



1 **Venusian Habitable Climate Scenarios: Modeling Venus**
2 **through time and applications to slowly rotating**
3 **Venus-Like Exoplanets**

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9 **Key Points:**

- 10 • Venus could have had habitable conditions for nearly 3 billion years.
11 • Surface liquid water is required for any habitable scenario.
12 • Solar insolation through time is not a crucial factor if a carbonate-silicate cycle
13 is in action.

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Abstract

One popular view of Venus' climate history describes a world that has spent much of its life with surface liquid water, plate tectonics, and a stable temperate climate. Part of the basis for this optimistic scenario is the high deuterium to hydrogen ratio from the Pioneer Venus mission that was interpreted to imply Venus had a shallow ocean's worth of water throughout much of its history. Another view is that Venus had a long lived (~ 100 million year) primordial magma ocean with a CO_2 and steam atmosphere. Venus' long lived steam atmosphere would sufficient time to dissociate most of the water vapor, allow significant hydrogen escape and oxidize the magma ocean. A third scenario is that Venus had surface water and habitable conditions early in its history for a short period of time ($< 1\text{Gyr}$), but that a moist/runaway greenhouse took effect because of a gradually warming sun, leaving the planet desiccated ever since. Using a general circulation model we demonstrate the viability of the first scenario using the few observational constraints available. We further speculate that Large Igneous Provinces and the global resurfacing 100s of millions of years ago played key roles in ending the clement period in its history and presenting the Venus we see today. The results have implications for what astronomers term "the habitable zone," and if Venus-like exoplanets exist with clement conditions akin to modern Earth we propose to place them in what we term the "optimistic Venus zone."

Plain Language Summary

We have little data on our neighbor Venus to help us understand its climate history. Yet Earth and Venus are sister worlds: they initially formed close to one another, and have nearly the same mass and radius. Despite the differences in their current atmospheres and surface temperatures, they likely have similar bulk compositions, making comparison between them extremely valuable for illuminating their distinct climate histories. We analyze our present data on Venus alongside knowledge about Earth's climate history to make a number of exciting claims. Evaluating several snapshots in time over the past 4+ billion years, we show that Venus could have sustained liquid water and moderate temperatures for most of this period. Cloud feedbacks from a slowly rotating world with surface liquid water reservoirs were the keys to keeping the planet clement. Contrast this with its current surface temperature of 450 degrees and an atmosphere dominated by Carbon Dioxide and Nitrogen. Our results demonstrate that it was not the gradual warming of the sun over the eons that contributed to Venus' present hot-house state. Rather we speculate that large igneous provinces and the global resurfacing 100s of millions of years ago played key roles in ending the clement period in its history.

1 Introduction

The case for prolonged habitability of Venus has been made by a number of authors in recent years (e.g., Grinspoon & Bullock, 2007; Way et al., 2016). If so, then if habitability was widespread and persisted over geological timescales (e.g. 10^9 years), it is possible that organisms were capable of filling a large variety of environmental niches as occurred on Earth via evolutionary processes. This has led to speculation about possible remaining life-friendly niches (e.g., Morowitz & Sagan, 1967; Grinspoon, 1997; Cockell, 1999; Schulze-Makuch & Irwin, 2002; Schulze-Makuch et al., 2004; Dartnell et al., 2015; Limaye et al., 2018). These possibilities alone provide sufficient motivation to understand whether early Venus was habitable or not. However, Venus is also interesting from the perspective of the large number of rocky exoplanets discovered to date inside the traditional inner edge of the habitable zone but far enough from their host stars to maintain an atmosphere (Kane et al., 2019). If some of these planets in close proximity to their host stars have long periods of habitability it may overturn traditional notions of the habitable zone (e.g., Kopparapu et al., 2013) and influence target selection

64 for characterization of the atmospheres of these planets. Hence, what appears to be a
65 modern Venus-like world close to its parent star might host surface liquid water. We refer
66 to such habitable worlds as residing in an “Optimistic Venus Zone,” a subset of the
67 planets in the “Venus Zone” described in Kane et al. (2014).

68 There are three primary requirements for the early habitability of Venus. The first
69 is that temperatures were low enough that liquid water was capable of condensing on
70 the surface of Venus. The second is that Venus had a sufficient inventory of water to create
71 the conditions believed necessary for the rise of life on terrestrial worlds (e.g., Brack
72 et al., 2010). Third, volatile cycling and a geologically active surface must exist to regulate
73 the climate as possibly supported by the work of Noack et al. (2012). For surficial
74 water reservoirs most research has focused on the poorly constrained measurement of
75 the D/H ratio of $\sim 150 \pm 30$ (times Earth’s value) by the Pioneer Venus Mission (<https://nssdc.gsfc.nasa.gov/plan>)
76 (Donahue et al., 1982, 1997). A number of other ground based measurements have been
77 made (De Bergh et al., 1991; Marcq, 2006; Bjoraker et al., 1992; Fedorova et al., 2008)
78 that have generally revealed values of D/H greater than 100. For a recent review of D/H
79 measurements of Venus see Section 6.3 of Marcq et al. (2017). We review the work of
80 authors who have considered Venus’ climate evolution below.

81 The first simple gray radiative transfer calculations of Venus’ climate history suggested
82 an early runaway greenhouse effect (e.g., Ingersoll, 1969; Rasool & de Bergh, 1970).
83 Any water would have remained in vapor form throughout Venus’ early history until its
84 loss to space via photodissociation processes (Goody & Walker, 1972; Walker, 1975). CO₂
85 would continue to be outgassed and would accumulate (offset by atmospheric loss over
86 the last 4.5Gyr) to the values we see today.

87 Pollack (1971) used a similar model but with non-gray radiative transfer. This was
88 probably the first work to show that Venus could have hosted liquid water on its surface
89 in its early history, but they also demonstrated that a runaway greenhouse was possible
90 as well. The difference between the two outcomes was *mostly* due to the specified cloud
91 cover and water vapor content of the atmosphere. With 50% cloud fraction the result
92 was usually a runaway greenhouse. For an N₂ dominated atmosphere with a modern Earth
93 water vapor profile, 100% cloud cover, and a 30% less luminous sun (akin to ~ 4.5 Ga)
94 he achieved surface temperatures ~ 300 K. These purely radiative models did not include
95 convection, which would limit the water vapor mixing ratio by precipitation and limit
96 the lapse rate of temperature to the moist adiabatic value (Rampino & Caldeira,
97 1994) and lessen the strength of the greenhouse effect to some degree. Likewise, they did
98 not include an atmospheric general circulation that would determine the cloud fraction
99 self-consistently. Donahue and Pollack (1983) replicated some aspects of the Pollack (1971)
100 work and again showed that clement conditions were possible in Venus’ early history.

101 Kasting et al. (1984) used a 1-D radiative-convective model to demonstrate that
102 a planet with an insolation $S_0X=1.45$ times that of modern Earth ($1973 \text{ W m}^{-2} \sim 3.8$ Gya
103 at Venus’ orbit) would have temperatures $\sim 100^\circ \text{ C}$ for a 2 bar N₂ H₂O atmosphere. This
104 atmosphere contained a wet stratosphere, and thus photodissociation of H₂O and subsequent
105 escape of hydrogen (e.g., Goody & Walker, 1972) could explain the long-term loss
106 of Venus’ primordial ocean and its present dry conditions. In a subsequent study
107 Kasting (1988) claimed that clouds would cool the atmosphere enough to keep it in a
108 moist greenhouse state, rather than the runaway state. At the same time the Kasting
109 (1988) surface temperature for his “Early Venus” (Figure 7) cloud-free model was ~ 500 K
110 and hovered just at the margin of the moist vs. runaway greenhouse states. Thus it was
111 clear early on that maintaining liquid water on ancient Venus required high fractional
112 cloud cover, but whether physical mechanisms exist to produce these cloudy conditions
113 was not addressed. These early Venus habitability scenarios would have taken place within
114 the first billion years of the planet’s evolution limiting the possibility for complex life to
115 evolve. Venus would have subsequently entered a moist/runaway greenhouse and left us
116 with the state it is in today. For the curious reader Bullock and Grinspoon (1999) pro-

117 vide a nice review of the literature on the possibility of an early habitable Venus. Note
118 that in a previous work (Way et al., 2016) it was shown that if habitable conditions were
119 possible in early Venus' history that it likely remained so and that increasing solar in-
120 solation through time is not a deciding factor. In the present work we find the same, and
121 hence believe that limited (in time) early habitability models are not supported by our
122 simulations.

123 Grinspoon and Bullock, in a number of abstracts (e.g., Grinspoon & Bullock, 2003;
124 Grinspoon, 2008), described a Venus climate evolution scenario with long-lived surface
125 habitability consistent with that described later in this paper. This work was never pub-
126 lished, though, so the details of their calculations and the justifications for their con-
127 jectures are not documented.

128 The first three-dimensional (3-D) General Circulation Model (GCM) exploration
129 of issues relevant to ancient Venus was that of Abe et al. (2011). Their study focused
130 on land planets (analogous to the planet Arrakis of *Dune*) with no permanent surface
131 water reservoirs, only limited fixed amounts of ground water. In some scenarios they found
132 that “in principle Venus could have been a habitable land planet as recently as 1 billion
133 years ago.” Their conclusions resulted from the limited water vapor in the atmosphere
134 caused by the modest subsurface reservoir, its limited greenhouse effect, and the albedo
135 of the planet. Thereafter Leconte et al. (2013) used a 3-D GCM to look at climate sce-
136 narios for Gl581 c and HD 85512 b. HD 85512 b orbits a K-dwarf with a synchronous
137 rotation period of ~ 58 days and receives approximately 1.86 times the insolation of present
138 day Earth. This is slightly less than present day Venus' insolation (1.9), but its host star
139 is a K-dwarf rather than a G-dwarf and hence the peak of its Planck blackbody spec-
140 trum is at longer wavelengths. They state, “if not for its thick atmosphere, Venus' cli-
141 mate would be very close to the one of HD 85512 b.” They modeled a dry planet akin
142 to that of Abe et al. (2011) using an N_2 dominated atmosphere with 376ppm CO_2 , and
143 no water vapor. The model produced large temperate regions for a number of different
144 atmospheric pressures (See Leconte et al. (2013) Figures 1, 2) for their land planet setup,
145 but they also found similar behavior in their ‘collapse’ scenario with different atmospheric
146 water vapor profile amounts.

147 A rationale for high albedo cloud cover on ancient Venus was first presented by Yang
148 et al. (2014). Yang et al. (2014) was a large parameter study looking at the inner edge
149 of the habitable zone around solar type stars using The National Center for Atmospheric
150 Research (NCAR) Community Atmosphere Model (CAM) with comprehensive atmo-
151 spheric physics but a thermodynamic ocean. A thermodynamic ocean, also known as a
152 slab or mixed-layer ocean, is typically of limited depth (<100 meters) with a prescribed
153 horizontal ocean heat transport, or no horizontal heat transport at all. The latter im-
154 plies that the temperature of each ocean grid cell is determined solely by the atmosphere
155 and incident sunlight directly above it. The shallow depth reduces the time lag between
156 solar forcing and ocean response. See Way et al. (2017) Section 2.2.2 for details. In essence
157 they stepped the sidereal rotation rate of an Earth-like world from 1 to 256 sidereal days.
158 At the same time they increased the insolation from that of modern Earth to as much
159 as 2.6 times the modern Earth for their most slowly rotating world of 256 sidereal days
160 period. One of their key conclusions was that the slowest rotators would have had a day-
161 night general circulation that would generate an optically thick contiguous cloud bank
162 in the substellar region. This cloud deck would greatly increase the planetary albedo,
163 keeping the surface temperature moderate even for quite high values of insolation. These
164 results were later confirmed with a completely different 3-D GCM with a fully-coupled
165 dynamic ocean (Way et al., 2018). Yang et al. (2014) also included a simulation with mod-
166 ern Venus orbital parameters, spin rate, insolation while using modern Earth topogra-
167 phy and land/ocean mask, but again with a thermodynamic ocean. These studies pro-
168 vide a possible rationale for the cloud cover needed to produce temperate surface con-
169 ditions as first postulated by Pollack (1971). Finally Way et al. (2016, hereafter Paper

170 I) took things one step further by exploring different topographies, insulations and ro-
171 tation rates to put tighter constraints on possible habitable conditions for ancient Venus.

172 In this paper, we extend those parameter studies to consider a wider variety of plan-
173 ets and attempt to justify our modeling assumptions in light of possible scenarios for Venus’
174 evolution. Sections 2–5 review existing observational constraints and hypotheses about
175 the composition, thickness, and evolution of Venus’ atmosphere and water history; its
176 surface and interior; and its rotation and obliquity. In Section 6 we draw upon this in-
177 formation to inform a series of 3-D global climate model simulations to illustrate pos-
178 sible scenarios for an early habitable Venus that transitioned to its current inhospitable
179 state relatively late in its history. We discuss the implications of our results for the de-
180 sign of future missions to Venus and for the potential habitability of exoplanets inside
181 the inner edge of the traditional “habitable zone” in Section 7. Finally, recently published
182 complimentary work by Weller and Kiefer (2019) supports many of our conclusions.

183 2 Atmospheric Composition and Pressure

184 Assuming that Venus and Earth formed from the same parts of the protoplanetary
185 disk and thus with similar compositions (e.g., Raymond et al., 2004), we are guided by
186 the history of Earth whose early atmosphere was likely CO₂-rich and possibly cool, but
187 not frozen due to the faint young sun for the late Hadean and early Archean (4.2–4.0 Ga)
188 (e.g., Owen et al., 1979; Kasting, 1993; Zahnle & Sleep, 2002; J. W. Valley et al., 2002;
189 Zahnle, 2006; Zahnle et al., 2010; Kunze et al., 2014; Catling & Kasting, 2017; Krissansen-
190 Totton et al., 2018; Mello & Friaça, 2019; S. J. Mojzsis et al., 2019). The picture of a
191 “Cool Early Earth” promoted in the early 2000s by J. W. Valley et al. (2002); J. Val-
192 ley (2005) and with more recent zircon data (Valley et al., 2014) may also be applica-
193 ble to Venus’ early evolutionary history if Venus survived its magma ocean (MO) phase
194 with some liquid surface water. Chassefière et al. (2012) was probably the first to make
195 the comparison of Hadean Earth with Venus, calling his hypothesis a “cool early Venus.”

196 The prospects for a cool early Earth and Venus hypothesis have improved in re-
197 cent years, while the likelihood for their end in a cataclysmic “late instability” known
198 as the Late Heavy Bombardment (LHB) may have started to fall out of favor. The LHB
199 is important because its intensity was imagined so intense that the surface temperatures
200 of Earth would return to those before the cool early Earth period, perhaps evaporating
201 the oceans and/or raising surface temperatures to values exceeding 100°C. However, in
202 recent years a number of studies have started to question the timing and strength of the
203 LHB (e.g. Boehnke & Harrison, 2016; Zellner, 2017; Morbidelli et al., 2018; Quarles &
204 Kaib, 2019; S. J. Mojzsis et al., 2019). For a brief overview of the latest work on the LHB
205 see Mann (2018) and Voosen (2020). An “early instability” has started to gain favor in
206 what is termed an “accretion tail” scenario. This accretion tail scenario has recently been
207 promoted in the works of Morbidelli et al. (2018, Figure 1) and S. J. Mojzsis et al. (2019,
208 Figure 1). In these works a late or early heavy bombardment is triggered by a period
209 of orbital ‘instability’ experienced by the giant planets as described in de Sousa Ribeiro
210 et al. (2020). More recent work by de Sousa Ribeiro et al. (2020) demonstrates that the
211 instability has a median timescale of 36.78–61.5 *Myr* and is within 136 *Myr* in 75% of
212 their cases. de Sousa Ribeiro et al. (2020) say the timing of an early instability fits in
213 nicely with the survival of the Patroclus–Menoetius Jupiter Trojan as a primordial bi-
214 nary from the Kuiper Belt (Nesvorný et al., 2018) and is a good match to other solar
215 system properties (e.g. Clement et al., 2019), but they explicitly discuss the many lim-
216 itations of their model. In the spirit of an early instability and accretion tail scenario,
217 Figure 4 of S. J. Mojzsis et al. (2019) suggests that Venus accreted less than 0.01 wt%
218 between 4.3 and 4.1 Ga, about the same as Earth. As S. J. Mojzsis et al. (2019) notes,
219 “Results show that an abating impact flux from late accretion is inadequate to steril-
220 ize the surface zone.” These works demonstrate the possibility that the cool early Earth
221 continued through the originally proposed time period of the LHB from \sim 3.8–4.1 Ga.

