

# Interplanetary Mesoscale Observatory (InterMeso): A mission to untangle dynamic mesoscale structures throughout the heliosphere

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## 27 Abstract

28 Mesoscale dynamics are a fundamental process in space physics, but fall within an observational gap  
29 of current and planned missions. Particularly in the solar wind, measurements at the mesoscales  
30 (100's  $R_E$  to a few degrees heliographic longitude at 1 au) are crucial for understanding the  
31 connection between the corona and an observer anywhere within the heliosphere. Mesoscale  
32 dynamics may also be key to revealing the currently unresolved physics regulating particle  
33 acceleration and transport, magnetic field topology, and the causes of variability in the composition

34 and acceleration of solar wind plasma. Studies using single-point observations do not allow for  
 35 investigations into mesoscale solar wind dynamics and plasma variability, nor do they allow for the  
 36 exploration of the sub-structuring of large-scale solar wind structures like coronal mass ejections  
 37 (CMEs), co-rotating/stream interaction regions (CIR/SIRs), and the heliospheric plasma sheet.

38 To address this fundamental gap in our knowledge of the heliosphere at these scales, the  
 39 Interplanetary Mesoscale Observatory (InterMeso) concept employs a multi-point approach using  
 40 four identical spacecraft in Earth-trailing orbits near 1 au. Varying drift speeds of the InterMeso  
 41 spacecraft enable the mission to span a range of mesoscale separations in the solar wind, achieving  
 42 significant and innovative science return. Simultaneous, longitudinally-separated measurements of  
 43 structures co-rotating over the spacecraft also allow for disambiguation of spatiotemporal variability,  
 44 tracking of the evolution of solar wind structures, and determination of how the transport of energetic  
 45 particles is impacted by these variabilities.

46

## 47 **1 Introduction**

48 The solar wind is a multi-scale and highly dynamic system with interplay between the micro-, meso-,  
 49 and macro-scales. Decades of single-point observations have led to great insight into the large-scale  
 50 variability in the solar wind and its transient phenomena, such as coronal mass ejections (CMEs),  
 51 stream and co-rotating interaction regions (SIR/CIRs), and solar energetic particle (SEP) events  
 52 associated with flares and/or CMEs. These observations have revealed differences between high-speed  
 53 streams originating from coronal holes, typical slow solar wind, and Alfvénic slow solar wind (see  
 54 *D’Amicis et al.*, 2019) and established the current paradigm under which the solar wind is interpreted  
 55 today.

56 Previous studies from the Solar Terrestrial Relations Observatory (STEREO) mission and other  
 57 serendipitous multi-point observations have allowed investigations of spatial variations of the solar  
 58 wind and transients over large distances. For instance, the radial evolution of the expanding solar wind  
 59 has been studied statistically using observations from Helios (*Perrone et al.*, 2018), while the radial  
 60 evolution of CIRs has been studied within the orbit of Earth (e.g., *Richter & Luttrell*, 1986; *Schwenn*,  
 61 1990; *Jian et al.*, 2008; *Allen et al.*, 2021a, b) and between Earth and Mars (e.g., *Geyer et al.*, 2021).  
 62 Multi-point observations of CIRs, and their associated energetic particles, over large longitudinal  
 63 separations have also revealed significant temporal evolution of the structures as they corotate over  
 64 10’s of degrees longitude (*Mason et al.*, 2009; *Jian et al.*, 2019; *Allen et al.*, 2021a). Additionally,  
 65 radial (e.g., *Burlaga et al.*, 1981; *Liewer et al.*, 2020) and longitudinal (e.g., *Kilpua et al.*, 2009;  
 66 *Farrugia et al.*, 2011; *Kollhoff et al.*, 2021) studies of CME structures have found significant variations  
 67 over these large separations.

68 On smaller, ion kinetic scales (<1000 km), the solar wind and transients have been found to be highly  
 69 turbulent and structured, with clear signs of coupling processes over a large range of spatiotemporal  
 70 scales (e.g., *Bandyopadhyay et al.*, 2018; *Roberts et al.*, 2020). While the Magnetospheric Multiscale  
 71 (MMS) and Cluster missions, designed to target kinetic scales, have demonstrated the importance and  
 72 richness of small-scale and highly dynamic plasma processes, these scales are “on the receiving end” of  
 73 the turbulent cascade that is driven by the large-scale solar wind structures (*Verscharen et al.*, 2019).  
 74 As such, the intermediate scale, between the small, kinetic scale dynamics explored by MMS and  
 75 Cluster, and the larger scale structuring and variations observed by STEREO and multi-mission  
 76 comparisons, represents the critical scale needed to understand cross-scale processes in the solar wind.

77 This intermediate scale – the mesoscale – is crucial for understanding the connection of the corona to  
 78 an observer anywhere within the heliosphere, as well as for revealing the currently unresolved physics  
 79 regulating particle transport, magnetic field topology, and the variability in composition and  
 80 acceleration of solar wind plasma.

81 The mesoscale solar wind currently falls within a gap both observationally and in spatiotemporal scales  
 82 of current simulations, and as such is a critical missing link in our fundamental understanding of the  
 83 heliosphere. A new mission targeting this critical gap of mesoscale dynamics would enable  
 84 investigations into how solar sources imprint themselves into the solar wind at 1 au and beyond via  
 85 mesoscale structuring, how mesoscale variability evolves as it propagates from the Sun, and how the  
 86 intrinsic structure of the solar wind impacts the structure of transients and particle acceleration and  
 87 transport. The mesoscale regime of the solar wind may likely be the missing piece to long sought  
 88 questions of sources of solar wind, particle acceleration, and particle transport. Gaining the ability to  
 89 probe the mesoscale solar wind will allow for leaps in our understanding of these outstanding  
 90 questions.

91

## 92 **2 InterMeso Science Objectives and Motivation**

93 To address the overarching science objective of *Investigating the fundamental mesoscale nature of*  
 94 *the variable solar wind and its impacts on particle acceleration, evolution, and transport*, the  
 95 Heliophysics community requires a mission to: (1) Characterize and identify the origin of the  
 96 mesoscale variability of the background solar wind and transient solar wind structures and (2)  
 97 Characterize and understand the impact of these mesoscale variations on particle acceleration and  
 98 transport.

### 99 **2.1 Objective 1: Characterize and identify the origin of the mesoscale variability of the** 100 **background solar wind and transient solar wind structures**

101 The solar surface exhibits structure on granule and supergranule scales (shown by the model results  
 102 in Figure 1a). As solar magnetic structures convect out with the solar wind, they can undergo  
 103 meandering due to footpoint motion, reconnection, and, possibly more importantly at 1 au, stochastic  
 104 motion and turbulent evolution (e.g., *Borovsky*, 2008; *Ashraf & Li*, 2019; *Bian & Li*, 2021). As such,  
 105 once flux tubes reach 1 au, they can be tangled into a complex meso-structure (illustrated in Figure  
 106 1b) (*Borovsky*, 2008). These processes and structuring can lead to various effects, such as “dropout”  
 107 phenomena in energetic particles (*Mazur et al.*, 2000). However, the fundamental structuring and  
 108 coherence of flux tubes in interplanetary space remain largely unknown.

109 Moving from the large-scale structures to mesoscale and smaller ranges, the nature of injection range  
 110 fluctuations in the solar wind is not well understood, nor is the transition at the break point between  
 111 the injection range and the inertial range (see the review by *Verscharen et al.*, 2019). Persistent,  
 112 large-scale structures may transition to dynamic structures at scales near  $10^9 - 10^{10}$  m (Figure 2).  
 113 However, this critical transitional scale between injection range and inertial ranges has been elusive  
 114 from a single-spacecraft vantage point, as these observations are limited in the ability to distinguish  
 115 temporal variations and those resulting from convection of spatially-variable solar wind over an  
 116 observer. While the currently planned HelioSwarm mission will, for the first time, robustly explore  
 117 the spatiotemporal transition between the inertial range and dissipation range, the injection-to-inertial  
 118 range transition will remain unexplored. HelioSwarm is a NASA Heliophysics Medium-Class  
 119 Explorer (MIDEX) mission currently in development that aims to study solar wind turbulence (at

120 scales of 50 - 3000 km), and consists of nine spacecraft (one hub and eight node spacecraft) in a lunar  
 121 resonant orbit (see <https://eos.unh.edu/helioswarm/mission> for additional details). Only through  
 122 multi-point observations with separations on the order of this critical scale can we robustly explore  
 123 this fundamental transition in the solar wind.

