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Experimental Cavity Pressure Distributions at Supersonic Speeds

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Scientific and Technical Information Office

Summary

An experimental investigation has been conducted to define pressure distributions for rectangular cavities over a range of free-stream Mach numbers and cavity dimensions. These pressure distributions together with schlieren photographs were used to define the critical values of cavity length-to-depth ratio $(l/h)_{cr}$ that separate open type cavity flows from closed type cavity flows. For closed type cavity flow, the shear layer expands over the cavity leading edge and impinges on the cavity floor, whereas for open type cavity flow, the shear layer bridges the cavity. The tests were conducted by using a flat-plate model that permitted the cavity length to be remotely varied from 0.5 to 12 in. Cavity depths and widths were varied from 0.5 to 2.5 in. The flat-plate boundary layer approaching the cavity was turbulent and had a thickness of approximately 0.2 in. at the cavity front face for the test range of Mach number from 1.50 to 2.86. Values of $(l/h)_{cr}$ obtained when decreasing cavity length were generally less than those obtained when increasing cavity length. Values of $(l/h)_{cr}$ ranged from 10 to 13 for the present tests. A large improvement in the correlation of measured cavity centerline pressure distributions for cavities of various depths was obtained when both the cavity widthto-depth ratio (w/h) and length-to-depth ratio (l/h)were held constant rather than l/h alone. The effects of cavity width on the cavity pressure distributions were much greater for cavities having closed or transitional flow fields than cavities having open flow fields. Decreasing cavity width resulted in a reduction in $(l/h)_{cr}$. Three-dimensional effects in the form of large lateral pressure gradients occurred on the rear faces of the cavities that had closed cavity flow fields.

Introduction

Numerous investigations have been conducted over the past several decades to investigate the flow fields over cavities and to define the resulting local pressure distributions and acoustic levels within the cavities (e.g., refs. 1 through 6). These investigations have been conducted over a speed range from subsonic through hypersonic Mach numbers. The results obtained at supersonic speeds are particularly important for application to cavities on contemporary and future aircraft and missile configurations capable of sustained supersonic flight speeds. Some examples of requirements for cavities on these configurations consist of weapon bays for high-speed military aircraft and recessed areas on wrap-around-fin missiles that contain the fins before they are deployed. Existing data available in the literature show that cavity flow fields can occur that result in large local turning angles of the shear layer over the cavity; this gives rise to large cavity drag levels (e.g., refs. 7 and 8) as well as large impact pressures on components within the cavity. Such cavity flow fields can also result in adverse separation characteristics for a store being launched from the cavity (e.g., refs. 9 and 10). Large fluctuating pressure levels can also occur in cavities, which sometimes are severe enough to cause component failure of hardware within the cavity (ref. 11).

In general, data available in the literature show that at supersonic speeds, there are two fundamentally different types of cavity flow fields which have been classified as open and closed cavity flows. The type of flow field appears to be primarily a function of cavity length-to-depth ratio (l/h). As illustrated in figure 1, for values of l/h > 13, the cavity flow field is generally of the closed flow type. For this case, the shear layer expands over the cavity leading edge, impinges on the cavity floor and exits ahead of the rear face. Typical cavity floor pressure distributions for this case consist of low pressures occurring in the expansion region behind the front face followed by an increase in pressure and a pressure plateau occurring in the impingement region. Further downstream, as the shear layer approaches the cavity rear face, the pressure levels again increase and reach a maximum value just ahead of the rear face. The local flows over the cavity front and rear faces for the closed cavity flow field are very similar to the flows over rearward-facing and forward-facing steps, respectively. At $l/h \approx 12$, the cavity flow field is on the verge of changing from closed cavity flow to open cavity flow (decreasing l/h) and is referred to as "transitional cavity flow." For this case, the shear layer turns through an angle to exit from the cavity coincident with impinging on the cavity floor resulting in the impingement shock and the exit shock collapsing into a single wave. The corresponding pressure distribution shows that the extent of the plateau pressures in the impingement region has diminished and the pressure increases uniformly from the low values in the region aft of the front face to the peak values ahead of the rear face. For l/h < 10, the high pressures ahead of the rear face venting into the low pressure region downstream of the front face cause the shear layer to flow over or bridge the cavity. This type flow field is generally referred to as "open cavity flow." The pressure coefficients over the cavity floor are slightly positive and relatively uniform with the exception of a small adverse gradient occurring ahead of the rear face that is associated with the shear layer impinging on the outer edge of the rear face.

Because of the large differences in the flow fields for open and closed cavity flows and the resulting varied loadings on the cavities and their contents, it is very important to be able to define the l/hboundary that separates the two types of flow fields, which generally is referred to as the critical value of l/h or $(l/h)_{cr}$. The complexity of the cavity flow field limits the applications of current computational methods for determining local flow conditions in the cavity and therefore experimental techniques are generally relied upon to obtain this information. Unfortunately, the data available in the literature for a particular investigation are generally limited to a single Mach number and a small range of geometric variables. The purpose of the present investigation is to provide critical values of l/h and cavity pressure distributions from a single investigation for a range of supersonic Mach numbers, cavity lengths, cavity depths, and cavity widths. The tests were conducted by using a model that permitted the cavity length to be remotely varied from 0.5 to 12 in. which greatly facilitated determining $(l/h)_{cr}$. Cavity depths and widths were varied from 0.5 to 2.5 in. The boundary layer approaching the cavity was turbulent and had a thickness of approximately 0.2 in. at the cavity front face for the range of test Mach numbers from 1.50 to 2.86.

Symbols

 C_p

h	cavity depth, in.
k	height of roughness used for boundary- layer transition, in.
l	cavity length, in.
l_D, l_F	separation distances downstream of a rearward-facing step and upstream of a forward-facing step, respectively, in.
$(l/h)_{cr}$	value of l/h that separates open type cavity flow from closed type cavity flow
M_{∞}	free-stream Mach number
p	measured surface pressure, lb/ft^2
p_t	free-stream stagnation pressure, lb/ft^2
p_{∞}	free-stream static pressure, lb/ft^2
q_∞	free-stream dynamic pressure, lb/ft^2
R	free-stream unit Reynolds number per foot
_	_

pressure coefficient, $\frac{p - p_{\infty}}{q_{\infty}}$

 T_t free-stream stagnation temperature, °R

- w cavity width, in.
- x_1, x_2, x_3 axial surface distance on forward plate, cavity floor, and rear plate as defined in figure 3, in.
- y_1, y_2 surface distances on cavity front face and cavity rear face as defined in figure 3, in.
- z lateral surface distance on cavity rear face as defined in figure 3, in.
- δ boundary-layer thickness, in.

Abbreviations:

- FF cavity front face
- FL cavity floor
- FP forward plate ahead of cavity
- Loc location
- Orif orifice
- RF cavity rear face
- RP rear plate downstream of cavity

Apparatus and Test Conditions

Wind Tunnel and Test Conditions

The tests were conducted in the low Mach number test section of the Langley Unitary Plan Wind Tunnel (UPWT). This facility is a variablepressure continuous-flow wind tunnel with two test sections that permit a variation in Mach number from 1.50 to 4.60.

Ahead of each test section is an asymmetric nozzle that permits a continuous variation in Mach number from 1.50 to 2.90 in the low Mach number test section and from 2.30 to 4.60 in the high Mach number test section. The test sections are approximately 7 ft long and have a square cross-sectional area of approximately 16 ft². A complete description of the facility is given in reference 12.

The tests were conducted at zero angle of attack for the test conditions shown in the following table:

M_{∞}	p_t	T_t	R	k
1.50	1051	585	$2 imes 10^6$	0.0128
2.16	1349	585	$2 imes 10^6$.0128
2.86	1934	585	$2 imes 10^6$.0215

Models and Instrumentation

Shown in figure 2 are drawings and photographs of the cavity model assembly. The model consisted of a sting-mounted flat plate 41.9 in. long and 34.0 in.

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wide that housed a cavity with the cavity forward face being located 10.4 in. downstream of the flatplate leading edge. In the region ahead of the cavity, the flat-plate leading edge had a sweep angle of 0° that provided a two-dimensional boundary layer approaching the cavity. The outboard sections of the leading edge were swept 30° to reduce the plate planform area in order to reduce tunnel starting loads. Sweeping the leading edge also positioned the tip vortices downstream to minimize the effect of these vortices on the cavity flow field. Starting loads were further reduced by sweeping the plate trailing edge to reduce planform area. The leading-edge wedge angle (5°) was sufficiently small that supersonic attached flow was maintained at the leading edge throughout the test range of Mach number.

Cavity length was remotely controlled by a slidingblock assembly that formed the rear face of the cavity. (See fig. 2(a).) Cavity depth was varied by positioning the cavity floor at the desired locations relative to the flat-plate surface. Depths of 0.5, 1.0, 2.0, and 2.5 in. were tested. Individual cavity rearface blocks for each cavity depth were constructed and attached to the sliding-block assembly to maintain the remote positioning feature of this assembly. Figure 2(b) is a photograph of the model installed in the low Mach number test section of the UPWT. Shown in figure 2(c) is a photograph of typical block inserts that were installed in the cavity to vary its width. These inserts were constructed to vary width for several depths and lengths as shown in the following table:

l	h	l/h	w	w/h
12	1.0	12	2.5	2.5
			2.0	2.0
			1.5	1.5
			1.0	1.0
			.5	.5
6	1.0	6	2.5	2.5
			2.0	2.0
			1.5	1.5
		-	1.0	1.0
			.5	.5
6	0.5	12	2.5	5.0
			2.0	4.0
			1.5	3.0
			1.0	2.0
			.5	1.0
3	0.5	6	2.5	5.0
			2.0	4.0
			1.5	3.0
			1.0	2.0
			.5	1.0

A boundary-layer transition strip was applied to the flat-plate leading edge to ensure fully developed turbulent flow on the plate surface at the cavity front face for all test conditions. Two different roughness sizes were used to cover the test Mach number range. At the lower test Mach numbers of 1.50 and 2.16, the transition strip consisted of randomly distributed No. 50 sand elements (0.0128 in. nominal height) in a 0.06-in. band applied 0.40 in. behind the flat-plate leading edge measured in a streamwise direction. At the maximum test Mach number of 2.86, the transition strip consisted of individually placed No. 35 sand elements (0.0215 in. nominal height) arranged in a line that also was 0.4 in. behind the leading edge measured in a streamwise direction. The elements were spaced approximately 0.09 in. between centers.

The model was instrumented with 84 pressure orifices with locations as defined in figure 3. Most of the orifices (orifices 1-72) were located along the plate longitudinal centerline. The remaining orifices were located in lateral rows on the rear face with the number of rows depending on the cavity depth as shown in figure 3.

The pressures were measured by using electrical transducers connected to a pressure scanning system. A total of four scanners were used with tubing from 21 orifices connected to each transducer. Three reference pressures were also connected to each scanner to provide transducer calibration for each test point. The maximum reference pressure for each scanner was selected to approximately match the maximum anticipated pressure to be measured for the particular group of tubes connected to the scanner. The reference pressures and tunnel free-stream pressures were measured independently by precision mercury manometers. The pressure measurements were reduced to coefficient form and are presented in tables I through V.

Accuracy

Accuracy of the system for measuring the cavity pressures is better than 1 percent of the fullscale range of the electrical transducers; this includes all errors of linearity, hysteresis, and repeatability. Transducers with a maximum range of 5.0 lb/in² and 7.5 lb/in² were used, and accuracies in C_p resulted as follows:

	Transducer	ΔC_p	for M_{∞}	of —
Orifice	range, lb/in^2	1.50	2.16	2.86
1-30, 58-61, 69-72, 81-84	5.0	± 0.016	± 0.016	± 0.019
31-57, 62-68, 73-80	7.5	± 0.024	± 0.025	±0.029

The accuracy of the precision mercury manometers with which the reference pressures and the tunnel stagnation pressure were measured was $\pm 0.0035 \text{ lb/in}^2$.

The results of a test section calibration (ref. 12) indicate that the maximum variation in free-stream Mach number in the vicinity of the model installation was ± 0.02 for the test range of Mach number.

Based on pretest calibrations, the cavity slidingblock assembly could be positioned to a given value of x_2 with an accuracy of ± 0.005 in.

