

THE EARTH'S LONGITUDE GRAVITY FIELD AS SENSED BY THE DRIFT OF THREE SYNCHRONOUS SATELLITES
by C. A. Wagner
Goddard Space Flight Center
Greenbelt, Md.


NASA TN D-3557

# THE EARTH'S LONGITUDE GRAVITY FIELD AS SENSED BY THE DRIFT OF THREE SYNCHRONOUS SATELLITES 

By C. A. Wagner

Goddard Space Flight Center Greenbelt, Md.


#### Abstract

One hundred and fifty-two orbits of three synchronous communications satellites have been analyzed for sensitivity to earth longitude gravity components through third order which are in resonance with them. Eighty-seven orbits are of Syncom 2 with inclination between $32^{\circ}$ and $33^{\circ}$ and distributed (but not uniformly) in mean geographic longitude between $66^{\circ}$ and $305^{\circ}$. Nineteen orbits, distributed with fair uniformity between $173^{\circ}$ and $180^{\circ}$, are of the nearly geostationary Syncom 3. Forty-six orbits are of the nearly geostationary Early Bird satellite between $330^{\circ}$ and $332^{\circ}$ longitude. These orbits were calculated without consideration of resonant gravitational effects.

The orbit data was reduced to give a set of essentially nine well determined long term longitude accelerations for these satellites between $66^{\circ}$ and $332^{\circ}$. From this reduced acceleration record, after extensive testing, four earth longitude gravity harmonics of second and third order appear to be well discriminated. These harmonics with their standard errors, for which adjustments for sun and moon effects and the probable influence of neglected higher order earth gravity have been made, are $$
\begin{aligned} \mathrm{J}_{22}= & -(1.816 \pm 0.020) \times 10^{-6} \text { (This corresponds to a difference in major and minor axes of the } \\ & \text { earth's elliptical equator of } 69.4 \pm 0.8 \text { meters.) } \\ \lambda_{22}= & -(15.4 \pm 0.3)^{\circ} \\ \mathrm{J}_{33}= & -(0.171 \pm 0.017) \times 10^{-6} \\ \lambda_{33}= & (24.9 \pm 3.3)^{\circ} . \end{aligned}
$$


In addition to these harmonics, a third pair ( $\mathrm{J}_{31}, \lambda_{31}$ ) was poorly discriminated from the limited acceleration record. The data shows tentatively that

$$
\begin{aligned}
& J_{31}=-\left(1.4_{-0.2}^{+1.0}\right) \times 10^{-6} \\
& \lambda_{31}=-(168 \pm 26)^{\circ} .
\end{aligned}
$$

Tests of the satellite data were also made to try and reveal the influence of resonant fourth order earth gravity. These tests were inconclusive.

The above results show that an equatorial 24 -hour satellite can be in uncontrolled long term east-west equilibrium at only the following four longitude locations:
$\lambda_{1}=76.7 \pm 0.8^{\circ}$ (dynamically stable east-west equilibrium)
$\lambda_{2}=161.8 \pm 0.7^{\circ}$ (statically stable east-west equilibrium)
$\lambda_{3}=-108.1 \pm 1.0^{\circ}$ (dynamically stable east-west equilibrium)
$\lambda_{4}=-12.2 \pm 0.7^{\circ}$ (statically stable east-west equilibrium).

According to the analysis of 24 -hour satellite drift thus far, the maximum longitude acceleration due to earth gravity which could be experienced by the nearly geostationary satellite is

$$
\ddot{\lambda}=-(1.83 \pm 0.05) \times 10^{-3} \text { degrees } / \text { day }^{2}
$$

at about $118^{\circ}$ east of Greenwich. To correct continuously for this east-west acceleration would require, conservatively, a velocity increment of $\Delta V=6.38 \mathrm{ft} / \mathrm{sec} / \mathrm{year}$.

## CONTENTS

Page
Abstract ..... ii
INTRODUCTION ..... 1
Experiment Plan and Error Sources ..... 1
Error Control ..... 2

1. REDUCTION OF THE BASIC ORBIT DATA IN NINE 24-HOUR SATELLITE FREE DRIFT ARCS FOR LONG TERM LONGITUDE ACCELERATIONS ..... 4
Arc 1, Syncom 2, 18 August 1963-18 November 1963 ..... 4
Arc 2, Syncom 2, 28 November 1963-18 March 1964 ..... 6
Arc 3, Syncom 2, 18 March 1964-25 April 1964 ..... 10
Arc 4, Syncom 2, 25 April 1964-4 July 1964 ..... 15
Arc 5, Syncom 2, 4 July 1964-19 February 1965 ..... 17
Arc 6, Syncom 3, 31 October 1964-21 December 1964 ..... 23
Arc 7, Syncom 3, 14 January 1965-16 March 1965 ..... 29
Arc 8, Syncom 2, 25 February 1965-10 May 1965 ..... 35
Arc 9, Early Bird, 23 April 1965-21 June 1965 ..... 39
2. SYNTHESIS OF THE LONGITUDE ACCELERATION RECORD TO REVEAL COMPONENTS IN THE EARTH'S LONGITUDE GRAVITY FIELD ..... 51
3. EAST-WEST EQUILIBRIUM LONGITUDES AND MAXIMUM EAST-WEST STATION KEEPING REQUIREMENTS FOR THE GEOSTATIONARY SATELLITE ..... 66
DISCUSSION ..... 67
CONCLUSIONS ..... 69
ACKNOWLEDGMENTS ..... 71
References ..... 71
Appendix A - Basic Orbit Data Used in This Report ..... 73
Appendix B - Earth Gravity Potential and Force Field Used in This Report: Comparison with Recent Investigations ..... 77
Appendix C - Preliminary Maximum Longitude Accelerations on 24-Hour Satellites Due to the Resonant Gravity Harmonics of the Earth through Fourth Order ..... 85
Appendix D - The Approximate Longitude Excursion of a Slowly Drifting 24-Hour Satellite ..... 89
Appendix E - The Secular Accelerations on Syncom 24-Hour Satellites Due to Particle Atmospheric Drag and Solar Radiation Pressure ..... 95
Appendix F - Average Second Order Resonant Gravity Fields on the Geostationary Satellite ..... 101
Appendix G-List of Symbols ..... 103

# THE EARTH'S LONGITUDE GRAVITY FIELD AS SENSED bY THE DRIFT OF THREE SYNCHRONOUS SATELLITES 

by<br>C. A. Wagner<br>Goddard Space Flight Center

## INTRODUCTION

This report summarizes the results of a two year investigation of the "resonant" earth gravity drift of the world's first three operational synchronous satellites. Previous reports (References 1 through 6) in this series have dealt with the accelerated drift theory for 24-hour satellites and its applications to the reduction of range and range rate orbit data returned from Syncom 2 during its "free" gravity drift over Brazil and the Pacific Ocean.

The specific objective of this summary report is to determine those longitude dependent components of the earth's gravity field which can be fairly said to explain the longitude acceleration record of these satellites (Syncom 2 [1963 31A], Syncom 3 [1964 47A], and Early Bird [1965 28A]) from August 1963 to June 1965. The previous investigations by the author cited above and earlier pioneer theoretical studies such as References 7, 8, 9, and 10 have established beyond reasonable doubt now that the principal long term longitude disturbance on the 24 -hour satellite arises from second order earth longitude gravity, associated with the ellipticity of the earth's Equator. The studies in References 1, 4, 5, and 6 indicated that earth longitude gravity effects of higher than second order on the 24 -hour satellite are at maximum, about an order of magnitude less than the maximum second order effect. Frick and Garber (Reference 9) and the author in the present report and in the simulation studies of References 3,5 , and 6 have shown the long term sun and moon gravity effects on the near circular orbit synchronous satellite to be about two orders of magnitude less than the maximum second tesseral effect for periods of record in excess of about two months.

## Experiment Plan and Error Sources

The basic orbit data of this gravity determining experiment are ascending Equator crossings, orbit vectors and subsatellite points which were reported by the Tracking and Data Systems Directorate of GSFC and the Communications Satellite Corporation without consideration of longitude gravity. Generally, the experiment plan and error control aspects follow that for the early Syncom 2 drift
analyses implicit in References 3,5, and 6. This plan is designed to insure as accurate a result as possible with the limited orbit data available.

The gravity experiment is carried out in two stages. In the first stage the basic orbit data is reduced by an appropriate model, and long term (periods greater than a day) longitude accelerations of the satellites are derived. These are called the "measured" accelerations. In this gravity experiment one hundred and fifty-two independently determined orbits of Syncom 2 and 3 and Early Bird have been reduced to give essentially nine well determined long term longitude accelerations for these satellites. These reduced accelerations are the basic, measured data for the second stage of the experiment. The second stage tests these reduced accelerations for sensitivity to the earth's longitude gravity which is assumed to be responsible for them. In the second stage of the experiment, then, we have a set of "actual" data which consists of measured longitude accelerations (long term) due to the following causes:

1. Physical agencies
a. True earth longitude gravity
b. Sun, moon, planetary and earth zonal gravity
c. Other assumed (or proven) negligibly small physical agencies (nongravitational) such as:
(1) micrometeorite drag or impact
(2) random or systematic outgassing from the satellite control jets
(3) solar wind and radiation pressure
(4) magnetic field interactions.
2. First stage experiment-induced accelerations (errors). Compared to accelerations from true earth longitude gravity, presumably small errors in "measured" accelerations will be due to errors in
a. the basic data (orbit determination) from which the accelerations were deduced
b. the model used to derive the "measured" accelerations from the basic orbit data (first stage model error)

In the second stage of the experiment two error sources are separated. The "basic data error" (in the "measured" accelerations) is considered to be the sum of all the acceleration producing sources in the measured data except true earth longitude gravity. The "model error" in the second stage is considered to be due entirely to the necessarily limited earth model which can be assumed to explain the limited number of measurements in the experiment.

## Error Control

First Experiment Stage (Accelevation Analysis)
The aim of this stage of the experiment is to obtain long term satellite accelerations as "basic data" which are as free from nonearth longitude gravity effects as possible. A number of
nongravitational disturbances on the satellite have been computed theoretically and shown to be negligibly small in this experiment (Appendix E). The others are assumed to be negligibly small also (Reference 3).

In one or two instances basic orbit data has been used at the beginning or end of a long free gravity drift arc which is known to have been disturbed by commanded control gas jet pulsing. In these instances it is evident by inspection that the use of this data does not significantly disturb the previous or following free drift record. This "overlap" data, when used, reduces significantly the standard experiment error over the limited record of that drift arc.

In order to keep the long term sun and moon (principally moon) disturbances almost negligibly small, it was found that free drift arc lengths of about two months or more were required. It also turned out that this was the length of the typical Syncom drift period between orbit corrections. It was also generally necessary to consider a time period of at least two months to obtain reasonably small standard errors for the "measured" Syncom accelerations. Thus, the length of time and record were the major means of error control in the first stage (acceleration analysis) of the experiment.

The model error in the acceleration analysis stage was also under somewhat independent control. The chief consideration for model error control was the mean drift rate. In those arcs (arcs $1,2,6,7,8$, and 9 ) where the mean drift rate was $\pm 0.1$ degree/day or less, the geographic longitude excursion was limited to less than $10^{\circ}$. It was found that with this "slow drift" regime a simple polynomial of third degree in the time could adequately describe the longitude drift of the ascending Equator crossing (Appendix D and Reference 3). In those arcs (arcs 3, 4, and 5) where the mean drift rate was greater than $\pm 0.1$ degree/day ("fast drift" regime) it was found that a simple three parameter function of the longitude of the ascending Equator crossing of the satellite could adequately describe (for a reasonable arc length) the change of the drift rate of these crossings. This was in accord with the use of the energy integral of the gravity drift of the 24 -hour satellite in a simple second order field (see Equation 8 and Reference 2). For arc lengths of up to about $50^{\circ}$, the three parameter (second order) model appeared theoretically adequate to reproduce true gravity accelerations within a reasonably small standard experimental error which includes sun and moon "gravity noise". However, even for somewhat longer arc lengths, simulations show that the simple three parameter velocity model still gives sufficiently good results to be utilized without adjustment (Table 11). This fact depends on the evident overbearing strength of the second order gravity field compared to higher orders. Nevertheless "velocity arcs" of as short a length as possible (to give reasonable standard errors) were chosen to provide a longitude-acceleration survey of as great an extent as possible. This was necessitated, finally, by the method of acceleration analysis. In order to keep the model prejudice in the measurement of the acceleration to as low a level as possible and also to utilize only the best determined statistic in each arc, only a single acceleration near the center of each arc was finally chosen to represent the mean acceleration for that arc. This being the case, it was found that only by breaking up a long drift arc such as arc 5 into smaller, approximately $50^{\circ}$ subarcs, could a reasonably extensive, precise and unprejudiced longitude survey be made with the data at hand.

Error control in this stage was primarily directed toward the refinement of the earth model which reflects the acceleration data. For additional error control, however, experiments with some adjustment of the acceleration data itself for sun, moon and first stage model errors on the basis of simulated 24 -hour trajectories have also been made. In addition, various data weighting schemes and random acceleration choices based on the standard acceleration errors in the drift arcs have been tried as error control methods in arriving at a final reasonable earth gravity synthesis as seen by the data. At each stage of the experiment the experimental error was checked or verified externally by simulations of 24-hour satellite drift in a sun, moon, and/or full earth field, numerically calculated, with trajectory conditions close to those actually experienced in the various arcs.

We now proceed with the first experiment stage and analyze the 24 -hour orbit data reported by GSFC and Comsat for what it reveals in terms of long term accelerations according to the "resonant gravity" models for slow and fast drift regimes previously discussed.

## 1. Reduction of the basic orbit data in nine 24-HOUR satellite free dRIFT ARCS FOR LONG TERM LONGITUDE ACCELERATIONS

## Arc 1, Syncom 2, 18 August 1963-18 November 1963

Syncom 2, the world's first operating synchronous communications satellite, was launched into orbit in late July 1963 and reached station over Brazil in mid-August 1963. The orbit inclination was close to $33^{\circ}$ and the ascending Equator crossing was near $55^{\circ} \mathrm{W}$ moving at less than 0.1 degree/ day eastward after the last corrective thrust was applied going into free drift arc 1 , on 18 August 1963. From this date to 28 November 1963 , the "figure of 8 " ground track of Syncom 2 drifted freely without orbit correction from $55^{\circ} \mathrm{W}$ to $59^{\circ} \mathrm{W}$ when on-board jet pulsing was applied to virtually stop the westward movement of the track. The details of this accelerated free gravity drift are presented in Table 1 and Figure 1.

The acceleration of the ascending Equator crossing in arc 1 was determined from a fit of the orbit data according to the third order polynomial,

$$
\begin{equation*}
L=a_{1}+a_{2} t+a_{3} t^{2}+a_{4} t^{3} \tag{1}
\end{equation*}
$$

which is shown to apply for slow drift regimes in Appendix D and in the simulations in Tables 1S, $1 \mathrm{~S} / 1$, etc. The longitude, L , is the ascending Equator crossing east of an arbitrary base longitude, which, for computational accuracy, is preferably near the center of the drift arc. The time, $t$, is an arbitrary base time which is also preferably located near the center of the arc for computational accuracy.

It is possible to evaluate semi-empirically the error due to sun and moon gravity and model bias implicit in calculating the satellite accelerations due to the earth's longitude gravity in this slow drift arc according to Equation 1. In the second stage of the experiment (gravity synthesis) we will assume long term 24 -hour satellite resonant earth gravity acceleration is derivable from the harmonic expansion through fourth order,

$$
\begin{equation*}
\ddot{\lambda}=-12 \pi^{2} \sum_{n=2}^{4} \sum_{m=1}^{n} F_{n m} \sin m\left(\lambda-\lambda_{n m}\right) F(i)_{n m}, \frac{\text { rad }}{(\text { sid. day })^{2}} \tag{2}
\end{equation*}
$$

$$
\text { for } n-m \text {, even }
$$

where

$$
\begin{align*}
& F_{22} F(i)_{22}=\frac{6 J_{22}}{a_{s}{ }^{2}}\left[\left(\cos i_{s}+1\right) / 2\right]^{2},  \tag{3}\\
& F_{31} F(i)_{31}=\frac{-3 J_{31}}{2 a_{s}{ }^{3}}\left\{\frac{\left(1+\cos i_{s}\right)}{2}-5 \sin ^{2} i_{s}\left(1+3 \cos i_{s}\right)\right\},  \tag{4}\\
& F_{33} F(i)_{33}=\frac{45 J_{33}}{a_{s}{ }^{3}}\left\{\left(\cos i_{s}+1\right) / 2\right\}^{3},  \tag{5}\\
& F_{42} F(i)_{42}=\frac{-15 J_{42}}{a_{s}{ }^{4}}\left\{\frac{\left(1+\cos i_{s}\right)^{2}}{4}-7 \sin ^{2} i_{s} \cdot \cos i_{s}\left(1+\cos i_{s}\right)\right.  \tag{6}\\
& 4 \tag{7}
\end{align*},
$$

See Equation 65 in Reference 2 for the derivation of Equation 2 above.
The symbol $\lambda$ represents the longitude location of the ascending Equator crossing (or mean daily longitude location) for the 24 -hour satellite of reasonably small eccentricity and drift rate (Reference 2). The symbols $a_{s}$ and $i_{s}$ are the "synchronous" semimajor axis (in earth radii) and inclination of the satellite's orbit. The significance of the gravity constants $J_{n m}, \lambda_{n m}$ is contained in Appendix B and explained in further detail in References 2 and 11. In evaluating the first stage model error we calculate numerically, particle trajectories closely parallel to the actual one in free drift, including as many relevant perturbation effects as desired. (In References 3, 5, and 6 these simulated trajectories clearly show the necessity of considering earth longitude gravity in the long term orbit determination for the 24 -hour satellite.) Then, a direct comparison of the longitude acceleration measured in the simulated trajectory with the theoretical 24 -hour satellite resonant gravity acceleration as given by Equation 2 gives an estimate of the bias error from the effects neglected in the real trajectory analysis. These errors include sun and moon gravity accelerations and model error implicit in the limited accuracy of the longitude Equation 1.

Ideally, this numerical approach to assessing nonresonant gravity effects should involve an attempt to duplicate as closely as possible the real trajectory. This is to avoid the criticism that
the nonresonant effects in the simulated trajectory may not be exactly or even nearly the same as in the actual trajectory. (Compare, for example, the two trajectory results in Table 4S.) It can be appreciated that the full process of such duplication, or closest duplication (in a least squares sense, for example), must involve adjustment of at least the six initial trajectory parameters, as well as the earth longitude gravity constants, the earth radius, and the principal gravity constant. Such an analysis is considerably beyond the scope of this one. Still, extensive experience with a more limited numerical approach to this problem appears to show that such bias errors can be fairly accurately determined without precise duplication, at least under a reasonably wide range of longitude gravity constants.

In this approach only the semimajor axis of the actual 24 -hour orbit at the beginning of each arc was adjusted so that, together with a fixed nominal second order longitude gravity field, the simulated trajectory gives close longitude-time congruence with the actual. Data and results from such trajectories are found in Tables $1 \mathrm{~S}, 2 \mathrm{~S}, 4 \mathrm{~S}$, etc. The second order earth longitude field was specified as $\mathrm{J}_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18.0^{\circ}$ in conformance with the 24 -hour gravity results of Reference 3. At a later stage in the analysis, a more accurate set of longitude gravity harmonics through third order was derived. Simulated trajectories utilizing this set with the same initial elements as before were calculated and analyzed. Data and results from these trajectories are found in Tables $1 \mathrm{~S} / 1,2 \mathrm{~S} / 1,4 \mathrm{~S} / 1$, etc. Except for arc 4 S versus $4 \mathrm{~S} / 1$, the bias results of thesetwo parallel simulated trajectories are consistent. The principal conclusion from the study of the simulated trajectories (not all of which are reported here) is that when arc lengths in excess of two months are considered, the cumulative model bias acceleration errors are within $\pm 0.03 \times 10^{-5}$ radians/sid. day ${ }^{2}$, RMS (see Table 11). In Section 2 the net gravity effect of these model bias errors in each acceleration measurement is shown to be small. In fact, there is evidence that the actual data analyzed without bias adjustment gives longitude gravity with greater precision than with such adjustment. The implication is that over the limited acceleration record of the experiment the bias error acts to cancel more often than not the random observation error attributable to the imprecise orbit determination.

The principal results of the acceleration analysis on the actual data in Table 1 and the simulated data in Tables 1 S and $1 \mathrm{~S} / 1$ are listed in Tables 10 and 11 in the next section.

## Arc 2, Syncom 2, 28 November 1963-18 March 1964

On 28 November 1963, on-board jets were fired to virtually stop the westward drift of Syncom 2 which had built up over the previous three months due to earth longitude gravity. From 28 November 1963 to 18 March 1964 Syncom 2 was allowed to drift freely about $8^{\circ}$ in mean longitude from $59^{\circ} \mathrm{W}$ to $67^{\circ} \mathrm{W}$ under the westward accelerating influence of earth longitude gravity. The orbit inclination during this time was about $32.8^{\circ}$. The details of this drift are presented in Table 2 and Figure 2. The length of the drift record in this arc should be long enough to extract two well defined acceleration values by considering a $t^{4}$ fit (see Appendix D). Unfortunately the orbit determination errors in this arc appear to be too great to allow this finer discrimination.

Syncom 2 Osculating Elements at the First Ascending Equator Crossing Past the Tracking Epoch and Related Data for Free Drift Are 1*.

| Orbit <br> Number <br> 1- | Tracking Epoch** (yr-mo-day-hr-min, UT) | Semimajor Axis, a (earth radii) | Inclination, i <br> (degrees) | Right Ascension of the Ascending Node (degrees) | Eccentricity | Argument of Perigee (degrees) | Mean <br> Anomaly <br> (degrees) | Time from January 0.0, 1963 (days) | Geographic <br> Longitude of First <br> Ascending <br> Equator <br> Crossing <br> After <br> Tracking <br> Epoch, <br> $\lambda$ <br> (degrees) | (1) <br> Time from January 276.5057, 1963, t (days) | (2) Longitude of Ascending Equator Crossing East of -56.5170, L (degrees) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 63-8-18-1-30.0 | 6.6105587 | 33.120 | -42.357 | . 00023 | 37.901 | -37.890 | 230.1302 | -55.004 | -46.3755 | 1.513 |
| 2 | 63-8-22-6-12.14 | 6.6105498 | 33.081 | -42.444 | . 00024 | 53.529 | -53.508 | 235.1159 | -54.893 | -41.3898 | 1.624 |
| 3 | 63-8-26-17.0 | 6.6105779 | 33.091 | -42.553 | . 00018 | 58.903 | -58.888 | 239.1047 | -54.847 | -37.4010 | 1.670 |
| 4 | 63-8-31.0 | 6.6107824 | 33.062 | -42.526 | . 00022 | 45.989 | -45.974 | 243.0937. | -54.803 | -33.4120 | 1.714 |
| 5 | 63-9-3-13-23.0 | 6.6105663 | 33.082 | -42.635 | . 00026 | 42.792 | -42.773 | 247.0824 | -54.792 | -29.4233 | 1.725 |
| 6 | 63-9-5.0 | 6.6110747 | 33.064 | -42.639 | . 00016 | 22.930 | -22.924 | 248.0796 | -54.753 | -28.4261 | 1.764 |
| 7 | 63-9-9.0 | 6.6107958 | 33.048 | -42.728 | . 00015 | 15.313 | -15.308 | 252.0685 | -54.774 | -24.4372 | 1.743 |
| 8 | 63-9-12-2.0 | 6.6110077 . | 33.078 | -42.788 | . 00021 | 31.588 | -31.577 | 256.0575 | -54.812 | -20.4482 | 1.705 |
| 9 | 63-9-17-2.0 | 6.6108515 | 33.040 | -42.850 | . 00020 | 29.394 | -29.384 | 261.0438 | -54.862 | -15.4619 | 1.655 |
| 10 | 63-9-20-2.0 | 6.6109260 | 33.009 | -42.919 | . 00020 | 39.841 | -39.830 | 264.0356 | -54.932 | -12.4701 | 1.585 |
| 11 | 63-9-27-2.0 | 6.6111697 | 33.031 | -43.018 | . 00020 | -24.116 | 24.109 | 271.0165 | -55.037 | - 5.4892 | 1.480 |
| 12 | 63-10-1-2.0 | 6.6107666 | 33.023 | -43.112 | . 00020 | 8.796 | - 8.793 | 275.0056 | -55.142 | - 1.5001 | 1.375 |
| 13 | 63-10-8-2.0 | 6.6114380 | 33.013 | -43.233 | . 00025 | 11.624 | -11.619 | 281.9872 | -55.519 | 5.4815 | . 998 |
| 14 | 63-10-14-2.0 | 6.6113174 | 32.979 | -43.201 | . 00030 | 23.884 | -23.872 | 287.9715 | -55.740 | 11.4658 | . 777 |
| 15 | 63-10-22-2.0 | 6.6116625 | 32.994 | -43.411 | . 00029 | -23.336 | 23.325 | 295.9503 | -56.171 | 19.4446 | . 346 |
| 16 | 63-10-30.0 | 6.6113312 | 32.946 | -43.449 | . 00024 | - 9.384 | 9.380 | 303.9297 | -56.658 | 27.4240 | - . 141 |
| 17 | 63-11-6.0 | 6.6119680 | 32.952 | -43.683 | . 00031 | 15.139 | -15.130 | 310.9116 | -57.253 | 34.4059 | - . 736 |
| 18 | 63-11-12-5.0 | 6.6117819 | 32.919 | -43.703 | . 00030 | 25.807 | -25.792 | 316.8964 | -57.713 | 40.3907 | -1.196 |
| 19 | 63-11-18-13.0 | 6.6120265 | 32.925 | -43.877 | . 00019 | 15.413 | -15.409 | 322.8811 | -58.280 | 46.3754 | -1.763 |
|  |  | Average: $6.611113=\mathrm{a}_{\mathrm{S}}$ | Average: $33.024=\mathrm{i}$ |  |  |  |  |  |  |  |  |

*One earth radius $=6378.388 \mathrm{~km}$; the earth gravity constant used in the trajectory program, $\mu_{e}=3.986267 \times 10^{5} \mathrm{~km}^{3} / \mathrm{sec}^{2}$
The osculating elements and the equator crossing data were derived from the satellite vectors reported by the Tracking and Data Systems Directorate of NASA-GSFC in Table A1. The trajectory generator, called "ITEM" (Interplanetary Trajectory by an Encke Method) at GSFC, for this derivation used the same earth, moon, and sun model as the original orbit (vector) determination program (see Table A1).
**The tracking epoch refers to the epoch of the satellite vectors reported in Table A1.
Results of least squares fit of data in (1) and (2) above according to the theory of Equation 1 :

$$
L=a_{1}+a_{2} t+a_{3} t^{2}+a_{4} t^{3}
$$

$a_{1}=1.2720 \pm 1.032 \times 10^{-2}$ degrees
$a_{2}=-(3.4899 \pm 0.0589) \times 10^{-2}$ degrees $/ \mathrm{sol}$. day
$a_{3}=-(6.4947 \pm 0.0938) \times 10^{-4}$ degrees $/$ sol. day ${ }^{2}$
$a_{4}=-(1.765 \pm 3.795) \times 10^{-7}$ degrees $/$ sol. day ${ }^{3}$
Standard error of estimate $=0.02825$
$\ddot{\lambda}$ (with minimum standard error) $=-(2.253 \pm 0.0325) \times 10^{-5} \mathrm{rad} /$ sid. day ${ }^{2}$, at $\mathrm{t}=-0.674$ days, $\mathrm{L}=1.296^{\circ}, \lambda=-55.22^{\circ}$, on January 275.832, 1963. See Figure 1 .

Ascending Equator Crossing Orbit Data from a Simulated Syncom 2 Trajectory for Free Drift Arc 1
with Earth Longitude Gravity Through Second Order.*

| Orbit Number 1S- | Tracking Epoch (yr-mo-day-hr-min UT) | Semimajor Axis, a (earth radii) | $\begin{aligned} & \text { Inclination, } \\ & \text { i } \\ & \text { (degrees) } \end{aligned}$ | Right <br> Ascension of the Ascending Node (degrees) | Eccentricity | Argument of Perigee (degrees) | Mean Anomaly (degrees) | Time from 1963.0 (days) | Geographic <br> Longitude of the <br> Ascending <br> Equator <br> Crossing, <br> $\lambda$ <br> (degrees) | (1) Time from January 276.5056, 1963, $t$ (days) | (2) <br> Longitude of the <br> Ascending Equator Crossing East of $-56.537^{\circ}$, L (degrees) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 63-8-18-1.5 | 6.6106431 | 33.120 | -42.358 | . 00023 | 37.945 | -37.930 | 230.1302 | -55.006 | -46.3754 | 1.531 |
| 2 | 63-8-22-6-12.14 | 6.6105732 | 33.102 | -42.423 | . 00024 | 53.526 | -53.505 | 235.1161 | -54.888 | -41.3895 | 1.649 |
| 3 | 63-8-26-17.0 | 6.6105273 | 33.095 | -42.488 | . 00019 | 57.150 | -57.131 | 239.1047 | -54.814 | -37.4009 | 1.723 |
| 4 | 63-8-31.0 | 6.6107816 | 33.091 | -42.539 | . 00017 | 44.072 | -44.063 | 243.0935 | -54.765 | -33.4121 | 1.772 |
| 5 | 63-9-3-13-23.0 | 6.6108968 | 33.077 | -42.585 | . 00021 | 48.646 | -48.632 | 247.0824 | -54.736 | -29.4232 | 1.801 |
| 6 | 63-9-5.0 | 6.6108259 | 33.072 | -42.601 | . 00020 | 54.853 | -54.836 | 248.0796 | -54.731 | -28.4260 | 1.806 |
| 7 | 63-9-9.0 | 6.6106796 | 33.061 | -42.673 | . 00014 | 61.617 | -61.604 | 252.0685 | -54.724 | -24.4371 | 1.81 .3 |
| 8 | 63-9-12-2.0 | 6.6109677 | 53.058 | -42.729 | . 00020 | 38.963 | -38.952 | 256.0575 | -54.750 | -20.4481 | 1.787 |
| 9 | 63-9-17-2.0 | 6.6110100 | 33.040 | -42.793 | . 00026 | 52.421 | -52.402 | 261.0438 | -54.801 | -15.4618 | 1.736 |
| 10 | 63-9-20-2.0 | 6.6109082 | 33.031 | -42.844 | . 00023 | 59.346 | -59.327 | 264.0356 | -54.840 | -12.4700 | 1.697 |
| 11 | 63-9-27-2.0 | 6.6111994 | 33.024 | -42.949 | . 00018 | 46.804 | -46.791 | 271.0166 | -54.986 | - 5.4890 | 1.551 |
| 12 | 63-10-1-2.0 | 6.6112339 | 33.009 | -43.002 | . 00021 | 55.489 | -55.469 | 275.0058 | -55.098 | - 1.4998 | 1.439 |
| 13 | 63-10-8-2.0 | 6.6112347 | 32.996 | -43.128 | . 00016 | 44.518 | -44.505 | 281.9870 | -55.348 | 5.4814 | 1.189 |
| 14 | 63-10-14-2.0 | 6.6113821 | 32.980 | -43.208 | . 00025 | 53.893 | -53.871 | 287.9712 | -55.616 | 11.4656 | . 921 |
| 15 | 63-10-22-2.0 | 6.6113844 | 32.970 | -43.346 | . 00017 | 58.412 | -58.401 | 295.9501 | -56.029 | 19.4445 | . 508 |
| 16 | 63-10-30.0 | 6.6114845 | 32.952 | -43.461 | . 00019 | 67.886 | -67.871 | 303.9293 | -56.530 | 27.4237 | . 007 |
| 17 | 63-11-6.0 | 6.6117444 | 32.947 | -43.581 | . 00020 | 42.270 | -42.258 | 310.9113 | -57.051 | 34.4057 | -. 514 |
| 18 | 63-11-12-5.0 | 6.6117173 | 32.932 | -43.668 | . 00023 | 58.345 | -58.325 | 316.8960 | -57.537 | 40.3904 | -1.000 |
| 19 | $63-11-18-13.0$ | 6.6117764 | 32.930 | -43.770 | . 00017 | 57.820 | -57.803 | 322.8809 | -58.067 | 46.3753 | -1.530 |
|  | Average: $6.611104=\mathrm{a}_{\mathrm{s}}$ | Average: $33.026=\mathrm{i}_{\mathrm{s}}$ |  |  |  |  |  |  |  |  |  |

* Computed by ITEM with gravity constants the same as in Table A 1 with the addition of earth constants, $\mathrm{J}_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18^{\circ}$.

Results of least squares fit of data in (1) and (2) above according to the theory of Equation 1:

$$
\begin{aligned}
& \mathrm{L}_{1}=\mathrm{a}_{1}+\mathrm{a}_{2} \mathrm{t}^{\mathrm{t}}+\mathrm{a}_{3} \mathrm{t}^{2}+\mathrm{a}_{4} \mathrm{t}^{3} \\
& \mathrm{a}_{1}=(1.3885 \pm 0.00152) \text { degrees } \\
& \mathrm{a}_{2}=-(3.289 \pm 0.00866) \times 10^{-2} \text { degrees } / \text { day } \\
& \mathrm{a}_{3}=-(6.451 \pm 0.0138) \times 10^{-4} \text { degrees } / \text { day }{ }^{2} \\
& \mathrm{a}_{4}=-(7.564 \pm 5.577) \times 10^{-8} \text { degrees } / \text { day }{ }^{3}
\end{aligned}
$$

Standard error of estimate $=0.004151$ degree
$\ddot{\lambda}($ measured $)=-(2.2389 \pm 0.0048) \times 10^{-5} \mathrm{rad} / \mathrm{sid}^{2}$. day $^{2}$, for $\mathrm{t}=-0.674$ day, $\mathrm{L}=1.4104^{\circ}, \lambda=-55.127^{\circ}$
$\ddot{\lambda}($ theoretical, from Equation 2$)=-(2.2211) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, for $\mathrm{a}_{\mathrm{s}}=6.611104$ earth radii, $\mathrm{i}_{\mathrm{s}}=33.026^{\circ}, \lambda=-55.127^{\circ}, \mathrm{J}_{22}=-1.68 \times 10^{-6}, \lambda_{22}=.18 .0^{\circ}$.
Estimate of measured bias due to sun-moon perturbations and $J_{22}$ model error (exclusive of higher order longitude gravity effects) in $\ddot{\lambda}$ at $t=-0.674$ day in Syncom 2 arc 1 on
January 275.832, 1963: Bias = theoretical -measured
$=-(2.2211) \times 10^{-5}+2.2389 \times 10^{-5}=+(0.0178) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$.

Ascending Equator Crossing Data from a Simulated Syncom 2 Trajectory for Free Drive Arc 1
Computed by ITEM with Earth Longitude Gravity through Third Order.*

*Gravity constants of this trajectory are the same as those in Table A1 with the addition of the earth constants:

$$
\mathrm{J}_{22}=-1.8 \times 10^{-6} \quad \lambda_{22}=-15.35^{\mathrm{o}} \quad \mathrm{~J}_{33}=-0.16 \times 10^{-6} \quad \lambda_{33}=-24^{\mathrm{o}} \quad \mathrm{~J}_{31}=-1.5 \times 10^{-6} \quad \lambda_{31}=0^{\circ},
$$

(see Figure Bl for the significance of the constants).
The initial elements of this trajectory, aside from those listed for orbit $1 S / 1-1$, are the same as those in orbit $1 S-1$ (Table 1S).
Results of least squares fit of data in (1) and (2)above according to the theory of Equation (for arc $1 \mathrm{~S} / 1$ ):
$L=a_{1}+a_{2} t+a_{3} t^{2}+a_{4} t^{3}$
$\mathrm{a}_{1}=(1.2439 \pm 0.00151)$ degrees
$a_{2}=-(3.2953 \pm 0.00864) \times 10^{-2}$ degrees $/$ day
$a_{3}=-(6.4319 \pm 0.0138) \times 10^{-4}$ degrees $/$ day $^{2}$
$a_{4}=-(2.94 \pm 5.57) \times 10^{-8}$ degrees $/$ day $^{3}$

## 3tandard error of estimate $=0.00414$ degree

$\ddot{\lambda}$ (measured, with minimum standard error) $=-2.2328 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, at $\lambda=-55.13^{\circ}$ on $\mathrm{t}^{\prime}=275.8318^{\circ}$ January 1963.
$\lambda$ (theoretical, from Equation 2) $=-2.2185 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, for $\mathrm{a}_{\mathrm{s}}=6.611$ earth radii, $i_{\mathrm{s}}=33.026^{\circ}, \lambda=-55.13^{\circ}, \mathrm{J}_{22}-\mathrm{J}_{31}$ as noted.
Estimate of acceleration bias at $\lambda=-55.13^{\circ}$ in arc $1 \mathrm{~S} / 1=\ddot{\lambda}$ (theoretical) $-\ddot{\lambda}$ (measured) $=+0.0143 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$.


Figure 1-Measured and simulated orbit data at ascending Equator crossings in free drift arc 1 (Syncom 2).

The principal results of the acceleration analysis on the actual data in Table 2 and the simulated data in Tables 2 S and $2 \mathrm{~S} / 1$ are listed in Tables 10 and 11 at the end of the next section.

## Arc 3, Syncom 2, 18 March 1964-25 April 1964

On 18 March 1964 the westward drift of Syncom 2 was speeded to 1.3 degrees/day by on-board gas jet pulsing. Between 18 March and 25 April, Syncom 2 drifted rapidly (in a "fast drift" regime) from $67^{\circ} \mathrm{W}$ to $116^{\circ} \mathrm{W}$. The orbit inclination during this period was about 32.7 .

Syncom 2 Osculating Elements at the First Ascending Equator Crossing Past the Tracking Epoch in Free Drift Arc 2.*

| Orbit Number 2- | Tracking Epoch (yr-mo-day=hr-min UT) | Seminajor Axis, a (earth radii) | Inclination, i (degrees) | Right <br> Ascension of the Ascending Node (degrees) | Eccentricity | Argument of Perigee (degrees) | Mean <br> Anomaly <br> (degrees) | Time from 1963.0 <br> (days) | Geographic <br> Longitude <br> of First <br> Ascending <br> Equator <br> Crossing <br> After <br> Tracking <br> Epoch $\lambda$ <br> (degrees) | (1) $\begin{gathered} \text { Time from } \\ \text { January } \\ 384.2230,1963, \\ t \\ \text { (days) } \end{gathered}$ | (2) <br> Longitude of the Ascending Equator Crossing East of $-62.649^{\circ}$, L (degrees) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 63-11-28-1.0 | 6.6103981 | 32.920 | -44.040 | . 00006 | -153.442 | 153.440 | 332.8558 | -59.161 | -51.3672 | 3.488 |
| 2 | 63-12-4.0 | 6.6112988 | 32.892 | -44.090 | . 00019 | 46.962 | - 46.949 | 338.8395 | -59.227 | -45.3835 | 3.422 |
| 3 | 63-12-10.0 | 6.6108885 | 32.881 | -44.138 | . 00009 | - 2.249 | 2.249 | 344.8230 | -59.243 | -39.4000 | 3.406 |
| 4 | 63-12-16-17.0 | 6.6109177 | 32.873 | -44.266 | . 00010 | 95.745 | - 95.744 | 350.8067 | -59.392 | -33.4163 | 3.257 |
| 5 | 63-12-23-19.0 | 6.6110073 | 32.795 | -44.242 | . 00019 | 81.870 | - 81.852 | 358.7851 | -59.465 | -25.4379 | 3.184 |
| 6 | 64-1-6-17.0 | 6.6111007 | 32.867 | -44.456 | . 00014 | 14.514 | - 14.511 | 371.7509 | -60.155 | -12.4721 | 2.494 |
| 7 | 64-1-9-6.0 | 6.6113769 | 32.857 | -44.539 | . 00013 | 5.238 | - 5.237 | 374.7431 | -60.360 | - 9.4799 | 2.289 |
| 8 | 64-1-15-18.0 | 6.6117781 | 32.810 | $-44.580$ | . 00025 | 60.341 | - 60.322 | 381.7246 | -60.607 | - 2.4984 | 2.042 |
| 9 | 64-1-20-21.0 | 6.6115475 | 32.825 | -44.713 | . 00016 | 68.659 | - 68.647 | 386.7119 | -61.112 | 2.4889 | 1.537 |
| 10 | 64-1-29-20.0 | 6.6120395 | 32.856 | -44.797 | . 00029 | 40.492 | - 40.473 | 395.6892 | -61.863 | 11.4662 | . 786 |
| 11 | 64-2-5-16.0 | 6.6119118 | 32.800 | -44.925 | . 00026 | 49.760 | - 49.738 | 401.6738 | -62.350 | 17.4508 | . 299 |
| 12 | 64-2-10-19.0 | 6.6123174 | 32.832 | -45.026 | . 00024 | 37.741 | - 37.727 | 407.6588 | -62.958 | 23.4358 | - . 309 |
| 13 | 64-2-17-17.0 | 6.6120695 | 32.760 | -45.134 | . 00023 | 54.393 | - 54.376 | 414.6413 | -63.632 | 30.4183 | -. 983 |
| 14 | 64-2-25-19.0 | 6.6124562 | 32.767 | -45.182 | . 00036 | 52.806 | - 52.775 | 422.6218 | -64.543 | 38.3988 | -1.894 |
| 15 | 64-3-4-23.0 | 6.6123649 | 32.723 | -45.386 | . 00017 | 49.417 | - 49.405 | 430.6019 | -65.452 | 46.3789 | -2.803 |
| 16 | 64-3-10-13.0 | 6.6124118 | 32.747 | -45.391 | . 00023 | 25.910 | - 25.899 | 435.5902 | -66.136 | 51.3672 | -3.487 |
|  |  | Average: $6.6116178=a_{s}$ | Average: $32.825=\mathrm{i}_{\mathrm{s}}$ |  |  |  |  |  |  |  |  |

*Computed by ITEM with gravity constants the same as in Table $A 1$ with the addition of earth constants, $\mathrm{J}_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18^{\circ}$.
Results of least squares fit of data in(1)and(2)above according to the theory of Equation 1 :

$$
L=a_{1}+a_{2} t+a_{3} t^{2}+a_{4} t^{3}
$$

$a_{1}=1.7349 \pm 2.273 \times 10^{-2}$ degrees
$a_{2}=-(7.118 \pm 0.126) \times 10^{-2}$ degrees $/$ sol. day
$a_{3}=-(6.6165 \pm 0.1650) \times 10^{-4}$ degrees $/$ sol. day ${ }^{2}$
$a_{4}=(1.492 \pm 0.636) \times 10^{-6}$ degrees $/$ sol. day ${ }^{3}$
Standard error of estimate $=6.041 \times 10^{-2}$ degrees
$\ddot{\lambda}$ (with minimum standard error) $=-(2.291 \pm 0.0572) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, at $t=0.412 \mathrm{day}, \mathrm{L}=1.706^{\circ}, \lambda=60.94^{\circ}$. (See Figure 2)
See Table 1 For Additional Notes.