124 In addition to the currently unconstrained, fundamental mesoscale structuring of the solar wind, sub-  
 125 structuring of transient events is also not well understood. The small number of studies that have  
 126 utilized fortuitous, but sporadic, multi-mission conjunctions to investigate the mesoscale structuring  
 127 of CME shocks have shown that CME-associated shocks and magnetic ejecta have smaller-scale  
 128 structuring, although the degree of such structuring is largely not understood (e.g., *Bale et al.*, 1999;  
 129 *Knock et al.*, 2003; *Pulupa & Bale*, 2008; *Koval & Szabo*, 2010, *Lugaz et al.*, 2018). As shown in  
 130 Figure 3a, the various scale lengths of the solar wind, CME sheath, and CME ejecta are expected to  
 131 be different from one another (*Ala-Lahti et al.*, 2020). Additionally, comparisons of shocks between  
 132 the Advanced Composition Explorer (ACE) and Wind missions have found that energetic particle  
 133 time-intensity profiles often change over mesoscales indicating important effects of mesoscale  
 134 structuring on particle acceleration and transport (Figure 3b and 3c, *Neugebauer et al.*, 2006).  
 135 Observations such as these indicate that the “large-scale-only” view of the solar wind is an  
 136 incomplete picture of the fundamental structure of the solar wind and the important processes that  
 137 define its evolution. Dedicated multi-point observations are needed to reveal these fundamental  
 138 physical processes.

## 139 **2.2 Objective 2: Characterize and understand the impact of these mesoscale variations on** 140 **particle acceleration and transport**

141 Particle acceleration is fundamentally dependent upon local conditions in the acceleration regions  
 142 and, as such, mesoscale variations in the solar wind and along/within transients will affect  
 143 acceleration processes. For example, a CME expanding into solar wind streams with spatial pressure  
 144 variations may lead to deformation of its structure (e.g., *Owens et al.*, 2017), resulting in spatially  
 145 varying shock structure in the mesoscale range. Such mesoscale variations along a shock surface lead  
 146 to localized differences in shock parameters affecting particle acceleration, and so could potentially  
 147 explain the variations observed in Figure 3b-c. Detailed observations at multiple points along such  
 148 structures are required to determine the mesoscale effect on CME particle acceleration.

149 Radio wave observations have demonstrated the presence of mesoscale variations in shock structure  
 150 (e.g., *Bale et al.*, 1999; Figure 4c); however, it remains unknown to what degree the mesoscale  
 151 structuring affects particle acceleration. While current models can reproduce large-scale variation in  
 152 event-integrated particle spectra along an interplanetary shock (Figure 4a-b), which have been  
 153 observed by conjunctions of spacecraft at large separations (e.g., *Hu et al.*, 2018), these approaches  
 154 are unable to capture the effects of mesoscale structuring or short time cadence evolution (Figure 4d).  
 155 For example, an event-integrated spectra includes particles accelerated from various locations along  
 156 the shock structure, but the seed populations and freshly accelerated ions will be more localized to  
 157 the mesoscale variations of the structure. Better understanding the role that mesoscale structure plays  
 158 in controlling the variability of particle acceleration along an interplanetary shock is critical for  
 159 investigating the causes of observed variation between events and potential seed populations.

160 Additionally, our understanding of flare acceleration of energetic electrons will be greatly advanced  
 161 through a robust exploration of mesoscale particle variations. For instance, *Li et al.* (2020) found that  
 162 the injection timing of energetic electrons compared to hard X-ray observations suggests the presence  
 163 of two distinct electron populations in an impulsive solar energetic electron event (Figure 5). The

164 presence or absence of these electron populations can differentiate between different types of  
165 magnetic reconnection at the flare site (e.g., interchange reconnection versus reconnection between  
166 two closed field lines). However, the determination of the path length that the energetic electrons  
167 traversed from the flare to the observer is the primary source of ambiguity in these observations,  
168 which may be a result of mesoscale variations impacting particle transport. For example, field line  
169 meandering and/or stochastic motion can lead to uncertainties of where a source region maps to of up  
170 to  $10^\circ$  in longitude at 1 au (Figure 6b, *Bian & Li, 2021*). Observing energetic electrons at multiple  
171 spacecraft with a close separation ( $<10^\circ$  total span) will allow for the determination of acceleration  
172 time scales of electrons at solar flares, through fractional dispersion analysis (FVDA, *Zhao et al.,*  
173 *2019*) or other analysis techniques, and so for differentiating among proposed reconnection  
174 mechanisms.

175 The complexities of transport between the corona and 1 au are apparent in impulsive SEP events. For  
176 example, when the two STEREO spacecraft were within  $34^\circ$  heliographic longitude of one another in  
177 2014 (Figure 7a), a flare event occurred at an active region near the footpoints of both spacecraft  
178 (Figure 7b). During this event, as shown in Figure 7c, STEREO-A, despite the footpoint being  
179 slightly further from the active region, observed a higher peak in 55-65 keV electron intensity with  
180 an onset time about an hour earlier, and with higher anisotropy, than STEREO-B (*Klassen et al.,*  
181 *2016*). Because transport of flare-accelerated energetic electrons depends on background solar wind  
182 structure, particle acceleration at flare sites, and the mechanisms driving cross-field diffusion, this  
183 event presents a clear example of the importance of mesoscale variations on the transport of energetic  
184 electrons from the flare site in the corona to 1 au. As such, a robust multi-point investigation at  
185 mesoscale separations is required to disentangle the complexities of particle transport in the inner  
186 heliosphere.

187 For particle transport from local acceleration sites associated with transients, mesoscale variations  
188 may also play an important role. As energetic particles propagate along a field line away from an  
189 acceleration site, such as from reverse shocks at CIRs, various processes (e.g., adiabatic cooling and  
190 particle scattering) leading to hardening of the lower-energy spectra are thought to occur (e.g., *Fisk &*  
191 *Lee, 1980*). However, this spectral hardening has not been observed as often, or as significantly, as  
192 expected (e.g., *Mason & Sanderson, 1999; Desai et al., 2020; Allen et al., 2021a; Joyce et al., 2021*).  
193 To explain this discrepancy, several theories have been developed, such as compressive, non-shock  
194 related acceleration (e.g., *Giacalone et al., 2002; Fisk & Gloeckler, 2006; Ebert et al., 2012; Chen et*  
195 *al., 2015*) and/or modifications to the magnetic topology, such as sub-Parker spirals (e.g., *Murphy et*  
196 *al., 2002; Schwadron, 2002; Schwadron et al., 2020*).

197 While these other processes may explain the observed weak modulation of the particle spectra, a CIR  
198 observation utilizing a fortuitous multi-spacecraft vantage point seems to suggest strong path length-  
199 dependent modulation (*Zhao et al., 2016; Figure 8a-b*). In this event, a CIR shock was observed at  
200 STEREO-B when it was  $\sim 23^\circ$  from Earth, however the shock was not observed at ACE, and the CIR  
201 evolved too significantly relative to STEREO-A for comparison. STEREO-B and ACE  
202 measurements indicate strong path-dependent modulation of the observed particle spectra (Figure  
203 8b). However, the required assumption of negligible temporal evolution of the CIR between  
204 observations does not fully capture the possible variations of the shock and dynamics regulating  
205 particle acceleration and transport, adding ambiguity to the findings. Constraining the true path  
206 length to the reverse shock and spatiotemporal variations of shock-related acceleration is difficult  
207 with single point observations. Only simultaneous, multi-point observations from different points  
208 within a CIR and its rarefaction region allows for differentiation between path-length dependent  
209 modulation and competing transport and acceleration processes.

## 210 **2.3 Significance, Impact, and Timeliness**

211 Fundamentally understanding the mesoscale (100's  $R_E$  to a few degrees heliographic longitude at 1  
 212 au, i.e., in-between the inertial and injection ranges) structure of the solar wind and transients and its  
 213 subsequent effects on particle acceleration and transport will be paradigm shifting in our insight into  
 214 the heliosphere, as mesoscale dynamics are vital for resolving long-standing questions of the  
 215 community. As such, a new mission with an enabling, multi-point architecture is essential to address  
 216 these objectives. InterMeso will fill a critical observational gap in the current Heliophysics System  
 217 Observatory (HSO) as mesoscale solar wind structure and dynamics falls between the global scales  
 218 studied by the in situ instrumentation on STEREO (*Kaiser et al.*, 2008) and occasional opportune  
 219 conjunctions within the HSO, and the kinetic-scales unlocked by MMS (*Burch et al.*, 2016), Cluster  
 220 (*Escoubet et al.*, 2001), and the HelioSwarm mission (illustrated in Figure 9).

