

1            **Asymmetric Warming/Cooling Response to CO<sub>2</sub>**  
2            **Increase/Decrease Mainly Due to Non-Logarithmic**  
3            **Forcing, not Feedbacks**

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8            **Key Points:**

- 9            • The global surface temperature responds asymmetrically to increased and decreased  
10            CO<sub>2</sub> levels, in both abrupt and transient cases
- 11            • Effective climate sensitivity is higher with warming (2×, 4×, 8×CO<sub>2</sub>) than with  
12            cooling (1/2×, 1/4×, 1/8×CO<sub>2</sub>), in two different coupled models
- 13            • The non-logarithmic nature of the CO<sub>2</sub> forcing is primarily responsible for the asym-  
14            metry, not the radiative feedbacks

**Abstract**

We explore the CO<sub>2</sub> dependence of effective climate sensitivity ( $S_G$ ) with symmetric abrupt and transient CO<sub>2</sub> forcing, spanning the range  $1/8\times$ ,  $1/4\times$ ,  $1/2\times$ ,  $2\times$ ,  $4\times$ , and  $8\times$ CO<sub>2</sub>, using two state-of-the-art fully coupled atmosphere-ocean-sea-ice-land models. In both models, under abrupt CO<sub>2</sub> forcing, we find an asymmetric response in surface temperature and  $S_G$ . The surface global warming at  $8\times$ CO<sub>2</sub> is more than one third larger than the corresponding cooling at  $1/8\times$ CO<sub>2</sub>, and  $S_G$  is CO<sub>2</sub> dependent, increasing non-monotonically from  $1/8\times$ CO<sub>2</sub> to  $8\times$ CO<sub>2</sub>. We find similar CO<sub>2</sub> dependence in the transient runs, forced with  $-1\%yr^{-1}$ CO<sub>2</sub> and  $+1\%yr^{-1}$ CO<sub>2</sub> up to  $1/8\times$ CO<sub>2</sub> and  $8\times$ CO<sub>2</sub>, respectively. The non-logarithmic radiative forcing – not the changing feedbacks – primarily explains the dependence of  $S_G$  on CO<sub>2</sub>, particularly at low CO<sub>2</sub> levels. The changing feedbacks, however, explain  $S_G$ 's non-monotonic behavior.

**Plain Language Summary**

Equilibrium climate sensitivity (ECS) is the global mean warming after doubling CO<sub>2</sub> concentrations from those of the year 1850. Since CO<sub>2</sub> levels will likely surpass a doubling, it is crucial to know whether the amount of warming per CO<sub>2</sub> doubling (which we refer to as the effective climate sensitivity,  $S_G$ ) is constant with each CO<sub>2</sub> doubling or whether it changes. Necessary conditions for constant  $S_G$  are 1) the radiative forcing introduced to the climate system from each CO<sub>2</sub> doubling is constant and 2) the net radiative feedback does not change with CO<sub>2</sub> levels. Current literature shows that  $S_G$  will increase in a warmer world because the radiative feedback will change. We here investigate  $S_G$  in both warmer and colder worlds, and confirm that  $S_G$  increases at higher CO<sub>2</sub> concentrations. However, we show that changes in the radiative forcing with each CO<sub>2</sub> doubling are mainly responsible for  $S_G$  increase with CO<sub>2</sub>, not feedback changes.

## 1 Introduction

Equilibrium climate sensitivity (ECS) is the global mean surface temperature change after the doubling of CO<sub>2</sub> concentrations from pre-industrial (PI) levels. ECS is perhaps the most important metric in climate science, and it has been extensively investigated in the literature (Sherwood et al., 2020). An important question is whether the amount of warming for each CO<sub>2</sub> doubling (which we refer to as the effective climate sensitivity,  $S_G$ ) is constant or not (i.e., whether it is CO<sub>2</sub> dependent). Necessary conditions for a constant  $S_G$  are 1) that the radiative forcing of the climate system for each CO<sub>2</sub> doubling is constant and 2) that the net radiative feedback does not change with CO<sub>2</sub> levels. This question has been investigated in many modeling studies (Meraner et al., 2013; Mauritsen et al., 2019; Sherwood et al., 2020; Bloch-Johnson et al., 2021), which have reported that  $S_G$  is indeed CO<sub>2</sub> dependent. Most of these studies find that  $S_G$  increases at higher CO<sub>2</sub> levels and that the change in feedbacks, not the change in CO<sub>2</sub> radiative forcing, is the primary driver of  $S_G$  CO<sub>2</sub> dependence.

An alternative approach to using climate models to investigate the dependency of  $S_G$  on CO<sub>2</sub> is to seek observational constraints from reconstructions of past climates. In particular, most studies conclude that  $S_G$  inferred from paleoclimate records does depend on CO<sub>2</sub> (Caballero & Huber, 2013; Anagnostou et al., 2016; Shaffer et al., 2016; Friedrich et al., 2016; Farnsworth et al., 2019; Zhu et al., 2019; Anagnostou et al., 2020), although a few studies disagree (e.g., Martínez-Botí et al. (2015)). An ideal period to study the  $S_G$  from past climate is the Last Glacial Maximum (LGM), approximately 21 kyr ago, when the Earth was roughly 6K colder than PI conditions (Tierney et al., 2020). The LGM period is of particular interest because the climate system was in a quasi-equilibrium state, the climate forcings were large, and the surface temperature reconstructions are relatively well-constrained (Zhu & Poulsen, 2021). However, when considering the LGM and other periods in Earth's past, one needs to account for how the feedbacks in those past climate states differ

65 from the feedbacks operating in the modern state: hence the challenge in using paleoclimate-  
66 based estimates to constrain  $S_G$ .

