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Technical Report

JUNCTURE STRESS FIELDS
IN MULTICELLULAR SHELL STRUCTURES

Final Report
Nine Volumes

Vol. VII Buckling Analysis of Segmental Orthotropic
Cylinders Under Uniform Stress Distribution

by

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Contract NAS 8-11480 to National Aeronautics and Space Administration
George C. Marshall Space Flight Center, Huntsville, Alabama

FOREWORD

This report is the result of a study on the buckling of segmental orthotropic cylinders under uniform axial compression. Work on this study was performed by staff members of Lockheed Missiles & Space Company in cooperation with the George C. Marshall Space Flight Center of the National Aeronautics and Space Administration under Contract NAS 8-11480. Contract technical representative was H. Coldwater.

This volume is the the seventh of a nine-volume final report of studies conducted by the department of Solid Mechanics, Aerospace Sciences Laboratory, Lockheed Missiles & Space Company. Project Manager was K. J. Forsberg; E. Y. W. Tsui was Technical Director and Professor D. O. Brush of the University of California at Davis was a Consultant for the work.

The nine volumes of the final report have the following titles:

- Vol. I Numerical Methods of Solving Large Matrices
- Vol. II Stresses and Deformations of Fixed-Edge Segmental Cylindrical Shells
- Vol. III Stresses and Deformations of Fixed-Edge Segmental Conical Shells
- Vol. IV Stresses and Deformations of Fixed-Edge Segmental Spherical Shells
- Vol. V Influence Coefficients of Segmental Shells
- Vol. VI Analysis of Multicellular Propellant Pressure Vessels by the Stiffness Method
- Vol. VII Buckling Analysis of Segmental Orthotropic Cylinders under Uniform Stress Distribution
- Vol. VIII Buckling Analysis of Segmental Orthotropic Cylinders under Non-uniform Stress Distribution
- Vol. IX Summary of Results and Recommendations

SUMMARY

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This volume presents a theoretical solution to the buckling problem of orthotropic cylindrical panels under uniform axial compression. The panels are simply supported along the curved edges and elastically supported along the straight edges. There is included a digital computer program that was used to obtain design curves which give non-dimensional buckling loads for a wide range of geometric parameters for the cylindrical panels. Derivation of the expressions for the bending influence coefficients of orthotropic rectangular plates is also included.

Author

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NOTATION

$A_1 = d_1 h_1; A_2 = d_2 h_2$	cross-sectional areas of stringer and ring
$D = Et^3/12(1 - \nu^2)$	flexural rigidity of shell
$\hat{e}_1, \hat{e}_2$	eccentricities of stringer and ring
E, E_1, E_2	moduli of elasticity of the shell, stringer, and ring
$G = E/2(1 + \nu)$	shear modulus of shell
I_1, I_2	moments of inertia about the centroidal axis of stringer and shell; ring and skin
$\bar{I}_1, \bar{I}_2$	moments of inertia about the neutral axis of stringer and ring
J_1, J_2	torsional constants for the cross section of stringer and ring
K	buckling coefficient
L	length of panel
$\hat{M}(\), \hat{N}(\)$	moments and stress resultants
$\hat{N}_{cr}$	critical buckling load
n	number of half-waves in the axial direction
P_x	total axial buckling load of the panel
R	radius of cylinder
r, s	spacings of rings and stringers
t	thickness of shell
$\hat{U}$	strain energy
$\hat{u}, \hat{v}, \hat{w}$	middle-surface displacement components in $\hat{x}, \hat{y}$, and $\hat{z}$ directions

$\hat{x}, \hat{y}, \hat{z}$	curvilinear coordinates for the middle-surface of shell
$\hat{V}$	total potential energy
$\hat{W}$	potential energy of external load
ν, ν_1, ν_2	Poisson's ratios of shell, stringer and ring
$\epsilon(\), \gamma(\)$	direct and shear strains
$\chi(\)$	change of curvature or twist of middle-surface of shell
$(\),_x$	$\partial(\)/\partial x$
Ω	angle subtending the width of panel (see Fig. 1)
$\theta = \hat{y}/R$	angular coordinate

$$\rho = 1 + (1 - \nu^2) \left(\frac{E_1}{E} \right) \left(\frac{d_1}{s} \right) \left(\frac{h_1}{t} \right) \left[\left(\frac{h_1}{t} \right)^2 + \frac{3(h_1/t + 1)^2}{(1 - \nu^2)(E_1/E)(d_1/s)(h_1/t) + 1} \right]$$

$$\gamma^2 = \frac{1}{12} \left(\frac{t}{R} \right)^2 \quad \text{constant}$$

Section 1 INTRODUCTION

When multiceullular configurations are employed as fuel tanks or interstages of space vehicles, buckling of the cylindrical panels under axial compression may be critical as far as the load-carrying capacity of the structure is concerned. In general, the loads on such panels are distributed in a nonuniform manner. However, the critical uniform loading may serve as a valuable guide in the design of multicellular shell structures. Besides local buckling (between stiffeners) two different buckling modes are possible: a panel buckling mode and a general instability mode. We define as panel buckling here a mode in which the displacement pattern is either symmetrical or antimetrical about each of the lines of intersection between adjacent panels. In the general instability mode, on the other hand, the wavelength is larger than the width of one panel. For structures having twenty cells or less, it is evident that this deformation pattern produces very large axial membrane stresses in the cylindrical panels. Therefore, general instability will not be critical and because satisfactory methods are available for the local buckling mode, only the panel buckling mode will be considered here.

It is believed that the buckling of monocoque simply supported cylindrical panels was first studied by S. Timoshenko (Ref. 1). S. C. Redshaw (Ref. 2) extended the analysis to include a few cases of different boundary conditions along the generators. Results for the simply supported case agreed fairly well with tests made by J. S. Newell and W. H. Gale (Ref. 3). However, for the cases of fixed straight edges, the correlation was poor. A general expression for the critical compressive stress for slightly curved sheets with equal elastic restraints against rotation along the unloaded edges was obtained by E. Z. Stowell (Ref. 4).

Using a simplified version of the buckling determinant derived by W. Flügge (Ref. 5) for orthotropic cylinders, D. D. Dschou (Ref. 6) investigated the buckling of axially

compressed segmental stiffened cylindrical shells. The general expression for the critical stress, however, is for the simple support boundary condition only, and eccentricity effects were not considered. Since then, the problem has received little attention (Refs. 7 and 8).

It is apparent that the available solutions cannot be directly applied to determine the critical compressive stresses of cylindrical segments subjected to the boundary conditions along the generators, which are applicable to the cylindrical panels in multicellular structures. The purpose of the present work is, therefore, to perform the necessary theoretical investigations required to formulate theory and to develop an analytical approach for predicting the axial buckling loads for both isotropic and orthotropic segmental cylindrical shells with displacements satisfying appropriate boundary conditions along the straight edges. Effects of eccentricities of stiffeners and bending of the web plates will also be examined.

Section 2

THEORETICAL FORMULATION

2.1 BASIC EQUATIONS

The analysis is based on the following assumptions:

- Stiffeners are located along coordinate lines.
- Strains in the stiffeners vary linearly with the radial coordinate and are compatible with those in the skin at the surfaces of contact.
- Stiffeners are under a uniaxial state of stress.
- In-plane shear forces are entirely carried by the skin.
- Stiffeners are so closely spaced that the effective width of the skin equals the spacing of the stringers or rings.

The strain-energy expression given in Sec. 2.2 for an orthotropic segmental cylinder as shown in Fig. 1 is based on the following elastic and kinematic relations (Refs. 9 and 10):

- Elastic relations

For the skin:

$$\sigma_x = \frac{E}{1 - \nu^2} [\epsilon_x + \nu\epsilon_\theta + z(\chi_x + \nu\chi_\theta)] \quad (2.1a)$$

$$\sigma_\theta = \frac{E}{1 - \nu^2} [\epsilon_\theta + \nu\epsilon_x + z(\chi_\theta + \nu\chi_x)] \quad (2.1b)$$

For the stiffeners:

$$\sigma_x = E_1(\epsilon_x + z\chi_x) \quad (2.2a)$$

$$\sigma_\theta = E_2(\epsilon_\theta + z\chi_\theta) \quad (2.2b)$$

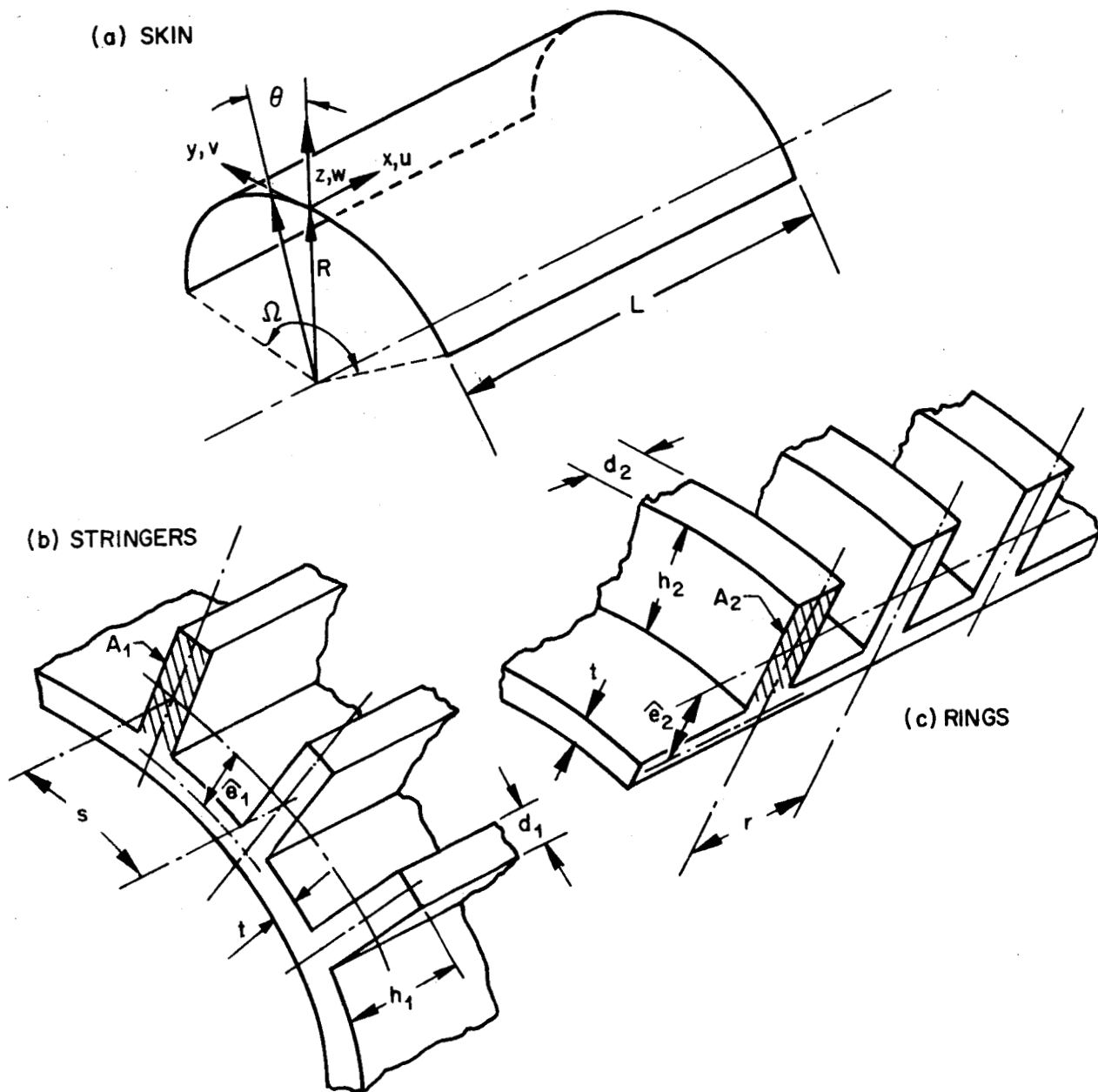


Fig. 1 Geometry of an Orthotropic Cylindrical Panel

- Middle-surface kinematic relations

Strain displacement:

$$\epsilon_x = u_{,x} + (w_{,x})^2/2 \quad (2.3a)$$

$$\epsilon_\theta = v_{,\theta} + w + (w_{,\theta})^2/2 \quad (2.3b)$$

$$\gamma_{x\theta} = u_{,\theta} + v_{,x} + (w_{,x})(w_{,\theta}) \quad (2.3c)$$

Curvature displacement:

$$\chi_x = -w_{,xx} \quad (2.4a)$$

$$\chi_\theta = -w_{,\theta\theta} \quad (2.4b)$$

$$\chi_{x\theta} = -w_{,x\theta} \quad (2.4c)$$

The nondimensional variables are given by

$$x = \hat{x}/R, \quad \theta = \hat{y}/R, \quad z = \hat{z}/R \quad (2.5a)$$

$$u = \hat{u}/R, \quad v = \hat{v}/R, \quad w = \hat{w}/R \quad (2.5b)$$

In terms of these relationships the corresponding stress resultants and moments are as follows:

$$N_x = \alpha_3 \epsilon_x + \nu \epsilon_\theta + \alpha_1 \chi_x \quad (2.6a)$$

$$N_\theta = \alpha_4 \epsilon_\theta + \nu \epsilon_x + \alpha_2 \chi_\theta \quad (2.6b)$$

$$N_{x\theta} = N_{\theta x} = \frac{1 - \nu}{2} \gamma_{x\theta} \quad (2.6c)$$

$$M_x = \alpha_5 \chi_x + \nu \gamma^2 \chi_\theta + \alpha_1 \epsilon_x \quad (2.6d)$$

$$M_\theta = \alpha_6 \chi_\theta + \nu \gamma^2 \chi_x + \alpha_2 \epsilon_\theta \quad (2.6e)$$

$$M_{x\theta} = \left[\gamma^2 + \frac{1 + \nu}{2tR^2} \frac{J_1}{(1 + \nu_1)s} \left(\frac{E_1}{E} \right) \right] (1 - \nu) \chi_{x\theta} \quad (2.6f)$$

$$M_{\theta x} = \left[\gamma^2 + \frac{1 + \nu}{2tR^2} \frac{J_2}{(1 + \nu_2)r} \left(\frac{E_2}{E} \right) \right] (1 - \nu) \chi_{\theta x} \quad (2.6g)$$

The nondimensional stress resultants and moments are defined as follows:

$$N_{()} = \frac{1 - \nu^2}{Et} \hat{N}_{()} \quad (2.7a)$$

$$M_{()} = \frac{1 - \nu^2}{EtF} \hat{M}_{()} \quad (2.7b)$$

and

$$\alpha_1 = \left(\frac{\hat{e}_1}{R} \right) \left(\frac{E_1}{E} \right) \frac{A_1 (1 - \nu^2)}{st} \quad (2.8a)$$

$$\alpha_2 = \left(\frac{\hat{e}_2}{R} \right) \left(\frac{E_2}{E} \right) \frac{A_2 (1 - \nu^2)}{rt} \quad (2.8b)$$

$$\alpha_3 = 1 + \left(\frac{E_1}{E}\right) \frac{A_1(1 - \nu^2)}{st} \quad (2.8c)$$

$$\alpha_4 = 1 + \left(\frac{E_2}{E}\right) \frac{A_2(1 - \nu^2)}{rt} \quad (2.8d)$$

$$\alpha_5 = \gamma^2 + (1 - \nu^2) \left(\frac{E_1}{E}\right) \frac{\bar{I}_1 + A_1 \hat{e}_1^2}{stR^2} \quad (2.8e)$$

$$\alpha_6 = \gamma^2 + (1 - \nu^2) \left(\frac{E_2}{E}\right) \frac{\bar{I}_2 + A_2 \hat{e}_2^2}{rtR^2} \quad (2.8f)$$

$$\alpha_7 = 2\gamma^2 + \frac{1 - \nu^2}{2tR^2} \left[\frac{J_1}{(1 + \nu_1)s} \left(\frac{E_1}{E}\right) + \frac{J_2}{(1 + \nu_2)r} \left(\frac{E_2}{E}\right) \right] \quad (2.8g)$$

$$\gamma^2 = \frac{1}{12} \left(\frac{t}{R}\right)^2 \quad (2.8h)$$

2.2 STRAIN-ENERGY EXPRESSIONS

The total potential energy of a stiffened cylindrical panel under axial load may be expressed as

$$\hat{V} = \hat{U} + \hat{W} \quad (2.9)$$

where $\hat{U}$ is the strain energy of the skin, stringers, and rings and $\hat{W}$ is the potential energy of the external load. From the relations in Sec. 2.1, it can be shown the strain energy of a cylindrical panel (Fig. 1) is given by the following equation.

$$\begin{aligned}
U = & \int_0^{L/R} \int_{-\Omega/2}^{\Omega/2} \left(\epsilon_x^2 + \epsilon_\theta^2 + 2\nu \epsilon_x \epsilon_\theta + \frac{1-\nu}{2} \gamma_{x\theta}^2 \right) dx d\theta \\
& + \frac{1}{12} \left(\frac{t}{r} \right)^2 \left[\chi_x^2 + \chi_\theta^2 + 2\nu \chi_x \chi_\theta + 2(1-\nu) \chi_{x\theta}^2 \right] dx d\theta \\
& + \frac{1-\nu^2}{st} \frac{E_1}{E} \left\{ A_1 \left[\epsilon_x^2 + \frac{2\hat{e}_1}{R} \epsilon_x \epsilon_x + \left(\frac{\hat{e}_1}{R} \right)^2 \chi_x^2 \right] + \bar{I}_1 \chi_x^2 + \frac{J_1}{2(1+\nu_1)} \chi_{x0}^2 \right\} dx d\theta \\
& + \frac{1-\nu^2}{rt} \frac{E_2}{E} \left\{ A_2 \left[\epsilon_\theta^2 + \frac{2\hat{e}_2}{R} \epsilon_\theta \epsilon_\theta + \left(\frac{\hat{e}_2}{R} \right)^2 \chi_\theta^2 \right] + \bar{I}_2 \chi_\theta^2 + \frac{J_2}{2(1+\nu_2)} \chi_{x0}^2 \right\} dx d\theta \quad (2.10)
\end{aligned}$$

where

$$U = \frac{2(1-\nu^2)}{EtR^2} \hat{U}$$

The potential energy of the external load is

$$\hat{W} = -\frac{\hat{N}_x}{2} \int_0^{L/R} \int_{-\Omega/2}^{\Omega/2} u_{,x} dx d\theta \quad (2.11)$$

Equation (2.10) becomes in terms of the nondimensional parameter α_j

$$\begin{aligned}
U = & \int_0^{L/R} \int_{-\Omega/2}^{\Omega/2} \left[\alpha_3 \epsilon_x^2 + \alpha_4 \epsilon_\theta^2 + 2\alpha_1 \epsilon_x \chi_x + 2\alpha_2 \epsilon_\theta \chi_\theta + 2\nu \epsilon_x \epsilon_\theta + \frac{1-\nu}{2} \gamma_{x\theta}^2 \right. \\
& \left. + \alpha_5 \chi_x^2 + \alpha_6 \chi_\theta^2 + 2\nu \chi_x \chi_\theta + (\alpha_7 - 2\nu \gamma^2) \chi_{x0}^2 \right] dx d\theta \quad (2.12)
\end{aligned}$$

2.3 BUCKLING EQUATIONS

The equations for the state of neutral equilibrium and the associated boundary conditions can be obtained by equating to zero the first variation of the potential energy [Eq. (2.10)] with respect to each of the displacement components u, v, w (i.e., $\delta\hat{V} = 0$) and by linearizing the resulting equations in accordance with the adjacent-equilibrium theory of stability. The linearized equations can then be expressed in terms of pre-buckling stress resultants, incremental displacement components, and other parameters. Suitable mathematical manipulation and substitution lead to the following stability equations for axial compression:

$$\alpha_3 u_{,xx} + \frac{1-p}{2} u_{,\theta\theta} + \frac{1+p}{2} v_{,x\theta} + p w_{,x} - \alpha_1 w_{,xxx} = 0 \quad (2.13a)$$

$$\frac{1+p}{2} u_{,x\theta} + \frac{1-p}{2} v_{,xx} + \alpha_4 v_{,\theta\theta} + \alpha_4 w_{,\theta} - \alpha_2 w_{,\theta\theta\theta} = 0 \quad (2.13b)$$

$$N_x w_{,xx} + p u_{,x} - \alpha_4 (v_{,\theta} + w) - \alpha_2 w_{,\theta\theta} - \alpha_1 u_{,xxx} - \alpha_2 (v_{,\theta\theta\theta} + w_{,\theta\theta}) + \alpha_5 w_{,xxxx} + \alpha_6 w_{,\theta\theta\theta\theta} + \alpha_7 w_{,xx\theta\theta} = 0 \quad (2.13c)$$

Specializing these equations for a homogeneous isotropic shell yields the well-known Donnell equations.

2.4 BOUNDARY CONDITIONS

The natural or dynamic boundary conditions of the problem under consideration can be obtained from the expression for the first variation of the strain energy. These are as follows:

- At the curved edges

$$\alpha_3 u_{,x} + p(v_{,\theta} + w) - \alpha_1 w_{,xx} = 0 \quad (2.14a)$$

$$u_{,\theta} + v_{,x} = 0 \quad (2.14b)$$

$$\alpha_5 w_{,xxx} + \nu \gamma^2 w_{,x\theta\theta} - \alpha_1 u_{,xx} + (\alpha_7 - 2\nu\gamma^2) w_{,x\theta} = 0 \quad (2.14c)$$

$$\alpha_5 w_{,xx} + \nu \gamma^2 w_{,\theta\theta} - \alpha_1 u_{,x} = 0 \quad (2.14d)$$

● At the straight edges

$$u_{,\theta} + v_{,x} = 0 \quad (2.15a)$$

$$\alpha_4 (v_{,\theta} + w) + \nu u_{,x} - \alpha_2 w_{,\theta\theta} = 0 \quad (2.15b)$$

$$\alpha_6 w_{,000} + \nu \gamma^2 w_{,xx} - \alpha_2 (v_{,\theta\theta} + w_{,\theta}) + (\alpha_7 - 2\nu\gamma^2) w_{,xx\theta} = 0 \quad (2.15c)$$

$$\alpha_6 w_{,\theta\theta} + \nu \gamma^2 w_{,xx} - \alpha_2 (v_{,\theta} + w) = 0 \quad (2.15d)$$

For simplicity in the analysis, it is assumed here that the panels are simply supported at the curved edges.

$$\alpha_3 u_{,x} + \nu (v_{,\theta} + w) - \alpha_1 w_{,xx} = N_x = 0 \quad (2.16a)$$

$$v = 0 \quad (2.16b)$$

$$w = 0 \quad (2.16c)$$

$$\alpha_5 w_{,xx} + \nu \gamma^2 w_{,\theta\theta} - \alpha_1 u_{,x} = M_x = 0 \quad (2.16d)$$

Since the radial webs in multicellular shell structure elastically restrain the deformations along the line of intersection between two cylindrical panels, boundary conditions for both antimetric and symmetric buckling can be expressed in terms of the influence coefficients of the web plates. The displacement components parallel and perpendicular to the plane of the web plates (Fig. 2) are respectively given by

$$\bar{w} = w \cos \phi - v \sin \phi \quad (2.17a)$$

$$\bar{v} = w \sin \phi + v \cos \phi \quad (2.17b)$$

where

$$\phi = \begin{Bmatrix} \frac{\Omega - \mathcal{D}}{2} \\ \frac{-\Omega + \mathcal{D}}{2} \end{Bmatrix} \quad \text{at} \quad \theta = \begin{Bmatrix} \Omega/2 \\ -\Omega/2 \end{Bmatrix}$$

The boundary conditions for the antimetric mode are

$$u = 0 \quad (2.18a)$$

$$\bar{w} = 0 \quad (2.18b)$$

$$\bar{v} - c_{11}(N_\theta \cos \phi + V_\theta \sin \phi)/2 - c_{12}M_\theta/2 = 0 \quad (2.18c)$$

$$w_{,\theta} - c_{21}(N_\theta \cos \phi + V_\theta \sin \phi)/2 - c_{22}M_\theta/2 = 0 \quad (2.18d)$$

In the above equations, N_θ and M_θ are given by Eqs. (2.6b) and (2.6c), and V_θ is the shear stress resultant.

