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CHARTS FOR CONICAL AND TWO-DIMENSIONAL

OBLIQUE-SHOCK FLOW PARAMETERS IN HELIUM AT

MACH NUMBERS FROM ABOUT 1 TO 100

By Arthur Henderson, Jr., and Dorothy O. Braswell

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SUMMARY

Charts are presented for the determination of certain flow properties in the flow field and on the surface of cones and wedges at Mach numbers from about 1 to 100 in a perfect gas with a ratio of specific heats of 5/3. In addition, a table of the isentropic subsonic flow parameters is included.

INTRODUCTION

The present charts supplement the work of references 1 and 2. Reference 1 presents ideal-gas flow parameters for helium. It includes tables of the one-dimensional isentropic flow relations and the normal shock relations for Mach numbers from 1 to 40.

For Mach numbers of 12, 16, 20, and 24, reference 1 also presents charts of the shock angle, Mach number behind the shock, and staticpressure ratio across the shock for two dimensionally deflected flow (weak shock solution only). In addition, charts of the shock angle, surface velocity ratio, and ratio of surface static pressure to freestream static pressure for conical flow are presented.

Reference 2 presents charts of the isentropic and normal shock relations at Mach numbers up to 20 and charts of several two-dimensional oblique-shock flow parameters at Mach numbers from about 5 to 25. Reference 2 has received only limited distribution.

The present paper presents charts of several flow parameters for both conical and two-dimensional oblique flow in helium at Mach numbers from about 1 to 100. A table of the isentropic subsonic flow parameters in helium is also given inasmuch as they are not presented elsewhere in the literature.

SYMBOLS

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А	cross-sectional area of stream tube or channel
a	speed of sound
cp	specific heat at constant pressure
cv	specific heat at constant volume
М	Mach number
р	pressure
đ	dynamic pressure
т	absolute temperature
v	velocity of flow
V _{max}	maximum velocity obtained by expanding to zero pressure
$\gamma = c_p / c_v$	
δ	flow-deflection angle across oblique shock
e	shock wave angle
μ	Mach angle
ν	Prandtl-Meyer angle
ρ	density
σ	semivertex angle of cone
Subscripts:	
с	cone surface
n	normal shock
t	stagnation conditions

1 immediately upstream of shock

2 immediately downstream of shock

Superscript:

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condition where local velocity equals speed of sound

RESULTS AND DISCUSSION

The conical-flow results were computed on an IBM 704 electronic data processing machine by the procedure outlined in reference 3. A cone-surface Mach number was assumed for various values of σ and the flow field was iterated out toward the shock to determine the shock angle ϵ and free-stream Mach number M_1 . Figures 1 and 2 present the variations of ϵ and M_c , respectively, with M_1 for various values of σ . Figures 3, 4, and 5 present the variations of $p_{t,c}/p_{t,1}$, p_c/p_1 , and ρ_1/ρ_c , respectively, with M_1 .

By using the results shown in figures 1 and 2, the results shown in figure 3 were obtained from

$$\frac{p_{t,c}}{p_{t,l}} = \left[\frac{(\gamma+1)M_{l}^{2}\sin^{2}\epsilon}{(\gamma-1)M_{l}^{2}\sin^{2}\epsilon+2}\right]^{\frac{\gamma}{\gamma-1}} \left[\frac{\gamma+1}{2\gamma M_{l}^{2}\sin^{2}\epsilon-(\gamma-1)}\right]^{\frac{1}{\gamma-1}}$$
(1)

The curves of figure 4 were obtained from

$$\frac{p_{c}}{p_{1}} = \frac{p_{c}}{p_{t,c}} \left(\frac{p_{1}}{p_{t,1}}\right)^{-1} \frac{p_{t,c}}{p_{t,1}}$$
(2)

where

$$\frac{P_{c}}{P_{t,c}} = \left(1 + \frac{\gamma - 1}{2} M_{c}^{2}\right)^{-\frac{\gamma}{\gamma - 1}}$$
(3)

$$\frac{\mathbf{p}_{1}}{\mathbf{p}_{t,1}} = \left(1 + \frac{\gamma - 1}{2} \mathbf{M}_{1}^{2}\right)^{-\frac{\gamma}{\gamma - 1}}$$
(4)

