

Experimental Investigation of Distributed Sand-Grain Roughness Effects on Transition Onset and Turbulent Heating Augmentation at Mach 6

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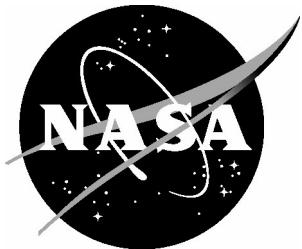
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Abstract

An experimental investigation of distributed sand-grain surface roughness effects on boundary-layer transition and turbulent heating has been performed at hypersonic test conditions. Two representative entry vehicle geometries, a sphere-cone aeroshell and a spherical-cap aeroshell, were considered. Cast ceramic models of each geometry were fabricated with distributed sand-grain roughness patterns of different heights that simulated an ablated thermal protection system. Wind tunnel testing was performed at Mach 6 over a range of Reynolds numbers sufficient to produce laminar, transitional, and turbulent flow. Aeroheating and boundary-layer transition onset data were obtained using global phosphor thermography. The experimental heating data are presented herein, as are comparisons to laminar and turbulent smooth-wall heat transfer distributions from computational flow field simulations. The boundary-layer transition data were found to correlate with a functional representation developed in prior roughness studies, although the data scatter was greater owing to the height variability of distributed sand-grain roughness.

Nomenclature

Symbols

D	ASTM mesh particle diameter
H_0	tunnel total (reservoir) enthalpy
H_{AW}	adiabatic wall enthalpy
H_k	roughness height enthalpy
H_w	wall enthalpy
H_{300K}	wall enthalpy at 300 K
h	measured heat-transfer film coefficient
h_{FR}	Fay-Riddell theory heat-transfer film coefficient
h	measured roughness height
h_{hex}	hexcomb cell height
h_{nom}	nominal ASTM mesh particle roughness height
h_{mean}	measured mean roughness height
h_{PVxx}	measured peak-to-valley roughness height for $xx\%$ exceedance height
h_{RMS}	measured root-mean-square roughness height
k	actual roughness height
k_{PVxx}	actual peak-to-valley roughness height for $xx\%$ exceedance height
M_e	boundary-layer edge Mach number
M_∞	free stream Mach number
p	effective roughness height transfer parameter
q	heat transfer rate
r_{stag}	approximate radius of edge Mach stagnation cutoff region
R	model radius
R_N	model nose radius
R_C	model corner radius
Re_{k+}	roughness height Reynolds number
Re_θ	boundary-layer momentum thickness Reynolds number
Re_∞	free stream unit Reynolds number
s	surface running length to edge Mach cutoff
s_0	surface running length from stagnation point
T_e	boundary-layer edge temperature
$T_{w,AVG}$	average wall temperature
T_∞	free stream temperature
U_e	boundary-layer velocity
U_τ	smooth-wall friction velocity
U_∞	free stream velocity
x, y, z	Cartesian coordinates
X_{TR}	transition correlation disturbance parameter
Y_{TR}	transition correlation flow field parameter
z_i	local surface height measurement
β_N	model spherical nose included angle
β_{POHL}	modified Pohlhausen parameter
δ	boundary layer thickness
ϕ	streamline angular coordinate identifier
ϕ_{corr}	heating calibration corrector factor

λ_{POHL}	Pohlhausen parameter
ν_e	boundary-layer edge kinematic viscosity
μ_w	wall viscosity
ρ_e	boundary-layer edge density
ρ_∞	free stream density
θ	boundary-layer momentum thickness
σ_h	roughness height standard deviation

Subscripts and Superscripts

∞	wind tunnel free stream condition
0	wind tunnel stagnation or reservoir condition
e	boundary layer edge condition
k	roughness height condition
PV	peak-to-valley surface roughness distance
RMS	root mean square
TR	transition location
w	model wall condition

Acronyms

ASTM	American Society for Testing and Materials
CFD	Computational Fluid Dynamics
IHEAT	Imaging for Hypersonic Experimental Aerothermodynamic Testing
LAL	Langley Aerothermodynamic Laboratories
LAURA	Langley Aerothermodynamic Upwind Relaxation Algorithm
OML	Outer Mold Line
TPS	Thermal Protection System

Introduction

This report serves to document an experimental dataset of distributed sand-grain surface roughness effects on boundary-layer transition and heating augmentation. The data were obtained through hypersonic wind tunnel testing of two representative entry vehicle aeroshell geometries with roughness patterns that simulated those produced by the ablation of a monolithic thermal protection system (TPS). This report represents a reference document that can be used as the basis for future detailed analysis of the heat-transfer distributions and boundary-layer transition onset locations measured in the test program. This study is a direct follow-on to a previous study (Refs. [1–2]) of ablated TPS with hexcomb pattern roughness effects on the same aeroshell geometries in which the transition criterion applied to the current data was developed. Another related study (Refs. [3–4]) on the effects of distributed sand-grain surface roughness – on hemispherical nose tips, rather than aeroshell geometries – provided the basis for the distributed sand-grain roughness height analysis presented herein.

Background

“Roughness” is a generic term in aerospace literature that encompasses many types of surface features that deviate from that of a smooth outer mold line (OML) surface. Roughness can be divided into two general types, discrete and distributed. Discrete roughness (Figure 1) includes surface features such as: a) protruding compression pads or recessed cavities at mechanical attachment points; b) steps or gaps between

heat shield tiles or blocks resulting from differential ablation of the TPS and the filler material between them; and c) physical damage to a TPS. Distributed roughness (Figure 2) includes features such as: a) regular patterns resulting from ablation of hexcomb-structure TPS; b) deflections of a flexible TPS over its support structure when subjected to aerodynamic loading; c) random “sand-grain” features resulting from ablation of a monolithic TPS; or d) the texture of overlapping fibers on a woven TPS.

Data on the effects of surface roughness are valuable because the roughness of an entry vehicle’s TPS can promote earlier boundary-layer transition and produce higher turbulent heating (and shear) levels than would be expected based on an idealized, smooth-surface analysis. However, due to the complexities of roughness effects, a vehicle’s TPS is typically designed using analytical, computational, and/or experimental techniques based on the assumption of a smooth surface. The effects of roughness on the aerothermodynamic environment are then included through approximate engineering correlations and methods.

The purpose of this test program was to obtain data on the effects of distributed sand-grain type roughness. Examples of ablating TPS that produce such patterns include the PICA material used on the Mars Science Laboratory and Mars 2020 missions’ entry vehicle heat shields and on the Stardust comet sample return mission. The data obtained in this test program are intended for use in the development and/or validation of engineering correlations for the effects of distributed sand-grain roughness on boundary-layer transition and turbulent heat transfer for such vehicles. These data can also serve as the basis for development and/or validation of higher-fidelity, numerical flow-field simulation models for roughness effects.

Experimental Tools and Methods

Wind Tunnel Models

Model Geometries

Roughness effects data were obtained on two representative entry vehicle geometries: a sphere-cone geometry (Figure 3) and a spherical-cap geometry (Figure 4). The sphere-cone geometry is representative of the Mars Viking – Mars Pathfinder – Mars Exploration Rover – Mars Phoenix – Mars Science Laboratory – Mars 2020 family of entry vehicles used in NASA’s robotic Mars exploration missions. The spherical-cap geometry is representative of the Mercury – Gemini – Apollo – Orion family of entry vehicles employed in NASA’s crewed space program. Geometry parameters for both model configurations are listed in Table 1. Multiple wind tunnel models of each configuration were fabricated with a range of distributed sand-grain roughness patterns, as detailed below.

Model Fabrication

Models with a wide range of distributed surface roughness heights were fabricated for this study. The fabrication process for a rough-surface model follows that for smooth models, as documented in Ref. [5], with additional steps to add roughness to the surface. The first step in fabrication of a smooth-surface model is the production of a rapid-prototype pattern of the geometry using a 3D wax printing machine. Investment casting is then used to make a mold from the pattern. A thin-shell silica ceramic model is then slip cast from the mold, dried and sintered. The shell is then backfilled with a hydraulically setting magnesia ceramic for strength and support and mounted on a stainless-steel support sting. Finally, the model is coated with a mixture of phosphors that luminesce under ultraviolet lighting and fiducial marks are placed at specified locations on the surface for image registration.

To fabricate a model with surface roughness, an adhesive coating is applied to the smooth wax pattern and then the pattern is placed into a container filled with precision-manufactured, spherical glass particles. The particles adhere to the wax pattern to form a distributed, sand-grain type roughness over the entire surface. The pattern may then be hand-worked if necessary to remove any obvious surface irregularities (i.e., clumps of glass spheres). A ceramic model with surface roughness can then be fabricated from the roughened pattern following the remaining steps detailed above for smooth-surface models.

The distributed sand-grain surface roughness was created using precision-manufactured, spherical glass particles ranging in size from 2.5 mil to 68.9 mil¹. The spherical glass particle diameters are specified according to an ASTM standard (Ref. [6]) that defines the mesh sieve opening size through which the particle can pass, hence the ‘ASTM-XXX’ mesh nomenclature for the models. A nominal roughness height for each ASTM mesh size is defined based on the assumption that for a uniform array of roughness elements (the spherical particles) in contact with each other, the roughness height is equal to the vertical distance from the top of an element to the point of contact with the adjacent element – that is the roughness height is equal to the glass particle radius. This height measurement is referred to as the nominal “peak-to-valley” roughness. The characterization of the actual “as-built” roughness height, which differs from the specified nominal value, will be presented below in the “Roughness Height Characterization” section. A listing of the nominal ASTM roughness sizes used in this study is given in Table 2.

Photographs of sample cast ceramic sphere-cone and spherical-cap roughness models with inset closeup images of the roughness at the nose are provided in Figure 5 and Figure 6, respectively. Although the image contrast and resolution are generally not sufficient to display the smaller roughness heights, all the roughness element heights did produce measurable effects on boundary-layer transition and surface heating.

Roughness Height Characterization

Surface roughness scan data acquisition and processing

The surface roughness data presented previously in Table 2 represent ideal values based on the sizes of the particles used in the fabrication of each model. However, the characterization of the actual “as-built” roughness was more complex and was based on a statistical analysis of the surface height distributions. The differences between the ideal and as-built roughness geometries are illustrated in Figure 7. In the ideal geometry, the roughness is defined by perfectly formed hemispherical elements in a single, flat, layer. However, the process of binding the glass particles to the surface, forming the mold, and then casting and coating the ceramic model introduces random imperfections in the surface.

To determine the as-built roughness characterization parameters, laser scans were made of 4-in. x 4-in. square flat sample plates for each of the ASTM roughness sizes to obtain a data cloud of *x-y-z* points. Margins of 0.5 in. on the sides of the plates were specified to avoid any edge effects and so the actual scan area was a 3 in. by 3 in. square. These flat plates were used in place of the actual model geometries owing to the difficulty of performing a scan over a curved surface. The stated, ideal scan resolution of the system was $\sim \pm 2.00$ mil ($\sim \pm 0.05$ mm), however, the actual achieved resolution was approximately ± 4.00 mil ($\sim \pm 0.1$ mm). The data cloud was then triangulated to form a continuous surface representation and the height (*z*) coordinate was shifted to put the average height of all points at zero.

Profile line-cuts were extracted from the global data sets at various stations to determine the height distribution of scan data points. The global surface scan data and representative profile line-cuts are shown

¹ The “mil” unit (0.001 inch) will be used frequently in the discussion of roughness heights rather than inches or SI units in deference to historical literature on surface roughness.

for each ASTM mesh size sample in Figure 8 – Figure 13. In the line plots, the dashed blue lines represent the nominal diameter and height of a roughness element, and the symbols represent the scan data point locations. As can be seen from these figures, several data points were obtained on each roughness element, for the larger roughness sizes, while for the smaller roughness sizes, the data point spacing was on the order of the roughness element diameter (e.g., 4.17 mil diameter for 140-Mesh and 2.48 mil diameter for the 240-Mesh). Therefore, while data and analyses for these smaller mesh sizes will be presented herein, these data are not considered to be quantitatively reliable.

Statistical Analysis of Roughness Data

The scan data were processed to determine several parameters, including the root-mean-square surface height distribution (h_{RMS}), local peak-to-valley roughness element heights (h_{PV}), and the actual effective peak-to-valley roughness height (k_{PV}).

The root mean square height distribution h_{RMS} is simply a function of the distribution of all height measurement points on the *entire* sample surface as per Eq. (1), where n is the number of points and z_i is the point height above a reference plane. This RMS value does not provide a direct characterization of the heights of the roughness elements, since sample points do not necessarily coincide with the maximum or minimum of an individual roughness element. Nevertheless, the RMS height is a relatively straightforward quantity to define and measure and is frequently reported in roughness studies. A functional relationship between RMS and peak-to-valley heights will be demonstrated subsequently.

$$h_{RMS} = \sqrt{\frac{1}{n} \sum_i (z_i)} \quad (1)$$

Previous studies of distributed roughness effects (e.g., Refs. [7–8]) have identified the peak-to-valley roughness height h_{PV} as a key parameter in the correlation of roughness data effects. This conclusion was confirmed in the more recent tests conducted in both wind tunnels (Refs. [3–4]) and ballistics ranges (Ref. [9]). The peak-to-valley roughness height represents the difference between the minimum and maximum surface heights of adjacent roughness elements, which for the idealized surface illustrated in Figure 7 would simply be the radius of the spherical roughness element.

To determine the h_{PV} of the actual roughness sample plates, peaks and valleys of adjacent roughness elements were determined manually from examination of the profile line-cut data. The selection of which points along the profiles of Figure 8 – Figure 13 actually represented the peaks and valleys of a roughness element, as opposed to minor variations over an element, was highly subjective and was additionally hindered by the resolution of the scans for the smallest ASTM mesh sample plates. Furthermore, while a fixed ASTM bead size was employed to create each roughness pattern, the application of roughness elements to the wax patterns, fabrication of molds, and final casting of the ceramic models introduced random deviations from the nominal peak-to-valley roughness height. A final complication is the difference between the peak-to-valley heights measured along an arbitrary profile line – which may not pass through the maximum and minimum of an element – as illustrated in Figure 14 – and the true peaks and valleys of individual elements.

These issues were dealt with using a statistical “exceedance” value to represent the random distribution of heights and a semiempirical correction factor to account for the differences between the peak-to-valley heights measured along an arbitrary profile and the actual peak-to-valley heights.

For each profile, the exceedance, which is defined as the percentage of data points in a set greater than a specified value, was computed from the database of peak-to-valley measurements. The exceedance height distributions obtained from scans of each sample plate are plotted in Figure 15. Also shown in this figure are Gaussian curve fits to these data of the form given by Eq. (2), where h_{mean} is the mean of the measured peak-to-valley heights and σ_h is the standard deviation.

$$\% \text{exceedance fit} = 100 \times \exp \left\{ -0.5 \left[\frac{(h_{PV} - h_{\text{mean}}/2)}{\sigma_h} \right]^2 \right\} \quad (2)$$

The exceedance height distributions are replotted in terms of the normalized roughness height (h_{PV}/h_{mean}) in Figure 16. As seen in these two figures, the exceedance distributions approximate the Gaussian distribution that would be expected from a large sample set with random deviations.

A relationship between the actual peak-to-valley height of a roughness element and that measured along a profile line (which does not necessarily pass through the center of the element) was developed by Dirling (Ref. 10) in an analysis for simple geometric elements including hemispheres, cones, and rectangles. Dirling's function for converting between a statistical mean of the measured height, h_{mean} , and the actual height, k , using a transfer function, p is given by Eq. (3). For the current analysis, where the roughness elements are represented as a tightly packed array of nominally hemispherical elements, the $\pi/4$ value of the transfer function is employed. Presumably, transfer functions could also be determined for more complex roughness shapes such as rods, fences, weaves, or honeycombs, but data would be required to validate such functions.

$$h_{\text{mean}} = k \times p; \text{where } p = \begin{cases} \pi/4 & \text{hemispherical element} \\ 0.5 & \text{conical element} \\ 1 & \text{rectangular element} \end{cases} \quad (3)$$

For the purposes of this study, it was assumed that this relationship for the mean value also held for any arbitrary exceedance percentile value, allowing for estimation of the true peak-to-valley exceedance height values from the heights measured along the profile lines. Values of the 30th and 50th percentiles have typically been reported in roughness literature, and herein values for the 50th, 30th, 15th and 5th percentiles are given. These values were derived from the measured roughness height distributions of Figure 15 - Figure 16 and are listed in Table 3.

The relationship between the estimates for the actual, as-built, heights to the nominal heights for the 50th, 30th, 15th and 5th percentile exceedances is shown in Figure 17. For the larger ASTM Mesh sizes (10, 20 and 40-mesh), the ratios of k_{PV}/h_{nom} are approximately the same for each exceedance percentile, varying from ~ 0.5 for the 50th percentile to ~ 1.0 for the 5th percentile. However, the ratio of actual to nominal height begins to increase rapidly with decreasing mesh size for the smaller ASTM mesh sizes (80, 140, and 230-Mesh). For the smallest, ASTM 230-Mesh, the ratios of actual to nominal height varies from ~ 1.6 for the 50th percentile to 2.7 for the 5th percentile.

Given that the fabrication process results in a roughness height distribution with a Gaussian shape, the estimates for the as-built to nominal ratios for the larger ASTM mesh sizes seems to be reasonable. However, absent any other information, the estimated peak-to-valley heights for the smaller ASTM mesh sizes would appear to be much larger than the nominal values. As noted earlier though, the nominal

roughness sizes for these models were on the order of the ideal image resolution of the scanning system used to obtain the data. Thus, while the results may indicate that the smaller mesh size model roughness heights were larger than intended, the fidelity of the scan data set was insufficient to definitively quantify the smaller mesh size heights. Instead, a rough estimate for the as-built heights for the smaller ASTM samples was made by fitting a curve through the higher ASTM sample ratios and extrapolating the results to the lower ASTM values to provide a corrected estimate for the peak-to-valley heights. Table 4 provides a summary of these original and corrected values of the 50% exceedance heights, as well as the nominal and RMS values for each ASTM mesh size.

For purposes of comparison with prior studies in which only RMS values are cited, it is useful to provide a relationship between the RMS and peak-to-valley exceedance values. It has been shown (e.g., Refs. [7–8]) that an approximate linear relationship holds between RMS and peak-to-valley exceedance parameters that allows for conversion between the two types of measurements. A plot of the current RMS roughness vs. actual peak-to-valley roughness data from each sample plate is presented in Figure 18 along with a linear correlation for each exceedance value. Depending on the exceedance value, the ratio of k_{PV}/h_{RMS} was found to vary from approximately 2.8 to 4.8.

These results can be compared to those from Jackson in Ref. 8, where it was stated that the “significant” peak-to-valley roughness was equal to 3.6 times the RMS value. Unfortunately, the term “significant” was not explicitly defined, although in Ref. [7], Batt later concluded that the “significant” value was equal to the 30th percentile exceedance height. Reexamination of the limited roughness profile data available from Ref. [8] suggests that Batt’s conclusion was tenuous, and that Jackson’s “significant” values could just as easily be equated to the 50th percentile roughness height. However, it is possible that Batt had access to more of the original data set than has been published and drew conclusions based on those data. Regardless of the definition of “significant” for the Ref. [8] data set, the cited value of 3.6 for the ratio falls within the range of exceedance values for the current data set, which indicates that the approximate relationship between RMS and peak-to-valley exceedance values is valid.

Wind Tunnel Test Facility

Facility Description

Hypersonic wind tunnel testing of the roughness models was performed in the NASA Langley Aerothermodynamics Laboratory (LAL) 20-Inch Mach 6 Air Tunnel. This wind tunnel is described in brief below and more detailed information on the LAL facilities can be found in Refs. [11–12].

The 20-Inch Mach 6 Air Tunnel (Figure 19 – Figure 20) is a blow-down facility in which heated, dried, and filtered air is used as the test gas. The tunnel has a two-dimensional contoured nozzle that opens into a 20.5 in. × 20.0 in. test section. The tunnel is equipped with a bottom-mounted injection system with a -5-deg to +55-deg pitch range and ±5-deg yaw range that can transfer a model from a sheltered model box to the tunnel centerline in less than 0.5 sec. Run times of up to 15 minutes are possible in this facility, although for the current aeroheating study, run times of only a few seconds were required. The nominal reservoir conditions of this facility produce perfect-gas free stream flows with Mach numbers between 5.8 and 6.1 and unit Reynolds numbers of $0.5 \times 10^6/\text{ft}$ to $8.3 \times 10^6/\text{ft}$. With its wide Reynolds number operating range capable of producing laminar, transitional, or turbulent flow on most geometries, this tunnel is primarily used for heat-transfer and boundary-layer transition studies.

Facility Operating Conditions

Data were obtained in Tests 7036 and 7057 in the 20-Inch Mach 6 Air Tunnel at six unit-Reynolds numbers from $Re_\infty = 2.1 \times 10^6/\text{ft}$ to $8.1 \times 10^6/\text{ft}$ with nominal free stream conditions as per Table 5. All sphere-cone data were obtained at $\alpha = 16$ deg and all spherical-cap data were obtained at $\alpha = 28$ deg. Full run matrices for the two tests are given in Table 6 and Table 7, respectively. Entries in these tables are sorted first by roughness height, then by free stream unit Reynolds number. These angles of attack and free stream conditions were selected for continuity with the hexcomb roughness dataset presented in Refs. [1–2].

Free stream velocity (U_∞), density (ρ_∞), temperature (T_∞), unit Reynolds number (Re_∞), and Mach number (M_∞) are provided in these tables. Additionally, an average model surface temperature (T_w), enthalpy difference (ΔH_{tot}), and reference heat transfer film-coefficient value (h_{FR}) are provided. The temperature is the average over the model surface when the thermographic phosphor image was obtained and is provided because boundary-layer transition is known to be sensitive to wall temperature. The enthalpy term is defined as the difference $H_0 - H_{300K}$ between the free stream total enthalpy and the wall enthalpy at cold wall (300 K) conditions. The film coefficient is the value from the Fay-Riddell theory (Ref. [13]) at cold wall conditions, where the radius is the nose radius of the model geometry.

Experimental Data

Data Acquisition and Reduction

Aeroheating data were obtained using the two-color, relative-intensity, global phosphor thermography method (Ref. [14]) and reduced using the IHEAT (Imaging for Hypersonic Experimental Aerothermodynamic Testing) code (Refs. [15–16]). In this method, a model is illuminated by ultraviolet light sources that induce temperature-dependent fluorescence of the phosphor coating. Fluorescent intensity images of a model are taken in the tunnel before and during a run using a three-color, charge-coupled device camera and the images are processed to determine heat-transfer distributions. The intensity data are then converted to temperatures using pretest calibrations of the data acquisition system.

Heat-transfer film coefficients are determined by assuming a step function in the film coefficient from the prerun temperature to the run temperature, which corresponds to a parabolic temperature-time history. The heating data are typically reported in terms of the ratio h/h_{FR} where the heat-transfer film coefficient, h , is defined in terms of enthalpy as:

$$h = q/\Delta H_{TOT} = q/(H_{AW} - H_w) = q/(H_0 - H_w) \quad (4)$$

In the calculation of the heat-transfer film coefficient, it is assumed that for a blunt-body, the adiabatic wall enthalpy H_{AW} is equal to the free stream total enthalpy of the tunnel, H_0 , and the wall enthalpy H_w is determined from the surface temperature at each image pixel. This heat transfer coefficient definition provides a theoretically near-constant value over the course of a run since the decrease in time of the heat transfer rate in the numerator as the model surface becomes hotter is balanced by the decrease of the enthalpy-difference term in the denominator.

Data Mapping and Presentation

The two-dimensional (2-D) image data output from IHEAT (Figure 21) for each run were transformed

to account for optical perspective effects and mapped to a three-dimensional (3-D) CAD surface of the wind tunnel model (Figure 22). To accomplish this mapping, perspective, translational, and rotational transformations were first performed on the 3-D CAD surface until its 2-D projection matched that of the 2-D image data. The image data were then assigned transformed (x, y, z) coordinates based on interpolation between the image and projected surface geometry. Finally, the transformation was inverted to obtain an orthographic, 3-D heating distribution map of the experimental data.

An additional data manipulation was performed to extract the streamline-based heating distributions from the mapped wind tunnel data set. These streamline-based data sets are used in boundary-layer transition analyses and for comparisons of Reynolds number and roughness effects. For each run, streamlines were defined based on the boundary-layer edge velocity vectors from the computed flow fields (to be discussed in the next section). Thirty-six streamline termini were established at locations spaced in 10-deg increments around the circumference of the geometry and the streamlines were then traced backward from each terminus toward the flow field stagnation point. Each streamline is identified by the angular location, ϕ , of its terminus. The resulting streamlines are shown in Figure 23 for the sphere-cone geometry and in Figure 24 for the spherical-cap geometry. The geometric (x, y, z) coordinates along each streamline were then interpolated onto the 3-D mapped image and h/h_{FR} heating data were extracted along each streamline in terms of s_0/R , which is the normalized streamline distance from the stagnation point. Additionally, the predicted flow field quantities (boundary-layer height, momentum-thickness Reynolds number, etc.) were also extracted along these streamlines and combined with the wind tunnel data set to enable transition onset analyses.

One additional complication needs to be noted with respect to extraction of data along streamlines. The extraction algorithm tended to fail near the stagnation point where the velocity vectors approached zero; essentially, the physical location became indeterminate, resulting in unreliable path-lengths through the stagnation region. This problem was resolved by stopping the reverse tracing of the streamlines from the outer edge of the model toward that stagnation point at the location where the edge Mach number, M_e , dropped below 0.025. The “true” streamline length value, s_0 was determined from an estimate of the physical length from the M_e cutoff to the stagnation region as a function of the streamline terminus angular location ϕ and the approximate radius of the stagnation region, r_{stag} :

$$s_0 = s + \Delta s \quad (5)$$

$$\Delta s = \cos(2\phi) \times r_{stag}/3 + r_{stag}, \text{ where } r_{stag} \cong 0.003 \text{ m} \quad (6)$$

In the body of this report, plotted data will be shown in terms of s/R , as that is the quantity in which the data were extracted along streamlines. The estimated actual distance, s_0/R , can be determined using Eqs. (5) and (6).

The mapped data from all runs are collected in the Appendices and presented therein as large, high-resolution images. These images are ordered by model geometry, roughness height, and Reynolds numbers. Smaller images will be shown in the body of the report along with streamline-based heating distributions.

Phosphor Thermography Data Quality

An important factor that influences the quality of phosphor thermography data quality is the local surface

inclination at a given point on the model with respect to both the camera and the UV lights. Phosphor thermography provides the best results when the surface to be imaged is normal to the camera, which reduces perspective distortion and image smearing, and when the surface is well illuminated, which induces the best temperature response of the phosphor coating. Because of the three-dimensional nature of a wind tunnel model, the entire surface of a model cannot be optimally imaged, or in some cases cannot even be viewed. For blunt bodies such as those in this test, the windward centerline region of the model – the ‘bottom’ of the model with respect to the view orientation - is the area where the data quality is most affected. This situation is illustrated for a simple hemispherical model in Figure 25. Because of this limitation, windward region data are only regarded as qualitative, not quantitative. Although image data from this region will be shown, quantitative plotted data and transition location data will not be provided for the streamlines originating from the $\phi = 160\text{-deg}$ through 200-deg termini.

Heat Transfer Data Uncertainty

The experimental uncertainty for convective heat transfer measurements on a *smooth, blunt body geometry* model in the 20-Inch Mach 6 Air Tunnel is quantified as a function of net uncertainties resulting from: the data acquisition method ($\pm 10\%$); flow quality and test-condition repeatability ($\pm 5\%$); and the accuracy of the 3D mapping process ($\pm 10\%$), which results in an overall root-sum-squared value of $\pm 15\%$. Experience with this technique indicates that these values are usually conservative and agreement between laminar, smooth-wall measured and predicted heating levels is generally well within this range.