221 The goal of InterMeso, to resolve the critical physics and consequences of the mesoscale solar wind  
 222 and transients, is also particularly timely for the next decade. With the continued operation of  
 223 ground-based solar observatories (e.g., Daniel K. Inouye Solar Telescope, DKIST; *Tritschler et al.*,  
 224 2016) and solar/heliographic imaging satellite missions, InterMeso measurements will have  
 225 complementary remote sensing observations of the solar footpoints of the spacecraft. The upcoming  
 226 discoveries from the Polarimeter to Unify the Corona and Heliosphere (PUNCH;  
 227 <https://punch.space.swri.edu/>) mission will enable a better understanding of the initial mesoscale  
 228 structuring of the coronal young solar wind, and so will feed directly into the interpretations of  
 229 mesoscale variability at 1 au from InterMeso, furthering our insight into the evolution of this  
 230 mesoscale variability between the solar source and Earth at 1 au. Understanding of the transition  
 231 from the inertial to dissipation range of plasma turbulence from HelioSwarm, coupled with  
 232 understanding of the injection to inertial range from InterMeso, would for the first time provide a  
 233 broad picture of plasma variability in the mesoscale solar wind. Additionally, the European Space  
 234 Agency (ESA) Vigil mission to L5 ([https://www.esa.int/Space\\_Safety/Vigil](https://www.esa.int/Space_Safety/Vigil)), in partnership with the  
 235 National Oceanic and Atmospheric Administration (NOAA), is planned to launch in 2027 and will  
 236 include a heliospheric imager that can provide broader context for the in situ observations of the  
 237 InterMeso mission. The potential synergies between InterMeso and upcoming missions (e.g.,  
 238 HelioSwarm, PUNCH, and Vigil) motivate such an architecture within the next decade.

## 239 **3 The InterMeso Concept**

### 240 **3.1 InterMeso Science Traceability**

241 The overarching science goals of InterMeso flow into eight targeted science questions, shown in the  
 242 Science Traceability Matrix (STM) provided in Table 1. Addressing these questions requires  
 243 simultaneous, multi-point observations with inter-spacecraft separations ranging from  $\sim 0.5$  Mkm to  
 244  $10^3$ 's Mkm (i.e., 100's  $R_E$  to a few degrees in heliographic longitude at 1 au, see Figure 2).  
 245 Maintaining the constellation at a near-1 au heliocentric distance enables observations of both well-  
 246 formed and still-steepening shocks of CMEs and SIR/CIRs.

247 All InterMeso spacecraft require measurements of the low-energy, bulk solar wind plasma. This  
 248 includes sub-minute cadence observations of the bulk proton population for speeds ranging from 250  
 249 – 1000 km/s and densities from  $1 - 100 \text{ cm}^{-3}$ . This range of measurements with necessary accuracy  
 250 can be accomplished with modern Faraday Cup instrumentation (e.g., *Case et al.*, 2020). Thermal ion  
 251 (1-10's keV/e) composition measurements on a less than few minute cadence will be critical to  
 252 disambiguate changes in flux tube composition. This will require the ability to distinguish between

253 various species and charge states including H, He<sup>2+</sup>, C<sup>5-6+</sup>, N<sup>5-6+</sup>, O<sup>6-7+</sup>, and Fe<sup>6-12+</sup> and can be  
 254 satisfied by modern iterations of the Ulysses & ACE Solar Wind Ion Composition Spectrometer  
 255 (SWICS) instrument design (*Gloeckler et al.*, 1992). Low-energy electron observations spanning  
 256 from a few eV to tens keV, i.e., including the core, halo, and strahl populations, are also needed and  
 257 be achieved with the inclusion of a Parker Solar Probe (PSP) Solar Probe Analyzers-Electrons  
 258 (SPAN-E)-type instrument (*Kasper et al.*, 2016).

259 Additionally, suprathermal and energetic ion/electron measurements are required to explore  
 260 anisotropies and, importantly, differentiate various ion species for investigating mass-per-charge-  
 261 dependent processes of particle acceleration and transport. Suprathermal ion (5 – 100 keV/nuc) and  
 262 energetic ion (0.1 – 10 MeV/nuc) measurements with mass determination (i.e., H, <sup>3</sup>He, <sup>4</sup>He, C, N, O,  
 263 Fe) are required on a cadence of several minutes and tens of minutes, respectively, which is within  
 264 the capabilities of current generation versions of the Solar Orbiter Suprathermal Ion Spectrograph  
 265 (SIS; *Rodríguez-Pacheco et al.*, 2020; *Wimmer-Schweingruber et al.*, 2021) instrument. Energetic  
 266 electron observations (few tens keV to few MeV) are needed with capabilities of determining first-  
 267 order anisotropy on a few second cadence to characterize mesoscale variations in SEP events, which  
 268 can be achieved with Solar Orbiter Energetic Particle Telescope (EPT)-like instrumentation  
 269 (*Rodríguez-Pacheco et al.*, 2020; *Wimmer-Schweingruber et al.*, 2021).

270 To probe mesoscale magnetic structuring and impacts of plasma waves and turbulence on particle  
 271 transport, observations of the magnetic field and radio emissions released from flare events and shock  
 272 acceleration are required. Vector magnetic field measurements must have a full-scale range of at least  
 273 +/- 100 nT with 0.1 nT or better resolution at cadences of 16 – 64 vectors/sec, obtainable by current  
 274 fluxgate instrumentation (e.g., PSP/FIELDS; *Bale et al.*, 2016). Radio wave observations ranging  
 275 from 10<sup>-18</sup> – 10<sup>-12</sup> V<sup>2</sup>/m<sup>2</sup>/Hz over the frequency range of 0.1 kHz to 20 MHz, achievable with  
 276 STEREO-SWAVES (*Bougeret et al.*, 2008) or Solar Orbiter Radio Plasma Waves (RPW;  
 277 *Maksimovic et al.*, 2020)-type instrumentation, are also needed. Accurate timing and characterization  
 278 of electron acceleration near the solar surface necessitates observations of hard X-rays with an  
 279 indirect imager (e.g., *Krucker et al.*, 2020) along with radio wave observations to determine both the  
 280 source location and electron characteristics. The X-ray instrument must measure hard X-ray spectra,  
 281 images, and time series with an energy range between 5 – 100 keV, and is achievable with Solar  
 282 Orbiter Spectrometer Telescope for Imaging X-rays (STIX)-type instrumentation (*Krucker et al.*,  
 283 2020).

284 As such, the baseline payload for this mission requires each spacecraft of the constellation to be  
 285 outfitted with: **(1)** bulk solar wind instrument, **(2)** thermal ion composition instrument, **(3)**  
 286 suprathermal ion composition telescope, **(4)** energetic ion composition telescope, **(5)** low energy  
 287 electron instrumentation, **(6)** energetic electron telescopes, **(7)** DC vector magnetic field sensor, and  
 288 **(8)** high-frequency electric field radio wave instrumentation. Additionally, only one of the spacecraft  
 289 must be equipped with **(9)** a hard X-ray indirect imager due to the relatively close separation needed  
 290 for the InterMeso spacecraft to span the mesoscale regime.

291 Closure of the science objectives given in Table 1 requires simultaneous multipoint observations with  
 292 longitudinal separations that span the range of mesoscale dynamics in the solar wind at 1 au. Figure  
 293 10 illustrates the ability to reconstruct longitudinal variations following either a Gaussian (top) or  
 294 sigmoidal (bottom) distribution by taking a representative distribution (dashed line) and placing 1000  
 295 randomly-placed constellations of equally-spaced observing points within the domain. For the  
 296 Gaussian reconstructions, a single spacecraft can only result in a single value, two spacecraft can  
 297 provide a linear variation, three spacecraft provide a quadratic, while four and more spacecraft can

298 constrain different orders of Gaussian fits. Similarly, the sigmoidal distribution can be represented by  
 299 a single value, linear fit, quadratic fit, or sigmoidal fit. As demonstrated in Figure 10, increasing the  
 300 number of spacecraft increases the ability to discern characteristic longitudinal profiles relevant to  
 301 structures in the solar wind, particularly Gaussian-type distributions (i.e., the spread of energetic  
 302 particles from an acceleration region) and sigmoidal-type variations (i.e., current sheets or changes in  
 303 topology across flux tubes). Distinguishing more complex topological variabilities, such as “ripples”  
 304 in a large-scale shock structure (e.g., *Bale et al.*, 1999), or better constraining the non-planarity and  
 305 radius of curvature of interplanetary shocks (e.g., *Neugebauer & Giacalone*, 2005), also benefit from  
 306 increasing numbers of longitudinally-separated observations.