67 While modeling and paleoclimatic evidence suggest that  $S_G$  depends on  $\text{CO}_2$ , a system-  
68 atic exploration of the symmetry over a wide range of  $\text{CO}_2$  forcing has yet to be performed.  
69 The question thus remains: is the climate system response symmetric across a broad range of  
70 positive (warm) and negative (cold)  $\text{CO}_2$  forcings? The question of symmetry was examined  
71 recently by Chalmers et al. (2022), who compared  $1/2\times$  and  $2\times\text{CO}_2$  simulations performed  
72 with the CESM1-CAM5 model, and found that global surface temperatures warm 20% more  
73 than they cool. Roughly 50% of this asymmetry was shown to derive from an asymmetry  
74 in  $\text{CO}_2$  radiative forcing; the rest was associated with differences in feedbacks which, inter-  
75 estingly, were found not to be related to clouds. Whether this result holds over a broader  
76 range of  $\text{CO}_2$  forcing, and whether it is model dependent remains an open question.

77 We here address these questions using a much broader range of both abrupt and tran-  
78 sient  $\text{CO}_2$  forcings, and do so with two different climate models. Specifically,  $\text{CO}_2$  is varied  
79 from  $1/8\times$  to  $8\times\text{PI}$  values, to test the  $\text{CO}_2$  symmetry of the climate system response to  
80 comparable increased and decreased  $\text{CO}_2$ . While we are not the first ones to perform such  
81 symmetric  $\text{CO}_2$  runs (Hansen et al., 2005; Colman & McAvaney, 2009; Russell et al., 2013;  
82 Chalmers et al., 2022), here we explore 1) a larger  $\text{CO}_2$  range than previously considered,  
83 2) we do so using two different fully coupled climate models and, most importantly, 3) we  
84 perform the experiments with both abrupt and transient  $\text{CO}_2$  runs.

85 Overall we confirm the asymmetric response in surface temperature: the climate system  
86 warms *more* with consecutive  $\text{CO}_2$  doublings ( $2\times$ ,  $4\times$ , and  $8\times\text{CO}_2$ ) than it cools with  
87 consecutive  $\text{CO}_2$  halvings ( $1/2\times$ ,  $1/4\times$ , and  $1/8\times\text{CO}_2$ ). This asymmetry is also reflected in  $S_G$ ,  
88 which *increases* at higher  $\text{CO}_2$  concentrations, consistent with previous studies. Surprisingly,  
89 we find that the non-logarithmic dependence of  $\text{CO}_2$  radiative forcing (i.e., the fact that  $\text{CO}_2$   
90 radiative forcing increases more rapidly than the log of the  $\text{CO}_2$  concentration) is primarily  
91 responsible for this asymmetric response, and not the changes in radiative feedbacks.

## 2 Methods

### 2.1 Models Used

We use two fully coupled atmosphere-ocean-sea-ice-land models: the large ensemble version of the Community Earth System Model (CESM-LE) and the NASA Goddard Institute for Space Studies Model E2.1-G (GISS-E2.1-G). CESM-LE comprises the Community Atmosphere Model version 5 (CAM5, 30 vertical levels), and parallel ocean program version 2 (POP2, 60 vertical levels) with approximately  $1^\circ$  horizontal resolution in all model components (Kay et al., 2015). GISS-E2.1-G is a 40-level atmospheric model with a resolution of  $2^\circ \times 2.5^\circ$  latitude/longitude, coupled to a  $1^\circ$  horizontal resolution 40-level GISS Ocean v1 (GO1) (Kelley et al., 2020). This configuration of the GISS model contributed to the CMIP6 project under the label “GISS-E2-1-G”. We show CESM-LE results in the main text, and some GISS-E2.1-G results in supplementary information (SI) to corroborate CESM-LE findings.

### 2.2 Abrupt $n \times \text{CO}_2$ Experiments

We perform a series of abrupt  $\text{CO}_2$  forcing runs using both models, subject to  $1/8 \times$ ,  $1/4 \times$ ,  $1/2 \times$ ,  $2 \times$ ,  $4 \times$ , and  $8 \times \text{CO}_2$  forcings, with all other trace gases, ozone concentrations, aerosols, and other forcings fixed at PI values. Following CMIP6 protocol for  $4 \times \text{CO}_2$  runs, we integrate all runs to 150 years starting from PI conditions. We contrast these to a PI control run to calculate the response.

For each model, we estimate the effective radiative forcing (ERF) with a companion series of  $\text{CO}_2$  experiments, as per Forster et al. (2016), with prescribed PI sea surface temperatures (SSTs) and sea-ice concentrations (SICs). These experiments are 30-year-long. We calculate ERF as the difference between the global mean net top of the atmosphere (TOA) flux between PI and  $n \times \text{CO}_2$  in these prescribed SSTs and SICs experiments. We do not here adjust for land warming simply because, in our ERF calculations, the surface

117 temperature response in the fixed SSTs and SICs simulations is minimal (Smith et al., 2020),  
 118 but we have verified that the adjustment does not change our results (see Figure S3).

