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# RADIO BLACKOUT ALLEVIATION AND PLASMA DIAGNOSTIC RESULTS FROM A 25 000 FOOT PER SECOND BLUNT-BODY REENTRY

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## CONTENTS

I

SUMMARY	1
INTRODUCTION	1
DESCRIPTION OF PAYLOAD	3
General Characteristics	3
Water Injection System	4
X-Band Telemetry System	5
Electrostatic Probe Experiment	5
Thermocouple Probe	6
RESULTS AND DISCUSSION	7
Trajectory and Tracking	7
Signal Blackout	7
X-Band Telemetry	8
Water Injection	8
Electrostatic Probe	9
Thermocouple Probe	11
CONCLUDING REMARKS	1 <b>1</b>
TABLES	13
FIGURES	17
APPENDIX A – DESCRIPTION OF RAM C-I VHF TELEMETRY	
ANTENNA SYSTEMS	45
APPENDIX B – DESIGN AND DEVELOPMENT OF C-BAND AND	
X-BAND RAM C-I ANTENNAS	60
APPENDIX C – INTERPRETATION OF RAM C-I ELECTROSTATIC	
PROBE DATA	81
REFERENCES	89

# RADIO BLACKOUT ALLEVIATION AND PLASMA DIAGNOSTIC RESULTS FROM A 25 000 FOOT PER SECOND BLUNT-BODY REENTRY

I

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#### SUMMARY

On October 19, 1967, RAM C-I was launched from the NASA Wallops Station aboard a four-stage Scout vehicle. The threefold purpose of the flight experiment was to determine the effectiveness of water injection as a communications blackout suppressant on a spacecraft entering the earth's atmosphere at earth-orbital velocities of about 25 000 ft/sec (7620 m/sec), to test the effectiveness of an X-band high-frequency telemetry system in reducing blackout time, and to make electrostatic probe measurements of flow-field characteristics. Water injection reduced the radio-frequency signal attenuation caused by the plasma surrounding the RAM C-I at altitudes between 260 000 and 135 000 ft (79.2 and 41.1 km). The electrostatic probe measurements were made continuously during the portion of the flight from 290 000 to 200 000 ft (88.4 to 61.0 km) and were used to infer electron density distributions in the plasma sheath. The X-band (9210 MHz) telemetry signal was received down to 144 000 ft (43.9 km), approximately 120 000 ft (36.6 km) after VHF blackout. Other transmission frequencies were 225.7 MHz, 259.7 MHz, and 5700 MHz.

#### INTRODUCTION

Disruption of communications with a spacecraft traveling at a high velocity through the earth's atmosphere is due to the plasma (ionized gas) formed about the spacecraft by aerodynamic heating. The dominant chemical kinetic processes in the plasma play an important role in determining the degree of electromagnetic interference and vary as the velocity and bluntness of the reentering spacecraft. For a given body bluntness, three velocity regimes can be considered (ref. 1). The low-velocity regime centers about 18 000 ft/sec (5486 m/sec) and singular molecular ionization is the important chemical kinetic factor. The medium-velocity region centers about 26 000 ft/sec (7925 m/sec) and atomic ionization and charge transfer processes become important. The high-velocity



regime is above 30 000 ft/sec (9144 m/sec) and the atomic ionization, photoionization, and electron-impact ionization processes and multiple-charged ions are all important.

In the Langley Research Center Project RAM (Radio Attenuation Measurements), plasma characteristics and radio interference alleviation techniques have been studied and tested in both ground facilities and flights. Results are documented in part in references 1 to 17; other pertinent information is given in references 18 to 24. Previous RAM flight experiments have included water injection during reentry at 18 000 ft/sec (5486 m/sec) from a hemisphere-cone body (ref. 11) and at 24 000 ft/sec (7315 m/sec) from a Gemini spacecraft (ref. 17).

RAM C-I, a hemisphere-cone probe with a 6-in. (15.24 cm) nose radius, was launched from the NASA Wallops Station on October 19, 1967, aboard a Scout vehicle and had a nominal reentry velocity of 25 000 ft/sec (7620 m/sec). During payload reentry, two techniques for alleviating blackout were tested (water injection and X-band signal transmission) and electrical characteristics of the plasma sheath were measured.

Water was injected into the flow field from the spacecraft as a means of reducing radio signal attenuation during reentry. Water injection into the flow field utilizes the water droplets as recombination surfaces for positive ions and free electrons. Theoretical and experimental background for both the mechanism and mechanics of liquid injection may be found in references 10 to 14.

An X-band (9210 MHz) telemetry system was used as a means of reducing the radio blackout period. This telemetry system transmitted a microwave carrier which was above the critical plasma frequency during a large portion of the reentry. A discussion of this technique may be found in references 1 and 5.

Plasma positive ion density was measured as a function of standoff distance from the spacecraft surface at a fixed body station by using an electrostatic probe fin. From these measurements, electron density concentrations in the flow field of the spacecraft were inferred.

The effectiveness of the two blackout alleviation techniques and the plasma-sheath diagnostic measurements are presented in this report. Also included are a description of the RAM C-I VHF telemetry antenna systems in appendix A by Thomas G. Campbell, a discussion of the design and development of C-band and X-band antennas in appendix B by Fred B. Beck, and an interpretation of RAM C-I electrostatic probe data in appendix C by W. Linwood Jones, Jr.

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#### DESCRIPTION OF PAYLOAD

#### General Characteristics

A four-stage solid-fuel Scout (Scout S-159) was the launch vehicle for the RAM C-I. (See fig. 1.) The payload and the fourth stage were spin stabilized at three revolutions per second just prior to fourth-stage ignition. A protective heat shield for the payload and fourth-stage motor was ejected during coasting flight between second-stage burnout and third-stage ignition.

The payload configuration (fig. 2) had a 6-in.-radius (15.24 cm) hemispherical nose followed by a  $9^{\circ}$  half-angle cone. The overall length of the payload was 51 in. (129.5 cm). The cone base diameter was 24 in. (61.0 cm) and the weight at launch was 267 lb (121.1 kg). The heat protection material covering the hemisphere nose was a charring ablator. The material was selected (ref. 25) in order that the integrity of the water injection nozzles, drilled through the material, would be maintained during ablation. The afterbody was covered with teflon.

The two major systems onboard the RAM C-I were the experiment system and the instrumentation system. The experiment system consisted of (1) water injection mechanism, (2) X-band telemetry, and (3) electrostatic probes. The instrumentation system included various sensors, payload performance instruments, a tape recorder, and two VHF telemeters. Figure 2, a cutaway drawing of the payload, shows the relative location of the various components. All instrumentation performance was nominal and is not discussed in this report.

A perspective sketch of the RAM C-I and coordinates of the important functional parts are given in table I. Indicated in the sketch are the respective locations (three) of the water injection nozzles, the VHF and X-band telemetry antennas, the C-band beacon antenna, and the electrostatic and thermocouple probe fins. One water injection site was in the stagnation region and the other two were located diametrically at the points of tangency between the hemisphere and cone section. Two 259.7-MHz VHF cavitybacked slot antennas, for the real-time telemetry system, were also located diametrically, aft of and in line with the water injection sites. The real-time 259.7-MHz VHF telemetry link used a 5-W FM transmitter. A 225.7-MHz VHF antenna, also excited by a 5-W FM transmitter, was just forward of the payload base. The 225.7-MHz system was also used as a telemetry link transmitting the same data as the 259.7-MHz system but with a 45-sec delay. This delayed telemetry was accomplished by a continuous loop tape recorder in the payload which stored real-time information for 45 sec before transmittal and supplied the data after the payload emerged from blackout. A detailed description of the two VHF telemetry antenna systems is given in appendix A.

The 9210-MHz X-band telemetry antenna system was a four-horn array. The X-band horns were positioned at  $90^{\circ}$  intervals on the periphery of the payload with a  $30^{\circ}$  offset from the slot antennas. The 5700-MHz C-band beacon antenna was located behind the 225.7-MHz VHF antenna. The C-band and X-band antenna systems are discussed in appendix B. The two retractable fins were diametrically positioned at the aft end of the payload and in line with the water injection sites. One fin contained eight electrostatic probes for measuring the plasma ion current. The other fin was equipped with three thermocouples for correlation of the fin leading-edge heating. Both fins were retracted at 176 000 ft (53.6 km) during reentry for payload stability considerations. Detailed descriptions of the two onboard blackout alleviation systems and the flow-field diagnostic experiment are presented in the following sections.