307 While additional spacecraft can increase the science return of InterMeso, four points of measurement  
 308 within relevant scales provide the highest benefit-to-cost for expected variability of mesoscale  
 309 phenomena with features manifesting as either a Gaussian or sigmoidal distribution (shown in Figure  
 310 10). With fewer than four spacecraft, the ability to reconstruct the distribution becomes significantly  
 311 deteriorated. However, while a constellation with more than four spacecraft can serve to hone the  
 312 resulting reconstructions of longitudinal variation in mesoscale phenomena, the added value,  
 313 particularly when including available mission resources and cost, diminishes with additional  
 314 spacecraft. As such, the InterMeso mission architecture targets four spacecraft as a baseline to  
 315 determine spatiotemporal variations and elucidate mesoscale structure.

### 316 **3.2 InterMeso Spacecraft, Orbit, and Separation Scales**

317 Key to the science of InterMeso, the mission requires multiple, longitudinally-separated spacecraft in  
 318 the solar wind near 1 au with varying inter-spacecraft separations over the course of the mission. In  
 319 family with spacecraft missions in development and in prime operations, the prime mission phase of  
 320 InterMeso is designed to last three years in order to sufficiently sample the mesoscale range in Table  
 321 1, but with systems that can remain operational likely for several years of extended mission  
 322 operation. The flexible mission design of InterMeso currently targets a 2033 launch to coincide with  
 323 the ascending phase of the solar cycle but is adjustable via daily launch opportunities throughout the  
 324 decade.

325 A design of the InterMeso constellation is shown in Figure 11. To reduce cost and complexity in  
 326 launch, the four InterMeso spacecraft can fit on a single Evolved Expendable Launch Vehicle  
 327 (EELV) Secondary Payload Adapter (ESPA)-Grande equivalent and within a standard 5m fairing  
 328 (shown in both the stowed and deployed configuration in Figure 11). All four spacecraft are launched  
 329 together into heliocentric, Earth-trailing orbit, and all with the same initial drift rate before  
 330 maneuvers separate the spacecraft. Figure 12a shows the initial drift rate set by the launch vehicle  
 331 (green line) as well as the drift rates of each InterMeso spacecraft throughout the mission. This  
 332 trajectory also maintains an inter-spacecraft separation of  $<0.01$  au at all times while keeping the  
 333 entire constellation between a heliographic distance of  $\sim 0.97$  and 1.05 au over the course of the prime  
 334 mission.

335 These drift rates allow the spacecraft to slowly separate and allow sampling of the solar wind at inter-  
 336 spacecraft separations spanning the full range of mesoscale dynamics (as illustrated in Figure 2).  
 337 Figure 12b illustrates the cumulative combined days at various inter-spacecraft separations for both  
 338 the three-year prime mission (black line) and for the inclusion of a two-year extended mission (red-  
 339 dotted line). The inclusion of the 2-year extended mission phase, especially, provides additional  
 340 opportunity for sampling mesoscale structuring near the transition from the inertial to the injection  
 341 range; however, this additional sampling is not required for addressing the primary science goals.

## 342 4 Summary

343 The mesoscale solar wind (100's  $R_E$  to a few degrees heliographic longitude at 1 au) remains elusive  
344 in our observations, modeling frameworks, and understanding. In order to finally close this  
345 fundamental gap in our knowledge of the solar wind, new multi-point missions are required with  
346 mesoscale separations. The InterMeso mission concept provides one such solution for addressing the  
347 unresolved physics of mesoscale structure of the solar wind and transients.

348 Consisting of four near-identical spacecraft operating independent of one another and outfitted with a  
349 suite of particle and field instrumentation, InterMeso is designed to robustly explore the mesoscale  
350 structuring of the solar wind and its consequences on particle acceleration and transport. Operation  
351 over the three-year prime mission lifetime allow for sampling over the full range of mesoscale  
352 variability in the solar wind. The science objectives addressed by the architecture of InterMeso is  
353 achievable in the next decade, enabling advances on long-outstanding questions in heliophysics.

## 354 5 Author Contributions

355 All authors performed research on this mission concept and contributed information for the  
356 manuscript.

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## 359 7 Conflict of Interest

360 The authors declare that the research was conducted in the absence of any commercial or financial  
361 relationships that could be construed as a potential conflict of interest.

