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Modeling the Transport Processes of a Pair of Solar Energetic Particle Events Observed by *Parker*
Solar Probe near Perihelion

2 LULU ZHAO,¹ MING ZHANG,¹ AND DAVID LARIO²

3 ¹*Department of Aerospace, Physics and Space Sciences, Florida Institute of Technology, Melbourne, FL, USA*

4 ²*NASA Goddard Space Flight Center, Heliophysics Science Division, Greenbelt, MD, USA*

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6 ABSTRACT

7 We present model calculations of the transport processes of solar energetic particles in
8 the corona and interplanetary medium for two events detected by Parker Solar Probe
9 near its second perihelion on 2019 April 2 and April 4. In the 2019 April 2 event,
10 the <100 keV proton differential intensity measured by the IS \odot IS EPI-Lo instrument
11 increased by more than a factor of 10 above the pre-event intensity, whereas the \sim 1
12 MeV proton differential intensity detected by the EPI-Hi instrument did not show any
13 intensity enhancement. In the 2019 April 4 event, the \sim 1 MeV proton intensity showed
14 an increase of more than a factor of 100 above the pre-event intensity, but the <100 keV
15 proton intensity enhancement was rather low and gradual. The observational fact that
16 only the 2019 April 4 event displayed a high-energy intensity enhancement indicates that
17 the associated acceleration process in the second event was more energetic than the first
18 event. However, the gradual and low enhancement of the low-energy proton intensity in
19 the second event is at odds with this scenario. In this paper, we investigate the injection

Corresponding author: Lulu Zhao
lzhao@fit.edu

and transport processes of protons in the corona and interplanetary magnetic fields with numerical model simulations. Our model calculations suggest that the gradual and low rise of the low-energy protons in the April 4 event was probably due to the different diffusion conditions in these two events and the fact that the spacecraft was moving away from the source region during these two events.

Keywords: Sun: corona – Sun: particle emission – Sun: magnetic fields – interplanetary medium

1. INTRODUCTION

On 2019 April 2 and April 4, *Parker Solar Probe* (*PSP*) detected a pair of solar energetic particle (SEP) events near its second perihelion at about ~ 0.17 AU from the Sun. These two events were small, and neither of them were detected by spacecraft near 1 AU. In the 2019 April 2 event, the <100 keV proton intensity measured by the Low-Energy Energetic Particle Instrument (EPI-Lo) of the Integrated Science Investigation of the Sun (IS \odot IS) experiment onboard *PSP* (McComas et al. 2016) increased by more than a factor of 10 above the pre-event measured intensities, whereas the ~ 1 MeV proton intensity measured by High-Energy Energetic Particle Instrument (EPI-Hi) of IS \odot IS did not show any noticeable enhancement. By contrast, in the 2019 April 4 event, the ~ 1 MeV proton intensities detected by EPI-Hi did increase by almost a factor of 100 above the pre-event intensities, while the <100 keV proton intensities from EPI-Lo showed only a low gradual increase delayed with respect to the onset of the higher-energy protons.

Figure 1 shows an overview of this time interval. The first three panels show, from top to bottom, (a) the heliocentric radial distance of *PSP*, R_{PSP} ; (b) the radial speed V_r of *PSP*; and (c) the azimuthal angular speed ω of *PSP*. The horizontal blue dashed line in Figure 1 (c) indicates the sidereal rotation angular speed of the sun. The vertical dotted line indicates the perihelion passage when R_{PSP} minimized, V_r became zero and ω maximized. Figure 1 (d) shows the intensity-time profiles as detected by the Time-of-Flight (ToF) system of EPI-Lo at three different energies (top three traces) and by the Low Energy Telescope (LET1) of EPI-Hi at about ~ 1 MeV (black trace).

46 In all these traces, 60-point moving averages of minute-cadence data are plotted. Each trace is
 47 labeled with the energy assigned by the IS \odot IS team under the assumption that particle intensities
 48 are dominated by protons. Figure 1 (e) shows 1-minute averages of the proton solar wind speed as
 49 measured by the Solar Probe Cup (SPC) of the Solar Wind Electrons Alphas and Protons (SWEAP)
 50 experiment on board *PSP* (Kasper et al. 2016). The red trace is a 60-point moving average applied
 51 to the 1-min averages. Figure 1 (f) shows 1-minute averages of the magnetic field magnitude as
 52 measured by the magnetometers of the FIELDS suite of instruments (Bale et al. 2016) on board
 53 *PSP*. The red trace is a 60-point moving average applied to the 1-min averages. Finally, Figure 1 (g)
 54 shows the magnetic field magnitude scaled to 1 AU, assuming a $\sim R^2$ dependence.

55 The two SEP events shown in Figure 1 were first presented by McComas et al. (2019) (see their
 56 Figure 4) and studied in detail by Roelof et al. (2019) and Leske et al. (2020). These authors associ-
 57 ated the origin of these two events with a single active region visible from the Ahead spacecraft of the
 58 *Solar Terrestrial Relations Observatory* (i.e. *STEREO-A*, hereafter *STA*). The Extreme Ultraviolet
 59 Imager (EUVI; Wuelser et al. 2004) of the Sun-Earth Connection Coronal and Heliospheric Investiga-
 60 tion (SECCHI; Howard et al. 2008) on board *STA* detected small brightness surges moving southward
 61 from this active region (see Figure 10 in Leske et al. 2020) that most likely were at the origin of the
 62 two SEP events (Roelof et al. 2019). The two purple arrows in Figure 1 (d) indicate the time when
 63 the most prominent surges at the origin of the two events occurred (i.e., 09:05 UT on 2019 April
 64 2 and 03:30 UT on 2019 April 4. Note that these times are the times of the *STA*/SECCHI/EUVI
 65 observation, and we have subtracted the light transit time from the sun to 1 AU to plot the arrows
 66 in Figure 1 (d)). The active region origin of these EUV surges was later numbered NOAA Active
 67 Region (AR) 12738 when it rotated into the view of Earth. With respect to *STA*, AR 12738 was
 68 located at N13E55 on April 2 and N13E38 on April 4, whereas *PSP* was $\sim 8^\circ$ and $\sim 35^\circ$ westward of
 69 *STA* on April 2 and April 4, respectively. The same active region was also the source of much larger
 70 SEP events observed later by *PSP* on April 20-21 (Wiedenbeck et al. 2020; Schwadron et al. 2020).

71 Apart from the different spectral signatures of both events (i.e. the event on April 2 observed mostly
 72 at low (<200 keV) proton energies whereas the event on April 4 being more prominent at >1 MeV

73 proton energies), the event on April 2 displayed enhanced abundances of heavy ions, whereas heavy
 74 ion (above He) counts were almost absent on April 4 (see [Roelof et al. 2019](#), and also Figure 4 in
 75 [Leske et al. \(2020\)](#)). The two SEP events were highly anisotropic as observed by *PSP*, with intensities
 76 measured in the apertures of the *IS \odot IS* telescopes looking sunward much larger than those in the
 77 opposite direction throughout the event ([Leske et al. 2020](#)). An initial ratio between intensities in
 78 the anti-sunward and sunward directions along the instantaneous magnetic field direction of $\sim 20:1$
 79 was reported in the April 2 event by [Roelof et al. \(2019\)](#) and in the April 4 event of $\sim 30:1$ by [Leske](#)
 80 [et al. \(2020\)](#).

81 Several factors might have played a role in the different properties of these two SEP events, from
 82 different properties of the parent solar eruptions to different transport conditions for the protons to
 83 propagate from their sources up to *PSP*. In this paper, we analyze the possibility that the signatures
 84 of the time intensity profiles in the April 4 event were due to the injection and transport processes
 85 in the proximity of the parent active region as well as the effects that the spacecraft’s motion had in
 86 shaping the time profiles of the SEP events.

87 2. MODEL SETUP

88 In this work, the energetic particle transport processes are calculated using the time-forward coro-
 89 nal and interplanetary transport model developed by [Zhang & Zhao \(2017\)](#). This model solves
 90 the stochastic differential equations corresponding to the focused transport equation to follow ener-
 91 getic particles’ propagation in the corona and interplanetary magnetic fields. All physical processes,
 92 including convection, adiabatic deceleration, magnetic focusing, drifts, parallel, and perpendicular
 93 diffusion are included (see [Zhang & Zhao \(2017\)](#) for details of the model). We point out that the
 94 particle acceleration processes in the source region are not explicitly modeled here. We only consider
 95 the transport processes of the particles in the solar corona and interplanetary space once they are
 96 released from the injection source region. The coronal magnetic fields between the solar surface (1
 97 R_s) and the solar wind source surface located at $2.5 R_s$ are calculated by using the potential field
 98 source surface (PFSS) model ([Altschuler & Newkirk 1969](#); [Schatten et al. 1969](#); [Schrijver & Derosa](#)
 99 [2003](#)) based on synoptic magnetograms obtained by the Global Oscillation Network Group (GONG)

of the National Solar Observatory (NSO) (<https://gong.nso.edu/data/magmap/archive.html>). For the April 2 event, we use the synoptic magnetogram at 09:04 UT on April 2, whereas for the April 4 event we use the synoptic magnetogram at 00:14 UT on April 4. A Parker nominal interplanetary magnetic field (IMF) is then assumed beyond $2.5 R_s$, with a matching boundary with PFSS at $2.5 R_s$. The solar wind speed is assumed to be zero below the source surface. The radial dependence of the solar wind speed above the source surface is calculated based on the empirical electron density formula of [Leblanc et al. \(1998\)](#) using the conservation of mass flow. The asymptotic solar wind speed at large enough radial distance is set to be 350 km s^{-1} as measured in situ by *PSP*.

Protons at energies 40 keV and 1 MeV are injected uniformly and isotropically on the open field lines at three different heliocentric radial distances of $1.05 R_s$, $1.5 R_s$, and $2.0 R_s$. The effect of the magnetic fluctuations on SEP propagation is treated as particle pitch-angle diffusion using the quasi-linear formulation ([Jokipii 1966](#); [Schlickeiser 2002](#)). We set the pitch-angle diffusion coefficient equivalent to a constant radial mean free path of 0.1 AU ($\sim 20 R_s$) for protons at energy of 250 keV around 1 AU. The radial mean free path below the radial distance of $20 R_s$ is set to be 1/4 times smaller ([Zhang & Zhao 2017](#)). And the radial mean free path of 40 keV and 1 MeV protons is determined by assuming a $v^{1/3}$ dependence as a result of a Kolmogorov turbulence spectrum ([Kolmogorov 1941](#)). The perpendicular diffusion is treated as a result of the supergranular motion on the photosphere ([Zhang & Zhao 2017](#); [Zhao & Zhang 2018](#)). In [Zhang & Zhao \(2017\)](#) and [Zhao & Zhang \(2018\)](#), the perpendicular diffusion coefficient equivalent to a 10% (κ_p) of the supergranular diffusion rate was found to be sufficient to explain the widely spread SEP events that occurred on 2010 February 7 and 2011 November 3. In this work, we adopt the same rate, i.e., 10% of the supergranular motion for the perpendicular diffusion calculation.

