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MANIFESTATION OF THE FERMI RESONANCE IN
SURFACE POLARITON SPECTRA

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Translation of "Proyavleniye razonansa fermi v spektrakh
povekhnostnykh poliaritonov". Ukrainsky Fizichyy Zhurnal,
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Abstract:
The method of disturbed full internal reflection (DFIR) is used to detect and interpret the resonance splitting of the surface polariton. The new effect in the spectra of oscillatory SP which are reflected by the DFIR method in Otto geometry was experimentally recorded. It is concluded that the resonance splitting of the dispersion branch of SP may serve as an effective method for detecting weak oscillations and for measuring their parameters.

MANIFESTATION OF THE FERMI RESONANCE IN SURFACE POLARITON SPECTRA


The method of disturbed full internal reflection (DFIR) was used to detect and interpret the resonance splitting of the surface polariton (SP) associated with the surface of single crystal α-LiI0₃ in the range of 790-810 cm⁻¹. The splitting is caused by the resonance interaction of SP with the weak volume polar oscillation of 801.5 cm⁻¹ in the contacting medium--a cholesteryl palmitate (CP) crystal.

We experimentally recorded the new effect in the spectra of oscillating SP which are reflected by the DFIR method in Otto geometry [1]. The effect consists of the emergence of a characteristic splitting in the dispersion branch of the SP, which is caused by resonance interaction of the SP genetically tied with one of the contacting media, with the weak polar oscillation of another dielectric medium. This effect is, to a certain degree, analogous to the polariton Fermi resonance--the resonance splitting of the disperse branches of volume polaritons (or anisotropic polar phonons) [2] interacting with weak polar oscillations of one or another medium. However, in our case, we are speaking of SP and the interaction of excitations associated with various contacting media due to the penetration of the SP field from one medium to another. This effect may be called the surface-polariton Fermi resonance. We are speaking, according to [3], of a quasiresonance interaction (in this case by means of an electromagnetic field) of the strong polar oscillation of one medium, caused by the SP, with the weak polar oscillation of another. We will note that work [4] examined the interaction of strong polar oscillations of various media by

*Numbers in margin indicate pagination of foreign text.
means of the SP field. However, this case differs from the case examined in this work.

![Figure 1](image)

The standard technology for the DFIR method was used for measuring the SP (for more detail see in [5]). Single crystalline hexagonal \(\alpha\)-LiI\(_3\) which was cut perpendicular to the optical axis was used as the dielectric in which the SP arose. The spectral interval of 750-850 cm\(^{-1}\) was studied, in which, as we know [5], there is an SP branch caused by intensive polar valent oscillations of the crystal with symmetry types A and E \((\omega^A = 795 \text{ cm}^{-1}, \omega^E = 769 \text{ cm}^{-1})\), with corresponding electrostatic splitting comprising 22 and 79 cm\(^{-1}\)). A CP liquid crystal was introduced into the gap between the \(\alpha\)-LiI\(_3\) crystal and the DFIR element (in this case, a half-cylinder made of KRS-5 crystal). The thickness of the CP crystal layer was approximately 1.5 \(\mu\)m. In the IR absorption spectrum for the volume CP, weak polar oscillation is registered without any great difficulty, with \(\omega_t = 801.5 \text{ cm}^{-1}\) (the force of the oscillator was \(s \sim 10^{-2}\), half-width of the corresponding IR-band was about 3 cm\(^{-1}\)).

The figure presents the obtained experimental points corresponding to the course of the dispersion branches of SP. The accuracy of determining \(\omega_p\) corresponding to the minimum reflection coefficient is indicated in the figure by vertical segments (spectral width of the spectrophotometer opening amounted to 1 cm\(^{-1}\)). The computed dispersion branches are also contained here: the broken line
shows the dispersion branch without consideration for oscillation \( \omega_t \), the solid lines—with consideration of the latter. Computa-
tions were performed on the study of SP in \( \alpha\)-LiIO\textsubscript{3} (polar axis of
crystal is perpendicular to its surface) in accordance with the
formula applicable for this geometry:

\[
k_{\parallel}^2 = -\frac{\omega_p^2}{c^2} \frac{\varepsilon_\infty (\varepsilon_{LL} - \varepsilon)}{\varepsilon_\infty^2},
\]

where \( \omega_p \) and \( k_{\parallel} \) are the frequency and wave vector of SP. \( \varepsilon_{LL} \) are
the main values of the tensor of dielectric permeability of \( \alpha\)-LiIO\textsubscript{3}
in a system of coordinates whose \( z \) axis is directed along the polar
axis of the crystal (their value is borrowed from [6], \( \varepsilon \) is the
dielectric permeability of the CP). The latter was approximated by
the standard expression (for isolated polar oscillation in an iso-
tropic medium):

\[
\varepsilon = \varepsilon_\infty \left[ 1 + (\omega_p^2 - \omega_1^2)/(\omega_p^2 - \omega^2) \right].
\]

while parameters \( \varepsilon_\infty \) and \( \omega_1 \) were selected from the condition of the
best correspondence with experimental data. It was achieved at \( \omega_1 = 803 \text{ cm}^{-1}, \varepsilon_\infty = 2.9 \), and in this case
\( \varepsilon = \varepsilon_\infty \approx 2\varepsilon_\infty (\omega_t - \omega) \omega_1 \approx 10^{-2}. \)
The physical sense of parameters \( \varepsilon = \varepsilon_\infty \) and \( \omega_1 \) is evident.

As we see, the interaction of SP with oscillation \( \omega_t \) causes a
rather large (in our case 2-3 cm\textsuperscript{-1}) characteristic splitting of the
disperse branch of the SP, and the experimental results corresponded
satisfactorily with the theoretical. We will note that the given
effect was theoretically predicted earlier in [7] using the example
of isotropic contacting media.

A double structure is observed in the DFIR spectrum, correspond-
ing to the two dispersion branches of the SP. In the range of
\( \omega_p(k_t) < \omega_t (\omega_p(k_t)) > \omega_t \) the low frequency (high frequency) component is
dominant, while near point \( \omega_p = \omega_t \) both components are commensurate
in intensity. Consequently, upon passage through the resonance, the
components seem to exchange intensities. A single band is observed
far from \( \omega_t \) in the SP spectrum.
The resonance splitting of the dispersion branch of SP may serve as a new effective method for detecting weak oscillations and for measuring their parameters. Here it is significant that it may be used for studying layers or films of the liquid or solid phase.

REFERENCES


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