Results and Discussion

Schlierens

Effects of l/h. Typical schlieren photographs from the present tests showing both open and closed cavity flow fields are presented in figures 4 and 5 for cavities with depths of 0.5 and 1.0 in., respectively. For a given cavity depth, the cavity length-to-depth ratio was varied by remotely changing the cavity length. The photographs presented in figures 4(a) and (b)are for values of l/h of 24.0 and 16.0 and illustrate closed cavity flow fields. Coincident impingement and exit shocks occurred at l/h = 11.6 (fig. 4(c)), corresponding to a transitional cavity flow field as illustrated in figure 1. At l/h = 11.2 (fig. 4(d)), the impingement shock is no longer apparent; this indicates an open type cavity flow with this type of flow field existing for all values of l/h < 11.2(figs. 4(e) and (f)). The photographs presented in figure 5 show that for a cavity depth of 1.0 in., closed cavity flow occurred for l/h > 10.5, transitional flow occurred for $l/h \approx 10.5$, and open cavity flow occurred for l/h < 10.5.

Critical 1/h values. The remotely controlled sliding-block feature of the present cavity model not only expedited pressure data acquisition for a wide range of cavity lengths but also facilitated determination of the critical values of l/h. The procedure for determining these critical values of l/h was as follows. Initially, with the cavity flow field being of the closed flow type, the cavity length was decreased until the flow field changed to open cavity flow as determined by the abrupt disappearance of the combined impingement-exit shock in the schlieren. This critical value of l/h was recorded and identified as being associated with decreasing l/h. After further decreasing the cavity length to well within the open flow region, the cavity length was then increased until the flow field changed back to closed cavity flow as determined by the sudden reappearance of the combined impingement-exit shock. This critical value of l/h was also recorded and identified as being associated with increasing cavity length. The cavity length

was then increased to well within the closed cavity flow region before repeating the above procedure. The procedure was repeated four times, and averages of these four critical l/h values were determined for both increasing and decreasing cavity length. The maximum variation of l/h from this mean value was approximately ± 0.12 .

Presented in figure 6 are results that were obtained throughout the range of test Mach number for a cavity depth of 0.5 in. The data show that the critical values of l/h obtained by increasing cavity length were greater than those obtained by decreasing cavity length and that the magnitude of this hysteresis effect increased with increasing Mach number. The crosshatched region between these two curves indicates the uncertainty level for determining if the length-to-depth ratio of a fixed geometry cavity at a constant Mach number is critical. For values of l/h above the crosshatched region, the fixed geometry cavity should have closed cavity flow; for values below the crosshatched region, it should have open cavity flow; and for values that fall within the crosshatched region, the flow could be of either type. Also shown in figure 6 are estimated critical l/h values that are assumed to correspond to cavities having lengths that are equal to the sum of the separation distances downstream of a rearward-facing step and upstream of a forward-facing step with the step heights equal to the cavity depths as shown in the following equation:

$$(l/h)_{cr} = (l_D + l_F)/h$$
 (1)

where l_D and l_F are the separation distances behind a downstream-facing step and ahead of a forwardfacing step, respectively. Values of l_D and l_F were obtained from reference 1. The rationale behind such an estimate is that when the sum of the two separation distances is equal to the cavity length, then the high pressure ahead of the cavity rear face will vent into the low pressure region behind the front face and will cause the flow to be forced out of the cavity. This estimate underpredicted the measured critical values throughout the test Mach number range.

Centerline Pressure Distributions

Effects of l/h. Shown in figure 7 are typical cavity centerline pressure distributions for open, transitional, and closed cavity flows for a range of cavity depths and free-stream Mach numbers. Results are presented for the plate surface ahead of the cavity, the cavity front face, the cavity floor, the cavity rear face, and the plate surface downstream of the cavity. Presented in figure 7(a) are data for $M_{\infty} = 1.50$ and h = 0.5 in. The pressure distributions for l/h = 24.0are for closed cavity flow, for l/h = 13.0 for transitional flow just prior to changing from closed to open flow, for l/h = 12.6 for transitional flow just after changing, and for l/h = 8.0 and 1.0 for open cavity flow. The pressure distributions on the cavity floor for the different types of cavity flow fields are consistent with the hypothetical distributions shown in figure 1 and discussed in the introduction. On the forward plate ahead of the cavity, the pressure coefficients are essentially constant at a value of zero; this indicates that any disturbances created by the cavities are not propagated upstream. The pressure measurements on the cavity front face for l/h > 1.0are invariant with y_1/h ; however, the magnitudes of the pressure measurements are sensitive to the type of cavity flow field and are essentially equal to the pressure level at the most forward instrumented station on the cavity floor. In general, for the range of l/h shown, the pressure coefficients on the front face increase with decreasing l/h with the greatest changes occurring for values of l/h at which the flow switches from open to closed cavity flow. On the cavity rear face, large pressure gradients exist and large variations in pressure levels occur with varying l/h. These large gradients, in contrast to the almost constant pressures on the front face, result from the fact that the rear face is exposed to the approaching high energy flow similar to a forward-facing step, whereas the front face is exposed to an almost quiescent region similar to a rearward-facing step. Peak pressures on the rear face for a given value of l/h occurred at the outer edge of the rear face with the exception of the case of transitional cavity flow (l/h = 13.0)where a minimum pressure occurred in this region. This trend is observed through the test range of Mach number (figs. 7(a) through (c)). With increasing y_2/h from the outer edge of the cavities $(l/h \neq 13)$, the pressures decrease to a minimum value at approximately mid-depth followed by an increase in pressure with further increases in y_2/h toward the cavity floor. The maximum values near the cavity floor are approximately equal to the peak values on the cavity floor for those cases where a pressure orifice was located at $x_2/l = 1$. On the rear plate downstream of the cavity, large pressure gradients occurred for those cavities having the larger values of l/h. The large gradients are associated with closed cavity flow fields and occur in a region of flow separation downstream of the outer edge of the rear face that is formed as the flow exits from the cavity and fails to expand around the 90° corner. For the cavities with open cavity flow fields, the flow essentially bridges the cavity, resulting in minimal separation at the rear corner and hence only small pressure gradients in this region.

The pressure distributions for the 0.5-in-deep cavities presented in figures 7(b) and (c) for Mach numbers 2.16 and 2.86 are somewhat similar to the results shown for M = 1.50. One of the most noticeable effects of increasing Mach number is to reduce the magnitude of the peak pressures. Also at M = 2.86, transition from closed to open cavity flow occurred when decreasing l/h from 11.6 to 11.2 as compared with 13 to 12.6 for the two lower Mach numbers. This trend is consistent with data obtained from the schlieren tests presented in figure 6. The data obtained on the rear plate shown in figure 7 indicate that increasing Mach number results in an increase in the extent of the separation region downstream of the rear face. This trend was also observed in the schlieren system as evidenced by a downstream movement of the reattachment shock with increasing Mach number.

Shown in figures 7(d), (e), and (f) are cavity pressure distributions for the 1-in-deep cavity at Mach numbers 1.50, 2.16, and 2.86, respectively. For this cavity depth, the maximum cavity length of 12 in. limits the maximum value of l/h to 12 and therefore only transitional and open cavity flow fields would be expected. As discussed previously in the introduction and as shown by the data from the 0.5-indeep cavity in figures 7(a), (b), and (c), the extent of the flow impingement plateau pressures for the transitional flow field diminished; this resulted in monotonically increasing pressures in this region. The pressure distributions presented in figures 7(d)and (e) show that this is also true for the 1-in-deep cavity at $M_{\infty} = 1.50$ and 2.16. At $M_{\infty} = 2.86$, however, the floor pressure distributions for l/h = 10.5(fig. 7(f)) are very similar to pressure distributions shown previously for closed cavity flow in that the plateau pressures occur over a significant range of x_2/l in the flow impingement region. Since the flow has changed to open cavity flow at l/h = 10.0, it is not clear why the distributions at l/h = 10.5 are not more representative of transitional cavity flow. Another unanticipated variation in the pressure distributions for the 1-in. cavities with l/h = 12.0 and 10.5occurred on the rear face when increasing Mach numbers from 1.50 to 2.16 as may be seen by comparing figures 7(d) and (e). With increasing Mach number, a large increase in pressure level and pressure gradient occurred as compared with a decrease in pressure level shown for the 0.5-in. cavity (figs. 7(a) and (b)). These large pressures also occur on the rear face at $M_{\infty} = 2.86$ (fig. 7(f)). The flow field associated with these large pressures also results in a bow shock at the outer edge of the rear face as can be seen in figures 5(a) and (b). This bow shock was not apparent for the 0.5-in. cavity flow field (fig. 4(c)). The pressure distributions on the forward plate, front

face, and rear plate of the 1-in. cavity through the test range of Mach number and l/h are similar to the results shown for the 0.5-in. cavity. Also, the pressure distributions on the cavity floor and rear face of the 1-in. cavity with open cavity flow $(y_2/h \leq 11.2)$ are similar to the results obtained for the 0.5-in. cavity.

Pressure distributions for the 2.0-in-deep cavity were only obtained at Mach numbers 1.50 and 2.16, and these results are shown in figures 7(g) and (h). The maximum value of l/h that could be obtained at this depth was 6 and therefore all the pressure distributions shown are for open cavity flow. Generally, the trends of the variation of C_p with l/h and Mach number that are shown are similar to the open cavity flow results shown previously for the 0.5- and 1.0-indeep cavities; however, the peak pressure magnitudes on the cavity floors and rear faces are greater than obtained for the more shallow cavities. This trend is consistent with an observation from reference 1 where it was found that as the ratio δ/h increases ($\delta \approx \text{Con-}$ stant with Mach number for present tests), pressure gradients are smoothed out presumably because of the decreased momentum transfer to the cavity.

For the deeper cavities of the present tests (h = 2.0 and 2.5 in.) more pressure instrumentation is available on the cavity floor for a given value of l/h simply because l is greater and therefore more pressure orifices are exposed. This more detailed instrumentation on the cavity floor indicates that a different type of flow field occurs for the smallest value of l/h (l/h = 1, figs. 7(g) through (j)), as compared with l/h = 3.0 and 6.0. This effect of l/h was not apparent for the more shallow cavities (h = 0.5 and 1.0) because of the reduced number of orifices on the cavity floor. For the deeper cavities with l/h = 1, the data in figures 7(g) through (j) show that a much smaller peak pressure occurs on the cavity floor ahead of the rear face as compared with l/h = 3.0 and 6.0. Additionally for l/h = 1, lower pressures occurred on the cavity rear face. This change in the pressure distributions may be associated with the flow restructuring from a two-vortex scheme to a single-vortex scheme as observed in reference 13 by flow visualization techniques. For values of l/h ranging from 5.0 to 2.5, Shchukin observed two vortices, as shown in the top of sketch A, of approximately the same size. The rear vortex had considerably greater circulation intensity and its center was located somewhat above the midsection of the cavity. For $l/h \approx 2.0$, the flow pattern restructured to form one vortex as shown in the bottom of sketch A. The one-vortex pattern was retained with increasing cavity depth to l/h = 1. Although pressure distributions are not presented in reference 13, the authors state that for the single-vortex case, the pressure



patterns become more symmetrical about the cavity midlength; this trend is consistent with the present data. Also, heat-transfer distributions presented in reference 13 show a large reduction in heat transfer ahead of the rear face for the single-vortex case which could in part be caused by a pressure reduction in this region as measured in the present tests.

Shown in figure 8 are the variations with l/h of the pressure coefficients in the outer-edge regions of the front and rear cavity faces to further illustrate the effect of the cavity flow field on the cavity pressure distributions. The data for the 0.5- and the 1.0-in. cavities (figs. 8(a) and (b)) clearly illustrate the increase in pressure on the front face and the decrease in pressure on the rear face that occurs as the flow changes from closed to open cavity flow $(l/h \approx 10 \text{ to } 13)$. The data also show that for all cavity depths, a decrease in pressure occurs on both the front and rear faces at the very low values of l/h, which is much more pronounced at the larger cavity depths. This decrease in pressure could be associated with the cavity flow restructuring from a two-vortex scheme to a one-vortex scheme discussed previously.

