

LOW SPEED AERODYNAMIC CHARACTERISTICS OF A VECTORED THRUST V/STOL TRANSPORT WITH TWO LIFT/CRUISE FANS

By Dirk J. Renselaer

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A wind tumnel test was conducted to obtain power-on low-speed characteristics of a twin fan vectored thrust V/STOL transport aircraft. longitudinal as well as some lateral-directional data were analyzed. Hover, STOL, and conventional flight modes were investigated.
The scope of the investigation included the determination of STOL characteristics in and out of ground effect, hover characteristics with and without bank angle, roll control effectiveness as affected by ground proximity and aircraft attitude. The study also included various means to improve the lifting capability of the aircraft such as by application of fuselage strakes, exhaust vanes capable of shifting the thrust vector aft, and external flap blowing for STOL performance.
A mmber of theoretical equations for longitudinal and lateral-directional characteristics were derived and results conpared with test data.

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# LOW SPEED AERODYNAMIC CHARACTERISTICS OF A VECTORED THRUST V/STOL TRANSPORT WITH TWO LIFT/CRUISE FANS 

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## SUMMARY

A wind thinel test was conducted to obtain lnw speed characteristics in and out of ground effect of a twin fan vectored thrust V/STOL transport. Hover, STOL and comventional flight modes were investigated using a $10 \%$ scale model with thrust nozzles vectored to various angles. Primarily longitudinal data were obtained, as well as some lateral-directional data.

In hover, only minor interference lift losses were encountered in ground effect with wings level. Larger losses were encountered in a banked condition both with and without roll control application.

Lift characteristics at forward speed were found to be essentially those of the power-off characteristics plus the direct thrust contribution with only minor negative interference. The negative interference was found to be mach less than that of a previously tested VTOL model with twin exhausts (ref. 4).

Various means to improve the lift were investigated. In hover, simple exhaust nozzle wall extensions, strakes, and a strake box proved beneficial. At forward speed strakes were also beneficial, and rotating the nozzles fully aft for flap blowing may be advantageous below $\mathrm{C}_{\mathrm{T}}=2$. However, a hybrid configuration with half the nozzles deflected downward and half deflected aft was not promising to improve lift.

## Introduction

Considerable interest exists in the design of a small transport aircraft with V/STOL capability for application in $t=1980$ 's. A viable propulsion concept for such an aircraft consists of using two lift fans the exhaust of which can be vectored aft for cruise, and downward for VTOL and STOL operation. Such a concept is proposed in reference 1.

The design of such an aircraft requires detailed knowledge of propulsion induced aerodynamic interference effecis in hover and in the low speed regime in and out of ground effect. The nature of the interference is such that heavy reliance on empirical data is required.

The present report includes test data for hover and transition of a ten percent scale wind tumnel model. Primarily longitudinal data were acquired, but some lateral and directional data were also obtained with emphasis on obtaining power effects. The tests were conducted in the V/STOL tumnel of the Langley Research Center of the National Aeronautics and Space Administration.

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SYMBOLS

| A | Total nozzle exit area |
| :---: | :---: |
| $\mathrm{A}_{1}$ | Nozzle exit area of one fan (i.e., exit area of two swivel nozzles) |
| AC | Aerodynamic Center |
| $\mathrm{a}_{\mathrm{N}}$ | Distance from center of fan inlet to the c.g. (of each fan), see figure 75a |
| AR | Aspect ratio |
| ${ }^{\text {a }}$ R | Moment arm of ram drag with respect to c.g., in pitch, $\alpha=0$ (fig. 49a) |
| ${ }^{3} \mathrm{~T}$ | Mament arm of thrust vector with respect to c.g. in pitch (fig. 49a) |
| b | Wing span |
| c | MAC of wing |
| $\bar{c}_{H}$ | MAC of horizontal tail |
| $\check{c i}_{V}$ | MAC of vertical tail |
| ${ }^{\text {c }}$ | Drag coefficient, D/qS |
| $\mathrm{C}_{\text {Daero }}$ | Coefficient of aerodynamic forces in drag |
| c.g. | Center of gravity (assumed to be located at 25\% ¢) |
| ${ }^{\text {CDIOT }}$ | Coefficient of total drag, identical to $C_{D}$ |
| $\mathrm{C}_{\mathrm{L}}$ | Lift coefficient, L/qS |
| $\mathrm{C}_{\text {Laero }}$ | Lift coefficient of circulation lift |
| $\mathrm{C}_{\text {LTOT }}$ | Coefficient of total lift, identical to $\mathrm{C}_{\mathrm{L}}$ |
| $\mathrm{C}_{\ell}$ | Rolling moment coefficient, $\mathrm{MR}_{\mathrm{R}} / \mathrm{qSb}$ |
| $\mathrm{C}_{\mathrm{m}}$ | Pitching mument coefficient (bout $25 \%$ c) , M/qSb |
| $\mathrm{Cm}_{\text {maero }}$ | Coefficient of aerodynamic portion of pitching moment |
| $\mathrm{C}_{\mathrm{mR}}$ | Pitching moment due to ram drag |
| Cmiot | Total pitching moment coefficient, identical to $\mathrm{C}_{\mathrm{m}}$ |
| $\mathrm{C}_{n}$ | Yawing moment coefficient. $\mathrm{N} / \mathrm{qSb}$ |
| $\mathrm{C}_{T}$ | Thrust coefficient, T/qS |
| $\mathrm{C}_{\mathrm{Y}}$ | Side force coefficient, Y/qS |
| D | Drag force |
| $\Delta D_{\text {int }}$ | Interference drag, i.e. aerodynamic drag with power, minus aerodynamic drag without power |
| $\mathrm{d}_{\text {e }}$ | Diameter of equivaler! circular exhaust area of one fan |
| $\mathrm{D}_{\mathrm{R}}$ | Inlet momentum drag |
| EBF | Externall; blown flaps |
| f | Ratio of inlet mass flow/exhaust mass flow |
| FRP | Fuseiage reference plane |
| $\mathrm{F}_{\mathrm{X}}$ | Force in x -direction in body axis system |
| $\mathrm{F}_{\mathbf{z}}$ | Force in 2 -direction in body axis system $\mathrm{fl}^{\text {dec }}{ }^{2}$ |
| g | Gravitational constant, $9.81 \mathrm{~m} / \mathrm{sec}^{2}\left(32.2 \mathrm{ft} / \mathrm{sec}^{2}\right)$ |


| h | Distance between ground and center of exhaust plane of nozzles when $\delta_{N}=90^{\circ}$. At other nozzle angles the nozzle height is defined as if th. nozzles were at $90^{\circ}$. In banked conditions $h$ is the average between left and right hand fan nozzles. |
| :---: | :---: |
| $\mathrm{i}_{\mathrm{H}}$ | Horizontal tail incidence with respect to FRP; positive trailing edge down |
| IGE | In ground effect |
| L | Total lift (stability axis system) |
| LE | Leading edge |
| $\Delta L_{\text {int }}$ | Interference lift due to power |
| $\ell$ | Longitudinal distance between nacelle inlets ar.f moment reference point |
| $\ell 1$ | Lateral moment arm of left hand ram drag with respect tc c.g. (fig. 75a) |
| $\ell_{2}$ | Lateral moment arm of right hand ram drag with respect to c.g. (fig. 75a) |
| ly | Same distance as $\ell_{1}$ or $\ell_{2}$, except that $\beta=0$ |
| $l_{2}$ | Vertical distance of ram drag vector above c.g. (fig. 49a) |
| M | Pitching moment |
| MAC | Mean aerodynamic chord |
| $\mathrm{M}_{\mathrm{R}}$ | Rolling moment (fig. 1) |
| m | Inlet mass flow per second |
| N | Yawing moment |
| NS | Nacelle station. (Stations are called out in incues) |
| OGE | Out of ground eifect |
| OH | Overhang of flap shroud |
| op | Subscript denoting zero power |
| p | Subscript denoting a total effect due to power |
| $\mathrm{P}_{\mathrm{T}}$ | Total pressure |
| $\mathrm{P}_{\mathrm{T}_{\mathrm{f}}}$ | PT at location of fan exit (figure 4) |
| $\mathrm{P}_{0}$ | Free stream static pressure |
| 9 | Free stream dynamic pressure |
| RPM | Revolutions per minute |
| r | Longitudinal displacement of nozzle thrust with respect to nozzle center (fig. 49a) |
| S | Wing area of reference wing |
| STOL | Short takeoff and landing |
| T | Total nozzle exhaust thrust (both rans) |
| $\mathrm{T}_{1}$ | Nozzle exhaust thrust of one fan. Also nozzle exhaust thrust of left fan |
| $\mathrm{T}_{2}$ | Nozzle exhaust thrust of right fan |


| $v$ | Free stream velocity |
| :---: | :---: |
| Vj | Exhaust velocity |
| $v_{\text {TIP }}$ | Fan tip velocity based on rudius 7.63 cm ( 3.0 in ) |
| VTOL | Vertical takeotf and landing |
| $\left(v / v_{j}\right)_{e}$ | Effective velucity ratio |
| w | Width of swivel nozzle exhaust |
| $w_{i}$ | Vertical downwash velocity in stability axis system |
| WCP | Wing chord plane |
| WRP | Wing reference plane |
| WS | Wing station, lateral d : sance from plane of symmetry (called out in inches) |
| $\mathbf{x}$ | Forward direction in body axis system (fig. 1) |
| $\mathbf{x}_{0}$ | Fuselage station, positive aft (called out in inches from fuselage nose) |
| Y | Side force (fig. 1) |
| $y$ | Lateral dimensions in body axis system |
| $\Delta y$ | Lateral distance between the centers of the exits of the two nozzles of one nacelle (fig. 6b) |
| $z$ | Upward direction in body axis system (fig. 1). Also distance downstream behind nozzle exit |
| $z_{0}$ | Vertical distance above FRP (called out in inches) |
| $\alpha$ | Angle of attack of FRP |
| $\beta$ | Sideslip angle (fig. 1) |
| $\gamma$ | Angle of inclination of nozzle side wall (fig. 6b) |
| $\delta$ | Ratic of atmospheric pressure/pressure at sea level standard day |
| $\delta \mathrm{F}$ | Fla. angle |
| $\delta^{\mathrm{N}}$ | Nozzle $\operatorname{mgle}$ ( $0^{\circ}$ full aft, $90^{\circ}$ full down in body axis system |
| $\varepsilon$ | Downwa a angle with respect to free stream direction |
| $\theta$ | Ratio of absolute temperatures at actual atmespheric condition/sea level staindard condition. Also pitch angle, positive nose up |
| $\theta \mathrm{j}$ | Exhaust flow angle with respect to FRP |
| $\rho$ | Atmospheric der.sity |
| $\rho_{j}$ | Density of jet exhaust |
| 0 | Bank angle (fig. l) |
| $\dot{\omega}$ | Inlet weight flow per second |
| $\mu$ | Tip speed ratio, vTIP/v |
| $\eta$ | Turning efficiency (see sketch on page 10) |
| $\chi$ | Angle, defined in fig. 75a |

## MODEL AND APPARATUS

The model represents a 108 scaling of a twin lift/cruise fan aircraft with fans located in wing nacelles. Photographs of the model are presented in figure 2 and a three-view is given in figure 3.

Each of the two model nacelles carries a tip turbine powered TD-457 fan from the Tech Development Company as shown in figure 4. The exit diameter of the fan including the tip turbine exhaust is 15.25 cm ( 6.0 in ). The drive air is supplied internally. The fan axis of rozation is drooped $5.5^{\circ}$ with respect to the fuselage reference plane.

Each nacelle carries two independently rotatable exhaust nozzles (fig.5). The nozzles can be rotated from a full aft cruise position to a forward position for braking. The axis of rotation of the nozzles is inclined with a double angle. Viewed along the fan axis, the axis of rotation is laterally inclined $21^{\circ}$ from horizontal, lying in plane A-A of figure 6 a and within the plane A-A the axis of rotation is rotated aft by $19^{\circ}$.

