

**An Efficient and Accurate Algorithm for Computing Grid-Averaged Solar
Fluxes for Horizontally Inhomogeneous Clouds**

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1 **Abstract**

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A computationally efficient method is presented to account for the horizontal cloud

inhomogeneity by using a radiatively equivalent plane parallel homogeneous (PPH) cloud. The

algorithm can accurately match the calculations of the reference (rPPH) independent column

approximation (ICA) results, but use only the same computational time required for a single

plane parallel computation. The effective optical depth of this synthetic sPPH cloud is derived by

exactly matching the direct transmission to that of the inhomogeneous ICA cloud. The effective

scattering asymmetry factor is found from a pre-calculated albedo inverse look-up-table that is

allowed to vary over the range from -1.0 to 1.0. In the special cases of conservative scattering

and total absorption, the synthetic method is exactly equivalent to the ICA, with only a small bias

(about 0.2% in flux) relative to ICA due to imperfect interpolation in using the look-up tables. In

principle, the ICA albedo can be approximated accurately regardless of cloud inhomogeneity.

For a more complete comparison, the broadband shortwave albedo and transmission calculated

from the synthetic sPPH cloud and averaged over all incident directions, have the RMS biases of

0.26% and 0.76%, respectively, for inhomogeneous clouds over a wide variation of particle size.

The advantages of the synthetic PPH method are that (1) it is not required that all the cloud

subcolumns have uniform microphysical characteristic, (2) it is applicable to any 1D radiative

transfer scheme, and (3) it can handle arbitrary cloud optical depth distributions and an arbitrary

number of cloud subcolumns with uniform computational efficiency.

Keywords: cloud inhomogeneity, radiative transfer, solar radiation.

1 **1. Introduction**

2 Clouds have very important effect on the Earth’s radiation budget of the atmosphere, thus
3 making clouds one of the fundamental issues in the study and modeling of the climate. Accurate
4 computation of radiative fluxes and absorption in clouds is needed to assess their impact on
5 climate, but rigorous radiative transfer computations (e.g., Hansen and Travis, 1974) are only
6 feasible in the framework of plane-parallel homogeneous (PPH) geometry. In contrast, surface
7 and space observations show clouds to be intrinsically inhomogeneous. Even stratus clouds that
8 appear to be uniform, are not actually homogeneous (e.g., King et al., 2013; Oreopoulos and
9 Cahalan, 2005; Zhang et al., 2019, Madhavan et al., 2016), and no terrestrial cloud can adhere
10 strictly to the horizontally unbounded PPH assumption.

11 It is well known that simply using the average value of the optical depth in PPH radiative
12 transfer calculations for a horizontally inhomogeneous cloud field will tend to overestimate the
13 reflectivity and underestimate the transmission (Cahalan et al., 1994a; Carlin et al., 2002; Bauml
14 et al., 2004). To address this problem, Cahalan et al. (1994a) employed the simple expedient of
15 an effective thickness approach (ETA), where ETA utilizes an effective cloud optical thickness
16 in the radiation computation that is smaller than the PPH thickness. However, the ETA reduction
17 factor has to be determined empirically and this method holds only for a rather narrow range of
18 optical depths and solar zenith angles (Barker, 1996).

19 Hogan et al. (2016) presented a method to approximate the 3D horizontal inhomogeneity
20 effects by adding extra terms to the two-stream equations to represent lateral transport between
21 clear and cloudy regions. The cloud inhomogeneity in this method is simply represented as
22 “Tripleclouds”: clear, thin cloud, and thick cloud (Shonk and Hogan, 2008). Other studies have
23 used a probability density function (PDF) representation of cloud optical depth (Barker, 1996;

1 Oreopoulos and Barker, 1999; Kato et al., 2005). For a Gamma distribution, Barker (1996)
2 showed that the domain-averaged fluxes can be computed as a function of two parameters,
3 expressing the PDF with a two-stream radiative transfer scheme. This approach is, however,
4 limited to the simple two-stream radiative transfer algorithm, and also requires about four times
5 more computational effort than the standard two-stream scheme (Barker, 1996). This approach
6 also assumes that clouds will have uniform microphysical characteristics with a uniform single
7 scattering albedo and asymmetry factor. But, perhaps as a more significant concern, the two-
8 stream approach has substantial intrinsic radiative transfer uncertainty issues that are on the order
9 of 10% compared to more precise radiative transfer modeling schemes (Barker et al., 2015; King
10 and Harshvardan, 1986).

11 A commonly used method of estimating solar fluxes for horizontally inhomogeneous
12 clouds is the independent column approximation (ICA). In this approach, the cloud field is
13 partitioned into a number of subcolumns to account for the inhomogeneity. The ICA method
14 applies a PPH radiative transfer model to each column, and then averages over all columns to
15 obtain the domain mean fluxes. Cahalan et al. (1994b) demonstrated that the ICA fluxes could be
16 in excellent agreement with full 3D Monte Carlo calculations. Neu et al (2007) proposed a cloud
17 quadrature approach dividing the cloud field into four representative independent columns based
18 on cloud optical depth. This approach could achieve an RMS error of less than 4% in actinic
19 fluxes but the instantaneous error could be large if the cloud heterogeneity is high, or if the solar
20 zenith angle is large (Prather. 2015).

21 The ICA approach can be considered as an accurate radiative transfer scheme to the
22 extent that the independent columns are a good representation for inhomogeneous cloud field.
23 But its computation time for radiative transfer increases proportionally with the number of

1 independent cloud columns. To improve the computation efficiency, several simplified methods
2 have been developed (e.g., Evans, 1993; Barker, 1996; Gabriel and Evans, 1996; Oreopoulos and
3 Barker 1999; Cairns et al. 2000; Kato et al., 2005). Some of these methods (e.g., Borde and Isaka
4 1996; Cairns et al. 2000) account for the effect of cloud inhomogeneity by rescaling the bulk
5 cloud scattering radiative parameters. As a result, the computation of radiative fluxes for the
6 inhomogeneous cloud cases can thus be transformed into those for an equivalent single column
7 PPH cloud. There is also the Monte Carlo Independent Column Atmosphere (McICA) method
8 (Pincus et al., 2003) that randomly selects a sub-column as a representative of clouds in a climate
9 GCM grid box. This stochastic approach is computationally efficient but it requires large domain
10 averaging to reduce the sampling error.

11 As a specific example of cloud heterogeneity treatment in a climate GCM, in the Cairns
12 et al. (2000) Monte Carlo based parameterization, the GCM grid-box-mean homogeneous cloud
13 radiative parameters (optical depth, single scattering albedo, and asymmetry parameter) are
14 transformed in accord with a cloud heterogeneity parameter $V = \exp(\delta^2)$, where δ is the *log*
15 standard deviation of the cloud (microscopic and/or macroscopic) density distribution. This
16 parameterization is used in the GISS ModelE climate GCM (Schmidt et al., 2006) utilizing the
17 Rossow et al. (2002) ($4^\circ \times 5^\circ$ monthly-mean) empirical determination of the Cairns et al. (2000)
18 cloud heterogeneity parameter V , based on the ISCCP D1 cloud climatology. The GISS ModelE
19 radiation model utilizes the Single Gauss Point (SGP) doubling/adding method as its radiative
20 transfer solver, using a look-up table to select its input radiative parameters (Hansen et al., 1983)
21 so as to precisely reproduce the planetary albedo for any solar zenith angle and cloud optical
22 depth for the case of conservative scattering.

1 In this paper we present a deterministic synthesizing approach to define a radiatively
2 equivalent PPH cloud to represent an inhomogeneous cloud field, and to thus provide for
3 efficient and accurate computation of solar fluxes for inhomogeneous clouds. Different from
4 previous deterministic methods, the single scattering optical properties are not required to be
5 uniform across the cloud subcolumns, and the subregion cloud optical depth distribution can be
6 arbitrary. This method is applicable to the utilization of sophisticated 1D radiative transfer
7 algorithms via look-up tables, which, in a sense, serve as the reference standard for comparison.
8 Section 2 describes this synthesizing approach in detail and its use of direct and inverse look-up
9 tables that also take on a radiative transfer solver role. Section 3 demonstrates the utility of this
10 methodology, displaying through comparisons with ICA results the numerical errors and biases
11 that arise for different examples of cloud heterogeneities, with concluding remarks contained in
12 section 4.

13 **2. Methodology**

14 In the ICA scheme the regional, or grid-box-mean cloud fluxes represent the summation
15 of individual cloud fluxes computed separately over all subregions of the grid-box. For a cloud
16 field divided into N subcolumns, the ICA method conducts plane-parallel radiative transfer
17 calculations for each column individually and then averages the subcolumn results to obtain the
18 grid-box regional mean. Thus for example, the ICA cloud albedo is expressed as

$$A_{ICA}(\mu_0) = \sum_{i=1}^N f_i A_i(\tau_i, \mu_0) \quad (1)$$

19 where μ_0 is the cosine of the solar zenith angle. f_i , A_i and τ_i are the column area fraction, the PPH
20 albedo, and the optical depth for the i th subcolumn, respectively. The ICA transmittance and
21 absorptance for a multiple-column inhomogeneous cloud are calculated in the same manner as
22 defined for the albedo.

