

Gamma Ray Bursts and their links with Supernovae and Cosmology

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Abstract

Gamma-ray bursts are the most luminous explosions in the Universe, whose origin and mechanism is the focus of intense interest. They appear connected to supernova remnants from massive stars or the merger of their remnants, and their brightness makes them temporarily detectable out to the largest distances yet explored in the Universe. After pioneering breakthroughs from space and ground experiments, their study is entering a new phase with observations from the recently launched *Fermi* satellite, as well as the prospect of detections or limits from large neutrino and gravitational wave detectors. The interplay between such observations and theoretical models of gamma-ray bursts is reviewed, as well as their connections to supernovae and cosmology.

1 Introduction

Roughly once a day, somewhere within our Hubble horizon a Gamma-ray burst (GRB) occurs, which for the next few seconds or tens of seconds completely overwhelms the gamma-ray flux from the rest of the Universe, including the Sun. In fact, the GRB prompt electromagnetic energy output during tens of seconds is comparable to that of the Sun over $\sim \text{few} \times 10^{10}$ years, or to that of our entire Milky Way over a few years; and their X-ray and optical afterglow over the first day after the outburst can outshine the brightest quasars, as well as supernovae, making them potentially important probes of the distant Universe. Since the discovery of their X-ray afterglows by the *Beppo-SAX* satellite in 1997 and the subsequent detection of their optical counterparts, we have measured these objects out to the farthest cosmological distances. Thanks to triggers and measurements from the *Swift* [1] and *Fermi* [2] [3] satellites, we have now detailed multi-wavelength data for many hundreds of bursts, and redshifts for over 200 of them, and this data set will continue to grow with the continuation of *Swift* and *Fermi*, and the possible upcoming *SVOM* mission [4].

GRBs are thought to arise either when a massive star ($\gtrsim 25M_{\odot}$) undergoes core collapse, or possibly when a double neutron star or a neutron star and a black hole binary merges [5]. The first scenario applies to so-called *long* GRBs (LGRBs), whose γ -ray light-curve lasts for $t_{\gamma} \gtrsim 2$ s, while the second scenario is the likeliest one so far for *short* GRBs, whose γ -ray light-curves last $t_{\gamma} \lesssim 2$ s [6] (for the latter, this short duration refers to photons at $\varepsilon \gtrsim 100$ keV; some “short” bursts, at softer energies,

have tails lasting as much as 100 s [7, 8]). In either scenario, it appears inevitable that a compact core object of a few to several solar masses forms, whose radius is of order of the Schwarzschild radius for this mass, $r_g \sim 10^6 (M/3M_\odot)$ cm, over a timescale comparable to a few dynamic (free-fall) times, which is likely to be a black hole. Accretion of residual infalling gas leads, if the core is fast rotating (guaranteed for a binary), to an accretion disk whose inner radius is $r_0 \sim 3r_g \sim 10^7$ cm, and the typical variability timescale of accretion is $t_0 \sim (2GM/r_0^3)^{-1/2} \sim 10^{-3}$ s. The bulk of the gravitational energy, of order a solar rest mass or 10^{54} erg is, as in SNe, rapidly radiated as thermal neutrinos ($E_{\nu,th} \sim 10$ MeV), and some amount is radiated as gravitational waves. A smaller fraction, of order $E_j \sim 10^{51} - 10^{52}$ erg is converted into a fireball of equivalent blackbody temperature $T_0 \sim \text{few MeV}$. This energy eventually emerges as the burst, in the form of a jet. However, for purposes of the dynamics one can generally use (see below) the isotropic equivalent energy $E_0 = E_j(4\pi/\Omega_j) \sim 10^{53} E_{53}$ erg, which with a nominal total burst duration $t_b \sim 10$ s implies a nominal isotropic equivalent luminosity $L_\gamma \sim E_0/t_b = 10^{52} L_{52}$ erg s $^{-1}$ (if most of that energy is emitted as γ -rays). The number density of photons at r_0 is roughly given by $L_\gamma = 4\pi r_0^2 c n_\gamma \epsilon$ where $\epsilon \sim kT \sim \text{MeV}$, and the “compactness parameter (roughly the optical depth of a photon with energy $\gtrsim m_e c^2$ against $\gamma\gamma \rightarrow e^+e^-$ pair production) is

$$\ell' \sim \tau_{\gamma\gamma} \sim n_\gamma \sigma_T r_0 \sim \frac{\alpha \sigma_T L_\gamma}{4\pi r_0 c \epsilon} \sim 10^{15}, \quad (1)$$

where σ_T is the Thomson cross section and α is the fraction of the luminosity above $m_e c^2$. This creates a fireball of gamma-rays, electron-positron pairs and hot baryons, where most of the entropy and pressure is in the photons and leptons. The optical depth is huge, and the radiation pressure far exceeds gravity, so the fireball expands and becomes relativistic. A simple lower limit on the expansion bulk Lorentz factor follows from the observations of photons up to $\gtrsim \text{GeV}$ energies, in some bursts. Such photons, trying to escape the source, would collide against softer photons and pair produce, $\gamma + \gamma \rightarrow e^+ + e^-$, degrading the spectrum to $\lesssim 0.5$ MeV. However the pair production threshold is angle dependent, and pair production is avoided if

$$\epsilon_1 \epsilon_2 (1 - \cos \theta) \leq 2m_e^2 c^4. \quad (2)$$

In a relativistically moving jet, causality implies that only photons within angles $\theta \lesssim 1/\Gamma$ can interact, so with $\epsilon_1 \sim 30$ GeV, $\epsilon_2 \sim \text{MeV}$, and $\cos \theta \sim 1 - \theta^2/2$, we see that this implies

$$\Gamma \gtrsim \sqrt{(\epsilon_1/m_e c^2)(\epsilon_2/m_e c^2)}/2 \gtrsim 10^2. \quad (3)$$

A more general constraint on Γ is obtained by considering the typical photon spectral distribution in GRBs, which is a broken power law “Band” function $n(\epsilon) \propto \epsilon^{-\beta} \text{ph cm}^{-3} \text{MeV}^{-1}$, where $\beta \simeq 1$ or $\beta \simeq 2$ for ϵ below or above a break frequency $\epsilon_{br} \sim \text{MeV}$ [9, 10]. The $\gamma\gamma$ optical depth at each energy $\geq m_e c^2$ depends on the optical depth to target photons at $\leq m_e c^2$ satisfying the threshold condition (2), in the jet comoving frame. This optical

depth can be shown to be $\propto \Gamma^{-6}$, so for increasingly high Γ the source becomes optically thin to increasingly higher energy photons. The result is that typical GRBs, even if the highest energy photons observed are only 100 MeV, require bulk Lorentz factors in excess of $\Gamma \sim 250$ [11].

If the entire burst energy is released impulsively, injecting an energy E_0 in a timescale t_0 inside a radius r_0 , with the numbers comparable to those above, the initial entropy per baryon is $\eta \sim E_0/M_0 c^2$, where M_0 is the baryon load of the fireball. If the pressure is mainly due to radiation and pairs, the inertia is due to baryons (“baryonic dynamics” regime) the bulk Lorentz factor initially accelerates as $\Gamma(r) \sim (r/r_0)$, e.g. [12]. After the baryons have become non-relativistic in their own frame, the expansion changes to a coasting behavior at a saturation radius $r_{sat} \sim r_0 \eta$, and the fireball continues to expand freely with $\Gamma \simeq \eta \simeq \text{constant}$. The observationally estimated values are $\Gamma \sim \eta \sim 10^2 - 10^3$, so the baryon load is typically $10^{-5} - 10^{-6} M_\odot$. The behavior is similar if the energy and mass input is spread out over accretion times (i.e. outflow feeding or ejection times) of $t_b \sim 10 - 100$ s, as inferred for “long” GRB. On the other hand, if the fireball pressure, or rather stress tensor, is dominated by magnetic fields, the dynamic behavior is different; depending on the symmetries of the fields, the acceleration behavior can range from $\Gamma \propto r^{1/3}$, e.g. [13, 14, 15] to $\Gamma \propto r^\eta$ where $1/3 \leq \eta \lesssim 2/3$, at least when the outflow is one-dimensional [16, 17]. This regime is referred to as magnetically dominated, or Poynting dominated dynamics.

