

Phased Array Beamforming and Imaging in Composite Laminates Using Guided Waves

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ABSTRACT

This paper presents the phased array beamforming and imaging using guided waves in anisotropic composite laminates. A generic phased array beamforming formula is presented, based on the classic delay-and-sum principle. The generic formula considers direction-dependent guided wave properties induced by the anisotropic material properties of composites. Moreover, the array beamforming and imaging are performed in frequency domain where the guided wave dispersion effect has been considered. The presented phased array method is implemented with a non-contact scanning laser Doppler vibrometer (SLDV) to detect multiple defects at different locations in an anisotropic composite plate. The array is constructed of scan points in a small area rapidly scanned by the SLDV. Using the phased array method, multiple defects at different locations are successfully detected. Our study shows that the guided wave phased array method is a potential effective method for rapid inspection of large composite structures.

Keywords: guided waves; phased arrays; anisotropic composites; array imaging; multiple defects; rapid inspection

1. INTRODUCTION

Advanced composite materials will likely play an important role in NASA's development of next generation vehicles for long duration space missions in order to enable lightweight space structures. NASA is currently investigating large scale composites for human spaceflight, such as a composite crew module and a composite fuel tank for rockets [1]. Rapid inspection techniques for detecting and quantifying damage in large composites are critical for ensuring operability and safety of composite structures [2, 3]. Though various damage inspection methods have been developed for metallic structures, reliable and efficient evaluations of large composite structures are not yet well established [3, 4].

For rapid damage inspection, an image of the structure being inspected often gives an efficient solution that quickly identifies and locates defects. Among various imaging methods, the guided wave phased array imaging is attractive since it uses sensors that are placed closed to each other in a compact format, steers the outputs of all sensors in a desired direction, and scans the entire structure like a radar [5], allowing for large area inspection with limited access. Intensive study has been conducted on the guided wave phased array beamforming and damage detection on metallic plate like structures [5-33]. Some researchers have started investigating the phased arrays for anisotropic composite materials [34-39]. It has been found that the guided wave complexity caused by anisotropic and inhomogeneous properties in composite materials makes the traditional metal-based methods inappropriate and sometimes even misleading [40-42].

In this paper, we investigate guided wave beamforming in anisotropic laminated composite plates. Based on the classic delay-and-sum principle, a generic formula of phased array beamforming in anisotropic composite plates is developed, in which the direction dependent guided wave properties are considered. Moreover, the array beamforming and imaging are performed in frequency domain where the guided wave dispersion effect has been considered. The presented phased array method is implemented with a non-contact scanning laser Doppler vibrometer (SLDV). For the proof of concept, laboratory tests are performed using rectangular grid arrays made of scanning points of the non-contact SLDV for detecting multiple surface defects. The results show that multiple defects at various directions can be successfully detected using the non-contact SLDV array. The remainder of this paper is organized as follows: Section 2 presents the formulation of phased array beamforming in anisotropic composite plates; Section 3 presents the

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implementation of phased arrays by using a non-contact SLDV for detecting multiple defects in a composite plate; Section 4 concludes the paper with novelties, discussions and planned future work.

2. GUIDED WAVE BEAMFORMING IN COMPOSITE LAMINATES

This section presents the guided wave beamforming in anisotropic composite laminates. The generic beamforming formulas have been derived based on the classic delay-and-sum principle. Compared to the array beamforming of circular wavefront guided waves in isotropic materials, the noncircular wavefront characteristic of guided waves in anisotropic materials is considered in the generic beamforming formulas.

2.1 Guided waves in composite laminates

When a guided wave with frequency ω and wavenumber \mathbf{k} is generated from a source at the coordinate origin O in a composite plate, the wave arriving at the location \mathbf{x} that is far away from the source (Figure 1a) can be expressed as [42-44],

$$u(t, \mathbf{x}) = Ae^{j(\omega t - \mathbf{k} \cdot \mathbf{x})} \quad (1)$$

where A is the amplitude, assumed to be independent of wave frequency. With the geometric relation illustrated in Figure 1a, we have,

$$\mathbf{k} \cdot \mathbf{x} = |\mathbf{k}| |\mathbf{x}| \cos \beta = k(\gamma) |\mathbf{x}| \cos \beta \quad (2)$$

with β being the angle between the wave propagation and wavenumber \mathbf{k} . Hence,

$$u(t, \mathbf{x}) = Ae^{j[\omega t - k(\gamma) |\mathbf{x}| \cos \beta]} \quad (3)$$