However, it is assumed that the distributed sand-grain roughness introduces additional uncertainties to the heat-transfer measurements. These roughness elements produce very detailed heating patterns due both to their three-dimensional shape and their influence on transition locations. In many cases, these patterns were smaller than the resolution of the camera system; thus, a measurement of heating over a roughness element is in effect, a spatial average, rather than a point measurement. Quantification of such errors on a macro-scale is not possible because of the localized and position/height dependency of each roughness element, but they are probably on the order of $\pm 10\text{--}20\%$. Taken together with the smooth-wall uncertainty, the uncertainty in distributed sand-grain roughness heating is estimated to be in the $\pm 18\text{--}25\%$ range

Calibration Correction for Heat Transfer Data

A central premise in the analysis of wind tunnel heating data is that for a given Mach number, the normalized heat transfer film coefficient, h/h_{FR} , at any point on a geometry remains constant with varying Reynolds number at perfect gas conditions for attached flow over a blunt body. This behavior is demonstrated through CFD simulations for a 2-inch diameter hemisphere over the current range of test conditions. As shown in Figure 26, laminar simulations using the LAURA code (see section below on Computational Tools and Methods) predict a constant value of $h/h_{FR} = 1.06$ at the hemisphere stagnation point for all test condition Reynolds numbers. The fact that the ratio is not exactly 1 is due to the differences between a modern CFD prediction for the film coefficient at perfect-gas wind tunnel conditions and the semiempirical Fay-Riddell correlation for the film coefficient based on approximate boundary-layer solutions for reacting-gas flight conditions. That the two predictions are so close is a testament to the utility of the original Fay-Riddell method that was developed in the 1950s.

While the CFD predictions do indeed demonstrate a constant value of h/h_{FR} for the wind tunnel conditions, the same behavior was not observed during the test program; in fact, a dependency on free stream Reynolds number was noted in the experimental data. This dependency is illustrated by stagnation point heating data from pretest checkout and calibration runs on a phosphor-coated, 2 inch radius hemisphere shown in Figure 27 for Test 7036 and in Figure 28 for Test 7057. Instead of a constant value

for h/h_{FR} , the measured values at lower Reynolds numbers were observed to be at or below the predicted value from the computations, while at higher Reynolds numbers the measured values were greater than the predictions. Second-order polynomial fits to the measured heating values that reflect these variations are also shown in the figures.

There are several potential sources of uncertainty which could be producing this dependency including: variations in the bulk materials used to cast the ceramic wind tunnel models; the consistency of the thermographic phosphor mixture used to coat the models; the fidelity of the phosphor intensity/temperature calibrations; degradation of the UV lighting or imaging camera; and/or the flow quality of the wind tunnel. Unfortunately, it was beyond the scope of this study to resolve whether any, or all, of these factors influenced the experimental data.

Because the differences in predicted and measured stagnation point film-coefficient ratios fell within the estimated uncertainty of $\pm 18\text{--}25\%$ range cited in the previous section, these results were considered “acceptable” from an experimental perspective. However, since the differences can be represented by a bias function that depends on Reynolds number, as opposed to a random dispersion, an additional data processing step was conducted to correct the heating data based on the hemisphere calibration run data. The original data were modified using the polynomial curve fits as per Eqs. (7) and (9) and all data and results presented herein reflect this calibration correction.

$$(h/h_{FR})_{exp,corr} = (h/h_{FR})_{exp} \times (h_{CFD}/h_{FR})_{hemi-stag} \times \phi_{corr} \quad (7)$$

where:

$$(h_{CFD}/h_{FR})_{hemi-stag} = 1.06 \quad (8)$$

$$\begin{aligned} \phi_{corr} &= 1/(A + Bx + Cx^2), \text{ and } x = Re_\infty/1,000,000 \\ \text{Test 7036: } A &= 0.8279 \quad B = 0.09120 \quad C = -7.583 \times 10^{-3} \\ \text{Test 7057: } A &= 0.9971 \quad B = 0.03163 \quad C = -2.034 \times 10^{-3} \end{aligned} \quad (9)$$

Computational Tools and Methods

Flow field solutions were generated using the LAURA (Langley Aerothermodynamic Upwind Relaxation Algorithm) code. LAURA (Refs. [17–18]) is a three-dimensional, structured-grid, finite-volume solver that includes perfect-gas and nonequilibrium chemistry options, a variety of turbulence models, and ablation and radiative transport capabilities. LAURA solutions were used for comparisons of predicted heating levels with the measured data and to define the streamlines along which to extract the mapped experimental data, as described above.

Solutions were computed on multiblock grids of each geometry with a smooth (no roughness elements) outer mold line. Grid adaption was performed to align the grid outer boundary with the bow shock and to cluster cells near the surface to produce wall cell Reynolds numbers on the order of 1 to 10. Free stream conditions were set to the nominal wind tunnel conditions for each operating point as given in Table 5. For these wind tunnel conditions, the perfect-gas air option was used. Both laminar and turbulent solutions were

generated. Turbulent cases were computed using the Cebeci-Smith algebraic model with fully turbulent flow over the entire geometry. Because the computations were performed on a smooth geometry, they are not quantitatively applicable to the actual wind tunnel tests performed on rough-surface models with heating augmentation but are still provided for qualitative comparisons.

For the wall temperature boundary condition, a change in the normal practice for wind tunnel simulations of setting this value to a “cold-wall” ambient temperature (because of the small variation in heat-transfer coefficient with temperature) was employed. Literature on roughness effects indicates a dependence of transition onset location on the ratio of boundary-layer edge temperature to wall temperature T_e/T_w . To approximately account for this effect (which was expected to be small, but non-negligible, for these test conditions), the computations were performed using a uniform “hot-wall” wall temperature set to the average of the measured surface temperature on the model. These values varied between ~ 325 K to 410 K, depending on roughness height and Reynolds number.

The flow field solutions also provided boundary-layer parameters that can be used in the correlation of transition and heating augmentation data. Centerline profiles of selected boundary-layer parameters are presented to provide insight into the range of the test data and potential relevance to flight missions.

The ratio of the physical roughness height to that of the boundary-layer, k/δ , has a first-order influence on transition onset. Centerline distributions of k/δ are presented in Figure 29 for the sphere-cone geometry and in Figure 30 for the spherical-cap geometry for the range of roughness heights and free stream Reynolds number conditions. In these figures, k is the corrected 50% exceedance values from Table 4 and δ is the physical height of the boundary layer. Values of k/δ varied over two orders of magnitude depending on roughness height and Reynolds numbers. The smallest k/δ values are well within the boundary layer, while the highest exceed the boundary layer height.

The turbulent roughness height Reynolds number, Re_{k+} as per Eq. (10), can be used as a correlation parameter for turbulent roughness heating augmentation. Centerline distributions of Re_{k+} are presented in Figure 31 and Figure 32 for the sphere-cone geometry and spherical-cap geometries, respectively, for the range of roughness heights and free stream Reynolds number conditions. Values of Re_{k+} also vary over two orders of magnitude, indicative of laminar flow at the lowest levels and roughness-augmented turbulent flow at the highest levels.

$$Re_{k+} = \rho_w U_\tau k_{PV50} / \mu_w, \text{ where } U_\tau = \sqrt{\tau_w / \rho_w} \quad (10)$$

Experimental Data Analysis

Transition Onset Location Definition

From a flow physics standpoint, transition onset is defined as the point where smooth, laminar flow in the boundary layer begins to break down into small eddies. This location can, in theory, be determined through flow field imaging and/or diagnostic techniques (e.g., high-frequency pressure measurements, laser velocimetry) to determine when fluctuations in a quantity of interest, such as the mean velocity, exceeded a specified criterion. However, in this study, the only measurements are of the surface temperature and (through data reduction) the surface heating. For such measurements, the differences in temperature or heating levels between laminar flow and transitional flow can be too subtle at some conditions to permit

precise definition of the onset location. This measurement is more difficult when the local roughness height is small, in which case, the change from laminar to transitional/turbulent flow is gradual, but easier when the local roughness height is large, and the transition length is very short.

In lieu of a precise measurement of the transition location, transition onset is defined herein through a common approach in which an “effective” or “apparent” onset point is determined through the “tangent-slope-intercept” method. As shown in Figure 33, using sample data for a hemisphere from Ref. 4, the effective transition onset location is identified as the point where a line drawn tangent to the slope of the heat-transfer distribution curve through the transition region intercepts the nominal, laminar level. While this method does not necessarily identify the precise location at which fluctuations in the boundary-layer flow begin, it is consistent with common practice for determining the roughness-induced transition location via surface-based measurement techniques. This method also permits a more consistent means of identifying a relevant transition parameter, since identification of the small rise in heating levels at the actual transition onset location would be highly susceptible to error through surface measurement techniques alone. If the data herein are compared to other datasets, then it will be necessary to ensure that the same definition of transition onset is applied to ensure consistency.

While this effective transition onset location is easy to define in principle, in practice, there can still be considerable uncertainty in determining the effects of roughness on transition location because of the difficulty in precisely defining the relevant roughness height. Consider again the distributed roughness hemisphere example, where now the heating data are shown in image form (Figure 34) rather than a plotted distribution. In the ideal case, where the surface roughness was invariant over the entire model surface, the transition onset location would be at a constant streamline length around the circumference of the model since the flow field is axisymmetric. However, as shown in Figure 34, there are clearly circumferential variations in the transition onset location. These variations occur because the local roughness height at any given location can vary from the nominal value for many reasons, including: the fidelity of the model fabrication process; the uniformity in application of the phosphor coating; and damage to the coating due to handling of the model or particle impacts during testing.

For an axisymmetric flow such as that over the example hemisphere, it was possible to reduce the uncertainty in the transition onset location by averaging multiple onset locations around the body to determine a mean value. This approach was followed in Ref. [4]. However, for the three-dimensional flow fields produced by the geometries considered herein, that was not possible. Thus, transition correlations drawn from this dataset can be expected to have greater scatter than those that were derived from Ref. [4].

Reynolds Number Effects on Heating and Transition

The effects of Reynolds number on the heating levels and boundary-layer transition onset locations are illustrated for each distributed sand-grain roughness pattern in Figure 35 – Figure 48 for the sphere-cone geometry and in Figure 49 –Figure 62 for the spherical-cap geometry. Two figures are provided for each case: in the first figure, global heating images are shown for each Reynolds number, ordered left-to-right, top-to-bottom in terms of increasing Reynolds number; in the second figure, line plots of h/h_{FR} vs. s/R are shown, ordered left-to-right, top-to-bottom in terms of streamline angular coordinate. For brevity and clarity, all the streamlines are not plotted in these figures. Instead, streamlines are shown at 30-deg increments from 0 deg to 150 deg. As noted earlier, the data for streamlines between 160 deg and 200-deg are considered qualitative, not quantitative, and were thus omitted. Data for streamlines from 210 deg to 360 deg are nominally symmetric with the data from 0 deg to 150 deg, although in practice, model surface irregularities can cause asymmetric behavior. Such local asymmetries can be observed in the images that

accompany the line plots.

In these line-plots, the CFD predictions for smooth-wall, laminar and turbulent heating levels are also shown. Because the laminar heat-transfer film coefficient ratio, h/h_{FR} , remains nearly constant with Reynolds number, only the lowest Reynolds number laminar prediction is shown for each case. However, since this invariance does not hold for turbulent flow, turbulent predictions are shown for each Reynolds number. As noted previously, simulations for these cases were treated as fully turbulent flow over the entire geometry. These turbulent predictions are shown for qualitative comparisons only, since the actual transition occurred at different locations for each test condition / model geometry and since the turbulent heating levels were augmented above smooth-wall levels by the surface roughness.

Roughness Height Effects on Heating and Transition

The same data are shown in the next group of figures, but they are reordered to show the effects of the distributed sand-grain surface roughness height on transition and heating at each Reynolds number. The sphere-cone data are shown in Figure 63 – Figure 74 and the spherical-cap data are shown in Figure 75 – Figure 86. Two figures are provided for each case: in the first figure, global heating images are shown for each Reynolds number, ordered left-to-right, top-to-bottom in terms of increasing roughness height; in the second figure, line plots of h/h_{FR} vs. s/R are shown, ordered left-to-right, top-to-bottom in terms of streamline angular coordinate. As with the Reynold number effects figure set, both laminar and turbulent CFD heating predictions are shown in each line plot. It is assumed that facility noise effects on transition are minimal in these data because the surface OML roughness features (step or gaps) promote a “bypass transition” mode (Ref. [19]) that is separate from the small disturbance growth modes of conventional stability theory analyses.

General Reynolds Number and Roughness Height Trends

In these line plots for Reynolds number and roughness height effects, the laminar CFD predictions allow for baseline assessment of the computational accuracy through comparisons with the low Reynolds number, small roughness height cases for which transition did not occur. In general, good agreement between data and predictions was observed for all laminar cases. However, the turbulent predictions are shown only for illustrative purposes since the fully-turbulent, smooth-wall computations do not account for roughness effects on the transition location or heating augmentation above smooth-surface levels.

Reynolds-number and roughness-height effects on transition and heating follow expected trends. As Reynolds number is increased, the transition onset location moves upstream toward the stagnation point of the model. The transition onset location also moves upstream as roughness height is increased and the measured rough-wall turbulent heating levels grow increasingly higher than the predicted smooth-wall turbulent heating levels.

Roughness Heating Augmentation

In this report, analysis of the heating augmentation due to the roughness patterns is limited to the expected observation that heating levels increase with roughness height. This limitation is due to the complexities of the problem and the goal of quickly releasing this data set as a basis for further analysis. For any given roughness height / Reynolds number / body-point location, the heating augmentation with respect to smooth-wall laminar or turbulent predictions can be determined through reference to the data and figures presented herein. However, the development of engineering correlations or numerical models for simulation of these data depend on not just modeling the effects of roughness on heating, but also modeling

the effects of roughness on transition onset; that is, it is not possible to accurately predict heating levels without being able to first predict the transition onset location. Implementation of a transition model into a CFD code and generation of transitional flow field solutions will be deferred to future in-depth analyses.

Transition Onset Data and Correlation

As per discussion above, transition locations (or lack thereof) were determined along each of 36 different streamlines for every run using the tangent-slope-intercept method. Discounting the windward streamline data from the $\phi = 160\text{-deg}$ through 200-deg rays with poor viewing angles and lighting, 2604 data points (2 model geometries \times 6 free stream Reynolds numbers \times 7 roughness heights \times 31 rays) on the state of the boundary layer were obtained. Tabulations of these transition onset data for each streamline are given in Table 8 – Table 14 for the sphere-cone geometry and in Table 15 – Table 21 for the spherical-cap geometry. Tabular entries are only provided for the 1720 streamlines along which transition was noted. For reference, the boundary-layer momentum thickness, Re_θ , value at transition is also listed. This quantity, along with the roughness height, is a first-order factor in the correlation of roughness transition data.

A correlation for the effects of roughness on boundary-layer transition was developed originally in Ref. [3] based on a survey of transition data on hemispherical geometries with distributed sand-grain roughness as per Eq. (11) through Eq. (13).

$$Y_{TR} = 165 \times (X_{TR})^{-0.5} \quad (11)$$

where

$$Y_{TR} = (Re_\theta)_{TR} \quad (12)$$

$$X_{TR} = \left[\left(\frac{kT_e}{\theta T_w} \right) \left(\frac{H_e}{H_k} \right)^{-1} (M_e)^{-0.5} \right]_{TR} \quad (13)$$

$$k = \begin{cases} k_{PV50} & \text{sand-grain roughness} \\ h_{hex} & \text{hexcomb pattern depth} \end{cases} \quad (14)$$

This original correlation was subsequently modified in Ref. [1] to account for the effects of varying pressure gradient in the non-axisymmetric flows over the current sphere-cone and spherical-cap geometries with patterned hexcomb roughness as per Eqs. (15) through (20). In this correlation, the pressure gradient effects are included through a modified Pohlhausen parameter, β_{POHL} , which is derived from the original Pohlhausen parameter, λ_{POHL} .

$$Y_{TR} = 171.4 \times (X_{TR})^{-0.6299} \quad (15)$$

where

$$Y_{TR} = (Re_\theta)_{TR} \quad (16)$$

$$X_{TR} = \left[\left(\frac{kT_e}{\theta T_w} \right)^{0.45} \left(\frac{H_e}{H_k} \right)^{-1.8} (M_e)^{-0.6} (\beta_{POHL})^{-0.5} \right]_{TR} \quad (17)$$

$$\log(1/\beta_{POHL}) = \bar{\lambda}_{POHL} \quad (18)$$

$$\bar{\lambda}_{POHL} = \min[10, \max(-10, \lambda_{POHL})] \quad (19)$$

$$\lambda_{POHL} = \frac{dU_e}{ds} \frac{\delta^2}{v_e} \quad (20)$$

In the Pohlhausen parameter, the pressure gradient is related to the velocity through the Bernoulli equation, which is approximately valid for the low Mach number boundary-layer conditions of a blunt body. The boundary-layer height, δ , and dynamic viscosity, v_e , terms provide a consistent nondimensionalization of the velocity gradient over a wide range of conditions. However, to be used in the power-law form of this correlation, the pressure gradient term must have a positive value, and must also be limited to prevent computational overflow or underflow at the aeroshell shoulder or stagnation point. Thus, the original Pohlhausen parameter is modified by *min* and *max* limiters and redefined through the *log* function to ensure a positive value.

The utility of this correlation is demonstrated through the plots in Figure 87 for the sphere-cone and spherical-cap hexcomb roughness datasets of Refs. [1–2] and the hemisphere sand-grain roughness dataset of Refs. [3–4]. The majority of transition data points from these two different types of distributed roughness tested in multiple wind tunnels and ballistics ranges fell within a $\pm 20\%$ uncertainty bounds of the correlation.

The current test program permits evaluation of this correlation against distributed sand-grain roughness data from three-dimensional geometries that produce varying pressure gradients. Comparisons of these new data with the correlation are given in Figure 88. While this new distributed sand-grain roughness data set follows the same trend line as the correlation, it exhibits considerably more scatter than earlier data sets, with the majority of points bounded by a higher range of $\pm 50\%$.

However, there is a key characteristic of the current data set that differs from prior data sets, which is the variability of roughness height over the surface of a wind tunnel model. As discussed earlier, the roughness elements in the current data set were found to exhibit an approximately Gaussian height distribution due to the method in which the roughness elements were created. The height values used in this analysis represent the 50% exceedance value, which means that half of the roughness elements have a larger height than the specified value and half have a smaller height. In contrast, the roughness heights for the prior hexcomb studies varied less because they were specified the hexcomb sizes were defined uniformly in the CAD design of the models and were fabricated as such. And for the hemispherical sand-grain roughness data sets, the transition data were based on analyses of multiple streamlines around the circumference of each model – all which experienced the same pressure gradient due to the axisymmetric nature of the flow – which had the effect on normalizing the data for the effects of varying roughness height.

Given these differences in the definition and reporting of the roughness heights between these studies, it is thus not surprising that the current data exhibit much more scatter than the earlier data. The current distributed sand-grain roughness data set is in fact, probably a more “flight-realistic” representation of roughness transition characteristics. The TPS of an actual flight vehicle is likely to exhibit more variability in roughness due to manufacturing defects, damage during handling and flight preparation, and random variations in material response. A real-world example of nonuniform transition response is illustrated in Figure 89 by a post-flight photograph (from Ref. [20]) of the Orion EFT-1 flight test mission heat shield. While the green-outlined areas that identify damage during recovery of the heatshield can be disregarded, the red-outlined areas represent isolated transition wedges caused by nonuniformities in the surface of the heatshield.

Since the new transition data follow the same trend as the earlier data, albeit with much more scatter, no modifications will be made to the transition of Eqs. (15) through (20). However, these wind tunnel data – and the flight data from EFT-1 – illustrate the need to consider large margins on the transition predictions to account for potential variations in roughness heights from an expected nominal value.

Summary

The effects of distributed sand-grain surface roughness patterns simulating that of a heat shield with an ablated TPS on hypersonic boundary-layer transition and turbulent heating have been investigated through wind tunnel testing of two representative entry vehicle geometries. Heating and transition onset data were obtained at Mach 6 over a range of roughness heights and free stream Reynolds numbers sufficient to produce laminar, transitional, and turbulent flow. Boundary-layer transition onset locations were tabulated, and heating distributions were provided in both line plot and image forms. Measured heating levels were found to increase with both Reynolds number and roughness height. The transition onset location trend was found to agree with a correlation developed from prior studies, although with greater scatter due to the variability of the sand-grain roughness heights over the surfaces of the wind tunnel models.

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Table 1. Model geometry parameters.

Model geometry	Model radius, R		Nose radius, R_N		Corner radius, R_c		Nose included angle, β		R_N/R	R_c/R
	in.	m	in.	m	in.	m	deg.			
Sphere-cone	3.000	0.0762	1.500	0.0371	0.1500	0.00381	20.00	0.5	0.050	
Spherical-cap	3.000	0.0762	7.200	0.1829	0.3000	0.00762	23.04	2.4	0.100	

Table 2. ASTM mesh parameters.

ASTM Mesh Designation	Spherical glass particle diameter, D (mil) (mm)		Nominal element height, h_{nom} (mil) (mm)	
ASTM-10	68.90	1.7501	34.45	0.8750
ASTM-20	33.58	0.8529	16.79	0.4265
ASTM-40	16.73	0.1798	8.37	0.0899
ASTM-80	7.09	0.1062	3.54	0.0531
ASTM-140	4.17	0.0630	2.09	0.0315
ASTM-230	2.48	1.7501	1.24	0.8750

Table 3. Roughness data from sample plate scans.

Model ID	Nominal element height h_{nom} (mil)	Measured RMS height h_{RMS} (mil)	50 % exceedance height k_{PV50} (mil)	30 % exceedance height k_{PV30} (mil)	15 % exceedance height k_{PV15} (mil)	5 % exceedance height k_{PV05} (mil)
10-Mesh	34.45	7.873	21.48	25.78	30.15	36.03
20-Mesh	16.79	2.380	8.49	10.19	14.18	16.93
40-Mesh	8.37	1.430	3.83	5.13	6.23	6.86
80-Mesh	3.54	1.181	3.40	4.38	5.09	6.47
140-Mesh	2.09	1.102	2.73	3.38	3.96	4.65
230-Mesh	1.24	0.683	1.94	2.30	2.71	3.32

Table 4. Model roughness information.

ASTM Mesh	Nominal height		Measured RMS height		Original 50% exceedance		Corrected 50% exceedance	
	h (mil)	h (mm)	h_{RMS} (mil)	h_{RMS} (mm)	k_{PV50} (mil)	k_{PV50} (mm)	k_{PV50} (mil)	k_{PV50} (mm)
10	34.45	0.8750	7.873	0.2000	21.48	0.5456	N/A	N/A
20	16.79	0.4265	2.380	0.0605	8.49	0.2156	N/A	N/A
40	8.37	0.2126	1.430	0.0363	3.83	0.0973	N/A	N/A
80	3.54	0.0899	1.181	0.0300	3.40 [#]	0.0864 [#]	2.63	0.0668
140	2.09	0.0531	1.102	0.0280	3.21 [#]	0.0815 [#]	2.15	0.0546
230	1.24	0.0315	0.683	0.0173	1.94 [#]	0.0493 [#]	1.29	0.0328
Smooth	N/A	N/A	0.500	0.0127	1.46 [#]	0.0371 [#]	0.98	0.0249

Values corrected due to lack of scan precision

Table 5. Nominal 20-Inch Mach 6 Air Tunnel Conditions.

Re_∞ (1/ft)	Re_∞ (1/m)	M_∞	T_∞ (K)	ρ_∞ (kg/m ³)	U_∞ (m/s)	ΔH (J/kg)	h_{FR} (kg/m ² -s)	
							sphere-cone	spherical-cap
2.051E+06	6.725E+06	5.968	62.94	3.203E-02	948.4	2.121E+05	1.144E-01	2.536E-01
2.992E+06	9.816E+06	5.998	63.26	4.666E-02	955.1	2.187E+05	1.401E-01	3.072E-01
4.966E+06	1.629E+07	6.030	63.72	7.741E-02	962.9	2.266E+05	1.824E-01	3.992E-01
6.489E+06	2.129E+07	6.042	63.48	1.008E-01	962.1	2.256E+05	2.078E-01	4.555E-01
7.199E+06	2.362E+07	6.047	63.48	1.118E-01	962.7	2.261E+05	2.189E-01	4.800E-01
8.136E+06	2.669E+07	6.034	59.21	1.224E-01	924.9	1.860E+05	2.180E-01	4.781E-01

Table 6. 20-Inch Mach 6 Air Tunnel Test 7036 run matrix.

Run	Geometry	ASTM Mesh	α (deg)	Re_{∞} (1/ft)	M_{∞}	T_{∞} (K)	ρ_{∞} (kg/m³)	U_{∞} (m/s)	ΔH (J/kg)	h_{FR} (kg/m²-s)	$T_{w,avg}$ (K)
50	sphere-cone	smooth	16	2.02E+06	5.96	62.9	3.161E-02	948.0	2.116E+05	2.505E-01	331
51	sphere-cone	smooth	16	2.96E+06	6.00	63.4	4.629E-02	956.1	2.198E+05	3.062E-01	342
52	sphere-cone	smooth	16	4.96E+06	6.03	63.7	7.733E-02	962.4	2.262E+05	3.990E-01	359
53	sphere-cone	smooth	16	6.47E+06	6.04	63.6	1.007E-01	962.7	2.262E+05	4.554E-01	371
54	sphere-cone	smooth	16	7.23E+06	6.05	63.4	1.122E-01	961.5	2.248E+05	4.799E-01	378
55	sphere-cone	smooth	16	8.13E+06	6.03	59.2	1.224E-01	924.8	1.859E+05	4.778E-01	368
74	sphere-cone	230	16	2.02E+06	5.97	63.2	3.164E-02	949.8	2.136E+05	2.512E-01	333
75	sphere-cone	230	16	2.98E+06	6.00	63.4	4.660E-02	956.1	2.198E+05	3.073E-01	344
76	sphere-cone	230	16	4.99E+06	6.03	63.6	7.772E-02	961.2	2.249E+05	3.994E-01	368
77	sphere-cone	230	16	6.45E+06	6.04	63.7	1.005E-01	963.8	2.274E+05	4.556E-01	387
78	sphere-cone	230	16	7.19E+06	6.05	63.5	1.118E-01	963.0	2.265E+05	4.800E-01	395
79	sphere-cone	230	16	8.21E+06	6.03	59.0	1.232E-01	922.4	1.834E+05	4.779E-01	380
80	sphere-cone	140	16	2.01E+06	5.97	63.4	3.159E-02	951.8	2.157E+05	2.516E-01	332
81	sphere-cone	140	16	2.97E+06	6.00	63.3	4.636E-02	955.1	2.187E+05	3.061E-01	342
82	sphere-cone	140	16	4.99E+06	6.03	63.6	7.771E-02	961.6	2.253E+05	3.995E-01	364
83	sphere-cone	140	16	6.48E+06	6.04	63.6	1.007E-01	962.7	2.262E+05	4.555E-01	387
84	sphere-cone	140	16	7.16E+06	6.05	63.7	1.115E-01	964.5	2.282E+05	4.803E-01	394
85	sphere-cone	140	16	8.15E+06	6.03	59.2	1.227E-01	924.8	1.859E+05	4.784E-01	381
87	sphere-cone	80	16	2.01E+06	5.97	63.4	3.151E-02	951.8	2.157E+05	2.513E-01	331
88	sphere-cone	80	16	2.97E+06	6.00	63.7	4.648E-02	958.4	2.223E+05	3.078E-01	343
89	sphere-cone	80	16	4.98E+06	6.03	63.7	7.770E-02	962.5	2.262E+05	3.999E-01	374
90	sphere-cone	80	16	6.51E+06	6.04	63.5	1.012E-01	962.0	2.255E+05	4.561E-01	394
91	sphere-cone	80	16	7.15E+06	6.05	63.8	1.113E-01	964.9	2.286E+05	4.802E-01	400
92	sphere-cone	80	16	8.14E+06	6.03	59.3	1.225E-01	925.1	1.862E+05	4.782E-01	386
93	sphere-cone	40	16	2.01E+06	5.97	63.2	3.148E-02	950.4	2.143E+05	2.507E-01	332
94	sphere-cone	40	16	2.97E+06	6.00	63.7	4.649E-02	958.5	2.224E+05	3.079E-01	350
95	sphere-cone	40	16	4.95E+06	6.03	63.9	7.733E-02	964.4	2.282E+05	3.999E-01	389
96	sphere-cone	40	16	6.48E+06	6.04	63.5	1.008E-01	962.1	2.256E+05	4.553E-01	406
97	sphere-cone	40	16	7.23E+06	6.05	63.4	1.122E-01	961.8	2.252E+05	4.801E-01	412
98	sphere-cone	40	16	8.20E+06	6.03	59.1	1.231E-01	923.3	1.844E+05	4.784E-01	394
99	sphere-cone	20	16	2.01E+06	5.97	63.4	3.157E-02	951.4	2.153E+05	2.514E-01	350
100	sphere-cone	20	16	2.98E+06	6.00	63.5	4.653E-02	956.5	2.203E+05	3.072E-01	375
101	sphere-cone	20	16	4.92E+06	6.03	63.8	7.679E-02	963.4	2.272E+05	3.981E-01	405
102	sphere-cone	20	16	6.47E+06	6.04	63.5	1.005E-01	961.9	2.254E+05	4.546E-01	415
103	sphere-cone	20	16	7.24E+06	6.05	63.3	1.123E-01	960.7	2.241E+05	4.798E-01	414
104	sphere-cone	20	16	8.17E+06	6.03	59.0	1.227E-01	922.5	1.835E+05	4.770E-01	404
107	sphere-cone	10	16	2.01E+06	5.97	63.3	3.153E-02	951.1	2.150E+05	2.512E-01	374
106	sphere-cone	10	16	2.97E+06	6.00	63.7	4.646E-02	958.6	2.225E+05	3.078E-01	393
105	sphere-cone	10	16	4.98E+06	6.03	63.9	7.776E-02	964.2	2.280E+05	4.010E-01	409
108	sphere-cone	10	16	6.45E+06	6.04	63.7	1.004E-01	964.1	2.277E+05	4.555E-01	393
109	sphere-cone	10	16	7.21E+06	6.05	63.4	1.119E-01	961.6	2.250E+05	4.795E-01	375
110	sphere-cone	10	16	8.20E+06	6.03	59.0	1.232E-01	922.9	1.840E+05	4.782E-01	407

Table 7. 20-Inch Mach 6 Air Tunnel Test 7057 run matrix.

