

1 **Asymmetric Warming/Cooling Response to CO₂**
2 **Increase/Decrease Mainly Due to Non-Logarithmic**
3 **Forcing, not Feedbacks**

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8 **Key Points:**

- 9 • The global surface temperature responds asymmetrically to increased and decreased
10 CO₂ levels, in both abrupt and transient cases
- 11 • Effective climate sensitivity is higher with warming (2×, 4×, 8×CO₂) than with
12 cooling (1/2×, 1/4×, 1/8×CO₂), in two different coupled models
- 13 • The non-logarithmic nature of the CO₂ forcing is primarily responsible for the asym-
14 metry, not the radiative feedbacks

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15 **Abstract**

16 We explore the CO₂ dependence of effective climate sensitivity (S_G) with symmet-
 17 ric abrupt and transient CO₂ forcing, spanning the range $1/8\times$, $1/4\times$, $1/2\times$, $2\times$, $4\times$, and
 18 $8\times\text{CO}_2$, using two state-of-the-art fully coupled atmosphere-ocean-sea-ice-land models. In
 19 both models, under abrupt CO₂ forcing, we find an asymmetric response in surface tem-
 20 perature and S_G . The surface global warming at $8\times\text{CO}_2$ is more than one third larger
 21 than the corresponding cooling at $1/8\times\text{CO}_2$, and S_G is CO₂ dependent, increasing non-
 22 monotonically from $1/8\times\text{CO}_2$ to $8\times\text{CO}_2$. We find similar CO₂ dependence in the transient
 23 runs, forced with $-1\%\text{yr}^{-1}\text{CO}_2$ and $+1\%\text{yr}^{-1}\text{CO}_2$ up to $1/8\times\text{CO}_2$ and $8\times\text{CO}_2$, respectively.
 24 The non-logarithmic radiative forcing – not the changing feedbacks – primarily explains the
 25 dependence of S_G on CO₂, particularly at low CO₂ levels. The changing feedbacks, however,
 26 explain S_G ’s non-monotonic behavior.

27 **Plain Language Summary**

28 Equilibrium climate sensitivity (ECS) is the global mean warming after doubling CO₂
 29 concentrations from those of the year 1850. Since CO₂ levels will likely surpass a doubling,
 30 it is crucial to know whether the amount of warming per CO₂ doubling (which we refer to
 31 as the effective climate sensitivity, S_G) is constant with each CO₂ doubling or whether it
 32 changes. Necessary conditions for constant S_G are 1) the radiative forcing introduced to the
 33 climate system from each CO₂ doubling is constant and 2) the net radiative feedback does
 34 not change with CO₂ levels. Current literature shows that S_G will increase in a warmer
 35 world because the radiative feedback will change. We here investigate S_G in both warmer
 36 and colder worlds, and confirm that S_G increases at higher CO₂ concentrations. However,
 37 we show that changes in the radiative forcing with each CO₂ doubling are mainly responsible
 38 for S_G increase with CO₂, not feedback changes.

39 1 Introduction

40 Equilibrium climate sensitivity (ECS) is the global mean surface temperature change
41 after the doubling of CO₂ concentrations from pre-industrial (PI) levels. ECS is perhaps
42 the most important metric in climate science, and it has been extensively investigated in
43 the literature (Sherwood et al., 2020). An important question is whether the amount of
44 warming for each CO₂ doubling (which we refer to as the effective climate sensitivity, S_G)
45 is constant or not (i.e., whether it is CO₂ dependent). Necessary conditions for a constant
46 S_G are 1) that the radiative forcing of the climate system for each CO₂ doubling is constant
47 and 2) that the net radiative feedback does not change with CO₂ levels. This question has
48 been investigated in many modeling studies (Meraner et al., 2013; Mauritsen et al., 2019;
49 Sherwood et al., 2020; Bloch-Johnson et al., 2021), which have reported that S_G is indeed
50 CO₂ dependent. Most of these studies find that S_G increases at higher CO₂ levels and that
51 the change in feedbacks, not the change in CO₂ radiative forcing, is the primary driver of
52 S_G CO₂ dependence.

53 An alternative approach to using climate models to investigate the dependency of S_G on
54 CO₂ is to seek observational constraints from reconstructions of past climates. In particular,
55 most studies conclude that S_G inferred from paleoclimate records does depend on CO₂
56 (Caballero & Huber, 2013; Anagnostou et al., 2016; Shaffer et al., 2016; Friedrich et al.,
57 2016; Farnsworth et al., 2019; Zhu et al., 2019; Anagnostou et al., 2020), although a few
58 studies disagree (e.g., Martínez-Botí et al. (2015)). An ideal period to study the S_G from
59 past climate is the Last Glacial Maximum (LGM), approximately 21 kyr ago, when the
60 Earth was roughly 6K colder than PI conditions (Tierney et al., 2020). The LGM period
61 is of particular interest because the climate system was in a quasi-equilibrium state, the
62 climate forcings were large, and the surface temperature reconstructions are relatively well-
63 constrained (Zhu & Poulsen, 2021). However, when considering the LGM and other periods
64 in Earth's past, one needs to account for how the feedbacks in those past climate states differ

65 from the feedbacks operating in the modern state: hence the challenge in using paleoclimate-
66 based estimates to constrain S_G .

67 While modeling and paleoclimatic evidence suggest that S_G depends on CO₂, a systematic
68 exploration of the symmetry over a wide range of CO₂ forcing has yet to be performed.
69 The question thus remains: is the climate system response symmetric across a broad range of
70 positive (warm) and negative (cold) CO₂ forcings? The question of symmetry was examined
71 recently by Chalmers et al. (2022), who compared $1/2\times$ and $2\times$ CO₂ simulations performed
72 with the CESM1-CAM5 model, and found that global surface temperatures warm 20% more
73 than they cool. Roughly 50% of this asymmetry was shown to derive from an asymmetry
74 in CO₂ radiative forcing; the rest was associated with differences in feedbacks which, interest-
75 ingly, were found not to be related to clouds. Whether this result holds over a broader
76 range of CO₂ forcing, and whether it is model dependent remains an open question.

77 We here address these questions using a much broader range of both abrupt and tran-
78 sient CO₂ forcings, and do so with two different climate models. Specifically, CO₂ is varied
79 from $1/8\times$ to $8\times$ PI values, to test the CO₂ symmetry of the climate system response to
80 comparable increased and decreased CO₂. While we are not the first ones to perform such
81 symmetric CO₂ runs (Hansen et al., 2005; Colman & McAvaney, 2009; Russell et al., 2013;
82 Chalmers et al., 2022), here we explore 1) a larger CO₂ range than previously considered,
83 2) we do so using two different fully coupled climate models and, most importantly, 3) we
84 perform the experiments with both abrupt and transient CO₂ runs.