1 While the ICA approach is frequently used to model radiative transfer for inhomogeneous
2 cloud configurations, computation time can be a significant issue if the number of subcolumns is
3 large. Here we propose an approach that combines all of the cloud subcolumns into a radiatively
4 equivalent single column synthetic PPH cloud (labeled hereafter as sPPH for clarity), such that
5 the ICA fluxes can be reproduced by a single plane parallel computation. To demonstrate the
6 basic concept of this synthesizing approach, we start with a heterogeneous cloud comprised of
7 four columns. The optical depths for the four columns are taken to be [0.3, 3, 10, 30], and the
8 corresponding asymmetry parameters for the particle scattering functions are taken to be,
9 respectively, [0.86, 0.84, 0.82, 0.80]. For simplicity, and not as a requirement, we assume
10 conservative scattering and equal area fractions for all columns. For uniform solar incidence on
11 such a 4-column cloud, the column-mean direct transmissivity (T_{dir}) is simply

$$T_{dir} = \frac{1}{4} \sum_{i=1}^4 e^{-\tau_i/\mu_0} \quad (2)$$

12 It then follows that the effective optical depth for a sPPH cloud having the same direct
13 transmission T_{dir} is

$$\tau_e = -\mu_0 \log(T_{dir}) \quad (3)$$

14 This derivation of the effective optical depth has also been adopted by Prather (2015). After the
15 effective optical depth of the sPPH cloud is determined, the next step is to determine an
16 asymmetry factor such that the sPPH cloud albedo is equal to the ICA computation. All of these
17 properties depend on solar zenith angle. To illustrate, for a solar zenith angle of 60°, the effective
18 optical depth of the single column sPPH cloud is 1.19, and the ICA albedo for the entire 4-
19 column cloud is 0.4843. Through radiative transfer iteration, the effective asymmetry factor for
20 the sPPH cloud that would reproduce the ICA column-mean albedo, is 0.232. In this way, a

1 synthetic sPPH cloud is defined that is radiatively equivalent to the heterogeneous four-column
2 cloud. In practice, this effective asymmetry factor for the sPPH cloud would be retrieved from a
3 pre-calculated *inverse* look-up table.

4 For our radiative transfer solver forward calculations, the forward look-up table for the cloud
5 albedo is tabulated for a range of incident zenith angles (μ_0), optical depth (τ), asymmetry factor
6 (g), and single scattering albedo (ϖ). In this table, there are a total of $30 \times 32 \times 21 \times 18$ grid points,
7 representing 30 grids in μ_0 from 0.02 to 1.0, 32 in τ from 0.05 to 210, 21 in g from -0.95 to 0.95,
8 and 18 in ϖ from 0.0 to 1.0. Because cloud albedo and transmittance approach asymptotes as
9 τ increases, values beyond $\tau=210$ are considered as unchanged. Since the radiative fluxes are
10 nonlinear functions of all four variables, the grid points are unevenly distributed. For example,
11 half of ϖ grids are in the range of 0.9 to 1.0.

12 As a general recipe, for a cloud with N columns, the complete processes to find the effective
13 optical properties of the equivalent sPPH cloud is to:

- 14 1) Calculate the effective optical depth for the sPPH cloud according to equation (3).
- 15 2) Obtain the effective single scattering albedo, $\varpi_e = \frac{\sum_{i=1}^N f_i \tau_i \omega_i}{\sum_{i=1}^N f_i \tau_i}$
- 16 3) Obtain the albedo, $A_i(\mu_0, \tau_i, \omega_i, g_i)$, of each column using the forward look-up-table.
- 17 4) Calculate the ICA albedo, A_{ICA} , using equation (1).
- 18 5) Find the effective asymmetry factor, g_e , corresponding to the ICA column-mean albedo,
19 using the inverse look-up-table tabulated as a function of albedo, μ_0 , τ_e , and ϖ_e .

20 Through this process of these five steps, a radiatively equivalent sPPH cloud to the N
21 columns of inhomogeneous cloud is found, as defined by the effective sPPH cloud optical
22 properties τ_e , g_e and ϖ_e .

1 We now apply this procedure to obtain the effective sPPH optical properties for the
2 example of the 4-column cloud described above. Figure 1 compares the albedo and transmission
3 of the synthetic sPPH cloud with those for reference rPPH ICA results for different incident
4 angles (the upper panel). To avoid the algorithm errors inherent in 2-stream radiative transfer
5 schemes (Barker et al., 2015; Jin et al., 2019), we use a 16-stream discrete ordinate algorithm for
6 all (rPPH cloud) radiative transfer and look-up table calculations. The black asterisks represent
7 the rPPH ICA results and the red diamonds represent the results calculated from the sPPH cloud.
8 The albedo and transmission calculated based on the simple mean cloud optical depth and mean
9 asymmetry factor (weighted by subcolumn optical depth) are also plotted for comparison (the
10 blue pluses). The lower panel shows the corresponding effective optical depth (black) and the
11 asymmetry factor (red) derived for the synthetic sPPH cloud.

12 Figure 1 shows that both the ICA albedo and transmission are reproduced accurately by
13 the simple one column of sPPH cloud. The RMS errors are only 0.25% and 0.24% for the albedo
14 and transmission, respectively, and are due to imperfect look-up-table interpolation. However,
15 for this specific example, the radiative transfer computation time using the synthetic sPPH cloud
16 is reduced by a factor of 4 compared with the full ICA approach. In general, the computation
17 time to be saved depends on the total number of cloud subcolumns. The more subcolumns, the
18 more time is saved. Computations that simply utilize the grid-box-mean cloud optical properties
19 might be equally as efficient in computing time, but as shown in figure 1, doing so would result
20 in a major loss of accuracy by greatly overestimating the albedo and underestimating the
21 transmission, compared to the synthetic sPPH approach.

22 It should be noted that the Henyey-Greenstein single scattering phase function is equally
23 well-defined for both positive and negative values of the asymmetry parameter. Accordingly, the

1 effective asymmetry factor for the synthetic sPPH cloud can be negative because the asymmetry
 2 factor is used here as a radiative parameter that simply regulates the relative amount upward and
 3 downward scattering, and is not used, or intended, as a physical representation of real cloud
 4 properties, for which negative values of asymmetry factor are not typically encountered in
 5 nature. Thus for example, the effective asymmetry factor needed to reproduce the ICA albedo is
 6 -0.567 when $\mu_0=0.1$ in figure 1.

7 **3. Application to General Heterogeneous Clouds**

8 Inhomogeneity in cloud distribution is often described in terms of cloud optical depth τ .
 9 Cloud heterogeneity can be in various formats, as there is no single formula to define all possible
 10 distributions. However, previous studies have shown that the τ distribution in many cloud fields
 11 can often be approximated by means of a Gamma distribution (e.g., Barker, 1996; Morrison and
 12 Gettelman, 2008; Oreopoulos and Barker, 1999; Oreopoulos and Cahalan, 2005; Zhang et al.,
 13 2019). The Gamma distribution function is expressed as

$$14 \quad p(\tau) = \frac{1}{\beta \Gamma(\nu)} \left(\frac{\tau}{\beta}\right)^{\nu-1} e^{-\tau/\beta} = \frac{1}{\Gamma(\nu)} \left(\frac{\nu}{\tau_m}\right)^{\nu} \tau^{\nu-1} e^{-\nu\tau/\tau_m} \quad (4)$$

15 where τ_m is the average cloud optical depth, and ν is the gamma distribution shape parameter.
 16 $\Gamma(\nu)$ is the Gamma function and $\nu = (\tau_m/\sigma)^2$ in which σ is the standard deviation of τ . Smaller
 17 ν indicates a higher degree of inhomogeneity. $\beta = \tau_m/\nu$ is the scale parameter. The larger the
 18 scale parameter, the more spread out the distribution.

19 Figure 2 shows an example of the Gamma distribution with mean optical depth $\tau_m=10$
 20 for four different shape parameters (ν), in which $\nu=8$ represents the least heterogeneity, and
 21 $\nu=0.5$ represents the most heterogenous case. The horizontal axis depicts the cloud optical depth
 22 τ in logarithmic scale. Different bars represent the probability or cloud fraction in a different τ

1 bin. The bin width is uniformly 0.5 in logarithmic scale. Each bin center corresponds to the mean
2 τ of the bin, which is 0.05 for the left-most bin, and 148 for the right-most bin in Figure 2. As
3 ν decreases, the distribution range of τ increases. A broad range of cloud inhomogeneities can be
4 approximated by different combinations of ν and τ_m in the distribution function of Equation (4).

5 We now apply the synthetic approach to inhomogeneous clouds with four different shape
6 parameters as shown in figure 2, and with different mean optical depth τ_m . Each bin in figure 2
7 is considered as one cloud subcolumn, and the probability for each bin is the cloud fraction for
8 the subcolumn. Clearly, the cloud field with smaller ν will have more cloud columns because of
9 wider variation range of τ . In our actual calculations, we ignore those columns with PDF value
10 (cloud fraction) less than 0.01, or τ less than 0.28. Under these conditions, the total number of
11 cloud columns is 4 for $\nu = 8$ and 12 for $\nu = 0.5$ for the example shown in figure 2. It should be
12 noted that this synthetic method could be applied to any τ distribution. We use the Gamma
13 distribution here because it is convenient and representative.

