

# **Forecasting Rates of Volcanic Activity on Terrestrial Exoplanets**

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**Abstract:** Similar to the planets and moons in our solar system, terrestrial exoplanets may be shaped by volcanism and tectonics. The magnitudes and rates of geological activity on terrestrial exoplanets will be intimately linked to their sizes and internal heating rates, and can either facilitate, or preclude, the existence of habitable environments. In order to place bounds on the potential for such activity, we estimate total internal heating rates for 52 exoplanets, with masses and radii up to  $\sim 8M_E$  and  $2R_E$ , respectively, assuming that internal heating is drawn from both radiogenic and tidal sources. We then compare these internal heating rates to those of the bodies in our solar system in an attempt to constrain the rates of volcanic activity on extrasolar worlds. We find that all of the exoplanets surveyed are likely to exhibit volcanic activity at their surfaces. We also find that at least 25% of the planets in our study may be extrasolar ocean worlds, the majority of which may contain internal oceans beneath layers of surface ice. These planets may be similar in structure to the icy moons of the giant planets and may have persistent cryovolcanic activity at their surfaces. Volcanic activity on exoplanets could be detected by next-generation space telescopes in transit spectra. In the case of planets with densities and/or effective temperatures that are consistent with  $H_2O$ -rich compositions, spectral identification of excess water vapor and other molecules that are explosively vented into space during cryovolcanic eruptions could serve as a way to infer the presence of subsurface oceans, and therefore indirectly assess their habitability. Considering the implications for habitability, our results suggest that characterizing Earth-like exoplanets in terms of the potential for geological activity at their surfaces should be a priority in the coming years.

## 1. Introduction

While in some cases volcanism might induce habitable environments on rocky exoplanets [e.g., *Ramirez and Kaltenegger, 2017*], excessive volcanism [*Demory et al., 2015*] or recurrent, extreme tectonic events (e.g., see [*Hurford et al., 2019*] would produce unstable surface environments that could render exoplanets incapable of hosting life [*Jackson et al., 2008*]. The frequency and relative strength of geological activity on rocky extrasolar planets will therefore directly affect their habitability. In order to gauge the habitability of terrestrial exoplanets, it is therefore necessary to constrain the amount of geological activity that may be occurring at their surfaces. Here, we estimate total internal heating rates of 52 terrestrial exoplanets and use these results to constrain rates of volcanic activity at their surfaces, employing rates of volcanic activity on the planets and moons in our solar system as a baseline.

## 2. Methods

All of the geological activity on the planets and moons in our solar system is driven by internal heating. On Earth, the Moon, and the terrestrial planets, geological processes are driven by radiogenic heating. However on the moons of the giant planets such as Io, Enceladus and Europa, geological processes are primarily driven by tidal heating [*Peale et al., 1979; Cassen et al., 1979; Hurford et al., 2007; Meyer and Wisdom, 2007; Nimmo et al., 2007; Roberts and Nimmo, 2008; Tobie et al., 2008*]. Terrestrial exoplanets may also experience geological activity as a result of tidal heating from their host stars [*Jackson et al., 2008; Henning and Hurford, 2014; Driscoll and Barnes, 2015; Barr et al., 2018; Hurford et al., 2019; Makarov et al., 2018*] as well as due to the decay of radioactive elements in their interiors [*Frank et al., 2014*]. We therefore assume that geological activity on terrestrial exoplanets is driven by both tidal and radiogenic heating.

If all energy imparted to a planet by tidal dissipation is dispersed as heat as the planet orbits

its parent star, then the amount of heat imparted to the planet by tidal heating can be expressed as:

$$H_{Tidal} = \frac{21}{2} \frac{k_2 \omega^5 R_P^5 e^2}{GQ} \quad (1)$$

[*Roberts and Nimmo*, 2008; *Quick and Marsh*, 2015; *Barr et al.*, 2018] where  $k_2$  is the degree 2 love number, which describes how the planet responds to the tide raised on it by its primary.  $k_2$  ranges from 0 for a completely rigid planet, to 1.5 for a planet that is entirely fluid and therefore has a significant tidal response to the gravitational tug of its star. The quantity  $\omega \approx \frac{2\pi}{T}$  is the planet's orbital mean motion, with orbital period  $T$ .  $R_P$  represents the planets' radius, and  $e$  is its eccentricity. Here  $G$  is the gravitational constant, and  $Q$  is the quality factor, which represents the fraction of energy that is dissipated as heat within the planet per orbital cycle.  $Q$  can range from 1 to  $1 \times 10^6$  depending on a planet or moon's composition and internal structure [*Goldreich and Soter*, 1966; *Ojakangas and Stevenson*, 1989; *Henning and Hurford*, 2014]. A significant portion of tidal energy is dissipated as heat in the interiors of planets with high  $Q$ -values.

*Frank et al.* [2014] provides the radiogenic heating rate, per kg of mantle mass, (henceforth  $H_{Radiogenic}^M$ ) in W/kg, for terrestrial exoplanets as a function of age. Assuming that each planet's mantle makes up 84% of its total mass, as is the case for Earth [*Stacey and Davis*, 2008], the volume of each planet's mantle,  $V_{Mantle}$ , can be expressed as:

$$V_{mantle} = 0.84V_P = 0.84 \times \frac{4}{3} \pi R_P^3 \quad (2)$$

If we assume that the density of each exoplanet's mantle,  $\rho_{Mantle} = 4000 \text{ kg/m}^3$ , identical to Earth's mantle, [*Stacey and Davis*, 2008], then the mass of each exoplanet's mantle may be approximated by:

$$M_{mantle} = 0.84 \times \frac{4}{3} \pi R_P^3 \rho_{mantle} \quad (3)$$

Upon obtaining  $H_{Radiogenic}^M$  from *Frank et al.* [2014] and assuming that the age of each exoplanet is identical to the average estimated age of its host star, the total radiogenic heating rate for each planet may be expressed as:

$$H_{Radiogenic} = H_{Radiogenic}^M \times M_{mantle} \quad (4)$$

### 3. Results & Discussion

To demonstrate the utility of (1) and (4)  $H_{Radiogenic}$ ,  $H_{Tidal}$ , and  $H_{Total}$ , are listed along with physical and orbital parameters, for 52 terrestrial exoplanets in Table 1. We considered planets with  $M_P \leq 8 M_{Earth}$  and  $R_P \leq 2R_{Earth}$ , which ensures an Earth-like, rather than a sub-Neptune-like, composition [*Stevenson et al.*, 1982; *Borucki et al.*, 2011; *Fabrycky et al.*, 2014]. In addition, we have used the relationship:  $H_{Total} = H_{Tidal} + H_{Radiogenic}$  and assume that the planets don't have substantive bound atmospheres so that their average surface temperatures,  $T_S$ , may be taken to be equal to their effective temperatures,  $T_{EFF}$ . We have also substituted  $k_2 = 0.3$ , consistent with past studies of the terrestrial planets and icy moons in our solar system [*Kozai*, 1968; *Jackson et al.*, 2008; *Quick and Marsh*, 2015; *Hurford et al.*, 2019], and  $Q = 100$ , commensurate with past studies of tidal dissipation in our solar system's moons [e.g., *Peale et al.*, 1979; *Cassen et al.*, 1979; *Chen et al.*, 2014; *Quick and Marsh*, 2015] and the conservative end of studies of tidal dissipation in terrestrial exoplanets [*Henning and Hurford* 2014; *Tamburo et al.*, 2018], into (1). For comparison, Table 1 also includes  $H_{Total}$ ,  $H_{Radiogenic}$  and  $H_{Tidal}$ , along with physical parameters for the planets and moons in our solar system. In most cases, internal heating rates and physical

parameters for the bodies in our solar system were extracted from the literature (Table 1 Refs. column). In cases where internal heating rates were not documented in the literature, they were calculated using the formulae introduced above. As is the case for the exoplanets in our study, these calculations assumed that solar system bodies are not fully deformable so that values for  $H_{Tidal}$  that are reflective of bodies with homogenous interiors and non-zero rigidities were adopted. (e.g., see Table 3 of *Chen et al.* [2014]). Physical and orbital parameters for each exoplanet were extracted from the NASA Exoplanet Archive and the references listed in Table 1.

