

Unexpected Nonlinear Effects in Superconducting Transition-Edge Sensors

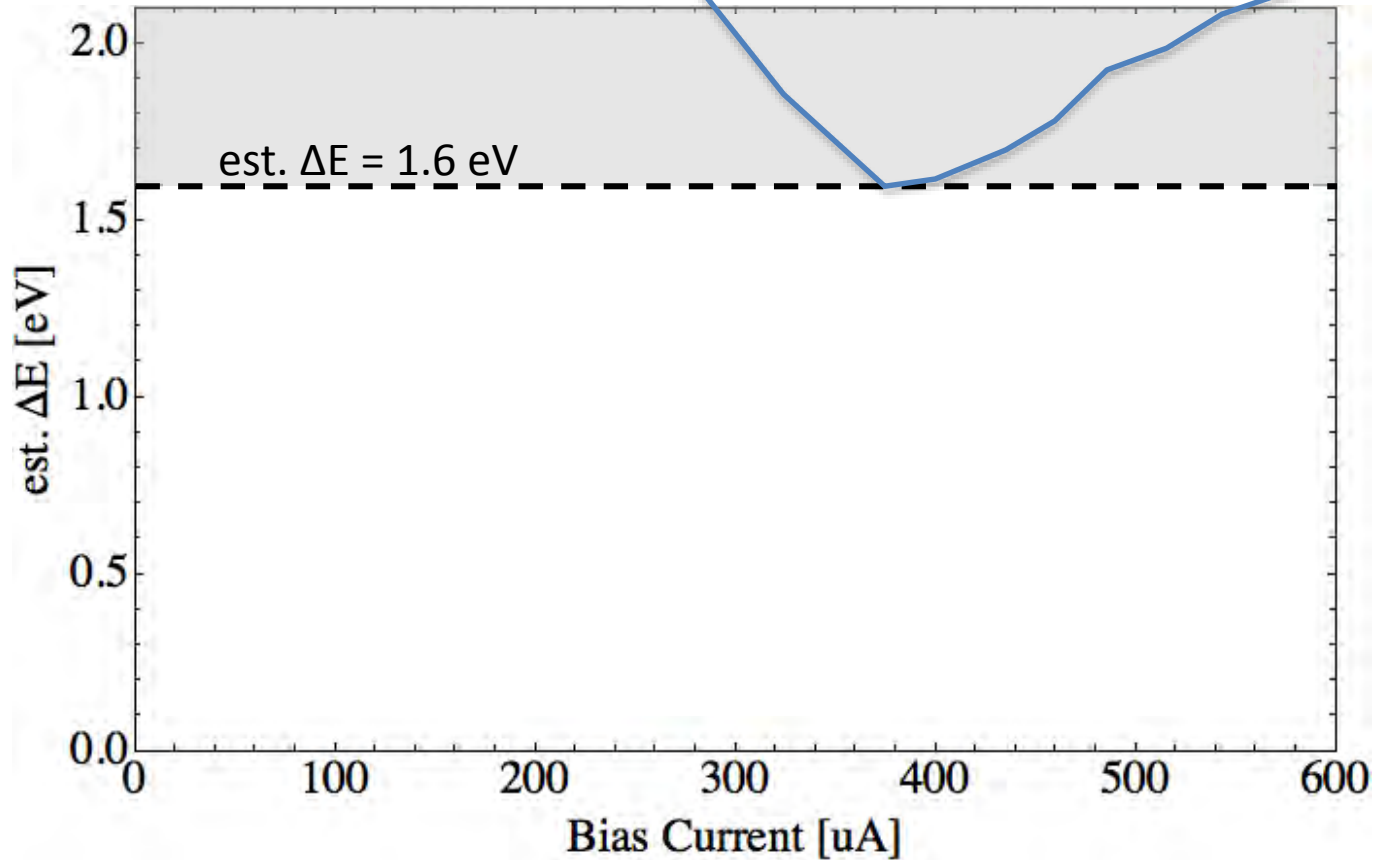
Jack Sadleir

Code 553

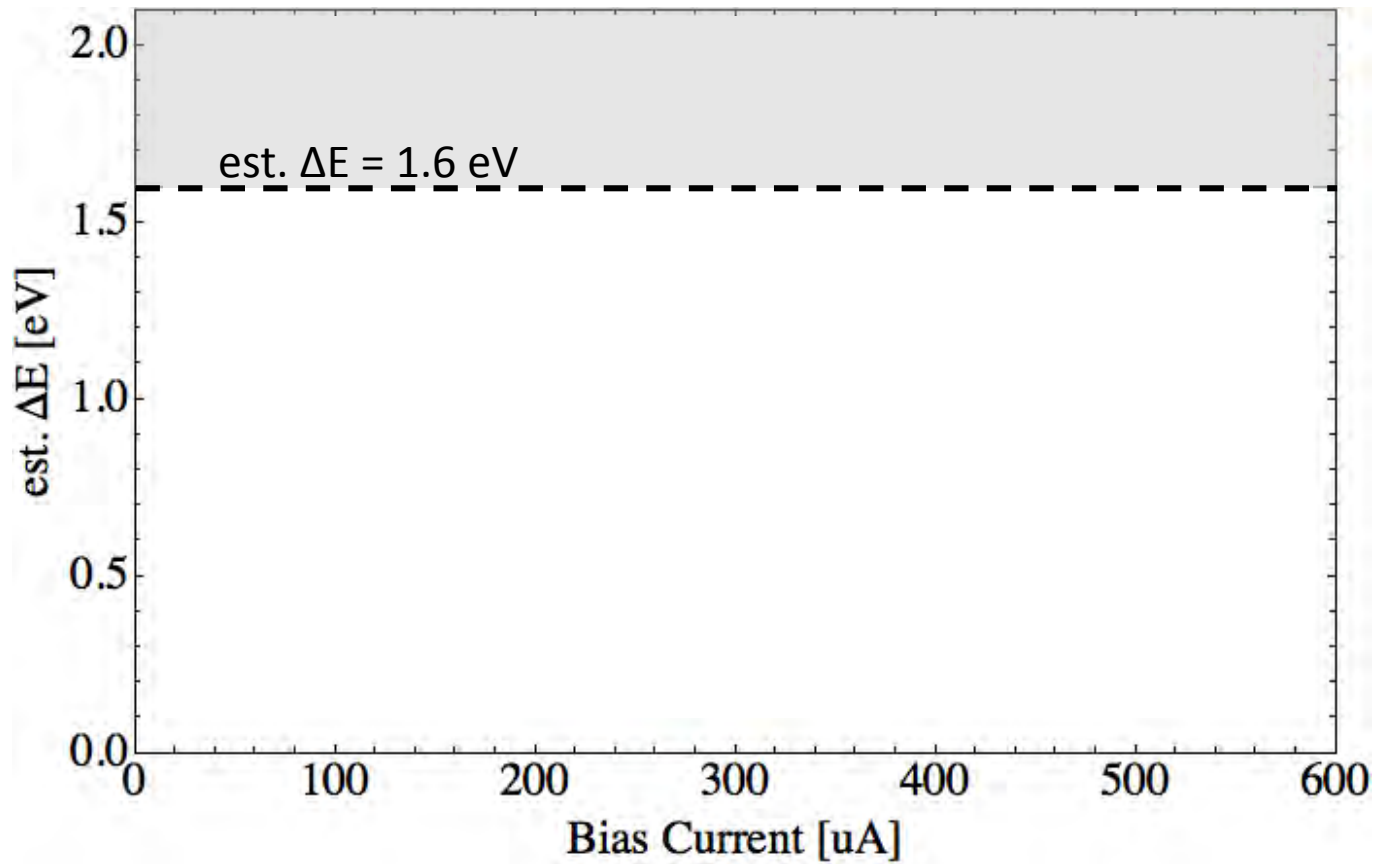
NASA Goddard Space Flight Center

6 keV

*World Record Energy
Resolution TES
Fabricated in the DDL
at NASA GSFC*

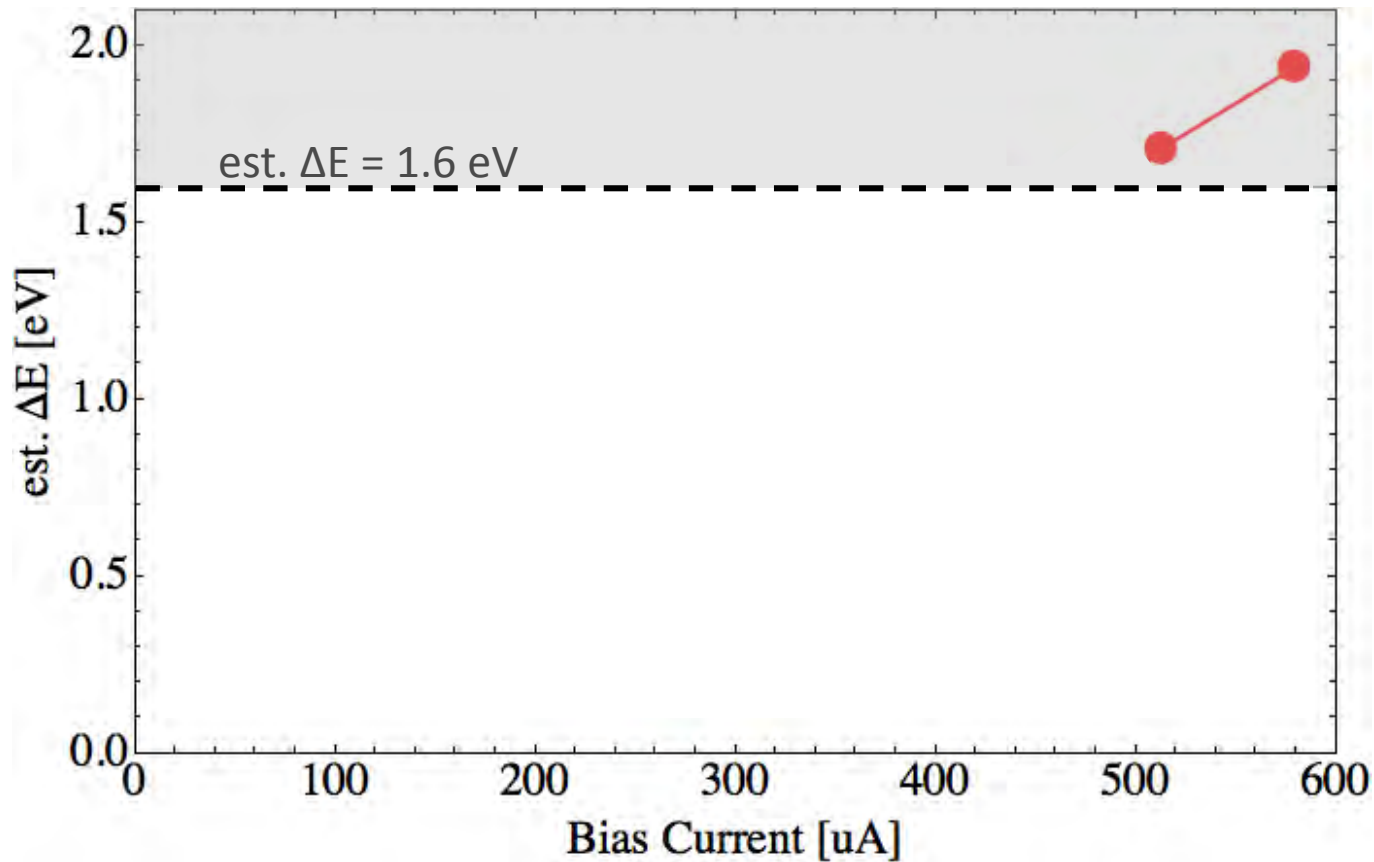


6 keV

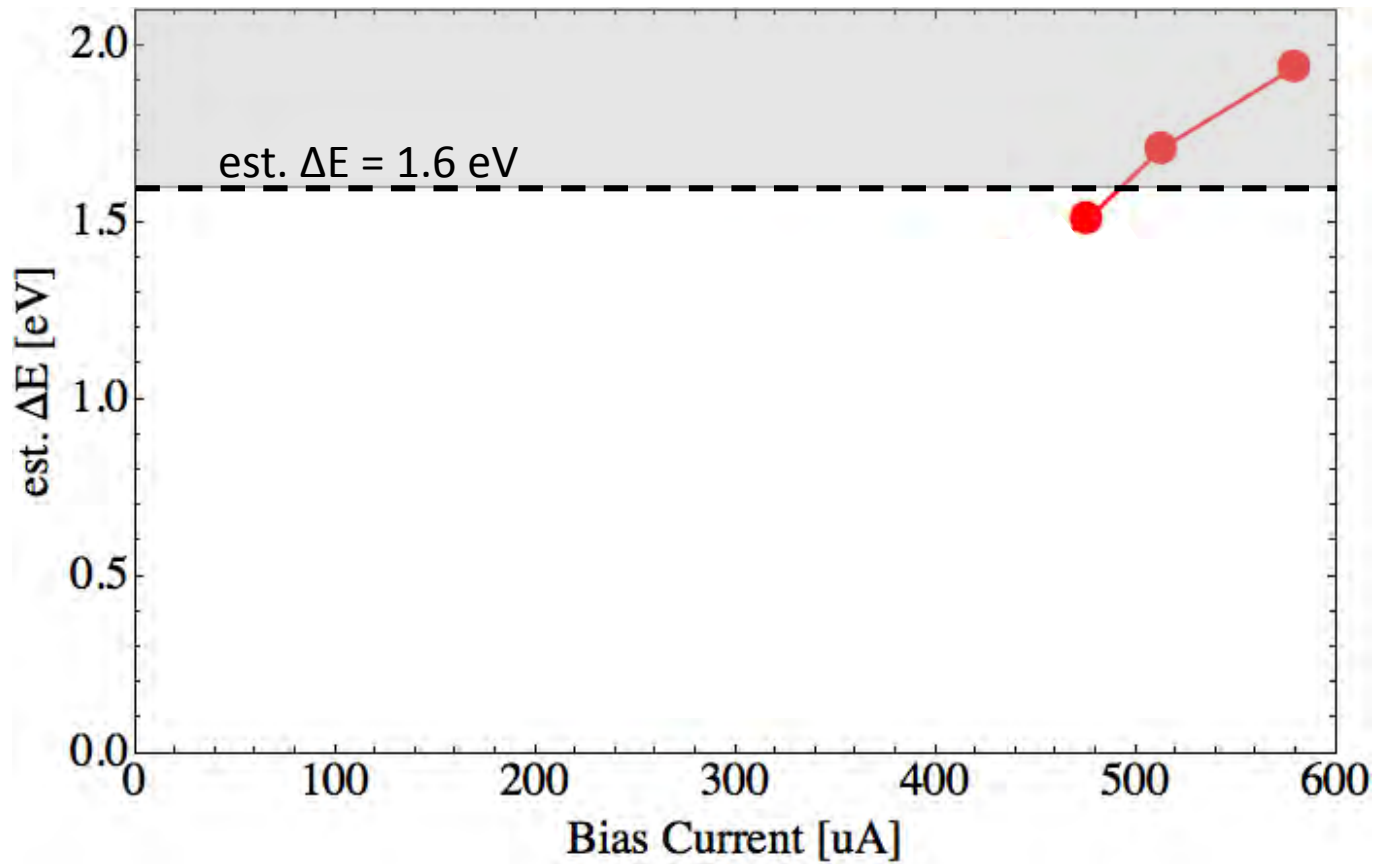


6 keV

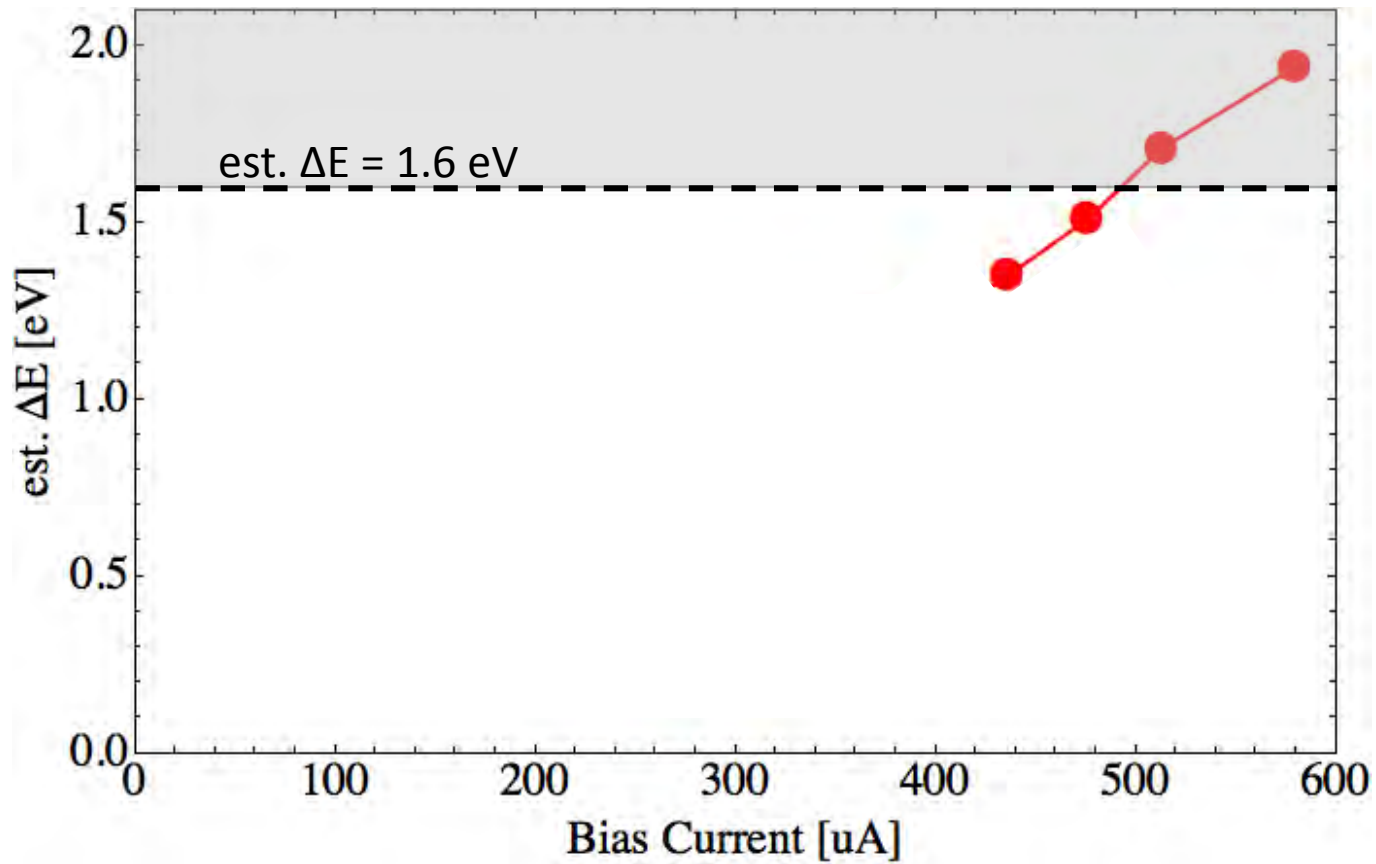
TES Fabricated in the
DDL at NASA GSFC
(But not specifically
designed for this
application)