14 Figure 3 shows the biases in albedo and transmission calculated with the synthetic sPPH
15 cloud from the ICA calculations. Again, albedo and transmission for each cloud subcolumn and
16 the rPPH cloud are calculated using the 16-stream radiative transfer scheme. The number of
17 cloud subcolumns and the fraction of each column for ICA are generated as shown in figure 2.
18 The horizontal and vertical coordinates in each panel represent the cosine of the solar zenith
19 angle (μ_0), and the cloud mean optical depth in logarithmic scale ($\log(\tau_m)$), respectively.
20 Different rows depict a different heterogeneity parameter. The left panels show the relative
21 biases of albedo in per cent, and the right panels show the transmission biases. Because the
22 transmission at large optical depth or large incident angle could be near zero, where tiny
23 differences could appear very large on a relative scale, the transmission bias is presented as

1 actual differences times a factor of 100. The dotted white lines are contours of bias. Conservative
2 scattering and a Henyey-Greenstein cloud scattering asymmetry factor of 0.86 are assumed here.
3 Regardless of the degree of inhomogeneity, figure 3 shows that the relative albedo biases for all
4 incident angles are less than 0.5% and that the transmission differences are within ± 0.003 .

5 Figure 4 compares the spherical albedo and transmission between the ICA approach and
6 the synthetic PPH approximation. The spherical albedo (transmission) can be regarded as the
7 averaged albedo (transmission) over all incident angles. Different panels represent a different
8 cloud inhomogeneity. In each panel, the black asterisk is for ICA calculation and the red
9 diamond is for that of the synthetic PPH approximation. The horizontal axis represents the
10 logarithmic cloud mean optical depth, $\log(\tau_m)$. Again, both the ICA albedo and transmission
11 based on the synthetic PPH cloud agree well with the intensive ICA computations in all cases.
12 The RMS errors in both the spherical albedo and transmission are 0.0004. Again, this small error
13 is simply due to imperfect interpolation of look-up-table variables, which is the only error source
14 when the scattering is conservative.

15 For conservative scattering, transmission has the same biases as the albedo but with
16 opposite direction. However, as shown in figure 5, this does not apply for non-conservative
17 scattering cases, in which case, a correct transmittance requires not only accurate reflection but
18 also accurate in-cloud absorption. Figure 5 is a counterpart of figure 3 but with particle single
19 scattering albedo of 0.98. For this absorptive case, the albedo bias is as small as in conservative
20 scattering, but the transmission bias becomes much larger. Generally, the transmission is
21 positively biased and the bias increases systematically with the optical depth. It also increases as
22 the cloud inhomogeneity increases. Under non-conservative scattering conditions, scaling the
23 single scattering albedo is required to reduce the bias. We have found that the transmission bias

1 can be significantly reduced by applying the following scaling factor to the effective single
 2 scattering albedo (ω_e)

$$3 \quad C = e^{-\left[\omega_e^{0.4}(1-\omega_e)^{0.8}\left(\frac{3.2}{(\tau_e+2)^{1.2}}+0.1\right)\left(3.5\left(1-\frac{\tau_e^0}{\tau_{mn}}\right)^{1.3}+0.1\right)S(6-1.5\tau_a)\right]} \quad (5)$$

4 where τ_{mn} is the mean cloud optical depth averaged in logarithmic scale and τ_e^0 is the effective
 5 optical depth of the synthetic PPH cloud for a $\mu_0=1$ solar zenith angle. $\tau_a = \tau_e(1-\omega_e)$ is the
 6 absorptive optical depth of the synthetic cloud. S represents the Sigmoid function, which is
 7 $S(x) = \frac{1}{1+e^{-x}}$. For either conservative scattering ($\omega_e=1.0$) or total absorption ($\omega_e=0$), the
 8 synthetic approach is accurate, and thus there is no correction required. In these special cases, the
 9 scaling factor C from equation (5) is 1.0, indicating no adjustment applied to ω_e . As τ_a increases,
 10 the Sigmoid function rapidly approaches zero, in which case, no adjustment is required either.

11 Figure 6 shows the same computation as figure 5, but with the adjustment of equation (5)
 12 applied to the effective single scattering albedo of the sPPH cloud. In comparing figure 6 to
 13 figure 5, the bias in albedo is found to be similar without much effect, but the bias in
 14 transmission is reduced significantly.

15 Figure 7 compares the spherical albedos and transmissions for this absorptive condition
 16 ($\omega_e=0.98$) with different panels representing a different inhomogeneity. In each panel, the black
 17 line is for rPPH ICA calculations, the blue is for the synthetic sPPH approach without adjustment
 18 of ω_e , and the red with adjustment of ω_e . The results show that the transmission can be well
 19 simulated after the ω_e adjustment except for the highly heterogenous clouds with large mean
 20 optical thickness, while the albedo is accurately simulated regardless of the adjustment.

21 The computations thus presented have demonstrated the methodology of the synthetic
 22 sPPH approach but have used the artificially assumed cloud optical properties and uniform

1 microphysical cloud characteristics. In more realistic situations, not only the optical depth, but
2 the single scattering albedo and the asymmetry factor can also be arbitrarily distributed within
3 the different subcolumns. This can be readily handled using the synthetic sPPH method.

4 We now perform further tests of the algorithm using the optical properties based on real
5 cloud particle size and phase. The particle size is allowed to vary randomly in different columns.
6 The variation range of the effective cloud particle radius is from 7 μm to 80 μm . The cloud
7 composition is considered to be water if the radius is less than 17 μm , otherwise it is treated as
8 ice. For ice cloud optical properties, we use the surface-roughened aggregate model used in the
9 MODIS Collection 6 cloud product (Platnick et al., 2017). For liquid cloud, the spherical water
10 particle model developed by Hu and Stamnes (1993) is used. The same heterogeneous cloud
11 fields as were used for Figures 3-7 are used to calculate the broadband shortwave fluxes and to
12 evaluate biases in the calculated solar radiation. The incident solar spectrum is taken from
13 MODTRAN 1 cm^{-1} resolution data (Berk et al., 2008). Spectral integrals are performed over the
14 0.3 to 5.0 μm range, but without including Rayleigh scattering, aerosols, or gaseous absorbers.

15 Figure 8 is the counterpart of figure 6 but shows the biases in shortwave albedo and
16 transmission for inhomogeneous clouds with nonuniform microphysical properties. Figure 9
17 compares the spherical (or all-angle averaged) albedo and transmission calculated from the
18 synthetic sPPH cloud, along with those from the reference rPPH ICA computation. The results in
19 figure 8 indicate that the biases are within $\pm 0.7\%$ for the broadband shortwave albedo and
20 ± 0.015 in transmission for all incident angles and heterogeneity cases. The relative RMS biases
21 for the spherical albedo and transmission are 0.26% and 0.76%, respectively.

22 If the water and ice clouds are not mixed between different columns, the biases are even
23 smaller than those shown in figures 8 and 9. We show in Figures 10 and 11 only the spherical

1 shortwave albedo and transmission for water cloud and ice cloud, respectively. In these
2 calculations, the effective particle radius was allowed to change randomly in the 17 to 80 μm
3 range for ice cloud, and in the range of 7 to 17 μm for water cloud columns. The relative RMS
4 biases for the spherical albedo and transmission of water clouds are 0.24% and 0.48%,
5 respectively. Similarly, the RMS biases of albedo and transmission for ice clouds are 0.25% and
6 0.52%, respectively. In comparing figures 10 and 11 with figure 9, not only are the biases
7 smaller, but the albedo and transmission lines, representing their variation with the mean cloud
8 optical depth, are also smoother. This is because there is only one cloud phase and thus a smaller
9 variation in the optical properties among different columns. In comparing figure 10 with figure
10 11, we see that the water cloud has a higher albedo and a lower transmission than the ice cloud
11 for the same optical depth for all heterogeneous cases.

12 Table 1 summaries the RMS biases of the synthetic PPH method in the spherical
13 shortwave albedo and transmission for the different cloud inhomogeneities and particle size
14 ranges. The numbers in parenthesis represent the relative biases in per cent. The bias in spherical
15 albedo is generally small (less than 0.4%) regardless of the cloud heterogeneity and particle size
16 variation. The bias in transmission is larger but still less than 1%, and in most cases, tends to
17 show somewhat of an increase as the variation range of the cloud particle size increases.

18 We have found that occasionally the ICA albedo can become so high that the effective
19 asymmetry factor interpolated from the inverse look-up table approaches the bounds of the table.
20 This usually occurs for extreme near-horizontal incidence conditions (cosine of solar zenith
21 angle less than 0.1) for some non-conservative scattering cases. When this happens, we decrease
22 the effective single scattering co-albedo ($1-\omega_e$) by 10%, which increases ω_e by the amount of
23 $0.1(1-\omega_e)$. We then go back to the look-up table to find the new effective asymmetry factor.

1 **5. Discussion and Conclusions**

2 We describe a new method for computing the mean radiative fluxes for inhomogeneous
3 cloud fields. A radiatively equivalent sPPH cloud can be found to represent the multiple column
4 horizontally inhomogeneous cloud, being transformed into the more tractable plane parallel
5 cloud framework. In ICA, a horizontally inhomogeneous cloud region is divided into a number
6 of subregions and radiative transfer computation is applied to each subregion individually. Using
7 the synthetic sPPH cloud, the radiative transfer computation for all subcolumns is transformed
8 into a single column regardless of the total number of subcolumns that may be used initially to
9 represent the cloud inhomogeneity. The effective optical depth of the synthetic sPPH cloud is
10 derived by matching the direct transmission to that of the inhomogeneous cloud field, which is
11 calculated as the subcolumn average in ICA. The effective scattering asymmetry factor, which is
12 allowed to vary from -1.0 to 1.0, is found from a pre-calculated inverse albedo look-up-table.

