A RAMANUJAN-TYPE MEASURE FOR THE ASKEY-WILSON POLYNOMIALS

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Abstract

A Ramanujan-type representation for the Askey-Wilson q-beta integral, admitting the transformation \( q \to q^{-1} \), is obtained. Orthogonality of the Askey-Wilson polynomials with respect to a measure, entering into this representation, is proved. A simple way of evaluating the Askey-Wilson q-beta integral is also given.

1 Introduction.

The Askey-Wilson polynomials \( p_n(x; a, b, c, d|q) \) [1], which have already become classical, represent a five-parameter system of polynomials. They satisfy the orthogonality relation

\[
\int_{-1}^{1} p_m(x; a, b, c, d|q) p_n(x; a, b, c, d|q) w(x; a, b, c, d|q) dx = \delta_{mn} I_n(a, b, c, d|q)
\]  

(1.1)

with respect to the absolutely continuous measure \( d\alpha(x) = w(x)dx \), with the weight function

\[
w(x; a, b, c, d|q) = \frac{1}{\sin \theta \prod_{j=a,b,c,d} h(\cos \theta, v; q)}
\]

\[
\times \frac{h(\cos 2\theta, 1; q)}{x = \cos \theta,}
\]

(1.2)

\[
h(a, b; q) = \prod_{j=0}^{\infty} (1 - 2abq^j + b^2 q^{2j}).
\]

As special and limiting cases, the Askey-Wilson polynomials contain many known systems of polynomials (see, for example, [2]). In particular, the choice of the parameters \( a = -b = \sqrt{\beta}, c = -d = \sqrt{q\beta} \), leads to the continuous q-ultraspherical polynomials \( C_n(x; \beta|q) \) [3], i.e.,

\[
p_n(x; \sqrt{\beta}, -\sqrt{\beta}, \sqrt{q\beta}, -\sqrt{q\beta}|q) = \frac{(\beta^2; q)_{2n}(q; q)_n}{(\beta, \beta^2; q)_n} C_n(x; \beta|q),
\]

(1.3)
where we have used the standard notation of the theory of $q$-special functions

$$(a; q)_n = \prod_{j=0}^{n-1} (1 - aq^j), \quad (a_1, ..., a_k; q)_n = \prod_{j=1}^{k} (a_j; q)_n. \quad (1.4)$$

In turn, from $C_n(x; \beta|q)$ one can obtain the continuous $q$-Hermite polynomials $H_n(x|q) = (q; q)_n C_n(x; 0|q)$, the Gegenbauer (ultraspherical) polynomials $C_n^\lambda(x) = \lim_{\lambda \to 1} C_n(x; q^\lambda|q)$, and also the Chebyshev polynomials of the first and second kinds, $T_n(x)$ and $U_n(x)$, by taking the limit $\beta \to 1$ or by putting $\beta = q$ in $C_n(x; \beta|q)$, respectively.

The key ingredient of the original proof of the orthogonality (1.1), which led to the discovery of the Askey-Wilson system of polynomials (see the discussion of this point in [4]), was the evaluation of the Askey-Wilson $q$-beta integral:

$$I_0(a, b, c, d|q) = \int_{-1}^{1} w(x; a, b, c, d|q) dx = \frac{2\pi(abcd; q)_\infty}{(q, ab, ac, ad, bc, bd, cd; q)_\infty}, \quad \text{max}_{v=a,b,c,d}|v| < 1, \quad |q| < 1. \quad (1.5)$$

The integral (1.5) has acquired its name because in a special case, when the parameters $a = q^{\alpha+1/2}$, $b = -q^\beta+1/2$, and $c = -d = q^{1/2}$, the $q \to 1^-$ limit of $I_0(a, b, c, d|q)$ is the beta function (or Euler's integral of the first kind)

$$\int_{-1}^{1} (1 - x)^\alpha(1 + x)^\beta dx = 2^{\alpha+\beta+1} B(\alpha + 1, \beta + 1) = 2^{\alpha+\beta+1} \frac{\Gamma(\alpha + 1)\Gamma(\beta + 1)}{\Gamma(\alpha + \beta + 2)}. \quad (1.6)$$

A nonstandard form of the orthogonality on the full real line for the continuous $q$-Hermite polynomials $H_n(\sin \kappa x|q)$, $q = \exp(-2\kappa^2)$, was considered in [5]. It turned out that if one uses the modular transformation and the periodicity property of the $\vartheta$-function appearing in the weight function for these polynomials, the finite interval of orthogonality can be transformed into an infinite one. This technique is of interest both from a mathematical point of view and from the point of view of possible applications in theoretical physics, beginning with a number of problems, related with $q$-oscillators (see the review [6]).

