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X-615-69-542

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NASA TM X- 63804

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DECEMBER 1969

N70-17393

(THRU)

(ACCESSION NUMBER)

(CODE)

(PAGES)

29

63804

(CATEGORY)

(NASA CR OR TMX OR AD NUMBER)



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**GREENBELT, MARYLAND**

To be submitted to J. Geophys. Res. for publication.



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**SEQUENCE OF DIFFUSE PLASMA RESONANCES  
OBSERVED ON ALOUETTE II IONOGRAMS**

by

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## ABSTRACT

An investigation of more than 1400 Alouette II ionograms indicates that a sequence of diffuse plasma resonances exists and that these resonances are grouped according to the following plasma conditions:

- (1)  $f_{D1}$  is observed between  $1.3 f_H$  and  $2f_H$  when  $2.2 < f_N/f_H < 3.6$ ,
- (2)  $f_{D2}$  is observed between  $2.4 f_H$  and  $3f_H$  when  $3.8 < f_N/f_H < 4.8$ ,
- (3)  $f_{D3}$  is observed between  $3.5 f_H$  and  $4f_H$  when  $4.5 < f_N/f_H < 5.8$ ,
- (4)  $f_{D4}$  is observed between  $4.5 f_H$  and  $5f_H$  when  $5.5 < f_N/f_H < 6.8$ ,

where  $f_{Dn}$  is the frequency of the diffuse plasma resonance,  $f_H$  is the electron cyclotron frequency, and  $f_N$  is the plasma frequency. When these conditions are satisfied, the corresponding  $f_{Dn}$  resonances are observed nearly 100% of the time. Diffuse resonances could not be identified when  $f_N < 1.8f_H$  or when  $f_N > 6.8f_H$ ; the upper limit may be due to insufficient frequency resolution. The diffuse resonance  $f_D$ , reported earlier by other workers as occurring at or near  $1.5f_H$ , corresponds to  $f_{D1}$  of this sequence. A nearly linear increase of  $f_{Dn}$  with increasing  $f_N/f_H$  is observed for  $n=1$  and  $n=2$  with dips near  $f_N/f_H=3$  and  $f_N/f_H=4$ , respectively; the correlation for  $n$  higher than 3 is not as apparent. The present observations cannot be explained in terms of pure Bernstein-mode electrostatic waves.

## INTRODUCTION

Nelms and Lockwood (1966) identified a diffuse resonance, which covers a relatively wide frequency range and shows somewhat similar pattern to a spread echo, in the Alouette II data. They used the notation  $f_D$  for this resonance and found the following approximate relationships:  $f_D \approx 1.5 f_H$  and sometimes  $f_D \approx f_T/2$  or  $f_D \approx f_N/2$ , depending on the local plasma conditions, where  $f_H$  is the electron gyrofrequency,  $f_T$  is the upper hybrid resonance frequency and  $f_N$  is the electron plasma frequency. Under certain plasma conditions, no simple relationship could be identified between  $f_D$  and the other principal resonances (see Calvert and McAfee (1969) for a review of these resonances); no explanation or hypothesis was presented for the observational results.

The purpose of the present paper is to report that a sequence of diffuse resonances (which includes the  $f_D$  resonance mentioned above) is observed on the Alouette II ionograms, and to show that all of these resonances are related in a definite manner to local plasma conditions. The work is based on an analysis of Alouette II data selected from the first 7 months of its life time. The time interval was separated into three periods in order to obtain a suitable variation of the plasma parameter  $f_N/f_H$  at three telemetry stations chosen to cover the latitude range from  $-50^\circ\text{S}$  to  $65^\circ\text{N}$ .

## OBSERVED DATA

The present study was based upon Alouette II data corresponding to an altitude range of 500 km to 3000 km, to a latitude range of  $-50^\circ\text{S}$  to  $65^\circ\text{N}$  (near 1000 km altitude), and to two

seasons (summer and spring). This coverage was obtained by using the data from the following three telemetry stations: (1) Santiago (November 29, 1965 to March 7, 1966), (2) Ottawa (April 5, 1966 to April 26, 1966), and (3) Quito (April 17, 1966 to June 19, 1966).