#### Water Injection System

The water injection system was designed to provide a programed variation of flow rates and jet velocities in order to place specific amounts of water throughout part of the radio-frequency (RF) attenuating layer of the flow field. The experiment was designed so that flow rate levels for the water injection system would provide for electron density reductions in the RF attenuating layer of the flow field, which varies with altitude.

The water injection system is shown in block diagram form in figure 3(a). It consists basically of a gas pressurizing system, a water reservoir, 11 solenoid valves, and an electronic programer. Prior to separation of the payload from the fourth stage, the system conditions were as follows: the nitrogen supply bottle was sealed off from the pressure reservoir by the supply valves; the water reservoir, which held 4.2 lb (1.905 kg) of water, was sealed from the gas side by the neoprene bladder and from the other side by the water dump valve; the water lines to the solenoid valves were empty and all of the valves were closed. When the payload had descended to an altitude of 410 000 ft (125.0 km), a command signal from Bermuda initiated the electronic programer action. Ten sec later two solenoid valves were opened to allow the water supply lines to the solenoid valves to be bled down to the ambient pressure, which was essentially vacuum. The valves were then closed and the water dump valve was opened to allow the water to fill the lines. The water pressure at that time was 2 psia  $(13.79 \text{ kN/m}^2)$  due to the initial condition in the pressure reservoir. The cycling of the solenoid valves was then begun with the sequencing shown in figure 3(b). As shown, a cycle of the valves consisted of 7 pulses and the cycle was repetitive every 4 sec. The valve-on times were 230 msec and the valves were opened at 0.5-sec intervals. Valves 5, 6, and 7 (at stagnation injection site) were consecutively opened and allowed increasingly larger amounts of water to be injected into the flow field. Next, the pairs of valves from the side locations (valves 1, 2, 3, and 4) were consecutively opened and also provided increasing flow rates.

Figure 4 shows the orifice configurations for the various flow rates. All orifices were 0.018 in. (0.4577 mm) in diameter and the increase in flow rate was accomplished by increasing the number of orifices used. To increase both the rate of flow and the penetration of the injection into the flow field as a function of decreasing altitude, the two pressure supply valves were opened by the programer, one at 250 000 ft (76.2 km) and the other at 176 000 ft (53.6 km). The gas from the nitrogen supply bottle was thus fed through the valve orifices at a predetermined rate and thereby the head pressure on the water was increased as the payload descended. The resultant flow rates as obtained from flight data are given in table II, with the altitude noted for the start of each pulse. (Cycle 1 had no flow due to a slow fill rate in the lines.)

#### X-Band Telemetry System

A block diagram of the X-band telemetry system which was used to demonstrate the effectiveness of higher transmission frequencies in reducing blackout time, is shown in figure 5. All input data transmitted by this system were also transmitted by the VHF telemetry links. Each of the 60 input channels was commutated 15 times per second by a solid-state electronic commutator-encoder. The outputs from the encoder were pulse position modulated (PPM) to the X-band transmitter and pulse duration modulated (PDM) to the VHF transmitter. The X-band transmitter had a peak power output of 500 W with a pulse width of 1.0  $\mu$ sec and a pulse repetition frequency of 1800 pulses per second. The power was divided and transmitted through a four-horn antenna array (described in appendix B) providing a symmetrical radiation pattern about the payload. Since the payload was spin stabilized during reentry, a symmetrical pattern would allow signalstrength interpretations of RF attenuation effects independent of vehicle spin motions.

#### Electrostatic Probe Experiment

The electrostatic probe experiment was designed to measure the positive ion density profiles through and beyond the flow-field boundary layer of the plasma surrounding the spacecraft as a function of distance normal to the payload surface. The leading-edge dimension of the probe was designed to be free molecular for altitudes above 180 000 ft (54.9 km). The probe body was designed to extend through the peak or maximum plasma density region of the flow field so that an experimental determination of peak density as a function of altitude could be made. Measurements were made at a fixed body station on the aft portion of the spacecraft.

The exterior physical configuration of the electrostatic probe fin, a sectional view of the leading edge, and standoff dimensions of the ion collectors are given in figure 6(a). The probe was essentially a rake consisting of eight individual sensors which were made from iridium strips. The sensors, or ion collectors, extended 0.010 in. (0.254 mm)

beyond the wedge-shaped leading edge. Larger iridium pieces mounted parallel to the leading edge on both sides of the wedge functioned as electron collectors. Iridium was used because of its high melting temperature, high electronic work function, and negligible oxidation property. The 0.010-in. radius of the beryllium oxide leading edge was followed by a  $60^{\circ}$  wedge, and the entire probe configuration was swept back at a  $45^{\circ}$  angle with respect to the payload surface.

Each ion collector was electrically insulated with beryllium oxide, which was used because of its high temperature insulation properties. The two electron collectors on either side of the wedge were electrically common. Referenced to the electron collectors, a fixed bias of -5.0 V was applied to all ion electrodes so that in the presence of a plasma the positive ions were collected. The resultant ion current on each of the eight probes was commutator sampled and then conditioned by a logarithmic amplifier for telemetering. The logarithmic amplifier had a current range from  $10^{-7}$  A to  $10^{-3}$  A; this current range corresponded to an electron density range from  $10^{8}$  electrons/cm<sup>3</sup> to  $10^{12}$  electrons/cm<sup>3</sup>.

When immersed in a flowing plasma, the collectors appear to the charged particles to be 0.010-in.-diameter (0.254 mm) cylindrical wires due to the directed ion velocity; therefore, a modified cylindrical probe theory can be used in the interpretation (ref. 15). The basic theory and a sample calculation of the data reduction procedure for interpretation of the electrostatic probe data are given in appendix C. Although only positive ions are measured by the electrostatic probes, electron densities can be inferred because equal distribution of the electrons and positive ions is assumed.

### Thermocouple Probe

The purpose of the thermocouple probe was to determine the altitude at which insulation degradation for the electrostatic probes began. The thermocouple probe fin was mounted on the payload diametrically opposite the electrostatic probe fin. Both the thermocouple and electrostatic probes were in line with the water injection sites to insure similar heating environments.

The geometry of the thermocouple probe fin shown in figure 6(b) was the same as that of the electrostatic probe fin. Instead of electrodes, however, three thermocouples were imbedded 0.025 in. (0.635 mm) from the leading edge of the wedge as shown in the sectional view of the probe leading edge. The thermocouples were platinum-platinum 13 percent rhodium and were located 0.787, 1.575, and 2.36 in. (2, 4, and 6 cm) from the payload surface. The useful measurement range of the thermocouples was from  $0^{\circ}$  F to  $3100^{\circ}$  F ( $255^{\circ}$  K to  $1978^{\circ}$  K).

#### **RESULTS AND DISCUSSION**

#### **Trajectory and Tracking**

RAM C-I was launched from the NASA Wallops Station at 17:33:00 GMT on October 19, 1967, on an azimuth of  $109^{\circ}$ . The launch vehicle flew a suborbital trajectory and reached an apogee of 750 000 ft (229 km) prior to driving the payload back into the atmosphere at a flight-path angle of  $-15^{\circ}$ . Maximum payload velocity was 25 165 ft/sec (7670 m/sec). Reentry was north of Bermuda, the primary tracking station. Midway through the data period the payload was 100 n. mi. (185.2 km) from the island. The RAM C-I ground track is shown in figure 7.

The variation of altitude and velocity with time for the reentry is shown in figure 8. The indicated events are the command signal from Bermuda which started the programer at 410 000 ft (125.0 km), fourth-stage separation from the payload at 360 000 ft (109.7 km), beginning of water injection at 273 000 ft (83.2 km), retraction of the two probe fins at 176 000 ft (53.6 km), and water cutoff at 110 000 ft (33.5 km). A detailed payload velocity history during reentry from 500 000 ft (152.4 km) down to 20 000 ft (6.10 km) in 5000-ft (1.524 km) altitude increments is given in table III.

Backup tracking was provided by the USNS Twin Falls Victory shown in figure 7. VHF telemetry coverage was provided by Bermuda, the USNS Twin Falls Victory, the USNS Range Recoverer, NASA aircraft 432 and 438, and an Eastern Test Range (ETR) aircraft. The X-band telemetry receiving stations were located at Bermuda.