We note that with  $H_{Total}$  on the order of  $10^{14}$  W (Table 1), Io is the most volcanically active body in our solar system [*Moore*, 2003; *Lopes et al.*, 2004; *McEwen et al.*, 2004]. Further, volcanic and tectonic activity is prevalent on Earth, which has  $H_{Total} \sim 4 \times 10^{13}$  W (Table 1) [*Turcotte*, 1995; *Turcotte and Schubert*, 2002], primarily due to radiogenic sources [*Schubert et al.*, 1997; *Turcotte and Schubert*, 2002]. Volcanic and tectonic activity was also widespread on early Venus [*Turcotte*, 1995; *Basilevsky et al.*, 1997; *Schubert et al.*, 1997], which may have had  $H_{Total}$  as large as  $2 \times 10^{14}$  W [*Turcotte*, 1995]. Jupiter’s icy moon Europa, which has  $H_{Total} \sim 1$  TW [*Chen et al.*, 2014; *Quick and Marsh*, 2015], is known to be cryovolcanically and tectonically active [*Fagents*, 2003; *Kattenhorn and Hurford*, 2009; *Prockter and Patterson*, 2009; *Roth et al.*, 2014a; *Sparks et al.*, 2016; 2017; *Prockter et al.*, 2017; *Quick et al.*, 2017b; *Jia et al.*, 2018], and explosive cryovolcanism, in the form of cold, geyser-like plumes, occurs on both Europa and Saturn’s moon Enceladus (Fig. 1) [*Porco et al.*, 2006; *Spencer et al.*, 2009; *Roth et al.*, 2014a; *Sparks et al.* 2016; 2017]. Effusive cryovolcanism, in which slurries of cryogenic fluids quiescently erupt, has been observed on both Europa [*Fagents*, 2003; *Miyamoto et al.*, 2005; *Parro et al.*, 2016; *Prockter et al.*, 2017; *Quick et al.*, 2017b] and Neptune’s moon Triton (Fig. 2) [*Croft et al.*, 1995]. We have utilized total internal heating rates, and the associated geological activity on these planets and moons, as a baseline from which the expected rates of volcanic activity on terrestrial exoplanets can be inferred.

### 3.1 Exoplanet Activity

In the context of this study we first consider 55 Cancri e, which observational analysis has suggested is a highly volcanic lava world that is completely melted on its day side [*Demory et al.*, 2015; 2016], and Corot-7b, which has previously been characterized as a super-Io [*Barnes et al.*, 2010]. Utilizing (1) to calculate the tidal heating rate of 55 Cancri e returns  $H_{Tidal} = 4.8 \times 10^{21}$  W (Table 1), while (3) returns  $M_{Mantle} = 2.54 \times 10^{25}$  kg. If we assume that each planet’s age is identical to the average estimated age of its host star, then the 10.2 Gyr age estimate for 55 Cancri [*von Braun et al.*, 2011] can be utilized to obtain  $H_{Radiogenic}^M = 4.49 \times 10^{-12}$  W/kg for 55 Cancri e [*Frank et al.*, 2014]. Multiplying this value by  $M_{Mantle}$ , as in (4), returns  $H_{Radiogenic} = 1.14 \times 10^{14}$  W (Table 1). In the case of Corot-7b, (3) returns  $M_{Mantle} = 1.4 \times 10^{25}$  kg. We have use the average estimated age of its host star, i.e., 1.32 Gyr [*Barros et al.*, 2014], to obtain  $H_{Radiogenic}^M = 1.59 \times 10^{-11}$  W/kg [*Frank et al.*, 2014]. Multiplying this value by  $M_{Mantle}$  as in (4), returns  $H_{Radiogenic} = 2.17 \times 10^{14}$  W for Corot-7b. If  $e = 0$  for Corot-7b [*Queloz et al.*, 2009; *Barros et al.*, 2014; *Stassun et al.*, 2017],  $H_{Tidal} = 0$  and  $H_{Total} = H_{Radiogenic} = 2.17 \times 10^{14}$  W (Table 1).  $H_{Total}$  for Io is between  $1 \times 10^{14}$  W and  $2 \times 10^{14}$  W [*Veeder et al.*, 1994; *Spencer et al.*, 2000]. This suggests that 55 Cancri e and Corot-7b are both likely to be very volcanically active. As suggested in previous studies, 55 Cancri e likely exhibits extreme volcanism and may have a molten surface. Identical analyses for L98-59b,c & d [*Kostov et al.*, 2019], assuming that L98-59 and its orbiting planets are at least 1 Gyr old, returns  $H_{Radiogenic} = 3.33 \times 10^{13}$  W,  $1.6 \times 10^{14}$  W, and  $2.52 \times 10^{14}$  W, respectively for each planet, with calculated  $H_{Tidal}$  and  $H_{Total}$  for each world being on the order of  $10^{17}$  W. A simple comparison

of  $H_{Total}$  for Io with  $H_{Radiogenic}$  for the aforementioned planets (Table 1) makes it clear that even in the absence of tidal heating, 55 Cancri e, Corot-7b, and L98-59 b, c, and d are all likely to exhibit extreme volcanism at their surfaces.

The results for  $H_{Total}$  presented in Table 1 can be employed to qualitatively assess the potential for volcanic activity on terrestrial exoplanets compared to the rates at which these processes occur on the planets and moons in our solar system. Owing to their large surface areas through which heat will escape, small planets will have much lower internal temperatures and may cool much faster than large planets [Stevenson, 2003]. This will ultimately result in the cessation of geological activity, including volcanism, at their surfaces. In order to constrain the magnitude of volcanic activity occurring on the remainder of the planets in our study, we have plotted  $H_{Total}$  as a function of surface area to volume ratio ( $SA/V = 3/R$ ) for the planets and moons in our solar system and 51 terrestrial exoplanets, in Fig. 3. The pink line in Fig. 3 connects Earth, Enceladus, and Europa, all of which are replete with geological activity, including (cryo)volcanism and tectonics. Exoplanets that lie along this line are likely to exhibit similar rates of geological activity as these bodies and may contain internal oceans or magma layers. Exoplanets that lie above this line are likely to display pervasive volcanism and tectonics at their surfaces. Depending on where they plot, rates of volcanism on these exoplanets may be similar to rates of volcanism on Io, Early Venus, or Corot-7b. Rates of volcanism on these bodies may be so high as to be detectable in transit spectra [Kaltenegger et al., 2010; Quick et al., 2017a] or in thermal emission observations [e.g., see Demory et al., 2015].

