

1 **Changes in Tropical Clouds and Atmospheric Circulation**
2 **Associated with Rapid Adjustment Induced by Increased**
3 **Atmospheric CO₂ – A Multiscale Modeling Framework Study**
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28 **Abstract**

29
30 The radiative heating increase due to increased CO₂ concentration is the primary source
31 for the rapid adjustment of atmospheric circulation and clouds. In this study, we investigate the
32 rapid adjustment resulting from doubling of CO₂ and its physical mechanism using a multiscale
33 modeling framework (MMF). The MMF includes an advanced higher-order turbulence closure in
34 its cloud-resolving model component and simulates realistic shallow and deep cloud climatology
35 and boundary layer turbulence. The rapid adjustment over the tropics is characterized by 1)
36 reduced ascent and descent strengths over the ocean, 2) increased lower tropospheric stability
37 (LTS) over the subsidence region, 3) shoaling of planetary boundary layers over the ocean, 4)
38 increased deep convection over lands and shift of cloud coverage from the ocean to lands, and 5)
39 reduced sensible (SH) and latent heat (LH) fluxes over the oceanic deep convective regions.
40 Unlike conventional general circulation models and another MMF, a reduction in the global-
41 mean shortwave cloud radiative cooling is not simulated, due to the increase in low clouds at
42 lower altitudes over the ocean, resulting from reduced cloud-top entrainment due to strengthened
43 inversion. Changes in regional circulation play a key role in cloud changes and shift of cloud
44 coverage to lands. Weaker energy transport resulting from water vapor and cloud CO₂ masking
45 effects reduces the upward motion and convective clouds in the oceanic regions. The ocean-land
46 transports are linked to the partitioning of surface SH and LH fluxes that increases humidity over
47 lands and enhances deep convection over the tropical lands.

48 **Keywords:** Rapid atmospheric adjustment . Tropical cloud changes . Rapid adjustment
49 mechanism . Low-level clouds . Increased atmospheric CO₂ . Multiscale modeling framework

51 **1 Introduction**

52 The climate response to increased CO₂ concentration in the atmosphere involves direct
53 and indirect effects; the direct effect is rapid adjustments to increased radiative heating while the
54 indirect effect is the slow response to the CO₂-caused change of surface air temperature (SAT)
55 (e.g., Andrews et al. 2010). The slow response is also referred to as the “temperature mediated”
56 change. Rapid adjustments act in a time scale of days and weeks, which is much shorter than that
57 of the slow response (Dong et al. 2009; Kamae and Watanabe 2012; Bony et al. 2013). As long
58 as the CO₂ concentration continues to rise, rapid adjustments will be a crucial contributor to the
59 climate change. It is important to understand rapid adjustments separately from slow response
60 because the changes due to rapid adjustments are often of opposite sign from those of the slow
61 response (e.g., Yang et al. 2003; Andrews et al. 2010; Bony et al. 2013; Sherwood et al. 2015).
62 For example, the global-mean precipitation decreases due to abrupt CO₂ increase but it increases
63 over the following years as the global-mean SAT increases in coupled ocean-atmosphere general
64 circulation model (GCM) simulations.

65 Changes in the atmospheric thermodynamic structure and regional circulation in response
66 to increased CO₂ concentration can cause changes in clouds and the associated changes in the
67 cloud radiative effects (CREs) at the top of the atmosphere (TOA). CREs are defined as the
68 differences in radiative fluxes between the clear and all skies. Zelinka et al. (2013) found that the
69 global-mean CRE change is $\sim 0.5 \text{ W m}^{-2}$ from doubling of CO₂ concentration in five GCM
70 simulations with prescribed sea surface temperature (SST), primarily from shortwave radiation.
71 However, the signs of CRE change vary among GCMs (Gregory and Webb 2008; Webb et al.
72 2013). The reason for the uncertainty in the global-mean CRE change is due to the uncertainty of
73 rapid cloud response, especially the response of low-level clouds (Andrews et al. 2012). The

74 diversity in the lower-tropospheric specific humidity due to changes in the regional atmospheric
75 circulations can explain most of the spreads in rapid cloud response and the associated CRE
76 change among GCMs (Kamae and Watanabe 2012; Kamae et al. 2015).

77 A plausible rapid adjustment mechanism that is applicable to the tropical ocean can be
78 described as follows. As the temperature over the oceanic regions increases in the troposphere,
79 the lower tropospheric stability (LTS) increases while the surface turbulent heat fluxes decrease.
80 The drying of the lower troposphere (in terms of relative humidity) reduces cloud condensate at
81 the cloud top, while the stabilization of the lower troposphere reduces the cloud top entrainment,
82 resulting in downshifts of the low clouds and shoaling of the planetary boundary layer (PBL)
83 (Wyant et al. 2012; Kamae et al. 2015). A consequence of this shoaling is a reduction in
84 shortwave (SW) cloud radiative cooling; that is, a positive change in the SW CRE in the
85 subsidence regions of the tropics and subtropics where low clouds are abundant. However,
86 conventional GCMs with parameterized physical processes produce diversity of low cloud
87 simulations (Soden and Vecchi 2011; Vial et al. 2013) and thereby contribute to the uncertainty
88 of rapid cloud adjustment and its physical mechanism.

89 In GCM simulations with prescribed SSTs, the global SAT cannot be held to fixed values
90 although both SST and sea ice concentration are fixed. The smaller heat capacity of the lands
91 allows the SAT to increase more than over the ocean due to the increased downward infrared
92 radiation from the heated troposphere. Such temperature contrasts between the land and oceanic
93 regions alter the regional atmospheric circulation and may shift the cloud cover at all levels from
94 the ocean to lands, which leads to the increase of precipitation over lands (Kamae and Watanabe
95 2013; Wyant et al. 2012; Zelinka et al. 2013; Kamae et al. 2016; Xu et al. 2017). Although the
96 regional circulation response to CO₂ direct radiative forcing is dominated by the CO₂ increase

97 over lands (Shaw and Voigt 2016), the changes in surface sensible heat fluxes due to increased
98 CO₂ concentration can also modulate changes in the regional atmospheric circulation and the
99 hydrological cycle (Kamae et al. 2015; Shaw and Voigt 2016; Xu et al. 2017). Additionally, the
100 plants react to increased CO₂ by closing the stomata and reducing the evapotranspiration of
101 moisture into the atmosphere. Such a response occurs at the same time scale as cloud adjustment,
102 therefore, should also impact the adjustment processes (Doutriaux-Boucher et al. 2009; Dong et
103 al. 2009; Andrews et al. 2012).

104 The direct CO₂ radiative forcing can be weakened by the masking effect of clouds and
105 water vapor in the deep convective regions but enhanced in the drier subtropics (Gregory and
106 Webb 2008; Merlis 2015). The differential heating resulting from the masking effect reduces
107 energy transport from the convective regions to the subsidence regions or higher latitudes and
108 thus reduces the regional atmospheric circulations. The extent of the weakening of regional
109 atmospheric circulation is related to the dependency of CO₂'s radiative forcing on the
110 climatology of clouds and humidity. Improving the representations of both shallow and deep
111 convective cloud processes is thus important for accurately capturing such masking effects and
112 the associated changes in regional atmospheric circulations, which can lead to a better
113 understanding the physical mechanisms of rapid cloud responses.

114 The complexity of subgrid effects associated with clouds, convection, precipitation and
115 radiation is the primary obstacle to improving model physical parameterizations in conventional
116 GCMs (Randall et al. 2003). The multiscale modeling framework (MMF) proposed by
117 Grabowski (2001) and Khairoutdinov and Randall (2001) is an attractive tool because it
118 explicitly simulates the largest and most organized circulations within deep convective systems
119 using a cloud-system resolving model (CRM) within each grid column of the global model. In

120 the past decade, MMF has been used to perform climate change simulations (Wyant et al. 2006,
121 2012; Arnold et al. 2014; Bretherton et al. 2014; Stan and Xu 2014; Xu and Cheng 2016; Xu et
122 al. 2017). The effective climate sensitivity (ECS) of MMFs is respectively 1.5 K in Wyant et al.
123 (2006), 2.1 K in Bretherton et al. (2014) and 2.0 K in Xu and Cheng (2016) assuming a CO₂
124 doubling forcing of 3.7 W m⁻² (Myhre et al. 1998), compared to 2.1-3.0 K for AMIP_4K
125 (Atmospheric Model Intercomparison Project +4 K SST) simulations by conventional GCMs
126 (Ringer et al. 2014).

127 In this study, we adopt the fixed SST approach using an MMF with an advanced higher-
128 order turbulence closure in its CRM component (Cheng and Xu 2011; Xu and Cheng 2013a).
129 This MMF is capable of simulating realistic shallow and deep cloud and water vapor climatology
130 and boundary layer turbulence (Xu and Cheng 2013a, b; Cheng and Xu 2013a, b; Painemal et al.
131 2015). We will separately examine the rapid adjustment over the tropical lands and ocean
132 resulting from doubling CO₂ levels except for the precipitation response. Xu et al. (2017)
133 compared the hydrological responses to both CO₂ increase and SST perturbation for two MMFs
134 with or without the higher-order turbulence closure. Similar to conventional GCMs, the land
135 surface temperature in MMF may vary regionally, resulting in differential changes in regional
136 circulation between lands and the ocean. Previous studies noticed less mid-tropospheric
137 subsidence over the tropical lands that may allow more convection and cloud formation in the
138 fixed SST experiments, compared to a fully coupled atmosphere-ocean GCM simulation
139 (Andrews et al. 2012; Wyant et al. 2012). Wyant et al. (2012) analyzed cloud and circulation
140 response of an MMF simulation to a CO₂ quadrupling, using a pair of 2-yr MMF simulations.
141 They found a shift of cloud cover and mean ascent from ocean to land regions, with little net
142 global change in cloud cover, and the shoaling of boundary-layer clouds in the subsidence

143 regions (by ~80 m), which was also identified in conventional GCMs (Kamae and Watanabe
144 2013).

145 The objectives of this study are twofold: one is to understand rapid adjustments resulting
146 from doubling of CO₂ in the MMF; and the other is to examine the rapid adjustment mechanisms
147 operating over lands and the ocean as well as different circulation regimes over the tropical
148 ocean. The results from this MMF will be further contrasted to those examined from
149 conventional GCMs and a MMF with a low-order turbulence closure in its CRM component.

