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Closed-Form Equations for the Lift, Drag, and Pitching-Moment Coefficients of Airfoil Sections in Subsonic Flow

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CLOSED-FORM EQUATIONS FOR THE LIFT, DRAG, AND PITCHING-

MOMENT COEFFICIENTS OF AIRFOIL SECTIONS IN

SUBSONIC FLOW

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SUMMARY

Closed-form equations for the lift, drag, and pitching-moment coefficients of two-dimensional airfoil sections in steady subsonic flow are obtained from published theoretical and experimental results. A turbulent boundary layer is assumed to exist on the airfoil surfaces. The effects of section angle of attack, Mach number, Reynolds number, and the specific airfoil type are considered. The equations are applicable through an angle-ofattack range of -180° to +180°; however, above about -20°, the section characteristic, are assumed to be functions only of angle of attack. A computer program is presented which evaluates the equations for a range of Mach numbers and angles of attack. Calculated results for the NACA 23012 airfoil section are compared with experimental data.

NOMENCLATURE

с	airfoil chord, m
c d	drag coefficient
c _F	friction drag coefficient
^c f	mean skin friction coefficient
c _k	lift coefficient
c (max	maximum lift coefficient (first lift-curve peak)
° m	pitching-moment coefficient about quarter-chord
$\mathbf{c}_{\mathbf{m}_0}$	pitching-moment coefficient at $\alpha = 0$
°s	pressure (form) drag coefficient
K	correlation parameter for airfoil drag due to lift



airfoil section perimeter, m

M Mach number

L

1 c

R

drag divergence Mach number M_{DD}

moment divergence Mach number MMD

Mach number at which trend of increasing lift-curve slope with increas-Μ1 ing M reverses; also, highest Mach number for which (). Is invariant with M

Revuolds number RN

Reviolds number based on free-stream velocity RN

mean value of airfoil pressure coefficient S_A

- airfoil section maximum thickness, m t
 - airfoil section angle of attack, deg
 - cutoff a, above which the slope of cd with M is assumed to be constant, deg -

a corresponding to highest value of M_{DD}, deg

^oPDD ° P MD a corresponding to highest value of MMM, deg

reference angle about which dc_d/dM vs a curve is assumed to be a mirror image, deg

 $\frac{\frac{e}{max}}{de_e/da} + a_0, deg$ airfoil section stall angle, approximated as astall

angle of attack for zero lift, deg 10

ratio of specific heats

INTRODUCTION

Piloted simulation of helicopters or other rotorcraft in real time is paced by the speed of solution of the rotor dynamic equations of motion. One method of increasing computation speed is the use of a hybrid (analog plus digital) computer, with all rotor celculations being accomplished in the analog computer. Unfortunately, this may preclude the normal practice of providing rotor blade airfoil section characteristics in tabular form, since the time required for the very large number of table lookups may negate the expected speed advantage. Therefore, it was determined to derive closed-form equations for the lift, drag, and pitching moment of airfoil sections which

would exhibit, insofar as possible and practical, the effects of angle of attack, Mach number, Reynolds number, and the specific section type. Published airfoil data were used to establish the form of equations which both reflect general crends and accommodate specific section type characteristics where possible. The coefficients of the equations must be obtained from experimental data for the particular airfoil of interest, but the data requirements are not as extensive as for construction of airfoil data tables typically used in helicopter rotor analysis programs.

The equations presented herein are equally applicable to fixed-wing problems, and should be useful in any case where it is either impractical or impossible to use an airfoil section data table. Conversely, if a particular data table is desired but not available, the equations provide a means of generating the table (with angle of attack and Mach number as parameters) from limited input data.

ASSUMPTIONS AND LIMITATIONS

The equations presented herein were derived from experimental airfoil section data obtained under conditions of two-dimensional, steady-state flow. The airfoil surface condition is assumed to be representative of in-service conditions; that is, smooth, but with a turbulent boundary layer over essentially the entire airfoil. Considering the combined effects of manufacturing defects, service wear, environmental deposits, high Reynolds number and freestream turbulence, it is reasonable, if slightly conservative, to assume a fully turbulent boundary layer for all full-scale rotors and wings for which no special laminar flow apparatus is provided (ref. 1). The assumption of a turbulenc boundary layer establishes the form of the incompressible drag coefficient (appendix B).

The subject equations are applicable only for flight conditions in which the free-stream Mach number is less than one. The Reynolds number, based on section chord, is assumed to be greater than 1×10^6 . Certain experimental data for the particular airfoil type under consideration must be available. Two examples of such data are: (1) moment coefficient at low Mach number and zero angle of attack; and (2) at least two values of drag divergence Mach number as a function of section angle of attack. All such requirements are presented in a subsequent section.

Equations are provided in this report for section lift. drag, and pitching moment coefficient through an angle of attack range of $\pm 180^\circ$. However, above about $\pm 20^\circ$, these coefficients are functions only of angle of attack, and are independent of Mach number, Reynolds number, and the airfoil type. This simplification is a consequence of the very small amount of experimental data available at high angle of attack.

It is assumed that the equations presented in this report will be applied only to airfoils whose characteristics follow the same general trends as the experimental data used to derive the equations (see the appendixes). For

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example, the drag divergence Mach number should be at least approximately linear with section angle of attack.

AIRFOIL SECTION FORCE AND MOMENT EQUATIONS

Equations for the lift, drag, and pitching moment of airfell sections through '180° angle of attack, derived in appendixes V. B. and C. respectively, are presented below. The equations are principally empirical in nature, and in the low angle of attack range require as input experimental data for the particular airfoll type. These requirements are discussed in a subsequent section. Assumptions and limitations are given in the preceding section. To facilitate reference to the appropriate section of the appendixes, equation numbers from the appendixes are retained here.

Lift Coefficient

Small angle of attack
$$(\alpha \ge \alpha_{stall})$$
-

$$c_{g} = \frac{dc_{g}}{d\alpha} (\alpha - \alpha_{0}) , \qquad |\alpha| \ge |\alpha_{stall}|$$
(A1)

$$\left(\frac{\mathrm{d}e_{y}}{\mathrm{d}\alpha}\right) = \left(\frac{\mathrm{d}e_{y}}{\mathrm{d}\alpha}\right)_{\mathrm{ine}} \left\{\mu + \frac{t/e}{1+t/e} \left[\mu(\mu-1) + 0.6(\mu^{2}-1)\right]\right\}$$
 (A2)

where

$$M = \frac{1}{\sqrt{1 - M^2}}$$

If $M \ge M_1$, use

$$\frac{de_{x}}{du} = \left(\frac{de_{x}}{du}\right)_{inc} \left\{ \mu + \frac{t/e}{1 + t/e} \left[\mu(\mu - 1) + 0.6(\mu^{2} - 1)^{2} \right] \right\} - (0.45)(M - M_{1})$$
(A3)

in lieu of equation (A2), but set $dc_{\ell}/d\alpha = 0.05$ as a lower limit.

 $\alpha_0 = (\alpha_0)_1 , \quad 0 \le M \le M_1$ (A4)

$$u_0 = (\alpha_0)_1 - \frac{(\alpha_0)_1 - (\alpha_0)_2}{M_1 - M_2} (M_1 - M), \quad M > M_1$$
 (A5)

$$c_{\ell_{\text{max}}} = C_1 + C_2 M + C_3 M^2 + C_4 M^3 + C_5 M^4 + (C_6 + C_7 M^{C_8}) \sin(C_9 + C_{10} M)$$
 (A6)

For negative angle of attack, the form of c_{fmax} is assumed to be given by equation (A6). The sign of c_{fmax} will be negative, and for cambered sections the constant term c_1 will have a different magnitude also.

