## NASA Technical Paper 1348



# Sonic-Boom Minimization With Nose-Bluntness Relaxation 

Christine M. Darden



$\qquad$

## 



르…............-
$\qquad$
$\qquad$
O-


$\qquad$
$\qquad$

## --an



# Sonic-Boom Minimization With Nose-Bluntness Relaxation 

Christine M. Darden<br>Langley Research Center<br>Hampton, Virginia

National Aeronautics
and Space Administration
Scientific and Technical Information Office

## SUMMARY

A method which provides sonic-boom-minimizing equivalent area distributions for supersonic cruise conditions is described. This work extends previous analyses to permit relaxation of the extreme nose bluntness which is required by conventional low-boom shapes and includes propagation in a real atmosphere. The procedure provides area distributions which minimize either shock strength or overpressure. The minimizing Whitham F-functions for the pressure signatures are also included.

## INTRODUCTION

Sonic-boom levels produced by future generation supersonic transports must be considered if their supersonic flight paths, and thus their economic success, are not to be severely restricted. It is generally believed that the best approach to this problem is to include low-boom constraints early in the design process. The procedure described herein provides this constraint in the form of an equivalent area distribution which depends on the design conditions of aircraft weight (lift) and length, and on cruise Mach number and altitude.

Sonic-boom-minimization studies, by Jones, Seebass, and George for uniform and isothermal atmospheres have been reported in references 1 to 6. These analyses have been extended to a real atmosphere in reference 7. All these studies have indicated that the minimizing area distribution is characterized by a bluntness in the leading portion of the area distribution, i.e., the aircraft nose region. Because extreme nose bluntness produces large drag, a method of relaxing the bluntness requirement is needed to offer the opportunity for compromise between blunt-nose, low-boom and sharp-nose, low-drag configurations.

This analysis extends the work of Seebass and George (ref. 6) and Darden (ref. 7) to include nose-bluntness relaxation. Because questions still remain regarding the type of pressure signature and the level of overpressure which are acceptable, either a minimum-shock or a minimum-overpressure signature may be produced for given cruise conditions and nose length.

This analysis has been incorporated into a computer program which is described in the appendix of this paper.

SYMBOLS

Although values are given in both SI and U.S. Customary Units, the calculations for the investigation were made in U.S. Customary Units.

Notation used in the computer printout is given in parentheses.


| X, Y | axial distances |
| :---: | :---: |
| $Y_{f}(\mathrm{YF})$ | width of forward spike in F-function |
| $Y_{r}(T)$ | position of rear area balancing |
| $z(r)$ | vertical distance of signal from airplane axis |
| $\alpha$ | advance of acoustic rays |
| $B$ | $=\left(M^{2}-1\right)^{1 / 2}$ |
| Bz (BR) | shifting factor for axial position |
| $\gamma$ | ratio of specific heats, 1.4 for air |
| $\Gamma$ | $=\frac{Y+1}{2}$ |
| $\delta$ | Dirac delta function |
| $\eta$ | fraction of balancing line slope used in defining F-function |
| $\lambda$ (LAM) | axial position at which F-function becomes negative |
| $\mu$ | Mach angle, $\sin ^{-1} \frac{1}{M}$ |
| $\xi$ | dummy integration variable |
| $\rho$ | density |
| Subscripts: |  |
| a | ambient |
| f | front shock |
| 9 | ground |
| h | altitude of initial waveform |
| 0 | minimum-overpressure signature |
| r | rear shock |
| S | minimum-shock signature |
| Y | axial position |

vertical distance of signal from airplane axis
First and second derivatives with respect to distance are denoted by single and double primes, respectively.

## BACKGROUND

Sonic-boom-minimization studies are based primarily on currently accepted sonic-boom prediction methods. These methods, outlined in figure 1, are extensions of work done by Whitham (ref. 8), Walkden (ref. 9), and Hayes (ref. 10). From the complex aircraft configuration and its lift distribution, an equivalent area distribution is defined. This equivalent area distribution then defines the Whitham "F-function" through the following integral relation:

$$
\begin{equation*}
F(y)=\frac{1}{2 \pi} \int_{0}^{y} \frac{A_{e} e^{\prime \prime}}{(y-\xi)^{1 / 2}} d \xi \tag{1}
\end{equation*}
$$

This F -function represents a distribution of sources which causes the same disturbances as the aircraft at some distance from the aircraft.

Because the pressure signal propagates at the local speed of sound, and each point of the signal advances according to its amplitude, the signal is distorted at the ground and, theoretically, can be multivalued. The physically unrealistic multiple values of pressure are, however, eliminated by the introduction of shocks. Shock location, based on the observation that for weak disturbances the shock bisects the angle between two merging characteristic lines, is determined by a balancing of the signature areas within loops on either side of the shocks. This procedure is demonstrated by the shaded areas of the distorted signal in figure 1.

