



A LATTICE BOLTZMANN METHOD FOR TURBOMACHINERY SIMULATIONS

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October 2002



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Outline

- Lattice Boltzmann Method
- Objectives
- Current LB model
- Simulation of cascades and result
- Parallel computing and result
- Conclusion remarks



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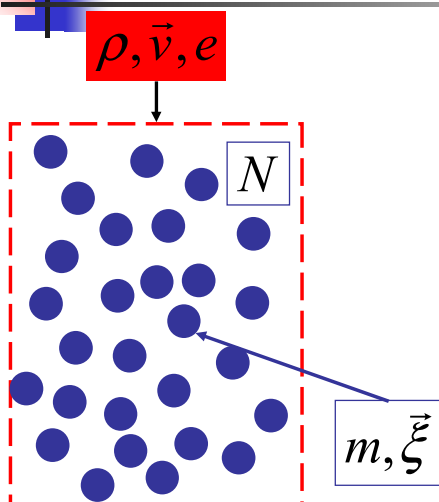


Introduction

- Lattice Boltzmann (LB) Method is a relatively new method for flow simulations.
- The start point of LB method is statistic mechanics and Boltzmann equation.
- The LB method tries to set up its model at molecular scale and simulate the flow at macroscopic scale
- LBM has been applied to mostly incompressible flows and simple geometry



Statistic Mechanics

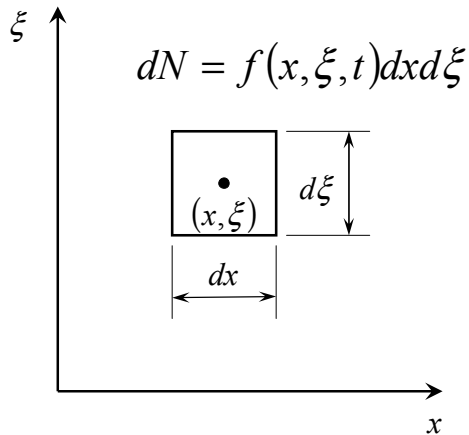


- Statistic mechanics views fluid as a collection of particles.
- The properties of the fluid are determined by the average properties of the particles in the collection

A small domain near point x

Boltzmann Equation

1) Distribution Function



- A phase space consists of both location and velocity of particles is introduced.
- The distribution function is defined as the density of the number of particles at point

(x, ξ)

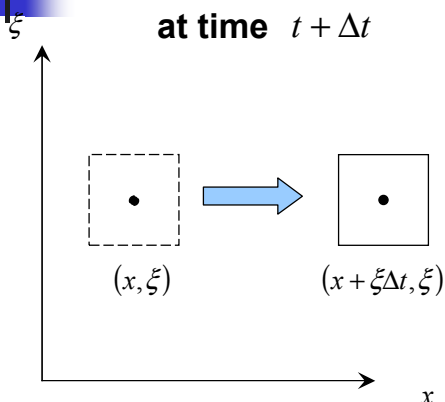
Phase space for 1D problem



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Boltzmann Equation

2) Time Evolution Of Distribution Function



- Assume that there is no external force and no collision between the particles, the velocity of the particles will not change

The two domain have the same number of particles

$$f(x + \xi \Delta t, \xi, t + \Delta t) dx d\xi = f(x, \xi, t) dx d\xi$$

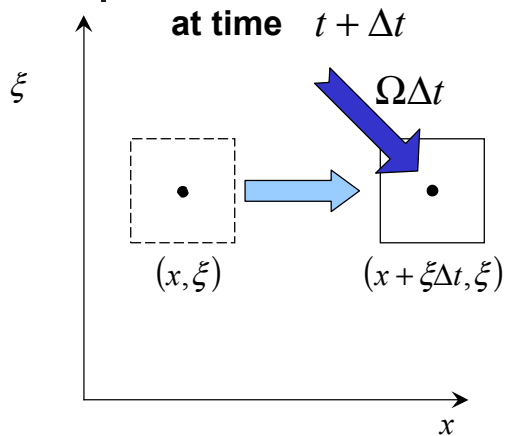
$$f(x + \xi \Delta t, \xi, t + \Delta t) = f(x, \xi, t)$$



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Boltzmann Equation

3) Collision Term



- Collisions between particles change their velocities, and make them move in and out of the domain
- A collision term describes the net increase of the density of the number of particles in the domain due to the collision

$$f(x + \xi\Delta t, \xi, t + \Delta t) = f(x, \xi, t) + \Omega\Delta t$$

Boltzmann Equation

4) BGK Collision Term

- One of the simplest collision models is the Bhatnagar, Gross and Krook (**BGK**) simplified collision model

$$\Omega = -\frac{1}{\tau} (f - f^{(eq)})$$

- The **BGK** model is widely used in **LB** models



Boltzmann Equation

5) Boltzmann Equation

- The Boltzmann equation in finite difference form:

$$f(\vec{x} + \vec{\xi}\Delta t, \vec{\xi}, t + \Delta t) = f(\vec{x}, \vec{\xi}, t) + \Omega\Delta t$$

- By Taylor expansion, the above equation can be written in the differential form

$$\frac{\partial f}{\partial t} + \vec{\xi} \cdot \vec{\nabla} f = \Omega$$

This is the Boltzmann Equation



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The Macroscopic Properties

$$Y(\vec{x}, t) = \frac{\int \eta(\vec{\xi}) f(\vec{x}, \vec{\xi}, t) d^3 \xi}{\int f(\vec{x}, \vec{\xi}, t) d^3 \xi}$$

- The Macroscopic properties are determined by the average value of properties of the particles

density

$$\rho = m \int f(\vec{x}, \vec{\xi}, t) d^3 \xi$$

momentum

$$\rho \vec{v} = m \int \vec{\xi} f(\vec{x}, \vec{\xi}, t) d^3 \xi$$

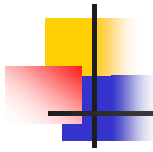
thermal energy

$$\rho \varepsilon = \frac{m}{2} \frac{D_f}{D} \int |\vec{\xi} - \vec{v}|^2 f(\vec{x}, \vec{\xi}, t) d^3 \xi$$



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The Enskog-Chapman Expansion

- It has been shown that the Euler equation and Navier-Stokes equation are the zeroth-order and first order approximations of the Boltzmann equation, respectively.,
- That is to say, the Boltzmann equation describes the fluid phenomena in a more accurate way than tradition fluid dynamics does

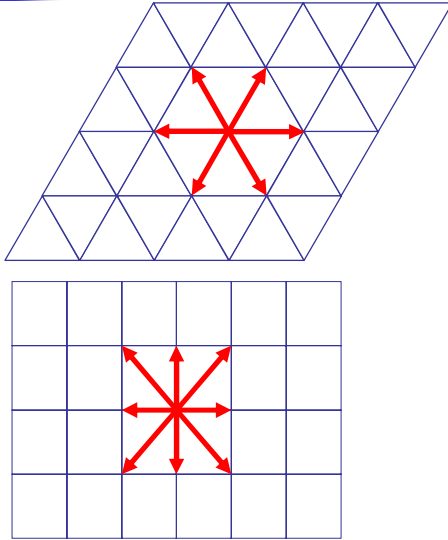


Lattice Boltzmann Method

- Lattice Boltzmann Method can be reviewed as a numerical method to solve the Boltzmann equation
- In LB method, the phase space is discretized.
- In a LB model, the velocity of a particle can only be chosen from a velocity set, which has only a finite number of velocities

Lattice Boltzmann Method

1) Velocity Set and Grid



- For the convenience of computation, the position space is discretized in such a way that the particles travel with one of the velocity in the velocity set will arrive at a correspondent node at next time step.