The purpose of this article is to discuss the applicability of this idea to the more general case, i.e. to the Askey-Wilson $q$-beta integral (1.5) [7, 8]. To simplify consideration it will be assumed in Sections 2-4 that $|v| < 1$, $v = a, b, c, d$, and that the parameter $q = \exp(-2\kappa^2)$ satisfies the requirement $0 < q < 1$. The possibility of extending these results to other values of the parameters is discussed in Section 5.

2 A Ramanujan-type representation for the $q$-beta integral.

From the point of view of symmetry the parametrization $x = \sin \varphi$ is most convenient; it corresponds to the change of variable $\theta = \frac{\pi}{2} - \varphi$, $-\frac{\pi}{2} \leq \varphi \leq \frac{\pi}{2}$ in formula (1.2). Consequently, the left
The numerator of (1.5) is equal to

\[ I_0(a, b, c, d|q) = \int_{-\pi/2}^{\pi/2} \frac{h(-\cos 2\varphi, 1; q)}{\prod_{\nu=a,d,c,d} h(\sin \varphi, v; q)} d\varphi. \] (2.1)

Comparison of the numerator

\[ h(-\cos 2\varphi, 1; q) = \prod_{j=0}^{\infty} (1 + 2q^j \cos 2\varphi + q^{2j}) \]

of the integral (2.1) with Jacobi's expression for the theta-function \( \vartheta_2(z, q) \equiv \vartheta_2(z|\tau) \), \( q = \exp(\pi i \tau) \) as an infinite product [9]

\[ \vartheta_2(z, q) = 2q^{1/4}(q^2; q^2)_\infty \cos \frac{z}{2} \prod_{j=1}^{\infty} (1 + 2q^2 j \cos 2z + q^4 j), \] (2.2)

shows that

\[ h(-\cos 2\varphi, 1; q) = \frac{2 \cos \varphi}{q^{1/8}(q, q)_\infty} \vartheta_2(\varphi, q^{1/2}) \] (2.3)

and therefore

\[ I_0(a, b, c, d|q) = \frac{2}{q^{1/8}(q, q)_\infty} \int_{-\pi/2}^{\pi/2} \frac{\vartheta_2(\varphi, q^{1/2}) \cos \varphi}{\prod_{\nu=a,b,c,d} h(\sin \varphi, v; q)} d\varphi. \] (2.4)

With the aid of the modular transformation [9]

\[ \vartheta_2(z|\tau) = \exp\left(\frac{-iz^2}{\tau}\right) \vartheta_4(z, q^{1/2}) e^{-\pi z^2/\tau}, \quad \tau = \frac{i\kappa^2}{\pi}, \] (2.5)

and the change of variable \( \varphi = \kappa x \), the integral (2.4) can be written as

\[ I_0(a, b, c, d|q) = \frac{2\sqrt{\pi}}{q^{1/8}(q, q)_\infty} \int_{-\pi/2}^{\pi/2} \frac{\vartheta_4(\frac{\pi i x}{\kappa}, e^{-\pi^2/\kappa^2}) e^{-x^2} \cos \kappa x}{\prod_{\nu=a,b,c,d} h(\sin \kappa x, v; q)} dx. \] (2.6)