13 We examined the biases of the synthetic sPPH method by comparing to the rPPH results
14 from ICA. For the special cases of conservative scattering (or total absorption), the synthetic
15 method is exact, and the radiative fluxes calculated would be equal to the ICA results, with the
16 small biases (about 0.2%) in the sPPH cases arising entirely due to imperfect interpolation in
17 using the forward look-up table. In principle, since this is primarily a function of the look-up
18 tables, the ICA albedo can be accurately approximated by the single column synthetic sPPH
19 cloud, regardless of cloud inhomogeneity. However, to obtain accurate transmission in
20 absorptive cases, a scaling factor is required to adjust the effective single scattering albedo. We
21 have tested the synthetic sPPH approach on cloud fields with different inhomogeneities and a
22 wide range of cloud particle properties. The broadband shortwave albedo and transmission agree
23 well with the corresponding rPPH ICA computations. Averaged over all incident directions, the

1 shortwave albedo and transmission calculated from the synthetic sPPH cloud have RMS errors of
2 0.26% and 0.76%, respectively, for inhomogeneous clouds over a wide range of particle size
3 with (R_e) ranging from 7 to 80 μm . The accuracy would be even higher for narrower cloud
4 microphysical property variation.

5 Compared with many approximate methods for modeling cloud inhomogeneities, the
6 advantages of the synthetic sPPH method developed here are that (1) it is not required that all the
7 cloud subcolumns have uniform microphysical characteristic (i.e., the same single scattering
8 albedo and scattering asymmetry factor), (2) it can be applied to sophisticated radiative transfer
9 schemes, (3) it can handle arbitrary cloud optical depth distributions and an arbitrary number of
10 subcolumns, and (4) it is straightforward to implement.

11 Although the synthetic sPPH method is computationally efficient and accurate, there are
12 some limitations. The computation accuracy of this method has been evaluated against the rPPH
13 ICA calculations and its viability rests on the assumption that the ICA approach is a reliable
14 reference point. Therefore, potential 3D radiative interactions between the different cloud
15 subcolumns is not included in the computations presented. Furthermore, the sPPH approach has
16 currently been applied only to single layer cloud cases. Future work should consider cloud
17 overlapping and its effect on atmospheric heating rate.

18

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Table 1. RMS Errors in Spherical Albedo and Transmission

Cloud Type (R_e range)	RMS Albedo Error				RMS Transmission Error			
	$v=8.0$	$v=4.0$	$v=1.0$	$v=0.5$	$v=8.0$	$v=4.0$	$v=1.0$	$v=0.5$
Water (7-17 μm)	0.0004 (0.10)	0.0008 (0.23)	0.0010 (0.30)	0.0010 (0.32)	0.0015 (0.25)	0.0020 (0.34)	0.0041 (0.65)	0.0034 (0.52)
Ice (17-80 μm)	0.0011 (0.25)	0.0011 (0.26)	0.0008 (0.22)	0.0010 (0.27)	0.0028 (0.55)	0.0026 (0.51)	0.0027 (0.50)	0.0030 (0.53)
Mixed (7-80 μm)	0.0011 (0.26)	0.0011 (0.26)	0.0008 (0.21)	0.0010 (0.28)	0.0046 (0.89)	0.0033 (0.65)	0.0034 (0.62)	0.0048 (0.83)

1

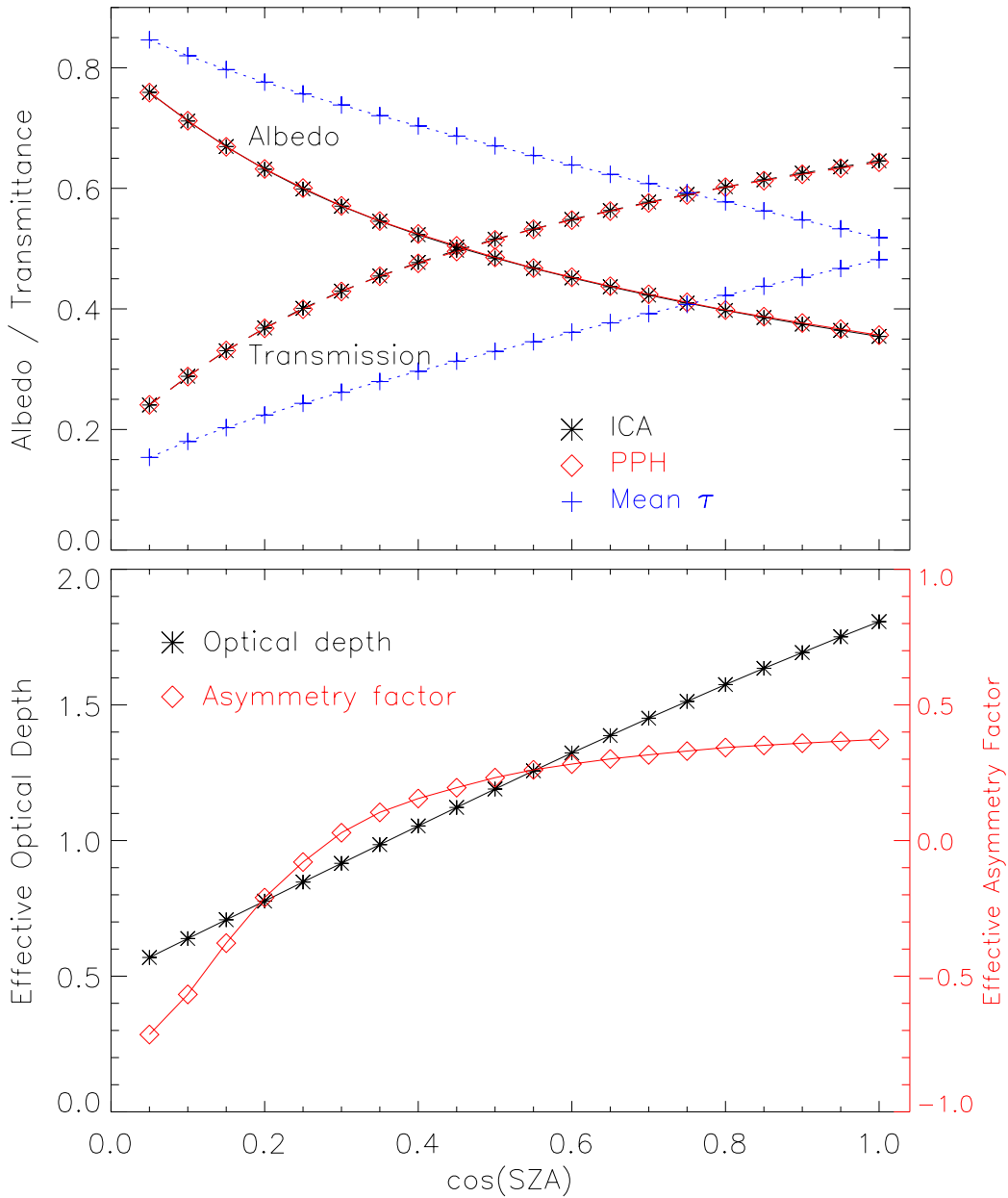


Fig. 1. Upper panel compares the albedo and transmittance of a 4-column cloud (see text for description) between the ICA method (black) and the synthetic PPH approach (red). The blue pluses represent the albedo and transmission calculated based on the simple mean cloud optical properties. Lower panel shows the corresponding effective optical depth and scattering asymmetry factor for the synthetic PPH cloud.