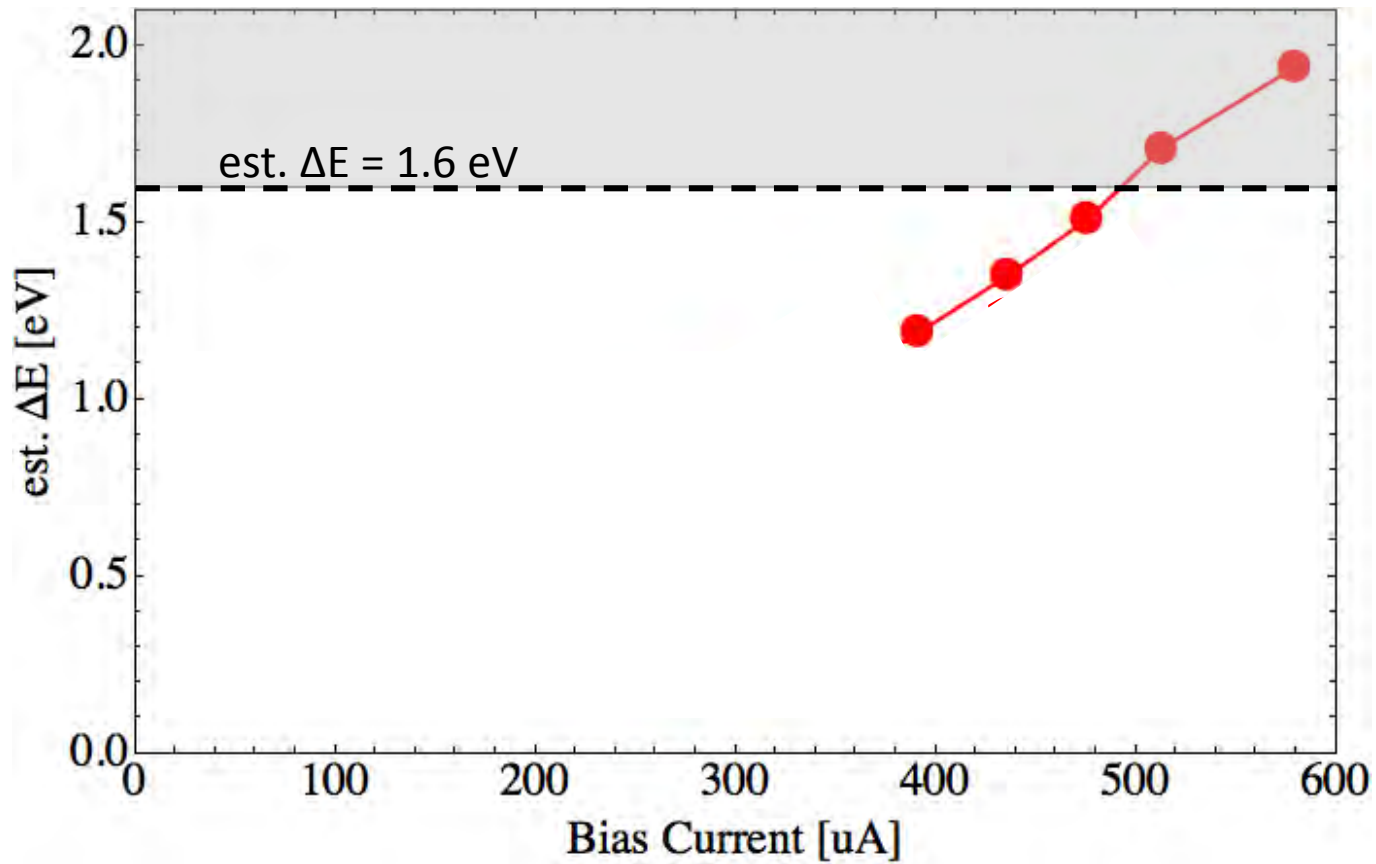
6 keV



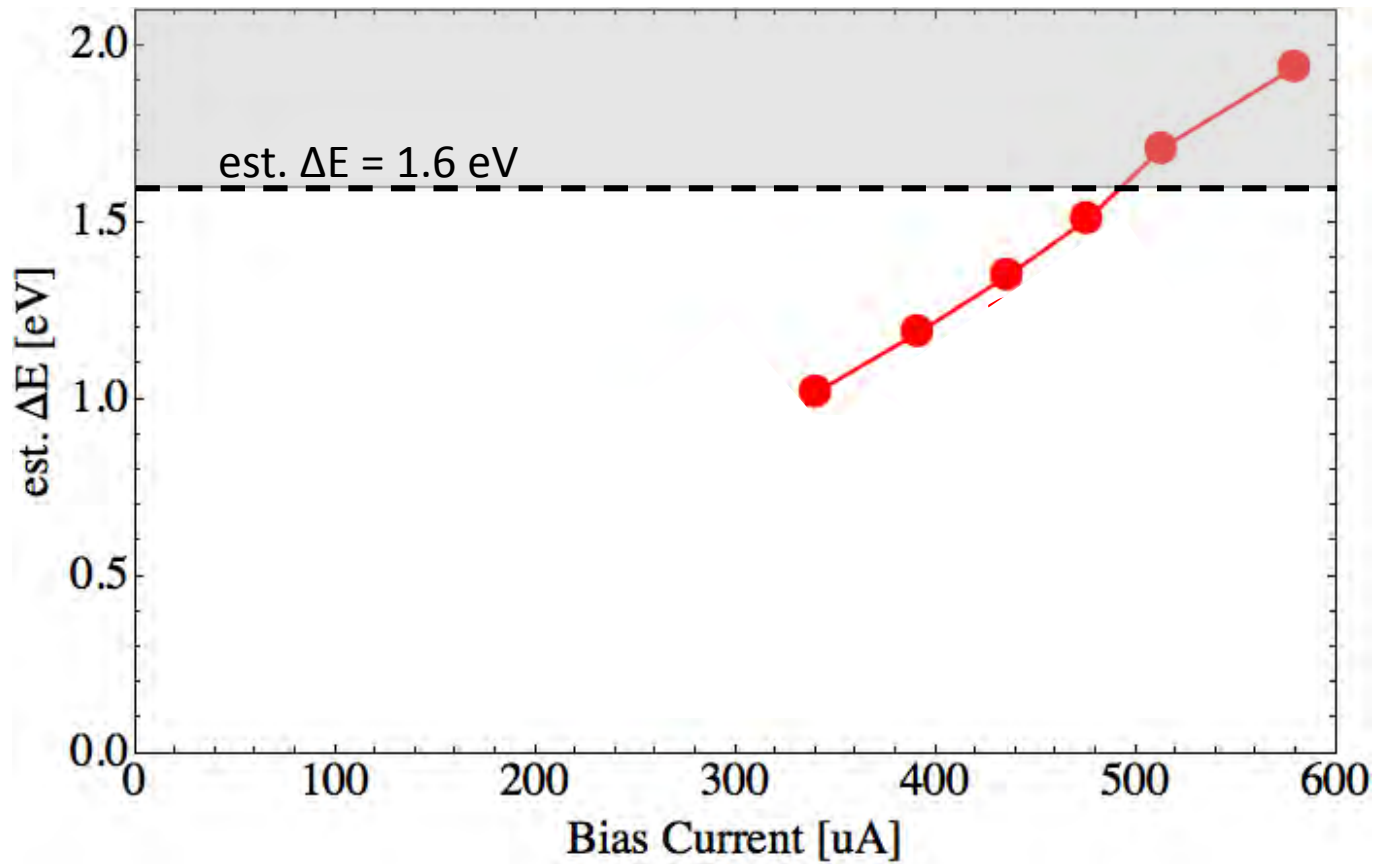
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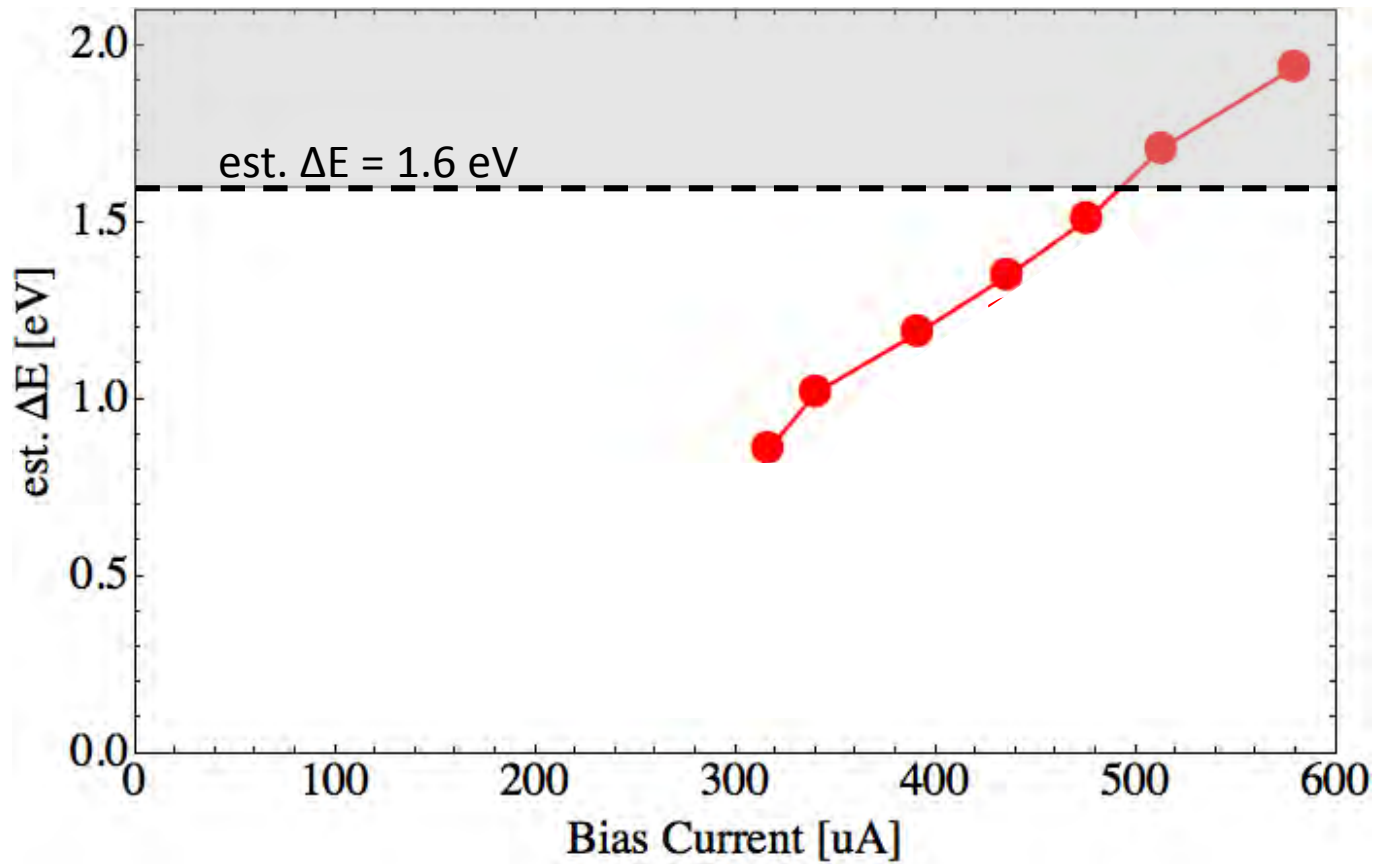
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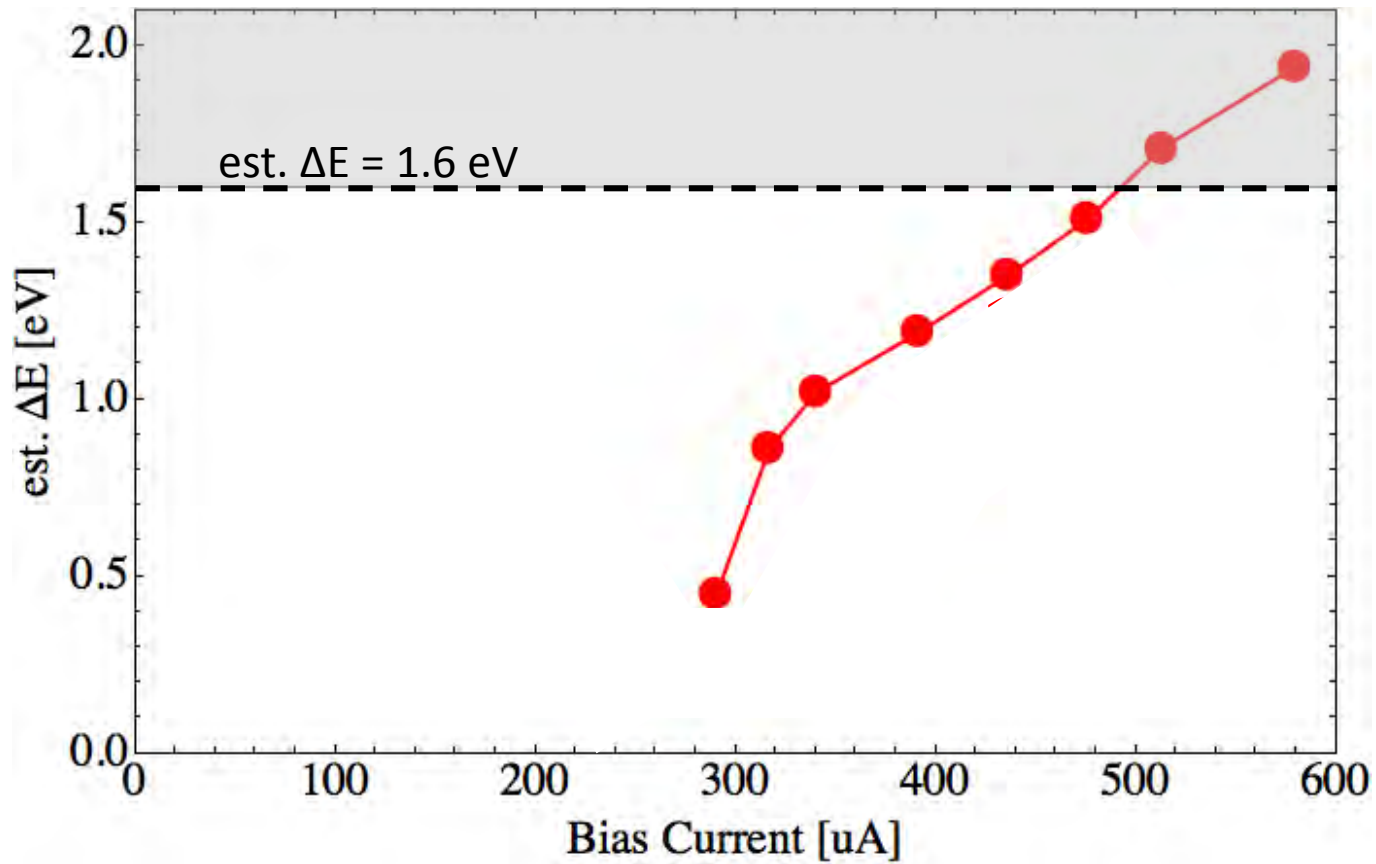
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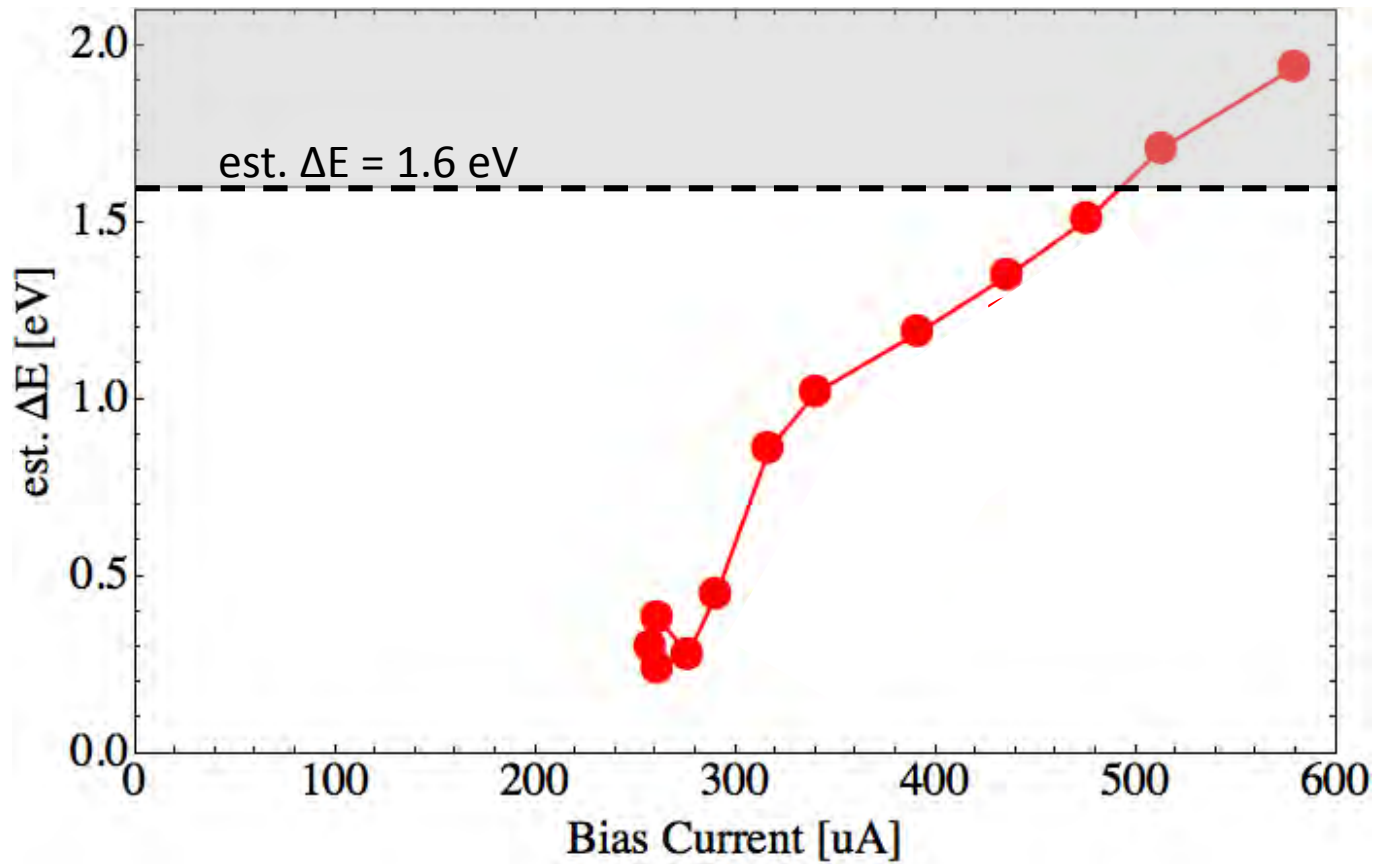
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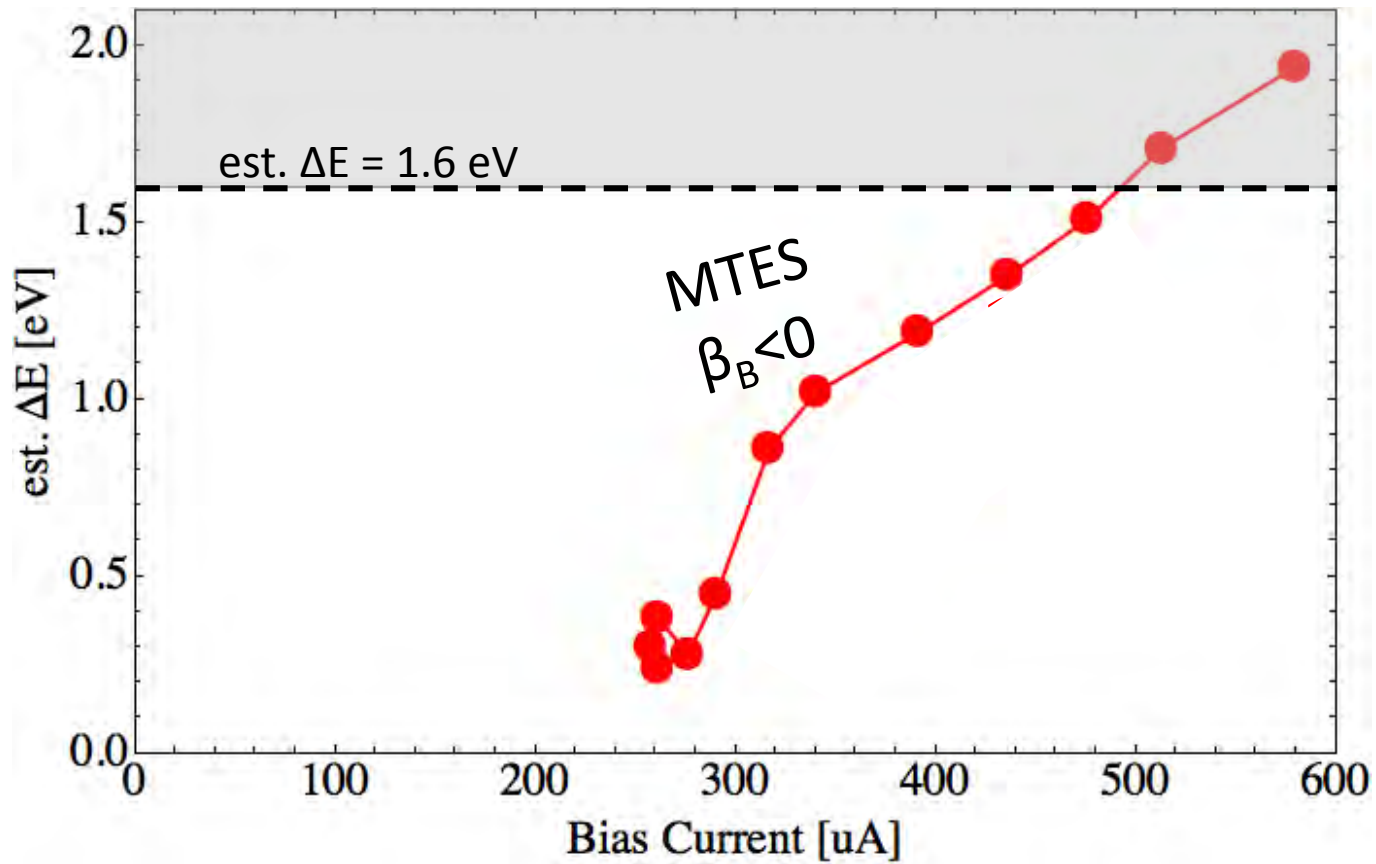
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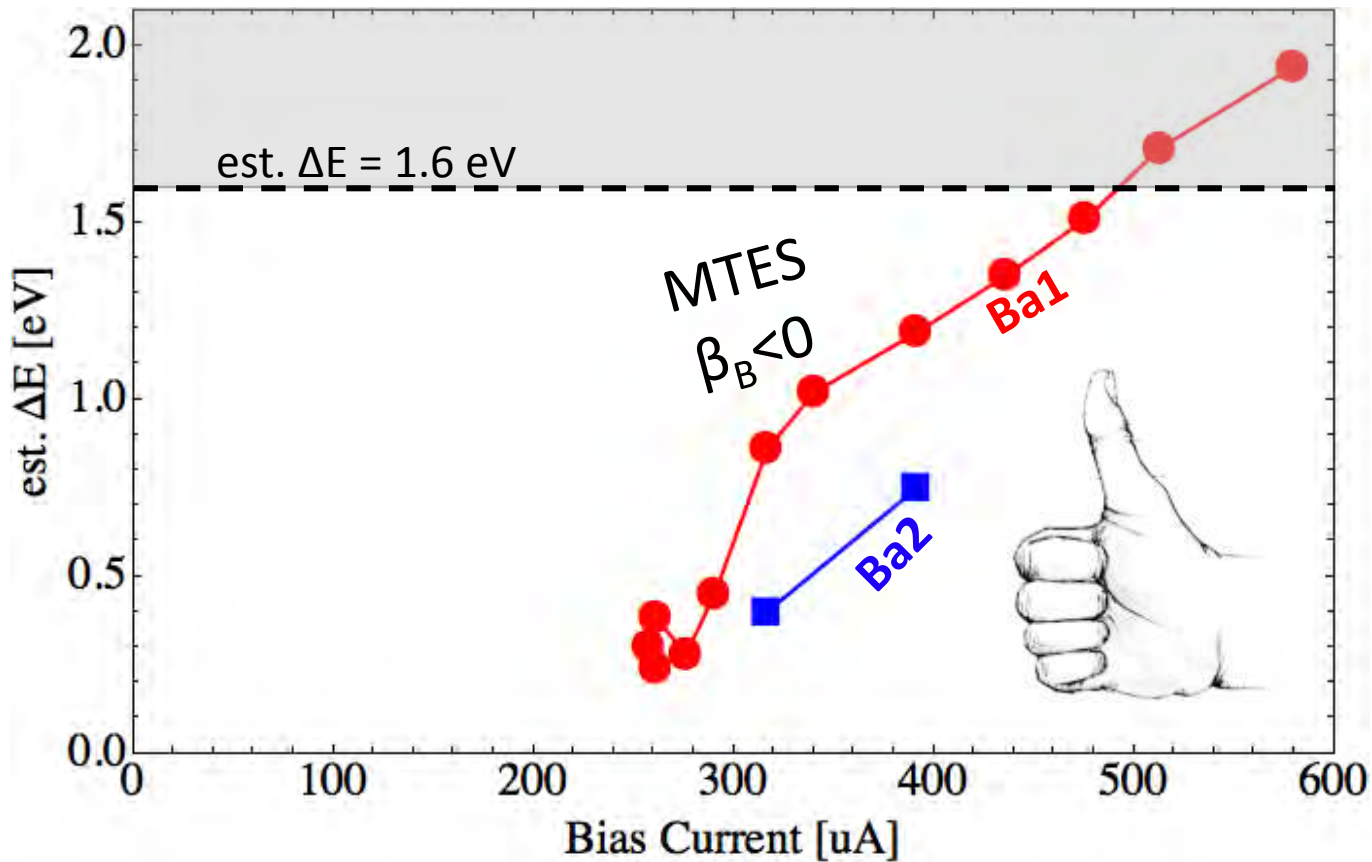


6 keV

CAUTION:

- *If a Linear detector...*
- *(Estimated)*

$$\text{est. } \Delta E = \text{true } \Delta E$$



What is “nonlinearity”? the “nonlinear-model” ... field “nonlinear dynamics”

Nonlinear \rightarrow Everything Else (*a very ∞ set ;-)*)

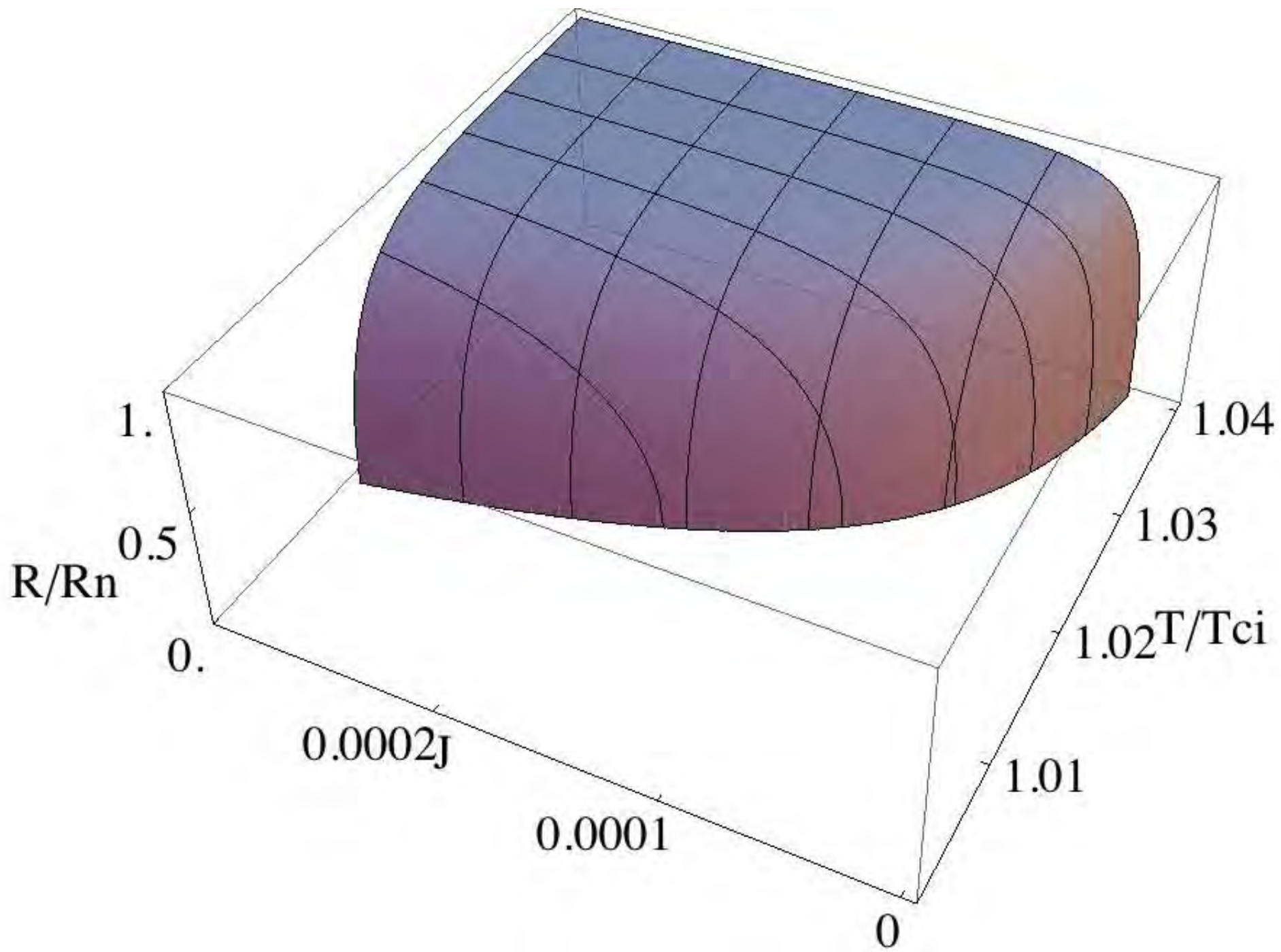
- **Nonlinearity ubiquitous** in nature:
 - e.g. 60 Hz harmonics pickup at 120 Hz, 180 Hz,.... requires nonlinearity in the system.
- **Linearity ubiquitous** in our mathematical description of nature:
 - often a good approximation to real physical systems.
 - It is mathematically easier
 - Linear Tools:
 - Superposition Principle
 - Transform methods, transfer functions etc.

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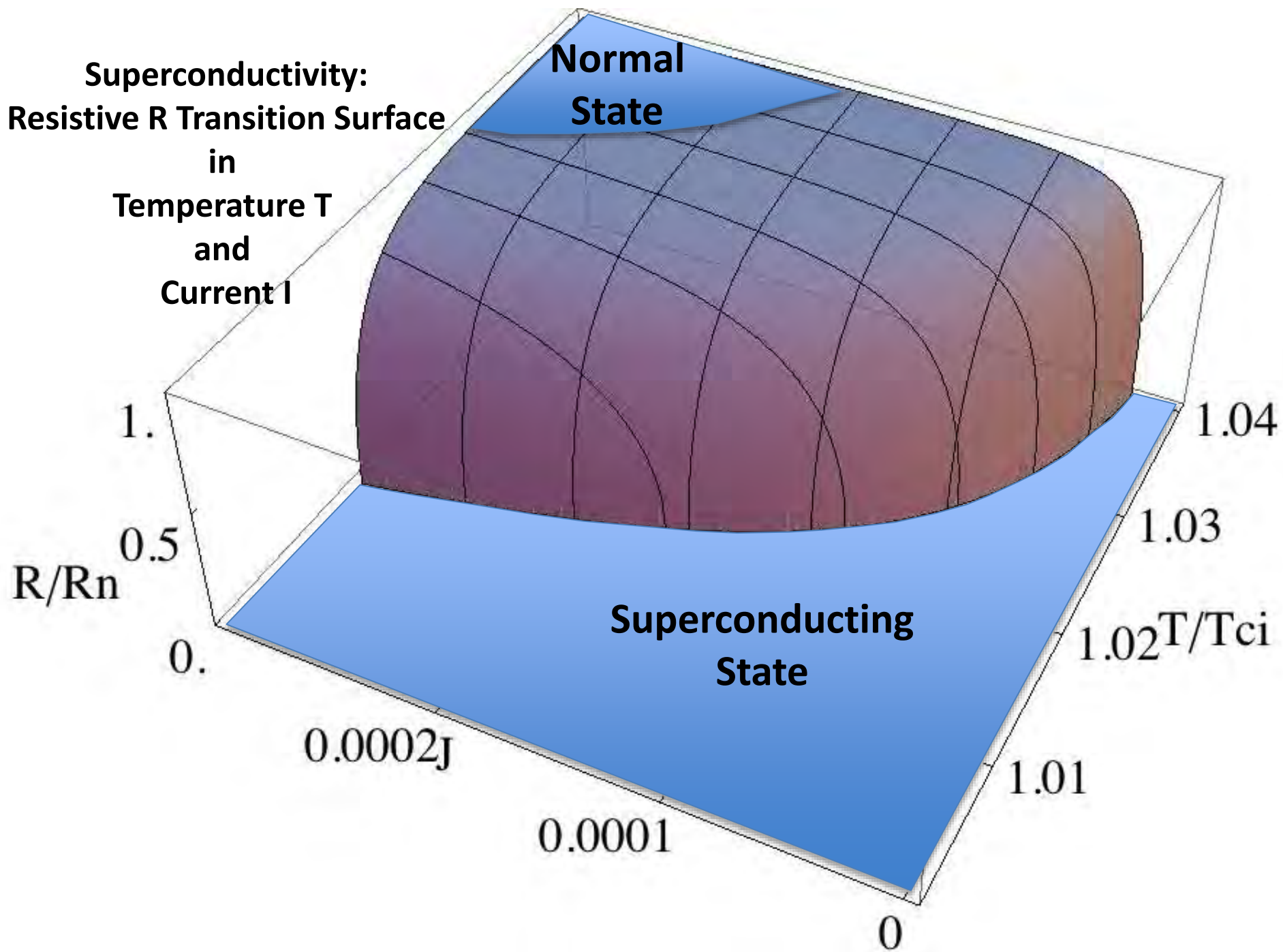
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*Known or
Expected
Nonlinearities*

- If our thermistor TES sensor...
 - If $R(T) \rightarrow R(T,J)$.
 - Then **Nonlinear** system of Diff Eqs.
 - (Can approximate by a linearized system of Diff Eqs.)
- “**Nonlinear**” if
 - $R(T) \rightarrow$ anything other than a single straight line
 - $R(T,J) \rightarrow$ “ “ “ a single plane.
 - $R(T,J,B) \rightarrow$ “ “ “ a single hyperplane.



