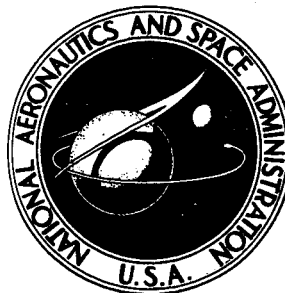


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CORRECTION OF BALLOON X-RAY ASTRONOMY DATA FOR THE EFFECTS OF ATMOSPHERIC ATTENUATION, K X-RAY ESCAPE, AND ENERGY RESOLUTION

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Prepared by
MASSACHUSETTS INSTITUTE OF TECHNOLOGY
Cambridge, Mass.

for



CORRECTION OF BALLOON X-RAY ASTRONOMY DATA FOR
THE EFFECTS OF ATMOSPHERIC ATTENUATION,
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Abstract

An x-ray flux from outside the atmosphere determines the counting rate of an x-ray detector beneath several g/cm^2 of atmosphere, but a unique inverse relationship does not exist. Therefore x-ray astronomy data analysis requires the computer simulation of the processes which transform assumed x-ray energy spectra into detectors' counting rates. The programs used by the author, E. Allen Womack, Jr., and Harvey D. Tananbaum are presented and are applied to a specific example. A set of conventions for presenting one's results as an x-ray spectrum above the atmosphere is also explained.

I. Introduction

There is no one-to-one correspondence between the counting rate of an x-ray detector beneath several g/cm² of atmosphere and the x-ray flux at the top of the atmosphere. This is primarily due to two effects which cause pulse heights from a detector not to be uniquely related to the energies of the x-rays which produced them. The first is the sum of statistical fluctuations in light photon and photoelectron production, usually lumped under the familiar term "energy resolution". The second, the escape of K x-rays from such high Z elements as Iodine or Xenon in the detector, causes a pulse to indicate about 30 keV less energy than the original photon actually had. Results of calculations on this effect have been presented by Stein and Lewin¹.

Another effect of the same sort is the partial deposition of energy by x-rays which Compton scatter in the detector and then leave. This is important above 100 keV and is extremely important for energies of a few MeV. This effect is neglected here because we are mainly concerned with energies below 100 keV. A more important effect is the increase in the detector's total absorption probability (referred to here as efficiency) due to the increased path length of an x-ray which has Compton scattered. The data^{2,3} of the author, E. Allen Womack, Jr., and Harvey D. Tananbaum have had such large statistical errors above 50 keV (where this effect is most important) that we have neglected this effect for the time being. Calibration with radioactive isotopes having known intensity ratios for their gamma ray lines will be used when the need arises.

In reference 3 the author and Harvey D. Tananbaum reported a significant increase in the x-ray flux from Cygnus XR-1 over a period of eight months. One of the main purposes of this report is to allow other investigators to compare their data with ours with the assurance that agreement or disagreement is not due to differences in data analysis or data presentation.

This report has four parts. The first part is a description and listing of a program, DETMD, which generates thermal bremsstrahlung and power law x-ray spectra, attenuates them by the atmosphere, multiplies them by the detector's efficiency, shifts x-rays in energy to account for K x-ray escape, and smears the result with a gaussian, energy dependent resolution function. The inputs to DETMD are tables of detector efficiency, K x-ray escape probabilities, attenuation coefficients for air, x-ray spectrum parameters, g/cm² of air considered, detector energy resolution parameters, and pulse height analyzer energy thresholds. Its purpose is to allow the user to calculate the ratio of x-ray fluxes in several energy bands to counting rates in the corresponding pulse height bands.

The second part consists of the programs we used to generate the tables of air attenuation coefficients, efficiency, and K x-ray escape probabilities. The efficiency and K x-ray escape program, NAIEF, applies only to our 1 mm thick NaI crystal with a .005 inch beryllium window. The user will have to modify it or start from scratch to generate efficiency tables for his own detector. The program is included here to allow others to check our analysis and to generate sample input data for program DETMD.

In the third part is presented a self consistent set of conventions for presenting one's results as a graph of an x-ray source's spectrum (above the atmosphere), thereby allowing accurate visual comparison of different experimenters' results.

The computer programs were written in Fortran IV for the IBM 7044, but they should run equally well on an IBM 7094 equipped with the operating system IBSYS. For operation on the IBM 360 some numbers will have to be truncated to 7 decimal digits. The Fortran logical designations of the input-output devices are: 5, card reader; 6, printer; and 7, card punch. As originally written the programs contained many statements which caused the generation of graphs of the results on a Calcomp

digital incremental plotter. This is the cause of some structures in the programs which seem to be awkward or unnecessary. It is also the cause of the small step size (0.1 keV) used in all the calculations.

II. Program DETMD

The assumed differential number spectra above the atmosphere are

$$\left(\frac{dn}{dE}\right)_0 = A \left(\frac{E_0}{E}\right) e^{-(E-E_0)/kT} \quad \text{or} \quad A \left(\frac{E_0}{E}\right)^N \quad \text{if}$$

NTYPE = 1 or 2 respectively.

A = ANORM, E_0 = EREF, kT = EKT, N = POWER

The spectrum at the detector is $\left(\frac{dn}{dE}\right)_1 = \left(\frac{dn}{dE}\right)_0 e^{-\mu(E)X}$

μ = the attenuation coefficient of air = ATTEN

X = g/cm² of air along line of sight = GRAMS.

The spectrum of energy losses in the detector is

$$\begin{aligned} \left(\frac{dn}{dE}(E)\right)_2 &= \left(\frac{dn}{dE}(E)\right)_1 [\varepsilon(E) - p_\alpha(E) - p_\beta(E)] \\ &+ \left(\frac{dn}{dE}(E + E_{K\alpha})\right)_1 p_\alpha(E + E_{K\alpha}) \\ &+ \left(\frac{dn}{dE}(E + E_{K\beta})\right)_1 p_\beta(E + E_{K\beta}) \end{aligned}$$

ε = the efficiency of the detector, i.e. the probability that an incident photon will interact = EFFIC

p_α = the probability that an incident photon will give rise to a K_α photon which escapes from the detector = ESCKA

p_β = the probability that an incident photon will give rise to a K_β photon which escapes from the detector = ESCKB

$E_{K\alpha}$, $E_{K\beta}$ = energies of K_α and K_β photons = 32.5, 28.6 keV for iodine.