In Figures 1 (a) - (d), examples of plasma resonances are indicated on the conventional Alouette II ionogram format (Nelms et. al., 1966). The signals of interest are labelled  $f_{D1}$ ,  $f_{D2}$ ,  $f_{D3}$  and  $f_{D4}$ ; the subscript 1 to 4 indicates that each resonance exists between the cyclotron harmonics  $nf_H$  and  $(n+1)f_H$ , where  $n = 1$  to 4. The diffuse plasma resonance  $f_{D1}$  (Figure 1a) corresponds to  $f_D$  which was initially investigated by Nelms and Lockwood (1966). The frequency and time duration of this resonance strongly depends on the plasma parameter  $f_N/f_H$ . The pattern of the  $f_{D1}$  resonance in Figure 1a is typical for the case when  $f_{D1}/f_H = 1.7$ . Under these conditions the resonant time duration is considerably less than for the condition  $f_{D1}/f_H \approx 1.5$  where the duration time is usually from 10 to 20 msec. The  $f_{D2}$  resonance (Figure 1b) has a pattern similar to that of  $f_{D1}$  in that it also has a frequency spectrum broader than the typical spectra of the principal plasma resonances. The  $f_{D3}$  resonance shown in Figure 1c suggests a double peak structure; the  $f_{D2}$  resonance appears on the same ionogram. An example of  $f_{D4}$  is given in Figure 1d. For higher values of  $n$ , the diffuse resonances are usually at frequencies greater than 2.0 MHz, where the frequency resolution is poorer. Also for higher values of  $n$  the spectrum tends to be broader.

The  $f_{Dn}$  resonances are called diffuse resonances in view of their broad frequency spectrum and as stated earlier they occur between  $nf_H$  and  $(n+1)f_H$ . This definition of  $f_{Dn}$  does not include the wide band signals which are seen below the gyro-frequency (Barry, et. al., 1967), and it does not include the broad responses frequently observed near the local Z-mode frequency. The plasma resonances which are observed at frequencies between  $nf_H$  and  $(n+1)f_H$  but occur above the upper hybrid resonance frequency (Warren and Hagg, 1968) are also excluded from this definition.

From a total of 147 satellite passes observed during the above mentioned time intervals, 1477 ionograms were examined for the presence of  $f_{Dn}$  diffuse resonances. On about 900 of these ionograms  $f_N/f_H < 1.8$  and no  $f_{Dn}$  resonance was observed. On the remaining ionograms, where  $f_N/f_H > 1.8$ ,  $f_{Dn}$  diffuse resonances were nearly always observed.

The scaled frequency of each of these diffuse resonances was taken as the center frequency of the diffuse resonance pattern except for very asymmetrical resonances; in this case, the frequency of the maximum duration portion of the resonance was measured. The plasma parameter  $f_N/f_H$  was obtained by scaling the frequencies of the resonances observed at  $f_N$  and  $f_H$  on each ionogram. On a few ionograms, however, where the frequencies of the  $f_N$  or  $f_H$  resonance could not be obtained directly,  $f_H$  was determined from the  $2f_H$  or  $3f_H$  resonance and  $f_N$  was determined from the observed cut-off value  $f_X$  of the extraordinary wave reflection trace.

#### A SEQUENCE OF DIFFUSE RESONANCES $f_{Dn}$

The  $f_{Dn}$  resonances observed on the Quito data are presented

in Figure 2, where  $f_{Dn}/f_H$  is plotted versus  $f_N/f_H$ . This presentation reveals a sequence of  $f_{Dn}$  resonances which consists of four groups defined by the conditions  $n \leq f_{Dn}/f_H < (n+1)$  for  $n = 1$  to 4, with clear gaps of distribution between each group. Similar resonances corresponding to  $n = 5$  could not be definitely identified, but this result is attributed to observational difficulties. Sometimes, the  $f_{Dn}$  resonances are accompanied by subsidiary resonances in each group. These resonant frequencies, which are plotted as open circles and x's in Figures 2 and 3, form different branches which are separated from the main branch (plotted as solid points). Nelms and Lockwood (1966) also reported the occasional occurrence of subsidiary resonances associated with their  $f_D$  resonance. Two subsidiary branches are observed with  $f_{D1}$ , one on the high frequency side of the main branch and one on the low frequency side. These subsidiary branches are not seen simultaneously, and they do not occur over the full range of  $f_N/f_H$  values for which the main  $f_{D1}$  resonance is observed. Only one subsidiary branch is observed with  $f_{D2}$  and it is on the high frequency side of the main branch. Again, the subsidiary branch does not extend over the full range of  $f_N/f_H$  values for which the main branch is observed. The rate of occurrence is higher for  $f_{D2}$  than that is for  $f_{D1}$ . Also, the subsidiary resonances belonging to  $f_{D2}$  have a more irregular pattern than those belonging to  $f_{D1}$  and sometimes it becomes difficult to discriminate them from the main branch resonance. This difficulty of discrimination tends to increase with increasing  $n$  (see for example  $f_{D3}$  in Figure 1c).