Hence we will begin our most ancient Venus simulations with CO₂ dominated atmospheres at ~ 4.2 Ga (during the J. W. Valley et al. (2002) cool early Earth period) that evolves over time to N₂ dominated atmospheres. If the LHB were a real event we would want to begin our simulations at 3.8 Ga, rather than 4.2 Ga, as others have speculated previously (e.g., Rampino & Caldeira, 1994; Lammer et al., 2018). There are still many unanswered questions regarding the early post-MO history of the atmospheres of Venus, Earth and Mars. Even though we have far more information to discern Earth's early post-MO atmosphere it remains a complicated story yet to be fully resolved (e.g., Hirschmann, 2012; Nikolaou et al., 2019) and may depend on atmospheric pressure more than previously assumed (e.g., Gaillard & Scaillet, 2014; Bower, Dan J. et al., 2019).

The carbonate-silicate cycle (e.g., Walker et al., 1981; Stewart et al., 2019) is the key to keeping most of Earth's CO₂ locked up in rocks for much of its history rather than in the atmosphere as on present day Venus. The carbonate-silicate cycle on Earth functions via subductive-type plate tectonics, the presence of a hydrosphere and continental crust. For Earth there are several lines of evidence to suggest these may go back to the Hadean (e.g., S. Mojzsis et al., 2001; Hopkins et al., 2008; Harrison, 2009; Korenaga, 2013; Harrison et al., 2017; O'Neill et al., 2017; Rozel et al., 2017; Kusky et al., 2018; Korenaga, 2018; Maruyama et al., 2018, and references therein). Dehant et al. (2019) reviews the literature for a later beginning of plate tectonics on Earth (Section 3.2). The requirements for how plate tectonics begins are still not fully understood and hence remains an active area of research (e.g., Lenardic et al., 2019). If Venus had a similar early atmospheric and interior evolution to that of Earth then early volatile cycling via some form of plate tectonics is a viable hypothesis. In fact recent work demonstrates convection regimes like that of plate tectonics in Venus' recent history (Gillmann & Tackley, 2014; Davaille et al., 2017). Dehant et al. (2019) reviews plate tectonic mechanisms that may not require as much water as previously believed, which may be relevant if Venus retained some water but less than Earth. Hence in our scenario a carbonate silicate cycle is hypothesized for Venus after the magma ocean phase and well before its resurfacing period to keep CO₂ largely sequestered in crustal carbonates as on modern Earth.

We are motivated to choose 1 bar atmospheres for our epochs of interest based on geological atmospheric pressure proxies for Earth that imply an atmospheric density of ~ 0.25 –1.1 bar for most of the past 4 billion years (Nishizawa et al., 2007; Goldblatt et al., 2009; D. Som S. Catling et al., 2012; S. Som et al., 2016; Marty et al., 2013; Avicé et al., 2018). However, some models of Nitrogen cycling imply that the atmospheric pressure could have been higher in the Archean (Johnson & Goldblatt, 2015; Mallik et al., 2018). Differences in N₂ of factors of a few have relatively small net impacts on climate, primarily due to small decreases/increases in the pressure broadening of CO₂ lines and partly offsetting decreases/increases in Rayleigh scattering. Given the likely similar geochemistry and volatile histories of Venus and Earth (e.g., Ringwood & Anderson, 1977; Lécuyer et al., 2000; Treiman, 2009; Chassefière et al., 2012; Rubie et al., 2015), a similar evolution in their early composition and pressure histories is plausible. Although whether they did indeed start out with similar volatile inventories is still an open research question (e.g. Horner et al., 2009).

Additional work demonstrates that the present day CO₂ and N₂ inventories of Earth and Venus may be similar (Donahue & Pollack, 1983; Goldblatt et al., 2009) if Venus is more degassed than Earth as Donahue and Pollack (1983) speculated. Early work by Rasool and de Bergh (1970) and Kasting (1988) estimated nearly as much CO₂ was locked up in carbonate rocks on Earth as exist in the atmosphere of Venus. Since there is no carbonate-silicate cycle active on Venus today most of the CO₂ that would otherwise be locked up in the interior is in the atmosphere, having degassed over the past several hundred million years (at least). As well, nearly 4 times as much N₂ is found in Venus' present day atmosphere as in Earth's when scaled by planetary mass. Earth's internal N₂ bud-

274 get is estimated at $\sim 7 \pm 4$ times that of the atmospheric mass (Johnson & Goldblatt,
275 2015).

276 However, Argon (Ar) modeling studies (Kaula, 1999; O'Rourke & Korenaga, 2015)
277 imply that Venus is less degassed than Earth, although see Donahue and Pollack (1983)
278 and Halliday (2013) for alternative explanations for the current ^{40}Ar abundance differ-
279 ence between Earth & Venus. Watson et al. (2007) claims that ^{40}Ar is not a reliable in-
280 dicator of degassing, although "the interpretation of their data is controversial" accord-
281 ing to O'Rourke and Korenaga (2015). Halliday (2013) also mentions that Venus' radio-
282 genic ^{40}Ar implies Venus is less degassed than Earth, but he notes that Venus' atmospheric
283 non-radiogenic ^{36}Ar has concentrations roughly two orders of magnitude greater than
284 Earth's and that "this is hard to explain unless it (Venus) is heavily degassed with a larger
285 inventory of primordial volatiles." In further support of a degassed Venus Halliday (2013)
286 notes that the similar Earth and Venus budgets of C and N support a degassed Venus
287 while helping to explain the high ^{20}Ne and ^{36}Ar budgets of Venus. One of the largest
288 stumbling blocks in deciding whether the ^{40}Ar modeling inference is correct is a proper
289 measurement of K/U on Venus, which is presently highly unconstrained. For this rea-
290 son modeling studies generally use Earth values. Namiki and Solomon (1998) use He to
291 confirm the ^{40}Ar estimates, but they require a Venus in-situ mission to make the appro-
292 priate Th and U abundance measurements to characterize the geochemistry necessary
293 for He production. One also needs to consider He escape, an under-explored area of re-
294 search given that it depends not only on the ability of He to remain charged or neutral,
295 but also whether Venus had a past magnetic field and how present day measurements
296 can distinguish in situ He in the upper atmosphere from solar wind deposition.

297 Another limiting factor in comparing Earth and Venus is the lack of good constraints
298 on the bulk water content of Earth. Some papers refer only to surface water amounts
299 of 0.023 wt% of the Earth's total mass, but there are compelling arguments for possi-
300 bly deep reservoirs within the lower mantle or core (e.g., Raymond et al., 2006; Schmandt
301 et al., 2014; Genda, 2016; Ikoma et al., 2018). For Venus, interior water amount estimates
302 are mostly speculation and are restricted to planetary formation studies (e.g., Ikoma et
303 al., 2018).

304 **3 Venus' Early Evolution & Evidence for Water**

305 Venus & Earth likely received similar initial water inventories during their forma-
306 tion histories, as has been shown in a number of works (e.g., Raymond et al., 2006). It
307 is also tied to their composition in general, as discussed above in Section 2.

308 The Pioneer Venus mass spectrometer measured a very high D/H ratio of 150 ± 30
309 times that of terrestrial water (Donahue et al., 1997) in a trapped droplet of sulfuric acid.
310 It is the only such published in-situ measurement. Other non in-situ measurements have
311 been made as noted in Section 1, and work continues apace (e.g. Tsang et al., 2017). The
312 original Venus D/H discovery paper by Donahue et al. (1982) was titled "Venus was Wet:
313 A measurement of the Ratio of Deuterium to Hydrogen." This was a tantalizing prospect,
314 but limited by our knowledge of atmospheric escape processes (Donahue & Pollack, 1983;
315 Donahue & Russell, 1997; Donahue, 1999) and the D/H of delivered materials over the
316 aeons.

317 It is possible that the D/H ratio is not an indicator of large amounts of water in
318 Venus' ancient history. Grinspoon (1993) pointed out that a short residence time for wa-
319 ter in the present atmosphere of Venus works against the primordial ocean hypothesis.
320 Grinspoon (1993) also noted that updated theoretical calculations at that time that im-
321 plied higher deuterium escape efficiency put constraints on the D/H source water of 10-
322 15. That would rule out source material such as meteorites, comets and dust particles
323 with high D/H ratios (e.g., Irvine et al., 2000; Charnley & Rodgers, 2009). As well, mea-

324 surements of D/H and Xenon isotopes in Comet 67P (Altwegg et al., 2015; Marty et al.,
325 2017) imply that Earth’s ocean has a much lower contribution from cometary objects
326 than previously thought. This would also likely rule out a large cometary contribution
327 to the high D/H ratio measured on Venus.

328 A number of authors have tried to model changes in the D/H fractionation over
329 time (Kasting & Pollack, 1983; Gurwell & Yung, 1993; Gurwell, 1995; Hartle et al., 1996)
330 to put some constraints on when the water was lost. Given the lack of data from Venus
331 it is equally difficult to constrain or move these models forward.

332 The possibility that the high D/H ratio implies long-lived surface water is also lim-
333 ited by in-situ measurements. There is some circumstantial evidence of past surface wa-
334 ter from surface emissivity observations from the Galileo NIMS instrument (Hashimoto
335 et al., 2008) and the Venus Express VIRTIS instrument (Mueller et al., 2008). These ob-
336 servations may imply that the highland “tessera” regions are mostly composed of felsic
337 rocks, and if they are indeed granitic they would have required surface water to form (Campbell
338 & Taylor, 1983). M. S. Gilmore et al. (2015); M. Gilmore et al. (2017) find that at least
339 one tessera region observed with VIRTIS (Alpha Regio) appears to be more felsic than
340 surrounding plains. This also suggests that these older stratigraphic units (Ivanov & Basilevsky,
341 1993; M. S. Gilmore et al., 1997) are granitic crustal remnants, but recent work by Wroblewski
342 et al. (2019) shows that parts of the Ovda Regio highland tessera are not in fact of granitic
343 origin. However, it is not yet possible to generalize the work of Wroblewski et al. (2019)
344 to the entirety of tessera.

345 Nikolayeva (1990) and Shellnutt (2019) analyzed surface rock measurements from
346 Venera 8. As Shellnutt (2019) explains “it is possible that the Venera 8 probe encoun-
347 tered a fragment of crust that resembles a terrestrial greenstone belt.” Zolotov et al. (1997)
348 and Johnson and Fegley (2000) have also demonstrated that signatures of water in hy-
349 drous minerals may persist on the surface of Venus for long periods even after the sur-
350 face morphology has changed. This motivates an in-situ mission to Venus to search for
351 such materials, if they exist. Watson et al. (2007) tried to demonstrate that the ^{40}Ar in
352 Earth’s atmosphere is related to the hydration of the oceanic lithosphere consisting of
353 relatively Ar-rich olivine and orthopyroxene. If the results from Watson et al. (2007) are
354 correct, (and there is skepticism (Ballentine, 2007)), this would lead one to believe that
355 the ^{40}Ar in Venus’ atmosphere today implies that water oceans could have persisted for
356 some time.

357 An outstanding unsolved and understudied problem is what happened in the epoch
358 of Venus’ MO as it cooled, as this may greatly affect the long-term water inventory of
359 the planet. The timescale of the MO crystallization could be of order a few million years
360 (*Myr*) as for Earth (e.g., Katyal et al., 2019; Nikolaou et al., 2019) or greater than 100
361 *Myr* (Hamano et al., 2013; Lebrun et al., 2013). The longevity of the MO and associ-
362 ated hot steam and CO_2 atmosphere is vital to understanding the volatile history of Venus
363 (e.g. Salvador et al., 2017). If the MO and steam atmosphere persist too long then much
364 of the primordial water inventory of Venus could have been lost in its very early history
365 from a stronger solar wind (Chassefière, 1997; Lichtenegger et al., 2016). An attractive
366 feature of the extended MO hypothesis is that it naturally solves the problem of the lack
367 of oxygen in the present day Venusian atmosphere. This would be accomplished by se-
368 questrating the O_2 left behind by H_2O dissociation in the magmatic crust and upper man-
369 tle (e.g. Lebrun et al., 2013; Gillmann et al., 2009; Lichtenegger et al., 2016; Lammer
370 et al., 2018). We discuss how large quantities of O_2 can be lost after a significant period
371 of habitability in bullet 6 of Section 8, in lieu of early MO losses. It is possible that the
372 high D/H ratio we see today (Donahue et al., 1982, 1997) is a relic of the early MO pe-
373 riod. If the MO cooled quickly, then there was an opportunity to build up a surface ocean
374 and atmosphere as is believed to have happened in Earth’s early history. The question
375 is whether Venus’ surface conditions as a result of its closer proximity to the Sun would
376 prevent the condensation of water on its surface or not. The answer is more complicated

377 than it may seem since water can condense under hot high pressure multi-bar atmospheres.
378 Matsui and Abe (1986) allow for temperatures up to 600K, while later work by Liu (2004)
379 allow temperatures approaching 720K.

380 Additionally the answer may reside in the planet's rotation history, what role clouds
381 played, and the outgassing rates of H₂O and CO₂. As we will show in Section 5 it is possible
382 for Venus to reach a tidally locked state in less than a few hundred *Myr* using constant
383 phase lag dissipation theory, suggesting that the planet's rotation rate could have
384 been slow early on. As shown in previous work (Way et al., 2016, 2018) as long as a planet
385 is in the slowly rotating regime (length of day greater than ~16 Earth sidereal days) its
386 climate dynamics work to allow liquid water to persist on the surface for insulations up
387 to ~ 2.6 times that of present day Earth. This is due a large contiguous dayside cloud
388 deck that significantly increases the planetary albedo as discussed in Section 4.

389 The timing of the MO termination is critical in more than one way. If the steam
390 and CO₂ atmosphere cooled sufficiently for MO crystallization to occur by the time of
391 the Late Veneer (also referred to as "Late Accretion") then even if Venus lost most/all
392 of its primordial H₂O through escape processes (Gillmann et al., 2009; Hamano et al.,
393 2013; Lichtenegger et al., 2016) there may have been a second chance to obtain a sur-
394 face ocean, albeit a shallow one. Recent work by Greenwood et al. (2018) implies that
395 Earth may have received as much as 30% of its H₂O inventory in post-accretion impact
396 delivery, consistent with research that shows that the entire H₂O budget cannot come
397 from the late veneer (Morbidelli & Wood, 2015). Halliday (2013) concludes that if ve-
398 neers were common they should be proportional to planetary mass, and hence Venus would
399 have received a percentage of late veneer H₂O similar to that of Earth. If Venus was left
400 dry after a long-lived magma ocean phase (Hamano et al., 2013), then this amount of
401 H₂O veneer also fits within the error bounds of Venus' measured D/H ratio (Donahue
402 et al., 1982, 1997). It should be noted that the work of Greenwood et al. (2018) can also
403 fit within the Ruthenium studies of Fischer-Gödde and Kleine (2017). For a contrary point
404 of view see Gillmann et al. (2019), who claim that most of the late veneer impactors would
405 have been Enstatite/ordinary chondrites which are water-poor, as opposed to water-rich
406 carbonaceous chondrites that would have been a mere 0-2% of the total chondrite de-
407 livery. These contrary points of view come about because different geochemical measure-
408 ments give different answers as pointed out in a number of recent works (e.g. Albarède,
409 2009; Fischer-Gödde & Kleine, 2017; Dauphas, 2017; McCubbin & Barnes, 2019; Zahnle
410 et al., 2019). There are two other important caveats to consider regarding the late ve-
411 neer. First the water content depends upon the composition of the accreting bodies. For
412 example, if the late veneer was made up of a few large bodies then the variations could
413 have been greater than if it was due to a large collection of smaller bodies. Second, the
414 definition of the late veneer is important since it is typically associated with the accre-
415 tion of bodies after the last giant impact. Jacobson et al. (2017) has suggested that the
416 last giant impact on Venus could have been much earlier than on Earth and this obvi-
417 ously affects the composition of objects making up the late veneer.