In practice, the outflow is inferred to be jet-like, rather than isotropic, with an average solid angle $\langle \Omega_j \rangle / 4\pi \sim 1/500$ or $\langle \theta_j \rangle \sim 1/30$ [18, 7]. In the case of core collapse (“long”) GRBs, this can be due to the outer parts of the star providing a massive barrier, which is best pierced along the centrifugally lightened rotation axis, along which the fireball escapes. The stellar envelope provides a sideways pressure which channels the jet. However, as long as the jet opening angle θ_j exceeds $1/\Gamma \sim 10^{-2}$ (i.e. $\gtrsim 0.5^\circ$), which is generally the case, the expansion occurs as if it were isotropic: causality prevents the gas to have any knowledge of what happens outside an angle $1/\Gamma$. For compact binary mergers, the data on jet opening angles is much sparser, but the average value may not be too different [19].

2 Observations

The *Swift* mission, launched in November 2004, finds bursts and observes the prompt phase with the Burst Alert Telescope (BAT). The afterglow is then observed with the X-Ray Telescope (XRT) and the UV Optical Telescope (UVOT). Measurements of the redshift and studies of host galaxies are typically done with large ground-based telescopes which receive immediate alerts from the spacecraft when GRBs are detected. *Swift* has, by far, the largest number of well-localized bursts, afterglow observations and redshift determinations. As of 1 April 2012, BAT has detected 669 GRBs (annual average rate of ~ 90 per year). Approximately 80% of the BAT-detected GRBs have rapid repointings (the remaining 20% have spacecraft constraints that prevent rapid slewing).

Of those, virtually all long bursts observed promptly have detected X-ray afterglow. Short bursts are more likely to have negligible X-ray afterglow, fading rapidly below the XRT sensitivity limit. The fraction of rapid-pointing GRBs that have UVOT detection is $\sim 35\%$. Combined with ground-based optical observations, about $\sim 60\%$ of *Swift* GRB have optical afterglow detection. There are so far (mid 2012) about 200 *Swift* GRBs with redshifts, compared with 41 in the pre-*Swift* era. The redshift distribution of *Swift* GRBs is shown in Fig. 2.

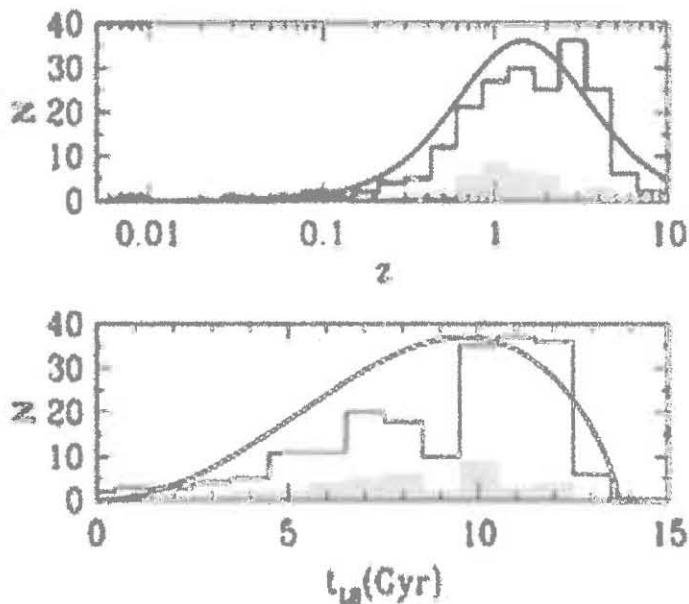


Figure 1: Redshift distribution and cosmic look-back time of GRBs. The *Swift* GRBs are in blue, the pre-*Swift* GRBs in yellow and the co-moving volume of the universe is the red curve. The GRBs roughly follow the co-moving volume [7].

The *Fermi* mission, launched in June 2008, has two instruments, the Gamma-ray Burst Monitor (GBM) and the Large Area Telescope (LAT). The GBM has scintillation detectors and covers the energy range from 8 keV to 40 MeV. It measures spectra of GRBs and determines their position to $\sim 5^\circ$ accuracy. The LAT is a pair conversion telescope covering the energy range from 20 MeV to 300 GeV. It measures spectra of sources and positions them to an accuracy of $< 1^\circ$. The GBM detects GRBs at a rate of ~ 250 per year, of which on average 20% are short bursts. The LAT detects bursts at a rate of ~ 8 per year.

3 The photon spectrum: standard picture

The expansion converts internal energy into bulk kinetic energy, so that the gas cools in its own rest frame and soon becomes an inefficient radiator. In the absence of dissipation,

at the photospheric radius r_{ph} where the flow becomes optically thin to scattering (which for baryon dominated dynamics usually occurs above the saturation radius r_{sat}) the escaping radiation would carry only a small fraction of the burst kinetic energy, and might be expected to have a quasi-black-body spectrum [20, 21]. This motivated the fireball shock model, where the bulk kinetic energy is reconverted by shocks into random particle energy, and thence into non-thermal radiation, at radii beyond the scattering photosphere, where the flow is optically thin; this could be at either external shocks [22, 23], where the jet interacts with external (e.g. interstellar) matter, or at internal shocks [24] occurring within the jet, at radii intermediate between the photosphere and the external shock radius. These radii can be expressed as

$$\begin{aligned} r_{ph} &\simeq (L\sigma_T/4\pi m_p c^3 \eta^3) \sim 4 \times 10^{12} L_{\gamma,52} \eta_{2.5}^{-3} \text{ cm}, \\ r_{dis} &\simeq \Gamma^2 c t_v \sim 3 \times 10^{13} \eta_{2.5}^2 t_{v,-2} \text{ cm}, \\ r_{dec} &\simeq (3E_0/4\pi n_{ext} m_p c^2 \eta^2)^{1/3} \sim 2 \times 10^{17} E_{53}^{1/2} n_0^{-1/2} \eta_2^{2/3} \text{ cm}. \end{aligned} \quad (4)$$

Here the photospheric radius assumes baryonic dynamics [24] (for magnetic dynamics see [15]). The dissipation (internal shock) radius r_{dis} follows from the relativistic relation between observer time t , the radius r and Lorentz factor, $r \simeq \Gamma^2 c t$. Considering two shells of matter ejected at time intervals comparable to the variability timescale of ejection t_v , with Lorentz factors differing by $\Delta\Gamma \sim \Gamma$, and the deceleration radius r_{dec} where external shocks start follows from the energy conservation assumption, $E_0 \simeq (4\pi/3) r_{dec}^3 n_{ext} m_p c^2 \Gamma^2$, when the swept up matter has been shock-heated to an energy comparable to the explosion energy (with $\Gamma \sim \eta$, e.g. [12] for details).

In the (collisionless) shocks, the particles are reheated to thermal energies comparable to the pre-shock relative kinetic energies per particle. The internal shocks are semi-relativistic (since $\Delta\Gamma \sim \Gamma$, the relative Lorentz factor $\Gamma_{rel} = (1/2)[(\Gamma_1/\Gamma_2) + (\Gamma_2/\Gamma_1)]$ is semi-relativistic), and this results in a shock luminosity $L_{sh} = \epsilon_{sh} L_0$ where L_0 is initial kinetic luminosity $L_0 = E_0 \eta / t_b \dot{M} c^2$ and the shock dissipation efficiency $\epsilon_{sh} \lesssim 0.1$. Particles repeatedly bounce across the shock, scattered by magnetic irregularities whose energy density can be assumed to build up to some degree of equipartition with the proton thermal energy, so the comoving magnetic field is $B'^2/8\pi \sim \epsilon_B 4\Gamma_{rel}^2 m_p c^2$ (primed quantities are in comoving frame), where $\epsilon_B \leq 1$. The repeated crossing Fermi accelerates the particles to a relativistic power law [22, 24]; the minimum electron (comoving) random Lorentz factor is $\gamma_{e,m} \sim \epsilon_e (m_p/m_e) \Gamma_{rel} \gg 1$ (for internal shocks $\Gamma_{rel} \sim 1$, while for external $\Gamma_{rel} \sim \eta$) and one expects a power law $N(\gamma_e) \propto \gamma_e^{-p}$ above that, with $p \sim 2 - 2.3$. One also expects something similar for the protons in the flow.

i) Internal shock prompt radiation: The electrons in the internal shock will emit synchrotron and inverse Compton (IC) radiation, leading to non-thermal broken power law photon spectrum, roughly similar to the observed “Band” spectra [23, 24]. For reasonable values of $L, \eta, t_v, \epsilon_B, \epsilon_e$ the synchrotron peak energy (observer frame) corresponding to the minimum $\gamma_{e,m}$ are comparable to the observed Band spectral break

energies,

$$\varepsilon_{sy,m} \sim \varepsilon_{br} \sim 1 \epsilon_B^{1/2} \epsilon_e^{3/2} \frac{L_{\gamma,52}^{1/2}}{\eta_{2.5}^2 t_{v,-2}} \text{ MeV.} \quad (5)$$

