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Explicit Robust Schemes for Implementation of a Class of Principal Value-Based Constitutive Models: Symbolic and Numeric Implementation

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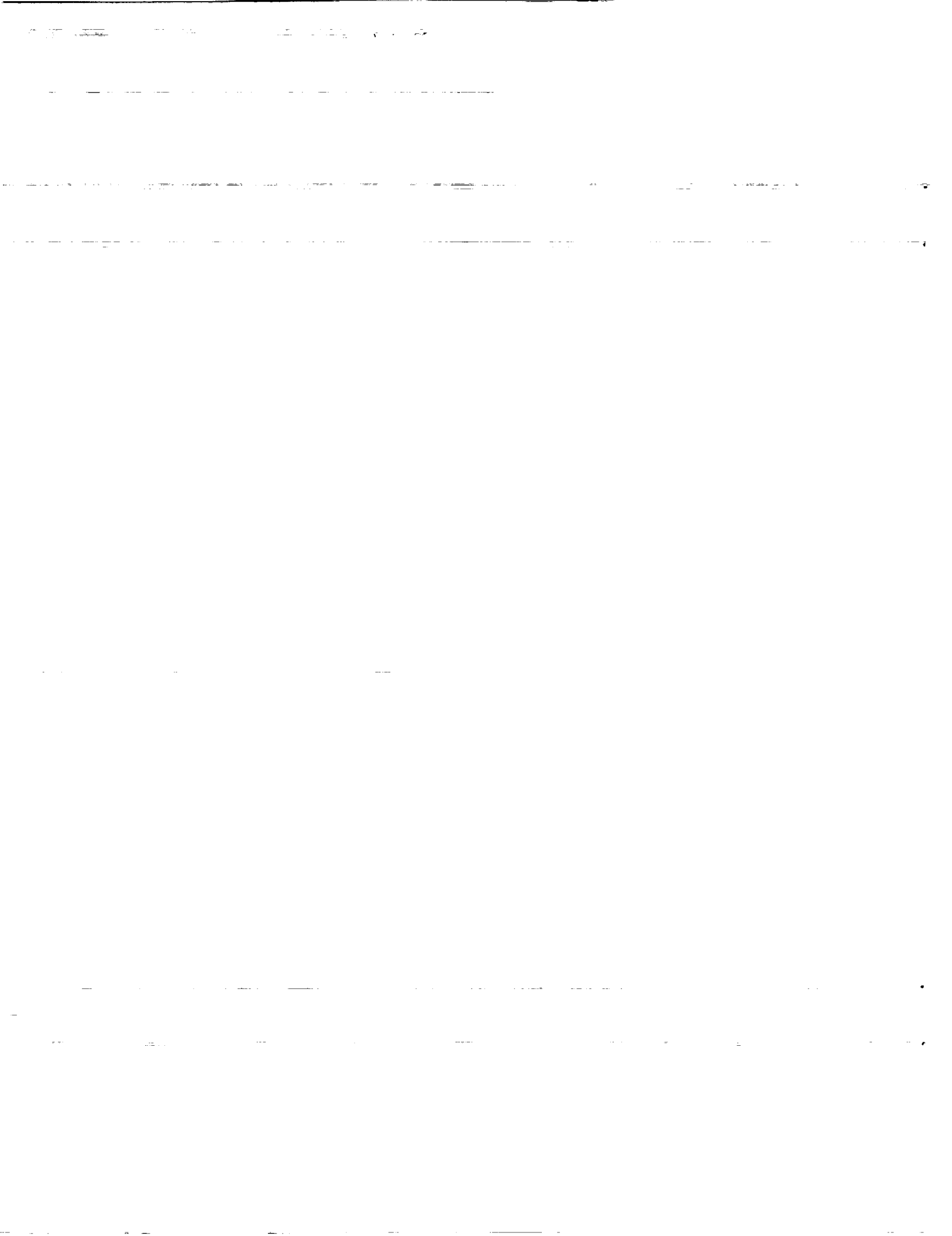
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Explicit Robust Schemes for Implementation of a Class of Principal Value-Based Constitutive Models: Symbolic and Numeric Implementation

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Abstract

The issue of developing effective and robust schemes to implement a class of the Ogden-type hyperelastic constitutive models is addressed. To this end, special purpose functions (running under MACSYMA) are developed for the symbolic derivation, evaluation, and automatic FORTRAN code generation of explicit expressions for the corresponding stress function and material tangent stiffness tensors. These explicit forms are valid over the entire deformation range, since the singularities resulting from repeated principal-stretch values have been theoretically removed. The required computational algorithms are outlined, and the resulting FORTRAN computer code is presented.

1 Introduction

To a great extent, constitutive models of the so-called generalized Rivlin-Mooney type [1,2] (i.e., with the stored strain energy density written as a polynomial function in terms of the deformation invariants) have dominated the phenomenological theory of isotropic hyperelasticity [1-6]. Such models dominate the related computational literature on finite-strain elasticity [7-9] as well. Recently

though, alternative representations in terms of the principal stretches have become increasingly popular in nonlinear finite element analyses [6,8,10]. However, from the viewpoint of numerical implementation, the use of these models presents a number of unique and difficult problems, which do not arise in classical representations using the strain invariants. The main difficulty is that (in addition to being reasonably complicated functions of the strain components) taken separately, the main constituents of the deformation tensor (i.e., principal values and associated eigenvectors) are, in general, not uniquely defined and continuously differentiable functions. A careful consideration is thus called for in implementing constitutive models formulated in terms of these principal-strain measures; this was the main problem addressed by Saleeb and Arnold [11]. They bypassed the difficulty entirely by resorting to explicit derivations of appropriate forms of the material tangent-stiffness matrices, which are valid for the entire deformation range. The explicit expressions they developed [11] were for two specific forms of the Ogden-type, strain-energy functions, which actually encompass many of the popular representations currently in use for rubber materials. Results were obtained by simply applying a systematic limiting procedure for one type of tensor-valued function and its spectral representation.

Symbolic computation specializes in exact computation with numbers, formulas, vectors, matrices, equations and the like. Numerical computation, on the other hand, uses floating-point numbers to compute approximate solutions to problems of practical interest. The two approaches are complementary and, when combined into an integrated form, can be very powerful in engineering applications. In particular, application of symbolic manipulation can provide significant incentive for the development of new constitutive theories and their applications, for example, finite element. Recently, a problem-oriented, self-contained, symbolic expert system, named SDICE (see [12-13]), was developed; it is capable of efficiently deriving, in analytical form, potential based constitutive models whose representations are in terms of the *classical* invariant formulation [14-15]. In addition, the FORTRAN code associated with the resulting analytical expressions can be automatically generated.

The objective of the present paper is to discuss three special purpose functions (SDIFF, SDIFFEV, and TEMPLATE) running under DOE MACSYMA [16]. These three functions have been developed to allow the derivation and automatic FORTRAN code generation of *alternative* potential based constitutive models composed of principal values and their associated eigenvectors, as discussed in reference 11. All three functions are written at the MACSYMA command level. In the future, these functions will be integrated into the collection of special purpose functions known as SDICE. This paper begins by reviewing *highlights* of the previous theoretical development and discussing the associated computer algorithm for the derivation of the explicit expressions for the second Piola Kirchhoff stress tensor S_{ij} and the material moduli tensor D_{ijkl} . The paper concludes with the evaluation of a separable strain energy function, similar to that discussed in reference 11, and its associated FORTRAN

source code generation.

2 Background

The theoretical development of *singularity-free* representations for principal value-based constitutive models has been discussed at length in reference 11. For brevity, we will confine our discussion, for illustrative purposes, to hyperelastic isotropic materials whose strain energy function W can be taken to have the following separable functional dependence:

$$W = W(\lambda_i) = \sum_{n=1}^p \bar{a}_n (\lambda_1^{\alpha_n} + \lambda_2^{\alpha_n} + \lambda_3^{\alpha_n}) \quad (1)$$

where λ_i represents the principal values of the right Cauchy-Green deformation tensor C_{ij} . Denoting n_i ($i = 1, 2$, or 3) to be the associated eigenvectors of C_{ij} , we can define

$$C_{ij} = \sum_{l=1}^3 \lambda_{(l)} N_{ij}^{(l)} \quad (2)$$

where $N_{ij}^{(l)}$ is defined as

$$N_{ij}^{(l)} = n_i^l n_j^l \quad (3)$$

and is often referred to as the (orthogonal) *eigenprojection* operator related to the associated eigenvectors of C_{ij} .

Equation (2) is valid for the case when all three eigenvalues, λ_i , are distinct. However, for the case when two eigenvalues are the same (i.e., double coalescence $\lambda_1 \neq \lambda_2 = \lambda_3 = \lambda$) we have

$$C_{ij} = (\lambda_1 - \lambda) N_{ij}^{(1)} + \lambda \delta_{ij} \quad (4)$$

And for the case of triple coalescence ($\lambda_1 = \lambda_2 = \lambda_3 = \lambda$), we have

$$C_{ij} = \lambda \sum_{l=1}^3 N_{ij}^{(l)} = \lambda \delta_{ij}. \quad (5)$$

Similarly, by manipulating equations (2) and (4), we can obtain explicit expressions for $N_{ij}^{(r)}$ in terms of C_{ij} :

$$N_{ij}^{(r)} = \frac{1}{(\lambda_r - \lambda_s)(\lambda_s - \lambda_t)} [(C_{ij} - \lambda_s \delta_{ij})(C_{ij} - \lambda_t \delta_{ij})] \quad (6)$$

and

$$N_{ij}^{(r)} = \frac{1}{(\lambda_r - \lambda)}(C_{ij} - \lambda \delta_{ij}) \quad (7)$$

In the preceding equations, the r, s , and t are any cyclic permutation of (1, 2, or 3). These definitions, equations (2) through (7), will prove very useful in obtaining the pertinent singularity-free directional derivatives of both the strain-energy potential function W and the stress function $S_{ij} = S_{ij}(C_{ij})$.

The explicit singularity-free expressions for the second Piola Kirchhoff stress tensor $S_{ij}(C_{ij})$ are defined as

$$S_{ij} = 2 \frac{\partial W}{\partial C_{ij}} \equiv S_{ij}(C_{ij}) \quad (8)$$

and those for the material moduli tensor $D_{ijkl}(C_{ij})$ are obtained by applying the directional derivative formula to S_{ij} , that is

$$D_{ijkl} = 2 \frac{\partial S_{ij}}{\partial C_{kl}} = 4 \frac{\partial^2 W}{\partial C_{ij} \partial C_{kl}} \equiv D_{ijkl}(C_{ij}) \quad (9)$$

As a result, the explicit expressions of the functional dependence of tensors S_{ij} and D_{ijkl} on C_{ij} can be obtained directly for the following three cases: 1) all three eigenvalues are distinct; 2) a single singularity ($\lambda_1 \neq \lambda_2 = \lambda_3 = \lambda$, i.e., double coalescence) is present; or 3) a double singularity ($\lambda_1 \neq \lambda_2 = \lambda_3 = \lambda$, i.e., triple coalescence) is present.

3 Computer Algorithm

The objective of the present study was to construct three special purpose functions (SDIFF, SDIFFEV, and TEMPLATE) written at the MACSYMA command level that can, respectively,

- (1) Derive explicit expressions for the stress tensor S_{ij} (eqs. (8)) and material tensor D_{ijkl} (eqs. (9)) given three, one, or no distinct eigenvalues
- (2) Evaluate symbolically the expressions generated by SDIFF for a given strain-energy function W
- (3) Evaluate the expressions generated by SDIFF and use the built-in MACSYMA function `gentran` to automatically generate the associated FORTRAN code needed to evaluate the expressions numerically for a given function, W .

