

# Internal thermal fluctuation noise in Mo/Au TES's

Nick Wakeham

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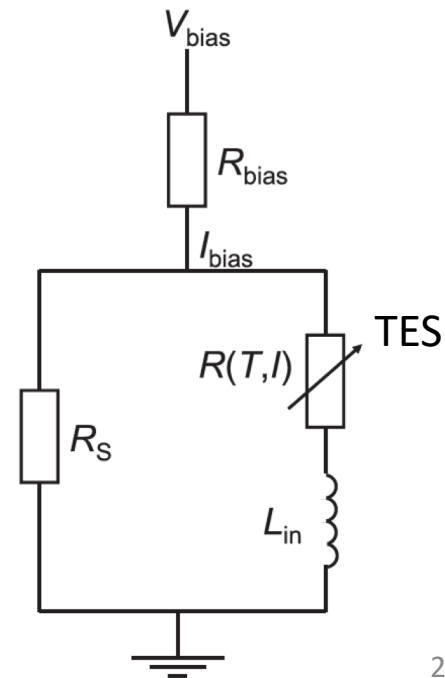
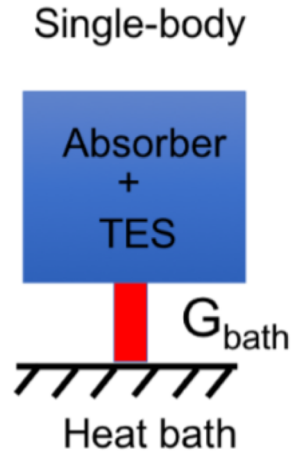
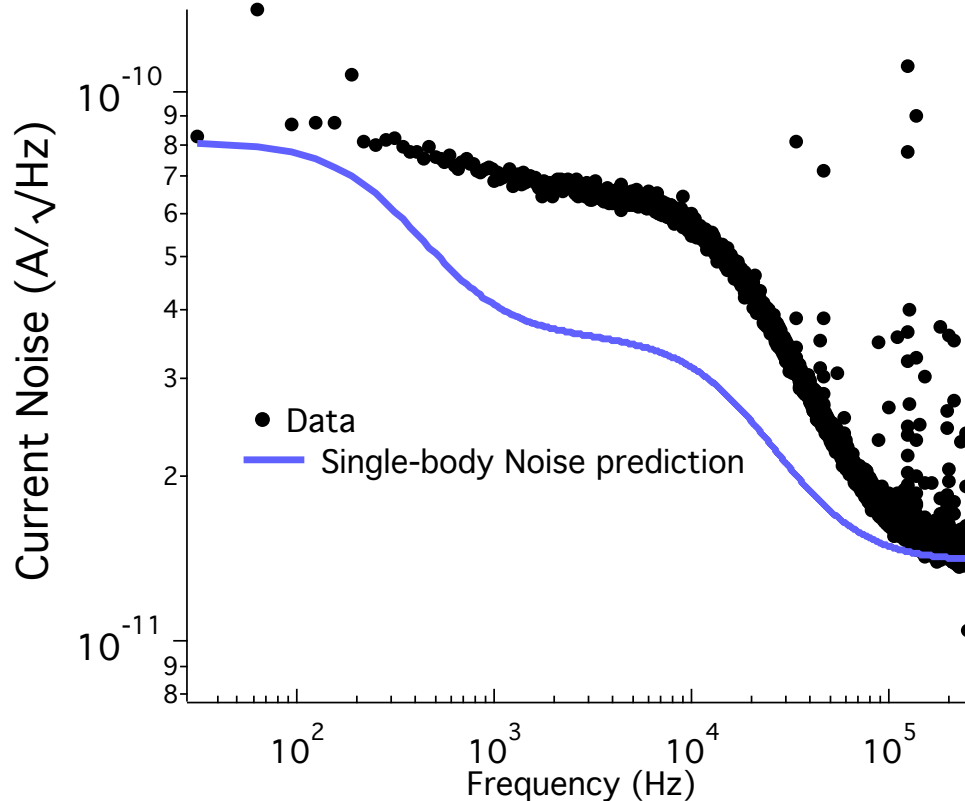
Joseph Adams, Simon Bandler, Sophie Beaumont, James Chervenak, Aaron Datesman, Megan Eckart, Fred Finkbeiner, Ruslan Hummatov, R. Kelley, Caroline Kilbourne, Antoine Miniussi, F. Porter, John Sadleir, Kazuhiro Sakai, Stephen Smith, Edward Wassell



# Noise in excess of single-body model

Simplest model of TES is a single body connected to bath by thermal conductance  $G_{\text{bath}}$

Measured noise is in excess of single body prediction

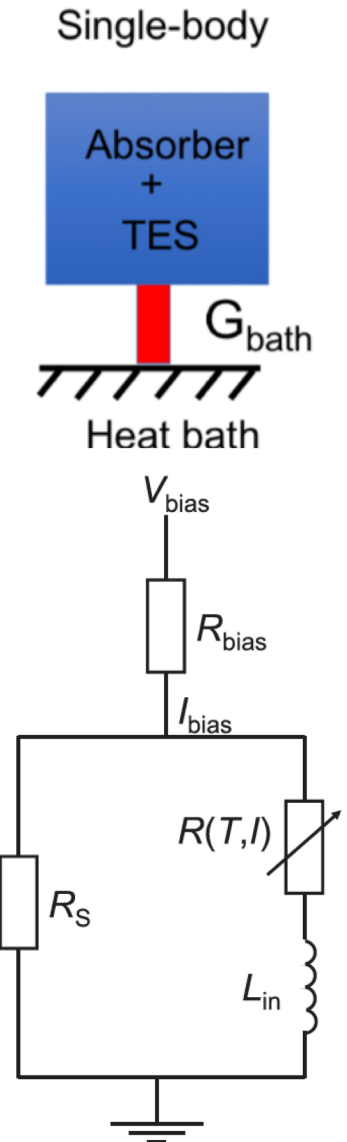
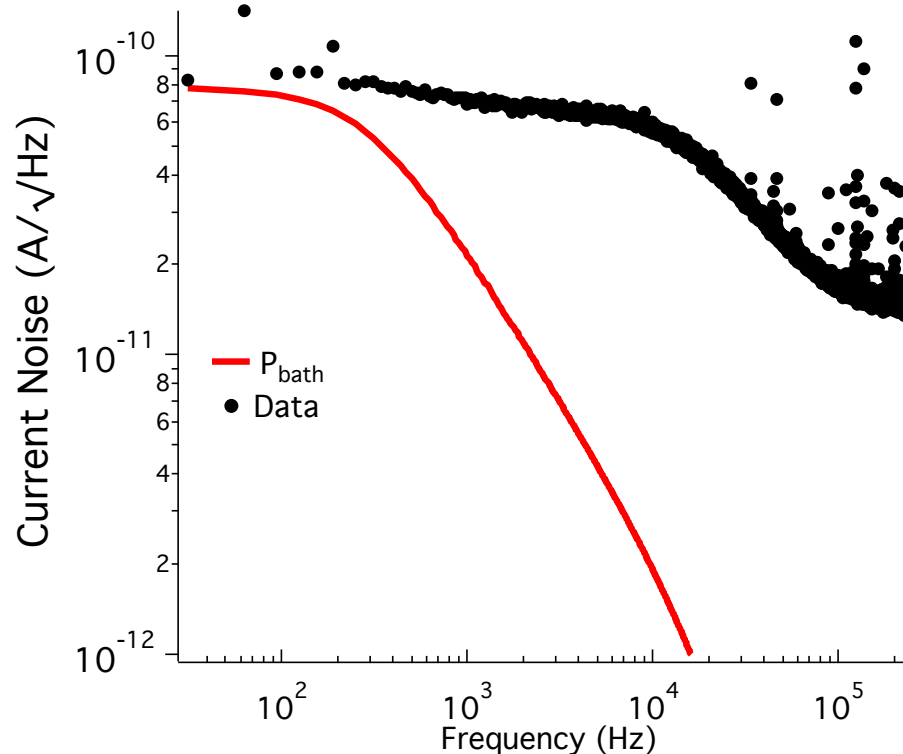


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3 contributions to current noise:

1. Thermal fluctuation noise between single body and bath -  $P_{\text{bath}}$

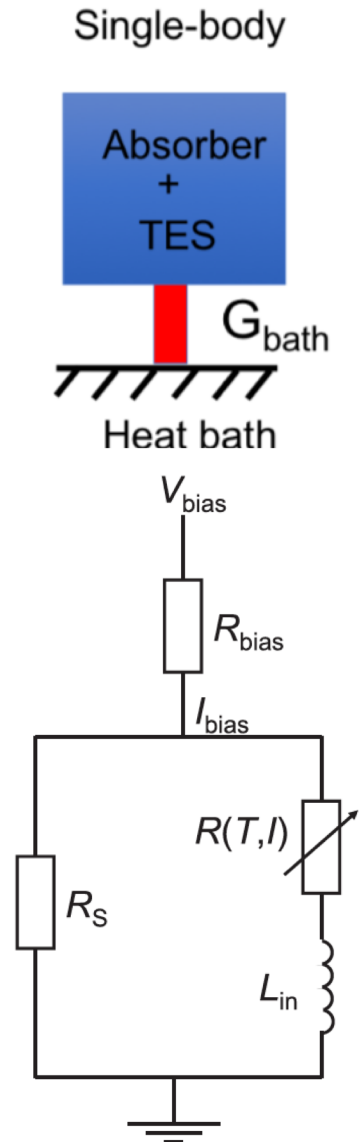
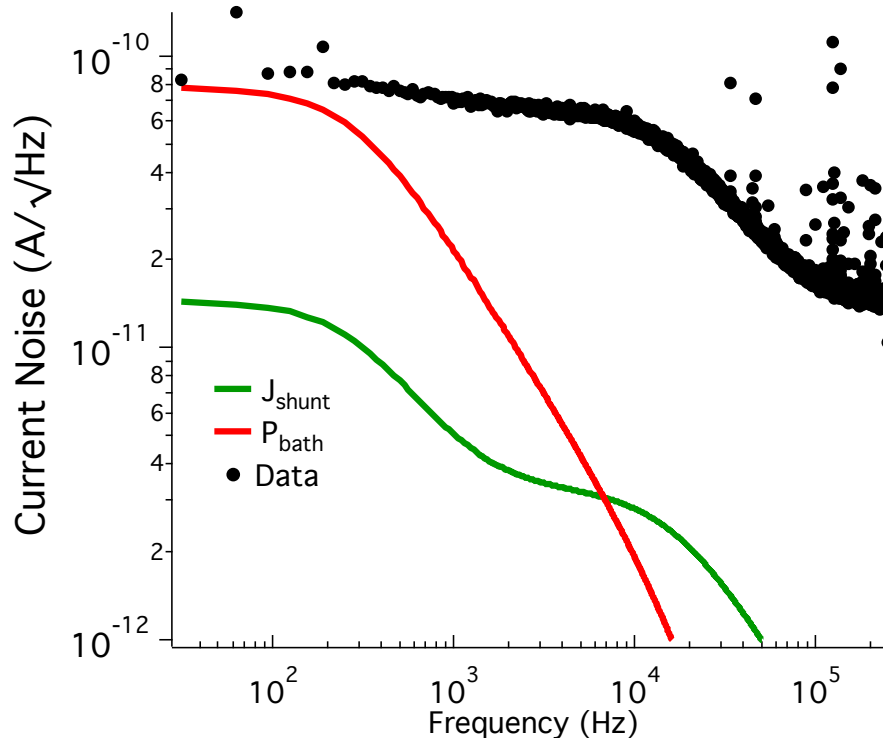


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Simplest model of TES is a single body connected to bath by thermal conductance  $G_{\text{bath}}$

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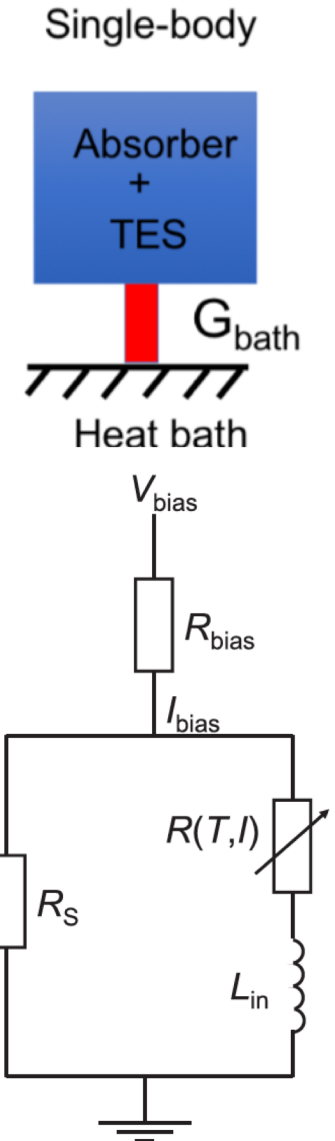
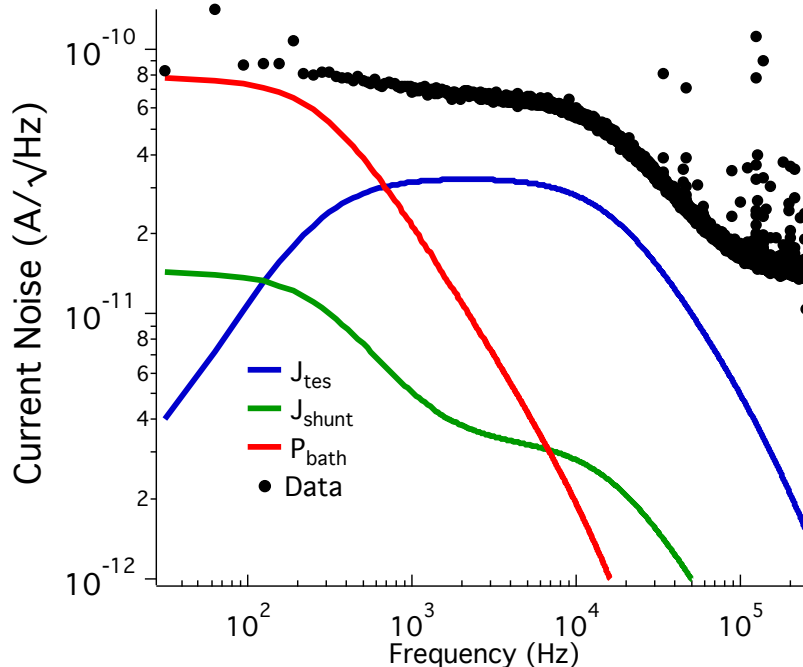
# Noise in excess of single-body model

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3 contributions to current noise:

1. Thermal fluctuation noise between single body and bath
2. Johnson noise of shunt resistor
3. (non-equilibrium) Johnson noise of TES  $J_{\text{tes}}$

K. D. Irwin Nuclear Instruments and Methods in Physics Research A 559 (2006) 718–720



# Noise in excess of single-body model

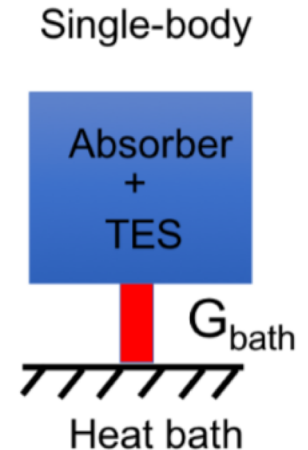
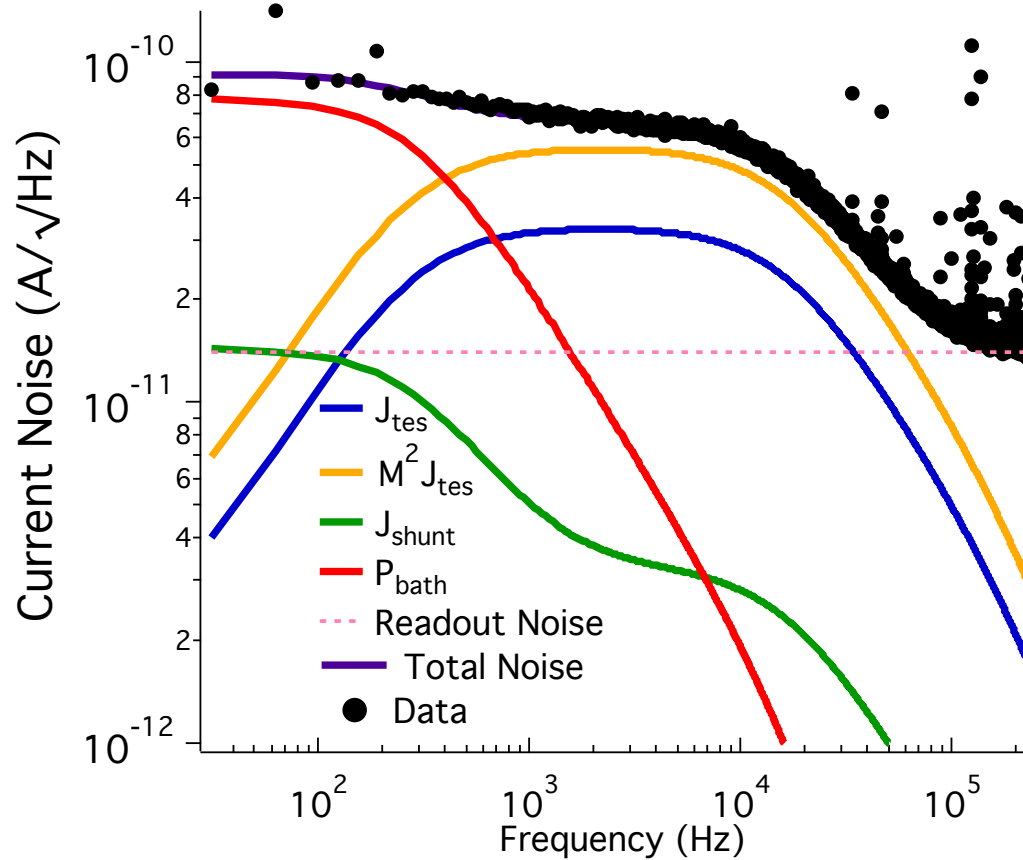
Can fit to measured data by adding another term  $M^2 J_{tes}$

*Smith et al. J. Appl. Phys. 114, 074513 (2013);*

$$V_n = \sqrt{(4K_B TR)(1 + 2\beta)(1 + M^2)}$$

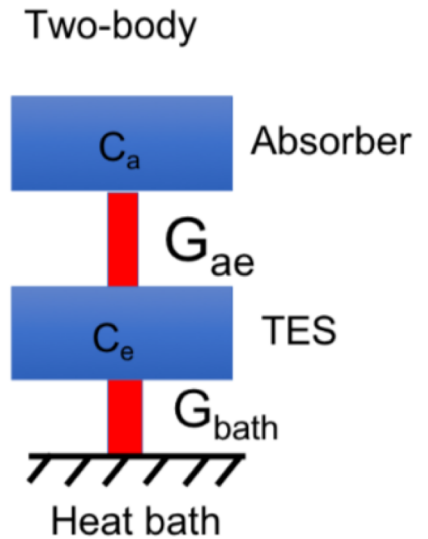
$M^2$  then parameterizes the magnitude of the excess noise in the device.