The curves of figure 5 were obtained from

$$\frac{\rho_{1}}{\rho_{c}} = \frac{\rho_{1}}{\rho_{t,1}} \left(\frac{\rho_{c}}{\rho_{t,c}} \right)^{-1} \frac{\rho_{t,c}}{\rho_{t,1}}$$
(5)

where

$$\frac{\rho_{\rm c}}{\rho_{\rm t,c}} = \left(1 + \frac{\gamma - 1}{2} M_{\rm c}^2\right)^{-\frac{1}{\gamma - 1}} \tag{6}$$

$$\frac{\rho_{l}}{\rho_{t,l}} = \left(1 + \frac{\gamma - 1}{2} M_{l}^{2}\right)^{-\frac{1}{\gamma - 1}}$$
(7)

$$\frac{\rho_{t,c}}{\rho_{t,l}} = \frac{P_{t,c}}{P_{t,l}}$$
(8)

The two-dimensional results were calculated from the oblique-shock equations which may be found in reference 4, for instance. Figures 6 to 10 present charts of ϵ , M_2 , $p_{t,2}/p_{t,1}$, p_2/p_1 , and ρ_1/ρ_2 , respectively, plotted against M_1 for various values of flow-deflection angle δ .

It is of interest to note that the minimum density ratio on the cone surface (fig. 5) is about 0.239 as compared with the value of 0.25 for two-dimensional flow. For comparison, it is noted that the minimum density ratio at the stagnation point behind a normal shock at $M = \infty$ is about 0.229.

Only weak shock solutions are presented for the cones, since after the occurrence of shock detachment the mathematical solution is physically unrealistic.

Both weak and strong shock solutions are presented for the twodimensional flow results. Although the two-dimensional results apply to conditions immediately behind the shock, when the shock is attached (weak shock solution), the results are also applicable on the surface of a wedge whose angle is identical to the flow-deflection angle across the shock, since the flow field is then uniform. Conditions immediately behind the two-dimensional shock apply to both two- and three-dimensional flow for both the weak and strong shock solutions.

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Often the compression ratio p_2/p_c , the ratio of the pressure immediately behind the conical shock to the pressure on the cone surface, is desired. This ratio can be obtained from

$$\frac{\mathbf{p}_2}{\mathbf{p}_c} = \frac{\mathbf{p}_2}{\mathbf{p}_1} \left(\frac{\mathbf{p}_c}{\mathbf{p}_1} \right)^{-1}$$
(9)

The ratio p_c/p_1 is given directly in figure 4, and the ratio p_2/p_1 is determined from figure 9. The value of δ to be used in figure 9 is the flow-deflection angle immediately behind the conical shock and is determined as follows: For given σ and M_1 , the cone shock angle ϵ is determined from figure 1. The flow-deflection angle δ immediately behind this shock is found from figure 6 for the given value of M_1 and the determined cone shock angle ϵ .

The scales on the charts have been chosen to allow the ordinates to be read to an accuracy of at least 0.5 to 1.0 percent at Mach numbers from 1 to 40, and at least 1.0 to 2.0 percent at Mach numbers from 40 to 100.

In general, the results presented here will be used in conjunction with the isentropic and normal shock relations which may be found in references 1 or 2 for Mach numbers up to 40. At Mach numbers above 40, the hypersonic approximation applied to the isentropic and normal shock relations yields the following expressions which may be used to a high degree of accuracy. In each case, the second form of the equation applies for $\gamma = 5/3$.

$$\frac{p}{p_{t}} = \left(\frac{\gamma - 1}{2} M^{2}\right)^{-\frac{\gamma}{\gamma - 1}} = \frac{15 \cdot 5884}{M^{5}}$$
(10)

$$\frac{\rho}{\rho_{\rm t}} = \left(\frac{\gamma - 1}{2} \,{\rm M}^2\right)^{-\frac{1}{\gamma - 1}} = \frac{5 \cdot 1962}{{\rm M}^3} \tag{11}$$

