REGIONAL SPHERICAL HARMONIC MAGNETIC MODELING FROM NEAR-SURFACE AND SATELLITE-ALTITUDE ANOMLAIES



Abstract

Hyung Rae Kim¹, Ralph R. B. von Frese², and Patrick T. Taylor³ tt. of Geoenvormental Sciences, Kongju National University, Shinkwan-dong, 182, Gongju, congchungNam-Do, Korea (kimhr@kongju.ac.kr); tool of Earth Sciences, 275 Mendenhall Lab., The Ohio State Univ., Columbus, Ohio, USA netary Geodynamics Lab, Goddard Space Flight Center, NASA, Greenbelt, Maryland, USA



e compiled near-surface data and satellite crustal magnetic measured data are modeled with a regionally concentrated spherical harmonic presentation technique over Australia and Antarctica. obal crustal magnetic anomaly studies have used a spherical harmonic analysis to represent the Earth's magnetic crustal field. This global approach, however is best applied where the data are formly distributed over the entire Earth. Statellite observations generally meet this requirement, but unequally distributed data cannot be easily adapted in global modeling. Even for the satellite servations, due to the errors spread over the globe, data smoothing is inevitable in the global spherical harmonic presentations. In addition, global high-resolution modeling requires a great numbe global spherical harmonic coefficients for the regional presentation of crustal magnetic anomalies, whereas a lesser number of localized spherical coefficients will satisfy. We compared methods in th global and regional approaches and for a case where the errors were propagated outside the region of interest. For observations from the upcoming Swarm constellation, the regional modeling I allow the production a lesser number of spherical coefficients that are relevant to the region of interest.

Introduction

Studies on geomagnetism have used a spherical harmonic analysis as a tool of representing Earth's magnetic fields both of internal and of external. Therefore, geomagnetic features have been modeled by spherical harmonics in a global scale. However, this global approach is best applied where the data are uniformly distributed over the entire Earth such as polar-orbiting satellite measurements. Until now, these satellite observations are good for regional or larger-scale studies due to the limited resolutions at such altitudes. For the detailed features of the magnetic anomalies, unequally distributed data measured at the near-surface are useful but cannot be easily adopted in global modeling. A spherical cap analysis (Haines, 1985) has been used as an alternative in regional modeling for geopotential fields and recently this technique was well updated by Thebault et al. (2006) and Thebault (2008). In this study, we introduce a more straightforward regional analysis technique as a magnetic field modeling that has been effectively implemented to model satellite gravity field data for small cap regions of the Earth and Moon (Han et al., 2008). We tested to compare between the methods in both global and regional approaches and a case that the errors are propagated outside the region of interest.

Theory of Spherical Slepian basis functions

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Slepian (1983) showed that optimally concentrated, band-limited and time-limited signal is considered by maximizing the energy inside the region (or interval) to the whole area. This concentration problem turned out to be the association of the eigenvalues and eigenvectors of an integral matrix form of spherical harmonics. Such basis functions are orthogonal each other and satisfy the same differential equation and boundary condition as the spherical harmonics on a sphere. These are in real the linear combinations of conventional spherical harmonic functions only concentrating the area of interest.

Given a band-limited signals $f(\theta, \varphi)$ on a sphere in terms of the usual Spherical Harmonics with a relevant set of coefficients, these signals also can be represented by a set of Slepian basis functions and expansion coefficients, respectively, as follows:

$$\Gamma(\theta, \varphi) = \sum_{n=0}^{N} \sum_{m=-n}^{n} g_{nm} Y_{nm}(\theta, \varphi) = \sum_{k=1}^{(N+1)^2} c_k s_k(\theta, \varphi)$$
(1)

n=0 m=-n k=1where θ and ϕ are colatitude and longitude of the sphere, respectively. $Y_{not}(\theta,\phi)$ is spherical harmonic function of degree n and order m, $g_{m,i}$ is the SH coefficient of the same degree and order, respectively. These signals can also be represented by the k-th Slepian function, s_k and its expansion coefficients, c_k , respectively. N is the highest degree of spherical harmonic expansion used. The k-th Slepian function is the expression of linear combination of spherical harmonic functions so that it can be rewritten as;

$$s_k(\theta,\varphi) = \sum_{n=0}^{N} \sum_{m=n}^{n} s_{k,nm} Y_{nm}(\theta,\varphi)$$
(2)

The new spherical harmonic coefficients, s_{klm} of new basis function, s_k (θ, φ) is now determined by maximizing the energy concentration over the area of interest to the energy over the entire sphere, which is as follows:

$$\mathbf{u}_{k} = \int_{0}^{2\pi} \int_{0}^{\theta_{0}} s_{k}^{2}(\theta, \varphi) \sin \theta d\theta d\varphi \Big/ \int_{0}^{2\pi} \int_{0}^{\pi} s_{k}^{2}(\theta, \varphi) \sin \theta d\theta d\varphi$$
(3)

where λ_e is the ratio of energy concentrated within the spherical cap (θ_a) compared to the entire sphere . Grunbaum et al. (1982) and Simons and Dahlen (2006) showed that the solution to this problem turned out an eigen problem in a spectral domain. With using (2) and orthonormality properly, this can be straightforwardly expressed as:

$$\sum_{n=0}^{N} \sum_{m=-n}^{n} D_{nmn'm'} s_{k,nm} = \lambda_k s_{k,n'm'}$$
(4)

$$D_{nmn'm'} = \int_0^{2\pi} \int_0^{\theta_0} Y_{nm}(\theta, \varphi) Y_{n'm}'(\theta', \varphi') \sin \theta d\theta d\varphi$$
(5)

The $\lambda_{\rm s}$ and $s_{\rm Lim}$ should be determined by the eigenvalues and eigenvectors of $D_{\rm ma,n}$ with the geometric bounds of the area of interest (or radius of spherical cap). This area can also have irregular bounds on the sphere such as Australia or Africa continents. It is noted that the total number of $(N+1)^2$ eigenvalues and eigenvectors exist but the necessary ones to optimally represent the maximum energy of signals over the area of interest are much fewer. That is, some of the largest eigenvalues (close to 1) and the corresponding eigenvectors are chosen to use as a set of basis functions concentrated within the area of interest.



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