

1 **Decadal and annual variations in meteoric flux from Ulysses, Wind, and**  
2 **SOFIE observations**

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16 **Main Points:**

17 1) SOFIE, Wind, and Ulysses give consistent estimates of the meteoric influx at Earth.

18 2) Annual and decadal variations in Wind interstellar dust observations agree with IMEX  
19 simulations.

20 3) Both interstellar and interplanetary dust are correlated to the 22-yr. solar magnetic cycle.

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22 Keywords: Meteoric influx, SOFIE, Wind, Ulysses, interstellar dust, meteoric smoke

23 **Abstract.** Our solar system is filled with meteoric particles, or cosmic dust, which is either  
24 interplanetary or interstellar in origin. Interstellar dust (ISD) enters the heliosphere due to the  
25 relative motion of the sun and the interstellar flow. Interplanetary dust (IPD) comes primarily  
26 from asteroid collisions or comet sublimation, and comprises the bulk of material entering  
27 Earth's atmosphere. This study examines variations in ISD and the IPD flux at Earth using  
28 observations from three different satellite techniques. First are size-resolved in situ meteoroid  
29 detections by the Ulysses spacecraft, and second are in situ indirect dust observations by Wind.  
30 Third are measurements of meteoric smoke in the mesosphere by the Solar Occultation For Ice  
31 Experiment (SOFIE). Wind and Ulysses observations are sorted into the interstellar and  
32 interplanetary components. Wind ISD show the anticipated correlation to the 22-yr. solar  
33 magnetic cycle, and are consistent with model predictions of ISD. Because Wind does not  
34 discriminate particle size, the IPD measurements were interpreted using meteoric mass  
35 distributions from Ulysses observations and from different models. Wind observations during  
36 2007-2020 indicate a total meteoric influx at Earth of 22 metric tons per day ( $\text{t d}^{-1}$ ), in reasonable  
37 agreement with long-term averages from SOFIE ( $25 \text{ t d}^{-1}$ ) and Ulysses ( $32 \text{ t d}^{-1}$ ). The SOFIE and  
38 Wind influx time series both show an unexpected correlation to the 22-yr. solar cycle. This  
39 relationship could be an artifact, or may indicate that IPD responds to changes in the solar  
40 magnetic field.

41 **1. Introduction**

42           The solar system is filled with meteoric particles, or cosmic dust, which is interplanetary  
43 or interstellar in origin. Interplanetary dust (IPD) comes primarily from asteroid collisions or  
44 comet sublimation, and is typically bound to solar orbits on the ecliptic plane. Interstellar dust  
45 (ISD) enters the heliosphere due to the relative motion of the sun within the local interstellar  
46 cloud [Krueger *et al.*, 2019; Sterken *et al.*, 2019]. Meteoroids are constantly entering Earth’s  
47 atmosphere, with larger IPD particles (roughly 1 - 100  $\mu\text{m}$  radius) dominating the mass influx.  
48 Frictional heating during entry vaporizes a fraction of the particles at altitudes from  $\sim 80$  to 100  
49 km, and ablation products combine to form nanometer sized meteoric smoke particles that reside  
50 in the mesosphere and stratosphere [Plane *et al.*, 2012; Hervig *et al.*, 2017]. Smoke in the  
51 mesosphere is reduced during polar summer due to transport by the mesospheric meridional  
52 circulation, as indicated by models [Megner *et al.*, 2008; Bardeen *et al.*, 2008] and satellite  
53 observations [Hervig *et al.*, 2009]. The annual smoke variation occurs despite the seasonal  
54 change in meteor influx at polar latitudes, where the highest influx occurs in summer as shown  
55 by radars [Janches *et al.*, 2004; Singer *et al.*, 2004; Reid *et al.*, 2006]. A comprehensive  
56 assessment of the spatial and temporal dependence in meteoric influx can be found in Fentzke *et*  
57 *al.* [2009]. Recent estimates of the meteoric influx into Earth’s atmosphere are from  $\sim 15$  - 60  
58 metric tons per day ( $\text{t d}^{-1}$ ) [e.g., Carrillo-Sánchez *et al.*, 2020], an improvement over the  
59 bewildering range of previous decades ( $1$  -  $300 \text{ t d}^{-1}$ ) [e.g., Plane *et al.*, 2012]. The meteoric  
60 influx at Earth has implications for atmospheric chemistry, aerosol processes, and ocean  
61 productivity [e.g., Rudraswami *et al.*, 2021], motivating further improvements to the  
62 understanding of influx and its variability.

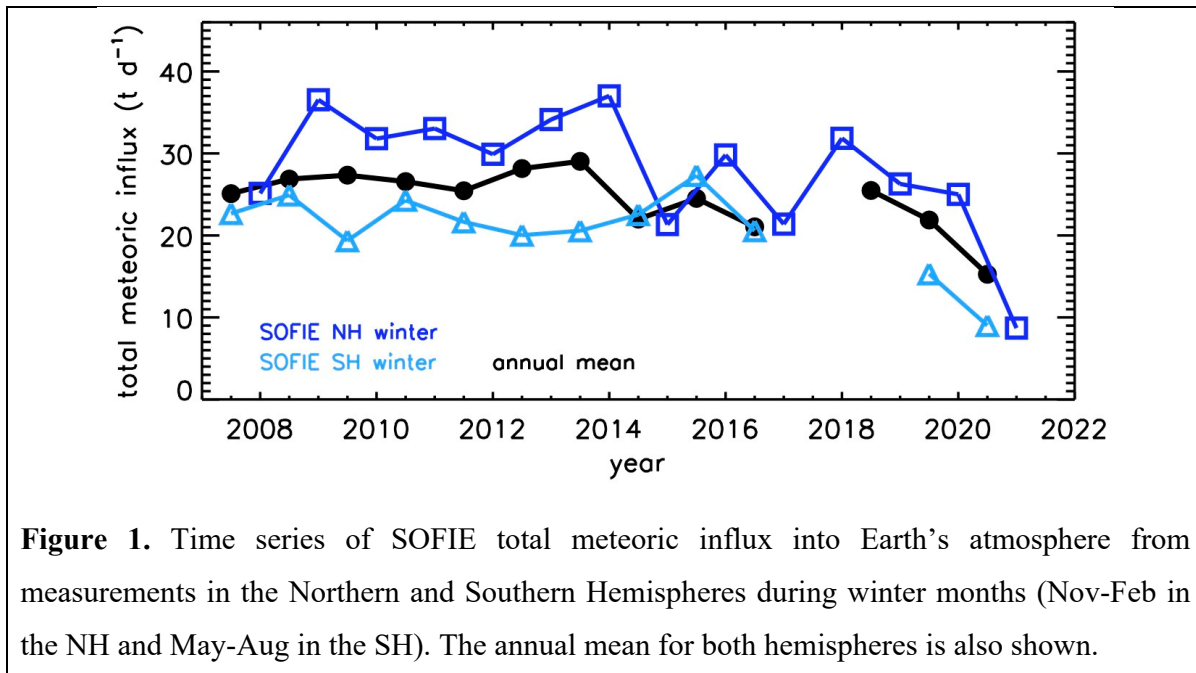
63           The present study examines meteoric flux in the near-Earth environment using  
64 observations from Ulysses, Wind, and the Solar Occultation For Ice Experiment (SOFIE). The  
65 Wind and Ulysses spacecraft offer long-term records of in situ dust measurements, which are  
66 related here to meteoric smoke measurements from the SOFIE satellite instrument. Interpreting  
67 the Ulysses, Wind, and SOFIE measurements requires an understanding of the meteoric mass  
68 distribution, and dust enhancement due to Earth’s gravity and size. The Ulysses and Wind  
69 observations contain both IPD and ISD, and separating these is important for estimating the  
70 meteoric influx at Earth. The Wind ISD results are validated through comparisons with ISD  
71 simulations from the Interplanetary Meteoroid environment for EXploration (IMEX) model  
72 [*Sterken et al.*, 2015; *Strub et al.*, 2019].