Using Eq.(1), for a source located at location \mathbf{p}_m , the wave resulted at the location \mathbf{x} is,

$$u_m(t, \mathbf{x}) = Ae^{j[\omega t - \mathbf{k} \cdot (\mathbf{x} - \mathbf{p}_m)]} \quad (4)$$

In anisotropic composite laminates, guided wave parameters such as wavenumbers, phase velocities and group velocities are direction dependent, due to the direction dependent physical properties of composite materials [40-42]. Figure 1b plots the wavenumber curve $k(\gamma)$. As illustrated in Figure 1b, the wavenumber vector \mathbf{k} is perpendicular to the wave front and the group velocity vector \mathbf{c}_g is orthogonal to the wavenumber curve $k(\gamma)$ [40-42]. The angle γ of the wavenumber vector \mathbf{k} is referred to as wavenumber angle. The angle θ of the group velocity vector \mathbf{c}_g is referred to as group velocity angle (or energy propagation angle). The angle β between wavenumber angle γ and energy propagation angle θ is referred to as skew angle, with the relation $\beta = \gamma - \theta$. It can be seen when \mathbf{c}_g is not parallel to \mathbf{k} , the skew angle β is not zero and hence the wave energy propagation direction is not perpendicular to the wave front.

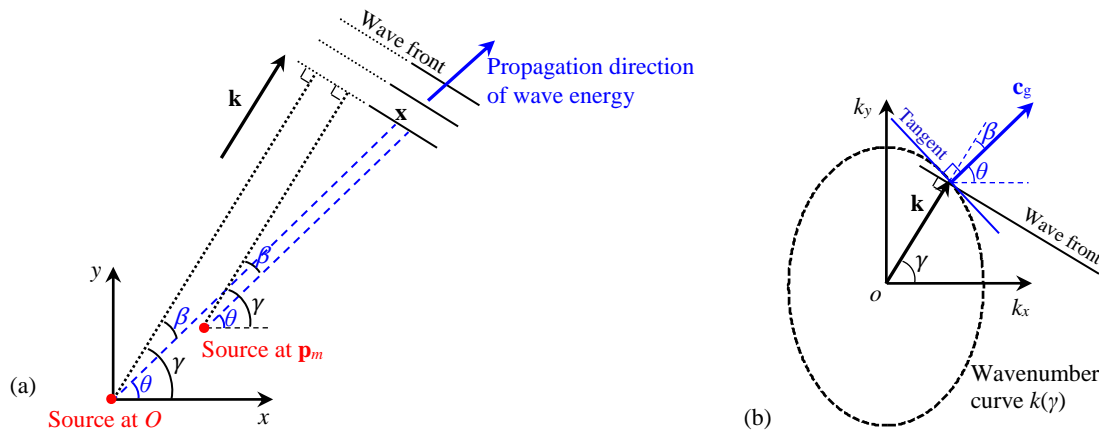


Figure 1 Schematics of geometric relations of guided waves in composite laminates: (a) the geometric relation of GW (in the far field) generated by sources at different locations; (b) wavenumber curve. The geometric relation between the wavenumber vector \mathbf{k} and the group velocity vector \mathbf{c}_g is given in figure b.

2.2 Delay-and-sum beamforming

Consider an array with M identical elements located at $\{\mathbf{p}_m\}$ ($m=0, 1, 2, \dots, M-1$) which are geometrically close to each other. The phase center is defined as the origin O of the Cartesian coordinate system, i.e., $\frac{1}{M} \sum_{m=0}^{M-1} \mathbf{p}_m = \mathbf{0}$. Each element serves as a wave source. When all elements generate waves with frequency ω and wavenumber vector \mathbf{k} simultaneously, using Eq. (4) the total output (synthesized wave) of the array at location \mathbf{x} can be derived as,

$$z(t, \mathbf{x}) = \sum_{m=0}^{M-1} A e^{j[\omega t - \mathbf{k} \cdot (\mathbf{x} - \mathbf{p}_m)]} = u(t, \mathbf{x}) \sum_{m=0}^{M-1} e^{j\mathbf{k} \cdot \mathbf{p}_m} \quad (5)$$