Detailed results of the  $f_{D1}$  and  $f_{D2}$  resonance scaling for the combined Quito and Santiago data are shown in Figure 3. When  $f_N/f_H < 1.8$  there is no  $f_{D1}$  resonance. The conditions  $1.8 < f_N/f_H < 2.2$  and  $3.6 < f_N/f_H < 3.8$  define transition regions where the  $f_{D1}$  and  $f_{D2}$  resonances, respectively, are not constantly observed. In the range  $2.2 < f_N/f_H < 3.6$  and  $3.8 < f_N/f_H < 4.8$ ,  $f_{D1}$  and  $f_{D2}$ , respectively, are observed almost 100% of the time independently of the satellite altitude, latitude, local time, and the season. The main branch resonance shows a functional relationship with  $f_N/f_H$ ; the quantity  $f_{Dn}/f_H$  varies almost linearly with  $f_N/f_H$ . There is however, a clear dip around  $f_N/f_H = 3$  for  $f_{D1}$ , and an indication of a dip around  $f_N/f_H = 4$ , for  $f_{D2}$ . The plasma conditions under which the  $f_{D3}$  and  $f_{D4}$  resonances (Figure 2) can be observed nearly 100% are the following:  $4.5 < f_N/f_H < 5.8$  for  $f_{D3}$ , and  $5.5 < f_N/f_H < 6.8$  for  $f_{D4}$ . A correlation between  $f_{Dn}/f_H$  and  $f_N/f_H$  is not apparent for these two groups.

Although the study of Ottawa data was limited to about 70 ionograms, nearly all of these exhibited  $f_{D1}$  resonances clearly split into a double peak structure. These separate peaks, as observed on the Ottawa data, are plotted as open squares and solid dots in Figure 4. There is a remarkable agreement between the average location of the  $f_{D1}$  main resonance from the Santiago and Quito data and the midpoint between the separate peaks of the Ottawa data. This frequency splitting of  $f_{D1}$  is independent

of  $f_N/f_H$ , and appears to be dependent on the satellite latitude.

## DISCUSSIONS

The  $f_{Qn}$  resonances observed at frequencies between the electron cyclotron harmonic values can be explained as the result of electrostatic plasma waves, propagating perpendicular to the magnetic field (the so-called Bernstein-mode waves), which are initiated by the transmitted sounder pulse (Warren and Hagg, 1968). The occurrence of the  $f_{Dn}$  resonance in groups corresponding to  $n = f_{Dn}/f_H = (n+1)$  with clear gaps between each group, and the dependency of  $f_{Dn}$  on  $f_N/f_H$ , suggests a similarity between the  $f_{Dn}$  and  $f_{Qn}$  resonances. The  $f_{Dn}$  resonances, however, differ from the  $f_{Qn}$  resonances in that the  $f_{Qn}$  resonances are observed only when  $f_{Qn} > f_T$  (Warren and Hagg, 1968), whereas the  $f_{Dn}$  resonances are observed only when  $f_{ZS} > f_{Dn}$ , (see Figure 2), which implies  $f_T > f_{Dn}$  since  $f_T > f_{ZS}$ , where  $f_T$  is the upper hybrid resonance frequency and  $f_{ZS}$  is the Z-mode frequency at the satellite level.

From the theoretical point of view, there is a clear difference between the  $f_{Dn}$  and  $f_{Qn}$  resonances in that the dispersion curves of the exact Bernstein-mode waves do not indicate a region of zero group velocity corresponding to the  $f_{Dn}$  resonances, whereas they do for the  $f_{Qn}$  resonances. Thus, if the  $f_{Dn}$  resonances are to be attributed to plasma waves of low group velocity, which remain in the vicinity of the satellite, then plasma waves other than the pure Bernstein-mode waves must be involved.

## SUMMARY

The results of the present data analysis of the  $f_{Dn}$  diffuse

resonances observed by Alouette II are the following:

1. A sequence of diffuse plasma resonances  $f_{Dn}$  which consists of several groups of corresponding  $n$ , is observed. Each group covers a relatively wide frequency range, and the individual groups are clearly separated from each other in frequency.
2. Sometimes, the  $f_{Dn}$  resonances are accompanied by subsidiary resonances which lead to additional branches separated from the main  $f_{Dn}$  branch on the  $f_{Dn}/f_H$  versus  $f_N/f_H$  diagram. Two subsidiary resonances, corresponding to a given  $n$  value, are never observed simultaneously, and the subsidiary resonances occur less frequently than the main branch resonances. There are two subsidiary branches in the  $f_{D1}$  group, one on the high frequency side and the other on the low frequency side of the main branch. In the  $f_{D2}$  group, there is only one subsidiary branch which is on the high frequency side of the main branch.
3. Occasionally, the main  $f_{D1}$  resonance is split into two peaks. This splitting appears to be related to latitude rather than to  $f_N/f_H$ .
4. In the  $f_{D1}$  and  $f_{D2}$  groups,  $f_{Dn}/f_H$  increases almost linearly with  $f_N/f_H$  with an indication of dips near  $f_N/f_H = 3$  for  $f_{D1}$ , and near  $f_N/f_H = 4$  for  $f_{D2}$ .