Continuous trajectory data were obtained at the Bermuda radar installation by the AN/FPS-16 and AN/FPQ-6 radars. Figures 9(a) and 9(b) show portions of the smoothed radar signal-strength data obtained. The FPS-16 radar was in the beacon track mode initially with the FPQ-6 radar operating in the skin track mode. As the payload beacon signal was first attenuated by the reentry plasma at 195 000 ft (59.4 km), the skin track return signal was enhanced. When the FPS-16 radar lost the beacon signal at 180 000 ft (54.9 km), the FPQ-6 radar continued tracking in the skin track mode. The skin return signal remained enhanced until the payload velocity decreased; this decrease caused a reduction of the plasma which, in turn, reduced the skin return signal and allowed the beacon signal to again be received upon emergence from blackout. Detailed C-band beacon track records are given in figure 9(c). A separate study of the tracking radar results has been made and is documented in reference 16.

#### Signal Blackout

Onset and emergence of signal blackout for the RAM C-I reentry are shown in figure 10. The payload velocity at the onset of blackout was about 25 000 ft/sec (7620 m/sec) for all frequencies. The first signal attenuated was the 259.7 MHz VHF

telemetry from the two slot antennas, followed by the 225.7 MHz VHF signal, the 5700 MHz C-band signal, and the 9210 MHz X-band signal. Table IV gives the altitudes for blackout of the various signal frequencies at the different receiving stations. The altitude at onset of attenuation and the altitude at emergence from blackout are also included.

#### **X-Band Telemetry**

The X-band telemetry system provided real-time data for two-thirds of the time period that the VHF telemetry signal was blacked out. The X-band signal did not black out until 144 000 ft (43.9 km) as compared with 260 000 ft (79.2 km) for the VHF signal (fig. 10). The blackout velocity was about the same for both signal frequencies, approximately 25 000 ft/sec (7620 m/sec). These results obtained with the X-band telemetry system show a significant improvement in operational time and altitude range for space-craft reentering the earth's atmosphere.

#### Water Injection

The effect of water injection on the two VHF telemetry signals, 225.7 MHz and 259.7 MHz, is shown in figures 11 to 15. The received signal strength in decibels is plotted above the noise level as a function of altitude for the five VHF telemetry stations. Signal-strength recovery during the prime data period can be compared with the water injection rates given in table II. Injection was from 273 000 ft (83.2 km) down to 111 000 ft (33.8 km) for a range of flow rates from 0.01 lb/sec to 1.55 lb/sec (4.54 g/sec to 703 g/sec). The last effect of water injection on VHF signal was noted at an altitude of approximately 150 000 ft (45.7 km). When the experiment was designed, it was expected that the largest water injection flow rate from the side (flow rate 4) would provide recovery throughout the entire altitude range of injection. Experimentally, however, only marginal recovery was observed for VHF signal down to 150 000 ft (45.7 km). At altitudes above 200 000 ft (61.0 km), water effects are relatively good.

The signal-strength records (figs. 11 to 15) have been faired to partially eliminate antenna pattern and payload spin effects, which are not readily known nor easily analyzed. For comparison, an unfaired telemetry record of the signal strength received at Bermuda is shown in figure 14(b). Recovery spikes at the beginning and end of the injection pulses are sometimes quite large. It is not known whether these spikes are due to optimum water penetration at these particular instances which allows maximum free-electron quenching or are due to a disruption of the flow field caused by the injected water stream. A spiking phenomenon was also observed during the water injection experiment on Gemini 3 (ref. 17).

The X-band recovery levels are shown in figure 16 and water injection effects were observed down to 129 000 ft (39.3 km). The effect of water injection was observed in the radar cross section of the payload and wake as a variation of intensity of the signal received by the skin track C-band radar. These data are evaluated in reference 16.

VHF antenna transmission characteristics are shown in figures 17(a) and (b). Figure 17(a) shows the effect of water injection on the forward and reflected power levels for the 259.7-MHz VHF system, a sharply tuned narrow-bandwidth slot antenna. During injection the reflected power is reduced and the forward power is increased; this indicates that the effects of the plasma on electromagnetic absorption and reflection are being reduced during the presence of water in the flow field. Figure 17(b) shows the forward and reflected power levels for the 225.7-MHz system, a broadly tuned widebandwidth ring antenna. Figure 17(c) shows forward and reflected power levels for the X-band system. An analysis of these power measurements (VSWR) will contribute to an improved understanding of the flow field.

The electrostatic probe data presented in the next section indicate that the attenuating layer of the flow field surrounding RAM C-I extended beyond the ends of the probes. Thus, it is likely that the water was not distributed far enough into the flow field to reduce the outer levels of the electron density so that the propagation of the signal through the flow field was possible.

#### Electrostatic Probe

Approximately 1 sec, corresponding to an altitude change of from 237 000 ft (72.2 km) to 231 000 ft (70.4 km), of the electrostatic probe data is shown in figure 18 where the collected ion current is plotted as a function of altitude. The eight ion collectors are sampled consecutively with one sweep of the eight collectors requiring about 0.026 sec. Thus, an ion distribution profile is obtained in 26 msec, which represents about 8 percent of one revolution of the payload. An in-flight calibration of the electronic system occurs after each three sweeps. The calibration levels are indicated in the figure (100%, 75%, 50%, 25%, 0% of full scale) and correspond to the scale markings on the ordinate axis.

The effects on the current distributions of flow rates 3 and 4 of water injection cycle 3 are also shown in figure 18. Flow rate 3 reduces the ion currents by approximately 1 decade, whereas the larger flow rate 4 reduces the ion currents by approximately 2 decades. The currents returned to their preinjection values very quickly when injection was stopped.

A history of electron density computed from measurements of probes 1, 5, and 8 is presented in figure 19. From a comparison of the curve slopes an insight into electrondensity distribution rates may be gained. A gradual increase in electron density for probes 1 and 5 is noted from 280 000 ft to 267 000 ft (85.3 km to 81.4 km), with a greater increase in density from 267 000 ft to 257 000 ft (81.4 km to 78.3 km). The buildup of electrons between 267 000 ft and 257 000 ft corresponds to the beginning of complete signal loss from the 259.7-MHz forward VHF slot antennas and to the initial attenuation of the 225.7-MHz aft VHF ring antenna. Critical electron densities for the 259.7- and 225.7-MHz frequencies are  $8.2 \times 10^8$  and  $6.3 \times 10^8$  electrons/cm<sup>3</sup>, respectively. Below 257 000 ft, electron density curves derived from all probes indicate only a gradual increase in magnitude. However, from 280 000 ft to 257 000 ft (85.3 km to 78.3 km), probe 8 curve shows an almost constant rate of increase in density and has a greater slope than those for probes 1 and 5; this indicates that peak or maximum electron density moves in an outboard to inboard direction since the standoff distance from the payload surface of probes 1, 5, and 8 is increasing as shown in figure 6(a).

Water injection and spacecraft motions, such as yawing, pitching, spinning, and coning, complicate the analysis of the electron density distribution rates. Figure 19 indicates that the probe density measurements include payload motion effects as evidenced by the null occurring during a period of no injection between cycle 2 and cycle 3 of the water injection (255 000 ft (77.7 km)). This null is certainly not attributable to injection. Likewise, the nulls occurring during injection cycle 2 are most likely due to payload motion effects. These effects are shown by the difference between the deep null during flow rate 2 and a mild depression during flow rate 3, although flow rate 3 is the larger of the two. A 3-rps oscillatory variation of electron density, which is not readily apparent from the figure, is definitely present in the measured density fluctuations. This variation is most likely due to the fact that the payload is at an angle of attack (unsymmetrical flow field about payload) while spinning at a 3-rps rate through the flow field. The probes, which are fixed on the payload, are therefore measuring a varying electron density at the spin rate because of the asymmetry.

During injection cycle 3, the sequential increase in flow rates are readily seen as corresponding decreases in electron density. Distribution profiles during the side injection portion of cycle 3 are shown in figure 20. The curve labeled "No flow" represents the average distribution during the no-injection periods of injection cycle 3. The data points for each profile represent instantaneous interpolations of electron density from plots such as those in figure 19. If a calculation of the attenuation change at VHF is made by applying plane-wave—homogeneous-plasma slab theory in increments to the electron density profile changes in figure 20 for flow rates 2, 3, and 4, there is good agreement between the resultant integral attenuation change and the observed attenuation change of figures 11 to 15.