Eq. (5) assumes the randomized kinetic luminosity of internal shocks L_{sh} to be related to their γ -ray luminosity through $L_\gamma = \epsilon_e L_s$, which is true in the fast cooling regime where cooling time is shorter than the dynamic time $t'_{sync} \ll t'_{dyn}$ [25], which in internal shocks is true [26]. In this fast cooling regime, for a typical Fermi electron index $p \simeq 2$ the photon spectral index above ε_{br} is expected to be $n(\varepsilon) \propto \varepsilon^{-(p/2)-1} \propto \varepsilon^{-2}$ ph cm⁻³ s⁻¹, as typical for the high energy branch of the canonical Band spectrum. The synchrotron model predicts below $\varepsilon_{sy,m}$ a low energy branch $n(\varepsilon) \propto \varepsilon^{-2/3}$, which through superposition of maxima for various parameters could fit the observed Band average low energy branch $n(\varepsilon) \propto \varepsilon^{-1}$ [23] (but flatter spectra are a problem see below).

ii) External shock radiation and afterglow: At r_{dec} the relativistic ejecta has used up about half its initial energy in sweeping up an amount $M_{sw} \sim M_{ejecta}/\eta$ of external material, driving a forward shock into the external gas and a reverse shock into the ejecta. This occurs at an observer time

$$t_{dec,0} \sim r_{dec}/(2c\eta^2) \simeq 10 (E_{53}/n_0)^{1/3} \eta_{2.5}^{-8/3} \text{ s.} \quad (6)$$

The forward shock is initially highly relativistic, $\Gamma_{sh} \sim \eta$, so $\Gamma_{rel,fs} \sim \eta$ and the synchrotron spectrum is in the hard X-rays or gamma-rays. The reverse shock builds up slowly and for usual conditions becomes semi-relativistic, $\Gamma_{rel,rs} \sim 1$ at the deceleration time, when it has crossed the ejecta. For this reason its $\gamma_{e,M}$ is smaller than that of the forward shock electrons, and the reverse shock spectrum peaks in the optical or UV [27, 28]. Beyond r_{dec} the expansion continues but it is increasingly slowed down due to increasing amount of swept up matter. In the adiabatic approximation the bulk Lorentz factor changes from being $\sim \eta \sim \text{constant}$ to a power law decline behavior given by $E_0 \propto r^3 \Gamma^2$, or

$$\Gamma \sim \eta (r/r_{dec})^{-3/2}. \quad (7)$$

In both the forward and the reverse shock one expects again Fermi acceleration of electrons to a power law distribution leading to synchrotron and inverse Compton radiation, but the synchrotron break energy becomes softer in time as the Doppler boost decreases in accordance with eq.(7). This leads to an afterglow [28] progressing from X-rays through optical to radio lasting from minutes to days to months, with fluxes decaying as power laws in time. This prediction was indeed confirmed by observations with the Beppo-SAX satellite of the X-ray [29] and with ground telescopes of the optical [30] afterglow of GRB 970228, soon followed by the first confirmation of a cosmological redshift [31] and a radio detection [32] for GRB 970508. The amount of data on, and understanding of, afterglows has since increased enormously, see e.g. [7, 8].

The external reverse shock gas, most luminous at $t \sim t_{dec}$, is in pressure equilibrium with forward shock gas, and having a higher particle density and smaller energy per electron than the forward shock, its synchrotron spectrum peaks in the O/UV. This was

predicted to lead to an observable prompt optical emission [27, 28], later detected with robotic ground telescopes such as ROTSE [33] triggered by spacecraft, the number of such detections being now several dozen [7].

The external forward shock is expected to also give rise to an IC component, in particular a synchrotron-self-Compton (SSC) from upscattering its own synchrotron photons [27, 34], which would appear in the GeV range. Such GeV emission was detected already by EGRET [35], and more recently by the Fermi LAT, e.g. [36]. This is discussed in §5.

4 Prompt MeV emission: issues and developments

Issues arise with the radiation efficiency of internal shocks, which is small in the bolometric sense (5-10%), unless the different shells have widely differing Lorentz factors [37, 38, 39]. The MeV efficiency is also substantially affected by IC losses [40, 41], in the BATSE range being typically $\sim 1-5\%$, both when the MeV break is due to synchrotron [42, 37, 43] and when it is due to inverse Compton [44].

The synchrotron interpretation of the GRB radiation is the most attractive; however, a number of effects can modify the simple synchrotron spectrum. One is that the cooling could be rapid, i.e. when the comoving synchrotron cooling time $t'_{sy} = 9m_e^3 c^5 / 4e^4 B'^2 \gamma_e) \sim 7 \times 10^8 / B'^2 \gamma_e$ s is less than the comoving dynamic time $t'_{dyn} \sim r/2c\Gamma$, the electrons cool down to $\gamma_c = 6\pi m_e c / \sigma_T B'^2 t'_{dyn}$ and the spectrum above $\nu_c \sim \Gamma(3/8\pi)(eB'/m_e c)\gamma_c^2$ is $F_\nu \propto \nu^{-1/2}$ [25, 45].

The radiative efficiency issue has motivated investigating various alternatives, e.g. relativistic turbulence in the emission region [46, 47]. This assumes that relativistic eddies with Lorentz factors $\gamma_r \sim 10$ exist in the comoving frame of the bulk $\Gamma \gtrsim 300$ flow, and survive to undergo at least γ_r changes over a dynamic time, leading both to high variability and better efficiency. Various constraints may however pose difficulties [48], while numerical simulations [49] indicate that relativistic turbulence would lead to shocks and thermalization, reducing it to non-relativistic.

The synchrotron spectral interpretation faces a problem from the observed distribution of Band low energy spectral indices β_l (where $N_e \propto \varepsilon^{\beta_l}$ below the spectral peak), which has a mean value $\beta_l \sim -1$, but for a fraction of bursts this slope reaches values $\beta_l > -2/3$ which are incompatible with a low energy synchrotron asymptote $\beta_l = -2/3$ [50]. Possible explanations include synchrotron self-absorption in the X-ray [51] or in the optical range up-scattered to X-rays [44], low-pitch angle scattering or jitter radiation [52, 53], or time-dependent acceleration [54], where low-pitch angle diffusion might also explain high energy indices steeper than predicted by isotropic scattering.

Pair formation can become important [24, 40, 41] in internal shocks or dissipation regions occurring at small radii, since a high comoving luminosity implies a large comoving compactness parameter $\ell' \gg 1$. Pair-breakdown may cause a continuous rather than an abrupt heating and lead to a self-regulating moderate optical thickness pair plasma at sub-relativistic temperature, suggesting a comptonized spectrum [45]. Copious pair

formation in internal shocks may in fact extend the photosphere beyond the baryonic photosphere value (4). Generic photosphere plus internal shock models [55, 56, 57] which includes the emission of a thermal photosphere as well as a non-thermal component from internal shocks outside of it, subject to pair breakdown, which can produce both steep low energy spectra, preferred breaks and a power law at high energies. A moderate to high scattering depth can lead to a Compton equilibrium which gives spectral peaks in the right energy range [58]. Pair enrichment of the outflow (due to back-scatter $\gamma\gamma$ interactions) can in general affect both the radiative efficiency and the spectrum [59, 60, 61, 62, 63].

4.1 Photospheric models

In the synchrotron interpretation the observed peak frequency is dependent on the bulk Lorentz factor, which may be random, and since the observed peaks appear to concentrate near 0.2-1 MeV [7], the question can be posed whether this is indeed due to synchrotron, or to some other effect. An alternative is to attribute a preferred peak to a black-body at the comoving pair recombination temperature in the fireball photosphere [64]. In this case a steep low energy spectral slope is due to the Rayleigh-Jeans part of the photosphere, and the high energy power law spectra and GeV emission require a separate explanation [55]. A related explanation has been invoked [65], considering scattering of photospheric photons off MHD turbulence in the coasting portion of the outflow, which up-scatters the adiabatically cooled photons up to the observed break energy and forms a power law above.

