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# 1999 NASA High-Speed Research Program Aerodynamic Performance Workshop 

Volume I-Configuration Aerodynamics
Edited by
David E. Hahne
Langley Research Center, Hampton, Virginia

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## PREFACE

The High-Speed Research Program sponsored the NASA High-Speed Research Program Aerodynamic Performance Review on February 8-12, 1999 in Anaheim, Califomia. The review was designed to bring together NASA and industry High-Speed Civil Transport (HSCT) Aerodynamic Performance technology development participants in areas of: Configuration Aerodynamics (transonic and supersonic cruise drag prediction and minimization) and High-Lift. The review objectives were to: (1) report the progress and status of HSCT aerodynamic performance technology development; (2) disseminate this technology within the appropriate technical communities; and (3) promote synergy among the scientist and engineers working HSCT aerodynamics. The HSR AP Technical Review was held simultaneously with the annual review of the following airframe technology areas: Materials and Structures, Environmental Impact, Flight Deck, and Technology Integration. Thus, a fourth objective of the Review was to promote synergy between the Aerodynamic Performance technology area and the other technology areas within the airframe element of the HSR Program.

The work performed in the Configuration Aerodynamics (CA) element of the High-Speed Research Program during 1998 was presented in the following sessions:

Propulsion Integration
Analysis Methods
Design Optimization
Testing
The work performed in the High Lift (HL) element of the High-Speed Research Program during 1998 was presented in the following sessions:

High-Lift Configuration Development Tools and Methods Development
The proceedings for the Aerodynamic Performance Annual Review are published in three volumes:

Volume I, Parts 1 and 2 Configuration Aerodynamics<br>Volume II, Parts 1 and $2 \quad$ High Lift<br>AP Review Chairperson: David Hahne NASA Langley Research Center

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Geometry-Driven Mesh Deformation

oulline
An outline of the material to be presented. Note that the examples section consisted of a
video presentation that is not included in this documentation.


- Motivation
- General Strategy
- Details of Methodology
- Examples
- Future Work
At each design step during the application of an aerodynamic shape optimization method,
it must be possible to obtain a mesh representing the current design changes and the
gradients of the cost function. In an adjoint based method, the determination of changes in
the cost function require that the variation in the mesh metrics be computed. Since it would
be difficult to obtain an explicit relationship between arbitrary surface changes and
variations in mesh metrics for complex configurations, these quantities are calculated
through finite differences. This requires the mesh to be regenerated a number of times
proportional to the number of design variables. The use of hyperbolic or elliptic iterative
mesh generation techniques to construct a mesh for a single design variable perturbation
becomes computationally prohibitive and may approach the cost of a flow solution.
Additionally, no automated method exists for generating multiblock meshes about complex
three-dimensional configurations.

Here, these difficulties are overcome through the use of a mesh perturbation technique. In
this approach, a high quality mesh appropriate for the flow solver is first generated by any
available procedure prior to the start of the design. This initial mesh becomes the basis for
all subsequent meshes which are obtained by analytic perturbations.
Applications for a mesh perturbation scheme can be found in any area which wishes to take advantage of an existing mesh in order to avoid the laborious and time intensive process of regenerating a grid system. Rapid mesh morphing permits the reuse of existing mesh topologies on similar configurations as well as analysis and design accounting for static

- Demands Arise From:
- Aerodynamic Shape Optimization of Complex Configurations
- Reuse of Existing Mesh Topologies on New Configurations
- Analysis and Design Accounting for Static Aeroelastics - Demands Arise From:
- Aerodynamic Shape Optimization of Complex Configurations
- Reuse of Existing Mesh Topologies on New Configurations
- Analysis and Design Accounting for Static Aeroelastics
- Limitations of Manual Approach:
 - Grid Generation
* Calendar-Intensive
* User-Intensive
* Difficult to Automate - Grid Generation
* Calendar-Intensive
* User-Intensive
* Difficult to Automate
$\infty$ - Geometry Treatment
* Lofting
* Intersections/Trimming
* Constraints
Structured Multi-Block Volume Mesh
The surface and symmetry planes of a typical multiblock volume mesh are illustrated
in this slide. In practice, a multiblock volume mesh for a complex configuration requires a
skilled grid generator one to three weeks to complete as well as related CAD support. In the
event of design changes, the mesh perturbation scheme allows one to circumvent this
labor-intensive grid generation process by recycling the previous multiblock mesh through
an automated procedure of surface and volume mesh warping.



# Direct Mesh Perturbation Approach 



> An advantage, however, of using the mesh points as design variables is that, when combined with an analytical mesh mapping transformation, the calculation of the gradient can be performed without explicitly computing the variations in the mesh metrics. Unfortunately, obtaining such a general mapping transformation increases in difficulty with added geometric complexity.
Direct Mesh Perturbation Approach

- Wing Vertical Translation Relative to Body Deform fuselage shape?
Re-spline fuselage shape?

Geometry Kernel Approach

At any particular point in the design process, changes in the underlying geometry are relayed to the mesh surfaces by first intersecting all of the geometric entities to construct a set of parametric surfaces, or panels, representing the modified complete configuration. Using a predetermined parametric mapping of the surface mesh points to the paneling of the initial configuration, changes between the modified and initial panelings at the CFD mesh points are evaluated and the surface mesh perturbed.

[^1]Geometry Kernel

- Use geometry components
- Apply design changes to components
- Impose constraints on components - Intersect components and panel the Deflected CFD Mesh $\stackrel{\text { 䍚 }}{4}$




General Strategy
In the method developed in this research, an underlying geometry database consisting of a set of simple geometric entities, such as wings and bodies, is input to the design algorithm along with the multiblock mesh used for the flow and adjoint calculations. An additional pre-processed input consists of the location of each surface mesh point in the initial mesh determined in relation to the parametric patch representation of the initial geometry and is termed the ( $u, v$ ) map.
Phich are
PERTURB then evaluates the linear and nonlinear constraints on these primary entities.
Next, an automated paneling method, AEROSURF, intersects the geometry components
and constructs parametric patches for the exterior geometry. AEROSURF is a generalized
paneling engine able to handle variable topologies of wings, fuselages, nacelles, pylons,
and diverters in an automated and efficient manner.


[^2]

Starting Point:

- Initial 3D Stru
- Basic Geomet
- Apply Shape Changes and Constraints to Geometry Components

Sections

- Initial 3D Structured Mesh (GRIDGEN)
- Basic Geometry Components as Defining
PERTURB:

0
-

- Initial 3D Structured Mesh (GRIDGEN)
- Basic Geometry Components as Defining Sections
- PERTURB:
- Apply Shape Changes and Constraints to Geometry Components

MAPXYZ:
- Use Parametric Patches to Update CFD Surface Mesh WARP-MB:
- Morph CFD Volume Mesh to Follow Surface Mesh
- 

Mesh Deformation Flow Chart
In this slide, the mesh deformation scheme implemented in the SYN107-MB design code
is represented as a flow diagram illustrating the functionality and sequential relation of the
individual modules as well as the module dependency upon the input data. The four
internal SYN107-MB procedures comprising the geometry engine and mesh perturbation
scheme are shown along the centerline of the diagram and are invoked sequentially from
left to right: PERTURB, AEROSURF, MAPXYZ, and WARP-MB. Additionally, the
three pre-defined inputs and their entry points to the procedures are indicated: an initial set
of geometry components, a multiblock volume mesh, and the $(u, v)$ map relating the surface
mesh points to the parametric paneling of the initial geometry. Given a set of geometry components, the PERTURB module applies the design variables to
affect changes in the geometry. AEROSURF then intersects the modified geometry entities
and produces a parametric paneling of the new configuration. The (u,v) map is then used
by MAPXYZ to propagate changes in the paneled geometry to CFD surface mesh. Finally,
WARP-MB completes the method by warping the volume mesh to follow the modified
surfaces.


A typical set of geometry components for the TCA-6 wing, fuselage, nacelles, and 100 points per section while the fuselage has 200 sections of 100 points per section and the nacelles have 178 sections of 101 points per section.

Design Variables and Constraints
The PERTURB module is responsible for the application of design variables to the
unintersected geometry components. The design variables consist of analytic shape
functions including Hickes-Henne sine bumps, Wagner functions, and various polynomial
and exponential functions. These functions allow the designer to make changes in wing
planform and sectional shapes such as: taper, dihedral, sweep, span, camber, and thickness.
Additionally, some of the specialized functions allow control of flaps, slats, leading edge
droop and trailing edge angle. Design variables applied to the body allow for changes in
length, camber, sectional area, radius, and, through the application of sine bumps, local
surface shape.

[^3]Geometry Intersecting and Paneling
After PERTURB has completed applying design changes to the underlying geometry
entities, the AEROSURF module first intersects the components then panels the exterior
surfaces with a set of parametric patches. AEROSURF is a robust and automated tool
which has been designed to handle general aircraft topologies containing features such as
fuselage, wings, canards, empennage, nacelles, pylons, and diverters. AEROSURF is able
to panel diverse topologies as the HSCT configurations, business jets, high-, mid-, and
low- mounted wings, multiple nacelles, and multiple lifting surfaces. In order to panel the
HSCT geometry correctly, AEROSURF has been modified to handle cases where the wing
drops below the fuselage lower surface and the nacelles touch the wing trailing edge.
Since the paneling procedure is repeated once for each design variable at each design cycle, special attention has been paid to developing an efficient routine. The paneling of the HSCT wing/body/nacelles/ diverters configuration requires 1.1 CPU seconds on an 195 MHz Origin 2000 processor while a full configuration business jet requires 0.4 CPU seconds on a similar processor.

## 



- Intersects Geometry Components (Simple Connectivity Scheme)
- Panels Exterior Surfaces with Parametric Patches
- Variable Topology
$\quad$ * Multiple Lifting Surfaces
$\quad$ * Multiple Nacelles
- Capabilities
$\quad$ * Handles Wing Dropping Below Body
* Handles Nacelle Touching Wing
- Automated
- Efficient
* HSCT wing/body/nacelles/diverters: $1.1 \mathrm{O} 2 \mathrm{~K} \mathrm{CPU} \mathrm{sec}$.
* Business Jet wing/body/nacelle/plyon: $0.4 \mathrm{O} 2 \mathrm{~K} \mathrm{CPU} \mathrm{sec}$.

0
AEROSURF Parametric Surface Paneling
The AEROSURF parametric surface paneling for the TCA-6 wing, body, nacelles, and
diverters configuration shown in the slide consists of 39 panels with approximately 120,000
points. Since AEROSURF is an automated procedure, the paneling topology will remain
consistent for reasonable design changes. It is this topological consistency which allows
one to treat the panels as a parametric space and map changes in the parametric surfaces to
changes in the CFD surface mesh.

$$
1
$$

Surface Preprocessing and Mesh Updating
In order to map changes in the parametric panels to the CFD surface mesh point locations,
one must have knowledge of where the surface points lie in the parametric space. Since the
AEROSURF panelings are toplogically consistent, the determination of this (u,v)
parametric map can be done once as a preprocessing step. A utility, UVMAP, has been
developed which finds the locations of the initial CFD surface mesh points on the initial
AEROSURF geometry paneling using a shortest distance projection method.

> Within the design code, the MAPXYZ module is responsible for updating the surface meshes to reflect the changes in the parametric geometry paneling. By evaluating the original and modified geometry panelings at the predetermined ( $u, v$ ) values, one can determine the required motion of each surface point and create an updated set of CFD surface meshes.


- UVMAP: - Run AEROSURF on Initial Geometry
- Prepare a CFD Surface Mesh (GRIDGEN or Equivalent)
- For Each CFD Surface Mesh Point, Calculate a (u,v) Pair Within
Some Patch (Shortest Distance Projection Method) MAPXYZ:
- Perturb the Geometry Components
- Reintersect and Repanel via AEROSURF
- Apply (u,v) Map to Evaluate New CFD Surface Mesh Points
Updated CFD Surface Mesh

[^4]
Volume Mesh Deformation
Once the surface meshes have been modified to follow the desired geometry changes,
WARP-MB perturbs the remaining portion of the mesh. In order to perturb the multiblock
mesh, two capabilities are required. First, the block corners, edges, and faces must be
moved in a manner that follows the desired surface changes and simultaneously retains
mesh continuity throughout the domain. The second requirement is to move all the points
interior to each block such that the spacing distributions and smoothness of the original
mesh are retained.
The methodology used to achieve the requirement of maintaining continuity in the blocking
structure is given as follows:

1. All faces that are directly affected by the design variables (active faces) are explicitly
2. In order to maintain mesh quality, internal faces that have been flagged by the user to move in conjunction with the active faces are perturbed using weighted contributions from the active faces.
3-4. All corners and edges that touch an active face, either in the same block or in an perturbed faces are adjusted using a TFI/Arc length based algorithm.
N
Volume Mesh Deformation (cont.)
Note that much of the mesh, especially away from the surfaces, will not require mesh
perturbations and thus may remain fixed through the entire design process. Close to the
surfaces, any blocks will either contain an active face or touch a block which contains an
active face, either by an edge or by a corner. As the design variations affect the active
faces, the above scheme ensures that the entire mesh will remain attached along block
boundaries. Added complexity is required to accomplish steps 3 and 4 since the
connectivity at the various edges and corners must be indicated somehow. Currently,
pointers to and from a set of master edges and master corners are determined as a
pre-processing step. During the design calculation, perturbations to any edges and corners
are fed to these master edges and master corners which in turn communicate these to all
connected edges and corners.
Since four of the required steps (1,5-7) during the warping procedure are independent of
operations on other blocks, WARP-MB has been highly parallelized. This, combined with
the algebraic nature of the deformations, results in a scheme able to process 500,000 mesh
points per second on an Origin 2000 processor.


Updated 3-D Mesh The result of transforming the TCA- 6 baseline mesh into the NCV geometry through the
use of the geometry-driven multiblock mesh deformation scheme is shown and may be
compared with the earlier slide depicting the baseline volume mesh.

Future Work



- Enhance Current Capabilities
- Automatic Interior Boundary Control
- Subfacing
- Adapt as a Stand-Alone Structured Mesh Utility
- Develop Unstructured Mesh Deformation Capability
- Upgrade AEROSURF as a CAD-Based Geometry Kernel
$8=72=7 \pi$
HSCT High Speed Aerodynamics - BCAG WBS 4.3.1.2
Aerodynamic Design Optimization Capability


February 10, 1999
Anaheim, CA


[^5]

 This Is the Outline
 progress in the three main areas: Methods Refinements, Point Design, and Multi-Point Design, are discussed. The presentation is then concluded with a brief overview of the remaining work in 1999.
This Is the Outline

- Objectives
- Personnel

Plans

 - Evaluate multi-point optimization
- Continue development of underlying point design

The particular multi-point optimization problem being pursued by these efforts is illustrated on this figure. The single TCA design mission is shown as a trace of altitude verses time. Spotted on the mission are the three flight acceleration condition which represents an evaluation of the initial climb through the least favorable thrust to drag segment, commonly referred to as the transonic pinch point. The second point is a representative midpoint of the supersonic cruise segment, and the final point is the corresponding representative midpoint of the terminal subsonic cruise segment.


Time $\longrightarrow$

 during 1997. Additional guidance in these upgrades was gained from the results of the multi-site, optimized TCA cross check analysis exercise.
Starting in 1998, work began on further extending the method to incorporate a three surface configuration and the target multi-point optimization problem. This presentation covers much of that work.
Viscous HSR optimization efforts began in 1997 with a demonstration of an inlet flow quality design on the NCV using OVERFLOW and a process developed under IR\&D in 1996. The 1997 work showed that inlet flow quality could be improved without degrading the improvements produced by the inviscid TRANAIR process. Given the current process, and computational resource limitations, this capability is considered to be a viscous tune-up to a larger inviscid optimization in areas where a small number of variables can work a viscous issue.
Based on this success, the viscous OVERFLOW optimization capability was directed at multi-point optimization of the TCA wing leading edge in 1998. That work is described in this presentation.


## February '99 HSR Airframe Technical Review

## HSCT High Speed Aerodynamics - BCAG

## There <br> Objectives <br> <br> Are Multiple

 <br> <br> Are Multiple} he goals are shown in this figureThe six objectives which support the goals are shown in this figure. They represent the simple steps required for
the target multi-point evaluation and the complementary viscous efforts.
The three TI-prescribed conditions must be realized as stand alone point design capabilities. With these three point design capabilities, a baseline sequential design can be created. This baseline, against which the multi-point optimization will be compared, is composed of a full point design optimization at the supersonic cruise condition, followed by flap deflection optimizations at the two other conditions.

The three stand alone point design capabilities must be integrated into the existing TRANAIR multi-point capability. The principle characteristic of this step is an extension to allow for different outer mold line (OML) geometries at the three conditions. Following this step, the final simultaneous multi-point optimization can be produced. The final step in this exercise is to perform the comparison with the sequential baseline.

Concurrent with this inviscid effort, is the complementary viscous multi-point work. It is assumed that a viscous tune-up of an inviscid multi-point optimization should most likely also be multi-point.

- Develop the 3 conditions as point designs
- Produce a baseline sequential optimization result
- Develop the 3 condition multi-point capability
- Produce a multi-point optimization result
- Compare the sequential and multi-point results
- Explore N-S multi-point capabilities
Q_eøz/F' February'g9 HSR Airframe Technical Review
Many People Contribute
This presentation represents the work of many people. The full list of contributors are shown on this figure. The
 figure.
The two viscous multi-point optimization cases and the work on collarless OVERFLOW grids are the work of Gordon Blom.
Flap effects correlation analysis with wind tunnel data for Mach number 1.2 and the initial efforts at flap deflection optimization at Mach number 1.1 were performed by Rusty Conner.
Josh Elliott produced the Mach number 0.90 spillage calculations, the work on transpiration modeling
improvement, and the initial Mach number 0.90 point design.
Chris Vegter is responsible for much of the work reported here. Specifically, he produced the flap correlation and optimization work at Mach number 0.90, the development of spilling inlet modeling, the initial trimmed supersonic cruise point design, and the initial two condition TRANAIR multi-point optimization. He also has picked up the remaining optimization work at Mach number 1.1 and has been preparing to begin the initial three condition TRANAIR multi-point optimization.
The OVERFLOW reference viscous solution for TRANAIR boundary layer calibration and the implementation of enthalpy damping in viscous OVERFLOW are the work of Anutosh Moitra.
Again, everyone listed on this figure has contributed to the work that is being reported. The specific progress described below is directly or indirectly the result of the efforts of everyone shown.

Eric Adamson
Chet Nelson
Steve Chaney
Anutosh Moitra
Kevin Mejia
Bob Patton
Bob Kulfan
Progress
The progress of the past year can be divided into the three reporting categories: Methods Refinements, Point
Design, and Multi-Point Design. Descriptions of progress in these three areas follow in that order.
The inviscid optimization efforts have narrowed to a specific problem characterized by a single configuration, a
single variable set, and a single process. All evaluations are made with longitudinally trimmed configurations which include the fundamental propulsion forces and prescribed full scale viscous drag terms. When appropriate, inlet spillage is also included in the aerodynamic model. The cases and results presented here are based on this focused
approach, whether explicitly noted or not.



- Flap deflection optimization
- Improved transpiration modeling
- Viscous modeling

Optimization of flap deflections is an important enabling capability for multi-point optimization. The most important aspect, however, is the effect on drag, both as an increment and as an absolute. Drag is the aspect of flap deflection optimization that appears in the objective function. Flap deflection values that result from drag optimization are of secondary interest.
The last year's efforts can be divided into two main areas. The first was establishing the correlation between the
wind tunnel results produced during LaRC 16' test 484 and inviscid TRANAIR analysis. For both Mach number 0.90
and 1.20 , the effects of leading and trailing edge flap deflections on drag showed good correlation. When further
quantified with the weighting factors used for multi-point evaluation and optimization, the correlation is sufficiently
accurate to support multi-point design.
It is interesting to note that the correlation of flap deflection angles was also encouraging. Agreement on implied optimum leading edge flap deflections was very good. Predictions of optimum trailing edge flaps were less accurate, as should be expected for inviscid analysis, although as noted above, the drag increments for optimum trailing edge flaps were sufficiently accurate.
All the correlation work was and the subsequent optimization efforts are limited to the flaps of the outboard wing only.
Once the validity of flap deflection optimization was established with the correlation exercise, the capability was applied to all the appropriate cases. As shown here, that included the trimmed baselines at Mach number 0.90 and 1.10, point design demonstrations at these two off-design conditions, and an initial two condition multi-point optimization.
Flap Deflections Are Optimized
- (Inviscid) Correlations with LaRC 16' 484
- Outboard flaps only
- Used in ...
- Trimmed baselines, Mach $0.9 \& 1.1$
- Mach 0.9 point design
- Mach 1.1 point design
- 2 condition multi-point (Mach $2.4 \& 0.9$ )

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February '99 HSR Airframe Technical Review
Inlet Spillage is Modeled
HSCT High Speed Aerodynamics - BCAG
 surface Mach number contours with and without spillage are compared. The two solutions were obtained at the same
angle of attack, that of the spilling model at the design condition shown.
Inlet spilling has a substantial effect on the flow fields under the wing and about the nacelles. Both the compressions and expansions are stronger in the presense of inlet spilling. Increased inlet lip expansion is also very evident.
The force and moment increments between the two cases (spilling minus not) are
The effects of inlet spillage on the external flow at Mach number 0.90 are shown on this figure. Lower wing
ang
The flap deflections shown do not represent optimally trimmed configurations. They are, however, representative of optimally trimmed states and should not invalidate these increments.



## Mach 1.1 Spillage

Similarly, the effects of inlet spillage on the external flow at Mach number 1.10 are shown on this figure. Lower at the same angle of attack, that of the spilling model at the design condition shown.
Inlet spilling appears to have an even more substantial effect on the flow fields under the wing and about the nacelles at this Mach number. Again, both the compressions and expansions are stronger in the presense of inlet spilling. The compression field ahead of the nacelles actually starts with a full span, almost unswept shock.
Significant changes to the outboard wing are also indicated. Increased inlet lip expansion is also very evident.
The force and moment increments between the two cases (spilling minus not) are $=0.0036$
$=-0.00032$
The moment is taken about the aft C.G. limit, station 2281.92 and waterline 218.
The flap deflections shown represent the optimally trimmed spilling configuration.


Trimmed baselines are the result of a small constrained drag optimization which is outlined on this figure. The
variables are canard and horizontal tail incidences, C.G. longitudinal location, and configuration angle of attack. For
the two off-design conditions, leading and trailing edge flaps become active variables. The cases are constrained to
the design lift coefficient and zero net pitching moment. Reasonable approximations of the propulsion forces are also
included in the force and moment balance. The optimization problem solves for the minimum gross thrust
configuration, subject to the constraints.
To date, every case using the TI-specified tri-surface limits has driven the C.G. to the aft limit.
An image of the trim problem, both the force and moment balance, and the variables, is shown on the next figure.
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- 4 (or 6) variable drag optimization
- Canard incidence
- Horizontal tail incidence
- C. G. location
- Angle of attack
- Leading edge flap deflection (Mach $0.9 \& 1.1$ )
- Trailing edge flap deflection (Mach $0.9 \& 1.1$ )

[^6]
The propulsion forces are also included in the lift and pitching moment balances.
An additional variable shown here is nozzle pitch angle. The current implementation assumes that this is a single value for all engines. An additional simplifying assumption is that the thrust vector can change without changing the nozzle exterior contours. The nozzle pitch angle variable is active during general optimization, but not when calculating the trimmed baseline starting configurations.


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## $p_{0}$

The principle effort on cycle time reduction was to improve the transpiration modeling used to calculate optimization sensitivities by TRANAIR. Two transpiration models are currently used. Most of the design process uses two-term transpiration modeling, the most accurate, but limited to small movements. A separate optimization problem is posed for wing twist and shear using one-term transpiration modeling to efficiently allow the large ment characteristic of those variables starting from linearly designed HSCT configurations. The current standard process is to start with a two-term cycle with all variables active, and follow with the one-term cycle, just for wing twist and shear. The process then uses three additional full variable set, two-term cycles. Some of the benefit of these subsequent two-term cycles appears to be recovering from the inaccuracies of the one-term cycle.
One way to improve the efficiency of the TRANAIR design process would be to develop a new mixed
transpiration model which matches the modeling to the variable movements, so that dedicated one-term cycles are unnecessary. This is the approach that was taken this year. The specific implementation uses one-term modeling for wing twist and shear movements and two-term modeling for wing camber and thickness. The new model has been shown to be an improvement on the 3-dimensional test cases checked. An example is shown on the next figure.
Eliminating the one-term cycle, and perhaps one of the subsequent two-term cycles, would reduce the flow time required for the current process by twenty to forty percent.
Transpiration Modeling Was Improved - Mixed transpiration model developed

- Modeling matched to movement
- 1-term -> wing twist \& shear
- 2-term $\rightarrow$ wing camber \& thickness
- Could eliminate 1 or 2 TRANAIR cycles
- $20-40 \%$ flow time reduction
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HSCT High Speed Aerodynamics - bcaG
The three efforts to increase the impact of viscous modeling in optimization are listed on this figure. The first
of the nacelle/diverter region at the cruise condition. This year a reference viscous solution for a simplified problem
was created. The simplified problem is an inboard TCA nacelle/diverter installed on a flat plate at zero angle of
attack. A viscous OVERFLOW solution for this case was produced and then boundary layer parameters were
extracted from the OVERFLOW viscous layer. This reference boundary layer data is currently being used to calibrate
the TRANAIR coupled boundary layer on the same simplified problem. The second effort was aimed at the development of a collarless OVERFLOW grid approach. This technique would
need to use collar grids and the accompanying interactive Pegsus step. A collarless TCA wing/body grid was
eloped with featured matched overlapping grid cells. An OVERFLOW solution was obtained for this grid, but it
found that there were small differences in overall forces and moments, relative to the standard grid solution, and
The flow details in the wing/body intersection region were also different. It was concluded that this particular
roach was not sufficiently accurate to produce the finite difference sensitivities needed to support optimization.
Thert was to increase the efficiency of OVERFLOW analysis through the implementation of an enthalpy damping technique. Building up to increasingly complex configurations, the work concluded this year with a demonstration of a two to three improvement in the convergence rate of a TCA wing/body/nacelle/diverter configuration case at the cruise condition.
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- OVERFLOW data for calibration of TRANAIR coupled
boundary layer
- Collarless OVERFLOW grid
- Not successful
- OVERFLOW implementation of enthalpy damping
- 2-3 convergence improvement, w/b/n/d, Mach 2.4

Point Design(s) Have Continued
- "Tri-surface TCA" configuration
- TCA wing/body/nacelles/diverters
- PTC canard/empennage
- Longitudinally trimmed
- Same as baselines
- 3 conditions
Supersonic cruise
- Transonic acceleration





[^7]

[^8]




[^9]

A full variable set trimmed optimization of the tri-surface TCA at the subsonic cruise condition was started this
past year, and the case is almost complete. Some characteristics and intermediate results are shown on this figure.
The condition is that specified by TI for the multi-point problem except that the lift coefficient has been reduced so
that the configuration angle of attack is closer to that of the supersonic cruise condition. This consistency in flight
attitude at the two cruise points would be a characteristic of a design, more well balanced than the TCA. Maintaining
this geometric relationship between the two flight conditions was judged to be more important than strict adherence to
the design mission lift coefficient at the subsonic cruise condition. Note that inlet spillage is included in the flow
model. As indicated previously, the case uses the common variable set and constraints.

[^10]

- 441 variables
- 34 k constraints
- 2.3 TRANAIR cycles
- 4.1 counts of thrust improvement
- Inviscid analysis
- Relative to trimmed baseline
- Shape changes dissimilar to cruise point designs

An upper surface view of the intermediate result is shown on this figure. No apparent shape changes are visible in
this view. The transition from inboard leading edge to outboard leading edge flap is shown, and illustrates that this is
modeled without the unporting of the actual system. The same modeling approximation is used for the trailing edge
flap also.
Two geometric results of the optimization thus far are, one, an increase in leading edge camber on the inboard
wing, and two, a general thickening of the outboard wing. The latter occurs primarily in the inspar region, and
appears to be reducing the flap hinge turning angles.
There is very little change to the body, or the wing twist and shear.

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This figure is the corresponding lower surface view of the intermediate subsonic point design geometry. As with
the upper surface, little change is apparent. There are, however, small changes occurring about the nacelle/diverter
regions in apparent response to the pressure fieldsamics - BCAG
be similar tho what occurs at the supersonic cruise condition in terms of ond cause and effect, but are at different appear to to
since the compressions and expansions driving the features are at different locations.


The third and final multi-point optimization flight condition is the transonic acceleration point outlined on this figure. This TI -specified condition represents the transonic pinch point of the initial climb out and, like the subsonic cruise condition, has significant inlet spillage.
The variable and constraint sets are the same as with the other cases.
This case is ready except for the trimmed baseline starting point, which is currently in work.


Initial efforts on multi-point optimization were started and/or completed this past year, and are outlined on this figure. Two cases of viscous multi-point design using OVERFLOW were accomplished. The cases demonstrated the feasibility and potential benefits of, and resources necessary for this viscous multi-point technology.
The initial TRANAIR HSCT multi-point optimization was also completed. This two condition developmental case exercised and demonstrated the process which will be employed for the upcoming three condition case.
Finally, preparations for the culminating three condition TRANAIR multi-point optimization exercise were started.
These three area of progress are next discussed in greater detail.


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\begin{aligned}
& \text { February 'g9 HSR Airframe Technical Review } \\
& \text { HSct High Speed Aerodynamics - bcaG } \\
& \text { Two closely related viscous multi-point optimization cases were accomplished this past year and are outlined on } \\
& \text { this figure. The cases used viscous OVERFLOW as the analysis tool. The optimization technique is the same hybrid } \\
& \text { genetic approach used in previous work on inlet flow quality. The configuration used for both cases is the TCA wing } \\
& \text { /body, and the designs were limited to the leading edge region of the inboard wing. } \\
& \text { The initial case is a three condition subsonic cruise design. The multiple conditions are used to optimize the wing } \\
& \text { leading edge over a wide range of lift coefficients. } \\
& \text { The second case is a two condition optimization based on flight conditions in the both the supersonic and subsonic } \\
& \text { cruise regimes. This case is closer to the type of multi-point problem being pursued with the inviscid TRANAIR } \\
& \text { process. }
\end{aligned}
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Viscous Multi-Point Was Demonstrated

- TCA wing/body configuration
- Inboard wing leading edge designs
- First case
- 3 conditions ( $C_{\text {L's }}$ ) at Mach 0.9
- Second case
- 2 conditions : Mach $2.4 \& 0.9$

HSCT High Speed Aerodynamics - BCAG
Case 1
Viscous multi-point numerical optimization was tested by using a hybrid genetic optimization driver connected to the OVERFLOW CFD solver. The TCA wing/body inboard wing leading edge was modified to minimize an objective function that required three CFD solutions at Mach number 0.90 for different angles of attack.
The figure shows the wing leading edge surrounded by a grid of boxes in red. These boxes define a region where free-form deformation (FFD) methods were used to modify the geometry. Four design variables controlled the movement of FFD control points within each box. The inboard wing between WBL 98.74 and WBL 471.56 forward of 2 percent of chord was modified. In order to avoid a large drag penalty at Mach number 2.40 that might be produced by an expected increase in leading edge camber, the leading edge point was held fixed and only the shape modified.



Multi-Point Objective Function
The figure graphically depicts the function being minimized. An objective function is formed from the integral shown by using three OVERFLOW solutions and fitting a cubic spline to the resultant lift and drag. The OVERFLOW solutions were for angles of attack of two, five, and eight degrees at Mach number 0.90.


Objective Function History
The figure shows a history of the optimization process. The objective function value is shown for each evaluation in sequence. The TCA baseline value is 0.004191 and the optimum value is 0.004176 for a reduction of 0.000015 . The optimization required 99 function evaluations each using three OVERFLOW solutions. Each OVERFLOW solution was a restart from a TCA baseline solution and required approximately 55 minutes of NAS C90 CPU time.



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The figure shows the wing leading edge airfoil section geometry for the TCA and for the optimized configuration at
several representative spanwise locations. Note that the leading edge points were not moved and the upper and lower
surfaces aft of the leading edge were recontoured in minimizing the objective function.

Case 2

The figure shows the region of the inboard wing that was redesigned. Six design variables were used to modify the leading edge geometry. The design variables provided for leading edge shape variation as well as leading edge camber variation by moving the control point locations of each of the FFD boxes.


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WBL 203

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Two condition viscous multi-point optimization

Flaps up TCA wing/body



February '99 HSR Airframe Technical Review
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The initial TRANAIR multi-point HSCT optimization was completed this year and is outlined on this figure. As noted previously, this case was also based on the common ti surfa
 tri-surface TCA condition, and the subsonic cruise condition is the tri-surface TCA condition. The objective function weighting factors were based on preliminary recommendations, rather than the final TI-specified values, $6: 1$ rather than 1:0.3.
As was the case with the cruise point design, over half of the 445 variables were lost due to active constraints and variable bounds. Only three TRANAIR cycles were completed, which was sufficient to verify the mechanics of the process. As a result, the slight decrease in thrust reduction relative to the cruise point design is inconclusive. This effect of multi-point will be evaluated with this coming year's three condition exercise.
It is encouraging that a drag improvement was found at the subsonic cruise condition. The magnitude of the improvement is not a final result, however. As noted above, this was a process development case, and is only . representative. A more valid evaluation will be made with the three condition multi-point optimization exercise.
Relative to cruise point designs, no new geometric effects of optimization were observed on this case.


[^12]
> design described earlier in this presentation, is still occurring.



[^13]| ne. | February '99 HSR Airframe Technical Review |
| :---: | :---: |
| Hscr High |  |
| Three | ndition TRANAIR Shakedown Case Has Almost Started |

[^14]
The remaining activities of this subtask are described on the following figures. The work is divided into the completion of the multi-point evaluation and the other remaining activities.

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February 'g9 HSR Airframe Technical Review
Ther High speed Aerodynamics - bcag
basic evaluation requires three main pieces. First, there are the common starting points, referred to as the trimmed
baseline configurations. These are not the baselines against which the results of multi-point optimization will be
compared, but will be used as baselines for the sequential optimizations. They are also the common starting points for
both the sequential and multi-point optimizations. The second main part is the sequential optimization. This is
composed of first, a full variable set, cruise point optimization. The result of this point design then passes through a
trim drag optimization with flaps at the two other flight conditions. These three results, combined in the weighted
multi-point objective function, constitute the sequential baseline against which multi-point will be evaluated. The
third main part is the simultaneous multi-point optimization. It also starts from the trimmed baseline configuration
and the three flight conditions are inherently combined in the weighted objective function.
Through coordination with the other CA organization participating in the multi-point evaluation exercise, NASA Ames, the task has been expanded to two separate, but related, multi-point problems. The problem described in this presentation is referred to as the extended problem in this context. A slightly simpler problem, representing the common capabilities and/or obtainable capabilities of both sites, is referred to as the core problem. The extended problem will allow each site to demonstrate local strengths and the largest possible optimization increment. The core problem will provide a more rigorous comparison of the two optimization processes.
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\begin{aligned}
& \text { As with past cruise point design exercises, the final evaluation of the multi-point optimization results will be } \\
& \text { viscous analysis. As a minimum, this will require } N-S \text { analyses of nine configurations, the three trimmed baseline } \\
& \text { configurations, the three sequential results and the three simultaneous results. Resource issues limit this step to only } \\
& \text { one of the two problems. The current plan is to provide viscous verification of the extended problem. } \\
& \text { There are two obvious open issues with viscous verification of inviscid, trimmed, multi-point results. First, } \\
& \text { optimum inviscid flap deflections are not optimum viscous flap settings, particularly for the trailing edge flap. } \\
& \text { Second, a trimmed inviscid configuration will not be trimmed in viscous analysis. Resolution of these issues will be } \\
& \text { made as data is obtained. Additional viscous analyses may be required. }
\end{aligned}
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Multi-Point Exercise Will Be Completed


- Extended problem only

- Design review



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High Speed Aerodynamics, Long Beach
Le-Time Reduction using
Geojoe Kuruvila
Robert P. Narducci


Design Cycle-Time Reduction using TLNS3D-Adjoint
Significant cycle-time reduction has been achieved by using the TLNS3D-Adjoint code for the aerodynamic shape optimization of High Speed Civil Transport (HSCT) configurations. Using the adjoint of TLNS3D (Euler) flow analysis code, aerodynamic gradients are computed, both accurately and efficiently, for hundreds of design variables. Results obtained by using TLNS3D-Adjoint in the aerodynamic shape optimization of HSCT Technology Concept Airplane (TCA) configurations are presented.
Outline
High Speed Aerodynamics, Long Beach



- Summary
Objective

[^16]$$
{ }_{\text {Objective }}^{\text {High Speed Aerodynamics, Long Beach }}
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- Reduce the cycle-time for HSCT aerodynamic
shape optimization
- Develop an adjoint code for TLNS3D (Euler)
- Compute aerodynamic gradients accurately and
efficiently
- Use hundreds of design variables
TLNS3D-Adjoint
TLNS3D-Adjoint is a suite of three codes which, computes the aerodynamic gradients of design variables, both accurately and efficiently, using the adjoint method.
The first code, TLNS3D, is the standard version of the NASA developed flow analysis code. Although it can analyze fully turbulent viscous flows, the adjoint of only its inviscid (Euler) component has been created. Given the grid, $X$, around an airplane configuration and the freestream conditions, the Euler equations along with the boundary conditions, $R(Q, X)=0$, are solved, using TLNS3D, to obtain the flow solution $Q$ and consequently the cost function $F(Q, X)=0$. The drag, lift, moments, etc. are examples of the cost function.
The second code, ATLNS3D, solves the adjoint of the Euler equations and the corresponding boundary conditions for the userspecified cost function, $F$. In addition to the grid, $X$, and the freestream conditions, ATLNS3D requires the flow solution, $Q$, to obtain the adjoint solution, $\lambda$. The third code, GTLNS3D, uses the flow and adjoint solutions along with the grid sensitivity, $\partial X / \partial b_{i}$ to compute the gradient of the
cost function, $\partial F / \partial b_{i}$, with respect to each design variable. GTLNS3D is executed as many times as the number of design variables.
TLNS3D - Adjoint

TLNS3D / ATLNS3D
In ATLNS3D, the discretization of the adjoint equations together with the dissipation terms is consistent with that of TLNS3D (Euler).
Both TLNS3D and ATLNS3D use central-difference discretization and blended $2^{\text {nd }}$ and $4^{\text {th }}$-order dissipation. They use Runge-Kutta
time-stepping to advance the solution and multigrid to accelerate the convergence to steady state. Both support structured multiblock
patched grids and they can be executed on either a serial or a parallel machine.
TLNS3D / ATLNS3D
 - Central-difference discretization
- Blended 2 nd and 4th-order dissipation
- Runge-Kutta time-stepping
- Multigrid acceleration
- Multiblock patched grids
- Serial and Parallel (MPI) versions
AEROSHOP Framework
AEROSHOP (AEROdynamic SHape OPtimization) is a Boeing - Long Beach developed shape optimization system. It features a modular structure and runs on many platforms including the Cray C-90, J90, T3E, the SGI Origin 2000, and clusters of Hewlett Packard workstations. In a top-level viewpoint, the design problem definition is input to the AEROSHOP core executable. The optimizer within the core is MDOPT, a Boeing proprietary version of NPSOL. When the optimizer needs to perform a search direction or a line search calculation, the core relinquishes control to AEROSHOP scripts. The scripts are UNIX shell commands designed to perform a function. At the end of a script's execution, control of the optimization is returned to the core. The scripts can be custom-built to incorporate well established, heritage technologies or new, state-of-the-art technologies. The modularity of the AEROSHOP scripts also allows codes to take advantage of parallel environments.
In the adjoint approach, any time the flow solution is required, the evaluation script is called to execute TLNS3D. When AEROSHOP needs the gradients, the gradient script first calls the evaluation script to execute TLNS3D and then calls auxiliary scripts to execute the adjoint module, ATLNS3D and the gradient module, GTLNS3D.
AEROSHOP Framework
$\longrightarrow \ldots \ldots$ Migh Speed Aerodynamics, Lono Beach

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TLNS3D TLNS3D is the standard version of the NASA developed flow analysis code. It is efficient, versatile and easy-to-use and is widely
used in the aerospace industry in a production environment.
For grid topologies where all the blocks are point-matched, TLNS3D reads the grid-file and the map-file where the boundary
conditions are specified. The flow solver parameters, like the freestream conditions, number of iterations, the CFL number, the
Runge-Kutta coefficients and the dissipation coefficients are read from a third file. For grid topologies where there are patched
interfaces, two additional files that specify the interpolation coefficients are also read. When the specified number of iterations are
completed, the flow solution is output into a file.


| Grid/Patch Information |
| :---: |
| Boundary Conditions |
| Flow Solver Parameters |

- 

ATLNS3D
ATLNS3D is similar to TLNS3D in structure, versatility and ease-of-use. Its performance and computer resource requirements are similar to that of TLNS3D.
In addition to all the input that TLNS3D reads, ATLNS3D needs the flow solution obtained from TLNS3D. The cost function, for which the adjoint solution is sought, is also specified via an input file. Note that, the CFL number, the Runge-Kutta coefficients and the dissipation coefficients used are identical for both TLNS3D and ATLNS3D.
ATLNS3D

$\longrightarrow$ High Speed Aerodynamics, Long Beach
GTLNS3D
GTLNS3D computes the gradients of the cost function with respect to each of the design variables. It is executed as many times as the number of design variables.

