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OPTIMIZATION OF SELF-ACTING THRUST BEARINGS FOR LOAD CAPACITY AND STIFFNESS

by Bernard J. Hamrock Lewis Research Center Cleveland, Ohio 44135

NATIONAL AERONAUTICS AND SPACE ADMINISTRATION • WASHINGTON, D. C. • AUGUST 1970



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1.	Report No. NASA TN D-5954	2. Government Acc	ession No.	3. Recipie	nt's Catalo	og No.		
4.	Title and Subtitle OPTIMIZATION OF SELF.	ST BEARINGS	5. Report I Augus	Date t 1970				
	FOR LOAD CAPACITY AN		6. Perform	ing Organi	zation Code			
7.	Author(s) Bernard J. Hamrock		8. Performing Organization Report No E-5648					
9.	Performing Organization Name and Lewis Research Center	Address	1	10. Work Ur 129-0	uit No. 3			
	National Aeronautics and S	Space Administr	ation	11. Contrac	t or Grant	No.		
	Cleveland, Ohio 44135]1	13. Type of	Report an	d Period Covered		
12.	Sponsoring Agency Name and Address National Aeronautics and Space Administration			Technical Note				
	washington, D.C. 20040			14. Sponsor	ing Agenc	y Code		
15.	5. Supplementary Notes							
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17.	Key Words (Suggested by Autho	18. Distribution Stat	B. Distribution Statement					
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Springfield, Virginia 22151

OPTIMIZATION OF SELF-ACTING THRUST BEARINGS FOR LOAD CAPACITY AND STIFFNESS by Bernard J. Hamrock Lewis Research Center

SUMMARY

A linearized analysis of both infinite- and finite-width (rectangular) step thrust bearings was performed. Dimensionless load capacity and stiffness of the finite-width bearing are expressed in terms of a Fourier cosine series. The dimensionless load capacity and stiffness are a function of the dimensionless bearing number Λ , the pad length-to-width ratio λ , the film thickness ratio k, the step location parameter ψ , and the feed groove parameter η .

The equations obtained in the analysis were verified. The assumptions imposed were substantiated by comparison of the results with an existing exact solution for the infinite-width bearing. A digital computer program was developed which determines optimal bearing configuration for maximum load capacity or stiffness. Simple design curves are presented. Results are shown for both compressible and incompressible lubrication. Through a parameter transformation, the results are directly usable in designing optimal step sector thrust bearings.

INTRODUCTION

One of the first to apply the step film to a gas-lubricated bearing was Kochi (ref. 1). For the infinitely wide single thrust bearing, Kochi obtained an exact numerical solution. The expressions for the pressure were contained in a set of transcendental equations. Graphical methods were used to obtain the results.

In practical applications, one must use finite-width bearings. The finite-width step thrust bearing can appear in the shape of a rectangular pad or as a sector. For both the rectangular and the sector step thrust bearings, there is a definite need to know the optimal step configurations for maximum load capacity or maximum stiffness.

Ausman (ref. 2) in 1961 analyzed the gas lubricated step sector thrust bearing. He

applied linearization assumptions to the Reynolds equation, thereby enabling the pressure to be determined. With the pressure known, the load capacity was obtained. The expression for the load capacity appeared in terms of eigenvalues and Bessel functions. Ausman's results do not lend themselves readily to obtaining optimal step configurations for maximum load capacity or maximum stiffness when various bearing operating conditions are considered. The reason for this is the way in which parameters were made dimensionless and the nature of the resulting equations.

In this report, a rectangular step thrust bearing is analysed. Linearization assumptions comparable to those imposed by Ausman (ref. 2) are used. The sector bearing results are obtained directly from the rectangular step bearing results since curvature effects are shown to be very small. Because of the simplified nature of the resulting equations, a computer program was developed which optimizes step parameters for maximum load capacity or maximum stiffness for a wide range of bearing operating conditions. Results are shown for both compressible and incompressible lubrication. Therefore, the objective of this report is to present easily usable design information for finding optimal step bearings of rectangular or sector shape. The results are valid for a wide range of operating conditions.

SYMBOLS

A, B, D, E	integration constants
AA, BB, CC, DD, EE, FF, GG, HH	constants defined in eqs. (A1) to (A8)
b	width of rectangular thrust bearing
С	film thickness in ridge region
G	constant
h	film thickness
I _m	Fourier coefficient
J	separation constant
К	dimensionless stiffness, $-C(\partial W/\partial C)$
k	film thickness ratio, $(C + \Delta)/C$
L	$l_r + l_s + l_g$
l	length of ridge, step, or feed groove region depending on subscript
М	last odd positive integer used in evaluation of Fourier cosine series

m	odd positive integers
N	number of pads placed in overall length
Р	dimensionless pressure, $(p - p_a)/p_a$
р	pressure
p _a	ambient pressure
Q	mass flow rate
R _i	inner radius of sector thrust bearing
R _o	outer radius of sector thrust bearing
s	film thickness ratio as defined by Kochi (ref. 1), $1/(k-1)$
U	velocity of bearing surface
w	dimensionless load capacity of finite-width bearing, $w/p_a b(l_r + l_s + l_g)$
W_{∞}	dimensionless load capacity of infinite-width bearing, $w_{\infty}/p_a(l_r + l_s + l_g)$
w	load capacity
w _∞	load capacity per unit width
x	dimensionless length coordinate, x/b
\mathbf{X}_{∞}	x/L
x	coordinate in direction of motion
Y	dimensionless width coordinate, y/b
У	coordinate in direction of width of bearing
β	dimensionless bearing number used by Kochi (ref. 1), $3\mu U(l_s + l_r)/p_a \Delta^2 = \Lambda/(k-1)^2$
Δ	depth of step
η	feed groove parameter, $(l_r + l_s)/(l_r + l_s + l_g)$
Λ	dimensionless bearing number, $6\mu \text{Ub} / \text{p}_a \text{C}^2$
Λ_{∞}	$6\mu U(l_s + l_r + l_g)/p_a C^2$
λ	ratio of length to width of pad, $(l_r + l_s + l_g)/b$
μ	viscosity of fluid
^ξ r	$\sqrt{\left(\Lambda/2\right)^2 + m^2 \pi^2}$

$$\xi_{\rm s} = \sqrt{\left(\Lambda/2k^2\right)^2 + m^2\pi^2}$$

 ρ_{a} mass density of lubricant

- ψ step location parameter, $l_s/(l_r + l_s + l_g)$
- ω angular velocity

Subscripts:

- g denotes feed groove region
- r denotes ridge region
- s denotes step region
- ∞ denotes infinite-width analysis

BEARING DESCRIPTION

Sketch (a) shows the rectangular step thrust bearing to be studied. The ridge region



is where the film thickness is C, and the step region is where the film thickness is $C + \Delta$. The feed groove is the deep groove separating the end of the ridge region and the beginning of the next step region. Although not shown in this sketch, the depth of the feed groove is orders of magnitude deeper than the film thickness C. A pad is defined as the section that includes a ridge, a step, and feed groove regions. The length of the feed groove is small relative to the length of the pad. It should be noted that each pad acts independently since the pressure profile is broken at the lubrication feed groove.

