NASA/TM-1999-208855



# High Temperature Mechanical Characterization and Analysis of $A1_2O_3/A1_2O_3$ Composites

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February 1999

# Acknowledgments

We would like to thank Mr. John A. Setlock of Case Western Reserve University (CWRU) for the processing of the specimens, scanning electron microscopy, and fiber test data. Also, we would like to thank Dr. Ali Sayir of CWRU for the high temperature fiber moduli data used in this paper.

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# $\label{eq:high-temperature-mechanical-characterization} High-Temperature-Mechanical-Characterization and Analysis of Al_2O_3/Al_2O_3 \ Composites$

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#### Abstract

Sixteen ply unidirectional zirconia coated single crystal  $Al_2O_3$  fiber reinforced polycrystalline  $Al_2O_3$  was tested in uniaxial tension at temperatures to 1400°C in air. Fiber volume fractions ranged from 26 to 31%. The matrix has primarily open porosity of approximately 40%. Theories for predicting the Young's modulus, first matrix cracking stress, and ultimate strength were applied and evaluated for suitability in predicting the mechanical behavior of  $Al_2O_3/Al_2O_3$  composites.

The composite exhibited pseudo tough behavior (increased area under the stress/strain curve relative to monolithic alumina) from 22° to 1400°C. The rule-of-mixtures provides a good estimate of the Young's modulus of the composite using the constituent properties from room temperature to approximately 1200°C for short term static tensile tests in air. The ACK theory provides the best approximation of the first matrix cracking stress while accounting for residual stresses at room temperature. Difficulties in determining the fiber/matrix interfacial shear stress at high temperatures prevented the accurate prediction of the first matrix cracking stress above room temperature. The theory of Cao and Thouless, based on Weibull statistics, gave the best prediction for the composite ultimate tensile strength.

#### 1. Introduction

Ceramics in general are extremely brittle, have low strain tolerance, and exhibit a wide variation in ultimate strength. The observed scatter in strength is caused by an abundance of imperfections, i.e., flaws, that are a result of material processing. Over the years, the strength reliability of monolithic ceramics have improved as better processing techniques have evolved. However, as Dev (1992), Taylor (1991), and Zweben (1998) indicate, the brittle failure characteristics of these materials make them acceptable in only a limited range of applications. Even in the limited structural application of monolithic ceramics to turbines in automotive turbochargers, turbo efficiency is sacrificed for structural reliability as noted by Yoshida and Kokji (1989). In an effort to increase ceramic toughness and strength, ceramic matrix composites with various reinforcements are being developed. These developments are covered briefly by King (1989) and Levine (1992). These composites may include multiple phases or matrices with particulates, whiskers, or continuous fibers. In the gas turbine industry, ceramic matrix composites (CMCs) are particularly attractive since they have the potential to replace nickel based superalloys in various hot section components (Dix and Petty (1990) and Constance (1990)). The primary attribute of CMCs relative to nickel based superalloys is the ability of CMCs to be used well beyond current turbine service temperatures, as well as to withstand more severe operating environments. This would enable engines to be operated at higher temperatures with near-stoichiometric combustion without cooling air requirement penalties as noted by Drascovich (1993). Increasing operating temperature is a classic approach for improving turbine efficiency. Zweben (1998) notes that the low density and high temperature properties of CMCs make them potential materials in hot structures and propulsion systems of reusable launch vehicles.

Another attraction of ceramic composites is their relative low density, which is, as Holmes and Wu (1995) point out, typically 65% to 75% lower than conventional superalloys. One can not overemphasize the fact that weight is a critical design facet for gas turbines utilized in aeropropulsion. Lastly, CMCs offer the potential of increased durability, relative to superalloys, at the high operating temperatures. This would result in increased time between engine overhauls reducing operating costs.

In virtually all ceramic matrix composite systems the goal of the materials scientist is to apply a closing pressure on existing matrix crack surfaces and to impart a tortuous fracture path by crack deflection. This results in an increase in the apparent toughness of the material as noted by Warren (1992). Unlike polymer and metal matrix composites, the fiber/matrix interface in a fiber reinforced ceramic composite must be relatively weak. Optimization of the interface prevents matrix cracks from propagating through the fibers while still providing load transfer. As a result, unbroken fibers bridge a propagating matrix crack, which increases the composite work of fracture. In essence, the fiber/matrix interface has to be strong enough to allow load transfer and retain acceptable strength in the transverse direction, but the interface must also allow debonding as a crack passes around the fiber.

There are many ceramic matrix composite systems being investigated and developed, as indicated by Sheppard (1992) and Studt (1991). One of these systems consists of zirconia coated continuous single crystal alumina fibers and an alumina matrix ( $Al_2O_3/Al_2O_3$ ). As an oxide, these composites don't have the oxidation problems of silicon based ceramics.

The path to successful commercialization of the CMC systems mentioned above must include the characterization and evaluation of engineering design properties. As Duffy and Gyekenyesi (1995) point out this requires characterization of mechanical and thermal properties. It is essential to characterize the creep behavior of these materials as well as ascertain how temperature affects their fast fracture. Motivated by the general lack of high quality test data required by the design engineer, this paper presents high-temperature, fast fracture data of the aforementioned  $Al_2O_3/Al_2O_3$  composite. This data was acquired using a high temperature tensile testing facility assembled by Gyekenyesi (1998) at NASA LeRC. The properties of these material systems were obtained in the primary fiber direction.

The objective of this paper is to determine selected CMC mechanical properties and predict the mechanical behavior with existing theoretical models. Flat  $Al_2O_3/Al_2O_3$  composite specimens were tested. The continuous monofilament fibers in the composite were oriented in the primary direction, which presents an upper bound for strength, stiffness, and work of fracture properties for this composite system. For this study unidirectional  $Al_2O_3/Al_2O_3$  composites were tested from room temperature to 1400°C (2550°F) in air. The experimental data were compared with theoretical predictions for stiffness, proportional limit or first matrix cracking stress, and ultimate strength.

# 2. Equipment

The tensile test equipment used for this experiment was setup to test flat ceramic matrix composites at temperatures to  $1500^{\circ}$ C (2730°F) in air. Most of the equipment was previously described by Gyekenyesi (1998), with the exception of a new furnace and grips that have been incorporated since then. A brief summary is presented here.

The universal tensile testing frame has a mechanical central screw actuator. The machine's actuator has a 100 kN (22 kip) load capacity and a displacement range of 100 mm (4 in.). It is digitally controlled, allowing closed-loop control with load, strain, or crosshead displacement.

A 50 kN (11 kip) load cell was used with the frame. The load cell's output was calibrated such that 10 volts corresponded with a 20 kN (4.5 kips) load. This output was used as input for the computerized data acquisition system.

Water cooled, hydraulically actuated wedge grips were used to hold the specimen in the test frame. These grips are mounted to a rigid load train with a permanent alignment device. The grip faces have a grit surface. Thermocouples are fastened to the grip faces to monitor the temperature near the ends of the specimen.

A low contact force mechanical extensometer was utilized for the high temperature tests to measure axial strain. This device employs a variable capacitor to convert displacements into electrical signals. Low contact forces minimize the bending in test specimens. The silicon carbide rods with knife edges are gas cooled along with the body of the extensometer. The gage length is fixed at 25 mm (0.98 in.).

A clip-on extensioneter was used for the room temperature tests to measure axial strain. The gage works in conjunction with the frame's electronic controls. As a result, the frame's controls are used to electronically calibrate the gage and process its signal. The gage has a fixed gage length of 12.7 mm (0.5 in.) with a 12.7 mm (0.5 in.) extender for a total gage length of 25.4 mm (1.0 in.). The maximum extension of this gage is  $\pm 1.27$  mm ( $\pm 0.05$  in.). This results in a maximum strain range of  $\pm 5.0\%$ . Rubber bands were used to hold the gage to the specimen.

Specimen heating was accomplished with the use of an induction heating system and a recrystalized silicon carbide cylindrical susceptor. The susceptor has a 51 mm (2.00 in.) inner diameter with a 60 mm (2.38 in.) outer diameter. The furnace is similar to the system presented by Worthem and Lewinsohn (1991) with the exception that the power supply is a slightly smaller bench top model. The specimen is passed through open slots in the side of the susceptor. The extensometer rods pass through another slot. The ends of the cylinder are plugged with silicon carbide disks. A thermocouple, placed near the specimen surface is used to monitor the gage section temperature.

