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The Inter-Amexianninstitute for Space Scierree Education 1963 Summer Conference for Science- Teachers
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ff 653 July 65

## Reentry

## by

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## Introduction

Any manned space mission is composed of several distinct phases: the launch phase, the accomplishment of mission objectives phase, and the reentry phase. We are concerned here with the final or reentry phase of the space mission. It is during this period that the uncorrectable errors accrued during the space mission must be overcome in order to safely land the reentry vehicle at the desired landing site.

During the launch phase a vast amount of kinetic energy must be imparted to the space vehicle in order that it may escape from the earth and accomplish its mission. In returning to earth after completion of the mission, this same amount of kinetic energy must be dissipated in some manner. There are two ways in which this may be accomplished: rocket braking or atmospheric braking. The use of rocket braking allows the space vehicle to reenter the atmosphere at low speeds, thus essentially eliminating the reentry problem. However, due to the huge weignt penaiities associaíed with rocket oraking this method of energy dissipation is not practical. Atmospheric braking, wherein the vehicle kinetic energy is dissipated by aerodynamic drag as the vehicle travels through the atmosphere, is much more favorable as may be seen in Figure 1.

In this figure the vehicle uses aerodynamic braking below satellite velocity and the weight ratio presented is the
ratio of the vehicle weight at any velocity to the vehicle weight for reentry at satellite velocity. As shown, atmospheric braking, where the increasing weight is due to the additional thermal protection required as the reentry velocity increases, is much more favorable than rocket braking.

It should be noted that the aerodynamic or ablation curve shown here is quite conservative since it is assumed that the vehicle has a lift-drag ratio capability of one and that the stagnation point heating applies over the entire vehicle. Thus, the ablation weights shown are much higher than is actually the case. But, even under the worst possible conditions aerodynamic, ablation cooled braking is far superior to rocket braking.

Therefore, since minimum weight must be the prime consideration, atmospheric braking is required to dissipate the kinetic energy attained during a space mission. The reentry of a vehicle into the earth's atmosphere at high velocities thus becomes a major problem. In studying reentry one immediately can forsee several broad problem areas associated with the safe return of a manned vehicle from space to a desired landing site on the earth's surface. These are:
A. Deceleration Loads
B. Reentry Corridor Width
C. Aerodynamic Heating

1. Convective
2. Radiative
D. Range Control
3. Longitudinal
4. Lateral

For a given space mission, the above problem areas define the reentry vehicle design. Thus we have the Mercury, Gemini, Apollo, and Dynasoar reentry vehicles (each designed for a different mission).

It is the purpose here to define the reentry equations of motion and investigate each of the reentry problem areas in some detail. Means of alleviating these problems by proper vehicle design will be demonstrated. A range of initial reentry velocity from satellite velocity to a reentry velocity of $100,000 \mathrm{ft} / \mathrm{sec}$ is chosen to illustrate the increase in reentry vehicle refinement required by increasing reentry velocity.

Equations of Motion
Consider a vehicle at some point along its reentry path with the aerodynamic and gravitational forces acting as shown.


Writing the equation of motion along the flight path one obtains:

$$
\begin{equation*}
\frac{m d V}{d t}=-D-m g \sin \varphi \tag{1}
\end{equation*}
$$

or

$$
\begin{equation*}
\frac{1}{g} \frac{d V}{d t}=\frac{-e V^{2}}{2\left(\frac{W}{C_{D} A}\right)}-\sin \gamma \tag{2}
\end{equation*}
$$

The equation of motion perpendicular to the flight path is:

$$
\begin{equation*}
\frac{-m V^{2}}{R_{c}}=-m V^{2}\left(\frac{\cos r}{r}-\frac{1}{V} \frac{d Y}{d t}\right)=L-m g \cos r \tag{3}
\end{equation*}
$$

or

$$
\begin{equation*}
\left.\frac{d \gamma}{d t}=\frac{g}{V}\left\{\frac{\rho V^{2}}{2\left(\frac{W}{C_{D} A}\right.}\right) \frac{L}{D}-\cos \gamma\left(1-\frac{V^{2}}{r g}\right)\right\} \tag{4}
\end{equation*}
$$

Additional equations of interest are:

$$
\begin{equation*}
\frac{d h}{d t}=V \sin \gamma \tag{5}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{d\left(\frac{R}{r_{e}}\right)}{d t}=\frac{V}{r} \cos r \tag{6}
\end{equation*}
$$

Equations 2, 4, 5, and 6 must then be solved simultaneously to achieve a solution. In general a high speed digital computer is required to solve these equations of motion. Analytical solutions are available however for certain maneuvers of interest which will be discussed in more detail subsequently.