3. 2019 APRIL 2 EVENT

On 2019 April 2, the $<100 \text{ keV}$ proton intensity detected by EPI-Lo started to increase at $\sim 10:00$ UT and remained elevated for about one day, whereas the $\sim 1 \text{ MeV}$ proton intensity detected by EPI-Hi did not show a visible enhancement above the instrument sensitivity (Figure 1 (d)). At 11:00 UT on 2019 April 2, the *PSP*'s Heliocentric Earth Equatorial (HEEQ) coordinates, radial distance, co-

127 latitude, and longitude, were 0.19 AU, 93° , and 89° , respectively. By assuming nominal Parker spiral
 128 IMF lines, we calculate the footpoints on the $2.5 R_s$ source surface of the field lines connecting with
 129 *PSP* and with other spacecraft distributed in the heliosphere. The solar wind speed is assumed to be
 130 350 km s^{-1} for *PSP* and 400 km s^{-1} for other spacecraft. The magnetic field line footpoints of the
 131 spacecraft on different solar radius from $1 R_s$ to $2.5 R_s$ are then calculated using the PFSS coronal
 132 magnetic field model.

133 The synoptic map we used to calculate the coronal magnetic field is built from GONG magne-
 134 tograms up to 09:04 UT April 2. The coordinates of the *PSP* IMF line footpoints are summarized
 135 in the first row of Table 1. In particular, we show the co-latitude (HEEQ co-Lat) and the lon-
 136 gitude (HEEQ Lon + 60°) of the magnetic field line footpoints on the $2.5 R_s$ source surface in
 137 columns 5-6, and on the photosphere in columns 7-8. Since the coordinates of the spherical harmon-
 138 ics used to calculate the coronal magnetic field are HEEQ co-latitude and HEEQ longitude + 60°
 139 (<https://gong.nso.edu/data/magmap/QR/bqc/>), the latitudes and longitudes in Table 1 and in the
 140 following figures are HEEQ co-Lat and HEEQ Lon+ 60° .

141 The loci and footpoints of different spacecraft using nominal Parker IMF lines and PFSS magnetic
 142 field lines at 11:00 UT on 2019 April 2 are plotted in the left panels of Figure 2. Green and red areas
 143 represent the anti-sunward and sunward open field regions, respectively. The cross symbols represent
 144 the magnetic field line footpoints of the spacecraft on the $2.5 R_s$ source surface and the solid circles
 145 represent the footpoints on spheres at, from top to bottom, $1.05 R_s$ (a), $1.5 R_s$ (b), and $2 R_s$ (c),
 146 respectively. The black curve is the separatrix of the opposite magnetic field polarities at $2.5 R_s$.
 147 The location of the active region on the solar surface is indicated by the magenta diamond. *PSP*
 148 is represented by the letter P in orange color. Letters A, and E represent *STA* and the *Advanced*
 149 *Composition Explorer* (*ACE*; located at the L1 Sun-Earth Lagrangian point). The footpoints of
 150 magnetic field lines connecting to STA and ACE are shown to indicate how other spacecraft are
 151 approximately located with respect to the active region.

152 Usually, the size of the energetic particle injection sources in SEP events can be estimated from the
 153 size of the parent coronal mass ejection (CME) or the associated EUV waves (Zhang & Zhao 2017).

154 An angular radius of 45° and 80° has been assumed to be the injection size in the 2010 February
 155 7 (Zhang & Zhao 2017) and 2011 November 3 (Zhao & Zhang 2018) SEP events. However, the
 156 events observed by *PSP* on April 2 and April 4 were too weak to estimate the injection size from
 157 coronagraph observations, since both were associated with small surges observed in the EUV images
 158 from *STA*. The dashed circles in Figure 2 represent a region around the parent active region with
 159 an angular radius of 60° . An injection region with this relatively small radius would allow *PSP* to
 160 detect particles without the need to invoke cross-field diffusion. By examining *STA*/EUVI images,
 161 the origin of the event on April 2 was found to be associated with brightenings and south-westerly
 162 surges observed moving out from the active region 12738 (Roelof et al. 2019). For some of these
 163 surges, accompanying EUV waves can be seen moving outward from the active region mainly in
 164 the southwestern direction extending over several tens of degrees, which might indicate a continued
 165 extended particle injection. Therefore, in the simulation, the injection size is assumed to have an
 166 angular radius of 60° to the west of the active region and an angular radius of 50° to the east of the
 167 active region.

168 Particles are released onto open field lines and propagate within the corona. Open magnetic field
 169 lines emanate from the active region and expand in both longitude and latitude, allowing the particles
 170 to spread toward *PSP*. The projection of the injection locations on the source surface at $2.5 R_s$ are
 171 plotted as green and red dots in the left panels of Figure 3. Green and red colors represent the anti-
 172 sunward and sunward magnetic field polarities. All other symbols are the same as in Figure 2. As
 173 shown in panels (a) and (b) of Figure 3, with an injection height of $1.05 R_s$ and $1.5 R_s$, particles are
 174 injected directly on the magnetic field lines connected to *PSP*. Therefore, *PSP* is expected to detect
 175 enhancement with a prompt onset. However, if protons were injected at a higher radial distance of
 176 $2.0 R_s$ as shown in panel (c), *PSP* is not directly connected to the particle injection region and will
 177 not detect any particle enhancement, under the assumption that particles strictly follow the magnetic
 178 field lines. When including perpendicular diffusion or random walk of field lines, the particles do
 179 not just follow the average magnetic field lines depicted by the PFSS model but they can spread to
 180 broader regions (Zhang & Zhao 2017; Zhao & Zhang 2018). Furthermore, magnetic field turbulence

181 can scatter particles back and forth and cause them spending more time inside the complex solar
 182 coronal magnetic fields.

183 A total number of more than 2 million protons with an energy of 40 keV are injected uniformly and
 184 simultaneously, and their motion is followed for 100 minutes in the time-forward simulation. The
 185 left panels of Figure 4 show the relative intensity of all the escaping protons in the reference frame
 186 co-rotating with the sun 64 minutes after the initial injection. Here, we consider escaping protons
 187 whose final location is above the source surface where field lines are considered completely open.
 188 And the intensity level is plotted in log scale. The panels (a), (b), and (c) correspond to the three
 189 injection heights of $1.05 R_s$, $1.5 R_s$, and $2.0 R_s$, specified in Figure 2. The intensity level in the field
 190 line connecting to *PSP* is similar in panels (a) and (b) because of the direct and uniform injection of
 191 particles onto these field lines. The intensity shown in panel (c) is much lower than in (a) and (b).
 192 This is because all the particles reach the magnetic field footpoint of *PSP* by cross-field diffusion,
 193 which is a slow process.

194 Using the longitude-latitude distributions shown in the left panels of Figure 4 as an inner bound-
 195 ary (at 0.05 AU from the sun), we performed a time-backward stochastic calculation of particles'
 196 propagation. This simulation allows us to conveniently calculate the SEP intensity time profile at
 197 the location of *PSP* (Zhang et al. 2009). The radial and perpendicular mean free paths are set to be
 198 the same values as the parameters used in the time-forward coronal transport model, i.e. a constant
 199 radial mean free path of 0.1 AU for 250 keV protons scaled with particle speed as $\lambda_r \sim v^{1/3}$, and 10%
 200 of supergranular diffusion as perpendicular diffusion. The injection time profile is assumed to be
 201 (Reid 1964; Qin et al. 2006; Zhao et al. 2016)

$$Q = Q_0 \frac{1}{t} \exp\left(-\frac{\tau_r}{t} - \frac{t}{\tau_d}\right), \quad (1)$$

202 where τ_r and τ_d indicate the rise and decay timescales of the injection time profile, respectively. These
 203 rise and decay timescales are estimated from the time intensity profiles of the energetic particles shown
 204 in Figure 1 and can depend on the energy of the particles. Different values of τ_r and τ_d would affect
 205 the rise and decay rate of the intensity profile. Q_0 is the normalization constant.

Figure 5 shows the 40 keV proton intensity-time profiles at *PSP* assuming both τ_r and τ_d equal to 0.2 day. Green, red, and blue colors correspond to the three different injection heights. The anti-sunward (outward with a pitch angle of 0° with respect to the magnetic field) and sunward (inward with a pitch angle of 180° with respect to the magnetic field) intensity profiles are plotted in solid and dashed curves, respectively. The difference between the $1.05 R_s$ and $1.5 R_s$ injection is small as expected. The intensity level is lower for the $2 R_s$ injection (blue traces), which is consistent with the distributions of SEPs shown in the left panels of Figure 4. A relatively large outward-to-inward intensity ratio (with a ratio of more than 100) is seen at the beginning of the event and gradually decreases (to ~ 1.5), which is consistent with the observation by Roelof et al. (2019). Figure 6 plots the calculated time intensity profiles assuming different radial mean free paths. The larger the value of the mean free path, the larger the outward/inward ratio (shown by the dotted curves in Figure 6), owing to the less parallel diffusion processes undergone by the particles throughout their transport. The missing of the 1 MeV protons in the April 2 event could be due to the lack of particle acceleration to this high energy by the associated weak eruption. Another reason for the missing of high-energy particles is that the source region of the high-energy particles is far away from the footpoint of *PSP*. We will examine this possibility in the next section when discussing the transport processes in the April 4 event.

4. 2019 APRIL 4 EVENT

On 2019 April 4, the ~ 1 MeV proton intensity at *PSP* showed an enhancement by a factor of more than 100 above the pre-event intensity, indicating a relatively stronger event than on 2019 April 2. However, the enhancement of the lower energies (< 100 keV) particle intensities was gradual and low (see Figure 1 and also Figure 4 of McComas et al. (2019)). Leske et al. (2020) analyzed the EUVI 195 Å images from *STA* and suggested that the possible sources of energetic particles were southwest directed surges of plasma injection originated from the same active region as associated with the origin of the 2019 April 2 event.