85 Overall we confirm the asymmetric response in surface temperature: the climate system
86 warms *more* with consecutive CO₂ doublings ($2\times$, $4\times$, and $8\times$ CO₂) than it cools with
87 consecutive CO₂ halvings ($1/2\times$, $1/4\times$, and $1/8\times$ CO₂). This asymmetry is also reflected in S_G ,
88 which *increases* at higher CO₂ concentrations, consistent with previous studies. Surprisingly,
89 we find that the non-logarithmic dependence of CO₂ radiative forcing (i.e., the fact that CO₂
90 radiative forcing increases more rapidly than the log of the CO₂ concentration) is primarily
91 responsible for this asymmetric response, and not the changes in radiative feedbacks.

92 **2 Methods**93 **2.1 Models Used**

94 We use two fully coupled atmosphere-ocean-sea-ice-land models: the large ensemble
 95 version of the Community Earth System Model (CESM-LE) and the NASA Goddard Insti-
 96 tute for Space Studies Model E2.1-G (GISS-E2.1-G). CESM-LE comprises the Community
 97 Atmosphere Model version 5 (CAM5, 30 vertical levels), and parallel ocean program version
 98 2 (POP2, 60 vertical levels) with approximately 1° horizontal resolution in all model com-
 99 ponents (Kay et al., 2015). GISS-E2.1-G is a 40-level atmospheric model with a resolution
 100 of $2^{\circ} \times 2.5^{\circ}$ latitude/longitude, coupled to a 1° horizontal resolution 40-level GISS Ocean
 101 v1 (GO1) (Kelley et al., 2020). This configuration of the GISS model contributed to the
 102 CMIP6 project under the label “GISS-E2-1-G”. We show CESM-LE results in the main text,
 103 and some GISS-E2.1-G results in supplementary information (SI) to corroborate CESM-LE
 104 findings.

105 **2.2 Abrupt $n \times \text{CO}_2$ Experiments**

106 We perform a series of abrupt CO_2 forcing runs using both models, subject to $1/8 \times$,
 107 $1/4 \times$, $1/2 \times$, $2 \times$, $4 \times$, and $8 \times \text{CO}_2$ forcings, with all other trace gases, ozone concentrations,
 108 aerosols, and other forcings fixed at PI values. Following CMIP6 protocol for $4 \times \text{CO}_2$ runs,
 109 we integrate all runs to 150 years starting from PI conditions. We contrast these to a PI
 110 control run to calculate the response.

111 For each model, we estimate the effective radiative forcing (ERF) with a companion
 112 series of CO_2 experiments, as per Forster et al. (2016), with prescribed PI sea surface
 113 temperatures (SSTs) and sea-ice concentrations (SICs). These experiments are 30-year-
 114 long. We calculate ERF as the difference between the global mean net top of the atmosphere
 115 (TOA) flux between PI and $n \times \text{CO}_2$ in these prescribed SSTs and SICs experiments. We
 116 do not here adjust for land warming simply because, in our ERF calculations, the surface

117 temperature response in the fixed SSTs and SICs simulations is minimal (Smith et al., 2020),
 118 but we have verified that the adjustment does not change our results (see Figure S3).

119 **2.3 Transient Experiments**

120 In addition to the abrupt CO₂ runs, we also perform transient CO₂ runs with the
 121 CESM-LE model. We start from PI conditions (same as in the abrupt CO₂ forcing), and
 122 we increase CO₂ at +1%yr⁻¹ for the “warm” case for 215 years (slightly above 8×CO₂) and
 123 -1%yr⁻¹ for the “cold” case for 215 years (slightly below 1/8×CO₂). We estimate transient
 124 effective radiative forcing as in the abrupt experiments, by running companion simulations
 125 with specified SSTs and SICs set to PI values (Forster et al., 2016), while ramping up CO₂
 126 at rates of +1%yr⁻¹ and -1%yr⁻¹. We contrast all variables to PI values to compute the
 127 response.

128 **2.4 Climate Sensitivity & Feedbacks**

We define effective climate sensitivity S_G as the x-intercept of the Gregory regression
 (Gregory et al., 2004) for each abrupt $n \times$ CO₂ run using the following equation:

$$S_G = \left| \frac{F_{y-\text{int}}(n \times \text{CO}_2)}{\lambda(n \times \text{CO}_2) \cdot \log_2 n} \right| \quad (1)$$

129 We find the radiative forcing $F_{y-\text{int}}$ as the y-intercept and the net feedback parameter λ
 130 as the slope from the Gregory regression (see Figure S1) where we regress the net TOA
 131 radiative imbalance against the global mean surface temperature response for years 1-150.
 132 In order to compare S_G for different CO₂ doubling / halving, we divide by $\log_2 n$ (assuming a
 133 logarithmic CO₂ forcing) and take the absolute value in Equation 1. Note that our definition
 134 of the effective climate sensitivity S_G is a generalization of the more common definition of
 135 effective climate sensitivity (which is typically defined as per Equation 1 but with $n = 2$).
 136 To check for the possibility that λ and S_G may be strongly affected by the “pattern effect”,
 137 we have repeated the calculations by regressing years 21-150 only, and our main results were
 138 not changed.

To calculate the individual feedbacks λ_i , we use radiative kernels (K_x) from both Pendergrass et al. (2018) and Huang et al. (2017) to quantify the sensitivity of TOA radiation imbalance (ΔR) to changes in surface and atmospheric temperature (T), water vapor (q), and surface albedo (α) (Soden et al., 2008; Shell et al., 2008). For each year of the 150-year experiment, we multiply the spatially-resolved kernels by the climate field anomalies ($R_x = K_x \cdot \Delta x$, where x is T, q, α), and then vertically integrate (for atmospheric temperature and water vapor) up to the tropopause. We define the tropopause as 100 hPa at the equator, 300 hPa at the poles, and in between, it varies by the cosine of the latitude (Soden & Held, 2006). Lastly, we regress these quantities on the surface temperature response to find the radiative feedbacks as the regression slope. The cloud feedbacks are computed via the residual method (Soden & Held, 2006) as follows. First, we subtract effective radiative forcing and the temperature, water vapor, and surface albedo radiative fluxes from the TOA net radiative flux, resulting in $\Delta R_{\text{cloud}} = \Delta R - \text{ERF} - \sum \Delta R_x$. Then, we regress ΔR_{cloud} onto ΔT_s anomalies and define the corresponding slope as the cloud feedback. Lastly, we find shortwave (SW) and longwave (LW) components of the cloud feedback by considering the radiative changes in LW and SW components separately.