Large angle of attack- For positive angles of attack,

$$c_{\chi} = 0.813 + \frac{22 - \alpha}{22 - \alpha}$$
, $\alpha_{stall} < \alpha < 22^{\circ}$, (A16)

where

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$$x_{\text{stall}} = \frac{c_0}{dc_0/du} + \alpha_0$$
 (A17)

$$c_{g} = 1.1 - 1.78[0.01745(\alpha) - 0.7853]^{2}$$
, $22^{\circ} \le \alpha < 90^{\circ}$ (A8)

$$e_{g} = -1.1 + 1.78[0.01745(\alpha) - 2.356]^{2}$$
, $90^{\circ} \le \alpha < 160^{\circ}$ (A9)

 $c_{g} = -0.763$, $160^{\circ} \le a < 172.5^{\circ}$ (A10)

$$c_{g} = -5.82[\pi - 0.01745(\alpha)]$$
, $172.5^{\circ} \le \alpha \le 180^{\circ}$ (A11)

For negative angles of attack:

$$c_{1} = -0.813 + \frac{\binom{0}{\max} + 0.813}{22 + \alpha} + \frac{(22 + \alpha)}{\max} + \frac{-22^{\circ}}{\cos} + \alpha + \frac{1}{\max} + \frac{(A18)}{\max}$$

where

$$a_{\text{neg stall}} = \frac{\begin{pmatrix} c_{\beta} \\ max \end{pmatrix}_{-\alpha}}{\frac{dc_{\beta}}{d\alpha} + \alpha_{0}}$$
(A19)

$$c_q = -1.1 + 1.78[0.01745|\alpha| - 0.7853]^2$$
, $-90^\circ < \alpha \le -22^\circ$ (A12)

$$c_{g} = 1.1 - 1.78[0.01745]\alpha - 2.356]^{2}$$
, $-160^{\circ} < \alpha \leq -90^{\circ}$ (A13)

$$\gamma = 0.763$$
, $-172.5^{\circ} < \alpha \leq -160^{\circ}$ (A14)

$$c_{g} = 5.82[\pi - 0.01745|\alpha|]$$
, $-180^{\circ} \le \alpha \le -172.5^{\circ}$ (A15)

Prag Coefficient

Small angle of attack
$$(|\alpha| \le \alpha_{stall})^{-}$$

$$M_{DD} = A + b \alpha , \qquad = \alpha_{P}$$
(B6)

$$M_{DD} = C + D, \qquad = 0$$

$$M_{\rm DD} = C + D C, \qquad (B7)$$

 $M_{DD} \ge 0.3$ for any α

For Mach number less than the drag divergence Mach number,

$$e_d = (e_d)_{M=1}$$
, $M = M_{DD}$, $(t = -stall)$ (B1)

$$(e_{d})_{M=0} = (e_{f}) \left[S_{A} \frac{i}{e} \left(1 + \frac{e_{S}}{e_{F}} \right) + \frac{K(0,0)745[e^{4}/2]^{2}}{(e_{f})} \right]$$
(B3)

$$\Gamma_{f} = \left(\frac{0.455}{\log - R_{N_{eff}}}\right)^{1.45}, \qquad (B4)$$

$$R_{N_{eff}} = R_{N_{eff}} \left(\frac{1}{2} \frac{L}{c}\right) \cdot \overline{S_{A}} , \qquad (B^{5})$$

For Mach number greater than the drag divergence Mach number:

$$c_d = (c_d)_{M=0} + \frac{dc_d}{dM} (M - M_{DD}), \quad M > M_{DD}, \quad \alpha = t_{stall}$$
 (B2)

$$\frac{dc}{dM} = A + B\alpha + C\alpha^2 + D\alpha^3 , \qquad \alpha_R \le \alpha \le \alpha_c$$
(B8)

$$\frac{dc_{d}}{dM} = \left(\frac{dc_{d}}{dM}\right)_{\alpha = \alpha_{c}}, \quad \alpha > \alpha_{c}$$
(B9)

For angles of attack less (more negative) than the reference angle $\alpha_{\mbox{R}},$ substitute

$$\alpha' = |\alpha| + 2\alpha_{\rm R}, \quad \alpha < \alpha_{\rm R} \tag{B10}$$

in equation (B8). Also the negative angle at which dc_d/dM becomes constant is

$$\alpha_{\mathbf{c}}^{\dagger} = -\alpha_{\mathbf{c}}^{\dagger} + 2\alpha_{\mathbf{R}}^{\dagger}$$
(B11)

Large angle of attack-

$$c_{d} = 0.219 - \frac{0.219 - (c_{d})}{15 - (s_{t} + 1)} (15 - (s_{t} + 1)), \quad (s_{t} + 11) = (B13)$$

where

$$a_{stall} = \frac{c_0}{dc_0/d\alpha} + \alpha_0$$
(B14)

$$c_d = 2.18(|\sin \alpha|)^{1+2}$$
, $15^\circ \pm |\alpha| \pm 180^\circ$ (B12)

For negative angle of attack, c_{\max} in equation (B14) is replaced by $(c_{\max})_{-\alpha}$.

Pitching Moment Coefficient

Small angle of attack ($|x| \leq 20^\circ$)-

$$M_{\rm MD} = A + B_{\rm M} + D_{\rm MD}$$
 (C2)

$$M_{\rm MD} = C + D_{\rm A}, \qquad A > A_{\rm P} \qquad (C3)$$

 $M_{\rm MD} \ge 0.3$ for any α

$$c_{\rm m} = c_{\rm m} + \frac{dc_{\rm m}}{d_{\rm c}} (\alpha)$$
, $M = M_{\rm MD}$, $|\alpha| = 4$ stall (C1)

$$c_{\rm m} = -0.077 + \frac{-\frac{1}{3} \text{stall}}{20 - \alpha} (20 - \alpha), \quad 20^{\circ} > \alpha > \alpha \text{stall}$$
 (C4)

$$c_{\rm m} = (c_{\rm m})_{\rm MD} - \frac{(c_{\rm m})_{\rm MD} + 0.077}{0.95 - m_{\rm MD}} (M - M_{\rm MD}), M \sim M_{\rm MD}$$
 (C5)

For negative angles of attack:

$$c_{\rm m} = 0.077 - (c_{\rm m})_{\alpha} \frac{\text{stall}}{\text{stall}} (20 + \alpha) , \quad 20^{\circ} > |\alpha| > |\alpha| \frac{\text{stall}}{\text{stall}} (C6)$$

$$c_{\rm m} = (c_{\rm m})_{\rm MD} - \frac{(c_{\rm m})_{\rm MD} - 0.077}{0.95 - M_{\rm MD}} (M - M_{\rm MD}), M > M_{\rm MD}$$
 (C7)

where α is negative.

