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Earth Models Consistent with Geophysical Data

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ABSTRACT

A suite of the most recently available geophysical data are inverted by an improved Monte Carlo procedure. The data are derived from surface waves for oceanic paths, eigenvibrations of the earth, elastic wave travel time and $dt/d\Delta$ data, mass and moment of inertia of the earth. A low velocity zone is required for the suboceanic mantle as is a high density lithosphere. The high density is related to eclogite fractionation from the underlying, partially molten asthenosphere in a process involving the creation and spreading of the lithosphere. If the asthenosphere is pyrolite or peridotite then an increase of mean atomic weight across the transition zone seems required. Fairborn's new $dt/d\Delta$ data for the lower mantle seem to show a higher shear velocity gradient than previously supposed. If correct, a compensatory lower density gradient is required. This may indicate a depletion of iron with depth in the lower mantle. The density at the top of the core is surprisingly well constrained to the range 9.9-10.2 gm/cc, a value appropriate for a mixture of iron and about 15 wt % silicon.

I. Introduction

A major goal of geophysics is to uniquely specify the distribution of two elastic velocities and density with depth in the earth and to relate these distributions to variations in composition, phase, and temperature in the interior. Impediments which block these achievements are many. It has not been proved in the mathematical sense that a unique solution can be obtained, although BACKUS and GILBERT (1968) have shown that under certain circumstances stable weighted averages (over depth) can be calculated. Furthermore, the data set available for recovering earth structure is incomplete and imprecise. Finally the equations of state available for interpreting elasticity and density distributions in terms of composition, state, etc. are tentative ones based on uncertain theories and assumptions and limited laboratory data.

Despite these difficulties it may be possible even with the presently available data to make some meaningful statements about the interior. In this paper we explore this possibility using the most recent and best available data in a Monte Carlo inversion procedure. Our results and conclusions supersede those presented in earlier papers (PRESS, 1968 a & b) because we are able to fit models to new, more extensive and accurate data with greater speed and better precision.

II. Method

The Monte Carlo method uses random selection to generate large numbers of models in a computer, subjecting each model to a test against geophysical data. Only those models are retained whose properties fit the data within a prescribed tolerance. The procedure offers the advantage that successful models are found without bias, preconceived ideas or uncertain assumptions of equations of state or composition. If the program is efficient so that a

very large number of models are examined, the retained models can be considered as representative of the family of successful models which fit the data. When the successful models fall in a narrow band, geophysically meaningful conclusions can often be reached despite our inability to specify a single unique model. Under certain conditions (BACKUS and GILBERT 1968) a single successful model can provide unique, local averages of density or velocity.

We have modified the Monte Carlo procedure reported last year (PRESS, 1968 a) speeding up the process by 1-2 orders of magnitude. This improved efficiency enabled us to find a larger number of successful models fitting a more extensive suite of data with better precision. The flow diagram of the currently used system (fig. 1) is printed with each run of the program and provides diagnostics so that controlling constants can be set for maximum efficiency. The figure shows the diagnostics following a run of 3347 seconds on an IBM 360-65 computer which yielded 11 successful models at a cost of about \$10 per successful model. SLMD is the random selection procedure for compressional velocity (α), shear velocity (β) and density (ρ). TTT is a test of the model against observed travel times using BULLEN'S (1961) method in which the earth is treated as a multi-layered sphere, the velocity varying according to a power law within each layer. VRPR uses a table of variational parameters (WIGGINS 1968) stored permanently in a data cell of the computer to test the perturbation of the eigenperiods due to velocity, density or core radius perturbations. This test is made after the selection of density and velocity models, the latter in order to eliminate early in the process those models which cannot be brought into agreement with eigenperiod data by density perturbations. MASMOM tests each density model against mass and moment of inertia of the earth. The flow diagram shows branching

according as the several tests are passed or failed. Each box shows the number of times the corresponding step was repeated, the average time and total time for each step and the percentage of model passing. Thus the time distribution over the various components of the program is available for adjustment of input constants for maximum efficiency and insight is provided as to how the various geophysical constraints figure in the elimination of models. A key requirement of the Monte Carlo method is that the selection procedure produce an unbiased, representative suite of models for examination. Figures 2, 3, and 4 show a run in which 25 models were generated, bypassing tests against geophysical data. It is seen that the velocity and density space between the permissible bounds is nearly uniformly filled. Since millions of earth models were generated and examined in this study it would be surprising if continued operation of the program would produce a successful model significantly different from those presented later in this paper. Some additional features of the program and procedures used are as follows:

1. The earth is assumed to be spherically symmetrical and isotropic, with an oceanic crust-upper mantle structure. The radius of the core is selected randomly for each model in the range 3473 ± 25 km.
2. Although α, β, ρ could be varied at 88 points in the earth, we chose the time saving device of randomly varying 19 points, (see fig. 5, section D for their location) obtaining the remaining values by linear interpolations.
3. The fluid core was assumed to be adiabatic. The density selection procedure for the mantle below 1000 km eliminated models with extreme density gradients. The gradients in α, β, ρ were restricted to a maximum number of reversals in sign (typically 2 or 4) to restrict the complexity of models.
4. The rigidity for the inner core only affects the mode ${}_0S_2$. Although zero rigidity was assumed in this paper, the

systematic, negative residuals for ${}_0S_2$ found in our models will be used to infer a rigidity for the inner core.

5. Several exact calculations of eigenperiods were made to check the accuracy of the variational parameter method. The differences were small and well within the uncertainty of the data.

III. The Data

Successful models were required to fit the following data:

1. Earth mass, $M=5.976 \times 10^{27}$ grams; dimensionless moment of inertia $I/Ma^2 = .3308$

2. Compressional velocity distribution in the mantle fixed very close to the models determined by JOHNSON (1969) and FAIRBORN (1969), based on $dt/d\Delta$ analyses of array data. P and PcP travel times fit the latest data (HERRIN et al. 1968) to ± 1 sec. The compressional velocity distribution in the core was fixed to the recent model of HUSEBYE and TOKSÖZ (1969).

3. Shear velocities below 800 km were restricted to lie within the narrow bounds reported by FAIRBORN (1969) who used Monte Carlo methods to interpret travel time and $dt/d\Delta$ data obtained from the Large Aperture Seismic Array (LASA). Wider bounds were used above 800 km. Travel times of S and ScS were required to fit FAIRBORN'S data to within ± 5 sec at 10 distances between 25° and 100° and a single failure was sufficient to reject a model. These travel time data primarily constrain the mantle below 800 km.

4. Eigenperiods tested were ${}_0S_0, {}_0S_2$ through ${}_0S_{22}, {}_1S_2, {}_1S_3, {}_1S_5, {}_1S_6, {}_1S_8, {}_1S_{12}, {}_2S_4, {}_2S_6, {}_2S_{10}$; toroidal oscillations tested were ${}_0T_3, {}_0T_{21}, {}_cT_2$ not being used because of its uncertain value; models were also required to fit surface wave phase velocities for predominantly oceanic paths as follows: Rayleigh waves in the period range 125-325 seconds (BEN-MENACHEM, 1965); Love waves in the period range 80-340 seconds (TOKSÖZ and ANDERSON, 1966). We used the eigenperiod data as reviewed and summarized by DERR (1969)

except as follows: S_2 - 3229.0 seconds, S_{11} - 537.5 seconds. The uncertainty in the eigenperiod and dispersion data was taken to be $\pm 0.4\%$ due to asphericity, rotational splitting and experimental errors (DAHLEN 1968). An error analysis of the oceanic surface wave data indicates that an accuracy better than 1% was achieved. Comparison with phase velocities for other oceanic paths verified this for Love and Rayleigh waves. The fit of S_0 was required to be within $\pm 0.1\%$. Actually the final models fit most of the data to about half these tolerances. Figure 5 shows the computed eigenperiods and the residuals for a typical model.

IV. Results

The results reported here supersede our earlier conclusions (PRESS 1968) because of the new and more extensive data set inverted in this paper. The effects of lateral variation were reduced by deriving higher mode data from oceanic surface wave phase velocities. Moreover the new procedure enabled us to find a much larger number of successful models and therefore a more representative selection from the set of successful models.

The shear velocity and density distribution are plotted in figs. 6, 7, and 8 and are also tabulated (together with the fixed compressional velocity distribution) in figs. 9-13.

V. The Upper Mantle Under Oceans

Without exception every successful model contains a low velocity zone for shear waves which centers at depths between 150 and 250 km. If the lid of this zone is characterized by $\beta > 4.5$ km/sec, then its thickness is 50-100 km. We failed to find a single model without a low velocity zone despite a special search in which 162,000 monotonic shear velocity models were examined. A low velocity zone seems required by our data since essentially every possible model without it was examined and eliminated. Nevertheless, HADDON and BULLEN (1969) reported a successful monotonic model, probably because: (1) they only use modes through $n=44$, whereas our data go to $n=105$; (2) our Love wave phase velocities trend towards lower values than the HB data (see fig. 14).

The several mechanisms which might account for the low velocity

partial (grain boundary) melting for the following reasons: (1) shear velocity and Q are sensitive to the presence of small amounts of melt along grain boundaries; (2) data presented at this conference by several investigators show that the temperature at which melting begins in the wet state for candidate upper mantle mineral assemblages is sufficiently low to be reached by the geotherm for most thermal models of the earth; (3) the partial melting product of candidate mineral assemblages can account for basaltic vulcanism; (4) a partially molten, low strength zone would serve to mechanically decouple the lithosphere from the underlying mantle as is required by some proposed mechanisms for the spreading sea floor.

The density values shown in figs. 7 and 15 fill the entire permissible range at the M-discontinuity indicating a lack of constraint by the geophysical data. However, the initial density gradients are all positive and in the vicinity of 100 km all the values fall in the narrow band $3.5-3.6 \text{ gm/cm}^3$ in the upper part of the permissible range. As a check on this result a special search was made without success to find models with densities below 3.4 gm/cm^3 in this depth range. For additional confirmation of this result we applied the BACKUS and GILBERT (1968) δ -ness criterion using weighting functions computed for our data by WIGGINS (1969). According to Backus and Gilbert, if the weighting functions are concentrated over narrow depth intervals, a stable local average can be obtained from a single model. Using this procedure every one of the models yielded an average density in the range $3.5-3.6 \text{ g/cm}^3$ for the depth interval 75-125 km. Presumably the average density near 100 km is uniquely determined in the sense that any model computed from our data set should give the same value.

Unfortunately the density resolution deteriorates below 100 km as can be seen by the wider band of Monte Carlo solutions. At

300 km the resolving length inferred from the \int -ness criterion is 200 km. One might argue on physical grounds that the lower density solutions should be favored below 150 km because of the low shear velocities. This implies a density reversal from the lithosphere to the asthenosphere ($3.5\text{--}3.6 \text{ gm/cm}^3$ at 100 km to $3.3\text{--}3.5 \text{ gm/cm}^3$ at 300 km).

More complex models were found involving two low velocity or two low density zones in the upper mantle. However, these models yield the same indication of high density near 100 km.

The indicated density for the lithosphere near 100 km is so high as to narrow the range of its possible composition to an eclogitic facies. This follows if the selection is made from the current petrologic hypotheses for the constitution of the upper mantle. In fig. 15 densities computed by CLARK and RINGWOOD (1964) for a mantle composed of pyrolite (peridotite or dunite would give about the same values) and eclogite. Only the eclogite model is consistent with our results between 80 and 150 km. Either model is acceptable above this region and the pyrolite model is weakly favored near 300 km. A more extended discussion of these results can be found in another paper (PRESS 1969) where a hypothesis is proposed in which eclogite fractionation from the underlying, partially molten asthenosphere is involved in the creation and spreading of the sub-oceanic, rigid, lithospheric plate. BIRCH (1969) also interpreted these results to imply an eclogitic composition.

VI. The Transition Zone

Seismic array data have been used recently to establish rapid velocity changes near 400 and 700 km (see for example JOHNSON, 1967). These results have been incorporated in our models by fixing the compressional velocity and narrowing the range of permissible shear velocities at these depths to conform to the rapid increases, as seen in fig. 6. Although no such restrictions were placed on

the density values the rapid increase in density across the transition zone is evident on all models in fig. 7. This increase is due to compression, and to phase changes and possibly to composition changes. Phase transitions are ^{also}inferred from the laboratory verification of the olivine-spinel phase change at pressures corresponding to depths near 400 km and by the theoretical and experimental indications for a post-spinel phase transformation. (See for example, D.L. ANDERSON (1967) or H. FUJISAWA (1968)).