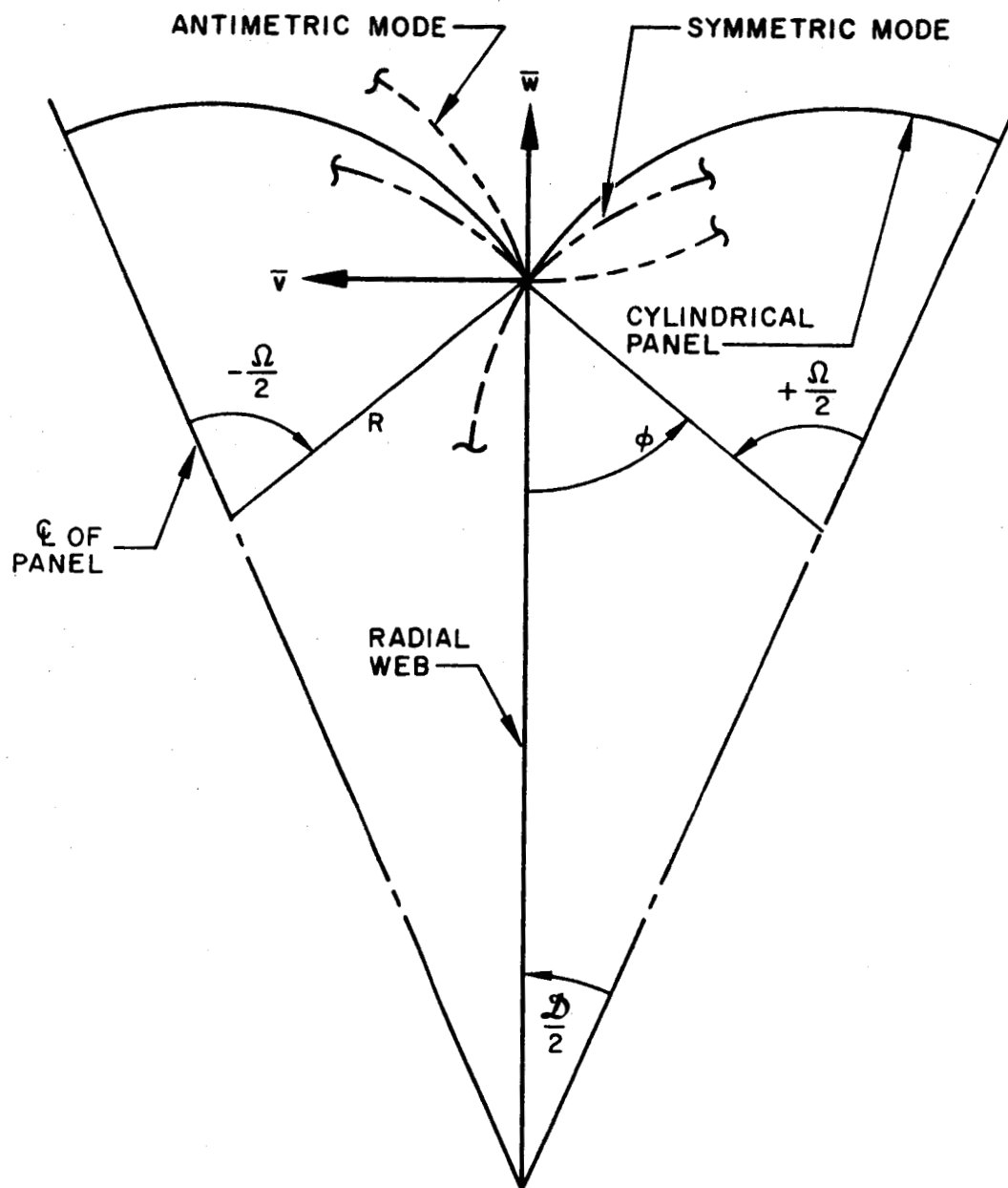


Fig. 2 Buckling Modes of Panel

which may be expressed in terms of the displacement components as follows:

$$V_{\theta} = (\alpha_7 - \nu\gamma^2)w_{,\theta xx} + \alpha_6 w_{,\theta\theta\theta} - \alpha_2(v_{,\theta\theta} + w_{,\theta}) \quad (2.19)$$

The coefficients c_{11} , c_{12} , c_{21} , and c_{22} are the bending influence coefficients of the web plates. A method to determine these coefficients is given in Appendix A.

For the symmetrical mode of panel buckling the boundary conditions at the straight edges are

$$u - c_{33}N_{x\theta}/2 = 0 \quad (2.20a)$$

$$\bar{v} = 0 \quad (2.20b)$$

$$\bar{w} - c_{44}(N_{\theta} \sin \phi + V_{\theta} \cos \phi)/2 = 0 \quad (2.20c)$$

$$w_{,\theta} = 0 \quad (2.20d)$$

where c_{33} and c_{44} are the stretching influence coefficients of the web plates. For the multicellular shell structures under consideration, it has been found numerically that the stretching flexibility of the web plates, especially when they are stiffened, is very small.

2.5 SOLUTION

For the case of uniform axial compression, Eqs. (2.13) are linear differential equations with constant coefficients. The simple support boundary conditions for the curved edges as given by Eqs. (2.16) are satisfied by the following displacement functions:

$$u = \cos anx \sum_{j=1}^8 A_j \exp [\lambda_j(\theta/\Omega)] \quad (2.21a)$$

$$v = \sin anx \sum_{j=1}^8 B_j \exp [\lambda_j(\theta/\Omega)] \quad (2.21b)$$

$$w = \sin anx \sum_{j=1}^8 C_j \exp [\lambda_j(\theta/\Omega)] \quad (2.21c)$$

where

$$a = \pi(R/L)$$

Substitution of the displacement functions into the equations of equilibrium [Eqs. (2.13)] yields the following characteristic equation:

$$\begin{aligned} & \left\{ \left(\frac{1-\nu}{2} \right) \left[(\alpha_1)^2 - \alpha_3 \alpha_5 \right] a^8 \right\} n^8 + \left[(1-\nu) \nu \alpha_1 a^6 \right] n^6 + \left\{ \left(\frac{1-\nu}{2} \right) \left(\nu^2 - \alpha_3 \alpha_4 \right) a^4 \right\} n^4 \\ & + \left\{ \left[\alpha_5 (\alpha_3 \alpha_4 - \nu) + \left(\frac{1-\nu}{2} \right) \alpha_3 \alpha_7 - (\alpha_1)^2 \alpha_4 \right] a^6 n^6 + (1-\nu) \left[\alpha_2 \alpha_3 + \alpha_1 \alpha_4 \right] a^4 n^4 \right\} (\tilde{\lambda})^2 \\ & + \left\{ \nu(1-\nu) \alpha_2 a^2 n^2 + \left[(\nu - \alpha_3 \alpha_4) \alpha_7 - \left(\frac{1-\nu}{2} \right) (\alpha_3 \alpha_6 + \alpha_4 \alpha_5) - (1+\nu) \alpha_1 \alpha_2 \right] a^4 n^4 \right\} (\tilde{\lambda})^4 \\ & + \left\{ \left[\alpha_6 (\alpha_3 \alpha_4 - \nu) + \left(\frac{1-\nu}{2} \right) \alpha_4 \alpha_7 - (\alpha_2)^2 \alpha_3 \right] a^2 n^2 \right\} (\tilde{\lambda})^6 + \left\{ \left(\frac{1-\nu}{2} \right) \left[(\alpha_2)^2 - \alpha_4 \alpha_6 \right] \right\} (\tilde{\lambda})^8 \\ & - N_{x_0} \left\{ \left(\frac{1-\nu}{2} \right) \alpha_3 a^6 n^6 + (\nu - \alpha_3 \alpha_4) a^4 n^4 (\tilde{\lambda})^2 + \left(\frac{1-\nu}{2} \right) \alpha_4 a^2 n^2 (\tilde{\lambda})^4 \right\} = 0 \end{aligned} \quad (2.22)$$

where

$$N_{x_0} = \frac{1 - \nu^2}{Et} \frac{P_x}{\Omega R \left[1 + \frac{\hat{e}_1/R}{1 + (E/E_1)st/(1 - \nu^2)A_1} \right]}$$

and

$$\tilde{\lambda}_j = \frac{1}{\Omega} \lambda_j$$

Equations (2.13) are identically satisfied if λ_j ($j = 1, 2, \dots, 8$) in Eqs. (2.22) are chosen to be the eight roots of this characteristic equation. The A_j , B_j and C_j in Eqs. (2.22) are constants of integration. For each given value of λ_j , A_j and B_j may be expressed in terms of C_j by use of the in-plane equilibrium equations [Eqs. (2.13a,b)]. These constants are governed by the homogeneous equation system which is obtained if Eqs. (2.22) are substituted into the eight boundary conditions for the straight edges (four at each edge). The critical value of the load is found as the lowest value of N_{x_0} for which the determinant of the coefficients of this equation system vanishes.

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Section 3

NUMERICAL ANALYSIS

The critical load of an axially compressed stiffened cylindrical panel may be obtained by use of the computer program which is discussed in Appendix B. This program is based on the theoretical formulation as described in the previous section. It should be noted, particularly in the preliminary design stage, that the inverse problem is frequently posed. That is, the axial load and sometimes material or certain geometrical parameters are given and it is the task of the stress analyst to determine the panel geometry. It is certainly possible that this problem may be solved by the application of the computer program to a number of cases in which one or more parameters have been varied. In many cases it may also be possible to devise an optimizing computer program for the inverse problem. However, because of the many alternatives that are open, the usefulness of such a program will necessarily be somewhat restricted. In view of these considerations, it is obvious that the computer program would be a very valuable tool if a set of design curves were provided. Unfortunately, the large number of independent parameters for stiffened panels makes a complete parametric study impossible. Therefore, the numerical analysis presented here is primarily intended to help the designer to obtain a reasonably close estimate of the critical loads. With such an estimate available, the determination of more accurate dimensions by use of the computer program is expedited.

Table 1 shows some results from a study of the influence of the angle ϕ on the buckling load. It was found that within the range of values applicable to multicellular design, ϕ has little effect on the critical load. Consequently, for the case in which the stiffness of the web plate can be omitted, the usual simple support boundary conditions give a good approximation for the panels in the shell structures under consideration.

Table 1
INFLUENCE OF ANGLE ϕ ON BUCKLING LOAD

ϕ (deg)	Buckling Load $N_{cr} \left(\frac{E}{10^7} \right)$	Remarks
0	11, 170	$R/t = 242$
7.5	11, 265	$L/R = 1.5$
		$d_1/R = d_2/R = 0.006$
		$h_1/R = h_2/R = 0.017$
11.25	11.313	$E_1/E = E_2/E = 1$
		$\nu = 0.3$
		$\Omega = 30 \text{ deg}$

Figure 3 shows critical loads as a function of the panel angle, Ω , for an external stringer stiffened panel. The nondimensional buckling curves are for the following four different boundary conditions along the straight edges:

PS1	$u = 0 ; v = 0 ; w = 0 ; M_{\theta} = 0$
PC1	$u = 0 ; v = 0 ; w = 0 ; w_{,\theta} = 0$
PS2	$u = 0 ; N_{\theta} = 0 ; w = 0 ; M_{\theta} = 0$
PC2	$u = 0 ; N_{\theta} = 0 ; w = 0 ; w_{,\theta} = 0$

In view of Fig. 3, the curve PS2 gives the critical load for the case in which the web plate stiffness is omitted, and curve PC1 corresponds to the case of rigid web plate. It is evident that for sufficiently narrow panels the influence of the web plate stiffness may be significant. It can be observed that to restrain tangential displacements is more effective than to restrain rotations at the edge (PS1-curve is well above the PC2-curve).

The critical load corresponding to a series of values of the web plate stiffness has been computed for a panel which is stiffened by an internal waffle pattern. The results

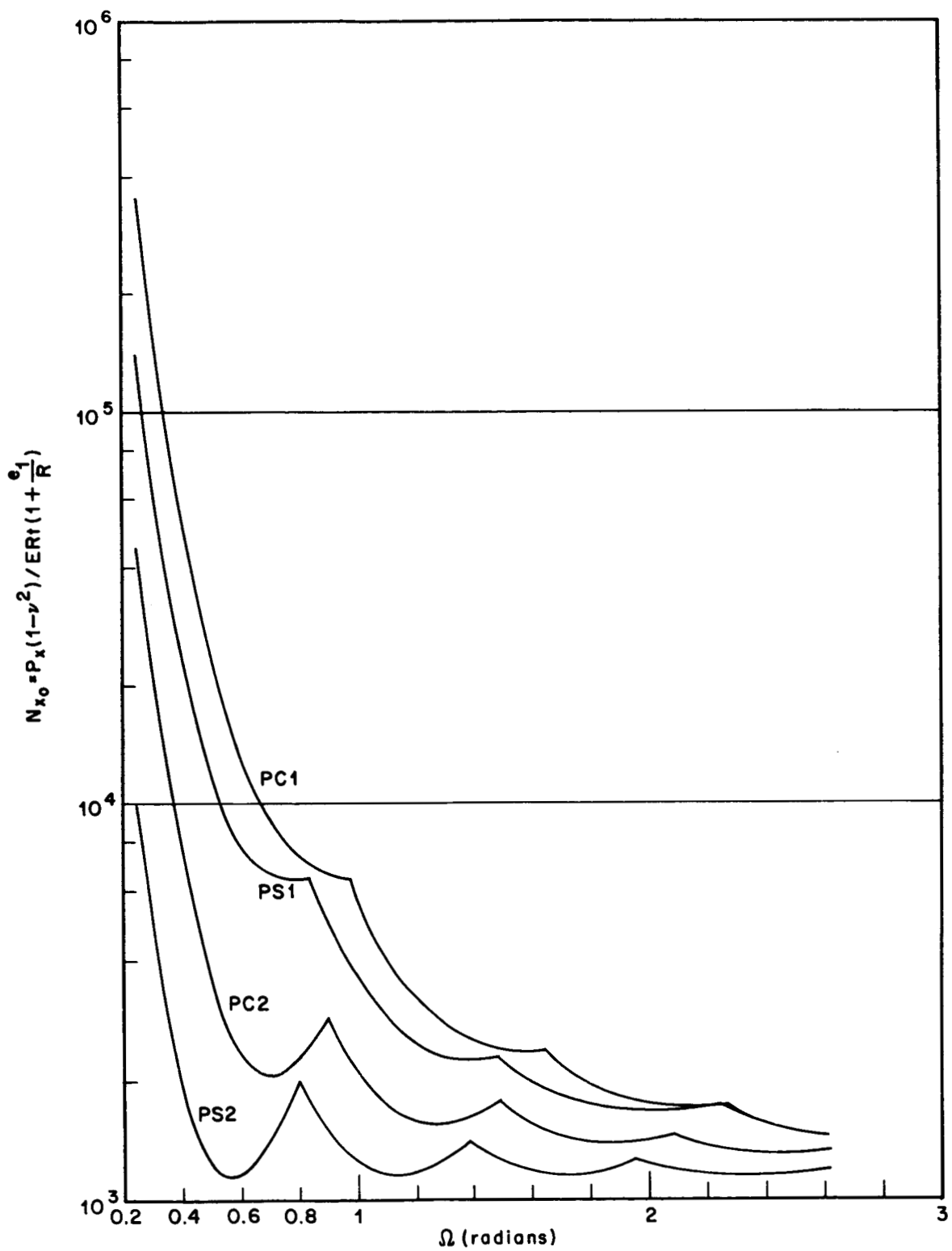


Fig. 3 Effects of Ω on Buckling Stress

are shown in Fig. 4. The web plate here is stiffened with the same waffle pattern as that for the cylinder but the stiffener height was varied. It is seen that the critical load does not increase significantly unless the plate is very stiff in comparison to that of the cylindrical panel. It appears that in most cases the simple support conditions should lead to a good approximation.

If the simply supported panel is entirely or primarily stiffened in the circumferential direction, the critical load would be the same as that obtained from a complete cylinder. For panels which are essentially stringer stiffened or with approximately the same stiffening in both directions, an estimate of the buckling load can be obtained by use of Figs. 5 through 16. For the case of simple support, Figs. 5 through 10 give buckling coefficients as a function of the angle Ω for various values of the parameters ρ , L/R , and R/t . Figures 11 through 16 give corresponding results for waffle stiffened cylindrical panels. Results are presented for external as well as internal stiffening. The graphs apply strictly only for rectangular stiffeners but they will yield reasonable buckling loads for other stiffener shapes.

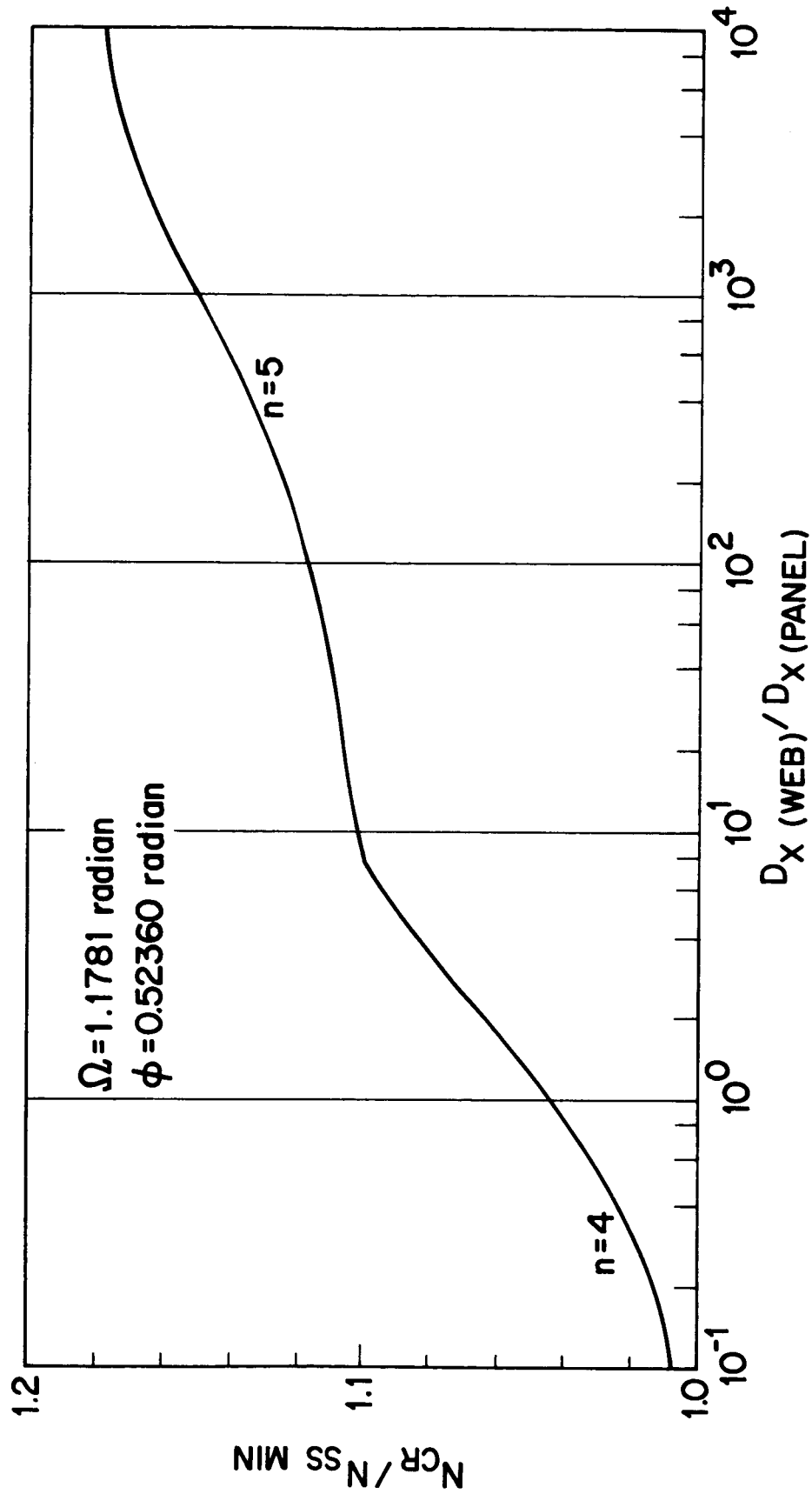


Fig. 4 Effect of Web Plate Stiffness on Panel Buckling (Antisymmetric Mode)

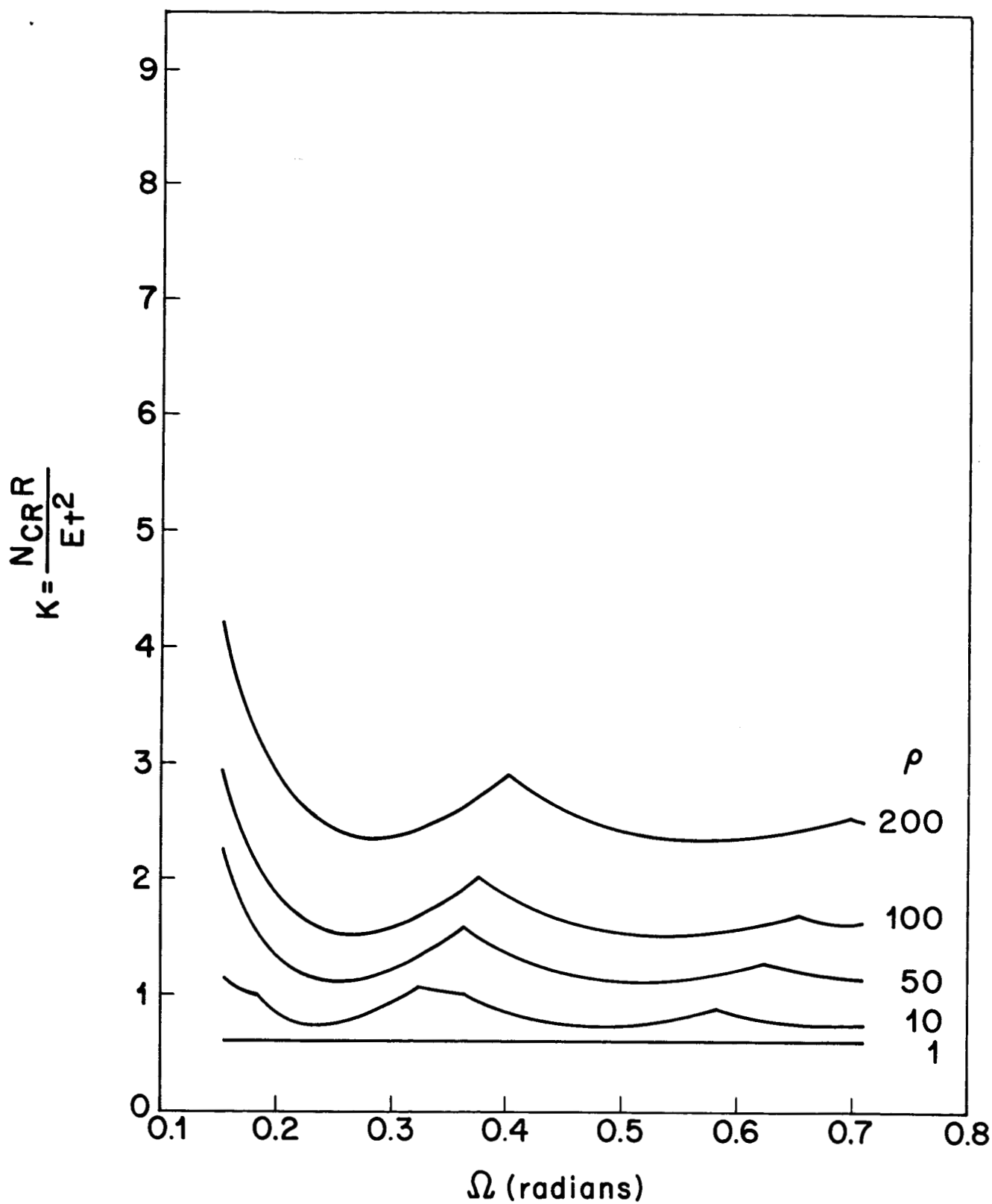


Fig. 5a Buckling Coefficients for Panels With Internal Stringers - $L/R = 0.5$;
 $R/t = 400$