In the transient runs, we estimate the net feedback parameter λ_{tr} following Rugenstein and Armour (2021) (see λ_{eff1pct} in their Figure 1d) with the expression:

$$\lambda_{\text{tr}} = -\frac{\text{ERF}(t) - \Delta R(t)}{\Delta T_s(t)} \quad (2)$$

$\Delta R(t)$ is the net TOA radiative imbalance, and $\Delta T_s(t)$ is the global mean surface temperature response in the transient runs at year t . $\Delta R(t)$ and $\Delta T_s(t)$ are 30-year moving averages of the respective terms. Note that we use different definitions for the feedback parameter in the abrupt and transient simulations.

159 **3 Results**160 **3.1 Abrupt CO₂ Experiments**

161 We start by examining the global mean surface temperature response ($|\Delta T_s|$) timeseries
 162 for the abrupt CO₂ runs (Figure 1). We contrast – in panels a, b, and c – the timeseries
 163 of each corresponding “warm” ($2\times$, $4\times$, and $8\times$ CO₂) and “cold” simulation ($1/2\times$, $1/4\times$,
 164 and $1/8\times$ CO₂) by taking the absolute value of the response from PI: note that the $|\Delta T_s|$
 165 in the “warm” case is always stronger than the “cold” case. In particular, we find 20%
 166 more warming at $2\times$ than cooling at $1/2\times$ CO₂ (Figure 1a), 15% more at $4\times$ than $1/4\times$ CO₂
 167 (Figure 1b), and 41% more at $8\times$ than $1/8\times$ CO₂ (Figure 1c). The asymmetry in $|\Delta T_s|$ is
 168 amplified at higher CO₂ forcing, and largest in the $1/8\times$ CO₂ vs. $8\times$ CO₂ case (Figure 1c).
 169 The asymmetry is reduced at $4\times$ CO₂ vs. $1/4\times$ CO₂ due to changes in ocean heat transport
 170 which result in a formation of the North Atlantic Warming Hole in this model at $4\times$ CO₂
 171 (see more details in Mitevski et al. (2021)).

To quantify the timescale of the asymmetry in $|\Delta T_s|$ between “warm” and “cold” cases,
 we define the asymmetry between “warm” and “cold” cases as

$$\Delta_a X = |\Delta X(\text{warm})| - |\Delta X(\text{cold})| \quad (3)$$

172 where X is any climate variable (e.g., T_s), and subscript a refers to “asymmetry” (Figure 1d).
 173 In particular, we find that the asymmetry emerges rapidly in the first ten years (e.g., 90%
 174 at $8\times$ CO₂). Relative to the (slower) response associated with SST-driven feedbacks, the
 175 asymmetry appears quickly, suggesting that it might be due to radiative changes.

176 Next, we calculate effective climate sensitivity S_G from the Gregory regression (Equa-
 177 tion 1), and plot it as percentage change from $2\times$ CO₂ (black line, Figure 2a). S_G is CO₂
 178 dependent and increases with CO₂ concentration: at $1/8\times$ CO₂, it is more than 20% *lower*
 179 than $2\times$ CO₂ values, and at $8\times$ CO₂, it is around 5% *higher* than at $2\times$ CO₂. CO₂ dependent
 180 S_G is possible if either the effective radiative forcing (ERF) or the net feedback parame-

181 ter (λ) change with CO₂. To individually test the relative importance of ERF and λ , we
 182 calculate the climate sensitivity in two different ways.

First, to examine the dependence of climate sensitivity on ERF, we calculate climate sensitivity as S_F using the expression:

$$S_F = \left| \frac{\text{ERF}(n \times \text{CO}_2)}{\lambda(2 \times \text{CO}_2) \cdot \log_2 n} \right| \quad (4)$$

183 where ERF is derived from the $n \times \text{CO}_2$ fixed SSTs and SICs runs, and λ (slope from Gregory
 184 Regression) is held constant at the $2 \times \text{CO}_2$ value. As seen in Figure 2a, we find that S_F
 185 (blue line) changes in tandem with S_G (black line), which reinforces the fact that changes
 186 in ERF explain the changes in S_G .

Second, to assess whether changes in feedback strength also contribute to S_G , we calculate climate sensitivity as S_λ :

$$S_\lambda = \left| \frac{\text{ERF}(2 \times \text{CO}_2)}{\lambda(n \times \text{CO}_2)} \right| \quad (5)$$

187 where λ is calculated at each $n \times \text{CO}_2$ and ERF is held constant at $2 \times \text{CO}_2$ value. As seen in
 188 Figure 2a, S_λ (red) changes in the opposite direction than S_G (black) for CO₂ values lower
 189 than $2 \times \text{CO}_2$. This suggests that changes in λ are not the main driver of the S_G dependence
 190 on CO₂. However, it is important to note that for CO₂ values higher than $2 \times \text{CO}_2$, we
 191 find λ non-monotonically increasing to $8 \times \text{CO}_2$, which can be linked to the corresponding
 192 non-monotonic behavior of S_G . We find qualitatively similar results using the GISS-E2.1-G
 193 model (Figure S2a), confirming that ERF is the primary driver of the dependence of S_G on
 194 CO₂.

195 Next, we correlate S_G with $1/\lambda$ (Figure 2c) and ERF (Figure 2d) across all abrupt CO₂
 196 experiments from $1/8 \times$ to $8 \times \text{CO}_2$ to examine whether feedbacks or forcing better correlate
 197 with changes in S_G . Overall, we find little correlation between S_G and $1/\lambda$ ($r=-0.44$) and
 198 a very strong correlation between S_G and ERF ($r=0.91$). Similarly, a high correlation
 199 between S_G and ERF is found in the GISS-E2.1-G model (Figure S2d). This strengthens our
 200 conclusions from Figure 2a that the changes in ERF are driving the S_G increase. However,

201 if one considers warm cases, one sees a strong correlation between S_G and $1/\lambda$, as indicated
202 earlier. This is in agreement with previous studies (Meraner et al., 2013; Bloch-Johnson et
203 al., 2021), which reported that feedback changes are important for the dependence of S_G
204 on CO₂. However, over a broad range of CO₂ forcing, including colder climates, that is not
205 the case: changes in ERF are more important than feedback changes.