Effects on electron density during one stagnation sequence of water injection occurring over a 1947-ft (593 m) altitude range is shown in figure 21. Almost a 2-decade

change in electron density due to water injection is observed. Injection began at 226 690 ft (69.1 km) and actually stopped at 225 250 ft (68.7 km).

A family of electron density profiles during no injection over an altitude range from 280 000 to 185 000 ft (85.3 to 56.4 km) is shown in figure 22. The profile at an altitude of 280 000 ft corresponds to the beginning of the electrostatic probe data period. (At altitudes of 195 000 ft (59.4 km) and 185 000 ft (56.4 km), the outermost collectors have saturated the measuring system.) The profiles inferred from the electrostatic probe measurements indicate a much thicker boundary layer than that used as the design criterion. From the general slope of the profiles, the outer edge of the boundary layer appears to be beyond a standoff distance of 7 cm.

#### Thermocouple Probe

A knowledge of the aerodynamic heating on the probe is necessary for valid probe current interpretation. Thermocouple probe data for the RAM C-I are presented in figure 23. Data from thermocouples 7, 8, and 9 show a small initial increase in heating. Thereafter data from thermocouple 9 indicate a sharp rise in heating commencing at about 230 000 ft (70.1 km) and the heating rate of thermocouple 8 changes abruptly at about 204 000 ft (62.2 km). Electrostatic probe data are considered to be useful only for comparative purposes and not for absolute values once the local probe temperature exceeds  $1000^{\circ}$  F (811<sup>o</sup> K) due to probable thermal degradation of the beryllium oxide insulator. Finally, at 189 500 ft (57.8 km), data from thermocouple 7 also indicate an abrupt rise in heating rate. Effects of water injection can also be seen in figure 23, particularly for altitudes below 210 000 ft (64.0 km).

#### CONCLUDING REMARKS

A reentry flight experiment has been conducted with the RAM C-I at about 25 000 ft/sec (7620 m/sec) to determine the reduction in radio-frequency signal attenuation due to injection of water into the plasma sheath surrounding the RAM C-I and to measure plasma properties by using electrostatic probes. From the data obtained, the following remarks can be made:

1. Water injection into the plasma surrounding the payload reduced the attenuation of signals radiated from the RAM C-I at frequencies of 225.7 MHz, 259.7 MHz, 5700 MHz, and 9210 MHz.

Effects of water injection were observed on signal-strength measurements down to 150 000 ft (45.7 km) for the VHF frequencies and to 130 000 ft (39.6 km) for the X-band frequency.

Signal recovery levels due to water injection did not meet the design estimates apparently because the attenuating portion of the flow field was thicker than expected, as evidenced by the electrostatic probe measurements at altitudes above 200 000 ft (61.0 km).

Both VHF and X-band VSWR measurements show plasma effects on the antennas, and a change in the effects during water injection is observed.

2. The X-band telemetry system provided real-time data during two-thirds of the time period that the VHF signal was blacked out at a reentry velocity of about 25 000 ft/sec (7620 m/sec) and thereby confirmed the worth of high-frequency transmission as a means of reducing the time interval of communications blackout.

The X-band signal was received to a 115 000-ft (35.0 km) lower altitude than were the VHF signals. The blackout altitudes were 260 000 ft (79.2 km) for VHF and 144 000 ft (43.9 km) for X-band frequency.

3. Electrostatic probes were effective in measuring the plasma properties.

Preliminary data reduction indicated that the region of high electron density in the flow field was thicker than expected.

During the period of valid electrostatic probe measurements, good agreement was obtained between the measured levels of electron density and the VHF signal attenuation data, both with and without water injections.

The effects of vehicular motion on the plasma density levels were also observed.

#### Langley Research Center,

National Aeronautics and Space Administration, Langley Station, Hampton, Va., October 28, 1969. TABLE I.- POSITION OF FUNCTIONAL PARTS ON THE RAM C-I CONFIGURATION



	Part	x,	φ,	x/D
Number	Function	in. (cm)	deg (a)	(b)
	Water injection nozzle:			
1	Stagnation	0 (0)		0
2	Lateral	6.0 (15.2)	0	.48
3	Lateral	6.0 (15.2)	180	.48
	Antenna:	•		
4	VHF slot (259.7 MHz)	29.6 (75.2)	0	2,36
5	VHF slot (259.7 MHz)	29.6 (75.2)	180	2.36
6	X-band horn (9210 MHz)	32.6 (82.8)	60	2.60
7	X-band horn (9210 MHz)	32.6 (82.8)	150	2.60
8	X-band horn (9210 MHz)	32.6 (82.8)	240	2.60
9	X-band horn (9210 MHz)	32.6 (82.8)	330	2.60
10	VHF ring (225.7 MHz)	43.0 (109.2)		3.42
11	C-band horn (5700 MHz)	46.6 (118.4)	30	3.71
	Probe:			
12	Electrostatic	48.6 (123.4)	0	3.87
13	Thermocouple	48.6 (123.4)	180	3.87

 $^{a}$ Center-line location of parts.

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 $^{b}$ Nose diameter D is 12.56 in. (31.90 cm).

TABLE II RAM C-I	WATER	INJECTION	FLOW	RATES
------------------	-------	-----------	------	-------

	Valve	for -	Flow	v rate	Altit	uđe		Valve	for -	Flo	w rate	Altit	ude	
Cycle	Stagnation injection	Side injection	lb/sec	kg/sec	ft	km	Cycle	Stagnation injection	Side injection	lb/sec	kg/sec	ft	km	
1	5	[	0	0	[	[	5	5		0.05	0.023	200 490	61.1	1
	6		0	0				6		.15	.068	197 236	60.1	1
	7		0	0			1	7		.28	.127	193 982	59.1	1
		1	0	0			lí –		1	.10	.045	190 724	58.1	1
		2	0	0	1				2	.24	.109	187 462	57.1	1   -
)		3	0	0					3	.65	.295	184 202	56.1	ι]
		4	0	0					4	1.01	.458	180 943	55.2	3
2	5		0	0	278 908	85.0	6	5		0.06	0.027	174 451	53.2	2
	6		0	0	275 639	84.0		6		.18	.082	171 212	52.2	2
	7		.03	.014	273 027	83.2		7		.36	.163	167 968	51.2	:
		1	.01	.004	269 115	82.0			1	.13	.059	164 722	50.2	:
		2	.04	.018	265 866	81.0			2	.32	.145	161 466	49.2	
		3	.10	.045	262 616	80.0			3	.90	.408	158 199	48.2	4
		4	.15	.068	259 363	79.0			4	1.42	.644	154 937	47.2	
3	5		0.02	0.009	252 861	77.1	7	5	ĺ	0.07	0.032	148 484	45.2	Í
	6		.07	.032	249 598	76.1		6		.23	.104	145 267	44.3	
	7	ļ	.14	.064	246 328	75.1		7		.43	.195	142 058	43.3	
		1	.06	.027	243 048	74.1			1	.15	.068	138 864	42.3	
		2	.14	.064	239 771	73.1			2	.35	.159	135 684	41.4	
		3	.40	.181	236 486	72.1	]		3	.98	.444	132 511	40.4	
		4	.66	.299	233 205	71.1			4	1.55	.703	129 532	39.5	
4	5		0.04	0.018	226 690	69.1	8	5	l	0.08	0.036	123 137	37.5	
	6		.12	.054	223 442	68.1		6	ĺ	.24	.109	120 058	36.6	
	7		.24	.109	220 185	67.1		7		.43	.195	117 010	35.7	
		1	.08	.036	216 915	66.1	[	[	1	.15	.068	113 987	34.7	
(	ĺ	2	.21	.095	213 635	65.1			2	.35	.159	110 986	33.8	
		3	.55	.249	210 348	64.1			3	0	0			ĺ
		4	.89	.404	207 054	63.1			4	0	0			