150 **2 Model and experiment design**

151 Both the fixed-SST perturbation and coupled slab-ocean (full ocean) methods have been
152 used in conventional GCMs and MMF to study the response of clouds and atmospheric
153 circulations to CO₂ increase (e.g., Gregory and Webb 2008; Andrews et al. 2012; Wyant et al.
154 2012; Bretherton et al. 2014). However, the former method is more feasible than the latter
155 method for MMF, due to the lengthy spin-up time of an ocean model and computational cost of
156 an MMF relative to that of a conventional GCM (Wyant et al. 2012). Both methods seem to
157 agree on the estimated radiative forcings from the same GCM (Gregory and Webb 2008; Bala et
158 al. 2010; Colman and McAvaney 2011; Andrews et al. 2012) although for the same method the
159 spreads of the cloud radiative forcings are significant among GCMs (Kamae et al. 2015).

160 In this study, super-parameterized Community Atmosphere Model (SPCAM) MMF is
161 used. This MMF includes an intermediately prognostic higher-order turbulence closure (IPHOC;
162 Cheng and Xu 2006) in its CRM component, hereafter, referred to as SPCAM-IPHOC (Cheng
163 and Xu 2011; Xu and Cheng 2013a). The host GCM is Community Atmosphere Model (CAM)
164 version 3.5 (Collins et al. 2006). It has a horizontal grid spacing of $1.9^\circ \times 2.5^\circ$ and there are 32
165 layers in the vertical with 12 of them below 700 hPa. The high resolution below 700 hPa is used

166 to better resolve the structures of stratocumulus clouds, compared to 6 layers below 700 hPa in
167 the SPCAM configuration (i.e., with low-order turbulence closure in its CRM component) used
168 in Wyant et al. (2012) and Bretherton et al. (2014). The embedded CRM is the System for
169 Atmospheric Modeling (SAM) (Khairoutdinov and Randall 2003) with IPHOC. The CRMs have
170 the same vertical levels as the host GCM. All CRMs have 32 grid columns with 4 km of
171 horizontal grid spacing. Cloud microphysics and radiation are parameterized at the CRM scale.
172 Tendencies of heat and moisture from the CRM scale communicate to the large scale via the
173 GCM. The dynamical core provides the large-scale advective tendencies to the CRMs. The sub-
174 CRM-scale variability is parameterized with IPHOC (Cheng and Xu 2006), as described below.

175 The MMF was forced by specifying climatological SST and sea ice distributions from
176 Hadley Centre Sea Ice and Sea Surface Temperature dataset (HadISST) (Rayner et al. 2003) with
177 monthly-mean annual cycles. The MMF simulation was integrated for 10 years and 3 months
178 (Xu and Cheng 2013a). The results from the last nine years are analyzed in this study. This
179 simulation is referred to as control.

180 In the sensitivity experiment, the CO₂ concentration of present-day climate is
181 instantaneously doubled (Hansen et al. 1984) while the rest of experiment design are identical to
182 those of the control experiment. As mentioned earlier, the SST and sea ice are fixed but land
183 surface temperature is allowed to change in both experiments. This experiment was also
184 integrated for 10 years and 3 months. As in control, the results from the last nine years are
185 analyzed. This simulation is referred to as 2xCO₂.

186 The sub-CRM-grid-scale variability is represented by IPHOC. As detailed in Cheng and
187 Xu (2006, 2008), IPHOC assumes a joint double-Gaussian distribution of liquid water potential
188 temperature, total water, and vertical velocity. The properties of the double-Gaussian probability

189 density function (PDF) are determined from the first-, second-, and third-order moments of the
190 variables given above. And the PDF is used to diagnose cloud fraction and grid-mean liquid
191 water mixing ratio, as well as the buoyancy terms and fourth-order terms in the equations
192 describing the evolution of the second- and third-order moments. Xu and Cheng (2013a, b),
193 Cheng and Xu (2013a, b) and Painemal et al. (2015) extensively evaluated the control simulation
194 described above. A major benefit of the IPHOC scheme is that the MMF is able to simulate
195 optically thin clouds and realistic global cloud coverage and the stratocumulus in the subsidence
196 regions as in observations. Due to its prediction of three skewness variables, shallow cumulus
197 and its transition to stratocumulus are well simulated in this MMF. Further, the vertically-
198 integrated water vapor is highly corrected (>0.98) with observations with a root-mean-square
199 error less than 10% of the global mean (Xu and Cheng 2013a).

200 **3 Results**

201 In the results presented below, we will first discuss the geographic distributions and the
202 global means of selected variables. We will then focus on statistics computed separately over
203 tropical land and ocean regions (30°S-30°N), as well large-scale circulation regimes. Results
204 from this MMF will be compared with previous studies using SPCAM (Wyant et al. 2012;
205 Bretherton et al. 2014) and conventional GCMs with similar experimental configurations (e.g.,
206 Andrews et al. 2012; Kamae and Watanabe 2012; Zelinka et al. 2013; Kamae et al. 2015).

207 **3.1 Global change in surface air temperature (SAT), cloud cover and cloud radiative effects**

208 The global mean tropospheric radiative cooling rate is reduced as the CO₂ concentration
209 doubles, especially in the middle troposphere of the Tropics between 800 and 500 hPa (Figure
210 1b). The global-mean SAT increases by 0.14 K as the land surfaces heat up (Figure 1a), with the
211 tropical-land SAT increasing by 0.38 K (Table 1). The global-mean SAT increase is less than

212 half of the values reported in similar fixed-SST MMF studies with quadrupling CO₂
213 concentration; 0.30 K in Wyant et al. (2012) for 2-yr runs, 0.40 K for years 2-10 and 0.49 K for
214 years 2-30 of long integrations in Bretherton et al. (2014). The corresponding tropical-land SAT
215 increases are 0.50 K, 0.77 K and 0.80 K. The weak sensitivity of SAT increase to the integration
216 length after 10 years justifies the use of 10-yr integration performed in this study.

217 The spatial patterns of the land SAT increases are consistent with Wyant et al. (2012) and
218 Bretherton et al. (2014); that is, the SAT increases over most continental areas except for small
219 regions over southern/equatorial Africa, northern South America and western Australia (Figure
220 1a). The decrease over the Siberian region was not simulated by SPCAM, but by a conventional
221 GCM (Andrews et al. 2012). Significant changes in regional circulation are likely responsible for
222 these SAT decreases.

223 A major focus of this study is the rapid cloud adjustment. It is important to understand
224 the vertical-horizontal distributions of cloud changes between the 2xCO₂ and control
225 experiments. The changes in low-, middle-, high-level and total cloud amounts are shown in
226 Figure 2. The low-level clouds are defined as those with tops between the surface and 700 hPa,
227 middle-level clouds between 700 hPa and 400 hPa, and high-level clouds between 400 hPa and
228 the model top. A maximum vertical overlap of cloud fractions between the chosen pressure
229 ranges is used to obtain cloud amount for each CRM grid column. Then, cloud amounts of all
230 CRM columns are horizontally averaged to obtain cloud amount over a GCM grid. The CRM
231 provides cloud fraction at each CRM gridpoint from the double-Gaussian PDF.

232 The global-mean total cloud amount increases by 0.62%, which is contributed by both
233 low- and high-level clouds (0.38% and 0.43%, respectively). One can see the large meridional
234 and zonal variations of cloud amount changes with local values up to 6% (Figure 2). Most of the

235 low-level cloud increases are contributed from the extratropical oceans (Figure 2a), which is the
236 largest meridional variation signal. The features described below contribute to the zonal
237 variations. The decreases of low-level cloud amount over lands such as South and North
238 America and Europe correspond well to the SAT increases, as may be explained by the lower
239 relative humidity (RH) due to surface temperature increase. Low-level cloud amount increases
240 over the rest of land areas where surface temperature experiences small changes (e.g., South Asia,
241 Africa and western Australia). A unique feature of the MMF simulations is the increase of low-
242 level clouds over the subtropics, in particular, the southeast Pacific and southwest and northeast
243 Atlantic Oceans, which will be further examined later. These increases also contribute to the
244 global-mean increase, which was not simulated in SPCAM (Wyant et al. 2012; Bretherton et al.
245 2014) or conventional GCMs (Zelinka et al. 2013).

246 The global mean middle-level cloud amount decreases by 0.16%, due to the stabilization
247 in the mid-troposphere caused by the CO₂ radiative heating (Figure 1b). This reduction is
248 simulated in all models (e.g., Colman and McAvaney 2011; Wyant et al. 2012; Zelinka et al.
249 2013), but the magnitude is smaller in this MMF. This is related to the fact that middle-level
250 cloud amount increases in some regions (Figure 2b) where there is weak large-scale vertical
251 motion in the control experiment (Figure 3a) such as Africa (except near the Equator), southern
252 Asia and southern Australia. The magnitudes of subsidence are reduced over the same regions in
253 the 2xCO₂ experiment (Figure 3b), resulting in the increase of both low- and middle-level clouds
254 (Figures 2a and 2b). This leads to an increase in albedo. But the small SAT increases over parts
255 of these regions are likely related to the increased CO₂ masking effect (Figure 1a).

256 The increase of high-level cloud amount is concentrated over the latitudinal bands over
257 the edges of the Tropics and the high latitudes. This meridional variation signal is more

258 pronounced than that of the zonal variation. The latter is associated with the cloud increases over
259 the land regions such as Africa, southern Asia, Australia and equatorial South America (Figure
260 2c), which are well correlated with the increases in upward vertical motion (Figure 3). The zonal
261 and meridional variations can be explained by circulation changes resulting from CO₂ and water
262 vapor masking effects. These effects act to oppose both the zonally symmetric Hadley circulation
263 and zonally asymmetric Walker circulation (Merlis 2015). The energy transport from tropical
264 convective regions to higher latitudes is reduced as the middle troposphere over the edges of the
265 subtropics becomes warmer from the CO₂ radiative heating so that the upper troposphere is
266 destabilized. But the transport from the land regions to the adjacent oceans increases as the land
267 SAT increases (Figure 1a). The weaker cloud masking effects over the subtropical oceans also
268 enhances the zonal energy transport from their adjacent lands. This, in turn, impacts the low
269 clouds over the subtropics by lowering their tops (Wyant et al. 2012; Kamae and Watanabe 2013)
270 but low-level cloud amount increases are also simulated in the 2xCO₂ experiment.