### 119 **2.3 Transient Experiments**

120 In addition to the abrupt CO<sub>2</sub> runs, we also perform transient CO<sub>2</sub> runs with the  
 121 CESM-LE model. We start from PI conditions (same as in the abrupt CO<sub>2</sub> forcing), and  
 122 we increase CO<sub>2</sub> at +1%yr<sup>-1</sup> for the “warm” case for 215 years (slightly above 8×CO<sub>2</sub>) and  
 123 -1%yr<sup>-1</sup> for the “cold” case for 215 years (slightly below 1/8×CO<sub>2</sub>). We estimate transient  
 124 effective radiative forcing as in the abrupt experiments, by running companion simulations  
 125 with specified SSTs and SICs set to PI values (Forster et al., 2016), while ramping up CO<sub>2</sub>  
 126 at rates of +1%yr<sup>-1</sup> and -1%yr<sup>-1</sup>. We contrast all variables to PI values to compute the  
 127 response.

### 128 **2.4 Climate Sensitivity & Feedbacks**

We define effective climate sensitivity  $S_G$  as the x-intercept of the Gregory regression  
 (Gregory et al., 2004) for each abrupt  $n \times \text{CO}_2$  run using the following equation:

$$S_G = \left| \frac{F_{y-\text{int}}(n \times \text{CO}_2)}{\lambda(n \times \text{CO}_2) \cdot \log_2 n} \right| \quad (1)$$

129 We find the radiative forcing  $F_{y-\text{int}}$  as the y-intercept and the net feedback parameter  $\lambda$   
 130 as the slope from the Gregory regression (see Figure S1) where we regress the net TOA  
 131 radiative imbalance against the global mean surface temperature response for years 1-150.  
 132 In order to compare  $S_G$  for different CO<sub>2</sub> doubling / halving, we divide by  $\log_2 n$  (assuming a  
 133 logarithmic CO<sub>2</sub> forcing) and take the absolute value in Equation 1. Note that our definition  
 134 of the effective climate sensitivity  $S_G$  is a generalization of the more common definition of  
 135 effective climate sensitivity (which is typically defined as per Equation 1 but with  $n = 2$ ).  
 136 To check for the possibility that  $\lambda$  and  $S_G$  may be strongly affected by the “pattern effect”,  
 137 we have repeated the calculations by regressing years 21-150 only, and our main results were  
 138 not changed.

139 To calculate the individual feedbacks  $\lambda_i$ , we use radiative kernels ( $K_x$ ) from both  
 140 Pendergrass et al. (2018) and Huang et al. (2017) to quantify the sensitivity of TOA radia-  
 141 tion imbalance ( $\Delta R$ ) to changes in surface and atmospheric temperature ( $T$ ), water vapor  
 142 ( $q$ ), and surface albedo ( $\alpha$ ) (Soden et al., 2008; Shell et al., 2008). For each year of the 150-  
 143 year experiment, we multiply the spatially-resolved kernels by the climate field anomalies  
 144 ( $R_x = K_x \cdot \Delta x$ , where  $x$  is  $T, q, \alpha$ ), and then vertically integrate (for atmospheric temper-  
 145 ature and water vapor) up to the tropopause. We define the tropopause as 100 hPa at the  
 146 equator, 300 hPa at the poles, and in between, it varies by the cosine of the latitude (Soden  
 147 & Held, 2006). Lastly, we regress these quantities on the surface temperature response to  
 148 find the radiative feedbacks as the regression slope. The cloud feedbacks are computed via  
 149 the residual method (Soden & Held, 2006) as follows. First, we subtract effective radiative  
 150 forcing and the temperature, water vapor, and surface albedo radiative fluxes from the TOA  
 151 net radiative flux, resulting in  $\Delta R_{\text{cloud}} = \Delta R - \text{ERF} - \sum \Delta R_x$ . Then, we regress  $\Delta R_{\text{cloud}}$   
 152 onto  $\Delta T_s$  anomalies and define the corresponding slope as the cloud feedback. Lastly, we  
 153 find shortwave (SW) and longwave (LW) components of the cloud feedback by considering  
 154 the radiative changes in LW and SW components separately.

In the transient runs, we estimate the net feedback parameter  $\lambda_{\text{tr}}$  following Rugenstein  
 and Armour (2021) (see  $\lambda_{\text{eff1pct}}$  in their Figure 1d) with the expression:

$$\lambda_{\text{tr}} = -\frac{\text{ERF}(t) - \Delta R(t)}{\Delta T_s(t)} \quad (2)$$

155  $\Delta R(t)$  is the net TOA radiative imbalance, and  $\Delta T_s(t)$  is the global mean surface tempera-  
 156 ture response in the transient runs at year  $t$ .  $\Delta R(t)$  and  $\Delta T_s(t)$  are 30-year moving averages  
 157 of the respective terms. Note that we use different definitions for the feedback parameter  
 158 in the abrupt and transient simulations.

### 3 Results

#### 3.1 Abrupt CO<sub>2</sub> Experiments

We start by examining the global mean surface temperature response ( $|\Delta T_s|$ ) timeseries for the abrupt CO<sub>2</sub> runs (Figure 1). We contrast – in panels a, b, and c – the timeseries of each corresponding “warm” ( $2\times$ ,  $4\times$ , and  $8\times\text{CO}_2$ ) and “cold” simulation ( $1/2\times$ ,  $1/4\times$ , and  $1/8\times\text{CO}_2$ ) by taking the absolute value of the response from PI: note that the  $|\Delta T_s|$  in the “warm” case is always stronger than the “cold” case. In particular, we find 20% more warming at  $2\times$  than cooling at  $1/2\times\text{CO}_2$  (Figure 1a), 15% more at  $4\times$  than  $1/4\times\text{CO}_2$  (Figure 1b), and 41% more at  $8\times$  than  $1/8\times\text{CO}_2$  (Figure 1c). The asymmetry in  $|\Delta T_s|$  is amplified at higher CO<sub>2</sub> forcing, and largest in the  $1/8\times\text{CO}_2$  vs.  $8\times\text{CO}_2$  case (Figure 1c). The asymmetry is reduced at  $4\times\text{CO}_2$  vs.  $1/4\times\text{CO}_2$  due to changes in ocean heat transport which result in a formation of the North Atlantic Warming Hole in this model at  $4\times\text{CO}_2$  (see more details in Mitevski et al. (2021)).