## 73 **2. SOFIE observations**

74           SOFIE has conducted solar occultation measurements from the Aeronomy of Ice in the  
75 Mesosphere (AIM) satellite since 2007 [*Russell et al.*, 2009]. The measurements are used to  
76 retrieve vertical profiles of temperature, five gases (O<sub>3</sub>, H<sub>2</sub>O, CO<sub>2</sub>, CH<sub>4</sub>, and NO), polar  
77 mesospheric cloud (PMC) extinction at 11 wavelengths, and meteoric smoke extinction at three  
78 wavelengths (330, 867, and 1037 nm). SOFIE observes primarily polar latitudes, with the  
79 exception of 2017 - 2019 when orbital progression caused an excursion through the tropics and a  
80 change from sunsets in the Southern Hemisphere (SH) to the Northern Hemisphere (NH) (vice  
81 versa for sunrises, see *Hervig et al.*, 2021 for details). The current SOFIE data is version 1.3  
82 which is available online ([sofie.gats-inc.com](http://sofie.gats-inc.com)).

83           SOFIE smoke measurements have been used to characterize the variation of smoke in  
84 height and time, and revealed the chemical composition of smoke [*Hervig et al.*, 2009; 2017].  
85 The smoke extinctions used here are monthly zonal means, avoiding summer measurements

86 when PMCs contaminate the smoke signal [Hervig *et al.*, 2012]. Extinction is converted to  
 87 volume density for a smoke composition of olivine ( $\text{Mg}_{2x}\text{Fe}_{2-2x}\text{SiO}_4$ ,  $x = 0.4$ ), which is optically  
 88 detected by SOFIE. Volume density is then used to derive the ablated meteoric influx through  
 89 comparisons with smoke simulations from the Whole Atmosphere Community Climate Model  
 90 (WACCM) [Bardeen *et al.*, 2008; Hervig *et al.*, 2017, 2021]. SOFIE results during 2007-2020  
 91 indicate a global mean ablated influx into Earth's atmosphere of  $7 \pm 2$  metric tons per day ( $\text{t d}^{-1}$ ).  
 92 Since only  $\sim 30\%$  of incoming meteoroids are ablated [Carrillo-Sánchez *et al.*, 2020], the  
 93 corresponding total influx (ablated plus surviving material) is  $25 \pm 7 \text{ t d}^{-1}$ . The influx versus time  
 94 from SOFIE observations is shown in Figure 1, where the results indicate year-to-year variations  
 95 and greater influx in the NH than in the SH. The hemispheric difference is still not understood,  
 96 but could indicate an asymmetry in meteoric influx that is not represented in WACCM. The  
 97 results below consider the meteoric influx from SOFIE as the average of the NH and SH values.

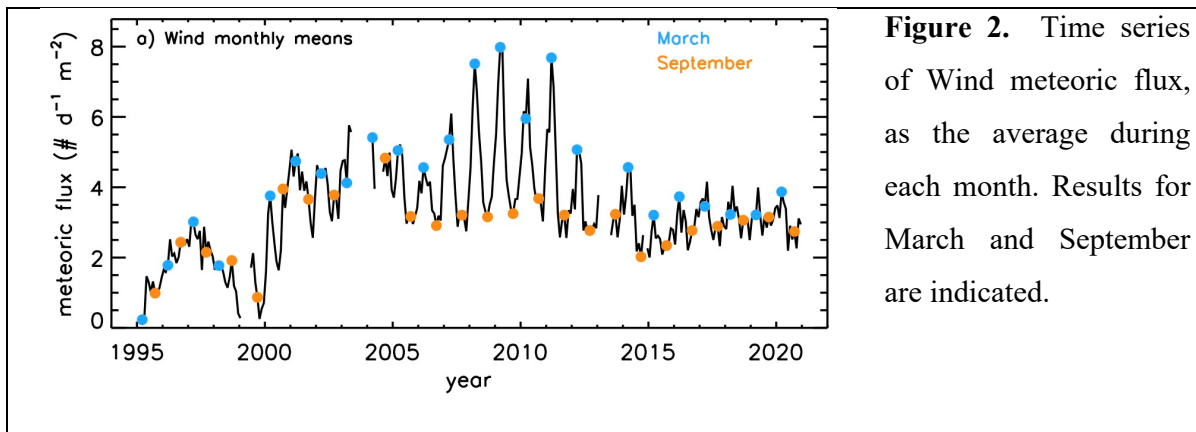


### 98 3. Wind observations

99           The Wind spacecraft was launched in 1994 to quantify the dynamics of the Sun-Earth  
100 system [Wilson *et al.*, 2021]. Wind operates within 1° of the ecliptic plane and has orbited the  
101 first Lagrange point (L1) since July 2004 ( $\sim 1.5 \times 10^6$  km sunward from Earth). Prior to 2004,  
102 Wind conducted a variety of orbital maneuvers, including petal orbits through the  
103 magnetosphere, lunar flybys, and an excursion to the second Lagrange point [Malaspina and  
104 Wilson, 2016; Wilson *et al.*, 2021]. Many of these periods are not useful for dust measurements,  
105 and were screened from the analyses presented here.

106           Meteoroids are detected when they collide with Wind and a fraction of the spacecraft  
107 body is vaporized and ionized [Mann, 2019]. The resulting plasma perturbs the electric potential  
108 of spacecraft surfaces [Shen *et al.*, 2021], which is observed by the WAVES electric field  
109 antennas and recorded by the Time Domain Sampler (TDS) [Bougeret *et al.*, 1995]. A similar  
110 approach has been used for dust measurements by other spacecraft including Voyager [Gurnett  
111 *et al.*, 1983] and the Mars Atmosphere and Volatile Evolution Mission (MAVEN) [Andersson *et al.*,  
112 2015]. The Wind dust detector area is the cross-sectional area of the cylindrical spacecraft  
113 body (1.8 m height  $\times$  2.4 m diameter), or 4.3 m<sup>2</sup>. Wind is estimated to be sensitive to meteoroids  
114 with radii ( $r$ ) of 0.1 - 11  $\mu\text{m}$ , or  $10^{-14}$  to  $10^{-8}$  g in mass ( $m$ ) for a dust density of  $\rho = 2.65 \text{ g cm}^{-3}$ ,  
115 but cannot resolve the size of individual impactors. Malaspina *et al.* [2014] noted that the lower  
116 and upper mass bounds are uncertain by a factor of 10 or more, due to observational  
117 uncertainties and assumptions in the measurement interpretation. Wind reports the number of  
118 dust detections per day, which represents particles with  $m$  from  $10^{-14}$  to  $10^{-8}$  g. The Wind dust  
119 observations are discussed in detail by Malaspina *et al.* [2014], Meyer-Vernet *et al.* [2014],  
120 Kellogg *et al.* [2016], and Wood *et al.* [2015]. Malaspina and Wilson [2016] describe the