The Lattice Boltzmann Method

2) Macroscopic Properties

- In the LB method, the macroscopic properties are evaluated through the weight summation

$$Y = \int \eta(\vec{\xi}) f^{(eq)}(\vec{x}, \vec{\xi}, t) d\vec{\xi} = \sum_{\alpha} W_{\alpha} \eta(\vec{\xi}_{\alpha}) f^{(eq)}(\vec{x}, \vec{\xi}_{\alpha}, t)$$

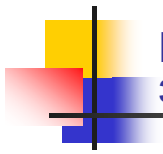
$$\rho = m \sum_{\alpha} f_{\alpha}^{(eq)}$$

$$\rho \vec{v} = m \sum_{\alpha} \vec{\xi}_{\alpha} f_{\alpha}^{(eq)}$$

$$\rho \epsilon = \frac{m}{2} \sum_{\alpha} |\vec{\xi}_{\alpha} - \vec{v}|^2 f_{\alpha}^{(eq)}$$

where

$$f_{\alpha}^{(eq)}(\vec{x}, \vec{\xi}_{\alpha}, t) = W_{\alpha} f^{(eq)}(\vec{x}, \vec{\xi}_{\alpha}, t)$$



Lattice Boltzmann Method

3) Time Evolution Equation

- Boltzmann equation in finite difference form with the **BGK** collision term

$$f(\vec{x} + \vec{\xi}_\alpha \Delta t, \vec{\xi}_\alpha, t + \Delta t) = f(\vec{x}, \vec{\xi}_\alpha, t) - \frac{\Delta t}{\tau} (f(\vec{x}, \vec{\xi}_\alpha, t) - f^{(eq)}(\vec{x}, \vec{\xi}_\alpha, t))$$

where $\vec{\xi}_\alpha$ is one of velocities in the velocity set

- Set $\Delta t = \tau$

$$f(\vec{x} + \vec{\xi}_\alpha \Delta t, \vec{\xi}_\alpha, t + \Delta t) = f^{(eq)}(\vec{x}, \vec{\xi}_\alpha, t)$$



The Lattice Boltzmann Method

4) The Calculation

$$Y(\vec{x}, t) = \sum_{\alpha} \eta(\vec{\xi}_\alpha) f_{\alpha}^{(eq)}(\vec{x} - \vec{\xi}_\alpha \Delta t, \vec{\xi}_\alpha, t - \Delta t) = \sum_{\alpha} Y_{\alpha}(\vec{x} - \vec{\xi}_\alpha \Delta t, t - \Delta t)$$

- The calculation of LB method can be viewed as following steps
 - 1) a node dividing its macroscopic quantities into parts
 - 2) the node sending parts of the macroscopic quantities to corresponding nodes
 - 3) a node receiving the parts of the macroscopic quantities sent by nearby nodes, adding them together and obtaining the macroscopic quantities for next time step



Advantages of LB method

- **Start from a clear and direct physical picture at molecular level;**
- **Algorithm is simple and straightforward;**
- **Natural parallel scheme;**
- **Easy to incorporate the physical phenomena at molecular level, possible of modeling fluid phenomena that can not be modeled with the traditional fluid mechanics.**



Motivation

- **There are very few reports about the successful LB simulations of the real life problems**
- **Successful extension of LB method to the simulation of turbomachinery can show the maturity of LB method and the promise of this method in the simulation of real life problem**
- **The parallel nature of LB method make it a possible high performance solver for turbomachine simulations**



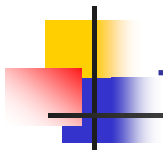
Challenges

1. **Most of LB models can only simulate the flow with a small Mach number**
2. **Most of past LB simulations are laboratory type simulations with simple computational domain and boundary**
3. **The proper LB model for simulation should be developed and necessary techniques should be developed**



The Current Work

- **Develop a compressible LB model.**
- **Introduce a boundary condition that allows accurate turbine blade simulation**
- **Develop mesh treatment for the irregular computational domain**



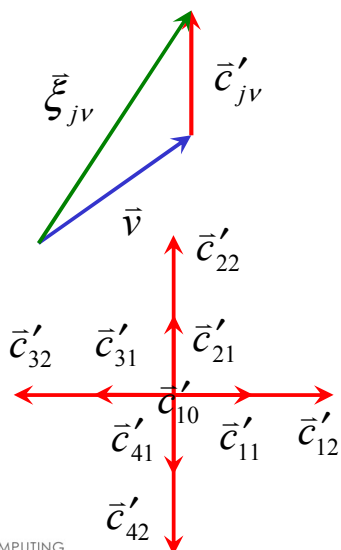
The Current Work

- Simulations carried out for three different cascades. It is the first time that turbomachines have been simulated by a LB model.
- The parallel performance of the LB model has been tested



Current Model

1) The Velocity Set

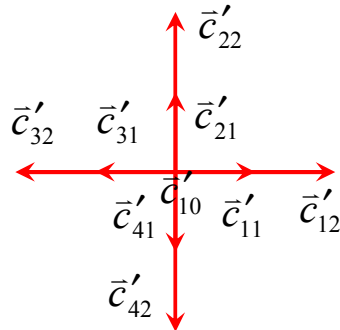


- A velocity in the velocity set consists of three parts.
- First the macroscopic velocity has been included explicitly into the microscopic velocity.
- Second, there is a set of diffusion velocity



Current Model

2) The Diffusion Velocity



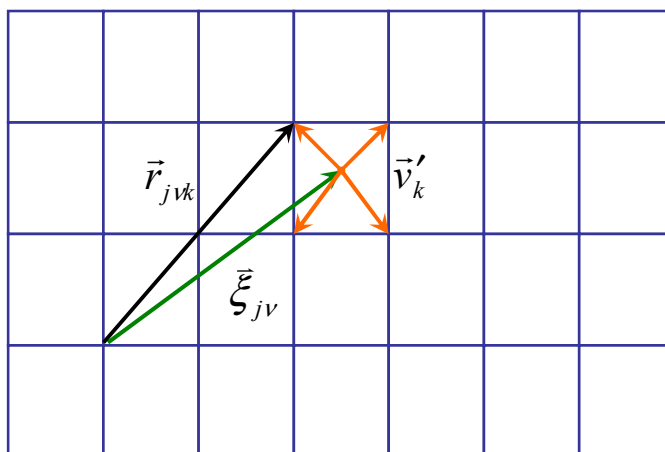
- There are 3 levels of diffusion velocity in current model
- The module of the diffusion velocities are determined by macroscopic quantities

$$c'_v = \begin{cases} 0 & \text{for } v = 0 \\ \text{int}(\sqrt{D(\gamma-1)\rho e/(\rho - b_0 d_0)}) & \text{for } v = 1 \\ \text{int}(\sqrt{D(\gamma-1)\rho e/(\rho - b_0 d_0)}) + 1 & \text{for } v = 2 \end{cases}$$



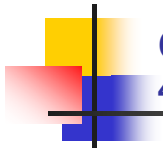
Current Model

3) The Velocity Set



- A third set of velocities \vec{v}'_k is introduced to carry the particles to the nearest vertex nodes

$$\vec{r}_{jvk} = \vec{v} + \vec{c}'_{jv} + \vec{v}'_k$$



Current Model

4) The Microscopic Quantities

- Mass, only one species is considered, is a constant

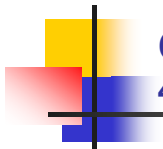
$$m = 1$$

- Velocity, consists of macroscopic velocity and a diffusion velocity

$$\vec{\xi}_{jv} = \vec{v} + \vec{c}'_{jv}$$

- Total energy is the same for all particles

$$\zeta = \frac{1}{2}v^2 + e$$



Current Model

4) Modification of Microscopic Quantities

- For the purpose of recovering the correct Navier-Stokes equation, a correction term has been introduced

$$\chi_{jvk} = \frac{\rho}{\rho - b_0 d_0} \frac{D}{2c_v'^2} (\vec{c}'_{jv} \cdot \vec{v}'_k)$$