In order to lowate the nozzle exhaust with respect to the model more conveniently, the nozzle angles were measured in front and side view, and the locations of the exsaust centers were measured longitudinally, vertically, and laterally using the definitions given in figure 6 b . Results are:

*measured from photographs, positive with exhaust directed downward.

* measured from photographs, positive where flow is directed inward towards the nacelle plane of symmetry.

| Nominal Nozzle <br> Angle, <br> ON | Fuselage Station <br> (Deg) | Vertical Station <br> $z_{0}$ |  | Distance $\Delta y$ Between <br> the two Nozzle Centers <br> of one Nacelle |
| :--- | :---: | :---: | :---: | :---: |
| (cm) |  |  |  |  |

The exhaust shape of each nozzle is rectangular (fig.5). The width is 6.22 cm ( 2.45 in ), height 12.45 cm ( 4.9 in ), and the exhaust area is $77.4 \mathrm{~cm}^{2}$ ( $12.0 \mathrm{in}^{2}$ ). The upper and lower walls of each nozzle are parallel. The side walls have a converging angle of $88.33^{\circ}$ as indicated in figure 7.

The figure also illustrates the internal duct shape of the nacelles. Fillets are used to eliminate sharp corners. The duct opening at the plane of rotation is circular and has a diameter of 12.45 cm ( 4.9 in ).

The reference wing planform is trapezoidal (fig. 3). The reference wing has an area of $0.387 \mathrm{~m}^{2}\left(4.17 \mathrm{ft}^{2}\right)$ and a mean aerodynamic chord of 22.75 cm ( 8.96 in ). The true wing planform includes an unswept portion with a constant chord between the fuselage and the nacelle, up to $32 \%$ semispan. The chord length of this section is $35.1 \mathrm{~cm}(13.82 \mathrm{in})$ and the airfoil is a NACA 64 A 412 section extending from the true leading edge to the true trailing edge. Outboard of the nacelles, from $43.3 \%$ to $97.5 \%$ semispan, a $17 \%$ thick supercritical airfoil is used with a camber of $3.15 \%$ and a leading-edge radius of $4.5 \%$ chord. Between the $32 \%$ and $43.3 \%$ semispan stations a transition exists between the comventional inboard wing panel and the supercritical outboard panel. The transition is made using straight lines between points of equal percent chord station.

The incidence of the airfoil reference line on the inboard panel with respect to the fuselage is $5.3^{\circ}$, constant up to a lateral wing station of $43.3 \%$ semispan. The outboard wing panel has a twist and thickness such that all constant percent chord stations fall laterally on a straight line. The incidence at $97.5 \%$ is $2.4^{\circ}$. The incidence is obtained by rotating the airfoil at the $97.5 \%$ semispan station about its $50 \%$ chord $s^{+}$ation; no dihedral is used at this chord station. Due to the twist, the quarter chord element lies below the wing reference plane and has a slight negative dihedral.

The model has a double slotted flap over $70 \%$ of the wing span. The flap chord inboard of 43.38 semispan is constant in absolute value, whereas a constant percentage wing chord is used between 43.3 and $70.0 \%$ semispan. Outboard of the $43.3 \%$ semispan station the detail flap dimensions are as shown in figure 8. A photograph of the flap near the nacelle afterbody is prescnted in figure 9.

Drooped and extended ailerons, from $70 \%$ to about $100 \%$ semispan, are used for one run. Geometric details are given in figure 10. The urper surface inclination of the extension is $20^{\circ}$ with respect to the airfoil chord.

Slotted Krueger flaps are installed inboard of the nacelles with chordwise dimensions as given in figure 11. The span is $9.9 \mathrm{~cm}(3.9 \mathrm{in})$ for each. No lateral gap exists between the flaps and the fuselage sidewall, however a gap exists between the flap and the sidewall of the nacelle because of the curved wing-nacelle fairings. The flap deflection is 60 degrees with respect to the wing chord.

Plain Krueger flaps are attached to the wing leading edge outboard of the nacelle fairings, figure 11. Flap chord is $15 \%$ of the wing chord, and the average deflection is 65 degrees with respect to the wing chord.

Nacelle strakes for the outboard sides of the nacelles are located as shown in figure 12. They were used on only a few test runs.

A few runs were made with exhaust vanes installed at the nozzle exits. A geometric description is given in figure 13.

Fuselage lower surface strakes are shown in figure 14a. The fuselage strakes consist of two walls parallel to the fuselage center line, and, in case of the "strake box", a laterally connecting wall in the front and rear of the two parallel walls.

Nozzle exhaust fences, consisting of an extension of the side wall of the nozzle so as to shield the nacelles from the exhaust flow, are shown in figure 14b.

Landing gear geometry is presented in figure 15. No associated gear doors were simulated.

Fan inlet covers (plugs) are illustrated in figure 16. Their shape is such that the forward portion of the nacelle represents a smooth streamline body.

External exhaust rake geometries and their location with respect to the left hand nacelle are presented in figure 17.

Each nacelle contains a built-in internal total pressure rake just aft of the fan stator (fig. 4). Six total pressure probes are manifolded together to produce one average total pressure for each nacelle. Two static pressure orificies are located in the nacelle wall in the same plane as the leading edge of the total pressure probes. The two static pressures were also manifolded together to produce one average static pressure for each nacelle.

Unless specifically noted otherwise, the configuration tested was equipped with slotted inboard and plain outboard Krueger flaps, but had no aileron droop or extensions, no fences or strakes, no landing gear, no inlet covers, and no external pressure rakes installed.

A boundary layer trip was applied $3.2 \mathrm{~cm}(1.25 \mathrm{in})$ behind the wing and tail leading edges and behind the fuselage and nacelle nose. Number 60 grit was applied $1.6 \mathrm{~mm}(1 / 16 \mathrm{in})$ wio.

The model was tested in the V/STOL facility of the National Aeronautics and Space Administration's Langley Research Center (tests 148, 149, and 152). Wall corrections were applied to all data at forward speed out of ground effect, based on H. Heyson's methods. In ground effect no wall corrections were applied because they were very small.

Calibration tests were made in the wind tumel as well as in the adjacent static test facility to determine static thrust by mas iorce measurements. Most of the calibration tests were made with the .dcelles and the fuselage shielded against the exhaust flow (reference 2). Shielding is necessary so that the forces of the model to be measured are ideally only those of the exhaust thrust. Without shields, entrainment flow and the exhaust plume generate interference forces on the madel.

The shielding used is schematically shown in figure 18a. The wing and other surfaces were deleted or substituted for by minimum structure to minimize forces on the model from interference. The fans were not operated simultaneously to avoid the generation of a "fountain" flow. Also, some calibration tests were made with the nacelle shields removed, but with the fuselage still shielded; see figure 18b. Additional details of the calibration setup are found in reference 2.

The inlet flow rate was calibrated under static conditions. For this calibration, the model inlet was removed and a bellmouth inlet was attached to the front of the fan. Total pressure in front of the bellmouth entrance (equal to room ambient pressure) and the static pressure in the bellmouth throat were measured.

## CALIBRATION RESULTS

The thrust of each of the two nozzles of each nacelle has a component in the axial and normal direction, as well as a small component in the lateral direction because of the inclinations of the axes of rotation of the nozzles (fig. 6a). The calibrated thrust of the exhaust nozzles, however, is defined as being only the resultant force in the $\mathrm{X}-\mathrm{Z}$ plane (parallel to plane of symmetry), leaving out any lateral components that may exist at each swivel nozzle position in this definition; see fig. 6b. The calibrated thrust so defined is presented as the sum of the thrusts of both nozzles per fan.

The calibrated thrust of the right hand fan, corrected for temperature and pressure, is presented in figure 19. The data are presented versus the square of RPM. Such a presentation facilitates detection of deviations from a first order linear theory that treats the fan as if it were a propeller. The figure pertains to conditions out of ground effect. Similar data are given in figure 20 in ground effect, with and without bank angle. Results show that the fan thrust decreases somewhat with closer proximity to the ground. Maximum scatter at high RPM values is in the order of 8 N (approximately 2 lbs.).

The calibrated thrust of the left hand fan is similar except that the ground effect is not as pronounced. Thrust data for this fan are presented in reference 2.

Comparison of the data with shields (as in figure 18a) and without shields (as in figure 18b) yields the exhaust interference on the nacelle. Results are given in figure 21. It is seen that, out of ground effect, the nacelle generates a "suck down" of about $5 \%$ at a nozzle angle of $90^{\circ}$. In ground effect this loss is increased to about $10 \%$

Extaust flow turning angles are shown in figure 22 a and b for the fans with calibration shields. It is seen that the turning angles almost coincide vii: the nozzle rotation angles. Average turning angles of the right and ceft hand fan together are listed in the following table:

| Nozzle angle, $\delta_{N}$ | $(\mathrm{deg})$ | 0 | 23 | 62 | 80 | 90 | 103 | 120 |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Jet de Zlection angle, $\theta_{\mathrm{j}}$ | $(\mathrm{deg})$ | 3.0 | 26.8 | 62.5 | 81.9 | 91.5 | 105.1 | 122.7 |

An assessment of the location of the exhaust thrust vector is made using moment data from the fews with calibration shields. Pitching moment data about $\bar{c} / 4$ are plotted in figure 23 for various RPM's in and out of ground effect. Computed moments of thrust times moment arm to the geometric center of the exhaust are also indicated for comparison. The measured moments out of ground effect show more negative values except at $\delta \mathrm{N}=120^{\circ}$, indicating that the thrust is generally located bohand the center line. The average distance $r$ behind the center is for the left and right hand fan:

| $\delta_{\mathrm{N}}$ | $80^{\circ}$ | $90^{\circ}$ | $103^{\circ}$ | $120^{\circ}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\mathrm{r} / \mathrm{C}$ | 0.0828 | 0.0727 | 0.0247 | -0.0388 |

In ground effect the measured moments are more negative, indicating that the center of thrust is located further behind the geometric nozzle center.

Lateral flow angularities are measured from rake surveys of the outboard nozzle of the left hand fan at distances of $0.40 \mathrm{~d}_{\mathrm{e}}$ and $1.12 \mathrm{~d}_{\mathrm{e}}$ downstream of the nozzle. On one occasion measurements were also taken at $0.03 \mathrm{~d}_{\mathrm{e}}$ distance from the outboard nozzle of the right hand fan. Data are presented in figure 24, indicating only a slight lateral inclination of the exhaust flow.

Additional exit flow characteristics are given in figure 25 where the pressure ratio at the fan exit is given as a function of thrust. The pressures are based on averages of six total pressure rakes and two static orifices, each manifolded together. The data are given only for nozzle angles of 90 and 103 degrees; at other angles oil was suspected to have accumulated in the pressure lines. Results are used in an assessment of -in thrust increasc with forward speed during the discussion of transition characteristics.

Inlet flow data are presented in figure 26 with correction factors for temperature, $\theta$, and pressure, $\delta$, which are unity for standard atmosphere. By first approximation, the inlet flow follows a linear variation with the square root of the thrust. Theoretically this is seen as follows. Assuming a uniform distribution of the flow across the fan exit area:

$$
\begin{aligned}
& \left(\rho_{j} A_{1} v_{j}\right) v_{j}=T_{1} \\
& \left(\rho_{j} A_{1} v_{j}\right)^{2}=\left(\rho_{j} A_{1}\right) T_{1}
\end{aligned}
$$

Herein, the mass flow leaving the fan is equal to $\left(\rho_{j} A_{1} v_{j}\right)$ and consists of the iniet mass flow and the mass added to the system by the fan turbine supply air. A factor $f$ is now defined to exclude the supply air, so that the inlet mass flow of one fan is

$$
\dot{m}_{1}=f \cdot\left(\rho_{j} A_{1} v_{j}\right)=\frac{\dot{e}_{1}}{g}
$$

and thus the weight flow per fan becomes

$$
\dot{\omega}_{1}=\left(f g \sqrt{\rho_{j} A_{1}}\right) \sqrt{T_{l}}
$$

In the present case, figure 26 yields
and thus

$$
\frac{\dot{\omega}_{1}}{\sqrt{T}}=\left(\mathrm{fg} \sqrt{\rho_{\mathrm{j}}}{ }_{1}\right) \approx 1.11 \frac{\mathrm{~N} / \mathrm{sec}}{\sqrt{\mathrm{~N}}}\left(=0.524 \frac{\mathrm{lb} / \mathrm{sec}}{\sqrt{1 \mathrm{~b}}}\right)
$$

$$
\mathrm{f}=\frac{\left(\mathrm{fg} \sqrt{\rho_{\mathrm{j}} \mathrm{~A}_{1}}\right)}{\mathrm{g} \sqrt{\overline{\rho A}_{1} \frac{\overline{\rho_{j}}}{\rho}}}=0.815 \sqrt{\frac{\rho}{\rho_{j}}}
$$

## HOVER DATA AND STATIC TURNING CHARACTERISTICS

Tail-off longitudinal forces and mements in hover are shown in figure 27a for various distances of the model above the ground. The distance, $h$, is the height of the exit center of the nozzle if it were deflected to an angle of $90^{\circ}$ even though the actual nozzle angle may be different. The quantity $d_{e}$ is the diameter of an equivalent circular exhaust area equal to the exhaust areas of the two swivel nozzles fed by one fan. The graphs cover a range of $h / d_{e}$ values up to 6 . Numerical values for out of ground effect, where $h / d_{e}=\infty$, are included in the upper right hand corner of each graph.