Large angle of attack-

$$c_{\rm m} = -9.00802 \ (x - 20) - 0.077 \ , 20^{\circ} \ z = 67^{\circ}$$
 (C3)

$$c_{\rm m} = -0.619[\sin(0.0260\tau - 1.26\tau)]^{(1-2)\tau}, \quad 67^{\circ} \le \tau \ge 162^{\circ}$$
 (C9)

$$c_m = -0.00838(\tau - 162) - 0.320$$
, $162^\circ = -170^\circ$ (C10)

$$c_{\rm m} = 0.0387 \ (a - 170) - 0.387 \ (170^{\circ} - 120^{\circ}) \ (11)$$

For negative angles of attack:

$$c_{\rm m} = 0.00802(|x|^2 - 20) + 0.077, -67^\circ < x < -20^\circ$$
 (C12)

$$c_{\rm m} = 0.619[\sin(0.0260|\alpha| - 1.26)]^{0.398}$$
, $-162^{\circ} \le \alpha < -67^{\circ}$ (C13)

$$c_{\rm m} = 0.00838(|\alpha| - 162) + 0.320$$
, $-176^{\circ} \le \alpha \le -162^{\circ}$ (C14)

$$c_{\rm m} = -0.0387 (|\alpha| - 170) + 0.387 , -180^{\circ} \leq \alpha < -170^{\circ}$$
 (C15)

INPUT DATA REQUIREMENTS

The input data that are necessary for evaluation of the lift, drag, and pitching moment equations can be obtained from airfuil section test results. The specific input variables, and the equations in which they are used, are listed in table 1. (It should be noted that S_A , L/c, c_S/c_F , and K, which are used in the drag equations (B3) and (B5), can be obtained from figs. 7-10, appendix B, respectively, ac functions of section thickness t/c.)

Minimum experimental data (for a particular airfoil) necessary to evalmate the inputs listed in table 1 are:

 c_2 vs α for several M from zero to M_2 c_d vs M for several α c_m vs M for several α

The effort required for data preparation will be reduced if the following curves are also available:

The c_{φ} vs a curves should extend to at least the stall angle. No test data past stall are required because the airfoil section characteristics at Lich angles of attack are assumed to be independent of section type. The appendixes provide a guide to the use of the above data in the preparation of input.

EXAMPLE CASE

The equations presented above were used to calculate lift, drag, and pitching moment coefficients for the NACA 23012 airfoil section. Input data were obtained from wind tunnel test results for the 23012 given in references 2-4, and are listed in table 2. The calculations were done by a digital computer program written to evaluate the equations over a range of Mach number and angle of attack. A FORTRAN listing of the program is presented in appendix D. Section coefficients were calculated through an angle of attack range of -180° to $+180^{\circ}$, and through a Mach number range of 0.0 to 0.9. From these results, c, cd, and cm at M = 0.1, for $\alpha = -20^{\circ}$ to $+180^{\circ}$, are plotted in figure 1.

Calculated and measured (ref. 2) section aerodynamic characteristics at low angle of attack for three Mach numbers, are compared in figure 2. Overall good agreement was realized, although some details of the 23012 behavior were not reproduced, since the equations are intended to be sufficiently general to represent most airfoil sections. No effort was made to tailor the inputs to improve the correlation. Experimental data for the 23012 at high angle of attack were not available. Figure 3 presents measured and calculated section coefficients as a function of Mach number. In this case, the experimental data are from reference 4. Again, generally good agreement was found.

CONCLUDING REMARKS

The airfoil section equations presented in this report were developed principally by fitting curves to a relatively limited set of experimental data. Some accuracy was sacrificed to obtain general applicability to a range of airfoil types. The utility of such equations can be judged only in the context of their intended use. If high accuracy is not required, these equations provid: the ability to calculate airfoil section aerodynamic coefficients at any angle of attack, over a wide range of subsonic Mach numbers.

The principal application of the equations will be in fixed or rotary wing computations for which it is not practical to use graphical or tabulated section data, or for which a data table must be generated from limited input. The equations also provide a method for calculating section lift, drag, and pitching moment at very high angles of attack (i.e., $-180^\circ \le \alpha \le 180^\circ$). Experimental data of this type are available for very few airfoils. Another important application is in the prediction of section characteristics at high subsonic Mach numbers in cases where such data are incomplete.

APPENDIX A

DELIVATION OF LIFT COFFFICIENT EQUATIONS.

SMALL ANGLE OF ATTACK

In the section angle of attack range $0 \le |x| \le 22^\circ$, effects of compressibility, Reynolds number, and the specific airfoil section of interest are considered. For angles of attack below the stall, the section lift curve is assumed to be linear:

$$e_{i} = \frac{de_{i}}{dx} (x - x_{i}) , \qquad x \ge \frac{1}{2} \operatorname{stall}$$
 (A1)

Lift-Curve Slope

Mach number and thickness effects- Kaplan's rule (ref. 3) is used to represent the effects of compressibility and airfoil section thickness on bitt curve slope:

$$\left(\frac{\mathrm{d}e_{ij}}{\mathrm{d}x}\right)_{\mathrm{comp}} = \left(\frac{\mathrm{d}e_{ij}}{\mathrm{d}x}\right)_{\mathrm{fine}} \left\{ u + \frac{t/c}{1+t/c} \left[u(u-1) + \frac{1}{4} (u+1)(u^2-1)^2 \right] \right\}$$

where

$$m = \frac{1}{\sqrt{1 - M^2}}$$

 $\frac{c}{a}$ = section thickness to chord ratio

$\gamma = 1.4$ for air

so that

$$\left(\frac{\mathrm{d}c_{\ell}}{\mathrm{d}\alpha}\right)_{\mathrm{comp}} = \left(\frac{\mathrm{d}c_{\ell}}{\mathrm{d}\alpha}\right)_{\mathrm{inc}} \left\{ \mu + \frac{t/c}{1+t/c} \left[\mu(\mu-1) + 0.6(\mu^2-1)^2 \right] \right\}$$
(A2)

The utility of this relation was investigated by comparing calculated and measured (ref. 2) lift-curve slopes for several airfoils over a range of Mach numbers, as shown in table 3. Lift-curve slope at M above 0.3 was obtained from the measured value at M = 0.3 and the ratio

$$\frac{(dc_{\ell}/d\alpha)_{M>0.3}}{(dc_{\ell}/d\alpha)_{M=0.3}},$$

calculated with equation (A2). The difference between the calculated and the measured lift-curve slope was 10% or less electric for one back such as ter even airfoil.

For the airfoils considered in table 3, and other airfoils procented in reference 2, the trend of increasing lift-burve slope with Mach number reverses above approximately M = 0.8. However, there is no discernible puttern to the rate of decline. References 5 and 6 present "synthesized" airfoil section data (for the NACA 0012 and 0015, respectively) which do exhibit a smooth decline of lift-curve slope above the reversa. This smoothness is due to the manner in which the airfoil characteristics were derived: an iteration between successive assumption of airfoil data values and comparison of calcalated with measured rotor performance. The average mate of decline of lift-curve slope for the two sets of section data is

Let M at which the trend of the lift-curve slope reverses be M_1 . Then, for M above M_1 , the lift-curve slope will be

$$\frac{dc_{c}}{dx} = \left(\frac{dc_{c}}{dx}\right)_{c \text{ omp}} - (0.45) (M - M_{1}), \text{ per deg}$$

$$\frac{dc_{c}}{dx} = \left(\frac{dc_{c}}{dx}\right)_{inc} \left\{ - \frac{t/c}{1 + t/c} \left[\mu(x - 1) + 0.6(- - 1)^{2} \right] \right\} - (0.45) (M - M_{1})$$
(A3)

It is necessary to set a lower limit for equation (A3). Again, the data available (e.g., ref. 2) are not sufficiently regular to provide a trend. Therefore, the lift-curve slope for the highest Mach number presented was measured for several airfoils in reference 2. The average value is 0.05. Thus the lower limit of equation (A3) is assumed to be

$$\frac{dc}{d\alpha} \ge 0.05$$
 per deg

<u>Reynolds number effects</u>- Airfoil section lift-curve slope is a weak function of R_N up to 1×10^5 or 2×10^6 , and is essentially independent of R_N for higher values (refs. 2 and 3). Therefore, for most uses, lift-curve slope caa be assumed to be constant with R_N . The value of incompressible lift-curve slope used in equation (A3) should be selected with consideration for the likely R_N range to be encountered. This value is also a function of the specific airfoil section considered.