418 At the same time, work by Gillmann et al. (2009), Morbidelli et al. (2000), and Raymond
419 et al. (2006) shows that Venus' initial water inventory at formation could be as much
420 as two terrestrial ocean's worth while large planetary embryos could deliver much more
421 within 200 *Myr* of formation. If true, and if the magma ocean lifetime on Venus was shorter
422 rather than longer, then our estimates of the water content on Venus from Pioneer Venus
423 D/H ratios (Donahue et al., 1982, 1997) should be more toward the higher end, ~ 16%
424 of a present day Earth's ocean (Donahue et al., 1997). However, it is not clear whether
425 Venus' primordial water content can readily be constrained by the D/H ratio (Grinspoon,
426 1987; Grinspoon, 1993) due to a lack of knowledge of sources and sinks over the lifetime
427 of the planet.

428 Another hypothesis (e.g., Rampino & Caldeira, 1994) states that because of its prox-
429 imity to the Sun, Venus could never condense water on its surface and hence its surface

430 temperature has always been 300K or higher (see Figure 1 in Rampino & Caldeira, 1994)
 431 and that most of this water was lost by photodissociation (Goody & Walker, 1972). At
 432 the same time the lack of water prevents silicate rock weathering (on Earth this removes
 433 CO₂ from the atmosphere), hence the CO₂ builds up in the atmosphere driving temper-
 434 atures ever higher due to the greenhouse effect as seen today.

435 **4 Surface History, Impactors and Climate Evolution**

436 Understanding the surface history of Venus is crucial to constraining any theory
 437 of its long-term climate evolution. Smrekar et al. (2018) reviews the literature on Venus'
 438 internal structure and dynamics. In this section we mostly focus on implications for the
 439 surface features we see today and how those might be consistent with a hypothesis for
 440 the long-term habitability of Venus and a transition to a more recent (~ 1 Gyr) hothouse
 441 state.

442 Up to 80% of the Venus surface has volcanic plains and tectonic structures emplaced
 443 over a relatively short geological interval as determined from crater counts (Ivanov & Head,
 444 2013, 2015). The cratering record seen in the plains regions imply surface ages ranging,
 445 for example, from ~ 180 Ma (Bottke et al., 2016), to ~ 300 Ma (Strom et al., 1994) to ~ 750 Ma
 446 (McKinnon et al., 1997). The relative youth of most of Venus' surface may be the re-
 447 sult of a large scale lithospheric overturn known as the Global Resurfacing Event (GRE),
 448 or it may be due to the latest GRE in a long sequence of episodic resurfacing events (e.g.,
 449 Turcotte, 1993; Strom et al., 1994). For example, Kaula (1999) constructed a simple model
 450 with outgassing events staggered at time periods of 4.1, 3.8, 3.5, 3.1, 2.6, 2.1, 1.5 and
 451 0.7 Ga constrained by ⁴⁰Ar measurements. The other hypothesis for the young surface
 452 of Venus is from continuous volcanic resurfacing (e.g. Basilevsky et al., 1997; Bjonnes
 453 et al., 2012; King, 2018).

454 The highland tesserae may be one of the keys to understanding this history. They
 455 are of particular interest because they may contain information about past crustal dif-
 456 ferentiation and other processes prior to the loss of any surface water. Some crater age
 457 estimates from the Magellan Mission imply that the tesserae are $\sim 40\%$ older than the
 458 plains (Ivanov & Basilevsky, 1993; M. S. Gilmore et al., 1997). However, Strom et al.
 459 (1994) did not agree with this conclusion. Additional work by V. Hansen and López (2010)
 460 points to the possibility that the Ribbon Tessera Terrain are older than the surface units
 461 identified with the GRE. Later analysis by Ivanov and Head (2013) implied that tessera
 462 are the oldest stratigraphic unit and that they were created near the beginning of Venus'
 463 surface observable history during the “tectonically dominated regime.”

464 However, *how* the large basaltic plains were emplaced remains controversial. A num-
 465 ber of authors (e.g., Herrick, 1994; Strom et al., 1994; Basilevsky & Head, 1996) postu-
 466 lated a nearly global ($\sim 80\%$) geologically instantaneous (10-100Myr) thick (>1 km) de-
 467 position of basaltic material from volcanic type outflows (GIBVO) that would have buried
 468 older craters we cannot observe today (akin to the GRE mentioned above). The outflow
 469 depth requirements are determined by the size of the largest impact craters that would
 470 have to be completely covered. However, as Ivanov and Head (2013) point out it is pos-
 471 sible that the cratering record previous to GIBVO could have also been erased in some
 472 manner. The GIBVO model was later augmented and became known as the global stratig-
 473 raphy hypothesis (e.g., Basilevsky & Head, 1996; Basilevsky et al., 1997; Basilevsky &
 474 Head, 1998; Head & Basilevsky, 1998). Yet another hypothesis to explain the Venus sur-
 475 face record was initially put forward by Phillips et al. (1992) and is termed the Equilib-
 476 rium Resurfacing Model (ERM). In this model the number of craters observed on Venus
 477 today is the result of an equilibrium between constant crater formation (via impacts) and
 478 the removal of such craters via on-going tectonic or volcanic methods. Monte Carlo cal-
 479 culations by Bullock et al. (1993) and Strom et al. (1994) demonstrated why the ERM
 480 was not feasible. Strom et al. (1994) decided that the GIBVO was a better fit to their

481 data, while Bullock et al. (1993) preferred a longer timeline of 550 *Myr*. More recent Monte
482 Carlo calculations by Bjonnes et al. (2012) show that the ERM is able to fit the obser-
483 vations.

484 V. Hansen and Young (2007) strove to demonstrate why none of these hypothe-
485 ses fit all available observational constraints. V. Hansen and Young (2007) then proposed
486 what they termed the Spatially Isolated Time-Transgressive Equilibrium Resurfacing (SPIT-
487 TER) hypothesis to explain more of the observational constraints. It is not clear that
488 the Venus geological community has settled on any of these hypotheses. Perhaps one of
489 the largest problems with the global lava hypothesis is the timescale, volume and depth
490 of the basaltic flows required, none of which have been observed on any present or pre-
491 viously active volcanic body in the solar system (including in Earth's past). The largest
492 known outflow to date in Earth's history is the mid-Cretaceous Superplume (Larson, 1991),
493 which is small by comparison to those envisioned to describe Venus' resurfacing. At the
494 same time the superplume hypothesis for Venus is compelling as large amounts of CO₂
495 could have been released at the same time as the plume event (Caldeira & Rampino, 1991).
496 Large overturn events have been proposed as an explanation for Venus' present surface
497 state, but in such a scenario it is possible to sequester large amounts of CO₂ in fresh flood
498 basalt outflows due to enhanced planetary weatherability (e.g., Godd ris et al., 2003; Cox
499 et al., 2016). Large Igneous Provinces (LIPs), on the other hand, can release copious amounts
500 of CO₂ sequestered in some sedimentary materials (e.g., Ganino & Arndt, 2009) while
501 avoiding the sequestration issues of a large overturn event. LIPs have been proposed as
502 an explanation for Venus' present day state as we will discuss below.

503 Previous simulations by Way et al. (2016) showed that Venus could have had tem-
504 perate conditions for nearly 2 billion years providing it had a shallow ocean of 310m in
505 depth, slow rotation rate, and modern orbital elements. Venus might even have experi-
506 enced more stable conditions than Earth in its early history since studies by Correia
507 and Laskar (2001) and J. W. Barnes et al. (2016) have shown that low obliquity states
508 (like that of modern Venus) may be stable over billions of years and we know that the
509 much shorter Milankovich cycles have had a strong influence on Earth's climate through
510 time. Deitrick et al. (2018) reviews the influence of such cycles on the climate of Earth
511 and possible influences on exoplanets. In addition, Weller et al. (2018) has also shown
512 from geological models that early Venus could have avoided glaciations more easily than
513 early Earth, which experienced several partial or total snowball periods in its history.
514 If long-term stable surface conditions are a requirement for life, Venus might have been
515 more stable and allowed primitive life to fill more ecological niches more quickly than
516 on Earth. This gives rise to the possibility that life may still exist in Venus' upper at-
517 mosphere (Limaye et al., 2018).

518 Ernst et al. (2017) speculate that "On Venus, voluminous LIP volcanism produced
519 high levels of CO₂ that led to run-away greenhouse effect, and high levels of SO₂ that
520 caused acid rain," but with little supporting evidence. Bullock and Grinspoon (2001)
521 present a similar hypothesis that involves outgassing of SO₂ and H₂O that eventually
522 drive the planet, over 100s of *Myr*, into a runaway greenhouse state, but do not men-
523 tion CO₂. If Venus had LIP volcanism then CO₂ as well as SO₂ can be outgassed if trapped
524 in sediments in the crust as is seen on Earth (e.g., J. W. Head III & Coffin, 1997; V. L. Hansen,
525 2007; Ernst et al., 2017; Ernst & Youbi, 2017). Hence if Venus had an earlier epoch of
526 liquid water habitability then it is logical to assume that CO₂ would have been trapped
527 in the crust of the planet in the same way it is trapped on Earth today and LIP volcan-
528 ism would have been the means to release that CO₂ into the atmosphere.

529 However, as noted in Macdonald and Wordsworth (2017) when the surface tem-
530 perature is warmer ($T > 300\text{K}$, see their Figure 2) more water vapor is injected into the
531 stratosphere, which stabilizes the lapse rate. Such warm climates (as seen in the Venus
532 models herein) would prevent the largest plumes from injecting SO₂ into the stratosphere,
533 allowing CO₂ warming without offsetting cooling by H₂SO₄ aerosols.

534 Another well known mechanism to get Venus from a cool clement state to its present
535 day hot and dry state was proposed by a number of authors (e.g. Ingersoll, 1969; Kast-
536 ing & Pollack, 1983; Kasting et al., 1984; Kasting, 1988; Taylor & Grinspoon, 2009) who
537 speculated that water loss via upper atmospheric dissociation and then hydrogen escape
538 would have eventually made the planet dry. Then, as stated in Taylor and Grinspoon
539 (2009) “With the loss of water, the removal mechanism for CO₂ would be eliminated,
540 and carbonate rocks on the surface would presumably eventually be subducted and lost
541 to thermal decomposition, with the CO₂ being irreversibly returned to the atmosphere
542 through outgassing.” This model fits in with more recent research by R. D. Wordsworth
543 (2016a) who states that the oxygen left over would eventually find its way to oxidize the
544 mantle and change its redox state, allowing for enhanced nitrogen outgassing which is
545 compatible with the nearly 3 bars of N₂ we see in Venus’ atmosphere today (also see re-
546 view by Lammer et al. (2018)). However, an alternative hypothesis is proposed by Gillmann
547 et al. (2009), who suggest that the oxidation of the mantle occurred in the first 100 *Myr*
548 of Venus’ history. They assume the surface was never cool enough to allow liquid wa-
549 ter to condense. The water would again be photodissociated and the hydrogen would
550 have been lost to space (Lichtenegger et al., 2016). The leftover oxygen would have dis-
551 solved in the magma ocean.

552 Genda and Abe (2005) have proposed that the lack of water on Venus and in Venus’
553 protoplanetary impactors in its early history (in contrast to that of Earth and its wa-
554 ter rich impactors) would explain differences in most of the noble gas abundances be-
555 tween Venus and Earth because oceanic protoplanets would enhance atmospheric loss,
556 implying that Venus’ original noble gas abundant proto-atmosphere survived to present
557 day on Venus, unlike that of Earth. A lack of water being detrimental to subductive plate
558 tectonics (see Section 8). Sakuraba et al. (2019) have also attempted to get the presently
559 observed nitrogen and noble gas abundances via impact degassing and atmospheric ero-
560 sion (also see work by, Pham et al., 2011), but unlike Genda and Abe (2005) they be-
561 lieve late accretion may have further influenced the atmosphere of Venus.

562 More recent work by Gillmann et al. (2016) show that large impactors (400-800km
563 in diameter) can cause atmospheric erosion and escape and deposit energy in the crust
564 and mantle. They believe the latter can cause a thermal anomaly in the crust and man-
565 tle triggering large scale volcanic events at the impact region and the antipode. This in
566 turn may deplete the upper mantle of volatiles and lead to water loss in the early atmo-
567 sphere, or conversely provide a volatile heavy atmosphere with extreme temperatures
568 for billions of years. In a sense this is similar to a theory by Davies (2008) who propose
569 a mega-collision (akin to that of the Earth’s moon-forming impact) to dry out the in-
570 terior of the planet. But thus far no large Venus impactor simulations have been utilized
571 to examine such a scenario, as has been done for Earth’s moon-forming collision (e.g.,
572 Canup, 2004).

573 To summarize, a number of mechanisms exist by which early Venus could have con-
574 densed liquid water on its surface. The key ingredient is that it must have been cool enough
575 for long enough in its early history. As shown by Yang et al. (2014); Way et al. (2016,
576 2018), the rotation rate of a planet greatly affects its climate dynamics. Specifically, for
577 very slow rotation a large contiguous water cloud forms at the substellar point, increas-
578 ing the Bond albedo markedly and keeping surface temperatures moderate for insola-
579 tion values up to nearly three times that of modern Earth’s 1361 W m⁻². In Paper I Way
580 et al. (2016) we demonstrated that early Venus could have had consistently habitable
581 conditions throughout its early history if it began with sufficiently slow rotation. In the
582 next section we review what is understood about the possible evolution of Venus’ spin-
583 orbit state.

584 In our scenario, early Venus’ has the earliest consistent liquid water habitability
585 in the solar system followed by Earth and then Mars. This is a broader statement of the
586 Faint Young Sun Paradox (FYSP), the challenge of explaining how early Earth, not to

587 mention Mars, could have been warm and wet early in their histories when the Sun was
588 25-30% dimmer than today (e.g., Feulner, 2012). There is still debate in the ancient Earth
589 GCM community about the actual composition and thus temperature of early Earth's
590 atmosphere given observational proxies for CO₂ that span orders of magnitude, though
591 models suggest that the range encompasses several viable scenarios (e.g., Charnay et al.,
592 2013; Wolf & Toon, 2013; Kunze et al., 2014; Le Hir et al., 2014; Charnay et al., 2017;
593 Krissansen-Totton et al., 2018). These GCM studies and most proxies (e.g., Spencer, 2019)
594 are from the Archean rather than the late Hadean, but there is some evidence that hab-
595 itable surface conditions existed well back into the Hadean (e.g., Harrison, 2009; Arndt
596 & Nisbet, 2012).

597 It is interesting to note that recent atmospheric pressure proxies from the late Archean
598 imply an atmospheric pressure less than half that of today (D. Som S. Catling et al., 2012;
599 S. Som et al., 2016). Atmospheres thinner than modern Earth's are less likely to avoid
600 snowball conditions, yet the literature above notes that there is geological evidence that
601 Earth was not in a snowball state during much of the late Archean that the pressure prox-
602 ies correspond to. Regardless, for this reason we feel it is necessary to explore the pos-
603 sibilities of lower atmospheric surface pressures in Venus' climatic history as described
604 for Simulations 26-30 in Section 6.

605 The FYSP for Mars remains difficult to resolve (e.g., R. D. Wordsworth, 2016b)
606 partly due to the fact that 3-D GCMs have traditionally struggled to consistently sus-
607 tain large-area liquid water conditions over millions of years (e.g., Goldblatt et al., 2009;
608 Kasting, 2010; Kienert et al., 2012; Feulner, 2012; Haqq-Misra et al., 2008) without snow-
609 ball type conditions. Long-standing solutions involving large amounts of atmospheric CO₂
610 are inconsistent with unobserved carbonate deposits expected from such CO₂ dominated
611 atmospheres (Shaw, 2018) and are insufficient in isolation to produce above-freezing con-
612 ditions. One possible solution to the lack of surface carbonates was proposed by Kasting
613 (2012). Other solutions to Mars' FYSP exist that involve H₂ with CO₂ as the background
614 gas (e.g. R. Wordsworth et al., 2017; Ramirez, Kopparapu, Zugger, et al., 2014; Ramirez
615 & Craddock, 2018; Haberle et al., 2019), although presently there appears to be little
616 consensus in the community.