Similar to the April 2 event, we calculate the magnetic field line footpoints of each spacecraft, and the results are shown in the second row of Table 1 and in the right panels of Figure 2. The longitudinal

233 separation between the active region and the footpoint of the magnetic field line connecting to *PSP*
 234 on the solar wind source surface at $2.5 R_s$ (orange crosses in Figure 2) was $\sim 81^\circ$, which is slightly
 235 larger than the $\sim 75^\circ$ estimated for the April 2 event. If we assume the same SEP injection region,
 236 the distance between the magnetic field line footpoint of *PSP* on the source surface and the edge of
 237 the SEP source region is larger in the April 4 event. However, the longitudinal separation between
 238 the magnetic field line footpoint of *PSP* on the photosphere and the active region as derived from
 239 the PFSS model of the April 4 event is just 39° , which is much smaller than the 54° obtained for
 240 the April 2 event. Whereas the *PSP* footpoint at $2.5 R_s$ was more separated from the SEP injection
 241 region in the April 4 event than in the April 2 event, the footpoints on the photosphere were closer
 242 for the April 4 than for the April 2 event. An energy-dependence of particles injection height could
 243 explain the differences between the two events. Assuming the same injection size as for the first event,
 244 because the distance between the magnetic field line footpoint of *PSP* and the particle source region
 245 changes dramatically with height, the SEP intensity-profile at *PSP* would be different for the two
 246 events, since the resulting SEP intensities depend sensitively on the particle injection heights as well
 247 as the details of the coronal magnetic field where particles initiate their propagation. Near-the-Sun
 248 particle measurements, such as those provided by *PSP*, are more sensitive to those injection and the
 249 initial transport processes than measurements at 1 AU.

250 Viewed from Earth, the active region corresponding to these two SEP events was on the backside of
 251 the sun. Therefore, the magnetogram measurements used by GONG to derive the coronal magnetic
 252 field configurations are many days old, and hence the PFSS results we use might contain significant
 253 uncertainties. To estimate how much uncertainty could be in our calculated magnetic field line
 254 footpoints using the PFSS model, we use two other times, 00:14 UT on 2019 March 22 and 00:14 UT
 255 on 2019 April 11. The first date, 2019 March 22, was the last day when the parent active region was
 256 still within the $\pm 60^\circ$ field of view from Earth before April 4. The second date, 2019 April 11, was
 257 the first day when the parent active region rotated back to the $\pm 60^\circ$ Earth's field of view. In Table 1
 258 (rows 3 and 4) we have considered the same spatial location of *PSP* as on 2019 April 4 (columns
 259 1-3) and compute field line footpoints at $2.5 R_s$ (columns 5-6) and $1 R_s$ (columns 7-8) using the

260 time-displaced PFSS results. The left (right) panels of Figure 7 show the loci and footpoints of each
 261 spacecraft using the March 22 (April 11) magnetogram in the same format as Figure 2.

262 Figure 7 shows that, in contrast to the results shown in Figure 2, the magnetic field line footpoints
 263 of *PSP* on the photosphere lie outside the injection region for all injection heights. In addition, the
 264 magnetograms obtained on March 22 (left panels in Figure 7) and April 4 (right panels in Figure 4)
 265 yield that the *PSP* footpoints on the source surface lie in the northern hemisphere (i.e. north of the
 266 neutral line represented by the black wavy line), while the footpoint is in a region of the opposite
 267 magnetic polarity if the magnetogram on April 11 is used (right panel in Figure 7). Comparing with
 268 the in-situ measurement of the magnetic fields made by FIELDS/*PSP* (Bale et al. 2016), the polarity
 269 of the magnetic field was inward as predicted when using the magnetogram on April 11. Here, we do
 270 not conclude that the magnetogram on April 11 is more accurate than those on March 22 or April 4
 271 since the *PSP*'s footpoint location is very close to the neutral sheet in the earlier two magnetograms,
 272 especially in that on April 4. Using two PFSS coronal magnetic fields constructed before and after
 273 the parent eruption of several SEP events, Zhao & Zhang (2018) examined the impact of the small
 274 scale variations of the coronal magnetic fields on the transport of energetic particles. They found
 275 that the final spatial distribution of energetic particles show only small variations as long as the
 276 effect of perpendicular particle diffusion from the random walk of field lines is faster than 10% of
 277 the supergranular diffusion rate, and therefore the arrival of SEPs at a spacecraft is not very much
 278 affected by the change of small-scale coronal magnetic fields. In the following calculation, we use the
 279 coronal magnetic fields reconstructed from the April 4 magnetogram and evaluate the results with
 280 the uncertainties of the magnetic fields in the corona.

281 Since in the 2019 April 4 event, both the high-energy (>1 MeV) and the low-energy (<100 keV)
 282 particles were detected, we examine here both 40 keV and 1 MeV protons. To account for the
 283 predominantly southwest surges suggested to be responsible for this event (Leske et al. 2020), the
 284 area of particle injection is set to be the same as in the April 2 event, i.e. 60° in angular radius to
 285 the west of the active region and 50° to the east of the active region. The parallel and perpendicular
 286 diffusion is set to be the same as in the case of the April 2 event.

287 The latitude-longitude distributions of escaping 40 keV protons are shown in the right panels of
 288 Figure 4. Since the properties of the latitude-longitude distributions of 40 keV and 1 MeV are similar,
 289 we only show here the distribution of 40 keV protons. From top to bottom, the height of particle
 290 injection is set at $1.05 R_s$, $1.5 R_s$, and $2 R_s$. As expected, the intensity at the *PSP* magnetic field
 291 line footpoint in panel (c) is much lower than these in panels (a) and (b).

292 The intensity profiles obtained from the time-backward simulations are calculated and shown in
 293 Figure 8. The inner boundary location, the diffusion parameters, τ_r and τ_d are assumed to be the
 294 same as in the April 2 event. The top panel is for 40 keV protons and the bottom panel for 1
 295 MeV protons. Similar to Figure 5, green, red, and blue colors represent the three different particle
 296 injection heights. The anti-sunward (outward) and sunward (inward) intensity profiles are plotted in
 297 solid and dashed curves, respectively. The intensity level of 1 MeV proton is determined by assuming
 298 a spectral index of -4.3 using a power law fit to the observations between ~ 1 and ~ 8 MeV in [Leske
 299 et al. \(2020\)](#). The overall shape of the calculated 1 MeV proton intensity-time profile agrees with
 300 the observation except that the decay is longer than the observation. The longer decay phase of the
 301 1 MeV profile is probably due to the larger than actual decay timescale τ_d assumed in the injection
 302 profile specified in Equation 1 (see discussion below). Despite of the small difference between the
 303 observation and calculation for the 1 MeV intensity profile, the calculated intensity profile for 40 keV
 304 is completely different from the observation. The overall observed intensity is orders of magnitude
 305 lower than the simulation results. Figure 8 shows that the intensity level is lower when the injection
 306 height is set at $2 R_s$ because particles are not directly injected onto the magnetic field line connecting
 307 to *PSP*. Therefore, one of the possible reasons for the low intensity level of the low-energy particles
 308 in April 4 event is that *PSP* was not magnetically connected to the source region. However, as
 309 shown in the top panel of Figure 8, with a perpendicular diffusion equals to 10% of the supergranular
 310 motion, the ~ 40 keV proton intensity profile detected by *PSP* still rises faster and higher than the
 311 observation. Therefore, the perpendicular diffusion rate assumed in the April 4 event is probably too
 312 large compared to the distance between the source region and the *PSP*'s magnetic footpoint.

313 To simulate the effect of a decreased perpendicular diffusion, we performed another set of calcula-
 314 tions assuming a smaller perpendicular diffusion rate, i.e., $\kappa_p = 1\%$ of the supergranular motion. The
 315 latitude-longitude distributions of 40 keV protons assuming different perpendicular diffusion rates,
 316 κ_p , are shown in Figure 9. In the top panel, the perpendicular diffusion rate is 10%, whereas in
 317 the bottom panel a lower perpendicular diffusion coefficient of 1% the supergranular diffusion rate
 318 is assumed. A smaller spread of particles assuming a 1% perpendicular diffusion rate is obtained as
 319 expected. In Figure 9, the particle injection height is set to be $1.5 R_s$ for both panels. Calculations
 320 with other values for the particle injection heights show similar results, so they are not shown here.

321 Whereas a decrease in the perpendicular diffusion rate might contribute to a smaller spread of
 322 low-energy particles and a delay in their release from the corona, it is interesting to investigate the
 323 effect produced by the spacecraft’s motion relative to the source of particles. Figure 1 (c) shows that
 324 the angular speed of the spacecraft increased from April 2 to April 4, and moved much faster than
 325 the sidereal solar angular speed of $14.7^\circ/\text{day}$ (indicated by the horizontal blue dashed line), especially
 326 on April 4. The supercorotation of *PSP* with respect to the sun started around 09:00 UT on April
 327 2 and reached its maximum by the end of April 4. On April 2, the near corotation of *PSP* with the
 328 source region facilitates the arrival of particles at the spacecraft, even when the particle injection was
 329 weak and lasted for a short time interval because *PSP* was continuously sampling the same field lines
 330 filled with low-energy particles. On April 4, *PSP* was continuously moving away from the source
 331 region. The supercorotation motion of the spacecraft makes it more difficult for low-energy particles
 332 to diffuse across the IMF to reach the magnetic field line that connects to *PSP* at each time.

333 In order to examine the effects of the decreased perpendicular diffusion and the motion of spacecraft
 334 on particle intensity measured at *PSP*, we plot in Figure 10 the different intensity profiles of the
 335 outward propagating 40 keV (top panel) and 1 MeV (bottom panel) protons for the April 4 event.
 336 τ_r and τ_d are set to be 0.2 day for both energies. Solid (dashed) curves represent the cases when the
 337 supercorotation of the *PSP* is (not) considered. The effects of the supercorotation on the intensity
 338 profiles are clearly seen. As time goes by, *PSP* moves further away from the source region due to
 339 the supercorotation. Therefore, the calculated intensities with supercorotation are lower compared to

340 those without supercorotation. The injection height is set to be $1.5 R_s$ for the red curves and $2 R_s$ for
 341 the blue and green curves. The perpendicular diffusion rate for the red and blue curves is assumed to
 342 be 10% of the supergranular motion and 1% for the green curves. The effects of the supercorotation
 343 are more pronounced when the injection height is higher and the perpendicular diffusion rate is
 344 smaller. This is because the longitudinal spread of particles is smaller and the magnetic footpoint of
 345 *PSP* on the source surface ($2.5 R_s$) is closer to the edge of the source region, as shown in Figure 4
 346 and Figure 9. Hence, when the supercorotation is considered, the magnetic footpoint connecting to
 347 *PSP* continuously moves into regions with lower intensity. In general, to which extent the particles'
 348 intensity is affected by the spacecraft's movement depends on the particles' distribution around the
 349 spacecraft's magnetic footpoint in the source region.