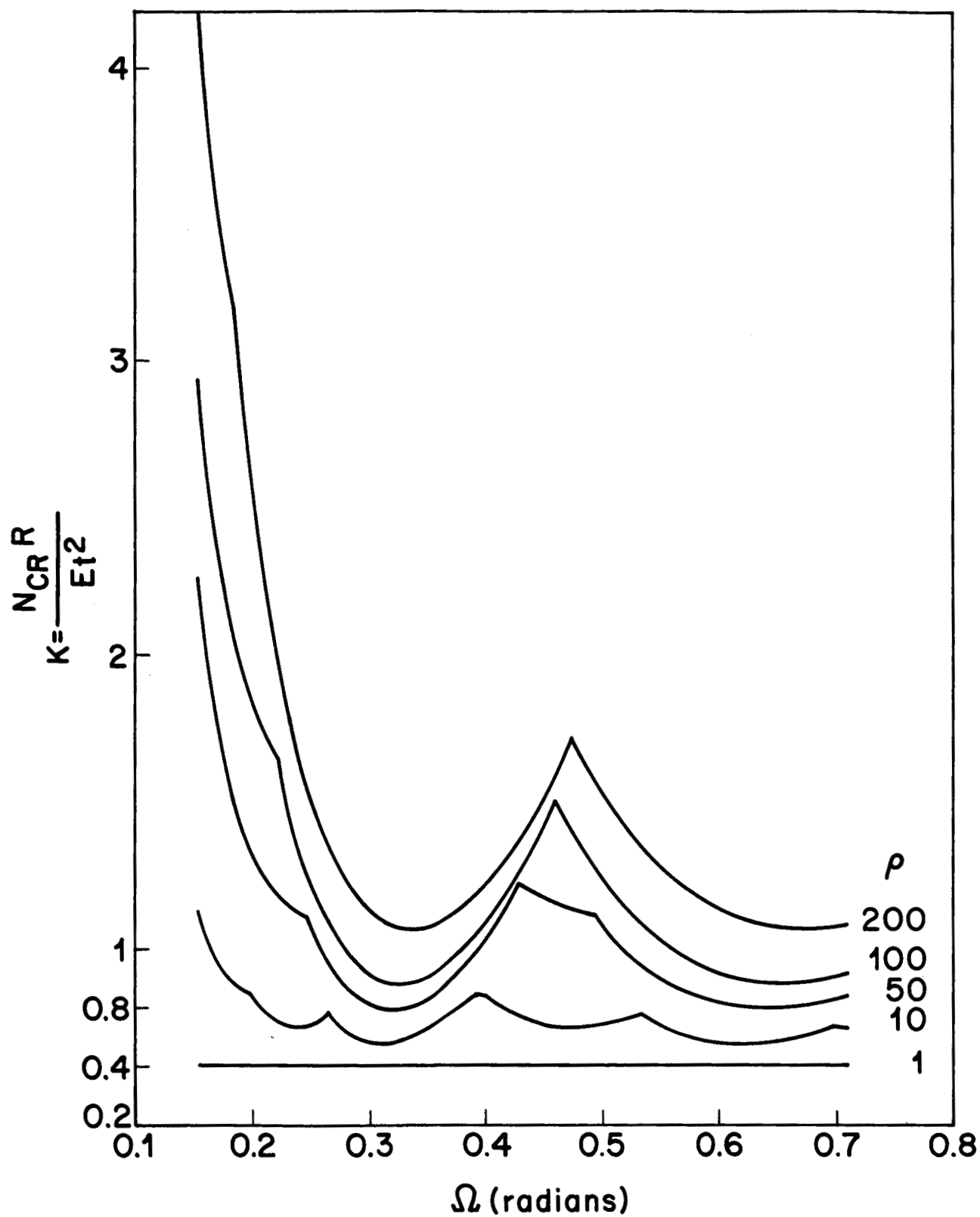


Fig. 5b Buckling Coefficients for Panels With Internal Stringers - $L/R = 1.0$;
 $R/t = 400$

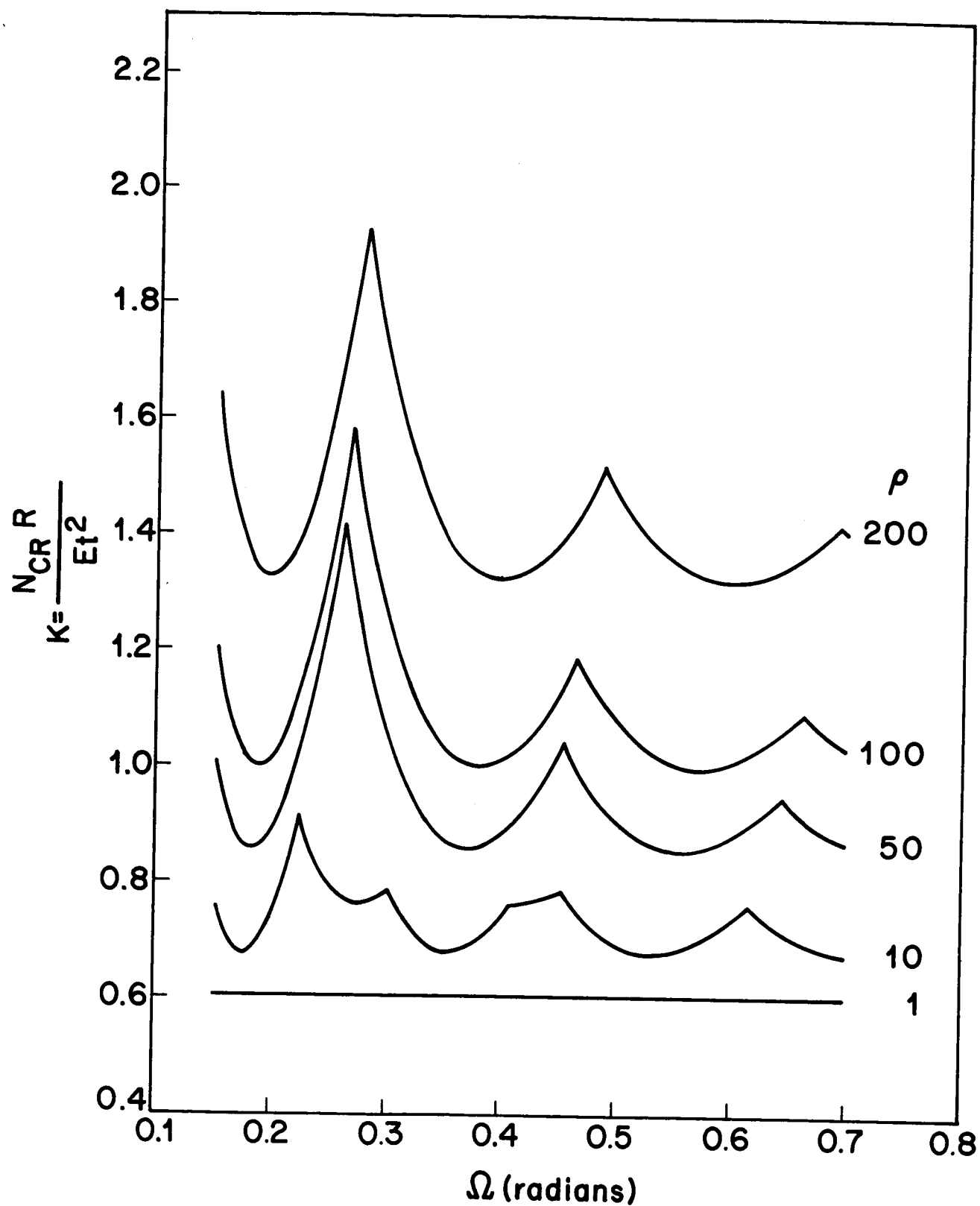


Fig. 6a Buckling Coefficients for Panels With Internal Stringers - $L/R = 0.5$;
 $R/t = 1000$

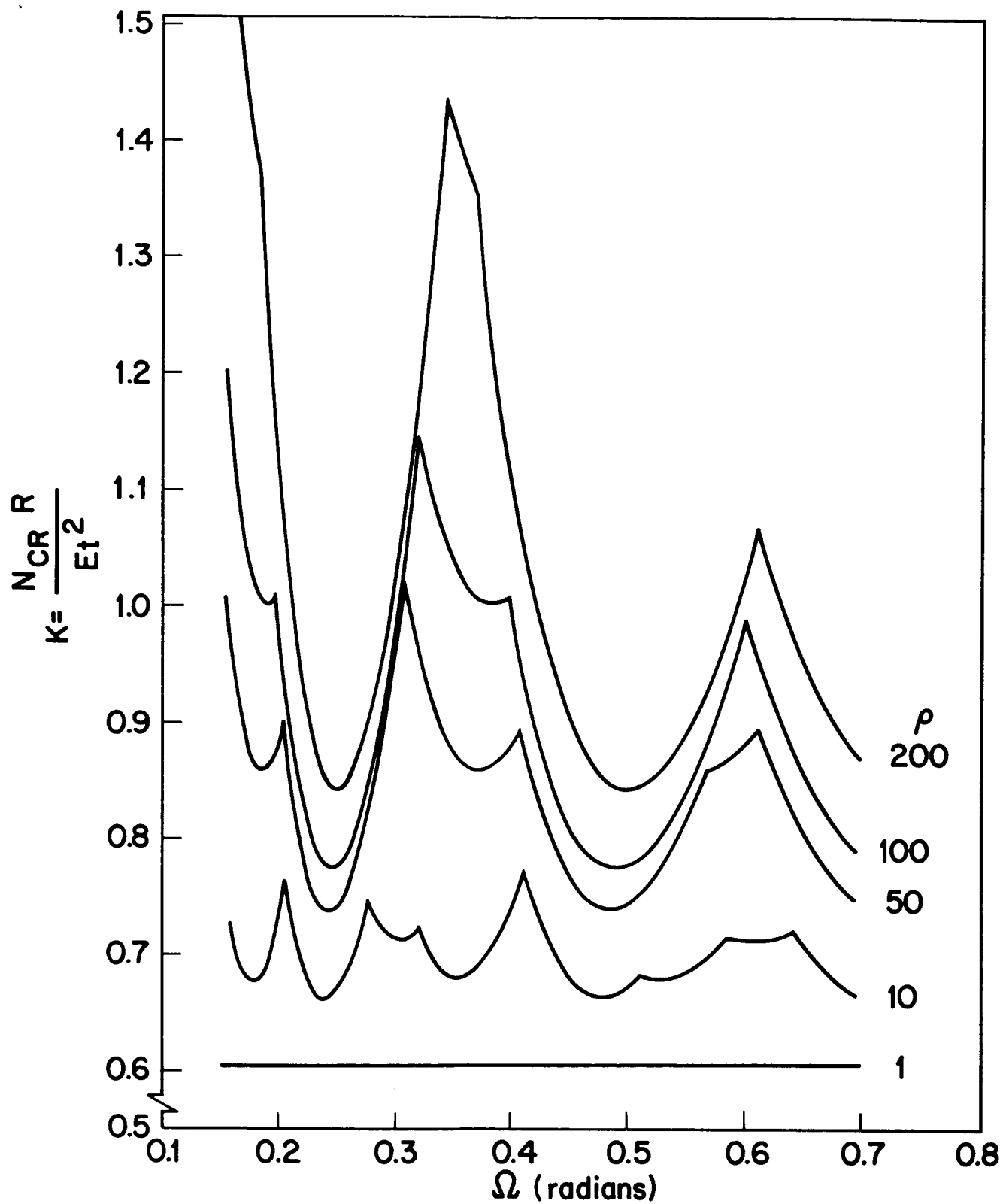


Fig. 6b Buckling Coefficients for Panels With Internal Stringers - $L/R = 1.0$;
 $R/t = 1000$

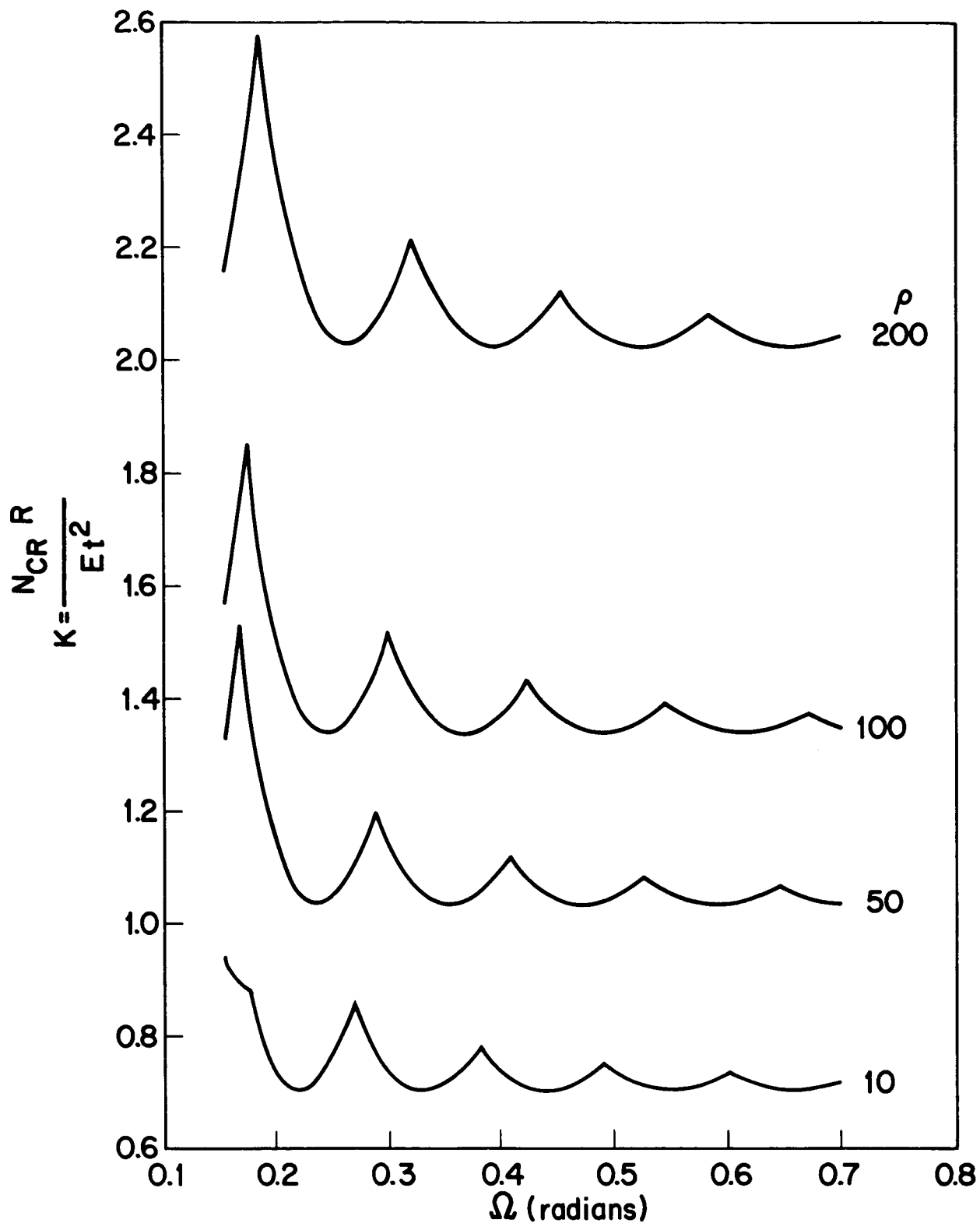


Fig. 7a Buckling Coefficients for Panels With Internal Stringers - $L/R = 0.25$;
 $R/t = 2000$

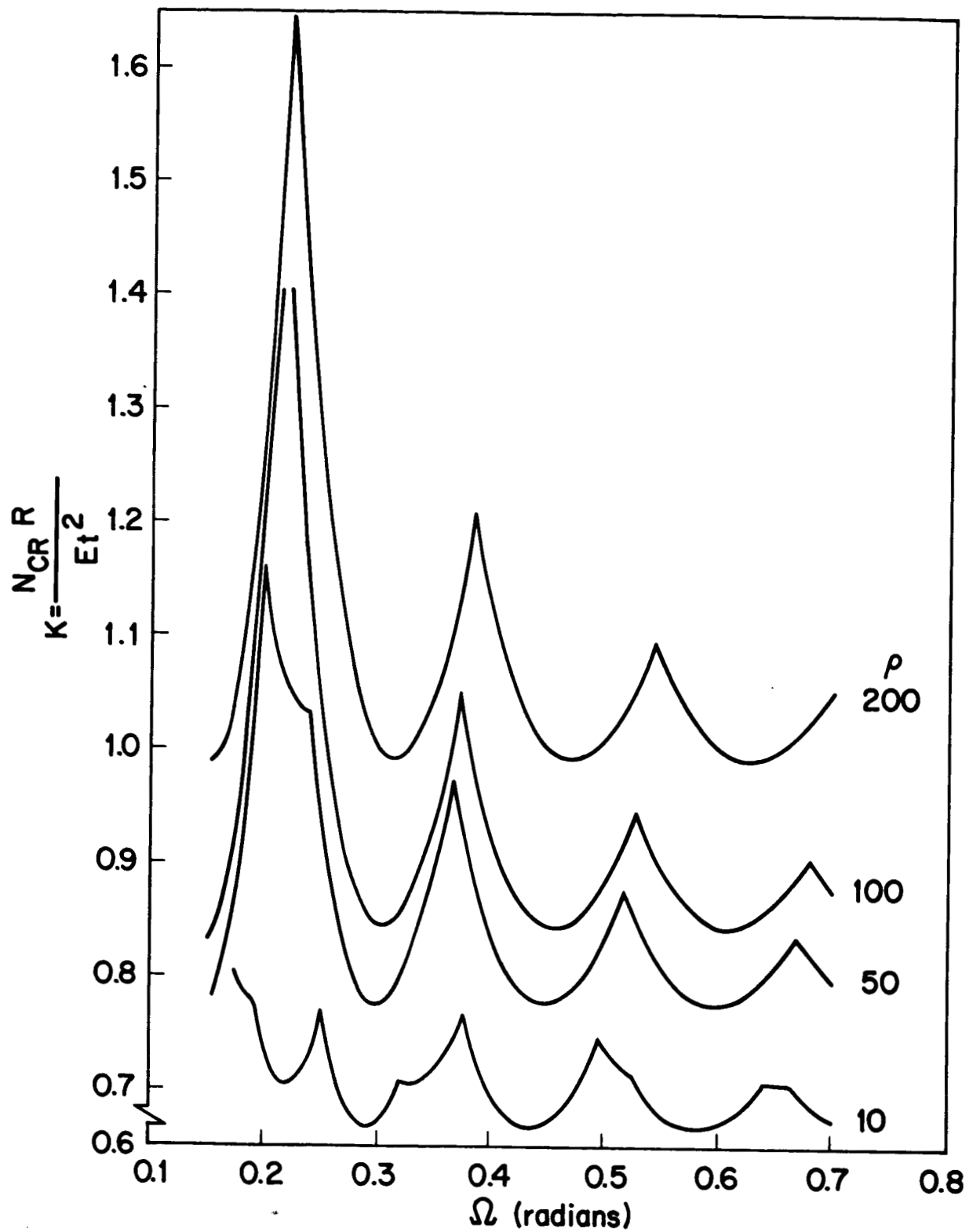


Fig. 7b Buckling Coefficients for Panels With Internal Stringers - $L/R = 0.5$;
 $R/t = 2000$

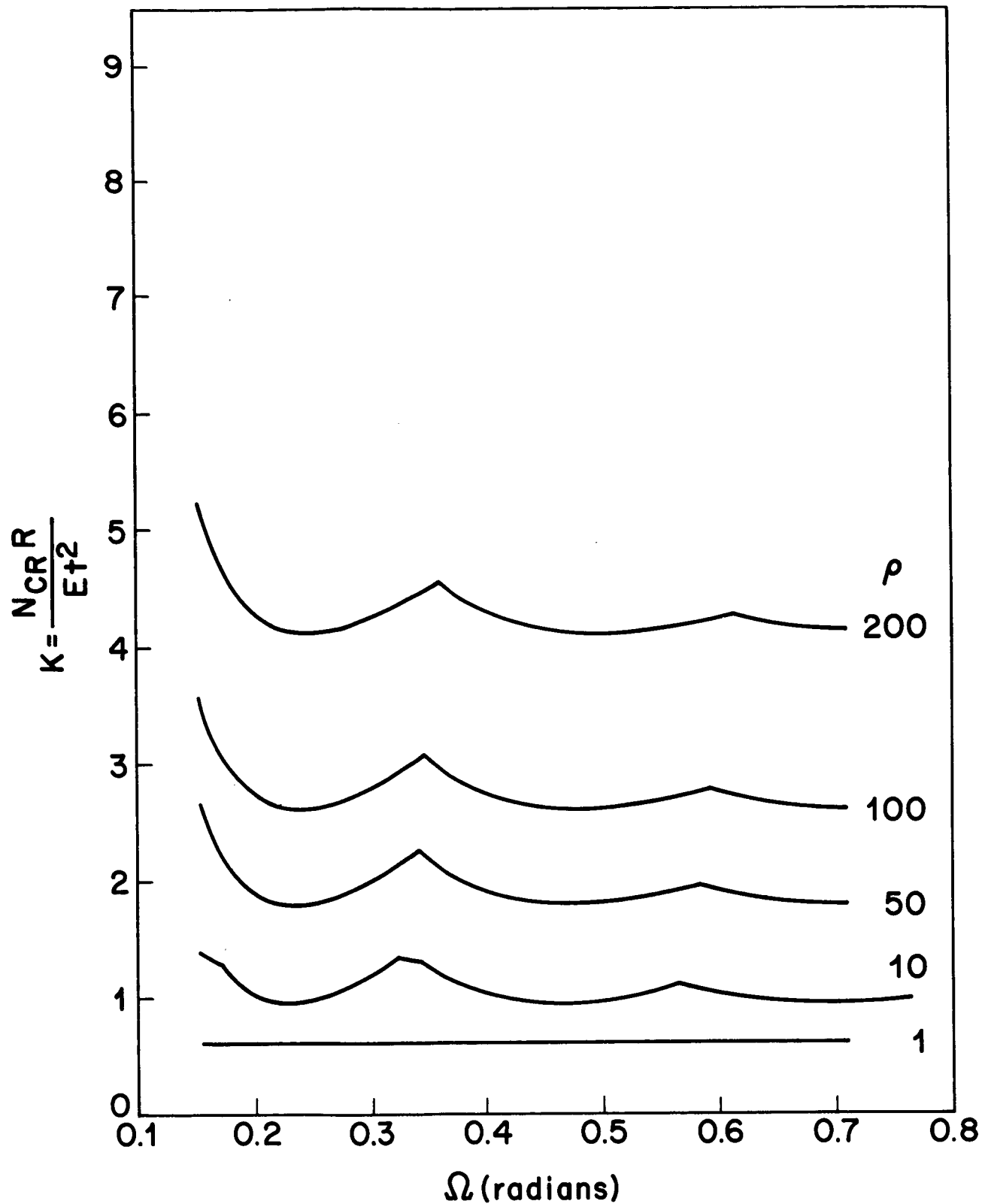


Fig. 8a Buckling Coefficients for Panels With External Stringers - $L/R = 0.5$;
 $R/t = 400$