617 **5 Rotation and Obliquity evolution**

618 To the best of our abilities we would like to constrain the obliquity and rotational
619 history of Venus to better constrain these important inputs for climate models. This is
620 limited by the absence of any direct information about Venus' initial rotation and obliq-
621 uity and the fact that impacts likely play a significant role in the early rotational his-
622 tory of the terrestrial planets (e.g., Lissauer & Kary, 1991; Dones & Tremaine, 1993).
623 On Earth a variety of means exist to obtain some constraints using dynamical model-
624 ing combined with geological data when available (e.g., Hays et al., 1976; Park & Her-
625 bert, 1987; Imbrie et al., 1992; Matthews et al., 1997; Petit et al., 1999; Pälike & Shack-
626 leton, 2000; Pälike et al., 2004; Olsen et al., 2019) and there has been modest success
627 doing the same for Mars (e.g., Cutts & Lewis, 1982; Laskar et al., 2002, 2004; Byrne, 2009;
628 Dickson et al., 2015; Bierson et al., 2016). For Earth, an additional constraint is provided
629 by the Moon, which has predictably affected the evolution of Earth's rotation and damped
630 obliquity excursions over its history (Zahnle & Walker, 1987; Lissauer & Chambers, 2011).
631 However, until and unless geological observables become available to constrain dynam-
632 ical models, only plausible scenarios for the rotational and obliquity history of Venus can
633 be defined.

634 Hoolst (2015) summarizes much of the literature on the rotational evolution of Venus
635 throughout its history. We summarize some of the work on this subject below and add
636 some additional estimates. First we look at the history of studies of the possible spin evo-
637 lution of Venus.

638 In the 1960-70s several authors investigated the possibility that Venus' rotation pe-
 639 riod was correlated with its synodic period (Goldreich & Peale, 1966; Gold & Soter, 1969,
 640 1979). Goldreich and Peale (1966) states, "the presence of the Earth may have stabilized
 641 the sidereal rotation period of Venus at the value of 243.16 days retrograde." An equi-
 642 librium between the atmospheric and body tide of Venus was first proposed by Gold and
 643 Soter (1969) to explain Venus' non-synchronous rotation period, based on the incorrect
 644 belief at that time that Venus always showed the same face at each inferior conjunction
 645 with Earth as proposed by Goldreich and Peale (1966).

646 The first work to analytically look at Venus' rotation rate and the role of atmo-
 647 spheric tides was by Ingersoll and Dobrovolskis (1978) who extended the earlier work of
 648 Lord Kelvin Thomson (1882), Chapman and Lindzen (1970) and Munk and MacDon-
 649 ald (1960). They mention that "Venus probably originated with a retrograde rotation
 650 in order to have evolved to the current retrograde state." In the 1980s this work was fur-
 651 ther extended in a series of papers (Dobrovolskis & Ingersoll, 1980; Dobrovolskis, 1980,
 652 1983). It was clear that Venus' rotation rate was probably determined by a balance be-
 653 tween the solid body tidal dissipation and the thermal tides of its thick atmosphere with
 654 the sun. Core-mantle friction (CMF) can also play an important role in slowing the spin
 655 rate of Venus, as first explored by Goldreich and Peale (1970). Goldreich and Peale (1970)
 656 were also the first to demonstrate that core-mantle viscous coupling can drive the obliq-
 657 uity to 0° when less than 90° and to 180° if it is greater than 90° over time.

658 This remained the state of understanding of Venus' rotational history until the early
 659 2000s when the long-term evolution of its spin state of Venus was investigated in a se-
 660 ries of papers by Correia and Laskar (2001); Correia et al. (2003); Correia and Laskar
 661 (2003), who suggested that Venus may have rotated faster in the past, and possibly pro-
 662 grade. It also became clear that at faster spin rates CMF plays an important role in slow-
 663 ing the rotation of the planet, but less so at slower spin rates.

664 Once a planet is spinning more slowly CMF may play an important role in obliq-
 665 uity variations (e.g., Correia et al., 2003). Correia and Laskar (2001) explored a num-
 666 ber of stable obliquity and spin states of Venus while more recent work by J. W. Barnes
 667 et al. (2016) has investigated how stable the obliquity of Venus might be though time.

668 The work on the thermal tides of Venus had led researchers to assume that its ef-
 669 fects would be minor (as it is for Earth) for atmospheres of modest density (e.g. 1 bar).
 670 However, more recent work by Leconte et al. (2015) has demonstrated that thermal tides
 671 arising from even 1 bar atmospheres can be significant depending on the distance to the
 672 host star and the host star's mass. Leconte et al. (2015) show that even if modern Venus
 673 had a 1 bar atmosphere the tidal torques would still be quite significant.

674 R. Barnes (2017) used an equilibrium tide model with a constant phase lag (CPL)
 675 to find that Earth could have ended up tidally locked today (after 4.5Gyr) had it started
 676 with a rotation rate of 3 Earth Days or longer (the latter more likely if Earth had no satel-
 677 lite). We have applied the same Equilibrium Tide Model (EqTide <https://github.com/RoryBarnes/EqTide>)
 678 from R. Barnes (2017) to Venus to explore how long it would take Venus to reach a tidally
 679 locked state only from solid body tides. As shown in Figure 1 using CPL theory we find
 680 that Venus could have been tidally locked within $684 Myr$ if it started with a prograde
 681 rotation period of 3 Earth days and zero obliquity. Unfortunately the EqTide model we
 682 utilize does not support retrograde spin states, but we expect the differences to be mi-
 683 nor. We will continue to explore these issues in a future work using the simulator vplanet
 684 (R. Barnes et al., 2019) once this functionality is added. Figure 1 gives further exam-
 685 ples for CPL and Constant Time Lag (CTL) theory results using EqTide. For input pa-
 686 rameters we assume that the tidal dissipation factor $Q=12$ and Love number of degree
 687 2 $k_2=0.3$. These are the same numbers used for the modern Earth in R. Barnes (2017).
 688 Recent work by Henning and Hurford (2014) demonstrates that our choice for Q may
 689 not be unreasonable for Venus. Henning and Hurford (2014) give estimates of Q for Earth-

690 like planets (see their Fig 15, top-center-row plot) with orbital periods from 0 to 200 days.
 691 Venus' 224 d period is slightly outside the range they explore (but can be anticipated
 692 from the trend visible in their figure). Our assumption of $Q=12$ is not far off the Henning
 693 and Hurford (2014) 'Warm Earth 2' estimate in their Fig 15. As an aside, Q and k_2 are
 694 poorly constrained for present-day Venus. We have even fewer constraints on these val-
 695 ues for an ancient Venus, but perhaps those values would be more Earth-like than present
 696 day Venus. For example, present day Venus' time lag may not be the same as Earth's
 697 because of higher internal temperatures (MacDonald, 1962; Henning & Hurford, 2014).
 698 Historically Goldreich and Soter (1966) estimated that $Q < 17$ for Venus, Lago and Cazenave
 699 (1979) had values up to $Q \sim 40$ while Leconte (2018) estimate $Q \sim 100$.

700 More recently Venus' tidal love number was estimated by Konopliv and Yoder (1996)
 701 using Magellan and Pioneer Venus Orbiter data to be $k_2 = 0.295 \pm 0.066$ implying the core
 702 is liquid (Yoder, 1997). Work by Zhang (1992); Xia and Xiao (2002) have estimated $k_2 = 0.18$
 703 ~ 0.26 . A smaller value ($k_2 = 0.17$) would imply a solidified iron core which is not con-
 704 sistent with Konopliv and Yoder (1996). Modeling work by Dumoulin et al. (2017) are
 705 consistent with the work of Konopliv and Yoder (1996) as well as our own modeling choices
 706 (discussed above) of $Q=12$ and Love number of degree 2 $k_2 = 0.299$ (see Table 3 in Dumoulin
 707 et al. (2017)). Regardless, if one uses higher values of Q and/or lower values of k_2 for
 708 ancient Venus it is sufficient to say that equilibrium tide theory predicts that the CPL
 709 and CTL for Venus estimates for tidal locking will be longer than those presented in Fig-
 710 ure 1. The values in Figure 1 then represent *lower limits* to tidal locking for a given start-
 711 ing rotational period. As a caveat there is a debate in the dynamics community about
 712 the appropriateness of the CPL and CTL approaches (Efroimsky & Williams, 2009; Efroim-
 713 sky & Makarov, 2013; Touma & Wisdom, 1994; Greenberg, 2009). so these tidal lock-
 714 ing timescales should be viewed with some caution in the context of the CPL and CTL
 715 models used herein. Of course we do not take into account magnetic braking to see how
 716 the Sun's natural spin-down might affect the tidal evolution of Venus, nor do we assume
 717 that Venus' orbital characteristics would have changed over the timescale of our calcu-
 718 lations, the latter being one of the criticism when applying CPL/CTL to evolving sys-
 719 tems (Efroimsky & Williams, 2009; Efroimsky & Makarov, 2013).

720 Recent work (Green et al., 2019) has investigated the influence of a hypothetical
 721 shallow ocean on Venus (water equivalent layers of ~ 330 meters deep and 830m) using
 722 present day topography and a range of initial rotation periods. The most dissipative sce-
 723 nario predicts a slow down of 72 days per million years. The latter result may be sur-
 724 prising until one recalls that tidal dissipation in Earth's oceans is larger than that of the
 725 Earth's solid body tides (e.g., Munk & MacDonald, 1960).

726 It is likely that Venus was initially a prograde spinning body like the other 3 ter-
 727 restrial planets in our solar system. The prograde hypothesis goes back at least to (Alfvén,
 728 1964) and more recent work by Lissauer and Kary (1991); Dones and Tremaine (1993);
 729 Kary and Lissauer (1995) would also support the idea of a primordial prograde Venus,
 730 barring the effects of a late large impactor as discussed in those works and that below.
 731 From that starting point we find the following perhaps the most compelling answer to
 732 Venus' present day spin state. As shown above there are models that can drive the planet
 733 toward a tidally locked state rather quickly. Core-mantle friction damps obliquity per-
 734 turbations which drive the spin rate to sub-synchronous prograde rotations. Then at-
 735 mospheric tides would reverse the spin to a retrograde equilibrium. These atmospheric
 736 tides continue to prevent the planet from being tidally locked and that is the state the
 737 planet has been in since that time as shown in some of the work of Correia and Laskar
 738 (2001).

739 Yet it has long been speculated that Venus' current retrograde rotation state is the
 740 result of a large impactor early in its history (McCord, 1968; Singer, 1970; French & Singer,
 741 1971; Counselman, 1973; Burns, 1973; Ward & Reid, 1973; Harris, 1978; Alemi & Steven-
 742 son, 2006; Davies, 2008). The large impactor hypothesis may also explain a possibly very

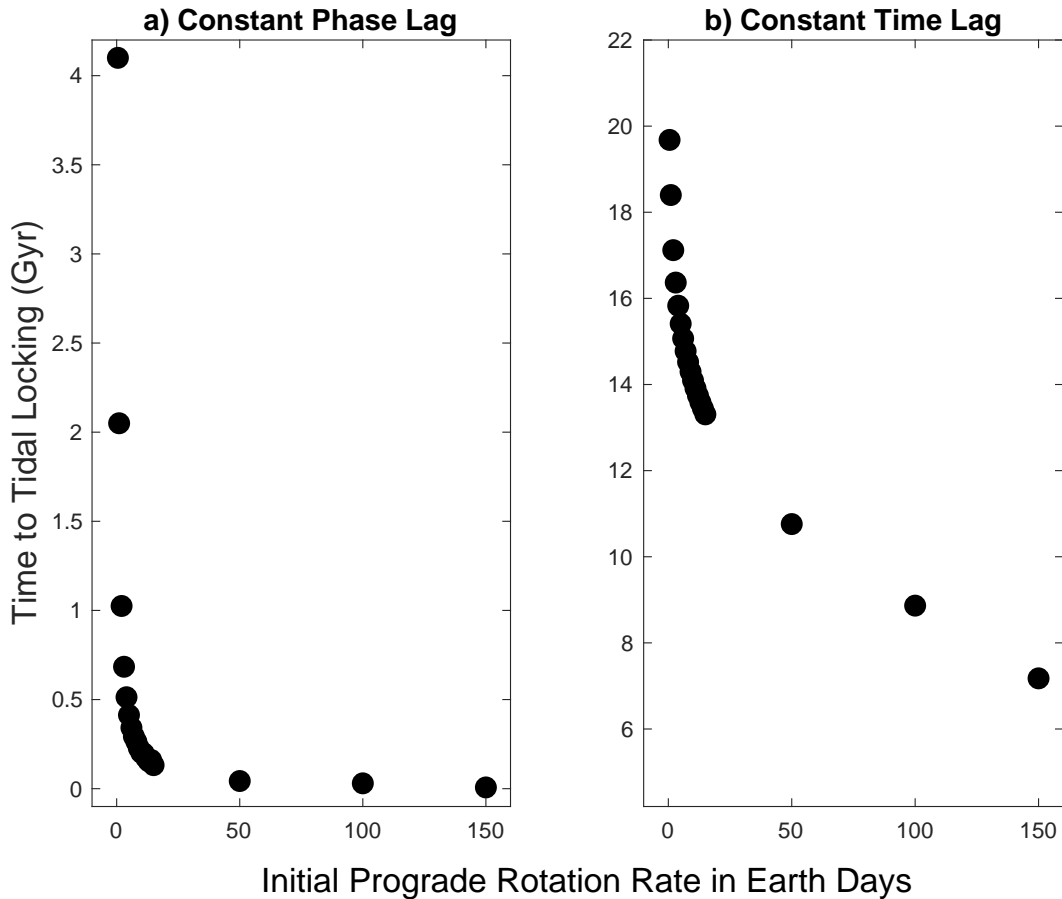


Figure 1. Venus tidal locking timescales using equilibrium tide constant phase (a) and time (b) lag theory. Note that the y-axes have different limits.

743 dry Venus interior, a lack of oxygen in the atmosphere and reconcile the ^{40}Ar results that
 744 imply it is less degassed than Earth (Davies, 2008). However, as pointed out in Ward
 745 and Reid (1973) it is possible for a small impactor (less than 1% of the Moon’s mass)
 746 to drive Venus retrograde if the planet has already spun down considerably due to tidal
 747 dissipation with the Sun (Figure 1).

748 Unfortunately, there is little hope that we will ever truly know the rotation rate
 749 of Venus through time without a way to either measure its “Geological Orrery” as on
 750 Earth (Olsen et al., 2019) and tie that to dynamical models, or find evidence that an im-
 751 pactor played a role in its rotational and hence geochemical evolution as seen in inves-
 752 tigation on Earth related to late accretion and its effect on different isotope abundances
 753 (e.g., Varas-Reus et al., 2019). In fact Brassier et al. (2016) and S. J. Mojzsis et al. (2019)
 754 prefer the hypothesis that the Earth’s Late Veneer was mainly delivered by a single Pluto-
 755 or Ceres-sized impactor. Hence if a larger object was involved in the late evolution of
 756 Venus’ spin or obliquity it may be possible to detect its geochemical fingerprints in a fu-
 757 ture in-situ mission.

758 For *most* of our climate simulations we assume, rightly or wrongly as discussed above,
 759 that Venus has had the same retrograde rotation and its present day obliquity for the
 760 past 4.2×10^9 years. We have included a few faster rotation rates that approach the ‘fast
 761 rotators’ described in Yang et al. (2014); Way et al. (2018), meaning sidereal day lengths
 762 of 16 and 64 times that of present day Earth.

6 Methods

All our simulations use ROCKE-3D (Way et al., 2017) a three-dimensional (3-D) General Circulation Model (GCM) developed at the NASA Goddard Institute for Space Studies (GISS). Radiative transfer in ROCKE-3D uses k-coefficients derived from the HITRAN 2012 line list, as well as the MT-CKD 3.0 water vapor continuum and CO₂ collisionally-induced absorption and sub-Lorentzian line shapes, as described in references cited in (Del Genio, Way, et al., 2019). HITRAN 2012 is accurate for temperatures below 350 K, as shown by (Kopparapu et al., 2013). As we will discuss below, a small number of our simulations exhibit a growing radiation imbalance with time, indicative of continually rising temperatures and a transition to a runaway greenhouse. If our objective were to determine the threshold for a runaway, we would need to use the more comprehensive HITEMP line list for these simulations, as discussed by (Kopparapu et al., 2013). Our purpose, though, is simply to identify such cases and exclude them from further analysis.

Most simulations use modern Venus' current orbital parameters, slow retrograde rotation period (-243 Earth sidereal days in length) and orbital period (224 sidereal days). In Paper I (plotted herein with ID = B) we looked at a faster rotation period (16 x modern Earth's sidereal day length) to see how the planet's early climate might have responded, and we also look at 16 & 64 day retrograde rotation periods in Venus' early history in this work. Our focus is on changing insolation, topography, land/sea mask, surface water availability and atmospheric constituents. We motivate our choices below.