**Superconductivity:
Resistive R Transition Surface
in
Temperature T
and
Current I**



**Normal
State**

**Superconducting
State**

R/R_n
1.
0.5
0.

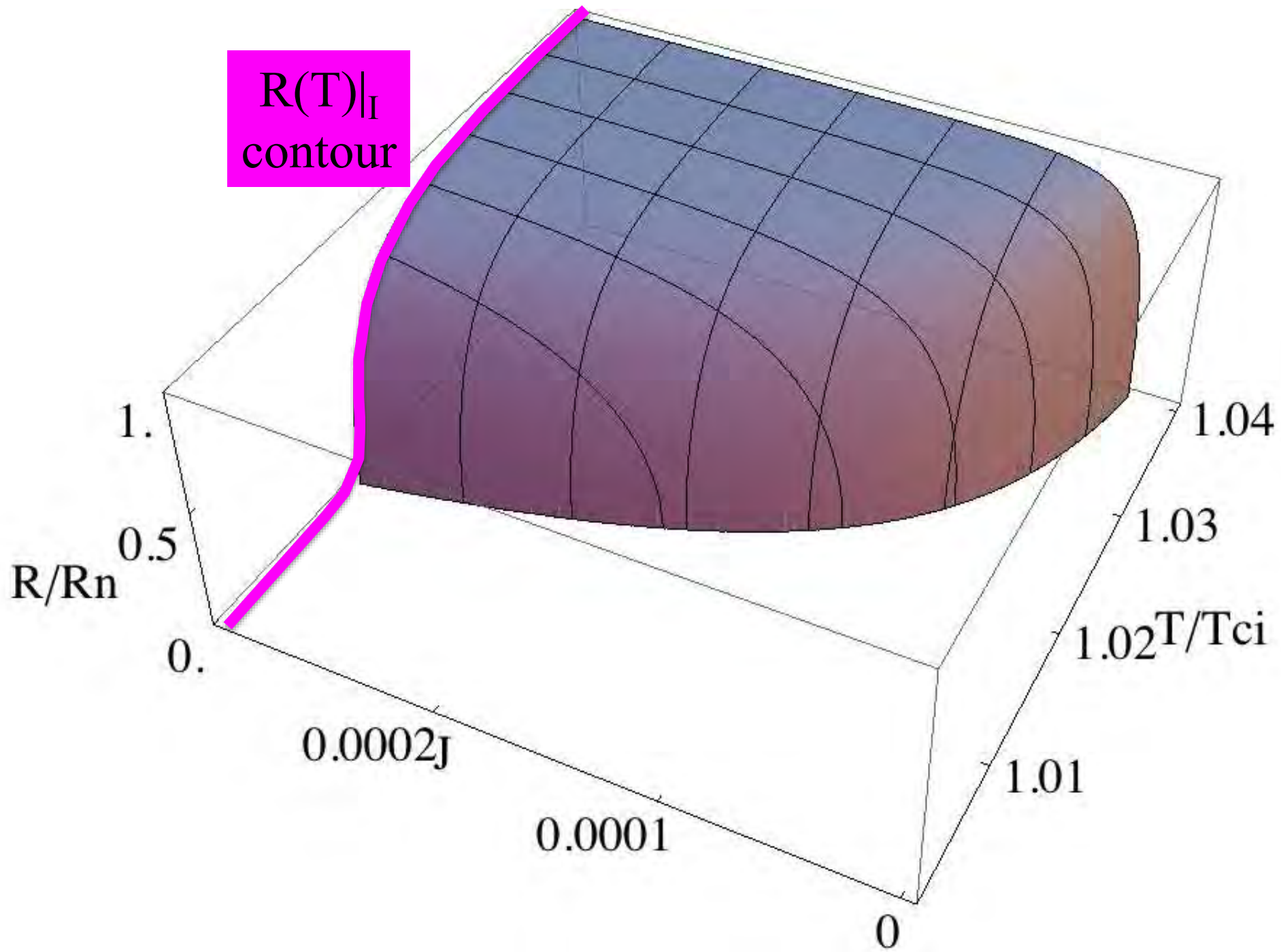
0.0002J

0.0001

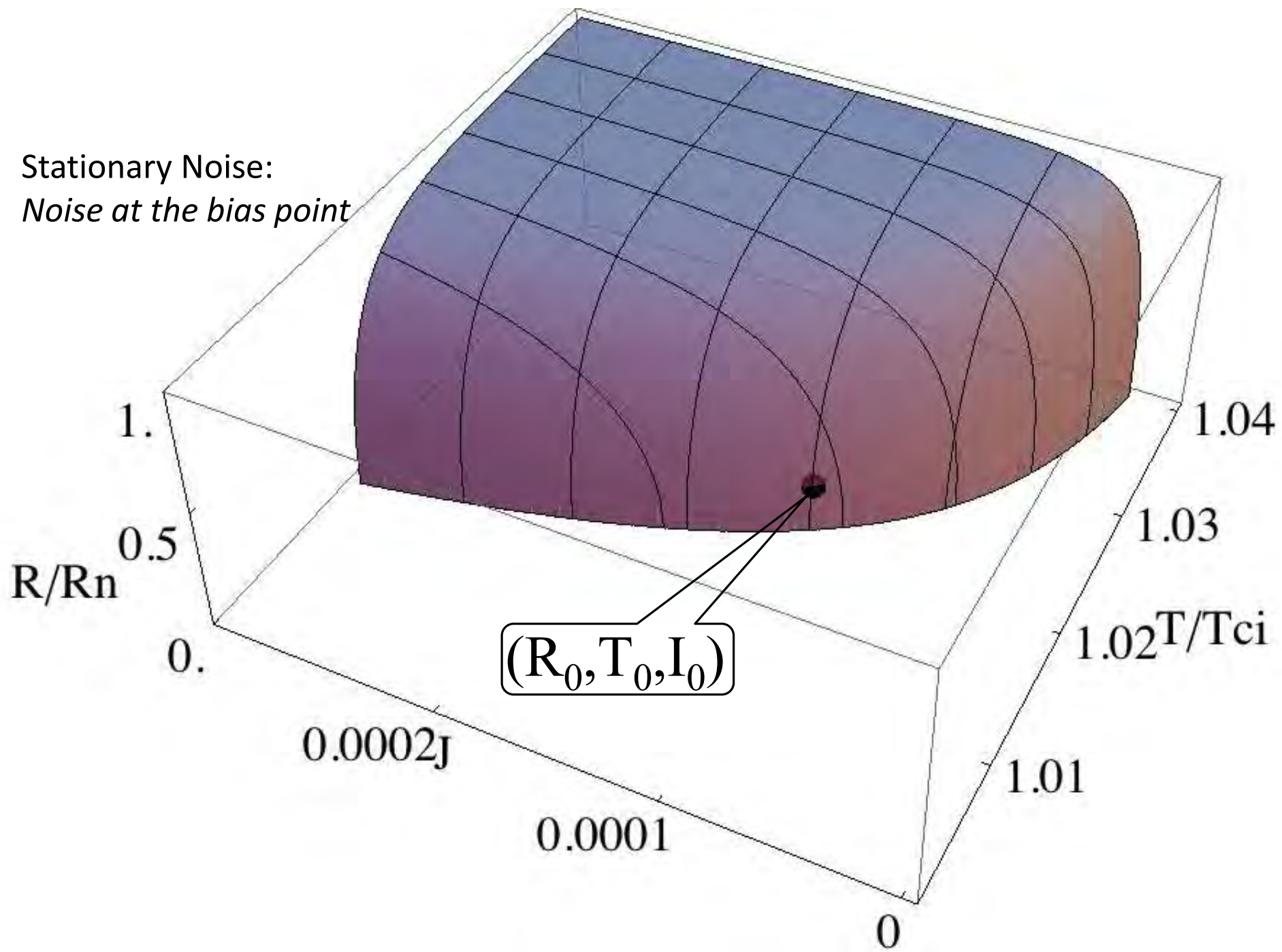
0

1.04
1.03
1.02
 T/T_{ci}
1.01

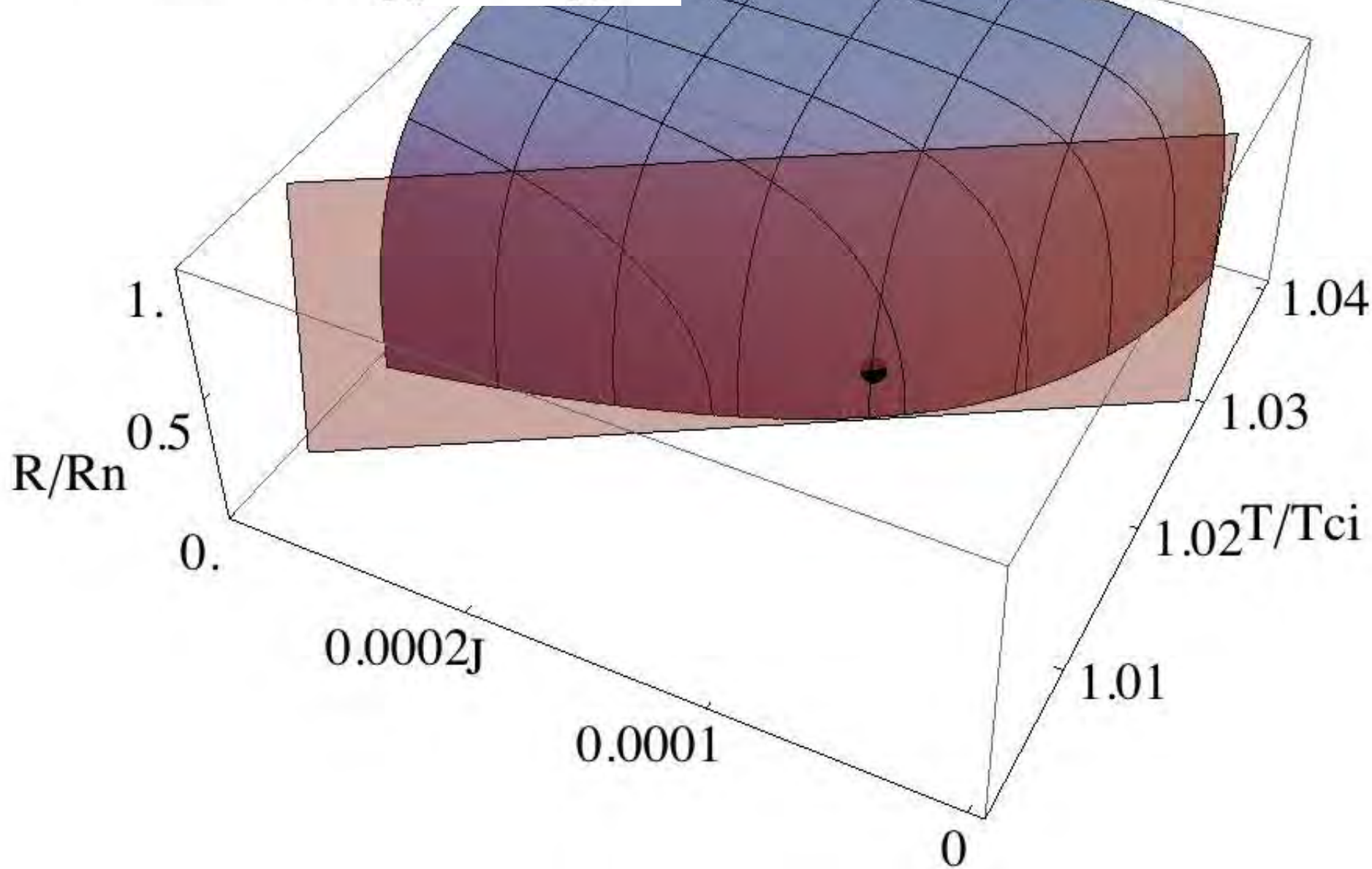
$R(T)|_I$
contour



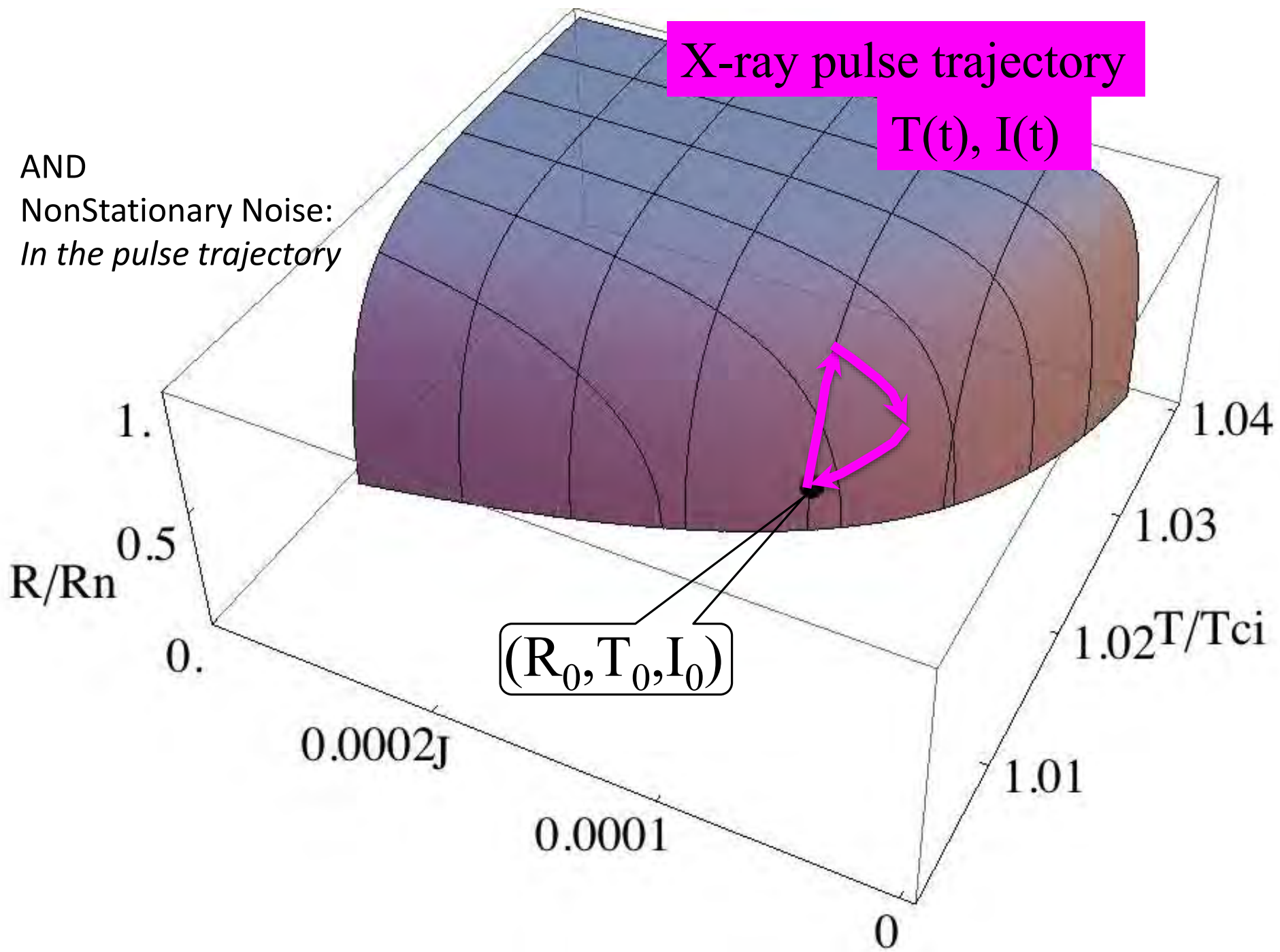
Stationary Noise:
Noise at the bias point



$$R(T, I) \approx R_0 + \frac{\partial R}{\partial T} \delta T + \frac{\partial R}{\partial I} \delta I$$



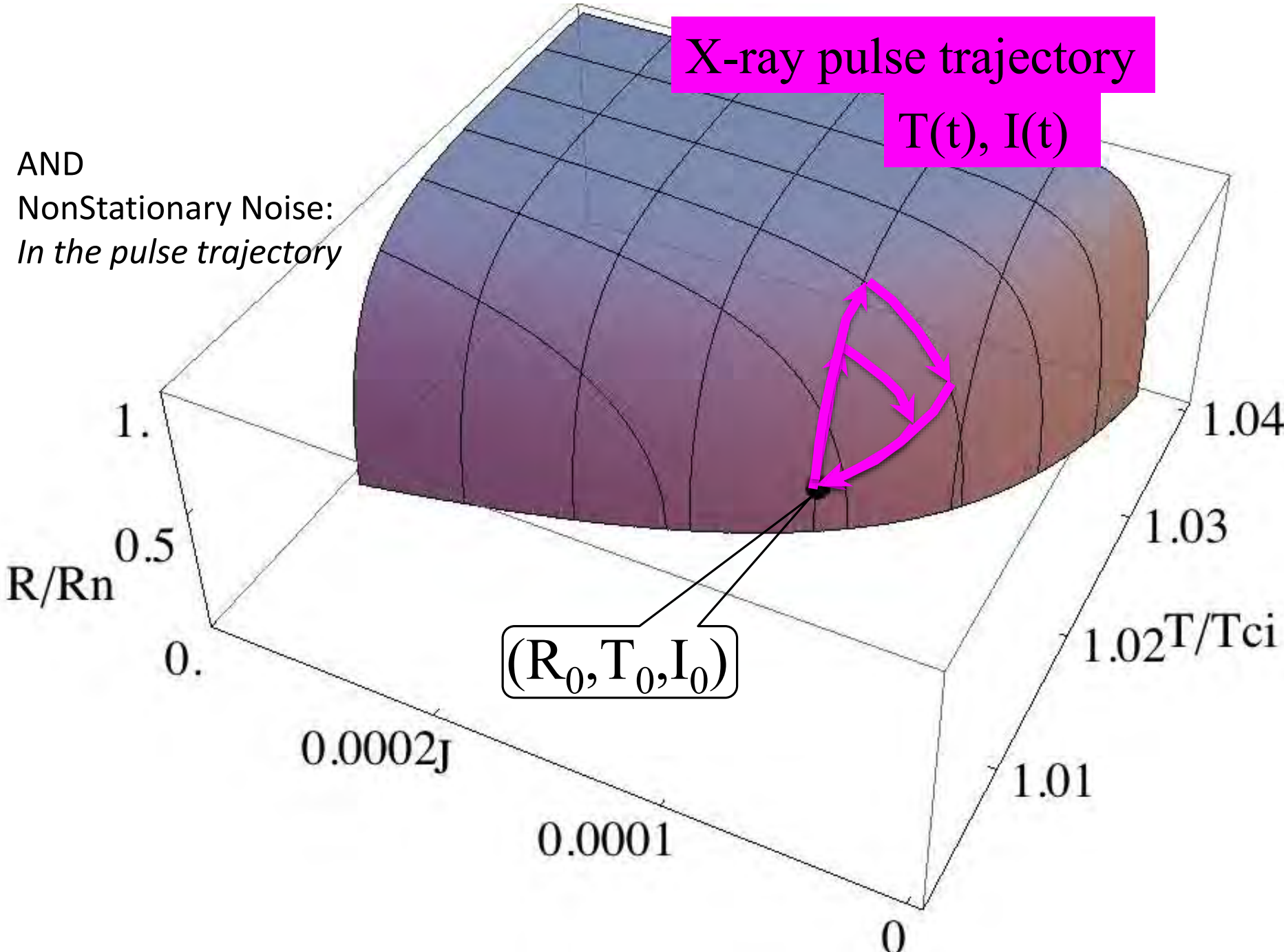
AND
NonStationary Noise:
In the pulse trajectory



X-ray pulse trajectory

$T(t), I(t)$

AND
NonStationary Noise:
In the pulse trajectory



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*“Unexpected
Nonlinearities”*

In steady state with standard DC biased TES operation.
The TES current is wildly oscillating in time at high frequencies and we only measure (*and are only aware of*) the time averaged current in the TES.

Conclusions: ...*Paradigm Shifting*

- In steady state with standard DC biased TES operation. The TES current is wildly oscillating in time at high frequencies and we only measure (*and are only aware of*) the time averaged current in the TES.
- The equations governing this time response is **nonlinear** time-dependent diff eqs.
 - **KEY POINT:** This time dependent current exists when the TES is only DC biased, NO time dependent inputs what so ever. It comes about from the intrinsic physics governing the superconducting state of the TES.

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- The measured resistive transition surface is not simply a function of $R(T,J,B)$. But is also a function of the electric circuit. $R(T,J,B,L,R_{sh})$.
 - → Problem compartmentalizing: → *Same TES measured in slightly different circuits will appear to have a different resistive transition surface. TEST in real setup early!!!*

Conclusions: ...*Paradigm Shifting*

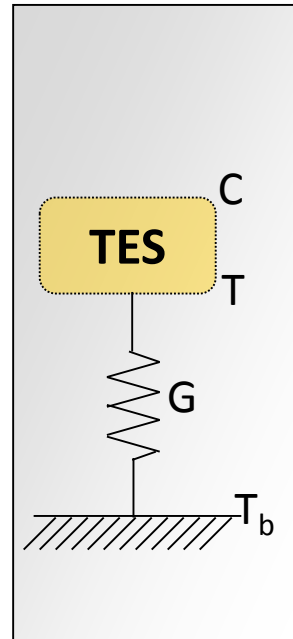
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- Circuit and TES parameters can have values such that these time dependent solutions become multivalued in time.
 - *What does that mean?* Can mean excess noise or fine structure of the time averaged resistive transition surface.

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- Circuit and TES parameters can have values such that these time dependent solutions become multivalued in time.
 - *What does that mean?* Can mean excess noise or fine structure of the time averaged resistive transition surface.
- The biased TES waiting to detect a photon (μ -calorimeter) or flux of photons (μ -bolometers) can itself act as a radiation source.
 - Tricky when an array of exquisitely sensitive micro-wave radiation detectors are themselves sources of microwave radiation.
 - This can lead to radiation resistance and fine structure in the resistive transition.
 - Possible cross talk between pixels in an array.
 - The time dependence of the current can take on pure sin waves and also very nonsinusoidal forms (variable harmonic content)
 - The fundamental frequency of these oscillations changes with bias voltage.
 - Makes the prospects of FDM TES arrays with an AC-biased TES challenging. Structure, sensitivity, and cross-talk.