Planets that lie on or above the blue stippled line are likely to exhibit more subdued geological activity at their surfaces. Similar to the Moon, they may occasionally experience tectonic activity in the form of earthquakes [Watters et al., 2019]. Although widespread volcanism may not occur on these planets, effusive volcanism in the form of (cryo)lava flows may occasionally manifest as it has recently on the Moon [Garry et al., 2012; Braden et al., 2014; Qiao et al., 2017; ], Venus [Smrekar et al., 2010] and Jupiter’s moon Ganymede [Head et al., 1998; Schenk et al., 2001]. These planets may also experience internal convection on a regional scale similar to Triton [Schenk and Jackson, 1993; Croft et al., 1995], and may contain global oceans in their interiors. Pluto, which sits just below the blue stippled line, is known to have internal convection, albeit localized to only a few regions on the surface [McKinnon et al., 2016; Trowbridge et al., 2016], and may also contain a subsurface ocean [Hammond et al., 2016; Nimmo et al., 2016]. It must be noted, however, that the maintenance of an ocean within Pluto may only be possible if the top of the ocean is capped by a layer of insulating clathrates [Kamata et al., 2019]. Consideration of the incorporation of clathrates and other insulating materials into planetary interiors is beyond the scope of this work. However, Pluto’s internal state suggests that exoplanets with similar heating rates and surface area to volume ratios might also be able to maintain subsurface liquid reservoirs in special cases where substantial amounts of insulating materials have been incorporated into their interiors. Conversely, exoplanets with lower internal heating rates and significantly larger surface area to volume ratios would plot substantially below the blue stippled line (Fig. 3). These exoplanets would be geologically inactive today. Figure 3 illustrates that the majority of bodies that lie below the blue dotted line have: small radii, on the order of hundreds of km, large surface area to volume ratios, and experience negligible tidal heating. Note that as a result of its large tidal heating rate,  $H_{Total}$  for 55 Cancri e is so large that it could not be plotted with the other worlds in Fig. 3.

### 3.2 Extrasolar Ocean Worlds

Several of the terrestrial planets considered in our study may be more specifically described as ocean planets. Ocean planets are a class of low-density, terrestrial exoplanets with substantial water layers that may be common throughout the galaxy [Kuchner, 2003; Léger *et al.*, 2004; Ehrenreich and Cassan, 2007; Sotin *et al.*, 2007; Fu *et al.*, 2010]. These planets may exist in one of a variety of climactic states, including, ice-free, partially ice covered, and completely frozen [Budyko, 1969; Sellars, 1969; Tajika, 2008]. In the context of this study, planets that have conditions which are favorable to the maintenance of liquid water on their surfaces or in their interiors are assumed to be ocean planets. Conditions which may be favorable to the maintenance of liquid water include  $T_{EFF} < 373$  K, bulk densities below  $5000 \text{ kg/m}^3$  or, in the case of high-gravity planets (Table 1) with potentially high surface pressures (e.g., surface pressures between  $10^6$  and  $10^7$  Pa),  $T_{EFF}$  below the critical point (i.e.,  $T_{EFF} < 647$  K). The effective temperatures (Table 1) and proposed internal structures of Trappist-1c & d [Barr *et al.*, 2018] suggest that they could be ocean planets with surfaces that are covered in liquid water. This could be especially true of Trappist-1d [Dobos *et al.*, 2019], as its low density is suggestive of a substantial fraction of volatiles (Table 1).  $H_{Total}$  for these planets is  $3 \times 10^{16}$  W and  $3 \times 10^{13}$  W, respectively (Table 1). This indicates that they are likely to have extreme volcanism at their surfaces, with Trappist-1c being more volcanically active than Corot-7b (Fig. 3), in agreement with the results of previous analyses by Dobos *et al.* [2019].

We refer to exoplanets with ice-covered surfaces that overlie internal oceans [Kuchner, 2003; Ehrenreich *et al.*, 2006; Ehrenreich and Cassan, 2007; Tajika, 2008; Yang *et al.*, 2017] as *cold ocean planets*. The internal structures of cold ocean planets may resemble the internal structures of our solar system’s icy moons (Fig. 4) [Ehrenreich *et al.*, 2006; Sotin *et al.*, 2007; Vance *et al.*, 2007; Fu *et al.*, 2010; Henning and Hurford, 2014; Vance *et al.*, 2015; Noack *et al.*, 2016; Luger *et al.*, 2017; Barr *et al.*, 2018], and they may exhibit similar geological activity at their surfaces, including ice tectonics [Fu *et al.*, 2010; Levi *et al.*, 2014] and cryovolcanism [Levi *et al.*, 2013; Quick *et al.*, 2017a; Barr *et al.*, 2018; Quick and Roberge, 2018]. In the absence of an atmosphere, Earth’s effective temperature is 255 K [Sagan and Mullen, 1972], which allows for the maintenance of liquid water at the surface. We therefore assume that cold ocean planets have  $T_{EFF} < 255$  K. Several of the planets plotted in Fig. 3 have  $T_{EFF} < 255$  K, and/or densities that are  $\leq 3500 \text{ kg/m}^3$  (Table 1), commensurate with the range of density values for our solar system’s icy moons [Jacobson *et al.*, 1992; Anderson *et al.*, 1997; Hussmann *et al.*, 2006]. Based on their surface temperatures and bulk densities, we have designated these exoplanets as cold ocean planets.

The effective temperatures, densities, and total internal heating rates of Kepler 138-b & d, Trappist-1e, f, g & h, Kepler 60-c & d suggest that these planets could be cold ocean planets that have persistent cryovolcanic activity at their surfaces (Table 1 & Fig. 3). Although LHS 1140 b’s high-density is indicative of an iron-rich terrestrial planet [Dittmann *et al.*, 2017a; Ment *et al.*, 2019], its 235 K surface temperature [Ment *et al.*, 2019] suggests that any water on its surface would be in a frozen state. Further, its estimated total internal heating rate of almost  $2 \times 10^{14}$  W may be high enough for it to maintain an internal ocean and cryovolcanic eruptions at the surface. While cryovolcanic activity on Kepler 138-b and Trappist 1-g & h may be more subdued than what has been observed on Enceladus and Europa (Fig. 1), their plotted positions Fig. 3 suggest that they may undergo periodic resurfacing by cryolava flows and/or experience periodic outgassing of volatiles similar to planet Venus (e.g., see [Bondarenko *et al.*, 2010; Smrekar *et al.*, 2010; Shalygin *et al.*, 2015]). Thus, cryolava flows on the surfaces of these planets may be similar to what has been observed in the walled plains units of Neptune’s moon Triton (Fig. 2), albeit more

widespread. Explosive cryovolcanism, similar to what has been observed on Enceladus and Europa (Fig. 1) [Porco *et al.*, 2006; Spencer *et al.*, 2009; Roth *et al.*, 2014a; Sparks *et al.*, 2016; 2017], may occur at comparable magnitudes on Kepler 138d, Trappist-1f and LHS 1140 b.

According to Fig. 3, rates of cryovolcanism on Trappist-1e and Kepler 60 d may be similar to rates of volcanism on Jupiter's moon Io, while Kepler 60c is likely to exhibit extreme cryovolcanism at rates much higher than the silicate volcanism that occurs on Corot-7b. Based on their internal heating rates and surface area to volume ratios compared to the planets and moons in our solar system (Fig. 1), all of the aforementioned exoplanets may contain extensive reservoirs of liquid water, possibly in the form of global oceans beneath layers of surface ice. As is the case for Europa [Crawford and Stevenson, 1988; Fagents, 2003] and Enceladus [Manga and Wang, 2007; Běhouňková *et al.*, 2015], cryovolcanic eruptions on these worlds may provide a pathway by which ocean water, or the contents of discrete water pockets within their icy crusts, reaches the surface. It must be noted that in the case of Trappist-1f, g, & h,  $H_{\text{Radiogenic}}$  is greater than, or equal to,  $H_{\text{Tidal}}$  (Table 1). Hence, even if Trappist-1f & h are too far away from their host star to experience substantial tidal heating [Dobos *et al.*, 2019], they are still likely to be geologically active and able to maintain subsurface oceans. On the other hand, our analyses suggest that in the absence of tidal heating, Trappist-1g may be geologically dead (see Table 1).