### **ACKNOWLEDGEMENTS**

The author wishes to express his sincere thanks to Dr. R. F. Benson for suggesting the investigation of the diffuse resonance problem and for his stimulating discussions during this study. The author is also grateful for helpful discussions with Mr. J. E. Jackson and Dr. S. J. Bauer.

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## FIGURE CAPTIONS

**Fig 1.** Ionograms illustrating diffuse plasma resonances:

- (a)  $f_{D1}$  resonance observed at Quito on April 22, 20:40:16 UT, 1966 (1007 km,  $7.0^{\circ}$  S geographic latitude,  $4.68^{\circ}$  N dipole latitude;  $f_N/f_H = 3.17$ ),
- (b)  $f_{D2}$  resonance observed at Santiago on January 28, 18:44:50 UT, 1966 (517 km,  $37.91^{\circ}$  S geographic latitude,  $36.16^{\circ}$  S dipole latitude;  $f_N/f_H = 4.01$ ), (c)  $f_{D3}$  resonance observed at Quito on May 26, 17:57:37 UT, 1966 (606 km,  $20.09^{\circ}$  N geographic latitude,  $31.38^{\circ}$  N dipole latitude;  $f_N/f_H = 5.21$ ), (d)  $f_{D4}$  resonance observed at Quito on May 26, 17:57:05 UT, 1966 (590 km,  $17.99^{\circ}$  N geographic latitude,  $29.27^{\circ}$  N dipole latitude;  $f_N/f_H = 6.21$ ).

**Fig. 2.** Scaling results of the  $f_{Dn}$  resonances for the Quito data. Open circles and x's indicate resonances subsidiary to the main resonance (solid points). The histogram at the top of the figure indicates the number of ionograms scaled vs.  $f_N/f_H$ . The normalized local Z - mode frequency  $f_Z/f_H$  is also indicated.

**Fig. 3.** Scaled  $f_{D1}$  and  $f_{D2}$  resonance frequencies vs.  $f_N/f_H$  as observed on the Quito and Santiago data (the Quito data of Fig. 2 are included in this figure). See caption with Fig. 2.

Fig. 4. The scaled  $f_{D1}$  resonance frequency vs.  $f_N/f_H$  for the Ottawa data. The open squares and solid points represent the two peaks observed on the main resonances; the x's represent subsidiary resonances. The three island-like figures represent the boundary of the  $f_{D1}$  data points from Quito and Santiago as plotted in Fig. 3. A histogram showing the number of ionograms scaled vs.  $f_N/f_H$  is given at the top of the figure.

