

Particle Simulation

HIWC Modeling Insights

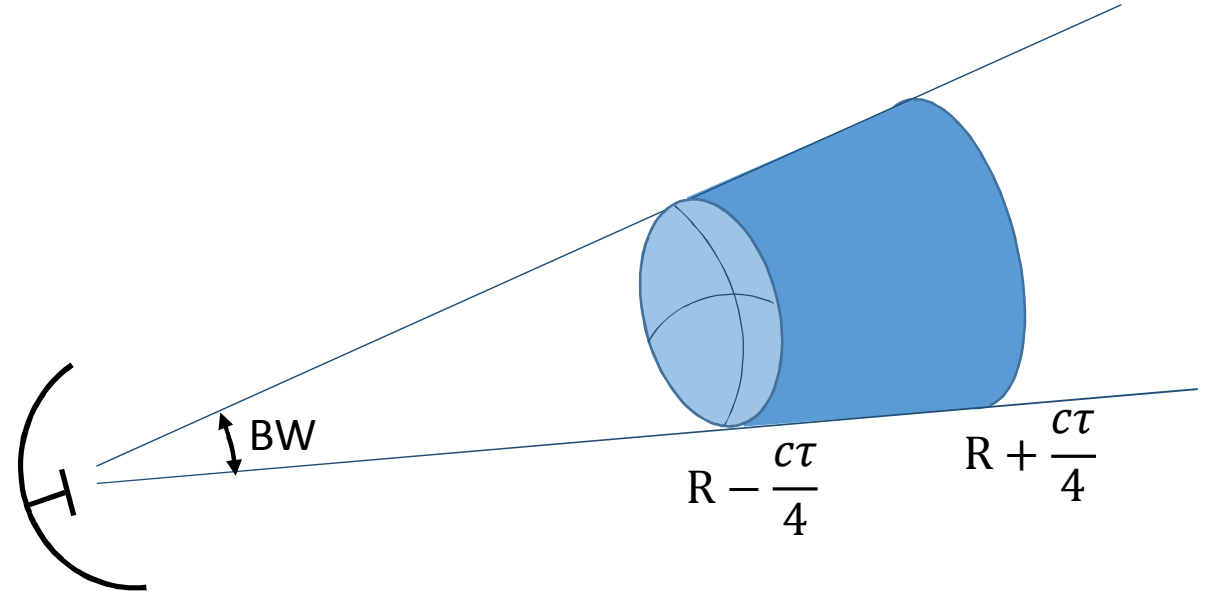
30 March 2020

Agenda

- ADWRS modeling approach
- Detailed particle model
 - Effects of angle resolution
 - Effects of range resolution
 - Random vs stochastic processes
 - Simulating HIWC
- Findings and Results
- Conclusions and next steps

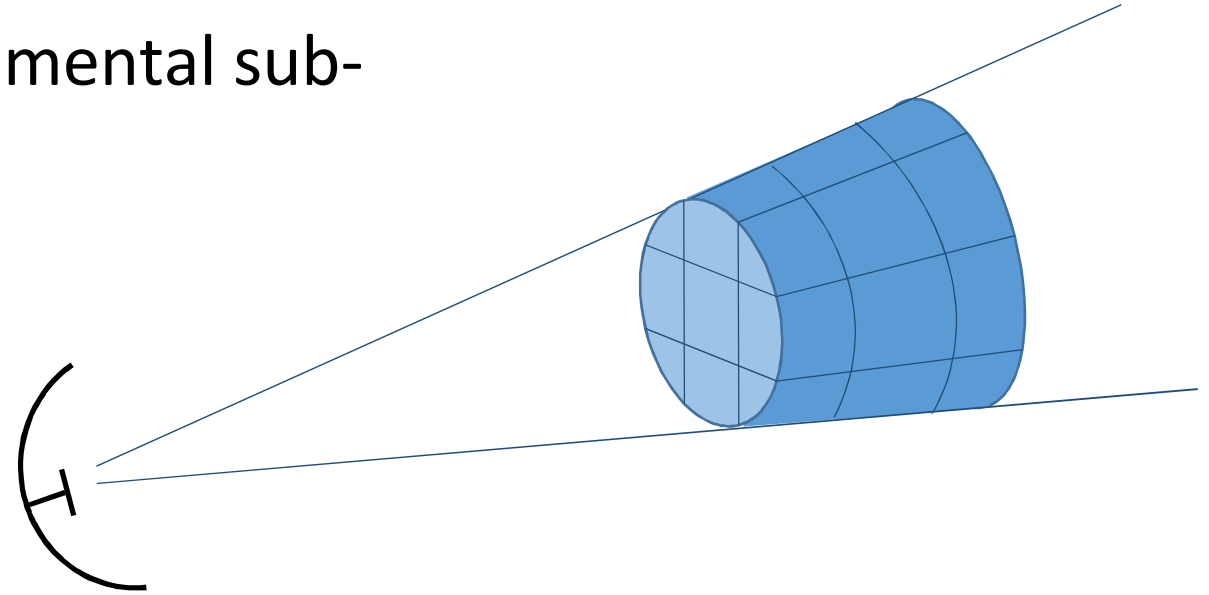
Radar Resolution Volume

- Beam width (BW)
- Range (R)
- Pulse with (τ)
- Homogeneous cloud
 - Particle uniformly distributed throughout the volume



ADWRS Weather Modeling

- Divide resolution volume into incremental sub-volumes (cells)
 - $\Delta\text{Azimuth}$, $\Delta\text{Elevation}$, ΔRange
- $\Delta V_i = R_i^2 \Delta A z_i \Delta E l_i \Delta R_i$
- $\eta_i = \frac{\pi^5}{\lambda^4} |K_W|^2 Z_i$
- $\sigma_i = \eta_i \Delta V_i$
- $A_i = \sqrt{\frac{P_t \lambda^2 G_i^2 \sigma_i}{(4\pi)^3 R_i^4}}$
- $v(t) = \sum_i A_i e^{-j\left(\bar{\phi}_i + \frac{4\pi v_i t}{\lambda}\right)}$