Angle of Zero Lift

Mach number effects- The variation of airfoil section angle of zero lift, α_0 , is small with M and may be neglected, until M reaches a high subsonic

value. Then, for most airfoils, a_0 decreases in magnitude with further increases in M. Table 4 presents measured a_0 values for M above M_1 , where M_1 is the highest M in the experimental data for which a_0 remains at a constant value. (The Mach number M_1 is also the Mach number at which the trend of increasing lift-curve slope with increasing M reverses.) The data in table 4 provide no consistent trend. For six of the sections, there is a strong decrease in the magnitude of a_0 with M, for M above M_1 . However, for three of the sections, the decline of a_0 is slight, while for two other sections, the decline of a_0 is reversed as M continues to increase. Since a general expression for a_0 was desired, it was assumed that the variation of section angle of zero lift with M can be adequately represented by a straight line above M_1 :

$$\alpha_{j} = (\alpha_{0})_{1}, \quad 0 \leq M \leq M_{1}$$
 (A4)

$$\alpha_0 = (\alpha_0)_1 - \frac{(\alpha_0)_1 - (\alpha_0)_2}{M_1 - M_2} (M_1 - M) , \quad M > M_1$$
 (A5)

Mach number M_2 is some convenient $M > M_1$ for which α_0 can be determined from the experimental section data. Note that for symmetrical airfoil sections, $\alpha_0 = 0$ even for high Mach numbers. Thus, for symmetrical sections, $\alpha_0 = (\alpha_0)_1 = (\alpha_0)_2 = 0$.

<u>Reynolds number effects</u>. The angle of zero lift (low Mach number) of an airfoil section is determined by the camber. The extensive section data presented in reference 3 indicate that α_0 is not significantly affected by Reynolds number.

Maximum Lift Coefficient

<u>Mach number effects</u>- In the angle of attack range below the first lift peak, the variation of airfoil section maximum lift coefficient, c_{max} , with M may be categorized as: (1) throughout the range of interest, c_{max} decreases with M, or (2) in part of the range, c_{max} increases with M (but decreases otherwise). Experimental data for both types are shown by the solid curves in figure 4, and both can be fitted by polynomials of the form

$$c_{2_{max}} = C_1 + C_2 M + C_3 M^2 + C_4 M^3 + C_5 M^4 + (C_6 + C_7 M^{C_8}) sin(C_9 + C_{10} M)$$
 (A6)

Fits for two specific airfoil sections are shown by the dashed lines in figure 4. For the V23010 - 1.58 (type 1 above), only the first three terms of equation (A6) were required. However, for the VR-7 (type 2 above), all ten coefficients are nonzero. In both cases, the functions were obtained by use of a least-squares curve-fit program.

<u>Reynolds number effects</u>- Figure 5, from reference 2, presents the combined effects of M and F_N for two airfoil sections. It can be seen that the effect of increasing R_N is to shift the complete $c_{\ell_{\text{max}}}$ vs M curve upwards, with the curves tending to collapse together above about $R_N = 3 \times 10^6$ in one case, and above about $R_N \approx 6 \times 10^6$ for the other. Therefore, the constant C_1 in equation (A6) above will be determined (from low M experimental data for the particular section) as a function of the expected R_N range.

Negative angle of attack- Experimental data for c_{\max} at negative angle of attack, as a function of M, are not available for most airfoll sections. It is assumed that the variation of c_{\max} with M at negative a has the same form as for positive A. The sign of c_{\max} is negative, of course, and for cambered sections the constant term C_1 in equation (A6) will have a different magnitude as well as a different sign. The C_1 term is obtained from the experimental lift curve at low M. If no data are available for negative c_{\max} max

$$\left(c_{\hat{x}_{\max}}\right)_{=\alpha} = -\left[\left(c_{\hat{x}_{\max}}\right)_{+\alpha} - 2\left(\Delta c_{\hat{x}}\right)_{\alpha=0}\right]$$
(A7)

That is, for a cambered section, negative c_{\max} is lower in magnitude than the positive c_{\max} by an amount double the lift increment due to the camber. The accuracy of this estimate was evaluated by comparing it with measured values for several airfoll sections, as presented in table 5. The error is less that 10% in six of the seven cases studied.

LARGE ANGLE OF ATTACK

Very little experiments' section lift coefficient data for angle of attack more than a few degrees above the stall angle are available. Section lift coefficients through '180° for the NACA 0012 and the NACA 63A012 are presented in figure 6. A curve fit of these data in four segments is given by

$$c_{0} = 1.1 - 1.78[0.01745(\alpha) - 0.7853]^{2}, \quad 22^{5} \le \alpha \le 90^{\circ}$$
 (A8)

$$c_{\Lambda} = -1.1 + 1.78[0.01745(\alpha) - 2.356]^2$$
, $90^{\circ} \le \alpha \le 160^{\circ}$ (A9)

$$c_s = -0.763$$
, $160^\circ \le \alpha \le 172.5^\circ$ (A10)

$$c_{g} = -5.82[n - 0.01745(a)]$$
, $172.5^{\circ} \le a \le 180^{\circ}$ (A11)

Equations (A8) through (A11) are taken from reference 7, except that (A8) is begun at $\alpha = 22^{\circ}$ rather than 16° since this yields better agreement with the experimental data in figure 6.

It is assumed that airfoil section lift coefficient at large angles of attack is an odd, symmetric function about $\alpha = 0$ (even for cambered sections). Thus c_{β} for large negative angles is given by the above equations, except that the sign of c_{β} is reversed:

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$$c_{1} = -1.1 + 1.78[0.01745]a] = 0.78531^{1}, -90^{\circ} \le a \le -22^{\circ}$$
 (A12)

 $\frac{1}{2} \rightarrow -\frac{1}{2}$

$$c_1 = 1.1 - 1.78[0.01745]a[-2.356]^{\circ}, -160^{\circ} < a \ge -90^{\circ}$$
 (A13)

$$c_{\rm f} = 0.763$$
, $-172.5^{\circ} + a_{\rm f} = -160^{\circ}$ (A14)

$$c_{1} = 5.82[\pi - 0.01745[\alpha]], -180^{\circ} \le \alpha \le -172.5^{\circ}$$
 (A15)

Equations (A8) through (A15) are assumed to hold for all airfoil sections, regardless of camber, thickness, M, or R_{N} .

Due to the meager amount of test data for $|c_{\chi}|$ above stall, the section lift coefficient is assumed to be a straight line between commax and $x = 22^{\circ}$. Then, from equation (A8),

$$c_{1}^{c} = 0.813 + \frac{c_{1}^{c}}{22 - \alpha} (22 - \alpha), \quad \alpha = 3 + \frac{c_{1}^{c}}{22 - \alpha} (A16)$$

where it is assumed that $c_{i_{max}}$ occurs at

.

$$\frac{a_{stall}}{de_{c}/da} + a_{0}$$
 (AP7)

If the particular airfoil section exhibits gradual stall characteristics, equation (A17) does not accurately predict the angle at which e_{max} occurs. However, equation (A17) may still be used in equation (A16), because the purpose of the latter is to approximate the lift curve in the region between the (assumed) linear lift-curve region and the high angle of attack region $(\alpha \ge 22^\circ)$. For negative angles of attack,

$$c_{\chi} = -0.813 + \frac{\left(c_{\chi}\right)_{-\chi} + 0.813}{22 + \alpha_{\text{neg stall}}} (22 + \alpha) , -22^{\circ} \le \alpha \le \alpha_{\text{neg stall}} (A18)$$

where

$$u_{\text{neg stall}} = \frac{\left(c_{\ell}\right)_{-\alpha}}{dc_{\ell}/d\alpha} + u_{\ell}$$
(A19)

APPENDIX B

DERIVATION OF DRAG COEFFICIENT EQUATIONS

SMALL ANGLE OF ATTACK

For airfoil section angle of attack less than the stall angle, section drag coefficient is essentially constant with Mach number below the drag divergence Mach number M_{DD} . Above M_{DD} , the drag itses very steepty. Thus it is assumed that

$$e_d = (e_d)_{M=0}$$
, $M \ge M_{DD}$, $\alpha \ge \alpha_{stall}$ (B1)

$$e_d = (e_d)_{M=0} + \frac{m}{dM} (M - M_{DD}), \quad M \ge M_{DD}, \quad \Delta \ge \Delta_{stall}$$
 (82)

where $(c_d)_{M=0}$, dc_d/dM , and M_{DD} are functions of section angle of attack.