350 In Figure 11, we plot the calculated time intensity profiles of the outward propagating 40 keV and
 351 1 MeV protons for the two events. The 40 keV profiles are plotted in red and the 1 MeV profiles
 352 are plotted in blue. The relative motion of *PSP* with respect to the source has been taken into
 353 consideration for both events indicated by the solid colored curves. The dashed curves are when
 354 *PSP* motion is not considered. The time intensity profiles of the 37 keV and 1.09 MeV protons
 355 observed by *PSP* are plotted in black. In the April 2 event, the calculated time intensity profiles are
 356 normalized with respect to the observation by setting the calculated peak intensity of the 40 keV
 357 protons considering the *PSP* motion (represented by the solid red curve) to be the same as the peak
 358 intensity of the 37 keV protons in the observation (represented by the black curve). In the April
 359 4 event, the number of injected 40 keV protons are assumed to be the same as the April 2 event.
 360 The intensity level of 1 MeV proton is calculated by assuming a spectral index of -4.3 (Leske et al.
 361 2020). Both solid and dashed curves correspond to an injection time profile follows Equation 1 with
 362 $\tau_r = \tau_d = 0.2$ day; whereas the dotted blue curves have a $\tau_d = 0.08$ day with the consideration of *PSP*
 363 motion (only plotted for 1 MeV protons). When we assume $\tau_d = 0.2$ day, the simulated 1 MeV proton
 364 profiles show a decay longer than that observed. In contrast, if $\tau_d = 0.08$ day, the simulated 1 MeV
 365 proton profiles (dotted lines) fit the decay for the April 4 event. As references, in this figure, we also
 366 plot the calculated time intensity profiles of 1 MeV protons shown as the blue curves for the April

367 2 event. Similar to the April 4 event, the solid blue curve assumes the injection time profile follows
 368 Equation 1 with $\tau_r=\tau_d=0.2$ day; whereas the dotted blue lines assume $\tau_d=0.08$ day and consider the
 369 *PSP* motion. The solid and dotted blue the dotted curves show only a small increase barely visible
 370 above the pre-event intensity level. The parallel diffusion parameters in both the April 2 and the
 371 April 4 event are set to be the same as those in Figure 5 and Figure 8.

372 The injection height in calculating the 40 keV proton intensity in April 2 event and the 1 MeV
 373 proton intensity in April 4 event is assumed to be $1.5 R_s$. In calculating the 40 keV intensity profile
 374 of April 4 event and 1 MeV intensity profile of April 2 event, the injection height is assumed to be
 375 $2 R_s$. The choices of $1.5 R_s$ injection height and $2 R_s$ injection height for different cases is to represent
 376 the direct and indirect injection of particles onto the *PSP*'s magnetic footpoints. The perpendicular
 377 diffusion is assumed to be 10% of the supergranular motion in calculating the 40 keV proton intensity
 378 in the April 2 event and 1% of the supergranular motion elsewhere. Note that as long as particles
 379 are injected directly onto the magnetic footpoints of *PSP*, the effect of the perpendicular diffusion
 380 rate on the time-intensity profiles is negligible.

381 Although the *PSP* angular speed is smaller in time period of the April 2 event than that in the
 382 April 4 event, the reduction in particle intensity due to the supercorotation effect (i.e. the difference
 383 between the solid and dashed lines in Figure 10) seems to be larger in the April 2 event compared
 384 to that in the April 4 event. This is because the extent to which the particle intensity is reduced
 385 depends not only on the angular speed of the spacecraft moving away from the active region but also
 386 on the distribution of particles around the injection particle source region. Although the particles are
 387 initially injected on the open field regions uniformly, their spatial distributions are changed by the
 388 complex coronal magnetic fields and diffusion processes. In the April 2 event, the intensity gradient
 389 around the *PSP*'s magnetic footpoint is larger than that in the April 4 event. Therefore, as *PSP*
 390 moves away from the injection source, the field line connecting to *PSP* at each moment is less and
 391 less populated with energetic particles in the April 2 event than in the April 4 event. Figure 11 shows
 392 that if energetic particles are not injected directly onto the footpoint of *PSP* (i.e. injection height

at $\geq 2 R_s$), a weaker perpendicular diffusion rate together with the supercorotation of the spacecraft lead to a delayed and slower onset of 40 keV protons, comparable to the observations.

5. DISCUSSION AND CONCLUSION

We have studied a pair of SEP events observed by *PSP* when it passed its second perihelion using a coronal and interplanetary particle transport model. Both events were small in intensity and size and they were not observed by spacecraft at 1 AU. Persistence of magnetic-field-aligned ion streaming anisotropies away from the sun in these two events (McComas et al. 2019) indicates that protons arrived onto the field lines connected to *PSP*. The first event was weak, with only the low-energy protons showing intensity enhancements. The second event was more energetic, with the intensity of >1 MeV protons increasing by more than a factor of 100 above pre-event intensities. However, the enhancement of the <100 keV protons in the second event was relatively weak. If the particles were injected from the same source region in the two events, the distance between the magnetic field line footpoint of *PSP* and the source region changed, not only because of the change of the *PSP* location relative to the parent active region, but also because the variations in the complex coronal magnetic field configuration.

From previous studies, the calculated magnetic field line footpoints strongly depends on the details of the coronal magnetic models. If the magnetic field line of *PSP* is close to the edge of the SEP source region, which is likely to be the case for these two events, a slight variation of the footpoint location would affect the intensity profile observed by *PSP* significantly. In this work, a nominal Parker IMF model with a constant radial solar wind speed of 350 km s^{-1} as obtained from observations is used to trace the magnetic connection from the *PSP*'s location back to the solar wind source surface. From the source surface to the photosphere, we use a coronal magnetic field derived from the PFSS method with the NSO/GONG synoptic maps built from magnetogram measurements to trace the field line to the photosphere. The synoptic maps obtained from NSO/GONG magnetograms are only up-to-date when the measured region is facing Earth within approximately $\pm 60^\circ$ (Nitta & Derosa 2008; Schrijver & Title 2011; Riley et al. 2006). However, the active region responsible for producing and releasing SEPs in the two events was on the backside of the sun as viewed from Earth. Therefore,

420 the uncertainty of the magnetogram close to the active region could be considerable. We also traced
421 the magnetic footpoints of *PSP* using magnetograms at two very different times, March 22 and April
422 11, corresponding to the last and first appearances of the active region within the $\pm 60^\circ$ field of view
423 from Earth before and after the events. The longitudinal distances between the magnetic field line
424 footpoint of *PSP* on the photosphere and the active region were $\sim 71^\circ$ and $\sim 87^\circ$ in these two
425 cases, which are larger than the $\sim 39^\circ$ obtained using the magnetogram on April 4. Furthermore,
426 the coronal magnetic field may vary even if it is reconstructed from the same source of photospheric
427 magnetic field measurements using different field models. Such a situation could be particularly
428 compromised if the field lines go through regions of weak magnetic fields where the field structure is
429 sensitive to the field models and data input.

430 Despite all these uncertainties, the PFSS model could be a good approximation of the global coronal
431 magnetic field structures for these two events, when the solar activity is low and the impact of the non
432 potential transient events and active regions is small (Bale et al. 2019). If the height and size of the
433 particles' injection is fixed, different configurations of the coronal magnetic fields determine whether
434 energetic particles are released directly onto the magnetic field lines that connect to *PSP*. When
435 the energetic protons are not directly injected onto the connected magnetic field lines, the efficiency
436 of the cross-field diffusion determines the particle enhancement profiles. Moreover, since *PSP* was
437 continuously moving away from the source region during the time period of these two events, the
438 profiles of the low-energy particles will be affected more than the profiles of the high-energy profiles
439 because it takes more time for the low-energy particles to reach the observation site. The gradual
440 and low-intensity of the < 100 keV protons in the second event can then be explained by an indirect
441 injection, a decreased perpendicular diffusion strength and the supercorotation of the spacecraft with
442 respect to the parent active region.

443 The supercorotation of *PSP*, that will repeat near its perihelia throughout the *PSP* mission, rep-
444 resent a challenge to correctly interpret energetic particle data from this spacecraft. During these
445 periods, the magnetic connection between *PSP* and the particle sources continuously changes, fol-
446 lowing a pattern opposite to what we have been habituated to interpret observations from 1 AU.

447 Whereas the injection of low-energy particles from a particle source is usually more abundant than
448 that of high-energy particles, the transport conditions of these particles to reach a spacecraft that
449 moves away from the particle source might lead to situations when intensity enhancements are more
450 significant at high than at low energies. The occurrence of more intense events throughout solar cy-
451 cle 25, with the capability to fill broad extents of the inner heliosphere, will most likely obscure this
452 effect, and hence the importance to analyze and understand these first small SEP events observed
453 by *PSP*.

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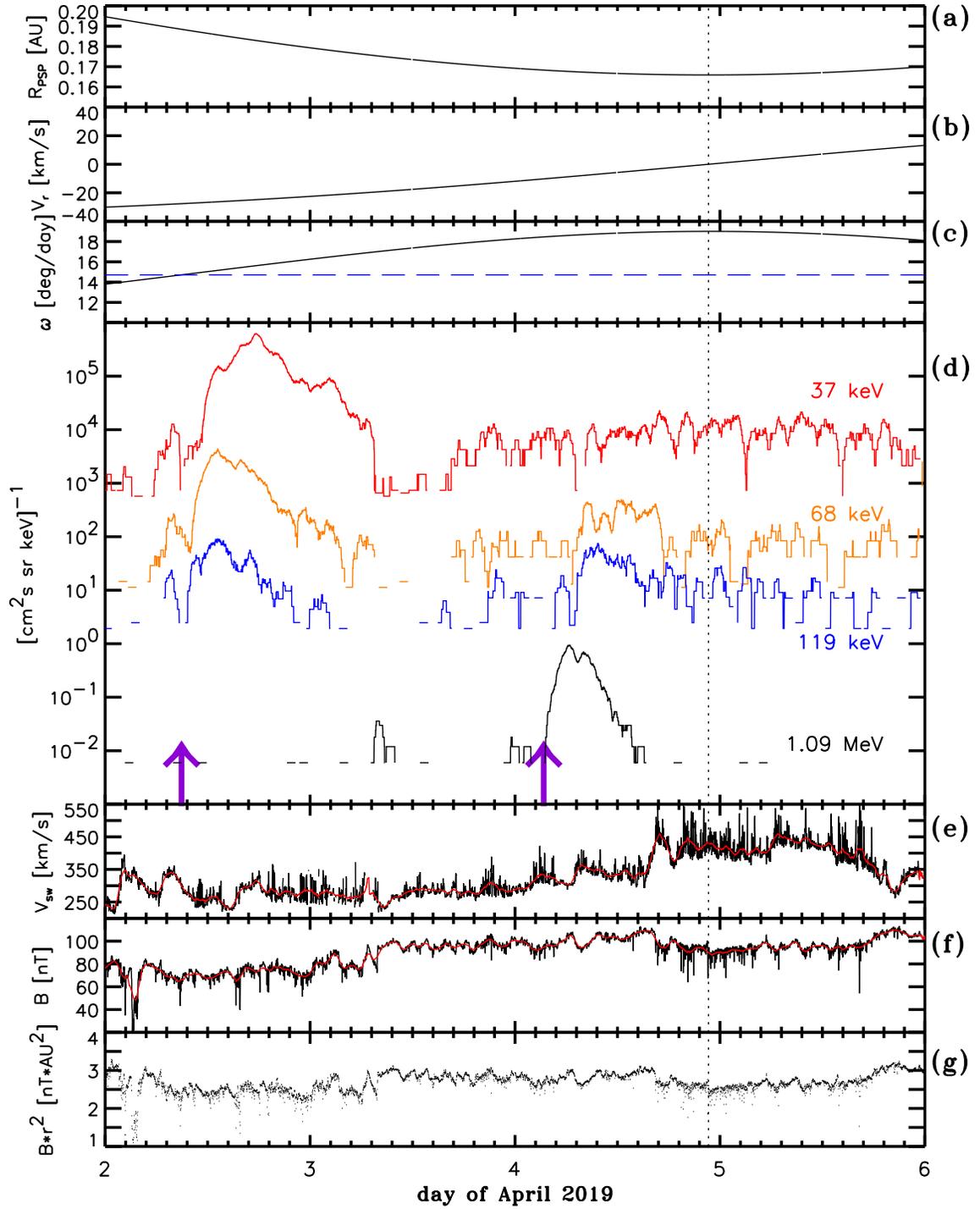