## TABLE III.- ALTITUDE, VELOCITY, AND TIME OF RAM C-I REENTRY TRAJECTORY

Altitude		Velocity		Time from lift-off.		Altitude		Velocity		Time from lift-off.
ft	km	ft/sec	m/sec	sec		ft	km	ft/sec	m/sec	sec
$500 \times 10^{3}$	152.4	23 004	7012	355.98		$255  imes 10^3$	77.7	25 121	7657	393.77
495	150.9	23 279	7095	356.83		250	76.2	25 135	7661	394.54
490	149.4	23 567	7183	357.64		245	74.7	25 164	7670	395.30
485	147.8	23 875	7277	358.42		240	73.2	25 143	7664	396.07
480	146.3	24 166	7366	359,19		235	71.6	25 131	7660	396.83
475	144.8	24 476	7460	359.98		230	70.1	25 139	7662	397.59
470	143.2	24 696	7527	360.77		225	68.6	25 148	7665	398.36
465	141.7	24 816	7564	361.57		220	67.0	25 165	7670	399.13
460	140.2	24 872	7580	362.34		215	65.5	25 156	7668	399.89
455	138.7	24 866	7579	363.11		210	64.0	25 145	7664	400.65
450	137.2	24 863	7578	363.87		205	62.5	25 123	7657	401.41
445	135.6	24 872	7580	364.64		200	61.0	25 102	7651	402.18
440	134.1	24 892	7587	365.40		195	59.4	25 115	7655	402.94
435	132.6	24 902	7590	366.18		190	57.9	25 108	7653	403.71
430	131.1	24 920	7596	366.95		185	56.4	25 076	7643	404.48
425	129.5	24 926	7597	367.73		180	54.9	25 064	7640	405.25
420	128.0	24 916	7594	368.49		175	53.3	25 034	7630	406.02
415	126.5	24 914	7594	369.25		170	51.8	25 009	7623	406.79
410	125.0	24 924	7597	370.00		165	50.3	24 967	7610	407.56
405	123.4	24 941	7602	370.75		160	48.8	24 923	7596	408.32
400	121.9	24 965	7609	371.52		155	47.2	24 829	7568	409.09
395	120.4	24 982	7614	372.30		150	45.7	24 752	7544	409.86
390	118.9	24 971	7611	373.08		145	44.2	24 663	7517	410.64
385	117.3	24 967	7610	373.85		140	42.7	24 532	7477	411.42
380	115.8	24 969	7610	374.61		135	41.4	24 384	7432	412.20
375	114.3	24 973	7612	375.37		130	39.6	24 157	7363	412.99
370	112.8	24 986	7616	376.14		125	38.1	23 882	7279	413.80
365	111.2	24 991	7617	376,92		120	30.0	23 563	7182	414.61
300	109.7	24 990	7019	377.70		110	30.0	23 202	2020	410.43
250	106.2	25 000	7696	270.92		105	22.0	22 103 99 176	6750	410.20
345	105.2	25 019	7620	370.00		100	30.5	22 170 91 455	6520	419.00
340	103.6	25 023	7627	380.75		95	29.0	20 446	6232	418 91
335	102.1	25 031	7629	381 52		90	27.4	19 483	5938	419.87
330	100.6	25 037	7631	382 29		85	25.9	18 226	5555	420.87
325	99.1	25 026	7628	383.07		80	24.4	16 697	5089	421.96
320	97.5	25 025	7628	383.86		75	22.9	14 935	4552	423.16
315	96.0	25 042	7633	384.63		70	21.3	12 942	3945	424.51
310	94.5	25 065	7640	385.39		65	19.8	10 624	3238	426.12
305	93.0	25 098	7650	386.14		60	18.3	8 000	2438	428.22
300	91.4	25 094	7649	386.88		55	16.8	5 706	1739	430.88
295	89.9	25 087	7646	387.64		50	15.2	3 441	1049	434.69
290	88.4	25 074	7642	388.40		45	13.7	1 906	581	440.26
285	86.9	25 082	7645	389.17		40	12.2	890	271	448.98
280	85.3	25 095	7649	389,93		35	10.7	721	220	459.34
275	83.8	25 103	7651	390.70		30	9.1	785	239	467.88
270	82.3	25 094	7649	391.46		25	7.6	744	227	475.40
265	80.8	25 096	7649	392.23	ĺ	20	6.1	907	276	484.45
260	79.2	25 113	7654	393.00						

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Station	Signal frequency,	Altitude at of atte	t beginning nuation	Altit signal	ude at blackout	Altitude at emergence from blackout		
		ft	km	ft	km	ft	km	
Bermuda	225.7	267 000	81.4	257 000	78.3	77 000	23.5	
	259.7	280 000	85.3	262 000	79.9	74 000	22.6	
	5700	195 000	59.4	181 000	55.2	75 000	22.9	
	9210	182 000	55.5	144 000	43.9	72 000	21.9	
NASA air-	225.7	274 000	83.5	256 000	78.0	87 000	26.5	
craft 438	259.7	273 000	83.2	263 000	80.2	79 000	24.1	
USNS Range	225.7	279 000	85.0	260 000	79.2	78 000	23.8	
Recoverer	259.7	280 000	85.3	267 000	81.4	70 000	21.3	

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## TABLE IV.- SIGNAL BLACKOUT HISTORY



Figure L- RAM C-I on launch pad.

L-67-220



Figure 2.- RAM C-I configuration showing experiment and instrumentation systems.



(b) Water injection cycle.

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Figure 3.- Block diagram of the RAM C-I water injection system and a typical water injection cycle.

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(a) Sectional view of typical nozzle insert with single orifice.



(b) Stagnation injection orifice layout.





Figure 4.- Orifice configurations for the water injection system.



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Figure 5.- Block diagram of X-band telemetry system.



(b) Thermocouple probe configuration.

Figure 6.- Body configurations of the electrostatic and thermocouple probes. All linear dimensions are in inches (centimeters).



Figure 7.- RAM C-I ground track with relative locations of tracking ships and aircraft.

![](_page_27_Figure_0.jpeg)

Figure 8.- RAM C-I reentry trajectory.

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![](_page_28_Figure_0.jpeg)

![](_page_28_Figure_1.jpeg)

![](_page_28_Figure_2.jpeg)

(b) AN/FPQ-6 C-band radar. (Automatic skin tracking; smoothed AGC data.)

Figure 9.- Strength of C-band radar signals received at Bermuda.

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![](_page_29_Figure_0.jpeg)

(c) Detailed beacon signal strength.

Figure 9.- Concluded.

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![](_page_30_Figure_0.jpeg)

Figure 10.- RAM C-I altitude-velocity reentry profile showing onset and end of RF signal blackout.

![](_page_31_Figure_0.jpeg)

Figure 11.- RAM C-I water injection effects on VHF telemetry signal strength received aboard the USNS Twin Falls Victory.

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![](_page_32_Figure_0.jpeg)

Figure 12.- RAM C-I water injection effects on VHF telemetry signal strength received aboard NASA aircraft 438.

![](_page_33_Figure_0.jpeg)

Figure 13.- RAM C-I water injection effects on VHF telemetry signal strength received aboard ETR aircraft number 1.

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![](_page_34_Figure_0.jpeg)

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![](_page_34_Figure_1.jpeg)

Figure 14.- Records of VHF telemetry signal strengths received at Bermuda during RAM C-I water injection period.

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![](_page_35_Figure_0.jpeg)

(b) Unfaired data. (LCP, left circular polarization; RCP, right circular polarization.)

Figure 14.- Concluded.


Figure 15.- RAM C-I water injection effects on VHF telemetry signal strength received aboard USNS Range Recoverer.



(a) Faired pulse-by-pulse data.

Figure 16.- Records of X-band telemetry signal strengths received at Bermuda during RAM C-I water injection period.



(b) Faired peak-pulse data.

Figure 16.- Concluded.



(a) 259.7-MHz system.

Figure 17.- Forward and reflected power of VHF and X-band telemetry systems as a function of altitude and water injection.



(b) 225.7-MHz system.

Figure 17.- Continued.



(c) 9210-MHz system.

Figure 17.- Concluded.



Figure 18.- Effects of water injection on electrostatic probe current measurements during injection cycle 3.



Figure 19.- Electron density of the RAM C-I flow field at standoff distances of 1, 4.5, and 7 cm as inferred from electrostatic probes 1, 5, and 8, respectively.



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Figure 20.- Electron density distribution profiles for various injection flow rates of cycle 3.