For present-day supersonic aircraft, the propagation distance (altitude) and the pattern of shock coalescence have been such that, at ground level, only two shocks remain in the signature. These two shocks have a linear pressure variation between them (as seen in fig. 2), thus the name "far-field N-wave." For this wave, the shape of the generating aircraft has an effect on the magnitude but no effect on the shape of the resulting signature. For a sufficiently long and slender aircraft, it was found that the ground wave form may possibly not have attained its N -shape remaining instead a mid-field wave. Because the shape of the aircraft does influence the shape of a mid-field pressure signature, aircraft shaping has now become a powerful tool in reducing the sonic boom.

The F-functions for the lower bound of an N-wave (refs. 1 and 2) and for the lower bound of the bow shock in a mid-field signature (refs. 3 and 4) led to the form of the minimizing F-function assumed by Seebass and George for the entire signature. Lung (ref. 11) later proved this to be the minimizing form by "bang-bang" control theory.

Common to all of these minimizing forms of the F-function is a Dirac delta function at $y=0$. This infinite impulse corresponds to an infinite gradient at the nose of the equivalent area distribution. Aircraft designed to match these area distributions generally have extremely blunt nose shapes which result in substantial drag. This result, though seemingly paradoxical, can be explained by the created shock-attenuation pattern in which the shock strengths, and therefore the drag, are greatest near the aircraft. Because of special shaping and area growth, secondary shocks from other aircraft components do not overtake and enhance the bow shock during the propagation of the wave front. (See fig. 3.) The net result of this process is weaker shocks at larger distances because of attenuation, but an overall increase in drag. In contrast, notice that the bow shock is weaker on the sharp-nosed aircraft, but coalescence with stronger shocks from rearward components causes a much stronger shock at mid- and far-field distances. If an aircraft configuration is to incorporate boom-minimization features and not suffer aerodynamic penalties which inordinately affect its flight efficiency, range, payload, productivity, and so forth, then means must be provided for including nose-bluntness sensitivity in design trade-off studies.

## ANALYSIS

The following is a brief description of the Seebass-George minimization scheme with modifications to provide propagation through the real atmosphere (ref. 7) and control over the bluntness of the area distribution and thereby the drag of the configuration. The assumed form for the class of minimizing F-functions is shown in figure 4. In mathematical terms it may be expressed as:

$$
\begin{align*}
& F(y)=2 y H / Y_{f}  \tag{2a}\\
& F(y)=C\left(2 y / Y_{f}-1\right)-H\left(2 y / y_{f}-2\right)  \tag{2b}\\
& F(y)=B\left(y-Y_{f}\right)+C  \tag{2c}\\
& F(y)=B\left(y-y_{f}\right)-D \tag{2d}
\end{align*}
$$

$$
\begin{array}{r}
\left(0 \leqq y \leqq Y_{f} / 2\right) \\
\left(y_{f} / 2 \leqq y<y_{f}\right) \\
\left(y_{f} \leqq y<\lambda\right) \\
(\lambda \leqq y<l)
\end{array}
$$

In these equations $H, B, C, D$, and $\lambda$ are unknown coefficients which are determined by the cruise conditions of the aircraft, by nose length, by the prescribed ratio of bow to rear shock, and by the signature parameter being minimized. Types of signatures studied include flat-topped signatures in which overpressure is minimized with $B=0$ and signatures in which $F(y)$ is allowed to rise between $y_{f}$ and $\lambda$ with a resulting minimum shock followed by a pressure rise as illustrated in figure 4. The value of $B$ in this form of $F(y)$ may range between 0 and $S$. Recall that the Whitham function $F(y)$ represents the shape characteristics of the pressure signature and is defined in reference 10 and equation (1) in terms of the equivalent area distribution as:

$$
F(y)=\frac{1}{2 \pi} \int_{0}^{y} \frac{A_{e} "}{(y-\xi)^{1 / 2}} d \xi
$$

Crucial to this minimization technique is the fact that equation (1) is an Abel integral equation which may be inverted to give the function $A_{e}$ in terms of the $F$-function. When this function is evaluated at $l$, the result is:

$$
\begin{equation*}
A_{e}(l)=4 \int_{0}^{l} F(y)(l-y)^{1 / 2} d y \tag{3}
\end{equation*}
$$