Figure 1. Overview of the *PSP* measurements from 2019 April 2 to 2019 April 6. From top to bottom, (a) *PSP* heliocentric radial distance, (b) *PSP* radial speed, and (c) *PSP* azimuthal angular speed. The dashed blue horizontal line in panel (c) plots the sidereal rotation speed of the Sun. (d) Proton intensity-time profiles at four different energies measured by EPI-Hi and EPI-Lo. (e) Solar wind speed measured by *PSP*/SWEAP. (f) Magnitude of the magnetic field as measured by *PSP*/FIELDS, and (g) magnetic field magnitude scaled to 1 AU, assuming a $\sim R^2$ dependence.

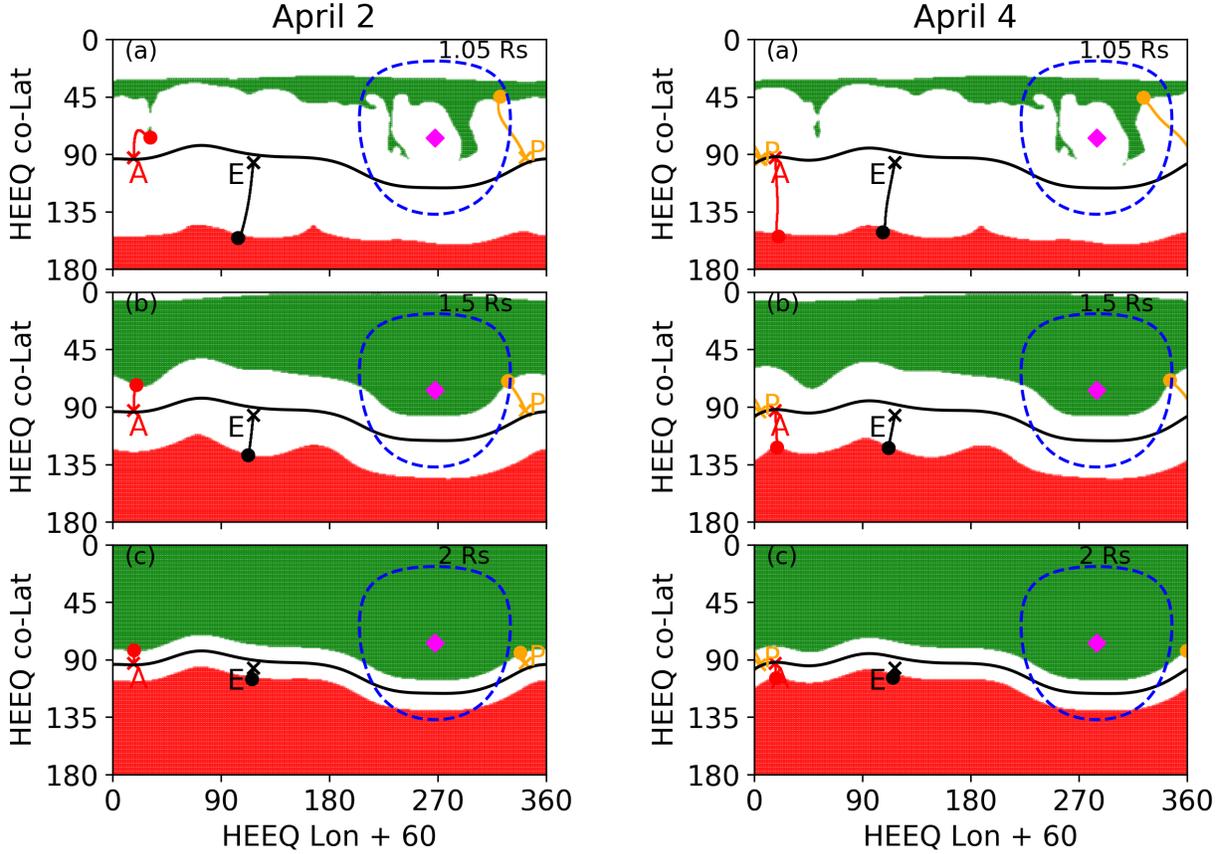


Figure 2. Loci of open fields calculated from the PFSS coronal magnetic field model using the synoptic maps at 09:04 UT 2019 April 2 (left) and 00:14 UT 2019 April 4 (right). The green dots are Anti-sunward and the red dots are sunward open fields on spheres of radius $1.05 R_s$ (panel (a)), $1.5 R_s$ (panel (b)) and $2 R_s$ (panel (c)). The magenta diamond symbol is the estimated active region location. The black curve is the separatrix between the opposite magnetic field polarities. The magnetic field line footpoints of spacecraft on the $2.5 R_s$ solar wind source surface are indicated by the cross symbols, and the solid circles are their footpoints on $1.05 R_s$ (panel (a)), $1.5 R_s$ (panel (b)), and $2 R_s$ (panel (c)). The spacecraft *STA* is plotted in red, *ACE* in black, and *PSP* in orange. The dashed circle represents the area of 60° angular radius centered at the active region.

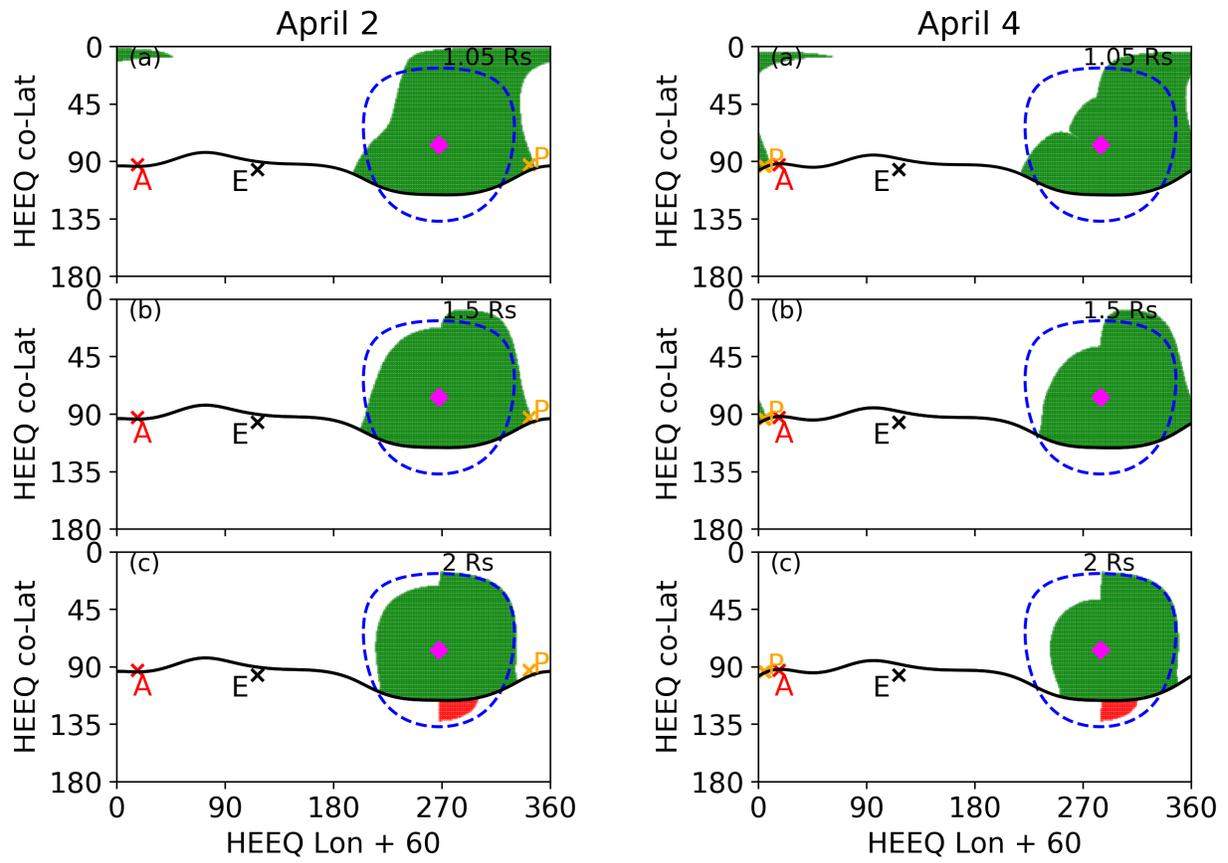


Figure 3. The same as Figure 2 except all the green (anit-sunward) and red (sunward) dots are the projection of all the injection locations onto the solar wind source surface of $2.5 R_s$.