To quantify the timescale of the asymmetry in  $|\Delta T_s|$  between “warm” and “cold” cases, we define the asymmetry between “warm” and “cold” cases as

$$\Delta_a X = |\Delta X(\text{warm})| - |\Delta X(\text{cold})| \quad (3)$$

where  $X$  is any climate variable (e.g.,  $T_s$ ), and subscript  $a$  refers to “asymmetry” (Figure 1d). In particular, we find that the asymmetry emerges rapidly in the first ten years (e.g., 90% at  $8\times\text{CO}_2$ ). Relative to the (slower) response associated with SST-driven feedbacks, the asymmetry appears quickly, suggesting that it might be due to radiative changes.

Next, we calculate effective climate sensitivity  $S_G$  from the Gregory regression (Equation 1), and plot it as percentage change from  $2\times\text{CO}_2$  (black line, Figure 2a).  $S_G$  is CO<sub>2</sub> dependent and increases with CO<sub>2</sub> concentration: at  $1/8\times\text{CO}_2$ , it is more than 20% *lower* than  $2\times\text{CO}_2$  values, and at  $8\times\text{CO}_2$ , it is around 5% *higher* than at  $2\times\text{CO}_2$ . CO<sub>2</sub> dependent  $S_G$  is possible if either the effective radiative forcing (ERF) or the net feedback paramete-



181 ter ( $\lambda$ ) change with  $\text{CO}_2$ . To individually test the relative importance of ERF and  $\lambda$ , we  
 182 calculate the climate sensitivity in two different ways.

First, to examine the dependence of climate sensitivity on ERF, we calculate climate sensitivity as  $S_F$  using the expression:

$$S_F = \left| \frac{\text{ERF}(n \times \text{CO}_2)}{\lambda(2 \times \text{CO}_2) \cdot \log_2 n} \right| \quad (4)$$

183 where ERF is derived from the  $n \times \text{CO}_2$  fixed SSTs and SICs runs, and  $\lambda$  (slope from Gregory  
 184 Regression) is held constant at the  $2 \times \text{CO}_2$  value. As seen in Figure 2a, we find that  $S_F$   
 185 (blue line) changes in tandem with  $S_G$  (black line), which reinforces the fact that changes  
 186 in ERF explain the changes in  $S_G$ .

Second, to assess whether changes in feedback strength also contribute to  $S_G$ , we calculate climate sensitivity as  $S_\lambda$ :

$$S_\lambda = \left| \frac{\text{ERF}(2 \times \text{CO}_2)}{\lambda(n \times \text{CO}_2)} \right| \quad (5)$$

187 where  $\lambda$  is calculated at each  $n \times \text{CO}_2$  and ERF is held constant at  $2 \times \text{CO}_2$  value. As seen in  
 188 Figure 2a,  $S_\lambda$  (red) changes in the opposite direction than  $S_G$  (black) for  $\text{CO}_2$  values lower  
 189 than  $2 \times \text{CO}_2$ . This suggests that changes in  $\lambda$  are not the main driver of the  $S_G$  dependence  
 190 on  $\text{CO}_2$ . However, it is important to note that for  $\text{CO}_2$  values higher than  $2 \times \text{CO}_2$ , we  
 191 find  $\lambda$  non-monotonically increasing to  $8 \times \text{CO}_2$ , which can be linked to the corresponding  
 192 non-monotonic behavior of  $S_G$ . We find qualitatively similar results using the GISS-E2.1-G  
 193 model (Figure S2a), confirming that ERF is the primary driver of the dependence of  $S_G$  on  
 194  $\text{CO}_2$ .

195 Next, we correlate  $S_G$  with  $1/\lambda$  (Figure 2c) and ERF (Figure 2d) across all abrupt  $\text{CO}_2$   
 196 experiments from  $1/8 \times$  to  $8 \times \text{CO}_2$  to examine whether feedbacks or forcing better correlate  
 197 with changes in  $S_G$ . Overall, we find little correlation between  $S_G$  and  $1/\lambda$  ( $r=-0.44$ ) and  
 198 a very strong correlation between  $S_G$  and ERF ( $r=0.91$ ). Similarly, a high correlation  
 199 between  $S_G$  and ERF is found in the GISS-E2.1-G model (Figure S2d). This strengthens our  
 200 conclusions from Figure 2a that the changes in ERF are driving the  $S_G$  increase. However,

201 if one considers warm cases, one sees a strong correlation between  $S_G$  and  $1/\lambda$ , as indicated  
202 earlier. This is in agreement with previous studies (Meraner et al., 2013; Bloch-Johnson et  
203 al., 2021), which reported that feedback changes are important for the dependence of  $S_G$   
204 on  $\text{CO}_2$ . However, over a broad range of  $\text{CO}_2$  forcing, including colder climates, that is not  
205 the case: changes in ERF are more important than feedback changes.