Upon substituting the minimizing form of the $F$-function into equation (3) and integrating, the following equation for the development of cross-sectional area is obtained:

$$
\begin{align*}
A_{e}(x)= & \frac{32}{15} \frac{H}{Y_{f}} x^{5 / 2}+1\left(x-\frac{Y_{f}}{2}\right)_{15}^{8}\left(x-\frac{Y_{f}}{2}\right)^{3 / 2}\left[\left(\frac{3 y_{f}}{2}+2 x\right)\left(\frac{1}{Y_{f}}\right)(2 C-4 H)\right. \\
& +5(2 H-C)]+1\left(x-y_{f}\right) 4\left(x-Y_{f}\right)^{3 / 2}\left[\left(\frac{2 C}{y_{f}}\right)\left(-\frac{2}{15}\right)\left(3 y_{f}+2 x\right)+\frac{2}{3} C\right. \\
& \left.+\frac{4}{15}\left(\frac{H}{y_{f}}\right)\left(3 y_{f}+2 x\right)-\frac{4}{3} H+\frac{2}{15} B\left(3 y_{f}+2 x\right)-\frac{2}{3} B Y_{f}+\frac{2}{3} C\right] \\
& -1(x-\lambda) \frac{8}{3}(x-\lambda)^{3 / 2}(C+D) \tag{4}
\end{align*}
$$

where $1(x-2)$ is the Heaviside unit step function. A typical form of the resulting area distribution, $A_{e}$, is seen in figure 4.

If the effects of aircraft wake and engine exhaust are neglected, and, if the aircraft cross-section area is zero at its base, then the area at $l$ is entirely due to cruise lift or

$$
\begin{align*}
A_{e}(l)= & \frac{W}{\rho u^{2}}=\frac{32}{15} \frac{H}{Y_{f}} l 5 / 2+\frac{8}{15}\left(l-\frac{Y_{f}}{2}\right)^{3 / 2}\left[\left(\frac{3 y_{f}}{2}\right)+2 l\left(\frac{1}{y_{f}}\right)(2 C-4 H)+5(2 H-C)\right] \\
& +4\left(l-Y_{f}\right)^{3 / 2}\left[\left(\frac{2 C}{Y_{f}}\right)\left(-\frac{2}{15}\right)\left(3 y_{f}+2 l\right)+\frac{2}{3} C+\frac{4}{15}\left(\frac{H}{Y_{f}}\right)\left(3 y_{f}+2 l\right)-\frac{4}{3} H\right. \\
& \left.+\frac{2}{15} B\left(3 y_{f}+2 l\right)-\frac{2}{3} B y_{f}+\frac{2}{3} C\right]-\frac{8}{3}(l-\lambda)^{3 / 2}(C+D) \tag{5}
\end{align*}
$$

The first constraint imposed upon $F(y)$ is that the front area balance must occur at $y=Y_{f}$, where $Y_{f}$ is the first point at which $F(y)=C$, that is,

$$
\begin{equation*}
\int_{0}^{Y_{f}} F(y) d y=G=\frac{\alpha_{Y f}}{2} F\left(y_{f}\right)=\frac{\alpha_{Y f}}{2} c \tag{6}
\end{equation*}
$$

For the real atmosphere, the advance (ref. 9) of any point of the signal is given by

$$
\begin{equation*}
\alpha_{y}=\frac{\Gamma M_{h}^{3} F(y)}{(2 \beta)^{1 / 2}} \int_{0}^{z} \frac{p_{h}}{p}\left(\frac{\rho a_{h}}{\rho_{h} a}\right)^{1 / 2}\left(\frac{A_{h}}{z_{h} A}\right)^{1 / 2} \frac{M}{\beta} d z \tag{7}
\end{equation*}
$$

where $z$ is the vertical distance of the signal from the aircraft axis, $z_{h}$ is the vertical distance of the initial waveform from the aircraft axis, $A$ and $A_{h}$ are ray-tube areas determined from

$$
\begin{equation*}
\frac{A_{h}}{z_{h} A}=\left[M_{h}\left(1-\frac{1}{M_{z}^{2}}\right)^{1 / 2} \int_{0}^{z} \frac{d z}{\left(M_{z}^{2}-1\right)^{1 / 2}}\right]^{-1} \tag{8}
\end{equation*}
$$

and $\Gamma=(\gamma+1) / 2$. The initial waveform must be defined away from the aircraft, because the $F$-function only can represent the body shape accurately at several body lengths away, and the acoustical theory used in describing the propagation of the signal fails near the aircraft. The slope of the balancing line is proportional to the reciprocal of the advance at any point of the signal; thus,

$$
\begin{equation*}
S=\frac{(2 \beta)^{1 / 2}}{\Gamma M_{h}^{3} \int_{0}^{h h} \frac{p_{h}}{p}\left(\frac{\rho_{h}}{\rho_{h} a}\right)^{1 / 2}\left(\frac{A_{h}}{z_{h} A}\right)^{1 / 2} \frac{M}{\beta} d z}=\frac{F(y)}{\alpha_{y}} \tag{9}
\end{equation*}
$$