- Mass

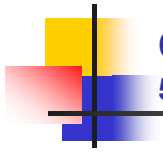
$$m_{jvk} = 1 - \lambda_{jvk}$$

- Velocity

$$\vec{\xi}_{jvk} = \vec{v} + \vec{c}'_{jv} - \lambda_{jvk} \vec{v}$$

- Total energy

$$\zeta_{jvk} = (1 - \lambda_{jvk}) \left(\frac{1}{2}v^2 + e \right)$$



Current Model

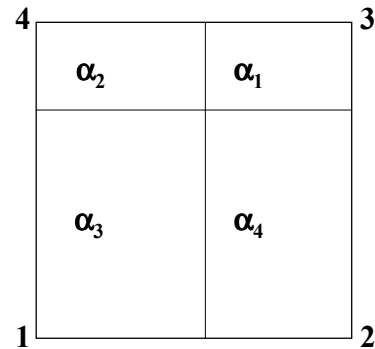
5) The Equilibrium Distribution Function

$$f_{jvk}^{(eq)} = f_{vk}^{(eq)}$$

$$f_{vk}^{(eq)} = \alpha_k d_v$$

$$d_1 = \frac{(\rho - b_0 d_0) c_2'^2 - D(\gamma - 1) \rho e}{b_1 (c_2'^2 - c_1'^2)}$$

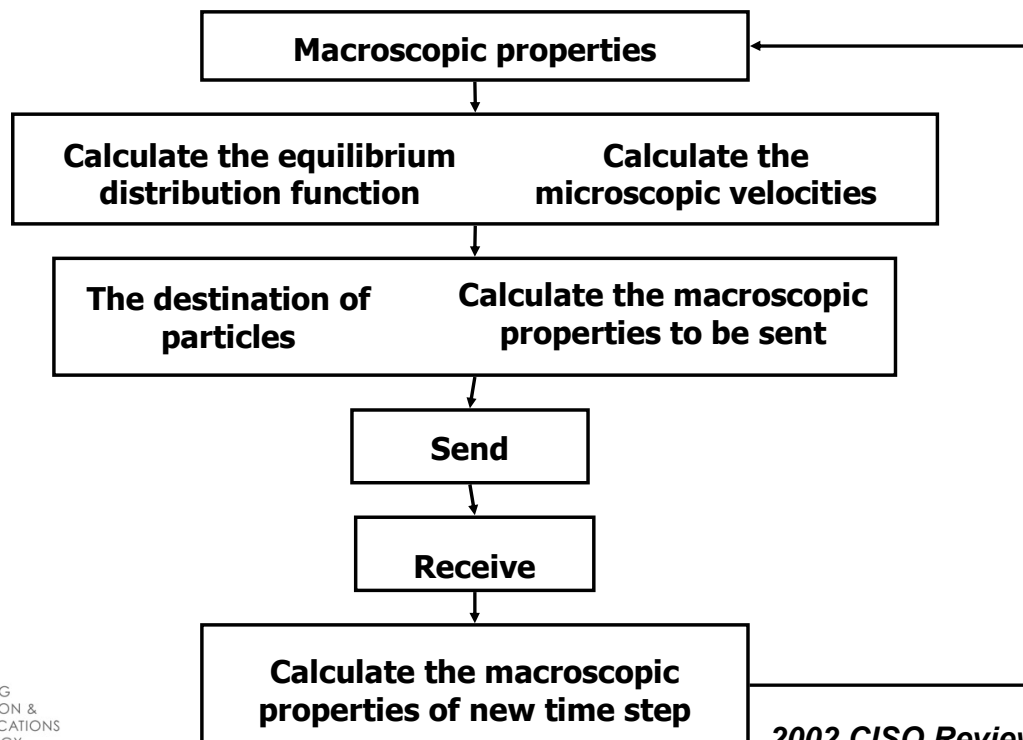
$$d_2 = \frac{D(\gamma - 1) \rho e - (\rho - b_0 d_0) c_1'^2}{b_2 (c_2'^2 - c_1'^2)}$$



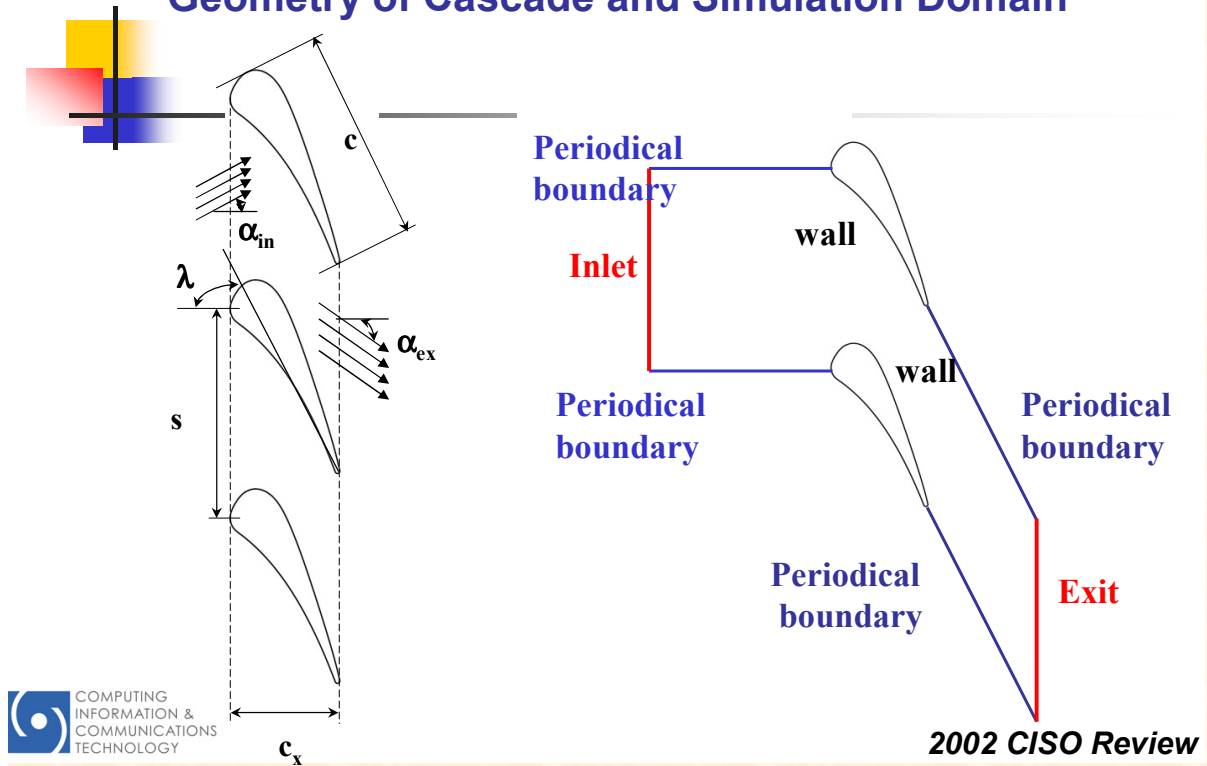
$$\alpha_1 = |u_3' v_3'| \quad \alpha_2 = |u_4' v_4'|$$