206 Given the aforementioned importance of ERF in driving the changes in S_G , we next
207 look in more detail at ERF, calculated from fixed SSTs and SICs runs, following Forster et
208 al. (2016), from $1/8\times$ to $8\times$ CO₂ (dark blue bars, Figure 2b). If ERF were scaled simply
209 with the logarithm of CO₂ concentration, then the dark blue bars would be identical for
210 all CO₂ values. However, we see that ERF grows more than logarithmically with CO₂. We
211 find a similar but weaker non-logarithmic behavior in the instantaneous radiative forcing
212 (IRF) reported in Byrne and Goldblatt (2014), which we obtain by linearly interpolating
213 their line-by-line radiative calculations (SI file “text03.txt” in Byrne and Goldblatt (2014))
214 and plot with light blue bars in Figure 2b. We also compare our ERF calculations with
215 the proposed stratospherically adjusted radiative forcing fit in Etminan et al. (2016) for the
216 warming case only (since it is not valid for low CO₂ values), and it appears both are in
217 agreement.

218 A limitation to our ERF calculation approach is that we only fix the SSTs and SICs
219 in the simulation, but not the land temperatures. Fixing the land temperatures has been
220 shown to increase ERF in warmer climates even more than when only SSTs and SICs are
221 fixed (Andrews et al., 2021). To account for this, we removed the land and sea-ice warming
222 effects in our ERF calculations, following Equation 1 in Hansen et al. (2005) as shown in
223 Figure S3, and found that the correction (dashed blue lines) leads, if anything, to a stronger
224 non-logarithmic ERF. Hence, incorporating fixed land temperatures leads to ERF increasing
225 even more rapidly than the log of CO₂ concentration; this strengthens our argument that
226 the S_G dependence on CO₂ is due to non-logarithmic CO₂ radiative forcing.

227 Next, we perform a standard decomposition of λ into individual radiative feedbacks λ_i .
 228 The summation of individual feedbacks ($\sum \lambda_i$) is shown in Figure 3a (blue). $\sum \lambda_i$ follows
 229 closely the net feedback calculated from the Gregory regression (black). We perform the
 230 decomposition using two radiative kernels from Pendergrass et al. (2018) and Huang et al.
 231 (2017), and we find minimal sensitivity to the choice of kernel (Figure S4). The individual
 232 feedbacks, plotted as differences from $2\times\text{CO}_2$ values, from the Pendergrass et al. (2018)
 233 kernels are shown in Figure 3b. We see a clear signal in the lapse rate feedback, which
 234 weakens the net feedback in the “cold” case and strengthens it in the “warm” case. The
 235 longwave cloud feedback has clear global surface temperature dependence, increasing with
 236 CO_2 monotonically for all CO_2 values. However, in general, we find no clear pattern in the
 237 changes in individual feedbacks that would sufficiently explain the overall feedbacks CO_2
 238 dependence. In addition, the changes in feedbacks in the GISS-E2.1-G model (Figure S5)
 239 are qualitatively different from those in the CESM-LE model (Figure 3). Since our models
 240 do not agree on the changes in individual feedbacks across the CO_2 range, and since we
 241 showed that feedback changes are strongly not correlated with changes in S_G (Figure 2c),
 242 we do not explore further the mechanisms driving feedback changes in the individual models.

243 3.2 Transient CO₂ runs

244 The abrupt CO₂ forcing runs show that the effective climate sensitivity increases with
 245 CO₂, and that the non-logarithmic nature of the ERF is largely responsible for this behavior.
 246 Now we seek to determine whether the same behavior is also seen in runs with transient
 247 CO₂ forcing, which are much more realistic. Our transient runs are forced, starting from
 248 PI, with CO₂ concentrations increasing at the rate of $1\%\text{yr}^{-1}$ and decreasing at $1\%\text{yr}^{-1}$. As
 249 seen in Figure 4a, the surface temperature response $|\Delta T_s|$ is stronger in the warming (red)
 250 than in the cooling (blue) case. Note that the responses computed from the last 50 years
 251 of the abrupt simulations at the corresponding CO₂ value (dots) are a good predictor of
 252 the response in the transient runs, demonstrating that the results of the abrupt runs carry

253 over to the transient runs. Together with the surface temperature, ERF also changes more
 254 rapidly in the warming than the cooling experiments, as seen in Figure 4b.

255 Next, we explore how the transient feedbacks (λ_{tr} , see Equation 2) change in the “warm”
 256 and “cold” cases (Figure 4c). The feedbacks timeseries are noisy at the beginning of the
 257 simulation, but in the last thirty years, the warm case shows 10% weaker (more positive)
 258 feedbacks compared to the cold case. The 10% difference indicates that S_G in the warming
 259 case should be higher than in the colder case. However, a robust difference in feedbacks
 260 only appears around year 130, whereas the $|\Delta T_s|$ asymmetry emerges much earlier, around
 261 year 60. This difference in the temporal evolution of the feedbacks, relative to the evolution
 262 of the forcing and S_G , adds additional strong evidence that the feedbacks are not driving
 263 the $|\Delta T_s|$ asymmetry.

264 Finally, as for the abrupt CO₂ runs, we correlate the asymmetry in global mean surface
 265 temperature response $\Delta_a T_s$ and effective radiative forcing $\Delta_a \text{ERF}$ (Figure 4d). We find
 266 a correlation of $r=0.96$, suggesting that the asymmetric changes in ERF drive the $|\Delta T_s|$
 267 asymmetry between the “cold” and “warm” cases. As we can see in Figure 4c, the transient
 268 feedbacks are contributing to the $|\Delta T_s|$ asymmetry at the end of the run, but their impact
 269 is much smaller than the one from ERF.

270 4 Summary and Discussion

271 We have explored the effective climate sensitivity (S_G) dependence on CO₂ with abrupt
 272 and transient CO₂ experiments spanning the range $1/8\times$ to $8\times$ CO₂ using two distinct CMIP-
 273 class climate models. First, we have found a considerable asymmetry in surface temperature
 274 response, with the climate system warming more than cooling for identical factors used to
 275 increase and decrease the CO₂ concentration, starting from a pre-industrial climate. Second,
 276 we showed that the asymmetry is due to the non-logarithmic nature of CO₂ radiative forcing,
 277 not the feedback changes. Upon decomposing the total feedback into individual feedbacks,

278 we found no simple explanation relating specific feedback changes to the changes in S_G
279 across the $1/8\times$ to $8\times\text{CO}_2$ forcing range examined in this study.