All simulations discussed below are outlined in Table 1. Our focus is on the retention and stability of surface liquid water on Venus over time. Therefore we simulate four types of planets with surfaces that differ in the amount of water they contain and how that water is allowed to interact with the atmosphere. ROCKE-3D allows for 3 types of surface water: Soil moisture at and beneath the surface with no standing bodies of water; "dynamic" lakes whose depth and area vary with time and that can appear or disappear as the competition between precipitation and evaporation dictates; and deeper oceans with permanent boundaries and an effectively infinite source of water for the atmosphere. All planets with oceans are fully dynamic. For more details on the capabilities of such oceans see Way et al. (2017).

- Arid Venus: This planet has modern Venus topography, but only contains 20cm of water in the subsurface soil layers, soil consisting of 100% sand, and no surface standing water at the start of the simulation. The atmosphere is initialized with zero water vapor and an isothermal temperature profile at 300K. This initial condition is similar to that of Kodama et al. (2019); Abe et al. (2011) who attempt to limit the amount of water vapor in the atmosphere (a strong greenhouse gas) and subsequently push the inner edge of the habitable zone farther inward. However, Kodama et al. (2019); Abe et al. (2011) use modern Earth's rotation rate for all their experiments.
- 10m-Venus: Uses modern Venus topography and places a 10 meter liquid water-equivalent layer in the lowest lying topographic areas. These are treated by the model as lakes, which have no circulation. The soil is a 50/50 sand/clay mix as used in Yang et al. (2014); Way et al. (2016, 2018).
- 310m-Venus: Similar to 10m-Venus, except with a 310 meter water equivalent layer again spread in the lowest lying regions. This is the same topography used in Way et al. (2016) simulations A,B and D.
- 158m-Aqua: This is a simple aquaplanet configuration that is commonly used in the exoplanet community. It uses a fixed 158 meter deep ocean, which corresponds to the bottom of the fifth layer of the ROCKE-3D ocean model. It is a bit shallower than the mean depth of the 310m-Venus ocean, and therefore comes into equilibrium a bit faster while still having a similar heat capacity, while including hor-

815 izontal heat transport as well as wind-driven and thermohaline overturning cir-
 816 culations.

817 • 310m-Earth: Similar to 310m-Venus, but using a modern Earth-like land/sea mask
 818 with a 310m deep bathtub dynamic ocean (i.e., every ocean grid cell is of a fixed
 819 depth of 310m). We call this an Earth-like land/sea mask since it is not exactly
 820 modern Earth, but has some modest changes as shown in Way et al. (2018) Fig-
 821 ure 8.

822 The five planets above are then given four types of atmospheres and four differ-
 823 ent insulations as described below:

- 824 • Simulations 1-5: These have a 10 bar 100% CO₂ atmosphere using a solar spec-
 825 trum and insolation from 4.2Ga from the work of Claire et al. (2012). CO₂ was
 826 probably the dominant gas in Earth's early atmospheric evolution (e.g., Kasting,
 827 1993). We pick atmospheric pressures of 1 bar (see next bullet point) and 10 bar
 828 to cover the Kasting (1993) ranges (see their Fig 2). Our 10 bar results either equi-
 829 librate at a temperature beyond that at which our radiative transfer is accurate,
 830 or do not reach equilibrium and the temperatures attained at the time the exper-
 831 iments were terminated are already beyond the upper limits of our radiation ta-
 832 bles. We report the results of these experiments in Table 2 below simply as a guide
 833 for future research, but we exclude them from our analysis in Figs. 2-8. The 10
 834 bar simulations use a modern Venus rotation rate and obliquity.
- 835 • Simulations 6-10: Similar to Simulations 1-5, but these use a 1 bar 97% CO₂ and
 836 3% N₂ atmosphere at 4.2Ga.
- 837 • Simulations 11-15: As in Simulations 6-10 but with a rotation period of -16 side-
 838 real Earth days to place the planet on the edge of the fast rotator regime as de-
 839 scribed in Yang et al. (2014); Way et al. (2018). This allows us to explore the pos-
 840 sibility that the planet was rotating more quickly in its early history than today.
 841 The choice of a retrograde rotation rate was chosen to be consistent with the present
 842 day retrograde rotation, but unpublished simulations with prograde rotation rates
 843 with these values produce very similar temperatures. Note that the work of Correia
 844 and Laskar (2001, 2003) indicate that prograde rotation rates of 16 days for Venus
 845 put its spin axis (obliquity) in a possibly chaotic regime. However, other work by
 846 J. W. Barnes et al. (2016) indicate that low obliquity retrograde rotation rates gen-
 847 erally have more stable spin axes. Even if there are spin axis variations on geo-
 848 logical timescales, it is not possible for us to model those here given that ROCKE-
 849 3D simulations are limited to < 10,000 years in length.
- 850 • Simulations 16-20: As in Simulations 6-10 but with a rotation period of -64 side-
 851 real Earth days. This allow us to explore the possibility that the planet was ro-
 852 tating somewhat more quickly in its early history but still in the slowly rotating
 853 dynamical regime. Again, prograde rotation rates were also used in unpublished
 854 results and have similar global surface temperature values.
- 855 • Simulations 21-25: These simulations use an atmospheric composition and pres-
 856 sure very similar to modern Earth, namely an N₂-dominated atmosphere with 400ppmv
 857 CO₂ and 1ppmv CH₄ with a 1013mb surface pressure. They also use a solar spec-
 858 trum and insolation at 2.9Ga from Claire et al. (2012). The rotation rate is the
 859 same as modern Venus.
- 860 • Simulations 26-30: Similar to Simulations 21-25, but with a lower atmospheric sur-
 861 face pressure of 250mb. This is again in the interest of comparative climatology
 862 since the Archean atmospheric pressure proxy work of D. Som S. Catling et al.
 863 (2012); S. Som et al. (2016) suggests that Earth may have had a surface pressure
 864 similar to 250mb at this time.
- 865 • Simulations 31-35: Similar to Simulations 21-25, but now using a solar spectrum
 866 and insolation from 0.715Ga from the work of Claire et al. (2012).

- 867 • Simulations 36-40: Again, similar to Simulations 21-25, but now using a modern
868 solar spectrum and insolation.
- 869 • Simulations 41-45: Similar to Simulations 21-25, but now using a modern solar
870 spectrum, but with insolation set to 1.26 times Venus' present day insolation (2.4
871 times modern day Earth's insolation) to test the boundaries of the inner edge of
872 the habitable zone as in Way et al. (2018). Two of these simulations are also out
873 of radiation balance and trending toward a runaway greenhouse state, and thus
874 we do not analyze them further.

875 Most simulations except 158m-Aqua use a fixed ground albedo of 0.2 (thermal con-
876 ductivity = $0.26 \text{ W m}^{-1} \text{ K}^{-1}$) and 50/50 mix of sand/clay soil following the work of Yang
877 et al. (2014); Way et al. (2016, 2018). The Arid-Venus simulation uses the same albedo
878 (0.2), but utilizes a 100% sand soil, rather than the sand/clay mix in other simulations.
879 The advantage of using sand is that it more quickly loses and absorbs water. This al-
880 lows the ground hydrology to come into balance more quickly than other soil types. This
881 is because in the Arid-Venus simulations we are focused on water availability to/from
882 the atmosphere from/to the soil and hence the amount of total water vapor acting as
883 a greenhouse gas in the atmosphere.

884 7 Results and Discussion

885 Simulations 1-5, all with a 10 bar pure CO_2 atmosphere and 4.2 Ga insolation, are
886 uniformly uninhabitable regardless of the surface water reservoir and topography (see
887 Table 2 in the Supplementary Information). The driest planet (Arid Venus) does reach
888 equilibrium, but with a surface temperature of 262°C , well above the accuracy limits of
889 the radiation parameterization used by ROCKE-3D. The other four planets are also well
890 above 100°C at the point at which they were terminated and are not converging to equi-
891 librium. Given the greater water reservoirs in these simulations, they are likely to be ap-
892 proaching a runaway greenhouse state.

893 In Figure 2 we show several different possible evolutionary scenarios for Venus de-
894 rived from the other experiments in Table 1. In all such scenarios we assume that Venus
895 had surface liquid water in varying amounts at model start, as described in Section 6.
896 The colors in this figure differentiate groups of simulations with different insolation, ro-
897 tation, surface pressure, and/or atmospheric compositions, while the numbers 1-5 and
898 corresponding symbols for each color delineate the range of climates obtained for differ-
899 ent surface water reservoir and topography assumptions.

900 7.1 4.2Ga

901 Since there are major uncertainties about what Venus' initial rotation rate was (See
902 Section 5) we explored early post-magma ocean scenarios at 4.2Ga with three different
903 retrograde initial rotation periods in the left hand part of Figures 2,3,4: -16 days (gray;
904 experiments 6-10), -64 days (magenta; experiments 11-15) and -243 days (red; exper-
905 iments 16-20). Each assumes a 1 bar CO_2 -dominated atmosphere. As one would expect
906 from the studies of Yang et al. (2014); Way et al. (2018) the faster spin rate simulations
907 generally have higher temperatures because of the cloud processes discussed in those pa-
908 pers, but almost all of them reach equilibrium at a habitable global mean surface tem-
909 perature. However, the clouds also differ to some degree because of water availability.
910 Contrary to the work of Abe et al. (2011); Kodama et al. (2019) the Arid-Venus cases
911 all have higher surface temperatures than their counterparts. This is because those pre-
912 vious works used modern Earth's rotation rate, whereas the cloud processes on these slower
913 rotating worlds better regulate the climate, more so the more water that is available for
914 cloud formation. This analysis is backed up by Figure 3 where we plot the shortwave cloud
915 radiative forcing (SWCRF). The Arid-Venus simulations have the smallest (in magni-

Table 1. Experiments

ID	Topography ^a	Epoch Ga	Insolation ^b S0X/W m ⁻²	P ^c bar	Spin days	N ₂ ppmv	CO ₂ ppmv	CH ₄ ppmv	Soil Type ^d
01	Arid-Venus	4.2	1.396/1913.6	10	-243	0	1000000	0	S
02	10m-Venus	"	"	"	"	"	"	"	S/C
03	310m-Venus	"	"	"	"	"	"	"	S/C
04	158m-Aqua	"	"	"	"	"	"	"	-
05	310m-Earth	"	"	"	"	"	"	"	S/C
06	Arid-Venus	4.2	1.396/1913.6	1	-16	43000	970000	0	S
07	10m-Venus	"	"	"	"	"	"	"	S/C
08	310m-Venus	"	"	"	"	"	"	"	S/C
09	158m-Aqua	"	"	"	"	"	"	"	-
10	310m-Earth	"	"	"	"	"	"	"	S/C
11	Arid-Venus	4.2	1.396/1913.6	1	-64	43000	970000	0	S
12	10m-Venus	"	"	"	"	"	"	"	S/C
13	310m-Venus	"	"	"	"	"	"	"	S/C
14	158m-Aqua	"	"	"	"	"	"	"	-
15	310m-Earth	"	"	"	"	"	"	"	S/C
16	Arid-Venus	4.2	1.396/1913.6	1	-243	43000	970000	0	S
17	10m-Venus	"	"	"	"	"	"	"	S/C
18	310m-Venus	"	"	"	"	"	"	"	S/C
19	158m-Aqua	"	"	"	"	"	"	"	-
20	310m-Earth	"	"	"	"	"	"	"	S/C
21	Arid-Venus	2.9	1.47/2001.0	1	-243	1012599	400	1	S
22	10m-Venus	"	"	"	"	"	"	"	S/C
23	310m-Venus	"	"	"	"	"	"	"	S/C
24	158m-Aqua	"	"	"	"	"	"	"	-
25	310m-Earth	"	"	"	"	"	"	"	S/C
26	Arid-Venus	2.9	1.47/2001.0	0.25	-243	1012599	400	1	S
27	10m-Venus	"	"	"	"	"	"	"	S/C
28	310m-Venus	"	"	"	"	"	"	"	S/C
29	158m-Aqua	"	"	"	"	"	"	"	-
30	310m-Earth	"	"	"	"	"	"	"	S/C
31	Arid-Venus	0.715	1.71/2358.9	1	-243	1012599	400	1	S
32	10m-Venus	"	"	"	"	"	"	"	S/C
33	310m-Venus	"	"	"	"	"	"	"	S/C
34	158m-Aqua	"	"	"	"	"	"	"	-
35	310m-Earth	"	"	"	"	"	"	"	S/C
36	Arid-Venus	0.0	1.9/2601.0	1	-243	1012599	400	1	S
37	10m-Venus	"	"	"	"	"	"	"	S/C
38	310m-Venus	"	"	"	"	"	"	"	S/C
39	158m-Aqua	"	"	"	"	"	"	"	-
40	310m-Earth	"	"	"	"	"	"	"	S/C
41	Arid-Venus	Future	2.4/3266.0	1	-243	1012599	400	1	S
42	10m-Venus	"	"	"	"	"	"	"	S/C
43	310m-Venus	"	"	"	"	"	"	"	S/C
44	158m-Aqua	"	"	"	"	"	"	"	-
45	310m-Earth	"	"	"	"	"	"	"	S/C
D ^d	310m-Venus	2.9	1.47/2001.0	1	-16	1012599	400	1	S/C

^aTopography: Arid-Venus=Only Ground Water, no surficial reservoirs, 20cm water in soil, with modern Venus topography; 10m-Venus=10m Water Equivalent Layer (WEL) spread in lowest elevations as lakes with modern Venus topography; 310m-Venus=310m deep ocean with modern Venus topography; 310m-Earth=Modern Earth-like topography with 310m deep ocean; 158m-Aqua=158m deep aquaplanet.

^bInsolation: S0X = multiple of amount that Earth receives today in insolation (S0=1361 W m⁻¹).

^c Pressure in bar. ^d S=100% Sand, S/C=50/50% Sand/Clay, - = Not Applicable, 100% ocean.

^d Simulation D from Paper 1 (Way et al., 2016). Most similar to ID 8.

916 tude) values, because a drier planet has less reflective clouds with less condensed water.
917 In Figure 4 we show the percentage of high level clouds (PCLDH), the dominant of the
918 three cloud types (high, medium, low) in Table 2. Here the distinction between the Arid-
919 Venus simulations and the others is not consistent across the different rotation periods,
920 suggesting that in some cases middle and/or low level clouds make important contribu-
921 tions to SWCRF.

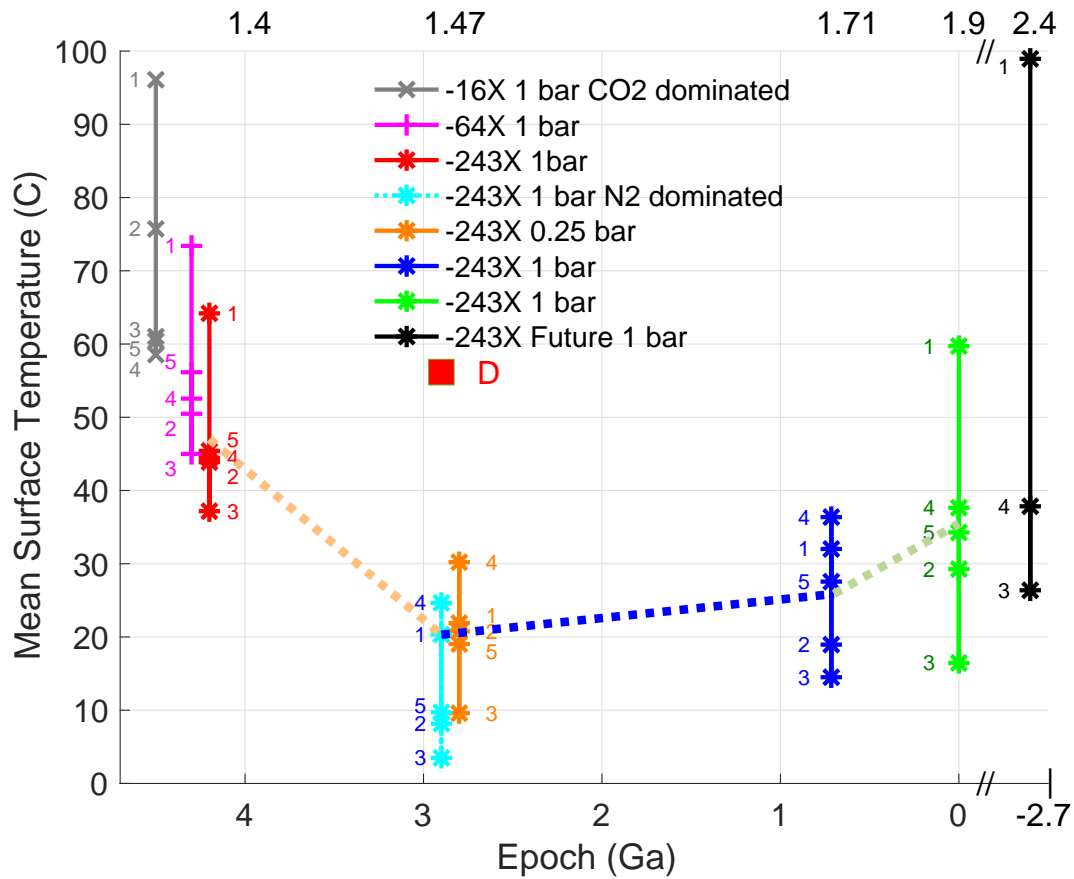


Figure 2. Atmospheric evolutionary scenarios for Venus. Top x-axis is insolation relative to modern Earth ($1.4 = 1.4 \times 1361 \text{ W m}^{-2}$). Note that the gray, magenta and red data in the left-most section of this plot are all for 4.2 Ga simulations. They are separated purely for visual effect and do not reflect differences in epoch or insolation. The same is true for the cyan, and orange data, all of which correspond to 2.9 Ga. The numbers oriented vertically along each set of simulations correspond to the different water reservoir/topography types: 1=Arid-Venus, 2=10m-Venus, 3=310m-Venus 4=158m-Aqua, 5=310m-Earth. See Table 2.