Effects of w/h. Presented in figure 9 are summary plots showing cavity centerline pressure distributions for cavities of different depths at constant or approximately constant values of l/h. The data are presented for the test range of Mach number for values of l/h representative of transitional and open cavity flow. These data indicate that for the transitional cavity flow field, l/h in itself is not a satisfactory correlation parameter. The lack of correlation of data is particularly obvious on the cavity rear face at the higher Mach numbers. This trend may be partially due to a three-dimensional effect created as a result of the cavity width not being scaled properly as depth was varied. The cavity width for the data presented in figure 9 was held constant as depth varied; therefore, the scaling parameter w/h varied as h varied as shown in the figure. Data were obtained, however, for different cavity widths at selected cavity lengths and these data provide an opportunity to examine the effect of w/h. Presented in figures 10 and 11 are data obtained at different depths but constant values of l/h and w/h. The data are again presented for both transitional and open cavity flow and show that much better correlation of the results is obtained when holding both l/h and w/h constant for the different cavity depths than was obtained by holding only l/h constant. There is, however, some lack of agreement of the data along the rear portion of the cavity floor, the rear face, and the rear plate. The disagreement on the floor and rear face generally consists of an increase in pressure level with an increase in cavity depth. This disagreement could result from the variation of the parameter δ/h that occurs with varying h, since for the present test, δ remained approximately constant as discussed previously. The trend of the variation of C_p with δ/h on the floor and rear face consisting of a decrease in pressure with increasing δ/h is consistent with previously published data (ref. 1) and is attributed to the decreased momentum transferred to the cavity with increasing δ/h .

A complete set of data obtained at the various cavity widths is presented in figures 12 and 13 for cavity depths of 0.5 and 1.0 in., respectively. The data are presented with w/h as a parameter at constant values of l and h. The data generally show that the greatest effect of cavity width occurs for the cavities having closed cavity flow fields or for the few cases of open cavity flow where decreasing cavity width caused the flow field to change from open to closed cavity flow. The flow field changed from open to closed flow when decreasing cavity width from w/h = 4 to 3 for the 0.5-in-deep cavity with l/h = 12at both Mach numbers 1.50 and 2.16 (figs. 12(a)and (b)), respectively. Therefore, one effect of decreasing cavity width is to reduce the critical value of l/h, since as shown in figure 6, $(l/h)_{cr} = 13$ for the cavity with w/h = 5 as compared with $(l/h)_{cr} = 12$ for the cavities with $3 \le w/h \le 4$. For the 1.0-in-deep cavity with l/h = 12 (figs. 13(a) and (b)), w/h was only varied from 2.5 to 0.5, and for this range of w/h the cavity flow field remained the closed flow type. The data show, however, that increasing w/hresulted in pressure distributions in the flow impingement region changing from distributions typical of closed cavity flow (w/h = 0.5) to distributions typical of transitional cavity flow (w/h = 2.5). It is quite possible, therefore, that if the range of w/h for the 1.0-in. cavity had been extended from 2.5 to 5, the flow field may have changed from closed flow to open flow and comparable values of $(l/h)_{cr}$ would have been measured for both the 0.5-in. and 1.0-in. cavities.

Results presented in figures 12 and 13 show that the effects of cavity width on the pressure distributions for the cavities having open cavity flow fields were relatively small compared with those for the cavities with closed cavity flows. For open cavity flow, increasing cavity width generally resulted in an increase in pressures on the cavity rear face and on the rear portion of the cavity floor.

Lateral Pressure Distributions

All the pressure distributions presented to this point have been along the cavity centerline and therefore do not give any indication of the lateral variations that may occur as a result of three-dimensional effects created because the cavity had finite width. An example of the complexity of such a flow field is shown in figure 14, reproduced from reference 8. For the closed cavity flow field, a pair of vortices are formed at the outer edges of the cavity side walls as the flow expands into the cavity. At the cavity rear face, the vortices are well developed and are observed to continue downstream of the cavity. The impingement of these vortices on the cavity rear face would be expected to create lateral pressure gradients in this region and in particular toward the outer edge of the rear facc. Also, as the cavity width decreases, there may be an interaction between the vortices which would further complicate the lateral pressure gradients. The vapor screen photographs shown in figure 14 for the case of open cavity flow show a less complicated flow field as there are no apparent vortices or shock waves.

As mentioned in the section "Models and Instrumentation," a lateral array of pressure orifices were installed on the cavity rear face of the present model, the rear face being selected since this was the region believed to have the maximum pressures as well as maximum pressure gradients. Measured pressures from the lateral row of orifices closest to the outer edge of the cavity rear face $(y_2 = 0.25 \text{ in.})$ are presented in figure 15 for cavity depths ranging from 0.5 in. to 2.5 in. Data are presented through the test range of Mach number for both open and closed cavity flow fields. The data presented in figures 15(a) and (b) for $l/h \ge 12$ are for closed cavity flow and show large lateral pressure gradients with the locations of the peak pressure ranging from the longitudinal centerline to the most outboard instrumentation location. The distributions are relatively symmetrical about the longitudinal centerline. The large pressure gradients and the location of the peak pressures are probably associated with the impingement of the edge vortices on the rear face.

Data presented in figure 15 for 3 < l/h < 6 are for open cavity flow, and even for this type flow field, large lateral pressure gradients occur although the magnitudes are considerably less than obtained for closed cavity flow. Also, within this range of l/h for the cavities with depths of 2 and 2.5 in. (figs. 15(c)and (d)), the distributions were in some cases unsymmetrical about the cavity centerline. Results presented in figure 15 for l/h = 1 show relatively uniform lateral pressure distributions of small magnitudes which are in all cases symmetrical about the cavity centerline. These differences in the lateral pressure distribution compared with the results shown for $3 \leq l/h \leq 6$ further substantiate the previous discussion concerning changes in the cavity flow field when decreasing l/h to values less than approximately 3.

Shown in figure 16 is the effect of cavity width on the rear face lateral pressure distributions. These results were only obtained for the 0.5- and 1.0-in-deep cavities and for values of l/h of 6 and 12. When the cavity flow field was of the open flow type, the cavity lateral pressure distributions were relatively insensitive to variations in cavity width as shown for l/h = 6 in figures 16(a) and (b) and for l/h = 12 and w/h > 4 in figure 16(a). Pressure levels representative of closed type flow fields were measured for the 1-in-deep cavities at l/h = 12 for the test range of w/h from 2.5 to 1.0 (fig. 16(b)). These data show that decreasing w/h results in a decrease in the pressure levels and a reduction in the lateral pressure gradient about the centerline. Pressure levels indicative of closed cavity flow were also measured for the 0.5-in-deep cavities at values of $w/h \leq 3$; however, for these reduced width cavities, insufficient instrumentation was available to determine the effect of width on the lateral pressure gradients.

Summary of Results

An experimental investigation has been conducted to define cavity pressure distributions for a range of free-stream Mach numbers and cavity geometries. These pressure distributions together with schlieren photographs were used to define the critical values of cavity length-to-depth ratio $(l/h)_{cr}$ that separate open type cavity flows from closed type cavity flows. For closed type cavity flow, the shear layer expands over the cavity leading edge and impinges on the cavity floor, whereas for open type cavity flow, the shear layer bridges the cavity. The tests were conducted by using a flat-plate model that permitted the cavity length to be remotely varied from 0.5 to 12 in. Cavity widths and depths were varied from 0.5 to 2.5 in. The plate boundary layer approaching the cavity was turbulent and had a thickness of approximately 0.2 in. at the cavity front face for the test range of Mach number from 1.50 to 2.86. The results from these tests are summarized as follows:

(1) Critical cavity length-to-depth ratios obtained when decreasing cavity length were generally less than obtained when increasing cavity length. The magnitude of this hysteresis effect increased with increasing Mach number. Values of $(l/h)_{cr}$ ranged from 10 to 13 for the present tests.

(2) For the transitional cavity flow field, measured pressures in the flow impingement region on the cavity floor downstream of the front face increased monotonically with surface length with no apparent plateau-pressure region through the test range of Mach number for the 0.5-in-deep cavity and at Mach numbers 1.50 and 2.16 for the 1.0-in-deep cavity.

(3) A reduction in pressures occurred on both the front and rear faces of the cavity for values of cavity length-to-depth ratio (l/h) less than approximately 3, which corresponds to the approximate value of l/h where previous investigators have observed a restructuring of the cavity flow from a twovortex scheme to a one-vortex scheme.

(4) A large improvement was obtained in the correlation of measured cavity centerline pressure distributions for cavities of various depths when both the cavity width-to-depth ratio (w/h) and l/h were held constant rather than l/h alone.

(5) Decreasing cavity width resulted in a reduction of $(l/h)_{cr}$.

(6) The effects of cavity width on the cavity pressure distributions were much greater for cavities having closed or transitional flow fields than cavities having open flow fields.

(7) Three-dimensional effects in the form of large lateral pressure gradients occurred on the rear faces of the cavities that had closed cavity flow fields. These large gradients also occurred for the cavities with open cavity flow fields when $3 \leq l/h \leq 6$. These gradients were generally symmetrical about the cavity longitudinal centerline for the closed cavity flows but in some cases they were asymmetrical for the open cavity flows. Relatively small lateral pressure gradients were measured on the rear faces of the cavities having l/h = 1, and these gradients were generally symmetrical about the cavity longitudinal centerline.

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OF POOR QUALITY

Table I. Pressure Coefficients for h = 0.5 in. and w = 2.5 in.

(a) $M_{\infty} = 1.50$

					<u> </u>					C _p	for 1/h	of -]
Orif	Loc	24	20	16	14	13	12.532	12	11	10	9	8	7.5	7	6	5	4.8	4	3	2	1.5	1
1 2 3 4	FP ↓ ♥	.0016 .0048 .0121 .0105	.0022 .0047 .0103 .0102	.0025 .0028 .0107 .0111	.0016 .0050 .0107 .0097	.0032 .0053 .0115 .0099	.0009 .0029 .0098 .0100	.0015 .0039 .0100 .0105	.0007 .0052 .0112 .0120	.0018 .0044 .0102 .0120	.0011 .0047 .0104 .0129	.0021 .0051 .0106 .0126	.0013 .0058 .0107 .0118	.0017 .0047 .0124 .0129	.0007 .0048 .0101 .0108	.0021 .0053 .0109 .0107	.0022 .0051 .0111 .0129	.0012 .0037 .0100 .0118	0002 .0037 .0097 .0110	.0009 .0046 .0098 .0112	.0020 .0043 .0101 .0117	.0010 .0051 .0095 .0111
5 6	FF FF	1963 1944	1968 1938	1968 1935	1912 1879	1504 1478	0312 0343	0135 0156	.0024	.0119	.0167 .0156	.0194 .0180	.0228	.0255	.0317	.0397 .0368	.0434 .0396	.0531 .0499	.0705	.0656	.0615 .0771	.0540 .0760
64 65 66 67 68	RF	.6270 .5677 .5891 .6752 .7125	.6349 .5866 .6029 .6562 .6566	.6295 .5939 .6000 .6122 .5726	.6079 .5754 .5792 .5781 .5266	.5512 .5136 .5158 .5181 .4794	.2717 .2440 .2432 .2751 .3013	.2349 .2070 .2057 .2327 .2666	.2033 .1729 .1712 .1981 .2484	.1838 .1536 .1489 .1811 .2397	.1747 .1433 .1399 .1798 .2514	.1746 .1418 .1347 .1746 .2524	.1688 .1370 .1289 .1712 .2535	.1665 .1302 .1247 .1670 .2556	.1608 .1243 .1148 .1663 .2635	.1662 .1249 .1165 .1700 .2807	.1686 .1240 .1142 .1683 .2809	.1679 .1253 .1126 .1647 .2796	.1496 .1106 .0942 .1287 .2354	.1091 .0725 .0528 .0690 .1766	.0719 .0637 .0480 .0353 .1059	.0807 .0676 .0485 .0535 .1487
69 70 71 72	RP	1844 2465 2665 0899	1720 2418 2713 0997	1430 2255 2490 0903	1241 2057 2091 0811	1221 1899 1629 0695	0732 0812 0491 0274	0604 0636 0386 0209	0482 0474 0292 0170	0366 0343 0210 0109	0293 0278 0154 0064	0250 0258 0131 0049	0224 0247 0117 0048	0221 0238 0119 0039	0181 0206 0097 0033	0166 0193 0094 0032	0170 0203 0108 0048	0157 0193 0083 0029	0138 0178 0065 0015	0027 0077 0014 .0030	.0011 0026 .0020 .0035	.0007 0005 .0035 .0040
73 74 75 76	k₽	.5232 .7164 .6681 .5333	.5843 .7675 .7270 .5930	.6274 .8005 .7381 .6408	.6249 .7660 .7140 .6351	.5921 .6917 .6339 .5958	.2866 .2779 .2673 .2771	.2298 .2351 .2250 .2238	.1806 .1899 .1841 .1756	.1380 .1671 .1568 .1348	.1054 .1535 .1412 .0952	.0926 .1472 .1391 .0759	.0857 .1414 .1363 .0692	.0755 .1341 .1286 .0631	.0613 .1267 .1210 .0514	.0554 .1288 .1216 .0501	.0544 .1252 .1197 .0479	.0527 .1273 .1212 .0462	.0651 .1233 .1047 .0689	.0435 .0901 .0488 .0780	.0411 .0554 .0523 .0414	.0449 .0504 .0511 .0450
$\begin{array}{c} 11\\ 13\\ 14\\ 15\\ 16\\ 17\\ 18\\ 20\\ 22\\ 23\\ 24\\ 25\\ 26\\ 27\\ 28\\ 29\\ 32\\ 33\\ 34\\ 27\\ 28\\ 29\\ 32\\ 33\\ 34\\ 35\\ 36\\ 37\\ 38\\ 39\\ 40\\ 41\\ 42\\ 43\\ 44\\ 45\\ 46\\ 47\\ 48\\ 49\\ 50\\ 51\\ 52\\ 55\\ 55\\ 55\\ 57\\ \end{array}$		2252 2437 2355 2110 1719 1257 .0047 .0828 .0963 .0963 .0963 .0973 .0983 .0953 .0953 .0953 .0953 .0953 .0953 .0953 .0571 .0550 .0571 .0550 .0571 .0550 .05510 .0496 .0494 .0425 .0608 .1471 .2545 .3254 .3256 .3254 .3254 .3256 .3254 .32566 .32566 .32566 .325666666666666666666666666666666666666	2257 2447 2370 2116 1720 1245 .0057 .00566 .0056 .0056 .0056 .0056 .0056 .0056	- 2254 - 2254 - 2430 - 2237 - 2124 - 1722 - 1253 - 0784 - 0043 - 0043 - 0043 - 0046 - 0087 - 0084 - 0087 -	2196 2381 2093 1711 1258 .0050 .0714 .0958 .1183 .1418 .1697 .2062 .2451 .2864 .4284 .4284 .4599 .4763 .4961 .5104 .6407	1776 1950 19518 1817 1187 1187 0822 .1164 .00322 .1164 .1468 .1792 .2105 .2453 .2453 .2453 .2453 .2453 .4435 .4435 .4435 .4570 .5783	0376 0454 0445 04490 0500 0369 0129 .0250 .0424 .0608 .10783 .0990 .1168 .1306 .1492 .1793 .1993 .2958	0196 0270 0230 0316 0328 00259 0059 0256 .0406 .0555 .0701 .0844 .0992 .1121 .1243 .1388 .1707 .2532	0032 0085 0111 0133 0157 0153 0152 0013 0268 0397 0519 0634 0744 0847 0920 0964 2223	.0070 .0031 .0012 -00042 -0042 -0045 .0045 .0045 .0388 .0484 .0585 .0646 .0671 .0771 .1192	.0133 .0103 .0089 .0064 .0045 .0123 .0045 .0123 .0312 .0383 .0431 .0489 .0601 .1073 .1949	.0175 .0144 .0135 .0012 .0089 .0175 .0325 .0379 .0517 .1040 .1929	.0195 .0165 .0158 .0123 .0125 .0102 .0113 .0204 .0328 .0482 .0993 .1896	.0220 .0173 .0160 .0136 .0135 .0134 .0143 .0143 .0143 .0428 .0955 .1850	.0243 .0187 .0172 .0159 .0151 .0130 .0128 .0302 .1801	.0293 .0221 .0190 .0172 .0134 .0147 .0276 .1877	.0293 .0223 .0193 .0159 .0123 .0154 .0382	.0327 .0243 .0170 .0118 .0227 .0823 .1917	.0421 .0197 .0208 .0735 .1695	.0418 .0561 .1233	.0489 .0769	.0662