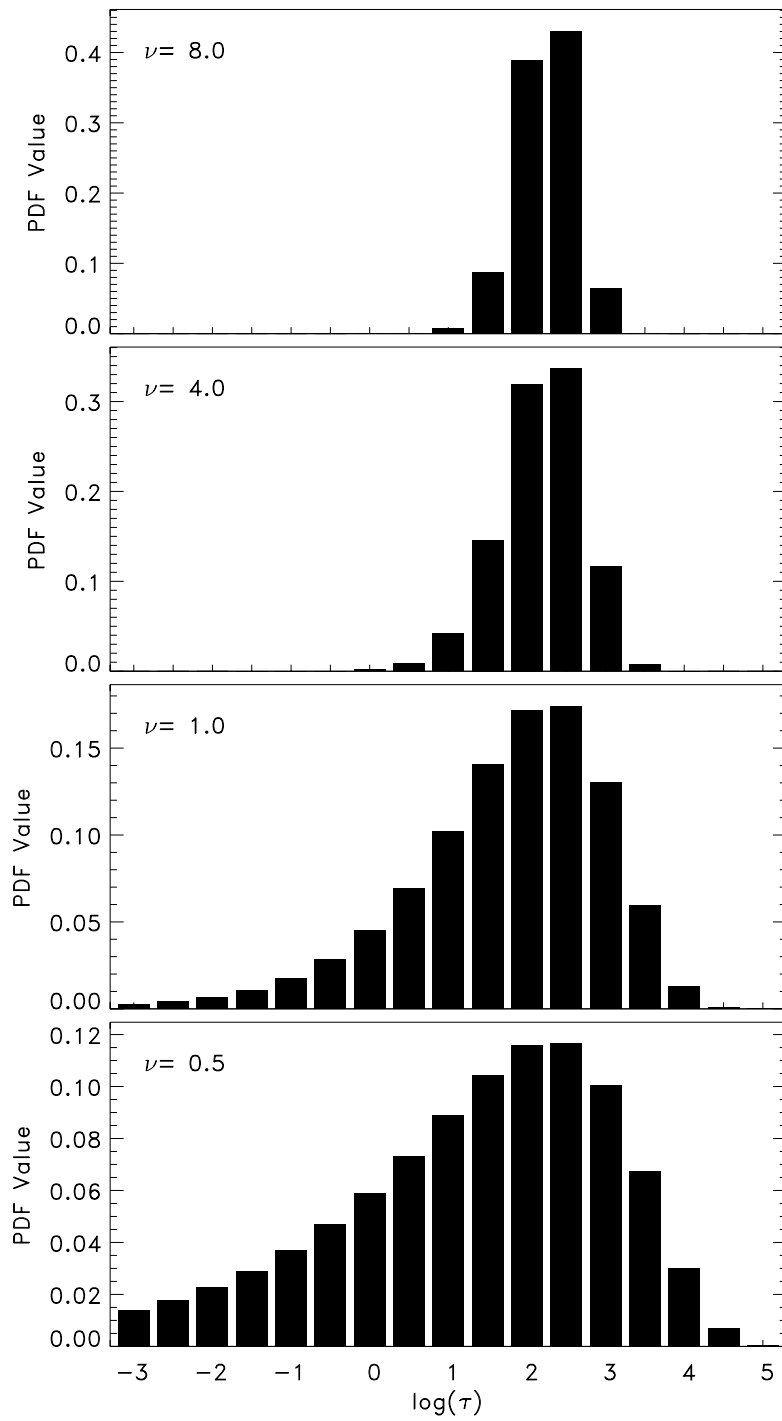
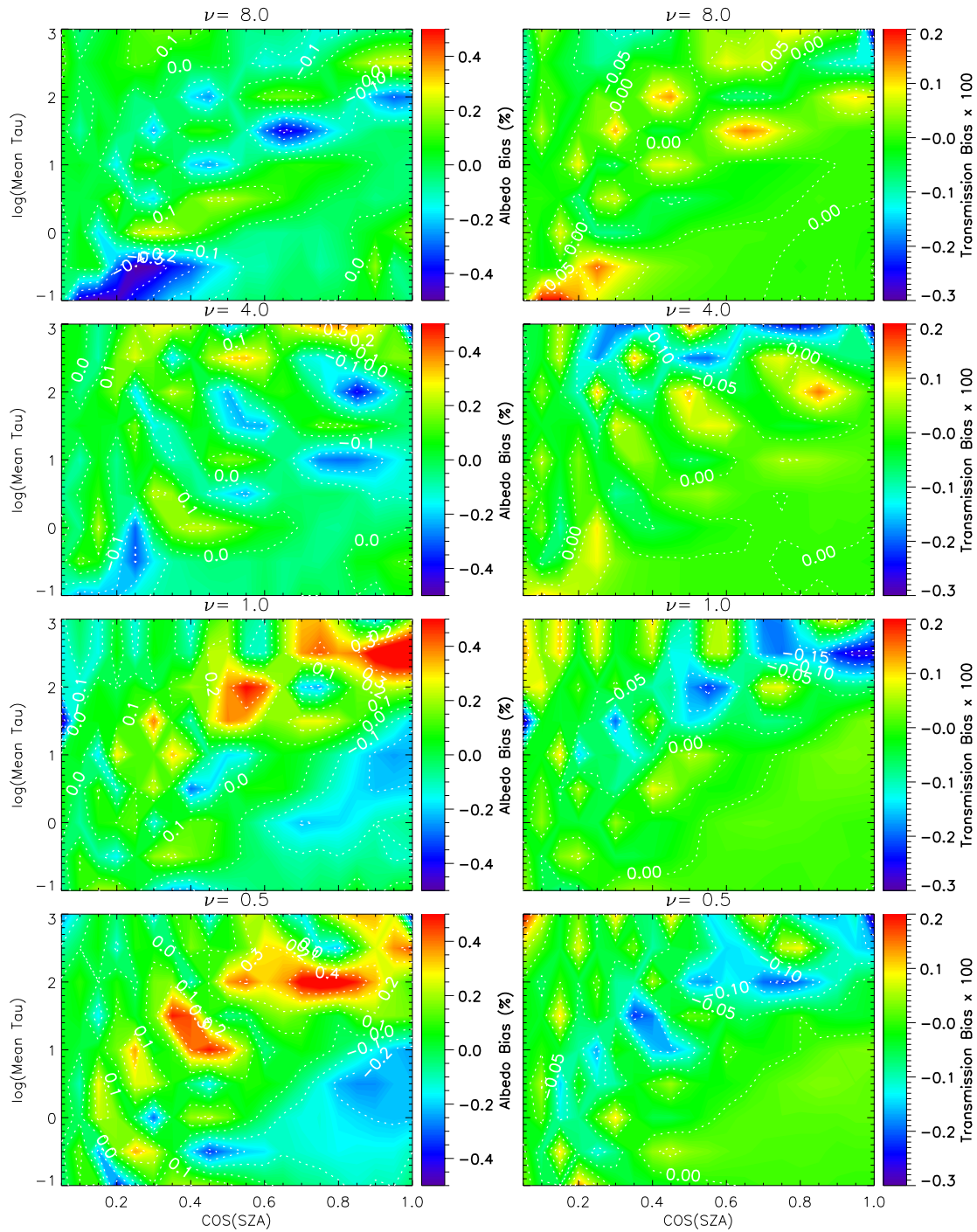


Fig. 2. An example of Gamma distributions with mean optical depth $\tau_m=10$ for four different shape parameters, ν . Heterogeneity increases from top ($\nu=8$) to bottom ($\nu=0.5$). A different bar represents the probability or cloud fraction in a different $\log(\tau)$ bin of 0.5 in width.

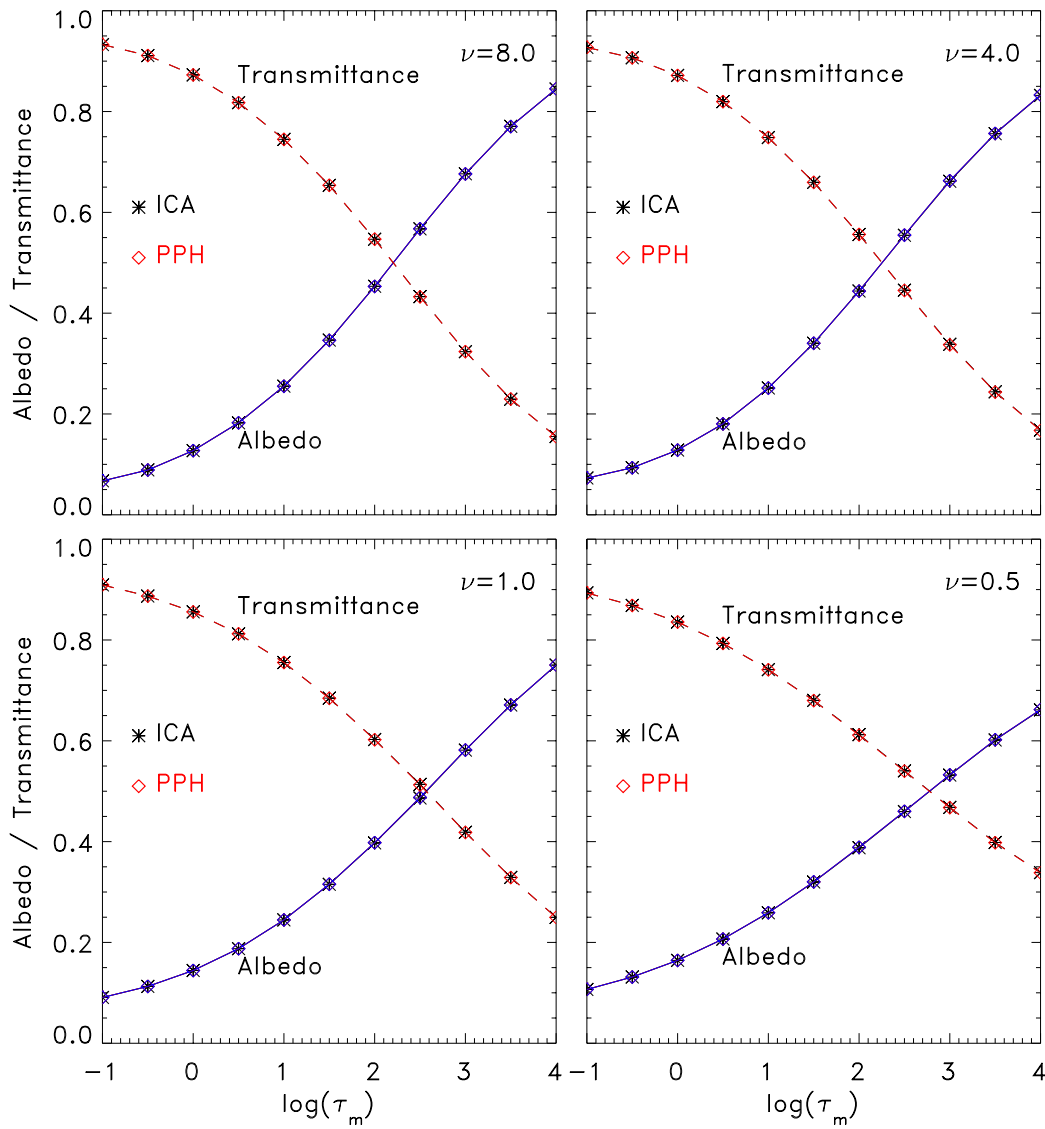
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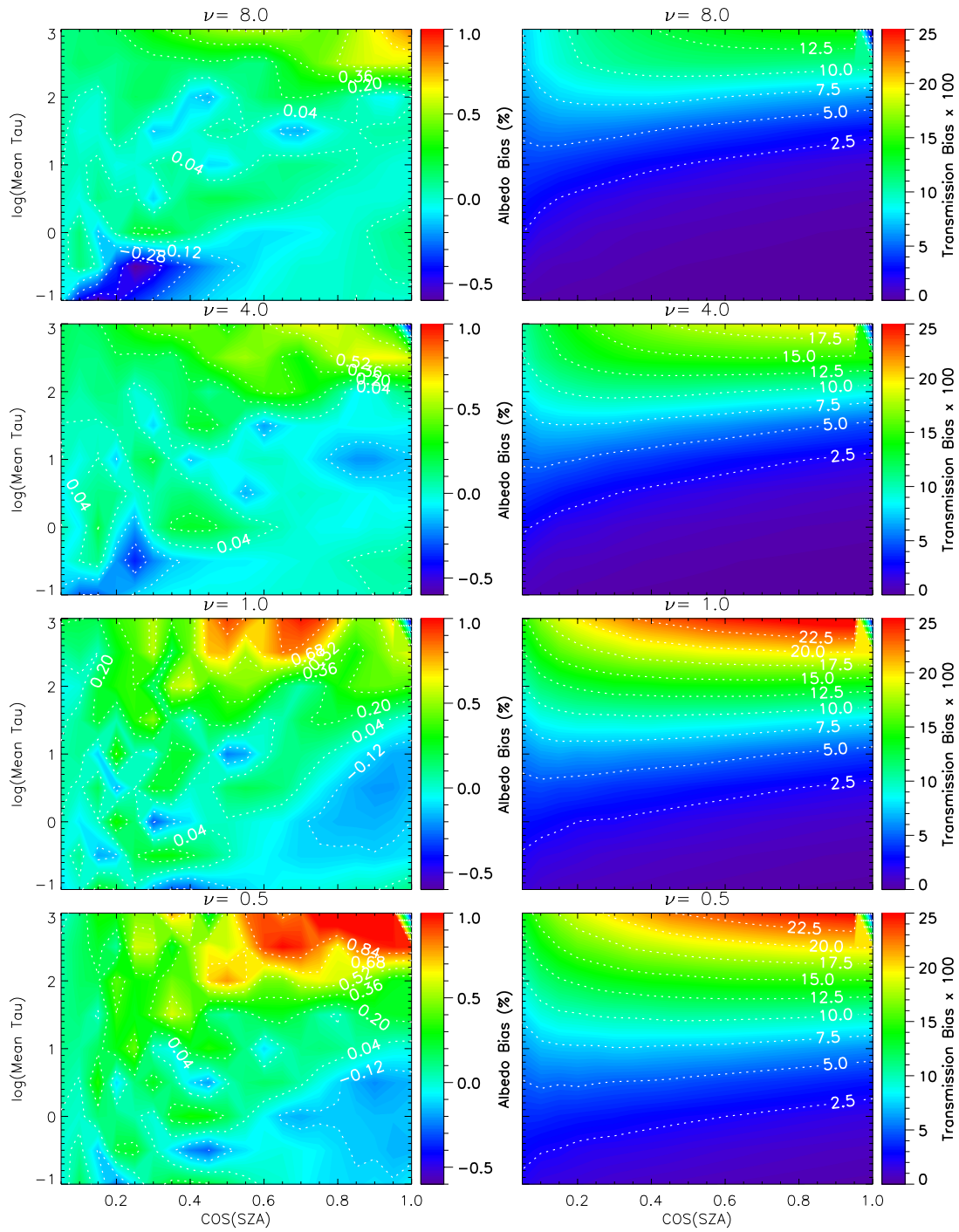
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Fig. 3. Biases in albedo (left) and transmission (right) calculated with the synthetic PPH cloud from the ICA calculations. The horizontal and vertical coordinates in each panel represent the cosine of zenith angle and the cloud mean optical depth ($\log(\tau_m)$), respectively. The number of cloud subcolumns and the fraction of each column are generated as in figure 2.