206 Given the aforementioned importance of ERF in driving the changes in  $S_G$ , we next  
207 look in more detail at ERF, calculated from fixed SSTs and SICs runs, following Forster et  
208 al. (2016), from  $1/8\times$  to  $8\times\text{CO}_2$  (dark blue bars, Figure 2b). If ERF were scaled simply  
209 with the logarithm of  $\text{CO}_2$  concentration, then the dark blue bars would be identical for  
210 all  $\text{CO}_2$  values. However, we see that ERF grows more than logarithmically with  $\text{CO}_2$ . We  
211 find a similar but weaker non-logarithmic behavior in the instantaneous radiative forcing  
212 (IRF) reported in Byrne and Goldblatt (2014), which we obtain by linearly interpolating  
213 their line-by-line radiative calculations (SI file “text03.txt” in Byrne and Goldblatt (2014))  
214 and plot with light blue bars in Figure 2b. We also compare our ERF calculations with  
215 the proposed stratospherically adjusted radiative forcing fit in Etminan et al. (2016) for the  
216 warming case only (since it is not valid for low  $\text{CO}_2$  values), and it appears both are in  
217 agreement.

218 A limitation to our ERF calculation approach is that we only fix the SSTs and SICs  
219 in the simulation, but not the land temperatures. Fixing the land temperatures has been  
220 shown to increase ERF in warmer climates even more than when only SSTs and SICs are  
221 fixed (Andrews et al., 2021). To account for this, we removed the land and sea-ice warming  
222 effects in our ERF calculations, following Equation 1 in Hansen et al. (2005) as shown in  
223 Figure S3, and found that the correction (dashed blue lines) leads, if anything, to a stronger  
224 non-logarithmic ERF. Hence, incorporating fixed land temperatures leads to ERF increasing  
225 even more rapidly than the log of  $\text{CO}_2$  concentration; this strengthens our argument that  
226 the  $S_G$  dependence on  $\text{CO}_2$  is due to non-logarithmic  $\text{CO}_2$  radiative forcing.

227 Next, we perform a standard decomposition of  $\lambda$  into individual radiative feedbacks  $\lambda_i$ .  
 228 The summation of individual feedbacks ( $\sum \lambda_i$ ) is shown in Figure 3a (blue).  $\sum \lambda_i$  follows  
 229 closely the net feedback calculated from the Gregory regression (black). We perform the  
 230 decomposition using two radiative kernels from Pendergrass et al. (2018) and Huang et al.  
 231 (2017), and we find minimal sensitivity to the choice of kernel (Figure S4). The individual  
 232 feedbacks, plotted as differences from  $2\times\text{CO}_2$  values, from the Pendergrass et al. (2018)  
 233 kernels are shown in Figure 3b. We see a clear signal in the lapse rate feedback, which  
 234 weakens the net feedback in the “cold” case and strengthens it in the “warm” case. The  
 235 longwave cloud feedback has clear global surface temperature dependence, increasing with  
 236  $\text{CO}_2$  monotonically for all  $\text{CO}_2$  values. However, in general, we find no clear pattern in the  
 237 changes in individual feedbacks that would sufficiently explain the overall feedbacks  $\text{CO}_2$   
 238 dependence. In addition, the changes in feedbacks in the GISS-E2.1-G model (Figure S5)  
 239 are qualitatively different from those in the CESM-LE model (Figure 3). Since our models  
 240 do not agree on the changes in individual feedbacks across the  $\text{CO}_2$  range, and since we  
 241 showed that feedback changes are strongly not correlated with changes in  $S_G$  (Figure 2c),  
 242 we do not explore further the mechanisms driving feedback changes in the individual models.

### 243 3.2 Transient $\text{CO}_2$ runs

244 The abrupt  $\text{CO}_2$  forcing runs show that the effective climate sensitivity increases with  
 245  $\text{CO}_2$ , and that the non-logarithmic nature of the ERF is largely responsible for this behavior.  
 246 Now we seek to determine whether the same behavior is also seen in runs with transient  
 247  $\text{CO}_2$  forcing, which are much more realistic. Our transient runs are forced, starting from  
 248 PI, with  $\text{CO}_2$  concentrations increasing at the rate of  $1\%_{\text{yr}^{-1}}$  and decreasing at  $1\%_{\text{yr}^{-1}}$ . As  
 249 seen in Figure 4a, the surface temperature response  $|\Delta T_s|$  is stronger in the warming (red)  
 250 than in the cooling (blue) case. Note that the responses computed from the last 50 years  
 251 of the abrupt simulations at the corresponding  $\text{CO}_2$  value (dots) are a good predictor of  
 252 the response in the transient runs, demonstrating that the results of the abrupt runs carry

253 over to the transient runs. Together with the surface temperature, ERF also changes more  
 254 rapidly in the warming than the cooling experiments, as seen in Figure 4b.

255 Next, we explore how the transient feedbacks ( $\lambda_{tr}$ , see Equation 2) change in the “warm”  
 256 and “cold” cases (Figure 4c). The feedbacks timeseries are noisy at the beginning of the  
 257 simulation, but in the last thirty years, the warm case shows 10% weaker (more positive)  
 258 feedbacks compared to the cold case. The 10% difference indicates that  $S_G$  in the warming  
 259 case should be higher than in the colder case. However, a robust difference in feedbacks  
 260 only appears around year 130, whereas the  $|\Delta T_s|$  asymmetry emerges much earlier, around  
 261 year 60. This difference in the temporal evolution of the feedbacks, relative to the evolution  
 262 of the forcing and  $S_G$ , adds additional strong evidence that the feedbacks are not driving  
 263 the  $|\Delta T_s|$  asymmetry.