922 **7.2 2.9Ga**

923 Here we plot two different sets of simulations for N_2 -dominated atmospheres: 1 bar
 924 (cyan; simulations 21-25) and 250mb (orange; simulations 26-30). In effect these portray
 925 representative possible atmospheres for an ancient Venus with liquid water that has evolved
 926 from an early CO_2 -dominated atmosphere to a more Earth-like composition via the carbonate-
 927 silicate cycle feedback that is believed to regulate CO_2 on planets with liquid water. In
 928 both cases the Aqua-158m simulations have the highest mean surface temperatures with
 929 the Arid-Venus a close second, but all 10 simulations have moderate surface tempera-
 930 tures fairly similar to modern Earth. However as for simulations 6-20 the shortwave cloud
 931 radiative forcing is again the smallest for the Arid-Venus simulations (Figure 3) while
 932 also having less high cloud in (Figure 4) than the simulations with more surface water.
 933 Unsurprisingly, the thin 250mb atmospheres (simulations 26-30) have cooler surface tem-
 934 peratures in Figure 2. Simulations 27 & 28 have lower mean surface temperatures than
 935 modern Earth. The surface temperature field for simulation 28 is plotted in Figure 5 for
 936 reference. It exhibits fairly uniformly warm oceans, a signature of slowly rotating plan-

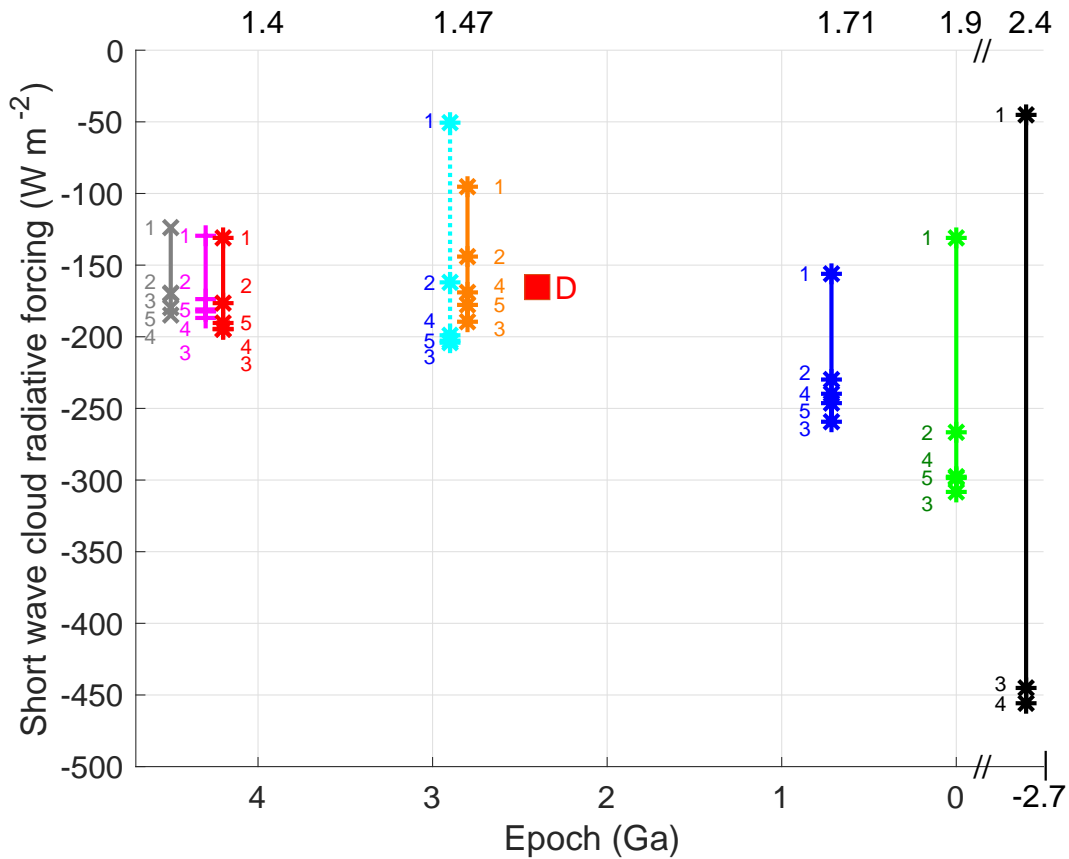


Figure 3. Shortwave cloud radiative forcing for the same evolutionary scenarios, defined as the difference between the solar radiation actually absorbed by the planet and how much would be absorbed if clouds were transparent. This is an estimate of how effective the clouds are at shielding the planet from the star’s intense radiation.

937 ets. Continental temperatures are cooler, slightly below freezing on average, due to night-
 938 time cooling that offsets daytime warming. We also plot Simulation D from Paper I. The
 939 other 3 simulations from Paper I have similar values to their corresponding simulations
 940 herein. Simulation D is an N₂-dominated atmosphere, but is otherwise similar to ID 8
 941 (a CO₂ dominated atmosphere) in Tables 1 and 2. It has a lower mean surface temper-
 942 ature than ID 8 as expected, but is significantly higher than the other simulations with
 943 larger rotation periods (ID=21–30) at 2.9Ga. It has less short wave cloud radiative forc-
 944 ing (Figure 3) and lower percentage of high level clouds (Figure 4) compared to the other
 945 310m-Venus simulations at 2.9Ga. This is expected given its faster rotation period, stronger
 946 Coriolis force, and less contiguous clouds at the substellar point as discussed in Paper
 947 I.

948 7.3 0.715Ga

949 This epoch captures a possible final habitable phase on Venus, if the thick CO₂ at-
 950 mosphere we see today was created by volcanic emissions during the global resurfacing
 951 event(s). The spread in surface temperatures between simulations remains about the same
 952 as in previous epochs, but for all surface types the temperature is warmer than at 2.9
 953 Ga because of the brighter Sun. Again, the Aqua-158m is has the highest surface tem-
 954 peratures with the Arid-Venus close behind as in the 2.9Ga epoch. The SWCRF is some-

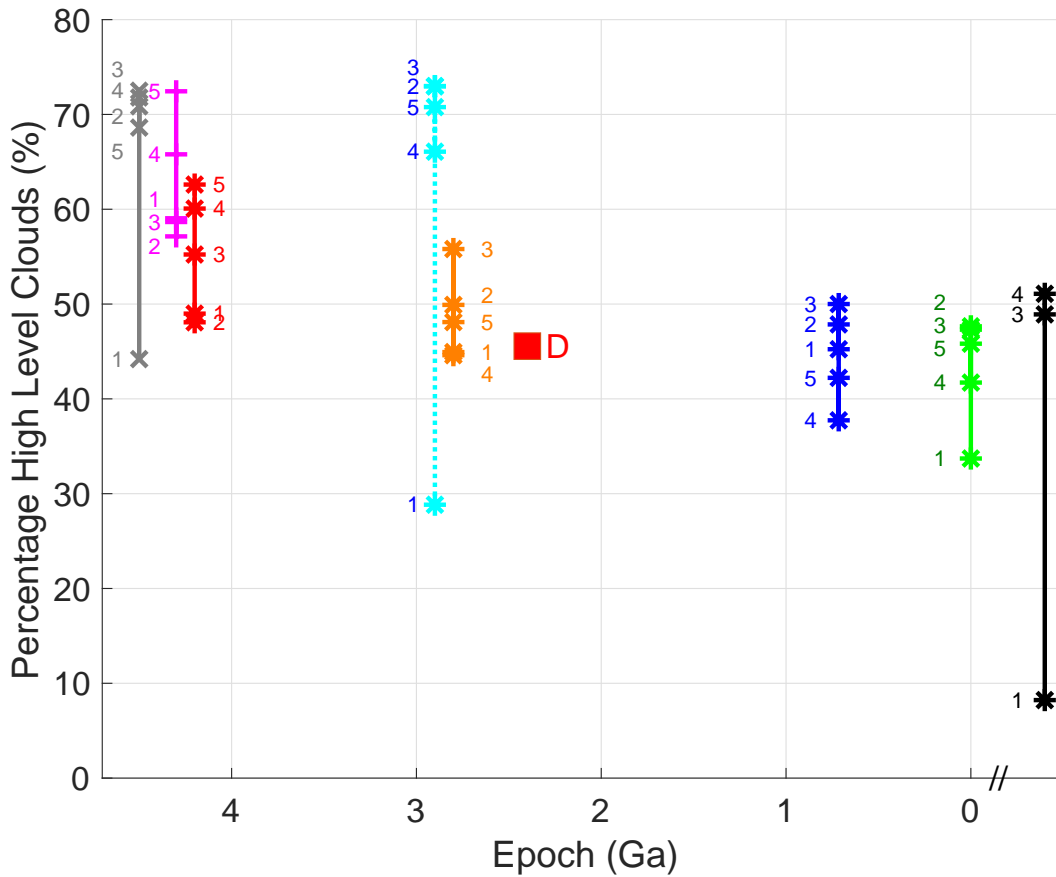


Figure 4. Percentage areal coverage of high level clouds (PCLDH). In Way et al. (2018) it was shown that PCLDH plays a key role in the cloud albedo feedback for slow rotators.

955 what larger in general, due to both the stronger insolation and slightly reduced high cloud,
 956 but again the Arid-Venus has the smaller effect. The spread in high level clouds has shrunk
 957 considerably from the 1 bar simulations at 2.9Ga.

958 7.4 Present Day

959 This suite of simulations at Venus' present day insolation are designed to demon-
 960 strate that even under today's Sun the slow rotation cloud feedback effect would have
 961 remained strong as Venus' atmospheric pressure and composition remain unchanged. This
 962 points to the idea that it was not an increase in insolation that drastically changed Venus'
 963 clement climate of earlier epochs, but rather something else, which we speculate to be
 964 multiple/simultaneous large igneous provinces. The Arid-Venus simulations again have
 965 the highest temperatures and corresponding smallest SWCRF. This is more along lines
 966 of what we saw with the simulations at 4.2Ga, with climate forcing by a stronger Sun
 967 replacing climate forcing by a thicker greenhouse gas atmosphere as the primary reason
 968 for a warm climate.

969 7.5 Future

970 Our last set of simulations at insolation values 2.4 times that of present day Earth
 971 are meant to show how long a temperate Venus-like world could have remained habit-
 972 able for a given surface type. Our 10m-Venus and 310m-Earth simulations are not in equi-

Simulation 28: 310m-Venus : 2.9Ga : -243X : 250mb : N2 dominated

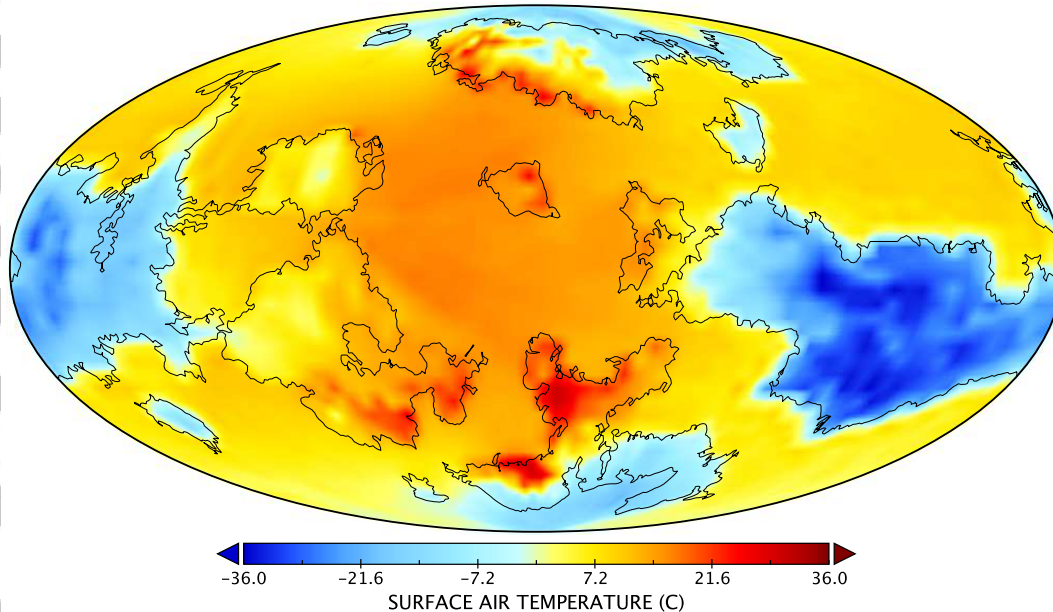


Figure 5. Simulation 28: mean surface temperature over 1/6th of a diurnal cycle. The sub-stellar point is centered over the middle of the plot.

973 librium and so are not plotted. The Arid-Venus simulation is at nearly 100° Celsius. It
 974 appears to be approaching radiative equilibrium, but the simulation crashed after 20 years
 975 so it is difficult to be certain. At this point the cloud/albedo feedback for the Arid-Venus
 976 case has decreased to Earth-like values (the SWCRF is a mere -50 W m^{-2}). This sim-
 977 ulation has the lowest value of PCLDH, which makes it hard to counter the increased
 978 insolation at this time in order to keep mean surface temperatures below the boiling point
 979 of water.

980 7.6 General Trends

981 A few relatively consistent trends are apparent from our simulations. First, the Arid-
 982 Venus simulations tend to have the highest surface temperatures, smallest values of SWCRF
 983 and lowest percentages of PCLDH. In many cases the 10m-Venus simulations are next,
 984 but not always. This may point to the fact that even with 10m of available water the
 985 cloud/albedo feedback is generally effective at shielding this slowly rotating world from
 986 the intense solar radiation at all epochs when considering what modern Earth receives.
 987 The simulations with generally higher water availability and similar percentages of land-
 988 to-sea (310m-Venus and 310m-Earth) tend to cluster together in Figures 2,3,4. The 158m-
 989 Aqua simulations seem to float in between, perhaps because of the lower surface albedo
 990 in combination with the cloud albedo feedback.