For a photosphere occurring at $r < r_{sat}$, which in a baryon-dominated model requires high values of η , the radiative luminosity in the observer frame is undiminished, since $E'_{rad} \propto r^{-1}$ but $\Gamma \propto r$ so $E_{rad} \sim \text{constant}$, or $L_{ph} \propto r^2 \Gamma^2 T'^4 \propto \text{constant}$, since $T' \propto r^{-1}$. However for the more moderate values of η the photosphere occurs at $r > r_{sat}$, and whereas the kinetic energy of the baryons is constant $E_{kin} \sim E_0 \sim \text{constant}$ the radiation energy drops as $E_{rad} \propto (r/r_{sat})^{-2/3}$, or $L_{ph} \sim L_0 (r_{ph}/r_{sat})^{-2/3}$ [23, 55]. This weakening of the photospheric luminosity leads again to a lowered efficiency, as well as a lower peak energy than observed. However, if the photosphere is dissipative (due to shocks or other dissipation occurring at or below the photosphere) then a high efficiency is regained, and the thermal peak photon energies are in the range of observed Band peaks [66]. An important aspect is that Compton equilibrium of internal shock electrons or pairs with photospheric photons lead to a high radiative efficiency, as well as to spectra with a break at the right preferred energy and steep low energy slopes [66, 67, 68]. It also leads to possible physical explanations for the Amati [69] or Ghirlanda [70] relations between spectral peak energy and burst fluence [66, 71].

4.2 Magnetic models

An alternative set of models for the prompt emission assume that this is due to magnetic reconnection or dissipation processes, or else to the external shock. Magnetic models fall into two categories, one where baryons are absent or dynamically negligible, at least initially [72, 13, 73, 74], and another where the baryon load is significant, although dynamically sub-dominant relative to the magnetic stresses [65, 75, 63]. These scenarios would in all cases still lead to an external shock, whose radius would be again given by r_{dec} in equ. (4), with a standard forward blast wave, but possibly a weaker or absent reverse shock [75, 28], due to the very high Alfvén (sound) speed in the ejecta. For the same reason, internal shocks may be prevented from forming in magnetized outflows. However, this depends on the magnetization parameter σ ; if not too large, reverse shocks [76, 77, 78] or internal shocks might still form [79], although with different strengths and radiation characteristics. In fact, “internal” dissipation regions may form due to magnetic reconnection, at radii comparable but differing from r_{dis} of eq.(4), where electric fields due to reconnection (instead of a Fermi mechanism) leads to particle acceleration, and a high radiative efficiency is conceivable.

A hybrid dissipation model, entitled ICMART [80] involves a hybrid magnetically dominated outflow leading to semi-relativistic turbulent reconnection. Here a moderately magnetized $\sigma = (B'^2/4\pi\rho'c^2) \lesssim 100$ MHD outflow undergoes internal shocks as $\sigma \rightarrow 1$, leading to turbulence and reconnection which accelerates electrons at radii $r \gtrsim 10^{15}$ cm. These involve fewer protons than usual baryonic models, hence less conspicuous photospheres, and have significant variability, and the efficiency and spectrum are argued to have advantages over those in the usual synchrotron internal shock models.

The baryon-free Poynting jet models resemble pulsar wind models, except for being jet-shaped, as in AGN baryon-poor models. The energy requirements of GRB (isotropic-equivalent luminosities $L_\gamma \gtrsim 10^{52}$ erg s⁻¹) require magnetic fields at the base in excess of $B \sim 10^{15}$ G, which can be produced by shear and instabilities in an accreting torus around the black hole (BH). The energy source can be either the accretion energy, or via the magnetic coupling between the disk and BH, extraction of angular momentum from the latter occurring via the Blandford-Znajek mechanism [81]. The stresses in this type of model are initially magnetic, involving also pairs and photons, and just as in purely hydro baryon-loaded models they lead to an initial Lorentz factor growth $\Gamma \propto r$ up to a pair annihilation photosphere [82]. This provides a first radiation component, typically peaking in the hard X-ray to MeV, with upscattering adding a high energy power law. Internal shocks are not expected beyond this photosphere, but an external shock provides another IC component, which reaches into the GeV-TeV range.

The baryon-loaded magnetically dominated jets have a different acceleration dynamics than the baryon-poor magnetic jets or the baryon dominated hydrodynamic jets: whereas both the latter accelerate initially as $\Gamma \propto r$ and eventually achieve a coasting Lorentz factor $\Gamma_f \sim L_\gamma/\dot{M}c^2$, the baryon-loaded magnetically dominated jets have a variety of possible acceleration behaviors, generally less steep than the above. In the simplest treatment of a homogeneous jet with transverse magnetic field which under-

goes reconnection, the acceleration is $\Gamma \propto r^{1/3}$ [73, 15], while in inhomogeneous jets where the magnetic field and the rest mass varies across the jet the average acceleration ranges from $\Gamma \propto r^{1/3}$ to various other power laws intermediate between this and $\Gamma \propto r$ [83, 14]. Few calculations have been made [84] of the expected (leptonic) spectral signatures in the simpler magnetized outflow photospheres, typically in a one-zone steady state approximation, showing that a Band-type spectrum can be reproduced.

5 GeV-TeV phenomenology and models

The first *Fermi* GRB observations, starting in late 2008, soon yielded a number of surprises. One of the first bright objects showing radically new features was GRB 080916C [36], in which the GeV emission started only with a second pulse, which was delayed by ~ 4 s relative to the first pulse, which visible only in MeV (Fig. 5).

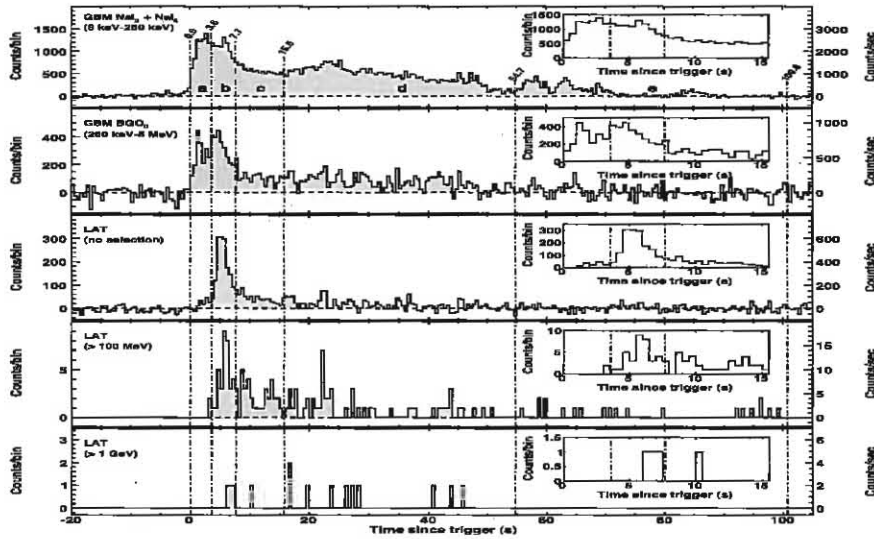


Figure 2: Light curves of GRB080916C showing the GBM (top two curves) and the LAT (bottom three curves) energy ranges [36].

The spectra of GRB 080916C consisted of simple Band-type broken power laws, the first pulse having a soft high energy index disappearing at GeV, but the second and subsequent pulses having harder high energy indices reaching well into the GeV range. There was no evidence for a second spectral component (such as expected from inverse Compton or hadronic effects). The peak energy of the Band function evolved from soft to hard and back to soft, but in this as well as in other *Fermi* LAT bursts, the GeV emission persisted in afterglows typically lasting $\gtrsim 1000$ s. On the other hand, in a few bursts, such as GRB090902B [85], a second spectral component did indeed appear, at 5σ significance, and also a lower energy power law extension whose significance is lower but suggestive. Another burst with a high energy second component was GRB 090926A [86], this one showing a clear cut-off or turnover to the high energy power law (Fig. 5).

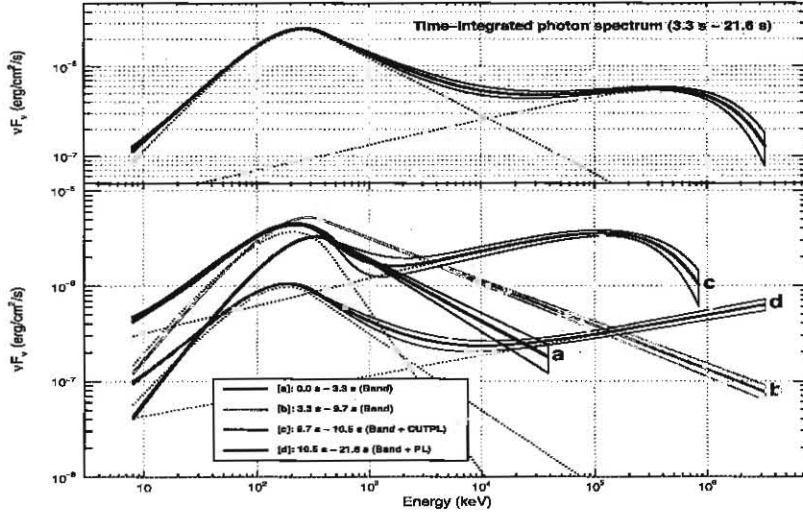


Figure 3: Spectra of GRB090926A from *Fermi* at four different time intervals, a= [0.0-3.3s], b= [3.3-9.7s], c= [9.7-10.5s], d= [10.5-21.6s] [86].