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$$\frac{\mathrm{T}}{\mathrm{T}_{\mathrm{t}}} = \left(\frac{\gamma - 1}{2} \,\mathrm{M}^2\right)^{-1} = \frac{3.0000}{\mathrm{M}^2} \tag{12}$$

$$\frac{a}{a_{t}} = \left(\frac{\gamma - 1}{2} M^{2}\right)^{-1/2} = \frac{1.7321}{M}$$
(13)

$$\frac{q}{p_{t}} = \frac{\gamma}{\gamma - 1} \left(\frac{\gamma - 1}{2} M^{2} \right)^{-\frac{1}{\gamma - 1}} = \frac{12.9905}{M^{3}}$$
(14)

$$\frac{A^{*}}{A} = \left(\frac{\gamma + 1}{2}\right)^{\frac{\gamma + 1}{2(\gamma - 1)}} M\left(\frac{\gamma - 1}{2} M^{2}\right)^{-\frac{\gamma + 1}{2(\gamma - 1)}} = \frac{16.0000}{M^{3}}$$
(15)

$$\frac{V}{V_{\text{max}}} = 1 - \frac{1}{(\gamma - 1)M^2} = 1.0000 - \frac{1.5000}{M^2}$$
(16)

$$\frac{\mathbf{v}}{\alpha^{\star}} = \left(\frac{\gamma + 1}{\gamma - 1}\right)^{1/2} \left[1 - \frac{1}{(\gamma - 1)M^2}\right] = 2.0000 - \frac{3.0000}{M^2}$$
(17)

$$\mu = \frac{57 \cdot 2958}{M}$$
(18)

$$v = 57.2958 \left(1.5708 - \frac{3.0000}{M} \right)$$
(19)

$$M_{2,n} = \left(\frac{\gamma - 1}{2\gamma}\right)^{1/2} = 0.4472$$
 (20)

$$\frac{p_{2,n}}{p_{1}} = \left(\frac{2\gamma}{\gamma+1}\right) M_{1}^{2} = 1.2500 M_{1}^{2}$$
(21)

$$\frac{\rho_{2,n}}{\rho_{1}} = \frac{\gamma + 1}{\gamma - 1} = 4.0000$$
(22)

$$\frac{T_{2,n}}{T_{1}} = \frac{2\gamma(\gamma - 1)}{(\gamma + 1)^{2}} M_{1}^{2} = 0.3125 M_{1}^{2}$$
(23)

$$\frac{P_{t,2,n}}{P_{t,1}} = \frac{\gamma + 1}{\gamma - 1} \left[\frac{(\gamma + 1)^2}{2\gamma(\gamma - 1)M_1^2} \right]^{\frac{1}{\gamma - 1}} = \frac{22.8970}{M_1^3}$$
(24)

$$\frac{P_{t,2,n}}{P_{1}} = \frac{\gamma + 1}{2} \left[\frac{(\gamma + 1)}{4\gamma} \right]^{\frac{1}{\gamma - 1}} M_{1}^{2} = 1.4689 M_{1}^{2}$$
(25)

References 1 and 2 do not present the subsonic isentropic flow relations for helium. Inasmuch as these relations are of value in the stagnation region of blunt bodies and in the subsonic portion of boundary layers, they are given in table I.

Langley Research Center, National Aeronautics and Space Administration, Langley Field, Va., March 8, 1961.