Run	Geometry	ASTM Mesh	α (deg)	Re_∞ (1/ft)	M_∞	T_∞ (K)	ρ_∞ (kg/m ³)	U_∞ (m/s)	ΔH (J/kg)	h_{FR} (kg/m ² -s)	$T_{w,avg}$ (K)
7	spherical-cap	smooth	28	2.04E+06	5.97	62.6	3.178E-02	946.0	2.094E+05	9.900E-02	332
8	spherical-cap	smooth	28	3.01E+06	6.00	63.0	4.676E-02	953.4	2.168E+05	1.212E-01	342
9	spherical-cap	smooth	28	5.03E+06	6.03	63.4	7.821E-02	960.8	2.243E+05	1.583E-01	357
10	spherical-cap	smooth	28	6.52E+06	6.04	63.3	1.011E-01	960.6	2.239E+05	1.799E-01	366
11	spherical-cap	smooth	28	7.28E+06	6.05	63.0	1.126E-01	959.1	2.223E+05	2.188E-01	364
12	spherical-cap	smooth	28	8.02E+06	6.04	59.5	1.208E-01	928.0	1.892E+05	2.176E-01	356
13	spherical-cap	230	28	2.02E+06	5.97	62.8	3.152E-02	947.8	2.113E+05	1.141E-01	329
14	spherical-cap	230	28	2.99E+06	6.00	63.1	4.665E-02	954.6	2.181E+05	1.400E-01	338
15	spherical-cap	230	28	4.94E+06	6.03	63.9	7.702E-02	964.5	2.283E+05	1.822E-01	357
16	spherical-cap	230	28	6.47E+06	6.04	63.5	1.006E-01	962.6	2.261E+05	2.077E-01	374
17	spherical-cap	230	28	7.20E+06	6.05	63.4	1.118E-01	961.9	2.253E+05	2.187E-01	386
18	spherical-cap	230	28	8.20E+06	6.03	59.0	1.230E-01	923.0	1.839E+05	2.181E-01	375
19	spherical-cap	140	28	2.04E+06	5.97	62.7	3.180E-02	946.8	2.102E+05	1.145E-01	329
20	spherical-cap	140	28	3.01E+06	6.00	62.9	4.687E-02	952.8	2.161E+05	1.401E-01	340
21	spherical-cap	140	28	4.95E+06	6.03	63.8	7.711E-02	963.6	2.273E+05	1.821E-01	384
22	spherical-cap	140	28	6.56E+06	6.04	63.1	1.016E-01	959.6	2.228E+05	2.080E-01	400
23	spherical-cap	140	28	7.15E+06	6.05	63.6	1.111E-01	964.1	2.276E+05	2.187E-01	405
24	spherical-cap	140	28	8.16E+06	6.03	59.0	1.225E-01	923.3	1.842E+05	2.178E-01	390
25	spherical-cap	80	28	2.04E+06	5.97	62.7	3.188E-02	947.3	2.107E+05	1.147E-01	330
26	spherical-cap	80	28	3.00E+06	6.00	63.1	4.676E-02	954.5	2.180E+05	1.402E-01	359
28	spherical-cap	80	28	4.97E+06	6.03	63.6	7.742E-02	961.8	2.254E+05	1.821E-01	398
29	spherical-cap	80	28	6.48E+06	6.04	63.4	1.006E-01	961.7	2.251E+05	2.075E-01	409
30	spherical-cap	80	28	7.19E+06	6.05	63.5	1.116E-01	963.1	2.266E+05	2.189E-01	422
31	spherical-cap	80	28	8.06E+06	6.04	59.3	1.214E-01	926.2	1.874E+05	2.176E-01	398
32	spherical-cap	40	28	2.05E+06	5.97	62.7	3.198E-02	947.3	2.108E+05	1.149E-01	338
33	spherical-cap	40	28	2.99E+06	6.00	63.2	4.665E-02	955.3	2.188E+05	1.402E-01	366
34	spherical-cap	40	28	4.97E+06	6.03	63.8	7.745E-02	963.7	2.274E+05	1.825E-01	402
35	spherical-cap	40	28	6.45E+06	6.04	63.6	1.003E-01	963.2	2.268E+05	2.076E-01	415
36	spherical-cap	40	28	7.19E+06	6.05	63.5	1.116E-01	962.9	2.264E+05	2.189E-01	418
37	spherical-cap	40	28	8.08E+06	6.04	59.4	1.217E-01	926.8	1.880E+05	2.180E-01	400
55	spherical-cap	20	28	2.04E+06	5.97	62.7	3.178E-02	947.3	2.107E+05	1.145E-01	342
56	spherical-cap	20	28	3.07E+06	6.00	62.4	4.755E-02	948.7	2.117E+05	1.403E-01	361
57	spherical-cap	20	28	4.98E+06	6.03	63.8	7.766E-02	963.4	2.271E+05	1.827E-01	393
58	spherical-cap	20	28	6.45E+06	6.04	63.6	1.003E-01	963.9	2.275E+05	2.077E-01	407
59	spherical-cap	20	28	7.22E+06	6.05	63.4	1.120E-01	962.6	2.261E+05	2.192E-01	414
60	spherical-cap	20	28	8.10E+06	6.04	59.2	1.219E-01	925.5	1.865E+05	2.178E-01	394
42	spherical-cap	10	28	2.01E+06	5.97	62.9	3.141E-02	948.1	2.117E+05	1.140E-01	359
43	spherical-cap	10	28	3.00E+06	6.00	63.1	4.672E-02	954.6	2.181E+05	1.401E-01	382
44	spherical-cap	10	28	4.94E+06	6.03	63.9	7.710E-02	964.6	2.285E+05	1.823E-01	408
45	spherical-cap	10	28	6.46E+06	6.04	63.6	1.004E-01	963.4	2.270E+05	2.077E-01	416
46	spherical-cap	10	28	7.19E+06	6.05	63.5	1.116E-01	962.8	2.262E+05	2.188E-01	411
47	spherical-cap	10	28	8.16E+06	6.03	59.1	1.226E-01	923.9	1.849E+05	2.180E-01	396

Table 14. Test 7036 sphere-cone smooth-OML transition locations.

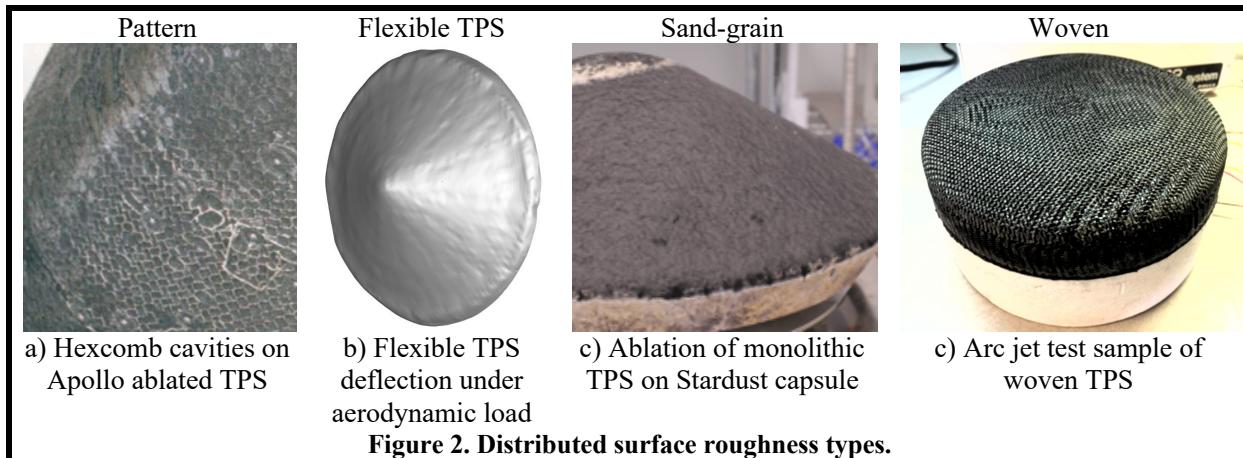
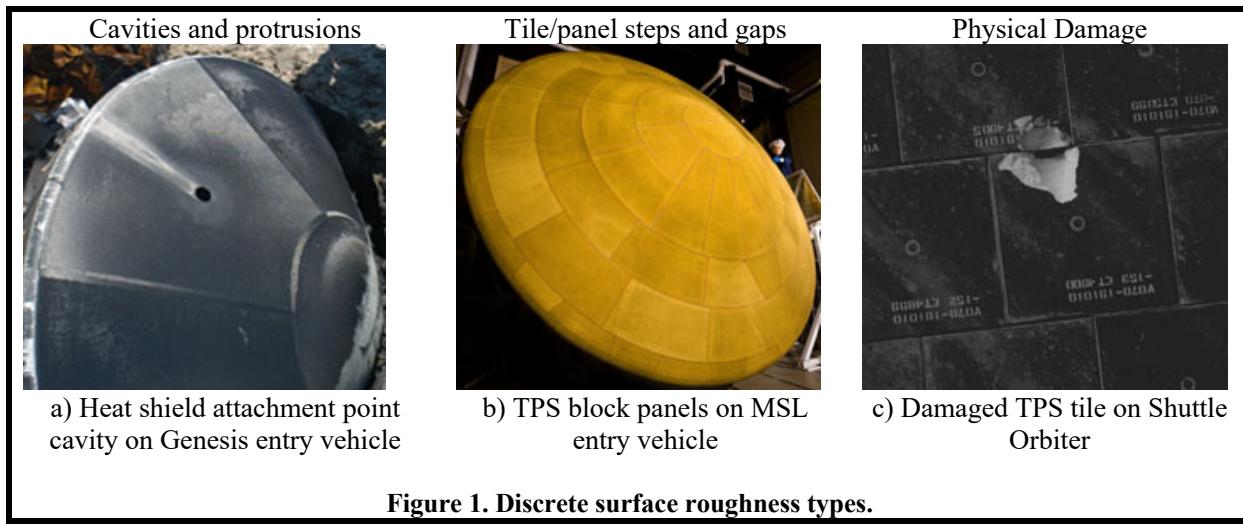
Model	Run	Ray	s ₀ /R	Re _θ
Smooth	53	0	0.883	246.0
Smooth	53	10	0.971	246.6
Smooth	53	20	0.853	216.9
Smooth	53	30	1.070	234.8
Smooth	53	40	0.733	182.1
Smooth	53	50	0.792	183.2
Smooth	53	60	0.773	174.1
Smooth	53	70	0.778	167.6
Smooth	53	80	0.961	174.8
Smooth	53	280	0.987	175.0
Smooth	53	290	0.876	179.3
Smooth	53	300	0.767	174.0
Smooth	53	310	0.846	190.8
Smooth	53	320	0.896	208.7
Smooth	53	330	0.988	226.8
Smooth	53	340	0.830	213.7
Smooth	53	350	0.977	247.1
Smooth	54	0	0.831	251.8
Smooth	54	10	0.849	244.1
Smooth	54	20	0.819	222.3
Smooth	54	30	0.824	213.9
Smooth	54	40	0.689	184.4
Smooth	54	50	0.717	178.7
Smooth	54	60	0.695	170.4
Smooth	54	70	0.724	168.0
Smooth	54	80	0.792	169.4
Smooth	54	90	0.855	169.2
Smooth	54	270	0.914	173.0
Smooth	54	280	0.788	168.1
Smooth	54	290	0.719	167.8
Smooth	54	300	0.684	167.8
Smooth	54	310	0.675	172.5
Smooth	54	320	0.758	197.3
Smooth	54	330	0.823	213.5
Smooth	54	340	0.638	183.5
Smooth	54	350	0.870	247.4
Smooth	55	0	0.824	258.0
Smooth	55	10	0.842	250.3
Smooth	55	20	0.809	228.0
Smooth	55	30	0.777	209.9
Smooth	55	40	0.645	177.4
Smooth	55	50	0.650	173.2
Smooth	55	60	0.588	156.1
Smooth	55	70	0.714	169.3
Smooth	55	80	0.740	166.5
Smooth	55	90	0.807	170.9
Smooth	55	270	0.837	172.5
Smooth	55	280	0.728	165.0
Smooth	55	290	0.696	165.4
Smooth	55	300	0.671	172.4
Smooth	55	310	0.619	167.2
Smooth	55	320	0.690	188.0
Smooth	55	330	0.788	211.6
Smooth	55	340	0.828	232.0
Smooth	55	350	0.822	247.4

Table 20. Test 7057 spherical-cap 230-mesh transition locations.

Model	Run	Ray	s ₀ /R	Re _θ	Model	Run	Ray	s ₀ /R	Re _θ
230-Mesh 14	320	1.166	153.8		230-Mesh 18	110	0.733	142.1	
230-Mesh 15	0	1.055	185.5		230-Mesh 18	120	0.736	136.0	
230-Mesh 15	10	0.980	176.4		230-Mesh 18	250	0.786	149.4	
230-Mesh 15	20	1.383	225.9		230-Mesh 18	260	0.708	142.9	
230-Mesh 15	50	1.120	188.6		230-Mesh 18	270	0.717	148.2	
230-Mesh 15	320	0.822	149.7		230-Mesh 18	280	0.748	156.6	
230-Mesh 15	340	0.958	172.3		230-Mesh 18	290	0.755	160.6	
230-Mesh 16	0	0.932	188.6		230-Mesh 18	300	0.781	167.6	
230-Mesh 16	10	0.912	186.5		230-Mesh 18	310	0.726	161.2	
230-Mesh 16	20	0.920	187.5		230-Mesh 18	320	0.706	160.2	
230-Mesh 16	30	0.867	176.4		230-Mesh 18	330	0.935	202.5	
230-Mesh 16	40	0.886	177.1		230-Mesh 18	340	0.613	147.3	
230-Mesh 16	50	0.879	174.3		230-Mesh 18	350	0.847	189.7	
230-Mesh 16	60	1.072	198.7						
230-Mesh 16	70	1.018	188.6						
230-Mesh 16	80	0.987	179.5						
230-Mesh 16	90	1.102	183.6						
230-Mesh 16	100	0.874	154.2						
230-Mesh 16	260	0.735	136.1						
230-Mesh 16	270	0.931	167.7						
230-Mesh 16	290	1.133	202.6						
230-Mesh 16	300	1.138	206.6						
230-Mesh 16	310	1.044	198.9						
230-Mesh 16	320	0.757	156.9						
230-Mesh 16	330	0.989	196.4						
230-Mesh 16	340	0.836	174.2						
230-Mesh 16	350	0.913	186.5						
230-Mesh 17	0	0.777	172.3						
230-Mesh 17	10	0.746	166.4						
230-Mesh 17	20	0.650	148.6						
230-Mesh 17	30	0.675	152.2						
230-Mesh 17	40	0.754	165.6						
230-Mesh 17	50	0.800	171.8						
230-Mesh 17	60	0.817	170.6						
230-Mesh 17	70	0.750	156.5						
230-Mesh 17	80	0.716	148.0						
230-Mesh 17	90	0.934	175.4						
230-Mesh 17	100	0.849	157.9						
230-Mesh 17	110	0.747	140.1						
230-Mesh 17	120	0.746	135.3						
230-Mesh 17	250	0.802	146.9						
230-Mesh 17	260	0.719	140.4						
230-Mesh 17	270	0.721	145.3						
230-Mesh 17	280	0.781	159.1						
230-Mesh 17	290	0.802	165.0						
230-Mesh 17	300	0.798	167.8						
230-Mesh 17	310	0.733	159.8						
230-Mesh 17	320	0.718	159.4						
230-Mesh 17	330	0.960	200.9						
230-Mesh 17	340	0.644	147.6						
230-Mesh 17	350	0.860	185.5						
230-Mesh 18	0	0.768	174.1						
230-Mesh 18	10	0.729	168.9						
230-Mesh 18	20	0.621	148.6						
230-Mesh 18	30	0.652	151.5						
230-Mesh 18	40	0.749	167.6						
230-Mesh 18	50	0.787	171.8						
230-Mesh 18	60	0.799	170.7						
230-Mesh 18	70	0.740	158.9						
230-Mesh 18	80	0.709	150.2						
230-Mesh 18	90	0.932	180.0						
230-Mesh 18	100	0.844	162.0						

Table 21. Test 7057 spherical-cap smooth-OML transition locations.

Model	Run	Ray	s₀/R	Re_θ
Smooth	10	310	1.393	245.5
Smooth	10	320	1.371	245.4
Smooth	10	330	1.390	247.7
Smooth	10	340	1.336	249.4
Smooth	11	0	1.311	262.9
Smooth	11	10	1.344	266.2
Smooth	11	20	0.902	197.0
Smooth	11	30	1.040	217.9
Smooth	11	40	1.170	235.2
Smooth	11	50	1.152	231.5
Smooth	11	60	1.158	225.5
Smooth	11	70	1.106	214.0
Smooth	11	300	1.124	221.1
Smooth	11	310	1.080	220.2
Smooth	11	320	1.025	213.4
Smooth	11	330	1.044	218.7
Smooth	11	340	1.115	231.9
Smooth	11	350	1.261	253.5
Smooth	12	0	1.296	266.1
Smooth	12	10	1.330	270.6
Smooth	12	20	0.818	187.9
Smooth	12	30	1.154	239.9
Smooth	12	40	1.318	260.9
Smooth	12	50	1.373	265.6
Smooth	12	300	1.177	234.7
Smooth	12	310	1.246	251.4
Smooth	12	320	1.052	223.2
Smooth	12	330	1.099	232.4
Smooth	12	340	1.103	236.5
Smooth	12	350	1.250	258.1



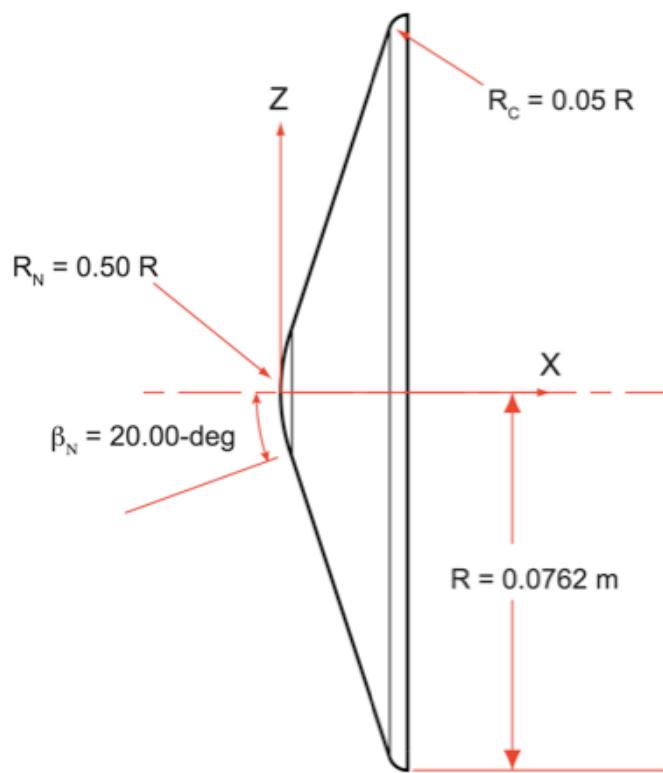


Figure 3. Sphere-cone geometry.

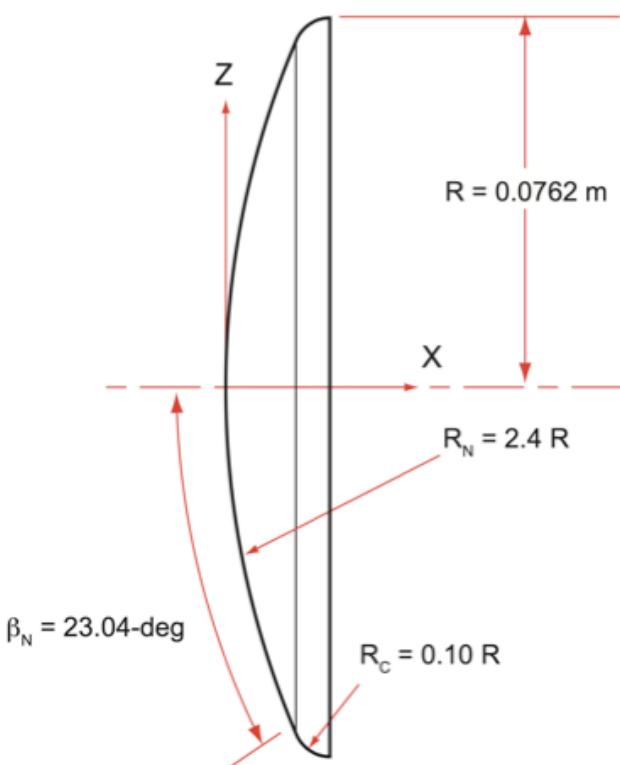


Figure 4. Spherical-cap geometry.

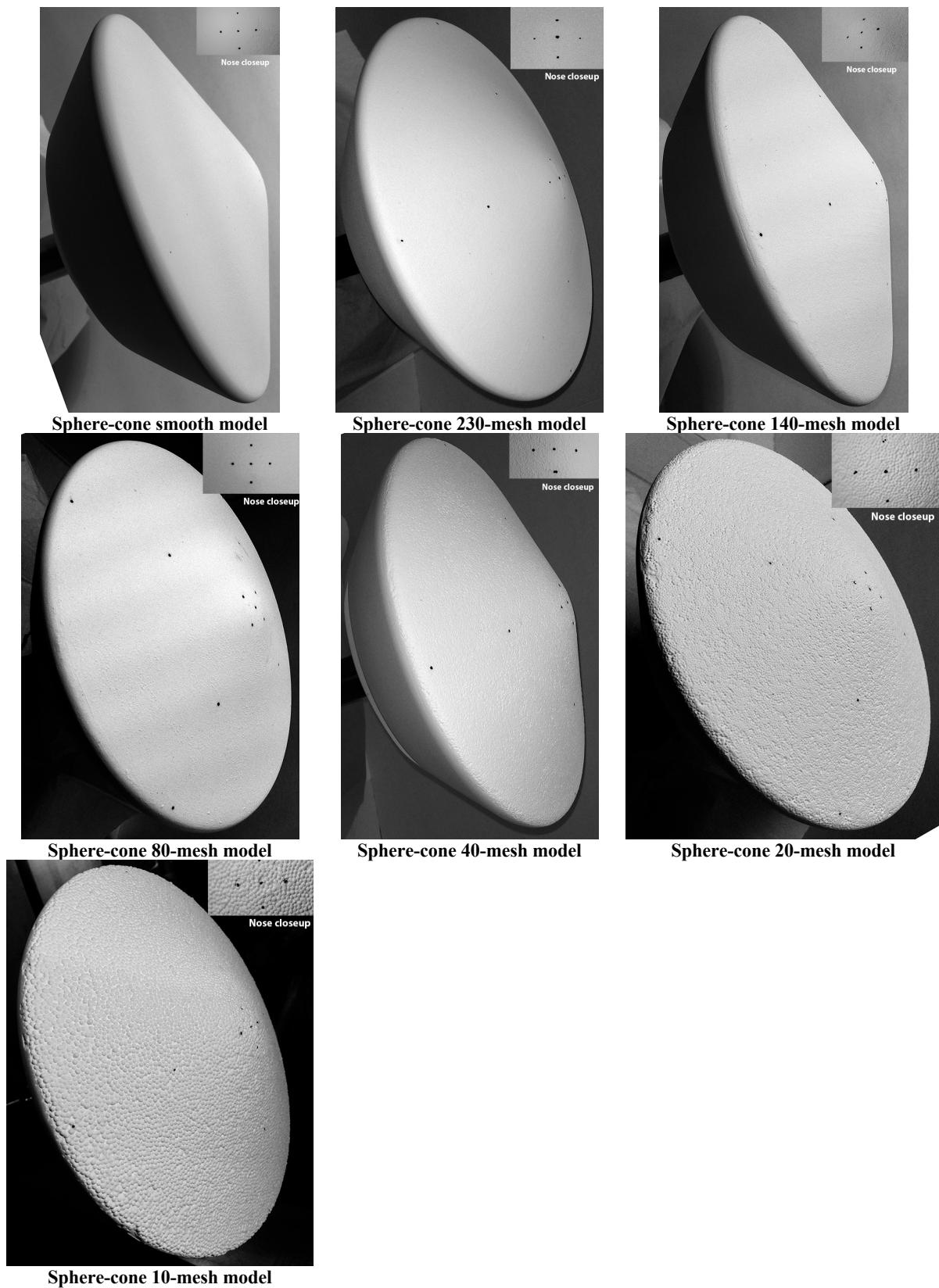
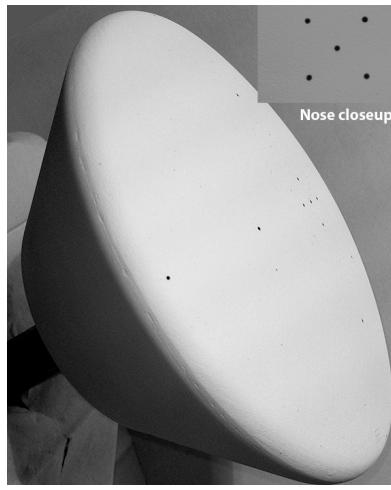


Figure 5. Sphere-cone model photographs



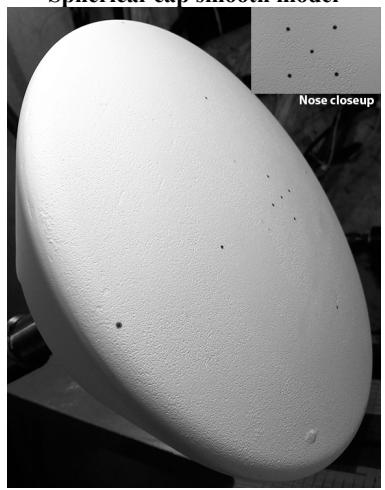
Spherical-cap smooth model



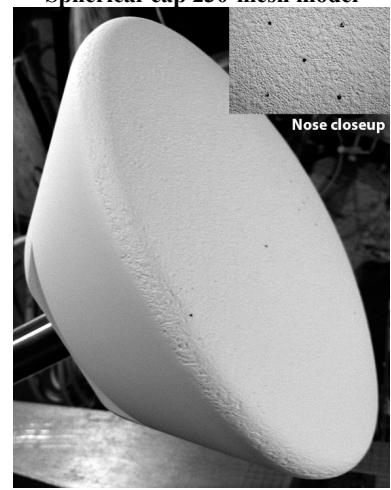
Spherical-cap 230-mesh model



Spherical-cap 140-mesh model



Spherical-cap 80-mesh model



Spherical-cap 40-mesh model



Spherical-cap 20-mesh model



Spherical-cap 10-mesh model

Figure 6. Spherical-cap model photographs

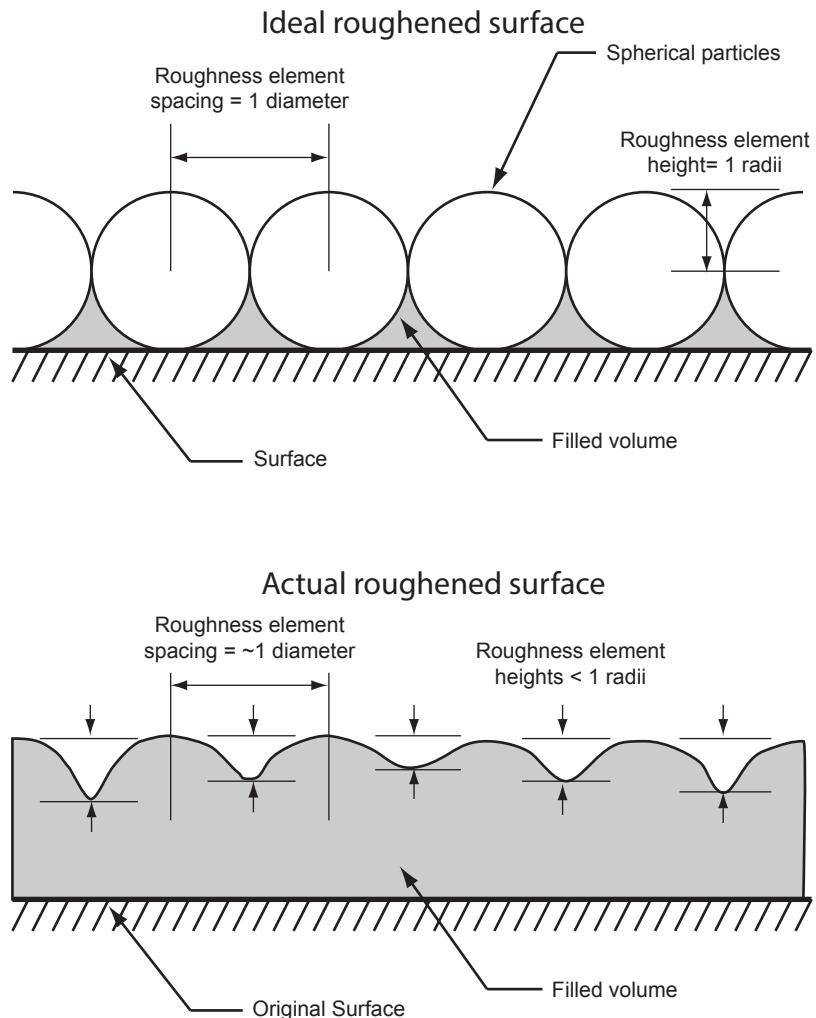


Figure 7. Illustration of ideal and actual surface roughness.