991 Even without a transition to a runaway greenhouse, prior water loss due to the on-
 992 set of a “moist greenhouse” state may have been important to Venus’ evolution. A num-
 993 ber of simulations in Table 2 contain stratospheric water concentrations (Q_{top}) greater
 994 than $3 \times 10^{-3} \text{ kg kg}^{-1}$, the traditional Kasting et al. (1993) limit for onset of the moist
 995 greenhouse. However, recent work by Chen et al. (2019) for M-star planets has demon-
 996 strated that previous work may have overestimated water loss rates. Hence we should
 997 exercise more caution in using $3 \times 10^{-3} \text{ kg kg}^{-1}$ as a hard value for the moist greenhouse
 998 until similar models are applied to G-star planets. Column Q_{surf} is a check on the amount

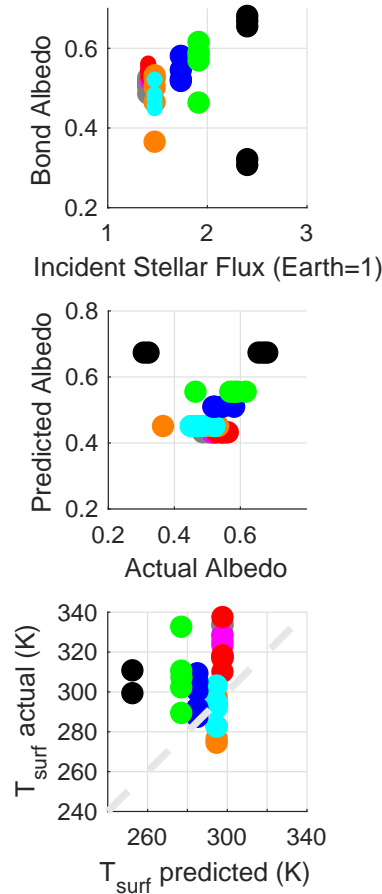


Figure 6. Upper panel: Bond albedo vs. incident solar flux for the planets in Figure 2. Middle panel: Predicted versus actual Bond albedos for the same set of planets after Del Genio, Kiang, et al. (2019). Lower panel: Predicted vs. observed surface temperatures after Del Genio, Kiang, et al. (2019).

999 of water vapor at the surface of the model. ROCKE-3D runs with a fixed molecular mass
 1000 at model start and ignores the spatially/temporally variable mass of water in calculat-
 1001 ing pressure gradients, so it is important to keep track of whether water becomes a non-
 1002 negligible fraction of atmospheric mass (e.g., 20% of the total or more) as the dynam-
 1003 ics in the model will begin to be outside an acceptable range. Only in simulation 45 does
 1004 this value go over the 20% limit.

1005 Figure 6 shows calculations for what exoplanet astronomers might find for a popu-
 1006 lation of “exo-Venuses,” some of them habitable and some not, in future observations.
 1007 We use an ensemble of ROCKE-3D simulations of a variety of rocky planet types from
 1008 which predictors for Bond albedo and surface temperature have been derived using in-
 1009 solation and star temperature as inputs (Del Genio, Kiang, et al., 2019, hereafter DG19).
 1010 Figure 6 applies the predictor to our Venus evolutionary scenarios to determine the pre-
 1011 dictability of albedo and surface temperature. In general Bond albedo increases with in-
 1012 solation in the Venus simulations (upper panel), the exceptions being 4 of the Arid-Venus
 1013 cases with limited surface water, fewer clouds and thus lower albedos than our other sim-
 1014 ulations. The DG19 predictor (middle panel) works well for all but these 4 cases, since
 1015 it predicts a high cloud-controlled albedo for the wetter planets that have more and/or
 1016 thicker clouds than others. For surface temperature (bottom panel), the predictor tends

1017 to underestimate the actual temperatures by roughly 20° or less in most cases, but by
 1018 up to 50° for the hotter, drier, marginally habitable Arid-Venus cases.

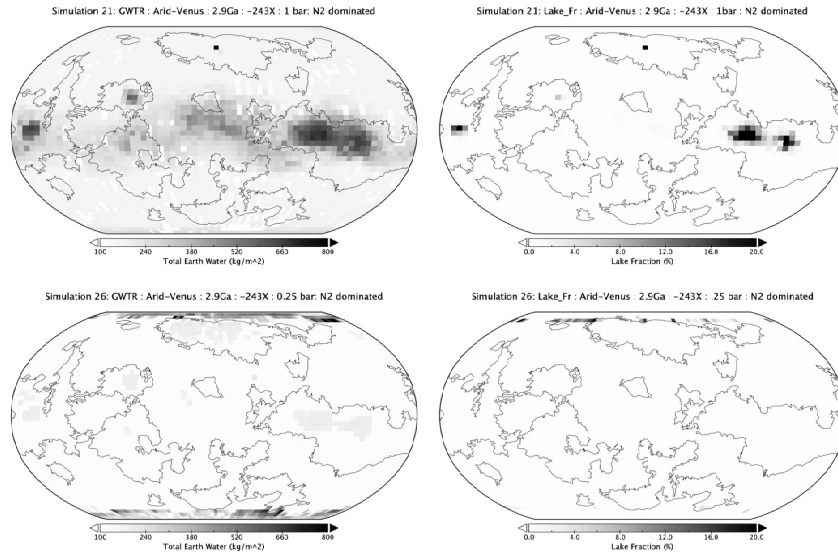


Figure 7. GWTR left panels, LakeFR right panels for Top to Bottom: Simulation 21 and 26 (Arid-Venus).

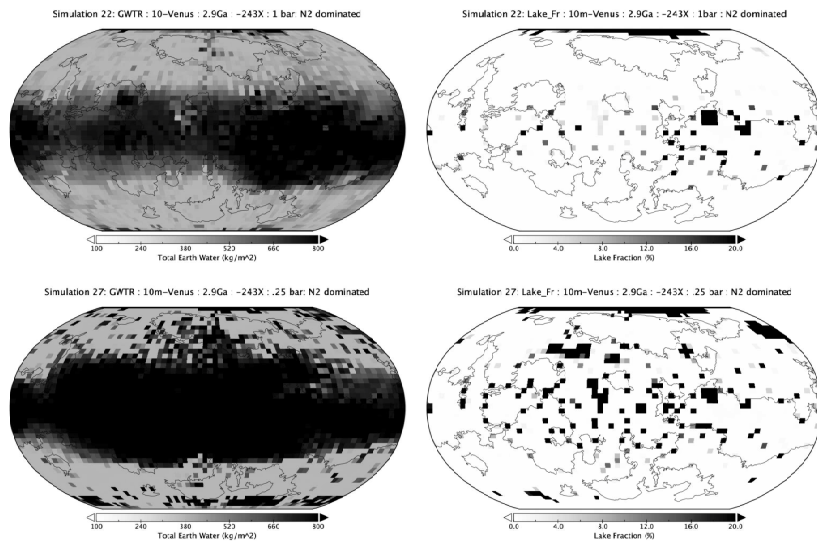


Figure 8. GWTR left panels, LakeFR right panels for Top to Bottom: Simulation 22 and 27 (10m-Venus).

1019 Figure 7 shows (left panels) the vertically integrated soil moisture and (right panels)
 1020 (right panels) lake fraction for two of our Arid-Venus simulations at 2.9 Ga with Earth-like atmos-
 1021 spheres: Experiment 21 (1 bar, top panels) and Experiment 26 (250 mb, bottom pan-
 1022 els). The Arid-Venus cases are of particular interest because they are initialized with a
 1023 spatially uniform subsurface soil water reservoir and no standing water bodies, and they
 1024 then equilibrate to a heterogeneous distribution of surface and subsurface water depend-
 1025 ing on the climate and circulation and thus the local precipitation-evaporation compe-

1026 titation. The 1 bar atmosphere is typical of the behavior of most of the experiments in Ta-
 1027 ble 1: In equilibrium, subsurface water collects primarily in the equatorial region where
 1028 rising motion and precipitation is prevalent during the daytime, and especially in high-
 1029 land regions such as Aphrodite Terra. Lakes (which are not present in the initial condi-
 1030 tion) also form preferentially over the equatorial highlands. The 250 mbar simulation
 1031 is an outlier, with subsurface water and lakes arising primarily at high latitudes. This
 1032 is reminiscent of what is observed for the methane-ethane lakes on Titan, which Mitchell
 1033 (2008) is able to reproduce in a GCM when a limited subsurface methane reservoir is
 1034 assumed.

1035 Figure 8 shows the same quantities for the analogous 10-m Venus simulations (Ex-
 1036 periments 22 and 27). These differ from the Arid-Venus cases not only because the planet
 1037 contains more water, much of it in surface lakes rather than subsurface soil moisture, but
 1038 also because the lakes fill the lowlands at the start of each simulation rather than be-
 1039 ing distributed uniformly across the planet. With a larger water reservoir than that for
 1040 the Arid-Venus planets, soil moisture collects throughout the tropics in the equilibrated
 1041 climate, but still with a slight preference for the highland regions even though there is
 1042 no standing water in the highlands in the initial condition. But unlike the Arid-Venus
 1043 planets, soil moisture also collects at the poles, both for the 1 bar and the 250 mb at-
 1044 mosphere planets. Likewise, lakes in both simulations form in both the tropics and pol-
 1045 ar region, more so over the highlands for the 1 bar atmosphere but fairly uniformly dis-
 1046 tributed in longitude for the 250 mb atmosphere.

1047 8 Conclusion

1048 Whether Venus' original water survived its initial MO stage, or whether significant
 1049 water was delivered afterwards, is unknown. It is therefore worth having a theoretical
 1050 framework that considers the possibility of an early habitable Venus as a starting point
 1051 for designing future observing strategies that might shed light on Venus' past. In this
 1052 spirit, we envision the following possible climatic evolution for Venus and provide Fig-
 1053 ure 9 as a guide:

- 1054 1.) Toward the end of the accretion period (~ 4.2 Ga) Venus would have cooled rapidly
 1055 as did Earth, as shown in the work of J. W. Valley et al. (2002). This would allow sur-
 1056 face water to condense and early oceans to form. Early oceans in turn could create sig-
 1057 nificant tidal dissipation that would spin down Venus' rotation rate on a relatively short
 1058 time scale as shown in work by (Green et al., 2019) and described in Section 5 above.
 1059 Solid body dissipation may have also been effective, see Section 5. Slow rotation com-
 1060 bined with an early ocean would provide the necessary ingredients for a dayside cloud
 1061 deck to emerge, shielding the planet from high insolation and allowing at least some of
 1062 the initial surface water to survive despite the planet being well inside the conventional
 1063 inner edge of the habitable zone, as shown in the work of Yang et al. (2014); Way et al.
 1064 (2018).
- 1065 2.) The carbonate-silicate cycle in concert with interior volatile cycling would al-
 1066 low CO_2 draw-down while N_2 was outgassed, eventually reaching a balance producing
 1067 an N_2 -dominated atmosphere with trace amounts of CO_2 over gigayears with pressures
 1068 ranging from several bars to hundreds of millibars.
- 1069 3.) We propose that any stable Venusian climate period came to an end at some
 1070 period of time (e.g. 0.2Ga to 3 Ga) before the global resurfacing event. As mentioned
 1071 previously surface age estimates for Venus range from as young as ~ 180 Ma (Bottke et
 1072 al., 2016) to 750Ma (McKinnon et al., 1997). We suggest that the ignition of multiple
 1073 large igneous provinces (LIPs) became active around that time (over a period of 10s or
 1074 100s of millions of years). This would not have been the global LIP as proposed by (López
 1075 et al., 1998) as it is not necessary and seems to have little support in the community.

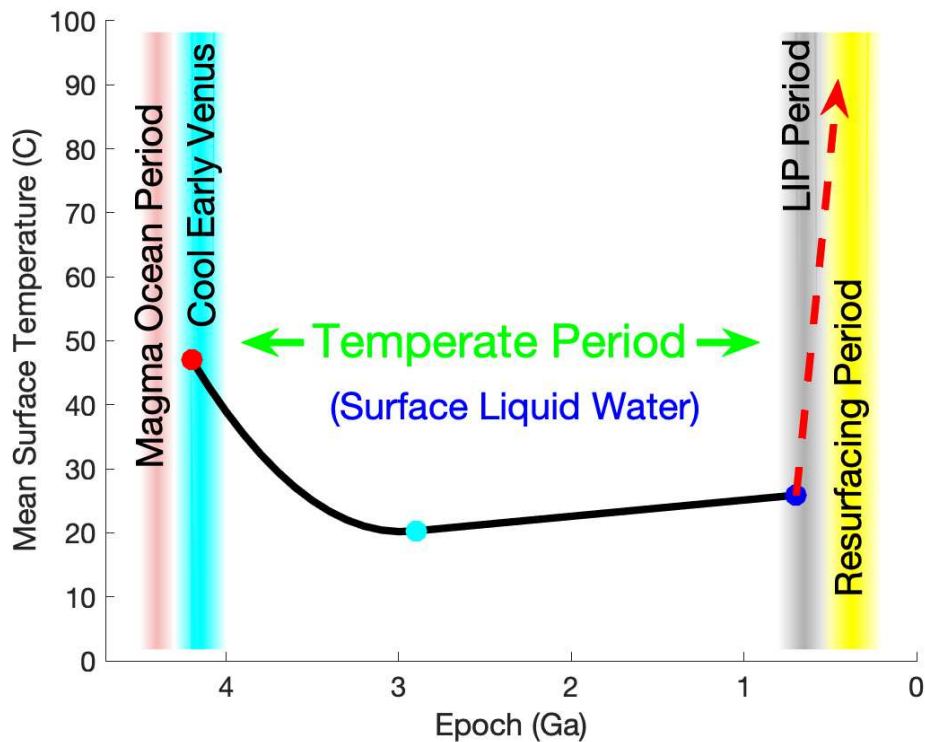


Figure 9. Graphical representation of Venus' possible climate history. The three data points represent the 1 bar atmospheres modeled at those points in time. The red dashed arrow to the right represents the transition to a moist/runaway greenhouse and eventually to Venus' present day surface temperature and atmospheric density.

1076 It is likely that the eruption of LIPs throughout Earth's history (Ernst, 2014; Ernst et
 1077 al., 2019) is a random stochastic process. This may imply that multiple large scale LIPs
 1078 have not occurred simultaneously on Earth by purely random chance, which is fortuitous
 1079 for life as we know it today. Venus may not have been as fortunate. Unfortunately, lit-
 1080 tle is known about Venus' interior structure today, much less its initial state and sub-
 1081 sequent evolution, so the question remains as to whether deterministic evolutionary pro-
 1082 cesses in the interiors of Venus-like planets will inevitably lead to catastrophic changes
 1083 and uninhabitable end states.

1084 4.) Multiple large scale LIPs would have warmed the world markedly via the re-
 1085 lease of large amounts of CO_2 via one or more mechanisms and its greenhouse effect on
 1086 the climate (e.g., Ogden & Sleep, 2012) over 10s or 100s of millions of years. It should
 1087 be noted that degassing ~ 90 bars worth of CO_2 via this mechanism would be difficult
 1088 and is an open research topic.

1089 5.) This warming could have hastened a movement toward the onset of a moist (Kasting
 1090 et al., 1993) and possibly a runaway greenhouse state in which much of the surface wa-
 1091 ter was lost via photodissociation accompanied by hydrogen escape and an oxidized
 1092 surface. Recall that present day Venus D/H measurements (if they are accurate) imply
 1093 .6 to 16% of Earth's present day surficial water stores (Donahue et al., 1997). Hence the
 1094 timescale of total loss would be at least an order of magnitude faster than that described
 1095 in Kasting et al. (1993), although there is still uncertainty about the stratospheric con-
 1096 ditions under which a moist greenhouse occurs (e.g., Chen et al., 2019) and when a run-

1097 away greenhouse is achieved (e.g., Kasting & Ackerman, 1986; Goldblatt et al., 2013; Ramirez,
 1098 Kopparapu, Lindner, & Kasting, 2014). Work by Grinspoon (1993) on the D/H ratio
 1099 may also support this hypothesis as they state: “Thus $(D/H)_{obs}$ may be the isotopic sig-
 1100 nature of a catastrophic resurfacing in the past 0.5–1 Gyr.”

1101 6.) A major issue with any proposed evolutionary scenario with long lived surface
 1102 water reservoirs then arises: what happened to the large quantities of oxygen expected
 1103 to be left over from a Venusian ocean? This is the oxygen that would have remained in
 1104 the atmosphere after the photo-dissociation of H_2O and the loss of the hydrogen (Watson
 1105 et al., 1981), but not oxygen, via atmospheric escape. In fact large quantities of abiot-
 1106 ically produced O_2 (100s to 1000s of bar) left over in such a scenario has been proposed
 1107 as an observational signature in exoplanetary atmospheres on planets that have lost their
 1108 oceans (Luger & Barnes, 2015). For our Arid-Venus scenario it may be possible to lose
 1109 much of the oxygen via a combination of atmospheric escape (Persson et al., 2018) and
 1110 absorption by a surface like that of present day Venus (e.g., M. Gilmore et al., 2017). How-
 1111 ever, recently submitted work by Persson et al. (2020) demonstrates that 0.3 m of a global
 1112 equivalent layer of water could have been lost via atmospheric escape alone in the past
 1113 ~ 4 Ga. Hence the 0.2 m global equivalent layer of water in our Arid-Venus scenario fits
 1114 within this framework. Yet work by Persson et al. (2018); Masunaga et al. (2019); Pers-
 1115 son et al. (2020) and their estimates for O^+ escape rates shows that atmospheric escape
 1116 alone is not sufficient to remove larger reservoirs of oxygen left over from the oceans in
 1117 our other non-Arid-Venus simulations. It should be noted that the escape estimates from
 1118 Persson et al. (2020) are distinctly lower than previous work by Chassefière (1996, 1997)
 1119 where they are mainly concerned with escape during a more active younger sun. The work
 1120 of Abe et al. (2011) also gives much higher escape rates estimating that an entire Earth
 1121 Ocean’s volume could be lost in 600Myr to 14Myr depending on how active the sun is.
 1122 Some caveats go with this work in that the Venus atmosphere simulated for our temper-
 1123 ate period (more akin to an Archean Earth atmosphere than modern day Earth or Venus)
 1124 is N_2 dominated with 400ppmv CO_2 and 1ppmv CH_4 and is very different from that of
 1125 modern Venus with thermospheric and exospheric temperatures likely to be distinct and
 1126 possibly affecting escape rates (Airapetian et al., 2017).