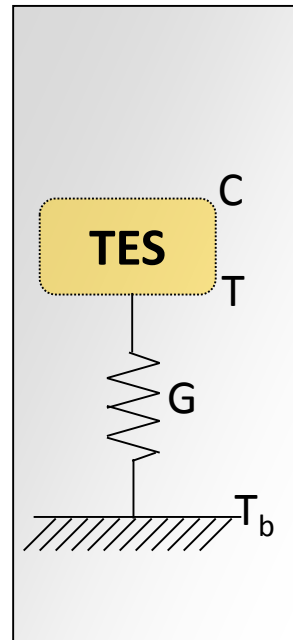
DONE

Thermal Circuit

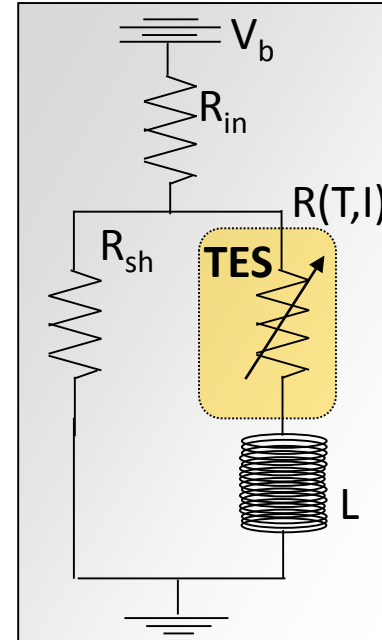


ELECTROTHERMAL MODEL

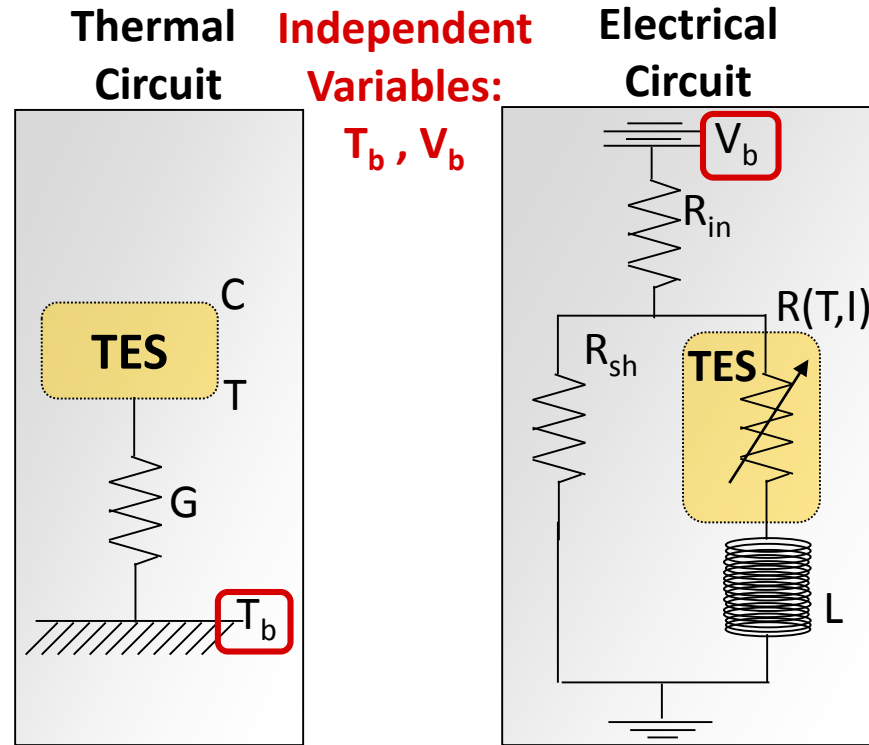
**Thermal
Circuit**



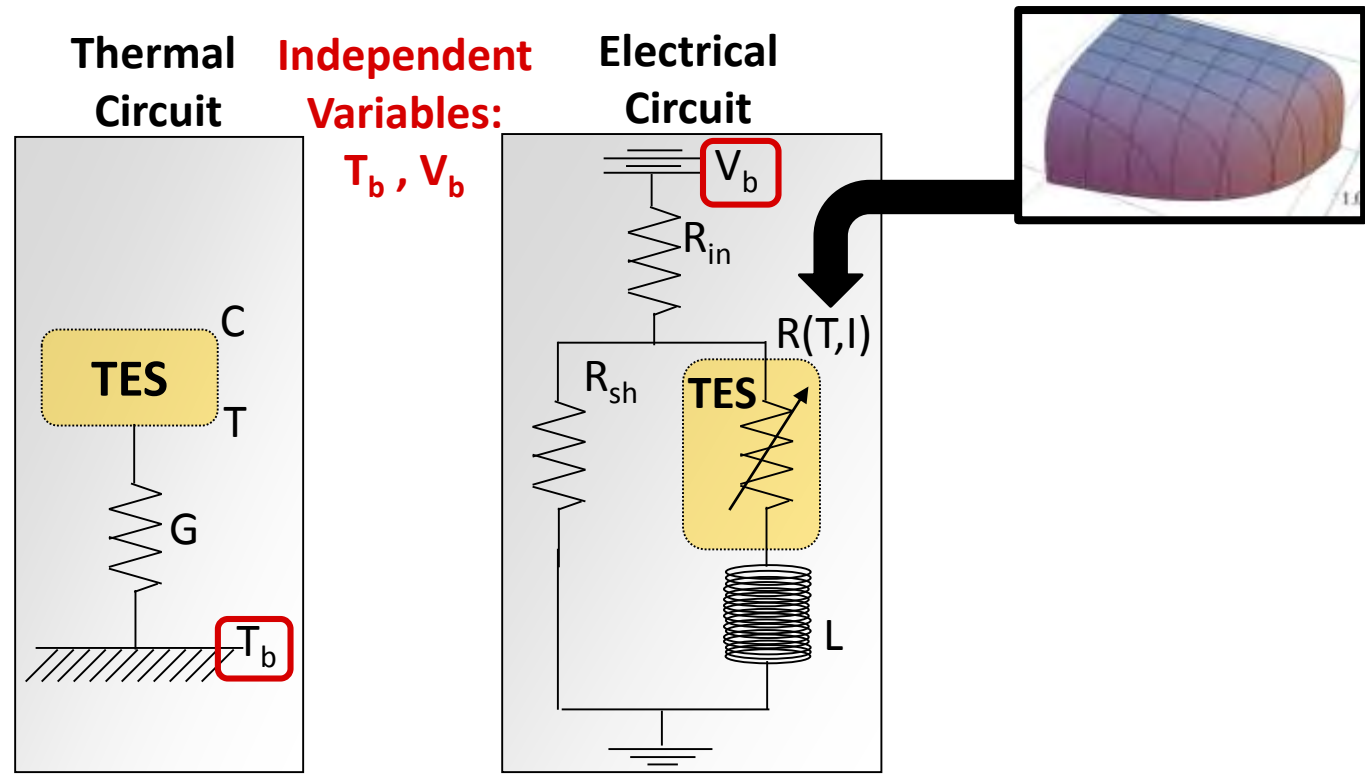
**Electrical
Circuit**



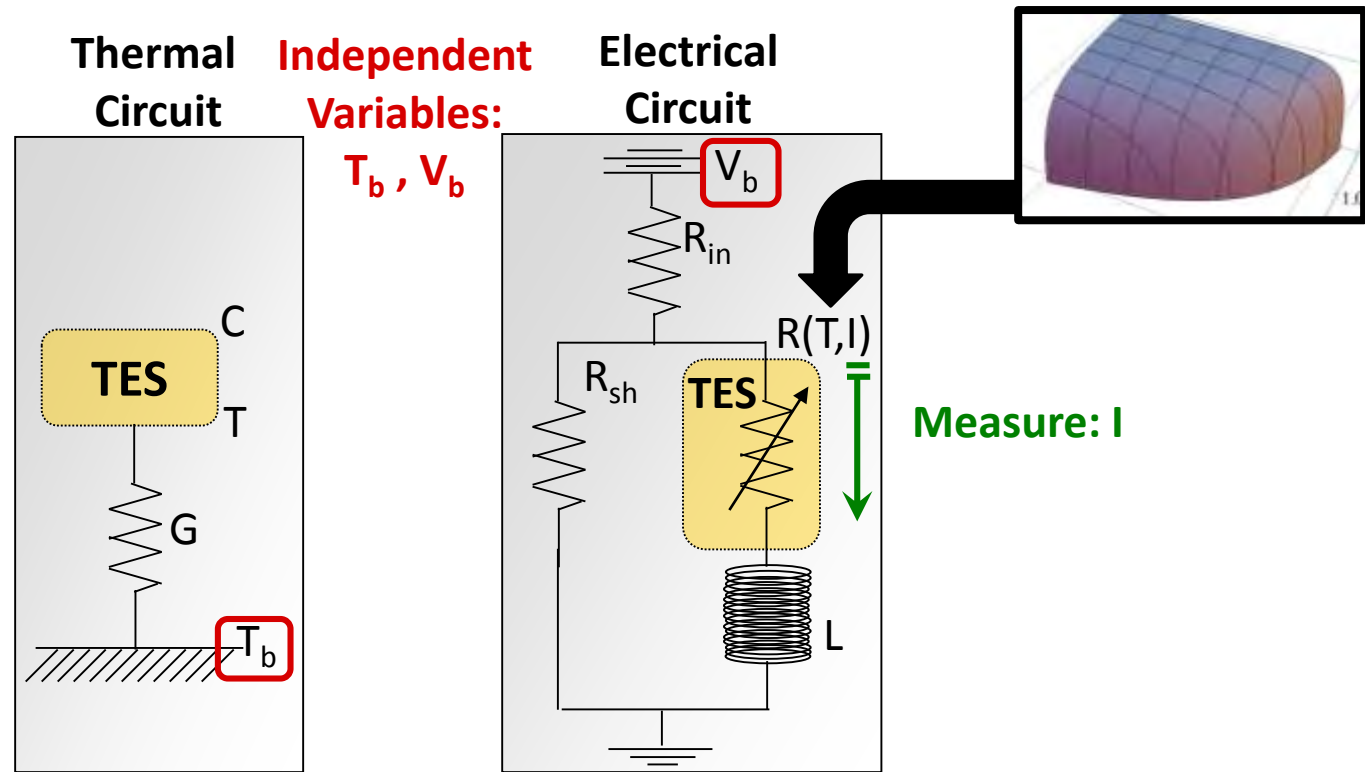
ELECTROTHERMAL MODEL



ELECTROTHERMAL MODEL

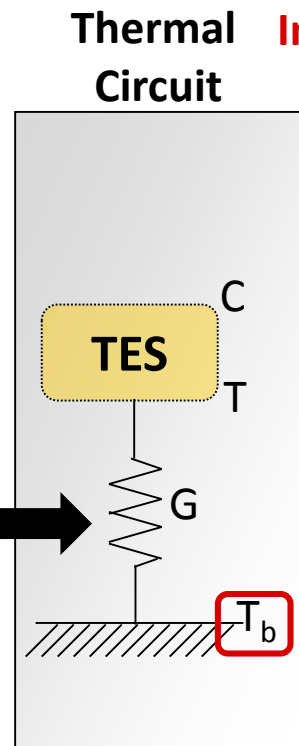


ELECTROTHERMAL MODEL

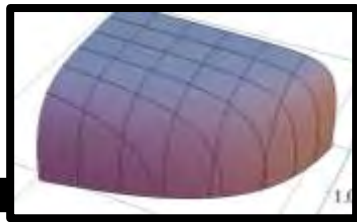
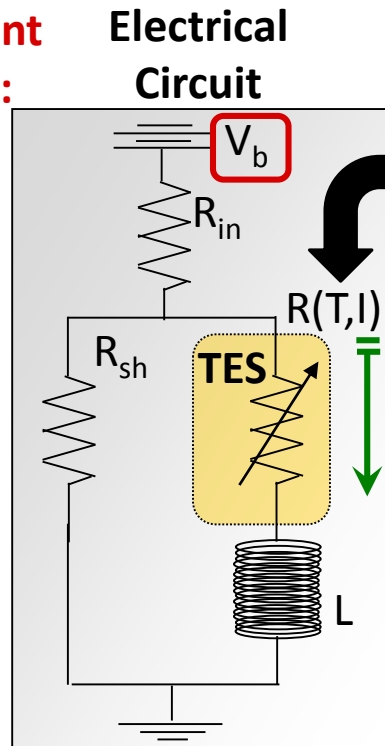


ELECTROTHERMAL MODEL

- TES on a: ...
- (1) solid substrate
 - (2) thin isolated membrane
 - (3) thin perforated membrane
 - (4) island suspended with long thin legs



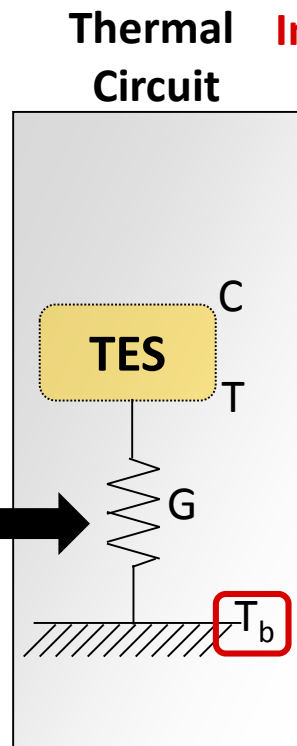
Independent Variables:
 T_b, V_b



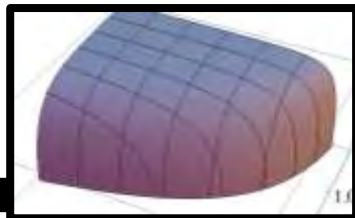
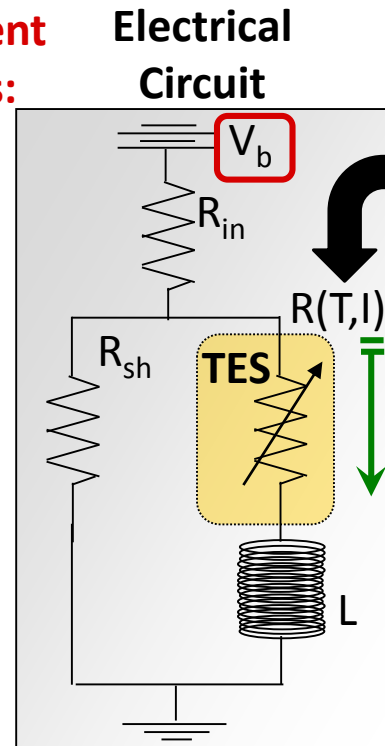
Measure: I

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Measure: I

Thermal power balance

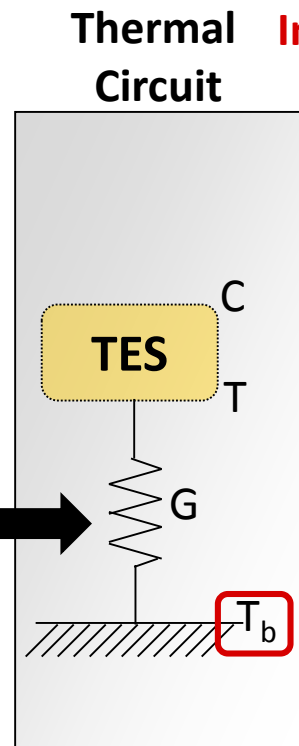
$$I^2 R = \frac{G}{nT^{(n-1)}} (T^n - T_b^n)$$

Electrical steady state

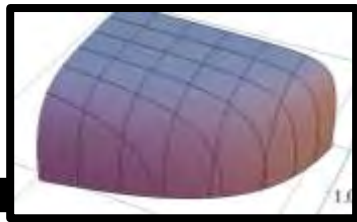
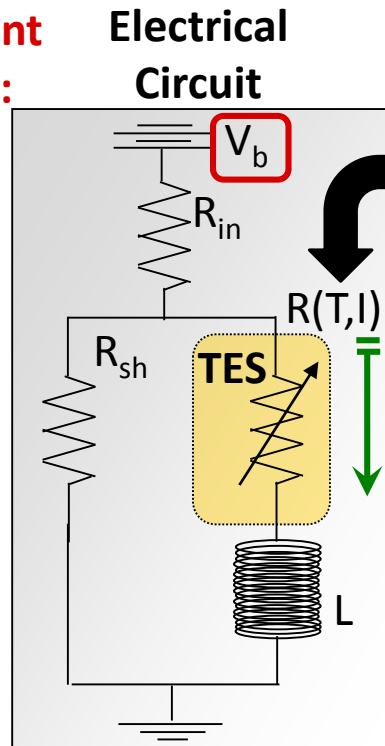
$$I = \frac{V_b}{R_{in}} \frac{R_{sh}}{R + R_{sh}}$$

IN STEADY STATE

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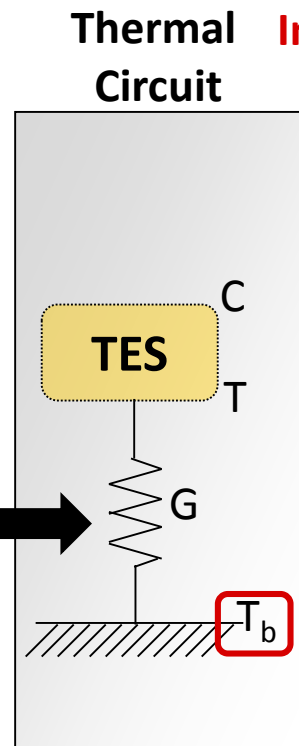
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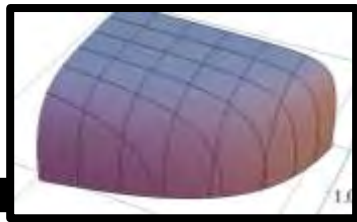
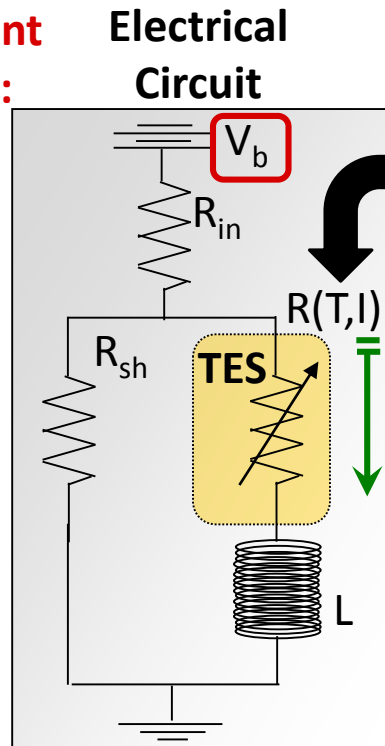
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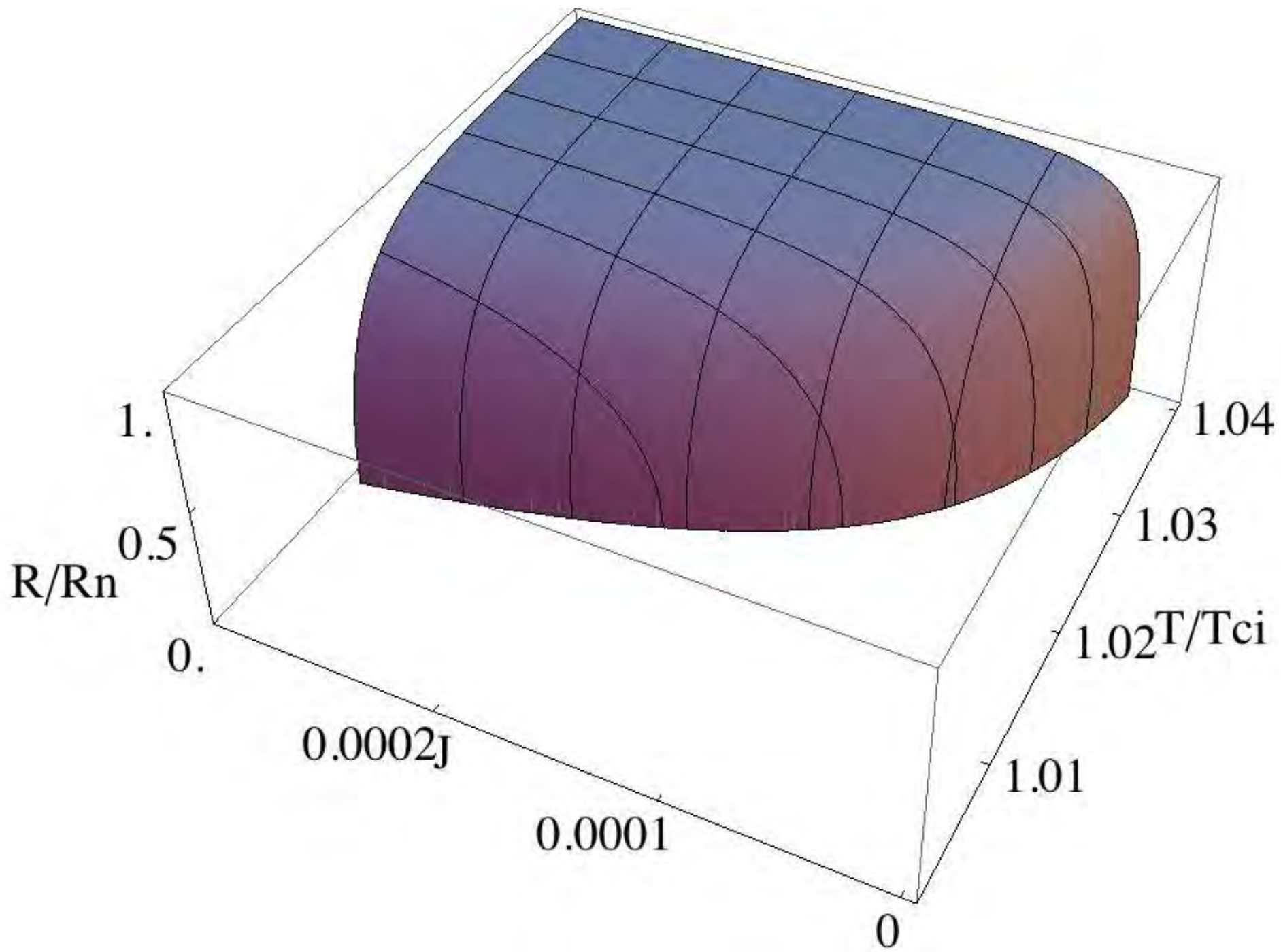
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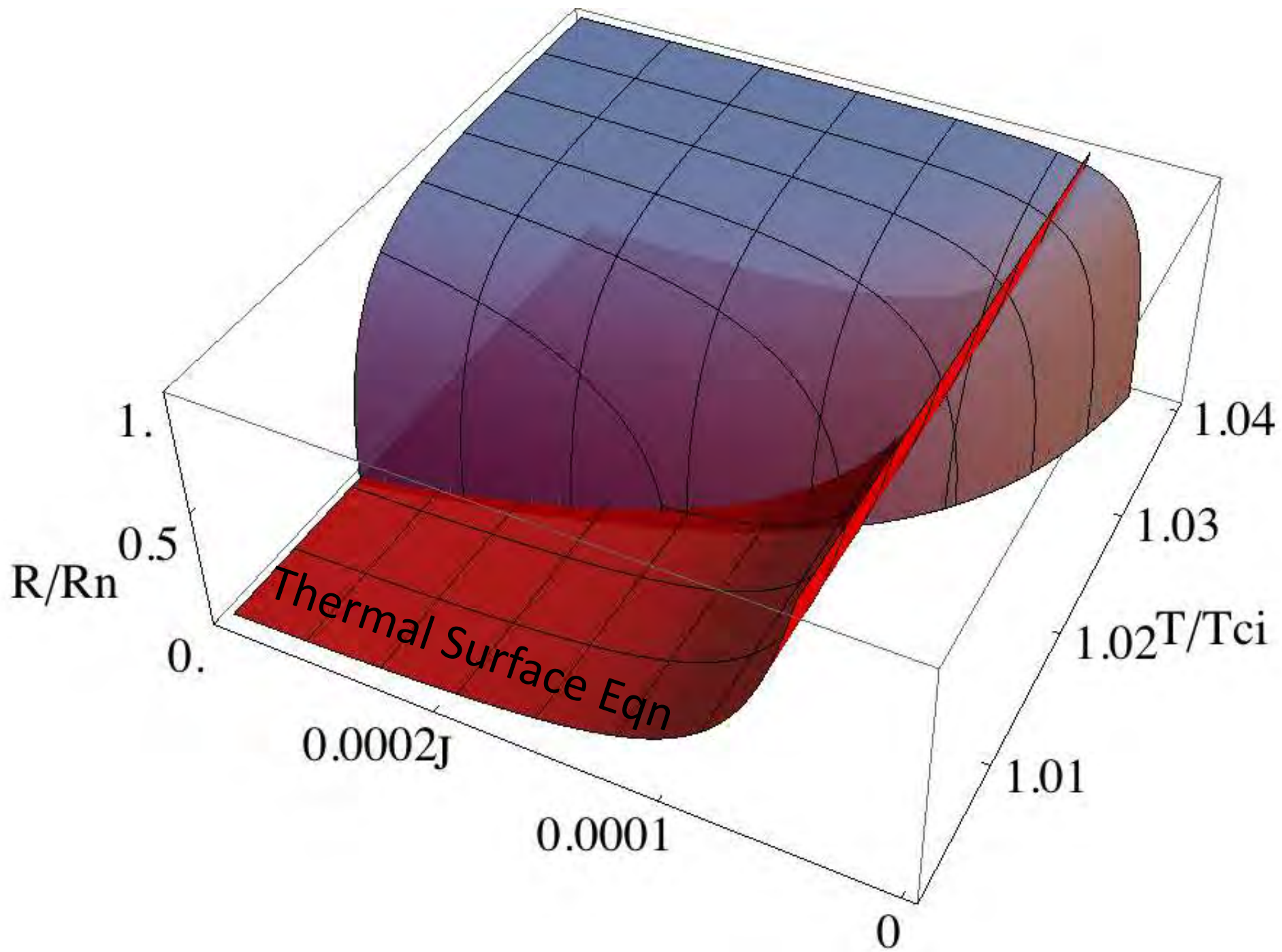
$$R = \frac{G}{I^2 n T^{(n-1)}} (T^n - T_b^n)$$