264 Finally, as for the abrupt CO<sub>2</sub> runs, we correlate the asymmetry in global mean surface  
 265 temperature response  $\Delta_a T_s$  and effective radiative forcing  $\Delta_a \text{ERF}$  (Figure 4d). We find  
 266 a correlation of  $r=0.96$ , suggesting that the asymmetric changes in ERF drive the  $|\Delta T_s|$   
 267 asymmetry between the “cold” and “warm” cases. As we can see in Figure 4c, the transient  
 268 feedbacks are contributing to the  $|\Delta T_s|$  asymmetry at the end of the run, but their impact  
 269 is much smaller than the one from ERF.

## 270 4 Summary and Discussion

271 We have explored the effective climate sensitivity ( $S_G$ ) dependence on CO<sub>2</sub> with abrupt  
 272 *and* transient CO<sub>2</sub> experiments spanning the range  $1/8\times$  to  $8\times$ CO<sub>2</sub> using two distinct CMIP-  
 273 class climate models. First, we have found a considerable asymmetry in surface temperature  
 274 response, with the climate system warming more than cooling for identical factors used to  
 275 increase and decrease the CO<sub>2</sub> concentration, starting from a pre-industrial climate. Second,  
 276 we showed that the asymmetry is due to the non-logarithmic nature of CO<sub>2</sub> radiative forcing,  
 277 not the feedback changes. Upon decomposing the total feedback into individual feedbacks,

278 we found no simple explanation relating specific feedback changes to the changes in  $S_G$   
279 across the  $1/8\times$  to  $8\times\text{CO}_2$  forcing range examined in this study.

280 Most studies to date have focused on the role of feedbacks in explaining the dependency  
281 of  $S_G$  on  $\text{CO}_2$ , with relatively little attention placed on radiative forcing. Indeed, consistent  
282 with these studies, we found that for warmer climates ( $> 2\times\text{CO}_2$ ), feedbacks are important  
283 for determining the changing behavior of  $S_G$  with  $\text{CO}_2$ . However, by considering a broader  
284 range of  $\text{CO}_2$  forcings, we have shown here that for cases in which  $\text{CO}_2$  concentrations are  
285 less than PI values, non-logarithmic ERF is the primary driver of  $S_G$  changes. Our goal  
286 here has been to isolate the role of  $\text{CO}_2$  alone, and we have set all other forcings to PI  
287 values. Needless to say, we have ignored the “slow” feedbacks present in cold climates (e.g.,  
288 the LGM), such as the formation of land ice sheets.

289 The results with our abrupt runs have been shown to be robust with two climate  
290 models for simulations up to 150 years. One may argue that our runs are not equilibrated,  
291 and we agree with that caveat. However, we have found that the asymmetry and the key  
292 role of ERF are also robustly seen in the transient runs. Because of this, we expect that  
293 prolonging the abrupt simulation for more than 150 years will yield similar results. In  
294 any case, it will be important to repeat similar experiments with longer simulations as in  
295 LongRunMIP (Rugenstein et al., 2019) to confirm that this asymmetry is still present at  
296 long times closer to equilibration. Finally, our findings indicate that future studies should  
297 place more emphasis on accurately quantifying the changes in effective radiative forcing  
298 when studying the effective climate sensitivity dependency on  $\text{CO}_2$ . The feedbacks appear  
299 unable to explain the cooling phase.

## 300 5 Open Research

301 The CESM-LE model data can be obtained at [https://doi.org/10.5281/zenodo](https://doi.org/10.5281/zenodo.5725084)  
302 [.5725084](https://doi.org/10.5281/zenodo.5725084) and GISS-E2.1-G model data at <https://doi.org/10.5281/zenodo.3901624>.