For the rear area balancing to occur between points $l$ and $Y_{r}$, then

$$
\begin{equation*}
\int_{l}^{Y_{r}} F(y) d y=\frac{1}{2}\left[B\left(l-Y_{f}\right)-D+F\left(y_{r}\right)\right]\left(y_{r}-l\right) \tag{10}
\end{equation*}
$$

where $Y_{r}$ is the unknown second intersection point of the rear area balancing line with $F(y)$. If a cylindrical wake is assumed, $F(y)$ and its integral for
$y>l$ can be expressed in terms of $F(y)$ for $y<l$ according to reference 1 as follows:

$$
\begin{align*}
& F(y)=-\frac{1}{\pi(y-l)^{l / 2}} \int_{0}^{l} \frac{(l-\xi)^{1 / 2}}{y-\xi} F(\xi) d \xi  \tag{11}\\
& \int_{l}^{y_{r}} F(y) d y=-\frac{2}{\pi} \int_{0}^{l} F(\xi) \tan ^{-1}\left(\frac{y_{r}-l}{l-\xi}\right)^{1 / 2} d \xi \quad(y>l) \tag{12}
\end{align*}
$$

It is necessary to define $F(y)$ and its integral for $y>l$ in this way for optimization problems, since aircraft geometry and thus the Whitham function can be varied arbitrarily only in the range $0 \leqq y \leqq l$. The constraint on the ratio of shocks is given by

$$
\begin{equation*}
\frac{P_{f}}{P_{r}}=\frac{C}{D-B\left(l-Y_{f}\right)+F\left(Y_{r}\right)} \tag{13}
\end{equation*}
$$

To insure that $Y_{r}$ is an intersection point of $F(Y)$ and the balancing line, then

$$
\begin{equation*}
F\left(y_{\Gamma}\right)=S\left(y_{Y}-l\right)+B\left(l-y_{f}\right)-D \tag{14}
\end{equation*}
$$

and the slope of $F(y)$ at $Y_{r}$ must be less than $S$.
Solving the system of equations ((5), (6), (10), (13), and (14)) provides values for the constants $H, C, D, \lambda$, and $Y_{r}$. One difficulty of this minimization approach is that, as yet, there is no precise definition of the rise time that will allow the ear to detect only the bow shock and not to be sensitive to the peak as well. With the values of the coefficients known, the minimizing F-function and the area distributions may be determined by equations (2) and (4).

To convert $F(y)$ into pressure near the airplane, the following equation is used:

$$
\begin{equation*}
\left(\frac{P}{P_{a}}\right)_{h}=\frac{\gamma M^{2} F}{\left(2 \beta z_{h}\right)^{1 / 2}} \tag{15}
\end{equation*}
$$

and on the ground

$$
\begin{equation*}
P_{g}=\left(\frac{A_{h}}{A_{g}}\right)^{1 / 2}\left(\frac{\rho_{g} a_{g}}{\rho_{h} a_{h}}\right)^{1 / 2} P_{h} \tag{16}
\end{equation*}
$$

Finally, by using a ground reflection factor $K$, the pressure perturbations are converted into the ground signature by

$$
\begin{equation*}
\Delta \mathrm{p}_{\mathrm{g}}=\mathrm{P}_{\mathrm{g}} \mathrm{~K} \tag{17}
\end{equation*}
$$

The ground-level signature is presented by the signature shown in figure 4.
By taking the limit as $y_{f} \rightarrow 0$ and by using L'Hospital's rule, the expression for area development given above (eq. (4)) reduces to the following delta function form as given by Seebass and George (ref. 6):

$$
\begin{equation*}
A_{e}(x)=4 G x^{1 / 2}+\frac{16}{15} B x^{5 / 2}+\frac{8}{3} C x^{3 / 2}-1(x-\lambda) \frac{8}{3}(x-\lambda)^{3 / 2}(C+D) \tag{18}
\end{equation*}
$$

SPECIAL CASES
Special types of pressure signatures occur for large or small values of the weight parameter $\tilde{w}$, where $\tilde{w}=\frac{\beta W}{\rho u^{2}} \frac{1}{B\left(l-y_{f}\right)^{5 / 2}}$. These special cases include signatures with no bow shock, signatures with no shocks, and signatures in which the rear shock is less than the bow shock without restriction (fig. 5). Values of $\tilde{w}$ may be checked to indicate when the special calculations are necessary.