The pulse height spectrum is

$$\left(\frac{dn}{dE}(E)\right)_3 = \int_0^{\infty} \left(\frac{dn}{dE}(E')\right)_2 \frac{e^{-\frac{(E'-E)^2}{2\sigma(E')^2}}}{\sqrt{2\pi}\sigma(E')} dE'$$

where $\sigma(E') = aE' + b\sqrt{E'}$

$a = \text{ASIG}$, $b = \text{BSIG}$

The four spectra are printed on four separate pages. Each is integrated over 1 keV bands, 5 keV bands, and eight other bands whose boundaries are the inputs CALIB*THRESH. Following the listing of program DETMD are some of its results for $kT = 4.3$ keV and the author's detector parameters. The first page of results is $\left(\frac{dn}{dE}\right)_0$ for $A = 1.0$ and $E_0 = 20$ keV. The next three pages are the corresponding $\left(\frac{dn}{dE}\right)_1$, $\left(\frac{dn}{dE}\right)_2$, and $\left(\frac{dn}{dE}\right)_3$ for 4.0 g/cm² of air along the line of sight. The last three pages are $\left(\frac{dn}{dE}\right)_3$ for 4.5, 5.0, and 5.5 g/cm² of air. The following example illustrates how these results are used. Absorption by 5.0 g/cm² of air along the line of sight is assumed. One counts $n_{ab} = .002$ pulses/cm²sec which have pulse heights corresponding to energies between $E_a = 23.4$ and $E_b = 46.7$ keV. If one assumes the spectrum above the atmosphere had the form $\frac{K}{E} e^{-E/kT}$ with $kT = 4.3$ keV then the number of incident photons with energies between E_a and E_b was $N_{ab}(kT) = R_{ab}(kT) n_{ab} = .01666$, where

$$R_{ab} = \frac{\int_{E_a}^{E_b} \left(\frac{dn}{dE}\right)_0 dE}{\int_{E_a}^{E_b} \left(\frac{dn}{dE}\right)_3 dE} = \frac{1.19006 + .19147 + .03155}{.11938 + .04034 + .00999} = 8.33$$

(see underlined numbers in the computer printout)

The best value of the coefficient K is

$$K = \frac{A E_0 e^{\frac{E_0}{kT}} n_{ab}}{\int_{E_a}^{E_b} \left(\frac{dn}{dE} \right)_3 dE} = \frac{1 \times 20 e^{\frac{20}{4.3}} \times .002}{.16971} = 24.68$$

FORTRAN SOURCE LIST

ISN SOURCE STATEMENT

```

0 $IBFIC DETMD
1   DIMENSION ATEN(1900), EFFIC(2000), ESCKA(2000), ESCKB(2000),
1   GAUSS(2000), COUNT(2000), CTSA(40), CTSB(8), THRESH(9), Y(2000)
2   NPOINT = 2000
3   DO 10 INIT = 1, 2000
4   GAUSS(INIT) = EXP(-((FLOAT(INIT - 1)/500.0)**2)/2.0)
5   10 CONTINUE
7   DO 12 JFIX = 1, 100
10  EFFIC(JFIX) = 1.0
11  ESCKA(JFIX) = 0.0
12  ESCKB(JFIX) = 0.0
13  12 CONTINUE
15  READ(5, 15) (EFFIC(JA), JA = 101, 2000), (ESCKA(JB), JB = 101,
1  2000), (ESCKB(JC), JC = 101, 2000)
32  15 FORMAT(10F7.5)
33  READ(5, 20) ATEN
34  20 FORMAT(10F7.4)
35  50 READ(5, 52) NDSPEC, NCATM, NDRESL, NDRSH, NDCALB, NDQUIT
44  52 FORMAT(6I1)
45  IF(NDQUIT.NE.0) GO TO 1000
50  IF(NDSPEC.NE.1) GO TO 60
53  READ(5, 54) NTYPE, EREF, ANORM, TORPWR
55  54 FORMAT(I1, 3F7.3)
56  60 IF(NDATM.NE.1) GO TO 65
61  READ(5, 63) GRAMS
62  63 FORMAT(F7.4)
63  65 IF(NDRESL.NE.1) GO TO 75
66  READ(5, 68) ASIG, BSIG
67  68 FORMAT(2F10.5)
70  75 IF(NDRSH.NE.1) GO TO 80
73  READ(5, 78) THRESH
74  78 FORMAT(6F10.4/3F10.4)
75  80 IF(NDCALB.NE.1) GO TO 85
100 READ(5, 83) CALIB
101  83 FORMAT(F10.5)
102  85 CONTINUE
103 IF(NTYPE.EQ.2) GO TO 110
106 EKT = TORPWR
107 DO 105 JD=2,NPCINT
110 EN = FLOAT(JD - 1)/10.0
111 Y(JC) = ANORM*EXP(-(EN-EREF)/EKT)/(EN/EREF)
112 105 CONTINUE
114 Y(1)=0.0
115 GO TO 140
116 110 POWER = -TORPWR
117 DO 115 JE=2,NPOINT
120 EN = FLOAT(JE-1)/10.0
121 Y(JE) = ANORM*((EN/EREF)**POWER)
122 115 CONTINUE
124 Y(1)=0.0
125 140 CONTINUE
126 GM = 0.0
127 ASSIGN 145 TO JRETRN
130 JRET = 145
131 GO TO 900