## VARIATION OF ENTRY VELOCITY WITH ENTRY ANGLE

- In analyzing the reentry phase of a space mission it is generally assumed that the initial reentry velocity is invariant with the initial reentry angle. To determine the applicability of this assumption let us consider the return of a vehicle from a circular orbit about the earth. For the condition of retrofire along the orbital track it may be shown that the initial reentry velocity and angle are related by the expression:

$$
\begin{equation*}
v_{1}^{2}=2 r_{0}^{2} g_{0}\left\{\frac{1}{r_{1}}-\frac{r_{1} \cos ^{2} r_{1}-r_{2}}{r_{1}^{2} \cos ^{2} r_{1}-r_{2}^{2}}\right\} \tag{7}
\end{equation*}
$$

```
where 1 = initial entry conditions
    o = earth surface conditions
    2 = point of retrofire conditions
```

Equation 7 is plotted in Figure 2 for a range of initial entry angles from 0 degrees to 12 degrees which should encompass the reentry corridor boundaries for vehicles of interest. As shown, initial reentry velocity is essentially invariant for return from orbital altitudes greater than about 2,000 miles. Therefore, we may say that, if the apogee altitude of the space mission is greater than 2,000 miles, the initial reentry velocity is independent of the initial reentry angle.

## ATMOSPHERIC MANEUVERS

A vehicle reentering the earth's atmosphere following a deep space mission must be capable of aerodynamic maneuvering since deviations from the desired entry conditions will occur due to guidance and control system inaccuracies encountered during the mission.

The trajectories traversed during reentry by an uncontrollable and a controllable vehicle are illustrated in Figure 3. As shown, a vehicle incapable of aerodynamic maneuvers will, in general, skip outside of the atmosphere. Maneuverability is required in order to reduce the maximum deceleration loads to tolerable levels and control the range traversed so as to land at some preselected site.

The "g" control maneuver requires that the reentry vehicle have the capability to vary the angle of attack from the angle for maximum lift to that for zero lift. To achieve range control, the vehicle must have the capability of either angle of attack variation or roll angle varlation such that the vehicle lift-drag ratio may be varied or modulated during the reentry period. Of the two maneuvers the range control maneuver is of the greater importance since the "g" control maneuver (discussed subsequently) would probably be used only in an emergency condition. The primary maneuvers considered for the range control problem are presented in Figure 4.

In general, the range control maneuver is initiated after the region of peak deceleration load and aerodynamic heating rate has been passed. Approximately minimum ranges are attainable by the constant heating rate and constant "g" maneuvers and maximum ranges by the constant altitude, equilibrium glide, and constant L/D-skip maneuvers. As shown, the constant $\dot{q}$ trajectory may not be maintained for a long period of time since the deceleration load is continually increasing during
this maneuver and will exceed acceptable limits. It is therefore not a practical maneuver, and is not considered further. The constant " $g$ " maneuver is one in which the deceleration load is maintained at a constant level by roll control. By maintaining the deceleration load at a comparatively high level, the vehicle's kinetic energy is rapidly dissipated and minimal ranges are obtained. An analytical solution is available for this maneuver with the assumption of an exponential densityaltitude relationship given by:

$$
\begin{equation*}
\frac{\rho}{\rho_{0}}=e^{-\beta h} \tag{8}
\end{equation*}
$$

where $\rho_{0}=$ sea level density and $\beta=$ scale height,$\frac{1}{23,500} \mathrm{ft}^{-1}$.

The reentry equations of motion of interest are:

$$
\begin{equation*}
\frac{1}{g} \frac{d V}{d t}=-\frac{G}{\sqrt{1+(L / D)^{2}}}-\sin r, \tag{9}
\end{equation*}
$$

where $G=$ constant deceleration load

$$
\begin{equation*}
\frac{d\left(\frac{R}{r_{e}}\right)}{d t}=\frac{V}{r} \cos \gamma \tag{i0}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{d h}{d t}=-V \sin \gamma \tag{11}
\end{equation*}
$$

Solution of equations 9, 10, and 11 with the aid of equation 8 yields the following result for the longitudinal range traversed during the maneuver.

$$
\begin{equation*}
\frac{R}{r_{e}}=\frac{\sqrt{1+(L / D)^{2}}}{G}\left\{\frac{\bar{v}_{1}^{2}-\bar{V}_{2}^{2}}{2}-\frac{\left(h_{1}-h_{2}\right)}{r_{e}}\right\} \tag{12}
\end{equation*}
$$

The term $\frac{\left(h_{1}-h_{2}\right)}{r_{e}}$ may be neglected in equation (12) as it is quite small in comparison to the velocity terms. Also, the maneuver end velocity, $\mathrm{V}_{2}$, may be shown to be given by

$$
\begin{equation*}
\overline{\mathrm{v}}_{2}^{2}=\frac{2}{\beta r_{e}} \frac{(1+G)}{\sqrt{1+(L / D)^{2}}} \tag{13}
\end{equation*}
$$

The constant altitude maneuver is initiated at the bottom of the pullout by a vehicle initially entering the atmosphere with positive Lift. The vehicle maintains a constant altitude flight path by either pitch or roll modulation.