Lastly, a computerized data acquisition and analysis system was used to collect the test data. The software also was used to determine various material properties from the collected data.

#### 3. Test Specimen Configuration

The  $Al_2O_3/Al_2O_3$  composite specimens had a fiber volume fraction ranging from 26% to 31%. Fibers consisted of uniaxially aligned single crystal alumina fibers. The matrix was polycrystalline alumina with primarily open porosity of 40%. The specimen geometry consisted of 16 plies. The fabrication of the composite is summarized by Jaskowiak, et al. (1997).

The C-axis sapphire fibers have a mean diameter of 132 microns (5.2 mils). An unstabilized zirconia coating of 1.5 to 2.0 µm (59-79 µin.) was applied to the fibers to act as an interface between the fibers and the matrix. The zirconia coating has a similar porosity to that of the alumina of the matrix. The resulting volume fraction of the interface is slightly greater than 1% of the composite.

The composite specimens used in this research were flat, constant thickness, straight sided coupons with bonded tabs on either end. Glass fiber/epoxy composite tabs were used for room temperature tests, and carbon fiber/polyimide composite tabs were utilized in high temperature tests. These tabs provided a compliant layer between the specimen and the grip faces. The tabs were bonded to the specimen using structural film adhesives. The structural film adhesives require curing temperatures of 120°C (250°F) and 180°C (350°F) for the room temperature and high temperature specimens, respectively.

Specimen lengths were 152 mm (6.0 in.). The specimen width was approximately 13 mm (0.5 in.) for all tests. The thickness was approximately 1.7 mm (0.07 in.). Tensile coupons were cut from a panel with a diamond impregnated abrasive wheel.

The specimen geometry was designed based on the American Society for Testing and Materials (ASTM) Standard C 1359-96. The specimen geometry is shown in Figure 1. ASTM Standard C 1359-96 is a recently developed standard test method for obtaining tensile properties of continuous fiber-reinforced ceramics with rectangular cross-sections at high temperatures. The configuration has been used successfully in the past with failures usually occurring within the gage section. In this research effort the carbon fiber/polyimide composite tabs were 32 mm (1.25 in.) long, 13 mm (0.5 in.) wide, and 1.5 mm (0.060 in.) thick.

In addition to the composite specimens, monolithic alumina coupons were fabricated. These specimens had the same geometry as the composite coupons, that is, straight sided flat coupons with bonded tabs on each end. Each coupon was 152 mm (6.0 in.) long with a width of 13 mm (0.5 in.).

#### 4. Procedure

Due to the limited availability of ceramic composite specimens, only two to four coupons were used for each of the test conditions.

The room temperature tests utilized clip-on gages. The clip-on gages were shunt calibrated using the tensile testing frame's controls. The gage was calibrated such that 5.0 percent strain resulted in a 10 volt output. The gage was used Figure 1. Flat straight sided specimen with the 25 mm (1.0 in.) gage length. The high with bonded tabs.



temperature tests used the capacitive extensometers. The mechanical extensometer was manually calibrated such that 10 volts corresponded to 1.0 percent strain.

Tests were performed on a universal, digitally controlled, tensile testing frame with an electric actuator. They were conducted in the load control mode which provides a constant change in load with respect to time. The loading rate was fixed at 1.0 kN/min. (4.4 kips/min.). The 1.0 kN/min. (4.4 kips/min.) provides a fast enough strain rate for the specimen to limit the effects of creep and oxidation, but slow enough to manually monitor and respond to any problems (i.e. with gripping or extensometry) that may arise during a test. The load cell output was shunt calibrated with the output set at 20 kN equal to 10 V.

The specimen was loaded by passing it through the slot in the furnace then pulling the furnace between the grips while holding the specimen with tongs. The grips were clamped onto the specimen and the frame was placed in load control to maintain a no load condition. Next, the extensometer was inserted into the hot zone without contacting the specimen. The furnace was brought up to the desired test temperature at a rate of  $30^{\circ}$ C per minute (54°F/min.). Following stabilization of the test coupon temperature, the extensometer was brought into contact with the sample and the tensile test was started.

Once the tensile tests were completed several fractured composite specimens were observed with a scanning electron microscope (SEM) to determine the average matrix crack spacing. Matrix cracks and their locations within the gage length were recorded. The mean crack spacing and its standard deviation were determined.

# 5. Theory

Considerable effort has been applied to the mathematical modeling of the mechanical behavior of ceramic matrix composites. A complete understanding of the mechanical behavior of these composites is necessary if designers are to make use of them for high temperature engine applications. Important properties include the stiffness, first matrix cracking stress, ultimate strength, and work of fracture from room temperature to at least 1200°C.

# 5.1. Modulus

One of the basic properties of interest is the elastic modulus. The longitudinal Young's modulus of the composite is related to the matrix and fiber moduli, and their respective volume fractions by the rule of mixtures as presented in the following equation

$$E_{c} = V_{f}E_{f} + V_{m}E_{m}$$
<sup>(1)</sup>

where E is the modulus, V is the volume fraction, and subscripts c, f, and m refer to the composite, fiber, and matrix, respectively.

#### 5.2. Matrix Cracking and Interfacial Shear Properties

In addition to studying the composite modulus, the proportional limit, in the direction of the fibers, needs to be investigated. It is assumed that the proportional limit on a stress/strain curve and the first matrix cracking stress are the same. This is a common assumption, as noted by Woodford and his associates (1993). The first matrix crack is taken by definition as the first through the cross section crack, wherein only the fibers are left to carry the total composite load. Any minor cracking within the composite before this condition or load is reached, is ignored in this definition. It is assumed that the fiber failure stress and strain are greater than the matrix failure stress and strain, respectively. The first matrix cracking stress is the onset of permanent major damage in a composite making it a critical design parameter. Shimansky (1989) describes the cracking process for CMCs with an increasing load as usually initiating with microcracks within an amorphous region at the fiber/matrix interface. The microcracks coalesce upon further increase in the applied load, forming major matrix cracks perpendicular to the loading axis that eventually traverse the whole composite cross-section. As the applied load continues to increase the matrix cracks at other points along the load axis. Finally, the matrix becomes saturated with regularly spaced parallel cracks. Once the composite matrix is saturated with cracks the remaining fiber/matrix interface area with each matrix segment is insufficient to transfer an adequate load to cause the matrix to fracture into smaller segments.

The first matrix cracking stress is dependent on many parameters. A key parameter is the fiber/matrix interfacial shear strength. Kerans (1989) has noted that the interfacial shear strength is difficult to characterize and may vary with location. In this work, the fiber/matrix interfacial shear strength is determined from matrix crack spacing measurements in a composite that has been loaded to or near its ultimate strength. Loading a specimen near the ultimate strength leads to matrix crack saturation. The average crack spacing for each specimen is used to determine the mean fiber/matrix interfacial shear strength,  $\tau$ , using the equation

$$\tau = \frac{\beta R V_m E_m \sigma_y}{2 V_f E_c x}$$
(2)

where  $\beta$ =1.337 from Kimber and Keer (1982), R is the fiber radius,  $\sigma_y$  is the composite stress where matrix cracking initiates, and x is the length over which the additional load sustained by the fibers at the crack is transferred back to the matrix. This equation is derived from a summation of forces within the composite as presented by Aveston, et al. (1971). Aveston, et al. showed that the crack spacing is between x and 2x. Kimber and Keer (1982) demonstrated analytically that the crack spacing was closer to 1.337x. This equation is deterministic since it is assumed that the composite stress where matrix cracking initiates, that is  $\sigma_y$ , is constant.