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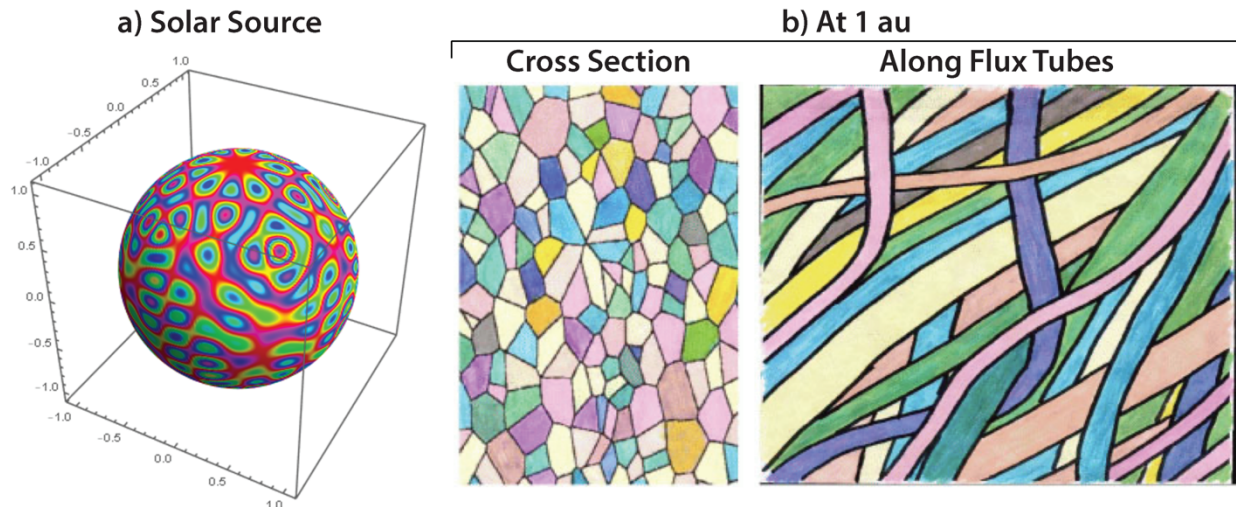
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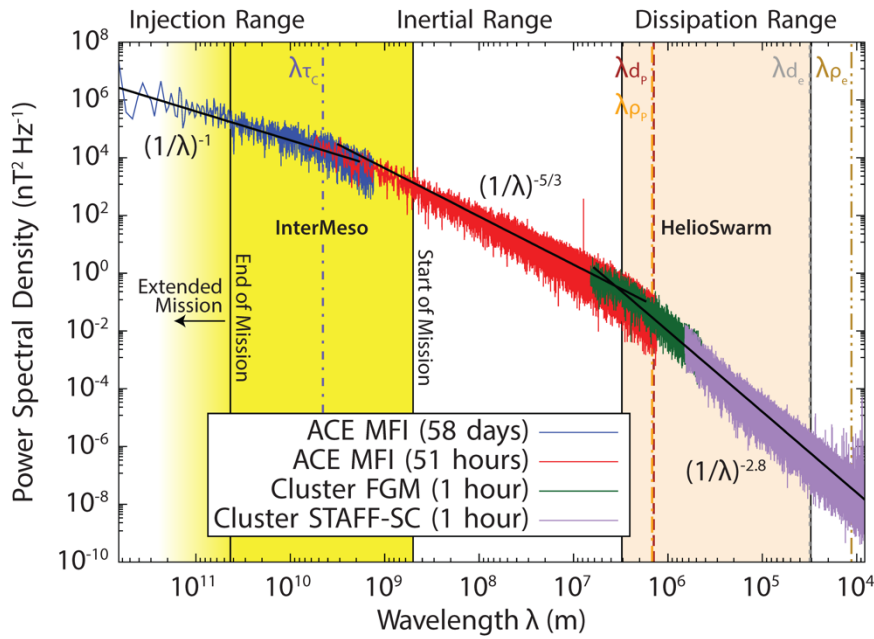
Goal	Obj.	Targeted Science Questions	Science Requirement	Measurement Requirements		
				Instruments Required	Inter-S/C Separations	Supporting Observations & Models
Investigating the fundamental mesoscale nature of the variable solar wind and impacts on particle acceleration, evolution, and transport.	1. Characterize and identify the origin of the mesoscale variability of the background solar wind and transient solar wind structures	1.1 What do the predominant scale sizes of the mesoscale solar wind and IMF reveal about the origins of these structures?	Conduct correlation analysis between spacecraft to determine the relevant scale sizes of the solar wind and fluxtubes. Need the ability to discern when changing between flux tubes: proton entropy, ion composition (tied to source), strahl electron properties.  Compare these scale sizes to expected solar variability (i.e., from granule/supergranule) vs. variability caused by dynamics in the inner heliosphere.	1, 2, 5, 7, 8	~5.5e5 km (flux tube cross-section)  ~2e6 - 6.5e6 km (granule/supergranule expansion at 1 au)	Solar Magnetogram DKIST (and other remote solar observatories)  Simulation advancements to resolve mesoscales Connectivity tools
		1.2 How do large-scale, persistent structures transition into dynamic turbulence?	Measure the azimuthal turbulence both within and across flux tubes encompassing the inertial and injection range.  Distinguish periods dominated by Alfvénic and slow-mode-like fluctuations.  Determine potential differences in azimuthal turbulence in longitudinally-separated parcels of plasma with composition measurements to identify different solar sources.	1, 2, 5, 7, 8	~4e5 km during calibration  1e6 - 1e9 km	PUNCH as input Inputs from L1 & L5  Global MHD & flux rope simulations
		1.3 How do mesoscale variabilities of the solar source manifest variations in seed particle properties and composition at 1au?	Determine the thermal/suprathermal ion spectra for main species (H, He, C, O, Fe) at high time cadence under both quiescent solar wind and across boundaries.  Use locally accelerated particles in ESP events or at CIRs to understand the role of seed particles in suprathermal-to-energetic ion composition.  Constrain mesoscale variations in these populations compared to solar wind structures (1.1). Determine if the mesoscale structuring of background solar wind maps to variability in seed particle populations. Determine if this points to variability of seed particle production in different solar structures (i.e., granules), or if this is from a broader reservoir.	1, 2, 3, 4, 5, 7	~5.5e5 km (flux tube cross-section)  ~2e6 - 6.5e6 km (granule/supergranule expansion at 1 au)	Solar imaging Solar magnetograms  CIR modeling (data assimilation into MHD)
		1.4 How does the mesoscale solar wind imprint itself on solar wind transients and interplanetary shock structures?	Determine morphological variability of transient structures (such as IP shocks, CMEs) expanding into the mesoscale solar wind. Determine if the variability across transient structures and shocks match that of the background solar wind (1.1) or if there are higher order perturbations or constraints.  Need ability to discern connection to source (shock or the Sun) through strahl/energetic electrons.	1, 2, 5, 6, 7, 8, 9	~5.5e5 km (flux tube cross-section)  ~2.2e6 km to ~6.5e6 km (granule/supergranule expansion at 1 au)	L5-based & STEREO heliographic imager  Ground-based radio arrays Simulation advancements to resolve mesoscales
	2. Characterize and understand the impact of these mesoscale variations on particle acceleration and transport	2.1 How does mesoscale solar wind variability and intermittency impact variability in particle acceleration?	Compare mesoscale variations in the suprathermal and energetic particle range to possible variations in wave instability occurrence predictions such as lower-hybrid waves, magnetosonic waves, ion-acoustic waves, electron cyclotron drift instability, and Langmuir waves.  Compare mesoscale variations in suprathermal and energetic particles to evidence of solar wind intermittency.	1, 2, 3, 4, 5, 6, 7, 8	~5.5e5 km (flux tube cross-section)  ~2e6 km - 6.5e6 km (granule/supergranule expansion at 1 au)	Simulation advancements to resolve mesoscales
		2.2 How does mesoscale structuring of solar wind transients and interplanetary shocks affect particle acceleration?	Determine mesoscale variations in longitudinally spread signatures of shock drift acceleration, diffusive shock acceleration, and fast Fermi processes to constrain the extent that mesoscale structuring may modulate these processes.  Need ability to discern plasma variations across shocks, and measure variations in A/Q for energetic ions.	1, 2, 3, 4, 5, 6, 7, 8	~3e5 - 1e7 km (nominal IP shock radius of curvature)  1's deg (CIR interfaces, connection to reverse shocks)	L5-based & STEREO heliographic imager Ground-based radio arrays Simulation advancements to resolve mesoscales
		2.3 How do mesoscale variations of the background solar wind impact the connectivity of an observer to the solar source or acceleration site?	Ability to examine the timing of hard X-rays and the timing of in situ electrons to determine transit time from most probable flare location and time.  Determine mesoscale variations in the injection timing and path length of energetic particles from source to observer and compare this to changes in particle population properties.  Determine mesoscale variability in ion drop out events.	1, 2, 3, 4, 5, 6, 7, 8, 9	1's deg (CIR interfaces, connection to reverse shocks)	Solar Magnetogram  Solar remote sensing observations Connectivity Tools Transport models
		2.4 How do mesoscale dynamics such as field line meandering and stochastic motion affect diffusion of particles in the solar wind?	Ability to detect energetic electrons with refined time resolution at multiple spacecraft separated in longitude.  Determine mesoscale variability in ion drop out events.	1, 2, 3, 4, 5, 6, 7, 8, 9	1's deg (full span ~10 deg)	Stochastic motion models with particle transport

554 Table 1: InterMeso STM. Instruments: (1) Bulk Solar Wind, (2) Thermal Ion Composition, (3)  
 555 Suprathermal Ion Composition, (4) Energetic Ion Composition, (5) Low-Energy Electrons, (6)  
 556 Energetic Electrons, (7) Magnetic Field, (8) Radio Waves, and (9) Hard X-rays.