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Drag Coefficient at low M

Section drag coefficient for smooth airfolls at low Mach number, with fully turbulent boundary layer (the important case for most practical applications), may be estimated (ref. 1) by

$$\left(c_{d}\right)_{M+s} = \left(c_{t}\right) \left[S_{A} \left[c_{t} \left(1 + \frac{c_{s}}{c_{F}} \right) + \frac{\kappa (0, 01/4S_{1} \alpha_{s}^{2})^{2+1}}{c_{t}} \right] \right]$$
(B3)

Equation (63) was derived for symmetrical airfoil sections; however, as shown by experimental data presented in reference 3, camber has a small effect on the minimum drag coefficient. Further, an inspection of the section data in reference 3 shows that minimum e_d occurs at approximately $a \neq 0$, for both symmetrical and cambered sections. Therefore, as discussed in reference 1, the e_d vs a relation obtained for a symmetrical airfoil may be applied with good accuracy to an airfoil with the same thickness distribution but a cambered mean line. The factors S_A , L/e, e_g/e_F , and K are constants for a specific airfoil section, and are obtained from graphs provided in reference 1. Those graphs are reproduced here as figures 7 through 10. The furbulent skin triction drag coefficient is represented by

$$\frac{c_1}{t} = \frac{0.455}{\left(\log - R_{N_{\text{eff}}}\right)^{1.58}}$$
(B5)

reference 10. The effective R_N (ref. 1) is given by

$$R_{N_{\text{eff}}} = R_{N_0} \left(\frac{1}{2} \frac{L}{c} \right), S_A , \qquad (B5)$$

where R_{N_D} is based on free-stream velocity and airfoil chord, except where it is specifically taken as 6×10^6 (third term of eq. (B3)).

Drag Divergence Mach Number

The drag divergence Mach number of an airfoil section is defined as the Mach number for which $(dc_d/dM) = 0.1$ as airspeed is increased at coastant section angle of attack. Measured drag divergence Mach numbers for several airfoil sections are presented in figure 11. The M_{DD} data points in figure 11 have been fitted with straight lines, or combinations of two straight lines in the cambered airfoil cases. Correlation is very good for the NACA 4-digit, 5-digit, 65-series, and Wortmann airfoils, throughout the α range for which data were available. However, the data for the NACA 64A-series airfoils bave a shift to different straight-line segments above about 5° angle of attack. Also, the NACA 64A-series data appear to have an abrupt fluctuation at about $\alpha = 10^\circ$. It was not established whether a general trend of 6-series airfoils is represented by these data (itso note that the NACA 65)-215 data are well represented by a straight line). Since a general expression was desired, it was assumed that M_{DD} can be represented by equations of the form

$$M_{\rm DD} = A + B\sigma_{\rm s}, \quad \alpha \ge \alpha_{\rm p} \tag{B6}$$

$$M_{\rm DD} = C + D\alpha , \qquad \alpha < \alpha_{\rm p}$$
 (B7)

where a_{PDD} is the peak of the M_{DD} data. For symmetrical sections, ${}^{d}P_{DD}$ is zero, C = A, and D = -B. Study of experimental data in references 3 and 4 indicates that M_{DD} is never less than 0.3, regardless of angle of attack.

Slope of c_d Curve Above M_{DD}

The slope of the c_d curve above M_{DD} was measured from experimental data for several airfoil sections, and the results are presented in Figure 12. The curves plotted in the figure were obtained from least-squares curve fits of the experimental data. The 4-digit series airfoils are easily represented by curves of the form

$$\frac{de}{dM} = A + B\alpha + C\alpha^2$$

However, the trend of (dc_d/dM) is more complex for the other two sections studied. Curve fits of the form

$$\frac{dc}{dM} = A + B\alpha + C\alpha^2 + D\alpha^3$$

were required to duplicate the rise and then leveling-off of the experimental data. Note that the quadratic equations will turn upwards, and the cubic

equations downward, at high α . Therefore, it is necessary to use a cutoff, α_c , above which (dc_d/dM) is assumed to be constant;

$$\frac{dc}{dM} = A + Bc + Ca^2 + Da^3, \quad a = a_c$$
(B8)

$$\frac{dc}{dM} = \begin{pmatrix} dc \\ \tilde{d} \\ \tilde{$$

Experimental data for (dc_d/dM) through a significant range of negative angle of attack are not available. Therefore, it is assumed that the dc_d/dM curve is a mirror image about a reference angle, a_R . For example, for the NACA 2312 section shown in figure 12, the negative data point is chosen as a_R , and so $a_R = -1^{\circ}$. For symmetrical sections, $a_R = 0$. Thus for species less (more negative) than a_R , substitute

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$$a^* = [a] + 2a_p , \quad a \le a_p \tag{B10}$$

for a in equation (B8). Also, the negative angle at which (dc_d/dM) becomes constant is

$$\alpha_{\alpha}^{*} = -\alpha_{\alpha} + 2\alpha_{R} \tag{B11}$$

LARGE ANGLE OF ATTACK

Experimental section drag coefficients through "180° are presented for two NACA airfoil sections in figure 13. These data are companion to the lift coefficient data in figure 6. The curve fit of the data shown in figure 13 is given by

$$c_{\rm A} = 2.18 (|\sin a|)^{1+7}, \quad 15^{\circ} \le |a| \le 180^{\circ}$$
 (B12)

The lower limit of 15° for equation (B12) is arbitrary, but reflects the fact that the effects of section type and of Mach number are significant only for small α . Equation (B12) was obtained from reference 1, but the constant factor has been increased slightly to provide better correlation with the experimental data in figure 1?. Equation (B12) is assumed to be applicable to all airfoil sections, regardless of section thickness, Reynolds number or Mach number.

If section angle of attack is greater than α_{stall} , but less than 15°, the drag coefficient is assumed to be a straight line between α_{stall} and 15°:

$$c_d = 0.219 - \frac{(c_d)_{astall}}{15 - |a|} (15 - |a|), \quad a_{stall} \le |a| \le 15^{\circ}$$
 (813)

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where α_{stall} is assumed to be the angle for $c_{e_{max}}$:

$$\alpha_{stall} = \frac{c_{s}}{dc_{s}/dx} + c_{0}$$
(B14)

For negative angle of attack, c_{max} in equation (B14) is replaced by $(c_{max})_{-\alpha}$.