271 Before discussing the changes in CREs due to cloud adjustment, the effective radiative
272 forcing (ERF) is examined, which is the change in net TOA downward radiative flux without any
273 changes in global-mean SAT under the conditions of doubled CO₂ (Hansen et al. 2005). The
274 ERF is 3.85 W m⁻² for SPCAM-IPHOC, 3.48 W m⁻² for SPCAM (Bretherton et al. 2014) and
275 3.77 ± 0.45 W m⁻² for 13 CMIP5 (Coupled Model Intercomparison Project Phase 5) models
276 (Kamae and Watanabe 2012), respectively. Therefore, the ERFs of both SPCAM and SPCAM-
277 IPHOC lie within the range of CMIP5 models although they differ by 0.37 W m⁻². As shown in
278 Table 1, the changes in LW CRE are negative for SPCAM-IPHOC (-0.47 W m⁻²), SPCAM (-
279 0.88 W m⁻²) and CMIP5 models (-0.76 ± 0.11 W m⁻²). But the change in SW CRE for SPCAM-
280 IPHOC has an opposite sign (-0.39 W m⁻² vs. 0.32 W m⁻² for SPCAM and 0.62 ± 0.24 W m⁻² for

281 CMIP5 models). This result is related to the large increases in low-level clouds over the mid- and
282 high-latitude and subtropical oceans, resulting in stronger radiative cooling. Consequently, the
283 change in net CRE for SPCAM-IPHOC (-0.86 W m^{-2}) is more negative, compared to SPCAM ($-$
284 0.56 W m^{-2}) and CMIP5 models ($-0.15 \pm 0.26 \text{ W m}^{-2}$) (Kamae et al. 2015). A majority of CMIP5
285 and earlier-generation models simulated positive SW CRE changes, which were consistent with
286 reduced low-level clouds (Gregory and Webb 2008; Kamae and Watanabe 2012).

287 The global-mean CRE changes contain both cloud adjustment and instantaneous cloud
288 masking effects. Utilizing an estimate of cloud masking effects from a conventional GCM (-0.56
289 W m^{-2} ; Andrews et al. 2012), cloud adjustment for SPCAM-IPHOC is -0.30 W m^{-2} , compared to
290 0.0 W m^{-2} for SPCAM. The difference in cloud adjustment between the two MMFs can largely
291 explain the difference in ERF (0.37 W m^{-2}).

292 The magnitudes of local changes in SW CRE can be as large as 10 W m^{-2} while those of
293 LW CRE can be as large as 5 W m^{-2} . The local changes in net CRE are similar to those of SW
294 CRE except for smaller magnitudes (Figure 4). Due to decreases of total cloud amount,
295 especially those of low-level clouds (Figures 2a and 2d), the net CRE changes are positive over
296 South and North Americas, and parts of Europe. Increases of low-level clouds in parts of the
297 subtropical and mid- and high-latitude oceans are associated with net CRE decreases (more
298 cooling). As the large-scale ascent is weakened (Figure 3) and clouds (especially, low-level
299 clouds) decrease (Figure 2), net CRE increases are simulated in tropical convective regions,
300 where net CRE increases due to cloud adjustment can be $1\text{-}2 \text{ W m}^{-2}$ higher than shown in Figure
301 4c due to strong CO_2 cloud masking effects (Andrews et al. 2012). Overall, the net CRE changes
302 are more pronounced over lands than over the ocean due to changes in regional circulation and
303 the associated shift of clouds from the ocean to lands (Wyant et al. 2012; Shaw and Voigt 2016).

304 **3.2 Changes in the thermodynamic and dynamic environments**

305 The rapid adjustment mechanism discussed in the introduction is characterized by 1)
306 reduction in the circulation strength, 2) stabilization of the lower troposphere, 3) reduction in
307 surface turbulent heat fluxes, 4) shoaling of the marine boundary layer and 5) reduction in SW
308 cloud radiative cooling (Colman and McAvaney 2011; Andrew et al. 2012; Kamae and
309 Watanabe 2012, 2013; Wyant et al. 2012; Kamae et al. 2015). A key aspect of this mechanism is
310 the decrease in the boundary-layer/cloud top entrainment, which is contributed through
311 strengthened trade inversion (Wyant et al. 2012). Unlike SPCAM and conventional GCMs, the
312 change in net CRE is negative, i.e., more cooling, in the southeast subtropical Pacific and
313 Atlantic regions, as well as the midlatitude storm track regions (Figure 4c). And the low-level
314 cloud amounts in these three regions increase rather than decrease in conventional GCMs
315 (Gregory and Webb 2008; Colman and McAvaney 2011; Kamae and Watanabe 2012; Zelinka et
316 al. 2013) except for one of the two GCMs in Watanabe et al. (2012). How much do the
317 remaining aspects of this mechanism differ from conventional GCMs and SPCAM? To answer
318 this question, we examine a few environmental variables, with an emphasis on the changes in
319 three regions; i.e., the tropical and subtropical oceanic subsidence regions (Figure 3a), the
320 tropical convective regions and the tropical lands. The changes in these variables are the primary
321 interest although the climatology of selected variables from the control experiment is also shown
322 in Figures 3 and 5-7.

323 The PBL depth/height is diagnosed from GCM grid variables based upon the bulk
324 Richardson number (Holtslag and Boville 1993). Change in the PBL depth is affected by
325 changes in both the stability and surface virtual heat flux. The lower tropospheric stability (LTS)
326 is defined as the difference in potential temperature between 700 hPa and the surface (Klein and

327 Hartmann 1993). Over the tropical and subtropical oceanic subsidence regions, there are small
328 decreases (and no increases anywhere) in the PBL depth, typically less than 40 m (Figure 5b).
329 The shoaling of marine boundary layer is thus consistent with previous studies (Wyant et al.
330 2012; Kamae and Watanabe 2013), but the largest decreases in the PBL depth occur near the
331 edges of the subtropics (30°S or 30°N). The decrease averaged over 30°S-30°N is ~20 m,
332 compared to 80 m reported in Wyant et al. (2012) for quadrupling CO₂. Note that 80 m was an
333 average over a certain LTS range. The small PBL depth decreases are accompanied by increased
334 low-level cloud amounts (Figure 2a) and strengthened trade inversions (Figure 5d). This result
335 can be explained by the high vertical resolution used in this MMF that reduces artificial
336 entrainment associated with a coarse resolution and limits the jump of cloud top to a smaller
337 height interval (Cheng et al. 2010; Xu and Cheng 2013a). The sharper inversion can lead to
338 simulation of more stratocumulus (with lower top) than cumulus clouds over these regions.

339 Changes in the large-scale subsidence ($\Delta\omega$) can impact the LTS and RH through the
340 subsidence warming (Figure 3b-d). In general, changes in ω (at 850, 700 and 500 hPa) are more
341 variable than those of the LTS. The large-scale subsidence in the strong subsidence regions
342 (ω_{500}) (Figure 3a), i.e., the southeast subtropical Pacific and Atlantic regions, increases slightly
343 in the 2xCO₂ experiment but its location shifts southwestwards from the strong subsidence
344 regions where the LTS increases (Figure 5d). The RH decreases over the subtropical low cloud
345 regions and the edges of the subtropical oceans except for the surface (Figures 6b-d). This is
346 directly related to the CO₂ radiative heating as cloud masking effects are weak there. Kamae and
347 Watanabe (2012) related the intermodel (conventional GCM) differences in the SW CRE change
348 due to rapid cloud adjustment to those in RH of the subsidence regions. However, the RH
349 decreases do not cause a reduction in low-level cloud amount there, due perhaps to a better

350 simulation of subgrid-scale turbulence and cloud processes within the embedded CRMs than
351 with parameterizations in conventional GCMs, which rely on RH to parameterize cloud amount.
352 The RH decreases are also consistent with increased surface latent heat fluxes (Figure 7d).

353 Over most of the tropical lands where surface sensible heat fluxes increase (Figure 7b),
354 the PBL depths increase, up to 100 m. The exceptions are the South and central Africa, Australia
355 and part of northwest South America where sensible heat fluxes decrease but latent heat fluxes
356 increase (Figure 7d). The LTS decreases are due to the strong increases in the SAT discussed
357 earlier (Figure 1a). Large-scale ascent is enhanced over most of the tropical lands. This suggests
358 that the regional atmospheric circulation changes are an important component of rapid
359 adjustment, agreeing with previous studies (e.g., Wyant et al. 2012; Shaw and Voigt 2016). The
360 RH increases, except for below 700 hPa over North America and parts of South American, are
361 associated with increased deep convection and enhanced large-scale ascent (Figure 3).

362 Over tropical deep convective regions (Figure 3a), neither the LTS (Figure 5b) nor the
363 PBL depth (Figure 5d) changes much as both are more relevant measures of variability for
364 shallow clouds than deep convection. Large-scale ascent is more likely slightly reduced than
365 enhanced, in particular, the Indian Ocean and northwestern Pacific and western Atlantic (Figures
366 3b-d). Since RH changes are closely connected to $\Delta\omega$, there are more areas with decreases than
367 increases in RH and the absolute magnitudes of RH changes increase as the height increases, as
368 in those of $\Delta\omega$. Additionally, the LH fluxes are reduced in the convective regions to balance the
369 decreased precipitation (Figure 7c). The LH flux change plays an important role in the
370 hydrological sensitivity, which was discussed in Xu et al. (2017).

371 As discussed above, regional circulation changes are a driver for cloud adjustment and
372 changes. A statistical description of circulation changes can be helpful. The ω_{500} PDFs for both

373 control and 2xCO₂ experiments show similar skewed distributions with a peak at ~ 20 hPa day⁻¹
 374 (Figure 8a). Convective regimes correspond to $\omega_{500} < 0$ whereas subsidence regimes correspond
 375 to $\omega_{500} \geq 0$. The differences in ω_{500} PDFs between the two experiments confirm the weakening
 376 of circulations discussed earlier. This is evidenced by increases in power for ω_{500} between -60
 377 and 20 hPa day⁻¹ but decreases in power for the moderate and strong subsidence regimes ($\omega_{500} >$
 378 20 hPa day⁻¹) and smaller decreases for the strongly convective regimes (Figure 8b). Over the
 379 tropical lands, there is a clear shift in power from the subsidence to convective regimes, which
 380 agrees with Wyant et al. (2012). Over the tropical ocean, the weakening is indicated by increases
 381 in power for the weakly convective and subsidence regimes and by decreases for the strongly
 382 convective ($\omega_{500} < -40$ hPa day⁻¹) and subsidence ($\omega_{500} > 25$ hPa day⁻¹) regimes. The decrease
 383 for the strong subsidence regimes was, however, not simulated by Wyant et al. (2012).

