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SIMULTANEOUS OPTICAL AND RADIO OBSERVATIONS OF NP 0532*

Simultaneous optical and radio observations of the Crab Nebula
pulsar, NP 0532, were obtained at Lick Observatory and Stanford Uni-
versity respectively, on March 13 and 15. While exact mechanisms
responsible for the radio and optical emissions are still not clear,
the detection of a time interval (if any) between the optical and radio
pulses would be an important test of the various emission theories[1,2] Our
data show, after correction for interstellar dispersion, that the optical
and radio pulses are emitted simultaneousl, with a phase uncertainty of
± 6 milliseconds which results almost entirely from the uncertainty in the
value of the interstellar dispersion. The measured phase delay before cor-
rection is accurate to approximately 250 microseconds.

A block diagram of the experiment is shown in Figure 1. Pulses from
an ITT FW 130 photomultiplier operated at the prime focus of the 120-inch
telescope were counted in a Hewlett-Packard Type 5401A multi-channel ana-
lyzer. Each scan through the 1024-channel memory was initiated by an ex-
ternal trigger signal set at the pulsar repetition rate. The rate of scan
was variable, and was set as fast as 10 μs per channel. The trigger sig-
nal was provided by a frequency synthesizer and scaler; it was possible to
match the apparent period of the pulsar to better than 1 part in 10^7, allow-
ing long integration times without degrading the pulse shape. The trigger
signal was transmitted over a VHF radio link to the 150-foot parabola at
Stanford, 48 km to the west. A second synthesizer at Stanford was adjusted
to the same frequency and phase as the first. Its scaled output was used
to trigger a Hewlett-Packard Type 5480A analog signal averager, which ana-
lyzed 424 MHz radio pulses, operating in much the same way as the multichannel

Pages-9
code 1
CRD1227
AT 30

analyzer at Lick. We estimate that the two trigger signals were syn-
chronized to ± 50 µs over any 1-hour period after compensation for travel
time and other systematic effects.

Figure 2 shows the results of one observation on March 15. The
optical pulses were measured in white light with a 5.5 arcsec diaphragm.
The analyzer had a channel width of 50 µs; the trace is the sum of 2000
scans. Smearing of the pulses due to a slight drift of the trigger signal
is about 60 µs, so that the pulse shapes shown are essentially unaffected
by the equipment. The half-power widths of the main pulse and subpulse are
1.5 and 2.6 ms; the ratio of their peak amplitudes is 3.5.

The radio pulses were observed at 424.00 MHz in a 600 kHz bandwidth.
The tracing in Figure 2 is an average of about 22,000 pulses, indicating
the weakness of the radio signal; the signal-to-noise ratio is low partly
because the Crab Nebula is itself a very strong continuous radio source.
Dispersion across the receiver bandwidth and a 1 ms time constant on the
receiver output have broadened the pulses by about 5 ms, so that little
information on the intrinsic widths is available; the observed widths do
not disagree with the optical widths. The ratio of main pulse to subpulse
peak amplitude, however, is about 1.2 instead of the optical value of 3.5.
Observations at Arecibo[3] give about the same ratio for the radio pulses;
the disagreement with the optical value may be due to polarization effects,
as both the Arecibo and Stanford antennas are circularly polarized.

The positions of the optical and radio pulses in Figure 2 do not rep-
resent the time relationship at the source because the radio pulses have
been delayed considerably by the interstellar medium. A pulse observed at
frequency ν will have been delayed by Δt seconds, where   .

$$\Delta t = \frac{1}{\nu^2} \left(\frac{\int n_e dl}{7.4378 \times 10^{14}} \right)$$

Here ν is in MHz, and the integrated electron density is in electrons/cm^2.
For NP 0532, the integrated density is $(1.755 \pm 0.007) \times 10^{20}$ el/cm^2 (ref. 3);
a pulse at 424.00 MHz is delayed 1312.49 ms, or more than 39 periods. The
optical pulse is delayed about 2×10^{-13} sec. Because the optical and radio
pulses may not come from the same regions there is a possible ambiguity of
an integral number of periods in determining the delay between the optical
and radio pulses. In what follows we will take this number to be zero, thus
making the assumption that the various emission centers are less than 10,000
km apart.