Figure 21.- Effect of stagnation injection at flow rate 5 of cycle 4 on electron density distribution profiles at various altitudes.



Figure 22.- Electron density distribution profiles at various altitudes. No water injection.



Figure 23.- Variation of thermocouple probe leading-edge temperature with altitude.

# DESCRIPTION OF RAM C-I VHF TELEMETRY ANTENNA SYSTEMS

By Thomas G. Campbell

# Introduction

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The RAM C-I had two VHF telemetry systems for the transmission of "direct" time and "delayed" time data. The antennas designed for these telemetry systems consisted of two diametrically opposed cavity-backed slot antennas for the transmission of directtime data and an array of circumferential slot antennas that transmitted the delayed-time data during the playback portion of the flight. This appendix presents a brief discussion of the design and development of these antenna systems and the final performance characteristics of each system.

The design objectives for the delayed-time telemetry antenna system were as follows:

(1) Optimize the pattern and radiation efficiency for the playback portion of the trajectory

(2) Provide broadband impedance characteristics to minimize plasma-sheath nearfield effects

(3) Reduce the possibility of electromagnetic interference when low-level components are used

The design objectives for the direct-time telemetry antenna system were as follows:

(1) Use thin rectangular slot antennas for comparison purposes (This type of antenna had been used on previous RAM experiments.)

(2) Optimize the radiation efficiency while providing broadside radiation from  $30^{\circ}$  to  $140^{\circ}$  off the payload nose (Typical antenna gain should be  $\geq -5$  dB for this electrically small antenna.)

(3) Reduce the possibility of electromagnetic interference when low-level components are used

# Delayed-Time Telemetry Antenna

A circumferential slot array was designed for the RAM C-I delayed-time telemetry antenna at a frequency f of 225.7 MHz to achieve end-fire (off the nose and tail)

radiation pattern characteristics. These characteristics were to optimize tracking and data acquisition during reentry and the playback phase of the trajectory.

The design of circumferential slot arrays has been discussed extensively in the past with respect to achieving omnidirectional pattern characteristics in the roll plane, but achieving end-fire or broadside radiation with polarization versatility is another problem. For this particular payload application, end-fire radiation was desired and obtained; however, the array could very easily have been altered to provide broadside radiation instead.

The antenna design was obtained by first testing two diametrically opposed circumferential slots on the coated conical configuration shown in figure A1. The slots were located in the aft portion of the nose cone to optimize the electrical length of each slot and to maximize the diametric spacing so that coupling would be reduced. The body circumference ka (where  $k = \frac{2\pi}{\lambda}$ ,  $\lambda$  is wavelength in free space, and a is radius of nose cone) at the selected location is  $1.34\lambda$  and the diameter d is  $0.425\lambda$ . The two diametrically opposed slots were therefore spaced  $0.425\lambda$  apart and the pattern control obtained agrees with the theory pertaining to an array of linear elements spaced the same distance apart. The two slots were fed in parallel with a half-wave-line balun which forced the potential at the feed point of each slot to be equal and opposite. A balanced-to-unbalanced impedance ratio of 4 is obtained and a 50-ohm impedance match was obtained at the paralleled input by adjusting the feed-point locations for each slot. The total slot length could also be adjusted by relocating shorting bolts across two circular rings inside the payload. The resulting slot lengths L were  $0.334\lambda$ .

It is interesting to note that these two slots mounted on the cone provided radiation pattern control similar to that of an array of two linear  $\frac{1}{2}\lambda$  elements with a spacing of approximately  $\frac{1}{2}\lambda$  (ref. 26, pp. 279-300 and 471). The pattern results for the two slots on the cone are shown in figure A2(a) and A2(b). In figure A2(a) the slots are fed out of phase, that is, the center conductor of the feed line for one slot is connected to the forward side of the slot and the center conductor for the second slot is connected to the aft side of the slot. In figure A2(b) the pattern is for end-fire radiation and this pattern is obtained by merely reversing the feed polarity of the second slot. Therefore, the center conductor for each slot was connected to the forward side of the slot.

Since end-fire radiation was desired for this particular payload, two additional slots were added to the array to optimize the polarization of the end-fire radiation. The additional slots were also diametrically opposed and fed with a half-wave-line balun. The resulting pattern is essentially the superposition of the fields from the two pairs of slots and the measured results are shown in figure A2(c). The capability for end-fire circular polarization is evident in the pattern amplitudes of the orthogonal linear polar-ization. The sense of the circular polarization was easily altered by adjusting the phase

angle  $\phi$  between slot pairs and the results are presented in figure A3. For  $\phi = 0^{\circ}$ , the time-phase angle by which the vertical component of the electric field leads the horizontal component of the electric field is  $90^{\circ}$  and, when the nose cone is viewed off the nose, the sense of the circular polarization is left-hand circular or a clockwise wave approaching. As the payload is viewed from the rear, the sense of the circular polarization is right-hand circular or a clockwise wave receding. As  $\phi$  is increased, the sense is completely reversed with right-hand circular off the nose and left-hand circular off the tail. The same polarization versatility could be achieved for broadside radiation by changing the polarity of the slots accordingly. The patterns in figure A3 are also shown for several roll angles  $\Phi$ , and no severe pattern changes were observed.

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The final radiation characteristics were somewhat different from those shown in figure A3 due to a change in the polarization capability of the main tracking ship. Therefore, the phase variation between slot pairs was changed to produce a more linearly polarized signal off the nose and tail of the payload. The final impedance, VSWR, and typical radiation pattern characteristics of the delayed-time telemetry antenna are shown in figures A4(a), A4(b), and A5, respectively.

# Direct-Time Telemetry Antenna

Two diametrically opposed, axial, cavity-backed slot antennas were used to transmit the direct-time telemetry data during the RAM C-I flight experiment. Slot antennas were used on previous RAM experiments and, therefore, they were used on this RAM payload for comparative purposes. Two antennas were used (locations are shown in fig. A1 and in sketch in table I) so that signal variation due to payload spin could be minimized. In order to prevent corona, the slot apertures were pressurized instead of being encapsulated. The slots were pressurized by placing a thin phenolic cover over the apertures. The cover contained a valve that would close when the ambient pressure dropped below 3.5 psi (24.132 kN/m<sup>2</sup>).

Since an encapsulating dielectric is usually lossy, pressurizing the slots increased the radiation efficiency. Pressurizing the slots also simplified the task of tuning the cavity-backed slot antennas because there was no need to compensate for an encapsulating dielectric. In figure A1 can be seen the pressurizing cover over the slot aperture.

The free-space impedance and voltage standing wave characteristics of the directtime telemetry antennas are shown in figure A6. As expected, these antennas are narrow bandwidth and experienced a near-field detuning effect when they were surrounded by the plasma sheath.

Typical radiation pattern characteristics of the direct-time telemetry antenna are shown in figure A7. On the basis of these radiation pattern characteristics, it was concluded that this antenna system would satisfy the experiment objective.



Figure A1.- RAM C-I configuration showing experiment and instrumentation systems.

APPENDIX A

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Figure A2.- Radiation pattern diversity capabilities of the circumferential slot array. Elevation plane.



(c) Four slots fed in phase.

Figure A2.- Concluded.

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Figure A3.- Polarization diversity capabilities of the end-fire slot array. Elevation plane.



Figure A3.- Continued.



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Figure A3.- Concluded.



(a) Impedance.

Figure A4.- Final performance characteristics of delayed-time telemetry antenna (free space).

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(b) Voltage standing wave ratio.

Figure A4.- Concluded.



Figure A5.- Equatorial plane radiation patterns of delayed-time telemetry antenna (free space). Isotropic level at the 6-dB circle;  $\theta = 90^{\circ}$ ; f = 225.7 MHz.

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(a) Impedance.

Figure A6.- Final performance characteristics of direct-time telemetry antenna (free space).



(b) Voltage standing wave ratio.

Figure A6.- Concluded.

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Figure A7.- Equatorial plane radiation pattern of direct-time telemetry antenna (free space).  $\theta = 90^{\circ}$ ; f = 259.7 MHz; peak gain, -4 dB.