$$\alpha_3 = |u_1' v_1'| \quad \alpha_4 = |u_2' v_2'|$$

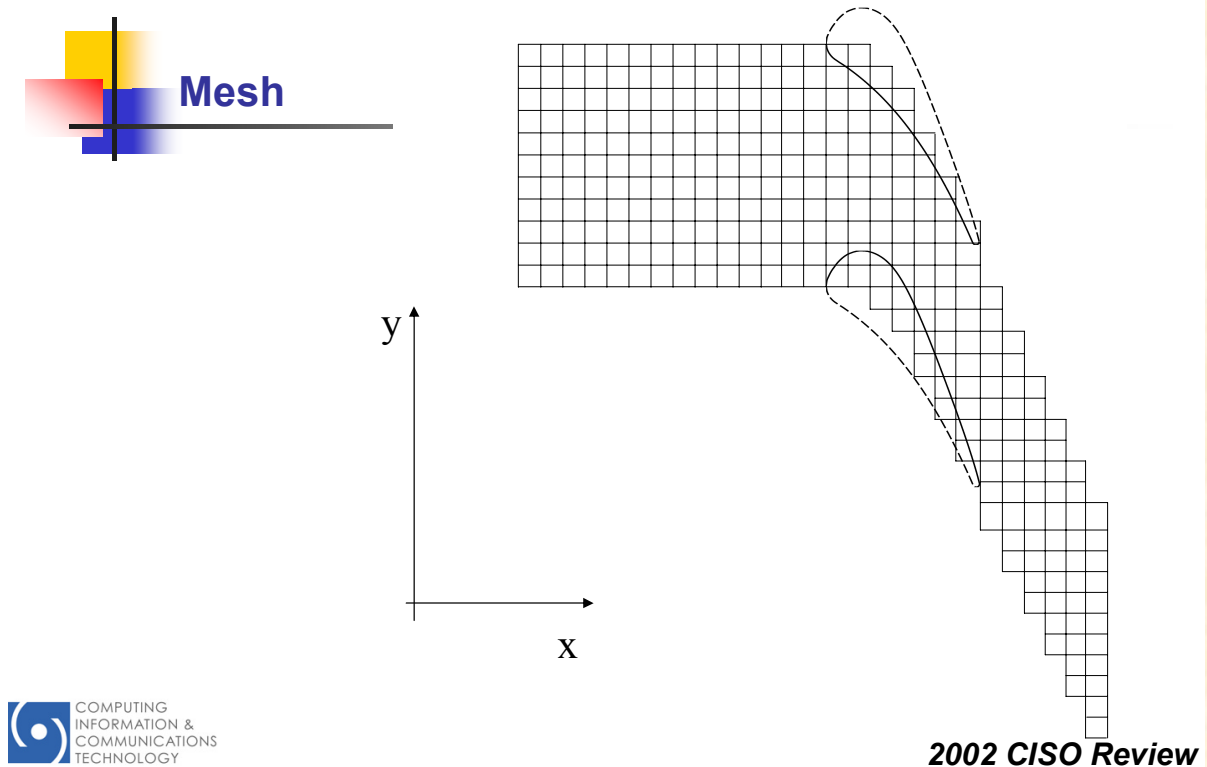
The Flow Chart Of Computation

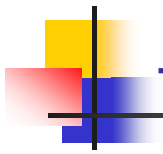


Geometry of Cascade and Simulation Domain

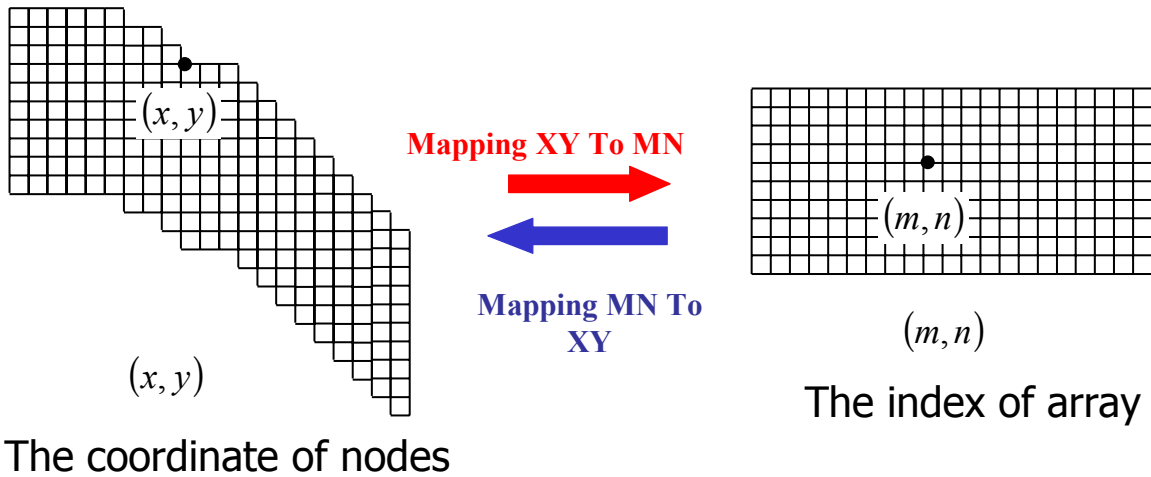


Mesh

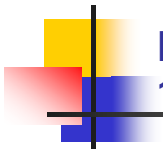




The Mapping Between Coordinates and Indexes

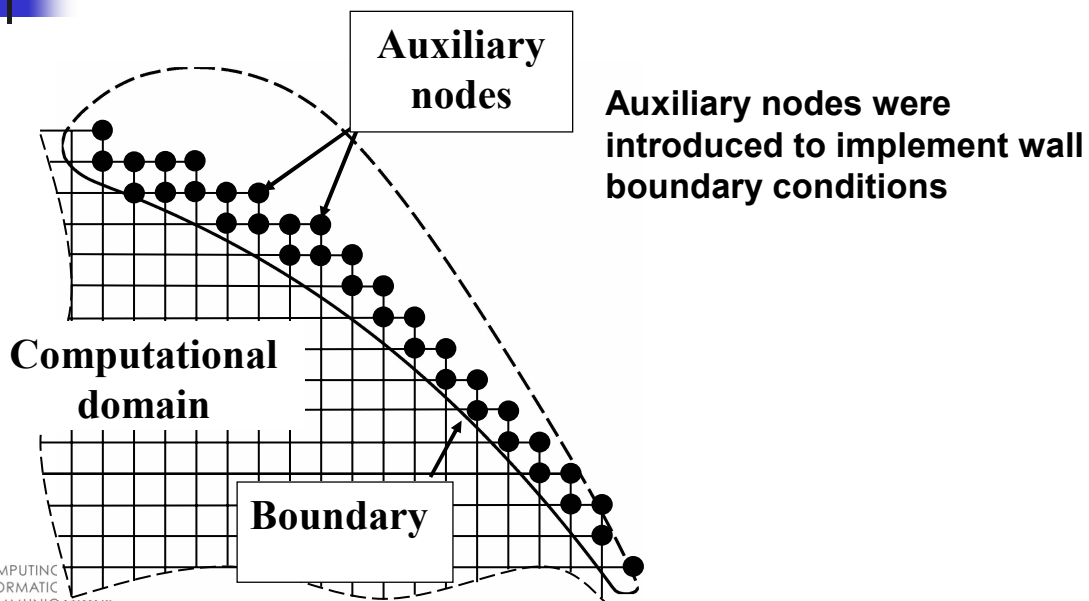


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Implement of Boundary Condition