LINEARIZATION ASSUMPTIONS

The Reynolds equation for the steady-state isothermal gas-lubricated thrust bearing can be written as

$$\frac{\partial}{\partial \mathbf{x}} \left(\mathbf{p} \ \frac{\partial \mathbf{p}}{\partial \mathbf{x}} \right) + \frac{\partial}{\partial \mathbf{y}} \left(\mathbf{p} \ \frac{\partial \mathbf{p}}{\partial \mathbf{y}} \right) = \frac{6\mu \mathbf{U}}{\mathbf{h}^2} \ \frac{\partial \mathbf{p}}{\partial \mathbf{x}}$$
(1)

Expanding and rearranging terms results in

$$\frac{\partial^2 p}{\partial x^2} + \frac{\partial^2 p}{\partial y^2} - \frac{6\mu U}{ph^2} \frac{\partial p}{\partial x} = -\frac{1}{p} \left[\left(\frac{\partial p}{\partial x} \right)^2 + \left(\frac{\partial p}{\partial y} \right)^2 \right]$$
(2)

In order to obtain an analytic solution to the foregoing equation, linearization assumptions will be imposed. The first linearization assumption states that the right side of equation (2) is zero. The validity of this assumption was shown by Scheinberg (ref. 3): For $U \rightarrow \infty$, the terms in equation (2) other than the one containing U become insignificant. Therefore, they can be neglected. Since the right side of equation (2) becomes zero for the two limiting cases, he concludes that it is a reasonable approximation to neglect these terms for all values of U. A second and final linearization assumption required is that p, where it appears as a coefficient, be replaced by the ambient pressure p_a . When these assumptions are applied, equation (2) becomes

$$\frac{\partial^2 \mathbf{p}}{\partial \mathbf{x}^2} + \frac{\partial^2 \mathbf{p}}{\partial \mathbf{y}^2} = \frac{6\mu U}{\mathbf{p}_2 \mathbf{h}^2} \frac{\partial \mathbf{p}}{\partial \mathbf{x}}$$
(3)

PRESSURE ANALYSIS OF FINITE-WIDTH BEARING

From equation (3), the linearized Reynolds equation can be written separately for the ridge and step regions of the finite-step thrust bearing:

$$\frac{\partial^2 \mathbf{p_r}}{\partial \mathbf{x}^2} + \frac{\partial^2 \mathbf{p_r}}{\partial \mathbf{y}^2} = \frac{6\mu U}{\mathbf{p_a} C^2} \frac{\partial \mathbf{p_r}}{\partial \mathbf{x}}$$
$$\frac{\partial^2 \mathbf{p_s}}{\partial \mathbf{x}^2} + \frac{\partial^2 \mathbf{p_s}}{\partial \mathbf{y}^2} = \frac{6\mu U}{\mathbf{p_a} (C + \Delta)^2} \frac{\partial \mathbf{p_s}}{\partial \mathbf{x}}$$

The subscripts r, s, and g refer to the ridge region (see sketch (a)), the step region, and the feed groove region, respectively. With x = bX, y = bY, $p_r = p_a(P_r + 1)$, and $p_s = p_a(P_s + 1)$, the foregoing equations can be rewritten as

$$\frac{\partial^2 \mathbf{P}_{\mathbf{r}}}{\partial \mathbf{X}^2} + \frac{\partial^2 \mathbf{P}_{\mathbf{r}}}{\partial \mathbf{Y}^2} = \Lambda \frac{\partial \mathbf{P}_{\mathbf{r}}}{\partial \mathbf{X}}$$
(4)

$$\frac{\partial^2 \mathbf{P}_s}{\partial \mathbf{X}^2} + \frac{\partial^2 \mathbf{P}_s}{\partial \mathbf{Y}^2} = \frac{\Lambda}{\mathbf{k}^2} \frac{\partial \mathbf{P}_s}{\partial \mathbf{X}}$$
(5)

where

$$\Lambda = \frac{6\mu \,\mathrm{Ub}}{\mathrm{p_a C}^2}$$

and

 $k = \frac{C + \Delta}{C}$

Using a separation-of-variables technique on equations (4) and (5) gives

$$\mathbf{P}_{\mathbf{r}} = \exp\left(\frac{\Lambda \mathbf{X}}{2}\right) \left[\mathbf{A}_{\mathbf{r}} \exp\left(\mathbf{X} \sqrt{\frac{\Lambda^2}{4} + \mathbf{J}_{\mathbf{r}}^2}\right) + \mathbf{B}_{\mathbf{r}} \exp\left(-\mathbf{X} \sqrt{\frac{\Lambda^2}{4} + \mathbf{J}_{\mathbf{r}}^2}\right) \right] \left[\mathbf{D}_{\mathbf{r}} \sin(\mathbf{J}_{\mathbf{r}} \mathbf{Y}) + \mathbf{E}_{\mathbf{r}} \cos(\mathbf{J}_{\mathbf{r}} \mathbf{Y}) \right]$$
(6)

$$P_{s} = \exp\left(\frac{\Lambda X}{2k^{2}}\right) \left[A_{s} \exp\left(X\sqrt{\frac{\Lambda^{2}}{4k^{4}} + J_{s}^{2}}\right) + B_{s} \exp\left(-X\sqrt{\frac{\Lambda^{2}}{4k^{4}} + J_{s}^{2}}\right)\right] \left[D_{s} \sin(J_{s}Y) + E_{s} \cos(J_{s}Y)\right]$$
(7)

