Electromagnetic Pumps for Liquid Metal-Fed Electric Thrusters*

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Abstract

Prototype designs of two separate pumps for use in electric propulsion systems with liquid lithium and bismuth propellants are presented. Both pumps are required to operate at elevated temperatures, and the lithium pump must additionally withstand the corrosive nature of the propellant. Compatibility of the pump materials and seals with lithium and bismuth were demonstrated through proof-of-concept experiments followed by post-experiment visual inspections. The pressure rise produced by the bismuth pump was found to be linear with input current and ranged from 0-9 kPa for corresponding input current levels of 0-30 A, showing good quantitative agreement with theoretical analysis.

Nomenclature

B, B - magnetic induction [T]
F, P - force [N]
I - pump current [A]
\( I_{sp} \) - specific impulse [s]
j - pump current density [A/m²]
l - pump channel length [m]
P - pressure [Pa]
s - pump channel height [m]
t - time [s]
w - pump channel width [m]
x - integration variable [m]

I. INTRODUCTION

Metallic propellants have been used in almost all major categories of electric propulsion (EP) devices. These have included lithium-fed arcjets¹ and magnetoplasmadynamic (MPD) thrusters,² mercury,³,⁴ cesium,⁵,⁶ and indium-fed⁷ electrostatic thrusters (both ion engines and field emission electric propulsion),
and bismuth-fed Hall thrusters.\textsuperscript{8,10}

Present interest in liquid metal-fed electric propulsion systems has been spurred by recently published data showing promising performance in systems operating on liquid-metal propellants. For example, a lithium-fed MPD thruster operating at 500 kW demonstrated over 50% efficiency with 500 hours of electrode erosion-free operation.\textsuperscript{11} Also, a bismuth-fed Hall thruster demonstrated high-power operation (140 kW) while maintaining high performance (8000s $L_p$ at greater than 70% anode efficiency) at 10 kV discharge voltage.\textsuperscript{8}

The use of a metallic propellant in a plasma thruster can result in systems and logistics-level advantages over gas-fed thrusters. For example, a gas (e.g. supercritical xenon) must be stored in high-pressure tanks to obtain a high propellant density, while a metallic propellant may be stored as either a solid or dense liquid, implying a reduction in the structural mass of the tank. In addition, vaporized metallic propellants can be condensed using simple, water-cooled plates to obtain equivalent pumping speeds of millions of liters/sec, making high-power, long-duration tests possible in present vacuum facilities.

An important component in a liquid metal-fed thruster is the mechanism that applies pressure to force propellant from a reservoir to the thruster. While there are many ways to pressurize the propellant, we chose to develop electromagnetic (EM) pumps to exploit the conductive nature of the propellant. The corrosive nature of the liquid lithium and bismuth environments severely restricted the range of materials that could be used. Both electrically insulating and conducting materials were employed in the construction of our pumps, which led to the additional challenge of joining and sealing vastly differing materials. To our knowledge this paper represents the first discussion of EM pumps where the main body was fabricated from an insulating material. Electromagnetic pump designs for both a lithium-fed and a bismuth-fed system are described, and both qualitative and quantitative data are presented to demonstrate the performance of the pumps.

In Sec. II we briefly review liquid-metal EM pumping technologies, specifically highlighting instances where EM pumps have been used in electric propulsion systems. We describe the general theoretical aspects of pump design and provide design details for both the lithium and bismuth pumps in Sec. III. The experiments performed to demonstrate pump operation and quantify performance are discussed in Sec. IV. Results obtained from our experiments are presented in Sec. V.
II. BACKGROUND

Much effort was expended on the development of EM pump technologies in the mid-1950s for use in liquid-metal cooling loops for nuclear fission reactors.\textsuperscript{12-14} This technology was widely deployed or considered for deployment in both ground-based\textsuperscript{5} and space-based\textsuperscript{16} reactor systems. Additional uses for EM pumps include the movement and control of seawater, molten metals, and various other conducting liquid metals employed in laboratory research\textsuperscript{17}.

