

Near-Infrared Spectral Geometric Albedos of Charon and Pluto: Constraints on Charon's Surface Composition

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The spectral geometric albedos of Charon and Pluto are derived at near-infrared wavelengths (1.4–2.5 μm) from measurements obtained in 1987. Comparisons of these to theoretical calculations are used to place constraints on the identity and relative abundances of surface ices on Charon. These comparisons suggest that widespread regions of pure CH_4 ice do not occur on Charon and that if CH_4 is abundant on Charon then it is large grained (≈ 5 mm) and is likely mixed at the granular level with H_2O ice, and possibly CO_2 ice. © 1996 Academic Press, Inc.

INTRODUCTION

Pluto and its moon Charon may represent two huge icy planetesimals that condensed from the solar nebula in a region where $T < 50$ K (Owen *et al.* 1993) and escaped accretion by the giant planets. They may therefore provide a sampling of the primitive stage in the transition from interstellar materials to the planets and satellites in our Solar System. The Pluto–Charon system also has the highest ratio of the radius of the satellite to the radius of the planet (≈ 0.5), about a factor of 2 greater than that of the Earth–Moon system (≈ 0.27). This may indicate that Pluto and Charon have undergone significantly different internal evolutionary sequences. For example, the Moon is depleted in volatiles relative to the Earth. This is believed to be due to the extensive melting of the lunar surface and interior early in its history (e.g., Warren 1985). Thus, insight into the original composition and subsequent evolu-

tion of Pluto and Charon may be provided by knowledge of their current surface composition.

While Pluto and Charon are in close proximity, they appear to have distinctly different surface compositions. Recent near-infrared (near-IR) spectroscopy of Pluto has revealed that the surface is dominated by N_2 ice with minor amounts of CH_4 and CO ices (Owen *et al.* 1993). Near-IR spectra of Charon have been interpreted as indicating the surface is dominated by H_2O ice (Buie *et al.* 1987, Marcialis *et al.* 1987), although recent theoretical calculations suggest a significant non- H_2O ice may be present as well (Roush 1994). Roush (1994) found that several calculated spectra of H_2O – CO_2 , H_2O – CH_4 , and H_2O – CO_2 – CH_4 mixtures could reproduce the shape of the observed Charon spectrum within the errors associated with the Charon data. Comparisons of the calculated spectral geometric albedo (p_λ) of these mixtures to the broadband $p_{2.2\mu\text{m}}$ of Charon (Bosh *et al.* 1992) eliminated only a few of the mixtures from further consideration (Roush 1994). Greater constraints could be placed on the surface composition if p_λ were known for Charon at higher spectral resolution over a broader range of near-IR wavelengths.

In this paper we derive near-IR values of p_λ for Charon from previous observations and provide a comparison to the calculated p_λ from Roush (1994). In the next section we discuss the telescopic observations, outline our assumptions, and present the values of p_λ we derive for Charon. In the following section, we compare our derived values of p_λ to the previous calculations of Roush (1994). Finally, we summarize our results and conclusions.

† Deceased June 13, 1994.

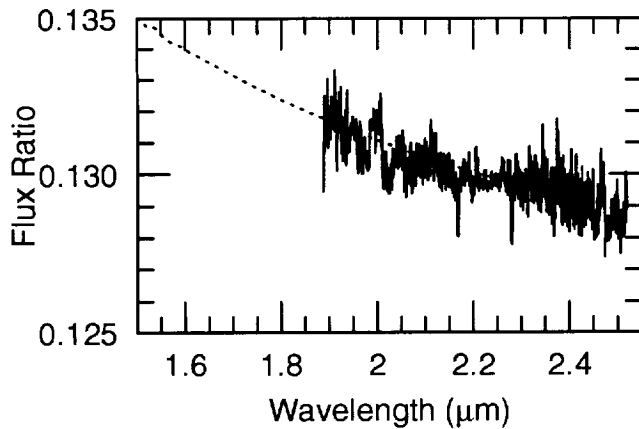


FIG. 1. The flux of SAO120599 relative to BS5384 (solid line) and the flux of a 6030-K blackbody relative to a 5950-K blackbody, scaled to the observational data (dotted line). The blackbody ratio is used to extrapolate the observational data to 1.65 μm .