$\qquad$ obtained from ATLNS3D

execur mas
GTLNS3D $\ldots$ High Speed Aerodynamics, Long Beach
Grid Sensitivity Options
The sensitivity of the grid to each design variable can be obtained using finite-differences or computed analytically. Generally, it is
tedious to differentiate the grid generation process by hand. Therefore, an automatic differentiation tool like ADIFOR (Automatic
Differentiation for FORTRAN) is used to obtain analytical sensitivities. Although analytical grid sensitivities are accurate, they are
more expensive to compute than finite-difference sensitivities.
Since the aerodynamic shape optimization seeks to find the best OML (Outer Mold Line) of an airplane, it maybe sufficient to use
only the sensitivity of the grid representing the airplane surface. Some evidence supporting this is discussed later in the paper. The
CPU time saved, by using surface sensitivities, can be significant when hundreds of design variables are used.
Grid Sensitivity Options
Grid Sensitivity Options

## 

Finitedifference

| Volume grid sensitivity | Need step-size | Using ADIFOR |
| :---: | :---: | :---: |
| Surface grid sensitivity | Need step-size | Using ADIFOR |
|  | Cost, Accuracy |  |

1
Cost Functions
TLNS3D-Adjoint currently supports the following cost functions. In addition to finding the gradients of drag, $D$, lift, $L$, and

Computer Platforms
TLNS3D-Adjoint has been tested on Cray C-90 serial, SGI Origin 2000 parallel, and HP Exemplar parallel computers. In the parallel mode, it uses the Message Passing Interface (MPI) protocols to communicate among processors. The MPI version of TLNS3DAdjoint has also been successfully executed on a cluster of HP workstations.
Computer Platforms
High Speed Aerodynamics, Long Beach

- Serial Machines
- Cray C-90
- Parallel Machines
- SGI Origin 2000
- HP Exemplar
- Workstations
- HP Cluster



## Ю!า

Gradients of D/L obtained using TLNS3D-Adjoint, for 26 typical design variables used in the shape optimization of an HSCT, are
compared with those obtained using finite-differences. The first 18 design variables were located on the wing and the rest on the
fuselage. These design variables included twist, camber and thickness of the wing, cross-sectional shape and camber of the fuselage
and position of the wing on the fuselage. The finite-difference gradients were computed, using CFL3D flow analysis code, with
carefully selected step-sizes and converged solutions. The grids, corresponding to the change in each design variable, were obtained
using the grid perturbation code, FlexMesh.
For the adjoint gradient, both TLNS3D and ATLNS3D solutions were converged by over six orders of magnitude. The grid
sensitivities corresponding to the change in each design variable were obtained by finite-differences. The perturbed grids were
obtained using CSCMDO.
In spite of the many differences, from flow solvers to grid perturbation schemes, the gradients computed using the adjoint and the
finite-difference methods agree well.

TCA Wing/Body Optimization
The first application presented is the optimization of the TCA wing/body configuration at $\mathrm{M}=2.4$. The objective was to minimize the ratio of drag-to-lift ( $\mathrm{D} / \mathrm{L}$ ) at a specified lift. The design-space was made up of 400 geometric functions and the angle-of-attack. These geometric functions included twist, camber, thickness, etc., of the wing, cross-sectional shape, camber, e., of the fusela , These constraints were obtained from the TCA Configuration Description Document.
$\underline{\text { TCA Wing/Body Optimization }}$ High Speed Aerodynamics, Long Beach $^{\text {TC }}$
(4)




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## History

The convergence of the analysis (Euler) and the adjoint equations, for the TCA wing/body configuration at $\mathrm{M}=2.4$, are shown. For a grid that has $193 \times 33 \times 73$ points in the chordwise, normal, and spanwise directions, respectively, the residuals of both the systems decreased by over six orders of magnitude in about 200 iterations. The convergence rates of the analysis and adjoint equations are comparable. Note that, the solver parameters like the CFL number, the Runge-Kutta coefficients, and the dissipation coefficients, used for the solution of both these systems, are identical.
Convergence History
High Speed Aerodynamics, Long Beach
TCA Wing/Body, Grid: 193x33x73, $M_{\infty}=2.4, \alpha=3.5^{\circ}$

200
150
100
Iteration
(ןenpisəy) 607
TCA
Optimization History The optimization history shows that the objective function decreased by about $5 \%$ in 24 design iteration. The drag decreased by about 5 counts and the lift constraint was satisfied. Note that the drag shown is only the pressure drag. In this optimization, only the surface grid sensitivities were used to compute the gradients. These sensitivities were computed analytically by differentiating the tools for surface grid perturbation using ADIFOR (Automatic Differentiation of FORTRAN)
An identical optimization exercise was performed, using the sensitivity equations method*. The sensitivity equations were obtained by automatically differentiating the CFL3D analysis code using ADIFOR. In this study, the analytical volume grid sensitivities were used to compute the gradients. In spite of these differences, the final design was similar in shape and performance to the one obtained using TLNS3D-Adjoint with surface grid sensitivities. This indicates that, in shape optimization problems, surface grid sensitivities are sufficient to improve the performance of configurations. Whether this argument holds true for more complex wing/body/nacelle/diverter configurations is yet to be demonstrated.
Computing only surface grid sensitivities requires considerably less CPU hours than computing volume grid sensitivities, especially for large number of design variables. Typically, the cost of computing volume grid sensitivities, using the geometry and grid perturbation tools at Boeing, Long Beach, is about twice that of computing only the surface grid sensitivities.
Optimization History
$=$ High Speed Aerodynamics, Long Beach
$3 \times 33 \times 73, M_{\infty}=2.4$
55 Constraints
TCA Wing/Body, Grid: $193 \times 1$ Design Variables,


000I $\times 7 / \square$
CPU Time / Memory on Cray C-90
This wing/body optimization was performed on von Neumann, the Cray C-90 at NAS. The CPU time and the memory, required to obtain a converged solution, for both TLNS3D (analysis) and ATLNS3D (adjoint) codes, are comparable. The gradient module, GTLNS3D, required about 30 seconds to compute the gradient of each design variable. Therefore, the CPU time for computing the gradients of all the design variables exceeded 3 hours. Yet, the total cost of 24 design iterations was only 136 hours, about 15 times less than an optimization using the finite-difference method.
Of the 30 seconds used by GTLNS3D, 28 were spent in computing the analytical surface grid sensitivity of each design variable. If the volume grid sensitivities were used, the computational cost would have doubled.

Beach


* Estimated CPU time for finite-difference method > 2000 hours
TCA Wing/Body/Nacelle/Diverter Optimization
The second application presented is the optimization of the TCA wing/body/nacelle/diverter configuration at $\mathbf{M}=2.4$. The objective was to minimize the pressure drag at a specified lift. The design variables were 319 geometric functions and the angle-of-attack. These geometric functions included twist, camber, thickness, etc., of the wing, cross-sectional shape, camber, etc., of the fuselage, and nacelle shear variables. In addition to the lift, 84 geometric parameters of the configuration were constrained. These constraints were obtained from the TCA Configuration Description Document.
TCA Wing/Body/Nacelle/Diverter Optimization

Optimization History
This TCA wing/body/nacelle/diverter configuration was modeled using a 19 -block patched grid. The starting point of the
optimization was a previously optimized configuration called TCA Cycle 3 . Starting with the TCA baseline, the Cycle 3
configuration was obtained by a sequence of optimizations that included both finite-difference and adjoint methods*. The
optimization history shows that, the decrease in drag from Cycle 3 is about 2 counts and the lift constraint is satisfied. Overall, this
configuration has about 7.3 counts less drag than the baseline. Note that the drag shown is only the pressure drag. In this
optimization, volume grid sensitivities were obtained using finite-differences.
Optimization History
High Speed Aerodynamics Lon Beach


## 00K points, $M_{\infty}=2.4$ 85 Constraints



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In the third application shown here, the TCA wing/body/flaps configuration at $\mathrm{M}=0.9$ is optimized for minimum pressure drag at the angle-of-attack. In addition to the lift, 55 leading-edge spar-thickness constraints were also imposed.
TCA Leading-Edge/Flaps Optimization - Wing/body/flaps at $M_{\infty}=0.9$

- Minimize pressure drag
- $\quad$ Design variables
$\quad-22$ Leading-edge shape functions
$\quad-5$ Leading- and trailing-edge flaps
$\quad-$ Angle-of-attack
- $\quad$ Constraints
$\quad-55$ Leading-edge spar thickness
$\quad-$ Lift - Wing/body/flaps at $M_{\infty}=0.9$
- Minimize pressure drag
- $\quad$ Design variables
$\quad-22$ Leading-edge shape functions
$\quad-5$ Leading- and trailing-edge flaps
$\quad-$ Angle-of-attack
- $\quad$ Constraints
$\quad-55$ Leading-edge spar thickness
$\quad-$ Lift - Wing/body/flaps at $M_{\infty}=0.9$
- Minimize pressure drag
- Design variables
$\quad-22$ Leading-edge shape functions
$\quad-5$ Leading- and trailing-edge flaps
$\quad-$ Angle-of-attack
- $\quad$ Constraints
$\quad-55$ Leading-edge spar thickness
$\quad-$ Lift
Convergence History

[^17]Convergence History
Wing/Body/Flaps, Grid: $113 \times 33 \times 89, M_{\infty}=0.9, \alpha=4.5^{\circ}$
 500
Optimization History
The starting point of this optimization was the TCA baseline with optimized flaps. The optimization history shows that the decrease in drag from the starting point is about 2.5 counts and the lift constraint is satisfied. Note that the drag shown is only the pressure drag. Volume grid sensitivities obtained using finite-differences were used in this optimization.
Optimization History чэeәg 6ио7 'sэ!ueuरpoлә৮ pəəds 4б!н $13 \times 33 \times 89, M_{\infty}=0.9$
56 Constraints
0.184


### 0.0080



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Wall Clock Time / Memory on SGI Origin 2000
The leading-edge/flaps optimization was performed on Turing, the SGI Origin-2000 at NAS, using 14 processors with the parallel
 solution, for both TLNS3D and ATLNS3D codes, are comparable. GTLNS3D was executed on all 14 processors to obtain the gradients of 14 design variables at a time. The total cost of the 10 design iterations was only 30 hours, about 2.5 times less than an optimization using the finite-difference method. This factor is not as large as the previous wing/body optimization, since the number of design variables is considerably fewer in this case.
Wall Clock Time / Memory on SGI Origin 2000
Wing/Body/Flaps, Grid: $113 \times 33 \times 89, M_{\infty}=0.9$

TCA Multipoint Optimization
 The objective was to minimize the weighted sum of the pressure drags at the two Mach numbers. The design variables included 22 leading-edge thickness and droop function, the 5 leading and trailing-edge flaps and the angle-of-attack. While the flaps were optimized only at Mach 0.9, the leading-edge thickness and droop functions, and the angle-of-attack were optimized at both the Mach numbers. Lift and 55 leading-edge spar thickness were constrained.
$\underline{\underline{T C A}}_{\text {TCA Multipoint Optimization }}^{\text {High Speed Aerodynamics, Long Beach }}$
sea

2.4)




spar thicknes
55 Leading-edge
Lift

- Angle-of-attack
- Constraints

1
- 

TCA Multipoint Optimization
The analysis of the wing/body configurations of the baseline TCA is compared with the ones obtained from the single point
optimization at Mach 0.9 and the multipoint optimization at Mach 0.9 and 2.4 . While the single point optimization reduced the drag at
Mach 0.9 by about 2.5 counts, the drag at Mach 2.4 went up by about 1.5 counts. In the multipoint optimization, the drag at Mach 0.9
was reduced by about 1.5 counts, without any penalty at Mach 2.4 .
TCA Multipoint Optimization

Drag Polars of W/B Configurations Resulting from Single and Multipoint Optimizations


Technology Transfer

[^18]
This technology being used in other Boeing
programs
TLNS3D-Adjoint successfully used to improve
performance of a space vehicle

- Wing/body/tail at $M_{\infty}=1.2$
- Cost function: Function of pitching moment
Summary
The development of TLNS3D-Adjoint (Euler) is complete. It computes the aerodynamic gradients for large number of design
variables, both accurately and efficiently. The memory requirements of the adjoint (ATLNS3D) and the gradient (GTLNS3D)
modules are comparable to that of the analysis (TLNS3D) module. By using TLNS3D-Adjoint in AEROSHOP, significant reduction
in design cycle-time has been realized.
TLNS3D-Adjoint has been applied to several aerodynamic shape optimization problems at supersonic and subsonic Mach numbers. It
has been applied to HSR as well as non-HSR problems.


Progress Towards Viscous Design Optimization
Using Automatic Differentiation


High Speed Research Airframe Review
Anaheim, California
February 8-11, 1999
P. Sundaram and Shreekant Agrawal
Configuration Aerodynamics
Phantom Works, The Boeing Company
Long Beach, CA
Acknowledgment
Several people helped during the course of this investigation. Arthur Taylor of Old Dominion
University was stationed at Boeing Long Beach from January to June 1998 under a subcontract through
Eagle Aeronautics, a small disadvantaged business located at Hampton, Virginia. He was primarily
assisting in the development of the ADIFOR versions of the BPW-LB grid tools. Alan Carle of Rice
University was instrumental in developing the ADJIFOR version of CFL3D code as part of a subcontract
to Rice University.

Terminology ADIFOR and ADJIFOR are preprocessors used to obtain the computer generated derivative
versions of FORTRAN function codes. ADIFOR was originally developed by a team of researchers from
Argonne National Laboratory and Rice University. They include Chris Bischof, and Paul Hovland of
Argonne National Laboratory and Alan Carle of Rice University. More recently, the ADJIFOR
development is actively pursued by Alan Carle and Mike Fagan of Rice University. Although ADJIFOR
and reverse mode ADIFOR are interchangeably used in the present context, it is important to note the
difference that ADJIFOR is specifically developed for iterative solvers while reverse mode ADIFOR is
applicable for a more general class of problems.


Background
Design cycle time in a nonlinear shape optimization process is mainly dependent on the time taken for the objective function sensitivity calculation. Considerable effort has been spent in the past few years to develop efficient techniques for calculating the sensitivities. Primary among them is the method of deriving the adjoints by posing the original continuous form of the problem as a calculus of variation problem. This requires long and tedious analytical derivations and hand-differentiation of the underlying partial differential equations. The adjoint variables for the Euler equations already derived and codes based on this are available. For Navier-Stokes equations, the theory of adjoints is much more complicated and becomes almost intractable. More so, for different turbulence models the handdifferentiation method of adjoints is impractical.

[^19]\[

$$
\begin{aligned}
& \text { Background } \\
& \text { Fast and accurate sensitivities calculation is vital for cycle } \\
& \text { time reductunamics, Long Beach } \\
& \text { Hand-differentiated adjoenor shape optimization code for N-S codes takes too long } \\
& \text { ADIFOR 2.0 from Argonne National Lab. and Rice University } \\
& \text { can provide accurate sensitivities of any FORTRAN code } \\
& \text { Alliance between LaRC, BPW-Long Beach, and Rice to } \\
& \text { advance the ADIFOR technology for shape optimizations }
\end{aligned}
$$
\]



 inclusion of these details will produce a more realistic representation of the local flow at the correct flight conditions. Although it is understood that it is an expensive proposition to perform Navier-Stokes optimization of the full configuration, it is also necessary to obtain all the available performance benefits in a realistic flow-field. Also, the Navier-Stokes design might produce a more desirable smoother geometry compared to the earlier Euler designs. The application of several technology advancements in the area of gradient calculations combined with the parallel processing might render this ambitious undertaking realizable. With this motivation, several new technologies are used in this task.


Outline
Outline
High Speed Aerodynamics, Long Beach
Objective
This chart outlines the objectives of the present task
High Speed Aerodynamics, Long Beach
Objective

$$
\begin{aligned}
& \text { Develop technology and tools required for efficient } \\
& \text { nonlinear viscous optimization using Navier-Stokes solver } \\
& \text { Apply the technology to full configuration design }
\end{aligned}
$$

Approach The nonlinear shape optimization scheme used at BPW-LB requires flow sensitivities of an
objective function to drive the optimizer. The approach chosen by the BPW-LB team is to obtain the
viscous flow sensitivities using the automatic differentiation (AD) approach. The CFL3D Navier-Stokes
equations solver is the chosen flow solver. Parallel processing is needed not only to yield the large
computer resources required, but also for achieving fast turn-around time of the design cycles.
Approach
High Speed Aerodynamics, Long Beach

- Apply ADIFOR to obtain computer generated sensitivity
codes used in nonlinear shape optimization at BPW-LB
- Validate ADIFOR / ADJIFOR-based codes for calculating
accurate grid and flow sensitivities
- Perform HSCT shape optimization with large number of
design variables using CFL3D Euler and N-S sensitivities
- Utilize parallel processing for fast turn-around time
Automatic Differentiation Using ADIFOR One of the major accomplishments in the area of application software technology that has
differentiation has a large spectrum of applications for problems arising in numerical weather prediction, nuclear engineering, and numerical optimization. Several tools have been developed for automatic
 The Automatic Dlfferentiation of FORtran (ADIFOR) software has been under development since the early 90's jointly at Argonne National Laboratory and Rice University with the support of DOE and NASA.
The ADIFOR preprocessor augments a FORTRAN function code and generates another
FORTRAN program that can compute both the original function and gradients. The preprocessor generates analytical chain-rule differentiation based on the user-specified independent and dependent variables. Since the differentiation is analytical, no step-size is required which is a major advantage compared to the divided-difference method of calculating gradients. Since ADIFOR performs analytical differentiation, it can use single-precision arithmetic to reduce the memory requirement on some computers. Because ADIFOR is memory-intensive, this property is very attractive. Also, ADIFOR can be easily applied to most FORTRAN code of any size. Last, but not the least, it is a software available open to public.

[^20]Automatic Differentiation Using ADIFOR

- Augments FORTRAN code to generate a derivative code
- Accurate analytical sensitivity calculations
- Uses chain-rule differentiation
- No step-size (an issue in finite-difference approach)
- Can use single-precision to reduce memory requirement
- Easily applied to any FORTRAN code
- ADIFOR is available in public domain
Forward Mode ADIFOR

> In the forward mode ADIFOR, or simply ADIFOR, the program flow maintains the function program
> flow direction and calculates the directional derivative of the dependent variables w.r.t. the independent variables.

> The main drawback of the ADIFOR software is the large computer resource requirement
> compared to some of the state-of-the-art sensitivity calculation methods such as the hand-differentiated
adjoint codes for sensitivity determination. The memory and CPU time for ADIFOR-processed code are
directly proportional to the number of design variables used. As a result, the computer resource
requirements for ADIFOR derivative code can quickly become staggering for large problems with
hundreds of design variables. However, significant developments are taking place to allay these
concerns. The reverse-mode ADIFOR is becoming practical for large problems. These discussions will
be alluded to later in the presentation.

Quatine.
Forward Mode ADIFOR in Shape Optimization The application of ADIFOR for a sequence of function codes is shown in this chart. Since gradient
calculation is a linear operation, it can be mathematically represented by matrix multiplication. The
gradient information is passed between codes through a sequence of matrix multiplication of a Seed
Matrix (SM). The SM consists of $p$ rows and $q$ columns where $p$ and $q$ are the independent and
dependent variables, respectively. The BPW-LB grid tools have several components. They all have
been ADIFOR-ed. As the seed matrix passes through each code, it gets accumulated with the gradient
information for that segment of the function.
ADIFOR and F-D Objective-Function (D/L) Sensitivity This chart shows the comparison of the ADIFOR sensitivities of the TCA baseline for the (D/L)
objective function with the finite-difference sensitivities. The finite-difference sensitivities have been
computed with a careful control of step size and a detailed step-size convergence study. The finite-
difference study was done on the CRAY C90 system in double precision. In contrast, the ADIFOR
sensitivities have been obtained on the SGI Origin 2000 system in single precision. The comparison
between the two is very good. When double precision was used for ADIFOR sensitivity calculation, the
minor discrepancies that are seen had completely disappeared.
ADIFOR and F-D Objective-Function (D/L) Sensitivity
High Speed Aerodynamics, Long Beach
TCA Wing/Body, FlexMesh, CFL3D, Euler, $M_{\infty}=2.4, \alpha=3.518^{\circ}, C_{D v}=41.35 \mathrm{cts}$.

ADIFOR Mode Reverse
The reverse mode ADIFOR uses adjoint principles to compute the gradients. The function and
gradient calculation steps are reversed and the derivatives are accumulated from the output to the input.
In this mode, the gradient of the output for unit value of the input is computed. For iterative solution
processes, the intermediate steps have to be stored to generate the adjoints. As a result, large storage
space is required for adjoint solutions of iterated solvers. If the number of dependent variables is much
smaller than the independent variables, the reverse mode is significantly faster. In CFL3D, if $p>25$, the
reverse mode is faster. For shape design problems with large number of design variables, the reverse
mode method of calculating the gradients is the best choice.
Reverse Mode ADIFOR

$$
\begin{aligned}
& \text { Uses Adjoint principles to compute gradients } \\
& \text { Accumulates derivatives from the output to the input } \\
& \text { Reverses the program flow and stores intermediate values } \\
& \text { that nonlinearly impact final result } \\
& \text { - Need huge storage space for large codes } \\
& \text { If } q \ll p \text {, reverse mode is significantly faster } \\
& \text { In CFL3D, } \\
& \text { - If } p<25 \text {, use forward mode }
\end{aligned}
$$


CPU Time for Forward and Reverse Mode ADIFOR Geom1, QGRID, and FlexMesh are the standard BPW-LB grid tools used in the design process.
Geom1 is the geometry kernel that modifies the baseline surface points based on a set of design
variables; QGRID uses surface point distribution and generates the surface grid; and FlexMesh perturbs
the baseline volume grid for a given surface grid modifications.
The chart below shows the CPU time comparison for computing the gradients using forward and
reverse mode ADIFOR. Note that the reverse mode ADIFOR CPU times are for 401 DVs. The
advantage of using the reverse mode ADIFOR is very clear.


|  | CPU Time (seconds) |  |
| :---: | :---: | :---: |
|  | Reverse-mode <br> (401 DV) | Forward-mode <br> (1 DV) |
| Geom1 | 17.2 | 6.1 |
| QGRID | 135 | 5.4 |
| FlexMesh | 423 | 143 |

ADJIFOR for Iterative Solvers
The mathematical theory of the application of reverse mode ADIFOR approach for iterated solvers
is described elsewhere in the literature. However, it is important to point out here that the adjoints for
the iterated solution schemes can use a simpler approach, denoted as ADJIFOR, of avoiding the
storage of all the forward mode iterations to evaluate the adjoint gradients typically used in reverse
mode ADIFOR. In the ADJIFOR appraoch, the converged solution to the problem at a given point is
first obtained and the logged iterated information for the last few cycles is used to construct the adjoint
variable derivatives. This derivative information is used for all the adjoint iterations to calculate the
objective function gradient information. Since the number of iterations for which the arithmetic
operations have to be logged is only the last three or four cycles, the information to be stored is
manageable. This ADJIFOR procedure for iterated solvers is very attractive for calculating the
gradients.
$\xlongequal{\text { ADJIFOR for Iterative Solvers }} \quad$ High Speed Aerodynamics, Long Beach


- Reverse-mode ADIFOR is too expensive for iterative solvers
- Careful algorithm reversal is required
- Generate ADJIFOR code that only affects the derivatives
- Use function solution at the start of the iterative process
- Iterate to ensure only the derivative convergence
- Mathematical similarity to discrete adjoint derivatives
- Contraction mapping theorem states, for $\left\|J_{s_{Q n}}\right\|_{1}<1$
- Need to store only 4 reverse-mode steps for $J_{s_{Q n}}$
- Derivative of objective function can be written as:
$V^{\prime}=L J_{F} J_{s_{Q n}} \ldots \ldots J_{s_{Q n}} J_{G} R$
ADIFOR/ADJIFOR Euler Sensitivities Comparison
This chart compares the objective function (D/L) sensitivities of the TCA W/B configuration. Both
ADIFOR and ADJIFOR-processed CFL3D code is used for calculating the sensitivities. The design
variabbe set chosen is a representative subset of the usual design variables used for shape optimization.
Excellent comparison between the forward mode and reverse mode adjoint sensitivities is observed for
all design variables except a few fuselage shape variables. This difference is attributed to the
inadequate convergence of the incremental iterative (II) version of the CFL3D forward mode ADIFOR
iterations. The sensitivities calculated using the black-box differentiated version of CFL3D showed
identical results as the ADJIFOR adjoint results.

Comparison of Sensitivity Calculation Methods The table below compares the various sensitivity calculation methods. Although the method of
divided difference is the simplest, its main problem is the need for knowing the proper step-size to
calculate the correct gradient values. The choice of improper step-size may yield a very erroneous
gradient value. ADIFOR does not suffer from this inaccuracy; but its computer resource requirement
increases in direct proportion with the number of independent variables. The adjoint methods have the
power of providing the derivatives with very little memory as well as execution time. However, the
determination of the adjoint variables for the problem as well as mathematically deriving the expressions
for Navier-Stokes equations is a very time consuming, labor intensive, and tedious task. ADJIFOR, on
the other hand, has the advantage of the adjoint-based approach so that the CPU time for derivative
calculations is independent of the number of design variables. Also, the adjoint code can be quickly
generated. However, the memory and disc storage requirements for ADJIFOR are large due to the data
logging requirement of iterative solvers. This is not a major concern for large parallel platforms.


[^21]The nonlinear shape optimization performs better with increasing number of design variables.
Since the ADIFOR-processed derivative codes require memory and CPU time proportional to the
number of design variables, there is a major bottleneck in using this sensitivity calculation approach for
large problems using sequential computers. However, distributed parallel processing using large
number of processors is a good source of large amount of memory and CPU. For example, the 256-
processor SGI Origin 2000 system at NAS has a total memory of 64 gigabytes and terabytes of disc
storage. Hence using parallel systems for ADIFOR, the large memory requirement is not a major
concern.

[^22]- Typically, nonlinear optimization shows more benefits with
larger number of design variables (DV)
- In ADIFOR, memory and CPU time increases with DVs
- Forward mode ADIFOR-ed CFL3D Euler with $0.5 \times 10^{6}$ grid
- CPU time for 1 DV gradient $2.5 \times$ CPU for function
- CPU time for 50 DVs gradients $\sim 63 \times$ CPU for function
- Memory increase with DV requires parallel processing
- Parallel processing enables fast turn-around time


The importance of utilizing parallel processing for ADIFOR sensitivity calculations was highlighted
 approach, each processor computes the grid sensitivities of a design variable. To compute the flow sensitivities, this approach is a very poor choice since it requires the function evaluation for each design variable. Instead, an alternate way is used for the forward mode ADIFOR sensitivity calculation. First, the parallel version of CFL3D is utilized so that the split volume grid can be mapped onto each processor. This code is black-box ADIFOR differentiated to obtain the forward mode ADIFOR version of the CFL3D code. Next, the back-box differentiated ADIFOR code is modified to include the incremental iterative (II) enhancement of Arthur Taylor. This served as the workhorse ADIFOR-ed CFL3D code for the design iterations during the TCA wing/body design. The seed matrix that contains the grid sensitivities of all the design variables is also split along the same way as the original grid is split and fed to the ADIFOR-processed CFL3D code. Since the grid is split, the sensitivity calculations of a group of design variables can fit into each processor. Several such
 processors. This method enables a fast turn-around time of the gradient computation. For the ADJIFOR adjoint gradient computation, the grid is split and mapped on to a set of
processors similar to function evaluation. However, the increased memory requirement for ADJIFOR
requires the grid to be split further to accommodate the increase in the memory requirement of the
ADJIFOR-processed code compared to the function code. Since the gradients are computed based on
the last four iterations of the converged function solution, large log files are stored. However, since log
files are written and read by each processor in parallel, the per-processor load is reasonable. files are written and read by each processor in parallel, the per-processor load is reasonable.
$\xlongequal{\text { Parallel Flow Sensitivities }} \begin{gathered}\text { High Speed Aerodynamics, Long Beach }\end{gathered}$

ADIFOR/ADJIFOR Euler Design Resource Needs
The resource requirements of the ADIFOR and ADJIFOR-processed CFL3D sensitivity code for a
$0.48 \times 10^{6}$-point Euler grid is shown in this chart. It can be seen that the total memory requirement of the
ADIFOR sensitivity calculation is nearly 40 gigabytes using single precision. Also, the CPU time for
obtaining the Euler flow sensitivities of the 401 design variables is nearly 450 single-processor hours on
the SGI Origin 2000 system. This is significantly large compared to the grid sensitivity calculation.
However, making use of large pool of processors, the wall-time to obtain the gradients can be reduced.

 for ADJIFOR is an order-of-magnitude smaller. Second, the CPU hours required for ADJIFOR gradients is nearly 25 times lower than that is required for ADIFOR. The wall time for adjoint gradient calculations performed for this grid using a 24 -processor Origin 2000 system is only 0.6 hour. This is a major advantage of using the ADJIFOR approach. Although the disc space requirement for logging the
intermediate steps during the adjoint calculation in ADJIFOR is nearly 16 gigabytes, it is not a concern.
ADIFOR/ADJIFOR Euler Design Resource Needs

- ADIFOR and ADJIFOR-processed CFL3D Euler code
- 401 DV TCA W/B Euler design with ~480,000 grid points
- DVs mapped onto PEs for parallel grid sensitivities
- Split grid mapped onto PEs for parallel flow sensitivities

Shape Optimization Using ADJIFOR Sensitivities
to determine the optimum direction.

AEROSHOP Framework
AEROSHOP is a Unix script-driven optimization environment. The heart of the AEROSHOP
optimizer environment is the NPSOL optimizer. The other segments of AEROSHOP are the gradient
generation and performance evaluation. The gradient-generation module built inside AEROSHOP is
based on finite difference. However, interfaces are provided to replace this with the customer-preferred
gradient modules such as ADIFOR or ADJIFOR and hand-differentiated adjoint.



## ADJIFOR-Based N-S Design

The next step in the grand scale of activity is to extend the ADJIFOR methodology to perform
viscous optimization. At this stage, the advantage of using ADJIFOR approach in contrast to the hand-
differentiated method of calculating the gradients is evident. The ADJIFOR adjoint code for Navier-
Stokes sensitivities was obtained in a quick and direct way by simply passing the CFL3D Navier-Stokes
code through the ADJIFOR preprocessor instead of the CFL3D Euler code used for the Euler
sensitivities.
To perform the Navier-Stokes wing/body design, the TCA baseline Euler wing/body design grid
was reclustered for the flight Re $=212 x 10^{6}$ with $y^{+}=5$ near the wall. Care was taken to ensure that the
viscous drag results of this design grid matches closely with the TCA baseline Navier-Stokes analysis
grid values. The resulting Navier-Stokes design grid had $0.89 \times 10^{6}$ points. The same set of design
variables and constraints used for the earlier Euler design were chosen for the Navier-Stokes design as
well.
Since this is the first time that the ADJIFOR technology is used to perform Navier-Stokes design, a
systematic study is warranted at the early stages to address the issues related to the iterative adjoint
procedure to calculate the derivatives. Also, the number of adjoint iterations required for the Navier-
Stokes gradient convergence should be established, prior to the beginning of design cycles. To this
effect, many adjoint gradient runs were made. From these evaluations, it was concluded that a large
number of adjoint iterations, comparable to the number of iterations required for function convergence, is
also required for gradient convergence. Another observation was that although accurate sensitivities are
obtained with single precision in ADJIFOR, due to the very tight spacing of the grid close to the wall,
double-precision arithmetic is required.

ADJIFOR-Based N-S Design



N-S Design Resource Requirements

$$
\begin{aligned}
& \text { The large size grid combined with the requirement of double precision arithmetic pose a burden on } \\
& \text { the computer memory requirement for Navier-Stokes ADJIFOR sensitivities. In addition, the hundreds of } \\
& \text { adjoint iterations needed for gradient convergence implies large CPU time for sensitivity calculations } \\
& \text { during the optimization. However, on parallel systems this is not a major concern. Using } 72 \text { processors } \\
& \text { of the SGI Origin parallel system, it is possible to obtain the converged adjoint gradients for the } \\
& \text { wing/body design grid in less than } 2.3 \text { hours of wall time. The ratio of CPU time for adjoint to analysis } \\
& \text { for the Navier-Stokes ADJIFOR sensitivities remained to be } 20 \text { to } 25 \text { as in the case of Euler ADJIFOR } \\
& \text { sensitivities. Indeed, very large log files ( } \sim 57 \text { gigabytes) are created during the iterative adjoint } \\
& \text { calculation process. It is anticipated that the available CPU and memory resources on the NAS Origin } \\
& 2000 \text { parallel system can easily handle up to } 3 \times 10^{6} \text { grid points. It is expected that the full configuration } \\
& \text { viscous multipoint design grid could fit within the available resources at NAS. }
\end{aligned}
$$

Requirements
High Speed Aerodynamics, Long Beach

N-S Design Resource

Results
computer resources.
Results

The design iteration history for various aerodynamic and optimization parameters for the TCA wing/body design with 401 design variables and 55 constraints are shown below. The chosen objective function for the minimization problem is drag/lift. The $\mathrm{G}_{\text {vio }}$ shown is the ratio of the constraint violation with over the constraint value. The first 20 design iterations were performed without any design parameter modification. As can be seen, the design progressed smoothly. At that time, these iterations were performed using the forward mode ADIFOR sensitivities that required more than a couple of month's time since the ADJIFOR adjoint technology was not available. It is interesting to see that switching to the ADJIFOR sensitivities, these steps can be completed within a couple of days. The ADJIFOR-processed CFL3D adjoint code became available when this design was completed. The last couple of design cycles were the only ones performed using the ADJIFOR Euler sensitivities. The comparison of the ADIFOR and ADJIFOR gradients were already shown in an earlier chart. Since the constraints were not being satisfied, the design variables were scaled after 20 iterations which = TCA $\stackrel{0}{\square}$ ation.
Design with ADIFOR-Based Euler Sensitivities

Design with ADIFOR-Based Euler Sensitivities
The variations of aerodynamic quantities with respect to the design iterations are shown in the
typical lift curve and drag polar form. The lift constraint is also shown as the dotted line.

ADIFOR/ADJIFOR Optimized TCA W/B Configuration The ADIFOR/ADIFOR optimized wing/body configuration is shown in this perspective view. A
significant spanwise waviness is seen that was absent in the TCA baseline linear design. A similar
waviness has also been obtained in other similar nonlinear Euler designs. Another observation to be
made is the nose droop. This was also absent in the baseline TCA configuration.

ADJIFOR-Processed CFL3D Navier-Stokes (B-L) Sensitivities

> The comparison of the sensitivities of the drag/lift objective function for the 401 design variable problem using the ADJIFOR-processed CFL3D Euler and Navier-Stokes adjoint codes are shown here. Important to observe are the large gradients of the fcwing design variable and the large fluctuations in the fuselage camber sensitivities for different number of Navier-Stokes adjoint iterations. The insets show that acceptable gradient convergence has been reached for around 600 adjoint iterations. As a result, all the wing/body design iterations based on the Navier-Stokes sensitivities use 700 adjoint iterations. Other design variables do not show a major change between Euler and Navier-Stokes
> sensitivities in this scale.
ADJIFOR-Processed CFL3D Navier-Stokes (B-L) Sensitivities (D/L)


ADJIFOR-Processed CFL3D Navier-Stokes (B-L) Sensitivities
This chart shows the close-up view of the wing design variables shown in the previous chart. It is
seen that, the Euler wing design variable sensitivities show a slightly different trend compared to the
Navier-Stokes sensitivities. Also, the number of Navier-Stokes adjoint iterations required for gradient
convergence of the wing design variables is not as large as the fuselage camber design variables. It is
noted that a mixed ADIFOR and ADJIFOR gradient calculation scheme is probably more economical.
Conclusions and Future Work
$\xlongequal{\text { Conclusions and Future Work }} \begin{aligned} & \text { High Speed Aerodynamics, Long Beach }\end{aligned}$

ADIFOR-processed codes provide accurate sensitivities
The W/B design technology demonstrator using CFL3D
Euler ADIFOR gradients improved TCA (L/D) by 5\%
Parallel processing enables design cycle time reduction
Parallel processing eliminates ADIFOR memory limitations
and ADJIFOR disc space issues
CFL3D Navier-Stokes ADJIFOR gradients have been
rapidly and successfully computed
Navier-Stokes designs for TCA wing/body and
wing/body/nacelle/diverter are underway

Technology Development for a Multipoint
Optimization Process for an HSCT
Parducci, James Hager, Eric Unger, Geojoe Kuruvila,
P. Sundaram, Peter Hartwich, Grant Martin,
Raul Mendoza, Alan Arslan, Shreekant Agrawal
The Boeing Company, Phantom Works - Long Beach
High Speed Research Airframe Review
Anaheim, California
February 8-11, 1999


## Introduction

Previous non-linear aerodynamic optimization work has focused mainly on the most important leg of the HSCT mission, the supersonic cruise. The benefits of this single-point optimization has not yet been fully realized and it is possible that with further work the goals of the HSR program may be met. However, consideration of the supersonic cruise performance may not be adequate to obtain a truly optimized HSCT since the airplane flies through diverse flowfields to accomplish each leg of its mission

Much of the technology development in 1998 has been to increase the fidelity of the analysis model, increase the number of degrees-of-freedom of the design, and reduce the cost of optimization. While these elements oppose each other, they are critical for a practical and successful multipoint optimization process.

This paper begins by assessing the progress of nonlinear aerodynamic shape optimization towards meeting program goals set for the Technology Concept Airplane (TCA). This section establishes that program goals have not been met and that steps towards reaching these goals may be made through a multipoint approach. The next section defines two approaches for multipoint optimization. The first involves a series of single-point optimizations where the OML is defined by supersonic cruise considerations, and flap schedules are determined in subsequent optimizations at off-design Mach numbers. The second approach involves concurrent assessments of the design performance at all critical Mach numbers.

The next section documents the optimization technology development in 1998. This includes enhancements to grid generation, configuration modeling, flow analysis, gradient calculations, and design variable definitions. Intermediate results that demonstrate these technologies conclude this paper.
outline

$$
\begin{aligned}
& \text { - Current Status of Aerodynamic Optimization } \\
& \text { - Multipoint Optimization Processes } \\
& \text { - Technology Developments in } 1998 \\
& \text { - Intermediate Results } \\
& \text { - Conclusions }
\end{aligned}
$$

Current Status - Supersonic
The objectives of this chart is to map out a feasible design space for the TCA at supersonic cruise and to determine where current variants of the TCA come to lie in this design space. The basic idea for designing this carpet plot is the assumption that the variation of drag with lift can be approximated by a parabola:
$C_{D}=C_{D_{0}}+K E \times C_{L}^{2}$
(1)
Straightforward differentiation of this formula gives the relationship between (L/D) max as a function of minimum drag and shape factor KE. Minimum drag is primarily driven by skin friction, by volume, and by nacelle installation. Primarily the effects of twist, camber and thickness distributions of the wing and of the fuselage (without twist, of course) determine the shape factor KE; the smaller KE, the better the performance of the aircraft. The estimates for the upper and lower bounds for these two parameters are taken from Kulfan.' The maximum lift-to-drag ratio for the Concorde is provided to put current results in perspective.
The maximum lift-to-drag ratio for three configurations at supersonic cruise Mach number is indicated in the feasible design space. A linear optimization process produced the baseline TCA. The TCA Cycle 2 and the NCV variants of the baseline TCA are the result of nonlinear design processes. Differences in the design processes and in the sets of constraints led to the different performance levels for the two nonlinear design variants of the TCA.
The three data points indicate that current designs emphasize reducing the shape factor, KE. Furthermore, assuming that Kulfan's estimates for the upper and lower bounds for minimum drag and shape factor KE are correct, it appears that the current design processes have come close to exhausting their capacity for further performance improvements through primarily optimizing thickness distributions, twist, and camber. On the other hand, minimum drag has suffered as a consequence of the optimization processes and is in danger of dropping below acceptable limits for $C_{D, 0}$. These results suggest that future improvements to design optimization methods should be aimed at reducing installation drag and exploiting any beneficial thrust and trim effects.
Current Status - Supersonic
High Speed Aerodynamics, Long Beach
$M_{\infty}=2.4, A R=2$

$$
i
$$



$$
\frac{0.5}{\sqrt{K E \times C D_{0}}}
$$

$(L / D)_{\text {max }}$

## Subsonic <br> Current Status

 This chart makes an attempt at quantifying the impact of flaps on the aerodynamic performance of a TCA wing/body configuration at transonic cruise Mach numbers. To this end, variations of a metric called suction parameter are plotted against lift coefficient. The six different curves pertain to a hypothetical case of minimum drag due to elliptic loading, a technical projection ( $=95$ percent of ideal suction), the TCA without deployed transonic flaps, the TCA with the leading- and trailing-edge flaps set to achieve the performance maxima as determined in computational and experimental parametric studies, and the TCA with all-wing flaps (inboard and outboard) set for maximum performance as determined by a nonlinear flap schedule optimization method. The suction parameter is defined as $\left.-\left.C_{o}\right|_{\text {elliptic }}\right)$
is the minimum drag due to elliptic load distribution, and

$$
\begin{aligned}
& +C_{L}^{2} /\left(C_{L} \tan ^{-1} \alpha\right)=C_{D_{0}}+C_{L}^{2} / C_{L_{\alpha}} \\
& \text { dicated in the chart, achieving maximum performance (i.e., } S=1 \text { ) requires the } \\
& C_{D_{\text {actual }}}=C_{D_{\text {elliptic }}}
\end{aligned}
$$

The TCA without deployed flaps turns in the worst performance. While the flap settings as determined in the parametric studies yield better peak performances for $0.04<\mathrm{C}_{\mathrm{L}}<0.16$, the optimized flap schedule surpasses the performance of all alternate flap settings at the transonic cruise condition ( $\mathrm{M}_{\infty}=0.9, \mathrm{C}_{\mathrm{L}}=0.18$ ). As a matter of fact, the optimized flap schedule lets the TCA perform at 93 percent of the technical projection, which, in turn, maintains that 95 percent of the ideal suction can be realized. This means that while there might be still further improvement achieved through refined nonlinear optimization of the transonic flap schedule, the current optimization results are already quite close to the practical limit.
For completeness, let it be mentioned that the estimation of the suction parameter based on experimental values is rather sensitive to uncertainties in determining minimum drag. This explains why the peak performance of some of the TCA wing/body configurations with transonic flaps exceeds the projected value and in one case even appears to match ideal performance; the experimental minimum

[^23]Current Status - Subsonic
Current Status - Subsonic
Current Status - Subsonic


## Process

 The multipoint optimization process shown below is an extension of the single-point methodologies used to perform the supersonic cruise optimization and the transonic flap optimization. The process can be thought to consist of 4 steps. The first two are start-up procedures; the third consists of optimization with $A E R O S H O P$, and the fourth is an analysis of the results produced in step three. The first step involves careful consideration of the design objectives and the available resources. It involves defining design variables, figure of merit (i.e., the objective function). Other considerations involved in step 1 include a decision of the preferred method for flow analysis and sensitivities, and the preferred computing platform.
 step 1 . These include generating the baseline grid, defining grid perturbation procedures, compiling computer codes, and generating baseline data (skin friction, spillage, etc.).
The third step is the heart of the optimization problem; it is the execution of the AEROdynamic SHape OPtimization system (AEROSHOP). The system solves general optimization problems. Within the context of multipoint optimizations, AEROSHOP executes surface perturbation, grid generation, and CFD analysis codes for each point of the design conditions. The data is passed to a design
 design is produced at the conclusion of this step. Optimizations may take advantage of shortcomings in the evaluation model or the
 plots of the optimization history, a refined CFD analysis, and communication to experts in other disciplines. Shortcomings are often addressed and the process is repeated.

