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ULTRAVIOLET PHOTOMETRY OF THE ECLIPSING VARIABLE CW CEPHEI

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16. Abstract

An extended series of photometric observations was made of the eclipsing variable CW Cephei using the Wisconsin instrument on OAO 2. Approximate elements derived solely from the eclipse depths and shape of the secondary are in satisfactory agreement with those found using ground-based observations. However, persistent asymmetries and anomalous light variations, all larger than the expected experimental error, were also found; subsequent ground-based observations show H_α entirely in emission indicating the presence of an extended gaseous system surrounding one or both of the components. Consistent solutions utilizing all data at all wavelengths were not found.

In addition to the lightcurve analysis, a detailed comparison was made of the flux distribution of the binary relative to that for the nominally unreddened stars δ Pic, B1III, and η Aur, B3V, to investigate the effects of interstellar extinction. The observations for these comparison stars were obtained near to the time of the binary observations, thereby minimizing spurious results due to instrumental changes. The resultant extinction curves, normalized at a wavelength of 3330Å, show a relatively smooth increase with decreasing wavelength; no conspicuous hump near a wavelength of 2200Å is evident. If, in addition, it is assumed that $E(B - V) = 0.68$, a smaller amplitude is found for the curves than that given by Stecher for the Perseus region.

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INTRODUCTION

In mid-April 1971, the OAO 2 spacecraft began to be stable for extended observing periods while pointed toward the region around Cepheus. As a result, it became feasible to obtain complete ultraviolet lightcurves of the eclipsing binary CW Cep. Petrie's (1947) observations of CW Cep show it to be a double-line spectroscopic binary consisting of two early B stars of nearly equal mass but with a differential spectroscopic luminosity $\Delta m = 0.30 \pm 0.03$. It is designated as no. 69 (bright) by Blaauw, Hiltner, and Johnson (1959) in their photometric study of the Cep OB3 association. They list for it $B - V = 0.41$ and $U - B = -0.52$; the application of the Q method results in unreddened colors of $(B - V)_0 = -0.27$, $(U - B)_0 = -1.04$, and hence a color excess $E(B - V) = 0.68$. Its revised MK classification is B1.5Vn according to Garrison (1970), but the intermediate band and β index photometry of Crawford and Barnes (1970) and McNamara (1966) leads to a somewhat earlier classification. These more recent classifications and that based on the earlier photometry are in substantial agreement. Membership in the association appears certain, based on the binary's location in the color-magnitude plane after corrections are made for interstellar absorption using $R = 3.0$. By implication, one can expect that the binary is young; i.e., of the order of the expansion age of the association, which is given by de Vegt (1966) to be approximately 1.6×10^6 yr. Abrami and Cester (1960) obtained approximate orbital solutions from two-color photoelectric lightcurves. They found that primary eclipse is a partial transit and they confirmed that the components, though quite similar in surface brightness, are distinctly different in size. A complete rediscussion of this system based on new multicolor photoelectric observations is being prepared by Il-Seong Nha (1971). Of some interest is his discovery of rapid apsidal motion.

DATA REDUCTION

The present observations were made with the four photometers in the University of Wisconsin instrumentation on OAO 2. A total of 49 data points, each the mean of six separate measurements of the binary, were obtained over 29 orbits. Dark and calibration readings were obtained at least once per orbit for each photometer. In addition, six orbits were dedicated to sky measurements distributed over the 9 days of observing. The majority of the orbits between 12280 and 12411 were used for this program. The 10-arcmin field aperture was decentered from the nominal position of the binary by 2.5 arcmin to avoid photometric contamination by a nearby fainter B star; sky readings were taken 30 arcmin north of this nominal pointing. No attempt was made to obtain data with the shortest wavelength filter, and all data from the longest wavelength filter were rejected because of contamination by the unavoidable inclusion of a faint late-type star in the aperture. The other short-wavelength data appear to be free of this contamination because the eclipse depths appear to be normal. All data were reduced in the manner suggested by Code.[1] The flux F measured through filter i with photometer j relative to that at an effective wavelength λ_{eff} of 3330Å is given by the equation

$$\log \frac{F_i}{F_{3330}} = \frac{(\mathrm{Star}_i - \mathrm{Dark}_j) - (\mathrm{Sky}_i - \mathrm{Dark}_j)}{\mathrm{Calibration}_j - \mathrm{Dark}_j} - \log \frac{(\mathrm{Star}_1 - \mathrm{Dark}_1) - (\mathrm{Sky}_1 - \mathrm{Dark}_1)}{\mathrm{Calibration}_1 - \mathrm{Dark}_1} + \log \Delta_i \, ,$$