These special purpose functions contain a list of built-in MACSYMA instructions (`factor`, `expand`, `ev`, `ratsubst`, `diff`, `limit`, and `for-loops`, to name a few) arranged in a specific algorithmic order. Each function, then, can be thought of as a macro command.

3.1 SDIFF(case)

Issuing the command SDIFF invokes the following algorithm (consisting of 15 steps) for automatic derivation of S_{ij} and D_{ijkl} . In this context, $\text{case} \equiv 1$ indicates that all three eigenvalues are distinct; $\text{case} \equiv 2$ indicates that only one is distinct; and $\text{case} \equiv 3$, that none are distinct.

To obtain S_{ij} ,

- (1) Differentiate W with respect to C_{ij} (see eq. (8))

$$S_{ij} = \sum_{l=1}^3 2 \frac{\partial W}{\partial \lambda_{(l)}} \frac{\partial \lambda_{(l)}}{\partial C_{ij}} \quad (10)$$

- (2) Apply the special directional derivative rules obtained from equation (2), that is,

$$N_{ij}^{(l)} = \frac{\partial \lambda_{(l)}}{\partial C_{ij}} \quad (11)$$

whose value is given in equation (6).

- (3) Obtain typical scalar derivatives by using the built-in **diff** command:

$$s(\lambda_{(l)}) = 2 \frac{\partial W}{\partial \lambda_{(l)}} \quad (12)$$

- (4) Multiply the results $s(\lambda_{(l)})$ and $N_{ij}^{(l)}$, then sum and factor out coefficients of like terms (i.e., $C_{ik}C_{kj}$, C_{ij} , and δ_{ij}), thereby obtaining the functional dependence of S_{ij} on C_{ij} . In the case of three distinct eigenvalues,

$$S_{ij} = aC_{ik}C_{kj} + bC_{ij} + c\delta_{ij} \quad (13)$$

where δ_{ij} is the second order identity tensor and

$$a = -m[s(\lambda_1)(\lambda_2 - \lambda_3) + s(\lambda_2)(\lambda_3 - \lambda_1) + s(\lambda_3)(\lambda_1 - \lambda_2)] \quad (14)$$

$$b = m[s(\lambda_1)(\lambda_2^2 - \lambda_3^2) + s(\lambda_2)(\lambda_3^2 - \lambda_1^2) + s(\lambda_3)(\lambda_1^2 - \lambda_2^2)] \quad (15)$$

$$c = -m[s(\lambda_1)\lambda_2\lambda_3(\lambda_2 - \lambda_3) + s(\lambda_2)\lambda_3\lambda_1(\lambda_3 - \lambda_1) + s(\lambda_3)\lambda_1\lambda_2(\lambda_1 - \lambda_2)] \quad (16)$$

and

$$m = \frac{1}{(\lambda_1 - \lambda_2)(\lambda_2 - \lambda_3)(\lambda_3 - \lambda_1)} \quad (17)$$

To obtain D_{ijkl} :

(5) Differentiate S_{ij} with respect to C_{kl} (see eq. (9)):

$$\begin{aligned} D_{ijkl} = & 2\left\{a\left[\frac{1}{2}(\delta_{ik}\delta_{ml} + \delta_{il}\delta_{mk})C_{mj} + \frac{1}{2}C_{im}(\delta_{jk}\delta_{ml} + \delta_{jl}\delta_{mk})\right]\right. \\ & + \sum_{r=1}^3 \frac{\partial a}{\partial \lambda_r} \frac{\partial \lambda_r}{\partial C_{kl}} C_{im} C_{mj} + b\left[\frac{1}{2}(\delta_{ik}\delta_{jl} + \delta_{il}\delta_{jk})\right] \\ & \left. + \sum_{r=1}^3 \frac{\partial b}{\partial \lambda_r} \frac{\partial \lambda_r}{\partial C_{kl}} C_{ij} + \sum_{r=1}^3 \frac{\partial c}{\partial \lambda_r} \frac{\partial \lambda_r}{\partial C_{kl}} \delta_{ij}\right\} \end{aligned} \quad (18)$$

(6) Apply the special directional derivative rule

$$N_{ij}^{(l)} = \frac{\partial \lambda^{(l)}}{\partial C_{ij}} \quad (19)$$

(7) Obtain the nine scalar derivatives,

$$\frac{\partial a}{\partial \lambda_r}, \frac{\partial b}{\partial \lambda_r}, \frac{\partial c}{\partial \lambda_r} \quad (20)$$

of equations (14) to (16) for $r = 1, 2$, and 3 .

(8) Substitute the preceding expressions and group-like terms, thus giving

$$\begin{aligned} D_{ijkl} = & 2a_1 C_{kl}^2 C_{ij}^2 + 2a_2 (C_{kl} C_{ij}^2 + C_{kl}^2 C_{ij}) + 2a_3 (\delta_{kl} C_{ij}^2 + C_{kl}^2 \delta_{ij}) \\ & + 2a_4 (C_{kl} C_{ij}) + 2a_5 (\delta_{ik} C_{jl} + C_{ik} \delta_{jl} + \delta_{kl} C_{ij} + C_{kl} \delta_{ij}) \\ & + 2a_6 (\delta_{ik} \delta_{jl} + \delta_{kl} \delta_{ij}) \end{aligned} \quad (21)$$

(9) For comparison of equation (21) to the forms described in reference 11, section 4, we make use of the symmetry properties of C_{ij} and δ_{ij} , and define two second order symmetric tensors, P and Q ,

$$P_{ijkl}(G, H) = G_{ik} H_{jl} + G_{il} H_{jk} \quad (22)$$

$$Q_{ijkl}(G, H) = G_{ik} H_{jl} + G_{ij} H_{jk} + G_{jl} H_{ik} + G_{jk} H_{il} \quad (23)$$

such that upon substitution we obtain

$$\begin{aligned} D_{ijkl} = & a_1 P(C_{kl}^2, C_{ij}^2) + a_2 [P(C_{kl}^2, C_{ij}) + P(C_{kl}, C_{ij}^2)] \\ & + a_3 [Q(C_{kl}^2 \delta_{ij}) + P(\delta_{kl}, C_{ij}^2)] + a_4 P(C_{kl}, C_{ij}) \end{aligned}$$

$$+ a_5[Q(C_{kl}, \delta_{ij}) + Q(\delta_{kl}, C_{ij})] + 2a_6 I_{ijkl} \quad (24)$$

where a_1, a_2, \dots, a_6 are as defined in reference 11 and the preceding equation (eq. (24)) is directly comparable to equations 4.6a in reference 11. Note that

$$I_{ijkl} = \frac{1}{2}[\delta_{ik}\delta_{jl} + \delta_{il}\delta_{jk}] \quad (25)$$

$$C_{ij}^2 = C_{im}C_{mj} \quad (26)$$

in the foregoing expressions.

Next, given the case of nondistinct eigenvalues, for example, case II when $(\lambda_1 \neq \lambda_2 = \lambda_3 = \lambda)$, or case III when $(\lambda_1 = \lambda_2 = \lambda_3)$, we must

- (10) Remove the singularity (case II) or singularities (case III) by defining an appropriate "path" for taking the limit of a, b, c and $C_{ik}C_{kj}$ in equations (13); that is,

- For case II
 $\lambda_1, \lambda_2 = \lambda + \Delta, \lambda_3 = \lambda - \Delta$
- For case III
 $\lambda_1 = \lambda, \lambda_2 = \lambda + \Delta, \lambda_3 = \lambda - \Delta$

- (11) Substitute the preceding eigenvalues into the expressions for $a, b,$ and c in equations (13), and take the limit of the numerator and denominator of $a, b,$ and c as $\Delta \rightarrow 0$.
- (12) If both limits are zero, apply l'Hospital's rule recursively to the now equivalent one dimensional problem. For example, given case II, we obtain

$$\lim_{\Delta \rightarrow 0} a(\Delta) = \frac{1}{(\lambda_1 - \lambda)^2} [s(\lambda_1) - s(\lambda) - (\lambda_1 - \lambda)s'(\lambda)] \quad (27)$$

$$\lim_{\Delta \rightarrow 0} b(\Delta) = \frac{1}{(\lambda_1 - \lambda)^2} [-2\lambda[s(\lambda_1) - s(\lambda)] + (\lambda_1^2 - \lambda^2)s'(\lambda)] \quad (28)$$

$$\lim_{\Delta \rightarrow 0} c(\Delta) = \frac{1}{(\lambda_1 - \lambda)^2} [\lambda^2 s(\lambda_1) + \lambda_1(\lambda_1 - 2\lambda)s(\lambda) - \lambda_1\lambda(\lambda_1 - \lambda)s'(\lambda)] \quad (29)$$

where

$$s'(\lambda_r) = \frac{\partial s(\lambda_r)}{\partial \lambda_{(r)}} = \frac{2\partial^2 W}{\partial \lambda_{(r)} \partial \lambda_{(r)}} \quad (30)$$

(13) Simplify $C_{ik}C_{kj}$ by using the definition of C_{ij} and N_{ij} , that is,

$$C_{ij}^2 = \lambda_1 N_{ij}^{(1)} + \lambda_2 N_{ij}^{(2)} + \lambda_3 N_{ij}^{(3)} \quad (31)$$

In addition, by using $\delta_{ij} = N_{ij}^{(1)} + N_{ij}^{(2)} + N_{ij}^{(3)}$ and equation (4), for case II we obtain

$$C_{ik}C_{kj} = \frac{1}{(\lambda_1 - \lambda)} [(\lambda_1^2 - \lambda^2)C_{ij} + (\lambda^2\lambda_1 - \lambda\lambda_1^2)\delta_{ij}] \quad (32)$$

and with equation (5) for case III we have

$$C_{ik}C_{kj} = \lambda\delta_{ij}. \quad (33)$$

(14) Substitute the limiting values of a,b,c and $C_{ik}C_{kj}$ into equations (13) and group like terms to obtain the modified stress function, S_{ij} , and the \bar{a} and \bar{b} values for case II:

$$S_{ij} = \bar{a}C_{ij} + \bar{b}\delta_{ij} \quad (34)$$

where

$$\bar{a} = \frac{s(\lambda_1) - s(\lambda)}{(\lambda_1 - \lambda)} \quad (35)$$

and

$$\bar{b} = -\frac{[s(\lambda_1)\lambda - s(\lambda)\lambda_1]}{(\lambda_1 - \lambda)} \quad (36)$$

For case III,

$$S_{ij} = s(\lambda)\delta_{ij} \quad (37)$$

(15) Repeat steps 5 through 10, but now use the appropriate modified stress function. For case II, this results in,

$$D_{ijkl} = 2\left(\frac{\partial \bar{a}}{\partial \lambda_1} N_{ij}^{(1)} + \frac{\partial \bar{a}}{\partial \lambda} N_{ij}^{(2)}\right)C_{kl} + 2\bar{a}\delta_{ijkl} + 2\left(\frac{\partial \bar{b}}{\partial \lambda_1} N_{ij}^{(1)} + \frac{\partial \bar{b}}{\partial \lambda} N_{ij}^{(2)}\right)\delta_{kl} \quad (38)$$

And for case III,

$$D_{ijkl} = 2\frac{\partial \bar{a}}{\partial \lambda_p} \frac{\partial \lambda_p}{\partial C_{kl}} \delta_{ij} \quad (39)$$

where the special derivative rule of equation (7) is now used.

The value in automating the foregoing procedure is evident: not only does this special purpose function relieve the user of the tedious manual derivation process but it also ensures analytical accuracy. This was illustrated prior to the publication of reference 11 in that a number of errors in the hand derivation were detected, verified and corrected. Furthermore, as will be discussed in a sequel paper [17], this automated derivation procedure facilitated the generalization of the preceding expressions to the *general nonseparable* case, which to the author's knowledge, has eluded researchers to date. Also, it should be apparent that this derivation process needs to be executed only once. However, with each new definition of W evaluation of $s(\lambda_{(I)})$ and $s'(\lambda_{(I)})$ is required in order to specialize the needed coefficients; for example, a, b and c , and a_1, a_2, \dots, a_6 . As a consequence, this motivated the development of SDIFFEV, as described in the next section.