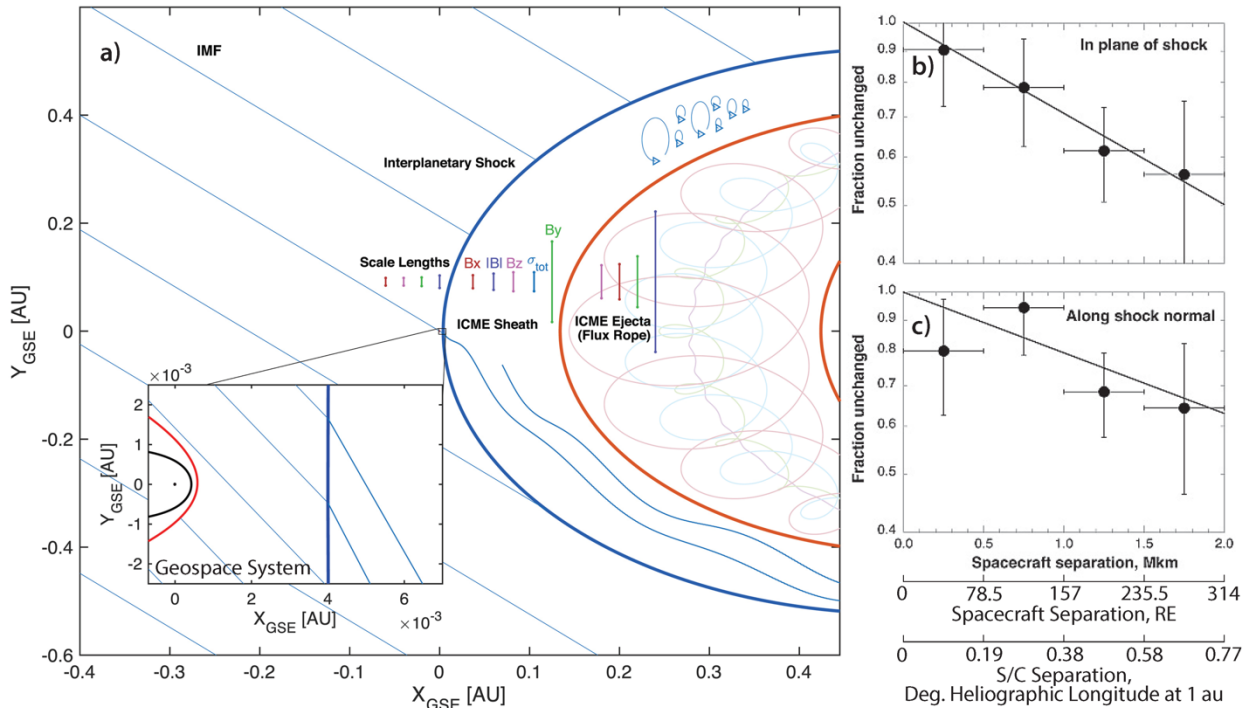
557 **11 Figures**



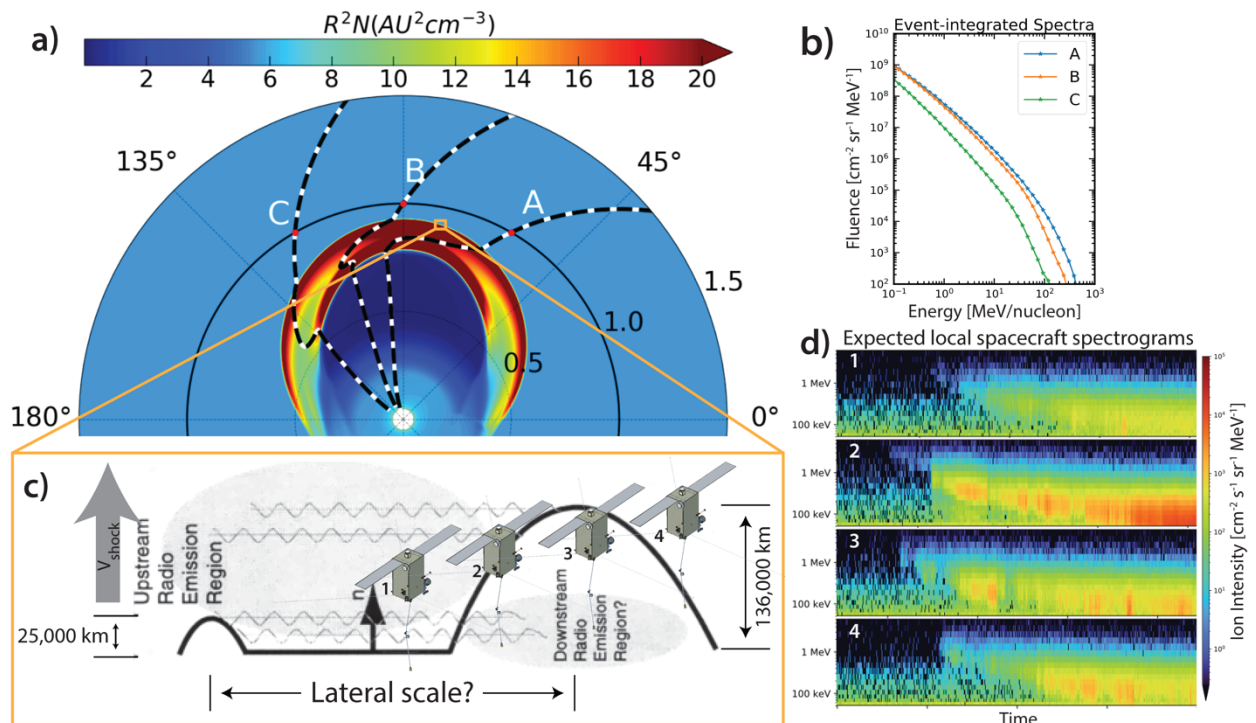
558 Figure 1: (a) Modeled representation of typical topology of magnetic structures on the solar surface  
 559 (adapted from *Bian and Li*, 2021). (b) As these structures convect outward to 1 au, they can undergo  
 560 processes such as expansion and stochastic meandering, but may preserve their characteristics as  
 561 granule to supergranule relics on the Sun (cartoon adapted from *Borovsky*, 2008). The nature of flux  
 562 tube structure at 1 au remains a fundamental open question.  
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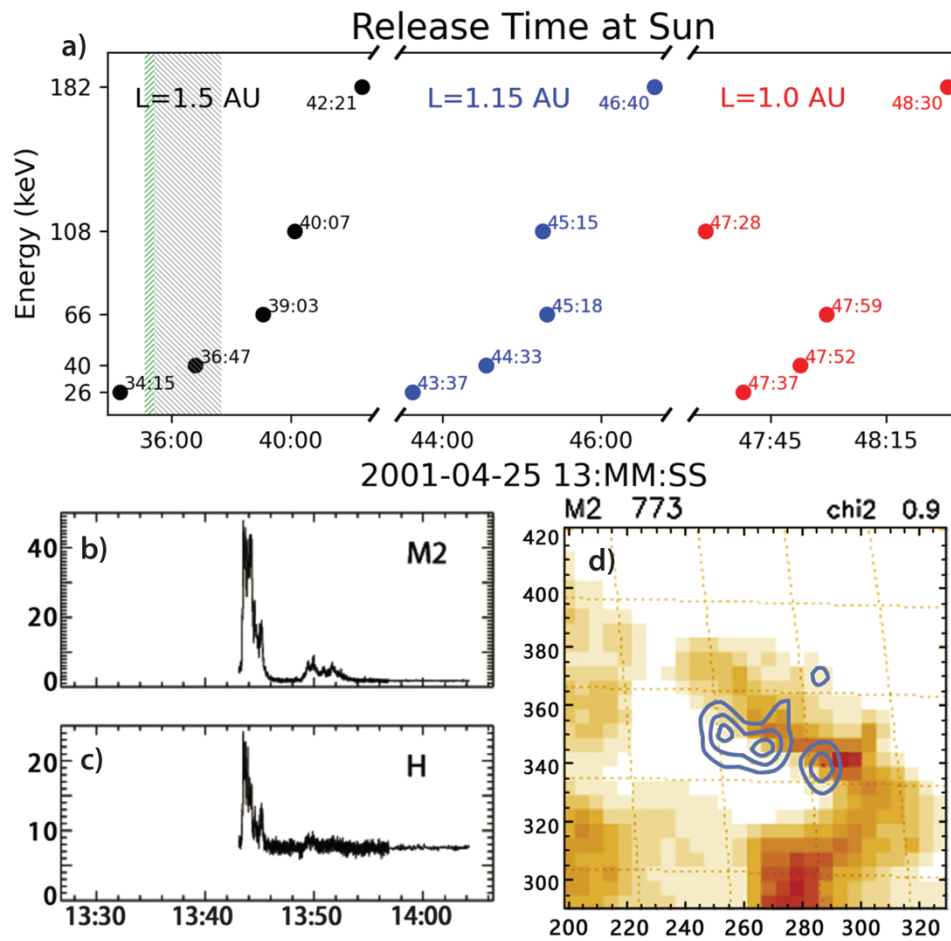
565  
 566 Figure 2: The power spectral density of magnetic field fluctuations at time of solar wind plasma beta  
 567 near 1 au shows two spectral breaks (see *Verscharen et al.*, 2019). The injection range is  
 568 characterized by a power spectral density slope of  $(1/\lambda)^{-1}$ , while the inertial range has a slope of  $(1/\lambda)^{-$   
 569  $^{5/3}$  in the spacecraft frame. The critical frequencies marked are the correlation length ( $f_{\tau_c}$ ), proton  
 570 inertial length ( $f_{dp}$ ) and ion-gyro scale ( $f_{pp}$ ), and the electron inertial length ( $f_{de}$ ) and electron-gyro  
 571 scale ( $f_{pe}$ ). While the upcoming HelioSwarm mission will explore the transition between the inertial and  
 572 dissipation ranges (tan region), InterMeso, for the first time, will reveal the spatiotemporal  
 573 dynamics between the injection and inertial range (yellow region). Adapted from *Verscharen et al.*  
 574 (2019).  
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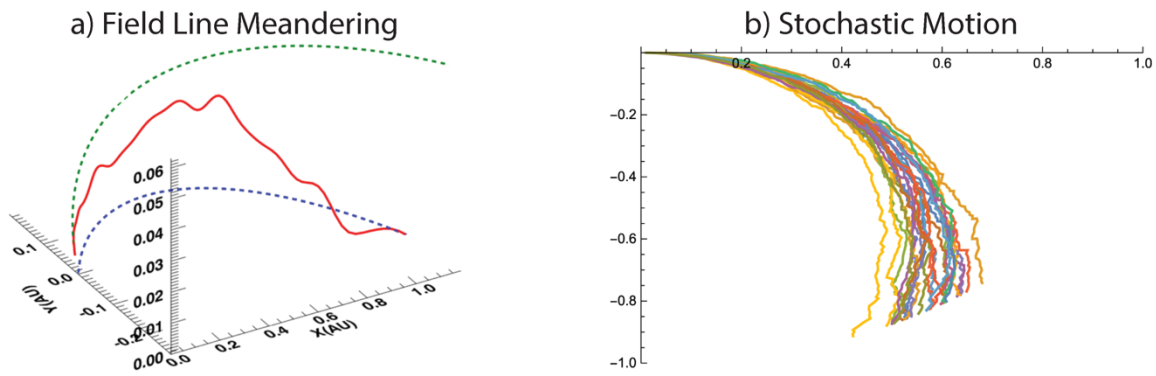
576  
 577 Figure 3: (a) Relevant scale sizes of CME structures at 1 au (from Ala-Lahti et al., 2020). The  
 578 fraction of 20-126 keV ion time intensity profiles that remain unchanged between two observing  
 579 points as a function of spacecraft separation both (b) in the plane of the shock and (c) along the shock  
 580 normal shows distinct variability that may be related to the variable scale sizes of CME sub-  
 581 structuring (from Neugebauer et al., 2006). The mesoscale structuring and subsequent effects on  
 582 particle acceleration remain outstanding open questions in Heliophysics.  
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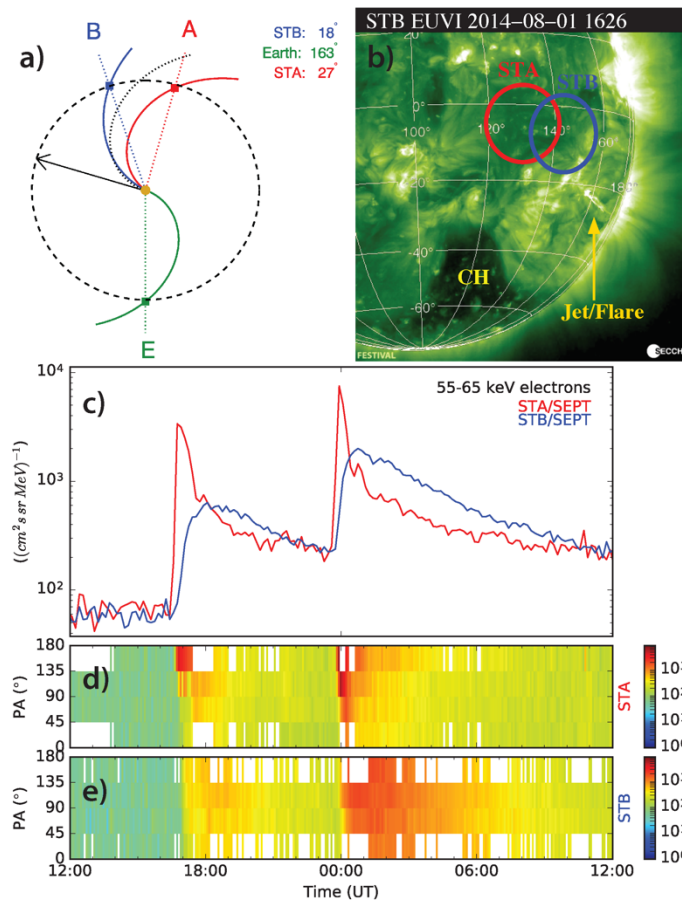
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 585 Figure 4: (a) Simulation of large-scale CME structure near 1 au yielding (b) longitudinal variations in  
 586 event-integrated particle spectra. (c) Mesoscale structuring of a CME shock inferred from radio wave  
 587 observations (adapted from *Bale et al.*, 1999) with (d) potential observations from four spacecraft  
 588 highlighting variations and impacts on particle dynamics. Mesoscale structuring of CMEs and effects  
 589 on particle acceleration are poorly understood and cannot be addressed by previous, current, or  
 590 planned missions.  
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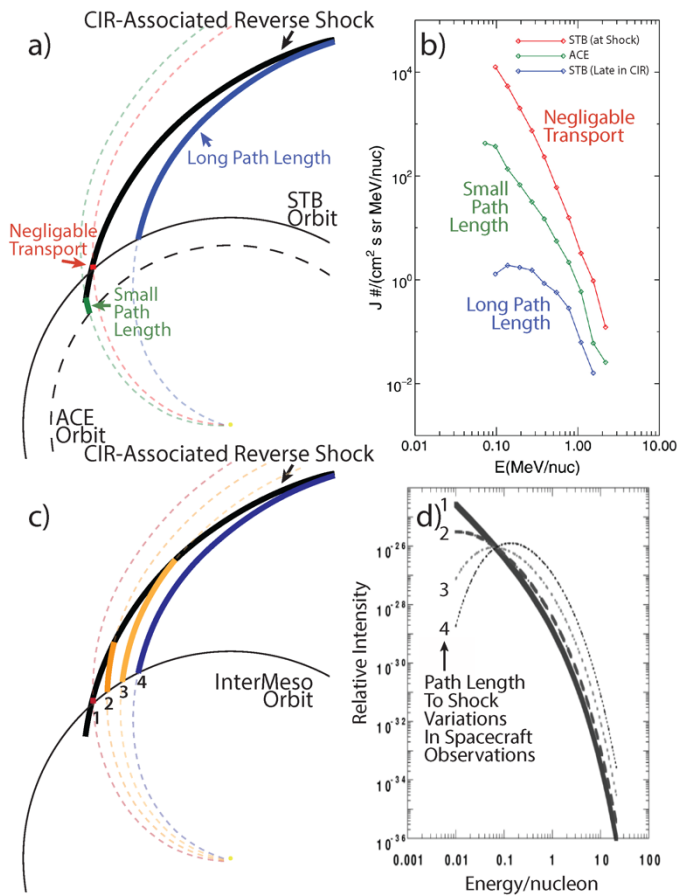
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 593 Figure 5: (a) Release times of energetic electrons for various path lengths and energies (dots) with  
 594 onset-to-peak (peak-to-end) phases of hard X-rays represented by the green and grey dashed regions.  
 595 (b, c, d) Yohkoh/HXT hard X-ray observations of the flare related to the energetic electron  
 596 observations with red contouring (d) from GOES soft X-ray. Adapted from *Li et al. (2020)*.  
 597 Comparisons between multi-point observations of hard X-ray and energetic electrons are needed to  
 598 constrain coronal dynamics.  
 599



