ELECTRON PARAMAGNETIC RESONANCE STUDY
OF ALINEMENT INDUCED BY MAGNETIC FIELDS
IN TWO SMECTIC-A LIQUID CRYSTALS
NOT EXHIBITING NEMATIC PHASES

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SUMMARY

Using vanadyl acetylacetonate (VAAC) as a paramagnetic probe, we have studied the molecular ordering in two smectic-A liquid crystals that do not display nematic phases. We have attained reproducible alignment by slow cooling through the isotropic smectic-A transition in dc magnetic fields of 1.1 and 2.15 teslas. The degree of order attained is small for a smectic-A liquid crystal, being $\sigma_Z = -0.1$. Measurements have been made of the variation of the average hyperfine splitting of the aligned samples as a function of orientation relative to the dc magnetic field of the spectrometer. This functional dependence is in agreement with the theoretical prediction except where the viscosity of the liquid crystal becomes large enough to slow the tumbling of the VAAC, as indicated by asymmetry in the end lines of the spectrum.

INTRODUCTION

The technique of electron paramagnetic resonance (EPR) has previously been used to determine the ordering properties of both nematic liquid crystals (refs. 1 and 2) and smectic liquid crystals which have a nematic phase (refs. 3 and 4). The technique consists of doping the liquid crystal with a small amount ($10^{-3}$ m) of the nearly planar paramagnetic molecule, vanadyl acetylacetonate (VAAC). This molecule exhibits a broad, eight-line spectrum with axial $g$ and hyperfine tensors. From the measurement of the average hyperfine splitting of the spectrum, it is possible to determine the orientation of the probe molecules. From this orientation the degree of order of the molecules of the liquid crystal is inferred.

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The dc magnetic field of the EPR spectrometer is large enough to aline the molecules in a nematic liquid crystal; that is, the field makes the preferred direction of the liquid-crystal molecules uniform throughout the sample. The degree of order can be determined as a function of temperature throughout the nematic phase. A magnetic field will not aline the molecules in the smectic phase of a liquid crystal. However, if the smectic liquid crystal displays a nematic phase, the molecules can be alined in this phase and then slowly cooled into the smectic phase while in the magnetic field (ref. 3). Since the magnetic field will not realine the molecules in the smectic phase, it is possible to study the EPR spectra as a function of sample orientation relative to the field (refs. 3 and 4).

In this study a dc magnetic field was used to aline two smectic-A liquid crystals that do not display a nematic phase. Foex (ref. 5) has stated that an alined smectic phase could be obtained by cooling the isotropic melt in a magnetic field, though Luckhurst and Sanson (ref. 6), more recently, were unable to duplicate this claim with fields up to 0.65 tesla. We have been able to obtain alinement in fields of 1.1 and 2.1 teslas. We have also measured the effect on the hyperfine splitting of sample orientation relative to the magnetic field.

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**EXPERIMENT**

The two liquid crystals used were 4-hexyloxybenzylidene-4'-aminobenzoate ethyl ester (HBABE) and 4-octyloxybenzylidene-4'-aminobenzoate ethyl ester (OBABE). The HBABE displayed a smectic-A phase from 69° to 96° C, while the OBABE displayed a smectic-A phase from 75° to 98° C. The materials were synthesized and purified for us by Dr. D. L. Fishel of Kent State University.

The details of the sample preparation were described previously (ref. 1). Briefly, a small amount of VAAC (10⁻³ m) was added to each liquid-crystal sample, and dissolved oxygen was removed by evacuation during melting. EPR spectra were obtained by using a modified Varian E-12 X-band spectrometer. Samples were heated by using a modified Varian 4547 variable temperature accessory, and temperatures were monitored with a 36-gauge copper-constantan thermocouple inserted in the nitrogen stream that heated the sample.

The smectic-A liquid crystals were ordered by slow cooling through the isotropic-
to-smectic-A transition while in a magnetic field. Fields of 0.33, 1.1, and 2.1 teslas were used. Cooling rates of 1/8°C per minute were employed. The EPR spectra of the aligned smectic sample were recorded for every 10° of rotation relative to the magnetic field of the spectrometer (0.33 T).

**THEORY AND RESULTS**

The theory required for interpreting the results of this study has been described previously (ref. 3). The variation of the average hyperfine splitting of the spectra \( \langle a \rangle \) with the angle of rotation relative to the dc magnetic field \( \alpha \) is given by the expression

\[
\langle a \rangle = a + \frac{b\sigma_Z}{3} (3 \cos^2 \alpha - 1)
\]

where \( a \) is the isotropic hyperfine splitting (-0.107 T; ref. 7) and \( b \) is the anisotropic splitting (-0.01215 T). The ordering parameter \( \sigma_Z \) for \( \alpha = 0 \) is obtained from the expression

\[
\sigma_Z = \frac{\langle a \rangle_0 - a}{2(a - A_\perp)}
\]

where \( A_\perp \) is the perpendicular component of the axial hyperfine tensor (ref. 7) (-0.00665 T) and \( \langle a \rangle_0 \) is the experimentally determined average hyperfine splitting after cooling through the isotropic-to-smectic-A transition. It should be noted that the pseudosecular terms have been neglected in deriving equation (1). This results in a maximum difference from the exact equation (ref. 8) of less than 1 percent at \( \alpha = 45^0 \). At \( \alpha = 0^0 \) and \( 90^0 \) the two equations are equivalent.