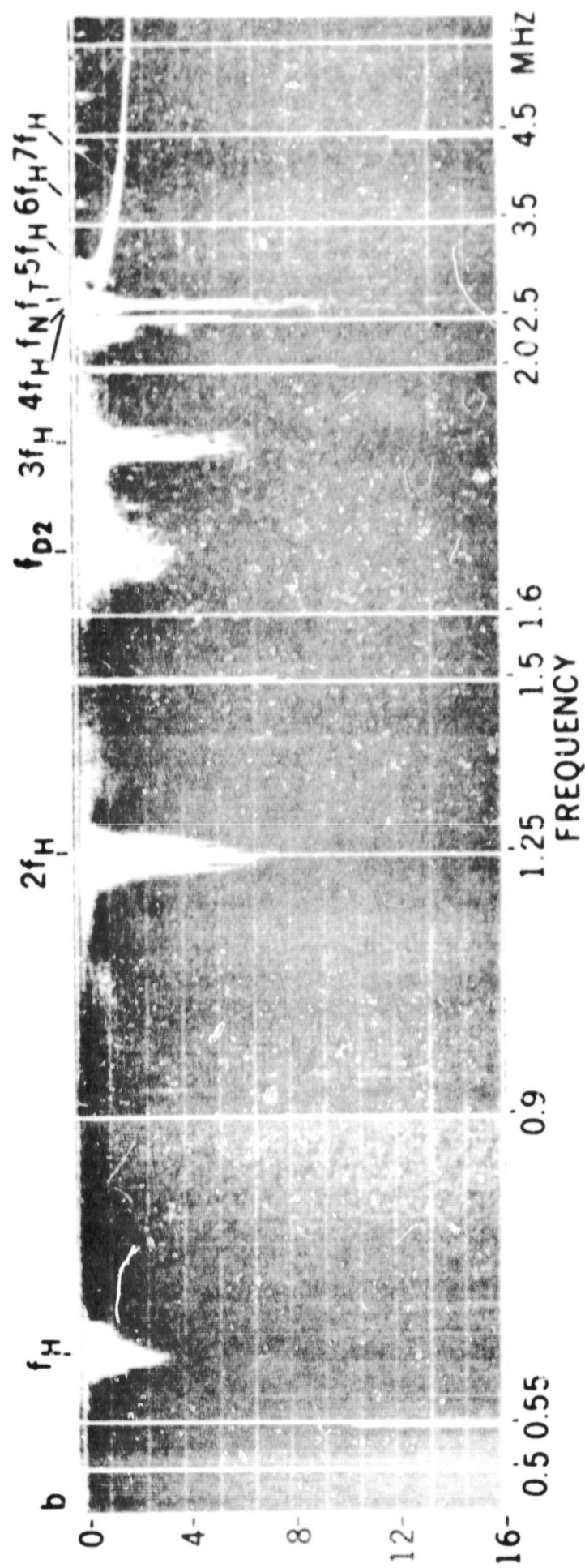
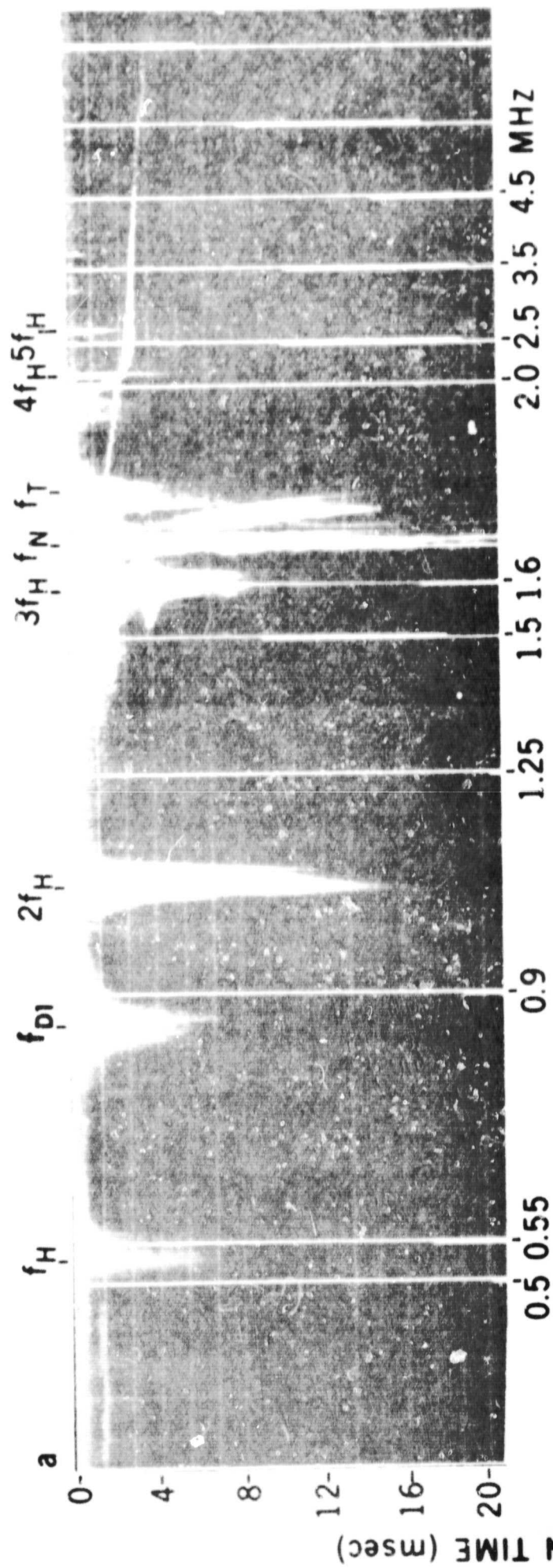