Using the expansion

\[ \vartheta_4(z, q) = \sum_{k=-\infty}^{\infty} (-1)^k q^{k^2} e^{2ikz} \] (2.7)

and taking into account the uniform convergence of the series (2.7) in any bounded domain of values of \( z \) [9], we substitute (2.7) into (2.6) and integrate this series termwise, i.e.,

\[ I_0(a, b, c, d|q) = \frac{2\sqrt{\pi}}{q^{1/8}(q, q)_\infty} \sum_{k=-\infty}^{\infty} (-1)^k \int_{-\pi/2\kappa}^{\pi/2\kappa} \frac{e^{-(x+\pi/\kappa)^2} \cos \kappa x dx}{\prod_{\nu=a,b,c,d} h(\sin \kappa x, v; q)}. \] (2.8)
The change of variable \( x_k = x + \frac{\pi}{\kappa} k, \quad x_k^{\min} = \frac{\pi}{\kappa}(k - \frac{1}{2}) \leq x_k \leq \frac{\pi}{\kappa}(k + \frac{1}{2}) = x_k^{\max} \) and an account for the relation \( x_k^{\max} = x_k^{\min} \) allows to sum the right-hand side of (2.8) with respect to \( k \) and represent (2.8) in the form

\[
I_0(a, b, c, d|q) = \frac{2\sqrt{\pi}}{q^{1/8}(q; q)_\infty} \int _0^\infty \frac{e^{-x^2} \cos \kappa x dx}{\prod _{v=a,b,c,d} h(\sin \kappa x, v; q)}. \tag{2.9}
\]

Thus, combining formulas (1.5) and (2.9) yields the following representation for the Askey-Wilson q-beta integral [7]

\[
\tilde{I}_0(a, b, c, d|q) \equiv \int _{-\infty} ^\infty \rho(\kappa x; a, b, c, d|q) e^{-x^2} \cos \kappa x dx = \frac{\sqrt{\pi} q^{1/4}(abcd; q)_\infty}{(ab, ac, ad, bc, bd, cd; q)_\infty}, \tag{2.10}
\]

where, in accordance with the definition (1.2),

\[
\rho(x; a, b, c, d|q) = \prod _{v=a,b,c,d} h^{-1}(\sin x, v; q) = \prod _{v=a,b,c,d} e_q(i v e^{-i\theta})e_q(-i v e^{i\theta}), \tag{2.11}
\]

and \( e_q(z) = (z; q)^{-1} \) is the q-exponential function [2].

We note that each factor \( h^{-1}(\sin \kappa x, v; q), \; v = a, b, c, d, \) in the integrand (2.10) is represented as

\[
(\sin \kappa x, v; q) = \sum _{n=0} ^\infty (iv)_n \sum _{k=0} ^n \frac{(-1)^k \exp[-i(n-2k)\kappa x]}{(v; q)_k(q; q)_{n-k}}, \tag{2.12}
\]

if one uses the generating function for the continuous q-Hermite polynomials \( H_n(x|q) \)

\[
(te^{i\theta}, te^{-i\theta}; q)^{-1} = \sum _{n=0} ^\infty \frac{H_n(\cos \theta|q)}{(q; q)_n} t^n, \quad |t| < 1, \tag{2.13}
\]

and their explicit representation [2]

\[
H_n(\cos \theta|q) = \sum _{k=0} ^n \left[ \begin{array}{c} n \\ k \end{array} \right]_q e^{i(n-2k)\theta}, \tag{2.14}
\]

where the symbol \( \left[ \begin{array}{c} n \\ k \end{array} \right]_q \) denotes the q-binomial coefficient [2]. Therefore the integration over \( x \) in (2.10) is reduced to the Fourier transformation formula for the ground state of the linear harmonic oscillator

\[
\frac{1}{\sqrt{2\pi}} \int _{-\infty} ^\infty \exp(-x^2/2 + ixy) dx = \exp(-y^2/2). \tag{2.15}
\]

An explicit evaluation of the nonstandard form of the Askey-Wilson q-beta integral (2.10) will be discussed in greater detail in Section 4.

