Radiation Damage and Defect Behavior in Ion-Implanted, Lithium Counterdoped Silicon Solar Cells

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ABSTRACT

Boron-doped silicon n^+p solar cells were counterdoped with lithium by ion implantation and the resultant n^+p cells irradiated by 1 MeV electrons. Performance parameters were determined as a function of fluence and a Deep Level Transient Spectroscopy (DLTS) study was conducted in order to correlate defect behavior with cell performance. It was found that the lithium counterdoped cells exhibited significantly increased radiation resistance when compared to boron doped control cells. Isochronal annealing studies of cell performance indicate that significant annealing occurs at 100°C. Isochronal annealing of the deep level defects showed a correlation between a single defect at $E_v+.43$ eV and the annealing behavior of short circuit current in the counterdoped cells. It was concluded that the annealing behavior was controlled by dissociation and recombination of this defect. The DLTS studies also showed that counterdoping with lithium eliminated at least three deep level defects and resulted in three new defects. It was speculated that the increased radiation resistance of the counterdoped cells is due primarily to the interaction of lithium with oxygen, single vacancies and divacancies and that the lithium-oxygen interaction is the most effective in contributing to the increased radiation resistance.

INTRODUCTION

In the past, extensive studies have been conducted on p^+n silicon solar cells in which lithium was used as the n-dopant [1,2]. Although some advantage, under 1 Mev electron irradiation was found at elevated temperatures, in general, the cells emanating from this terminated program were found to exhibit the same tolerance to 1 MeV electron irradiation as conventional n+p silicon solar cells. We report here results on lithium counterdoped n^+p silicon solar cells in which lithium is introduced into the boron doped p-region in small enough quantities so that, despite the compensating effects of lithium, the cell base remains p-type. This procedure was

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followed in order to exploit the increased radiation resistance of p-type over n-type silicon [3]. The present work is an extension of our earlier work on electron-irradiated low resistivity, lithium counterdoped n^+p silicon cells [4]. It differs from our earlier work in the use of higher resistivity silicon and the use of ion implantation in introducting the lithium. We have, in addition, investigated the radiation induced deep level defects by Deep Level Transient Spectroscopy (DLTS) [5].

EXPERIMENTAL

The cells were fabricated from 1 ohm-cm boron-doped float zone silicon. All cells were 2 x 2 cm, 250 micrometers thick with no anti-reflection coating. The cell's n^+ region was formed by phosphorus ion implantation while the lithium was introduced by implantation of lithium ions. Electron beam annealing was used to selectively anneal the cell's n^+ region after ion-implantation [6]. Lithium concentrations were determined by four point resistivity measurements at the back surface and C-V measurements at the junction. For comparison purposes, n^+p control cells were fabricated by phosphorus ion implantation into the boron doped 1 ohm-cm float zone silicon. Cell characteristics are shown in Table I. After fabrication, the cells were irradiated by 1 MeV electrons and solar cell parameters determined using an AMO xenon-arc solar simulator. DLTS measurements were performed on small area portions of the cells (0.01 to 0.03 cm^2) using a 30 MHz capacitance bridge and boxcar averager [5]. To prevent extraneous annealing, the cells were immersed in liquid nitrogen between irradiations, measurements and isochronal anneals.

RESULTS

From Table I it is seen that the pre-irradiation power output levels are less for the lithium counterdoped cells. However, with irradiation, the effects of lithium are such that the output power of the counterdoped cells eventually exceeds that of the control cells (fig. 1). This can also be seen from Fig. 2 which shows normalized maximum power for all cells after completion of the irradiations.

The effects of isochronal annealing are shown in Fig. 3 where it is seen that significant cell recovery occurs for $T \leq 100°$ C. It is noted here that, at these temperatures, thermal degradation of space solar array components is minimal. Hence the present cells with further optimization could be candidates for components of an annealible array.

Lithium gradients are shown in Fig. 4. In all cases, the gradients are such that the lithium concentration is greatest at the rear of the cell. Previous data and calculations on lithium gradients for p$^+$n cells indicate a correlation between cell performance and gradient [7,8]. However this is not the case for the present cells.