```

FORTRAN SOURCE LIST DETMO

ISN	SOURCE STATEMENT
132	145 CONTINUE
133	GM = GRAMS
134	DO 155 JG = 1, 100
135	Y(JG) = 0.0
136	155 CONTINUE
140	DO 160 JH = 101, NPOINT
141	Y(JH) = Y(JH)*EXP(-GRAMS*ATTEN(JH-100))
142	160 CONTINUE
144	ASSIGN 190 TO JRETRN
145	JRET = 190
146	GO TO 900
147	190 CONTINUE
150	DO 210 JI = 1, NPOINT
151	IF(JI-286.LT.1) GO TO 205
154	Y(JI-286) = Y(JI-286) + ESCKA(JI)*Y(JI)
155	IF(JI-325.LT.1) GO TO 205
160	Y(JI-325) = Y(JI-325) + ESCKB(JI)*Y(JI)
161	205 Y(JI) = Y(JI)*(EFFIC(JI) - ESCKA(JI) - ESCKB(JI))
162	210 CONTINUE
164	235 CONTINUE
165	ASSIGN 240 TO JRETRN
166	JRET = 240
167	GO TO 900
170	240 CONTINUE
171	DO 255 JCT = 1, NPOINT
172	COUNT(JCT) = 0.0
173	255 CONTINUE
175	DO 270 KE = 1, NPOINT
176	E = FLOAT(KE-1)/10.0
177	SIGMA=ASIG*E+BSIG*SQRT(E)
200	QKSIG = 50.0/SIGMA
201	YADJ = 0.1*Y(KE)/(2.5065*SIGMA)
202	DO 265 KEPRIM = 1, NPOINT
203	NGAUS = IFIX(FLOAT(1ABS(KE-KEPRIM))*QKSIG) + 1
204	IF(NGAUS.GT.2000) GO TO 260
207	COUNT(KEPRIM) = COUNT(KEPRIM) + YADJ*GAUSS(NGAUS)
210	GO TO 265
211	260 IF(KEPRIM.GT.KE) GO TO 270
214	265 CONTINUE
216	270 CONTINUE
220	DO 275 JJ = 1, NPOINT
221	Y(JJ) = COUNT(JJ)
222	275 CONTINUE
224	ASSIGN 305 TO JRETRN
225	JRET = 305
226	GO TO 900
227	305 CONTINUE
230	GO TO 50
231	900 IF(NTYPE.EQ.2) GO TO 905
234	WRITE(6, 901) TCRPWR
235	GO TO 906
236	901 FORMAT(1H1, 40HTHERMAL BREMSSTRAHLUNG SOURCE WITH KT = , F7.4, 1 4H KEV)
237	903 FORMAT(1H , 21HAFTER ATTENUATION BY , F7.4, 15H G/CM**2 OF AIR 1 / 1H , 32HINCIDENT SPECTRUM NORMALIZED TO , F7.4,

FORTRAN SOURCE LIST DETMD

ISN	SOURCE STATEMENT
	2 12H PER KEV AT , F7.4, 4H KEV)
240 905	WRITE(6, 904) TCRPWR
241 904	FORMAT(1H1, 48HPOWER LAW DIFFERENTIAL SPECTRUM WITH EXPONENT = 1 , F7.4)
242 906	WRITE(6, 903) GM, ANDRM, EREF
243	IF(JRET.EQ.145.OR.JRET.EQ.190) WRITE(6, 908)
246	IF(JRET.EQ.240) WRITE(6, 910)
251	IF(JRET.EQ.305) WRITE(6, 912)
254	IF(JRET.EQ.305) WRITE(6, 914) ASIG, BSIG
257 908	FORMAT(1H0, 46HBEFORE CORRECTIONS FOR K ESCAPE AND RESOLUTION)
260 910	FORMAT(1H0, 29HAFTER CORRECTION FOR K ESCAPE)
261 912	FORMAT(1H0, 45HAFTER CORRECTIONS FOR K ESCAPE AND RESOLUTION)
262 914	FORMAT(1H ,8HSIGMA = ,F7.5,5H+E + ,F8.5,14H*SQRT(E) KEV)
263	DO 915 KBIN=1,8
264 915	CTSB(KBIN) = 0.0
266	DO 916 JBIN = 1, 40
267 916	CTSA(JBIN) = 0.0
271	MBIN = 1
272	DO 918 KNRG = 1, NPOINT
273	NBIN = (KNRG-1)/50 + 1
274	CTSA(NBIN) = CTSA(NBIN) + 0.1*Y(KNRG)
275	ENERGY = 0.1*FLOAT(KNRG-1)
276	IF(ENERGY.LT.CALIB*THRESH(1)) GO TO 918
301	IF(ENERGY.GE.CALIB*THRESH(MBIN+1)) MBIN = MBIN + 1
304	CTSB(MBIN) = CTSB(MBIN) + 0.1*Y(KNRG)
305 918	CONTINUE
307	WRITE(6, 920)
310 920	FORMAT(1H0, 6HENERGY)
311	LINMX = (NPOINT-1)/100 + 1
312	DO 921 LIN = 1, LINMX
313	LENRGY = 10*(LIN-1)
314	LSTAR1 = 1 + 100*(LIN-1)
315	LSTCP = LSTART + 90
316	WRITE(6, 919) LENRGY, (Y(LA), LA = LSTART, LSTOP, 10)
323 919	FORMAT(1H , 13, 10F10.5)
324 921	CONTINUE
326	WRITE(6, 123) CTSA, CALIB, THRESH, CTSB
327 123	FORMAT(1H0, 20HCCUNTS IN 5 KEV BINS /1H0, 10F10.5 /1H , 1 10F10.5/1H , 10F10.5/1H ,10F10.5/ 1H0, 33HCOUNTS IN BINS WITH T 2HRESHOLDS = , F7.4, 6H TIMES /1H , 9F8.3/1H0,8F8.5)
330	GO TO JRETRN, (145, 190, 240, 305)
331 1000	CONTINUE
332	END

INPUT DATA CARDS FOR PROGRAM DETMD (FOR SAMPLE RESULTS PRESENTED IN THIS REPORT)

190 CARDS CONTAINING EFFIC FROM 10.0 TO 199.9 KEV

190 CARDS CONTAINING ESCKA FROM 10.0 TO 199.9 KEV

190 CARDS CONTAINING ESCKB FROM 10.0 TO 199.9 KEV

190 CARDS CONTAINING ATTEN FROM 10.0 TO 199.9 KEV

111110

1 20.000 1.000 4.30

4.00

0.07 0.29

16.60 23.40 30.30 37.50 46.70 54.20

66.50 81.00 101.00

1.0000

010000

4.50

010000

5.00

010000

5.50

000001

BEFORE CORRECTIONS FOR K ESCAPE AND RESOLUTION

[illegible]

COUNTS IN 5 KEV BINS

[illegible]

COUNTS IN BINS WITH THRESHCLDS =	1.0000 TIMES	
33.000	37.500	46.700
32.000	40.000	54.200
31.000	42.500	66.500
30.000	45.000	81.000
29.000	47.500	101.000

8.09347 1.19006 0.19147 0.03155 0.00294 0.00049 0.00002 0.00000

BEFORE CORRECTIONS FOR K ESCAPE AND RESOLUTION

COUNTS IN 5 KEV BINS

COUNTS IN BINS WITH THRESHOLDS = 1.0000 TIMES

[illegible]