The lift characteristics are given in this figure in terms of the lift/ thrust ratio, $\mathrm{L} / \mathrm{T}$, where T is the total calibrated exhaust thrust of both fans. A high RPM case ( $100 \%$ T) and a lower RPM case ( $80 \% \mathrm{~T}$ ) are shown. The lift loss due to aerodynamic interference is 5 to $6.7 \%$ out of ground effect and 10 to $14 \%$ in ground effect. The larger lift loss pertains to the lower RPM case. An almost identical lift loss for high RPM was found in figure 21 for the nacelle alone which suggests, at this nozzle angle, that the nacelle is primarily responsible for most of the lift loss for the whole model.

Similar data are given in figure $27 b$ at various nozzle angles, $\delta_{N}$, but all at high RPM. The data show that L/T decreases from 0.98 to 0.88 in going from $\delta_{\mathrm{N}}=80^{\circ}$ to $\delta_{\mathrm{N}}=100^{\circ}$, out of ground effect, and that ground effect decreases L/T further by approximately $6 \%$ at $h / d_{e}=1.0$. These are readings at zero pitch angle; in ground effect a decrease in $L / T$ is found with increasing pitch angle (fig. 27c, 27d).

In an effort to reduce the lift loss due to ground effect, a strake box was tested on the bottom of the fuselage. The box consisted of sidewalls and a front and rear wall (see fig. 14a), and is designed to capture a "fountain" flow between the two fans. The box changed the ground effect by $17 \%$ from a lift loss to a lift gain at $h / d_{e}=1.0$. The chaige was less at greater ground distances (fig. 28a).

Another sidewall which is formed by an extension of nozzle walls as illustrated in figure 14b, also helps reduce the loss due to ground effect by at least $5 \%$ throughout the range of $h / d_{e}$ values (fig. 28b). This may be caused by moving the region of flow mixing between the exhaust and surrounding air further downstream, thereby reducing the velocity oi entrained air at the lower surface of the nacelle which, in turn, decreases the suction forces at that surface. Application of these sidewalls reduces the lift loss due to interference to only $5 \%$ of the nozzle thrust at $h / d_{e}=1.0$.

A sumary of minimun values of $L / T$ encountered at a nozzle angle of 90 degrees is included in the following table (in addition to asymmetric conditions which will be discussed further below):


Pitching moment characteristics are also presented in the figures. The data show a pitch-up due to ground effect. Interference effects in ground proximity were analyzed in figure 27c, showing that the interference effects increase the stability in pitch. In this figure the computed curves without interfereace effects are based on $\theta_{j}=90^{\circ}$ and on a varying value of $\mathrm{r} / \tilde{\mathrm{c}}$ as a function of nozzle attitude as derived from figure 23 for close ground proximity.

The effect of the horizontal tail on pitching moment in hover, in or out of groumd effect, is negligible; see figure 29.

The discussion of the above hover data pertains to symmetric flow conditions. However, with the aircraft banked $10^{\circ}$ (no strake box, $80 \%$ thrust setting) a minimum value of $\mathrm{L} / \mathrm{T}=0.80$ is reached as compared to 0.86 without bank (see fig. 30a and 27a) indicating that the lift loss in bank is larger. It is believed that the fountain existing between the exhausts of the two fans is oblique; it escapes the fuselage bottom somewhat, thereby generating less positive pressure which, in turn, cannot overcome the suckdown, see sketch:


Using the strake box to capture the fountain increases the lift to $\mathrm{L} / \mathrm{T}=$ 0.84 in the banked condition, as compared to 0.80 without strake box, see figure 30a.

Application of roll control in the banked condition, in the direction against the bank, reduces $L / T$ to 0.72 at $h / d_{e}=1.0$ without the strake hox (fig. 30b). This may be due to an increased obliqueness of the fountain when control is applied. The amount of control is based on an approximate $20 \%$ fan thrust increase on one side and about $20 \%$ decrease on the other. The strake box limits the lift loss to $\mathrm{L} / \mathrm{T}=0.78$ as shown in the same figure. Results are included in above table.

This trend of the above lift characteristics has also been shown in previous tests of a twin fan VTOL model, using water as a test medium rather than air (see ref. 3). Results uf this previous test were used in selecting test runs for the present test.

Roll control characteristics are also affected by ground effect. Basic data for wings level, without strakes are given in figure 30c, and are analyzed in figure 30d. The control effectiveness is reduced to $81 \%$ at the critical distance of $h / d_{e}=1.6$. This is with respect to the rolling moment computed from the thrust differential. Use of strakes changes the effectiveness to $72 \%$ as shown in the same figure. In the banked condition at a height of $\mathrm{h} / \mathrm{d}_{\mathrm{e}}=1.0$ the roll control moment remaining is about 61\%, as compared to $81 \%$ wings level, and is roughly unaffected by the strake box (fig. $30 \mathrm{~b}, \mathrm{~d}$ ). The percentage improves rapidly with increasing ground distance.

The effect of the landing gear on hover characteristics is presented in figure 31. The effect is negligible.

Static turning characteristics are given in figure 32, based on data obtained at the highest RPM values tested, and pertain to conditions out of ground effect with the landing flap deflected and no strakes or nozzle wall extensions. The thrust efficiency and turning angle are defined in the following sketch, where $F_{x}$ and $F_{z}$ are measured model forces with interference effects, and where $T$ is the total calibrated fan thrust.


The thrust efficiency decreases by 48 while increasing the nozzle angle from $90^{\circ}$ to $123^{\circ}$, and increases $6 \%$ in going from $90^{\circ}$ to $23^{\circ}$. The turning angles are essentially equal to those of the calibration data with shields installed (fig. 22), i.e., interference effects do not change the turning angles significantly.

Included in figure 32 is a case with zero nozzle deflection where the exhaust impinges on the flap (without EBF deflector). As expected, the thrust efficiency is considerably lower, and the turning angle significantly different from zero.

## LONGITUDINFI POWER-OFF DATA

Power-off data were measured in conjunction with power-on data so that it is possible to isolate power effects. The power-off data are given in figures 33 and 34 for flaps up and down, with and without fan inlets covered, all without ground effect. The data are pres ${ }^{\prime}$ nted in the stability axis system.

Flaps up, without leading edge devices and with the landing gear retracted, a tail-off CLMAX of 1.32 is obtained (fig. 33, fans windmilling). The increase in the maximum lift coefficient due to flap deflection and full snan leading-edge devices is $\Delta C_{\text {LMAX }}=1.24$, bringing the CLMAX to 2.56 (see sumary in figure 35). The tumnel dynamic pressure in this condition was $q=2769 \mathrm{~N} / \mathrm{m}^{2}\left(60 \mathrm{lb} / \mathrm{ft}^{2}\right)$.

Inlet covers produced, in all cases, a decrease of CLMAX (fig. 34 and 35). In a number of runs, covers were used for drag investigations, showing a reduction in drag.

Flaps-up, the tail-off aerodynmic center for low positive angles of attack is located at $12 \%$ MAC, and tail-on at $49 \%$ (fig. 33a). The stall is stable with tail on and off, although a destabilizing trend exists at high angles of attack. Flaps-down, the tail-off derodynamic center is located at about $14 \% \mathrm{MAC}$, and teil-on at $36 \%$ (fig. 3jc). Addition of the flaps made the model stall unstable with and without tail.

Additional power-off characteristics, such as downwash and tail effectiveness are included in summary plots in the subsequent section.

LONGITUDINAL DATA FOR TRANSITION OUT OF GROUND EFFECT

In the present report, transition is defined as the speed regime between hover and conventional wing-borne flight.

## Reference Thrust

Either the thrust coefficient $C_{T}=T /(q S)$, or its inverse $q S / T$, is used in the present report to describe the power effects. Herein, the static value of the thrust, $T$, is used and was compluted from calibration data using measured RPM values and applying appropriate corrections for atmospneric temperature and pressure.

The static thrust, rather than the exhaust thrust existing at forward speed, is used because the change of this thrust with speed is small. An indication of the change with speed may be obtained from figure 36. The figure shows the increase of the pressure ratio at the fan exit when RPM is held constant while the tunnel dynamic pressure is increased. This pressure ratio increase can be compared with the static exhaust thrust increase with pressure ratio obtained during calibration, see figure 25. Such a comparison (not explicitly shown here) indicates that a speed increase between the static condition and $\mathrm{C}_{\mathrm{T}}=2$ results in an exhaust thrust increase of approximately $2 \%$.

## Inlet Effects

Inlet effects are included in the power-on test data given. However, the ram drag portion of the inlet effects can be isolated using the inlet mass flow information given in figure 26 and using a procedure discussed in the drag section of this report. The remainder of the inlet effects are considered here as part of aerodynamic interference effects. Also the analysis plots include the ram drag effects unless it is specifically stated that they are excluded.

## Lift

Test data for swivel nozzle deflections of $0^{\circ}, 23^{\circ}, 62^{\circ}, 80^{\circ}, 90^{\circ}$, $103^{\circ}$, and $120^{\circ}$ are given in figures 37 through 44 for conditions out of ground effect.

A discussion of the maximum lift is given in the following. CLMAX is defined as the lift coefficient where the first reversal in the lift curve slope occurs, even though a higher value in $\mathrm{C}_{\mathrm{L}}$ may be reached at higher angles of attack.

Maximum lift information for zero nozzle angle is correlated in figures $45 a$ and 45 b , showing that the $\mathrm{C}_{\mathrm{MAX}}$ increment due to power in the flaps-up case consists primarily of the vertical component of the thrust vector. However, flaps-down, an additional maximum lift increment exists beyond this vector component because an externally blown flap (EBF) principle is involved, rather than a vectored thrust principle. With this principle the exhaust jet impinges on the flap and leaves the flap trailing edge similar to a jet sheet. This generates supercirculation analogous to a jet flap, and is responsible for the additional increase of $\mathrm{C}_{\text {LMAX }}$ with $\mathrm{C}_{\mathrm{T}}$.

The figure also shows the effect of leading-edge devices on the maximum lift at this nozzle angle. The increase in $C_{\text {LMAX }}$ from these devices is only slightly affected by $\mathrm{C}_{\mathrm{T}}$.

Maximun lift information for other nozzle angles, flaps down, is sum with CT is shown in all cases. For these larger nozzle angles this is primיrily caused by the direct thrust contribution to lift.
$\therefore$ more detailed discussion of the lift characteristics now follows. First, the direct thrust contribution is subtracted to icolate the aerodynamic circulation lift, Claero. This is carried out for various nozz ${ }^{\circ}$. angles and as a function of angle of attack in figures 46 a through 46 g , ${ }^{\text {sing }}$

$$
C_{L_{\text {aero }}}=C_{L}-C_{T} \sin \left(\theta_{j}+a\right)
$$

Herein, $C_{L}$ is the coefficient of the total measured and $\theta_{j}$ is the jet deflection angle.