OBSERVATIONS AND DERIVATION OF CHARON'S SPECTRAL GEOMETRIC ALBEDO

Observations. We base our study on the observations of Charon originally obtained by Buie *et al.* (1987) (referred to hereafter as BEA). The details regarding their acquisition, reduction, and original interpretation are presented in BEA. At the same time BEA also obtained spectral observations of Pluto. For completeness, and to provide another entry in the record of Pluto observations, we derive p_λ for Pluto below. However, because the focus of this paper is Charon we do not include any discussion or comparisons to the Pluto data. Both of these data were presented as differences in stellar magnitudes, relative to the standard star SAO120599. Digital versions of all the data were kindly provided by Marc Buie (Buie, personal communication, 1993).

Derivation of stellar magnitudes. To derive p_λ for Pluto and Charon, the near-IR magnitudes of the standard star must be determined. SAO120599 (hereafter referred to as SAO) is listed as a G0V stellar type with a visual magnitude of 8.414 ± 0.031 (Turon *et al.* 1992). A literature search for color properties of this star in the near-IR was unsuccessful. In their derivation of the near-IR fluxes of Pluto and Charon, Marcialis *et al.* (1987) assumed the near-IR colors of SAO were "solar-like." Because this star is presently close to Pluto in the sky ($\alpha = 14^\circ 41' 18.77''$; $\delta = +03^\circ 26' 56.5''$; J2000), it is useful to determine its near-IR spectral energy distribution for future observations. Instead of assuming color properties for SAO, spectral observations of SAO and BS5384 (hereafter referred to as BS) were obtained in May 1995. Figure 1 shows the flux of SAO relative to BS as a function of wavelength in

the region 1.88 to 2.52 μm . The H and K photometric magnitudes of BS are well documented (Gezari *et al.* 1993) and after reviewing the various values we adopt $m_{\text{BS,H}} = 4.74 \pm 0.03$ and $m_{\text{BS,K}} = 4.68 \pm 0.03$. The magnitude of the object of interest can be determined via

$$m_o = -2.5 \log_{10} \left(\frac{f_o}{f_s} \right) + m_s, \quad (1)$$

where m_o and m_s are the magnitudes of the object and the standard being used, respectively, and the quantity in the brackets is the flux of the object relative to the standard. At 2.2 μm , the monochromatic wavelength of the K photometric filter, the value of $f_{\text{SAO}}/f_{\text{BS}}$ is 0.130 ± 0.003 , giving $m_{\text{SAO,K}} = 6.90 \pm 0.03$. Poor seeing conditions and other observational priorities prevented observations of SAO at shorter wavelengths. However, the BEA data extend to wavelengths shorter than those of the observational data, so in order to evaluate these data we extrapolate the data in Fig. 1 to the H photometric region (1.65 μm). We use a ratio of blackbodies representing the effective temperature (T_e) of SAO and BS. SAO is a G0V stellar type, so we use $T_e = 6030$ K (Lang 1991) and BS is a G1V stellar type, so $6030 > T_e > 5860$ K (Lang 1991). We require the blackbody ratio to reproduce the shape of the observational data. The dotted line in Fig. 1 is a curve corresponding to $(f_{6030}/f_{5950}) - C = f_{\text{obs}}$, where the subscripts denote the black body temperature and C is a constant used to scale the flux ratio to the observational data ($C = 0.89821$). The dotted line is used to extrapolate the observed flux ratio to 1.65 μm and using this value (0.134 ± 0.003) yields $m_{\text{SAO,H}} = 6.92 \pm 0.03$. The magnitude of the object (Pluto