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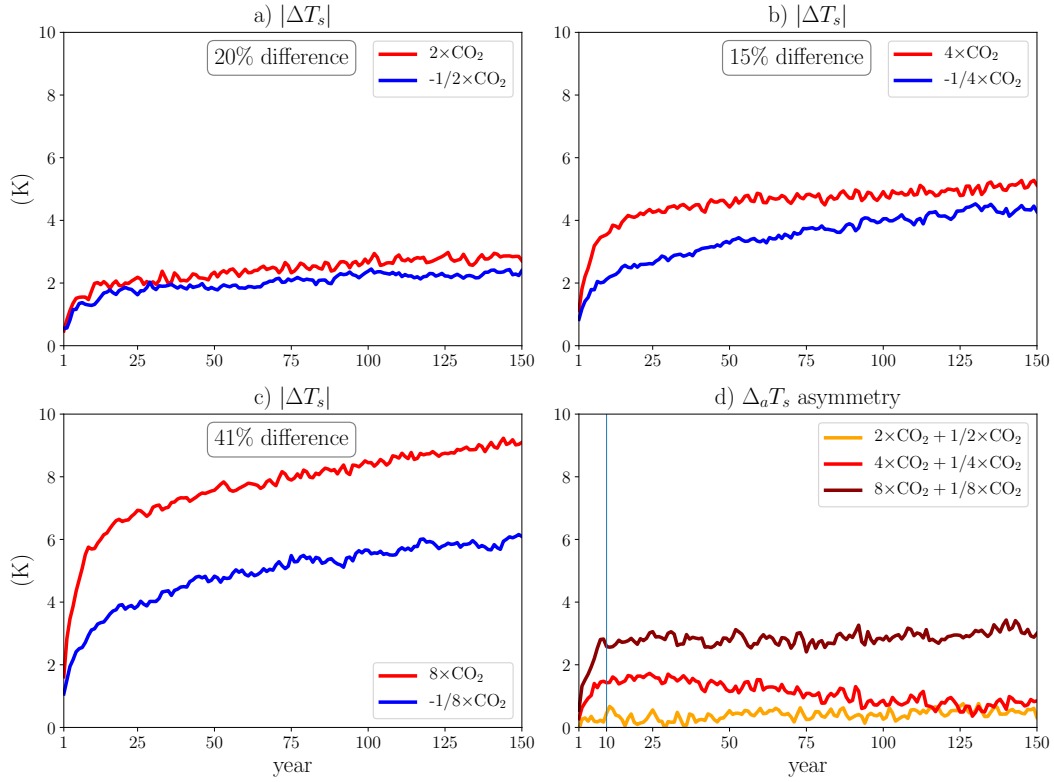
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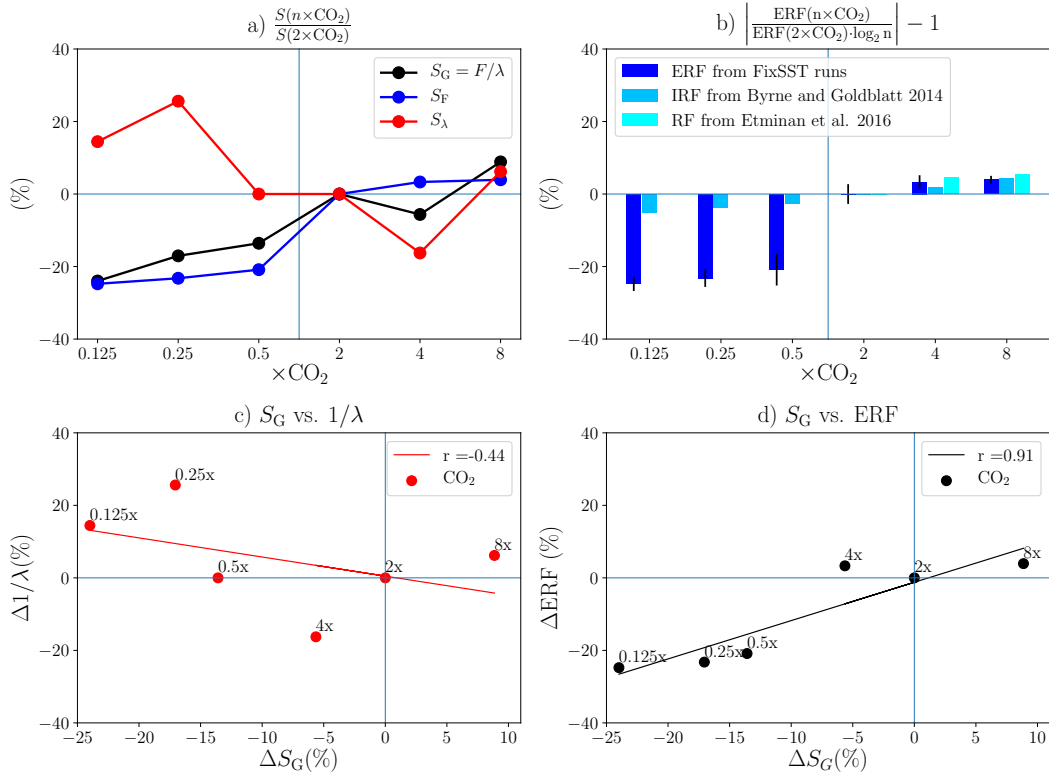
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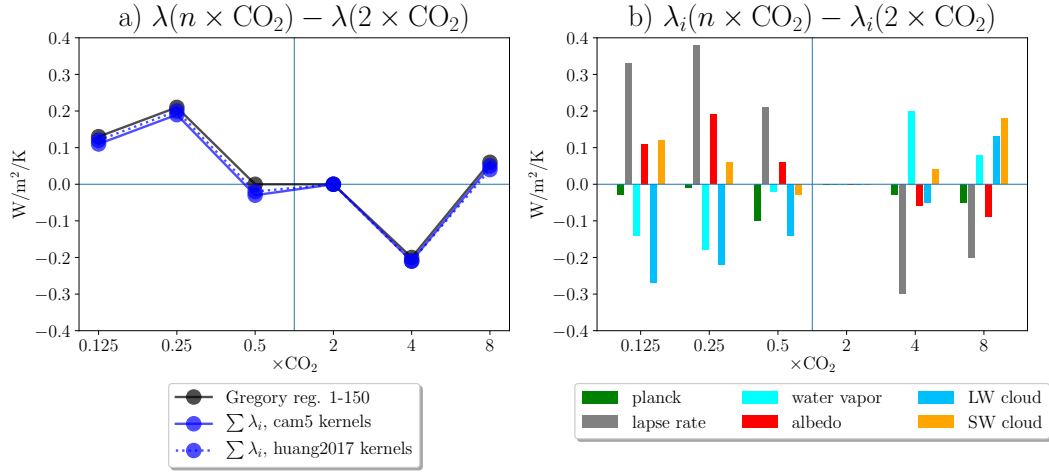
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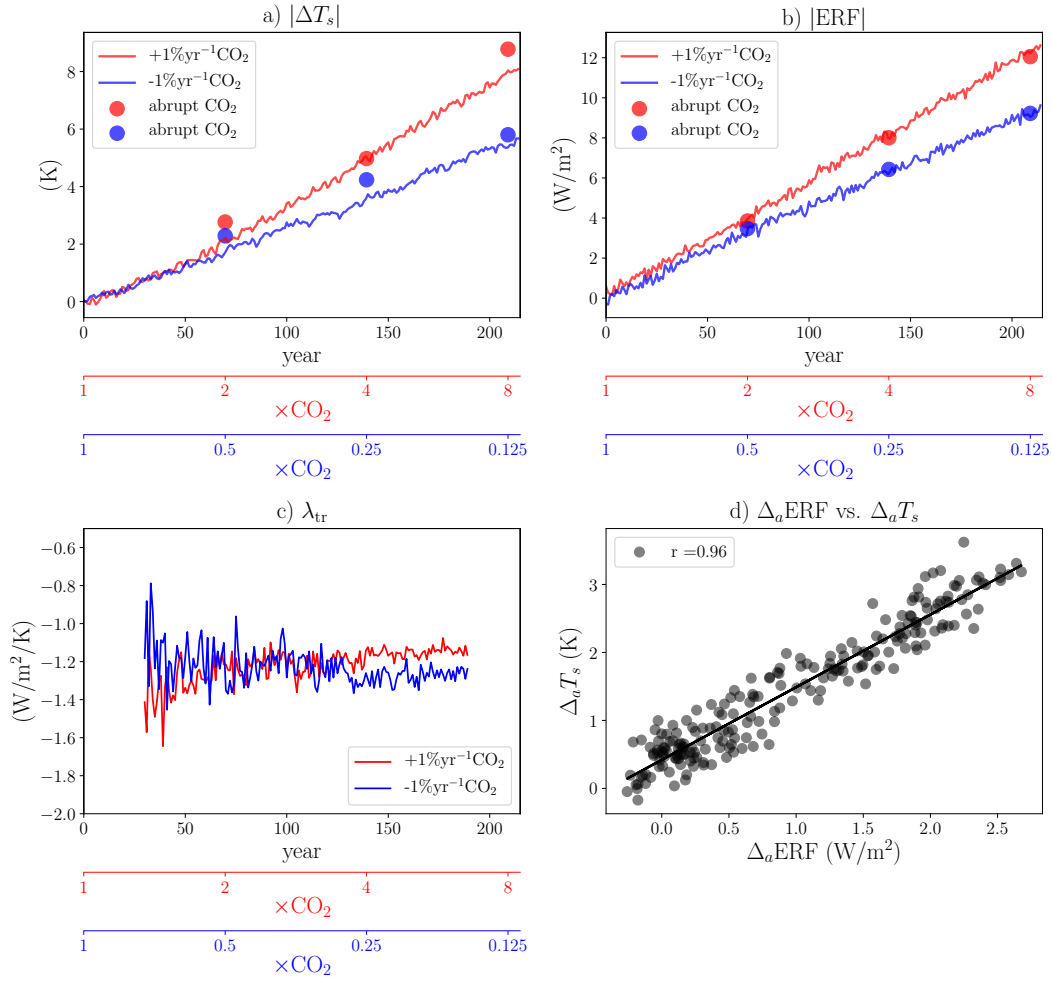
**Figure 1.** Timeseries of surface temperature response ( $|\Delta T_s|$ ) for abrupt CO<sub>2</sub> runs with CESM-LE model. a)  $2\times\text{CO}_2$  and  $1/2\times\text{CO}_2$ , b)  $4\times\text{CO}_2$  and  $1/4\times\text{CO}_2$ , c)  $8\times\text{CO}_2$  and  $1/8\times\text{CO}_2$  runs, and d) surface temperature asymmetry ( $\Delta_a T_s$ ) between “warm” and “cold” cases.