Osculating Elements at the First Ascending Equator Crossing Past the Syncom 2 Tracking Epoch, and Related Data
in a Simulated Syncom 2, Arc 2 Trajectory with Earth Longitude Gravity through Second Order.*

| Orbit <br> Number 2S- | Tracking Epoch (yr-mo-day-hr-min UT) | Semimajor Axis, a (earth radii) | Inclination, <br> i <br> (degrees) | Right Ascension of the Ascending Node (degrees) | Eccentricity | Argument of Perigee (degrees) | Mean Anomaly (degrees) | Time from 1963.0 (days) | Geographic Longitude of first Ascending Equator Crossing After Tracking Epoch, $\lambda$ (degrees) | (1) Time from January $384.2230,1963$, $t$ (days) | (2) Longitude of the Ascending Equator Crossing East of $-62.649^{\circ}$, L (degrees) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 63-11-28-1.0 | 6.6105430 | 32.920 | -44.040 | . 00005 | -153.540 | 153.543 | 332.8558 | -59.160 | -51.3672 | 3.489 |
| 2 | 63-12-4.0 | 6.6109361 | 32.920 | -44.128 | . 00008 | - 18.896 | 18.894 | 338.8393 | -59.201 | -45.3837 | 3.448 |
| 3 | 63-12-10.0 | 6.6108628 | 32.906 | -44.213 | . 00004 | 34.301 | - 34.307 | 344.8229 | -59.270 | -39.4001 | 3.379 |
| 4 | 63-12-16-17.0 | 6.6110089 | 32.907 | -44.304 | . 00004 | - 66.460 | 66.491 | 350.8066 | -59.386 | -33.4164 | 3.263 |
| 5 | 63-12-23-19.0 | 6.6111138 | 32.891 | -44.404 | . 00001 | - 72.413 | 72.423 | 358.7851 | -59.615 | -25.4379 | 3.034 |
| 6 | 64-1-6-17.0 | 6.6113789 | 32.870 | -44.582 | . 00007 | 25.588 | - 25.594 | 371.7507. | -60.185 | -12.4723 | 2.464 |
| 7 | 64-1-9-6.0 | 6.6113491 | 32.867 | -44.629 | . 00004 | 8.904 | - 8.904 | 374.7428 | -60.341 | - 9.4802 | 2.308 |
| 8 | 64-1-15-18.0 | 6.6117630 | 32.861 | -44.709 | . 00008 | - 20.008 | 20.007 | 381.7246 | -60.760 | - 2.4984 | 1.889 |
| 9 | 64-1-20-21.0 | 6.6115831 | 32.844 | -44.778 | . 00002 | - 7.961 | 7.969 | 386.7117 | -61.093 | 2.4887 | 1.556 |
| 10 | 64-1-29-20.0 | 6.6119796 | 32.829 | -44.893 | . 00013 | 14.624 | - 14.621 | 395.6888 | -61.801 | 11.4658 | . 848 |
| 11 | 64-2-5-16.0 | 6.6118246 | 32.811 | -44.984 | . 00006 | 25.370 | - 25.371 | 401.6735 | -62.313 | 17.4505 | . 336 |
| 12 | 64-2-10-19.0 | 6.6121671 | 32.804 | -45.058 | . 00009 | - 13.385 | 13.385 | 407.6585 | -62.879 | 23.4355 | - . 230 |
| 13 | 64-2-17-17.0 | 6.6120119 | 32.778 | -45.157 | . 00002 | 24.230 | - 24.248 | 414.6411 | -63.594 | 30.4181 | -. 945 |
| 14 | 64-2-25-19.0 | 6.6124152 | 32.762 | -45.263 | . 00013 | 18.256 | - 18.252 | 422.6215 | -64.512 | 38.3985 | -1.863 |
| 15 | 64-3-4-23.0 | 6.6122805 | 32.739 | -45.392 | . 00006 | 26.790 | - 26.792 | 430.6020 | -65.490 | 46.3790 | -2.841 |
| 16 | 64-3-10-13.0 | 6.6126075 | 32.731 | -45.455 | . 00009 | - 0.415 | 0.415 | 435.5900 | -66.150 | 51.3670 | -3.501 |
|  |  | $\begin{gathered} \text { Average: } \\ 6.611614=\mathrm{a}_{\mathrm{s}} \end{gathered}$ | $\begin{gathered} \text { Average: } \\ 32.84=\mathrm{i}_{\mathrm{s}} \end{gathered}$ |  |  |  |  |  |  |  |  |

*Computed by ITEM with gravity constants the same as in Table A1 with the addition of the earth constants: $\mathrm{J}_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18.0^{\circ}$.
Results of least squares fit of data in (1) and (2) above according to the theory of Equation 1:
$L=a_{1}+a_{2} t+a_{3} t^{2}+a_{4} t^{3}$
$a_{1}=1.7226 \pm 2.223 \times 10^{-3}$ degrees
$a_{2}=-(6.7983 \pm 0.01237) \times 10^{-2}$ degrees $/$ sol. day
$\mathrm{a}_{3}=-(6.5640 \pm 0.01614) \times 10^{-4}$ degrees $/ \mathrm{sol}$. day ${ }^{2}$
$\mathrm{a}_{4}=-(1.493 \pm 62.19) \times 10^{-9}$ degrees $/$ sol. day ${ }^{3}$
Standard error of estimate $=5.908 \times 10^{-3}$ degrees
$\ddot{\lambda}$ (measured $)=-(2.2787 \pm 0.00560) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, at $t=0.412$ day and $\lambda=-60.954^{\circ}$
$\ddot{\lambda}\left(\right.$ theoretical, from Equation 2) $=-2.3061 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, for $\mathrm{a}_{\mathrm{s}}=6.61161$ earth radii, $\mathrm{i}_{\mathrm{s}}=32.84^{\circ}, \lambda=-60.954^{\mathrm{o}}, \mathrm{J}_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18.0^{\circ}$
Estimate of acceleration bias due to sun-moon perturbations and $J_{22}$ model error (exclusive of higher order longitude gravity effects), in $\boldsymbol{\lambda}^{\text {at }} \mathrm{t}=0.412$ day in Syncom arc 2 :
Bias $=$ theoretical - measured $=-\left(2.3061 \times 10^{-5}\right)+\left(2.2787 \times 10^{-5}\right)=-(0.0274) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day $^{2}$

Table 2S/1
Ascending Equator Crossing Data from a Simulated Syncom 2 Trajectory for Free Drift Arc 2, Computed by ITEM with Earth Longitude Gravity through Third Order*.

| Orbit Number 2S/1- | Tracking Epoch (yr-mo-day-hr UT) | Semimajor Axis, a (earth radii) | $\begin{gathered} \text { Inclination, } \\ \text { i } \\ \text { (degrees) } \end{gathered}$ | Time from 1963.0 (days) | Geographic <br> Longitude <br> of the <br> Ascending <br> Equator <br> Crossing, <br> $\lambda$ <br> (degrees) | $\begin{gathered} \text { © } \\ \text { Time From } \\ \text { January } \\ \text { 384.2227, 1963, } \\ \text { t } \\ \text { (days) } \end{gathered}$ | (2) <br> Longitude of Ascending Equator Crossing East of $-62.556^{\circ}$, L (degrees) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 63-11-28-1.0 | 6.6105430 | 32.920 | 332.8558 | -59.160 | -51.3669 | 3.396 |
| 2 | 63-12-4.0 | 6.6109340 | 32.920 | 338.8393 | -59.200 | -45.3834 | 3.356 |
| 3 | 63-12-10.0 | 6.6108585 | 32.906 | 344.8229 | -59.268 | -39.3998 | 3.288 |
| 4 | 63-12-16-17.0 | 6.6110021 | 32.907 | 350.8066 | -59.382 | -33.4161 | 3.174 |
| 5 | 63-12-23-19.0 | 6.6111039 | 32.891 | 358.7851 | -59.605 | -25.4376 | 2.951 |
| 6 | 64-1-6-17.0 | 6.6113633 | 32.870 | 371.7506 | -60.162 | -12.4721 | 2.394 |
| 7 | 64-1-9-6.0 | 6.6113321 | 32.867 | 374.7427 | -60.315 | - 9.4800 | 2.241 |
| 8 | 64-1-15-18.0 | 6.6117426 | 32.861 | 381.7245 | -60.723 | - 2.4982 | 1.833 |
| 9 | 64-1-20-21.0 | 6.6115603 | 32.844 | 386.7116 | -61.048 | 2.4889 | 1.508 |
| 10 | 64-1-29-20.0 | 6.6119519 | 32.829 | 395.6886 | -61.737 | 11.4659 | . 819 |
| 11 | 64-2-5-16.0 | 6.6117933 | 32.811 | 401.6733 | -62.234 | 17.4506 | . 322 |
| 12 | 64-2-10-19.0 | 6.6121317 | 32.804 | 407.6583 | -62.785 | 23.4356 | - . 229 |
| 13 | 64-2-17-17.0 | 6.6119720 | 32.778 | 414.6408 | -63.479 | 30.4181 | - . 923 |
| 14 | 64-2-25-19.0 | 6.6123694 | 32.762 | 422.6211 | -64.368 | 38.3984 | -1.812 |
| 15 | 64-3-4-23.0 | 6.6122286 | 32.739 | 430.6016 | -65.315 | 46.3789 | -2.759 |
| 16 | 64-3-10-13.0 | 6.6125514 | 32.731 | 435.5895 | -65.952 | 51.3668 | -3.396 |
|  |  | Average: $6.6116=a_{s}$ | Average: $32.84=i_{\mathrm{S}}$ |  |  |  |  |

(see Figure 2)
*Gravity constants of this trajectory are the same as that in Table A1, with the addition of the earth constants:

$$
\begin{aligned}
& \mathrm{J}_{22}=-1.8 \times 10^{-6}, \lambda_{22}=-15.35^{\circ} \\
& \mathrm{J}_{33}=-0.16 \times 10^{-6}, \lambda_{33}=24.0^{\circ} \\
& \mathrm{J}_{31}=-1.5 \times 10^{-6}, \lambda_{31}=0.0^{\circ}
\end{aligned}
$$

(See Figure B1 for the significance of these constants). The initial elements of this trajectory, aside from those listed for orbit $2 \mathrm{~S} / 1-1$, are the same as those in orbit $2 \mathrm{~S}-1$ (Table 2S).
Results of least squares fit of data in (1) and (2) according to the theory of Equation 1: (for arc $2 \mathrm{~S} / 1$ )

$$
\begin{aligned}
& \mathrm{L}=\mathrm{a}_{1}+\mathrm{a}_{2} \mathrm{t}+\mathrm{a}_{3} \mathrm{t}^{2}+\mathrm{a}_{4} \mathrm{t}^{3} \\
& \mathrm{a}_{1}=(1.6705 \pm 0.00214) \text { degrees } \\
& \mathrm{a}_{2}=-(6.6252 \pm 0.0119) \times 10^{-2} \text { degrees } / \text { sol. day } \\
& \mathrm{a}_{3}=-(6.3450 \pm 0.0155) \times 10^{-4} \text { degrees } / \text { sol. day }{ }^{2} \\
& \mathrm{a}_{4}=(7.13 \pm 5.98) \times 10^{-8} \text { degrees } / \text { sol. day }{ }^{3}
\end{aligned}
$$

Standard error of estimate $=0.00569$ degree
$\ddot{\lambda}$ (measured, with minimum standard error) $=-2.2024 \times 10^{-5} \mathrm{rad} / \mathrm{sid}^{\text {. day }}{ }^{2}$, at $\lambda=-60.91^{\circ}$ on $\mathrm{t}^{\prime}=384.6346$ January 1963
$\ddot{\lambda}($ theoretical, from Equation 2$)=-2.2330 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, for $\mathrm{a}_{\mathrm{s}}=6.6116$ earth radii, $\mathrm{i}_{\mathrm{s}}=32.84^{\circ}, \lambda=-60.91^{\circ}, \mathrm{J}_{22}-\mathrm{J}_{31}$ as noted.
Estimate of acceleration bias in arc $\mathrm{S} 2 / 1=\ddot{\lambda}($ theoretical $)-\ddot{\lambda}($ measured $)=-0.0306 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$


Figure 2-Measured and simulated orbit data at ascending Equator crossings in arc 2 (Syncom 2).

We distinguish two drift regimes for the 24 -hour satellite on the basis of mean drift rate as discussed in the introduction. If the mean drift rate in an arc is less than $\pm 0.1$ degree/day, the ground track of the synchronous satellite is somewhat arbitrarily said to be in a slow drift regime. In this regime, it has been found in practice (and reconciled theoretically in Appendix D; see also Reference 3) that the mean geographic longitude can be expressed as a low degree polynomial in the time with constant coefficients related to the strength and orientation of the underlying longitude gravity field. If the mean drift rate is greater than $\pm 0.1$ degree/day, the satellite is said to be in a "fast drift" regime. In this regime the energy, or first, integral of Equation 2 has been found to give the best representation of the motion as a function of the underlying longitude gravity field. In particular, it has been found that for a realistic earth, if the drift excursion in the arc is limited to about $50^{\circ}$, only the second order $\mathrm{H}_{22}$ harmonic terms in Equation 2 need be retained for a sufficiently accurate representation. With this limitation, the energy integral can be shown to be expressable as

$$
\begin{equation*}
(\dot{\lambda})^{2}=c_{0}+c_{22} F\left(i_{s}, a_{s}\right)_{22} \cos 2 \lambda+S_{22} F\left(i_{s}, a_{s}\right)_{22} \sin 2 \lambda \mathrm{rad} / \mathrm{sid} . \mathrm{day}^{2} \tag{8}
\end{equation*}
$$

where

$$
C_{22}=J_{22} \cos 2 \lambda_{22}
$$

and

$$
\begin{equation*}
S_{22}=J_{22} \sin 2 \lambda_{22} . \tag{9}
\end{equation*}
$$

Equation 9 gives the dependence of the determinable constants in Equation 8 in terms of the strength $\left(J_{22}\right)$ and orientation $\left(\lambda_{22}\right)$ of the $\left(\mathrm{H}_{2}\right)$ earth gravity harmonic. The orbit constant $F\left(i_{s}, a_{s}\right){ }_{22}$ is given as

$$
\begin{equation*}
\mathbf{F}\left(i_{s}, a_{s}\right)_{22}=18\left[\pi\left(\cos i_{s}+1\right) / a_{s}\right]^{2}, \tag{10}
\end{equation*}
$$

where $a_{s}$ is in units of earth radii (see Equations 76 and 77 in Reference 2).

In References 5 and 6 it was shown that in the fast drift regime, for excursions of about 10 days or $10^{\circ}$ or less, sufficient accuracy is maintained if the drift rate is calculated simply from the difference of successive longitudes and assigned to the midlongitude. Unfortunately, in arc 3 the orbit determination appears to be so poor, or the record so brief, that utilization of only a single Equator crossing in each determined orbit has not proved adequate to give even one well determined acceleration for this arc. Some improvement in acceleration discrimination has resulted from the utilization of successive crossings in each orbit to provide additional independent drift velocity data. (See the discussion in Reference 6.) But even with this extra data and also with the use of an estimated crossing just prior to the jet pulsing on 25 April initiating arc 4 , the best standard error in the acceleration for this arc is close to $100 \%$ of the measured value. The details of this drift are presented in Table 3. The principal results of the acceleration analysis on the actual data in Table 3 are listed in Table 10 in the next section. The data was so inconclusive that the "best" measured acceleration in this arc was ignored in the final gravity synthesis.

## Arc 4, Syncom 2, 25 April 1964-4 July 1964

On 25 April 1964, on-board jet pulsing slowed the westward drift of Syncom 2 from - 1.3 degree/ day to -0.8 degree/day. The ascending Equator crossing at this time was at $116^{\circ} \mathrm{W}$. The orbit inclination was about $32.6^{\circ}$. From 25 April to 7 July 1964, the "figure of $8^{\prime \prime}$ ground track of Syncom 2 moved from $116^{\circ} \mathrm{W}$ to $164^{\circ} \mathrm{W}$ in free gravity drift. The principal effect in this two month arc was a deceleration of the drift rate, due to resonant earth longitude gravity, from $\mathbf{- 0 . 8 1}$ degree/day to -0.75 degree/day. The drift regime is "fast" and the details of the long term acceleration analysis on the measured and simulated ascending Equator crossings according to Equation 8 are presented in Table 4 and Figure 3 and summarized in Tables 10 and 11.

A number of gravity drift simulated trajectories for this arc were calculated with initial semimajor axis, Equator crossing longitude and earth longitude gravity as variables. The results of these (Tables $4 S$ and $4 S / 1$ ) were only fairly conclusive as to the exact magnitude of long term acceleration bias due to sun and moon gravity and model error in this arc. It would seem that the

Table 3
Syncom 2 Osculating Elements at the First Ascending Equator Crossings Past the Tracking Epoch and Related Data for Free Drift Arc 3. ${ }^{\dagger}$

| Orbit <br> Number <br> 3 - | Tracking Epoch (yr-mo-day-hr UT) | Semimajor Axis, a (earth radii) | $\begin{gathered} \text { Inclination, } \\ \text { i } \\ \text { (degrees) } \end{gathered}$ | Right <br> Ascension <br> of the <br> Ascending <br> Node <br> (degrees) | Eccentricity | Argument of Perigee (degrees) | j | Mean <br> Anomaly <br> (degrees) | $\begin{gathered} \text { Time from } \\ 1964.0, \\ \mathrm{~T} \\ \text { (days) } \end{gathered}$ | $\begin{gathered} \Delta \mathrm{T}= \\ \mathrm{T}_{j+1}-\mathrm{T}_{\mathrm{j}} \\ \text { for }{ }^{[\text {bracketed }]} \\ \text { Data: } \\ \mathrm{T}_{\mathrm{j}}^{\prime}-\mathrm{T}_{\mathrm{j}} \\ \text { (days) } \end{gathered}$ | (1) <br> Longitude of the Ascending Equator Crossing, $\lambda$ (degrees) | $\begin{gathered} \Delta \lambda= \\ \lambda_{j+1}-\lambda_{j} \\ \text { for [bracketed] } \\ \text { Data: } \\ \eta_{i}-\lambda_{j} \\ \text { (degrees) } \end{gathered}$ | (2) <br> $\Delta \mathrm{V} / \Delta \mathrm{T}, \dot{\lambda}$ <br> (degrees/day) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 64-3-18-2.0 | 6.6266333 | 32.682 | -45.512 | . 00205 | -173.534 | 1 | 173.508 | 78.5721 |  | $\begin{array}{cc} -67.623 \\ {[-68.28]^{*}} \end{array}$ |  | [-1.3082] |
| 2 | $64-3-18-2.0$ | 6.6266800 | 32.681 | -45.529 | . 00203 | $-173.162$ | $1{ }^{\prime}$. | 173.134 | 79.5730 | 1.0009 | $\begin{gathered} -68.932 \\ (-71.55)^{* *} \end{gathered}$ | $\begin{aligned} & -1.309 \\ & -7.848 \end{aligned}$ | $(-1.3068)$ |
| 3 | 64-3-24-13.0 | 6.6266362 | 32.724 | -45.544 | . 00190 | -169.992 | 2 | 169.954 | 84.5774 | 6.0053 | $\begin{gathered} -75.471 \\ {[-76.11]} \end{gathered}$ |  | $[-1.2838]$ |
| 4 | 64-3-24-13.0 | 6.6266143 | 32.720 | -45.558 | . 00190 | $-170.431$ | $2^{\prime}$ | 170.394 | 85.5782 | 1.0008 | $\left(\begin{array}{cc} -76.756 \\ (\quad 81.31) \end{array}\right.$ | - 1.285 | $(-1.2959)$ |
| 5 | 64-4-1-22.0 | 6.6269968 | 32.685 | -45.699 | . 00204 | $-171.037$ | 3 | 171.000 | 93.5847 | 9.0073 | $\left\lvert\, \begin{gathered} -87.144 \\ {[-87.81]} \end{gathered}\right.$ | -11.673 | [-1.3339] |
| 6 | 64-4-1-22.0 | 6.6270420 | 32.684 | -45.714 | . 00204 | $-170.705$ | $3^{\prime}$ | 170.667 | 94.5856 | 1.0009 | $\left(\begin{array}{c} -88.479 \\ (-91.12) \end{array}\right.$ | - 1.335 | $(-1.3242)$ |
| 7 | 64-4-7-15.0 | 6.6269419 | 32.701 | -45.831 | . 00197 | $-172.877$ | 4 | 172.849 | 99.5900 | 6.0053 | $\begin{array}{r} -95.096 \\ {[-95.75]} \end{array}$ | - 7.952 | $[-1.3093]$ |
| 8 | 64-4-7-15.0 | 6.6269010 | 32.697 | -45.846 | . 00198 | -173.240 | $4^{\prime}$ | 173.213 | 100.5909 | 1.0009 | $\begin{aligned} & -96.406 \\ & (-99.03) \end{aligned}$ | - 1.310 | $(-1.3109)$ |
| 9 | 64-4-13-19.0 | 6.6264550 | 32.599 | -45.990 | . 00213 | -172.507 | 5 | 172.475 | 105.5950 | 6.0050 | $\begin{array}{r} -102.968 \\ {[-103.62]} \end{array}$ | - 7.872 | $[-1.3001]$ |
| 10 | 64-4-13-19.0 | 6.6264975 | $32.599$ | -46.008 | . 00212 | -172.055 | $5^{\prime}$ | 172.022 | 106.5958 | 1.0008 | $\begin{gathered} -104.269 \\ (-109.58) \end{gathered}$ | $-1.3011$ | $(-1.3215)$ |
| 11*** | 64-4-25-2.0 | $$ | 32.603 <br> Average: <br> $32.67=i_{s}$ |  |  |  | 6 |  | 115.6026 | 10.0076 | -116.193 | -13.225 |  |

$*\left(\lambda_{j}^{\prime}-\lambda_{j}\right) / 2$; for [bracketed] data.
** $\left(\lambda_{j+1}-\lambda_{j}\right) / 2$ for other data.
***Data estimated for crossing just prior to Epoch 64-4-25-2.0 hour from arc 4, orbir 4-1, and arc 3, orbir 3-9.
${ }^{\dagger}$ See notes in Table 1.
Results of least squares fit of (bracketed) data in (1) and (2) above according to the drift theory of Equation 8:
$(\dot{\lambda})^{2}=C_{0}+C_{22} F\left(i_{s}, a_{s}\right) \cos 2 \lambda+S_{22} F\left(i_{s}, a_{s}\right) \sin 2 \lambda$
$C_{22}=-(1.01 \pm 3.12) \times 10^{-6}$
$S_{22}=(5.85 \pm 6.75) \times 10^{-7}$
Standard error of estimate $=1.215 \times 10^{-5} \mathrm{rad} / \mathrm{sid}^{2}$. day ${ }^{2}$
$\ddot{\lambda}($ with minimum standard error $)=-(0.897 \pm 0.888) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, at $\lambda=-88.0^{\circ}$, (see Figure 10$)$.
difference between the acceleration biases in trajectories 4 S and $4 \mathrm{~S} / 1$ can be accounted for by the neglect, in the analysis in Table $4 \mathrm{~S} / 1$, of higher order earth longitude gravity which was responsible for the fine details of the drift. Analysis of long drift arcs (in excess of $30^{\circ}$ ) in various realistic earth longitude gravity fields has shown that the order of magnitude of this higher order earth gravity bias in an analysis with only second order gravity is about $0.03 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$.

Evidently, an accurate assessment of the total model bias in a long, fast drift arc can only be made in a numerical analysis which includes at least some third order gravity effects. The most accurate discrimination of the model bias in such a regime would involve a third order reduction of the drift as well as its inclusion in the trajectory generator. Unfortunately, the parallel analysis of the actual orbits with a third order energy integral (similar to Equation 8) appears to lose considerable acceleration discrimination over the second order analysis of the third order trajectories in arcs 4 and 5 , due to the large observation errors present in the limited actual data.

After many techniques were tried, the straightforward numerical evaluation of the model bias in the gravity experiment on the data presented here seems to be adequate to the observational precision of that data. Other more analytical and more lengthy iterative techniques have been used by Allan (Reference 12) to evaluate sun, moon, and model bias in Syncom 2 drift data (see Discussion.) The present analysis on the limited nine arc record points to the conclusion that the actual unadjusted accelerations as a whole are, if anything, better measures of earth gravity effects, more precise than bias adjusted measurements on any basis (see Table 12.) Perhaps in the future, when a greater proportion of augmentation (rather than cancellation) of error data is received and processed, it will prove more than academic to examine these bias removal techniques with thoroughness.

## Arc 5, Syncom 2, 4 July 1964-19 February 1965

On 4 July 1965 the westward drift of Syncom 2, at a mean longitude of $171^{\circ} \mathrm{W}$, was slowed from -0.75 degree/day to -0.5 degree/day by ground commanded on-board jet pulsings. For the next $7-1 / 2$ months the satellite, as far as can be determined, drifted freely in the gravity fields of the earth, sun, and moon. The details of this drift as derived from orbits for Syncom 2 determined at GSFC are presented in Table 5 and Figure 4. Tables $5 S$ and $5 S / 1$ give the results of closely paralleling simulated trajectories numerically calculated in the presence of earth longitude gravity through second and third order. The second order trajectory uses gravity constants which were derived from earlier, more limited Syncom 2 data (Reference 3). The constants in the third order trajectory represent a best estimate at an intermediate stage in this analysis. The $\mathrm{J}_{31}, \lambda_{31}$ constants in these "ITEM" computed trajectories are best estimates in a private communication from W. M. Kaula in October 1964.

Greatest precision and fidelity to true earth effects at this stage in the reduction of the actual "noisy" data appears to be preserved when a single acceleration is calculated from a second order longitude gravity model extending over as short an arc as feasible. Model bias effects appear to be most accurately assessed by paralleling these reductions on a third order numerically calculated trajectory which closely follows the actual drift.

Syncom 2 Osculating Elements at the First Ascending Equator Crossing Past the Tracking Epoch, and Related Data for Free Drift Arc 4.*

| Orbit <br> Number <br> 4 - | Syncom 2 <br> Tracking Epoch <br> I (yr-mo-day-hr UT) | Semimajor Axis, a (earth radii) | $\begin{gathered} \text { Inclination, } \\ \text { i } \\ \text { (degrees) } \end{gathered}$ | Right <br> Ascension of the Ascending Node (degrees) | Eccentricity | Argument of Perigee (degrees) | Mean <br> Anomaly <br> (degrees) | Time from 1964.0, t (days) | $\begin{gathered} \Delta t= \\ t_{j+1}-t_{j} \\ \text { (days) } \end{gathered}$ | (1) <br> Longitude of the <br> Ascending Equator Crossing, $\lambda$ (degrees) | $\begin{gathered} \Delta \lambda= \\ \lambda_{j+1}-\lambda_{j} \\ \text { (degrees) } \end{gathered}$ | (2) $\begin{gathered} \frac{\Delta \lambda}{\Delta t}, \dot{\lambda} \\ \text { (degrees/day) } \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 64-4-25-2.0 | 6.6208880 | 32.602 | -46.131 | . 00120 | -168.118 | 168.090 | $\begin{gathered} 116.6039 \\ (118.6030) \end{gathered}$ |  | $\begin{aligned} & -117.182 \\ & (-118.81)^{* *} \end{aligned}$ |  | (-0.8134) |
| 2 | 64-4-28-15.0 | 6.6208306 | 32.594 | -46.142 | . 00116 | -163.317 | 163.279 | $\begin{gathered} 120.6020 \\ (124.1001) \end{gathered}$ | 3.9981 | $\begin{gathered} -120.434 \\ (-123.26) \end{gathered}$ | -3.252 | $(-0.8086)$ |
| 3 | 64-5-5-16.0 | 6.6206780 | 32.560 | -46.274 | . 00117 | -165.263 | 165.229 | $\begin{gathered} 127.5982 \\ (131.0963) \end{gathered}$ | 6.9962 | $\begin{aligned} & -126.091 \\ & (-128.91) \end{aligned}$ | -5.657 | $(-0.8072)$ |
| 4 | 64-5-12-16.0 | 6.6206500 | 32.626 | -46.424 | . 00118 | -159.744 | 159.697 | $\begin{gathered} 134.5943 \\ (137.5927) \end{gathered}$ | 6.9961 | $\begin{aligned} & -131.738 \\ & (-134.11) \end{aligned}$ | -5.647 | $(-0.7926)$ |
| 5 | 64-5-19-14.0 | 6.6205248 | 32.577 | -46.452 | . 00113 | -163.604 | 163.567 | $\begin{gathered} 140.5910 \\ (144.0889) \end{gathered}$ | 5.9967 | $\begin{gathered} -136.491 \\ (-139.25) \end{gathered}$ | -4.753 | $(-0.7898)$ |
| 6 | 64-5-25-15.0 | 6.6205311 | 32.600 | -46.614 | . 00122 | -162.158 | 162.115 | $\begin{gathered} 147.5867 \\ (151.5844) \end{gathered}$ | 6.9957 | $\begin{gathered} -142.016 \\ (-145.13) \end{gathered}$ | -5.525 | $(-0.7791)$ |
| 7 | 64-6-2-21.0 | 6.6204625 | 32.576 | -46.678 | . 00123 | -163.458 | 163.418 | $\begin{gathered} 155.5820 \\ (159.0798) \end{gathered}$ | 7.9953 | $\begin{gathered} -148.245 \\ (-150.94) \end{gathered}$ | -6.229 | $(-0.7711)$ |
| 8 | 64-6-9-21.0 | 6.6198435 | 32.580 | -46.763 | . 00119 | -159.072 | 159.023 | $\begin{gathered} 162.5776 \\ (166.0753) \end{gathered}$ | 6.9956 | $\begin{gathered} -153.639 \\ (-156.30) \end{gathered}$ | -5.394 | $(-0.7604)$ |
| 9 | 64-6-16-15.0 | 6.6199213 | 32.565 | -46.874 | . 00118 | -162.909 | 162.870 | $\begin{gathered} 169.5729 \\ (173.0705) \end{gathered}$ | 6.9953 | $\begin{gathered} -158.958 \\ (-161.59) \end{gathered}$ | -5.319 | $(-0.7518)$ |
| 10 | 64-6-23-15.0 | 6.6201029 Average: $6.6204433=a_{\mathrm{S}}$ | $\begin{aligned} & 32.562 \\ & \hline \text { Average: } \\ & 32.584=\mathrm{i}_{\mathrm{s}} \end{aligned}$ | -46.991 | . 00122 | -161.998 | 161.955 | 176.5680 | 6.9951 | -164.217 | -5.259 |  |

Results of least squares fit of (bracketed) data in (1) and (2) above according to the drift theory of Equation 8:
$(\dot{\lambda})^{2}=C_{0}+C_{22} F\left(i_{5}, a_{8}\right) \cos 2 \lambda+S_{22} F\left(i_{s}, a_{s}\right) \sin 2 \lambda$
$C_{22}=-(1.4347 \pm 0.0788) \times 10^{-6}$
$S_{22}=(0.8114 \pm 0.2823) \times 10^{-6}$
Standard error of estimate $=1.152 \times 10^{-6} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$
$\lambda$ (with mimimum standard error) $=(2.138 \pm 0.0842) \times 10^{-5} \mathrm{rad} / \mathrm{sid}^{2}$ day ${ }^{2}$, at $\lambda=-140.00^{\circ}$ (see Figure 3)
*See notes in Table 1.
** $\Delta \lambda / 2$ for (bracketed) longitudes

Table 4S-A
Osculating Elements at the First Equator Crossing Past the Syncom 2 Tracking Epoch, and Related Data in Two Simulated Syncom 2 Arc 4 Trajectories with Earth Longitude Gravity through Second Order.*

|  |  |  |  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Orbit <br> Number 4S-A - | Syncom 2 <br> Tracking Epoch (yr-mo-day-hr UT) | Semimajor Axis, a (earth radii) | Inclination, i (degrees) | Right <br> Ascension of the Ascending Node (degrees) | Eccentricity | Argument of Perigee (degrees) | Mean <br> Anomaly <br> (degrees) | $\begin{gathered} \text { Time from } \\ 1964.0, \\ \mathrm{t} \\ \text { (days) } \end{gathered}$ | $\begin{gathered} \Delta t=\underset{\text { (days) }}{t_{j+1}}-t_{j} \\ \end{gathered}$ | Longitude of the Ascending Equator Crossing, $\lambda$ (degrees) | $\begin{gathered} \Delta \lambda=\lambda_{j+1}-\lambda_{j} \\ \text { (degrees) } \end{gathered}$ | (2) <br> $\Delta v / \Delta t, \dot{\lambda}$ <br> (degrees/day) |
| 1 | 64-4-25-2.0 | 6.6206462 | 32.602 | -46.131 | . 00120 | -168.120 | 168.093 | $\begin{gathered} 116.6039 \\ (118.6029) \end{gathered}$ |  | $\begin{gathered} -117.182 \\ (-118.81)^{* *} \end{gathered}$ |  | (-.8167) |
| 2 | 64-4-28-15.0 | 6.6205913 | 32.601 | -46.201 | . 00123 | -166.322 | 166.289 | $\begin{gathered} 120.6019 \\ (124.1001) \end{gathered}$ | 3.9980 | $\begin{array}{r} -120.447 \\ (-123.29) \end{array}$ | -3.265 | $(-.8138)$ |
| 3 | 64-5-5-16.0 | 6.6207620 | 32.592 | -46.293 | . 00120 | -165.756 | 165.722 | $\begin{gathered} 127.5983 \\ (131.0964) \end{gathered}$ | 6.9964 | $\begin{gathered} -126.141 \\ (-128.97) \end{gathered}$ | -5.694 | (-.8087) |
| 4 | 64-5-12-16.0 | 6.6204145 | 32.582 | -46.404 | . 00127 | -164.655 | 164.617 | $\begin{gathered} 134.5945 \\ (137.5929) \end{gathered}$ | 6.9962 | $\begin{array}{r} -131.799 \\ (-134.21) \end{array}$ | -5.658 | $(-.8046)$ |
| 5 | 64-5-19-14.0 | 6.6205767 | 32.572 | -46.495 | . 00115 | -167.061 | 167.031 | $\begin{gathered} 140.5913 \\ (144.0893) \end{gathered}$ | 5.9968 | $\begin{gathered} -136.624 \\ (-139.40) \end{gathered}$ | -4.825 | $(-.7940)$ |
| 6 | 64-5-25-15.0 | 6.6202807 | 32.566 | -46.613 | . 00122 | -165.700 | 165.665 | $\begin{gathered} 147.5872 \\ (151.5849) \end{gathered}$ | 6.9959 | $\begin{array}{r} -142.179 \\ (-145.320 \end{array}$ | -5.555 | (-.7856) |
| 7 | 64-6-2-21.0 | 6.6203666 | $32.557$ | $-46.711$ | $.00119$ | -165.422 | 165.387 | $\begin{gathered} 155.5825 \\ (159.0803) \end{gathered}$ | 7.9953 | $\begin{aligned} & -148.460 \\ & (-151.18) \end{aligned}$ | $-6.281$ | $(-.7765)$ |
| 8 | 64-6-9-21.0 | 6.6201344 | 32.551 | -46.829 | . 00122 | -163.006 | 162.966 | $\begin{gathered} 162.5781 \\ (166.0759) \end{gathered}$ | 6.9956 | $\begin{array}{r} -153.892 \\ (-156.58) \end{array}$ | -5.432 | $(-.7682)$ |
| 9 | 64-6-16-16.0 | $6.6200798$ | $32.535$ | -46.914 | . 00115 | -166.916 | 166.887 | $\begin{gathered} 169.5736 \\ (173.0713) \end{gathered}$ | 6.9955 | $\begin{aligned} & -159.266 \\ & (-161.91) \end{aligned}$ | -5.374 | $(-.7571)$ |
| 10 | 64-6-23-15.0 | $\begin{aligned} & \frac{6.6199550}{\text { Average: }} \\ & 6.6203807=a_{s} \end{aligned}$ | $\begin{array}{\|l} 32.533 \\ \hline \text { Average: } \\ 32.57=\mathrm{i}_{\mathrm{s}} \end{array}$ | -47.020 | . 00120 | -164.171 | 164.134 | 176.5689 | 6.9953 | -164.562 | -5.296 |  |

*Computed by ITEM with gravity constants the same as in Table A1 with the addition of earth constants: $\mathrm{J}_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18.0^{\circ}$.
** $\Delta \lambda / 2$ for (bracketed) longitudes.
Results of least squares fit of (bracketed) data for arc S4-A in (1) and (2) above according to the drift theory of Equation 8:
$(\dot{\lambda})^{2}=C_{0}+C_{22} F\left(i_{2}, a_{s}\right) \cos 2 \lambda+S_{22} F\left(i_{s}, a_{s}\right) \sin 2 \lambda$
$C_{0}=1.7987 \times 10^{-4} \mathrm{rad}^{2} / \mathrm{sid} . \mathrm{day}^{2}$
$C_{22}=-(1.3234 \pm 0.0390) \times 10^{-6}$
$S_{22}=(1.0770 \pm 0.1372) \times 10^{-6}$
Standard error of estimate $=5.685 \times 10^{-7} \mathrm{rad}^{2} / \mathrm{sid}$. day ${ }^{2}$
$\ddot{\lambda}$ (measured, at $\left.\lambda=-140.0^{\circ}\right)=(2.051 \pm 0.042) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day $^{2}$
$\ddot{\lambda}$ (theoretical, from Equation 2) $=2.078 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, where $a_{s}=6.6203807$ earth radii, $i_{s}=32.57^{\circ}, \lambda=-140.0^{\circ}, J_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18.0^{\circ}$.
Bias $($ theoretical-measured $)=2.078 \times 10^{-5}-2.051 \times 10^{-5}=+0.027 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$

Osculating Elements at the First Equator Crossing Past the Syncom 2 Tracking Epoch, and Related Data in Two Simulated Syncom 2 Arc 4 Trajectories with Earth Longitude Gravity through Second Order.*

| Orbit <br> Number <br> 4S-B - | Syncom 2 <br> Tracking Epoch (yr-mo-day-hr UT) | Semimajor Axis, a (earth radii) | $\begin{gathered} \text { Inclination, } \\ \text { i } \\ \text { (degrees) } \end{gathered}$ | Right <br> Ascension of the Ascending Node (degrees) | Eccentricity | Argument of Perigee (degrees) | Mean Anomaly (degrees) | Time from 1964.0, t (days) | $\begin{gathered} \Delta \mathrm{t}=\mathrm{t}_{\mathrm{i}+1}-\mathrm{t}_{\mathrm{i}} \\ \text { (days) } \end{gathered}$ | (1) <br> Longitude of the Ascending Equator Ctossing, $\lambda$ (degrees) | $\begin{gathered} \Delta \lambda=\lambda_{i+1}-\lambda_{i} \\ \text { (degrees) } \end{gathered}$ | (2) $\Delta \lambda / \Delta t, \dot{\lambda}$ (degrees/day) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 64-4-25-2.0 | 6.6206118 | 32.602 | -46.131 | . 00120 | -163.674 | 163.635 | 116.6039 |  | $\begin{gathered} -117.160 \\ (-118.79) \end{gathered}$ |  | (-.8139) |
| 2 | 64-4-28-15.0 | 6.6205573 | 32.601 | -46.201 | . 00124 | -161.997 | 161.953 | 120.6018 | 3.9979 | $\begin{gathered} -120.414 \\ (-123.25) \end{gathered}$ | -3.254 | $(-.8110)$ |
| 3 | 64-5-5-16.0 | 6.6207281 | 32.592 | -46.292 | . 00120 | -161.312 | 161.268 | 127.5981 | 6.9963 | $\begin{aligned} & -126.088 \\ & (-128.91) \end{aligned}$ | -5.674 | $(-.8060)$ |
| 4 | 64-5-12-16.0 | 6.6204353 | 32.582 | -46.420 | . 00125 | -159.660 | 159.611 | 134.5943 | $6.9962$ | $\begin{aligned} & -131.727 \\ & (-134.13) \end{aligned}$ | -5.639 | $(-.8019)$ |
| 5 | 64-5-19-14.0 | 6.6205429 | 32.572 | -46.495 | . 00115 | -162.421 | 162.381 | 140.5910 | 5.9967 | $\begin{gathered} -136.536 \\ (-139.30) \end{gathered}$ | -4.809 | $(-.7913)$ |
| 6 | 64-5-25-15.0 | 6.6202474 | 32.566 | -46.613 | . 00122 | -161.330 | 161.285 | 147.5869 | 6.9959 | $\begin{gathered} -142.072 \\ (-145.20) \end{gathered}$ | -5.536 | (-.7827) |
| 7 | 64-6-2-21.0 | 6.6203333 | 32.557 | -46.711 | . 00119 | -160.950 | 160.906 | 155.5821 | 7.9952 | $\begin{gathered} -148.330 \\ (-151.04) \end{gathered}$ | -6.258 | (-.7739) |
| 8 | 64-6-9-21.0 | 6.6201010 | 32.551 | -46.829 | . 00122 | -158.629 | 158.578 | 162.5777 | 6.9956 | $\begin{gathered} -153.744 \\ (-156.42) \end{gathered}$ | -5.414 | (-.7653) |
| 9 | 64-6-16-15.0 | 6.6200466 | 32.535 | -46.914 | . 00115 | -162.291 | 162.252 | 169.5732 | 6.9955 | $\begin{gathered} -159.098 \\ (-161.74) \end{gathered}$ | -5.354 | (-.7544) |
| 10 | 64-6-23-15.0 | 6.6199205 <br> Average: <br> $6.6203524=a_{S}$ | $\frac{32.533}{\text { Average: }} \begin{aligned} & 32.57=\mathrm{i}_{\mathrm{s}} \end{aligned}$ | -47.020 | . 00120 | -159.742 | 159.695 | 176.5684 | 6.9952 | -164.375 | -5.277 |  |

*See notes in Table 4S-A.
Results of least squares fit of (bracketed) data for arc 4 S - B in (1) and (2) above according to the drift theory of Equation 8:
$(\dot{\lambda})^{2}=C_{0}+C_{22} F\left(i_{s}, a_{s}\right) \cos 2 \lambda+S_{22} F\left(i_{s}, a_{s}\right) \sin 2 \lambda$
$C_{22}=-(1.3250 \pm 0.0388) \times 10^{-6}$
$S_{22}=(1.0787 \pm 0.1378) \times 10^{-6}$
Standard error of estimate $=5.669 \times 10^{-7} \mathrm{rad}^{2} /$ sid. day $^{2}$
$\tilde{\lambda}$ (measured, at $\left.\left.\lambda=-140.0^{\circ}\right)=(2.054) \pm 0.0413\right) \times 10^{-5} \mathrm{rad} / \mathrm{sid}^{2}$ day $^{2}$
$\ddot{\lambda}$ (theoretical, from Equation 2 ) $=2.078 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, where $\mathrm{a}_{s}=6.620352$ earth radii, $\mathrm{i}_{\mathrm{s}}=32.57^{\circ}, \lambda=140.0^{\circ}, \mathrm{J}_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18.0^{\circ}$.
Bias (theoretical-measured) $=2.078 \times 10^{-5}-2.054 \times 10^{-5}=+0.024 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$.