A significant advance from *Fermi* LAT was the discovery of the first GeV *short* burst, GRB 090510 [87], whose general behavior (including a GeV delay) was qualitatively similar to that of long bursts. Several more short bursts have been discovered since with the *Fermi* LAT.

The *Fermi* -LAT extended emission, if one ignores various details, has a relatively simple interpretation in terms of conventional forward shock leptonic synchrotron models (i.e. relying on accelerated electrons or e^+e^- pairs) [88, 89]. Such models provide a natural delay between an assumed prompt MeV emission (assumed implicitly to come from, e.g. internal shocks or other “inner” mechanisms) and the GeV emission from the external shock, which starts after a few seconds time delay. However, taking into account more carefully the constraints provided by the Swift MeV and X-ray observations, and considering carefully the accompanying inverse Compton (IC) scattering and Klein-Nishina effects, it is clear that at least during the prompt emission, there must be a subtle interplay between the shorter lasting mechanism providing the MeV radiation and the mechanism or emission region responsible for the bulk of the longer lasting GeV radiation [90, 91, 92]. One general shortcoming of these early studies was a postponement of addressing the interaction of the GeV emission with a specific, self-consistent model of the prompt emission, including the radiative inefficiency in an implicit internal shock assumption.

A resolution of this problem is possible if the prompt MeV Band spectrum is due to an efficient dissipative photosphere (baryonic, in this case) with an internal shock upscattering the MeV photons at a lower efficiency, giving the delayed GeV spectrum [93]. Alternatively, for a magnetically dominated outflow, where internal shocks may not occur, an efficient dissipative photospheric Band spectrum can be up-scattered by the external shock and produce the observed delayed GeV spectrum [94]. Depending on the parameters, the combined spectrum can look like a two-component or a single

Band spectrum (Fig. 4). On the other hand, a delayed GeV spectrum can also be

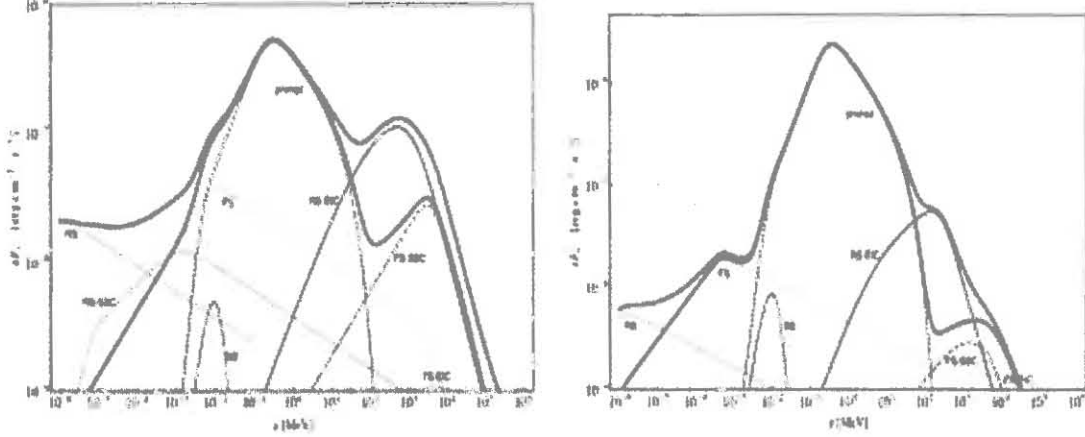


Figure 4: A magnetically dominated leptonic model where the MeV Band spectrum is due to photospheric emission, there are no internal shocks, and the external reverse and forward shock upscatter the MeV photospheric spectrum into the GeV range. Parameters are typical for *Fermi* LAT GRBs, but in some cases lead (left) to a two-component spectrum, while in others (right panel) it can be fitted as a single Band spectrum extending to the GeV range [94].

expected in hadronic models, which assume the co-acceleration, along with the electrons, of protons which undergo electromagnetic cascades and synchrotron losses along with their secondaries [95, 96, 97, 98] see §6.

The *LAT* data show that a fraction of GRB are emitting (in their own rest frame) photons in the energy range of at least up to 30 – 90 GeV. A partial list of *Fermi* LAT detections [99] of maximum observed photon energies and redshifts ($E_{\gamma,obs}, z$) is (13.2, 4.35), (7.5, 3.57), (5.3, 0.74), (31.3, 0.90), (33.4, 1.82), (19.6, 2.10), (2.8, 0.897), (4.3, 1.37). This list shows that (i) even $z > 4$ bursts can produce $E_{\gamma} > 10$ GeV photons at the observer, and (ii) some $z \sim 1$ bursts can produce $E_{\gamma} > 30$ GeV photons at the observer. This is highly encouraging for the planned large Cherenkov Telescope Array (CTA), as described in recent reviews [100, 101]. The CTA detection rate is estimated [100] to be 0.7 – 1.6 per year, based on the rate of Swift triggers (while GBM triggers on *Fermi* are more frequent, their positional accuracy is poorer). This rate is affected by uncertainties in the fraction of bursts which emit in the GeV range, relative to those emitting below 100 MeV [102, 103]. E.g., as of February 2011, in 2.5 years, *Fermi* LAT detected 4 bursts at energies > 10 GeV (or 20 at > 0.1 GeV) out of some 700 bursts detected by *Fermi* GBM at $E < 100$ MeV. This very small fraction of the total ($\lesssim 1\%$) of course is in part due to the size constraints under which space detectors must operate.

In the standard internal shock model of prompt emission, the intra-source $\gamma\gamma$ absorption typically prevents photons in excess of a few GeV to emerge [40, 41], unless the bulk Lorentz factor is above ~ 700 [104]. For photospheric models of the prompt emission,

e.g. [105], photons in excess of 10 GeV can escape the source from radii $r_{\gamma\gamma} \sim 10^{15}$ cm, and such radii are also inferred phenomenologically from one-zone analyses of the Fermi data on GRB. However, most of the GeV emission occurs during the afterglow, which is good for ground-based TeV Cherenkov telescopes, whose reaction time can be slower. Indeed, the GeV emission can last up to ~ 1000 s, far more than the $\sim 2 - 50$ s of the MeV emission. In the standard external shock scenario, the compactness parameter is smaller than in the internal shock, and inverse Compton scattering is expected to lead to multi-GeV and TeV photons [75, 34], the details depending on the electron distribution slope and the radiative regime (e.g. slow or fast cooling). This scenario is thought to be responsible for the afterglows of GRB [28], and is also thought to be responsible for the extended GeV emission observed by LAT so far [88, 89, 106, 107], etc. Of course, propagation in the intergalactic medium from high redshifts leads to additional $\gamma\gamma \rightarrow e^\pm$ interaction with the extragalactic background light, or EBL [108, 109, 110], the threshold for which depends on the photon energy and the source redshift.

Thus, if TeV emission is produced, it is mainly expected to be detectable from $z \lesssim 0.5$, while the 10-30 GeV emission should be (and is) detectable from higher redshifts. Thus, the GeV detectability is dictated by the source physics, the source rate and the immediate source environment. The source rate, based on MeV observations, is well constrained [7], while the near-source environment effects can be reasonably parametrized (e.g. [111]). The source physics, however, has large uncertainties. E.g. in an external shock model the simple synchrotron self-Compton (SSC) model can be additionally complicated by the scattering of photons arising at other locations well inside the external shock, e.g. from the photosphere [93], or from an inner region energized by continued central engine activity [112]. Similar uncertainties about the soft photon source and location would affect hadronic cascade models. The observed GBM high energy spectral slopes are in many cases steep enough not to expect much GeV emission from their extrapolation [99], while in other cases the LAT spectrum shows a cutoff or turnover, e.g. in GRB 090926B [86]. Nonetheless, all things considered, the estimate [100] of 0.7 – 1.6 CTA detections per year appears to be a conservative lower limit.

6 Hadronic models

If GRB jets are baryonic, or magnetically dominated but with non-negligible baryon load, the charged baryons should be co-accelerated with the electrons in any shocks or reconnection zones, and hadronic processes would lead to both secondary high energy photons and neutrinos. Monte Carlo codes have been developed to model hadronic effects in relativistic flows, including p, γ cascades, Bethe-Heitler interactions, etc. E.g. [113, 96] used such a code to calculate the photon spectra from secondary leptons resulting from hadronic interactions following proton acceleration in the same shocks that accelerate primary electrons in GRBs. The code uses an escape probability formulation to compute the emerging spectra in a steady state, and provides a detailed quantification of the signatures of hadronic interactions, which can be compared to those arising from

purely leptonic acceleration. Spectral fits of the *Fermi* LAT observations of the short GRB 090510 were modeled by [96] as electron synchrotron for the MeV component and photohadronic cascade radiation for the GeV distinct power law component.