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Table I. Continued

(b) $M_{\infty} = 2.16$

										C _p for	l∕h of	-								
Urit	LOC	24	20	16	14	13	12.532	2 12	11	10	8	7	6	5	4.8	4	3	2	1.5	1
1	FP	.0042	.0033	.0039	.0021	.0038	.0030	.0030	.0035	.0027	.0022	.0018	.0018	.0019	.0018	.0017	.0018	.0017	.0014	.0018
3		.0034	.0026	.0033	.0012	.0028	.0027	.0024	.0024	.0019	.0017	.0017	.0019	.0010	.0008	.0008	.0008	.0009	.0009	.0008
. 4		.0057	.0057	.0057	.0040	.0057	•0054	.0048	•0047	•0050	•0045	.0042	•0040	.0037	.0038	.0038	.0039	.0038	.0032	.0035
5 6	FF FF	1748 1765	1748 1766	1748 1766	1745 1760	1692 1713	0287	0129 0149	.0008 0010	.0069	•0134 •0120	.0177 .0167	.0223	.0262 .0243	•0268 •0246	.0314 .0297	.0327	.0349 .0388	.0304 .0409	.0274 .0431
64	RF	.5168	.5197	.4950	.4652	.4421	.1984	.1704	.1458	.1354	.1245	.1146	.1056	.0978	•0960	.0865	.0696	.0529	.0469	.0417
66		.4612	.4803	.4711	.4420	.4151	.1790	.1491	.1251	.1134	.1012	.0900	.0809	.0741	+0639	.0560	.0394	.0250	.0338	.0207
67		.5458	.5127	.4591	.4246	.4116	.2202	.1877	1598	.1523	.1386	.1287	.1189	.1043	•0997	.0864	.0739	•0447	.0330	.0242
68	<u> </u>	.6146	• 5468	.43/8	.4024	.4097	.2780	.24/7	•2265	.2257	.2236	.2200	•2112	.1950	.1899	.1698	.1/89	.1500	.1451	.1075
69 70	RP	.0236	.0352	.0379 0247	.0338 0241	.0344 0220	0073	0066	0054	0042	0013	0005	.0004	.0027	.0026 0058	.0039 0035	.0042 0006	.0062 .0026	.0063	.0068 .0039
71	[0887	0778	0687	0655	0630	0339	0269	0194	0138	0088	0071	0065	0058	0061	0046	0022	.0012	.0025	.0028
		1210	1150	0900		0005		0210				0040		0050		0039	0023	0002		.0012
73		.4328	.5042	.5829	• 55/3 • 5445	.5032	-1867 -1852	.1409	•1062 •1353	.0/69	.0480	.0388	+0307	.0230	.0215	.0155	.0133	.0121	.0145	.0178
75		.5561	. 5936	.5571	.5092	.4776	.1778	.1516	+1282	.1111	.0970	.0848	•0740	.0646	.0634	.0535	.0319	.0215	.0192	.0208
//		.424/	.4964	.5629	•5283	.4818	•1848	.1342	.1025	.0763	.0426	.0328	.0249	.01/2	.0152	.0123	.0270	.0132	.0156	.0182
	FL	1825	1828	1827	1819	1780	0326	0175	0038	.0033	+0109	.0133	.0157	.0168	.0170	.0201	.0230	.0222	.0193	.0379
14		1573	1575	1577	1574	1568	0395	0245	0095	0002	.0075	.0093	.0107	.0116	.0112	.0062	.0046	.0599		
15		1186	0820	1190	1191	1202	0419	0267	0117	0022	+0060	.0082	•0104	.0081	.0066	.0036	.0353			
17		0498	0499	0508	0509	0509	0374	0245	0114	0042	.0041	.0065	.0066	.0052	.0063	.0463				
18	1	0233	0243	0245	0246	0223	0292	0188	0084	0025	.0047	.0078	.0067	.0139	.0222	.0984				
22		.0379	.0398	.0408	.0429	.0756	.0180	.0189	.0205	.0202	.0200	.0255	.1226							
23		.0461	.0478	.0490	•0562	.1016	·0304	.0290	•0278	.0264	•0228	.0655								
25		.0547	.0577	.0579	.1009	.1537	.0556	.0393	.0333	.0378	.0712	.1505								
26		.0566	•0592	.0593	1339	.1745	•0678	.0594	-0511	.0416	.1399									
28		.0555	.0590	.0594	.1938	.1943	.0792	.0665	.0576	.0443										
29		.0547	•0573	.0610	-2142	.2309	.1008	.0848	.0637	.0836										
32		.0524	•0522 •0499	.1820	•2648 •2892	.2990	-1248 -1260	.0942	.1634											
34		.0488	.0495	.2545	.3120	.3357	.1342	.1878												
35	11	.0489	+0492	.2731	.3344	.3620	•2251													
37		.0479	-0482	.3106	-3879															
38		.0467	.04/1	.3348	.4806															
40		.0436	.0661	.3792																
41		.0423	·1569 ·2466	.4040																
43		.0392	•2874																	
44.		.0371	•3146 •3370																	
46		.0340	.3605																	
47	11	.0320	• 37 59 • 3785																	
49		.0935	.4007																	
50		.2213	•5453																	
52		.3101																		
53		.3367																		
55		3570																		
56	1	.3423																		
L3/		1.3/41																		

Table I. Concluded

($[\mathbf{c}]$	M_{∞}	=	2.	86
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					C _p	for 2/h	of -			
Orif	Loc	24	20	16	12	11.6	11.2	8	4	2
1	FP	.0018	.0011	.0001	.0019	.0005	.0018	.0011	.0015	.0016
2		.0027	.0007	.0010	.0007	.0018	.0014	0007	.0007	0008
3		.0020	.0001	.0003	0004	~.0005	.0011	0009	.0009	0001
4	1	.0006	.0007	.0014	.0006	.0007	.0008	0002	.0007	0002
5 6	FF FF	1263 1284	1262 1285	1264 1284	1262 1284	1259 1281	0071 0080	.0088 .0069	.0209 .0198	.0201 .0195
64	RF	.3596	.3786	.3932	.3305	.3222	.1076	.0832	.0533	.0300
65		.3379	.3568	.3684	.3095	.3006	.0950	.0705	.0443	.0219
66		.3289	.3459	.3509	.29/1	•28/1	.0954	.0691	.0368	.0134
68		.3639	.3303	.3409	.2890	•2688	.1881	.1768	.1315	.1217
69	RP	.0522	.0559	.0506	.0431	.0426	.0057	.0050	.0060	.0052
70	Ĩ	.0122	.0143	.0150	.0121	.0140	0065	~.0069	0004	.0017
71		0170	0138	0116	0103	0090	0125	0085	0038	.0021
72	. 🕴	0381	0344	0294	0281	0274	0129	0085	0047	0004
73	RF	.5328	.6675	.9431	.6835	.6239	.0801	.0335	.0096	.0074
74		.4422	.4216	.3613	.3220	.3161	.1033	.0692	.0399	.0114
75	j.	.3872	.3796	.3405	.3031	.3061	.0955	•0662	.0336	.0124
76		.5293	.6476	.7903	.5783	.5351	.0787	.0324	.0103	.0082
11	FL	1272	1276	1273	1275	1272	0121	.0068	.0125	.0134
13		1233	1234	1232	1234	1231	0140	.0039	.0083	.0178
14		0986	09/8	0987	0989	0990	0161	.0024	.0029	.0355
16		0468	0690	0471	0489	- 0478	- 0173	- 0000	.0000	
17		0306	0310	0311	0318	0306	0145	0008	.0314	
18		0160	0156	0159	0167	0141	0096	.0004	.0608	
20		.0021	.0023	.0014	.0098	.0337	.0014	.0061		
22		.0148	.0138	.0128	.0929	.1212	.0173	.0130		
23		.0195	.0171	.0175	1324	.1483	.0236	.0138		
24		.0229	.0202	.0206	.15/4	•1612	.0301	.0209		
26		.0255	.0250	.0241	1835	1783	.0337	.0323		
27		.0284	.0271	.0275	.1874	.1893	.0465	.0702		
28		.0299	.0281	.0289	.1961	.2001	.0490			
29		.0304	.0279	.0325	.2067	.2180	.0508			
32		.0314	.0295	.1730	.2563	.2778	.1265			
33		.0305	.0279	.1961	.2984	.3248				
34		.0277	.0283	.2106	.3375					
35		.0281	.02/1	•2190						
30		.0274	-0207 -025×	.2320						
38		.0268	.0276	.2474						
39		.0268	.0492	.2719						
40		.0267	.1134	.2973						
41		.0242	.1727	.3508						
42		.0244	.2032	.3998						
43		.0232	.21/1							
44		.0237	.2313							
46		.0243	.2622							
47		.0279	.2791							
48		.0657	.2993							
49		.1496	.3355							
50		.1919	.3871							
52		.2279								
53		.2430								
54		.2605								
55		.2646								
56	¥	.2760								
57	1	.3069								

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Table II. Pressure Coefficients for h = 1.0 in. and w = 2.5 in.