In most composite systems there is a coefficient of thermal expansion (CTE) mismatch between the fibers and the matrix. This results in residual stresses within the composite at temperatures other than the processing temperature. The following equation from Budiansky, et al. (1986) is used to determine the axial residual stress within the matrix:

$$\sigma_{\mathrm{m}\alpha} = \left[ \mathrm{E}_{\mathrm{m}} \frac{\phi_2}{\phi_1} \right] \left[ \frac{\mathrm{E}_{\mathrm{f}}}{\mathrm{E}_{\mathrm{c}}} \right] \left[ \frac{\mathrm{V}_{\mathrm{f}}}{1 - v_{\mathrm{m}}} \right] \varepsilon_{\mathrm{T}}$$
(3)

where:

$$\boldsymbol{\mathcal{E}}_{\mathrm{T}} = (\boldsymbol{\alpha}_{\mathrm{f}} \cdot \boldsymbol{\alpha}_{\mathrm{m}}) \Delta \mathrm{T}$$
(4)

$$\Delta T = (T - T_{\text{proc}})$$
<sup>(5)</sup>

$$\phi_{1} = 1 - 0.5 \left[ \frac{1 - 2v}{1 - v} \right] \left[ 1 - \frac{E_{c}}{E_{f}} \right]$$
(6)

$$\phi_2 = 0.5 \left[ 1 + \frac{\mathrm{E_c}}{\mathrm{E_f}} \right] \tag{7}$$

The variables  $\alpha_f$  and  $\alpha_m$  are the coefficients of thermal expansion for the fiber and matrix, respectively. The processing temperature is  $T_{proc}$  and the test temperature is T. The Poisson's ratio, v, for the fiber and the matrix is assumed to be the same.

Aveston, et al. (1971) derived an equation for predicting the composite stress where matrix cracking initiates in brittle matrix composites. This is commonly referred to as the ACK theory and the result is presented in equation 8. The parameters include the fiber/matrix interfacial shear strength  $\tau$ , the matrix fracture surface energy  $\gamma_m$ , the fiber radius R, the constituent volume fractions V, and constituent moduli E.

$$\sigma_{y} = \left[\frac{12\tau\gamma_{m}E_{c}^{2}E_{f}V_{f}^{2}}{E_{m}^{2}V_{m}R}\right]^{\frac{1}{3}}$$
(8)

Note that the above equation is independent of crack size. It is a discrete model that was derived using an energy balance approach. The ACK theory uses the change in energy states within the composite from just prior to matrix crack initiation to just after the crack propagates completely through the matrix. Important assumptions included are: (1) the fiber failure strain is greater than the matrix failure strain, (2) a frictional interfacial bond between the fiber and the matrix exists, and (3) the fibers can bear the load without any support from the matrix.

Chawla (1993) pointed out some limitations or problems with the ACK theory. First, the theory indicates that the matrix strain to failure, or the first matrix cracking stress in equation 8, goes to zero as the fiber volume fraction goes to zero, whereas, the limit should be the strain or stress to failure of the monolithic matrix material. The model predicts that the strain to failure increases with increasing fiber/matrix interfacial shear strength. Unfortunately, the theory does not account for the limitation where fiber/matrix interfacial debonding and/or sliding does not occur. The lack of relative displacement at the fiber/matrix interface prevents fiber bridging of the matrix crack, resulting in linear elastic behavior.

Another equation for predicting the first matrix cracking stress was developed by Marshall and Cox (1988) using the crack closure pressure suggested by McCartney (1987). This approach was applied by Chulya, et al. (1991) to SiC/RBSN to calculate the first matrix cracking stress. Marshall and Cox (1988) used linear elastic fracture mechanics and assumed a single crack, loaded in mode I, propagating through a semi-infinite medium. The traction from the bridging fibers was superimposed on the crack model as closing pressure. The resulting equation is similar to the ACK results. Marshall and Cox also assumed a weak frictional fiber/matrix interface with bridging fibers. The resultant equation for predicting the composite stress where steady state matrix cracking begins within the lamina is

$$\sigma_{y} = \left[\frac{6\tau\gamma_{m}E_{c}^{2}E_{f}V_{f}^{2}}{E_{m}^{2}V_{m}R}\right]^{\frac{1}{3}}$$
(9)

Combining equations 3 and 9 leads to an equation which predicts the composite cracking stress with residual stress effects included. The result is

$$\sigma_{y\alpha} = \sigma_{y} - \sigma_{m\alpha} \frac{E_{c}}{E_{m}}$$
(10)

Danchaivijit and Shetty (1993) and Budiansky, et al. (1986) have pointed out that the ACK result, as well as equation 9 above, lead to lower bound predictions.

# 5.3. Fiber Properties

As mentioned earlier, ceramic matrix composites contain fibers that have a higher failure strain than the matrix. As a result, the composite's mechanical properties are fiber dominated near the material's ultimate strength. This section discusses the behavior of ceramic fibers in preparation for analyzing the ultimate strength of unidirectional composites.

Most brittle materials, including ceramic fibers, show a large variation in their tensile strength. Therefore, it is appropriate to apply statistical techniques for predicting their reliability and probability of failure. Statistical analysis allows the designer to use data generated from a small sample to systematically predict the stochastic response of complex structures.

Bergman (1984) reported that the cumulative distribution function proposed by Weibull (1939) is the most useful for characterization of brittle materials. Weibull analysis is based on the weakest link theory, where failure is assumed to occur at the largest flaw within the material. The sapphire fibers in this study consist of single crystals eliminating creep due to sliding grains but according to Bunsel and Berger (1997) the fibers are not without defects. Bunsel and Berger (1997) observed bubbles in the fibers most probably due to meniscus at the point of fiber growth and convection in the melt during fiber production. The Weibull theory is purely statistical in nature. It should be noted that a single flaw population and a strength that is independent of time will be assumed. The three parameter Weibull cumulative distribution function for the failure probability of ceramic fibers is

$$P_{f} = 1 - e^{-\int_{V} \left(\frac{\boldsymbol{\sigma} - \boldsymbol{\sigma}_{t}}{\boldsymbol{\sigma}_{0}}\right)^{m} dV}$$
(11)

where  $P_f$  is the probability of failure by fracture at a given stress,  $\sigma$ , m is the shape factor, known as the Weibull modulus, V is the stressed volume,  $\sigma_0$  is the scale factor, and  $\sigma_t$  is the location parameter defined as the threshold stress below which the probability of failure is zero. The scale factor,  $\sigma_0$ , has the dimension of stress times (volume)<sup>1/m</sup> for  $P_f$  to be dimensionless. The above equation can be simplified by taking the conservative assumption of setting the threshold stress to zero, assuming uniaxial fiber stress acting in a material volume with only internal imperfections. Consequently, equation 11 can be expressed as

$$P_{f} = 1 - e^{-V \left(\frac{\sigma}{\sigma_{0}}\right)^{m}}$$
(12)

The above equation is linearized by taking the natural logarithm twice, resulting in the following equation:

$$y = \ln\left[\ln\left(\frac{1}{1 - P_{f}}\right)\right] = \ln(V) + m \ln\left(\frac{\sigma}{\sigma_{0}}\right)$$
(13)

This equation can be plotted on a Weibull plot of y versus  $\ln \sigma$  wherein the Weibull modulus, m, is the slope of the plotted line. The Weibull characteristic strength may be substituted for the volume and Weibull scale parameter to simplify the above equations. The Weibull characteristic strength,  $\sigma_{\theta}$ , is a function of the Weibull scale parameter,  $\sigma_{\sigma}$ ; the stressed volume, V; and the Weibull modulus, m; as shown in the following equation:

$$\sigma_{\theta} = \frac{\sigma_{o}}{V^{\frac{1}{m}}}$$
(14)

Equation 14 is valid for the condition where the stress is evenly distributed over the volume, V, as is the case in uniaxial tension. Substituting equation 14 into equations 12 and 13 we get

$$P_{f} = 1 - e^{\left(\frac{\sigma}{\sigma_{\theta}}\right)^{ff}}$$
(15)

$$y = \ln \left[ \ln \left( \frac{1}{1 - P_{f}} \right) \right] = m \ln \left( \frac{\sigma}{\sigma_{\theta}} \right)$$
(16)

Pai and Gyekenyesi (1988) and Mahfuz, et al. (1997) report that the least squares analysis and maximum likelihood method are the most popular techniques for estimating Weibull parameters from experimental data. The least squares method offers simplicity with the estimated Weibull modulus being the slope of the best fit line by linear regression on a Weibull plot of y versus  $\ln(\sigma)$ . Bergman (1984) also states that the use of the least squares method implies that the ln  $\sigma$  values follow a Gaussian distribution around the line obtained from equation 13. The maximum likelihood technique is a nonlinear model offering efficiency and is better suited to model uniaxial strength data of brittle ceramics. ASTM (1995) has a standard (Designation: C 1239-94a) utilizing the maximum likelihood technique for estimating the Weibull parameters. The maximum likelihood method will be used for this work. The likelihood function, from the ASTM (1995) standard, for the two-parameter Weibull distribution, with a single-flaw population or uncensored data set is

$$= \prod_{n=1}^{n_{tot}} \left(\frac{m}{\sigma_{\theta}}\right) \left(\frac{\sigma_{n}}{\sigma_{\theta}}\right)^{m-1} e^{\left(\frac{\sigma_{i}}{\sigma_{\theta}}\right)^{m}}$$
(17)

where n is the rank of a specimen and  $n_{tot}$  is the total number of specimens in a sample.