It is seen from Eq. (5) that the synthesized wave $z(t, \mathbf{x})$ is an amplification of the wave $u(t, \mathbf{x})$ emitted from the Origin. The amplification is controlled by the exponential component $\sum_{m=0}^{M-1} e^{j\mathbf{k} \cdot \mathbf{p}_m}$ in Eq. (5). Therefore, by adjusting the component $\sum_{m=0}^{M-1} e^{j\mathbf{k} \cdot \mathbf{p}_m}$, we can control the amplification. One way to adjust the component $\sum_{m=0}^{M-1} e^{j\mathbf{k} \cdot \mathbf{p}_m}$ is applying phase delays to all exponents. For example, to maximize the amplification $\sum_{m=0}^{M-1} e^{j\mathbf{k} \cdot \mathbf{p}_m}$ in a specific direction θ_s , we can apply phase delays $\Delta_m(\theta_s)$ to each exponent to make each exponent zero,

$$\sum_{m=0}^{M-1} e^{j[\mathbf{k} \cdot \mathbf{p}_m - \Delta_m(\theta_s)]} = \sum_{m=0}^{M-1} e^{j0} = M, \text{ where } \Delta_m(\theta_s) = \mathbf{k} \cdot \mathbf{p}_m \quad (6)$$

Using phase delays $\Delta_m(\theta_s)$, the amplification is maximized at the direction θ_s , i.e., generating a directional “beam”. θ_s is also known as steering angle. Eq. (6) shows that phase delays $\Delta_m(\theta_s)$ depend on the m^{th} element’s position vector \mathbf{p}_m and wavenumber vector \mathbf{k} .

In anisotropic composites, the wavenumber vector \mathbf{k} depends on wave frequency ω and wavenumber angle γ_s . In addition, the geometry relation between wavenumber angle γ_s and steering angle θ_s is $\gamma_s = \theta_s + \beta_s$, where β_s is referred to as skew angle. Hence, phase delays $\Delta_m(\theta_s)$ in Eq.(6) can be further expressed as,

$$\Delta_m(\theta_s) = \mathbf{k}(\omega, \theta_s + \beta_s) \cdot \mathbf{p}_m \quad (7)$$

It can be seen from Eq.(7) that the phase delay includes a frequency dependent term, $\mathbf{k}(\omega, \theta_s + \beta_s)$. Hence, it is frequency dependent and has the advantage of compensating dispersion effects (a.k.a. frequency dependent wave properties) by using frequency-related components during the delaying [14], [45]. This is otherwise not readily achievable in the commonly used time delay methods [5, 11, 34].

In addition to phase delays, weighting factors w_m can also be applied to delayed waves to further control the quality of beamforming [43, 46]. In summary, with phase delays and weighting factors, the beamforming is represented as,

$$z(t, \mathbf{x}) = u(t, \mathbf{x}) \sum_{m=0}^{M-1} w_m e^{j[\mathbf{k} \cdot \mathbf{p}_m - \Delta_m(\theta_s)]} = u(t, \mathbf{x}) \sum_{m=0}^{M-1} w_m e^{j[\mathbf{k} \cdot \mathbf{p}_m - \mathbf{k}(\omega, \theta_s + \beta_s) \cdot \mathbf{p}_m]} \quad (8)$$

2.3 Array beamforming in a CFRP composite plate

To give a direct view and better understanding of the array beamforming in anisotropic composites, wavefields are simulated by using equations given in subsection 2. The simulation specimen is an 8-ply $[0_2/90_2]_s$ CFRP composite plate. The material properties for a single ply are given in Table 1. The A_0 mode at 90 kHz is chosen for the simulation (its wavenumber curve is plotted in Figure 2).