Since acceleration as well as cascade development can take some time, in principle even one-zone models might result in GeV-MeV photon delays. E.g. [95] assumes for GRB 090510 the prompt MeV to be electron synchrotron and the GeV to be proton synchrotron, whose cooling time cranks down the photon energy into the GeV range on a few second delay timescale, the electron plus proton synchrotron merging into a single Band function with the approximate spectral slope of the GeV photons. A more recent one-zone hadronic calculation [98] shows that even when proton synchrotron is not important, hadronic cascade development leads to a second GeV component, with time delays comparable to the observed ones (Fig. 5, right panel). Similar delays can, however, be also obtained in purely leptonic two-zone photosphere plus external shock models [114] (Fig. 5, left panel).

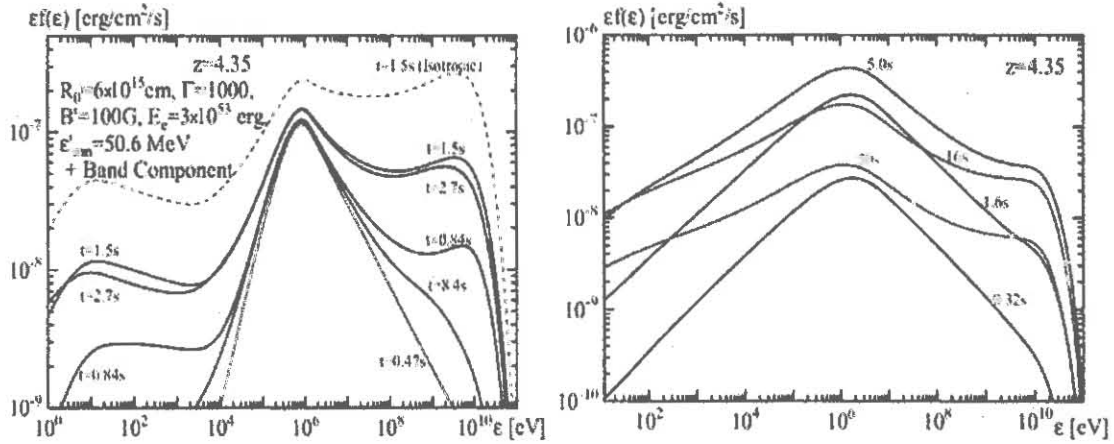


Figure 5: Temporal evolution of the observable spectral photon flux for typical *Fermi* LAT parameters, from Monte Carlo simulations. Left: a purely leptonic two-zone model with a photospheric (MeV) Band component and upscattering into the GeV range by a shock further out [114]. Right: a one-zone hadronic model, where electron synchrotron produces the Band MeV spectrum and hadronic cascade secondaries produce the GeV spectrum, as well as a low energy component [98].

Hadronic interactions can also have implications for a low energy photon power law below the Band function, perhaps resulting in a GRB optical prompt flash, as discussed by [115]. For the usual Band MeV spectrum produced by conventional leptonic mechanisms, the acceleration of hadrons leads to secondaries whose radiation produces both a high energy “extra” GeV component and a prompt bright optical emission from secondary synchrotron. This might explain, e.g. the observed “naked eye” 5th magnitude flash of GRB 080319B, e.g. [116].

Hadronic binary collisions in baryon-loaded jets can also be important, both for efficient energy dissipation and for shaping the photon spectrum. This is because the baryons will consist of both protons (p) and neutrons (n), especially if heavy elements are photo-dissociated. The protons are coupled to the radiation during the acceleration phase but the neutrons are carried along only thanks to nuclear (p, n) elastic collisions, whose characteristic timescale at some point becomes longer than the expansion time. At this point the p and n radial relative drift velocity v approaches c , leading to the collisions becoming inelastic, $p + n \rightarrow \pi^+, \pi^0$, in turn leading to positrons, gamma-rays and neutrinos [117]. Such inelastic (p, n) collisions can also arise in jets where the bulk Lorentz factor is transversely inhomogeneous [118], e.g. going from large to small as the angle increases, as expected intuitively from a jet experiencing friction against the surrounding stellar envelope. In such cases, the neutrons from the slower, outer jet regions can diffuse into the faster inner regions, leading to inelastic (p, n) and (n, n) collisions resulting again in pions. An interesting consequence of either radial or tangential (n, p) drifts is that the decoupling generally occurs below the scattering photosphere, and the resulting positrons and gamma-rays deposit a significant fraction of the relative kinetic energy into the flow, reheating it [105]. Internal dissipation below the photosphere has been advocated, e.g. [119] to explain the MeV peaks as quasi-thermal photospheric peaks [120, 121], while having a large radiative efficiency. Such internal dissipation is naturally provided by (p, n) decoupling, and numerical simulations [105] indicate that a Band spectrum and a high efficiency is indeed obtained, which remains the case even when the flow is magnetized up to $\varepsilon_B = 2$ [122], while keeping the dynamics dominated by the baryons. These numerical results were obtained for nominal cases based on a specific radial (n, p) velocity difference, although the phenomenon is generic.

The photon spectral signatures of a magnetically dominated, baryon loaded leptonic plus hadronic GRB model involving nuclear collisions has been calculated by [123]. This uses a realistic transverse structure of a fast core-slow sheath. The analytical results indicate that the transverse neutron collisions become most effective, resulting in GeV photons at radii from which the observer-frame time delay relative to the photospheric MeV photons is appropriate to explain the observed *Fermi* time lags. The purely leptonic (SSC, EIC) time delays and spectral components of such a baryon-loaded magnetic model, in the absence of drifts and transverse gradients, have been calculated by [124], leading to delays in the range observed by *Fermi*.

A hadronic model which attempts to self consistently produce the GeV radiation, the MeV Band spectrum and the low energy (optical) power law is discussed in [125]. The protons accelerated in the shocks or magnetic reconnection regions result in hadronic cascades which produce, as in [115], the GeV and optical power laws, while the cooled leptonic secondaries are re-accelerated via Fermi 2nd order mechanism in the MHD turbulent waves produced by the same shocks or reconnection regions, leading self-consistently to an MeV Band spectrum.

7 Gravitational waves from GRB

GRB may be also sources of gravitational waves (GWs). The most likely such sources are short GRBs [126], if these indeed arise from merging compact objects [7]. The rates in advanced LIGO and VIRGO may be at least several per year [127]. Long GRBs are more speculatively as sources, since in the favored core collapse scenario the collapse may be more chaotic [128]. They may, nonetheless, be weakly detectable as GW sources, especially if the core collapse breaks up into substantial blobs [129], or if they go through a magnetar phase leading to a bar [130]. More recent, detailed numerical calculations of collapsar (long) GRBs lead to GW prospects which range from pessimistic [131] to modest [132].

8 High energy neutrinos from GRBs

High energy ($10^9 \text{ eV} \lesssim E_\nu \lesssim 10^{18} \text{ eV}$) neutrinos may be expected from baryon-loaded GRBs if sufficient protons are co-accelerated in the shocks [133]. The most widely considered paradigm involves proton acceleration and $p\gamma$ interactions in internal shocks, resulting in prompt $\sim 100 \text{ TeV}$ HENUs [134, 135]. Other interaction regions considered are external shocks, with $p\gamma$ interactions on reverse shock UV photons leading to EeV HENUs [136]; and pre-emerging or choked jets in collapsars resulting in HENU precursors [137]. Also, for baryonic dominated GRBs, a neutrino component may arise from photospheric $p\gamma$ and pp interactions [138, 139]. An EeV neutrino flux is also expected from external shocks in very massive Pop. III magnetically dominated GRBs [140]. Current IceCube observations [141, 142, 143] are putting significant constraints on the internal shock neutrino emission model, with data from the full array still to be analyzed. One caveat is that, since the above analysis, several groups [144, 145, 146] have recalculated the GRB internal shock neutrino production in greater detail, including multi-pion and Kaon production in the $p\gamma$ interactions, and allowing for various astrophysical uncertainties including different values of the Lorentz factor and the accelerated proton to electron ratio $L_p/L_e = 1/f_e$. The conclusion from these revised calculations is that the current IceCube (IC40+IC59) measurements need to be extended for another 4 to 9 years for obtaining a strong constraint.