But what is the origin of this excess noise term?



# Two-body model

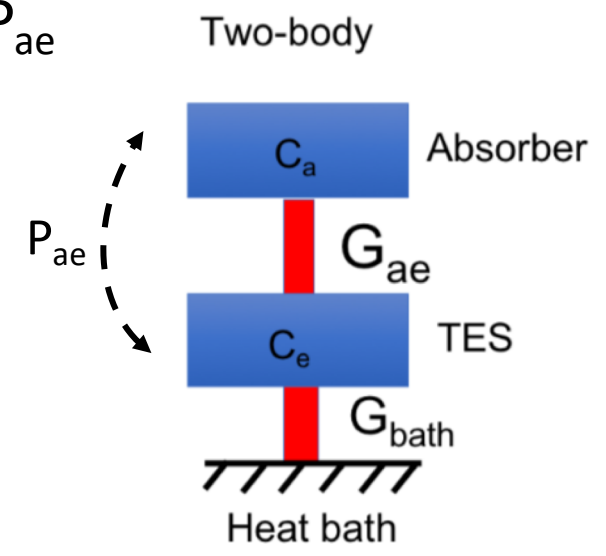
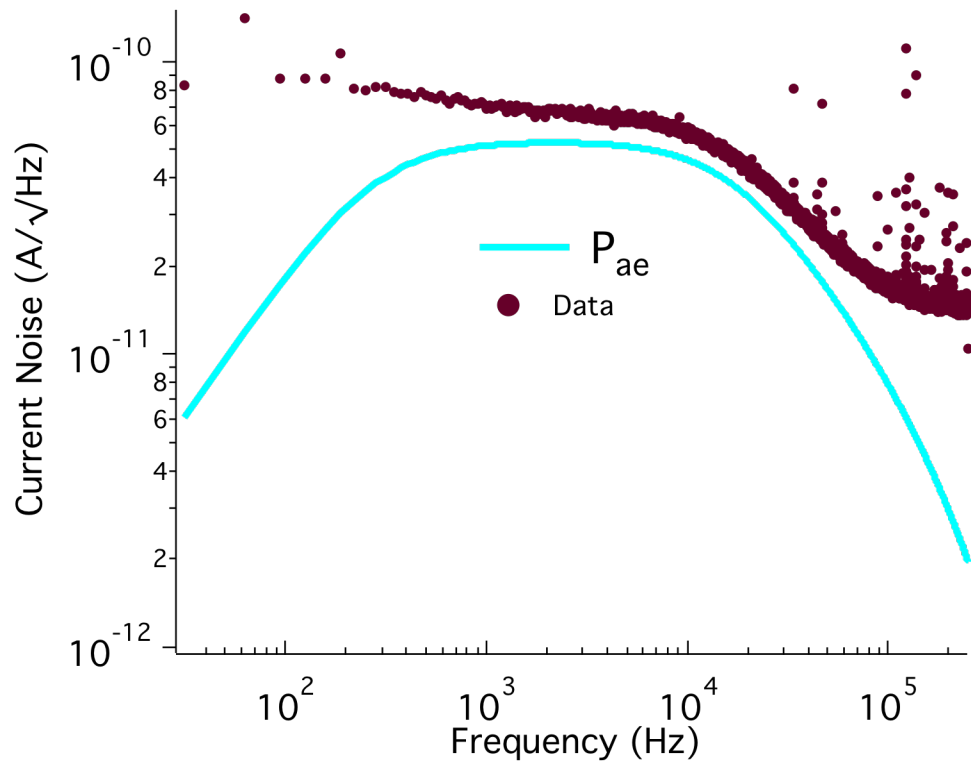
In two-body model, absorber and TES (for example) are separated by thermal conductance  $G_{ae}$



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- Internal thermal fluctuation noise between the two bodies. -  $P_{ae}$





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No new theory here. It has been described by many others...

*M. A. Lindeman, Ph.D. thesis, UC Davis, 2000;*

*Lindeman et al. Review of Scientific Instruments 75, 1283 (2004)*

*I. J. Maasilta AIP Advances 2, 042110 (2012);*

*E Figueroa-Feliciano et al. Journal of Applied Physics 99, 114513 (2006);*

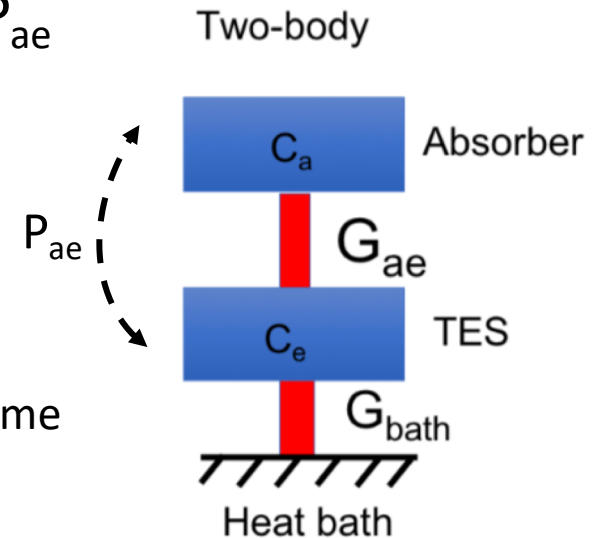
*.....et al.*

This additional noise has also been found to be significant in some cases

*H. F. C. Hoevers et al. Appl. Phys. Lett. 77, 4422 (2000);*

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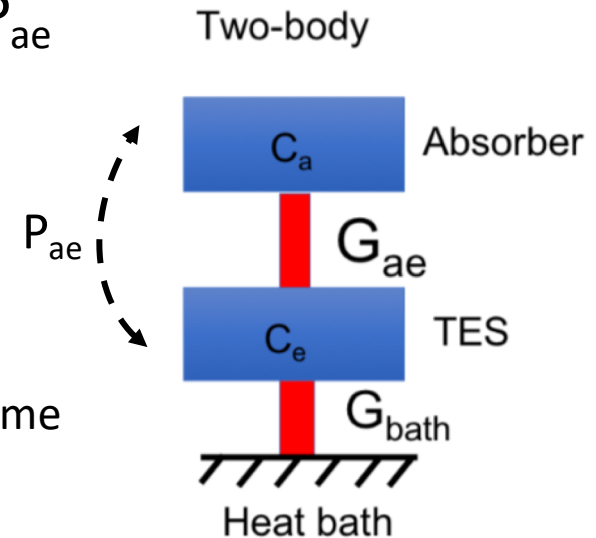
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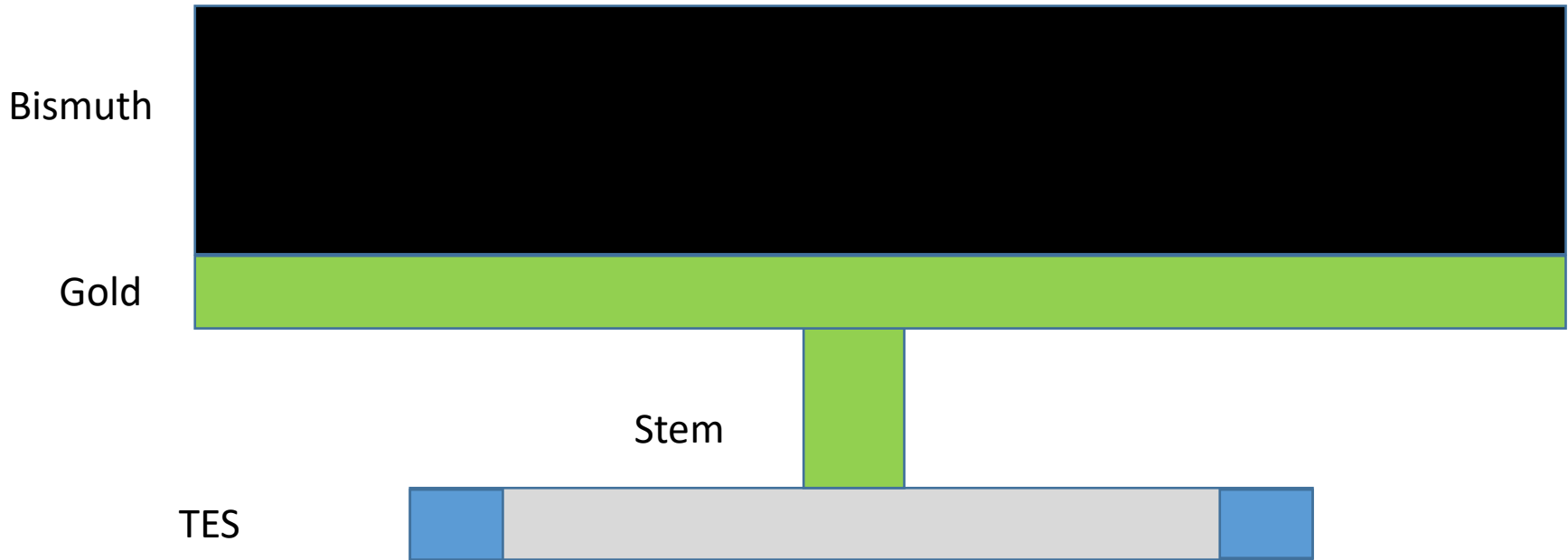
*.....et al.*



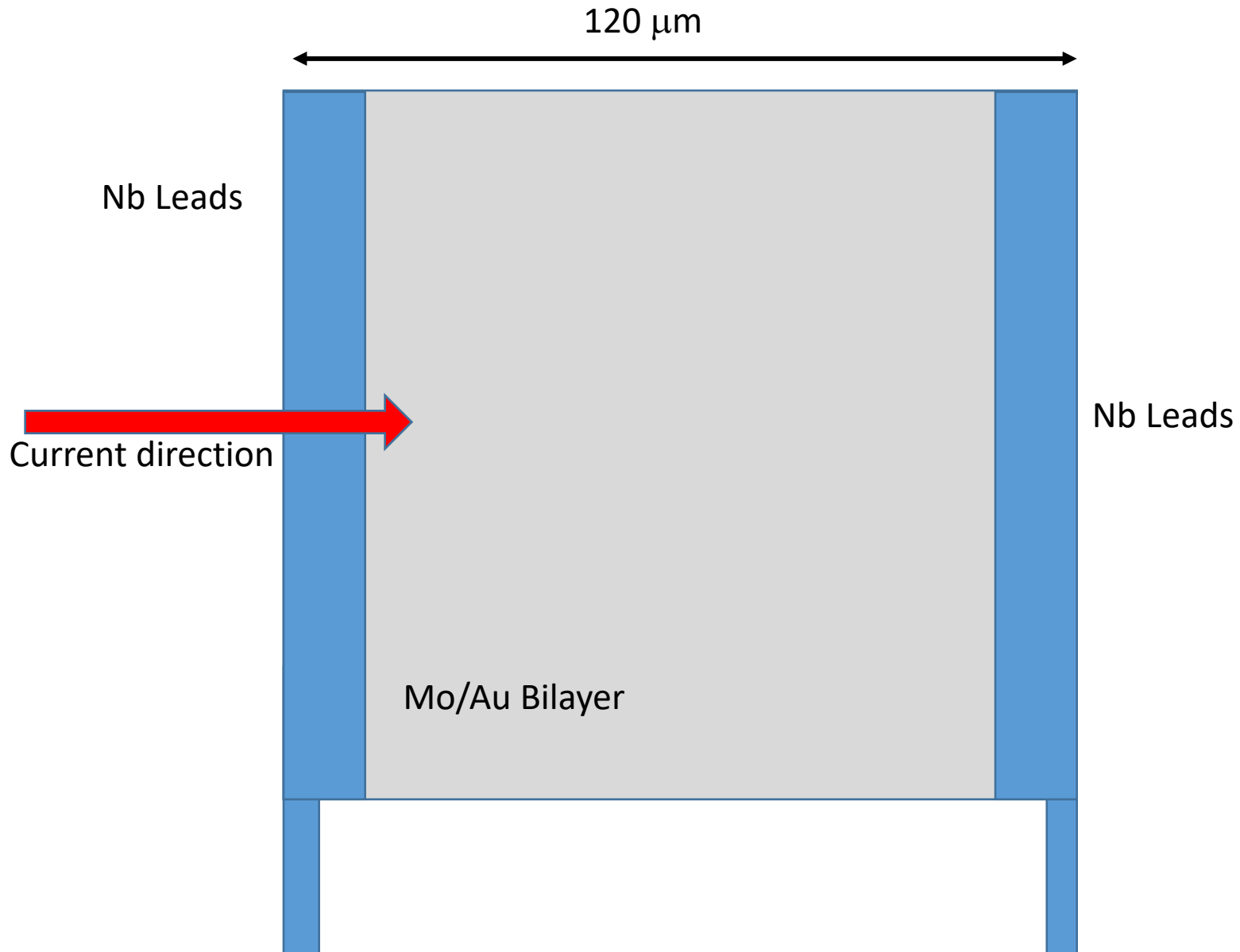
**Central question:**

**In latest NASA Mo/Au bilayer TESs is  
internal thermal fluctuation noise  
significant?**

# TES design - Side view

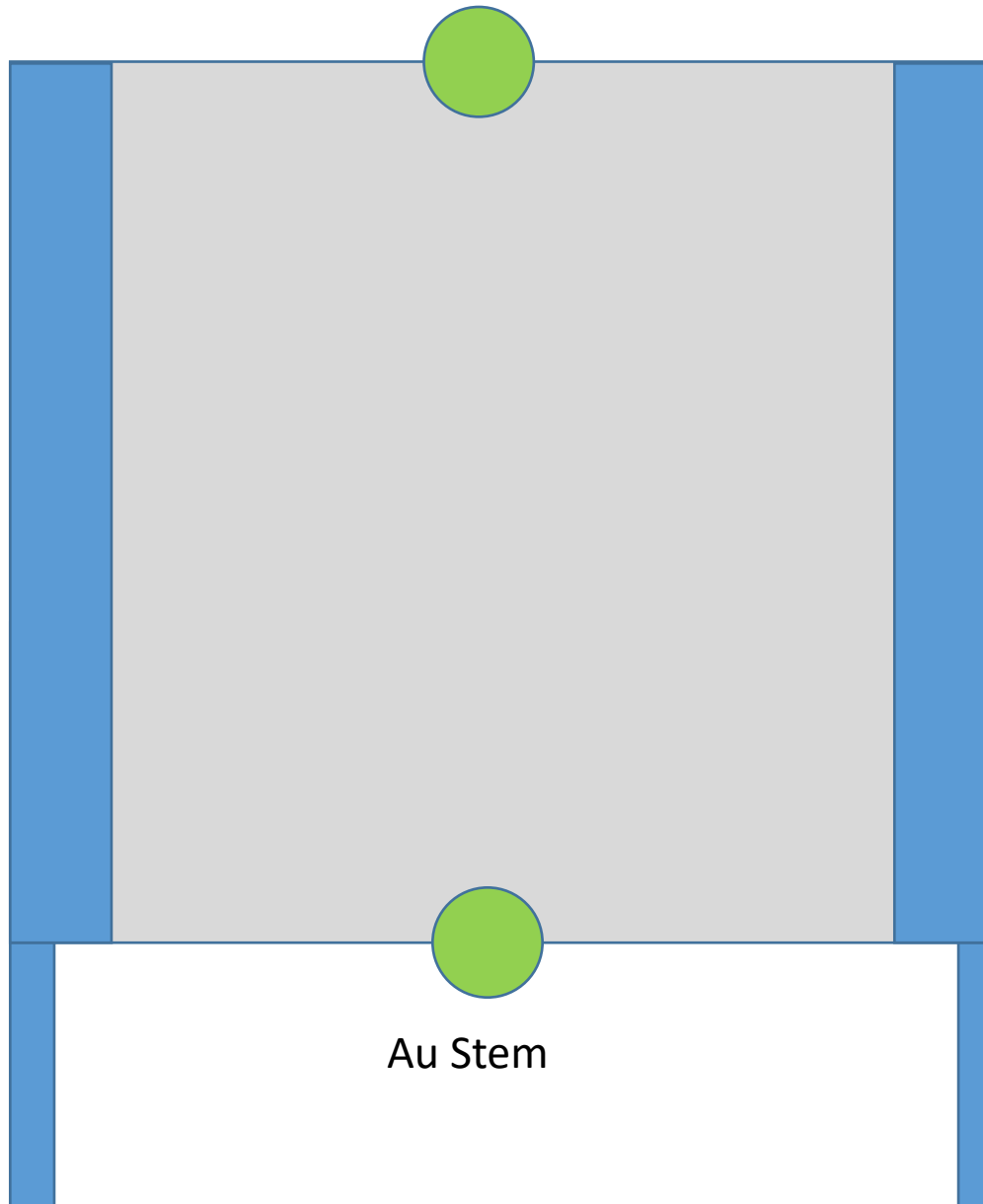


# TES design – Top View



# Top view

Absorber stems are pillars



We have measured complex impedance and noise



Extract transition properties and estimate noise contributions

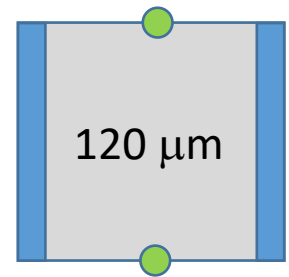
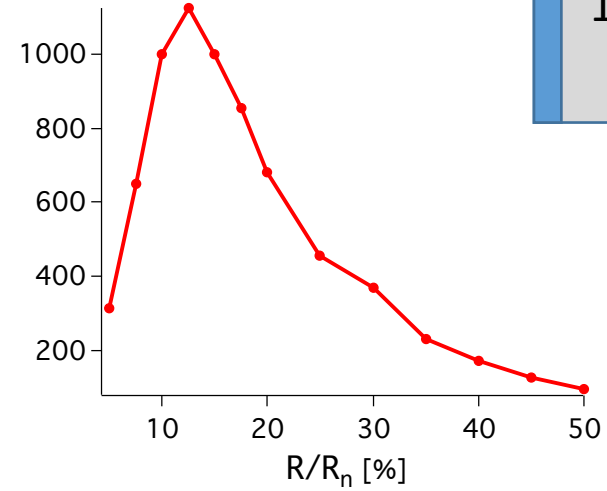
# Single-body