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REFERENCES

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TABLE I. - SUBSONIC ISENTROPIC FLOW RELATIONS FOR HELIUM

M	p/p _t	۵/۹t	T/T _t	a/a _t	q/p _t	A/A*	v/v _{max}	V/a
0	1.0000	1.0000	1.0000	1.0000	0	0	0	0
.01	.9999	.9999	1.0000	1.0000	.0001	.0178	.0058	.0115
.02	.9997	.9998	.9999	.9999	.0003	.0355	.0115	.0231
.03	.9993	.9996	.9997	.9999	.0007	.0533	.0173	.0346
.04	.9987	.9992	.9995	.9997	.0013	.0710	.0231	.0462
.05	•9979	•9988	•9992	•9996	.0021	.0887	.0289	.0577
.06	•9970	•9982	•9988	•9994	.0030	.1064	.0346	.0692
.07	•9959	•9976	•9984	•9992	.0041	.1240	.0404	.0808
.08	•9947	•9968	•9979	•9989	.0053	.1416	.0461	.0923
.09	•9933	•9960	•9973	•9987	.0067	.1591	.0519	.1038
.10	.9917	.9950	•9967	•9983	.0083	.1766	.0576	.1153
.11	.9900	.9940	•9960	•9980	.0100	.1940	.0634	.1268
.12	.9881	.9928	•9952	•9976	.0119	.2113	.0691	.1382
.13	.9861	.9916	•9944	•9972	.0139	.2285	.0748	.1497
.14	.9839	.9903	•9935	•9967	.0161	.2457	.0806	.1611
.15	.9815	•9889	.9926	•9963	.0184	.2627	.0863	.1726
.16	.9790	•9873	.9915	•9958	.0209	.2796	.0920	.1840
.17	.9763	•9857	.9905	•9952	.0235	.2965	.0977	.1954
.18	.9735	•9840	.9893	•9946	.0263	.3132	.1033	.2067
.19	.9705	•9822	.9881	•9940	.0292	.3298	.1090	.2181
.20	•9674	.9803	.9868	•9934	.0322	.3463	.1147	.2294
.21	•9642	.9783	.9855	•9927	.0354	.3626	.1204	.2407
.22	•9608	.9763	.9841	•9920	.0388	.3788	.1260	.2520
.23	•9572	.9741	.9827	•9913	.0422	.3948	.1316	.2633
.24	•9536	.9719	.9812	•9905	.0458	.4108	.1372	.2745
.25	•9498	•9695	.9796	.9897	.0495	.4265	.1429	.2857
.26	•9458	•9671	.9780	.9889	.0533	.4421	.1484	.2969
.27	•9417	•9646	.9763	.9881	.0572	.4575	.1540	.3081
.28	•9379	•9620	.9745	.9872	.0613	.4727	.1596	.3192
.29	•9332	•9594	.9727	.9863	.0654	.4878	.1651	.3303
.30	.9288	.9566	•9709	.9853	.0697	.5027	.1707	.3413
.31	.9242	.9538	•9690	.9844	.0740	.5174	.1762	.3524
.32	.9195	.9509	•9670	.9834	.0785	.5320	.1817	.3634
.33	.9147	.9479	•9650	.9823	.0830	.5463	.1872	.3743