(a) $M_{\infty} = 1.50$

					<u> </u>			<u> </u>	C _p for	l/h of	-					· .	
Orif	LOC	12	11.5	11	10	9	8	7.5	5 7	6	5	4.8	4	3	2	1.5	1
1	FP	.0017	.0042	.0064	.0061	.0043	.0041	.0065	.0042	.0028	.0041	.0045	.0051	.0046	.0039	.0027	.0041
2		.0036	.0068	.0082	.0078	.0069	.0053	.0093	.0061	.0071	.0056	.0071	.0077	.0056	.0072	.0056	.0070
4	¥	.0095	.0117	.0127	.0113	.0126	.0105	.0125	.0101	.0111	.0112	.0112	.0121	.0104	.0116	.0108	.0112
5	FF	1869	1779	0045	.0087	.0127	.0146	.0179	.0208	.0272	.0342	.0386	.0514	.0708	.0587	.0346	.0320
6	1	1858	1772	0055	.0074	.0120	.0146	.0173	.0190	.0253	.0340	.0365	.0493	.0674	.0457	.0286	.0228
<u> </u>	<u> </u>	1040	1/39	0094		.0124	.0140	.0185	.0190	.0203	.0341	•0300	.04/4	.0390		.0277	.0324
62	RF	.6862	.6589 .6368	.2527	.2336	.2412	.2735	.3018	.3027 .2192	·2933 ·2020	.2481	.2451	.2515	.2106	.1520	.0709	.0428
64		.7594	.7228	.2256	.2028	.2109	.2497	.2651	.2867	•2629	.2191	.2133	.2138	.1565	.0875	.0415	.0145
65		.8162	.7586	.2491	.2411	.2486	.2955	.3149	.3342	.3269	.2774	.2689	.2707	.1996	.1064	.0487	.0223
67	V	.8666	.7998	.3164	.3412	.3696	.4317	.4631	.4916	.4888	.4337	.4316	.4456	.3391	.2253	.1354	.0740
69	RP	3399	3292	1026	0695	0623	0654	0686	0682	0678	0620	0598	0557	0403	0209	0081	.0026
70		3357	3228	0792	~.0516	0465	0567	0647	0708	0681	0536	0484	0489	0362	0177	0038	.0027
72	🕴	1510	1414	0269	0126	0058	0016	0076	0064	0089	0080	0066	0059	0043	0007	.0039	.0062
73	RF	.7180	.7002	.2904	.1880	.1547	.1449	.1400	.1364	.1183	.1102	.1019	.0988	.0898	.0487	.0263	.0268
74	11	.8993	.8534	.3076	.3066	.3180	.3564	.3742	.3961	.3726	.3061	.2990	.3088	.2717	.1569	.0608	.0302
75		.8484	•7958	.3147	.2935	.3098	.3538	.3692	.3913	.3732	.3063	.1062	.3173	.2743	-1565	.0607	.0307
77	11	.7082	•6843	.2650	.2320	.2367	.2661	.2744	.2841	.2677	.2247	.2171	.2211	.1848	.1148	.0749	.0452
78		.6736	•6530 6414	.2417	.2156	.2188	.2462	.2584	.2685	•2514 2516	.2140	.2100	.2129	.1803	·1292	.0779	.0451
80	♥	.6704	•6520	.2570	.2302	.2256	.2544	.2606	.2703	.2568	.2168	.2121	.2206	.1868	.1105	.0718	.0459
11	FL	1925	1856	0136	.0034	.0105	.0146	.0147	.0175	.0233	.0281	.0305	.0411	.0492	.0313	.0240	.0302
13		2000	1924	0183	.0014	.0072	.0132	.0110	.0146	•0187	.0163	.0166	.0218	.0298	0083	.0158	.0356
15		1991	1932	0194	0005	.0074	.0123	.0108	.0141	.0157	.0195	.0205	.0243	.0419	.0338	.0552	.0443
16		1991	1926	0193	0002	.0066	.0106	.0080	.0109	.0165	.0214	.0227	.0299	.0287	.0562	.0802	
18		2011	1929	0189	0010	.0067	.0113	.0082	.0110	.0150	.0209	.0226	.0310	0012	.1438		
19		1995	1929	0196	0017	.0058	.0077	.0067	.0089	.0156	.0225	.0253	.0284	.0146			
20		1965	1896	0193	0027	.0043	.0072	.0051	.0088	.0158	.0225	.0235	.0163	.0591			
23		1475	1438	0156	.0015	.0055	.0054	.0109	.0121	.0207	.0204	.0180	.0366				
24		0912	1186	0166	0005	.0045	.0044	.0105	.0139	.0222	.0168	.0158	.0831				
26		0609	0620	0150	0011	.0036	.0074	.0139	.0182	.0250	.0240	.0412	.3195				
28		.0028	0022	0133	0008	.0038	.0099	.0166	.0204	.0271	.0489	.0758					
29		.0336	.0274	0076	.0029	.0078	.0167	.0235	.0292	.0376	.1999	.3264					
30		.0625	.0561	0037	.0042	.0102	.0212	.02/8	.0351	.04/1	.3150						
32		.1061	.1109	.0101	.0133	.0199	.0302	.0411	.0465	.1129							
33	11	.1255	.1336	.0167	.0184	.0230	.0353	.0465	.0521	.2491							
35		.1540	.1732	.0312	.0281	.0347	.0455	.0534	.0783								
36		.1649	.1918	.0391	.0342	.0383	.0473	.0583	.1219								
38		.1849	.2330	.0485	.0390	.0430	.04/8	.1158	.3700								
39		.2002	·2595	+0652	.0503	.0481	.0688	-2264									
40		.2525	.2876	.0714	.0569	.0502	.2135	.3201									
42		-2862	.3503	.0897	.0645	.0519	.3342										
43		.3608	.3/81	.1039	.0673	.1010											
45		.3971	•4403	.1094	.0658	.1937											
40		.4270	-4820	.1185	.0778	.2936											
48		.4799	•5059	.1222	.1116	•											
50		.5115	•5241 •5391	.1206	.1927												
51		.5322	• 55 38	.1233													
52		.5474	.5719	.1451													
54		.5837	.5391	.2871													
55		-5784	·5996														
57	1	.6102	.,,														

Table II. Continued



(b) $M_{\infty} = 2.16$

[T							c	n for l	/h of -							
Orif	Loc	12	11.5	11	10.5	10	9	8	7.5	7	6	5	4.8	4	3	2	1.5	1
1	FP	.0016	.0019	.0023	.0022	.0019	.0014	.0019	.0019	.0016	.0014	.0016	.0020	.0019	.0018	.0021	.0020	.0017
3		.0017	.0025	.0031	.0026	.0024	.0021	.0025	.0024	.0019	.0016	.0022	.0021	.0027	.0022	.0022	.0018	.0021
4	V	.0030	.0039	.0041	.0040	•0039	.0037	.0038	.0043	.0037	.0034	.0032	.0038	.0041	.0040	.0031	.0036	.0033
5	FF	1711	1709	1704	1708	.0073	.0085	.0104	.0116	.0123	.0142	.0178	.0187	.0230	.0262	.0175	.0132	.0131
6		1729	1723	1719	1723	.0069	.0083	.0103	.0114	.0123	.0147	.0171	.0185	.0237	.0249	.0161	.0117	.0087
<u>⊢́</u>	<u> </u>	.104/		1057	1045				.0114					.0224		.0158		.01/0
62	RF	-8155	.8135	.7977	.7739	.1448	.1331	-1304 0×48	.1278	.1227	.1140	.1082	.1072	.1090	•0846 0463	.0472	.0350	.0242
64		.9787	.9693	.9269	.8651	.1299	.1199	.1159	.1139	.1079	.0965	.0864	.0834	.0818	.0529	.0216	.0185	.0033
65		1.1413	1.1256	1.0754	1.0064	.1619	.1535	.1493	-1449	.1389	.1250	.1110	-1081 1489	.1065	.0701	.0291	.0209	.0076
67	V.	1.4494	1.4052	1.3808	1.2885	.2707	.2700	.2702	.2630	.2524	.2297	.2023	.1960	.1880	.1420	.1059	.0781	.0485
69	RP	0797	0721	0671	0563	0171	0150	0143	0142	0172	0181	0129	0120	0063	0034	.0023	.0052	.0077
70		1168	1129	1114	1014	0276	0237	0222	0217	0239	0234	0194	0188	0156	0104	0005	.0025	.0061
72	¥	1765	1762	1759	1632	0153	0125	0106	0094	0115	0112	0078	0081	0080	0043	.0015	.0023	.0034
73	RF	.7729	.8126	.8355	.8337	.1006	.0716	.0612	.0584	.0550	.0512	.0454	.0426	.0408	.0229	.0146	.0063	.0113
74		1.2261	1.2325	1.1970	1.1318	.1970	.1741	.1628	.1574	.1519	.1401	.1364	.1388	.1594	.1113	.0341	.0188	.0153
75		.7515	1.1990	1,1357	1.0669	.1843	•1610 •0678	.1543	.1473	.1425	.1322	.1294	.1297	.1501	.1045	.0339	.0196	.0157
17		.7991	.8118	.8036	.7778	.1415	.1246	.1169	.1129	.1080	.0994	.0925	.0923	.0963	.0629	.0337	.0270	.0244
78		.7574	.7541	.7606	.7348	.1342	.1197	.1122	.1088	.1046	.0972	.0911	.0902	.0927	.0677	.0442	.0354	.0242
80	¥	.7784	.7783	.7845	.7473	.1331	.1160	.1101	.1073	.1027	.0943	.0883	.0863	.0892	.0598	.0324	.0257	.0246
11	FL	1749	1748	1744	1747	.0044	.0076	.0091	.0101	.0108	.0133	.0158	.0160	.0203	.0205	.0132	.0097	.0135
14		1835	1835	1823	1834	.0025	.0060	.0050	.0041	.0030	.0033	.0049	.0053	.0071	.0123	0034	.0094	.0246
15		1844	1848	1836	1846	.0022	.0061	.0050	.0048	-0042	.0044	.0057	.0066	.0096	.0164	0023	.0278	
17		1780	1784	1772	1785	.0020	.0053	.0048	.0047	.0049	.0056	.0080	.0078	.0115	0021	.0077	.0372	
18		1679	1682	1670	1683	.0016	.0048	.0046	.0054	.0063	.0079	.0098	.0109	.0154	0062	.0537		
19		1358	1532	1351	1532	.0016	.0039	.0040	.0055	.0063	.0088	.0104	.0113	.0137	.0059			
22		0975	0962	0980	0922	.0042	.0073	.0091	.0105	.0089	.0105	.0165	.0163	.0027	.1066			
23		0761	0746	0759	0705	.0032	+0058	.0083	.0097	.0082	.0102	.0148	.0108	.0135				
25		0331	0316	0330	0277	.0021	.0048	.0077	.0099	.0087	.0113	.0050	.0078	.0965				
26		0134	0121	0133	0081	.0016	.0043	+0079	.0101	.0092	.0114	.0099	.0187	.1368				
28		.0199	.0214	.0203	.0254	.0024	.0044	.0085	.0119	.0116	.0088	.0409	.0829					
29		.0334	.0346	.0339	.0394	.0031	.0048	.0096	.0130	.0125	.0065	.0976	.1393					
31		.0587	.0592	.0596	.0627	.0057	.0073	.0132	.0160	.0167	.0233	.1303						
32		.0666	.0675	.0686	.0747	.0078	.0090	.0148	.0173	-0159	.0427							
34		.0776	.0791	.0823	.1067	.0121	.0133	.0185	.0183	.0149	.1444							
35		.0816	.0825	.0912	.1325	.0140	.0155	.0198	.0177	.0275								
37		.0842	.0803	.1264	.2137	.0206	.0179	.0189	.0201	.10485								
38		-0885	.1001	.1638	.2597	.0237	.0216	.0213	.0504	.1550								
40		.0915	.1236	.2111	.3047 .3484	.0266 .0296	.0229 .0235	.0297	.1074 .1606									
41		.1141	2142	.3113	.3927	.0312	.0226	.1048										
42		.1542	.2725	.3562	.4324	.0336	.0248	.1639										
44		.2915	.3660	.4353	.5047	.0350	.0547											
45		.3444	.3983 .4324	•4786 •5134	•5322 •5528	•0340 •0367	.1097											
47		.4057	.4701	.5442	•5645	.0444	- 10/ 4											
48		.4360	.5129	.5740	•5682	•0653												
50		.4983	.5681	•5782	.5255	.1774												
51		.5327	.5866	•5348	+6853													
53		.5803	.5240	.6980														
54		.5694	.5127	.9239														
56		.4930	.9606															
57		.7150																