APPENDIX C

DERIVATION OF MOMENT COEFFICIENT EQUATIONS

SMALL ANGLE OF ATTACK

The variation of airfoil section pitching moment (about the quarter-chord) with Mach number is small below the moment divergence Mach number, and may be neglected in that region. Also, the moment is essentially independent of Reynolds number. For most airfoils, the slope of the moment coefficient curve is approximately constant (often zero) with angle of attack until the stall angle is reached, whereupon c_m breaks sharply. Therefore, it is assumed that

$$r_{\rm m} = c_{\rm m_0} + \frac{dc_{\rm m}}{d_{\rm A}} (\alpha) , \quad M \leq M_{\rm MD} , \quad |\alpha| \leq \alpha_{\rm stall}$$
 (C1)

where c_{m_0} is c_m at $\alpha \neq 0$. At positive stall, c_m breaks in the negative (nose-down) c_m direction. At negative stall, the reverse is true.

P. C. Law

Moment Divergence Mach Mumber

An inspection of experimental airfoil section pitching moment coefficient plotted against Mach number (e.g., ref. 4) shows that, at a fixed angle of attack, c_m is essentially constant with M until a certain M is reached, whereupon c_m diverges rapidly with further increases in M. A pitching moment divergence Mach number is defined such that M_{MD} is the Mach number for which $|dc_m/dM| = 0.5$ as airspeed is increased at constant section angle of attack. Using this definition, M_{MD} as a function of angle of attack was obtained from experimental data for several airfoil sections, and is plotted in figure 14. Although there is some scatter, the data are reasonably well fitted by straight lines, or combinations of straight lines. Therefore, M_{MD} will be represented by equations of the form

$$M_{MD} = A + B\alpha$$
, $\alpha \ge \alpha_P$ (C2)

 $M_{MD} = C + D\alpha$, $\alpha < \alpha P_{ID}$ (C3)

where $\alpha_{P_{MD}}$ is the peak of the data. For symmetrical sections $\alpha_{P_{MD}} = 0$, C = A, and D = -B Experimental data in references 3 and 4 indicate that M_{MD} is always greater than about 0.3, for any angle of attack.

Above Stall or MMD

There are very few data available for pitching moment above stall or M_{MD} , and no general trends are discernable. It is assumed that c_m varies linearly with a between α_{stall} and 20° (assumed to be the start of the large α region). A linear relation of c_m with M is also assumed between M_{MD} and M = 0.95. At M = 0.95 it is assumed that c_m has reached the same large-a-region value of -0.077 as at $\alpha = 20^\circ$. The selection of M=0.95 for this relation is arbitrary. Thus,

$$c_{\rm m} = -0.077 + \frac{\frac{(\rm cm)}{\alpha} + 0.077}{20 - \alpha} (20 - \alpha), \quad 20^{\circ} > \alpha > \alpha \\ \text{stall} \quad (C4)$$

$$c_m = (c_m)_{MD} - \frac{(c_m)_{MD} + 0.077}{0.95 - M_{MD}} (M - M_{MD}), M > M_{MD}$$
 (C5)

Note that the condition $a \ge a_{stall}$, $M \ge M_{MD}$ may occur, in which case equation (C4) is used to evaluate $(c_m)_{MD}$ in equation (C5). For negative angles of attack, these equations become

$$c_{m} = 0.077 - \frac{(c_{m})_{\alpha}}{20 + \alpha} \frac{stall}{stall} (20 + \alpha), \quad 20^{\circ} > |\alpha| > |\alpha|_{stall} (C6)$$

and

2

$$c_{\rm m} = (c_{\rm m})_{\rm MD} - \frac{(c_{\rm m})_{\rm MD} - 0.077}{0.95 - M_{\rm MD}} (M - M_{\rm MD}), M > M_{\rm MD}$$
 (C7)

where a_{stall} is negative.

LARGE ANGLE OF ATTACK

Experimental section pitching moment coefficients through +180° are presented in figure 15 for the NACA 63A012 and NACA 0012 airfoils. A curve through the data was determined in four sections, as shown in figure 15. The curve is given by

$$c_m = -0.00802 \ (\alpha - 20) - 0.077 \ , 20^\circ \le \alpha \le 67^\circ$$
 (C8)

$$c_m = -0.619[sin(0.0260a - 1.26)]^{0.396}$$
, $67^\circ < a \le 162^\circ$ (C9)

$$c_m = -0.00838(\alpha - 162) - 0.320$$
, $162^\circ < \alpha \le 170^\circ$ (C10)

$$c_m = 0.0387 (\alpha - 170) - 0.387$$
, $170^\circ < \alpha \le 180^\circ$ (C11)

The lower limit of 20° for these large-angle equations is arbitrary.

Section pitching moment at large angle of attack is assumed to be an odd, symmetric function about $\alpha = 0$, even for cambered sections. Thus c_m for large negative angles is given by the above equations, except that the sign of c_m is reversed (c_m is positive)

c _m •	0.00802()1	- 20)	+ 0.077 ,	-57° 2 1 2 -20°	((12))
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c _m	*	$0.619\{\sin(0.0260 \alpha - 1.26)\}^0$. 198 ,	-162° ≤ a × ~67°	(c13)
c m	-	$0.008_{38}(\alpha - 162) + 0.320$,	-170°	≤ a < -162°	(C14)
c _m	•	$-3.0387(\alpha - 170) + 0.387$,	-180°	$\leq \alpha \leq -170^{\circ}$	(C15)

APPENDIX D

COMPUTER PROGRAM FOR EVALUATION OF AIRFOIL

SECTION AERODYNAMIC CHARACTERISTICS

A digital computer program was written to allow rapid evaluation of the airfoil section force and moment coefficient equations presented in this report. The program calculates section lift, drag, and pitching moment coefficients for angles of attack from -180° to $+180^{\circ}$, and for a range of Mach numbers (M \leq 1.0) which may be selected by the user. Input parameters for the specific airfoil type under study are required. Sample input for the NACA 23012 airfoil section is presented in table 2. Calculated output is in the form of lift coefficient, drag coefficient, and pitching moment coefficient tables. In each table, the calculated coefficient is printed as a function of section angle of attack and of Mach number. Calculated aerodynamic coefficients for the NACA 23012 are plotted in figures 2 and 3. These data were calculated by the program as a result of the input data shown in table 2. A FORTRAN listing of the computer program is presented in this appendix.