384 The weakening of circulations can be looked in another way. We compute the average
 385 upward ($\bar{\omega}^\uparrow$) and downward ($\bar{\omega}^\downarrow$) pressure vertical velocities at four selected levels (850, 700,
 386 500 and 200 hPa) for the entire tropics, as well as their difference $I = \bar{\omega}^\downarrow - \bar{\omega}^\uparrow$, which represents
 387 the strength of tropical overturning circulations (Bony et al. 2013). The average changes in
 388 ascending and descending areas of the control experiment ($\Delta\bar{\omega}^\uparrow$ and $\Delta\bar{\omega}^\downarrow$) are also computed
 389 (Table 3), using the monthly composite data to exclude the changes due to interannual
 390 fluctuations. The largest weakening of $\bar{\omega}^\uparrow$ occurs at 700 and 500 hPa, but that of $\bar{\omega}^\downarrow$ occurs at
 391 850 and 700 hPa. The relative weakening, $\Delta I/I$, is similar for all four levels, ranging from 0.059
 392 at 200 hPa to 0.064 at 700 hPa. The magnitude of the weakening over the tropical ocean is larger
 393 than that of the entire tropics (Bony et al. 2013). This is evident from Figure 8 for ω_{500} .

394 **3.3 Rapid cloud adjustment over the tropical lands and ocean and circulation regimes**

395 We first examine the mean properties over the tropical lands and ocean (Table 1). The
396 mean low-, middle- and high-level cloud amounts over the tropical lands are significantly higher
397 (0.2 to 1.6%, absolute percentage) in the 2xCO₂ experiment compared to the control experiment,
398 so are cloud liquid water path (LWP) and ice water path (IWP). According to the definitions of
399 LWP and IWP, which are the vertical integrals and horizontal averages of cloud-water and
400 cloud-ice mixing ratios, respectively, changes in cloud amount also contribute to those of LWP
401 and IWP. Over the tropical ocean, low-level cloud amount does not change, high-level cloud
402 amount increases (0.7%) but middle-level cloud amount decreases (-0.2%). Since both LWP and
403 IWP decrease and the decrease in LWP (-2.1%) is larger than the amount that is attributed to the
404 decrease in middle-level clouds, oceanic clouds are optically thinner in terms of both in-cloud
405 LWP and IWP in the 2xCO₂ experiment. SW cloud radiative cooling over the tropical ocean is
406 reduced by 0.3 W m⁻² due to optically thinner clouds, but enhanced by 0.9 W m⁻² over the
407 tropical lands due to the increase of total cloud amount (especially, high-level clouds) and
408 optically thicker clouds (Table 1). LW cloud radiative heating also has opposite signs over the
409 two regions, i.e., -0.8 W m⁻² for the ocean and 0.5 W m⁻² for lands, reflecting the shift of large-
410 scale ascent and deep convection from the ocean to lands.

411 How well do these results compare to those of SPCAM averaged over the same 9-year
412 period (Bretherton et al. 2014) after linearly scaling the 4xCO₂ changes to 2xCO₂ changes? The
413 signs of the changes agree between the two MMFs in ω_{500} , middle-level cloud amount, LWP,
414 IWP, precipitation, TOA and surface SW and LW radiative fluxes, SW and LW CREs and SH
415 for the tropical lands and the tropical ocean (Table 2). The disagreement appears in low-level
416 clouds for the tropical lands (+0.2% for SPCAM-IPHOC vs -0.5% for SPCAM), high and total
417 cloud amounts over the tropical ocean (+0.6-0.7% for SPCAM-IPHOC vs. -0.4% for SPCAM).

418 This disagreement does not alter the signs of CRE and TOA and surface radiative flux changes,
419 due to compensation of increases in cloud amount with decreases in in-cloud cloud optical depth
420 simulated in SPCAM-IPHOC. Therefore, the results from both MMFs largely support the
421 conceptual diagram of rapid tropical cloud changes proposed by Wyant et al. (2012).

422 Both MMFs have the same amount of precipitation decrease on the tropical ocean (-0.15
423 mm day⁻¹). But SPCAM-IPHOC simulates a larger precipitation increase over the tropical lands
424 than SPCAM (0.14 vs. 0.06 mm day⁻¹). This suggests that convection is more strongly enhanced
425 over lands in SPCAM-IPHOC. This result is explained by different partitioning of SH and LH
426 flux changes (Xu et al. 2017); i.e., small increases for both in SPCAM-IPHOC versus large
427 increase in SH but large decrease in LH in SPCAM (Table 2). DeAngelis et al. (2016) found that
428 different SH and LH partitioning in conventional GCMs also changes the ocean-land transports
429 and the intensity of convection over both the ocean and lands. The stronger land convection in
430 SPCAM-IPHOC also causes appreciable differences in the magnitudes (ω_{500} , LWP, IWP, LH,
431 SW and LW CRE, TOA SW and LW, surface LW) and the signs (low- and high-level, and total
432 cloud amounts, SH and surface SW) of the tropical-mean changes between the two MMFs.

433 Next, vertical profiles of a few variables are obtained for the tropical ocean and tropical
434 lands as in previous studies (Wyant et al. 2012; Kamae and Watanabe 2013) to further explain
435 the rapid cloud adjustment mechanism. We also obtain the mean vertical profiles for sub-regions
436 and circulation regimes. For the tropical lands, the desert region over Africa is separated from
437 the rest of tropical lands due to smaller cloud masking effects there. For the tropical ocean,
438 vertical profiles for five circulation regimes are further examined. The five circulation regimes
439 with equal size of population (20%) are obtained, based upon the distribution of annual-mean
440 ω_{500} of the control experiment (not shown). The ω_{500} ranges are $\omega_{500} < -21$ hPa day⁻¹ for “0%-

441 20%”, $-21 \leq \omega_{500} < -4$ hPa day⁻¹ for “20%-40%”, $-4 \leq \omega_{500} < 10$ hPa day⁻¹ for “40%-60%”, 10
442 $\leq \omega_{500} < 21$ hPa day⁻¹ for “60-80%” and $\omega_{500} \geq 21$ hPa day⁻¹ for “80%-100%.” The first two
443 correspond to strongly (“0%-20%”) and weakly (“20%-40%”) convective regimes while the last
444 three correspond to weak, moderate and strong subsidence regimes, respectively.

445 There are large contrasts in the response to doubling of CO₂ between the tropical ocean
446 and lands (Figure 9) and between desert and non-desert regions. The potential temperature
447 increases ($\Delta\theta$) over lands are 0.1-0.3 K higher than over the ocean while they are ~0.5 K higher
448 over deserts than over non-deserts (Figure 9b). The maximum increase ($\Delta\theta$) over the ocean
449 appears above the marine boundary layer top (~800 hPa) and $\Delta\theta$ decreases as the height
450 decreases. The negative $\Delta\theta$ in the upper troposphere can be explained by a reduced meridional
451 heat transport and less condensational heating resulting from weaker deep convection. The
452 increased stability between the surface and ~800 hPa is consistent with the increased low-level
453 clouds simulated in SPCAM-IPHOC (Figure 9d). Over lands, $\Delta\theta$ increases in the lower
454 troposphere (0.2-0.8 K) are much larger than in the upper troposphere (0.1-0.2 K), especially
455 over deserts. The increased instability over lands corresponds well with enhanced upward motion
456 (up to -2.5 hPa day⁻¹). Over deserts, $\Delta\omega$ is “bottom heavy,” which is related to the large
457 instability increase in the lower troposphere (Figure 9b). However, the upward motion over the
458 ocean is slightly reduced (~0.5 hPa day⁻¹), especially in the upper troposphere (Figure 9c).

459 Over the tropical ocean, sum of cloud water and cloud ice mixing ratios ($q_c + q_i$), i.e., in-
460 cloud value multiplied by cloud fraction, increases over a thin layer immediately below the
461 boundary layer top (~825 hPa) but decreases rapidly above the boundary layer top. The
462 magnitude of the largest decrease (at 760 hPa) is twice as large as that of the largest increase (at
463 860 hPa). This feature is related to the shoaling of the boundary layer. The decrease associated

464 with weakened deep convection is much smaller above 600 hPa. Conventional GCMs also show
465 an increase of marine boundary cloud fraction at below ~900 hPa (Zelinka et al. 2013) instead of
466 below 825 hPa in this study. This difference may be related to either stronger downward
467 radiative heating resulted from quadrupling CO₂ or artificially strong entrainment resulted from
468 the coarser resolution in the lower troposphere used in GCMs.

469 Over the tropical lands, $q_c + q_i$ increases throughout the troposphere except for the lower
470 troposphere of deserts (Figure 9d), indicating that deep convection is mostly enhanced. Over
471 deserts, low-level clouds are reduced, which is related to the drier lower troposphere in terms of
472 RH (Figure 9f), but deep convection is slightly enhanced. The latter is related to the slight
473 increases of upward motion in the upper troposphere. Over non-deserts, both RH and $q_c + q_i$
474 increase over the entire troposphere, indicating that both low- and high-level clouds are enhanced.
475 In the lower troposphere, there are larger moisture increases (Figure 9e) and smaller temperature
476 increases (Figure 9b) over non-deserts, compared to over deserts. These results are related to the
477 efficient turbulent mixing simulated in this MMF. This is also supported by the small increases
478 in surface LH and SH fluxes, compared to SPCAM (Table 2). Thus, SPCAM-IPHOC simulates
479 enhanced low clouds, except for those below 980 hPa, while SPCAM simulates reduced low
480 clouds, as the land surface warms up less over non-deserts.

481 Both specific humidity (Δq) and RH (ΔRH) decrease above the shoaled boundary layer
482 top of the tropical ocean. The RH increase below the shoaled boundary layer top is likely related
483 to reduced cloud top entrainment due to strengthened inversion. To further understand this, the
484 RH changes (ΔRH) are partitioned into two components; i.e., one is due to temperature change,
485 $\Delta RH(\Delta T)$, and the other due to specific humidity change, $\Delta RH(\Delta q)$, so that $\Delta RH \approx \Delta RH(\Delta T) +$
486 $\Delta RH(\Delta q)$, following Kamae and Watanabe (2012). The temperature and specific humidity are

487 set to be the values in the control experiment to calculate $\Delta RH(\Delta q)$ and $\Delta RH(\Delta T)$, respectively.
488 As seen from Figure 9g and 9h, $\Delta RH(\Delta T)$ and $\Delta RH(\Delta q)$ over the tropical ocean have opposite
489 signs in the vertical except between 600 and 800 hPa, with $\Delta RH(\Delta T)$ dominating in the upper
490 troposphere and $\Delta RH(\Delta q)$ in the lower troposphere. Between 800 hPa and 600 hPa, both act to
491 reduce RH. In contrast, vertical variations of ΔRH over lands are mostly explained by $\Delta RH(\Delta T)$
492 except for the middle and upper troposphere over deserts.

493 Figure 10c shows that the weakening of regional circulations over the tropical ocean seen
494 from Figure 9c is achieved through a large reduction in the ascent strength of convective regimes
495 (up to 2 hPa day⁻¹) and a small reduction in the descent strength in the low and middle
496 troposphere of moderate and strong subsidence regimes (up to 1 hPa day⁻¹). The ascent strength
497 in the low and middle troposphere of the strongly convective regime is reduced more greatly
498 than that of the weakly convective regime, but the descent strength of the moderate subsidence
499 regime is reduced more greatly than that of the strong subsidence regime.