## Table 4S/1

Osculating Elements at the First Ascending Equator Crossing Past the Syncom 2 Tracking Epoch, and Related
Data in a Simulated Syncom 2 Arc 4 Trajectory with Earth Longitude Gravity through Third Order.*

| Orbit Number 4S/1 - | Tracking Epoch (yr-mo-day-hr UT) |  | Inclination, i (degrees) | $\begin{gathered} \text { Time from } \\ 1964.0, \\ \mathbf{t} \\ \text { (days) } \end{gathered}$ | $\left\lvert\, \begin{gathered} \Delta t=t_{i+1}-\mathrm{t}_{\mathrm{i}} \\ \text { (days) } \end{gathered}\right.$ | (1) <br> Longitude of the Ascending Equator Crossing, $\lambda$ (degrees) | $\begin{gathered} \Delta \lambda=\lambda_{j+1}-\lambda_{j} \\ \text { (degrees) } \end{gathered}$ | (2) <br> $\Delta \lambda / \Delta t, \dot{\lambda}$ (degrees/day) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 64-4-25-2.0 | 6.6206454 | 32.602 | 116.6039 |  | -117.182 |  |  |
|  |  |  |  |  |  | (-118.81)** |  | (-.8167) |
| 2 | 64-4-28-15.0 | 6.6205886 | 32.601 | 120.6019 | 3.9980 | $\begin{aligned} & -120.447 \\ & (-123.29) \end{aligned}$ | -3.265 | $(-.8134)$ |
| 3 | 64-5-5-16.0 | 6.6207557 | 32.592 | 127.5983 | 6.9964 | $\begin{aligned} & -126.138 \\ & (-128.96) \end{aligned}$ | -5.691 | $(-.8080)$ |
| 4 | 64-5-12-16.0 | 6.6204594 | 32.583 | 134.5945 | 6.9962 | $-131.791$ | -5.653 |  |
| 5 | 64-5-19-14.0 | 6.6205639 | 32.572 | 140.5912 | 5.9967 | $\begin{array}{r} (-134.20) \\ -136.611 \end{array}$ | -4.820 | (-.8038) |
|  |  |  |  |  |  | $(-139.38)$ | . 8.80 | (-.7929) |
| 6 | 64-5-25-15.0 | 6.6202643 | 32.566 | 147.5871 | 6.9959 | $\begin{aligned} & -142.158 \\ & (-145.29) \end{aligned}$ | -5.547 | (-.7840) |
| 7 | 64-6-2-21.0 | 6.6203440 | 32.557 | 155.5824 | 7.9953 | $\begin{aligned} & -148.426 \\ & (-151.13) \end{aligned}$ | -6.268 | (-.7743) |
| 8 | 64-6-9-21.0 | 6.6201058 | 32.551 | 162.5780 | 6.9956 | $\begin{gathered} -153.843 \\ (-156.52) \end{gathered}$ | -5.417 | (-.7655) |
| 9 | 64-6-16-15.0 | 6.6200444 | 32.535 | 169.5734 | 6.9954 | $\begin{array}{r} -159.198 \\ (-161.84) \end{array}$ | -5.355 | (-.7539) |
| 10 | 64-6-23-15.0 | 6.6199117 | 32.533 | 176.5686 | 6.9952 | -164.472 | -5.274 | (-.7539) |
|  |  | Average: $6.6204=\mathrm{a}_{\mathrm{s}}$ | Average: $32.57=\mathrm{i}_{\mathrm{S}}$ |  |  |  |  |  |

*Gravity constants of this trajectory computed by ITEM are the same as those in Table A1, with the addition of these earth constants:
$\mathrm{J}_{22}=-1.8 \times 10^{-6}, \lambda_{22}=-15.35^{\circ}$
$J_{33}=-0.16 \times 10^{-6}, \lambda_{33}=24.0^{\circ}$
$\mathrm{J}_{31}=-1.5 \times 10^{-6}, \lambda_{31}=0.0^{\circ}$
See Figure B1 for the significance of these constants. The initial elements of this trajectory, aside from those listed for orbit $4 \mathrm{~S} / 1-1$, are the same as those in orbit S4A-1 (Table 4S).
** $\left(\lambda_{i+1}-\lambda_{j} / 2\right)$ for (bracketed) longitudes
Results of least squares fit of the (bracketed) data in (1) and (2) above according to the theory of Equation 8:
$(\dot{\lambda})^{2}=C_{0}+C_{22} F\left(i_{s}, a_{s}\right) \cos 2 \lambda+S_{22} F\left(i_{s}, a_{s}\right) \sin 2 \lambda_{22}$
$\mathrm{C}_{0}=1.786 \times 10^{-4} \mathrm{rad}^{2} /$ sid. day ${ }^{2}$
$C_{22}=-(1.3905 \pm 0.0386) \times 10^{-6}$
$\mathrm{S}_{22}=(1.1362 \pm 0.1362) \times 10^{-6}$
Standard error of estimate $=5.627 \times 10^{-7} \mathrm{rad}^{2} / \mathrm{sid}^{\text {day }}{ }^{2}$
$\ddot{\lambda}($ with minimum standard erros $)=2.156 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, measured, at $\lambda=-140.0^{\circ}$
$\ddot{\lambda}$ (theoretical, from Equation 2) $=2.163 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, where $\mathrm{a}_{\mathrm{s}}=6.6204$ earth radii, $\mathrm{i}_{\mathrm{s}}=32.57^{\circ}, \lambda=-140.0^{\circ}$, and $\mathrm{J}_{22}-\mathrm{J}_{31}$ as noted
Estimate of acceleration bias at $\lambda=-140.0^{\circ}$ in arc S4/1 $=\ddot{\lambda}($ theoretical $)-\ddot{\lambda}($ actual $)=+0.007 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \mathrm{day}^{2}$.


Figure 3-Measured and simulated orbit data at ascending Equator crossings in arc 4 (Syncom 2).

In line with this reasoning and also with the object in mind of obtaining as wide and precise a longitude survey as possible with the single-acceleration reduction technique, arc 5 and one simulated arc 5 trajectory were split up into 18 consecutive sub-arcs. (See Tables 10 and 11 in the next section). The first such sub-arc spans the first 9 (bracketed) longitude-drift rate data points in Tables 5 and $5 \mathrm{~S} / 1$. Each succeeding overlapped sub-arc drops the leading point of the previous one and adds a new point consecutively until the entire arc is covered. Each of these sub-arcs are analyzed for acceleration in exactly the same way as in arc 4, except as noted below.

Between 17 November 1964 and 10 January 1965 no orbits for Syncom 2 were calculated. The range and range rate transponder on board the satellite was not used then to conserve battery power, as the satellite was spending a considerable time in the earth's shadow in this period. As a result, the first few new orbits in January 1965 were particularly ill determined. In addition extensive data testing has shown that a few orbits in November 1964 also gave poorly defined single Equator crossing information.

On the other hand, it appeared that such suspect single crossing orbits gave reasonably good drift rates from successive crossings. This is equivalent to saying that the satellite position at epoch was poorly determined for these orbits ( 21,24 , and 27 in Table 5 ), but the semimajor axis was reasonably well determined. The reverse is the usual situation. In Table 5 it will be seen that the orbits mentioned above have been utilized for drift rate information without association with neighboring orbits. While the mean crossing longitudes for these are probably in error by about $0.2^{\circ}$, it is sufficiently accurate in this reduction where rate information demands the greater precision.

It is regretted that the standard errors in the semimajor axes and longitude locations reported by GSFC for these and other Syncom 2 and Syncom 3 orbits did not always agree with the evident errors revealed by the full arc analysis and simulations in this investigation. An effort was made to get better (smoother) information by using the complete record in these arcs without exception. Various a priori data weighting schemes were tried on the basis of these reported errors without noticeable gain in accuracy. Simple data rejection on the basis of a $3 \sigma$ criteria, after a trial reduction, was used in this and other arcs to arrive at the final acceptable data record.

In the future, further smoothing of the data in all the Syncom arcs may be possible by a separate analysis of the semimajor axis drift according to the resonant formulations in References 2 and 3. (See Discussion).

## Arc 6, Syncom 3, 31 October 1964-21 December 1964

Syncom 3, the first geostationary satellite was launched in August 1964 and reached station in the Pacific over the International Date Line in October. From 31 October to 21 December 1964, the satellite was permitted to drift from $180^{\circ}$ to $178^{\circ}$ at less than 0.1 degree/day westward without correction maneuvers. The details of this free gravity accelerated drift are found in Table 6 and Figure 5 and summarized in Tables 10 and 11 in the next section.

The actual data analysis for acceleration of the geostationary or slow moving Syncom 3 drift in this arc follows the $t^{3}$ fit technique used in arcs 1 and 2 for Syncom 2. While the inertial location of Syncom 3's Equator crossing (determined by the crossing time) was poorly defined (the orbit being nearly equatorial), the geographic location had good definition in this arc since the subsatellite point was nearly stationary. This is brought out best by the low value of the standard error of the estimate of the Equator crossing drift under a $t^{3}$ formulation, compared to the Syncom $2 \operatorname{arcs} 1$ and 2.

The poor definition of the inertial location of Syncom 3 in arc 6 has caused some difficulty in finding an acceptably close parallel simulated trajectory. Comparison of the results of the two simulations in Table $6 S$ shows the discrepancy even half a day can make in the sun, moon, and model bias errors. However, it is noted that in arc 6, as in arcs 1 and 2, there is good agreement in the bias results in the two simulations from identical initial elements between second order and

Table 5
Syncom 2 Osculating Elements at the First Ascending Equator Crossings Past the Tracking Epoch and Related Data in Free Drift Arc 5*.

| $\begin{gathered} \text { Orbit } \\ \text { Number } \\ 5- \end{gathered}$ | Tracking Epoch (yr-mo-day-hr UT) | $\begin{gathered} \text { Semimajor Axis, } \\ a \\ \mathbf{a} \\ \text { (earth radii) } \end{gathered}$ | $\begin{gathered} \text { Inclination, } \\ \text { i } \\ \text { (degrees) } \end{gathered}$ | Right Ascension of the Ascending Node (degrees) | Eccentricity | Argument of Perigee (degrees) | Mean Anomaly (degrees) | $\begin{gathered} \text { Time from } \\ \text { I964.0, } \\ \mathbf{t} \\ \text { (days) } \end{gathered}$ | $\begin{gathered} \Delta t=t_{t_{1}}^{\mathbf{t}_{1+1}-t_{1}}(\mathrm{days}) \end{gathered}$ | (1) <br> Longitude of the Ascending Equator Crossing, $\lambda$ (degrees) | $\begin{gathered} \Delta \lambda= \\ \lambda_{j+1}-\lambda_{1} \\ \text { (degres) } \end{gathered}$ | $\underset{\text { (degrees/day) }}{\Delta \Delta / \Delta t, \dot{\lambda}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 64-7-4-2.0 | 6.6165754 | 32.541 | -47.118 | . 00086 | -154.569 | 154.527 | 186.5599 |  | $\begin{gathered} -171.261 \\ (-171.97)^{* *} \end{gathered}$ |  | (-.4734) |
| 2 | 64-7-7-3.0 | 6.6169477 | 32.537 | -47.161 | . 00094 | -150.702 | 150.649 | 189.5555 | 2.9956 | $\begin{array}{r} -172.679 \\ (-174.36) \end{array}$ | - 1.418 | (-.4805) |
| 3 | 64-7-13-17.0 | 6.6166919 | 32.500 | -47.211 | . 00079 | -153.849 | 153.810 | 196.5455 | 6.9900 | $\begin{array}{r} -176.038 \\ (-177.91) \end{array}$ | - 3.359 | (-.4692) |
| 4 | 64-7-21-21.0 | 6.6164887 | 32.515 | -47.363 | . 00087 | -147.458 | 147.404 | 204.5337 | 7.9882 | $\begin{gathered} -179.791 \\ (178.81) \end{gathered}$ | - 3.753 | (-.4677) |
| 5 | 64-7-27-16.0 | 6.6165999 | 32.500 | -47.422 | . 00084 | -150.162 | 150.114 | 210.5249 | 5.9912 | $\begin{array}{r} 177.407 \\ (175.79) \end{array}$ | - 2.802 | (-.4624) |
| 6 | 64-8-3-17.0 | 6.6166087 | 32.476 | -47.555 | . 00081 | -145.700 | 145.648 | 217.5144 | 6.9895 | $\begin{array}{r} 174.175 \\ (172.59) \end{array}$ | - 3.232 | (-.4548) |
| 7 | 64-8-11-1.0 | 6.6163047 | 32.445 | -47.606 | . 00077 | -149.337 | 149.292 | 224.5039 | 6.9895 | $\begin{array}{r} 170.996 \\ (169.43) \end{array}$ | - 3.179 | (-.4494) |
| 8 | 64-8-17-19.0 | 6.6164211 | 32.408 | -47.700 | . 00089 | -145.420 | 145.362 | 231.4932 | 6.9893 | $\begin{array}{r} 167.855 \\ (166.27) \end{array}$ | - 3.141 | (-.4522) |
| 9 | 64-8-25-10.0 | 6.6164814 | 32.443 | -47.739 | . 00089 | -138.379 | 138.311 | 238.4828 | 6.9896 | $\begin{array}{r} 164.694 \\ (163.13) \end{array}$ | - 3.161 | (-.4466) |
| 10 | 64-9-1-10.0 | 6.6163260 | 32.397 | -47.933 | . 00080 | -144.436 | 144.382 | 245.4718 | 6.9890 | $\begin{gathered} 161.573 \\ (159.54) \end{gathered}$ | - 3.121 | (-.4521) |
| 11 | 64-9-9-14.0 | 6.6160374 | 32.372 | -48.100 | . 00090 | -145.744 | 145.686 | 254.4580 | 8.9862 | $\begin{array}{r} 157.510 \\ (156.17) \end{array}$ | - 4.063 | (-.4487) |
| 12 | 64-9-15-12.0 | 6.6165134 | 32.360 | -48.148 | . 00083 | -140.900 | 140.840 | 260.4489 | 5.9909 | $\begin{array}{r} 154.822 \\ (153.46) \end{array}$ | - 2.688 | (-.4562) |
| 13 | 64-9-22-10.0 | 6.6164083 | 32.302 | -48.256 | . 00087 | -141.927 | 141.866 | 266.4398 | 5.9909 | $\begin{array}{r} 152.089 \\ (150.48) \end{array}$ | - 2.733 | (-.4605) |
| 14 | 64-9-29-6.0 | 6.6165573 | 32.327 | -48.344 | . 00083 | -140.203 | 140.143 | 273.4294 | 6.9896 | $\begin{array}{r} 148.870 \\ (147.29) \end{array}$ | - 3.219 | (-.4534) |
| 15 | 64-10-6-5.0 | 6.6164264 | 32.321 | -48.414 | . 00081 | -135.504 | 135.440 | 280.4189 | 6.9895 | $\begin{gathered} 145.701 \\ (144.05) \end{gathered}$ | - 3.169 | (-.4733) |
| 16 | 64-10-13.0 | 6.6167308 | 32.312 | -48.565 | . 00090 | -135.824 | 135.753 | 287.4085 | 6.9896 | $\begin{array}{r} 142.393 \\ (140.50) \end{array}$ | - 3.308 | (-.4746) |
| 17 | 64-10-20-16.0 | 6.6163209 | 32.289 | -48.662 | . 00090 | -139.171 | 139.104 | 295.3969 | 7.9884 | $\begin{gathered} 138.602 \\ (137.16) \end{gathered}$ | - 3.791 | (-.4817) |
| 18 | 64-10-26-16.0 | 6.6169645 | 32.250 | -48.815 | . 00072 | -133.261 | 133.202 | 301.3881 | 5.9912 | $\begin{array}{r} 135.716 \\ (134.21) \end{array}$ | - 2.886 | (-.5012) |
| 19 | 64-11-2-5.0 | 6.6164857 | 32.249 | -48.864 | . 00079 | -139.157 | 139.097 | 307.3799 | 5.9918 | 132.713 | - 3.003 |  |
|  |  | $\begin{array}{\|l\|} \hline \text { Average: } \\ 6.6165206=\mathrm{a}_{\mathrm{S}} \\ \hline \end{array}$ | $\begin{aligned} & \text { Average: } \\ & 32.397=i_{\mathrm{s}} \end{aligned}$ | arc 5A data |  |  |  |  |  | (130.41) |  | (-.512 |
| 20 | 64-11-11-2.0 | 6.6171112 | 32.287 | -48.990 | . 00080 | -128.016 | 127.945 | 316.3677 | 8.9878 | 128.107 | - 4.606 |  |
| 21 | 64-11-17-6.0 | 6.6169067 | 32.230 | -49.018 | . 00081 | -137.326 | 137.264 | 322.3588 |  | $\begin{aligned} & 125.370 \\ & (125.1) \end{aligned}$ |  | (-.5177) |
| 22 | 64-11-17-6.0 | 6.6168678 | 32,229 | -49.038 | . 00083 | -136.946 | 136.882 | 323.3575 | 0.9987 | $\begin{gathered} 124.853 \\ (112.0)^{\dagger} \end{gathered}$ | - 0.517 | (-.5570) |
| 23 | 65-1-10-6.0 | 6.6173810 | 32.131 | -49.815 | . 00066 | -120.051 | 119.987 | 376.2941 | 59.9264 | 94.729 | -33.378 |  |
| 24 | 65-1-13-16.0 | 6.6180715 | 32.164 | -49.942 | . 00071 | -130.815 | 130.754 | 380.2891 |  | $\begin{array}{r} 92.461 \\ (92.16) \end{array}$ |  | (-.6057) |
| 25 | 65-1-13-16.0 | 6.6181446 | 32.163 | -49.954 | . 00070 | -128.832 | 128.771 | 381.2880 | 0.9989 | $\begin{array}{r} 91.856 \\ (91.45) \end{array}$ | - 0.605 | (-.5977) |
| 26 | 65-1-20-12,0 | 6.6185305 | 32.117 | -50.105 | . 00060 | -319.867 | 139.823 | 387.2814 | 10.9873 | 88.162 | -6.567 |  |
| 27 | 65-1-27-4-5.0 | 6.6181841 | 31.956 | -49.833 | . 00066 | 160.176 | -160.151 | 393.2771 |  | $\begin{array}{r} 84.089 \\ (83.79) \end{array}$ |  | (-.6087) |
| 28 | 65-1-27-4-5.0 | 6.6182309 | 31.955 | -49.845 | . 00065 | 160.915 | -160.891 | 394.2760 | 0.9989 | $\begin{aligned} & 83.48 \\ & (80.21)^{\dagger} \end{aligned}$ | - 0.608 | $(-.6127)$ |
| 29 | 65-2-16-4-5.0 | 6.6183288 | 32.111 | -50.406 | . 00061 | -127.648 | 127.592 | 413.2537 | 25.9723 | 72.248 | -15.914 |  |
| 30 | 65-2-19-23.0 | Average: $6.617425=a_{s}$ | Average: $32.16^{\circ}=i_{s}$ | arc 5B data data in orbi | (includes ts 5-17, |  |  | 417.2496 | 3.9959 | $\begin{aligned} & (71.02) \\ & 69.783^{\dagger t} \end{aligned}$ | - 2.465 | (-.6169) |
|  |  | Average: $6.617=\mathrm{a}_{\mathrm{s}}$ | Average: $32.33^{\circ}=\mathrm{i}$ | $\begin{aligned} & 18,19) \\ & \text { full are } 5 \text { da } \end{aligned}$ | ata |  |  |  |  |  |  |  |

[^0]*See notes in Table 1.
** $\left(\lambda_{j+1}-\lambda_{j}\right) / 2$ for (bracketed) Iongitudes between [bracketed] or un-[bracketed] data.
${ }^{\dagger}$ Data adjusted for long arc effect (Reference 5).
${ }^{\dagger} \dagger$ This data is from an orbit not listed in Table A1 for which no elements are available but which, when calculated at GSFC, appeared to show little mean motion effects from commanded but poorly executed jet pulsings just prior to the listed epoch.
Results of least squares fit of (bracketed) data for the full arc 5 in (1) and (2) according to the drift theory of Equation 8:
$(\dot{\lambda})^{2}=C_{0}+C_{22} F\left(i_{s}, a_{s}\right) \cos 2 \lambda+S_{22} F\left(i_{s}, a_{s}\right) \sin 2 \lambda$
$C_{0}=8.789 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day $^{2}$
$C_{22}=-(1.6261 \pm 0.0288) \times 10^{-6}$
$\mathrm{S}_{22}=(1.0500 \pm 0.0448) \times 10^{-6}$
Standard error of estimate $=1.415 \times 10^{-6} \mathrm{rad} / \mathrm{sid}$. day $^{2}$
$\ddot{\lambda}$ (with minimum standard error) $=-(2.295 \pm 0.0397) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, at $\lambda=134.0^{\circ}$, (see Figure 4)
Results of least squares fit of bracketed data for arc SA in (1) and (2) above according to the drift theory of Equation 8:
$(\dot{\lambda})^{2}=C_{0}+C_{22} F\left(i_{8}, a_{s}\right) \cos 2 \lambda+S_{22} F\left(i_{s}, a_{8}\right) \sin 2 \lambda$
$C_{0}=8.7719 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \mathrm{day}^{2}$
$C_{22}=-(1.5996 \pm 0.1465) \times 10^{-6}$
$S_{22}=(1.0667 \pm 0.1161) \times 10^{-6}$
Standard error of estimate $=1.465 \times 10^{-6} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$
$\ddot{\lambda}$ (with minimum standard error) $=-(0.1991 \pm 0.0661) \times 10^{-5} \mathrm{rad} / \mathrm{sid}^{\text {d }}$ day ${ }^{2}$, at $\lambda=161.0^{\circ}$
Results of least squares fit of bracketed data for arc $5 B$ in (1) and (2) according to the drift theory of Equation 8:
$(\dot{\lambda})^{2}=C_{0}+C_{22} F\left(i_{s}, a_{s}\right) \cos 2 \lambda+S_{22} F\left(i_{s}, a_{s}\right) \sin 2 \lambda$
$C_{22}=-(1.5746 \pm 0.1653) \times 10^{-6}$
$\mathrm{S}_{22}=(1.0523 \pm 0.1126) \times 10^{-6}$
Standard error of estimate $=1.603 \times 10^{-6} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$
$\ddot{\lambda}($ with minimum standard error $)=-(2.389 \pm 0.0724) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, at $\lambda=106.0^{\circ}$

Osculating Elements at Ascending Equator Crossings Past the Syncom 2 Tracking Epochs and Related Data in a Simulated Syncom 2 Arc 5 Trajectory with Earth Longitude Gravity through Second Order*.


[^1]* Computed by ITEM with gravity constants the same as in Table Al, with the addition of the earth constants:

$$
\begin{gathered}
\mathrm{J}_{22}=-1.68 \times 10^{-6} \\
\lambda_{22}=-18^{\circ} \\
* *\left(\lambda_{j+1}-\lambda_{j}\right) / 2 \text { for (bracketed) longitudes between [bracketed] or un-[bracketed] data. }
\end{gathered}
$$

***Data adjusted for long arc effects (Reference 5) (see Table 11 for summary of sub-arc analysis).
Results of least squares fit of bracketed data for arc $5 S-A$ in (1) and (2) above according to the drift theory of Equation 8:

$$
\begin{aligned}
(\dot{\lambda})^{2} & =C_{0}+C_{22} F\left(i_{s}, a_{s}\right) \cos 2 \lambda+S_{22} F\left(i_{s}, a_{s}\right) \sin 2 \lambda . \\
C_{0} & =8.5845 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \text { day }^{2} \\
C_{22} & =-(1.4109 \pm 0.0293) \times 10^{-6} \\
S_{22} & =(1.0378 \pm 0.0233) \times 10^{-6}
\end{aligned}
$$

Standard error of estimate $=2.939 \times 10^{-7} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$
$\ddot{\lambda}$ (measured at $\left.\lambda=161.0^{\circ}\right)=-(0.0702 \pm 0.0132) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$
$\ddot{\lambda}$ (theoretical, according to Equation 2), $=-(0.0809) \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \mathrm{day}^{2}$,
for $a_{s}=6.616475$ earth radii, $i_{s}=32.409^{\circ}, J_{22}=-1.68 \times 10^{-6}, \Lambda_{22}=-18.0^{\circ}, \lambda=161.0^{\circ}$
Estimate of measured bias in $\vec{\lambda}$ at $161.0^{\circ}$ in Syncom 2 arc 5A:

$$
\begin{aligned}
\text { Bias } & =\text { theoretical-measured } \\
& =-(0.0809) \times 10^{-5}+(0.0702) \times 10^{-5} \\
& =-(0.0107) \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \text { day }^{2}
\end{aligned}
$$

Results of least squares fit of bracketed data for arc $5 S-B$ in (1) and (2) above according to the drift theory of Equation 8:

$$
\begin{aligned}
& (\dot{\lambda})^{2}=C_{0}+C_{22} F\left(i_{s}, a_{s}\right) \cos 2 \lambda+S_{22} F\left(i_{s}, a_{s}\right) \sin 2 \lambda \\
& C_{22}=-(1.4219 \pm 0.0550) \times 10^{-6} \\
& S_{22}=(0.9296 \pm 0.0392) \times 10^{-6}
\end{aligned}
$$

Standard error of estimate $=5.152 \times 10^{-7} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$
$\ddot{\lambda}$ (measured at $\left.\lambda=106.0^{\circ}\right)=-(2.132 \pm 0.0240) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$
$\ddot{\lambda}$ (theoretical, according to Equation 2 ),$=-(2.154) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day $^{2}$,
for $a_{s}=6.61739$ earth radii, $i_{s}=32.20^{\circ}, J_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18.0^{\circ}, \lambda=106.0^{\circ}$
Estimate of measured bias in $\ddot{\lambda}$ at $\lambda=106.0^{\circ}$ in Syncom 2 arc $5 B$, due to sun-moon perturbations and $J_{22}$ model error exclusive of higher order longitude gravity:

$$
\begin{aligned}
\text { Bias } & =\text { theoretical-measured } \\
& =-(2.154) \times 10^{-5}+(2.132) \times 10^{-5} \\
& =-(0.022) \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \text { day }{ }^{2}
\end{aligned}
$$

Results of least squares fit of bracketed data for the full arc $5 S$ in (1) and (2) above according to the theory of Equation 8:

$$
\begin{aligned}
& (\dot{\lambda})^{2}=C_{0}+C_{22} F\left(i_{s}, a_{8}\right) \cos 2 \lambda+S_{22} F\left(i_{s}, a_{s}\right) \sin 2 \lambda \\
& C_{22}=-(1.3664 \pm 0.00725) \times 10^{-6}, S_{22}=(1.0008 \pm 0.0115) \times 10^{-6}
\end{aligned}
$$

Standard error of estimate $=3.566 \times 10^{-7} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$
$\ddot{\lambda}$ (measured at $\lambda=134.0$ ) $=-(19344 \pm 0.0100) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$
$\ddot{\lambda}$ (theorerical, $\left.i_{s}=32.33^{\circ}, J_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18.0^{\circ}, \lambda=134.0^{\circ}\right)=-(1,3237) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, for $\mathrm{a}_{\mathrm{s}}=6.6169$ earth radii.
Estimate of measured bias in $\ddot{\lambda}$ at $\lambda=134.0^{\circ}$ in Syncom 2, arc 5 , due to sun-moon perturbations and $J_{22}$ model error exclusive of higher order longitude gravity:

Bias $=$ Theoretical-measured

$$
=-(1.9237) \times 10^{-5}+1.9344 \times 10^{-5}=+0.0107 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \mathrm{day}^{2}
$$

Table $5 \mathrm{~S} / 1$
Osculating Elements at Ascending Equator Crossings Past the Syncom 2 Tracking Epochs and Related Data in a Simulated Syncom 2 Arc 5 Trajectory with Earth Longitude Gravity through Third Order*.


[^2]$J_{31}=-1.5 \times 10^{-6}, \lambda_{31}=0.0^{\circ}$. The initial elements of this trajectory, aside from those listed for orbit 5S/1-1, are che sarae as chose in orbit $55-1$ (Table $5 S$ ).
${ }^{* *}\left(\lambda_{j+1}-\lambda_{j}\right) / 2$ for (bracketed) longitudes between [bracketed] or ur- [bracketed] data.
**Data adjusted for long arc effects (Reference 5 )
Results of least squares fit of (bracketed) data for the full are $55 / 1$, in (1) and (2) above, according to Equation 8:
$(\dot{\lambda})^{2}=C_{0}+C_{22} F\left(i_{i}, z_{2}\right)_{22} \cos 2 \lambda+S_{22} F\left(i_{z}, a_{2}\right)_{22} \sin 2 \lambda$
$C_{0}=8.797 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \mathrm{day}^{2}, \mathrm{C}_{22}=-(1.6095 \pm 0.0108) \times 10^{-6}, \mathrm{~S}_{22}=(1.0981 \pm 0.0169) \times 10^{-6}$
Standard error of estimate $=5.318 \times 10^{-7} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, (see Table 11 for complete 18 sub atcs analysis)


Figure 4-Measured and simulated orbit data at ascending Equator crossings in arc 5 (Syncom 2).
full third order trajectories. In this case the initial elements of orbit 6S-B were chosen to give the best representation of the time characteristics (and thus the more faithful bias estimate) over the whole arc. Orbit 6S-A (with GSFC reported "mean" elements for epoch 6S-A-1) appears to give a better overall representation of the inclination and eccentricity in arc 6.

## Arc 7, Syncom 3, 14 January 1965-16 March 1965

Between 20 December 1964 and 14 January 1965 a number of orbit change maneuvers were performed on Syncom 3, repositioning the mean longitude of the satellite back near the International Date Line and increasing the orbit inclination to about $1^{\circ}$. Between 14 January and 30 January 1965 a number of attitude and inclination change maneuvers were performed which apparently had little effect on the mean motion of the satellite. Between 30 January and 16 March 1965 the control jets

Table 6
Syncom 3 Osculating Elements at the First Ascending Equator Crossing Past the Tracking Epoch and Related Data for Free Drift Arc 6*.

| Orbit Number 6 - | Tracking Epoch (hr-mo-day-hr UT) | Semimajor Axis, a (earth radii) | Inclination, i (degrees) | Right <br> Ascension of the <br> Ascending Node (degrees) | Eccentricity | Argument of Perigee (degrees) | Mean <br> Anomaly <br> (degrees) | Time from 1964.0, t (days) | Geographic Longitude of the First Ascending Equator Crossing After the TrackingEpoch $\lambda$ (degrees) | (1) Time from January $331.2059,1964$, $t$ (days) | (2) <br> Longitude of the Ascending Equator Crossing East of 179.017 ${ }^{\circ}$, L (degrees) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 64-10-31-2.0 | 6.6119343 | . 097 | -149.455 | . 00026 | - 45.737 | 45.719 | 305.9723 | 180.219 | -25.2336 | 1.202 |
| 2 | 64-11-3-13-18.0 | 6.6116163 | . 066 | 135.469 | . 00005 | - 13.946 | 13.948 | 308.7567 | 180.028 | -22.4492 | 1.011 |
| 3 | 64-11-7-8.0 | 6.6116896 | . 057 | 165.876 | . 00008 | - 23.052 | 23.052 | 312.8308 | 179.737 | -18.3751 | . 720 |
| 4 | 64-11-16-3.0 | 6.6115890 | . 038 | 106.668 | . 00012 | 151.303 | -151.297 | 321.6437 | 179.191 | - 9.5622 | . 174 |
| 5 | 64-11-24-3.0 | 6.6115203 | . 080 | 112.569 | . 00008 | 43.507 | - 43.514 | 329.6393 | 178.798 | $-1.5666$ | - . 219 |
| 6 | 64-11-30-10-35.0 | 6.6114447 | . 090 | 68.895 | . 00001 | 140.529 | -140.585 | 335.5028 | 178.509 | 4.2969 | -. 508 |
| 7 | $64-12-8-12-45.0$ | 6.6111913 | . 261 | 54.788 | . 00017 | -137.715 | 137.706 | 344.4400 | 178.171 | 13.2341 | -. 846 |
| 8 | $64-12-15-12.0$ | 6.6113693 | .127 | 74.467 | . 00003 | $-157.902$ | 157.913 | 351.4760 | 177.968 | 20.2701 | -1.049 |
| 9 | 64-12-21-9.0 | 6.6110148 | . 204 | 66.016 | . 00009 | - 8.119 | 8.118 | 356.4394 | 177.814 | 25,2335 | $-1.203$ |
|  |  | Average: $6.611474=a_{\mathrm{s}}$ | Average: $0.113=\mathrm{i}_{\mathrm{s}}$ |  |  |  |  |  |  |  |  |

*See notes in Table 1.
Results of least squares fit of data in (1) and (2) according to the theory of Equation 1:
$L=a_{1}+a_{2} t+a_{3} t^{2}+a_{4} t^{3}$
$\mathrm{a}_{1}=-(3.080 \pm 0.067) \times 10^{-1}$ degrees
$\mathrm{a}_{\mathbf{2}}=-(4.629 \pm 0.073) \times 10^{-2}$ degrees/day
$\mathrm{a}_{3}=(4.915 \pm 0.170) \times 10^{-4}$ degrees $/ \mathrm{day}^{2}$
$a_{4}=-(2.12 \pm 1.40) \times 10^{-6}$ degrees $/$ day ${ }^{3}$
$\ddot{\lambda}($ with minimum standard error $)=+(1.707 \pm 0.0591) \times 10^{-5} \mathrm{rad} / \mathrm{sid}^{\text {id }}$ day ${ }^{2}$, at $\mathrm{t}=-0.0332$ days, $\mathrm{t}^{\prime}=331.1727 \mathrm{January} 1964, \mathrm{~L}=-0.310^{\circ}, \lambda=178.707^{\circ}$. (see Figure 5)

Table 6S-A
Osculating Elements at the First Ascending Equator Crossing Past the Syncom 3 Tracking Epoch and Related Data,
in Two Simulated Syncom 3 Arc 6 Trajectories with Earth Longitude Gravity through Second Order.*

| Orbit Number 6S-A- | Syncom 3 <br> Tracking Epoch (yr-mo-day-hr- UT) | Semimajor Axis, a (earth radii) | Inclination, i (degrees) | Right <br> Ascension <br> of the <br> Ascending Node (degrees) | Eccentricity | Argument of Perigee (degrees) | Mean <br> Anomaly <br> (degrees) | Time from $\begin{gathered} \text { 1964.0, } \\ t \\ \text { (days) } \end{gathered}$ | Geographic Longitude of the First Ascending Equator Crossing After the Tracking Epoch, $\lambda$ (degrees) | $(1)$ Time From 330.9708 January 1964, $t$ (days) | (2) <br> Longitude of the Ascending Equator Crossing East of $179.003^{\circ}$, L (degrees) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 64-10-31-2.0 | 6.6118485 | . 066 | 37.760 | . 00011 | -130.250 | 130.248 | 305.4935 | 180.275 | -25.4773 | 1.272 |
| 2 | 64-11-3-13-18.0 | 6.6119239 | . 069 | 37.793 | . 00004 | -133.732 | 133.752 | 308.4860 | 180.053 | -22.4848 | 1.050 |
| 3 | 64-11-7-8.0 | 6.6117741 | . 081 | 39.858 | . 00006 | 153.835 | -153.831 | 312.4816 | 179.779 | -18.4892 | 0.776 |
| 4 | 64-11-16-3.0 | 6.6116883 | . 095 | 51.145 | . 00010 | -162.428 | 162.426 | 321.4898 | 179.226 | - 9.4810 | 0.223 |
| 5 | 64-11-24-3.0 | 6.6114084 | . 123 | 56.365 | . 00014 | -179.030 | 179.030 | 329.4837 | 178.790 | - 1.4871 | -0.213 |
| 6 | 64-11-30-10-35.0 | 6.6115566 | . 132 | 60.210 | . 00005 | -165.418 | 165.417 | 335.4787 | 178.495 | 4.5079 | -0.508 |
| 7 | 64-12-8-12-45.0 | 6.6112925 | . 161 | 64.207 | . 00011 | 144.315 | -144.310 | 344.4663 | 178.125 | 13.4955 | -0.878 |
| 8 | 64-12-15-12.0 | 6.6114751 | . 174 | 66.575 | . 00007 | -178.506 | 178.505 | 351.4544 | 177.880 | 20.4836 | -1.123 |
| 9 | 64-12-21-9.0 | 6.6111823 | . 198 | 69.060 | . 00014 | 158.894 | -158.891 | 356.4480 | 177.731 | 25.4772 | -1.272 |
|  |  | Average: $6.611572=a_{s}$ | Average: $0.122=\mathrm{i}_{\mathrm{s}}$ |  |  |  |  |  |  |  |  |

Results of least squares fit of data in (1) and (2) according to the theory of Equation 1, for arc 6S-A
$L=a_{1}+a_{2} t+a_{3} t^{2}+a_{4} t^{3}$
$a_{1}=-(0.29025 \pm 0.00217)$ degrees
$a_{2}=-(4.947 \pm 0.0234) \times 10^{-2}$ degrees $/$ day
$a_{3}=(4.438 \pm 0.0546) \times 10^{-4}$ degrees $/ \mathrm{day}^{2}$
$a_{4}=-(6.082 \pm 4.435) \times 10^{-7}$ degrees $/$ day $^{3}$
Standard error of estimate $=3.938 \times 10^{-3}$ degrees
$\ddot{\lambda}$ (measured) $=(1.5394 \pm 0.0190) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, at $\lambda=178.703^{\circ}$ and $\mathrm{t}^{\circ}=331.1727$ January 1964
$\ddot{\lambda}($ theoretical $)=+(1.5036) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$ for $\mathrm{a}_{2}=6.611572$ earth radii, $\mathrm{i}_{2}=0.122^{\circ}, \lambda=178.703^{\circ}, \mathrm{J}_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18.0^{\circ}$
-Computed by ITEM with gravity constants the same as in Table A1, with the addition of the earth constants: $\mathrm{J}_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18^{\circ}$.