No Bow Shock
For no bow shock, the expression for total area reduces to the $y_{f}=0$ form (eq. (18)) automatically. Using the facts that $G=\alpha C^{2}$ and $\lambda=l$ at the minimum length necessary for no shock, the resulting quadratic in $C$ may be solved explicitly to give

$$
\begin{equation*}
C=\frac{B l}{\eta}\left[-\frac{2}{3} \pm\left(\frac{4}{9}-\frac{8 \eta}{15}+\frac{\eta}{2} \frac{v}{B l^{5 / 2}}\right)\right]^{1 / 2} \tag{19}
\end{equation*}
$$

where

$$
\begin{equation*}
v=\frac{\beta W}{\rho u^{2}} \quad B=S \eta \quad s=\frac{1}{\alpha} \quad(0 \leqq \eta \leqq 1) \tag{20}
\end{equation*}
$$

For no bow shock ( $C=0$ ) equation (19) reduces to

$$
\begin{equation*}
\tilde{w}=\frac{V}{\mathrm{~B} l^{5 / 2}} \leqq \frac{16}{15} \tag{21}
\end{equation*}
$$

in the minimum shock signature.

No Shocks
To eliminate both shock waves, $G$ and $C$ are zero and there must be no discontinuity in $F$ at $l$. In fact, both $F(y)$ and $F^{\prime}(y)$ must be continuous at $l$ and thus $A_{e}{ }^{\prime}(x)$ and $A_{e}{ }^{\prime \prime}(x)$ must be zero there, since $A_{e}$ is constant for $x>2$. With no shocks,

$$
\begin{equation*}
A_{e}(l)=\frac{16}{15} B q^{5 / 2}-\frac{8}{3} D(l-\lambda)^{3 / 2} \tag{22}
\end{equation*}
$$

By taking the first and second derivatives and solving the two resulting equations simultaneously

$$
\begin{align*}
& \lambda=\frac{2}{3} l  \tag{23}\\
& D=\frac{2 \eta}{3 \alpha} \tag{24}
\end{align*}
$$

Substituting these values back into the expression for $A_{e}$ gives

$$
\begin{equation*}
\tilde{\mathrm{w}}=\frac{\mathrm{V}}{\mathrm{~B} l^{5 / 2}} \leqq 0.47407 \tag{25}
\end{equation*}
$$

## No Restriction on Rear Shock

Values of $\lambda<l$ in the minimizing $F$-function result because of restrictions on the rear shock. For short lengths, high weights, etc., or other combinations which give large values of $\tilde{w}$, the rear shock is less than the bow shock
without restriction. In these instances $\lambda=1$ and the system of equations is solved without restricting the ratio of rear to front shock:

$$
\begin{array}{ll}
F(y)=\frac{2 y h}{y_{f}} & \left(0 \leqq y \leqq y_{f} / 2\right) \\
F(y)=C\left(\frac{2 y}{y_{f}}-1\right)-h\left(\frac{2 y}{y_{f}}-2\right) & \left(y_{f} / 2<y<y_{f}\right) \\
F(y)=B\left(y-y_{f}\right)+C & \left(y_{f} / 2 \leqq y \leqq \eta\right) \tag{26c}
\end{array}
$$

Substituting this expression for $F$ into equation (3) and integrating gives a quadratic equation in $C$ which can be solved explicitly. Values for $C$, $h$, $\lambda$, $l$, and $B$ are now established. For the rear area balance to occur at $y_{r}$

$$
\begin{equation*}
\int_{l}^{t} F(y) d y=\frac{1}{2}\left[B\left(l-y_{f}\right)-D+F\left(y_{r}\right)\right]\left(y_{r}-l\right) \tag{27}
\end{equation*}
$$

also, $Y_{r}$ must be an intersection point for $F(y)$ and balance point, or

$$
\begin{equation*}
F\left(y_{r}\right)=s\left(y_{r}-l\right)+B\left(l-y_{f}\right)-D \tag{28}
\end{equation*}
$$

Combining these equations

$$
\begin{equation*}
Y_{r}-Z=\frac{\int_{l}^{y_{r}} F(y) d y}{\frac{1}{2}\left[2 F\left(y_{r}\right)-S\left(y_{r}-1\right)\right]} \tag{29}
\end{equation*}
$$

which may be solved iteratively for $Y_{r}$ with

$$
\begin{equation*}
F\left(y_{r}\right)=-\frac{1}{\pi\left(y_{r}-l\right)^{1 / 2}} \int_{0}^{l} \frac{(l-\xi)^{1 / 2}}{y_{r}-\xi} F(\xi) d \xi \tag{30}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{l}^{y_{r}} F(y) d y=-\frac{2}{\pi} \int_{0}^{l} F(\xi) \tan ^{-1}\left(\frac{y_{r}-l}{l-\xi}\right)^{1 / 2} d \xi \tag{31}
\end{equation*}
$$