The estimates of the Weibull modulus and the characteristic strength are determined by taking the partial derivatives of the natural logarithm of the likelihood function, equation 17, with respect to m and  $\sigma_{\theta}$  and equating the resulting expressions to zero. Following are the resulting equations, for an uncensored sample, which are to be solved numerically.

$$\frac{\sum_{n=1}^{n_{\text{tot}}} \sigma_n^m \ln(\sigma_n)}{\sum_{n=1}^{n_{\text{tot}}} \sigma_n^m} - \frac{1}{n_{\text{tot}}} \sum_{n=1}^{n_{\text{tot}}} \ln(\sigma_n) - \frac{1}{m} = 0$$

$$\sigma_{\theta} = \left[ \left( \sum_{n=1}^{n_{\text{tot}}} \sigma_n^m \right) \frac{1}{n_{\text{tot}}} \right]^{\frac{1}{m}}$$
(18)

A probability estimator is used to give the probability of failure for each failure stress. These probability values are used to calculate y for a given corresponding stress. The measured fracture strengths of the fibers are ranked in ascending order. The following estimator is used to calculate the failure probability:

$$P_n(\sigma_n) = \frac{n - 0.5}{n_{tot}}$$
(20)

where n is the rank of the specimen data point and  $n_{tot}$  is the sample size. Other common probability estimators were discussed by Bergman (1984). The above

probability estimator, which gives an average value of the empirical density function about the corresponding stress, is most useful for small samples of less than 50 as discussed by Bergman (1984). Also, the probability estimator of equation 20 is part of the ASTM (1995) standard, mentioned above.

# 5.4. Ultimate Strength

The ultimate strength of the sapphire fiber reinforced polycrystalline alumina composites is primarily dependent on fiber properties. The strength properties of single crystal alumina fibers, as a function of temperature, were measured in separate tests.

A simple approximation for predicting the composite ultimate tensile strength utilizes the rule of mixtures and the mean fiber strength. For the unidirectional lamina, assuming that the in-situ fiber and the independent fiber strengths are identical and that the matrix carries no load, we have

$$\sigma_{\rm cu} = V_{\rm f} \sigma_{\rm fu} \tag{21}$$

where:

 $\sigma_{cu}$  - the composite ultimate strength  $\sigma_{fu}$  - the mean fiber ultimate strength

It is further assumed here that all the fibers are intact until just prior to the composite ultimate loading.

The assumption of uniformly strong fibers would be the ideal situation, but in practice the fibers tend to fail sequentially starting with the weakest fiber, until the applied load cannot be supported leading to total fracture. Since the fibers are brittle and exhibit stochastic behavior, it is more appropriate to apply statistics to determine their ultimate strength. Duffy, et al. (1991) have pointed out that the strength distribution of the fibers needs to be incorporated into an analytical model for predicting the ultimate strength of the composite.

Equation 22 from Curtin, et al. (1993) and Curtin (1993) has been used to determine the composite's ultimate strength in terms of fiber properties.

$$\sigma_{cu} = V_{f} \left(\frac{2}{m+2}\right)^{\frac{1}{m+1}} \left(\frac{m+1}{m+2}\right) \left[\frac{\sigma_{fu}^{m} \tau L_{f}}{R \ln(2)}\right]^{\frac{1}{m+1}}$$
(22)

where:

 $L_f$  - fiber gage length at which the strength was determined m - Weibull modulus of fiber  $\sigma_{fu}$  - mean fiber strength

It is of interest to compare the above result to that found for a "dry" bundle containing only fibers and no matrix. The composite strength based on bundle ultimate strength,  $\sigma_{cub}$ , can be calculated, per Evans, et al. (1995), from

$$\sigma_{\rm cub} = V_{\rm f} \sigma_{\rm fu} \left( \frac{L_{\rm f}}{L_{\rm c}} \right)^{\frac{1}{\rm m}} e^{-\frac{1}{\rm m}}$$
(23)

where  $L_c$  is the fiber bundle gage length and  $L_f$  is the fiber length used to determine  $\sigma_{fu}$ and m. Equation 23 uses the rule of mixtures to estimate the composite strength,  $\sigma_{cub}$ , with the product of the fiber volume fraction and the fiber bundle strength,  $\sigma_{fub}$ .

Evans (1989), as well as Evans and Marshall (1989), presented a model for predicting the composite ultimate strength based on weakest link statistics, incorporating the fiber Weibull modulus, m. The model is a modified bundle failure analysis which assumes failed fibers have no load bearing ability. The model of the modified bundle failure theory is presented in the following equation:

$$\sigma_{\text{cubm}} = V_{f}\sigma_{\text{fub}}e^{-\frac{1\left(1-\frac{\pi}{R\sigma_{\text{fub}}}\right)^{m+1}}{(m+1)\left[1-\left(1-\frac{\pi}{R\sigma_{\text{fub}}}\right)^{m}\right]}}$$
(24)

The fiber bundle strength,  $\sigma_{\rm fub},$  is determined by iteratively solving the following equation:

$$\left(\frac{R\sigma_{fub}}{\tau x}\right)^{m+1} = \frac{A_{o}}{2\pi RL_{c}} \left(\frac{R\sigma_{o}}{\tau x}\right)^{m} \left[1 - \left(1 - \frac{\tau x}{R\sigma_{fub}}\right)^{m}\right]^{-1}$$
(25)

where  $\tau$  is the interfacial shear stress as defined by equation 2, x is the saturated matrix crack spacing, R is the fiber radius, L<sub>c</sub> is the composite gage length, and A<sub>o</sub> is an area normalizing factor. The scale parameter,  $\sigma_o$ , is defined, according to Chulya, et al. (1991), by the following equation:

$$\sigma_{o} = \frac{\sigma_{fu}}{\Gamma\left(1+\frac{1}{m}\right)} \left(2\pi RL_{f}\right)^{\frac{1}{m}}$$
(26)

Here  $L_f$  is the fiber gage length and  $\Gamma$  is Euler's gamma function, defined as:

$$\Gamma\left(1+\frac{1}{m}\right) = \int_{0}^{\infty} t^{\left(1+\frac{1}{m}\right)^{-1}} e^{-t} dt$$
(27)

Cao and Thouless (1990) made an attempt to predict the ultimate strength of a ceramic composite with the application of two parameter Weibull statistics. Their theory assumed that the matrix is saturated with cracks. As a result, the initial linear elastic behavior and the nonlinear deformation associated with matrix cracking are not incorporated. Another simplifying assumption is that upon fracture of a fiber anywhere within the gage length of the composite, the fiber is unable to carry any

load. Given the assumptions, the following equation is used to predict the ultimate strength of a ceramic composite:

$$\sigma_{cu} = V_f \Sigma \left( \frac{\Sigma R}{m(m+1)\tau L_c} \right)^{\frac{1}{m}} e^{-\frac{1}{m}}$$

(28)

where:

$$\Sigma = \left[\frac{A_{o}\sigma_{o}^{m}\tau(m+1)}{2 \pi R^{2}}\right]^{\frac{1}{m+1}}$$

(29)

As with the other theories in this section, this theory is based on fiber statistics, primarily the Weibull modulus and the scale parameter, and the variables are as defined in the statistical failure theories above.

