

MASS SPECTROMETRIC DETERMINATION OF THE DISSOCIATION ENERGY OF VANADIUM AND CHROMIUM DICARBIDE AND VANADIUM TETRACARBIDE

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16. Abstract

The Knudsen effusion method was used in conjunction with a double-focusing mass spectrometer to study the vaporization of the vanadium-carbon and chromium-carbon systems over the temperature ranges 2417 to 2603 and 2083 to 2176 K, respectively. The molecular species VC_2 , VC_4 , and CrC_2 were identified, and experimentally determined third-law enthalpies were combined with published thermodynamic data to yield the following dissociation energies:

| $VC_2(g) = V(g) + C_2(g)$ | $D_0^0(V-C_2) = 570 \pm 20 \text{ kJ/mole}$ |
|--|--|
| $VC_4(g) = V(g) + 2C_2(g)$ | $D_0^0(C_2^-V^-C_2) = 1193\pm22 \text{ kJ/mole}$ |
| $\operatorname{CrC}_2(g) = \operatorname{Cr}(g) + \operatorname{C}_2(g)$ | $D_0^0(Cr-C_2) = 445 \pm 18 \text{ kJ/mole}$ |

Atomization energies were also calculated. Results are discussed in terms of the metal oxide - metal dicarbide bond strength analogy.

| 17. Key Words (Suggested by Author(s)) Vanadium and chromium carbides; High- temperature mass spectrometry; Knudsen effusion; Dissociation energies; Vapori- zation of chromium carbide; Vaporization of vanadium carbide; Vanadium dicarbide; Vanadium tetracarbide; Chromium dicarbide | | 18. Distribution Statement Unclassified - unlimited | | | |
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SUMMARY

The Knudsen effusion method was used in conjunction with a double-focusing mass spectrometer to study the vaporization of the vanadium-carbon and chromium-carbon systems over the temperature ranges 2417 to 2603 and 2083 to 2176 K, respectively. The molecular species VC_2 , VC_4 , and CrC_2 were identified, and experimentally determined third-law enthalpies were combined with published thermodynamic data to yield the following dissociation energies:

$$\begin{array}{lll} VC_2(g) = V(g) + C_2(g) & D_0^O(V-C_2) = 570 \pm 20 \ \text{kJ/mole} \\ \\ VC_4(g) = V(g) + 2C_2(g) & D_0^O(C_2-V-C_2) = 1193 \pm 22 \ \text{kJ/mole} \\ \\ CrC_2(g) = Cr(g) + C_2(g) & D_0^O(Cr-C_2) = 445 \pm 18 \ \text{kJ/mole} \end{array}$$

Atomization energies were also calculated. Results are discussed in terms of the metal oxide - metal dicarbide bond strength analogy.

INTRODUCTION

Only a few vaporization studies of the vanadium-carbon (V-C) and chromium-carbon (Cr-C) systems have been reported in the literature (refs. 1 to 5). Two of these studies suggest or present evidence for the existence of gaseous metal-carbon molecular species. Drowart et al. (ref. 1) gave no details of their study but did report an atomization energy of 1167±21 kilojoules per mole for VC₂. Doan (ref. 2) used spark-source mass spectrometry to study the Cr-C system and reported peaks that correspond to the CrC⁺ and CrC⁺₂ molecular ions. His ratios of metal carbide to metal for these two species were

 1×10^{-5} and 3×10^{-4} , respectively. The other investigators (refs. 3 to 5) reported vapor pressures of the metal atoms over various metal-carbon condensed phases.

Recent identifications of gaseous dicarbide and tetracarbide molecules of high stability in equilibrium over condensed phases of various rare-earth-carbon systems (refs. 6 to 11) and transition metal-carbon systems (refs. 12 and 13) prompted the present investigation of the V-C and Cr-C systems. Knudsen effusion molecular-beam mass spectrometric methods appeared to be ideally suited to these studies. The objectives were to use this method to identify metal-carbide molecules and to measure their dissociation energy. The study was limited in both cases to compositions of condensed phase metal carbide in equilibrium with excess carbon at each temperature.