A DLTS spectrum, showing both minority and majority carrier recombination centers is shown in Fig. 5 for the lithium counterdoped cells. From the figure, four deep level defects are detected with energy levels as shown. One result of lithium addition is the formation of new defects. This can be seen from Table II where the energy levels and capture cross sections of the counterdoped cells can be compared to those of the boron doped control cells. Isochronal annealing date for the defects in the counterdoped cells are shown in Fig. 6. Comparison with Fig. 3 shows a correlation between the isochronal annealing behavior of the defect at Ev + 0.43 eV and recovery of short circuit current.

DISCUSSION

Cell Performance After Irradiation

The significant improvement over the control cells, after irradiation, observed in the present case is much greater than the slight improvement previously noted in our earlier work on lithium counterdoped n$^+$p cells [4]. These latter cells were processed from 0.35 ohm-cm float zone and Czochralski grown silicon the results indicating that float zone was preferrable and that positive or zero lithium gradients were preferrable to negative lithium gradients. Our present use of float zone silicon and the avoidance of negative lithium gradients follows the recommendations emanating from our earlier results [4]. The present much greater improvement in radiation resistance could possibly be due to the use of higher resistivity silicon and/or the different processing technique used. With respect to processing, we recall that previously the lithium was applied as a paste to the cell's back surface followed by heating to drive in the lithium [4]. Concerning the absence of a definite correlation between gradient and performance, we note that the gradients found in the present n$^+$p cells are at least an order of magnitude lower than those cited for the p$^+$n cells where dependence of cell output on lithium gradient was previously reported [7,8].

Annealing

Comparison of Figs. 3 and 6 shows a strong correlation between the annealing behavior of the defect at Ev+0.43 eV and cell performance. A similar correlation is observed for diffusion length in Fig. 7. These data indicate that the Ev+0.43 defect is dominant in controlling the annealing behavior of the present counterdoped cells.

Interaction with Lithium

Little is known concerning the composition of the defects which correspond to the energy levels and capture cross sections of the deep level defects detected by DLTS in the counterdoped cell. However, much more is known concerning the defects in the boron doped control cell. In addition, there is a background of information, obtained by other techniques, on the composition of defects in lithium-doped silicon. For example, it is known that lithium combines with divacancies [9] oxygen [10] and substitutional boron [10] in silicon. This information can be used to speculate on the interactions involving lithium which lead to the increased radiation resistance we observe in the present counterdoped cells. Therefore in the remainder of this discussion we consider the defects in the boron doped cells and the possible interactions with lithium which could lead to the changes observed in the counterdoped cells.

Defect at E_c–0.27 eV

This defect has been tentatively identified as a complex of interstitial oxygen and interstitial boron [11]. A later investigation tends to conform the identification as a boron-oxygen complex [12]. Since the complex is positively charged [11], we assume no interaction with lithium which takes the form Li$^+$ in silicon. However lithium is known to combine with interstitial oxygen and substitutional boron in silicon [10]. Of these latter two interactions it is more likely that lithium in combining with interstitial oxygen would tend to inhibit formation of this defect.

Defect at Ev + 0.23 eV

This defect has been identified as the divacancy [11,13]. It has also been established that lithium forms complexes with divacancies [9]. Hence it is concluded that this defect is altered on counterdoping predominantly by the complexing of lithium with divacancies.

Defect at Ev + 0.33 eV

This defect has been alternately identified, from DLTS data, as a vacancy-oxygen-carbon complex [11] or as a carbon interstitial-carbon substitutional pair [14]. In both cases, the DLTS peak anneals out at $T \sim 400°$ C [11,14]. It is well to note that the DLTS data, by itself, does not suffice to identify the atomic constituents

of a defect. Other data, for example EPR, is usually required to identify a specific complex. In this connection, it is significant to note that the EPR spectrum, associated with the carbon-carbon pair, anneals out at $T \sim 300°$ C [15] while the EPR spectrum associated with the vacancy-oxygen-carbon defect anneals out at $T \sim 400°$ C [16]. Hence the EPR data favors the vacancy-oxygen-carbon identification. Since this defect is positively charged [16] it is unlikely that it would form complexes with lithium. Also there is no evidence that lithium interacts with carbon in silicon. Hence it appears likely that lithium interacts with oxygen [10] and vacancies [17] to alter the structure of this defect.