The densities and/or effective temperatures of Trappist-1b, Kepler 93 b, Kepler 138 c, Kepler 114 c, Kepler 80 e, LHS 1140 c, Pi Mensae c, and L 98-59 d are such that they could also be ocean planets or cold ocean planets. If surface pressures on these planets are high enough, liquid water could be maintained at their surfaces at the critical point, or they could contain high pressure ices. These planets are labeled as candidate ocean worlds in Table 1. With the exception of Kepler 114 c, for which we were unable to obtain  $H_{\text{Total}}$  due to the unknown age of its host star and its unknown eccentricity, estimated internal heating rates for all of these planets suggest that they are all likely to be geologically active (Fig. 3). Of note is that Pi Mensae c and L98-59 d likely receive enough internal heating from radiogenic sources alone to be as volcanically active as Io (Table 1 and Fig. 3), while LHS 1140 c may receive enough internal heating from radiogenic sources alone to be as geologically active as Earth. *Ment et al.* [2019] report an eccentricity less than 0.31 for LHS 1140c. Using 0.31 as an upper bound for eccentricity in (1) results in tidal and total internal heating rates for this planet that are several orders of magnitude greater than the tidal and total internal heating rates for Corot-7 b (Table 1 & Fig 3). Such high internal heating rates would suggest that LHS 1140 c is a lava world instead of a candidate ocean world. However, owing to uncertainties in  $e$  [Ment *et al.*, 2019] the estimated  $H_{\text{Total}}$  for this planet represents a maximum value that may change once  $e$  is better constrained. Bearing this mind, we contend that LHS 1140 c may be an ocean world that contains exotic forms of ice on its surface or within its interior. Due to its high  $T_{\text{EFF}}$  and low density, Pi Mensae c could also contain exotic forms of ice.

### 3.3 Model Limitations

Utilizing (1) to determine  $H_{\text{Total}}$  for Io and Enceladus where Jupiter and Saturn serve as the primaries from which tidal heating is sourced, returns total internal heating rates of  $2 \times 10^{13}$  W and  $3.2 \times 10^8$  W, respectively, which are 1-2 orders of magnitude less than their actual values of  $1 \times 10^{14}$  W, and  $1.6 \times 10^{10}$  W, respectively [Veeder *et al.*, 1994; Spencer *et al.*, 2000; Schubert *et al.*, 2004; Howett *et al.*, 2011]. Enceladus may be considered an outlier as unlike Io where  $H_{\text{Total}}$  is representative of global heating rates, the bulk of Enceladus' internal heating is concentrated at its south pole [Howett *et al.*, 2011]. Nevertheless, this discrepancy suggests that substituting  $k_2 = 0.3$  and  $Q = 100$  in (1) returns conservative values for exoplanet tidal heating rates, and by extension,

for the magnitude of geological activity that may be occurring at the surfaces of the exoplanets we have considered. Furthermore, we have not considered additional heating sources that might be contributed by dynamical planet-planet or planet-moon interactions within exoplanetary systems. We have also neglected to consider cases in which the dynamical histories and internal structures of planetary bodies could preclude substantial internal heating. This is the case for Saturn’s moon Mimas. Although Mimas is known to be a geologically dead world [Rhoden *et al.* 2017; Kirchoff *et al.*, 2018], utilization of (1) returns a tidal heating rate that suggests that the small moon should be more geologically active than Enceladus. Neveu and Rhoden [2019] suggest that the absence of current geological activity on Mimas may be due to a loss of radiogenic heating early in its evolution. These authors assert that loss of radiogenic heating precluded Mimas from having a dissipative interior. Such considerations for Earth-like exoplanets are outside the scope of this work. Nevertheless, the agreement of our results with past studies that considered the magnitude of geological activity on exoplanets (e.g., [Jackson *et al.*, 2008; Barnes *et al.*, 2010; Demory *et al.*, 2015; 2016; Dobos *et al.*, 2019]) illustrates that tidal heating rates obtained using (1), and radiogenic heating rates extracted from Frank *et al.* [2014], can be reliably employed to place conservative estimates on the expected magnitude of internal heating, and by extension, volcanic activity, on solid exoplanets. Our calculations reveal that the majority of exoplanets considered in this study are likely to exhibit volcanic activity, and may also exhibit active tectonics, at levels that are at least on par with, if not greater than, the most geologically active bodies in our solar system. Hence from the standpoint of comparative exoplanetology, in which the internal heating rates and geological activity of the planets and moons in our solar system are used as a baseline, our results regarding the expected magnitude of geological activity on terrestrial exoplanets are robust.

In order to ensure maximum accuracy when constraining the amount of geological activity on the exoplanets listed in Table 1, specific  $k_2$  and  $Q$  values must be determined for each planet. This will require a new modeling approach that would allow us to constrain the relative amounts of ice, metal, and rock likely to be present in the planets considered. This approach was investigated in Barr *et al.* [2018] and could be utilized to confirm the terrestrial status of the largest planets in our study, including, Kepler 60c & d, LHS 1140 b, Kepler 138 c, and Pi Mensae c. All of these planets have radii greater than  $1.6R_{Earth}$ , which could be indicative of a sub-Neptune-like composition [Rogers *et al.*, 2015]. Such an approach will be the focus of a forthcoming manuscript.

#### 4. Conclusions

Based on estimated internal heating rates, all of the exoplanets surveyed here are likely to be geologically active. Comparing calculated internal heating rates of each exoplanet in our study to the estimated internal heating rates of the planets and moons in our solar system indicates that all 52 exoplanets we considered are likely to exhibit moderate to extreme rates of volcanism at their surfaces. In addition, out of the 52 exoplanets that were surveyed, 13 (~ 25%) have effective temperatures and/or densities that are consistent with them being candidate ocean planets. 8 out of these 13 planets (~ 61%) have effective temperatures and/or densities that are indicative of them being cold ocean planets. Hence, about 15% of all of the planets surveyed in this study may be cold ocean planets with internal structures that are similar to the moons of our solar system’s giant planets, and similar geological activity at their surfaces. If conditions at the surfaces of candidate ocean planets Trappist-1b, Kepler 93 b, Kepler 138 c, Kepler 114 c, Kepler 80 e, LHS 1140 c, Pi Mensae C and L98-59 d are such that they can maintain liquid water or high-pressure ices, then 21/52 or 40% of the planets surveyed here may be ocean planets. Of note is that these totals do not include cold exoplanets such as OGLE 2005-BLG-390-Lb [Beaulieu *et al.*, 2006], MOA-2007-

BLG-192 [Bennett *et al.*, 2008], or OGLE 2016-BLG-1195-Lb [Shvartzvald *et al.*, 2017] which were detected by gravitational microlensing. These worlds may be cold ocean planets in their own rights, perhaps containing ice-covered surfaces and internal oceans [Ehrenreich *et al.*, 2006; Ehrenreich and Cassan, 2007; Bond *et al.*, 2017].