3.2 SDIFFEV(case, W)

The function SDIFFEV symbolically evaluates the explicit expressions for the stress function S_{ij} and material moduli tensor D_{ijkl} , which were generated by SDIFF and stored in a LISP [18] level disk file. Only the coefficients of these expressions need be changed when a different strain-energy function is specified. The evaluation algorithm is illustrated here in pseudo code:

```

SDIFFEV(case, W)
IF (diff(W,  $\lambda_1$ ,  $\lambda_2$ ), diff(W,  $\lambda_2$ ,  $\lambda_3$ ), diff(W,  $\lambda_3$ ,  $\lambda_1$ ))=0 THEN
Display message: W is separable.
SEP = 1
ELSE Display message: W is non separable. SEP = 2
ENDIF
IF case=1 THEN .
Call Subroutine A
ELSE IF case = 2 THEN
. Call Subroutine B
ELSE IF case = 3 THEN
. Call Subroutine C
END IF
End

Subroutine A
IF SEP = 2 THEN
Do loop i = 1, 6
a[i] = ea[i] (ea[i] are the coefficients of tensor D stored on the disk file produced
by SDIFF(1))
End loop
ELSE IF SEP = 1 THEN

```

```

s[2,1] = s[3,1] = s[3,2] = 0
ENDIF
Do loop i = 1, 6
a[i] = ev(ea[i])
End loop
Do loop i = 1, 3
s[i] = 2*diff(W,λi,1)
s[i,i] = 2*diff(W,λi,2)
IF SEP = 2 THEN
Do loop j = 1, 3
s[i,j] = diff(W,λi, λj,2)
End loop
ENDIF
End loop
Call OPTION
Return End

```

```

Subroutine B
W = ev(W,λ3 = λ2)
IF SEP = 2 THEN
Do loop i = 1, 3
b[i] = eb[i] (eb[i] are the coefficients of tensor D stored on the disk file produced
by SDIFF (2))
End loop
ELSE IF SEP = 1 THEN
s[2,1] = 0
Do loop i = 1, 6
b[i] = ev(eb[i])
End loop
Do loop i = 1, 2
s[i] = 2*diff(W,λi,1)
s[i,i] = 2*diff(W,λi,2)
IF SEP = 2 THEN
Do loop j = 1, 2
s[i,j] = diff(W,λi, λj,2)
End loop
ENDIF
End loop
Call OPTION
Return

```

Subroutine C

```
W = ev(W,  $\lambda_3 = \lambda_2 = \lambda_1$ )
s[1]=2*diff(W, $\lambda_1$ ,1)
s[1,1]=2*diff(W, $\lambda_1$ ,2)
Call OPTION
Return
```

Subroutine OPTION

Display the formulae $S[i,j]$ and $D[i,j,k,l]$. Then, ask if user wants to see the symbolic form for the given function W , the intermediate step evaluations, and the derivatives of W .

```
READ(type y, or n to the question)
DISPLAY the options user may choose
Return
```

3.3 TEMPLATE ()

The function **TEMPLATE** is similar to the function **SDIFFEV** in that both will evaluate the explicit expressions obtained from **SDIFF**. As a result neither can be employed unless preceded by an invocation of **SDIFF**. **TEMPLATE**, however, will automatically generate the associated FORTRAN source code needed to evaluate the expressions numerically for a given potential function W . Code generation is accomplished by utilizing the built-in MACSYMA function **gentran**, and a number of template files. The template files can be thought of as a framework for the generation of four basic FORTRAN subroutines (i.e., the main driving routine **COMPSD** and the three subroutines – one each for case I, case II, and case III) and numerous functions. Appendix A contains the template file for the main driving routine **COMPSD**. This subroutine is constructed for easy implementation into a finite element code; the input requirements are the strain tensor C_m (denoted as *cmu*) and its associated eigenvalues (i.e., $\lambda_1, \lambda_2, \lambda_3$ denoted by *gl1*, *gl2*, and *gl3* respectively), and the outputs are the stress tensor S_n (denoted as *s*), and the material moduli tensor D_{nm} (denoted as *d*). Here, n and m run from 1 to 6. The only automated code generation required is that for the subroutines **COMPSD1**, **COMPSD2**, and **COMPSD3**. These codes are generated by issuing the command **<gentranin>**. The subroutines **COMPSD1**, **COMPSD2**, and **COMPSD3** are associated with case I ($\lambda_1 \neq \lambda_2 \neq \lambda_3$), case II ($\lambda_1, \lambda = \lambda_2 = \lambda_3$), and case III ($\lambda = \lambda_1 = \lambda_2 = \lambda_3$), described in Section 2.0. The template files corresponding to these three cases are shown, respectively in appendixes B, C, and D. Note that in these routines, most of the FORTRAN code is automatically generated, since it pertains to the definition of coefficients a, b, c ; a_1, a_2, \dots, a_6 , and the first (s_1, s_2, s_3 , see eq. (12)) and second scalar (s_{11}, s_{22}, s_{33} , see eq. (30)) derivatives of the strain energy function W . The

gentran commands are enclosed by double inequality signs, that is, $\ll \gg$. Finally, all functions that are associated with a given case have been included in the corresponding appendix.

4 Example

As an example, consider the case in which the strain energy function W of equation (1) consists of only two terms; that is,

$$W = x1(gl1^{y1} + gl2^{y1} + gl3^{y1}) + x2(gl1^{y2} + gl2^{y2} + gl3^{y2}) \quad (40)$$

where $x1$, $x2$, $y1$, and $y2$ are material coefficients and $gl1 = \lambda_1$, $gl2 = \lambda_2$, and $gl3 = \lambda_3$. After defining W , we can symbolically obtain the analytical expressions for S_{ij} and D_{ijkl} (given the case of three distinct eigenvalues) by merely issuing the command

sdiffov(1, W);

at the MACSYMA command level. Case II or III can just as easily be obtained by substituting a 2 or 3 in place of the 1 in this command. The resulting output is shown in appendix E where the expressions for the coefficients a, b, c and $a1, a2, \dots, a6$ could be further simplified and manipulated, if desired, by using other MACSYMA built-in functions. Typically, however, the analyst will ultimately desire a FORTRAN code for the resulting expressions in order to solve a given structural problem using the foregoing constitutive model. This code, described in the previous section, can easily be obtained by issuing the command

template();

at the MACSYMA command level. The generated FORTRAN code will then be stored in a file named temp.f. The automatically generated FORTRAN code for the above example is shown in appendix F.

5 Summary of Results

Taken separately, the main constituents of the deformation tensor (i.e., principal values and associated eigenvectors) are, in general, not uniquely defined and continuously differentiable functions. Careful consideration is thus called for in implementing constitutive models formulated in terms of these principal-strain measures. This difficulty can be entirely bypassed by resorting to explicit symbolic derivations of appropriate forms of the material tangent-stiffness matrices, which are valid over the entire deformation range. Furthermore, to enhance effective utilization and implementation of the present results, automatic FORTRAN generation of these explicit expressions has been pursued and

achieved. As a result, three special purpose functions(SDIFF, SDIFFEV and TEMPLATE), running under MACSYMA, have been developed and verified.

References

- [1] Green, A.E. and Adkins, J.E.: Large Elastic Deformations. 2nd Edition. Oxford Univ. Press, Clarendon, 1970.
- [2] Rivlin, R.S.: Large Elastic Deformations of Isotropic Materials; Part III: Experiments on the Deformation of Rubber. Philo. Trans. Roy. Soc., London, Ser. A, Vol. 243, pp.251-288, 1951.
- [3] Rivlin, R.S.: "Forty Years of Nonlinear Continuum Mechanics", Proc. IX Int., Congress on Rheology, Mexico, 1984.
- [4] Treloar, L.R.G.: The Physics of Rubber Elasticity, 3rd Ed., Oxford Univ. Press., U.K., 1975.
- [5] Ogden, R.W., "Recent Advances in the Phenomenological Theory of Rubber Elasticity," J.Rubber Chem. Tech., Vol. 59, pp. 361-383, 1986.
- [6] Finney, R.H., and Kumar, A.: "Development of Material Constants for Non-linear Finite Element Analysis," J. Rubber Chem. Tech., Vol.. 61, pp. 879-891, 1988.
- [7] Oden, J.T.: Finite Elements of Nonlinear Continua, McGraw Hill, New York, 1972.
- [8] Sussman, T.; and Bathe, K.J.: A Finite Element Formulation for Nonlinear Incompressible Elastic and Inelastic Analysis, Comp. Struct., Vol. 26, pp. 357-409, 1987.
- [9] Simo, J.C.; Taylor, R.L; and Pister, K.S.: "Variational and Projection Methods for the Volume Constraint in Finite Deformation Elastoplasticity", Comp. Meth. Appl. Mech. Engng., Vol. 51, pp. 177-208, 1985.
- [10] Chang, T.Y., Saleeb, A.F., and Li, G.: "Large Strain Analysis of Rubber-Like Materials by a Mixed Finite Element", Computational Mech., Vol. 8, No. 4, pp. 221-233, 1991.
- [11] Saleeb.A.F.; and Arnold, S.M.: Explicit Robust Schemes for Implementation of a Class Of Principle Value-Based Constitutive Models: Theoretical Development. NASA TM 105345, 1991
- [12] Arnold, S.M.; and Tan, H.Q.: "Symbolic Derivation of Potential Based Constitutive Equations", Computational Mech., Vol. 6, pp. 237-246, 1990.

- [13] Arnold, S.M.; Tan, H.Q.; and Dong, X.: Application of Symbolic Computations to The Constitutive Modeling of Structural Materials. Symbolic Computations and Their Impact on Mechanics, Noor, A.K., Elishakoff, I. and Hulbert, G., eds., PVP-Vol. 205, ASME, pp.215-229.,1990.
- [14] Malvern, L.E.: Introduction to the Mechanics of a Continuous Medium. Prentice-Hall, 1969.
- [15] Chen, W.F.; and Saleeb, A.F.: Constitutive Equations For Engineering Materials Volume 1: Elasticity and Modeling, John Wiley & Sons, 1982.
- [16] MATHLAB Group: MACSYMA Reference Manual. Version 10. Laboratory for Computer Science, Massachusetts Institute of Technology, Cambridge, MA, 1984.
- [17] Arnold, S.M.; et al.: Explicit Robust Schemes for Implementation of General Principal Value-Based Constitutive Models, submitted to Comp. and Struct., 1993.
- [18] Wilensky, R.: LISPcraft. W.W. Norton Co, 1984.