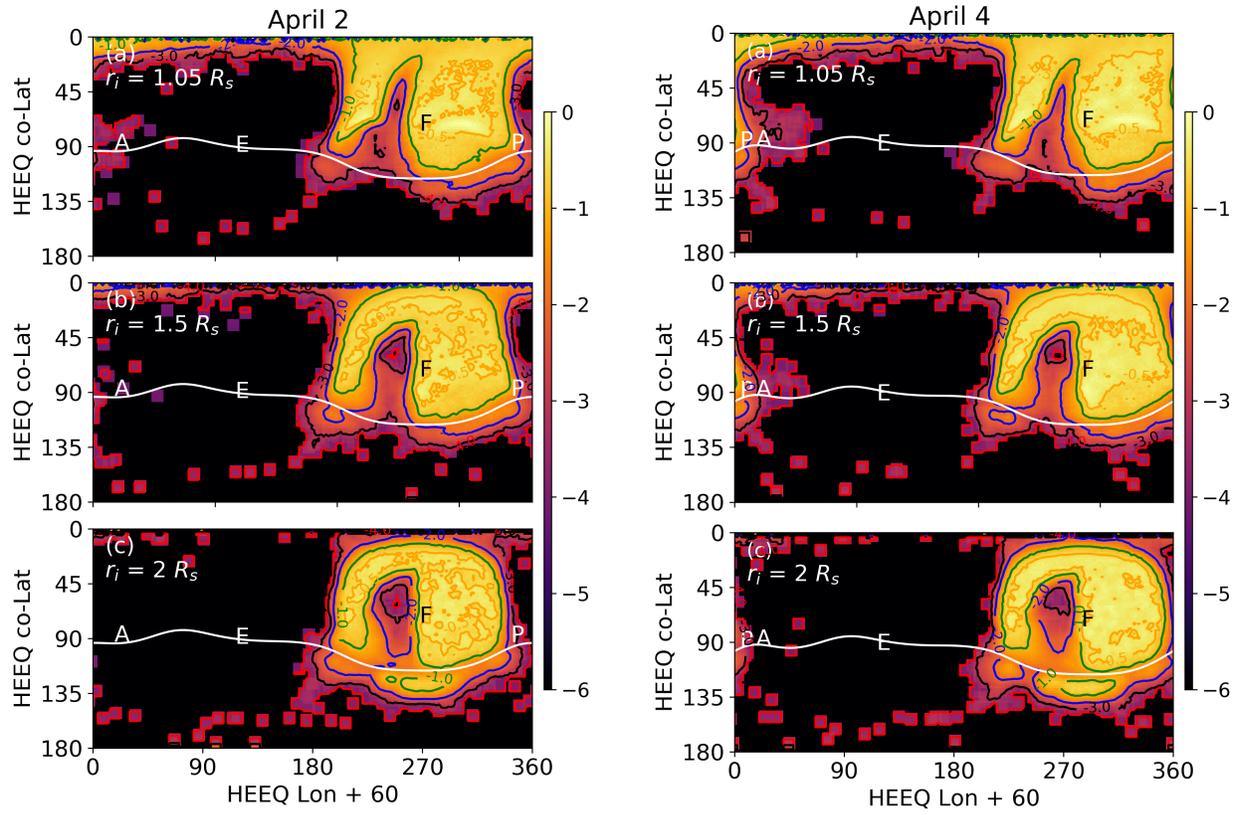


Figure 4. The magnetic field line footpoint distribution of the escaping proton locations on the solar wind source surface ($2.5 R_s$) at the time 64 minutes after injection for the April 2 event (left) and the April 4 event (right) event. The particle injection height is $1.05 R_s$, $1.5 R_s$, and $2 R_s$, for panel (a), (b), and (c) respectively. The white curve is the separatrix between the opposite magnetic field polarities. Letters A, B, E, and P mark the footpoint location of spacecraft on the source surface.

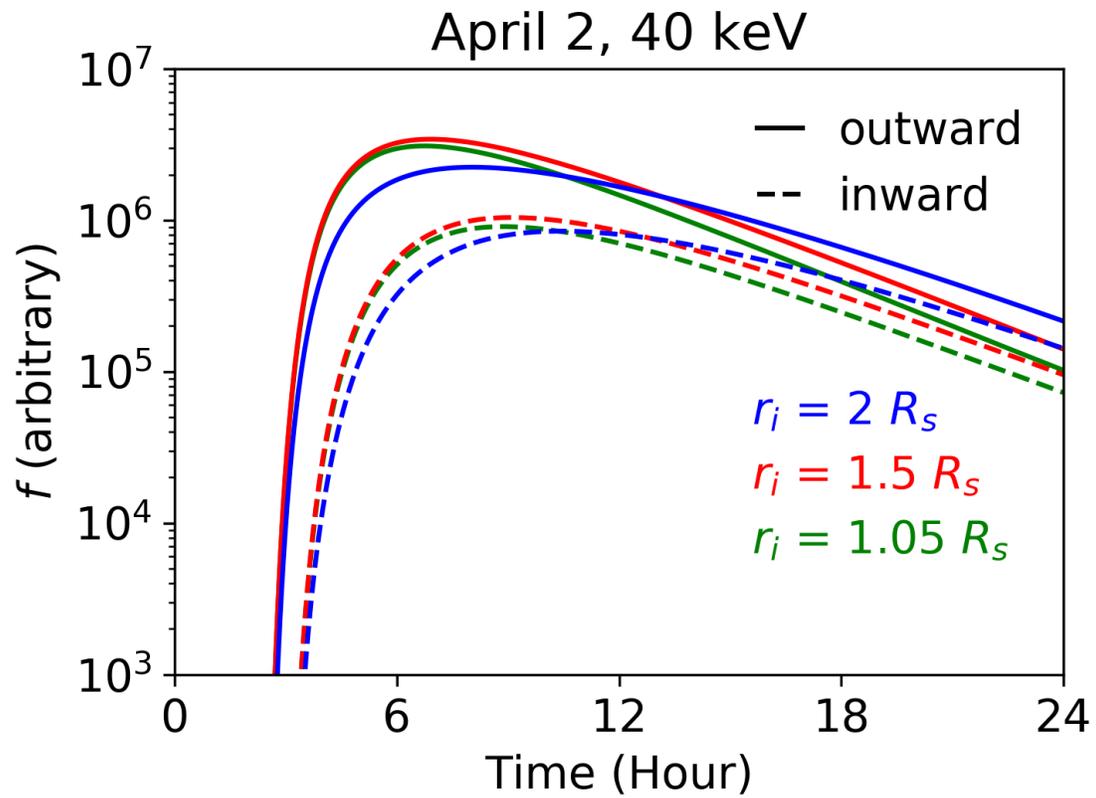


Figure 5. The calculated time intensity profiles of 40 keV protons for the April 2 event. Green, red, blue colors represent different injection heights. The intensity profiles of outward and inward moving particles are plotted in solid and dashed curves respectively.

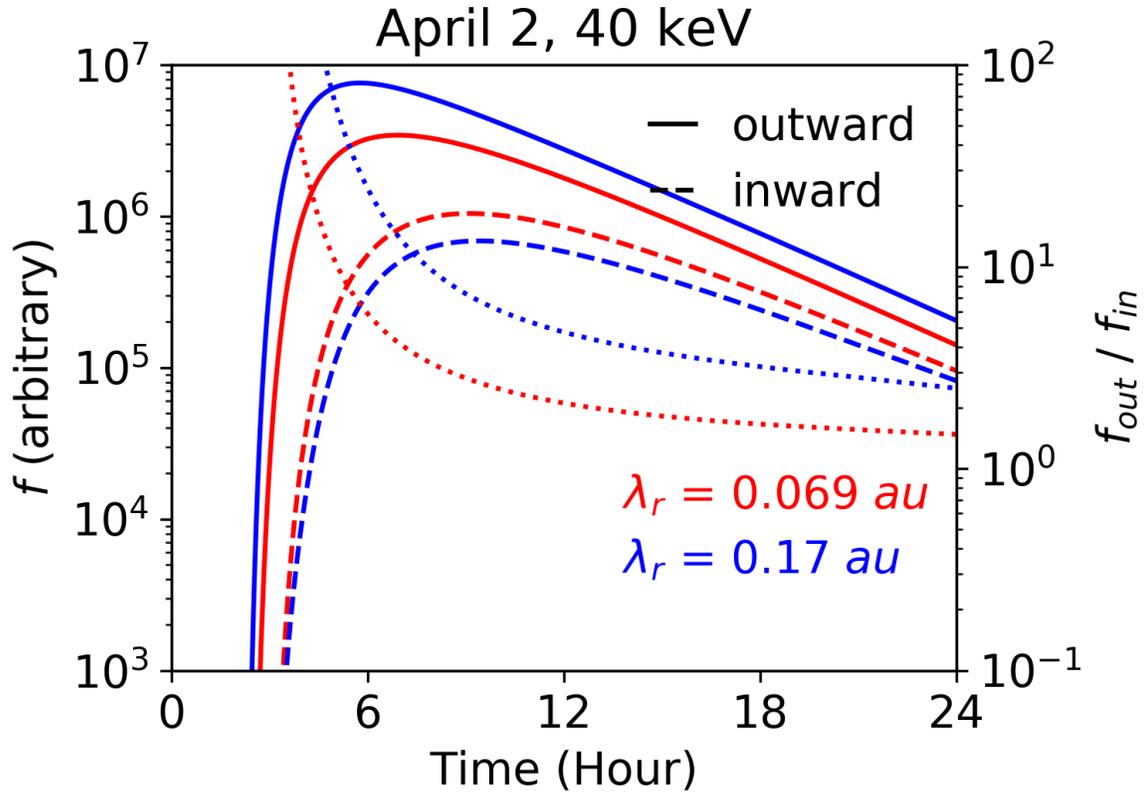


Figure 6. The calculated time intensity profiles of 40 keV protons for the April 2 event. Red, blue colors indicates different radial mean free paths. The intensity profiles of outward and inward moving particles are plotted in solid and dashed curves. The dotted curves plot the ratio of the outward and inward moving particle intensity.