The sum of the lift and centrifugal force is thus maintained equal to the vehicle weight. Eventually, the velocity decreases to the point where sufficient lift cannot be generated to satisfy this equality. A constant L/D trajectory is then flown to the landing site. For the constant-altitude maneuver the reentry equations of motion reduce to:

$$
\begin{align*}
& \frac{L}{W}=1-\bar{V}^{2}  \tag{14}\\
& \frac{d V}{d t}=-\frac{\rho V^{2} g}{2\left(\frac{W}{C_{D}^{A}}\right)} \tag{15}
\end{align*}
$$

$$
\begin{equation*}
\frac{d\left(\frac{R}{r_{e}}\right)}{d t}=\frac{V}{r_{e}} \tag{16}
\end{equation*}
$$

If the maneuver is controlled by roll angle modulation at constant drag coefficient, the longitudinal range may be shown to be given by:

$$
\begin{equation*}
\frac{R}{r_{e}}=-\frac{\left.2^{\left(\frac{W}{C_{D} A}\right.}\right)}{\rho r_{e^{g}}} \ln \frac{V_{2}}{V_{1}} \tag{17}
\end{equation*}
$$

where $V_{1}$ and $V_{2}$ are the velocities at the beginning and ending of the maneuver. Control of this maneuver by pitch modulation does not, in general, result in an analytic solution.

The equilibrium gilde maneuver is an approximation to the constant L/D maneuver. It is initiated at the point on the constant L/D pullup trafectory defined by the condition:

$$
\begin{equation*}
\rho=\frac{\left(1-\bar{v}^{2}\right)}{V^{2}} \quad\left(\frac{2 W}{S}\right) \quad\left(\frac{1}{C_{L}}\right) \tag{18}
\end{equation*}
$$

where $C_{L_{\text {e.g. }}}$ is a constant.
It is then assumed that the flight path angle is negligibly small and that the equality of aquation (14) nolds. If the maneuver is initiated at velocities greater than local sateilite velocity ( $\bar{V}>1$ ),$C_{L_{e . g}}$. is negative and the altitude increases (o decreases) with decreasing velocity. Note that at $\overline{\mathrm{V}}=1$ an infinite altitude is required. For velocities less than the local satellite value, $C_{L_{e}}$.g. is positive and the altitude
decreases ( $\rho$ increase ) with decreasing velocity. Obviously some transition maneuver is required in the region of $\overline{\mathrm{V}}=1$ to transfer from the negative equilibrium glide to the positive equilibrium glide maneuver.

It is assumed here that a minimum dynamic pressure of 10 psf is required for aerodynamic maneuvering. Therefore, a maximum range transition maneuver would be one which is carried out at a constant dynamic pressure of 10 psf . This combination of negative equilibrium glide, constant $q=10$ psf transition, and positive equilibrium glide should yield approximately maximum range for a wholly atmospheric maneuver.

The equations of motion for this case are given by equations 14 and 15 and may be shown to give the following expression for longitudinal range.

$$
\begin{equation*}
\frac{R}{r_{e}}=\frac{L}{D}\left\{1+\frac{1}{2} \ln \left[\frac{\overline{\mathrm{~V}}_{1}^{2}-1}{\left(\overline{\mathrm{~V}}_{2}^{2}-1\right)^{2}}\right]\right\} \tag{19}
\end{equation*}
$$

where $\overline{\mathrm{V}}_{2}=\sqrt{1+\frac{10}{\left(\frac{W}{C_{L}{ }^{A}}\right)}}$

The ifnal maximum range maneuver considered here is the true constant L/D maneuver which generally involves skipping outside the atmosphere. In this maneuver no control over the vehicle trajectory is available except that allowed by trimming the vehicle at a desired value of $L / D$ prior to initiating reentry. This value is then maintained throughout the reentry period.

## DECELERATION LOADS

The undershoot boundary; generally the steepest allowable descent into the atmosphere, is generally defined from consideration of both aerodynamic heating and deceleration loads. It is assumed here that this boundary may be defined by man's tolerance to deceleration loading only since improved technology can alleviate the heating problem but can do little to increase man's ability to withstand high deceleration or "g" loads.

The maximum deceleration load obtained during reentry on a vehicle with (L/D) $\max =1$ is presented in Figure 5 as a function of the initial reentry angle for several values of initial reentry velocity. It is assumed here that the crew can withstand no more than 12 g 's without serious damage. This value then defines the undershoot boundary. The overshoot boundary, indicated by the dashed line, is presented to demonstrate the reduction in the reentry corridor with increasing velocity.

It is of interest to consider the maximum deceleration load obtained at this overshoot boundary. This gives the minimum "g" tolerance required by crew members to safely return from a space mission. These are shown in Figure 6 as affected by initial reentry velocity and vehicle lift-drag ratio capability. As is to be expected, increasing the vehicle L/D capability decreases the overshoot " g " load for a given value of initial velocity. Also, note that the maximum deceleration load at the overshoot boundary increases with increasing initial entry velocity. Therefore, if the undershoot boundary
is "g" limited, a limiting velocity may be obtained where the undershoot and overshoot boundaries cross. For instance, if the undershoot boundary is defined by $G_{\max }=12 \mathrm{~g}$ 's, a limiting velocity of $93,300 \mathrm{fps}$ is obtained for a vehicle having infinite lift-drag ratio capability. Thus, entry at velocities in excess of $93,300 \mathrm{fps}$ is not possible unless man's tolerance to deceleration loads may be extended above 12 g's.