Substituting back into equation (28) gives the required value for $D$.
Values of $\tilde{W}$ producing the special cases are shown in figure 5. Typical values giving these extreme $\tilde{w} ' s$ at $M=2.7, \quad z=18288 \mathrm{~m}(60000 \mathrm{ft})$, $W=272155 \mathrm{~kg}(600000 \mathrm{lb})$ would be $l=33.5 \mathrm{~m}(110 \mathrm{ft})$ for no restriction on the rear shock and $l=167.64 \mathrm{~m}(550 \mathrm{ft})$ for no bow shock. These values may vary slightly with $Y f$ as seen.

TYPICAL RESULTS

Except as indicated otherwise, all results shown are for the following conditions: Mach number, 2.7; altitude, $18288 \mathrm{~m}(60000 \mathrm{ft})$; weight, 272155 kg ( 600000 lb ) ; equivalent length, $91.44 \mathrm{~m}(300 \mathrm{ft})$; reflection factor, 2.0 ; ratio of bow to rear shock pressure, 1 ; and $B=0.5 S$. The value of the front balancing point Yf determines the length of the conical nose region of the equivalent area distribution and is the controlling factor for drag when all of the other flight conditions remain constant. Referring again to figure 4, this position is indicated on the $F$-function and the corresponding area distribution.

The effect that varying $y_{f}$ has on the equivalent area distributions for both the minimum-shock and minimum-overpressure signatures is seen in figure 6 . Each of the distributions was determined from the inversion formula which defines the effective area corresponding to the minimizing F-function for a given value of $y_{f}$. Values of the overpressure and impulse which correspond to these distributions have been inserted for reference. Recall that the impulse is the integral of the positive portion of the pressure signature, as indicated in the signatures at the top of figure 6. Note that the lowest values of overpressure and impulse for these conditions occur when the area distribution has an infinite gradient at the nose. As yf increases, the longer conical nose region, generated to reduce drag, also causes an increase in overpressure and impulse. Thus, extensive tradeoff studies between drag and boom levels would be necessary in any aircraft design studies. The ratio of the increase in overpressure and impulse for both types of signatures as a function of the ratio of nose length to airplane length $y_{f} / l$ is seen more easily in figure 7. With increasing values of $Y_{f}$, note that the area under the peak portion of the $F$-function also increases but not in the same proportions. Therefore, to achieve the necessary total area, the level of the constant portion of the F -function also must increase, thereby producing a higher level of $\Delta \mathrm{p}$ and I . For illustrative purposes, the nose spike of the F-function for $Y_{f}=0$ has been drawn with a finite width. For this condition, the spike is defined mathematically as a pulse of zero width and infinite height which nevertheless has a finite area. In these figures, $\Delta p$ is the level of overpressure at the bow shock and represents the initial "bang" heard by the ear when a sonic boom occurs.

The variations of overpressure level and the corresponding drag levels are seen in figure 8. To develop the corresponding variation of drag with $\mathrm{Y}_{\mathrm{f}}$, the assumption is made that necessary configuration changes would be confined to the fuselage forebody itself. The increment in drag produced by these changes was calculated using the near-field program of reference 12 and then applied to drag values for a typical arrow-wing SST configuration. As seen in figure 8, there is a sharp decline in drag as the blunt nose is converted to a conical region. Thus, the boom levels could be reduced significantly without prohibitive drag penalties by defining the proper ratio $y_{f} / 2$.

The actual trend of overpressure with length is seen in figure 9. The lowest curve represents the $y_{f}=0$ case, the higher ones increasing with $y_{f}$. In all instances observe that overpressure levels decrease with length. Though these results are for minimum-overpressure signatures, similar trends are found to exist for the minimum-shock signatures.

## CONCLUDING REMARKS

The analysis for a sonic-boom minimization method has been presented. The method, which represents an extension of previous work, provides the minimizing equivalent area distribution with relaxation of the nose-bluntness requirement for given cruise conditions and includes propagation in the standard atmosphere. The minimizing Whitham $F$-functions for the minimum-overpressure or minimum-shock signature are also included.

Langley Research Center
National Aeronautics and Space Administration
Hampton, VA 23665
November 13, 1978

## APPENDIX

## COMPUTER PROGRAM DESCRIPTION AND FLOW CHART

This program calculates the Whitham $F$-function and the corresponding equivalent area distribution which produce the minimum-overpressure or minimum-shock signature. The solution is based on five governing equations: total area growth (eq. (5)), front area balancing point (eq. (6)), rear area balancing point (eq. (10)), ratio of front to rear shock (eq. (13)), and the area balancing line at the rear balancing point (eq. (14)). These five equations reduce to two nonlinear equations in two unknowns which are solved iteratively using the NewtonRaphson and a combination of the secant and bisection methods.