Electrical steady state

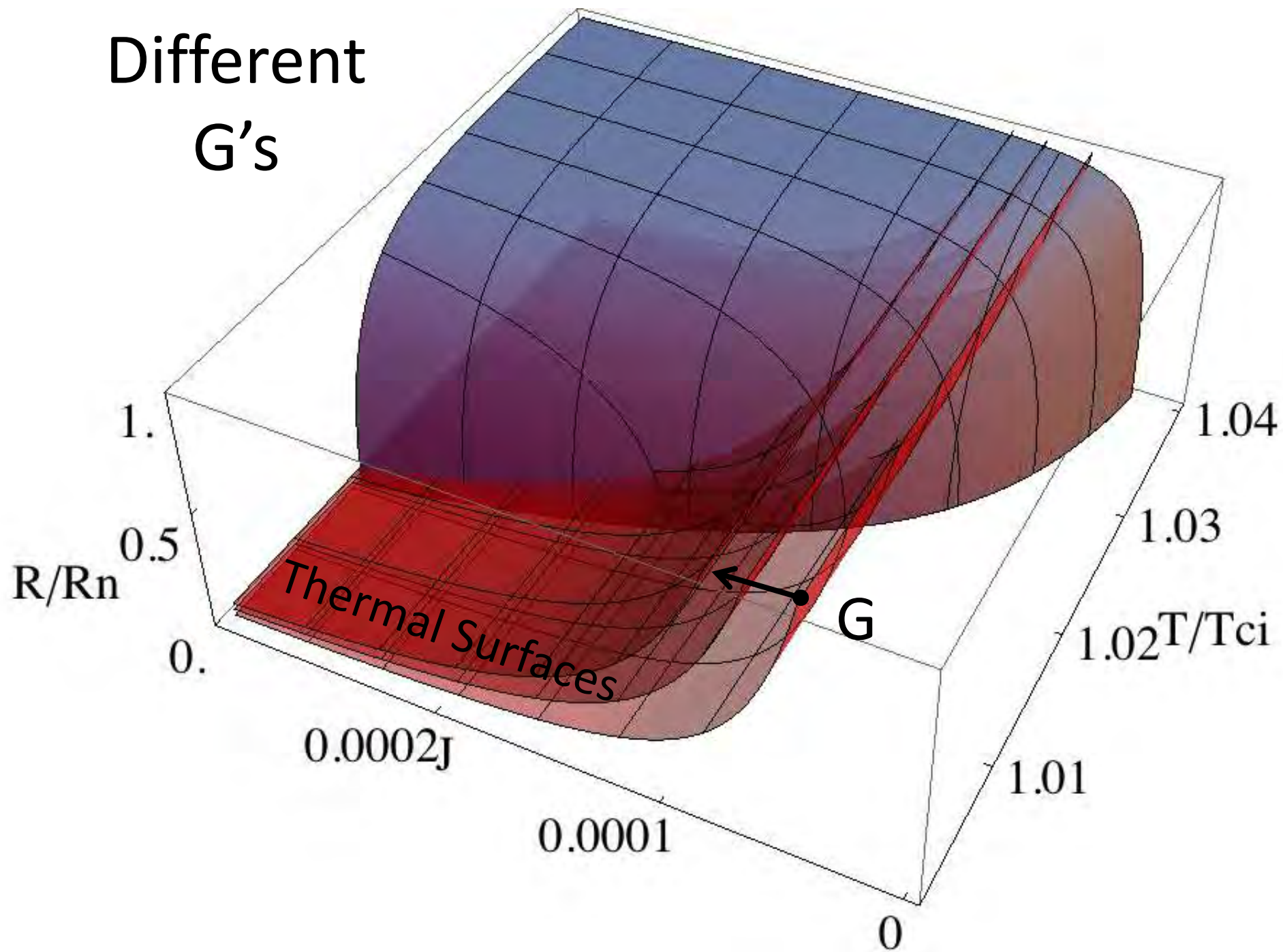
$$I = \frac{V_b}{R_{in}} \frac{R_{sh}}{R + R_{sh}}$$

$$R = \frac{V_b}{R_{in}} \frac{R_{sh}}{I} - R_{sh}$$

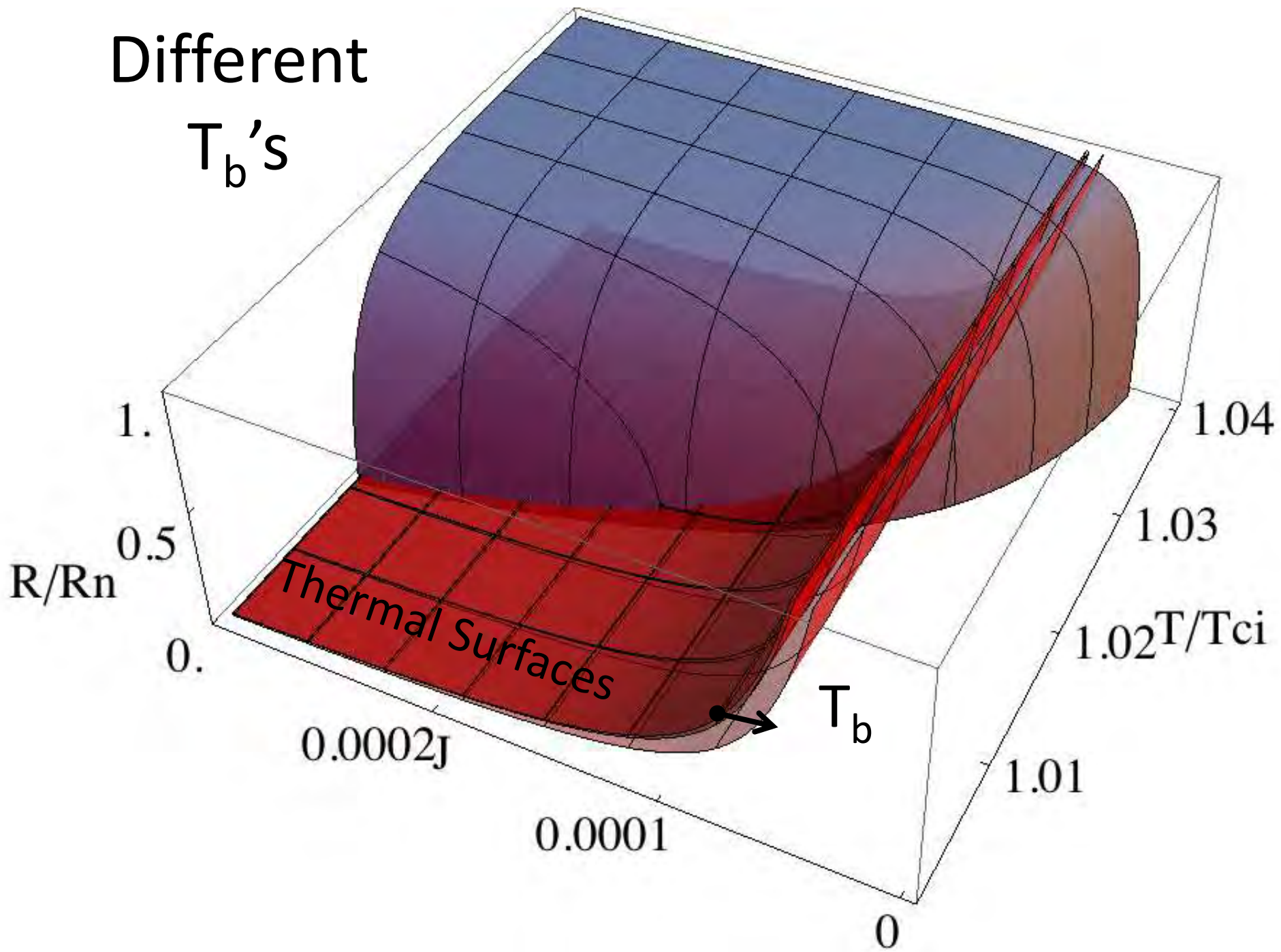


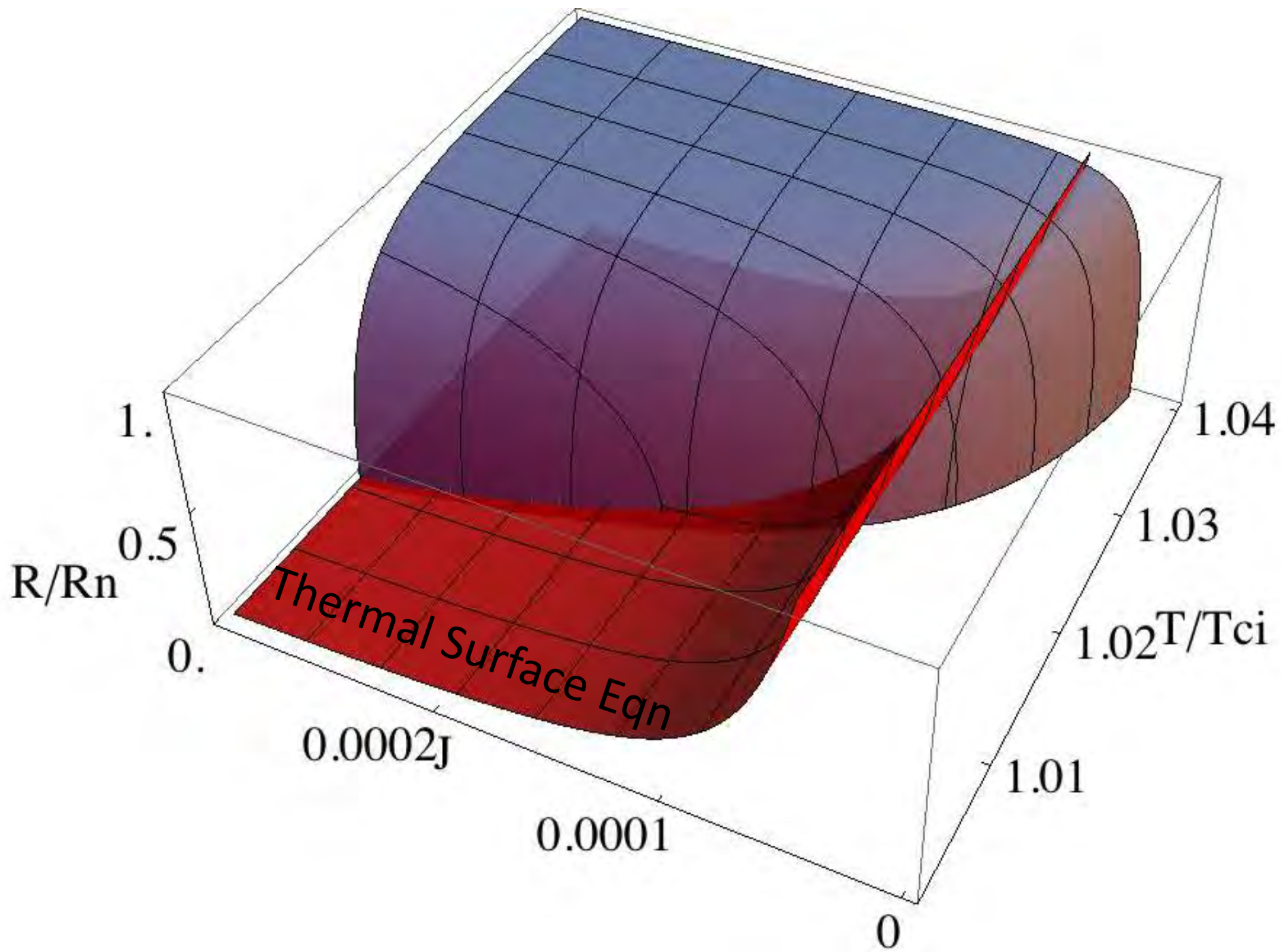


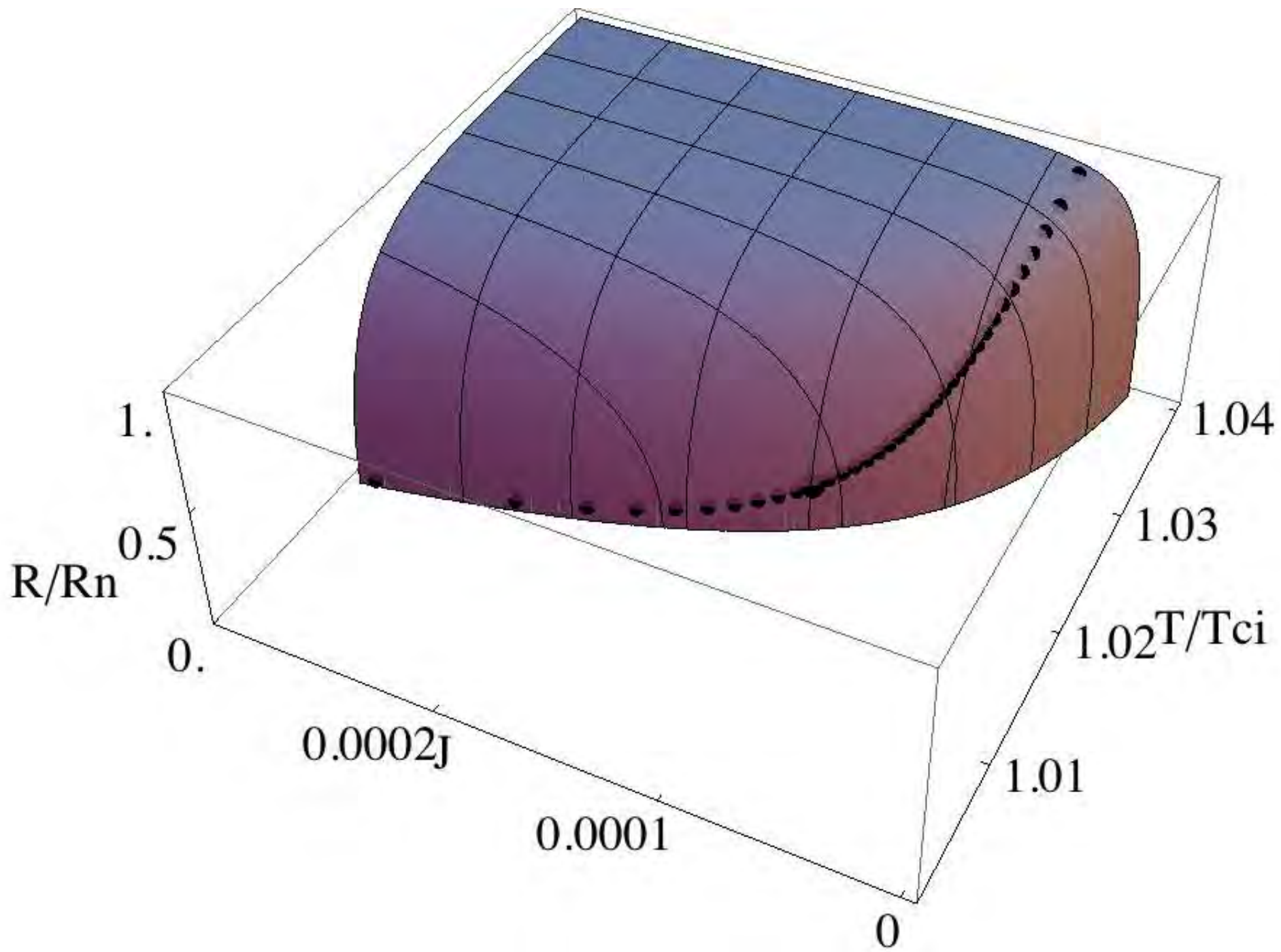
Different G's

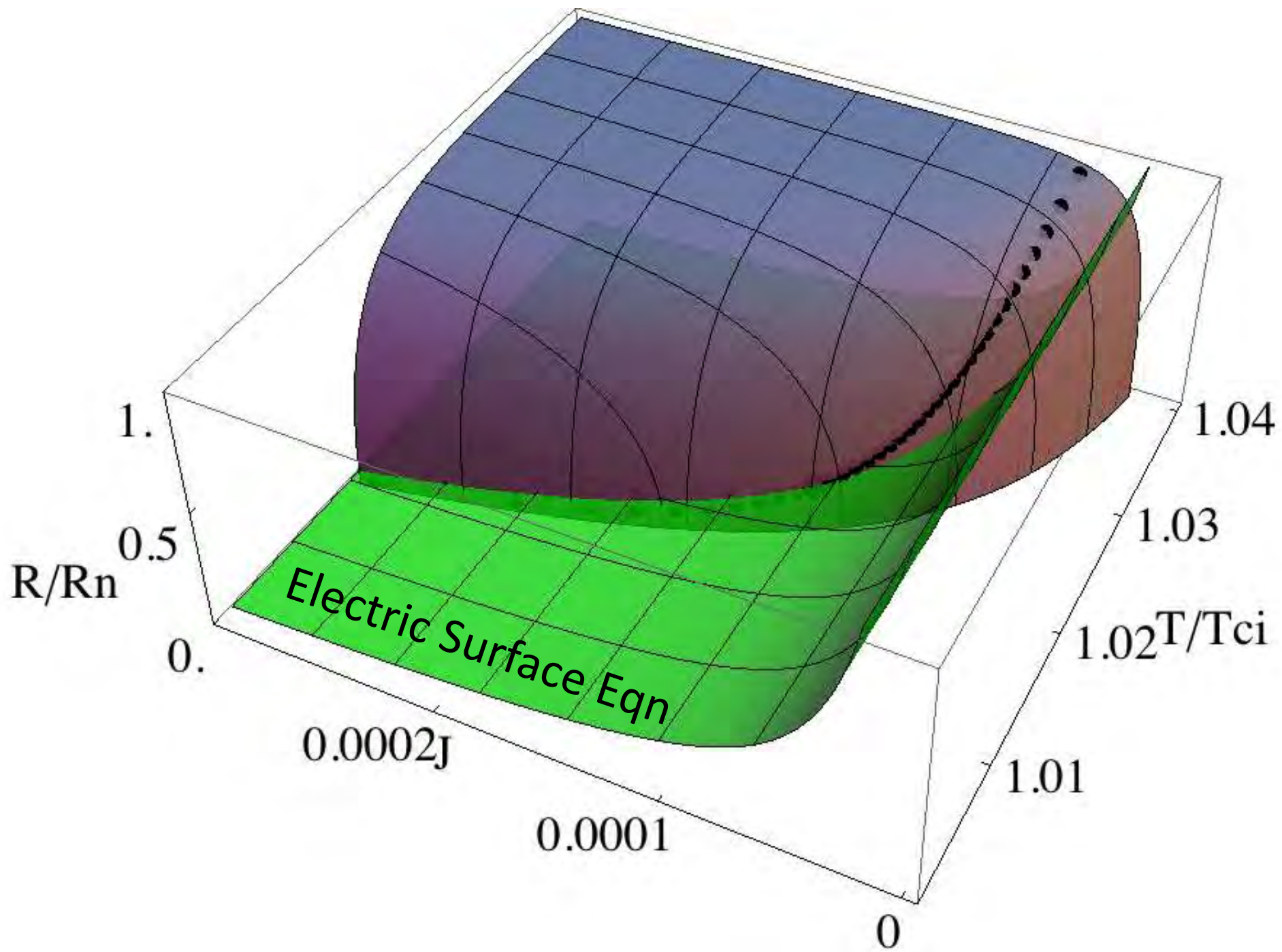


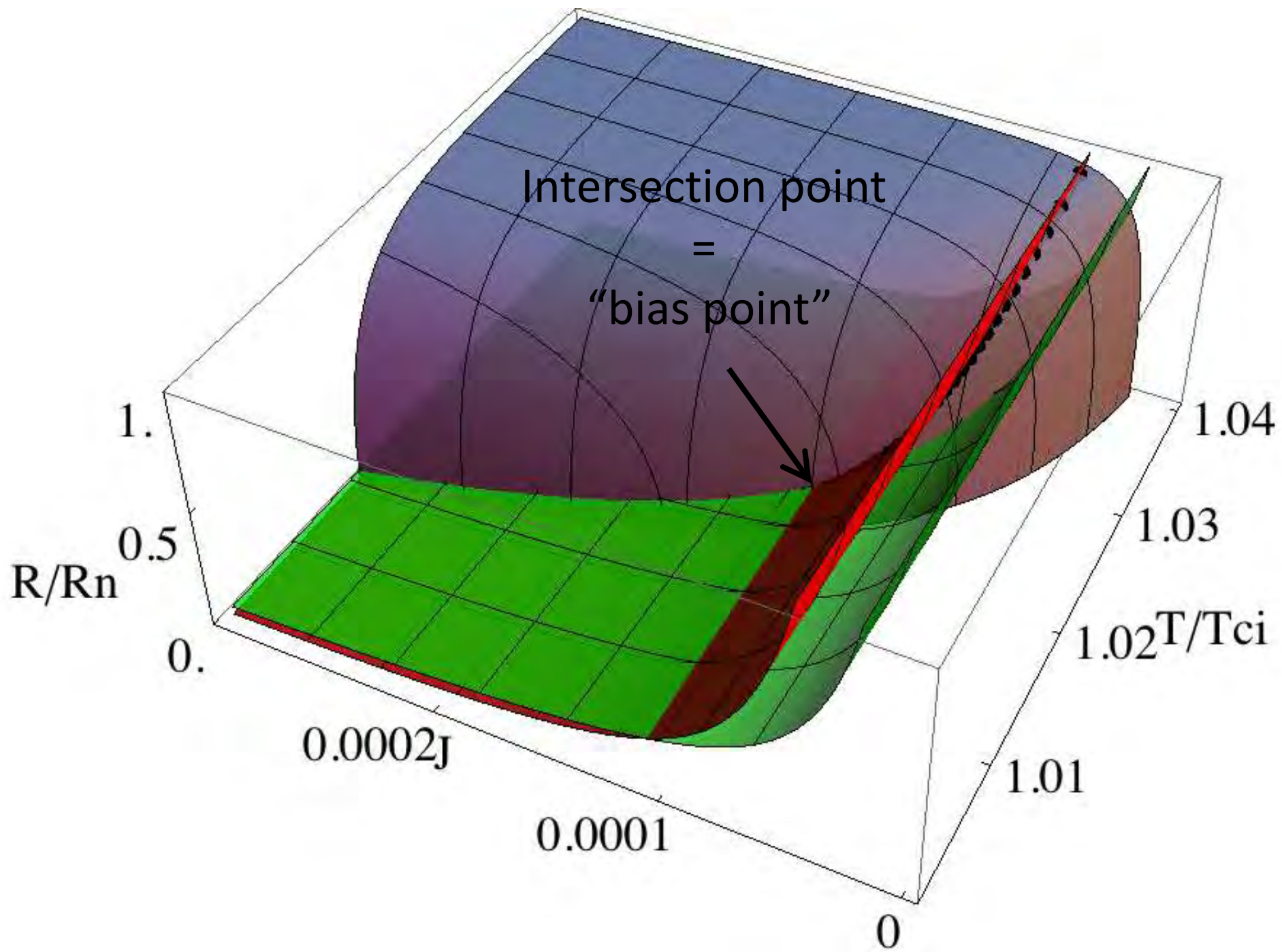
Different
 T_b 's



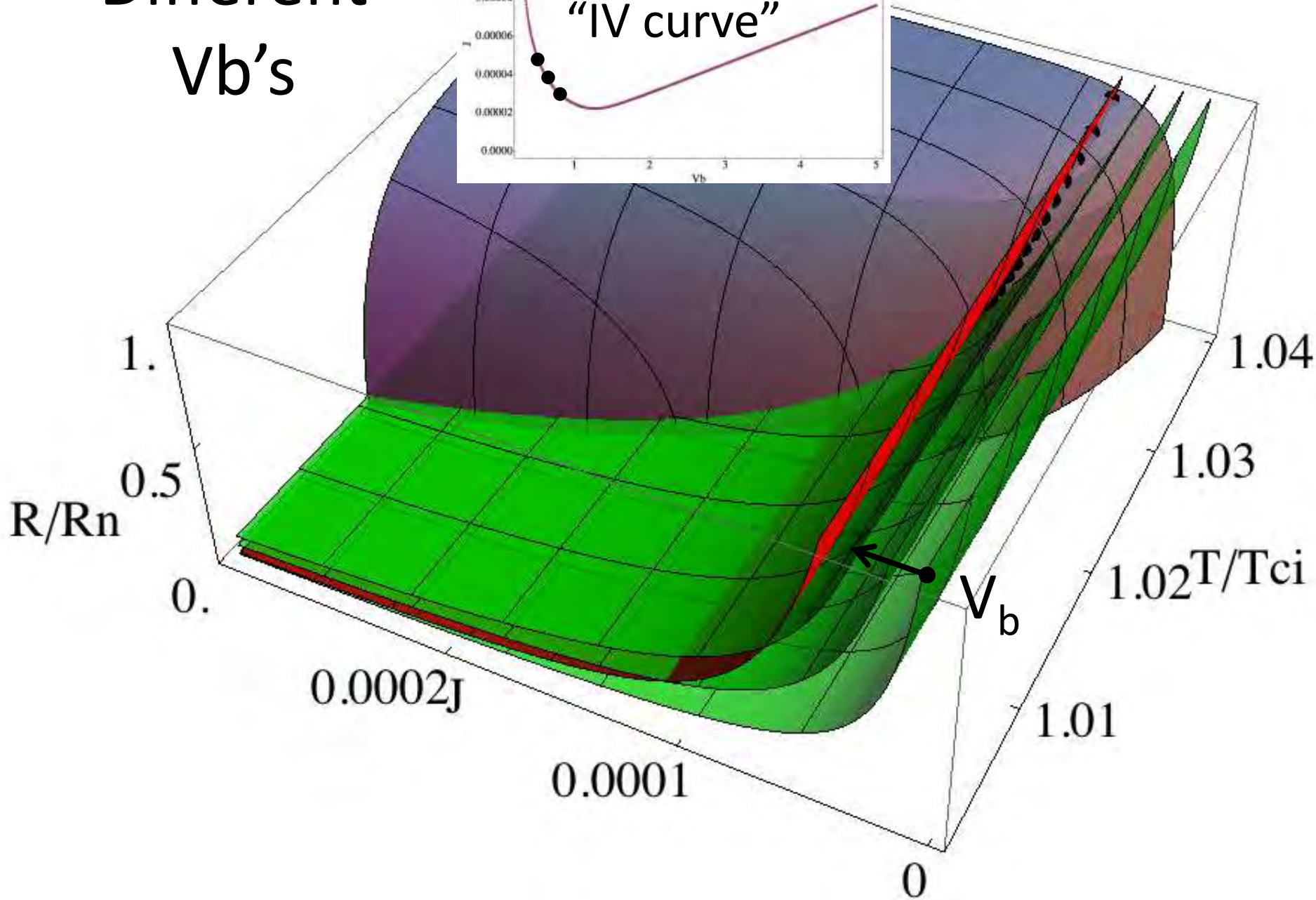
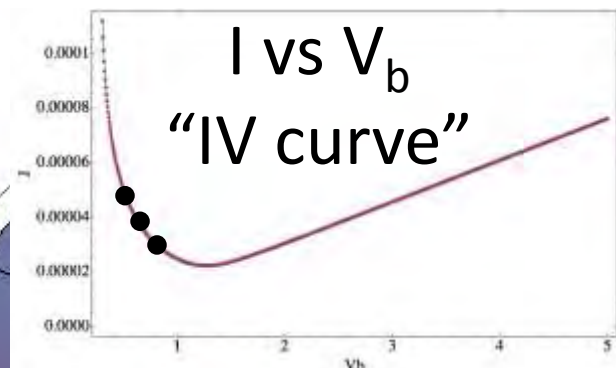




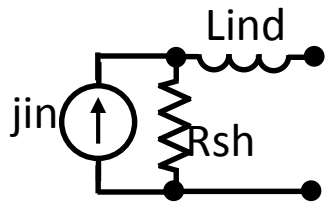




Different
 V_b 's



Bias Circuit



R_{sh}, L

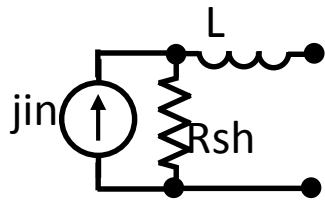
TES



$R(T, j, B)$

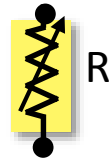
Separable?

Bias Circuit



R_{sh}, L

TES



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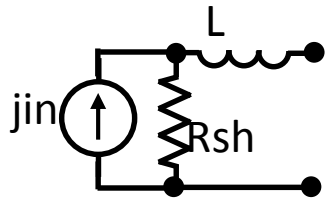
Does changing R_{sh} and L have any direct impact on the TESs $R(T, j, B)$?

Of course changing R_{sh} and L impacts:

- bias stability conditions
- time constants
- the j_{in} needed to get the the same R etc.

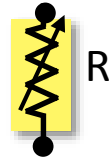
Can we separate the TES transition from the bias circuit?

Bias Circuit



R_{sh}, L

TES



$R(T, j, B)$

Separable?