Summary of Defect Interactions With Lithium

From the preceding it is suggested that lithium interacts with oxygen, single vacancies and divacancies to alter the structures of the deep level defects seen in the present boron doped cells. The question as to which one of these interactions is most effective in contributing to the increased radiation resistance of the counterdoped cells is open to speculation. However, we note that of the three deep level defects affected by lithium in boron doped silicon, the boron-oxygen complex is the only one whose production rate increases as both cell resistivity and radiation resistance decrease [11]. This and the high minority carrier capture cross section of the $E_C-0.27$ defect in the boron doped cell suggest that the lithium-oxygen interaction is most effective in contributing to the increased radiation resistance observed in the present lithium counterdoped cells.

CONCLUSIONS

As a result of this work, it is concluded that;

-Lithium counterdoping results in significant increases in radiation resistance when compared to the 1 ohm-cm boron doped control cells.

-Performance of the counterdoped cells can be improved by annealing at $100°$C.

-The defect at $E_v + 0.43$ eV in the counterdoped cells is dominant in controlling the annealing behavior of the counterdoped cells.

-The increased radiation resistance of the counterdoped cells is primarly due to the interaction of lithium with oxygen, single vacancies and divacancies. It is suggested that the lithium-oxygen interaction is the most effective in contributing to the increased radiation resistance.

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TABLE I. – PRE-IRRADIATION CELL CHARACTERISTICS

Resistivity[a]	Li gradient, cm^{-4}	I_{sc}, mA	V_{oc}, mV	P_{max}, mW	FF, percent
1	---------	97.1	595	44	76.1
1.8	1.6×10^{17}	98.3	540	39.5	74.4
1.4	5.2×10^{16}	100.4	494	33	66.5
1.7	2.2×10^{17}	100.8	508	36.2	70.7
1.5	2.6×10^{16}	96.2	541	39.8	76.4
1.7	2.1×10^{17}	101.3	505	36.5	71.3
1.4	1.2×10^{17}	100.1	555	41.6	74.9

[a]Except for control: measured at back contact after introduction of lithium.

All cells 2 x 2 cm, 250 μm thick: no AR coating.