TABLE I	SUBSONIC	ISENTROPIC	FLOW	RELATIONS	FOR	HELIUM	-	Continued

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м	p/p _t	ρ/ρ _t	T/T _t	a/a _t	q/Pt	A/A [*]	v/v _{max}	V/a
0.34	0.9098	0.9449	0.9629	0.9813	0.0876	0.5604	0.1926	0. 38 53
.35	.9048	.9417	.9608	.9802	.0924	.5744	.1981	.3961
.36	.8997	.9385	.9586	.9791	.0972	.5881	.2035	.4070
.37	.8944	.9353	.9564	.9779	.1020	.6016	.2089	.4178
.38	.8891	.9319	.9541	.9768	.1070	.6150	.2143	.4286
•39	.8837	.9285	.9517	.9756	.1120	.6280	.2197	.4393
•40	.8782	.9250	.9494	.9744	.1171	.6409	.2250	.4501
•41	.8726	.9215	.9469	.9731	.1222	.6536	.2303	.4607
•42	.8669	.9179	.9445	.9718	.1274	.6660	.2356	.4713
•43	.8611	.9142	.9419	.9705	.1327	.6783	.2409	.4819
.44	.8553	.9105	.9394	.9692	.1380	.6902	.2462	.4924
.45	.8493	.9067	.9368	.9679	.1433	.7020	.2515	.5029
.46	.8433	.9028	.9341	.9665	.1487	.7136	.2567	.5134
.47	.8373	.8989	.9314	.9651	.1541	.7248	.2619	.5238
.48	.8311	.8949	.9287	.9637	.1596	.7359	.2671	.5341
•49	.8249	.8909	.9259	.9622	.1650	.7468	.2722	.5441
•50	.8186	.8869	.9231	.9608	.1705	.7574	.2774	.5547
•51	.8123	.8827	.9202	.9593	.1761	.7678	.2825	.5649
•52	.8059	.8786	.9173	.9578	.1816	.7779	.2876	.5751
•53	.7995	.8786	.9144	.9562	.1871	.7878	.2926	.5852
•54	. 7930	.8701	.9114	•9547	.1927	.7975	.2976	•5953
•55	. 7865	.8658	.9084	•9531	.1983	.8069	.3026	•6053
•56	. 7799	.8615	.9054	•9515	.2038	.8161	.3076	•6153
•57	. 7733	.8571	.9023	•9499	.2094	.8250	.3126	•6252
•58	. 7667	.8526	.8992	•9482	.2149	.8337	.3175	•6350
.59	.7600	.8482	.8960	.9466	.2205	.8422	.3224	.6449
.60	•7533	.8437	.8929	.9449	.2260	.8503	.3273	.6546
.61	•7465	.8391	.8897	.9432	.2315	.8583	.3322	.6644
.62	•7398	.8346	.8864	.9415	.2370	.8660	.3370	.6740
.63	•7330	.8300	.8832	.9398	.2424	.8736	.3418	.6837
.64	.7262	.8253	.8799	•9380	•2479	.8809	.3466	.6932
.65	.7194	.8207	.8766	•9362	•2533	.8878	.3513	.7027
.66	.7125	.8160	.8732	•9345	•2586	.8947	.3561	.7122
.67	.7057	.8113	.8698	•9327	•2640	.9012	.3608	.7216

М	p/p _t	ρ/ρ _t	T/T _t	a/a _t	q∕p _t	A/A*	v/v _{max}	V/a
0.68	0.6988	0.8065	0.8665	0.9308	0.2693	0.9075	0.3654	0.7309
.69	.6919	.8018	.8630	.9290	.2745	•9136	.3701	.7402
.70	.6851	.7970	.8596	.9271	.2797	•9195	.3747	.7494
.71	.6782	.7922	.8561	.9253	.2849	•9252	.3793	.7586
.72	.6713	.7873	.8527	.9234	.2900	•9306	.3839	.7677
•73	.6645	.7825	.8492	.9215	.2951	.9358	.3884	.7768
•74	.6576	.7776	.8456	.9196	.3001	.9408	.3929	.7858
•75	.6508	.7728	.8421	.9177	.3051	.9455	.3974	.7948
•76	.6439	.7679	.8386	.9157	.3099	.9501	.4018	.8036
•77	.6371	.7630	.8350	.9138	.3148	.9544	.4062	.8125
.78	.6303	.7581	.8314	.9118	.3195	.9585	.4106	.8212
.79	.6235	.7532	.8278	.9098	.3243	.9623	.4150	.8299
.80	.6167	.7482	.8242	.9078	.3289	.9661	.4193	.8386
.81	.6099	.7433	.8205	.9058	.3334	.9696	.4236	.8472
.82	.6032	.7383	.8169	.9038	.3380	.9728	.4279	.8558
.83	.5964	•7334	.8133	.9018	.3424	•9759	.4321	.8643
.84	.5897	•7284	.8096	.8998	.3467	•9787	.4364	.8728
.85	.5831	•7235	.8059	.8977	.3511	•9815	.4405	.8811
.86	.5764	•7185	.8022	.8957	.3552	•9840	.4447	.8895
.87	.5698	•7136	.7985	.8956	.3594	•9863	.4489	.8977
.88	.5632	. 7086	. 7948	.8915	.3634	.9884	.4529	.9059
.89	.5567	. 7037	. 7911	.8894	.3675	.9903	.4570	.9140
.90	.5502	. 6987	. 7874	.8874	.3714	.9920	.4611	.9222
.91	.5437	. 6938	. 7837	.8853	.3752	.9936	.4651	.9303
.92	.5372	. 6888	. 7800	.8853	.3789	.9949	.4691	.9381
•93	.5308	.6839	.7762	.8810	.3826	.9961	.4730	.9461
•94	.5245	.6789	.7725	.8789	.3862	.9972	.4770	.9540
•95	.5181	.6740	.7687	.8768	.3896	.9981	.4809	.9618
•96	.5119	.6691	.7650	.8746	.3931	.9987	.4848	.9695
•97	.5056	.6642	.7612	.8725	.3964	.9987	.4886	.9773
.98	.4994	.6593	•7575	.8703	•3997	•9997	.4924	.9848
.99	.4933	.6544	•7537	.8682	•4029	•9999	.4962	.9925
1.00	.4871	.6495	•7500	.8660	•4059	1.0000	.5000	1.0000