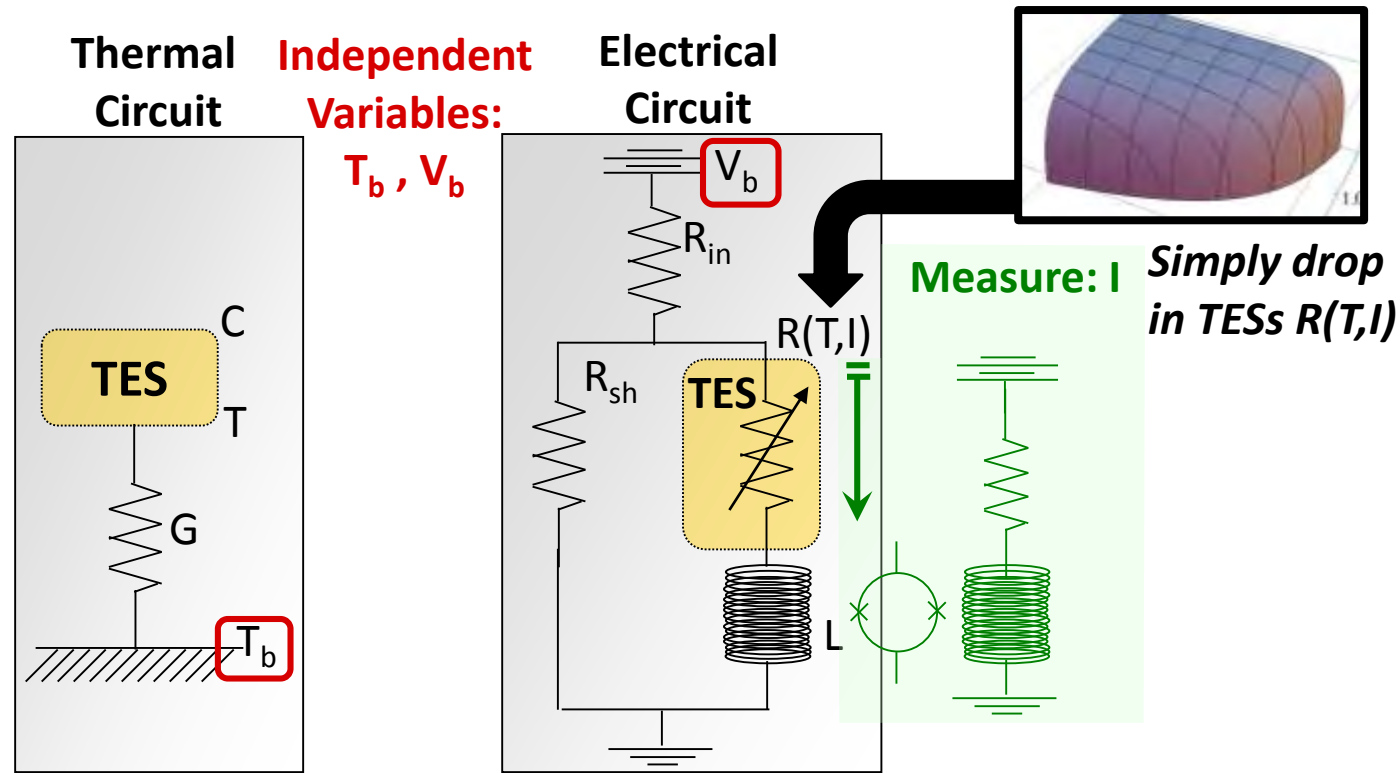
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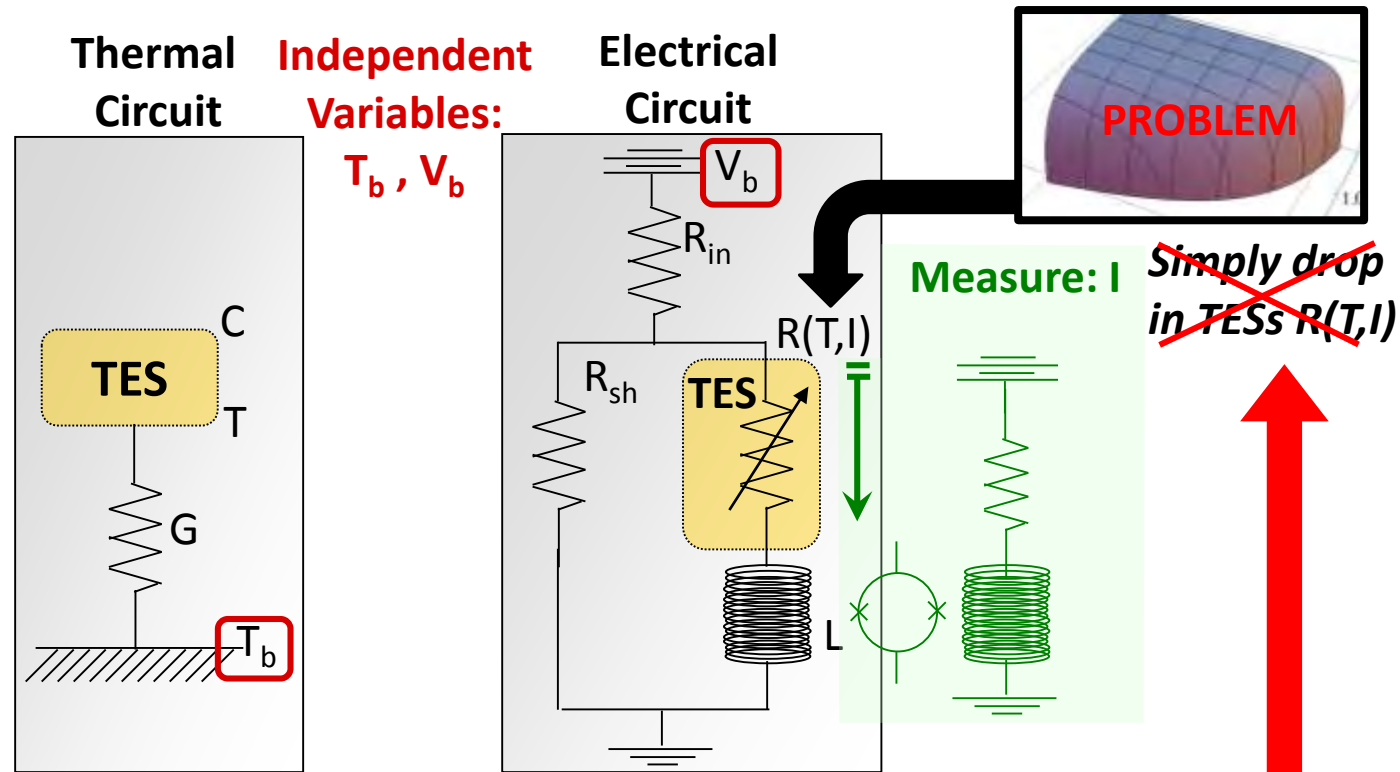
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ANSWER: generally NO.



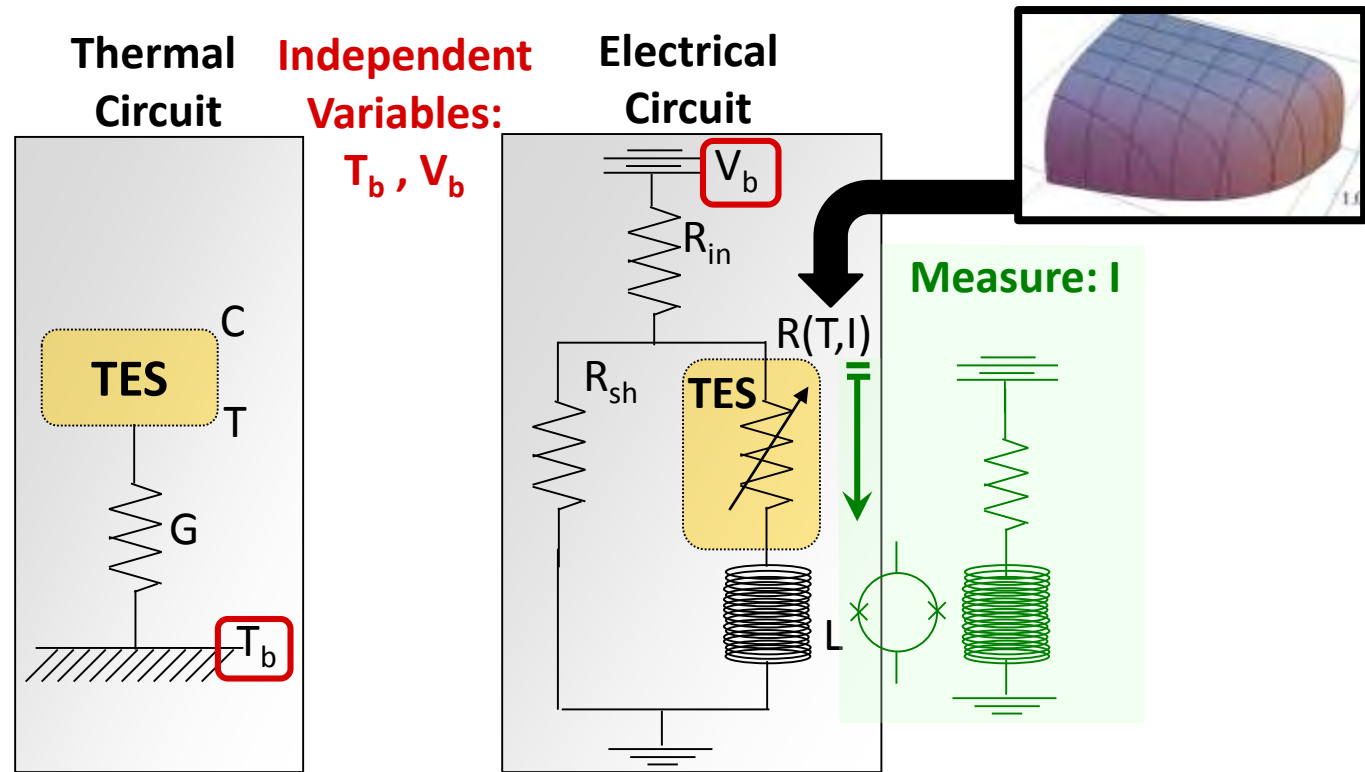
ELECTROTHERMAL MODEL



ELECTROTHERMAL MODEL

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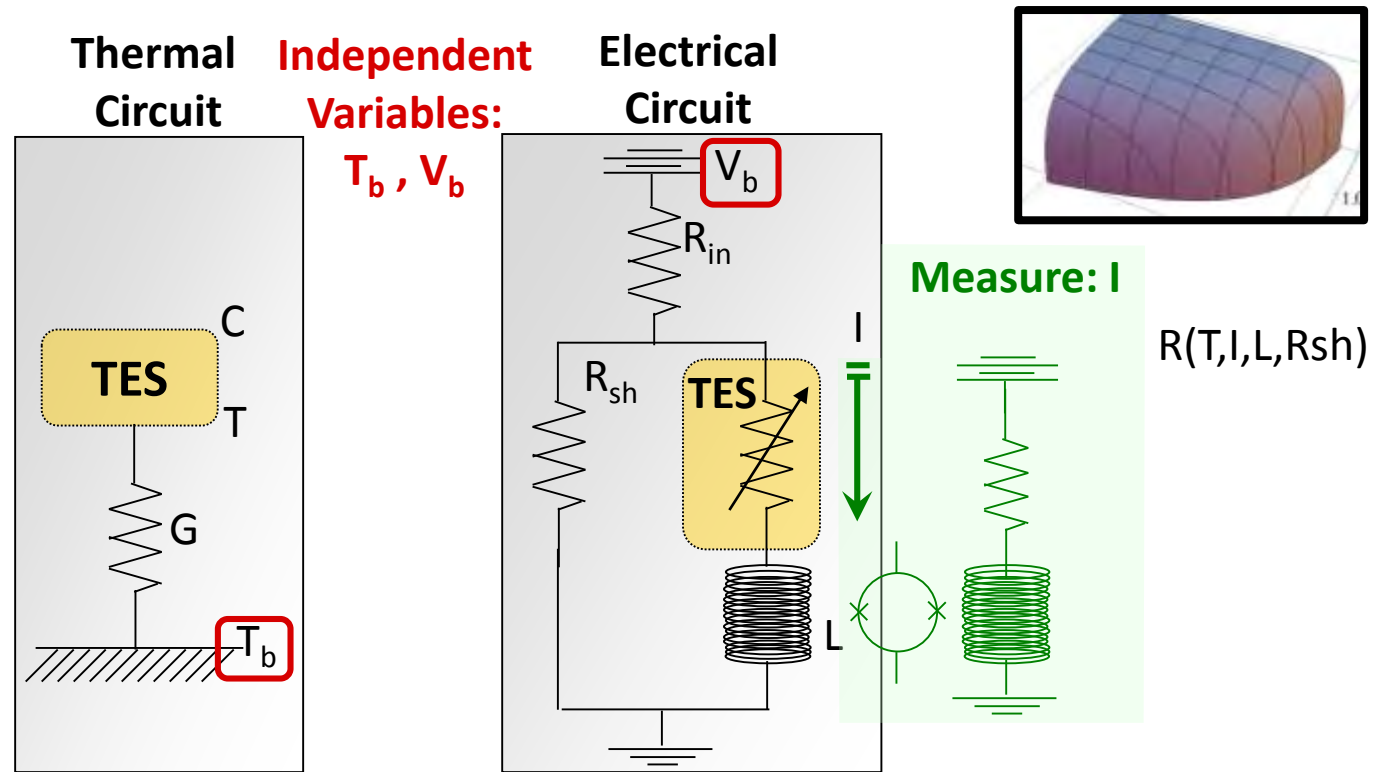
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ELECTROTHERMAL MODEL

Can we separate the TES transition from the bias circuit?

ANSWER: **generally NO.**

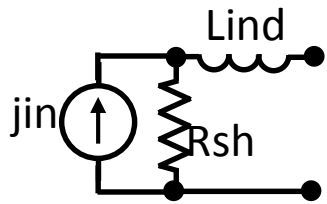


ELECTROTHERMAL MODEL

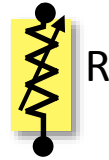
Can we separate the TES transition from the bias circuit?

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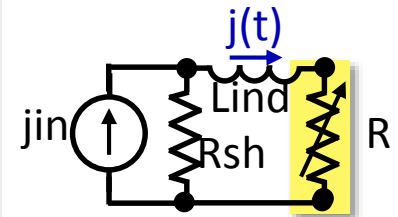
Bias Circuit



TES



Bias Circuit + TES



(1) $j(t)=?$

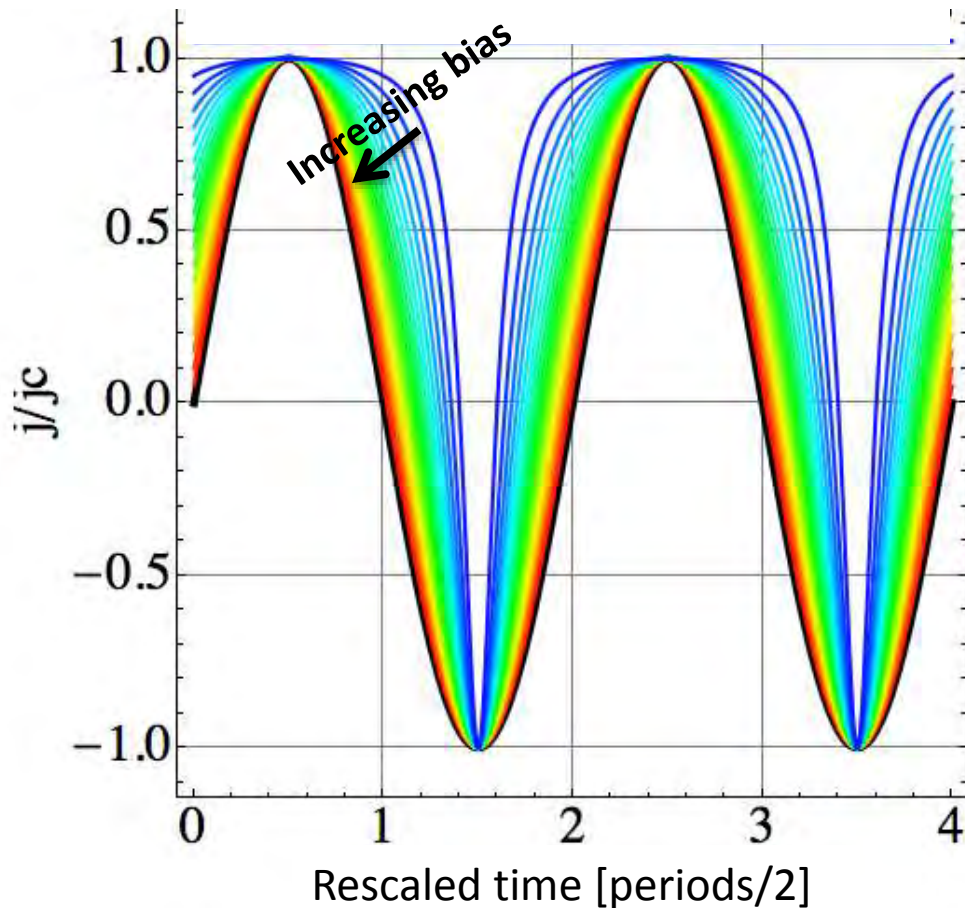
What is the time dependence of the current j ?

(2) $\langle j(t) \rangle$ vs $j_{in} = ?$

What is the shape of the time averaged IV curve?

J-TES model

Time rescaled so we can compare shape in time (harmonic content).



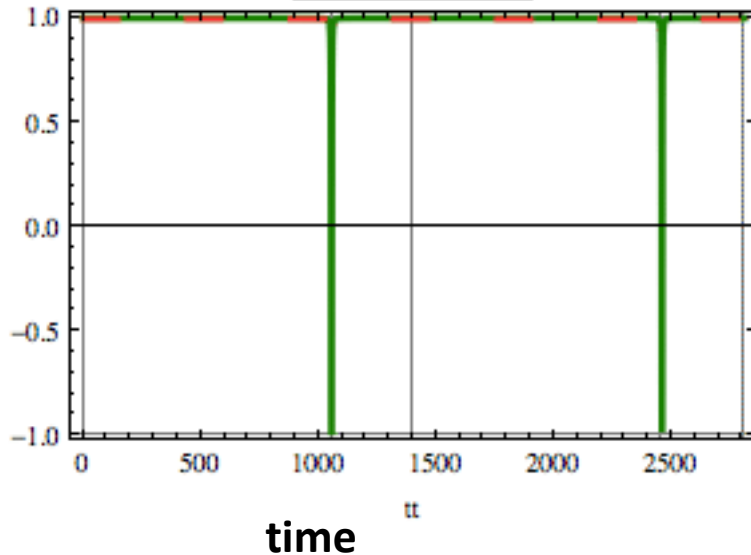
L = 0 limit

Increasing bias $j_{\text{in}} = j_{\text{in}}/j_c$

J-TES model

L = 0 limit

jjin	1.00001
jjavg	0.995538



Green: TES current versus time

Red Dashed: time averaged TES current

Movie evolves with increasing bias current jjin.

Movie rescales time as the bias is increased so two periods are contained in time

Bias jjin:

Small:

slow spikes (+jc to -jc)

Large:

fast sinusoidal (+jc to -jc)

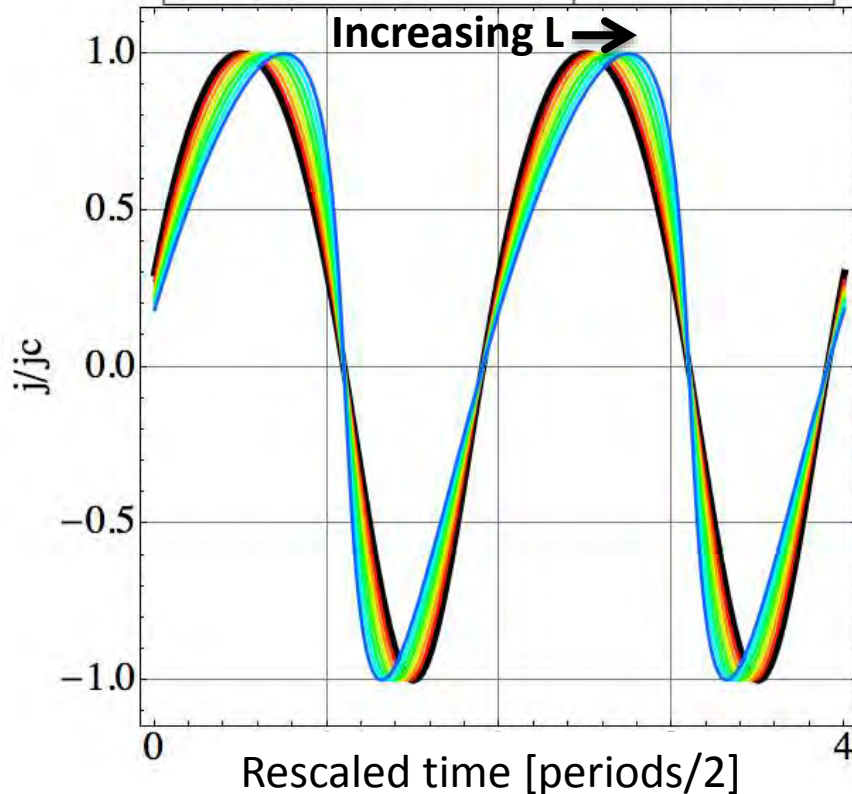
J-TES model

Finite Inductance Effect

**Closed form solution found!*

Large j_{in}

$j_{jin} = \frac{j_{in}}{j_c}$	3.33333
{Lmin,Lmax,Lstep}	{0, 0.7, 0.1}



J-TES model

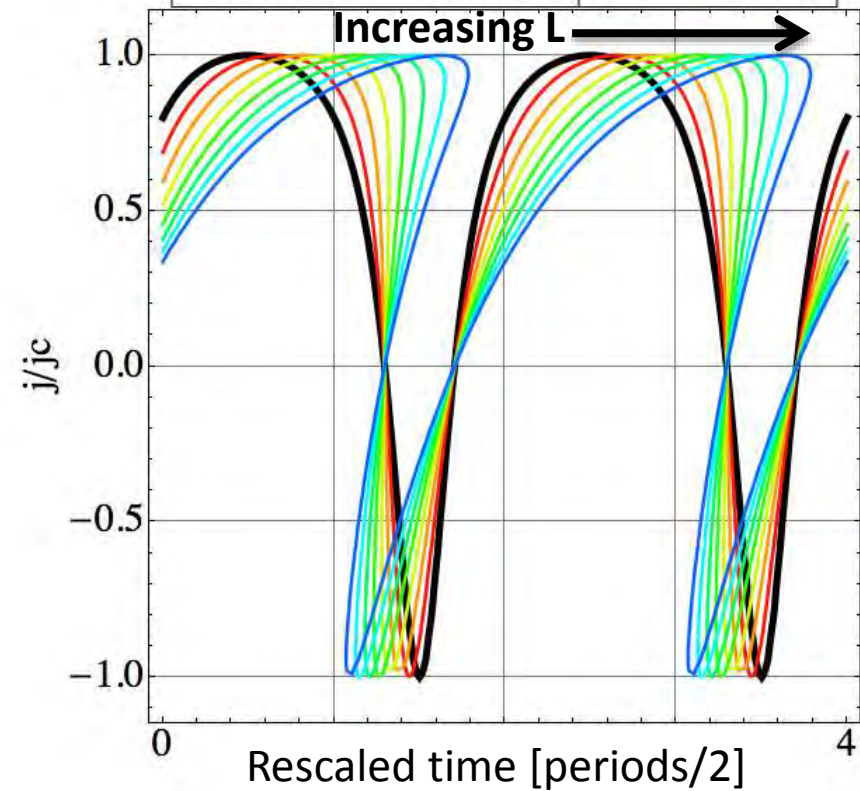
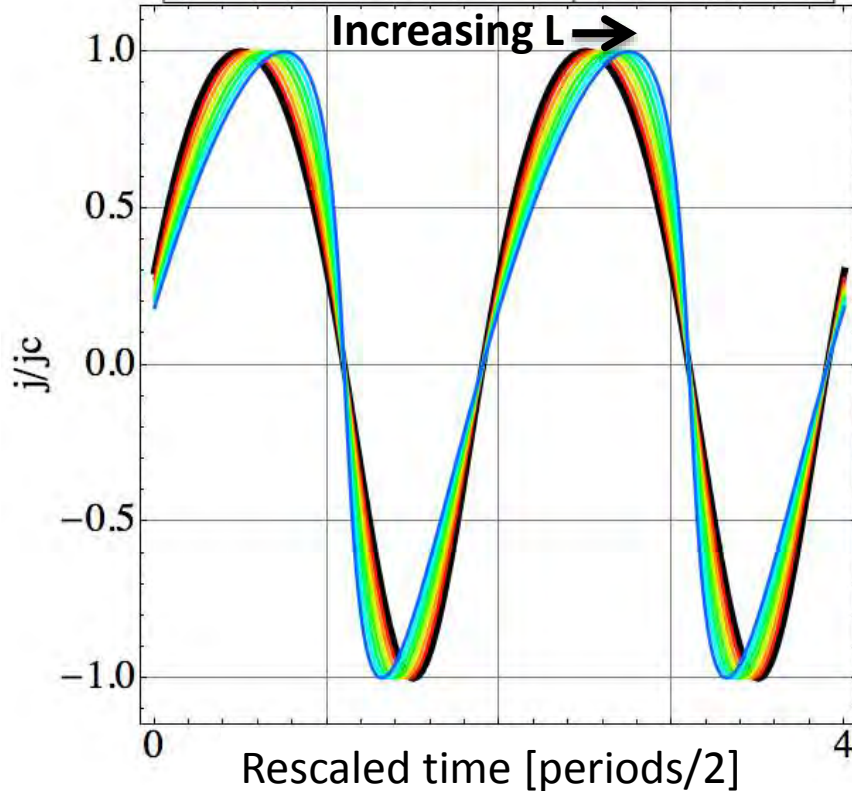
Finite Inductance Effect

Large j_{in}

Small j_{in}

$j_{jin} = \frac{j_{in}}{j_c}$	3.33333
$\{L_{min}, L_{max}, L_{step}\}$	$\{0, 0.7, 0.1\}$

$j_{jin} = \frac{j_{in}}{j_c}$	1.25
$\{L_{min}, L_{max}, L_{step}\}$	$\{0, 0.7, 0.1\}$