TABLE I
Magnitudes and Spectral Geometric Albedos of
Charon and Pluto

Wavelength (μm)	m_c	p_c	m_p	p_p
1.504	14.97 ± 0.21	0.23 ± 0.05	12.30 ± 0.03	0.77 ± 0.04
1.595	15.00 ± 0.14	0.22 ± 0.03	12.72 ± 0.03	0.52 ± 0.03
1.684	14.69 ± 0.10	0.30 ± 0.03	12.80 ± 0.03	0.49 ± 0.03
1.771	14.57 ± 0.13	0.33 ± 0.05	12.74 ± 0.03	0.51 ± 0.03
1.858	14.62 ± 0.19	0.30 ± 0.06	12.42 ± 0.03	0.65 ± 0.04
1.944	14.96 ± 0.17	0.22 ± 0.04	12.36 ± 0.03	0.69 ± 0.04
2.028	15.42 ± 0.19	0.14 ± 0.03	12.42 ± 0.03	0.65 ± 0.04
2.111	14.97 ± 0.11	0.22 ± 0.03	12.74 ± 0.03	0.48 ± 0.03
2.193	14.64 ± 0.09	0.30 ± 0.03	13.16 ± 0.04	0.33 ± 0.02
2.274	14.71 ± 0.08	0.28 ± 0.03	13.60 ± 0.03	0.22 ± 0.01
2.353	14.89 ± 0.15	0.24 ± 0.04	13.58 ± 0.05	0.22 ± 0.01
2.432	14.92 ± 0.22	0.23 ± 0.05	13.00 ± 0.04	0.38 ± 0.02
2.509	14.85 ± 0.67	0.24 ± 0.15	12.71 ± 0.08	0.50 ± 0.04

or Charon, m_o) is determined via

$$m_o = \Delta + m_{SAO}, \quad (2)$$

where Δ is the magnitude difference as presented in BEA and m_{SAO} is the magnitude of SAO. We used $m_{SAO,H}$ for wavelengths $\leq 1.8 \mu\text{m}$ and $m_{SAO,K}$ for wavelengths $> 1.8 \mu\text{m}$. H and K solar magnitudes of -28.188 ± 0.030 and -28.248 ± 0.030 , respectively, were obtained from Campins *et al.* (1985). Rearranging Eq. (1), the ratio of the flux from an object to that of the Sun (at 1 AU) can be expressed as

$$\frac{F_o}{F_s} = 10^{-0.4(m_o - m_s)}, \quad (3)$$

Derivation of spectral geometric albedos. The spectral geometric albedo is calculated at each wavelength via

$$\rho_\lambda = \frac{R^2 d^2}{r^2} \frac{F_o}{F_s}, \quad (4)$$

where R is the heliocentric distance (in AU), d is the geocentric distance (in AU), r is the radius of the object (in AU), and F_o/F_s is determined via Eq. (3). Values used in Eq. (4) included radii of 1164 ± 22.9 and 621 ± 20.6 km (converted to AU), for Pluto and Charon respectively (Young and Binzel 1994), and 29.6991 and 28.7411 AU for the calculated helio- and geocentric distances on April 23, 1987, when the BEA observations were made at a rotational phase angle (ϕ) of 0.75. At $\phi = 0.25$, Charon is

between the Earth and Pluto, while at $\phi = 0.75$ Pluto is between the Earth and Charon. At $\phi = 0.5$ and 1.0 (or 0.0) Charon is at southern and northern elongation in its orbit, respectively. The derived magnitudes and spectral geometric albedos are tabulated for each wavelength in Table I, and the spectral geometric albedos are shown graphically in Fig. 2a.

COMPARISONS TO PREVIOUS STUDIES

Fink and DiSanti (1988) and Marcialis *et al.* (1992) found that $p_\lambda \sim 0.45$ near $1 \mu\text{m}$ for Charon ($\phi = 0.75$). In this paper we find that $p_\lambda \approx 0.20 \pm 0.05$ near $1.5 \mu\text{m}$, indicating a significant difference from the $1\text{-}\mu\text{m}$ value, although this is not surprising since H_2O ice has an absorption near $1.5 \mu\text{m}$. Narrowband filter observations of Charon ($\phi = 0.75$, Marcialis *et al.* 1987) were used to derive values of p_λ for Charon at four wavelengths assuming a radius of 593 km (Marcialis *et al.* 1992). Using a radius of 602 km for Charon, Bosh *et al.* (1992) report the $p_{2.2 \mu\text{m}}$ of Charon at rotational phase angles (ϕ) of 0.42 and 0.06. All of these values have been scaled to the Charon radius used here, 621 km, and are compared in Fig. 2b. Inspection of Fig. 2b indicates that the p_λ derived here are comparable, within the errors, to the values derived by Marcialis *et al.* (1992) with the notable exception of the $1.5\text{-}\mu\text{m}$ value. It is difficult to ascribe the magnitude of this discrepancy to the different assumptions regarding the near-IR colors of SAO. In fact, Marcialis *et al.* (1987) note some concerns related to the data point at $1.5 \mu\text{m}$. The $2.2\text{-}\mu\text{m}$ value derived here is higher than either value presented by Bosh *et al.* (1992). Convolution of the data presented