PROCEDURE

The Knudsen cell mass spectrometer system and experimental method used in this study were the same as described previously (refs. 6, 12, and 14). The tungsten Knudsen cell had a 0.076-centimeter-diameter zero-length (knife edge) orifice. The cell was fitted with a graphite liner made from Ultra Carbon UFS^a graphite rod. A 0.25-centimeter-diameter hole in the wall of the liner was alined to the effusion orifice.

For the study of the V-C system, a homogeneous mixture of graphite (Ultra Carbon UFS-4) and vanadium, in the atomic ratio of 2.1 to 1.0, was loaded into a new cell liner. The vanadium was 99 percent pure as obtained from Cerac Inc. The mixture was heated in the Knudsen cell to 2330 K. Eight hours were required to reach this temperature with the pressure maintained at less than 1×10^{-5} torr. The cell was kept at temperature for an additional 16 hours before measurements were started. At the conclusion of the experiment, X-ray diffraction analysis showed the residue to be a mixture of graphite and cubic vanadium carbide with a lattice parameter of 416.5±0.2 picometers. According to Storms (ref. 4), this lattice parameter corresponds to a composition of approximately VC_{0.88}, which is close to the upper phase limit for the VC phase.

For each of the two experiments (I and II) on the Cr-C system, the cell liners were loaded with a homogeneous mixture of Cr_3C_2 and graphite (Ultra Carbon UFS-4) in the weight ratio of 10 to 1. The Cr_3C_2 , of unspecified purity, was purchased from Astro Metals Inc. Spectrographic analysis for metallic impurities revealed moderate amounts of aluminum plus trace amounts of boron, copper, iron, vanadium, nickel, silicon, and

^aCertain commercial products are identified by name in this report to specify the experiment completely. In no case does such identification imply a recommendation or endorsement by NASA, nor does it imply that the products are the best available for the purpose.

titanium. Each $\operatorname{Cr}_3\operatorname{C}_2$ -C mixture was heated in a Knudsen cell to 1400 K over a period of 24 hours prior to the start of the ion intensity measurements. For the last 10 hours of the preheating period, the pressure was maintained at less than 2×10^{-6} torr. X-ray diffraction analysis of the sample residue from experiment II showed that the solid phase was a mixture of graphite, $\operatorname{Cr}_3\operatorname{C}_2$, and a small amount of $\operatorname{Cr}_7\operatorname{C}_3$.

The ion species were identified by their mass-to-charge ratio (m/e) and isotopic abundances. A movable slit between the ion source and the Knudsen cell was used as a shutter. The shutter could be used to block the molecular beam, to allow optimum passage of the beam, or to scan the beam to study its profile. The shutter was used to establish which species were originating from the Knudsen cell. The shutter effect is the difference in measured ion intensity with the shutter in the open and closed positions. All intensity measurements were made with ionizing electron energies of 20 eV and an anode current of 150 microamperes. Appearance potentials were measured by recording ion intensity against electron energy (ionization efficiency curve) and applying the linear extrapolation method.

RESULTS AND CALCULATIONS

The only pertinent ion species detected were V^+ , VC_2^+ , VC_4^+ , Cr^+ , and CrC_2^+ . A careful search at the maximum temperature of each experiment failed to reveal any monocarbides or higher carbides other than those just noted.

In the vanadium experiment, ion intensities were measured at 13 temperatures in the range from 2417 to 2603 K. Below 2417 K, the intensity of the carbide species was too low to measure, and above 2603 K, the intensity of V^+ became high enough to indicate that the pressures were approaching the point where the conditions of effusive flow no longer applied. The shutter effect for V^+ , VC_2^+ , and VC_4^+ was 100 percent.