For propellant control in electric propulsion applications, EM pumps possess several advantages over conventional mechanical pumps. They have no moving parts, thus eliminating mechanical friction losses and the need for bearings and rotating seals. This is especially advantageous if the propellant is highly reactive. Propellant pressurization is controlled by adjusting the input current and, as a result, control can be executed with a high degree of fidelity over the entire pump operating range. EM pumps are reliable, possessing no parts that are particularly failure-prone, and can be scaled to small sizes. Consequently, these pumps should prove well suited as flow controllers in electric propulsion feed systems.

While EM pumps have not been used as the primary propellant feeding mechanism in electric thrusters, they have been used as an auxiliary system in conjunction with elastic diaphragms in mercury ion engines.\textsuperscript{3,4,18} In these systems, the EM pump acted as a mechanism to provide fine-adjustment of the mercury pressure head at the propellant vaporizer. One such pump produced a pressure differential of 0.6 atm at a 20 A current level.\textsuperscript{4}

The primary designs used for electromagnetically pumping liquid metals are direct-current (DC) and alternating-current (AC) conduction pumps and linear induction pumps. In conduction pumps, current is directly passed through the fluid and a magnetic field is applied perpendicular to both the current and flow vectors. In induction pumps, the current and magnetic field in the fluid are both induced by externally flowing currents that are time varying. The external currents oscillate so as to drive an induced, travelling magnetic wave through the fluid. We chose to develop DC conduction pumps for EP applications due to their relatively simple and easily scalable design.
III. ENGINEERING DESIGN

We proceed with a brief discussion of the theoretical aspects of DC EM conduction pump design and then describe the two EM pumps constructed for this study.

A. EM Pump Theory

An idealized schematic of a DC conduction pump is presented in Fig. 1. The liquid metal flows through a channel of width \( w \), length \( l \), and height \( s \). The magnetic field \( B \) and current \( I \) are perpendicular to the flow (and each other) and are uniform over the indicated faces. When current flows between the electrodes, a Lorentz body force is exerted on the fluid in the direction of the flow vector. Integration of the body force is written as

\[
|F| = \int j \times B \, d^3x,
\]

where \( j \) is the current density (equal in magnitude to \( I/(l \, s) \) when \( |j| \) is uniform). Performing the integration yields

\[
F = \frac{I}{l \, s} B \, l \, s \, w = I \, B \, w,
\]

for uniform \( B \) and \( j \). The force causes a fluid pressure increase to develop which is equal to

\[
P = \frac{F}{s \, w} = \frac{IB}{s}.
\]

![FIG. 1: Idealized schematic of an electromagnetic pump.](image)
In Eq. (3) the pressure developed by the pump is only a function of the channel height, magnetic field strength, and total current. Material properties, specifically the conductivity of the fluid, are important only when considering the electric potential required to drive the current \( I \) through the pump. The relationship in Eq. (3) is plotted in Fig. 2 as a function of the parameters \( B/s \) and \( I \). In practice, the maximum value of \( B/s \) that can be achieved is \( \sim 10 \) T/cm. If we limit the current to below \( \sim 100 \) A to avoid excessive Ohmic heating losses, the pump pressure that can be developed in an optimized pump is roughly \( 10^9 \) Pa (1 atm).

**B. Lithium Pump Design**

In the present study, we chose to fabricate our lithium EM pump from ceramic to minimize stray conduction losses. To our knowledge, this is the first use of an insulating material for the body of an alkali metal EM pump. While the previous analysis shows that the current required to attain a given pump pressure decreases for increasing values of \( B/s \), the materials employed place certain practical limits on the maximum attainable value. The brittle nature of ceramic materials led to the choice of a relatively large value for \( s \) to minimize the potential for cracking in the walls separating the magnets from the fluid flow channel.

Aluminum nitride (AlN) was chosen as the insulating body material and while it is difficult to machine, it can survive the corrosive nature of the molten lithium environment (\( \geq 250 \) C). A major pump fabrication issue was the method by which metallic electrodes and liquid feed lines were joined to the ceramic body.
FIG. 3: A) Photograph of assembled lithium EM pump (with a U.S. quarter to show the relative scale) and B) a CAD drawing with the pump body removed to show the pump's inner detail.