Depending on their total internal heating rates, they could also be geologically active. Our results suggest that a significant number of exoplanets that have been previously classified as terrestrial planets may instead be extrasolar oceans worlds that contain significant amounts of water and may exhibit cryovolcanism and/or ice tectonics at their surfaces. In our solar system, cryovolcanism on the moons of the giant planets serves as an important process that transports liquid water, energy and organics between their interiors and surfaces [Fagents, 2003; Manga and Wang, 2007; Lopes *et al.*, 2013; Postberg *et al.*, 2018]. In some cases, the circulation of cryomagmatic fluids could create transient habitable niches in the interiors of ice-covered worlds [Ruiz *et al.*, 2007]. Additionally, recent studies suggest that even planets that are mostly ice-covered might have substantial amounts of unfrozen land near their equators where life could flourish [Paradise *et al.*, 2019]. The significant number of candidate cold ocean planets listed in Table 1 suggests that it is important to consider the possibility of habitable environments on terrestrial exoplanets that exist beyond the snowline. Cryovolcanic eruptions on cold ocean planets could be detected by next-generation telescopes as transient or periodic excesses in H<sub>2</sub>O, O<sub>2</sub>, and/or H in transit spectra that display spatial variability [Quick *et al.*, 2017a], or are localized to one hemisphere, as is the case for the cryovolcanically active moons in our solar system [Hurford *et al.*, 2007; Hedman *et al.*, 2013; Roth *et al.*, 2014a,b; Rhoden *et al.*, 2015; Sparks *et al.*, 2016; 2017; Teolis *et al.*, 2017]. If traces of cryovolcanic activity on exoplanets could be detected in transit spectra, this activity could be used as an indicator of which planets have substantial amounts of internal energy and water, both of which are necessary ingredients for life.

**Figure 1.** Explosive cryovolcanism, in the form of geyser-like plumes, has been detected on several of the icy moons of the giant planets in our solar system. **(a)** Plumes erupting at the south pole of Saturn’s moon Enceladus, as imaged by the Cassini spacecraft. Image Credit: NASA/JPL-Caltech/SSI. **(b)** The Hubble Space Telescope’s STIS instrument detected plumes at the south pole of Jupiter’s moon Europa. The pixilated blocks represent the locations where water vapor ejected during eruptions was detected spectroscopically. Image Credit: NASA/ESA/L. Roth/SWRI/University of Cologne.

**Figure 2. (a)** This smooth, semi-circular feature on the surface of Jupiter’s moon Europa is believed to be a cryolava flow that was emplaced during an effusive cryovolcanic eruption on the icy moon. **(b)** Cryolava flows on Neptune’s moon Triton. Given their densities, effective temperatures, and the estimated magnitude of total internal heating for each, the surfaces of Trappist-1g and Trappist-1h may be covered in cryolava flows.

**Figure 3.** Based on rates of volcanic activity on Earth, and the rates of cryovolcanic activity on Europa and Enceladus, all bodies on or above the pink line connecting these worlds are likely to have high rates of (cryo)volcanic activity. Worlds on or above the blue dotted line will also be volcanically active, albeit at lower rates. Planetary bodies below the blue dotted line may be geologically dead. Based on this plot, the high rates of internal heating on cold, low-density planets such as Trappist-1h and Kepler 138 b suggest that they are likely to exhibit cryovolcanism similar to what has been observed on Europa and Enceladus. Commensurate with previous investigations, our results suggest that CoRoT-7b’s high internal heating rates would induce volcanic activity of a magnitude similar to what has been observed on Jupiter’s moon Io.

**Figure 4.** The internal structures of cold ocean planets may be similar to the internal structures of our solar system’s icy moons. The interior structure of Trappist-1h **(a)** [Barr *et al.*, 2018; Luger *et al.*, 2017] may be similar to the interior structure of icy bodies like Jupiter’s moon Europa **(b)** [Schubert *et al.*, 2004; 2009].

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**Table 1. Solar System and Extrasolar Planetary Parameters**

World	$R_p$ (R <sub>E</sub> )	$M_p$ (M <sub>E</sub> )	$\rho$ (kg/m <sup>3</sup> )	$g$ (m/s <sup>2</sup> )	$T_{\text{EFF}}/T_s$ † (K)	System Age (Gyr)	$H_{\text{Radiogenic}}$ (W)	$H_{\text{Tidal}}$ (W)	$H_{\text{Total}}$ (W)	Type	Refs.
Earth	1	1	5500	9.8	255 ( $T_{\text{EFF}}$ ) 288 ( $T_s$ )	4.5	$4 \times 10^{13}$	# —	$4.7 \times 10^{13}$	rocky	Turcotte, 1995; Turcotte and Schubert, 2002
Venus	0.82	0.95	5240	8.87	260 K ( $T_{\text{EFF}}$ ) 737 ( $T_s$ )	4.5	$2.91 \times 10^{13}$	—	$2.91 \times 10^{13}$	rocky	Turcotte, 1995; Schubert et al. 1997
Early Venus	0.82	0.95	5240	8.87	?	4	$2 \times 10^{14}$	—	$2 \times 10^{14}$	rocky	Turcotte, 1995
Io	0.286	0.015	3528	1.8	110	4.5	$3.7 \times 10^{11}$	$1 \times 10^{14}$	$1 \times 10^{14}$	rocky	Yoder and Peale, 1981; Schubert et al., 2004; Hussmann and Spohn, 2004
Europa	0.245	0.008	3000	1.31	100	4.5	$2.1 \times 10^{11}$	$1 \times 10^{12}$	$1.21 \times 10^{12}$	icy	Schubert et al., 2004; Hussmann and Spohn, 2004; <sup>a</sup> Chen et al., 2014; Quick and Marsh, 2015
Enceladus	0.0395	$1.8 \times 10^{-5}$	1610	0.11	75	4.5	$2.73 \times 10^8$	$1.6 \times 10^{10}$	$1.63 \times 10^{10}$	icy	Howett et al., 2011; <sup>a</sup> Chen et al., 2014
Triton	0.21	0.00359	2060	0.78	38	4.5	$7.12 \times 10^{10}$	$6.6 \times 10^{10}$	$1.37 \times 10^{11}$	icy	Gaeman et al., 2012; <sup>a</sup> Chen et al., 2014 and refs. therein
Callisto	0.378	0.018	1834	1.235	126	4.5	$3.2 \times 10^{11}$	$3.3 \times 10^9$	$3.2 \times 10^{11}$	icy	Moore and Schubert, 2003; <sup>a</sup> Chen et al., 2014 and (1)
Ganymede	0.413	0.025	1942	1.43	117	4.5	$4.67 \times 10^{11}$	$1.1 \times 10^{10}$	$4.78 \times 10^{11}$	icy	Moore and Schubert, 2003; Schubert et al., 2004; <sup>a</sup> Chen et al., 2014 and (1)
Titan	0.404	0.0225	2575	1.35	94	4.5	$4.11 \times 10^{11}$	$8.75 \times 10^{10}$	$4.99 \times 10^{11}$	icy	<sup>a</sup> Chen et al., 2014; Schubert et al., 1986; and (1)
Mercury	0.38	0.055	5427	3.7	440	4.5	$2.25 \times 10^{12}$	$1.4 \times 10^9$	$2.25 \times 10^{12}$	rocky	Peplowski et al., 2011; Makarov & Efroimsky, 2014; Ogawa, 2016**
Mars	0.533	0.107	3934	3.72	210	4.5	$2.74 \times 10^{12}$	$1 \times 10^9$	$2.74 \times 10^{12}$	rocky	Parro et al., 2017; Manga et al., 2019
The Moon	0.273	0.012	3344	1.62	215(equator) 104 (poles)	4.5	$3 \times 10^{11}$	—	$3 \times 10^{11}$	rocky	Siegler and Smrekar, 2014; Paige et al., 2016; Williams et al., 2017