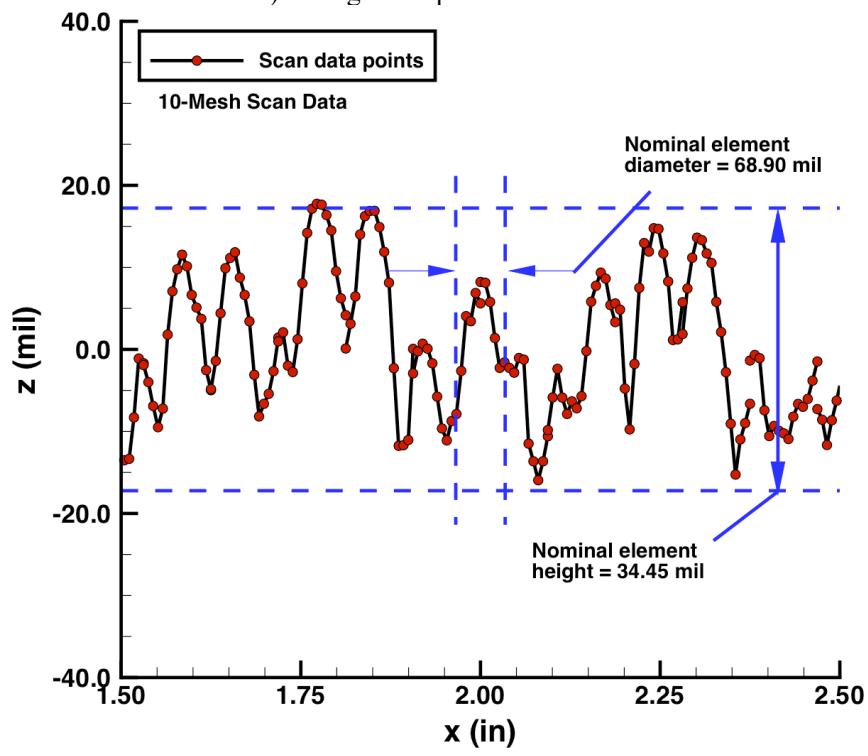
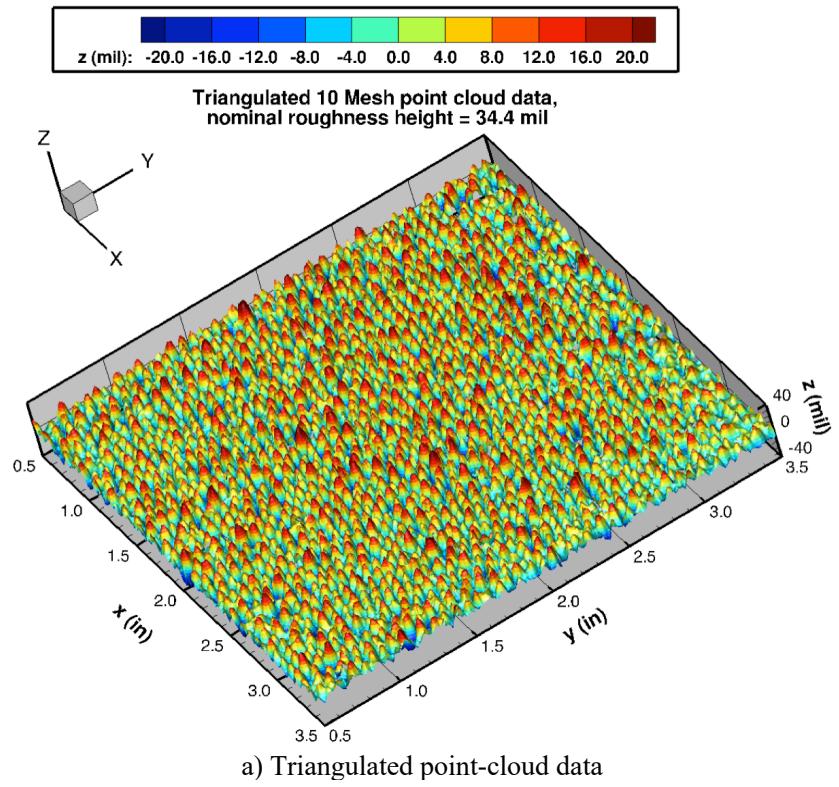


Figure 8. 10-Mesh sample plate scan data.

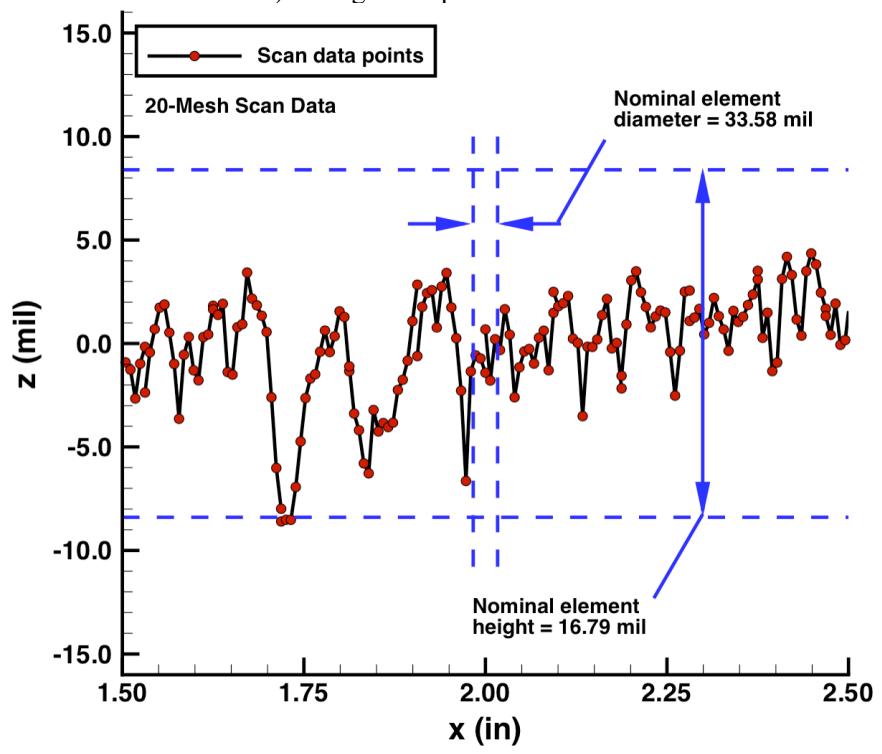
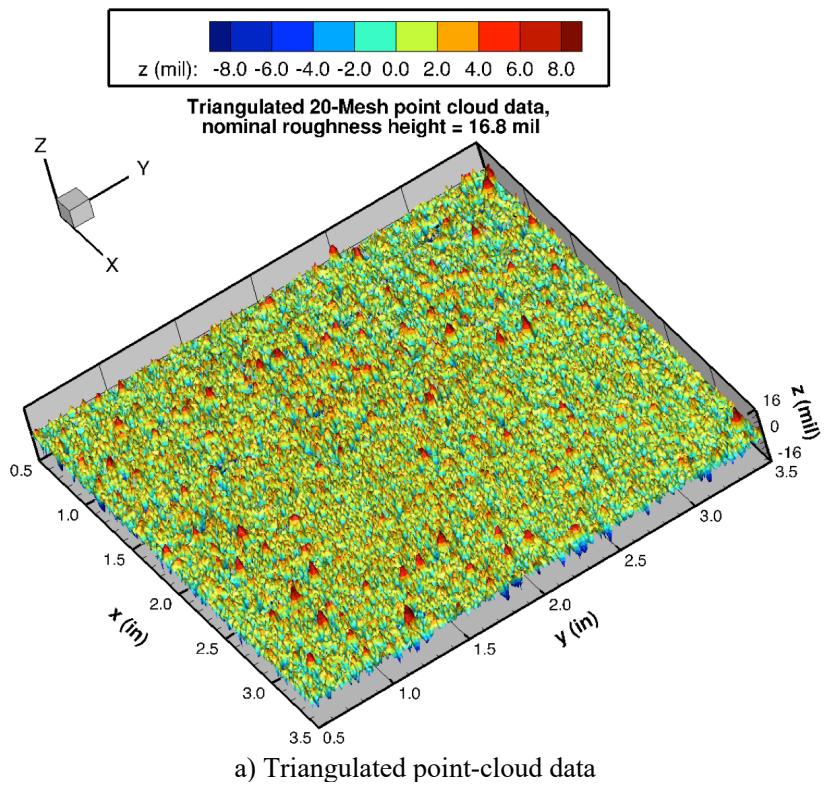


Figure 9. 20-Mesh sample plate scan data.

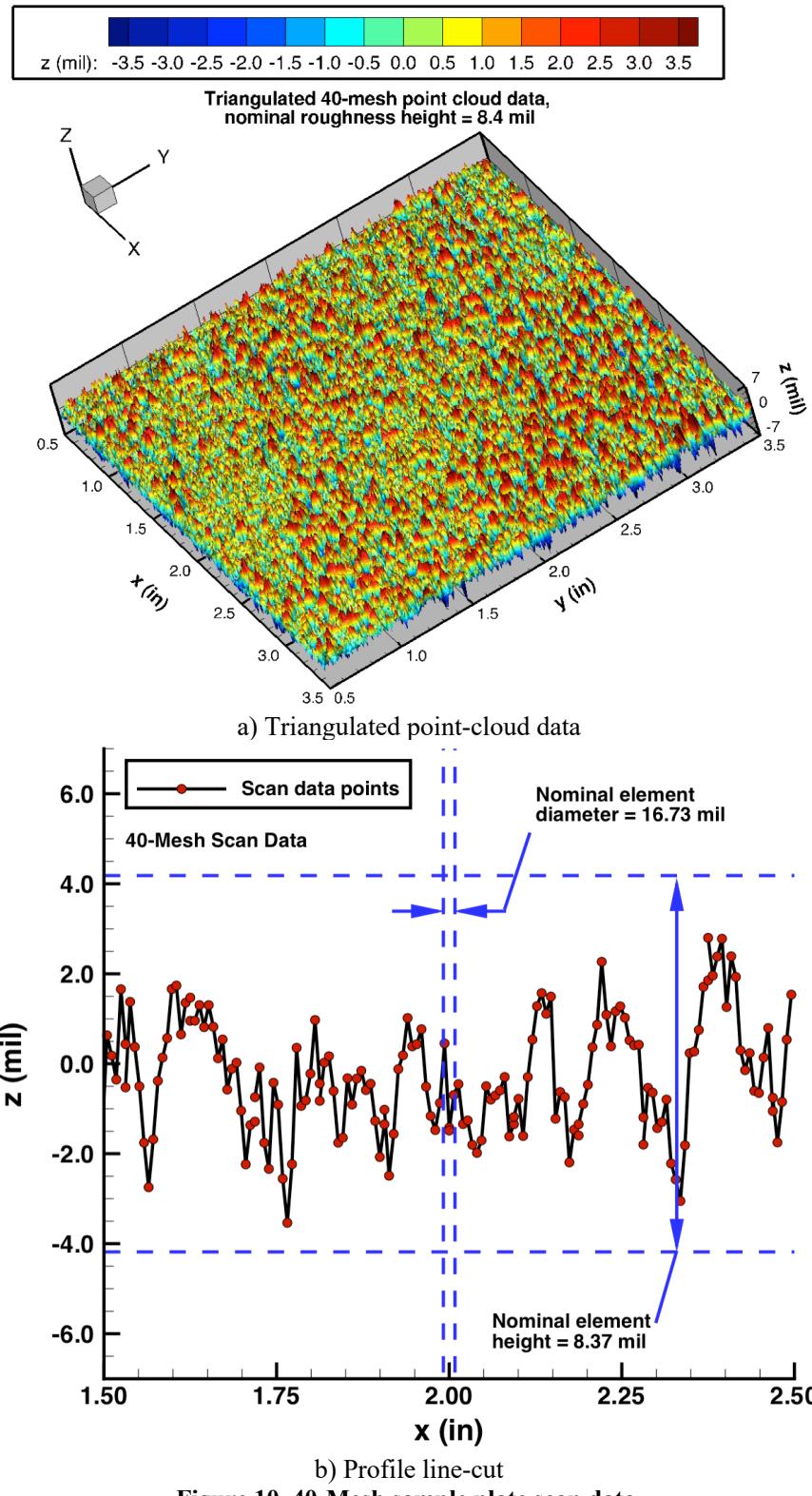


Figure 10. 40-Mesh sample plate scan data.

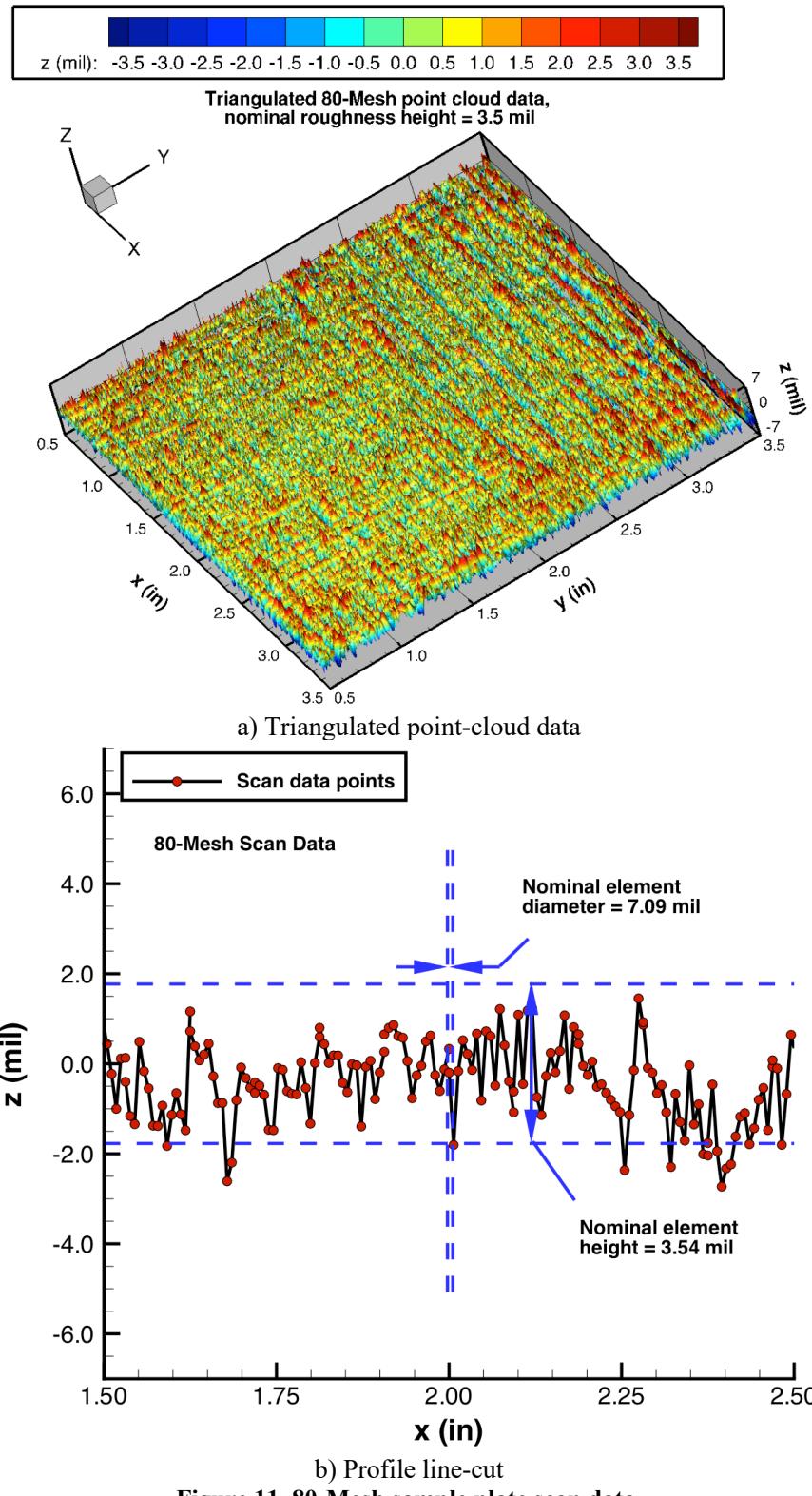


Figure 11. 80-Mesh sample plate scan data.

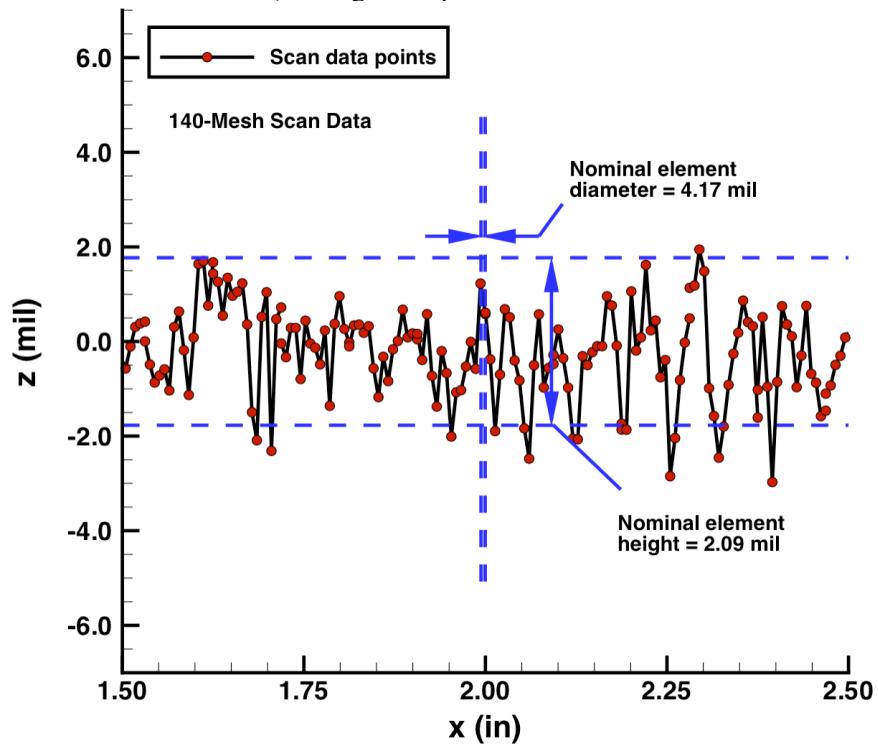
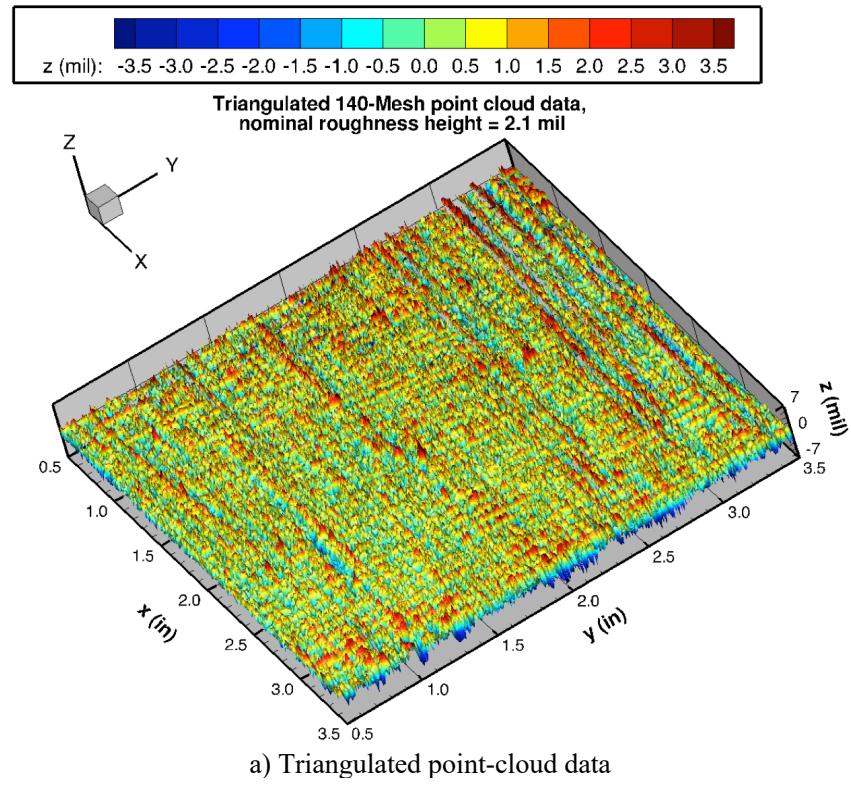


Figure 12. 140-Mesh sample plate scan data.

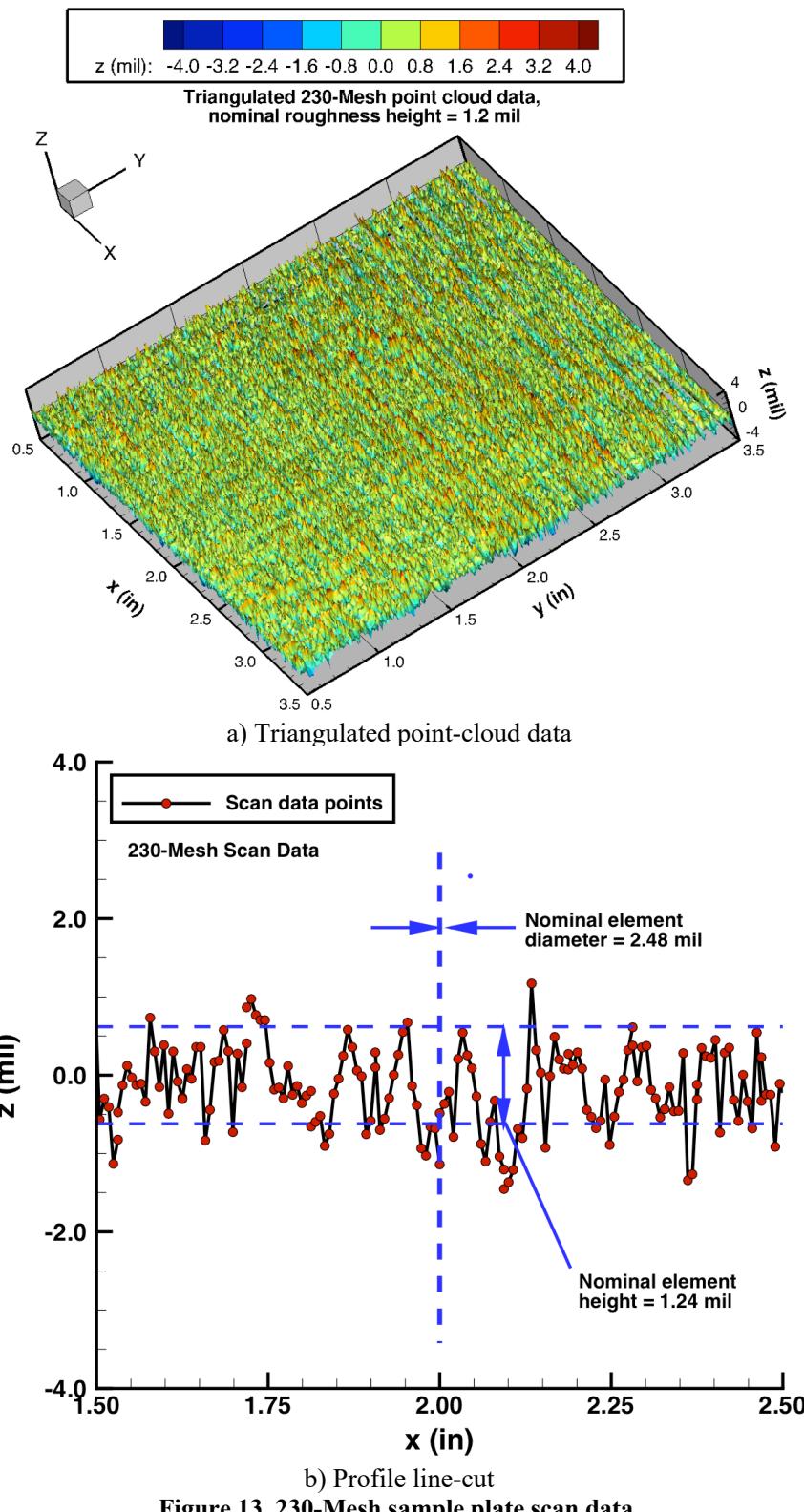


Figure 13. 230-Mesh sample plate scan data.