To further illustrate the reduction of maximum deceleration loads by increased vehicle $L / D$ capability, $G_{\max }$ is presented in Figure 7 in terms of initial entry angle for vehicles reentering the earth's atmosphere at escape speed (36,500 fps). As shown, the effect of increasing vehicle $L / D$ capability is to reduce the maximum deceleration load for a given initial reentry angle. Maximum benefits occur by increasing L/D from 0 to 0.5 with little advantage to increasing a vehicle's liftdrag ratio capability above 1.

The comparitively high reentry velocities required by trips to the outer planets or "fast" trips to nearby planets may necessitate development of modulation techniques to reduce the peak "g" loads and achieve acceptable reentry corridor widths. One such method, originated by Grant, is presented in Figure 8. In this method, the vehicle is considered to operate on the illustrated drag polar where:

$$
C_{L}=\left(C_{D_{\max }}-C_{D_{\min }}\right) \sin ^{2} \alpha \cos \alpha
$$

and

$$
\begin{equation*}
C_{D}=C_{D_{\min }}+\left(C_{D_{\max }}-C_{D_{\min }}\right) \sin ^{3} \alpha \tag{20}
\end{equation*}
$$

In this method the vehicle initially operates at maximum lift coefficient. When the deceleration loads have reached a specified level, the angle of attack is modulated towards zero so as to maintain the deceleration load at a constant value. This results in a decreased value of $G_{\max }$ as is illustrated in Figure 8 for the particular case of a vehicle with $(L / D)_{\max }=1 / 2$ entering the atmosphere at escape speed. Modulation from $C_{L}$ to (L/D) $\max$ is seen to decrease $G_{\max }$ from 12.7 to 10 "g's" with only a slight effect on the altitude and range at pullout. A further reduction in peak "g" to 7 "g's" is attainable by modulation from $C_{L_{\max }}$ to $C_{D_{\min }}$. This greatly affects the ranging capability of the vehicle as indicated by the altituderange curves of Figure 8. Increased modulation causes the vehicle to dig deeper into the atmosphere resulting in pullout at lower altitudes and velocities, thereby decreasing the vehicle's range capability.

Large increases in the reentry corridor width are attainable through the use of this technique of peak "g" reduction.

## Reentry Corridor Width

Reentry corridor width is defined as the difference between the perigee altitudes of the overshoot and undershoot trajectories neglecting atmospheric effects and considering the
earth as a point mass as shown by Figure 9. The overshoot and undershoot trajectories are thus simple conics defined by the initial entry velocity, $V_{1}$, and the angles $\gamma_{1_{0}}$ and $\gamma_{1_{u}}$. The angles $\gamma_{1_{0}}$ and $\gamma_{1_{u}}$ are obtained by definition of the corridor boundaries for a particular vehicle lift-drag ratio and ballistic parameter, $\frac{W}{C_{D}{ }^{A}}$.

In the previous section it was stated that the undershoot boundary is defined by the maximum deceleration load attained during reentry. It is felt that $G_{\max }=12$ "g's" is a reasonable undershoot limit for manned reentry vehicles. Therefore, the undershoot boundary is defined as that trajectory for which a vehicle reentering the atmosphere with a positive value of $L / D$ will receive a maximum deceleration load of 12 "g's".

It is much more difficult to define the overshoot boundary. However, for general purposes, it shall be defined as that trajectory for which the vehicle requires negative maximum lift capability to maintain constant altitude at the bottom of the pullout with initial entry at positive lift. Several other definitions will be discussed subsequently.

Utilizing the above boundary definitions, the overshoot and undershotot boundaries have been determined for a vehicle with $(L / D)_{\max }=1$ and are shown in Figure 10. Here, overshoot and undershoot initial reentry angles are presented in terms of initial reentry velocity. Note that at satelife velocity $\gamma_{1_{0}}=0^{\circ}$ since the centrifugal force is initially equal to the vehicle weight requiring that the vehicle remain within the
atmosphere. Also, a limit velocity of $83,000 \mathrm{fps}$ (corresponding to zero corridor width) is obtained using the present boundary definitions.

The reentry corridor width may be indicated by the parameter, $\Delta r_{i}=r_{i_{u}}-\gamma_{1_{0}}$. The effect of vehicle L/D capability on this parameter is shown in Figure ll. As is to be expected, large gains in corridor width are obtained by increasing the vehicle L/D capability from 0 to 1 . Note again that the limiting velocity is given by the initial reentry velocity for which $\Delta \gamma_{1}=0$; and that safe entry is impossible for values of $V_{1}$ greater than $93,300 \mathrm{fps}$ under the constraints of the present corridor boundary definitions.