The temperature range for the two chromium experiments was from 2083 to 2176 K. In the first chromium experiment, ion intensities were measured at 11 temperatures, and in the second experiment, intensities were measured at seven temperatures. The high and low limits of the temperature range studied were dictated by the same factors that prevailed in the vanadium experiment. The shutter effect for Cr^+ was 100 percent, but for CrC_2^+ the shutter effect was between 80 and 100 percent. Ion intensity values were derived by making measurements with the shutter open and closed and taking differences. It was not possible to identify the nonshutterable background peak at m/e = 76 that interferred with the CrC_2^+ peak, and the resolution of the mass spectrometer was not high enough to separate it from the CrC_2^+ .

The measured ion currents I_i are related to partial pressures P_i at temperature T by the equation

$$P_{i} = \frac{kI_{i}TE_{i}}{\sigma_{i}\gamma_{i}n_{i}}$$
(1)

where k is the sensitivity constant for the apparatus, E_i is the experimentally determined factor to correct ion intensities measured at the working voltage to the maximum of the ionization efficiency curve (ratio of intensity at maximum to intensity at measurement electron energy), σ_i is the relative maximum electron impact ionization cross section, γ_i is the relative multiplier gain, and n_i is the fractional isotopic abundance of the species i.

For a reaction of the form

$$M(g) + xC(s) = MC_{x}(g)$$
(2)

the equilibrium constant $K_{\mathbf{p}}$ is given as

$$K_{P} = \frac{P_{MC_{x}}}{(P_{M})(a_{C})^{X}}$$
(3)

where a_{C} is the activity of the solid graphite phase. Reactions of this type are termed pressure independent since K_{P} is directly proportional to the ratio of the pressures of the gaseous products and reactants. Therefore, for such reactions,

$$K_{\mathbf{p}} = \frac{I_{1}E_{1}\sigma_{2}\gamma_{2}n_{2}}{I_{2}E_{2}\sigma_{1}\gamma_{1}n_{1}}$$
(4)

where 1 and 2 refer to the products and reactants, respectively, and the activity of graphite is taken as 1.

Ionization efficiency curves could only be obtained for V⁺ and VC₂⁺ because the other ion intensities were too low to obtain meaningful results. The measured appearance potential for VC₂⁺ was 8.6±0.5 eV, with the 6.74-eV value for V⁺ (ref. 15) used as the internal standard. The factor E_i was required to correct ion intensities obtained with a working ionizing energy of 20 eV to the maximum of the ionization efficiency curve. The factor E_i was measured directly for V⁺ and VC₂⁺, but for VC₄⁺, $E_{VC_4}^+ = E_{VC_2}^+$ was assumed. For chromium, $E_{CrC_2}^+ = E_{VC_2}^+$ was assumed.

Relative maximum ionization cross sections σ_i for atoms were taken from Mann (ref. 16), and the cross sections for VC₂, VC₄, and CrC₂ were calculated by summing the respective metal atom cross section with the cross section of the appropriate number

of carbon atoms. All carbide ion intensities were too low to measure the absolute gain of the electron multiplier, and, therefore, the relative gains γ_i for the various species were taken to be the same as those measured in a previous work (ref. 12) on the titanium-carbon system.

Three pressure-independent reactions were considered. These reactions and the corresponding calculated third-law heats are

$$V(g) + 2C(s) = VC_2(g)$$
 $\Delta H_0^0 = 254.1 \pm 18.0 \text{ kJ/mole}$ (5)

$$V(g) + 4C(s) = VC_4(g)$$
 $\Delta H_0^0 = 455.6 \pm 18.8 \text{ kJ/mole}$ (6)

$$Cr(g) + 2C(s) = CrC_2(g) \qquad \Delta H_0^0 = 379 \pm 16 \text{ kJ/mole}$$
 (7)

The value for chromium was obtained by combining the results of experiments I and II. Unit activity for C(s) was assumed because the graphite liner always assured an excess of carbon. Values of the various parameters just discussed are listed in table I. Equilibrium constants for reactions 5 to 7 are listed in tables II to IV for each measurement temperature.