APPENDIX A: Template File Associated With COMPSD The Main Driver Routine

```

c *****
c      This is the template subroutine to calculate
c      tensor S and D. inputs are eigenvalues gl1,gl2,gl3,
c      and cmu(6). cmu is assumed to be engineering strain(e),
c      e.g. the Cauchy-green deformation tensor cm(3,3) is related
c      to cmu(6) in the following fashion:
c      cm(1,1)=cmu(1), cm(2,2) = cmu(2), cm(3,3) =cmu(3),
c      cmu(4)=2*cm(1,2), cm(5)=2*cm(2,3),cmu(6)=2*cm(1,3).
c      The outputs are the second order tensor S(6)
c      and forth order tensor D(6,6) are related in the
c      following way:
c      S=D*C
c      S(1,1) = S(1)
c      S(2,2) = S(2)
c      S(3,3) = S(3)
c      S(1,2) = S(4)
c      S(2,3) = S(5)
c      S(3,1) = S(6)
c      C(1,1) = C(1)
c      C(2,2) = C(2)
c      C(3,3) = C(3)
c      C(1,2) = C(4)
c      C(2,3) = C(5)
c      C(3,1) = C(6)
c
c      subroutine compsd(gl1,gl2,gl3,cmu,s,d)
c      real*8 gl1,gl2,gl3,ts(3,3),td(3,3,3,3)
c      real*8 delt(3,3),delt4(3,3,3,3),s(6),d(6,6)
c      real*8 cmu(6),cm(3,3)
c
c      converts cmu(6) to matrix cm(3,3) in a way that
c      cm(1,2)=cm(2,1)=cmu(4),cm(2,3)=cm(3,2)=cmu(5),
c      cm(1,3)=cm(3,1)=cmu(6).
c

```

```

do 5 i=1,3
  do 5 j=1,3
    if (i.eq.j) then
      iq=i
      cm(i,j)=cmu(iq)
    else if (i.ne.j) then
      if((i+j).eq.3) iq=4
      if((i+j).eq.4) iq=6
      if((i+j).eq.5) iq=5
      cm(i,j)=cmu(iq)/2
    end if
  continue
5 continue
c
c   Initiates the second identity tensor delt(3,3) which
c   is a 2X2 identity matrix.
c
do 6 i=1,3
  delt(i,i)=1.0
6  continue
c
c   Computes the forth order identity tensor delt4(3,3,3,3)
c   by definition.
c
do 7 i=1,3
  do 7 j=1,3
    delt4(i,j,i,j)=delt(i,i)*delt(j,j)+delt(i,j)*delt(j,i)
    delt4(i,j,j,i)=delt4(i,j,i,j)
7  continue
c
c*****
c   For different eigenvalues gl1,gl2,gl3 the computation
c   is different. case1 is gl1#gl2#gl3 call subroutine comsd1.
c   case2 is gl3=gl2#gl1 or gl1=gl3#gl2 or gl1=gl2#gl3 then
c   call subroutine compsd2. case3 is gl1=gl2=gl3 call subroutine
c   compsd3.
c*****
c   if ((gl1.ne.gl2).and.(gl2.ne.gl3).and.(gl1.ne.gl3)) then
c   call compsd1(gl1,gl2,gl3,delt,delt4,cm,ts,td)
c   else if((gl2.eq.gl3).and.(gl1.ne.gl3)) then

```

```

      call compsd2(g11,g12,delt,delt4,cm,ts,td)
    else if((g11.eq.g12).and.(g13.ne.g12)) then
      g11=g13
      call compsd2(g11,g12,delt,delt4,cm,ts,td)
    else if((g11.eq.g13).and.(g12.ne.g13)) then
      g11=g12
      g12=g13
      call compsd2(g11,g12,delt,delt4,cm,ts,td)
    else
      call compsd3(g11,delt,delt4,ts,td)
    end if
c
c      Rewrite the tensor ts(i,j) td(i,j,k,l)to S(i) and D(i,j)
c      respectively by using the symetric property.
c      converts ts(3,3) s(6) and td(3,3,3,3) to D(6,6)
c
      do 8 i=1,3
        do 8 j=i,3
          if (i.eq.j) iq=i
        if (i.eq.1.and.j.eq.2) iq=4
        if (i.eq.2.and.j.eq.3) iq=5
        if (i.eq.1.and.j.eq.3) iq=6
          s(iq)=ts(i,j)
        continue
      8 continue
      do 9 i=1,3
        do 9 j=i,3
          d(i,j)=td(i,i,j,j)
        continue
      9 continue
      do 10 i=1,3
        d(i,4)=td(i,i,1,2)+td(i,i,2,1)
        d(i,5)=td(i,i,2,3)+td(i,i,3,2)
        d(i,6)=td(i,i,3,1)+td(i,i,1,3)
      10 continue
      d(4,4)=(td(1,2,1,2)+td(1,2,2,1)+td(2,1,1,2)+td(2,1,2,1))/2.
      d(4,5)=(td(1,2,2,3)+td(1,2,3,2)+td(2,1,2,3)+td(2,1,3,2))/2.
      d(4,6)=(td(1,2,1,3)+td(1,2,3,1)+td(2,1,1,3)+td(2,1,3,1))/2.
      d(5,5)=(td(2,3,2,3)+td(2,3,3,2)+td(3,2,2,3)+td(3,2,3,2))/2.
      d(5,6)=(td(2,3,1,3)+td(2,3,3,1)+td(3,2,1,3)+td(3,2,3,1))/2.

```

```

d(6,6)=(td(3,1,1,3)+td(3,1,3,1)+td(1,3,1,3)+td(1,3,3,1))/2.
do 11 i = 1,6
  do 11 j = 1,6
    d(i,j) = d(j,i)
11  continue
c
c  prints out the inputs gl1,gl2,gl3,cmu(6) and outputs S and D
c
  print*, 'gl1=', gl1
  print*, 'gl2=', gl2
  print*, 'gl3=', gl3
  print*, 'Input tensor C(6):'
  print*, (cmu(i), i = 1,6)
  print*, "second order tensor S(6):"
  print*, (s(i), i=1,6)
  print*, "The forth order tensor D(6,6):"
  do 101 i=1,6
    print*,(d(i,j),j=1,6)
101  continue
  return
end
c
  subroutine compsd1(gl1,gl2,gl3,delt,delt4,cm,ts,td)
<<
  gentranin("case11.tem")$
>>
  subroutine compsd2(gl1,gl2,delt,delt4,cm,ts,td)
<<
  gentranin("case22.tem")$
>>
  subroutine compsd3(gl1,delt,delt4,ts,td)
<<
  gentranin("case3.tem")$
>>

```

```

c
c      This subroutine computes P and Q forth order tensors
c      which we define in tensor D.
c
      subroutine pqcom(cm1,cm2,p,q)
      real*8 cm1(3,3),cm2(3,3), p(3,3,3,3),q(3,3,3,3)
      do 100 i=1,3
        do 100 j=1,3
          do 100 k=1,3
            do 100 l=1,3
              p(i,j,k,l)=cm1(i,k)*cm2(j,l)+cm1(i,l)*cm2(j,k)
              q(i,j,k,l)=p(i,j,k,l)+cm1(j,l)*cm2(i,k)+cm1(j,k)*cm2(i,l)
100      continue
      return
      end

c
c      This subroutine computes matrix product cmXcm.
c
      subroutine product(mat1,cmm)
      real*8 mat1(3,3),cmm(3,3),sum
      do 25 i=1,3
        do 25 j=1,3
          sum=0.0
          do 26 k=1,3
            sum=sum+mat1(i,k)*mat1(k,j)
26      continue
          cmm(i,j)=sum
25  continue
      return
      end

```

APPENDIX B: Template File Associated With COMPSD1 Valid For Three Distinct Eigenvalues

```

real*8 gl1,gl2,gl3,ts(3,3),td(3,3,3,3)
real*8 cm(3,3),delt(3,3),delt4(3,3,3,3),p(3,3,3,3)
real*8 q(3,3,3,3),cmm(3,3),p1(3,3,3,3),p21(3,3,3,3)
real*8 p31(3,3,3,3),q11(3,3,3,3),q12(3,3,3,3),p22(3,3,3,3)
real*8 q21(3,3,3,3),q22(3,3,3,3),a,b,c,a1,a2,a3,a4,a5,a6

c
c      Obtains cmm(3,3)=cm(3,3)*cm(3,3) from subroutine product
c      call product(cm,cmm)
c      Uses the formula we derived in code to compute second order
c      tensor ts(3,3).
c
c<<

gentran(for i:1 thru 3 do
  (for j:1 thru 3 do
    (ts[i,j]:a(gl1,gl2,gl3)*cmm[i,j]+b(gl1,gl2,gl3)
      *cm[i,j]+c(gl1,gl2,gl3)*delt[i,j]))))$

>>
c
c      Call subroutine to compute all the functions we defined
c      when we derived forth order tenosor td, namely P(i,j,k,l)
c      and Q(i,j,k,l) which are the functions of cm(3,3) and
c      the matrix product cmm(3,3).
c

call pqcom(cmm,cmm,p1,q)
call pqcom(cmm,cm,p21,q)
call pqcom(cm,cmm,p22,q)
call pqcom(cm,cm,p31,q)
call pqcom(cmm,delt,p,q11)
call pqcom(delt,cmm,p,q12)
call pqcom(cm,delt,p,q21)
call pqcom(delt,cm,p,q22)

```

```

c
c      Computes forth order tensor td(i,j,k,l)
c
c
<<
gentran(for i:1 thru 3 do
  (for j:1 thru 3 do
    (for k:1 thru 3 do
      (for l:1 thru 3 do
        (td[i,j,k,l]:a1(gl1,gl2,gl3)*p1[i,j,k,l]+a2(gl1,gl2,gl3)
        *(p21[i,j,k,l]+p22[i,j,k,l])+a4(gl1,gl2,gl3)*p31[i,j,k,l]
        +a3(gl1,gl2,gl3)*(q11[i,j,k,l]+q12[i,j,k,l])+
        a5(gl1,gl2,gl3)*(q21[i,j,k,l]+q22[i,j,k,l])+
        a6(gl1,gl2,gl3)*delt4[i,j,k,l]))))$
      >>
    return
  end
  >>
c
c      a,b,c,a1-a6 are the coefficients we derived in code.
c
c
<<
gentran(a(gl1,gl2,gl3):=block(type(function,a),
                              type("real*8",gl1,gl2,gl3),
                              type("real*8",a,s1,s2,s3),
                              a:eval(ta)))$
  >>
  <<
gentran(b(gl1,gl2,gl3):=block(type(function,b),
                              type("real*8",b,gl1,gl2,gl3),
                              type("real*8",s1,s2,s3),
                              b:eval(tb)))$
  >>
  <<

gentran(c(gl1,gl2,gl3):=block(type(function,c),
                              type("real*8",c,gl1,gl2,gl3),
                              type("real*8",s1,s2,s3),
                              c:eval(tc)))$
  >>

```

```

<<
gentran(a1(gl1,gl2,gl3):=block(type(function,a1),
                                type("real*8",a1,gl1,gl2,gl3),
                                type("real*8",s1,s2,s3,s11,s22,s33),
                                a1:eval(ta1)))$
>>
<<
gentran(a2(gl1,gl2,gl3):=block(type(function,a2),
                                type("real*8",a2,gl1,gl2,gl3),
                                type("real*8",s1,s2,s3,s11,s22,s33),
                                a2:eval(ta2)))$
>>
<<
gentran(a3(gl1,gl2,gl3):=block(type(function,a3),
                                type("real*8",a3,gl1,gl2,gl3),
                                type("real*8",s1,s2,s3,s11,s22,s33),
                                a3:eval(ta3)))$
>>
<<
gentran(a4(gl1,gl2,gl3):=block(type(function,a4),
                                type("real*8",a4,gl1,gl2,gl3),
                                type("real*8",s1,s2,s3,s11,s22,s33),
                                a4:eval(ta4)))$
>>
<<
gentran(a5(gl1,gl2,gl3):=block(type(function,a5),
                                type("real*8",a5,gl1,gl2,gl3),
                                type("real*8",s1,s2,s3,s11,s22,s33),
                                a5:eval(ta5)))$
>>
<<
gentran(a6(gl1,gl2,gl3):=block(type(function,a6),
                                type("real*8",a6,gl1,gl2,gl3),
                                type("real*8",s1,s2,s3,s11,s22,s33),
                                a6:eval(ta6)))$
>>