Multipoint Design Process
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Multipoint Strategies
Two strategies are currently implemented for multipoint design at Phantom Works. In the first, the optimization is performed in
several steps. The first step involved optimization for supersonic cruise. Here, the OML is redefined in wing twist, thickness, camber,
and shear, fuselage camber, area distribution and shape, and nacelle shape. The supersonic cruise optimization is followed by several
independent, single-point flap optimizations at conditions away from the supersonic cruise point. In the second approach, the OML
and flap settings are determined simultaneously for all relevant conditions with each condition having its own flap schedule. Flaps are
retracted at the supersonic cruise condition.
The assumption in approach 1 is that performance levels away from the $M_{\infty}=2.4$ are not heavily dependent on the geometric features
that resulted from the supersonic cruise point optimization. The assumption implies that whatever drag penalty might be incurred,
away from supersonic cruise, is small or can be regained with optimized flap deflections. The second approach removes the major
assumption made in the first approach. Here, the optimization is carried out by integrating input from all relevant flight conditions.
Conditions away from the supersonic cruise point can influence the wing and fuselage OML.
Multipoint Strategies
High Speed Aerodynamics, Long Beach


## Aeroshop Framework

Multipoint optimization can be performed within the framework of the AEROdynamic SHape OPtimization (AEROSHOP) system. AEROSHOP is the preferred environment for shape optimization with nonlinear aerodynamic considerations at Boeing - Long Beach. It features a modular structure ideal for testing new technologies and runs on many platforms including the Cray C-90, J90, T3E, the SGI Origin 2000, and clusters of Hewlett Packard Workstations. In a top-level viewpoint, the design problem definition is input to the AEROSHOP core executable. The optimizer within the core is MDOPT, a Boeing proprietary version of NPSOL. When the optimizer

 and the second returns the gradient of the evaluation with respect to each of the design variables. Here, an evaluation consists of the objective function and all constraints. At the end of a script's execution, control of the optimization is returned to the core. The scripts can be custom-built to incorporate well established, heritage technologies or new, state-of-the-art technologies. The modularity of the AEROSHOP scripts also allows for codes to take advantage of parallel environments.
The dependence of the gradient script to the evaluation script can be strong or weak. In a finite-difference approach, the connection between the two scripts is strong as the gradient script will execute the evaluation script many times to generate the data needed for the difference. In an adjoint approach, the dependence is weaker. The gradient script will call the evaluation script to solve the flowfield equations, but then will execute a new branch to solve adjoint-related codes.
AEROSHOP Framework


Configuration Modeling
The success of an optimization problem is strongly dependent on the analysis model, the objective function, and the constraints. The optimizer may exploit weaknesses in any of these pieces of the multipoint model. These elements, as applied to the TCA, are the next topics addressed.
The technology developed at Phantom Works is focused on solving a multipoint optimization problem of an HSCT similar to the layout of the Technology Concept Airplane (TCA). The major components under consideration are the wing (W), body (B), nacelles
 been generated for full configuration optimization in addition to certain configurations containing subsets of the above mentioned components. These sub-configurations include W/B, W/B/F, W/B/N/D, W/B/N/D/C, W/B/N/D/E, W/B/N/D/F, W/B/N/D/E/F, and W/B/N/D/C/F. When nacelles and diverters are included in the configuration, it is not possible to deflect the inboard and middle trailing-edge flaps. When the W/B/F configuration is considered, additional details are provided to model the gap between the wing and the flaps at the trailing edge. Configurations may differ among evaluation points. This allows flaps to be deflected at low speed, transonic, and low supersonic conditions, while being retracted at the supersonic cruise condition.


* Does not include inboard \& middle trailing-edge flaps


## Function

 The objective function is the heart of any optimization problem. An improperly constructed objective function will inevitably lead to a

 that seems to indicate that the cost of an aircraft is proportional to its take-off gross weight (TOGW). In a single discipline optimization, modeling the TOGW is impossible since all the elements contributing to its evaluation are simply not available. In aerodynamic shape optimization, the best that can be done is to assume the TOGW is proportional to the fuel consumption and thus proportional to the thrust. It is on this premise that we base the objective function on a weighted sum of the thrust evaluated at key points along the mission profile,

## $f=\sum \omega_{i} C_{T_{i}}$

$f=\sum_{i} \omega_{i} C_{T_{i}}$
The weights associated with each element of the mission and which elements to include in the evaluation of (1) is an area of little understanding. It seems natural to exploit the Technology Integration team's analysis of the TCA mission profile that suggests the dominant legs of the mission are the supersonic cruise ( $\mathrm{M}_{\infty}=2.4, \mathrm{C}_{\mathrm{L}}=0.1$ ), the transonic cruise ( $\mathrm{M}_{\infty}=0.9, \mathrm{C}_{\mathrm{L}}=0.18$ ), and the supersonic acceleration ( $\mathrm{M}_{\infty}=1.1, \mathrm{C}_{\mathrm{L}}=0.15$ ). The weights associated with these segments are $1,0.3$, and 0.15 , respectively.
The thrust is computed according to the free-body diagram shown in the figure below. The net thrust is comprised of the engine thrust, $\mathrm{C}_{\mathrm{T}}$, the nacelle-inlet ram drag, $\mathrm{C}_{\text {Dram }}$, and aerodynamic drag, $\mathrm{C}_{\text {Daero }}$. The engine thrust acts at and along the centroid of the CA reference-nacelle nozzle, and the magnitude is the same for all nacelles. (The CA reference nacelle is the one in the IGES file that the CA community uses. It does not have an inlet spike or true nozzle.) The engine thrust is calculated to balance the forces in the drag direction. The ram-drag acts at the centroid of the CA reference-nacelle inlet in the drag direction, and is a constant. The components of the aerodynamic drag include pressure, viscous, spillage, and excressence drag and are assumed to act at the aerodynamic center. Engine thrust, ram drag, and aerodynamic drag are included when evaluating trim conditions
The Objective Function
The objective function is a weighted sum of the
aerodynamic performance at each design point $P=\omega_{1} P_{1}+\omega_{2} P_{2}+\ldots+\omega_{n} P_{n}$
where $\omega_{i}$ are the weighting factors and $P_{i}$ is typically
$C_{D}, D / L$, or $C_{T}$

- Thrust is computed to balance drag

Constraints
The constraint package used during the single-point optimization for Cycle 2 has been relatively unmodified. Constraints that may be
enforced for a multipoint design include those indicated in the figure below, though it is unlikely that inlet flow quality will be
exercised. Geometric constraints with a dependence on the flow condition, such as the in-flight deck incidence, will be enforced at the
supersonic cruise condition, and may be relaxed at other conditions.
Constraints
Speed Aerodynamics, Long Beach
Lift Conditions
Pitching Moments
Inlet Flow
Quality
Fuel Volume
Design Variables
Most design variables are formulated as shape function perturbations to an underlying geometry. Shape functions can allow for a smooth progression from the initial geometry to an optimized shape, though when many shape functions are used in a single optimization it often becomes necessary for a post-optimization smoothing procedure. Shape functions have an advantage over NURBS-based design by reducing the number of design variables required to define an optimized shape. Formulating the design variables as perturbations assumes that the initial design is good and the optimum shape is close. Perturbations offer the advantage that large changes to the design are not possible. This avoids wild designs that are often undesirable, difficult to grid, and hard to analyze. Design variables which are applied as absolute values (i.e., are not perturbations) include the angle-of-attack, flap deflections, and canard and empennage incidences.
Changes to the configuration can occur on the wing, fuselage, and nacelles. The wing can change in twist, camber, thickness, shear, and flap deflection. The fuselage can change in camber, cross-sectional area, and cross-sectional shape.

Nacelle/Diverter Shaping Capability Design variables applied to the nacelles preserve the internal cross-sectional area, but allow the outer mold line to change in thickness and camber. The inlet, nozzle or entire nacelle may change in pitch and yaw. Variations to the empennage are limited to pitch. The canard can change in pitch and dihedral. Changes in horizontal and vertical placement of the canard and empennage are restricted due to its adverse impact on other disciplines. Shown below are the programmed degrees-of-freedom for changes to the nacelle shape.
Nacelle/Diverter Shaping Capability
Shearing Capability $\begin{aligned} & \text { High Speed Aerodynamics, Long Beach } \\ & \text { Outer Surface Design }\end{aligned}$



$\alpha_{i}=$ Shear in pitch variable for segment $i$
$\beta_{i}=$ Shear in yaw variable for segment $i$
* Complete nacelle/diverter yaw variable
W/B/N/D Geometry/Grid Generation
A top-level diagram of the geometry and grid generation in AEROSHOP is shown below. Presented in its basic form, the geometry
generation module takes baseline geometry and input design variables and then perturbs the W/B surface. Next, the nacelles are rolled
and pitched to follow the lower surface of the wing (assuming a W/B/N/D configuration is being optimized). The final geometry step
is the generation of a new diverter between the new wing and the translated and rotated nacelles.
The grid generation module uses the current defining geometry and creates a CFD surface grid on the W/B using the QGRID code.
This surface grid is then passed to either CSCMDO or FlexMesh to create a multiblock W/B surface and to perturb the multiblock
volume grid. Again, if only a W/B configuration is being optimized, the volume is single block, and the nacelle/diverter work is
unnecessary.

Canard and Empennage Integration
A top-level diagram of the system that integrates/perturbs the canard and empennage grid blocks to an existing W/B/N/D
configuration is shown. The first step in this process is to integrate the new canard to the current configuration. The canard can be
modified in both pitch and dihedral and the system maintains the relative position of the canard on the fuselage. The new canard
surface grid, along with a baseline W/B/N/D/C grid and the current configuration's W/B/N/D grid are given to CSCMDO to produce a
new W/B/N/D/C surface and volume grids.
The next step in this process is similar to the canard module with the integration of the vertical and horizontal tail surfaces to the fuselege aftbody. Once again, the system maintains the relative position of the tails to the fuselage. These new empennage surfaces along with the baseline $\mathrm{W} / \mathrm{B} / \mathrm{N} / \mathrm{D} / \mathrm{C} / \mathrm{E}$ grid and the previously created current $\mathrm{W} / \mathrm{B} / \mathrm{N} / \mathrm{D} / \mathrm{C}$ grid are again given to CSCMDO to produce a new W/B/N/D/C/E surface and volume grids for analysis.
Full Configuration Grid
The figure below shows the CFD surface grid for the W/B/N/D/C/E configuration with every-other point removed for clearity. Also shown are detailed views for the canard and empennage regions.
Ful| Configuration Design Grid

Every-other point shown

Flow Analysis The AEROSHOP system is connected to two different flow solvers at this time. The first of these is NASA Langley's well known CFL3D. This code was selected due to its ability to work on both point-matched and patched multiblock grids. CFL3D will work on both serial and parallel platforms providing AEROSHOP with even more flexibility. The code has been differentiated with ADIFOR in both the forward and reverse modes to provide flow sensitivities to AEROSHOP (in addition to a finite-difference option). In addition to these design-oriented enhancements, CFL3D now has a wall-function method to provide Navier-Stokes results on coarser meshes more suitable to the design environment.
The second solver that has been connected to AEROSHOP is the TLNS3D code, also from NASA Langley. Implementation of this code was delayed due to its previous inability to work with patched-multiblock grids (a handicap that has recently been lifted). TLNS3D also functions on serial and parallel platforms, thus offering the same flexibilty as CFL3D. The code offers very fast and generally more robust solutions than CFL3D, but it has not been processed through ADIFOR for sensitivities. There is however, a very efficient hand-differenced adjoint code developed for TLNS3D that has been used extensively in AEROSHOP for Euler-based optimzation.
Flow Analysis for Optimization
$\longrightarrow$ High Speed Aerodynamics, Long Beach

Gradient Calculations
Automatic Differentiation
This chart shows a three-step process to obtain analytic gradients using automatic differentiated FORTRAN codes. The design variables are passed through surface and volume grid generation tools to produce a grid suitable for analysis. The differentiated CFL3D code operates on the grid to produce a flow solution and an adjoint solution. The next step involves generating grid sensitivities. These establish how the grids change with perturbations in the design variables. The flow, adjoint, and grid sensitivitives are passed along to the next step that produces the sensitivities of aerodynamic quantities with respect to design variables.
TLNS3D-Adjoint Gradients
Beach

No


TLNS3D-Adjoint
The adjoint method provides a means to compute sensitivities in an efficient and accurate manner. Details of the method can be found in many papers such as Reuther et al. ${ }^{2}$ and Kuruvila et al. ${ }^{3}$ In essence, the method requires two solutions for a gradient of an
 adjoint solutions are required for gradients of lift and drag. In addition to the flow and adjoint solutions, the TLNS3D-Adjoint formulation requires gradients of the grid with respect to the design variables. This can be computed by finite-difference or analytically. Finite-difference grid sensitivities are easy to compute, but can be inaccurate if a poor step size is chosen. Analytic sensitivities are not dependent on the step size, but can be difficult to obtain. The ADIFOR tool is ideal for obtaining analytic grid sensitivities.

Depicted below is the implementation of the TLNS3D-Adjoint system with the AEROSHOP framework for a single-point design. AEROSHOP relinquishes control to a script that drives TLNS3D-Adjoint when sensitivities are required. This script first generates a baseline grid. Next it computes a flow solution using TLNS3D. In a parallel environment, TLNS3D can run on a user-defined number of processors, typically between 10 and 30 . The next step involves obtaining solutions for the adjoint equations. This step is repeated for each aerodynamic quantity required (i.e., $\mathrm{C}_{\mathrm{L}}, \mathrm{C}_{\mathrm{D}}$, and $\mathrm{C}_{\mathrm{M}}$ ). The adjoint problem is similar in size to the flow problem and thus can be run on the same number of processors as was used to compute the flow solution. The next step involves obtaining the grid sensitivities. In the "production" version of AEROSHOP, the grid generator is applied to each perturbation of the design variables for a finite-difference estimation of the grid sensitivities. On a serial machine, this process is done sequentially and is the time dominating step in the optimization for problems involving large numbers of design variables. On a parallel machine, the grid sensitivities are done in blocks of $m$ number of design variables. For efficient use of the computer, the number of processors used here should be close in value to the number used to compute the adjoint and flow solutions. The final step is the calculation of the sensitivities. For multipoint optimization, each step is repeated over the number of evaluation points.
Automatic Differentiation
$\ldots$ High Speed Aerodynamics, Long Beach

Volume Perturbation

Intermediate Results

## Optimization with Analytic Derivatives

A head-to-head comparison of the analytic derivative methods, CFL3D-ADIFOR and TLNS3D-Adjoint, was performed using a W/B optimization of the TCA. The objective of the exercise was to compare the methods in accuracy, usability, and potential for growth. The optimization problem was set up identically in both cases; each started from the same configuration (TCA), used 401 design variables, and was formulated to minimize $\mathrm{D} / \mathrm{L}$ at $\mathrm{M}_{\infty}=2.4$ for a given lift condition. The problem was constrained with the set of constraints used during Cycle 2 optimization. Sensitivities of the constraints, in both cases, were obtained using the ADIFORed processed constraint package.
Differences in the optimization were limited to the sensitivity method for aerodynamic quantities and the flow solver of performance evaluations. The ADIFOR optimization used CFL3D and the hand-differentiated adjoint used TLNS3D, thus making the analysis consistent with the sensitivity method.
The final results of the optimization are plotted on the polar below. The accuracy of the gradients produced by the methods seems to be comparable as both optimizations produced substantial improvements in the design. The CFL3D-ADIFOR procedure produced sensitivities in less wall clock time since it was able to make use of the parallel environment on the NAS Origin 2000. At the time of the optimizations, TLNS3D-Adjoint was restricted to serial computers and significant time was required to obtain grid sensitivities. The TLNS3D-Adjoint optimization has the advantage of running on a single platform and operated in a relatively hands-off manner. The line searches for the CFL3D-ADIFOR optimization was performed on the Cray C-90. This process required file transfers and thus was prone to human errors. Work is underway to eliminate the dependency on the C-90. Both methods have proven to be viable candidates for future optimizations.
Current State of TCA W/B Design


- Two optimizations were performed
using analytic derivatives of the
objective function and constraints
- Case 1 computed flow sensitivities
with ADIFORed-CFL3D
- Case 2 computed flow sensitivites
with TLNS3D-Adjoint
- Constraint sensitivities were
computed with ADIFOR-based code
- Optimizations used 401 design
variables and 55 constraints
Intermediate Results

The removal of the restriction to configurations griddable with a single-block grid provided the first major benefit to the $A E R O S H O P$ optimization process. This technology was demonstrated in a multipoint environment to simultaneously reduce the supersonic and transonic cruise point drag. At the supersonic condition, the $\mathrm{W} / \mathrm{B} / \mathrm{N} / \mathrm{D}$ configuration was modeled while the transonic configuration was limited to $\mathrm{W} / \mathrm{B} / \mathrm{F}$. The exercise produced a total improvement of 5.7 dragcounts at the supersonic condition as predicted by Euler analyses.
The development of the TLNS3D-Adjoint system benefited the optimization process in two ways. First, the flow sensitivities were computed with greater accuracy and second, they are computed with greater speed. The first implementation of the TLNS3D-Adjoint system was limited to optimizations involving single-block grids and thus the test case was restricted to the W/B configuration. Here, the latest geometry (from the $W / B / N / D$ multipoint optimization) was used as the initial geometry with the nacelle removed. The nacelles were reinstalled after the optimization. The optimization effort further reduced the W/B/N/D Euler drag by 0.9 counts. The reduction can be attributed to greater accuracy in the gradients and the larger degree-of-freedom resulting from the hundreds of design variables used in the optimization.
The final point shown on the figure documents the benefit of two technologies; the nacelle kinking capability, and the extension of the TLNS3D-Adjoint system to configurations gridded with multiple blocks. The optimization exercised the nacelle shaping optimization that included inlet, nozzle, and the entire nacelle alignment in pitch and yaw. Optimization including the nacelle shape was not available at the time of this optimization. The technologies resulted in a 7.3 count improvement in Euler drag over the TCA Baseline configuration.


Drag Coefficient,
- Phase 1: W/B with nacelle effects;
Intermediate Results

[^24]Baseline TCA W/B/N/D/E with Thrust, CFL3D Euler, $M_{\infty}=2$
$\operatorname{Re}_{\mathrm{c}}=212 \times 10^{6}\left(C_{D v}=45.63 \mathrm{cts}.\right), C_{L,}$ w/Thrust $=0.1, C_{D, R a m}=267.1 \mathrm{cts}$., $Z_{C G}$


Intermediate Results Trim Conditions, $\theta_{T}$ variation
A parametric study was performed to evaluate the trim drag benefit of the TCA W/B/N/D/E configuration with thrust as a function of
the thrust-vector pitch angle, $\theta_{\mathrm{T}}$. For this study, $\theta_{\mathrm{T}}$ is the same for the inboard and outboard engines and is independent of the nozzle
angle. The baseline thrust-vector pitch angles are $\theta_{\mathrm{T}}$, inboard $=-3.11^{\circ}$ and $\theta_{\mathrm{T}}$, outboard $=-2.15^{\circ}$ for an average $\theta_{\mathrm{T}}=-2.9^{\circ}$. Pitching
the thrust-vector up reduces $\theta_{\mathrm{T}}$ and $\mathrm{i}_{\mathrm{H}}$, and increases $\mathrm{C}_{\text {Drimim. }}$ The minimum-thrust condition is near $\theta_{\mathrm{T}}=0.8^{\circ}$.
Sequential Multipoint Optimization
The sequential procedure for multipoint design (outlined in an earlier section) has been applied to four supersonic cruise point designs.
Three of these designs correspond to the Cycle 2 configurations produced by NASA Ames, BCAG, and Phantom Works. The fourth is
a configuration representing the latest design from Phantom Works (Cycle 3). In this procedure, the OML is frozen with the exception
of leading- and trailing-edge flaps. The optimization is done as a single-point design at $\mathrm{M}_{\infty}=0.6,0.9$ and 1.1 . The results of the
optimizations are shown below as drag decrements with respect to the TCA baseline configuration with all flaps retracted. This
exercise demonstrates that off-design Mach number performance is not seriously compromised as a result of supersonic cruise point
design. This method, however, does not measure the benefits of a simultaneous multipoint design approach.
Sequential Multipoint Results
$\square \longrightarrow$ High Speed Aerodynamics, Long Beach

## Drag Reduction Relative to TCA Baseline with no Flaps Results Based on CFL3D Euler Analysies on the Design Grid



Intermediate Results One area of potential benefit to the overall performance of the TCA is the optimization of the leading edge. During take-off, climb, and transonic cruise, a blunt leading edge can benefit the performance by keeping the flow attached. At the supersonic cruise condition, a sharp leading edge is preferred to reduce wave drag. The purpose of this exercise is to optimize the leading edge for optimum performance.
In the first attempt, a single-point optimization was performed at the transonic cruise condition ( $\mathrm{M}_{\infty}=0.9, \mathrm{C}_{\mathrm{L}}=0.18$ ). Here the transonic performance could be improved by modifying the thickness and camber of the leading edge or by deflecting the leading-edge flaps. Trailing-edge flaps and the angle-of-attack were other variables in the design problem to assist in maintaining the lift constraint. The result of the optimization was an increase in thickness of the leading edge near the break, and a droop in leading edge over the inboard section of the wing. The part-span inboard flap was deflected to $15^{\circ}$. This led to a 2.5 count reduction in drag at the design condition, but a significant 1 count increase at the supersonic cruise condition.
The second attempt repeated the optimization in a multipoint setting. Here, an objective function based on the drag at the supersonic and transonic cruise condition was constructed,
$f=\left.C_{D}\right|_{M_{\infty}=2.4, C_{L}=0.09}+\left.0.3 C_{D}\right|_{M_{\infty}=0.9, C_{L}=0.18}$
This function, evaluated for the baseline TCA and the configuration obtained from the single-point optimization, is 86.6 and 86.9 , respectively. The function after multipoint optimization was reduced to 86.2 with a small reduction in supersonic cruise drag ( $\sim 0.1$ count) and $3 / 4$ count reduction in transonic cruise drag. The reduction in the objective function is small because the configuration changes occur only at the leading edge ( $10 \%$ of the chord) and because each configuration has optimized flaps at the transonic condition (flaps are retracted at $\mathrm{M}_{\infty}=2.4$ ). compromising the supersonic performance.




## Optimization <br> Leading-Edge

High Speed Aerodynamics, Long Beach
Drag Polars of TCA W/B Configurations Resulting from Single \& Multipoint Optimizations

Conclusions
This paper documents the need for further nonlinear aerodynamic shape optimization to meet the performance expectations of the
HSR program. The optimization may be done as a single point of multipoint process. Multipoint optimization is a more general
optimization and may produce a better overall design. Single-point optimization is cheaper, but does not offer any promises on
performance away from the design condition. Multipoint optimization is costly and thus motivated the development of certain
technologies to increase model fidelity and degrees-of-freedom while decreasing the cost of optimization.
The technologies developed at Phantom Works have focused on meeting these goals. Model fidelity has increased with the extention
of grid perturbation tools to multiblock grids. Increases in degrees-of-freedom and the reduction of optimization costs have been
accomplished with the development of TLNS3D-Adjoint and the portability of $A E R O S H O P$ to parallel environments. Sensitivity
accuracy and quick turn-around time have been improved with sensitivities obtained through ADIFOR.
These technologies have been demonstrated in several applications to deliver tangible improvements to the TCA configuration.
Conclusions
High Speed Aerodynamics, Long Beach
References

[^25]
# Accounting for Laminar Run \& Trip Drag in Supersonic Cruise Performance Testing 

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This work was done to support the the WT Database and WT Data Corrections programs under the Test Program and Techniques approach for the Configuration Aerodynamics Technology Development.


## Background:

One of the goals of the High Speed Research Program is to achieve new levels of precision and accuracy, particularly in supersonic performance testing.

The picture shown below is that of the $1.675 \%$ Baseline TCA Model 2 b . This is our baseline performance model which has been tested numerous times in the past couple of years.


## Background (continued):

As part of the goal to obtain increased precision and accuracy, it is important to reduce the statistical uncertainty in the wind tunnel data. This is achieved by taking between 3-6 repeat measurements of lift and drag over a range of angles-of-attack.

The plot on the left shows an example of a drag polar for the $1.675 \%$ Baseline TCA Model. The squares are the actual lift and drag measurements. The data are represented by fitting a least-squares polynomial curve through the points.

However, we are primarily interested in determining the drag at cruise lift conditions.

## HSR Annual Airframe Technical Review <br> Motivation

- Why do we care so much about drag?

1 drag count ( 0.0001 ) $\sim 6,000 \mathrm{lb}$. take-off weight

- Ultimate program goal is to measure drag to within $\pm 0.5$ count with $80 \%$ confidence


## Motivation:

Why do we care so much about this cruise drag?
The reason is that, at supersonic cruise, each drag count costs approximately 6000 lb in take-off weight, which is equivalent to about 30 passengers. Therefore, there is a significant economic incentive to both know what the drag is and to try to reduce it.

The ultimate goal of the Configuration Aerodynamics Test Team is to obtain drag to within $+/-0.5$ counts with $80 \%$ confidence at supersonic cruise.

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## Sub-scale Modeling Issues

- Cannot directly simulate flight conditions
- Flight $\mathrm{Re}=209$ million; boundary layer transition near LE
- Wind Tunnel $\operatorname{Re}=6.4$ million; transition is delayed
- Mixed laminar/turbulent skin friction is too low


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## Sub-scale Modeling Issues:

Part of the difficulty with wind tunnel testing is that we cannot directly simulate flight conditions.

For example, at flight, the Reynolds number for the TCA configuration is 209 million. At this Reynolds number, the boundary layer transition from laminar to turbulent flow is near the leading edge over most of the wing. On the other hand, a much lower Reynolds number is attained in the wind tunnel. The Reynolds number for the $1.675 \%$ TCA Model in the NASA Langley UPWT is equal to 6.4 million. For this low Reynolds number condition, the boundary layer transition is delayed over most of the wing. This is demonstrated in the flow visualization image below, which shows the extent of laminar flow on the model.

As a result of the mixed laminar/turbulent flow, the skin friction is lower than that for a flow which is fully-turbulent from the leading edge.


## Forcing Early Transition

- Standard approach is to add roughness near wing leading edge
- This does not completely solve the problem
- Transition may not move all the way forward, especially in supersonic flow
- Roughness can create additional drag
- Analysis of wind tunnel drag data requires accounting for laminar run \& trip drag


## Forcing Early Transition:

The standard experimental approach to solve this problem is to add roughness (e.g., grit, glass beads, epoxy dots) near the leading edge of the wing in order to induce early transition.

However, this does not always resolve the difficulty. First of all, the transition line may not move all the way forward to the trip location, especially in supersonic flow. Second, applying roughness to the model can create additional drag.

Therefore, it is important that the analysis of the wind tunnel data takes into account all of the laminar run which exists behind the leading edge and any additional trip drag that occurs.


## Not a New Problem

- Braslow (Langley, 1960s)
- Advocated extrapolating to zero trip height
- Daugherty (Ames, 1970s)
- Observed constant drag for a range of trip sizes
- Vaucheret (France, 1960s)
- Applied laminar run correction
- Corrected drag was constant for small trips

In all cases, prompt transition was obtained for sufficiently large roughness

## Not a New Problem:

The whole issue involved with tripping the flow and correcting the wind tunnel data is not a new problem in wind tunnel testing.

Braslow worked at NASA Langley in the 1960s to develop a method to estimate trip drag. He advocated using an extrapolation technique. This method extrapolates a curve through the drag data as a function of trip height in order to estimate the drag for zero trip height. This technique will be further examined in Slide 11.

Daugherty conducted a series of wind tunnel tests at NASA Ames in the 1970s on the XB-70 model. He observed constant drag over a range of trip sizes. Some of his results are presented on the next slide.

Vaucheret conducted tests on a delta wing/body model in France during the 1960s. Based on sublimation images, he observed some laminar flow on the model. Therefore, he first applied a laminar run correction to his drag data. Once he made the correction, he found that the drag was constant for small trips. Some data from his work are shown in Slide 10.

In all cases, it is believed that all of the researchers used airfoil sections with sharp leading edges which produced gentle pressure gradients. Therefore, they were able to obtain prompt transition for sufficiently large roughness sizes.

## Jim Daugherty -XB-70 Test



Daugherty, et al. (Ames)

- $\mathrm{M}=2.53, \mathrm{Re}=4\left(10^{6}\right) / \mathrm{ft}$, sharp LE sections
- Transition observed at or near the grit strip
- Drag "plateau" observed; zero trip drag for $k$ up to about 0.013 in

Jim Daugherty -XB-70 Test:

For further information, please refer to the following references:

James C. Daugherty, "Wind-Tunnel/Flight Correlation Study of Aerodynamic Characteristics of a Large Flexible Supersonic Cruise Airplane (XB-70-1), I-Wind-Tunnel Tests of a 0.03 -Scale Model at Mach Numbers from 0.6 to 2.53," NASA Technical Paper 1514, November 1979.

Henry H. Arnaiz, John B. Peterson, Jr., and James C. Daugherty, " WindTunnel/Flight Correlation Study of Aerodynamic Characteristics of a Large Flexible Supersonic Cruise Airplane (XB-70-1), II-A Comparison Between Characteristics Predicted from Wind-Tunnel Measurements and Those Measured in Flight," NASA Technical Paper 1516, 1979.


## Xavier Vaucheret



- $\mathrm{M}=1.65, \operatorname{Re}=9.8\left(10^{6}\right) / \mathrm{m}, \alpha=0^{\circ}$, symmetrical sections
- Trip consisted of a single row of glass beads
- Drag "plateau" found after correcting for laminar flow


## Xavier Vaucheret:

For more details of Vaucheret's work and results for Mach 2.34, please refer to the following reference:

Xavier Vauchert, "Artificial Initiation of the Supersonic Transition," Translation of "Declenchement Artificiel de la Transition en Supersonique," La Recherche Aerospatiale, No. 120, pp. 25-32, 1967, NASA TT F-11, 558.


## Drag vs Trip Height:

The Configuration Aerodynamics Test Team attempted to follow Braslow's approach in their initial analysis of the wind tunnel data. The first step is to plot the drag as a function of trip height for those values where prompt transition is believed to occur. This is shown in the plot on the left. Then, a curve is extrapolated through the drag data to zero trip height. The trip drag is computed as the difference between the measured drag and the extrapolated drag.

For this data, it is not clear which extrapolation technique is the correct one to use. If a quadratic extrapolation is made, the trip drag is equal to 3.3 counts. On the other hand, if a linear extrapolation is applied, then the trip drag becomes 5.8 counts. The difference between these two methods- 2.5 counts-is unacceptably large.

At this point, it was important to determine a way to resolve this uncertainty in the trip drag.


## Test \#1679:

In order to resolve the uncertainty in estimating trip drag, a more careful examination of the results from Test \#1679 was undertaken. This test was the first entry of the $1.675 \%$ Baseline TCA Model 2 b in the Unitary Pressure Wind Tunnel Test Section \#2 at NASA Langley Research Center. The other collaborators involved in this wind tunnel test were Mina Cappucio and Kevin Mejia.

It is important to note that the inboard section of this model has a blunt leading edge, whereas the outboard section has a sharp leading edge. The pressure distribution around the blunt leading edge is different from that around the sharp leading edge. The pressure distribution in the outboard section with the sharp leading edge is more similar to that obtained in the previous historical tests.

## W/B Trip Drag Study

- Trip height on wing varied from $\mathrm{k}=0.006^{\prime \prime}$ to 0.017 "
- Trips were epoxy disks with diameter 0.050", applied 0.6" from LE
- Each trip configuration was carefully applied and measured for uniformity
- Measure forces \& moments at $\mathrm{Ma}=2.4$
- Obtain series of sublimation images at cruise ( $\alpha=3.5^{\circ}$ ) until stable pattern emerges


## W/B Trip Drag Study:

In Test \#1679, a trip drag study was conducted on the wing/body configuration. The trip height was varied from 0.006 inches to 0.017 inches. The trips were epoxy disks with a constant diameter equal to 0.050 inches. The disks were applied at 0.6 inches from the wing leading edge. The picture in the upper right corner shows an example of the trip disks on the model near the leading-edge break.

Each trip configuration was carefully applied and measured in order to maintain uniformity. For each trip configuration, forces and moments were measured at the supersonic cruise Mach number, which is equal to 2.4 . In addition, a series of sublimation images at the cruise angle-of-attack were obtained until a stable pattern emerged in the sublimation material.


## Sublimation Process

- Model is positioned with wings vertical
- Cameras are mounted on each side of test section
- Upper \& lower surfaces of wing are sprayed with sublimation material
- Pictures are taken at 30 sec intervals
- Allows a complete record of the development of the transition front


## Sublimation Process:

To obtain the sublimation images, the model is rotated $90^{\circ}$ so that the wings are vertical.

Cameras are mounted on each side of the test section to capture the upper and lower surfaces of one wing at the same time.

The upper and lower surfaces of the wing are sprayed with the sublimation material, which is a mixture of fluorene in Genesolv.

The test section is closed and the tunnel is brought up to the supersonic cruise condition.

Then, pictures are take every 30 seconds until a stable pattern emerges and most of the material has sublimed.

This allows a complete record of the development of the transition location.


## Sublimation Reveals Transition:

This slide gives examples of sublimation images for the free transition configuration at four different times.

The transition location is apparent in these images because the material sublimes more rapidly in areas of turbulent flow.

After 5 minutes, the initial development of the transition location appears. As more time elapses, the transition location becomes more evident. The last two images demonstrate that the transition location remains stable with time.

## Locating Transition

Transition location was measured as a percentage of local chord from the digitized images


## Locating Transition:

The sublimation images are digitized and read into a software program like Adobe Illustrator to help locate the transition line on the model for each trip configuration.

First, the transition location is drawn on the wing, shown as a thick purple line in the picture above. Next, the green guidelines are used to measure the transition location as a percentage of the local chord for various spanwise stations.

This procedure is followed for all of the different trip height configurations.


## Transition Location-Upper:

This viewgraph gives a summary of the transition location for all of the dot heights tested on the upper surface. The thick black line represents the location of the free transition line. Note that transition is significantly delayed over most of the wing without any trip disks applied to the wing.

In the outboard section, which has sharp leading-edge airfoil sections, there is an abrupt forward movement of the transition location with any application of the epoxy disks, even for the smallest disk height. For disk heights of 0.010 inches and greater, the transition moves forward to the disk location.

On the other hand, the inboard section behaves differently. The blunt leading edge makes it difficult to trip the flow and transition is delayed for all disk heights. In fact, for the smallest disk height, the transition location is not changed from that for free transition, except in the region near the leadingedge break in the planform. Prompt transition is not attained for any of the trip configurations, even with the largest trip height of 0.017 inches.

The dissimilar behavior of the transition line in the inboard and outboard sections may be due to differences in the pressure gradient and boundary layer development between the two sections.


## Transition Location--Lower:

The summary of the transition locations for the different trip heights on the lower surface of the wing is shown here.

The outboard section of the lower surface behaves similarly to that of the upper surface.

However, the transition location in the inboard section displays different behavior than on the upper surface. First of all, the free transition location is closer to the leading edge than observed on the upper surface. Furthermore, the application of most disk heights moves the transition forward to the trip location.

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## Methodology

- Laminar run correction
- Compressible flat-plate skin friction code developed
- Compute $\mathrm{C}_{\mathrm{F}}$ for both mixed laminar/turbulent and fully-turbulent flow from LE
- Difference between these is the laminar run correction, $\Delta \mathrm{C}_{\text {Dlam }}$

> Once corrected, the data provided guidance for extrapolating to zero trip height, resolving the uncertainty in $\Delta C_{D r r i p}$

## Methodology:

Once it became evident that a significant amount of laminar flow existed, particularly on the inboard, upper, wing surface, it was clear that the wind tunnel data needed to be corrected for laminar run before the data could be analyzed for trip drag.

In order to make this laminar run correction, co-author Robert Kennelly wrote a FORTRAN code that estimates the skin friction for a compressible flow over a flat plate. This code was then used to compute the skin friction for a mixed laminar/turbulent flow using the measured transition locations and for a flow that is fully-turbulent from the leading edge. The difference between these two results is the laminar run correction.

Once the drag data was corrected and plotted against trip height, the proper extrapolation of the data to zero trip height became apparent, which resolves the uncertainty in the trip drag.

## CDf: New Code, Old Method



- Flat plate, compressible, several laminar and turbulent $\mathrm{C}_{\mathrm{f}}$ methods
- Iteratively matches laminar \& turbulent momentum thicknesses at a specified x-location
- Drag areas ( $\mathrm{D} / \mathrm{q}$ ) are summed over all elements; normalized by Sref to yield $C_{D}$ for configuration


## CDf: New Code, Old Method:

This FORTRAN code has been run on a number of platforms and is available through Robert Kennelly at rakennelly @ mail.arc.nasa.gov.


## Drag vs. Trip Height:

The reader will now be led through a step-by-step procedure for correcting the wind-tunnel data.

First, the drag data for $\alpha=3.5^{\circ}$ is obtained by interpolating the 7th-order polynomial curves which represent the wind tunnel data. The drag data are plotted as a function of trip height for all values tested. The uncorrected drag data are shown as blue squares.

Next, the laminar run correction is added to the drag data for each disk height. The laminar run correction is largest for free transition and is equal to 5.8 counts. For a trip height of 0.012 inches, the laminar run correction is 1.7 counts. The corrected data are represented by the black circles.

The data, corrected for laminar run, fall into two regions that will be described further in the next slide.


## Drag vs. Trip Height:

First, a "drag plateau" is defined using data from the first two points-free transition and 0.006 inches. The plateau is only tentatively identified since drag data in between the two endpoints are not available from Test \#1679. However, some flattening of the drag curve is expected as zero trip height is approached. In addition, the existence of a drag plateau is consistent with Braslow's analysis of earlier subsonic data and other researchers' supersonic data, including those of Daugherty and Vaucheret.

The remaining data varies linearly with trip height. For trip heights that are approximately equal to the boundary layer height or greater, the drag is expected to be proportional to the disk frontal area. Since the diameter remains constant in this study and only the height is varying, a linear variation of the drag with trip height is appropriate.


## Final Drag vs. Trip Height:

The trip drag is computed by taking the difference between the linear fit of the data and the value of the drag plateau.

The trip drag is subtracted from each of the data points that lie above the plateau value. The trip drag for the disk height equal to 0.012 inches is 1.9 counts, which is close to the value of the laminar run correction for the same disk height.

The final corrected drag for all disk heights are represented by the solid black circles. All of the corrected drag data are now consistent with the value of the drag plateau.

Therefore, for this configuration at the cruise angle-of-attack, the plateau value is the fully-corrected cruise drag that the Configuration Aerodynamics Test Team has sought.


## Comparison with CFD Results:

In this slide, the final corrected wind tunnel data are compared to three different CFD solutions near the cruise angle-of-attack.

The corrected wind tunnel data are in excellent agreement with the lift and drag results from UPS (Scott Lawrence, NASA Ames Research Center) and CFL3D (Chih Fang Shieh, Boeing-Phantom Works, Long Beach). Both of these codes used the Baldwin-Lomax turbulence model.

On the other hand, the drag computed by OVERFLOW (Steve Chaney, Boeing Commercial Airplane Group, Seattle) is about two counts lower than the wind tunnel data. This is most likely due to the anomalous behavior of the Spalart-Allmaras turbulence model which exhibits a "laminar-flow effect" near the leading edge of the wing.

## HSR Annual Airframe Technical Review <br> Nacelle Trip Drag Study

- Constant wing trip height: 0.012 "
- Three nacelle trip heights (inside \& out) tested: $0.008^{\prime \prime}, 0.012 "$, and $0.014 "$
- Note: large number of trip disks on nacelles (49\% of wing total)
- Measure W/B/N/D forces \& moments
- Sublimation images obtained at $\alpha=3.5^{\circ}$


## Nacelle Trip Drag Study:

In addition to the wing/body trip drag study, a reduced trip drag study was conducted on the nacelles. For this nacelle study, the trip disk height on the wing was kept constant at 0.012 inches.

Three different trip heights were applied on the nacelles, both internally and externally, at 0.6 inches downstream from the leading edges of the nacelles. This results in a large number of disks on the nacelles-the total number of trip disks on the nacelles is about $49 \%$ of the total number on the wing.