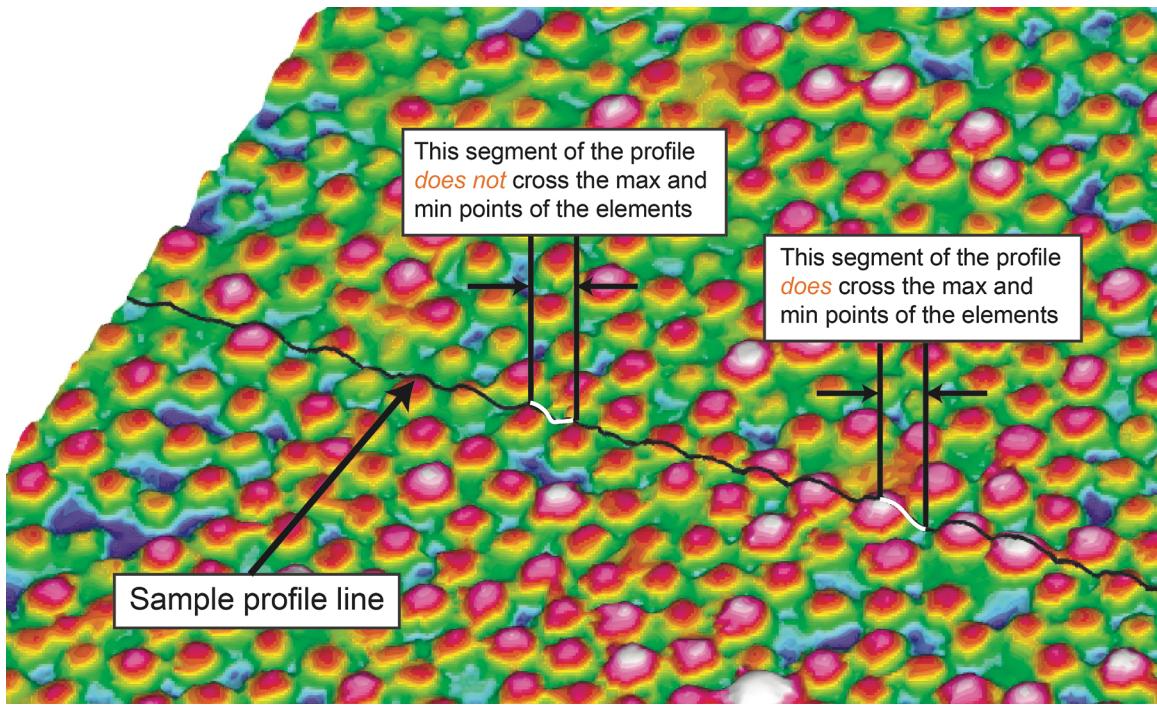


Figure 14. Profile alignment with roughness elements.

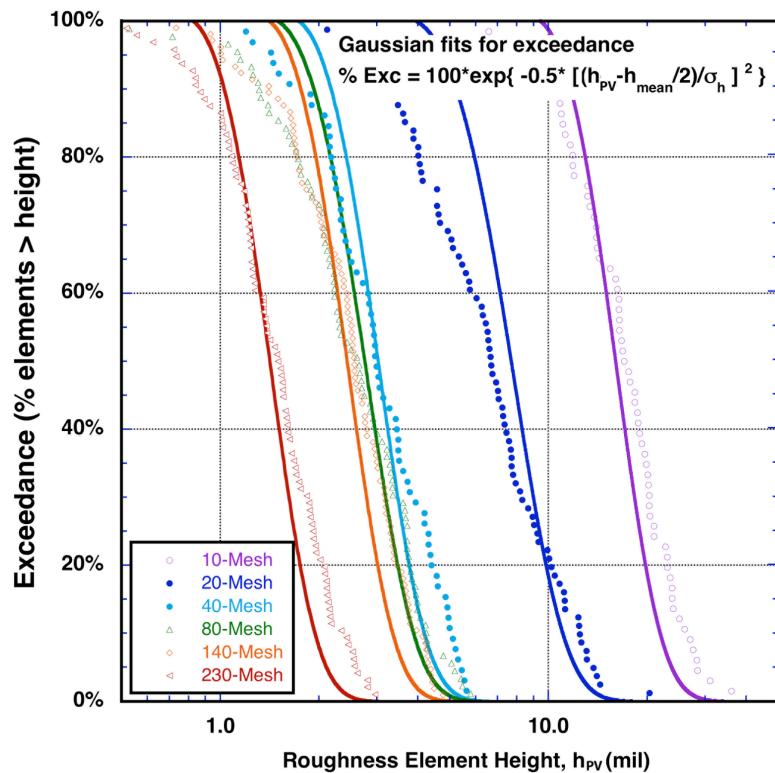


Figure 15. Roughness height probability of exceedance distributions.

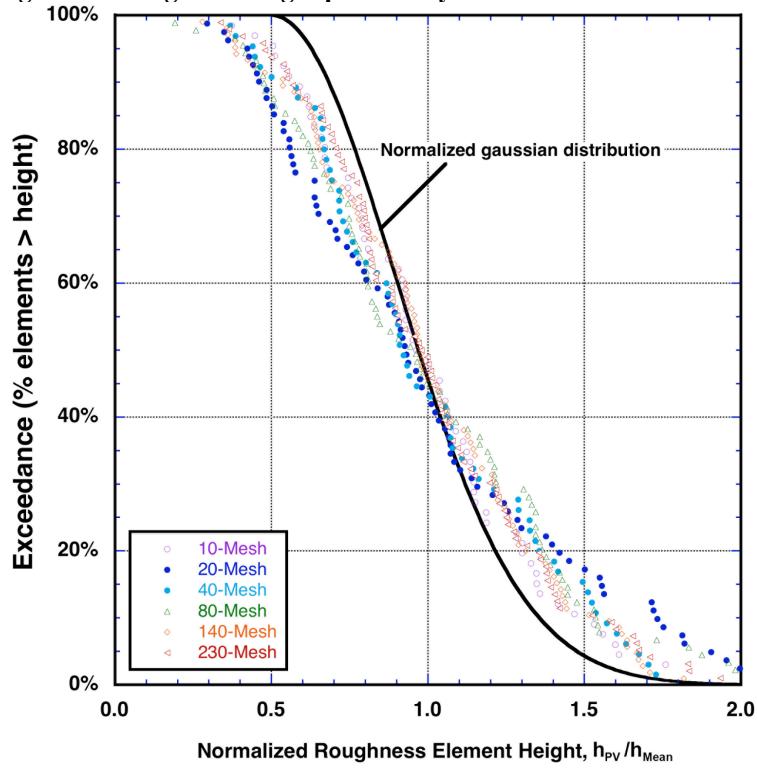


Figure 16. Normalized exceedance distributions.

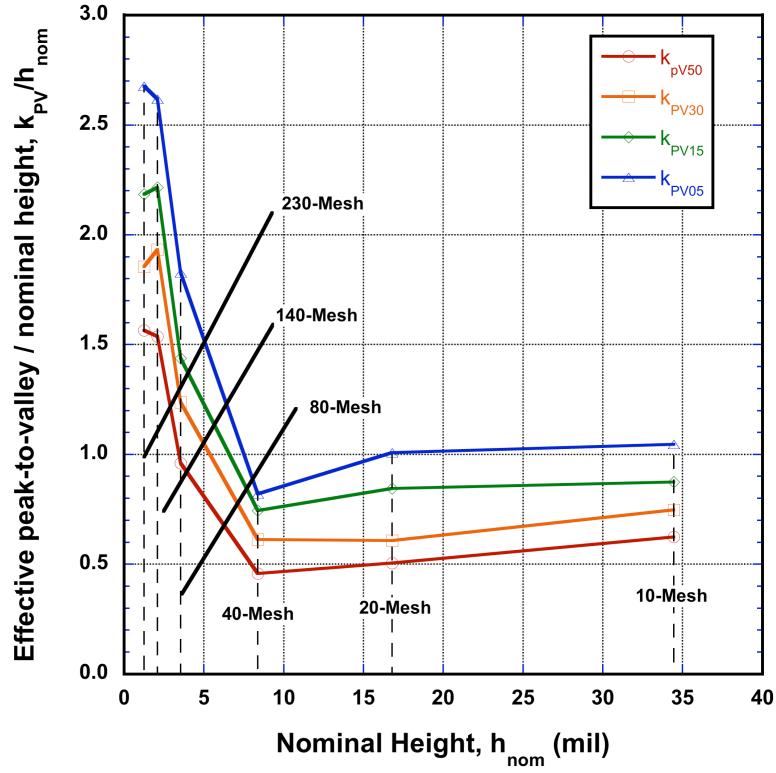


Figure 17. Comparison of effective and nominal roughness heights.

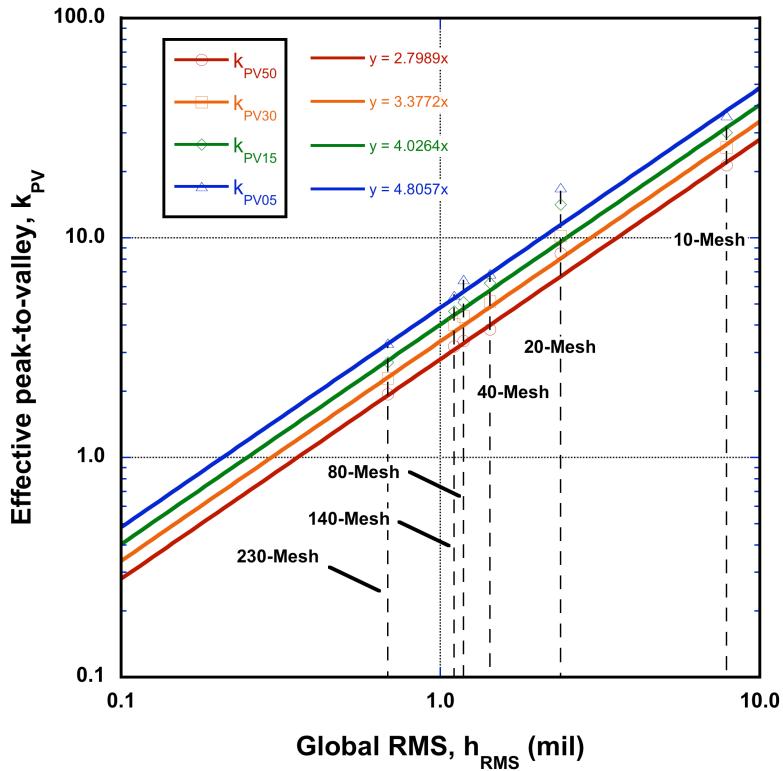


Figure 18. Relationship between effective roughness heights and measured RMS heights.

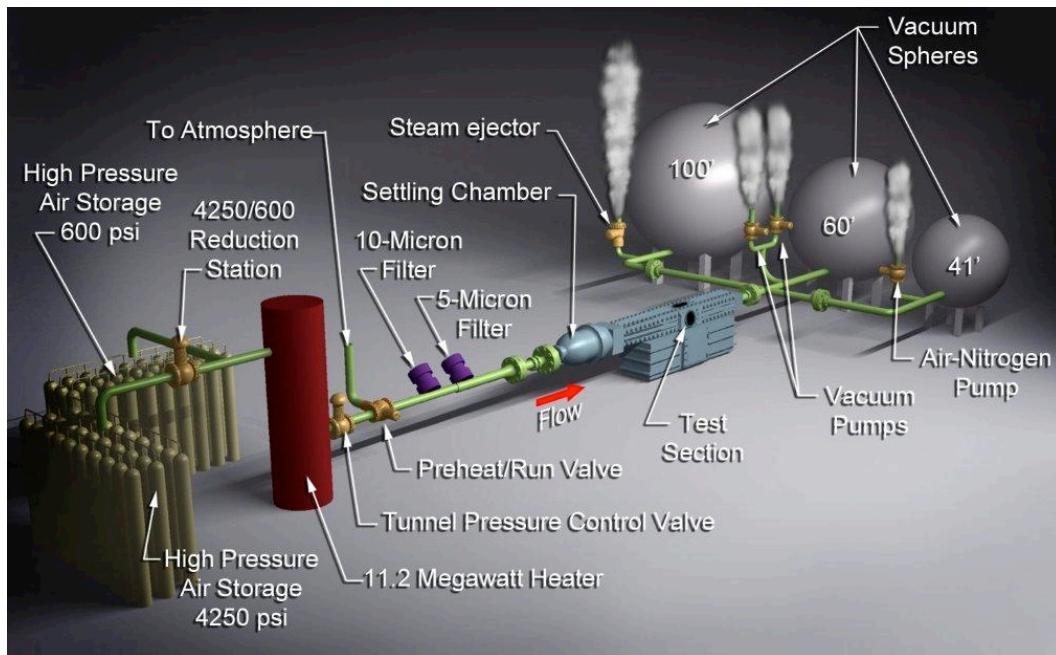


Figure 19. Schematic of Langley Research Center 20-Inch Mach 6 Air Tunnel.



Figure 20. Langley Research Center 20-Inch Mach 6 Air Tunnel test section.

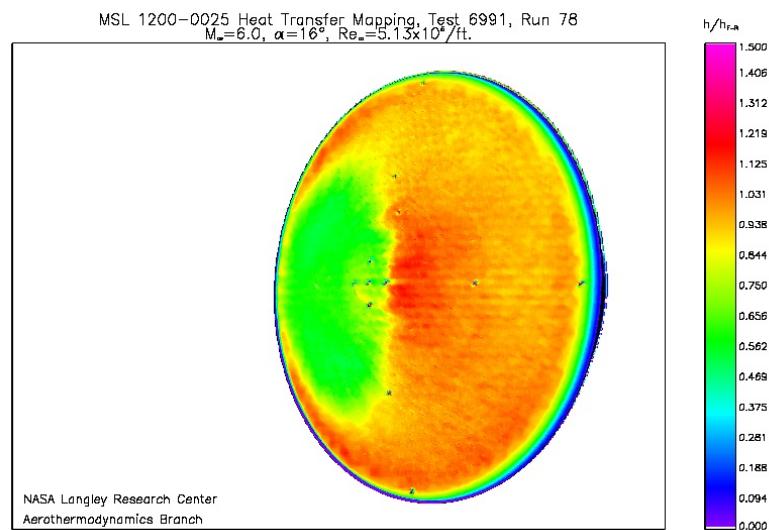


Figure 21. Sample phosphor thermography 2-D image data.

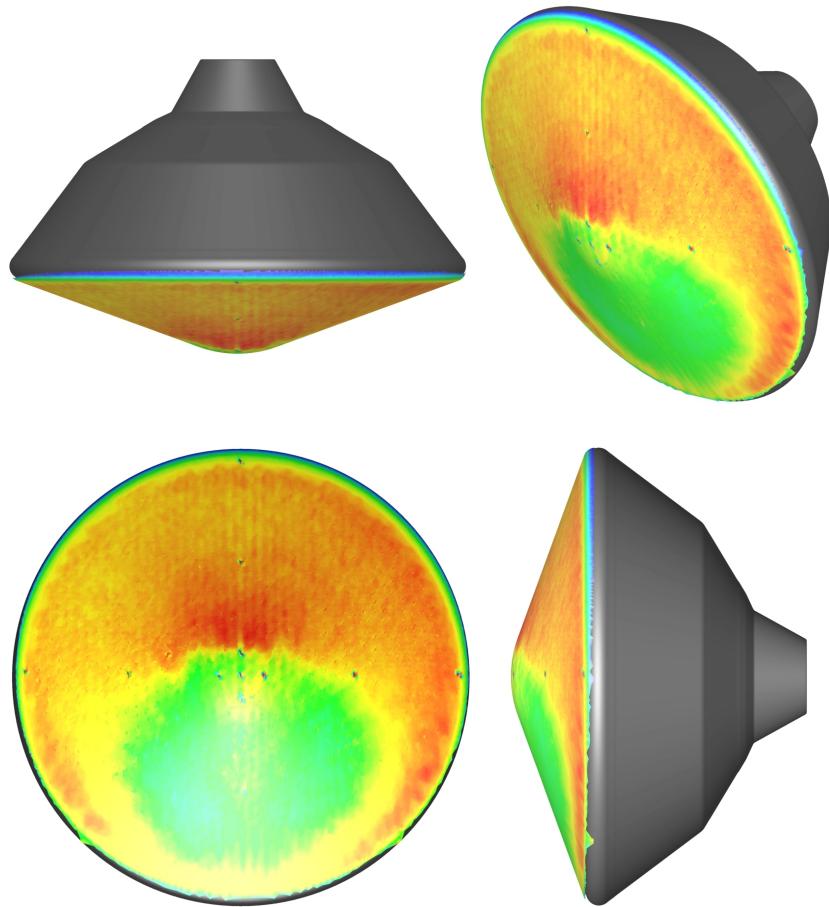


Figure 22. Sample 3-D mapping of phosphor thermography data.

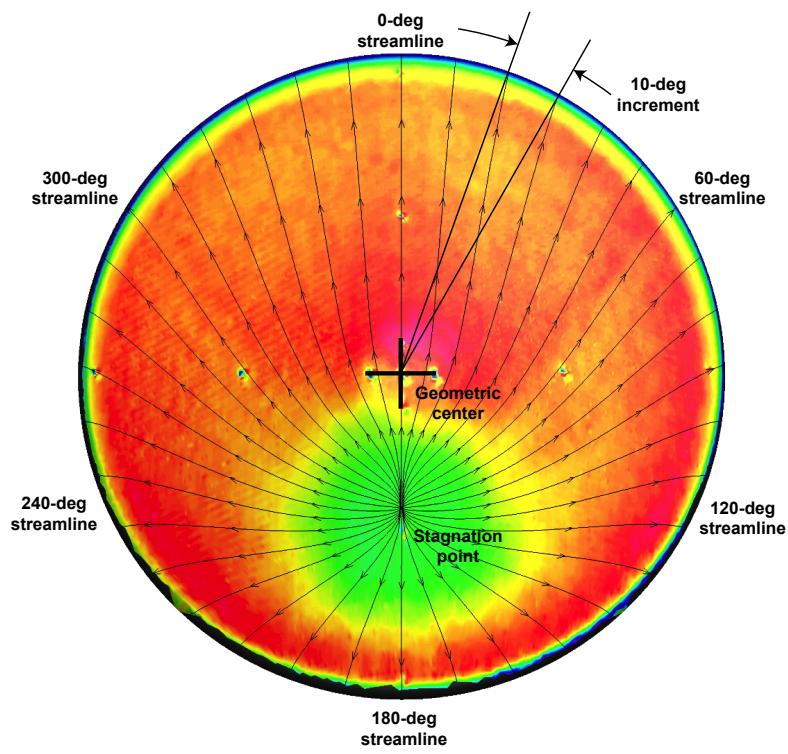


Figure 23. Streamlines for data extraction on sphere-cone geometry.

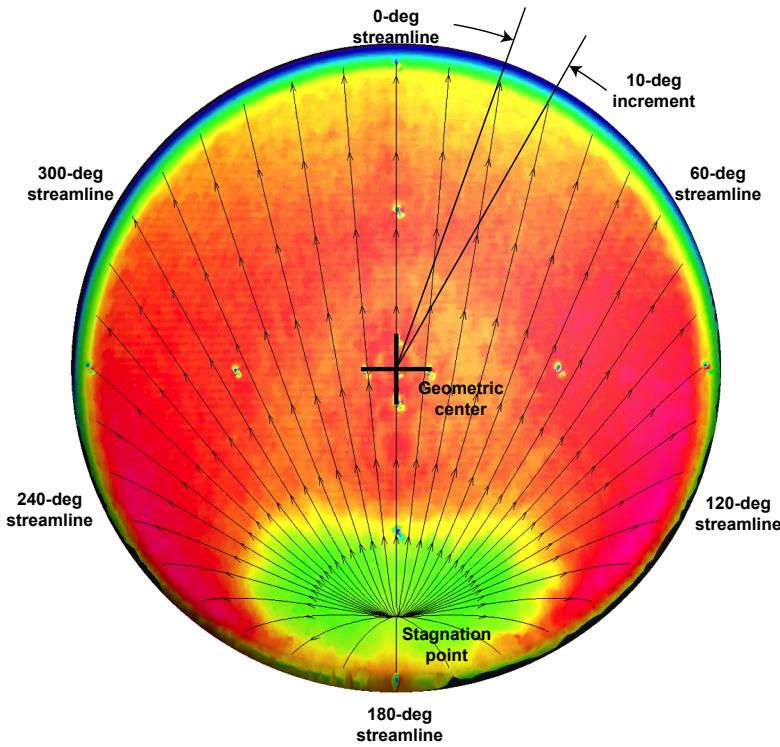


Figure 24. Streamlines for data extraction on spherical-cap geometry.

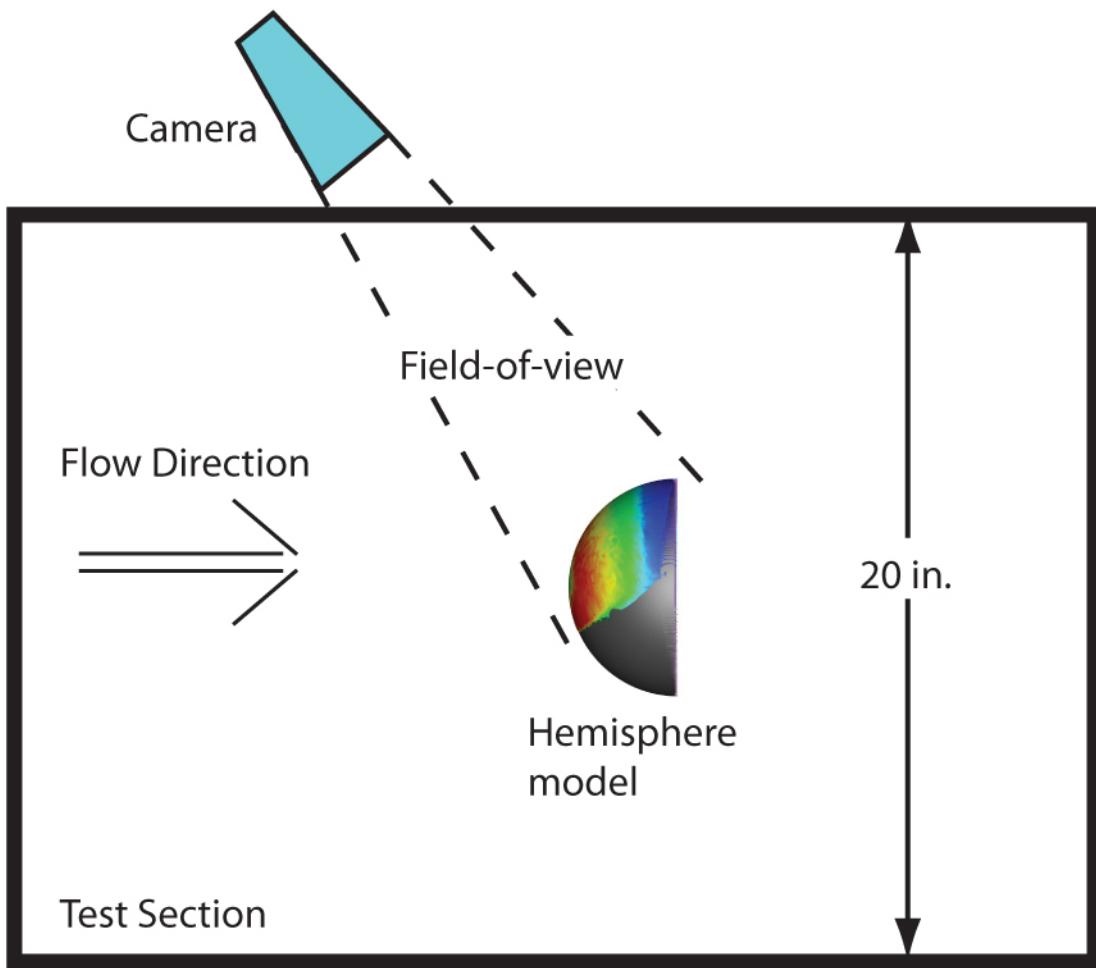


Figure 25. Illustration of camera field-of-view for hemisphere model in 20-Inch Mach 6 Air Tunnel.

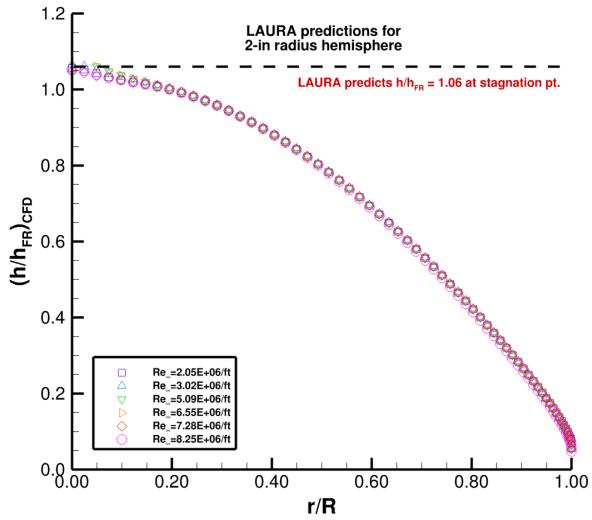


Figure 26. CFD predictions for hemisphere heating at wind tunnel conditions.

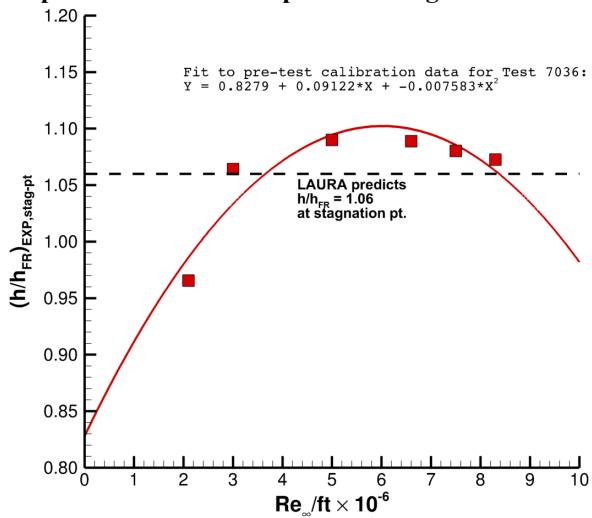


Figure 27. Measured stagnation point heating for pretest calibrations for Test 7036.

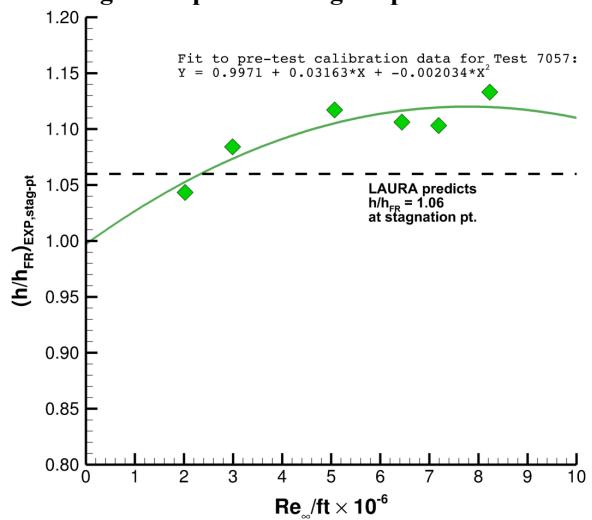


Figure 28. Measured stagnation point heating for pretest calibrations for Test 7057.

