RESULT OF MONTE-CARLO SIMULATION OF ELECTRON-PHOTON CASCADES IN LEAD AND LAYERS OF LEAD-SCINTILLATOR

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Results of Monte-Carlo simulation of electromagnetic cascade development in lead and lead-scintillator sandwiches are analysed. It is demonstrated, that the structure function for "core approximation" is not applicable in the case in which the primary energy is higher than 100 GeV. The simulation data has shown that introduced inhomogeneous structure of chamber gives subsequent reduction of secondary particles.

Introduction

In this paper we present results of detailed simulations , including the LPM effect, of electromagnetic cascades in lead and lead-scintillator sandwiches which are often used in emulsion chamber experiments. The following processes were taken into account in calculations: pair production, bremstrahlung including LPM effect with energy dependent cross - sections /Fig.1/, ionisations loss, Compton effect, single and multiple Coulomb scattering, the electron and positron inelastic scattering, photoelectric effect, annihilation of the positron.



Fig.1.a The product of differential cross-section of bremstrahlung and v, where $v=E/E_0$. b The differential cross-section of pair production, where $u=E/E_0$.

The simulations were performed:

- in lead, for primary photons of energy 100 GeV, 10 TeV and 100 TeV and for primary electrons of energy 100 GeV,
 in sandwiches of lead and scintillator: for primary elec-
- in sandwiches of lead and scintillator: for primary electrons of energy 100 GeV /Tab.1,No.1/, primary photons of energy 10 TeV and 100 TeV /Tab.1,No.2/.

absorber		Pb	Sc	Pb	Sc	Pb	Sc	Pb	Sc	Pb	Pb
thickness of aborb. /cm/	12	1 4	•5 2	•5 1	•5 2	•5 1	15 90	•5 3	•5 2	•5 1	•5 1

Tab.1 The structure of chambers investigated in our simulation. No.1: this sandwich is similar to that used in FNAL experiment/1/.No.2: this one is similar to the big gamma-hadron detector used in cosmic rays research /for instance Pamir Collaboration [2] /.

General procedure was following.All particles with energies exceeding certain value of $k=E_0/E_k$ were followed individually. Particles below E_k down to the threshold energy were followed using the so-called "thinning" method. The value of k was usually taken to be 1000 or 100 and the threshold energy was $2 m_e c^2$.

Longitudinal and lateral development of the electromagnetic caseades

In the emulsion chamber technique, the energy of an incoming electron or photon is measured by the behaviour of cascade showers near the shower axis. In "core approximation" is assumed that the lateral structure functions depend on E_0r/K_s and s only and not of separate variables E_0 and r." In this paper we analyse validity of this statement for wide range of energies and for different structure of chambers." In Fig.2 one can see the transition curves for primary electrons and photons of energy 100 GeV for shower radii 100 um together with the cascades obtained in chamber C1 of FNAL experiment."



Fig.2 Cascades curves of elec trons in photon and electron initiated showers for radius 100 um.

Our results are presented separately for cascade passing only throught lead and for cascades initiated in lead-scintillator / sandwich No. 1/. As can be seen the introduced scintillator plays important role, giving subsequent reduction of secondary particle numbers. The transition curve obtained for sandwiche's structure shows similar tendence as is seen in experiment, i.e. the smaller electron size than in pure lead chamber. We conclude that the agreement with experiment can be achieved by emploing as an absorber in our simulations the precise structure of FNAL sandwiches.

Let's compare the shape of transition curves when the energy of primary particles increases!

In Fig. 3a we present the transition curves with Ethr=2mec² $E_0 = 10$ TeV and $E_0 = 100$ TeV for pure lead chamber and for for sandwich No.2. Fig. 3b gives the transition curves for "core approximation" i.e. for shower radius 100 um and for the same primary energies of particles. For these cascades the LPM effect plays an important role in transition curves as the energy increases giving a deeper position of shower maximum and smaller size at maximum. In Fig.3 one can see that even primary photons of energy 10 TeV can penetrate through all chamber /sandwich No.2/. This effect is stronger when the primary energy of photon increases. The most of EAS experiments have one or more layers of carbon: this enhances the detector capability to discriminate hadrons from photons. and electrons. The insteresting results of the present simulations is that the number of particles after penetration of



the thick /90 cm/ scintillator and 3 cm lead increases if the lead follow this scintillator. Furthermore the numbers

Fig. 3 Cascade curves of electrons in photon initiated shower.

a. for all particles with threshold 2 mec^2 , b. the same threshold but in radius 100 um.

of electrons after first 4-5 cm in lead is comparable with ones after 90 cm of scintillator and 3 cm of lead for primary photons of energy 100 TeV. A detailed simulation would be necessary to quantify this effect with respect to different experiment.

The comparison of the simulation transition curves for different energies shows that validity of statement "core approximation" for all range of energies and depths is not good - Fig.4 .





Our results show that applied "core approximation" for lower energies than 100 GeV is rather good, but for higher energies LPM effect makes this relation wrong."

On the basis of presented transition curves we estimated the energy of primary particle for different depths of development cascades according to "core approximation". It is

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shown in Tab.2. The value of energy determinated from our data changes from a value lower than the energy of primary particle at low s to the value greater than ones at higher s. In the other words, our results show that the method of describing lateral distributions of electrons by "core approximation" leads to underestimation of the energy of the primary particle at low value of s and to overestimation at higher s.

	S	0.5	0.6	0.7	0.8	0,9	1.0	1.1	1.2	1.3	1.4
a	t /cm/	2.9	3.6	4.5	5.3	6.3	7.3	8.5	9.7	11.2	12.8
	E _e /TeV/	5.0	7.1	9.2	10.1	12.8	15.0	17.0	22.3	26.6	31.0
b	t/cm/	6.6	8.3	10.1	12.0	14.2	16.6	19.2	22.1	25.3	29.0
	E _e /TeV/	7.0	21.0	40.0	77.0	107.0	170.0	280.0	370.0	670.0	1000.

Tab.2 The comparison of the energy of primary photon estimated according to "core approximation". The energy of primary photon is : a.10 TeV and b.100 TeV.

As one can see, these discrepancies depend on the depth of the observation level and on the primary energy of particles.

<u>Conclusions</u>

1. It was stated that inhomogeneous structure of chamber has an effect on the transition curves. It appeared that numbers of secondary particles / for constant radii and constant primary photon energy / is determined by thickness, the kind of absorbers and their place in sandwiches.

2. It is not possible to describe the lateral distributions of particles with "core approximation" independent of the primary particle energy because the LPM effect plays an important role in transition curves. The method of describing lateral distributions of electrons by the "core approximation" used in chamber experiments lead to underestimation of the primary photon energy at low values of s and overestimation at higher s

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References

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