Once again, the forces and moments were measured at the supersonic cruise Mach number, 2.4, and a series of sublimation images were obtained at the cruise angle-of-attack.


## Nacelle Sublimation:

This is an example of a sublimation image on the nacelles for a trip height of 0.014 inches. The nacelle transition location is not as easily discernible as the wing transition is in these images. In addition, only one of the four external surfaces of each nacelle is visible. Therefore, it is difficult to locate the transition line on the complete external surfaces of the nacelles due to poor visibility in the sublimation region and around the nacelles. As a result, the laminar run correction will assume that transition occurs at the trip location.


## WBND Drag Corrections:

The plot on the left shows the drag coefficient for the W/B/N/D configuration at the cruise angle-of-attack as a function of the nacelle trip height. In addition, the complete laminar run and trip drag corrections for the W/B/N/D configuration are shown on this plot. Both the laminar run and trip drag corrections consist of a separate wing and a separate nacelle contribution.

The wing contribution to the laminar run correction is constant for all nacelle trip heights and is obtained from the wing/body study for a disk height equal to 0.012 inches. The laminar run contribution from the nacelles is based on the assumption that transition occurs at the trip location for all trip heights. Therefore, the nacelle laminar run contribution is also constant. This laminar run correction only takes into account the external surfaces of the nacelles because the internal skin friction, which includes the assumption of laminar flow between the leading edge and the trip location, is already removed from the wind tunnel data. The total laminar run correction ( 2.0 counts) is added to the drag data for each trip height.

Similar to the laminar run correction, the wing contribution to the trip drag is based on the value of the wing/body trip drag for a disk height of 0.012 inches. The nacelle contribution is computed by scaling the corresponding trip drag for the wing/body by 0.49 , which accounts for both the internal and external dots. The combined trip drag value is subtracted from the drag for each disk height.

Once these corrections are applied, the final values agree with each other.


## Comparison with CFD results:

The plot on the left shows the drag polar including both final corrected wind tunnel and CFD data for the W/B/N/D configuration near the cruise condition. The comparison of the wind tunnel data for the W/B/N/D configuration to the lift and drag coefficients computed by CFL3D using the Baldwin-Lomax turbulence model show good agreement.

The OVERFLOW results with the Spalart-Allmaras turbulence model are about five drag counts lower than the corrected wind tunnel data. It is believed that this under prediction of the drag is, once again, due to the turbulence model, but the CFD solutions for the complete configuration have not been closely examined.

## Work in ProgressA Closer Look at the B'Layer

- Why is the inboard, upper surface so difficult to trip?
- Why do we see prompt transition on the outboard wing?
- Approach: Study boundary layers on flat plate and wing (CFD results provided by Scott Lawrence/Ames)


## Work in Progress-A Closer Look at the Boundary Layer:

The authors are now going to switch directions from focusing on methodology to attempting to understand why a modified approach to trip drag is required. The main question that comes to mind is "why is the inboard, upper surface so difficult to trip?", whereas, prompt transition on the outboard wing is easy to obtain. Our initial approach to answering this question is to study the boundary layers on both flat plates and the baseline TCA wing using CFD results provided by Scott Lawrence at NASA Ames Research Center.


## Laminar Boundary Layer Profiles:

The plot on the left gives laminar boundary layer profiles computed by UPS at 0.6 inches downstream of the leading edge on a flat plate and on the baseline TCA wing at both an inboard and outboard section. The UPS solutions for the TCA were obtained at $3.6^{\circ}$ angle-of-attack. Also shown on the plot, as a solid black line, is the incompressible Blasius solution for a flat plate. The heights of the trip disks applied to the wing in Test \#1679 are indicated on the plot to show the disk heights relative to the boundary layer thickness.

The first observation in comparing the Blasius solution to the UPS solution for the flat plate is that compressibility thickens the boundary layer. This effect would require a larger disk height to trip the flow. Next, examining the differences in the boundary layer profile between the flat plate and TCA wing suggests that the gentle, favorable pressure gradient in the outboard section tends to slightly thin the boundary layer. This can be seen by comparing the solid, blue line with the dashed, green line. The opposite effect is observed for the inboard section where the boundary layer is thicker than on the flat plate (compare the dashed, red line to the solid, blue line). This may be a result of the attachment line boundary layer on the blunt leading edge. This increase in the boundary layer thickness implies that a larger trip height is required in the inboard section, but doesn't completely explain the why even the largest disk height does not cause prompt transition.


## Mach and $\mathrm{Re}_{y}$ in Boundary Layer:

The plot on the left gives the distribution of Mach number through the laminar boundary layers computed by UPS on a flat plate and the TCA wing. This plot shows that the flow at the top of all of the trip disks that were tested is supersonic.

The plot on the right provides the distribution of Reynolds number based on local flow properties with the length scale equal to the height off the surface. This Reynolds number is being examined because the value of the Reynolds number at the trip height, referred to $\mathrm{Re}_{\mathrm{k}}$, has been used by researchers in the past to predict the required trip height that will cause boundary layer transition. This plot shows that the value of the inboard trip height Reynolds number is lower than that of the flat plate and the outboard section. This further implies that a larger disk height is needed to trip the flow.

## Summary

- Improved laminar run \& trip drag corrections for supersonic cruise performance testing
- Careful flow visualization revealed delayed transition, even for the largest disks
- New code developed for laminar run corrections
- Corrected data could be analyzed for trip drag
- Despite initial appearances, our data are consistent with (most) previous results
- Our approach extends prior work, taking into account the effects of blunt leading edges


## Summary:

In summary, an improved laminar run and trip drag correction methodology for supersonic cruise performance testing was derived. This method required more careful analysis of the flow visualization images which revealed delayed transition particularly on the inboard upper surface, even for the largest trip disks. In addition, a new code was developed to estimate the laminar run correction. Once the data were corrected for laminar run, the correct approach to the analysis of the trip drag become evident.

Although the data originally appeared confusing, the corrected data are consistent with previous results. Furthermore, the modified approach, which was described in this presentation, extends prior historical work by taking into account the delayed transition caused by the blunt leading edges.

## HSR Annual Airframe Technical Review <br> Follow-up "Grit Drag" Test in 4/99

- Repeat portions of Test \#1679
- Investigate shorter trip disks to verify that a drag plateau exists
- Study Re \& angle-of-attack effects
- Can we improve the present method?
- Are there more efficient and/or effective trips?
- Intentionally contaminate attachment line
- Different trip configurations on outboard wing
- Is it really necessary to trip the flow at all?

Follow-up Grit Drag Test in 4/99:
A follow-up grit drag test is planned for April 1999 at NASA Langley's UPWT Test Section \#2. First, portions of Test \#1679 will be repeated to confirm the general trends in the wing/body data which were observed during that test and reported in this presentation. Shorter trip disks will be tested on the wing and the resulting data will be examined to verify the existence of a drag plateau. In addition, some testing will be done to study the effects of varying Reynolds number and angle-of-attack on the wing transition location.

However, the main purpose of the follow-up test will be to try to improve the present method, which is time-consuming in both the testing and analysis aspects. First, more efficient and/or effect trip mechanisms will be examined. For the inboard section, intentionally contaminating the attachment line will be attempted. In the outboard region, different trip configurations will be tested. Second, it may be possible that for the baseline TCA configuration, it is not really necessary to trip the flow at all. If the free transition data corrected for laminar run provides the final answer, then testing without the application of trip disks is possible as long as other aerodynamic characteristics are not adversely affected.

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This test was done for Configuration Aerodynamics in support of WBS 4.3.1.1, Analysis and Testing Methods.
Technology Development
Session 4: Testing


There were two objectives for this test. First, was to assess the reasons why there is approximately 1.5 drag counts (cts) discrepancy between measured and computed drag improvement of the Non-linear Cruise Validation (NCV) over the Technology Concept Airplane (TCA) wing body (WB) configurations. The Navier-Stokes (N-S) pre-test predictions from Boeing Commercial Airplane Group (BCAG) show 4.5 drag cts of improvement for NCV over TCA at a lift coefficient (CL) of 0.1 at Mach 2.4. The pre-test predictions from Boeing Phantom Works - Long Beach, BPW-LB, show 3.75 drag cts of improvement. BCAG used OVERFLOW and BPW-LB used CFL3D. The first test entry to validate the improvement was held at the NASA Langley Research Center (LaRC) UPWT, test number 1687. The experimental results showed that the drag improvement was only 2.6 cts, not accounting for laminar run and trip drag. This is approximately 1.5 cts less than predicted computationally. In addition to the low Reynolds Number (RN) test, there was a high RN test in the Boeing Supersonic Wind Tunnel (BSWT) of NCV and TCA. BSWT test 647 showed that the drag improvement of NCV over TCA was also 2.6 cts , but this did account for laminar run and trip drag. Every effort needed to be done to assess if the improvement measured in LaRC UPWT and BSWT was correct.

The second objective, once the first objective was met, was to assess the performance increment of NCV over TCA accounting for the associated laminar run and trip drag corrections in LaRC UPWT. We know that the configurations tested have laminar flow on portions of the wing and have trip drag due to the mechanisms used to force the flow to go from laminar to turbulent aft of the transition location.
Objectives

- Assess the reasons why there is $\sim 1.5 \mathrm{ct}$ discrepancy
between measured and computed drag improvement of
NCV over TCA

| - OVERFLOW: |
| :--- |
| - CFL3D: |$\quad \Delta \mathrm{CD}=4.5 \mathrm{cts}$

- Test 1687:
- Assess the performance increment of NCV over TCA with
associated laminar run and trip drag corrections.

This is a photograph of the TCA WB configuration tested. The tripping mechanisms used on the wing upper and lower surfaces are 0.05 -inch diameter epoxy discs spaced 0.2 -inches apart, center to center, and are 0.012 inches high. They are placed 0.6 inches aft of the leading edge in the streamwise direction on all wing surfaces. The same tripping mechanisms, of the same diameter and height, are also placed 1 inch aft of the nose, but spaced 0.1 inches apart.

The TCA WB shown here is the solid wing model known as TCA 2 b . The TCA 2 b model has been tested in the LaRC UPWT three (3) times in test section 2. It also was tested in BSWT as part of the high RN test of the NCV model and as part of an Internal Research and Development (IR\&D) test.


This is a photograph of the NCV WB configuration tested. The exact same trip configuration used on TCA 2 b was used for the NCV model. This trip configuration is called the baseline trip configuration. Any other trip configuration is considered an alternate trip configuration. Alternate trip configurations were also applied to the NCV model to assess if it might be more effective to trip the flow and cause the drag improvement to increase. Alternate trip configurations were also used to assess the trip drag correction. A similar study of alternate trip was done very extensively on TCA 2 b during test 1679 .

NCV WB


This is the drag polar from the BCAG OVERFLOW calculations showing 4.5 cts drag improvement of NCV over TCA WB at cruise CL of 0.1 at wind tunnel RN of 4 million/ft at Mach 2.4. The Spalart-Allmaras turbulence model was used in OVERFLOW for the N-S calculations.

This is the drag polar of the BPW-LB CFL3D calculations showing 3.75 cts drag improvement of NCV over TCA WB as cruise CL of 0.1 at wind tunnel RN of 4 million/ft. The Baldwin-Lomax turbulence model was used with CFL3D for the N-S calculations.


This figure shows drag versus lift for NCV and TCA WB as measured during test 1687. The plot on the left is a close-up of the drag polars around minimum drag for the two configurations. The plot on the right is a close-up around cruise. The solid curve is for TCA and the dashed curve is for NCV. The curves represent the mean, using a $7^{\text {th }}$ order polynomial fit, through all runs within test 1687 of the same configuration. Test 1687 shows the drag improvement for NCV over TCA WB is only 2.6 cts . This is approximately 1.5 cts less than the calculated pre-test predictions.

[^26]$\begin{array}{ll}\angle 6 \cdot \varepsilon & 0+2 \\ 06 \cdot \varepsilon & 0+Z \\ N y & H O \forall W\end{array}$


## $\Delta \mathrm{CD}=-2.60$ cts @ $\mathrm{CL}=0.1$



TCA vs NCV WB
Langley UPWT Test 1687

The approach taken to meet the objectives of the test is to measure forces and moments on both the TCA and NCV; get off body flow characteristics using Laser Vapor Screen (LVS) for both configurations; get TCA and NCV surface flow characteristics using UV oil and mini tufts; and test NCV with alternate trip configurations.
Approach
WB Forces and Moments


- Alternative trip configurations

Test procedures were carefully looked at and followed throughout the test to make sure the most accurate data possible was obtained to reduce errors that might have contributed to the measured drag improvement during test 1687. Angle of attack is one of the biggest influences on drag. An observation made during the test just prior to this one was that the adapter between the sting and strut needed to be replaced. The hardware set up produced an error in alpha that was not repeatable. In addition to installing a new adapter, brass buffers were added on the retaining area of the sting. The fit between hardware pieces uses a cylindrical fit. The buffers allow for the errors in alpha to be reduced when the system is under load.

Another component of angle of attack is the droop angle. This is the angle incurred by the weight of the model on the hardware support system. This angle gets measured for both the upright and inverted runs. The angle should be equal unless there is slop in the system or if the fit of the model on the balance is not perfect. For TCA, the upright is different from the inverted droop angle. The values of $-2.92^{\circ}$ upright and $-2.54^{\circ}$ inverted were consistent with previous tests. For NCV, the upright and inverted droop angles were the same. The value of $-2.83^{\circ}$ upright and inverted was different from test 1687 . It is felt that the difference between tests is due to the hardware change. While for TCA, the difference between upright and inverted is tied up in the fit of the model on the balance.

An additional component of angle of attack is the balance to model misalignment angle. The model was designed such that the balance is installed in the model at a $2^{\circ}$ angle down at the front of the balance. Models are rarely built such that this angle is met to within $0.001^{\circ}$ accuracy. Test 1687 used the theoretical value for the balance to model misalignment angle. Test 1703 used the NASA Ames Research Center (ARC) QA measurements of $-2.0061^{\circ}$ for TCA and $-2.0086^{\circ}$ for NCV. What these new values ultimately effected was the computed stream angle correction. After test 1703 while QA'ing TCA and NCV, the balance to model misalignment angle was measured. There was a larger difference in the measurement of the TCA than for the NCV. The TCA LaRC QA measurement was smaller by $0.0166^{\circ}$. The NCV LaRC QA measurement was $0.002^{\circ}$ smaller. These new measurements, if applied to the data, would not affect the data because the stream angle correction would compensate for the difference.

Another large influence to drag is the actual instrument used for measuring forces and moments. Balance UT65A was used for all HSR tests in the LaRC UPWT for TCA and NCV. The balance was monitored throughout the test. It showed shifts in normal, side, and axial force gages. Most of the shifts were within tolerances required by HSR. There was a time these shifts were outside of HSR tolerances and it was found that the UV oil damaged the moisture proofing on the gages. After this discovery, the balance was removed for the gages to be moisture proofed and for the rest of the test the balance was protected during the UV oil runs.

Another large correction to drag is the pressure acting on the balance by the cavity in the aftbody. The aftbody for the NCV configuration has an unusual shape due to the optimization process. The TCA aftbody cross section is circular. It was thought that the pressures across the aftbody shape might not be constant. Additional pressures were measured just behind the balance and just ahead of the fuselage exit at various radial locations.

The test was heavily supported on and off site. A large amount of flow visualization data was acquired. This required good documentation of the results. Images were acquired using Hasselblad and high resolution digital cameras. In addition to still images, standard and digital video cameras were used for moving images. The images, plots, data, and notes were posted on ADAPT daily. ADAPT is a secure website being used within the HSR community. No data was ever passed to others within HSR in a non-secure way.

The reason for posting all this information daily is so that others working on this test that couldn't be on site could participate during the test. Telecons were held daily between experimentalists and computational analysts. The team worked together to analyze the results and help make any changes necessary in the test plan.

When ever an interesting feature was detected that hadn't been looked at as part of the pretest predictions, a computational solution or analysis was done to help the team assess what was observed.

This test could not have been successful as it was without the teamwork of NASA and Boeing personnel.

The results from the test are presented as outlined for the rest of the presentation. First, the short and long term repeatability of the data will be shown. Second, the mini trip drag study that was done resulted in a laminar run and trip drag corrections to be made. Finally flow visualization images will be shown for the UV oil and LVS runs.
$\quad$ Outline of Results

- Short \& Long Term Repeatability
- Laminar Run and Trip Drag Corrections
- uv Oil
- Laser Vapor Screen

This figure shows the short-term repeatability of the lift and pitching moment for TCA. These are mean curves through all the repeat runs for this configuration during the test. Short-term repeatability was very good for lift and pitching moment.


This figure shows the short-term repeatability of drag for TCA. A $7^{\text {th }}$ order polynomial fit was put through the repeat runs to get a mean curve. The left plot shows a close-up around minimum drag. The right plot shows a close-up around the cruise point. The minimum to maximum range of drag at a CL of 0.1 shows a repeatability of less than $\pm 0.5 \mathrm{cts}$. This is within the maximum tolerance required by HSR for supersonic cruise performance out of a wind tunnel test.





Repeatability in $\mathrm{CD}< \pm 0.5$ cts



This figure shows the long-term repeatability of the lift and pitching moment for TCA. These are mean curves through all the repeat runs for this configuration during each test. Long-term repeatability was very good for lift and pitching moment between tests 1687 and 1703. It was slightly worse between tests 1703 and 1679. Test 1679 was the very first time the TCA- 2 b model was tested. The change in angle of attack between tests 1703 and previous tests at the CL of 0.1 is $0.0025^{\circ}$ for test 1687 and $0.02^{\circ}$ for test 1679.


This figure shows the long-term repeatability of drag for TCA. These are mean curves through all the repeat runs for this configuration using a $7^{\text {th }}$ order polynomial fit through all the data during each test. The left plot shows a close-up around minimum drag. The right plot shows a close-up around the cruise point. Long term repeatability was very good for drag between tests 1687 and 1703. It was worse between tests 1703 and 1679. The change in drag between tests 1703 and previous tests at the CL of 0.1 is 0.1 cts less for test 1687 and 0.7 cts less for test 1679.


This figure shows the short-term repeatability of the lift and pitching moment for NCV. These are mean curves through all the repeat runs for this configuration during the test. Short-term repeatability was very good for lift and pitching moment.


This figure shows the short-term repeatability of drag for NCV. A $7^{\text {th }}$ order polynomial fit was put through the repeat runs to get a mean curve. The left plot shows a close-up around minimum drag. The right plot shows a close-up around the cruise point. The minimum to maximum range of drag at a CL of 0.1 shows a repeatability of less than $\pm 0.25 \mathrm{cts}$. This is within the maximum tolerance required by HSR for supersonic cruise performance out of a wind tunnel test.

This figure shows the long-term repeatability of the lift and pitching moment for NCV. These are mean curves through all the repeat runs for this configuration during each test. Long-term repeatability was very good for lift and pitching moment between tests 1687 and 1703. The change in angle of attack between tests 1703 and 1687 at the CL of 0.1 is $0.025^{\circ}$ for test 1687.
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 $\frac{\text { SYMBOL }}{---\ldots .}$

Through All Repeat Runs

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This figure shows the long-term repeatability of drag for NCV. These are mean curves through all the repeat runs for this configuration using a $7^{\text {th }}$ order polynomial fit through all the data during each test. The left plot shows a close-up around minimum drag. The right plot shows a close-up around the cruise point. Long term repeatability was very good for drag between tests 1687 and 1703. The change in drag between tests 1703 and 1687 at the CL of 0.1 is 0.4 cts less for test 1687 .


This figure shows the drag improvement for NCV over TCA for test 1703. NCV is 2.3 cts less than TCA. Recall test 1687 showed 2.6 cts improvement for NCV over TCA.






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\begin{array}{ll}
\text { SYMBOL } & \begin{array}{l}
\text { CONFIGURATION } \\
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\end{array} \\
\text { TCA 2b WB } \\
\text { NCV WB }
\end{array}
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Langley UPWT Test 1703

In an attempt to check if the tripping mechanisms used were the right ones for NCV, CFD pre-test calculations were done to assess the boundary layer thickness and the location of the baseline trip configuration with respect to the pressure distribution. The long dashed vertical lines are the trip discs located 0.6 inch aft of the leading edge in the streamwise direction. The dash-dot vertical lines represent the trip discs to be located at 1 inch aft of the leading edge. The top two plots are of the pressure distribution on the upper wing surface at two spanwise locations close to side of body. It was felt that the baseline location was not tripping the flow effectively because they were ahead of the pressure peak. The bottom two plots are of the boundary layer and displacement thickness. At 0.6 inch aft of the leading edge, the baseline trip discs were 0.012 inches high. This discs are taller than the boundary layer thickness. At 1 inch aft of the leading edge, it was felt that 0.015 -inch high discs would be a good starting point to evaluate the trip effectiveness since this is the height of the boundary layer thickness. It was also evident that the trip configuration only needed to be replaced on the inboard upper wing surface. The baseline trip configuration was effective everywhere else.

In addition to the baseline trip configuration and 0.015 -inch high discs located 1 inch aft of the leading edge, three other trip configurations were investigated as part of the mini trip drag study on the NCV model. The trip discs were totally removed from the upper and lower wing surfaces to get the data for free transition. The other trip heights investigated on the inboard upper wing surface were 0.012 and 0.016 inch high discs.


The 0.016 inch high discs on the inboard upper wing surface were the most effective, but still did not force the flow to go turbulent just aft of the discs. There still was a small amount of laminar run past the disc location. The drag improvement of NCV over TCA using these discs on the NCV model increased the improvement only by 0.17 cts from the baseline trip configuration for a total drag improvement of 2.57 cts . This data is still not corrected for the laminar run and the trip drag.


To be able to make the laminar run correction, a series of images were taken during the sublimation flow visualization runs. Sublimation is a material that sublimes on the area of the wing that is turbulent. The material that remains represent $s$ the areas of laminar run. These images are scanned and read into a program like Adobe Illustrator. The transition line is laid out on the image and points representing the line as a percentage of the wing chord are outputted. This data is used in a boundary layer code that was developed at NASA Ames Research center by Robert Kennelly. The code outputs the skin friction drag based on flat plate boundary layer theory on a fully turbulent and mixed laminar-turbulent wing flow. The difference between these outputs is the drag due to the laminar run that gets added to the total uncorrected drag.

I will be showing the images of some of the sublimation flow visualization runs on the NCV model. A trip drag study was done on TCA during test 1679. The results of the laminar run and trip drag analysis from the data acquired during test 1679 will be applied to the TCA data from this test.

This image is the transition line laid out for the free transition case on the upper wing surface of the NCV model. The points represent the transition line for the boundary layer code to compute the laminar run correction.


This image is from the sublimation run on the upper wing surface. The trip configuration here is 0.016 -inch high discs located 1 inch aft of the leading edge inboard and 0.012 inch high discs located 0.6 inch aft of the leading edge outboard. To help determining the location of transition, two discs were removed on wither side of a disc to create a lone disc in regions of the wing. These lone discs help during a sublimation run to determine where transition occurs. A $15^{\circ}$ turbulent wedge is created where the flow transitions from laminar to turbulent flow. These wedges can be seen readily wherever a lone disc exists. This helps the transition line to be laid out on the image on the wing surface. The transition line can sometimes be difficult to observe due to not being able to distinguish laminar from turbulent flow. The best attempt was made to lay out where turbulent flow starts.


This is a compilation chart of the transition line for each of the trip configurations tested for the upper wing surface. For each transition line, a corresponding laminar run correction is listed in the legend. An interesting observation is that the baseline trip was just as effective as the alternate trip location except for a small region of the wing near side of body. The difference between the baseline and the alternate trip laminar run is less than 0.2 cts.
UPWT Test \#1703


Sublimation flow visualization runs were also done on the lower wing surface for the 2 trip configurations tested. This is an image of the transition line laid out for the free transition case.

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\mathrm{NCV}
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Transition on Lower Wing

This is an image of the transition line laid out for the baseline trip configuration on the lower wing surface.

NCV Transition for $\mathrm{k}=0.012 / \mathrm{x}=0.6$ on Lower Wing


This compilation chart show the difference between free transition and the baseline trip configuration. Note the amount of laminar run correction made for the lower wing surface for each configuration.


This plot shows the uncorrected and laminar run corrected drag at the lift coefficient the sublimation runs were acquired. The procedure on how trip drag is determined is described in the paper that proceeded this one by Aga Goodsell and Robert Kennelly. Due to this being a mini trip drag study, not enough trip configurations were tested to be able to determine where the drag plateau occurs. An estimation was made that the plateau occurs at zero trip height and that the rest of the data needed to cross the plateau at a trip height of 0.007 inches. These estimations came about from the trip drag study done on TCA during test 1679. A linear fit is put through the data for the configurations tested with trip discs. The difference between the linear fit and the plateau is the trip drag correction.
NCV WB Mini Trip Drag Study

| Average, uncorrected | - Plateau, first point |
| :--- | :--- |
| Corrected for laminar run | -Trip drag contribution (linear) |


upper inboard wing $k @ x=1 \mathrm{in}$, in
8

This plot shows the drag corrections made for laminar run and trip drag for each configuration tested. These corrections are applied to the whole drag polar. The laminar run correction is added, while the trip drag correction is subtracted.
NCV WB Mini Trip Drag Study @CL=0.0875
Laminar Run and Trip Drag Corrections


After making these corrections, the drag improvement of NCV over TCA only improved by 0.2 cts more than that determined using the uncorrected data. This applies to all tests and all trip configurations. For the baseline trip configuration the improvement of NCV over TCA is 2.8 cts . For test 1703 , the improvement is 2.5 cts . The difference between the 2 tests is wrapped up in the long-term repeatability. For the alternate trip location and the highest height, the improvement is 2.8 cts . The difference between the baseline and the alternate trip configuration is wrapped up in the cumulative error in laminar run, trip drag, and data repeatability.

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& \text { Laminar Run and Trip Drag Corrected } \\
& \text { Drag Improvement of NCV Over TCA }
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| Trip Configuration | Uncorrected Drag <br> Improvement | Corrected Drag <br> Improvement |
| :---: | :---: | :---: |
| Test 1687 <br> Baseline Trip | 2.6 | 2.79 |
| Test 1703 <br> Baseline Trip | 2.3 | 2.49 |
| Test 1703 <br> Trip Height $=0.016$ " at X=1" <br> on Inboard Upper Wing | 2.57 | 2.8 |

This is a summary image of the UV oil run done on TCA at the supersonic cruise angle of attack. The upper most image is of the upper wing surface. Just under the upper wing is the lower wing image from the UV oil run. At the bottom left is side of body near the wing apex and on the right is looking at the forebody upper surface. Hasselblad cameras were mounted on both windows. This image makes up 2 runs. One run was with the model mounted wings vertically to get the upper surfaces. The second run was with the model mounted wings horizontally to get side of body. The images depicts the surface flow on the model. The flow is very clean without any separation.
?
Forebody
$\begin{array}{r}\text { II } \\ \frac{9}{2} \\ \frac{1}{2} \\ \hline\end{array}$

TCA UV


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 1This is a compilation of images from UV oil runs on the NCV model at the supersonic cruise angle of attack. Here you will notice that the upper wing surface has some distinct features representing vortices. There seems to be two traces of the vortices on the wing. One goes streamwise over the inboard nacelle location and the other over the outboard nacelle location. The one that passes over the inboard nacelle location seems to be emanating from the wing body junction. The one that passes over the outboard nacelle location seems to come from the same region but moves span wise as it goes downstream. The vortices will be evident when the LVS images are seen.

The flow over the lower wing surface, side of body, and upper forebody look clean. There is no indication of separation in any of these regions.

NCV

For every UV oil run, images for a maximum of three (3) angles of attack could be obtained. This figure shows the UV oil results for the cruise and two (2) higher angles of attack for the NCV model. The two (2) vortices that seem to pass over the upper wing surface become more evident at the higher angles of attack. They become stronger and cause flow separation on the wing surface.


This figure show the UV oil results on the NCV model at the same three (3) angles of attack. The flow turns more and the attachment line moves as the angle of attack increases.


Post test, a comparison was made between the UV oil images and the CFD calculations that were done at NASA Ames Research Center by Scott Lawrence using UPS. The top figure compares the fully turbulent solution from UPS with the UV oil image. The bottom figure compares it with a fully laminar solution. It was determined from this that the original thought that the upper wing surface had laminar flow on it was incorrect. The flow is turbulent except ahead of the trip mechanisms and slightly behind.
NCV UPPER WING SURFACE FLOW COMPARISON
UPS Solutions, AOA $=4.5, \mathrm{Re}=4 \mathrm{Mft}$

0.015 UW Trip Dots



The UPS solution was also compared to the UV oil images for side of body. It was feared that the flow over the NCV forebody was laminar. It was hard to see during the sublimation runs. The top figure compares the fully laminar UPS solution with the UV oil image. The bottom two (2) figures compare fully turbulent UPS solutions with the UV oil image. The turbulence models being compared are Spalart-Allmaras and Baldwin-Lomax. The flow over the forebody is more like the fully turbulent UPS solutions.
NOV SURFACE FLOW COMPARISON
UPS SOLUTIONS, AOA $=4.5, \operatorname{Re}=4 \mathrm{Mft}$
(FINE GRID,FUSELACE)

This figure shows the set up used for the Laser Vapor Screen runs. The top left image shows the Spectra-Physics laser used. The right image shows a light box attached to the tunnel window. The model was translated axially so that the laser light sheet was at a different location on the model for tracking any shock and vortices. A Hasselblad camera was mounted in a housing to the tunnel ceiling and was remotely triggered in the control room to take the still image. A mini video camera was in a housing mounted to the hardware holding onto the model. The two (2) cameras can be seen in the bottom left image.


The model was painted black with reflective targets placed 4 inches apart from the aftbody forward. These targets were used to stop the axial movement of the model when the laser light sheet passed to take the still images. A video was recorded from the mini video camera signal of the entire LVS run. Any shocks and vortices were tracked for every angle of attack during the LVS run.
Model LVS Targets

## $X=1715^{\prime \prime} \quad X=2193^{\prime \prime}$


-Targets 4 inches Apart from Fuselage Base
-Left Hand Upper Side of Model Painted Flat Black

During the LVS run water was injected in the tunnel to create condensation to help better pick out shocks and vortices. This figure compares the LVS image between TCA and NCV at airplane station 2671 inches. The very dark area is the wing behind the light sheet. Above the light sheet close to the leading edge, a compression shock can be seen sitting on the wing. There is also a horseshoe vortex at side of body on the NCV model. It is hard to see it in this figure but the video clearly shows it. The TCA does not have a horseshoe vortex as strong as on the NCV as observed from the mini video camera.


This figure compares the LVS run for TCA and NCV at the aftbody location, which is downstream of the wing trailing edge. This figure helps show that the NCV model has these two (2) vortices that pass over the upper wing surface. The NCV wake shows traces of two (2) vortices. The TCA wake is clean.


This figure compares TCA and NCV LVS results at an angle of attack of $1.5^{\circ}$ above the cruise angle of attack for each configuration at airplane station 2671 inches. For TCA, the image is for $5^{\circ}$, and for NCV the image is for $6^{\circ}$. The dark spots above the light sheet for NCV are the horseshoe vortex and the vortex that passes over the inboard nacelle location. These cannot be seen for TCA.


This figure shows the wake behind the wing trailing edge for TCA and NCV at $1.5^{\circ}$ higher than the cruise angle of attack. The NCV wake clearly show three (3) vortices: a horseshoe vortex at side of body, a vortex that passes over the inboard nacelle location, and a vortex that passes over the outboard nacelle location. The TCA wake is just starting to show very weak vortices over the wing.


The forces and moments acquired on the NCV and TCA models are as good as can be with the angle of attack and balance instrumentation used during the test. Both the low and high RN tests did not measure the calculated drag improvement of NCV over TCA. Both tests show that NCV is approximately 2.8 cts better than TCA.

Both UV oil and LVS are good flow visualization tools for acquiring surface and off body flow characteristics. The two techniques go together well. LVS backs up what is seen in the UV oil images.

The only items that are needed to close-out the experimental portion of the NCV testing are to get surface pressure measurements, CFD solutions of as-measured (QA) geometry, and validate effect of balance bore angle measurements by LaRC on experimental results.
Flowfield Studies for the TCA/NCV
High Speed Aerodynamics, Long Beach
Configurations
Raul Mendoza, Chih Fang Shieh, and P. Sundaram
Configuration Aerodynamics
The Boeing Company
Phantom Works, Long Beach
High Speed Research Airframe Review
Anaheim, California
February 8-11, 1999
Outline
This paper will follow the standard outline shown here. First, an introduction will be given,
and the objective of this work will be stated. Then, the approach will be described and some
results presented. Finally, a summary and some conclusions will be offered.

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\begin{aligned}
& \text { - Background } \\
& \text { - Objective } \\
& \text { - Approach } \\
& \text { - Results } \\
& \text { - Summary and conclusions }
\end{aligned}
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Background The Technology Concept Airplane (TCA) was developed using linear design and analysis
 Boeing Phantom Works (BPW) at Long Beach performed non-linear aerodynamic shape





 design.

 configuration on the order of 4 to 4.5 counts over the TCA W/B at cruise conditions; but the improvement measured in the wind tunnel was on the order of only 2 to 2.5 counts. Even though the experimental drag reduction for the W/B/N/D configuration was also smaller than expected,

 on the W/B configurations.

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Background - TCA designed with linear methods

- Three TCA Cycle 2 configurations developed though
non-linear cruise-point optimization
- One optimized design selected for experimental
validation: NCV
- Wind-tunnel data for W/B configurations did not show
predicted performance gain of NCV over TCA
- Predicted $\Delta \mathrm{C}_{\mathrm{D}}: 4-4.5$ cts.
- Measured $\Delta \mathrm{C}_{\mathrm{D}}: 2-2.5$ cts.
Background: Drag Polars for TCA and NCV W/B
This figure compares predicted and measured drag polars for the TCA and the NCV W/B
 Navier-Stokes computations with the Baldwin-Lomax turbulence model show an expected drag improvement for the NCV W/B configuration of about 3.7 counts over the TCA W/B. However, (uncorrected) experimental data obtained during Test 1703 at NASA Langley's Unitary Plan Wind Tunnel (UPWT) give a drag reduction of only 2.3 counts. The experimental curves shown in this plot and in subsequent figures were obtained by fitting a cubic spline though all the repeat runs.
Predicted and Measured Drag Polars for TCA and NCV W/B

Predictions at High Re In order to eliminate trip drag uncertainties while minimizing the amount of laminar flow on
the models, both the TCA and NCV configurations were tested in the Boeing Supersonic Wind
Tunnel (BSWT) during a high Reynolds number entry (Test 647). For a Reynolds number of
Re $e_{c}=22.27 \times 10^{6}$, CFL3D Navier-Stokes computations performed on a dense grid with more than
11 million points predicted a drag improvement of the NCV W/B over the TCA W/B of
approximately 4.3 counts (as shown in the figure, the CFD data were extrapolated to a C $\mathrm{C}_{\mathrm{L}}$ of 0.1
using the local slope of the experimental curves). The drag improvement measured in the BSWT
tunnel at these conditions, however, was only 2.8 counts. This is 1.5 counts less than the
expected gain, based on uncorrected data.
Drag Predictions for TCA and NCV W/B Models at High Re

Reduction

 $R e_{c}=6.36 \times 10^{6}$, based on the mean aerodynamic chord of the $1.675 \%$-scale models), OVERFLOW Spalart-Allmaras calculations (performed at BCAG) predict a 4.5-count drag



 entries: Test 1687, 2.6 counts; and Test 1703, 2.3 counts. Also shown in this figure are predictions and wind-tunnel test results for a unit Reynolds number of $14 \times 10^{6} / \mathrm{ft}$ (or $\left.\operatorname{Re}_{\mathrm{c}}=22.27 \times 10^{6}\right)$. For this condition, the drag reduction predicted by both OVERFLOW and



 trip disks located 0.6" downstream from the leading edge), while the BSWT data correspond to
 rigid, theoretical configurations. (However, the Spalart-Allmaras turbulence model in the OVERFLOW code has a built-in transition model that simulates laminar flow with reduced skin friction immediately downstream of the attachment line at low Reynolds numbers.)
NCV Drag Reduction relative to TCA for W/B Configurations
CFD vs. Experiment
$M_{\infty}=2.4, C_{L}=0.1(1.675 \%-$ scale models $)$



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Objective The objective of several Computational Fluid Dynamics (CFD) studies performed at BPW
over the past few months was to identify and analyze potential factors that may have contributed
to the discrepancies between the predicted and the measured performance improvement of the
NCV W/B model.

- Analyze factors that may have contributed to discrepancies
between the predicted and the measured aerodynamic
performance of the NCV model
Approach Various hypotheses regarding the NCV performance shortfall have been suggested. In
order to test some of these hypotheses, several CFD activities were performed at Boeing
Phantom Works: flow analysis and flow visualization of the TCA and NCV W/B configurations,
grid density effects on Navier-Stokes solutions, analysis of the as-built NCV geometry to quantify
the effects of model fidelity, aeroelastic assessment of the TCA wind-tunnel model, boundary-
layer stability analysis, and mixed laminar/turbulent flow computations to estimate laminar run
corrections and trip drag.


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& \text { - TCA/NCV flow analyses and flow visualization } \\
& \text { - Grid density effects on CFL3D Navier-Stokes results } \\
& \text { - Effects of model fabrication imprecision on NCV } \\
& \text { - Effects of aeroelastic deformation on TCA } \\
& \text { - Boundary layer stability analysis } \\
& \text { - Mixed laminar/turbulent flow analyses } \\
& \text { - Laminar run corrections } \\
& \text { - Trip drag }
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 the wind-tunnel flow over the forward section of the NCV fuselage might not be fully turbulent as assumed in most Navier-Stokes computations. To try to answer this question, several CFL3D fully turbulent and fully laminar flow solutions were obtained. This figure compares surface particle traces on the TCA and NCV fuselage forbodies. With fully turbulent flow (BaldwinLomax), the patterns on the NCV are similar to those on the TCA. However, with fully laminar flow, the NCV shows an additional attachment line and a larger separation line than the TCA. It was argued that there could be a problem with laminar flow separation affecting the NCV forebody in the wind tunnel. However, comparisons of CFD surface particle traces with flow visualization results (ultra violet oil and colored oil) from the NCV flow diagnostics test (UPWT 1703) indicated that the flow on the NCV forebody appeared to be, in fact, fully turbulent and did not trigger any massive laminar flow separation that could cause an increase in the NCV drag.

TCA
(Fully Turbulent, B-L)
TCA
(Fully Laminar)
NCV
(Fully Turbulent, B-L)
NCV
(Fully Laminar)
3

Trip Dot Location
NCV Crossflow Velocity Vectors The Laser Vapor Screen (LVS) visualization technique that was employed during the NCV
flow diagnostics test to visualize off-body flow detected a series of vortices on the NCV model.
Of particular interest was a large vortex that originated near the wing apex and traveled
downstream along the wing/body junction. This vortex was stronger at high angles-of-attack, but
it was present even at an angle of $4.5^{\circ}$, which is near the design condition for the NCV
configuration. The turbulent CFL3D solution indicates that the vortex does exist at these
conditions, as illustrated in this figure by the crossflow velocity vectors at fuselage station
FS2100. Also shown in the figure are off-body particle traces that point to the wing apex as the
origin of the vortex. The computational results indicated that this vortex is indeed enhanced after
FS1500, as detected in the experiment. Neither the wind-tunnel LVS results nor the CFL3D
solutions showed any evidence of this vortex pattern forming on the TCA model at the design
conditions.