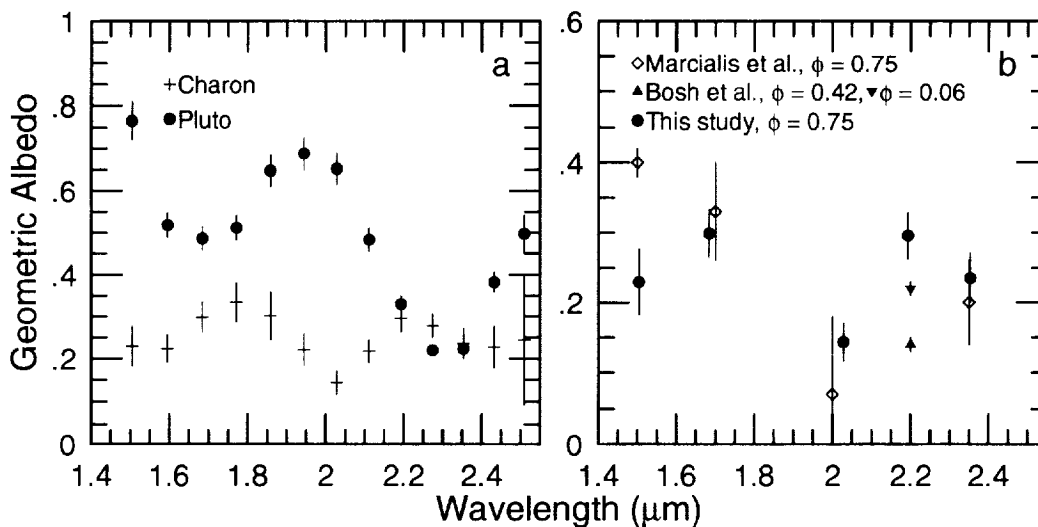


FIG. 2. (a) Spectral geometric albedo of Charon (+) using a radius of 621 ± 20.6 km, and Pluto (●), using a radius of 1164 ± 22.9 km. Both data sets were derived using the data of Buie *et al.* 1987). (b) Comparison of the spectral geometric albedo values derived here from the Buie *et al.* (1987) data (●) to values previously reported for Charon: ▲ and ▼ from Bosh *et al.* (1992); and ◇ from Marcialis *et al.* (1992).

here, using the nominal K photometric filter response given in the IRTF photometry manual (Tokunaga 1988), yields a broadband $p_{2.2 \mu\text{m}}$ of 0.23 ± 0.03 for $\phi = 0.75$ and thus is comparable to the value given by Bosh *et al.* (1992) for $\phi = 0.06$. Buie and Shriver (1994) present differential magnitudes ($m_c - m_p$) for six filter band passes, including the standard H and K photometric filters, as a function of Charon's orbital phase, but did not obtain any data at $\phi = 0.75$ (at 180° East Longitude, using their terminology). In the H filter the differential magnitudes of obtained by Buie and Shriver (1994) range from 1.5 to 2.9 while in the K filter the range is from 1.56 to 2.4. Convolution of the magnitude values in Table I with the nominal IRTF filter responses yields differential magnitudes of 2.17 ± 0.14 and 1.90 ± 0.12 at H and K, respectively. These values fall within the range of values defined by the Buie and Shriver data (1994).

COMPOSITIONAL IMPLICATIONS OF CHARON'S SPECTRAL GEOMETRIC ALBEDO

Roush (1994) calculated the p_λ using scattering theory for several mixtures of volatile ices representing the surface of Charon. These models included two- and three-component spatial and intimate mixtures consisting of H_2O , CO_2 , and CH_4 ices. Here we compare p_λ to some of these mixtures to see if we can further constrain the surface composition of Charon.

A comparison of calculated spatial mixtures to the Charon data is provided in Fig. 3a. The solid line in Fig. 2a is the best comparison to the Charon data and it corresponds to a spatial mixture of H_2O and CO_2 , having grain sizes of 200 and 50 μm and relative areal coverages of 80 and 20%, respectively. The other two curves appear to be too bright at the shortest wavelengths, suggesting that the presence of significant occurrences of discrete patches of CH_4 ice are unlikely on Charon.