Three successful simultaneous observations of optical and radio pulses
were made, each lasting about 30 minutes. The results are summarized in
Table 1. The line headed Observed Radio Delay gives the apparent time be-
tween the peak of the main optical pulse and the peak of the main radio
pulse; corrections for small phase differences in the trigger signals,
travel times of the trigger and pulsar signals over the 48 km baseline, and
propagation delays in the receiving systems have been made. The probable
errors quoted for the observed delay are due almost entirely to errors in
measuring the position of the weak radio pulse; the error quoted for March
13 is larger due to a lower signal-to-noise ratio at that time. The error
in the calculated delay is due to uncertainty in the integrated electron
density.

A weighted average of the data indicates that the radio pulse arrives
sooner than expected by 0.63 ms, with an experimental error of $\pm$ 0.25 ms
and a systematic error of ± 6.0 ms due to the uncertain electron density.

The difference in arrival times is obviously not significant in view of the total error involved; at present we can only conclude that the radio and optical pulses are emitted within 5 ms of each other. However, when more accurate values for the integrated electron density become available, the observed delay values are sufficiently precise to determine time intervals at the source as small as 0.25 ms. Time delays of this magnitude are important for radio emission mechanisms of the type proposed by Gold[2], in which the observed pulses are produced at a transverse distance from a neutron star where the peripheral speed is near the speed of light. For NP 0532, this transverse distance is 1600 km. The observed time delay of less than 6 ms requires a transverse distance between optical and radio emission centers of less than 1800 km; if the delay proves to be less than 0.25 ms, the distance would be less than 75 km.

Alternatively, we can assume that the optical and radio pulses are emitted simultaneously. We then can derive an accurate value for the integrated electron density from the observed pulse arrival times. The delays in Table 1 give a value of $(1.7542 \pm .0003) \times 10^{20}$ el/cm^2.

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[1]Chiu, H.-Y., Canuto, V., and Fassio-Canuto, L., Nature, 221, 529 (1969).

[2]Gold, T., Nature, 221, 25 (1969).

[3]Comella, J. M., Craft, H. D., Lovelace, R. V. E., Sutton, J. M., and Tyler, G. L., Nature, 221, 453 (1969).

FIGURE CAPTIONS

Figure 1. Block diagram of the simultaneous experiment.

Figure 2. Simultaneous observations on March 15, 2000-2030 PST. (a) 5 minute average of white light optical pulses; (b) 25 minute average of 424 MHz radio pulses. The abscissa is the time after a trigger signal at each site. The dispersion delay and bandwidth broadening of the radio pulses have not been removed. S = subpulse; M = main pulse.

TABLE I - SUMMARY OF OBSERVATIONS

Date	Mar 13	Mar 15	Mar 15
Time (PST)	2000-2020	2000-2030	2100-2130
Apparent Period (ms)	33.09876	33.09885	33.09887
Radio Frequency (MHz)	423.28	424.00	424.00
Observed Radio Delay (ms)	26.13 ±0.60	21.09 ±0.40	20.48 ±0.40
Calculated Delay (ms)	1316.96 ±6.0	1312.49 ±6.0	1312.49 ±6.0
Calculated Delay minus 39 Periods (ms)	26.11 ±6.0	21.63 ±6.0	21.63 ±6.0
Difference (ms)	+0.02 ±6.0	-0.54 ±6.0	-1.15 ±6.0

PHOTOMULTIPLIER AT 120" PRIME FOCUS

PULSAR RECEIVER $f_0 = 424\,MHz$

SIGNAL AVERAGER

150' PARABOLA

SYNTHESIZER $f \approx 3.0212\,MHz$

SCALER $\div 10^5$

OSCILLO-SCOPE

VHF RECEIVER

48 km

STANFORD UNIVERSITY

PULSE SHAPING

MULTI-CHANNEL ANALYZER

PRINTER

VHF TRANSMITTER

SCALER $\div 10^5$

SYNTHESIZER $f \approx 3.0212\,MHz$

LICK OBSERVATORY