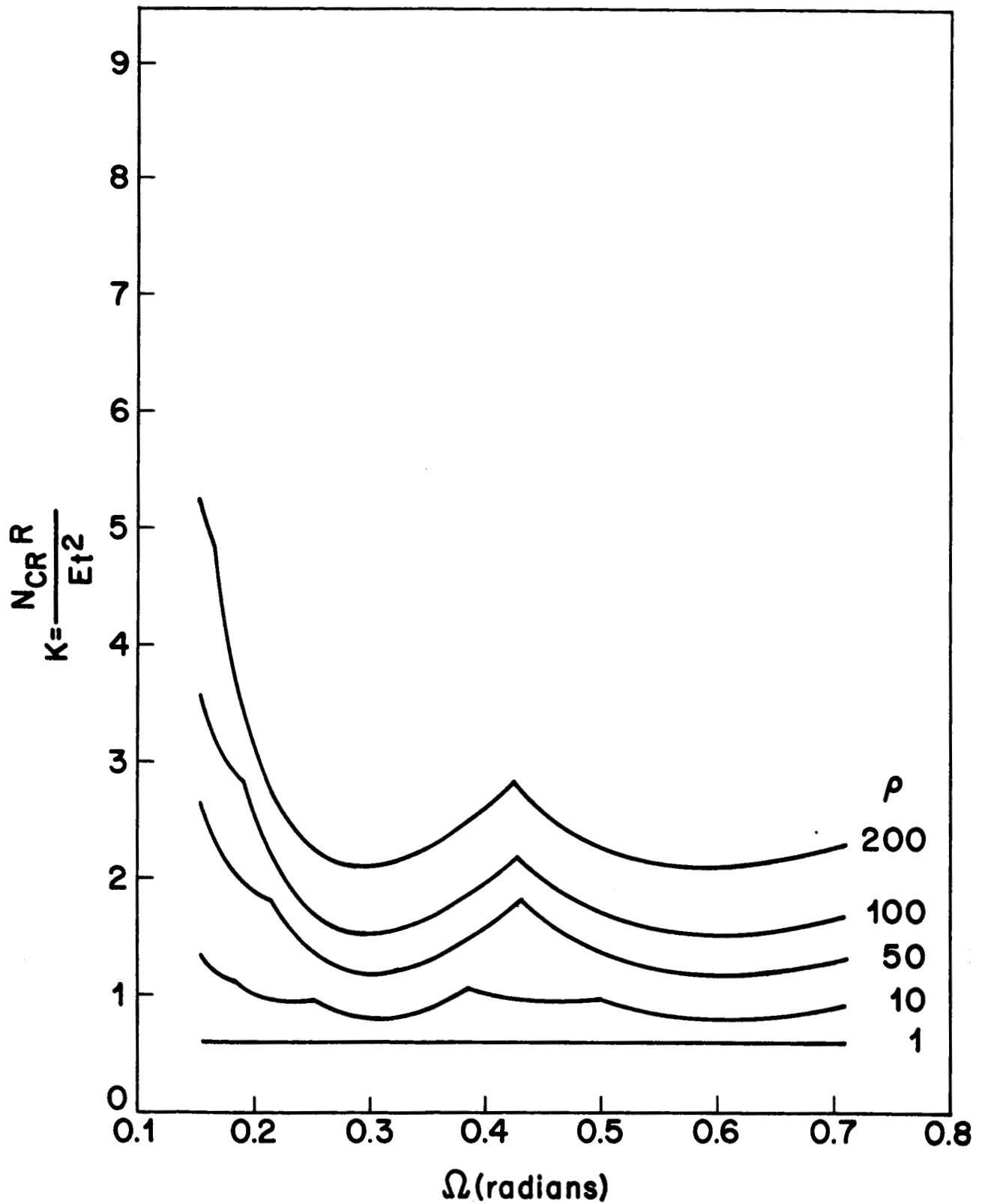


Fig. 8b Buckling Coefficients for Panels With External Stringers - $L/R = 1$;
 $R/t = 400$

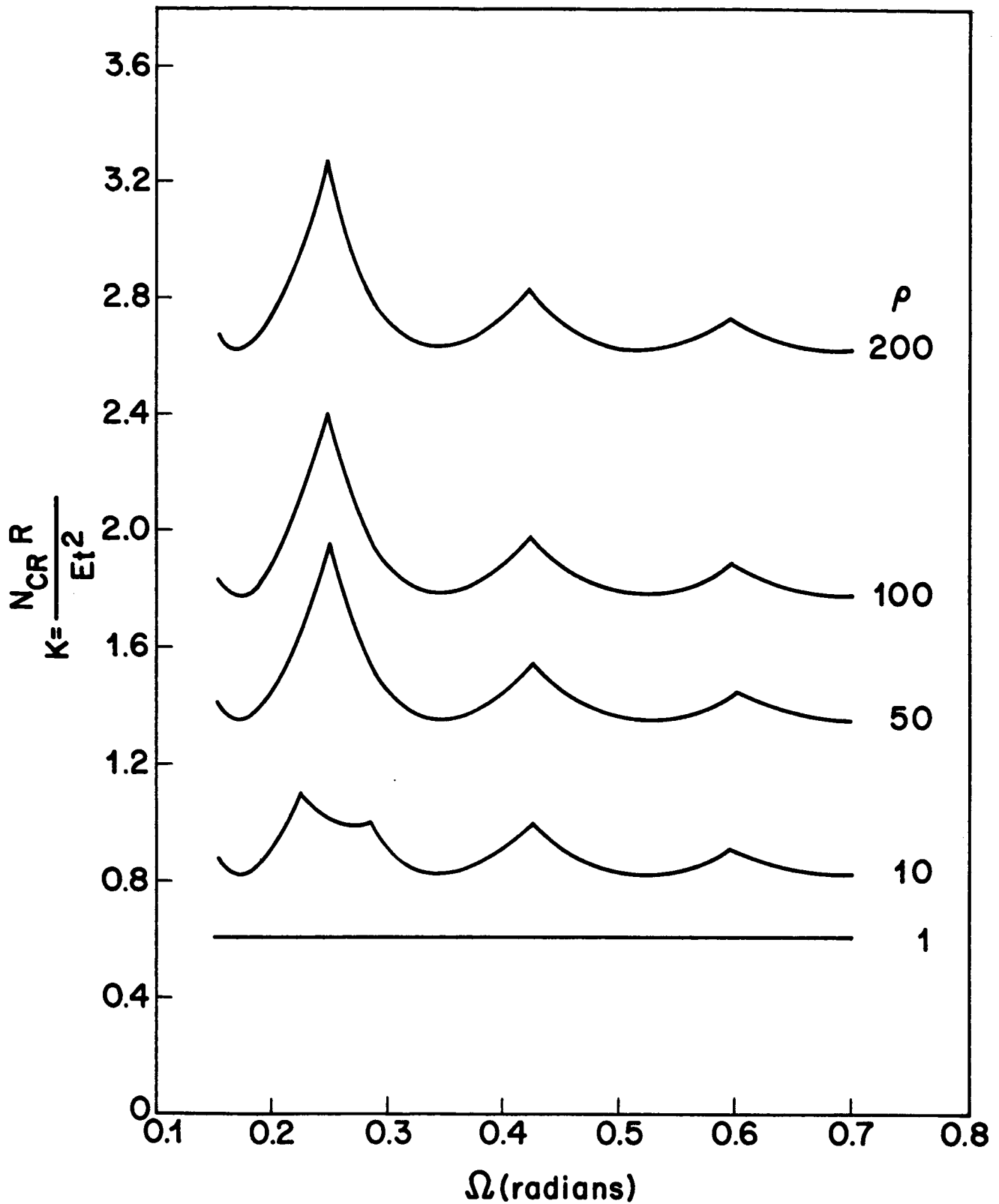


Fig. 9a Buckling Coefficients for Panels With External Stringers - $L/R = 0.5$;
 $R/t = 1000$

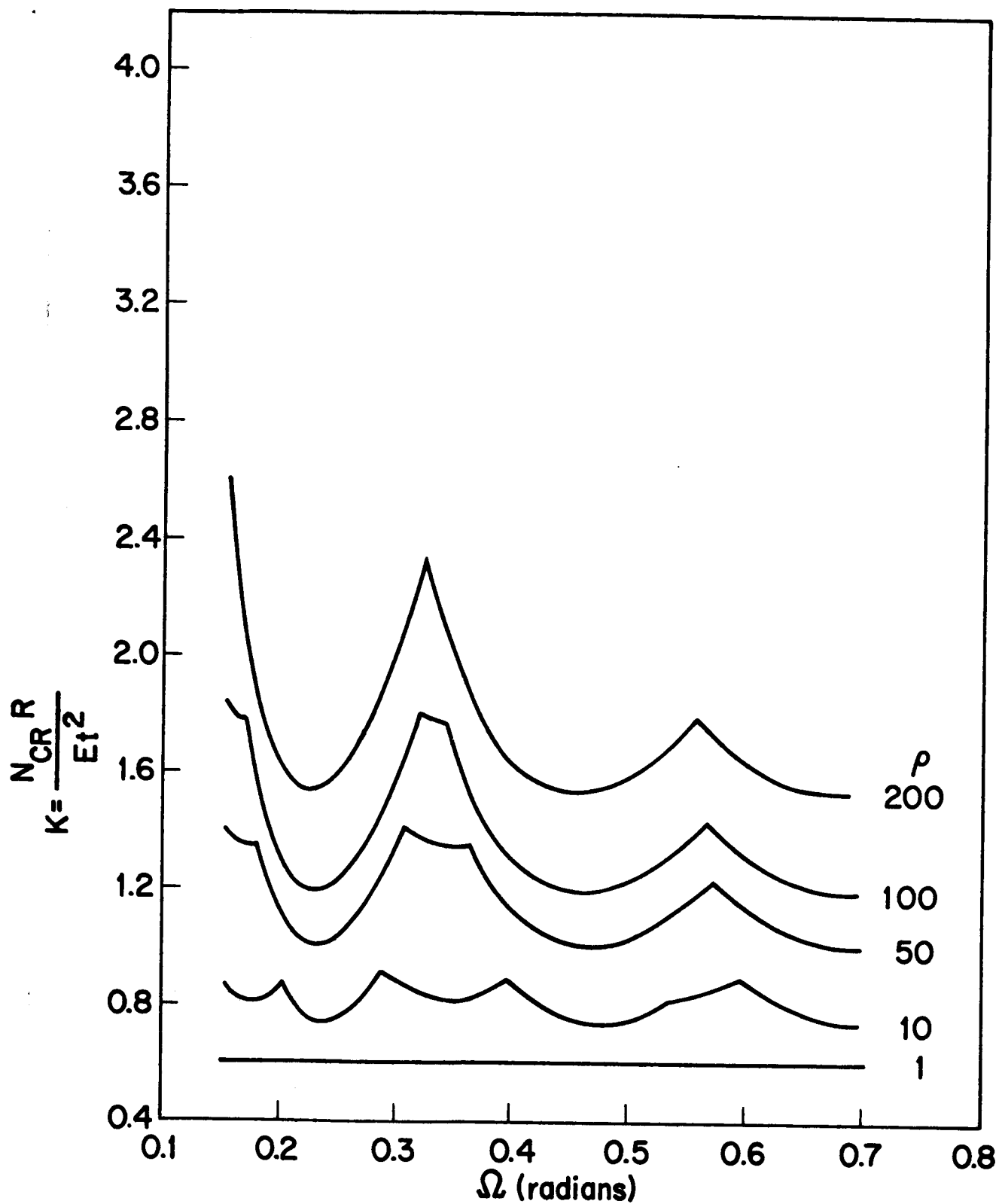


Fig. 9b Buckling Coefficients for Panels With External Stringers - $L/R = 1.0$;
 $R/t = 1000$

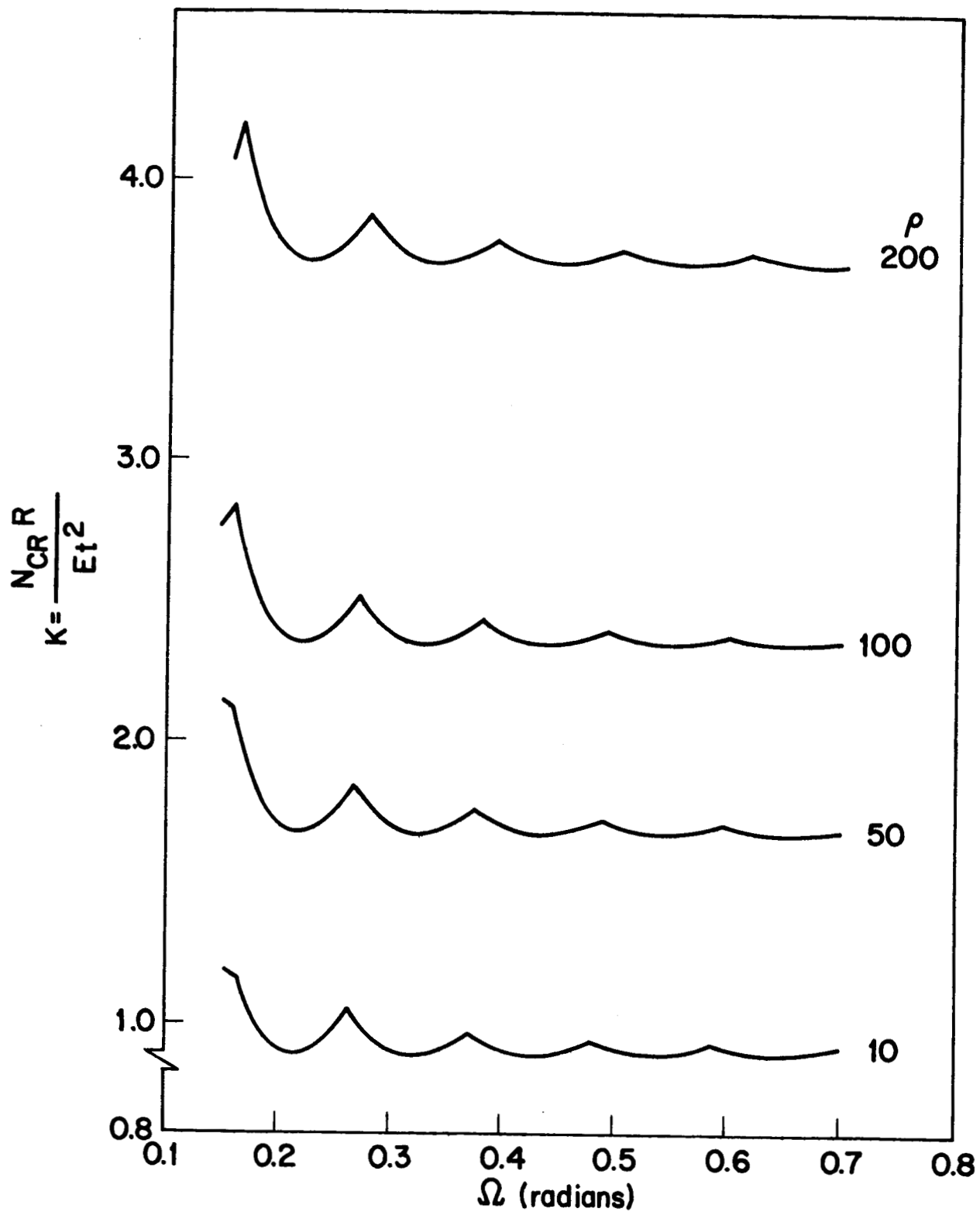


Fig. 10a Buckling Coefficients for Panels With External Stringers - $L/R = 0.25$;
 $R/t = 2000$

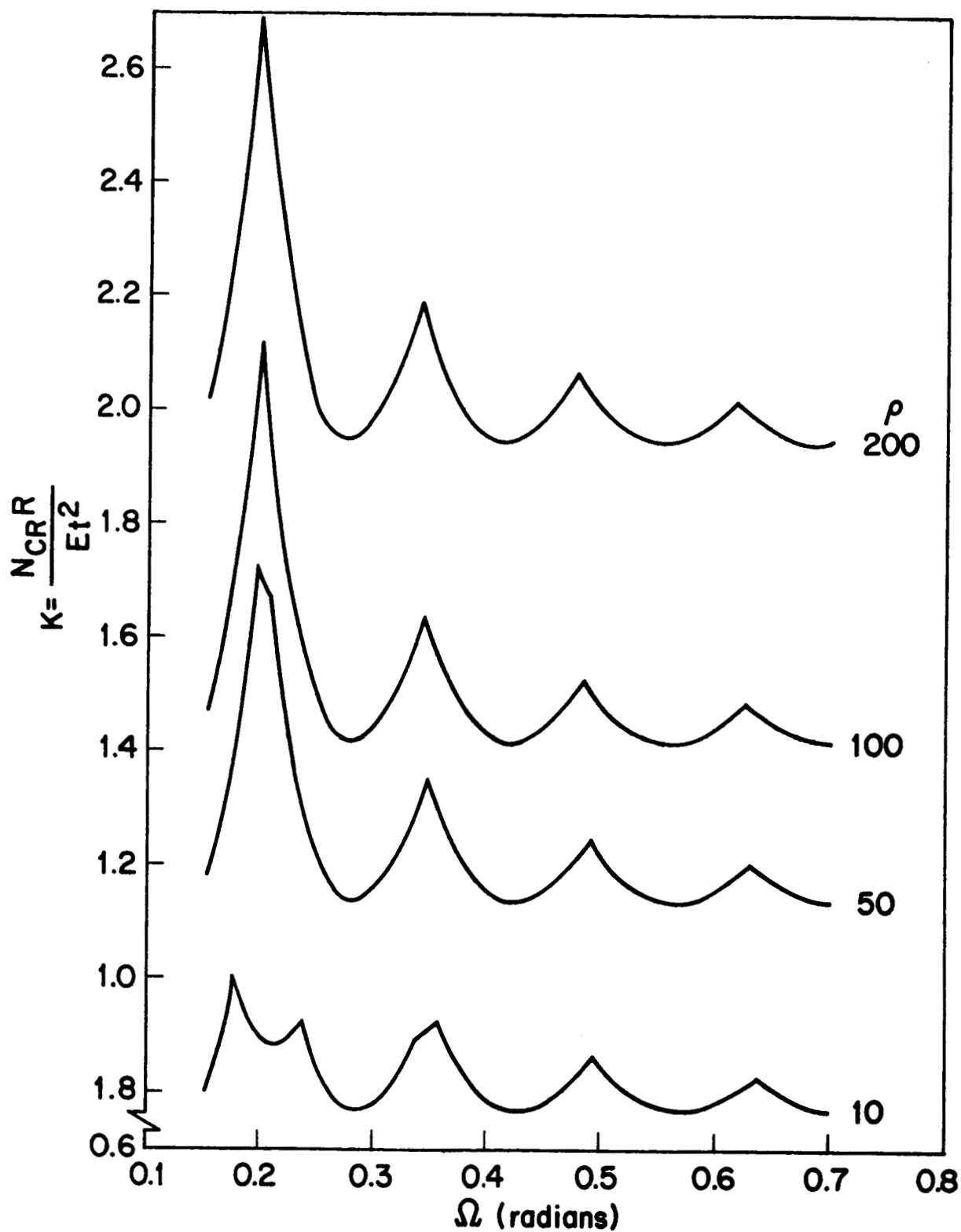


Fig. 10b Buckling Coefficients for Panels With External Stringers - $L/R = 0.5$;
 $R/t = 2000$

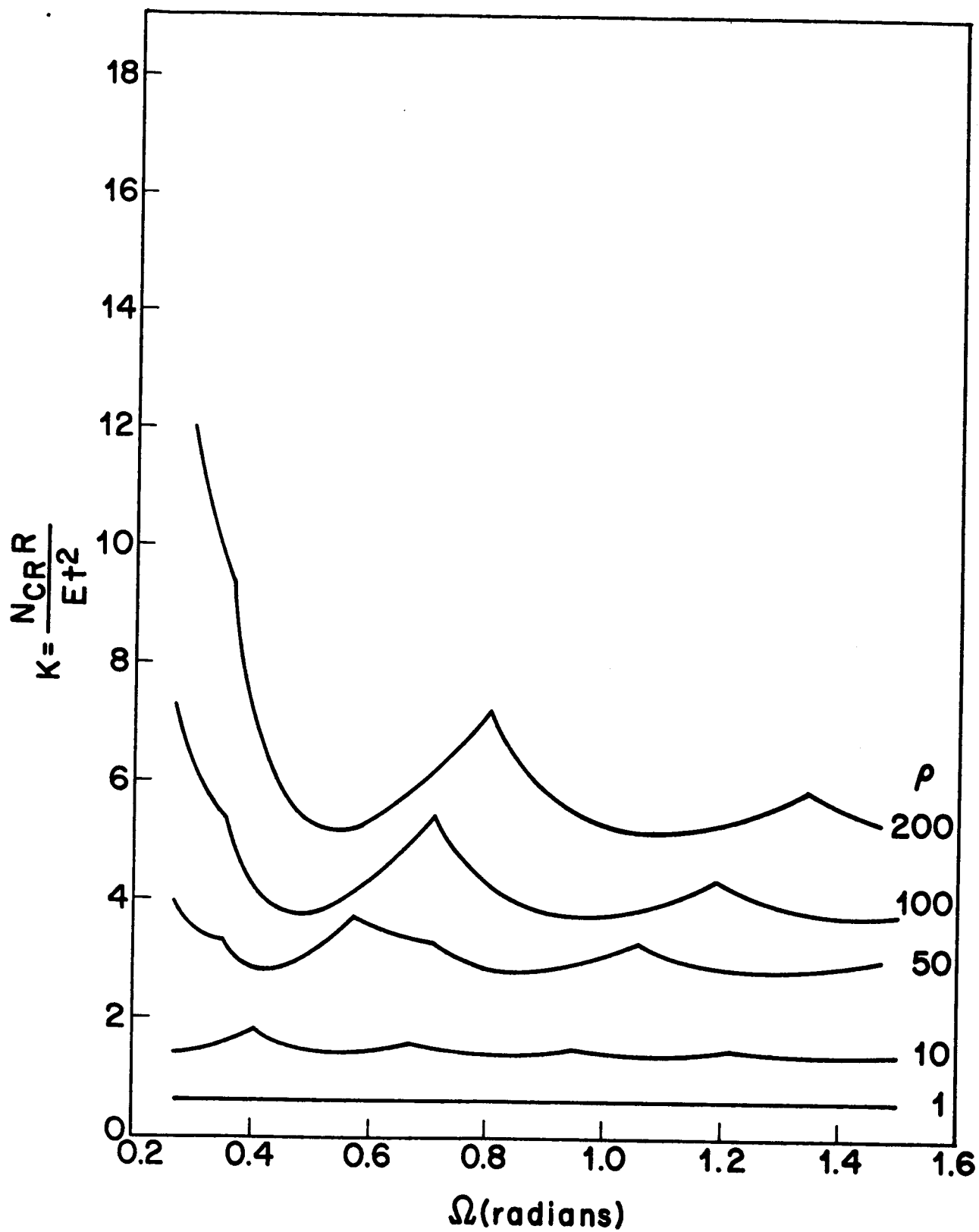


Fig. 11a Buckling Coefficients for Panels With Internal Waffles - $L/R = 0.5$;
 $R/t = 400$