Figure 1



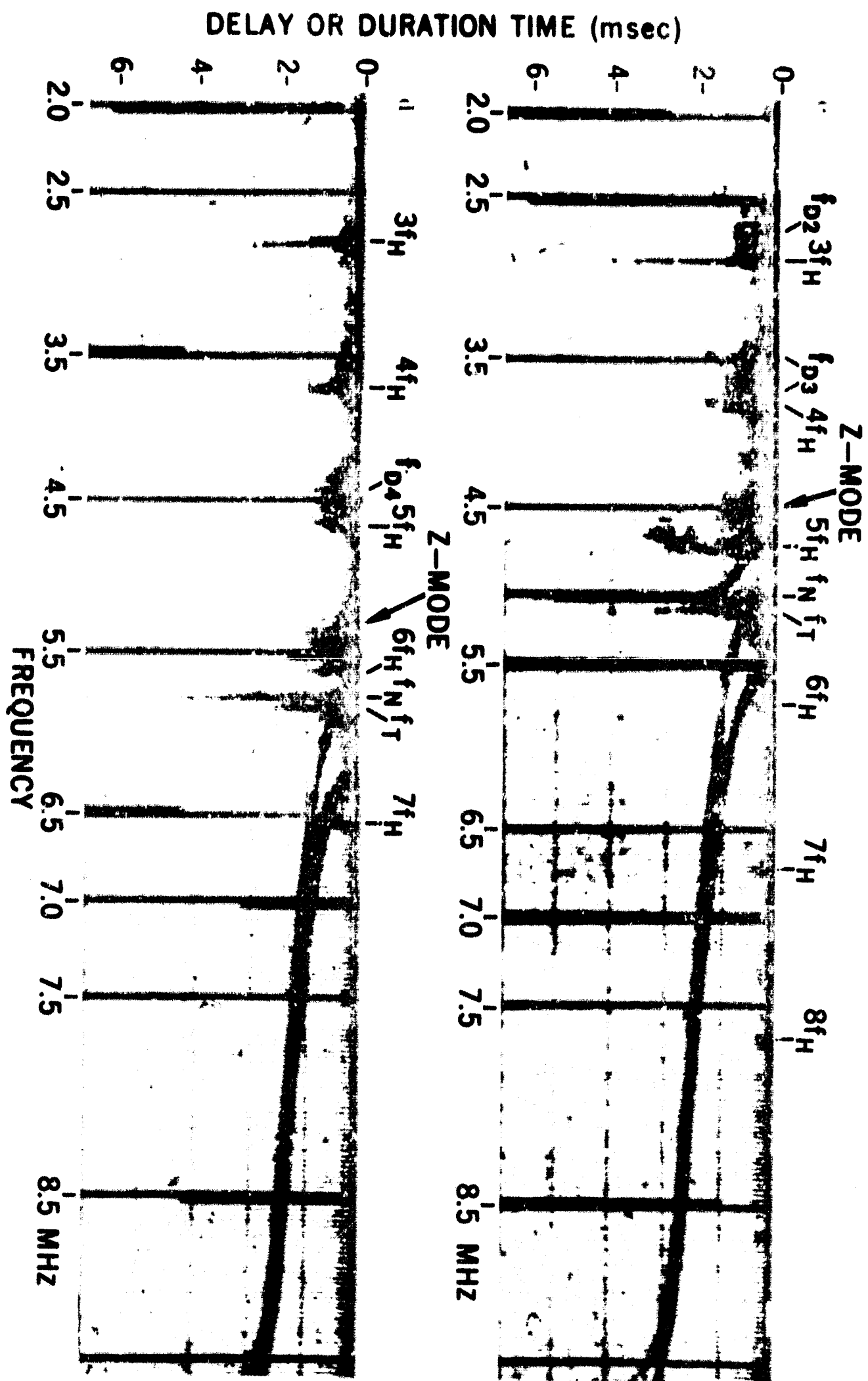


Figure 1A