The problem was made worse by the fact that the interface must be hermetically sealed to avoid liquid leaks. We chose to employ Inconel c-ring seals (compression gaskets) at all sealed joints.

The lithium EM pump was comprised of the following additional components: tungsten electrodes, 316L stainless steel liquid feed lines, samarium cobalt magnets, and an iron magnetic yoke. The yoke, electrodes, and feed line end caps were designed to also serve as a means for compressing the c-ring seals. The assembled pump is shown in Fig. 3A and a solid model with the pump body removed to show the inner detail is found in Fig. 3B. The magnet separation in this design was 6.4 mm and the measured value of $B$ was 0.45 T. The magnet employed was much wider than the diameter of the flow channel and much longer than the electrode width. Consequently, we expect the $B$ field is constant over the region where current passes through the liquid metal in agreement with the assumptions made in the derivation of Eq. (3).

C. Bismuth Pump Design

Like the lithium pump, the bismuth EM pump was also fabricated from an insulating material to minimize stray conduction losses. While molten bismuth at $\sim 300^\circ$C is not as corrosive as lithium, it was still a challenge to identify an insulating material that could be used for the pump body. We chose to use glass-mica ceramic (also known as macor) because it is nonreactive at liquid bismuth temperature levels. Relative to the AlN used in the lithium pump, it is also inexpensive and easily machinable. Samarium cobalt magnets were again employed in this pump, but relative ease of machining allowed for a magnet separation (2.9 mm) that was considerably smaller than in the lithium pump case resulting in an on-axis magnetic field strength of
FIG. 4: A) Photograph of assembled bismuth EM pump and B) a CAD drawing sectioned to reveal the pump’s inner detail.

0.64 T. Again, the magnets were significantly larger than the dimensions of the flow channel and electrodes to maintain a constant B field in the pump.

A photograph of the assembled bismuth EM pump is shown in Fig. 4A and a sectioned-view CAD model in Fig. 4B reveals the inner details of the pump’s construction. The pump was comprised of the following additional components: Inconel electrodes and c-ring compression seals, 316L stainless steel liquid feed lines, and an iron magnet yoke. The electrodes were bonded to the macor body using Ceramabond 569 high temperature alumina adhesive.

IV. TEST APPARATUS

In this section, we describe two types of experiments that were conducted to demonstrate EM pump operation, verify the design approach, and quantify pump performance.

A. Open-Loop Apparatus

Qualitative, proof-of-concept experiments were performed using an open-loop flow apparatus. These tests were carried out to demonstrate that the pumps could contain liquid metals at the required operating temperature levels and electromagnetically push the fluids from a reservoir into a catch basin. The flow loops used for the lithium and bismuth pumps are shown in Fig. 5. Since the EM pumps were situated above their respective reservoirs, it was necessary to ‘prime’ the pumps at the beginning of a test by pressurizing
the liquid metal in the reservoir using argon. Testing was conducted inside a vacuum chamber to minimize convective heat losses from the system and keep the liquid metals from excessively oxidizing.

**B. Pressure Measurement Apparatus**

A test apparatus was constructed to facilitate direct, quantitative measurement of the EM pump pressure as a function of applied current. From Eq. (3) we observe that pressure developed by the pump is independent of the conductive fluid so long as sufficient voltage is available to drive current through the pump. Consequently, we employed liquid gallium as a substitute fluid to eliminate the hazards and difficulties associated with handling and testing lithium and bismuth.

A photograph and schematic illustration of the apparatus used to measure the pressure rise across the pump are shown in Fig. 6. The bottoms of the two liquid metal reservoirs were connected through the EM pump and both reservoirs were initially pressurized using argon. Pressure in each reservoir was monitored...
using an MKS 221 differential capacitance manometer and the valves (V₁, V₂, V₃) in the system allowed for pressurization and evacuation of the reservoirs and pressure equalization between the two reservoirs.