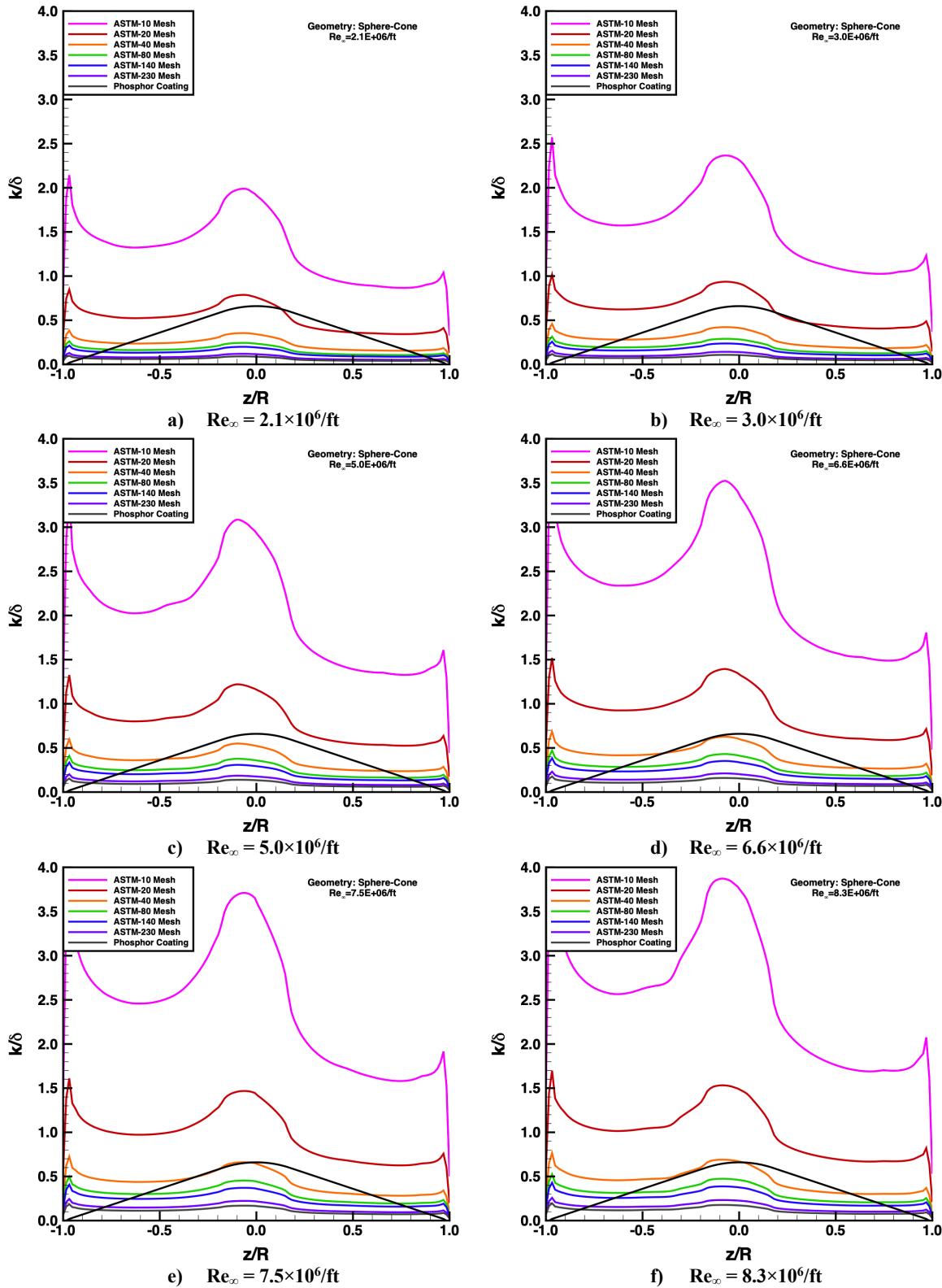


Figure 29. Centerline profiles of roughness effects on k/δ , sphere-cone geometry.

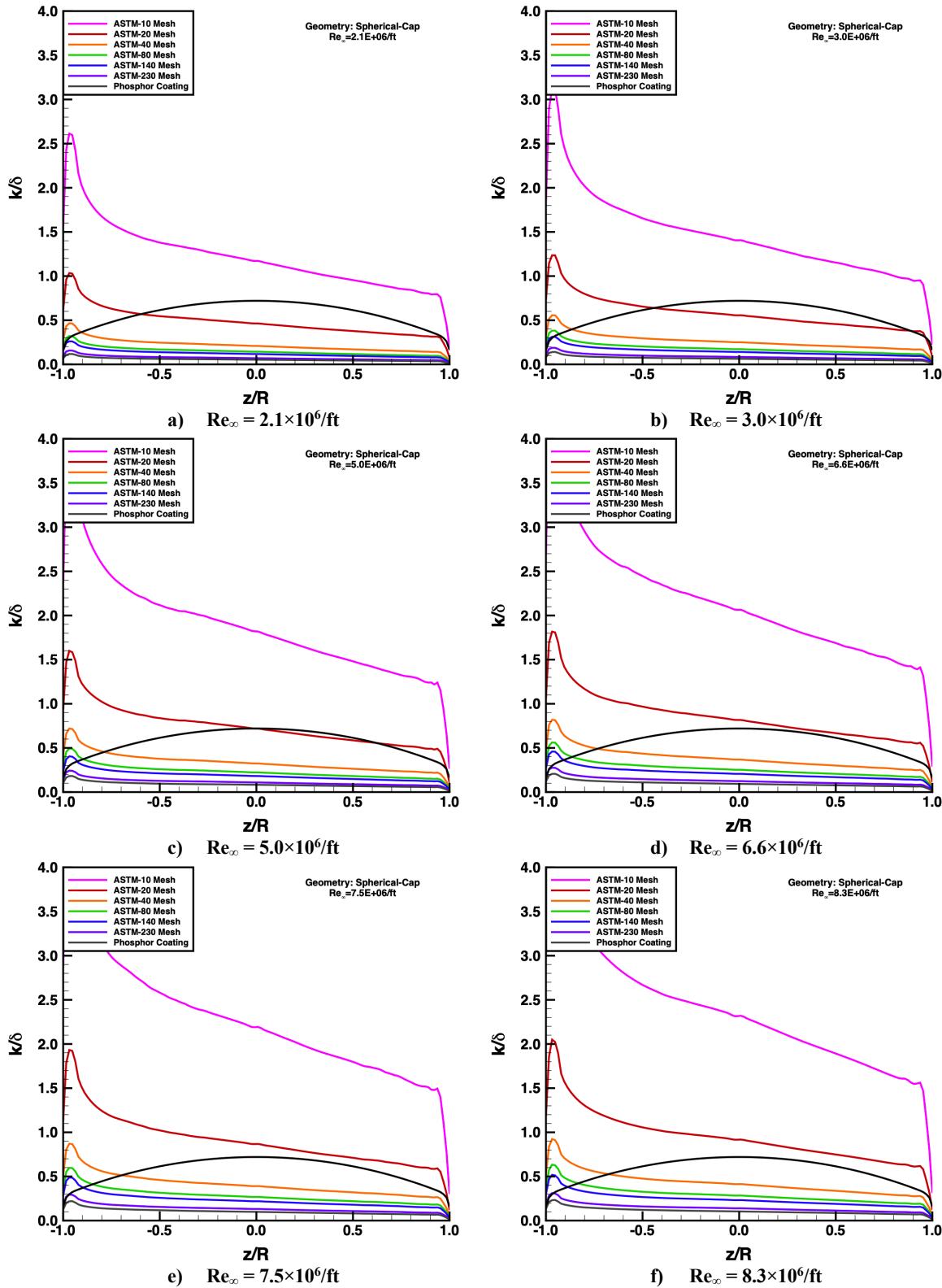


Figure 30. Centerline profiles of roughness effects on k/δ , spherical-cap geometry.

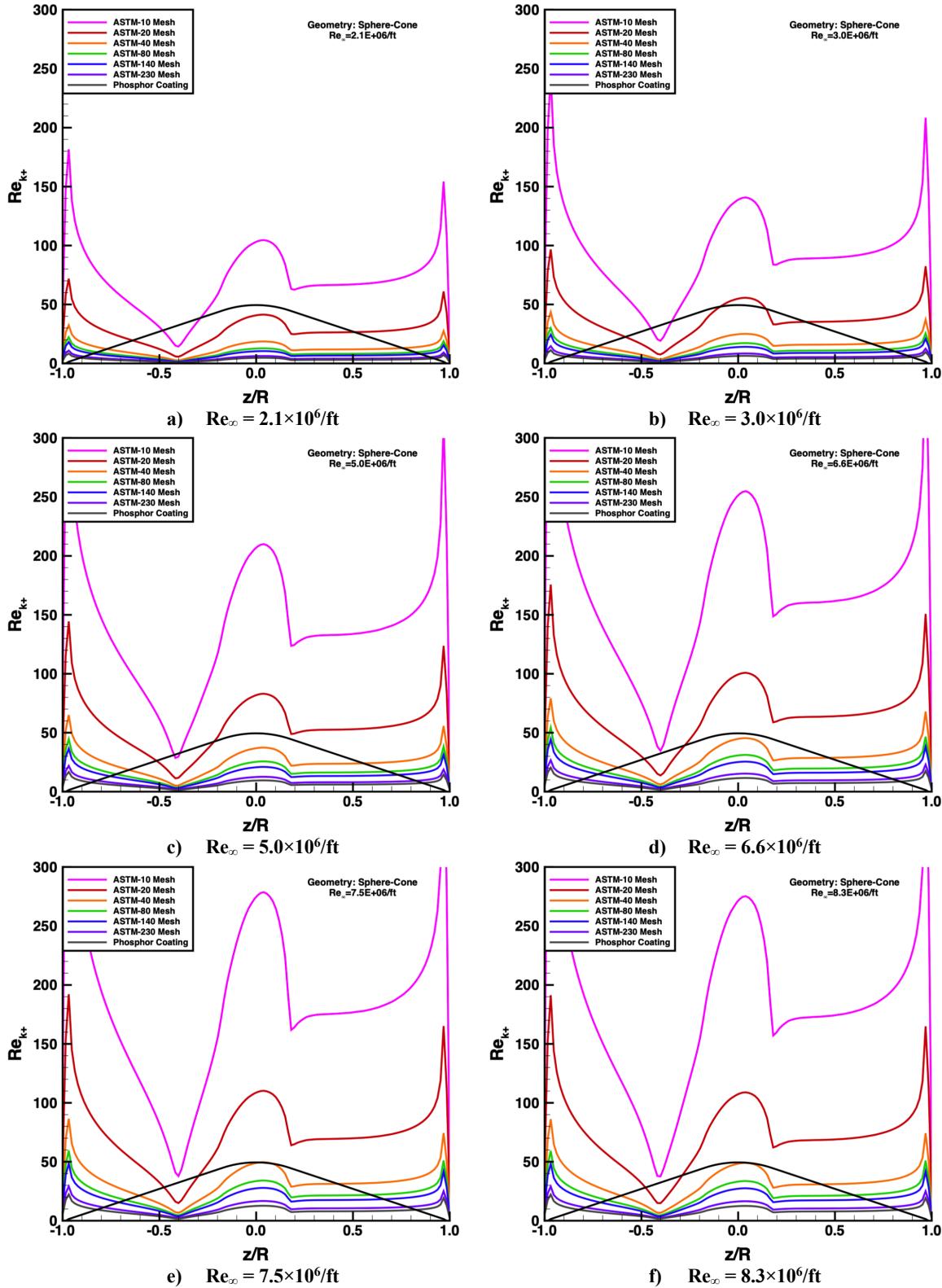


Figure 31. Centerline profiles of roughness effects on Re_{k+} , sphere-cone geometry.

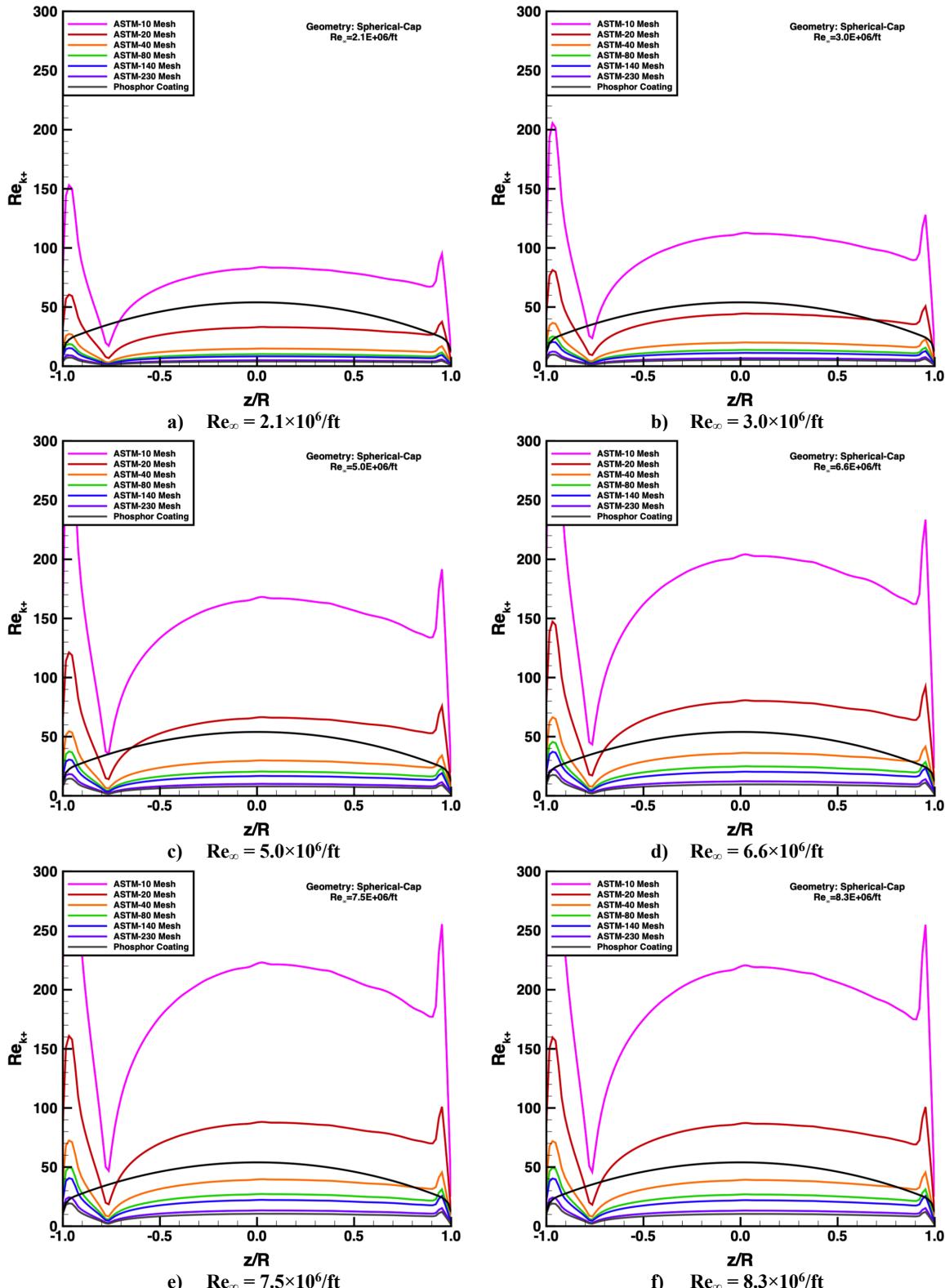


Figure 32. Centerline profiles of roughness effects on Re_{k+} , spherical-cap geometry.

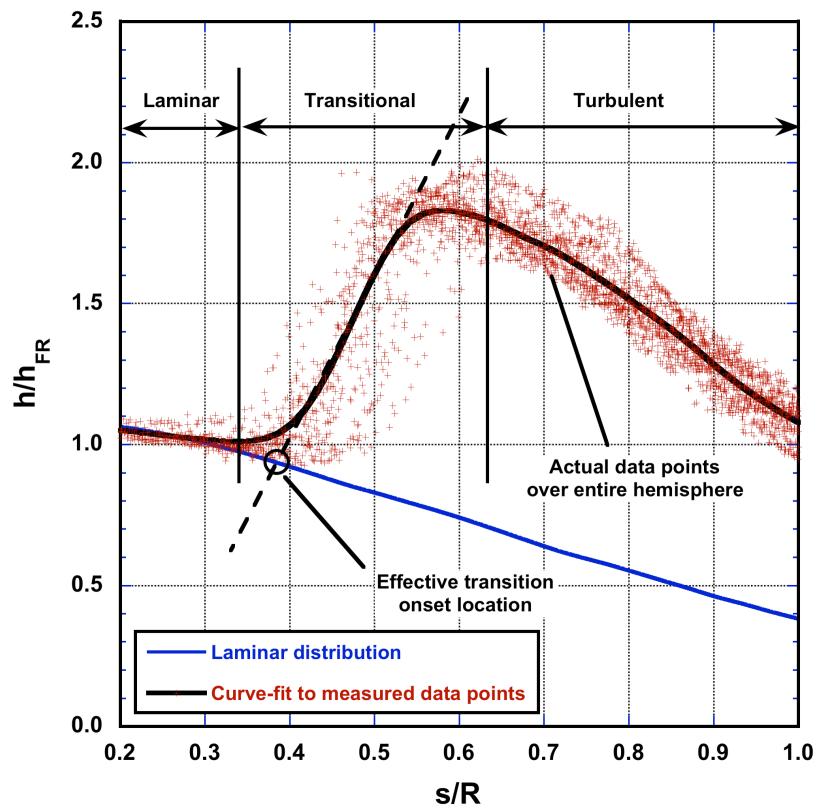


Figure 33. Tangent-slope-intercept method for determination of effect transition onset location.

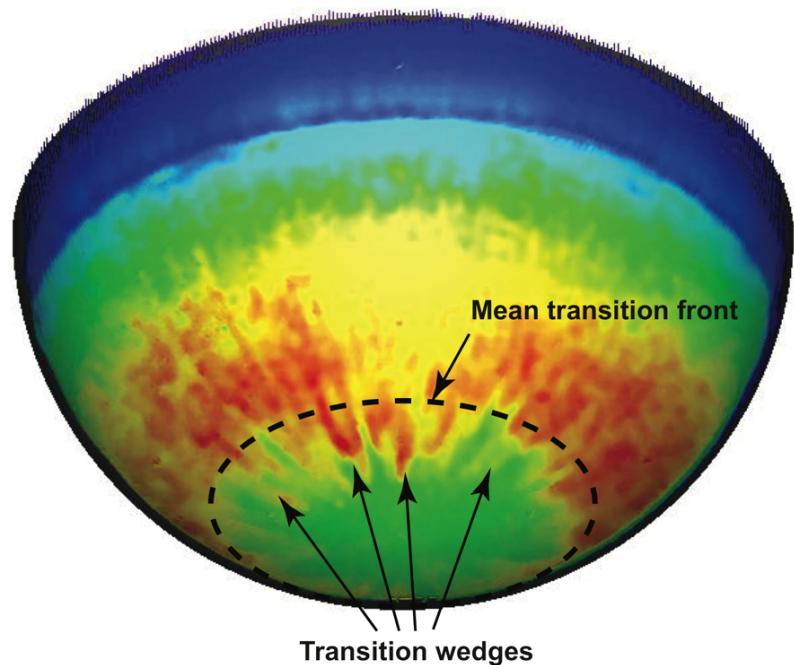


Figure 34. Comparison of irregular transition wedges vs. mean transition front.

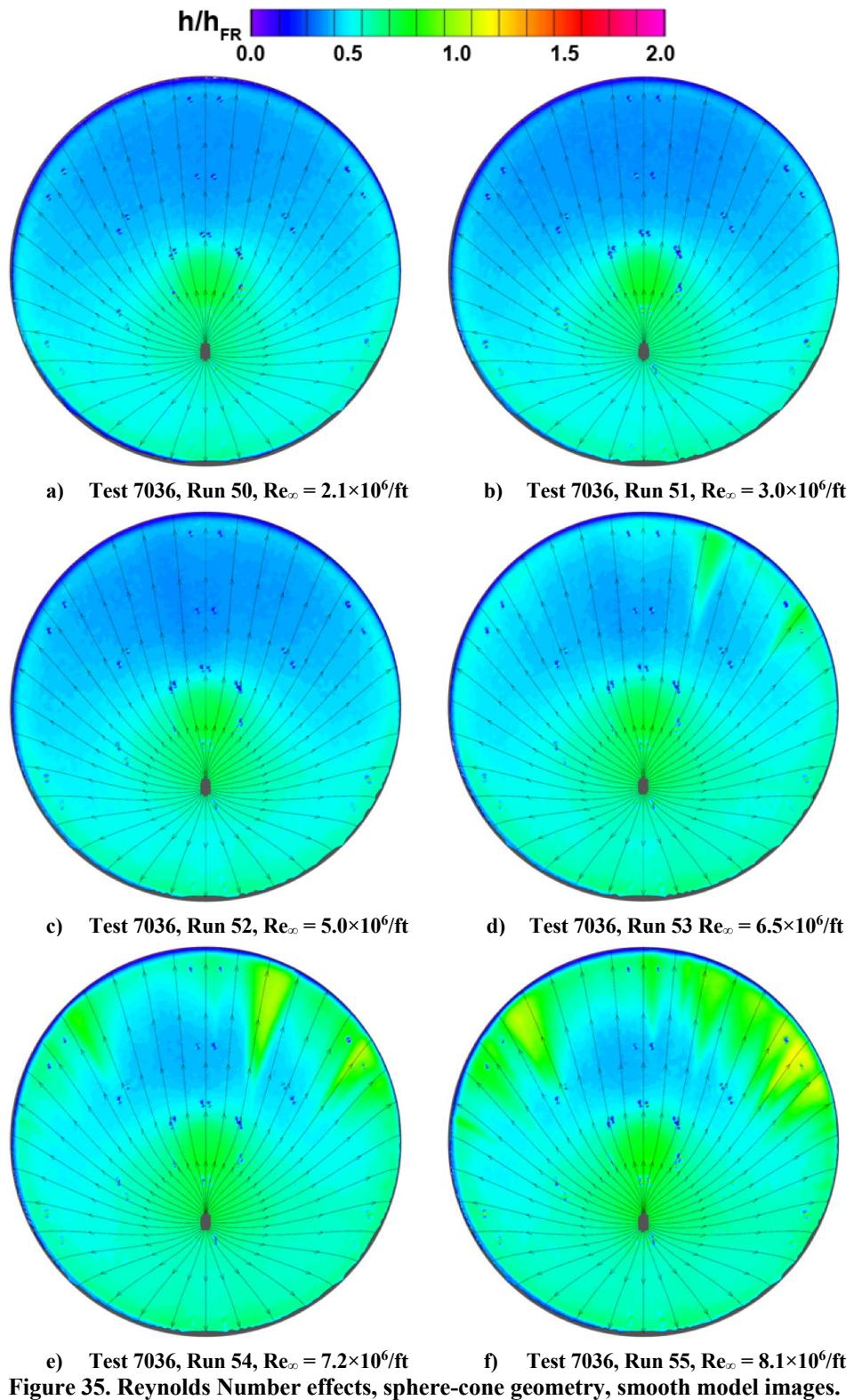


Figure 35. Reynolds Number effects, sphere-cone geometry, smooth model images.

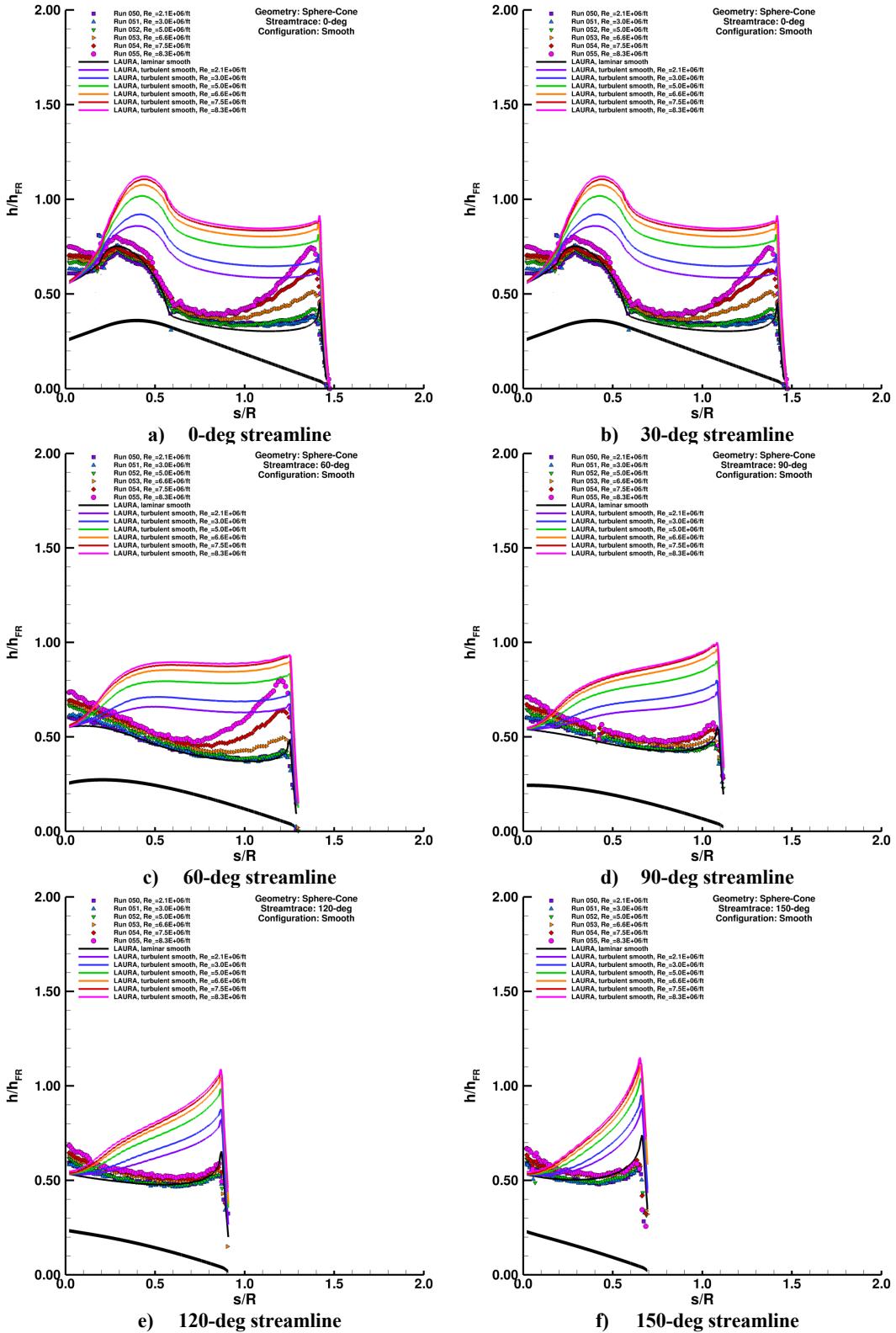


Figure 36. Reynolds number effects, sphere-cone geometry, smooth model plots.