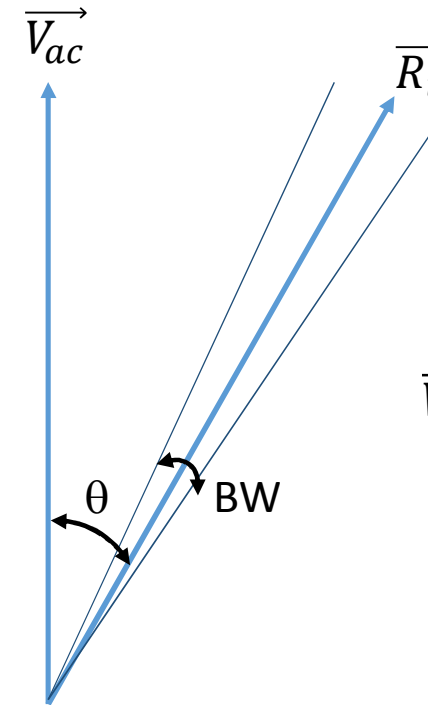
Z_i and v_i interpolated from TASS database
 G_i from antenna pattern

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Particle Radar Return Signal

- $\sigma_i = \frac{\pi^5}{\lambda^4} |K_W|^2 D_i^6$
- $A_i = \sqrt{\frac{P_t \lambda^2 G_i^2 \sigma_i}{(4\pi)^3 R_i^4}}$
- $v(t) = \sum_i A_i e^{-j\left(\frac{4\pi R_i}{\lambda} + 2\pi f_i t\right)}$
- $f_i = \frac{2}{\lambda} \frac{dR_i}{dt}$