NCV Crossflow Velocity Vectors
High Speed Aerodynamics, Long Beach
CFL3D N-S (B-L), Fully Turbulent, $M_{\infty}=2.4, C_{L}=0.089, R_{c}=6.36 \times 10^{6}$
NCV Crossflow Velocity Vectors
High Speed Aerodynamics, Long Beach
CFL3D N-S (B-L), Fully Turbulent, $M_{\infty}=2.4, C_{L}=0.089, R_{c}=6.36 \times 10^{6}$
Fuselage Stations


TCA and NCV Pressure Distributions Computed off-surface pressure contours for the TCA and NCV wing/body configurations
are shown in this figure. At FS1500, the NCV results show a stronger flow compression than the
TCA as the flow moves upward from the wing surface in the wing/body junction region. At
FS2100, a spanwise-compression wave is observed near $33 \%$ semi-span. During UPWT Test
1703 , a standing shock was clearly observed on the NCV inboard upper wing surface instead of
the compression predicted by CFL3D at the cruise angle-of-attack. It is plausible that the
presence of this shock could increase the drag of the NCV model in the wind tunnel, adversely
impacting the NCV performance improvement. On the other hand, only a very weak shock was
detected during the Laser Vapor Screen (LVS) runs on the TCA wing near cruise conditions.
In order to resolve the shock wave issue, and perhaps be able to capture the strength of
the detected vortices more accurately, a grid refinement study was performed next for both the

## TCA and NCV Pressure Distributions


High Speed Aerodynamics, Long Beach

Beach
Grid Density Effects on N-S Solutions To investigate the effect of grid density on drag predictions, refined grids were generated,
and new CFL3D Navier-Stokes (Baldwin-Lomax) solutions were computed for both the TCA and
NCV W/B configurations. The new grids had eight times the number of points as the original
grids: the number of grid points was doubled in each of the three computational coordinates
compared to the original grids. Thus, the total number of grid points increased from 1.5 million to
11.5 million. Care was taken to ensure that the grid points inside the boundary layer were also
doubled, while maintaining the original $y^{+}$values next to the surface. The surface grids were
refined using a linear interpolation of the original grids. Thus, the surface representation of the
refined grids retained the same fidelity as the original grids.
In this figure, the original drag predictions are compared with the refined-grid Navier-Stokes
results for the TCA and NCV W/B configurations. It can be seen that the refined-grid predictions
fall almost on top of the corresponding original-grid polar, and therefore do not significantly
change the predicted drag increment between the TCA and NCV configurations $(3.3$ counts at
$C_{L}=0.09, M_{\infty}=2.4$ and $\left.\mathrm{Re}_{\mathrm{C}}=6.36 \times 10^{6}\right)$.
Effect of Grid Density on Drag Predictions for TCA and NCV W/B CFL3D Navier-Stokes (Baldwin-Lomax)
$M_{\infty}=2.4, \operatorname{Re}_{c}=6.36 \times 10^{6}$


Model Fabrication Imprecision Effects
(NCV)
$S$
$\vdots$
$\vdots$
The impact of model fidelity on drag predictions was also examined. Detailed
measurements of the NCV wind-tunnel model were used to create new computational grids. Noticeable differences between the theoretical NCV definition and the Quality Assurance (QA)
 impact of these imprecisions, CFL3D Navier-Stokes computations for the as-built geometry were performed and compared to the original drag predictions for $M_{\infty}=2.4$ and $\operatorname{Re}_{c}=6.36 \times 10^{6}$. The



 the left half is approximately 0.3 counts. Most of the drag increase is due to an increase in pressure drag, as the viscous drag of the as-built model is nearly the same as that of the asdesigned configuration.
A similar analysis of the TCA as-built model will be performed later when that the QA

[^27]Effect of Fabrication Imprecision on Performance of NCV W/B
High Speed Aerodynamics, Long Beach CFL3D N-S (Baldwin-Lomax)
$M_{\infty}=2.4, \mathrm{Re}_{\mathrm{c}}=6.36 \times 10^{6}, 93 \times 241 \times 65 \mathrm{C}-\mathrm{O}$ Grid

0.0120
Drag Coefficient, $C_{D}$
Fabrication Imprecision Effects: $\mathbf{C}_{\mathbf{P}}$ Distributions
This figure shows that the $\mathrm{C}_{\mathrm{p}}$ distributions of the NCV theoretical and QA models are fairly
similar. The most significant difference between the surface pressures of these two models is the
presence of a suction peak on the lower surface of the NCV QA geometry near the leading edge
of the right inboard wing (red dashed curve).
Pressure Distributions for the NCV W/B Model
Beach

 C

CFL3D Navier-Stokes (Baldwin-Lomax)
$M_{\infty}=2.4, C_{L}=0.089, \operatorname{Re}_{c}=6.36 \times 10^{6}$




$\begin{array}{lllllllll}0.15 \\ \mathrm{C}_{0} .0 & 0.2 & 0.4 & 0.6 & 0.8 & 1.0\end{array}$

$\begin{array}{lllllllllll}0.15 & & & & & & & \\ C_{p} & 0.0 & 0.4 & 0.6 & 0.8 & 1.0\end{array}$

Reduction This bar chart shows the impact of model fidelity on the CFD solutions for the NCV W/B. It
can be seen that after NCV model imprecision effects are taken into account, the predicted NCV
drag improvement over the TCA decreases from 3.3 to 2.6 counts at $M_{\infty}=2.4, R_{c}=6.36 \times 10^{6}$,
and $C_{L}=0.09$ (this is approximately the lift coefficient generated by the W/B configurations at
their respective cruise angle-of-attack). The NCV drag reduction measured during UPWT Test
1703 with the baseline trip configuration ( $k=0.012^{\prime \prime} @ x=0.6^{\prime \prime}$ from the leading edge), also
shown in this figure, is 2.1 counts for these conditions. Therefore, once the CFL3D predictions
account for NCV QA effects, the discrepancy between CFD and wind tunnel reduces from 1.2 to
only 0.5 counts. As stated before, the TCA as-built model is yet to be analyzed to determine the
exact impact of model imprecision on the CFD predictions.
NCV Drag Reduction relative to TCA for W/B Configurations


Aeroelastic Deformation Effects
(TCA)
Aeroelastic Deformation Effects (TCA) While a model undergoes deformation in the wind-tunnel due to aerodynamic loads, most
CFD analyses assume a rigid configuration. So, in order to investigate the effects of static
aeroelasticity on model performance, the deformations of the TCA Model 2 a W/B/N/D
configuration were measured at several angles-of-attack in the Unitary Plan Wind Tunnel (Test
1671 ) at $\mathrm{M}_{\infty}=2.4$ and $\mathrm{Re}_{\mathrm{c}}=6.36 \times 10^{6}$. The deformations (vertical wing displacement and wing
twist) measured for each of these angles-of-attack were then lofted onto the baseline TCA OML,
and CFL3D Navier-Stokes analyses were performed on each of the deformed geometries.
Changes in wing camber and deformation of the fuselage were assumed to be negligible.
The next four charts show the effects of aeroelastic deformation on the performance of the
TCA W/B/N/D configuration. Wind-tunnel data are compared with Navier-Stokes results for the
deformed configurations as well as computations for the rigid model.
The agreement between the predicted lift coefficient and test data improves when model
deflections are included in the computations, as the outboard wing unloads reducing the overall
lift. Due to the unloading of the outboard wing, there is a nose-up pitching moment. As a result,
the agreement between the computed and the measured pitching moments improves significantly
when the aeroelastic deformation of the wind-tunnel model is taken into account. However, the
drag polar shows that the impact of aeroelasticity on the aerodynamic performance of the model
is negligible, since there is no significant change in drag at a given lift coefficient. The unloading
of the outboard wing is evident in the pressure cuts shown ( $\eta=67.9 \%$ and $\eta=85.0 \%$, for
instance).




Effect of Deformation on TCA Model 2a W/B/N/D
CFL3D N-S (B-L) Solutions and Wind Speed Aerodynamics, Long Beach
$M_{\infty}=2.4, \operatorname{Re}_{c}=6.36 \times 10^{6}, \alpha=3.5^{\circ}\left(C_{L} \approx 0.1\right)$







Boundary Layer Stability Analysis
Boundary Layer Stability Analysis

 turbulent state (depending on the Reynolds number, among other factors). It has been suggested that the flow transition location on the TCA wing has been different than the transition





 layer trips, there is still a laminar flow correction that needs to be applied to the data. The first step in correcting the test data for a proper comparison to CFD results is to determine the transition location of the flow. For this purpose and in order to obtain some
 the TCA and NCV wings were obtained through boundary layer stability analyses. The predictions were based on linear stability theory and the envelope method, in which the flow is
 value of $e^{N}$, where $N$ is in the range of 8 to 9 . The linear stability equations were solved using

 the mean flow cannot change abruptly in the streamwise or crossflow directions. The boundary ayer profiles required by eMalik3d were generated from CFL3D N-S laminar flow solutions. The computations were performed at $M_{\infty}=2.4, \operatorname{Re}_{c}=6.36 \times 10^{6}$, and $C_{L}=0.09$.

• CFL3D
$\quad$ - N-S laminar boundary-layer profiles
• eMalik3d

- Linear stability
- Compressible flows
- General swept wings
$\quad$ - Parallel-flow assumption
- Envelope $\left(\mathrm{e}^{\mathrm{N}}\right)$ method
$\quad$ - Transition @ $\mathrm{N}=8$ or 9
BL Stability Analysis: N Factors The next two figures show the growth of the amplification factor, N , on the upper surface of
 predicted to grow faster than the $N$ factors on the TCA, except near the wing break ( $\eta=50 \%$, for instance) where the opposite is true. The baseline trip dot location ( 0.6 " from the wing leading edge) is also marked in these figures. It can be seen that at this location the N factors on the TCA wing have a value of less than 4. The N factors on the NCV wing at the same location are only slightly higher.
Boundary Layer Stability Analysis for the TCA W/B Configuration Neach
$N$ Factor


0.10 $00^{\circ}$

Boundary Layer Stability Analysis for the NCV W/B Configuration
dynamics, Long Beach


N Factor


1
 the TCA and the NCV were determined by NASA Ames researchers from sublimation images obtained during UPWT Test 1679 and UPWT Test 1703, respectively. No trip dots were placed on the wings for these particular runs. Two sets of data from eMalik3d are presented in each figure, showing the locations on the wing surface where the amplification factor has reached values of 8 and 9
It can be seen that the agreement between the stability analysis predictions and the experimentally-determined transition locations is quite good for the TCA model on the mid-wing рعч!ן predicts the transition location on the TCA wing to occur considerably further downstream than measured in the wind tunnel. There are several factors that may contribute to the disagreement.



 code. Furthermore, in the wind tunnel, the flow on the wing is contaminated by the turbulent wedge

 from the Navier-Stokes computations, in which the flow is assumed to be completely laminar.
The stability calculations on the NCV wing also correlate well with experiment, including the area near the wing/body junction. However, the good agreement observed between analysis and experiment in this region may be considered somewhat fortuitous, since some of the assumptions in the eMalik3d linear stability code are violated in this region of highly three-dimensional flow.
Boundary Layer Transition Location for the TCA W/B Configuration
High Speed Aerodynamics, Long Beach
$M_{\infty}=2.4, \operatorname{Re}_{c}=6.36 \times 10^{6}, C_{L}=0.09$

Boundary Layer Transition Location for the NCV W/B Configuration

$Q$ EOENNE
BL Stability Analysis: Transition Locations The next two figures compare transition locations for both models on the same plot. The
sublimation data shown in the first figure indicate that the experimental boundary layer transition
locations for the TCA and NCV configurations are rather similar everywhere on the wing upper
surface, including the area close to the wing/body junction. This was not expected, since the
models (and, therefore, the flows) are quite different from each other. The linear stability results,
on the other hand, show some differences between the predicted transition locations for the two
models, with the most significant differences occurring close to the wing/body junction, as seen in
the next figure.
Experimental BL Transition Location for TCA and NCV W/B

Predicted BL Transition Location for TCA and NCV W/B eMalik3d with CFL3D N-S (Laminar)
$\mathrm{M}_{\infty}=2.4, \mathrm{Re}_{\mathrm{c}}=6.36 \times 10^{6}, \mathrm{C}_{\mathrm{L}}=0.09$

BL Stability Analysis: $C_{p}$ Distributions The surface pressure associated with the laminar flow solutions used in the stability
analyses do suggest that the flow on the inboard wing upper surface of the NCV should transition
before the flow on the TCA. This figure shows the computed laminar pressure distributions for
the TCA and NCV W/B configurations, along with wind-tunnel pressures measured during UPWT
Test 1671 for the TCA Model 2a. The experimental pressures were obtained with 0.012 "-high trip
dots on the model. It can be seen that on the mid-span and outboard regions, the pressure
distributions for both models are fairly similar. However, on the inboard region of the wing
( $\eta=19.9 \%$, for example), the TCA and NCV pressures are quite different. The NCV, for
instance, displays a peaky distribution, with a stronger adverse pressure gradient close to the
leading edge that could promote transition of the boundary layer. The TCA model, on the other
hand, shows a milder pressure gradient. Thus, based exclusively on pressure distributions, it is
reasonable to expect the flow on the inboard portion of the NCV wing to transition sooner than
the flow on the TCA model (as predicted by the eMalik3d code, which assumes parallel flow and
does not take into account factors such as freestream turbulence, flow contamination from the
fuselage, or wall roughness).


Laminar Run Correction Once the boundaries of laminar flow have been defined (based on stability analysis or
sublimation photographs), an estimate of the laminar run corrections can be obtained. The first
step in this process would be to compute the fully turbulent drag, $\mathrm{C}_{\mathrm{D}}$ tur, of a model using CFD
Navier-Stokes solutions (equivalent flat-plate estimates for skin friction can also be used). Next,
drag predictions for mixed laminar/turbulent flows, $\mathrm{C}_{\mathrm{D}}$ mix, are obtained using the predicted or
measured transition locations. An estimate of the laminar-flow drag correction, $\mathrm{C}_{\mathrm{D}}$ cor, can then
be obtained by subtracting the combined laminar/turbulent-flow drag predictions from the
computed fully turbulent flow drag values. Note that if skin-friction estimates from flat-plate theory
are used instead of N-S solutions, any pressure drag contribution is assumed to be negligible.
In the figure, CFL3D results are shown for the TCA W/B configuration near the design
condition. In this case, experimentally-determined transition locations from UPWT Test 1679
were used to perform the combined laminar/turbulent flow analyses. The laminar run corrections
are seen to vary from 5.5 counts (for free-transition runs) to 1.9 counts (for runs with $0.014 "$-high
trip dots).
Laminar Run Correction
High Speed Aerodynamics, Long Beach ient drag, $C_{D \text { tur }}$, from CFD or flat-plate theory
inar/turbulent drag, $C_{D \text { mix }}$, using predicted (stability
Laminar Run Corrections for TCA W/B

Laminar Run Correction for TCA and NCV W/B This figure shows CFL3D-based results from BPW along with the flat-plate estimates from
NASA Ames. In general, the laminar-flow drag corrections obtained from CFD are approximately
0.2 counts higher than the flat-plate estimates for both the TCA and NCV configurations (except
for the free-transition case where this pattern reverses). The transition locations used to perform
the combined laminar/turbulent flow computations were defined by NASA Ames personnel from
sublimation testing with different trip dot heights. The trip dots on the TCA wing were located 0.6
in. from the leading edge (UPWT Test 1679), while the NCV results presented here correspond
to test runs for which the dots on this configuration were placed 1.0 in. downstream of the leading
edge (UPWT Test 1703). A laminar run correction for the NCV with the baseline trip
configuration ( $k=0.012$ " @ 0.6" from the leading edge) was also obtained ( 1.96 cts.), even
though it is not shown in the figure. The corresponding laminar flow correction for the TCA was
found to be almost identical ( 1.97 cts.), as seen in the figure.
Laminar Run Correction for TCA and NCV W/B

Free-Transition Data Corrected for Laminar Flow This chart illustrates the process for correcting free-transition data for laminar run in order to make a consistent comparison between the CFD fully turbulent drag predictions, $\mathrm{C}_{\mathrm{D} \text { tur }}$, and windtunnel measurements. The approach is to add the predicted laminar-flow drag correction, $\mathrm{C}_{\mathrm{D} \text { cor }}$, to the free-transition test data, $\mathrm{C}_{\mathrm{D} \text { mix }}$, as shown in this table. Since the wind-tunnel data are freetransition data, there are no trip drag corrections to apply. Once the laminar run adjustment is

 values shown in this table have been interpolated to a lift coefficient of 0.09). However, the agreement is not as good for the NCV model: CFL3D results for the theoretical configuration

 counts. After applying laminar run corrections to both models, the NCV improvement reduces to
 stated before, the CFL3D-predicted performance improvement for the NCV W/B at this lift coefficient is 3.32 counts, so that the difference between the CFD predictions (for the asdesigned models) and the corrected experimental data becomes 1.91 counts.

[^28]Free-Transition Data Corrected for Laminar Flow
High Speed Aerodynamics, Long Beach
Laminar run corrections obtained from CFL3D $N$-S solutions
$M_{\infty}=2.4, R e_{c}=6.36 \times 10^{6}, C_{L}=0.09$

|  |  | $\mathrm{C}_{\text {D tur }}$ | $\mathrm{C}_{\mathrm{D} \text { mix }}$ |  | $\mathrm{C}_{\text {D cor }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| TCA | CFL3D N-S (B-L) <br> UPWT Test 1679 | $\begin{aligned} & 120.44 \\ & 120.50 \end{aligned}$ | $\begin{aligned} & -\quad 114.92 \\ & =\quad 114.98 \end{aligned}$ | $=$ $+$ | $\begin{aligned} & 5.52 \\ & 5.52 \end{aligned}$ |
| NCV | CFL3D N-S (B-L) UPWT Test 1703 | $\begin{aligned} & 117.12 \\ & 119.09 \end{aligned}$ | $\begin{aligned} & -\quad 110.98 \\ & =\quad 112.95 \end{aligned}$ | $\begin{aligned} & = \\ & + \end{aligned}$ | $\begin{aligned} & 6.14 \\ & 6.14 \end{aligned}$ |
|  |  | $\mathrm{C}_{\mathrm{D}}$, TCA | $\mathrm{C}_{\mathrm{D}}$, NCV | Drag | Reduction |
| Wind-Tunnel Test |  | 120.50 | - 119.09 | $=$ | 1.41 |
| CFL3D N-S (B-L) |  | 120.44 | - 117.12 | $=$ | 3.32 |
| Test - CFD |  | 0.06 | 1.97 |  | -1.91 |

All drag values given in counts
Laminar Run Corrections: TCA-to-NCV Drag Reduction The data presented in the previous slide is repeated in this figure, which shows a bar chart
comparing CFD predictions with wind-tunnel measurements. The predicted NCV drag
improvement is seen to decrease by 0.69 counts after NCV fabrication imprecision effects are
accounted for in the CFLSD solutions (as explained earlier). The measured NCV drag
improvement over the TCA, on the other hand, decreases by 0.62 counts (from 2.03 to 1.41 cts.)
after the free-transition test data are adjusted for laminar run. As a result, the difference
between the CFD-predicted and the measured performance improvement for the NCV W/B over
the TCA W/B is still about the same, after corrections have been applied. Note that the
experimental free-transition TCA-to-NCV drag reduction was computed using data from two
separate tests: UPWT 1679 (TCA) and UPWT 1703 (NCV), since the NCV model was not tested
during Test 1679 and free-transition data for the TCA model was not acquired during Test 1703 .


After Corrections

Before Corrections
NCV Drag Reduction relative to TCA for W/B Configurations

Trip Drag The laminar run corrections computed in the previous section can now be used to estimate
the trip drag associated with different trip dot heights. First, the computed laminar-flow drag
corrections, $\mathrm{C}_{\mathrm{D}}$ cor, are added to the corresponding wind-tunnel drag values, $\mathrm{C}_{\mathrm{D}}$ w-v, for different
trip-dot heights. Since the free-transition drag corrected for laminar run requires no further
corrections, this is the experimental fully turbulent drag (which is assumed to be the "true" drag of
the model). The trip drag for each dot height, $k$, can now be estimated as the difference between
the trips-on data corrected for laminar run and the free-transition data that has been adjusted for
laminar flow, as illustrated in the figure. Shown here are trip drag corrections for the TCA W/B,
as obtained from CFL3D computations with transition locations defined from UPWT 1679 data at
cruise conditions.
Trip Drag
High Speed Aerodynamics, Long Beach

- Laminar run corrections applied to wind-tunnel data for each dot height, $k$ :
- Corrected free-transition data, $C_{D \text { tur }}(k=0)$, assumed to be "true" drag of model
- Trip drag: $C_{D \text { trip }}=C_{D \text { tur }}(k>0)-C_{D \text { tur }}(k=0)$

- Laminar run $C_{D \text { tur }}=C_{D w-t}+C_{D \text { cor }}$
Trip Drag for TCA and NCV W/B The TCA and NCV trip drag for different dot heights is shown in this figure. Both, CFD
results and flat-plate estimates are presented. It can be seen that the CFL3D trip drag
computations are about 0.3 to 0.6 counts higher than the flat-plate estimates. For a typical trip
dot height of 0.012 ", the trip drag estimates for the TCA and NCV configurations are
approximately 2.2 and 1.6 counts, respectively (based on the laminar-flow correction from the
CFL3D solutions). Note that the trip dots on the TCA model are located 0.6 " from the wing
leading edge, whereas the dots on the NCV wing are 1" downstream of the leading edge.
Previous trip drag estimates obtained using techniques such as the variable trip height and the
variable Reynolds number methods were as high as 5 counts.
Trip Drag for TCA and NCV W/B


1703) $\forall O \perp$ әцł ฝəへо before and after corrections have been applied to the data. The experimental drag values
 baseline trip configuration ( $k=0.012 " @ x=0.6$ " from the leading edge) on both models. The details on how the corrected data were obtained are also given in the figure. At the given

 NCV W/B configurations with $k=0.012$ " is approximately 2.2 counts. Therefore, since the corrections for the NCV data are about the same as the corrections applied to the TCA data, the

 decreases by 0.7 when the NCV as-built model is analyzed, the difference between CFD predictions and wind-tunnel measurements now becomes only 0.5 counts (in terms of TCA-toNCV drag reduction).

Recall that the wind-tunnel data corrections for the TCA were derived using data from
UPWT Test 1679, while the corrections for the NCV data were obtained from UPWT Test 1703. However, as long as the conditions remain the same, it is assumed that data corrections derived from one test can be applied to data from another test.
NCV Drag Reduction relative to TCA for W/B Configurations
High Speed Aerodynamics, Long Beach
CFD (CFL3D N-S, B-L) vs. Experiment (UPWT 1703, $\mathrm{k}=0.012^{\prime \prime}$ at $\mathrm{x}=0.6^{\prime \prime}$ )
$C_{L}=0.09$

Data Corrections: TCA-to-NCV Drag Reduction (UPWT 1687) Experimental data from UPTW Test 1687 with the baseline trip configuration are presented
in this chart. Since the conditions are the same as for the previous set of data (UPWT 1703), the
data corrections are assumed to be the same. Hence, again, the measured NCV drag
improvement is essentially the-same before and after corrections are applied. However, the NCV
drag improvement observed during this wind-tunnel test is 0.34 counts higher than the
improvement measured during UPWT Test 1703 . Therefore, the agreement between the
CFL3D-predicted improvement (corrected for NCV model imprecisions) and data from UPTW
Test 1687 reduces to just over 0.1 drag counts.
NCV Drag Reduction relative to TCA for W/B Configurations
CFD (CFL3D N-S, B-L) vs. Experiment (UPWT 1687, $\mathrm{k}=0.012^{\prime \prime}$ at $\mathrm{x}=0.6^{\prime \prime}$ )

Summary and Conclusions While a single definite reason for the NCV performance shortfall was not found, several factors have been identified as possible culprits and others have been ruled out.
Factors that have been found to have no impact on the discrepancies between the predicted and the measured performance improvement of the NCV W/B over the TCA W/B are the following:

- Laminar flow separation on the NCV forebody: On-body flow visualization data obtained
 the flow over the NCV forebody is fully turbulent (as modeled in most CFD simulations), ruling out
 shortfall.
Grid density effects on CFD solutions: The effect of grid density on the CFD drag predictions for both models was investigated by increasing the number of points in the computational grid from 1.5 million to 11.5 million. The refined-grid Navier-Stokes solutions performed at $R e_{c}=6.36 \times 10^{6}$ did not significantly change the original predictions at $R e_{c}=6.36 \times 10^{6}$.
- Aeroelastic effects on drag: Results from Navier-Stokes analyses of the TCA W/B/N/D
model with aeroelastic deformations indicate that static aeroelasticity affects lift and pitching moment, but it does not have a significant impact on the drag polar of the TCA configuration.
Summary and Conclusions
Beach
(Cont'd)
Conclusions Summary and
 TCA-to-NCV drag improvement are the following:
- Model fabrication imprecision: The impact of model fidelity on drag predictions was
examined by performing Navier-Stokes analyses of the NCV QA model. It was found that the
predicted drag of the as-built geometry is approximately 0.7 counts higher than the drag of the as-
designed W/B configuration at cruise. The TCA as-built model will have to be analyzed in order
to assess the full impact of model fidelity on the CFD predictions.








 CFD predictions and wind-tunnel data is just over 0.1 count.
Shock on the NCV wing upper surface: Off-body flow visualization data obtained from laser vapor screen testing on the NCV, detected the presence of a shock on the upper surface of
 shock. It is possible that this shock could adversely affect the NCV performance.
Summary and Conclusions (Cont'd)
- Partial contributors
- Model fabrication imprecision
- NCV QA has 0.7 cts. more drag than theoretical design
- TCA QA must be analyzed to assess full impact
- Laminar flow and trip drag
- Laminar run corrections applied to free-transition data decreas
measured TCA-to-NCV drag improvement by 0.6 cts.
- For baseline HSR trip configuration $\left(\mathrm{k}=0.01 \mathrm{~m}^{\prime \prime}\right.$ @ $\mathrm{x}=0.6$ "),
laminar run and trip drag corrections do not change measured
improvement significantly
- Shock on NCV wing upper surface
- Observed in wind tunnel
- Not predicted by CFD

Two important unresolved issues remain:
- Uncertainty of wind-tunnel measurements: It has been suggested that the uncertainty in
wind-tunnel measurements is as big as 2 or even 2.5 counts. If this were the case, the wind-
tunnel drag uncertainty band would be large enough to include the CFD predictions. But an
extensive test techniques program to reduce such a large uncertainty would need to be
conducted.
- Accuracy of CFD predictions: The accuracy of the CFD computations has not been
quantified. Perhaps the CFD tools used in the HSR program cannot accurately predict wind-
tunnel flows over configurations like the TCA or NCV. In this case, a test program for
validation/calibration of CFD tools would be appropriate.
$\xrightarrow{\text { Summary and Conclusions (Cont'd) }}$
Unresolved issues
- Uncertainty of wind-tunnel measurements
- Accuracy of CFD predictions
TUNNEL
FTBODY CLOSURE WIND-
NG, DATA ANALYSIS, AND
MPUTATIONAL RESULTS

Jerry Allen, LaRC
Grant Martin, BPW
Paul Kubiatko, BPW

Configuration Aerodynamics

Supersonic Aftbody Closure Wind-Tunnel
This paper reports on the model, test, and results from the Langley Supersonic Aftbody Closure wind tunnel test. Each of the authors prepared a specific section of the paper. Jerry Allen from Langley prepared the opening section describing the test and model. Grant Martin from Boeing Phantom Works prepared the section on the performance results. Paul Kubiatko, also from Boeing Phantom Works, prepared the section on the stability and control results. Each author presented his section during the oral presentation of the paper.
Configuration Aerodynamics Technology Development
This presentation was the fourth paper in the Testing Session (Session 4) held on Thursday, February 11,
1999. The work included in this paper supports the three activities that are shaded in this chart.
Configuration Aerodynamics Technology Development
Session 4: Testing Session

## Demonstrate Significant L/Dmax Gains


Goals
Objectives
Challenges
Approaches

Design


| Tech. Baseline |
| :---: |
| Development | T

WT Data Corrections



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Multi-Point Optimization
S\&C CFD Predictions
Program
Project Background:
This project is an experimental evaluation of the $1.5 \%$ Technology Concept Aircraft (TCA) aftbody closure
model (Model 23) in the Langley Unitary Plan Wind Tunnel. The baseline TCA design is the result of a
multidisciplinary, multipoint optimization process and was developed using linear design and analysis methods,
supplemented with Euler and Navier-Stokes numerical methods. After a thorough design review, it was decided to
use an upswept blade attached to the forebody as the mounting system. Structural concerns dictated that a wingtip
support system would not be feasible. Only the aftbody part of the model is metric. The metric break was chosen
to be at the fuselage station where prior aft-sting supported models had been truncated. Model 23 is thus a
modified version of Model 20. The wing strongback, flap parts, and nacelles from Model 20 were used, whereas
new aftbodies, a common forebody, and some new tails were fabricated. Microcraft of Hampton was tasked with
the design and fabrication of the new hardware.

[^29]

Test Objectives
The six objectives of this test listed in this figure can be grouped into the following four areas: 1) determine and tail control effectiveness, 3) provide a database for CFD code validation and correlation with other models, and 4) establish and validate the aftbody closure test technique at supersonic speeds.
Test Support NASA Langley served as the lead aero organization for this test, and was provided on-site support by
Boeing-Phantom Works (BPW) and Boeing-Seattle (BCAG). Jerry Allen served as focal point for NASA Langley
and was on-site for the 7.5-week duration of the test. Support from Boeing was provided by David Blake (BPW),
Paul Kubiatko (BPW), Grant Martin (BPW), Christine Titzer (BCAG), Greg Stanislaw (BCAG), and Servando
Flores (BCAG). David Tuttle of NASA Langley served as the lead test engineer throughout the test. The UPWT
technicians listed in this figure made model changes and kept the tunnel running.


- Researchers:
Comparison of Aftbody Configurations
This figure shows comparison drawings of the three aftbody configurations used in this test - flared,
baseline closed and modified closed. Also indicated are which parts were new for this test and which came from
the preexisting sting-mounted model (Model 20). Basically, all three configurations used the same forebody, wing
section, and nacelles, but had separate aftbody transition sections and metric aftbodies. The flared configuration
used the existing tails from Model 20 , whereas the closed configurations needed new tails to match their aftbody
shapes. All three configurations used the same internal balance to measure force and moment data. There were
actually two models of the baseline aftbody. In addition to the force and moment model, a separate pressure-
instrumented model was built. This model had 47 pressure taps located at 12 longitudinal stations on the aftbody
to measure detailed pressure distributions on this configuration.



Model Installation in UPWT
A side view drawing of the UPWT test setup of the flared aftbody configuration is shown in this figure. The tems outlined in red are part of the permanent model support equipment in the test section. This figure shows the forward blade attachment sting that was used for all aftbody configurations to support the model throughout this test. For some of the flared aftbody runs a dummy sting (not shown in this figure) was used to simulate the aftsting mounted Model 20 setup.
Flared Aftbody Configuration in UPWT
A side-view photograph of the flared aftbody configuration with dummy sting attached is shown in this
figure. The front end of the dummy sting is located inside the base cavity but does not make contact with the
model so that the aftbody balance measurements would not be compromised. As shown in this figure, the dummy
sting is held in place by attachment of its downstream end to the permanent model support equipment of the tunnel.

Dummy Sting for Flared Aftbody
 sting on the forebody part of the model. Because it did not rotate with the model, the dummy sting could be used
only at zero sideslip angle. This figure shows photographs of the flared aftbody configuration with and without the
dummy sting attached. Note the pressure tube running along the dummy sting into the base cavity of the model.
For the majority of the test this was the only method used to measure the pressures inside the cavity. Thus base
cavity pressures were not measured on runs where the dummy sting was not used. This deficiency was addressed
toward the end of the test, as will be discussed in a later figure.

Baseline Closed Tail
A photo of the baseline closed aftbody is shown in this figure. Since the closed configurations had no base
cavity, there was never any need to use the dummy sting with these configurations. This figure also highlights the
metric break which isolated the aftbodies from the rest of the configuration. The white area seen in the gap at the
metric break is a very thin piece of Teflon tape that was inserted in groves that were cut in the facing surfaces.
This tape was designed to inhibit airflow into the metric break area while not fouling the balance measurements.
This tape was used on all configurations for force and moment runs.

Comparison of Aftbodies
Side-by-side photographs of the three aftbody configurations are shown in this figure in both top and side
views. No tails are attached on either aftbody; however, the modified aftbody is in its tails-off running
configuration - a vertical-tail-off block is attached and the horizontal tail mounting holes have been filled. The
basic purpose of the modified aftbody was to provide larger flat surfaces on the sides of the aftbody to prevent
unporting of the horizontal tails. The differences in the two closed aftbodies seems small from the top view of the
aftbodies in this figure, but is clearly seen in the side view.
Comparison of Aftbodies
Top View Side View



Baseline

Comparison of Baseline and Modified Aftbodies
The unporting of the horizontal tails on the two closed aftbodies is illustrated in the drawings in this figure,
where the tails are shown at +4 deg. incidence. The top-view drawings at the bottom of this figure show the gap
that is created between the tail root chord and the baseline aftbody fuselage. The size of this gap increases with tail
incidence angle due to the curvature of the aftbody. The gap is been to be virtually eliminated for the modified aftbody due to the larger flat surface on the fuselage.

Comparison of Baseline and Modified Aftbody
Horizontal Tail $\delta_{\mathrm{HT}}=4^{\circ}\left(2^{\circ} \mathrm{FRP}\right)$ $\checkmark$



Baseline


Modified
Top View of Hor. Tail LE (zoomed)


Modified
Top View of Hor. Tail TE
Baseline




## Q ETDEINE

Aftbody Streamwise Sections
This figure shows longitudinal cuts at the centerline of both the flared and closed aftbodies to show how the balance was attached to the metric and non-metric parts of the model. Note that the drawing is not to scale downstream of the metric break (the base cavity on the flared aftbody is much deeper than it appears here). There are three distinct internal cavities on this model on which pressures could act. One is from the nose of the model to the upstream face of the balance mounting block. The second is in the metric break region. And the third is in the base cavity region (flared aftbody only). The first cavity lies entirely in the nonmetric part of the configuration; therefore the pressures in this cavity have no effect on the balance measurements. Pressures on the later two cavities, however, act on the metric aftbody and therefore have to be accounted for when correcting the measured drag. In this test the pressure in the metric-break cavity was measured for all runs; however, the base cavity pressures on the flared aftbody were measured initially only when the dummy sting was used. This created an issue that had to be resolved to get accurate drag data on the flared aftbody. This issue was addressed by having several runs at the end of the test that were dedicated to measuring the base cavity pressures without the dummy sting. This was done as shown in the next figure.
Aftbody Streamwise Sections

Measuring Base Cavity Pressures without Sting
This figure shows a rear view of the flared aftbody configuration for the dedicated base pressure runs. Four
pressure tubes were run down the outside of the fuselage to the base of the model and routed into the base cavity.
Since these tubes bridge the metric gap, valid balance data were not recorded for these runs. Base cavity pressures
were thus measured without the dummy sting being present.
Cavity Pressures without Sting
Measuring Base

Internal Cavity Pressures
A plot of the measured pressures in the internal cavities of the flared aftbody at zero sideslip angle are shown
in this figure. The pressures in the balance (metric break) cavity are near freestream static pressure, and thus have little effect on the measured drag. The pressures in the base cavity, however, are about $50 \%$ lower than freestream static. The corrections needed for these pressures had a large effect on the final drag data. In addition, this figure shows that there were small but measurable effects of the dummy sting on the base cavity pressures. The most ideal way of obtaining accurate drag on the flared aftbody configuration without the dummy sting, therefore, would be to measure base cavity pressures simultaneously with the balance measurements for all test conditions. Modifications to the model have been made to do this and will be incorporated into the upcoming test on this model at the 16-FT Transonic Tunnel at Langley.

Test Highlights
A summary of the test conditions are listed in this figure. The test was conducted in June of 1998 in the high

[^30]Performance Results
The performance portion of this report will be presented in the order listed above.

Performance Results

W/B Configuration
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$$
\begin{aligned}
& \text { The following charts document the repeatability of the } 1.5 \% \text { TCA Model } 23 \text { with the baseline aftbody. } \\
& \text { Lift versus angle-of-attack and drag versus angle-of-attack are presented. Throughout the range of angles-of- } \\
& \text { attack, the lift is repeatable within } \pm 0.0001 \text {. The drag is generally repeatable within } \pm 0.2 \text { counts at a given } \\
& \text { angle-of-attack. Near the cruise condition, the drag repeatability improves to within } \pm 0.1 \text { count, which is } \\
& \text { equivalent to } 2.5 \% \text { of the total drag on the aftbody. }
\end{aligned}
$$




The next few charts present a build-up for the baseline aftbody. Six configurations are shown: W/B,
$\mathrm{W} / \mathrm{B} / \mathrm{N} / \mathrm{D}, \mathrm{W} / \mathrm{B} / \mathrm{N} / \mathrm{D} / \mathrm{V}, \mathrm{W} / \mathrm{B} / \mathrm{N} / \mathrm{D} / \mathrm{H}, \mathrm{W} / \mathrm{B} / \mathrm{V} / \mathrm{H}$, and $\mathrm{W} / \mathrm{B} / \mathrm{N} / \mathrm{D} / \mathrm{V} / \mathrm{H}$. It is evident that without the horizontal
stabilizers, the aftbody generates little or no lift. Adding nacelles to the W/B/E configuration increases the
aftbody lift-curve slope, while adding the nacelles to the W/B configuration shows little sensitivity. The
$\mathrm{W} / \mathrm{B} / \mathrm{N} / \mathrm{D} / \mathrm{H}$ configuration generates the most lift on the aftbody at $\alpha=3.5^{\circ}$. The addition of the vertical
stabilizer at $\alpha=3.5^{\circ}$, reduces the $\mathrm{C}_{\mathrm{L}}$ by 0.0002 and 0.0003 for the $\mathrm{W} / \mathrm{B} / \mathrm{N} / \mathrm{D}$ and the $\mathrm{W} / \mathrm{B} / \mathrm{N} / \mathrm{D} / \mathrm{H}$
configurations, respectively.
TCA Baseline Aftbody LIft

Angle-of-Attack, $\alpha\left({ }^{\circ}\right)$
TCA Basel ine Aftbody Drag
1.58 TCA Model 23

The aftbody drag for the six baseline configurations tested are presented below. At $\alpha=3.5^{\circ}$, the aftbody drag for the W/B and W/B/N/D configurations are approximately 4.0 counts. Addition of the nacelles
on the W/B/E increases drag by 0.7 counts at $\alpha=3.5^{\circ}$. The addition of the vertical stabilizer increases the drag by 4.2 and 4.4 counts for the $\mathrm{W} / \mathrm{B} / \mathrm{N} / \mathrm{D}$ and $\mathrm{W} / \mathrm{B} / \mathrm{N} / \mathrm{D} / \mathrm{H}$ configurations, respectively. The horizontal stabilizers increase the aftbody drag by 8.6 and 8.8 counts for the W/B/N/D and W/B/N/D/V configurations, respectively. The minimum aftbody drag occurs at $\alpha \cong 1.5^{\circ}$ for all the horizontal stab-on configurations.
 to the total aftbody drag.
TCA Baseline Aftbody Drag

Pitching Moment
Kpoq7ゴ
TCA Baseline
All configurations without the horizontal stabilizers have a small negative (nose-down) pitching
moment, while the horizontal stab-on runs show a stable aftbody up to $9^{\circ}$ angle-of-attack. The addition of the
nacelles on the $\mathrm{W} / \mathrm{B} / \mathrm{E}$ decreases the $\mathrm{C}_{\mathrm{M}}$ slope producing a slightly more stable aftbody. The vertical
stabilizer increases the pitching moment for both the W/B/N/D and W/B/N/D/H configurations.

Comparison of Baseline and Flared Aftbody Lift
1.5\% TCA Model $23, \mathrm{~W} / \mathrm{B} / \mathrm{N} / \mathrm{D} / \mathrm{V} / \mathrm{H}$ Configuration, $\mathbf{i}_{\mathrm{H}}=$
The baseline and flared abodes are compared in this section. The lift curve comparison is shown
below. The flared aftbody has a higher lift coefficient throughout the angle-of-attack range.

Comparison of Baseline and Flared Aftbody Drag
$1.5 \% \mathrm{TCA} \operatorname{Model} 23, \mathrm{~W} / \mathrm{B} / \mathrm{N} / \mathrm{D} / \mathrm{V} / \mathrm{H}$ Configuration $i_{H}=0^{\circ}$
At $\alpha=3.5^{\circ}$, the flared aftbody drag is $\sim 2$ counts greater than the baseline. The difference increases
with angle-of-attack.

Comparison of Baseline and Flared Aftbody Pitching Moment
$1.5 \%$ TCA Model $23, \mathrm{~W} / \mathrm{B} / \mathrm{N} / \mathrm{D} / \mathrm{V} / \mathrm{H}$ Configuration $i_{\mathrm{H}}=0^{\circ}$
The baseline aftbody shows a considerable reduction in longitudinal stability compared to the flared aftbody.

aftbody Closure Effects
This section presents the corrected and uncorrected lift curve, drag polar, and pitching moment for the 1SəL) LMA

 closure effect on the lift curve is seen above. The lift-curve slope is reduced slightly. At $\alpha=3.5^{\circ}$, the
 was obtained at $3 \times 10^{6} / \mathrm{ft}$. The increment due to aftbody geometry is shown to significantly reduce the longitudinal stability of the TCA configuration.
TCA Lift Coefficient Corrected for Aftbody Closure Effects $M_{\infty}=2.4,1.5 \%$ TCA Model, W/B/N/D/V/H configuration, $\mathrm{i}_{\mathrm{H}}=0^{\circ}$

$\begin{array}{llllll}7 & 8 & 9 & 10 & 11 & 12\end{array}$ ${ }^{6}\left({ }^{6}\right)$

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$\pm$ -0.100
TCA Drag Coefficient Corrected for Aftbody Closure Effects

$$
\begin{array}{l}1.5 \% \mathrm{TCA} \text { Model, } \mathrm{W} / \mathrm{B} / \mathrm{N} / \mathrm{D} / \mathrm{V} / \mathrm{H} \text { configuration, } \mathrm{i}_{\mathrm{H}}=0^{\circ} \\ \mathrm{M}_{\infty}=2.4, \mathrm{Re}=4 \times 10^{6} / \mathrm{ft}\end{array}
$$


Angle-of-Attack, $\alpha\left({ }^{\circ}\right)$

## 

Effects

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| Angle-of-Attack, | 6 |  |  |
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Wo ‘quemow 6u!upuld
Aftbody Contribution to Full Aircraft Forces at $1.5 \%$ TCA
Model 23, Baseline Aftbody on $W / B / N / D / V / H$ Configuration $i_{H}=0$. The baseline aftbody drag and lift contribution to the full configuration is $10 \%$ and $3 \%$, respectively. The full
configuration forces were obtained from a previous sting-mounted test at the UPWT (Test 1677) with the
aftbody closure effects obtained in Test 1707 corrected to the data.
Aftbody Contribution to Full Aircraft Forces $\left(\alpha=3.5^{\circ}\right)$

$$
M_{\infty}=2.4, \beta=0^{\circ}, \mathrm{I}_{\mathrm{H}}=0^{\circ}, \text { Re }=4 \times 10^{6} / \mathrm{ft}
$$

(Model 23) $\quad \begin{gathered}\text { Full Aircraft }(\mathrm{W} / \mathrm{B} / \mathrm{N} / \mathrm{D} / \mathrm{V} / \mathrm{H}) \\ (\text { Model } 20+\text { Model } 23)\end{gathered}$

Comparison of Baseline and Modified Aftbody Lift
$1.5 \%$ TCA Model $23, \mathrm{~W} / \mathrm{B} / \mathrm{N} / \mathrm{D} / \mathrm{V} / \mathrm{H}$ Configuration $\mathrm{i}_{\mathrm{H}}=0$.
The baseline and modified afbodies are compared in this section. The lift curve comparison is shown
below. The modified aftbody has a higher lift coefficient throughout the angle-of-attack range. At $\alpha=3.5^{\circ}$, the
modified aftbody $\mathrm{C}_{\mathrm{L}}$ is 0.0005 higher than the baseline aftbody.



 negative (nose-down) CM.
Pitching Moment from the TCA Baseline and Modified Aftbodies

Trim drag is the drag of the aircraft at $\mathrm{C}_{\mathrm{M}}=0$ minus the drag of the aircraft where the horizontal tail has zero lift. During Test 1707, trim drag studies were planned for both the baseline and modified aftbodies. The baseline aftbody trim study was dropped from the run schedule due to time limitations in the wind tunnel. The chart above shows the process of determining trim drag for the modified aftbody configuration. The full configuration force and moments were obtained from a previous sting-mounted test at the UPWT (Test 1677) using the flared aftbody geometry. Differences between the flared and modified aftbodies were obtained and applied to the existing data. The trim drag was determined to be -0.4 counts for the modified aftbody.