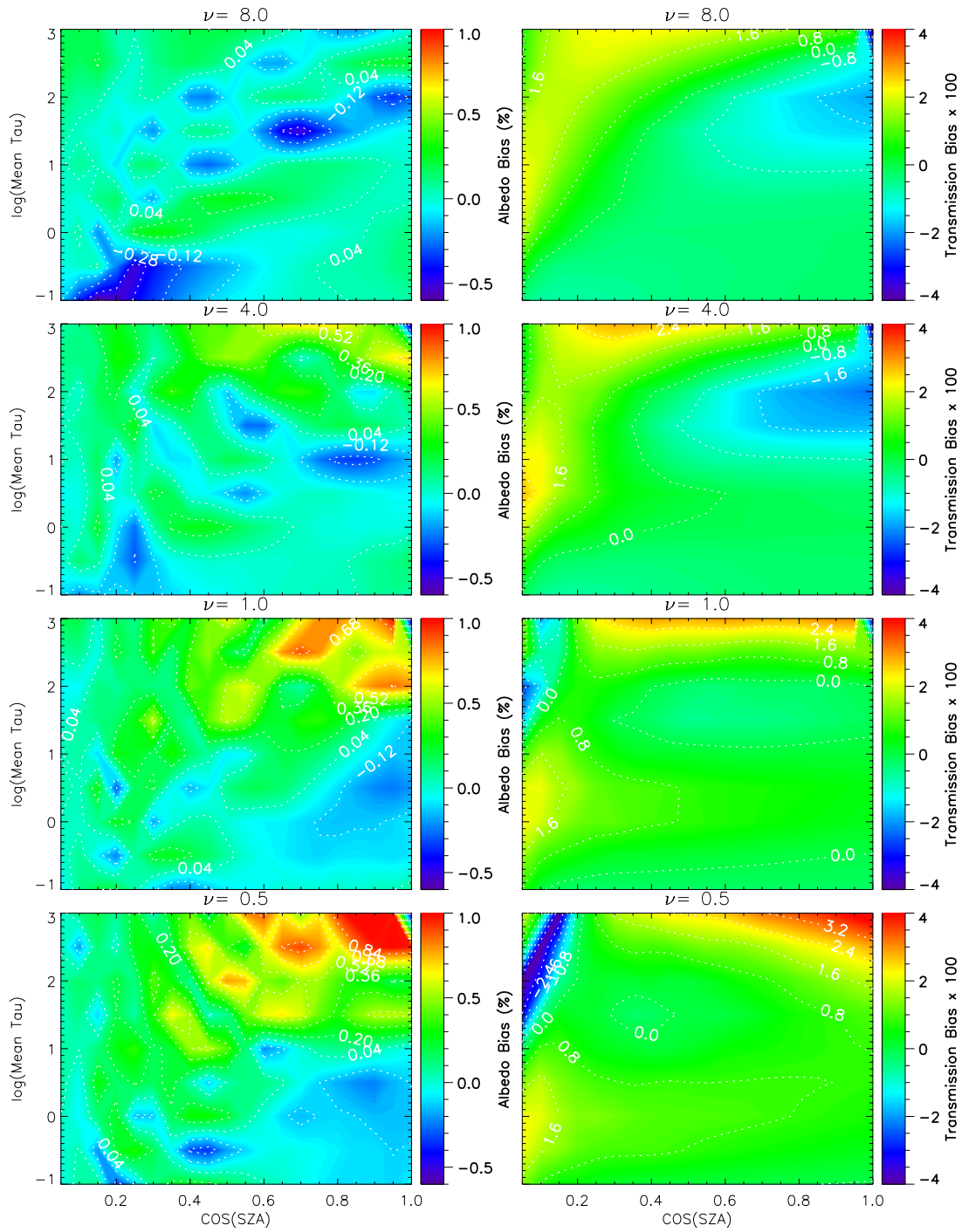
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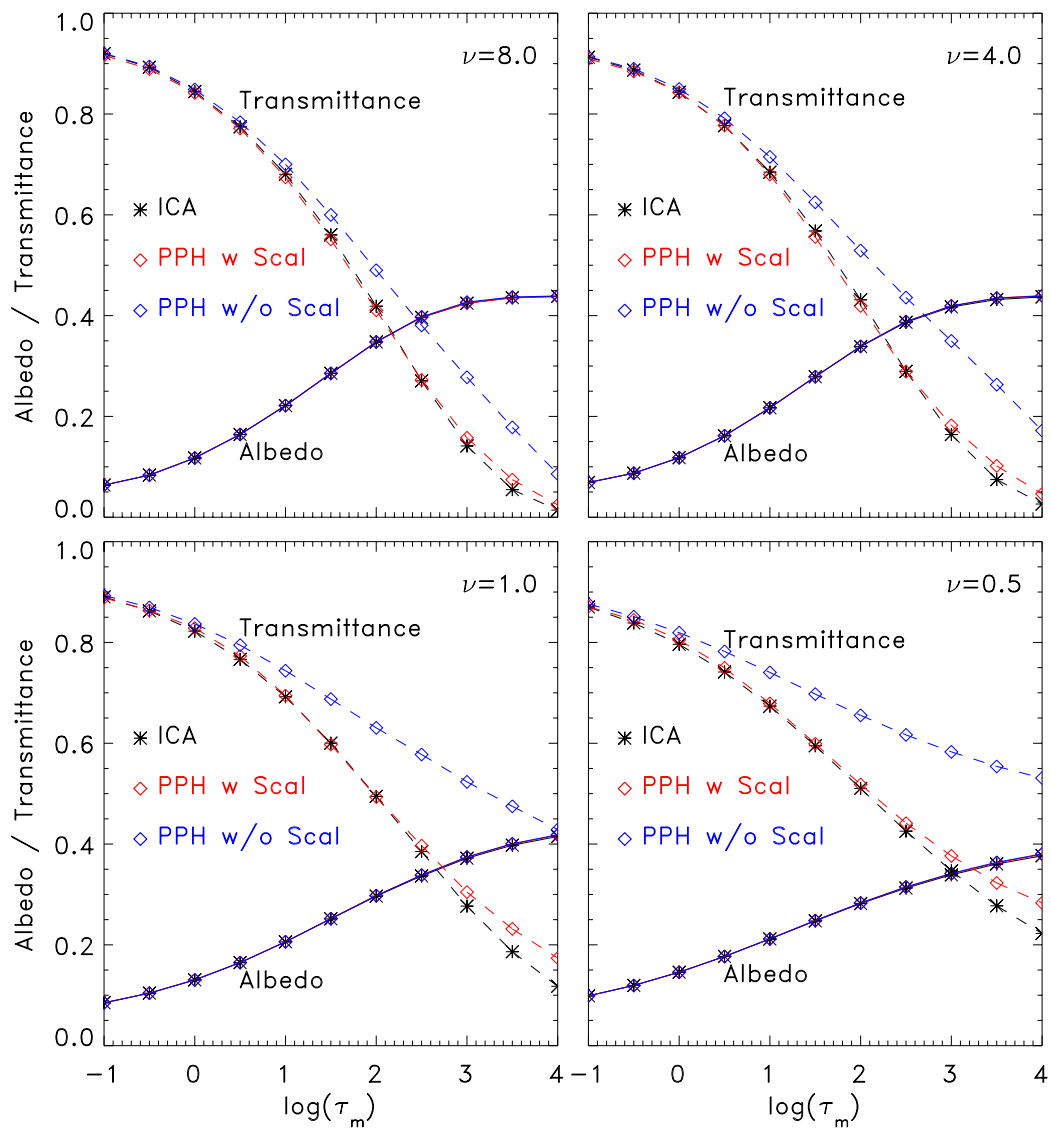
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 2 Fig. 4. Comparison of spherical albedo and transmission between the ICA approach (black) and the synthetic PPH approximation (red). A different panel represents a different cloud inhomogeneity. τ_m is the mean cloud optical depth.