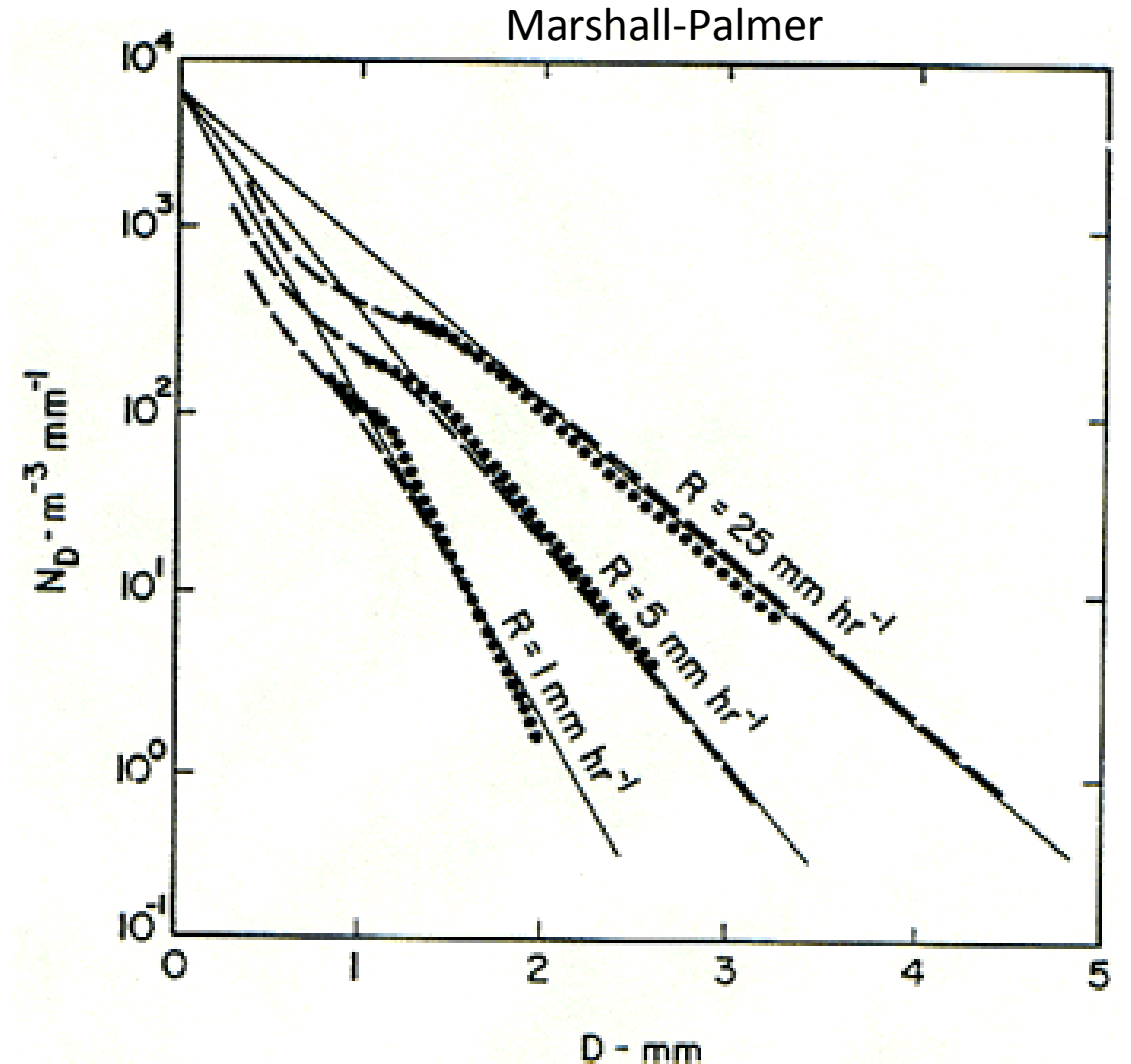


$$\vec{V}_{ac} \cdot \vec{R}_i = |\vec{V}_{ac}| |\vec{R}_i| \cos \theta_i$$

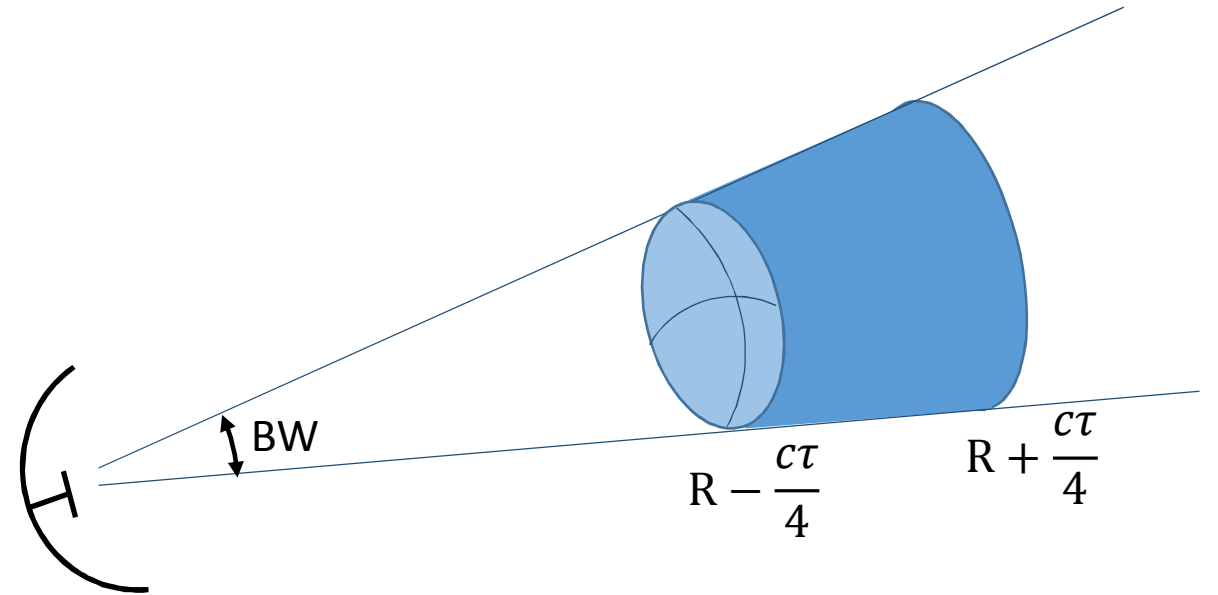
$$\frac{dR_i}{dt} = V_{ac} \cos \theta_i$$

Detailed Particle Model

- Particle/Drop size distribution
 - $N(D) = N_0 e^{-\Lambda D}$
- Number of particles of a given size in a unit volume
- Theoretical/Truth information
 - Mass = $\left(\frac{\pi\rho}{6}\right) \int_0^\infty D^3 N(D) dD = \frac{\pi\rho N_0}{\Lambda^4}$
 - RRF = $\int_0^\infty D^6 N(D) dD = \frac{N_0(6!)}{\Lambda^7}$



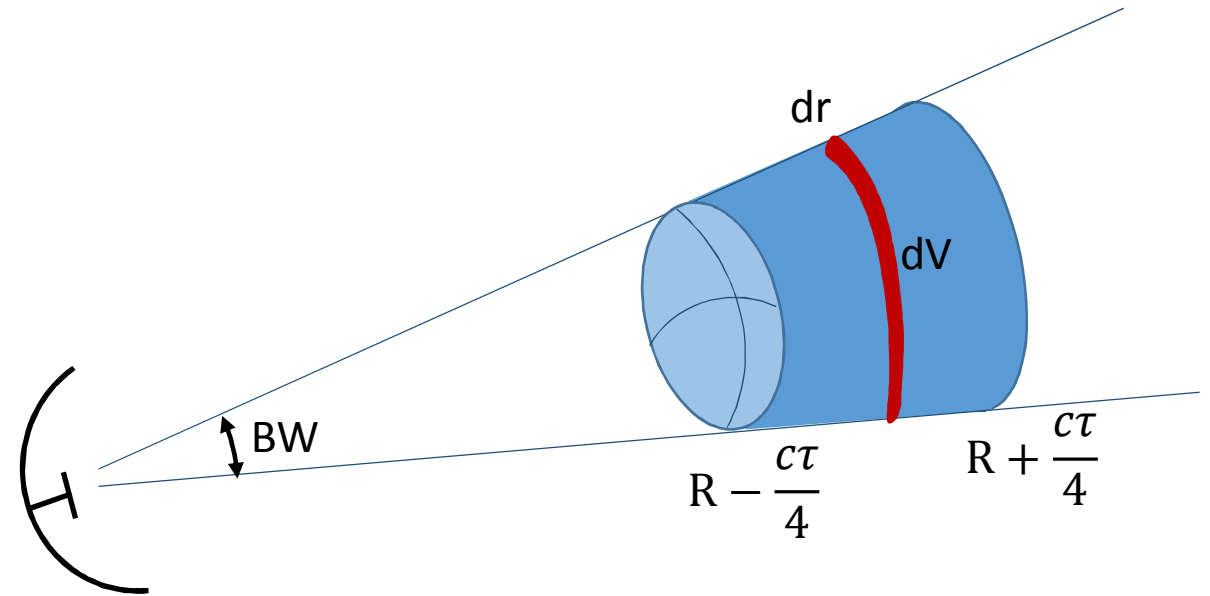
Resolution Volume



- $V = \int_{R - \frac{c\tau}{2}}^{R + \frac{c\tau}{2}} \int_0^\pi \int_0^{2\pi} r^2 g(\theta, \phi) \sin \theta d\theta d\phi dr$
- Accounting for Gaussian antenna gain pattern, $g(\theta, \phi)$
 - Doviak and Zrnic (4.3)
- $V = R^2 \frac{\pi BW^2 c\tau}{16(\ln 2)}$

Volume Distribution

- Area of a spherical cap
 - $A = 2\pi r^2(1 - \cos \theta)$
- Differential volume
 - $dV = 2\pi r^2(1 - \cos BW)dr$
- Distribution of volume in resolution cell
 - $V(r) = \int_{R - \frac{c\tau}{4}}^r dV = \frac{2}{3}\pi(1 - \cos BW) \left[r^3 - \left(R - \frac{c\tau}{4} \right)^3 \right]$
- Inverse transform sampling
 - Generate random range resulting in uniform particle density
 - $r = \sqrt[3]{U(0,1) \left[\left(R + \frac{c\tau}{4} \right)^3 - \left(R - \frac{c\tau}{4} \right)^3 \right] + \left(R - \frac{c\tau}{4} \right)^3}$