Figure 3a plots the wavefield calculated using Eq. (1), in which the amplitude $A(\omega)$ is assumed to be unity. This wavefield is for the 90 kHz A_0 mode generated by a point source at the coordinate origin in the $[0_2/90_2]_s$ CFRP plate. As shown in the wavefield in Figure 3a, waves propagate to all directions from the actuator at the coordinate origin. Figure 3b plots the synthesized wavefield calculated using Eq. (8). This wavefield represents guided waves generated from a rectangular grid array (16×16 elements with the array spacing 2 mm), when the wave energy steering angle θ_s is 45° . Compared to the wavefield in Figure 3a, the synthesized wavefield in Figure 3b generated from the array shows directional guided waves in the direction $\theta=45^\circ$.

Table 1 Material properties

ρ (kg/m ³)	E_1 (GPa)	E_2 (GPa)	E_3 (GPa)	G_{12} (GPa)	G_{13} (GPa)	G_{23} (GPa)	ν_{12}	ν_{13}	ν_{23}	Thickness (mm)
1570	171.4	9.08	9.08	5.29	5.29	2.80	0.32	0.32	0.34	0.1068

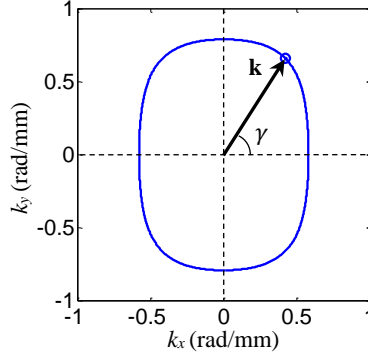


Figure 2 The wavenumber curve of A_0 mode at 90 kHz in the composite laminate with $[0_2/90_2]_s$ layup

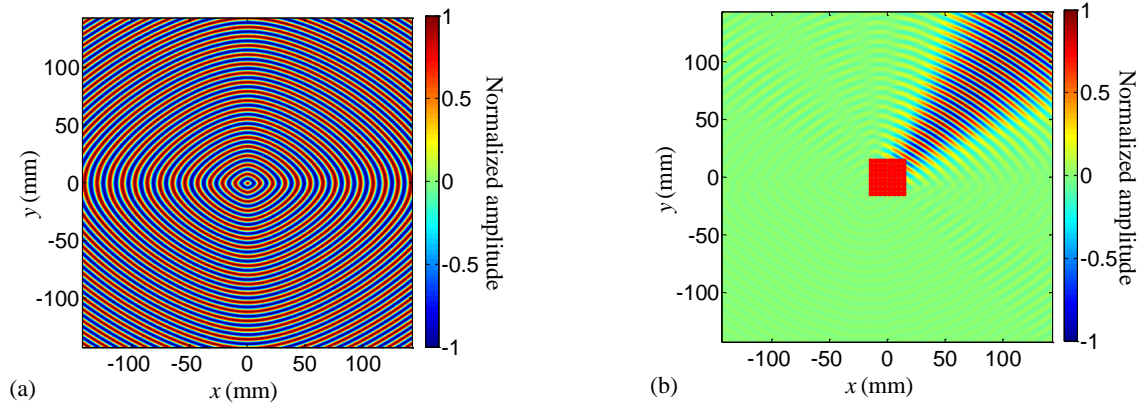


Figure 3 Simulated wavefields of the 90 kHz A_0 mode in a composite plate with $[0_2/90_2]_s$ layup: (a) generated by a single point source at the origin, (b) generated by rectangular grid array (16×16 elements) when the wave steering angle θ_s is 45° .

3. DETECTION OF MULTI-DEFECTS IN A CFRP COMPOSITE PLATE

In this section, guided wave beamforming and damage detection are implemented for a laminated composite plate using rectangular grid arrays. A hybrid measurement PZT-SLDV system consisting of a surface bonded PZT wafer and a non-contact scanning laser Doppler vibrometer (SLDV) is used for guided wave actuation and sensing [47]. Phased arrays are constructed by selected scanning points of the SLDV. Four quartz rods are surface bonded on a composite plate at various angular positions as defects to be detected by the phased array method.