Table II. Concluded

(c) $M_{\infty} = 2.86$

					C _p	for 2/H	n of -			
Orif	Loc	12	11	10.5	10	8	6	4	2	1
1	FP	.0009	.0031	.0023	.0040	.0030	.0020	.0034	.0010	.0027
2		0008	.0007	.0015	.0017	.0011	.0016	.0010	.0010	.0001
3	1 V	- 0009	0001	0004	.0021	.0011	.0011	.0020	.0012	.0007
					.0020	.0021	.0010	.0010		
5	FF	1309	1290	1287	.0069	.0074	.0090	.0137	.0144	.0065
6	¥	1325	1306	1305	.0059	.0076	.0088	.0146	.0123	.0059
<u> </u>	-	1209	1209	12/2	.0025	.0071	.009	.0135	.0104	.0122
62	ĸf	.6664	.7220	.7117	.0879	.0785	.0703	.0590	.0234	.0112
63		.5763	.6007	.6136	•0643	.0536	.0456	.0359	.0084	.0039
65		.8805	.9595	.9937	.0975	.0948	.0305	.0558	.0035	0023
66		1.0317	1.1860	1.1944	.1333	.1387	.1168	.0799	.0287	.0070
67		1.1824	1.3475	1.3945	.1744	.1955	.1634	.1148	.0766	.0298
69	RP	.0467	.0424	.0424	0015	0011	0002	.0028	.0041	.0056
70	ΙĨ	0159	0252	0264	0104	0102	0117	0095	0006	.0054
71		0573	0660	0671	0104	0100	0123	0097	0017	.0010
- 12	Ľ		0889	088/	0152	-+0036	0097	00//	0018	.0002
73	RF	.8363	.8409	.8561	.0809	.0508	.0325	.0210	.0173	.0031
74		1.3366	1.4297	1.4459	.1259	.0989	.0845	·0808	.0210	·0058
76		.8434	.8082	.8172	.0837	.0421	.0281	.0215	.0160	.0023
77		.6911	.7163	.7193	.0846	.0666	.0564	.0439	.0206	.0093
78		.6328	.6615	.6813	.0803	.0675	.0569	.0460	.0230	.0101
80	🕴	.7101	.7310	.7383	.0823	.1490	.1426	.1313	.1159	.1095
	ļ									
11	FL	1326	1304	1310	.0008	.0060	.0085	.0112	.0099	.0082
14		1365	1342	1340	0012	.0031	0002	.0018	0025	.0105
15		1344	1323	1322	0010	.0031	0003	.0048	0040	
16		1275	1255	1255	0025	.0038	.0027	.0061	.0023	
11/		1029	1007	1009	0029	.0035	-0035	.0073	.0292	
19	11	0893	0874	0872	0020	.0021	.0046	.0088		
20		0735	0715	0719	0024	.0009	.0063	.0065		
22		0412	0420	0416	.0004	.0041	.0072	.0042		
23		0302	0312	0192	0019	.0020	.0074	.0078		
25		0084	0096	0077	0002	.0028	.0078	.0625		
26		.0021	0003	.0015	0006	.0023	.0071	.0740		
27		.0160	.0072	.0082	0005	.0025	.0057			
29		.0230	.0218	.0212	.0005	.0030	.0041			
30		.0271	.0279	.0271	.0019	.0035	.0055			
31		.0303	.0305	.0293	.0038	.0050	.0268			
33		.0361	.0343	.0363	.0085	.0088	.0659			
34		.0357	.0354	.0417	.0111	.0105	.0905			
35		.0372	.0358	.0517	.0123	.0105				
37		.0342	.0459	.1150	.0140	.0110				
38		.0351	.0676	.1607	.0168	.0111				
39	1	.0345	.1160	.2044	.0188	.0156				
40		.0340	.2187	.2397	.0225	.0693				
42		.0486	.2625	.3161	.0222	.0993				
43		.1005	.2981	.3485	.0237					
44		.2399	.3634	.3834	.0244					
46	11	.2902	.3872	.4384	.0245					
47		.3213	.4161	.4468	.0282					
48		.3512	.4452	.44/4	.0418					
50		.3916	.4336	.4425	.1069					
51		.4157	.4118	.6663						
52		.4320	.4536	.8263						
54		.4141	.8066							
55		.3892								
56	🖌	.4421								
1 "	1	1.0309								

Table III. Pressure Coefficients for h = 2.0 in. and w = 2.5 in.

(a) $M_{\infty} = 1.50$

(b) $M_{\infty} = 2.16$

			(a) <i>M</i>	$\infty =$	1.50		÷.	11
	Γ.				C _p for	r l/h of	F _	· • • 5	*
Orif	Loc	6	5	4.8	4	3	2	1.5	1
1	FP	.0052	.0065	.0029	.0034	.0029	.0079	.0019	.0044
2		.0065	.0092	.0068	.0048	.0056	.0102	.0036	.0076
3	1	.0098	.0161	.0111	.0075	.0098	.0139	.0083	.0112
5	FF	.0273	.0430	.0410	.0450	.0810	.0917	.0496	.0434
6		.0259	.0400	.0389	.0386	.0748	.0746	.0398	.0287
7		.0269	.0405	-0376	.0380	.0691	+0617	.0222	.0108
9	۲	.0256	.0362	.0372	.0342	.0584	.0782	.0660	.0720
59	RF	.2719	.2678	.2670	.2520	.2682	.1896	.1269	.0647
60		.1883	.1685	.1729	.1539	.1397	.1023	.0632	.0476
61 62		.2747	.1487	.1538	.1263	.1037	.0388	.0212	.0153
64		.4459	.4381	.4252	.3807	.4232	.2619	.0944	.0406
65		.5055	4875	.4832	.4444	.4940	.3410	.1466	.0571
67	¥	.5705	.5516	.5529	.5253	.5813	.4461	.2487	.1093
69	עס	- 1049	- 1026	- 1010	- 1005	- 0932	- 0270	- 0211	- 0037
70	Ĩ	0665	0711	0675	0640	0696	0333	0119	0039
71		0204	0276	0257	0208	0303	0164	0006	0038
72	۲	0003	0099	0094	0032	0076	0041	.0005	0002
73	ĸF	.2830	.2954	. 3364	.3125	.3216	.2475	.1798	.1375
74		4576	-4739	.4118	.4676	.5860	-3885	.2032	.1002
76		.2904	.2702	.2488	.2804	.3373	.2498	.0685	.1442
77		.2649	.2645	.2599	.2441	.2399	.0682	.0342	.0051
78		.2488	.2270	.2134	.1655	.1972	-0506	.0234	.0049
80		.2664	.2709	.2763	.2415	.2587	.0708	.0017	.0001
81		.2754	.2677	.2579	.2687	.2592	.1345	.0633	.0311
82		.1995	.1910	.1868	.1778	.1737	.1078	.0500	.0333
84	¥	.2834	.2690	.2639	.2627	.2600	.1389	.0372	.0339
11	FL	.0254	.0274	.0302	.0291	.0452	.UShN	.0481	.0529
12	ī	.0224	.0182	.0217	.0193	.0280	.0471	.0320	.0290
13		.0215	.0148	.0170	.0147	.0196	.0404	.0088	.0219
15		.0163	.0097	.0094	.0095	.0143	.02/9	.0009	.0108
16		.0140	.0040	.0052	.0081	.0137	0070	.0095	.0519
17		.0145	.0067	.0032	.0081	.0189	0174	.0037	•0728
19		.0121	.0106	.0093	.0122	.0253	0248	.0280	.0/10
20		.0114	.0131	.0115	.0155	.0198	0241	.0836	
21		.0103	.0153	.0125	.0191	.0147	0128	.1699	
23		.0175	.0216	.0191	.0234	0006	-0192	.1455	
24		.0163	.0259	.0226	.0235	0047	.1456		
25		.0159	.0256	.0240	.0220	0087	.2511		
26 27		.0147	.0257	.0235	.0176	0098	.2314		
28		.0139	.0255	.0238	.0129	.0012			
29		.0129	.0223	.0230	.0085	.0138			
30		.0124	.0192	.0207	+0089	.0321			
32		.0112	.0164	.0121	.0092	.1470			
33		.0117	.0185	.0123	.0151	.3356			
34 35		.0121	.0160	.0115	.0167	.3316			
36		.0103	.0192	.0192	.0292				
37		.0110	.0196	.0219	.0394				
38		.0126	.0228	.0261	.0523				
40		.0143	.0317	.0382	.1434				
41		.0177	.0321	.0411	.2955				
42 43		+0193	.0370	+0467 -0558	.3267				
44		.0268	.0481	.0630					
45		.0309	.0542	.0779					
46		.0379	.0711	.1256					
48		.0541	.1732	.3836					
49		.0510	.3090						
50		.0562	.3281						
52		.0719							
53		.0759							
54		.0904							
56		.1712							
57	V	.3148							

		C _p for 2/h of -								
Orif	Loc	6	5	4.8	4	3	2	1.5	1	
1	FP	.0049	.0051	.0046	.0047	.0042	.0062	.0059	.0046	
2	1	.0051	.0051	.0042	.0049	.0042	.0054	.0064	.0043	
3	I V	.0034	.0042	.0029	+0031	.0031	-0037 0067	.0056	.0028	
5	FF	.0170	.0248	.0234	.0314	.0424	.0229	.0191	.0199	
7		.016/	.0235	.0224	.0295	.0397	.0210	.0175	.0100	
8	T	.0176	.0233	.0223	.0272	.0362	.0221	.0156	.0103	
59	RF	1349	1564	1445	1416	1313	.0474	.0313	.0327	
60	Ĩ	.0779	.0890	.0809	.0736	.0675	.0298	.0181	.0293	
61		.0578	.0664	.0559	.0473	.0288	.0156	.0063	.0098	
64		.2038	.2324	.2206	.2218	.1669	.0356	.0177	.0244	
65		.2510	.2920	.2836	.2895	.2387	.0534	.0279	.0412	
66 67	•	·3259	.3732	.3695	.3859	.3425	.0990	-0645	.0748	
			.4524	.4540	.4730	.4/10				
69	RP	0298	0270	0250	0250	.0105	.0056	.0029	.0109	
70		0362	0351	0326	0311	0061	0017	0026	.0043	
72	*	0267	0253	0235	0214	0157	0012	0029	.0002	
70					0.054	1000	0070	05/2	0710	
74	RF 	.3133	.3041	.2266	• 5258 • 5683	.3536	.0972	.0343	.0661	
75		.2218	.2296	.4881	.2933	.3237	.0956	.0357	.0679	
76		.1524	.1462	.3383	.3408	-2124	.1055	.0209	.0742	
78		.0702	.1568	.1460	.0781	.0392	.0127	.0090	.0104	
79		.0673	.0644	.0596	.0547	.0408	.0143	.0044	.0113	
80 81		.1326	.1343	.1053	.1300	.0620	.0182	.0032	.0256	
82		.0868	.1011	.0952	.0924	.0712	.0310	.0199	.0273	
83	¥	.0871	.1020	.0982	.0927	.0731	.0326	.0171	.0291	
		.114/	.1302	.1520	.1505	.0354	.0501	.0054	.0255	
11	FL	.0157	.0185	.0190	.0215	.0278	.0169	.0148	.0256	
13		.0139	.0146	.0147	.0154	.0216	.0151	.0123	.0091	
14		.0091	.0070	.0073	.0077	.0119	.0104	.0007	.0075	
15		.0067	.0036	.0044	.0051	.0094	.0054	0022	.0160	
17		.0043	.0010	.0024	.0042	.0136	0018	.0001	.0397	
18		.0037	.0006	.0036	.0053	.0173	0038	.0035	.0392	
19 20		.0032	.0021	.0037	.0068	-0209	0022	.0136	ŀ	
21		.0042	.0056	.0072	.0101	.0200	.0016	.0446		
22		.0039	.0127	.0129	.0159	.0174	.0120	.0402		
23		.0053	.0140	.0141	.0177	0011	.0246			
25		.0062	.0175	.0160	.0208	0078	.0667			
26		.0063	.0191	.0164	.0217	0122	.0594			
28		.0064	.0193	.0158	.0196	0076				
29		.0059	.0173	.0153	.0160	.0037				
30 31		.0068	.0171	.014/	.0104	.0212				
32		.0115	.0128	.0126	.0019	.0889				
33		.0115	.0138	.0125	.0024	.1728				
35		.0115	.0161	.0115	.0047	.1507				
36		.0117	.0189	.0142	.0112					
37		.0100	.0166	.0095	.0112					
39		.0090	.0154	.0077	.0328					
40		.0084	.0146	.0089	.0826				f	
41		.0082	.0133	.0156	.1932					
43		.0104	.0168	.0206						
44 45		.0118	.0209	.0275						
46		.0131	.0343	.0701						
47		.0116	.0501	.1472						
48 49		.0134	.1952	.21/4						
50		.0124	.1961							
51 52		.0137								
53		.0207								
54		.0290								
55 56		.0450							l	
57	V	.1709								
									- 1	

Table IV. Pressure Coefficients for h = 2.5 in. and w = 2.5 in.