#### 6. Results and Discussion

In this section the high temperature tensile test results are presented and compared with theory to predict the mechanical behavior of the unidirectional fiber reinforced  $Al_2O_3/Al_2O_3$  composite.

Table 1.	High Temperature Mechanical Tensile Properties
for Al <sub>2</sub> O <sub>3</sub>	$/Al_2O_3$ Composites.

Specimen No.	T (°C)	E <sub>c</sub> (GPa)	$\sigma_{y}  (\text{MPa})$	ε <sub>y</sub> (%)	$\sigma_{cu}  (\text{MPa})$	ε <sub>αι</sub> (%)
64a	22	146	274	0.190	425	0.400
64b	22	151	285	0.190	438	0.350
54a	22	167	308	0.190	362	0.320
54b	22	167	283	0.170	459	0.320
mean	22	158	288	0.185	421	0.348
std. dev.		11	14	0.010	42	0.038
SA167-1	800	168	142	0.080	149	0.082
SA167-2	800	138	206	0.135	206	0.135
SA167-3	800	152	147	0.092	147	0.092
mean	800	153	165	0.102	167	0.103
std. dev.		15	36	0.029	34	0.028
SA171-1	1000	108	47	0.452	89	0.047
SA171-2	1000	177	63	0.037	80	0.030
mean	1000	143	55	0.245	84	0.039
std. dev.		49	12	0.293	7	0.013
SA169-1	1200	128	84	0.054	84	0.054
SA169-2	1200	95	71	0.064	71	0.064
SA169-3	1200	110	103	0.080	103	0.080
mean	1200	111	86	0.066	86	0.066
std. dev.		17	16	0.013	16	0.013
SA180-1	1400	99	70	0.071	112	0.101
SA180-2	1400	106	74	0.071	120	0.097
SA180-3	1400	84	74	0.088	93	0.099
SA180-4	1400	106	52	0.047	95	0.027
mean	1400	99	67	0.069	105	0.081
std. dev.		10	11	0.017	13	0.036

T - test temperature

E<sub>c</sub> – Young's modulus of composite

 $\sigma_v$  – first matrix cracking stress

 $\varepsilon_y$  – strain at first matrix cracking stress

 $\sigma_{cu}$  – ultimate strength of composite

 $\varepsilon_{cu}$  – strain at ultimate strength of composite

A typical stress-strain curve for the composite specimens is presented in Figure 2. All the specimens exhibited pseudo-toughness, or an ability to sustain progressive damage, from room temperature to  $1400^{\circ}$ C (2550°C) in air relative to monolithic polycrystalline alumina. These temperatures were selected to represent potential service conditions. All the specimens tested shared linear elastic behavior, followed by matrix cracking producing a limited nonlinear stress versus strain region. The measured mechanical properties for all the short term static tensile tests for the composite are presented in Table 1.

#### 6.1. Modulus

Figure 3 and Table 1 show the Young's modulus of the composite as a function of temperature. The composite modulus decreases with increasing temperature. In

addition, Figure 3 illustrates the constituent moduli. The fiber moduli at  $22^{\circ}$ ,  $1000^{\circ}$ , and  $1400^{\circ}$ C ( $72^{\circ}$ ,  $1830^{\circ}$ , and  $2550^{\circ}$ F) were provided by Dr. Ali Sayir (1998) of Case Western Reserve University. The high temperature fiber modulus data was generated using the system described by Sayir, et al. (1994). It can be observed in Figure 3 that the single crystal alumina fiber modulus varies fairly linearly with respect to temperature from room temperature to  $1400^{\circ}$ C ( $2550^{\circ}$ F). This behavior is consistent with the observation made by Wachtman Jr. and Lam (1959) for single crystal alumina. Fiber moduli at  $800^{\circ}$  and  $1200^{\circ}$ C ( $1470^{\circ}$  and  $2190^{\circ}$ F) were derived using linear interpolation.

Monolithic polycrystalline alumina specimens were tested at room temperature and 1200°C (2190°F). The room temperature results were reported by Jaskowiak and Setlock (1994). The material is the same as the matrix in the  $Al_2O_3/Al_2O_3$  composite used for this work. The measured and interpolated moduli are presented in Figure 3. Wachtman Jr. and Lam (1959) showed that the modulus decreased linearly with respect to temperature from room temperature to approximately 950°C (1740°F). Above 950°C (1740°F) the modulus of the polycrystalline alumina dropped much more rapidly due to grain boundary softening and sliding according to Wachtman Jr. (1996). The percentage change in the modulus with respect to temperature presented by Wachtman Jr. and Lam (1959) was used to extrapolate moduli for the alumina used for the composite matrix at different temperatures.

The moduli of the constituents were used to determine the Young's modulus of the composite using the rule-of-mixtures, equation 1. The measured composite moduli and the calculated moduli are presented in Figure 3. There is excellent correlation between the measured and predicted values for the composite Young's modulus from room temperature to 1000°C (1830°F). At 1200°C (2190°F) the predicted value is greater than the mean measured modulus but within one standard deviation. Above 1400°C (2550°F) the rule of mixtures overestimates the modulus of the composite. The lower value of the measured Young's modulus at 1400°C (2550°F) may be due to changes in the zirconia interface between the fibers and the matrix. It is noted by Ryshkewitch and Richerson (1985) that the crystal structure of the zirconia starts to change from monoclinic to tetragonal at approximately 1000°C (1830°F) resulting in a volume reduction of about 9%. The volume reduction of the zirconia may be leading to debonding of the fiber/matrix interface. The fiber/matrix interfacial debonding would decrease or prevent load transfer between the fibers and the matrix which would reduce the fibers' contribution to the overall stiffness of the composite. A more in depth investigation of the interfacial properties at high temperatures may be a subject for future work.

# 6.2. Matrix Cracking and Interfacial Shear Properties

The effect of temperature on the proportional limit also was studied. It was assumed that the proportional limit, on a stress/strain curve, and the first matrix cracking stress were the same. The first matrix cracking stress is the onset of permanent damage in a composite, making it a critical design parameter.

The first matrix cracking stress is affected by various parameters. One of these parameters is the fiber/matrix interfacial shear strength. Other parameters are addressed later in this section. It has been noted by Kerans, et. al. (1989) that the interfacial shear strength is difficult to characterize with certainty. In this study the fiber/matrix interfacial shear strength is taken from the work of Jaskowiak, et al.

(1997) for room temperature. The high temperature interfacial shear stress is determined as a function of matrix crack spacing in a composite that has been loaded to the ultimate strength of the composite. Figure 4 illustrates the average matrix crack spacing as a function of temperature in air. Scatter was observed in the matrix crack spacing as indicated by the first standard deviation error bars in Figure 4. The matrix crack spacing was measured using a scanning electron microscope (SEM). It should be noted that it was very difficult to locate through-the-composite, fiber bridged, regularly spaced cracks within the Al<sub>2</sub>O<sub>3</sub>/Al<sub>2</sub>O<sub>3</sub> composite. Next, the average crack spacing for each specimen was used to determine the mean fiber/matrix interfacial shear strength,  $\tau$ , using equation 2. It is assumed that the matrix is characterized by a single

**Table 2.** Coefficients of Thermal Expansion for Sapphire Fibers and Polycrystalline Al<sub>2</sub>O<sub>3</sub> Matrix.

T (°C)	<sub>αf</sub> (10 <sup>-6</sup> /°C)	αm (10 <sup>-6</sup> /°C)
22	6.26	6.03
800	8.43	7.90
100	8.72	8.14
1200	8.89	8.30
1400	9.02	8.43

 $\alpha_{\rm f}$  - fiber coefficient of thermal expansion  $\alpha_{\rm m}$  - matrix coefficient of thermal expansion

assumed that the matrix is characterized by a single value for strength. Yang and Knowles (1992) have made an attempt to apply Weibull statistics to characterize the matrix crack spacing with limited success, but more in depth analysis is required.