The occurrence of composition changes in the transition zone are more difficult to establish. BIRCH (1961) used the velocity change $\Delta\alpha$, and the density change $\Delta\rho$ to separately estimate the effects of phase and composition change. Using $\Delta\alpha$ and $\Delta\rho$ values for each model between 333 and 871 km, allowing .36 gm/cc for compression and using Birch's values for $(\partial\alpha/\partial\rho)_m$, $(\partial\rho/\partial m)_{T,P}$, $(\partial\alpha/\partial m)_g$, the change in mean atomic weight Δm was computed across the transition zone for each model. Those models with reduced densities in the asthenosphere ($\rho < 3.4$ gm/cc) showed an increase of 1-2 units in m . Thus for an asthenosphere with $m \sim 21$, and $Fe/Fe + Mg \sim 0.1$, as would be the case for peridotite or pyrolite, the $Fe/Fe + Mg$ ratio would increase to 0.2 or 0.3 across the transition zone. On the other hand, no increase in m was found for those models with a high density asthenosphere ($\rho > 3.5$ gm/cc). If the entire upper mantle is closer to eclogite in its iron content no increase in the $Fe/Fe + Mg$ ratio seems to be required across the transition zone. In a recent paper D.L. ANDERSON (1968) proposed that $\Delta m \sim 1.5$ and BIRCH (1961) gave $\Delta m \sim 1.0$ for one model.

VII. The Lower Mantle

Our results for the lower mantle rest heavily on Fairborn's independent determinations of a band of shear velocity distributions consistent with $dt/d\Delta$ and travel time data obtained at LASL. The range of shear velocities permitted by Fairborn's results is quite

narrow, as can be seen in fig. 6. This enables us to use eigenperiod data to constrain the density in the lower mantle to a greater degree than was possible before. Fairborn's shear velocity envelope shows a higher gradient than has usually been assumed (e.g. when compared to the Gutenberg model) and this requires a compensatory reduction in the density gradient in order to fit the spheroidal eigenperiod data. The results are shown in figs. 6 and 7. The density is constrained surprisingly well, to within about 0.2 gm/cc for most of the lower mantle. The density gradient is less than the adiabatic gradient as can be seen by comparison with the lower bound which approximates an adiabatic gradient. Figure 16 shows our band of solutions plotted against D.L. ANDERSON'S (1968) theoretical calculations for density of the solid solution series forsterite-fayalite and his summary of shock wave data. The band of density solutions is discordant with respect to profiles of constant composition, suggesting a change of mean atomic weight from 22-23 at the top of the lower mantle to 20-22 at the bottom of the mantle. This implies a depletion of iron with depth with the Fe/Fe + Mg ratio going from 0.2-0.3 to 0.1-0.2. Although superadiabatic temperature gradients might also account for the smaller density gradient, the augmented shear velocity gradient argues against this.

This can also be seen in figure 17 where the bulk sound velocity and density values for each model are plotted. The figure also shows the shock wave values for Twin Sisters dunite ($m=20.9$) and hortonolite dunite ($m=25.1$), as reduced by AHRENS, ANDERSON and RINGWOOD (1969). Although the data are scanty, a reduction in mean atomic weight from 22-23 at 871 km to 20-21 at 2898 km is indicated. The discordance with lines of constant composition seems too large to be accounted for by superadiabatic temperature gradients. WANG (1969) also suggested that these data might indicate decreasing m in the lower mantle. If subsequent studies do not establish Fairborn's shear velocity distribution as a world wide phenomenon, this conclusion will

have to be changed.

VIII. The Mantle

The c - ρ graph illustrates the main features of the mantle discussed earlier. The olivine-spinel phase transformation is evidenced between 371 km and 421 km by models with increasing c and ρ . Between 421 km and 621 km models with large increases in ρ and with little change in c could be interpreted as the result of compression and increasing iron content, the two effects having the same sign for ρ and opposite signs for c . The increase in c and ρ between 621 km and 721 km implies a phase change as the major feature. Decomposition of the ferro-magnesium-aluminum silicate to close packed simple oxides, or transformations to structures such as ilmenite or perovskite have been suggested for this region (BIRCH 1952, D.L. ANDERSON 1967, RINGWOOD 1969).

The distributions for ϕ and k/μ in the mantle are given in figs. 18 and 19.

XI. The Core

Results for the core are shown in figs. 8 and 20. The assumption of adiabaticity in the fluid core and the constraints imposed by the data prescribe the densities to the surprisingly small range of .25 gm/cc. The δ -ness criterion also indicates high resolving power for density at the top of the core. Using the shock wave data of McQUEEN and MARSH (1966) and BALCHAN and COWAN (1966) we see that iron alloyed with a miscible, abundant element such as silicon (15 wt %) would account for the core densities. There is no control on the density in the inner core with the data used here. Changes in the core radius ranged from -3 km to + 10 km.

The requirement for finite rigidity of the inner core was evidenced in an interesting way. Our procedure neglects core

rigidity and the only mode affected by this assumption, ${}_0S_2$ showed negative residuals for every model. Using ALSOP'S (1963) correction for rigidity resulted in a reduction of the residuals, for ${}_0S_2$, the largest discrepancy being .15%. This agrees with D.L. ANDERSON'S conclusions concerning rigidity in the inner core (1969).

X. Discussion

The question of uniqueness arises in all discussions of internal earth structure. The Backus-Gilbert δ -ness criterion demonstrates how stable local averages can be computed for α , β and ρ using eigenperiod data. However with currently available data the resolving power is adequate at too few places in the earth. Also the procedure does not yet allow for travel time or $dt/d\Delta$ data which under certain circumstances have high resolving power, nor does it consider errors in data. The band of solutions provided by the Monte Carlo method, if sufficiently narrow and if derived without bias from a large and representative selection of models, can under certain circumstances lead to meaningful conclusions. Unfortunately one is never quite sure that continued search would not reveal models significantly different from those already found, vitiating the conclusions. The use of physical arguments, laboratory experiments, theoretical-empirical equations of state will eventually provide powerful constraints. However the poor state of knowledge of the behavior of materials at internal earth pressures and temperatures, though improving rapidly, is a severe, current limitation.

With regard to the present paper and the other studies cited, we believe the following results are firm if the assumptions and data are correct:

1. The low velocity zone for shear waves in the sub-oceanic upper mantle. The parameterization included sufficiently few elements so that all possible models without a low velocity zone

could be tested and eliminated.

2. The high density for the lithosphere near 100 km. The δ -ness criterion has high resolving power for density at this depth and the narrow spread of Monte Carlo solutions indicate that errors in surface wave data of about ~~10%~~^{10%} do not weaken the constraint. As mentioned earlier, we believe this accuracy was achieved.

3. The rapid velocity increase near 400 km and its association with the olivine-spinel transformation: directly obtainable from $dt/d\Delta$ data (with minor depth uncertainty); the phase transformation was experimentally verified in the laboratory and olivine is almost certainly a major constituent of the upper mantle.

4. The rapid velocity increase near 700 km: directly obtainable from $dt/d\Delta$ data (with some depth uncertainty).

5. The density at the top of the core is between 9.9 and 10.2 gm/cc. The δ -ness criterion shows high resolving power and the spread of Monte Carlo solutions is small.

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Figures

Figure 1 - Flow diagram of Monte Carlo program during a run in which 531, 881 density models and 5025 shear velocity models were tested and yielded 11 successful models.

Figure 2 - Twenty-five shear velocity models of the mantle not subjected to geophysical constraints to test distribution of randomly generated models.

Figure 3 - Unconstrained mantle densities (see Figure 2 caption).

Figure 4 - Unconstrained core densities (see Figure 2 caption).

Figure 5 - Results for a typical model. Section A line 1: mass and moment of inertia; lines 2-4: p, Δ , theoretical times, model times and residuals for S and ScS. Section B: model eigenperiods and residuals against observed eigenperiods. Section C: model printout. Section D: depths at which parameters were varied and corresponding Q values.

Figure 6 - Twenty-seven successful shear velocity models for the mantle. Ticks on upper and lower bounds show where parameter was randomly varied.

Figure 7 - Density in the mantle (see Figure 6 caption).

Figure 8 - Density in the core (see Figure 6 caption).

Figure 9 - Tabulated parameters of successful models. Change in core radius shown in lower right corner. Each model has fixed crustal layers for depth, alpha, beta and rho as follows: 0., 1.52, 0. 1.03; 3., 6.55, 3.73, 2.84; 10, 6.55, 3.73, 2.84.

Figure 10 - See Figure 9

Figure 11 - See Figure 9

Figure 12 - See Figure 9

Figure 13 - See Figure 9.

Figure 14 - Differences in Love wave phase velocity data used by Hadden and Bullen and by Press which accounts for latter's requirement of low velocity zone. Points show how models fit the data.

Figure 15 - Successful density models for the upper mantle plotted together with Clark and Ringwood models for pyrolyte and eclogite.

Figure 16 - D. L. Anderson's theoretical models and his summary of experimental data plotted with our density solutions shown by the shaded band.

Figure 17 - Bulk sound velocity-density plot for successful models of the mantle together with static and shock wave data for dunite and forsterite-fayalite.

Figure 18 - Seismic parameter ϕ for the mantle obtained from successful models.

Figure 19 - Ratio K/μ for the mantle obtained from successful models.

Figure 20 - Band of core densities from successful models together with shock wave density data for Fe, Ni and Fe+ 19.8 wt.% Si.