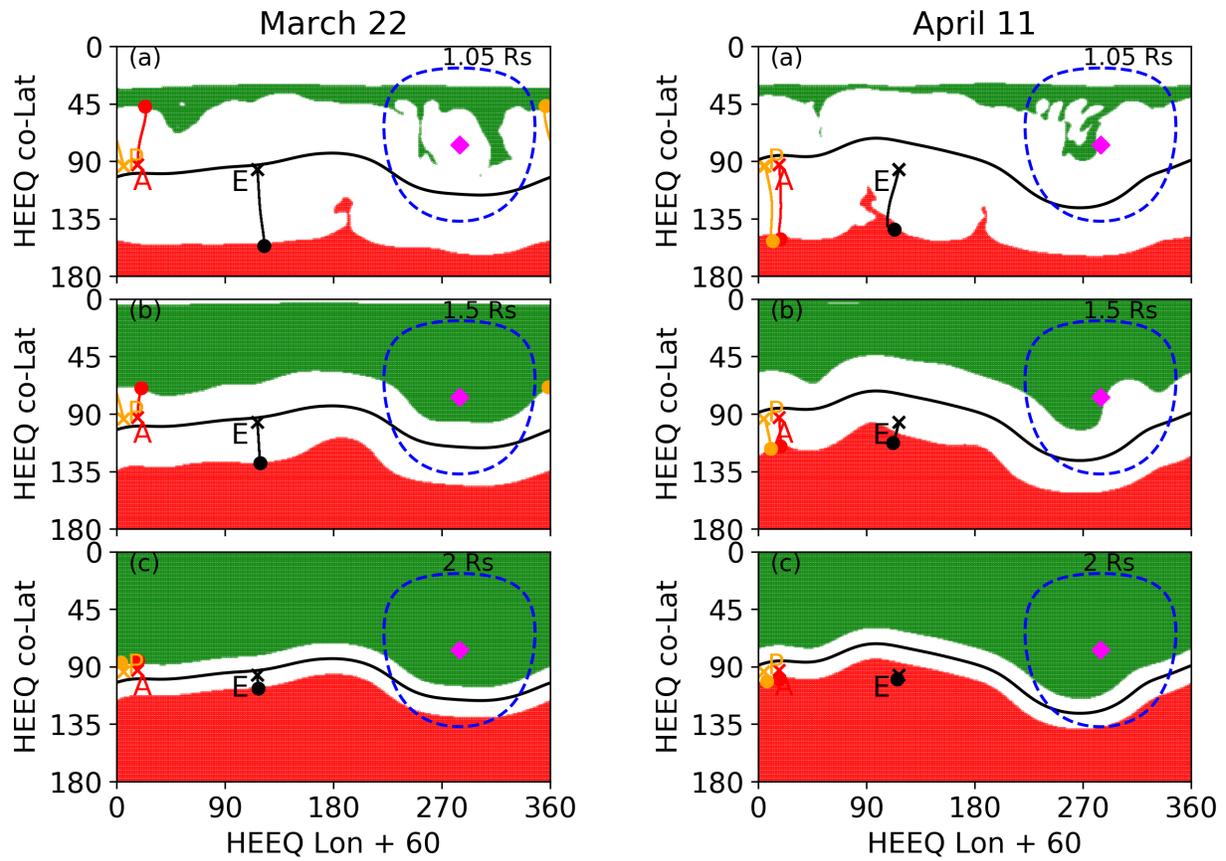


Figure 7. The same as Figure 2 except the coronal magnetic field model are taken from the magnetogram measurements at the time of 00:14 UT 2019 March 22 (left) and 00:14 UT 2019 April 11 (right).

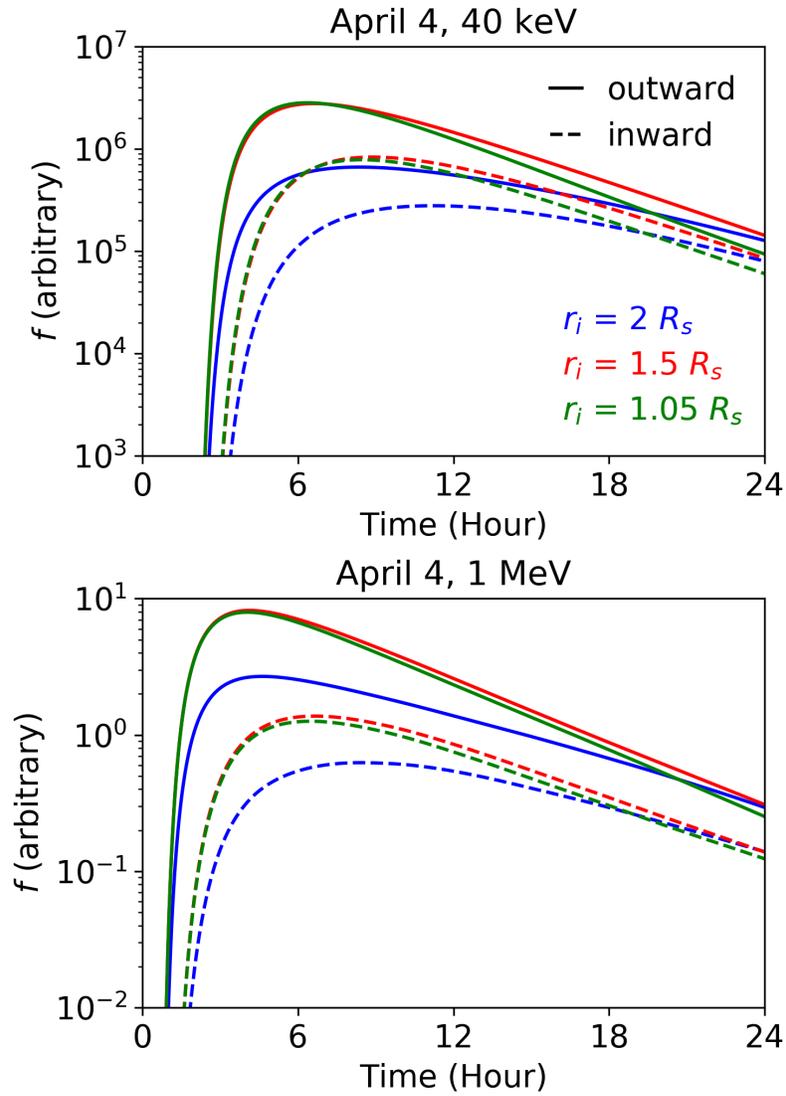


Figure 8. The calculated time intensity profiles of 40 keV (top) and 1 MeV (bottom) protons for the April 4 event. Green, red, blue colors represent different injection heights. The intensity profiles of outward and inward moving particles are plotted in solid and dashed curves.

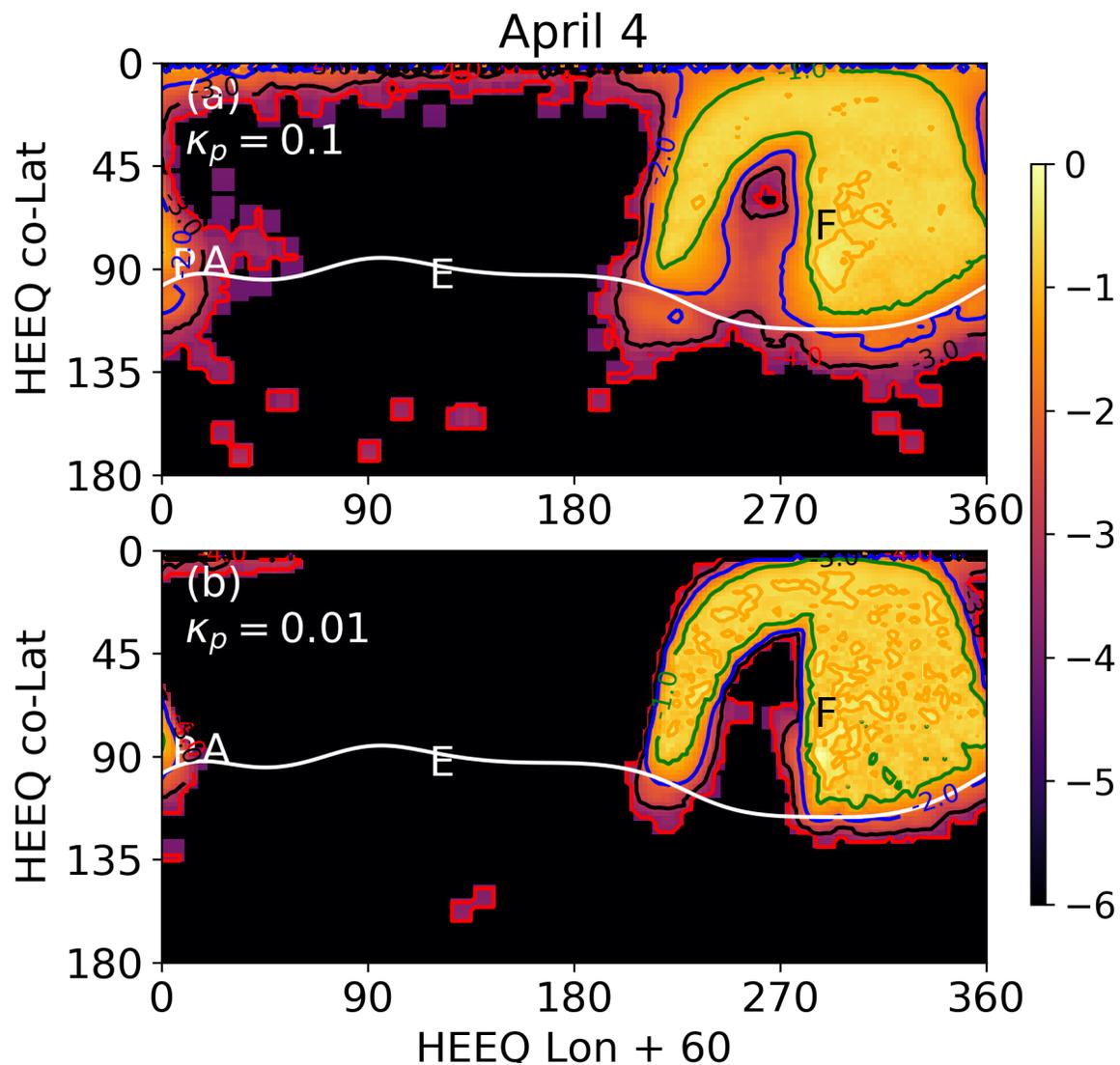


Figure 9. The magnetic field line footprint distribution of the escaping protons on the solar wind source surface ($2.5 R_s$) at the time 64 minutes after particle injection for the April 4 event with a diffusion coefficient of 10% (top) and 1% (bottom) the supergranular diffusion rate. The injection height is set to be $1.5 R_s$. All other symbols are the same as Figure 4.