As mentioned above, the weight function (1.2) with the parameters \( a = -b = \beta^{1/2}, \; c = -d = aq^{1/2}, \) corresponds to the continuous q-ultraspherical polynomials \( C_n(x; \beta|q). \) The relations [2]

\[
(a; q)_\infty = (a, aq; q^2)_\infty, \quad (a, -a; q)_\infty = (a^2, q^2)_\infty,
\]

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enable the representation (2.10) for this particular case to be simplified to

\[
\int_{-\infty}^{\infty} \frac{\exp(-x^2 + i\kappa x)dx}{(-\beta \exp(2i\kappa x), -\beta \exp(-2i\kappa x); q)_\infty} = \frac{\sqrt{\pi} q^{1/4}(\beta, q\beta; q)_\infty}{(\beta^2; q)_\infty}.
\] (2.16)

If one compares (2.16) with the Ramanujan integral \( q = \exp(-2k^2), |q| < 1 \) [10, 11]

\[
\int_{-\infty}^{\infty} e^{-x^2 + 2m-x} e_q(a^{1/2} e^{2i\kappa x})e_q(b^{1/2} e^{-2i\kappa x}) dx = \frac{\sqrt{\pi} e^{m^2}}{(q ab; q)_\infty} E_q(a q e^{2imk}) E_q(b q e^{-2imk}),
\] (2.17)

it is easy to verify that (2.16) agrees with (2.17) if one sets \( 2m = ik = i\kappa \) and \( a = b = -\beta q^{1/2} \).

3 Orthogonality of the Askey-Wilson polynomials with respect to the measure \( \rho(\kappa x; a, b, c, d | q) \).

A direct proof of the orthogonality for the Askey-Wilson polynomials

\[
\int_{-\infty}^{\infty} p_m(\sin \kappa x; a, b, c, d | q)p_n(\sin \kappa x; a, b, c, d | q) \rho(\kappa x; a, b, c, d | q) \exp(-x^2) \cos \kappa x dx =
\]

\[= \delta_{mn} \tilde{I}_n(a, b, c, d | q), \] (3.1)

with respect to the weight function appearing in the nonstandard integral representation (2.10), is analogous to the proof of eigenfunctions orthogonality for the Sturm-Liouville differential equation [12]. Indeed, the difference differentiation formula for the Askey-Wilson polynomials [1]

\[
\sin \kappa \partial_x p_n(\sin \kappa x; a, b, c, d | q) =
\]

\[= q^{-n/2}(1 - q^n)(1 - abcd q^{n-1}) \cos \kappa x p_{n-1}(\sin \kappa x; aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2} | q) \]

provides a lowering operator for these polynomials. To find a raising operator one can use the relation

\[
w(\sin \varphi; a, b, c, d | q) = \frac{2\partial_2(\varphi, q^{1/2})}{q^{1/8}(q; q)_\infty} \rho(\varphi; a, b, c, d | q), \] (3.3)

which follows from (1.2), (2.3) and (2.11), and write the difference equation for the Askey-Wilson polynomials [1] in the form

\[
\sin \kappa \partial_x \left[ \frac{\partial_2(\kappa x, q^{1/2})}{\cos \kappa x} \rho(\kappa x; aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2} | q) \sin \kappa \partial_x p_n(\sin \kappa x; a, b, c, d | q) \right] =
\]

\[= (1 - q^{-n})(1 - abcd q^{n-1}) \cos \kappa x \partial_2(\kappa x, q^{1/2}) \rho(\kappa x; a, b, c, d | q) p_n(\sin \kappa x; a, b, c, d | q). \] (3.4)
Now, using the difference differentiation formula (3.2) in the left-hand side of (3.4) and the periodicity property of the $\vartheta_2$-function [9],

$$\vartheta_2(z \pm \pi \tau, q) = q^{-1} \exp(\pm 2i\tau)\vartheta_2(z, q), \quad q = \exp(\pi i \tau), \quad (3.5)$$

we arrive at the raising operator

$$p_n(sin \kappa x; aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2}|q) = q^{\frac{1-m}{2}} \cos \kappa x \rho(\kappa x; \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}|q). \quad (3.6)$$