J-TES model

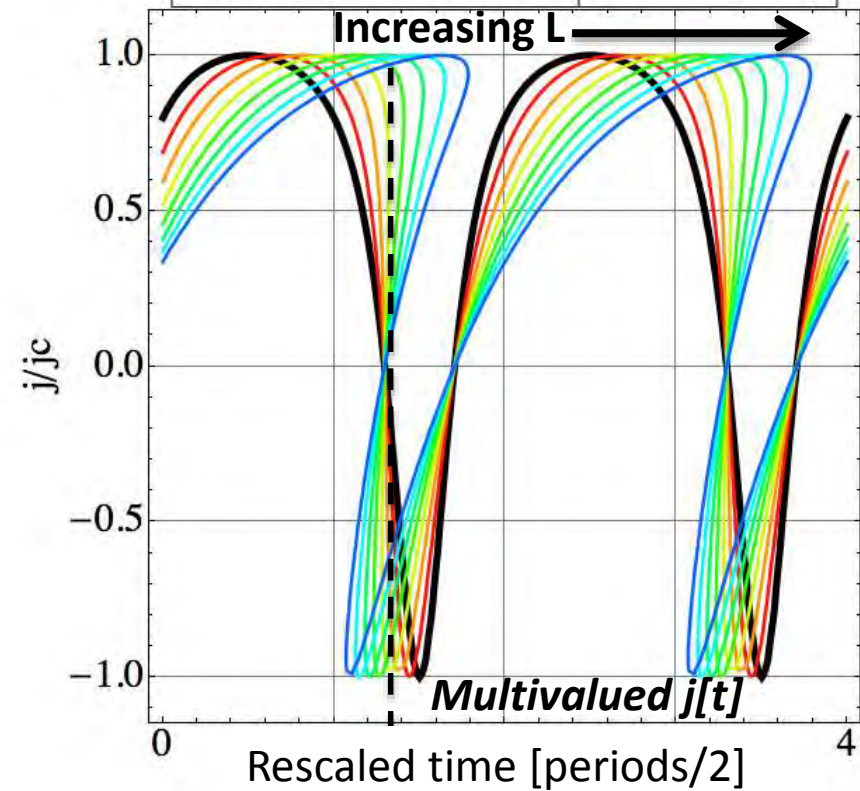
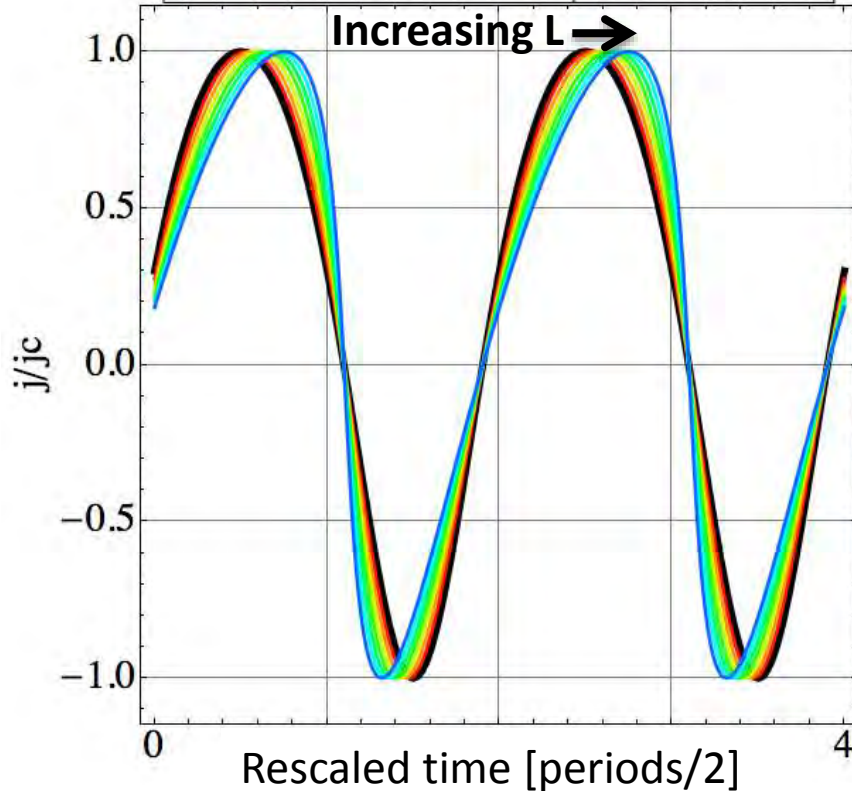
Finite Inductance Effect

Large j_{jin}

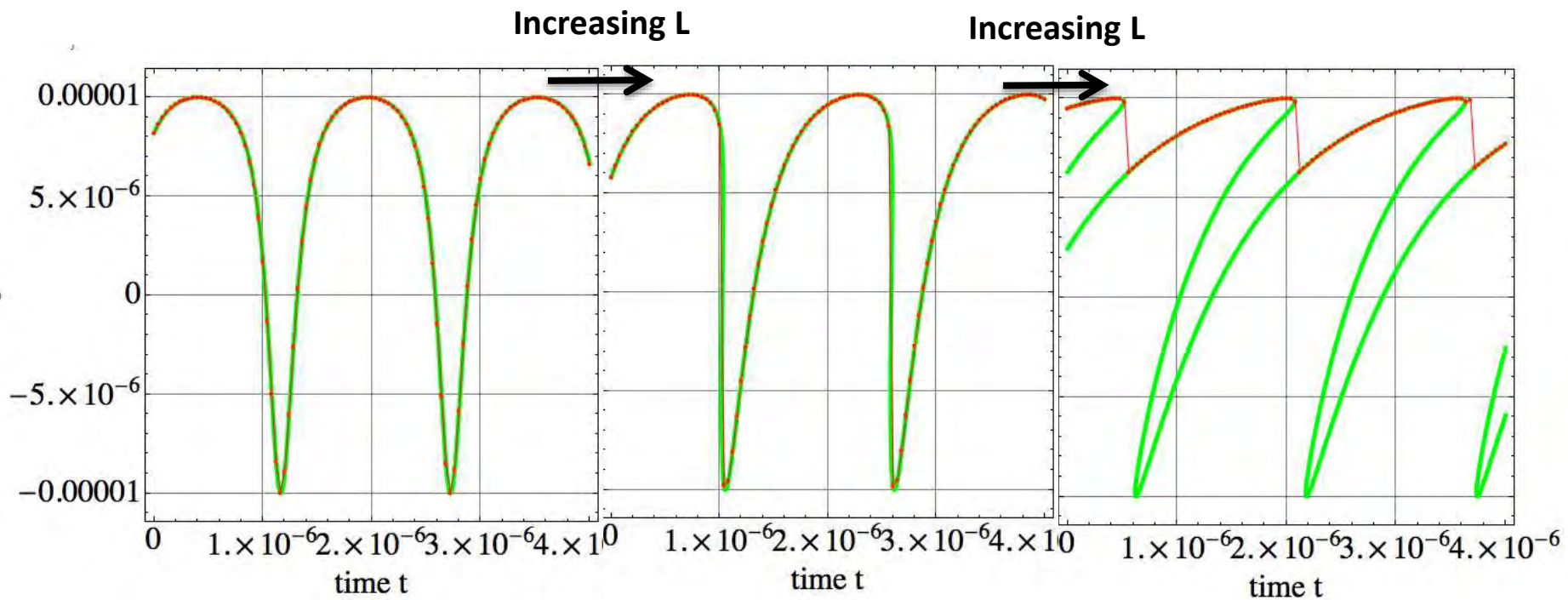
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$j_{jin} = \frac{j_{in}}{j_c}$	1.25
$\{L_{min}, L_{max}, L_{step}\}$	$\{0, 0.7, 0.1\}$



Current becomes multivalued with sufficiently large L

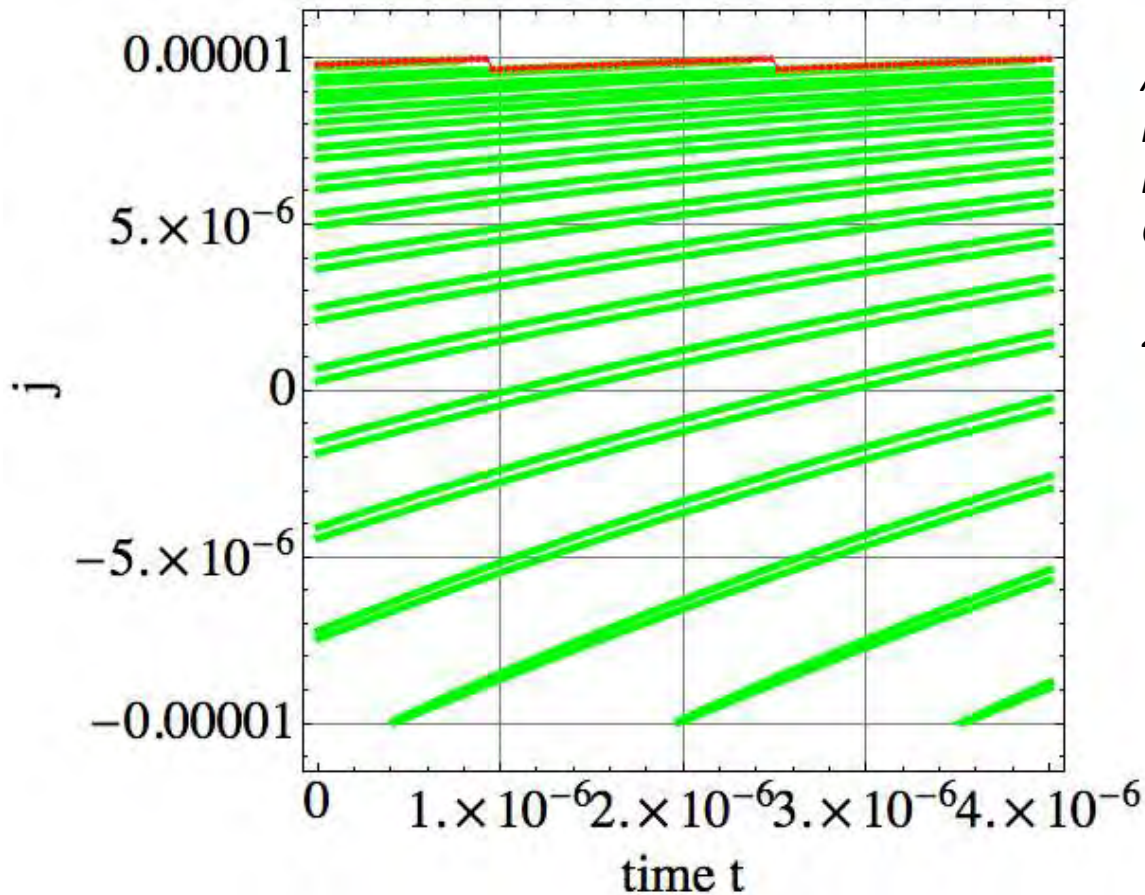


Red: solution

J-TES model

Large L behavior

Large L \rightarrow \sim Sawtooth $j[t]$



*Allowed current states
become very close together
in the large L limit.
Current jump between levels
 \rightarrow
2 level system noise*

Test my claim

Quantum wave function of the superconducting condensate

$$\Psi = |\Psi| e^{i \phi(r,t)}$$

First Josephson Equation: $J = J_c \sin \phi(t)$

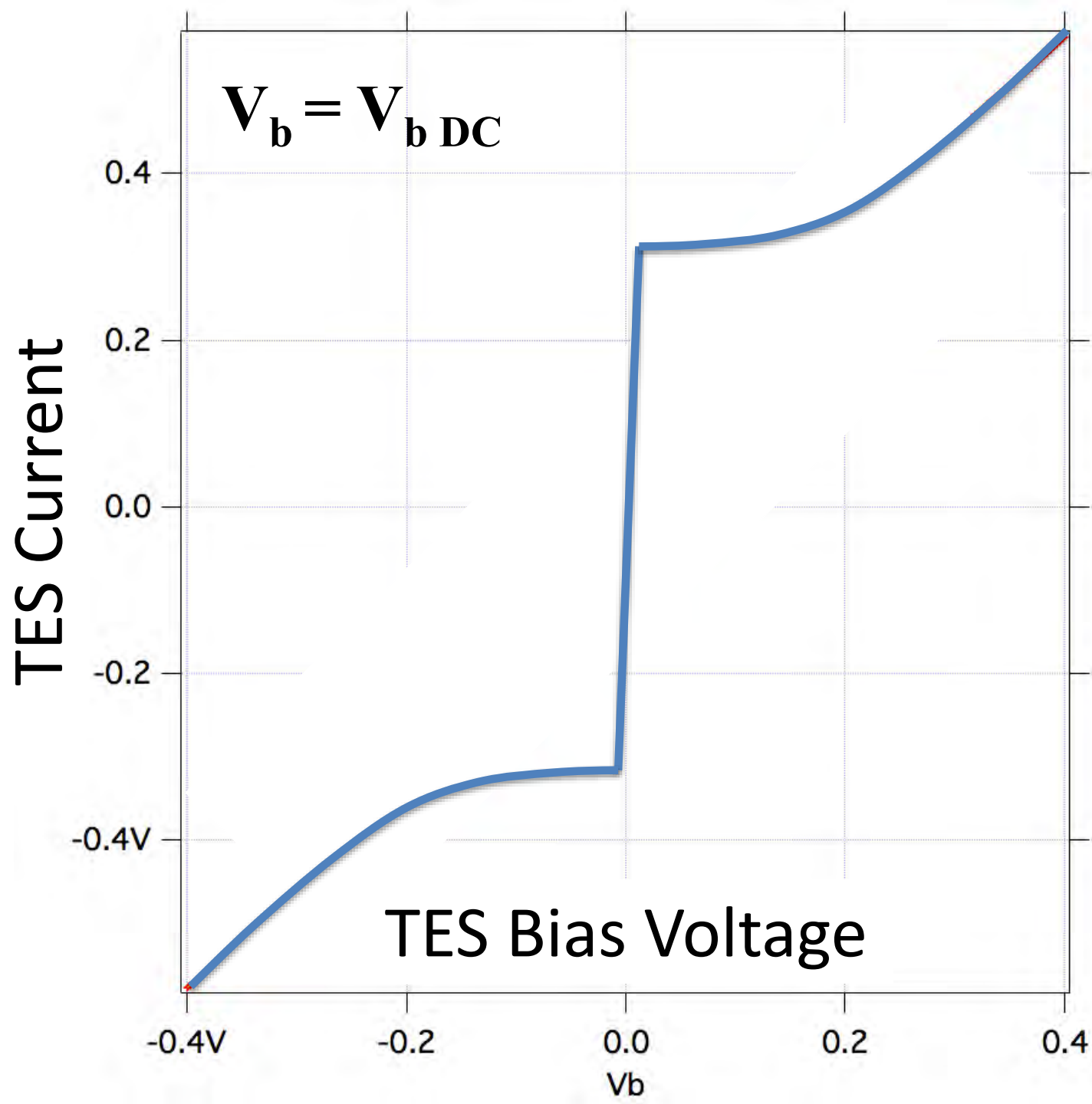
Second Josephson Equation: $\phi'(t) = 2\pi V / \Phi_0$

$$V_b = V_{b \text{ DC}} + \underline{v \sin (2\pi f t)}$$

Shapiro voltage steps: $\underline{V_n = n \Phi_0 f}$
(voltage to frequency transducer)

Fundamental:

1. Gauge Invariance
2. Energy Conservation.



$$V_b = V_{b \text{ DC}} + v \sin(2\pi f t)$$

Phase Locking

TES Current

-0.4V

TES Bias Voltage

-0.4V

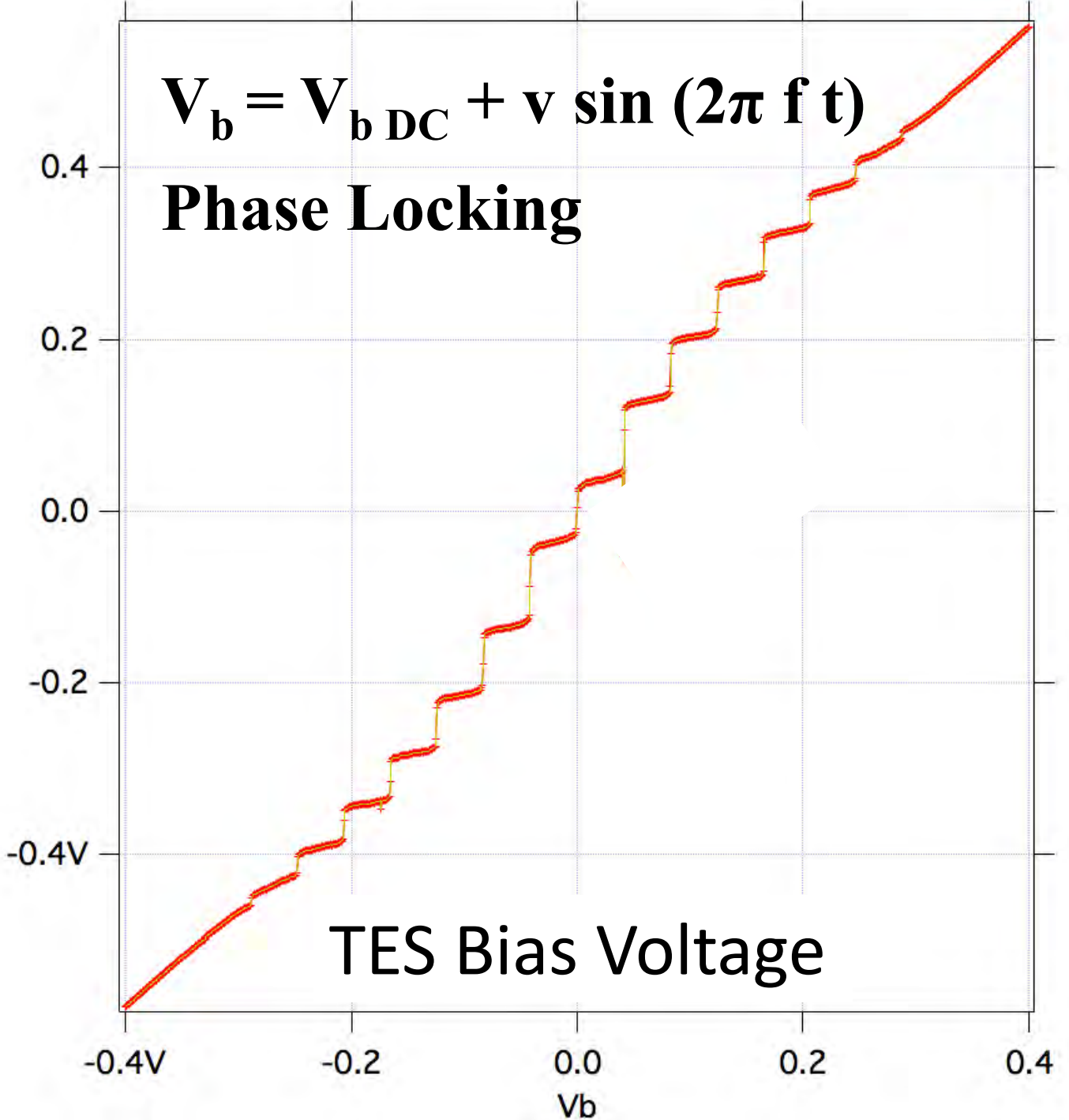
-0.2

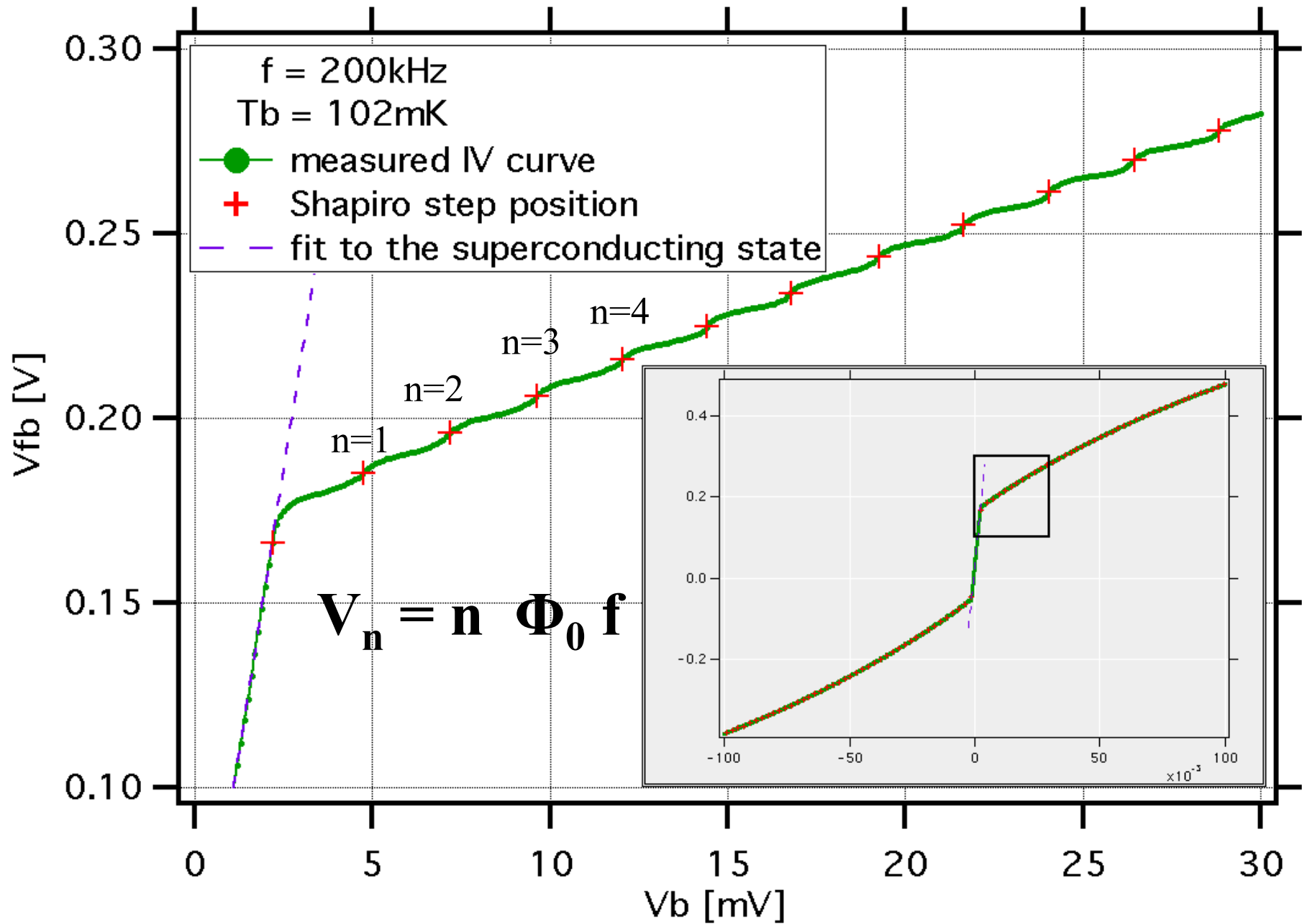
0.0

0.2

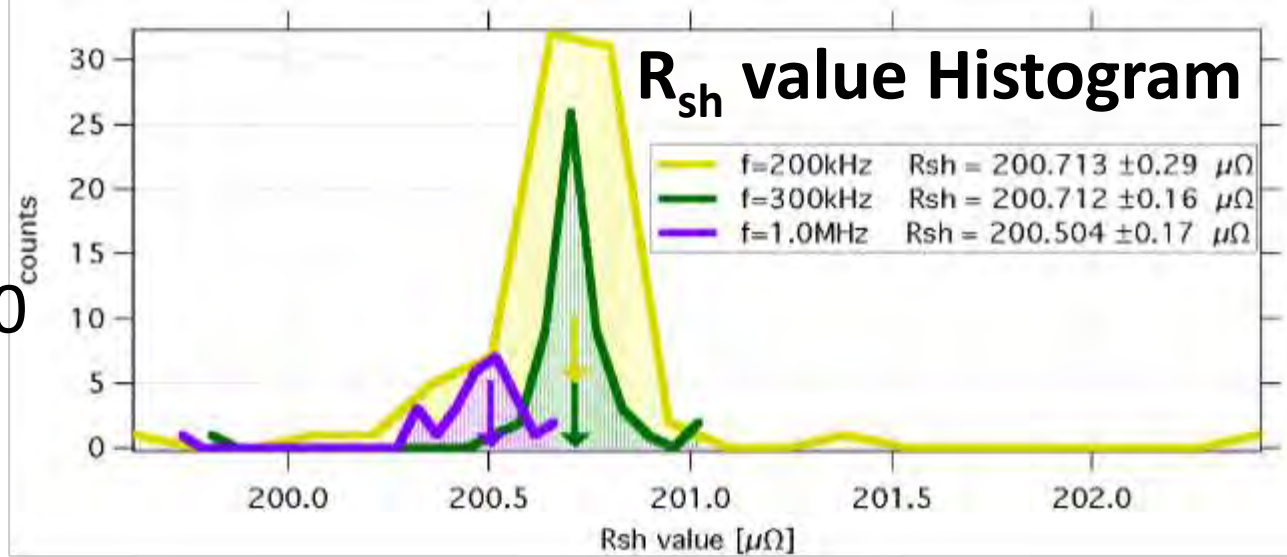
0.4

Voltage Steps
Observed from
f=200 kHz to
f=30 MHz

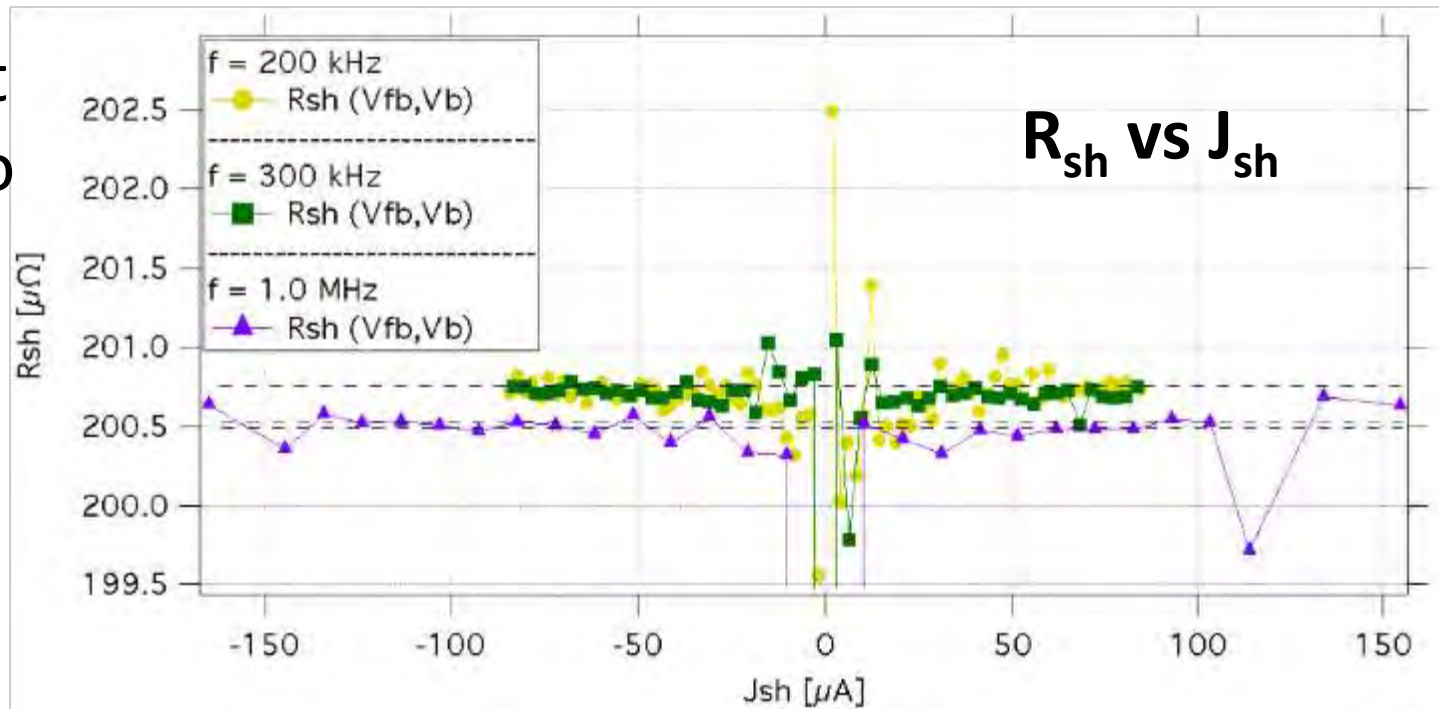




Measuring a 200 μOhm resistor to better than 1/1000 !