280 Most studies to date have focused on the role of feedbacks in explaining the dependency
281 of S_G on CO_2 , with relatively little attention placed on radiative forcing. Indeed, consistent
282 with these studies, we found that for warmer climates ($> 2 \times \text{CO}_2$), feedbacks are important
283 for determining the changing behavior of S_G with CO_2 . However, by considering a broader
284 range of CO_2 forcings, we have shown here that for cases in which CO_2 concentrations are
285 less than PI values, non-logarithmic ERF is the primary driver of S_G changes. Our goal
286 here has been to isolate the role of CO_2 alone, and we have set all other forcings to PI
287 values. Needless to say, we have ignored the “slow” feedbacks present in cold climates (e.g.,
288 the LGM), such as the formation of land ice sheets.

289 The results with our abrupt runs have been shown to be robust with two climate
290 models for simulations up to 150 years. One may argue that our runs are not equilibrated,
291 and we agree with that caveat. However, we have found that the asymmetry and the key
292 role of ERF are also robustly seen in the transient runs. Because of this, we expect that
293 prolonging the abrupt simulation for more than 150 years will yield similar results. In
294 any case, it will be important to repeat similar experiments with longer simulations as in
295 LongRunMIP (Rugenstein et al., 2019) to confirm that this asymmetry is still present at
296 long times closer to equilibration. Finally, our findings indicate that future studies should
297 place more emphasis on accurately quantifying the changes in effective radiative forcing
298 when studying the effective climate sensitivity dependency on CO_2 . The feedbacks appear
299 unable to explain the cooling phase.

300 5 Open Research

301 The CESM-LE model data can be obtained at [https://doi.org/10.5281/zenodo](https://doi.org/10.5281/zenodo.5725084)
302 .5725084 and GISS-E2.1-G model data at <https://doi.org/10.5281/zenodo.3901624>.

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315 **References**

- 316 Anagnostou, E., John, E. H., Babila, T. L., Sexton, P. F., Ridgwell, A., Lunt, D. J., ...
 317 Foster, G. L. (2020, Sep 07). Proxy evidence for state-dependence of climate sensitivity
 318 in the eocene greenhouse. *Nature Communications*, 11(1), 4436. doi: <https://doi.org/10.1038/s41467-020-17887-x>
 319
- 320 Anagnostou, E., John, E. H., Edgar, K. M., Foster, G. L., Ridgwell, A., Inglis, G. N., ...
 321 Pearson, P. N. (2016, May 01). Changing atmospheric co₂ concentration was the
 322 primary driver of early cenozoic climate. *Nature*, 533(7603), 380-384. doi: <https://doi.org/10.1038/nature17423>
 323
- 324 Andrews, T., Smith, C. J., Myhre, G., Forster, P. M., Chadwick, R., & Ackerley, D.
 325 (2021). Effective radiative forcing in a gcm with fixed surface temperatures. *Journal*
 326 *of Geophysical Research: Atmospheres*, 126(4), e2020JD033880. doi: <https://doi.org/10.1029/2020JD033880>
 327
- 328 Bloch-Johnson, J., Rugenstein, M., Stolpe, M. B., Rohrschneider, T., Zheng, Y., & Gregory,
 329 J. M. (2021). Climate sensitivity increases under higher co₂ levels due to feedback

- 330 temperature dependence. *Geophysical Research Letters*, 48(4), e2020GL089074. doi:
331 <https://doi.org/10.1029/2020GL089074>
- 332 Byrne, B., & Goldblatt, C. (2014). Radiative forcing at high concentrations of well-mixed
333 greenhouse gases. *Geophysical Research Letters*, 41(1), 152-160. doi: <https://doi.org/10.1002/2013GL058456>
- 335 Caballero, R., & Huber, M. (2013). State-dependent climate sensitivity in past warm
336 climates and its implications for future climate projections. *Proceedings of the Na-*
337 *tional Academy of Sciences*, 110(35), 14162–14167. doi: <https://doi.org/10.1073/pnas.1303365110>
- 339 Chalmers, J., Kay, J. E., Middlemas, E., Maroon, E., & DiNezio, P. (2022). Does disabling
340 cloud radiative feedbacks change spatial patterns of surface greenhouse warming and
341 cooling? *Journal of Climate*. doi: <https://doi.org/10.1175/JCLI-D-21-0391.1>
- 342 Colman, R., & McAvaney, B. (2009). Climate feedbacks under a very broad range of forcing.
343 *Geophysical Research Letters*, 36(1). doi: <https://doi.org/10.1029/2008GL036268>
- 344 Etminan, M., Myhre, G., Highwood, E. J., & Shine, K. P. (2016). Radiative forcing of
345 carbon dioxide, methane, and nitrous oxide: A significant revision of the methane
346 radiative forcing. *Geophysical Research Letters*, 43(24), 12,614-12,623. doi: <https://doi.org/10.1002/2016GL071930>
- 348 Farnsworth, A., Lunt, D. J., O'Brien, C. L., Foster, G. L., Inglis, G. N., Markwick, P.,
349 ... Robinson, S. A. (2019). Climate sensitivity on geological timescales controlled
350 by nonlinear feedbacks and ocean circulation. *Geophysical Research Letters*, 46(16),
351 9880-9889. doi: <https://doi.org/10.1029/2019GL083574>
- 352 Forster, P., Richardson, T., Maycock, A. C., Smith, C. J., Samset, B. H., Myhre, G., ...
353 Schulz, M. (2016). Recommendations for diagnosing effective radiative forcing from
354 climate models for cmip6. *Journal of Geophysical Research: Atmospheres*, 121(20),
355 12,460-12,475. doi: <https://doi.org/10.1002/2016JD025320>
- 356 Friedrich, T., Timmermann, A., Tigchelaar, M., Elison Timm, O., & Ganopolski, A. (2016).
357 Nonlinear climate sensitivity and its implications for future greenhouse warming. *Sci-*