Flow Solver and Computational Grid
The nonlinear flow solver, CFL3D, was used in simulations of flows over the TCA configuration. A
parallel version of CFL3D was used on the SGI Origin 2000. CFL3D uses an upwind Roe's flux-differencing
scheme to solve either the Euler or Navier-Stokes equations on patched, structured, multi-block grids. A
number of convergence accelerators are programmed in CFL3D to reduce CPU time. Here, grid sequencing
and multigrid were actively used. The Baldwin-Lomax turbulence model was used in all the calculations.
Flow Solver \& Computational Grid - CFL3D

- Baldwin-Lomax turbulence model
- Parallel version
- SGI Origin 2000
- TCA supersonic W/B/N/D/V/H grid
- 21 blocks
- 6.7 milion points
- $1: 1$ and patched block boundaries
- Strut not modeled in CFD

Predicted and measured drag on the modified aftbody for the $1.5 \%$ Model 23 at $\mathrm{M}_{\infty}=2.4, \alpha=3.5^{\circ}$, and
$\mathrm{Re}=4 \times 10^{6} / \mathrm{ft}$. CFL3D predicts 0.2 counts less aftbody drag than the wind tunnel test. Although the test data
has not been corrected for the trip drag, the correction is expected to be small compared to the total aftbody
drag. The CFD solutions do not model the forward swept strut used to support the wind-tunnel model in the
test section. The CFL3D N-S pressure distribution and surface streamlines in the aftbody region are also
shown. CFL3D N-S does not predict separation in the aftbody region. CFL3D N-S and test data show that
the modified aftbody has $\sim 0.5$ counts more drag than the baseline aftbody.

Comparison of CFL3D $\mathrm{N}-\mathrm{S}$ and Test Data Pressure Distributions
$1.5 \%$ TCA Model 23 , Baseline Aftbody
Predicted and measured pressures are compared for the baseline aftbody The CFL3D N-S surface
pressures agree well with the test data for WL 4.086, 4.586, 5.086, and the upper centerline station. The
comparison is not as good for the lower centerline pressures starting at MS 54 .

Comparison of CFL3D N-S and Test Data surface Streamlines

$$
1.5 \% \text { TCA Model 23, Baseline Aftbody }
$$

CFL3D N-S surface streamlines and UV oil are compared in the above chart for the W/B/N/D/V/H
configuration. Overall, there is good agreement between the computational and experimental surface flow
Aftbody guration
CA Baseline onfig $/{ }^{\circ}, i_{4}=0^{\circ}$ /ft, $\alpha=3.5^{\circ}$,
$\stackrel{0}{4}$
23, W/
Particle Traces on
1.5\% TCA Model 23
$M_{\infty}=2.4, \mathrm{Re}$


UV Oil

CFL3D N-S (B-L)
Summary of Performance Results
The test data exhibits excellent short-term repeatability. Overall, there was good agreement between
the predicted CFL3D forces and the test data for the 1.5\% TCA Model 23 . This agreement is extremely
important in order to validate the analysis tools used for the complex configurations. The trim drag was
determined for the modified aftbody configuration. The modified aftbody allowed the horizontal tail to remain
ported at higher raiil incidence angles compared to the baseline aftbody. The UV oil images and CFL3D N-S
surface streamlines showed no signs of separation on the aftbody. The wind-tunnel test established aftbody
closure test techniques.
Summary of Performance Results (cont.) - Both UV oil images and CFD results indicate attached
flow on the aftbody near supersonic cruise

- Established aftbody closure test techniques at
supersonic speeds
Stability and Control Test Objectives
This slide outlines the stability and control objectives of the aftbody closure test. First the longitudinal and directional stability levels of the flared and baseline closed aftbody will be determined so that an increment between the two can be derived. Similarly, the longitudinal and directional control effectiveness for the horizontal stabilizer, elevator, and rudder will be
 between the two can be derived.
These stability and control increments can be compared against those predicted by CFD codes, such as A502SC, to gage the ability of those codes to estimate aftbody closure effects.
 control database obtained for the TCA configuration with the flared aftbody. This corrected database can then be used for development of a non-linear aerodynamic simulation of the TCA configuration.
Stability and Control Test Objectives
- Validate Aftbody Longitudinal and Directional
$\frac{\text { Stability Levels: }}{\Delta \mathrm{Cm}_{\alpha}, \Delta \mathrm{Cn}_{\beta}}$
$(\Delta=$ Baseline Closed Aftbody - Flared Aftbody $)$
- Validate Control Effectiveness Levels:
$\Delta \mathrm{Cm}_{\mathrm{iH}}, \Delta \mathrm{Cm}_{\delta \mathrm{e}}, \Delta \mathrm{Cn}_{\delta \mathrm{r}}$
$(\Delta=$ Baseline Closed Aftbody - Flared Aftbody)
- Provide Database for CFD Code Validation and

Corrections to the TCA Aerodynamic Database:
$\mathrm{C}_{\mathrm{x}_{\text {corrected }}}=\mathrm{C}_{\mathrm{x}_{\text {uncorrected }}}+\Delta \mathrm{C}_{\mathrm{x}}$
(where x is any force or moment variable)

## Longitudinal Characteristics <br> Stability Comparison (ABB vs. ABF)

configuration.

TCA Pitching Moment Coefficient
Corrected for Aftbody Closure Effects
This figure shows the corrected and uncorrected pitching moment coefficient variation with
angle-of-attack for the TCA configuration at Mach 2.4. The uncorrected data was obtained
from a previous sting-mounted test at the UPWT (test \#1677) using the flared aftbody
geometry. Horizontal tail-on and off increments due to the difference between the flared and
closed aftbody geometries were obtained from the previous slide and applied to the
uncorrected data to obtain the corrected pitching moment curves shown here. The increment
due to aftbody geometry is shown to significantly reduce the horizontal tail-on longitudinal
stability of the TCA configuration.

 This slide shows the variation of pitching moment coefficient vith angle-of-attack and Horizontal stabilizer deflections tested include $0,+6$, and -6 degrees. Stabilizer effectiveness (incremental pitching moment per degree of stabilizer deflection) is fairly constant with angle-of-attack for either aftbody geometry; however, the stabilizer effectivenes for the ABB configuration is reduced relative to the ABF configuration. Again, the primary reason for the
 effectiveness for either aftbody geometry is shown to be fairly constant between positive and negative deflections.
Tails-On Longitudinal Characteristics
Stabilizer Control Effectiveness
(ABB vs. ABF)

Angle-of-Attack, $\alpha$ (deg.)

Tails-On Longitudinal Characteristics
Elevator Control Effectiveness
(ABB vs. ABF) This slide shows the variation of pitching moment coefficient vith angle-of-attack and
elevator deflection, at Mach 2.4 , for the ABB and ABF aftbody geometries. Elevator
deflections tested include $0,+10$, and +20 degrees. Elevator effectiveness (incremental
pitching moment per degree of elevator deflection) is fairly constant with angle-of-attack for
both aftbody geometries. Also, there is only a small difference in elevator effectiveness
between the two aftbody geometries. This is in contrast to the stabilizer effectiveness
comparison shown on the previous slide.
Tails-On Longitudinal Characteristics
Elevator Control Effectiveness
(ABB vs. ABF)

Angle-of-Attack, $\alpha$ (deg.)

Longitudinal Characteristics
Stability and Control Derivatives
(ABB vs. ABF)

$$
\begin{aligned}
& \text { This slide summarizes the longitudinal stability and control derivative comparison between } \\
& \text { the flared and baseline closed aftbody geometries at a representative angle-of-attack of } 3.5 \\
& \text { degrees. The significant reduction in tail-on stability, the moderate reduction in stabilizer } \\
& \text { effectiveness, and the small reduction in elevator effectiveness are easily observed for the } \\
& \text { baseline closed aftbody when compared to the flared aftbody. Stabilizer effectiveness is } \\
& \text { shown to be symmetric between positive and negative deflections for both aftbodies. Also, } \\
& \text { elevator effectiveness remains constant up to +20 degrees of deflection for both aftbodies. }
\end{aligned}
$$


Tails-On Directional Characteristics
Stability Comparison (ABB vs. ABF)

 yawing moment at any other $\beta$ and dividing the resulting yawing moment by the $\beta$. The data shows that the ABB aftbody, has significantly less directional stability than the ABF aftbody. This reduction in directional stability for the ABB aftbody is the result of a reduction in both the aftbody side area, as well as, the reduction in vertical tail carry-over lift.
TCA Yawing Moment Coefficient
Corrected for Aftbody Closure Effects
This figure shows the corrected and uncorrected yawing moment coefficient variation with
 uncorrected data was obtained from a previous sting-mounted test at the UPWT (test \#1677) using the flared aftbody geometry. Vertical tail-on and off increments due to the difference between the flared and closed aftbody geometries were obtained from the UPWT-1707 test
and applied to the uncorrected data to obtain the corrected yawing moment curves shown
 vertical tail-off directional stability of the TCA configuration, although the reduction for the vertical tail-on is more significant.


 aftbody geometries at several sideslip angles and at a representative angle-of-attack of 3.5 degrees. The moderate reduction in directional stability of the ABB aftbody compared to the ABF aftbody is shown at each of the sideslip angles tested. Also shown is the relative asymetry in measured directional stability between the $\beta=+3$ and $\beta=-3$ degrees. This is believed to be caused by asymmetries associated with the model and tunnel sideflow which was not determined during the test.
-On Directional Characteristics
(ABB vs. ABF)
(ABility Derivative Comparison
vitectionat staty Aftbody (ABB) Configurations $\mathrm{MaCh}=2.4,=3.5^{\circ}$, Re/ft $=\mathbf{4} \times 10^{6}$

Tails-On Directional Characteristics
Rudder Control Effectiveness at $\beta=0^{\circ}$
(ABB vs. ABF) The slide below shows the variation of yawing moment coefficient with angle-of-attack and rudder deflection ,at Mach 2.4, for the ABB and ABF aftbody geometries at a constant $\beta$ of 0 degrees. Rudder deflections tested include $0,+10$, and +20 degrees. Rudder effectiveness (incremental yawing moment per degree of rudder deflection) is fairly constant with angle-ofattack for either aftbody geometry; however, the rudder effectivenes for the ABB configuration is reduced relative to the ABF configuration. Again, the primary reason for the effectiveness reduction is the smaller carry-over area of the ABB configuration. The effectiveness reduction is most dramatic at the +20 degree rudder deflection.
Tails-On Directional Characteristics
Rudder Control Effectiveness at $\beta=0^{\circ}$
(ABB vs. ABF)

Angle-of-Attack, $\alpha$ (deg.)
Tails-On Directional Characteristics
Rudder Control Effectiveness at $\beta=-3^{\circ}$
(ABB vs. ABF) The slide below shows the variation of yawing moment coefficient with angle-of-attack and
rudder deflection ,at Mach 2.4, for the ABB and ABF aftbody geometries at a constant $\beta$ of -3
degrees. For positive rudder deflections, a negative sideslip angle will result in an increased
(more negative) yawing moment level due to positive directional stability. Rudder
effectiveness for the ABF aftbody remained fairly constant when compared to the level at $\beta=$
0 ; however, rudder effectiveness of the ABB aftbody increased significantly but still remained
below the level of the ABF aftbody.
Tails-On Directional Characteristics
Rudder Control Effectiveness at $\beta=-3^{\circ}$
(ABB vs. ABF)

$\alpha$ (deg.)

$$
\begin{aligned}
& \text { Tails-On Directional Characteristics } \\
& \text { Rudder Control Effectiveness at } \beta=+3^{\circ} \\
& \text { (ABB vs. ABF) }
\end{aligned}
$$

$$
\begin{aligned}
& \text { The slide below shows the variation of yawing moment coefficient with angle-of-attack and } \\
& \text { rudder deflection,at Mach } 2.4 \text {, for the ABB and ABF aftbody geometries at a constant } \beta \text { of }+3 \\
& \text { degrees. For positive rudder deflections, a positive sideslip angle will result in a decrease } \\
& \text { (more positive) yawing moment level due to positive directional stability. Rudder effectiveness } \\
& \text { for both the ABF and ABB aftbodies remained fairly constant when compared to the level at } \beta \\
& =0 \text { except the }+20 \text { degree rudder deflection on the ABF aftbody which appears to have less } \\
& \text { effectiveness at } \beta=+3 \text { than at } \beta=0 \text { degrees. }
\end{aligned}
$$



This slide summarizes the rudder effectiveness comparison between the ABB and ABF
aftbody geometries at several sideslip angles and at a representative angle-of-attack of 3.5
degrees. The significantly reduced rudder effectiveness of the ABB aftbody configuration at $\beta$
$=0$ degrees and the significant increase at $\beta=-3$ degrees are easily observed.
Stability and Control Summary of Results In summary, significant differences in longitudinal and direction stability and control
characteristics between the ABF and ABB aftbody geometries were measured. Correcting
the experimental data obtained for the TCA configuration with the flared aftbody to the
representative of the baseline TCA closed aftbody will result in a significant reduction in
longitudinal stability, a moderate reduction in stabilizer effectiveness and directional stability,
and a moderate to significant reduction in rudder effectiveness. These reductions in the
stability and control effectiveness levels of the baseline TCA closed aftbody are attributed to
the reduction in carry-over area.

# Aft Body Closure - Predicted Strut Effects at M=2.4 

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This paper reports the predicted $\mathrm{M}=2.4$ strut-interference effects on a closed aft body with empennage for the TCA baseline model. The strut mounting technique was needed in order to assess the impact of aft-end shaping, i.e. open for a sting or closed to better represent a flight vehicle. However, this technique can potentially lead to unanticipated effects that are measured on the aft body. Therefore, a set of computations were performed in order to examine the closed aft body with and without strut present, at both zero and non-zero angles of sideslip (AOS). The work was divided into a computational task performed by Javier A. Garriz, using an inviscid (Euler) solver, and a monitoring/reporting task done by John E. Lamar. All this work was performed during FY98 at the NASA Langley Research Center.


The work reported in this paper is associated with the highlighted boxes.

## Outline

- CFD Objectives and Geometrical Modeling
- Supersonic Results
- Conclusions
- Recommendations

This presentation has four parts. The first deals with the objectives of the CFD study and details the geometrical modeling employed. The second part shows the resulting pressures and forces/moments obtained at $\mathrm{M}=2.4$ with the strutoff and -on. The last two parts summarize the paper with some conclusions and provides recommendations for further studies.

## CFD Objectives and Geometrical Modeling

- Estimate strut interference on TCA airplane model at $\mathrm{M}=2.4,0.6,0.9$ and 1.2 for AOA $4^{\circ}$ and AOS of $0^{\circ}$ and $-3^{\circ}$ using an Euler solver (USM3Dns in its inviscid mode)
- Complete airplane - tip to tip - and supersonic strut modeled
- Strut added to airplane at AOS of $0^{\circ}$ and $-3^{\circ}$

The general goal of this task is to be able to estimate the strut interference effects on the TCA airplane model at both transonic and supersonic speeds. To date only the $\mathrm{M}=2.4$ computations have been completed at an angle-of-attack $(\mathrm{AOA})=4$ degrees and at $A O S=0$ and -3 degrees. These have been made using the USM3Dns code in its Euler or inviscid mode. An overall description of the code follows:
USM3Dns is a tetrahedral-based Euler and Navier-Stokes flow solver. Spatial discretization is accomplished by a cell-centered finite-volume formulation using an accurate linear reconstruction scheme and upwind flux differencing. Time is advanced by an implicit backward-Euler time-stepping scheme. Flow turbulence effects are modeled by the Spalart-Allmaras one-equation model, which is coupled with a wall function to reduce the number of cells in the sublayer region of the boundary layer.

Since the study involved sideslip, it was necessary to model the entire configuration and not take advantage of symmetry. Even though this was not required at AOS $=0$ degrees, for consistency full configuration modeling was maintained for all computational situations.
In addition, the supersonic strut was modeled, including the region where it abutted the lower surface of the model. All together about 1.3 million cells were used to model the complete airplane and strut combination.

During the computations, the model geometry was maintained at $\mathrm{AOS}=0$ degrees regardless of the onset flow. For AOS $=-3$ degrees the boundary conditions for the flow solution were changed on the model and the relative alignment angle (as viewed from the top) between the model and the strut was altered by moving the strut only the 3 degrees. This was done because the strut surface was represented by a fewer number of cells than the model.

## Perspective Top View of TCA



This figures shows a top view of the complete TCA aircraft model from a $3 / 4$ front perspective. The different colors on the various parts of the model are only to assist in keeping track of the geometrical elements during the grid generation portion of this task.


This figure shows the arrangement of the complete TCA model and the supersonic strut from a top view at an AOS $=0$ degrees. Again the various colors shown on the different surfaces are only used to keep track of the geometrical el ements during the grid creation process.


Side view of TCA model mounted atop the supersonic strut at $A O S=0$ degrees.


This figure is a lower surface view from a $3 / 4$ front perspective of the TCA model mounted atop the supersonic strut at $\mathrm{AOS}=0$ degrees. Some details of the strut system, both in terms of its cross sectional change with longitudinal distance and the face of the strut where attaches to the tunnel support sy stem, can be seen here.

## Alt. Perspective view of TCA UPWT Strut at AOS of $0^{\circ}$



This figure is another lower surface view of the TCA model and supersonic strut system at AOS $=0$ degrees, but this time looking forward from underneath the model and with a $3 / 4$ rear perspective.


This figure shows the arrangement of the complete TCA model and the supersonic strut from a top view at an AOS $=-3$ degrees. Again the various colors shown on the different surfaces are only used to keep track of the geometrical elements during the grid creation process.


This figure is a lower surface view from a $3 / 4$ front perspective of the juncture between the TCA model mounted and the supersonic strut at $\mathrm{AOS}=-3$ degrees. Note that at the top of the strut leading-edge there is an air gap between the strut and the model which does occur because of the 3 degree misaiignment. This gap was modeled.

# Supersonic Results 

- Surface Pressure Contours
- Forces/Moments
- Cross Sections and Pressures

The figures that follow present the $\mathrm{M}=2.4$ aerodynamic results for $\mathrm{AOA}=4$ degrees at AOS $=0$ and -3 degrees with the strut-off and -on. These figures include eight surface pressure contours, four for the upper and four for the lower; the overall forces/moments for the entire configuration, as well as the forces/moments for the aft body alone; and cross-sectional plots of local pressures and associated wing/body geometry at fixed Fuselage Stations. The figure presentation follows this order.


This figure shows the upper surface pressure coefficient contours for the TCA configuration without the supersonic strut at $\mathrm{M}=2.4, \mathrm{AOA}=4$ degrees, and AOS $=0$ degrees. The most noticeable phenomenon is that even at AOS $=0$ degrees the contours are not exactly symmetrical. This is due to the unstructured CFD code (USM3D) modeling the entire configuration and not taking advantage of geometrical symmetry. Consequently, one should expect small lateral side-force and moments at this attitude.

Along the inboard leading-edge and trailing aft over the outboard wing, a region of higher suction pressure is present.


This figure shows the upper surface pressure coefficient contours for the TCA configuration with the supersonic strut at $\mathrm{M}=2.4, \mathrm{AOA}=4$ degrees, and AOS $=0$ degrees. In comparison with the previous figure (strut-off), there are no striking differences seen on the upper surface associated with the strut, as would be expected under these conditions.


This figure shows the lower surface pressure coefficient contours for the TCA configuration without the supersonic strut at $\mathrm{M}=2.4, \mathrm{AOA}=4$ degrees, and AOS $=0$ degrees. Just as for the upper surface, these contours are not exactly symmetrical even at AOS $=0$ degrees. The most interesting features here are the shock-on-shock interactions between the nacelles, the relatively large suction values of pressure coefficient reached inboard of the inner nacelles, due to flow expansion, and the larger positive pressure coefficients near the outerwing leading edge.


This figure shows the lower surface pressure coefficient contours for the TCA configuration with the supersonic strut at $\mathrm{M}=2.4, \mathrm{AOA}=4$ degrees, and AOS $=0$ degrees. In comparison with the previous figure (strut-off), the nacelle shocks and outer-wing leading edge pressure coefficients appear the same. However, there are differences in the pressure coefficients in the streamwise region between the fuselage/wing juncture and the start of the wing outer panel. The presence of the strut is to cause this region of the wing (an inverted Vshape just inboard-and-downstream of the inboard leading-edges) to have a larger positive value than without the strut. Moreover, along the centerline and downstream of the wing/fuselage juncture -- where the strut footprint would be located (see figure 25) -- there is a region of higher suction pressure coefficient indicated on the model. This is unexpected result for a no-flow situation and therefore the integration of the pressure underneath the strut footprint is not performed. Another reason for omitting these force/moment contributions to the reported values is given on later on figure 21.


This figure shows the upper surface pressure coefficient contours for the TCA configuration without the supersonic strut at $\mathrm{M}=2.4, \mathrm{AOA}=4$ degrees, and AOS $=-3$ degrees. The effect of sideslip is clearly noted on the inboard wing where the higher suction region on the windward wing moves inboard and on the leeward wing moves outboard. Moreover, asymmetries are noted in local pressures over the entire length of the model, including higher suctions on the leeward portion of the horizontal tail.


This figure shows the upper surface pressure coefficient contours for the TCA configuration with the supersonic strut at $\mathrm{M}=2.4, \mathrm{AOA}=4$ degrees, and AOS $=-3$ degrees. In comparison with the previous figure (strut-off), there are no striking differences seen on the upper surface associated with the strut, as would be expected under these conditions.


This figure shows the lower surface pressure coefficient contours for the TCA configuration without the supersonic strut at $\mathrm{M}=2.4, \mathrm{AOA}=4$ degrees, and $\mathrm{AOS}=-3$ degrees. In comparison with figure 15 at $\mathrm{AOS}=0$ degree, the asymmetries associated with sideslip are present all along the configuration, but especially noticeable are the changes to the nacelle shocks and the pressures on the lifting surfaces.


This figure shows the lower surface pressure coefficient contours for the TCA configuration with the supersonic strut at $\mathrm{M}=2.4, \mathrm{AOA}=4$ degrees, and AOS $=-3$ degrees. In comparison with figure 19 (strut-off), the nacelle shocks and outer-wing leading edge pressure coefficients appear the same. However, there are differences in the pressure coefficients in the streamwise region between the fuselage/wing juncture and the start of the wing outer panel. In particular, the inverted V -shape noted in figure 16 at $\mathrm{AOS}=0$ degrees is more-or-less retained at sideslip. Hence, the strut tends to have a mitigating effect on the pressures in its immediate vicinity.

TCA Total Force Coefficients Estimated by Euler Solver (USM3D)

| $\mathrm{M}=2.4, \mathrm{AOA}=4^{\circ}$ |  |  |  |  |  |
| :--- | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{AOS}=0^{\circ}$ |  |  |  |  |
|  | Normal | Axial | Side |  |  |
| - Without strut | 0.11392 | 0.00104 | 0.00080 |  |  |
| - With strut | 0.11363 | 0.00145 | -0.00003 |  |  |
|  | AOS $=3^{\circ}$ |  |  |  |  |
|  |  |  |  |  |  |
|  | Normal | Axial | Side |  |  |
|  | 0.11277 | 0.00097 | 0.00035 |  |  |
| - Without strut | 0.11316 | 0.00147 | 0.01514 |  |  |

[Note that the forces/moments reported in figures 21 to 24 for the TCA airplane model do not contain those generated over the area associated with the strut footprint (see figure 25). This is done to keep the comparisons consistent with the minimum exposed model-surface-area, i.e. strut on.]

The strut effect on the normal- and side-force coefficients at both AOS of 0 and -3 degrees is seen to be small, and its effect on the axial force coefficient at either AOS value amounts to a "pressure drag" increase of about $\sim 5 \mathrm{xE}-04$ ( $\sim 5$ drag counts).
(Note that the side force coefficient should be zero at AOS $=0$ degrees and it isn't due to the unstructured grid modeling of the complete configuration.)

TCA Total Moment Coefficients
Estimated by Euler Solver (USM3D)

$$
\mathrm{M}=2.4, \mathrm{AOA}=4^{\circ}
$$

$$
\mathrm{AOS}=0^{\circ}
$$

- Without strut
- With strut

| Pitch | Roll | Yaw |
| :---: | :---: | :---: |
| -0.21010 | -0.00229 | -0.00059 |
| 0.00203 | -0.00001 | 0.0 |

$$
\operatorname{AOS}=3^{\circ}
$$

- Without strut

| Pitch | Roll | Yaw |
| :---: | :---: | :---: |
| -0.02056 | -0.04442 | 0.09141 |
| 0.00155 | 0.00157 | -0.00302 |

At both AOS of 0 and -3 degrees the strut effect on the rolling- and yawingmoment coefficients is seen to be small*; however, there is a destabilizing effect on the pitching moment coefficient to the extent that it goes from stable to slightly unstable.

* Note they should be zero as should all lateral moments values at AOS $=0$ degrees.

TCA Total and Aftbody Force Coefficients Estimated by Euler Solver (USM3D)

$$
\mathrm{M}=2.4, \mathrm{AOA}=4^{\circ}, \mathrm{AOS}=0^{\circ}
$$

Total Configuration

|  | Normal | Axial | Side |
| :--- | :---: | :---: | :---: |
| - Without strut | 0.11392 | 0.00104 | 0.00080 |
| - With strut | 0.11363 | 0.00145 | -0.00003 |

## Aftbody Only

|  | Normal | Axial | Side |
| :--- | :---: | :---: | :---: |
| - Without strut | 0.00075 | 0.00028 | 0.00001 |
| - With strut | 0.00047 | 0.00032 | -0.00003 |

The strut effect at AOS of 0 degrees on the total TCA configuration force coefficients are repeated from figure 21 for comparison with those for the aft body only. For both it is noted that the changes in aft body force coefficients mirror those of the total configuration. In particular, with the strut present there is less normal force and more axial force; moreover the increase in the latter or "pressure drag" amounts to about one-tenth ( $\sim 5 \mathrm{E}-05$ or $\sim 0.5$ drag counts) of the total increase ( $\sim 5 \mathrm{E}-04$ or $\sim 5$ drag counts).


The strut effect at AOS of 0 degrees on the total moments coefficients are repeated from figure 22 for comparison with those for the aft body only. The only consistent strut effect noted between the two sets of data is the slightly destabilizing effect on pitching moment coefficient. The other aft body moment coefficients should all be zero but are not due to the unstructured modeling of the complete configuration.


This figure shows the lower surface of the TCA airplane model, the footprint of the strut at $\mathrm{AOS}=0$ degrees, the moment reference point and the thirteen computational stations where the local pressure coefficients were computed. Not all thirteen locations of pressure data are presented, but enough of them are shown in order to provide understanding as to how the reported forces/moments are developed. Stations 2 to 7 bound the strut footprint on the model.

TCA Configuration Grid Cross-Section and Surface Cp Comparisons at Station 1 $M=2.4, \alpha=4.0^{\circ}, \beta=0.0^{\circ}$
USM3D Inviscid Flow Simulation

[Note for figures 26 to 37 that the view is looking upstream with the right hand side having the positive $y$ values.]
This figure shows the grid cross-section and the USM3Dns inviscid solution for surface pressure coefficients at station $1, \mathrm{M}=2.4, \mathrm{AOA}=4$ degrees and $\mathrm{AOS}=$ 0 degrees. This station was chosen because it is ahead of the strut, hence no effect is expected here and thus it can be used for validating the solutions. The pressure data show a negligible effect of the strut .

TCA Configuration Grid Cross-Section and Surface Cp Comparisons at Station 3
$M=2.4, \alpha=4.0^{\circ}, \beta=0.0^{\circ}$
USM3D Inviscid Flow Simulation


This figure shows the grid cross-section and the USM3Dns inviscid solution for surface pressure coefficients at station $3, \mathrm{M}=2.4, \mathrm{AOA}=4$ degrees and $\mathrm{AOS}=$ 0 degrees. This is the first station on the strut and shows there to be more positive pressure on the lower surface with the strut present, hence more normal force. Since this station is ahead of the moment reference point, it will contribute a more nose-up moment and be destabilizing.


This figure shows the grid cross-section and the USM3Dns inviscid solution for surface pressure coefficients at station $4, \mathrm{M}=2.4, \mathrm{AOA}=4$ degrees and $\mathrm{AOS}=$ 0 degrees. This station also shows there to be more positive pressure on the lower surface with the strut present, hence more normal force. Since this station is al so ahead of the moment reference point, it will contribute a more nose-up moment and be destabilizing.

TCA Configuration
Grid Cross-Section and Surface Cp Comparisons at Station 7
$M=2.4, \alpha=4.0^{\circ}, \beta=0.0^{\circ}$
USM3D Inviscid Flow Simulation



This figure shows the grid cross-section and the USM3Dns inviscid solution for surface pressure coefficients at station $7, \mathrm{M}=2.4, \mathrm{AOA}=4$ degrees and $\mathrm{AOS}=$ 0 degrees. This station shows that with the strut present there are two suction pressure spikes on the lower surface near $y=0$. However, the strut on pressures are more positive outboard of $y=+/-140$, hence the resulting normal force could be slightly less or the same-as the no-strut configuration. Since this station is still ahead of the moment reference point, the strut effect could be slightly stabilizing or produce no change.

TCA Conflguration Grid Cross-Section and Surface Cp Comparisons at Station 9 $M=2.4, \alpha=4.0^{\circ}, \beta=0.0^{\circ}$
USM3D Inviscid Flow Simulation



This figure shows the grid cross-section and the USM3Dns inviscid solution for surface pressure coefficients at station $9, \mathrm{M}=2.4, \mathrm{AOA}=4$ degrees and $\mathrm{AOS}=$ 0 degrees. This station shows that with the strut present there is less positive pressure on the lower surface, hence less normal force. Since this station is behind the moment reference point, it will contribute a small er nose-down moment and therefore be destabilizing.


This figure shows the grid cross-section and the USM3Dns inviscid solution for surface pressure coefficients at station $12, \mathrm{M}=2.4, \mathrm{AOA}=4$ degrees and AOS $=0$ degrees. No discernable differences are noted in the pressure coefficients so no incremental contribution to the normal force or pitching moment from the strut is expected on the empennage.

TCA Configuration Grid Cross-Section and Surface Cp Comparisons at Station 1 $M=2.4, \alpha=4.0^{\circ} \beta=-3.0^{\circ}$ USM3D Inviscid Flow Simulation


Note: Sting should read Strut

This figure shows the grid cross-section and the USM3Dns inviscid solution for surface pressure coefficients at station $1, M=2.4, A O A=4$ degrees and $A O S=$ -3 degrees. This station was chosen because it is ahead of the strut, hence no effect is expected here and thus it can be used for validating the solutions. The pressure data show a negligible effect of the strut .


This figure shows the grid cross-section and the USM3Dns inviscid solution for surface pressure coefficients at station $3, \mathrm{M}=2.4, \mathrm{AOA}=4$ degrees and $\mathrm{AOS}=$ -3 degrees. This is the first station on the strut and shows there to be more positive pressure on the lower surface with the strut present, hence more normal force. Since this station is ahead of the moment reference point, it will contribute a more nose-up moment and be destabilizing.
[Note that there will be small side force coefficient and rolling- and yawingmoment coefficient differences associated with the integration of the two different lateral, delta-pressure-coefficient distributions over the cross-sectional geometry for figures 33 to 37.]


This figure shows the grid cross-section and the USM3Dns inviscid solution for surface pressure coefficients at station $4, \mathrm{M}=2.4, \mathrm{AOA}=4$ degrees and $\mathrm{AOS}=$ -3 degrees. This station also shows there to be more positive pressure on the lower surface with the strut present, hence more normal force. Since this station is also ahead of the moment reference point, it will contribute a more nose-up moment and be destabilizing.

TCA Configuration Grid Cross-Section and Surface Cp Comparisons at Station 7 $M=2.4, \alpha=4.0^{\circ} \beta=-3.0^{\circ}$
USM3D Inviscid Flow Simulation



Note: Sting should read Stru

This figure shows the grid cross-section and the USM3Dns inviscid solution for surface pressure coefficients at station $7, \mathrm{M}=2.4, \mathrm{AOA}=4$ degrees and $\mathrm{AOS}=$ -3 degrees. This station shows that with the strut present there are two suction pressure spikes on the lower surface near $y=0$. However, the strut on pressures are more positive outboard of $y=-100$ and +150 , hence the resulting normal force could be slightly less or the same-as the no-strut configuration. Since this station is still ahead of the moment reference point, the strut effect could be slightly stabilizing or produce no change.

TCA Configuration
Grid Cross-Section and Surface Cp Comparisons at Station 9 $M=2.4, \alpha=4.0^{\circ} \beta=-3.0^{\circ}$
USM3D Inviscid Flow Simulation



Note: Sting should read Strut

This figure shows the grid cross-section and the USM3Dns inviscid solution for surface pressure coefficients at station $9, \mathrm{M}=2.4, \mathrm{AOA}=4$ degrees and $\mathrm{AOS}=$ -3 degrees. This station shows that with the strut present there is less positive pressure on the lower surface, hence less normal force. Since this station is behind the moment reference point, it will contribute a smaller nose-down moment and therefore be destabilizing.

TCA Conflguration
Grid Cross-Section and Surface Cp Comparisons at Station 12
$M=2.4, \alpha=4.0^{\circ} \beta=-3.0^{\circ}$
USM3D Inviscid Flow Simulation



Note: Sting should read Strut

This figure shows the grid cross-section and the USM3Dns inviscid solution for surface pressure coefficients at station $12, \mathrm{M}=2.4, \mathrm{AOA}=4$ degrees and AOS $=-3$ degrees. No discernable differences are noted in the pressure coefficients so no incremental contribution to the normal force or pitching moment from the strut is expected on the empennage.

## Conclusions

Inviscid (Euler) strut effects at $\mathrm{M}=2.4$ on the TCA configuration and aft body are:

- Negligible for normal force coefficient
- Increase pressure drag (axial force) coefficient
- Slightly pitch destabilizing

The USM3Dns code has been used in its inviscid (Euler) mode to provide solutions for the TCA model in free air and in the presence of a support strut at $\mathrm{M}=2.4$ at both Angles Of Sideslip (AOS) of 0 and -3 degrees. The strut effects on the TCA configuration and aft body were determined to be negligible on normal force, to increase the pressure drag by $\sim 5$ counts on the TCA configuration and $\sim 0.5$ counts on the aft body, and be slightly pitch destabilizing on both.

## Recommendations

- Compare Euler and Navier-Stokes solutions over the TCA aft body for particle traces, Cp , and force/moment
- Do transonic solutions of transonic strut and compete airplane prior to testing in 16FTT in FY99

It is recommended that viscous (Navier-Stokes) solutions be generated for comparison with those shown herein. The resulting comparisons should include particle traces, pressure coefficients, and forces/moments over the aft body.

It is also recommended that prior to the transonic wind-tunnel testing of the TCA model in the 16FTT during Spring 1999, transonic solutions, of at least the Euler type, be available to provide understanding of the anticipated strutinterference effects. Since the transonic strut is shaped differently than the supersonic one, there is a need for some new geometric modeling to be completed in order for this task to be preformed.

The task of developing an internal drag calibration process for HSCT nacelles falls under the category of wind tunnel corrections. For the most part, however, it has been funded out of the Propulsion Induced Effects (PIE) Program. This interesting funding arrangement makes sense because a HSCT nacelle calibration process is a critical requirement for any PIE testing. As the HSCT program has evolved, accurate experimental assessments of nacelle integration effects have become more critical, and nacelle internal geometries have become more complex. As a result, it was apparent that the program had outgrown the empirical nacelle internal skin friction corrections of the past. In the case of $\mathbf{M}=2.4$ cruise point experimental evaluation, the cost and accuracy of CFD based corrections may be sufficient to meet the programs internal drag correction needs. For off-design evaluations, the nacelle internal geometry becomes more complex (ramps \& plugs), the number of required CFD runs increase, and convergence time and uncertainty increase. All this adds up to a cost for CFD based corrections that is currently
 is to calibrate nacelles. Historically, nacelle calibration has been shown to be a reliable, accurate, and
 features in current HSCT geometries and the thrust-drag bookkeeping system, the program could not adopt "as-is" an existing calibration procedure from one of these programs.
Configuration Aerodynamics Technology Development
Demonstrate Significant L/Dmax Gains


Outline
This report documents the results of two tests that were run in the Boeing Flight Simulation Chamber
(FSC) to develop an HSCT specific calibration process. The report is divided into 6 parts. The first part
is a review of thrust-drag accounting systems. The choice of how internal drag can be obtained is
determined by the T-D system. The second part reviews the basic principles concerning how and why
calibration works. The third and fourth part review results of the first HSCT calibration test. This test
went a long way toward validating a HSCT methodology. The fifth and sixth parts review the results of the
second calibration test. This test was intended to use some available hardware to refine the handling of
Reynolds number effects and to explore the feasibility of calibrating with supersonic duct flow.
Before a HSCT specific internal drag calibration methodology can be developed, one must understand the HSCT definition of internal drag. Concurrent with this 1998 CA task was a Technology Integration (TI) Thrust-Drag (T-D) task. This task was chartered to review the existing documentation on the HSCT T-D bookkeeping system and compare that definition to the thrust and drag databases being generated with CFD and experiments. It also compared the HSCT T-D system to that of other commercial and military programs. The ultimate goal of the TI task was to figure out where there were mismatches between the WT \& CFD databases being generated and the T-D system that would be used to develop airplane performance. Recommendations could then be made as to whether to change the types of tests being planned, the types of data these tests would generate, and/or how these data went into airplane performance.

[^31]Review of Thrust-Drag Basic Principles


A Thrust-Drag Bookkeeping System is a formal agreement between Aerodynamics and Propulsion staffs on how to handle 3 key questions: What is the control volume, what is the reference condition and how are off-reference conditions to be handled.
The control volume definition answers the questions: Where does drag end and thrust begin? How does it vary with external and internal conditions?
The reference system definition forms the baseline from which propulsion airframe interactions can be incremented. The fact is, airframe and propulsion control volumes interact. Variable geometry airframes and propulsion systems further complicate matters. Airframe effects entering the propulsion control volume change thrust, and effects within the propulsion control volume change lift and drag. The interactions are termed Propulsion Induced Effects (PIE). PIE effects must be incremented from a reference condition/geometry that is consistent between Aerodynamcs (Aero) and Propulsion staffs.
The off-reference PIE effects system clarifies who bookkeeps what. When an off-reference effect is produced, it often can be termed thrust or drag. The T-D system negotiates these ambiguities into a constant framework.
(1)

A Thrust-Drag Bookkeeping System is a formal agreement
between Aerodynamics and Propulsion staffs on how to
handle 3 key issues.
Control Volume:

- In general, where does Drag end and Thrust begin?

Reference System: - What common configuration will be a baseline both Propulsion and Aero thrust and drag assessments?

Off-Reference PIE effects:

- For non-reference configuration, how will the resultant forces be
handled.

Q eneme February 1999 HSR Airframe Technical Review $\quad \therefore$

> In 1998, a comprehensive review of the HSCT Thrust-Drag Bookkeeping system was funded under the Technology Integration task. Because TI has the ultimate responsibility of incorporating the databases generated by each staff, it was natural that they should conduct such a review. The current T-D system has remained relatively unchanged since its first use on the old Boeing SST. The system has been
> acceptable for HSCT conceptual design studies. The second review was in 1995, when a downselect was required between the system used by Seattle HSCT and that used by Long Beach. The most recent review was meant to be more comprehensive. TI was to review the potential advantages and disadvantages of the current method and alternatives open for consideration. TI would make an assessment of how well current and future CFD and test databases would fit with current and alternate T-D systems. What TI found was that the current T-D system had issues with its control volume definition, reference system definition, and how PIE effects were handled. The conclusion was that the accurate determination of thrust in the current T-D system was expensive and error prone, that neither Aero or Propulsion are
> currently determining thrust accurately or consistently, and that Propulsion was ill equipped to change methodology in the short term.
Review of Thrust-Drag Basic Principles
TIR
TI Review of T-D status Found:

- The accurate determination of thrust in the current T-D
system is expensive and error prone
- Propulsion is currently not accurately determining thrust
- Aero is currently not accurately determining thrust
- CFD solutions are not calculating thrust (internal drag)
- Aero and Propulsion are assessing thrust differently
- In the short term Propulsion cannot change methodology
Qemene February 1999 HSR Airframe Technical Review
The accompanying figure shows the current HSR control volume definition of thrust. This definition is a variation of what textbooks refer to as "intrinsic thrust". An accurate assessment of drag under this definition requires instrumentation to determine the integrated pressure, weight flow, and velocity at location just upstream of the nacelle, at the nacelle highlight, and at the nacelle exit. In practice this would be costly, inaccurate, and problematic.