The boundary conditions can be written as

when X = 0

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(1)
$$\mathbf{P}_{\mathbf{s}} = 0$$
 when $\mathbf{X} = 0$
(2) $\mathbf{P}_{\mathbf{r}} = 0$ when $\mathbf{X} = \frac{l_{\mathbf{s}} + l_{\mathbf{r}}}{b} = \left(\frac{l_{\mathbf{s}} + l_{\mathbf{r}}}{l_{\mathbf{s}} + l_{\mathbf{r}} + l_{\mathbf{g}}}\right) \left(\frac{l_{\mathbf{s}} + l_{\mathbf{r}} + l_{\mathbf{g}}}{b}\right) = \eta \lambda$

(3)
$$\mathbf{P}_{\mathbf{r}} = \mathbf{P}_{\mathbf{S}} = \sum_{m=1, 3, \ldots}^{\infty} \mathbf{I}_{m} \cos(m\pi \mathbf{Y})$$
 when $\mathbf{X} = \frac{l_{\mathbf{S}}}{b} = \psi \lambda$

where $\ensuremath{I_m}$ is a Fourier coefficient

- (4) $\frac{\partial \mathbf{P}_{\mathbf{r}}}{\partial \mathbf{Y}} = \frac{\partial \mathbf{P}_{\mathbf{S}}}{\partial \mathbf{Y}} = \mathbf{0}$ when $\mathbf{Y} = \mathbf{0}$
- (5) $P_r = P_s = 0$ when $Y = \frac{1}{2}$
- (6) $Q_r = Q_s$ when $X = \psi \lambda$

With the use of boundary conditions (1) to (5), equations (6) and (7) can be written as

$$\mathbf{P}_{\mathbf{r}} = \sum_{m=1,3,\ldots}^{\infty} \left\{ \begin{array}{l} I_{m} \cos(m\pi \mathbf{Y}) \exp\left[\frac{\Lambda}{2} (\mathbf{X} - \psi\lambda)\right] \\ \exp\left(-\lambda\psi\xi_{\mathbf{r}}\right) - \exp\left[-\lambda\xi_{\mathbf{r}}(2\eta - \psi)\right] \end{array} \right\} \left\{ \exp\left(-\mathbf{X}\xi_{\mathbf{r}}\right) - \exp\left[-\xi_{\mathbf{r}}(2\lambda\eta - \mathbf{X})\right] \right\}$$

$$\mathbf{P}_{\mathbf{S}} = \sum_{\mathbf{m}=1,3,\ldots}^{\infty} \left\{ \begin{bmatrix} \mathbf{I}_{\mathbf{m}} \cos(\mathbf{m}\pi \mathbf{Y}) \exp\left[\frac{\Lambda}{2\mathbf{k}^{2}} (\mathbf{X} - \psi\lambda)\right] \\ \exp(-\psi\lambda\xi_{\mathbf{S}}) - \exp(\psi\lambda\xi_{\mathbf{S}}) \end{bmatrix} \exp(-\mathbf{X}\xi_{\mathbf{S}}) - \exp(\mathbf{X}\xi_{\mathbf{S}}) \end{bmatrix}$$
(9)

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(8)

where

$$\xi_{\mathbf{r}} = \sqrt{\left(\frac{\Lambda}{2}\right)^2 + \mathbf{m}^2 \pi^2}$$

and

$$\xi_{\rm s} = \sqrt{\left(\frac{\Lambda}{2{\rm k}^2}\right)^2 + {\rm m}^2\pi^2}$$

The linearized equations describing the mass flow across the ridge and step regions can be written as

$$Q_{r} = \frac{\rho_{a}}{p_{a}} \left(\frac{p_{r}UC}{2} - \frac{p_{a}C^{3}}{12\mu} \frac{\partial p_{r}}{\partial x} \right)$$
$$Q_{s} = \frac{\rho_{a}}{p_{a}} \left[\frac{p_{s}U(C + \Delta)}{2} - \frac{p_{a}(C + \Delta)^{3}}{12\mu} \frac{\partial p_{s}}{\partial x} \right]$$

These equations may be made dimensionless by letting $p_r = p_a(P_r + 1)$, $p_s = p_a(P_s + 1)$, and x = bX as was done for the Reynolds equations. Therefore,

$$Q_{r} = \frac{\rho_{a} p_{a} C^{3}}{12\mu b} \left[\Lambda(P_{r} + 1) - \frac{\partial P_{r}}{\partial X} \right]$$
$$Q_{s} = \frac{\rho_{a} p_{a} (C + \Delta)^{3}}{12\mu b} \left[\frac{\Lambda}{k^{2}} (P_{s} + 1) - \frac{\partial P_{s}}{\partial X} \right]$$

Making use of boundary conditions (3) and (6) gives

$$k^{3} \left(\frac{\partial \mathbf{P}_{s}}{\partial \mathbf{X}} \right)_{\mathbf{X}=\lambda\psi} - \left(\frac{\partial \mathbf{P}_{r}}{\partial \mathbf{X}} \right)_{\mathbf{X}=\lambda\psi} = \Lambda(k-1) \left[1 + \sum_{m=1,3,\ldots}^{\infty} \mathbf{I}_{m} \cos(m\pi \mathbf{Y}) \right]$$
(10)

With the use of equations (8) to (10), the Fourier coefficient I_m can be solved:

$$I_{m} = \frac{4(k-1)\sin\left(\frac{m\pi}{2}\right)}{m\pi\left(\frac{1-k}{2} + \frac{\xi_{s}k^{3}}{\Lambda}\left[\frac{1+\exp(-2\xi_{s}\lambda\psi)}{1-\exp(-2\xi_{s}\lambda\psi)}\right] + \frac{\xi_{r}}{\Lambda}\left\{\frac{1+\exp\left[-2\lambda\xi_{r}(\eta-\psi)\right]}{1-\exp\left[-2\lambda\xi_{r}(\eta-\psi)\right]}\right\}\right)}$$
(11)

LOAD ANALYSIS OF FINITE-WIDTH BEARING

The dimensionless load capacity for the ridge and step region is

$$W_{r} = \frac{W_{r}}{p_{a}bL} = \frac{2}{\lambda} \int_{0}^{1/2} \int_{\lambda\psi}^{\lambda\eta} P_{r} dX dY$$
$$W_{s} = \frac{W_{s}}{p_{a}bL} = \frac{2}{\lambda} \int_{0}^{1/2} \int_{0}^{\lambda\psi} P_{s} dX dY$$