1 Fig. 5. As in Figure 3 except this is for a nonconservative scattering with single scattering
 2 albedo of 0.98.



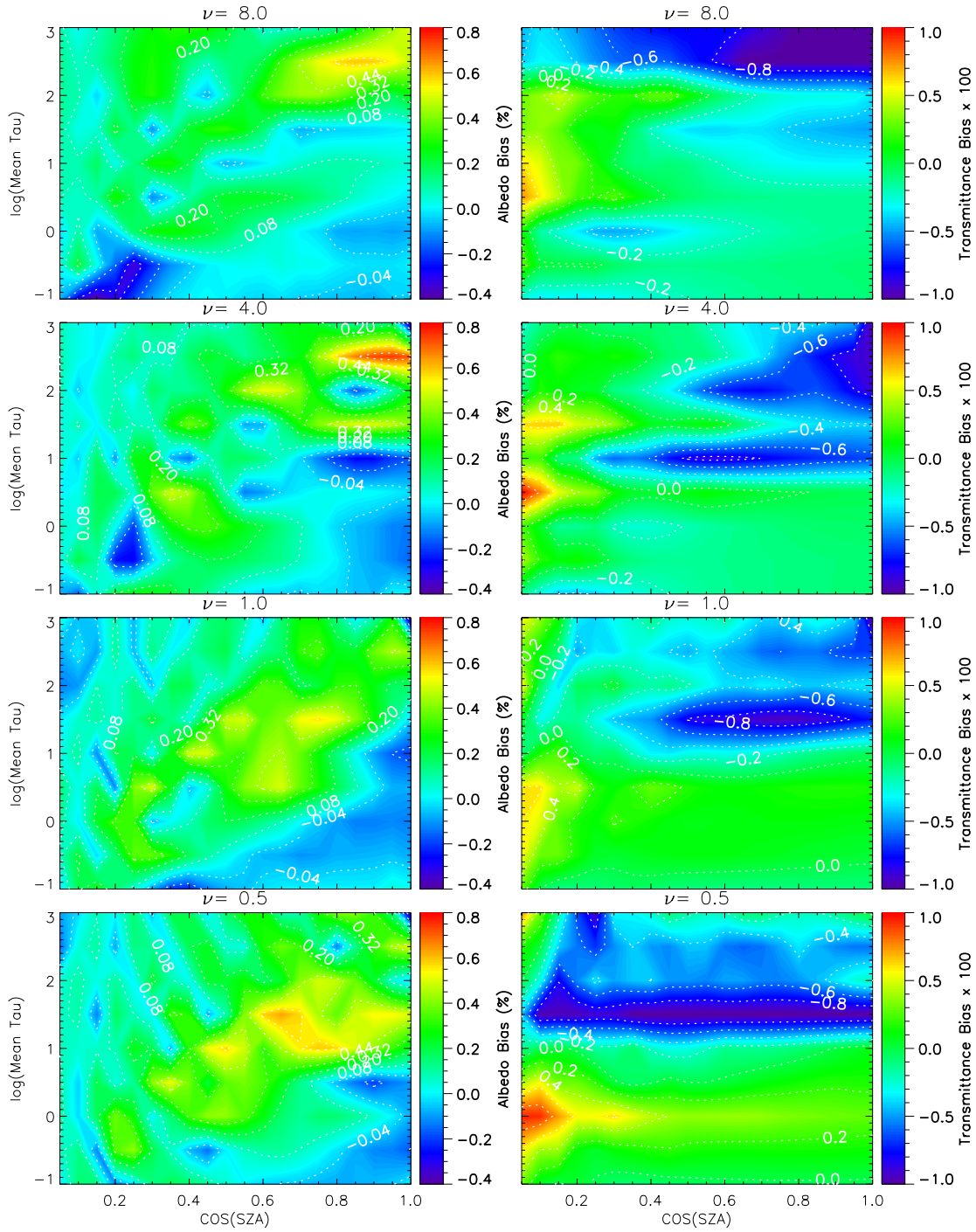
1 Fig. 6. As in Figure 5 except an adjustment of the effective single scattering albedo as in
 2 Equation (5) is applied to the PPH cloud.



1

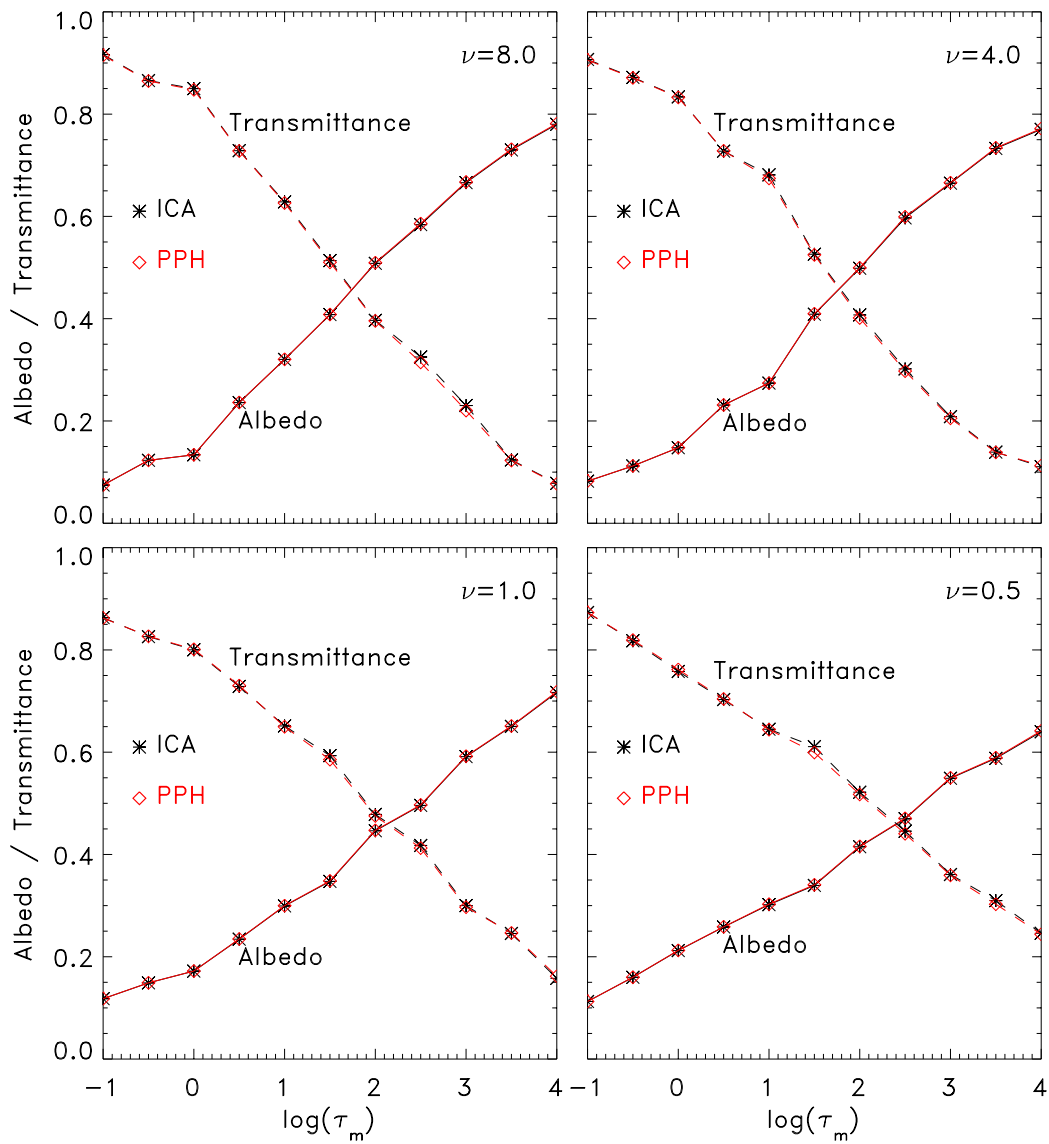
Fig. 7. Comparison of spherical albedo and transmission between the ICA approach (black), the synthetic PPH approximation (blue), and the synthetic PPH with scaling of the single scattering albedo (red). Cloud single scattering albedo is 0.98.