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 601 Figure 6: (a) Field line meandering (from *Ashraf and Li, 2019*) and (b) stochastic motion (from *Bian*  
 602 *and Li, 2021*) are important mechanisms that cause the location of energetic electrons reaching 1 au  
 603 to dramatically deviate from the nominal Parker Spiral. The relative contributions of these processes  
 604 on particle transport are unknown, but critical to understanding particle dynamics.  
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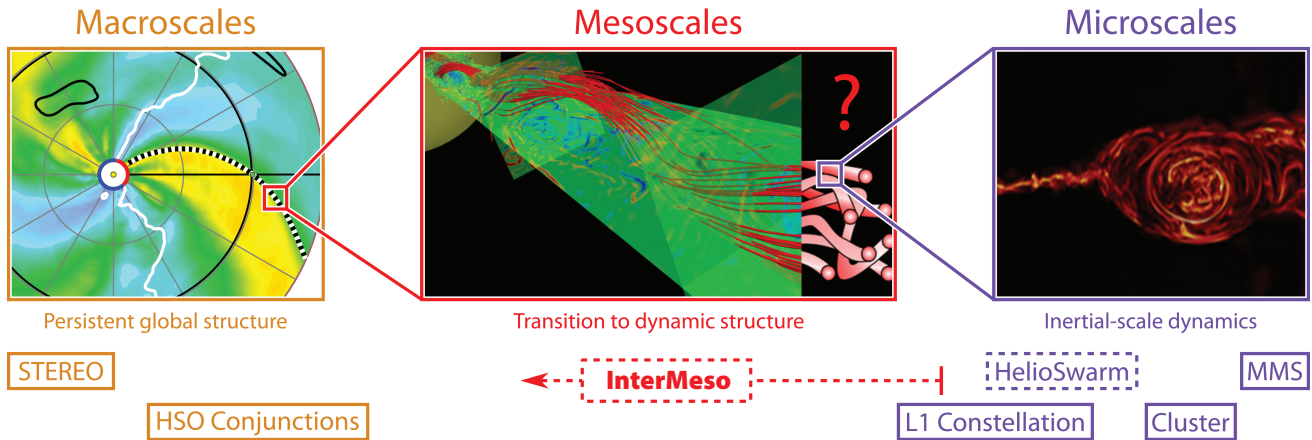