Substituting equations (8) and (9) into the foregoing equations and integrating gives

$$W_{\mathbf{r}} = \sum_{m=1,3,\ldots}^{\infty} \frac{2I_{m} \sin\left(\frac{m\pi}{2}\right)}{m^{3}\pi^{3}\lambda} \left(\frac{\Lambda}{2} + \frac{\xi_{\mathbf{r}} \left\{ 1 - 2 \exp\left[-\lambda(\eta - \psi)\left(\xi_{\mathbf{r}} - \frac{\Lambda}{2}\right)\right] + \exp\left[-2\lambda\xi_{\mathbf{r}}(\eta - \psi)\right] \right\}}{1 - \exp\left[-2\lambda\xi_{\mathbf{r}}(\eta - \psi)\right]} \right)$$

(12)

$$W_{s} = \sum_{m=1,3,\ldots}^{\infty} \frac{2I_{m}\sin\left(\frac{m\pi}{2}\right)}{m^{3}\pi^{3}\lambda} \left(\frac{\xi_{s}\left\{1-2\exp\left[-\lambda\psi\left(\xi_{s}+\frac{\Lambda}{2k^{2}}\right)\right]+\exp(-2\lambda\xi_{s})\right\}}{1-\exp(-2\lambda\psi\xi_{s})}\right)$$
(13)

The total dimensionless load supported by the rectangular step slider bearing is

$$W = \frac{W_r + W_s}{p_a Lb} = W_r + W_s$$
(14)

STIFFNESS ANALYSIS OF FINITE-WIDTH BEARING

The equation for the dimensionless stiffness is

$$\mathbf{K} = -\mathbf{C} \, \frac{\partial \mathbf{W}}{\partial \mathbf{C}}$$

With the use of equations (12) to (14), the foregoing equation becomes

$$K = \sum_{m=1,3,\ldots}^{\infty} \left\{ W(AA) - \frac{2I_{m}\sin\left(\frac{m\pi}{2}\right)}{m^{3}\pi^{3}\lambda} \left[N\left(1 - \frac{1}{k^{3}}\right) - \xi_{r}(BB) - \xi_{s}(CC) + \frac{\Lambda^{2}}{2\xi_{r}}(DD) + \frac{\Lambda^{2}}{2\xi_{s}k^{5}}(EE) \right] \right\}$$
(15)

where AA, BB, CC, DD, and EE are constants defined in the appendix.

Therefore, with equations (11) to (15), the dimensionless load capacity and stiffness for a self-acting gas-lubricated finite-width step thrust bearing is completely defined. From these equations, it is evident that the dimensionless load capacity and stiffness are functions of the following parameters:

Dimensionless bearing number

$$\Lambda = \frac{6\mu \text{ Ub}}{p_a \text{C}^2}$$

Length-to-width ratio of pad

$$\lambda = \frac{l_{\rm s} + l_{\rm r} + l_{\rm g}}{\rm b}$$

Film thickness ratio

$$k = \frac{C + \Delta}{C}$$

Step location parameter

$$\psi = \frac{l_{\rm s}}{l_{\rm s} + l_{\rm r} + l_{\rm g}}$$

Feed groove parameter

$$\eta = \frac{l_{s} + l_{r}}{l_{s} + l_{r} + l_{g}}$$

INFINITE-WIDTH-BEARING ANALYSIS

The linearized Reynolds equations for the ridge and step regions of an infinitely wide thrust bearing can be written in dimensionless form from equations (4) and (5) as

$$\frac{\partial^2 \mathbf{P}_{\mathbf{r},\infty}}{\partial \mathbf{X}_{\infty}^2} = \Lambda_{\infty} \frac{\partial \mathbf{P}_{\mathbf{r},\infty}}{\partial \mathbf{X}_{\infty}}$$
$$\frac{\partial^2 \mathbf{P}_{\mathbf{s},\infty}}{\partial \mathbf{X}_{\infty}^2} = \frac{\Lambda_{\infty}}{\mathbf{k}^2} \frac{\partial \mathbf{P}_{\mathbf{s},\infty}}{\partial \mathbf{X}_{\infty}}$$

where

$$\Lambda_{\infty} = \frac{6\mu U(l_{s} + l_{r} + l_{g})}{p_{a}C^{2}}$$

and

$$X_{\infty} = \frac{x}{l_{s} + l_{r} + l_{g}}$$

Proceeding in much the same manner as in the Finite-Analysis section of this report, one arrives at the following expressions for the dimensionless pressure for the ridge and step regions:

$$\mathbf{P}_{\mathbf{r},\infty} = \mathbf{G}\left\{\mathbf{1} - \exp\left[\Lambda_{\infty}(\mathbf{X}_{\infty} - \eta)\right]\right\}$$
(16)

$$\mathbf{P}_{\mathbf{s},\infty} = \mathbf{G} \left\{ \frac{1 - \exp\left[\Lambda_{\infty}(\psi - \eta)\right]}{1 - \exp\left(\frac{\Lambda_{\infty}\psi}{k^2}\right)} \right\} \left[1 - \exp\left(\frac{\Lambda_{\infty}\mathbf{X}_{\infty}}{k^2}\right) \right]$$
(17)

where

$$G = \frac{k - 1}{1 + k \exp\left(\frac{-\Lambda_{\infty}\psi}{k^2}\right) \left\{\frac{1 - \exp\left[\Lambda_{\infty}(\psi - \eta)\right]}{1 - \exp\left(\frac{-\Lambda_{\infty}\psi}{k^2}\right)\right\}}$$

With the pressure known, the resulting dimensionless load capacity of an infinitewidth step thrust bearing can be easily formulated as

$$W_{\infty} = G\left(\eta - \psi \left\{ \frac{1 - \exp\left(\frac{-\Lambda_{\infty}\psi}{k^{2}}\right) \exp\left[\Lambda_{\infty}(\psi - \eta)\right]}{1 - \exp\left(\frac{-\Lambda_{\infty}\psi}{k^{2}}\right)} \right\} + \frac{(k^{2} - 1)}{\Lambda_{\infty}} \left\{1 - \exp\left[\Lambda_{\infty}(\psi - \eta)\right]\right\}\right)$$
(18)