1) Wall Boundary - Auxiliary Nodes

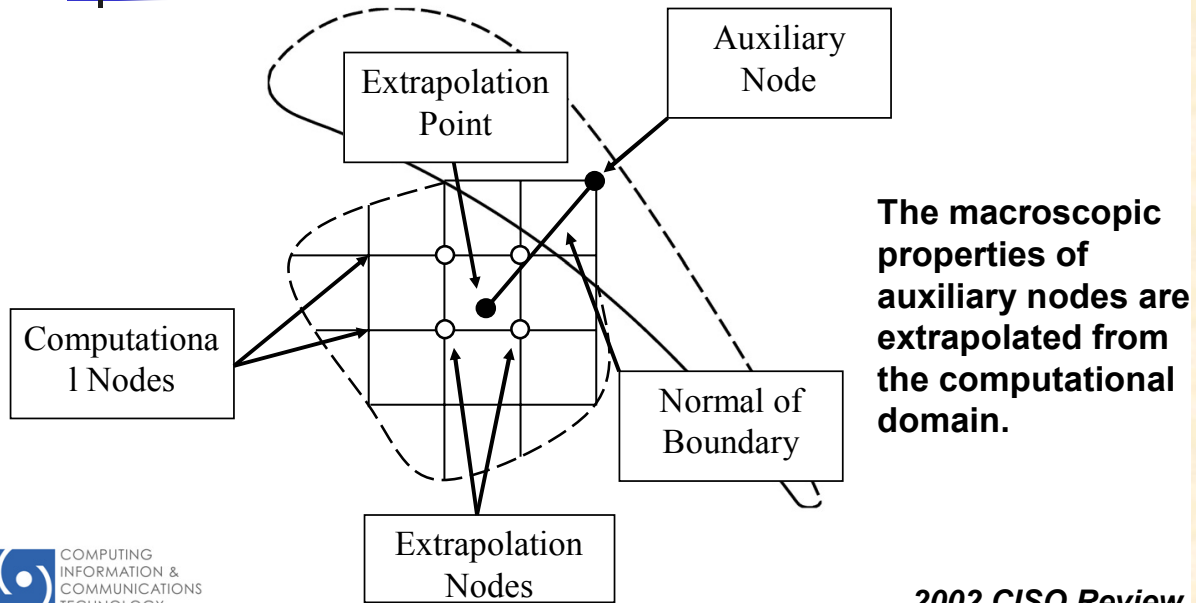


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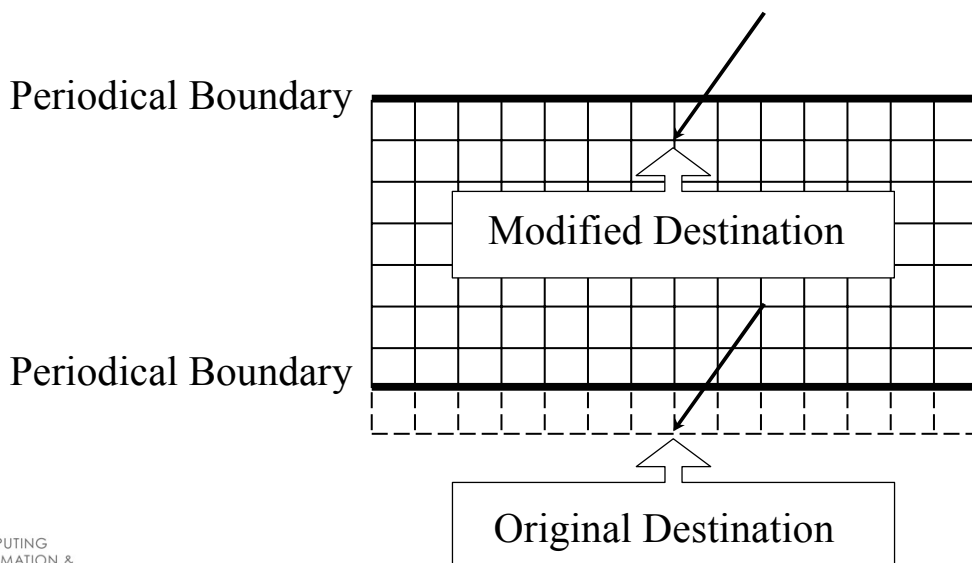
Implement of Boundary Condition

2) Wall Boundary - Extrapolation

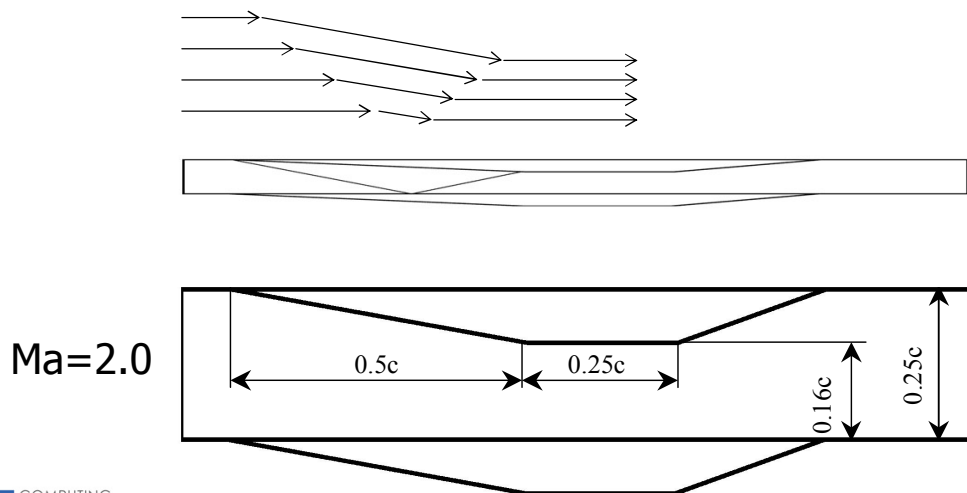


Implement of Boundary Condition

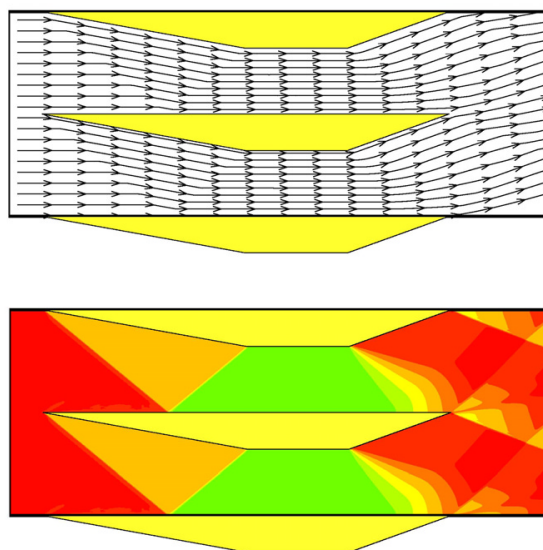
3) Periodical Boundary



Wedge Cascade Geometry



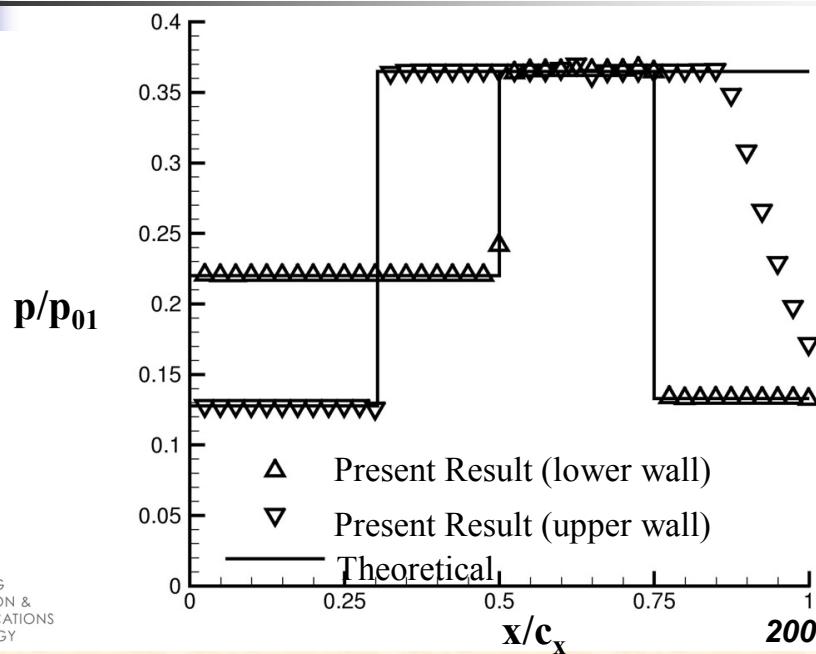
Wedge Cascade Result: Stream Line and Ma Contour





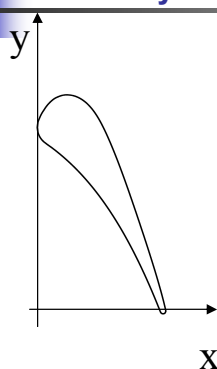
Wedge Cascade

Result: Static Pressure at Walls



C3X Cascade

Geometry and Inlet and Outlet Parameters



Stagger Angle **59.89°**

Chord **14.49cm**

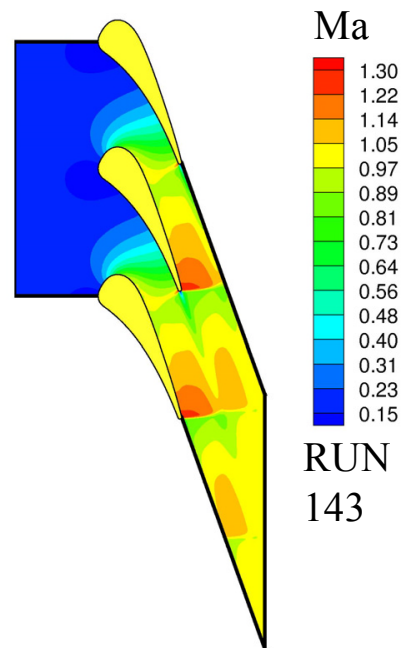
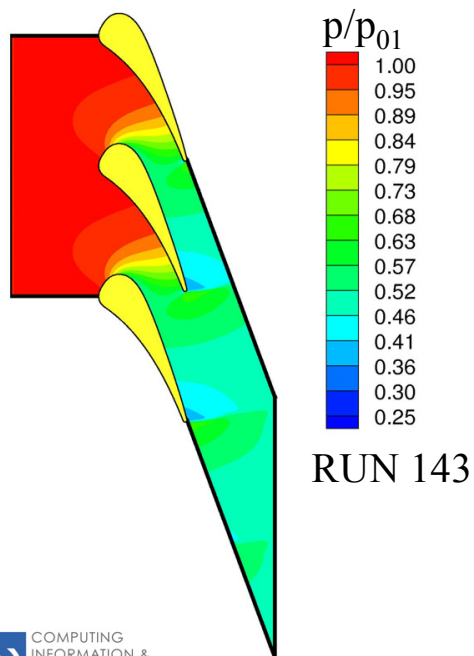
Spacing **11.77cm**