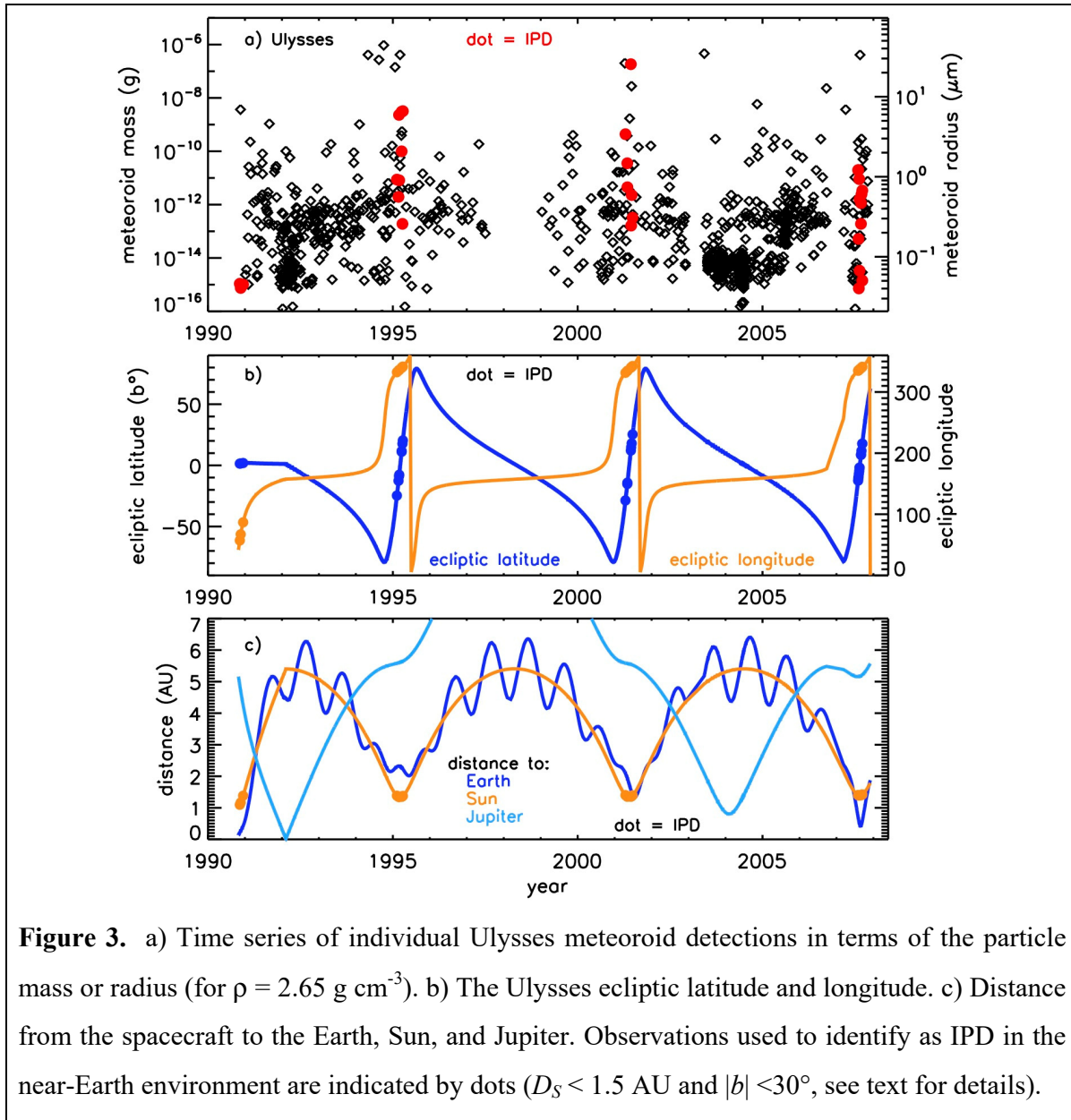
121 archived data which are available online (cdaweb.gsfc.nasa.gov). The Wind dust record is  
 122 summarized in Figure 2, where the monthly average meteoric flux is shown. For this work,  
 123 observations were ignored when there were known problems, or complications due to special  
 124 spacecraft maneuvers [see *Malaspina and Wilson, 2016; Wilson et al., 2021*]. The results  
 125 indicate pronounced decadal and annual variations (e.g., more dust in March than September),  
 126 which are investigated below.



#### 127 4. Ulysses Observations

128 Ulysses operated during 1990 - 2007, and was the first spacecraft to conduct a polar orbit  
 129 of the Sun. An in situ dust detector used impact ionization to measure the mass of individual  
 130 particles with  $m$  from  $\sim 10^{-16}$  -  $10^{-6}$  g [*Grün et al., 1992; Krüger et al., 2006; 2019*]. The detector  
 131 sensitive area is quoted as a maximum of  $0.02 \text{ m}^2$  by *Krüger et al. [2015]*, which is the value  
 132 used here. The Ulysses dust observations are illustrated in Figure 3a, where the mass and radius  
 133 of each particle are shown. The reported particle mass uncertainties are typically a factor of 5 to  
 134 10, due to measurement errors combined with uncertainties in the interpretation. Due to the high-  
 135 inclination polar orbit of the Sun, the Ulysses data set consists mostly of ISD. *Strub et al. [2015]*  
 136 describe the criteria for identifying ISD in Ulysses, and the present study considered these (in  
 137 reverse) to find IPD in the Ulysses record. Particles detected at high ecliptic latitudes ( $b$ , Figure

138 3b) are most likely interstellar in origin [*Krüger et al.*, 2006], since interplanetary dust is  
139 concentrated near the ecliptic plane [e.g., *Soja et al.*, 2019]. Exceptions can occur, however, as  
140 dust from Halley type comets and Oort cloud comets has been detected far above the ecliptic  
141 plane. IPD in the near-Earth environment was identified by considering 1) spacecraft - sun  
142 distances ( $D_S$ ) of less than 1.5 AU, and 2) relatively low ecliptic latitudes ( $|b| < 30^\circ$ ). These  
143 criteria limit the Ulysses observations to those near perihelion and also exclude the Jupiter flybys  
144 (see Figure 3). Note that filtering by rotation angle will not separate IPD from ISD because the  
145 interstellar flow direction was aligned with the heliocentric prograde motion (i.e., the dominant  
146 IPD direction) near perihelion [*Strub et al.*, 2015]. The selection criteria here yields dust  
147 observations at an average  $D_S$  of 1.36 AU. These observations occurred during four periods (see  
148 Figure 3) comprising a total of 0.76 years, which is the observing interval used for the Ulysses  
149 IPD flux quantities below.



**Figure 3.** a) Time series of individual Ulysses meteoroid detections in terms of the particle mass or radius (for  $\rho = 2.65 \text{ g cm}^{-3}$ ). b) The Ulysses ecliptic latitude and longitude. c) Distance from the spacecraft to the Earth, Sun, and Jupiter. Observations used to identify as IPD in the near-Earth environment are indicated by dots ( $D_S < 1.5 \text{ AU}$  and  $|b| < 30^\circ$ , see text for details).