Osculating Elements at the First Ascending Equator Crossing Past the Syncom 3 Tracking Epoch and Related Data, in Two Simulated Syncom 3 Arc 6 Trajectories with Earth Longitude Gravity through Second Order.*

| Orbit Number 6S-B- | Syncom 3 Tracking Epoch (yr-mo-day-hr UT) | Semimajor Axis, a (earth radii) | $\begin{gathered} \text { Inclination, } \\ \text { i } \\ \text { (degrees) } \end{gathered}$ | Right <br> Ascension of the Ascending Node (degrees) | Eccentricity | Argument of Perigee (degrees) | Mean Anomaly (degrees) | $\begin{gathered} \text { Time from } \\ 1964.0 \\ t \\ \text { (days) } \end{gathered}$ | Geographic Longitude of the First Ascending Equator Crossing After the Tracking Epoch, $\lambda$ (degrees) | $\begin{gathered} \text { © } \\ \text { Time From } \\ 330.9708 \\ \text { January 1964, } \\ \mathbf{t} \\ \text { (days) } \end{gathered}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 64-10-31-2.0 | 6.6119240 | . 097 | -149.455 | . 00026 | -45.719 | 45.698 | 305.9723 | 180.219 | -24.9985 | 1.216 |
| 2 | 64-11-3-13-18.0 | 6.6119574 | . 093 | -149.660 | . 00026 | -41.421 | 41.403 | $308.9642{ }^{\text { }}$ | 180.003 | -22.0066 | 1.000 |
| 3 | 64-11-7-8.0 | 6.6116738 | . 082 | -153.621 | . 00022 | -45.310 | 45.294 | 312.9430 | 179.735 | -18.0278 | 0.732 |
| 4 | 64-11-16-3.0 | 6.6117460 | . 077 | -168.885 | . 00028 | -18.764 | 18.754 | 321.8777 | 179.173 | - 9.0931 | 0.170 |
| 5 | 64-11-24-3.0 | 6.6115233 | . 070 | 167.455 | . 00022 | - 8.705 | 8.701 | 329.7916 | 178.729 | - 1.1792 | -0.274 |
| 6 | 64-11-30-10-35.0 | 6.6113283 | . 075 | 159.890 | . 00020 | 1.339 | - 1.338 | 335.7550 | 178.435 | 4.7842 | -0.568 |
| 7 | 64-12-8-12-45.0 | 6.6114224 | . 088 | 138.974 | . 00028 | 9.174 | - 9.169 | 344.6736 | 178.064 | 13.7028 | -0.939 |
| 8 | 64-12-15-12.0 | 6.6110825 | . 099 | 132.866 | . 00017 | 31.826 | -31.819 | 351.6382 | 177.805 | 20.6930 | -1.198 |
| 9 | 64-12-21-9.0 | 6.6114383 | . 118 | 124.209 | . 00025 | 27.235 | -27.225 | 356.6010 | 177.654 | 25.6302 | -1.349 |
|  |  | $\begin{gathered} \text { Average: } \\ 6.611566=\mathrm{a}_{\mathrm{S}} \end{gathered}$ | Average: $0.089=i_{\mathrm{s}}$ |  |  |  |  |  |  |  |  |

Results of least squares fit of data in (1) and (2) according to the theory of Equation 1, for arc 6 S -B
$L=a_{1}+a_{2}{ }^{t}+a_{3} \mathbf{t}^{2}+a_{4} \mathbf{t}^{3}$
$a_{1}=-(3.310 \pm 0.0237) \times 10^{-1}$ degrees
$a_{2}=-(5.082 \pm 0.0257) \times 10^{-4}$ degrees $/$ day
$a_{3}=(4.388 \pm 0.0604) \times 10^{-6}$ degrees $/$ day $^{2}$
$a_{4}=-(1.930 \pm 4.926) \times 10^{-7}$ degrees $/$ day $^{3}$
Standard error of estimate $=4.290 \times 10^{-3}$
$\ddot{\lambda}($ measured $)=(1.5231 \pm 0.0209) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$ at $\lambda=178.662^{\circ}$ and $\mathrm{t}^{\prime}=331.1727 \mathrm{January} 1964$
$\ddot{\lambda}$ (theoretical) $=1.5003 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$ for $\mathrm{a}_{2}=6.611566$ earch radii, $\mathrm{i}_{ \pm}=0.089^{\circ}, \lambda=178.662^{\mathrm{o}}, \mathrm{J}_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18.0^{\mathrm{o}}$
Bias $=$ theoretical-measured $=1.5003 \times 10^{-5}-1.5231 \times 10^{-5}=-(0.0228) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$
*Computed by ITEM with gravity constants the same as in Table A1, with the addition of the earth constants: $\mathrm{J}_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18^{\circ}$.

Table 6S/1
Ascending Equator Crossing Data From a Simulated Syncom 3 Trajectory for Free Drift Arc 6, Computed by ITEM With Earth Longitude Gravity through Third Order.*

| Orbit Number 6S/1 | Tracking Epoch (yr-mo-day-hr UT) | Semimajor Axis, <br> a <br> (earth radii) | $\begin{gathered} \text { Inclination, } \\ \mathbf{i} \\ \text { (degrees) } \end{gathered}$ | $\begin{gathered} \text { Time from } \\ 1964.0 \\ \text { (days) } \end{gathered}$ | Geographic <br> Longitude of the <br> Ascending <br> Equator <br> Crossing, <br> $\lambda$ <br> (degrees) | (1) Time from January $331.2865,1964$ $\mathbf{t}$ (days) | (2) of the Ascending Equator Crossing East of $179.002^{\circ}$, L (degrees) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 64-10-31-2.0 | 6.6119228 | . 097 | 305.9723 | 180.219 | -25.3142 | 1.217 |
| 2 | 64-11-3-13-18.0 | 6.6119523 | . 093 | 308.9642 | 180.004 | -22.3223 | 1.002 |
| 3 | 64-11-7-8.0 | 6.6116643 | . 082 | 312.9430 | 179.738 | -18.3435 | . 736 |
| 4 | 64-11-16-3.0 | 6.6117257 | . 078 | 321.8777 | 179.188 | -9.4088 | . 186 |
| 5 | 64-11-24-3.0 | 6.6114943 | . 070 | 329.7915 | 178.760 | - 1.4950 | - . 242 |
| 6 | 65-11-30-10-35.0 | 6.6112925 | . 075 | 335.7549 | 178.482 | 4.4684 | - . 520 |
| 7 | 64-12-8-12-45.0 | 6.6113766 | . 088 | 344.6733 | 178.141 | 13.3868 | - . 861 |
| 8 | 64-12-15-12.0 | 6.6110288 | . 099 | 351.6379 | 177.911 | 20.3514 | -1.091 |
| 9 | 64-12-21-9.0 | 6.6113790 | . 118 | 356.6007 | 177.784 | 25.3142 | -1.218 |
|  |  | Average: $6.611537=\mathrm{a}_{\mathbf{s}}$ | Average: $0.089=i_{s}$ |  |  |  |  |

Results of least squares fit of data in (1) and (2) above according to the theory of Equation 1:
$L=a_{1}+a_{2} t+a_{3} t^{2}+a_{4} t^{3}$
$a_{1}=-(3.1148 \pm 0.0233) \times 10^{-1}$ degrees
$a_{2}=-(4.8012 \pm 0.0253) \times 10^{-2}$ degrees $/$ day
$\mathbf{a}_{3}=(4.8565 \pm 0.0594) \times 10^{-4}$ degrees $/$ day $^{2}$
$a_{4}=-(1.362 \pm 4.856) \times 10^{-7}$ degrees $/$ day $^{3}$
Standard error of estimate $=0.00423$ degrees
$\ddot{\lambda}$ (with minimum standard error) $=1.6860 \times 10^{-5} \mathrm{rad} / \mathrm{sid}^{\text {. day }}{ }^{2}$, at $t=-0.0091$ day, $t^{\prime}=331.2774 \mathrm{Jan} .1964, \lambda=178.69^{\circ}$, measured $\ddot{\lambda}$ (theoretical from Equation 2, $\mathrm{i}_{\mathrm{s}}=0.089^{\circ}, \lambda=178.69^{\circ}, \mathrm{J}_{22}-\mathrm{J}_{31}$ as noted) $=1.6615 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, for $\mathrm{a}_{\mathrm{s}}=6.61154$ earth radii Estimate of acceleration bias at $\lambda=178.69^{\circ}$ in arc $6 S / 1=\ddot{\lambda}$ (theoretical) $-\ddot{\lambda}$ (measured) $=-0.0245 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$, day ${ }^{2}$
*Gravity constants of this trajectory are the same as that in Table A1, with the addition of the earth constants:

$$
\begin{aligned}
& \mathrm{J}_{22}=-1.8 \times 10^{-6}, \lambda_{22}=-15.35^{\circ} ; \mathrm{J}_{33}=-0.16 \times 10^{-6}, \lambda_{33}=24^{\circ} \\
& \mathrm{J}_{31}=-1.5 \times 10^{-6}, \lambda_{31}=0^{\circ} .
\end{aligned}
$$

The initial elements of this trajectory, aside from those listed for orbit $6 \mathbf{S} / 1-1$, are the same as those in orbit $65-\mathrm{B}-1$ (Table $6 S$ ).


Figure 5-Measured and simulated orbit data at ascending Equator crossings in free drift arc 6 (Syncom 3).
on Syncom 3 were inactive. During this two month period the mean longitude of the satellite moved westward from about $181^{\circ}$ to $173.5^{\circ}$ with a mean drift rate of about -0.1 degree/day. The details of this gravity accelerated drift are found in Table 7 and Figure 6 and summarized in Tables 10 and 11 in the next section.

The acceleration analysis during this relatively slow westward drift ${ }^{\prime}$ utilized the same $t^{3}$ fit method as in previous slow drift 24 -hour satellite arcs 1,2 , and 6 . The efficacy of this method is once again attested to by the small bias results evident in the close arc 7 simulated trajectories in Tables 7 S and $7 \mathrm{~S} / 1$. In these simulations, no attempt was made to break the arc at the end of January 1965 to account for the evident inclination change maneuver. Near-equatorial 24 -hour orbits, theoretically, suffer resonant gravity effects with small sensitivity to inclination changes (see Equations 3 through 7).

It is noted that the standard error of estimate for the $t^{3}$ longitude fit in this arc is over four times that in the arc 6 experiment for the same amount of data. The results of attempts to reduce this error by both a priori and a posteriori weighting have been inconclusive, in large part because of the scarcity of the data in this arc. In arc 2, which apparently suffers from similarly ill determined orbits, the standard error in the best measured acceleration is relatively small because more data is available in that arc than in arc 7. Nevertheless, the best measured acceleration in arc 7 seems to be much less in error from true earth resonant acceleration than its standard first stage experiment error (in Table 7) would indicate, (compare the results of the gravity synthesis in the next section). In the future, a more detailed analysis of the day-by-day drift of Syncom 3 in this arc, after the method used for Early Bird (arc 9), promises a far better discrimination of the acceleration, and with greater longitude separation from arc 6 than presently attained (see Discussion).

## Arc 8, Syncom 2, 25 February 1965-10 May 1965

Between 19 and 24 February 1965, gas jets were pulsed on board Syncom 2 to reorient the satellite and adjust its mean motion to as close to synchronous as possible. The last of this final series of Syncom 2 maneuvers took place on 24 February 1965 and left the satellites ground track with a westward drift rate of 0.05 degree/day (reduced from 0.6 degree/day westward on 19 February) at a mean longitude of $67.7^{\circ}$. Under the influence of earth longitude gravity, the ground track drift rate westward was further reduced till a momentarily stationary condition was reached in late May 1965 near $65.2^{\circ}$. The details of this gravity decelerated drift are found in Table 8 and Figure 7 and summarized in Tables 10 and 11 in the next section.

Once again, the slow drift " $t^{3}$ " theory was utilized for the acceleration analysis in this arc, both for the actual data and the closely parallel simulated data (see Tables 8 S and $8 \mathrm{~S} / 1$ ). The orbit determination over this slow drift arc appears to be about as good as in arc 1 (Syncom 2), but not as precise as in arc 6 (Syncom 3). The analysis of the two simulated trajectories with different earth models shows reasonably consistent model bias acceleration results for this arc considering the different lengths of these trajectories. The results of the parallel acceleration analysis on the simulated trajectories are summarized in Tables 10 and 11 in the next section.

According to the $\mathbf{2 4}$-hour satellite earth gravity drift theory summarized in this report (see Conclusions), a point of stable drift equilibrium exists for the Equatorial Geostationary Satellite, at about $77^{\circ}$ (and at about the same longitude for a $32^{\circ}$ inclined orbit satellite). Thus Syncom 2 is now probably forever in a long period oscillatory drift regime between mean longitudes of about $65^{\circ}$ and $89^{\circ}$. The exact description of this oscillation (with a period of about 2-1/2 years initially) depends on the long term change in the inclination of Syncom 2's orbit due to the gravitational attractions of the sun and moon. At present (summer 1965) the orbit inclination is about $31.7^{\circ}$ and is being reduced at about 0.8 degree/year (see Table 8S/1).

Syncom 3 Osculating Elements at the First Ascending Equator Crossing Past the Tracking Epoch,

| Orbit <br> Number 7- | Syncom 3 Tracking Epoch (yr-mo-day-hr-min UT) | $\begin{gathered} \text { Semimajor Axis, } \\ a \\ \text { (earth radii) } \end{gathered}$ | $\begin{gathered} \text { Inclination, } \\ \text { i } \\ \text { (degrees) } \end{gathered}$ | Right <br> Ascension of the Ascending Node (degrees) | Eccentricity | Argument of Perigee (degrees) | Mean Anomaly (degrees) | Time from 1965.0 (days) | Geographic <br> Longitude of the <br> Ascending Equator Crossing, $\lambda$ (degrees) | $\underset{\text { Time from }}{(1)}$ January 45.4138, 1965, t (days) | $\begin{gathered} \text { (2) } \\ \text { Longitude } \\ \text { of the } \\ \text { Ascending } \\ \text { Equator } \\ \text { Crossing } \\ \text { East of } \\ 177.164^{\circ} \mathrm{E}, \\ \mathrm{~L} \\ \text { (degrees) } \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1** | 65-1-14-23-30.0 | 6.6121095 | 1.137 | - 34.097 | . 00181 | -97.425 | 97.220 | 15.0855 | 180.803 | -30.3283 | 3.639 |
| 2 | 65-1-30-13-10.0 | 6.6124038 | . 173 | - 61.562 | . 00029 | -64.872 | 64.844 | 30.9720 | 178.543 | -14.4418 | 1.379 |
| 3 | 65-2-2-6.0 | 6.6129132 | . 101 | - 35.767 | . 00013 | -64.858 | 64.847 | 34.0364 | 178.124 | -11.3774 | . 960 |
| 4 | 65-2-9-11.0 | 6.6124858 | . 166 | - 81.989 | . 00025 | -76.958 | 76.934 | 40.8915 | 177.300 | - 4.5223 | . 136 |
| 5 | 65-2-16-12.0 | 6.6118440 | . 060 | -112.788 | . 00001 | -30.762 | 30.809 | 47.7895 | 176.446 | 2.3757 | - . 718 |
| 6 | 65-2-23.0 | 6.6123345 | . 138 | -100.918 | . 00019 | -82.895 | 82.879 | 54.8054 | 175.677 | 9.3916 | -1.487 |
| 7 | 65-3-2.0 | 6.6119364 | . 205 | -118.912 | . 00034 | -87.998 | 87.963 | 61.7382 | 175.021 | 16.3244 | -2.143 |
| 8 | 65-3-9.0 | 6.6123015 | . 074 | - 71.138 | . 00020 | -76.296 | 76.279 | 68.8539 | 174.130 | 23.4401 | -3.034 |
| 9 | 65-3-16.0 | 6.6120906 Average: $6.612269=\mathrm{a}_{\mathrm{S}}$ | $\begin{gathered} .355 \\ \hline \text { Average } \\ 0.268=i_{s} \end{gathered}$ | -105.212 | . 00070 | -96.150 | 96.072 | 75.7421 | 173.525 | 30.3283 | -3.639 |

[^3]Ascending Equator Crossing Data from a Simulated Syncom 3 Trajectory for Free Drift Arc 7,
Computed by ITEM in the Presence of Earth Longitude Gravity*.

| Orbit Number 7S | Syncom 3 Tracking Epoch (yr-mo-day-hr-min UT) | $\begin{aligned} & \text { Semimajor } \\ & \text { Axis, } \\ & \text { a } \\ & \text { (earth radii) } \end{aligned}$ | $\begin{gathered} \text { Inclination, } \\ \text { i } \\ \text { (degrees) } \end{gathered}$ | Right Ascension of the Ascending Node (degrees) | $\begin{gathered} \text { Eccentri- } \\ \text { city } \end{gathered}$ | Argument of Perigee (degrees) | Mean <br> Anomaly <br> (degrees) | $\begin{gathered} \text { Time from } \\ 1965.0 \\ \text { (days) } \end{gathered}$ | Geographic <br> Longitude of the <br> Ascending Equator Crossing, $\lambda$ (degrees) | (1) <br> Time From January 45.4138, 1965, $t$ (days) | (2) <br> Longitude of the Ascending Equator Crossing East of $177.164^{\circ} \mathrm{E}$, L (degrees) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 65-1-14-23-30.0 | 6.6127641 | . 200 | -64.894 | . 00030 | - 98.417 | 98.416 | 15.0024 | 180.836 | -30.4094 | 3.672 |
| 2 | $65-1-30-13-10.0$ | 6.6126603 | . 148 | -63.708 | . 00025 | -111.274 | 111.249 | 30.9658 | 178.628 | -14.4480 | 1.464 |
| 3 | 65-2-2-6.0 | 6.6125307 | . 140 | -65.126 | . 00024 | -117.018 | 116.994 | 33.9548 | 178.244 | -11.4590 | 1.080 |
| 4 | 65-2-9-11.0 | 6.6123492 | . 126 | -66.027 | . 00030 | -105.735 | 105.705 | 40.9356 | 177.363 | - 4.4782 | . 199 |
| 5 | 65-2-16-12.0 | 6.6122952 | . 100 | -67.046 | . 00021 | -111.429 | 111.408 | 47.9160 | 176.519 | 2.5022 | - . 645 |
| 6 | 65-2-23.0 | 6.6122513 | . 091 | -66.187 | . 00028 | -116.269 | 116.241 | 54.9015 | 175.710 | 9.4877 | -1.454 |
| 7 | 65-3-2.0 | 6.6121399 | . 071 | -68.756 | . 00025 | -119.829 | 119.804 | 61.8774 | 174.951 | 16.4636 | -2.213 |
| 8 | 65-3-9.0 | 6.6120920 | . 065 | -67.449 | . 00029 | -104.308 | 104.279 | 68.8639 | 174.210 | 23.4501 | -2.954 |
| 9 | 65-3-16.0 | 6.6119327 | . 044 | -70.595 | . 00023 | -115.776 | 115.755 | 75.8381 | 173.502 | 30.4243 | -3.662 |
|  |  | Average: $6.612335=a_{\mathrm{s}}$ | Average: $0.109=\mathrm{i}_{\mathrm{s}}$ |  |  |  |  |  |  |  |  |

${ }^{*} \mathrm{~J}_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18^{\circ}$, only earth longitude gravity used in this simulation. All orher gravity constants are the same as those used in Table Al.
Results of least squares fit of data in (1) and (2) above (Table 7S) according to the theory of Equation 1:

$$
\begin{aligned}
& L=a_{1}+a_{2} t+a_{3} t^{2}+a_{4} t^{3} \quad a_{1}=-(3.4626 \pm 0.0182) \times 10^{-1} \text { degrees } \quad a_{2}=-(1.1971 \pm 0.0017) \times 10^{-1} \text { degrees } / \text { day } \\
& \mathrm{a}_{3}=(3.8156 \pm 0.0374) \times 10^{-4} \text { degrees } / \text { day }{ }^{2} \quad \mathrm{a}_{4}=-(8.980 \pm 2.241) \times 10^{-7} \text { degrees } / \text { day }{ }^{3}
\end{aligned}
$$

Standard error of estimate $=0.00383$ degree
$\ddot{\lambda}$ (measured, at $176.871^{\circ}$ ) $=(1.3289 \pm 0.0129) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, at $t=-0.4507$ day, $\mathrm{t}^{1}=44.9631$ January 1965 .
$\ddot{\lambda}($ theoretical $)=1.3546 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, for $\mathrm{a}_{\mathrm{s}}=6.612335$ earth radii, $\mathrm{i}_{\mathrm{s}}=0.109^{\circ}, \lambda=176.871^{\circ}, \mathrm{J}_{22}=-1.68 \times 10^{-6}, \lambda_{22}-18.0^{0}$.
Bias $=$ Theoretical - Measured $=1.3546 \times 10^{-5}-1.3289 \times 10^{-5}=+0.0257 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day $^{2}$

Table 7S/1
Ascending Equator Crossing Data from a Simulated Syncom 3 Trajectory for Free Drift Arc 7, Computed by ITEM in the Presence of Earth Longitude Gravity*.

| Orbit Number 7S/1- | Syncom 3 <br> Tracking Epoch (yr-mo-day-hr-min UT) | $\begin{gathered} \text { Semimajor } \\ \text { Axis, } \\ \text { a } \\ \text { (earth radii) } \end{gathered}$ | $\begin{aligned} & \text { Inclination, } \\ & \mathbf{i} \\ & \text { (degrees) } \end{aligned}$ | $\begin{gathered} \text { Time from } \\ 1965.0 \\ \text { (days) } \end{gathered}$ | Geographic <br> Longitude of the <br> Ascending <br> Equator <br> Crossing, $\lambda$ <br> (degrees) | (1) Time From January 45.4138, 1965, t (days) | (2) <br> Longitude of the Ascending Equator Crossing East of $177.164^{\circ} \mathrm{E}$, L (degrees) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 65-1-14-23-30.0 | 6.6127641 | . 200 | 15.0002 | 180.836 | -30.4136 | 3.672 |
| 2 | 65-1-30-13-10.0 | 6.6126416 | . 148 | 30.9658 | 178.641 | -14.4480 | 1.477 |
| 3 | 65-2-2-6.0 | 6.6125089 | . 137 | 33.9547 | 178.262 | -11.4591 | 1.098 |
| 4 | 65-2-9-11.0 | 6.6123196 | . 126 | 40.9355 | 177.396 | - 4.4783 | . 232 |
| 5 | 65-2-16-12.0 | 6.6122585 | . 100 | 47.9158 | 176.571 | 2.5020 | - . 593 |
| 6 | 65-2-23.0 | 6.6122077 | . 091 | 54.9013 | 175.786 | 9.4875 | -1.378 |
| 7 | 65-3-2.0 | 6.6120897 | . 071 | 61.8771 | 175.054 | 16.4633 | -2.110 |
| 8 | 65-3-9.0 | 6.6120352 | . 065 | 68.8635 | 174.344 | 23.4497 | -2.820 |
| 9 | 65-3-16.0 | 6.6118698 | . 044 | 75.8376 | 173.671 | 30.4238 | -3.493 |
|  |  | Average: $6.6123=a_{s}$ | Average: $0.109=i_{s}$ |  |  |  |  |

${ }^{*} \mathrm{~J}_{22}=-1.8 \times 10^{-6}, \lambda_{22}=-15.35^{\circ}, \mathrm{J}_{33}=-.16 \times 10^{-6}, \lambda_{33}=+24^{\circ}, \mathrm{J}_{31}=-1.5 \times 10^{-6}, \lambda_{31}=0^{\circ}$, only earth longitude gravity used in this simulation. Initial elements as in orbit $7 \mathrm{~S}-1$ above. All other gravity constants as in Table Al.
Results of least squares fit of data in (1) and (2) above (Table 7S/1) according to the theory of Equation 1 :
$L=a_{1}+a_{2} t+a_{3} t^{2}+a_{4} t^{3} \quad a_{1}=-(3.0135 \pm .0178) \times 10^{-1}$ degrees $\quad a_{2}=-(1.1687 \pm 0.00166) \times 10^{-1}$ degrees $/$ day $a_{3}=(4.2397 \pm 0.0366) \times 10^{-4}$ degrees $/$ day ${ }^{2} \quad a_{4}=(9.599 \pm 2.192) \times 10^{-7}$ degrees $/$ day ${ }^{3}$
Standard error of estimate $=0.00375$ degree
$\ddot{\lambda}$ (measured, with minimum standard error) $=1.4726 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, at $\boldsymbol{t}=-0.4433$ days, $\lambda=176.915^{\circ}$.
$\ddot{\lambda}\left(\right.$ theoretical, from Equation 2 , for $a_{s}=6.6123$ E.R., $i_{s}=0.109^{\circ}, \lambda=176.915^{\circ}$ ) $=1.501 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$
Estimated bias in arc $\mathrm{S} 7 / 1$ at $\lambda=176.915^{\circ}=$ theoretical - measured $=1.501 \times 10^{-5}-1.473 \times 10^{-5}=$
$0.028 \times 10^{-5} \mathrm{rad} /$ sid. day ${ }^{2}$


Figure 6-Measured and simulated orbit data at ascending Equator crossings in drift arc 7 (Syncom 3).

## Arc 9, Early Bird, 23 April 1965 - 21 June 1965

The Communications Satellite Corporation's first 24-hour satellite, Early Bird, was launched in March 1965 and was brought to station at $30^{\circ} \mathrm{W}$ in late April 1965. Free gravity drift of this nearly geostationary satellite commenced on 23 April with the mean longitude at $30^{\circ} \mathrm{W}$ and the drift rate about +0.06 degree/day. Tracking has been maintained on nearly an around-the-clock basis since this time from the A.T. and T. facility at Andover, Maine. Figure 8 shows Early Bird reached a momentarily stationary configuration at $28^{\circ} \mathrm{W}$ in late June 1965.

Extremely fine precision in defining the long term drift of this satellite has been achieved by averaging the results of a large number of daily subsatellite position determinations converted directly from simultaneous range, azimuth and elevation fixes on Early Bird from Andover. The technique of utilizing this large amount of directly observed data almost every day in arc 9 can be followed in Table 9.

Syncom 2 Mean Elements at the Tracking Epoch, and Related Data for Free Drift Arc 8*.

| Orbit Number 8- | ```Tracking Epoch (yr-mo-day-hr- min, UT)``` | Semimajor Axis, a (earth radii) | $\begin{gathered} \text { Inclination, } \\ \text { i } \\ \text { (degrees) } \end{gathered}$ | Right <br> Ascension of the Ascending Node (degrees) | Eccentricity | Argument of Perigee (degrees) | Mean <br> Anomaly <br> (degrees) | Time from January $0.0,1965.0$ (days) | Geographic <br> Longitude <br> of the <br> Ascending <br> Equator <br> Crossing <br> After <br> Tracking <br> Epoch, <br> $\lambda$ <br> (degrees) | (1) Time from January $86.1612,1965$, $t$ (days) | $\begin{gathered} \text { (2) } \\ \text { Longitude } \\ \text { of } \\ \text { Ascending } \\ \text { Equator } \\ \text { Crossing } \\ \text { East of } \\ 66.634^{\circ} \\ \text { L } \\ \text { (degrees) } \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| **1 | 65-2-25.0 | 6.6115151 | 31.956 | -50.539 | . 00075 | -28.510 | 28.469 | 56.2412 | 67.733 | -29.9200 | 1.099 |
| **2 | 65-3-3.0 | 6.6112551 | 31.930 | -50.539 | . 00079 | -25.098 | 25.060 | 62.2257 | 67.423 | -23.9355 | . 789 |
| 3 | 65-3-6.0 | 6.6112849 | 31.939 | -50.630 | . 00076 | -26.952 | 308.359 | 65.2177 | 67.270 | -20.9435 | . 636 |
| 4 | 65-3-13.0 | 6.6113429 | 31.912 | -50.905 | . 00069 | -38.553 | 326.845 | 72.1986 | 67.007 | -13.9626 | . 373 |
| **5 | 65-3-29.0 | 6.6114858 | 31.955 | -50.921 | . 00077 | -18.507 | 18.479 | 88.1567 | 66.314 | 1.9955 | -. 320 |
| 6 | 65-4-5.0 | 6.6111445 | 31.830 | -51.104 | . 00069 | -19.366 | 329.612 | 95.1378 | 66.071 | 8.9766 | - . 563 |
| 7 | 65-4-12.0 | 6.6111334 | 31.848 | -51.053 | . 00073 | -11.053 | 328.269 | 102.1195 | 65.822 | 15.9583 | -. 812 |
| 8 | 65-4-19.0 | 6.6110585 | 31.867 | -51.414 | . 00065 | -29.162 | 353.146 | 109.0997 | 65.712 | 22.9385 | -. 922 |
| 9 | 65-4-26.0 | 6.6110267 | 31.787 | -51.420 | . 00063 | -22.995 | 353.717 | 116.0811 | 65.534 | 29.9199 | -1.100 |
|  |  | $\begin{gathered} \text { Average: } \\ 6.611250=\mathrm{a}_{\mathrm{s}} \end{gathered}$ | $\begin{array}{r} \text { Average: } \\ 31.892=\mathrm{i}_{\mathrm{s}} \end{array}$ | Data thr | ugh orbit 8 |  |  |  |  |  |  |
| 10 | 65-5-3.0 | 6.6110035 | 31.807 | -51.583 | . 00061 | -22.243 | 359.897 | 123.0619 | 65.404 | 36.9007 | -1.230 |
| 11 | 65-5-10.0 | 6.6109407 | 0.731 | -51.665 | . 00059 | -15.485 | 0.012 | 130.0428 | 65.288 | 43.8816 | -1.346 |
|  |  | Average: $6.611199=\mathrm{a}_{\mathrm{s}}$ | Average: $31.869=i_{s}$ | Data thr | h orbit | $1 .$ |  |  |  |  |  |

*Mean (Brouwer) elements and longitude data as reported by GSFC Tracking and Data Systems Directorate (except as noted). Longitude data also reported by GSFC is at the first ascending Equator crossing past the Syncom 2 tracking epoch.
**Osculating elements and longitude data at the first ascending Equator crossing past the tracking epoch (see notes in Table 1).
Results of least squares fit of data in (1) and (2) above acording to the theory of Equation 1: (through orbit 8-9).

$$
\begin{aligned}
& \mathrm{L}=\mathrm{a}_{1}+\mathrm{a}_{2} \mathrm{t}+\mathrm{a}_{3} \mathrm{t}^{2}+\mathrm{a}_{4} \mathrm{t}^{3} \quad \mathrm{a}_{1}=-(2.430 \pm 0.168) \times 10^{-1} \text { degrees } \quad \mathrm{a}_{2}=-(3.857 \pm 0.131) \times 10^{-2} \text { degrees } / \text { day } \\
& \mathrm{a}_{3}=(2.699 \pm 0.319) \times 10^{4} \text { degrees } / \text { day }^{2} \quad \mathrm{a}_{4}=(2.445 \pm 1.873) \times 10^{-6} \text { degrees } / \text { day }{ }^{3}
\end{aligned}
$$

Standard error of estimate $=0.0291$ degree
$\ddot{\lambda}($ with minimum standard error $)=(0.9566 \pm 0.1097) \times 10^{-5} \mathrm{rad} . / \mathrm{sid}^{2}$ day ${ }^{2}$, at $\mathrm{t}=0.7759$ days, $\mathrm{t}=86.9371$ January $1965, \mathrm{~L}=-0.273^{\circ}, \lambda=66.361^{\circ}$.
Revised data: (through orbit 8-11)

$$
\begin{aligned}
& \mathrm{L}=\mathrm{a}_{1}+\mathrm{a}_{2} \mathrm{t}+\mathrm{a}_{3} \mathrm{t}^{2}+\mathrm{a}_{4} \mathrm{t}^{3} \quad \mathrm{a}_{1}=-(2.403 \pm .159) \times 10^{-1} \text { degrees } \quad \mathrm{a}_{2}=-(3.741 \pm 0.0864) \times 10^{-2} \text { degrees } / \text { day } \\
& \mathrm{a}_{3}=(2.6082 \pm 0.2845) \times 10^{-4} \text { degrees } / \text { day }^{2} \quad \mathrm{a}_{4}=(5.471 \pm 9.487) \times 10^{-7} \text { degrees } / \text { day }{ }^{3}
\end{aligned}
$$

Standard error of estimate $=0.0280$ degree
$\ddot{\lambda}_{\text {(with minimum standard error) }}=(0.9500 \pm .0616) \times 10^{-5} \mathrm{rad} . / \mathrm{sid}^{2}$ day ${ }^{2}$, at $\mathrm{t}=7.8135$ days, $\mathrm{t}^{1}=93.9747$ January $1965, \lambda=66.115^{\circ}$.
(see figure 7)

Table 8S
Ascending Equator Crossing Data From a Simulated Syncom 2 Trajectory for Free Drift Arc 8, Computed by ITEM in the Presence of Earth Longitude Gravity*.

| Orbit <br> Number 8S - | Tracking Epoch (yr-mo-day UT) | Semimajor Axis, a (earth radii) | Inclination, i (degrees) | Right <br> Ascension of the Ascending Node (degrees) | Eccentricity | Argument of Perigee (degrees) | Mean <br> Anomaly <br> (degrees) | $\begin{gathered} \text { Time from } \\ 1965.0 \\ \text { (days) } \end{gathered}$ | Geographic <br> Longitude of the <br> Ascending Equator <br> Crossing, $\lambda$ <br> (degrees) | (1) Time from January $86.1612,1965$, $t$ (days) | (2) <br> Longitude of the Ascending Equator Crossing East of $66.634^{\circ}$, L (degrees) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 65-2-25.0 | 6.6114062 | 31.956 | -50.538 | . 00075 | -28.528 | 28.488 | 56.2412 | 67.735 | -29.9200 | 1.101 |
| 2 | 65-3-3.0 | 6.6114759 | 31.936 | -50.611 | . 00075 | -25.625 | 25.588 | 62.2254 | 67.456 | -23.9358 | . 822 |
| 3 | 65-3-6.0 | 6.6112778 | 31.925 | -50.663 | . 00070 | -27.225 | 27.188 | 65.2175 | 67.321 | -20.9437 | . 687 |
| 4 | 65-3-13.0 | 6.6114573 | 31.913 | -50.770 | . 00079 | -24.836 | 24.798 | 72.1990 | 66.993 | -13.9622 | . 359 |
| 5 | 65-3-29.0 | 6.6113880 | 31.871 | -51.015 | . 00075 | -25.549 | 25.512 | 88.1565 | 66.318 | 1.9953 | - . 316 |
| 6 | 65-4-5.0 | 6.6110762 | 31.853 | -51.142 | . 00070 | -29.586 | 29.547 | 95.1377 | 66.048 | 8.9765 | - . 586 |
| 7 | 65-4-12.0 | 6.6112985 | 31.839 | -51.239 | . 00076 | -20.941 | 20.910 | 102.1191 | 65.770 | 15.9579 | -. 864 |
| 8 | 65-4-19.0 | 6.6110384 | 31.826 | -51.369 | . 00071 | -26.757 | 26.721 | 109.1003 | 65.525 | 22.9391 | -1.109 |
| 9 | 65-4-26.0 | 6.6112287 | 31.816. | -51.462 | . 00074 | -25.117 | 25.081 | 116.0816 | 65.291 | 29.9204 | -1.343 |
|  |  | Average: $6.611289=a_{\mathrm{S}}$ | Average: $31.882=i_{s}$ |  |  |  |  |  |  |  |  |

${ }^{*} \mathrm{~J}_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18^{\circ}$ only earth longitude gravity used in this simulation. All other gravity constants as in Table A1.
Results of least squares fit of data in (1) and (2) above according to the theory of Equation 1:

$$
\begin{aligned}
\mathbf{L} & =\mathbf{a}_{1}+\mathbf{a}_{2} t+a_{3} \mathbf{t}^{2}+a_{4} \mathbf{t}^{3} \\
\mathbf{a}_{1} & =-(2.337 \pm 0.0336) \times 10^{-1} \text { degree } \\
\mathbf{a}_{2} & =-(4.126 \pm 0.0262) \times 10^{-2} \text { degrees } / \text { day } \\
\mathbf{a}_{3} & =(1.264 \pm 0.0638) \times 10^{-4} \text { degrees } / \text { day }^{2} \\
\mathbf{a}_{4} & =(4.151 \pm 3.743) \times 10^{-7} \text { degrees } / \text { day }^{3}
\end{aligned}
$$

Standard error of estimate $=0.00581$ degree
$\grave{\lambda}$ (measured) $=(0.4421 \pm 0.0219) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, for $t=0.7759 \mathrm{day}, \mathrm{t}^{\prime}=86.9371 \mathrm{Jan} .1965, \mathrm{~L}=-0.266^{\circ}, \lambda=66.368^{\circ}$
$\lambda$ (theoretical, from Equation 2, $\left.i_{1}=31.882^{\circ}, \lambda=66.368^{\circ}, J_{22}=-1.68 \times 10^{-6}, \lambda_{22}=-18.0^{\circ}\right)=(0.4560) \times 10^{-5} \mathrm{rad} / \mathrm{sid}^{2}$. day ${ }^{2}$, for a $=6.611289$ earth radii
Estimate of measured bias due to sun-moon perturbations and $\mathrm{J}_{22}$ model error (exclusive of higher order longitude gravity effects) in $\lambda$ at $t=0.7759$ days in Syncom 2 arc 8 , on January 86.9371, 1965: $\quad$ Bias = the oretical-measured

$$
\begin{aligned}
& =0.4560 \times 10^{-5}-0.4421 \times 10^{-5} \\
& =+0.0139 \times 10^{-5} \mathrm{rad} / \mathrm{sid} \mathrm{day}^{2}
\end{aligned}
$$

Table 8S/1
Ascending Equator Crossing Data From a Simulated Syncom 2 Trajectory for Free Drift Arc 8, Computed by ITEM in the Presence of Earth Longitude Gravity*.

| Orbit <br> Number <br> 8S/1 - | $\begin{gathered} \text { Tracking } \\ \text { Epoch } \\ \text { (yr-mo-day UT) } \end{gathered}$ | Semimajor Axis, <br> a <br> (earth radii) | $\begin{gathered} \text { Inclination, } \\ \text { i } \\ \text { (degrees) } \end{gathered}$ | Time from 1965.0 (days) | Geographic <br> Longitude of the <br> Ascending Equator Crossing, $\lambda$ (degrees) | (1) Time from January 86.1612, 1965, t (days) | (2) <br> Longitude of the Ascending Equator Crossing East of $66.634^{\circ}$, L (degrees) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 65-2-25.0 | 6.6114053 | . 956 | 56.2412 | 67.735 | -29.9200 | 1.101 |
| 2 | 65-3-3.0 | 6.6114581 | . 936 | 62.2254 | 67.460 | -23.9358 | . 826 |
| 3 | 65-3-6.0 | 6.6112516 | . 925 | 65.2174 | 67.331 | -20.9438 | . 697 |
| 4 | 65-3-13.0 | 6.6114112 | . 913 | 72.1989 | 67.023 | -13.9623 | . 389 |
| 5 | 65-3-29.0 | 6.6112959 | . 871 | 88.1561 | 66.438 | 1.9949 | - . 196 |
| 6 | 65-4-5.0 | 6.6109642 | . 853 | 95.1372 | 66.226 | 8.9760 | - . 408 |
| 7 | 65-4-12.0 | 6.6111669 | . 839 | 102.1184 | 66.018 | 15.9572 | - . 616 |
| 8 | 65-4-19.0 | 6.6108871 | . 827 | 109.0994 | 65.853 | 22.9382 | - . 781 |
| 9 | 65-4-26.0 | 6.6110582 | . 816 | 116.0804 | 65.711 | 29.9192 | - . 923 |
| 10 | 65-5-3.0 | 6.6107742 | . 805 | 123.0613 | 65.597 | 36.9001 | -1.037 |
| 11 | 65-5-10.0 | 6.6109303 | . 792 | 130.0422 | 65.492 | 43.8810 | -1.142 |
|  |  | Average: $6.611146=a_{s}$ | Average: $31.867=i_{s}$ |  |  |  |  |
| $\mathrm{J}_{22}=-1.8 \times 10^{-6} \quad \mathrm{~J}_{33}=-0.16 \times 10^{-6}$ |  |  | $\mathrm{J}_{31}=-1.5 \times 10^{-6}$ |  |  |  |  |
| $\lambda_{22}=-15.35^{\circ} \quad \lambda_{33}$ |  |  | $\lambda_{31}=0^{\circ}$ |  |  |  |  |

Only earth longitude gravity used in this simulation. Initial elements as in orbit $8 \mathbf{S}-1$ above. All other gravity constants as in Table A1.
Results of least squares fit of data in (1) and (2) above according to the theory of Equation 1 :

$$
\begin{aligned}
& L=a_{1}+a_{2} t+a_{3} t^{2}+a_{4} t^{3} \\
& a_{1}=-(1.2826 \pm 0.0336) \times 10^{-1} \text { degrees } \\
& a_{2}=-(3.4062 \pm 0.0183) \times 10^{-2} \text { degrees } / \text { day } \\
& a_{3}=(2.4390 \pm 0.0601) \times 10^{-4} \text { degrees } / \text { day }{ }^{2} \\
& a_{4}=(1.826 \pm 2.005) \times 10^{-7} \text { degrees } / \text { day }^{3}
\end{aligned}
$$

Standard error of estimate $=0.00591$ degree

$\ddot{\lambda}$ (theoretical, from Equation $\left.2, i_{s}=31.87^{\circ}, \lambda=66.255^{\circ}\right)=0.884 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$, for $a_{s}=6.61115$ earth radii.
Estimated bias in arc $8 S / 1$ at $\lambda=66.255^{\circ}$

$$
\begin{aligned}
& =\text { theoretical-measured } \\
& =0.884 \times 10^{-5}-0.862 \times 10^{-5}=+0.022 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \text { day }^{2}
\end{aligned}
$$



Figure 7-Measured and simulated orbit data in free drift arc 8 (Syncom 2).

In the first three columns of that table, the densely observed reference orbit ground track of Early Bird for 23 April is listed. The theory of 24 -hour satellite drift used in this report references such drift specifically to the ascending Equator crossing of the satellite (see References 2 and 3, for example). However, the results apply equally well to orbit averaged drift from any nodal argument. If the satellite suffered no orbit perturbations, an orbit period would be specified precisely by the time between two similarly directed latitude passes. Thus, we can measure the 24-hour orbit longitude drift by comparing a given measured longitude-latitude point in an actual orbit with its longitude in the reference orbit by way of the indicated latitude. Unfortunately, perturbations caused the inclination of Early Bird to grow by about $100 \%$ in the two months of arc 9. For a large number of measured ground points in this history of Early Bird, there was no clear reference latitude match to indicate orbit longitude drift. This problem was especially aggravated by the preponderance of Early Bird observations at the maximum north and south points. Fortunately, the observation time alone provides a second and always clear reference orbit match criteria, even in the presence of perturbations, providing the orbit period (nodal) is reasonably well known. In the case of Early Bird in the spring of 1965, the nodal period is obviously close to synchronous, or one revolution in about 4 minutes short of 24 hours. Thus, the reference orbit longitude which corresponds to $N$ number of orbit periods from the observed longitude occurs at a reference orbit time approximately 4 N minutes later in the day than the observed longitude for a nearly synchronous satellite. Many of the comparative longitude matchings in Table 9 were made solely on this daily orbit time basis, especially where the observations were in the maximum north and south regions. Near equatorial observations were generally matched by both methods and an average reference longitude is listed in Table 9 for these cases. It is noted that the observed reference orbit ground track is "biased" north by about $0.06^{\circ}$. Most of this bias is probably due to error in the Andover antenna altitude calibration. A similar bias in the azimuth calibration would shift the mean longitude of the satellite by a small constant amount, insignificant in this acceleration analysis.