TABLE I. - RELATIVE MULTIPLIER GAINS, CROSS SECTIONS, ION INTENSITY

FACTORS, AND APPEARANCE POTENTIALS FOR VANADIUM-

| Ion | Relative multiplier gain, ^γ i | Relative maximum electron impact ionization cross section, ^c i | Ion intensity cor- rection factor, E _i | Appearance potential, A.P., eV |
|--|--|---|---|--------------------------------------|
| $\begin{array}{c} v^+ \\ v c_2^+ \\ v c_4^+ \\ c r^+ \\ c r c_2^+ \end{array}$ | 1.0 b1.4 ^c 1.4 1.0 b1.4 | 5.58 9.04 12.50 5.45 8.91 | $1.00 \\ 1.36 \\ d_{1.36} \\ 1.00 \\ e_{1.36} $ | ^a 6.74 8.6±0.5 |

AND CHROMIUM-CONTAINING IONS

^aRef. 15 (standard).

 ${}^{b}\gamma_{V^{+}}/\gamma_{VC_{2}^{+}}$ and $\gamma_{Cr^{+}}/\gamma_{CrC_{2}^{+}}$ were taken to be same as that measured for $\gamma_{Ti^{+}}/\gamma_{TiC_{2}^{+}}$ (ref. 12).

(ref. 12). $^{\circ}\gamma_{VC_{2}^{+}/\gamma_{VC_{4}^{+}}}$ was taken to be same as that measured for $\gamma_{TiC_{2}^{+}/\gamma_{TiC_{4}^{+}}}$ (ref. 12).

^dAssume that $E_{VC_4^+} = E_{VC_2^+}$ (measured).

^eAssume that $E_{CrC_2^+} = E_{VC_2^+}$ (measured).

| Temperature, | Equilibrium | Log of equilibrium | Change of Gibbs free | Third-law |
|--------------|-----------------------|--------------------|--|--------------------------------|
| т, | constant, | constant, | energy function, | enthalpy, |
| к | к _р | $\log_{10} K_{P}$ | $-\Delta\left[(G_{T}^{O}-H_{0}^{O}/T)\right],$ | ΔH ₀ ^o , |
| | | | J/(deg)(mole) | kJ/mole |
| 2503 | 1.21×10 ⁻³ | -2.917 | 45.97 | 254.8 |
| 2568 | 1.80 | -2.745 | 45.72 | 252.4 |
| 2539 | 1.37 | -2.863 | 45.83 | 255.5 |
| 2547 | 1.69 | -2.785 | 45.80 | 252.5 |
| 2580 | 1.95 | -2.710 | 45.68 | 251.7 |
| 2549 | 1.59 | -2.799 | 45.79 | 253.3 |
| 2503 | 1.33 | -2.876 | 45.97 | 252.9 |
| 2488 | 1.22 | -2.914 | 46.02 | 253.3 |
| 2470 | 1.18 | -2.928 | 46.09 | 252.3 |
| 2452 | 8.41 $\times 10^{-4}$ | -3.075 | 46.15 | 257.5 |
| 2417 | 7.28×10^{-4} | -3.138 | 46.27 | 257.0 |
| 2557 | 1.50×10^{-3} | -2.824 | 45.76 | 255.2 |
| 2603 | 1.84×10^{-3} | -2.735 | 45.59 | 255.0 |
| | | | Average | ^a 254.1±1.9 |

TABLE II., - THIRD-LAW ENTHALPIES FOR REACTION $V(g) + 2C(s) = VC_2(g)$

^aError quoted is standard deviation of points. See text for overall estimated uncertainty.