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** BETA SLMD
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** 101- 600
** 601- 900
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** 1201-1500
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** 2401-2700
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** FAILED 110
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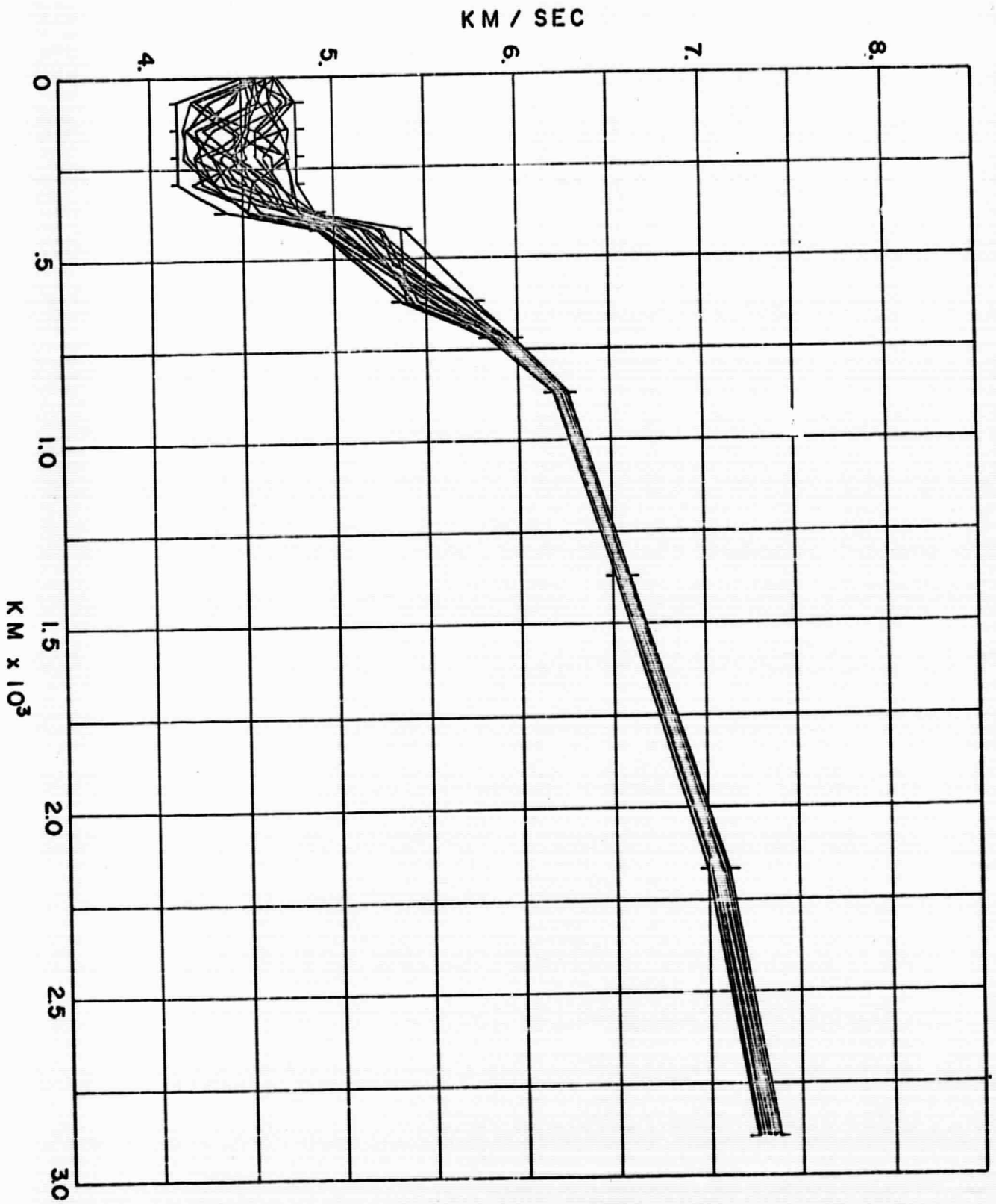
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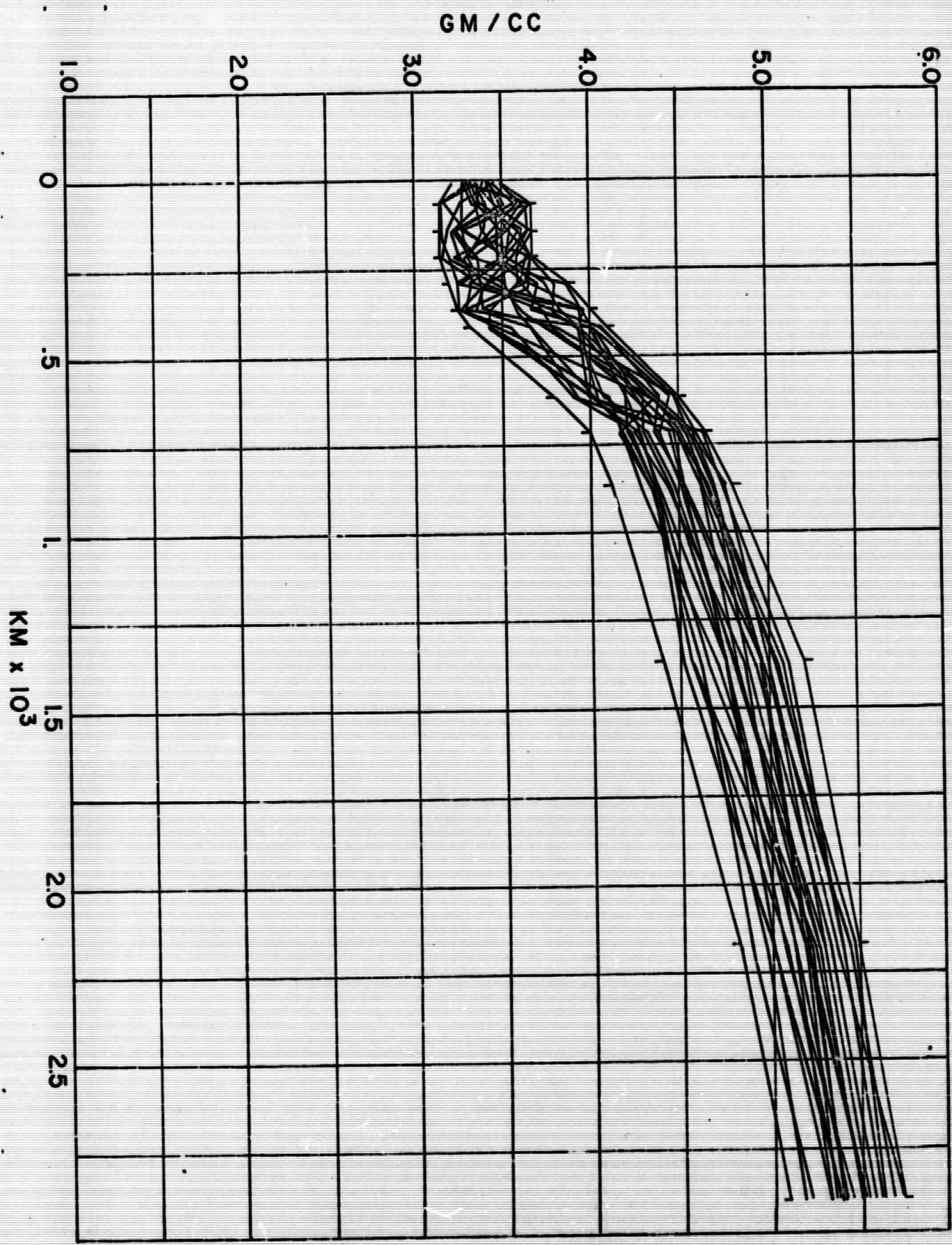
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BETA TTT	1561	1149	56	150	251	31	15	476	15	2	1319	0	0	0	0	0	0	0	0	0