Complex Impedance

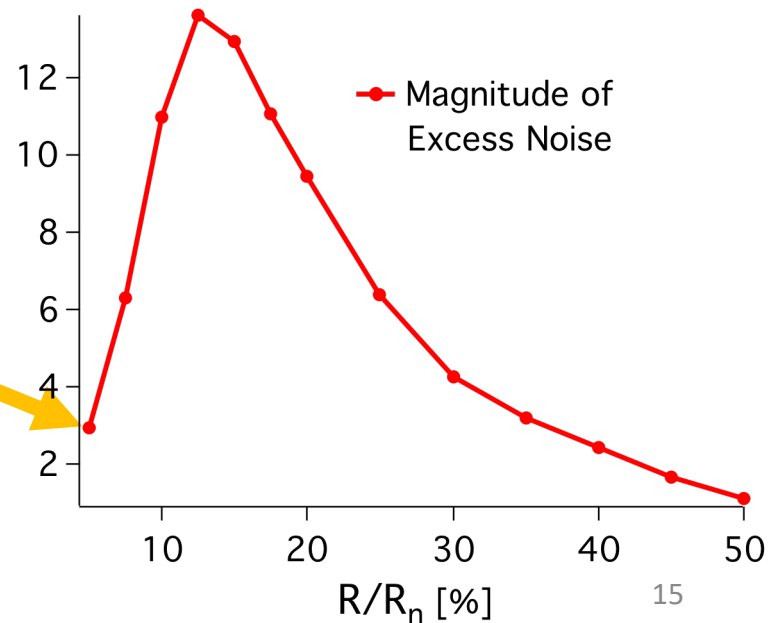
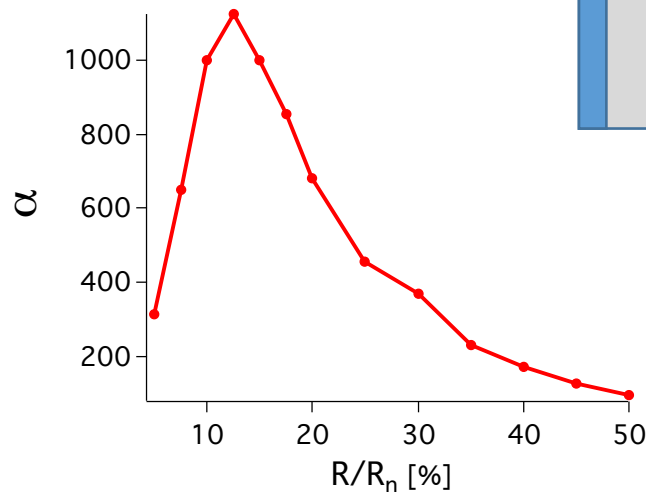
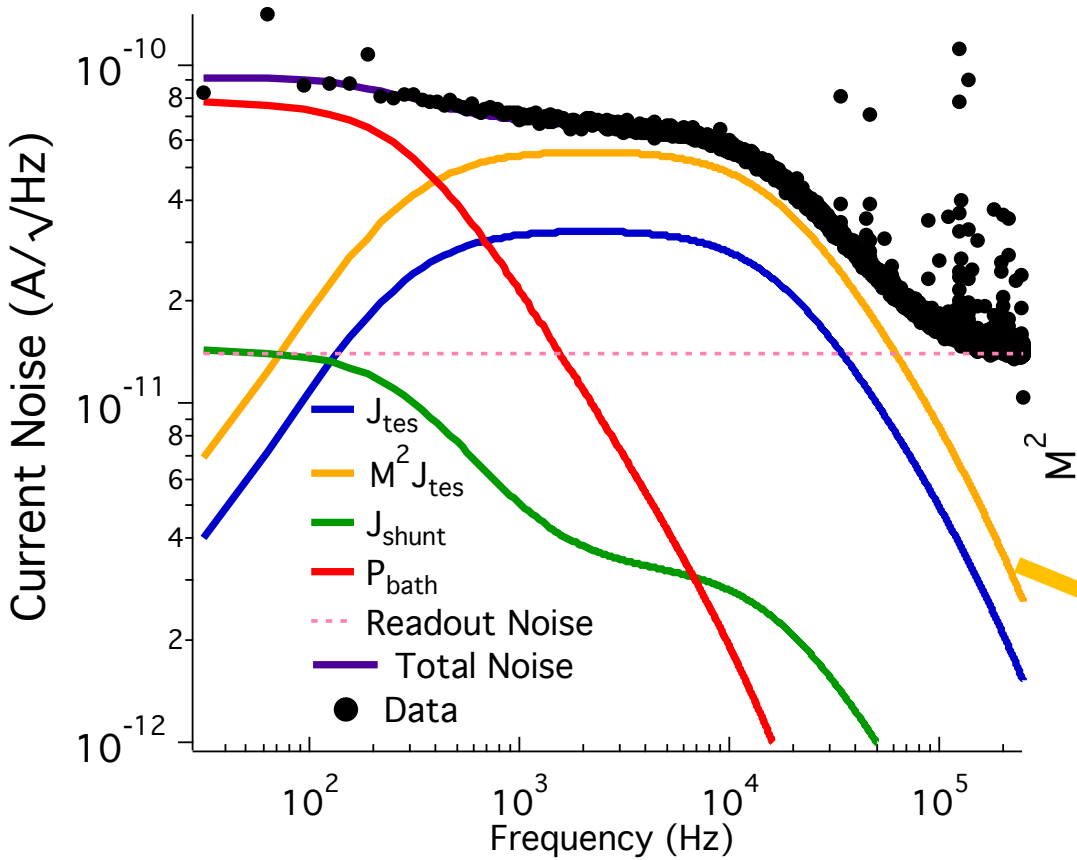
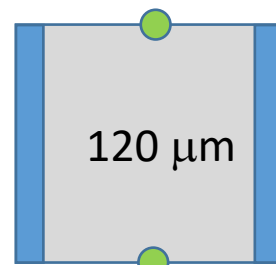
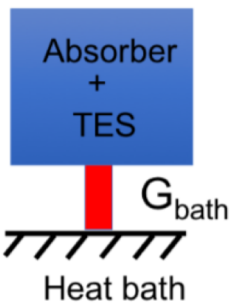
$\alpha = T/R \, dR/dT$

$\alpha$

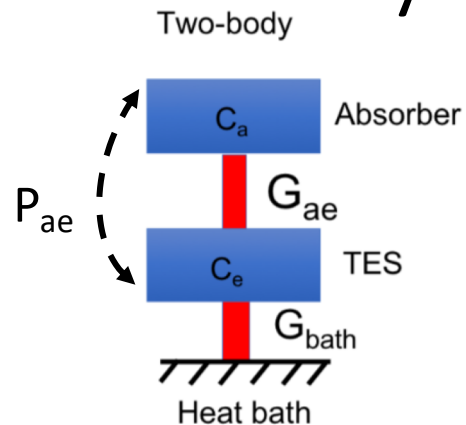


# Single-body

Single-body



# Two-body

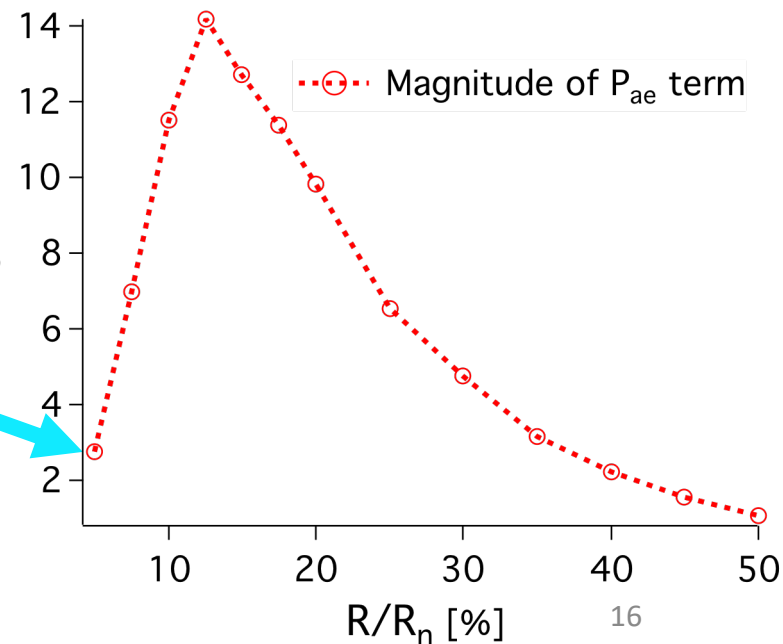
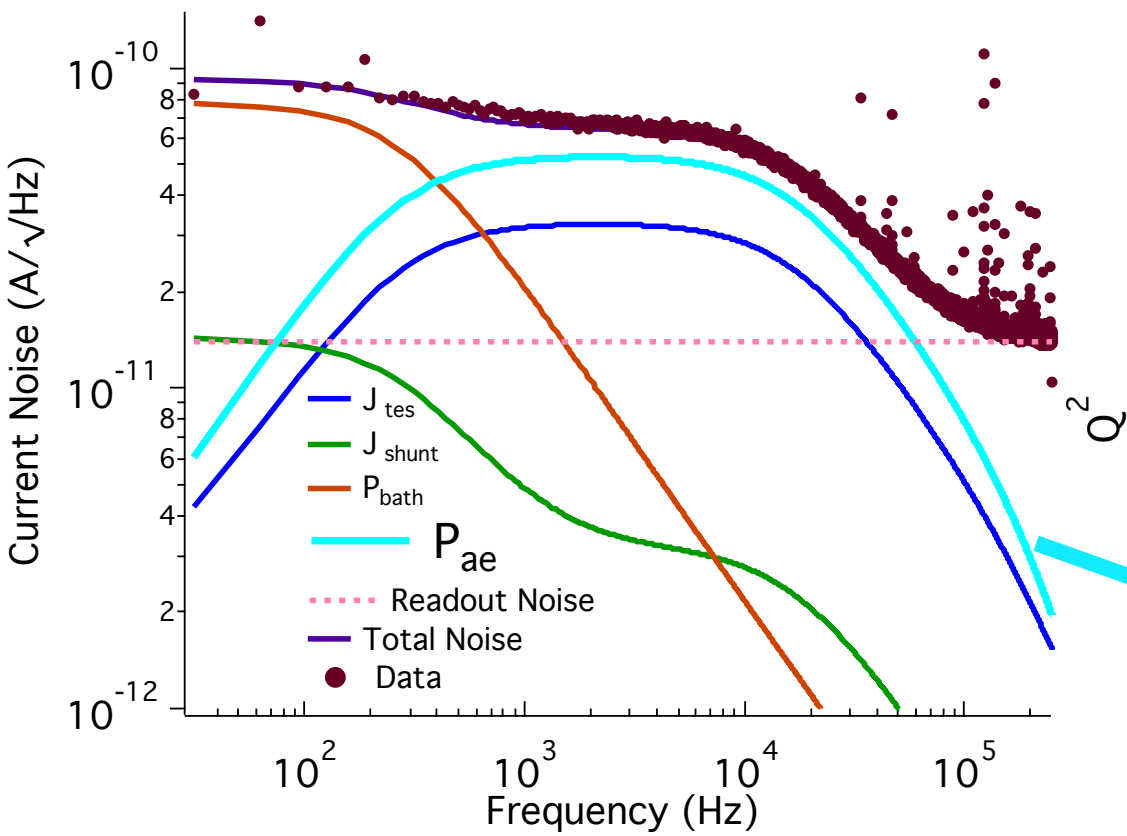


Attribute excess noise to internal thermal fluctuation noise -  $P_{ae}$

No need for additional Johnson noise

For convenience: Parameterize magnitude of  $P_{ae}$  term relative to the Johnson noise.

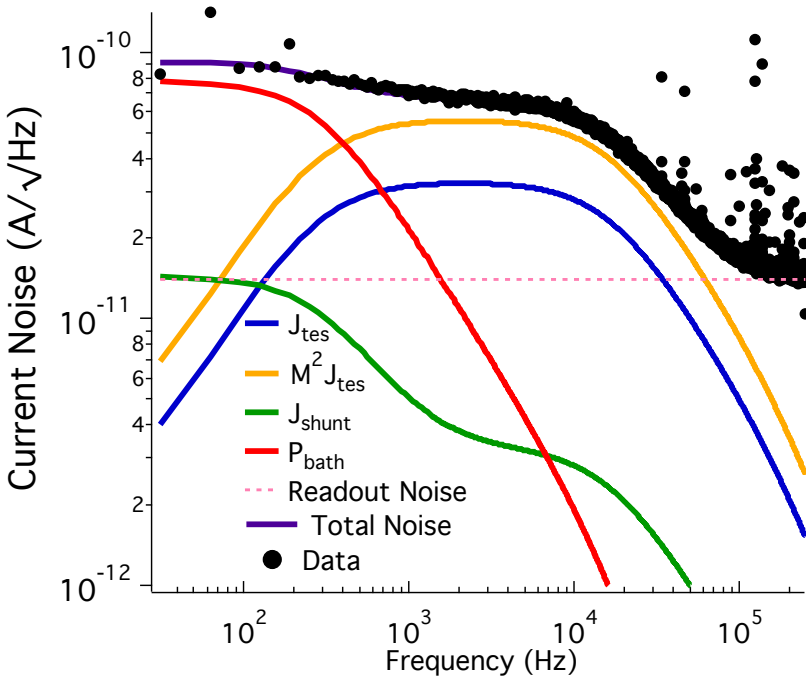
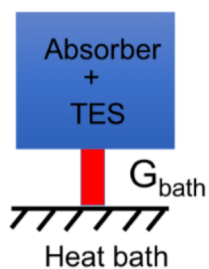
$$P_{ae} \approx Q^2 J_{tes}$$



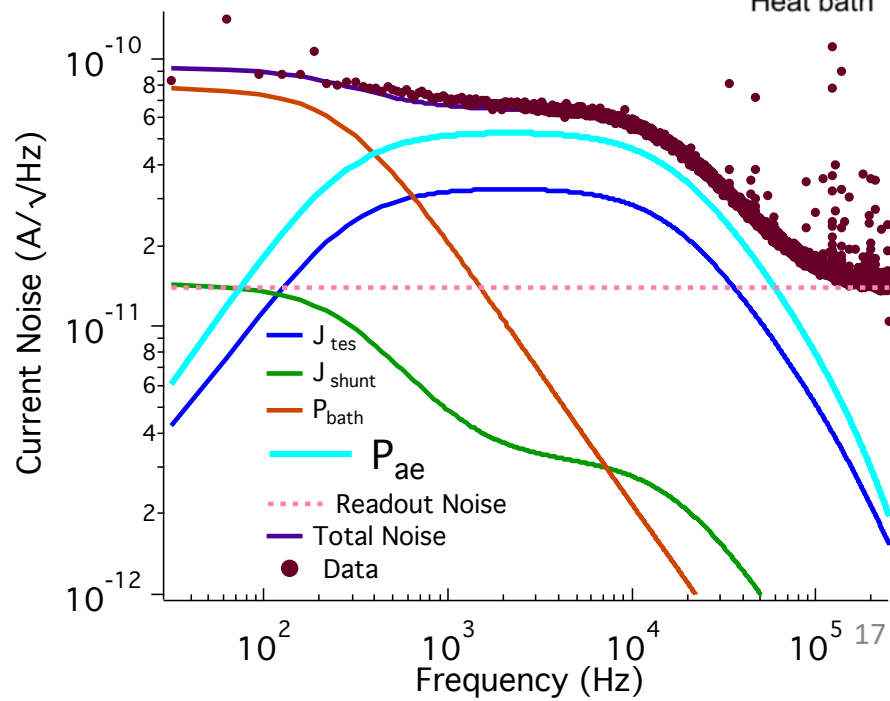
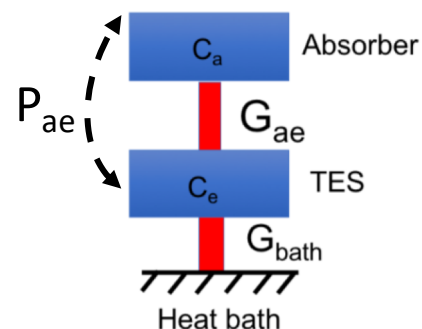