It is of interest to determine the effects of several boundary definitions on both the corridor boundaries and the actual corridor widths. Figures 12 and 13 illustrate such effects for a vehicle reentering the atmosphere at escape speeds. Three definitions of each boundary are considered. The undershoot boundaries are defined as follows:
$+L / D$, uncontrolled -- The vehicle enters the atmosphere with positive $L / D$ and maintains constant L/D to pullout where a range control maneuver may be initiated.

Modulated, $C_{L_{\max }}$ to (L/D) $\max$.- The modulation technique discussed in Figure 8 is utilized to increase $\gamma_{I_{u}}$. The vehicle angle of attack may be modulated only from that for $C_{L_{\max }}$ to that for $(L / D)_{\max }$.

Modulated, $C_{L_{\max }}$ to $C_{L}=0-$ The same as the above with the modulation capability extended to $c_{L}=0$.

As shown, modulation greatly increases the undershoot entry angle and the benefits of this maneuver increase with increasing vehicle L/D capability. Note that without modulation the maximum benefits of increasing $L / D$ are achieved at $L / D=1$.

The overshoot boundaries are defined as follows:
$+L / D$, uncontrolled ( $h_{\text {skip }}=400 \mathrm{mi}$.) -- The vehicle enters the atmosphere with positive L/D and maintains this value throughout the reentry period (no maneuver capability). A 400 mile maximum skip altitude is chosen as the limit so as to prevent penetration of the Van Allen radiation belts.
$+L / D$, controlled -- The vehicle enters the atmosphere with positive L/D and utilizes full negative lift capability to maintain constant altitude at the bottom of the pullout.
${ }^{-} \mathrm{C}_{\mathrm{L}_{\max }}$-- The vehicle enters the atmosphere with full negative lift capability and barely remains within the atmosphere. This is the absolute limit for wholly atmospheric reentry maneuvers.

The $+L / D$, uncontrolled overshoot boundary is shown to be impractical for significant values of L/D (L/D ) .2). The
other two definitions show little variation with lift-drag ratio with the $-C_{L_{\text {max }}}$ definition yielding approximately $1 / 2$ degree more corridor width than the $+L / D$, controlled definition.

A design reentry corridor should have flexible boundaries capable of extension for emergency reentry conditions. Therefore, the overshoot boundary is generally taken as the $+\mathrm{L} / \mathrm{D}$, controlled case and the undershoot boundary as the $+L / D$, uncontrolled case.

The midcourse guidance and control required for a space mission is specified by a combination of system guidance and control capabilities and reentry vehicle corridor width capabilities. Since system guidance and control capability is essentially an engineering or "hardware" problem, we shall be concerned in this study only with the vehicle corridor width capability.

The actual reentry corridor widths corresponding to the boundary definitions of Figure 12 are presented in Figure 13 for an initial entry velocity of $36,500 \mathrm{fps}$. At this reentry velocity the maximum corridor width for the $+L / D$, uncontrolled case is about 15 miles and becomes zero for a value of $L / D=.475$. Therefore, the vehicle must be controlled at the overshoot boundary for safe reentry of vehicles with significant L/D capability. Modulation at the undershoot boundary is seen to significantly increase corridor width. This method will be required at the high entry velocities associated with short interplanetary trips to insure sufficient reentry corridor widths from the guidance and control standpoint. Note also that by
choosing the $+L / D$, controlled definition for the overshoot boundary, an additional 10 miles of corridor is maintained in reserve for emergency conditions wherein the $-C_{L_{\max }}$ entry made may be utilized.

Increasing the reentry velocity above escape velocity results in decreasing corridor widths. Eventually, the point of zero corridor width will be reached. A combination of propulsive and aerodynamic braking will be required for such high initial reentry velocities.

RANGE CONTROL
Longitudinal
Range control is achieved by aerodynamic maneuvers initiated after pullout to zero flight path angle. A reentry vehicle must be capable of reaching the desired landing site after initial reentry from any point within the reentry corridor. That is, the maximum range attainable by the vehicle entering the atmosphere at the undershoot boundary must be at least equal to the minimum range attainable by entry at the overshoot boundary. In addition, some range overlap (see Figure 14) is desirable to offset errors in initial reentry time since time errors introduce range errors due to the earth's rotation. The degree of range overlap required is of course defined by the allowable mission time errors and the angle of the reentry plane with respect to the equatorial plane.

Before proceeding with a discussion of the range overlap capabilities of reentry vehicles, it is desirable to consider
the longitudinal range capabilities of a vehicle carrying out the aerodynamic maneuvers of Figure 4.

The range attainable by the constant $L / D$ reentry maneuver is shown in Figure 15 for an initial velocity of $36,500 \mathrm{fps}$ (escape speed). The boundaries considered in this figure are a 10 g undershoot boundary and a 400 mlle skip overshoot boundary. Range control is obtained by selecting the appropriate L/D for a given initial reentry angle. The slope of the lines of constant L/D indicate that large errors in range would occur for small errors in setting the trim angle of attack (reentry L/D). Note also that, for these boundary conditions, safe reentry is not available for values of L/D greater than . 475 . These limitations in range control and usable vehicle $L / D$ demonstrate that control over the reentry trajectory must be utilized for the safe reentry and landing of a manned vehicle returning from a space mission at escape or higher speeds.