where Δ_i is a correction factor. The values of Δ_i used are those suggested by Code but amended by Holm[2] to allow for the declining sensitivity of photometer 4. Linear interpolation between dark measurements gave the value appropriate for the time of the star or sky measurement. The means for all calibration measurements obtained during the experiment were used to normalize the data from each photometer. Heliocentric orbital phases corresponding to the mean GMT of each observation were computed using Nha's (1971) ephemeris,

$$\text{Julian date} = 2435373.4487 + 2^{\mathrm{d}}.72919396E \pm 0^{\mathrm{d}}.0256 \sin (0.07018E - 31^{\circ}.55),$$

where $E = 2083$ for the observed primary minimum. Approximate times of minima were determined graphically from the data, and they are in excellent agreement with this ephemeris. The observed epochs for primary and secondary eclipses are at Julian dates 2441058.270 and 2441054.128, respectively.

INTERSTELLAR EXTINCTION

To investigate the ultraviolet extinction of CW Cep, relative fluxes were found first by forming the means of data taken on April 10 and 15, 1971, when the system was at maximum light. The approximate effective wavelengths (expressed in angstroms) corresponding to each flux measurement, the data for CW Cep expressed in magnitudes, and similar data for two nominally unreddened stars are presented in the first four columns of Table 1.

The original data for the two comparison stars, supplied by the University of Wisconsin, were reduced in the manner described above. The stars η Aur, B3V, and δ Pic, B1.5, were selected to match, respectively, the original and the more recent spectral classifications of CW Cep. Just as important, these observations were obtained at approximately contemporaneous epochs, thus minimizing any effects of instrumental variation. The extinction law is defined here as the difference in magnitude at each effective wavelength between the binary and comparison star, normalized to unity color excess. It differs somewhat from the usual representation in that Δm is set equal to zero, arbitrarily, at $\lambda = 3330$Å. The differential magnitudes thus formed are listed separately in columns 5 and 6 of Table 1 while the "average" extinction curve found by Bless and Savage[3] from OAO 2 observations and that found for the Perseus region by Stecher (1969) are given in the last two columns. The data and

[1] A. D. Code, 1971, private communication.
[2] A. V. Holm, 1971, private communication.
[3] See Table 1, Footnote a, p. 3.

Table 1—Observed flux and extinction data.

λ_{eff} (Å)	Relative Flux (mag)			Extinction			
	CW Cep	δ Pic	η Aur	CW Cep − δ Pic	CW Cep − η Aur	Bless and Savage[a]	Stecher (1969)
3330	0	0	0	0	0	0	0
2980	−.03	−.46	−.36	.63	.49	.65	.8
2940	.06	−.12	−.34	.26	.59	.85	1.0
2460	.43	−1.16	−.80	2.34	1.81	2.34	3.0
2380	.59	−.96	−1.04	2.27	2.39	2.90	3.4
2040	.40	−1.26	−1.47	2.44	2.75	4.00	5.5
1920	.14	−1.93	−1.55	3.05	2.48	3.32	5.2
1680	−.02	−2.14	−1.68	3.11	2.44	2.48	3.8
1500	−.83	−2.60	−2.19	2.59	2.00	2.58	4.3
1380	−1.29	−2.75	−2.19	2.14	1.32	3.07	5.2

[a]R. C. Bless and B. D. Savage, 1972, "Ultraviolet Photometry From the Orbiting Astronomical Observatory. II. Interstellar Extinction." *Proc. Symp. OAO* (Amherst, Mass.), Aug. 22-23, 1971, to be published.

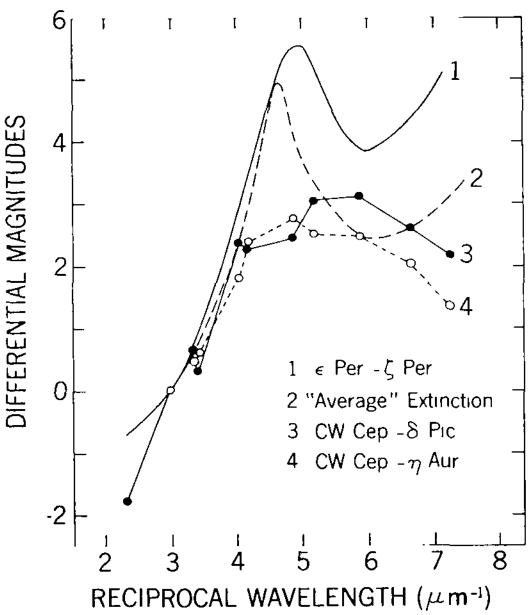

Figure 1—Ultraviolet extinction curves for CW Cep formed with δ Pic, curve 3, and η Aur, curve 4, as comparison stars. Curves labeled 2 and 1 refer respectively to the average extinction given by Bless and Savage and to the extinction in Perseus according to Stecher (1969).