(a)
$$M_{\infty} = 1.50$$

(b)
$$M_{\infty} = 2.16$$

				C _p for	ℓ/h of	-	
Orif	Loc	4.8	4	3	2	1.5	1
1	FP	.0030	.0058	.0057	.0064	.0052	.0070
2	Ĩ	.0049	.0085	.0082	.0093	.0080	.0098
3		.0075	.0112	+0105 0144	.0106	.0099	.0113
5	FF	.0300	.0424	.0734	.1238	.0474	.0430
7		.0288	.0388	.0571	.0725	.0265	.0176
8		.0272	.0370	.0502	.0849	.0237	.0128
10	🕴	.0276	.0323	.0476	.1213	.0421	.0241
58	RF	.2973	.2913	. 2883	.2431	.0810	.0775
59	Ĩ	.1939	.1708	.1349	.0915	.0288	.0230
60		.1951	.1662	.1233	.0368	.0089	.0137
63		.2827	.2473	.2299	.1033	.0129	.0139
62		.3328	.3090	.3047	.2096	.0254	.0224
65		.4701	.4521	.4857	.4876	.0741	.0856
66		.5211	.5129	.5520	.5730	.1282	.1584
6/	<u> </u>	.5604	.5397	.6088	.6465	.2143	.2623
69	RP	1095	1050	1034	0473	0095	0132
70		0751	0733	0714	0458	0168	0129
72	1	0096	0142	0092	0102	.0025	0033
73	RF	.3412	.4291	.2404	.2780	.2586	,1218
74	l'ī	.6081	.6596	.3231	.5215	.1877	.1517
75		.4768	.3892	-6224 5269	.5753	.1760	.1537
77		.3406	.2817	.2508	.1064	.0315	.0215
78		.3358	.3200	.2591	.1490	.0428	.0184
81		.2633	.3011	.3390	.0935	.0221	.0124
82		.2357	.1758	.1876	.0394	.0134	.0082
83	🕴	.2217	.2050	.1815	.0359	.0125	.0091
L	PI	02/0	0206	0.26.2	0004	0/00	04.22
12	Ĩ	.0233	.0220	.0209	.0779	.0400	.0200
13		.0208	.0152	.0133	.0672	.0262	.0099
15	11	.0191	.0139	.0083	.0480	0004	.0138
16		.0175	.0117	.0080	.0138	0102	.0229
18		.0169	.0135	.0038	0201	0161	.0439
19		.0169	.0137	.0119	0269	0113	.1140
20		.0161	.0142	.0165	0352	0049	.0956
22		.0131	.0202	.0290	0383	.0510	
23		.0143	.0204	-0298	0328	.0978	Ì
25		.0124	.0211	.0228	0003	.1020	
26		.0135	-0224	.0172	.0476		
28		.0138	.0233	.0122	.2186		
29		.0137	.0224	.0059	.2935		
30		.0160	.0207 .0188	.0091	.2651		
32		.0161	.0153	.0240			
33		.0162	.0159	.0403			
35		.0153	.0124	.0668			
36		.0123	.0144	.0868			
38		.0125	.0158	.11/8			
39		.0124	.0263	.3215			
40		.0119	.0336 .0396	.3268			
42		.0158	.0522				
43		.0208	.0706				
45		.0301	.0955				
46		.0362	.1159				
48		.04/6	.2208				
49		.0650	.3131				
50		.0793	. 3224				
52		.1043					
53		.1121					
55		.1522					
56		.2088					
<u> </u>	I.'_	1.3269					_

				C _p fo	or I/h a	f -	
Orif	Loc	4.8	4	3	2	1.5	1
	עט	0030	0036	0036	.0034	.0028	.0033
2	ï	.0025	.0034	.0032	.0032	.0026	.0029
3		.0018	.0026	.0024	-0029	.0024	.0026
4		.0036	.0045	.0041	.0040	.0041	.0033
5	FF	.0219	.0313	.0406	.0247	.0178	.0166
6		.0217	·0296	.0379	.0217	.0156	.0131
8		.0215	.0279	.0345	.0212	.0092	.0008
9		.0215	.0280	.0314	.0243	.0170	.0094
10		.0207	.0256	.0291	.0214	.0193	.0255
58	RF	.1690	.1705	.1424	.0549	.0341	.0329
59		.0744	.0683	.0430	.0226	.0161	.0209
61		.0688	.0588	.0240	.0217	.0042	.0031
63		.1266	.1169	.0706	.0160	.0070	.0091
62		.1853	.1/53	.1229	.0283	.0106	.0082
65		.3692	.4038	.3350	.0963	.0406	.0504
66		.4579	.5096	.4457	.1568	.0773	.0853
0/	<u> </u>	.5290	• 5645	.3613	.2357	.1445	.14.14
69	RP	0283	0310	.0119	.0072	.0041	.0081
70		0339	0366	0120	0040	.0018	.0001
72	🕴 -	0263	0304	0216	0075	.0007	0032
7.2	<u>,</u>	2062	2622	217/	1505	0757	0720
74		.3816	.3032	.4186	.1558	.0744	.0834
75		.4540	•4435	.4781	.1557	.0739	.0861
76		.2352	.3560	.3244	.1529	.0750	.0812
78		.1665	.1557	.0921	.0331	.0139	.0132
79		.1813	.1957	.0803	.0138	.0061	.0078
81		.1170	.1226	.0680	.0260	.0099	.0049
83		.0805	.0756	.0275	.0208	.0081	.0020
84	I.	.1041	.1123	.0493	.0253	.0097	.0052
11	FL	.0189	.0225	.0259	.0180	.0153	.0193
12		.0169	.0184	.0196	.0161	.0135	.0133
14		.0143	.0137	.0145	.0146	.0052	0003
15		.0102	.0099	.0087	.0112	0009	.0010
16		.0095	.0092	.0077	.0062	0051	.0106
18	11	.0090	.0084	.0117	.0035	0082	.0405
19		.0089	.0089	.0141	0013	0061	.0460
20		.0098	.0111	.0193	0030	.0060	.0420
22		.0117	.0140	.0264	.0078	.0235	
23		.0115	.0140	.0267	.0085	.0375	
25		.0118	.0136	.0216	.0077	.0404	
26		.0121	.0138	.0142	.0187		
28		.0127	.0132	0017	.0503		
29		.0125	.0145	0081	.0784		
30		.0121	.0159	0113	.0700		:
32		.0124	.0201	0110			
33		.0105	.0207	0033			
34		.0094	.0190	.0055			
36		.0091	.0158	.0198			
37		.0096	.0031	.0644			
38		.0113	0034	.1155			
40		.0167	0006	.1658			
41		.0187	.0026				
42		.0183	.0114				
44		.0155	.0291				
45		.0110	.0360				
47		.0039	.0679				
48		.0042	.1321				
49		.0063	.2092				
51		.0172	. 2000				
52		.0301					
54		.0524					
55		.0731					
57	1	.1218					
1	1 .						

			M=	1.50		M _{ap} = 2.16			
Orif	Loc		C _p fo	rwof	-		C _p for	w of -	
		2.0	1.5	1.0	0.5	2.0	1.5	1.0	0.5
1	FP	.0022	.0005	.0010	.0043	.0008	0000	.0050	.0065
2		.0052	.0043	.0049	.0063	.0002	0008	.0040	.0057
3		.0119	.0093	.0132	.0127	.0003	0010	.0041	.0052
4	V	.0130	.0121	.0127	.0145	.0025	.0013	.0063	.0075
5	FF	0225	1287	1390	1071	0176	1481	1358	0897
6	FF	0251	1288	1384	1084	0193	1437	1338	0915
64	RF	,2585	.4853	.4926	.3986	.1693	.4863	.4835	.3411
65		.2283	.4595	.4629	.3910	.1494	.4352	.4396	.3316
66		.2276	.4792	.5092	.4599	.1475	•4664	•5334	.3855
67	i L	.2521	.5160	.5910	.5794	.1723	.5710	.7298	.4985
68	, V	.2721	. 5099	.6302	.6974	.2131	•6856	.9109	.6278
69	RP	0929	2263	2884	1641	0202	0577	1335	1124
70		0739	1607	0986	0517	0392	1098	1213	0418
71		0367	0631	0501	0230	0331	1034	0464	0262
72	T	0192	0382	0285	0078	0236	0509	0249	0168
74	RF	.2688	.5529			.1688	.4992		
75	RF	.2782	.5221			.1654	.4690		
11	FL	0338	1478	1489	1141	0251	1559	1453	0920
13		0404	1515	1504	1172	0307	1616	1474	0950
14		0417	1517	1527	1156	0321	1462	1347	0894
15		0435	1478	1443	1070	0334	1184	1049	0731
16		0435	1306	1167	0788	0324	0874	0702	0480
17		0406	1009	0751	0341	0297	0559	0357	0181
18		0335	0620	0298	.0182	0238	0271	0049	.0105
20		0102	.0191	.0529	.0859	0067	0210	.0407	.0544
22		.0293	.09/8	.1072	.1087	.0186	-0686	.0732	.0776
23		.0400	.1319	1444	1009	.0291	10951	.0806	.07/8
24		0835	1020	•199/	1070	0517	1590	1164	.0702
25		1003	2361	-1/24	1306	.051/	1013	140	1002
27		.1165	. 2669	2649	1919	0727	2265	.1007	1531
28		.1321	3025	3125	2442	0821	2635	•2320 2847	2072
29		.1448	. 3401	. 3535	2897	08021	3010	3751	2662
32		.1672	.4039	4159	3395		3547	3751	2846
33	I L	.1851	.4051	.4090	.3478	.1167	.3598	.3492	2040
34		.2748	.5192	.5223	.4306	.1812	.5320	.5452	.3751

(a) h = 0.5 in., l/h = 12

Table V. Pressure Coefficients for Range of Cavity Widths

(b) h = 0.5 in., l/h = 6

			M =	1.50		M ₆₀ = 2.16					
Orif	Loc		C _p for	rwof–			C _p for	w of -			
		2.0	1.5	1.0	0.5	2.0	1.5	1.0	0.5		
1	FP	.0029	.0014	.0014	.0011	.0035	.0036	.0005	.0027		
2		.0057	.0041	.0053	.0048	.0029	.0025	0003	.0018		
3		.0120	.0117	.0115	.0109	.0035	.0031	0007	.0025		
4	V	.0136	.0134	.0127	.0115	.0058	.0055	.0022	.0054		
5	FF	.0282	.0260	.0230	.0200	.0187	.0159	.0136	.0141		
6	FF	.0273	.0249	.0242	.0191	.0185	.0156	.0133	.0132		
64	RF	.1700	.1750	.1345	.1054	.1137	.1218	.0802	.0619		
65		.1417	.1348	.0874	.0671	.0949	.0950	.0517	.0424		
66		.1302	.1390	.0974	.0657	.0839	.0959	.0589	.0357		
67	ΙL	.1606	.1914	.1714	.1200	.1163	.1424	.1198	.0695		
68	T	.2238	.2626	.2914	.2321	.1986	.2125	.2384	.1594		
69	КР	0278	0394	0135	.0005	0050	0218	0118	.0003		
70		0184	0163	0005	.0045	0100	0130	0023	.0049		
71		0026	0023	.0064	.0077	0044	0031	.0016	.0069		
72	V	.0028	.0015	.0068	.0072	0001	.0011	.0029	.0071		
74	RF	.1410	.0837			.0842	.0449				
75	RF	.1336	.0815			.0781	.0451				
11	FL	.0238	.0158	.0204	.0177	.0161	.0091	.0115	.0123		
13		.0229	.0156	.0177	.0150	.0140	.0091	.0093	.0102		
14		.0220	.0171	.0164	.0132	.0124	.0123	.0077	.0088		
15		.0205	.0164	.0133	.0092	.0107	.0125	.0064	.0071		
16		.0195	.0169	.0117	.0068	.0085	.0124	.0041	.0060		
17		.0195	.0201	.0083	.0068	.0079	.0139	.0017	.0050		
18		.0231	.0232	.0055	.0051	.0081	.0161	0003	.0021		
20		.0386	.0391	.0057	0007	.0190	.0266	0027	0043		
22	1	.1867	.2005	.1632	.1225	.1242	.1403	.1008	.0730		