TABLE II. – ENERGY LEVELS AND CAPTURE CROSS SECTIONS

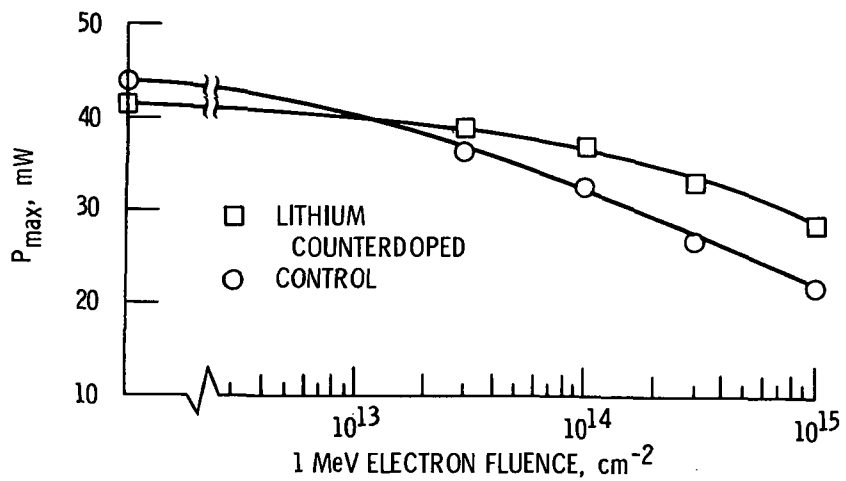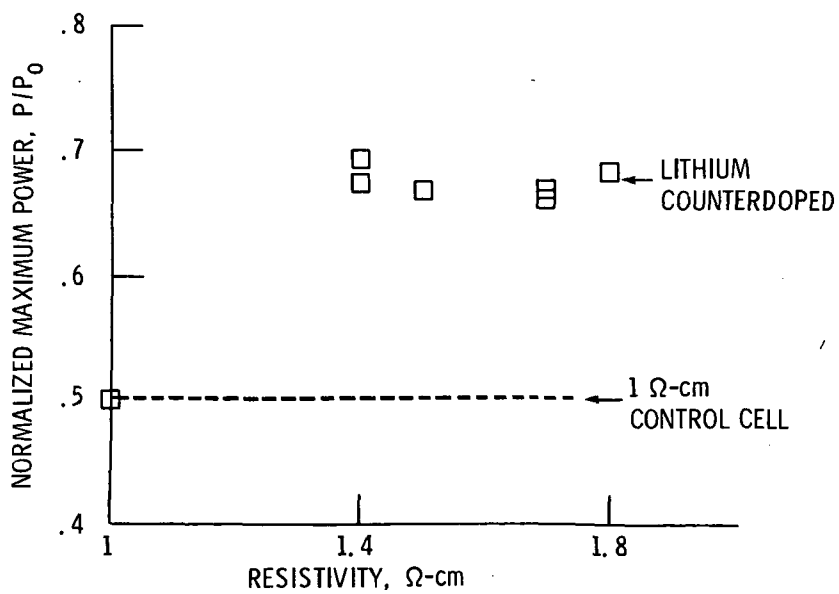|		1 OHM-cm boron doped				1.8 OHM-cm Li counterdoped			
Energy level, eV		$E_V+0.23$	$E_V+0.26$	$E_V+0.33$	$E_C-0.27$	$E_V+0.28$	$E_V+0.43$	$E_V+0.52$	$E_C-0.46$
Capture cross section, cm^2	σp	3×10^{-16}	4×10^{-17}	2×10^{-16}		8.5×10^{-16}	2×10^{-13}	1×10^{-14}	
	σN				3×10^{-13}				9.3×10^{-18}

Figure 1. - P_{max} versus 1 MeV electron fluence.

Figure 2. - Normalized maximum power for lithium
counterdoped and boron doped control cells;
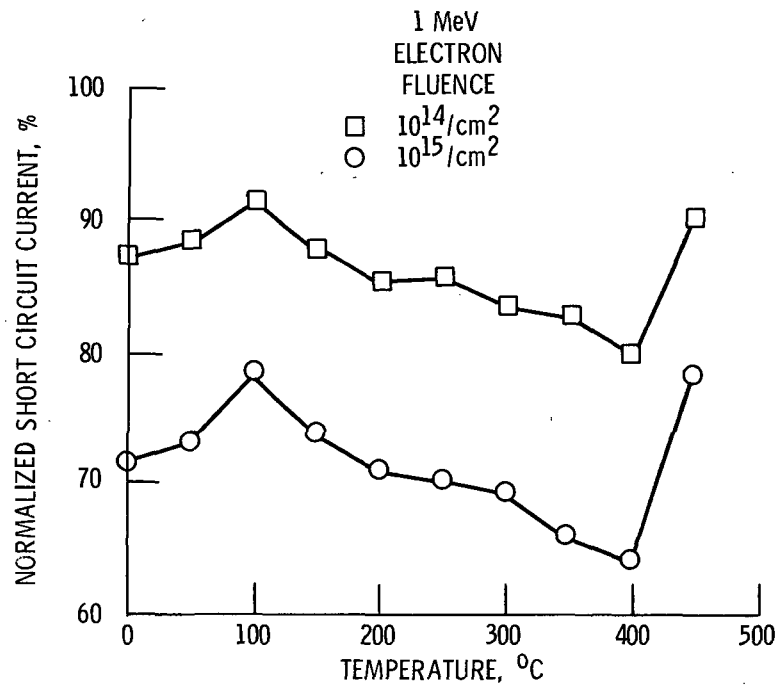$\Phi = 10^{15}/cm^2$.