9 Progenitors and Supernova Connection

Including the collimation correction, the GRB electromagnetic emission is energetically quite compatible with an origin in, say, either compact mergers of neutron star-neutron star (NS-NS) or black hole-neutron star (BH-NS) binaries [20, 147, 148, 149], or with a core collapse (hypernova or collapsar) model of a massive stellar progenitor [150, 151, 152], which would be related to but much rarer than core-collapse supernovae [5, 7]. While in both scenarios the outcome could be, at least temporarily, a massive fast-

rotating ultra-high magnetic field neutron star (a magnetar) [153, 14], the mass of the resulting central object exceeds substantially the Chandrasekhar mass and is expected to lead, sooner or later, to the formation of a central black hole. The latter will be fed through a (seconds to minutes) accretion episode from the surrounding disrupted core stellar matter, which provides the energy source for the ejection of relativistic matter responsible for the radiation. This inference appears to be confirmed by numerical simulations, for NS-NS or NS-BH mergers [154, 155, 156] as well as for collapsar models [152, 157, 158].

The above numerical simulations also indicate that (1) the compact object merger accretion disks are less massive and the accretion episode (when the disk is not highly magnetized) lasts less than a few seconds, compatible with the observations of the canonical short GRBs (SGRBs); and (2) in the collapsar models the accretion lasts for tens of seconds or more, compatible with the durations of canonical long GRBs (LGRBs).

The observations of LGRBs indicate that they are generally located in active star-forming environments, usually in blue, small or not too massive, gas-rich galaxies [159, 5, 7], which is where one expects massive stars to be present. Progenitor stars more massive than $\sim 25 - 28 M_{\odot}$, following core collapse, are expected to result in a black hole (BH) central remnant, either directly or through an intermediate neutron star (NS) phase [160, 161]. Such BH core collapse events, if the core is sufficiently fast rotating, can lead to a fall-back fed accretion disk powering a relativistic jet, which is able to escape a star of $\sim 10^{11}$ cm [157, 158, 162]. This radius corresponds to those of Wolf-Rayet (WR) stars, which are thought to arise from more massive $M > 25 M_{\odot}$ progenitors, whittled down by wind mass loss prior to core collapse. The high rotation rate, which favors the wind mass loss and also the formation of a longer lasting accretion disk, is thought to be enhanced when the star arises in a metal-poor environment [5], which in fact seems to characterize many LGRB host galaxies [163, 7].

The massive core collapse model of LGRBs is confirmed by the fact that LGRBs are, in some cases, demonstrably associated with type Ib/c supernovae, whose explosion is, to within errors, contemporaneous with the GRB [164, 165, 166, 167]. The SNe Ib/c are generally thought to have WR progenitors, whether they are associated with GRB or not; only a small fraction of order few % of SNe Ib/c appear to be associated with LGRBs [168, 169, 167]. However, the SNe Ib/c associated with GRBs, as well as a good fraction of those not associated, are classified as hypernovae (HNe), meaning that they have unusually broad spectral lines indicating a semi-relativistic envelope ejecta ($v/c \sim 1$), and inferred isotropic energies $E_{HN,iso} \sim 10^{52.5}$ erg, as opposed to the average SNe with $v/c \sim 0.1$ and $E_{SN,iso} \sim 10^{51}$ erg [170, 171, 172, 173, 174].

For short gamma-ray bursts (SGRB) the most widely favored candidates are mergers of neutron star binaries (DNS) or neutron star-black hole (BH) binaries, which lose orbital angular momentum by gravitational wave radiation and undergo a merger. This second progenitor scenario has only now begun to be tested thanks to the Swift detection of short burst afterglows [175, 176, 177, 178, 179]. The SGRBs are found both in evolved (elliptical) galaxies [175] and in galaxies with star formation [7, 180], in proportions

compatible with that expected for an old population such as neutron stars. While neutron stars are expected, and found, in young star-forming galaxies, massive young stars are not expected, and not found in old population ellipticals. And, indeed, no supernova has ever been found exploding at the same time and location as a SGRB [7]. Of course, neutron stars are the product of supernovae, which can have occurred much earlier than the burst, from progenitors whose initial mass was $8M_{\odot} \lesssim M_{*} \lesssim 25M_{\odot}$. A kick is imparted to the NS at the supernova event, so the NS can wander off significant distances from its birth site (many Kpc, even outside the host galaxy). If the NS was born in a binary and/or later became a binary, the time between the initial explosion and the eventual merger can range up to $10^8 - 10^9$ years, and is very unlikely to be less than 10^6 years [181, 182]. Only a handful of SGRBs have yielded reliable lightcurve breaks suitable for determining a jet opening angle. The latest measurements and comparison to previous data [19] indicate an average $\theta_j \sim 5^\circ$ (comparable to the LGRB average value, although there is one outlier at 25°). This is interesting, since for DNS mergers there is no stellar envelope (as for LGRBs) to provide collimation, at most there would be a wind; however, such narrow jets would be compatible with twisted or hoop-stress collimated MHD jets from DNS mergers [156, 183].

10 High redshift GRBs

Long GRBs are astonishingly bright, both in gamma-rays and at longer wavelengths. In the optical, typical brightnesses are ~ 18 th magnitude a few hours after the trigger (and some have been detected in the 5th-10th magnitude range seconds after trigger), while a Milky-Way-type galaxy has ~ 32 nd magnitude at a redshift $z = 8$. In fact, GRBs vie with galaxies for the record on the highest confirmed redshift measurements, e.g. GRB080913 at $z = 6.7$ [184], GRB090423 at $z = 8.2$ [185, 186] (through spectroscopy), while GRB 090429B has a photometric redshift of $z \simeq 9.4$ [187]. It is possible that even much more distant objects than these have already been detected in the gamma-ray and X-ray detectors of *Swift* and *Fermi*, although for such $z > 9$ objects a specific (optical/IR or other) redshift signature is extremely difficult and noisy, so redshift diagnostics are increasingly harder to obtain in this range. The above discoveries do, however, indicate that the prospect of eventually reaching into the realm of Pop. III objects is becoming increasingly realistic.

Population III GRBs at $10 \lesssim z \lesssim 20$ may result, as they do at lower redshifts, from massive, $M > 25 - 30M_{\odot}$ metal-poor stars whose core collapses to a black hole [188]. However, the mass of Pop. III stellar progenitors could be as high as $\sim 1000M_{\odot}$, leading to $100 - 500M_{\odot}$ black holes [189], although the Pop. III masses are a subject of debate (and could be much lower [190, 191]). For extremely massive black holes, the jets are likely to be Poynting-dominated, e.g. powered by the Blandford-Znajek mechanism. The expansion dynamics and the radiation arising from such very massive Poynting jet GRBs was discussed by [192]. At typical redshifts $z \sim 20$ this implies a “prompt” emission extending to $\lesssim 1$ day which should be detectable by *Swift* or *Fermi*, being

most prominent initially around 50 keV due to the jet pair photosphere, followed after a similar time interval by an external shock synchrotron component at a few keV and an inverse Compton component at $\gtrsim 70$ GeV [193]. Both the ‘prompt’ emission and the longer-lasting afterglows [193] of such Pop. III GRBs should be detectable with the BAT or XRT on *Swift* or the GBM on *Fermi*. On *Swift*, image triggers may be the best way to detect them, and some constraints on their rate are provided by radio surveys. They are expected to have GeV extensions as well, but redshift determinations need to rely on L-band or K-band spectroscopy.

11 The Cosmology Connection

Since GRBs are seen out to the largest redshifts yet measured, and for periods of hours to days they can outshine any other objects at those distances, their potential usefulness as tools for cosmology has been intriguing for some time.

The simplest way, to use them as distance markers, is unfortunately not straightforward. This is because they are not good “standard candles”, which could be used e.g. in a Hubble-type diagram to compare flux against redshift to compare against cosmological models to deduce a closure parameter, or an acceleration rate. Even the collimation-corrected average gamma-ray energy $E_{\gamma,j} \sim 10^{51}$ ergs has too much variance to be of direct use as a standard candle. The error is still almost twice as large as that obtained from SNIa, at least at redshifts $z \lesssim 1.5$, which is the most important region for dark energy studies. The hope is, still, that one or more of the various empirical correlations between observed spectral quantities could lead to a calibration, as the Phillips relation does for supernovae, which could turn GRBs into an effective distance ruler.