Results show that $\mathrm{C}_{\text {aero }}$ decreases as $\mathrm{C}_{\mathrm{T}}$ increases when the angle of attack is low, except for $\delta_{N}=0$ where the flap is externally blown. The difference between $\mathrm{C}_{\mathrm{T}}=0$ (fans windmilling) and finite $\mathrm{C}_{\mathrm{T}}$ values represents the interference effect of power on the circulation lift. Thus, except for $\delta_{N}=0$, the interference lift coeffizient is negative at low $\alpha^{\prime}$ s.

For a further discussion of the data it may be helpful to introduce a different form of presentation since the significance of th.e lift loss is not easy to assess immediately from above curves. For example, a large drop in $C_{L}$ may be relatively insignificant when multiplied by a very low dynamic pressure. To obtain a proper assessiment of a change in lift the actual lift value, rather than $C_{L}$, should be compared with a tangible quantity such as aircraft weight or nozzle thrust.

In the present report the static thrust is used for comparison, being always relatively high at low transition speeds:

$$
\frac{\Delta \mathrm{L}}{\mathrm{~T}}=\frac{\Delta \mathrm{C}_{\mathrm{L}} \mathrm{qS}}{\mathrm{~T}}=\frac{\Delta \mathrm{C}_{\mathrm{L}}}{\mathrm{C}_{\mathrm{T}}}
$$

The data are now also presented as a function of the inverse of $\mathrm{C}_{\mathrm{T}}$, i.e. $\mathrm{qS} / \mathrm{T}$, so that conditions for zero velocity can be shown graphically. A schematic of such a presentation is given in the fnllowing sketch for the tctal aircraft lift:


Graphs of this type give a clear visibility of the magnitude of lift change with speed suring takeoff, a characteristic which cannot be found in presenting $\mathrm{C}_{\mathrm{L}}$ versus $\mathrm{C}_{\mathrm{T}}$ because of the large changes in $\mathrm{C}_{\mathrm{L}}$ encountered when $q$ is increased at low speed. For example, when speed is increased from zero the large associated drop in $C_{L}$ does not even give an immediate impression as to whether L actually increases or decreases. However, when the data are expressed i:. terms of $\mathrm{L} / \mathrm{T}$ a change in 9 giv : a clear impression of the magnitude and sign of the lift change. Also, for a known thrust and given aircraft weight $W=L$, the dynamic pressure needed for flight equilibrium can immediately be determi, $n d$ in the latter format.

One reason for the selection of $\mathrm{qS} / \mathrm{T}$ as the abscissa is to preserve a ready determination of conventional coefficients from these graphs based on the following equation:

$$
C_{L}=-\frac{L}{q S}=\frac{(L / T)}{(q S / T)}
$$

For example, in a graph showing $1 / T$ versus $q S / T$, the $C_{L}$ val te can be determined by dividing the ordinate $\mathrm{L} / \mathrm{T}$ by the abscissa $\mathrm{qS} / \mathrm{T}$ as is illustrated in the above sketch (note also that $\mathrm{C}_{\mathrm{L}}=(\mathrm{I} / \mathrm{T}) \mathrm{C}_{\mathrm{T}}$ ).

Furthermore, the abscissa $q S / T$ can easily be related to the often quoted physically important velocity ratio $\left(v / v_{j}\right)_{e}$. This is shown as follows: The thrust is

$$
T=\left(v_{j} \rho_{j} A\right) v_{j}
$$

where $A$ is the total nozzle exhaust area when $T$ is the total nozzle thrust. Thus

$$
\frac{q S}{T}=\frac{\frac{1}{2} \rho v^{2} S}{\rho_{j} v_{j}{ }^{2} A}=\frac{S}{2 A}\left(\frac{\rho v^{2}}{\rho_{j} v_{j}{ }^{2}}\right)
$$

Defining

$$
\frac{\rho v^{2}}{\rho_{j} v_{j}^{2}}=\left(\frac{v}{v_{j}}\right)_{e}^{2}
$$

The desired relation becomes

$$
\left(\frac{v}{v_{j}}\right)_{e}=\sqrt{\frac{q S}{T}} \sqrt{\frac{2 A}{S}}
$$

With the geometric property of $A / S=0.080$ for the present model the relation reduces to

$$
\left(\frac{v}{v_{j}}\right)_{e}=0.40 \sqrt{\frac{q \mathrm{~S}}{T}}
$$

This yields the following numerical comparison between $q S / T$ and $\left(v / v_{j}\right)_{e}$ :

| $\mathrm{C}_{\mathrm{T}}$ | 0 | 0.5 | 1 | 2 | 4 | 8 | 16 | 50 | $\infty$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{qS} / \mathrm{T}$ | $\infty$ | 2 | 1 | 0.5 | $0.2 う$ | 0.125 | 0.0625 | 0.020 | 0 |
| $\left(\mathrm{v} / \mathrm{v}_{\mathrm{j}}\right)_{\mathrm{e}}$ | $\infty$ | 0.565 | 0.400 | 0.283 | 0.200 | 0.1415 | 0.100 | .0565 | 0 |

Applying the ahove rela. 1 hips the increase in maximurn lift in terms of LMAX/r versus qS;'T is presented in figure 47a, showing that with an increase in freestream dynamic presscre "? film lift capability of the
present aircr. ${ }^{f}$ increases immediately while keeping the thrust constant. Also, the lift at low angle of attack increases immediately with increase in forward speed (fig. 47b). An angle of attack of four degrees was chosen because at a lesser angle of attack the lift decreases significantly at very low speeds (fig. 42b). The lift presented is the total lift which includes the aerodynamic lift as affected by power, and the direct lift component from the exhaust thrust.

The above form of presentation leads to a different analysis procedure than before. The lift is broken up into various components. The various components are illustrated in figure 47 c , describing the case of $\delta_{\mathrm{N}}=90^{\circ}$ from figure 47b. In general, the total lift consists of the following three contributions:

$$
\mathrm{L}=\mathrm{C}_{L_{o p}} \mathrm{q} S+\mathrm{T} \sin \left(\theta_{j}+\alpha\right)+\Delta L_{i n t}
$$

or

$$
\frac{L}{T}=C_{L o p} \frac{q S}{T}+\sin \left(\theta_{j}+\alpha\right)+\frac{\Delta L_{i n t}}{T}
$$

Herein, CLop is the lift coefficient at zero power. It is seen that, given a certain thrust, the power-off-lift/thrust ratio (first term on the right in the equation) is proportional to $\mathrm{qS} / \mathrm{T}$ because of the dynamic pressure variation. The second term is a constant (assuming a constant angle of attack). In the figure the two terms together are represented as a straight line. An increase above this line indicates positive interference lift, $\Delta \mathrm{Lint}$, and a drop below this line represents a negative interference lift. In the case of figure 47c, the negative interference lift amounts to only about $7 \%$ of the thrist. The loss occurs at speeds where sufficient lift is generated to enable the aircraft to lift weights in excess of those at hover. In this figure, the power-off lift used was that for fans windmilling ( $\mathrm{CT}_{\mathrm{T}}=0$, $\alpha=4^{\circ}, C_{L}=2.00$, fig. 42a).

In figure 47d this interference lift is isolated and sumarized for various nozzle angles and for various angles of attack.

The above results reveal that the interference lift loss of the present model is significantly less than a previously tested model with twin exhausts (ref. 4). In that test, losses in excess of $25 \% \mathrm{~T}$ were encountered for a nozzle location slightly forward of $25 \% \mathrm{MAC}$ ( $\delta \mathrm{N}=90^{\circ}$, flaps down). It is surmised that in the present test the nacelle afterbody shape is such that not as much flow separation exists behind the exhaust pipes at forward speeds.

Drag
Similariy to the lift, the drag can also be broken up into various components. The total drag is

$$
C_{D_{\text {tot }}}=C_{D_{\text {aero }}}+C_{D_{R}}-C_{T} \cos \left(\alpha+\theta_{j}\right)
$$

Herein, $C_{D_{R}}$ is the ram drag, and the last term on the right hand side is the direct nozzle thrust component. In figures 48a through 48 g the aerodynamic drag, $\mathrm{CD}_{\text {aero }}$, is presented by subtracting the ram drag and the thrust component from the total measured drag.

The ram drag was computed using

$$
\mathrm{D}_{\mathrm{p}}=\frac{\dot{\omega} \mathrm{v}}{\mathrm{~g}}=\frac{\dot{\omega}}{\mathrm{g}} \sqrt{\frac{2 g}{\rho}}
$$

where, for two fans (see before):

$$
\dot{\omega}=2\left(\mathrm{fg} \sqrt{\rho_{\mathrm{j}} \mathrm{~A}_{1}}\right) \sqrt{\mathrm{T}_{1}}
$$

and where the subscript 1 denotes one fan. This yields for two fans:

$$
\mathrm{D}_{\mathrm{R}}=2\left(\mathrm{fg} \sqrt{\rho_{\mathrm{j}} \mathrm{~A}_{1}}\right) \sqrt{2 \mathrm{~T}} 1 \frac{\sqrt{q}}{\mathrm{~g} \sqrt{\rho}}
$$

and because $2 \mathrm{~T}_{1}=\mathrm{T}$ :

$$
\frac{\mathrm{D}_{\mathrm{R}}}{\mathrm{~T}}=\frac{2\left(\mathrm{fg} \sqrt{\rho_{\mathrm{j}} \mathrm{~A}_{1}}\right)}{\mathrm{g} \sqrt{\rho_{\mathrm{S}}}} \sqrt{\frac{\mathrm{~A}}{\mathrm{~S}}}
$$

Note that conversions between the forms of presentation can be carried out using:

$$
C_{D}=\frac{D}{q S}=\frac{D / T}{q S / T}=\left(\frac{D}{T}\right) C_{T}
$$

so that

$$
C_{D_{R}}=\frac{D_{R}}{q S}=\left(\frac{D_{R}}{T}\right) C_{T}=\frac{2\left(f g \sqrt{\rho_{j} A_{1}}\right)}{g \sqrt{\rho_{S}}} \sqrt{C_{T}}
$$

in which, from calibration results:

$$
\mathrm{fg} \sqrt{\rho_{\mathrm{j}} \mathrm{~A}_{1}}=1.11 \frac{\mathrm{~N} / \mathrm{sec}}{\sqrt{\mathrm{~N}}}\left(=0.524 \frac{\mathrm{lb} / \mathrm{sec}}{\sqrt{1 \mathrm{~h}}}\right)
$$

The above figure 48 gives the aerodynamic drag for $C_{T}=0$ (fan windmilling), as well as for finite $C_{T}$ values. An aerodynamic interference drag, $\Delta D_{\text {int }}$, is defined as the difference between the aerodynamic drag $\mathrm{a}_{\mathrm{L}} \mathrm{C}_{\mathrm{T}}=0$ and $\mathrm{C}_{\mathrm{T}}>0$ while the angle of attack is held constant. This drag, in terms of

$$
\frac{\Delta D_{\text {int }}}{T}=\Delta C_{D_{\text {int }}} \frac{q S}{T}
$$

is plotted separately in figure 48 h . When this drag is combined with the direct thrust term the sum represents the total power effect:

$$
\Delta D_{p}=\Delta D_{i n t}+T \cos \left(\alpha+\theta_{j}\right)
$$

This quantity, non-dimensionalized by $T$, is presented in figure 48i, along with the associated lift change due to power, all for zero angle of attack. The data show the lift and drag increment due to power as being relatively constant with speed over a fairly large speed range (a fact that could not have been visible if the data were plotted in customary coefficient form since constants divided by $q S$ show large variations).