150 **4. Meteoric mass distributions**

151 Interpreting the various measurements and relating them to each other requires an  
 152 understanding of the meteoric mass distribution. The mass of meteoroids spans many orders of  
 153 magnitude, with ISD ranging from  $10^{-16} - 10^{-10} \text{ g}$  ( $r \approx 0.1 - 2 \mu\text{m}$ , for  $\rho = 2.65 \text{ g cm}^{-3}$ ) and IPDs  
 154 spanning roughly  $10^{-15}$  to  $10 \text{ g}$  ( $r \approx 0.1 \mu\text{m}$  to  $1 \text{ cm}$ ) [e.g., Krüger *et al.*, 2019; Sterken *et al.*,

155 2015]. Visual meteors are roughly  $10^{-2}$  to  $10^3$  g ( $r \approx 0.1 - 5$  cm) but contribute little to the total  
 156 meteoric mass influx at Earth, and larger bodies ( $m > 1$  kg) appear only on geologic time scales.  
 157 *Grün et al.* [1985] (G85) described a meteoric mass distribution based on spacecraft in situ  
 158 observations, lunar crater analysis, and photometric measurements of the Zodiacal light. The  
 159 G85 expression yields the cumulative dust flux in free space,  $n_C(m)$  ( $\text{g m}^{-2} \text{s}^{-1}$ ), for a given  $m$  (g)  
 160 (i.e., the number of particles with mass  $> m$ ),

$$161 \quad n_C(m) = (2.2 \times 10^3 m^{0.306} + 15)^{-4.38} + 1.3 \times 10^{-9} (m + 10^{11} m^2 + 10^{27} m^4)^{-0.36} +$$

$$162 \quad 1.3 \times 10^{-16} (m + 106 m^2)^{-0.85} \quad (1)$$

163 The number of particles with a given  $m$ ,  $n(m)$ , is found from  $n(m_i) = n_C(m_i) - n_C(m_{i+1})$ . The  
 164 number distribution is easily converted to a mass distribution,  $f(m) = m n(m)$  [*Grün et al.*, 1985].

165 The G85 expression and observations from Wind and Ulysses describe the meteoric flux  
 166 in free space. Relating this to the flux into Earth's atmosphere requires consideration of Earth's  
 167 gravity and surface area. The focusing effect of planetary gravity on meteoric dust was described  
 168 by *Drolshagen et al.* [2017] as the enhancement factor

$$169 \quad H_F = v^2 / (v^2 - v_{esc}^2) \quad (2)$$

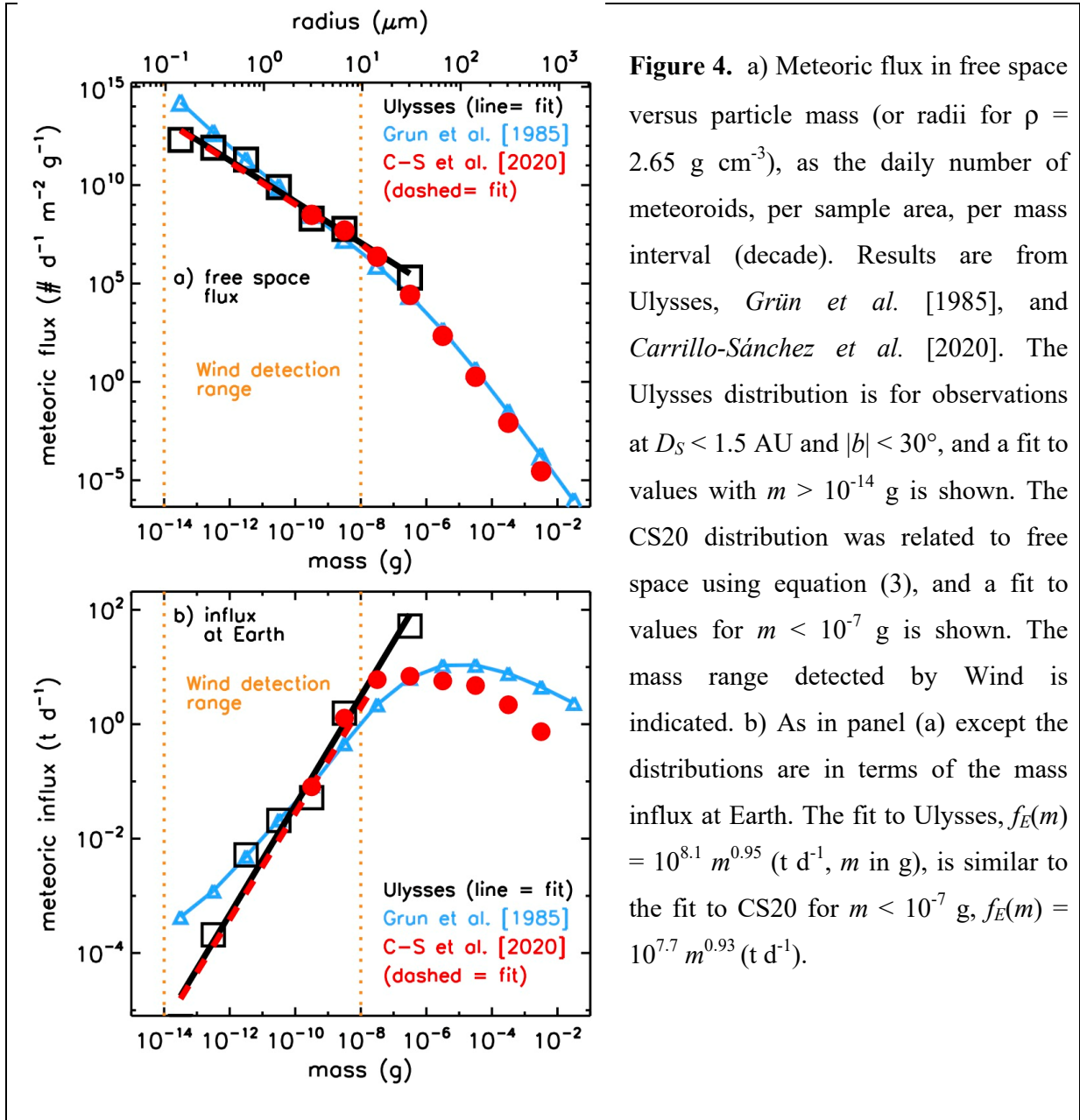
170 where  $v$  is the dust velocity (far from Earth) and  $v_{esc}$  is the escape velocity ( $11.1 \text{ km s}^{-1}$  for Earth)  
 171 [see also *Jones and Poole*, 2007]. The present study assumed a mean dust velocity of  $17 \text{ km s}^{-1}$   
 172 [e.g., *Borin et al.*, 2017] giving  $H_F = 1.74$ , which is the value used below. The statistical  
 173 uncertainty in  $H_F$  is  $\sim 20\%$ , for considering  $v$  in the range of  $14 - 28 \text{ km s}^{-1}$ . Earth's surface area is  
 174 calculated at the typical meteor ablation altitude of  $100 \text{ km}$  ( $S_E = 5.26 \times 10^{14} \text{ m}^2$ ). With these  
 175 considerations the total meteoric influx at Earth,  $F_E$ , is given by

$$176 \quad F_E = F H_F S_E \quad (3)$$

177 where  $F$  is the total flux in free space,  $F = \sum f(m)$ . Note that this expression is adaptable to the  
178 numeric or mass flux distributions as well (e.g.,  $f_E(m) = f(m) H_F S_E$ , assuming constant  $v$ ).