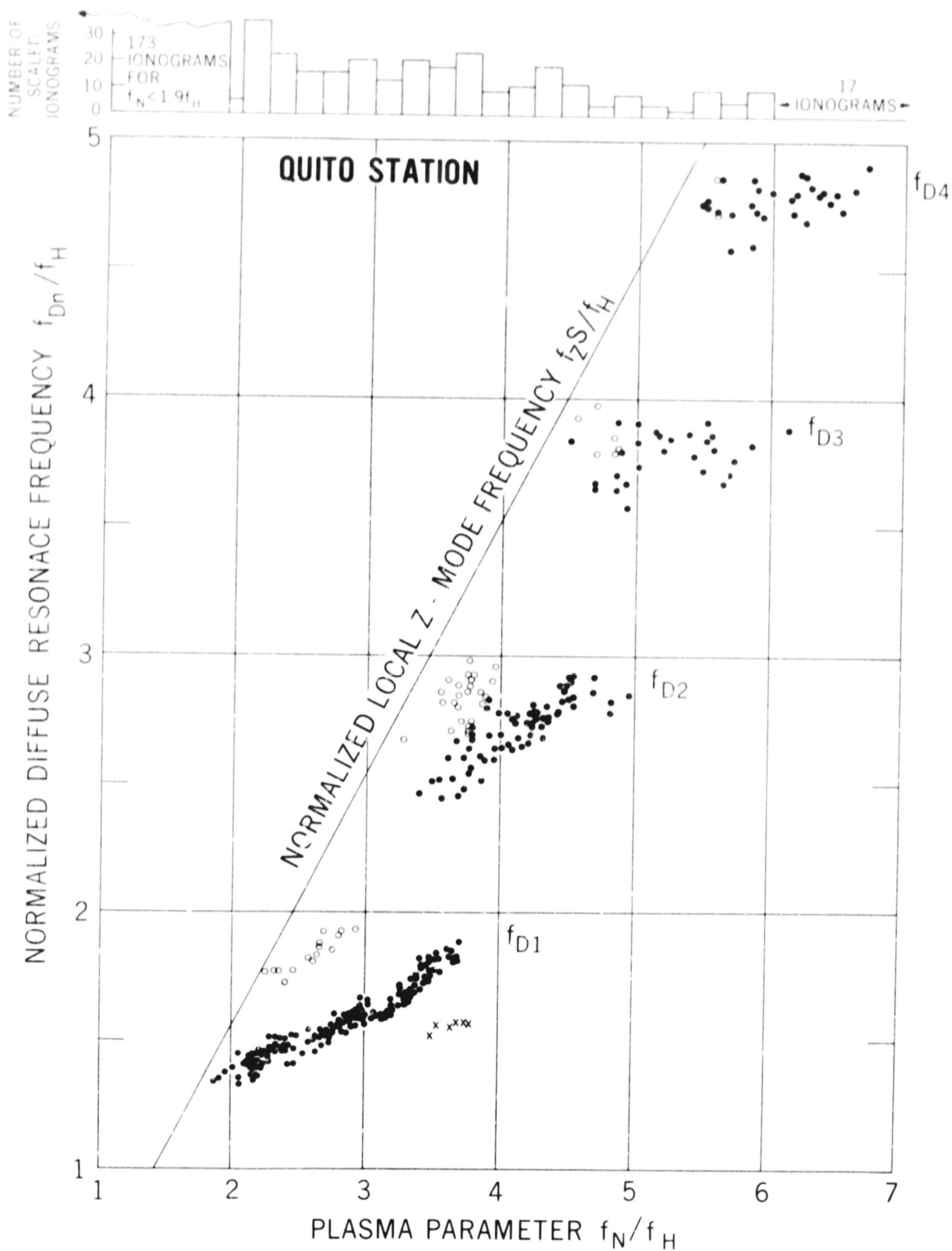


Figure 2

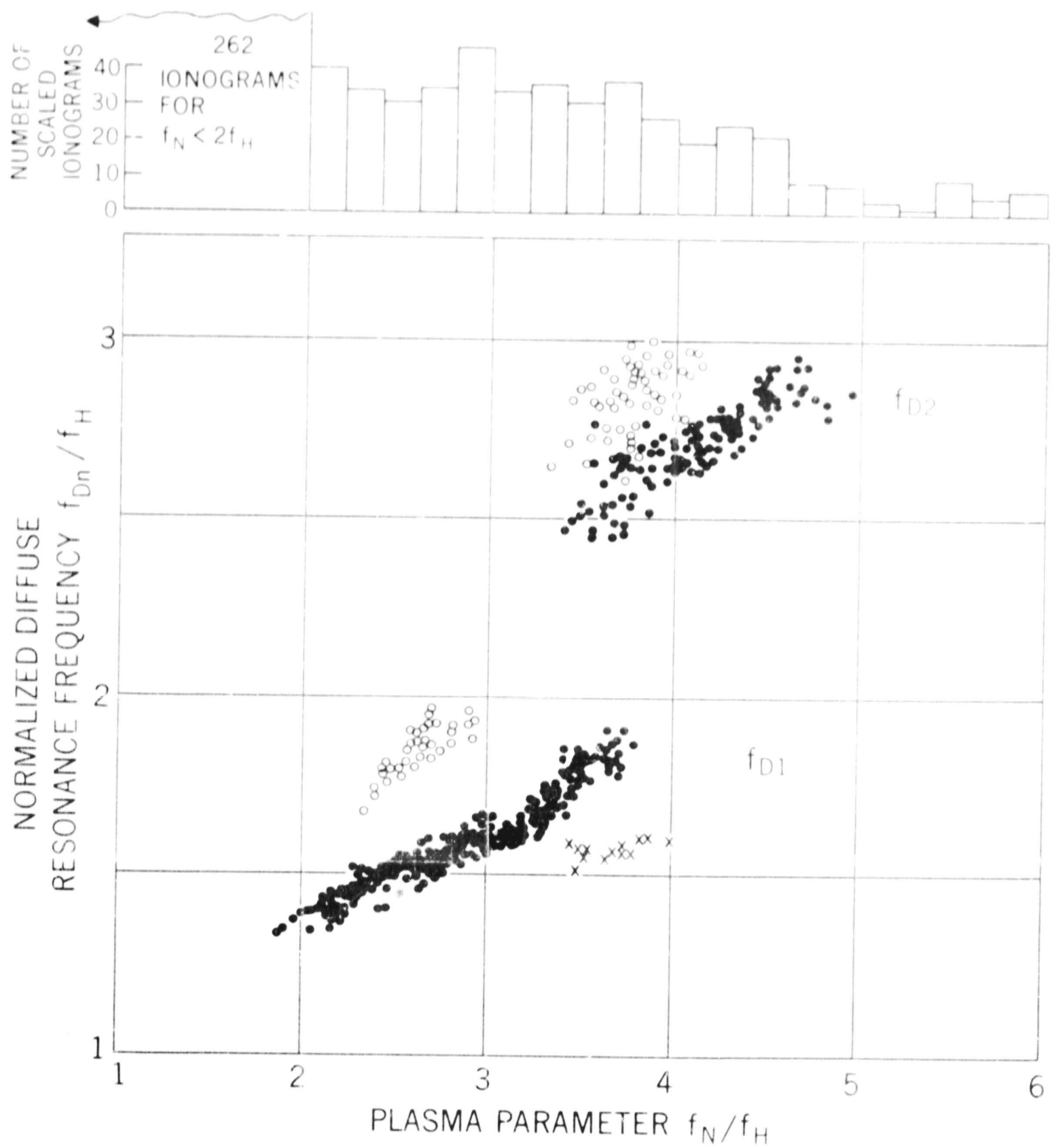