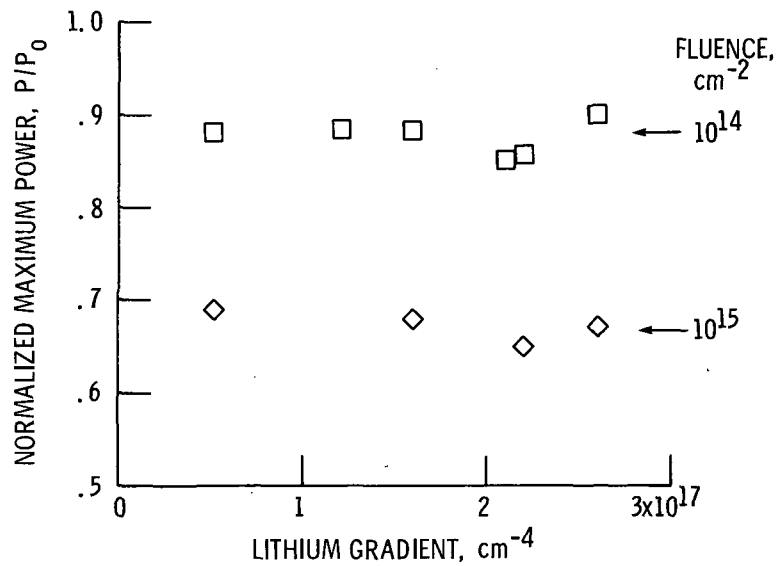Figure 3. - Isochronal anneal of lithium counterdoped silicon cells.

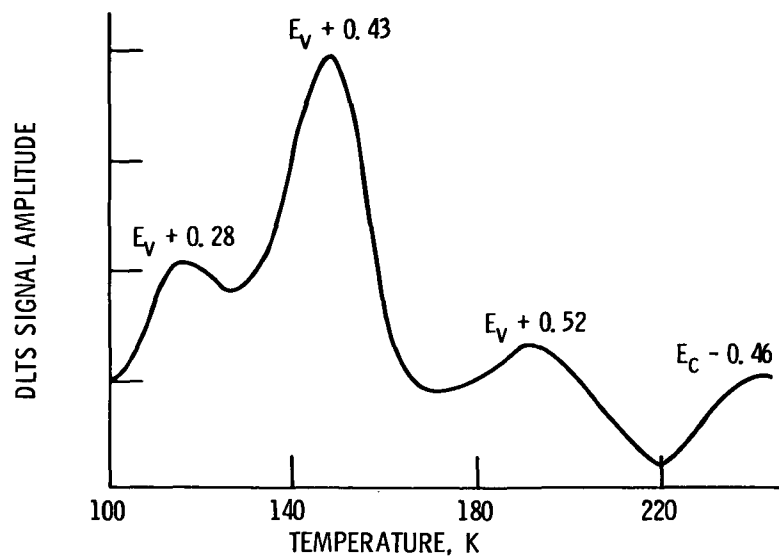Figure 4. - Lithium gradient versus normalized maximum power-lithium counterdoped cells.