# Single-body Vs Two-body

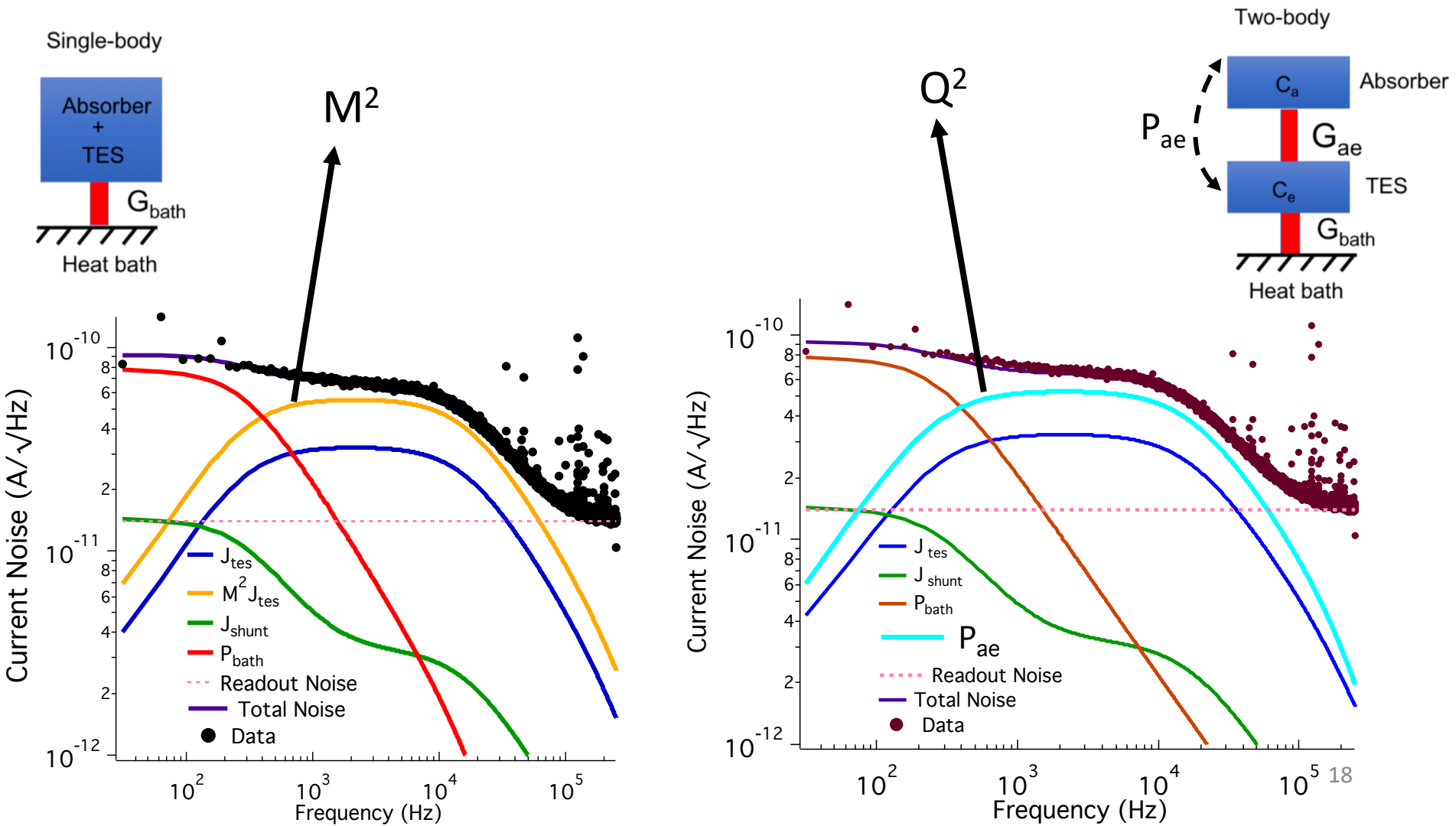
Single-body



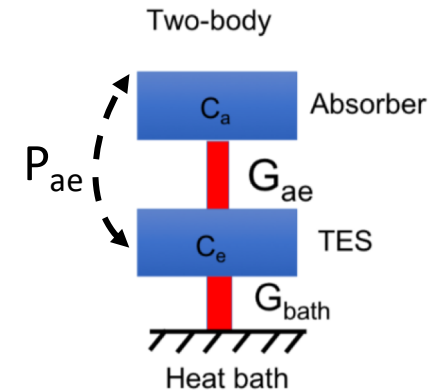
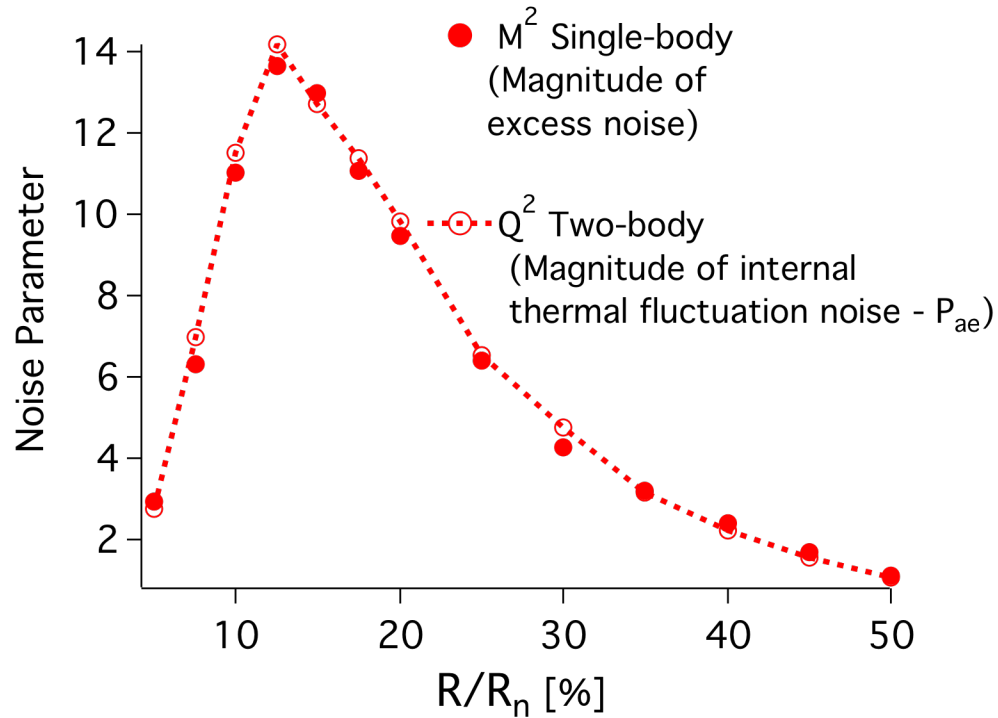
Two-body



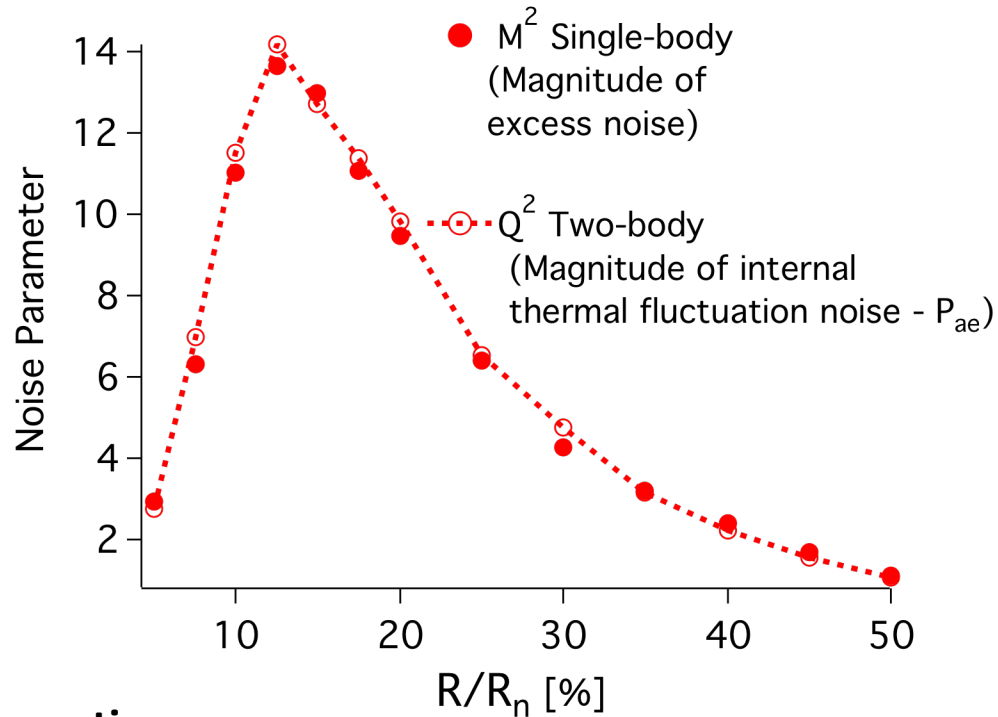
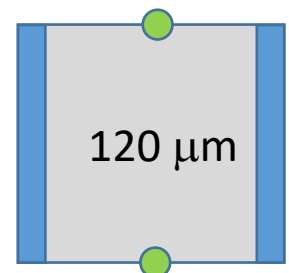
# Single-body Vs Two-body



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# Single-body Vs Two-body



## Two-body assumptions:

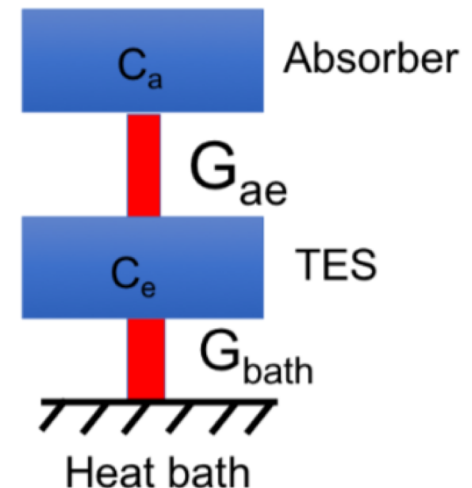
$C_e$  = BCS Predicted C of TES including jump from superconductivity (+ membrane). **Relatively insensitive.**

$C_a + C_e$  = measured heat capacity of device

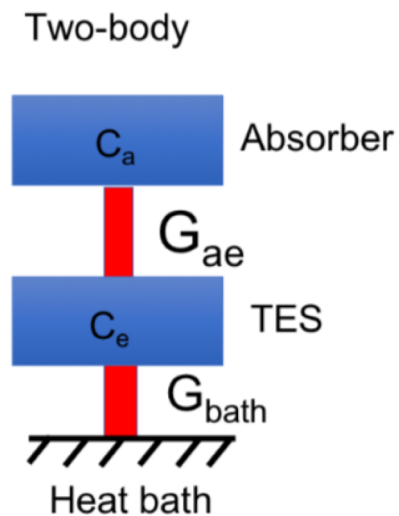
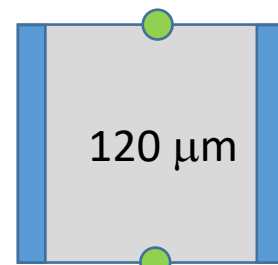
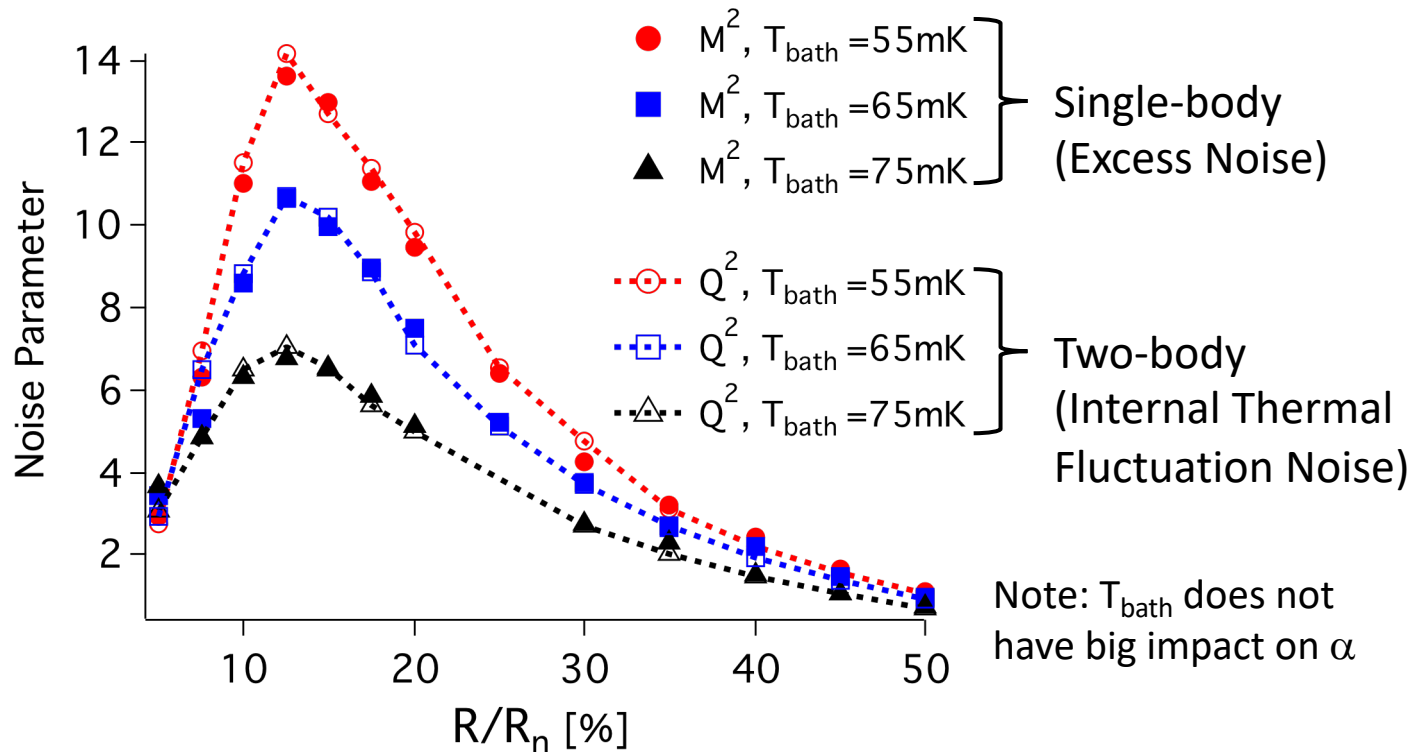
$G_{bath}$  is measured value

$G_{ae}$  = constant in  $R/R_n$

## Two-body



# Variation with $T_{\text{bath}}$



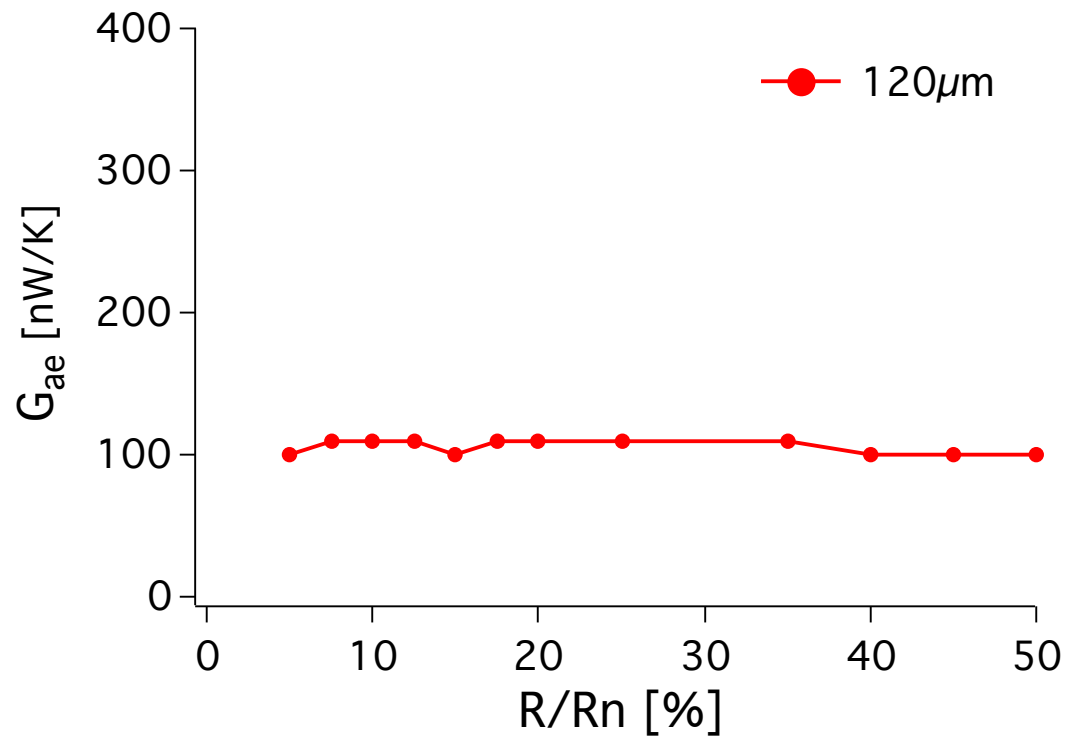
$G_{ae}$  is a constant for all points on graph.

Measured excess noise well described by two-body model with fixed parameters over wide range of  $T_{\text{bath}}$  and  $R/R_n$ .

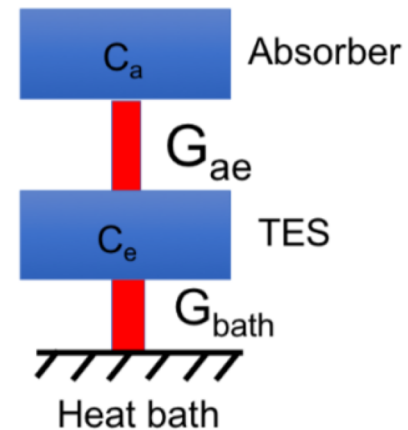
# Variable $G_{ae}$

Allow  $G_{ae}$  to float in fitting of two-body model.

Largely independent of bias point



Two-body

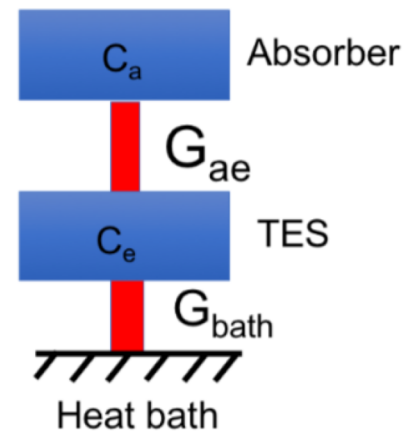
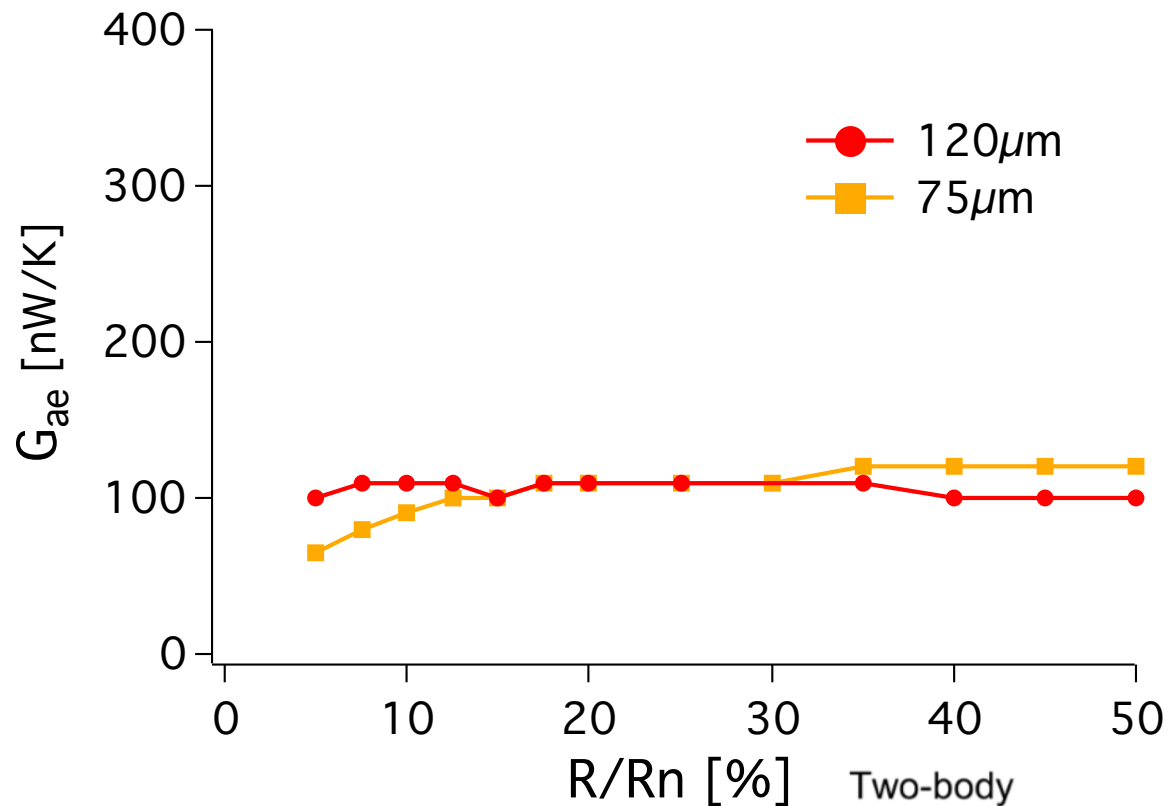


# Variable $G_{ae}$

Allow  $G_{ae}$  to float in fitting.