Figure 3

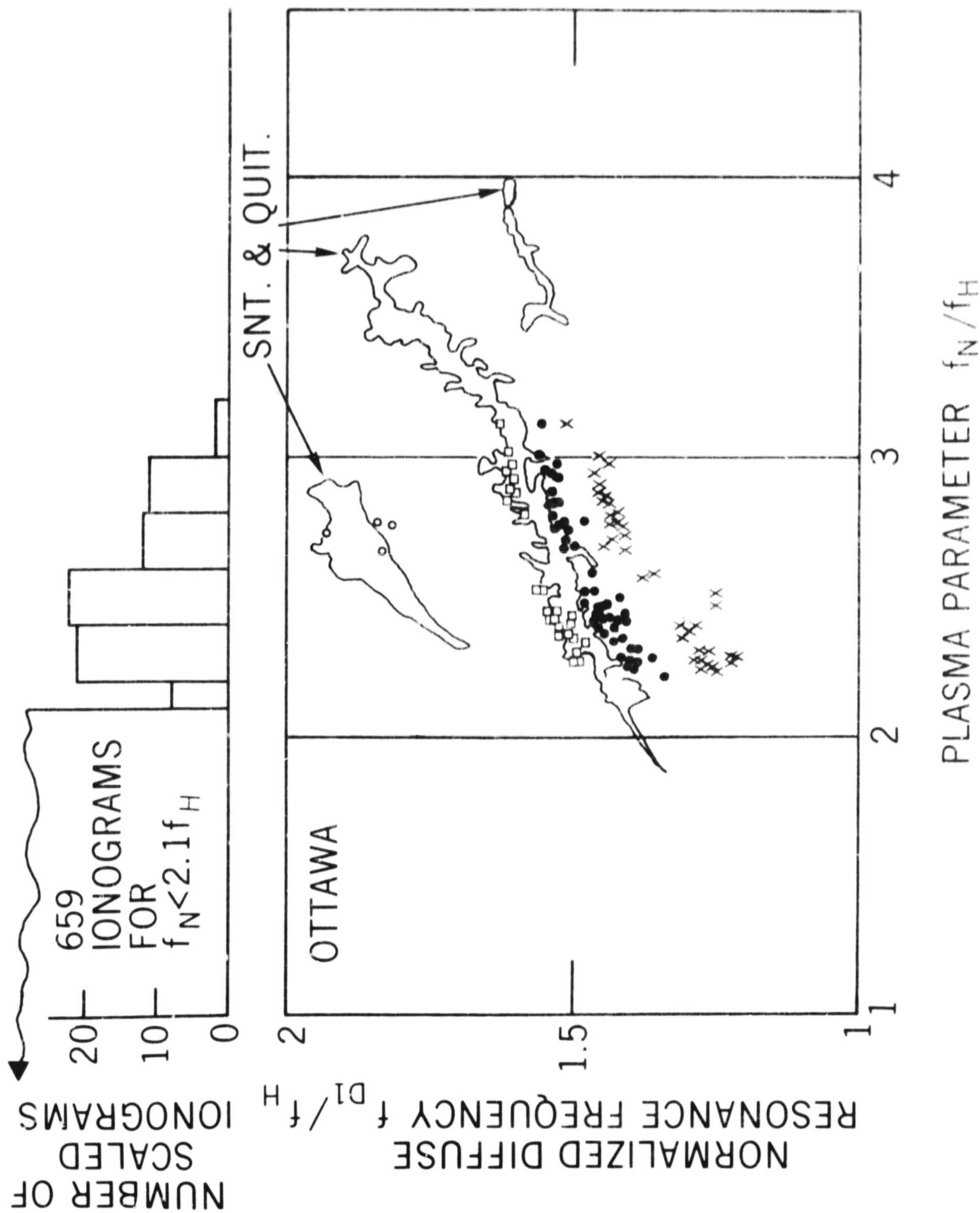


Figure 4