It thus becomes necessary to consider the longitudinal range attainable by the controlled atmospheric maneuvers of Figure 4. As an illustrative example we shall consider return from a lunar mission with a reentry vehicle having a maximum $L / D$ capability of $1 / 2$. The longitudinal ranges attainable by this venicie throughout the reentry corridor are presented in Figure 16. An entry angle of $-5.25^{\circ}$ represents the overshoot boundary and $7.6^{\circ}$, the undershoot boundary. The two limiting curves represent essentially the limiting ranges of which the vehicle is capable. The maximum range curve is obtained by utilizing a constant altitude maneuver at the bottom of the
initial reentry pullout. Constant altitude is maintained until sufficient kinetic energy has been dissipated such that, if a constant $L / D=1 / 2$ trajectory is then initiated, the vehicle will skip, outside the atmosphere to a maximum altitude of 400 miles. This maximum range maneuver is, however, critically dependent on the velocity and path angle at which the skip is initiated. Therefore, this maneuver is considered to be too susceptible to large uncorrectable range errors to be a reliable method of range control. The greatest range attainable by a wholly atmospheric maneuver is given by the equilibrium gilde curve, which gives ranges of about 6,000 miles near the undershoot boundary. This appears adequate for return from the lunar mission. The constant altitude curve is obtained by maintaining constant altitude at the bottom of the pullout for as long as possible. As shown, quite short ranges are obtained at the undershoot boundary. The minimum indicated ranges were obtained by the use of a constant 10 " $g$ " deceleration load maneuver.

The range overlap for this vehicle and mission varies from 400 to 12,000 miles depending on the maximum range maneuver utilized. The equilibrium glide maneuver, yielding 4,000 miles range overlap, is probably the best operational maneuver.

The effect of increasing the initial reentry velocity on the longitudinal range overlap is presented in Figure l7. Here, the maximum range is obtained by the constant altitude maneuver and the minimum range by the constant " $g$ " maneuver. The trends would be the same for the equilibrium glide maneuver although this maneuver is not considered here. As shown, increasing
either the initial reentry velocity or the vehicle $L / D$ capability yields an increased range overlap. Both results are to be expected. One effect of increasing velocity is to bring the corridor boundaries closer together, thus the initial maneuver conditions are more nearly the same for the maximum and minimum range maneuvers. The effect of increasing the vehicle L/D capability yields increased maneuverability and hence, longer ranges. Of course, this increased range overlap with increased reentry velocity may be offset by the increased range overlap requirements of the deep space missions associated with these reentry velocities.

It is necessary to determine the ability of pilots to fly the maneuvers of interest. Many pilot simulation studies have been carried out in an effort to define optimum methods of pllot control over the reentry trajectory. The ranges attainable by some of these are shown in Figure 18 for reentry at escape speed of a vehicle with a maximum $L / D$ capability of $1 / 2$. The reentry guidance and control techniques considered are: the reference trajectory technique, the repetitive prediction technique, and the pilot controlled technique. In the reference trajectory procedure the control feedbacks were developed for successful operation of the system. The repetitive prediction system utilized a rapidtime analog computer to predict the range capability from the present conditions. The pilot's intelligence and learning capabilities were used to provide the guidance logic and the control commands in the pilot controlled technique.

The repetitive prediction technique was shown to yield very good control of the initial peak deceleration and, hence, minimal ranges. The maximum range maneuver utilized here for the simulator studies is an arbitrary maneuver wherin a pullup, initiated upon entering the atmosphere, is terminated at an altitude of approximately 250,000 feet with a velocity of about $26,000 \mathrm{ft} / \mathrm{sec}$. This flight plan is most nearly approximated by the equilibrium glide maneuvers of the present study, indicated by the dashed line on Figure 18. The piloted maximum ranges are shown to be much less than the theoretical values. It appears, however, that with further system refinement and pilot schooling, ranges quite close to the theoretical values may be obtained.

Lateral
Significant lateral range capability is required of a vehcile returning from a deep space mission since the reentry plane angle may vary considerably from the nominal due to midcourse guidance corrections. In addition, time errors may introduce large lateral range requirements if the reentry plane is not the equatorial plane.

A vehicle's lateral range capability may be shown to increase with increasing entry velocity. Therefore, determination of a vehicle with sufficient $L / D$ capability to satisfy lateral range requirements for reentry at satellite velocity will apply for reentry at all higher velocities. The effect of the vehicle's maximum $L / D$ capability on its lateral range capability
is shown in Figure 19 for entry at satellite velocity. Lateral range is seen to increase rapidly with L/D capability reaching $1 / 4$ the earth's circumference at an L/D of about 3.5. This is the maximum lateral range capability which could be required since a vehicle reentering the atmosphere in the equatorial plane could reach either of the poles. Similarly, a vehicle reentering in a polar plane could reach any point on the earth by proper combination of the vehicle lateral and longitudinal range capabilities. A vehicle's design lateral range capability will, of course, depend on the allowable range errors introduced by the space vehicle's guidance and control systems.