Can then do same analysis on other devices

**75 $\mu$ m TES**

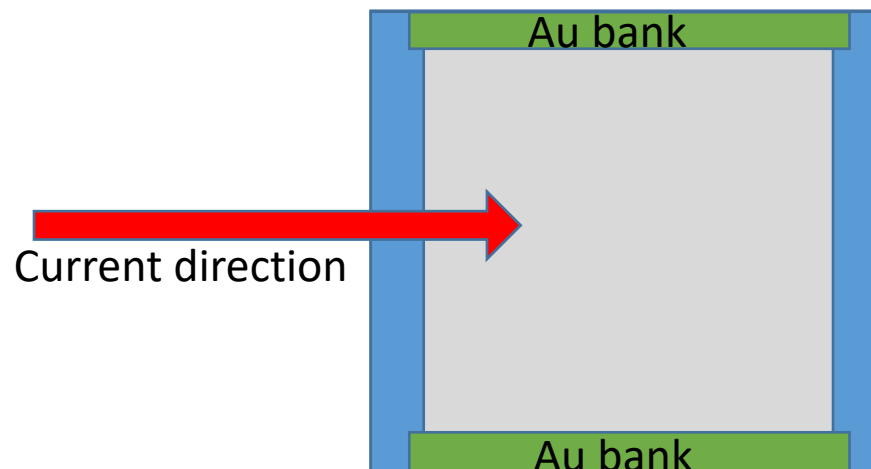
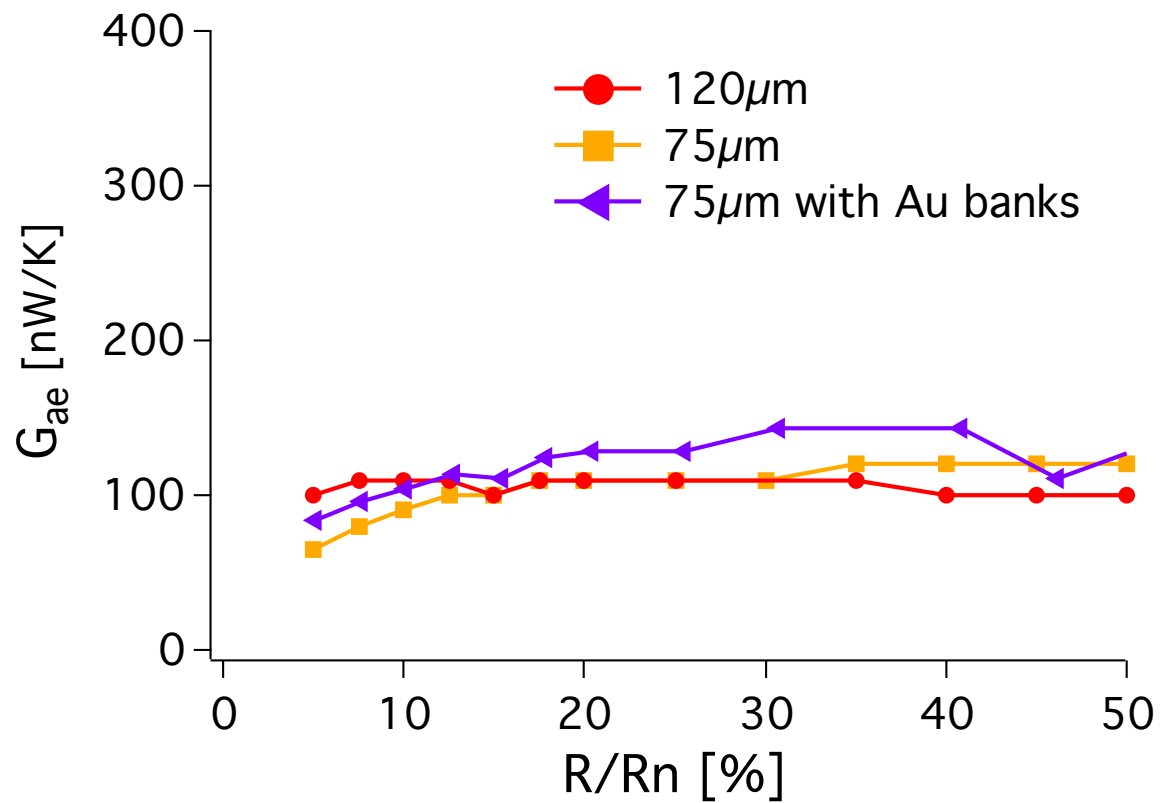


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## 75 $\mu\text{m}$ TES with Au banks



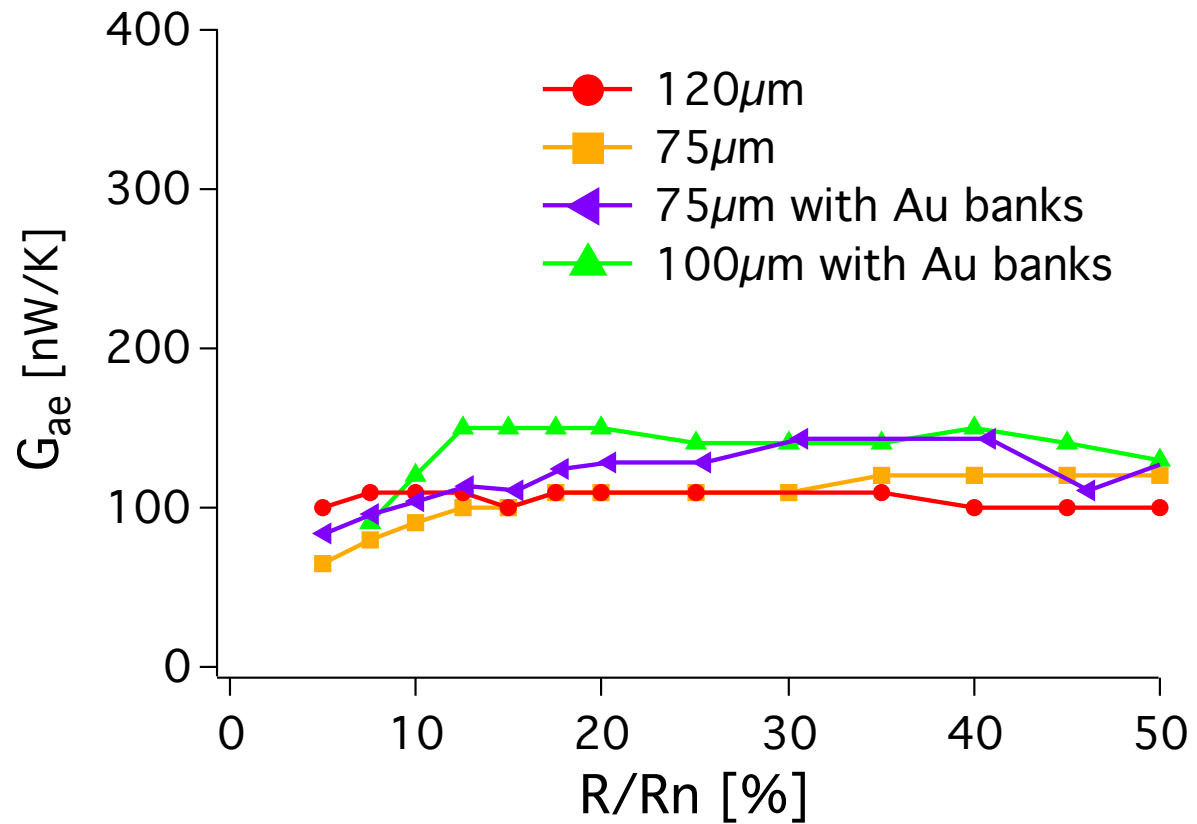


# Variable $G_{ae}$

Allow  $G_{ae}$  to float in fitting.

Can then do same analysis on other devices

**100 $\mu\text{m}$  TES with Au banks**



➤ Little variation in fitted value  $G_{ae} \sim 100$  nW/K

# Variable $G_{ae}$

## What is origin of finite $G_{ae}$ ?

Estimated G from stem pillars

90  $\mu\text{W/K}$

→ Too large

Estimated G from electron-phonon interaction 2 nW/K

→ Too small

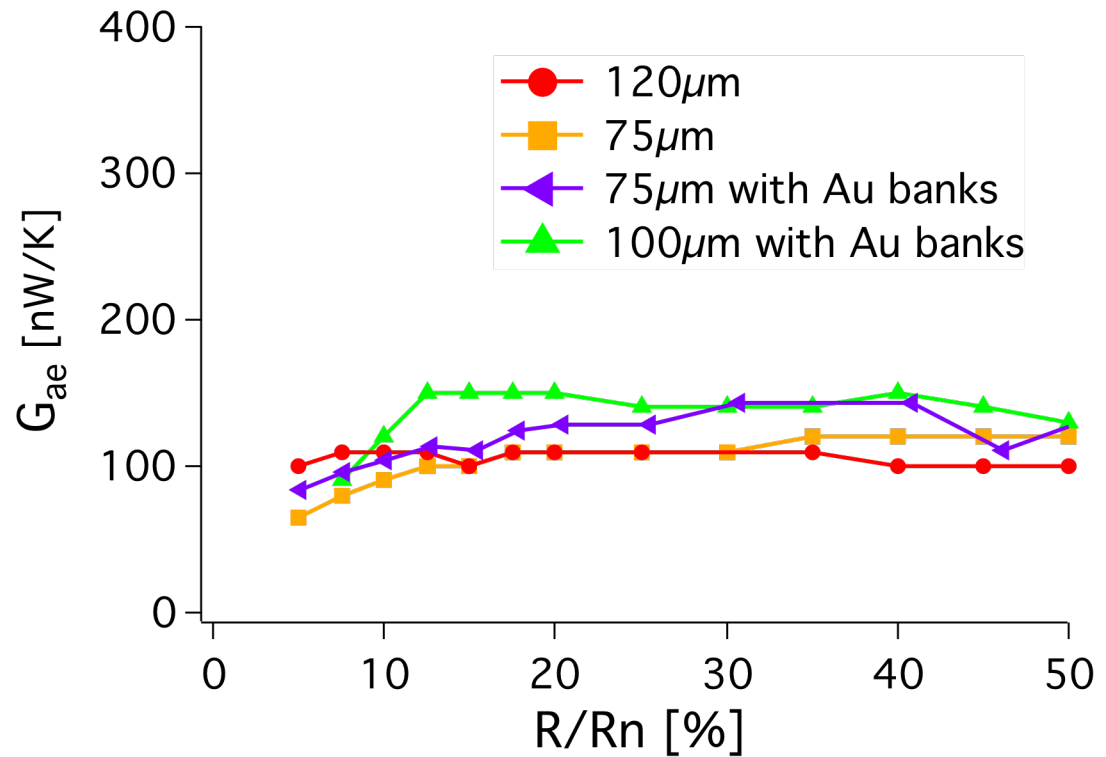
Estimated G of bilayer from Wiedemann-Franz law and measured sheet resistance

(50  $\text{m}\Omega/\square$ )

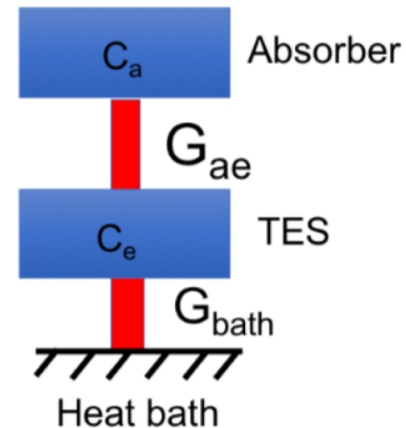
→ 50 nW/K.

**Suggests finite thermal conductance ( $G_{ae}$ ) is from finite resistance of bilayer itself**

Reported before in 200  $\text{m}\Omega$  Ti/Au bilayers



Two-body

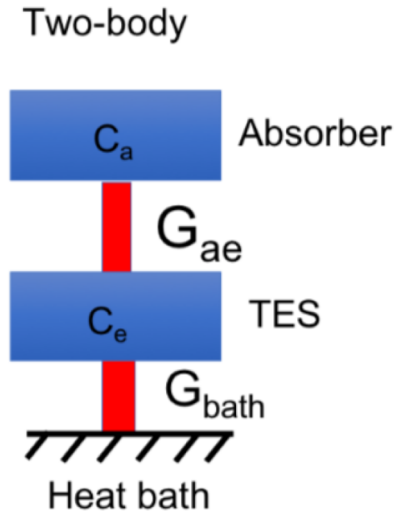


*H. F. C. Hoevers et al. Appl. Phys. Lett. 77, 4422 (2000);*

# Low R bilayer

If bilayer resistance is responsible for internal thermal fluctuation noise.

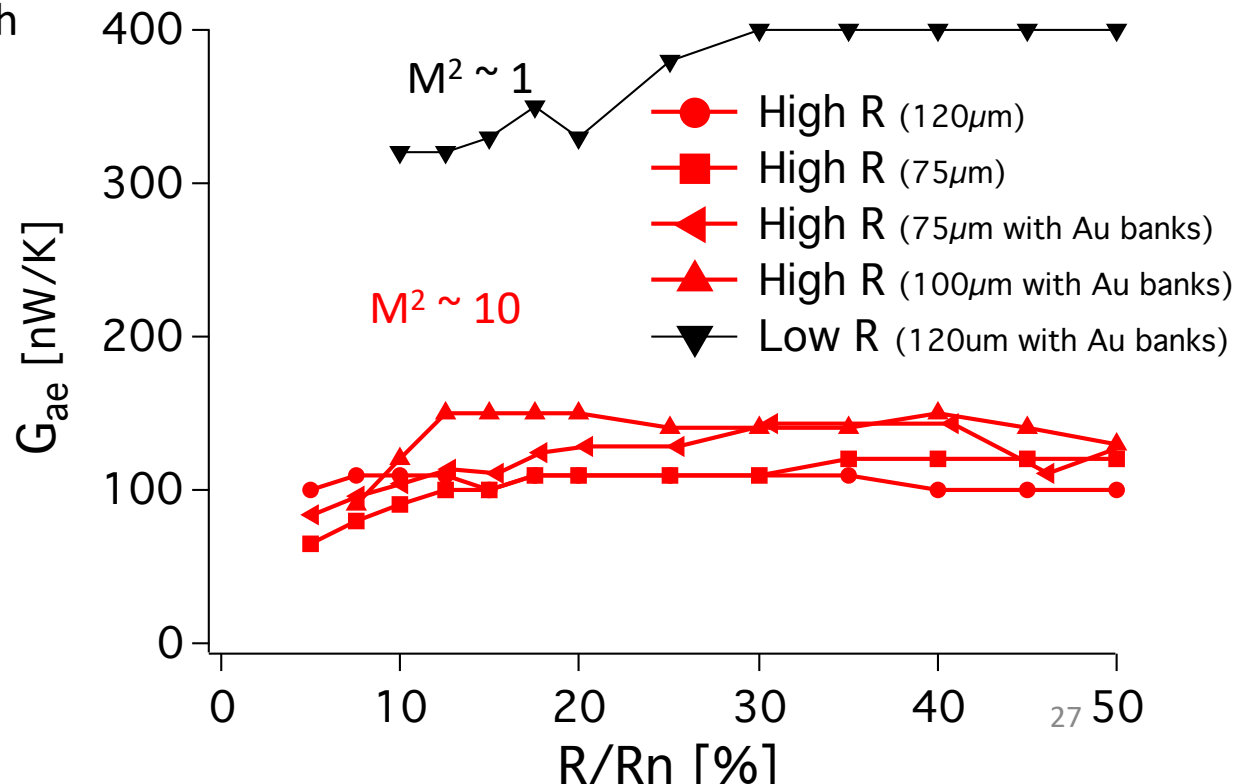
EXPECT: Lower  $R_{\square}$  bilayer  $\rightarrow$  Higher  $G_{ae}$



Measured and fitted devices with bilayer with factor  $\sim 4$  smaller  $R_{\square}$

Fitted  $G_{ae} \sim$  factor 4 larger

Clearly a crude estimate of thermal conductance of TES but captures coarse trends.



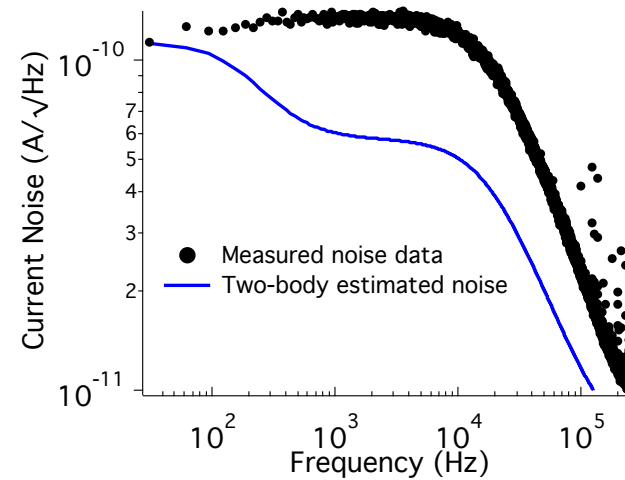
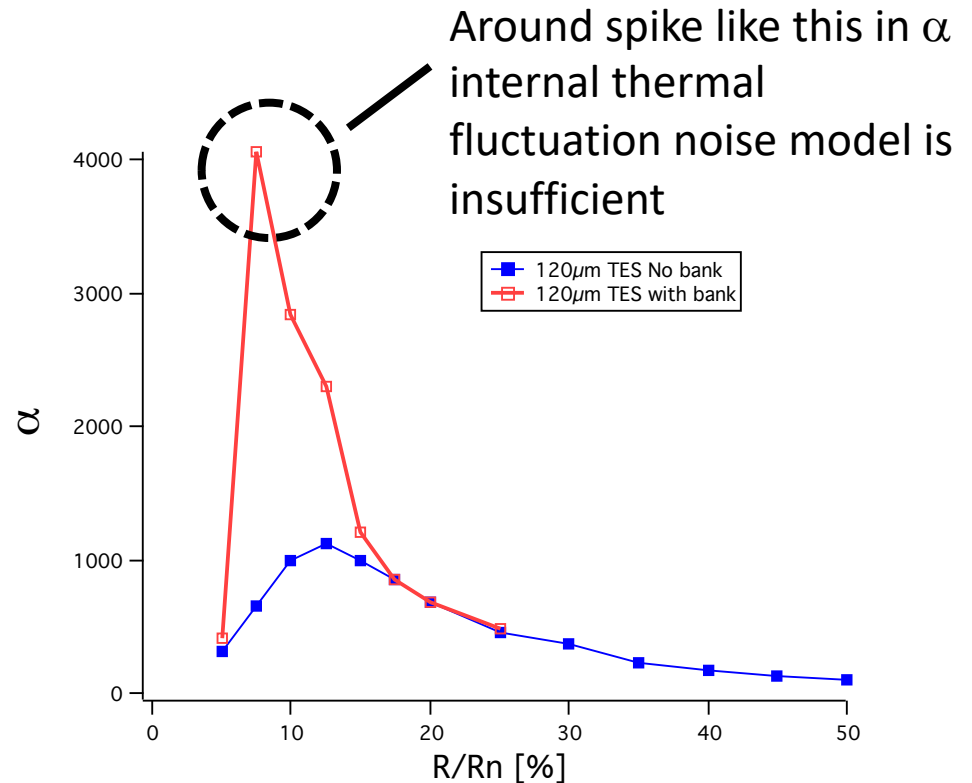
# Where two-body model doesn't work...