The dimensionless stiffness can be written as

$$K_{\infty} = -C \frac{\partial W_{\infty}}{\partial C} = \frac{W_{\infty}(HH)}{G} + G\left(-\psi(GG) + \frac{2(k-1)}{\Lambda_{\infty}} \left\{1 - \exp\left[-\Lambda_{\infty}(\eta - \psi)\right]\right\} - 2(k^{2} - 1)(\eta - \psi)\exp\left[-\Lambda_{\infty}(\eta - \psi)\right]\right)$$
(19)

where FF, GG, and HH are constants defined in the appendix. For the infinite-width analysis, the dimensionless load and stiffness are functions of the following parameters:

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Dimensionless bearing number for infinite-width bearing

$$\Lambda_{\infty} = \frac{6\mu \,\mathrm{U}(l_{\mathrm{s}} + l_{\mathrm{r}} + l_{\mathrm{g}})}{p_{\mathrm{a}} \mathrm{C}^2}$$

Film thickness ratio

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$$k = \frac{C + \Delta}{C}$$

Step location parameter

$$\psi = \frac{l_s}{l_s + l_r + l_g}$$

Feed groove parameter

$$\eta = \frac{l_{s} + l_{r}}{l_{s} + l_{r} + l_{g}}$$

VERIFICATION OF EQUATIONS

The equations for the dimensionless load capacity and stiffness for the finite- (eqs. (11) to (15)) and infinite- (eqs. (18) and (19)) width analyses were programmed on a digital computer. It should be recalled that linearization assumptions were imposed in order to obtain simplified results. Figure 1 shows that these assumptions are generally valid for the infinite-width bearing. This figure compares the results from the present work with Kochi's (ref. 1) exact solution. The agreement is good. Comparing equations (2) and (3) with and without the width coordinates y, one could further conclude that for a finite-width bearing the linearized analysis should be in good agreement with the exact results.

Table I shows that the solutions for the finite and infinite analyses approach each other when the length-to-width ratio approaches zero. The results, when 500 terms (M = 1001) are used in the Fourier cosine series, approach the infinite-width analysis much closer than when only 50 terms (M = 101) are used. Furthermore, the rate of convergence is much slower at large dimensionless bearing numbers $(\Lambda \rightarrow 500)$ than at smaller values of Λ . Note the decrease in dimensionless load capacity W when Λ increases from 100 to 500. This is due to the fact that the step parameters are held

constant; that is, the step parameters chosen are closer to the optimal for $\Lambda = 100$ than for $\Lambda = 500$.

Table II compares the dimensionless load capacity obtained from Ausman (ref. 2) with the present work for various dimensionless bearing numbers Λ and inner- to outer-radius ratios R_i/R_o . Ausman (ref. 2) considers curvature effects, whereas the present work does not. The equivalent length of a sector pad is assumed to be the arc length along the average radius. The width is the difference between inner and outer radii. For all inner- to outer-radius ratios, there is close agreement between the two analyses. Curvature effects are small. Therefore, the simplified equations of the present analysis are valid for evaluating the circular sector thrust bearing.

OPTIMIZING PROCEDURE

The problem as defined in the INTRODUCTION is to find the optimal step bearing for maximum load capacity or stiffness for various bearing numbers. This means, given the dimensionless bearing number Λ , finding the optimal length-to-width ratio λ , optimal film thickness ratio k, and optimal step location parameter ψ . The significance of the feed groove parameter η is much less than that of the other parameters. Therefore, for all evaluations, the feed groove parameter η will be set equal to 0.97.

The basic problem in optimizing λ , k, and ψ for maximum load and stiffness is essentially that of finding values of λ , k, and ψ that satisfy the following equations:

$$\frac{\partial \mathbf{W}}{\partial \lambda} = \frac{\partial \mathbf{W}}{\partial \mathbf{k}} = \frac{\partial \mathbf{W}}{\partial \psi} = \mathbf{0}$$
(20)

$$\frac{\partial \mathbf{K}}{\partial \lambda} = \frac{\partial \mathbf{K}}{\partial \mathbf{k}} = \frac{\partial \mathbf{K}}{\partial \psi} = \mathbf{0}$$
(21)

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The method used in solving equations (19) to (21) is the Newton-Raphson method for solving simultaneous equations. This method is described in reference 4 and in most other tests on numerical analysis.

Therefore, given the dimensionless bearing number Λ , the optimization computer program obtains optimum values of λ , k, and ψ for maximum dimensionless load capacity or stiffness. As a check on the optimization procedure, the following case was considered: $\Lambda = \lambda = 1 \times 10^{-5}$. This case approaches an infinitely-wide incompressibly lubricated step bearing for which the results are known. For this case, the computer program indicated that k = 1.866 and $\psi = 0.718$ were optimal for maximum load capacity. These results are in exact agreement with Archibald (ref. 5). For optimization of a step sector thrust bearing, parameters for the sector must be found that are analogous to those for the rectangular step bearing. Sketch (b) shows the transformation.



The following substitutions accomplish this transformation:

$$b \rightarrow R_{o} - R_{i}$$
$$N(l_{s} + l_{r} + l_{g}) \rightarrow \pi(R_{o} + R_{i})$$
$$U \rightarrow \frac{\omega}{2} (R_{o} + R_{i})$$

where N is the number of pads placed in the step sector. By use of the foregoing equations, the dimensionless bearing number can be rewritten as

$$\Lambda \rightarrow \frac{3\mu\omega \left(R_0^2 - R_1^2\right)}{p_a C^2}$$
(22)

The optimal number of pads to be placed in the sector is obtained from

$$N = \frac{\pi(R_{o} + R_{i})}{(\lambda)_{opt}(R_{o} - R_{i})}$$
(23)

where $(\lambda)_{opt}$ is the optimal value for the length-to-width ratio. The way $(\lambda)_{opt}$ is obtained is discussed in the next section. Since N will not be an integer normally, rounding it to the nearest integer is required.

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DISCUSSION OF RESULTS

Tables III and IV give optimal step parameters $(\psi, \lambda, \text{ and } k)$ for resulting maximum load capacity and stiffness. The differences between these tables are that table III optimizes with respect to load capacity, whereas table IV optimizes with respect to stiffness. The following observations can be made about both tables III and IV as Λ (or bearing speed) is increased:

(1) The length-to-width ratio λ increases; that is, the length of the pad increases relative to its width.

(2) The step location parameter ψ decreases; that is, the length of the step region decreases relative to the length of the pad.

(3) The film thickness ratio k increases; that is, the step depth increases relative to the clearance.