Referring to Figure 4 it can be observed that matrix crack spacing decreases with increasing temperatures up to approximately 1200°C (2190°F). At 1400°C (2550°F) there is a significant increase in the matrix crack spacing. The large increase in matrix crack spacing may be attributed to a volume reduction of the zirconia at the fiber/matrix interface as discussed in the previous section, 6.1, on the modulus. The resultant debonding between the fibers and the matrix prevents the fibers from enhancing the matrix fracture stress relative to monolithic alumina. Figure 5 shows the interfacial shear strength as a function of temperature in air. It is difficult to conclude how the interfacial shear stress varies with temperature but there is a general reduction going from room temperature to 1400°C (2550°F). In addition, it may be noted that the interfacial shear strength is rather weak ranging from 55 to 3 MPa (8.0 to 0.4 ksi). The interfacial shear stress can vary significantly within a composite due to variations in processing conditions and effects of neighboring fibers as proposed by Chulya, Gyekenyesi, and Bhatt (1991). Research is continuing to identify a test for accurately determining the interfacial shear strength between the fiber and the matrix of various composite materials. This includes tensile tests of single fibers coated with the matrix material by Morscher, Martinez-Fernandez, and Purdy (1994) and fiber push-out tests by Eldridge, Bhatt, and Kiser (1991). Eldridge and Ebihara (1994) and Eldridge (1995) developed a high temperature fiber push-out system capable of testing specimens up to 1100°C (2010°F) in a vacuum environment.

The coefficient of thermal expansion (CTE) for the matrix and fibers are slightly different as can be observed in Table 2. The CTE values were taken from the work of Gitzen (1970) for 0° single crystal alumina and polycrystalline alumina. This results in residual stresses within the composite at temperatures other than the processing temperature. The processing temperature, assumed to be approximately 1300°C (2400°F), is used as the reference temperature at which there are no residual stresses due to the CTE mismatch within the composite. Equation 3 is used to determine the residual stress within the matrix. The Poisson's ratio is taken from the work of Gitzen (1970) and is assumed to be the same for the fiber and the matrix at a mean value of  $v=v_f=v_m=0.22$ .

The stress values at the proportional limit are illustrated in Figure 6 and Table 1. In addition, Figure 6 presents predicted values from the ACK theory and the Marshall and Cox theory combined with the McCartney theory. The matrix fracture surface energy,  $\gamma_m$ , is assumed to be a conservative 5 J/m<sup>2</sup> (0.03 in·lb<sub>f</sub>/in<sup>2</sup>). The value for the matrix crack surface energy is taken from Wachtman Jr. (1996) who showed ceramics to typically vary from 5 to 50 J/m<sup>2</sup> (0.03 to 0.29 in·lb<sub>f</sub>/in<sup>2</sup>).

Looking at Figure 6 and Table 1 we find that the estimated values for the first matrix cracking stress by the ACK theory (equation 8), and the Marshall and Cox theory combined with the McCartney theory (equation 9) were conservative at room temperature while not accounting for residual stresses and assuming a frictional fiber/matrix interface. By accounting for residual stresses the ACK theory estimates the room temperature first matrix cracking stress fairly accurately. On the other hand, by accounting for the residual stresses the Marshall and Cox theory combined with the McCartney theory still results in a conservative prediction for the first matrix cracking stress. Above room temperature all the theories over estimate the first matrix cracking stress. This is probably due to the difficulty in determining the fiber/matrix interfacial shear stress. In addition, the matrix fracture surface energy may be reduced slightly from the assumed value as temperature is increased. The ACK theory, with the assumption of a frictional fiber/matrix interface and not accounting for residual stresses, is recognized as a lower bound as pointed out by Danchaivijit and Shetty (1993) and Budiansky, Hutchinson, and Evans (1986), although, the Marshall-Cox theory combined with McCartney's theory, and not accounting for residual stresses, results in a more conservative prediction.

It should be noted that this study is primarily concerned with the high temperature behavior of these composites, that is, approximately at potential operating temperatures up to  $1400^{\circ}$ C (2550°F). Wachtman Jr. (1996) has shown that many researchers have observed a drop in strength between approximately 400° to 900°C (750° to 1650°F). Future studies of this composite system may be made in the lower temperature regime.

# 6.3. Sapphire Fiber Properties

The single crystal alumina fibers were tested in an "as-received" condition. That is, they were taken directly off the spool on which they were delivered and then tested.

The properties including the ultimate strength of individual fibers, mean ultimate strength with a variance of one standard deviation, Weibull modulus, and Weibull characteristic strength as a function of temperature for the sapphire fibers are presented in Table 3. In addition, the mean fiber ultimate strengths with one standard deviation are plotted as a function of temperature in Figure 7. Figure 8 presents a typical two-parameter Weibull plots with  $Ln(Ln(1/(1-P_f)))$  versus the fiber ultimate strength for the 'as-received' sapphire fibers. The same figure illustrates the 90% confidence bounds as determined using the ASTM (1995) standard designation C 1239-94a. Figure 9 presents the Weibull modulus as a function of temperature with 90% confidence bounds for the sapphire fibers.

The Weibull modulus is determined by numerically solving equations 18 and 19, which are part of the maximum likelihood method.

The data from Table 3 and Figure 7 indicate a steady decrease in ultimate strength with increasing temperature from room temperature to 1500°C (2730°F). Weibull moduli range from a low of 10 to a high of 24 from room temperature to

n	22°C <sub>ofu</sub> (GPa)	900°C <sub>ofu</sub> (GPa)	1000°C <sub>σfu</sub> (GPa)	1100°C <sub>σfu</sub> (GPa)	1200°C <sub>σfu</sub> (GPa)	1300°C <sub>σfu</sub> (GPa)	1400°C <sub>σfu</sub> (GPa)	1500°C <sub>σfu</sub> (GPa)
1	3.18	1.02	0.97	0.86	0.88	0.75	0.67	0.64
2	3.16	1.02	0.94	0.82	0.82	0.73	0.64	0.63
3	3.05	1.02	0.93	0.79	0.81	0.71	0.62	0.62
4	2.98	0.97	0.89	0.77	0.80	0.70	0.61	0.62
5	2.84	0.94	0.88	0.77	0.75	0.69	0.60	0.62
6	2.80	0.93	0.82	0.76	0.74	0.69	0.59	0.62
7	2.68		0.82	0.76	0.73	0.68	0.58	0.57
8	2.50		0.70	0.73	0.73	0.67	0.55	0.56
9	2.45		0.69	0.73	0.72	0.67	0.53	0.56
10	2.27			0.70	0.70	0.65	0.48	0.55
mean <sub>ofu</sub> (GPa) =	2.79	0.99	0.85	0.77	0.77	0.69	0.59	0.60
standard deviation	0.31	0.04	0.10	0.05	0.06	0.03	0.06	0.04
*m =	9.7	23.9	9.5	14.3	12.0	20.4	11.4	18.8
5% bound for m	6.2	13.1	6.0	9.2	7.7	13.1	7.4	12.1
95% bound for m	15.4	45.8	15.6	22.6	18.9	32.4	18.1	29.8
*σ <sub>θ</sub> (GPa) =	2.91	1.00	0.89	0.79	0.79	0.71	0.61	0.61

**Table 3.** The Ultimate Strength, Mean Ultimate Strength, Weibull Modulus, and Weibull Characteristic Strength as a Function of Temperature of Sapphire Fibers.

environment: air

gage length: 25 mm

\*Unbiased values derived by using maximum likelihood method

with probability estimator  $P=(n-0.5)/n_{tot}$  as per ASTM standard

designation C 1239-94a

n - rank of the specimen data point

 $L_{\rm f}$  - specimen gage length

m - Weibull modulus

 $\sigma_{\theta}$  - characteristic strength

approximately 1500°C (2730°F) for the sapphire fibers. The scatter of the Weibull moduli as a function of temperature with the wide confidence bounds may be an indication that a larger sample of fibers is needed.

As noted in the section on matrix cracking, it should be pointed out that this study is primarily concerned with the high temperature behavior of these composites, that is, at potential operating temperatures up to approximately  $1400^{\circ}C$  (2550°F). Wachtman Jr. (1996) has shown that many researchers have observed a drop in strength between approximately  $300^{\circ}$  to  $600^{\circ}C$  (570° to  $1110^{\circ}F$ ) for single crystal alumina. Future studies of this composite system may be made in the lower temperature regime.