the two reference extinction curves are shown in Figure 1. It is immediately evident by comparing these curves that the shape of the extinction curve for CW Cep differs markedly from that appropriate to the Perseus region. It is also seen that the average extinction fits the observed data only poorly. Most conspicuous are the apparent absence of the 4.6 μm^{-1} peak and the turning down of the curve at the shortest wavelengths. The overall shape resembles that found by Bless and Savage for θ^1 and θ^2 Ori except that the amplitude, i.e., relative absorption, is greater for CW Cep in the midultraviolet range. No definite explanation can be offered for the unusual shape of this extinction curve. The fact that the photometric colors are consistent with the spectral type, and that the resultant position of the binary in the color-magnitude diagram of the association appears reasonable after the visible wavelength data are corrected for reddening and absorption with the usual relations, suggests that the absorbing medium is normal, at least in regard to its properties at visible wavelengths. In support of this view, spectra of CW Cep, obtained with the GSFC Cassegrain spectrograph and 36-in. telescope, show moderately strong absorption at a wavelength of 4430Å; the estimated equivalent width of 2Å to 3Å appears consistent with the $E(B - V) = 0.68$. On the other hand these same spectra show H_α entirely in emission, thus confirming the report of Wackerling (1970). The emission strength appears to vary from approximately 3Å to 5Å as a function of the orbital phase. With this evidence for the presence of an extended gaseous system about one or both components, one can speculate that the intrinsic flux distribution is abnormal in the sense that overlying Balmer emission could be modifying the observed ultraviolet flux distribution in a way similar to that found for Be stars. However, the net effect of this would be to exaggerate the apparent absorption at the shortest wavelengths where the Balmer emission would be the weakest which, of course, is contrary to what is observed. Furthermore, this emission as opposed to strategically located line emission cannot explain the absence of the 4.6 μm^{-1} peak. It should be noted that spectral scans of $\theta^1 + \theta^2$ Ori had revealed no emission at 4.6 μm^{-1}. It is likely that for CW Cep both conditions exist, i.e., the ultraviolet flux distribution may well be abnormal, and the interstellar medium or perhaps the immediate stellar environs may be peculiar. Like $\theta^1 + \theta^2$ Ori, CW Cep is a multiple system and, additionally, is a member of a young association. It would be of interest to determine if the infrared color excess is unusual as is reported for $\theta^1 + \theta^2$ Ori. Of course, data for other members of the Cep OB3 association would be valuable in elucidating this puzzle.

LIGHTCURVE ANALYSIS

The original objective was to verify the theoretical limb-darkening coefficients by solving the lightcurves. This could not be done because of instrumental problems that resulted in incomplete coverage and poor phase resolution. Hence, only approximate solutions were possible. Preprimary maximum is adequately observed as is the entire secondary eclipse, but the egress of primary eclipse is represented by few observations of low weight, and the postprimary maximum is essentially unobserved. Comparison stars were not used as controls, but no systematic differences in light levels were detected for the phases in the preprimary maximum for which redundant data were available and also for the shallow partial phases of primary ingress.

The lightcurves, consisting of the highest weight data corresponding to the largest signal-to-noise cases, are shown in Figure 2, where for simplicity the individual data points have been connected by straight lines. On the basis of visual inspection of the data, rectification in the main appeared unwarranted although some curvature in the maximum of the $\lambda = 1920$Å lightcurve seems to be present. One can show that, based on the solution given later in this paper, the effects of reflection and photometric ellipticity should be observed, and one must conclude that the scatter in the present data masks these effects. The proximity effects, if present, are small, but peculiar light variations are evident at the shoulders of both eclipses. Primary eclipse is grossly asymmetric but it must be recalled that egress is only poorly represented. Secondary eclipse changes shape with wavelength. To quantify the shapes of both eclipses, the shape parameter $\chi(n = 0.8)$, as defined by Russell and Merrill (1952), was determined from smooth curves passed through the data for the ingress of primary and the combined data from both branches of secondary. The shape parameters, the observed eclipse depths, and their mean values, which are used in a nomographic solution, are listed in Table 2. All observations are tabulated in the appendix. Because of the paucity of data and the peculiar light variation at the