In calculating the advance of the wave front the standard atmosphere has been used except for an area near the aircraft axis where it is necessary to assume a uniform atmosphere.

This program was coded for the Control Data 6600 computer in Fortran IV. The field length required is 100000 octal units. Program SEEB is available as LAR 11979 from

COSMIC
Suite 112 Barrow Hall
University of Georgia
Athens, GA 30602

The flow chart of program SEEB is on the following page.

## INPUT DESCRIPTION

| M | Cruise Mach number (1.25 to 3.0) |
| :---: | :---: |
| Z | Cruise altitude ( 20000 to $100000 \mathrm{ft} \mathrm{(6096} \mathrm{to} 30480 \mathrm{~m}$ ) ) |
| L | Airplane equivalent length ( 200 to $700 \mathrm{ft}(60.96$ to 213.36 m$)$ ) |
| WG | ```Gross weight of airplane (200 000 to 1 000 000 1b (90 720 to 453 600 kg))``` |
| YF | Balancing point in $F$-function to determine the front shock ( $0 \leqq y_{f} / l \leqq 0.2$ ) (Default $=0.0$ delta function) |
| IPRINT | $=0$ Iterations on Lambda and $T$ not printed. |
|  | $\begin{aligned} =1 & \text { Iterations printed. } \\ & \text { (Default }=0) \end{aligned}$ |
| RKF | Reflection factor. <br> (Default $=2.0$.) |



```
BLEN Factor determining where iteration on T begins. T = BLEN•L
    (Default = 3.0)
PRPF Ratio of rear to front shock
    (Default = 1.0)
WIWG
ISIG
IS = 0 Minimum-shock and minimum-overpressure signatures.
= 1 Minimum-overpressure signature only.
=2 Minimum-shock signature only.
    (Default = 0)
```

PERCEN

ICH

IYF $\quad=0$ Delta function used in F-function.
$=1$ Nonzero value of $Y_{f}$ used.
If the value of $y_{f}$ is less than 1 , the program automatically uses the delta function in the $F$-function.
(Default $=1$ )

## APPENDIX

OUTPUT DESCRIPTION
The output consists of a listing of input conditions, the solution coefficients, solution tables, and summary values. Codes exist so that only the pressure signature is printed, the $F$-function, area, and pressures are printed, or extensive printout occurs during iterations.

The axial position of the front balancing point is Yf.

The advance factor on the ground is $\alpha$.

Altitude is given in feet and kilometers.
Gross weight is given in pounds.
Cruise weight is given in pounds and kilograms.
WIWG $\quad=$ Ratio of beginning cruise weight to gross weight
$W$-TILDE $\quad=\tilde{w}=\frac{\beta W}{\rho u^{2}} \frac{\alpha}{\eta\left(2-Y_{f}\right)^{5 / 2}}$

Total Area $=\frac{\beta W}{\rho u^{2}}$
ISLOPE $\quad=$ Code for signature type
$0=$ Flat-top minimum overpressure
1 = Minimum shock

H, B, C, D, LAM, and Yf are coefficients in equations for F-function and area distribution. $\quad\left(Y_{r}=T ; \quad \lambda=L A M\right)$
$X \quad=$ Axial position for $F$-function and area distribution near aircraft.
$X-B R \quad=x-\beta z=$ Shifted axial position for pressure signature on the ground.

DELTA $P \quad=\Delta p=$ Overpressure levels corresponding to $x-\beta z$.
$P_{1}, P_{2} \quad=$ Overpressure levels in pressure signature. $P_{1}$ and $P_{2}$ at same $X$ - $B R$ indicate a shock.

## APPENDIX

## SAMPLE CASE



222222222222222222222222222222222222222222222222222222222222222222222222222

4444444444444444444444444444444444444444444444444444444444444444444444444
555 5555555555555555555555555555555555555555555555555555555555 55 55555555




12:
[I4M508]

For this sample case an "IS" value of zero was used to produce both a minimum-overpressure solution and a minimum-shock solution.
APPENDIX

$$
\begin{gathered}
\text { OUTPUT } \\
\text { MINIMUM-OVERPRESSURE SOLUTION }
\end{gathered}
$$

POSITIVE VALUE OF YF USED IN F-FUNCTION




 0

[^0]





[^1]