179 Ulysses observations in the near-Earth environment (see Section 4) were used to  
180 characterize the IPD mass distribution. The numeric flux distribution from Ulysses is compared  
181 to G85 in Figure 4a, along with results from *Carrillo-Sánchez et al.* [2020] (CS20) which  
182 describe  $f(m)$  at Earth based on observations and models. The CS20 distribution (their Figure 1a)  
183 was related to free space here using equation (3). The distributions from Ulysses and CS20 are in  
184 good agreement in the range where they overlap, and a log-linear fit to CS20 for  $m < 10^{-7}$  g is  
185 very close to the Ulysses values for smaller particles. The meteoric mass distributions from G85  
186 and Ulysses were related to Earth using equation (3), and are compared to the CS20 curve in  
187 Figure 4b. The distributions in terms of mass reveal some differences that are not apparent in the  
188 number distributions. In particular, both Ulysses and CS20 indicate lower influx than G85 for the  
189 smallest particles ( $m < 10^{-10}$  g), while none of the distributions agree for the largest meteoroids  
190 ( $m > 10^{-7}$  g). In the case of Ulysses, there are no observations of  $m > 10^{-6}$  g, although the  
191 meteoric flux in the interval  $10^{-7}$  to  $10^{-6}$  g is  $\sim 4$  times greater than for CS20 or G85. Note that the  
192 meteoric mass influx at Earth is dominated by particles in the range of  $10^{-8}$  to  $10^{-3}$  g. The sum of  
193 the distributions in Figure 4b gives the total meteoric influx at Earth, with  $F_E = 37$  t d $^{-1}$  for G85  
194 and 28 t d $^{-1}$  from CS20. Because Ulysses resolves the meteoroid mass, the individual  
195 observations can be added to give  $F_E = 32$  t d $^{-1}$ . This is slightly more accurate than summing the  
196 histogram values, which assign the average mass to each interval. The reported Ulysses particle  
197 mass uncertainties are typically a factor of 5 to 10, and this error dominates the uncertainty in  $F_E$ .  
198 Even with the advantage of counting statistics (error reduction by  $N^{1/2}$ ) the Ulysses  $F_E$

199 uncertainty is large at  $76 \text{ t d}^{-1}$ . The above estimates are in good agreement with SOFIE  
 200 observations which give  $F_E = 25 \pm 7 \text{ t d}^{-1}$  on average [Hervig *et al.*, 2021].



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The meteoric mass and number distributions (Figure 4) were used to interpret the Wind observations, which do not resolve the mass of an impactor, but rather indicate the total flux of

204 particles within the Wind mass detection range,  $N_W = \sum_{m_1}^{m_2} n(m)$ , with  $m_1 = 10^{-14}$  g and  $m_2 = 10^{-8}$   
 205 g.  $N_W$  can be computed from the meteoric size distributions (e.g., Figure 4a), and used to provide  
 206 and an estimate of the total influx at Earth from Wind,

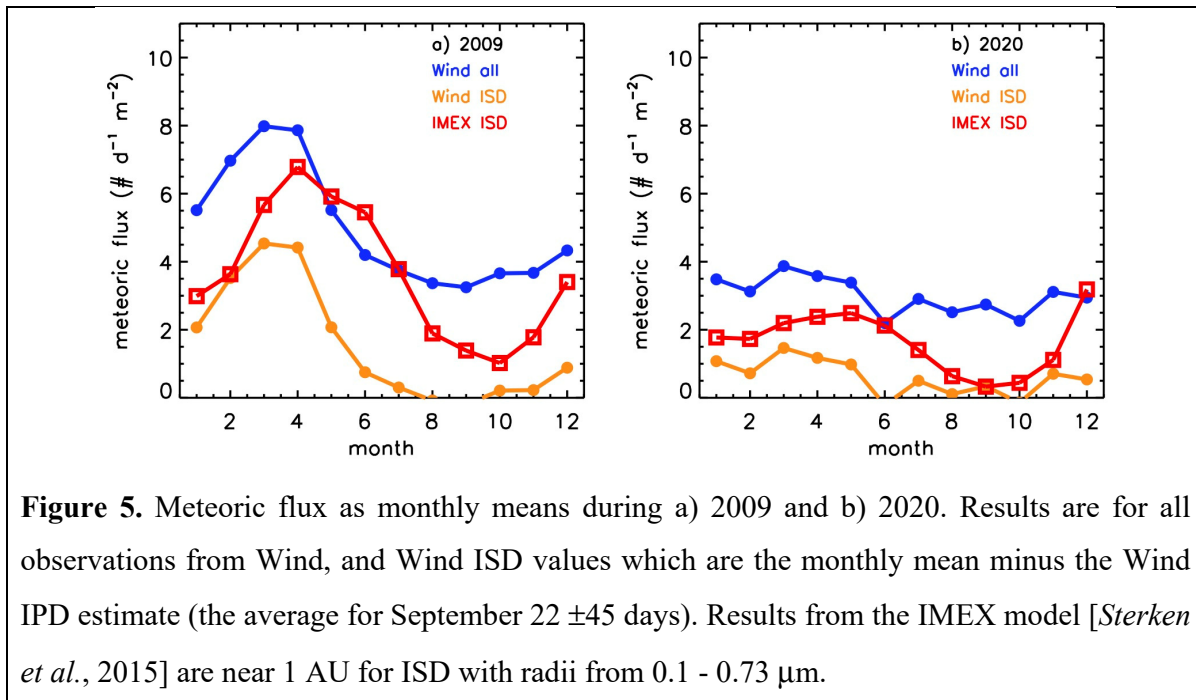
$$207 \quad F_E = N_W \sum f_E(m) / \sum_{m_1}^{m_2} n(m) = N_W C \quad (4)$$

208 In this expression  $N_W$  is the Wind observation, and  $C = \sum f_E(m) / \sum_{m_1}^{m_2} n(m)$  can be computed  
 209 from a given modeled or measured mass distribution. Note that the summation of  $f_E(m)$  is for the  
 210 relevant range of IPD (roughly  $10^{-9}$  to  $10^{-2}$  g). For the mass distributions in Figure 4,  $C = 2.1$  (t  
 211  $m^2$ ) for G85 and  $C = 11.3$  (t  $m^2$ ) for CS20. One error component in  $C$  arise from the factor of 10  
 212 uncertainties in the Wind detection limits ( $m_1$  and  $m_2$ ) [Malaspina *et al.*, 2014]. This error  
 213 component was estimated to be  $\sim 29\%$ , using the Ulysses mass distributions and perturbing  $m_1$   
 214 and  $m_2$  alternately by factors of 10 and 0.1. The total uncertainty in  $C$  from the G85 distribution  
 215 is difficult to estimate, but it is at least 35% due to the combined errors in  $H_F$  (20%) and in  $m_1$   
 216 and  $m_2$  (29%). The uncertainty in  $C$  from CS20 is 70% due to errors in  $F_E$ ,  $H_F$ , and  $m_1$  and  $m_2$ .  
 217 The value for Ulysses was determined directly by summing the individual observations, giving  $C$   
 218 = 9.9 (t  $m^2$ ). The uncertainty in  $C$  from Ulysses is close to a factor of  $\sim 3$ , due to the uncertainty  
 219 in  $F_E$  (discussed above) combined with the other terms. The value of  $C$  based on G85 is lower  
 220 than from Ulysses or CS20 because the G85 curve indicates more dust at the smallest sizes. For  
 221 comparison, the G85 distribution gives  $N_W = 21.7$   $d^{-1} m^{-2}$ , compared to 4.1  $d^{-1} m^{-2}$  from Ulysses  
 222 and 2.4  $d^{-1} m^{-2}$  from CS20. While the small particles contribute little to the total mass, they  
 223 contribute greatly to the total number of meteoroids (Figure 4). Equation (4) is used below with  
 224 Wind observations to examine the IPD influx at Earth.