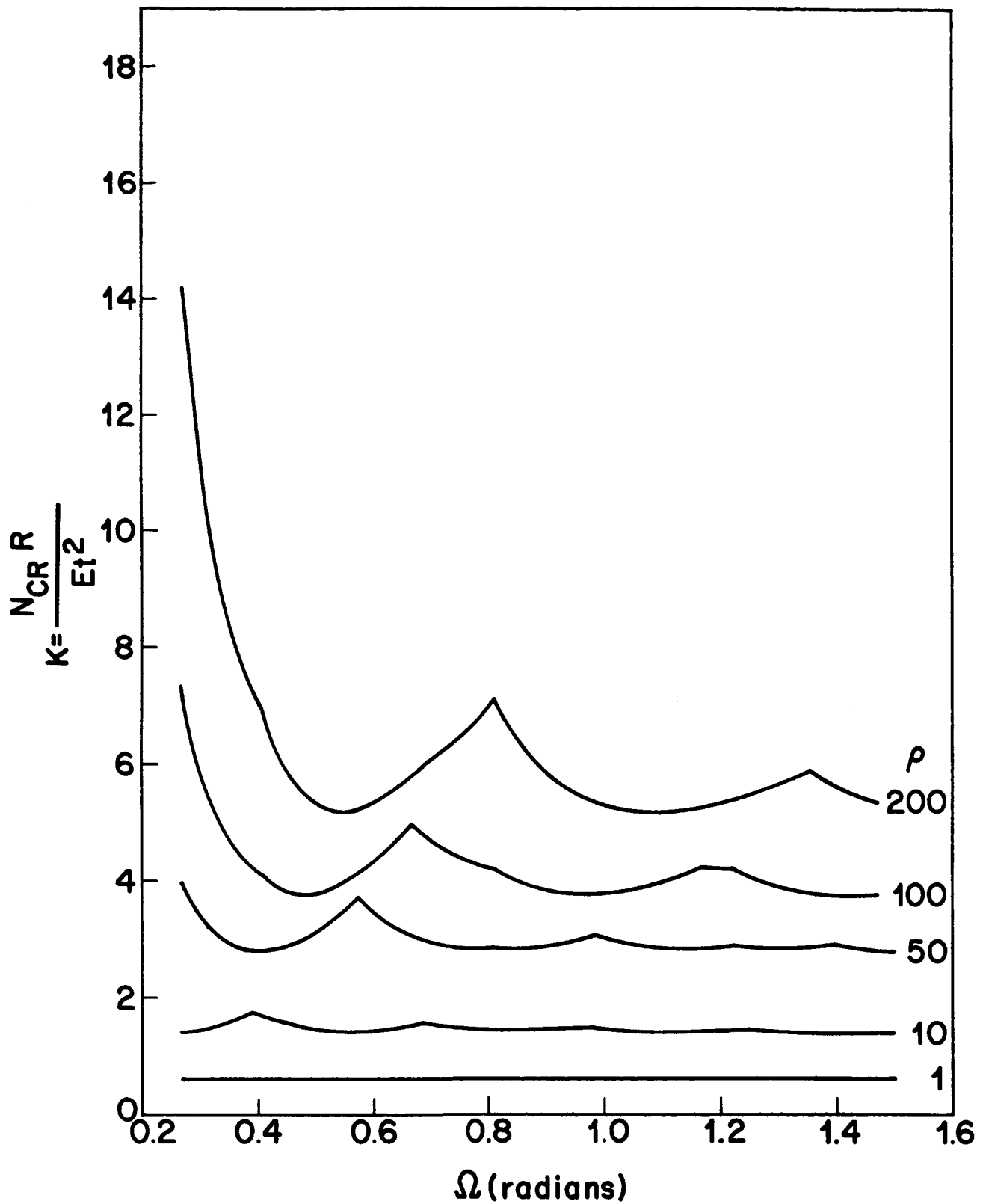


Fig. 11b Buckling Coefficients for Panels With Internal Waffles - $L/R = 1.0$;
 $R/t = 400$

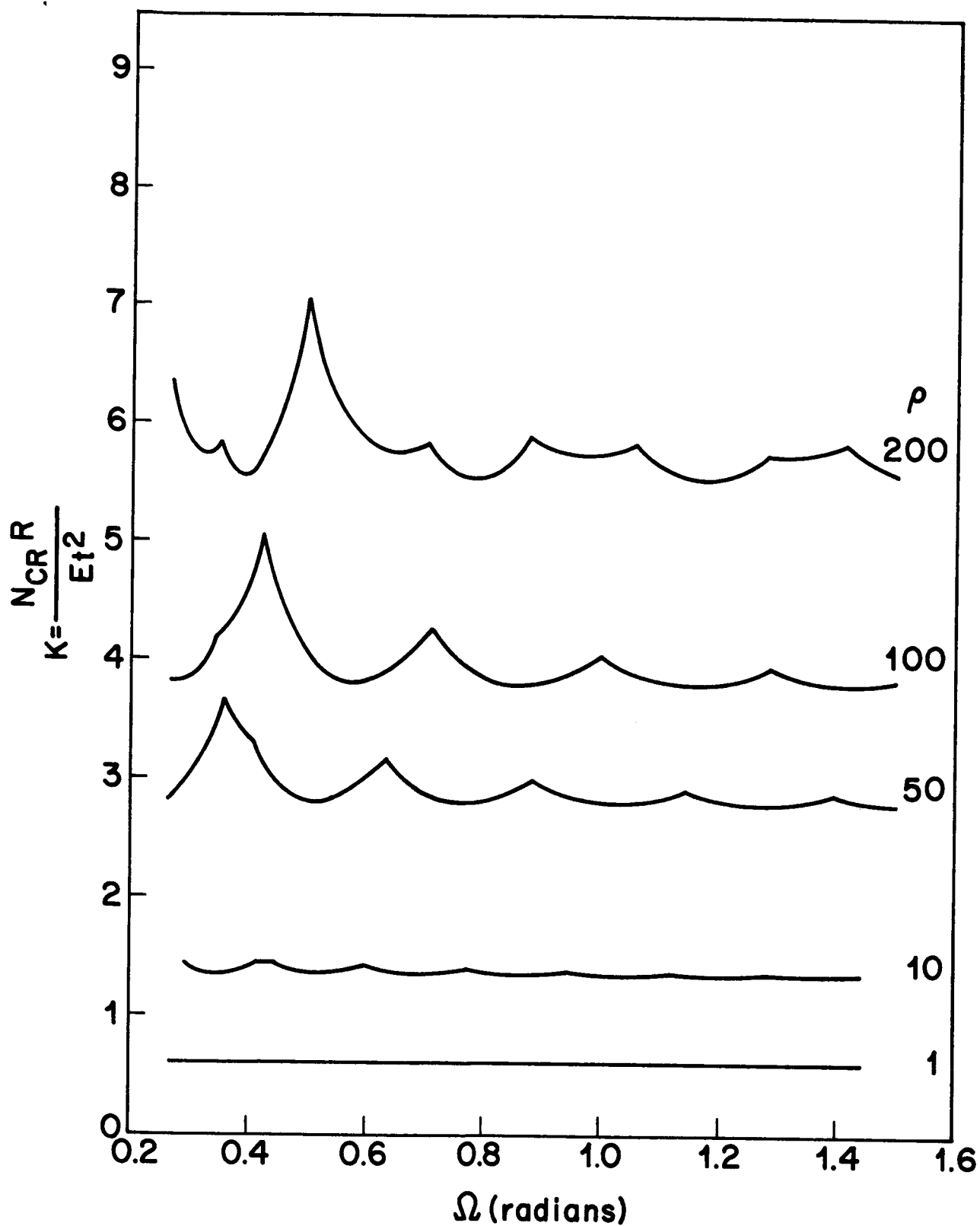


Fig. 12a Buckling Coefficients for Panels With Internal Waffles - $L/R = 0.5$;
 $R/t = 1000$

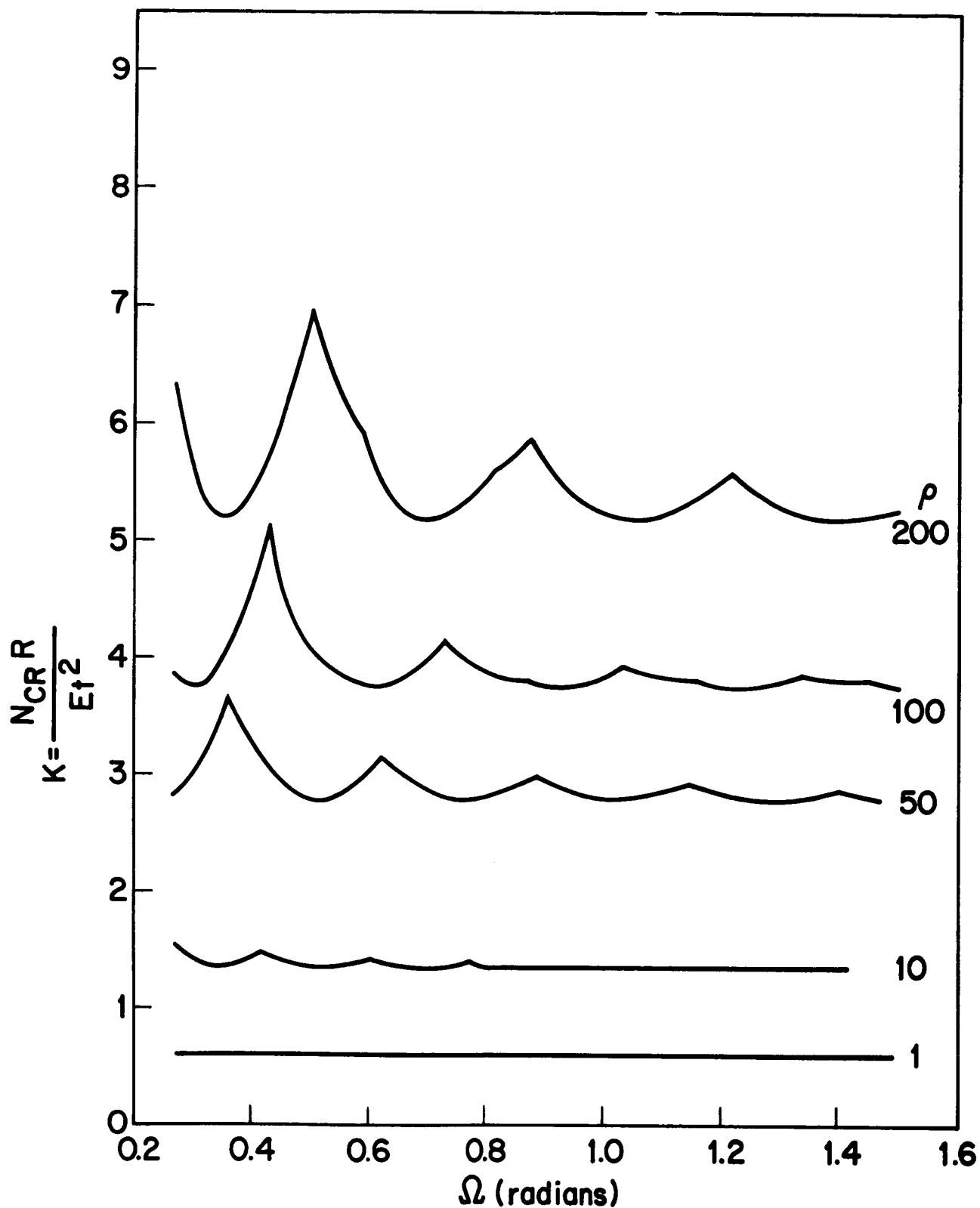


Fig. 12b Buckling Coefficients for Panels With Internal Waffles - $L/R = 1$;
 $R/t = 1000$

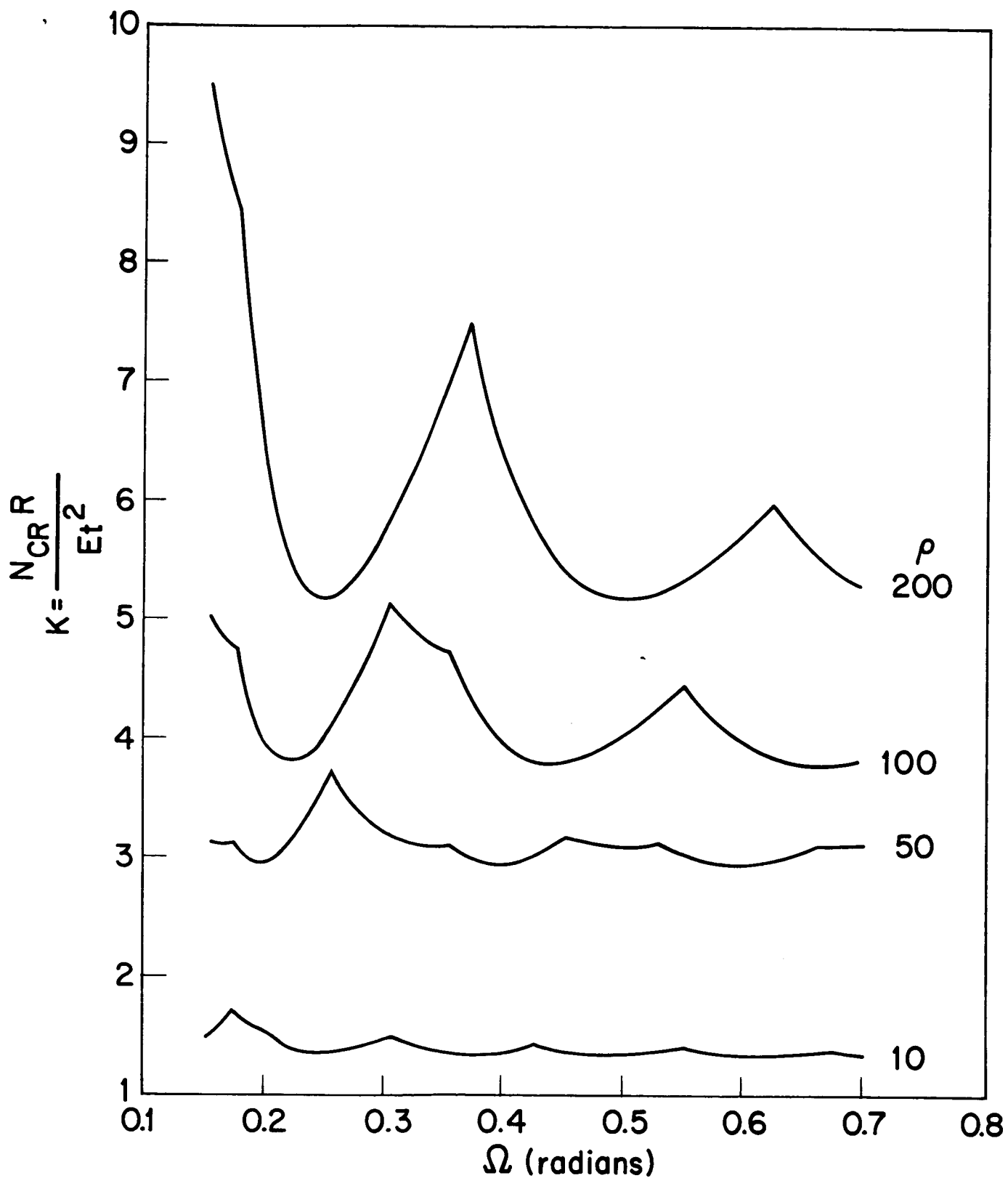


Fig. 13a Buckling Coefficients for Panels With Internal Waffles - $L/R = 0.25$;
 $R/t = 2000$

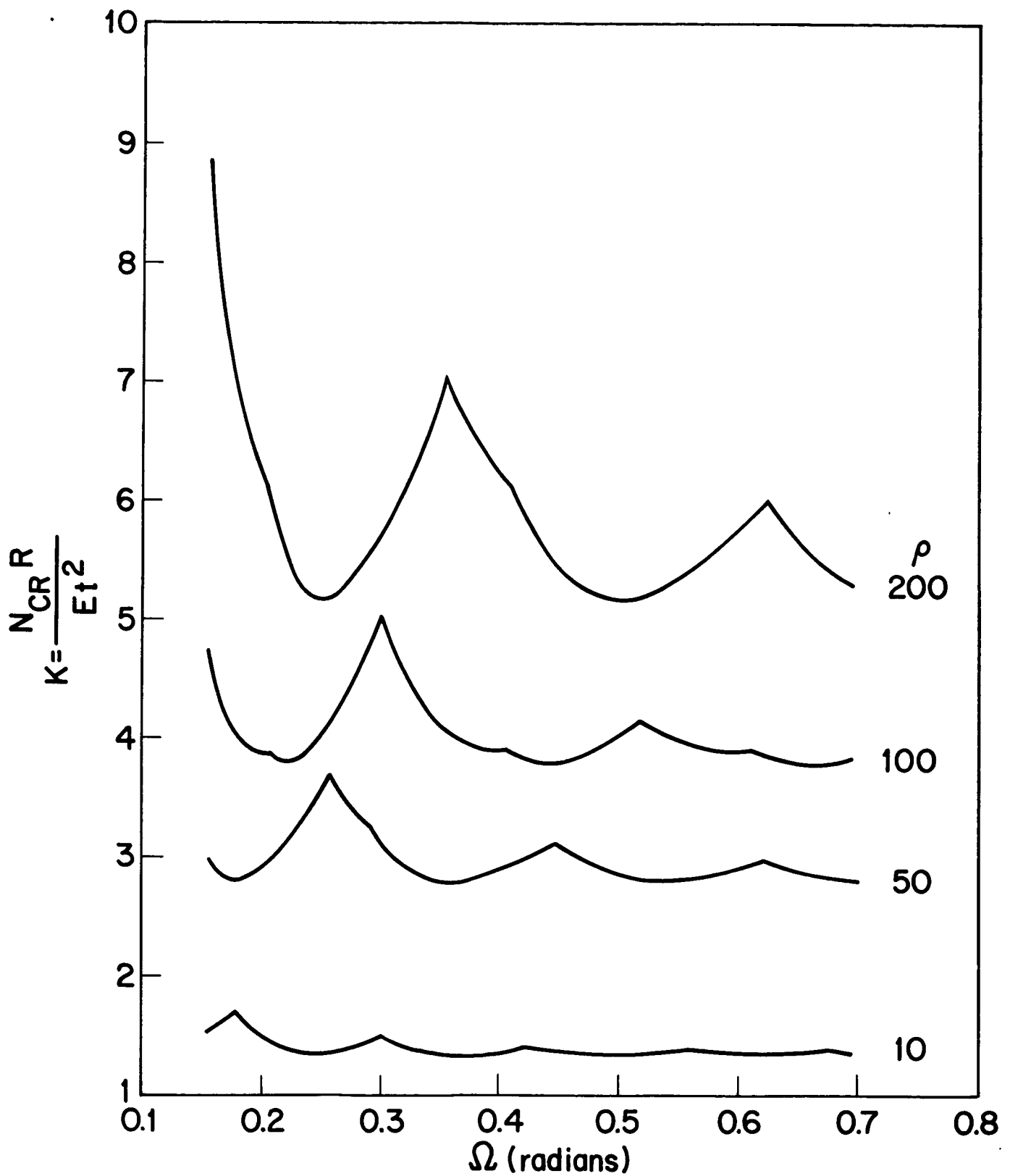


Fig. 13b Buckling Coefficients for Panels With Internal Waffles - $L/R = 0.5$;
 $R/t = 2000$

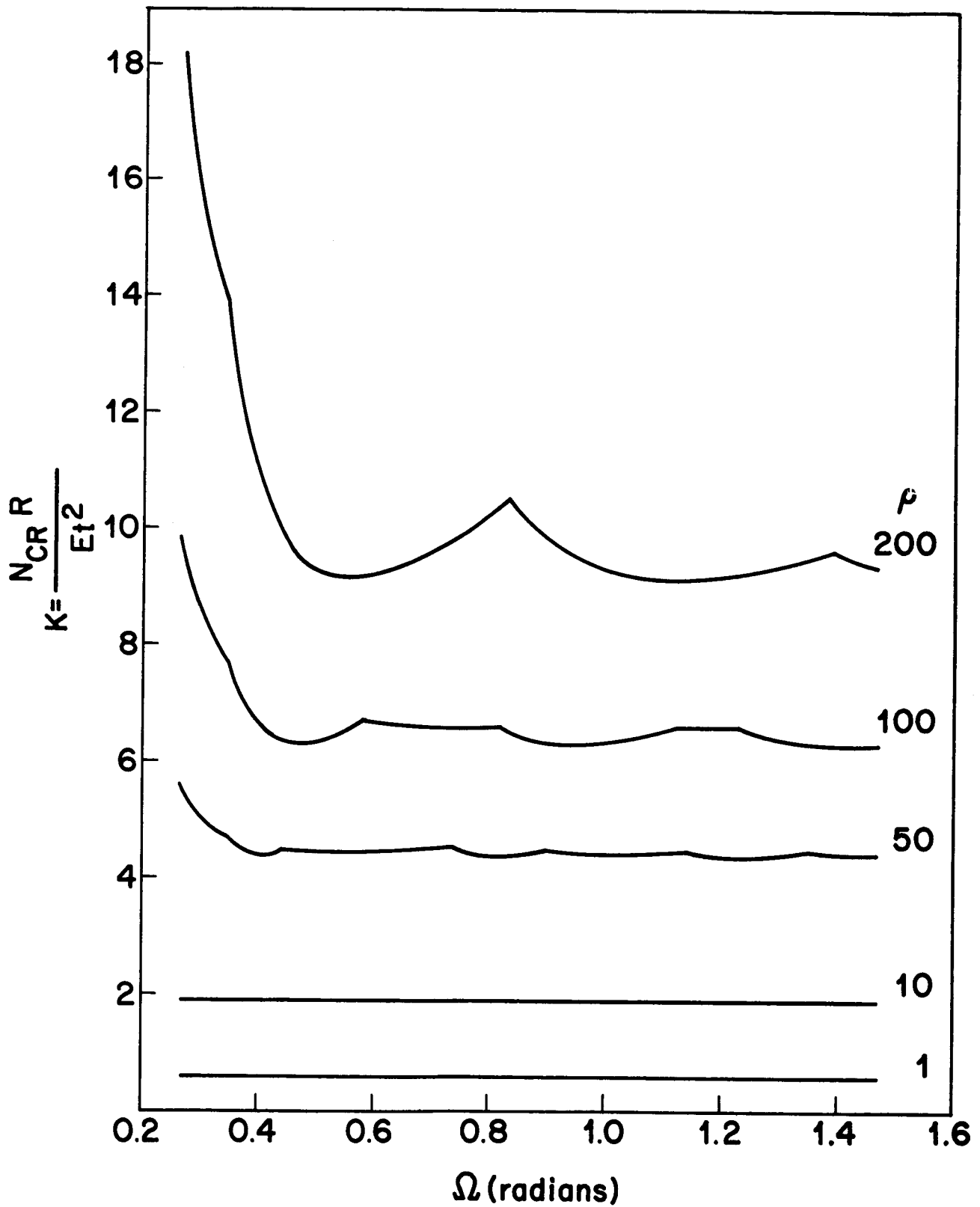


Fig. 14a Buckling Coefficients for Panels With External Waffles - $L/R = 0.5$;
 $R/t = 400$