```



```

c
c      The s1,s2,s3,s11,s22,s33 are derivatives of W
c
      function s1(gl1,gl2,gl3)
      <<cut(var);>>
<<
      gentran(type(  "real*8",s1,gl1,gl2,gl3),
                  s1:2*eval(diff(w,'gl1,1')))$
>>

      return
      end
c
      function s2(gl1,gl2,gl3)
      <<cut(var);>>
<<
      gentran(type(  "real*8",s2,gl1,gl2,gl3),
                  s2:2*eval(diff(w,'gl2,1')))$
>>

      return
      end
c
      function s3(gl1,gl2,gl3)
      <<cut(var);>>
<<
      gentran(type(  "real*8",s3,gl1,gl2,gl3),
                  s3:2*eval(diff(w,'gl3,1')))$
>>

      return
      end
c
      function s11(gl1,gl2,gl3)
      <<cut(var);>>
<<
      gentran(type(  "real*8",s11,gl1,gl2,gl3),
                  s11:2*eval(diff(w,'gl1,2')))$
>>

      return
      end
c

```

```

function s22(gl1,gl2,gl3)
<<cut(var);>>
<<
    gentran(type("real*8",s22,gl1,gl2,gl3),
              s22:2*eval(diff(w,'gl2,2)))$
>>
    return
end
c
function s33(gl1,gl2,gl3)
<<cut(var);>>
<<
    gentran(type( "real*8",s33,gl1,gl2,gl3),
              s33:2*eval(diff(w,'gl3,2)))$
>>
    return
end

```

APPENDIX C: Template File Associated With COMPSD2 Valid For Double Coalescence Case

```

c
c
real*8 gl1,gl2,ts(3,3),td(3,3,3,3)
real*8 cm(3,3),delt(3,3),delt4(3,3,3,3),p1(3,3,3,3)
real*8 q2(3,3,3,3),q1(3,3,3,3),p(3,3,3,3),q(3,3,3,3)
real*8 b1,b2,b3, abar,bbar

c
c Computes second order tensor ts(i,j) based on the formula
c derived in code.
c
c<<
gentran(for i:1 thru 3 do
  (for j:1 thru 3 do
    (ts[i,j]:abar(gl1,gl2)*cm[i,j]+bbar(gl1,gl2)*delt[i,j]))))$
>>
c
c Call subroutine to get P, Q which are defined in code.
c
call pqcom(cm,cm,p1,q)
call pqcom(cm,delt,p,q1)
call pqcom(delt,cm,p,q2)

c
c Computes tensor td(i,j,k,l).
c
c<<
gentran(for i:1 thru 3 do
  (for j:1 thru 3 do
    (for k:1 thru 3 do
      (for l:1 thru 3 do
        (td[i,j,k,l]:b1(gl1,gl2)*p1[i,j,k,l]+b2(gl1,gl2)*
          (q1[i,j,k,l]+q2[i,j,k,l])+b3(gl1,gl2)*
            delt4[i,j,k,l])))))))$
>>

return
end

```

```

c
c      abar,bbar are b1, b2, b3 functions derived in code.
c
<<
gentran(abar(gl1,gl2):=block(type(function,abar),
                             type("real*8",abar,gl1,gl2),
                             type("real*8", ss1,ss2),
                             abar:eval(abar)))$
>>
<<
gentran(bbar(gl1,gl2):=block(type(function,bbar),
                             type("real*8",bbar,gl1,gl2),
                             type("real*8", ss1,ss2),
                             bbar:eval(bbar)))$
>>
<<
gentran(b1(gl1,gl2):=block(type(function,b1),
                             type("real*8",b1,gl1,gl2),
                             type("real*8", ss1,ss2,ss11,ss22),
                             b1:eval(tb1)))$
>>
<<
gentran(b2(gl1,gl2):=block(type(function,b2),
                             type("real*8",b2,gl1,gl2),
                             type("real*8", ss1,ss2,ss11,ss22),
                             b2:eval(tb2)))$
>>
<<
gentran(b3(gl1,gl2):=block(type(function,b3),
                             type("real*8",b3,gl1,gl2),
                             type("real*8", ss1,ss2,ss11,ss22),
                             b3:eval(tb3)))$
>>
<<
neww:subst(['gl3='gl2],w)$
>>

```

```

c
c      ss1,ss2,ss11,ss22 are derivatives of W.
c
      function ss1(gl1,gl2)
      <<cut(var);>>
<<
      gentran(type("real*8",ss1,gl1,gl2),
              ss1:2*eval(diff(neww,'gl1,1')))$
>>
      return
      end
c
      function ss2(gl1,gl2)
      <<cut(var);>>
<<
      gentran(type("real*8",ss2,gl1,gl2),
              ss2:2*eval(diff(neww,'gl2,1')))$
>>
      return
      end
c
      function ss11(gl1,gl2)
      <<cut(var);>>
<<
      gentran(type("real*8",ss11,gl1,gl2),
              ss11:2*eval(diff(neww,'gl1,2')))$
>>
      return
      end
c
      function ss22(gl1,gl2)
      <<cut(var);>>
<<
      gentran(type("real*8",ss22,gl1,gl2),
              ss22:2*eval(diff(neww,'gl2,2')))$
>>
      return
      end

```

APPENDIX D: Template File Associated With COMPSD3
Valid For The Triple Coalescence Case

```

c
real*8 gl1,ts(3,3),td(3,3,3,3),delt(3,3),delt4(3,3,3,3)
real*8 cc1,abbar
c
<<
gentran(for i:1 thru 3 do
  (for j:1 thru 3 do
    (ts[i,j]:abbar(gl1)*delt[i,j]))))$
>>
<<
gentran(for i:1 thru 3 do
  (for j:1 thru 3 do
    (for k:1 thru 3 do
      (for l:1 thru 3 do
        (td[i,j,k,l]:cc1(gl1)*delt4[i,j,k,l])))))$
>>
return
end
<<
gentran(abbar(gl1):=block(type(function,abbar),
  type("real*8", abbar,gl1),
  abbar:eval(abbar)))$
>>
<<
gentran(cc1(gl1):=block(type(function,cc1),
  type("real*8", cc1,gl1),
  cc1:eval(cc1)))$
>>
<<
www:subst(['gl3='gl1, 'gl2='gl1],w)$
>>

```

```

c
function sss1(gl1)
<<cut(var);>>
<<
    gentran(type("real*8",sss1,gl1),
        sss1:2*eval(diff(www,'gl1,1')))$
>>
    return
end
c
function sss11(gl1)
<<cut(var);>>
<<
    gentran(type("real*8",sss11,gl1),
        sss11:2*eval(diff(www,'gl1,2')))$
>>
    return
end

```

APPENDIX E: Listing of MACSYMA Session Resulting From Issuing The SDIFFEV Command

c6) sdiffev(1,w);

w is a separable function.

$$w = (g_{l3} + g_{l2} + g_{l1}) x_2 + (g_{l3} + g_{l2} + g_{l1}) x_1$$

This is case 1 with distinct eigenvalues $g_{l1} \# g_{l2} \# g_{l3}$.

Please type y if your answer is yes, otherwise type n to skip it.

Do you want to display the second order tensor $s[i,j]$?

y;

$$s_{i,j} = a c_{i,k} c_{k,j} + c \delta_{i,j} + b c_{i,j}$$

Do you want to display $c[i,j]$ and $\delta[i,j]$?

y;

$$\delta_{i,j} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$c_{i,j} = g_{l3} n_{3,i} n_{3,j} + g_{l2} n_{2,i} n_{2,j} + g_{l1} n_{1,i} n_{1,j}$$

n_1, n_2, n_3 are eigenvectors associatedd with eigenvalues g_{l1}, g_{l2}, g_{l3} .

If $c[i,j]$ is given then the eigenvectors can be computed

Do you want to display a,b,c in s[i,j] form?

y;

$$a = - \frac{s3}{t31 (t31 - t21)} + \frac{s2}{t21 (t31 - t21)} - \frac{s1}{t21 t31}$$

$$b = - \frac{s3 (2 t31 - t21 - 2 gl3)}{t31 (t31 - t21)} + \frac{s1 (t31 + t21 + 2 gl1)}{t21 t31}$$

$$- \frac{s2 (t31 - t21 + gl2 + gl1)}{t21 (t31 - t21)}$$

$$c = - \frac{s3 (t31 - gl3) (t31 - t21 - gl3)}{t31 (t31 - t21)} - \frac{s2 (t21 - gl2) (t31 + gl1)}{t21 (t31 - t21)}$$

$$- \frac{s1 (t21 + gl1) (t31 + gl1)}{t21 t31}$$

Do you want to display t21,t31, s1,s2,s3?

y;

$$t21 = gl2 - gl1$$

$$t31 = gl3 - gl1$$

s1,s2,s3 are the first derivatives of W with respect to gl1,gl2,gl3.

$$s1 = gl1^{y2 - 1} x2 y2 + gl1^{y1 - 1} x1 y1$$

$$s2 = g_{12}^{y2-1} x_2 y_2 + g_{12}^{y1-1} x_1 y_1$$

$$s3 = g_{13}^{y2-1} x_2 y_2 + g_{13}^{y1-1} x_1 y_1$$

Do you want to display the forth order tensor d[i,j,k,l]?
y;

$$d_{i,j,k,l} = [a_{6} \text{delt4} + a_{3} (q_{2}(\text{delt}, c) + q_{2}(c, \text{delt})) + a_{5} (q_{1}(\text{delt}, c) + q_{1}(c, \text{delt})) + a_{1} p_{2}(c, c) + a_{2} (p_{2}(c, c) + p_{2}(c, c)) + a_{4} p_{4}(c, c)]$$

Do you want to display the functions p,q and delt4?
y;

$$p(g, h) = g_{i,k} h_{j,l} + g_{i,l} h_{j,k}$$

$$q(g, h) = g_{i,k} h_{j,l} + h_{i,k} g_{j,l} + 2 g_{i,l} h_{j,k}$$

$$\text{delt4} = \text{delt}_{i,k} \text{delt}_{j,l} + \text{delt}_{i,l} \text{delt}_{j,k}$$

Do you want to continue displaying a
1
y;

$$\begin{aligned}
 a = & \frac{2 s_3 (2 t_{31} - t_{21})}{t_{31}^3 (t_{31} - t_{21})} - \frac{2 s_1 (t_{31} + t_{21})}{t_{21}^3 t_{31}} - \frac{s_{33}}{t_{31}^2 (t_{31} - t_{21})} \\
 & - \frac{s_{22}}{t_{21}^2 (t_{31} - t_{21})} + \frac{2 s_2 (t_{31} - 2 t_{21})}{t_{21}^3 (t_{31} - t_{21})} - \frac{s_{11}}{t_{21}^2 t_{31}}
 \end{aligned}$$

Do you want to continue displaying a ?