| TABLE III | - THIRD-LAW | ENTHALPIES | FOR REACTION | V(g) + 4C(s) = | $VC_4(g)$ |
|-----------|-------------|------------|--------------|----------------|-----------|
|-----------|-------------|------------|--------------|----------------|-----------|

| Temperature, T, K | Equilibrium constant, K _P | Log of equilibrium constant, ^{log} 10 ^K P | Change of Gibbs free energy function, $-\Delta \left[(G_T^0 - H_0^0)/T \right],$ J/(deg)(mole) | Third-law enthalpy, △H ⁰ ₀ , kJ/mole |
|-------------------------|--|---|--|--|
| 2503 | 3.16×10 ⁻⁶ | -5,500 | 73.93 | 448.6 |
| 2568 | 3.18 | -5.498 | 73,72 | 459.6 |
| 2539 | 3.45 | -5.462 | 73.82 | 452.9 |
| - 2547 | 3.00 | -5.523 | 73.79 | 457.2 |
| 2580 | 4.61 | -5.336 | 73.68 | 453.7 |
| 2549 | 3.12 | -5.506 | 73.78 | 456.8 |
| 2503 | 2.48 | -5.606 | 73.93 | 453.7 |
| 2557 | 3.32 | -5.479 | 73.76 | 456.8 |
| 2603 | 3.97 | -5.401 | 73.61 | 460.8 |
| | | | Averag | $e^{a_{455.6\pm 3.7}}$ |

^aError quoted is standard deviation of points. See text for overall estimated uncertainty.

| Temperature, T, K | Equilibrium constant, K _P | Log of equilibrium constant, ^{log} 10 ^K P | Change of Gibbs free energy function, $-\Delta \left[(G_T^0 - H_0^0)/T \right],$ J/(deg)(mole) | Third-law enthalpy, ∆H ⁰ ₀ , kJ/mole | | | | |
|-------------------------|--|---|--|--|--|--|--|--|
| | Experiment I | | | | | | | |
| 2083 | 4.1 ₈ ×10 ⁻⁷ | -6.379 | 62.32 | 384 | | | | |
| 2119 | 7.61 | -6.119 | 62,24 | 380 | | | | |
| 2133 | 9.2 ₈ | -6.032 | 62.21 | 379 | | | | |
| 2111 | 6.8 ₉ | -6.162 | 62.26 | 381 | | | | |
| 2116 | 7.63 | -6.117 | 62.25 | 380 | | | | |
| 2154 | 1.02×10^{-6} | -5.991 | 62.17 | 381 | | | | |
| 2125 | $1.0_{5}^{-} \times 10^{-6}$ | -5.979 | 62.23 | 376 | | | | |
| 2124 | $6.5_5 \times 10^{-7}$ | -6.184 | 62.23 | 384 | | | | |
| 2146 | $1.0_1 \times 10^{-6}$ | -5.996 | 62.18 | 380 | | | | |
| 2166 | $1.2_{3} \times 10^{-6}$ | -5.910 | 62.14 | 380 | | | | |
| 2100 | $1.0_4 \times 10^{-6}$ | -5.983 | 62.28 | 371 | | | | |
| | | Experime | nt II | | | | | |
| 2101 | 9.2 ₇ ×10 ⁻⁷ | -6.033 | 62.28 | 374 | | | | |
| 2132 | 1.41×10^{-6} | -5.851 | 62.21 | 371 | | | | |
| 2154 | $9.2_0 \times 10^{-7}$ | -6.036 | 62.17 | 383 | | | | |
| 2188 | 1.79×10^{-6} | -5.747 | 62.10 | 377 | | | | |
| 2121 | $8.2_{5}^{-7} \times 10^{-7}$ | -6.084 | 62.24 | 379 | | | | |
| 2147 | $1.3_{4} \times 10^{-6}$ | -5.873 | 62.18 | 375 | | | | |
| 2176 | 1.25×10^{-6} | -5.903 | 62.11 | 381 | | | | |
| | | | Ave | erage ^a 379±4 | | | | |

TABLE IV. - THIRD-LAW ENTHALPIES FOR REACTION $Cr(g) + 2C(s) = CrC_2(g)$

^aError quoted is standard deviation of points. See text for overall estimated uncertainty.