Solidity **1.23**

Axial Chord **7.82cm**

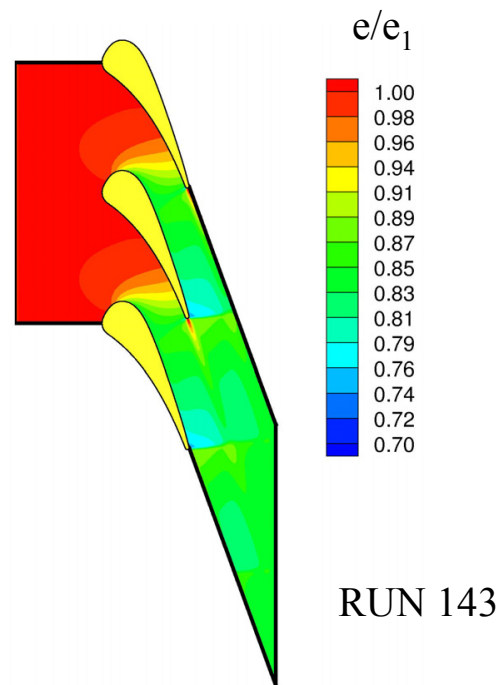
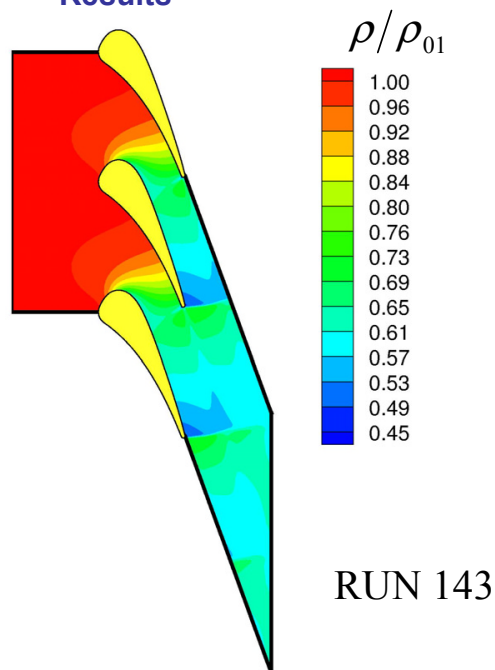
Run Number	α_{in}	p_{t1} (Pa)	T_{t1} (K)	Ma_1	Re_1	p_2/p_{t1}
RUN 143	0	7755	811	0.17	0.63×10^6	0.50
RUN 144	0	7889	815	0.16	0.63×10^6	0.59