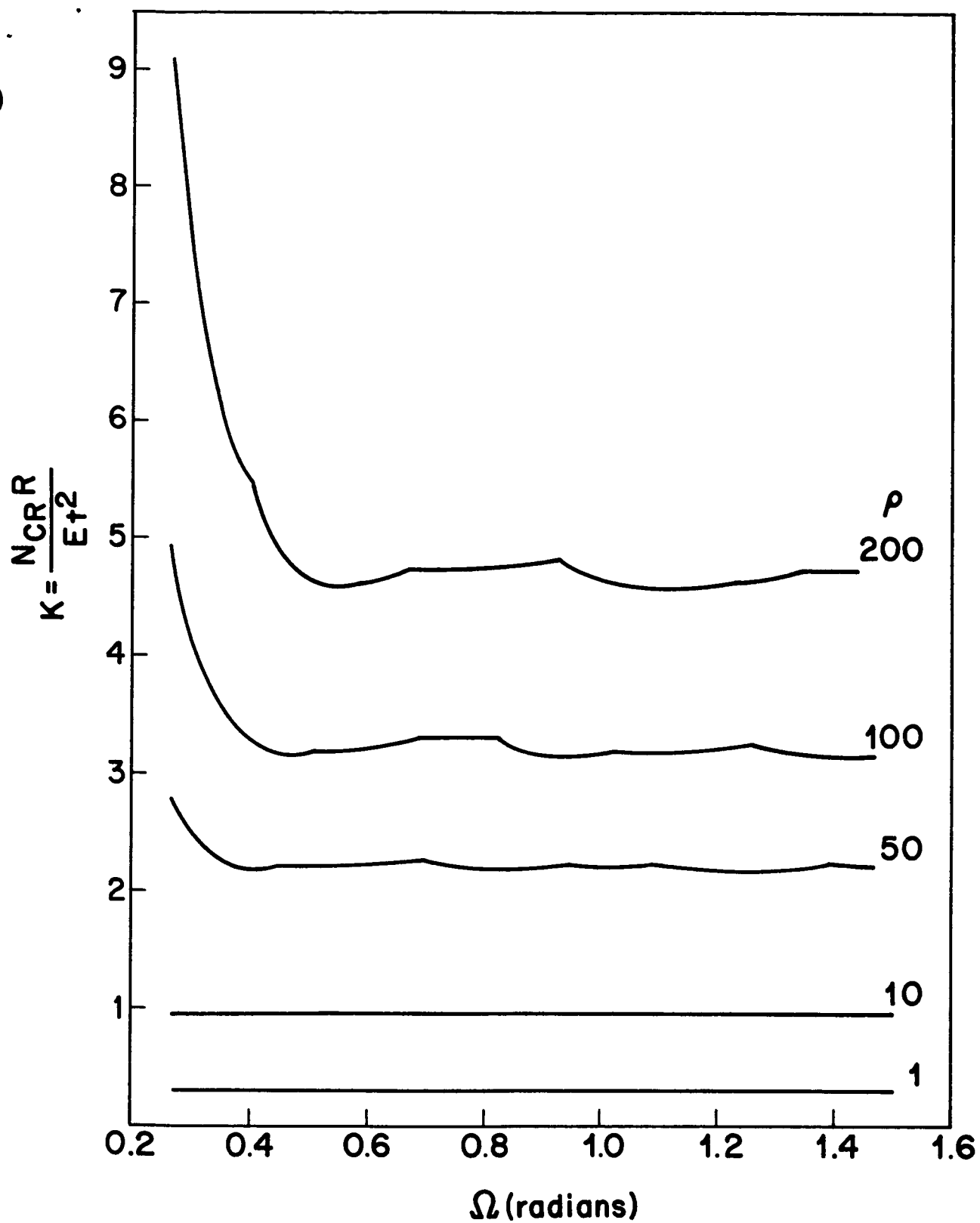


Fig. 14b Buckling Coefficients for Panels With External Waffles - $L/R = 1.0$;
 $R/t = 400$

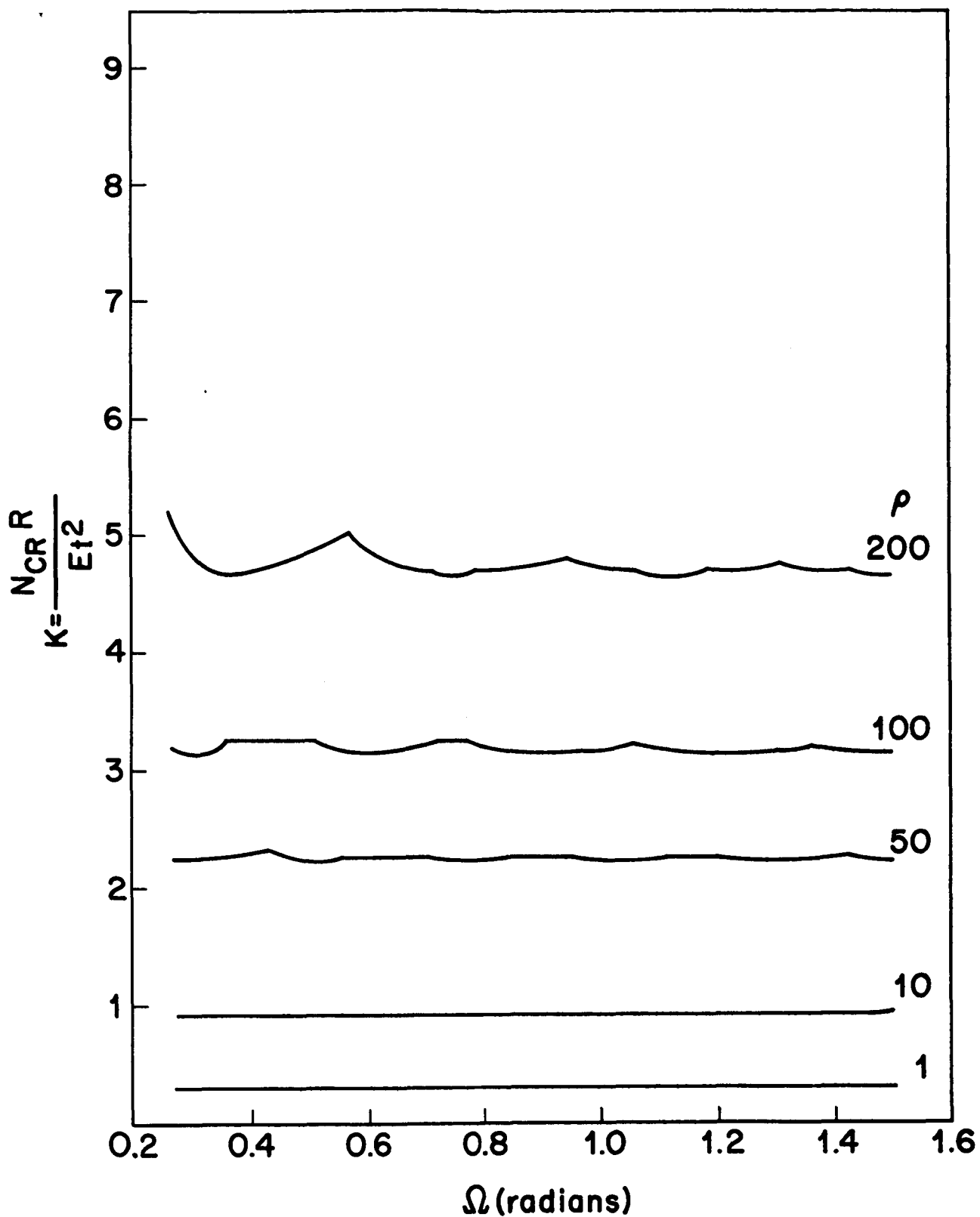


Fig. 15a Buckling Coefficients for Panels With External Waffles - $L/R = 1.0$;
 $R/t = 1000$

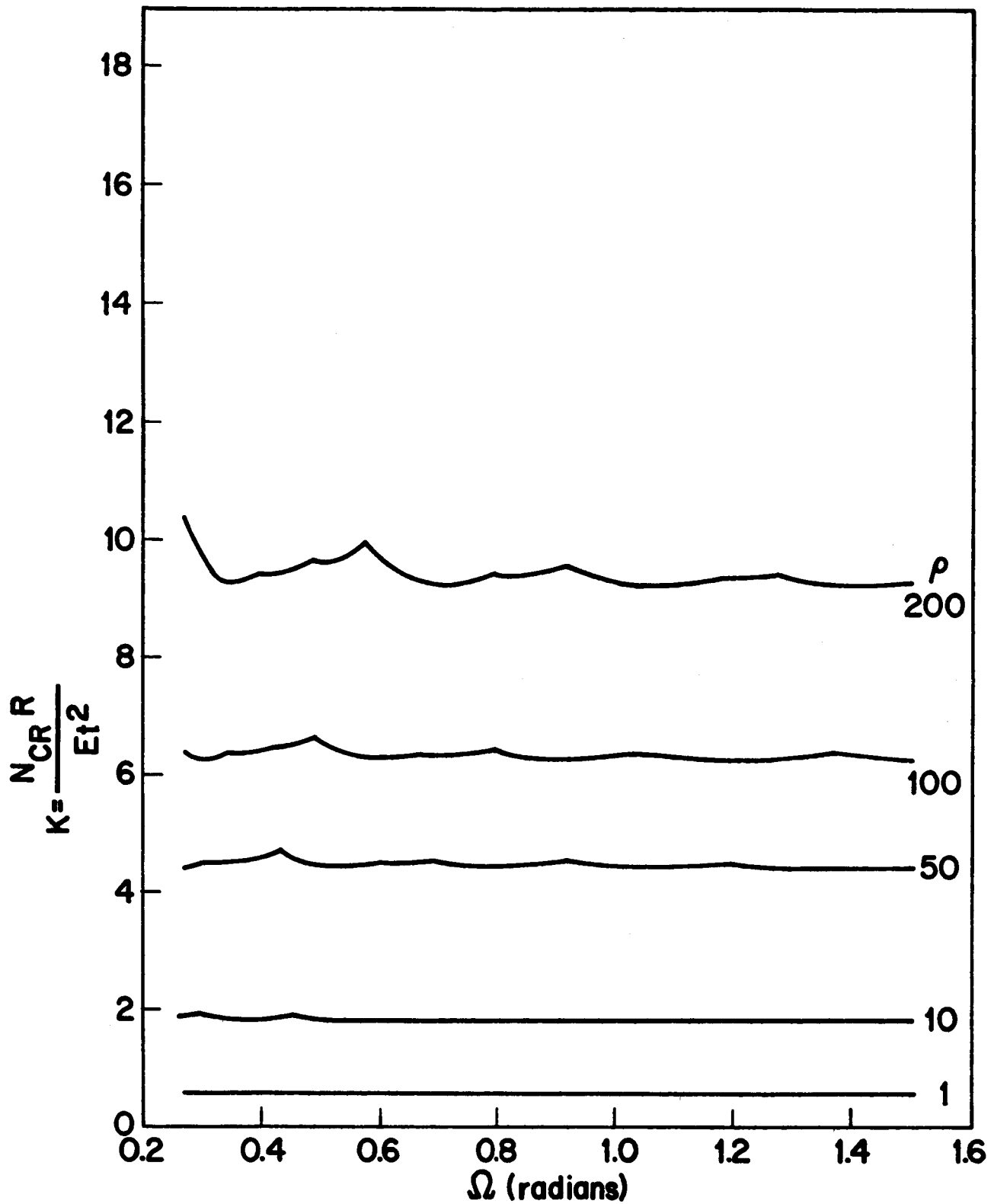


Fig. 15b Buckling Coefficients for Panels With External Waffles - $L/R = 1.0$;
 $R/t = 1000$

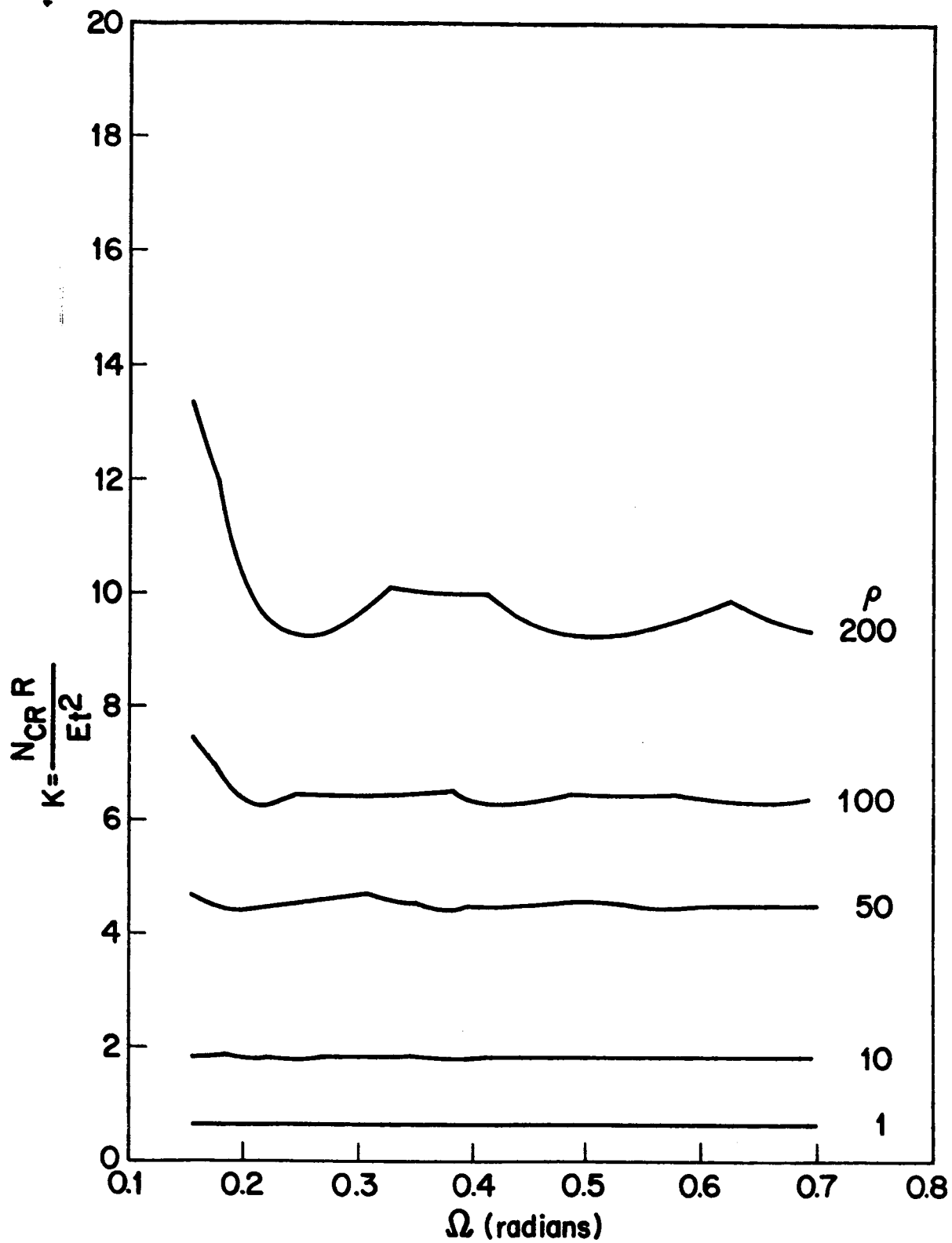


Fig. 16a Buckling Coefficients for Panels With External Waffles - $L/R = 0.25$;
 $R/t = 2000$

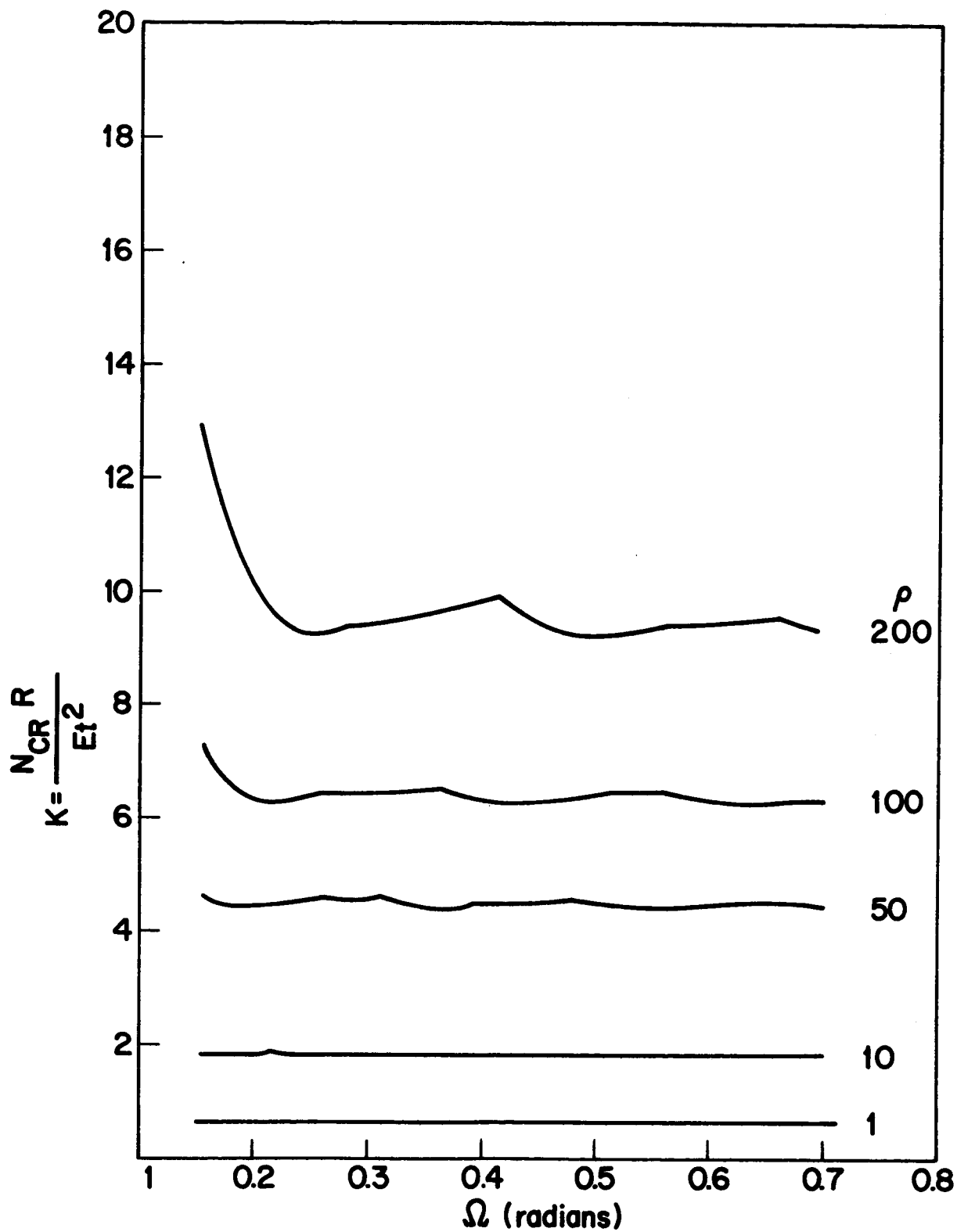


Fig. 16b Buckling Coefficients for Panels With External Waffles - $L/R = 0.5$;
 $R/t = 2000$

Appendix A INFLUENCE COEFFICIENTS FOR ORTHOTROPIC PLATES

An analysis of influence coefficients for orthotropic rectangular plates subjected to edge moments and shears is presented in this appendix. The influence coefficients are used to determine the elastic restraint provided to cylindrical panels supported along the straight edges by stiffened web plates (see Fig. 2. Vol. II).

The following equation which governs the bending of an orthotropic flat plate may be obtained through specialization of the equations for an orthotropic cylinder (see Sec. 2. Vol. II):

$$D_x w_{,xxxx} + 2H w_{,xxyy} + D_y w_{,yyyy} = 0 \quad (A. 1)$$

where

$$H = D_1 + 2D_{xy}$$

$$D_x = \frac{Et^3}{12(1 - \nu^2)} \left\{ 1 + \frac{E_1}{E} \left(\frac{d_1}{S_1} \right) \left(\frac{h_1}{t} \right) \left[\left(\frac{h_1}{t} \right)^2 + \frac{3[(h_1/t) + 1]^2}{(1 - \nu^2)(E_1/E)(d_1/S_1)(h_1/t) + 1} \right] (1 - \nu^2) \right\}$$

$$D_y = \frac{Et^3}{12(1 - \nu^2)} \left\{ 1 + \frac{E_2}{E} \left(\frac{d_2}{S_2} \right) \left(\frac{h_2}{t} \right) \left[\left(\frac{h_2}{t} \right)^2 + \frac{3[(h_2/t) + 1]^2}{(1 - \nu^2)(E_2/E)(d_2/S_2)(h_2/t) + 1} \right] (1 - \nu^2) \right\}$$

$$D_1 = \frac{\nu Et^3}{12(1 + \nu)}$$

$$D_{xy} = \frac{Et^3}{24(1 + \nu)} \left\{ 1 + \frac{1}{2} \left[\frac{(E_1/E)J_1}{(1 + \nu_1)S_1} + \frac{(E_2/E)J_2}{(1 + \nu_2)S_2} \right] (1 + \nu) \right\}$$

and S_i is the spacing of the stringers; subscripts 1 and 2 refer to the x- and y-directions; w is the lateral displacement of the plate.