0.32554 0.20907 0.05711 0.01200 0.00129 0.00023 0.00001 0.00000

THERMAL BREMSSTRAHLUNG SOURCE WITH KT = 4.3000 KEV
 AFTER ATTENUATION BY 4.0000 G/CM**2 OF AIR
 INCIDENT SPECTRUM NORMALIZED TO 1.0000 PER KEV AT 20.0000 KEV

AFTER CORRECTIONS FOR K ESCAPE AND RESOLUTION
 SIGMA = 0.07000*E + 0.29000*SQRT(E) KEV

ENERGY	0	0.00001	0.00036	0.00035	0.00035	0.00060	0.00101	0.00112	0.00096	0.00079	0.00076
10	0.00101	0.00182	0.00182	0.00360	0.00676	0.01148	0.01761	0.02458	0.03157	0.03771	0.04233
20	0.04505	0.04584	0.04584	0.04487	0.04251	0.03917	0.03525	0.03108	0.02693	0.02298	0.01936
30	0.01612	0.01330	0.01330	0.01088	0.00884	0.00714	0.00575	0.00462	0.00371	0.00298	0.00239
40	0.00192	0.00155	0.00155	0.00125	0.00101	0.00082	0.00066	0.00053	0.00043	0.00035	0.00028
50	0.00022	0.00018	0.00018	0.00015	0.00012	0.00009	0.00008	0.00006	0.00005	0.00004	0.00003
60	0.00002	0.00002	0.00002	0.00002	0.00001	0.00001	0.00001	0.00001	0.00001	0.00000	0.00000
70	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
80	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
90	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
100	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
110	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
120	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
130	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
140	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
150	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
160	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
170	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
180	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
190	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000

COUNTS IN 5 KEV BINS

0.00207	0.00463	0.03163	0.16655	0.21351	0.12690	0.05147	0.01766	0.00596	0.00204
0.00063	0.00023	0.00008	0.00002	0.00001	0.00000	0.00000	0.00000	0.00000	0.00000
0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000

COUNTS IN BINS WITH THRESHOLDS = 1.0000 TIMES

16.600	23.400	30.300	37.500	46.700	54.200	66.500	81.000	101.000
0.28597	0.19014	0.05672	0.01304	0.00173	0.00037	0.00002	0.00000	0.00000

THERMAL BREMSSTRAHLUNG SOURCE WITH KT = 4.3000 KEV
 AFTER ATTENUATION BY 4.5000 G/CM*2 OF AIR
 INCIDENT SPECTRUM NORMALIZED TO 1.0000 PER KEV AT 20.0000 KEV

AFTER CORRECTIONS FOR K ESCAPE AND RESOLUTION
 SIGMA = 0.07000*E + 0.29000*SQRT(E) KEV

ENERGY	0	10	20	30	40	50	60	70	80	90	100	110	120	130	140	150	160	170	180	190
	0.00001	0.00067	0.00158	0.01329	0.00168	0.00020	0.00002	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
	0.00031	0.00105	0.003294	0.01105	0.00136	0.00016	0.00002	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
	0.00030	0.00199	0.003299	0.00911	0.00110	0.00013	0.00001	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
	0.00052	0.00376	0.003190	0.00746	0.00089	0.00010	0.00001	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
	0.00087	0.01048	0.02738	0.00491	0.00058	0.00007	0.00001	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
	0.00097	0.01517	0.02449	0.00397	0.00047	0.00005	0.00001	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
	0.00068	0.02493	0.01857	0.00258	0.00031	0.00003	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
	0.00060	0.02885	0.01581	0.00208	0.00025	0.00003	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000

COUNTS IN 5 KEV BINS

0.00180	0.00387	0.01812	0.10929	0.15785	0.10138	0.04310	0.01523	0.00524	0.00182
0.00062	0.00021	0.00007	0.00002	0.00001	0.00000	0.00000	0.00000	0.00000	0.00000
0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000

COUNTS IN BINS WITH THRESHOLDS = 1.0000 TIMES

16.600 23.400 30.300 37.500 46.700 54.200 66.500 81.000 101.000

0.20063 0.15068 0.04782 0.01141 0.00154 0.00034 0.00002 0.00000

THERMAL BREMSSTRAHLUNG SOURCE WITH KT = 4.3000 KEV
 AFTER ATTENUATION BY 5.0000 G/CM*2 OF AIR
 INCIDENT SPECTRUM NORMALIZED TO 1.0000 PER KEV AT 20.0000 KEV

AFTER CORRECTIONS FOR K ESCAPE AND RESOLUTION
 SIGMA = 0.07000*E + 0.29000*SQRT(E) KEV

ENERGY	0	0.00001	0.00027	0.00026	0.00026	0.00045	0.00073	0.00085	0.00073	0.00059	0.00049
10	0.00049	0.00066	0.00116	0.00218	0.00389	0.00638	0.01310	0.00954	0.01310	0.01668	0.01986
20	0.00231	0.00383	0.00439	0.00629	0.00954	0.01310	0.01721	0.01935	0.01721	0.01503	0.01292
30	0.01097	0.00920	0.00764	0.00629	0.00515	0.00419	0.00276	0.00341	0.00276	0.00224	0.00181
40	0.00147	0.00119	0.00097	0.00079	0.00064	0.00052	0.00034	0.00042	0.00034	0.00028	0.00022
50	0.00018	0.00015	0.00012	0.00009	0.00008	0.00006	0.00004	0.00005	0.00004	0.00003	0.00003
60	0.00002	0.00002	0.00001	0.00001	0.00001	0.00001	0.00001	0.00001	0.00001	0.00000	0.00000
70	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
80	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
90	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
100	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
110	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
120	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
130	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
140	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
150	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
160	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
170	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
180	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
190	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000

COUNTS IN 5 KEV BINS

0.00156	0.00330	0.01080	0.07279	0.11739	0.08118	0.03612	0.01314	0.00461	0.00162
0.00056	0.00019	0.00006	0.00002	0.00001	0.00000	0.00000	0.00000	0.00000	0.00000
0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000

COUNTS IN BINS WITH THRESHOLDS = 1.0000 TIMES

16.600 23.400 30.300 37.500 46.700 54.200 66.500 81.000 101.000

0.14226 0.11938 0.04034 0.00999 0.00138 0.00030 0.00002 0.00000

AFTER CORRECTIONS FOR K ESCAPE AND RESOLUTION
SIGMA = 0.07000*E + J.29000*SQRT(E) KEV

COUNTS IN 5 KEV BINS

COUNTS IN BINS WITH THRESHOLD =	1.000 TIMES	
16.600 23.400 30.300 37.500	46.700 54.200	66.500 81.000 101.000

-	.10184	.-09485	0.-03467	0.00874	0.00123	0.00028	0.00000
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III. Programs Used to Generate Inputs to DETMD

AIRAT generates a 1900 entry table of attenuation coefficients for air between 10.0 and 199.9 keV. This table is called TABLE within AIRAT and is called ATTEN when used as input to DETMD. The inputs to AIRAT are energies and Compton attenuation coefficients for nitrogen, oxygen, and argon at these energies. The constituents of air correspond to the indices 1, 2, and 3 respectively.