Particle Simulation

- Start at smallest D
- While $N(D) > 0$
 - Compute $\sigma_i = f(D^6)$
 - Get random range, R_i
 - Compute amplitude, $A_i = f(\sigma_i, R_i)$
 - Get random angle
 - Compute radial velocity/Doppler, f_i
 - For pulses in CPI
 - Compute phase = $f(\text{PRI}, R_i, f_i)$
 - $v(t) = \text{sum complex returns for all } i$
 - Increment D
 - Compute $N(D)$

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Doppler Frequency Distribution

- Particles uniformly distributed in angle

- Let $X = U\left(\theta - \frac{BW}{2}, \theta + \frac{BW}{2}\right)$

- Let $Y = f = \frac{2V_{ac}}{\lambda} \cos X$

- Transformation of Random Variable

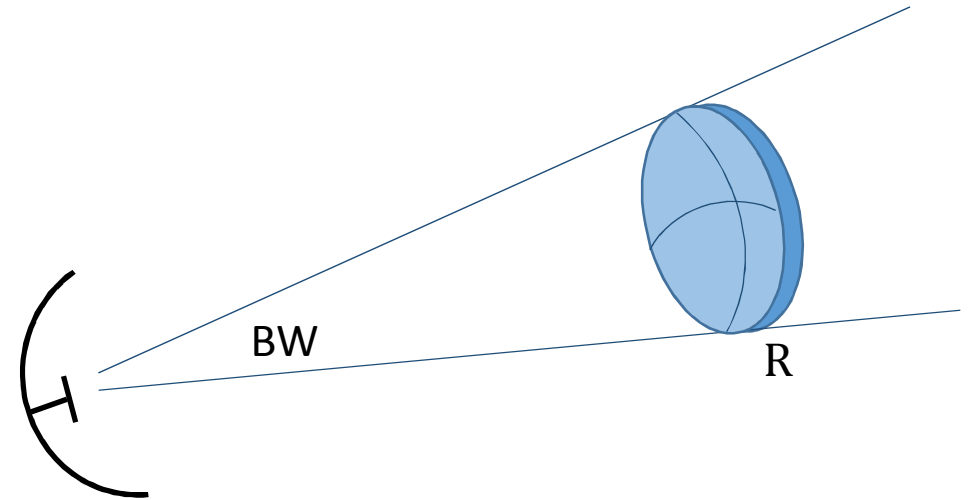
- Probability distribution function (PDF)

- $$p(f) = \frac{\lambda}{2BWV_{ac} \sqrt{1 - \left(\frac{f\lambda}{2V_{ac}}\right)^2}}, f_{\min} \leq f \leq f_{\max}$$

- $$f_{\min} = \frac{2V_{ac}}{\lambda} \cos\left(\theta - \frac{BW}{2}\right) \text{ \& } f_{\max} = \frac{2V_{ac}}{\lambda} \cos\left(\theta + \frac{BW}{2}\right)$$

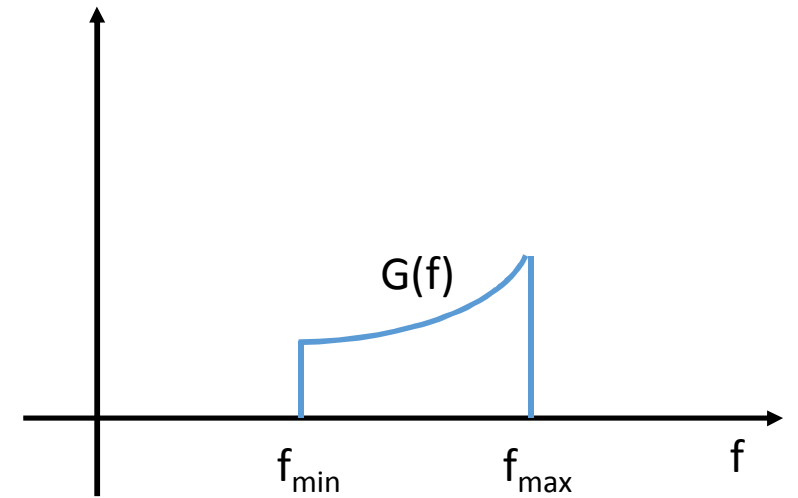
Amplitude Distribution

- $A = \sum_i A_i = \sqrt{P_t} \frac{\pi |K_w| G}{4\lambda} \sum_i \frac{D_i^3}{R_i^2}$
- Assume small range resolution
 - $R_i = R$ for all i
- $A = \sqrt{P_t} \frac{\pi |K_w| G}{4\lambda R^2} V \int_0^\infty D^3 N(D) dD$
- $A = \sqrt{P_t} \frac{3\pi^2 BW^2 c\tau |K_w| GN_0}{32(\ln 2)\lambda R^2 \Lambda^4}$



Frequency Domain

- $G(f) = Ap(f) = \frac{A\lambda}{2BWV_{ac} \sqrt{1 - \left(\frac{f\lambda}{2V_{ac}}\right)^2}}$



- Inverse Fourier transform

- $g(t) = \int_{-\infty}^{\infty} G(f) e^{2\pi jft} df$

- $g(t) = \frac{A\lambda}{2BWV_{ac}} \int_{f_{\min}}^{f_{\max}} \frac{e^{2\pi jft}}{\sqrt{1 - \left(\frac{f\lambda}{2V_{ac}}\right)^2}} df$