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An example of the issues associated with the reference definition is shown in the following figure. The reference nacelle definition differs both from what Aero tests and what a real airplane sees over much of it's mission. As an example, the internal geometry is different for all three cases. For the reference geometry is irrelevant). The WT geometry in comparison, has a constant area flow-through duct. Finally, the flight article has the real geometry which has complex geometry throughout. Each of these
geometries will have a different force in the lift direction. For the reference nacelle the force will be zero, the WT geometry will generate some force in the positive or negative lift direction, and the flight article might generate a large lift force in the positive direction. Is duct lift part of drag or thrust? Currently, Propulsion has no way to included lift in the engine deck and never runs an installed flow-thru nacelle so it can not be included in thrust. But Aero never runs the real geometry. The net result is that the Aero CFD solutions are generating a force that has no real home in the current T-D system.
The handling of off-reference PIE effects is further confused in the current system. All effects are broken down into installed versus isolated effects. This requires drag to be build up from 6 numbers: the flight conditions installed and isolated, the reference conditions installed and isolated, and the WT conditions installed and isolated. As one can see, the current system is hopelessly complicated.
Current T-D Reference System
.................:: Riverter diverts all B.L Reference Nacelle Instalation

Nominal Flight Conditions

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While the Technology Integration task was not directly related to the CA calibration methods task, it has been key to determining the direction of the calibration work. The T-D task indicated that there have been some major mismatches in the past on how Aero and Propulsion determine thrust and drag. These mismatches are not trivial and cannot be overcome without substantially more testing and analysis than is currently budgeted in the program management is prepared to spend. In short, , the study showed that it would be impractical to develop a calibration methodology that fully supports the current T-D system. As a result, a short term and a long term solution were proposed. Short term, Aero would incorporate into our drag model any forces that are a result of Aero/Propulsion mismatches. Aero could calculate thrust identical to Propulsion for our wind tunnel models and any forces above and beyond thrust are drag. This would solve the immediate problem of forces being double bookkept or systematically neglected. The long term solution proposed is to update the thrust drag bookkeeping system to something that could and would be supported by both experiment and CFD, and by Propulsion and Aero groups. The system that is proposed is "Net Standard Thrust" and is the industry standard. It is the recommended system by both AIAA and AGARD advisory groups. A nacelle internal drag calibration methodology can be readily adapted to this system. CFD codes can also be easily adapted to accommodate such a system.
Review of Thrust-Drag Basic Principles

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            Na
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        TI T-D status Review Recommendations:
    - Short term fix: Aero should calculate thrust like Propulsion
(right or wrong), and assume all residual forces are drag.
- Long term fix: Adopt a new T-D system that Aero and
Propulsion can and will follow
Q enenve February 1999 HSR Airframe Technical Review
The TI task recommended a change to a new T-D system consisting of a different control volume, reference configuration, and division of PIE effects. The following figure illustrates the recommended control volume for the program in the future. Textbooks refer to this system as Net Standard Thrust. This system requires that the instrumentation to determine the integrated pressure, weight flow, and velocity
only be at a single location. Ram drag is known from the exit weight flow and the freestream velocity.
The new reference configuration would be variable consisting of two or more configurations. As a minimum, there would be separate $M=2.4$ and $M=.90$ cruise references. These reference configurations would be chosen to match flight with the wind tunnel models as closely as possible.
Finally, the system of separating all PIE increments into isolated and installed effects would be
discontinued.
Recommended Thrust Definition For Future Applications
Q AOEINE

> Proposed Definition : Net Standard Thrust: (Industry Standard) $$
\begin{aligned} & \text { Thrust }= M_{\text {dot }} V_{\text {exit }}+A_{\text {exit }}\left(P_{\text {exit }}-P_{\text {inf }}\right) \\ &-M_{\text {dot }} V_{i n f} \\ & \text { (Gross Thrust) } \\ & \text { New T-D Sy Drag) } \\ & \text { Measurements at } 1 \text { Location }\end{aligned}
$$

So why calibrate nacelles, and if it is such a good idea why haven't we done it before? The underlying
assumption behind this work is that the program has reached the point where the accurate determination of internal drag is necessary and worth some expense. Historically HSCT nacelles were straight, constant area, circular ducts aligned with the local inlet flow. Most design work was focused exclusively at the $\mathrm{M}=2.4$ cruise condition. For this scenario, an empirical skin friction analysis very closely approximates the real nacelle internal drag. However, over the last several years many things have changed. The internal ducts are rarely straight, often not aligned with the local flow, and off-design effects are considered important. At all conditions other than the $M=2.4$ condition, a ramp and exit plug must be included in the duct to simulate proper propulsion induced effects. As a result, empirical skin friction analysis is often not only inaccurate, but may bias results with respect to one technology or another. Alternatives to empirical skin friction included CFD, direct force measurement, calculation from internal flow measurements, and calibration + internal flow measurements. Our recent experience with CFD has been mixed. We have had good results with CFD at the cruise condition but less so for transonic conditions. For a transonic WT test, we often want to know the installed internal drag over a broad range of angles of attack, for at four to eight Mach numbers, and for several configurations using the current state of the art, where the accuracy is debatable, the calendar time and cost for Navier-Stokes solutions have been prohibitivly high. There have been cases where 6 months of flow time were required for a matrix of solutions. CFD also has the problem of being run on the "desired" geometry/flow as opposed to WT geometry/flow. The actual model may not have the "desired" contour or smoothness and may have trips, partial laminar flow, and finite leading and trailing edge thickness. For some transonic conditions convergence may also be a problem. Direct force measurement (having a metric duct) is also an option. However, considering accuracy limitations caused by the limited size of HSCT nacelles and the potential cost of implementation, this is assumed to be an option of last resort. Calculation of internal forces from flow measurements is probably the most cost effective approach of the alternatives. However, historically uncalbrated flow measurements have yielded very poor accuracy. This leaves calibration as the most viable approach in terms of cost and accuracy.
General Internal Drag Objectives for HSCT Program
Q Boeine

[^32]February 1999 HSR Airframe Technical Review
The following figure illustrates how a "subsonic" calibration works. Pictured on the left is the mounting
arrangement (above) and internal flow characteristics (below) of a nacelle in the wind tunnel. Pictured on
the right is the mounting arrangement (above) and internal flow characteristics (below) of a nacelle in the
FSC. The objective of a calibration is to correlate the nozzle exit gross thrust and weight flow
characteristics with the nacelle internal flow. The nacelle exit flow (ie. exit Mach) is a function of both the
upstream flow conditions and exit conditions (free stream Mach). It is not necessary to simulate the
entire installed system to reproduce the same exit flow. All that is required is that the control volume just
upstream and downstream of the exit is reproduced. The FSC does this by applying a pressure
differential across the nacelle, thereby accelerating air though a converging bellmouth, through the
nacelle, and out the nozzle exit. In this way only the subsonic flow in the duct aft of the inlet throat is
simulated.
February 1999 HSR Airframe Technical Review The following figure illustrates the nominal "plate mount" testing arrangement in the Boeing Flight Simulation Chamber. The FSC is a 4 ft diameter by 12 ft long chamber which provides the capability to calibrate a blown or flow-through nacelle over the Mach range from 0 to 1.4 at Reynolds numbers fairly assembly consisting of a vertical model support structure supported on either side of the chamber by two 6 component balances. Each balance has a capacity of $\pm 100 \mathrm{lb}$. with a stated $2 \sigma$ accuracy of $\pm .1 \%$ of the applied load. An opening in the support frame midway between the balances allows the model to be mounted with the inlet open to room ambient pressure and the exhaust to be confined to the pressure enviroment of the chamber. The chamber is seal across the balance by a bellows. A pair of compensating bellows are incorporated to cancel out the large pressure force across the balance. This allows for smaller more accurate balances to be used. It also allows the balance to measure only the gross thrust produce by the test article. Two ejectors downstream of the chamber are used to maintain any desired chamber pressure from atmospheric down to 3 psi . Screens are installed in the chamber to break up the exhaust jet, prevent recirculation and entrainment and mix the flow before it exit through the low pressure Multiple Critical-flow Venturi (MCV) at the rear of the chamber. The MCV's provide an accurate reading of mass flow through the test article to a stated $2 \sigma$ accuracy of $\pm .08 \%$ of the applied load. The tank minimum pressure limits the maximum attainable ideal Mach number of the system to 1.4.

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The following picture is of the HSCT 2-D nacelle N31 installed in the FSC in nominal the plate mount

EINE February 1999 HSR Airframe Technical Review measures gross thrust by varying pressure ratio. In an ideal world, the bellmouth design would be so determine flow conditions in the duct. Now, suppose there were losses associated with the bellmouth. It is true that it would require a greater pressure differential to achieve the same gross thrust. But keep in mind, as a result the velocity in the duct must drop as well. This will show up in the duct instrumentation. The goal of calibration is to map the instrumentation reading to gross thrust, not facility pressure ratio.

Conclusions: - However, a bellmouth that provides uniform laminar flow is ideal
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While it is true that in theory, bellmouth losses should have a minimal effect on a calibration, it is good to know that historically Boeing's cubic bellmouth design has produced very near ideal flow. Pictured on the left is ratio of balance measured velocity versus the velocity of ideally (isentropic) expanded flow for a cubic calibration bellmouth/nozzle. The data were taken over 4 installations over 3 tests (1year). This ubic bellmouth's exit velocity is $99.5 \%$ of ideal. Pictured in the middle is the ratio of venturi measured. weight flow versus the weight flow of ideally (isentropic) expanded flow. The fact that the weight flow is $99 \%$ of ideal indicates that there has been very little boundary layer growth. Pictured on the right is the internal drag of this bellmouth based on an HSCT Sref and mass flow. Over most of the regime, the losses are less than .3cts.

Erne February 1999 HSR Airframe Technical Review challenges to adapting it to the HSCT program. HSCT nacelles are typically smaller because our models 747. Our thrust drag bookkeeping is unique. It is more consistent with preliminary design tools than experimental methods. Several facinties where we test at are not well matched in terms of Reynolds As a number and pressure ratio to available calibration facilities. The HSCT's inlet is mixed compression. As a result, the internal flow will be supersonic at the cruise condition. Lastly, an HSCT inlet is square. This presents a challenge of designing a bellmouth that can accelerate flow form $M=0$ to 1 while maintaining uniformity and laminar flow.
Challenges to HSCT Program Using Calibration
to Improve Current T-D Situation
(t)

If this test was successful, a second test exploring long term repeatability, Reynolds number matching,
and supersonic duct flow calibration techniques could be planned.
Two Tests Conducted to Resolve HSCT Calibration Challenges
Q ADEENE
Test 1: IRAD FSC + BSWT (Must be success for HSR PIE test)

- Resolve if Old T-D system is an obstacle
- Develop HSCT data reduction scheme
- Resolve if nacelle size is an obstacle
- Resolve if cost/productivity will be an issue
Test 2: FSC (Desirable for reduction in long term costs)
- Assess long term repeatability
- Explore Reynolds number matching
- Explore supersonic duct flow calibration technique

> To have a focused methods development task, cost and accuracy goals were developed. In this way approaches that could not meet either the cost or accuracy requirement could be immediately abandoned.

> The cost goals were based on comparisons with CFD, assuming Navier-Stokes could produce accurate WT matched forces on a flow-thru nacelle with a ramp and plug. If internal drag is to be known at our standard transonic Mach range, 24 runs (4 alphas at $M=.9, .95, .98,1.06,1.2,1.4$ ) are required. At $100 \mathrm{hrs} / \mathrm{run}$ and $\$ 40 / \mathrm{hour}$ that is $\$ 96 \mathrm{k}$. Assuming engineering costs are $\$ 16 \mathrm{~K}$ for 2 eng. months at $\$ 100 /$ hour, the total cost would be $\$ 112 \mathrm{~K}$ per geometry. Given some automation and realistic reductions in run time, the CFD cost should be able to be brought down by $50 \%-60 \%$. Consequently, a reasonable goal for calibration costs appears to be $<\$ 45 \mathrm{~K}$
The calibration accuracy goals were derived from the WT workshop testing accuracy requirements. The total airplane requirement is $\pm 1 \mathrm{ct}$ at $\mathrm{M}=2.4$ and $\pm 3 \mathrm{ct}$ transonically (and supersonic climb). Consequently, an acceptable level for each nacelle would translate to $\pm .1 \mathrm{ct}$ and $\pm .3 \mathrm{ct}$ repectively (note RSS error on 4 nacelles would be .2cts and .6cts respectively)


[^33]Q enene February 1999 HSR Airframe Technical Review
The following chart provides the definition of accuracy, error, precision, and bias that will be used
Defin
\[

$$
\begin{aligned}
& \text { - Accuracy - the closeness of agreement between a } \\
& \text { measured value and the truth. } \\
& \text { - Error - the difference between the measured value } \\
& \text { and the true value. } \\
& \text { - Precision - the random component of error. } \\
& \text { - Spot repeatability - samples within a data point } \\
& \text { - Short term repeatability - repeats within a test } \\
& \text { - Near term repeatability - repeat model builds within a test } \\
& \text { - Long term repeatability - repeat runs, different test } \\
& \text { - Bias - the systematic component of error. } \\
& \text { Error = Bias + Precision }
\end{aligned}
$$
\]

Test Results
Results Summary of First Test
(February/IRAD)
Results Summary of First Test (February/IRAD)
February 1999 HSR Airframe Technical Review
Once set on the path of using Net Standard Thrust, the primary unknowns are the nozzle exit conditions. Any cost associated to assessing the local or hilight flow conditions become non-value added. The objective of the calibration work conducted in 1998 was to develop a methodology the would support both the existing and proposed definition of thrust without incurring the added costs that the current system mandates. In this way, interim performance numbers could be generated until the program formally changes the thrust definition.
If a calibration is done to generate net standard thrust, there is a way that the same data can support HSCT's intrinsic thrust definition. If all the inlet unknowns that are not required by net standard thrust are obtained directly from Propulsion, there can be no mismatch between Aero and Propulsion on the assessment of thrust. This hybrid solution may have a bias error relative to the "true" value of intrinsic thrust. However, this bias will have an offsetting drag increment that will result in the correct aircraft performance.
QueganeDevelopment of Calibration Process for HSCT T-D system
A calibration process can be applicable to either T-D system

Q ecenv February 1999 HSR Airframe Technical Review
How much information does the instrumentation have to provide in order to calculate drag. The following
chart examines the instrumentation requirement based on 1-D flow.
Development of Calibration Process

| Given | $\mathrm{Me}=2.236\left[(\mathrm{Pe} / \mathrm{PT})^{-2 / 7}-1\right]^{(.5)}$ |
| :---: | :---: |
| Then | $\mathrm{Me}=\boldsymbol{f}(\mathrm{Pe} / \mathrm{PT})$ |
| Given | $\mathrm{Ve}=49.02{ }^{*} \mathrm{Me}$ * $\mathrm{TT}^{(.5) ~ *}\left(1+0.2{ }^{*} \mathrm{Me}^{2}\right)^{(-.5)}$ |
| Then | $\mathrm{Ve}=\boldsymbol{f}(\mathrm{Pe} / \mathrm{PT}, \mathrm{TT})$ |
| Given | Wdot $=0.91996$ * PT * TT ${ }^{(-0.5) ~ * ~} \mathrm{Ae}$ * $\mathrm{Cde}^{*} \mathrm{Me} *\left(1+0.2 * \mathrm{Me}^{2}\right)^{(-3)}$ |
| Given | Cde=f (Me, PT, TT) |
| Then | Wdot $=\boldsymbol{f}$ (Pe/PT, PT, TT) |
| Given | $\mathrm{Fg}=\mathrm{Wdot}^{*} \mathrm{Ve} / \mathrm{g}+\mathrm{Ae}^{*}$ (Pe-Pamb) |
| Then | $\mathrm{Fg}=\boldsymbol{f}$ ( $\mathrm{Pe} / \mathrm{PT}, \mathrm{PT}, \mathrm{Pamb}, \mathrm{TT}$ ) |

[^34]Q enene February 1999 HSR Airframe Technical Review

[^35]Problem: It is difficult to instrument for Pe

Conclusion: If PT, Pamb and TT are known, Fg and Wdot are also known (and Pe is redundant)

instrumentation. The following charts shows how, in theory, the static pressure at two stations provides as much information as the total pressure.

| Problem: PT rakes add flow distortion <br> Define a upstream duct location, called a charging station, preferably where Acharge > Aexit |  |
| :---: | :---: |
|  |  |
| Given | $\mathrm{Mc}=2.236\left[(\mathrm{Pc} / \mathrm{PT})^{-2 / 7}-1\right]^{(5)}$ |
| Then | $\mathrm{Mc}=\mathrm{f}$ (Pc/PT) |
| Given | (Wdot)e = (Wdot)c |
| Given | $\begin{aligned} \text { Wdot } & =0.91996 * \text { PT * TT }(-0.5) * \mathrm{Ae}^{*} \mathrm{Cde}^{*} \mathrm{Me} *\left(1+0.2 * \mathrm{Me}^{2}\right)^{(-3)} \\ & =0.91996 * \text { PT *TT }{ }^{(-0.5)} \text { Ac*Cdc }{ }^{(2 c}{ }^{*}\left(1+0.2 * \mathrm{Mc}^{2}\right)^{(-3)} \end{aligned}$ |
| Given | Cde=f(Me, PT, TT) and Cdc=f(Mc, PT, TT) |
| Then | Wdot $=\boldsymbol{f}(\mathrm{Pe} / \mathrm{PT}, \mathrm{PT}, \mathrm{TT})=\boldsymbol{f}(\mathrm{Pc} / \mathrm{PT}, \mathrm{PT}, \mathrm{TT})$ |
| Therefore | - $\mathrm{PT}=f(\mathrm{Pe} / \mathrm{Pc}, \mathrm{Pc}, \mathrm{TT})$ |
| Then | Wdot $=\boldsymbol{f ( P e / P c , ~ P c , ~ T T ) ~})$ |
| Then F | $\mathrm{Fg}=\boldsymbol{f}(\mathrm{Pe} / \mathrm{Pc}, \mathrm{Pc}, \mathrm{Pamb}, \mathrm{TT})$ |
| Conclusion: The static pressure at two stations, Pe and Pc, can substitute for knowing PT |  |

Q enenve February 1999 HSR Airframe Technical Review
Using the same logic that indicates knowledge of Pe is redundant if PTe is known, one can'deduce that Pe is also redundant if Pc is known. The following chart is the derivation that in principle, a single nacelle static pressure along with the free stream static pressure (or total pressure) are enough to characterize the internal flow. There are two basic scenarios. The first is described here, where Ac < Ae and the duct is choked. In this situation the duct total pressure and a charging station pressure are offset but track almost one for one. Consequently Pc/Pamb can be substituted for PTe/Pamb. If Ac~Ae, and the duct is unchoked, Pc~Pamb, so Pc/Pamb~1. For this scenario, the ratio of PTinf/Pc can be used instead. This is reasonable because if the duct is unchoked, PTinf >PTe (PTe includes duct viscous losses) and will track almost one for one.
Development of Calibration Process
Problem: It is difficult to instrument for Pe

Conclusion: If only Pc, Pamb and TT are known, Fg and Wdot
are also known (if solution is unique)
February 1999 HSR Airframe Technical Review The following figure illustrates the instrumentation suite selected for these tests. The preceding derivation indicates that a calibration can be accomplished by either static and total pressure instrumentation at a single station, static pressures at two stations, or static pressures at a single "charging station". Aside from fabrication costs (+2K/geometry), there doesn't seem to be any downsides to instrumented both the charging station and the exit for static pressure. Total pressure rakes, on the other hand, have a several issues. On nacelles of such small size, they are difficult and costly to install and tend to provide unsatisfactory readings of total pressure. This is because it is difficult to achieve high enough density in the rake while maintaining minimal rake induced distortion and blockage. In the end, the uncalibrated rake will never read the "true" average exit total pressure and will require a large calibration correction. The net result is that a rake seems to add less information than the uncertainty it creates. In the end, static instrumentation at two stations were selected because they provide the maximum data reduction possibilities (for little additional cost), and serve as an invaluable tool for duct flow validation/troubleshooting. The location of the charging station was just up stream of where the area growth of a plug would begin. The exit pressures were .1" upstream of the exit.
Development of Calibration Process for HSCT T-D system

Selected "Rakeless" Instrumentation suite: static's at 2 stations - Costly/Difficult to install rake in small model

- Difficult to get enough rake density 1 sqin duct
- Rake adds own uncertainty, PTexit/ PTrake $<93 \%$
Q enerne February 1999 HSR Airframe Technical Review
Developing a data reduction scheme for this test was challenging. While it would be desirable to strive for consistency between an HSCT methodology and that of another program, it is not obvious which program to select. It seems as if there is no shortage of parameters in use to normalize the calibration data. Some examples are

The following chart reviews one of several methodologies used by St. Louis on fighter programs and the issues associated with applying them to an HSCT.
Development of Calibration Process

| BOEINE Del |  |
| :---: | :---: |
| \%pewtivinuy |  |
| Military Group Process | HSCT Issues With This Process |
| Nacelle has total pressure rake PT, and exit pressure rake, PS | - $\quad$ Thrust $=\mathbf{F g}-$ Fram $_{\text {local }}-\mathrm{D}_{\text {add }}-$ Dpref <br> - Requires duct PT rake |
| Thrust $=$ Fg - Fram | - Requires exit static pressure |
| $\begin{aligned} \mathrm{CW}= & \text { Wdot }_{\text {actual }} / \text { Wdot }_{\text {rake-measured-calc }} \\ & =\mathbf{A 0}+\mathrm{A} 1^{*}(\text { PTR } / \text { Pinf })+\text { A }^{\star}(\text { PTR } / \text { Pinf })^{2}+ \\ & \text { A3 }^{*}(\text { PTR } / \text { Pinf })^{3} \end{aligned}$ | - Difficult to get an accurate reading of PT in a 1 " model |
| $\mathrm{CFG}=\mathrm{Fg}_{\text {actual }} / \mathrm{Fg}_{\text {rake-measured-calc }}$ | 7\% |
| $\begin{aligned} & =\mathrm{B0}+\mathrm{B} 1^{*}(\mathrm{PTR} / \text { Pinf })+\mathrm{B} 2^{*}(\mathrm{PTR} / \text { Pinf })^{2}+ \\ & \text { B3 }^{\star}(\text { PTR/Pinf })^{3} \end{aligned}$ | - Sparse rake may not fully capture Reynolds number variation |
| $\mathrm{Me}=\mathrm{f}(\mathrm{PT} / \mathrm{PS})$ |  |
| $\mathrm{Ve}=\mathrm{f}(\mathrm{Mc}, \mathrm{TT})$ | Advantages This Process |
| $\begin{aligned} & \text { Wdot }_{\text {ideal }}=f\left(\mathrm{Me}^{2}, \mathrm{PT}, \mathrm{TT}\right) \\ & \mathrm{Fg}_{\text {ldeal }}=\mathrm{Wdot}_{\text {ldeal }} / \mathrm{g} * \mathrm{Ve}+\mathrm{Ae}(\text { PS-Pinf }) \end{aligned}$ | - In theory, does not require Reynolds number correction |
| $\begin{aligned} & \text { CDint }=[\text { CW Wdot } \\ &-[\text { CFG Fg } \\ & \text { ldeal } \\ &/ \mathrm{g} * \text { Vinf }] / \mathrm{Q} / \text { Sref } \\ & \hline \text { Sref } \end{aligned}$ | - Captures any inlet shock PT losses |

Q enecine February 1999 HSR Airframe Technical Review
The following chart reviews one of several methodologies used by Seattle on transport programs and the issues associated with applying them to an HSCT.

$$
\begin{aligned}
& \text { HSCT Issues With This Process } \\
& \text { - Thrust = Fg - Framlocal }-D_{\text {add }}-D_{\text {Pref }} \\
& \text { - Inlet shocks may make PTduct<PTinf } \\
& \text { - For long ducts PTduct/PN not constant } \\
& \text { with Reynolds number (i.e. Re correction } \\
& \text { will be too large) }
\end{aligned}
$$

Commercial Group Process
Nacelle has static pressures, PN, at 1 location

Wcorr $=A 0+A 1^{*}(P T / P N)+A 2^{*}(P T / P N)^{2}+$
Wdot $=$ Wcorr ${ }^{*}(\mathrm{PT} / 14.696) /(\mathrm{TT} / 518.69)^{5}$
VH/Vinf=Mhi/Minf * $\left[\left(1+.2 \text { Minf }^{2}\right) /\left(1+.2 \mathrm{Mhi}^{2}\right)\right]^{5}$ $\begin{aligned} \mathrm{CV}= & \mathrm{Fg}_{\text {actual }} / \mathrm{Fg}_{\text {ideal }} \\ = & \mathrm{BO}+\mathrm{B1}^{*}(\mathrm{VH} / \text { Vinf })+\mathrm{B2}^{*}(\mathrm{VH} / \text { Vinf })^{2}+ \\ & \mathrm{B1}^{*}(\mathrm{VH} / \text { Vinf })^{3} \\ \text { CVNC }= & 1-(1-\mathrm{CV})(\text { Re/ft~FSC/Re/ft~test) })^{+.2}\end{aligned}$
CDint $=(1-C V N C) W d o t / g ~ * V i n f / Q / S r e f ~$

$$
\begin{aligned}
& \text { For ducts where Ae<< Aduct, favorable } \\
& \text { pressure gradient may reduce Reynolds } \\
& \text { number effect }
\end{aligned}
$$

$$
\begin{aligned}
& \text { Advantages This Process } \\
& \text { - Simplicity of application } \\
& \text { - few curvefits reduce lack of fit error } \\
& \text { - Requires only static pressures }
\end{aligned}
$$

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Given that there are infinite variations of a data reduction scheme, which is best for an HSCT? For this test several methods where examined for assessing duct total pressure. The knowns in a calibration are weight flow, gross thrust, TT, and the static pressure and area at two stations. This implies that there are too few equations and too many knowns. Consequently, the total pressure can be solved for independently 5 ways.

1) Solve for PT at the charging station using static pressure and weight flow
2) Solve for PT at the charging station using static pressure, gross thrust and weight flow 3) Solve for PT at the exit station using static pressure and weight flow
3) Solve for PT at the exit station using static pressure, gross thrust and weight flow
Solve for PT at the either station using both static pressures and weight flow

Interestingly, each of these PT's will be different. There is also an infinite combinations of pressure ratios and functions that these resulting PT's can be correlated against.

> A gross thrust coefficient was used to calibrate gross thrust. Coefficients based on measured gross thrust versus ideal (expanded to Pamb) and calculated (using ideal 2-D flow) gross thrust were looked at. Given there were 5 values of PT and two methods of gross thrust coefficient (Ideal \& calculated), there were 10 ways to assess the thrust coefficient. As with the value of PT, there was also any number of variables to correlate to. In the end, all these methods are self consistent. They were derived from calibration weight flow and gross thrust values and will therefore produce these same values. When these numbers are subjected to wind tunnel data, the results diverge slightly. At $M=.90$, the scatter between methods was up to 2 counts. Ultimately, the "best" methodology is one that maintains consistency in form and process to what Propulsion must use to calibrate thrust (for wind tunnel and flight test). In that way "thrust", which is a defined value, will be by definition accurate.
Development of Calibration Process
Dozens of data reduction schemes in use on other programs
Examined 5+ methods for determining duct total pressure Examined 10+ methods for normalizing gross Thrust
All methods are self consistent with calibration data
"Best method" is one that's most consistent with Propulsion
$\theta$
Q emerno February 1999 HSR Airframe Technical Review
The following figure shows an example of the correlation curves for the 5 difference values of PT and 9 different values of gross thrust coefficient.
Development of Calibration Process
Examined 5+ methods for determining duct total pressure

Q eacone February 1999 HSR Airframe Technical Review
Given all these available methods, which one is recommended? While there is no clear distinction among
the methods, it is reasonable to believe that getting a reliable reading of Pe where $\mathrm{M} \sim 1$ for most cases is
considerably more difficult than Pc. Therefore, it is recommended the for the near term, PT be derived
form Wdot and Pc, based only on Pc/Pamb. The issue of whether it is better to have gross thrust
coefficients based on measured gross thrust versus ideal (expanded to Pamb) and calculated (using ideal
2-D flow) gross thrust is less clear. For the near term, Ideal gross thrust will be the preferred method
because it is more consistent with Propulsions methodology. The following figure shows how the thrust
coefficient will be incorporated into the data. It is worth noting that the gross thrust in this methodology
is treated as a scalar quantity. This is because it is consistent with Propulsions current methodology.
Development of Calibration Process
Rewriting HSCT Thrust equations

Resulting Definition Becomes:
Thrust $=-$ CDint $=\mathbf{C g} \mathbf{F g}_{\text {calc/ideal }}-$ Wdot $_{\text {calculated }}{ }^{*} V_{\text {local }}-D_{\text {add }}-D_{\text {Pref }}$
 where
$\quad V_{\text {local }}=f\left(M_{\text {inf }}, T T\right)$
$D_{\text {add }}=f\left(M_{\text {inf }}\right.$, TTinf, PTinf, Wdot $\left.{ }_{\text {actual }}\right)$
$V_{c=f}\left(\mathrm{Mc}^{\prime}, \mathrm{TT}\right)$
Wdot $_{\text {caculated }}=f($ Mc, PTc, TT $)$
February 1999 HSR Airframe Technical Review
The following figure illustrates how the recommended process would be applied. PT is derived from Pc, and a function of Pc/Pamb. With Pc, PT, and TT, Wdot is calculated. Cg is also derived from the same function of Pc/Pamb. With PT and Pamb, FGideal is calculated. With Cg and FG 恠eal, FG is calculated. With Wdot, PTinf, TT, and Pinf, Vlocal, Dadd, Dpref are calculated. Given Wdot, Sref, Q, FG, Vlocal, Dadd, and Dpref internal drag can be finally calculated.

## Development of Calibration Process

Example of How Calibration Data is Used in WT Tunnel Corrections |  |
| :---: |
| Step 1: Calculate PTC/PC |
| from and Pamb/PN |
| Step 2: Calculate PTc from |
| PTC/PC and PC |
| Step 3: Calculate Cg from |
| Pamb/PC |
| Step 4: Calculate |
| MC,VEideal from |
| PC/PTC |
| Step 5: calculate Wdot |
| Step 6:Calculate Fgideal |
| Find 7: Calculate internal |
| drag |



$=-$

> One of the main objectives of the first test was to detirmine whether an HSCT model's small scale had exceeded the calibration facility's precision limits or whether the combination of repeatability and productivity would make the testing too costly. The short term force repeatability turned out to be less than $\pm .04$ cts. The balance checkloads were excellent, being within $\pm .012$ lbs ( $\sim .06 \mathrm{cts})$ of the calibration standard. The short term repeatability of weight-flow was < $\pm .03 c t s$ when converted to ram drag. The accuracy compared to the calibration standard is quoted to be $(2 \sigma) \pm .08 \%(\sim .04 c t s)$.
Assessment of Facilities Force/Weight flow Measurement
Force/Gross thrust:

- Short term force repeatability < $\pm .04 \mathrm{cts}$
- Checkload within $\pm .012$ lbs $(\sim .06 \mathrm{cts})$
Weight-flow/Ram Drag:
- Short term weight-flow repeatability < $\pm .03 \mathrm{cts}$
- Calibration accuracy $(2 \sigma) \pm .08 \%$ ( $\sim .04 \mathrm{cts})$
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The productivity was such that 2-3 installations/shift were possible (at $\$ 4 \mathrm{k} /$ shift). With the level of repeatability/accuracies seen in this test, there does not appear to be a big incentive to run multiple repeats within a test. Assuming new bellmouths are made ( $\$ 3 \mathrm{~K}$ ), two nacelles are instrumented and unı $\ddagger$ SOכ uо! approximately $\$ 10 \mathrm{~K}$ per geometry (less if multiple geometries are run). Nacelles were instrumented with 10 pressures at a cost of $\$ 2 \mathrm{~K} /$ nacelle. Future tests will most likely only have 8 . With the described data reduction methodology, the use of calibrated nacelles would not significantly reduce in-tunnel productivity because internal drag would be obtained with base drag data runs.
Assessment of Costs/Productivity
for HSCT Applications
Productivity:
- Over 3 shift test 6 Mach sweeps obtained
$-\quad$ 2-3+ geometries can be calibrated per shift
- Use of calibrated nacelles would not significantly reduce
in-tunnel productivity
- Internal drag obtained with base drag data
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This chart shows the results of the BSWT spillage runs when the nacelle internal drag was corrected with
the calibration data. 4 runs were made: Mach $=0.90$ and $M=1.2$ at $\sim 8 \%$ and $\sim 35 \%$ spillage. If Aero is
assuming responsibility for all spillage effects above and beyond what Propulsion accounts for in
additive drag, then spillage interference is the difference in nacelle increments between the $\sim 8 \%$ and
$\sim 35 \%$ spillage cases. That number is shown in the right hand axis. At $M=.9, C L=.24$, spillage is worth
$\sim 8 c t s$. At $M=1.2, C L=.14$, it is worth $\sim 4.0$ cts. This result is consistent with past results from the 1995
ARC111-119 spillage test and recent OVERFLOW results.

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The first FSC test was a success in all respects. A "proof of concept" low cost PIE test was conducted, and the results were consistent with expectations. Several "rakeless" data reduction schemes were developed for old/new T-D systems so that despite the current small nacelles, repeatability/accuracy and costs (productivity) should meet requirements for HSCT applications. Despite the success of this test, follow-on improvements are always possible. It was therefore recommended that a future test should be conducted to assess long term repeatability, explore Reynolds number matching capability and supersonic duct flow. In addition, CFD and test should always be complementary. Therfore some CFD should be run on some of the calibration cases to assess uncertainty and evaluate the "best" data reduction scheme.
Conclusions After First FSC/BSWT Test
Summary:
- A "proof of concept" low cost PIE test was conducted
- Several "rakeless" data reduction schemes developed for
old/new T-D system
- Despite nacelle size, repeatability/accuracy and costs
(productivity) should meet requirements for HSCT applications
Recommendations

[^36]Test Results Results Summary of Second Test
(October/HSR)
Results Summary of Second Test
(October/HSR)
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Scheduling issues required that the second test be broken down into two entries with another customer's
test in between. The following table outlines the runs made in the first entry of two.

| dCminatismot |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Series | Runs | Configuration | RN | Mach | Comment | Resultatend of series |
| 11 | 227-237 | 2* CUBIC | nom | 6-1.4 | repeat of lacility check | Data finally within historical trends |
| 12 | 240.247 | BM 1, N31.3 | nom | $6 \cdot 1.4$ | repeat of test-to-test check | CDTC slightly higher than series 2 |
| 13 | 256.266 | $2^{\prime \prime}$ CUBIC, $7^{*}$ STB | $\mathrm{pt}=14.7$ | 6-1.4 | confirm no moment/installation tare | oposed to +.02 cts , and drag is nigher by -.5 cts |
| 14 | 267-274 | $2^{*}$ CUBIC, $7^{\text {- }}$ ST8 | $\mathrm{PT}=23$ | 6-1.4 | confirm no moment/installation tare | down and CVN goes up only to the tire air level |
| Repeat of AF4145 contiguration with STB mounted |  |  |  |  |  |  |
| 15 | 283-333 | BM 3.1.7*STB, N31.3 | 4 | 6-1.4 | contirm no moment/installation tare |  |
| 16 | 283-333 |  | 4 | 6-1.4 | confirm no moment/installation tare |  |
| 17 | 283-333 | 8M3.1, $7^{\circ}$ STB, N31.3 | 4 | 6-1.4 | conlirm no moment/installation tare |  |
| 18 | 283-333 | ВM3.1, $\mathbf{7 C N S}^{\text {STB, N31.3 }}$ | 14 | 6-1.4 | contirm no moment/installation tare |  |
|  |  |  |  |  |  |  |
| 19 | 339-341 | BM 7.1, $\mathbf{7}^{\text {S }}$ STB | 7 | 2.4 | Install alum. $\mathrm{M}=2.4 \mathrm{C}-\mathrm{D}$ bellmouth | CDV -1 , but the CVN are low at $\sim .97$ |
| 20 | 342-385 | BM3.1, 7 STB, N31.3 | 4 | 8.1.4 | repeat of N31.3 BM 3.1 | Repeat multiple AN capability |
| 21 | 342-385 | BM3.1, $\mathbf{7}^{\text {S STB, }}$ N 31.3 | 4.5 | 8-1.4 | repeat of N31.3 8M3.1 | Repeat multiple RN capability |
| 22 | 342-385 | BM3.1, $7^{\circ}$ STB. ${ }^{\text {S }} 1.3$ | 5.5 | 8.1.4 | repeat of N31.3 8M3. 1 | Repeat multiple RN capability |
| 23 | 342-385 | 8M3.1, 7* STB, $^{\text {S }} 1.3$ | 7.2 | .8-1.4 | repeat of N31.3 8M3.1 | much scatter and drag still goes up with AN |
| 24 | 386.389 | BM7. $\mathbf{7 C O}^{\text {STB }}$ | 4,9,14 | 2.4 | Install $\mathrm{M}=2.4 \mathrm{C} \cdot \mathrm{D}$ bellmouth | indicate $\mathrm{M}=2.4$ |
| 25 | 392-403 | BM3.2.7* STB | 4.5 | nom | repeat of N31.3 BM3. 1 | isolated belmouth |
| 26 | $405 \cdot 407$ | BM6 | 6,9,12 | 2.1 | Install $\mathrm{M}=2.1 \mathrm{C}$ - bellmouth | indicate $M=2.1$ |
| 27 | 408.410 | BM5 | 6.9,12 | 1.8 | Install $M=1.8 \mathrm{C}$ - D bellmouth | indicate $M=1.8$ |
| 28 | 411.413 | BM 4 | 5.7.10 | 1.5 | Install $M=1.5 \mathrm{C} \cdot \mathrm{D}$ bellimouth | indicate $M=1.5$ |
| 29 | 415.423 | BM7. $\mathbf{7 *}^{\text {S }}$ STB | 4.9.14 | 2.4 | install $\mathrm{M}=2.4 \mathrm{C}$ - b bellmouth | indicate $M=2.4$ |
| 30 | 428.436 | BM7.t. N31.3 | 4,9,14 | 2.4 | install alum. $\mathrm{M}=2.4 \mathrm{C}-\mathrm{D}$ bellmouth | drag |
| 31 | 437-445 | BM4, N31.3 | 4,9,14 | 1.5 | instail $\mathrm{M}=1.5 \mathrm{C}-\mathrm{D}$ bellmouth |  |
| 32 | 446-462 | BM3.2, $7^{\text {S }}$ STB | 4 | 8-1.4 |  |  |
| 33 | 446-462 | BM $3.2,7^{*}$ STB | 4 | 8.1 .4 |  |  |
| 34 | 446-462 | BM 3.2, 7* STB | 4 | 8-1.4 |  |  |
| 35 | 446-462 | ВM3.2, $7^{\text {² }}$ STB | 4 | 8.1.4 |  | ellects and yet drag on an isolated bellmout is almost |
|  |  |  |  |  |  |  |
| 36 | 468-523 | BM 3.1, $7^{\circ}$ STB, 31.3 | 4 | .8-1.4 | repeat of N31.3 8M3.9 |  |
| 37 | 468.523 | BM3.1, $\mathbf{7}^{\text {C S STB, N31.3 }}$ | 4 | . $8-1.4$ | repeat of 31.3 8M3.1 |  |
| 38 | 468.523 | BM3.1, ${ }^{\text {P S STB, N31.3 }}$ | 4 | .8-1.4 | repeat of N31.3 BM3.1 |  |
| 39 | 468.523 |  | 4 | .8-1.4 | repeat of N31.3 BM3.1 | drag is still increasing with Reynotds number |
| 40 | 527-535 | BM5, 7" STB, N31.3 | 4,9,14 | 1.8 | Install $\mathrm{M}=1.8 \mathrm{C}-\mathrm{D}$ bellmouth |  |
| 41 | 536.544 |  | 4.9,14 | 2.1 | Ins tall $\mathrm{M}=2.1 \mathrm{C}$ - D bellmouth |  |
| 42 | 545-553 | BM7, 7* STB, N31.3 | 4,9,14 | 2.4 | Install $\mathrm{M}=2.4 \mathrm{C}$-D bellmouth |  |
| 43 | 554.562 | BM7.1, ${ }^{\text {² S STB, N31.3 }}$ | 4,9,14 | 2.4 | install alum. $\mathrm{M}=2.4 \mathrm{C}-\mathrm{O}$ bellmouth | nigher than skin friction. |
| 44 | 566-601 | BM3.1, $7^{*}$ STB, N31.3 R1.1 P2 | 4 | .8.1.4 | repeat of N31.3 BM3.1 |  |
| 45 | 566-601 | BM3.1, $\mathbf{7}^{\text {® }}$ STB, N31.3 R1.1 P2 | 4 | .8-1.4 | repeat of N31.3 BM3.1 |  |
| 46 | 586-801 | BM3.1, $7^{*}$ STB, N31.3 71.1 P2 | 4 | .8-1.4 | repeal of N31.3 BM3.1 |  |
| 47 | 566.601 | BM3.1, $\mathbf{7}^{\text { }}$ STB, N 31.3 P 1.1 P 2 | 4 | .8-1.4 | repeat of N31.3 BM3.1 | is basically insensitive to reynolds number |
| 48 | 602.642 | BM 3.1, $\mathbf{7}^{\text {C S STB, }}$ N31.3 trips off | 4 | .8-1.4 | repeat of N31.3 BM3.1 |  |
| 49 | 602.842 |  | 4 | .8-1.4 | repeal of N31.3 BM3.1 |  |
| 50 | 602-642 | BM 3.1, 7" STB. N31.3 trips off | 4 | .8-1.4 | repeat of N31.3 BM3.1 |  |
|  |  | BM 3.1, $7^{+}$STB, 3 31.3 trips off | 4 | 8-1.4 | repeat of N31.3 EM3.1 | that the trip on data, but it has less variability to RN. |
|  |  | End ol Entry ! |  |  |  | Total charged hours |

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This table outlines the runs made during the second entry of two.