2

y;

$$\begin{aligned}
 a = & \frac{s_1 (2 t_{31}^2 + 3 t_{21} t_{31} + 4 g_{11} t_{31} + 2 t_{21}^2 + 4 g_{11} t_{21})}{t_{21}^3 t_{31}} \\
 & - \frac{s_2 (2 t_{31}^2 - 3 t_{21} t_{31} + 4 g_{11} t_{31} - t_{21}^2 - 8 g_{11} t_{21})}{t_{21}^3 (t_{31} - t_{21})} \\
 & - \frac{s_3 (t_{31}^2 + 3 t_{21} t_{31} + 8 g_{11} t_{31} - 2 t_{21}^2 - 4 g_{11} t_{21})}{t_{31}^3 (t_{31} - t_{21})} \\
 & - \frac{s_{33} (2 t_{31} - t_{21} - 2 g_{13})}{t_{31}^2 (t_{31} - t_{21})} + \frac{s_{11} (t_{31} + t_{21} + 2 g_{11})}{t_{21}^2 t_{31}}
 \end{aligned}$$

$$+ \frac{s_{22} (t_{31} - t_{21} + g_{l2} + g_{l1})}{t_{21}^2 (t_{31} - t_{21})^2}$$

Do you want to continue displaying a ?
3

y;

$$\begin{aligned} a = & -s_3 (t_{31}^3 - 2 t_{21}^2 t_{31} - g_{l1} t_{31}^2 + t_{21}^2 t_{31} - 3 g_{l1} t_{21} t_{31} - 4 g_{l1}^2 t_{31}^2 \\ & + 2 g_{l1}^2 t_{21}^2 + 2 g_{l1}^3 t_{21}) / (t_{31}^3 (t_{31} - t_{21})^3) \\ & - s_1 (t_{21}^2 t_{31}^2 + 2 g_{l1}^2 t_{31}^2 + t_{21}^2 t_{31}^2 + 3 g_{l1}^2 t_{21} t_{31} + 2 g_{l1}^2 t_{31}^2 \\ & + 2 g_{l1}^2 t_{21}^2 + 2 g_{l1}^3 t_{21}) / (t_{21}^3 t_{31}^3) \\ & + s_2 (t_{21}^2 t_{31}^2 + 2 g_{l1}^2 t_{31}^2 - 2 t_{21}^2 t_{31}^2 - 3 g_{l1}^2 t_{21} t_{31} + 2 g_{l1}^2 t_{31}^2 + t_{21}^3 \\ & - g_{l1}^2 t_{21}^2 - 4 g_{l1}^3 t_{21}) / (t_{21}^3 (t_{31} - t_{21})^3) \\ & - \frac{s_{33} (t_{31} - g_{l3}) (t_{31} - t_{21} - g_{l3})}{t_{31}^2 (t_{31} - t_{21})^2} + \frac{s_{22} (t_{21} - g_{l2}) (t_{31} + g_{l1})}{t_{21}^2 (t_{31} - t_{21})^2} \\ & - \frac{s_{l1} (t_{21} + g_{l1}) (t_{31} + g_{l1})}{t_{21}^2 t_{31}^2} \end{aligned}$$

Do you want to continue displaying a ?

4

y;

$$\begin{aligned}
 a &= \frac{2 s_3 (t_{21} + 2 g_{11}) (t_{31}^2 + t_{21} t_{31} + 4 g_{11} t_{31} - t_{21}^2 - 2 g_{11} t_{21})}{t_{31}^3 (t_{31} - t_{21})^3} \\
 &- \frac{2 s_1 (t_{31} + t_{21} + 2 g_{11}) (t_{31}^2 + t_{21} t_{31} + 2 g_{11} t_{31} + t_{21}^2 + 2 g_{11} t_{21})}{t_{21}^3 t_{31}^3} \\
 &+ \frac{2 s_2 (t_{31} + 2 g_{11}) (t_{31}^2 - t_{21} t_{31} + 2 g_{11} t_{31} - t_{21}^2 - 4 g_{11} t_{21})}{t_{21}^3 (t_{31} - t_{21})^3} \\
 &- \frac{s_{33} (2 t_{31} - t_{21} - 2 g_{13})}{t_{31}^2 (t_{31} - t_{21})^2} - \frac{s_{11} (t_{31} + t_{21} + 2 g_{11})}{t_{21}^2 t_{31}^2} \\
 &- \frac{s_{22} (t_{31} - t_{21} + g_{12} + g_{11})}{t_{21}^2 (t_{31} - t_{21})^2}
 \end{aligned}$$

Do you want to continue displaying a ?

5

y;

$$\begin{aligned}
a = & s_1 \frac{(t_{21}^3 t_{31}^3 + 2 g_{11}^2 t_{31}^3 + t_{21}^2 t_{31}^2 + 6 g_{11} t_{21} t_{31}^2 + 6 g_{11}^2 t_{31}^2}{5} \\
& + t_{21}^3 t_{31} + 6 g_{11}^2 t_{21} t_{31} + 9 g_{11}^2 t_{21} t_{31} + 4 g_{11}^3 t_{31} + 2 g_{11}^3 t_{21} \\
& + 6 g_{11}^2 t_{21}^2 + 4 g_{11}^3 t_{21}) / (t_{21}^3 t_{31}) \\
& + s_3 \frac{(t_{21}^3 t_{31}^3 + 2 g_{11}^3 t_{31}^3 - 2 t_{21}^2 t_{31}^2 - 6 g_{11}^2 t_{21} t_{31}^2 - 3 g_{11}^2 t_{31}^2}{3} \\
& + t_{21}^3 t_{31} - 9 g_{11}^2 t_{21} t_{31} - 8 g_{11}^3 t_{31} + 2 g_{11}^3 t_{21} + 6 g_{11}^2 t_{21}^2 \\
& + 4 g_{11}^3 t_{21}) / (t_{31}^3 (t_{31} - t_{21})) - s_2 \\
& \frac{(t_{21}^3 t_{31}^3 + 2 g_{11}^3 t_{31}^3 - 2 t_{21}^2 t_{31}^2 + 6 g_{11}^2 t_{31}^2 + t_{21}^2 t_{31}^3 - 6 g_{11}^2 t_{21}^2 t_{31} - 9 g_{11}^2 t_{21} t_{31} + 4 g_{11}^3 t_{31} + 2 g_{11}^3 t_{21}^3 - 3 g_{11}^2 t_{21}^2 - 8 g_{11}^3 t_{21})}{(t_{21}^3 (t_{31} - t_{21}))} \\
& \frac{s_{33} (t_{31} - g_{13}) (t_{31} - t_{21} - g_{13}) (2 t_{31} - t_{21} - 2 g_{13})}{t_{31}^2 (t_{31} - t_{21})^2} \\
& + \frac{s_{11} (t_{21} + g_{11}) (t_{31} + g_{11}) (t_{31} + t_{21} + 2 g_{11})}{t_{21}^2 t_{31}^2}
\end{aligned}$$

$$-\frac{s_{22} (t_{21} - g_{l2}) (t_{31} + g_{l1}) (t_{31} - t_{21} + g_{l2} + g_{l1})}{t_{21}^2 (t_{31} - t_{21})^2}$$

Do you want to continue displaying a ?
6

y;

a =
6

$$-\frac{2 g_{l1} s_1 (t_{21} + g_{l1}) (t_{31} + g_{l1}) (t_{31}^2 + t_{21} t_{31} + g_{l1} t_{31} + t_{21}^2 + g_{l1} t_{21})}{t_{21}^3 t_{31}^3}$$

$$+ 2 g_{l1} s_2 (t_{21} + g_{l1}) (t_{31} + g_{l1}) (t_{31}^2 - 2 t_{21} t_{31} + g_{l1} t_{31} + t_{21}^2 - 2 g_{l1} t_{21}) / (t_{21}^3 (t_{31} - t_{21})^3) - 2 g_{l1} s_3 (t_{21} + g_{l1}) (t_{31} + g_{l1})$$

$$(t_{31}^2 - 2 t_{21} t_{31} - 2 g_{l1} t_{31} + t_{21}^2 + g_{l1} t_{21}) / (t_{31}^3 (t_{31} - t_{21})^3)$$

$$-\frac{s_{33} (t_{31} - g_{l3})^2 (t_{31} - t_{21} - g_{l3})^2}{t_{31}^2 (t_{31} - t_{21})^2} - \frac{s_{22} (t_{21} - g_{l2})^2 (t_{31} + g_{l1})^2}{t_{21}^2 (t_{31} - t_{21})^2}$$

$$-\frac{s_{11} (t_{21} + g_{l1})^2 (t_{31} + g_{l1})^2}{t_{21}^2 t_{31}^2}$$

Do you want to display s11?

y;

$$s_{11} = g_{11} \frac{y_2 - 2}{x_2 (y_2 - 1) y_2} + g_{11} \frac{y_1 - 2}{x_1 (y_1 - 1) y_1}$$

Do you want to display s22?

y;

$$s_{22} = g_{12} \frac{y_2 - 2}{x_2 (y_2 - 1) y_2} + g_{12} \frac{y_1 - 2}{x_1 (y_1 - 1) y_1}$$

Do you want to display s33?

y;

$$s_{33} = g_{13} \frac{y_2 - 2}{x_2 (y_2 - 1) y_2} + g_{13} \frac{y_1 - 2}{x_1 (y_1 - 1) y_1}$$

(d6) done

APPENDIX F: Listing of Automatically Generated FORTRAN Code

```

c      This is the template subroutine to calculate tensor S and D.
c      inputs are eigenvalues gl1,gl2,gl3, and cmu(6). cmu is assumed to be
c      engineering strain(e), e.g. the Cauchy-green deformation tensor
c      cm(3,3) is related to cmu(6) in the following fashion:
c
c      cm(1,1)=cmu(1), cm(2,2)=cmu(2), cm(3,3)=cmu(3), cmu(4)=2*cm(1,2),
c      cm(5)=2*cm(2,3),cmu(6)=2*cm(1,3).
c
c      The outputs are the second order tensor S(6) and forth order
c      tensor D(6,6) are related in the following way:
c
c      S=D*C
c
c      S(1,1) = S(1)
c      S(2,2) = S(2)
c      S(3,3) = S(3)
c      S(1,2) = S(4)
c      S(2,3) = S(5)
c      S(3,1) = S(6)
c
c      C(1,1) = C(1)
c      C(2,2) = C(2)
c      C(3,3) = C(3)
c      C(1,2) = C(4)
c      C(2,3) = C(5)
c      C(3,1) = C(6)
c
c      subroutine compsd(gl1,gl2,gl3,cmu,s,d)
c      real*8 gl1,gl2,gl3,ts(3,3),td(3,3,3,3)
c      real*8 delt(3,3),delt4(3,3,3,3),s(6),d(6,6)
c      real*8 cmu(6),cm(3,3)
c
c      converts cmu(6) to matrix cm(3,3) in a way that
c      cm(1,2)=cm(2,1)=cmu(4),cm(2,3)=cm(3,2)=cmu(5),
c      cm(1,3)=cm(3,1)=cmu(6)

```

```

c
do 5 i=1,3
  do 5 j=1,3
    if (i.eq.j) then
      iq=i
      cm(i,j)=cmu(iq)
    else if (i.ne.j) then
      if((i+j).eq.3) iq=4
      if((i+j).eq.4) iq=6
      if((i+j).eq.5) iq=5
      cm(i,j)=cmu(iq)/2
    end if
  5 continue
c
c   Initiates the second identity tensor delt(3,3) which
c   is a 2X2 identity matrix
c
do 6 i=1,3
  delt(i,i)=1.0
  6 continue
c
c   Computes the forth order identity tensor delt4(3,3,3,3)
c   by definition.
c
do 7 i=1,3
  do 7 j=1,3
    delt4(i,j,i,j)=delt(i,i)*delt(j,j)+delt(i,j)*delt(j,i)
    delt4(i,j,j,i)=delt4(i,j,i,j)
  7 continue
c
c*****
c   For different eigenvalues gl1,gl2,gl3 the computation is
c   different.
c   case1 is gl1#gl2#gl3 call subroutine comsd1.
c   case2 is gl3=gl2#gl1 or gl1=gl3#gl2 or gl1=gl2#gl3 then
c   call subroutine compsd2.
c   case3 is gl1=gl2=gl3 call subroutine ccmpsd3.
c*****
c