[^2]
## APPENDIX


DELTA P
$-.336928 \varepsilon-01$
$-.321526 \varepsilon-01$


80
$N N$

## APPENDIX







 $\underset{\sim}{\sim}$

## APPENDIX






89 $27741 t+00$
+03
+03
HE +CI

APPENDIX
MINIMUM-SHOCK STRENGTH SOLUTION
shock minimization in the real atmosphere after segbass and geqrge

a






$$
i
$$

$$
0
$$













## APPENDIX


UELTA P
$-.330955 E-01$
$-.315544 E-01$
$x-B R$
$144135 \mathrm{E}+04$
$.148195 \mathrm{E}+04$
$\underset{\sim}{\underset{4}{山}} \underset{\sim}{\sim}$
$u$
$-.171321 E-02$
$-.1633445-02$
$\begin{array}{ll}200 & .143500 t+04 \\ 201 & .147589 t+04\end{array}$

## APPENDIX

[^3]







## APPENDIX








## APPENDIX



$.820959016+00$ (LAM-YFI/(L-YF)
LOCATION OF FRONT
LOCATION OF REAR SHD
TIME OETWEEN SHOCK
CHARACTERISTIG OVER

## BOW SHOCK= 47.67PA



CHARACTEKSSTIC OVEKPQ: $103.20 P A$ 105.72METERS N UF REAR SHOCK=

BOW SHOCK= 47.67PA REAR SHUCK= 47.67PA
OISTANCE FRQM BOW SHOCK TO PMAX= 110.83 FEET
TIME FROM BOW SHOCK TU PMAX: 04 SEC

## REFERENCES

1. Jones, L. B.: Lower Bounds for Sonic Bangs. J. Roy. Aeronaut. Soc., vol. 65, no. 606, June 1961, pp. 433-436.
2. Jones, L. B.: Lower Bounds for Sonic Bangs in the Far Field. Aeronaut. Q., vol. XVIII, pt. 1, Feb. 1967, pp. 1-21.
3. Seebass, R.: Minimum Sonic Boom Shock Strengths and Overpressures. Nature, vol. 221, no. 5181, Feb. 15, 1969, pp. 651-653.
4. George, A. R.: Lower Bounds for Sonic Booms in the Midfield. AIAA J., vol. 7, no. 8, Aug. 1969, pp. 1542-1545.
5. George, A. R.; and Seebass, R.: Sonic Boom Minimization Including Both Front and Rear Shocks. AIAA J., vol. 9, no. 10, Oct. 1971, pp. 2091-2903.
6. Seebass, R.; and George, A. R.: Sonic-Boom Minimization. J. Acoust. Soc. America, vol. 51, no. 2 (pt. 3), Feb. 1972, pp. 686-694.
7. Darden, Christine M.: Minimization of Sonic-Boom Parameters in Real and Isothermal Atmospheres. NASA TN D-7842, 1975.
8. Whitham, G. B.: The Flow Pattern of a Supersonic Projectile. Commun. Pure \& Appl. Math., vol. V, no. 3, Aug. 1952, pp. 301-348.
9. Walkden, F.: The Shock Pattern of a Wing-Body Combination, Far From the Flight Path. Aeronaut. Q., vol. IX, pt. 2, May 1958, pp. 164-194.
10. Hayes, Wallace D.: Linearized Supersonic Flow. Ph. D. Thesis, California Inst. Technol., 1947.
11. Lung, Joseph Lui: A Computer Program for the Design of Supersonic Aircraft To Minimize Their Sonic Boom. M.S. Thesis, Cornell Univ., 1975.
12. Marconi, Frank; Salas, Manuel; and Yaeger, Larry: Development of a Computer Code for Calculating the Steady Super/Hypersonic Inviscid Flow Around Real Configurations. Volume I - Computational Technique. NASA CR-2675, 1976.


Figure 1.- Sonic-boom prediction methods.


Figure 2.- Pressure-signature propagation.
Sharp nose

High boom
low drag
Blunt nose
Low boom
high drag
Figure 3.- The low-boom, high-drag paradox.


Figure 4.- Illustration of the theoretical concepts of near-field sonic boom.

Figure.5.- Special cases shown as function of weight factor.





Figure 8.- Estimated drag increments associated with configuration changes for sonic-boom minimization. $M=2.7 ; h=18288 \mathrm{~m}(60000 \mathrm{ft}) ; W=272155 \mathrm{~kg}(600000 \mathrm{lb}) ; \quad 2=91.44 \mathrm{~m}(300 \mathrm{ft})$.

-


[^4]
[^0]:    

[^1]:    
    

[^2]:    
    

[^3]:    
    
    

[^4]:    * For sale by the National Tectnical Information Service, Springfield. Virginia 22161