# DESIGN AND DEVELOPMENT OF C-BAND AND X-BAND RAM C-I ANTENNAS

By Fred B. Beck

# Introduction

The major design considerations for the RAM C-I C-band and X-band antennas are physical size, pattern coverage, and gain. The limitations on physical size are surface area on the circumference of the payload and internal volume. Since the antenna surface area primarily determines the pattern coverage, the antenna apertures were designed to be a compromise of these two parameters. With the antenna aperture dimensions defined, the resultant pattern coverage and gain are primarily determined by leaky surface wave effects (ref. 27) from the ablation coating and efficiency of aperture excitation, respectively.

# Preliminary Antenna Design

The first step in the preliminary design procedure for the RAM C-I C-band and X-band antennas was to build a metal mock-up of the payload to be used in conjunction with various antenna test models. The next step was to determine the available payload space for mounting the antennas and their far-field pattern requirements. The final inputs needed for the preliminary design were trajectory information and antenna polar-ization requirements. The important trajectory parameter at this stage of the design was the aspect angle since this angle dictates the orientation of the far-field patterns. The aspect angle is formed by a line through the nose-to-tail axis of the payload and the line of sight from a receiving station. For the RAM C-I flight, the primary downrange receiving station is Bermuda.

Different test antenna models were designed within the preliminary guidelines of space and mounted in the RAM C-I metal mock-up in specified locations. Horn antennas were chosen rather than some other aperture-type antenna because of their simplicity and because they had been used on previous RAM flights. Since pattern coverage at X-band frequency was required over the entire 360° in the roll plane, an array of four horns was chosen. However, for C-band frequency, only one horn was required for continuous smoothed radar tracking.

After the preliminary design decision to use horns for both the C-band and X-band antenna systems was made, far-field radiation patterns were then measured to determine whether the coverage was compatible with anticipated flight conditions. Examination of these preliminary patterns for both circumferential and axial polarization clearly

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demonstrated that circumferential polarization would not be acceptable for the four-horn X-band array. The reason for the unacceptability of the X-band circumferential polarization was the occurrence of deep nulls around the entire 360<sup>o</sup> of the roll pattern. Although either vertical or horizontal polarization gave a satisfactory C-band pattern, the downrange C-band tracking radar could receive only vertical polarization, which dictated that the C-band payload polarization be circumferential.

Once the preliminary X-band and C-band antenna design, including the choice of polarization, was completed, effort was then channeled toward refinements of the basic design. In order to make these refinements, trajectory information and a payload mock-up including the teflon heat shield were required. By using the mock-up with heat shield, more meaningful pattern and standing wave ratio measurements were made and incorporated into a final design.

# **Trajectory Parameters**

The reentry trajectory parameters used in the final design of the C-band and X-band antennas are given in figure B1. The variation of altitude, velocity, elevation angle, aspect angle, and heat-shield ablation rate with time from lift-off in seconds is given in figure B2. From the standpoint of antenna design, the two most important of these parameters are aspect angle and heat-shield ablation rate. The aspect angle of the trajectory dictates the orientation of the antennas with respect to the spacecraft axis. Since the aspect angle is approximately  $90^{\circ}$  as the payload passes Bermuda, the antennas should have radiation patterns favoring the side rather than the nose or tail of the payload. The heat-shield ablation rate has an important effect on the matching of the antennas and on radiation patterns due to trapped energy in the ablative material.

# Antenna Voltage Standing Wave Ratio

A major contribution to a high VSWR on flight vehicles is frequently the impedance mismatch caused by the dielectric heat shield covering the antennas. The method commonly used to decrease the VSWR of horn antennas is to insert a proper size dielectric plug into the throat of the aperture. However, since the protective heat shield ablates during reentry, some detuning effects are still expected to occur. Therefore, the dielectric plug should be designed to allow some detuning before and after the ablation period. The heat-shield thickness on the RAM C-I before ablation is approximately 0.31 in. (0.79 cm) at the location of the C-band and X-band antennas. After ablation, the thickness is calculated to be approximately 0.11 in. (0.28 cm); this indicates that 0.2 in. (0.51 cm) of the heat shield will ablate during reentry.

The E-plane and H-plane views of the C-band horn antenna and the location of the teflon matching plug are shown in figure B3. The matched condition for the horn occurs

for a teflon thickness equal to a multiple of dielectric half-wavelengths. Since 0.2 in. (0.51 cm) of the heat shield ablates during reentry, the plug plus heat shield and bonding agent should be designed 0.1 in. (0.25 cm) greater than the matched condition.

The variation of the VSWR with frequency of the C-band horn antenna on the RAM C-I is presented in figure B4. Note that the VSWR over the frequency range from 5.40 GHz to 5.8 GHz is always less than the maximum desirable preflight value of 2 to 1. Since the horn has approximately 0.1 in. (0.25 cm) more teflon covering than is required for a matched condition, the horn is expected to become matched when 0.1 in. of the heat shield has ablated.

In designing the X-band matching plugs extra care was deemed necessary because 0.2 in. (0.51 cm) of teflon approaches  $\frac{1}{4}\lambda$  at the design frequency of 9.21 GHz. This thickness is a sufficient amount to mismatch the horns if the plugs are not designed properly. A drawing of the flight X-band horn antenna with matching plug is shown in figure B5.

The variation of VSWR with frequency of the four-horn-array X-band antenna on the RAM C-I is shown in figure B6. The total thickness of the matching plug, heat shield, and bonding agent on the flight payload is approximately 0.94 in. (2.39 cm). This thickness is approximately 0.06 in. (0.15 cm) greater than  $1\lambda$  in the dielectric. Since the X-band horns are designed primarily for the reentry data period and the total 0.2-in. (0.51 cm) heat-shield ablation does not all occur during this time, less overdesign for the matching plugs of the X-band horn is required than for the matching plugs of the C-band horn.

Once the C-band and X-band matching plugs are designed and incorporated into the flight payload, no further refinements should be required unless the far-field patterns have changed so that they are no longer acceptable for flight.

# Far-Field Patterns of C-Band Antenna

Antenna patterns were measured for different aspect angles varying from  $0^{\circ}$  to 180°. These patterns for various aspect angles are shown in figure B7. The payload top is at  $\phi = 0^{\circ}$  with all patterns measured over 360° in the roll plane of the model. Note that the level of the patterns increases with an increase in aspect angle up to 80° and then decreases as the aspect angle is increased further. The reason for the symmetry around an aspect angle of 80° instead of 90° is the approximate 10° half-angle of the cone-shaped payload. For tracking purposes, the patterns in figure B7 appear very adequate except for aspect angles near 0° and 180°, which are never encountered at Bermuda.

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The curve in figure B8 gives a better idea of how the maximum pattern level varies with aspect angle. The peak gain of the horn is 7.5 dB above the isotropic level.

# Far-Field Patterns of X-Band Antenna

Some typical X-band patterns measured on the RAM C-I are shown in figure B9. Each plot is a  $360^{\circ}$  roll-plane pattern for a specific aspect angle. The aspect angle, starting with  $0^{\circ}$ , is increased in increments of  $10^{\circ}$ . For all the roll-plane patterns, the payload top is at  $\phi = 0^{\circ}$ .

An important observation to make about most of the patterns in figure B9 is the number and magnitude of the peaks and nulls that occur in the roll plane. These variations are due to individual pattern overlap and leaky surface waves radiating around the periphery of the payload. If the payload was not covered with a layer of teflon, the individual pattern overlap would be the only contributor to the peaks and nulls. Also, the deep variations would occur primarily in the regions between any two adjacent horns. However, since a layer of teflon is attached for thermal protection, variations also occur which can be attributed to leaky surface waves. Leaky surface waves (ref. 28) can exist for a dielectric skin thickness of zero. The equation for the different modal cutoff thicknesses t (in wavelengths) is

$$t = \frac{n}{4\sqrt{\epsilon - 1}}$$

where n is an integer representing the mode number and  $\epsilon$  is the dielectric constant. Since the dielectric constant of teflon is 2.1, t = 0.238n. Due to the boundary conditions, only the even TM (transverse magnetic) and the odd TE (transverse electric) modes can exist. Therefore, at 9.21 GHz and a teflon skin thickness of 0.31 in. (0.79 cm), only the TM<sub>0</sub> and the TE<sub>1</sub> leaky wave modes can exist. The cutoff thickness for the TE<sub>1</sub> mode at 9.21 GHz is approximately 0.305 in. (0.77 cm). Thus, with more than 0.005 in. (0.013 cm) of heat-shield ablation, the TE<sub>1</sub> leaky surface wave mode can no longer be excited. The TM<sub>0</sub> mode never cuts off except for a heat-shield thickness of zero.