The x and y axes are taken so that the edges of the plate are at $x = 0$, $x = a$, $y = 0$, and $y = b$. The plate is assumed to be simply supported along the edges $y = 0$, $y = b$, and $x = a$. A moment M and a transverse shear force V are specified along the edge $x = 0$. These edge loadings may be represented by the following Fourier series:

$$M = \sum_{m=1}^{\infty} M_m \sin \frac{m\pi y}{b} \quad (\text{A. 2a})$$

$$V = \sum_{m=1}^{\infty} V_m \sin \frac{m\pi y}{b} \quad (\text{A. 2b})$$

Thus, the boundary conditions for the plate are

$$\underline{y = 0} \quad w = w_{,yy} = 0 \quad (\text{A. 3a})$$

$$\underline{y = b} \quad w = w_{,yy} = 0 \quad (\text{A. 3b})$$

$$\underline{x = a} \quad w = w_{,xx} = 0 \quad (\text{A. 3c})$$

$$\underline{x = 0} \quad M_x = M = \sum_{m=1}^{\infty} M_m \sin \frac{m\pi}{b} y = -D_x w_{,xx} - D_1 w_{,yy} \quad (\text{A. 3d})$$

$$V_x = V = \sum_{m=1}^{\infty} V_m \sin \frac{m\pi}{b} y$$

$$= -D_x w_{,xxx} - (H + 2D_{xy}) w_{,xyy} \quad (\text{A. 3e})$$

By use of the following Fourier series representation of the displacement function, the space variables in Eq. (A. 1) can be separated and the boundary conditions along the edges $y = 0$ and $y = b$ are satisfied:

$$w(x, y) = \sum_{m=1}^{\infty} w_m(x) \sin \frac{m\pi y}{b} \quad (\text{A. 4})$$

Insertion of Eq. (A. 4) into Eq. (A. 1) results in the following ordinary differential equation for the Fourier function w_m :

$$D_x \frac{d^4 w_m}{dx^4} - 2\alpha_m^2 H \frac{d^2 w_m}{dx^2} + \alpha_m^4 D_y w_m = 0 \quad (\text{A. 5})$$

where

$$\alpha_m = \frac{m\pi}{b}, \quad m = 1, 2, 3, \dots$$

Since Eq. (A. 5) is linear and has constant coefficients, we substitute

$$w_m = \sum_{i=1}^4 A_{im} e^{\lambda_i x} \quad (\text{A. 6})$$

into Eq. (A. 5) and obtain a characteristic equation of the fourth order. The roots of this characteristic equation are

$$\lambda_{1, 2, 3, 4} = \pm \alpha_m \left(\frac{H}{D_x} \pm \sqrt{\frac{H^2}{D_x^2} - \frac{D_y}{D_x}} \right)^{1/2} \quad (\text{A. 7})$$

The form of the solution of Eq. (A. 5) depends on the value of $H^2 - D_x D_y$, and thus the following three cases are possible:

Case 1 $H^2 > D_x D_y$

$$w_m = A_{1m} \cosh(\lambda_1 x) + A_{2m} \sinh(\lambda_1 x) + A_{3m} \cosh(\lambda_2 x) + A_{4m} \sinh(\lambda_2 x) \quad (A. 8)$$

where

$$\lambda_1 = \alpha_m \left(\frac{H}{D_x} + \sqrt{\frac{H^2}{D_x^2} - \frac{D_y}{D_x}} \right)^{1/2}$$

$$\lambda_2 = \alpha_m \left(\frac{H}{D_x} - \sqrt{\frac{H^2}{D_x^2} - \frac{D_y}{D_x}} \right)^{1/2}$$

Case 2 $H^2 = D_x D_y$

$$w_m = (A_{1m} + A_{2m} x) \cosh(\lambda x) + (A_{3m} + A_{4m} x) \sinh(\lambda x) \quad (A. 9)$$

where

$$\lambda = \alpha_m \sqrt{\frac{H}{D_x}}$$

Case 3 $H^2 < D_x D_y$

$$w_m = [A_{1m} \cos(dx) + A_{2m} \sin(dx)] \cosh(ex) + [A_{3m} \cos(dx) + A_{4m} \sin(dx)] \sinh(ex) \quad (A. 10)$$

where

$$c = \frac{\alpha_m}{\sqrt{2}} \sqrt{\beta^2 + \frac{H}{D_x}}$$

$$d = \frac{\alpha_m}{\sqrt{2}} \sqrt{\beta^2 - \frac{H}{D_x}}$$

$$\beta^2 = (D_y/D_x)^{1/2}$$

The integration constants A_{1m} to A_{4m} are to be determined such that the boundary conditions [Eqs. (A.3c) through (A.3e)] at the edges $x = 0$ and $x = a$ are satisfied. This yields

$$(A_m) = [R_m]^{-1} (P_m) \quad (A.11)$$

where

$$(A_m) = \begin{pmatrix} A_{1m} \\ A_{2m} \\ A_{3m} \\ A_{4m} \end{pmatrix}$$

$$(P_m) = \begin{pmatrix} 0 \\ 0 \\ M_{xm} \\ V_{xm} \end{pmatrix}$$

$$[R_m] = \begin{bmatrix} \cosh(\lambda_1 a) & \sinh(\lambda_1 a) & \cosh(\lambda_2 a) & \sinh(\lambda_2 a) \\ \lambda_1^2 \cosh(\lambda_1 a) & \lambda_1^2 \sinh(\lambda_1 a) & \lambda_2^2 \cosh(\lambda_2 a) & \lambda_2^2 \sinh(\lambda_2 a) \\ -D_x \lambda_1^2 + D_1 \alpha_m^2 & 0 & -D_x \lambda_2^2 + D_1 \alpha_m^2 & 0 \\ 0 & -D_x \lambda_1^3 + k \lambda_1 \alpha_m^2 & 0 & -D_x \lambda_2^3 + k \lambda_2 \alpha_m^2 \end{bmatrix}$$

for $H^2 > D_x D_y$, or

$$[R_m] = \begin{bmatrix} \cosh(\lambda a) & a \cosh(\lambda a) & \sinh(\lambda a) & a \sinh(\lambda a) \\ \lambda^2 \cosh(\lambda a) & \lambda^2 a \cosh(\lambda a) + 2\lambda \sinh(\lambda a) & \lambda^2 \sinh(\lambda a) & \lambda^2 a \sinh(\lambda a) + 2\lambda \cosh(\lambda a) \\ -D_x \lambda^2 + D_1 \alpha_m^2 & 0 & 0 & -D_x (2\lambda) \\ 0 & -D_x (3\lambda^2) + k \alpha_m^2 & -D_x \lambda^3 + k \lambda \alpha_m^2 & 0 \end{bmatrix}$$

for $H^2 = D_x D_y$, or

$$[R_m] = \begin{bmatrix} \cos(da) \cosh(ca) & \sin(da) \cosh(ca) & \cos(da) \sinh(ca) & \sin(da) \sinh(ca) \\ (c^2 - d^2) \cos(da) \cosh(ca) & (c^2 - d^2) \sin(da) \cosh(ca) & (c^2 - d^2) \cos(da) \sinh(ca) & (c^2 - d^2) \sin(da) \sinh(ca) \\ -2cd \sin(da) \sinh(ca) & -2cd \cos(da) \sinh(ca) & -2cd \sin(da) \cosh(ca) & -2cd \cos(da) \cosh(ca) \\ -D_x (c^2 - d^2) + D_1 \alpha_m^2 & 0 & 0 & -D_x (2cd) \\ 0 & -D_x d(3c^2 - d^2) + k d \alpha_m^2 & -D_x c(c^2 - 3d^2) + k c \alpha_m^2 & 0 \end{bmatrix}$$

for $H_2 < D_x D_y$

and

$$k = H + 2 D_{xy}$$

Through use of Eqs. (A. 11), (A. 8), (A. 9), or (A. 10), the edge deflection $w_m(0)$ and edge rotation $(d/dx)w_m(0)$ are next computed for the following values of M_m and V_m :

Case I $M_m = 0, V_m = 1$

Case II $M_m = 1, V_m = 0$

Finally, for each harmonic m of edge loading, the influence coefficients c_{ij} are obtained as follows:

From Case I

$$c_{11} = w_m(0) \quad (A. 12a)$$

$$c_{21} = \frac{d}{dx} w_m(0) \quad (A. 12b)$$

From Case II

$$c_{12} = w_m(0) \quad (A. 13a)$$

$$c_{22} = \frac{d}{dx} w_m(0) \quad (A. 13b)$$

This procedure has been programmed for the IBM 7094 digital computer and the influence coefficients so obtained were used to arrive at the results presented in Sec. 3.

Appendix B

DIGITAL COMPUTER PROGRAM

The solution presented in Sec. 2 was programmed for the digital computer and a listing of the FORTRAN II program is given in this appendix. Use of the program requires input of geometrical and elastic properties of the cylindrical panel, estimated ranges of critical load, number of waves in the axial direction, and a set of control variables. In the general case, such properties as moments of inertia of stiffeners are required in the input, but a separate branch is provided in the program for use in the case of rectangular stiffeners. If this branch is used the stiffener dimensions are read in and stiffnesses are computed internally.

The program computes the determinant of the stability matrix and can give the output in two different ways. A series of values of the determinant may be printed out together with the corresponding load levels. Thus, the user can find the critical load as the lowest load for which the determinant vanishes. The computer can also determine the value of the critical load by use of an iterative procedure.

At certain values of the load a change may occur in the number of real roots of the characteristic equation. Multiple roots exist at such values and hence two columns in the stability matrix will be identical. As this leads to a zero determinant independently of the boundary conditions at the straight edges, the corresponding value of the load is a false solution to the buckling problem. When a series of determinant values are printed, such a solution may be recognized in the output since the number of real roots is included. In case the iterative method is used to find the critical load, the program will ignore zero crossings of the determinant when a change takes place in the number of real roots within 10^{-4} times the load step. Printing out the values of the roots of the characteristic equation is an additional option.

The buckling mode, which is symmetrical around the line of intersection between adjacent panels, is expected to be less critical than the antimetrical mode even for the case without web plates. In the presence of web plates, their stiffness will practically prevent the in-plane displacements and therefore the antimetric mode becomes critical. For completeness, the branch for the symmetrical mode is included in the computer program. However, to avoid unnecessary input, this program branch assumes a negligible effect for the web plate, and a lower bound to the symmetrical buckling mode can be computed. For the antimetric mode, the case of zero web plate stiffness corresponds to infinite flexibility constants (c_{11} , c_{12} , c_{21} , and c_{22}). For other cases the coefficients may be computed according to the solution given in Appendix A.

Figure B-1 shows a flow chart of the computer program. Geometric and elastic properties needed as input are listed in Table B-1. Load range and load step size are defined by the input parameters listed in Table B-2. In addition some control variables are input parameters. They are listed in Table B-3. The complete sequence of input parameters and the corresponding "read-in-formats" are shown in Table B-4.

A sample problem is demonstrated in Table B-5 (input) and Table B-6 (output). The coefficients c_{11} , c_{12} , c_{21} , and c_{22} are computed according to the solution in Appendix A.

The critical load is given under the heading NXO in the output (Table B-6). This buckling load is computed for an internally waffle stiffened panel as shown in Fig. B-2.

The complete program is given in Table B-7.

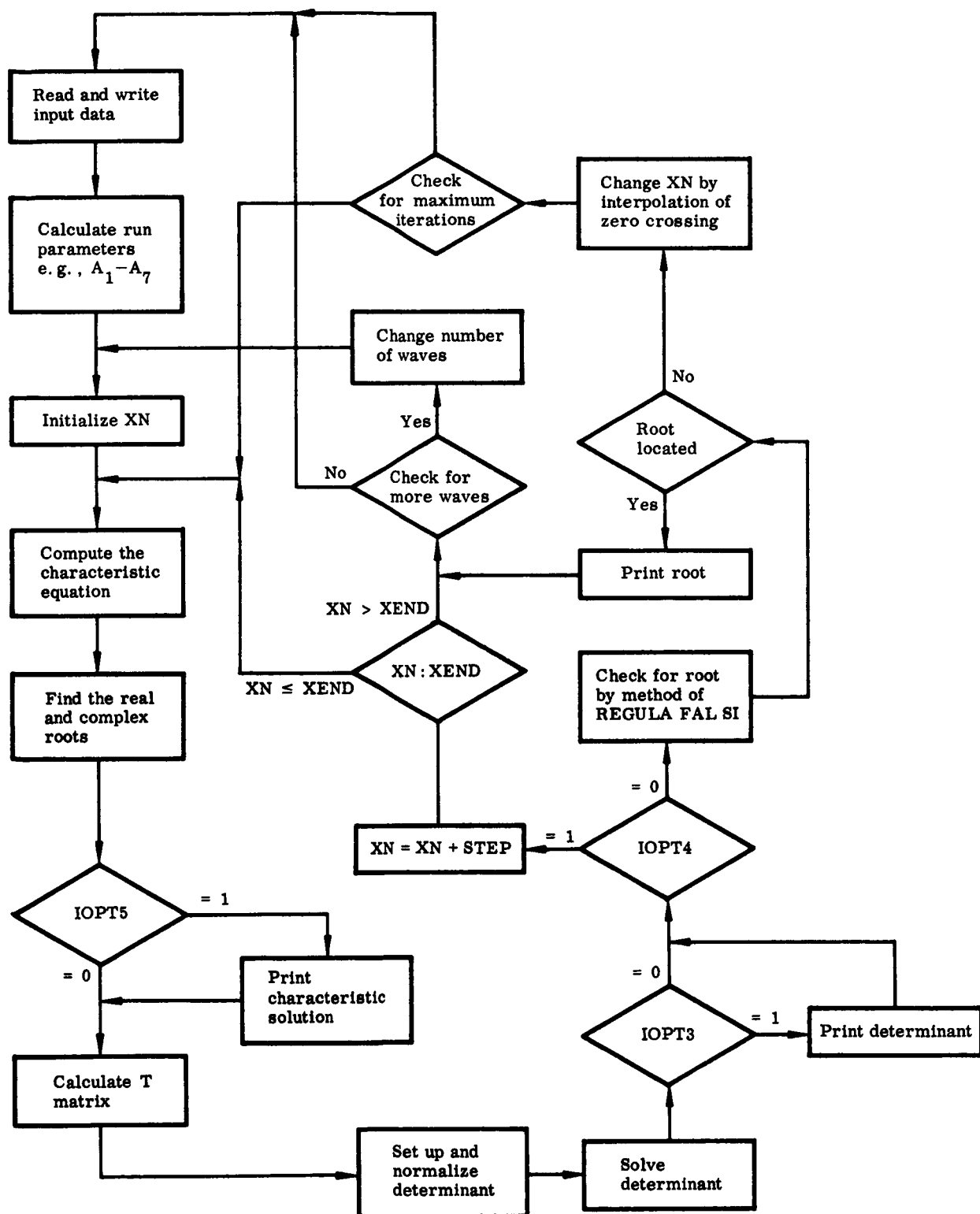


Fig. B-1 Flow Chart

Table B-1

GEOMETRIC AND ELASTIC PROPERTIES USED AS INPUT

	Computer Program Notation	Notation in Analysis (Sec. 2)	Item
For Rectangular Stiffeners	ZNU	ν	Poisson's ratio
	AL	L	Panel length
	AR	R	Panel radius
	AE	E	Young's modulus of skin
	AT	t	Thickness of skin
	DIS		Thickness of rectangular stringers
	D2		Thickness of rectangular rings
	H1		Height of rectangular stringers
	H2		Height of rectangular rings
	AAS		Spacing of rectangular stringers
	AB		Spacing of rectangular rings
	W	Ω	Angle subtending width of panel
	PHI	ϕ	Cell angle
For Arbitrary Stiffeners	XA	s	Spacing of arbitrary stringers
	XB	r	Spacing of arbitrary rings
	XE1	$\hat{e}_1$	Eccentricity of arbitrary stringers
	XE2	$\hat{e}_2$	Eccentricity of arbitrary rings
	XA1	A ₁	Area of arbitrary stringers
	XA2	A ₂	Area of arbitrary rings
	XI1	$\bar{I}_1$	Moment of inertia of arbitrary stringers
	XI2	$\bar{I}_2$	Moment of inertia of arbitrary rings
	XIT1	J ₁	Torsional constant for arbitrary stringers
	XIT2	J ₂	Torsional constant for arbitrary rings
	E1	E ₁	Young's modulus of arbitrary stringers
	E2	E ₂	Young's modulus of arbitrary rings
	CC ₁₁ , CC ₁₂ , CC ₂₁ , CC ₂₂	c ₁₁ , c ₁₂ , c ₂₁ , c ₂₂	Influence coefficients for web plate (included only if KC = 1)
	BN	n	Number of axial waves

Table B-2

LOAD RANGE AND STEP SIZE

Computer Program Notation	Item
XBEG	Starting value for axial load (compression corresponds to negative values). The numerical value should be chosen below the expected solution.
XEND	Maximum value of axial load. (The search for a solution will not be extended beyond this point.)
STEP	Increment of axial load used in search for solution.

Table B-3

CONTROL VARIABLES AND INPUT PARAMETERS

Computer Program Notation	Item
KC	$\begin{cases} = 0 & \text{for symmetrical mode} \\ = 1 & \text{for antisymmetrical mode} \end{cases}$
KN	Number of elements in the wave-number vector
LIM	Maximum number of iterations in the search for zero-crossing of determinant
ICARD1	$\begin{cases} = 0 & \text{The "option card" (Card No. 3) will not be included.} \\ & \text{Information for previous case applies} \\ = 1 & \text{Card No. 3 will be included} \end{cases}$
ICARD2	$\begin{cases} = 0 & \text{The "geometry card" (Card No. 4) not included} \\ = 1 & \text{(Card No. 4) will be included} \end{cases}$
ICARD3	$\begin{cases} = 0 & \text{The "\alpha-card" (Card No. 5) is not included} \\ = 1 & \text{(Card No. 5) will be included} \end{cases}$
ICARD4	$\begin{cases} = 0 & \text{The "load-card" (Card No. 6) is not included} \\ = 1 & \text{(Card No. 6) will be included} \end{cases}$
IOPT3	$\begin{cases} = 0 & \text{The determinant is not printed} \\ = 1 & \text{The determinant is printed} \end{cases}$
IOPT4	$\begin{cases} = 0 & \text{Regular falsi; search is used} \\ = 1 & \text{Determinant computed for several points} \end{cases}$

Table B-3 (cont'd)

Computer Program Notation	Item
IOPT5	$\begin{cases} = 0 & \text{Roots of characteristic equation are printed} \\ = 1 & \text{Roots of characteristic equation are not printed} \end{cases}$
IOPT7	$\begin{cases} = 0 & \text{External stiffeners} \\ = 1 & \text{Internal stiffeners} \end{cases}$
IOPT11	$\begin{cases} = 0 & \text{Stiffeners with rectangular cross section} \\ = 1 & \text{Stiffeners with any cross section} \end{cases}$

Table B-4

INPUT PARAMETER SEQUENCE AND CORRESPONDING READ-INFORMATS

Card	FORTTRAN Symbol	Format
1	Record	72H
2	ICARD1, ICARD2, ICARD3, ICARD4, KN, KC, LIM	4I1, 2I3, I4
3	BN(I)	6E12.8
4 ^(a)	IOPT3, IOPT4, IOPT5, IOPT7, IOPT11	20I1
5 ^(b)	ZNU, AL, AR, AE, AT	6E12.8
6a ^(c)	DIS, D2, H1, H2, AAS, AB, E1, E2, W	6E12.8
6b ^(d)	XA, XB, XE1, XE2, XA1, XA2, XI1, XI2, XIT1, XIT2, E1, E	6E12.8
7 ^(e)	XBEG, XEND, STEP	6E12.8
8a ^(f)	PHI	6E12.8
8b ^(g)	PHI, C ₁₁ , C ₁₂ , C ₂₁ , C ₂₂	6E12.8

- (a) Omitted unless ICARD1 = 1.
 (b) Omitted unless ICARD2 = 1.
 (c) Omitted unless ICARD3 = 1, for rectangular stiffeners.
 (d) Omitted unless ICARD3 = 1, for arbitrary stiffeners.
 (e) Omitted unless ICARD4 = 1.
 (f) For rigid webplate.
 (g) For the case without webplate for symmetrical mode or with finite stiffness webplate for antimetrical mode.

Table B-5

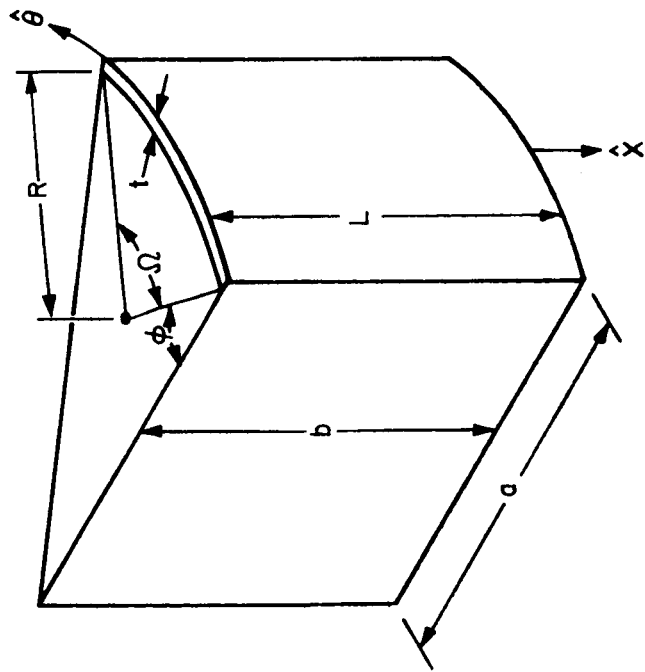
INPUT DATA IN REQUIRED FORMAT

MODULE		ORNL	ALDA	FAC.	JOB NUMBER	DATE OF REQUEST	DISPATCH NUMBER	PROGRAM	OPER. CONTROL NO.	DATE REQUEST	PRIORITY	PAGE	
PROGRAMMER		ORNL	ORNL	ORNL	ORNL	ORNL	ORNL	ORNL	ORNL	ORNL	ORNL	OF	OF
1	1	OCTOBER	1	1965	1								
1	1	1	300										
4			+01										
0	0	1	0										
3			+00	51477		+02	35598		+08	147			
2	2	1	+00	221		+00	603		+00	6821			
1			+08	1		+08	11781						+01
-	8		+04	-12		+05	-25						
5	2	3	6		25842557	-02	43124422	-03	43124422	-03	15400285	-03	

Table B-7

EXAMPLE: CYLINDRICAL PANEL WITH INTERNAL WAFFLE

D1= 0.221000E-00, D2= 0.221000E-00, H1= 0.603000E 00, H2= 0.603000E 00, A= 0.682100E 01, B= 0.682100E 01
INTERNAL STIFFENERS.
PARAMETER RANGE, -0.80000E 04 TO -0.12000E 05, STEP= -0.25000E 03
ANTIMETRIC MODE PW1= 5.235999E-01 C11= 2.5842557E-03 C12= 4.3124421E-04 C21= 4.3124421E-04 C22= 1.5400283E-04
ZNU= 0.306900E-00, L= 0.514770E 02, R= 0.355960E 02, N= 0.400000E 01, E= 0.100000E 08, T= 0.147000E-00
OMEGA = 0.117810E 01 E1 = 0.100000E 08 E2 = 0.100000E 08
ALPHA1 = ALPHA7 ARE
-0.127486E-02 -0.127486E-02 0.112094E 01 0.112094E 01 0.177343E-04 0.177343E-04 0.403729E-05
4.000000E 00 -5.9121157E 03 1.2453256E-26+ 60 DETP -1.8650651E-32+ 60 DETI



Panel	Web Plate
L = 51.477 in.	a = 90.765759 in.
R = 35.598 in.	b = L = 51.477 in.
$\Omega = 3\pi/8$	t = 0.147 in.
$\phi = 0.52360$ radian	
t = 0.147 in.	
$d_1 = d_2 = 0.221$ in.	$d_1 = d_2 = 0.221$ in.
$h_1 = h_2 = 0.603$ in.	$h_1 = h_2 = 0.603$ in.
s = r = 6.821 in.	$s_1 = s_2 = 6.821$ in.