Pluto	0.1868	$2.18 \times 10^{-3}$	1854	0.620	44	4.5	$5.32 \times 10^{10}$	—	$5 \times 10^{10}$	icy	Robuchon and Nimmo, 2011; McKinnon et al., 2016
Charon	0.095	$2.7 \times 10^{-4}$	1702	0.288	53	4.5	$4.4 \times 10^9$	—	$4.4 \times 10^9$	icy	Hussmann et al., 2006; Cook et al., 2007
Ceres	0.074	$1.5 \times 10^{-4}$	2160	0.28	150	4.5	$4.5 \times 10^9$	—	$4.5 \times 10^9$	ice/rock hybrid	Castillo-Rogez et al., 2019
Miranda	0.037	$1.1 \times 10^{-5}$	1200	0.079	60	4.5	$8.5 \times 10^7$	—	$8.5 \times 10^7$	icy	<sup>&amp;</sup> Chen et al., 2014
Ariel	0.091	$2.3 \times 10^{-4}$	1592	0.269	60	4.5	$3.6 \times 10^9$	—	$3.6 \times 10^9$	icy	<sup>&amp;</sup> Chen et al., 2014
Rhea	0.12	$3.9 \times 10^{-4}$	1233	0.264	99 (dayside) 53(nightside)	4.5	$3.3 \times 10^9$	—	$3.3 \times 10^9$	icy	<sup>&amp;</sup> Chen et al., 2014
Mimas	0.03	$6.3 \times 10^{-6}$	1148	0.064	64	4.5	$4 \times 10^7$	$2.9 \times 10^8$	$3.3 \times 10^8$	icy	<sup>&amp;</sup> Chen et al., 2014
Dione	0.09	$1.8 \times 10^{-4}$	1480	0.232	87	4.5	$2.46 \times 10^9$	—	$2.46 \times 10^9$	icy	<sup>&amp;</sup> Chen et al., 2014
55 Cancri e	1.91	8.08	6400	21.7	1958	10.2	$1.14 \times 10^{14}$	$4.28 \times 10^{21}$	$4.28 \times 10^{21}$	rocky	Demory et al., 2011; 2015; 2016
Corot 7 b	5.7	1.55	7500	23	1756	1.32	$2.17 \times 10^{14}$	0 ( $e = 0$ )	$2.17 \times 10^{14}$	rocky	Queloz et al., 2009; Barros et al., 2014; Stassun et al., 2017
GJ 1132 b	1.13	1.66	6300	12.9	529	5	$3.66 \times 10^{13}$	TBD: $e$ unknown	$3.66 \times 10^{13}$	rocky	Berta-Thompson et al., 2015; Bonfils et al., 2018; Dittmann et al., 2017b
HD 219134 b	1.6	4.74	6340	18	1015	11	$6.42 \times 10^{13}$	0 ( $e = 0$ )	$6.42 \times 10^{13}$	rocky	Gillon et al., 2017a
HD 219134 c	1.5	4.36	6950	18.7	782	11	$5.39 \times 10^{13}$	$2.17 \times 10^{16}$	$2.17 \times 10^{16}$	rocky	Gillon et al., 2017a
<sup>II</sup> HD 219134 f	1.31	unknown	unknown	unknown	522	11	$3.51 \times 10^{13}$	$1.42 \times 10^{14}$	$1.77 \times 10^{14}$	unknown	Gillon et al., 2017a
HD 3167 b	1.7	5.02	5600	17	1669	7.8	$9.49 \times 10^{13}$	0 ( $e = 0$ )	$9.49 \times 10^{13}$	rocky	Christiansen et al., 2017; Livingston et al., 2018
Kepler 10 b	1.45	4.6	8255	21	2130	10.6	$4.9 \times 10^{13}$	TBD: $e$ unknown	$4.91 \times 10^{13}$	rocky	Esteves et al., 2015
Kepler 21 b	1.61	5.09	6720	19	2025	3.03	$1.48 \times 10^{14}$	$2.6 \times 10^{17}$	$2.6 \times 10^{17}$	rocky	López-Morales et al., 2016
Kepler 36 b	1.46	4.45	7810	20	978	6.8	$6.6 \times 10^{13}$	$2 \times 10^{14}$	$2.7 \times 10^{14}$	rocky	Carter et al., 2012
Kepler 60 b	1.68	4.2	4620	14	unknown	5.1	$1.3 \times 10^{14}$	$9.4 \times 10^{14}$	$1.1 \times 10^{15}$	ocean planet	Gozdziowski et al. 2016; Jontoff-Hutter et al., 2016; Morton et al., 2016

Kepler 60 c	1.9	3.85	3060	10.4	unknown	5.1	$1.7 \times 10^{14}$	$3.81 \times 10^{15}$	$3.98 \times 10^{15}$	cold planet	Gozdzewski et al. 2016; Jontoff-Hutter et al., 2016; Morton et al., 2016
Kepler 60 d	1.99	4.16	2910	10.3	unknown	5.1	$1.97 \times 10^{14}$	$3.03 \times 10^{14}$	$5 \times 10^{14}$	cold ocean planet	Gozdzewski et al. 2016; Jontoff-Hutter et al., 2016; Morton et al., 2016
Kepler 62 c	0.54	unknown	unknown	unknown	578	7	$3.16 \times 10^{13}$	TBD: e unknown	$3.16 \times 10^{13}$	unknown	Borucki et al., 2013
Kepler 68 c	0.96	2.04	unknown	21.7	unknown	6.3	$1.94 \times 10^{13}$	0 (e = 0)	$1.94 \times 10^{13}$	unknown	Berger et al., 2018. Mills et al., 2019
Kepler 70 b	0.76	0.44	5528	7.5	unknown	unknown	TBD: system age unknown	TBD: e unknown	unknown	rocky	Charpinet et al., 2011
Kepler 70 c	0.87	0.66	5521	8.5	unknown	unknown	TBD: system age unknown	TBD: e unknown	unknown	rocky	Charpinet et al., 2011
Kepler 78 b	1.1	3.2	unknown	14	2250	0.75	$9.57 \times 10^{13}$	TBD: e unknown	$9.57 \times 10^{13}$	candidate rocky	Pepe et al., 2013; Stassun et al., 2017
Kepler 80 d	1.53	6.75	7040	28	720	2	$1.55 \times 10^{14}$	TBD: e unknown	$1.55 \times 10^{14}$	rocky	Muirhead et al., 2012; MacDonald et al., 2016
Kepler 80 e	1.6	4.1	3750	16	628	2	$1.8 \times 10^{14}$	TBD: e unknown	$1.8 \times 10^{14}$	candidate ocean planet	Muirhead et al., 2012; MacDonald et al., 2016
Kepler 93 b	1.6	3.2	4292	12	1037	6.6	$8.74 \times 10^{13}$	0 (e = 0)	$8.74 \times 10^{13}$	candidate ocean planet	Dressing et al., 2015; Stassun et al., 2017
Kepler 97 b	1.5	3.5	5440	16	unknown	8.4	$5.66 \times 10^{13}$	TBD: e unknown	$5.66 \times 10^{13}$	rocky	Marcy et al., 2014
Kepler 99 b	1.5	6.2	10900	28	unknown	1.5	$1.66 \times 10^{14}$	TBD: e unknown	$1.66 \times 10^{14}$	rocky	Marcy et al., 2014
Kepler 100 b	1.3	7.3	14250	42.5	unknown	6.5	$4.7 \times 10^{13}$	TBD: e unknown	$4.7 \times 10^{13}$	rocky	Marcy et al., 2014
Kepler 101 c	1.2	unknown	unknown	unknown	1412	5.9	$4.3 \times 10^{13}$	0 (e = 0)	$4.3 \times 10^{13}$	candidate rocky	Bonomo et al., 2014
Kepler 102 d	1.2	3.8	13270	28	unknown	1.4	$8.66 \times 10^{13}$	TBD: e unknown	$8.66 \times 10^{13}$	rocky	Marcy et al., 2014
Kepler 105 c	1.3	4.6	11200	27.3	997	3.5	$6.85 \times 10^{13}$	TBD: e unknown	$6.85 \times 10^{13}$	rocky	Everett et al., 2015; Jontoff-Hutter et al., 2016; Morton et al., 2016
Kepler 114 c	1.6	2.8	4039	11.3	508	2.7	$1.49 \times 10^{14}$	TBD: e unknown	$1.49 \times 10^{14}$	candidate ocean planet	Muirhead et al., 2012; Xie, 2014; Morton et al., 2016