TABLE I.- SUBSONIC ISENTROPIC FLOW RELATIONS FOR HELIUM - Concluded





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Figure 1.- Continued.

(b) $M_1 = 1$ to μ ; $\varepsilon = 12^0$ to 52^0



(c)
$$M_1 = 1$$
 to μ ; $\epsilon = 52^\circ$ to 90°





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Figure 1. - Continued.



Figure 1. - Continued.



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Figure 1. - Continued.



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Figure 1. - Continued.



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Figure 1. - Continued.



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i)
$$M_1 = 8$$
 to 12; $\epsilon = 60^{\circ}$ to 90°

Figure 1. - Continued.



Figure 1. - Continued.



(k)
$$M_1 = 12$$
 to 20; $\epsilon = 0^{\circ}$ to 60° .

Figure 1.- Continued.





Figure 1. - Continued.





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(m) $M_{\rm l}$ = 20 to 40; ϵ = 1⁰ to 10⁰.



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Figure 1.- Continued.





Figure 1.- Continued.



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Figure 1.- Continued.



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$$M_{\gamma} = 40$$
 to 100; $\epsilon = 0^{\circ}$ to 60° .

Figure 1. - Continued.



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Figure 1. - Concluded.





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Figure 2.- Continued.





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$$M_{\rm l} = 4$$
 to 10; $M_{\rm c} = 0$ to ¹



Figure 2.- Continued.

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(g) $M_1 = 10 \text{ to } 25; M_c = 4 \text{ to}$

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(h) $M_1 = 10 \text{ to } 25; M_c = 8 \text{ to } 18.$



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Figure 2.- Continued.

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$$M_1 = 25$$
 to 40; $M_c = 8$ to 1





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(n) $M_1 = 40$ to 100; $M_c = 0$ to 10.







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Figure 2.- Continued.

(p) $M_{\rm l}$ = 40 to 100; $M_{\rm c}$ = 30 to 70.



Figure 2.- Concluded.



Figure 3.- Variation of ratio of total pressure behind cone shock to free-stream total pressure with free-stream Mach number. (Dashed line denotes limit of weak shock solution.)



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Figure 3.- Continued.



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Figure 3.- Continued.

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1.10 L

14 5 0 5 0

6.5

7.5

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P_{t,1}

Figure 3.- Continued.



(j)
$$M_1 = 10 \text{ to } 25; \frac{P_{t,c}}{P_{t,1}} = 0.1 \text{ to } 0.5.$$

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Figure 3.- Continued.



Figure 3.- Continued.

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Figure 3.- Continued.





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(**p**)
$$M_{l} = 25$$
 to 40; $\frac{P_{t,c}}{P_{t,l}} = 0.1$ to 0.3.

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Figure 3.- Continued.

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Figure 3.- Continued.







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Figure 4.- Continued.









Figure 4.- Continued.








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(g)
$$M_1 = 10 \text{ to } 25; p_c/p_1 = 10 \text{ to 100.}$$





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Figure 4.- Continued.





Figure 4.- Continued.

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Figure 4.- Continued.



(m) $M_{\rm l}$ = 40 to 100; $p_{\rm c}/p_{\rm l}$ = 1 to 10.