- 358 ence *Advances*, 2(11). doi: <https://doi.org/10.1126/sciadv.1501923>
- 359 Gregory, J. M., Ingram, W. J., Palmer, M. A., Jones, G. S., Stott, P. A., Thorpe, R. B.,
360 ... Williams, K. D. (2004). A new method for diagnosing radiative forcing and
361 climate sensitivity. *Geophysical Research Letters*, 31(3). doi: <https://doi.org/10.1029/2003GL018747>
- 362
- 363 Hansen, J., Sato, M., Ruedy, R., Nazarenko, L., Lacis, A., Schmidt, G. A., ... Zhang, S.
364 (2005). Efficacy of climate forcings. *Journal of Geophysical Research: Atmospheres*,
365 110(D18). doi: <https://doi.org/10.1029/2005JD005776>
- 366 Huang, Y., Xia, Y., & Tan, X. (2017). On the pattern of co2 radiative forcing and poleward
367 energy transport. *Journal of Geophysical Research: Atmospheres*, 122(20), 10,578-
368 10,593. doi: <https://doi.org/10.1002/2017JD027221>
- 369 Kay, J. E., Deser, C., Phillips, A., Mai, A., Hannay, C., Strand, G., ... Vertenstein, M.
370 (2015, 09). The Community Earth System Model (CESM) Large Ensemble Project: A
371 Community Resource for Studying Climate Change in the Presence of Internal Climate
372 Variability. *Bulletin of the American Meteorological Society*, 96(8), 1333-1349. doi:
373 <https://doi.org/10.1175/BAMS-D-13-00255.1>
- 374 Kelley, M., Schmidt, G. A., Nazarenko, L. S., Bauer, S. E., Ruedy, R., Russell, G. L., ...
375 Yao, M.-S. (2020). Giss-e2.1: Configurations and climatology. *Journal of Advances*
376 in Modeling Earth Systems, 12(8), e2019MS002025. doi: <https://doi.org/10.1029/2019MS002025>
- 377
- 378 Martínez-Botí, M. A., Foster, G. L., Chalk, T. B., Rohling, E. J., Sexton, P. F., Lunt,
379 D. J., ... Schmidt, D. N. (2015). Plio-pleistocene climate sensitivity evaluated using
380 high-resolution CO2 records. *Nature*, 518(7537), 49-54. doi: <https://doi.org/10.1038/nature14145>
- 381
- 382 Mauritsen, T., Bader, J., Becker, T., Behrens, J., Bittner, M., Brokopf, R., ... Roeckner,
383 E. (2019). Developments in the mpi-m earth system model version 1.2 (mpi-esm1.2)
384 and its response to increasing co2. *Journal of Advances in Modeling Earth Systems*,
385 11(4), 998-1038. doi: <https://doi.org/10.1029/2018MS001400>

- 386 Meraner, K., Mauritsen, T., & Voigt, A. (2013). robust increase in equilibrium climate
387 sensitivity under global warming. *geophysical research letters*, 40(22), 5944-5948. doi:
388 <https://doi.org/10.1002/2013gl058118>
- 389 Mitevski, I., Orbe, C., Chemke, R., Nazarenko, L., & Polvani, L. M. (2021). Non-monotonic
390 response of the climate system to abrupt co2 forcing. *Geophysical Research Letters*,
391 48(6), e2020GL090861. doi: <https://doi.org/10.1029/2020GL090861>
- 392 Pendergrass, A. G., Conley, A., & Vitt, F. M. (2018). Surface and top-of-atmosphere
393 radiative feedback kernels for cesm-cam5. *Earth System Science Data*, 10(1), 317–
394 324. doi: <https://doi.org/10.5194/essd-10-317-2018>
- 395 Rugenstein, M., & Armour, K. C. (2021). Three flavors of radiative feedbacks and their im-
396 plications for estimating equilibrium climate sensitivity. *Geophysical Research Letters*,
397 48(15), e2021GL092983. doi: <https://doi.org/10.1029/2021GL092983>
- 398 Rugenstein, M., Bloch-Johnson, J., Abe-Ouchi, A., Andrews, T., Beyerle, U., Cao, L.,
399 ... Yang, S. (2019). Longrunmip: Motivation and design for a large collection of
400 millennial-length aogcm simulations. *Bulletin of the American Meteorological Society*,
401 100(12), 2551 - 2570. doi: <https://doi.org/10.1175/BAMS-D-19-0068.1>
- 402 Russell, G. L., Lacis, A. A., Rind, D. H., Colose, C., & Opstbaum, R. F. (2013). Fast
403 atmosphere-ocean model runs with large changes in co2. *Geophysical Research Letters*,
404 40(21), 5787-5792. doi: <https://doi.org/10.1002/2013GL056755>
- 405 Shaffer, G., Huber, M., Rondanelli, R., & Pepke Pedersen, J. O. (2016). Deep time evidence
406 for climate sensitivity increase with warming. *Geophysical Research Letters*, 43(12),
407 6538-6545. doi: <https://doi.org/10.1002/2016GL069243>
- 408 Shell, K. M., Kiehl, J. T., & Shields, C. A. (2008). Using the radiative kernel technique
409 to calculate climate feedbacks in ncar's community atmospheric model. *Journal of*
410 *Climate*, 21(10), 2269 - 2282. doi: <https://doi.org/10.1175/2007JCLI2044.1>
- 411 Sherwood, S. C., Webb, M. J., Annan, J. D., Armour, K. C., Forster, P. M., Hargreaves,
412 J. C., ... Zelinka, M. D. (2020). An assessment of earth's climate sensitivity using
413 multiple lines of evidence. *Reviews of Geophysics*. doi: <https://doi.org/10.1029/>

414 2019RG000678

415 Smith, C. J., Kramer, R. J., Myhre, G., Alterskjær, K., Collins, W., Sima, A., ... Forster,
416 P. M. (2020). Effective radiative forcing and adjustments in cmip6 models. *At-*

417 *mospheric Chemistry and Physics*, 20(16), 9591–9618. doi: [https://doi.org/10.5194/](https://doi.org/10.5194/acp-20-9591-2020)

418 acp-20-9591-2020

419 Soden, B. J., & Held, I. M. (2006). An assessment of climate feedbacks in coupled
420 ocean–atmosphere models. *Journal of Climate*, 19(14), 3354 - 3360. doi: <https://doi.org/10.1175/JCLI3799.1>

421 Soden, B. J., Held, I. M., Colman, R., Shell, K. M., Kiehl, J. T., & Shields, C. A. (2008).
422 Quantifying climate feedbacks using radiative kernels. *Journal of Climate*, 21(14),
423 3504 - 3520. doi: <https://doi.org/10.1175/2007JCLI2110.1>

424 Tierney, J. E., Zhu, J., King, J., Malevich, S. B., Hakim, G. J., & Poulsen, C. J. (2020,
425 Aug 01). Glacial cooling and climate sensitivity revisited. *Nature*, 584(7822), 569–573.
426 doi: <https://doi.org/10.1038/s41586-020-2617-x>