# 7.1.4 Ultimate Tensile Strength of Al<sub>2</sub>O<sub>3</sub>/Al<sub>2</sub>O<sub>3</sub>

The ultimate strengths of these composites are primarily dependent on fiber properties. As a result, the mechanical properties of the composite, near the ultimate strength, will be effected by the environment in a similar manner as the fiber constituent. The strength properties of the single crystal alumina fibers as a function of temperature were discussed in the previous section.

The ultimate tensile strengths of the composites are plotted as a function of temperature in Figure 10. The plot includes the mean ultimate strength with one standard deviation from tensile tests in air and theoretical curves using constituent properties. The same properties are presented in Table 4. To reiterate, the test temperatures were room temperature, 800°, 1000°, 1200°, and 1400°C. Figure 10 and Table 4 indicate a steady drop in the ultimate strength of the composite from room temperature to approximately 1000°C (1830°F). There is an almost insignificant increase in the ultimate strength from 1000°C (1830°F) to 1400°C (2550°F).

It can be observed in Figure 10 and Table 4 that the rule of mixtures, equation 21, over-predicts the ultimate strength considerably. It is only at 1400°C (2550°F) where the rule of mixtures prediction approaches the measured ultimate strength, but it is still higher than the experimentally observed values and the other theoretical predictions utilized in this section. It was mentioned in the previous chapter that this approach assumes that all the fibers are intact just prior to the composite ultimate strength. The rule of mixtures does not account for the brittle nature of the reinforcing fibers and the corresponding variance in the strength of the fibers. As a result, the rule of mixtures assumes ideal conditions, producing optimistic strength values relative to the experimentally derived measurements of the composites. Lastly, fiber damage, due to processing of the composite, may reduce the ultimate strength of the composite also.

Table 4. The Experime	ntal and Theor	etical Ultimate	e Tensile Strength	of $Al_2O_3/Al_2O_3$	as a Function of
Temperature in Air.					

Т	experimental	ROM	Curtin	Evans 1995	Evans 1989	Cao & Thouless
(°C)	(MPa)	(MPa)	(MPa)	(MPa)	(MPa)	(MPa)
22	421±42	728	704	656	584	523
800	167±34	324	311	303	270	262
1000	84±7	227	209	205	173	170
1200	86±16	207	205	190	167	162
1400	105±13	182	159	167	145	142

experimental - mean and one standard deviation values from tensile tests

ROM - rule of mixtures

Curtin - Curtin's theory incorporating rule of mixtures and Weibull statistics with L<sub>f</sub>=25 mm

Evans 1995 - Evans' theory using fiber bundle theory with no matrix material with  $L_f=L_c=25$  mm except  $L_c=38$  mm at room temperature

Evans 1989 - Evans' theory based on a modified fiber bundle theory

Cao & Thouless - Theory incorporating rule of mixtures and Weibull statistics

\* Theoretical results using fiber properties in air

Since the fibers are brittle and exhibit linear elastic behavior with a significant variance in strength, relative to metals or other common structural materials, it is more appropriate to apply statistics to determine the ultimate strength. Curtin's (1993) theory, equation 22, incorporates Weibull statistics to address the variance in the strength of the fibers. Data from Figure 5 and Table 3 were used with Curtin's theory. The fiber gage length,  $L_f$ , is 25 mm (0.98). The results are presented in Table 4 and Figure 10. Curtin's theory also produces optimistic values for the ultimate strength. Unfortunately, the theory predicts strengths significantly over the measured

ultimate strengths. Curtin's theory is only a slight improvement over the simple rule of mixtures for predicting the ultimate strength of the brittle fiber reinforced  $Al_2O_3$  composite system. As noted with the rule-of-mixtures, part of the discrepancy between Curtin's theory and the measured ultimate strength may be due to fiber damage during processing of the composites.

Evans, et al. (1995) tried to predict the ultimate strength of the composite using the rule of mixtures, the fiber bundle strength with Weibull statistics, and with the assumption that influence from the matrix was insignificant. The theory is presented in Equation 23. The results are shown in Figure 10 and Table 4. The mean fiber strengths and the fiber Weibull moduli were taken from Table 3. All gage lengths,  $L_f$  and  $L_c$ , are 25 mm except at room temperature where the composite gage length,  $L_c$ , assumed to be the same as the bundle gage length, is 38 mm. The theory overestimates the ultimate strength of the  $Al_2O_3/Al_2O_3$  composite at all temperatures but it shows a similar curve to Curtin's theory. The curves show a reduction in ultimate strength from room temperature to approximately 1000°C (1830°F) with a minimal change 1000°C (1830°F) to 1400°C (2550°F). Again, fiber damage due to processing may add to the discrepancy between the theory of Evans, et al. (1995) and the measured composite ultimate strengths.

Evans (1989) presented a model based on a modified fiber bundle theory. The theory is presented in equation 24. The results are presented in Figure 10 and Table 4. The required parameters of the fiber bundle strength,  $\sigma_{fub}$ , was solved iteratively using equation 25 and the scale parameter,  $\sigma_0$ , was determined using equation 26. The area normalizing factor,  $A_0$ , was equal to 1.0 m<sup>2</sup>. All gage lengths,  $L_f$  and  $L_c$ , are set at 25 mm (0.98 in.) except at room temperature where  $L_c=38$  mm (1.5 in.). The theory produces results which are better than Curtin's theory (1993); Evans', et al. (1995) theory; and the rule of mixtures. At room temperature the prediction is 39% greater than the experimentally derived data. Unfortunately, the Evans' modified fiber bundle theory over predicts by up to 106% over the measured ultimate strength at 800°C (1470°F). Fiber degradation, due to the processing conditions of the composite, may have added to the discrepancy between the theory and the measured composite ultimate strength.

Cao and Thouless (1990) also presented a theory for predicting the ultimate strength for a ceramic composite. It is based on two parameter Weibull statistics as described in the previous section. The theory was applied to the  $Al_2O_3/Al_2O_3$  system with the results presented in Figure 10 and Table 4. The same scale parameters,  $\sigma_0$ , calculated for the modified fiber bundle theory of Evans (1989), were used for this theory. The composite gage length, L<sub>c</sub>, is set at 38 mm (1.5 in.) for room temperature and 25 mm (0.98 in.) for temperatures greater than room temperature. At room temperature, the theory over predicted the composite ultimate strength by 24%. The theoretical predictions over estimate the ultimate strength up to 102% at 1000°C (1830°F). Of the theories evaluated in this paper, the theory from Cao and Thouless (1990) does offer the best predictions for the composite ultimate strength but the results are still overly optimistic relative to the experimentally derived ultimate strengths. The theory provides similar predictions to the Evans' (1989) theory for temperatures between 800°C (1470°F) to 1400°C (2550°F). Composite fabrication induced fiber damage may have reduced the measured composite ultimate strength also.

All the theories that were addressed in this section using the fiber properties overestimated the ultimate strength of the  $Al_2O_3/Al_2O_3$  composite. It was noted with each of the theories that the lower values of the experimentally obtained ultimate

strength measurements for the  $Al_2O_3/Al_2O_3$  composite may be due to fiber damage incurred during processing of the Al<sub>2</sub>O<sub>3</sub>/Al<sub>2</sub>O<sub>3</sub> or due to possible surface damage produced upon matrix fracture during the tensile test. Although as the test temperatures approach 1000°C (1830°F) and over, it is believed that the fiber strength degrades due to dislocation microplasticity, twinning, stresses due to thermal expansion anisotropy, and slow crack growth as noted by Wachtman Jr. (1996) for single crystal alumina. This degradation of intrinsic strength becomes more significant than the strength reduction due to surface damage. As a result, the predicted composite ultimate tensile strength, based on fiber properties, becomes more accurate when compared with the measured composite ultimate tensile strength. The application of the mean fiber strength and the fiber Weibull modulus does predict the ultimate strength of the composite with increased accuracy. Increasing the number of specimens per test should show more accurate results when applying statistics. In addition, a more accurate technique is needed to model the fiber damage due to the processing of the composite. Lastly, the mean strengths and Weibull moduli for the fibers were generated by testing individual fibers. Hill and Okoroafor (1995) have observed a reduction in strength of fiber bundles due to inter-fiber friction. This inter-fiber friction may be causing some of the reduction in strength of the experimentally derived composite data relative to the theoretical values derived from individually tested fibers. On an interesting note, Cox, Marshall, and Thouless (1989) found that the fracture of composites is not greatly influenced by the breadth of the fiber strength distribution. Although the theories incorporating the scatter in fiber strength improved the predictions of the composite ultimate strength, the conclusions of Cox, Marshall, and Thouless (1989) reinforces the idea that the lower strength of composites is most likely due to damaged fibers caused by the processing conditions of the composite.