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lr(ALFs[L.GE.MLFMLF.ANC.ALFSTL.LE.ALF(UT) UCCEMS=ALUM+BLUM#ALFSfL CuM45=L6F+{55a+kL6+{1.+C56F}+{FK+{.01745++85{ALB37L}}+*.237L}} IF(4LFA.LT.C.ù) ČL(M,N)=-C.El3+(ČLMAX+u.El3)/(22;+ALFuLM)=(22. 1130 IF(ALFA.Gc.0.0) (L(M.A)=(.E12+(CLPAX-(.E12)/(22.-ALFULM)*(22. IF (AMACH ... I. AMACHL) ALF Z=ALFZI-(ALFZI-ALFZZ)/(MPACHZ-AMACHZ) CLAmule(APu+TC/(1.+TC)*(AMU=(AMU-1.)+C.6*(APU**z-1.)**z))
If(AMAU+.GT.4MACH1) CLAmCL2-C.45*(AMACH-AMACH1) lf(acfslt.61.st/ful) cufitys=ACDM+cCOK#ALftLl+CCOM# ALfUlf##z LLMxX=CiN+UzN=bPaC++C3N=EPaCH==2+C4N=EMaCF==3+C5N=dMaCH==4 Cil49C0C 111C (LMAA=Ci+UZ=EMAC++C3=EMACH=+z+C4=EMACH==3+C5=EMACH==4 IF (AMACH.GT. LUMN) UD (M.N)=CCM2+CCCCM4 (AMACH-CCM4) ifialfSTL.GE.ALFPLCJ COMNS=AMCL+EMCC+ALFSTL afialfSTL.LT.ALFMCJ COMNS=CMCC+EMCC+ALFSTL 1 + (CON+C 7N*BMACH**CBN1*SIN(CSN+CLCN*BFACH) ł L +(Lo+C7+BMACH++C8)+SIN(C5+C1C+BMACH) IF(MALFA.UT.AUS(ALFCLM)) GL TC 113C IFLANLEA.UT.ABS(ALFSTL)) UC TC 1150 ANGLE CF ATTACK FLR MAX CL PAXIPUN LIFT CUEFFICIENT +L.UM#ALrSTL##2+UCUM#ALFSTL##3 Jf(MAALF-LE-JUMA) CU(M'A)=CUM2 IF (UMACF.LT..001) EMACH=.UCL ANGLE OF ZEAL LIFT IFIALFA.Cc.0.0) GU TC 111C LFILUMNS+LT .U.S. LLMASEU.S LIFI CLEFFICIENT CHAG CUCYFILIENT LL[M,N]=1,LA#(ALFA-ALF2) 0015300C 1120 ALFULMFULMAX/ULAFALFZ L * (AMACHI-AMACH) ALFSTLEALFULM EPACH=APACH 66 TU 1120 61 TU 1160 6C TU 1140 215=ALF2 1 -ALFAI L +ALFA} 1150 55000 0-1510000 001520000 021560000 001640000 001360300 001370000 001410000 031420000 UC154COC CC155CUCC 001651000 00166000 CCIEOCUC 0.0134000 L0135CUC UC13ECOC J 00 6 1 3 J D C UI146C0C 00147000 00148000 C()15000C 00157000 00159000 UC17100C 0.,133CUC 60140000 00143000 04,145COC UC158CUC UCIEICOC LC16200 00163000 UCIESCJC uC16360C 00170000 0 C I 7 2 CU C **JU174CUC** 00175600 00175000 0 C I 77 Cu C 0016700 0014425

[F[ALFA.GE.O.U] CM[M,N]=CP[M,N]-[CM[M,N]+O.C7]]/[U.95-AMUMN]+[AMACH-AMDMN] ORIGINAL PAGE IS OF POOR QUALITY JF14LF4+LT+C+U) [M(F+N)+CF(M+N)+(CF(F+N)+-U17)/(U+95-ANDMN)+(AMACH-AMDMN) ł CHANGE HIGH-ANGLE CCEFFICIENTS IC SUBSCHIPTER FLAM IFIALFSTL.GE.ALCUTF) CCCCMS=ACCM+BCDM+ALFPRM+CCCM+ALFPRM++2 IF(ALFSTL.LI.ALCUTF) CCDCPS=ACDP+ECDM+ALFCUT+CCCP+ALFCUT+*2 LL(M+N) = C. 219-(J. 219-CDSTL)/(15.- ABS(ALFSTL))+(15.-AALFA) IF (ALFA.6E.U.U) CM(M,A)=-U.077+1CPSTL+U.U77)/(2U.-ALFSTL) IF [ALFA.LI.C.O.) CP (M.N)=C.C77-(U.C77-CMSTL)/(20.+ALFSTL) 1 F (AMACF-LE DUMNS) CDSTL=CCM2S 1 F (AMACF-GT-UDMNS) CUSTL=CCM2S+DCCDMS+TAMACP-COMNS) IF (AALFF. UT. ABS (ALFSTL)) GC TC 1170 IF(ALFSTL.GE.ALFREF) GU TC 1165 INCREPENT MACH NUPBER ALFPKM=AES(ALFSIL)+2.C+ALFREF JFIAMACP.LE.AMDMN) GL TC 1200 IF (AALFA.LT.IS.) GC TC 1220 u TC izeC PUPENT CLEFFICIENT IF (34 LFA.46 .20.) GC TC 1200 IF(ALFA.E4.16U. GC TC 1225 66 11 124C ALCUTP=2.0*ALFREF-ALFCUT CP(M, N) = CPC+CCPCA+ALFA CMSTL =CMC+CCMCA#ALFSTL UC211CUC 12CC APACH=APACH+UPACH i f (4 m L F L . . T) +CLDM+ALFCUT++3 +DCUM*ALFCUT**3 IF(ANLFD.LT.20.) UL 4450 AFL APA EL 1270 A=1,NMAP +UCD##ALFPRF##3 CC 1230 A=1, AMAP CC 1235 N=1,NVNP CC(M.N)=LŪ(46,N) LE(**N)*CU(48,N) LL(N,N) = UHA **イト・ト・ト・ド・** * (20.-ALFA) *{<u.+ALFA} CL(3,N)=CLFA GC TU 1175 GL TU 1220 1175 0C214C0C 121C 1235 1240 1270 1165 1230 1225 1220 1250 1160 1173 ÚC21200ÚC CC12900CC CC15CC2CC 0020300000 CCZIOCUCC 002230000 00204000 002130001 06216666 00222000 20251222 CC 2 20 CO C 0018400C 0C186CJC 0C21700C 00187C0C 0021500 00174000 0C179CUC CCIENCOC 0019100 0016200 C 00183000 00185000 CC188CUC 00205000 CCZCECOC 00202000 002CEC0C UCZIAOUC **UC**221 LUC 0620068 0515050 0021610 0021620 0021630 UC21c40 CC1:455 0020816 0020832 0021425 29 I 0 I 0 0 0015320

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ŧ FLRMAT(IEL, 20%; "PITCHING MCMENT CCEFFICIEN! TABLE") IF (ALFA.GE.-25. ANC. ALFA.LT.-10.) ALFA=ALFA+1.5 FIALFA.LI.-25..UR.ALFA.GE.25.) ALFA-ALFA+4.0 F(ALFA.CL.10..ANC.ALFA.LT.25.) ALFA=ALFA.1.5 1234 HAIF (EJAUEC) PALFA(J), (CCJUINI, NELINANF) 1320 hRITc(6,5056) PALFA(J), (CL(J,N), N=1,NPAP) MRITE(C.+4050) PALFA(J), [CP(J,N), N=1,NANP) FCRMAT(IM1, 30x, "CRAG CCEFFICIENT TALLE") FCMAT(1H1, 30X, 'LIFT CCEFFICIENT TABLE') INCKEMENT ANGLE CF ATTACK AKITE(6,4030) (PMACH(N), N=1,NMNP) 4C3U FURMAT (11 , "MACH N.", (12F5.31) nfITE(6,4030) (PRACH(N), N=1,NMNP) 4050_ FURMAT (1H . Fo.L. 2X, (12F5.4)) IF(ALFA.LI.180.) GC TC 1300 PRINI CUT RESULTS 4640 FERMATIIN . "ALPHA") ARACHEAPACH+CPACH CC 131L A=1, AMAP PPACH(N)=APACH PALFA(P)=ALFA NRITE (014425) 06 1320 JEL 0 CC 1330 J=1,M NRITE (6 ,4060) hEITE(6.4070) 10 1340 UHI -1 NELTE (6 ,4040) NRITC (C, 4040) rkIT= { c , 404C) APACHESPACH S TUP r NC 00225000 1260 4 6 2 5 1310 4 010 4 (6 Ú 1340 002270000 002240000 0 C 2 2 8 00 C C 0034600 00236000 002386000 0022600 00229000 CC230C0C 00231000 00232000 0023300C 00235CU <u>C</u> 00240C0C JC241C0C 0 C 2 4 3 C U C 0024900C 00253C0C 000339000 3024200 0C247CUC 00244000 00245000 06246606 0 625200 002337000 00251000 **0025400C** JC248C0C 00250000 00255200 0022524 0022536 0022512