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GM / CC

10

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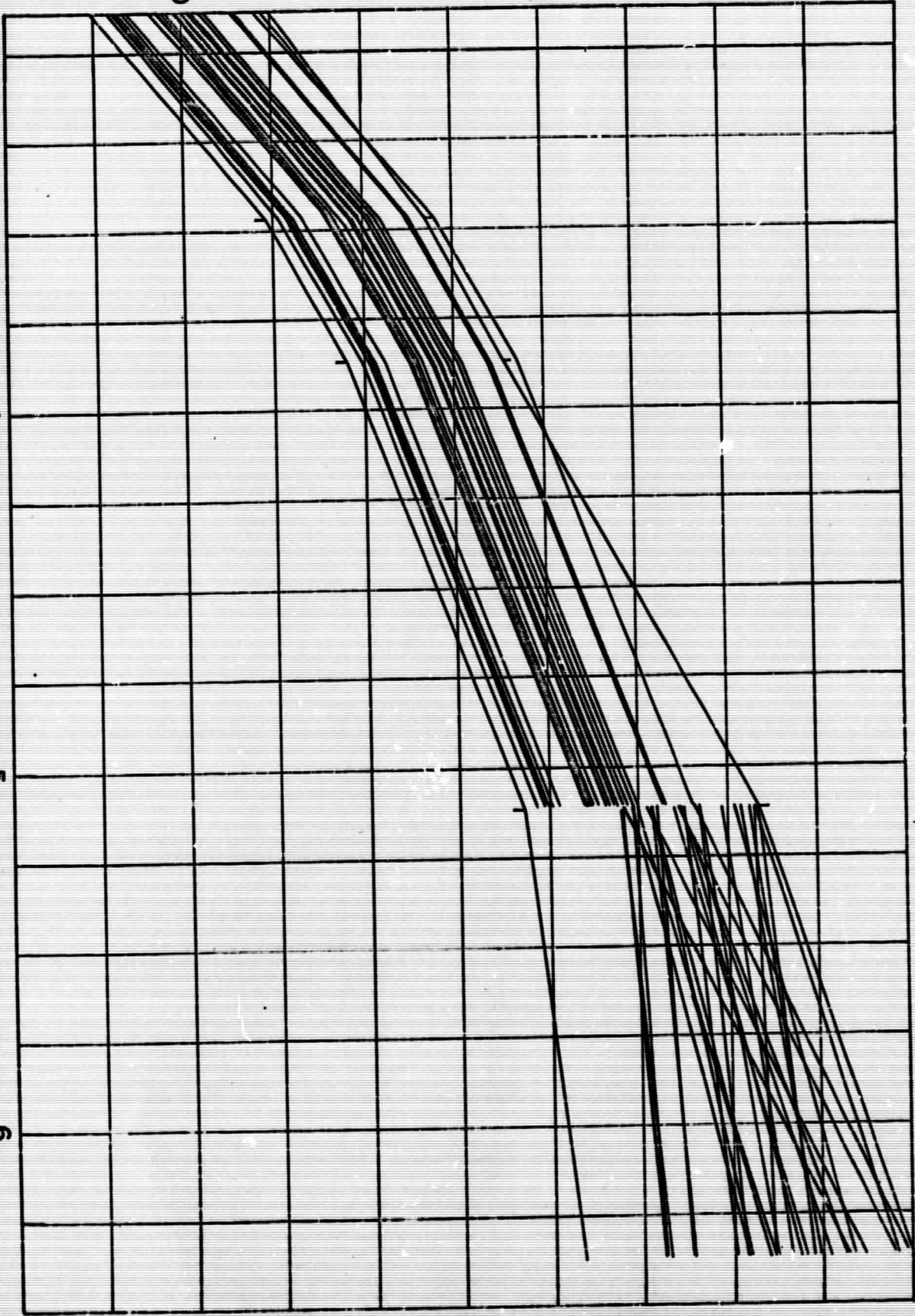
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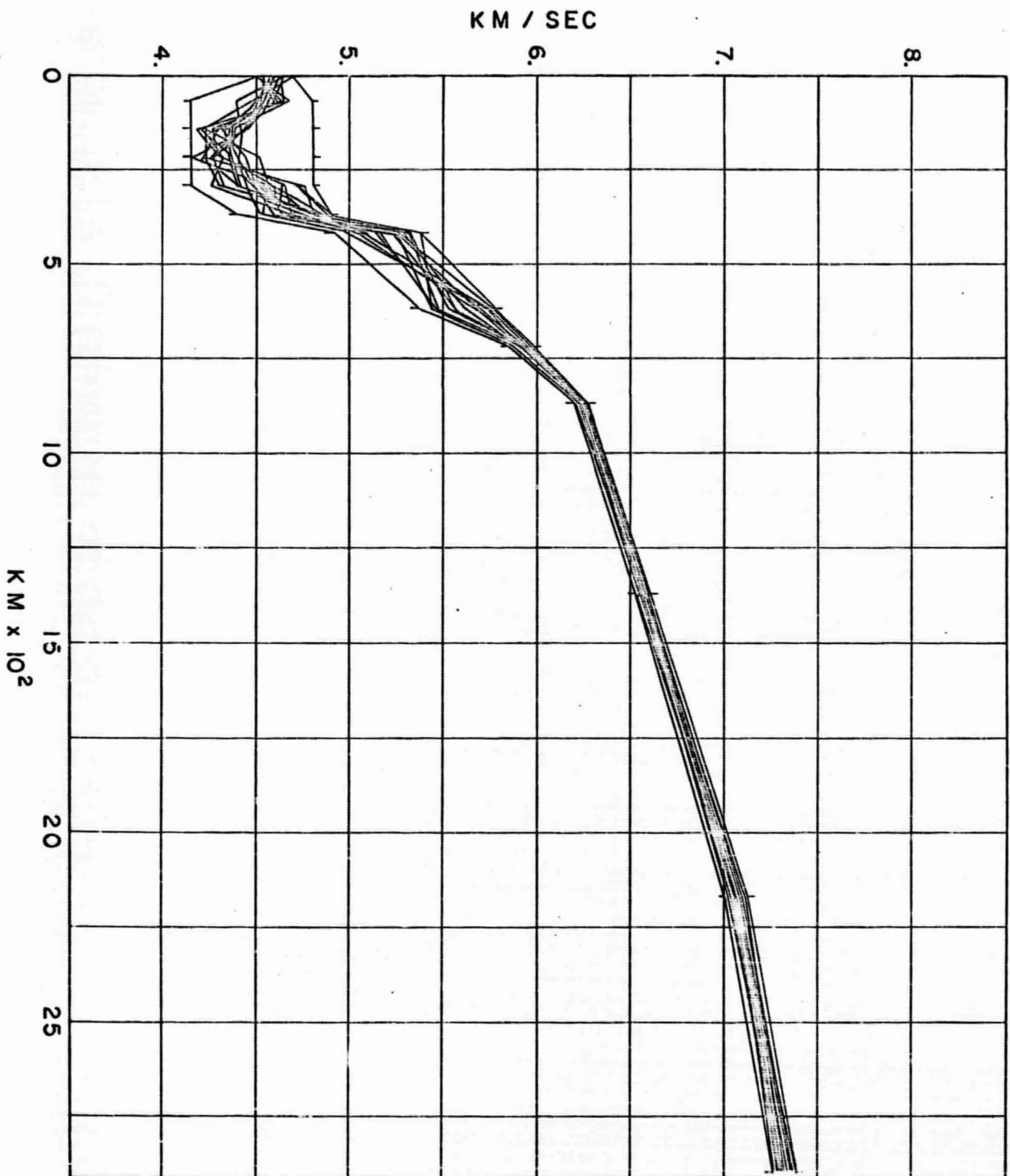
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 14.4 47.02 931.97 926.75 5.22 13.3 55.22 1045.72 1039.71 6.01 12.2 65.22 1173.89 1168.21 5.68
 11.1 72.63 1260.32 1256.00 4.32 10.0 83.69 1375.71 1374.27 1.44 8.9 94.14 1474.78 1470.89 3.89
 7.3 68.09 1264.38 1260.27 4.11
 0.5 2 3221.5 -7.5 4 1543.4 -3.6 6 963.7 -0.2 14 449.3 0.9 30 261.8 -0.7 50 178.7 0.5 70 134.0 -0.2
 * 12 504.2 -2.7 5 1189.4 0.4 7 813.0 1.3 8 709.2 1.6 9 635.7 1.7 10 581.4 1.4 11 538.9 1.4
 21 335.4 -0.4 13 474.7 0.6 15 427.0 0.8 16 407.3 0.8 17 389.8 0.4 18 374.1 0.7 19 347.1 -0.3
 36 229.9 -0.1 22 324.7 0.7 24 305.7 -0.0 26 289.2 -0.2 28 274.7 -0.3 32 250.2 -0.6 34 239.6 -0.4
 75 125.8 -0.5 19 360.0 -0.6 42 205.0 0.3 45 194.4 0.5 55 165.2 0.6 60 153.5 0.4 65 143.1 -0.1
 0 T 4 1308.4 4.4 8 738.2 1.2 11 575.9 0.2 40 201.1 -0.2 80 106.5 0.0 17 409.8 1.5 3 1708.5 0.5
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 13 505.7 1.0 14 477.3 0.3 16 429.9 0.3 18 391.7 1.5 19 375.2 0.7 20 360.1 -0.5
 21 346.2 -0.9 22 333.4 -0.9 23 321.6 0.1 25 300.3 0.1 27 281.7 0.1 30 257.8 -0.1 35 225.9 -0.3
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 0 R 1 1226.7 -0.9 3 1061.4 3.4 * 5 734.2 5.5 6 662.2 2.2 8 559.3 3.9 * 12 393.4 -3.4
 1 T 2 1473.5 1.5
 1 R
 2 T * 2 921.8 16.3 * 4 730.7 5.8 6 595.9 1.3 10 415.9 0.7
 2 S
 2 R
 1 6371.0 1.52 0.0 1.03 2 6368.0 1.52 0.0 4.58 3.48 * 3 6368.0 6.55 3.73 2.84 4 6361.0 6.55 3.73 2.84 *
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 17 6100.0 8.54 4.38 3.55 18 6075.0 8.58 4.40 3.56 19 6050.0 8.61 4.51 3.55 20 6025.0 8.65 4.61 3.54
 21 6000.0 8.68 4.71 3.53 22 5975.0 9.17 5.01 3.46 23 5950.0 9.67 5.31 3.39 24 5925.0 9.71 5.34 3.48
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 37 5600.0 11.05 6.06 4.50 38 5550.0 11.13 6.15 4.52 39 5500.0 11.21 6.23 4.55 40 5400.0 11.37 6.30 4.58
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 2 BETA REVERSALS BEFORE LAYERS 6 19 58 2898. 24.0
 2 RHO REVERSALS BEFORE LAYERS 22 24 59 2898. 54.8
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 9 71. 22.6 31 621. 23.0
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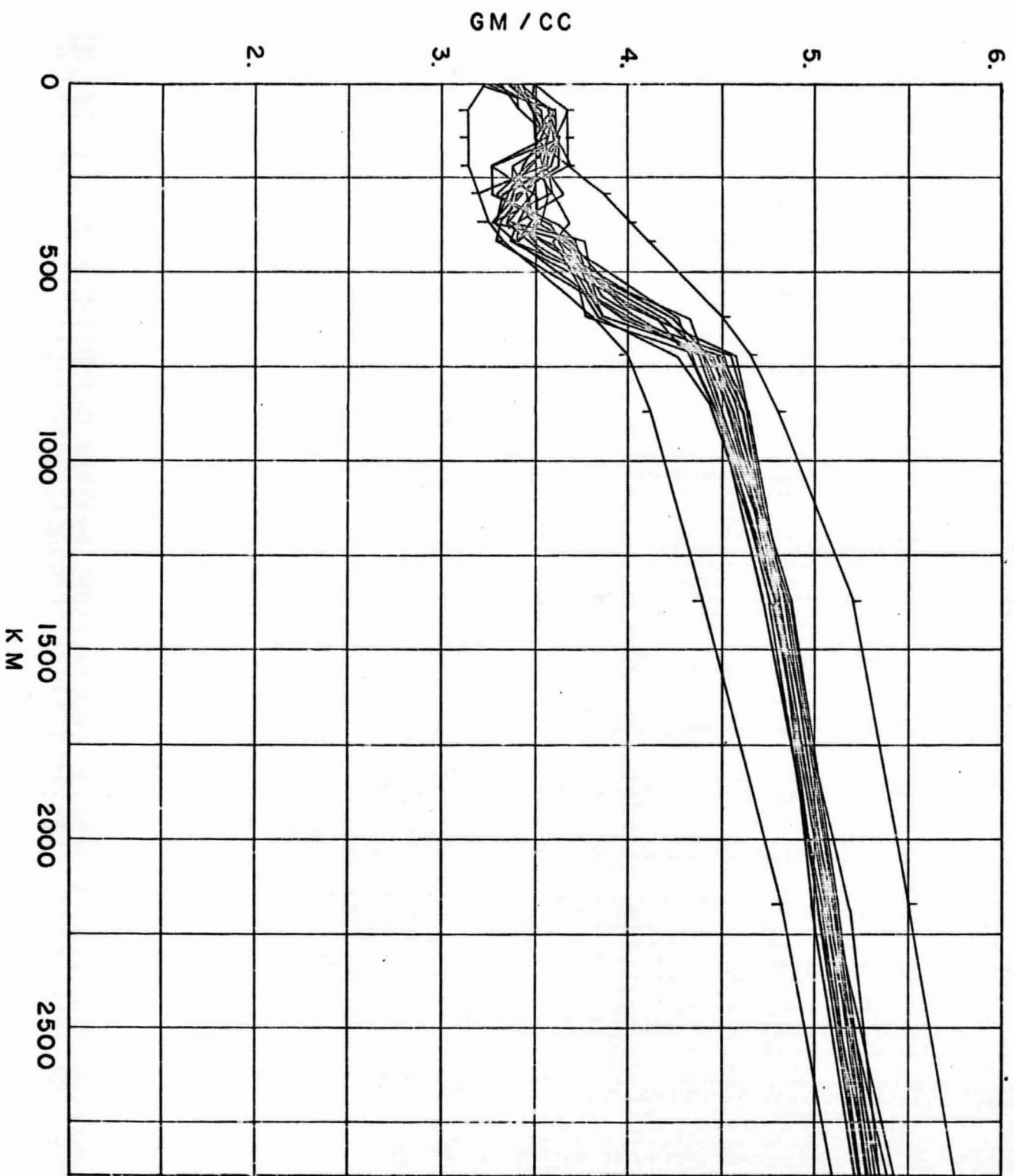
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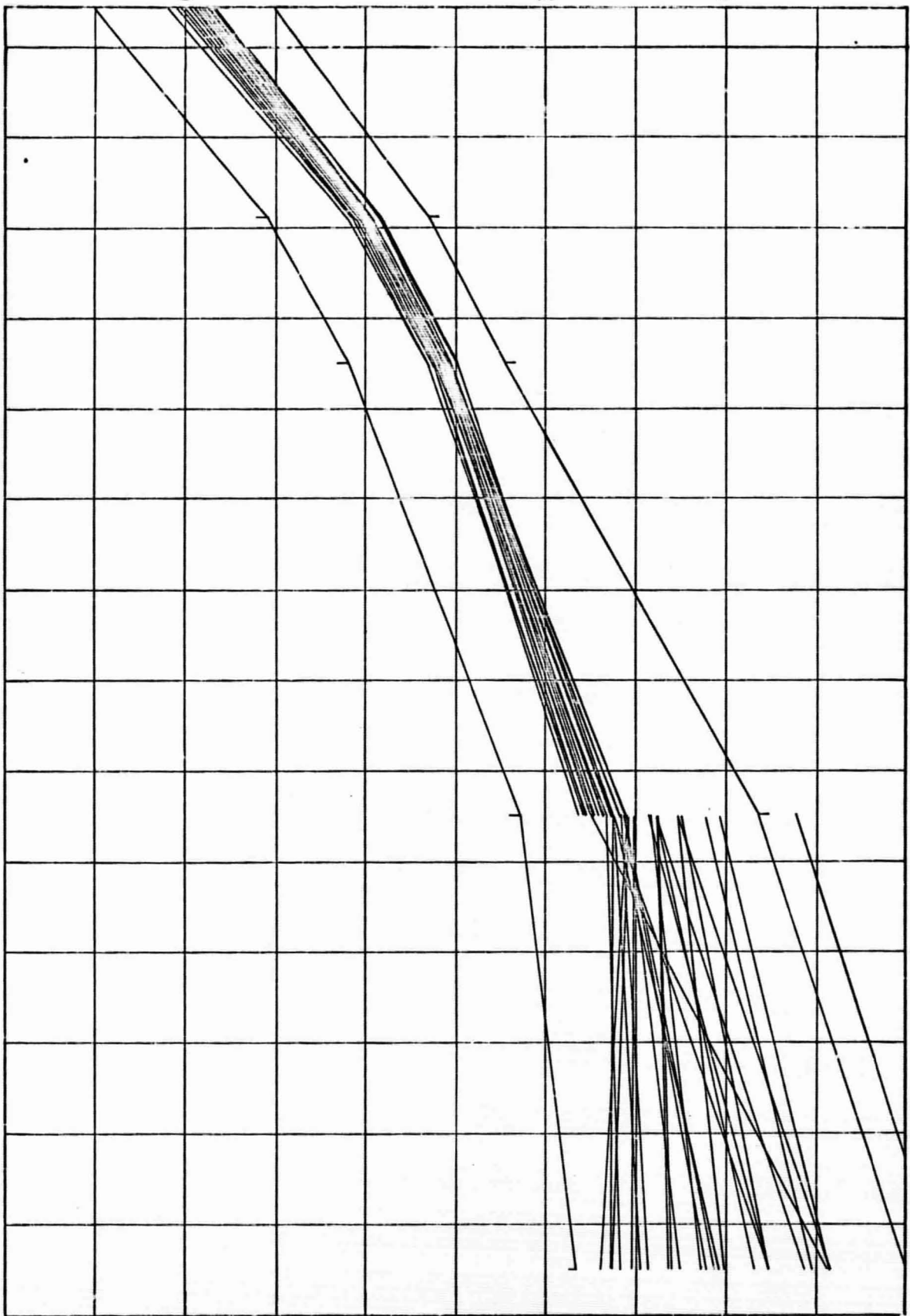
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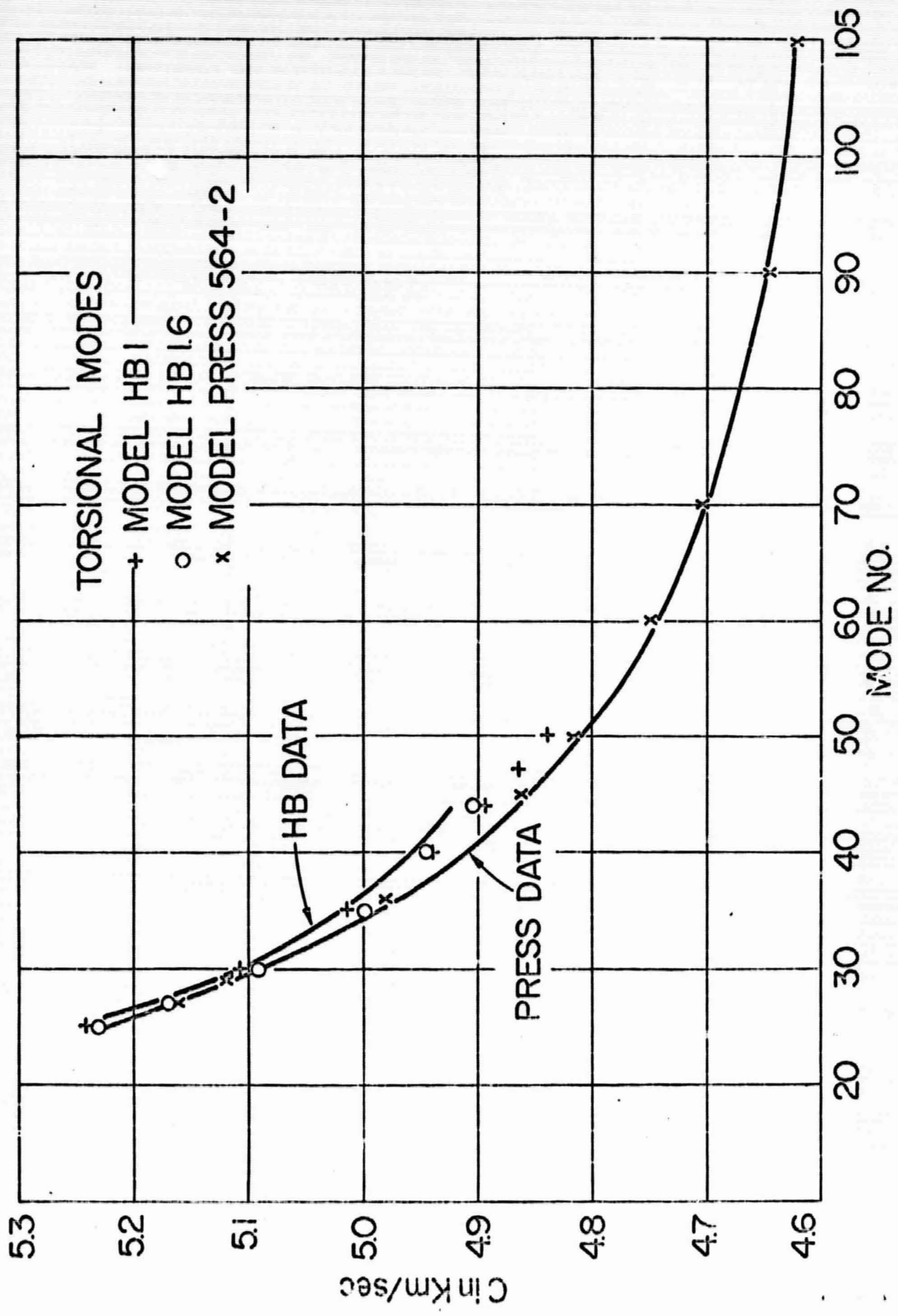