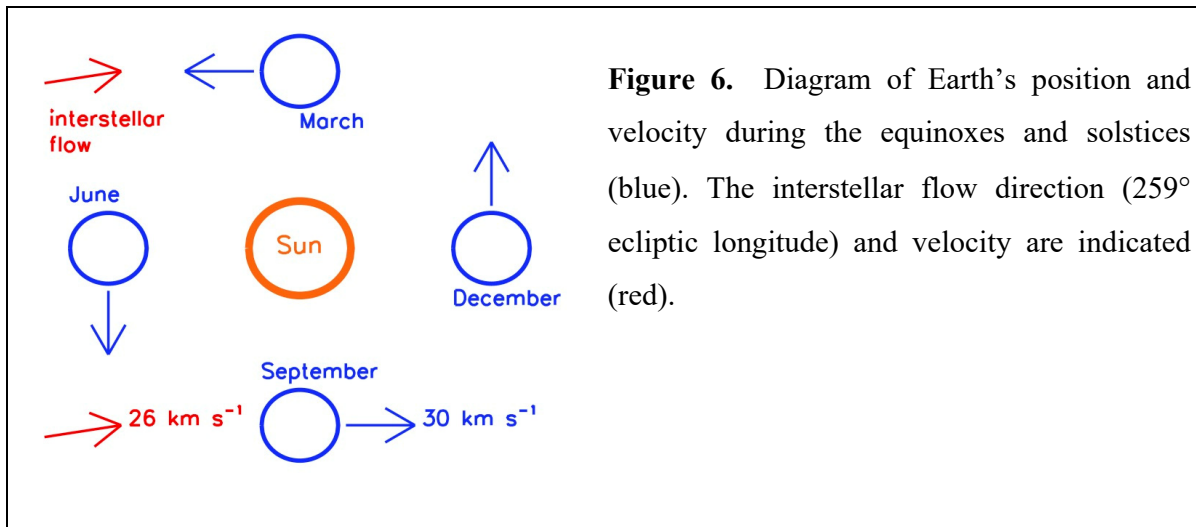
## 225 **5. Decadal and annual variations in meteoric flux**

226 Wind shows an annual cycle in meteoric flux, that varies in amplitude on a decadal time  
227 scale (e.g., Figure 2). The flux annual cycle is examined in greater detail in Figure 5, where the  
228 amplitude in 2009 is much larger than in 2020. This annual variation has been discussed by  
229 others [Kellog *et al.*, 2016; Malaspina *et al.*, 2014; Wood *et al.*, 2015; Zaslavsky *et al.*, 2012],  
230 who concluded that higher flux near the Vernal equinox is associated with ISD. The reason is  
231 that Earth's ram direction is into the interstellar flow in March, and away from it in September  
232 [see also Malaspina *et al.*, 2014]. Because the Earth's orbital velocity ( $\sim 30 \text{ km s}^{-1}$ ) is similar in  
233 magnitude to that of ISD ( $\sim 26 \text{ km s}^{-1}$ ), the ISD flux detected by Wind approaches zero during  
234 September, when the velocities are nearly parallel (Figure 6). Given this pattern, Wind  
235 observations near the autumn equinox should represent mostly IPD, and an estimate of the IPD  
236 flux was determined from Wind as the mean of observations surrounding the autumn equinox  
237 (September  $22 \pm 45$  days). It should be noted that the assumption of a constant ISD velocity here  
238 is a simplification [e.g., Sterken *et al.*, 2012], and that faster and slower ISD can exist. It is thus  
239 possible that Wind measurements in September contain a few ISD, although a more detailed  
240 treatment of ISD velocities is beyond the scope of this paper. The ISD flux from Wind was  
241 subsequently obtained by subtracting the IPD flux from the monthly mean values during the rest  
242 of the year (Figure 5). The approach taken here for isolating IPD and ISD in Wind observations  
243 is similar to that used by Wood *et al.* [2015]. The Wind ISD flux shows a strong annual variation,  
244 that is consistent with ISD flux simulations from the IMEX model [Sterken *et al.*, 2015; Krüger  
245 *et al.*, 2019]. The IMEX results shown here are for the near Earth environment (1 AU), and ISD  
246 with radii from 0.1 - 0.73  $\mu\text{m}$ . Note that both Wind and IMEX indicate a stronger spring ISD  
247 peak in 2009 than in 2020, which is explored below. The model peak occurs slightly later (April)

248 than in Wind (March), which could be related to the assumed interstellar flow direction ( $259^\circ$   
 249 ecliptic longitude) in IMEX. The model peak in ISD would occur at an earlier time if the  
 250 interstellar flow was from a higher longitude (roughly 1 day per degree). Indeed, a slightly  
 251 higher ecliptic longitude ( $274^\circ$ ) for ISD was considered by *Sterken et al.* [2014] in an  
 252 interpretation of Stardust in situ ISD measurements. In additions, the analysis of Wind and  
 253 Ulysses observations by *Wood et al.* [2015] also suggests that the interstellar flow comes from  
 254 ecliptic longitudes greater than  $259^\circ$ .