The average of the days estimated drift from the reference orbit is then applied to the ascending Equator crossing longitude in the reference orbit, to arrive at an estimated ascending Equator crossing longitude for that day (see Table 9). The subsequent acceleration analysis of this derived crossing data follows the same technique as used in the other slow drift arcs $(1,2$, 6, 7, and 8). The average semimajor axis for Early Bird during arc 9 was not observed but inferred from the closely fitting simulated trajectory of orbit $9 \mathrm{~S} / 1$ (Table $9 \mathrm{~S} / 1$ ). The north-south excursion in the observed reference orbit (Table 9) gave a nominal value of $0.14^{\circ}$ for the initial orbit inclination of Early Bird. The average orbit inclination in this arc was also inferred from the close simulated trajectory of orbit $9 \mathrm{~S} / 1$.

The results of the acceleration analysis on the actual data and on the simulated data are found in Tables $9,9 \mathrm{~S}$ and $9 \mathrm{~S} / 1$, and also summarized in Tables 10 and 11 in the next section.

Table 9
Ground Track and Related Ascending Equator Crossing Data for Free Drift Arc 9 (Early Bird)*


Results of least squares fit of data in (1) and (2) (pages 47-48) according to the theory of Equation 1:
$L=a_{1}+a_{2} t+a_{3} t^{2}+a_{4} t^{3}$
$a_{1}=(3.6180 \pm 0.01023) \times 10^{-1}$ degrees

$$
\begin{aligned}
& \mathrm{a}_{3}=-(4.127 \pm 0.0279) \times 10^{-4} \text { degrees } / \text { sid. day }{ }^{2} \\
& \mathrm{a}_{4}=(1.039 \pm 1.799) \times 10^{-7} \text { degrees } / \text { sid. day }{ }^{3}
\end{aligned}
$$

$a_{2}=(3.202 \pm 0.0107) \times 10^{-2}$ degrees $/$ sid. day
Standard error of estimate $=4.95 \times 10^{-3}$ degrees
$\ddot{\lambda}$ (with minimum standard error) $=-(1.4408 \pm 0.0097) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$ day ${ }^{2}$, at $t=-0.3003 \mathrm{sid}$. day, $t^{\prime}=22$ May $1965, \lambda=-28.703^{\circ}$.

- Ground track (subsatellite point) information on Early Bird, from Andover Maine Tracking, supplied through the Offices of Robert Greene, Comsat Corp., Washington, D.C.
**Calculated as $30.005^{\circ}$ minus average drift from the reference orbit.

Ground Track and Related Ascending Equator Crossing Data for Free Drift Arc 9 (Early Bird)*

| Ground Track Reference Orbit Data |  |  | $\begin{gathered} \text { (May 1965) } \\ \text { Time } \\ \text { Day Hr } \end{gathered}$ | Latitude (degrees) | Longitude <br> (degrees west) | Comparative Longitude in Reference Orbit (degrees west) | $\begin{gathered} \text { Drift } \\ \text { from } \\ \text { Reference } \\ \text { Orbit } \\ \text { (degrees } \\ \text { east) } \end{gathered}$ | ** <br> Estimated <br> Ascending Equator Crossing <br> Longitude <br> (degrees west) | Time of Ascending Equator Crossing During Day: Estimate (hours) | Time (sid. day integers from 23 April 1965) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Time <br> (April 1965) <br> Day Hr <br> 23 | Latitude (degrees) | Geographic Longitude (degrees west) |  |  |  |  |  |  |  |  |
|  19 <br>  19 <br>  19 <br>  20 <br>  20 <br>  20 <br>  21 <br>  21 <br>  22 <br>  23 <br> $23 / 24$ $24 / 0$ <br> 24 1 <br>  2 <br>  2 <br>  2 | . 10 | 29.97 | 46 | . 01 | 29.495 | 30.08 | . 585 |  |  |  |
|  | . 11 | 29.965 | 12 | -. 08 | 29.465 | 30.055 | . 590 |  |  |  |
|  | . 12 | 29.96 | 21 | . 185 | 29.385 | 29.96 | . 575 |  |  |  |
|  | . 13 | 29.955 | 24/0 | . 20 | 29.40 | 29.97 | . 570 |  |  |  |
|  | . 14 | 29.955 |  |  |  | avg: | . 580 | 29.425 | 15 | 11 |
|  | . 15 | 29.95 | 56 | . 00 | 29.445 | 30.08 | . 635 |  |  |  |
|  | . 16 | 29.955 | 12 | -. 09 | 29.42 | 30.05 | . 630 |  |  |  |
|  | . 17 | 29.955 | 24/6 | . 19 | 29.35 | 29.97 | . 620 |  |  |  |
|  | . 18 | 29.96 |  |  |  | avg: | . 630 | 29.375 | 15 | 12 |
|  | . 19 | 29.965 | 66 | . 005 | 29.40 | 30.08 | . 680 |  |  |  |
|  | . 18 | 29.97 |  |  |  | avg: | . 680 | 29.325 | 15 | 13 |
|  | . 17 | 29.975 | $7 \quad 9$ | -. 085 | 29.35 | 30.07 | . 720 |  |  |  |
|  | . 16 | 29.98 |  |  |  | avg: | . 270 | 29.285 | 15 | 14 |
|  | . 15 | 29.98 | $8 \quad 22$ | . 22 | 29.205 | 29.965 | . 760 |  |  |  |
|  |  |  |  |  |  | avg: | . 760 | 29.245 | 15 | 15 |
|  |  |  | 108 | -. 075 | 29.23 | 30.07 | . 840 |  |  |  |
|  |  |  | 24/0 | . 21 | 29.13 | 29.97 | . 840 |  |  |  |
|  |  |  |  |  |  | avg: | . 830 | 29.165 | 15 | 17 |
|  |  |  | 117 | -. 035 | 29.18 | 30.075 | . 895 |  |  |  |
|  |  |  | 24/0 | . 21 | 29.085 | 29.97 | . 885 |  |  |  |
|  |  |  |  |  |  | avg: | . 890 | 29.115 | 15 | 18 |
|  |  |  | $12 \quad 8$ | -. 065 | 29.14 | 30.07 | . 930 |  |  |  |
|  |  |  | 24/0 | . 215 | 29.045 | 29.97 | . 925 |  |  |  |
|  |  |  |  |  |  | avg: | . 930 | 29.075 | 15 | 19 |
|  |  |  | 138 | -. 075 | 29.105 | 30.065 | . 960 |  |  |  |
|  |  |  | 24/0 | . 205 | 29.01 | 29.965 | . 955 |  |  |  |
|  |  |  |  |  |  | avg: | . 960 | 29.045 | 15 | 20 |
|  |  |  | $14 \quad 8$ | -. 08 | 29.065 | 30.075 | 1.010 |  |  |  |
|  |  |  | 24/0 | . 21 | 28.97 | 29.97 | 1.00 |  |  |  |
|  |  |  |  |  |  | avg: | 1.005 | 29.000 | 15 | 21 |
|  |  |  | 15 8 | -. 09 | 29.02 | 30.075 | 1.055 |  |  |  |
|  |  |  | 24/0 | . 21 | 28.93 | 29.965 | 1.035 |  |  |  |
|  |  |  |  |  |  | avg: | 1.045 | 28.960 | 15 | 22 |
|  |  |  | 168 | -. 10 | 28.985 | 30.065 | 1.080 |  |  |  |
|  |  |  | 24/0 | . 21 | 28.895 | 29.97 | 1.075 |  |  |  |
|  |  |  |  |  |  | avg: | 1.080 | 28.925 | 15 | 23 |
|  |  |  | $18 \quad 8$ | -. 11 | 28.91 | 30.065 | 1.155 |  |  |  |
|  |  |  | 23 | . 225 | 28.81 | 29.97 | 1.160 |  |  |  |
|  |  |  |  |  |  | avg: | 1.155 | 28.850 | 14 | 25 |
|  |  |  | 198 | -. 12 | 28.87 | 30.065 | 1.195 |  |  |  |
|  |  |  | 24/0 | . 21 | 28.74 | 29.97 | 1.230 |  |  |  |
|  |  |  | 25/1 | . 18 | 28.79 | 29.975 | 1.185 |  |  |  |
|  |  |  |  |  |  | avg: | 1.190 | 28.815 | 14 | 26 |

Results of least squares fit of data in (1) and (2) (pages 47-48) according to the theory of Equation 1:
$L=a_{1}+a_{2} t^{t}+a_{3} t^{2}+a_{4} t^{3}$
$a_{1}=(3.6180 \pm 0.01023) \times 10^{-1}$ degrees

$$
\begin{aligned}
& a_{3}=-(4.127 \pm 0.0279) \times 10^{-4} \text { degrees }, ' \text { sid. day }{ }^{2} \\
& a_{4}=(1.039 \pm 1.799) \times 10^{-7} \text { degrees } / \text { sid. day }{ }^{3}
\end{aligned}
$$

$a_{2}=(3.202 \pm 0.0107) \times 10^{-2}$ degrees $/$ sid. day
Standard error of estimate $=4.95 \times 10^{-3}$ degrees
$\ddot{\lambda}($ with minimum standard error $)=-(1.4408 \pm 0.0097) \times 10^{-5} \mathrm{rad} / \mathrm{sid}^{2} \mathrm{day}^{2}$, at $\mathrm{t}=-0.3003$ sid. day, $\mathrm{t}^{\prime}=22 \mathrm{May} 1965, \lambda=-28.703^{\circ}$.
*Ground track (subsatellite point) information on Early Bird, from Andover Maine Tracking, supplied through the Offices of Robert Greene, Comsat Corp., Washington, D.C.
**Calculated as $30.005^{\circ}$ minus average drift from the reference orbit.

Table 9 (Cont.)
Ground Track and Related Ascending Equator Crossing Data for Free Drift Arc 9 (Early Bird)*

|  | Latitude (degrees) | Longitude (degrees west) | Computed <br> Longitude in Reference Orbit (degrees west) | Drift from Reference Orbit (degrees east) | Estimated Ascending Equator Crossing Longitude (degrees west) | Time of Ascending Equator Crossing During Day: Estimate (hours) | Time (sid. day integers from 23 April 1965) |  | Ascending Equator Longitude East of $\stackrel{-}{29.055}$ (degrees) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\begin{array}{cc}20 & 8 \\ & 24 / 0\end{array}$ | -. 12 | 28.835 | 30.065 | 1.230 |  |  |  | -29.5 | -. 950 |
|  | . 205 | 28.75 | 29.97 | 1.220 |  |  |  | -28.5 | -. 890 |
|  | -. 115 |  | avg: | 1.225 | 28.780 | 14 | 27 | -26.5 | -.780 |
| 218 |  | 28.80 | 30.065 | 1.265 |  |  |  | -25.5 | -. 725 |
|  |  |  | avg: | 1.265 | 28.740 | 14 | 28 | -24.5 | -. 665 |
| $\begin{array}{lc}22 & 2 \\ & 8 \\ & 24 / 0\end{array}$ | . 145 | 28.73 | 30.045 | 1.315 |  |  |  | -23.5 | -. 615 |
|  | -. 13 | 28.765 | 30.065 | 1.300 |  |  |  | -22.5 | -. 570 |
|  | . 205 | 28.685 | 29.975 | 1.290 |  |  |  | -21.5 | -. 515 |
|  | -. 135 |  | avg: | 1.300 | 28.705 | 14 | 29 | -20.5 | -. 475 |
| 238 |  | 28.73 | 30.065 | 1.335 |  |  |  | -18.5 | -. 370 |
|  |  |  | avg: | 1.335 | 28.670 | 14 | 30 | -17.5 | -. 320 |
| $24 \quad 8$ | -. 14 | 28.70 | 30.065 | 1.365 |  |  |  | -16.5 | -. 270 |
|  |  |  | avg: | 1.365 | 28.640 | 14 | 31 | -15.5 | -. 230 |
| $\begin{array}{rr}25 & 2 \\ & 8 \\ \\ & 23\end{array}$ | . 12 | 28.64 | 30.05 | 1.410 |  |  |  | -14.5 | -. 190 |
|  | -. 14 | 28.67 | 30.065 | 1.395 |  |  |  | -12.5 | -. 110 |
|  | . 23 | 28.58 | 29.97 | 1.390 |  |  |  | -11.5 | -. 060 |
|  |  |  | avg: | 1.400 | 28.605 | 14 | 32 | -10.5 | -. 020 |
| $\begin{array}{rr}26 & 8 \\ \\ \\ 23\end{array}$ | -. 14 | 28.635 | 30.065 | 1.430 |  |  |  | - 9.5 | +. 010 |
|  | . 225 | 28.555 | 29.97 | 1.415 |  |  |  | - 8.5 | +. 055 |
|  |  |  | avg: | 1.425 | 28.580 | 14 | 33 | - 7.5 | . 095 |
| $\begin{array}{cc}27 & 8 \\ & 25 / 1\end{array}$ | -. 14 | 28.605 | 30.065 | 1.460 |  |  |  | - 6.5 | . 130 |
|  | . 165 | 28.545 | 29.98 | 1.435 |  |  |  | - 4.5 | . 205 |
|  |  |  | avg: | 1.450 | 28.555 | 14 | 34 | - 3.5 | . 240 |
| $\begin{array}{lc}28 & 8 \\ & 24 / 0\end{array}$ | -. 145 | 28.585 | 30.065 | 1.480 |  |  |  | - 2.5 | . 275 |
|  | -. 205 | 28.51 | 29.975 | 1.465 |  |  |  | - 1.5 | . 315 |
|  |  |  | avg: | 1.475 | 28.530 | 14 | 35 | - 0.5 | . 350 |
|  |  |  |  |  |  |  |  | + 0.5 | . 385 |
| $\begin{array}{rr}29 & 8 \\ & 23\end{array}$ | -. 15 | 28.55 | 30.065 | 1.515 |  |  |  | + 1.5 | . 415 |
|  | . 225 | 28.485 | 29.97 | 1.485 |  |  |  | 2.5 | . 450 |
|  |  |  | avg: | 1.500 | 28.505 | 14 | 36 | 3.5 | . 475 |
| $\begin{array}{cc}30 & 8 \\ & 24 / 0\end{array}$ | -. 15 | 28.53 | 30.065 | 1.535 |  |  |  | 4.5 | . 500 |
|  | . 19 | 28.465 | 29.975 | 1.510 |  |  |  | 5.5 | . 525 |
|  |  |  | avg: | 1.525 | 28.480 | 14 | 37 | 6.5 | . 550 |
| $\begin{array}{lc} 31 & 8 \\ & 25 / 1 \end{array}$ | -. 15 | $28.50$ | 30.065 |  |  |  |  | 7.5 | . 575 |
|  | . 17 | 28.445 | 29.98 | 1.535 |  |  |  | 8.5 | . 600 |
|  |  |  | avg: | 1.550 | 28.455 | 14 | 38 | 9.5 | . 625 |
| (June 1965) |  |  |  |  |  |  |  | 10.5 | . 650 |

Results of least squares fit of data in (1) and (2) above according to the theory of Equation 1:
$L=a_{1}+a_{2} t+a_{3} t^{2}+a_{4} t^{3}$
$a_{1}=(3.6180 \pm 0.01023) \times 10^{-1}$ degrees $\quad a_{3}=-(4.127 \pm 0.0279) \times 10^{-4}$ degrees $/ \mathrm{sid}:$ day ${ }^{2}$
$a_{2}=(3.202 \pm 0.0107) \times 10^{-2}$ degrees $/$ sid. day
$a_{4}=(1.039 \pm 1.799) \times 10^{-7}$ degrees $/$ sid. day ${ }^{3}$
Standard error of estimate $=4.95 \times 10^{-3}$ degrees
$\ddot{\lambda}($ with minimum standard error $)=-(1.4408 \pm 0.0097) \times 10^{-5} \mathrm{rad} / \mathrm{sid}$ day ${ }^{2}$, at $t=-0.3003$ sid. day, $t^{\prime}=22 \mathrm{May} 1965, \lambda=-28.703^{\circ}$.
-Ground track (subsatellite point) informiation on Early Bird, from Andover Maine Tracking, supplied through the Offices of Robert Greene, Comsat Corp., Washington, D.C.
**Calculated as $30.005^{\circ}$ minus average drift from the reference orbit.

Table 9 (Cont.)
Ground Track and Related Ascending Equator Crossing Data for Free Drift Arc 9 (Early Bird)*


Results of least squares fit of data in (1) and (2) above according to the theory of Equation 1:
$L=a_{1}+a_{2} t+a_{3} r^{2}+a_{4} t^{3}$
$a_{1}=(3.6180 \pm 0.01023) \times 10^{-1}$ degrees

$$
a_{3}=-(4.127 \pm 0.0279) \times 10^{-4} \text { degrees } / \text { sid. day }{ }^{2}
$$

$a_{2}=(3.202 \pm 0.0107) \times 10^{-2}$ degrees $/ \mathrm{sid}$. day

Standard error of estimate $=4.95 \times 10^{-3}$ degrees
$\ddot{\lambda}($ with minimum standard ertor $)=-(1.4408 \pm 0.0097) \times 10^{-5} \mathrm{rad} / \mathrm{sid}^{\mathrm{s}}$ day ${ }^{2}$, at $t=-0.3003$ sid. days, $t^{\prime}=22 \mathrm{May} 1965, \lambda=-28.703^{\circ}$.
*Ground track (subsatellite point) information on Early Bird, from Andover Maine Tracking, supplied through the Offices of
Robert Greene, Comsat Corp., Washington, D.C.
**Calculated as $30.005^{\circ}$ minus average drift from the reference orbit.
***Data not used in acceleration analysis as inclusion gives unacceptably large residuals for this arc.

Ascending Equator Crossing Orbit Data from a Simulated Early Bird Trajectory for Free Drift Are 9 Computed by "ITEM" with Earth Longitude Gravity Through Third Order.*



Figure 8-Measured and simulated orbit data at ascending Equator crossings in free drift arc 9 (Early Bird).

## 2. SYNTHESIS OF THE LONGITUDE ACCELERATION RECORD TO REVEAL components in the earth's longitude gravity field

According to Equation 2, the long term resonant earth gravity accelerated longitude drift of the 24 -hour satellite is given very closely through fourth order, by

$$
\begin{align*}
\ddot{\lambda} & =-12 \pi^{2}\left[\frac{6}{a_{s}{ }^{2}} F(i)_{22}\left\{C_{22} \sin 2 \lambda-S_{22} \cos 2 \lambda\right\}\right. \\
& +\frac{45}{a_{s}{ }^{3}} F(i)_{33}\left\{C_{33} \sin 3 \lambda-S_{33} \cos 3 \lambda\right\}-\frac{3 F(i)_{31}}{2 a_{s}{ }^{3}}\left\{C_{31} \sin \lambda-S_{31} \cos \lambda\right\} \\
& \left.+\frac{420 F(i)_{44}}{a_{s}{ }^{4}}\left\{C_{44} \sin 4 \lambda-S_{44} \cos 4 \lambda\right\}-\frac{15 F(i)_{42}}{a_{s}{ }^{4}}\left\{C_{42} \sin 2 \lambda-S_{42} \cos 2 \lambda\right\}\right] \tag{11}
\end{align*}
$$



Figure 9 -Measured and geoid-predicted 24 -hour satellite longitude drift accelerations.
where $\ddot{\lambda}$ is the longitude acceleration in units of radians/sid. day ${ }^{2}$ and $a_{s}$ is the "synchronous" semimajor axis of the satellite in earth radii. The $F(i)_{n m}$ functions of the satellite inclination $i^{\prime} s$ are

$$
\begin{align*}
& F(i)_{22}=\left[\frac{1}{2}\left(\cos i_{s}+1\right)\right]^{2} \\
& F(i)_{33}=\left[\frac{1}{2}\left(\cos i_{s}+1\right)\right]^{3} \\
& F(i)_{31}=\left[\frac{1}{2}\left(\cos i_{s}+1\right)-\frac{5}{8} \sin ^{2} i_{s}\left(1+3 \cos i_{s}\right)\right]  \tag{12}\\
& F(i)_{44}=\left[\frac{1}{2}\left(\cos i_{s}+1\right)\right]^{4} \\
& F(i)_{42}=\left[\frac{1}{4}\left(\cos i_{s}+1\right)^{2}-\frac{7}{4} \sin ^{2} i_{s}\left(\cos i_{s}+1\right)\right] .
\end{align*}
$$

The $C_{n m}, S_{n m}$ gravity coefficients are given in terms of the $J_{n m}, \lambda_{n m}$ coefficients by

$$
\left.\begin{array}{l}
\mathrm{c}_{\mathrm{nm}}=\mathrm{J}_{\mathrm{nm}} \cos \mathrm{~m} \lambda_{\mathrm{nm}}  \tag{13}\\
\mathrm{~s}_{\mathrm{nm}}=\mathrm{J}_{\mathrm{nm}} \sin \mathrm{~m} \lambda_{\mathrm{nm}}
\end{array}\right\}
$$

In Equations 13, the $J_{n m}$ are all negative so that the $\lambda_{n m}$ gravity harmonic phase angles with respect to Greenwich are interpreted physically as in Appendix B. Thus the proper quadrant for $\lambda_{\mathrm{nm}}$ from the $\mathrm{C}_{\mathrm{nm}}$ and $\mathrm{S}_{\mathrm{nm}}$ is determined by

$$
\begin{equation*}
\lambda_{\mathrm{nm}}=\frac{1}{\mathrm{~m}} \tan ^{-1}\left[\frac{-\mathrm{S}_{\mathrm{nm}}}{-\mathrm{C}_{\mathrm{nm}}}\right] \tag{14}
\end{equation*}
$$

The $J_{n m}$ are given from the $C_{n m}$ and $S_{n m}$ by

$$
\begin{equation*}
J_{n m}=-\left[C_{n m}^{2}+S_{n m}^{2}\right]^{1 / 2} \tag{15}
\end{equation*}
$$

From a set of measured long term 24-hour satellite accelerations $\ddot{\lambda}$ at longitudes $\lambda$, with inclinations $i_{s}$ and semimajor axes $a_{s}$, it is possible to determine the $C_{n m}, S_{n m}$ in Equation 11 which best satisfy this set of accelerations. If the set of measurements numbers the same as the $C_{n m}, S_{n m}$
coefficients, then clearly the $\mathrm{C}_{\mathrm{nm}}$, $\mathrm{S}_{\mathrm{nm}}$ can be determined from them by a simple simultaneous solution of the set of Equations 11 with the specified measurements. However, such a straightforward solution for the underlying gravity field assumes that Equation 11 is an exact equation, if, for example, ten measurements are available. Actually, as discussed in the introduction, the measurements of $\ddot{\lambda}$ will be in error from true resonant gravity acceleration due to a number of sources.

A more realistic model from which to determine the dominating gravity effects from the measured accelerations is the general linear least squares model. In this model, we allow each of the measurements j to be conditioned by Equation 11 (with specified $\mathrm{C}_{\mathrm{nm}}$ ' s and $\mathrm{S}_{\mathrm{nm}}$ ' s as unknowns) with the addition of a small unknown error $\epsilon_{j}$. We require more measurements than unknowns in order to allow for these additional unknown $\epsilon_{\mathrm{j}}$ 's. In the least squares solution, the otherwise overdetermined $C_{n m}$ and $S_{n m}$ are adjusted so that the sums of the squares of the residuals of $\ddot{\lambda}$ (measured, or the left hand side of Equation 11) and $\ddot{\lambda}$ (theoretical, or the right hand side of Equation 11) are a minimum. These residuals are an estimate of the unknown $\epsilon_{j}$. Furthermore, if we assume that the $\epsilon_{j}$ are random normally distributed with mean of zero and constant variance $\sigma$, we can estimate $\sigma$ and make statistically significant statements about the likely variation of the $\mathrm{C}_{\mathrm{nm}}$ and $\mathrm{S}_{\mathrm{nm}}$ coefficients in any given test of the measured data according to Equation 11. (For general treatments of the least squares model, see References 13 and 14.) In Tables 10 and 11 we list the results of the acceleration analysis in the actual and simulated 24 -hour satellite gravity drift arcs. Since we have available less than ten independent acceleration measurements, we cannot hope to determine all ten resonant gravity harmonics from a solution of Equation 11. In addition, of course, we tacitly ignore in this analysis the effects on the data of the infinite set of earth resonant gravity harmonics of order higher than fourth. It seems evident, even before looking at the data, that from four to six harmonics is all that might reasonably be extracted from a least squares solution of Equation 1 (ignoring selected coefficients).

In Appendix $C$ we have calculated maximum resonant gravity effects from a recent (1965) geoid due to W. H. Guier (Reference 15). From this calculation, it is evident that $\mathrm{C}_{22}, \mathrm{~S}_{22}, \mathrm{C}_{33}$ and $\mathrm{S}_{33}$ should be the dominating harmonics on the three 24 -hour satellites in this study. The next most influential set of harmonics on all the satellites appears to be $C_{31}$ and $S_{31}$, although on Syncom 2, $\mathrm{C}_{44}$ and $\mathrm{S}_{44}$ might rival it in influence according to other recent geoids listed in Table B1.

It is interesting to compare the standard acceleration error in the actual experiment with these maximum theoretical effects. Except for arc 3 (not used), arc 7 and about half of the 18 sub-arcs in arc 5 , the standard acceleration error measures between $0.03 \times 10^{-5}$ and $0.10 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$. These levels are over an order of magnitude below the maximum theoretical $\mathrm{J}_{22}$ caused accelerations, and significantly below the theoretical maximum $\mathrm{J}_{33}$ effects on all the satellites. We are consequently encouraged to believe that the four harmonics $\mathrm{C}_{22}, \mathrm{~S}_{22}, \mathrm{C}_{33}$ and $\mathrm{S}_{33}$ should be the minimum yield from the synthesis (or explanation) of the measured accelerations in Table 10 according to the theory of Equation 11. The theoretical maximum $J_{31}$ effects appear to be at or just above the average noise level of this experiment, while the $\mathrm{J}_{42}$ and $\mathrm{J}_{44}$ effects appear to be somewhat below that level. While it would seem that not enough data is yet available to distinguish these harmonics with good precision, we are encouraged to hope that the apparently low errors in
the widely separated equatorial arcs 6 and 9 will allow at least a tentative discrimination of $C_{31}$ and $S_{31}$ (see Discussion).

We now proceed to test the actual data in Table 10 for sensitivity to $C_{22}, S_{22}$ alone and then together with $\mathrm{C}_{33}, \mathrm{~S}_{33}$, and finally in combination with $\mathrm{C}_{31}, \mathrm{~S}_{31}$ as well, according to a least squares model based on Equation 11 (see Table 12). The blanks in Table 12 indicate that particular harmonic was not considered in the test.

The first three tests in Table 12, with the unadjusted data weighted according to the measured standard errors in Table 10, shows the general trend of the results. The weighting scheme chosen for these tests assumes that samples with standard errors less than $0.05 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$ contain predominantly model bias errors and (in this unadjusted test) carry equal, unit weight. There is some justification for this assumption based on the bias results in Table 11. In this test, arc 5 was assigned a weight of 1.0 on the basis of the sum of the independent arcs 5 A and 5 B weights (column 3 of Table 10). This unit weight was divided among the 18 sub-arcs according to an independent arc 5 weighting scheme which gave unit weight to the best determined sample, that of sub-arc 5-18.

There are two strong conclusions which can be drawn at once from the first three tests in Table 12. The first is that the 24 -hour longitude coverage around the equator is now so complete that we can almost define the dominant $\mathrm{C}_{22}, \mathrm{~S}_{22}$ harmonics without regard for the presence of higher order effects. (See also tests 28-30.) Theoretically, since the potential is an infinite series of orthogonal Legendre functions, a least squares fit of the actual potential through all space with respect to any combination of potential harmonics with determinable coefficients will yield the true Legendre coefficients of the potential. It appears a reasonable conjecture that, because of the natural suppression of higher order effects, near convergence to the true low order gravity potential coefficients should be possible from a complete longitude survey at 24-hour altitudes.

The second conclusion from tests $\mathbf{1 - 3}$ is that 24 -hour satellite drift to date is now strongly sensitive to $\mathrm{C}_{33}$ and $\mathrm{S}_{33}$, or at least third order earth longitude gravity. The sensitivity to higher order gravity is shown by the over fivefold reduction in the standard error of these tests when higher order effects are allowed. In fact, the higher order tests bring the standard acceleration error of the test to the level of the individual acceleration standard errors.

In addition to this dramatic reduction of the residuals upon allowance for third order effects, we note a similarly large reduction in the standard errors of the $H_{22}$ coefficients. If we can attribute the residuals of the $\mathrm{H}_{22}, \mathrm{H}_{33}$ test 2 purely to "observation noise," then we would expect the inclusion of the $\mathrm{H}_{31}$ harmonic (test 3) to make negligible change in the previously determined coefficients. We could also expect no clear $\mathrm{H}_{31}$ result, as well as an increase in the test- $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ standard errors, because the degrees of freedom of this limited sample would have been reduced at no comparable improvement of the fit.

Longitude Accelerations in 24 -Hour Satellite Arcs 1 to 9

| (Satellite)-Arc | (1) <br> Longitude Acceleration Measured, $\pi$ $\left(1 \sigma^{5} \mathrm{rad} / \mathrm{sid}^{2} . \mathrm{day}^{2}\right)$ | Longitude, <br> $\lambda$ (degrees) | Semimajor Axis, $a_{s}$ (eatth radii) | $\begin{gathered} \text { Orbit } \\ \text { Inclination, } \\ i_{s} \\ \text { (degrees) } \end{gathered}$ | Standard <br> Errot of Longitude Acceleration, $o$ ( $10^{-5} \mathrm{rad} / \mathrm{sid}$. day $^{2}$ ) | (2) <br> Estimated Acceleration Model Bias* $\left(10^{-5} \mathrm{rad} / \mathrm{sid} . \mathrm{day}^{2}\right)$ | (1) + (2) Bias Adjusted Acceleration, $\dot{\lambda}$ (adj.) $\left(10^{-5} \mathrm{rad} /\right.$ sid. day $\left.{ }^{2}\right)$ | (3) <br> Weight of Sample (If $\sigma \leq 0.05 \times 10^{-5}$ $\mathrm{rad} / \mathrm{sid} . \mathrm{day}^{2}$, $\mathrm{wt} .=1.0$ ) | $\begin{gathered} \text { (4) } \\ \text { Weight of Sample } \\ \text { in Arc } 5 \\ \text { (If } \sigma \leq 0.616 \times 10^{-5} \\ \text { rad } / \text { sid. day }{ }^{2}, \\ w t .=1.0 \text { ) } \end{gathered}$ | Number of Orbits Considered in Acceleration Determination | $\begin{aligned} & \text { Arc Time } \\ & \text { (span) } \end{aligned}$ | Arc Time (months) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| (SYNCOM 2) 1 | -2.253 | - 55.22 | 6.611113 | 33.024 | 0.0325 | +. 015 | -2.238 | 1.000 |  | 19 | AUG-DEC 1963 | 3 |
| (SYNCOM 2) 2 | -2.291 | - 60.94 | 6.611618 | 32.825 | 0.0572 | -. 029 | -2.320 | . 765 |  | 16 | DEC-MAR 1963/64 | $31 / 2$ |
| (SYNCOM 2) ${ }^{* * * 3}$ | -0.897 | - 88.00 | 6.62675 | 32.67 | 0.888 |  |  | . 000 |  | 6 | MAR-APR 1964 | 1 |
| (SYNCOM 2) 4 | 2.138 | -140.00 | 6.620443 | 32.584 | 0.0842 | +. 015 | 2.153 | . 353 |  | 10 | APR-JULY 1964 | 2 |
| (SYNCOM 2) 5A | -0.199 | 161.00 | 6.616521 | 0.397 | 0.0661 | -. 011 | -. 210 | . 570 |  | 17 | JULY-NOV 1964 | $41 / 2$ |
| (SYNCOM 2) 5 | -2.295 | 134.00 | 6.617 | 0.33 | 0.0397 | +. 011 | -2.284 | 1.000 |  | 26 | JULY-FEB 1964/65 | $71 / 2$ |
| (SXNCOM 2) 5B | -2.389 | 106.00 | 6.617425 | 0.16 | 0.0724 | -. 022 | -2.411 | . 479 |  | 10 | NOV-FEB 1964/65 | $31 / 2$ |
| (SYNCOM 2) 5-1 | 1.066 | 175.5 | 6.6165 | 0.47 | 0.131 | - . 045 | 1.021 | **** | 0.252 | 10 | JULY-SEPT 1964 | 2 |
| (SXNCOM 2) 5-2 | 0.954 | 172.5 | 6.6165 | 0.45 | 0.0887 | -. 029 | 0.925 |  | 0.550 | 10 | JULY-SEPT 1964 | 2 |
| (SYNCOM 2) 5-3 | 0.672 | 169.0 | 6.6164 | 0.43 | 0.0926 | +. 006 | 0.678 |  | 0.501 | 10 | JULY-SEPT 1964 | 2 |
| (SYNCOM 2) 5-4 | 0.394 | 166.0 | 6.6164 | 0.40 | 0.0819 | - . 021 | 0.373 |  | 0.645 | 10 | JULY-SEPT 1964 | 2 |
| (SYNCOM 2) 5-5 | 0.016 | 163.0 | 6.6164 | 0.39 | 0.0822 | -. 006 | 0.010 |  | 0.640 | 10 | AUG-OCT 1964 | 2 |
| (SYNCOM 2) 5-6 | -0.157 | 160.0 | 6.6164 | 0.38 | 0.129 | -. 002 | -0.159 |  | 0.260 | 10 | AUG-OCT 1964 | 2 |
| (SYNCOM 2) 5-7 | -0.541 | 157.0 | 6.6164 | 0.37 | 0.167 | +. 019 | -0.522 |  | 0.155 | 10 | AUG-OCT 1964 | 2 |
| (SYNCOM 2) 5-8 | -0.788 | 153.5 | 6.6164 | 0.35 | 0.163 | -. 003 | -0.791 |  | 0.162 | 10 | AUG-OCT 1964 | 2 |
| (SYNCOM 2) 5-9 | -1.103 | 150.0 | 6.6165 | 0.34 | 0.160 | +. 036 | -1.067 |  | 0.169 | 10 | SEPT-NOV 1964 | 2 |
| (SYNCOM 2) 5-10 | -1.560 | 146.5 | 6.6166 | 0.32 | 0.185 | +. 034 | -1.526 |  | 0.126 | 10 | SEPT-NOV 1964 | 2 |
| (SYNCOM 2) 5-11 | -1.990 | 143.5 | 6.6167 | 0.31 | 0.181 | +. 030 | -1.960 |  | 0.132 | 10 | SEPT-NOV 1964 | 2 |
| (SYNCOM 2) 5-12 | -2.138 | 139.5 | 6.6167 | 0.30 | 0.213 | +. 018 | -2.120 |  | 0.095 | 10 | SEPT-NOV 1964 | 2 |
| (SYNCOM 2) 5-13 | -2.501 | 132.0 | 6.6169 | 0.22 | 0.167 | -. 033 | -2.534 |  | 0.154 | 10 | OCT-JAN 1964/65 | $31 / 2$ |
| (SYNCOM 2) 5-14 | -2.828 | 120.0 | 6.6171 | 0.22 | 0.122 | +. 037 | -2.791 |  | 0.290 | 10 | OCT-JAN 1964/65 | $31 / 2$ |
| (SYNCOM 2) 5-15 | -2.663 | 116.0 | 6.6172 | 0.20 | 0.102 | $+.025$ | -2.638 |  | 0.415 | 10 | OCT-JAN 1964/65 | $31 / 2$ |
| (SYNCOM 2) 5-16 | -2.584 | 112.0 | 6.6173 | 0.18 | 0.0907 | -. 005 | -2.589 |  | 0.525 | 10 | OCT-JAN 1964/65 | $31 / 2$ |
| (SYNCOM 2) 5-17 | -2.474 | 209.0 | 6.6174 | 0.16 | 0.0847 | -. 033 | -2.507 |  | 0.600 | 10 | NOV-FEB 1964/65 | $31 / 2$ |
| (SYNCOM 2) 5-18 | -2.278 | 104.5 | 6.6176 | 0.15 | 0.0656 | -. 041 | -2.319 |  | 1.000 | 10 | NOV-FEB 1964/65 | $31 / 2$ |
| (SYNCOM 3) 6 | 1.707 | 178.707 | 6.611474 | 0.113 | 0.0591 | -. 024 | 1.683 | 0.715 |  | 10 | NOV-DEC 1964 | 2 |
| (SYNCOM 3) 7 | 1.550 | 176.801 | 6.612269 | 0.268 | 0.175 | +. 027 | 1.577 | 0.082 |  | 9 | JAN-MAR 1965 | 2 |
| (SYNCOM 2) 8 | 0.950 | 66.115 | 6.611199 | 31.869 | 0.0616 | +. 018 | 0.968 | 0.660 |  | 11 | FEB-MAY 1965 | 21/2 |
| (EARLY BIRD) 9 | -1.441 | - 28.703 | 6.6105 | 0.200 | 0.010 | -. 031 | -1.472 | 1.000 |  | 46 | APR-JUNE 1965 | 1 |
| **5' | 1-2.528 | 129.00 | 6.617 | 32.33 | 0.11 |  |  |  |  | 26 | JULY-FEB 1964/65 | $71 / 2$ |

*From a compromise of the biases reported in Table 11.
**Independent full are 5 acceleration from drift rate data reduced by the $\mathrm{J}_{22}$ model, Equation B , determined from successive equator crossings in each orbit only (velocity data).
***Not used in the gravity synthesis.
$* * *$ In weighted analyses interdependent ares $5-1-5-18$ were given total wt. = 1.0 (replacing arcs 5 A \& 5 B), distributed according to column 4 weights. In "unweighted" analyses, arcs $5-1-5-18$ were given total weight $=2.0$ without predjudice,
all other independent arcs used being given wt. = 1.0.