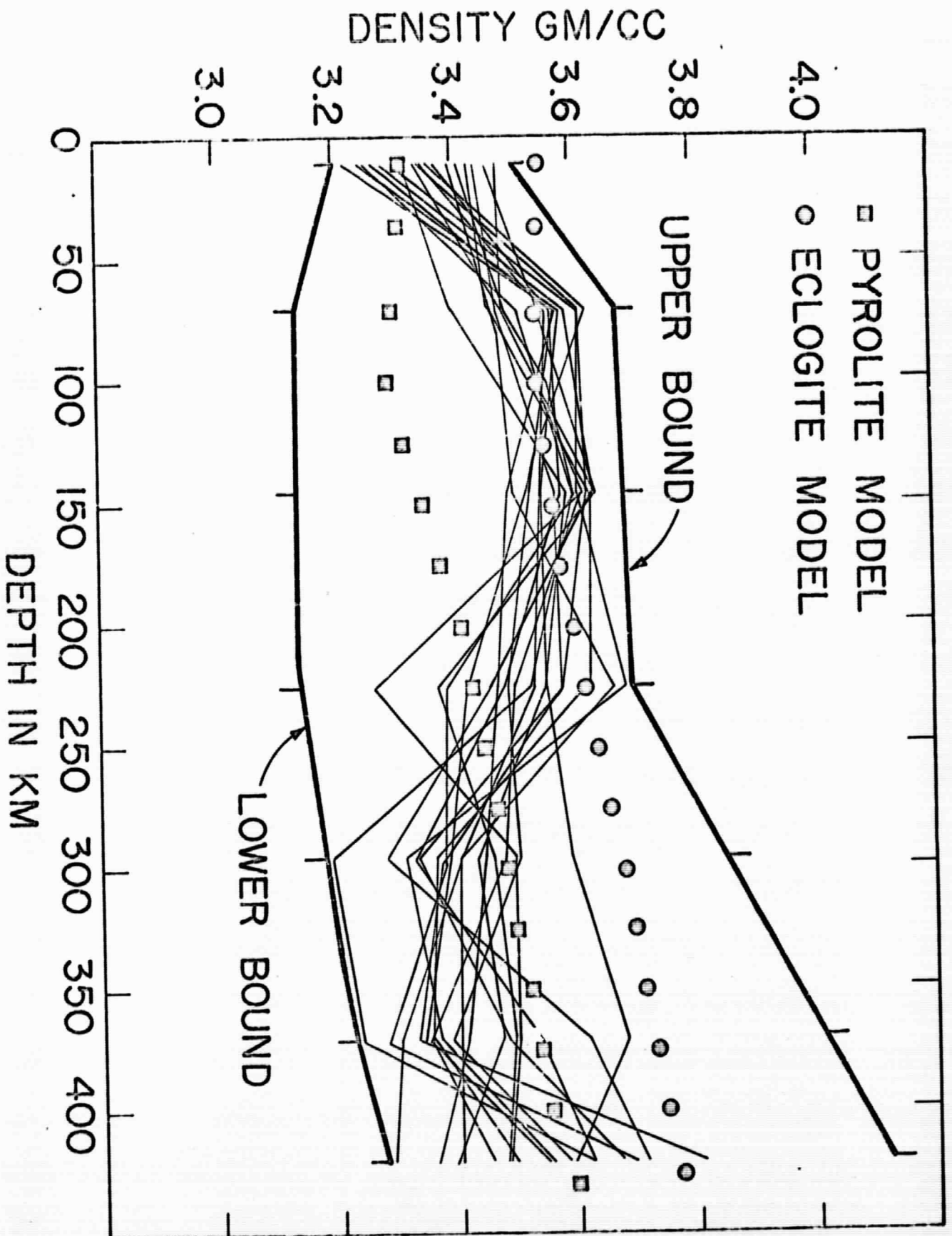
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71.	8.160	4.585	3.501	38.6	1.83	0.260	71.	8.160	4.585	3.572	38.6	1.81	0.269
146.	7.760	4.438	3.550	34.0	1.72	0.257	146.	7.760	4.451	3.531	33.8	1.71	0.255
221.	8.470	4.224	3.556	48.0	2.69	0.334	221.	8.470	4.209	3.498	48.1	2.72	0.334
295.	8.580	4.202	3.503	40.1	2.65	0.313	295.	8.580	4.412	3.504	47.7	2.45	0.320
371.	8.680	4.889	3.681	43.5	1.81	0.268	371.	8.680	4.947	3.888	44.0	1.87	0.273
421.	9.670	5.270	3.509	57.0	2.93	0.386	421.	9.680	5.124	3.582	58.5	2.73	0.385
621.	10.000	5.480	3.838	56.9	1.76	0.261	621.	10.000	5.473	4.036	57.1	1.77	0.263
721.	10.970	5.965	4.255	72.9	2.05	0.290	721.	10.970	5.930	4.448	73.5	2.09	0.294
871.	11.210	4.265	4.461	73.3	1.87	0.273	871.	11.210	4.207	4.548	74.3	1.81	0.279
1371.	12.020	6.571	4.962	96.9	2.01	0.287	1371.	12.020	6.589	4.747	86.6	1.80	0.285
2171.	13.030	7.057	5.116	103.4	2.08	0.292	2171.	13.030	7.078	5.053	103.0	2.06	0.291
2898.	13.770	7.277	5.272	119.0	2.25	0.306	2898.	13.770	7.270	5.390	119.1	2.25	0.307
2898.	7.960	0.0	10.000				2898.	7.960	0.0	10.087			
3471.	9.000	0.0	10.945				3471.	9.000	0.0	11.003			
5118.	10.150	0.0	12.389				5118.	10.150	0.0	11.424			
5118.	11.150	0.0	12.749				5118.	11.150	0.0	12.285			
6371.	11.100	0.0	13.231				6371.	11.100	0.0	12.745			
6.42													
10.	8.000	4.482	3.281	37.2	1.85	0.271	10.	8.000	4.529	3.354	36.7	1.70	0.276
71.	8.160	4.445	3.587	37.8	1.75	0.260	71.	8.160	4.529	3.571	39.3	1.93	0.276
146.	7.760	4.260	3.632	36.0	1.98	0.284	146.	7.760	4.308	3.525	35.5	1.67	0.277
221.	8.470	4.313	3.624	46.9	2.52	0.325	221.	8.470	4.439	3.533	45.5	2.31	0.311
295.	8.580	4.845	3.365	47.3	2.30	0.317	295.	8.580	4.404	3.402	46.7	2.31	0.311
371.	8.680	4.893	3.357	43.4	1.81	0.267	371.	8.680	4.698	3.209	45.9	2.08	0.293
421.	9.670	5.29	3.476	56.2	2.01	0.286	421.	9.670	5.229	3.286	57.1	2.00	0.293
621.	10.000	5.550	3.953	58.9	1.91	0.277	621.	10.000	5.641	3.980	57.6	1.91	0.287
721.	10.970	5.950	4.384	73.0	2.05	0.291	721.	10.970	5.874	4.561	76.3	2.15	0.290
871.	11.210	4.201	4.544	74.4	1.93	0.280	871.	11.210	6.212	4.640	74.2	1.92	0.278
1371.	12.020	6.571	4.821	96.9	2.01	0.287	1371.	12.020	6.541	4.832	87.4	2.04	0.289
2171.	13.030	7.051	5.102	103.5	2.09	0.293	2171.	13.030	7.052	5.088	103.9	2.08	0.293
2898.	13.770	7.262	5.244	117.3	2.17	0.300	2898.	13.770	7.336	5.324	117.9	2.19	0.302
2898.	7.960	0.0	10.135				2898.	7.960	0.0	10.027			
3471.	9.000	0.0	11.028				3471.	9.000	0.0	10.040			
3471.	4.600	0.0	11.428				3471.	4.600	0.0	11.361			
5118.	10.150	0.0	12.211				5118.	10.150	0.0	12.226			
5118.	11.150	0.0	12.370				5118.	11.150	0.0	12.730			
6371.	11.100	0.0	12.353				6371.	11.100	0.0	12.986			
3.38													
10.	8.000	4.540	3.298	36.3	1.74	0.256	10.	8.000	4.683	3.321	34.8	1.58	0.250
71.	8.160	4.533	3.612	39.2	1.91	0.277	71.	8.160	4.504	3.400	39.5	1.85	0.281
146.	7.760	4.427	3.674	36.1	1.74	0.259	146.	7.760	4.224	3.580	36.1	1.74	0.250
221.	8.470	4.267	3.600	47.7	2.64	0.332	221.	8.470	4.265	3.551	47.5	2.61	0.330
295.	8.580	4.820	3.318	46.4	2.27	0.308	295.	8.580	4.325	3.435	48.7	2.66	0.330
371.	8.680	4.781	3.366	44.9	1.96	0.282	371.	8.680	4.931	3.449	44.2	1.89	0.276
421.	9.670	5.226	3.455	57.1	2.00	0.294	421.	9.670	5.223	3.393	57.1	2.00	0.284
621.	10.000	5.541	4.195	59.1	1.92	0.278	621.	10.000	5.550	4.117	58.9	1.91	0.277
721.	10.970	5.961	4.403	73.1	2.06	0.291	721.	10.970	5.914	4.579	73.1	2.11	0.285
871.	11.210	4.256	4.521	73.5	1.88	0.274	871.	11.210	6.204	4.624	74.3	1.93	0.270
1371.	12.020	6.558	4.746	86.6	2.00	0.285	1371.	12.020	6.570	4.812	86.9	2.01	0.287
2171.	13.030	7.037	5.019	103.8	2.10	0.294	2171.	13.030	7.106	4.973	102.5	2.03	0.288
2898.	13.770	7.297	5.294	118.4	2.22	0.304	2898.	13.770	7.311	5.230	118.3	2.21	0.304
2898.	7.960	0.0	10.168				2898.	7.960	0.0	9.976			
3471.	9.000	0.0	11.098				3471.	9.000	0.0	10.968			
3471.	4.600	0.0	11.505				3471.	4.600	0.0	11.437			
5118.	10.150	0.0	12.336				5118.	10.150	0.0	12.426			
5118.	11.150	0.0	12.611				5118.	11.150	0.0	13.202			
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-2.37													
4.03													