**Figure 2.** Percent change (from  $2 \times \text{CO}_2$ ) for abrupt  $\text{CO}_2$  runs with CESM-LE model of: a) climate sensitivity as x-intercept of Gregory Regression (black,  $S_G$ ), as a function of ERF (blue,  $S_F$ ), and as a function of  $1/\lambda$  (red,  $S_\lambda$ ); b) effective radiative forcing (dark blue, ERF), instantaneous radiative forcing (IRF) fit from Byrne and Goldblatt (2014) (light blue), and stratospherically adjusted radiative forcing (RF) fit from Etminan et al. (2016) (cyan). c) Percent change of  $S_G$  vs.  $1/\lambda$  (red) and d)  $S_G$  vs. ERF (black).  $r$  is the Pearson correlation coefficient.



**Figure 3.** Feedbacks for abrupt  $CO_2$  runs with CESM-LE model are shown as a difference from  $2 \times CO_2$ . a) Total feedback calculated with Gregory Regression years 1-150 (black), Pendergrass et al. (2018) kernels for CESM1-CAM5 (blue solid), and Huang et al. (2017) kernels (blue dashed). b) Individual feedbacks calculated with Pendergrass et al. (2018) kernels.



**Figure 4.** Transient runs annual timeseries with CESM-LE of a) the absolute value of surface temperature response ( $|\Delta T_s|$ ), b) effective radiative forcing ( $|ERF|$ ), c) net feedback ( $\lambda_{tr}$ ), and d) correlation between asymmetries in  $\Delta_a T_s$  and  $\Delta_a ERF$ . Responses from abrupt simulations are shown as dots.