One such empirical correlation is between the photon spectral peak energy E_{pk} and the apparent isotropic energy E_{iso} [69, 194], namely $E_{pk} \propto E_{iso}^\alpha$ with $\alpha \sim 1/2$ (Amati relation). Other correlations are between E_{pk} and the peak luminosity L_{pk} [195] (Yonetoku relation); or for bursts where the jet opening angle is known, between E_{pk} and the collimation corrected gamma-ray energy of the jet $E_{\gamma,jet}$ [196, 197] (Ghirlanda relation); or between E_{pk} and $L_{\gamma,jet}$ [198]; or between E_{pk} , E_{iso} and the light curve break time t_{br} [199] (Liang-Zhang relation); or between the X-ray luminosity at break time $L_{X,br}$ and the break time t_{br} [200]. Of course, these correlations are of interest in themselves as possible constraints on the radiation mechanism or emission region, and various interpretations have been made, e.g. [201, 119, 71, 202]. However, for cosmology only the tightness of the empirical correlation is what matters. Observational biases could, of course, pose problems [203, 204], and circularity issues may be a concern; the latter, however can in principle be minimized by restricting oneself only to directly observed quantities [205].

While such correlations make it possible to extend Hubble diagrams up to redshifts $z \lesssim 8 - 9$ with GRBs, their usefulness for deriving cosmological parameters is still limited. The main reason is that the matter (dark plus baryonic) and the vacuum energy densities evolve as $\rho_m \propto (1+z)^3$, $\rho_v \propto (1+z)^{3(1+w)}$, respectively, where w is the

equation of state multiplier, $p = w\rho$, for vacuum $w = -1$ being the currently favored value. At present the vacuum energy dominates the dynamics, $\Omega_{V,0} = (\rho_{V,0}/\rho_{cr,0} \simeq 0.7$ while $\Omega_{m,0} \simeq 0.3$. Going back in time or up in redshift, the vacuum energy grows slowly (or not at all, if $w = -1$), while the matter energy grows fast, so that for redshifts $z > z_V = [(\Omega_{V,0}/\Omega_{m,0})^{-1/3w} - 1] \sim 0.5$ the vacuum energy becomes negligible compared to that of matter. Thus, the information about the vacuum-related acceleration and a vacuum equation of state is gleaned mainly from data at $z \lesssim 0.5$, where the SNIa with redshifts greatly outnumber the GRBs with redshift. Some recent papers using GRBs in Hubble diagrams are, e.g. [206] and [207], the latter including 109 GRBs with known redshifts calibrated with 567 SNIa. While strongly suggestive, and increasingly interesting, the sample is still small compared to SNe Ia, especially at low z , and the dispersion remains larger then for SNe, so the usefulness of GRBs as statistical distance indicators remains to be seen.

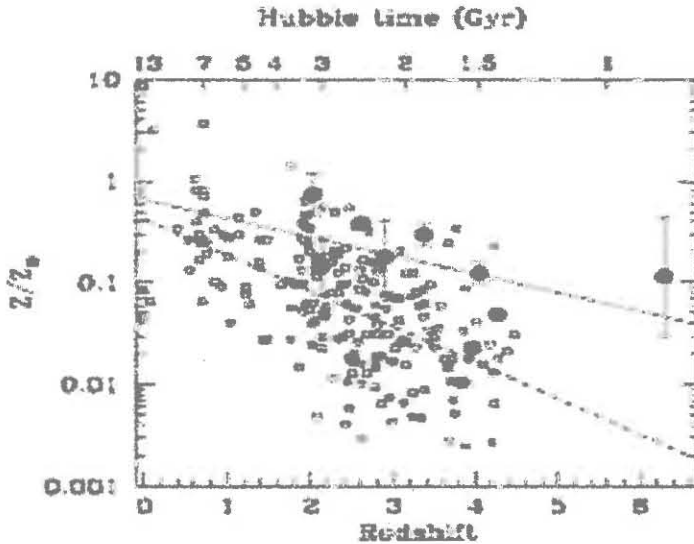
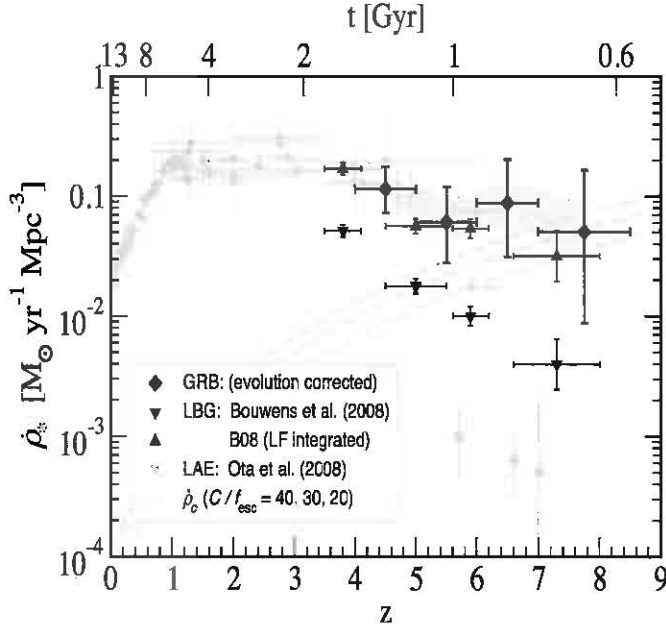


Figure 6: Redshift evolution of the metallicity relative to solar values, for GRBs shown with blue dots and QSOs show with open circles. The GRB metallicity is on average ~ 5 times larger than in QSO. These are based on damped Lyman alpha (DLA) spectral features. The upper horizontal x-axis indicates the age of the Universe (Hubble time) [208].

GRBs, however, are likely to be unique as beacons for probing the high redshift Universe. They are detectable with current gamma-ray, X-ray and infra-red detectors out to distances corresponding to the earliest stars formation epoch [209, 210], and they may provide possible redshift signatures [211, 212] extending into the $10 \lesssim z \lesssim 20$ range. Their strong, featureless power law continuum spectrum shining through the intergalactic medium and intervening young galaxies or protogalaxies provide a sensitive probe of the ionization state, velocity distribution and chemical composition at those redshifts [213, 214, 215]. An example is shown in Fig. 6, indicating the change of the metallicity (given by the oxygen to hydrogen abundance ratio) as a function of redshift. These abundances are determined from spectral absorption lines in the continuum radiation of GRBs and quasars. The GRB lines come mainly from the host galaxy gas in the star forming region where the explosion occurred, while the quasar lines arise in random intervening galaxies along the line of sight. One sees that the GRB provide information



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Figure 7: Cosmic star formation history [216]. Shown are the data compiled in [217] (light circles) and from Ly- α emitters (LAE) [218]. Also recent LBG data for two UV LF integrations: down to $0.2L_*$ (down triangles, [219]), and complete up to $z = 3$ (up triangles). Swift GRB-inferred rates are diamonds, the shaded band shows the range resulting from varying evolutionary parameters. Also shown is the critical ρ_* from [220] for $C/f_{esc} = 40, 30, 20$ (dashed lines, top to bottom), see also [221].

GRBs also provide an excellent tool for investigating the cosmic star formation rate (SFR) of the high redshift universe, and thereby also the rate of large scale structure (LSS) formation, out to (so far) redshifts in the 8-10 range. This is exemplified in Fig. 7, which shows the star formation rate determined through various techniques. The long GRBs are the endpoints of the lives of massive stars, and their rate is therefore approximately proportional to the star formation rate in general. However, at high redshifts the rates are very uncertain, and may be subject to various observational biases. There may also be evolutionary biases, such as a dependence of the long GRB formation on the metallicity of the host galaxy, which needs to be taken into account, e.g. [214, 215].

The most distant GRBs may also provide the only possible probes of the era of the first generation of (Population III) stars formed in the Universe [189, 192, 193, 222]. These relics of the infant universe could be the most sensitive probes of the redshift for the start of large scale structure formation, with significant implications for the properties of the dark matter.

12 The future

Both *Swift* and *Fermi* are likely to be functional and return GRB data for many years to come. They have orbital lifetimes extending beyond 2025 and no critical expendables.

The *SVOM* mission [4] is an approved Chinese-French mission to observe GRBs and their afterglow. It has a wide-field instrument to image the bursts and one to study the spectrum. The spacecraft rapidly slews like *Swift* to point X-ray and optical telescopes for afterglow observation. There are other concepts in consideration, such as *Lobster* [223], which performs the wide-field observations in the X-ray band suitable for high-redshift GRBs. In addition, combined with such electromagnetic detection programs, increasingly sensitive multi-messenger detection attempts will continue to be pursued using high energy neutrinos (§8) and gravitational waves (§7). These expanded observational efforts will require more detailed theoretical interpretation and models, extending well beyond what has been achieved so far. Based on past experience, GRBs may be counted on to provide further exciting surprises during the next decade.

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