Cannot fit within our assumptions around kinks

- Regions with rapid changes in  $\alpha$ .
- For example in this 120 $\mu\text{m}$  device with banks.

No stripes  $\rightarrow$  smoother transitions.

Perhaps in other devices small transition features may have given additional noise largely not present in our small no-stripe devices



# Conclusion

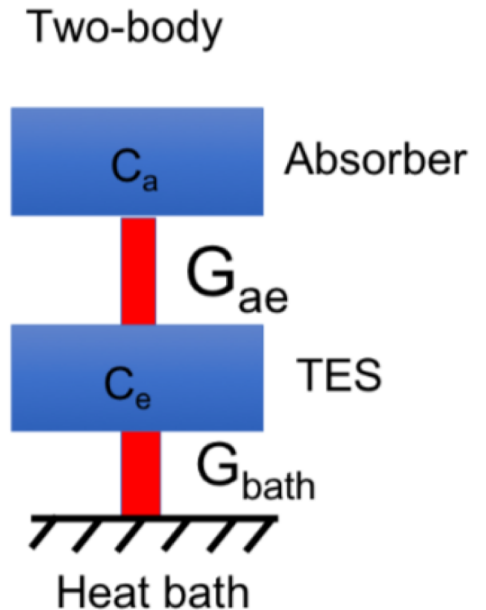
Fit measured noise in our “no-stripe” devices with a two-body model.

Internal thermal fluctuation noise appears to dominate excess noise

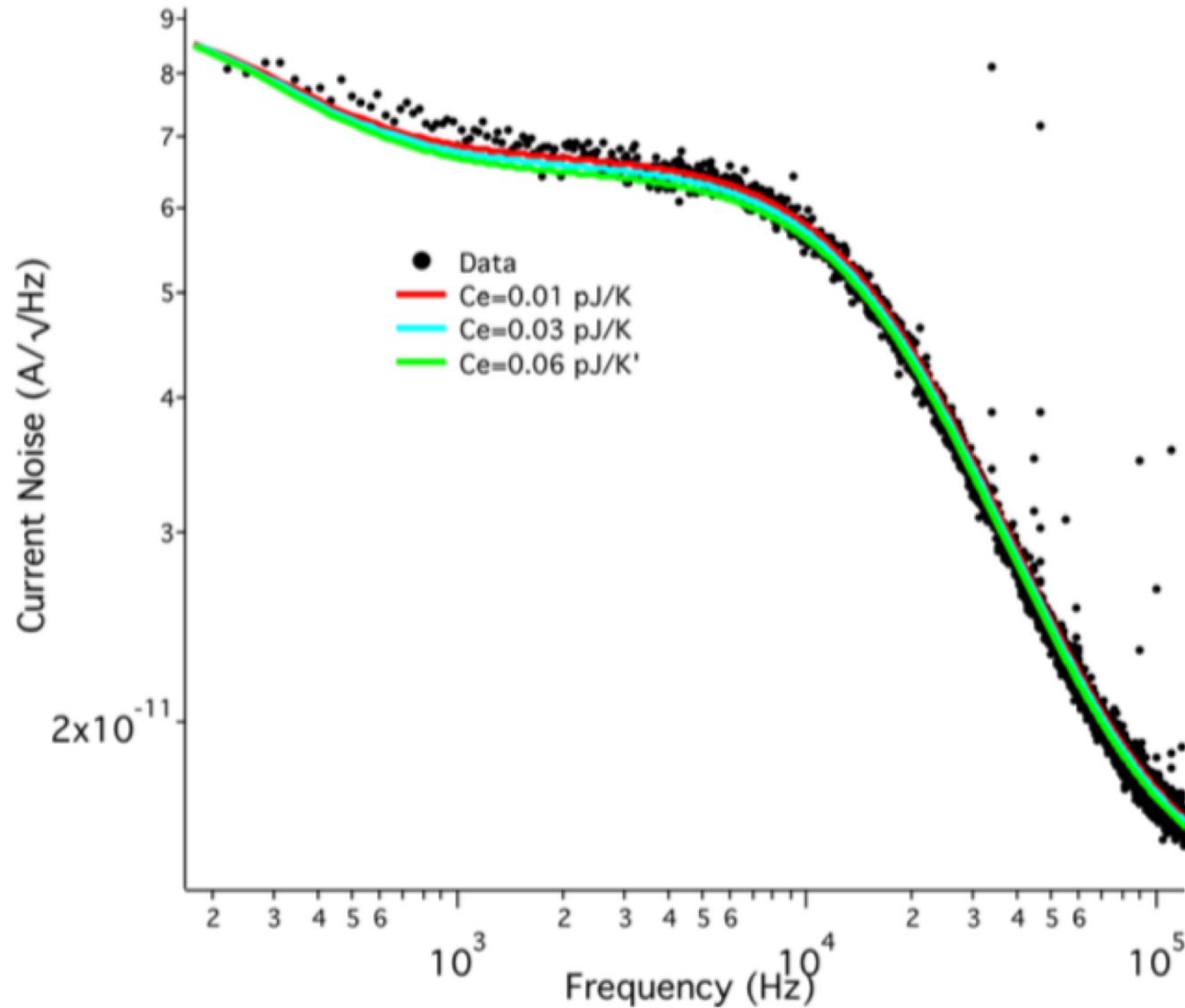
Finite thermal conductance responsible appears to be from resistance of the bilayer.

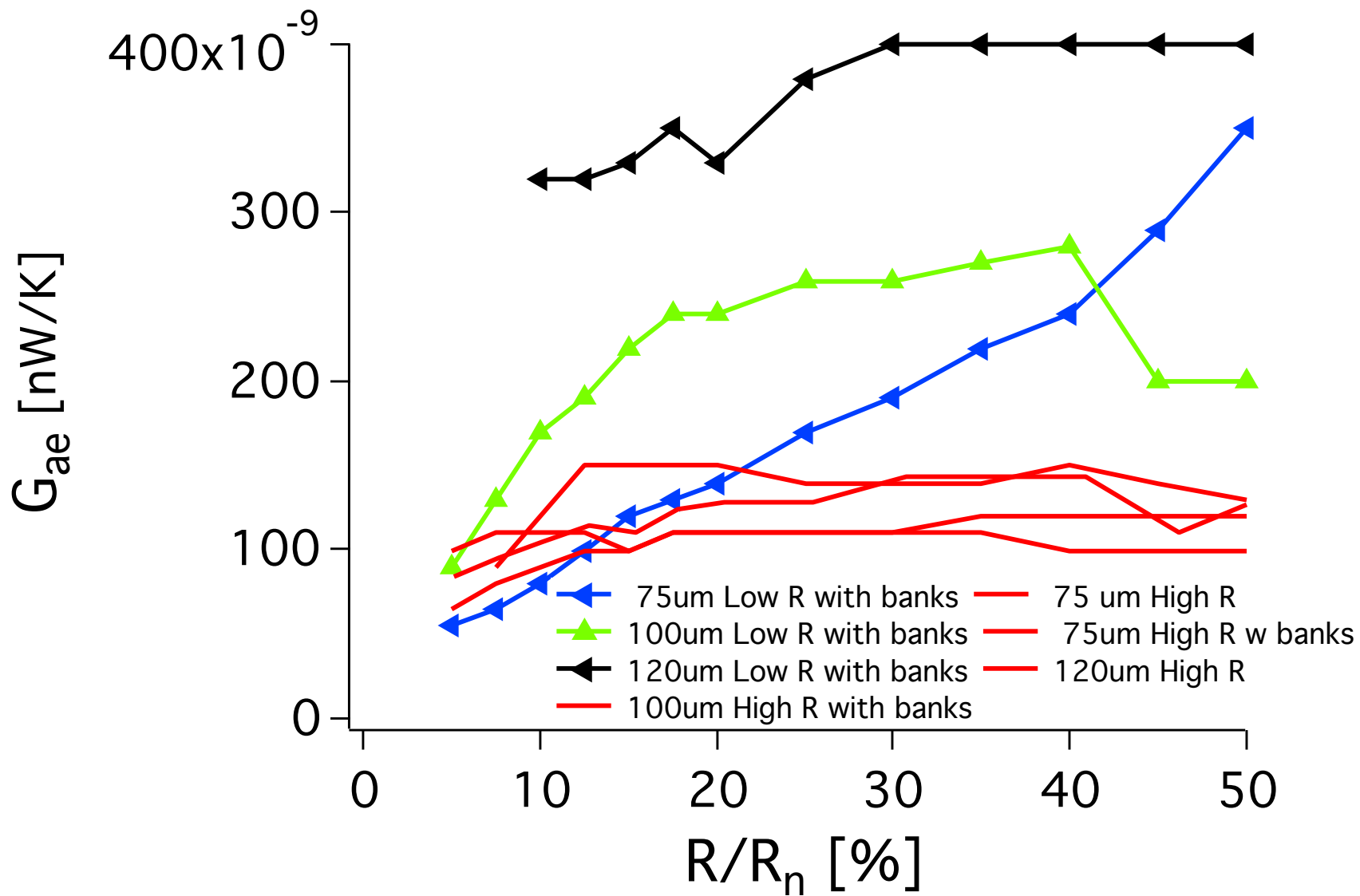
In regions with rapidly changing  $\alpha$  this model is insufficient.

It's likely there that an additional noise mechanism may be present in that case.



# Insensitivity to $C_e$

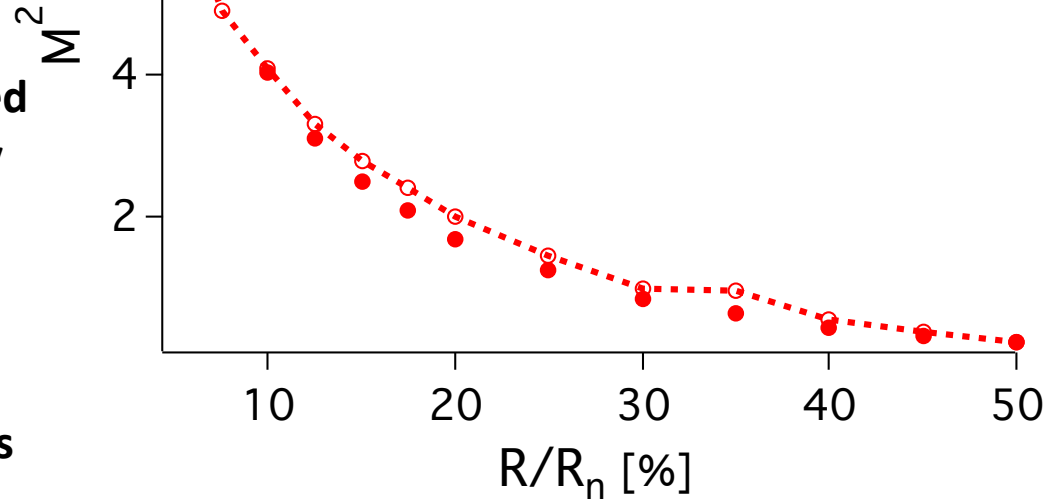




# 75 $\mu$ m TES

With constant  $G_{ae}$  through the transition.  $M^2$  calculated from two body model is only partial agreement with measured values.

If we allow  $G_{ae}$  to vary then we are able to fit at all points





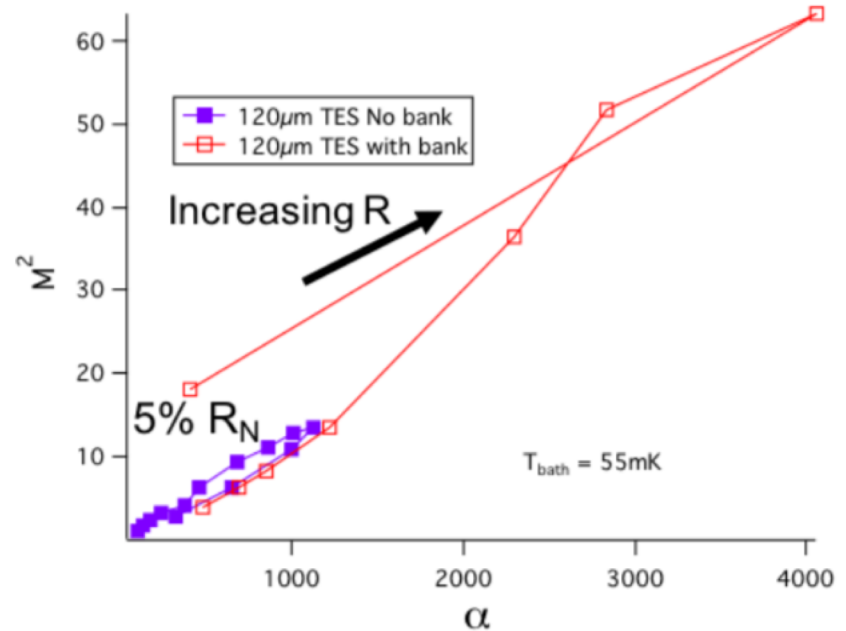
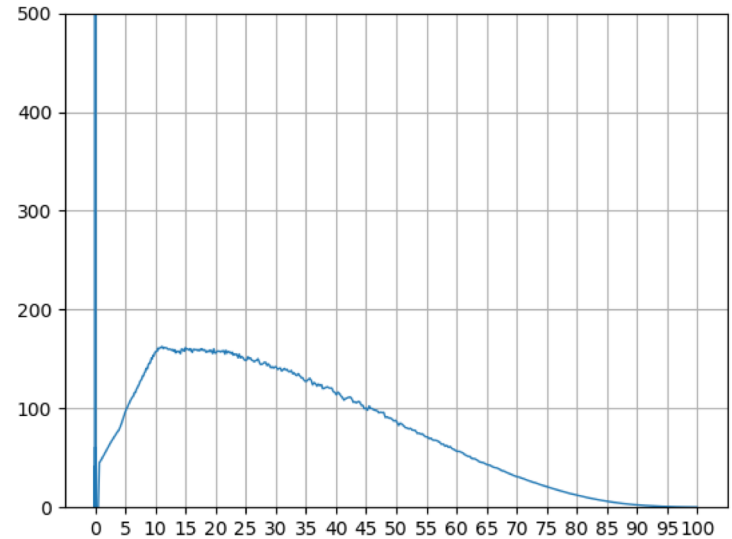
# Where this doesn't work

In kink-like region excess noise far exceeds expectation from our two-body model with reasonable value of  $G_{ae}$ .

In fact noise fits can not be good within our assumptions even with very low  $G_{ae}$ .

Speculate that in this region there is a separate noise term.

This other noise term may have dominated in older devices (e.g. with normal metal stripes)



# Where this doesn't work

Why does this fitting work in these devices?

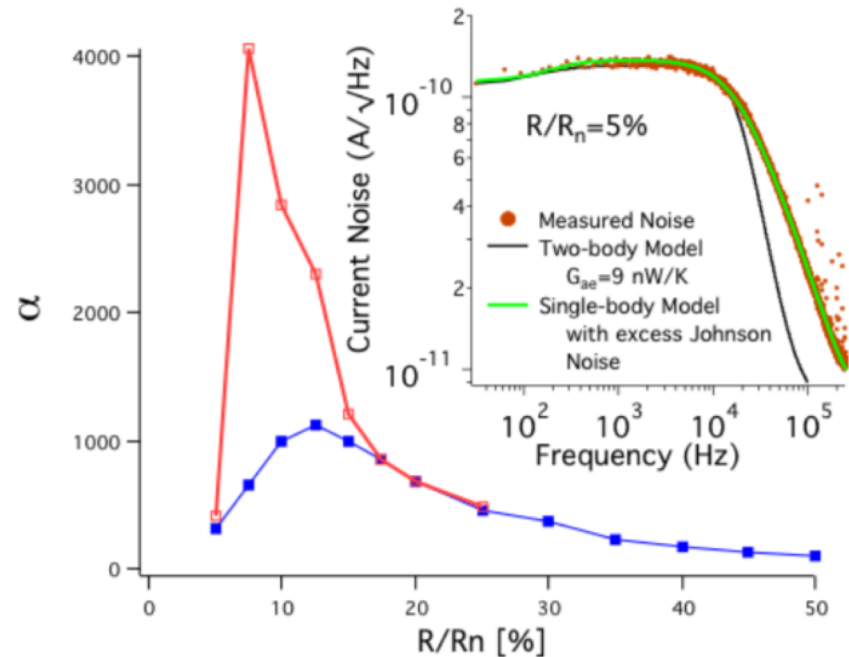
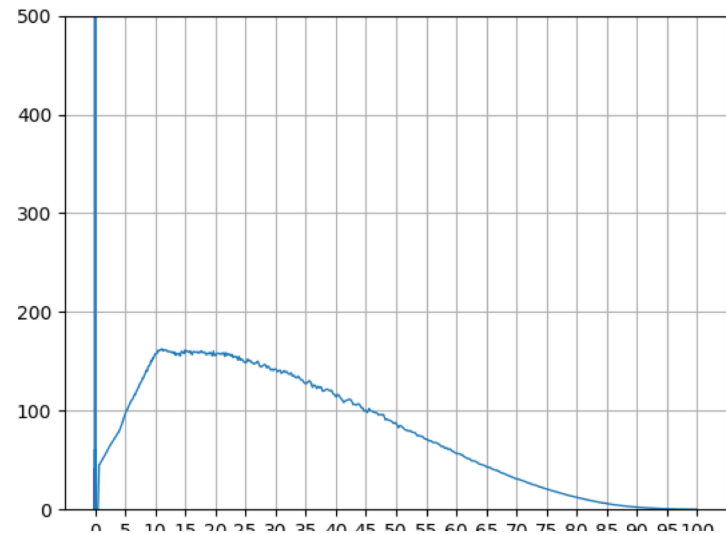
Removing stripes in general produces smoother transitions.