To illustrate the operation of the test apparatus, a notional depiction of the gas pressure readings in each reservoir is presented in Fig. 7. At time $t₀$ there is no current passing through the pump and the fluid in each reservoir is at an equilibrium level, yielding gas pressures in reservoirs 1 and 2 equal to $P₁(t₀)$ and $P₂(t₀)$, respectively. At time $t₁$ a fixed, steady current is applied to the EM pump and fluid begins to flow from reservoir 1 to reservoir 2, decreasing the pressure in the former and increasing it in the latter until a new
equilibrium is reached at time \( t_2 \). Unfortunately, the change in pressure at time \( t_2 \) is a composite of the EM pump pressure \( P_{EM} \) and the gravitational pressure change (elevation head) associated with the change in the relative heights of the two columns of liquid metal in the reservoirs. Our test apparatus was designed to circumvent this difficulty, using the subsequent steps to eliminate the need to measure any liquid metal height change. Opening valve \( V_2 \) at time \( t_3 \) further increases the pressure in reservoir 2, forcing the liquid metal back through the pump and into reservoir 1. At time \( t_4 \), the liquid metal column heights in the reservoirs are equal to their levels at \( t_0 \) and the pressure in reservoir 1 satisfies the condition \( P_1(t_4) = P_1(t_0) \). When the fluid column heights return to their initial levels, the excess pressure in reservoir 2 is solely balanced by the pressure applied by the EM pump:

\[
P_{EM} = P_2(t_4) - P_2(t_0).
\]
V. EXPERIMENTAL RESULTS

A. Lithium Open-Loop Tests

Approximately 100 g of lithium were loaded into the reservoir of the test apparatus shown in Fig. 5A. The reservoir was heated to ~300 °C while the EM pump was ~250 °C. Argon gas at a pressure of 15 Torr was introduced into the reservoir to prime the pump. After the electrodes in the pump were electrically bridged by liquid lithium, the pump current was set to 20 A. Lithium flow (at a rate of ~1 g/sec) commenced and liquid droplets were expelled into the catch basin. The EM pump was operated with molten lithium for approximately 100 sec during the test run, after which time the lithium supply in the reservoir was exhausted. The pump was exposed to liquid and then hot solid lithium for several hours while it cooled to room temperature.

Aside from demonstrating electromagnetic pumping with lithium, we also sought to validate the short-term compatibility of our pump materials and sealing techniques with the molten lithium environment (i.e. demonstrate that no immediate problems arose due to our choice of pump materials and seals). After the test run, the pump was disassembled and visually inspected. No evidence of material degradation was observed (see Fig. 8). The AIN body showed no visible signs of cracking or short-term failure and no lithium was evident outside of the compression seals. The metallic components of the pump also appeared to be undamaged. The long term corrosive effects of liquid lithium on the pump materials and seals were not evaluated, and would require extended life testing that was beyond the scope of the present study.

B. Bismuth Open-Loop Tests

The same test procedure was implemented for open-loop bismuth tests and similar, positive results were obtained. The main difference was that the bismuth reservoir had to be brought to a higher temperature (~325 °C) before flow could be initiated. Several tests (~10) were performed, with the pump remaining at elevated temperature for several hours during each test. No degradation of the macor or seals was evident through visual inspections after the tests and no leaks were found.
The pressure rise produced by bismuth EM pump design was evaluated over an applied current range from 10-30 A using the test apparatus shown in Fig. 6, and with liquid gallium substituted for liquid bismuth as discussed in Sect. IV.B. A time history of the raw pressure measurements in the two reservoirs for an applied EM pump current of 15 A is presented in Fig. 9. The results of our testing indicated that the liquid metal had a tendency to "stick" at low flow speeds, possibly due to oxidation layers forming on the gallium surface or reservoir walls impeding wetting and flow. A result of this impediment was that reservoir 1 was unable to return its initial (pressure) state in a monotonic manner. It was necessary to alternate between overpressurizing and underpressurizing reservoir 2 to maintain liquid flow, which caused the pressure in reservoir 1 to oscillate about the initial pressure with progressively decreasing amplitude. The smallest difference between the high and low values of the pressure in reservoir 2 were used to establish the uncertainty in the measurement of $P_{EM}$.