500 Although the five circulation regimes have various characteristics in the changes in
501 thermodynamic/cloud/radiative profiles between the 2xCO₂ and control experiments, the strong
502 subsidence regime is most distinct and a major contributor to the tropical-ocean mean cloud
503 changes (Figure 9d), for example, the increase of condensate below 800 hPa. Further, this regime
504 has the largest stability increase between the surface and 800 hPa (Figure 10b) and the largest
505 change in radiative heating/cooling rate, (Figure 10a) but condensate reduction is minimal above
506 800 hPa (Figure 10d). This result is related to the weak cloud masking effects above the PBL top.
507 For other four regimes, reduction in condensate appears in the layer above 840 hPa, but the
508 strongly convective regime shows the greatest reduction in the vertical extent and amount of
509 condensate (Figure 10d). Therefore, the large condensate reduction in the tropical-mean profile

510 between 840 and 600 hPa (Figure 9d), i.e., the shoaling of the PBL, is contributed by all regimes
511 except for the strong subsidence regime.

512 Although the vertical variations of moisture changes (Figures 10e-h) resemble those of
513 the tropical-mean changes (Figures 9e-h), the magnitudes of moisture changes have large
514 diversity in the middle/upper troposphere (<800 hPa) among the five regimes. The moisture
515 increases from the surface to 800 hPa are related to weakened vertical moisture transport, which
516 may be caused by strengthened inversion (Figure 10b) and/or reduced surface evaporation
517 (Figure 7d). Above 600 hPa, RH changes are small for all regimes except for the strong
518 subsidence regime, due to opposite effects and temperature and moisture changes. For the strong
519 subsidence regime, both $\Delta RH(\Delta T)$ and $\Delta RH(\Delta q)$ act to reduce RH throughout the troposphere.
520 This is only true for the layer between 800 and 700 hPa for the other four regimes. Above 500
521 hPa, the magnitudes of $\Delta RH(\Delta T)$ are slightly larger than those of $\Delta RH(\Delta q)$, which are related to
522 the monotonic increase of negative temperature changes.

523 **4 Summary and discussions**

524 In this study, we have investigated rapid adjustment resulting from doubling of
525 atmospheric CO₂ concentration and its physical mechanism using a multiscale modeling
526 framework (MMF). The MMF includes an advanced higher-order turbulence closure in its CRM
527 component. It simulates realistic shallow and deep cloud climatology and boundary-layer
528 turbulence, in comparison with conventional GCMs and SPCAM. This ability is important for
529 simulating rapid adjustment because of CO₂-induced cloud and water vapor masking effects that
530 depend on cloud climatology and their impacts on regional circulation (Merlis 2015).

531 Although the simulated global cloud distributions from this MMF show a decrease in
532 middle-level cloud amount as in conventional GCMs and SPCAM, low-level cloud amount

533 increases slightly, as opposed to decreases in conventional GCMs and SPCAM. This is due to
534 increases over the extratropical and subtropical oceans. As in conventional GCMs, high-level
535 cloud amount also increases, due to increases over the edges of the Tropics and several land
536 regions. The oceanic clouds in the 2xCO₂ experiment are optically thinner than those in the
537 control experiment. Their counterparts over lands are optically thicker. Both aspects agree with
538 conventional GCMs. However, the increases in low-level clouds simulated in this MMF result in
539 global-mean SW cloud radiative cooling instead of warming simulated by most conventional
540 GCMs and SPCAM. The change in net CRE is more negative than in previous studies (Gregory
541 and Webb 2008; Kamae et al. 2015). The cloud adjustment effect, which is the net CRE change
542 without cloud masking effect, differs by 0.30 W m⁻² between the two MMFs, which largely
543 explain their difference in effective radiative forcing (ERF).

544 This study adds distinct contributions to the rapid adjustment mechanism from different
545 circulation regimes. This modified mechanism based upon the results from this advanced MMF
546 is characterized by the following elements: 1) reduced ascent and descent strengths over the
547 ocean with the largest reduction for the strongly convective and moderate subsidence regimes,
548 respectively, 2) increased lower tropospheric stability (in particular, between the surface and 800
549 hPa) over the subsidence region, which strengthens the trade inversion and reduces the
550 entrainment despite an overall reduction in the subsidence strength, 3) shoaling of planetary
551 boundary layers over the ocean with the largest reduction in the PBL depth away from the
552 tropical and subtropical regions and the smallest reduction for the strong subsidence regime, 4)
553 increased strengths of ascent and deep convection over lands and shift of cloud coverage from
554 the ocean to lands, particularly over non-deserts, and 5) reduction in the SH and LH fluxes over
555 the oceanic deep convective regions but small increases of both SH and LH fluxes over lands.

556 In this study, a decrease in the global-mean SW cloud radiative cooling is not simulated,
557 which is a key outcome of the rapid adjustment mechanism in earlier studies with conventional
558 GCMs and SPCAM. The shoaling of the PBL is not as pronounced as in earlier studies, in
559 particular, over the strong subsidence regions where low-level clouds prevail. The MMF
560 simulates an increase in low-level cloud fraction/condensate at a slightly lower altitude. This is
561 related to the CO₂-induced radiative heating above the PBL top, which increases the stability and
562 reduces the entrainment. This means that the embedded CRM with IPHOC more likely simulates
563 stratocumulus clouds rather than cumulus clouds. The high vertical resolution in the lower
564 troposphere used in this MMF also helps to reduce artificial entrainment and minimizes the
565 extent of the PBL shoaling (Cheng et al. 2010; Xu and Cheng 2013a).

566 Changes in regional circulation play a key role in influencing the different cloud changes
567 between the ocean and lands. Weaker energy transport resulting from CO₂ cloud and water vapor
568 masking effects in the oceanic regions with strong large-scale ascent reduces the upward motion
569 and convective clouds. Over lands, large-scale upward motion is enhanced throughout the entire
570 troposphere, accompanied by increased moisture related to efficient turbulent mixing. Due to
571 differences in surface warming between desert and non-desert regions, low-level clouds over
572 non-deserts increases rather than decreases over deserts. Deep convection is also enhanced more
573 over non-deserts. Large increases of humidity over lands are related to the land-ocean transports
574 that are linked to the partitioning of surface SH and LH fluxes (DeAngelis et al. 2016). This
575 MMF simulates small increases in both SH and LH over the tropical lands, compared to large
576 increases in SH but large decreases in LH in SPCAM.

577 The rapid adjustment mechanism outlined in this study may need further investigation
578 because some of the differences from previous studies may be related to nonlinearity in the

579 responses between doubling and quadrupling of CO₂ and to the coarser vertical resolutions used
580 in other studies. However, the MMF results presented in this study will be helpful to further
581 advance our understanding of rapid adjustment simulated in conventional GCMs if the
582 differences between the parameterized and explicitly simulated cloud processes can be
583 understood.

584

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Table 1. The average values over the entire tropics, tropical ocean, tropical land and entire globe for the control experiment as well as the differences between the 2xCO₂ and control experiments. For ω_{500} , the percentage change is the strength of tropical overturning circulation shown in Table 3.

	Control				2×CO ₂ - Control				
	Tropics			Globe	Tropics			Globe	Tropics Total %
	Ocean	Land	Total		Ocean	Land	Total		
T _s (K)	299.22	296.19	298.42	287.85	0.00	0.38	0.10	0.16	0.03
ω_{500} (hPa day ⁻¹)	-1.76	3.21	-0.46	0.08	0.59	-2.16	-0.13	0.00	6.02
Low cloud (%)	43.6	20.6	37.5	45.1	0.0	0.2	0.0	0.38	0.12
Middle cloud (%)	8.4	11.8	9.3	17.2	-0.2	0.4	-0.1	-0.16	-0.63
High cloud (%)	43.9	25.9	39.1	29.2	0.7	1.6	1.0	0.43	2.43
Total cloud (%)	68.3	40.5	61.0	61.7	0.6	1.2	0.7	0.62	1.21
LWP (g m ⁻²)	111.4	70.5	100.7	98.2	-2.1	1.4	-1.2	0.13	-1.15
IWP (g m ⁻²)	39.2	33.2	37.6	48.3	-0.3	2.5	0.5	0.11	1.21
Rain (mm day ⁻¹)	4.02	2.67	3.67	2.86	-0.16	0.14	-0.08	-0.06	-2.23
LHF (W m ⁻²)	150.4	60.6	126.8	88.5	-2.3	0.4	-1.6	-0.95	-1.24
SHF (W m ⁻²)	15.6	48.9	24.4	23.4	-0.2	0.3	-0.1	-0.11	-0.33
SW CRE (W m ⁻²)	-55.8	-40.5	-51.8	-50.4	0.3	-0.9	0.0	-0.39	0.07
LW CRE (W m ⁻²)	26.6	20.4	25.0	22.9	-0.8	0.5	-0.4	-0.47	-1.75
SW TOA (W m ⁻²)	306.8	282.2	300.4	240.4	0.3	-0.7	0.0	-0.30	0.01
LW TOA (W m ⁻²)	259.7	264.6	261.0	240.3	-3.9	-5.0	-4.2	-3.55	-1.61
SW SFC (W m ⁻²)	213.9	187.3	207.0	161.8	0.1	-1.3	-0.2	-0.55	-0.12
LW SFC (W m ⁻²)	55.9	78.5	61.8	57.5	-2.5	-1.5	-1.7	-1.90	-2.81

Table 2. The average differences over the entire tropics, tropical ocean and tropical lands between the xCO₂ and control experiments from the ensemble of CMIP5 models (Kamae and Watanabe 2012), SPCAM (Bretherton et al. 2014) and SPCAM-IPHOC.