But sometimes see "kinks" low in transition

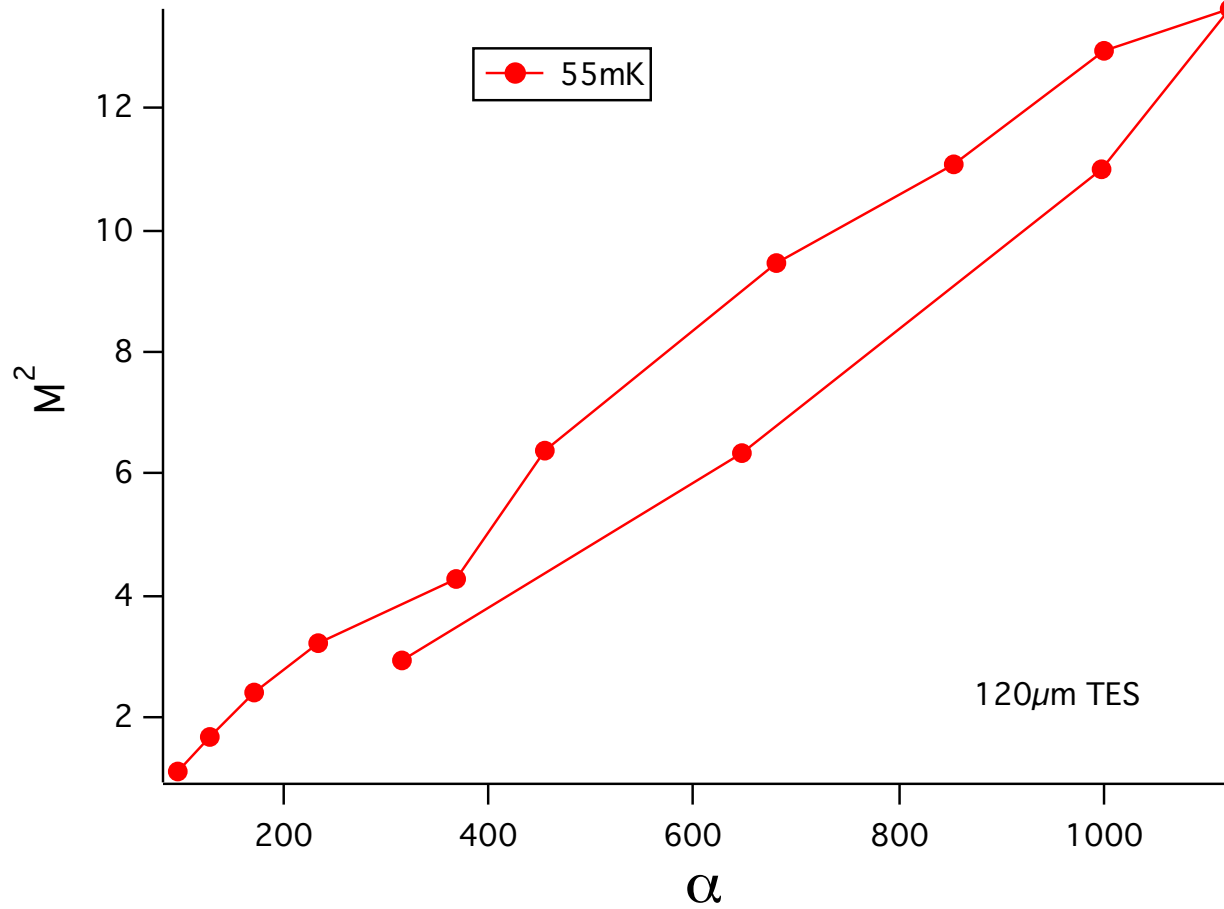
For example in this 120um device with banks.

Subtle features is  $\alpha_{IV}$

Dramatic spike in Alpha.

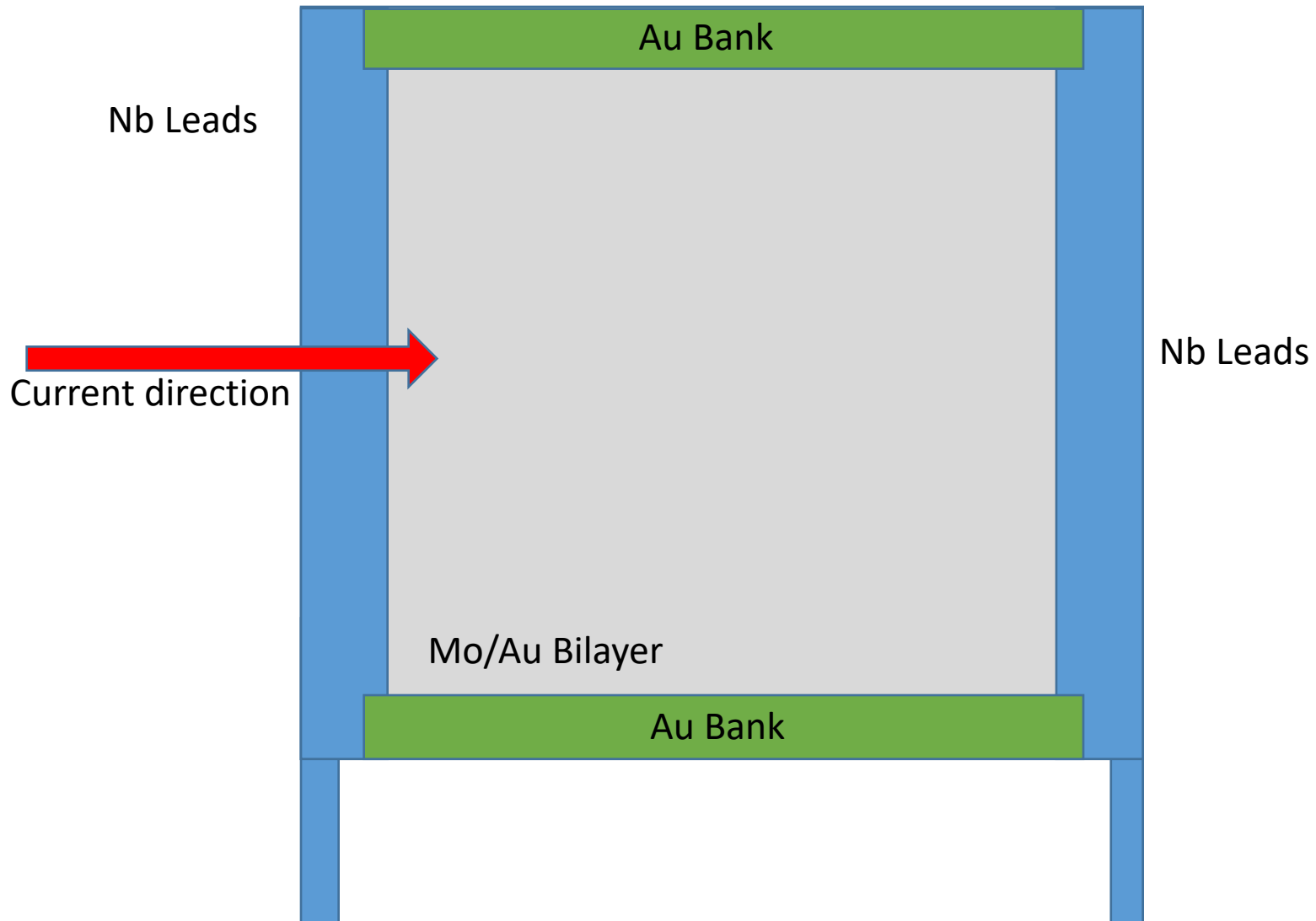


# Assumptions



# Our simple TES design

Some have Au banks



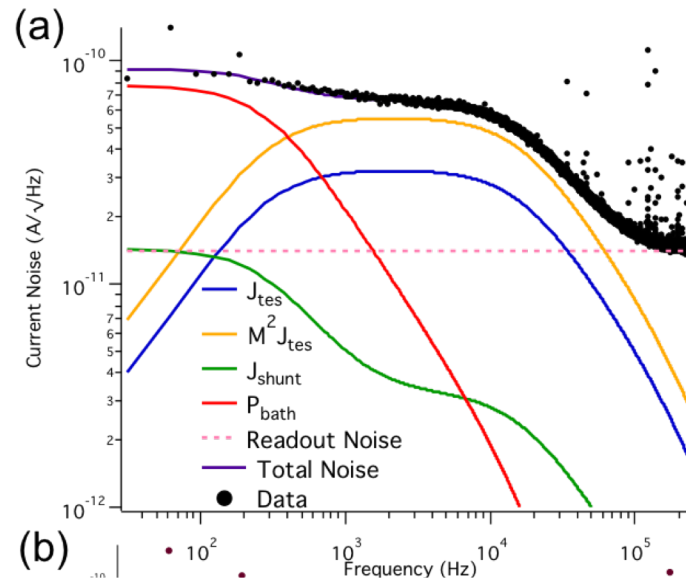
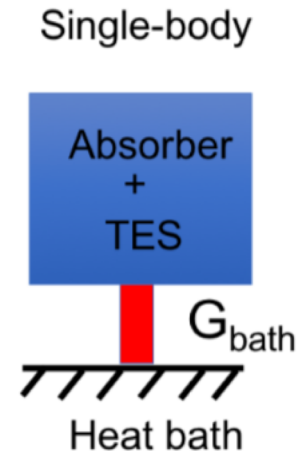
# Noise in excess of single-body model

What is the origin of this additional noise?

Is it related to the Johnson noise (higher order terms)?

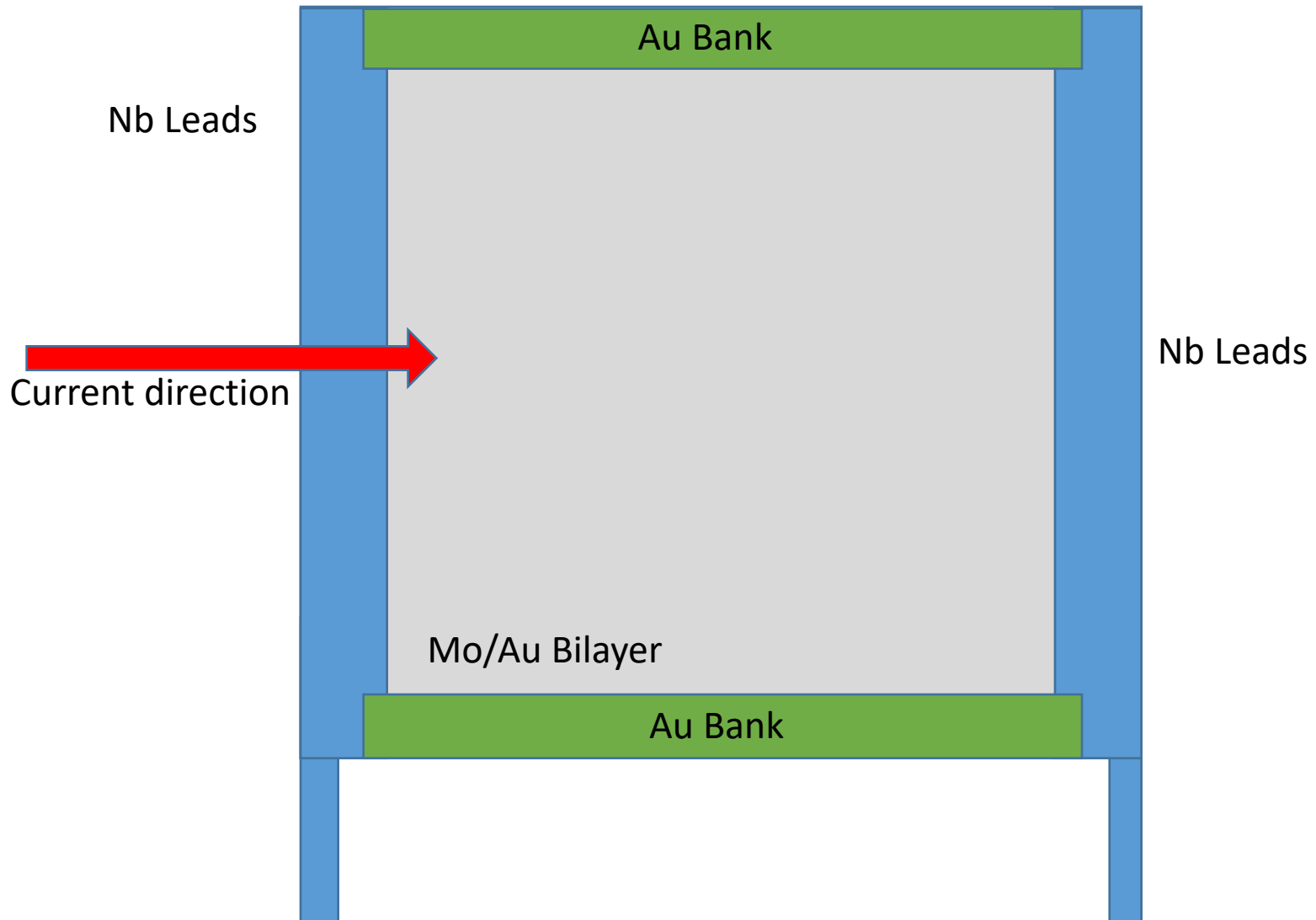
Is it related to phase-slip line behavior?

**Or is it additional thermal fluctuations noise not captured in single-body model?**



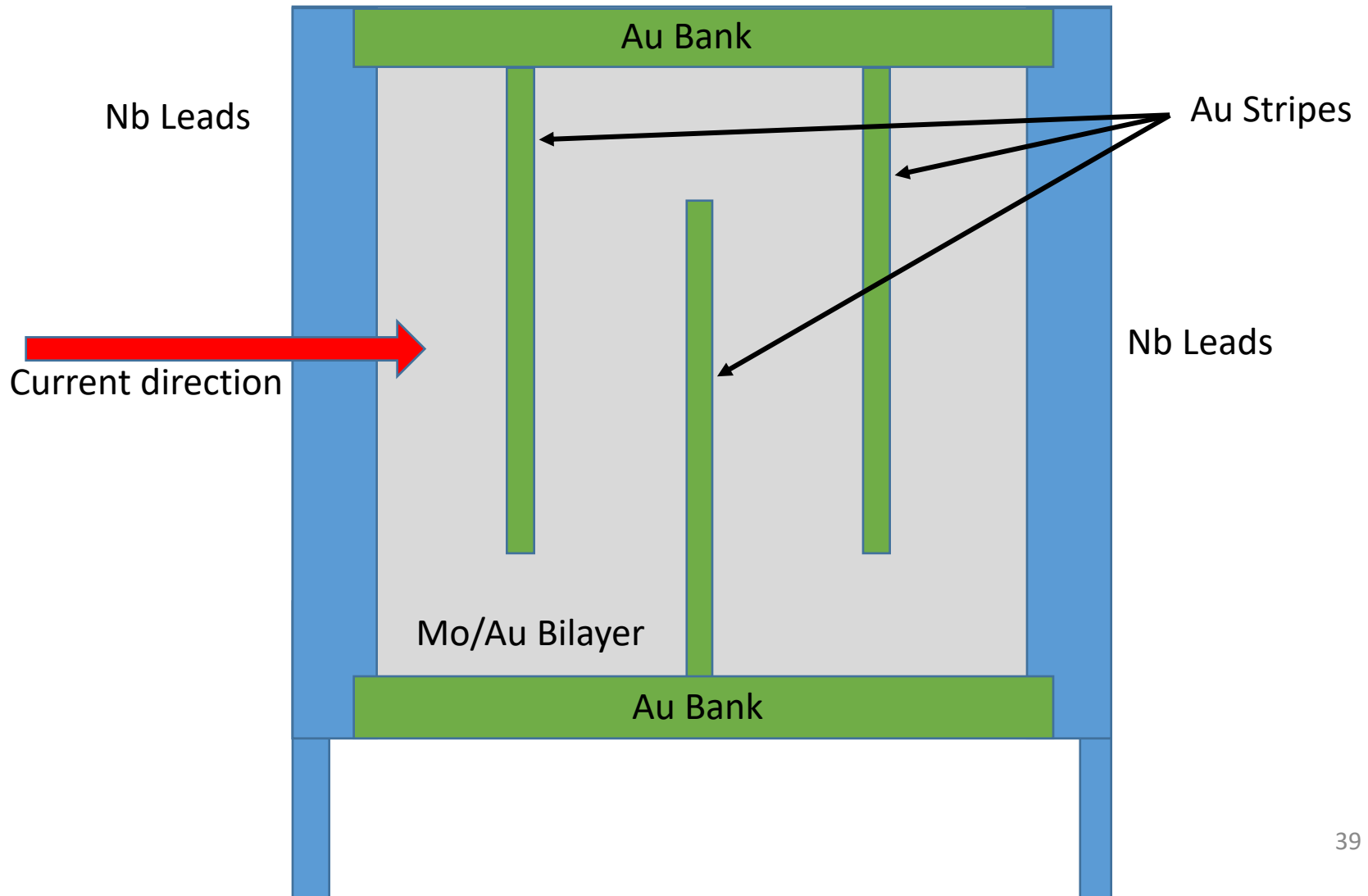
# TES design

Some have Au banks



# Typical TES design

Au stripes reduce unexplained (excess) noise



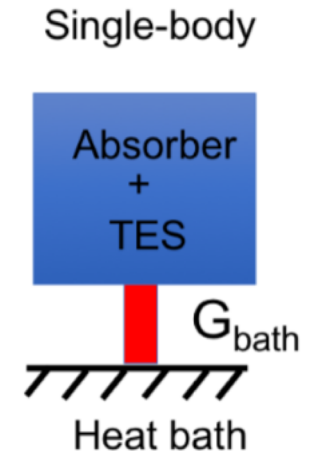
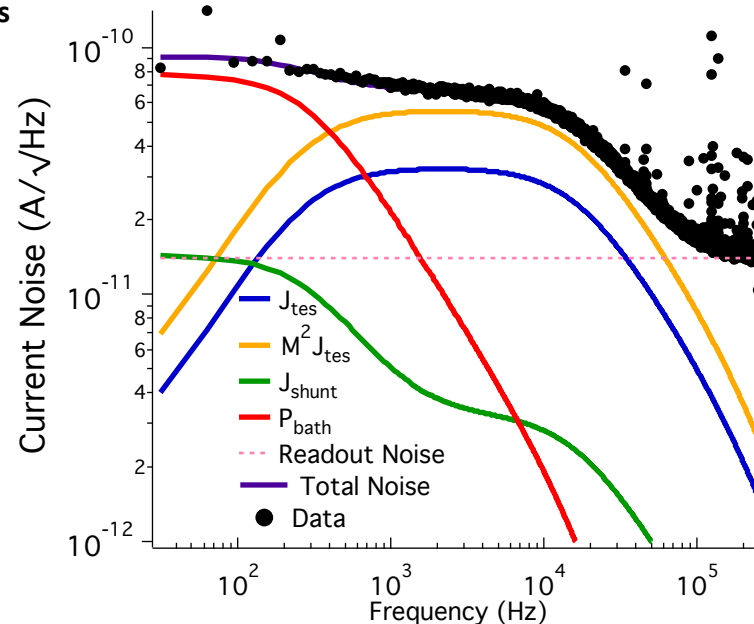
# Noise in excess of single-body model