The same comparison for intimate mixtures is presented in Fig. 3b. In this case, two mixtures appear to be comparable to the Charon data, the solid and dashed lines. The solid line corresponds to an intimate mixture of H_2O and CO_2 , both having a grain size of 100 μm and relative mass fractions of 50%. The dashed line corresponds to a mixture of H_2O , CO_2 , and CH_4 ices with grain sizes of 100, 100, and 5000 μm and relative mass fractions of 35, 35, and 30%, respectively. This suggests that if CH_4 ice is present on Charon, then it likely occurs as large particles mixed at the granular level with H_2O , and perhaps CO_2 , ices.

SUMMARY

Using the spectral reflectance measurements of Pluto and Charon obtained by Buie *et al.* (1987), and the radii

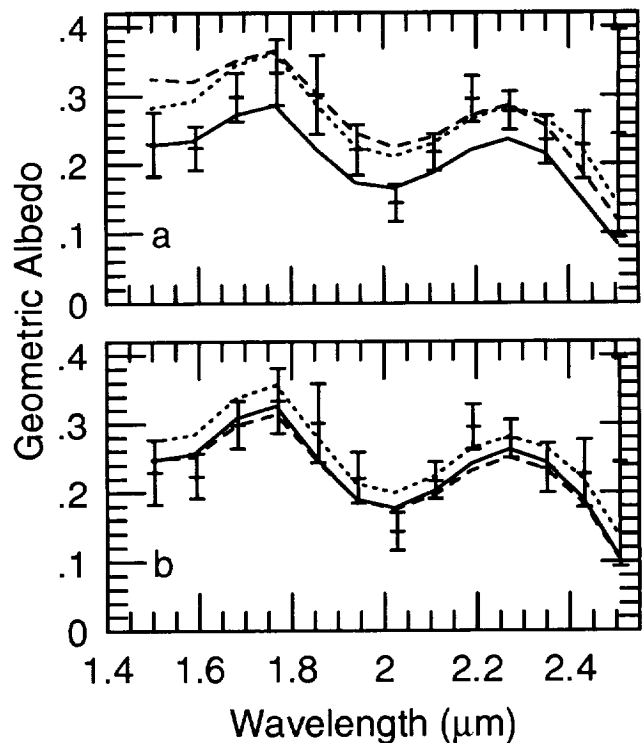


FIG. 3. (a) Comparison of derived spectral geometric albedo of Charon to theoretical calculations of spatial mixtures representing the surface of Charon. The solid line is a mixture of H_2O and CO_2 ices having grain sizes of 200 and 50 μm , and relative areal coverages of 80 and 20%, respectively. The dotted line is a mixture of H_2O and CH_4 ice having grain sizes of 50 and 25 μm , and relative areal coverages of 95 and 5%, respectively. The dashed line is a mixture of H_2O , CO_2 , and CH_4 ice having grain sizes of 100, 100, and 10 μm and relative areal coverages of 80, 19, and 1%, respectively. (b) Comparison of derived spectral geometric albedo of Charon to theoretical calculations of intimate mixtures representing the surface of Charon. The solid line is a mixture of H_2O and CO_2 ices each having a grain size of 100 μm , and relative mass fractions of 50%. The dotted line is a mixture of H_2O and CH_4 ice each having grain sizes of 50 μm , and relative mass fractions of 95 and 5%, respectively. The dashed line is a mixture of H_2O , CO_2 , and CH_4 ice having a grain sizes of 100, 100, and 5000 μm and relative mass fractions of 35, 35, and 30%, respectively.

determined by Young and Binzel (1994), we have derived the spectral geometric albedo of these two objects. We find that the values derived here are comparable to values previously reported at similar near-IR wavelengths (Bosh *et al.* 1992, Marcialis *et al.* 1992).

Comparison of the derived spectral geometric albedo of Charon to theoretical calculations of spatial mixtures of volatile ices suggest that widespread isolated patches of pure CH_4 do not occur on Charon. However, the presence of modest amounts ($\approx 20\%$) of isolated patches of CO_2 ice remains plausible. Comparisons involving theoretical calculations of intimate mixtures suggest significant

amounts of CO₂ (50 wt.%) and CH₄ (30 wt.%) ices may be present, provided the CH₄ ice grain size is quite large (≈ 5 mm). Further progress in determining the surface composition of Charon requires high spectral resolution observations of the satellite spatially resolved from Pluto.

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