Fig. B-2 Cylindrical Panel (Stiffeners not shown)

Table B-7

STUDY OF JUNCTURE STRESS FIELDS: STIFFENED PANEL UNDER UNIFORM AXIAL COMPRESSION

```

C CONTRACT NAS 8-11480 TO NASA G.C. MARSHALL SPACE FLIGHT CENTER
C HUNTSVILLE, ALABAMA
C BY SOLID MECHANICS, AEROSPACE SCIENCES LABORATORY, 52-20
C LOCKHEED MISSILES AND SPACE COMPANY, PALO ALTO CALIF
C STIFFENED PANEL, UNIFORM AXIAL COMPRESSION, REGULA FALSI
C DIMENSION RN(100)
C DIMENSION C(9)
C DIMENSION T(8,8),X(8),RU(8),RV(8),D1(1),DET(1),C0(5),A1SW(4)
C COMMON KEXP
C COMMON IOPT3,IOPT4,IOPT5
C IOPT3=0 DO NOT PRINT DET
C IOPT3=1 PRINT DET
C IOPT4=0 REGULA FALSI SEARCH FOR DET=0
C IOPT4=1 COMPUTE DET FOR SEVERAL POINTS
C IOPT5=0 DO NOT PRINT CHARACTERISTIC SOLUTION
C IOPT5=1 PRINT CHARACTERISTIC SOLUTION
C IOPT7=0 EXTERNAL
C IOPT7=1 INTERNAL
C IOPT11=0 RECTANGULAR CROSS-SECTION.
C IOPT11=1 ARBITRARY CROSS-SECTION
C KC=0 SYMMETRIC MODE
C KC=1 ANTIMETRIC MODE
C DELT = .00005
C READ INPUT TAPE 5,1020
C
10 CONTINUE
C READ INPUT TAPE 5,990,ICARD1,ICARD2,ICARD3,ICARD4,KC,LIM
C WRITE OUTPUT TAPE 6,1020
C ICARD1=0 SKIP OPTION CARD
C ICARD2=0 SKIP GEOMETRY CARD
C ICARD3=0 SKIP ALPHA CARD
C ICARD4=0 SKIP STEP SIZE CONTROL
C READ INPUT TAPE 5,1000,(RN(I),I=1,KN)
C KKK=1
C IF (ICARD1) 30,30,20
COMP0010
COMP0020
COMP0030
COMP0040
COMP0050
COMP0060
COMP0070
COMP0080
COMP0090
COMP0100
COMP0110
COMP0120
COMP0130
COMP0140
COMP0150
COMP0160
COMP0170
COMP0180
COMP0190
COMP0200
COMP0210
COMP0220
COMP0230
COMP0240
COMP0250
COMP0260
COMP0270
COMP0280
COMP0290
COMP0300
COMP0310
COMP0320
COMP0330
COMP0340

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```

20 CONTINUE
  READ INPUT TAPE 5,1010,I0PT3,I0PT4,I0PT5,I0PT7,I0PT11
30 IF (ICARD2) 50,50,40
40 CONTINUE
  READ INPUT TAPE 5,1000,ZNU,AL,AR,AE,AT
  Y1=(1.-ZNU)/2.
  Y2=(1.+ZNU)/2.
  W1=(1.-ZNU**2)/2.
  W2=(AT/AR)**2/12.
  W4=(1.-ZNU**2)
  ARL = 3.14159265*AR/AL
  XN = XN1S
50 IF (ICARD3) 80,80,60
60 IF (I0PT11) 70,70,160
70 CONTINUE
  READ INPUT TAPE 5,1000,D1S,D2,H1,H2,AAS,AB,E1,E2,W
  WRITE OUTPUT TAPE 6,1110,D1S,D2,H1,H2,AAS,AB
80 IF (I0PT11) 90,90,170
90 CONTINUE
  H1R=H1/AR
  H2R=H2/AR
  H1I=H1/AT
  H2I=H2/AT
  D1R = D1S*E1/(AB*AE)
  D2A = D2*E2/(AAS*AE)
  A1=W1*D1R*H1R*(H1I+1.)
  A2=W1*D2A*H2R*(H2I+1.)
  A3=1.+2.*W1*D1R*H1I
  A4=1.+2.*W1*D2A*H2I
  A5=W2+W1/6.*H1R**2*D1R*(4.*H1I+6.*3./H1I)
  A6=W2+W1/6.*H2R**2*D2A*(4.*H2I+6.*3./H2I)
  IF (D1S-H1) 100,110,110
100 *5=H1I*D1R*(71S/AR)**2
  GO TO 120
110 *5=H1I*D1R*H1R**2
120 IF (D2-H2) 130,140,140
130 *6=H2I*D2A*(72/AR)**2
  GO TO 130
140 *6=H2I*D2A*H2R**2

```

COMP0350
 COMP0360
 COMP0370
 COMP0380
 COMP0390
 COMP0400
 COMP0410
 COMP0420
 COMP0430
 COMP0440
 COMP0450
 COMP0460
 COMP0470
 COMP0480
 COMP0490
 COMP0500
 COMP0510
 COMP0520
 COMP0530
 COMP0540
 COMP0550
 COMP0560
 COMP0570
 COMP0580
 COMP0590
 COMP0600
 COMP0610
 COMP0620
 COMP0630
 COMP0640
 COMP0650
 COMP0660
 COMP0670
 COMP0680
 COMP0690
 COMP0700
 COMP0710
 COMP0720
 COMP0730

```

150 A7=2.*W2+Y1/3.*(W5+W6)
160 GO TO 180
160 CONTINUE
160 READ INPUT TAPE 5,1000,XA,XB,XE1,XE2,XA1,XA2,XI1,XI2,XIT1,XIT2,E1,
162,W
170 WRITE OUTPUT TAPE 6,1120,XA,XB,XE1,XE2,XA1,XA2,XI1,XI2,XIT1,XIT2
170 CONTINUE
XAI=XA*AI
RT=XB*AT
E1E = E1/AE
E2E = E2/AE
A1E=XA1*XE1
A2E=XA2*XE2
A1=W4*A1E/(AR*BT) * E1E
A2=W4*A2E/(AR*XAT) * E2E
A3=1.+W4*XAI/BT * E1E
A4=1.+W4*XA2/XAT * E2E
A5=W2+W4*(XI1+A1E*XE1)/(BT*AR**2) * E1E
A6=W2+W4*(XI2+A2E*XE2)/(XAT*AR**2) * E2E
A7=2.*W2+Y1/(AT*AR**2)*(XIT1/XR*E1E + XIT2/XA*E2E)
GO TO 200
180 IF (IOPT7) 200,200,190
190 A1=-A1
A2=-A2
WRITE OUTPUT TAPE 6,1030
200 CONTINUE
210 IF (ICARD4) 270,270,220
220 CONTINUE
READ INPUT TAPE 5,1000,XBEG,XEND,STEP
WRITE OUTPUT TAPE 6,1100,XREG,XEND,STEP
READ INPUT TAPE 5,1000,PHI,CC11,CC12,CC21,CC22
IF (KC) 230,230,240
230 WRITE OUTPUT TAPE 6,1130,PHI
GO TO 250
240 WRITE OUTPUT TAPE 6,1140,PHI,CC11,CC12,CC21,CC22
XBOT = (CC11*CC22-CC12*CC21)*2.
AXBOT = ABSF(XBOT)
CS = ABSF(CC22)/AXBOT
C4 = -ABSF(CC21)/AXBOT
COMP0740
COMP0750
COMP0760
COMP0770
COMP0780
COMP0790
COMP0800
COMP0810
COMP0820
COMP0830
COMP0840
COMP0850
COMP0860
COMP0870
COMP0880
COMP0890
COMP0900
COMP0910
COMP0920
COMP0930
COMP0940
COMP0950
COMP0960
COMP0970
COMP0980
COMP0990
COMP1000
COMP1010
COMP1020
COMP1030
COMP1040
COMP1050
COMP1060
COMP1070
COMP1080
COMP1090
COMP1100
COMP1110
COMP1120

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```

C5 = C4
C6 = ABSF(CC11)/AXBOT
C7=0.
C8=0.
250 FUDGE=2.*W1/(AE*AT)
XBEG=XBEG*FUDGE
XEND=XEND*FUDGE
STEP=STEP*FUDGE
260 CONTINUE
STEPS=STEP
XN1S=XN
270 CONTINUE
XN =XBEG
AN=BN(KKK)
WRITE OUTPUT TAPE 6,1070,ZNU,AL,AR,AN,AE,AT
WRITE OUTPUT TAPE 6,1080,W,E1,F2
WRITE OUTPUT TAPE 6,1090,A1,A2,A3,A4,A5,A6,A7
280 CONTINUE
IOUT = 15
290 CONTINUE
IF (IOUT) 300,300,320
300 CONTINUE
310 WRITE OUTPUT TAPE 6,1060
320 CONTINUE
STEP=STEPS
Y4=A3*A4-ZNU
Y6=12.
ANV = AN
AN = AN*ARL
AN2=AN**2
AN3=AN**3
AN4=AN2**2
AN5=AN2*AN3
AN6=AN3*AN3
AN8=AN4*AN4
ALIM=XEND
DEI=0.
DET(2)=0.
STEPS=STEP

```

```

COMP1130
COMP1140
COMP1150
COMP1160
COMP1170
COMP1180
COMP1190
COMP1200
COMP1210
COMP1220
COMP1230
COMP1240
COMP1250
COMP1260
COMP1270
COMP1280
COMP1290
COMP1300
COMP1310
COMP1320
COMP1330
COMP1340
COMP1350
COMP1360
COMP1370
COMP1380
COMP1390
COMP1400
COMP1410
COMP1420
COMP1430
COMP1440
COMP1450
COMP1460
COMP1470
COMP1480
COMP1490
COMP1500
COMP1510

```

COMP1520
COMP1530
COMP1540
COMP1550
COMP1560
COMP1570
COMP1580
COMP1590
COMP1600
COMP1610
COMP1620
COMP1630
COMP1640
COMP1650
COMP1660
COMP1670
COMP1680
COMP1690
COMP1700
COMP1710
COMP1720
COMP1730
COMP1740
COMP1750
COMP1760
COMP1770
COMP1780
COMP1790
COMP1800
COMP1810
COMP1820
COMP1830
COMP1840
COMP1850
COMP1860
COMP1870
COMP1880
COMP1890
COMP1900

```

K=0
K1=0
KOUNT=0
TH=(W-PHI)/2.
CTH=COSF(TH)
STH=SINF(TH)
AA = XN
C
  COMPUTE THE CHARACTERISTIC EQUATION
330 CONTINUE
  C(5) = Y1 * (-A3*A5+A1**2)*AN8+ (1.-ZNU)*ZNU*A1*AN6
  1 +Y1*(ZNU**2-A3*A4)*AN4 - XN*Y1*A3*AN6
  C(4) =(A5*Y4 + Y1*A3*A7 - A1**2*A4)*AN6 + ((1.-ZNU)*A2*A3
  1 +2.* Y1*A1*A4)*AN4
  C(3)=Y1*2.*ZNU*A2*AN2 - (1.+ZNU)* A1*A2*AN4
  1 - XN*Y1*A4*AN2 - (1.+ZNU)* A1*A2*AN4
  C(2) = (A6*Y4 + Y1*A4*A7 - A2**2*A3)*AN2
  C(1) = Y1*(A2**2 - A4*A6)
  DO 340 I=10,14
340 C(I)=0.
  CALL MULLER (C, 4, X, X(9))
  IF (IUP15) 380,380,350
350 WRITE OUTPUT TAPE 6,1040,(X(I),X(I+8),I=1,4)
  WRITE OUTPUT TAPE 6,980,(C(I),I=1,5)
  DO 360 I=1,5
360 C(I) = C(I)
  DO 370 I=1,4
  ANSW(I) = (0., 0.)
  DO 370 J=1,5
  1 ANSW(I) = X(I)*ANSW(I) + C(J)
  1 370 ANSW(I) = X(I)*ANSW(I),ANSW(I+4),I=1,4)
  WRITE OUTPUT TAPE 6,980,(ANSW(I),ANSW(I+4),I=1,4)
380 CONTINUE
  N12 = N11
  N11 = 0
  DO 410 I=1,4
  I1=5-I
  IF (ABSF(X(I1))-1.E6*ABSF(X(I1+8))) 400,390,390
390 N11 = N11+1
  1 400 Z = SUMF(X(I1))
  1 X (2*I1-1) = Z

```

```

1      X (2*11) = -Z
410    CONTINUE
DO 420 I=1,8
Z=X(I)
Z2=Z*Z
Z3=Z*Z2
Z4=Z2**2
Z5=Z*Z4
T20 = Y1*(A3*AN4 + A4*Z4) - Y4*AN2*Z2
RU(I) = (Y1*(A1*AN5 + ZNU*AN3) + A4*AN*Z2*(Y1-A1*AN2)
1      - Y2*A2*AN*Z4) / T20
RV(I) = -((Y2*((1.,0.)*A1*AN2) - A3*A4)*AN2*Z + (Y1*A4 + A2*A3
1      *AN2)*Z3 - Y1*A2*Z5) / T20
420    CONTINUE
GNA = - W2*ZNU*A7
DO 450 I=1,8
ELI = EXPF(.5,0.) * X(I) * W
PLI = (1.,0.)/ELI
OLB = (1.,0.) + X(I)*RV(I)
C11=A4*OLB=ZNU*AN*RU(I)-A2*X(I)**2
C22=A6*X(I)**2-AN2*ZNU*W2-A2*OLB
C11=C11*AE*AT /W4
C22=C22*AE*AT*AR /W4
C12=AE*AT*(A6*X(I)**3-A2*(RV(I)*X(I)**2+X(I))-(A7-ZNU*W2)*AN**2
1      *X(I)) /W4
C21=C12*CTH-C11*STH
C11=C11*CTH+C12*STH
IF (KC) 430,430,440
SYMMETRIC MODE
430    T(1,I)=ELI*X(I)
T(2,I)=RLI*X(I)
T(3,I)=ELI*(RV(I)*CTH+STH)
T(4,I)=RLI*(RV(I)*CTH-STH)
C15=BU(I)*AR*W4
C25= AE*AT*(RV(I)*AN+RU(I)*X(I))
T(5,I)=ELI*(C15*C7+C25)
T(6,I)=RLI*(C15*C7-C25)
T(7,I)=ELI*(CTH-RV(I)*STH)*C8-C21/AR)
T(8,I)=RLI*(CTH+RV(I)*STH)*C8-C21/AR)

```

COMP1910
 COMP1920
 COMP1930
 COMP1940
 COMP1950
 COMP1960
 COMP1970
 COMP1980
 COMP1990
 COMP2000
 COMP2010
 COMP2020
 COMP2030
 COMP2040
 COMP2050
 COMP2060
 COMP2070
 COMP2080
 COMP2090
 COMP2100
 COMP2110
 COMP2120
 COMP2130
 COMP2140
 COMP2150
 COMP2160
 COMP2170
 COMP2180
 COMP2190
 COMP2200
 COMP2210
 COMP2220
 COMP2230
 COMP2240
 COMP2250
 COMP2260
 COMP2270
 COMP2280
 COMP2290

```

C
I 440 GO TO 450
      ANTIMETRIC MODE
I      T(3,I)=ELI*(CTH-RV(I)*STH)
I      T(4,I)=RLI*(CTH+RV(I)*STH)
I      T(1,I)=BU(I)*ELI
I      T(2,I)=BU(I)*RLI
I      T(5,I)=ELI*(C11+C3*(RV(I)*CTH+STH)*AR-C4*X(I))
I      T(6,I)=RLI*(C11-C3*(RV(I)*CTH-STH)*AR+C4*X(I))
I      T(7,I)=ELI*(C22-C5*(RV(I)*CTH+STH)*AR+C6*X(I))
I      T(8,I)=RLI*(C22+C5*(RV(I)*CTH-STH)*AR-C6*X(I))
I 450 CONTINUE
      DDET=1.
      KEXP=0
      DO 540 I=1,8
        CALL NORM (T(1,I),T(1,I+8),DET2)
        DD=DET2
        460 IF (DD-1.E+10) 480,470,470
        470 DD=DD*1.E-10
        KEXP=KEXP+10
        GO TO 460
        480 IF (DD-1.E=10) 490,490,500
        490 DD=DD*1.E+10
        KEXP=KEXP-10
        GO TO 480
        500 CONTINUE
        510 DDET=DDET*DD
        520 IF (DDET-1.E+10) 540,530,530
        530 DDET=DDET*1.E-10
        KEXP=KEXP+10
        GO TO 520
        540 CONTINUE
I      D1=DET
        CALL PLXDET (T, T(1,9), 8, DET, DET(2))
        DET=DET*DDET
        DET(2)=DET(2)*DDET
        IF (IOPT3) 550,620,550
        550 IF (IOUT-50) 560,560,590
        560 IOUT=IOUT+1
        IF (IOPT5) 580,580,570

```

COMP2300
 COMP2310
 COMP2320
 COMP2330
 COMP2340
 COMP2350
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 COMP2660
 COMP2670
 COMP2680

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570 CONTINUE
    WRITE OUTPUT TAPE 6,1060
    IOUT=IOUT+3
580 CONTINUE
    GO TO 610
590 IOUT=3
    WRITE OUTPUT TAPE 6,1020
600 WRITE OUTPUT TAPE 6,1060
610 CONTINUE
    XNX = XN/FUDGE
    WRITE OUTPUT TAPE 6,1050,ANN,XNX,DET,KEXP,DET(2),KEXP,N11
620 IF (IOUT) 630,640,630
630 XN = XN + STEP
    IF ((XN-XEND)*STEP) 330,330,960
640 KK=KEXP1-KEXP
    AAZ=AA
    IF (KK) 650,670,650
650 IF (D1) 660,670,660
660 XD=10.**KK
    D1=D1*XD
    D1(2)=D1(2)*XD
    FOLD=FOLD*XD
    FTWO=FTWO*XD
    KEXP1=KEXP
670 IF (K) 650,680,860
680 IF (ABSF(DEF)-ABSF(DET(2))) 690,720,720
690 IF (D1(2)/DET(2)) 700,780,780
700 IF (N11-N12) 740,710,740
    C      IGNORE CROSSING IF NUMBER OF COMPLEX ROOTS HAS CHANGED
    C      ZERO CROSSING OF IMAGINARY PART
710 K=-1
    FOLD=D1(2)
    FNEW=DET(2)
    GO TO 900
720 IF (D1/DET) 730,780,780
730 IF (N11-N12) 740,770,740
    C      CHANGE STEPSIZE
740 COUNT=COUNT+1
    IF (COUNT-3) 750,750,760

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COMP2690
COMP2700
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COMP2800
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COMP2980
COMP2990
COMP3000
COMP3010
COMP3020
COMP3030
COMP3040
COMP3050
COMP3060
COMP3070

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750 AA=AA1
STEP=STEP/10.
I DET=D1
N11=N12
GO TO 780
760 KOUNT=0
STEP=STEPS
GO TO 780
C ZERO CROSSING OF REAL PART
770 K=1
FOLD=D1
FNEW=DET
GO TO 900
780 AA1=AA
AA=AA+STEP
KEXP1=KEXP
K1=K1+1
IF (K1-LIM) 800,800,790
790 WRITE OUTPUT TAPE 6,1150
GO TO 960
800 IF ((AA-ALIM)*STEP) 820,820,810
810 WRITE OUTPUT TAPE 6,1160
GO TO 960
820 CONTINUE
830 XN = AA
GO TO 350
840 WRITE OUTPUT TAPE 6,1170
GO TO 960
C SEARCH FOR DET=0
850 K=K-1
FNEW=DET(2)
GO TO 870
860 K=K+1
FNEW=DET
870 IF (FOLD-FNEW) 880,940,930
880 CONTINUE
890 IF (XARSF(K)-LIM) 900,900,840
900 STEP1= (AA-AA1)*FOLD/(FOLD-FNEW)
AA2=AA

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COMP3080
COMP3090
COMP3100
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COMP3120
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COMP3460

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AA=AA1+STEP1
FTWO=FNEW
IF (ABSF((AAZ-AA)/AA)-DELT) 940,940,910
910 IF (XMODF(ABSF(K),3)) 920,920,830
920 AA=(AA1+AA2)*.5
GO TO 830
930 STEP1=AA2-AA
C CROSSING BETWEEN AA AND AA2
AA1=AA
AA=AA2
FOLD=FNEW
FNEW=FTWO
KEXP1=KEXP
GO TO 880
C ROOT LOCATED
940 ROOT= AA1+(AA-AA1)*FOLD/(FOLD-FNEW)
XNX = ROOT/FUDGE
950 WRITE OUTPUT TAPE 6,1050,ANN,XNX,DET,KEXP,DET(2),KEXP
960 CONTINUE
IF (KKK-KN) 970,10,10
970 KKK=KKK+1
GO TO 270
980 FORMAT (1H0, 8E15.8)
990 FORMAT (4I1,2I3,I4)
1000 FORMAT (6E12.8)
1010 FORMAT (20I1)
1020 FORMAT (72H1
1
1030 FORMAT (30H0 INTERNAL STIFFENERS,
1040 FORMAT (41H0 SOLUTION OF THE CHARACTERISTIC EQUATION / (1P8E15.7))
1050 FORMAT (1P3E20.7, 1H+, I3, 1PE20.7, 1H+, I4, I20)
1060 FORMAT ( 1H0,12X, 1HN, 18X, 3HNX0,17X, 16X,4HDETR,21X,
1 4HDEII )
1070 FORMAT (6H0 ZNU=E15.6, 6H, L=E15.6, 9H, R=E15.6,4H, N=E15.6
1 , 4H, E=E15.6, 4H, T=E15.6 )
1080 FORMAT (9H0 OMEGA =, E15.6, 7H E1 =, E15.6, 7H E2 =, E15.6)
1090 FORMAT (24H0 ALPHA1 - ALPHA7 ARE, /7E15.6 )
1100 FORMAT (19H0 PARAMETER RANGE, F14.5,5H TO E14.5,7H, STEP=E14.5
1 )
COMP3470
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COMP3500
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COMP3790
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COMP3830
COMP3840
COMP3850

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1110 FORMAT (5H, D1=
1 5H, H2=E15.6, 5H, A= E15.6, 5H, B= E15.6 )
1120 FORMAT (120H
1A1
2 / 10E12.5 )
1130 FORMAT (15H SYMMETRIC MODE5X5HPHI =PE15,8)
1140 FORMAT(16H ANTIMETRIC MODE2X5HPHI =1PE14,7,2X4HC11=E14.7,2X4HC12=
1 E14.7,2X4HC21=E14.7,2X4HC22=E14.7)
1150 FORMAT ( 50H0 MAXIMUM NUMBER OF STEPS EXCEEDED
1160 FORMAT (50H0 UPPER LOAD LIMIT EXCEEDED
1170 FORMAT (50H0 REGULA FALSI DOES NOT CONVERGE
END
COMP3860
COMP3870
ACOMP3880
COMP3890
COMP3900
COMP3910
COMP3920
COMP3930
COMP3940
COMP3950
COMP3960
COMP3970

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