Third-law heats were calculated according to the relation

$$\Delta H_0^0 = -2.303 \text{ RT} \log_{10} K_P - T \Delta \left(\frac{G_T^0 - H_0^0}{T} \right)$$
(8)

where R is the gas constant and $\Delta \left[(G_T^0 - H_0^0)/T \right]$ is the change of the Gibbs free energy function for the reaction under consideration. The values of $\Delta \left[(G_T^0 - H_0^0)/T \right]$ and ΔH_0^0 are presented in tables II to IV.

The values of the free energy functions used for V(g) and Cr(g) were taken from Hultgren et al. (ref. 17), and those used for C(s) were taken from the JANAF tables (ref. 18). The heat contents and free energy functions for $VC_2(g)$ and $CrC_2(g)$ were calculated on the basis of estimated molecular parameters for an assumed linear asym-

| | Molecule | | | |
|---|---|---|--|--|
| | vc ₂ | vc4 | CrC ₂ | |
| Internuclear separation, pm | | | | |
| V-C | *162 | 162 | , | |
| Cr-C | | | *168 | |
| C-C | **131 | 131 | 131 | |
| Moment of inertia, g-cm ² | 15.74×10 ⁻³⁹ | 44.71×10 ⁻³⁹ | 16.59×10 ⁻³⁹ | |
| Force constants, N/m k_{V-C}^{\dagger} k_{Cr-C}^{\dagger} $k_{C-C}^{\ast\ast\ast}$ $k_{C-C}^{\ast\ast\ast}$ k_{δ}^{δ} , N·m/rad [‡] $f_{13, 22}^{\dagger}$ | 7. 36×10 ² 9. 25×10 ² 6. 70×10 ⁻¹⁹ | 7. 36×10^{2} 9. 25×10^{2} 1. 30×10^{2} | 5.82×10 ² 9.25×10 ² 6.70×10 ⁻¹⁹ | |
| Frequencies (degeneracy), cm^{-1} | | | | |
| ω_1 | 780(1) | 812(1) | 709(1) | |
| ω_2 | 494(2) | 1883(1) | 488(2) | |
| ω_3 | 1814(1) | 1013(1) | 1767(1) | |
| ω_4 | | 1841(1) | | |
| ω_5 | | 136(2) | | |
| ^ω 6 | | 570(2) | | |
| ω_{7} | | 468(2) | | |
| Ground-state statistical weight | "4 | "4 | # ₅ | |

TABLE V. - ESTIMATED MOLECULAR PARAMETERS FOR

VC₂, VC₄, AND CrC₂

*Calculated by use of modified Badger's rule (ref. 25).

** From C₂ (ref. 26).

[†]Force constants for V-C and Cr-C were assumed to be same as those of corresponding oxides (ref. 26). See ref. 27.

[§]See ref. 19.

^{II}Ground state of ${}^{4}\Sigma$ was assumed for VC₂ by analogy to VO (ref. 28). The same electronic contribution was assumed for VC₄. [#]Ground state of ${}^{5}\Sigma$ was assumed for CrC₂ by analogy with apparent

ground state for CrO (private communication with K. D. Carlson of Case Western Reserve University).

metric M-C-C molecule. For $VC_4(g)$, a linear symmetric C-C-V-C-C structure was assumed, and the method of Engler and Kohlrausch (ref. 19) was used to estimate fundamental vibration frequencies. All thermodynamic functions were calculated with a computer program (ref. 20) using the rigid rotator - harmonic oscillator approximation. Pertinent molecular parameters are listed in table V, and heat contents and free energy functions are given in tables VI and VII. The estimated total errors associated with the third-law heats ΔH_0^0 were obtained by combining the standard deviations from the mean of all data points listed in tables II to IV with estimated uncertainties in the values of the