```

```

      if ((gl1.ne.gl2).and.(gl2.ne.gl3).and.(gl1.ne.gl3)) then
        call compsd1(gl1,gl2,gl3,ae t,delt4,cm,ts,td)
      else if((gl2.eq.gl3).and.(gl1.ne.gl3)) then
        call compsd2(gl1,gl2,delt,delt4,cm,ts,td)
      else if((gl1.eq.gl2).and.(gl3.ne.gl2)) then
        gll=gl3
        call compsd2(gl1,gl2,delt,delt4,cm,ts,td)
      else if((gl1.eq.gl3).and.(gl2.ne.gl3)) then
        gll=gl2
        gl2=gl3
        call compsd2(gl1,gl2,delt,delt4,cm,ts,td)
      else
        call compsd3(gl1,delt,delt4,ts,td)
      end if

c
c      Rewrite the tensor ts(1,j) td(i,j,k,l) to S(i) and D(i,j)
c      respectively by using the symetric property.
c      converts ts(3,3) s(6) and td(3,3,3,3) to D(6,6)
c
      do 8 i=1,3
        do 8 j=i,3
          if (i.eq.j) iq=i
          if (i.eq.1.and.j.eq.2) iq=4
          if (i.eq.2.and.j.eq.3) iq=5
          if (i.eq.1.and.j.eq.3) iq=6
        s(iq)=ts(i,j)
      8      continue
c
      do 9 i=1,3
        do 9 j=i,3
          d(i,j)=td(i,i,j,j)
      9      continue
c
      do 10 i=1,3
        d(i,4)=td(i,i,1,2)+td(i,i,2,1)
        d(i,5)=td(i,i,2,3)+td(i,i,3,2)
        d(i,6)=td(i,i,3,1)+td(i,i,1,3)
      10      continue
c

```

```

d(4,4)= (td(1,2,1,2)+td(1,2,2,1)+td(2,1,1,2)+td(2,1,2,1))/2.
d(4,5)= (td(1,2,2,3)+td(1,2,3,2)+td(2,1,2,3)+td(2,1,3,2))/2.
d(4,6)= (td(1,2,1,3)+td(1,2,3,1)+td(2,1,1,3)+td(2,1,3,1))/2.
c
d(5,5)=(td(2,3,2,3)+td(2,3,3,2)+td(3,2,2,3)+td(3,2,3,2))/2.
d(5,6)=(td(2,3,1,3)+td(2,3,3,1)+td(3,2,1,3)+td(3,2,3,1))/2.
d(6,6)=(td(3,1,1,3)+td(3,1,3,1)+td(1,3,1,3)+td(1,3,3,1))/2.
c
do 11 i=1,6
  do 11 j=1,6
    d(i,j) = d(j,i)
11  continue
c
c  prints out the inputs gl1,c12,gl3,cmu(6)
c  and outputs S and D
c
print*, ' gl1=', gl1
print*, ' gl2=', gl2
print*, ' gl3=', gl3
print*, 'Input tensor C(6):'
print*, (cmu(i), i=1,6)
print*, "second order tensor S(6):"
print*, (s(i), i=1,6)
print*, "The forth order tensor D(6,6):"
do 101 i=1,6
  print*, (d(i,j),j=1,6)
101  continue
  return
end
c
subroutine compsd1(gl1,c312,gl3,delt,delt4,cm,ts,td)
c
real*8 gl1,gl2,gl3,ts(3,3),td(3,3,3,3)
real*8 cm(3,3),delt(3,3),delt4(3,3,3,3),p(3,3,3,3)
real*8 q(3,3,3,3),cmm(3,3),p1(3,3,3,3),p21(3,3,3,3)
real*8 p31(3,3,3,3),q11(3,3,3,3),q12(3,3,3,3),p22(3,3,3,3)
real*8 q21(3,3,3,3),q22(3,3,3,3),a,b,c,a1,a2,a3,a4,a5,a6

```

```

c
c Obtains cmm(3,3)=cm(3,3)*cm(3,3) from subroutine product
c
      call product(cm,cmm)
c
c Uses the formula we derived in code to compute second order
c tensor ts(3,3).
c
      do 25037 i=1,3
        do 25038 j=1,3
          ts(i,j)=a(gl1,gl2,gl3)*cmm(i,j)+b(gl1,gl2,gl3)*cm(i,j)
          c(gl1,gl2,gl3)*delt(i,j)
25038      continue
25037      continue
c
c      call subroutine to compute the functions we defined
c      when we derived forth order tensor td, namely P(i,j,k,l)
c      and Q(l,j,k,l) which are the functions of cm(3,3) and
c      the matrix product cmm(3,3).
c
      call pqcom(cmm,cmm,p1,q)
      call pqcom(cmm,cm,p21,q)
      call pqcom(cm,cmm,p22,q)
      call pqcom(cm,cm,p31,q)
      call pqcom(cmm,delt,p,q11)
      call pqcom(delt,cmm,p,q12)
      call pqcom(cm,delt,p,q21)
      call pqcom(delt,cm,P,q22)
c
c      computes forth order tensor td(i,j,k,l)
c
      do 25039 i=1,3
        do 25040 j=1,3
          do 25041 k=1,3
            do 25042 l=1,3
              td(i,j,k,l)=a1(gl1,gl2,gl3)*p1(i,j,k,l)+
              a2(gl1,gl2,gl3)*(p21(i,j,k,l)+p22(i,j,k,l))+
              a4(gl1,gl2,gl3)*p31(i,j,k,l)+a3(gl1,gl2,gl3)*
              (q11(i,j,k,l)+q12(i,j,k,l))+a5(gl1,gl2,gl3)*
              (q21(i,j,k,l)+q22(i,j,k,l))+a6(gl1,gl2,gl3)

```

```

                *delt4(i,j,k,l)
25042          continue
25041          continue
25040          continue
25039          continue
c
          return
          end
c
c          a,b,c,a1-a6 are the coefficients we derived in code.
c
      real*8 function a(gl1,gl2,gl3)
      real*8 gl1,gl2,gl3,s1,s2,s3
      a=-s1(gl1,gl2,gl3)/(-gl1+gl2)/(-gl1+gl3)+s2(gl1,gl2,gl3)/
      (-gl1+gl2)/(-gl2+gl3)-s3(gl1,gl2,gl3)/(-gl1+gl3)/(-gl2+gl3)
      return
      end
c
      real*8 function b(gl1,gl2,gl3)
      real*8 gl1,gl2,gl3,s1,s2,s3
      b=s3(gl1,gl2,gl3)*(gl1+gl2)/(-gl1+gl3)/(-gl2+gl3)-s2(gl1,gl2,
      gl3)*(gl1+gl3)/(-gl1+gl2)/(-gl2+gl3)+s1(gl1,gl2,gl3)*(gl2+gl3)
      /(-gl1+gl2)/(-gl1+gl3)
      return
      end
c
      real*8 function c(gl1,gl2,gl3)
      real*8 gl1,gl2,gl3,s1,s2,s3
      c=-s1(gl1,gl2,gl3)*gl2*gl3/(-gl1+gl2)/(-gl1+gl3)+gl1*s2(gl1,
      gl2,gl3)*gl3/(-gl1+gl2)/(-gl2+gl3)-gl1*s3(gl1,gl2,gl3)*gl2/
      (-gl1+gl3)/(-gl2+gl3)
      return
      end
c

```

```

real*8 function a1(g11,g12,g13)
real*8 g11,g12,g13,s1,s2,s3,s11,s22,s33

c
a1=- s11(g11,g12,g13)/(-g11+g12)**2/(-g11+g13)**2+2.0*
.      s2(g11,g12,g13)*(g11-2.0*g12+g13)/(-g11+g12)**3/
.      (-g12+g13)**3-s22(g11,g12,g13)/(-g11+g12)**2/(-g12+g13)**2
.      -s33(g11,g12,g13)/(-g11+g13)**2/(-g12+g13)**2-2.0*
.      s1(g11,g12,g13)*(-2.0*g11+g12+g13)/(-g11+g12)**3/
.      (-g11+g13)**3+2.0*s3(g11,g12,g13)*(-g11-g12+2.0*g13)/
.      (-g11+g13)**3/(-g12+g13)**3
return
end

c
real*8 function a2(g11,g12,g13)
real*8 g11,g12,g13,s1,s2,s3,s11,s22,s33
a2=s33(g11,g12,g13)*(g11+g12)/(-g11+g13)**2/(-g12+g13)**2+
. s22(g11,g12,g13)*(g11+g13)/(-g11+g12)**2/(-g12+g13)**2+
. s11(g11,g12,g13)*(g12+g13)/(-g11+g12)**2/(-g11+g13)**2
. -s3(g11,g12,g13)*(-2.0*g11**2-3.0*g11*g12-2.0*g12**2+
. 3.0*g11*g13+3.0*g12*g13+g13**2)/(-g11+g13)**3/(-g12+g13)**3
. -s2(g11,g12,g13)*(2.0*g11**2-3.0*g11*g12-g12**2+3.0*g11*
. g13-3.0*g12*g13+2.0*g13**2)/(-g11+g12)**3/(-g12+g13)**3+
. s1(g11,g12,g13)*(-g11**2-3.0*g11*g12+2.0*g12**2-3.0*g11*
. g13+3.0*g12*g13+2.0*g13**2)/(-g11+g12)**3/(-g11+g13)**3
return
end

c
real*8 function a3(g11,g12,g13)
real*8 g11,g12,g13,s1,s2,s3,s11,s22,s33
a3=-s11(g11,g12,g13)*g12*g13/(-g11+g12)**2/(-g11+g13)**2-g11*
. s22(g11,g12,g13)*g13/(-g11+g12)**2/(-g12+g13)**2-g11*
. s33(g11,g12,g13)*g12/(-g11+g13)**2/(-g12+g13)**2+
. s2(g11,g12,g13)*(g11**2*g12-2.0*g11*g12**2+g12**3+g11**2
. *g13-g11*g12*g13-2.0*g12**2*g13+g11*g13**2+g12*g13**2)/
. (-g11+g12)**3/(-g12+g13)**3-s1(g11,g12,g13)*(g11**3-2.0
. *g11**2*g12+g11*g12**2-2.0*g11**2*g13-g11*g12*g13+g12**2*g13
. +g11*g13**2+g12*g13**2)/(-g11+g12)**3/(-g11+g13)**3-s3(g11,
. g12,g13)*(g11**2*g12+g11*g12**2+g11**2*g13-g11*g12*g13+g12**2*
. g13-2.0*g11*g13**2-2.0*g12*g13**2+g13**3)/(-g11+g13)**3
. /(-g12+g13)**3

```

c

```

real*8 function a4(g11,g12,g13)
real*8 g11,g12,g13,s1,s2,s3,s11,s22,s33
a4=-s33(g11,g12,g13)*(g11+g12)**2/(-g11+g13)**2/(-g12+g13)**2-
. s22(g11,g12,g13)*(g11+g13)**2/(-g11+g12)**2/(-g12+g13)**2
. -s11(g11,g12,g13)*(g12+g13)**2/
. (-g11+g12)**2/(-g11+g13)**2+2.0*s2(g11,g12,g13)*(g11+g13)*
. (g11**2-g11*g12-g12**2+g11*g13-g12*g13+g13**2)/(-g11+
. g12)**3/(-g12+g13)**3-2.0*s1(g11,g12,g13)*(g12+g13)*(-g11**2-
. g11*g12+g12**2-g11*g13+g12*g13+g13**2)/(-g11+g12)**3/(-g11+g13
. )**3+2.0*s3(g11,g12,g13)*(g11+g12)*(-g11**2-g11*g12-g12**2+g11*
. g13+g12*g13+g13**2)/(-g11+g13)**3/(-g12+g13)**3
return
end