Preliminary experiments with these two liquid crystals indicated that cooling through the isotropic-to-smectic-A (\( i - S_A \)) transition in a 0.33-tesla field resulted either in no ordering or, sometimes, in a very small degree of ordering (\( \sigma_Z < -0.05 \)). The degree of ordering was not reproducible even though slow cooling rates were used. However, with 1.1- and 2.1-tesla fields, reproducible ordering was attained with slow cooling. The degree of order attained was the same for both the 1.1- and 2.1-tesla fields. With rapid cooling, no order was effected. Frequently, the spectra obtained after cooling in the 2.1-tesla field displayed small extra lines. This effect became progressively greater with sample use, and when the effect became pronounced, the sample was discarded.
Slow cooling of HBABE in a 1.1-tesla (or a 2.1-T) field resulted in a small degree of order: \( \sigma_Z = -0.087 \). The EPR spectra indicated that the viscosity of the liquid crystal was exerting only a moderate influence; that is, the lines were broadened but they were not asymmetric. In figure 1 we show the results of the rotation of the aligned sample in the 0.33-tesla dc field of the spectrometer. We have plotted the average splitting \( \langle a \rangle \) as a function of the angle of rotation relative to the magnetic field \( \alpha \). At \( 0^\circ \) and \( 180^\circ \) the average splitting is characteristic of the order for the given temperature \( (T = 91.5^\circ C) \), namely, \( \sigma_Z = -0.087 \). The curve is representative of the theoretical function obtained from equation (1) by substituting the experimental value of \( \sigma_Z \) and the reported hyperfine parameters (ref. 7). The agreement between the experimental points and the theoretical curve is within experimental error.

Slow cooling of OBABE in a 1.1-tesla (or a 2.1-T) field gave results very similar to those obtained with HBABE. The degree of order was a little higher, \( \sigma_Z = -0.115 \). Spectra were taken at three different temperatures within the smectic-A phase, namely, \( 96^\circ, 84^\circ, \) and \( 75^\circ C \). Spectra obtained at \( 96^\circ C \) show little viscous effects other than some line broadening and are similar to those obtained with HBABE. However, spectra obtained at \( 84^\circ C \) show slight asymmetry in the shape of the two end lines. This asymmetry indicates that the viscosity is significantly affecting the tumbling time of the VAAC molecules (ref. 2). This effect is even more pronounced in the spectra obtained at \( 75^\circ C \). Rotation experiments were performed at these same three temperatures. The results are presented in figure 2. The degree of order was the same at all tem-
temperatures, $\sigma_Z = -0.115$. Again, the curve is representative of the theoretical function obtained by substituting this value of $\sigma_Z$ along with the reported hyperfine parameters into equation (1). The data taken at $96^\circ$ C agree with the theoretical curve within experimental error. However, the data taken at $84^\circ$ and $75^\circ$ C fall below the theoretical curve at angles between $45^\circ$ and $135^\circ$.

Some visual observations of the liquid crystals were made during their cooling through the $i - S_A$ transition. The experiments were conducted in an auxiliary magnet capable of producing fields up to 2.15 teslas. For fields greater than 1 tesla and with slow cooling, the first visual indication that the $i - S_A$ transition had been reached was the appearance throughout the clear isotropic liquid of horizontal needles growing in from the inside surface of the quartz sample tube. This was followed by a gradual clearing until the entire sample was again quite clear. However, careful examination in oblique light revealed the presence of very fine horizontal rings visible at the tubing wall. Warming this sample up through the transition indicated that the process was reversible, though the needle-like phase was more pronounced and more opaque than during the cooling. For fields around 0.33 tesla or with fast cooling through the transition, the sample did not appear to be uniform throughout. There were several planes visible across the quartz tube, and they were not always horizontal. In addition, the samples did not become as clear, and different regions displayed different variations of clearing.
CONCLUDING REMARKS

Alignment of two smectic-A liquid crystals not displaying nematic phases has been attained in the experiments just described. However, the degree of order attained is quite small for a smectic-A liquid crystal. Apparently, the direction of the director is not uniform throughout the entire sample. That this is the case is especially evident from our visual observations for the fast-cooling experiments and for the experiments in the 0.33-tesla field. The appearance of small extra lines in the spectra obtained in some of the slow-cooling, high-field experiments may also arise from the nonuniformity of the director throughout the entire sample.

In our orientation studies the agreement between the experimental values and the theoretical curve is good except in the case of OBABE at the lower temperatures (fig. 2). The asymmetry in the end lines of the spectra obtained at these temperatures indicates that the tumbling of the VAAC molecule is being slowed by the higher viscosity of the liquid crystal. The effect becomes greater as the temperature is decreased. We feel the discrepancy between experiment and theory is associated with this viscous effect. This conclusion is in agreement with that postulated in one of our recent studies (ref. 3) involving two other smectic-A liquid crystals: 4-butyloxybenzylidene-4′-acetoaniline (BBAA) and 4-oxyloxy benzylidene-4′ ethylaniline. For these liquid crystals the viscous effects were much more pronounced and the discrepancy between experiment and theory was larger. Recently, Luckhurst and Setaka (ref. 9) repeated our experiments with BBAA, but used a nitroxide probe. The nitroxide probe tumbles rapidly ($\tau_c = 10^{-10}$ sec), and the effects associated with the viscosity are absent. They found very close agreement between their experimental points and their fitted equation. This result corroborates our contention that the discrepancy between experimental points and the theoretical curve in these orientation experiments arises from a viscous effect of the liquid crystal on the probe molecule, VAAC.

Lewis Research Center,
National Aeronautics and Space Administration,
Cleveland, Ohio, September 29, 1972,
502-01.

REFERENCES


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— National Aeronautics and Space Act of 1958

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