2



1 Fig. 8. As in Figure 6 but these are for solar radiation calculated with real cloud optical properties. The cloud particle effective radius is allowed to vary randomly from $7 \mu\text{m}$ to $80 \mu\text{m}$ among cloud subcolumns. Particle size less than $17 \mu\text{m}$ is assumed as water and as ice otherwise.

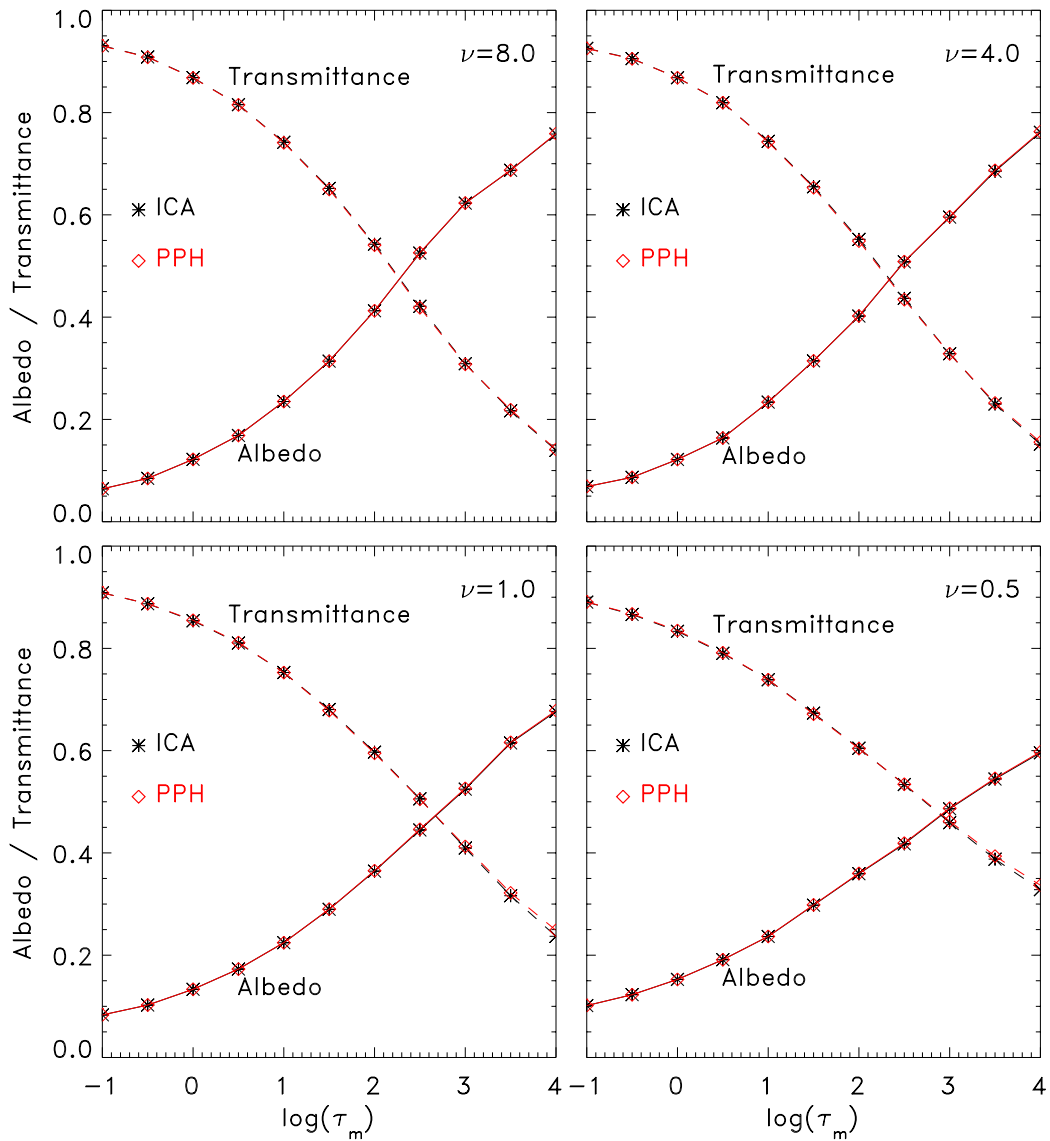
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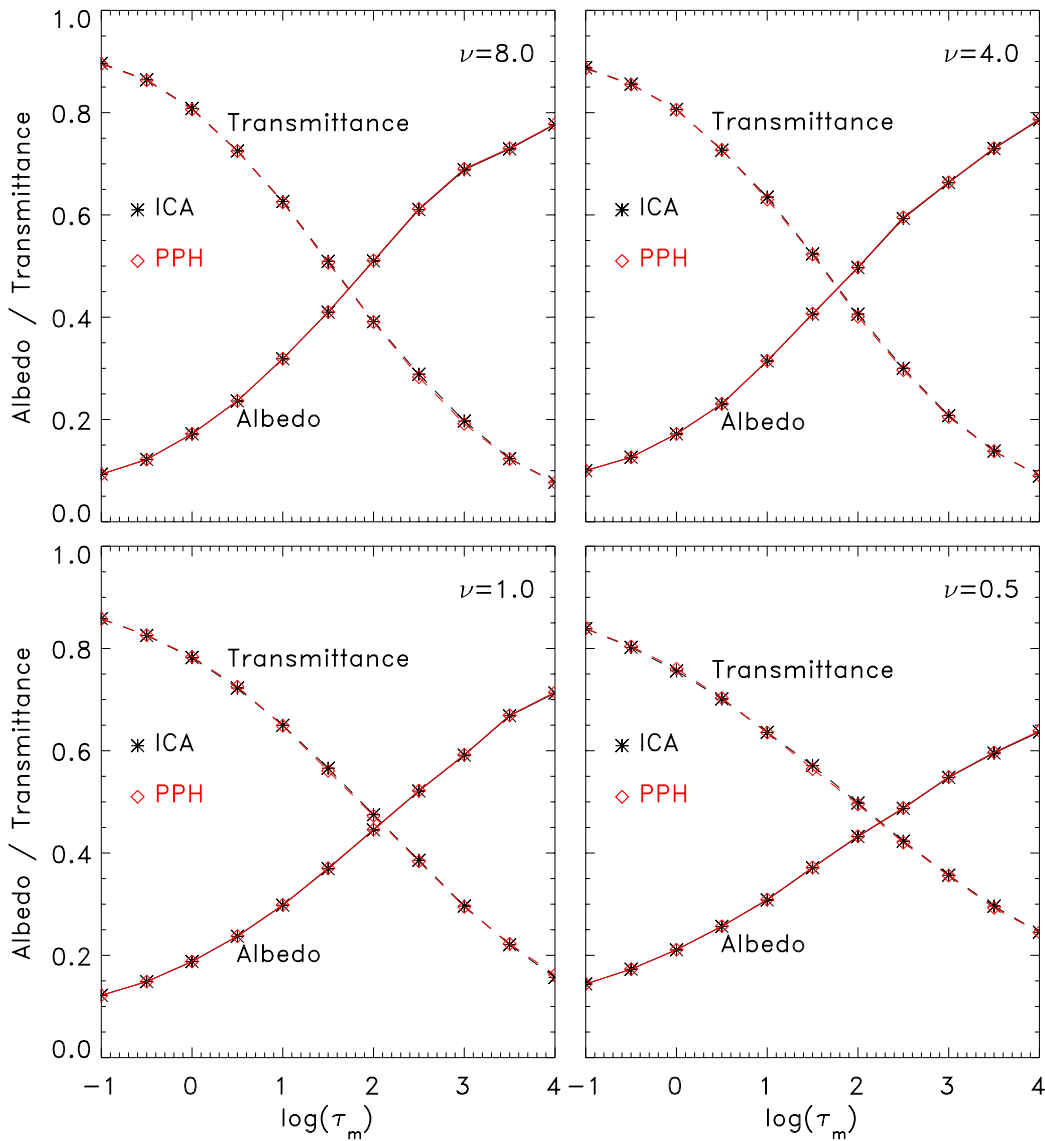
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Fig. 9. As in Figure 4 but for broadband shortwave albedo and transmission and the same clouds as in Figure 8.

2
3



1
 2 Fig. 10. As in Figure 9 except for water cloud only with particle size varies from 7 μm to 17 μm .



1

Fig. 11. As in Figure 9 except for ice cloud only with particle size varies from 17 μm to 80 μm .

2