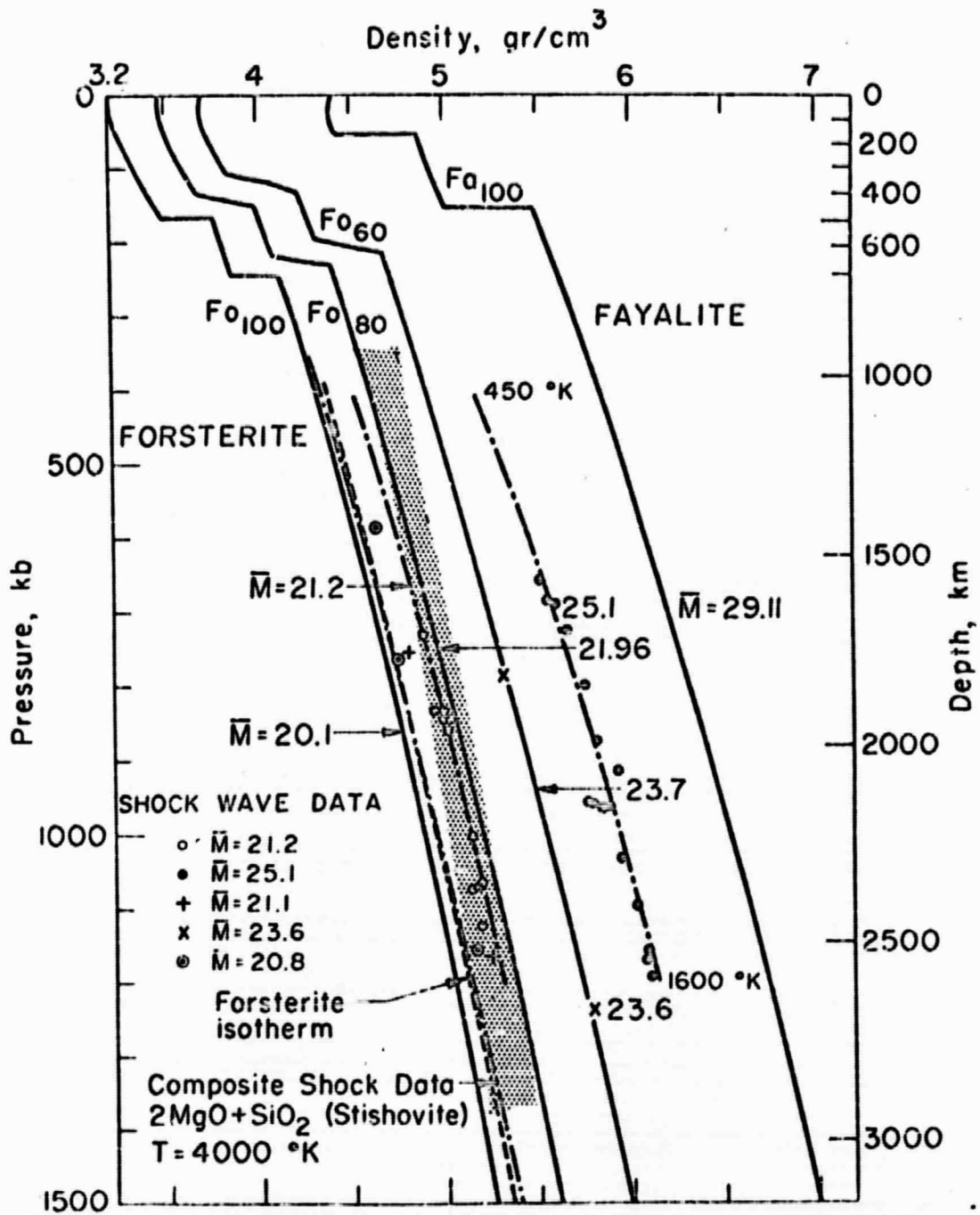
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71	9.160	4.603	3.562	39.6	1.84	0.282	71	8.000	4.572	3.384	35.2	1.65	0.245
146	8.470	4.783	3.532	46.6	2.47	0.322	146	7.760	4.808	3.516	39.6	1.80	0.280
221	8.580	4.403	3.585	47.8	2.68	0.321	221	8.470	4.783	3.532	46.6	2.44	0.310
721	9.680	4.572	3.525	45.7	2.05	0.291	721	8.580	4.528	3.258	46.6	2.23	0.285
421	9.670	5.312	3.305	55.0	1.98	0.286	421	8.680	4.686	3.572	45.9	2.08	0.293
621	10.000	5.561	4.110	58.8	1.90	0.276	621	9.000	5.150	3.644	58.0	2.18	0.361
721	10.970	5.970	4.110	72.2	2.03	0.280	721	10.000	5.445	4.264	60.5	2.04	0.280
871	11.210	6.232	4.549	73.0	1.90	0.274	871	10.970	5.922	4.337	73.3	2.07	0.292
1371	12.020	6.873	4.717	85.6	2.01	0.287	1371	11.210	6.236	4.444	73.8	1.90	0.274
2171	13.030	7.056	5.029	103.4	2.08	0.293	2171	12.020	6.580	4.803	85.6	1.90	0.285
2809	13.770	7.202	5.378	118.7	2.08	0.305	2809	13.030	7.042	5.108	103.2	2.00	0.294
2809	7.960	0.0	0.0	0.0	2.23	0.305	2809	13.770	7.394	5.271	117.9	2.19	0.307
3471	9.500	0.0	0.0	0.0	3.71	0.305	3471	2809	7.940	0.0	10.119	0.0	0.0
3871	10.150	0.0	0.0	0.0	7.77	0.305	3871	9.500	0.0	0.0	11.427	0.0	0.0
5118	11.150	0.0	0.0	0.0	7.75	0.305	5118	10.150	0.0	0.0	12.241	0.0	0.0
5118	11.150	0.0	0.0	0.0	3.75	0.305	5118	11.150	0.0	0.0	12.415	0.0	0.0
6371	11.100	0.0	0.0	0.0	7.77	0.305	6371	11.100	0.0	0.0	12.936	0.0	0.0
10	8.000	4.506	3.305	36.0	1.82	0.268	10	8.000	4.522	3.355	36.7	1.80	0.265
71	8.160	4.460	3.400	37.5	1.72	0.267	71	9.160	4.504	3.595	39.4	1.82	0.268
146	7.760	4.757	3.572	46.6	1.84	0.276	146	7.760	4.725	3.519	36.6	2.04	0.280
221	8.470	4.313	3.544	46.0	2.52	0.325	221	8.470	4.271	3.520	47.0	2.67	0.334
721	8.580	4.252	3.549	49.5	2.72	0.337	721	9.580	4.725	3.417	43.8	1.94	0.283
421	8.680	4.677	3.303	45.2	2.11	0.285	421	9.680	4.884	3.261	43.5	1.83	0.268
621	9.570	5.260	3.491	56.5	2.03	0.280	621	9.670	5.264	3.547	56.8	2.04	0.281
721	10.000	5.735	4.120	56.1	1.71	0.265	721	10.000	5.430	3.918	60.7	2.04	0.281
871	10.970	5.958	4.312	74.6	1.77	0.261	871	10.970	5.880	4.580	74.1	1.97	0.273
1371	11.210	6.197	4.467	74.5	1.94	0.280	1371	11.210	6.264	4.604	73.3	1.97	0.273
2171	13.030	6.585	4.777	86.7	2.05	0.286	2171	12.020	6.587	4.803	85.6	2.00	0.285
2809	13.770	7.067	5.102	103.2	2.07	0.292	2809	13.030	7.066	5.021	103.5	2.08	0.292
2809	7.960	0.0	0.0	0.0	2.25	0.307	2809	13.770	7.350	5.265	117.6	2.19	0.301
3471	9.500	0.0	0.0	0.0	7.77	0.307	3471	2809	7.960	0.0	10.076	0.0	0.0
3871	10.150	0.0	0.0	0.0	7.77	0.307	3871	9.500	0.0	0.0	10.946	0.0	0.0
5118	11.150	0.0	0.0	0.0	7.77	0.307	5118	10.150	0.0	0.0	12.244	0.0	0.0
5118	11.150	0.0	0.0	0.0	7.77	0.307	5118	11.150	0.0	0.0	12.556	0.0	0.0
6371	11.100	0.0	0.0	0.0	7.77	0.307	6371	11.100	0.0	0.0	13.220	0.0	0.0
71	8.000	4.610	3.359	35.7	1.68	0.261	71	8.000	4.638	3.213	35.3	1.64	0.247
146	9.160	4.501	3.482	39.6	1.95	0.281	146	8.160	4.593	3.594	39.5	1.82	0.268
146	7.760	4.276	3.563	36.6	2.04	0.280	146	7.760	4.275	3.581	36.6	2.06	0.285
221	8.470	4.513	3.542	44.4	2.10	0.302	221	8.470	4.220	3.508	47.0	2.48	0.334
721	8.580	4.553	3.316	45.0	2.20	0.303	721	9.580	4.717	3.337	43.0	1.84	0.263
421	8.680	4.643	3.201	46.4	2.14	0.300	421	9.680	4.955	3.281	43.0	1.84	0.263
621	9.570	5.036	3.306	59.8	2.37	0.315	621	9.670	5.141	3.636	56.3	2.20	0.303
721	10.000	5.672	4.322	57.1	1.77	0.263	721	10.000	5.450	3.927	60.3	2.00	0.280
871	10.970	5.952	4.395	73.1	2.04	0.281	871	10.970	5.919	4.494	73.3	2.08	0.285
1371	11.210	6.260	4.552	73.6	1.88	0.275	1371	11.210	6.264	4.644	73.8	2.08	0.275
2171	13.030	6.575	4.377	86.8	2.01	0.287	2171	12.020	6.585	4.808	85.7	2.00	0.284
2809	13.770	7.014	5.203	103.2	2.10	0.290	2809	13.030	7.029	5.028	104.1	2.00	0.284
2809	7.960	0.0	0.0	0.0	2.17	0.300	2809	13.770	7.356	5.235	117.6	2.17	0.306
3471	9.500	0.0	0.0	0.0	9.82	0.300	3471	7.960	0.0	0.0	10.136	0.0	0.0
3871	10.150	0.0	0.0	0.0	9.82	0.300	3871	9.500	0.0	0.0	11.702	0.0	0.0
5118	11.150	0.0	0.0	0.0	9.82	0.300	5118	10.150	0.0	0.0	12.360	0.0	0.0
5118	11.150	0.0	0.0	0.0	9.82	0.300	5118	11.150	0.0	0.0	12.619	0.0	0.0
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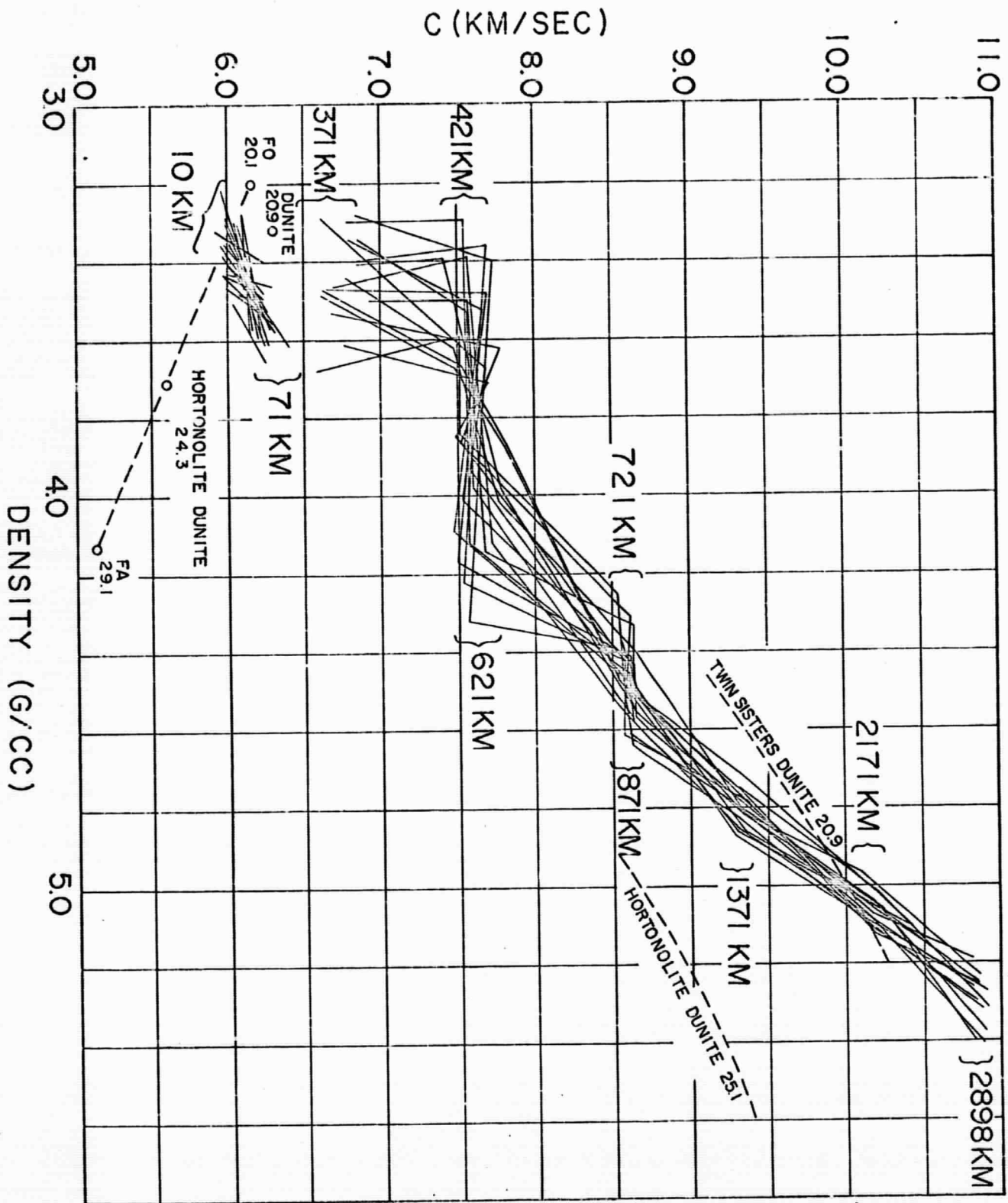
13.	8.000	4.574	3.419	36.1	1.77	0.257	10.	8.000	4.430	3.992	36.6	1.76	0.264	
14.	8.160	4.514	3.401	39.2	1.91	0.277	11.	8.160	4.490	3.575	39.6	1.96	0.282	
15.	7.760	4.311	3.421	35.4	1.91	0.277	12.	7.760	4.343	3.492	35.1	1.86	0.272	
16.	8.470	4.376	3.263	46.2	2.41	0.318	13.	8.470	4.350	3.472	46.5	2.44	0.321	
17.	8.580	4.356	3.508	45.3	2.21	0.304	14.	8.580	4.368	3.445	45.6	2.15	0.286	
18.	8.680	4.527	3.503	48.1	2.34	0.313	15.	8.680	4.620	3.498	46.9	2.20	0.302	
19.	8.270	4.520	3.278	57.3	1.83	0.296	16.	8.270	4.570	3.355	56.2	2.30	0.310	
20.	19.000	5.422	4.145	57.9	1.83	0.269	17.	19.000	5.070	3.325	56.7	1.74	0.259	
21.	19.970	5.869	4.332	74.6	2.17	0.300	18.	19.970	5.014	4.415	73.7	2.11	0.295	
22.	8.710	4.488	4.326	73.4	1.87	0.274	19.	8.710	4.605	4.074	73.4	1.88	0.274	
23.	12.020	4.507	4.868	73.4	1.99	0.284	20.	12.020	4.775	4.775	87.5	2.06	0.291	
24.	13.030	7.020	5.107	103.9	2.10	0.295	21.	13.030	7.085	5.098	102.9	2.05	0.290	
25.	7.370	7.375	5.764	117.1	2.15	0.299	22.	7.370	7.340	5.203	117.8	2.16	0.302	
26.	9.960	9.000	10.035				23.	9.960	10.031	10.031				
27.	9.600	9.0	11.379				24.	9.600	10.965	10.965				
28.	10.150	9.0	12.334				25.	10.150	11.403	11.403				
29.	11.150	9.0	12.723	5.70			26.	11.150	12.375	12.375	7.43			
30.	11.190	9.0	13.577				27.	11.190	12.374	12.374				
31.							28.			13.236				
32.	10.	8.000	4.587	3.432	35.9	1.71	10.	8.000	4.573	3.345	36.1	1.73	0.257	
33.	8.160	4.504	3.464	39.5	1.95	0.291	11.	8.160	4.540	3.611	39.1	1.90	0.276	
34.	7.760	4.338	3.620	35.1	1.87	0.273	12.	7.760	4.295	3.617	35.7	1.95	0.281	
35.	8.470	4.345	3.579	46.6	2.47	0.321	13.	8.470	4.243	3.503	47.7	2.65	0.332	
36.	8.580	4.522	3.366	46.4	2.27	0.308	14.	8.580	4.733	3.478	45.7	1.95	0.281	
37.	8.680	4.693	3.335	46.1	2.10	0.295	15.	8.680	4.822	3.348	44.3	1.91	0.277	
38.	8.270	5.676	5.678	59.0	2.28	0.309	16.	8.270	5.981	3.448	59.1	2.20	0.309	
39.	6.210	4.509	4.002	45.8	1.69	0.252	17.	6.210	4.604	4.012	54.8	1.75	0.260	
40.	7.210	10.970	4.386	73.1	2.05	0.291	18.	7.210	5.921	4.435	73.6	2.10	0.286	
41.	8.710	4.231	4.487	73.1	1.90	0.274	19.	8.710	5.195	4.563	74.5	1.94	0.260	
42.	12.020	6.559	4.642	87.1	2.03	0.288	20.	12.020	6.600	4.802	86.4	1.98	0.288	
43.	13.030	7.052	5.554	101.5	2.08	0.293	21.	13.030	7.005	5.047	104.4	2.13	0.297	
44.	2808.	13.770	7.371	5.242	117.2	2.16	22.	2808.	13.770	7.373	5.388	117.1	2.15	0.299
45.	2808.	7.960	9.0	10.150			23.	2808.	7.960	10.087	10.087			
46.	3871.	9.000	9.000	11.436			24.	3871.	9.000	11.002	11.002			
47.	5118.	10.150	9.0	12.219			25.	5118.	10.150	11.424	11.424			
48.	5118.	11.150	9.0	12.459	6.17		26.	5118.	11.150	12.443	12.443			
49.	5118.	11.190	9.0	12.881			27.	5118.	11.190	12.877	12.877	1.64		
50.							28.							
51.	10.	8.000	4.517	3.247	36.8	1.80	10.	8.000	4.625	3.437	35.5	1.66	0.249	
52.	8.160	4.612	3.614	38.2	1.80	0.265	11.	8.160	4.447	3.466	37.9	1.75	0.260	
53.	7.760	4.216	3.614	36.5	2.05	0.291	12.	7.760	4.191	3.547	35.8	2.00	0.286	
54.	8.470	4.301	3.532	47.1	2.54	0.326	13.	8.470	4.384	3.578	46.1	2.46	0.317	
55.	8.580	4.656	3.189	44.7	2.06	0.291	14.	8.580	4.424	3.342	44.5	2.43	0.310	
56.	4.680	4.689	3.237	44.5	2.09	0.294	15.	4.680	4.860	3.474	43.8	1.86	0.272	
57.	4.210	9.670	3.619	56.0	1.90	0.285	16.	4.210	9.670	3.671	56.3	2.02	0.297	
58.	6.210	10.000	4.014	60.1	2.01	0.297	17.	6.210	5.416	3.866	60.9	2.09	0.299	
59.	7.210	5.948	4.397	76.7	2.18	0.301	18.	7.210	5.958	4.445	73.0	2.04	0.290	
60.	8.710	11.210	4.235	73.8	1.90	0.274	19.	8.710	6.212	4.571	74.2	1.92	0.278	
61.	13.710	12.020	6.590	86.6	1.99	0.285	20.	13.710	6.549	4.792	87.3	2.04	0.289	
62.	21.710	13.030	5.095	103.5	2.08	0.293	21.	21.710	7.101	5.076	102.5	2.03	0.289	
63.	2808.	13.770	7.319	5.422	118.2	0.306	22.	2808.	7.349	5.293	117.6	2.19	0.301	
64.	2808.	7.960	9.0	10.035			23.	2808.	7.960	10.035	10.035			
65.	3871.	9.000	9.0	11.338			24.	3871.	9.000	11.430	11.430			
66.	5118.	10.150	9.0	12.227	10.33		25.	5118.	10.150	12.436	12.436	9.71		
67.	5118.	11.150	9.0	12.244			26.	5118.	11.150	12.613	12.613			
68.	5371.	11.190	9.0	13.547			27.	5371.	11.190	12.583	12.583			