Sketch (c) shows optimal step bearings for two extremes of dimensionless bearing numbers.



Figures 2 to 4 are obtained directly from the data presented in tables III and IV. Figure 2 shows the effect of Λ on λ , k, and ψ for the maximum load capacity condition for a range of Λ from 0 to 410. The optimal step parameters (λ , k, and ψ) are seen to approach an asymptote as the dimensionless bearing number Λ becomes small. That is, for small $\Lambda(\Lambda \leq 0.1)$ the optimal step parameters are not a function of Λ . In the incompressible solution of a step bearing, the right side of equations (4) and (5) are zero. Therefore, it must be concluded that the asymptotic values that the step parameters approach in figure 2 correspond to the incompressible solution. These asymptotes are $\lambda = 0.918$, $\psi = 0.555$, and k = 1.693. Figure 3 shows the effect of Λ on λ , k, and ψ for the maximum stiffness condition for a range of Λ from 0 to 410. As in figure 2, the optimal step parameters are seen to approach asymptotes as the incompressible solution is reached. The asymptotes are $\lambda = 0.915$, $\psi = 0.557$, and k = 1.470. Note that there is a difference in the asymptote for the film thickness ratio but virtually no change in λ and ψ when compared to those obtained from figure 2.

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Figure 4 shows the effect of dimensionless bearing number Λ on dimensionless load capacity and stiffness. The difference in these figures is that the optimal step parameters are obtained in figure 4(a) for maximum load capacity and in figure 4(b) for maximum stiffness. Also shown in these figures are values of K and W when the step parameters are held fixed at the optimal values obtained for the incompressible solution. A significant decrease between the solid and dashed lines in W or K does not occur until $\Lambda > 8$.

Figures 2 to 4 contain all the necessary information for the design of an optimal rectangular step thrust bearing for maximum load capacity or stiffness. With the dimensionless bearing number Λ , the optimal values of λ , ψ , and k can be obtained from figure 2 or 3 depending on whether or not maximum load or stiffness is a major consideration. From figure 4(a) or (b), the resulting values for the dimensionless load and stiffness can be obtained. Furthermore, from figures 2 to 4 and equations (22) and (23), the optimal step sector thrust bearing with load capacity or stiffness considered can be obtained. The dimensionless bearing number Λ is obtained from figure 2 or 3 depending on whether maximum load or stiffness is a major (22). With Λ known, the optimal values of λ , ψ , and k can be obtained from figure 2 or 3 depending on whether maximum load or stiffness is a major consideration. From equation (23), the optimal number of pads placed in a sector can be determined. Finally, from figure 4, the resulting values for the dimensionless load capacity and stiffness can be obtained.

SUMMARY OF RESULTS

A linearized analysis of both infinite-width and finite-width (rectangular) step thrust bearings was performed. Dimensionless load capacity and stiffness of the finite-width bearing are expressed in terms of a Fourier cosine series. The equations obtained in the analysis were verified. The assumptions imposed were substantiated by comparison of the results with an existing exact solution for the infinite-width bearing. A digital computer program was developed which determines optimal bearing configuration for maximum load capacity or stiffness. Simple design curves are presented. Results are

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shown for both compressible and incompressible lubrication. Through a parameter transformation, the results are directly usable in designing an optimal step sector thrust bearing.

Lewis Research Center, National Aeronautics and Space Administration, Cleveland, Ohio, May 15, 1970, 129-03.

APPENDIX - CONSTANTS OBTAINED IN EVALUATING DIMENSIONLESS STIFFNESS

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Constants developed when deriving the dimensionless stiffness expressed herein. Equations (A1) to (A5) are obtained from the finite-width analysis, and equations (A6) to (A8) are obtained from the infinite-width analysis.

$$AA = \frac{-m\pi I_{m}}{4(k-1)\sin\left(\frac{m\pi}{2}\right)} \left[\frac{2\lambda\psi\Lambda}{k^{2}} \left\{ \frac{\exp(-2\xi_{s}\lambda\psi)}{\left[1 - \exp(-2\lambda\xi_{s}\psi)\right]^{2}} \right\} + \left(\frac{3k^{2}\xi_{s}}{\Lambda} - \frac{\Lambda}{2k^{2}\xi_{s}} \right) \left[\frac{1 + \exp(-2\xi_{s}\lambda\psi)}{1 - \exp(-2\xi_{s}\lambda\psi)} \right] \right]$$

$$+\left(\frac{3\xi_{\mathbf{r}}}{\Lambda}-\frac{\Lambda}{2\xi_{\mathbf{r}}}\right)\left\{\frac{1+\exp\left[-2\lambda\xi_{\mathbf{r}}(\eta-\psi)\right]}{1-\exp\left[-2\lambda\xi_{\mathbf{r}}(\eta-\psi)\right]}\right\}+2\lambda\Lambda(\eta-\psi)\left(\frac{\exp\left[-2\xi_{\mathbf{r}}\lambda(\eta-\psi)\right]}{\left\{1-\exp\left[-2\xi_{\mathbf{r}}\lambda(\eta-\psi)\right]\right\}^{2}}\right)\right]$$
(A1)

$$BB = \frac{2\lambda\Lambda(\eta - \psi)}{\left\{1 - \exp\left[-2\lambda\xi_{\mathbf{r}}(\eta - \psi)\right]\right\}^2} \left[\frac{\Lambda}{2\xi_{\mathbf{r}}} \left(2 \exp\left[-2\lambda\xi_{\mathbf{r}}(\eta - \psi)\right] - \exp\left[-\lambda(\eta - \psi)\left(\xi_{\mathbf{r}} - \frac{\Lambda}{2}\right)\right]\right]$$

$$\times \left\{ 1 + \exp\left[-2\lambda\xi_{\mathbf{r}}(\eta - \psi)\right] \right\} + \left. \exp\left[-\lambda(\eta - \psi)\left(\xi_{\mathbf{r}} - \frac{\Lambda}{2}\right)\right] \left\{ 1 - \exp\left[-2\lambda\xi_{\mathbf{r}}(\eta - \psi)\right] \right\} \right]$$
(A2)

$$CC = \frac{2\lambda\psi\Lambda}{k^{3}\left[1 - \exp(-2\lambda\psi\xi_{s})\right]^{2}} \left(\frac{\Lambda}{2\xi_{s}k^{2}} \left\{2 \exp(-2\lambda\psi\xi_{s}) - \exp\left[-\lambda\psi\left(\xi_{s} + \frac{\Lambda}{2k^{2}}\right)\right]\right\}$$