| | | 2 | 7 |
|-------------------|--------------------|------------------------------------|--|
| | Temperature, T, | Heat content, $H_T^0 - H_0^0$, | Free energy function, - $(G_T^0 - H_0^0)/T$, |
| | ĸ | kJ/mole | J/(deg)(mole) |
| vc ₂ (| | | |
| ſ | 298.15 | 10.10 | 213.83 |
| | 2000 | 106.37 | 297.14 |
| I | 2100 | 112.46 | 299.75 |
| | 2200 | 118.57 | 302.25 |
| | 2300 | 124.69 | 304.65 |
| | 2400 | 130.81 | 306.96 |
| | 2500 | 136.95 | 309.19 |
| | 2600 | 143.09 | 311.35 |
| | 2700 | 149.24 | 313.43 |
| | 2800 | 155.40 | 315.44 |
| | 2900 | 161.56 | 317.39 |
| | 3000 | 167.72 | 319.29 |
| | | | |
| | 298.15 | 14.72 | 233.74 |
| | 2000 | 184.09 | 369.51 |
| | 2100 | 195.00 | 374.02 |
| | 2200 | 205.94 | 378.36 |
| | 2300 | 216.90 | 382.53 |
| | 2400 | 227.88 | 386.56 |
| | 2500 | 238.88 | 390.45 |
| | 2600 | 249.90 | 394.21 |
| | 2700 | 260.94 | 397.84 |
| | 2800 | 271.98 | 401.37 |
| | 2900 | 283.04 | 404.78 |
| | 3000 | 204 11 | 408 10 |

| TABLE VI | HEAT | CONTENTS | AND | FREE | ENERGY |
|----------|------|----------|-----|------|--------|
|----------|------|----------|-----|------|--------|

| ENERGY FONCTIONS FOIL CIC2(g) | | | |
|-------------------------------|--------------|-------------------|-------------------------|
| | Temperature, | Heat content, | Free energy function |
| | T, | $H_T^0 - H_0^0$, | - $(G_T^0 - H_0^0)/T$, |
| | K | kJ/mole | J/(deg)(mole) |
| | 298.15 | 10.19 | 216,43 |
| | 1800 | 94.78 | 294,89 |
| | 1900 | 100.86 | 297,75 |
| | 2000 | 106.95 | 300,49 |
| | 2100 | 113.05 | 303,10 |
| | 2200 | 119.17 | 305,62 |
| | 2300 | 125.29 | 308.03 |
| | 2400 | 131.43 | 310.36 |

137.57

2500

TABLE VII. - HEAT CONTENTS AND FREE

per mole. The estimated error for the values of the relative cross sections accounts for the metal-dicarbide molecules.

ENERCY ETIMOTIONS FOR Croc (a)

following: K_{P} , ±50 percent; T, ±10 K; and $\Delta \left[(G_{T}^{O} - H_{0}^{O})/T \right]$, ±6.3 joules per degree for a large part of the overall estimated error in K_{P} . The estimated error for $\Delta \left[(G_T^0 - H_0^0) / T \right]$ does not include a possible difference in ground-state statistical weight

The heat of formation of $C_2(g)$, $\Delta H_{0,f}^0 = 824.2\pm 8.4$ kilojoules per mole (ref. 18), was combined with the third-law heats of reactions (eqs. (5) and (7), respectively), to give the following dissociation energies:

$$VC_2(g) = V(g) + C_2(g)$$
 $D_0^0(V-C_2) = 570.1 \pm 19.9 \text{ kJ/mole}$ (9)

312.60

$$\operatorname{CrC}_{2}(g) = \operatorname{Cr}(g) + \operatorname{C}_{2}(g)$$
 $D_{0}^{0}(\operatorname{Cr-C}_{2}) = 445 \pm 18 \text{ kJ/mole}$ (10)

The atomization energies $D_{0, \text{ atom}}^{0}$ of V-C-C and Cr-C-C were obtained by combining, respectively, the heats of reactions (5) and (7) with the heat of formation of C(g), which is 709.5 ± 1.9 kilojoules per mole (ref. 18):

$$VC_2(g) = V(g) + 2C(g)$$
 $D_{0, atom}^0 = 1164.9 \pm 18.2 \text{ kJ/mole}$ (11)

$$CrC_2(g) = Cr(g) + 2C(g)$$
 $D_{0, \text{ atom}}^0 = 1040 \pm 16 \text{ kJ/mole}$ (12)