INTERP allows interpolation between three consecutive points on a smooth curve by finding coefficients A, B, and C of the quadratic $Ax^2 + Bx + C$ which fits the three points. It is used by AIRAT to find Compton attenuation coefficients in between the tabulated values it receives as input.

E1 generates values of the exponential integral $E_1(x) = \int_x^\infty e^{-t} \frac{dt}{t}$ using approximation formulae given in Reference 4.

E1 is used by NAIEF in calculating the probability that an isotropically emitted iodine K x-ray escapes from the uniformly thick sheet of NaI.

NAIEF calculates ϵ , p_α , and p_β for x-rays between 10 and 199.9 keV normally incident on a 1 mm. thick NaI crystal. The formulae used are

$$\begin{aligned} \epsilon(E) &= 1 - e^{-\mu x_0} \quad \text{where } x_0 = .367 \text{ g/cm}^2 \text{ and} \\ \mu &= 124 \left(\frac{E}{10} \right)^{-2.7778} \quad \text{for } 10 \leq E \leq 33.23 \text{ keV} \\ \mu &= 30.8 \left(\frac{E}{33.23} \right)^{-2.6596} \quad \text{for } 33.23 < E < 199.9 \text{ keV} \\ p_\alpha(E) &= \int_0^{x_0} dx \mu e^{-\mu x} \left(\frac{1}{2} \delta \omega f_\alpha \left(e^{-\mu_\alpha x} + e^{-\mu_\alpha (x_0 - x)} - \right. \right. \\ &\quad \left. \left. \mu_\alpha x E_1(\mu_\alpha x) - \mu_\alpha (x_0 - x) E_1(\mu_\alpha (x_0 - x)) \right) \right) \end{aligned}$$

$$p_{\beta}(E) = \int_0^{x_0} dx \mu e^{-\mu x} \frac{1}{2} \delta \omega f_{\beta} \left(e^{-\mu_{\beta} x} + e^{-\mu_{\beta} (x_0 - x)} - \right. \\ \left. \mu_{\beta} x E_1(\mu_{\beta} x) - \mu_{\beta} (x_0 - x) E_1(\mu_{\beta} (x_0 - x)) \right)$$

δ = DELTA = fraction of photoelectric processes taking place in the K shell = .875

ω = OMEGA = K fluorescence yield of iodine = .84

μ_{α} = AATTEN = $\mu(E_{K\alpha})$ = 6.62 cm²/g

μ_{β} = BATTEN = $\mu(E_{K\beta})$ = 4.66 cm²/g

f_{α} = ABRANC = fraction of K x-rays which are K_{α} x-rays = .7937

f_{β} = BBRANC = fraction of K x-rays which are K_{β} x-rays = .2063

FORTRAN SOURCE LIST

```

1SN      SOURCE STATEMENT

0 $IBFTC AIRAT
1      DIMENSION ETAB(11), SIGMA(3, 11), EPLOT(1900), TABLE(1900),
1      A(3), B(3), C(3), FRAC(3), CONV(3), POWER(3), COEFF(3)
2      READ(5, 25) ETAB, ((SIGMA(J,K), K = 1, 11), J = 1, 3)
13 25   FORMAT(11F6.1/11F6.2/11F6.2/11F6.2)
14      FRAC(1) = .7556
15      FRAC(2) = .2315
16      FRAC(3) = .0129
17      CONV(1) = .04301
20      CONV(2) = .03765
21      CONV(3) = .01508
22      POWER(1) = -3.28443
23      POWER(2) = -3.30707
24      POWER(3) = -3.071145
25      COEFF(1) = 77.0
26      COEFF(2) = 142.0
27      COEFF(3) = 4230.0
30      DE = 0.1
31      J = 1
32      DO 500 N = 1, 1900
33      EN = N
34      E = (EN - 1.0)*DE + ETAB(1)
35      EPLOT(N) = E
36      IF(E.LT.ETAB(J)) GO TO 35
41      DO 30 K = 1, 3
42      CALL INTERP(ETAB(J), ETAB(J+1), ETAB(J+2), SIGMA(K,J),
1 SIGMA(K, J+1), SIGMA(K, J+2), A(K), B(K), C(K))
43 30   CONTINUE
45      J = J+1
46 35   CONTINUE
47      TABLE(N) = 0.0
50      DO 100 M = 1, 3
51      TABLE(N) = TABLE(N) + FRAC(M)*CONV(M)*((COEFF(M)*(E/ETAB(1))
1 **POWER(M) + A(M)*E**2 + B(M)*E + C(M))
52 100  CONTINUE
54 500  CONTINUE
56      WRITE(6, 510)
57 510  FORMAT(1H1, 55HPHOTON ATTENUATION COEFFICIENTS IN CM**2/GM FOR DRY
1 AIR / 1H0, 6HENERGY )
60      DO 550 NA = 1, 1900, 10
61      IF(MOD(NA, 500).EQ.1.AND.NA.NE.1) WRITE(6, 515)
64      CONTINUE
65 515  FORMAT(1H1)
66      NB = NA + 9
67      WRITE(6, 520) EPLOT(NA), (TABLE(JA), JA = NA, NB)
74 520  FORMAT(1H , F4.0, 10F8.4)
75 550  CONTINUE
77      WRITE(7, 560) (TABLE(JB), JB = 1, 1900)
104 560  FORMAT(10F7.4)
105      CALL EXIT
106      END

```

INPUT DATA CARDS FOR PROGRAM AIRAT (ENERGIES AND COMPTON ATTENUATION
COEFFICIENTS OF NITROGEN, OXYGEN, AND ARGON AT THESE ENERGIES)