$$\times \left[1 + \exp(-2\lambda\psi\xi_{\rm s})\right] - \exp\left[-\lambda\psi\left(\xi_{\rm s} + \frac{\Lambda}{2k^2}\right)\right] \left[1 - \exp(-2\lambda\psi\xi_{\rm s})\right]$$
(A3)

$$DD = \frac{1 - 2 \exp\left[-\lambda(\eta - \psi)\left(\xi_{r} - \frac{\Lambda}{2}\right)\right] + \exp\left[-2\lambda\xi_{r}(\eta - \psi)\right]}{1 - \exp\left[-2\lambda\xi_{r}(\eta - \psi)\right]}$$
(A4)

$$EE = \frac{1 - 2 \exp\left[-\lambda\psi\left(\xi_{s} + \frac{\Lambda}{2k^{2}}\right)\right] + \exp(-2\lambda\psi\xi_{s})}{1 - \exp(-2\lambda\psi\xi_{s})}$$
(A5)

$$FF = \frac{2}{\left[1 - \exp\left(\frac{-\Lambda_{\infty}\psi}{k^2}\right)\right]^2} \left(\frac{\Lambda_{\infty}\psi}{k^3} \left\{1 - \exp\left[-\Lambda_{\infty}(\eta - \psi)\right]\right\} \exp\left(\frac{-\Lambda_{\infty}\psi}{k^2}\right)$$

$$-\Lambda_{\infty}(\eta - \psi) \left[1 - \exp\left(\frac{-\Lambda_{\infty}\psi}{k^2}\right) \right] \exp\left[-\Lambda_{\infty}(\eta - \psi)\right] \right)$$
(A6)

$$GG = \frac{2 \exp\left(\frac{-\Lambda_{\infty}\psi}{k^2}\right)}{\left[1 - \exp\left(\frac{-\Lambda_{\infty}\psi}{k^2}\right)\right]^2} \left(-\Lambda_{\infty}(\eta - \psi) \exp\left[-\Lambda_{\infty}(\eta - \psi)\right]\left[1 - \exp\left(\frac{-\Lambda_{\infty}\psi}{k^2}\right)\right]$$

$$+\frac{\Lambda_{\infty}\psi}{k^{3}}\left\{1 - \exp\left[-\Lambda_{\infty}(\eta - \psi)\right]\right\}\right)$$
(A7)

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$$HH = \frac{-G^{2}}{(k-1)} \left[1 + \exp\left(\frac{-\Lambda_{\infty}\psi}{k^{2}}\right) \left(k(FF) + \left(1 + \frac{2\Lambda_{\infty}\psi}{k^{2}}\right) \left\{ \frac{1 - \exp\left[-\Lambda_{\infty}(\eta - \psi)\right]}{1 - \exp\left(\frac{-\Lambda_{\infty}\psi}{k^{2}}\right)} \right\} \right) \right]$$
(A8)

REFERENCES

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- Kochi, Kikuo C.: Characteristics of a Self-Lubricated Stepped Thrust Pad of Infinite Width With Compressible Lubricant. J. Basic Eng., vol. 81, no. 2, June 1959, pp. 135-146.
- Ausman, J. S.: An Approximate Analytical Solution for Self-Acting Gas Lubrication of Stepped Sector Thrust Bearings. ASLE Trans., vol. 4, no. 2, Nov. 1961, pp. 304-313.
- Scheinberg, S. A.: Gas Lubrication of Sliding Bearings (Theory and Calculations). Friction and Wear in Machines, Inst. of Machine Sci., Acad. Sci. USSR, vol. 8, 1953, pp. 107-204.
- 4. Scarborough, James B.: Numerical Mathematical Analysis. Sixth ed., Johns Hopkins Press, 1966.
- Archibald, F. R.: A Simple Hydrodynamic Thrust Bearing. Trans. ASME, vol. 72, no. 4, May 1950, pp. 393-400.

TABLE I. - COMPARISON OF DIMENSIONLESS LOAD CAPACITY FOR INFINITE-WIDTH SOLUTION AND LIMITING CASE OF

FINITE-WIDTH SOLUTION FOR TWO-SERIES TRUNCATIONS

[Feed groove parameter, $\eta = 1.0$; step location parameter, $\psi = 0.45$; film thickness ratio, k = 2.0.]

Dimensionless	Dimensionless load	Limiting case of dimensionless load		Dimensionless	Limiting case of dimensionless stiffness	
bearing number,	capacity of infinite-	capacity of finite-width bearing,		stiffness of	of infinite-width bearing,	
Λ	width bearing,	lim W		infinite-width	lim K	
	W _∞	$\lambda \rightarrow 0$		bearing,	$\lambda \rightarrow 0$	
		M = 101 (50 terms)	M = 1001 (500 terms)	K∞	M=101. (50 terms)	M = 1001. (500 terms)
10 ⁻⁴	2.5516×10 ⁻⁶	2.5414×10 ⁻⁶	2.5505×10 ⁻⁶	4.1825×10 ⁻⁶	4.1659×10^{-6}	4.1808×10 ⁻⁶
10-3	2.5517×10 ⁻⁵	2.5416×10 ⁻²	2.5507×10 ⁻⁵	4.1831×10 ⁻⁰	4.1664×10 ⁻⁵	4.1814×10^{-5}
10 ⁻²	2.5533×10 ⁻⁴	2.5431×10 ⁻⁴	2.5522×10 ⁻⁴	4.1887 $\times 10^{-4}$	4.1720×10 ⁻⁴	4.1870×10 ⁻⁴
10 ⁻¹	2.5687×10^{-3}	2.5585×10^{-3}	2.5677×10^{-3}	4.2448×10 ⁻³	4.2279×10 $^{-3}$	4.2431×10 ⁻³
1	2.7185×10 ⁻²	2.7077×10^{-2}	2.7174×10 ⁻²	4.7869×10 ⁻²	4.7678×10^{-2}	4.7849×10 ⁻²
10	3.2354×10^{-1}	3.2226×10 ⁻¹	3. 2341×10^{-1}	5.8440 $\times 10^{-1}$	5.8208×10 ⁻¹	5.8416 $\times 10^{-1}$
100	5.7998×10^{-1}	5.7716×10^{-1}	5.7923×10 ⁻¹	5.6021 $\times 10^{-1}$	5.5891×10 ⁻¹	5.6091 $\times 10^{-1}$
500	5.5600×10 ⁻¹	5.4967 $\times 10^{-1}$	5.5165×10^{-1}	5.5200×10 ⁻¹	5.8944×10 ⁻¹	5.6141×10 ⁻¹