Time Domain

- Small Angle Approximation

- $G(f) = A \text{rect}(F)$

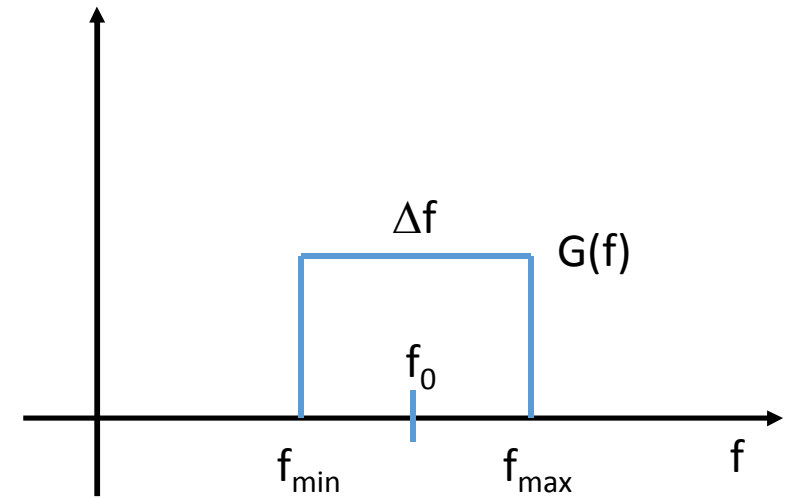
- $F = \frac{f-f_0}{\Delta f}$

- Inverse Fourier transform

- $g(t) = \int_{-\infty}^{\infty} G(f) e^{2\pi j f t} df$

- $g(t) = A \int_{-\frac{1}{2}}^{\frac{1}{2}} e^{2\pi j (F\Delta f + f_0)t} dF$

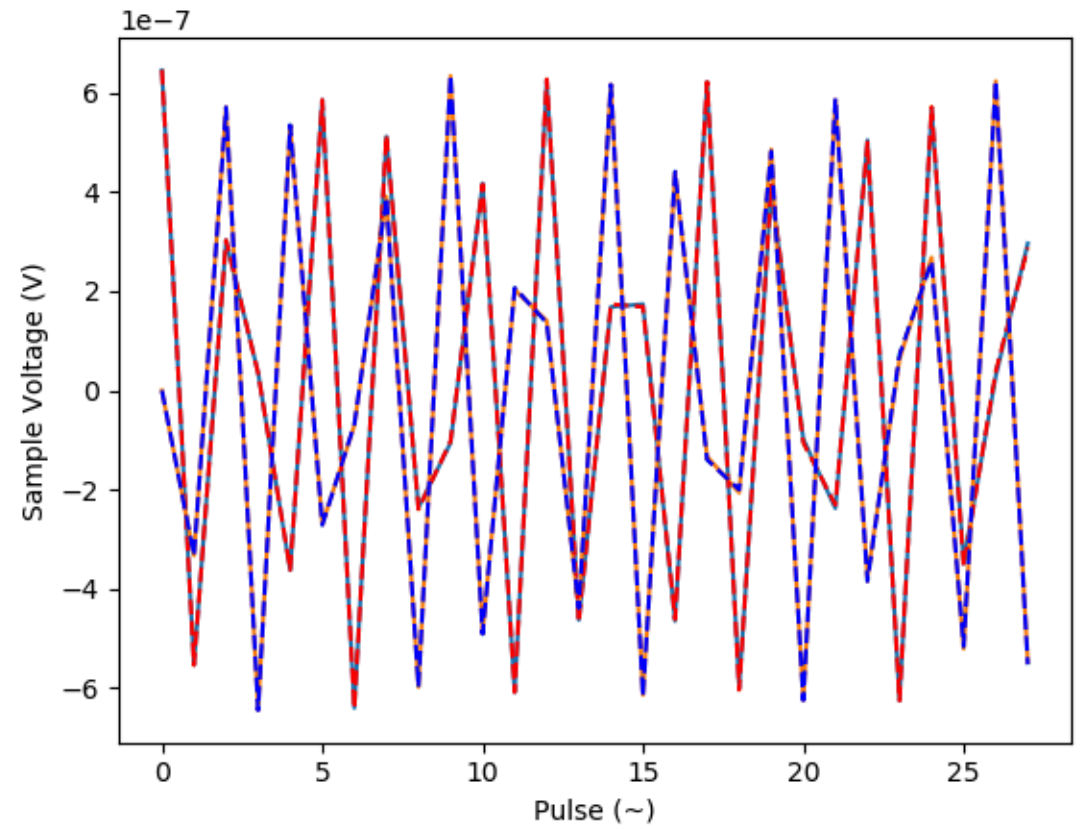
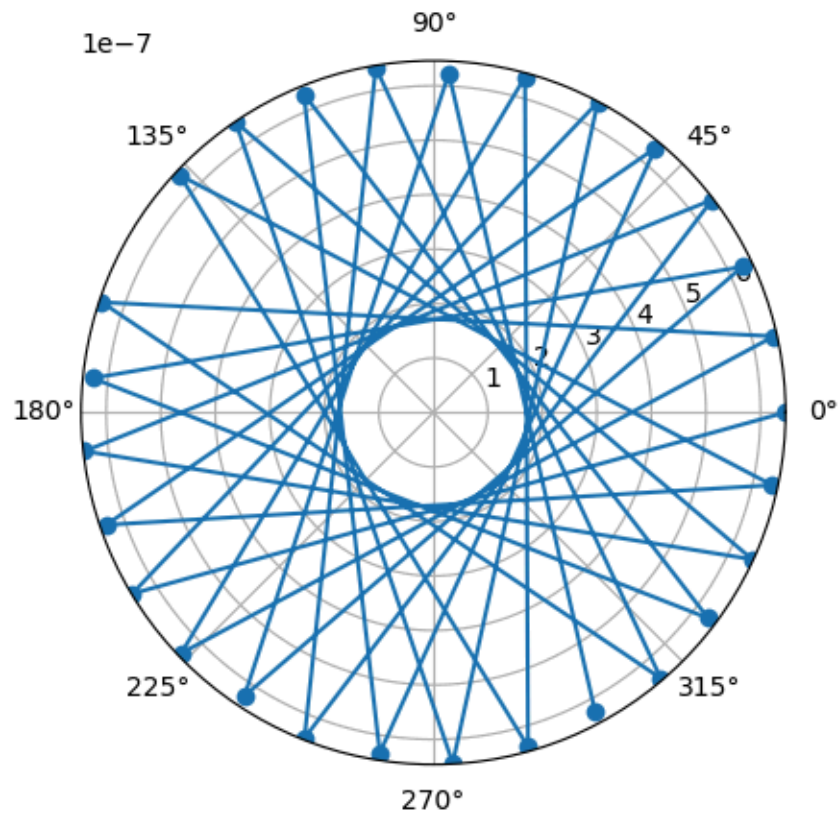
- $g(t) = A \text{sinc}(\Delta f t) e^{2\pi j f_0 t}$