As an example of the $L / D$ required from the lateral range standpoint, let us consider reentry from a polar orbit. The range required to reach a particular landing site either once or twice daily is presented in Figure 20. Also shown is the L/D required. Return to the continental United States may be accomplished twice daily by a vehicle with L/D of approximately 0.9 . Thus, the particular mission requirements are seen to reduce the L/D required for lateral range from 3.5 to 0.9 . This effect of specific mission requirements reducing the lateral range requirements may be expected to hold for the deep space missions with high reentry velocities. It cannot be stated at the present time exactly what lateral range capability will be required for reentry at high velocities. Therefore, no attempt is made here to define specific lateral range capabilities or requirements for these missions.

AERODYNAMIC HEATING

## Convective

Aerodynamic heating is the heating of a reentry vehicle due to the friction of the air as the vehicle passes through it. It may be divided into two components: convective and radiative. Convective heating is the dominant source of heating at the lower reentry velocities with radiative heating becoming the dominant source at the higher velocities.

The convective stagnation point heating rate equation may be written approximately as:

$$
\begin{equation*}
\dot{q}_{c}=\frac{K \sqrt{\rho} V^{3} \cdot 15}{\sqrt{R_{n}}} \tag{21}
\end{equation*}
$$

where $\rho=$ atmospheric density

$$
\begin{aligned}
& V=\text { velocity } \\
& R_{n}=\text { vehicle nose radius }
\end{aligned}
$$

The total stagnation point convective heat load is obtained by integration of equation (15) over the reentry time period. Note that convective heating is inversely proportional to the square root of the vehicle nose radius. For this reason, reentry vehicles operating in the range of entry velocities where convective heating dominates have blunt nose shapes. Notable examples are the Mercury, Gemini, and Apollo vehicles.

Since a knowledge of the vehicle shape is required for a complete heating analysis, we are concerned here only with the stagnation point heating loads so as to maintain the generality
of the study. The stagnation point heating is quite sufficient to indicate the effects of vehicle $L / D$ capabilities, reentry velocities, and atmospheric maneuvers on the aerodynamic heating problem. In particular, design of a reentry vehicle heat shield is dependent on the maximum heating rates and total heat loads expected to be encountered during reentry. These are, of course, the stagnation point values.

The effect of the particular atmospheric maneuver utilized during reentry on the maximum heating rate and total heat load is presented in Figure 21 for a vehicle with $L / D=1 / 2$ reentering the earth's atmosphere at escape speed. Since all the maneuvers considered here are initiated after pullout, the same value of maximum heating rate and maximum "g" load apply to each maneuver for the same initial reentry conditions. As is to be expected, the minimum range, constant "g" maneuver yields minimum total heat loads for a given initial entry condition. The maximum range, constant L/D, skipping maneuver yields the greatest total heat loads. Note that maximum total heat loads occur at the overshoot boundary with low values of $\dot{q}_{c_{m a x}}$ while minimum total heat loads and high values of $\dot{a}_{c_{\max }}$ occur at the undershoot boundary. This is due to the fact that at the undershoot boundary the vehicle dips deeper into the atmosphere into regions of higher atmospheric density than at the overshoot boundary. Since the vehicle is in a higher density region when the maneuver is initiated, it will, of course, dissipate its kinetic energy much faster than at the overshoot boundary resulting in lower total heat loads.

The maximum convective heating rates (undershoot) and the maximum convective total heat loads (overshoot) are presented in Figures 22 and 23 to demonstrate the effects of vehicle L/D capability and initial reentry velocity on these quanities. The expected results of increasing maximum heating rates and total heat loads with increasing vehicle L/D capability or initial reentry velocity are obtained. Increasing L/D capability yields increased values of $\dot{q}_{c_{\text {max }}}$ since the vehicle resultant force coefficient is reduced thereby causing the undershoot boundary pullout to occur at lower altitudes (higher density) for the same initial reentry velocity and angle. Increased total heat loads result due to the increased maneuverability of the higher L/D vehicle. Increased initial reentry velocity yields increased heating rates since the heating rate is more velocity dependent than density dependent as shown by equation (15). Also, the total heat load is increased due primarily to the greater kinetic energy of the higher velocity vehicles which must be dissipated within the atmosphere.

Obviously, the convective heating problem becomes more severe as more sophisticated space missions are undertaken.