1127 We propose that the large-scale resurfacing evident on Venus today, which took place
 1128 over 100s of millions of years, is a possible effective answer. It must be noted that this
 1129 would be separate from the earlier (in geological time) LIP scenario above. Since the Mag-
 1130 ellan mission it has been known that $\sim 80\%$ of the surface of Venus is relatively young,
 1131 with estimates ranging from 300-700Myr old as mentioned in Section 4. These newly
 1132 exposed basalts would be the ideal sink for large quantities of oxygen (possibly 100s of
 1133 bars) over 100s of millions of years. According to (Lécuyer et al., 2000) Venus would need
 1134 to oxydize a rock layer ~ 50 km deep to absorb an Earth Ocean’s worth of oxygen, and
 1135 they propose a mechanism for doing so while citing the earlier work of Pieters et al. (1986).
 1136 Note that none of the oceans proposed herein are close to an Earth’s ocean in volume,
 1137 hence the number could be much smaller. For example, work by Grinspoon (1993) and
 1138 J. W. Head et al. (1992) note that the volume of magma necessary to cover all pre-existing
 1139 craters would need to be a global layer ~ 10 km deep, and that would be sufficient for
 1140 the volume of oceans proposed in our models. Some fraction of the oxygen may actu-
 1141 ally be deep within Venus’ lithosphere and possibly even within its mantle. This may
 1142 be consistent with coronae-related subduction hypotheses (Sandwell & Schubert, 1992;
 1143 Davaille et al., 2017) and other ideas about downwelling-associated highlands (e.g. J. W. Head
 1144 et al., 1992, see Fig 2), where thicker crustal regions may exhibit orogenesis (e.g., Head
 1145 et al., 1990) and sinking of parts of the lithosphere into the mantle (e.g., Lenardic et al.,
 1146 1991; Bindschadler et al., 1992).

1147 7.) The loss of water would in turn change the planet from an initial subductive
 1148 or mobile plate tectonic mode to more of a stagnant lid mode (as on present day Venus
 1149 & Mars) since it is currently believed that water plays a key role in plate tectonics on

1150 Earth (e.g., Grove et al., 2012; Lécuyer, 2014). This scenario fits in very nicely with the
 1151 recent work of (Weller & Kiefer, 2019) who give a timescale of order 1Gyr for the tran-
 1152 sition from a mobile to a stagnant lid mode on Venus in their simplified model. With-
 1153 out a mechanism to efficiently cycle volatiles in a stagnant lid mode (e.g., Tosi et al., 2017;
 1154 Höning et al., 2019), outgassing would have continued without the major weathering and
 1155 subduction surface sinks that operate on Earth, hence CO₂ and N₂ would build up over
 1156 time to reach the levels we see on Venus today. Some studies have also shown that even
 1157 in a stagnant lid mode it is possible to cycle volatiles, possibly up to gigayears in time
 1158 (e.g., Foley & Smye, 2018; Godolt et al., 2019), but these mechanisms depend on the ini-
 1159 tial CO₂ budget and the retention of at least some water after cooldown.

1160 8.) This stagnant lid mode may then allow very large mantle upwelling and/or down-
 1161 welling centers that would produce some of the features we see on Venus' surface today
 1162 produced over hundreds of millions of years, as described most recently in the works of
 1163 e.g Rolf et al. (2018); Weller and Kiefer (2019).

1164 Our scenario can also fit within the Fortunian, Guineverian & Atlian periods pro-
 1165 posed in the works of Ivanov and Head (2015) and Airey et al. (2017), as it is not pos-
 1166 sible to constrain the start of the LIP period we propose with the data we have today.
 1167 Our LIP period could easily have concluded in the pre-Fortunian or Fortunian period
 1168 ~1.5Ga.

1169 One of the remaining quandaries in our hypothesis is the fact that the 92 bar at-
 1170 mosphere we see on Venus today must go back at least as far as the age of the visible
 1171 surface because there are fewer small craters (<35 km in diameter) to be seen in the Mag-
 1172 ellan data (Schaber et al., 1992, see Figure 2). Certainly smaller craters would be vis-
 1173 ible if the atmosphere had been significantly thinner in the lifetime of the observed sur-
 1174 face when atmospheric filtering of smaller impactors would have been less prevalent.

1175 For example, assuming the tesserae are the oldest stratigraphic units why are there
 1176 no small craters present if the present day atmosphere is not a primordial feature from
 1177 many eons ago? One resolution to this problem could be that the tesserae are not as old
 1178 as we think they are, and until we date these units and the basaltic flats we really won't
 1179 know. Secondly, as mentioned above, there is the possibility that there have been mul-
 1180 tiple resurfacing events and the tesserae are left-over from one of the previous events.
 1181 Neither is a terribly optimistic scenario if one is hoping that some of tesserae may be rem-
 1182 nant crust from a period of hosting surface water. Finally, a large impactor may be the
 1183 cause of the 'catastrophic' climate change we propose. This could have also played a role
 1184 in resetting the clock on the surface of Venus reconciling the lack of small craters. In this
 1185 scenario the LIP hypothesis plays a partial role in the evolution of Venus' climate. This
 1186 would be similar to what we have seen in the on-going debate regarding the K-Pg pe-
 1187 riod on Earth (e.g. Hull et al., 2020; Schoene et al., 2019; Sprain et al., 2019). Such an
 1188 impactor's imprint would have long been lost due to Venus' relatively young surface. Our
 1189 comments about impactors in Section 4 apply here as well.

1190 Clearly a great deal more modeling work and more observations are required to con-
 1191 firm or refute this hypothesis. Did Venus follow the 'canonical' path with Earth-like con-
 1192 ditions in it's early history and then experience a moist-runaway greenhouse due to in-
 1193 creasing solar insolation? Did it experience a longer period of habitability throughout
 1194 most of its history, with its demise and present hothouse state the consequence of a se-
 1195 ries of LIP events releasing CO₂ bound up in the crust as on Earth, and/or released from
 1196 the deep interior where CO₂ is more easily sequestered (e.g., Kuramoto & Matsui, 1996)?
 1197 Or did it become bone dry in an extended magma ocean phase in the first 100Myr, as
 1198 described in Hamano et al. (2013) for Type II planets?

1199 We believe the only scenario we can begin to rule out with the present work is the
 1200 'canonical path' since there is no evidence that an early period of habitability would have

1201 been affected by increasing solar luminosity in the first billion years. In essence, if Venus
1202 had habitable surface conditions with surface liquid water ~ 4 Ga then the same cloud
1203 albedo effect that allows such a scenario would continue for eons. On the other hand we
1204 will not be able to distinguish between the two remaining scenarios until we return to
1205 Venus to make proper noble gas and other elemental and isotopic measurements at the
1206 surface (Baines et al., 2013) and better constrain escape processes at the top of the at-
1207 mosphere through time. The latter will also rely upon how such gases escape from present
1208 day Earth given the possibility that Venus may have had a magnetic field in previous
1209 epochs, even if it is not clear how important the magnetic field is to escape processes in
1210 general (Gunell et al., 2018). Likewise, whether the actual evolution of the one Venus
1211 we can visit is the ultimate fate of all highly irradiated rocky planets or an accident of
1212 an evolutionary path that might have proceeded differently in other circumstances (e.g.
1213 Lenardic et al., 2016) is not known. The stakes are high for answering this question, since
1214 many exoplanets have been discovered in the “Venus zone” just inside the traditional
1215 inner edge of the habitable zones of other stars (Kane et al., 2014). Efforts to simulta-
1216 neously characterize the CO₂ concentrations and climates of a number of these exoplan-
1217 ets, combined with a focused observational strategy for unveiling the history of the “ex-
1218 oplanet next door” to Earth in our own solar system (Kane et al., 2019), will be our best
1219 chance to understand whether the envelope for habitability and the emergence of life is
1220 much broader than usually assumed.

Table 2. Results

ID ^a	Runtime years	Temp ^b C	Balance ^c W m ⁻²	Q _{top} ^d kg/kg	Q _{surf} ^e %	Albedo Planetary	Albedo Surface	Clouds High	Clouds Medium	Clouds Low
01	71	262	0.01	1.36e-05	0.08	54	29	54	5	0
02	111	151	27.21	2.67e-05	3.74	54	18	33	17	33
03	67	121	21.31	4.50e-04	7.87	52	11	45	23	37
04	52	120	51.28	1.74e-04	10.01	44	7	36	17	51
05	56	123	36.27	4.28e-04	9.16	49	10	46	24	44
06/01	100	96	0.06	1.73e-03	1.65	51	29	44	2	0
07/02	100	76	0.41	9.34e-03	7.70	53	16	71	32	21
08/03	53	61	5.92	6.49e-03	7.50	50	10	72	32	47
09/04	100	59	-0.48	6.23e-03	6.95	52	7	73	35	62
10/05	60	60	1.90	5.46e-03	7.19	52	9	69	37	53
11/01	200	73	0.05	9.03e-04	1.70	52	29	59	4	1
12/02	500	50	0.04	5.94e-04	2.07	55	16	57	10	10
13/03	200	45	0.87	1.56e-03	3.15	54	10	59	18	43
14/04	156	53	0.11	3.47e-03	4.81	51	7	66	25	60
15/05	200	56	0.44	4.58e-03	5.65	52	9	72	28	49
16/01	300	64	0.01	8.34e-04	0.80	53	28	49	4	1
17/02	500	44	0.06	5.68e-04	1.40	56	16	48	9	8
18/03	1000	37	0.13	8.18e-04	2.07	56	10	55	15	38
19/04	500	44	-0.50	1.59e-03	3.01	54	7	60	17	56
20/05	500	45	0.23	1.88e-03	3.15	54	9	63	19	44
21/01	500	22	0.02	2.52e-05	0.36	46	28	45	6	4
22/02	3000	21	-0.04	3.08e-05	0.61	47	17	50	10	12
23/03	1000	10	-0.05	2.53e-06	0.70	52	10	56	20	48
24/04	500	30	-0.12	3.51e-05	2.32	45	6	45	14	37
25/05	1000	19	0.05	5.87e-06	1.24	48	9	48	22	39
26/01	300	20	0.11	5.94e-03	0.53	37	29	29	5	9
27/02	1000	8	-0.14	9.98e-03	1.26	50	20	73	17	27
28/03	1000	3	-0.02	6.15e-03	1.94	52	13	73	8	13
29/04	1000	25	0.34	2.69e-02	6.39	47	6	66	1	0
30/05	500	10	0.08	8.58e-03	2.76	50	10	71	7	10
31/01	300	32	-0.03	1.12e-04	0.63	52	28	45	7	4
32/02	297	19	0.02	8.32e-05	0.87	2	1	48	20	26
33/03	1000	15	-0.03	1.18e-05	0.95	58	10	50	28	47
34/04	500	36	-0.03	6.41e-04	3.32	52	7	38	17	42
35/05	500	28	0.27	1.11e-04	2.04	55	9	42	26	40
36/01	50	60	-0.09	3.05e-03	0.58	46	29	34	2	1
37/02	2000	29	-0.01	4.33e-04	1.17	59	17	48	16	23
38/03	1000	16	0.12	3.52e-05	1.08	62	10	47	33	47
39/04	500	38	-0.07	1.37e-03	3.57	57	7	42	21	42
40/05	1000	34	0.35	1.10e-03	2.92	58	9	46	24	41
41/01	20	99	0.52	3.13e-03	0.43	32	29	8	0	0
42/02	8	288	90.21	1.00e-01	11.71	31	17	31	0	0
43/03	1000	26	0.01	3.35e-03	1.99	68	11	49	33	54
44/04	500	38	0.23	5.52e-03	3.75	67	8	51	44	54
45/05	35	83	32.64	1.05e-01	30.36	66	10	89	19	16
D	2000	56	0.19	5.31e-03	8.97	44	10	46	39	45

^aID: The colored numbers correspond to those in Figure 2.^bTemp: Surface temperature in Celsius.^cBalance: These numbers come from 50 year averages unless the Runtime (Column 2) is less than 150 years, then the average is 10 years.^dSpecific humidity in top layer of the atmosphere.^eSurface humidity as percentage of atmosphere.1221 **Appendix A Energy balance and temperature**

1222 Different simulations reach radiative balance sooner than others, while some never
1223 reach it at all. Herein we plot the energy balance (in units of W m⁻²) and surface tem-
1224 perature (in Celsius) as a function of simulation year. This should allow the reader to
1225 have a better grasp of which simulations are appropriate for a given interest.

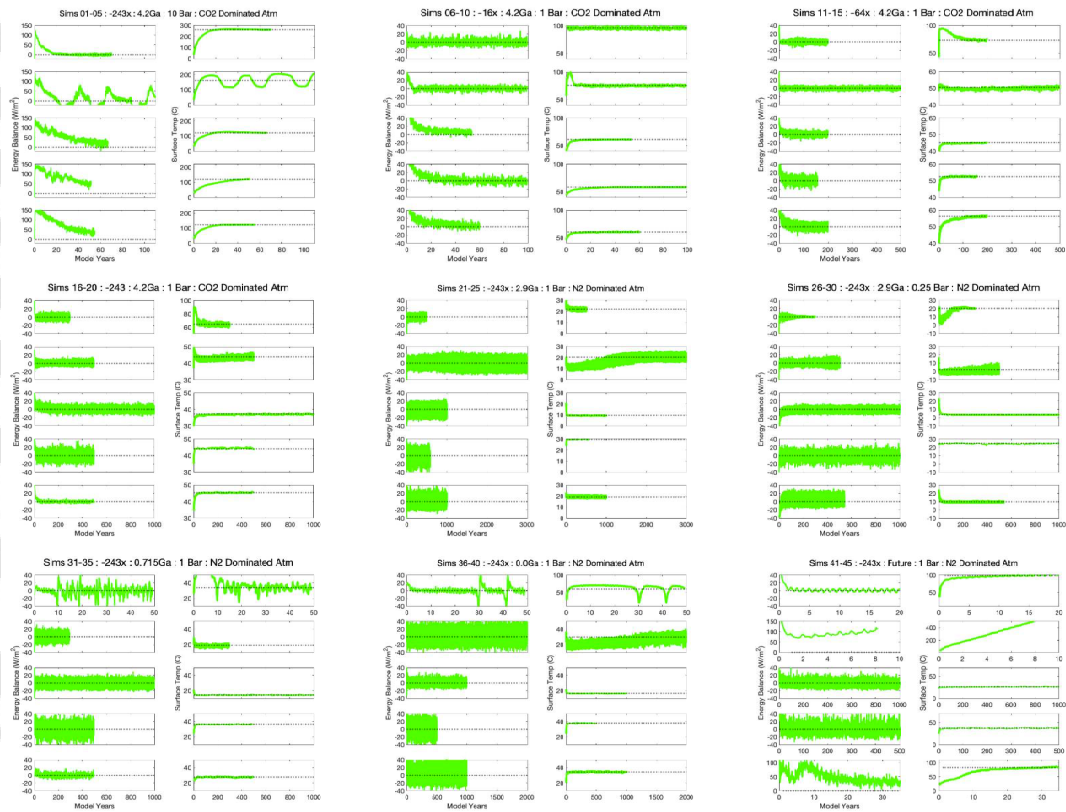


Figure A1. Energy Balance (also called Net Radiative Balance) in the left columns and surface temperature in the right columns as a function of simulation year for all simulations in this study. Note that not all limits on the x or y axes are the same.

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Acknowledgments

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1242

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