$$f_0 = (f_{\min} + f_{\max})/2$$

$$\Delta f = f_{\max} - f_{\min}$$

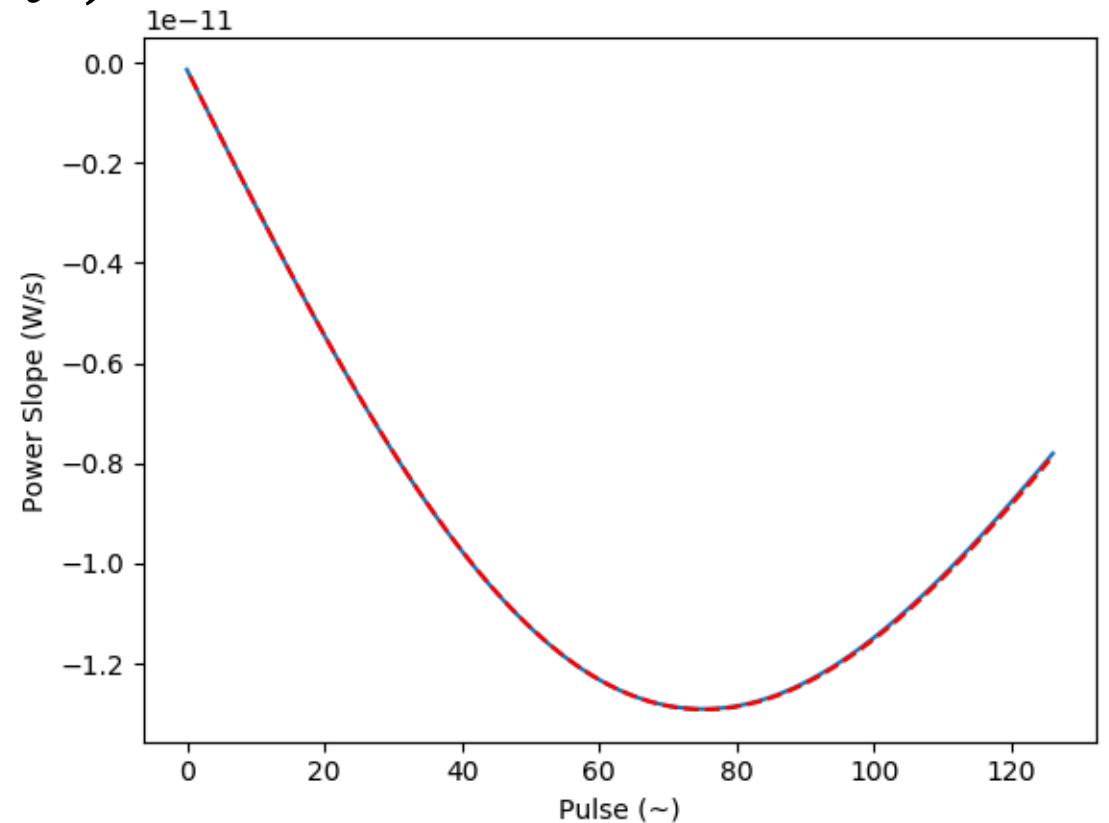
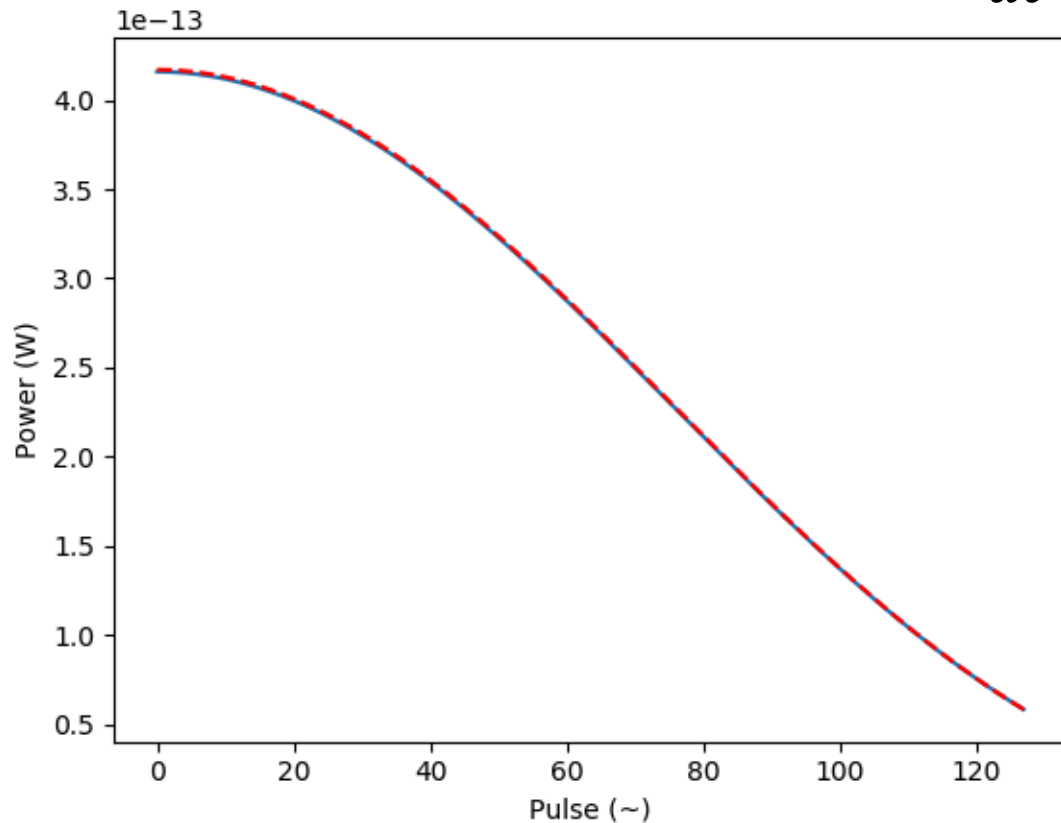
Particle Simulation Results