We are now in a position to give a direct proof of the orthogonality relation (3.1). We multiply both sides of the equality (3.6) by $p_m(sin \kappa x; a, b, c, d|q) \exp(-x^2)$ and integrate in $x$ over the full real line. As a result we obtain in the right-hand side,

$$q^{\frac{1-n}{2}} \int_{-\infty}^{\infty} p_m(sin \kappa x; a, b, c, d|q) p_n(sin \kappa x; a, b, c, d|q) \rho(\kappa x; a, b, c, d|q) e^{-x^2} \cos \kappa x \ dx \equiv$$

$$q^{\frac{1-n}{2}} I_{mn}(a, b, c, d|q). \quad (3.7)$$

The left-hand side

$$\int_{-\infty}^{\infty} dx p_m(sin \kappa x; a, b, c, d|q) e^{-x^2} (\sin 2\kappa x \cos \kappa \partial_x - \cos 2\kappa x \sin \kappa \partial_x)$$

$$\rho(\kappa x; aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2}|q) p_{n-1}(sin \kappa x; aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2}|q), \quad (3.8)$$

can be integrated by parts with the aid of (3.2) and the evident relations

$$\int_{-\infty}^{\infty} dx f(x) \cos \kappa \partial_x \varphi(x) = \int_{-\infty}^{\infty} dx \varphi(x) \cos \kappa \partial_x f(x), \quad (3.9)$$

$$\int_{-\infty}^{\infty} dx f(x) \sin \kappa \partial_x \varphi(x) = - \int_{-\infty}^{\infty} dx \varphi(x) \sin \kappa \partial_x f(x),$$

which apply to (3.8) because the function $\rho(\kappa x; aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2}|q)$ has no singularities inside of the strip $-\kappa \leq y \leq \kappa$, $-\infty < x < \infty$ in the complex plane $z = x + iy$. This leads to

$$q^{\frac{1-m}{2}} (1 - q^m)(1 - abcdq^{m-1}) I_{m-1n-1}(aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2}|q). \quad (3.10)$$

Equating the right-hand (3.7) and left-hand (3.10) sides thus yields

$$q^{\frac{m-n}{2}} I_{mn}(a, b, c, d|q) = (1 - q^m)(1 - abcdq^{m-1}) I_{m-1n-1}(aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2}|q). \quad (3.11)$$
We now interchange $m$ and $n$ in (3.11) and take into account that the integral $I_{mn}(a,b,c,d|q)$ is symmetric in $m$ and $n$ due to the definition (3.7), i.e.,

$$q^{n-m}I_{mn}(a,b,c,d|q) = (1-q^n)(1-abcdq^{n-1})I_{m-1n-1}(aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2}|q).$$

(3.11')

Finally, multiplying both sides of (3.11) by $(1-q^n)(1-abcdq^{n-1})$ and of (3.11') by $(1-q^m)(1-abcdq^{m-1})$ and subtracting the second expression from the first, we get

$$(q^{n-m} - q^{m-n})(1-abcdq^{m+n-1})I_{mn}(a,b,c,d|q) = 0.$$  

(3.12)

From (3.12) it follows that $I_{mn}(a,b,c,d|q) = \delta_{mn}I_n(a,b,c,d|q)$, confirming the orthogonality (3.1) of the Askey-Wilson polynomials for $m \neq n$ [8].

We note that as special and limiting cases, (3.1) contains the orthogonality relations for other known sets of polynomials, such as the continuous $q$-ultraspherical polynomials $C_n(x; \beta|q)$, the continuous $q$-Hermite polynomials $H_n(x; q) = (q;q)_nC_n(x; 0|q)$ (the corresponding special case of (3.1), when the all parameters $a, b, c, d$ are equal to zero, is considered in [5] ), the Chebyshev polynomials of the first and second kinds, $T_n(x)$ and $U_n(x)$, and so on.

4 Evaluation of the integrals $\tilde{I}_n(a, b, c, d|q)$.

Iterating the recurrence relation

$$\tilde{I}_n(a, b, c, d|q) = (1-q^n)(1-abcdq^{n-1})\tilde{I}_{n-1}(aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2}|q),$$

(4.1)

which follows from (3.11) or (3.11') when $m = n$, allows to express the normalization integrals $\tilde{I}_n(a, b, c, d|q)$, $n = 1, 2, ..., $ through a known value of the Askey-Wilson $q$-beta integral $\tilde{I}_0(a, b, c, d|q)$, i.e.

$$\tilde{I}_n(a, b, c, d|q) = \frac{(q, ab, ac, ad, bc, bd, cd; q)_n}{(1-abcdq^{2n-1})(abcd; q)_{n-1}} \tilde{I}_0(a, b, c, d|q).$$