	CMIP5 Ensemble			SPCAM			SPCAM-IPHOC		
	Ocean	Land	Tropics	Ocean	Land	Tropics	Ocean	Land	Tropics
ω_{500} (hPa day ⁻¹)	0.43	-0.96	0.08	0.60	-1.98	-0.06	0.59	-2.16	-0.13
Total cloud (%)	-0.4	-0.3	-0.4	-0.4	0.4	-0.2	0.6	1.2	0.7
Low cloud (%)				-0.1	-0.5	-0.2	0.0	0.2	0.0
Middle cloud (%)				-0.3	0.6	-0.1	-0.2	0.4	-0.1
High cloud (%)				-0.4	1.0	-0.1	0.7	1.6	1.0
LWP (g m ⁻²)				-2.7	0.9	-1.8	-2.1	1.4	-1.2
IWP (g m ⁻²)				-0.7	2.4	0.1	-0.3	2.5	0.5
Rain (mm day ⁻¹)				-0.16	0.06	-0.10	-0.16	0.14	-0.08
LHF (W m ⁻²)				-2.8	-2.7	-2.8	-2.3	0.4	-1.6
SHF (W m ⁻²)				-0.2	1.7	0.3	-0.2	0.3	-0.1
SW CRE (W m ⁻²)	0.7	0.4	0.6	1.1	-0.8	0.6	0.3	-0.9	0.0
LW CRE (W m ⁻²)	-1.1	-0.1	-0.8	-1.5	0.7	-0.9	-0.8	0.5	-0.4
SW TOA (W m ⁻²)				1.1	-0.4	0.7	0.3	-0.7	0.0
LW TOA (W m ⁻²)				-2.9	-5.0	-3.5	-3.9	-5.0	-4.2
SW SFC (W m ⁻²)				1.1	-0.7	0.7	0.1	-1.3	-0.2
LW SFC (W m ⁻²)				-1.0	-1.9	-1.2	-2.5	-1.5	-1.7

Table 3. Pressure vertical velocity (ω) and fractional area (σ) at 850, 700, 500 and 200 hPa levels for 30°S to 30°N that are obtained for different signs of vertical velocity, upward ($\bar{\omega}^\uparrow$, negative) and downward ($\bar{\omega}^\downarrow$, positive) of the control experiment. I represents the strength of tropical overturning circulation. Note that changes between the 2xCO₂ and control experiments ($\overline{\Delta\omega}^\uparrow$ and $\overline{\Delta\omega}^\downarrow$) are calculated over the same areas with upward and downward velocity of the control experiment. The monthly composite (averaged over the same month from years 2-9) data are used to eliminate the changes due to interannual fluctuations. Unit for pressure vertical velocity is hPa day⁻¹.

	850 hPa	700 hPa	500 hPa	200 hPa
$\bar{\omega}$	-0.62	-0.65	-0.47	-0.37
$\bar{\omega}^\downarrow$	23.43	23.26	22.00	12.41
$\bar{\omega}^\uparrow$	-32.00	-33.75	-34.60	-18.35
$I = \bar{\omega}^\downarrow - \bar{\omega}^\uparrow$	55.43	57.01	56.60	30.76
$\overline{\Delta\omega}^\downarrow$	-1.66	-1.73	-1.49	-0.78
$\overline{\Delta\omega}^\uparrow$	1.70	1.91	1.92	1.05
$\Delta I = \overline{\Delta\omega}^\downarrow - \overline{\Delta\omega}^\uparrow$	-3.36	-3.64	-3.41	-1.83
$\sigma(\omega^\downarrow)$	0.57	0.58	0.60	0.58
$\sigma(\omega^\uparrow)$	0.44	0.42	0.40	0.42

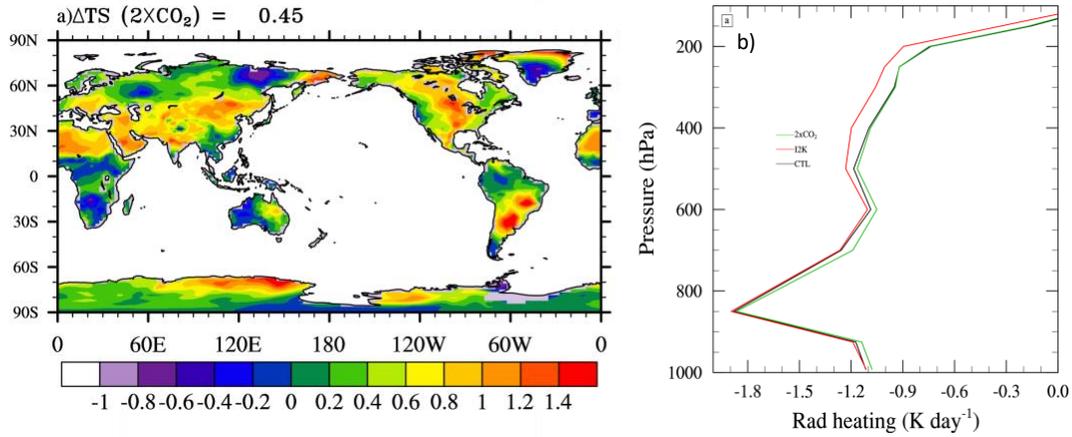
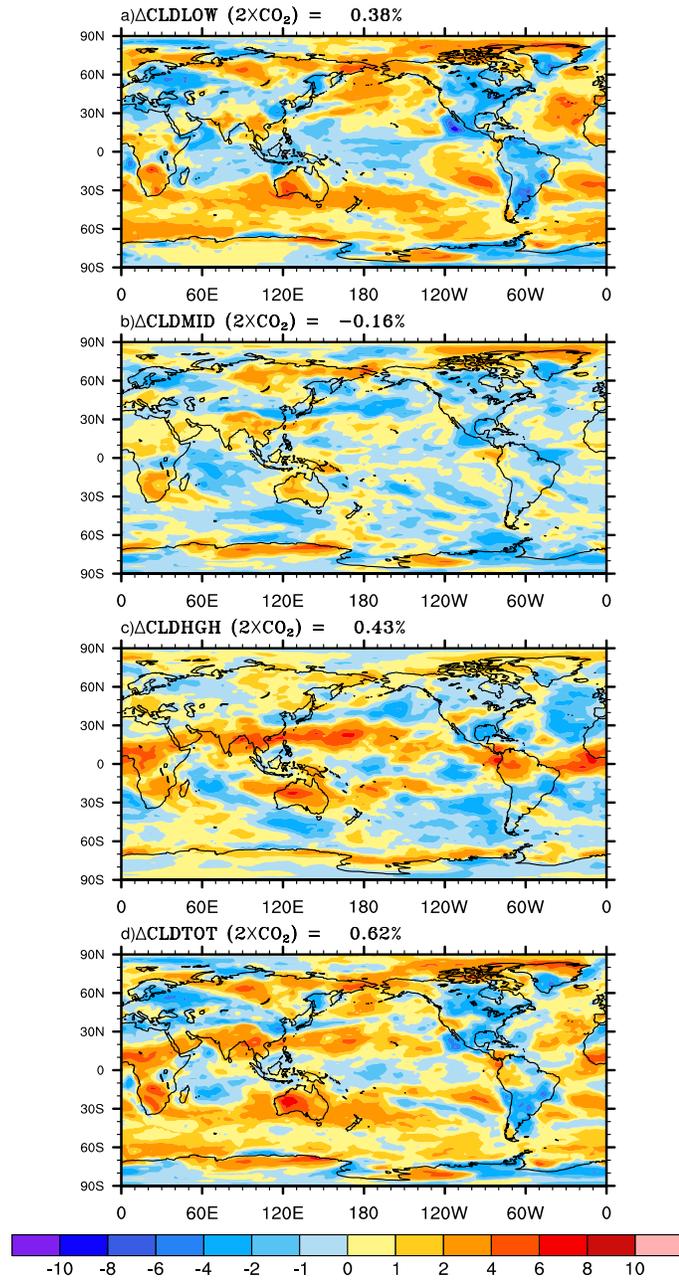


Figure 1. (a) Annual-averaged surface air temperature change between the $2\times\text{CO}_2$ and control experiments. (b) Annual- and tropical-mean profiles of radiative heating rates for the control (CTL; black), $2\times\text{CO}_2$ (green) and +2K SST (red) experiments. The +2K SST experiment is provided for reference although it is discussed in Xu and Cheng (2016) and Xu et al. (2017).

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744 **Figure 2.** Annual-averaged (a) low-, (b) middle-, high-level (c) and total cloud amount changes
 745 between the $2\times\text{CO}_2$ and control experiments.

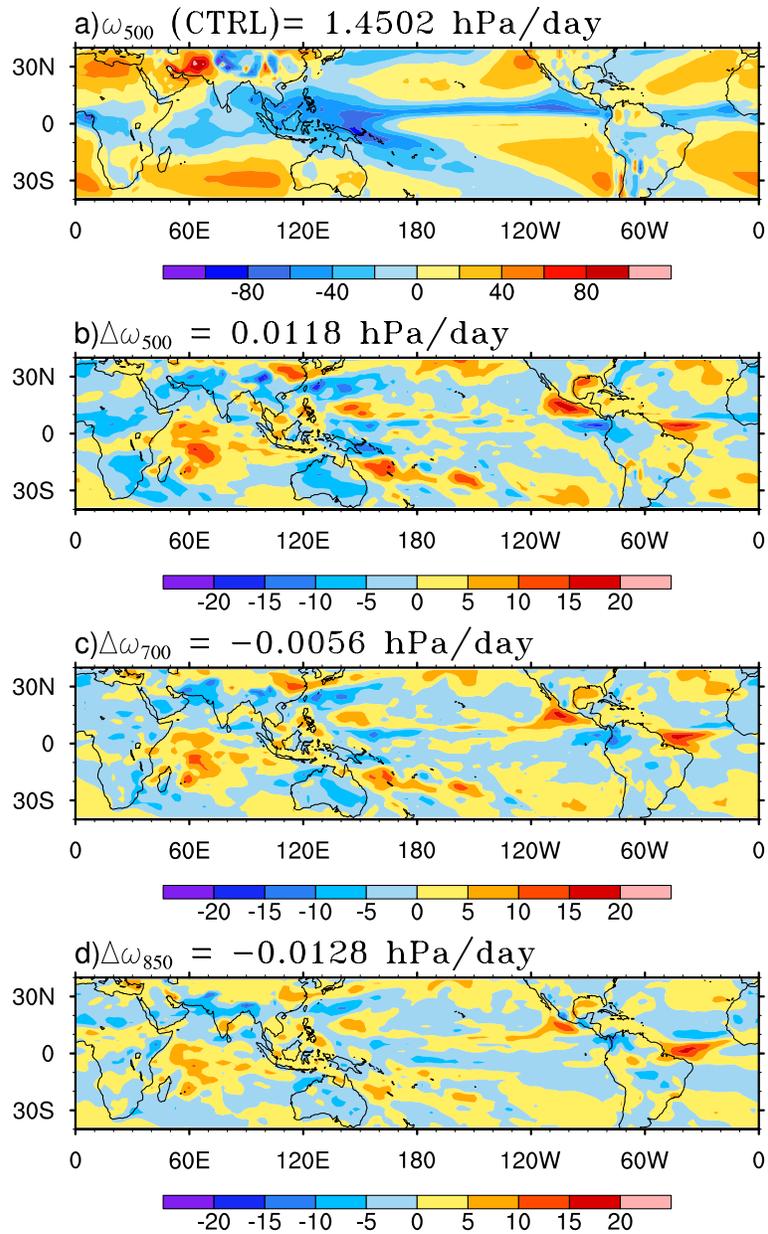
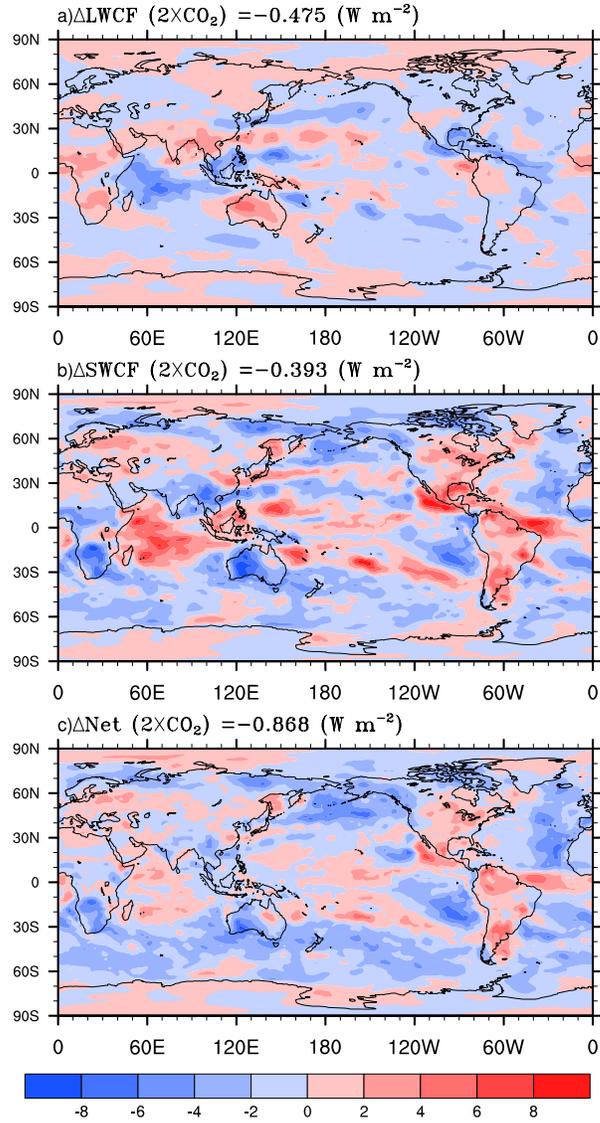