Kepler 138 b	0.7	0.19	3020	3.8	unknown	4.7	$9 \times 10^{12}$	$5.33 \times 10^{11}$	$9.5 \times 10^{12}$	cold ocean planet	Morton et al., 2016; Almenara et al., 2018
Kepler 138 c	1.7	5.2	6100	18.3	398	4.7	$1.3 \times 10^{14}$	$3.2 \times 10^{13}$	$1.6 \times 10^{14}$	candidate ocean planet	Muirhead et al., 2012; Morton et al., 2016; Almenara et al., 2018
Kepler 138 d	1.7	1.2	1360	4	335	4.7	$1.3 \times 10^{14}$	$1.5 \times 10^{13}$	$1.4 \times 10^{14}$	ocean planet	Muirhead et al., 2012; Morton et al., 2016; Almenara et al., 2018
Kepler 186 b	1.07	unknown	unknown	unknown	579	4	$3.6 \times 10^{13}$	TBD: e unknown	$3.6 \times 10^{13}$	unknown	Muirhead et al., 2012; Quintana et al., 2014; Torres et al., 2015
Kepler 186 c	1.25	unknown	unknown	unknown	470	4	$5.7 \times 10^{13}$	TBD: e unknown	$5.7 \times 10^{13}$	unknown	Muirhead et al., 2012; Quintana et al., 2014; Torres et al., 2015
Kepler 186 d	1.4	unknown	unknown	unknown	384	4	$8 \times 10^{13}$	TBD: e unknown	$8 \times 10^{13}$	unknown	Muirhead et al., 2012; Quintana et al., 2014; Torres et al., 2015
Kepler 186 e	1.27	unknown	unknown	unknown	unknown	4	$6 \times 10^{13}$	TBD: e unknown	$6 \times 10^{13}$	unknown	Quintana et al., 2014; Torres et al., 2015
Kepler 186 f	1.2	unknown	unknown	unknown	unknown	4	$4.7 \times 10^{13}$	$1.1 \times 10^9$	$4.7 \times 10^{13}$	unknown	Torres et al., 2015
Kepler 406 b	1.4	6.4	11820	30	unknown	5.8	$6.8 \times 10^{13}$	TBD: e unknown	$6.8 \times 10^{13}$	rocky	Marcy et al., 2014
Kepler 406 c	0.85	2.7	24390	37	unknown	5.8	$1.4 \times 10^{13}$	TBD: e unknown	$1.4 \times 10^{13}$	rocky	Marcy et al., 2014
Kepler 414 b	1.7	3.5	3845	11.7	unknown	5.5	$1.14 \times 10^{14}$	TBD: e unknown	$1.14 \times 10^{14}$	ocean planet	Hadden and Lithwick, 2014; Morton et al., 2016
L 98-59 b	0.8	0.5	5365	7.6	unknown	1 <sup>+</sup>	$3.33 \times 10^{13}$	$5.72 \times 10^{17}$	$5.72 \times 10^{17}$	rocky	Kostov et al., 2019
L 98-59 c	1.35	2.4	5358	12.9	unknown	1 <sup>+</sup>	$1.6 \times 10^{14}$	$5.38 \times 10^{17}$	$5.38 \times 10^{17}$	rocky	Kostov et al., 2019
L 98-59 d	1.57	3.4	4826	13.5	unknown	1 <sup>+</sup>	$2.52 \times 10^{14}$	$1.68 \times 10^{17}$	$1.69 \times 10^{17}$	candidate ocean planet	Kostov et al., 2019
LHS 1140 b	1.73	6.98	7500	23	235	5 <sup>A</sup>	$1.31 \times 10^{14}$	$\approx 6.15 \times 10^{13}$	$1.93 \times 10^{14}$	cold ocean planet	Ment et al., 2019
LHS 1140 c	1.3	1.81	4700	10.8	438	5 <sup>A</sup>	$5.34 \times 10^{13}$	$\approx 4.38 \times 10^{18}$	$4.38 \times 10^{18}$	candidate ocean planet	Ment et al., 2019
Pi Mensae c	2.04	4.82	2970	11.3	1169.8	2.98	$3 \times 10^{14}$	0 (e = 0)	$3 \times 10^{14}$	rocky	Huang et al., 2018
Trappist-1 b	1.07	0.85	3400	7.3	400	8	$2.3 \times 10^{13}$	$7.02 \times 10^{17}$	$7.02 \times 10^{17}$	candidate	Gillon et al., 2017b;

										ocean planet	Wang et al., 2017
Trappist-1 c	1.04	1.38	7630	12.6	342	8	$2.12 \times 10^{13}$	$3.13 \times 10^{16}$	$3.13 \times 10^{16}$	ocean planet	Wang et al., 2017
Trappist-1 d	0.76	0.41	3950	7	288	8	$8.29 \times 10^{12}$	$2.3 \times 10^{13}$	$3.13 \times 10^{13}$	ocean planet	Gillon et al., 2017b; Wang et al., 2017
Trappist-1 e	0.9	0.64	1710	7.7	251	8	$1.39 \times 10^{13}$	$3.84 \times 10^{13}$	$5.23 \times 10^{13}$	cold ocean planet	Gillon et al., 2017b; Wang et al., 2017
Trappist-1 f	1.03	0.67	1740	6.2	219	8	$2.05 \times 10^{13}$	$2.31 \times 10^{13}$	$4.36 \times 10^{13}$	cold ocean planet	Gillon et al., 2017b; Wang et al., 2017
Trappist-1 g	1.11	1.34	2180	10.7	199	8	$2.57 \times 10^{13}$	$5.77 \times 10^{11}$	$2.63 \times 10^{13}$	cold ocean planet	Wang et al., 2017
Trappist-1 h	0.72	0.086	1270	1.6	167	8	$6.99 \times 10^{12}$	$3.03 \times 10^{12}$	$1 \times 10^{13}$	cold ocean planet	Wang et al., 2017

In the case of non-transiting planets, detected solely by the radial velocity technique, where only planet mass in terms of  $m \sin i$  is known, we have assumed that planet mass,  $M_p$ , is equal to the minimum mass. Planet density,  $\rho$ , has been calculated for planets with an unknown bulk density using the relation:  $\rho = \frac{3M_p}{4\pi R_p^3}$ . Similarly, surface gravity,  $g$ , has been calculated for planets

for which surface gravity is unknown using  $g = \frac{GM_p}{R_p^2}$

<sup>†</sup> $T_S$  corresponds to the mean surface temperature of the planets and moons in our solar system.