Simplest model of TES is a single body connected to bath by thermal conductance  $G_{\text{bath}}$

3 noise sources:

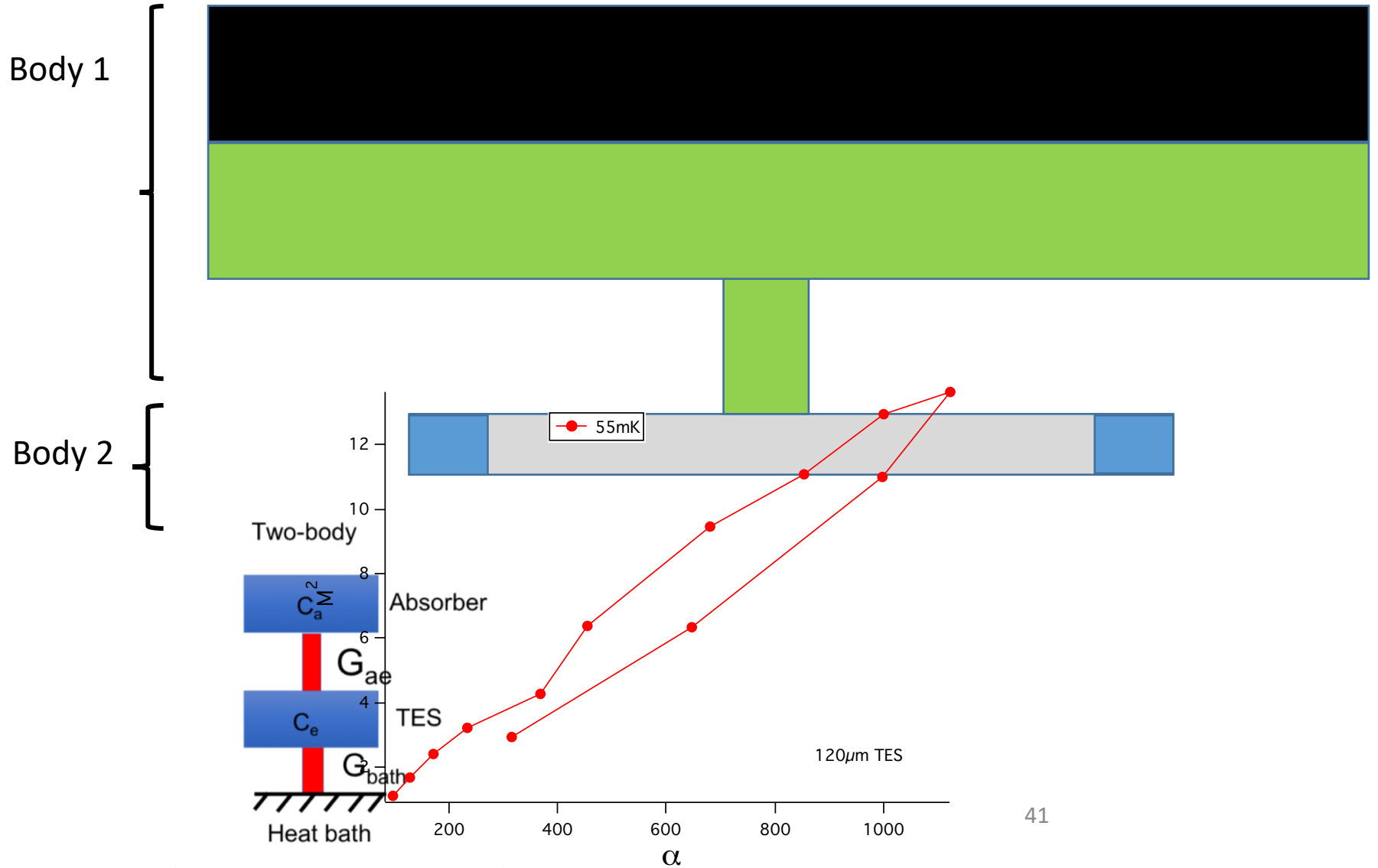
1. Johnson noise of TES
2. Johnson noise of shunt resistor
3. Thermal fluctuation noise between TES/Absorber and bath
4. **Quantify magnitude of excess noise by adding another TES**

Johnson noise term -  $M^2 J_{\text{tes}}$

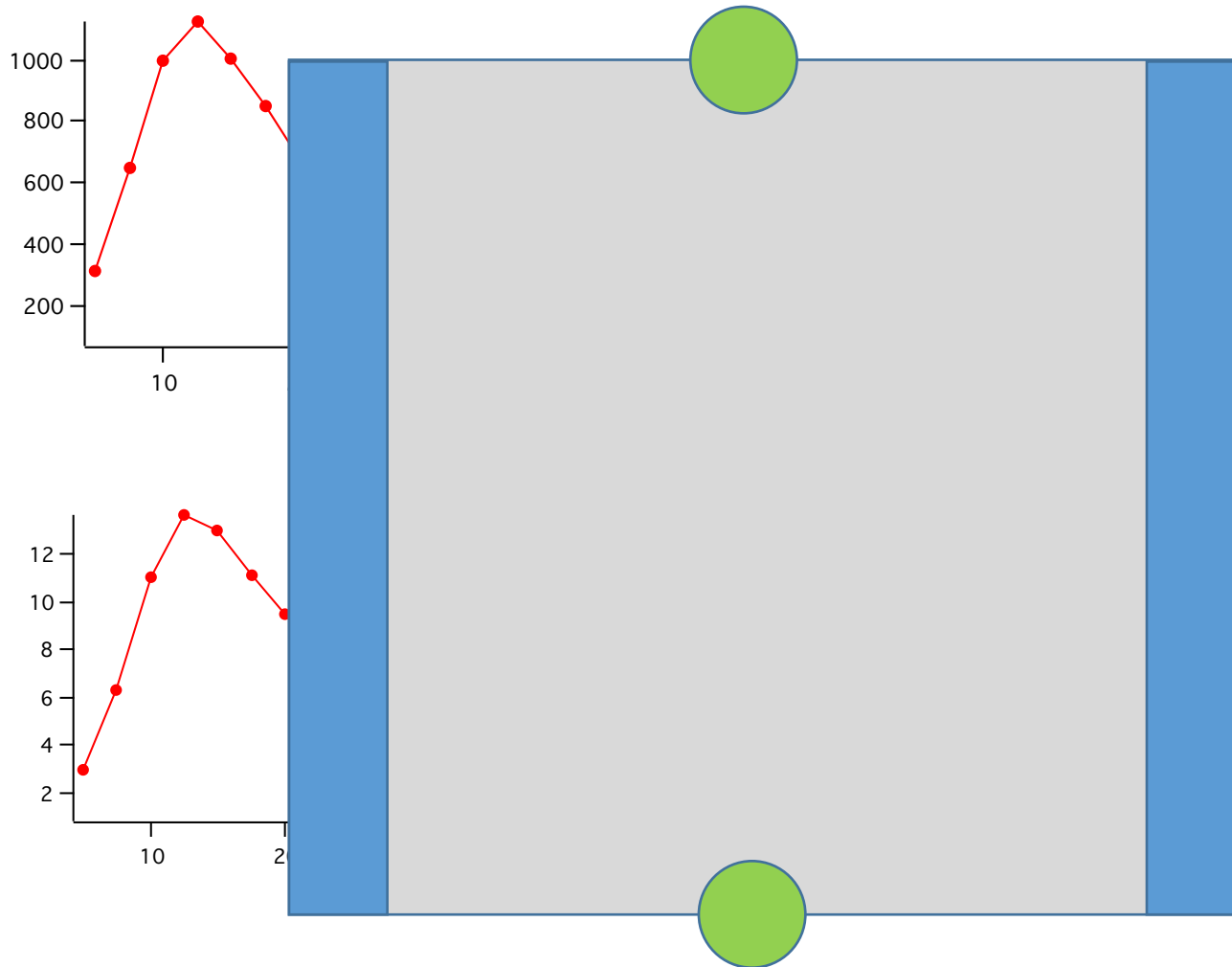




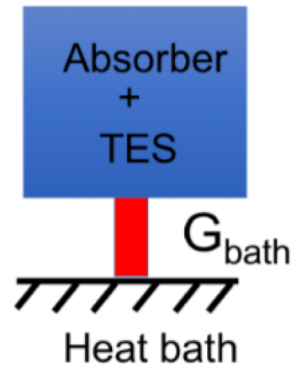
# Assumptions



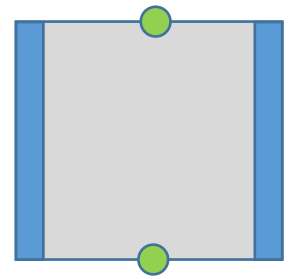
# Fitted $M^2$ for 120 $\mu$ m TES



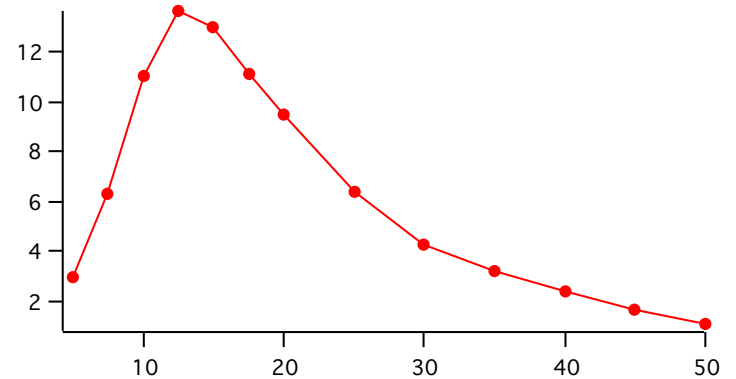
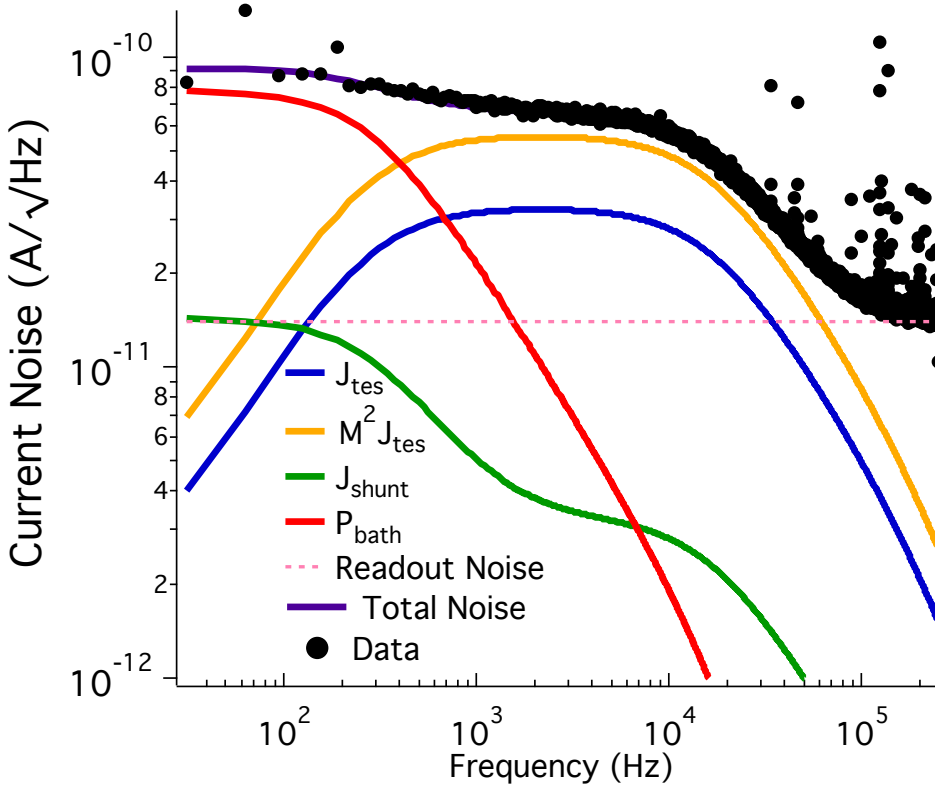
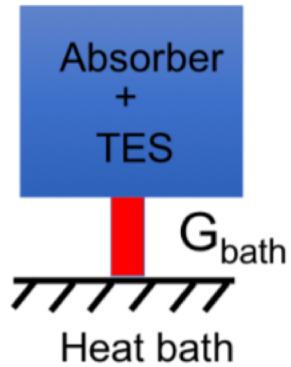
Single-body



# Single-body Vs Two-body



Single-body

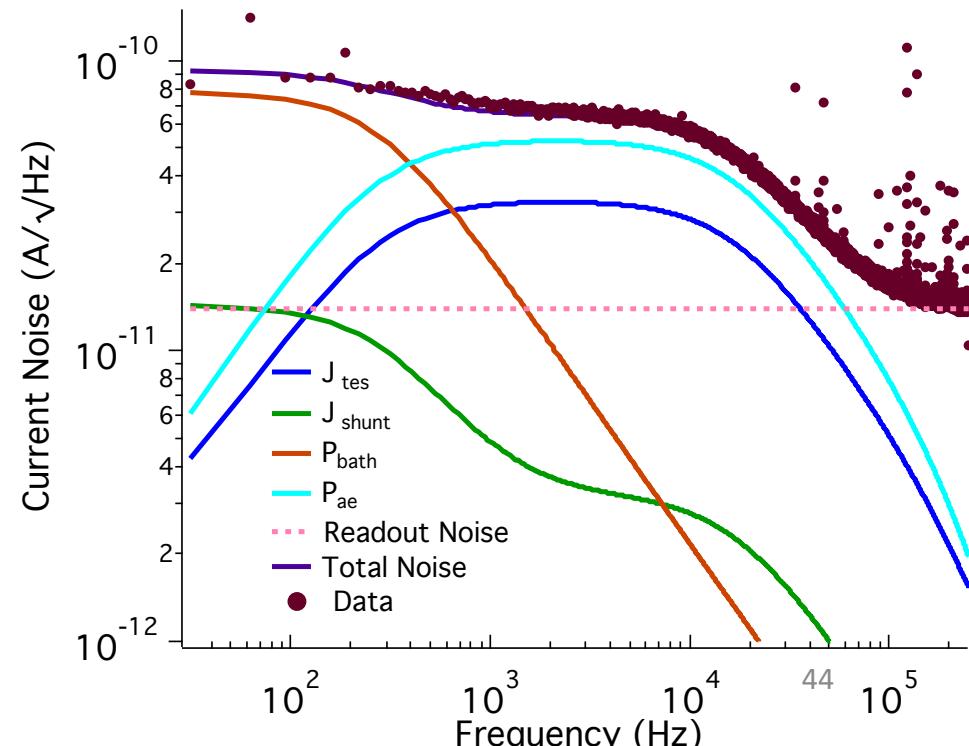
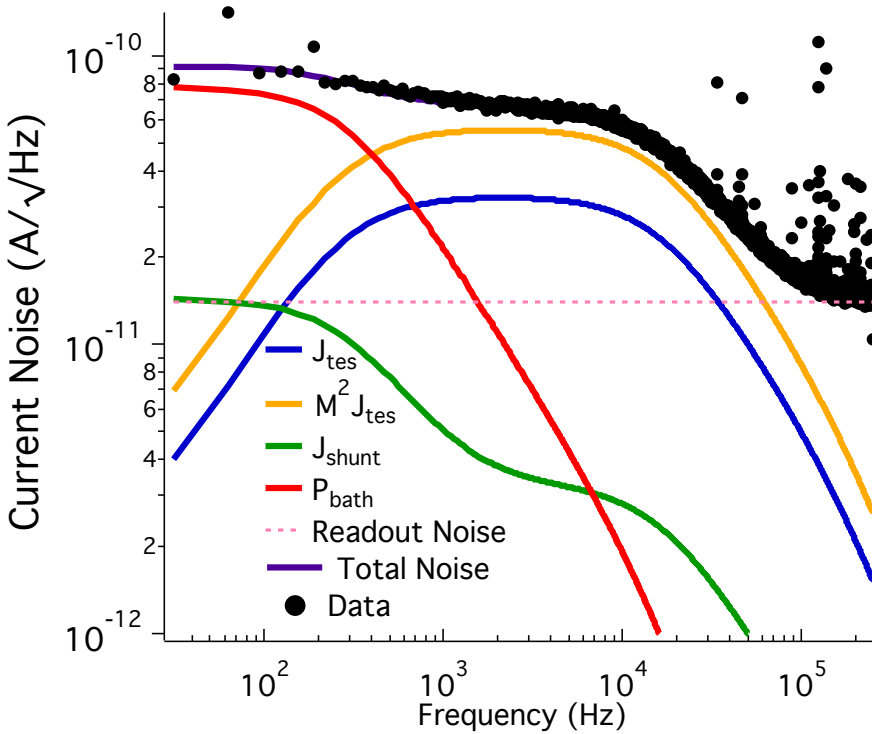


# Single-body Vs Two-body

Internal thermal fluctuation noise can also be used to fit data.

Frequency dependence very similar to  $J_{tes}$

Therefore, I will still quantify magnitude of this excess noise term with  $M^2$

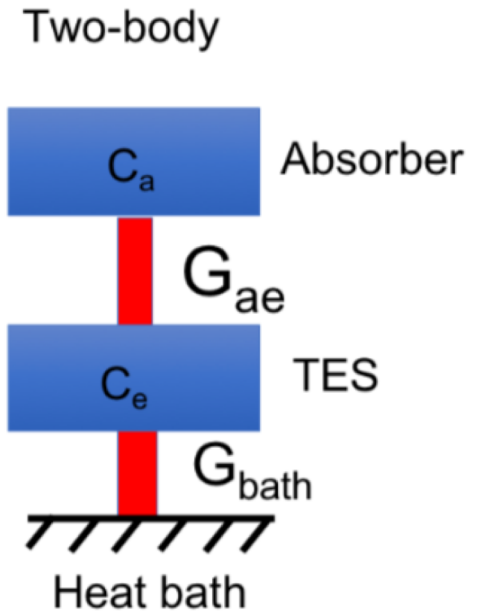


## Assumptions:

$C_e$  = BCS Predicted heat capacity of TES including jump from superconductivity (+ membrane). Relatively insensitive

$C_a + C_e$  = measured heat capacity of device

$G_{\text{bath}}$  is measured value



Fit for  $G_{ae}$

but initially assume  $G_{ae}$  is not a function of  $R/R_n$  or  $T_{\text{bath}}$