The effects of leaky surface wave radiation can also be observed in the crosspolarized component of the far-field pattern. Measured patterns depicting the crosspolarized components are shown in figure B10. These patterns were measured in the same manner as those in figure B9 except for the polarization change. The high magnitude of the patterns demonstrates that considerable cross-polarized energy is coupled out of the leaky surface waves; this should not cause undue concern, since it is possible to use the cross-polarized component to advantage. With a polarization diversity ground station, the cross-polarized component can be used to fill in many nulls in the dominantly polarized pattern. An illustration of what should be recorded with a polarization diversity system is shown in figure B11.

#### **Concluding Remarks**

The trajectory for the RAM C-I experiment was very compatible with reception of a strong signal from the payload during the prime data period of the mission. The high level of pattern reception was due primarily to an aspect angle near  $90^{\circ}$  for much of the flight during prime data transmission.

Problems encountered in antenna design were mostly those of a minor nature such as the X-band matching plug and effects of leaky surface waves on far-field patterns. Since the matching plugs were designed properly, no serious detuning due to heat-shield ablation was expected from either the X-band or C-band antennas unless considerably more than 0.2 in. (0.51 cm) of teflon is lost during reentry. The leaky surface wave problem was overcome by choosing the proper polarization for the payload. On the basis of the results obtained during this flight, the antenna designs were adequate.



Figure B1.- RAM C-I reentry trajectory.





Figure B2.- Variation of RAM C-I altitude, velocity, elevation angle, aspect angle, and ablation rate with time from lift-off.



Figure B3.- C-band horn antenna on RAM C-I.

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Figure B4.- Variation of voltage standing wave ratio with frequency of the C-band horn antenna.

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APPENDIX B






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Figure B6.- Variation of voltage standing wave ratio with frequency of four-horn-array X-band antenna.

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APPENDIX B

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Figure B7.- Measured C-band roll-plane antenna patterns for various aspect angles.





Figure B7.- Concluded.

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Figure B8.- Variation of maximum pattern level with aspect angle of C-band horn antenna.

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Figure B9.- Measured X-band roll-plane antenna patterns for various aspect angles.



Figure B9.- Continued.



Figure B9.- Concluded.



Figure B10.- Cross-polarized X-band roll-plane antenna patterns for various aspect angles.



Figure B10.- Continued.

78

APPENDIX B



Figure B10.- Concluded.



Figure B11.- Sum of dominant and cross-polarized far-field components as a function of roll angle for various aspect angles.

## INTERPRETATION OF RAM C-I ELECTROSTATIC PROBE DATA

By W. Linwood Jones, Jr.

#### Introduction

The positive ion current collected by the electrostatic probe is used to measure the free-electron density of the flow-field plasma for altitudes greater than 180 000 ft (54.9 km). In relation to the flowing plasma, each ion collector appears to be a cylindrical wire 0.010 in. (0.0254 cm) in diameter and 0.212 in. (0.5385 cm) long, which is inclined at an angle of  $45^{\circ}$  with respect to the plasma flow. A classical electrostatic probe theory modified to account for the directed ion flux resulting from the plasma flow is used to calculate the positive ion density. This theory has been experimentally verified to be applicable within a factor of  $\pm 2$  for the RAM C-I probe physical configuration at altitudes greater than 150 000 ft (45.7 km). (See ref. 15.)

## Theory

In the positive-ion saturation region of the probe voltage-current characteristic, the current is

$$I_{+} = f(n_{+}, T_{-}, v_{f}, m_{+}, V, A)$$
(C1)

where

I\_+collected currentn\_+positive-ion densityT\_-electron temperaturev\_fflow velocitym\_+ion mass..Vprobe potentialAprobe area

Since there are four unknown variables  $(n_+, T_-, v_f, m_+)$  and only one measured quantity  $(I_+)$ , the quantities  $T_-$ ,  $v_f$ , and  $m_+$  must be determined from theoretical flow-field calculations. The exact functional dependence of  $n_+$  on the other three unknowns is complicated; however, to a first-order approximation

$$n_+ \propto v_f \sqrt{\frac{m_+}{T_-}}$$
 (C2)

To illustrate the data reduction procedure, a sample calculation is presented.

## Sample Calculation

The following factors are given in the sample calculation:

Altitude 250 000 ft (76.2 km)

Ion collector

Measured current  $50 \times 10^{-6}$  A

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The first factor to be calculated is the probe potential relative to the plasma potential V which is obtained from the following equation:

$$V = \left| V_{O} - \frac{5kT_{-}}{e} \right|$$
(C3)

where

 $V_0$  applied bias, -5 volts

k Boltzmann constant

e electronic charge

The electron temperature is assumed to be equal to the local gas temperature. A plot of gas temperature as a function of altitude for the ion collector locations is shown in figure C1. These temperature data as well as the flow velocity data given in figure C2 were generated by a nonequilibrium, boundary-layer-corrected flow-field analysis of the RAM C-I trajectory. For an altitude of 250 000 ft (76.2 km) and ion collector 8, the electron temperature was calculated to be  $3600^{\circ}$  K. The probe potential V is there-fore 6.551 volts.

The next factor to be determined is the potential-to-kinetic energy ratio for ions  $\eta$ , which is calculated to be

$$\eta = \frac{\mathrm{eV}}{\mathrm{kT}} = 21.12 \tag{C4}$$

The ratio of ion-sheath radius to probe radius  $\gamma$  is estimated by manipulating the space-charge-limited diode equation in terms of the measured current, the probe potential, and the physical dimensions of the probe. The result is shown in figure C3 where  $(\gamma - 1)$  is plotted as a function of the parameter  $\Theta$ , which is defined as

$$\Theta = \frac{\mathrm{L}}{\mathrm{r}_{\mathrm{p}}} \frac{\mathrm{V}^{3/2}}{\mathrm{I}_{+}} \left( 1 + \frac{2.66}{\sqrt{\eta}} \right) \tag{C5}$$

where

L probe length, 0.5385 cm

For a measured current of 50  $\mu$ A, the value of  $\Theta$  is  $2.24 \times 10^7$ . The corresponding sheath-to-probe radius ratio  $\gamma$  is 2.25.

The modified normalized probe current R is now determined. Figure C4 is a plot of R as a function of the modified potential-to-kinetic energy ratio H, which is defined as

$$\mathbf{H} = \frac{\eta}{(1 + s^2)} \tag{C6}$$

The ion speed ratio s is

$$s = \frac{v_f \sin \theta}{v_+}$$
(C7)

where

 $v_f$  flow velocity,  $6.35 \times 10^5$  cm/sec

 $\theta$  angle between flow velocity and the cylindrical axis,  $45^{\circ}$ 

 $v_+$  ion velocity, calculated to be

$$v_{+} = \left(\frac{2kT_{-}}{m_{+}}\right)^{1/2} = 1.413 \times 10^{5} \text{ cm/sec}$$
 (C8)

for  $m_{+} = 49.88 \times 10^{-24}$  g (theoretical positive ion concentration 95% NO<sup>+</sup>). From equation (C7) s = 3.180, and thus the value of H is 1.906. The corresponding value of R is 1.61, determined from the following equation:

$$R = \frac{I_{+}}{I_{r}} \frac{1}{(1 + s^{2})^{1/2}}$$
(C9)

where

 $I_r = JA$ 

A probe area,  $0.0420 \text{ cm}^2$ 

J random current density expressed as

$$J = n_{+}e \frac{v_{+}}{2\sqrt{\pi}}$$
(C10)

Solving for  $n_+$  and equating  $n_+$  to  $N_e$  yields

$$N_{e} = \frac{2\sqrt{\pi}I_{+}}{\text{Rev}_{+}A(1+s^{2})^{1/2}} = 3.40 \times 10^{10} \text{ cm}^{-3}$$
(C11)



Figure C1.- Computed electron or gas temperature as a function of altitude for electrostatic probe locations.



Figure C2.- Computed flow velocity as a function of altitude for electrostatic probe locations.



Figure C3.-  $(\gamma - 1)$  as a function of  $\Theta$ .



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Figure C4.- Normalized probe current R as a function of modified potential-to-kinetic energy ratio H for various values of  $\gamma$ .

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