## Pitching Moment

The aerodynamic contribution in the tail-off pitching moment has been isolated according to equation

$$
C_{m}=C_{m_{\text {aero }}}-C_{m_{R}}-\frac{T}{q S}-\frac{a}{-}
$$

Herein, aT is the moment arn of the thrust wit? respect to the center of gravity as illustrated in figure 49a. The coefficient $C_{\text {InR }}$ is defined as the ram drag contribution to the pitching moment coefficient, i.e.

$$
\mathrm{C}_{\mathrm{m}_{\mathrm{R}}}=\frac{\mathrm{i}_{\mathrm{z}} \cdot \mathrm{D}_{\mathrm{R}}}{\mathrm{qSC}}
$$

The distance $\ell_{2}$ is the height of the center of the inlet above the center of gravity, and is (see fig. 49a for the definition of aT and $\ell$ ):

$$
f_{z}=l \alpha \quad a_{R}
$$

Because

$$
\mathrm{R}_{\mathrm{R}}=\frac{2\left(\mathrm{fg} \sqrt{\rho_{j} A_{1}}\right)}{g \sqrt{\rho_{S}}} \sqrt{\frac{q S}{T}} \cdot T
$$

the pitching moment coefficient contribution from the ram drag becomes

$$
C_{m_{R}}=\left(\frac{1}{\frac{L}{L}} a-\frac{a_{R}}{c}\right)^{2\left(f g \sqrt{\rho_{j} A_{1}}\right)} \sqrt{\frac{T}{P_{S}}}
$$

Results of $C_{\text {paero }}$ are presented in figures 49 b through d for various nozzle angles and CT values.

Similar to the presentation of aerodynamic lift and drag, the difference between the values for $C_{T}=0$ and $C T>0$ is due to the power effect on the aerodynanic monent, and is presented separately in figure 49f, nondimensionalized, using the form


In the numerical treatment, the following geometric quantities have been used:

$$
\begin{array}{ccccc} 
& \frac{\ell}{\bar{c}}=1.32, \frac{{ }_{\mathrm{a}}^{\mathrm{R}}}{} \\
\overline{\mathrm{c}} & =0.39 \\
\varepsilon_{N} & 80 & 90 & 103 & 120 \\
\frac{a \mathrm{~T}}{\overline{\mathrm{C}}} & 0.227 & 0,157 & 0.077 & -0.058
\end{array}
$$

The values for at/ $\bar{c}$ were determined using thrust inclinations $\theta_{j}$ (as opposed to $\delta N$ ) and using thrust displacements $r / \bar{C}$, both as tabulated in the calibration section.
Stability
Longitudinal tail-on stability characteristics as a function of the inverse of the thrust coefficient are presented in figure 50a, showing a forward shift of the aerodynamic center when speed is decreased (given a constant thrust level). For comparison, tail-off aerodynamic center locations are presented in figure 50 b , revealing a much lesser AC shift. The tail contribution in stability is given in figure 5la showing a large decrease of its effectiveness which is responsible for the forward shift in the tail-on case.

In general the pitching moment curves are quite nonlinear. The nonlinearity is probably due to early local wing upper surface flow separations at positive angles of attack, and early wing lower surface flow separations at negative angles. This may be caused by large flow nonuniformity due to power effects and by low Reynolds mmbers. In many instances it appeared impossible to read reliable slopes and corresponding points have been deleted from figure 50. Examples of cases where slopes have been read or where it was possible to isolate the tail contribution in stability are presented in figures $51 \mathrm{~b}, \mathrm{c}$, and d . These figures cover nozzle angles of zero and 90 degrees, and $\mathrm{C}_{\mathrm{T}}$ values of zero to about 8. The slopes were consistently read at $\alpha \geq 2$ degrees, and care was taken to exclude conditions where the tail was stalled.

Stability changes with power were also computed for comparison with test data. The computations were made for the tail-off case for selected nozzle angles, and results are included in figure 50 b . The computation is based on the assumption that the AC shift is caused only by inlet momentum effects. The equations are derived as follows: From

$$
c_{m_{R}}=\left(\frac{1}{c} \alpha-\frac{a_{R}}{\bar{c}}\right)=\frac{\left(\operatorname{tg} \sqrt{\rho_{j} A_{1}}\right)}{g \sqrt{\rho S}} \sqrt{\frac{T}{q S}} \frac{l}{c}
$$

it follows that

$$
\Delta \frac{\mathrm{dC}}{\mathrm{~m}} \mathrm{~d}_{\mathrm{d}}=2 \frac{\left(\mathrm{fg} \sqrt{\rho_{\mathrm{j}} \mathrm{~A}_{1}}\right)}{\mathrm{g} \sqrt{\rho_{\mathrm{S}}}} \sqrt{\frac{\mathrm{~T}}{\mathrm{aS}}} \frac{l}{\bar{c}}
$$

where $T$ is the total exhaust thrust of two fans. The lift curve slope without interference is

$$
\frac{\mathrm{dC}_{\mathrm{L}}}{\mathrm{~d} \alpha}=\frac{\mathrm{d}\left[\mathrm{~T} \sin \left(\theta_{j}+\alpha\right)+C_{\mathrm{LOp}} \mathrm{qS}\right]}{\mathrm{qS} \cdot d \alpha_{\mathrm{rad}}}=C_{\mathrm{L}_{\alpha_{\mathrm{op}}}}+\frac{\mathrm{T}}{\mathrm{qS}} \cos \left(\theta_{j}+\alpha\right)
$$

so that the aerodynamic center shift becomes (positive for a forward shift):

$$
\Delta \frac{d C_{m}}{d C_{L}}=\frac{\left.2 \frac{\left(f g \sqrt{P_{j} A} 1\right.}{}\right)}{g \sqrt{P S}} \cdot \frac{\ell}{\bar{c}} \sqrt{\frac{T}{a S}} C_{L_{\alpha_{o p}}}+\left(\frac{T}{Q S}\right) \cos \left(\theta_{j}+\alpha\right) \quad=\frac{\left(f g \sqrt{\rho_{j} A}\right)}{g \sqrt{\rho S}} \cdot \frac{1}{c} \sqrt{\frac{q S}{T}}
$$

In figure $50 b$ this shift was added to experimentally obtained power-off locations. (The power-off locations are indicated at $\mathrm{qS} / \mathrm{T}=\infty$ ).

The large difference between the above computed AC locations and the experimentally obtained locations shown in this figure for of $=90^{\circ}$ may result from a positive pressure field just ahead of the nozzle exhaust; see sketch below. However, no pressure measurements were made that could verify this.


Positive pressures
(+)

Note that moment characteristics presented in the form $M /(T \bar{c})$ can be converted to regular coefficient form using

$$
\begin{gathered}
C_{m}=\frac{M}{q S \bar{c}}=\frac{\left(\frac{M}{T \bar{c}}\right)}{\left(\frac{q S \bar{c}}{T \bar{c}}\right)}=\frac{\left(\frac{M}{T \bar{c}}\right)}{(q S / T)}=\left(\frac{M}{T \bar{c}}\right) \cdot C_{T} \\
C_{m_{\alpha}}=\frac{d\left(\frac{M}{T \bar{c}}\right)}{d \alpha} \cdot C_{T} \\
\frac{d C_{i n}}{d C_{L}}=\frac{C_{m_{\alpha}}}{C_{L_{\alpha}}}=\frac{d\left(\frac{M}{T \bar{c}}\right)}{\frac{d\left(\frac{L}{T}\right)}{d \alpha} \cdot C_{T}}=\frac{d\left(\frac{M}{T C}\right)}{d\left(\frac{L}{T}\right)}
\end{gathered}
$$

## Downwash and Tail Effectiveness

A downwash correlation is presented in figure 52a for various nozzle angles and $C_{T}$ values. They are generally obtained at intersections of tail off and tail-on pitching moments to avoid inaccuracies due to interpulations of nonlinear data.

The theoretical downwash shown in this figure is based on a
jet flap theory with full span blowing and with the downwash located at an infinite distance behind the wing. Calling $w_{i}$ the vertical velocity at infinity of a stream tube with diameter $b$ representing the virtual mass, the following is obtained:

$$
w_{i}\left(p v \pi \frac{b^{2}}{4}\right)+T \sin \epsilon=L
$$

or

$$
\frac{w_{i}}{v}\left(\frac{1}{2} \rho v^{2} \pi \frac{b^{2}}{S} \cdot s\right)+2 T \sin \epsilon=2 L
$$

where $\varepsilon$ is the downwash with respect to the free stream direction. Substituting for small angles:

$$
\begin{gathered}
\frac{w_{i}}{v}=\tan \epsilon \approx \epsilon \\
\sin \epsilon \approx \epsilon
\end{gathered}
$$

the above relation becomes

$$
\epsilon=\frac{2 C_{L}}{\pi A R+2 C_{T}}
$$

or, in non-conventional form:

$$
\epsilon=\frac{2 \mathrm{~L} / \mathrm{T}}{\left(\frac{\Phi}{\mathrm{~T}}\right) \pi \mathrm{AR}+2}
$$

Considerable scatter of experimental data points exists about this theoretical relation. lowever, the theory is a simplification and is not really applicable for vectored thrusts at discrete points along the span. Aiso, the average value of the downwash is somewhat larger than the theory indicates which may be due to the wing twist.

Perhaps a better correlation is obtained by removing the direct thrust contribution $\mathrm{T} . \sin \left(\delta_{\mathrm{N}}+\alpha\right)$ from the correlation, see figure 52 b . This curve leads to a downwash due to circulation lift, plus a downwash increment $\Delta_{\varepsilon}$ due to interaction of the jet exhaust with the surrounding flow, figure 52c, i.e.

$$
\epsilon=\left(\frac{d \epsilon}{d C_{L}}\right)_{\text {aero }} \cdot C_{L_{\text {aero }}}+\Delta \epsilon
$$

Herein, the circulation lift $C_{\text {Laero }}$ includes power effects on the circulation.

The tail effectiveness $\partial \mathrm{Om} / \partial \mathrm{i}_{\mathrm{H}}$ is given in figure 53 showing an increase with $\mathrm{CT}_{\mathrm{T}}$ as may be expected from an increased velocity due to entrainments into the exhaust flow.

## Aircraft Modifications

Modifications to the basic vectored thrust aircraft configuration were tested as part of this investigation. These modifications were an externally blown flap (EBF) with flaps extended but with $\delta \mathrm{N}=0$, a hybrid configuration (asymmetric nozzle deflections), a configuration using nozz]e exhaust vanes, and a configuration having a fuselage strake box. The EBF and hybrid are illustrated in figure 54 along with the vectored thrust concept. Detail information of the vanes and the strake box are given in figures 13 and 14a respectively. Basic data are in figures 38 a, 55,56 and 57 .

Test data pertaining to externally blown flaps (EBF) from figure 38 a are analyzed in terms of the lifting capability $\mathrm{L} / \mathrm{T}$ versus nozzle angle for conditions where a net forward thrust or zero thrust exists, see figure 58a. This is done to investigate only cunditions of interest for takeoff. At low nozzle angles, th ift is limited by $\mathrm{s}^{\text {tall }}$ while $\mathrm{D} / \mathrm{T}<0$. At high nozzle angles the lift is limited by $\mathrm{D} / \mathrm{T}=0$. Results are plotted in figure 58 b . At each point a different angle of attack exists. Without making a complete analysis it appears from figure 58b that at $\mathrm{C}_{\mathrm{T}}=4$ the vectored thrust concept ( $\delta_{\mathrm{N}}>20^{\circ}$ ) yields better lifting capabilities, whereas at $\mathrm{C}_{\mathrm{T}}=2$ the externally blown flap (EBF, $\delta_{N}<20^{\circ}$ ) shows promise (a complete analysis would require the application of flight safety margins not considered here). In selecting the points for zero or forward acceleration the ram drag was included in the drag.

Application of thrust deflectors for the EBF configuration did not increase the maximm lift (fig. 58c). The deflectors used were considered to be formed by double fold-out landing gear doors and may not have had the optimum shape, size, or location.

The hybrid configuration is defined as having half the nozzles deflected downward for lift and half deflected aft for external blowing and thrust. It yields an $\mathrm{L} / \mathrm{T}$ value of 1.33 at $\mathrm{C}_{\mathrm{T}}=3.19(\mathrm{D} / \mathrm{T}=0)$, but does not have enough thrust at $\mathrm{C}_{\mathrm{T}}=1.57$ to achieve $\mathrm{D} / \mathrm{T} \leq 0$, see figure 55 . Comparison with the EBF concept in figure 58b shows the hybrid configuration to be inferior.

