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### Development of an Algebraic Stress/Two-Layer Model for Calculating Thrust Chamber Flow Fields

By

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### ABSTRACT

Following the consensus of a workshop in Turbulence Modeling for Liquid Rocket Thrust Chambers, the current effort was undertaken to study the effects of second-order closure on the predictions of thermochemical flow fields. To reduce the instability and computational intensity of the full second-order Reynolds Stress Model, an Algebraic Stress Model (ASM) coupled with a two-layer near wall treatment was developed. Various test problems, including the compressible boundary layer with adiabatic and cooled walls, recirculating flows, swirling flows and the entire SSME nozzle flow were studied to assess the performance of the current model. Detailed calculations for the SSME exit wall flow around the nozzle manifold were executed. As to the overall flow predictions, the ASM removes another assumption for appropriate comparison with experimental data, to account for the non-isotropic turbulence effects.

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- Improve Predictive Capabilities of Turbulent Transport in Thrust Chamber
- Non-Isotropic and Compressibility Effects are the Focus of the Study
- Simplified Reynolds Stress Modeling
- Further Modeling in Turbulent Transport of Thermal Energy and Chemical Species -  $u'_iC'$  and  $u'_iT'_i$  etc.

Motivation and Objective
• Higher Order Models Are Desirable For Calculating Thrust Chamber
Flow Fields 1991 Thrust Chamber Turbulence Modeling Workshop
• To Develop a Simplified 2nd-Order Turbulence Model For Thrust Chamber
Flow Calculation Near wall treatment Efficiency and stability

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### APPROACH

- Modeling any unknown in terms of Reynolds stress, the mean strain PDE's for Reynolds stress  $\overline{u_i u_j}$  can be derived. rate etc. •
- Simplifications of the Differential Reynolds stresses Equations Algebraic Stress Model(ASM)
- Non-linear constitutive relations (Spezial)

APPROACH (DRS Equation)	
<ul> <li>Differential Reynolds Stress Equation</li> </ul>	
$\frac{D}{Dt} \frac{pu_i^* u_j^*}{Dt} = P_{ij} + D_{ij} + \pi_{ij} + C_{ij} - \varepsilon_{ij}$	
$P_{ij} = -\overline{\rho} \left[ \overline{u_i^{i} u_k^{i}} \frac{\partial \overline{u_j}}{\partial x_k} + \overline{u_j^{i} u_k^{i}} \frac{\partial \overline{u_j}}{\partial x_k} \right]$	production
$D_{ij} = \frac{\partial}{\partial x_k} \left[ \bar{\rho} \ \bar{u}_i^{T} \bar{u}_k^{T} + \delta_{ik} \ \bar{u}_j^{T} \bar{p}^{T} + \delta_{jk} \ \bar{u}_i^{T} \bar{p}^{T} - (\mu S_{ik} u_j^{T} + \mu S_{jk} u_i^{T}) \right]$	diffusion
$\pi_{ij} = p'(\frac{\partial u'_i}{\partial x_i} + \frac{\partial u'_j}{\partial x_i})$ pressure-9	train correlation
$C_{ij} = -\left[ \vec{u}_i \frac{\partial \vec{P}}{\partial x_j} + u_j \frac{\partial \vec{P}}{\partial x_i} \right]$	compressibility
$\varepsilon_{ij} = \mu \left[ S_{ik} \frac{\partial u_j^{"}}{\partial x_k} + S_{jk} \frac{\partial u_i^{"}}{\partial x_k} \right]$	dissipation

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**Turbulence model costants** 

$$C_{\epsilon 1}$$
  $C_{\epsilon 2}$   $C_k$   $C_\epsilon$   $C_1$   $C_2$ 

1.45 1.92 0.22 0.15 2.5 0.5



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**METHODOLOGIES**  

$$\begin{aligned}
\nu_{t} = u' l' = C_{\mu} \kappa^{\mu} l_{\mu} = C_{\mu} \frac{k^{2}}{\epsilon} = C_{\mu} \frac{l_{e}}{\kappa^{2} l_{e}} \kappa^{2}
\end{aligned}$$
within Intertia Sublayer  $\epsilon = C_{\mu}^{24} \frac{k^{2} l_{\mu}}{l_{\mu}} = \frac{k^{2} l_{e}}{l_{e}}$ 

$$l_{\mu} = C_{l} y \left[ 1 - \exp\left(-\frac{R_{e}}{A_{\mu}}\right) \right] \longrightarrow \text{ to be used in Eddy Viscosity}$$

$$l_{e} = C_{l} y \left[ 1 - \exp\left(-\frac{R_{e}}{A_{\mu}}\right) \right] \longrightarrow \text{ to be used in k-equation}$$

$$l_{e} = C_{e} r C_{\mu}^{24}$$
Matching at  $R_{e} = \frac{k^{2} l_{e}}{v}$ 



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# **Compressible Flat Plate Flow**



# **Compressible Flat Plate Flow**

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**Backward Facing Step** 

flow (9:1), with data from [31]

Fig. 5 Reynolds stress profiles for the backward-facing step turbulent

### **Confined Swirling Flows**









**Confined Swirling** 

**H'lows** 

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Fig. 9a Contour of Mach number for  $k - \epsilon$  with two-layer model.



Fig. 10a Contour of temperature for k- $\epsilon$  with two-layer model.

SSME Nozzles



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Fig. 9b Contour of Mach number for ASM with two-layer model.



Fig. 10b Contour of temperature for ASM with two-layer model.

### **8 - STEP REACTIONS**

		Α	Ν	E
$M + O_2$	==== 0 + 0	0.72000E+19	-1.0000	117908
$M + H_2$	==== H + H	0.55000E+19	-1.0000	103298
$M + H_2O$	==== H +OH	0.52000E+22	-1.5000	118000
O + H	==== OH	0.71000E+19	-1.0000	0.
$H_2O + OH$	==== H <sub>2</sub> O + H	0.58000E+14	0.0000	18000
$H_2 + OH$	$===H_{2}O + H$	0.20000E+14	0.0000	5166
0 <sub>2</sub> + H	==== OH + O	0.22000E+15	0.0000	16800
$H_{2} + O$	==== OH +H	0.75000E+14	0.0000	11099

 $\begin{aligned} k &= AT^N exp(-E_{RT}) \\ \text{with } k \text{ in } cm^3 \cdot mole^{-1} \cdot s^{-1} \text{ and } E \text{ in } cal \cdot mole^{-1} \end{aligned}$ 

From R.C.Rogers and Chinitz, 'Using a global hydrogen-air model in turbulent flow calcultions ', AIAA J., vol. 21, pp. 586-592, 1983.





CONTOUR OF TEMPERATURE



Figure A.5 Sample SSME Nozzle Flow Inputs and Results --- Turbulent, 8-Step Kinetics. Exp. ISP=453.3sec ISP = 452.78 sec.

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Fig. 12, SSME wall contour and geometry at exit

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Fig. 17(a), Pressure levels along the wall near the nozzle exit using the ASM and  $k-\epsilon$  models



Fig. 17(b), Effects of wall temperature on the wall pressure



Fig. 18, Contour of pressure using 75% of the chamber pressure level





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	Summaries
•	The Algebraic Stress Model Removes the Isotropic Turbulence Assumption for the Eddy Viscosity Type Models
•	Improved On the Reynolds Stresses Predictions
٠	The ASM Does Not Improve Too Much On SSME Nozzle & Outlet Flows
	RSM Other Mechanisms Shock- Boundary Layer Interactions Entrainment Issues
•	3-D Calculations Are Desirable

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