Figure 3. (a) Annual-averaged pressure vertical velocity at 500 hPa (ω_{500}) for the control experiment, and annual-averaged changes in (b) ω_{500} (c) ω_{700} , and (d) ω_{850} between the 2xCO₂ and control experiments.



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748 **Figure 4.** Annual-averaged (a) longwave, (b) shortwave and (c) net cloud radiative effect
 749 changes between the 2xCO₂ and control experiments.

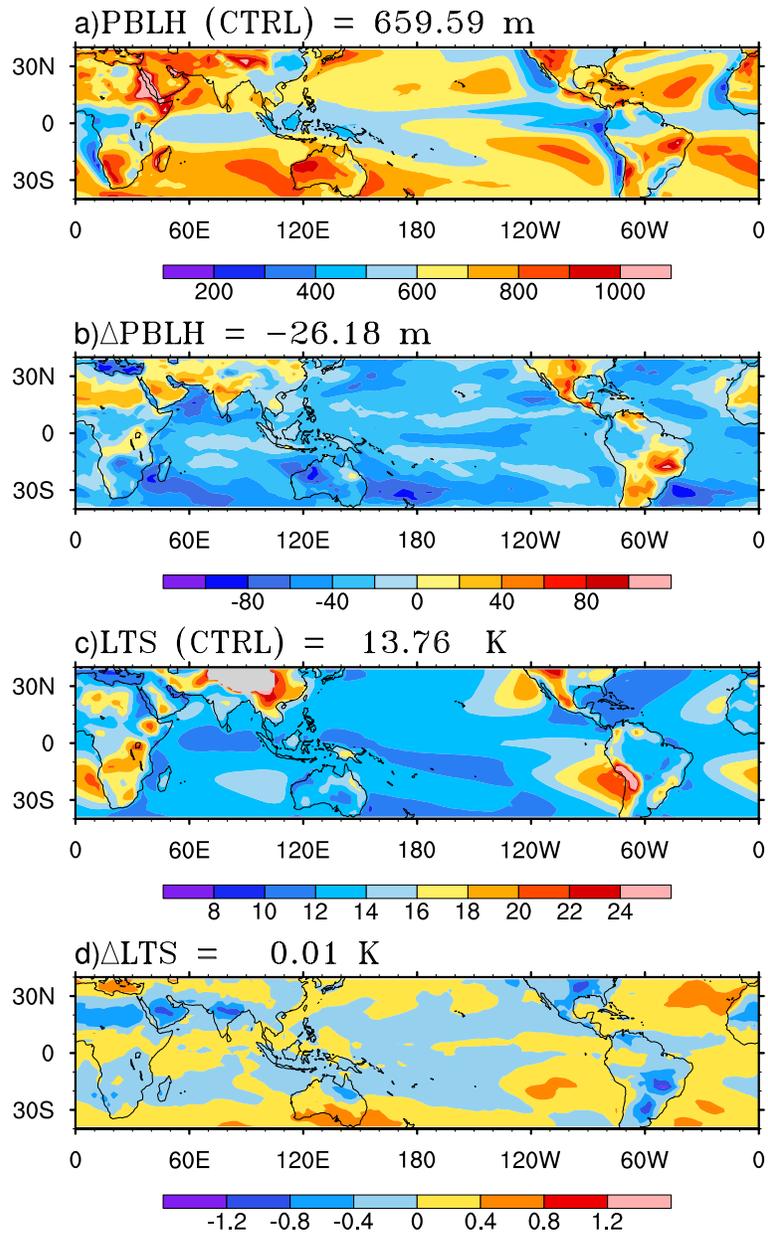


Figure 5. Annual-averaged (a) PBL height and (c) lower tropospheric stability (LTS) from the control experiment and their changes (b, d) between the 2xCO₂ and control experiments.

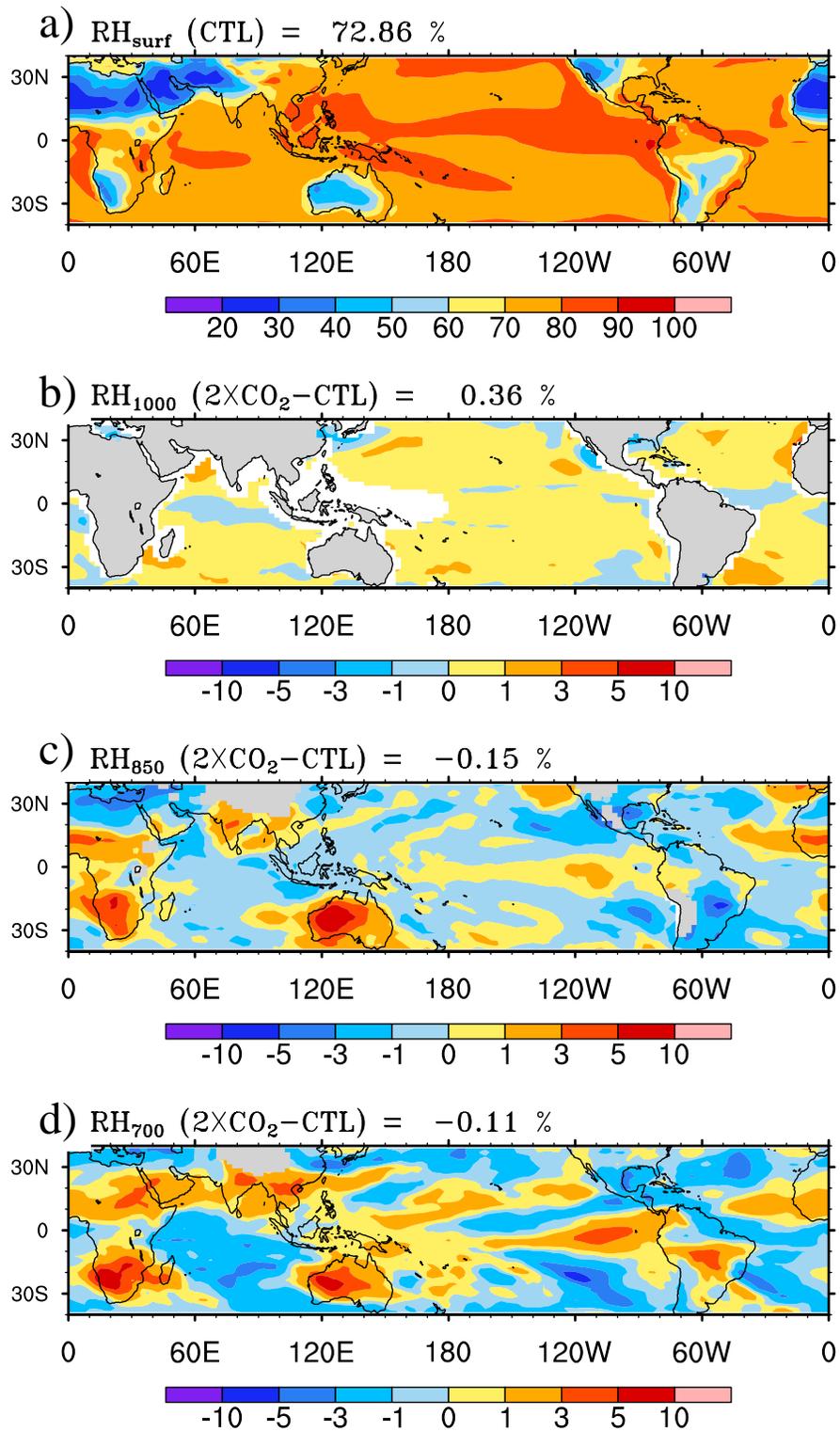


Figure 6. Annual-averaged relative humidity at first model level from the control experiment (a) and the changes at first model level (b), 850 hPa (c) and 700 hPa (d) between the $2\times\text{CO}_2$ and control experiments.