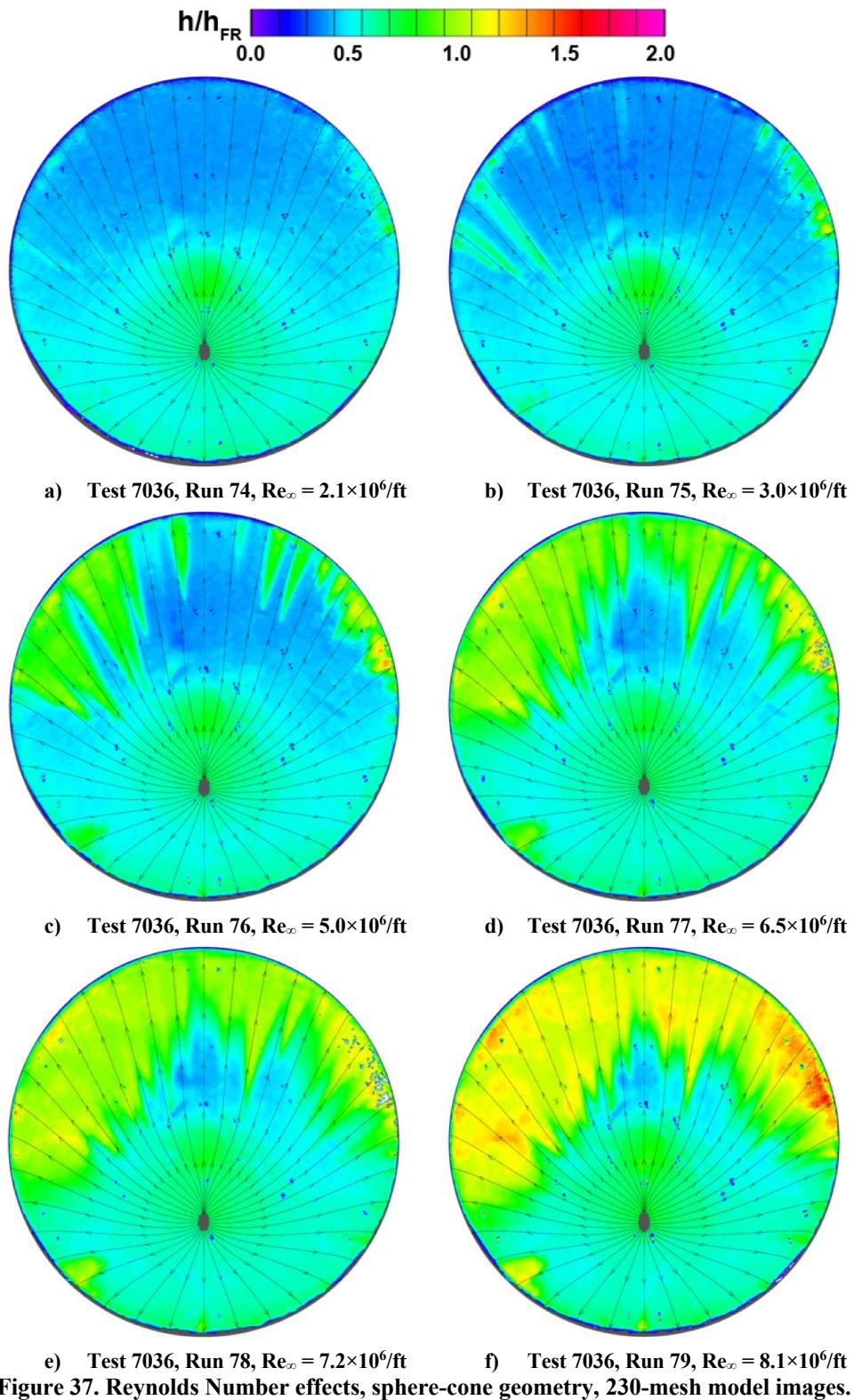


Figure 37. Reynolds Number effects, sphere-cone geometry, 230-mesh model images.

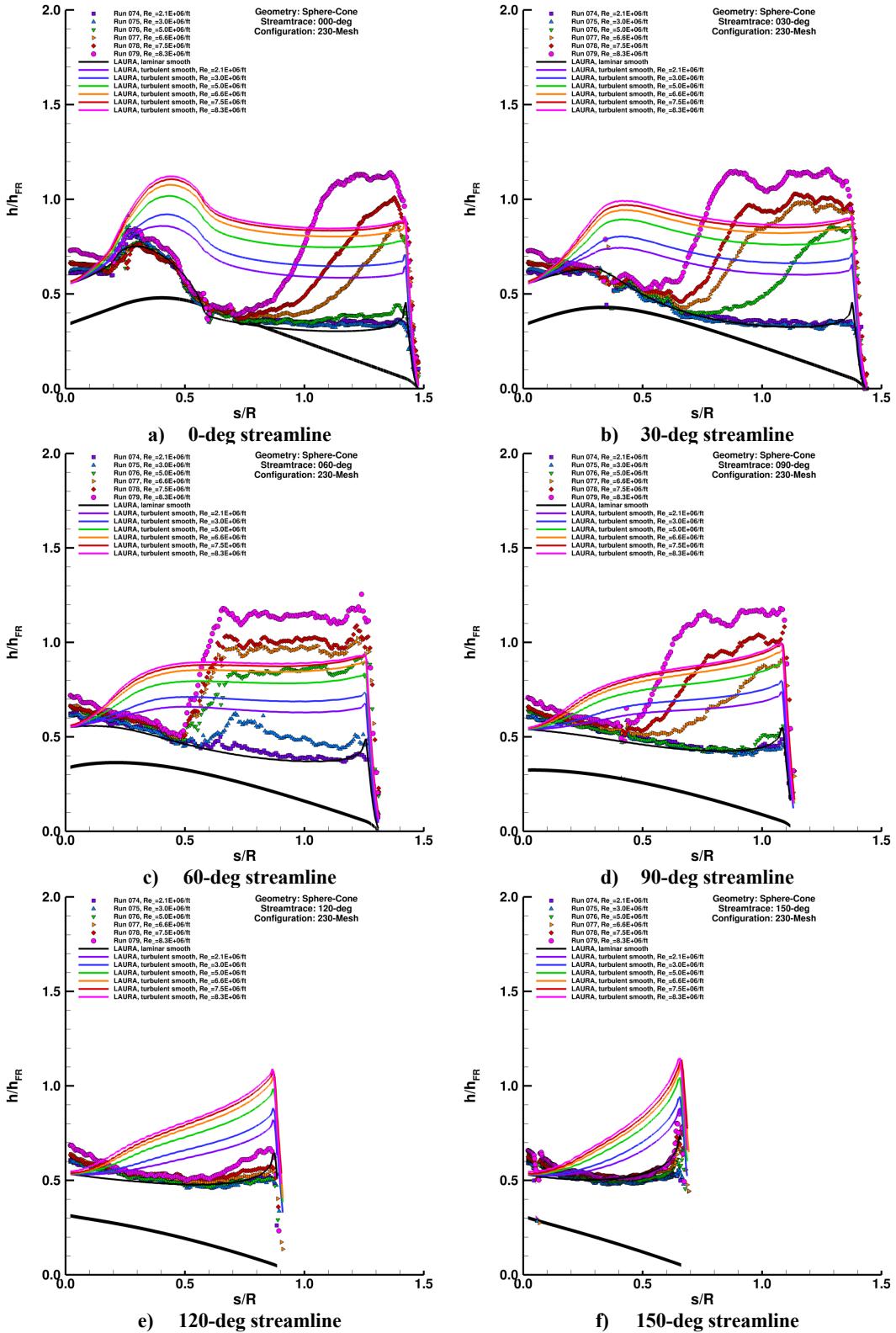


Figure 38. Reynolds Number effects, sphere-cone geometry, 230-mesh model plots.

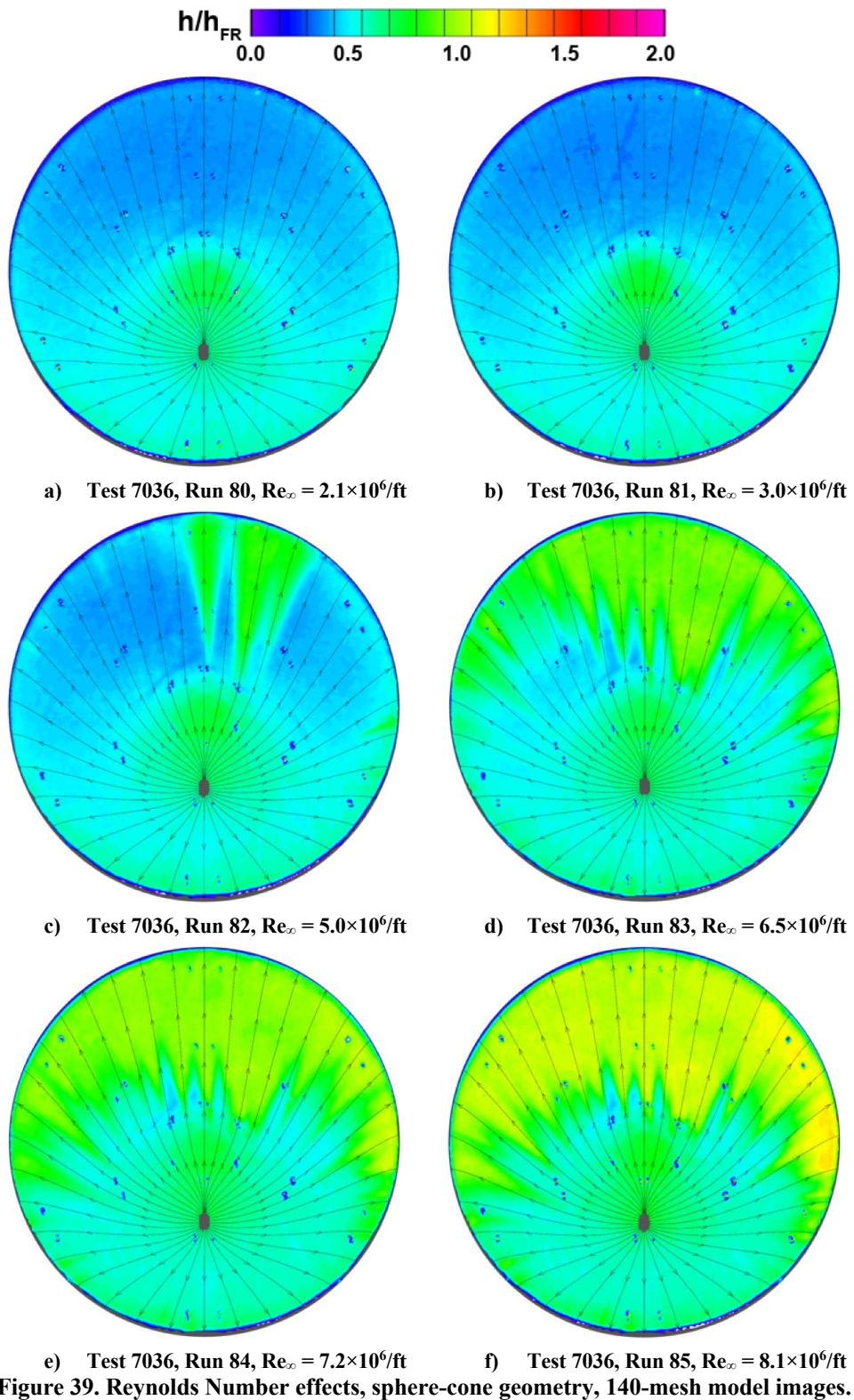


Figure 39. Reynolds Number effects, sphere-cone geometry, 140-mesh model images.

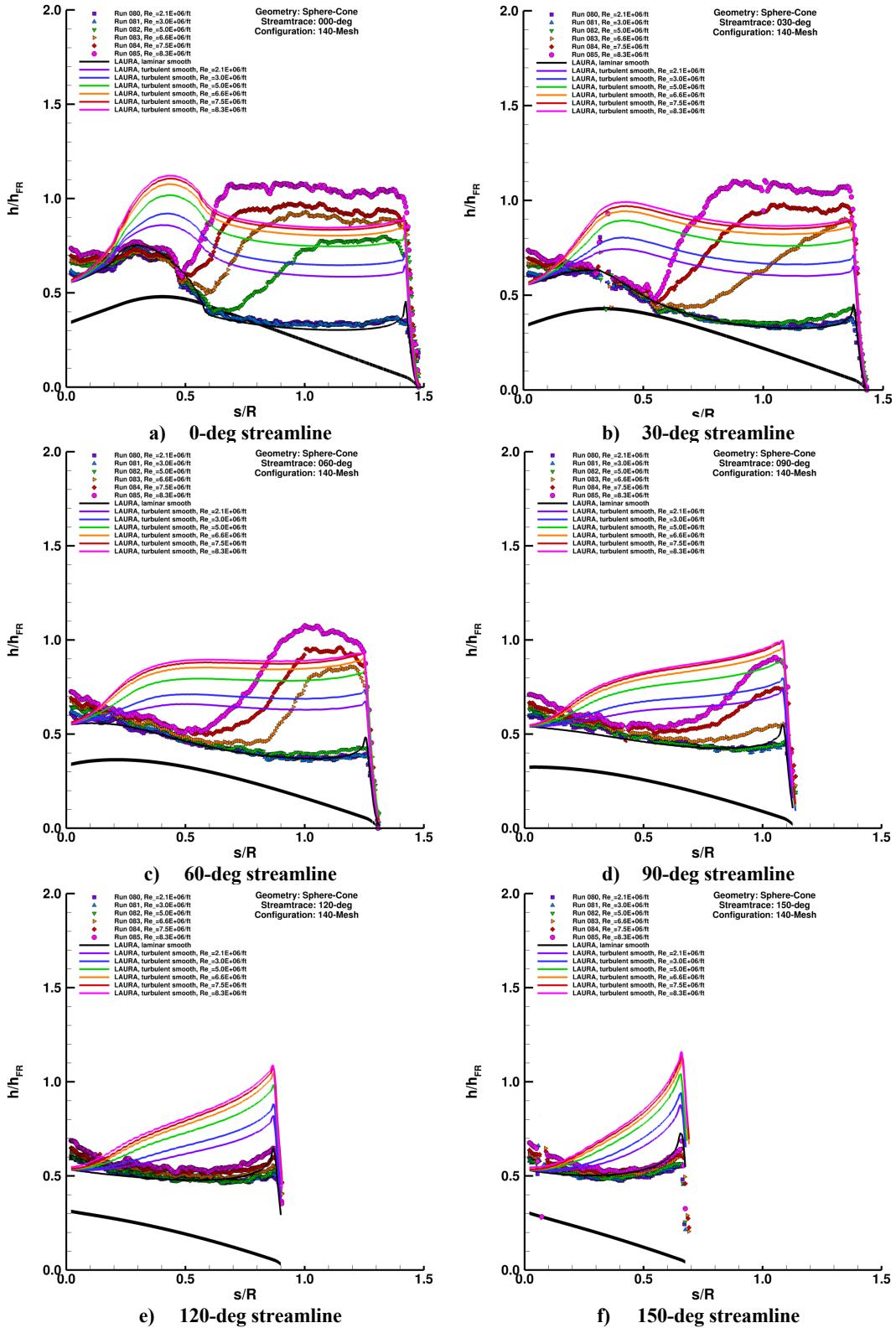


Figure 40. Reynolds Number effects, sphere-cone geometry, 140-mesh model plots.

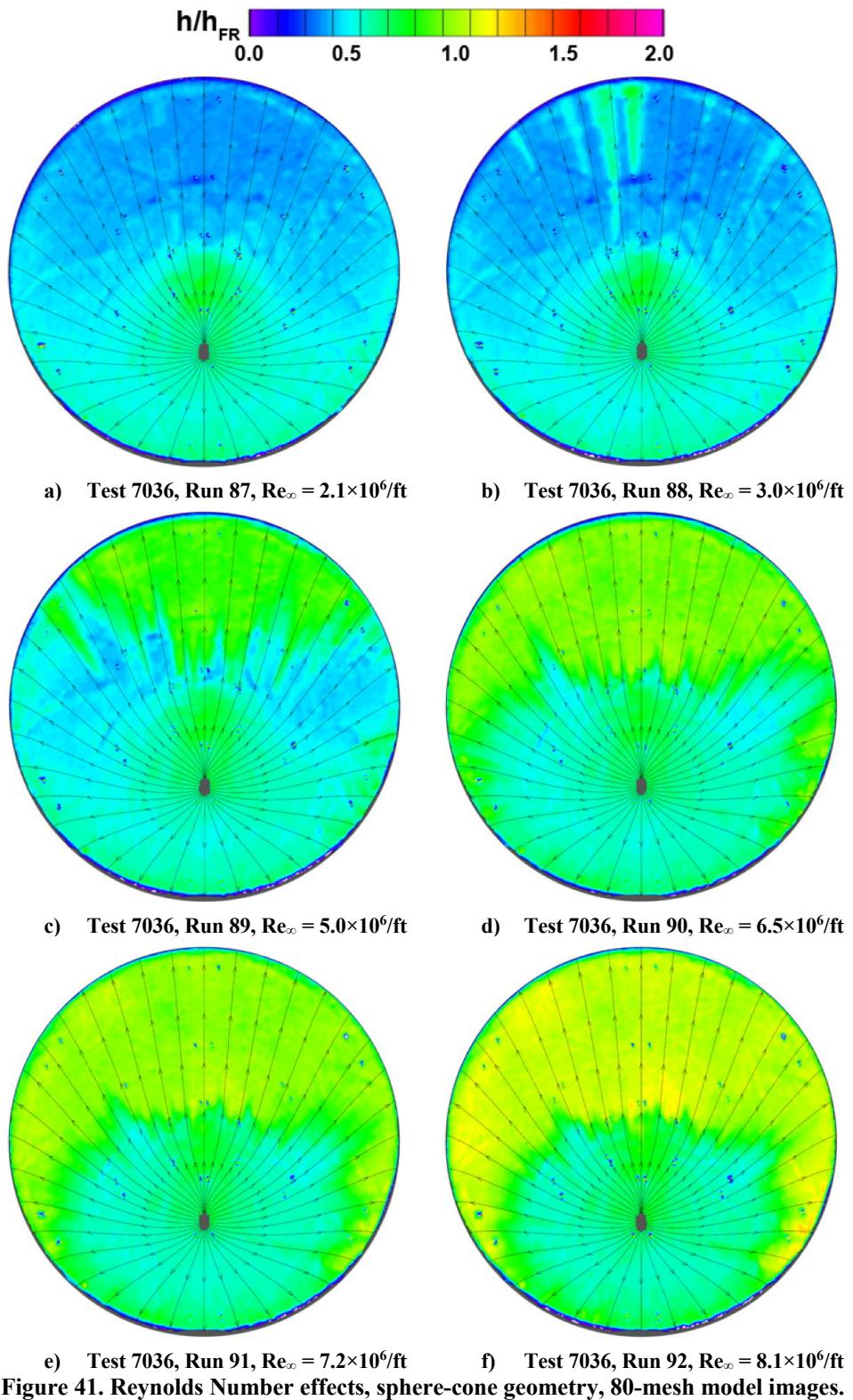


Figure 41. Reynolds Number effects, sphere-cone geometry, 80-mesh model images.

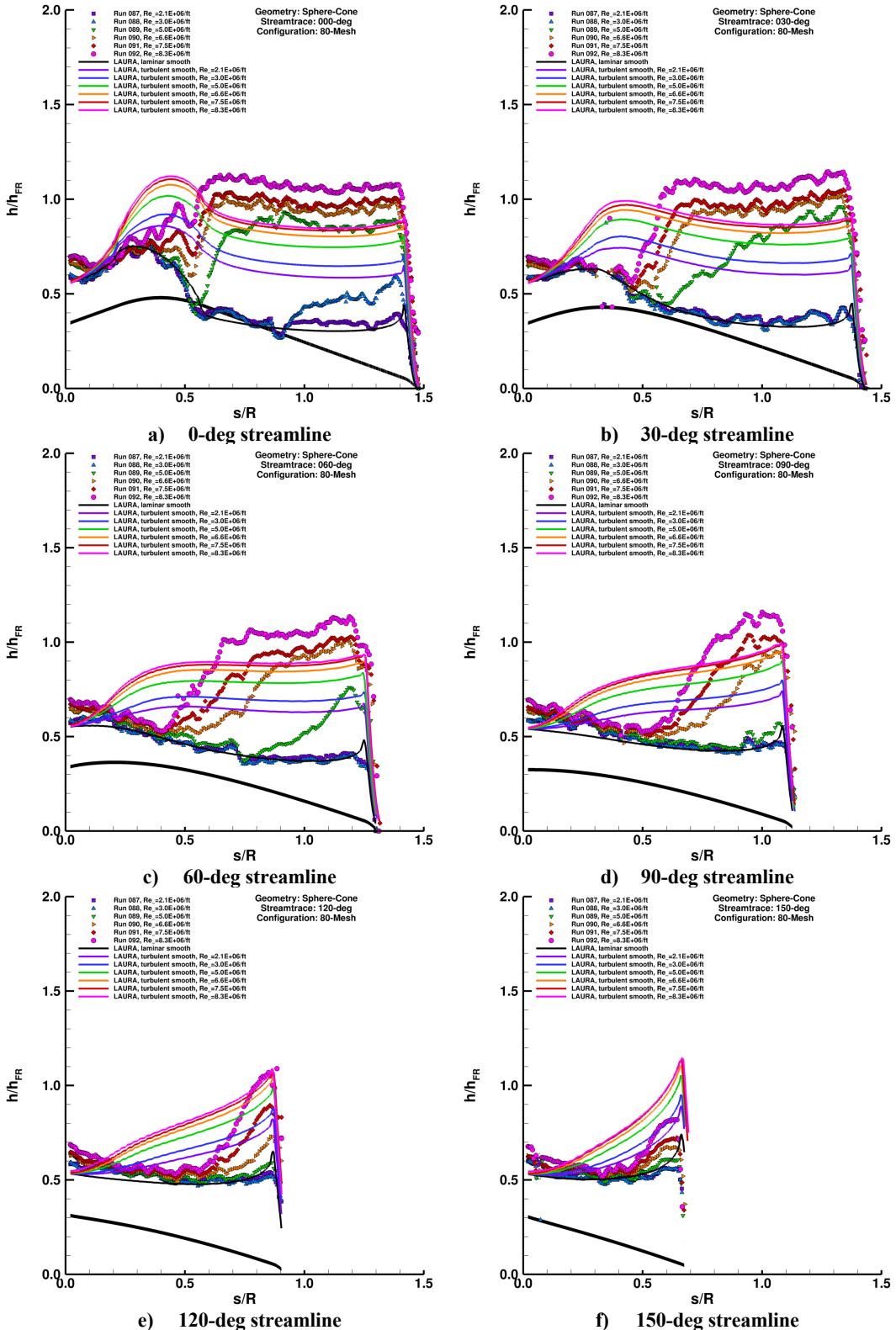


Figure 42. Reynolds Number effects, sphere-cone geometry, 80-mesh model plots.

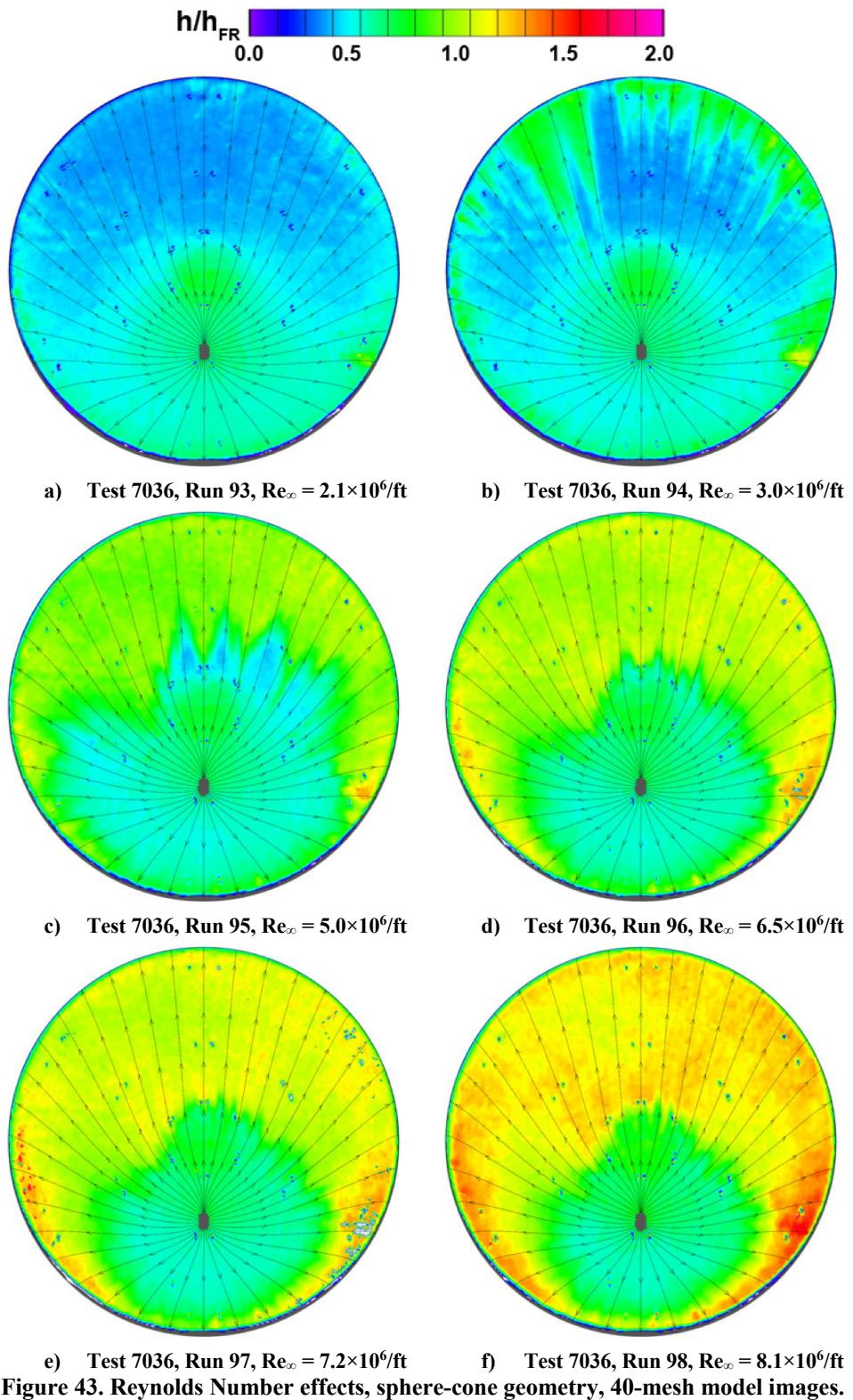


Figure 43. Reynolds Number effects, sphere-cone geometry, 40-mesh model images.

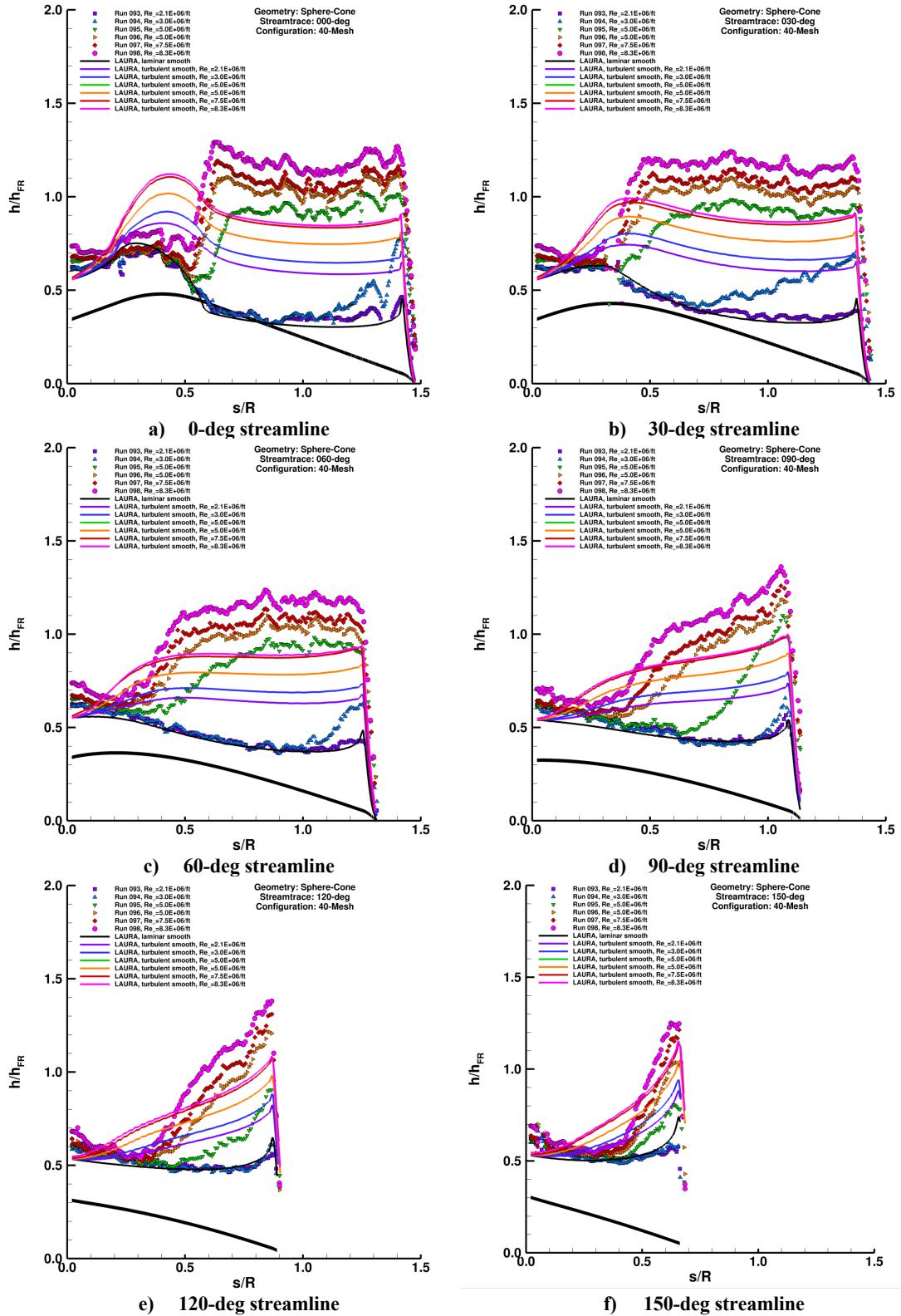


Figure 44. Reynolds Number effects, sphere-cone geometry, 40-mesh model plots.