(4.2)

It only remains to evaluate the integral $\tilde{I}_0(a, b, c, d|q)$ itself. To this end, having defined the symmetrical $\rho_+(x)$ and antisymmetrical $\rho_-(x)$ combinations with respect to the inversion $x \rightarrow -x$,

$$\rho_\pm(x; a, b, c, d|q) = \frac{1}{2} [\rho(x; a, b, c, d|q) \pm \rho(-x; a, b, c, d|q)],$$

(4.3)

it is convenient to rewrite (2.10) as

$$\tilde{I}_0(a, b, c, d|q) = \int_{-\infty}^{\infty} dx \exp(-x^2 + i\kappa x)\rho_+(\kappa x; a, b, c, d|q).$$

(4.4)

Let us carry out the replacements $v \rightarrow v \sqrt{q}$, $v = a, b, c, d$, and the subsequent shift of the variable of integration $x \rightarrow x + i\kappa$ in (4.4). (We remind that the function $\rho(\kappa z; aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2}|q)$ does not have singularities in the strip $-\kappa \leq y \leq \kappa$, $-\infty < x < \infty$ of the complex plane $z = x + iy$ ). Then, taking into account that in accordance with the definitions (1.2) and (2.11)

$$\rho(\kappa(x + i\kappa); aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2}|q) = \rho(\kappa x; a, b, c, d|q) \prod_{v=a,b,c,d} (1 + iv \exp(i\kappa x)),$$

(4.5)
we obtain

\[ \tilde{I}_0(aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2}|q) = (1 - s_2)\tilde{I}_0(a, b, c, d|q) + \]

\[ + s_4 \int_{-\infty}^{\infty} dx \exp(-x^2 + 3i\kappa x)\rho_+(\kappa x; a, b, c, d|q) - is_3 \int_{-\infty}^{\infty} dx \exp(-x^2 + 2i\kappa x)\rho_-(\kappa x; a, b, c, d|q), \]

where

\[ s_2 = ab + ac + ad + bc + bd + cd, \]

\[ s_3 = abc + abd + acd + bcd, \]

\[ s_4 = abcd. \]

It remains only to express the second and third integrals in the right-hand side of (4.6) in terms of \( \tilde{I}_0(a, b, c, d|q) \). To that end one can use the \( n = 1 \) case of (3.6)

\[ (\sin 2\kappa x \cos \kappa \partial_x - \cos 2\kappa x \sin \kappa \partial_x)\rho(\kappa x; aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2}|q) = \]

\[ = [(1 - s_4) \sin 2\kappa x + (s_3 - s_1) \cos \kappa x]\rho(\kappa x; a, b, c, d|q), \]

taking into account that \( p_0(x; a, b, c, d|q) = 1, \rho_1(x; a, b, c, d|q) = 2(1 - s_4)x + s_3 - s_1 \) and \( s_1 = a + b + c + d \). The symmetrization of (4.8) leads to the relations

\[ (\sin 2\kappa x \cos \kappa \partial_x - \cos 2\kappa x \sin \kappa \partial_x)\rho_\pm(\kappa x; aq^{1/2}, bq^{1/2}, c q^{1/2}, d q^{1/2}|q) = \]

\[ = (1 - s_4) \sin 2\kappa x \rho_\pm(\kappa x; a, b, c, d|q) + (s_3 - s_1) \cos \kappa x \rho_\mp(\kappa x; a, b, c, d|q). \]

Multiplying both sides of the equality (4.9) for the antisymmetrical combination \( \rho_-(\kappa x) \) by \( \exp(-x^2) \) and integrating over the variable \( x \) yields

\[ (1 - s_4) \int_{-\infty}^{\infty} dx \exp(-x^2 + 2i\kappa x)\rho_-(\kappa x; a, b, c, d|q) = i(s_1 - s_3)\tilde{I}_0(a, b, c, d|q). \]