11	10.	8.000	4.610	3.353	35.7	1.69	0.251
	71.	8.160	4.558	3.527	38.9	1.87	0.273
12	146.	7.760	4.428	3.559	34.1	1.74	0.259
	221.	8.470	4.151	3.580	48.8	2.83	0.342
13	296.	8.580	4.594	3.323	45.5	2.15	0.299
	371.	8.680	4.770	3.382	45.0	1.98	0.284
14	421.	9.670	5.050	3.762	59.5	2.33	0.312
	621.	10.000	5.740	3.849	56.1	1.70	0.254
15	721.	10.970	5.930	4.471	73.5	2.09	0.294
	871.	11.210	6.253	4.607	73.5	1.88	0.274
16	1371.	12.020	6.545	4.799	87.4	2.04	0.289
	2171.	13.030	7.025	5.007	104.0	2.11	0.295
17	2898.	13.770	7.377	5.238	117.1	2.15	0.299
	2898.	7.960	0.0	10.155			
18	3471.	9.000	0.0	11.076			
	3871.	9.600	0.0	11.501			
19	5118.	10.150	0.0	12.373	-3.41		
	5118.	11.150	0.0	12.483			
20	6371.	11.190	0.0	12.471			
21	10.	8.000	4.552	3.514	36.4	1.76	0.261
	71.	8.160	4.536	3.659	39.2	1.90	0.276
22	146.	7.760	4.332	3.598	35.2	1.88	0.274
	221.	8.470	4.377	3.572	46.2	2.41	0.318
23	296.	8.580	4.413	3.373	47.7	2.45	0.320
	371.	8.680	4.840	3.477	44.1	1.88	0.274
24	421.	9.670	5.318	3.684	55.8	1.97	0.283
	621.	10.000	5.515	4.041	59.4	1.95	0.281
25	721.	10.970	5.933	4.308	73.4	2.09	0.293
	871.	11.210	6.263	4.470	73.4	1.87	0.273
26	1371.	12.020	6.587	4.718	86.6	2.00	0.285
	2171.	13.030	7.048	5.086	103.5	2.08	0.293
27	2898.	13.770	7.343	5.341	117.7	2.18	0.301
	2898.	7.960	0.0	10.182			
28	3471.	9.000	0.0	11.076			
	3871.	9.600	0.0	11.476			
29	5118.	10.150	0.0	12.261	3.33		
	5118.	11.150	0.0	12.449			
30	6371.	11.190	0.0	12.300			
31	10.	8.000	4.601	3.289	35.8	1.69	0.253
	71.	8.160	4.557	3.542	38.9	1.87	0.273
32	146.	7.760	4.316	3.513	35.4	1.90	0.276
	221.	8.470	4.402	3.515	45.9	2.37	0.315
	296.	8.580	4.451	3.583	47.2	2.38	0.316
	371.	8.680	4.522	3.274	48.1	2.35	0.314
33	421.	9.670	5.252	3.686	56.7	2.06	0.291
	621.	10.000	5.660	3.769	57.3	1.79	0.264
34	721.	10.970	5.891	4.484	74.1	2.13	0.297
	871.	11.210	6.202	4.644	74.4	1.93	0.279
35	1371.	12.020	6.598	4.807	86.4	1.99	0.284
	2171.	13.030	7.047	5.051	103.6	2.00	0.293
36	2898.	13.770	7.281	5.218	118.9	2.24	0.306
	2898.	7.960	0.0	10.041			
37	3471.	9.000	0.0	10.996			
	3871.	9.600	0.0	11.450			
38	5118.	10.150	0.0	12.426	4.28		
	5118.	11.150	0.0	12.578			
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(KM/SEC)²