3.1 Experimental setup

Figure 4 plots a schematic of experimental setup for the PZT-SLDV phased array sensing. The test specimen is an 8-ply $[0_2/90_2]_s$ layup composite laminate with dimensions of $380 \times 380 \times 0.85$ (unit: mm). The material properties of a single ply are listed in Table 1. The plate center is set as the coordinate origin. Four defects (D_1 , D_2 , D_3 and D_4) are simulated by using four quartz rods with the same size bonded on the plate surface with the same distance 100 mm to the

plate center while different angles 0° , 45° , 90° and 135° . A PZT wafer (APC 851: 7 mm diameter, 0.2 mm thickness) is bonded on the plate center to generate guided waves. A 3-cycle toneburst with the frequency 90 kHz is generated from the function generator (model: Agilent 33522B), amplified to 30V by a voltage amplifier (model: Krohn-Hite 1506), and applied to the PZT wafer for excitation. An SLDV (model: Polytec PSV-400-M2) is used to acquire velocity components of guided waves at selected scanning points, based on Doppler Effect. The laser beam is set perpendicular to the plate such that the out-of-plane velocity of guided waves is measured.

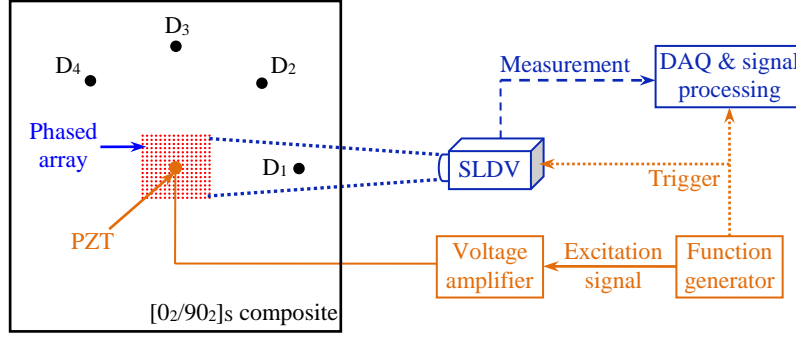


Figure 4. The schematic of experimental setup for multi-defects detection in a $[0_2/90_2]_s$ CFRP composite laminate.

3.2 Array imaging algorithms

The phased array is constructed of SLDV scanning points at selected locations $\{\mathbf{p}_m\}$ ($m=0, 1, 2, \dots, M-1$), whose phase center satisfies $\frac{1}{M} \sum_{m=0}^{M-1} \mathbf{p}_m = \mathbf{0}$. At the m^{th} array point (\mathbf{p}_m), the acquired signal by the SLDV is denoted as $v_m(t)$. The signal's frequency spectrum is obtained using the Fourier transform, as:

$$V_m(\omega) = F[v_m(t)] = \int_{-\infty}^{\infty} v_m(t) e^{-j\omega t} dt \quad (9)$$

Using the frequency spectrum $V_m(\omega)$, the synthesized frequency-space representation $Z(\omega, \mathbf{x})$ of the array [14] is:

$$Z(\omega, \mathbf{x}) = \sum_{m=0}^{M-1} w_m V_m(\omega) e^{j[-\varphi(\omega, \mathbf{x}) - \Delta_m(\omega, \mathbf{x})]} \quad (10)$$

where,

$$\Delta_m(\omega, \mathbf{x}) = \mathbf{k}(\omega, \gamma) \cdot \mathbf{p}_m, \text{ and } \varphi(\omega, \mathbf{x}) = -2\mathbf{k}(\omega, \gamma) \cdot \mathbf{x} \quad (11)$$

$\Delta_m(\omega, \mathbf{x})$ is the phase delay applied to the m^{th} array point for beamsteering, and $\varphi(\omega, \mathbf{x})$ represents the spatial phase shift. As guided waves travel from the PZT to the damage and then back to the array, they undergo a spatial phase shift $\varphi(\omega, \mathbf{x})$. Thus, $-\varphi(\omega, \mathbf{x})$ is applied in Eq. (10) in order to compensate for the spatial phase shift. $\mathbf{k}(\omega, \gamma)$ is the wavenumber vector at the frequency ω and wavenumber angle γ , and is obtained from wavenumber dispersion curves. γ is the wavenumber angle determined from the geometry relation as $\gamma = \theta + \beta$, where θ is the energy propagation angle and equals the angle of vector \mathbf{x} , i.e., $\theta = \angle \mathbf{x}$, and β is the skew angle corresponding to the energy propagation angle θ . Weighting factors w_m can be used to further control the quality of the beamforming such as the beam shape [43, 46]. In this study, weighting factors are set equal to one.