## Radiative

The radiative stagnation point heating rate equation may be written approximately as:

$$
\begin{equation*}
\dot{q}_{r}=K \rho{ }^{r} V^{s} R_{n} \tag{22}
\end{equation*}
$$

where the exponents $r$ and $s$ are dependent on the velocity regime in which the vehicle is flying. Radiative heating is considered
to be negligibly small at velocities less than $25,000 \mathrm{fps}$ and approaching that required to dissipate the entire kinetic energy at high velocities. Note that the radiative stagnation point heating is directly proportional to the vehicle nose radius. This indicates that pointed shapes are optimum from the radiative heating standpoint. Thus, once again, the designer is faced with a tradeoff problem. The vehicle nose shape must be designed from consideration of both convective and radiative heating and is quite mission dependent.

The total heat loads as obtained for a vehicle with a one foot nose radius and a maximum $L / D$ capability of one are presented in Figure 24 for two types of entry: maximum L/D and maximum lift coefficient. This figure indicates the superiority of entry at maximum lift coefficient from consideration of heating only. Of course, operation at $C_{L_{\max }}$ results in a loss of corridor width and range capability since the vehicle is operating at a value of $L / D$ less than the maximum value.

Of particular interest is the role played by radiative heating. For the case shown here, radiative heating is negligible in comparison to convective heating for initial entry velocities less than about $45,000 \mathrm{fps}$. As the reentry velocity is increased however, radiative heating quickly becomes the dominant heating factor. Therefore a completely different vehicle design is required for a reentry velocity of $70,000 \mathrm{fps}$ than for a reentry velocity of 40,000 fps.

## CONCLUDING REMARKS

The major problems in reentry--deceleration loads, corridor width, range control, and heating loads--have been discussed. The effects of reentry vehicle $L / D$ capability, initial reentry velocity, and atmospheric maneuvers on these problems have been demonstrated. No clear cut conclusions may be drawn here since the only intention has been to bring to the reader a better knowledge of what is meant by the word "reentry". Also, it has been the purpose here to point out the complexities of designing a reentry vehicle.

The effect of increased $L / D$ capability was shown to be advantageous from the standpoint of corridor width, deceleration loads, and range control, but most disadvantageous from the standpoint of aerodynamic heating. The role played by vehicle nose shape in reentry has been demonstrated by the opposing effects of convective and radiative heating. The necessity for a vehicle to have the capability of maneuvering within the atmosphere was shown. The effect of initial reentry velocity was to indicate that in all probability a specific reentry vehicle must be defined for each space mission or reentry velocity range. Thus, the trend of the past--the Mercury vehicle for the orbital mission and the Apollo vehicle for the lunar mission--will probably be continued in the future. Increased L/D capability will be demanded by the more stringent requirements of reentry at the hyperbolic velocities associated with interplanetary missions. Finally, propulsion may well be required to slow a vehicle returning from fast interplanetary trips to velocities for which the vehicle may safely reenter the atmosphere.

APPENDIX A

SYMBOLS

| A | vehicle reference area |
| :---: | :---: |
| $\mathrm{C}_{\mathrm{D}}$ | drag coefficient |
| $\mathrm{C}_{\mathrm{L}}$ | lift coefficient |
| $\mathrm{C}_{\mathrm{L}_{\mathrm{e}, \mathrm{~g}}}$ | lift coefficient used during the equilibrium glide maneuver |
| D | drag |
| $g$ | gravitational acceleration |
| h | altitude |
| L | Iift |
| L/D | lift-drag ratio |
| m | vehicle mass |
| R | longitudinal range |
| $\mathrm{R}_{\mathrm{c}}$ | radius of curvature of reentry flight path |
| r | radial distance from earth center |
| $\mathrm{r}_{\mathrm{e}}$ | earth radius |
| $t$ | time |
| v | velocity |
| w | vehicle weight |
| $\alpha$ | angle of attack |
| $\gamma$ | reentry angle |
| $\rho$ | atmospheric density |

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Figure 1. - Comparison of rocket and aerodynamic braking.


Figure 2. - Reentry velocity and angle for return from a circular orbit.
"Uncontrolled"
Fixed Attitude Fixed Altitude Constant L/D

Figure 4. - Atmospheric maneuvers.


## No







Figure 7. - Effect of vehicle L/D capability in maximum deceleration loads. $\left(V_{i}=36,500 \mathrm{fps}\right)$.

$\begin{aligned} & \text { Figure 8. - Modulated reentry } \\ & \gamma_{i}=-8.11 .\end{aligned}$
$\mathrm{v}_{\mathrm{i}}=36,500 \mathrm{fps}$.

figure 9. - Reentry corridor width.




- Figure 12. - Dffect of vehicle ID ospability and atmospheric maneureming on the xecatiy comidor bombaries,
$\left(Y_{2}=36,500 \mathrm{fP}\right)$


Figure 13. - Effect of vehicie $\mathrm{L} / \mathrm{D}$ capability and atmospheric maneuvering on the reentry corridor width.

$$
\left(\mathrm{v}_{\mathrm{i}}=36,500 \mathrm{fps}\right)
$$



Figure 15. - Longitudinal range attainable by the constant L/D maneuver.




Figure 19. - Vehicle maximum lateral range capability.