Power Fluctuations

$$P(t) = |g(t)|^2 = A^2 \text{sinc}^2(\Delta f t)$$

$$\bullet \frac{dP}{dt} = \left(\frac{2A^2}{t} \right) \text{sinc}(\Delta f t) [\cos(\pi \Delta f t) - \text{sinc}(\Delta f t)]$$



Assessment of HIWC PSD

- $\text{Mass} = \frac{\pi \rho N_0}{\Lambda^4}$

- $\text{RRF} = \frac{N_0(6!)}{\Lambda^7}$

- Consider HIWC and non-HIWC conditions with equivalent RRF

- $\frac{\text{Mass}_{HIWC}}{\text{Mass}_{non-HIWC}} = \left(\frac{N_{0HIWC}}{N_{0non-HIWC}} \right)^{\frac{3}{7}}$

Rate of Change in Received Power

- $\frac{P_{HIWC}}{P_{non-HIWC}} = \left(\frac{A_{HIWC}}{A_{non-HIWC}}\right)^2 = \left(\frac{N_{0HIWC}}{N_{0non-HIWC}}\right)^{\frac{3}{7}}$
- $\frac{\left(\frac{dP}{dt}\right)_{HIWC}}{\left(\frac{dP}{dt}\right)_{non-HIWC}} = \left(\frac{A_{HIWC}}{A_{non-HIWC}}\right)^2 = \left(\frac{N_{0HIWC}}{N_{0non-HIWC}}\right)^{\frac{6}{7}}$

Particle Simulation

- Non-HIWC
 - $N_0 = 8000 \text{ m}^{-3} \text{ mm}^{-1}$ (Marshall-Palmer)
 - RRF = 26 dBZ
 - Mass = 0.1 g m^{-3}
- HIWC
 - $1000N_0$
 - Mass = 2.1 g m^{-3}
- $\frac{Mass_{HIWC}}{Mass_{non-HIWC}} = (1000)^{\frac{3}{7}} = 19$
- $\frac{\left(\frac{dP}{dt}\right)_{HIWC}}{\left(\frac{dP}{dt}\right)_{non-HIWC}} = (1000)^{\frac{6}{7}} = 372$

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Random vs Stochastic Process

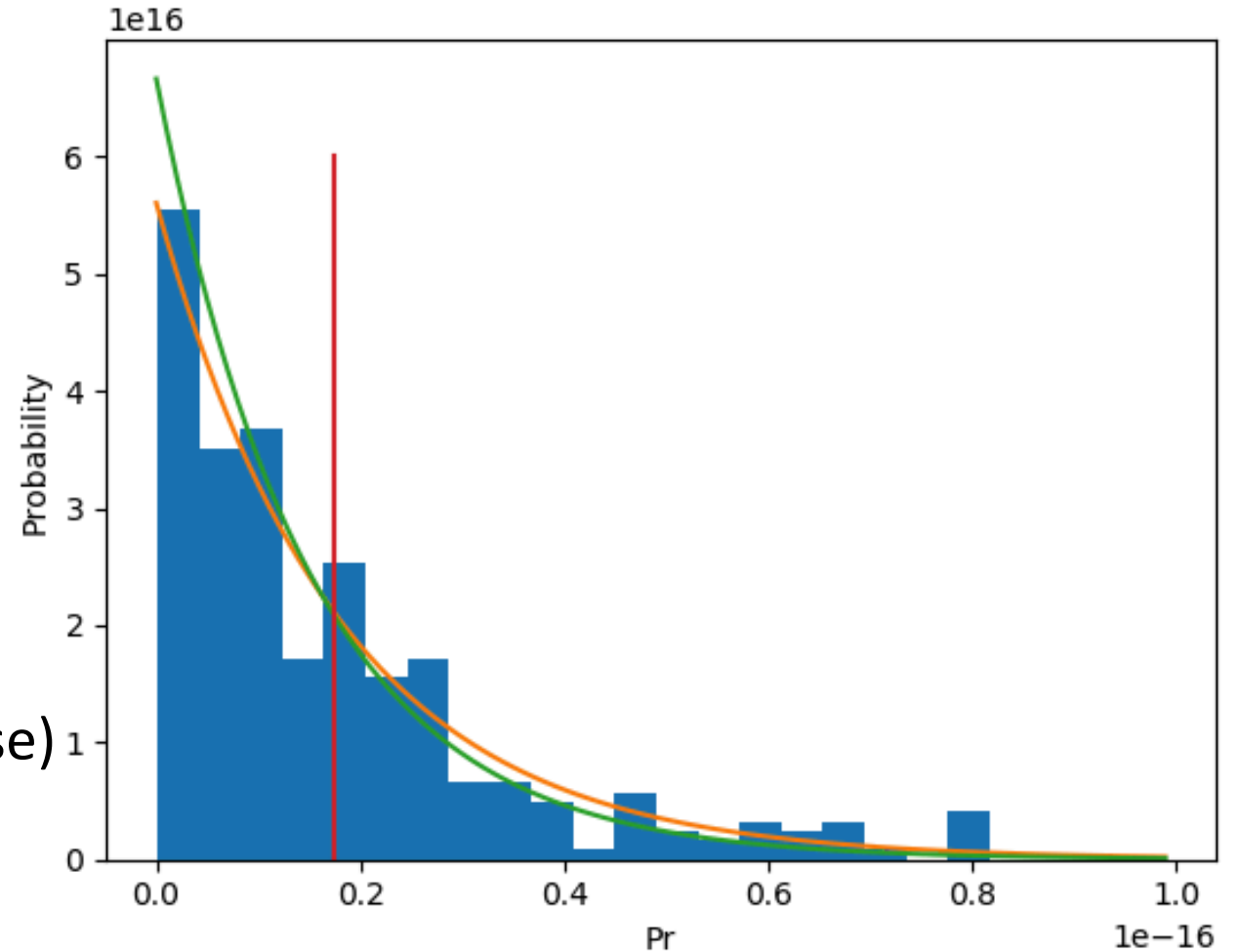
- Random Processes
 - Range bin to Range bin are independent samples
 - Independent resolution volumes
- Stochastic Processes
 - Random variation in time
 - Pulse to pulse within a CPI?
- PRI = $O(0.1 \text{ ms})$
 - At 400 kts, aircraft changes position by $> \lambda$ between pulses
 - However, relative particle motion small