427 Zhu, J., & Poulsen, C. J. (2021). Last glacial maximum (lgm) climate forcing and ocean
428 dynamical feedback and their implications for estimating climate sensitivity. *Climate
429 of the Past*, 17(1), 253–267. Retrieved from <https://cp.copernicus.org/articles/17/253/2021/> doi: <https://doi.org/10.5194/cp-17-253-2021>

430 431 Zhu, J., Poulsen, C. J., & Tierney, J. E. (2019). Simulation of eocene extreme warmth
432 and high climate sensitivity through cloud feedbacks. *Science Advances*, 5(9). doi:
433 <https://doi.org/10.1126/sciadv.aax1874>

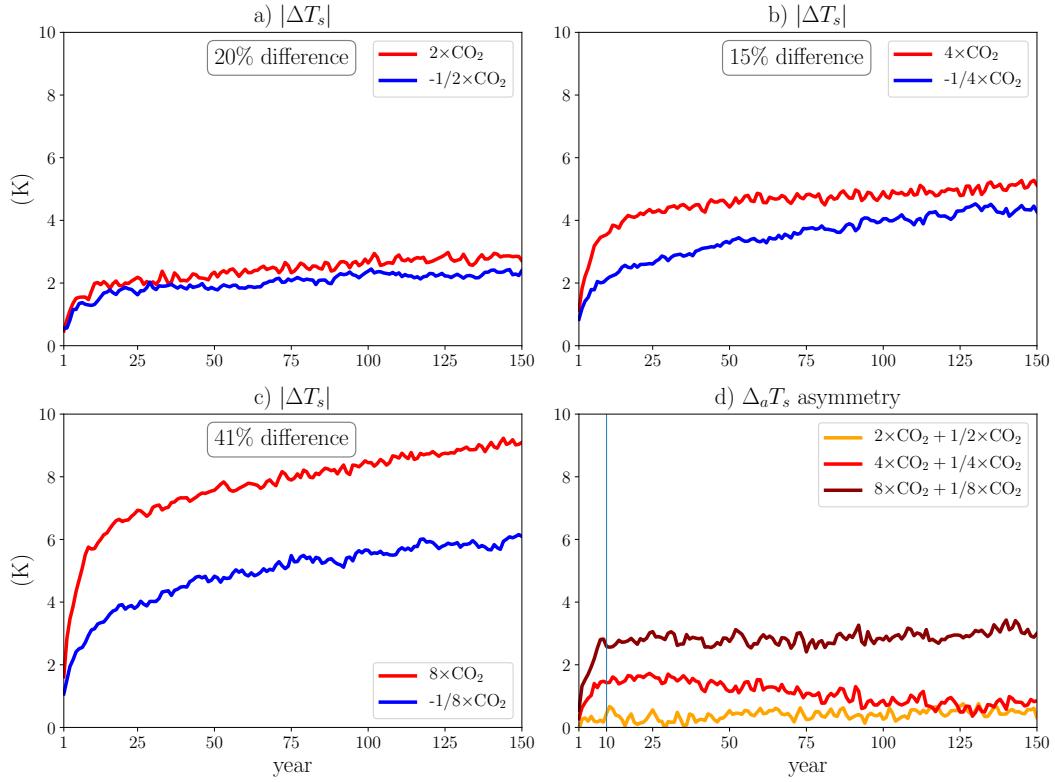


Figure 1. Timeseries of surface temperature response ($|\Delta T_s|$) for abrupt CO₂ runs with CESM-LE model. a) $2\times\text{CO}_2$ and $1/2\times\text{CO}_2$, b) $4\times\text{CO}_2$ and $1/4\times\text{CO}_2$, c) $8\times\text{CO}_2$ and $1/8\times\text{CO}_2$ runs, and d) surface temperature asymmetry ($\Delta_a T_s$) between “warm” and “cold” cases.

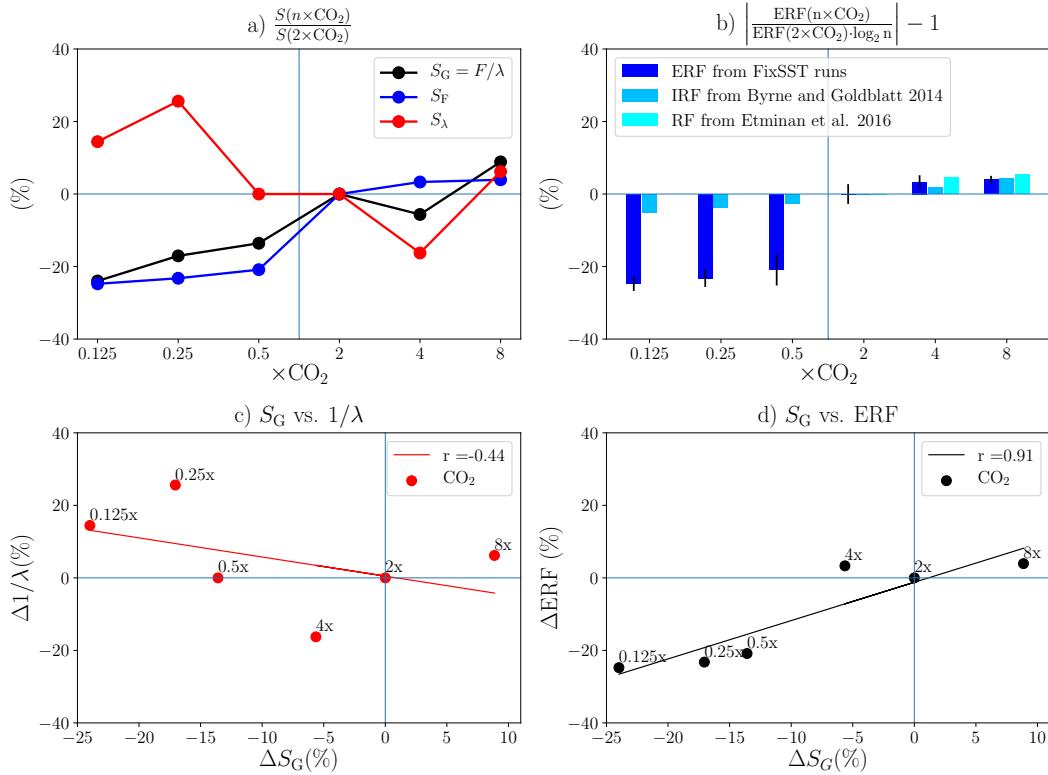


Figure 2. Percent change (from $2\times\text{CO}_2$) for abrupt CO_2 runs with CESM-LE model of: a)

climate sensitivity as x-intercept of Gregory Regression (black, S_G), as a function of ERF (blue, S_F), and as a function of $1/\lambda$ (red, S_λ); b) effective radiative forcing (dark blue, ERF), instantaneous radiative forcing (IRF) fit from Byrne and Goldblatt (2014) (light blue), and stratospherically adjusted radiative forcing (RF) fit from Etminan et al. (2016) (cyan). c) Percent change of S_G vs. $1/\lambda$ (red) and d) S_G vs. ERF (black). r is the Pearson correlation coefficient.