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1ABLE 1 D	ATA REQUIREMENTS
Input data	Equation numbers
	oeificient
(dc _o /da) inc	(A2), (A3)
t/c	(A2), (A3), also figs. 7-10
M ₁ , M ₂	(A3), (A4), (A5)
$(\alpha_0)_1, (\alpha_0)_2$	(A4), (A5)
ce coefficients	
$C_1 + C_{10}$ for +a case	(A6)
$\mathbf{C}_1 + \mathbf{C}_{10}$ for $-\alpha$ case	(A6)
Drag c	oefficient
M _{DD} coefficients A, B. C, D	(B6), (B7)
^a P _{nn}	(86), (87)
R _{N₀}	(85)
dc _d /dM coefficients A, B, C, D	(88)
°c	(B8), (B9), (B11)
a _R	(810), (811)
Moment of	coefficient'
M _{MD} coefficients: A, B, C, D	(C2), (C3)
^a P _{KD}	(C2), (C3 [°]
c _n	(C1)
dem/da	(C1)

Character and the state of the			01 1010 101		
Symbol	FORTRAN name	Value	Symbol	FORTRAN name	Value
	NMNP	10	L/c	RLC	2.035
	SMACH	0-	SA	SSA	1.18
	DMACH	.1	C _S /C _F	CSCF	.037
t/c	тс	.12	К	FK	1.55
RNo	RNO	8.10×10 6	Coeffic	ients for	M _{DD} :
(dc _l /da) inc	CLAI	.100/deg	A	AMDD	0.730
м ₁	AMACH1	.80	В	BMDD	0246
(a ₀) ₁	ALFZ1	-1.20 deg	с	Chidd	.830
M_2	AMACH2	.85	D	DMDD	.0246
(a ₀) ₂	ALFZ2	70 deg	ap	ALFPDD	-2.0°
Coefficients	for +c _{fm}	ax	ac.	ALFCUT	10.0°
C ₁	C1	1.622	⁽¹ R	ALFREF	-2.0°
<i>c</i> ₂	C2	. 337	Coeffic	ients for	dc _d /dM:
C ₃	C3	-2.315	A	ACDM	0.274
C4	C4	.0	B.	BCDM	.0253
C 5	C5	.0	С	CCDM	.00273
۲ _ΰ	C6	.0	D	DCDM	.000264
C 7	C7	.0	Coeffic	lents for	M _{MD} :
C ₈	C8	.0	A	AMMD	0.810
Cg	C9	.0	В	BMMD	026
C10	C10	.0	C	CMMD	.910
Coefficients	for -cl	r.:	D	DMMD	.026
c1	CIN	-1.200	α _{Ρνη}	ALFPMD	-2.0°
C ₂	C2N	250	ciao	СМО	010
C ₃	C 3N	1.716	dc _m /da	DCMDA	.0014/deg
C4	C4N	.0			
C5	C5N	.0			
C ₆	C6N	.0			
C7	C7N	.0			
C ₈	C8N	.0			
Cو	C9N	.0			
C ₁₀	CION	.0			

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	THE OUVERAL ATRIVIL SECTIONS							
Airfoil section	Mach number	Measured ⁽⁾ dc _y /da per deg	Calculated ¹ de _x /da per deg	Difference,				
NACA 0012	0.30 .40 .60	0.103	0.108	0.0				
	.80	.200	.194	-3.0				
NACA 23015	. 30	.100						
· · ·	.60	.105	.105	.0				
NACA 634012	.75	.155	.170	9.7				
INCA UJAUIZ	. 30	.100	.104	-1.9				
	.58 .74	.125	.121	-3.2				
VERTOL	. 30	.12?						
V23010-1.58	.40	.124	.129	4.0				
	.77	. 240	.206	-14.2				
VERTOL VR-7	.30	.110						
	. 62	.138	.139	-2.5				
	./5	.180	.178	-1.1				

 TABLE 3.- MEASURED AND CALCULATED LIFT-CURVE SLOPE

 FOR SEVERAL AIRFOIL SECTIONS

^aReference 2.

^bEquation (A2).

		mond or	L'ERU I	LITI -	
Airfoil section	Mla	(α ₀) ₁ , deg	м	deg	Reference
NACA 2409-34	0.80	-2.5	0.83	-0.5	3
NACA 4409-34	.70	-4.5	.80	-1.9	3
NACA 23012	.80	-1.2	.85	7	2
NACA 23015	.78	-1.2	.80 .83	9 2.0	2
NACA 64A(4.5)08	.81	-3.2	.85 .90 .96	-2.0 9 5	2
NACA 64A608	. 81	-4.5	.86 .90 .96	-2.5 9 5	2
NACA 64A312	.75	-2.3	.80 .85 .90	-2.0 .4 2	2
NACA 64612	.76	-4.5	.80 .85 .90	-2.6 4 3	2 .
V23010-1.58 with T.E. Tab	.77	.5	.82	.2	2
VR-7			.86	. 2	
with T.E. Tab	.75	-2.0	.82 .92	-1.1 -1.7	2
VR-8					
with T.E. Tab	.85	8	.90 .95	6 6	2

TABLE 4. - MEASURED MACH NUMBER EFFECT ON AIRFOIL SECTION ANGLE OF ZERO LIFT

 ${}^{\rm C}\!M_1$ is the highest Mach number in the experimental data for which α_0 remains at a constant value.

Airfoil Section	Measured c a		Measured	Calculated	Error,	
	+1		(Ac) , 1≡)	(c) max	• • • • • •	
NACA 1412 NACA 2412 NACA 2412 NACA 4412 NACA 63_1-212 NACA 63_1-412 NACA 65_1-212 NACA 65_1-412	1.60 1.70 1.64 1.58 1.72 1.45 1.65	-1.20 -1.08 78 -1.18 -1.00 -1.10 80	0.15 .26 .43 .23 .35 .13 .35	-1.30 -1.18 78 -1.12 -1.02 -1.19 95	8.3 9.3 0 -5.1 2.0 8.2 18.8	

TABLE 5.- MEASURED AND CALCULATED LOW MACH NUMBER COMBAX FOR NEGATIVE ANGLE OF ATTACK

Reference 3.

Equation (A7).



Figure 1.- Calculated lift, drag, and pitching-moment coefficients for the NACA 23012 airfoil section, at M = 0.1.



(a) Section lift coefficient.

Figure 2.- Calculated and measured aerodynamic characteristics for NACA 23012 airfoil section.



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Figure 2.- Continued.



ODA MEASURED, REF.2





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(a) Section normal torce coefficient;

Sigure 3. Calculated and measured derodynamic characteristics for XACA (1912) airford section.



(b) Section drag coefficient.

Figure 3.- Continued.





Figure 3.- Concluded.



(a) Airfoil section V23010-1.58.

(b) Airfoil section VR-7.

Figure 4.- Maximum lift coefficient.









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Figure 7.- Airfoil section calculated mean pressure coefficient.





Figure 8.- Airfoil section perimeter/chord ratio.







Figure 10.- Correlation parameter for airfoil section drag due to lift.

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MDD = M FOR WHICH dCu/dM = 0.1, AT CONSTANT α

(a) NACA symmetrical four-digit series airfoll sections.

(b) NACA cambered four-digit series airfoil sections.

Figure 11.- Measured airfoil section drag divergence Mach number.





(c) NACA five-digit series airfoil sections.

(d) NACA 63A symmetrical series airfoil sections.

Figure 11.- Continued.

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(e) NACA 64A series airfoil sections.

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(f) NACA 65 series and Wortmann airfoil sections.

Figure 11.- Concluded.

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Figure 12.- Slope of c_d curve above drag divergence Mach number.



Figure 13.- Airtoil section drag coefficient through 180°.

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