10.0	15.0	20.0	30.0	40.0	50.0	60.0	80.0	100.0	150.0	200.0
8.96	6.72	5.73	4.84	4.45	4.14	3.98	3.73	3.54	3.15	2.87
11.50	8.28	6.95	5.77	5.18	4.80	4.61	4.30	4.06	3.61	3.29
56.00	36.00	28.00	19.00	15.80	13.60	12.40	10.80	9.85	8.43	7.57

FORTRAN SOURCE LIST

ISN SOURCE STATEMENT

```

0 $IBFTC NAIEF
1   DIMENSION EXPTAB(2000), EFFIC(1900), ESCKA(1900), ESCKB(1900),
    1 FRACKA(200), FRACKB(200)
2   DO 5 JEX = 1, 2000
3     X = FLOAT(JEX - 1)/200.0
4     EXPTAB(JEX) = EXP(-X)
5   5 CONTINUE
C   ATTENUATION COEFFICIENTS FOR IODINE K ALPHA AND BETA
7     AATTEN = 6.62
10    BATTEN = 4.66
11    DELTA = .875
12    OMEGA = .84
C   RELATIVE NUMBERS OF K ALPHA AND BETA PHOTONS PRODUCED
13    ABRANC = .7937
14    BBRANC = .2063
15    DO 20 JDEPTH = 1, 200
16      DEPTH = .001835*FLOAT(JDEPTH)
C   .367 = THICKNESS OF NAI CRYSTAL CONSIDERED (IN G/CM**2)
17      HEIGHT = .367 - DEPTH
20      XFA = AMAX1(.000001, DEPTH*AATTEN)
21      XFB = AMAX1(.000001, DEPTH*BATTEN)
22      XRA = AMAX1(.000001, HEIGHT*AATTEN)
23      XRB = AMAX1(.000001, HEIGHT*BATTEN)
24      FRACKA(JDEPTH) = 0.5*DELTA*OMEGA*ABRANC*(EXP(-XFA) + EXP(-XRA) -
    1 XFA*E1(XFA) - XRA*E1(XRA))
25      FRACKB(JDEPTH) = 0.5*DELTA*OMEGA*BBRANC*(EXP(-XFB) + EXP(-XRB) -
    1 XFB*E1(XFB) - XRB*E1(XRB))
26  20 CONTINUE
30    DO 100 JENRGY = 1, 1900
31      ENRGY = 10.0 + FLOAT(JENRGY - 1)/10.0
32      EFFIC(JENRGY) = 0.0
33      ESCKA(JENRGY) = 0.0
34      ESCKB(JENRGY) = 0.0
35      IF(ENRGY.GT.33.23) GO TO 25
40      ATTEN = 124.0*(ENRGY/10.0)**(-2.7778)
41      EFFIC(JENRGY) = 1.0 - EXP(-.367*ATTEN)
42      GO TO 100
43  25 CONTINUE
44      ATTEN = 30.8*(ENRGY/33.23)**(-2.6596)
45      DATTEN = .001835*ATTEN
46      DO 80 KDEPTH = 1, 200
47        DEPTH = .001835*FLOAT(KDEPTH)
50        NLAMDA = DEPTH*ATTEN*200.0 + 1.0
51        IF(NLAMDA.GT.2000) GO TO 35
54        FLUX = EXPTAB(NLAMDA)
55        DEFFIC = DATTEN*FLUX
56        EFFIC(JENRGY) = EFFIC(JENRGY) + DEFFIC
57        ESCKA(JENRGY) = ESCKA(JENRGY) + DEFFIC*FRACKA(KDEPTH)
60        ESCKB(JENRGY) = ESCKB(JENRGY) + DEFFIC*FRACKB(KDEPTH)
61  35 CONTINUE
62  80 CONTINUE
64  100 CONTINUE
66      WRITE(6, 105)
67  105 FORMAT(1H1, 67HDEC 6, 1966          PHOTOELECTRIC EFFICIENCY OF 1 MM. T
    1 THICK NAI CRYSTAL          /1H , 75HIFRACTION OF NORMALLY INCIDENT PHOTO

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FORTRAN SOURCE LIST NAIEF

ISN	SOURCE STATEMENT
	2NS THAT ARE ABSORBED PHOTOELECTRICALLY) /1H0, 6HENERGY)
70	DO 120 NLINE = 1, 190
71	IF(NLINE.EQ.51.OR.NLINE.EQ.101.OR.NLINE.EQ.151) WRITE(6,110)
74	CONTINUE
75	110 FORMAT(1H1)
76	NENRGY = NLINE + 9
77	NSTART = (NLINE - 1)*10 + 1
100	NSTOP = NSTART + 9
101	WRITE(6, 115) NENRGY, (EFFIC(NA), NA = NSTART, NSTOP)
106	115 FORMAT(1H , 13, 10F7.4)
107	120 CONTINUE
111	WRITE(7, 121) EFFIC
112	121 FORMAT(10F7.5)
113	WRITE(6, 125)
114	125 FORMAT(1H1, 88HDEC 6, 1966 PROBABILITY THAT A PHOTON INCIDENT INORMALLY ON A 1 MM. THICK NAI CRYSTAL)
115	WRITE(6, 126)
116	126 FORMAT(1H , 79H PRODUCES AN IODINE K ALPHA X-RAY WH 1CH ESCAPES FROM THE CRYSTAL /1H0, 6HENERGY)
117	DO 130 MLINE = 1, 190
120	IF(MLINE.EQ.51.OR.MLINE.EQ.101.OR.MLINE.EQ.151) WRITE(6,110)