Longitude Accelerations in Simulated 24 -Hour Satellite Arcs $1 \mathrm{~S} / 1$ to $9 \mathrm{~S} / 1^{*}$.

| Arc | (1) <br> Longitude Acceleration, $\lambda$ $10^{-5} \mathrm{rad} / \mathrm{sid}^{2} \text { day }^{2} \text { ) }$ | Longitude, $\lambda$ (degrees) | $\begin{gathered} \text { Semimajor } \\ \text { Axis, } \\ a_{3} \\ \text { (earth radii) } \end{gathered}$ | Orbit Inclination, $i_{s}$ (degrees) | Standard Error of Longitude Acceleration, $\sigma t$ $\left(10^{-5} \mathrm{rad} / \mathrm{sid}\right.$. day $\left.^{2}\right)$ | (2) <br> Theoretical $\ddot{\lambda}$ for Given $\mathbf{a}_{s}, \mathbf{i}_{\mathbf{s}}$, and $\lambda^{* *}$ $\left(10^{-5} \mathrm{rad} / \mathrm{sid}\right.$. day $\left.^{2}\right)$ | Column (1) minus Column (2), Longitude Acceleration Model Bias $\left(10^{-5} \mathrm{rad} / \mathrm{sid}\right.$. day $\left.^{2}\right)$ | Bias from a Second Order Earth Longitude Gravity Trajectory $\dagger \dagger$ $\left(10^{-5} \mathrm{rad} / \mathrm{sid}\right.$. day $\left.^{2}\right)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1S/1 | -2.2328 | - 55.13 | 6.6111 | 33.03 | . 0048 | -2.2185 | +. 0143 | +. 018 |
| 2S/1 | -2.2024 | - 60.91 | 6.6116 | 32.84 | . 0056 | -2.2330 | -. 0306 | -. 027 |
| 4S/1 | 2.156 | -140.00 | 6.6204 | 32.57 | . 0413 | 2.163 | +. 007 | +. 025 |
| 6S/1 | 1.686 | 178.69 | 6.6115 | 0.09 | . 0209 | 1.6615 | -. 0245 | -. 023 |
| 7S/1 | 1.473 | 176.915 | 6.6123 | 0.11 | . 0129 | 1.501 | +. 028 | +. 026 |
| 8S/1 | 0.862 | 66.255 | 6.61115 | 31.87 | . 0219 | . 884 | +. 022 | +. 014 |
| 9S/1 | -1.349 | -28.69 | 6.6105 | 0.20 | . 0064 | -1.380 | -. 031 | -. 032 |
| 5S/1-1 | 1.1915 | 175.5 | 6.6164 | 32.47 | . 0305 | 1.146 | -. 0455 |  |
| 5S/1-2 | . 927 | 172.5 | 6.6164 | 32.46 | . 0304 | . 898 | -. 029 |  |
| 5S/1-3 | . 5885 | 169.0 | 6.6164 | 32.44 | . 0230 | . 594 | +. 0055 |  |
| 5S/1-4 | . 341 | 166.0 | 6.6164 | 32.42 | . 0232 | . 321 | -. 021 |  |
| 5S/1-5 | . 049 | 163.0 | 6.6164 | 32.40 | . 0266 | . 043 | -. 006 |  |
| 5S/1-6 | - . 237 | 160.0 | 6.6164 | 32.39 | . 0274 | - . 239 | -. 002 |  |
| 5S/1-7 | - . 540 | 157.0 | 6.6164 | 32.38 | . 0243 | - . 521 | +.019 |  |
| 5S/1-8 | - . 842 | 153.5 | 6.6164 | 32.36 | . 0248 | -. 845 | -. 003 |  |
| 5S/1-9 | -1.194 | 150.0 | 6.6165 | 32.35 | . 0361 | -1.1585 | +. 0355 |  |
| 5S/1-10 | -1.449 | 147.0 | 6.6165 | 32.34 | . 0427 | -1.415 | +. 034 |  |
| 5S/1-11 | -1.726 | 143.5 | 6.6166 | 32.32 | . 0420 | -1.6955 | +. 0305 |  |
| 5S/1-12 | -2.003 | 139.5 | 6.6167 | 32.31 | . 0392 | -1.985 | +. 018 |  |
| 5S/1-13 | -2.3835 | 132.0 | 6.6171 | 32.26 | . 0310 | -2.416 | -. 0325 |  |
| 5S/1-14 | -2.768 | 120.0 | 6.6172 | 32.25 | . 0266 | -2.731 | +. 037 |  |
| 5S/1-15 | -2.752 | 116.5 | 6.6173 | 32.23 | . 0272 | -2.727 | +. 025 |  |
| 5S/1-16 | -2.662 | 112.5 | 6.6174 | 32.22 | . 0380 | -2.667 | -. 005 |  |
| 5S/1-17 | -2.535 | 109.0 | 6.6175 | 32.19 | . 0404 | -2.568 | -. 033 |  |
| 5S/1-18 | -2.363 | 105.0 | 6.6176 | 32.16 | . 0350 | -2.404 | -. 041 |  |
| ***5S-A |  |  |  |  | . 0132 |  |  | -. 0107 |
| ***5S |  |  |  |  | . 0100 |  |  | +.0107 |
| ***5S-B |  |  |  |  |  |  |  |  |

*Trajectories computed numerically by ITEM in the presence of sun and moon gravity through third order earth longitude gravity field given by:

$$
\begin{aligned}
& \mathrm{J}_{22}=-1.8 \times 10^{-6}, \lambda_{22}=-15.35^{0} \\
& \mathrm{~J}_{33}=-0.16 \times 10^{-6}, \lambda_{33}=24^{0} \\
& \mathrm{~J}_{31}=-1.5 \times 10^{-6}, \lambda_{31}=0^{\circ} .
\end{aligned}
$$

Accelerations derived by a second order gravity drift model (see Tables 1S/1 through 9S/1)
**Computed from Equation 2 with the gravity and orbit constants above.
***Data from Table 5S; simulated trajectory with sun and moon gravity and second order earth longitude gravity
$\dagger+$ See Tables $1 \mathrm{~S}-9 \mathrm{~S}$.
+Data from arcs $1 \mathrm{~S}-9 \mathrm{~S}$ except for 5S/1-1 through 5S/1-18

Tests of 24 -Hour Satellite Accelerations for Sensitivity to Resonant Earth Gravity Harmonics.*
(All values in units of $10^{-6}$ except as noted)


## Comments

Reduced from unadjusted data: weighted according to $\sigma$ of measured accelerations: relative arc $5 \mathrm{wt} .=1.0$ : (see Table 10 and text): total samples $=100$ Reduced from unadjusted data: weighted according to $\sigma$ of measured accelerations: relative arc $5 \mathrm{wt}=1.0$ : (see Table 10 and text): total samples $=100$ Reduced from unadjusted data: weighted according to $\sigma$ of measured accelerations: relative arc $5 \mathrm{wt}=1.0$ : (see Table 10 and text): total samples $=100$ Reduced from bias adjusted data: weighted according to $\sigma$ of measured accelerations: relative arc 5 wt . $=1.0$ : (see Table 10 and text): total samples $=100$ Reduced from bias adjusted data: weighted according to $\sigma$ of measured accelerations: relative arc $5 \mathrm{wt} .=1.0$ : (see Table 10 and text): total samples $=100$ Reduced from bias adjusted data: welghted according to $\sigma$ of measured accelerations: relative arc $5 \mathrm{wt}=1.0$ : (see Table 10 and text): total samples $=100$ ( $\mathrm{A} / \mathrm{B} / \mathrm{C}$ ) reduced from unadjusted data: 3 random choices $(\mathrm{A} / \mathrm{B} / \mathrm{C})$ from normal distributions specified by $\ddot{\lambda}$ and $\sigma$ in Table 10 ( 9 data for each test): total samples $=9$ ( $\mathrm{A} / \mathrm{B} / \mathrm{C}$ ) reduced from unadjusted data: 3 random choices ( $\mathrm{A} / \mathrm{B} / \mathrm{C}$ ) from normal distributions specified by $\dot{\lambda}$ and $\sigma$ in Table 10 (same data as test 7 ): total samples =9 ( $\mathrm{A} / \mathrm{B} / \mathrm{C}$ ) reduced from unadjusted data: 3 random choices ( $\mathrm{A} / \mathrm{B} / \mathrm{C}$ ) from normal distributions specified by $\tilde{\lambda}$ and $\sigma$ in Table 10 (same data as test 7): total samples = 9 Reduced from unadjusted data: uses Table 10 values (unweighted), includes independent arc ${ }^{5}$ ' measurement: total samples $=10$
Reduced from unadjusted data: uses Table 10 values (unweighted), includes independent arc 5 ' measurement: total samples $=10$
Reduced from unadjusted data: uses Table 10 values (unweighted), includes independent arc 5 ' measurement: total samples $=10$
Reduced from unadjusted data: uses Table 10 values, unweighted, without arc $5-11$ : total samples $=100$
Reduced from unadjusted data: uses Table 10 values, unweighted, without arc 5-11: total samples $=100$
Reduced from unadjusted data: uses Table 10 values, unweighted, without arc $5-11$ : total samples $=100$
Reduced from simulated data in Table 11 (geoid: Wagner-Kaula combined 1964/65 through 3rd order), includes sun and moon effects; unweighted, without arc $S 5 / 1-11$ : total samples $=100$ Reduced from simulated data in Table 11 (geoid: Wagner-Kaula combined 1964/65 through 3rd order), includes sun and moon effects; unweighted, without arc $\operatorname{S5} / 1-11$ : total samples $=100$ Reduced from simulated data in Table 11 (geoid: Wagner-Kaula combined $1964 / 65$ through 3 rd order), includes sun and moon effects; unweighted, without arc $55 / 1-11$ : total samples $=100$ (Actual geoid harmonics: Wagner-Kaula combined (1964/65))
19 Reduced from simulated data calculated according to Equation 2 from harmonic coefficients of geoid of Kaula (1964); unweighted, without arc S5/1-11: total samples = 100
20 Reduced from simulated data calculated according to Equation 2 from harmonic coefficients of geoid of Kaula (1964); unweighted, without arc $55 / 1-11$ : total samples $=100$
21 Reduced from simulated data calculated according to Equation 2 from harmonic coefficients of geoid of Kaula (1964); unweighted, without arc S5/1-11: total samples $=100$ (Actual geoid harmonics: Kaula (1964), including: $J_{42}=-0.117 \times 10^{-6}, \lambda_{42}=42.3^{\circ}, \mathrm{J}_{44}=-0.0104 \times 10^{-6}, \lambda_{44}=14.5^{\circ}$ )
22 Reduced from simulated data calculated according to Equation 2 from harmonic coefficients of geoid of Guier (1965); unweighted, without arc $\mathbf{5 5} / 1-11$ : total samples $=100$ 23 Reduced from simulated data calculated according to Equation 2 from harmonic coefficients of geoid of Guier (1965); unweighted, without arc $\mathbf{S 5} / 1-11$ : total samples $=100$ 24 Reduced from simulated data calculated according to Equation 2 from harmonic coefficients of geoid of Guier (1965); unweighted, without arc S5/1-11: total samples $=100$ (Actual geoid harmonics: Guier (1965), including: $J_{42}=-0.193 \times 10^{-6}, \lambda_{42}=23.4^{\circ}, \mathrm{J}_{44}=-0.006 \times 10^{-6}, \lambda_{44}=34.5^{\circ}$ )
25 Reduced from simulated data calculated according to Equation 2 from harmonic coefficients of geoid of Wagner-Guier (1965); unweighted, without arc S5/1-11: total samples = 100
27
Reduced from simulated data calculated according to Equation 2 from harmonic coefficients of geoid of Wagner-Guier (1965); unweighted, without arc S5/1-11: total samples $=100$ (Actual geoid harmonics: Wagner-Guier (1965), including: $\mathrm{J}_{42}=-0.19 \times 10^{-6}, \lambda_{42}=42.3^{\circ}, J_{44}=-0.006 \times 10^{-5}, \lambda_{44}=34.5^{\circ}$ )
28 Reduced from unadjusted data in Table 10, unweighted, includes independent arc $5^{\prime}$ measurement; without ares 5-1, 5-10, 5-11: total samples = 111
29 Reduced from unadjusted data in Table 10 , unweighted, includes independent arc $5^{\prime}$ measurement; without ares $5-1,5-10,5-11$ : total samples $=111$
30 Reduced from unadjusted data in Table 10 , weighted according to $\sigma$ of measured accelerations: relative arc $5 \mathrm{wt}=1.0$ (see Table 10): total samples $=100$

The results of test 3 (see also tests 28-30) are strongly suggestive but not conclusive as to the sensitivity of 24 -hour satellite drift thus far to other than second and third order sectorial ( $\mathrm{n}=\mathrm{m}$ ) earth gravity. Test 3 (and almost all the other $\mathrm{H}_{22}, \mathrm{H}_{33}$, and $\mathrm{H}_{31}$ tests in Table 12) does produce small changes in the $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ coefficients and a slight reduction in the test standard error, as well as a fairly strong $\mathrm{C}_{31}$ coefficient with a standard error of about $20 \%$. These signs point to some significance in the $H_{31}$ results of test 3 . On the other hand, since the longitude and sample survey is limited in these tests, the small changes in the $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ coefficients might be the result of random "observation" or model bias error in the data. The slight decrease in the test standard error may be equally fortuitous. The increase in the standard errors of the $H_{22}$ and $H_{33}$ coefficients is the strongest indication that the $H_{31}$ harmonic is not yet clearly sensed by the 24-hour data so far. In the future, with more data of similar quality, we would expect the standard test error to remain in the vicinity of $(3-5) \times 10^{-7} \mathrm{rad} / \mathrm{sid}$. day $^{2}$ upon the inclusion of third and fourth order terms. At the same time, we would expect the standard errors of the lower order coefficients to decline consistent with coefficient convergence as the dominating higher order terms are brought into the synthesis one by one or in combinations.

We now want to determine as closely as possible the effects on these acceleration data reductions which can be attributed to sun and moon gravity as well as other model bias errors inherent in the method of acceleration analysis. First, we rerun tests $1-3$, adjusting the acceleration samples by the probable model biases determined from the simulated trajectories in the previous section (see Table 10). The results of these data adjusted tests (Table 12, 4-6) show minor and insignificant changes from the unadjusted results. In fact, the unadjusted accelerations appear to be even closer, on the average, to true resonant gravity accelerations, judging by the smaller standard errors throughout the unadjusted data tests. Evidently, the model errors in the actual data have acted to cancel the "observation" errors (in the orbit determinations) more often than not. As a result, the conclusions of this study have been drawn primarily from the unadjusted data.

Before proceeding with independent tests of the simulated data, we would like to make a few more tests of the actual data to better judge the true latitude in the harmonics which is allowed by the measurements. In these tests (7-12 in Table 12), we use the unweighted, unadjusted accelerations in Table 10 and replace the interdependent 18 sub-arcs of arc 5 by the independent measurements for arcs 5 A and 5B. In tests $10-12$ the independently determined arc 5 ' measurement is also included. Tests 7-9 each used three random samplings ( $A / B / C$ ) of a normal distribution given by the mean $\ddot{\lambda}$ and $\sigma$ values for these accelerations in Table 10. On the assumption that there is no bias in the accelerations and all the measured $\sigma^{\prime}$ s arise from random observation error, tests 7-9 give an example of the widest latitude permitted by the measurements. From the results of tests 4-6, there is no reason to expect randomly chosen bias adjusted data to yield significantly more divergent results. Indeed, since the sun, moon, and insufficient model introduce pseudo-random "noise" (of up to $\sigma / 2$ ) into the accelerations as well as biases (see Table 11), the "noise" level attributable to the observations alone should be somewhat less than the $\sigma^{\prime}$ s reported in Table 10. Thus, random tests 7-9 should be reasonably conservative as to the divergent results permitted by the data. In view of the apparent oddness of $\lambda_{31}$ that is implicit in the "mean" accelerations of Table 10 (see Conclusions), it is interesting and perhaps significant that two of three random
samplings in these tests ( A and C) yield best $\lambda_{3_{1}}$ 's in the neighborhood of $-80^{\circ}$. This is about $90^{\circ}$ from the best $\lambda_{31}$ determined from the "mean" accelerations. Test 9A, in fact, shows a standard test error as low as similar tests with the "mean" accelerations. The relatively great range of $H_{31}$ harmonics revealed in test 9 further emphasizes the tentative nature of the best reported values for these quantities. The random tests resulted in a far smaller divergence of $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ harmonics which increase our confidence in the best reported values for them (see Conclusions). Tests $10-12$ with unweighted data are consistent with the previous actual weighted data tests which used 18 sub-arcs of arc 5 .

Finally, a series of unweighted data tests was made using the dense longitude coverage of arc 5 provided by sub-arcs $5-1$ to $5-18$. In these (tests $13-15$ ), arc 5 was counted with relative weight $=2.0$ since the acceleration $\sigma^{\prime}$ s of independent arcs 5 A and 5 B were each near the average $\sigma$ for the independent arcs in Table 10. Tests 13-15 ignore the influence of arc 5-11 because of its excessive residual ( $4 \sigma$ ) when it is included in the tests for the harmonics to and through third order. This series of tests, relatively unprejudiced by arbitrary weighting and giving complete coverage to arc 5, was chosen as the basis for drawing final conclusions in this gravity experiment.

We have already computed the effects on the gravity synthesis of previously determined acceleration model biases in the various arcs (compare tests $1-3$ with tests $4-6$ ). There is another way to present this result which reveals the likely model errors in both stages of this experiment directly in terms of the harmonic coefficients and the test standard errors. We merely repeat precisely the same gravity synthesis on the simulated parallel acceleration data in Table 11 as we did in the preceding tests on the actual data in Table 10. However, the simulated 24 -hour trajectories summarized in Table 11 contained no effects from resonant earth gravity of higher than third order. To gage the likely effects of higher order gravity on the second and third order gravity syntheses, or rather to obtain a wider range of likely second stage model error in the actual data, we have also simulated satellite drift over four geoids taken from recent studies. In these simulations, no full gravity trajectories were calculated. Instead, simple drift accelerations from Equation 2 were computed for the satellite arcs in Table 11 according to the gravity constants of the four geoids. These accelerations were then combined in the same way as the actual data (reversing the solution of Equation 11 with selected coefficients ignored) to yield gravity fields whose bias is readily apparent.

The results of these parallel gravity syntheses on the simulated data are found in Table 12, tests 16-27. Two of the geoids chosen for these simulations come from single comprehensive studies of satellite perturbations. The geoid of Kaula (1964) (see Table B1) is derived from camera observations of five to ten medium altitude, medium and high inclination satellites. The geoid of Guier (1965) was determined from comprehensive, worldwide Doppler radar tracking of five medium altitude satellites (see Table B1). These two geoids are considered representative of the best available geodetic results from independent satellite tracking to date. The geoid labeled Wagner-Kaula combined 1964/65 (used in the ITEM trajectory computations) used $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ harmonic coefficients derived at an earlier stage of the present study. The $H_{31}$ coefficients are from an "average" geoid due to W. M. Kaula (private communication) recommended as a "first guess" in geodetic studies when
there is no better evidence of different values. The geoid labeled Wagner-Guier (1965) uses the $\mathrm{H}_{22}$, $\mathrm{H}_{33}$ values derived earlier in this study, $\mathrm{H}_{31}$ values near those finally chosen as best representing 24 -hour satellite drift to date, and $\mathrm{H}_{42}, \mathrm{H}_{44}$ values from the geoid of Guier (1965). This latter geoid was thought to be the most accurate known at the time this study was initiated (spring, 1965). These tests illustrate the kind of convergence to true harmonic values, which should be evident and apparently is, in the actual data reductions for $\mathrm{C}_{22}$ and $\mathrm{S}_{22}$, when higher order gravity is introduced without constraint into the tests. Test 18 shows clearly that sun, moon, and first stage model bias has almost negligibly small effect on these harmonic reductions. This is also evident from the low values of the acceleration biases in Table 10.

The harmonic biases (theoretical-measured values) shown by these simulated data tests are calculated in Table 13. Except for the effects on $\mathrm{H}_{3} 1$, they are all reasonably consistent. This illustrates, more than anything else, the agreement in the geoids themselves. As more harmonics are permitted, the sharp reduction in the $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ biases shows that, observation errors excluded, these harmonics should be essentially determined from the arcs in the data test through the third order. In Table 14 the model biases in the harmonics determined from the simulated arcs in Table 13 are added to the harmonics derived from the actual data (Table 12, tests 13-15) to arrive at a reasonable range of bias free harmonics at all stages in the reductions.

It appears noteworthy that all $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ harmonics except $\mathrm{C}_{33}$ are relatively unchanged through the reductions when they are corrected at each stage through third order by the average biases. This result, which could not be anticipated in advance (for two of the four geoid simulations), gives added assurance in both the overall quality of the basic data and the $\mathrm{H}_{22}, \mathrm{H}_{33}$ gravity field implicit in that drift data after only a third order reduction. As a further test of the conjecture that $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ are essentially determined by only a third order reduction of this wide coverage 24-hour data, we allow a fourth order harmonic into the reductions (tests 28 and 29, Table 12). Appendix C makes it fairly clear that $\mathrm{H}_{44}$ is the next strongest harmonic after $\mathrm{H}_{31}$ in its influence on the 24 -hour satellites in this study.

It is encouraging to find that these tests show essentially the same $H_{22}$ and $H_{33}$ results as previously in the all third order reductions (i.e., tests $13-15$ ). It is interesting that $\mathrm{J}_{44} \simeq-0.02 \times 10^{-6}$ in these tests since this checks reasonably well with recent results for this harmonic determined from lower altitude data (see Table B1). The fact that both $J_{44}$ and $\lambda_{44}$ determined from tests 28 and 29 are almost identical with and without the inclusion of $\mathrm{H}_{31}$ seems to be coincidental. The $C_{31}$ from test 29 seems unrealistically high from the results of this and other recent geodetic investigations from satellite motions (see Table B1). As a further confirmation of this, in test 30 we repeat the weighted data test 3 through third order with the inclusion of $H_{44}$. While all the other harmonics are virtually unchanged from test 29 , the $\mathrm{S}_{44}$ harmonic has increased significantly. With only 9 or 10 well determined 24 -hour satellite accelerations at this point in the analysis, it should not be surprising that only four resonant gravity coefficients appear well determined.

The best estimate of these coefficients is presented at the bottom of Table 14, including a very preliminary estimate of $\mathrm{H}_{31}$, as far as can be judged from the many gravity tests in this section. In general, mean values and standard errors were chosen together to encompass as widely as
possible the results of all the actual data tests in Table 12 and bias adjusted reductions in Table 14. The single reduction most heavily relied upon in this judgement was the average bias-added gravity synthesis (in Table 14) from unweighted basic acceleration data with dense arc 5 coverage. It appears that a more balanced estimate of $S_{33}$ should be $S_{33}=-(0.16 \pm 0.01) \times 10^{-6}$, since not one test showed $\mathrm{S}_{33}$ to be greater than $-0.17 \times 10^{-6}$. In this one instance, balance over all the tests was ignored in favor of the bias-adjusted results on the unweighted dense arc 5 coverage data reductions. The reason for this exception is to be found in the remarkably consistent bias-adjusted $\mathrm{S}_{33}$ ' S in Table 14. This seems to arise in large part from the consistency in $\mathrm{S}_{33}$ between $\mathrm{H}_{22}$, $\mathrm{H}_{33}$, and $\mathrm{H}_{22}, \mathrm{H}_{33}, \mathrm{H}_{31}$ reductions when arc 5 is densely covered (compare tests 14 and 15, dense coverage tests with tests 11 and 12 in Table 12). To guard against error in this somewhat special judgement of the test results, we have allowed for a $0.015 \times 10^{-6}$ standard error in $S_{33}$ which is probably higher than is strictly seen in this experiment.

The judgement of the $\mathrm{H}_{31}$ harmonics in Table 14 also calls for some comment. Again, a balanced view over all the tests was the criterion of choice. The results of the random sampling tests 7A, 7B, 7C (in Table 12) were strongly relied upon in the $H_{31}$ estimates. This was in spite of the fact that arc 5 was sparsely covered in these tests and there were only nine samples to test for six coefficients (independent arc 5 ' not being used, for example). In particular, test 7A shows that viewed probablistically, the basic data allows low overall acceleration residuals and $\mathrm{H}_{22}, \mathrm{H}_{33}$ values reasonably consistent with other 'best values' tests and, in addition, $\mathrm{H}_{31}$ values considerably divergent from the "best value" results. The standard errors of the $\mathrm{H}_{31}$ harmonics were, in fact, taken from test 7A. The best estimates of $\mathrm{H}_{31}$ were considered to be the last two bias adjusted values in Table 14. Together with the estimated standard errors, a fair measure of the range of $\mathrm{H}_{31}$ seen in these tests is covered. The standard error in $\mathrm{J}_{31}$ was not estimated from $\mathrm{S}\left(\mathrm{C}_{31}\right)$ and $\mathrm{S}\left(\mathrm{S}_{31}\right)$ through Equation 15, assuming uncorrelated coefficients (Reference 5, Appendix E), since $C_{31}$ and $S_{31}$ appear to be significantly correlated in these tests. Clearly $J_{31}$, seen in the experiment, is bounded between about $-10^{-6}$ and $-3 \times 10^{-6}$. The estimated $J_{31}$ deviation in Table 14 reflects this result.

Finally, in Figure 9, we display the measured (unadjusted) 24 -hour satellite accelerations (in solid) in this study (from table 10) and match them against the accelerations (in dots) from the geoid (at the bottom of Table 14) which the measured accelerations have determined in large part. The standard errors in the matching geoid accelerations (about $0.03 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day $^{2}$ for Syncom 2 arcs) were calculated from somewhat smaller $H_{n m}$ deviations than those finally chosen in Table 14. Actual geoid accelerations may differ from the best values calculated through Equation 2 with the best estimated coefficients in Table 14 by about $0.08 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$ for equatorial 24 -hour satellites. This number reflects only the long term acceleration uncertainty remaining in our knowledge of the effect of the earth's field on the distant synchronous satellite. But we know also from this study that, beyond two months of drift, long term sun and moon gravity effects may continue to be responsible for as much as $\pm 0.03 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$ in the drift acceleration of the 24 -hour satellite.

The solid curve in Figure 9 represents the best estimated geoid accelerations on the geostationary satellite around the equator, as determined from this study.

Table 13
Model Biases in Gravity Harmonics Reduced from Simulated 24 -Hour Satellite Drifts with Four Recent Geoids*.
(all bias values in units of $10^{-6}$ )

${ }^{\text {Data }}$ from least squares reductions of simulations over ares $1,2,4,5-1$ to $5-18,6,7,8$ and 9 : see Table 12 and text.
GEOIDS USED
(1) Wagner-Kaula combined 1964/65 (3 $3^{\text {rd }}$ order) plus sun, moon, and $1^{\text {tt }}$ stage experiment bias effects: see test 16 , Table 12
(2) Kaula 1964: see test 13, Table 12.
(3) Guier 1965: see test 22, Table 12
(4) Wagner-Guier 1965: see test 25 , Table 12
(5) Wagner-Kaula comb. (1964/65) plus sun and moon
(6) Kaula 1964
(7) Guier 1965
(8) Wagner-Guier 1965
(9) Wagner-Kaula comb. (1964/65) plus sun and moon
(10) Kaula 1964
(11) Guier 1965
(12) Wagner-Guier 1965

Table 14
Gravity Synthesis, from Actual 24-Hour Satellite Data and Likely Model Bias Errors. (all harmonic values in units of $10^{-6}$ except as noted)

| Reductions for |  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :--- | :--- | :--- | :--- |
| Harmonics $H_{n m}\left(\mathrm{C}_{n \mathrm{~m}}, \mathrm{~S}_{\mathrm{nm}}\right)$ | $\mathrm{C}_{22}$ | $\pm \mathrm{S}\left(\mathrm{C}_{22}\right)$ | $\mathrm{S}_{22}$ | $\pm \mathrm{S}\left(\mathrm{S}_{22}\right)$ | $\mathrm{C}_{33}$ | $\pm \mathrm{S}\left(\mathrm{C}_{33}\right)$ | $\mathrm{S}_{33}$ | $\pm \mathrm{S}\left(\mathrm{S}_{33}\right)$ | $\mathrm{C}_{31}$ | $\pm S\left(\mathrm{C}_{31}\right)$ |


| $\mathrm{H}_{22}$ only | (1) | -1.525 |  | 1.005 |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ | (2) | -1.565 |  | . 930 |  | -. 047 |  | -. 164 |  |  |  |  |  |
| $\mathrm{H}_{22}, \mathrm{H}_{33}$ and $\mathrm{H}_{31}$ | (3) | ${ }^{-1.554}$ | . 005 | . 935 | . 006 | -. 029 | . 006 | -. 164 | . 004 | 1.72 | .4 | . 02 | . 2 |
| $\mathrm{H}_{22}$ only ${ }^{\text {a }}$ |  | (-1.569 |  | . 940 |  |  |  |  |  |  |  |  |  |
| $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ |  | $\{-1.577$ |  | . 925 |  | -. 065 |  | -. 165 |  |  |  |  |  |
| $\mathrm{H}_{22}, \mathrm{H}_{33} \text { and } \mathrm{H}_{31}$ | (4) | $\left\{\begin{array}{c}-1.558 \\ (-1.568)\end{array}\right.$ |  | . 932 |  | -. 033 |  | -. 164 |  | 1.88 |  | . 09 |  |
| (Average) |  | (-1.568) | (.012)* | (.932) | (.013)* | (-.049) | $(.018) *$ | (-.165) | (.003)* | (1.88) | (.6)* | (.09) | (.3)* |
| $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ | (5) | -1.584 |  | . 932 |  | -. 070 |  | -. 166 |  |  |  |  |  |
| $\mathrm{H}_{22}, \mathrm{H}_{33}$ and $\mathrm{H}_{31}$ | (6) | -1.560 |  | . 932 |  | -. 030 |  | -. 165 |  | 2.17 |  | . 08 |  |
| $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ | (7) | -1.581 |  | . 927 |  | -. 071 |  | -. 164 |  |  |  |  |  |
| $\mathrm{H}_{22}, \mathrm{H}_{33}$ and $\mathrm{H}_{31}$ | (8) | -1.558 |  | . 949 |  | -. 030 |  | -. 168 |  | 2.60 |  | -. 33 |  |
| $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ | (9) | -1.583 |  | . 923 |  | -. 079 |  | -. 166 |  |  |  |  |  |
| $\mathrm{H}_{22}, \mathrm{H}_{33}$ and $\mathrm{H}_{31}$ | (10) | -1.556 |  | . 924 |  | -. 036 |  | -. 161 |  | 1.35 |  | . 30 |  |
| $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ | (11) | -1.559 |  | . 928 |  | -. 039 |  | -. 163 |  |  |  |  |  |
| $\mathrm{H}_{22}, \mathrm{H}_{33}$ and $\mathrm{H}_{31}$ | (12) | -1.557 |  | . 923 |  | -. 036 |  | -. 161 |  | 1.39 |  | . 30 |  |
|  | (13) | -1.56 | . 02 | . 93 | . 02 | -. 045 | . 030 | -. 165 | . 015 | 1.4 | 1.2 | . 3 | . 6 |
|  |  | $\mathrm{J}_{22}$ -1.816 | $\mathrm{S}\left(\mathrm{J}_{22}\right)$ .02 | $\lambda_{22}$ -15.40 | $\mathrm{S}\left(\lambda_{22}\right)$ $.32^{\circ}$ | $\mathrm{J}_{33}$ -.171 | $\mathrm{S}\left(\mathrm{J}_{33}\right)$ .017 | $\lambda_{33}{ }^{\circ}$ 24.92 | $\mathrm{S}\left(\lambda_{33}\right)$ $3.3^{\circ}$ | J 31 -1.4 | ${ }_{+0.2}^{\mathrm{S}\left(\mathrm{J}_{31}\right)^{* *}}$ | ${ }^{\lambda_{31}}{ }^{\text {-167 }}{ }^{\circ}$ | ${ }_{\text {S }}^{S\left(\lambda_{31}\right)}$ |
|  | (14) | -1.816 | . 02 | -15.40 | $.32^{\circ}$ | -. 171 | . 017 | $24.92^{\circ}$ | $3.3^{\circ}$ | -1.4 | +0.2 | $-167.9^{\circ}$ | $25.8{ }^{\circ}$ |

*Maximum RMS bias dev. from average biases in Table 13.
**Estimated from a range of $\left(\mathrm{C}_{31}^{2}+\mathrm{S}_{31}^{2}\right)^{1 / 2}$ values in Tables 12,13 and 14 .

## GEOIDS USED

$$
\begin{array}{ll}
\overline{C_{22}}(\text { best })=2.42 \pm .03 \times 10^{-6} & \overline{C_{33}}(\text { best })=0.322 \pm .215 \times 10^{-6} \\
\overline{\mathrm{~S}_{22}}(\text { best })=-1.44 \pm .03 \times 10^{-6} & \left.\overline{S_{33}} \text { (best }\right)=1.183 \pm .108 \times 10^{-6}
\end{array}
$$

## 3. EAST-WEST EQUILIBRIUM LONGITUDES AND MAXIMUM EAST-WEST STATION KEEPING REQUIREMENTS FOR THE GEOSTATIONARY SATELLITE

This study confirms the results of many recent satellite-geoid reductions (see Table B1 and Reference 1) that the dominance of the $\mathrm{J}_{22}$ harmonic in the earth's field establishes only four narrow longitude zones at which a geostationary satellite may be placed and kept for long periods of time without the necessity of east-west station keeping. As a study of Figure 9 shows, initially geostationary satellites placed near geoid acceleration zeros over the Indian Ocean and the eastern Pacific will be forever trapped in the field under the influence only of the small longitude dependent components of the earth's gravity potential. Similar satellites placed near the east-west acceleration zeros over the Atlantic and western Pacific will tend to drift away from their initial positions, but initially only very slowly, at least due to perturbations of the earth's field (see also Reference 1 and 8). The present study of the drift of three 24 -hour satellites over a period of two years indicates these east-west equilibrium points in the earth's field at synchronous altitudes are at:

$$
\begin{align*}
& \lambda_{1}=76.7 \pm 0.8^{\circ} \text { (dynamically stable equilibrium longitude) } \\
& \lambda_{2}=161.8 \pm 0.7^{\circ} \text { (statically stable equilibrium longitude) } \\
& \lambda_{3}=-108.1 \pm 1.0^{\circ} \text { (dynamically stable equilibrium longitude) }  \tag{16}\\
& \lambda_{4}=-12.2 \pm 0.7^{\circ} \text { (statically stable equilibrium longitude). }
\end{align*}
$$

These longitudes and their standard errors have been calculated through Equation 2 from the geoid harmonics and their likely deviations at the bottom of Table 14. The derived geoid in Table 14 , reflecting the drift record of the 24 -hour satellites, is as free as we can make it from model bias. This includes higher order earth effects, as well as sun and moon gravity effects, and the error sources in the data reduction methods. Therefore these equilibrium longitude zones are believed to be absolute measures. Barring a much stronger higher order earth longitude dependent field than appears likely now, the true equilibrium zones should fall within the limits above.

To calculate the maximum east-west station keeping requirements for geostationary satellites, we compare Equation B4 with Equation C1. The longitude perturbing force per unit mass on the circular orbit equatorial 24 -hour satellite given in terms of the drift acceleration (in radians/sid. day ${ }^{2}$ ) it produces, is

$$
\begin{equation*}
\mathrm{F}_{\lambda}=\frac{\left(-\ddot{\lambda} \mu_{\mathrm{e}} / \mathrm{a}_{\mathrm{s}}{ }^{2}\right)}{12 \pi^{2}}=-0.00621 \ddot{\lambda} \mathrm{ft} / \mathrm{sec}^{2}, \tag{17}
\end{equation*}
$$

for $a_{s}=6.611$ earth radii, with $\ddot{\lambda}$ in units of rad/sid. day ${ }^{2}$. The maximum longitude acceleration which might be experienced by such a satellite is (from Figure 9)

$$
\begin{align*}
\ddot{\lambda}_{\max } & =-(3.18 \pm 0.08) \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \text { day }{ }^{2} \\
& =-(1.83 \pm 0.05) \times 10^{-3} \text { degrees } / \text { day }^{2}, \tag{18}
\end{align*}
$$

at $\lambda=118^{\circ} \mathrm{E}$ (over Indonesia). Using the conservative upper bound of Equation 18 in Equation 17, we calculate the maximum east-west station keeping requirements for the geostationary as (following Reference 3, p. 31):

$$
\begin{align*}
\Delta \mathrm{V}_{\mathrm{T} \cdot \max } & =\mathrm{F}_{\lambda, \max } \times \Delta \mathrm{T}(1 \mathrm{yr}) \\
& =0.00621 \times 3.26 \times 86,400(\mathrm{sec} / \text { day }) \times 365(\text { days } / \mathrm{yr}) \times 10^{-5} \\
& =6.38 \mathrm{ft} / \mathrm{sec}-\mathrm{yr}) . \tag{19}
\end{align*}
$$

Other longitudes where near maximum east-west station keeping requirements on geostationary satellites would exist, occur roughly halfway between $\lambda_{2}$ and $\lambda_{3}, \lambda_{3}$ and $\lambda_{4}$, and $\lambda_{4}$ and $\lambda_{1}$ in Equations 16.

## DISCUSSION

Perhaps the best check on the validity of the geodetic results of this study (in Figure 9 and Table 14) is a recent private communication from R. R. Allan to the author (June 1965). Mr. Allan of the Royal Aircraft Establishment in England reports the following gravity harmonics as seen by Syncom 2 drift in arcs 1-5:

$$
\begin{aligned}
& \mathrm{J}_{22}=-1.80 \times 10^{-6} \\
& \lambda_{22}=-15.0^{\circ} \\
& \mathrm{J}_{33}=0.178 \times 10^{-6} \\
& \lambda_{33}=24.7^{\circ} \\
& \mathrm{J}_{44}=-0.017 \times 10^{-6} \\
& \lambda_{44}=37.9^{\circ}
\end{aligned}
$$

Mr. Allan performed his independent reductions on essentially the same orbit data in arcs 1-5 as found in this report (in Appendix A). Allan apparently used a somewhat different analysis than here. He removed sun and moon perturbations by a semi-analytic technique and solved for all the harmonics directly from the drift in the arcs by an iterative method. His best results for $\mathrm{H}_{22}$ and $H_{33}$ are well within the standard deviations reported for the final values in Table 14 except for $\lambda_{22}$ which does not differ significantly.

The present study could come to no firm conclusion on $H_{44}$ except to estimate broadly that $0.01<\left|J_{44} \times 10^{6}\right|<0.03$ (see Conclusions). The best external check on the results, (at least for $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ ) is found in the recent geodetic reductions in Table B1. For the two dominant sectorials
through third order, the results of this study agree most closely individually and as a set with the Doppler - satellite geoids of Guier (1965) and Anderle (1965), and the camera - satellite geoid of Kaula (1964). The only surprising result of this study is the phase angle $\lambda_{31}$ of the lowest order "mixed" or tesseral harmonic, which was sensed to be almost $180^{\circ}$ from the consensus of recent observations (see Table B1). However, as pointed out in Section 2 of this report, the 24 -hour data thus far appears to allow for a considerably greater range of $\lambda_{31}{ }^{\prime} s$ than settled on in the final geoid values of Table 14. The random sampling tests on the 24 -hour data (Table 12) showed that $\lambda_{31}$ could be less than $90^{\circ}$ (west) without seriously affecting the $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ harmonics determined from the estimated "best" data. Since the inclusion of $\mathrm{H}_{44}$ into the gravity synthesis of that data did not seem to improve the situation with regard to $\mathrm{H}_{31}$, we must wait for more acceleration data, particularly from the equatorial 24 -hour satellites, to clear up the mystery of the apparent discrepancy in this harmonic. It should be emphasized again, however, that the reported values of $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ as a set, even without the inclusion of $H_{31}$, give nearly as good a reproduction of the "best measured" accelerations as the complete third order geoid reported. In fact, it is believed that the reported $\mathrm{H}_{22}$ and $H_{33}$ values are individually absolute within their stated deviations, and are not expected to vary significantly from these ranges when the 24 -hour data is complete enough to reveal fourth order gravity as well.

The greater part of the effort in this study has been to obtain long term accelerations with as small likely deviations as possible. As Appendix C shows, not a great deal could be expected of obtaining meaningful results on third or fourth order harmonics from the limited record, unless accuracies of the order of $0.1 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$.day ${ }^{2}$ in the longitude drift accelerations were obtained. For the most part this goal was met and exceeded. But for much of the Syncom 2 and 3 record it was not an easy task because the orbit determinations were often of much reduced quality as can be judged from a comparison of the standard test errors among arcs 1, 2, and 8. In the main, use of a single well determined longitude location for each reported orbit gave sufficiently precise accelerations for the Syncom arcs, provided the drift exceeded about two months and more than eight individual orbits were available for the arc. But in many Syncom free drift arcs the longitude data alone was not good enough to obtain sufficiently precise accelerations in spite of the arc length. This was especially true for arcs $3,5 B$ and 7 (see Table 7). For arc 3, the length was short of two months, but even with the use of two longitudes in each GSFC reported orbit to help determine the drift velocity changes, unacceptable precision of about $1 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$ was all that could be attained. In arc 5B (and the latter sub-arcs of arc 5) some of the single longitude estimations in November 1964 and January 1965 proved so poor they were discarded altogether as drift velocity indicators and replaced for this purpose by successive longitude estimations for those orbits.

In Reference 6 the use of such drift rate data from a single urbit was discussed. In theory it should provide a semi-independent determination of the accelerations in the Syncom arcs since an orbit is specified by both independent position and velocity information. However, except for isolated regions in the 24 -hour record, the velocity (or semimajor axis) measurements were not as "smooth" for the purposes of the analysis as the position (longitude) measurements. In theory this is understandable since the differences of two errored equator positions and crossing times over a small time span gives the velocity from one independent orbit determination. The satellite velocity
determined from two independent 'best" position measurements over a longer base time (between orbits) should have superior precision. For example, in arc 5', we could only measure the best acceleration from such single orbit velocity data with a precision of $0.11 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$ day ${ }^{2}$ compared to $0.4 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$ in arc 5 covering essentially the same Syncom 2 orbits (see also the results of the semimajor axis analysis in Reference 3). It appears on inspection that the mean semimajor axes for these orbits, reported by GSFC but not presented here, change with sufficient smoothness over arc 5 to permit independent acceleration determinations to be made from them within $0.1 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$ day ${ }^{2}$ accuracy. Alternately, one could also use for this purpose the average drift rate from three or four successive Equator crossings generated from a single set of vector elements in the presence of sun and moon and zonal gravity perturbations (such as in Appendix A). The "mean" semimajor axis implies such velocity smoothing over the perturbations for a number of days past the orbit epoch.

Both of these velocity-from-single-orbit approaches will be tried in the future on a more thoroughgoing basis than here to augment and refine the rather inadequate 24 -hour acceleration record presently available for the Syncom satellites. It is hoped that such an augmentation of the past record, together with new equatorial data from Syncom 3, Early Bird, and near future 24-hour satellites will produce a clear picture of $\mathrm{H}_{31}$ as well as show conclusive evidence of fourth order resonant earth gravity. The earth harmonics of higher than fourth order appear to be beyond reasonable discrimination from 24 -hour altitudes until far more data is available than is foreseen for the next few years.

With a few exceptions, noted in the tables of Section 1, the orbit data used in the arc analyses are believed to be substantially free from all but gravity perturbations (see also the Discussion in Reference 3). For example, in Appendix E, are calculated likely magnitudes for residual atmospheric drag and solar radiation pressure accelerations on the 24 -hour Syncom satellites. These are found to be entirely negligible compared to resonant earth gravity accelerations.