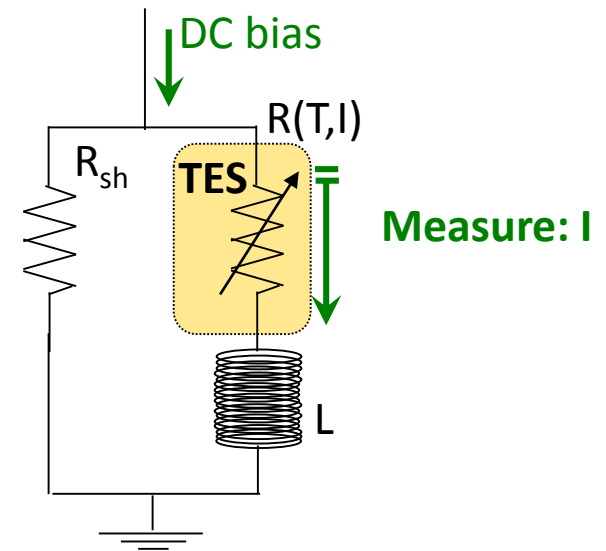
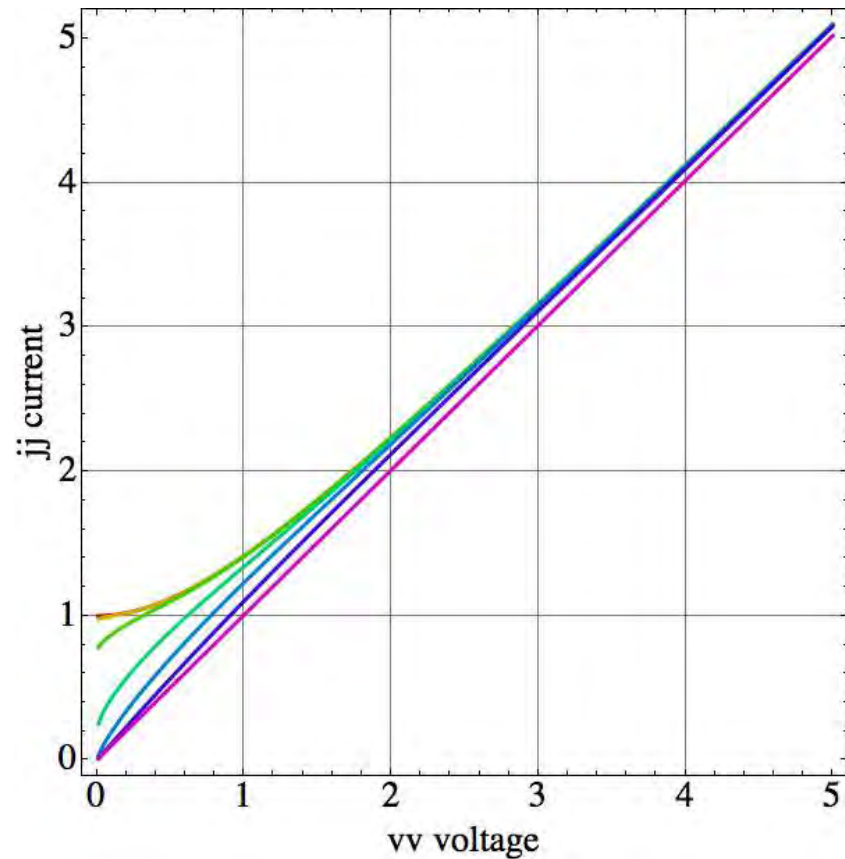
R_{sh} vs current is *flat* down to $5\mu\text{A} \rightarrow$ Ohmic!



Time Averaged IV curves

VRSJ-TES model

Time averaged IV curve.

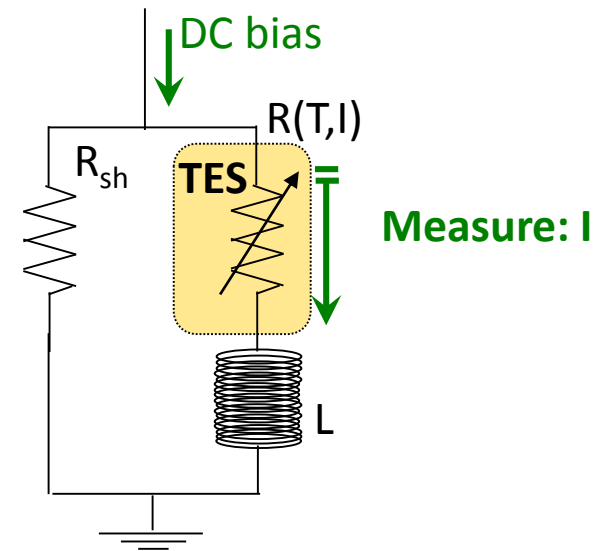
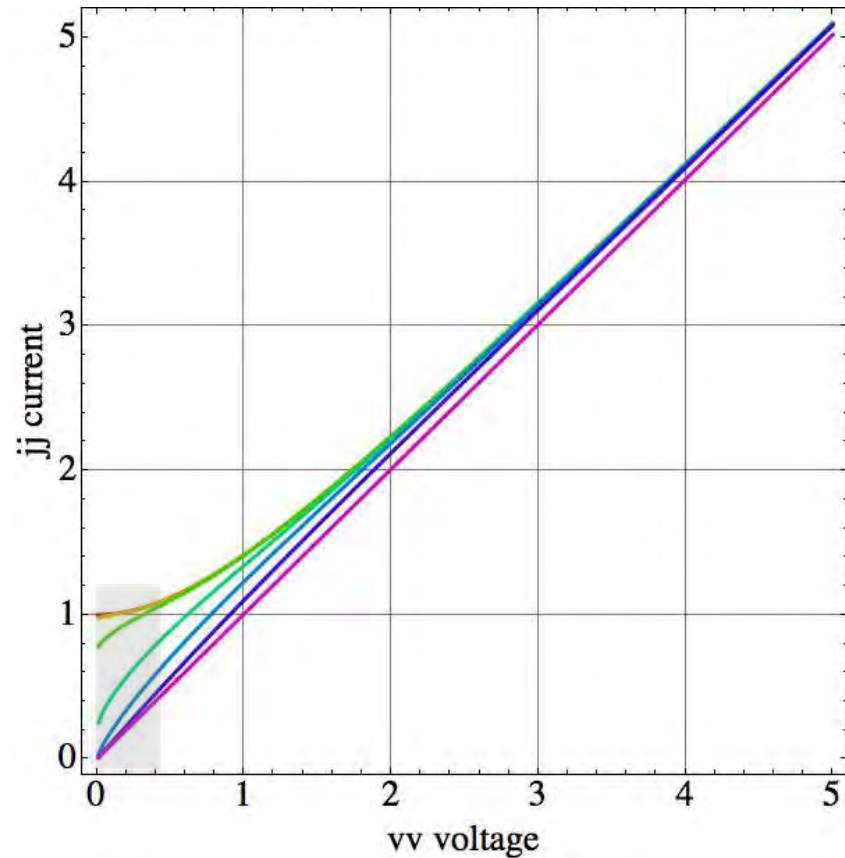


Only changing the circuit Inductor value
Everything else constant.

Time Averaged IV curves

VRSJ-TES model

Time averaged IV curve.

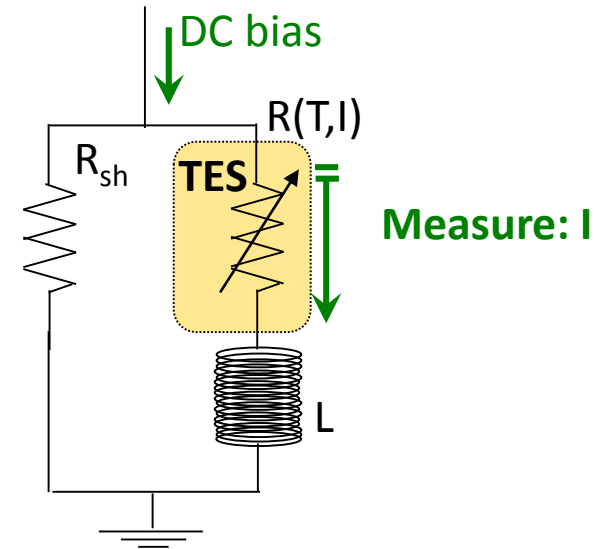
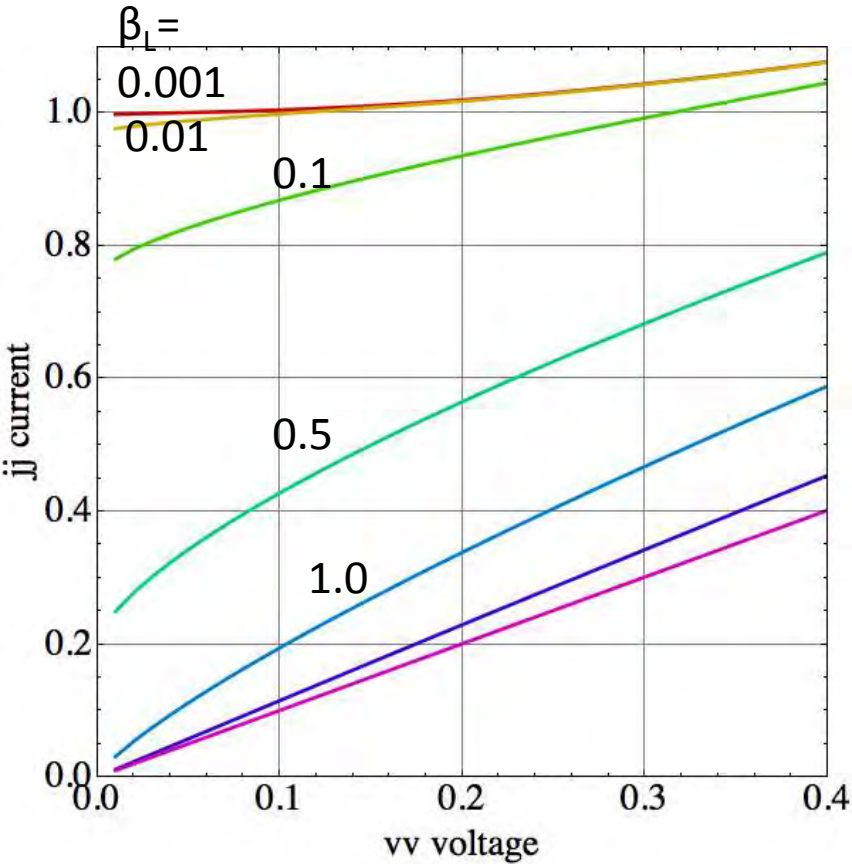


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Time Averaged IV curves

VRSJ-TES model

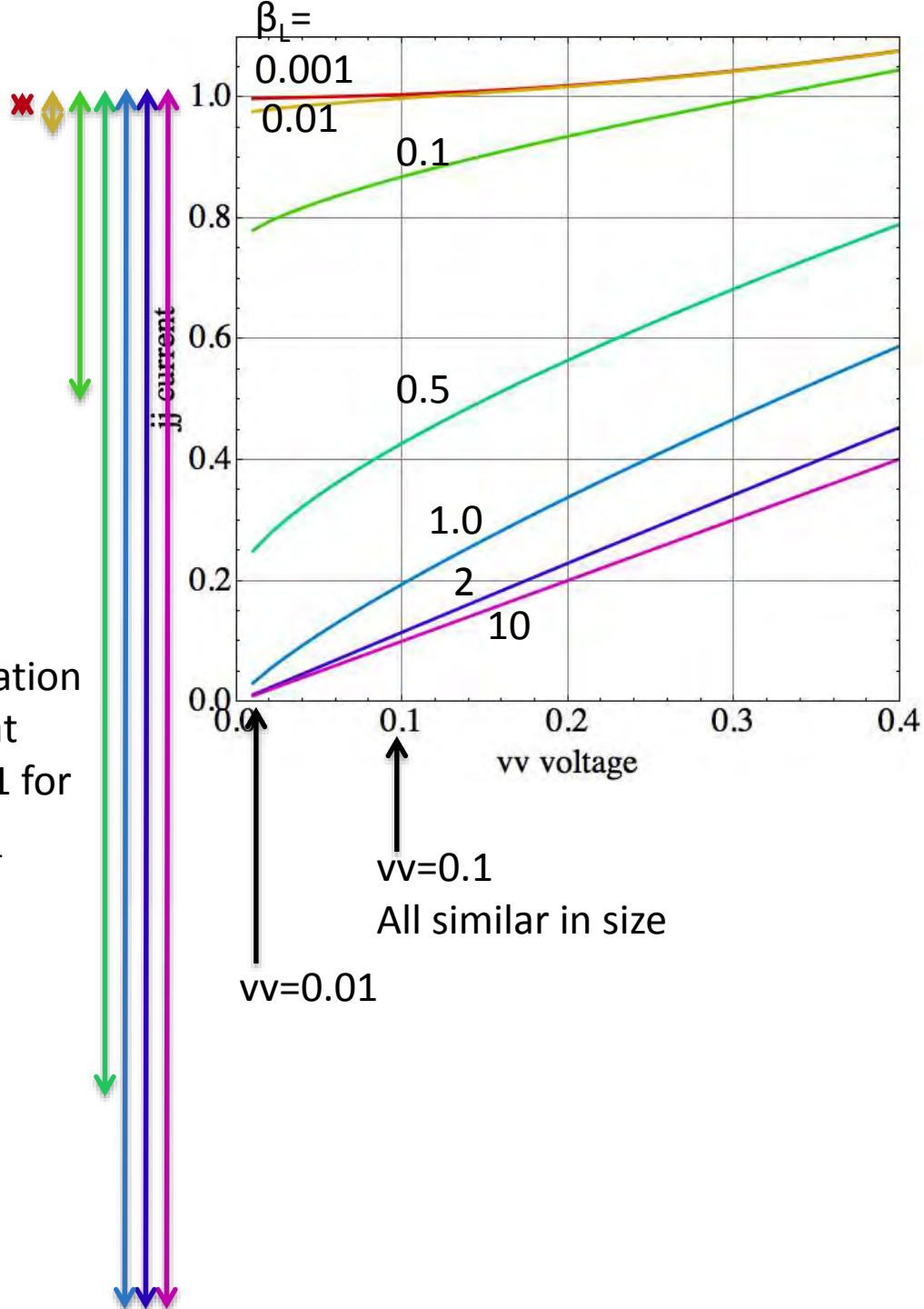
Time averaged IV curve.



Only changing the circuit Inductor value
Everything else constant.

jj oscillation amplitude

jj oscillation range at $v_v=0.01$ for each β_L value.

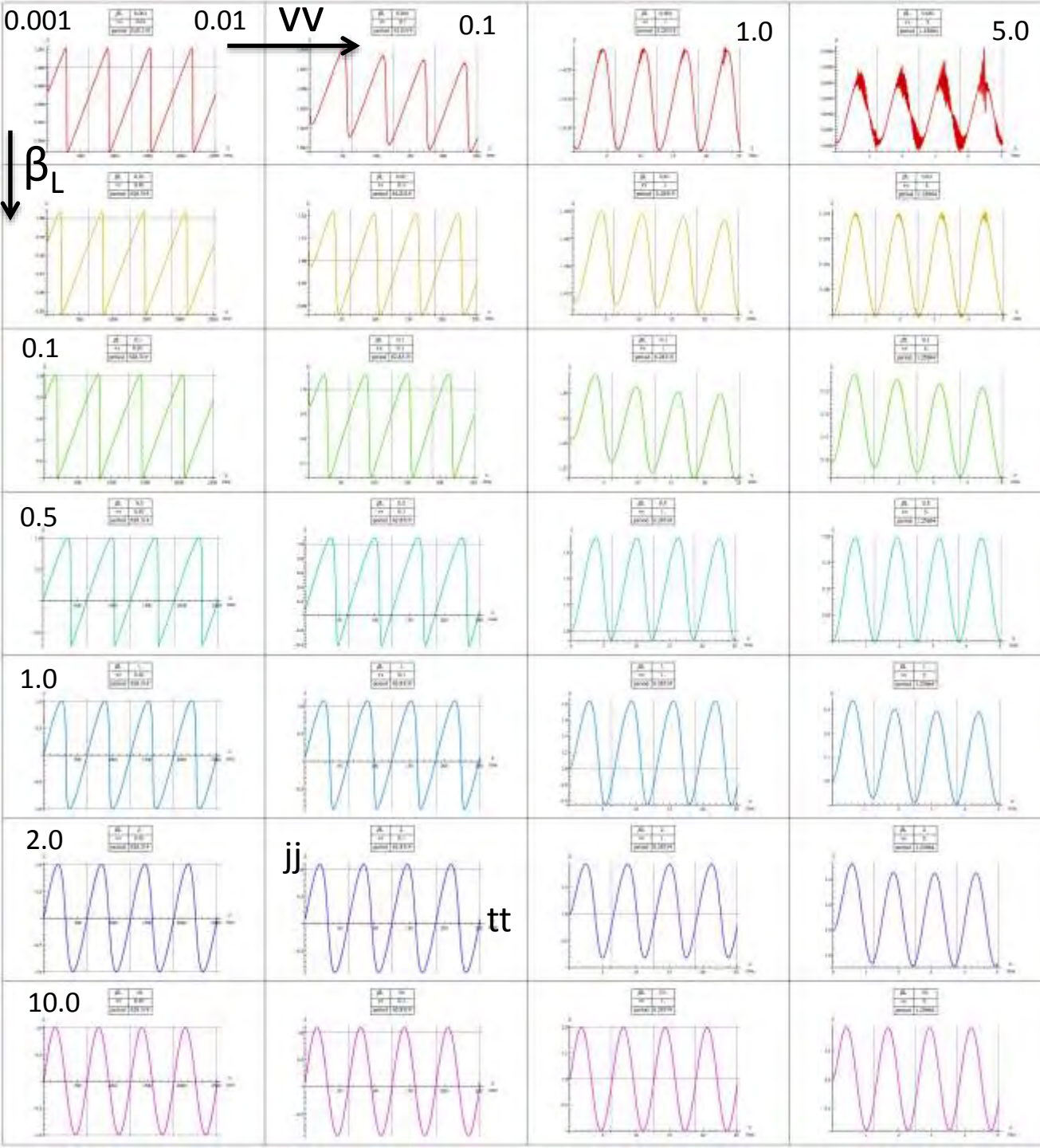


Conclusions: ...*Paradigm Shifting*

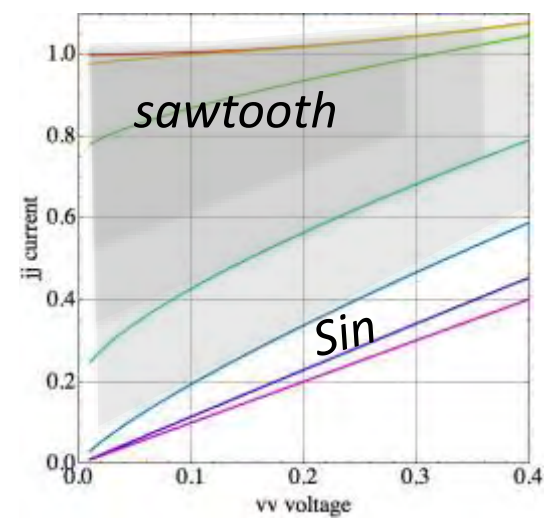
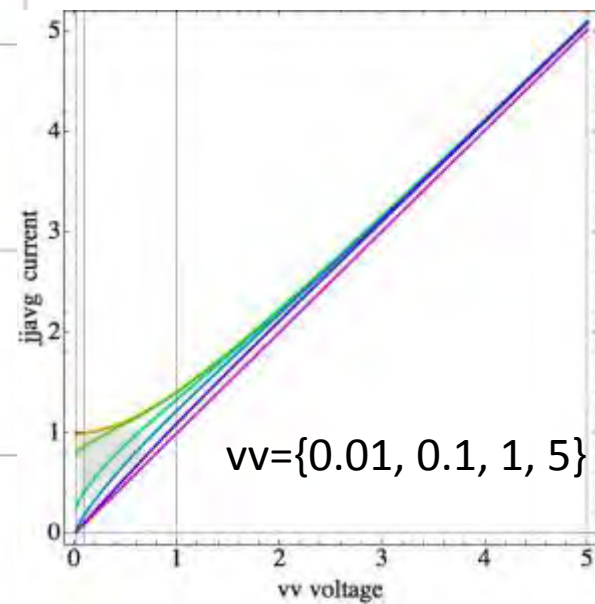
- In steady state with standard DC biased TES operation. The TES current is wildly oscillating in time at high frequencies and we only measure (*and are only aware of*) the time averaged current in the TES.
- The equations governing this time response is **nonlinear** time-dependent diff eqs.
 - **KEY POINT:** This time dependent current exists when the TES is only DC biased, NO time dependent inputs what so ever. It comes about from the intrinsic superconductivity physics governing TESs.
- The measured resistive transition surface is not simply a function of $R(T,J,B)$. But is also a function of the electric circuit. $R(T,J,B,L,R_{sh})$.
 - → Problem compartmentalizing: → *Same TES measured in slightly different circuits will appear to have a different resistive transition surface. TEST in real setup early!!!*
- Circuit and TES parameters can have values such that these time dependent solutions become multivalued in time.
 - *What does that mean?* Can mean excess noise or fine structure of the time averaged resistive transition surface.
- The biased TES waiting to detect a photon (μ -calorimeter) or flux of photons (μ -bolometers) can itself act as a radiation source.
 - Tricky when an array of exquisitely sensitive micro-wave radiation detectors are themselves sources of microwave radiation.
 - This can lead to radiation resistance and fine structure in the resistive transition.
 - Possible cross talk between pixels in an array.
 - The time dependence of the current can take on pure sin waves and also very nonsinusoidal forms (variable harmonic content)
 - The fundamental frequency of these oscillations changes with bias voltage.
 - Makes the prospects of FDM TES arrays with an AC-biased TES challenging. Structure, sensitivity, and cross-talk.

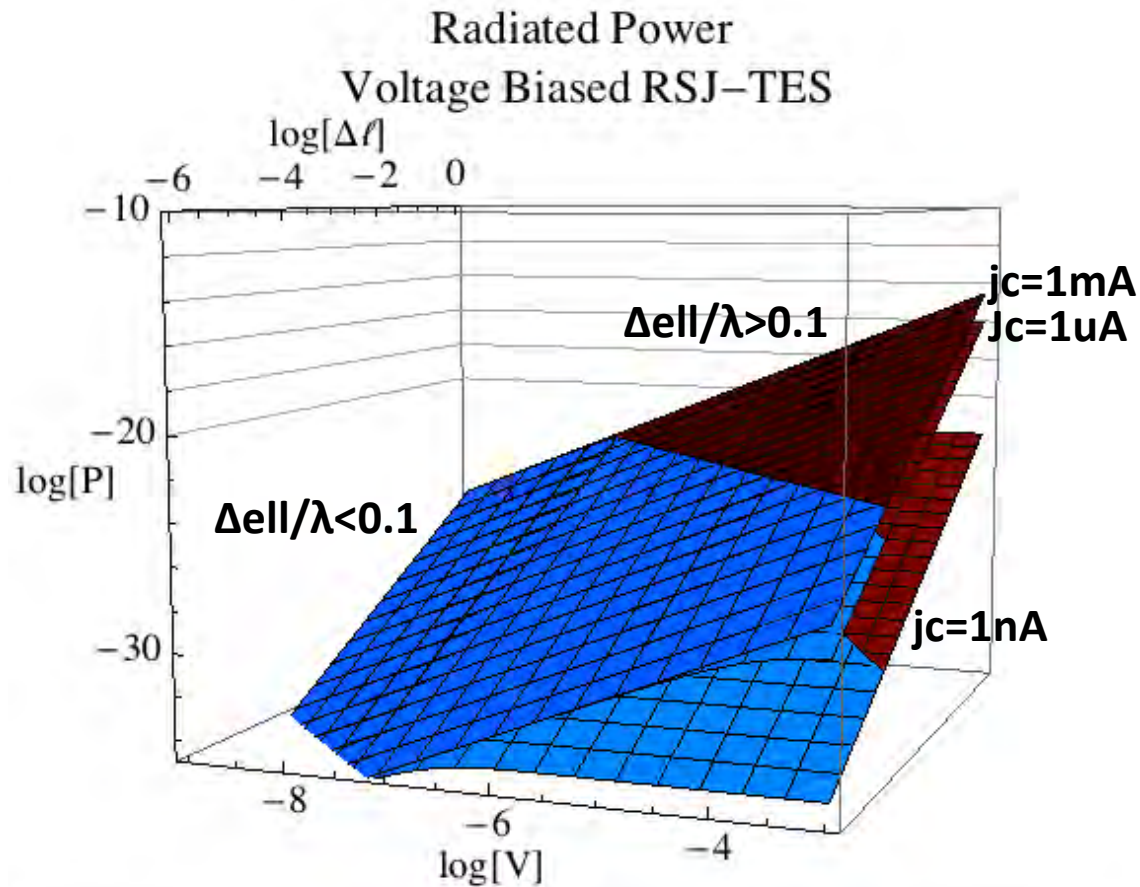
THE

END



$jj(tt)$





Power for relationship holds for:
 $\lambda \gg$ antenna length

Blue: $\Delta ell/\lambda < 0.1$
 (satisfied)

Red: $\Delta ell/\lambda > 0.1$
 Actual radiated power is larger than red surface

Rwsh	$\frac{1}{100}$
log [L]	-8
log [jc] {min,max,step}	{-9, -3, 3}

approximating the waveform as a pure sinusoid for the purposes of this