120

100

80

60

40

0

5

10

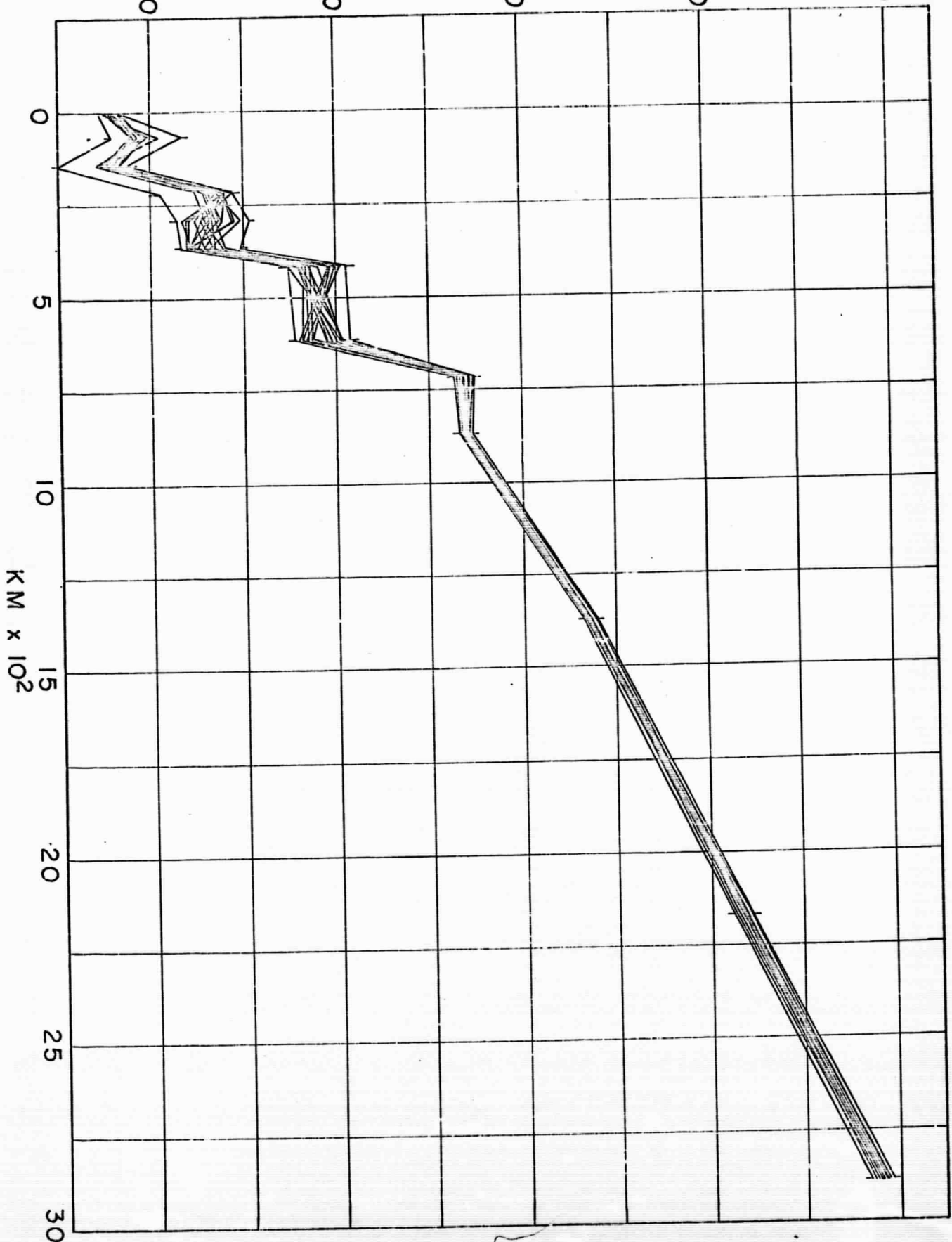
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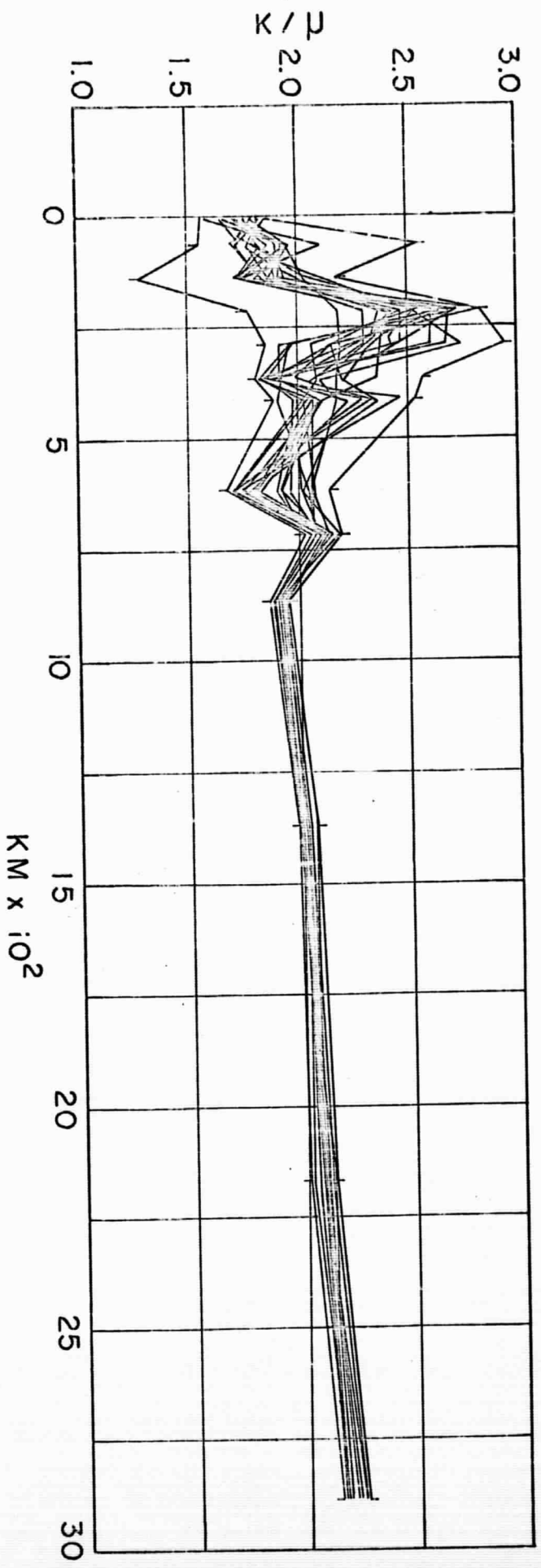
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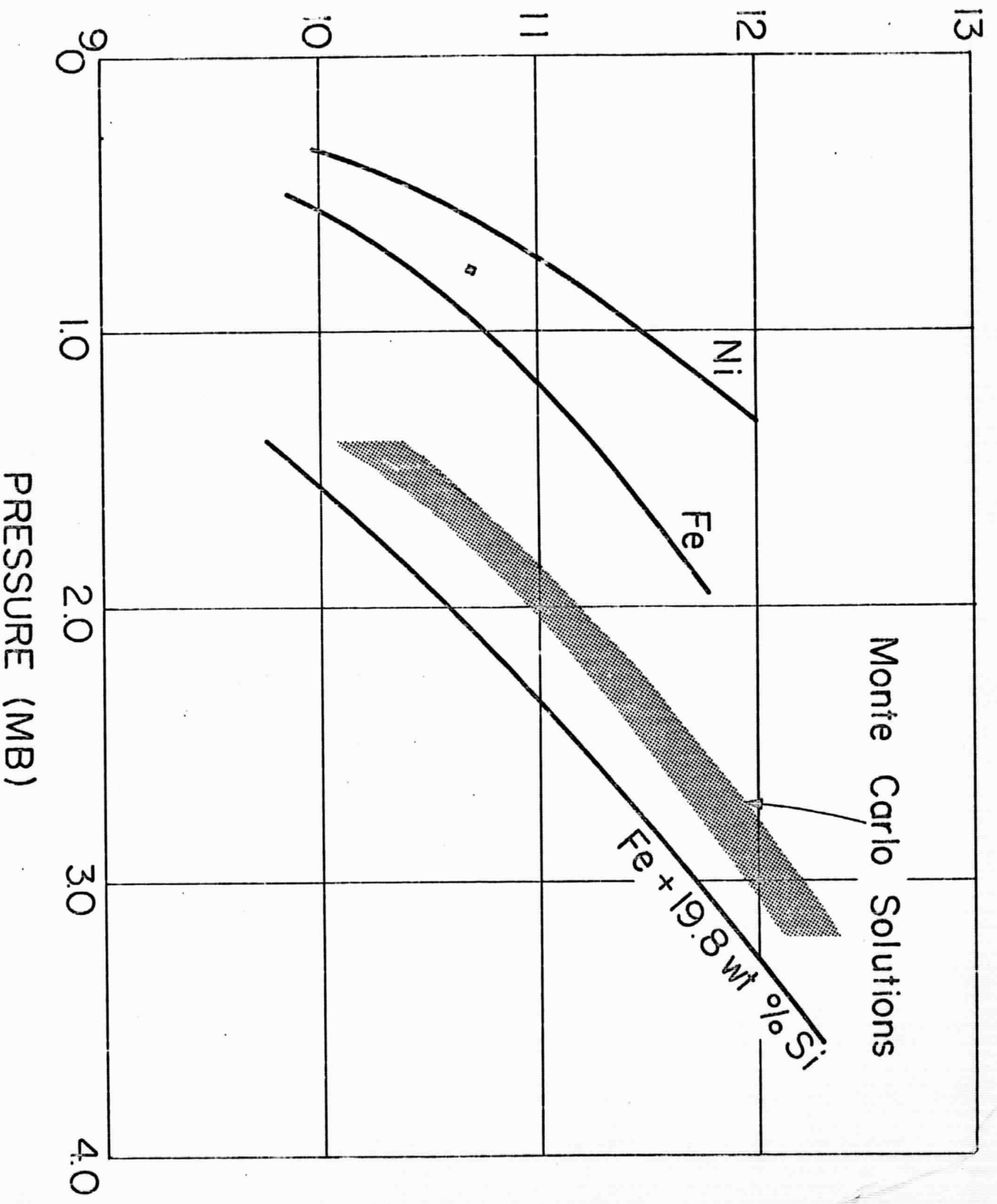
KM x 10²





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DENSITY G/CC



Monte Carlo Solutions

Ni

Fe

Fe + 19.8 wt % Si

PRESSURE (MB)