Figure 5. - DLTS spectrum of lithium counterdoped silicon cells after 1 MeV electron irradiation.
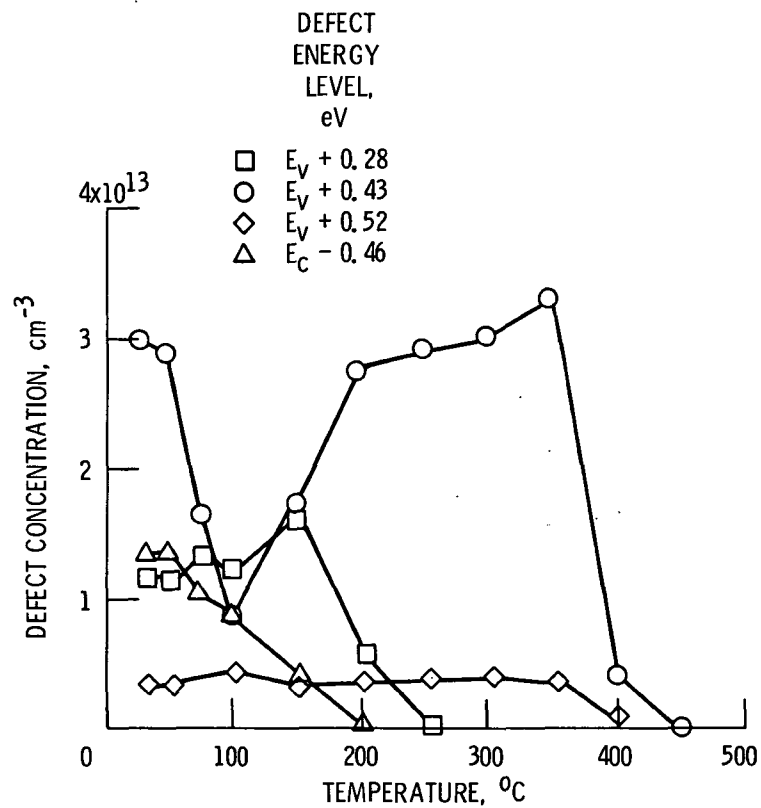

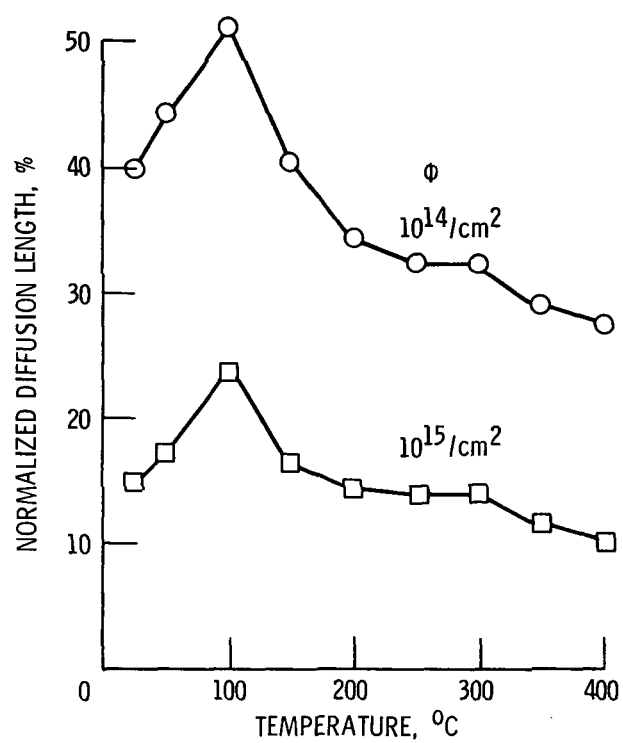Figure 6. - Isochronal anneal of defects in lithium counterdoped silicon solar cells using DLTS.

Figure 7. - Normalized diffusion length versus annealing temperature - lithium counter-doped cells.

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16. Abstract

Boron-doped silicon n^+p solar cells were counterdoped with lithium by ion implantation and the resultant n^+p cells irradiated by 1 MeV electrons. Performance parameters were determined as a function of fluence and a Deep Level Transient Spectroscopy (DLTS) study was conducted in order to correlate defect behavior with cell performance. It was found that the lithium counterdoped cells exhibited significantly increased radiation resistance when compared to boron doped control cells. Isochronal annealing studies of cell performance indicate that significant annealing occurs at $100^\circ C$. Isochronal annealing of the deep level defects showed a correlation between a single defect at $E_v+0.43$ eV and the annealing behavior of short circuit current in the counterdoped cells. It was concluded that the annealing behavior was controlled by dissociation and recombination of this defect. The DLTS studies also showed that counterdoping with lithium eliminated at least three deep level defects and resulted in three new defects. It was speculated that the increased radiation resistance of the counterdoped cells is due primarily to the interaction of lithium with oxygen, single vacancies and divacancies and that the lithium-oxygen interaction is the most effective in contributing to the increased radiation resistance.

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