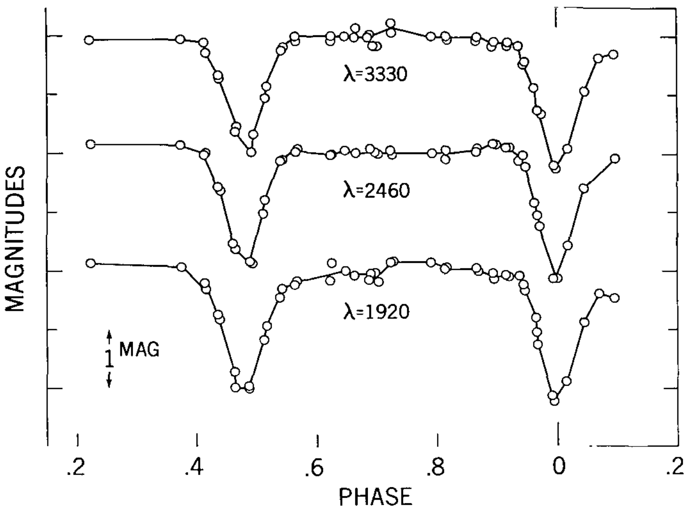

Figure 2—Representative lightcurves of CW Cep.

Table 2—Computed parameters.

Parameter	Wavelength							
	3330Å			2460Å		1920Å		
	p I	p II	s	p	s	p	s I	s II
Depth (mag)	0.473	0.455	0.380	0.454	0.377	0.451	0.395	0.360
Shape χ (n = 0.8)	.274	.273	.315	.251	.353	.249	.398	.440
Correction (mag)	0	.025	0	0	0	0	0	.04

p = primary; s = secondary; I = branch I; II = branch II.

shoulders, an uncertainty in the unit light level existed at several wavelengths for data taken during the eclipses. The values of the corrections listed in Table 2, which were applied to all eclipse data, are those that force the lightcurve to attain unit light near to external tangency. The correction in magnitudes was subtracted from each data point before the corresponding shape parameter was calculated.

The variation in the values of χ for secondary eclipse is quite large; the sense of the variation is that the branches steepen with decreasing wavelength. Even if one excludes the case where the corrections were applied, the variation is much larger than can be explained by changing the limb darkening of the eclipsed secondary components. Although the run of the eclipse is only poorly defined, the data at minimum for the λ = 1920Å lightcurve suggest that a total phase may exist. On the other hand, both the eclipse depths and the shape of primary remain constant. The depths of the eclipses are similar to those found by Abrami and Cester (1960), $\Delta m_p = 0.44$ and $\Delta m_s = 0.37$. This implies that, to a first approximation, the surface brightnesses of the components are alike.

In attempting a nomographic solution, it was found that the depths and line-shape relation for primary eclipse yielded no solution at any wavelength, whereas the secondary eclipse shapes all yielded solutions specifying primary eclipse to be a geometrically deep partial transit. The nomographic solution based on the mean depths and shape of secondary eclipse yielded the following elements: ratio of radii = 0.69, radius of primary = 0.26, orbital inclination = 83°, and luminosity of the primary = 0.7. The usual normalizations were adopted, the ratio of surface brightness was set equal to 0.9 (the value used by Abrami and Cester), and a limb darkening of 0.6 was assumed for both components. These elements agree in the main with those found by Abrami and Cester and, in view of the low precision and scarcity of the data, further refinement is considered unwarranted. The computed solutions of Abrami and Cester (1960) and the values given in this paper during the minima are shown in Figure 3 for λ = 2460Å.

DISCUSSION

The system of CW Cep is unusual in many respects. To recapitulate, the ultraviolet extinction is abnormal, emission in H_α is observed, peculiar light variations appear at the shoulders of the minima, and the shape of secondary eclipse varies with wavelength in the ultraviolet. The photometric solution implies an approximate difference in luminosity of Δm = 0.9, which is inconsistent with that determined spectroscopically. Both components should be well situated on the main sequence if the expansion age of the association applies, because the contraction to the main sequence is given by Iben (1965) for comparable mass stars as 1.5×10^5 yr, whereas the main-sequence lifetime is of the order of 10^7 yr. And yet from the ratio of the radii or the differential luminosity, either the primary has evolved away from the main sequence or the secondary is a subdwarf. On the basis of the empirical mass-radius relation given by Harris, Strand, and Worley (1963), Petrie's spectroscopic solution combined with the photometric results would indicate that the primary is evolved. These questions and particularly the problem of the variation of the eclipse shapes cannot be resolved solely with the present photometric data. If an extended atmosphere exists about either of the components, then one may question the applicability of the Russell-Merrill model and the elements derived thereby. It is easy to show that for a system undergoing eclipses where the geometric depth $p_0 = -1.0$, i.e., the