Using the inverse Fourier transform, the frequency-space representation $Z(\omega, \mathbf{x})$ can be transformed back to the time-space domain, as:

$$z(t, \mathbf{x}) = F^{-1}[Z(\omega, \mathbf{x})] = \frac{1}{2\pi} \int_{-\infty}^{\infty} Z(\omega, \mathbf{x}) e^{j\omega t} d\omega \quad (12)$$

where $z(t, \mathbf{x})$ is a synthesized time-space wavefield that represents the array beamforming in the time-space domain. An inspection image of the plate to detect/locate any existing damage is then acquired by building an intensity image with the pixel value at the location \mathbf{x} defined as:

$$I(\mathbf{x}) = |z(t=0, \mathbf{x})|^2 \quad (13)$$

3.3 Array imaging results

In this test, a rectangular phased array is used for damage detection. The array is constructed of 31×31 SLDV scan points with the spacing 2 mm and the array center at the origin. Using the phased array method presented in subsection 3.2, an intensity image can be generated with the pixel definition given in Eq. (13). Figure 5a plots the imaging result when the sizes of four defects are 4 mm (defects are simulated by four quartz rods with the same diameter of 4 mm and length of 10 mm). The defects show up as highlighted spots in the image providing a useful means to detect as well as to locate their positions. Hence the four defects are located by looking for the local maximum pixel value (normalized amplitude): D_1 at (99.0, 0.8) mm with errors of (1.0, 0.8) mm, D_2 at (74.4, 69.3) mm with errors of (3.7, 1.4) mm, D_3 at (0.2, 101.0) mm with errors of (0.2, 1.0) mm, and D_4 at (-71.9, 69.9) mm with errors of (1.2, 0.8) mm.

Besides the case where the all four defect are size 4 mm, other cases with defect sizes increased to 6, 8 and 10 mm are also investigated. The array imaging results are plotted in Figure 5b, c and d, which show that the four defects are successfully detected. For comparison purpose, the image amplitudes (local maximum normalized amplitude) at the four detected areas are extracted and plotted with respect to the damage size in Figure 6, showing increasing amplitude trends with the increase of damage size. Moreover, Figure 6 shows how the amplitude varies with the inspection angle. The plot for 0° has the highest amplitude, while the plot for 135° (close to the plot for 45°) gives the smallest amplitude. The damage at 0° has the highest amplitude due to the fact that the guided waves have the strongest propagation along this direction in the subject composite plate. In contrast, the amplitudes in 45° and 135° are very weak due to the fact that guided waves in these two directions are weak.