<sup>#</sup>Negligible tidal heating rate

<sup>\*</sup> $H_{Tidal}$  for Triton corresponds to obliquity tides in the subsurface ocean, which contribute more internal energy than solid body tides. See *Chen et al.* [2014].

<sup>\*\*</sup> $H_{Radiogenic}$  for Mercury calculated using internal heating rate from *Peplowski et al.* [2011] and mantle mass from *Ogawa et al.* [2016]

<sup>&</sup>We have assumed that moons do not have a fully deformable interior (i.e.,  $k_2 \neq 1.5$ )

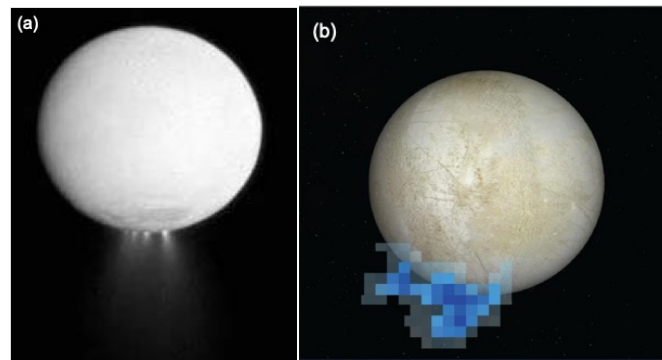
<sup>□</sup>Smallest possible value of  $R_p$  from *Gillon et al.* [2017]

<sup>+</sup>Minimum stellar age from Kostov et al., 2019

<sup>^</sup>Minimum stellar age from Ment et al., 2019

<sup>a</sup>Eccentricity is poorly constrained; maximum eccentricity values substituted into (1).  $H_{Tidal}$  values are preliminary.

**Figure 1**



**Figure 2**

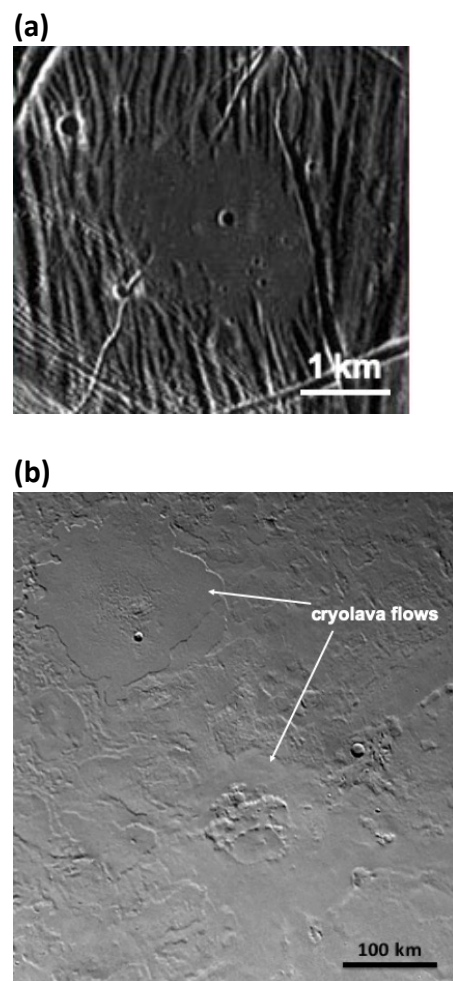
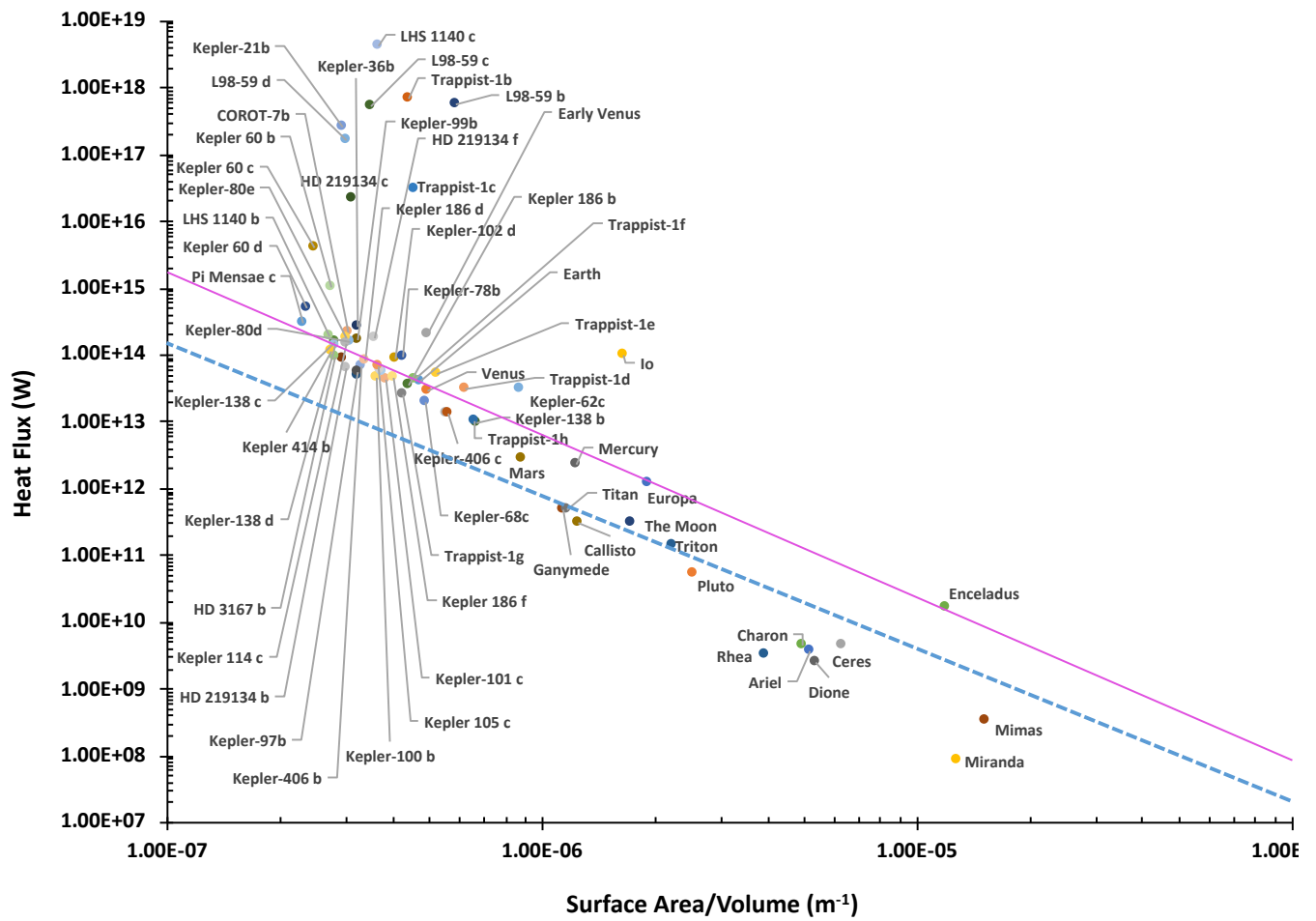


Figure 3



**Figure 4**