## CONCLUSIONS

From this comprehensive investigation of the long term gravity drift of three 24 -hour "synchronous" satellites over a period of two years (1963-1965), the following conclusions are drawn:

1. Virtually all of the east-west geographic acceleration of these satellites can be accounted for by the second and third order sectorial harmonics of the earth's gravitational field which resonate with them.
2. With due adjustment for small effects of sun and moon gravity and the neglect of likely higher order resonant earth gravity, these dominant sectorial harmonics are estimated to be

$$
J_{22}=-(1.816 \pm 0.020) \times 10^{-6}
$$

which corresponds to a difference in major and minor axes of $69.4 \pm 0.9$ meters in the earth's elliptical Equator, and

$$
\lambda_{22}=-(15.4 \pm 0.3)^{\circ},
$$

which is the longitude location of the major axis of the elliptical Equator, and

$$
\begin{aligned}
& J_{33}=-(0.171 \pm 0.017) \times 10^{-6} \\
& \lambda_{33}=24.9 \pm 3.3^{\circ} .
\end{aligned}
$$

3. The sectorial harmonics above, within their ranges, are believed to be absolute or true measures of those individual components of the earth's field.
4. A third pair of third order resonant earth harmonics was just evident but poorly discriminated from the limited acceleration record. The data shows tentatively

$$
\begin{aligned}
& J_{31}=-\left(1.4{ }_{-0.2}^{+1.0}\right) \times 10^{-6}, \\
& \lambda_{31}=-(168 \pm 26)^{\circ} .
\end{aligned}
$$

5. Tests of the complete 24 -hour satellite record for the effects of individual fourth order earth resonant gravity harmonics were inconclusive but gave some evidence that

$$
0.01<\left|J_{44} \times 10^{6}\right|<0.03
$$

6. The geoid resulting from this close study of operating 24 -hour satellites implies that an equatorial synchronous satellite can be in uncontrolled long term east-west equilibrium at only the following four longitude locations:

$$
\begin{aligned}
& \lambda_{1}=76.7 \pm 0.8^{\circ} \text { (dynamically stable east-west equilibrium) } \\
& \lambda_{2}=161.8 \pm 0.7^{\circ} \text { (statically stable east-west equilibrium) } \\
& \lambda_{3}=-108.1 \pm 1.0^{\circ} \text { (dynamically stable east-west equilibrium) } \\
& \lambda_{4}=-12.2 \pm 0.7^{\circ} \text { (statically stable east-west equilibrium). }
\end{aligned}
$$

7. The maximum long term longitude acceleration due to earth gravity which can be experienced by the nearly geostationary satellite, according to the 24 -hour acceleration record thus far, is conservatively $\ddot{\lambda}=-1.88 \times 10^{-3}$ degrees $/$ day $^{2}$, at about $118^{\circ}$ east of Greenwich.
8. To correct continuously for this east-west acceleration would require a velocity increment of

$$
\Delta V_{\max }=6.38 \mathrm{ft} /(\mathrm{sec}-\mathrm{yr}) .
$$

9. Further study of the drift of both present and near future 24 -hour satellites should be rewarded in a few years by the first unambiguous picture of the earth's resonant longitude gravity field through fourth order.

## ACKNOWLEDGMENTS

The author warmly appreciates the encouragement and fruitful exchange of ideas and information with R. R. Allan (Royal Aircraft Establishment, England), José Osório Pereira (Astronomical Observatory of the University of Monte Da Virgem, Portugal), Peter Musen (Goddard Space Flight Center, Greenbelt, Maryland), David Mott (University of New Mexico) and William Kaula (University of California at Los Angeles). The author is no less thankful for the more immediate assistance of Ken Squires, Fred Schaeffer, Fred Whitlock and Ed Monasterski (from the Theoretical Division of the Goddard Space Flight Center) in the laborious and so often frustrating numerical calculations that directly supported the results of this effort. Special thanks are also due Bob Green of the Comsat Corporation for supplying the excellent Andover, Maine Early Bird tracking data from which was calculated the finest observed 24 -hour satellite acceleration in this study.
(Manuscript received September 14, 1965)

## REFERENCES

1. Wagner, C. A., 'The Drift of a 24 -Hour Equatorial Satellite Due to an Earth Gravity Field Through Fourth Order," NASA Technical Note D-2103, February 1964.
2. Wagner, C. A., "The Drift of an Inclined Orbit 24 -Hour Satellite in an Earth Gravity Field Through Fourth Order," NASA Technical Note, in press, 1965 (G-666).
3. Wagner, C. A., "Determination of the Ellipticity of the Earth's Equator From Observations on the Drift of the Syncom 2 Satellite," NASA Technical Note D-2759, May 1965.
4. Wagner, C. A., "On the Probable Influence of Higher Order Earth Gravity on the Determination of Equatorial Ellipticity From the Drift of Syncom 2 Over Brazil," NASA Technical Note, in press, 1965 (G-664).
5. Wagner, C. A., "The Equatorial Ellipticity of the Earth From Syncom 2 Drift Over the Central Pacific," NASA Technical Note, in press, 1965 (G-665).
6. Wagner, C. A., "The Equatorial Ellipticity of the Earth as Seen From Syncom 2 Drift Over the Western Pacific," NASA Technical Note, in press, 1965 (G-663).
7. Williams, D. D., "Dynamic Analysis and Design of the Synchronous Communication Satellite," Hughes Aircraft Corp. Report TM 649, May 1960.
8. Blitzer, L., Boughton, E. M., Kang, G., and Page, R. M., "Effect of Ellipticity of the Equator on 24-Hour Nearly Circular Satellite Orbits," J. Geophys. Res. 67(1):329-335, January 1962.
9. Frick, R. H., and Garber, T. B., "Perturbations of a Synchronous Satellite Due to the Triaxiality of the Earth," Rand Corp. Memo. RM-2996-NASA, January 1962.
10. Musen, P., and Bailie, I. E., "On the Motion of a 24-Hour Satellite," J. Geophys. Res. 67(3):11231132, March 1962.
11. Wagner, C. A., "The Gravitational Potential and Force Field of the Earth Through Fourth Order," NASA Technical Note, in press 1965 (G-667).
12. Allan, R. R., "Even Tesseral Harmonics in the Geopotential Derived from Syncom 2," paper presented at the Second International Symposium on the Use of Artificial Satellites for Geodesy, Athens, April 1965.
13. Bowker, A. H., and Lieberman, G. J., "Engineering Statistics," Englewood Cliffs, N. J.: Prentice-Hall, 1959.
14. Smart, W. M., "Combination of Observations," New York: Cambridge University Press, 1958.
15. Guier, W. H., "Recent Progress in Satellite Geodesy," Johns Hopkins University Applied Physics Lab. Document TG-659, February 1959.
16. Kaula, W. M., "Theory of Satellite Geodesy," New York: Blaisdell, in press, 1965.

## Appendix A

## Basic Orbit Data Used in This Report

Vector and mean elements for Syncom 2 ( $\operatorname{arcs} 1,2,3,4,5$, and 8) and Syncom 3 (arcs 6 and 7) were used as basic data in this report and are reported in Tables $\mathbf{A 1}$ and A2 below. These elements were calculated by the GSFC Tracking and Data Systems Directorate from radar and minitrack observations on the satellite made over a period of about three days per orbit following the listed epochs. The orbit determination program used for the calculation of these elements employed a gravity-earth model with the following constants (see Appendix B):

$$
\begin{aligned}
\mu_{\text {earth }} & =3.98627 \times 10^{5} \mathrm{~km}^{3} / \mathrm{sec}^{2} \\
\mathrm{R}_{0} & =6378.388 \mathrm{~km} / \text { mean equatorial earth radius } \\
\mathrm{J}_{20} & =1082.21 \times 10^{-6} \\
\mathrm{~J}_{30} & =-2.29 \times 10^{-6} \\
\mathrm{~J}_{40} & =-2.10 \times 10^{-6} \\
\mu_{\text {sun }} & =332.490 \mu_{\text {earth }} \\
\mu_{\text {moon }} & =0.01229491 \mu_{\text {earth }}
\end{aligned}
$$

The Equator crossing data in Section 1 of this report was derived (mainly) by generating numerically a short trajectory from the vector elements of Tables A1 and A2 utilizing the gravityearth constants above. The absence of longitude earth gravity in the orbit determinations for the Syncom satellites generally limited the time for which data could be applied to each orbit to less than a week.

The numerical trajectory generator employed in deriving the Equator crossing data and in running the long simulated longitude gravity trajectories in Section 1, is called "ITEM" at Goddard Space Flight Center. Details of this generator can be found in GSFC Document X-640-63-71, "Interplanetary Trajectory Encke Method (ITEM) Program Manual", May 1963.

The basic subsatellite position data for arc 9 (Early Bird) was supplied by Robert H. Green of the Comsat Corporation (see Table 9). It was determined by simultaneous range-azimuth-elevation observations on Early Bird from the A. T. and T. tracking facility at Andover, Maine.

Table AI
Inertial Position and Velocity Coordinates for Syncom 2 and Syncom 3 as Reported by GSFC.*

| Tracking Epoch ( Y -mo-day-hr-min UT) | $\underset{\left(10^{4}\right.}{\mathrm{X} \mathrm{~km})}$ | $\stackrel{Y}{\left(10^{4} \mathrm{~km}\right)}$ | $\stackrel{\mathrm{Z}}{\left(10^{4} \mathrm{~km}\right)}$ | $\underset{(\mathrm{km} / \mathrm{sec})}{\dot{\mathrm{x}}}$ | $\begin{gathered} \dot{\mathbf{Y}} \\ (\mathrm{km} / \mathrm{sec}) \end{gathered}$ | $\begin{gathered} \dot{\boldsymbol{z}} \\ (\mathrm{km} / \mathrm{sec}) \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 763-8-18-1-30.0 | 1.8517253 | -3.6656408 | -0.95346425 | 2.5197630 | 0.87570544 | 1.5296730 |
| 63-8-22-6-12.14 | 3.8192813 | 0.19653916 | 1.7732339 | -0.64139671 | 2.8111679 | 1.0698498 |
| 63-8-26-17.0 | -3.9190365 | 0.71434463 | -1.3838910 | -0.030659063 | -2.7659652 | -1.3414586 |
| 63-8-31.0 | 1.2517473 | -3.8173186 | -1.2805028 | 2.7082374 | 0.42130417 | 1.3937010 |
| 63-9-3-13-23.0 | -2.6683093 | 3.2430659 | 0.37724660 | -2.0920997 | -1.5282086 | -1.6556391 |
| 63-9-5.0 | 1.5690433 | -3.7540418 | -1.1058861 | 2.6179496 | 0.66119258 | 1.4710835 |
| 63-9-9.0 | 1.8101775 | -3.6842867 | -0.96169979 | 2.5333080 | 0.84735057 | 1.5232783 |
| 63-9-12-2.0 | 3.4100141 | -2.4566117 | 0.33394973 | 1.4036497 | 2.1745471 | 1.6606540 |
| 63-9-17-2.0 | 3.5605032 | -2.1949356 | 0.52762629 | 1.1881299 | 2.3199229 | 1.6320863 |
| Arc 1 63-9-20-2.0 | 3.6381029 | -2.0310580 | 0.64240062 | 1.0551433 | 2.3993342 | 1.6084890 |
| 63-9-27-2.0 | 3.7828113 | -1.6282253 | 0.90343513 | 0.73324914 | 2.5577582 | 1.5413518 |
| 63-10-1-2.0 | 3.8398851 | -1.3915057 | 1.0449696 | 0.54792106 | 2.6322531 | 1.4928190 |
| 63-10-8-2 | 3.8997221 | -0.96844341 | 1.2765410 | 0.22281006 | 2.7325580 | 1.3929604 |
| 63-10-14-2.0 | 3.9114166 | -0.59213158 | 1.4569605 | -0.059266662 | 2.7895298 | 1.2936179 |
| 63-10-22-2.0 | 3.8655564 | -0.089237507 | 1.6821569 | -0.43011588 | 2.8231687 | 1.1399466 |
| 63-10-30.0 | 3.7935711 | -1.5763511 | 0.94831895 | 0.69205550 | 2.5804974 | 1.5229227 |
| 63-11-6.0 | 3.8724273 | -1.1762815 | 1.1818209 | 0.38097748 | 2.6938419 | 1.4335987 |
| 63-11-12-5.0 | 1.1446010 | 3.4554430 | 2.1294725 | -2.7244144 | 1.2840822 | -0.61733366 |
| 63-11-18-13.0 | -3.7038702 | -0.58064695 | -1.9332132 | 0.90759960 | -2.7961257 | -0.89794623 |
| /63-11-28-1.0 | 3.4781437 | 1.1376687 | 2.0944723 | -1.2933144 | 2.7059520 | 0.67780023 |
| 63-12-4.0 | 3.7151749 | 0.53514140 | 1.9201218 | -0.87415171 | 2.8043813 | 0.90960269 |
| 63-12-10.0 | 3.5731662 | 0.92615699 | 2.0379487 | -1.1461780 | 2.7500585 | 0.76033287 |
| 63-12-16-17.0 | 1.0227024 | -3.8691081 | -1.3292554 | 2.7458357 | 0.25931246 | 1.3585847 |
| 63-12-23-19.0 | 3.0804533 | -2.8787714 | 0.055186822 | 1.7485346 | 1.9038213 | 1.6650503 |
| 64-1-6-17.0 | 2.2436953 | -3.5177574 | -0.60704656 | 2.3262615 | 1.2062956 | 1.6089899 |
| 64-1-9-6.0 | -3.0926247 | 2.8665793 | -0.080953346 | -1.7324029 | -1.9159081 | -1.6667974 |
| Arc 2) 64-1-15-18.0 | 3.3299438 | -2.5658898 | 0.32807382 | 1.4797710 | 2.1321920 | 1.6489369 |
| 64-1-20-21.0 | 3.7030072 | 0.56000451 | 1.9369923 | -0.88759725 | 2.8080493 | 0.88485319 |
| 64-1-29-20.0 | 3.8204264 | 0.093177361 | 1.7811029 | -0.55721471 | 2.8372181 | 1.0470801 |
| 64-2-5-16.0 | 2.8698222 | -3.0876991 | -0.10282631 | 1.9209748 | 1.7308532 | 1.6640514 |
| 64-2-10-19.0 | 3.8650196 | -0.19769561 | 1.6738813 | -0.34675337 | 2.8365031 | 1.1356011 |
| 64-2-17-17.0 | 3.7268026 | -1.7536007 | 0.90292256 | 0.82362190 | 2.5381818 | 1.5281483 |
| 64-2-25-19.0 | 3.6533259 | 0.69908427 | 1.9851452 | -0.98120516 | 2.7965023 | 0.82109998 |
| 64-3-4-23.0 | 0.18113658 | 3.8087583 | 1.8019013 | -2.8347353 | 0.61701878 | -1.0178939 |
| 64-3-10-13.0 | 2.0644779 | -3.6124046 | -0.68632246 | 2.4052556 | 1.0737921 | 1.5862758 |
| /64-3-18-3.0 | -3.6249117 | 2.0233288 | -0.74902788 | -1.0365741 | -2.4353607 | -1.5693664 |
| 64-3-24-13.0 | 2.3283566 | -3.4999056 | $\sim 0.50715805$ | 2.2721038 | 1.2746548 | 1.6158241 |
| Arc 3 64-4-1-22.0 | 0.76296064 | 3.6490419 | 1.9857054 | -2.7806607 | 1.0205339 | -0.81923611 |
| 64-4-7-15.0 | 3.4461234 | -2.4088100 | 0.50885287 | 1.3277141 | 2.2415368 | 1.6144312 |
| 64-4-13-19.0 | 3.3872421 | 1.3385265 | 2.1522760 | -1.4069923 | 2.6725021 | 0.54080057 |
| ¢64-4-25-2.0 | -2.5603443 | 3.3390780 | 0.29994066 | -2.1318214 | -1.4887792 | -1.6428786 |
| 64-4-28-15.0 | 3.2628949 | -2.6693162 | 0.32101150 | 1.5457882 | 2.0856009 | 1.6366277 |
| 64-5-5-16.0 | 3.7257132 | -1.7632128 | 0.94048572 | 0.82412264 | 2.5449717 | 1.5038900 |
| 64-5-12-16.0 | 3.7475309 | -1.6867283 | 0.99296812 | 0.76273664 | 2.5725322 | 1.4891387 |
| Arc 4 64-5-19-14.0 | 2.7849331 | -3.1784434 | -0.10945226 | 1.9747957 | 1.6719452 | 1.6505955 |
| 64-5-25-15.0 | 3.4307022 | -2.4126224 | 0.53392194 | 1.3319460 | 2.2493385 | 1.6072617 |
| 64-6-2-21.0 | 1.7406646 | 3.1597779 | 2.1946670 | -2.5288299 | 1.6899106 | -0.43439828 |
| 64-6-9-21.0 | 1.6500071 | 3.2194181 | 2.1774956 | -2.5601432 | 1.6304633 | -0.47782909 |
| 64-6-16-15.0 | 3.5655116 | -2.1536100 | 0.72108034 | 1.1248629 | 2.3871212 | 1.5668312 |
| 64-6-23-15.0 | 3.6075927 | -2.0587512 | 0.78733380 | 1.0494437 | 2.4320607 | 1.5495642 |
| 64-7-4-2.0 | -3.2093967 | 2.7165420 | -0.32063943 | -1.5841089 | -2.0660850 | -1.6378223 |
| 64-7-7-3.0 | -3.6954496 | 1.7950587 | -0.94983861 | -0.85046154 | -2.5458831 | -1.5023067 |
| 64-7-13-17.0 | 3.7354217 | 0.41 .539202 | 1.9259618 | -0.76727529 | 2.8431082 | 0.87209834 |
| 64-7-21-21.0 | 0.68011374 | 3.6978808 | 1.9157389 | -2.7798310 | -0.96706741 | -0.88586596 |
| 64-7-27-16.0 | 3.8417868 | -0.11824881 | 1.7510381 | -0.39126192 | 2.8593634 | 1.0493355 |
| 64-8-3-17.0 | 3.4664817 | 1.1480683 | 2.1210692 | -1.2649191 | 2.7375684 | 0.58214636 |
| 64-8-11-1.0 | -3.3701691 | 2.4805650 | -0.51870386 | -1.3894809 | -2.2250586 | -1.6062097 |
| 64-8-17-19.0 | 1.6959231 | 3.2024028 | 2.1645995 | -2.5377008 | 1.6647695 | -0.48003497 |
| 64-8-25-10.0 | 1.5330007 | -3.8258893 | -0.91436182 | 2.5905352 | 0.67495556 | 1.5073075 |
| 64-9-1-10.0 | 1.7622665 | -3.7590544 | -0.76800569 | 2.5108491 | 0.85857627 | 1.5478150 |
| 64-9-9-14.0 | 3.8729406 | -0.56806750 | 1.5864037 | -0.064200095 | 2.8378419 | 1.1715632 |
| 64-9-15-12.0 | 3.4973065 | -2.2642849 | 0.69291203 | 1.2114165 | 2.3483095 | 1.5648130 |
| 64-9-22-10.0 | 2.3957332 | -3.4619283 | -0.32713684 | 2.2142000 | 1.3764322 | 1.6238887 |
| Arc 5<64-9-29-6.0 | -1.2057564 | -3.4918173 | -2.0387967 | 2.6768512 | -1.3377786 | 0.70292893 |
| 64-10-6-5.0 | -1.9027758 | -3.0642587 | -2.1871975 | 2.4522966 | -1.8112966 | 0.39976832 |
| 64-10-13.0 | -3.7909750 | 1.3890374 | -1.2159474 | -0.54220619 | -2.6929405 | -1.3843106 |
| 64-10-20-16.0 | 2.3418688 | 2.7061755 | 2.2406431 | -2.2408225 | 2.0926322 | -0.18928830 |
| 64-10-26-16.0 | 2.1737218 | 2.8588045 | 2.2200917 | -2.3278640 | 1.9849102 | -0.280.44993 |
| 64-11-2-5.0 | -1.0575392 | -3.5703165 | -1.9843669 | 2.7067256 | -1.2347065 | 0.77366371 |
| 64-11-11-2.0 | -3.2269498 | -1.5931919 | -2.1991397 | 1.5433675 | -2.6360297 | -0.35716605 |
| 64-11-17-6.0 | 0.44169064 | -3.9493268 | -1.4227193 | 2.7947443 | -0.14738721 | 1.2691720 |
| 65-1-10-6.0 | 1.8733482 | -3.7329199 | -0.61394837 | 2.4551255 | 0.97133874 | 1.5716495 |
| 65-1-13-16.0 | -0.024904471 | 3.9169200 | 1.5734233 | -2.8064159 | 0.44943751 | -1.1686893 |
| 65-1-20-12.0 | 3.2361444 | 1.5906403 | 2.1989944 | -1.5260383 | 2.6448765 | 0.33028481 |
| 65-1-27-4-5.0 | 0.42866892 | -3.9652039 | -1.3910958 | 2.7914381 | -0.14522612 | 1.2721610 |
| 65-2-2-13-30.0** | 1.8326846 | 3.1377111 | 2.1684747 | -2.4504221 | 1.7730384 | -0.49949115 |
| 65-2-16-4-5.0 | 0.94696655 | -3.9573969 | -1.1249388 | 2.7174166 | 0.24662704 | 1.4128023 |

*From the Data Systems and Tracking Directorate, computed from range and range rate and Minitrack Data by Robert Chaptick, Gerald Repass, and Carleton Carver, from an orbit determination program due to Dr. Joseph Siry (see text in Appendix A for carth-gravity constants used in this program.) The inertial orthogonal system $x, y, z$ has the $x$ axis pointing cowards the vernal equinox of epoch and the $z$ axis pointing towards the North Celestial Pole.

Table A2
Inertial Position and Velocity Coordinates or Mean Elements for Syncom 2 and Syncom 3 as Reported by GSFC*.

| Tracking Epoch (yrmo-day-hr-min UT) | $\underset{\left(10^{4} \mathrm{~km}\right)}{\mathrm{x}}$ | $\begin{gathered} y \\ \left(10^{4} \mathrm{~km}\right) \end{gathered}$ | $\begin{gathered} z \\ \left(10^{4} \mathrm{~km}\right) \end{gathered}$ | $\begin{gathered} \mathrm{k} \\ (\mathrm{~km} / \mathrm{sec}) \end{gathered}$ | $\begin{gathered} y \\ (\mathrm{~km} / \mathrm{sec}) \end{gathered}$ | $\begin{gathered} \dot{z} \\ (\mathrm{~km} / \mathrm{sec}) \end{gathered}$ | Semimajor Axis (earth radii) | Eccentricity | Inclination (degrees) | Mean Anomaly (degrees) | Argument of Perigee (degrees) | Right Ascension of the Ascending Node (degrees) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| '65-1-14-23.5 | 1.2675646 | -4.0179928 | -0.051959215 | 2.9359926 | 0.92060244 | -0.047842619 |  |  |  |  |  |  |
| 65-1-30-13-10.0 | -3.4792181 | 2.3828935 | -0.0058904292 | -1.7368621 | -2.5371972 | -0.0083528082 |  |  |  |  |  |  |
| 65-2-2-6.0 | 3.2099039 | 2.7368950 | 0.0073851190 | -1.9943285 | 2.3390669 | 0.0012901180 |  |  |  |  |  |  |
| 65-2-9-11.0 | -2.2119906 | 3.5906580 | -0.0049486959 | -2.6173528 | -1.6131485 | -0.0081798989 |  |  |  |  |  |  |
| Arc 65-2-16-12.0 | -3.3564297 | 2.5533931 | -0.0042520595 | -1.8613403 | -2.4471176 | -0.0086970709 |  |  |  |  |  |  |
| 7 7 65-2-23.0 | 3.5938828 | -2.2084394 | 0.0096558595 | 1.6094139 | 2.6188512 | 0.0026358856 |  |  |  |  |  |  |
| \|65-3-2.0 | 3.8120235 | -1.8069874 | 0.015075846 | 1.3165671 | 2.7773230 | -0.00057799012 |  |  |  |  |  |  |
| 65-3-9.0 | 3.9810232 | -1.3940918 | 0.0045517934 | 1.0162757 | 2.9011202 | 0.0024674641 |  |  |  |  |  |  |
| 65-3-16.0 | 4.1104734 | -0.95605039 | 0.026135089 | 0.69589465 | 2.9925915 | -0.00060597968 |  |  |  |  |  |  |
| 65-2-25.0 | -2.6260533 | -2.4306852 | -2.2283266 | 2.0528730 | -2.2888731 | 0.081154582 |  |  |  |  |  |  |
| 65-3-3.0 | -2.3397165 | -2.7254249 | -2.2051595 | 2.2300224 | -2.1051269 | 0.23910311 |  |  |  |  |  |  |
| 65-3-6.0 |  |  |  |  |  |  | 6.6112847 | . 00076 | 31.939 | 308.359 | 333.048 | 309.370 |
| 65-3-13.0 |  |  |  |  |  |  | 6.6113429 | . 00069 | 31.912 | 326.845 | 321.447 | 309.095 |
| Arc 65-3-29.0 | -0.85799562 | -3.6815475 | -1.8631946 | 2.7341821 | -1.0890101 | 0.89574242 |  |  |  |  |  |  |
| Arc 65-4-5.0 |  |  |  |  |  |  | 6.6111445 | . 00069 | 31.830 | 329.612 | 340.634 | 308.896 |
| 65-4-12.0 |  |  |  |  |  |  | 6.6111335 | . 00073 | 31.848 | 328.269 | 348.585 | 308.947 |
| 65-4-19.0 |  |  |  |  |  |  | 6.6110585 | . 00065 | 31.867 | 353.146 | 330.838 | 308.586 |
| 65-4-26.0 |  |  |  |  |  |  | 6.6110267 | . 00063 | 31.787 | 353.717 | 337.005 | 308.580 |
| 65-5-3.0 |  |  |  |  |  |  | 6.6110035 | . 00061 | 31.807 | 359.897 | 337.757 | 308.417 |
| 65-5-10.0 |  |  |  |  |  |  | 6.6109407 | . 00059 | 31.731 | 0.012 | 344.515 | 308.335 |

*From the Data Systems and Tracking Directocate (see note in Table AI). The mean elements, good for more than one orbit (smoorhing out the periodic sun, moon, and zonal gravity effects) are calculated from the vector elements of position and velocity according to a cheory due to D . Brouwer.
"

## Appendix B

## Earth Gravity Potential and Force Field Used in This Report: Comparison with Recent Investigations

The gravity potential used as the basis for the data reduction in this study is the exterior potential of the earth derived (in Reference B1 at the end of this appendix) for geocentric spherical coordinates referenced to the earth's spin axis and its center of mass (Equation B1). The infinite series of spherical harmonics is truncated after $J_{44}$ because the great height of the synchronous satellite makes it very insensitive to higher orders of earth gravity. The zonal gravity and other earth constants used in this study (and illustrated in Equation B1) are from Reference B2 and represent a somewhat outdated set used by GSFC in the orbit determination program for the Syncom satellites (see Appendix A). They are (with the corresponding mean equatorial radius):

$$
\begin{aligned}
\mathrm{R}_{0} & =6378.388 \mathrm{~km} / \text { mean equator earth radius (Reference B3) }, \\
\mu_{\text {earth }} & =3.98627 \times 10^{5} \mathrm{~km}^{3} / \mathrm{sec}^{2} \text { (Reference B3), } \\
\mathrm{J}_{20} & =1082.21 \times 10^{-6}, \\
\mathrm{~J}_{30} & =2.29 \times 10^{-6}, \\
\mathrm{~J}_{40} & =-2.10 \times 10^{-6} .
\end{aligned}
$$

Though these values individually are not the most accurate known to date (1965), they were chosen for the trajectory generations in this study to insure consistency with the published orbits. The longitude gravity reductions themselves are not significantly affected by the probable errors in these zonal gravity and other principal earth constants. The most accurate "zonal geoid" is still probably that of Kozai (1962) in Reference B4, with the following constants:

$$
\begin{aligned}
\mathrm{R}_{0} & =6378.2 \mathrm{~km}, \\
\mu_{\text {earth }} & =3.98603 \times 10^{5} \mathrm{~km}^{3} / \mathrm{sec}^{2}, \\
\mathrm{~J}_{20} & =1082.48 \times 10^{-6}, \\
\mathrm{~J}_{30} & =-2.56 \times 10^{-6}, \\
\mathrm{~J}_{40} & =-1.84 \times 10^{-6},
\end{aligned}
$$

plus higher order terms.
A good review of zonal gravity investigations from satellite observations through 1965 is to be found in Reference B5.


$$
\mathrm{V}_{\mathrm{E}}=\frac{\mu_{\mathrm{E}}}{\mathrm{r}}\left[1-\frac{\mathbf{J}_{20} \mathrm{R}_{0}{ }^{2}}{2 \mathbf{r}^{2}}\left(3 \sin ^{2} \phi-1\right)-3 \mathrm{~J}_{22} \frac{\mathbf{R}_{0}{ }^{2}}{\mathrm{r}^{2}} \cos ^{2} \phi \cos 2\left(\lambda-\lambda_{22}\right)\right.
$$



$$
-\frac{J_{30} R_{0}^{3}}{2 r^{3}}\left(5 \sin ^{3} \phi-3 \sin \phi\right)-\frac{J_{31} R_{0}^{3}}{2 r^{3}} \cos \phi\left(15 \sin ^{2} \phi-3\right) \cos \left(\lambda-\lambda_{31}\right)
$$


$-15 \mathrm{~J}_{32} \frac{\mathrm{R}_{0}{ }^{3}}{\mathrm{r}^{3}} \cos ^{2} \phi \sin \phi \cos 2\left(\lambda-\lambda_{32}\right)-15 \mathrm{~J}_{33} \frac{\mathrm{R}_{0}{ }^{3}}{\mathrm{r}^{3}} \cos { }^{3} \phi \cos 3\left(\lambda-\lambda_{33}\right)$


The earth-gravity field (per unit test mass), whose potential is Equation B1, is given as the gradient of B 1 , or

$$
\begin{equation*}
\overline{\mathbf{F}}=\hat{\mathbf{r}} \mathbf{F}_{\mathbf{r}}+\hat{\lambda} \mathbf{F}_{\lambda}+\hat{\phi} \mathbf{F}_{\phi}=\mathbb{V}_{E}=\hat{\mathbf{r}} \frac{\partial \mathbf{V}_{E}}{\partial r}+\frac{\hat{\lambda}}{\mathrm{r} \cos \phi} \frac{\partial V_{E}}{\partial \lambda}+\frac{\hat{\phi}}{\mathbf{r}} \frac{\partial \mathrm{V}_{E}}{\partial \phi} ; \tag{B2}
\end{equation*}
$$

or
$F_{r}=\frac{\mu_{\mathrm{E}}}{\mathrm{r}^{2}}\left\{-1+\left(\mathrm{R}_{0} / \mathrm{r}\right)^{2}\left[3 / 2 \mathrm{~J}_{20}\left(3 \sin ^{2} \phi-1\right)+9 \mathrm{~J}_{22} \cos ^{2} \phi \cos 2\left(\lambda-\lambda_{22}\right)\right.\right.$
$+2\left(R_{0} / r\right) J_{30}\left(5 \sin ^{2} \phi-3\right)(\sin \phi)+6\left(R_{0} / r\right) J_{31}\left(5 \sin ^{2} \phi-1\right) \cos \phi \cos \left(\lambda-\lambda_{31}\right)$
$+60\left(R_{0} / r\right) J_{32} \cos ^{2} \phi \sin \phi \cos 2\left(\lambda-\lambda_{32}\right) \cdot+60\left(R_{0} / r\right) J_{33} \cos ^{3} \phi \cos 3\left(\lambda-\lambda_{33}\right)$
$+5 / 8\left(\mathrm{R}_{0} / \mathrm{r}\right)^{2} \mathrm{~J}_{40}\left(35 \sin ^{4} \phi-30 \sin ^{2} \phi+3\right)$
$+25 / 2\left(R_{0} / r\right)^{2} J_{41}\left(7 \sin ^{2} \phi-3\right) \cos \phi \sin \phi \cos \left(\lambda-\lambda_{41}\right)$
$+75 / 2\left(R_{0} / \mathrm{r}\right)^{2} \mathrm{~J}_{42}\left(7 \sin ^{2} \phi-1\right) \cos ^{2} \phi \cos 2\left(\lambda-\lambda_{42}\right)$
$\left.\left.+525\left(R_{0} / r\right)^{2} \mathrm{~J}_{43} \cos ^{3} \phi \sin \phi \cos 3\left(\lambda-\lambda_{43}\right)+525\left(R_{0} / r\right)^{2} J_{44} \cos ^{4} \phi \cos 4\left(\lambda-\lambda_{44}\right)\right]\right\}$,
$F_{\lambda}=\frac{\mu_{E}}{r^{2}}\left(R_{0} / r\right)^{2}\left\{6 J_{22} \cos \phi \sin 2\left(\lambda-\lambda_{22}\right)+3 / 2\left(R_{0} / r\right) J_{31}\left(5 \sin ^{2} \phi-1\right) \sin \left(\lambda-\lambda_{31}\right)\right.$
$+30\left(R_{0} / \mathrm{r}\right) \mathrm{J}_{32} \cos \phi \sin \phi \sin 2\left(\lambda-\lambda_{32}\right)+45\left(\mathrm{R}_{0} / \mathrm{r}\right) \mathrm{J}_{33} \cos ^{2} \phi \sin 3\left(\lambda-\lambda_{33}\right)$
$+5 / 2\left(R_{0} / r\right)^{2} J_{41}\left(7 \sin ^{2} \phi-3\right) \sin \phi \sin \left(\lambda-\lambda_{41}\right)+15\left(R_{0} / r\right)^{2} J_{42}\left(7 \sin ^{2} \phi-1\right) \cos \phi \sin 2\left(\lambda-\lambda_{42}\right)$
$+315\left(R_{0} / r\right)^{2} J_{43} \cos ^{2} \phi \sin \phi \sin 3\left(\lambda-\lambda_{43}\right)$
$\left.+420\left(R_{0} / r\right)^{2} J_{44} \cos ^{3} \phi \sin 4\left(\lambda-\lambda_{44}\right)\right\}$,

$$
\begin{align*}
F_{\phi}=\frac{\mu_{E}}{r^{2}}\left(R_{0} / r\right)^{2}\{ & -3 J_{20} \sin \phi \cos \phi+6 J_{22} \cos \phi \sin \phi \cos 2\left(\lambda-\lambda_{22}\right) \\
& -3 / 2\left(R_{0} / r\right) J_{30}\left(5 \sin ^{2} \phi-1\right) \cos \phi+3 / 2\left(R_{0} / r\right) J_{31}\left(15 \sin ^{2} \phi-11\right) \sin \phi \cos \left(\lambda-\lambda_{31}\right) \\
& +15\left(R_{0} / r\right) J_{32}\left(3 \sin ^{2} \phi-1\right) \cos \phi \cos 2\left(\lambda-\lambda_{32}\right) \\
& +45\left(R_{0} / r\right) J_{33} \cos ^{2} \phi \sin \phi \cos 3\left(\lambda-\lambda_{33}\right)-5 / 2\left(R_{0} / r\right)^{2} J_{40}\left(7 \sin ^{2} \phi-3\right) \sin \phi \cos \phi \\
& +5 / 2\left(R_{0} / r\right)^{2} J_{41}\left(28 \sin ^{4} \phi-27 \sin ^{2} \phi+3\right) \cos \left(\lambda-\lambda_{41}\right) \\
& +30\left(R_{0} / r\right)^{2} J_{42}\left(7 \sin ^{2} \phi-4\right) \cos \phi \sin \phi \cos 2\left(\lambda-\lambda_{42}\right) \\
& +105\left(R_{0} / r\right)^{2} J_{43}\left(4 \sin ^{2} \phi-1\right) \cos { }^{2} \phi \cos 3\left(\lambda-\lambda_{43}\right) \\
& \left.+420\left(R_{0} / r\right)^{2} J_{44} \cos ^{3} \phi \sin \phi \cos 4\left(\lambda-\lambda_{44}\right)\right\} \tag{B5}
\end{align*}
$$

The actual sea-level surface of the earth is to be conceptualized through Equation B1 as a sphere of radius 6378 km , around which are superimposed the sum of the separate spherical harmonic deviations illustrated. To these static gravity deviations, of course, must be added a centrifugal earth-rotation potential at the earth's surface, to get the true sea-level surface.

Table B1 gives longitude coefficients for this earth-gravity field form as reported by geodesists from 1942 to 1965.

The longitude gravity represented in Equation B1 represents only that gravity determined in the final longitude geoid of this report (see Table 14). (The longitude gravity drift simulations used slightly different values than these as reported in Sections 1 and 2.) The coefficients of this longitude geoid are

$$
\begin{aligned}
& \mathrm{J}_{22}=-1.816 \times 10^{-6}, \\
& \lambda_{22}=-15.4^{\circ}, \\
& \mathrm{J}_{31}=-1.4 \times 10^{-6}, \\
& \lambda_{31}=-168^{\circ}, \\
& \mathrm{J}_{33}=-0.171 \times 10^{-6}, \\
& \lambda_{33}=+24.9^{\circ} .
\end{aligned}
$$

Other longitude coefficients from recent studies are reported in Table B1 for purposes of comparison. The maximum geoid heights and depressions illustrated in Equation B1 are proportional to the $\mathrm{J}_{\mathrm{nm}}$ amplitudes of the corresponding gravity harmonics.

Longitude Coefficients in the Earth's Gravity Potential $\left\{v_{E}=\frac{\mu}{r} \sum_{n=2}^{\infty} \sum_{m=0}^{\infty}\left[1-\left(R_{0} / r\right)^{n} P_{n}^{m}(\sin \phi) J_{n m} \cos m\left(\lambda-\lambda_{n m}\right)\right]\right\}^{1}$, as Reported 1942-1965 ${ }^{2}$.

| Longitude Geoid Reference | $J_{22}$ | $\lambda_{22}$ | $\mathrm{J}_{31}$ | $\lambda_{31}$ | $\mathrm{J}_{3}{ }^{2}$ | $\lambda_{32}$ | $\mathrm{J}_{33}$ | $\lambda_{33}$ | $\mathrm{J}_{4} 1$ | $\lambda_{41}$ | $\mathrm{J}_{42}$ | $\lambda_{42}$ | $\mathrm{J}_{43}$ | $\lambda_{43}$ | $\mathrm{J}_{44}$ | $\lambda_{44}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| (1) Wagner (1965a) ${ }^{3}$ | $-1.81 \times 10^{-6}$ | $-15.4^{\circ}$ | $-1.4 \times 10^{-6}$ | $-168.0^{\circ}$ |  |  | $-.171 \times 10^{-6}$ | $24.9^{\circ}$ |  |  |  |  |  |  |  |  |
| (2) Izsak (1965) ${ }^{4}$ | -1.57 | $-15.2^{\circ}$ | -1.72 | $-1.0{ }^{\circ}$ | $-.300 \times 10^{-6}$ | $-32.5^{\circ}$ | -. 199 | $32.4{ }^{\circ}$ | $-.519 \times 10^{-6}$ | $-134.0^{\circ}$ | $-.138 \times 10^{-6}$ | $35.5{ }^{\circ}$ | $-.0413 \times 10^{-6}$ | $-2.7^{\circ}$ | $-.0076 \times 10^{-6}$ | $27.2^{\circ}$ |
| (3) Guier (1965) ${ }^{3}$ | -1.72 | $-13.4{ }^{\circ}$ | -2.01 | $6.7^{\circ}$ | -. 477 | $-14.6{ }^{\circ}$ | -. 165 | $18.7^{\circ}$ | -. 679 | $-142.0^{\circ}$ | -. 193 | $23.4{ }^{\circ}$ | -. 0506 | $0.2^{\circ}$ | -. 0060 | $34.5{ }^{\circ}$ |
| (4) Anderle (1965) ${ }^{3}$ | -1.86 | $-16.0^{\circ}$ | -2.32 | $7.0^{\circ}$ | -. 455 | $-21.3^{\circ}$ | -. 242 | $23.5{ }^{\circ}$ | -. 723 | $-130.0^{\circ}$ | -. 158 | $33.7^{\circ}$ | -. 0626 | -4.3 ${ }^{\circ}$ | -. 0116 | $34.6{ }^{\circ}$ |
| (5) Kaula (1964) ${ }^{4}$ | -1.77 | $-18.2^{\circ}$ | -2.12 | $-5.4{ }^{\circ}$ | -. 379 | $10.5^{\circ}$ | -. 105 | $23.1{ }^{\circ}$ | -. 263 | $-239.0^{\circ}$ | -. 117 | $42.3^{\circ}$ | -. 0473 | $15.0^{\circ}$ | -. 0104 | $14.5{ }^{\circ}$ |
| (1966) | -1.82 | -14.9 | -2.27 | 5.5 | -. 36 | -21.9 | -. 194 | 22.6 |  |  | -. 184 | 31.2 |  |  | -. 0081 | 33.8 |
| (6) Uotila (1964) ${ }^{5}$ | -1.52 | $-36.5{ }^{\circ}$ | -0.685 | -81.0 $0^{\circ}$ | -. 409 | $-5.2^{\circ}$ | -. 398 | $19.5{ }^{\circ}$ | -. 238 | $-127.0^{\circ}$ | -. 211 | $14.6{ }^{\circ}$ | -. 082 | $-9.3{ }^{\text {a }}$ | -. 0142 | $-2.6{ }^{\circ}$ |
| (7) Kozal (1962) ${ }^{4}$ | -1.2 | $-26.4^{\circ}$ | -1.9 | $4.6{ }^{\circ}$ | -. 14 | $-16.8{ }^{\circ}$ | -. 10 | $42.6{ }^{\circ}$ | -. 52 | -122.5 | -. 062 | $65.2^{\text {c }}$ | -. 035 | $0.5^{\circ}$ | -. 031 | $14.9{ }^{\circ}$ |
| (8) Zhongolovitch (1961) ${ }^{5}$ | -5.95 | $-7.7^{\circ}$ | -2.21 | $-25.7^{\circ}$ | -. 628 | $-26.4^{\circ}$ | -. 54 | $13.0{ }^{\circ}$ | $-.78$ | -149.1 | -. 080 | $45.0^{\circ}$ | -. 051 | $-3.8^{\circ}$ | -. 0224 | $15.9{ }^{\circ}$ |
| (9) Jeffreys (1942) ${ }^{\text {3 }}$ | -4.1 | 0.0 | -2.1 | 0.0 | -. 66 | 0.0 | -. 24 | $33.3^{\circ}$ |  |  |  |  |  |  |  |  |

$I_{r}$ is the radial distance of the field point to the center of mass of the earth, $\mu$ the earth's Gaussian gravity constant $=3.9860 \times 10^{20} \mathrm{~cm}^{3} / \mathrm{sec}^{2}, \mathrm{R}_{0}$ the mean equatorial radius of the earth $=6378.2 \mathrm{~km}$. $\phi$ is the geocentric latitude of the field

$T_{n m t}=\frac{(-1)^{!}(2 \mathrm{n}-2 t)!}{2^{n} \mathrm{t}!(\mathrm{n}-\mathrm{t})!(\mathrm{n}-\mathrm{m}-2 \mathrm{t})!}$ (See Kaula, 1965 [Reference 16]). The longitude coefficients are those for which $\mathrm{m} \neq 0$.
 to the observed data.
'Satellite - radar geoid.
${ }^{4}$ Satellite-camera geoid.
${ }^{\text {Surface-gravimetric geoid. }}$

References and Notes to Longitude Geoid Data of Table B1

| Longitude Geoid | Reference | Notes |
| :---: | :---: | :---: |
| (1) | This report | Considers the drift of three 24 -hour satellites, two geostationary and one of medium inclination, with fair global longitude coverage. |
| (2) | Data due to I. Izsak quoted in: "Y. Kozai," Summary of Numerical Results Derived from Satellite Observations," paper presented at the 2nd International Symposium on the Use of Artificial Satellites for Geodesy, Athens, Greece, April, 1965. | Uses $10-15$ medium altitude, medium and high inclination satellites; data reduced from Baker-Nunn camera observations. |
| (3) | In: "Recent Progress in Satellite Geodesy," Johns Hopkins University, APL Report TG-659, Feb. 1965 | Uses data from about 5 "Transit" medium altitude, medium and high inclination satellites, reduced from Doppler observations. |
| (4) | Data due to Anderle quoted in the Y. Kozai paper above [for longitude geoid (2)] | Uses data from both "Transit" and "Anna" satellites; reduction uses cross track perturbations as well as along track. |
| (5) | Private communication to the author (July 1964) from W. M. Kaula | Uses about 10 medium altitude, medium and high inclination satellites. |
| (6) | Data due to Uotila quoted in a private communication to the author (July 1964) from W. M. Kaula | Believed by Kaula to be the most comprehensivecoverage gravimetric geoid to date (1964). |
| (7) | Private communication to the author (Oct. 1962) from Y. Kozai. | Uses about 5 medium altitude, medium inclination satellites. |
| (8) | Data due to Zhongolovitch quoted in: Y. Kozai, "Tesseral Harmonics of the Gravitational Potential of the Earth," Astronom. J. 66(7): Sept. 1961. | Recent Russian gravimetric geoid for comparison purposes. |
| (9) | Data due to Jeffreys quoted in the paper above [for longitude geoid (8)]. | Older gravimetric geoid for comparison purposes. |

## REFERENCES

B1. Wagner, C. A., "The Gravitational Potential of a Triaxial Earth," GSFC Document X-623-62-206, October 1962.