123	MENRGY = MLINE + 9
124	MSTART = (MLINE - 1)*10 + 1
125	MSTOP = MSTART + 9
126	WRITE(6, 115) MENRGY, (ESCKA(MA), MA = MSTART, MSTOP)
133	130 CONTINUE
135	WRITE(7, 121) ESCKA
136	WRITE(6, 125)
137	WRITE(6, 136)
140	136 FORMAT(1H , 78H PRODUCES AN IODINE K BETA X-RAY WHI 1CH ESCAPES FROM THE CRYSTAL /1H0, 6HENERGY)
141	DO 140 LLINE = 1, 190
142	IF(LLINE.EQ.51.OR.LLINE.EQ.101.OR.LLINE.EQ.151) WRITE(6,110)
145	LENRGY = LLINE + 9
146	LSTART = (LLINE - 1)*10 + 1
147	LSTOP = LSTART + 9
150	WRITE(6, 115) LENRGY, (ESCKB(LA), LA = LSTART, LSTOP)
155	140 CONTINUE
157	WRITE(7, 121) ESCKB
160	CALL EXIT
161	END

		FORTRAN SOURCE LIST	
ISN		SOURCE STATEMENT	
0		\$IBFTC E1	
1		FUNCTION E1(X)	
2		Y = ABS(X)	
3		X2 = Y*Y	
4		X3 = X2*Y	
5		IF(Y.GT.1.0) GO TO 10	
10		X4 = X2*X2	
11		X5 = X4*Y	
12		E1 = -ALOG(X) - .57721566 + .99999193*Y - .24991055*X2 +	
	1	.05519968*X3 - .00976004*X4 + .00107857*X5	
13		RETURN	
14	10	IF(Y.GT.10.0) GO TO 20	
17		E1 = EXP(-Y)*(X2 + 2.334733*Y + .2506211/(X3 + 3.330657*X2 +	
	1	1.681534*Y)	
20		RETURN	
21	20	E1 = EXP(-Y)*(X2 + 4.03640*Y + 1.15198)/(X3 + 5.03637*X2 +	
	1	4.19160*Y)	
22		RETURN	
23		END	

		FORTRAN SOURCE LIST	
ISN		SOURCE STATEMENT	
0		\$IBFTC INTERP	
1		SUBROUTINE INTERP(X1, X2, X3, Y1, Y2, Y3, A, B, C)	
2		X12 = X1**2	
3		X22 = X2**2	
4		X32 = X3**2	
5		DET = X12*(X2 - X3) - X1*(X22 - X32) + (X3*X22 - X2*X32)	
6		A = ((X2-X3)*Y1 + (X3-X1)*Y2 + (X1-X2)*Y3)/DET	
7		B = ((X32-X22)*Y1 + (X12-X32)*Y2 + (X22-X12)*Y3)/DET	
10		C = ((X3*X22-X2*X32)*Y1 + (X1*X32-X3*X12)*Y2 + (X2*X12-X1*X22)*Y3)	
	1	/DET	
11		RETURN	
12		END	

IV Data Presentation

A logically complete way to present the results of the foregoing analysis is to give the coefficient K found for each energy interval as a function of kT or N or other parameters. This is also a correct way if for some one value of the parameters K is nearly the same for all energy intervals. This is essentially the procedure followed by Clark⁵ in his measurement of the spectrum of the Crab Nebula between 15 and 62 keV. Since then, however, results of x-ray observations from balloons have been presented as energy spectra characteristic of the x-ray sources alone, with the effects of finite energy resolution and K x-ray escape supposedly removed.

Our method of spectrum presentation described below is one which approaches correctness as the range of K values for different energy bands approaches zero for some value of a spectrum parameter such as kT or N . Let us refer to this "best" value of the parameter as $\hat{kT}$. The ordinates in a graph of the x-ray differential number spectrum, $\frac{dn}{dEd\hat{kT}}$, are then taken to be $\frac{N_{ab}(\hat{kT})}{E_b - E_a}$, $\frac{N_{bc}(\hat{kT})}{E_c - E_b}$, etc. where $N_{ab}(kT)$ is as defined in section II of this report. For the abscissa corresponding to an energy band from E_j to E_k we pick the energy E' such that

$$K_{jk} \frac{e^{-E'/\hat{kT}}}{E'} = \frac{N_{jk}(\hat{kT})}{E_k - E_j}$$

$$\text{where } K_{jk} = \frac{AE_0 e^{E_0/\hat{kT}} n_{jk}}{\int_{E_j}^{E_k} \left(\frac{dn}{dE}(\hat{kT}) \right) dE}$$