C3X Cascade RUN 143 results



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C3X Cascade RUN 143 Results

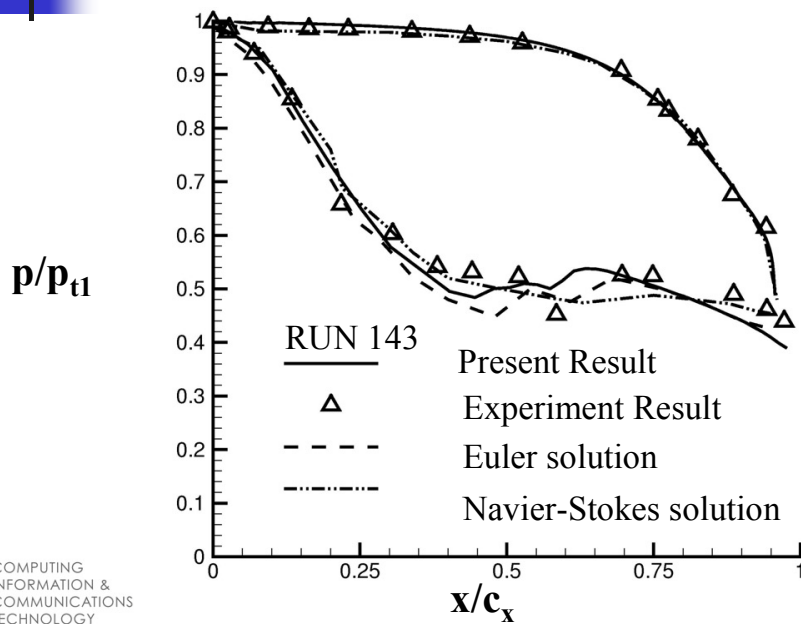


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C3X Cascade RUN 143

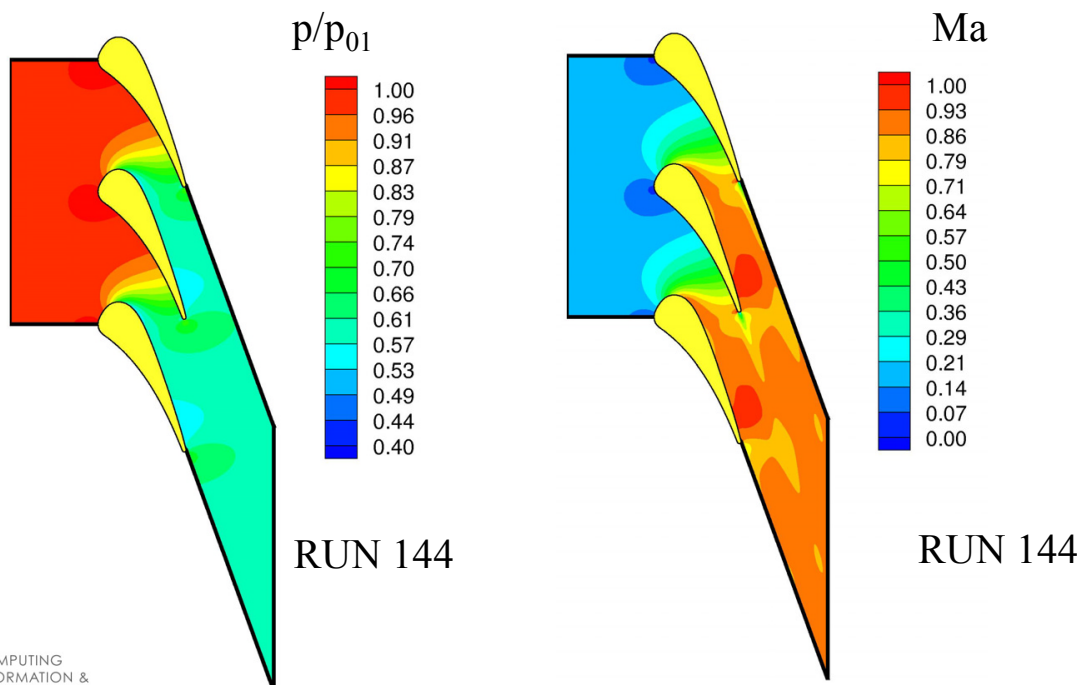
Results



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C3X Cascade RUN 144

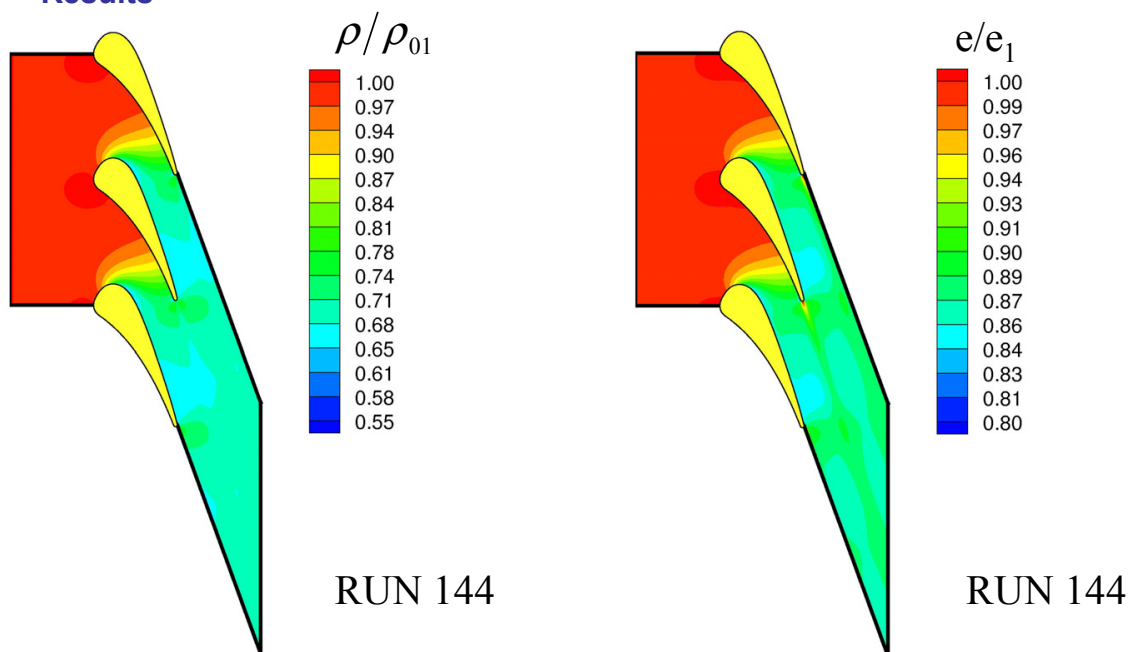
Results



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C3X Cascade RUN 144

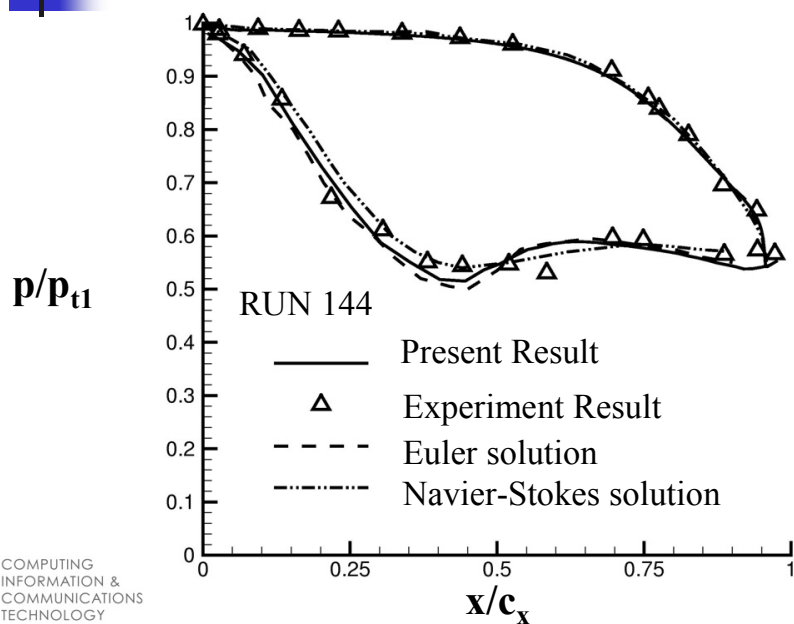
Results



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C3X Cascade RUN 144

Results

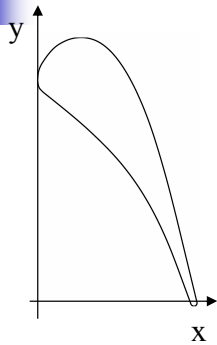


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VKI Cascade

Geometry and Inlet and Outlet Parameters



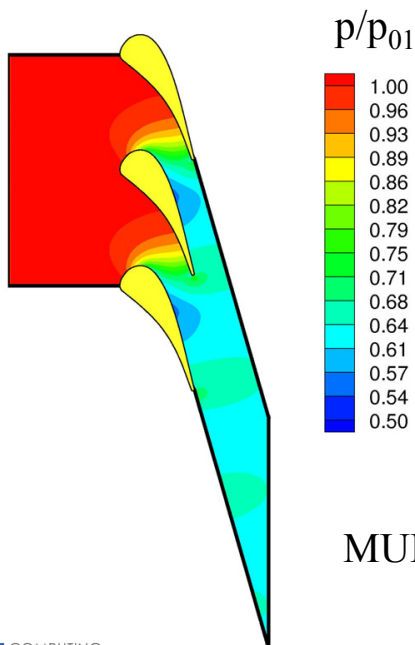
Stagger Angle	55.0°
Chord	67.646 mm
Spacing	57.50 mm
Solidity	1.1765
Axial Chord	36.98 mm

Run Number	α_{in}	p_{t1} (Pa)	T_{t1} (K)	Ma_1	Re_1	Ma_2
MUR 129	0	18200	409.20	0.15	0.27×10^6	0.84