Table V. Concluded

(c) h = 1.0 in., l/h = 12

(d) h = 1.0 in., l/h = 6

			M ₆₀ = 1	.50			M =	2.16	
Orif	Loc		C _p for	w of -			C _D for	w of -	
, , , , , , , , , , , , , , , , , , ,	[~~]	20	1 5	1.0		2.0	<u> </u>	1 0	0.5
		2.0	1.5	1.0	U.0	2.0	1.5	1.0	0.5
1 1	FP	.0041	.0058	.0055	.0057	.0005	.0033	.0013	.0004
2		.0079	.0108	.0073	.0063	.0009	.0046	.0003	0011
3	🖌	.0102	.0120	.0120	.0114	0006	.0028	.0009	0002
4		.0084	.0112	.0099	.0119	.0023	+0053	.0036	.0014
5	FF	1566	1355	1168	0798	-,1546	1316	-,1055	0737
6	111	1553	1349	1173	0798	1543	1314	1064	0741
7	🖠	1558	1361	1162	0795	1492	1285	1062	0738
	\vdash							·	
62	RF	.6075	.5619	.5127	.4083	.6610	5800	.4896	3434
63		.0050	.5577	.5354	.4234	-6242	.2519	.4921	.3441
64		.7434	.6840	6797	.5//2	.0131	-0838 -7702	• 56 55 5 - 650 7	-9120
66		.7968	.7303	.7296	.7695	1.0656	.8902	.7154	.7998
67	🕴	.8149	.7556	.7674	.8125	1.1958	.9960	.7893	.8997
<u>├</u>	<u> </u>	<u>├</u>							
69	RP	3705	3954	3645	2004	1100	1549	1929	1568
70		3342	2553	1768	0678	1455	1779	1375	0628
11	 ∦	- 1052	1302	0709	0417	1679	1493	0/12	05/1
· "	للكل	.1053	.0/2/	.0399	+0192		.0902	-0410	.0403
74	RF	.7792	.7374			1.0190	.8712		
75	11	.7180	.6718			.9475	.7726		
78		.5949	.5758		l	.6434	.6060		
79	L	.5768	.5751			.6255	.6025		
11	FI	- 1644	1407	-,1182	0815	-,1500	1375	1073	0750
13	l'í I	1717	1457	1200	0834	1655	1424	1091	0759
14	! 1	1704	1465	1200	0828	1652	- 1413	1099	0770
15		1694	1474	1205	0854	1649	1408	1105	0783
16	11	- 1683	1465	1224	0869	1654	- 1408	1115	0800
17		- 1705	1466	1227	0888	1641	1408	1110	0775
19		1700	1457	1212	0847	1480	1316	1054	0725
20	[]	1680	- 1434	1191	0807	1343	1219	1002	0667
22		1472	- 1255	1020	0591	0994	0908	0795	0523
23		1314	1101	0898	0469	0798	0727	0666	0416
24	11	1112	0923	0734	0306	0594	0533	0508	0275
25		0865	0684	0527	0087	0382	0332	0327	0132
20		0395	0421	0273	.0121	0194	0135	0145	-0027
28	11	0024	.0130	.0004	.0574	.0162	.0031	.0030	,0340
29		.0260	.0412	.0545	.0797	.0308	.0401	.0389	.0486
30	11	.0525	.0650	.0780	.0982	.0436	.0550	.0551	.0606
31	11	.0772	.0883	.0994	.1100	.0578	.0712	.0683	.0764
32		.0986	1097	.1175	.1169	•0682	-0826	.0804	.0849
1 33		.1150	.1214	.1259	.1181	.0767	.0915	•0896	.0910
35	11	1408	-1402	+1312	.1144	+U858 - 0412	.1054	.1036	.0937
36		.1531	.1453	.1332	.1087	.0990	.1107	.1071	.0929
37	11	.1608	.1471	.1303	.0994	.1026	.1135	.1091	.0901
38	11	.1674	.1479	.1253	.0916	.1051	.1136	.1085	.0853
39	11	.1793	.1499	1201	.0848	1075	.1138	.1076	.0800
40	11	1959	1605	.1143	.0767	.1141	•1135	.1033	.0725
41	11	.2453	.1911	.1751	.0720	1605	-1212	.0990	.0604
43		.2766	.2194	.1488	.0751	.2167	.1429	1054	.0557
44		.3127	.2566	.1861	.0894	.2690	.1956	.1258	.0578
45	11	.3511	.2968	.2295	.1151	.3155	.2532	.1721	•0792
46		.3833	3303	.2773	.1509	•3449	.3047	.2321	.1201
4/	11	-4125	- 3657	.3154	.1942	3657	•3372	•2809	•1666
40	11	.4405	. 4194	מסכני רניער	.2750		, 3094 , 3096	- 3720	.2009
50		.4733	.4424	.4077	.3072	.4403	.4201	.3699	.2606
51		.4882	-4581	.4261	.3374	.4628	.4441	.3868	.2805
52	11	.5016	4821	.4338	.3632	.4840	. 4636	.3979	-2944
53		.5098	.4945	• 4404	.3779	.5012	.4796	.4047	•3026
54		-5205	5035	.4520	.3823	.5107	.4867	.4192	.3043
55	11	-5322	.5052	-4592	.3841	.4955	.4786	+4253	.3032
57	♥	5443	.4936	•4662	• 3843 3000	4620	.46/3	.42/9	.3108
<u>,,</u>	1'	1.1443	+ 3148	.+000	10101	1.2010	. 3134	• 44 30	.34/0

			M., =	1.50		$M_{gg} = 2.16$				
Orif	Loc		C _p for	w of -		(C _p for w	1 of -		
		2.0	1.5	1.0	0.5	2.0	1.5	1.0	0.5	
,	FP	.0038	.0030	.0053	.0039	0004	.0001	.0031	,0007	
2	î	.0047	.0060	.0046	.0037	0019	0013	.0024	0011	
3		.0109	.0104	.0114	.0111	0008	0009	.0030	0003	
4	V	.0109	.0120	.0111	.0100	.0017	.0013	.0050	.0022	
5	FF	.0255	.0246	.0167	.0143	.0122	.0094	.0136	.0057	
6		.0256	.0253	.0165	.0145	.0122	.0091	.0130	.0061	
7	Į	.0259	.0235	.0170	.0147	.0121	.0088	.0131	.0054	
62	RF	.2395	.2116	.1846	.1316	.0843	.0762	.0661	.0585	
63	1	.1590	.1358	.1259	.0868	.0418	.0377	•0256	.0232	
64	1	.2333	.2093	.1764	.1227	.0665	.0494	.0280	.0243	
65		.3072	.2820	.2256	.1700	.0996	.0703	•0458	.0436	
66	11	.4069	.3893	.3074	.2426	.1599	.1177	.0871	.0890	
67	V	.5020	.5040	•4041	.3399	.2415	.1846	•1500	.1568	
69	RP	0639	0662	0544	0018	0198	0250	0067	0033	
47	1	0375	0111	0041	.0094	0179	0177	.0027	.0014	
71		0070	.0069	.0044	.0151	0105	0081	•0062	.0031	
72	♥	.0047	.0106	.0095	.0179	0062	0040	.0068	.0031	
74	RF	.2489	.1843			,1043	.1428			
75	11	.2312	.1782			.0687	.0383			
78		.2136	.2408			.0816	.0976			
79	I V	.2055	.2262			.0745	.0747			
11	FL	.0246	.0200	.0145	.0150	.0113	.0054	.0136	.0056	
13	í ī	.0200	.0121	.0099	.0146	.0096	0012	.0127	.0053	
14		.0183	.0125	.0109	.0141	.0090	0018	.0122	•0050	
15		.0173	.0154	.0095	.0136	.0082	.0000	.0115	.0052	
16	11	.0147	.0158	.0096	.0139	.0074	.0021	.0114	.0049	
17		.0146	.0153	.0062	.0137	.0067	.0040	.0108	.0040	
18		.0136	.0143	.0069	.0129	.0061	.0046	•0099	.0036	
19		.0129	.0132	+0052	.0116	.0055	.0042	.0090	.0026	
20	11	.0120	.0104	.0035	+0094	.0046	.0041	.0079	0022	
22		.0166	.0119	.0060	.0084	.0042	.0040	.0093	0051	
23	11	01/5	+0121	.0002	.0000	.0027	.0036	0053	00/1	
24	11	0145	0116	0100	.0111	- 0005	0012	.0000	0030	
25		01/2	0124	.0100	0092	- 0000	- 00012	- 0025	- 0016	
27		1 -0156	.0131	.0126	.0049	0029	0021	0073	0035	
28		.0186	.0167	.0212	- 0007	- 0081	- 0023	- 0110	- 0067	
29		.0179	-0188	.0212	0031	- 0105	- 0020	- 0138	- 0137	
30		.0206	.0219	.0367	0001	0112	0014	0147	0207	
31		.0352	0250	.0547	.0176	0050	.0040	0088	0144	
32	11	.0663	.0455	.0692	.0494	.0090	.0137	.0108	.0067	
33	11	.1737	.1346	.1265	.1064	.0612	.0538	.0596	.0496	
34	IV.	.3209	.2947	.2446	.1533	.1181	.1057	.0871	.0719	



Figure 1. Sketches of cavity flow field models.



Figure 2. Model description. Linear dimensions are in inches.



(b) Photograph of model installed in tunnel. h = 1.0 in.; w = 2.5 in.

Figure 2. Continued.



Figure 2. Concluded.

(c) Photograph of typical block inserts for varying cavity width. h = 1.0 in.



Figure 3. Pressure orifice locations. Linear dimensions are in inches.



(a) l/h = 24.0.



(b) l/h = 16.0.



(c) l/h = 11.6.

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Figure 4. Schlieren photographs of cavity flow field for h = 0.5 in. and $M_{\infty} = 2.86$.



(d) l/h = 11.2.



(e) l/h = 8.0.



(f) l/h = 2.0.

Figure 4. Concluded.

L-87-570



(a) l/h = 12.0.



(b) l/h = 10.5.



(c) l/h = 10.0.

Figure 5. Schlieren photographs of cavity flow field for h = 1.0 in. and $M_{\infty} = 2.86$.

L-87-571



(d) l/h = 8.0.



(e) l/h = 6.0.



(f) l/h = 2.0.

L-87-572

Figure 5. Concluded.

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Figure 7. Effect of cavity length-to-depth ratio on cavity centerline pressure distributions. w = 2.5 in.



(b) $M_{\infty} = 2.16$; h = 0.5 in.







(d) $M_{\infty} = 1.50; h = 1.0$ in.



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(b) h = 1.0 in.

Figure 8. Continued.



(c) h = 2.0 in.

(d) h = 2.5 in.



Figure 9. Effect of cavity depth on cavity centerline pressure distributions. w = 2.5 in.















(g) $M_{\infty} = 2.16; l/h = 4.$







Figure 9. Concluded.











Figure 10. Concluded.





Figure 11. Correlation of cavity pressure distributions based on l/h and w/h. w/h = 1.0.







(d) $M_{\infty} = 2.16$; l/h = 6. Figure 11. Concluded.



Figure 12. Effect of cavity width on cavity pressure distributions. h = 0.5 in.









Figure 13. Effect of cavity width on cavity pressure distributions. h = 1.0 in.

(a) $M_{\infty} = 1.50; l/h = 12.$









Figure 14. Cavity vapor screen photographs from reference 8. $M_{\infty} = 2.16$.



Figure 15. Lateral pressure distributions measured on cavity rear face. w = 2.5 in.

(a) h = 0.5 in.; $y_2/h = 0.5$.


Figure 15. Continued.

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Figure 15. Continued.



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Figure 15. Concluded.

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Figure 16. Effect of cavity width on cavity rear face pressure distributions.

(a) h = 0.5 in.; $y_2/h = 0.5$.

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i.



Figure 16. Concluded.