# 7. Summary and Conclusion

The composite system studied includes uniaxial single crystal alumina fiber reinforced polycrystalline alumina utilizing a porous  $ZrO_2$  fiber/matrix interface. The Al<sub>2</sub>O<sub>3</sub>/Al<sub>2</sub>O<sub>3</sub> composite was tensile tested at high temperatures in air from room temperature to 1400°C (2550°F).

The composite exhibited progressive fracture at all the test temperatures of the short term static tensile tests. The Young's modulus shows a slight decrease as the test temperature was increased from room temperature to  $1000^{\circ}$ C ( $1830^{\circ}$ F) but decreased at a greater rate above  $1000^{\circ}$ C ( $1830^{\circ}$ F). The rule of mixtures gives a good approximation of the Al<sub>2</sub>O<sub>3</sub>/Al<sub>2</sub>O<sub>3</sub> modulus as function of temperature using the constituent moduli up to approximately  $1200^{\circ}$ C ( $2190^{\circ}$ F). More information is required for predicting composite modulus above  $1200^{\circ}$ C ( $2190^{\circ}$ F).

The characterization of the fiber/matrix interfacial shear strength requires further study. The first matrix cracking stress decreases with respect to temperature from room temperature to 1000°C (1830°F) followed by a slight increase up to approximately 1200°C (2190°F). The first matrix cracking stress continues to decrease as temperature is increased from 1200°C (2190°F) to 1400°C (2550°F). The ACK theory (1971), with the assumption of a frictional fiber/matrix interface and accounting for residual stresses, produces a good approximation for the composite first matrix cracking stress at room temperature. Above the room temperature the ACK theory and the Marshall and Cox theory combined with the McCartney theory significantly overestimate the first matrix cracking stress whether the residual stresses are accounted for or not. This is believed to be, primarily, due to the difficulty in characterizing the fiber/matrix interfacial sheer stress above room temperature.

The ultimate strength of the composite showed a drop from room temperature to approximately 1000°C (1830°F) followed by a negligible increase in strength from 1000°C (1830°F) to 1400°C (2550°F). The theory of Cao and Thouless (1990) produced the best estimates of the composite ultimate strength relative to the theories evaluated in this work. Unfortunately, the Cao and Thouless theory over estimates the composite ultimate strength from 24 to 102% relative to the experimental values. There is a need to study the effects of composite processing on the fiber properties.

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Figure 2. Typical tensile stress-strain curve for unidirectional  $Al_2O_3/Al_2O_3$  composites with approximately 27% fiber volume fraction and monolithic  $Al_2O_3$ .



Figure 3. Tensile Young's modulus of  $Al_2O_3/Al_2O_3$  composite, sapphire fiber, and monolithic  $Al_2O_3$  vs. temperature. The mononlithic polycrystalline  $Al_2O_3$  has approximately the same porosity as the matrix in the composite.



Figure 4: Mean matrix crack spacing with one standard deviation versus temperature in air for unidirectional  $Al_2O_3/Al_2O_3$  composites when loaded in the primary direction.



Figure 5. Fiber/matrix interfacial shear strength versus temperature in air for Al<sub>2</sub>O<sub>3</sub>/Al<sub>2</sub>O<sub>3</sub> composites.



Figure 6. Average with one standard deviation, measured and predicted unidirectional  $Al_2O_3/Al_2O_3$  composite tensile first matrix cracking stresses as a function of temperature in air. Loading was along the primary direction.



Figure 7. Mean ultimate tensile strength with one standard deviation of single crystal alumina fibers as a function of temperature in air.



**Figure 8.** A typical plot for  $Ln(Ln(1/(1-P_f)))$  versus the ultimate tensile strength of individual single crystal alumina fibers. Ninety percent confidence bounds determined as per ASTM (1995) standard designation C 1239-94a.



Figure 9. Weibull modulus of single crystal alumina fibers as a function of temperature in air.



Figure 10. Average measured  $Al_2O_3/Al_2O_3$  composite ultimate tensile strength with one standard deviation relative to predictions of the composite ultimate strength using single crystal alumina fiber properties as a function of temperature. Loading was along the primary direction.

# **REPORT DOCUMENTATION PAGE**

Form Approved OMB No. 0704-0188

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1. AGENCY USE ONLY (Leave blank)	2. REPORT DATE	3. REPORT TYPE AND	D DATES COVERED					
4. TITLE AND SUBTITLE	5. FUNDING NUMBERS							
High Temperature Mechanica Composition	al Characterization and Analysi	is of A1 <sub>2</sub> O <sub>3</sub> /A1 <sub>2</sub> O <sub>3</sub>	NUL 522 21 12 00					
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John Z. Gyekenyesi and Mart	tha H. Jaskowiak							
7. PERFORMING ORGANIZATION NAM	IE(S) AND ADDRESS(ES)		8. PERFORMING ORGANIZATION REPORT NUMBER					
National Aeronautics and Spa	ce Administration							
Lewis Research Center			E-11562					
Cleveland, Ohio 44135–319	1							
9. SPONSORING/MONITORING AGEN	CY NAME(S) AND ADDRESS(ES)		10. SPONSORING/MONITORING AGENCY REPORT NUMBER					
National Aeronautics and Spa	ce Administration							
Washington, DC 20546-000	01		NASA TM—1999-208855					
John Z. Gyekenyesi, Universi NCC3–615); and Martha H. J organization code 5130, (216	John Z. Gyekenyesi, University of Akron, Civil Engineering Department, Akron, Ohio 44325 (work funded under NCC3–615); and Martha H. Jaskowiak, NASA Lewis Research Center. Responsible person, Martha H. Jaskowiak, organization code 5130, (216) 433–5515.							
12a. DISTRIBUTION/AVAILABILITY ST	ATEMENT		12b. DISTRIBUTION CODE					
Unclassified - Unlimited								
Subject Categories: 39 and 27								
This publication is available from	This publication is available from the NASA Center for AeroSpace Information. (301) 621–0390.							
13. ABSTRACT (Maximum 200 words)								
Sixteen ply unidirectional zirconia coated single crystal Al <sub>2</sub> O <sub>3</sub> fiber reinforced polycrystalline Al <sub>2</sub> O <sub>3</sub> was tested in uniaxial tension at temperatures to 1400 °C in air. Fiber volume fractions ranged from 26 to 31%. The matrix has primarily open porosity of approximately 40%. Theories for predicting the Young's modulus, first matrix cracking stress, and ultimate strength were applied and evaluated for suitability in predicting the mechanical behavior of Al <sub>2</sub> O <sub>3</sub> /Al <sub>2</sub> O <sub>3</sub> composites. The composite exhibited pseudo tough behavior (increased area under the stress/strain curve relative to monolithic alumina) from 22° to 1400 °C. The rule-of-mixtures provides a good estimate of the Young's modulus of the composite using the constituent properties from room temperature to approximately 1200 °C for short term static tensile tests in air. The ACK theory provides the best approximation of the first matrix cracking stress while accounting for residual stresses at room temperature. Difficulties in determining the fiber/matrix interfacial shear stress at high temperatures prevented the accurate prediction of the first matrix cracking stress above room temperature. The theory of Cao and Thouless, based on Weibull statistics, gave the best prediction for the composite ultimate tensile strength.								
14. SUBJECT TERMS			15. NUMBER OF PAGES					
Mechanical characterization;	36 16. PRICE CODE							
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