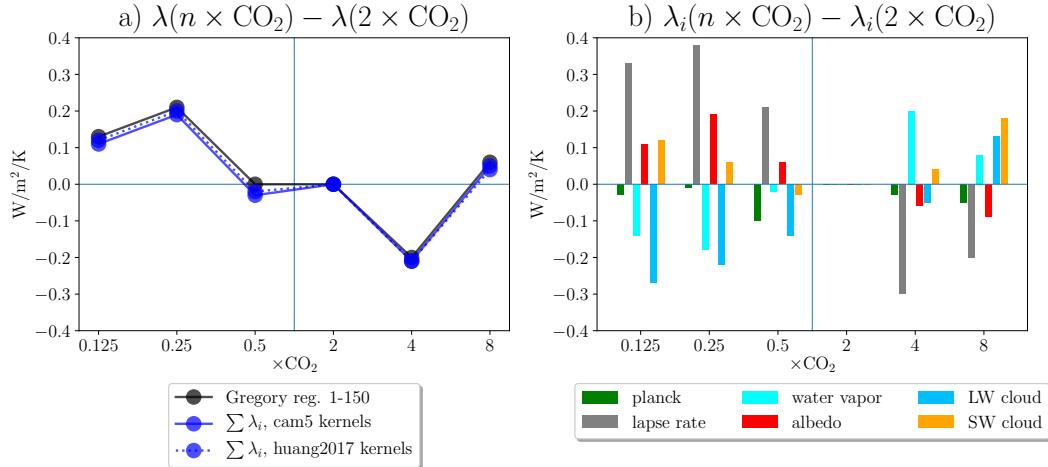


Figure 3. Feedbacks for abrupt CO_2 runs with CESM-LE model are shown as a difference from to $2 \times \text{CO}_2$. a) Total feedback calculated with Gregory Regression years 1-150 (black), Pendergrass et al. (2018) kernels for CESM1-CAM5 (blue solid), and Huang et al. (2017) kernels (blue dashed). b) Individual feedbacks calculated with Pendergrass et al. (2018) kernels.

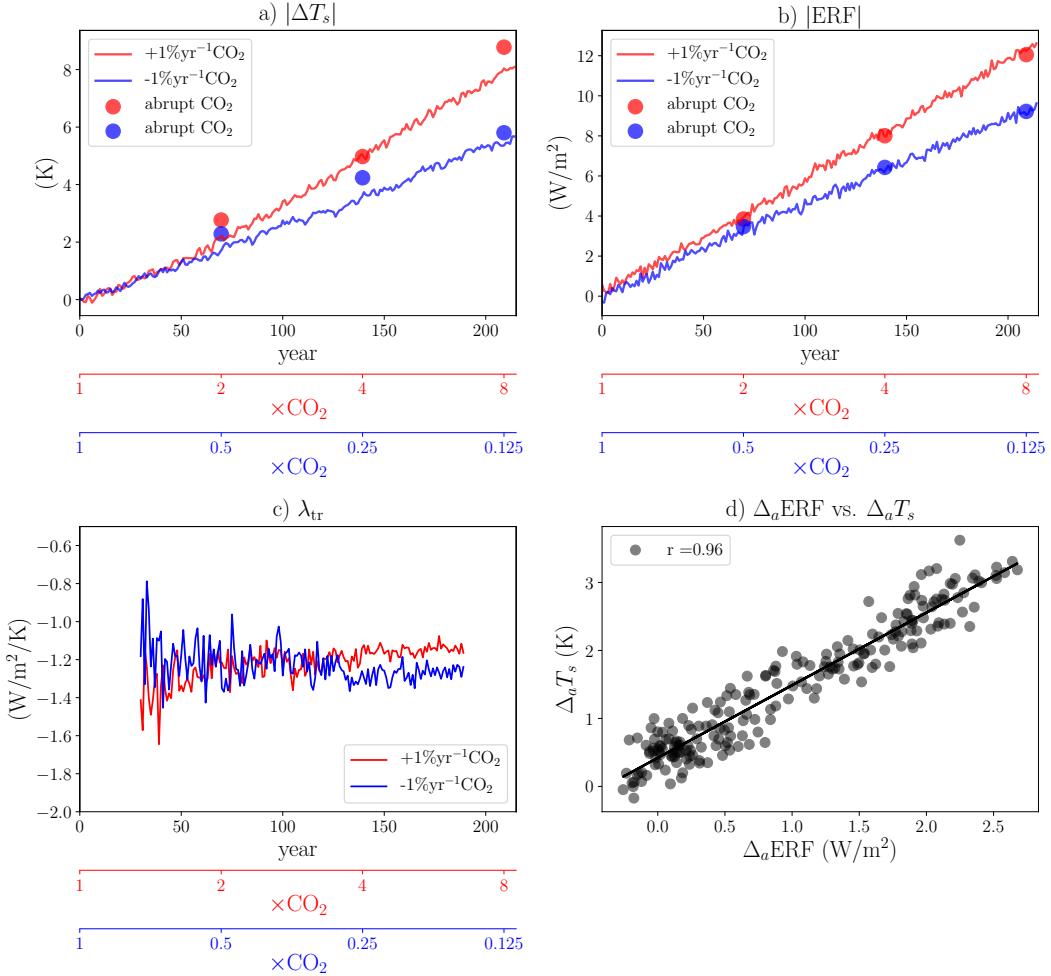


Figure 4. Transient runs annual timeseries with CESM-LE of a) the absolute value of surface temperature response ($|\Delta T_s|$), b) effective radiative forcing ($|\text{ERF}|$), c) net feedback (λ_{tr}), and d) correlation between asymmetries in $\Delta_a T_s$ and $\Delta_a \text{ERF}$. Responses from abrupt simulations are shown as dots.