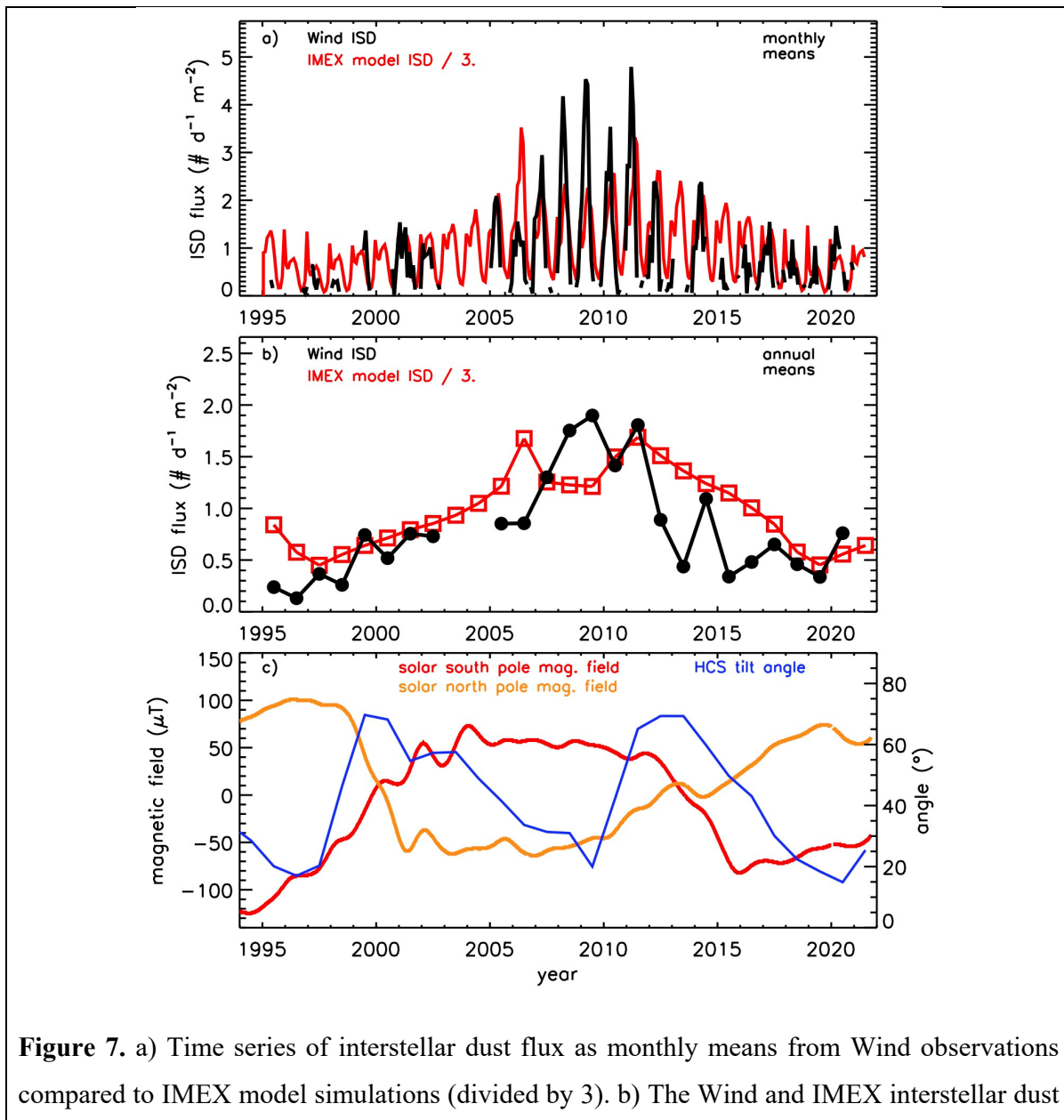
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257 Monthly mean ISD flux during 1995 - 2021 from Wind compare favorably to IMEX  
 258 simulations (Figure 7a), with both indicating a stronger annual variation and greater overall ISD  
 259 flux during 2004 - 2016. The trajectories of ISD in the solar system are controlled by solar  
 260 gravity, solar radiation pressure, and magnetic field strength. Theory predicts that ISD flux  
 261 should correlate to changes in the solar magnetic field (SMF), due to the Lorentz forces  
 262 experienced by charged particles in motion [e.g., Landgraf et al., 2000; Sterken et al., 2015]. The  
 263 SMF varies with a 22-yr periodicity (the Hale cycle, see Figure 7c), where the net effect on ISD  
 264 is a focusing of particles in the inner heliosphere during solar south pole positive phases.  
 265 Variations in the ISD flux are indeed coincident with changes in SMF, and linear regression to  
 266 annual means gives correlation coefficients ( $\rho$ ) of 0.74 for Wind and 0.67 for IMEX (Figure 7b).  
 267 The heliospheric current sheet (HCS) tilt was also examined, which is a representation of the  
 268 interplanetary magnetic field that exhibits an 11-yr cycle (in phase with Lyman- $\alpha$  flux). The  
 269 HCS is considered a factor in ISD trajectories [Sterken et al., 2015], although the Wind and  
 270 IMEX ISD are only weakly correlated with the HCS tilt ( $\rho = 0.1$  for Wind and 0.34 for IMEX).  
 271 The good agreement between Wind and IMEX confirms that Wind observations are dominated

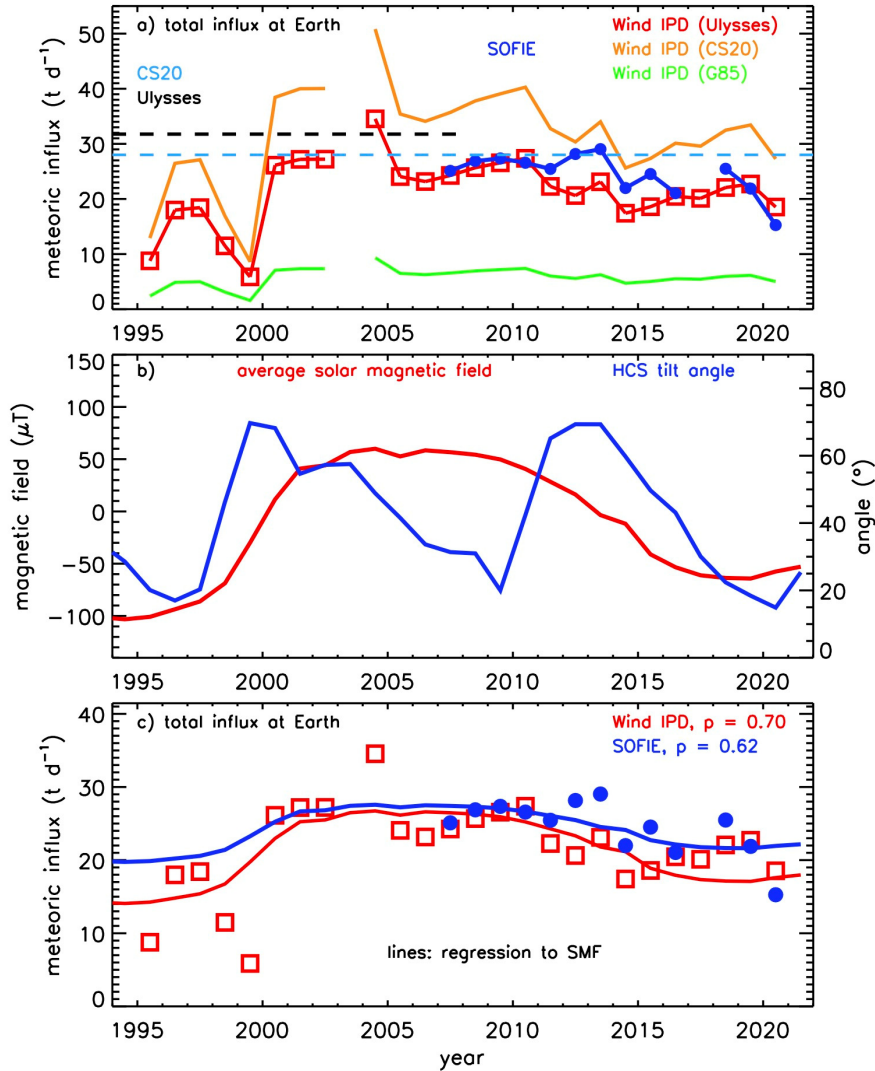
272 by ISD during spring. This agreement furthermore suggests that using Wind measurements near  
 273 the autumn equinox provides a good approximation of IPD, which is important for characterizing  
 274 the influx at Earth. Note that the Wind and IMEX results indicate an ISD influx at Earth (using  
 275 equation (3)) of less than  $0.01 \text{ t d}^{-1}$ , which is insignificant compared to the IPD influx of  $\sim 25 \text{ t d}^{-1}$   
 276 <sup>1</sup>.