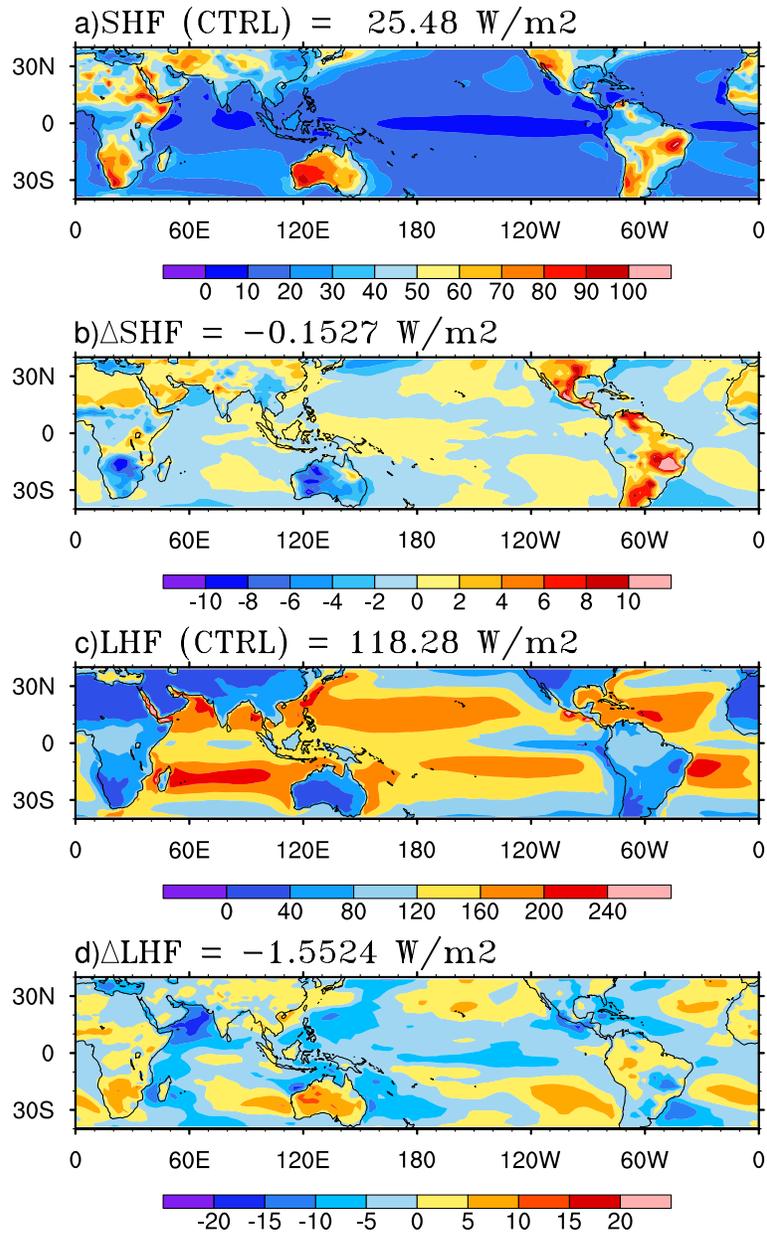


Figure 7. Annual-averaged (a) sensible heat (SH) and (c) latent heat (LH) fluxes from the control experiment and their changes (b, d) between the 2xCO₂ and control experiments.

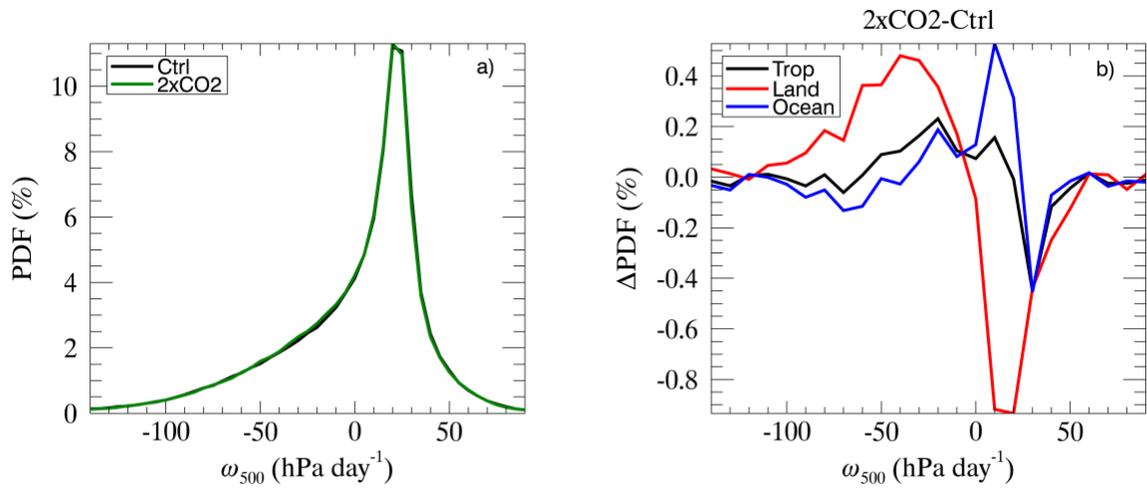


Figure 8. The frequency (%) of midtropospheric pressure velocity (ω_{500}) over the entire tropics (a) and the difference of frequency between the 2xCO₂ and control experiments (b) for the entire tropics (black), tropical land (land) and tropical ocean (blue). The bin size is 5 hPa day^{-1} .

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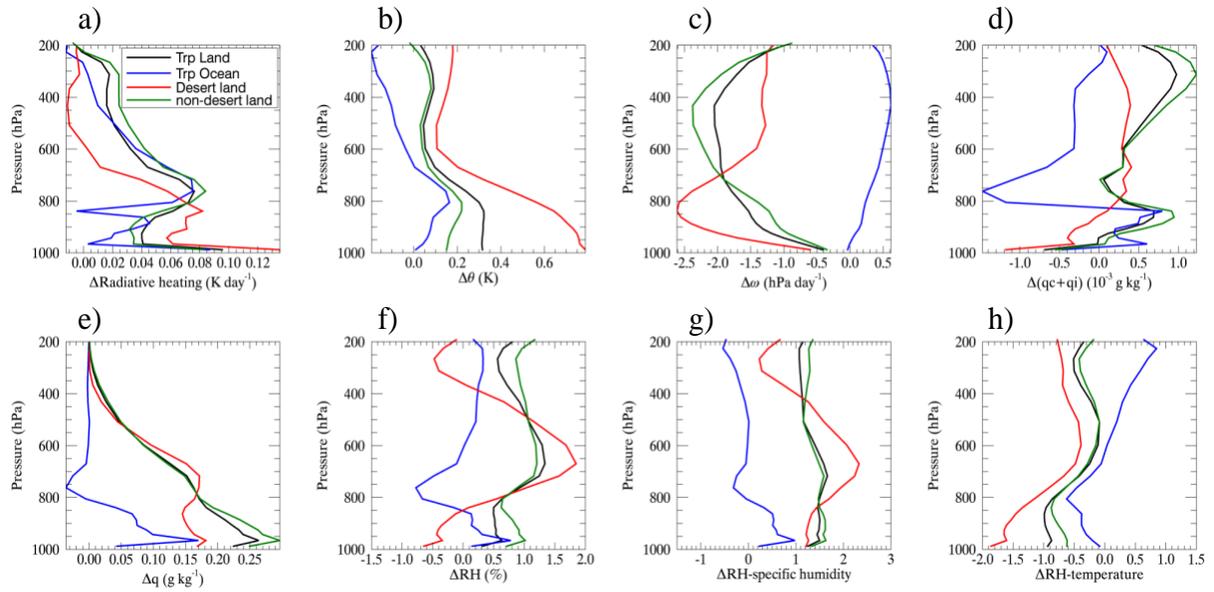
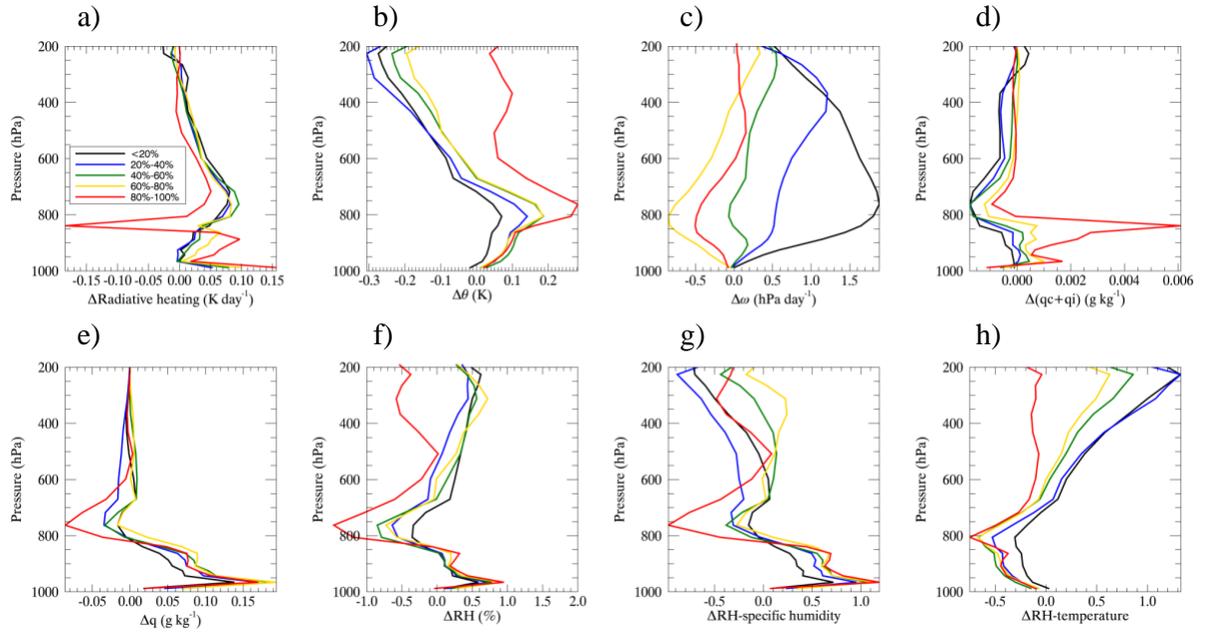


Figure 9. Vertical profiles of changes in (a) radiative heating rate, (b) potential temperature, (c) pressure vertical velocity, (d) sum of cloud water and cloud ice mixing ratio, (e) specific humidity, (f) relative humidity, (g) relative humidity due to changes in specific humidity and (h) relative humidity due to changes in temperature averaged over the tropical lands, tropical ocean, tropical deserts and tropical non-deserts between the 2xCO₂ and control experiments.



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Figure 10. Same as Figure 9 except for the tropical ocean profiles sorted by annual-mean ω_{500} , each color represents the average for 20% of ω_{500} values. The 20%, 40%, 60% and 80% thresholds of ω_{500} are roughly -24 hPa, -1 hPa, 15 hPa, and 23 hPa.