The effect of exhaust vanes as compared to regular nozzles without vanes is presented in figure 59 for flow angles of about $60^{\circ}$ and for flow angles of $90^{\circ}$ to $100^{\circ}$. The data are obtained from figure 56. Only a negligible affect on CL is shown for about $60^{\circ}$, but for about $100^{\circ}$ a decrease in lift is shown with vanes at low $\mu$, and an increase at high $\mu$. However, the comparison is made at equal fan operating conditions using the speed ratio $\mu=v_{T} / v$ rather than on the basis of equal $C_{T}$ values. This is done because it was not possible to calibrate the thrust with vanes because of their physical interference with calibration shields. In this way it is not necessary to know the accurate thrust for the presentation of the data.

Strake box effects are presented in figure 57, showing an increase in maximum lift, as well as an increase of lift at zero angle of attack. An increase in lift is even found for conditions of equal drag.

## LONGITUDINAL DATA FOR TRANSITION IN GROUND EFFECT

Longitudinal forces and moments in ground effect and as a function of power are presented in figures 60 through 65 for nozzle angles of $0^{\circ}$ through $103^{\circ}$. In addition, data for fuselage strakes are given in figure 66. All data pertain to uings level conditions.

Lift coefficients are presented in figures 67a, and $b$ for $h / d_{e}=1.0$ and $\alpha=0$, and in figure 67 c for $\alpha=4^{\circ}$. Comparison with figure 47 b for $\alpha=4^{\circ}$ shows the lift in ground effect to be lower than out of ground effect.

The addition of parallel strakes shows an increase in lift in all cases tested (fig. 66a, b, and c). A larger increase in lift is obtained using the strake box in close proximity to the ground, as indicated in the same figure.

The tail-off aerodynamic enter for $\delta_{N}=62^{\circ}$ is shifted aft by the ground (fig. 68). At $\delta_{N}=90^{\circ}$ the effect is difficult to determine at low speeds because of the small lift change involved when the aircraft is rotated.

Downwash data are summarized in figure 69, showing that $\varepsilon$ is reduced when the nozzle angle is increased from $0^{\circ}$ to $90^{\circ}$. The quanties $\Delta M$ and $\Delta C_{m}$ in this figure pertain to the difference in pitching moment due to the tail.

## LATERAL-DIRECTIONAL DATA AT FORWARD SPEEDS

Lateral-directional data were acquired on a limited basis using a relatively small range of sideslip angles. Larger sideslip angles required a different model support sting which was outside the scope of the present investigation. All data pertain to zero bank angle conditions.

Basic lateral-directional test data are presented in figure 70 for power-off conditions with zero nozzle angle, and in figure 71 with power and various nozzle angles ( $0^{\circ}, 62^{\circ}$, and $90^{\circ}$ ). They pertain to tail on and off conditions, all out of ground effect. Figure 72, however, contains basic data with ground effect. All data are given in the stability axis system.

The data are given in the customary coefficient form for $\mathrm{CT}_{\mathrm{T}}=0$ to about 4, as well as in the form non-dimensionalized by T or Tb for $\mathrm{C}_{\mathrm{T}}=2$ to about 50. Conversion between the two forms is readily accomplished by using:

$$
\begin{aligned}
& C_{n}=\frac{\left(\frac{N}{T b}\right)}{\left(\frac{q S}{i}\right)}=\left(\frac{N}{T b}\right) C_{T} \\
& C_{l}=\frac{\left(\frac{M_{R}}{T b}\right)}{\left(\frac{q S}{T}\right)}=\left(\frac{M_{R}}{T b}\right) C_{T} \\
& C_{y}=\frac{\left(\frac{Y}{T}\right)}{\left(\frac{q S}{T}\right)}=\left(\frac{Y}{T}\right) C_{T}
\end{aligned}
$$

Similarly, slopes with respect to sideslip angles are converted using

$$
\begin{aligned}
& C_{n_{\beta}}=\frac{{ }^{N} \beta}{\mathrm{~Tb}} \cdot \mathrm{C}_{\mathrm{T}} \\
& \mathrm{C}_{\mathrm{l}_{\beta}}=\frac{\mathrm{dM}_{\mathrm{R}} / \mathrm{d}_{\beta}}{\mathrm{T}_{\mathrm{b}}} \cdot \mathrm{C}_{\mathrm{T}} \\
& \mathrm{C}_{\mathrm{Y}_{\beta}}=\frac{\mathrm{Y}_{\beta}}{\mathrm{T}} \cdot \mathrm{C}_{\mathrm{T}}
\end{aligned}
$$

Analysis plots of directional stability, out of ground effect, are presented in figures 73 and 74. Tail-on as well as tail-off data are shown. In the case of the tail-off data both the horizontal and vertical tails are removed.

Generally, power decreases the directional stability when changing $\delta_{N}$ from $0^{\circ}$ to $62^{\circ}$. At $62^{\circ}$ even the tail-on aircraft becomes unstable. At this angle $\mathrm{C}_{n_{\beta}}$ can even become relatively large in the unstable direction when $C_{T}$ is high (fig. 74a, top). However, this value is large not so much because of a large instability, but because of the low dynamic pressure used for reference at the high CT values. To obtain a true impression of the magnitude of instability the value of ${C_{n}}_{\beta}$ should be multiplied by a term containing that dynamic pressure:

$$
N_{\beta}=\left(\frac{N_{\beta}}{q S b}\right) q S b
$$

When non-dimensionalized by thrust (which is always high at low speed), this becomes:

$$
\frac{N_{\beta}}{T b}=\left(\frac{N_{\beta}}{q S b}\right) \frac{q S b}{T b}=C_{n_{\beta}} \cdot \frac{q}{T / S}
$$

It is seen in figure 74a, bottom, $t_{1}$ at at $C_{T}=9$ the value of $N_{B}$ is only about -0.0001 (Tb) per deg, which represents only a slight instability in this tail-on case.

Contrary to the above characteristics fi:r $\delta_{N}=62^{\circ}$, at $\delta \mathrm{N}=90^{\circ}$ the aircraft is generally directionally stable, tail-on, and even stable at $C_{T}=4$ with the tail off; see figure 74b. Tests were repeated and verified these characteristics. It is surmised that the model exhausts act, in this case, as a wall in front of which a positive pressure may build up. This build-up may be greater at the windward side than at the lee side, which then contributes to a positive scability; see sketch below. However, no pressure measurements were made in this area of the model, nor were flow visualizations attempted that may have substantiated this pattern.

When the exhaust flow is bent backwards due to higher free stream dynamic pressure, such as at $C_{T}=2$, this phenomenon disappears. Similarly, when the flow is bent backwards due to ground proximity, this characteristic also disappears; see figure 74c.

ine above variation with $C_{T}$, out of ground effect, is relatively unaffected by the tail except for a general shift in stable direction. The tril contributions in directional stability are about the same for $\delta \mathrm{N}=90^{\circ}$ as ©or $\delta_{\mathrm{N}}=62^{\circ}$. The tail effectiveness increases moderately with $\mathrm{C}_{\mathrm{T}}$; see figure 74d.

A component breakdown, tail-off, for directional stability with power is now given for $\delta_{N}=62^{\circ}$ for analysis purposes. The stability can he broken down into a power-off portion, inlet momentum effects, and interference effects. 'The power-of portion is obtained directly from test data for CT $=0$ (fans windmilling). Expressed in coefficient form it was found that

$$
\mathrm{C}_{\mathrm{n}_{\beta_{\text {op }}}}=-0.0006 \mathrm{deg}^{-1}
$$

Thus, for any given thrust, the actual power-off yawing moment increases with $q$ according to

$$
\frac{{ }^{{ }^{N}} \boldsymbol{\beta}_{o p}}{\mathrm{~Tb}}=\mathrm{C}_{\mathrm{n}_{\beta_{o p}}} \cdot \frac{q}{\mathrm{~T} / \mathrm{S}}=-0.0006 \frac{q}{\mathrm{~T} / \mathrm{S}}
$$

The inlet momentum effect is derived theoretically. This is carried out as follows, using the geometric definitions of figure 75 a .

The moment arm of the center of the inlet of the left fan engine is:

$$
\ell_{1}=a_{N} \cos (X+B)
$$

Calling $D_{\mathrm{K}}$ the total inlet momentum drag of both fans, then the contribution of the left engine only is:

$$
\Delta N_{1}=-\frac{D_{R}}{2}-\ell=\frac{D_{R}}{2} a_{N} \cos (\chi+\bar{\beta})
$$

Similarly, for the right hand fan:

$$
\begin{gathered}
\ell_{2}=a_{N} \cos (\chi-\beta) \\
\Delta N_{2}=-\frac{D_{R}}{2} \ell_{2}=-\frac{D_{R}}{2} a_{N} \cos (x-\beta)
\end{gathered}
$$

The total moment increment $\angle \mathbb{N}$ due to ram drag becomes:

$$
\Delta N={ }_{2}^{D_{R}} a_{N}[\cos (x+\beta)-\cos (x-\beta)]
$$

so that (per radian):

$$
\Delta v_{\beta}=\frac{d \Delta N}{d \beta}=\frac{D_{R}}{2} a_{1}[-\sin (x+\beta)-\sin (x-\beta) \mid
$$

For $\beta=0$, this yields

$$
\Delta N_{\beta}=-D_{R} a_{N} \sin x=-D_{R} \ell
$$

Substituting

$$
D_{R}=\frac{2\left(\mathrm{rg}^{\left(\rho_{j}^{A}\right.}\right)}{g \sqrt{\rho^{\prime} S}} \sqrt{{ }_{T}^{q S}} T
$$

the increment in stability due to inlet momentum effect's becomes then

$$
\frac{\Delta N_{\beta}}{\mathrm{Tb}}=-2 \frac{\left(\mathrm{fg} \sqrt{\rho_{j} A_{1}}\right)}{\mathrm{g} \sqrt{\rho_{S}}} \frac{l}{\mathrm{~b}} \sqrt{\frac{q}{\mathrm{~T} / \mathrm{S}}} \mathrm{rad}^{-1}
$$

Herein is to be used:

$$
\left(\mathrm{fg} \sqrt{\rho_{\mathrm{j}} \mathrm{~A}_{1}}\right) \approx 1.11 \frac{\mathrm{~N} / \mathrm{sec}}{\sqrt{\mathrm{~N}}}\left(=0.524 \frac{\mathrm{lb} / \mathrm{sec}}{\sqrt{\mathrm{lb}}}\right)
$$

The power-off portion as well as the inlet momentum effect is plotted in figure 75b, both labeled "coimputed". Comparison of the sum of these
two portions with power-on test data shows that no significant difference exists between test and computation at that nozzle angle, indicating that interference effects on yawing moment are negligible in this tail-off case.