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Figure 4.- Continued.







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Figure 5.- Continued.



Figure 5.- Continued.

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(d) $M_1 = \frac{1}{4}$ to 10; $\rho_1 / \rho_c = 0.2$ to 0.6.









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Figure 5.- Continued.

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Figure 5.- Continued.



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Figure 5.- Continued

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flow-deflection angles. (Dashed line denotes boundary between weak and strong shock solu-Figure 6. - Variation of two-dimensional shock angle with free-stream Mach number for various tions. Weak shock solutions are below dashed line.)

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Figure 6.- Continued.

(b) $M_{1} = 1$ to μ ; $\epsilon = 12^{\circ}$ to 52°



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Figure 6.- Continued.



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Figure 6. - Continued.



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(g) $M_1 = 8$ to 12; $\epsilon = 1^{\circ}$ to 10° .

Figure 6. - Continued.





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(i)
$$M_1 = 8$$
 to 12; $\epsilon = 60^\circ$ to 90° .

Figure 6. - Continued.



Figure 6.- Continued.



Figure 6.- Continued.



(1)
$$M_1 = 12$$
 to 20; $\epsilon = 60^{\circ}$ to 90° .

Figure 6.- Continued.



m)
$$M_1 = 20 \text{ to } 40; \epsilon = 1^0 \text{ to } 10^1$$

Figure 6.- Continued.


(n)
$$M_1 = 20$$
 to 40; $\epsilon = 0^0$ to 60^0 .

Figure 6. - Continued.



Figure 6.- Continued.









Figure 6.- Continued.





Figure 6. - Concluded.



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(a) $M_1 = 1$ to 100; $M_2 = 0$ to 76.

Figure 7.- Variation of Mach number behind two-dimensional shock with free-stream Mach number (Dashed line denotes boundary between weak and strong shock solutions. Weak shock solution is above dashed line.) for various flow-deflection angles.



Figure 7.- Continued.

(b) $M_1 = 1 \text{ to } \mu; M_2 = 0.1 \text{ to}$

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Figure 7.- Continued.



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Figure 7.- Continued.



Figure 7.- Continued.

(f) $M_1 = 4$ to 10; $M_2 = 4.8$ to 6.8.





(g) $M_1 = 4$ to 10; $M_2 = 6.8$ to 10.8.





Figure 7.- Continued.

(h) $M_1 = 10 \text{ to } 25; M_2 = 0.1 \text{ to } 1.$



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Figure 7.- Continued.









Figure 7.- Continued.











$$M_1 = 25$$
 to 40; $M_2 = 0.1$ to





(n) $M_1 = 25$ to 40; $M_2 = 0$ to 4.



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 $M_1 = 25$ to 40; $M_2 = 4$ to 8.



Figure 7.- Continued.



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Figure 7.- Continued.



Figure 7.- Continued.









line denotes boundary between weak and strong shock solutions. Weak shock solution is above dashed line.) Figure 8.- Variation of ratio of total pressure behind a two-dimensional shock to free-stream (Dashed total pressure with free-stream Mach number for various flow-deflection angles.

(a)
$$M_1 = 1$$
 to 100; $\frac{P_{t,2}}{P_{t,1}} = 0$ to 1.







Figure 8.- Continued.

























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Figure 8.- Continued.









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Figure 8.- Continued.

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Figure 8.- Continued.



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Figure 8.- Continued.



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Figure 8.- Continued.





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Figure 8.- Concluded.







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(b)
$$M_1 = 1$$
 to μ ; $p_2/p_1 = 1$ to 10.





Figure 9.- Continued.



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Figure 9.- Continued.









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Figure 9.- Continued.



Figure 9.- Continued.

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line denotes boundary between weak and strong shock solutions. Weak shock solution is above (Dashed Figure 10.- Variation of ratio of free-stream static density to static density behind a twodimensional shock with free-stream Mach number for various flow-deflection angles. dashed line.)















Figure 10.- Continued.



Figure 10. - Continued.




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 $\frac{\sigma}{\rho} \frac{\sigma}{\rho}$

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