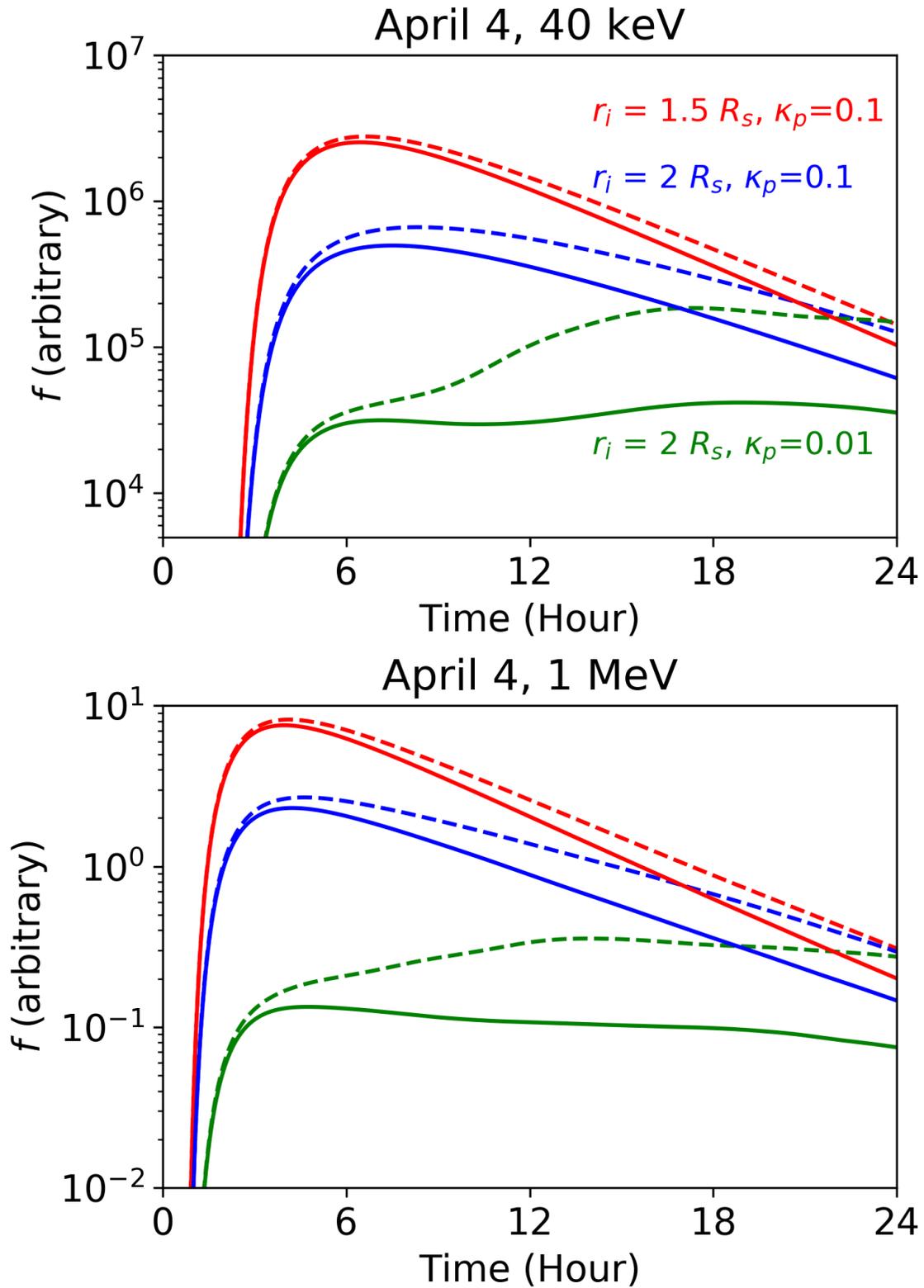


Figure 10. The calculated 40 keV (top) and 1 MeV (bottom) outward proton time-intensity profiles in the April 4 event. The solid (dashed) curves plot the time intensity profiles with (without) considering the supercorotation of *PSP*. The perpendicular diffusion is assumed to be 10% of the supergranular motion for the red and blue curves, and 1% for the green curves.

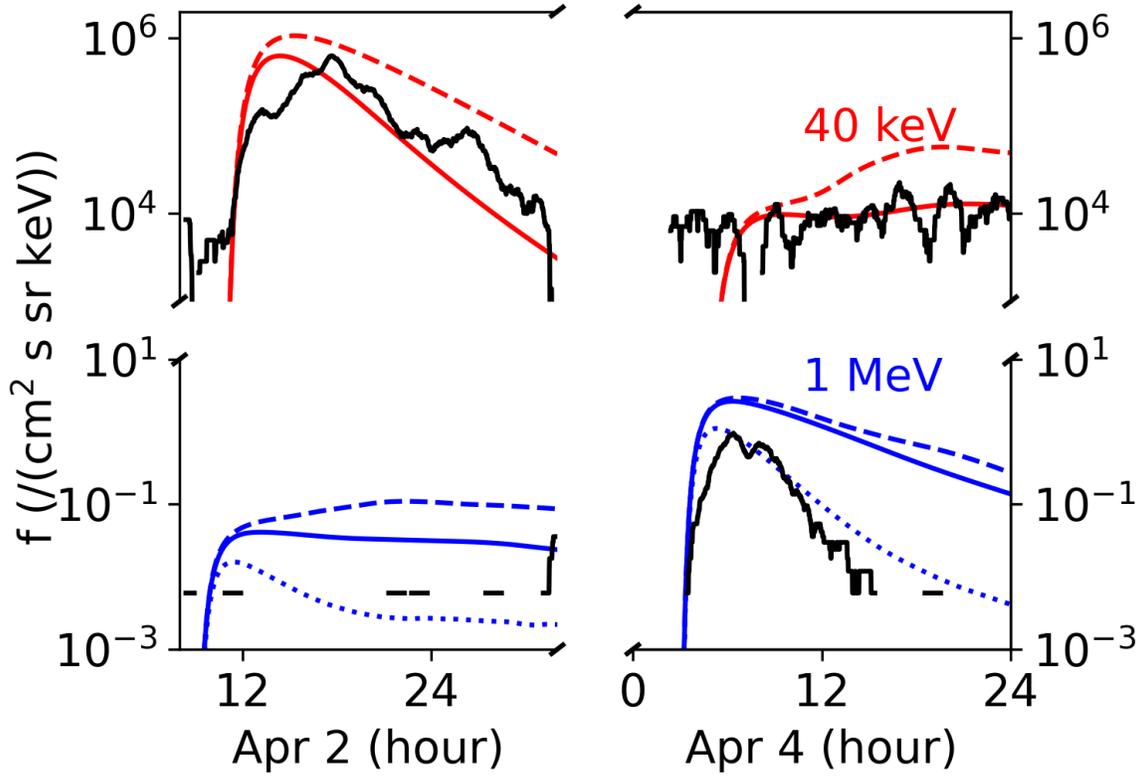


Figure 11. The outward propagating 40 keV and 1 MeV proton time-intensity profiles for the April 2 and the April 4 event. The *PSP* observations are plotted in black and the numerical calculations are plotted in red (40 keV) and blue (1 MeV). The solid and dotted curves plot the intensity profiles considering the supercorotation of *PSP*. The dashed curves do not consider the supercorotation. The solid and dashed curves correspond to the case when the decay time scale (τ_d) of the injection profile is 0.2 day. The dotted curves assume the decay time scale, τ_d , is 0.08 day (only plotted for 1 MeV protons).

Table 1. Magnetic Field Line Footpoint of *PSP*

Event	<i>PSP</i> location			$2.5 R_s$		$1 R_s$	
	r (AU)	co-Lat	Lon+60	co-Lat	Lon+60	co-Lat	Lon+60
April 2	0.19	93	331	93	342	45	322
April 4	0.17	94	355	94	5	45	324
March 22	0.17	94	355	94	5	47	356
April 11	0.17	94	355	94	5	153	12

NOTE—The magnetic field line footpoints of *PSP* in the 2019 April 2 and 4 events. Columns 2-4 show the spacecraft locations. Columns 5-6 show the magnetic field line footpoints on the source surface ($2.5 R_s$) that connect to the spacecraft. Columns 7-8 show the footpoints on the photosphere ($1 R_s$) obtained by tracing along the magnetic field line from $2.5 R_s$ to $1 R_s$ using the PFSS model. The longitude and latitude are given in the HEEQ coordinates in the unit of degree.

REFERENCES

- 469 Altschuler, M. D., & Newkirk, G. 1969, *Sol. Phys.* 496
470 9, 131 497
- 471 Bale, S. D., Goetz, K., Harvey, P. R., et al. 2016, 498
472 *Space Sci. Rev.*, 204, 49 499
- 473 Bale, S. D., Badman, S. T., Bonnell, J. W., et al. 500
474 2019, *Nature*, 576, 237 501
- 475 Howard, R. A., Moses, J. D., Vourlidas, A., et al. 502
476 2008, *Space Sci. Rev.*, 136, 67 503
- 477 Jokipii, J. R. 1966, *Astrophys. J.*, 146, 480 504
- 478 Kasper, J. C., Abiad, R., Austin, G., et al. 2016, 505
479 *Space Sci. Rev.*, 204, 131 506
- 480 Kolmogorov, A. N. 1941, *Dokl. Akad. Nauk.* 507
481 *SSSR*, 30, 301 508
- 482 Leblanc, Y., Dulk, G., & Bougeret, J. 1998, *Sol.* 509
483 *Phys.*, 183, 165 510
- 484 Leske, R. A., Christian, E. R., Cohen, C. M. S., 511
485 et al. 2020, *Astrophys. J.*, 246, 1 512
- 486 McComas, D. J., Alexander, N., Angold, N., et al. 513
487 2016, *Space Sci. Rev.*, 204, 187 514
- 488 McComas, D. J., Christian, E. R., Cohen, C. M., 515
489 et al. 2019, *Nature*, 576, 516
490 doi:10.1038/s41586-019-1811-1 517
- 491 Nitta, N. V., & Derosa, M. L. 2008, *Astrophys. J.* 518
492 673, 207 519
- 493 Qin, G., Zhang, M., & Dwyer, J. R. 2006, *J.* 520
494 *Geophys. Res.*, 111, A08101 521
- 495 Reid, G. C. 1964, *J. Geophys. Res.*, 69, 2659
- Riley, P., Linker, J. A., & Mikic, Z. 2006,
Astrophys. J., 653, 1510
- Roelof, E. C., Allen, R. C., Bale, S. D., et al.
2019, American Geophysical Union Abstract
- Schatten, K., Wilcox, J., & Ness, N. 1969, *Sol.*
Phys., 9, 131
- Schlickeiser, R. 2002, *Cosmic Ray Astrophysics*
(Springer Berlin Heidelberg)
- Schrijver, C. J., & Derosa, M. L. 2003, *Sol. Phys.*,
212, 165
- Schrijver, C. J., & Title, A. M. 2011, *J. Geophys.*
Res., 116, 1
- Schwadron, N. A., Bale, S., Bonnell, J., & et al.
2020, *ApJS*, 246, 33
- Wiedenbeck, M. E., Allen, R. C., Bale, S., & et al.
2020, *ApJS*, 246, 42
- Wuelser, J.-P., Lemen, J. R., Tarbell, T. D., et al.
2004, *Telesc. Instrum. Sol. Astrophys.*, 5171, 111
- Zhang, M., Qin, G., & Rassoul, H. 2009,
Astrophys. J., 692, 109
- Zhang, M., & Zhao, L. 2017, *Astrophys. J.*, 846,
107
- Zhao, L., & Zhang, M. 2018, *Astrophys. J. Lett.*,
859, L29
- Zhao, L., Zhang, M., & Rassoul, H. K. 2016,
Astrophys. J., 821, 62