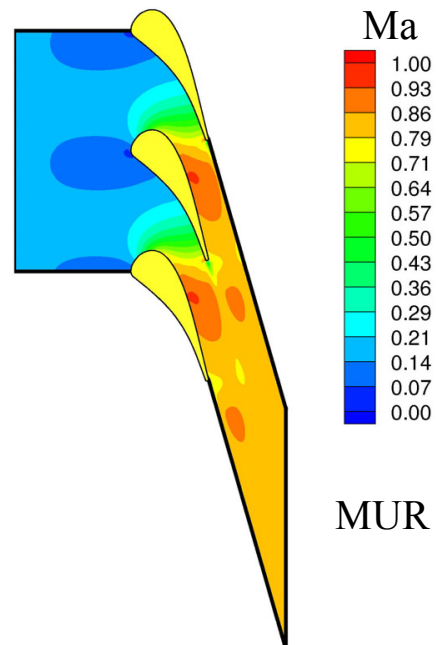


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VKI Cascade MUR 129 Results



MUR 129

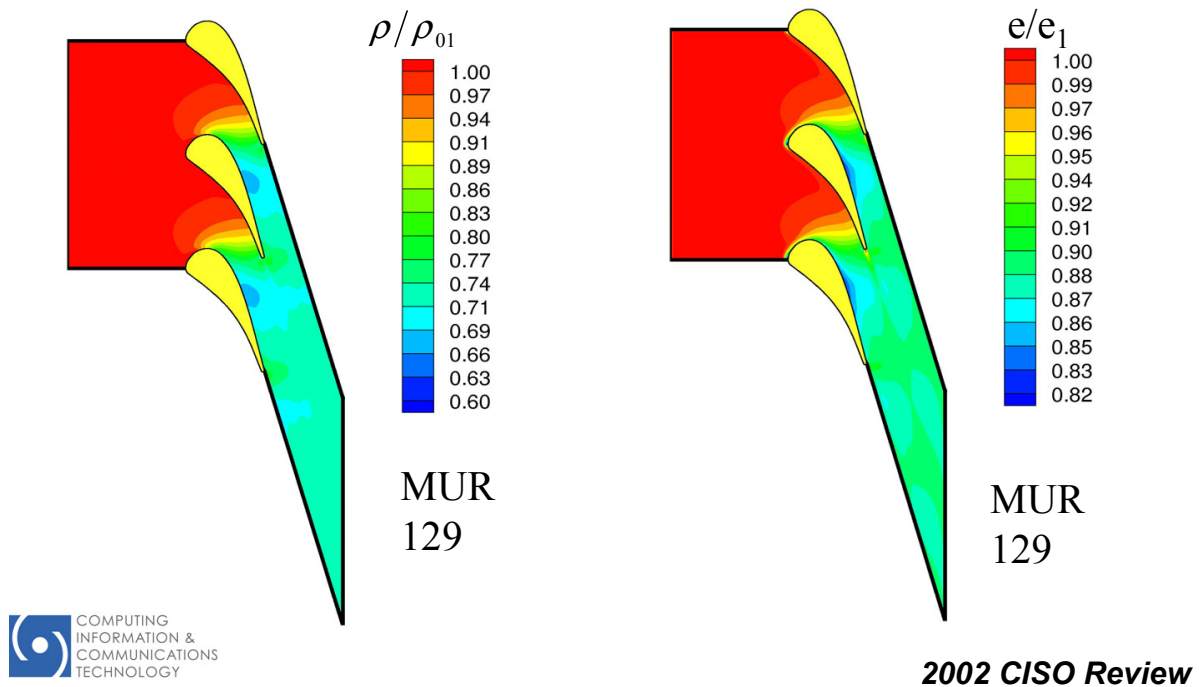


MUR 129



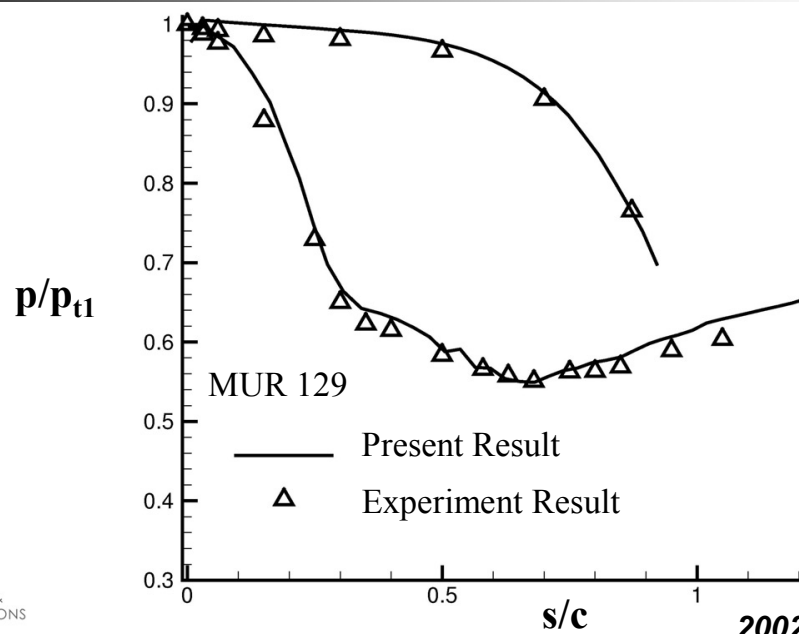
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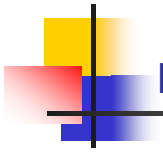
VKI Cascade MUR 129



VKI Cascade MUR 129

Results





Parallel Computing

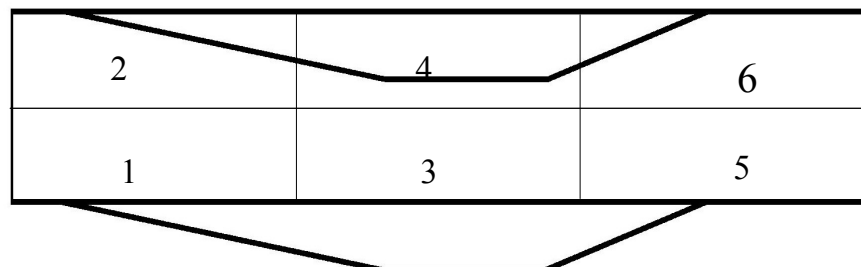
- The LB method is a natural parallel method;
- The LB model is explicit in time;
- The area of dependent domain of a node is determined by the magnitude of velocity set, which is usually a small number.



Parallel Computing

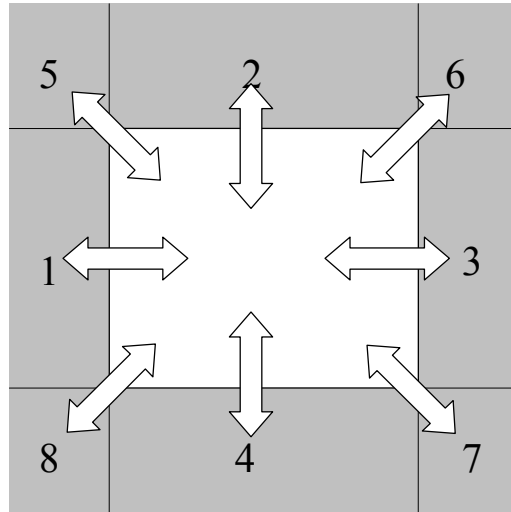
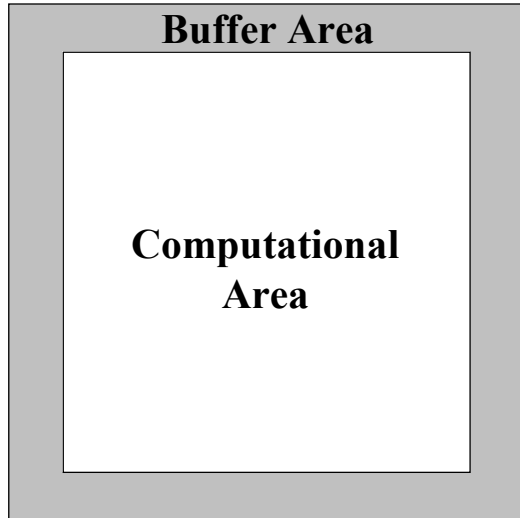
1) Division of Blocks

- The simulation of wedge cascade was parallelized to test the parallel performance of current LB model



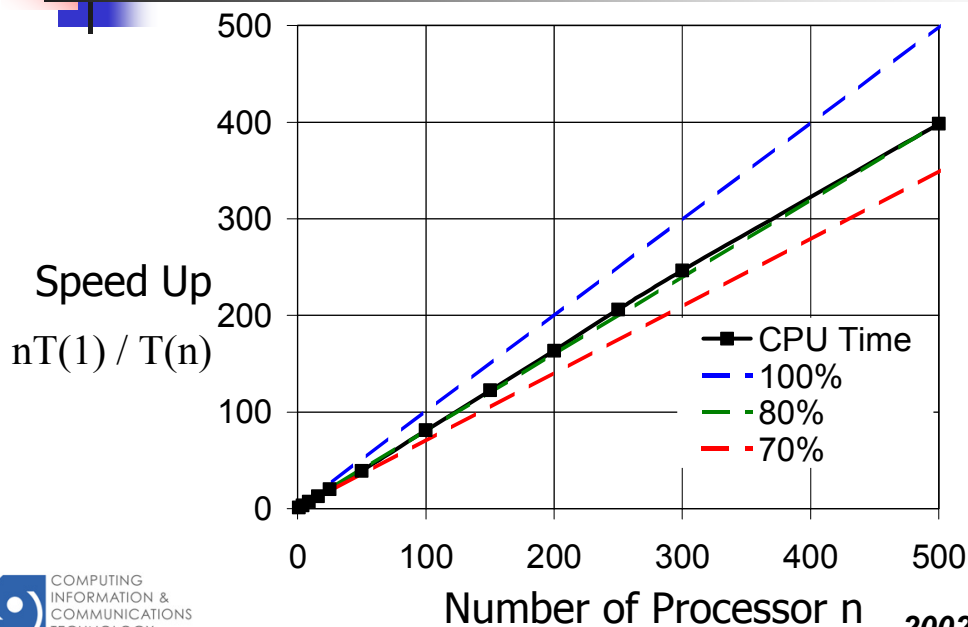
Parallel Computing

2) Information Exchange



Parallel computing

3) Result





Conclusion

- 1) A compressible LB model has been successfully developed for turbomachinery simulations.**
- 2) Successful simulation of cascades has been carried out and it is the first successful turbomachinery simulation by a LB model.**



Conclusion

- 3) A treatment of boundary condition in LB method has been introduced to the current compressible LB model.**
- 4) A new mesh treatment method has been devised in order to use regular mesh on a irregular geometry.**



Conclusion

- 5) The parallel efficiency of the new compressible LB model is studied. A linear efficiency has been demonstrated.**
- 6) The theoretical basis of the current model is analyzed in detailed.**