TABLE II. - COMPARISON OF DIMENSIONLESS LOAD CAPACITY

Inner- to	Dimensionless bearing number, Λ						
outer- radius	10	20	40	80	160		
ratio, R _i /R _o	Dimensionless load capacity of finite-width bearing, W (a)						
0.2	0.064 .063	0,141 .138	0.286 .280	0.470 .458	0.638		
0.3	0.059	0,131	0.270	0.457 .450	0.622		
0.4	0.053 .053	0.119 .118	0.248 .246	0.431 .429	0.602 .598		
0.5	0.046	0.103 ,102	0.219	0. 397 . 396	0.572		
0.6	0.038 .038	0.084 .084	0.184 .183	0. 349 . 348	0.530 .531		
0.7	0.029 .029	0.063 .063	0.141 .140	0.284 .284	0.466 .466		
0.8	0.019 .019	0.041 .041	0.091 .091	0.195 .196	0.363 .363		

OF AUSMAN (REF. 2) WITH PRESENT ANALYSIS

^aFirst value from ref. 2; second value from present analysis.

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TABLE III. - OPTIMAL STEP PARAMETERS FOR RESULTING MAXIMUM DIMENSIONLESS LOAD

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CAPACITY FOR VARIOUS DIMENSIONLESS BEARING NUMBERS

Dimensionless	Optimal value	Optimal value	Optimal value	Maximum value of	Dimensionless
bearing number,	for length-to-	of step location	of film	dimensionless load	stiffness,
Λ	width ratio	parameter,	thickness	capacity of finite-	К
	of pad,	$(\psi)_{ont}$	ratio,	width bearing,	
	(\lambda) opt	- Opt	^(k) opt	(W) _{max}	
10 ⁻⁵	0.918	0.555	1.693	1.181×10^{-7}	2.362×10 ⁻⁷
2.5×10 ⁻²	. 919	. 554	1.693	2.955 $\times 10^{-4}$	5.913×10^{-4}
5.0×10^{-2}	. 920	. 553	1.693	5.914 $\times 10^{-4}$	1.184×10^{-3}
0.1	. 922	. 552	1.693	1.184×10^{-3}	2.376 $\times 10^{-3}$
. 2	. 925	. 549	1.693	2.376×10 ⁻³	4.779×10^{-3}
. 4	. 933	. 544	1.694	4.779×10 $^{-3}$	9.669 $\times 10^{-3}$
. 8	. 948	. 533	1.696	9.670 $\times 10^{-3}$	1.979×10^{-2}
1.6	. 980	. 511	1.703	1.980×10^{-2}	4.145×10^{-2}
3.2	1.043	. 471	1.723	4.144×10^{-2}	9.006 $\times 10^{-2}$
6.4	1.145	. 412	1.790	8.903×10 ⁻²	1.974×10^{-1}
12.8	1.294	. 344	1.949	1.878×10 ⁻¹	3.854 $\times 10^{-1}$
25.6	1.575	. 271	2.240	3.651×10^{-1}	6.514×10 ⁻¹
51.2	2.037	. 204	2.698	6.492 $\times 10^{-1}$	1.003
102.4	2.710	. 151	3.359	1.072	1.456
204.8	3.642	. 110	4.270	1.674	2.034
409.6	4.901	. 080	5.501	2.502	2.770

[Resulting dimensionless stiffness also given when optimum step parameters are used.]

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TABLE IV. - OPTIMAL STEP PARAMETERS FOR RESULTING MAXIMUM DIMENSIONLESS

STIFFNESS FOR VARIOUS DIMENSIONLESS BEARING NUMBERS

[Resulting dimensionless load capacity also given when optimal step parameters are used.]

Dimensionless	Optimal value	Optimal value	Optimal value	Maximum value of	Dimensionless load
bearing number,	for length-to-	of step location	of film	dimensionless	capacity of finite-
Λ	width ratio	parameter,	thickness	stiffness,	width bearing,
3	of pad,	$(\psi)_{\text{opt}}$	ratio,	(K) _{max}	w
	$(\lambda)_{opt}$		(k) _{opt}		
10 ⁻⁵	0.915	0.557	1.470	2.550 $\times 10^{-7}$	1.115×10 ⁻⁷
2.5×10 ⁻²	. 917	. 555	1.471	6.334 $\times 10^{-4}$	2.789 $\times 10^{-4}$
5.0×10 ⁻²	. 919	. 554	1.471	1.278×10^{-3}	5.582 $\times 10^{-4}$
0.1	. 922	. 551	1.471	2.563×10 ⁻³	1.118×10^{-3}
. 2	. 929	. 546	1.472	5.152 $\times 10^{-3}$	2.245×10^{-3}
. 4	. 943	. 535	1.474	1.041×10^{-2}	4. 520×10 ⁻³
. 8	. 973	. 514	1.479	2.125×10^{-2}	9.164 $\times 10^{-3}$
1.6	1.035	. 474	1.494	4.432×10^{-2}	1.882×10^{-2}
3.2	1.153	. 408	1.537	9.561×10 ⁻²	3.953×10 ⁻²
6.4	1.353	. 328	1.642	2.072×10^{-1}	8.518 $\times 10^{-2}$
12.8	1.863	. 232	1.849	4.070×10^{-1}	1.757×10^{-1}
25.6	2.952	.148	2.191	7.036×10 ⁻¹	3.289×10^{-1}
51.2	5.035	.088	2.687	1.105	5.624 $\times 10^{-1}$
102.4	9.093	. 051	3.368	1.627	8.931×10 ⁻¹
204.8	17.172	.030	4.274	2.293	1.338

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Figure 1, - Comparison of present linearized results with Kochi's exact results for infinite-width step slider bearing when step location parameter is 0, 5.



Figure 2. - Effect of dimensionless bearing number on optimum parameters for maximum dimensionless load.



Figure 3. - Effect of dimensionless bearing number on optimal step parameters for maximum dimensionless stiffness.



(b) Maximum stiffness.

Figure 4. - Effect of dimensionless bearing number on dimensionless load capacity and dimensionless stiffness.

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