**Figure 7.** a) Time series of interstellar dust flux as monthly means from Wind observations compared to IMEX model simulations (divided by 3). b) The Wind and IMEX interstellar dust

flux as annual means. c) Solar magnetic field strength at both poles and the heliospheric current sheet (HCS) tilt angle, from the Wilcox Solar Observatory.

277 Meteoric influx at Earth was estimated from the Wind IPD flux determined as above. The  
278 Wind  $F_E$  estimates use a constant (equation (4)) that is based on integrating the mass  
279 distributions from either Ulysses, G85, or CS20 (see Section 4). Wind  $F_E$  determined from the  
280 yearly IPD flux values are compared to SOFIE, Ulysses, and CS20 in Figure 8a. Wind  $F_E$  using  
281  $C$  determined from G85 are much lower than when  $C$  from Ulysses or CS20, with the later  
282 providing Wind influx estimates that are in reasonable agreement with SOFIE, Ulysses, and  
283 CS20. Note that both Wind and SOFIE show similar year-to-year influx variations. The influx  
284 time series are compared to the solar magnetic field and HCS tilt angle in Figure 8b. Linear  
285 regression indicates a strong  $F_E$  - SMF correlation for both Wind ( $p = 0.70$ ) and SOFIE ( $p =$   
286  $0.62$ ). The correlation with HCS tilt angle is weaker, however, with  $p = 0.01$  for Wind and  $0.42$   
287 for SOFIE. The Wind and SOFIE  $F_E$  time series are shown in Figure 8c with a fit to the results  
288 based on linear regression to the SMF. These results suggest that changes in the SMF can alter  
289 the flow of IPD in the inner heliosphere. This result is not anticipated, however, because IPD  
290 have small charge-to-mass ratios ( $Q/m$ ), unlike ISD. While the IPD response to changing SMF  
291 would not be instantaneous, it is possible that a cumulative effect is realized after many orbits of  
292 the sun. This idea is supported somewhat by the apparent time delay in changing  $F_E$  with respect  
293 to SMF. Indeed, the highest  $F_E$  - SMF correlation is found for a time lag of 14 months, with  $p =$   
294  $0.76$  for Wind and  $0.70$  for SOFIE. The agreement between SOFIE and Wind concerning year-  
295 to-year and decadal changes in meteoric influx is encouraging, and furthermore suggests that  
296 variability in IPD appears in Earth's mesosphere.



**Figure 8.** a) Time series of total meteoric influx ( $F_E$ ) at Earth. Wind results are shown for conversion factors based on the dust distribution from either Ulysses, *Carrillo-Sánchez et al.* [2020] (CS20), or *Grün et al.* [1985] (G85). SOFIE results are the average for winter months near  $67^\circ N$  (Nov-Feb) and  $67^\circ S$  (Jun-Sep) [*Hervig et al.*, 2021]. The average influx based on Ulysses IPD observations during 1990 - 2007 is indicated ( $|b| < 30^\circ$  and  $D_s < 1.5$  AU). The CS20  $F_E$  based on observations of iron and sodium in the mesosphere and cosmic spherules collected at South Pole is shown. b) The North-South average solar polar magnetic field strength and HCS tilt angle, as annual means. c) SOFIE and Wind  $F_E$  ( $C$  from Ulysses), with regressions to the average SMF time series and regression coefficients as listed.

297 The median  $F_E$  from Wind during 2007 - 2020 for  $C$  determined from the Ulysses, CS20,  
 298 and G85 results are 22, 33, and 6 t d<sup>-1</sup>, respectively, each with a standard deviation of ~14%  
 299 during the time period. The median  $F_E$  based on the three Wind time series during 2007 - 2020  
 300 (i.e., Figure 8a) is 22 t d<sup>-1</sup>, with a standard deviation of 11 t d<sup>-1</sup> (50%) which is due mostly to the  
 301 spread in  $C$ . The total uncertainty in Wind  $F_E$  includes contributions from the error in  $H_F$  and  
 302 uncertainties in the IPD flux measurements, and is closer to 13 t d<sup>-1</sup> (60%). Table 1 summarizes  
 303 the meteoric influx estimates from this work and some recent publications.

**Table 1.** Total meteoric influx (ablated + surviving material) at Earth from different sources.

Source	Method	Meteoric influx (t d <sup>-1</sup> )
Wind	Satellite in situ dust detection (2007-2020), <i>this work</i>	22 ± 13
Ulysses	Satellite in situ dust detection (1990-2007), <i>this work</i>	32 ± 76
SOFIE	Satellite remote measurements of meteoric smoke in the mesosphere (2007-2021) [ <i>Hervig et al., 2021</i> ]	25 ± 7
Lidar and cosmic spherules	Lidar measurements of Fe & Na in the mesosphere and cosmic spherules collected at South Pole [ <i>Carrillo-Sánchez et al., 2020</i> ]	28 ± 16
LDEF	Long Duration Exposure Facility (LDEF) satellite in situ dust measurements (1984-1990) [ <i>Borin et al., 2017</i> ]	15 ± 3
Lidar	Lidar measurements of Na in the mesosphere [ <i>Gardner et al., 2014</i> ]	60 ± 16

## 304 5. Summary

305 This work examined meteoric influx using in situ dust detection by the Wind and Ulysses  
 306 spacecraft, and observations of meteoric smoke in the mesosphere by the SOFIE satellite  
 307 instrument. Wind does not resolve the mass of the detected meteoroids, but rather reports the  
 308 total number of particles with mass from 10<sup>-14</sup> - 10<sup>-8</sup> g. The Wind measurements were separated  
 309 into the interstellar and interplanetary components. The Wind ISD estimates are in good  
 310 agreement with simulations from the IMEX model, in terms of the both the annual and decadal  
 311 flux variations. The decadal ISD variation is correlated to the 22-year solar magnetic cycle, as

312 anticipated by theory and predicted by IMEX. The Wind IPD observations were related to the  
313 total meteoric influx at Earth using IPD mass distributions from Ulysses and previous  
314 publications. The resulting Wind influx estimates are in good agreement with SOFIE and  
315 Ulysses. The SOFIE and Ulysses influx time series show similar year-to-year and decadal  
316 changes in meteoric influx. The decadal influx variation exhibits an unanticipated correlation to  
317 the 22-year solar magnetic cycle. This relationship may be an artifact, or could indicate that  
318 changes in the solar magnetic field can alter the trajectories of interplanetary dust.

319 **Acknowledgements.** This work was funded in part by the AIM mission through NASA contract  
320 NAS5-03132. V. J. Sterken received funding from the European Union's Horizon 2020 research  
321 and innovation program under grant agreement N 851544. L. B. Wilson was partially supported  
322 by Wind MO&DA funds.

### 323 **Data Availability Statement**

324 SOFIE data are available online ([sofie.gats-inc.com](http://sofie.gats-inc.com)). WIND meteoric dust data are available  
325 online ([cdaweb.gsfc.nasa.gov](http://cdaweb.gsfc.nasa.gov)). Wilcox Solar Observatory data are available online  
326 (<http://wso.stanford.edu/#MeanField>). Ulysses dust observations are available online  
327 ([cdaweb.gsfc.nasa.gov](http://cdaweb.gsfc.nasa.gov)).

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