Now we multiply both sides of (4.9) for \( \rho_+(\kappa x; aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2}|q) \) by \( \exp(-x^2 + i\kappa x) \) and integrate over \( x \). Using (4.10), the result can be written as

\[ \int_{-\infty}^{\infty} dx \exp(-x^2 + 3i\kappa x)\rho_+(\kappa x; a, b, c, d|q) = \]

\[ = \left[ 1 - \frac{(s_3 - s_1)^2}{(1 - s_4)^2} \right] \tilde{I}_0(a, b, c, d|q) - \frac{1 - q}{1 - s_4} \tilde{I}_0(aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2}|q). \]
Substituting (4.10) and (4.11) into (4.6), we find

\[(1 - abcd)(1 - qabcd) \tilde{I}_0(aq^{1/2}, bq^{1/2}, cq^{1/2}, dq^{1/2}|q) = \]

\[(1 - ab)(1 - ac)(1 - ad)(1 - bc)(1 - bd)(1 - cd) \tilde{I}_0(a, b, c, d|q). \]

Since 0 < q < 1, by iterating equation (4.12) one can express the Askey-Wilson q-beta integral (2.10) with arbitrary parameters in terms of its value for vanishing parameters a, b, c, d, i.e.,

\[\tilde{I}_0(a, b, c, d|q) = \frac{(abcd; q)_\infty}{(ab, ac, ad, bc, bd, cd; q)_\infty} \tilde{I}_0(0, 0, 0, 0|q). \quad (4.13)\]

The value of \(\tilde{I}_0(0, 0, 0, 0|q)\) is easily found from (2.10) and (3.1) with the aid of the Fourier transformation formula (2.15) for the quadratically decreasing exponential function, i.e.,

\[\tilde{I}_0(0, 0, 0, 0|q) = \int_{-\infty}^{\infty} dx \exp(-x^2 + i\kappa x) = \sqrt{\pi} q^{1/8}. \quad (4.14)\]

Combining formulas (4.13) and (4.14) leads to

\[\tilde{I}_0(a, b, c, d|q) = \frac{\sqrt{\pi} q^{1/8}(abcd; q)_\infty}{(ab, ac, ad, bc, bd, cd; q)_\infty}, \quad (4.15)\]

which is the known value of the Askey-Wilson q-beta integral [1]

\[I_0(a, b, c, d|q) = \frac{2\sqrt{\pi}}{q^{1/8}(q; q)_\infty} \tilde{I}_0(a, b, c, d|q) = \frac{2\pi(abcd; q)_\infty}{(q, ab, ac, ad, bc, bd, cd; q)_\infty}. \quad (4.15')\]

Substituting (4.15) into (4.2), we finally obtain the explicit form for the normalization integral

\[\tilde{I}_n(a, d, c, d|q) = \frac{\sqrt{\pi} q^{1/8}(q; q)_n(abcdq^{n+1}; q)_\infty}{(1 - abcdq^{n+1})(abq^n, acq^n, adq^n, bdq^n, cdq^n; q)_\infty}. \quad (4.16)\]

The complications arising in the evaluation of the standard form of the Askey-Wilson q-beta integral (1.5) can be illustrated by the following short quotation from reference [4]: "This was surprisingly hard, and it has taken over five years before relatively simple ways of evaluating this integral were found".

5 The transformation \(q \to q^{-1}\).

It is necessary to emphasize that the nonstandard orthogonality relation (3.1) admits the transformation \(q \to q^{-1}\) [7, 8]. The standard form of the Askey-Wilson integral (1.5) does not in general have this property. Even in the simplest case of vanishing parameters \(a, b, c\) and \(d\), which corresponds to the continuous \(q\)-Hermite polynomials \(H_n(x|q)\), the definition of a weight function for the system of polynomials \(h_n(x; q) = i^{-n}H_n(ix|q^{-1})\) requires a special analysis [13, 14].
Since 

\[(z; q^{-1})_\infty = (qz; q)_{-1}, \tag{5.1}\]

the change \(q \to q^{-1}\) (i.e. \(\kappa \to i\kappa\)) in the function \(\rho(\kappa x; a, b, c, d|q)\) appearing in (2.10) and (3.1), transforms it into