Similar analyses were made for the side force derivative with respect to $B$ and for the dihedral effect. The side force characteristics are shown in figure 75c. It is seen that the total derivative is iarger than the sum of the power-off portion and the computed inlet momentum effect. This indicates the existence of a power-on interference effect in side forc

Herein, the computation of the power-off portion is made using

$$
\frac{\mathrm{Y}_{\boldsymbol{\beta}_{\mathrm{op}}}}{\mathrm{~T}}=\mathrm{C}_{y_{\beta}} \frac{\mathrm{qS}}{\mathrm{~T}}=\mathrm{C}_{\mathrm{y}_{\boldsymbol{\beta}}}\left(\frac{1}{\mathrm{C}_{\mathrm{op}}}\right)=-0.0092\left(\frac{1}{\mathrm{C}_{\mathrm{T}}}\right)
$$

and the inlet momentum effect $i$ s obtained using:

$$
\begin{gathered}
\Delta y=-D_{R} \cdot \beta \\
\frac{\Delta Y \beta}{T}=\frac{\Delta d Y / d \beta}{T}=-\frac{D_{R}}{T}=\frac{-2\left(f g \sqrt{\rho_{j} A_{l}}\right)}{g \sqrt{\rho \mathrm{~S}}} \sqrt{\frac{\mathrm{qS}}{\mathrm{~T}}} \mathrm{rad}^{-\mathrm{i}}
\end{gathered}
$$

The dihedral effect is presented in figure 75d for zero nozzle angle, nd in figure 75 e for nozzle angles of $62^{\circ}$ and $90^{\circ}$. It is seen that the ram drag contribution is relatively small, and that also the interference effects are small compared to the basic power-off dihedral effect. Herein, the power-off portion is computed from (assuming the thrust is finite and given)

$$
\left(\frac{r_{\mathrm{R}} / \mathrm{d} \beta}{\mathrm{~Tb}}\right)_{\mathrm{op}}=\frac{\mathrm{C}_{\mathrm{l}_{\mathrm{op}}} \mathrm{qSb}}{\mathrm{Th}}=\mathrm{C}_{\beta_{\mathrm{op}}}\left(\frac{1}{\mathrm{C}_{\mathrm{T}}}\right)
$$

while the inlet momentum effect is based on

$$
\frac{\Delta \mathrm{dM}_{\mathrm{R}} / \mathrm{d} \beta}{\mathrm{~Tb}}=\frac{1_{z} \Delta y_{\beta}}{\mathrm{Tb}}=\frac{-1_{z}}{\mathrm{~b}}\left(\frac{\mathrm{n}_{\mathrm{R}}}{\mathrm{~T}}\right)=\frac{-1_{z}}{\mathrm{~b}} \frac{2\left(\mathrm{fg}_{\mathrm{g}} \sqrt{\rho_{\mathrm{j}} \mathrm{~A}_{1}}\right)}{\mathrm{g} \sqrt{\rho \mathrm{~S}}} \sqrt{\frac{\underline{q}}{\mathrm{~T}}} \mathrm{rad}^{-1}
$$

Herein, $\ell_{2}$ is the location of the center of the inlets above the center of gravity. In the present case $\ell Z / b$ is equal to -0.0476 (at zero angle of attack).

## CONCUUSIONS

A powered wind tumel model of $3 \mathrm{~V} /$ STOL transport aircraft with two fans for vectored thrust was tested. The major conclusions of the data analysis are:

Interference lift losses in hover near the ground are about $10 \%$ of the nozzle thrust. This loss can be reduced to onily $5 \%$ when the nozzle exhaust walls are extended somewhat so as to shield the nacelles from exhaust flow interference.

The lift losses increase when the aircraft is banked, as well as when roll control is applied (in the direction against the bank) by increasing the nozzle thrust on one side and decreasing it on the other.

In the transition speed range interference lift losses of only $7 \%$ of the nozzle thrust were encountered ( $\delta_{\mathrm{N}}=90^{\circ}, \alpha=4^{\circ}$, no strakes) which is significantly less than was found with a previously tested twin-jet VIOL model (ref. 4). The loss can be further reduced by employing a strake box.

The lift characteristics can be enhanced by rotiting all nozzles fully aft for external flap blowing, provided that the speed is high enough so that CT is less than approximately 2. However, a hybrid configuration with half the nozzles deflected downward and half deflected aft was not as promising.

Power effects decreased the directional stability at a nozzle angle of $62^{\circ}$, but showed an increase in directional stability at $90^{\circ}$.

The twin fan vectored thrust concept is a feasible concept for a V/STOL 'ransport aircraft in terms of lift and pitching moment characteristics, rovided that the lift losses in hover during momentary banked conditions -an be accommodated.

The reduction of the lateral control effectiveness due to ground effect in hover is an important consideration; additional roll control tests with nozzle shielas and using flow visualization are recommended to investigate the prssibility of improving the roll control.

The adequacy of lateral-directional characteristics at forward speeds cannot be fully assessed from the data due to the limited sideslip angles imposed on the test. At low sideslip angles unknown flow phenomena were encountered at high nozzle angles, but the forces and moments were controllable.

Analy tical methods for the determination of the inlet momentum effects on the lur, itudinal and directional stability appeared to adequately describe the whole power effect for tail-off conditions at low nozzle angles. At high nozzle ar, les, however, additional power effects exist.

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Figure 1. - Axis system.

1sume: *. Whotograph of model,

(b) Sitle view

(c) One-fourth lower view

Figure 2. concludel.

Figure 3. - Three-view of model.

Figure 4. - Nacelle details in side view.

ligure 5. Photograph of nacelle details.

(a) Axes of rotation of nozzies.
ligure b. Nozale angles and nozale exnt locat lons.

(b) Definition of measured nozzle distances and angles.

Figure 6. - Concluded.


Figure 7. Ducting details of nozzle.

ligure 8. - llay geonetry


Tigure 9. Whotograph of flaps.


Figure 10. - Drooped and extended aileron.


Figure 11. - leading edge flaps.

figure $2 \therefore$ letaile of natelle st rahes.

!'gure 13. - Exhaust vane details.

(a) Fuselage strakes.

Figure 14. - Fuselage strakes and exhaust fences.

(b) Nozzle exhaust fences.
ifgure 14. - Concloded.


Figure 15. - Landing gear dimensions.


Figure 10. - Fan inlet covers.


Figure 17. - Whalust raken.

(a) Shields installed

(b) Shields removed

Figure 18. - Schematic of calibration test setup (not to scale).


Figure 19. - Cal ibrated thrust of right-hand fan (starboard), (Gil.
(N) (1b)


Figure 20. - Calibratcd thmst of right-hand fan in ground effect, $\delta_{N}=90^{\circ}$

(b) Bank angle $10^{\circ}$

Figure 20. Concluded.


Figure 21. - Nacelle interference on static thrust of right-hand fan.

(a) Right-hand fan.

Figure 22. - Calibrated flow turning chamateristas.

> h/de $=\infty$ Flain symbol
> ————— h/de $=1.0$ Flagged symbol

(b) Left-hand fan.

Figure $22 . \quad$ Concluded.

(a) Right-hand fan.

Figure 23.- Thrust versus pitching moment of fan during calibration.

hyure $2: 3$. Conchuded.


Figure 24. . . Vozz1c $\because t$ total pressure profile for $\delta_{N}=90^{\circ}$.


Figere ? i. - Thrust of right hand fan versus fan exit pressure ratio.


Figure 26. - Inlet flow rate of right-hand fan, out of ground effect.


Figure 27.- Tail-off longitudinal forces and moments in hover.

(c) Effect of pitch angle at $\mathrm{h} / \mathrm{d}_{\mathrm{e}}=1.0 . \delta_{\mathrm{N}}=90^{\prime \prime}$.

Figure 27.- Continued.

(d) Effect of pitch angle at $h / d_{e}=1.2$. Various nozzle angles.

Figure 27.- Concluded.

(a) Effect of strake box.

Figure 28. - Effect of shiekds on tail-off longitudinal forces and moments in hover. $\theta=0, \Phi=0, \delta \mathrm{~N}=.90^{\circ}$.



(b) Effect of nozzle wall extension.

Figure 28.- Concluded


Figure 29. - Tail effect on longitudinal forces and moments in hover. $\theta=0, \delta \mathrm{~N}=90^{\circ}$.

(a) Aircraft banked, no roll control input, $80 \%$ thrust setting, strake box on and off.

Figure 30.- Tail-off forces and moments in asymmetric hover condition. $\delta_{\mathrm{N}}=90^{\circ}, \theta=0$.

(b) Aircraft banked, with corrective roll control input. With and without strakes or strake box.

Figure 30.- Continued.

(c) Aircraft level, with roll control ipput. With and without parallel strakes.

Figure 30.- Contimed.

(d) Roll control effectiveness.

Figure 30. - Concluded.


Figure 31.- Effect of landing gear in hover. $\delta_{\mathrm{N}}=90^{\circ}$, tail off, $\theta \quad \mathrm{J}, \phi=0, \mathrm{~h} / \mathrm{d}_{\mathrm{e}}=1.2$.

| Sym | $\delta_{N}$ deg |  |
| :---: | :---: | :--- |
| $\square$ | 0 | With flap blowing |
| $\triangle$ | 23 | No flap blowing |
| $\square$ | 62 | No flap blowing |
| $O$ | 90 | No flap blowing |
| $\square$ | 120 | No flap blowing |


rigure 32. - Static turning ${ }^{\text {naracteris:ics }}$ OGE.

(a) Flaps up, tail on and off, fan inlets spen.

Figure 33.- Longitudinal power-cff forces and moments out of ground effect. $\delta_{N}=0$.

Figure 33.- Continued.




$$
\begin{array}{cc}
\text { Tail } & i_{H} \text {, deg } \\
0 \text { off } & - \\
\square \text { on } & -5 \\
\forall \diamond \text { on } & 0 \\
\Delta \text { on } & +10
\end{array}
$$

(c) Flaps down, tail on and off, fan inlets open.

Figure 33.- Concluded.


Figure 34.- Effect of fan inlet covers on power-off longitudinal forces and moments out of ground effect. $\delta_{\mathrm{N}}=0$.

(b) Flaps down.
Figure 34.- Concluded.




Figure 36. - Fan exit pressure ratio versus inverse of thrust coefficient.


Figure 37.- Flaps -up longitudinal forces and monents during
transition out of ground effect. on $=0$, tail off.

(b) With full span leading edge devices.

Figure 37.- Concluded.


| $\mathrm{C}_{\top}$ |  | Krueger flaps |
| :---: | :---: | :---: |
| 0 | 0 | full span |
| $\square$ | 1.92 | full span |
| $\cdots$ | 0 | inb'd only |
| $\pm$ | 1.88 | inb'd only |



(a) Effect of leading edge devir .., tail off.

Figure 38.- Flaps-down longitudinal forces and moren.is during transition out of ground effect. $\delta_{N}=$


|  |  | Nacelle |
| :--- | :--- | :--- |
|  | $C_{T}$ | Strakes |



(b) Effect of nacelle strakes. Tail off, full span Krueger flap. $C_{L}$ as variable.

Figure 38. - Continued.

(c) Effect of nacelle strakes, tail off, $\mathrm{L} / \mathrm{T}$ as variable.

Figure 28. - Continued.



| Tail $i_{H}$, deg |  |
| :--- | ---: |
| O off | - |
| $\square$ on | -5 |
| $\diamond$ on | 0 |
| $\triangle$ on | +10 |


(e) Tail effectiveness at $\mathrm{C}_{\mathrm{T}}=0$.

Figure 38.- Co. inued.




|  | Tail $i_{H}$, deg | $C_{T}$ |  |
| :--- | :--- | :--- | :--- |
| $O$ | off | - | 7.81 |
| $\square$ | on | -5 | 7.76 |
| 0 | on | +10 | 7.61 |




(h) Tail effectiveness at $\mathrm{C}_{\mathrm{T}}=7.7$.

Figure 38.- Concluied.


$\cdot \circ \varepsilon Z=$



(a) Tail off.
Figure 40.- Flaps-Cown longitudinal forces and moments during transition
out of ground effect. $\delta \mathrm{N}_{\mathrm{i}}=62^{\circ}$.



(b) Tail Effectiveness at $\mathrm{CT}=0$.

Figure 40.- Continued.


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コ

$$
\begin{aligned}
& \begin{array}{cc}
\text { Tail iH, deg } \\
0 & \text { off } \\
\square & - \\
\diamond \text { on } & 0 \\
\diamond \text { on } & +10 \\
\Delta \text { on } & +20
\end{array}
\end{aligned}
$$


Figure 40.- Continued.




(d) Tail effectiveness at $\mathrm{C}_{\mathrm{T}}=4.1$.
Figure 40.- Continued.


(e) Tail effectiveness at $\mathrm{C}_{\mathrm{T}}=8.4$

Figure 4l - Concluded.

(a) Tail off.
Figure 41.- Flaps-down $\begin{aligned} & \text { longitudinal forces and noments during transition } \\ & \text { out of gron ld effect. } \delta N=80^{\circ} .\end{aligned}$


$$
\begin{aligned}
& \text { ○｣ }
\end{aligned}
$$





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\begin{aligned}
& \text { 上ロ } \\
& \text { ○くな』1の }
\end{aligned}
$$





3.
 - 于fo tue $\frac{D}{T}$上 Figure 42.- Flaps-down longitudinal forces and moments during transition out of ground effect. $\delta_{N}=90^{\circ}$.
 ${ }_{\text {deg }}^{20} 0^{0.4} \frac{0}{T}_{0.8}^{0.8}$
(b) Tail on, $\mathrm{i}_{\mathrm{K}}=20^{\circ}$.
Figure 42.- Continued.





| $\quad$ Tail $i_{H}$, deg |  |
| :--- | :---: |
| $O$ off | - |
| [] on | 0 |
| $\triangle$ on | +10 |
| $\triangle$ on | +20 |

(c) Tail effectiveness at $C_{T}=0$.