The measured values of $P_{EM}$ as a function of pump current for several different trials are presented in Fig. 10. The data show that the pressure developed by the EM pump is linear with current, as indicated by the linear curve fit (dashed line) of the function

$$P_{EM} = a I,$$  \hspace{1cm} (5)
FIG. 9: Raw reservoir pressure data as a function of time obtained using the setup shown in Fig. 6 with the EM pump current set to 15 A at $t=40$ s using liquid gallium as a working fluid.

FIG. 10: Measured EM pump pressure versus current (the solid lines are theoretical curves computed using Eq. (3) and based on the values of $s$ given while the dashed line represents a linear curve fit of the data).

where $a$ is the fit coefficient. For the bismuth pump, this value was $336\pm 16$ Pa/A. As noted, this coefficient is independent of the type of liquid metal propellant used. The systematic uncertainties associated with the measurement instruments were small compared to the uncertainty inherent in the measurement procedure, and is therefore not included in the uncertainty analysis.

The pump's theoretical performance can be computed using Eq. (3), but there is some ambiguity in defining the channel height $s$. The pump was constructed by drilling an axial, 2.4 mm diameter hole through
The actual value of \( s \) is dependent on the (unknown) manner in which current distributes itself within the cylindrical channel. It is reasonable to assume that the effective value of \( s \) must be somewhere between 1.6 and 2.4 mm in our pump. In Fig. 10 we have plotted theoretical EM pump pressures computed using Eq. (3) for each of these extreme cases. The measured data show good quantitative agreement, lying roughly midway between the two theoretical curves.

The data demonstrate that the bismuth EM pump is capable of delivering a pressure of \( O(10^4) \) Pa at 30 A. Furthermore, the pump is seen to operate at a level consistent with the theoretical maximum performance, implying that losses associated with stray conduction currents are negligible in the present design. Finally, the linearity of the pump pressure data indicate that the effective value of \( s \) in the pump is relatively constant (linear scaling of current density in the channel) as input current is increased up to at least 30 A.

VI. CONCLUSIONS

The following major results and conclusions are reported in this study. Our EM pump designs were based around central, insulating ceramic bodies, which was a departure from all earlier liquid metal EM pump work. Material compatibility with liquid lithium led us to design an EM pump around an aluminum nitride body. The pump was fabricated and tested in the relevant environment, and subsequent disassembly of the pump revealed no short-term degradation of the AlN body, stainless steel feed lines, or tungsten electrodes, indicating that the pump would not immediately fail when brought into contact with liquid lithium. The metallic compression seals used to join the conducting and insulating components showed no short-term signs of leakage. An EM pump for use with hot liquid bismuth was fabricated from macor. The relative ease of machinability of macor compared to AlN allowed for a pump design with a much smaller magnet separation than that found in the lithium pump. This resulted in a greater magnetic field in the flow channel. The pump was successfully operated and showed no signs of degradation or leakage after exposure to liquid bismuth over several thermal cycles and \( O(10 \text{ hrs}) \) of operation. Quantitative measurements of EM pump pressure in the bismuth pump using liquid gallium as a substitute fluid show that the output was linear with current, with a pump pressure of \( \sim 10 \text{ kPa} \) produced at 30 A. The experimentally measured and theoretically
computed pressure values exhibit good quantitative agreement, indicating minimal stray conduction current losses and linear scaling of current density in the channel.

Acknowledgments

We gratefully acknowledge the contributions and support of our the MSFC technical support staff: Boris Stanojev, Doug Davenport, Tommy Reid, Doug Galloway, Keith Chavers, and Rondal Boutwell. We also extend our gratitude to Dr. Noah Rhys and students Amado DeHoyos, Chris Dodson, and Jeff Gross for their contributions to this work. This work was supported by NASA’s Exploration Systems Mission Directorate (Project Prometheus) and funded under contract NNM05AA25C managed by Randy Baggett and NAS7-03001 managed by John Warren.

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