\[
\rho(i\kappa x; a, b, c, d|q^{-1}) = \prod_{\nu=\alpha, \beta, \gamma, \delta} (ivqe^{\kappa x}, -ivqe^{-\kappa x}; q)_\infty = \prod_{\nu=\alpha, \beta, \gamma, \delta} E_q(ivqe^{\kappa x})E_q(-ivqe^{\kappa x}), \tag{5.2}\]

where \(E_q(z) = e_q^{-1}(-z) = (-z; q)_\infty\) [2]. Therefore, under the transformation \(q \to q^{-1}\), the orthogonality relation (3.1) for the Askey-Wilson polynomials with the parameter \(q < 1\) converts into the following orthogonality relation for the Askey-Wilson polynomials with \(q > 1\):

\[
\int_{-\infty}^{\infty} p_m(i \sinh \kappa x; a, b, c, d|q^{-1}) p_n(i \sinh \kappa x; a, b, c, d|q^{-1}) e^{-x^2} \cosh \kappa x dx =
\delta_{mn} f_n(a, b, c, d|q^{-1}) \tag{5.3}\]

The explicit form of \(f_n(a, b, c, d|q^{-1})\) is readily obtained from (4.16), upon making use of the formulas (5.1) and \((a; q^{-1})_n = (a^{-1}; q)_n(-a)^n q^{-n(n-1)/2}\) [2].

On the other hand, with the aid of the explicit representation for the Askey-Wilson polynomials [1, 2]

\[
p_n(\cos \varphi; a, b, c, d|q) = (ab, ac, ad; q)_n a^{-n} \Phi_3 \left[ q^{-n}, abcdq^{n-1}, iae^\varphi, -iae^{-i\varphi} \right]_{ab, ac, ad} \tag{5.4}\]

and the inversion formula (with respect to the transformation \(q \to q^{-1}\)) for the basic hypergeometric series \(\Phi_3\) (see [2], p.21, exercise 1.4(i)), it is easy to show that

\[
p_n(x; a, b, c, d|q^{-1}) = (-1)^n(abcd)^{-3/2n(n-1)} p_n(x; a^{-1}, b^{-1}, c^{-1}, d^{-1}|q). \tag{5.5}\]

Consequently, from (5.3) and (5.5) it follows the orthogonality relation

\[
\int_{-\infty}^{\infty} p_m(i \sinh \kappa x; a^{-1}, b^{-1}, c^{-1}, d^{-1}|q) p_n(i \sinh \kappa x; a^{-1}, b^{-1}, c^{-1}, d^{-1}|q) \rho(i\kappa x; a, b, c, d|q^{-1})\]

\[
e^{-x^2} \cosh \kappa x dx = \frac{(q, 1/ab, 1/ac, 1/ad, 1/bc, 1/bd, 1/cd, 1/d; q)_n}{(1-q^{2n-1}/abcd)(1/abcd; q)_{n-1}} f_0(a, b, c, d|q^{-1}) \delta_{mn} \tag{5.6}\]

for the Askey-Wilson polynomials with the parameters \(|v| > 1, v = a, b, c, d\) and \(0 < q < 1\). The value of the integral \(f_0(a, b, c, d|q^{-1})\) is simple to obtain from (4.15) by means of the formula (5.1).
6 Concluding remarks.

The orthogonality relations (3.1) and (5.6) are bound to be related by the Fourier transformation for the Askey-Wilson functions, analogous to the well-known transformation for the harmonic oscillator wave functions $H_n(x) \exp(-x^2/2)$ (or Hermite functions in the terminology of mathematicians [15, 16]) connecting the coordinate and momentum realizations in quantum mechanics. It should be interesting to compare this Fourier transformation with the $q$-transformations, that reproduce the Askey-Wilson polynomials [17, 18]. For the $q$-Hermite functions $H_n(\sin \kappa x|q) \exp(-x^2/2)$, $q = \exp(-2\kappa^2)$, which are the simplest case of the Askey-Wilson functions with vanishing parameters $a$, $b$, $c$, and $d$, such Fourier transformation has the form [5]

$$\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \exp(ixy - x^2/2)H_n(\sin \kappa x|q)dx = i^nq^{n^2/4}h_n(\sinh \kappa y|q)\exp(-y^2/2).$$

The general case needs to be analyzed in greater detail.

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