Fluctuation Echo from Randomly Distributed Scatterers

- Independent Samples
- Rayleigh Distribution of Signal Magnitude
 - Marshall and Hitschfield (1953)
 - Doviak and Zrnic (4.2)
- Exponential Power Distribution
 - $p(P) = \lambda e^{-\lambda P}$
 - Where rate parameter $\lambda = \frac{1}{\langle P_r \rangle}$
- Mean = $\frac{1}{\lambda}$
- Variance = $\frac{1}{\lambda^2}$
- $I_D = \frac{Var}{Mean} = \frac{1}{\lambda} = \text{mean}$
- $RIWC \approx \log_{10}(I_D) = \log_{10}(RRF)$

Particle Simulation

- 300 independent samples
- RIWC
 - I_D equivalent to mean RRF
 - Insensitive to IWC
- Flight data
 - I_D varies widely from mean RRF
 - Sensitive to IWC
 - Samples within a CPI (pulse-to-pulse) are not independent



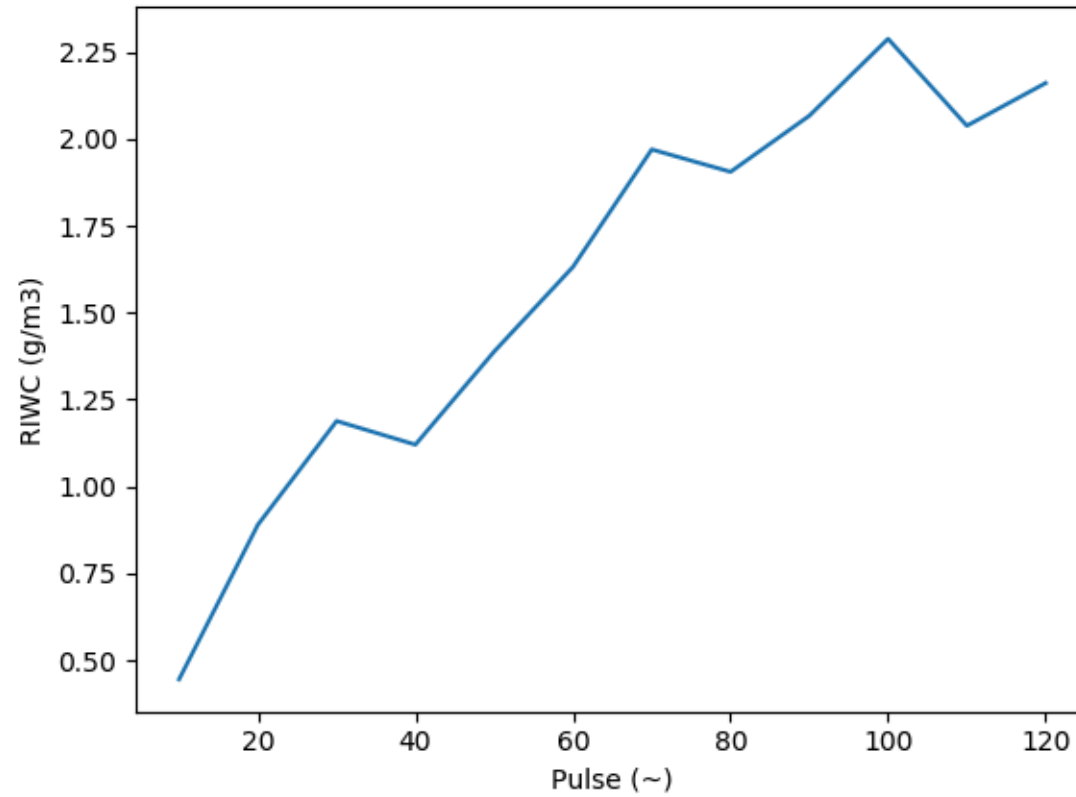
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Particle Simulation

- Full sized resolution volume
 - Particles randomly distributed in angle and range
 - Uniform density
- Monte-Carlo runs
 - Multiple CPI's in a measurement
 - Beamwidth averaging, spatial filtering
 - Mean values from n pulses in a CPI
 - Mean values from 10 CPI's

Effect of CPI Duration



Simulating HIWC Conditions

- Mean values from 8 pulses in a CPI
- Non-HIWC condition
 - RRF 25-29 dBZ, RIWC $\approx 0 \text{ g m}^{-3}$
- HIWC conditions
 - RRF 25-27 dBZ, RIWC $\approx 0 \text{ g m}^{-3}$

Findings and Results

- Detailed particle model verifies ADWRS approximation
 - Independent samples in random process follow exponential distribution
- Theoretical solution assuming small range extent
 - Suggests variation in stochastic process should increase with N_0
- Detailed particle model and ADWRS output
 - Coherency de-correlates with time (long duration CPI)
- Detailed particle model HIWC simulation
 - variation in stochastic process insensitive to N_0

Conclusions and Next Steps

- Doppler distribution due to beamwidth results in “slow” power fluctuations
 - Return signals decorrelate over long durations
- Uniformly distributed particles do not result in “rapid” power fluctuations with increasing mass
 - Return signals remain equally correlated regardless of PSD
- Observed power variation must come from another source
 - Non-uniformly distributed particles?
- Investigate effect of clustering