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 607 Figure 7: (a) The location of the STEREO spacecraft during a flare event and (b) the approximate  
 608 footpoint of both spacecraft relative to the flare site. (c-e) Energetic electron observations at both  
 609 spacecraft showing clear differences at 1 au over 10's degrees of longitude. Adapted from *Klassen et*  
 610 *al.* (2016). Transport of energetic electrons from flare sites to 1 au is complex, and requiring multi-  
 611 point measurements at mesoscale separations to constrain.  
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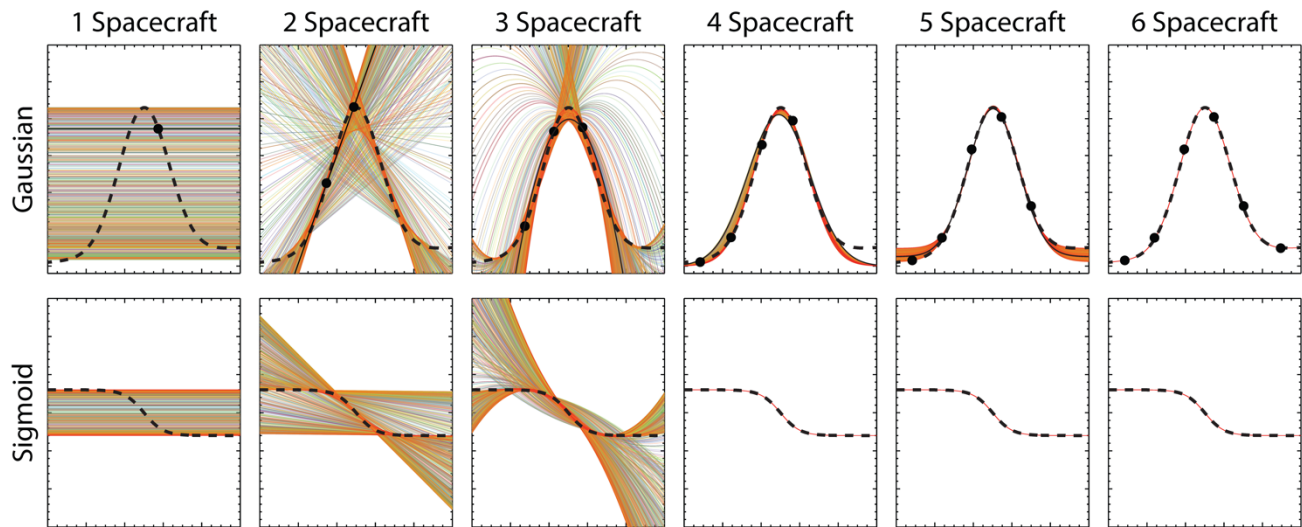
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 614 Figure 8: (a) Cartoon of the observation locations relative to a CIR-associated reverse shock inferred  
 615 from (b) energetic particle spectra (adapted from *Zhao et al.*, 2016). (c) Distributed observations at  
 616 mesoscale separations are needed to (d) constrain the transport effect on particle spectra without  
 617 ambiguities from temporal evolution (adapted from *Mason et al.*, 1999).  
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## Resolving mysteries of the mesoscale solar wind and transient structures

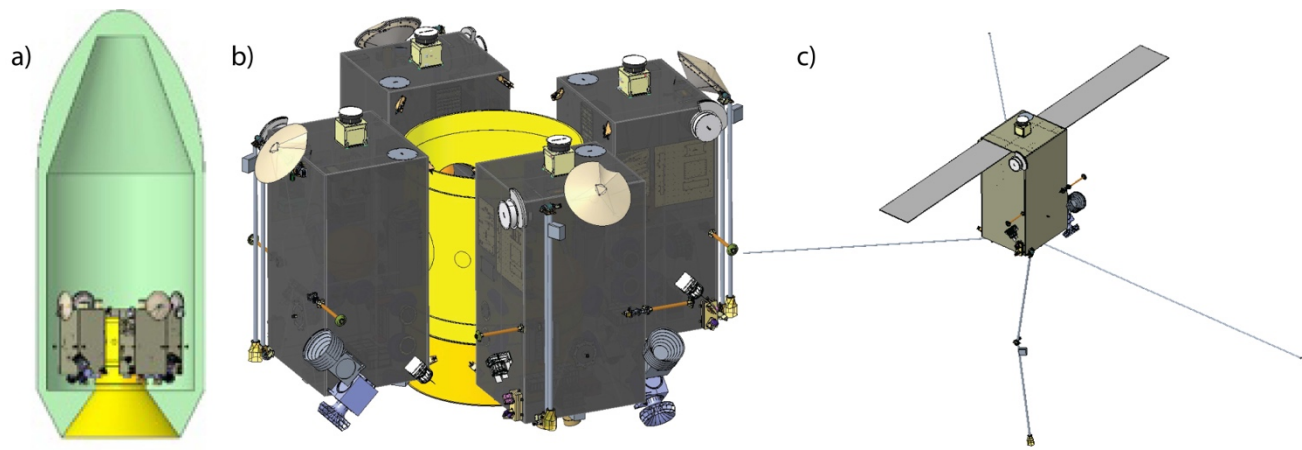


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Figure 9: Mesoscale dynamics are fundamental to solar wind dynamics, enabling interaction between macro and micro scale dynamics, but currently fall within an observational gap. Left panel from *Allen et al. (2020)*, middle panel adapted from *Maruca et al. (2021)*, and right panel from *Lazarian et al. (2020)*. InterMeso targets fundamental knowledge gaps of solar wind dynamics and structure at these critical scale lengths.

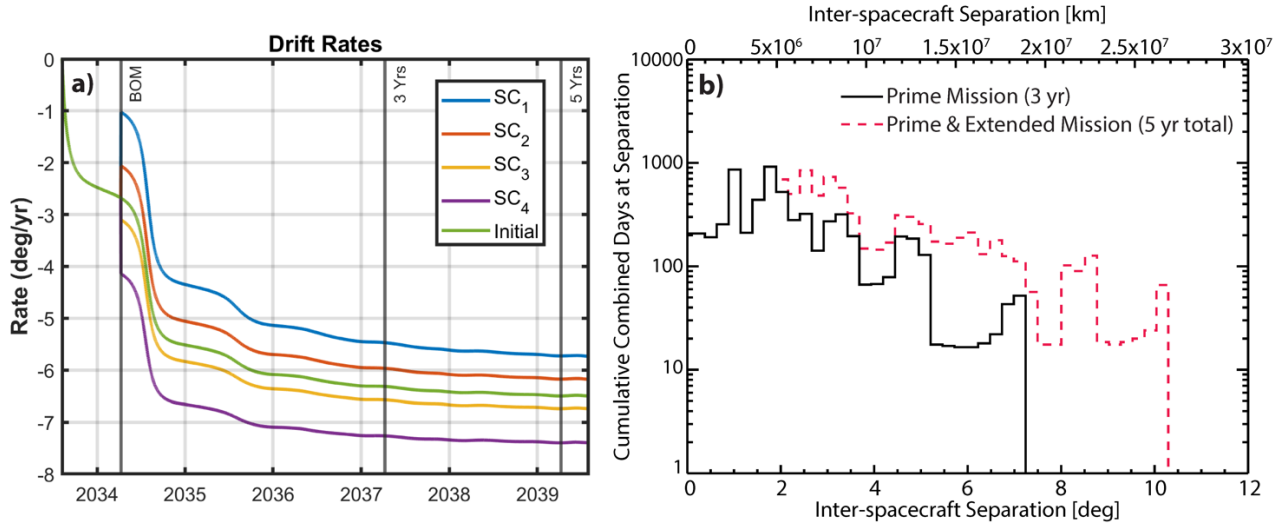


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 627 Figure 10: Randomized placement of a four-spacecraft constellation throughout both a Gaussian (top  
 628 row) and Sigmoidal (bottom row) distribution allow for different functional fits. Filled in circles in  
 629 the top row are examples of the 1000 random placements of spacecraft constellations resulting in the  
 630 shown colored fits. The solid black curve on the top row is the corresponding fit from the shown  
 631 example placements. The initial distribution the fits are attempting to recreate is shown by the dashed  
 632 curve. Four spacecraft are the optimal benefit-to-cost configuration for mesoscale variations.  
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635 Figure 11: Four “strawman” InterMeso spacecraft in a stowed configuration mounted on an ESPA-  
636 Grande equivalent viewed inside a 5-meter faring (a) and viewed without the faring (b). (c) A  
637 “strawman” InterMeso spacecraft in the deployed configuration.  
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Figure 12: (a) Change in Earth-relative drift rate over the course of the mission. The “initial” line represents the drift rate targeted by the launch vehicle. (b) Cumulative combined days the InterMeso spacecraft spend at different inter-spacecraft separations over the prime mission (black) plus extended mission (red). Cumulative combined days is the sum of the number of days spent at each separation for each spacecraft pair (i.e., spacecraft 1&2, 1&3, 1&4, 2&3, 2&4, and 3&4).