The heat of formation of $C_2(g)$ was combined with the third-law heat of reaction (eq. (6)) to give the dissociation energy of C_2 -V- C_2

$$VC_4(g) = V(g) + 2C_2(g)$$
 $D_0^0(C_2 - V - C_2) = 1192.8 \pm 22.2 \text{ kJ/mole}$ (13)

The atomization energy of $VC_4(g)$ was calculated by combining the heat of reaction (eq. (6)) with the heat of formation of C(g):

$$VC_4(g) = V(g) + 4C(g)$$
 $D_{0, atom}^0 = 2382.4 \pm 19.2 \text{ kJ/mole}$ (14)

DISCUSSION

Second-law treatments of the experimental data were not attempted because of the relatively short temperature ranges covered in each experiment. However, the absence of any severe trend with temperature in the third-law heats serves to indicate that no large systematic temperature error existed.

The value obtained for the atomization energy of $VC_2(g)$, 1165±18 kilojoules per mole, is in excellent agreement with that reported by Drowart et al. (ref. 1), 1167±21 kilojoules per mole. The results obtained for the Cr-C system, however, differ in two respects from findings reported by Doan (ref. 2). First, no monocarbide was detected, and, second, the ratio of CrC_2^+ to Cr^+ is smaller by a factor of 10^{-2} to 10^{-3} than that reported by Doan. These differences in results are not particularly disturbing because Doan's results were obtained by sparking, a process which is not well understood. The mechanism of vaporization and ion production by sparking is complex. The energy dissipated by the spark vaporizes, ionizes, and physically removes material from the electrodes. The neutral and ionic species formed do not necessarily represent an equilibrium vaporization process. On the other hand, the present results are derived from the Knudsen method, which samples an equilibrium process.

The comparisons of M-C₂ and M-O bond strengths that have previously been made for rare-earth-, Group IIIA-, Group IVA- (refs. 6 to 11, 21, and 22), and transition metal-carbon systems (ref. 12) can now be extended to include the dicarbides of vanadium and chromium. The comparison of M-C₂ and M-O bond strengths is made because it has been postulated that the bonding of the C_2^{-} group is similar (isoelectronic) to that of the O⁼ ion (refs. 21 and 22). The V-C₂ bond energy (570±20 kJ) follows the apparent trend that the oxide bond energy, $D_0^{0}(V-O) = 617\pm11$ kilojoules (ref. 23), is somewhat higher than the bond energy of the corresponding dicarbide. However, the difference between $D_0^{0}(V-C_2)$ and $D_0^{0}(V-O)$ is not as large as that of some of the other metal oxides and dicarbides, notably titanium (ref. 12), and some of the rare earths (refs. 6 and 21). The case of chromium seems to be an exception because the dicarbide bond energy (445±18 kJ) is slightly higher than that of the oxide, $D_0^{0}(Cr-O) = 423\pm29$ kilojoules (ref. 23). However, the overlapping uncertainties make the situation nebulous enough that no positive conclusion can be made.

It is interesting to note that the bond energies for C_2 -V- C_2 , $D_0^0 = 1193\pm22$ kilojoules, and O-V-O, $D_0^0 = 1213\pm42$ kilojoules (ref. 23), are comparable. This is in contrast to the case of titanium, where the O-Ti-O species, $D_0^0 = 1311\pm21$ kilojoules (ref. 24), is markedly more stable than C_2 -Ti- C_2 , $D_0^0 = 1214\pm23$ kilojoules (ref. 12).

The value of the atomization energy for $VC_4(g)$ obtained in this study indicates that this molecule is reasonably stable. The stabilities of VC_2 and VC_4 and the existence of stable oxides of niobium and tantalum (ref. 23) suggest that the corresponding carbides of niobium and tantalum might be stable molecules.

Lewis Research Center,

National Aeronautics and Space Administration, Cleveland, Ohio, January 14, 1970, 129-03.

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