```

c

```

real*8 function a5(g11,g12,g13)
real*8 g11,g12,g13,s1,s2,s3,s11,s22,s33
a5=g11*s33(g11,g12,g13)*g12*(g11+g12)/(-g11+g13)**2/(-g12+g13)
. **2+g11*s22(g11,g12,g13)*g13*(g11+g13)/(-g11+g12)**2/(-g12+g13)
. **2+s11(g11,g12,g13)*g12*g13*(g12+g13)/(-g11+g12)**2/(-g11+g13)
. **2+s3(g11,g12,g13)*(g11**3*g12+g11**2*g12**2+g11*g12**3+g11**
. 3*g13+g11**2*g12*g13+g11*g12**2*g13+g12**3*g13-2.0*g11**2*g13**
. 2-5.0*g11*g12*g13**2-2.0*g12**2*g13**2+g11*g13**3+g12*g13**3)/
. (-g11+g13)**3/(-g12+g13)**3-s2(g11,g12,g13)*(g11**3*g12-2.0*g11
. **2*g12**2+g11*g12**3+g11**3*g13+g11**2*g12*g13-5.0*g11*g12**2*
. g13+g12**3*g13+g11**2*g13**2+g11*g12*g13**2-2.0*g12**2*g13**2+
. g11*g13**3+g12*g13**3)/(-g11+g12)**3/(-g12+g13)**3+s1(g11,g12,
. g13)*(g11**3*g12-2.0*g11**2*g12**2+g11*g12**3+g11**3*g13-5.0*
. g11**2*g12*g13+g11*g12**2*g13+g12**3*g13-2.0*g11**2*g13**2+g11
. +g12*g13**2+g12**2*g13**2+g11*g13**3+g12*g13**3)/(-g11+g12)**3/
. (-g11+g13)**3
return
end

```

c


```

real*8 function a6(gl1,gl2,gl3)
real*8 gl1,gl2,gl3,s1,s2,s3,s11,s22,s33
a6=-s11(gl1,gl2,gl3)*gl2**2*gl3**2/(-gl1+gl2)**2/(-gl1+gl3)**2
. -gl1**2*s22(gl1,gl2,gl3)*gl3**2/(-gl1+gl2)**2/(-gl2+gl3)**2-gl1
. **2*s33(gl1,gl2,gl3)*gl2**2/(-gl1+gl3)**2/(-gl2+gl3)**2-2.0*gl1
. *s3(gl1,gl2,gl3)*gl2*gl3*(gl1**2+gl1*gl2+gl2**2-2.0*gl1*gl3-2.0*
. gl2*gl3+gl3**2)/(-gl1+gl3)**3/(-gl2+gl3)**3+2.0*gl1*s2(gl1,gl2,
. gl3)*gl2*gl3*(gl1**2-2.0*gl1*gl2+gl2**2+gl1*gl3-2.0*gl2*gl3+
. gl3**2)/(-gl1+gl2)**3/(-gl2+gl3)**3-2.0*gl1*s1(gl1,gl2,gl3)*gl2*
. gl3*(gl1**2-2.0*gl1*gl2+gl2**2-2.0*gl1*gl3+gl2*gl3+gl3**2)/
. (-gl1+gl2)**3/(-gl1+gl3)**3
return
end

```

c
c The s1,s2,s3,s11,s22,s33 are derivatives of W
c

```

function s1(gl1,gl2,gl3)
common /param/ x1,x2,y1,y2
real*8 x1,x2,y1,y2
real *8 s1,gl1,gl2,gl3
s1=2.0*(gl1**(-1+y1)*x1*y1+gl1**(-1+y2)*x2*y2)
return
end

```

c
function s2(gl1,gl2,gl3)
common /param/ x1,x2,y1,y2
real*8 x1,x2,y1,y2
real*8 s2,gl1,gl2,gl3
s2=2.0*(gl2**(-1+y1)*x1*y1+gl2**(-1+y2)*x2*y2)
return
end

c
function s3(gl1,gl2,gl3)
common /param/ x1,x2,y1,y2
real*8 x1,x2,y1,y2
real*8 s3,gl1,gl2,gl3
s3=2.0*(gl3**(-1+y1)*x1*y1+gl3**(-1+y2)*x2*y2)
return
end

c

```

function s11(gl1,gl2,gl3)
common /param/ x1,x2,y1,y2
real*8 x1,x2,y1,y2
real*8 s11,gl1,gl2,gl3
s11=2.0*(gl1**(-2+y1)*x1*(-1.0+y1)*y1+gl1**(-2+y2)*x2*
(-1.0+y2)*y2)
return
end

```

c

```

function s22 (gl1,gl2,gl3)
common /param/ x1,x2,y1,y2
real*8 x1,x2,y1,y2
real*8 s22, gl1,gl2,gl3
s22=2.0*(gl2**(-2+y1)*x1*(-1.0+y1)*y1+gl2**(-2+y2)*x2*
(-1.0+y2)*y2)
return
end

```

c

```

function s33(gl1,gl2,gl3)
common /param/ x1,x2,y1,y2
real*8 x1,x2,y1,y2
real*8 s33,gl1,gl2,gl3
s33=2.0*(gl3**(-2+y1)*x1*(-1.0+y1)*y1+gl3**(-2+y2)*x2*
(-1.0+y2)*y2)
return
end

```

c

```

subroutine compsd2(gl1,gl2,delt,delt4,cm,ts,td)
real*8 gl1,gl2,ts(3,3),td(3,3,3,3)
real*8 cm(3,3),delt(3,3),delt4(3,3,3,3),pl(3,3,3,3)
real*8 q2(3,3,3,3),q1(3,3,3,3),p(3,3,3,3),q(3,3,3,3)
real*8 b1,b2,b3, abar,bbar

```

```

c
c      Computes second order tensor ts(i,j) based on the formula
c      derived in code.
c
      do 25043 i=1,3
        do 25044 j=1,3
          ts(i,j)=abar(gl1,gl2)*cm(i,j)+bbar(gl1,gl2)*delt (i,j)
25044      continue
25043  continue
c
c      Call subroutine to get P, Q which are defined in code.
c
      call pqcom(cm,cm,p1,q)
      call pqcom(cm,delt,p,q1)
      call pqcom(delt,cm,p,q2)
c
c      Computes tensor td(i,j,k,l).
c
      do 25045 l=1,3
        do 25046 j=1,3
          do 25047 k=1,3
            do 25048 i=1,3
              td(i,j,k,l)=b1(gl1,gl2)*p1(i,j,k,l)+b2(gl1,gl2)*
                (q1(i,j,k,l)+q2(i,j,k,l))+b3(gl1,gl2)*delt4(i,j,k,l)
25048      continue
25047      continue
25046      continue
25045      continue
      return
      end
c
c      abar,bbar are b1, b2, b3 functions derived in code.
c
      real*8 function abar(gl1,gl2)
      real8 gl1,gl2,ss1,ss2
      abar=(-ss1(gl1,gl2)+ss2(gl1,gl2))/(-gl1+gl2)
      return
      end
c

```

```

      real*8 function bbar(gl1,gl2)
      real*8 gl1,gl2,ss1,ss2
      bbar=(-gl1*ss2(gl1,gl2)+ss1(gl1,gl2)*gl2)/(-gl1+gl2)
      return
      end

c
      real*8 function b1(gl1,gl2)
      real*8 gl1,gl2,ss1,ss2,ss11,ss22
      b1=2.0*ss1(gl1,gl2)/(-gl1+gl2)**3-2.0*ss2(gl1,gl2)/(-gl1+gl2)
      . **3+ss11(gl1,gl2)/(-gl1+gl2)**2+ss22(gl1,gl2)/(-gl1+gl2)**2
      return
      end

c
      real*8 function b2(gl1,gl2)
      real*8 gl1,gl2,ss1,ss2,ss11,ss22
      b2=-gl1*ss22(gl1,gl2)/(-gl1+gl2)**2-ss11(gl1,gl2)*gl2/
      . (-gl1+gl2)**2-ss1(gl1,gl2)*(gl1+gl2)/(-gl1+gl2)**3+ss2(gl1,
      . gl2)*(gl1+gl2)/(-gl1+gl2)**3
      return
      end

c
      real*8 function b3(gl1,gl2)
      real*8 gl1,gl2,ss1,ss2,ss11,ss22
      b3=2.0*gl1*ss1(gl1,gl2)*gl2/(-gl1+gl2)**3-2.0*gl1*ss2(gl1,gl2)
      . *gl2/(-gl1+gl2)**3+gl1**2*ss22(gl1,gl2)/(-gl1+gl2)**2+ss11(gl1,
      . gl2)*gl2**2/(-gl1+gl2)**2
      return
      end

c
c      ss1,ss2,ss11,ss22 are derivatives of W.
c
      function ss1(gl1,gl2)
      common /param/ x1,x2,y1,y2
      real*8 x1,x2,y1,y2
      real*8 ss1,gl1,gl2
      ss1=2.0*(gl1**(-1+y1)*x1*y1+gl1**(-1+y2)*x2*y2)
      return
      end

c

```

```

function ss2(gl1,gl2)
common /param/ x1,x2,y1,y2
real*8 x1,x2,y1,y2
real ss2,gl1,gl2
ss2=2.0*(2.0*gl2**(-1+y1)*x1*y1+2.0*gl2**(-1+y2)*x2*y2)
return
end

c

function ss11(gl1,gl2)
common /param/ x1,x2,y1,y2
real*8 x1,x2,y1,y2
real*8 ss11,gl1,gl2
ss11=2.0*(gl1**(-2+y1)*x1*(-1.0+y1)*y1+gl1**(-2+y2)*x2*
(-1.0+y2)
return
end

c

function ss22(gl1,gl2)
common /param/ x1,x2,y1,y2
real*8 x1,x2,y1,y2
real*8 ss22,gl1,gl2
ss22=2.0*(2.0*gl2**(-2+y1)*x1*(-1.0+y1)*y1+2.0*gl2**
(-2+y2)*x2*(-1.0+y2)*y2)
return
end

c

subroutine compsd3(gl1,delt,delt4,ts,td)
real*8 gl1,ts(3,3),td(3,3,3,3),delt(3,3),delt4(3,3,3,3)
real*8 cc1,abbar

c

do 25049 i=1,3
do 25050 j=1,3
ts(i,j)=abbar (gl1)*delt(i,j)
25050 continue
25049 continue

```

```

do 25051 i=1,3
  do 25052 j=1,3
    do 25053 k=1,3
      do 25054 l=1,3
        td(i,j,k,l)=cc1(gl1)*delt4(i,j,k,l)
25054      continue
25053    continue
25052  continue
25051 continue
c
  return
end
c
  real*8 function abbar(gl1)
  real*8 gl1
  abbar=sss1(gl1)
  return
end
c
  real*8 function cc1(gl1)
  real*8 gl1
  cc1=sss11(gl1)
  return
end
c
  function sss1(gl1)
  common /param/ x1,x2,y1,y2
  real*8 x1,x2,y1,y2
  real*8 sss1,gl1
  sss1=2.0*(3.0*gl1**(-1+y1)*x1*y1+3.0*gl1**(-1+y2)*x2*y2)
  return
end
c

```

```

function sss11(gl1)
common /param/ x1,x2,y1,y2
real*8 x1,x2,y1,y2
real*8 sss11,gl1
sss11=2.0*(3.0*gl1**(-2+y1)*x1*(-1.0+y1)*y1+3.0*gl1**
(-2+y2)*x2*(-1.0+y2)*y2)
return
end

c
c   This subroutine computes P and Q forth order tensors
c   which we define in tensor D
c

subroutine pqcom(cm1,cm2,p,q)
real*8 cm1(3,3),cm2(3,3),p(3,3,3,3),q(3,3,3,3)
do 100 i=1,3
  do 100 j=1,3
    do 100 k=1,3
      do 100 l=1,3
        p(i,j,k,l)=cm1(i,k)*cm2(j,l)+cm1(i,l)*cm2(j,k)
        q(i,j,k,l)=p(i,j,k,l)+cm1(j,l)*cm2(i,k)+cm1(j,k)
        *cm2(i,l)
100    continue
      return
    end

c
c   This subroutine computes matrix product cmXcm.
c

subroutine product(mat1,cmm)
real*8 mat1(3,3),cmm(3,3),sum
do 25 i=1,3
  do 25 j=1,3
    sum=0.0
    do 26 k=1,3
      sum=sum+mat1(i,k)*mat1(k,j)
    continue
    cmm(i,j)=sum
25  continue
  return
end

```

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