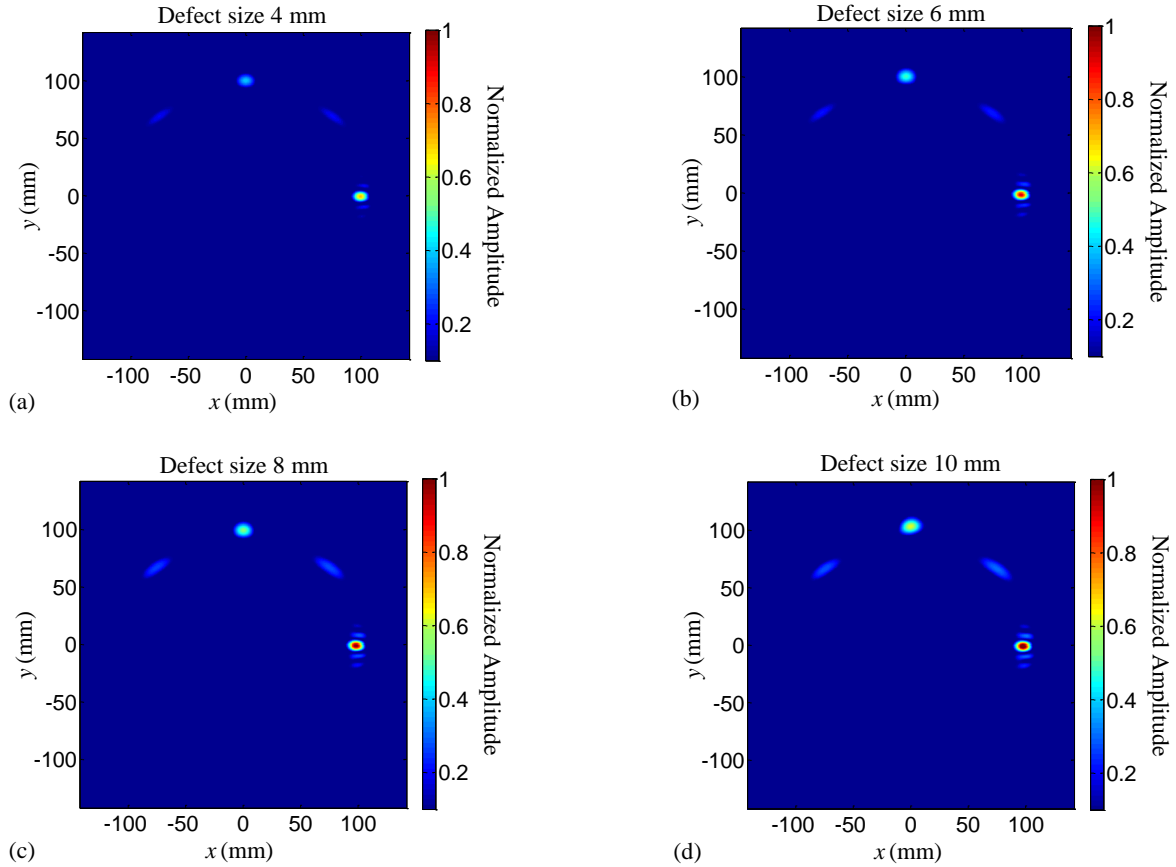


Figure 5. Array imaging results of the four defects. Sizes of defects are (a) 4 mm, (b) 6 mm, (c) 8 mm and (d) 10 mm.

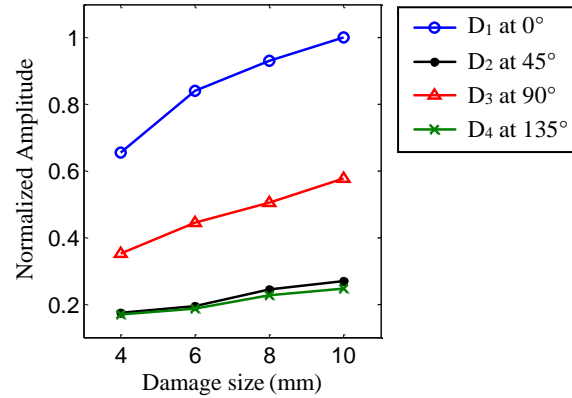


Figure 6. Image amplitude versus damage size. The four lines represent amplitudes at different inspection angles.

4. CONCLUSIONS

In this paper, we have presented a generic guided wave phased array beamforming method for rapid damage inspection in anisotropic composite laminates. The generic beamforming formula uniquely considers the direction-dependent guided wave properties and the energy skew effect; thus it allows for phased array beamforming in anisotropic composites without requiring any quasi-isotropic assumption. The phased array beamforming is performed in the frequency domain which offers the advantage of compensating the guided wave dispersion effect through the use of frequency-dependent wavenumber.

The presented phased array method is implemented with a non-contact SLDV system which allows for the flexibility of using selected points from the SLDV scan to construct phase arrays of various configurations. The high spatial resolution feature of the SLDV system also significantly reduces array element spacing that gives the potential of using guided waves with smaller wavelengths to detect smaller defects in composite structures. We have shown that the present method can detect multiple defects at various locations in an anisotropic CFRP composite plate.

Despite the success of phased array beamforming and multiple defect detection in a composite plate presented here, further studies are still needed. Besides the selected 90 kHz A_0 mode used in the current study, additional research needs to be performed with guided waves at other frequencies and/or different wave modes in highly anisotropic composite plates. The effectiveness of the present method should be also evaluated with realistic defects such as impact induced delamination in composite plates, and studies of the smallest detectable defect size should be performed. The potential of weighting factors to improve the beamforming quality could also be investigated.

5. ACKNOWLEDGEMENT

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