| Series | Runs | Configuration | RN | Mach | Comment | Result at end of series |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 52 | 655-665 | BM3.2 tree air | nom | nom |  | is BAD |
| 53 | 666-673 | BM3.1, N31.3 free air | nom | nom | repeat of test-to-test check | data hi by .3cts, flow viz indicates stron recirculation b |
| 54 | 678-685 | BM1, N31.3 free air | nom | nom | repeat of test-to-test check | Agrees with series 12 |
| 55 | 688-695 | BM3.2.1, N31.3 tree air | nom | nom | modify adaptor to reduce recirculation | data looks better, but not yet back to plate mount level |
| 56 | 696-700 | BM3.2.2, N31.3 tree air | nom | nom | modity adaptor to reduce recirculation | data looks still better, but not yet back to plate mount |
| 57 | 701-709 | BM3.2.2 free air | nom | nom | effects perf. | isolated bellmouth Matches original perf. |
| 58 | 713-721 | BM1, N31.3 R1.1 P2 TS2 free air | nom | nom | (TS2) free air | baseline nacelle with gaurenteed furbulant flow |
| 59 | 722-731 | BM1.1, N31.3 R1.1 P2 TS2 free air | nom | nom | (TS2) free air low mount tongue |  |
| 60 | 732-741 | BM1.2, N31.3 R1.1 P2 TS2 | nom | nom | (TS2) free air low mount and mid horiz |  |
| 61 | 751-787 | BM3.2.2, STB, N31.3 R1.1 P2 TS2 | 4 | 8-1.4 | (TS2) STB | Find out if all super trip cases collapse |
| 62 | 751-787 | BM3.2.2, STB, N31.3 R1.1 P2 TS2 | 4.5 | .8-1.4 | (TS2) STB |  |
| 63 | 751-787 | BM3.2.2, STB, N31.3 R1.1 P2 TS2 | 5.5 | .8-1.4 | (TS2) STB |  |
| 64 | 751-787 | BM3.2.2, STB, N31.3 A1.1 P2 TS2 | 7.2 | .8-1.4 | (TS2) STB |  |
| 65 | 794-807 | BM8, STB | 14.7 |  | 1" cubic, free air |  |
| 66 | 810-813 | BM7.0.1, STB, N31.3 R1.1 TS2 | 4,9,14 | 2.4 | Run N31.3 R1 with BM7 bellmouth | Ramp broke off |
| 67 | 815-817 | BM7.0.1, STB, N31.3 R1.1 TS2 | 4,9,14 | 2.4 | Run N31.3 R1 with BM7 bellmouth | wrong MCV code |
| 68 | 818-820 | BM7.0.1, STB, N31.3 R1.1 TS2 | 4,9,14 | 2.4 | Run N31.3 R1 with BM7 bellmouth | run supersonic case with ramps |
| 69 | 821 | BM7.0.1, STB, N31.3 R1.1 TS2 | 4,9,14 | 2.4 | Run N31.3 R1 with BM7 bellmouth | run supersonic case with ramps with "nickle" asymetrid |
| 70 | 823-825 | BM6.0.1, STB, N31.3 R1.1 TS2 | 6,9,12 | 2.1 | Run N31.3 R1 with BM6 bellmouth | run supersonic case with ramps |
| 71 | 826-828 | BM6.0.1, STB | 6,9,12 | 2.1 | isolated BM6 bellmouth |  |
| 72 | 829-832 | BM7.0.1, STB | 4,9,14 | 2.4 | isolated BM7 |  |
| 73 | 834-840 | BM3.2.2, STB, N36 R1 N3 TS2 | 4 | .8-1.4 | (TS2) STB |  |
| 74 | 841-846 | BM3.2.2. STB, N36 R1 N3 TS2 | 7 | .8-1.4 | (TS2) STB |  |
| 75 | 847-864 | BM3.2.2, STB, N36 R1 N2 TS2 | 4 | .8-1.4 | (TS2) STB |  |
| 76 | 847-864 | BM3.2.2, STB, N36 R1 N2 TS2 | 7 | .8-1.4 | (TS2) STB |  |
| 77 | 867.875 | BM3.2.3, STB | 4 | .8-1.4 | isolated BM3.2.3 | modified bellmouth |
|  | 876 | BM3.2.3, free alr, low viz | nom | 0.95 | isolated BM3.2.3 | modified bellmouth, free air, flow viz |
|  | 877 | BM3.2.2 free air, flow viz | nom | 0.95 | isolated BM3.2.2 | modified bellmouth, free air, flow viz |
|  | 877 | BM3.2 free air, flow Viz | nom | 0.95 | Isolated BM3.2 | modified bellmouth, free air, flow viz |
| 78 | 879-886 | BM3.2.4 free air | nom | .8-1.4 | isolated BM3.2.4 | modified bellmouth, free air |
|  | 887 | BM8, free air, flow viz | nom | 0.95 | isolated 1 inch cubic | fow viz |
| 79 | 888-894 | BM3.1.5 free air | nom | .8-1.4 | isolated BM3.1.5 | modifled bellmouth, tree air |

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In accordance with our objectives, the same model that was tested in the first test was re-tested in each of the second test's two entries. Long term force repeatability was under $\pm .16$ cts. Long term weight-flow repeatability was under $\pm .07 \mathrm{ct}$. It turns out that the random shifts in force correlated well with weight flow. Because drag is gross thrust minus ram drag, the net result is that the total variability on drag was less than $\pm .08 \mathrm{cts}$.
Assessment of Long term repeatability
for HSCT type nacelles
Force/Gross Thrust:

- Long term force repeatability $< \pm .16$ cts Weight-flow/Ram Drag:
- Long term weight-flow repeatability < $\pm .07 \mathrm{ct}$
*Random Shifts in Force correlated well with Weight flow so
uncertainty in Fnet < RSS(Fram + FG)
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The following figure shows the short and long term repeatability of gross thrust in the left hand axis and
of ram drag in the right hand axis.
The BSWT test that was to use this calibration data was run at a Re=7million/ft. The calibration data was
obtained at a Re~ 4 million/ft. A reasonable question that is often posed is whether the Re=4million
calibration data is still applicable at 7 million, and what corrections should be applied. Having a mismatch
between calibration Reynolds numbers and test is not unique to the HSCT program. Depending on the
wind tunnel facility(High Reynolds number or atmospheric) and the calibration facility (FSC or Thrust
stand) a Re mismatch of $2-4$ times exists for most programs. Other programs have evolved two distinct
approaches to handling this issue.
On commercial transports, nacelles are large and have a low length/diameter ratio (reiative to an fighter). As a result $\delta^{\star} \ll$ duct diameter. Consequently, the nacelle internal static's pressure coefficient is relatively insensitive to changes in Reynolds number. Therefore if internal drag ( $1-\mathrm{Cv}$ ) is correlated to the nacelle pressure ratio (PTinf/PN), a correction must be applied to (1-Cv) for Reynolds number. A derivation of a first-order correction for axisysmetric and 2-D nacelles is described in the next two charts.

$$
\begin{aligned}
& \text { On fighter programs, the duct diameters are less and I/d is larger }(\sim 10) \text {. The net result is that } \delta^{\star} \text { is not } \\
& \text { small relative to the duct diameter so duct pressures are sensitive to changes in Reynolds number. } \\
& \text { Consequently, a drop in Reynolds number will show up in the average PT. The resultant change to the } \\
& \text { local Mach due to boundary layer growth must also show up in the the static pressures. So it follows that } \\
& \text { a calibration at a single Reynolds number has broad applicability without any corrections for Reynolds } \\
& \text { number. }
\end{aligned}
$$

[^37]Quno We Need to Calibrate At The Reynolds Number We Test?

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The following figure shows the derivation of the first order Reynolds number correction for a axisymmetric nacelle.
Q Analytical Correction Available to Match FSC Data to Test Facility

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EINE February 1999 HSR Airframe Technical Review
The following figure shows the derivation of the first order Reynolds number correction for' 2 2-D nacelle.
ical Correction Available to Match FSC Data to Test Facility
Derivation of Standard Reynolds Number Correction For 2-D Nozzle

Quene February 1999 HSR Airframe Technical Review The following figure illustrates how Reynolds number matching can be achieved in the FSC. On a
standard thrust stand or in the nominal FSC arrangement, either the upstream or downstream pressure is
a variable, but never both. Consequently, the nozzle pressure ratio pressure is determined by a single
pressure, and Reynolds number is a fallout. In the FSC, PT is fixed at atmospheric and the pressure ratio
is varied by changing the downstream static pressure. On a thrust stand, it is PT that is the variable and
the downstream static pressure that is fixed at atmospheric pressure. By attaching the "Standard Thrust
Body" (STB) to the FSC upstream of the nacelle, uniform high pressure air can be supplied to the nacelle.
In the STB mounting arrangement, both the upstream and downstream pressures can vary independently.
The maximum pressure ratio is determined by a minimum tank pressure limit of 4 psi and a maximum
supply pressure of 300 psi.


The following picture is of the STB installed on the FSC. Note the high pressure air lines entering the STB can-like objects on right), to minimize the potential of a momentum tare.

The following plot shows the operating envelope of Mach and Reynolds number with the STB installed. It
is possible to expand the envelope with some re-engineering of the hardware to expand the
minimum/maximum tank pressures or supply pressures.
FSC/STB Operational Envelope

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The following is a picture of the STB as viewed looking upstream. In the foreground is the
instrumentation section that is placed between the STB and the test article. It consists of 4 PRT's to
$\rightarrow 5$

## FSC Basic Principles: Upstream Instrumentation

For Variable PT/RE Test Rig
$\rightarrow$

- Basic


The results from using the STB were mixed. The good news is that, in general, the set up appeared to work. Comparison for the same nacelle between the plate mount and the STB indicate that the STB data fall within the long-term scatter of the plate mount data. To check for a bias with PT, the cubic calibration nozzle on the plate mount was compared with the cubic on the STB at two Reynolds numbers. The cubic calibration nozzles are designed to be insensitive to Reynolds number. At both Reynolds numbers the STB data fall within the long-term scatter of the plate mount data.

The bad news was that the use of the STB degraded the short term precision two fold. This was because it was almost impossible to set PT precisely and hold it constant for more than a few seconds. At the same time, the nacelle is unstable because it's throat is highly Reynolds number dependent. In addition to the precision penalty, productivity was also worse by a factor of two to three. This was because set up times are longer, setting PT was harder, and one naturally wants to run more Reynolds numbers if the STB is installed. The net result is that if only atmospheric data are all that is required, using the STB could increase the test time by a factor of 6 and offers no addtional benefits.

The other curious result was the assessment of Reynolds number effects. Several configurations were run at 4 Reynolds numbers. All runs always fell within the long term data scatter of plate mount data. However, the visable trend indicated that drag was either constant or went up with increasing Reynolds number. This is opposite of what CFD would predict. There are several potential explanations for these results. It could be just scatter, or it could be all the data are corrupted by the bellmouth design (discussed later), or it could be that there is some bias with PT we failed to diagnose, or it could be a real effect. Running the STB presents a real paradox. While the technique is necessary to simulate variable Reynolds Number, the reduced precision makes it harder to measure an effect.

The following figure illustrates how there is statistically no difference between platemount and STB mount data for nacelle N31.3.
Assess facilities ability to match Reynolds number (from 16ft to Polysonic)
Good News: At Same RE, STB Data Within Plate Mount Repeatability

Data based on same instrumentation: MCV4 (aft tank venturi) and TTN (upstream TT)
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A cubic calibration nozzle is designed to be insensitive to Reynolds number. Verification of the design has historically been done by comparing FSC data and Thrust Stand data. Within experimental precision limits, the two data sets have agreed despite a factor of two difference in Reynolds number ( $\sim 4$ million vs 7 7million). To verify that the STB was not imparting a bias with PT setting, a 2" diameter cubic
calibration nozzle was run at two PTs (14.7 \& 23psi). The upper and lower bound of the data scatter for these two cases as well as a plate mount run are shown in the following figure. These data indicates that all data, regardless of mounting system or Reynolds number, produce the same results (within precision limits), and hence there is no system bias with PT. It should be noted that these data also illustrate the reduced precision of STB data.
Assess facilities ability to match Reynolds number
(from 16ft to Polysonic)
Good News: At Same RE, STB Data Within Plate Mount Repeatability


The following figure shows the long term precision limits for the STB mount. Three installations on the STB are within the limits of the plate mount. These data would imply that the STB degrades the short term precision by a factor of two, but also that the effect on long term precision is negligible.

Data based on same instrumentation: MCV4 (aft tank venturi) and TTN (upstream TT)

There are two interesting issues with the STB data that may be contributing to the precision losses. The
 tank total measured at the tank exit. This is shown in the following left hand plot. The flow through the nacelle is believed to be adiabatic and for the plate mount/free air cases the instrumentation seems to substantiate this. When the STB is installed the data shows a 3 degree temperature drop from upsteam of the nacelle to downstream.

The second issue concerns the different readings of mass flow rate from the critical flow ventrui upstream of the STB/nacelle (MCV1), and the one downstream of the nacelle/tank (MCV4). The maximum scatter of MCV1 and MCV4 is shown in the following right hand plot. The reduction in short term precision seems to be caused by system instability. The complete system consists of 3 choked venturi's. The first is MCV1. Then there is a supply line that is approximately 5" diameter that runs $\sim 12$ ' from MCV1 to the STB. For the Mach numbers of interest, the nacelle itself forms a choked venturi as its exit area fluctuates with

Reynolds number based on upstream PT. Aft of the nacelle is the 4' diameter, 12' long settling chamber. At the exit of the chamber is MCV4, the primary measurement critical flow venturi. In practice, we found it very difficult to set and hold the supply line PT upstream of the STB for more than a few seconds. At the low mass flow rates of HSCT nacelles (.2lbs/sec), there was a lag between when PT was set and when it equalized at the rakes in the STB. This upstream PT is what would determine boundary layer characteristics through the nacelle and therefore the mass flow rate through the nacelle. There was also a lag between when the STB PT equalized and the when the tank ambient pressure would equalize. The tank pressure forms the PT for the MCV4 venturi. At any given moment, the system could conceivable see 3 different mass flow rates: MCV1 (upstream venturi), nacelle mass flow (felt by balance), and MCV4 (downstream venturi). Because the volume of air in the supply line between MCV1 and the STB is
approximately 23 times less than that of the tank, it stands to reason that MCV1 should track the mass flow through the nacelle more closely and be less sensitive to supply pressure changes. However, MCV1 has never been as fully calibrated as MCV4 has for data accuracy.
STB's Reduced Precision is Caused By System Instability
DATA QUALITY ISSUES RUNNING WITH THE THRUST BODY
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The following figure illustrates the difficulty in interpreting the Reynolds number trends from the STB. Taken at one level, this plot says that any variation in drag with Reynolds number is within the long term data scatter and is therefore meaningless. On another level, it is recognized that the Reynolds number variations are taken on a single installation over a very short period so it is the short term precision that should be viewed for comparison. If this second case is closer to reality, the data indicate the gross thrust per unit of weight flow actually increases as Reynolds number is decreased.
 Depending on whether or not it is short or long term precision that is guiding the interpretation, a
completely opposite trend with Reynolds number is indicated by the data. Shown in the left hand figure is the trend with Reynolds number extracted from a linear theory code (H349) and OVERFLOW NS. Both indicate that a Reynolds number increase form 4 to 7 million/ft will produce a drag reduction of .4cts per/pair (2 nacelles). If the STB Re variation data is just scatter, then it can be assumed that the nominal drag at 4 million corresponds to the plate mount data. This data can then be adjusted using the 1 st order Reynolds number correction described previously. The resultant data also indicate that drag goes down by . 4 ct for a 4 to 7 million/ft Re change. On the other hand, if the STB data are good, then the trend actually is a drag increase of . 4 cts going from 4 to 7 million/ft.
If the STB data are not scatter, they could be caused by two effects. The first is a bias with PT. To change Reynolds number requires an increase in the upstream PT. If there is a systematic change in leak rate, pressure tare (through high pressure line), weight flow measurement or balance accuracy with change in PT, a misleading trend could be produced. All the checks available and the 2 " cubic runs seem to rule out the possibility, but something could have been missed.

[^38]
N31.3/N32.3 Flow Thru Nacelle Trip = .008" losses per unit of mass would actually be less.
The following figure provides some supporting data that this may actually be happening, and therefore the trend seen by the STB with Reynolds number might "real". Shown is a comparison between data from the same model with two different trip definitions. A large change in trip can dramatically change the boundary layer profile. In this way, a change in trip is very similar to a change in Reynolds number. Both configurations were run with the plate mount at a nominal Re=4 million/ft. Using this configuration, it can be said that these data have no hidden bias with respect to PT associated with the STB mount. One trip is the nominal trip or $.008^{\prime \prime}, .4$ " from the inlet. It was sized for BSWT at $\operatorname{Re}=7 \mathrm{million} / \mathrm{ft}$, so it is probably under-tripped for the FSC. Consequently, the duct most likely has a substantial length of laminar flow. The second trip, designated the super trip, has \#40 (.0165 in.) grit filling the space from the inlet leading edge all the way back to the nominal trip. There is no chance that this configuration has any laminar run. Moving left to right, the first chart shows that increased trip effectiveness reduces the weight flow (consistent with a large boundary layer). The second chart shows that supertrip is increasing thrust loss (Inlet thrust-exit thrust). Both charts are as one would expect. However, the third chart shows thrust loss as a function of inlet momentum. The supertrip case actually has less of a loss per/count of inlet momentum. The fourth chart shows exit momentum in counts versus inlet moment in counts. The fifth chart shows the same data as the fourth in lbs of thrust versus lbs of weight flow. These data clearly indicates this configuration has less reduced losses per lb of mass flow with increasing boundary layer. Therefore this configuration should show a reverse trend with Reynolds number when viewed by some variables.
There are two possible sources of thrust loss. There are PT losses in the boundary layer and flow nonuniformity in the core flow (pressure drag). Transonically, it is believed that the dominant effect for HSCT nacelles is boundary layer losses and that the core flow is fairly uniform. CFD solution have been
consistant with this belief. It follows that a decrease in Reynolds number will result in a larger boundary layer, hence larger losses, with little effect on flow non-uniformity. To have a reverse trend with Reynolds number would require that the flow non-uniformity losses be greater than the boundary layer losses per unit of mass flow. Therefore, as more boundary layer fills the duct there would be less core flow, so the
Variable Trip Data Suggest That The Reverse RE Trend May Not Be
Due to A STB Bias

N31.3/N32.3 Flow Thru Nacelle Trip = .008"
February 1999 HSR Airframe Technical Review
One fortuitous feature of the mounting design to be used with the STB was the use of an adapter to hold the bellmouth. In the nominal plate mount arrangement, the nacelle is mounted to the plate and the bellmouth to the nacelle. This was how the nacelle was mounted in the first test. With this arrangement, there is no easy way to run an isolated bellmouth without making a new plate. When the STB is used, an adapter is mounted in the plate, the bellmouths are attached to the adapter, and the nacelle is mounted off the bellmouth. It is relatively easy to run an isolated bellmouth. When the trends with Reynolds number appeared, one diagnostic was to look at an isolated bellmouth. The following plots show the isolated bellmouth results.
Moving from left to right, the first plot shows the balance based velocity versus ideal, the second shows the venturi based weight flow versus ideal, and the last show the thrust loss. The first thing that was apparent for these results was that no one had imagined that such a bad bellmouth design was possible. Instead of getting $>99 \%$ of ideal velocity and weightflow, this bellmouth was at $96 \%$ or worse. Instead of .8 cts of drag it had up to 2.8 cts . As mentioned earlier, past designs have been very successful. However those designs were all axisysmetric in nature. Other non-axisymmetric designs have been built ( F -22,
JSF), but those designs have not dhe calibration. This requires that there is sufficient instrumentation in the duct that can pick up the bellmouth induced effects and not include them when they do not appear in a wind tunnel test. There was a concern that with the limited static pressures in the HSCT nacelles, they would be unable to distinguish bellmouth induced effects from the nominal duct effects. To assess this, an isolated bellmouth was run post-test in OVERFLOW NS.
February 1999 HSR Airframe Technical Review The following figure shows a post-test OVERFLOW solution at an equivalent Mach=0.8. The PT contour plot indicates that there are no large PT losses so the drag of the bellmouth is associated primarily with distortion and not turbulent flow. The pressure and streamline data indicate that while the sidewalls are accelerating the flow uniformly, the three dimensional nature of the corners is producing a hyper acceleration. The result is a highly distorted flow that is substantially off ideal. Two other important facts are evident. First, the distortion is such that it will invariably propagate aft. The second is that the nacelle static pressures have little ability to compensate. This is because primary instrumentation at the charging station consists of only 4 nacelle static pressures located mid-wall on the four sides. The location and amount of instrumentation was selected based on transonic CFD solutions on the installed nacelle. These data indicated the duct static pressure profile to be very uniform. For such a scenario, a static pressure rake would be overkill. It would mostly cause more distortion than it would resolve.
For the calibration, instead of the bellmouth providing uniform flow at $M=.8$ and the pressures reading $M=.8$, CFD indicates the bellmouth is producing thrust more consistent with ideal flow at $M=.75$ with the pressures seeing $M=.82$. Clearly this has to introduce some bias into the system. With this much distortion it is possible that more gross thrust can be generated in the turbulent boundary layer than in the core flow. It is also possible that due to the boundary layer effects at the duct exit, the increased convergence makes the flow more uniform. In either case, this uncertainty makes it impossible to determine whether the reverse Reynolds number trends produced by the STB were the product of the bellmouth, data scatter, or some real effect.

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As mentioned earlier, subsonic transport programs apply Reynolds number corrections to calibration data where as fighter programs do not. Therefore, one of the principal objectives of this test was to verify whether some form of Reynolds correction would be required for an HSCT. With the resulting uncertainty in the STB data, the reverse trends, and the bellmouth issues, one might wonder if this entry got us any closer to knowing whether a Reynolds number correction is necessary for HSCT nacelles? The answer is yes. The main reason subsonic transports require a correction while fighters do not is that the instrumentation in a fighter flow-through nacelle tends to be sensitive to boundary layer changes.
The following charts show how trip effectiveness changes the static pressures inside an HSCT nacelle. Moving from left to right, the first chart shows the ratio of charging station static pressure to free stream static, the second shows the ratio of exit pressure static to free stream static, and the third shows the ratio of the two. For a nacelle where the boundary layer effects are small relative to the core flow, these three ratios should remain constant for the two trip definitions. If these ratios remained constant, a Reynolds number correction would be a requirement. There was a measurable change and because of this sensitivity, the static pressures become good indicators of changes in local Mach. This relation is shown in the last figure.
Nominal Plate Mount Data Should Be applicable Over Broad Re Range

The following figure illustrates how a "supersonic" calibration works. Pictured on the left is the mounting arrangement (above) and internal flow characteristics (below) of a nacelle in the wind tunnel. Pictured on the right is the mounting arrangement (above) and internal flow characteristics (below) of a nacelle in the
 thrust and weight flow characteristics with the nacelle internal flow. The nacelle exit flow (exit Mach) is a function of only the upstream flow conditions (Inlet local Mach). It is not necessary to simulate the entire installed system to reproduce the same exit flow. All that is required is that the control volume just upstream of the inlet is reproduced. If the effect of plume shape on gross thrust is also required, the static pressure downstream of the exit must also be reproduced. The FSC does this by applying a
bellmouth, followed by the nacelle, up to the nozzle exit. In this way entire nacelle internal flow aft of the inlet is simulated.
FSC Test Set Up For Supersonic Flow
Q botina


The following figure illustrates how supersonic flow is simulated in the FSC. The overall test setup is
identical to the variable Reynolds number transonic set up. Again the STB is used. By replacing the
converging bellmouth with a converging/diverging bellmouth, supersonic flow can be reproduced at the
inlet highlight.

February 1999 HSR Airframe Technical Review On one level, this entry was a success for verifying that supersonic bellmouth design is possible and
variable inlet Machs can be simulated. Inlet Machs of $1.6,1.8,2.1$, and 2.4 were simulated at 3 Reynolds
numbers a piece. Unfortunately, the same problems associated with the square bellmouth apply to the were all tested isolated and showed the same drag as the isolated subsonic. The supersonic bellmouths were fitted with static pressures just upstream of where the inlet face would be. These pressures were to be used in the data analysis to define the inlet Mach of the nacelle. What was found, as shown in the following figure, was that the wall static produced an indicated Mach very close to the ideal design Mach of each bellmouth. If the bellmouth was designed to supply $M=2.4$ flow, the statics indicated a mid-wall Mach of 2.399. However, the average Mach, as indicated from the manipulated balance data was substantially lower (maybe 2.25). This is a clear indication of the cornerflow distortion seen of the subsonic bellmouths propagating aft and it staying in the corners.
8 COEFNF
Pressure Ratios and Supersonic duct Flow
 vs. RN trend
Assess Ability to Achieve M=2.4


- Reynolds numbers per/Mach Issues:
- Static's indicate near ideal flow
- Balance \& venturi data indicates


February 1999 HSR Airframe Technical Review
Results Summary:
The second FSC test set had three goals: 1) Define long term repeatability, 2) Explore Reynolds number effects, and 3) Explore supersonic duct flow. The long term repeatability appears acceptable. The results of Reynolds number effects runs were inconclusive. The data from this test consistently showed a small reverse trend with Reynolds number. This effect could be real or it could be caused by random data scatter, a systematic bias for the STB mount with PT, or interactions with a bad bellmouth. Results from the 2"cubic nozzle appears to rule out bias, and plate mount runs with two trips implicate the bad bellmouth as a culprit. Unfortunately, there are insufficient data for a conclusive diagnosis. Productivity rates and precision limits on the STB mount indicate that the STB should be used only as a last result if no other way of correcting for Reynolds number effects is available. The two trip plate mount data indicated that the internal statics are responsive to changes in boundary layer profile. This implies that an HSCT calibration at a single Reynolds number (plate-mount) should have broad applicability across Reynolds numbers (similar to fighters).
The supersonic bellmouth work was also inconclusive. The data seemed to imply that the general experimental approach is feasible. Unfortunately, the distortion in the basic bellmouth compromised all the drag results. The basic bellmouth produced a vortex down the corners that is unrepresentative of wind tunnel conditions and cannot be picked up by simple internal instrumentation.
The first solution to most of these issues is to design a new bellmouth and verify the design in CFD before any new testing. Bellmouths are inexpensive $<\$ 2 \mathrm{~K}$, and it is believed round inlet bellmouths with short transition to the square inlet is all that is required to clean up the flow. The supersonic calibration should also be fixed with this new bellmouth design ( $<\$ 4 K$ ). Based on these data, the recommendation for the PIE test is to calibrate nacelles using the plate mount method. Follow-on CFD and uncertainty analysis would be useful to determine "best" data reduction scheme.
February 1999 HSR Airframe Technical Review This test was one of those tests when you find out that you did not even know what you did not know. In pretest discussions about risk, the bellmouth design was never considered a big issue. Now this appears to be the primary obstacle to accurate data. On the plus side, several successful CFD runs have already been made on low distortion bellmouths that fit our square inlet. Another lesson learned was to run the FSC with only one ejector. The FSC has two ejectors to suck down the tank. The low flow rates on our nacelles only require one to be run to achieve the appropriate tank pressures. In principle one or two ejectors should have no effect on data quality because the ejectors are downstream of the last MCV and
 accurate data acquisition difficult. The effect can be subtle, but it can reduce repeatability and impart biases to the balance load. In one case, in the beginning of the test, a faulty balance lead produced good data with low vibrations during checkloads and zeros, only to act up while the system is running with two ejectors.

[^39]Conclusions From HSCT FSC Entries
Lessons Learned - Square bellmouths are BAD ! BAD! BAD!

- It is possible to make a bellmouth so bad that it effects
calibration
- HSCT nacelles should always be run with only one ejector
(reduce vibrations)
- There is an unresolved bias in TTN vs. TTVC with the STB
- When using the STB, weight flow from upstream (MCV1) is
twice as precise as from tank (MCV4) but may introduce a
bias
February 1999 HSR Airframe Technical Review
Lessons learned-continued
If one truly wants to understand Reynolds number effects, the STB still offers a unique capability. To fully utilize that capability requires confidence that all potential bias issues have been resolved. To that end an "HSCT" calibration standard with known Reynolds number characteristics could flush out bias issues between CFD and the FSC. Such a standard might consist of a $1 \mathrm{in}^{2}$ cubic bellmouth followed by $\sim 11$ " Lastly, while the rakeless/single station data reduction scheme appears to work well, it remains unproven for all flow scenarios, and therfore, it is recommended that future designs continue to have static's at 2 stations.

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[^0]:    *Part 1 is presented under separate cover.

[^1]:    The important feature of this approach is that a set of simple geometric entities lies at the core of the entire design process. This technique retains the typical way in which
    aerodynamic vehicles are defined and provides strict control over how the surface
    intersections are treated. Furthermore, since the chosen design variables act directly on the
    geometric entities, these entities may be output for future analysis at the end of the design
    process. process.

[^2]:    Finally, the WARP-MB procedure is used to morph the CFD volume mesh to follow
    the updated surface meshes producing a volume grid to be used in calculating the adjoint and flow solutions.

[^3]:    Additionally, PERTURB is responsible for the imposition of constraints to the geometry entities. Examples of implemented constraints include the thickness for spars of hinges, cabin radius, fuel volume, spar carry-through, and spar straightness.

[^4]:    into
    geometry be compared with an earlier slide
    of surface meshes fr depicting the baseline mesh.

[^5]:    This figure is the Configuration Aerodynamics Program on a Page. As shown, the work reported here falls under both Cruise Point and Multi-Point Optimization.

[^6]:    Gross thrust as objective function
    -•

[^7]:    - Mach $=2.4, C_{L}=0.092$
    - 437 variables, 19 on bounds
    - 34 k constraints, 237 active
    - 4 TRANAIR cycles
    - 7.6 counts of thrust improve
    - Mach $=2.4, C_{L}=0.092$
    - 437 variables, 19 on bounds
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    - 437 variables, 19 on bounds
    - 34 k constraints, 237 active
    - 4 TRANAIR cycles
    - 7.6 counts of thrust improvement
    - Mach $=2.4, C_{L}=0.092$
    - 437 variables, 19 on bounds
    - 34 k constraints, 237 active
    - 4 TRANAIR cycles
    - 7.6 counts of thrust improve


    ## - Inviscid analysis

    - Relative to trimmed baseline
    - Shape changes similar to TCA optimization

[^8]:    The drag reduction of the initial trimmed cruise point design is
     the NCV, about one count of additional drag解 the TCA cruise point design, rather than an effect of including trim.

    The first bar from the left is the Ames 103 design. Of interest are both the total and the nacelle/diverter increment. Both are comparable to that achieved with this latest design. The larger nacelle/diverter increment, relative to the NCV , is one of the features of the Ames 103 which was addressed by the process and technology improvements
    mentioned above.

[^9]:    The corresponding lower surface view of the trimmed cruise point design is shown on this figure. The most apparent features in this view are the wing lower surface troughs associated with the nacelles and the bump between the troughs. These features are also present on the NCV . The sculpting present around the nacelle/diverters is not well resolved by this image.

[^10]:    As with the cruise point design, the thrust improvement shown is relative to the trimmed baseline, and based on inviscid TRANAIR analysis. As expected, the shapes produced by this case are not similar to those found in the cruise point case.

[^11]:    - 2 viscous demonstrations (N-S OVERFLOW)
    - 2 condition TRANAIR shakedown problem
    - Initial 3 condition TRANAIR problem

[^12]:    - 445 variables ( 29 on bounds) \& 34k constraints ( 214 active) - 3 TRANAIR cycles
    - 7.45 counts of thrust improvement
    - 7.4 counts at Mach 2.4
    - 0.3 counts at Mach 0.9
    - Shape changes similar to cruise point design
    - Tri-surface TCA configuration
    - Mach $2.4 / C_{L} 0.092$ \& Mach $0.9 / C_{L} 0.18$
    - $6: 1$ weighting in objective function

[^13]:    The initial three condition TRANAIR multi-point optimization case, outlined on this figure, is essentially ready except for the trimmed baseline starting points. The common confrguration, Ti-sp.

    Once fully underway, this problem will represent the multi-point case on which the final evaluation will be made. More details about the evaluation process, and how this case fits in the process, are outlined in the concluding section of this presentation.

[^14]:    - Tri-surface TCA configuration
    - Mach $2.4 / \mathrm{C}_{\mathrm{L}} 0.089$, Mach $0.9 / \mathrm{C}_{\mathrm{L}} 0.18$ \& Mach $1.1 / \mathrm{C}_{\mathrm{L}} 0.15$
    - 1:0.3:0.15 weighting in objective function
    - 451 variables \& 35 k constraints
    - Waiting for trimmed baseline starting points
    - Will begin with initial Mach 1.1 sensitivities

[^15]:    - Sensitivities to active constraints
    - Coupled viscous TRANAIR optimization
    - Other TRANAIR multi-point problems
    - Sensitivities to active constraints
    - Coupled viscous TRANAIR optimization
    - Other TRANAIR multi-point problems

[^16]:    Nonlinear aerodynamic shape optimization is essential in designing an economically viable High Speed Civil Transport (HSCT).
     time remains a challenge. In a typical optimization, significant cost is incurred in obtaining the gradient of the cost function, especially when finite-differences or sensitivity equations are used.

    It is well known that, by using the adjoint method, the gradients of hundreds of design variables can be computed accurately and efficiently. With design cycle-time reduction as a goal, an adjoint code was developed for TLNS3D (Euler), a versatile flow analysis code. Results of the aerodynamic shape optimization of HSCT Technology Concept Airplane (TCA) configurations, performed using this code, are presented in this paper.

[^17]:    The convergence of the analysis (Euler) and the adjoint equations, for the TCA wing/body/flaps configuration at $\mathrm{M}=0.9$, are shown. For a single-block grid with $13 \times 33 \times 89$ points in the chordwise, norma, and spanwise diretions, respectiv, he rady and equations are comparable. Note that, the solver parameters like the CFL number, the Runge-Kutta coefficients, and the dissipation coefficients, used for the solution of for both these systems, are identical.

[^18]:    Due to the efficiency and versatility of TLNS3D-Adjoint in computing aerodynamic gradients, it is being used in other Boeing programs. It has been successfully used to improve the performance of a space vehicle. This lifting-body type configuration was optimized at Mach 1.2. The cost function was a function of the pitching moment.

[^19]:    The ADIFOR-based methods, on the other hand, are straight-forward and are applied directly to
    the differential form of the equations of motion already available in the form of FORTRAN codes. The
    resulting computer-generated differentiated form of the existing function computer codes provide
    accurate analytical sensitivities. Also, the ADIFOR and ADJIFOR application procedures are the same
    for both the Euler and Navier-Stokes codes. This, it can be seen that the automated procedure for
    generating adjoint codes using ADJIFOR has several advantages over the hand-differentiated adjoint codes.

[^20]:    Several codes have been ADIFOR-processed to obtain sensitivities. These codes have varied sizes and applications. The ADIFOR software does not require large user intervention from function code to derivative code.

[^21]:    N: No. of DVs; C: CPU time for 1 analysis; M: No. of Grid points, $\lambda$ : words/grid point $\alpha, \beta$ : Savings factor due to the use of previous solution as initial guess
    $\gamma: \quad$ Cost of evaluating one gradient; $\sigma$ : for CFL3D, approx. $=20$ to 25

[^22]:     the Euler gradients of each design variable for a $0.48 \times 10^{6}$ point Euler grid. However, when 50 design
    variable sensitivities are computed simultaneously, the CPU time required is only 63 times the function sevel $1 i \mathrm{c}$. This is dual
     number of function evaluations in each processor that results in significant CPU savings. The use of large number of processors can also result in fast turn-around time.

[^23]:    drag values are most likely too low for these configurations by a few counts of drag.

[^24]:    Trim Conditions, $x_{C G}$ variation
    A parametric study was performed to evaluate the trim drag benefit of the TCA W/B/N/D/E configuration with thrust as a function of the axial CG location, $\mathrm{x}_{\mathrm{CG}}$. Trim drag is the difference in drag between the trimmed configuration and the configuration with no lift on the horizontal tail. Moving the CG aft reduces angle-of-attack, increases the tail incidence ( $\mathrm{i}_{\mathrm{H}}$ ), and reduces the thrust and trim drag (CDrim). At the aft-most CG location, the trim-drag benefit is 1.3 counts. This is close to the 1.0 count that the TI team has bookkept.

[^25]:    1. Kulfan, Robert, Configuration Aerodynamics Metrics Update, HSR Aerodynamics Performance Workshop Proceeding, February

    Reuther, J., Jameson, A., Farmer, J., Martinelli, L., Saunders, D., Aerodynamic Shape Optimization of Complex Aircraft
    3. Kuruvila, G., Sundaram, P., Hager, J.O., Aerodynamic Sensitivities Using Three Methods, HSR Airframe Technical Review Proceeding, February 9-13, 1998.

[^26]:    CONFIGURATION
    -

     | 0 |
    | :--- |
    | 3 |
    | 3 |
    | 3 | SYMBOL

[^27]:    geometry becomes available.

[^28]:    The computed laminar flow corrections are influenced, of course, by the definition of the
     likely to vary with lift. Therefore, strictly speaking, the laminar flow corrections shown here would only apply in the immediate vicinity of the given lift coefficient.

[^29]:    Presentation Outline:
    geometry and test highlights will be presented first. This will be followed by a summary of the performance results of the test, which will include comparisons with CFD and flow viz. The third section will be an analysis of the stability and control results. A summary will be given at the end of each data analysis section.

[^30]:    some limited UV and colored oil runs were made.

[^31]:    This review indicated that there were some commonly held misconceptions as to what the definition of internal drag was and what it meant relative to aircraft performance. The understanding among many people was that internal drag was the sum of all nacelle internal forces and that to correct wind tunnel model or CFD model forces, one simply removes the forces predicted on the internal surfaces. While this may be a useful approximation for a select few scenarios, for most cases, it is a gross oversimplification. Internal drag is negative net thrust. Because the Propulsion group never runs installed propulsion systems, the actual definition is a complex arrangement negotiated by the T-D document. In the case of aircraft performance, it turns out that it is more important that thrust and drag are assessed consistantly than "accurately". For example, if a T-D system lumps nacelle skin friction in with thrust, it decreases the thrust available per lb/fuel, but it also reduces the thrust required the same amount. If on the other hand neither staff includes nacelle skin friction, but calculates all other forces very accurately, a performance short fall will result.

[^32]:    Transonic calibration has been has been the method of choice for internal corrections on most major programs for several reasons. It's cost is reasonable and it is quick. Bellmouths can cost less than $\$ 2,000$. Test time can be as reasonable as $\$ 4000 /$ geometry and take as little time as a single shift. Such a calibration covers the nacelle for all Machs/angle of attack. Most importantly, the results are consistant: They are consistent with the as-fabricated part and can be completely consistent with the thrust methodology and how Propulsion calibrates nozzles.

    > Stokes cases (4 alphas X 6 Mach numbers). With transonic convergence times running at 100hrs/run and $\$ 40 / \mathrm{hr}$ the cost for computer time alone might be as high as $\$ 96,000$. Gridding and shepherding solutions through the queue could take months of flow time. Finally, the solution will not necessarily be consistent with the as-built geometry.

[^33]:    Calibration cost must be less $\$ 30 \mathrm{~K} /$ geometry
    Instrumentation cost must add less $\$ 15 \mathrm{~K} /$ nacelle
    Use of calibrated nacelles cannot significant reduce in-tunnel
    productivity
     Calibration error (precision + bias) must be $< \pm .1 \mathrm{ct} / \mathrm{nacelle}$ for
    $\mathrm{M}=2.0-2.5$
    Calibration error (precision + bias) must be $< \pm .3 \mathrm{ct} / \mathrm{nacelle}$ for
    $\mathrm{M}<2.0$
    Data reduction process should be simple enough to be repeatable
    by any HSR researcher
    -
    $\bullet$ $\qquad$ --
    -

[^34]:    F
    둗
    PT, Pamb and TT are known, Fg and Wdot are also known
    to calibrate measured values of Pe, PT, Pamb
    to measured values of Fg and Wdot
    Goal of FSC is $t$
    Conclusion: If $\mathrm{Pe}, \mathrm{P}$
    

[^35]:    For a long time propulsion staffs have known that it is not necessary to actually calculate exit gross thrust. In the first place, the calculation will invariably not agree with the test data (instrumentation bias, non-uniform flow, etc...). In the second place, actual thrust can more easily be approximated by some power series function of nozzle pressure ratio (PTnozzle/Pamb). This effectively makes static pressure instrumentation obsolete. The following chart illustrates that a function of NPR and direct calculation are one and the same.

[^36]:    Follow-on CFD and uncertainty analysis is needed to "best" data reduction scheme

[^37]:    So which is right for an HSCT? We have no reason to suspect that an HSCT should behave differently than a fighter model. But what if it doesn't and a correction is required? How would we go about developing such a correction? On an HSCT nacelle, skin friction changes by about .2cts per/nacelle going from 4-7million. The accuracy capability demonstrated in the first entry indicated that we should be able to resolve effects down to $\pm .08 \mathrm{cts}$. In addition, the Boeing FSC has the hardware available (in the form of the "standard thrust body") to test at a wide range of Reynolds numbers. Therefore the easiest way to resolve the uncertainty is to run the test.

[^38]:    A second possibility is that the effects are real, but there is something going on in the calibration that is not going on in CFD. One such feature is the bellmouth design. This will be discussed later, but the bellmouth is imparting an unusually high drag bias to the data (relative to historical trends for cubics). If the bellmouth imparts "too" much distortion, adding boundary layer (by reducing Reynolds number) could actually reduce drag. The data in the next figure seems to support this theory.

[^39]:    ! 1 (จヘJW) recommended that MCV4 remains the primary measurement of weight flow until more comparisons to MCV1 can be made.