Figure 42.- Continued.

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(d) Tail effect reness at $C_{T}=1.9$
Figure 42.- Continued.
ق

$\begin{array}{lcc}\text { Tail } i_{H} \text {, deg } c_{T} \\ \text { Ooff } & - & 3.73 \\ \text { Jon } & 0 & 3.73 \\ \bigcirc \text { on } & +10 & 3.77 \\ \triangle \text { on } & +20 & 3.73\end{array}$

(e) Tail effectiveness at $\mathrm{C}_{\mathrm{T}}=3.7$.
Figure 42.- Continued.




| Tail $i_{H}$, deg | $C_{T}$ |  |
| :---: | :---: | :---: |
| O off | - | 7.52 |
| $\square$ on | 0 | 7.63 |
| $\diamond$ on | +10 | 7.53 |
| $\Delta$ on | +20 | 7.58 |

(f) Tail effectiveness at $\mathrm{C}_{\mathrm{T}}=7.6$.

Figure 42.- Concluded.


(b) Tail on, $\mathrm{i}_{\mathrm{H}}=20^{\circ}$.
Figure 43.- Concluded.



(b) Tail effectiveness at $\mathrm{CT}=0$.

Figure 44.- Continued.






Figure 44. - Continued

Figure 44.- Continued.




|  | Tail | $i_{H}$, deg | $C_{T}$ |
| :---: | :---: | :---: | :---: |
| 0 | off | - | 7.50 |
| $\square$ | on | 0 | 7.32 |
| $\diamond$ | on | +10 | 7.29 |
| $\Delta$ | on | +20 | 7.38 |

Figure 44. - Concluded.

(a) Flaps up, $\delta_{N}=0$.


(b) Flaps down, $\delta_{\mathrm{N}}=0$.

Higure 4.5. : int

(c) Flaps down, varicus nozzle angles.

Figure 45. - Concluded.

(a) $C_{\text {Laero }}$ for $\delta_{N}=0^{\circ}$.

Figure 46. - Tail-off acrodynamic lift in coafficient form. Flaps down, (OFI.

(b) $\mathrm{C}_{\text {Laero }}$ for $\delta_{\mathrm{N}}=23^{\circ}$.

(c) $\mathrm{C}_{\text {Laero }}$ for $\delta_{\mathrm{N}}=62^{\circ}$.

Figure 46. - Continued.


Figure 46. - Continued.


Figure 46. - Concluded

(a) Maximum Total lift.

Figure 47. - Tail-off lift characteristics in terms of $1 / / \mathrm{I}$, flaps down, Ofil.

(b) Total lift at $a=4^{\circ}$.

Figure 47.-Cont inued.
$C_{L_{O P}}=L i f t$ coeff. at zero power
$\theta_{j}=$ Exhaust angle with respect to FRP

(c) Lift components for $\delta_{N}=90^{\circ}, \alpha=4^{\circ}$.

Figure 47. - Continued.

(d) Interference lift at various values of $\alpha$ and $\delta_{N}$.

Figure $4^{7}$. - Concluded.

(a) $C_{D_{\text {aero }}}$ at $\delta_{N}=0^{\circ}$.

Figure 48. - Tail-off drag components, flaps down, transition OGF.


Figure 48. - Contimuod.

(d) $C_{D_{\text {aero }}}$ for $\delta N=80^{\circ}$.

Figure 48. - Cont inued.


Figure is. Continued.

(f) $C_{D \text { aero }}$ for $\delta \mathrm{N}=103^{\circ}$.

Figure 48. - Continued.

(g) $\mathrm{CD}_{\text {aero }}$ for $\delta_{\mathrm{N}}=120^{\circ}$.

Figure 48. - Continued.

(h) Interference drag at various values of $\alpha$ and $\varepsilon_{N}$.

Figure 48. - Continued.

(i) Lift and drag increment due to power, $\alpha=0$.

Figure 48. - Concluded.

(a) Monent arms of thrust and ram drag.

Figure 49. - Analysis of pitching moments, Tail off, OGE.

(b) Aerodynamic pitching moment in coefficient form, $\delta_{N}=62^{\circ}$.

Figure 49. - Continued.

(c) Aerodynamic pitching moment in coefficient form, $\delta_{N}=90^{\circ}$.

Figure 49. - Cont inued.

(d) Aerodynamic pitching moment in coefficient form, $\delta_{N}=103^{\circ}$.

Figure 49. - Continued.

(e) Aerodynamic pitching moment in coefficient form, $\delta \mathrm{N}=120^{\circ}$.

Figure 49. - Continued.

(f) Interference pitching moment in ratio form.

Figure 49.- Concluded.


Figure 50. - Aerodynamic center location versus thrust coefficient. OCIE, $\alpha \geq 2^{\circ}$.


Figure 50. Concluded.

(a) A.C. shift at $\alpha \geq 2^{\circ}$.

Figure 51.- Effect of horizontal tail on longitudinal stability, OGE.




$$
\frac{2 C_{L} T O T}{\pi A R+2 C_{T}}
$$


(a) Comparison with theory.

Figure 52. - Downwash correlation, OCl:.

(b) Interference downwash.

|  | $C_{T} \approx$ |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| $\delta_{N}, \operatorname{deg}$ | 0 | 2 | 4 | 8 |
| 0 | $\square$ |  |  |  |
| 62 | 0 | $\rangle$ | $\diamond$ | $\diamond$ |
| 90 | 0 |  | 0 | 0 |
| 120 | $D$ |  | 0 | 0 |


(c) Downwash versus circulation lift.

Figure 52. - Concluded.



Figure 53. - Power-on tail effectiveness.

## Vectored thrust concept:



Externally blown flap (EBF) concept:


Hybrid concept:


Figure 54. - Shematics of the EBF, vectored thrust, and hohrid lift/propulsion concepts.




Figure 55. Flaps-down longitudinal forces and moments for nozzles deflected differentially (Hybrid). Transition out of ground effect. Tail off.

Figure 56. - liffect of exhaust vanes on longitudinal characteristics out of ground effect. Tail off.

Figure 57.- Effect of strake box on longitudinal characteristics out of ground effect.

(a) Conditions for comparison.

Figure 58. - Lift comparison between vectored thrust and EBF concepts. tail off, OGI:.

(b) Maximm lift with zero or forward acceleration.

Figure $58 . \quad$ Continued.

-igure 58. Concluded.


Figure 59. - liffect of exinast vanes on laft at $\alpha=1$. Transitun. (xil, tail off.


Figure 60.- Flaps-down longitudinal f., ces and moments during transition in ground effect. $\delta_{\mathrm{N}}=0^{\circ}$.


Figure 60.- Concluded.


Figure 61.- Flaps-down longitudinal forces and moments during transition in ground effect. Tail off, $a=0 . \delta_{N}=23^{\circ}$.

(a) Tail off, $\alpha=0^{\circ}$.

Figure 62.- Flaps-down longitudinal forces and moments during transition in ground effect. $\quad \delta_{N}=02^{\circ}$.

(b) Tail on, $\mathrm{iH}=+10^{\circ}, \alpha=0$.

Figure 62.- Continued.



Figure 63.- Flaps-down longitudinal forces and moments during transition in ground effect. $\delta_{N}=80^{\circ}, a=0^{\circ}$.



Figure 64.- Continued.

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(c) Effect of angle of attack, tail off.
Figure 64.- Concluded.




Figure 65.- Flaps-down longitudinal forces and moments during transition in ground effect. Tail off, $x=0$, $\delta_{N}=103^{\circ}$.

${ }^{C T}$
$\square \quad 2.03$
〉 3.95
$\triangle 8.05$
$\triangle 16.56$
D 52.35
(b) Force and moment ratios.

Figure 65.- Concluded.



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Nolll




Fuselage
Strakes $C_{T}$
O off 51.57
$\square$ parallel 50.46
Sbox $\quad 49.69$
(b) Nozz1e angle $80^{\circ}, \mathrm{C}_{\mathrm{T}}=50.6$.

Figure 66.- Continueu.


Fuselage Strakes $C_{T}$
O off 51.32
$\square$ parallel
49,12
$\bigcirc$ bux
47.83
(i) Nozzle angle $90^{\circ}, C_{T}=49.4$.

Figur. Du. - Conclixed.

| Sym | $\delta_{N}$, deg |  |
| :---: | :---: | :--- |
| $\square$ | 0 | With flap blowing |
| $\triangle$ | 23 | No flap blowing |
| $\bigcirc$ | 62 | No flap blowing |
| $D$ | 80 | No flap blowing |
| $O$ | 90 | No flap blowing |
| $D$ | 103 | No flap blowing |


(a) $C_{L}$ versus $C_{T}, \alpha=0$.

Figure 0. - Lift at low angle of attack. Tail off, flaps dow, transition IGI: $\left(\mathrm{h} / \mathrm{d}_{\mathrm{e}}=1.0\right)$.

| Sym | $\delta_{N}$, deg |  |
| :--- | :--- | :--- |
| $\square$ | 0 | With flap blowing |
| $\triangle \triangle$ | 23 | No flap blowing |
| $\Delta$ | 62 | No flap blowing |
| $\square$ | 80 | No flap blowing |
| $O$ | 90 | No flap blowing |
| $\square$ | 103 | No flap blowing |


(b) $1 / 1$ versus $1 / \mathrm{C}, \alpha=0$.

Higur: " ${ }^{-}$. Cur imad.


Figure 67. - Concluded.


Figure 68. - Ground effect on tail-off aerodynamic center location with power, $\delta_{N}=62^{\circ}, \because=2^{\circ}$.


Figure 69. - Donnuas: correlation in ground effect. $h / d_{e}=1.0, \boldsymbol{\alpha}=0$.



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-
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0
0

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\begin{aligned}
& \bigcirc \diamond \oplus \oplus \\
& \text { ᄃ } \\
& \stackrel{\oplus}{\oplus}
\end{aligned}
$$







Figure 71.- Continued.







(f) Effect of nozzle angle at $\mathrm{C}_{\mathrm{T}}=4.01$. Tail on.

Figure 71. - Concluded.
 $\beta$, de9

Figure 72.- Ground effect on lateral-directional forces and moments during
Figure $\begin{aligned} \text { trarsition. Flaps down, Nozzle angle } 90^{\circ} . a=0 . \quad \Gamma_{-2}=4.01\end{aligned}$


Figure 75. - Power effect on directional stability at $\boldsymbol{\delta}_{N}=0$. OGE, $\alpha=0, \beta=0$

(a) $\delta=62^{\circ}$.

Figure 74. - Power effect on directional stability, flaps down, OGE, $\alpha=0, \beta=0$.

(b) $\mathrm{N}_{\beta} /(\mathrm{Tb})$ versus $1 / \mathrm{C}_{\mathrm{T}}$ for $\delta_{\mathrm{N}}=90^{\circ}$.

(c) $\mathrm{C}_{\mathrm{n}_{\beta}}$ versus $\mathrm{C}_{\mathrm{T}}$ for $\delta_{\mathrm{N}}=90^{\circ}$.

Figure 74. - Continued.

(d) Effect of power on vertail tail effectiveness.

Figure 74. - Concluded.

(a) Description of moment arms and angles.

Figure 75. - Computation of lateral-directional derivatives. Flaps down, tail off, transition $0 G E, \boldsymbol{\alpha}=0$.

(b) Yawing moment derivatives for $\delta_{N}=62^{\circ}$.

(c) Side force derivatives.

Figure 75. - Continued.

(d) Rolling moment derivative for $\delta_{N}=0^{\circ}$.

(e) Rolling moment derivative for $\delta_{N}=62^{\circ}$ and $90^{\circ}$.

Figure 75. - Concluded.