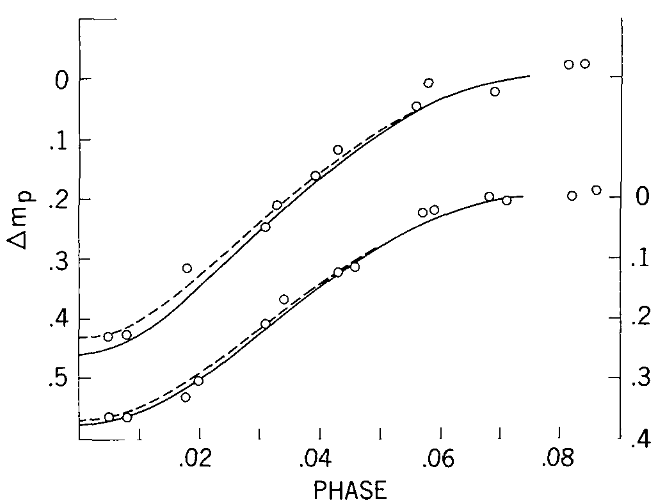

Figure 3—Computed solutions of Abrami and Cester (1960), represented by the dashed curve, and the present work, represented by the solid curve, are shown for λ = 2460Å.

case of internally tangent eclipses, a large change can be induced in the shape parameter by moderate changes in the radii. For example, a 10-percent change in the radius of the secondary star in CW Cep could account for the observed χ variation if the radius decreases with decreasing wavelength. This situation could obtain if the smaller secondary were surrounded by an optically thin shell that emits a Balmer continuum. An optically thick shell surrounding the larger (and presumably more evolved) star would produce the inverse effect. The phase of external tangency should vary, but the photometric difficulties at the shoulders will mask this effect. Unfortunately this picture of the binary aggravates the already unusual ultraviolet extinction by requiring the true extinction freed of the anomalous shell flux to turn down even more steeply at the shortest wavelengths. Additional observations, particularly scanner measures of the flux distribution, may resolve some of these problems.

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Appendix
Observed Data for CW Cep

Phase	m_λ			
	3330Å	2980Å	2460Å	1920Å
0.018	0.384	0.349	0.316	0.372
0.043	.191	.152	.120	.167
0.069	.076	.070	–	.064
0.094	.060	–	.019	.079
0.223	.016	.012	-.033	-.028
0.375	.010	.018	-.031	-.011
0.412	.020	-.012	.007	.040
0.415	.057	–	.000	.065
0.437	.133	.080	.114	.147
0.440	.142	.122	.123	.168
0.463	.328	.285	.306	.346
0.466	.305	.331	.333	.401
0.488	.399	.350	.364	.389
0.491	.337	.339	.366	.400
0.514	.210	.228	.208	.237
0.517	.164	.173	.158	.190
0.539	.047	.023	.026	.088
0.542	.037	.036	.022	.060
0.565	-.002	.019	-.002	.047
0.568	.021	.031	-.014	.039
0.621	-.002	-.010	.005	.027
0.624	.014	.011	.006	-.027
0.648	-.000	-.006	-.011	.000
0.662	.004	-.034	.000	.018
0.663	-.033	–	.004	–
0.687	.003	-.014	–	.029
0.689	-.010	.014	-.016	.007
0.698	.027	-.002	.002	.005
0.701	.030	.011	.002	.039
0.723	-.050	.009	-.009	-.031
0.726	-.016	-.013	.000	-.030
0.790	-.000	.013	.003	-.033
0.814	-.000	.015	.019	-.007
0.817	.005	.006	-.010	-.014
0.865	.015	.013	-.001	-.006
0.867	.003	.021	-.015	.001
0.891	.035	-.015	-.028	.006
0.893	.016	-.020	-.030	.034
0.916	.035	.000	-.021	.018
0.919	.020	.013	-.021	.017
0.936	.031	.034	.027	.013
0.942	.097	.088	.011	.042
0.944	.082	.065	.048	.071
0.961	.177	.184	.165	.158
0.967	.255	.220	.214	.212
0.969	.269	.241	.250	.247
0.992	.441	.421	.427	.416
0.995	.453	.372	.430	.437

NOTE–Heliocentric phases were computed with the ephemeris given in the text for primary minimum.

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