B2. Kozai, Y., "The Earth's Gravitational Potential Derived from Motions of Three Satellites," Astron. J. 66(1):8-10, February 1961.

B3. O'Keefe, J. A., Eckels, A., and Squires, R. K., "The Gravitational Field of the Earth," Astron. J. 64(7):245-253, September 1959.

B4. Kozai, Y., 'Numerical Results from Orbits," Smithsonian Inst. Astrophys. Obs. Spec. Rept. 101, July 31, 1962, pp. 1-21.

B5. Kozai, Y., "Summary of Numerical Results Derived from Satellite Observations," paper presented at the Second International Symposium on the Use of Artificial Satellites for Geodesy Athens, Greece, April 1965.

## Appendix C

## Preliminary Maximum Longitude Accelerations on 24-Hour Satellites Due to the Resonant Gravity Harmonics of the Earth through Fourth Order

The source for this appendix is Reference 15. The studies in Reference 1 and especially References 3 through 6 have already established the dominance of second order earth gravity (represented by the elliptical Equator) on the perturbed drift of the 24 -hour satellite. In order to establish a reasonable basis for higher order gravity tests, we calculate here maximum 24 -hour satellite accelerations due to resonant gravity harmonics through fourth order as well. The source for the higher order harmonics is Guier (1965), Reference 15 (see also Table B1), considered to be representative of the best of recent high order satellite-geoid determinations.

To cover the three satellites in this study, we calculate maximum drift accelerations for equatorial ( $\mathrm{i}_{\mathrm{s}}=0$, Syncom 3 and Early Bird) as well as moderately inclined orbit satellites ( $i_{s}=32.5^{\circ}$, Syncom 2).

## Equatorial Satellites (Syncom 3 and Early Bird)

From Equation 57B in Reference 1 (or Equation 66 in Reference 2) the long term longitude drift acceleration of the 24 -hour equatorial satellite is given through fourth order earth gravity as

$$
\begin{align*}
\ddot{\lambda}=-12 \pi^{2}\left(R_{0} / a_{s}\right)^{2}\{ & 6 J_{22} \sin 2\left(\lambda-\lambda_{22}\right)-\frac{3}{2}\left(R_{0} / a_{s}\right) J_{31} \sin \left(\lambda-\lambda_{31}\right)+45\left(R_{0} / a_{s}\right) J_{33} \sin 3\left(\lambda-\lambda_{33}\right) \\
& \left.-15\left(R_{0} / a_{s}\right)^{2} J_{42} \sin 2\left(\lambda-\lambda_{42}\right)+420\left(R_{0} / a_{s}\right)^{2} J_{44} \sin 4\left(\lambda-\lambda_{44}\right)\right\} \mathrm{rad} / \mathrm{sid} . \mathrm{day}^{2}, \tag{C1}
\end{align*}
$$

where the following set of harmonic constants are used:

$$
\begin{aligned}
& \mathrm{J}_{22}=-1.8 \times 10^{-6} \quad(\text { Wagner, preliminary estimate from this study }), \\
& \mathrm{J}_{31}=-2.0 \times 10^{-6} \quad(\text { Guier }(1965) ; \text { Reference } 15), \\
& \mathrm{J}_{33}=-0.17 \times 10^{-6} \quad(\text { Guier }(1965)), \\
& \mathrm{J}_{42}=-0.19 \times 10^{-6} \quad(\text { Guier }(1965)), \\
& \left.\mathrm{J}_{44}=-0.006 \times 10^{-6} \quad \text { Guier }(1965)\right) .
\end{aligned}
$$

Calculating only the maximum values of the terms in Equation C 1 for $\left(\mathrm{R}_{0} / a_{\mathrm{s}}\right)^{2}=0.0229$ with $a_{s}=6.61$ earth radii gives

$$
\begin{align*}
& \left|\ddot{\lambda}_{22}(\max )\right|_{\text {equator }}=2.92 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \text { day }^{2},  \tag{C2}\\
& \left|\ddot{\lambda}_{31}(\max )\right|_{\text {equator }}=0.12 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \text { day }^{2},  \tag{C3}\\
& \left|\ddot{\lambda}_{33}(\max )\right|_{\text {equator }}=0.31 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \text { day }^{2},  \tag{C4}\\
& \left|\ddot{\lambda}_{42}(\max )\right|_{\text {equator }}=0.018 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \text { day }^{2},  \tag{C5}\\
& \left|\ddot{\lambda}_{44}(\max )\right|_{\text {equator }}=0.016 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \mathrm{day}^{2} \tag{C6}
\end{align*}
$$

## $32.5^{\circ}$ Inclined Satellite (Syncom 2)

For this satellite the maximum values in Equations C2 and C6 should be multiplied by the inclination factors $\mathrm{F}_{\mathrm{nm}}(\mathrm{i})$ (Reference 2, Equation 67):

$$
\begin{array}{ll}
n, m & F_{\mathrm{nm}}(i), \\
2,2 & \frac{1}{4}(\cos i+1)^{2}, \\
3,1 & \frac{1}{2}(\cos i+1)-\frac{5}{8} \sin ^{2} i(1+3 \cos i), \\
3,3 & \frac{1}{8}(\cos i+1)^{3}, \\
4,2 & \frac{1}{4}(\cos i+1)^{2}-\frac{7}{4} \sin ^{2} i \cos i(\cos i+1), \\
4,4 & \frac{1}{16}(\cos i+1)^{4} . \tag{C11}
\end{array}
$$

Evaluating Equations C7-C11 at $i=32.5^{\circ}$ gives

$$
\begin{align*}
& \mathrm{F}_{22}(\mathrm{i})=0.850,  \tag{C12}\\
& \mathrm{~F}_{31}(\mathrm{i})=0.284,  \tag{C13}\\
& \mathrm{~F}_{33}(\mathrm{i})=0.785,  \tag{C14}\\
& \mathrm{~F}_{42}(\mathrm{i})=0.063,  \tag{C15}\\
& \mathrm{~F}_{44}(\mathrm{i})=0.723 . \tag{C16}
\end{align*}
$$

Equations C12-C16 as multiplying factors of Equations C2-C6 give the maximum accelerations on the Syncom 2 satellite due to the gravity harmonics through fourth order as

$$
\begin{align*}
& \mid\left.\ddot{\lambda}_{22}(\text { max })\right|_{\text {Syncom 2 }}=2.48 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \text { day }^{2},  \tag{C17}\\
& \left|\ddot{\lambda}_{31}(\mathrm{max})\right|_{\text {Syncom } 2}=0.03 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \text { day }^{2},  \tag{C18}\\
& \left|\ddot{\lambda}_{33}(\max )\right|_{\text {Syncom } 2}=0.24 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \text { day }^{2},  \tag{C19}\\
& \left|\ddot{\lambda}_{42}(\max )\right|_{\text {Syncom } 2=0.001 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \text { day }^{2}, ~}^{\text {, }}  \tag{C20}\\
& \left|\ddot{\lambda}_{44}(\max )\right|_{\text {syncom } 2}=0.012 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \mathrm{day}^{2} \tag{C21}
\end{align*}
$$

## Conclusions

Seven arcs of longitude-acceleration data ( $\operatorname{arcs} 1,2,4,5 \mathrm{~A}, 5^{\prime}, 5 \mathrm{~B}$ and 8 in Table 10) are available from Syncom 2 drift with standard errors from ( $0.03-0.11$ ) $\times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$. With this data alone, only both harmonic coefficients of $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ should be well discriminated. Additionally, one arc of longitude-acceleration data (arc 9) from Early Bird is available with a standard error of $0.01 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$. Sun, moon, and model bias accelerations average about $0.02 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$ for all of the above ares.

It is anticipated then, that if all the data without model bias adjustment are used, $\mathrm{H}_{22}$ and $\mathrm{H}_{33}$ should be well discriminated, and the effects of $\mathrm{H}_{31}$ should be marginally apparent. The separate effects of fourth order resonant earth gravity are probably at or below the average noise level of this experiment. Only an indication of the probable order of magnitude of fourth order effects should be realizable from the available limited acceleration record.

## Appendix D

## The Approximate Longitude Excursion of a Slowly Drifting $\mathbf{2 4 - H o u r ~ S a t e l l i t e ~}$

## Introduction

A very close approximation to the geographic drift excursion in a resonant gravity field of a 24-hour satellite follows the differential equation of motion, Equation 2, (see also Reference 2) and is, evidently, given by an elliptical integral such as that developed in Appendix $E$ of Reference 3. If the excursion itself is limited in extent, a simple closed form of this solution in terms of harmonic functions becomes applicable (References 2 and 9 ). Unfortunately this solution is essentially nonlinear in the unknown gravity constants. This makes the extraction of these constants from timelongitude data somewhat cumbersome statistically. If the time in the excursion is limited, a simple closed form of the elliptic integral solution is applicable in terms of a power series in the time (Reference 3). This solution is essentially linear in the gravity constants and will be developed here, by a Taylor series, to cover the excursion $\Delta \lambda$ from a longitude $\lambda_{0}$ where the drift rate $\dot{\lambda}_{0}$ is low but not zero (as in Reference 3). This solution (within appropriate limits on the excursion time $\Delta t$ ) should then be suited to describe the slow drift of the 24 -hour satellites in arcs $1,2,6,7$, 8 and 9 of Section 1.

## Development

Expanding the drift of the 24 -hour satellite from $\lambda_{0}$ (i.e., the longitude of the ascending Equator crossing) in a Taylor series in the excursion time $\Delta t$ gives

$$
\begin{equation*}
\Delta \lambda=\lambda-\lambda_{0}=\dot{\lambda}_{0} \Delta t+\frac{\ddot{\lambda}_{0} \Delta t^{2}}{2}+\frac{\lambda_{0}^{(3)} \Delta t^{3}}{6}+\frac{\lambda_{0}^{(4)} \Delta t^{4}}{24}+\cdots . \tag{D1}
\end{equation*}
$$

When resonant earth terms of higher than second order are ignored, Equation 2 gives

$$
\begin{equation*}
\ddot{\lambda}=-\mathrm{A}_{22} \sin 2 \gamma, \tag{D2}
\end{equation*}
$$

where $\gamma$ is the longitude of the 24 -hour satellite east of the minor axis of the elliptical Equator, and

$$
\begin{equation*}
A_{22}=-72 \pi^{2} \mathrm{~J}_{22}\left(\mathrm{R}_{0} / \mathrm{a}_{\mathrm{s}}\right)^{2}\left[\frac{1}{4}\left(\cos \mathrm{i}_{\mathrm{s}}+1\right)^{2}\right] \mathrm{rad} / \mathrm{sid} . \text { day }^{2} . \tag{D3}
\end{equation*}
$$

Differentiating Equation D2 with respect to time gives

$$
\begin{equation*}
\lambda^{(3)}=-2 \mathrm{~A}_{22} \dot{\gamma} \cos 2 \gamma=-2 \mathrm{~A}_{22} \lambda \cos 2 \gamma \tag{D4}
\end{equation*}
$$

Similarly, differentiation of Equation D4 shows that

$$
\begin{align*}
\lambda^{(4)} & =-2 A_{22} \ddot{\lambda} \cos 2 \gamma+4 A_{22}(\dot{\lambda})^{2} \sin 2 \gamma \\
& =A_{22} \sin 4 \gamma+4 A_{22}(\dot{\lambda})^{2} \sin 2 \gamma \tag{D5}
\end{align*}
$$

Substituting Equations D5, D4 and D2 at $\lambda=\lambda_{0}$ (or $\gamma=\gamma_{0}$ ) into Equation D1 gives the excursion from $\lambda_{0}$ to fourth order in $\Delta t$ as

$$
\begin{align*}
\Delta \lambda & =\dot{\lambda}_{0} \Delta t-\frac{A_{22}}{2}\left(\sin 2 \gamma_{0}\right) \Delta t^{2}-\frac{A_{22}}{3}\left(\cos 2 \gamma_{0}\right) \dot{\lambda}_{0} \Delta t^{3} \\
& +\frac{1}{24}\left[A_{22}^{2} \sin 4 \gamma_{0}+4 A_{22}\left(\dot{\lambda}_{0}\right)^{2} \sin 2 \gamma_{0}\right] \Delta t^{4}+\cdots \tag{D6}
\end{align*}
$$

The adequacy of a series such as Equation D6 to approximate the actual drift motion stemming from the basic differential equation is discussed at greater length in Reference 3.

A typical value for the "longitude noise" level (due to sun and moon effects and observational error) in the actual gravity experiment is $0.025^{\circ}$ (see Section 1). We wish to apply a polynomial fit to the actual data to the third degree in $\Delta t$, by way of Equation D6. Such a fit allows the acceleration in any slow drift arc to vary with time as the satellite samples different longitudes in the gravity field. By introducing an additional degree of freedom it also permits the precision of the acceleration determination to vary. Under normal circumstances the "best acceleration" measurement will occur with this fit near the "centroid" of the longitude-time data. This statistical result coincides with our intuition of where the best measured parameters of the drift should occur.

Let us assume that at $\Delta t=0, \lambda=\lambda_{0}, \gamma=\gamma_{0}$ the drift rate is $\dot{\lambda}_{0}=0.1$ degree/day. Then Equation D6, truncated at the $\Delta t^{3}$ term, will hold to an error of the order of $0.025^{\circ}$ at the end of an arc
time length $\Delta t$ if

$$
\begin{equation*}
\frac{1}{24}\left(A_{22}^{2} \sin 4 \gamma_{0}+4 A_{22}\left(\frac{0.1}{57.3}\right)^{2} \sin 2 \gamma_{0}\right) \Delta t^{4} \leq 0.025^{\circ} \tag{D7}
\end{equation*}
$$

with $A_{22}$ in units of rad/sid. day ${ }^{2}$. In Reference 3, $A_{22}$ for measured Syncom 2 drift over Brazil was

$$
A_{22}=23.2 \times 10^{-6} \mathrm{rad} / \mathrm{sid} . \text { day }^{2}
$$

corresponding to

$$
\mathrm{i}_{\mathrm{s}}=33^{\circ} \text { and } \mathrm{J}_{22} \text { (measured) }=-1.7 \times 10^{-6} .
$$

We wish to evaluate Equation D7 in the most conservative drift condition (giving the minimum $\Delta t$ for the inequality to apply). In Appendix $F$ (from Figure 9) it is estimated that for drift of an equatorial satellite over the Indian Ocean, $\mathrm{J}_{22}$ (measured) $\simeq-1.93 \times 10^{-6}$, giving the strongest longitude acceleration on the 24 -hour satellite. From Equation D3, the $\mathrm{A}_{22}$ measured for this satellite location would be

$$
\begin{equation*}
\mathrm{A}_{22}(\max )=23.2 \times 10^{-6} \times\left(\frac{1.94}{1.7}\right) \times\left(\frac{2}{\cos 33.0^{\circ}+1}\right)^{2}=31.3 \times 10^{-6} \mathrm{rad} / \mathrm{sid} . \text { day }^{2} \tag{D8}
\end{equation*}
$$

Equation D8 in Equation D7 gives the critical inequality as

$$
\begin{equation*}
\left(9.797 \times 10^{-10} \sin 4 \gamma_{0}+3.814 \times 10^{-10} \sin 2 \gamma_{0}\right) \Delta t^{4} \leq 0.01047 \tag{D9}
\end{equation*}
$$

with $\Delta t$ in units of days. The factor of $\Delta t^{4}$ in Equation D9 is maximum when $2 \gamma_{0} \doteq 48.7^{\circ}$, or the satellite is about $24.4^{\circ}$ east of one of the two minor equatorial axis longitudes. At this longitude the $\Delta t^{4}$ factor is $12.58 \times 10^{-10}$. From Equation D9, the minimum time for inclusion of the $\Delta t^{4}$ term in Equation D6 according to the above criteria is

$$
\begin{align*}
\Delta t_{\min }(\text { for fourth order term inclusion in slow } 24 \text {-hour satellite drift }) & =\left(0.0832 \times 10^{8}\right)^{1 / 4} \\
& = \pm 53.7 \text { days } \tag{D10}
\end{align*}
$$

The conservative time criteria in Equation D10 shows that Equation D6 should give an adequate description of the slow drift regime in all the 24 -hour satellite arcs of Section 1 . The excellent agreement of theoretical accelerations with measured results from $\Delta t^{3}$ analyses of the simulated data in these arcs confirms the adequacy of this description. (See the bias results in Table 11.)

Let the drift time from an arbitrary base time $t$ (i.e. the beginning of the year or the middle of the arc for statistical convenience) be given by $t$. Let $t=0$ be the time (with respect to the base time) when the satellite is at $\lambda=\lambda_{0}$ moving at drift rate $\dot{\lambda}_{0}$. Let the drift be given from an arbitrary base longitude by $L$. Let the longitude $\lambda_{0}$ with respect to the base longitude $\underline{L}$ be $L_{0}$. Then


Equations D11 into Equation D6 truncated at the $\Delta t^{3}$ term gives the longitude excursion $L$ as

$$
\begin{equation*}
L=L_{0}+\left(\dot{\lambda}_{0}\right) t-\left(\frac{A_{22}}{2} \sin 2 \gamma_{0}\right) t^{2}-\left(\frac{A_{22}}{3} \dot{\lambda}_{0} \cos 2 \gamma_{0}\right) t^{3} \tag{D12}
\end{equation*}
$$

Let

$$
\begin{align*}
& a_{1}=L_{0} \\
& a_{2}=\dot{\lambda}_{0} \\
& a_{3}=-\frac{A_{22}}{2} \sin 2 \gamma_{0}  \tag{D13}\\
& a_{4}=-\frac{A_{22} \dot{\lambda}_{0}}{3} \cos 2 \gamma_{0}
\end{align*}
$$

Equations D13 in Equation D12 gives the drift from the arbitrary longitude $\underline{L}$ in the terms of the time from an arbitrary base time $t$ as

$$
\begin{equation*}
L=a_{1}+a_{2} t+a_{3} t^{2}+a_{4} t^{3} \tag{D14}
\end{equation*}
$$

where the four arbitrary constants of Equation D14 serve to define the elements $L_{0}, \dot{\lambda}_{0}, A_{22}$ and $\gamma_{0}$ of the arc dynamics in terms of a purely second order resonant gravity drift, presumably dominant at 24-hour altitudes.

It is easily seen from Equation 2 that the effect of higher order resonant gravity on the slow drift expansion, Equation D6, is to merely alter the coefficients of the $\Delta t^{2}$ and higher order terms. The fundamental polynomial representation does not change.
i

## Appendix E

## The Secular Accelerations on Syncom 24 -Hour Satellites Due to Particle Atmospheric Drag and Solar Radiation Pressure

## Introduction

We have seen previously (Section 1) that the long term tracking of the three 24 -hour satellites (modeled after Syncom) has provided a discrimination of drift acceleration of the order of $0.1 \times 10^{-5}$ $\mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$. It will be instructive to compare this figure with the likely long term accelerations on these satellites due to particle-atmospheric drag and solar radiation pressure. In performing this calculation, we will follow the work of D. D. Williams in Reference 7. We shall assume that long term secular accelerations in drift from these sources will be negligible on the gravity analysis here if these accelerations are of the order of magnitude of $0.01 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$ or less.

## Atmospheric-Particle Drag

With respect to the drag caused by a continuum of nonmeteoritic particles in the upper atmosphere, Williams distinguishes two cases. The first is drag caused by particles of residual earth atmosphere, at rest but not rotating with respect to the earth. The second is drag caused by the interplanetary particles, at rest with respect to the sun in the earth's orbital path.

## Residual Atmospheric Drag

For a conservative Syncom spacecraft mass of only 0.75 slug and fully elastic collisions, Williams calculates the density of residual atmosphere at 24 -hour altitudes necessary to produce one degree of longitude drift in a year to be

$$
\rho=0.1826 \times 10^{-17} \mathrm{gm} / \mathrm{cc} \text { (with drift proportional to density). }
$$

An estimate of the order of magnitude of the residual atmosphere at 24-hour altitudes is (Reference E1 at the end of this appendix, pp. 2-8)

$$
\rho \doteq 10^{-21} \mathrm{gm} / \mathrm{cc}
$$

The acceleration, in radians per sidereal day ${ }^{2}$, necessary to give a drift of one degree in a year is (from $\Delta \lambda=1 / 2 \ddot{\lambda} \Delta t^{2}$ ):

$$
\ddot{\lambda}=\frac{2 \Delta \lambda}{\Delta \mathbf{t}^{2}}=\frac{2 \times 1^{\circ}}{57.3^{\circ} / \mathrm{rad} \times(366 \mathrm{sid} . \text { day })^{2}}=0.261 \times 10^{-6} \mathrm{rad} / \mathrm{day}^{2}:
$$

Thus, for a Syncom spacecraft (including empty apogee motor) of 2.34 slugs, the drift acceleration due to residual atmosphere drag should be of the order of

$$
\begin{aligned}
\ddot{\lambda}(\text { residual atmosphere }) & \simeq 0.261 \times 10^{-6} \times\left(10^{-21} / 0.1826 \times 10^{-17}\right) \times(0.75 / 2.34) \mathrm{rad} / \mathrm{sid} . \text { day }^{2} \\
& =0.46 \times 10^{-10}=0.0000046 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \text { day }^{2}
\end{aligned}
$$

Clearly, the long term acceleration effects of residual earth atmosphere on the 24-hour satellite gravity experiment are negligible.

## Interplanetary Particles (Non Meteoritic)

Again, for a Syncom spacecraft mass of 0.75 slug and fully elastic collisions, Williams calculates the density of (solar stationary) interplanetary particles necessary to produce one degree of longitude drift in a year to be

$$
\rho=0.2432 \times 10^{-18} \mathrm{slug} / \mathrm{ft}^{3}=0.1257 \times 10^{-18} \mathrm{gm} / \mathrm{cc}
$$

with the drift proportional to the density. From Reference E1, pp. 2-8, the interplanetary density is of the order of $10^{-22} \mathrm{gm} / \mathrm{cc}$. Thus, for a Syncom spacecraft of 2.34 slugs, the drift acceleration due to solar stationary interplanetary particles should be of the order of

$$
\begin{aligned}
\ddot{\lambda}(\text { interplanetary particles, non meteoritic }) & \simeq 0.261 \times 10^{-6} \times\left(10^{-22} / 0.1257 \times 10^{-18}\right) \times(0.75 / 2.34) \\
& =0.0000067 \times 10^{-5} \mathrm{rad} / \mathrm{sid} . \mathrm{day}^{2} .
\end{aligned}
$$

In summary, the effects of residual atmospheric and solar-stationary interplanetary particle drag on the 24 -hour satellite gravity experiment appear to be insignificant.

## Solar Radiation Pressure

In Reference E2, Appendix F, it was shown that the solar radiation force on Syncom was about five orders of magnitude less than the solar gravity force. But the conclusion does not follow that,
because of this, the solar radiation pressure perturbations are insignificant compared to solar gravity perturbations. The reason is that solar gravity perturbations of an earth satellite arise from the slightly different solar gravity accelerations experienced by the earth and its satellite. Solar radiation pressure caused accelerations of the earth, on the other hand, are extremely small because the earth presents to the sun a very small projected area/mass ratio. Thus between the two effects, we ought to compare the differences (between earth and satellite) of large accelerations due to solar gravity, with the small acceleration on the satellite alone due to the solar radiation pressure. It is by no means certain which effect will predominate for any given satellite. Experience with the Echo satellites has shown that for satellites with sufficiently large area/mass ratios, radiation pressure effects cannot be ignored.

Instead of making a direct comparison of solar gravity and radiation effects, since they do not act analogously (as was tacitly assumed in Reference E2), we shall merely calculate here the secular effect of radiation pressure pertinent to the 24 -hour satellite gravity experiment.

In Reference E2, Appendix F, it is calculated that in August 1963 the radiation force on Syncom 2 (for full radiant energy absorption) was $4.13 \times 10^{-7}$ pounds. With a spacecraft mass of 2.34 slugs, this gives an acceleration caused by radiation pressure of $4.13 \times 10^{-7} / 2.34=1.77 \times 10^{-7} \mathrm{ft} / \mathrm{sec}^{2}$. In Reference 7 on p. 13, for a spacecraft mass of 0.75 slug, the radiation acceleration is calculated as $4.30 \times 10^{-7} \mathrm{ft} / \mathrm{sec}^{2}$ for a fully reflecting equatorial spacecraft at the equinoxes. For a 2.34 slug equatorial Syncom, fully reflecting (with spin axis pointing north) at the equinoxes, the radiation acceleration would be $4.30 \times 10^{-7} \times(0.75 / 2.34)=1.38 \times 10^{-7} \mathrm{ft} / \mathrm{sec}^{2}$, since the radiation pressure acceleration is inversely proportional to the mass (at constant projected area). This latter acceleration actually represents the most conservative value for Syncom for our purposes since in this configuration the sun is in the orbit plane and no part of the solar pressure is wasted on "plane change effects".

Consider in Figure E1 the Syncom-earthsun orbital geometry in this latter configuration (from Figures 2-4, Reference 7), where $a$ is the earth rate about the sun, and $\theta_{\mathrm{s}}$ is the orbital argument of Syncom with respect to the moving earth-sun line. At $\mathrm{t}=0, \theta_{\mathrm{s}}=\theta_{0}$ which is the initial sun-satellite argument, for example, at the first ascending Equator crossing in an arc. Clearly, when $\theta_{\text {s }}$ is positive, solar pressure acts to retard the motion of the satellite. When $\theta_{\mathrm{s}}$ is negative, the pressure adds energy to the satellite. In fact, from Figure E1 the energy adding (or tangential) force $F$ (per unit mass)


Figure E1-Equatorial Syncom-earth-sun orbital configuration at the equinoxes, looking south (after Williams).
on Syncom due to solar pressure is

$$
\mathbf{F}=-\mathbf{F}_{\mathrm{s}} \cos \beta=-\mathbf{F}_{\mathrm{s}} \sin \theta_{\mathrm{s}}^{\prime} \doteq-\mathbf{F}_{\mathrm{s}} \sin \theta_{\mathrm{s}}
$$

since $\theta_{s}^{\prime} \doteq \theta_{\mathrm{s}}$ due to the small "parallax" of the sun from the 24 -hour orbit ( $\sim 1$ min of arc). From Figure E1 also

$$
\theta_{s}=\omega_{s} t-a t+\theta_{0}
$$

Thus

$$
\mathrm{F} \doteq-\mathrm{F}_{\mathrm{s}} \sin \left[\left(\omega_{\mathrm{s}}-\alpha\right) \mathrm{t}+\theta_{0}\right]=-\mathrm{F}_{\mathrm{s}}\left[\cos \theta_{0} \sin \left(\omega_{\mathrm{s}}-\alpha\right) \mathrm{t}+\sin \theta_{0} \cos \left(\omega_{\mathrm{s}}-\alpha\right) \mathrm{t}\right]
$$

The orbit averaged $\left(\omega_{s} t=2 \pi\right)$ solar pressure perturbing force (per unit mass) is then

$$
\begin{aligned}
\overline{\mathbf{F}} & =-\frac{\mathbf{F}_{\mathrm{s}}}{\mathbf{1} \operatorname{sid} . \operatorname{day}} \int_{0}^{2 \pi / \omega_{\mathrm{s}}}\left[\cos \theta_{0} \sin \left(\omega_{\mathrm{s}}-\alpha\right) \mathrm{t}+\sin \theta_{0} \cos \left(\omega_{\mathrm{s}}-\alpha\right) \mathrm{t}\right] \mathrm{dt} \\
& =\frac{+\mathbf{F}_{\mathrm{s}}}{\omega_{\mathrm{s}}-\alpha}\left|\cos \theta_{0} \cos \left(\omega_{\mathrm{s}}-\alpha\right) \mathrm{t}-\sin \theta_{0} \sin \left(\omega_{\mathrm{s}}-\alpha\right) \mathrm{t}\right|_{0}^{2 \pi / \omega_{\mathrm{s}}} \\
& =\frac{\mathbf{F}_{\mathrm{s}}\left[\cos \left(\frac{2 \pi \alpha}{\omega_{\mathrm{s}}}\right)-1\right] \cos \theta_{0}+\mathbf{F}_{\mathrm{s}}\left[\sin \left(\frac{2 \pi \alpha}{\omega_{\mathrm{s}}}\right)\right] \sin \theta_{0}}{\left(\omega_{\mathrm{s}}-\alpha\right)} \\
& \doteq \frac{\mathbf{F}_{\mathrm{s}}}{\omega_{\mathrm{s}}-\alpha}\left[\sin \left(\frac{2 \pi \alpha}{\omega_{\mathrm{s}}}\right)\right] \sin \theta, \text { for } \frac{2 \pi \alpha}{\omega_{\mathrm{s}}} \operatorname{sma} 11,
\end{aligned}
$$

with $\omega_{\mathrm{s}}$ and $\alpha$ in units of rad/sid. day. But $\omega_{\mathrm{s}}=2 \pi$ and $\alpha \doteq 2 \pi / 366=0.01715$. Therefore $2 \pi \alpha / \omega_{\mathrm{s}}$ $\doteq 0.01715$ radians and $\omega_{\mathrm{s}}-a \doteq 6.26 \mathrm{rad} / \mathrm{sid}$. day. Thus

$$
\bar{F}(\text { solar pressure }) \doteq \frac{F_{s} \times 0.01715 \sin \theta_{0}}{6.26}=2.74 \times 10^{-3} F_{\mathrm{s}} \sin \theta
$$

which is in units of acceleration $\left(F_{s}\right)$. It is assumed in these calculations that $F_{s}$ is a constant over the 24 -hour orbit. For $\mathrm{F}_{\mathrm{s}}=1.38 \times 10^{-7} \mathrm{ft} / \mathrm{sec}^{2}$.
$\overline{\mathrm{F}}_{\text {max }}($ solar pressure on Syncom $) \doteq 3.78 \times 10^{-10} \mathrm{ft} / \mathrm{sec}^{2}$, when $\theta_{0}=90^{\circ}$. From Equation 10 in Reference E2, the long term longitude acceleration on the 24 -hour satellite is given from $\overline{\mathrm{F}}$ by

$$
\ddot{\lambda}=\frac{-12 \pi^{2} \bar{F}}{\mu_{\mathrm{e}} / \mathrm{a}_{\mathrm{s}}{ }^{2}}=-\frac{12 \pi^{2} \overline{\mathrm{~F}}}{\mathrm{~g}_{\mathrm{s}}}\left(\mathrm{rad} / \mathrm{sid} . \text { day }^{2}\right)
$$

where $g_{s}$ is the radial gravity acceleration ( $\doteq 0.7355 \mathrm{ft} / \mathrm{sec}^{2}$ ) on Syncom. Thus
$\ddot{\lambda}_{\text {max }}$ (due to solar pressure on Syncom satellites) $\doteq \frac{+12 \pi^{2}}{0.7355} \times 3.26 \times 10^{-10}=0.00609 \times 10^{-5} \mathrm{rad} / \mathrm{sid}$. day ${ }^{2}$

It is evident that solar radiation pressure has negligible effect on the 24-hour satellite gravity experiment. The maximum total longitude excursion due to solar pressure in this conservative case accumulates over half a year and is approximately

$$
\Delta \lambda_{\max }=\frac{1}{2}\left(0.637 \times 6.09 \times 10^{-8}\right)(183)^{2} \times 57.3=0.0372^{\circ} .
$$

The factor 0.637 is the average of $\sin \theta_{0}$ for $0<\theta_{0}<\pi$, which is the range of $\theta_{0}$ over half a year.

## Summary

The effects of high altitude atmospheric particle drag and solar radiation pressure on the 24-hour satellite experiment should be entirely negligible.

## REFERENCES

E1. Koelle, H. H., "Handbook of Astronautical Engineering," New York: McGraw-Hill, 1961.
E2. Wagner, C. A., "Determination of the Triaxiality of the Earth from Observations on the Drift of the Syncom 2 Satellite," GSFC Document X-621-64-90, April 1964.

## Appendix F

## Average Second Order Resonant Gravity Fields on the Geostationary Satellite

In the near future, many 24 -hour equatorial satellites (mainly for communication purposes) will be placed and maintained in orbit at selected longitudes around the Equator. It will usually be necessary to keep these satellites within a longitude band about $20^{\circ}$ wide or less depending on the location of the ground stations and the sophistication of the transmission and tracking equipment. Since the $\mathrm{J}_{22}$ gravity field is dominant at 24 -hour altitudes, it may be convenient, in predicting the trajectory of these nearly "fixed" satellites between orbit corrections, to consider the average $\mathrm{J}_{22}$ field on the geostationary satellite over wide longitude arcs.

Average $\mathrm{J}_{22}$ gravity fields on the geostationary satellite are most naturally grouped into four longitude zones (of about $90^{\circ}$ each) surrounding the four equilibrium longitudes (see Figure 9 and Section 3). To determine these average fields we first find the average of the peak accelerations in Figure 9 (without sign) on either side of these equilibrium longitudes. The average $\mathrm{J}_{22}$ in the region between these relative maxima is then determined by solving Equation 2 for $\mathrm{J}_{22}$ with all other gravity coefficients zero, $i_{s}=0^{\circ}$, $a_{s}=6.611$ earth radii and $2\left(\lambda-\lambda_{22}\right)=90^{\circ}$. The effective $\lambda_{22}$ in this region is given by the corresponding equilibrium longitude $\lambda_{e}$ (see Section 3). The results of this calculation follow in Table F1 below.

Table F1

Average Second Order Resonant Gravity Fields on the Geostationary Satellite*.

| Longitude Region | $\bar{J}_{22}$ | $\bar{\lambda}_{22}$ <br> (degrees) | Equilibrium Longitude <br> (degrees) |
| :---: | :---: | :---: | :---: |
| $34^{\circ}<\lambda<118^{\circ}$ | $-1.93 \times 10^{-6}$ | -13.3 | 76.7 |
| $118^{\circ}<\lambda<204^{\circ}$ | $-1.87 \times 10^{-6}$ | -18.2 | 161.8 |
| $204^{\circ}<\lambda<300^{\circ}\left(-60^{\circ}\right)$ | $-1.73 \times 10^{-6}$ | -18.1 | $251.9(-108.1)$ |
| $-60^{\circ}<\lambda<34^{\circ}$ | $-1.79 \times 10^{-6}$ | -12.2 | -12.2 |

[^4]
## Appendix G

## List of Symbols

a, $a_{s} \quad$ Semimajor axis and synchronous semimajor axis of the orbit of a 24-hour satellite. (More specifically, $a_{s}$ is defined in the text and tables of Section 1 as the average semimajor axis of the satellite during a drift arc.)
$C_{n m}, S_{n m} \quad$ The cosine and sine parameters of the $n, m$ gravity harmonic.

F A gravity force per unit mass acting on a 24-hour satellite.
$F_{n m}$ (i) Inclination factor corresponding to the 24-hour drift caused by the $n$, $m$ resonant gravity harmonic.
$F_{22}\left(i_{s}, a_{s}\right)$ A "constant" of the 24-hour drift motion caused by the $H_{22}$ resonant gravity harmonic, a function of the average inclination and semimajor axis in the drift arc.
$H_{n m} \quad$ Specifying the $n$, m earth gravity harmonic.
i, $i_{s} \quad$ Orbit inclination and synchronous orbit inclination during a 24-hour drift arc.
$J_{n m}, \lambda_{n m} \quad$ Amplitude and geographic phase of the $n, m$ earth gravity harmonic.
$\mathrm{R}_{0} \quad$ The mean equatorial radius of the earth $(\doteq 6378.2 \mathrm{~km})$.
s( ) Standard error of the bracketed quantity ().
$\lambda, r, \phi \quad$ Geographic longitude, geocentric radius, and geocentric latitude of the 24-hour satellite position.
$\mu_{\mathrm{e}} \quad$ Earth's Gaussian gravity constant $\left(\doteq 3.986 \times 10^{5} \frac{\mathrm{~km}^{3}}{\mathrm{sec}^{2}}\right)$.


#### Abstract

"The aeronautical and space activities of the United States shall be conducted so as to contribute . . . to the expansion of human knowledge of phenomena in the atmospbere and space. The Administration shall provide for the widest practicable and appropriate dissemination of information concerning its activities and the results thereof."


## NASA SCIENTIFIC AND TECHNICAL PUBLICATIONS

TECHNICAL REPORTS: Scientific and technical information considered important, complete, and a lasting contribution to existing knowledge.

TECHNICAL NOTES: Information less broad in scope but nevertheless of importance as a contribution to existing knowledge.

TECHNICAL MEMORANDUMS: Information receiving limited distribution because of preliminary data, security classification, or other reasons.

CONTRACTOR REPORTS: Technical information generated in connection with a NASA contract or grant and released under NASA auspices.

TECHNICAL TRANSLATIONS: Information published in a foreign language considered to merit NASA distribution in English.

TECHNICAL REPRINTS: Information derived from NASA activities and initially published in the form of journal articles.

SPECIAL PUBLICATIONS: Information derived from or of value to NASA activities but not necessarily reporting the results of individual NASA-programmed scientific efforts. Publications include conference proceedings, monographs, data compilations, handbooks, sourcebooks, and special bibliographies.

Details on the availability of these publications may be obtained from:

SCIENTIFIC AND TECHNICAL INFORMATION DIVISION
NATIONAL AERONAUTICS AND SPACE ADMINISTRATION
Washington, D.C. 20546


[^0]:    See Footnotes Page 25.

[^1]:    *See Footnotes Poge 27.

[^2]:    ${ }^{*}$ Gravity constants of this crajectory (computed by ITEM) are the same ns those io Table A1, with the addition of the earth constants: $J_{22}=-1.8 \times 10^{-5}, \lambda_{22}=-15.35^{\circ}, J_{33}=-0.16 \times 10^{-6}, \lambda_{33}=24^{\circ}$,

[^3]:    ${ }^{* *}$ A number of attitude and orbit inclination change maneuvers were performed on Syncom 2 between 14 January and 30 January 1965. These did not appear to affect the mean motion of the satellite significantly.
    Results of least squares of data in (1) and (2) above according to the theory of Equation 1:

    $$
    \begin{aligned}
    & \mathbf{L}=\mathbf{a}_{1}+\mathbf{a}_{2} \mathbf{t}+\mathbf{a}_{3} \mathrm{t}^{2}+\mathbf{a}_{4} \mathbf{t}^{3} \\
    & \mathbf{a}_{\mathbf{1}}=-(0.4168 \pm 0.0244) \text { degrees }
    \end{aligned}
    $$

    $$
    a_{2}=-(0.1166 \pm 0.0023) \text { degree } / \text { day }
    $$

    $$
    \mathrm{a}_{3}=(4.409 \pm 0.505) \times 10^{-4} \text { degrees } / \text { day }^{2}
    $$

    $$
    \mathrm{a}_{4}=-(4.131 \pm 3.042) \times 10^{-6} \text { degrees } / \text { day }^{3}
    $$

    Standard error of estimate $=0.0515^{\circ}$
    $\ddot{\lambda}\left(\right.$ with minimum standard error) $=(1.550 \pm 0.175) \times 10^{-5} \mathrm{rad} / \mathrm{sid}^{\text {. day }}{ }^{2}$, at $\mathrm{t}=-0.4507$ day, $\mathrm{t}^{\prime}=44.9631$ January $1965, \mathrm{~L}=-0.3634^{\circ}, \lambda=176.801^{\circ}$ (see Figure 6 ).

[^4]:    *Representing the average effects of a geoid without higher order longitude gravity on a 24 -hour equatorial satellite ( $a_{s}=6.611$ earth
    radii) in the given longitude region radii) in the given longitude region.