Thus E' is the energy at which the spectrum equals its average value between E_j and E_k . In terms of K_{jk} the ordinate is

$$\frac{K_{jk}}{E_k - E_j} \int_{E_j}^{E_k} e^{-E/k\hat{T}} \frac{dE}{E}$$

Similarly in a graph of the differential energy spectrum,

$E \frac{dn}{dE dA dt}$, the ordinate is

$$\frac{K_{jk}}{E_k - E_j} \int_{E_j}^{E_k} e^{-E/k\hat{T}} dE$$

and the abscissa is the E' such that

$$K_{jk} e^{-E'/k\hat{T}} = \frac{K_{jk}}{E_k - E_j} \int_{E_j}^{E_k} e^{-E/k\hat{T}} dE$$

Through the point defined by ordinate and abscissa pass a horizontal bar extending from E_j to E_k and a vertical bar whose length represents the error in the determination of N_{jk} , atmospheric depth, and detector properties. The horizontal bar is not to be interpreted as an error bar or as a measure of energy resolution, but just shows the width of the pulse height channel which yielded the data. This is our intent in references 2 and 3 and we believe it is the intent of the other workers in this field. Our vertical error bar does not reflect any error in the process of finding $k\hat{T}$ or finding which parameter, kT , N , or some other, is most appropriate.

If we assume $k\hat{T} = 4.3$ keV and one standard deviation in $n_{ab} = .002$ pulses/cm²sec is .0008 pulses/cm²sec then the ordinate in

a graph of $\frac{dn}{dE dA dt}$ (Figure 1a) is

$$\frac{N_{ab}(\hat{kT})}{E_b - E_a} = \frac{.01666}{23.3} = 7.15 \times 10^{-4} \text{ photons/cm}^2 \text{ sec keV}$$

and the vertical error bar extends from 0.6 to 1.4 times 7.15×10^{-4} . So $K_{ab} = 24.68 \pm 9.87$ and the abscissa, $E' = 30.3$ keV. Note that $\hat{kT} = 4.3$ keV implies that the flux at the middle of the energy band, 35 keV, is very different from that at 30.3 keV. The curved line in Figure 1a is the assumed spectrum at the top of the atmosphere. Figure 2 shows this and the pulse height spectra our detector would produce in case there were 4.0 and 5.5 g/cm² of air along the line of sight. Figure 1B shows the differential energy spectrum $E \frac{dn}{dE dA dt}$. The ordinate is 1.97×10^{-2} keV/cm²sec keV and the abscissa is 30.7 keV.

Of course the one datum considered in this example in no way justifies the choice of $\hat{kT} = 4.3$ keV. The example was taken from an observation of Scorpius X-1. Previous observations of this source by other groups had established the value of $\hat{kT}$. The example was chosen because the steepness of the spectrum necessitates careful presentation of the spectral data and the weakness of the source necessitated consolidation of data over a wide energy range.

References

1. J.A. Stein and W.H.G. Lewin, Journal of Geophysical Research, 72, 383 (1967).
2. J.W. Overbeck, E.A. Womack, and H.D. Tananbaum, Astrophysical Journal, 150, 47 (1967).
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4. M. Abramowitz and I.A. Stegun, editors, Handbook of Mathematical Functions with Formulas, Graphs, and Mathematical Tables, Dover Publications, Inc., New York (1965), page 231.
5. G.W. Clark, Physical Review Letters, 14, 91 (1965).

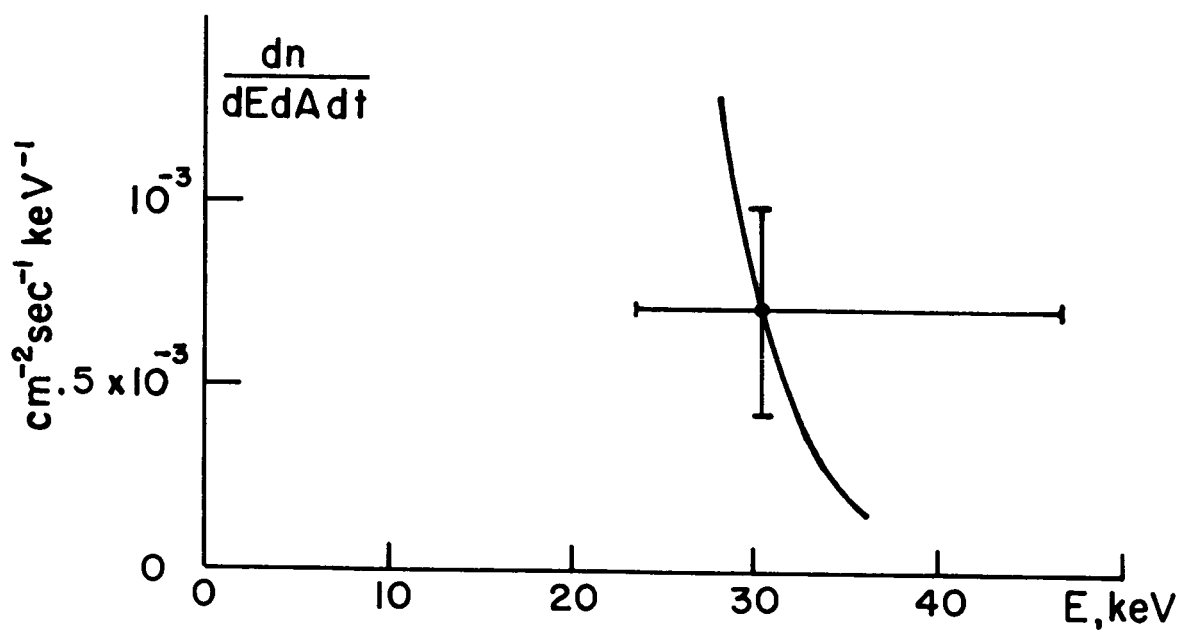


Figure 1a

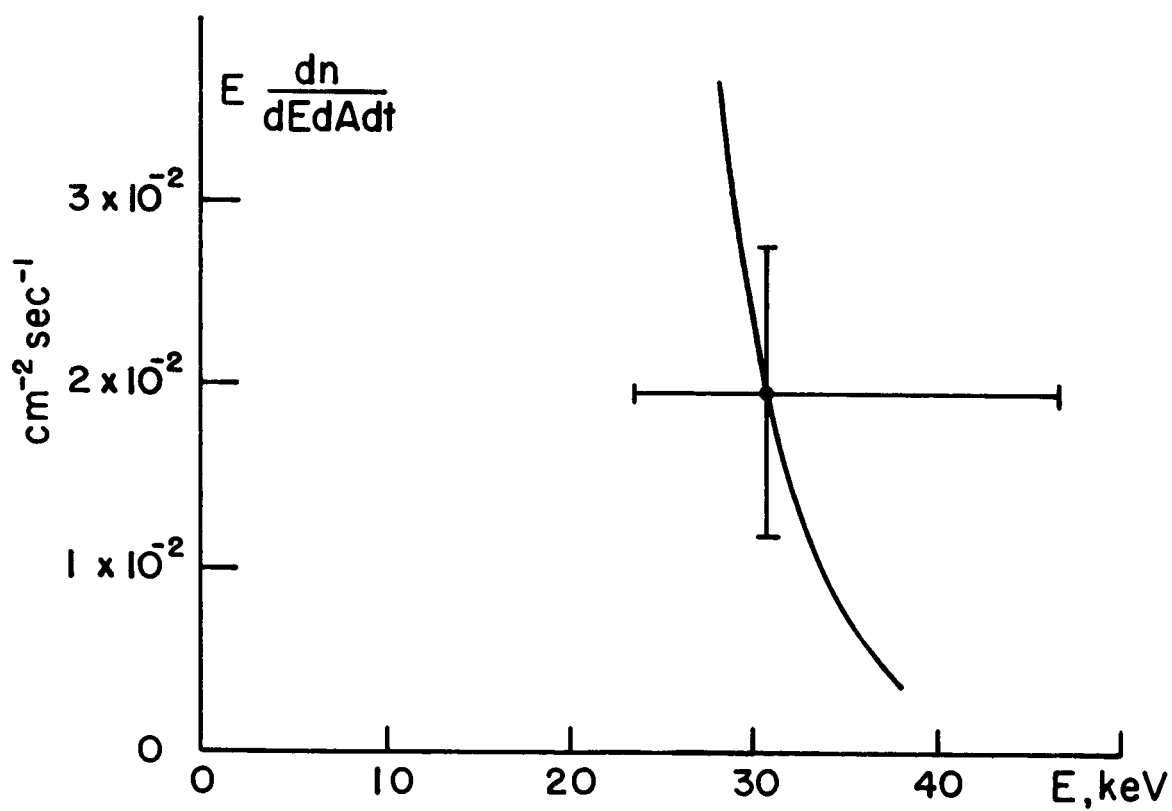


Figure 1b