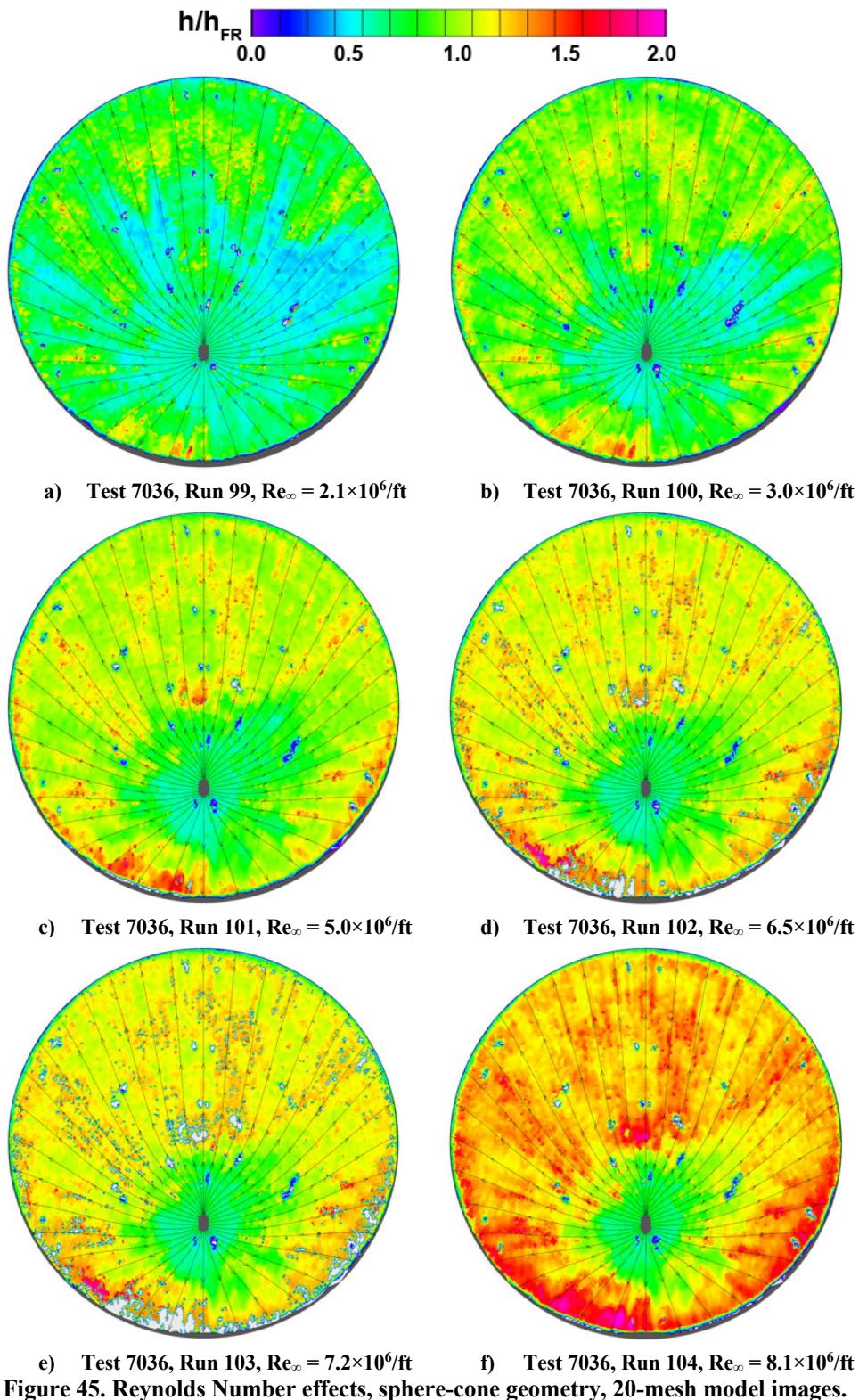


Figure 45. Reynolds Number effects, sphere-cone geometry, 20-mesh model images.

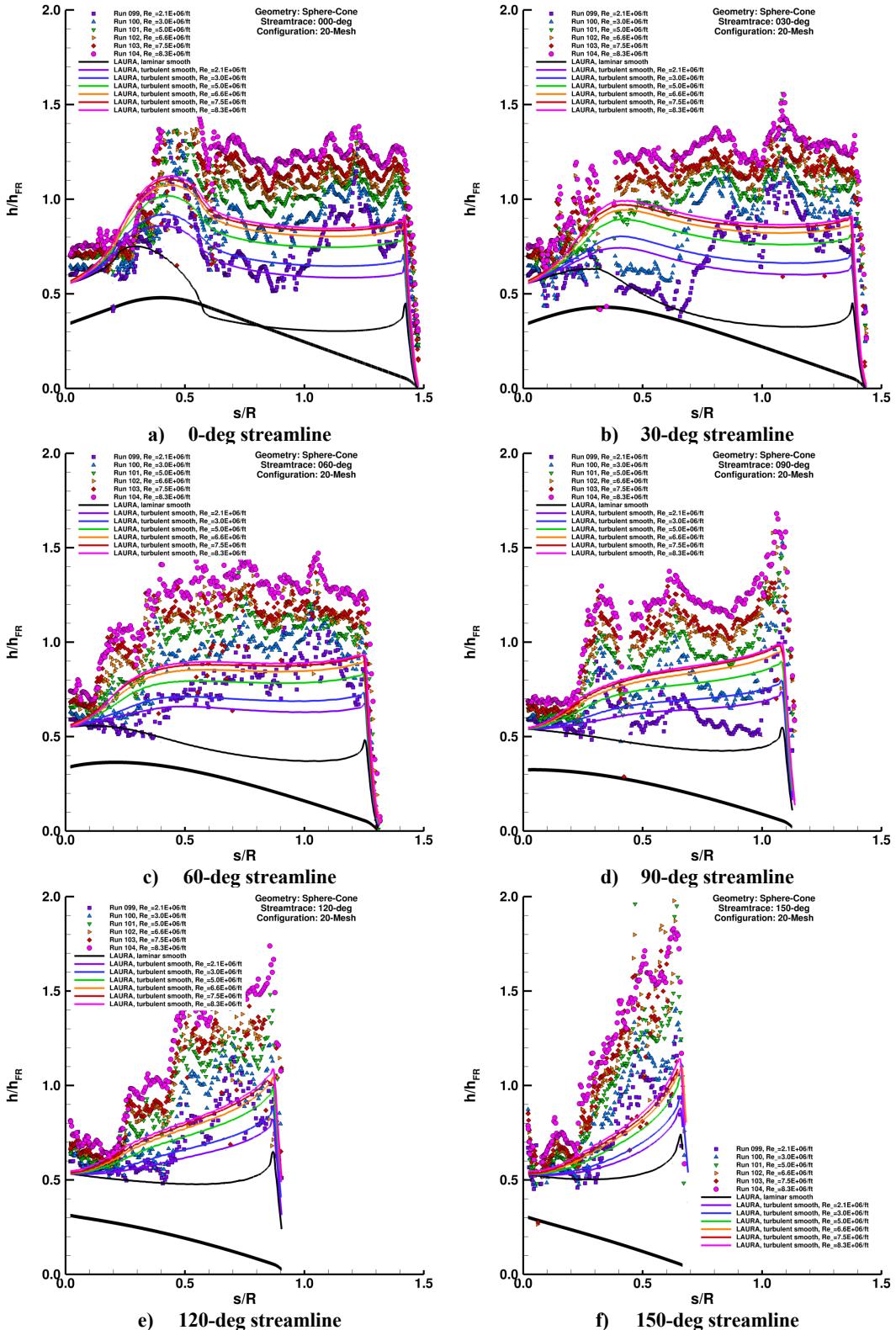


Figure 46. Reynolds Number effects, sphere-cone geometry, 20-mesh model plots.

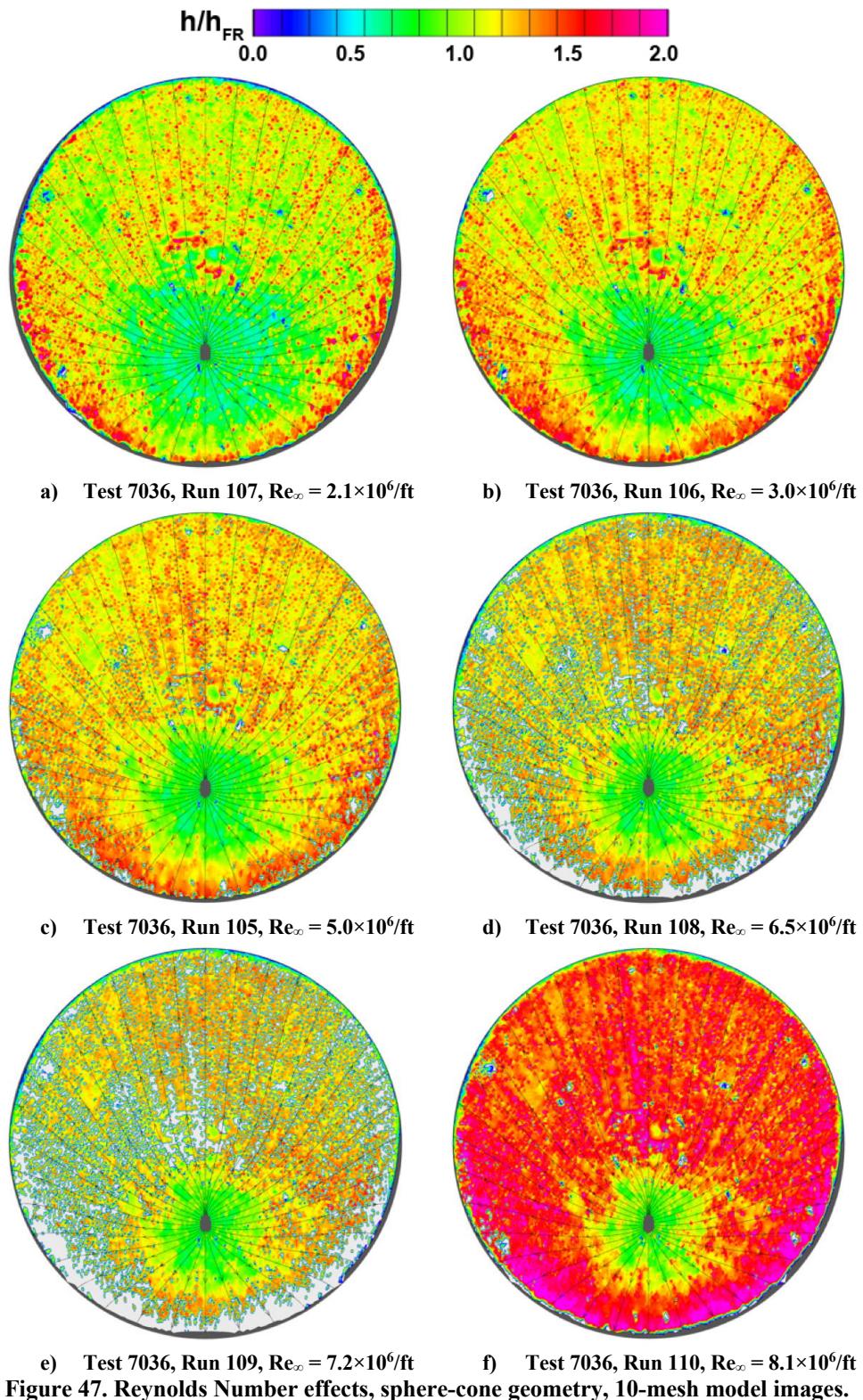


Figure 47. Reynolds Number effects, sphere-cone geometry, 10-mesh model images.

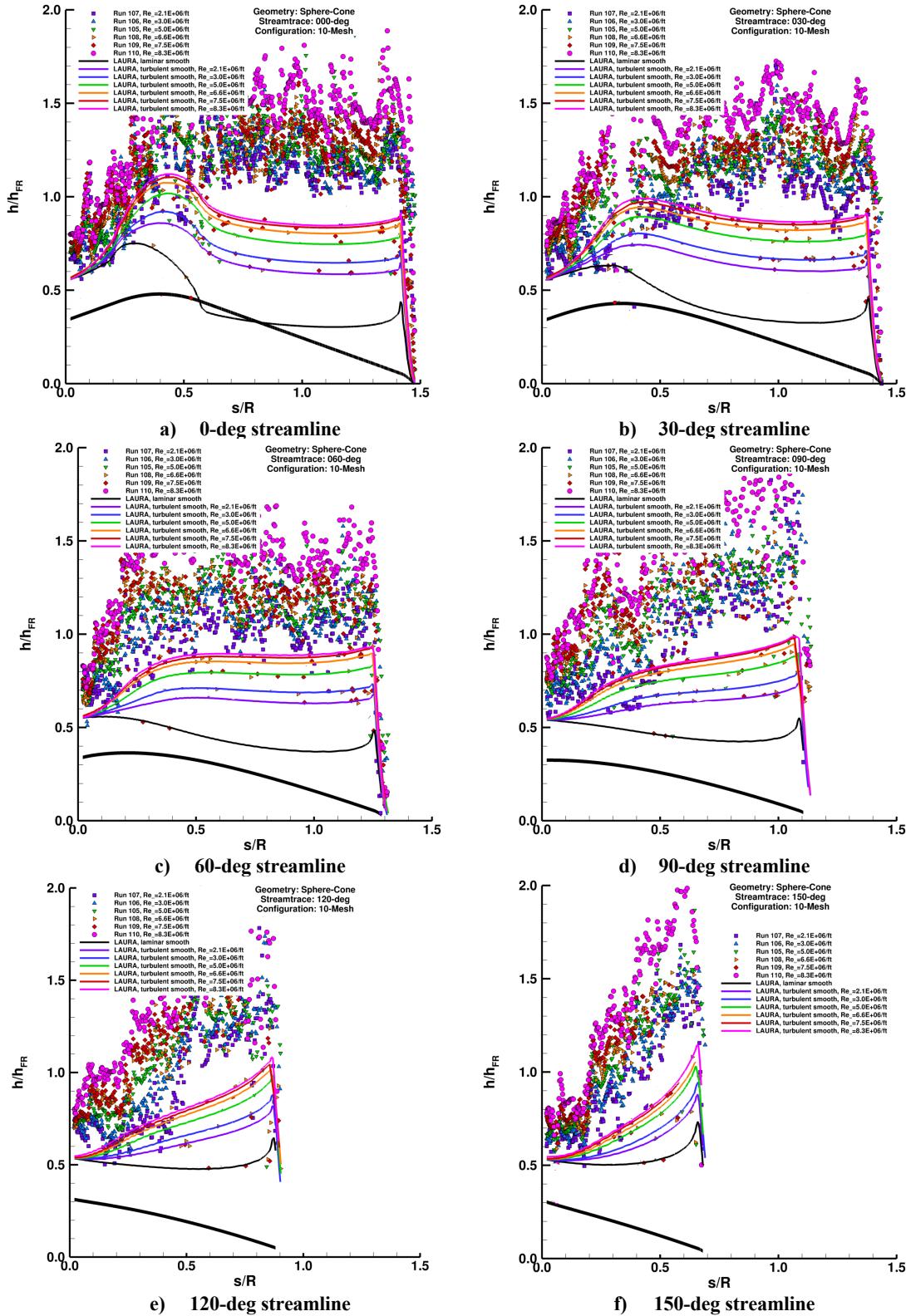


Figure 48. Reynolds Number effects, sphere-cone geometry, 10-mesh model plots.

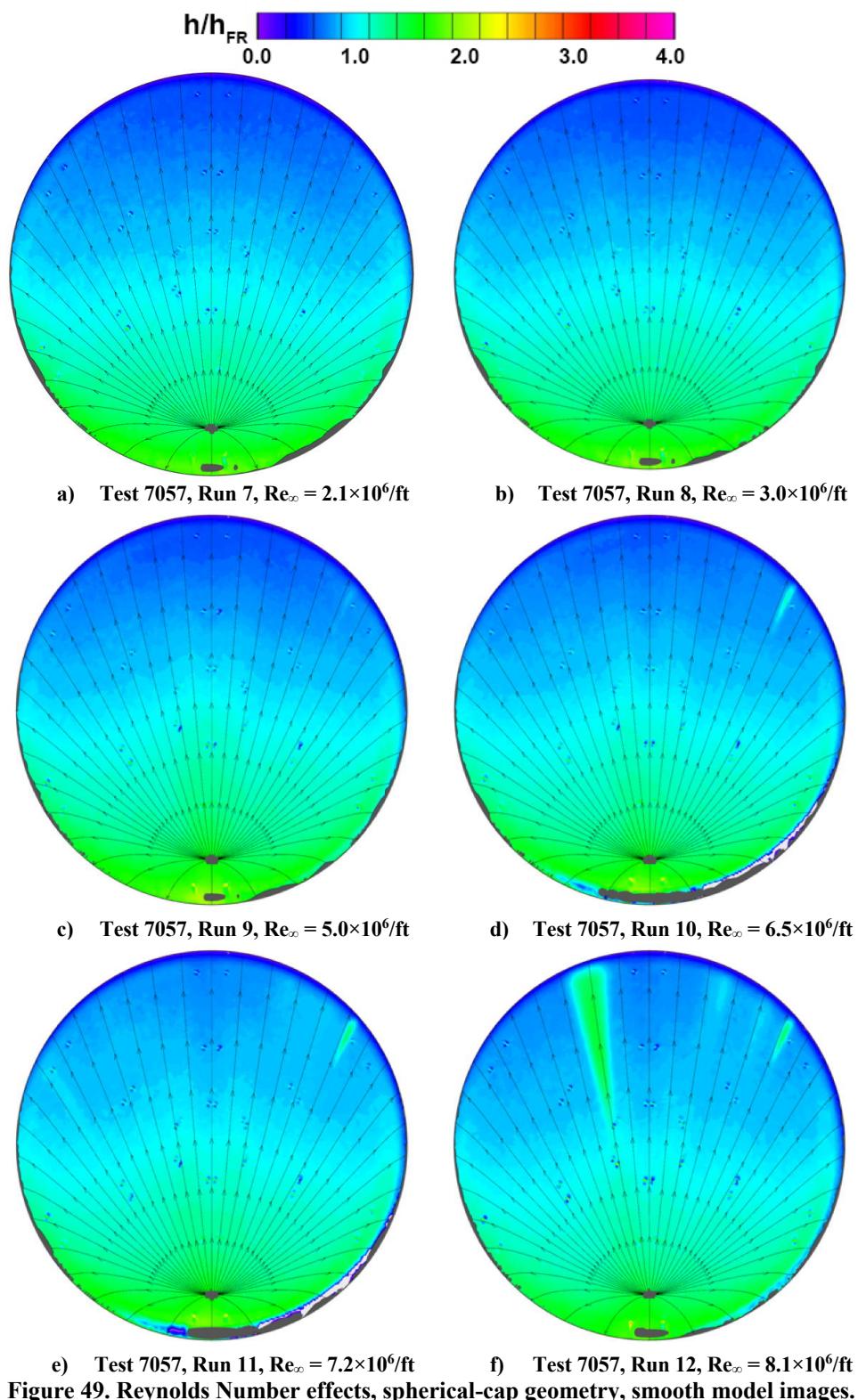


Figure 49. Reynolds Number effects, spherical-cap geometry, smooth model images.

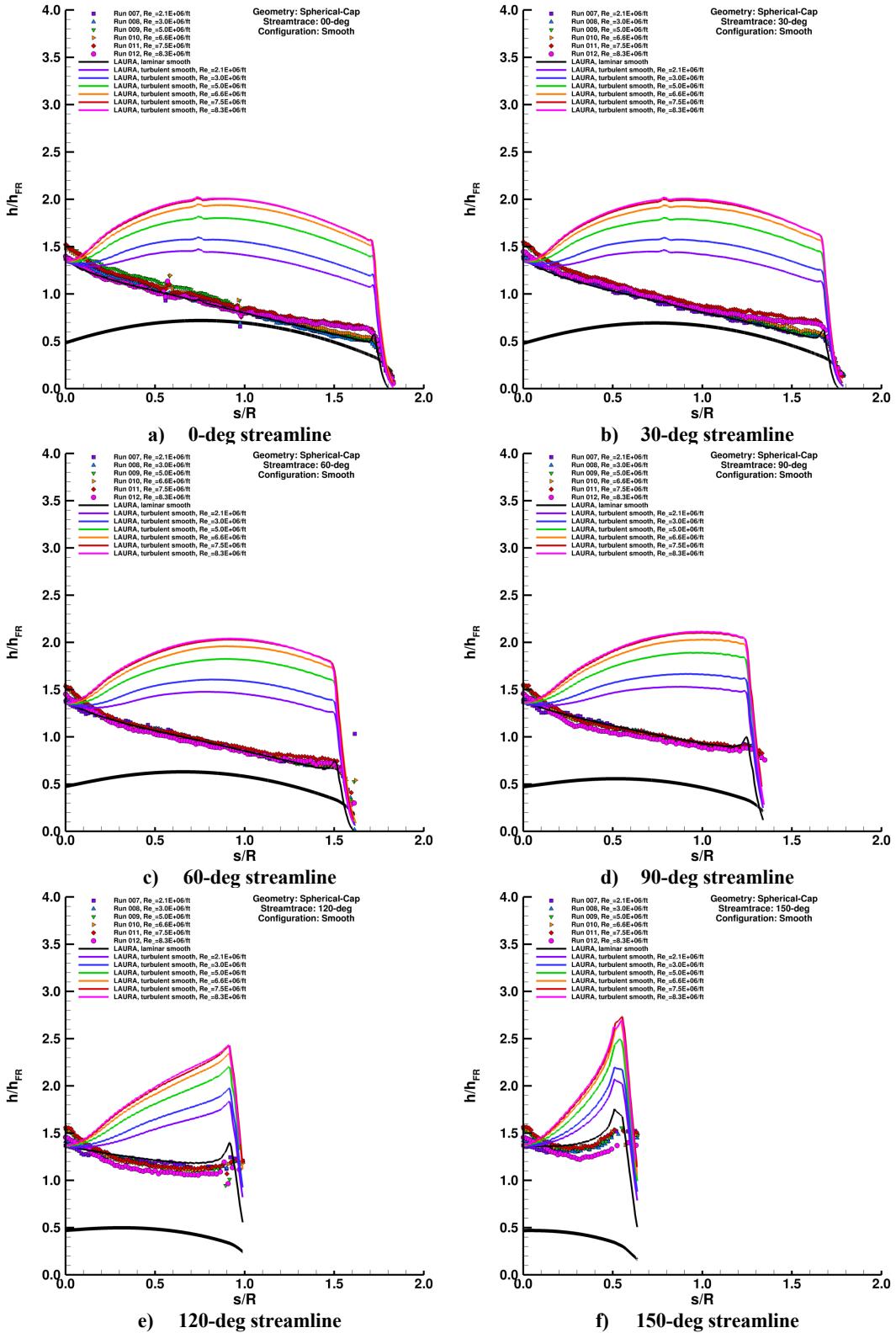


Figure 50. Reynolds number effects, spherical-cap geometry, smooth model plots.

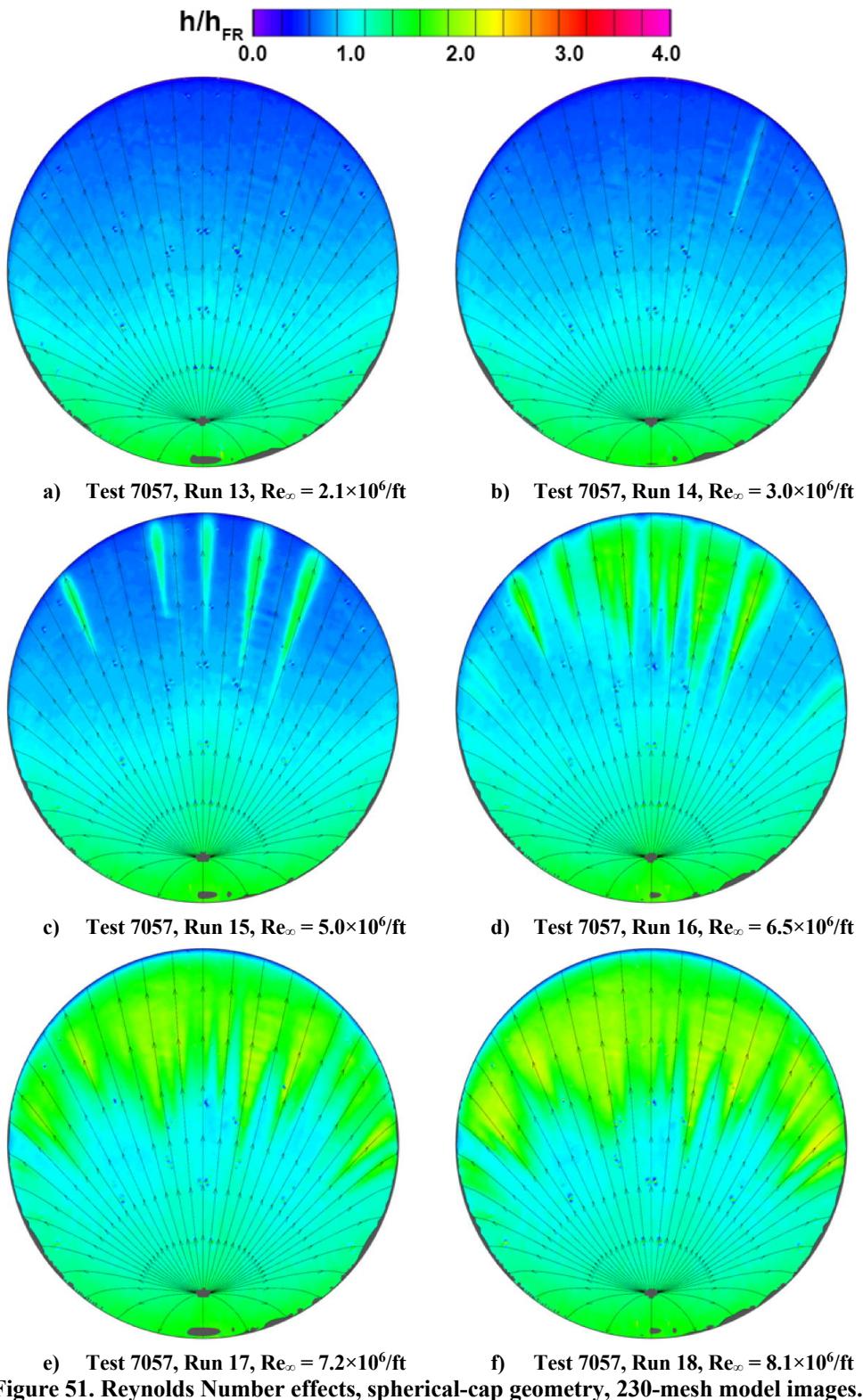


Figure 51. Reynolds Number effects, spherical-cap geometry, 230-mesh model images.

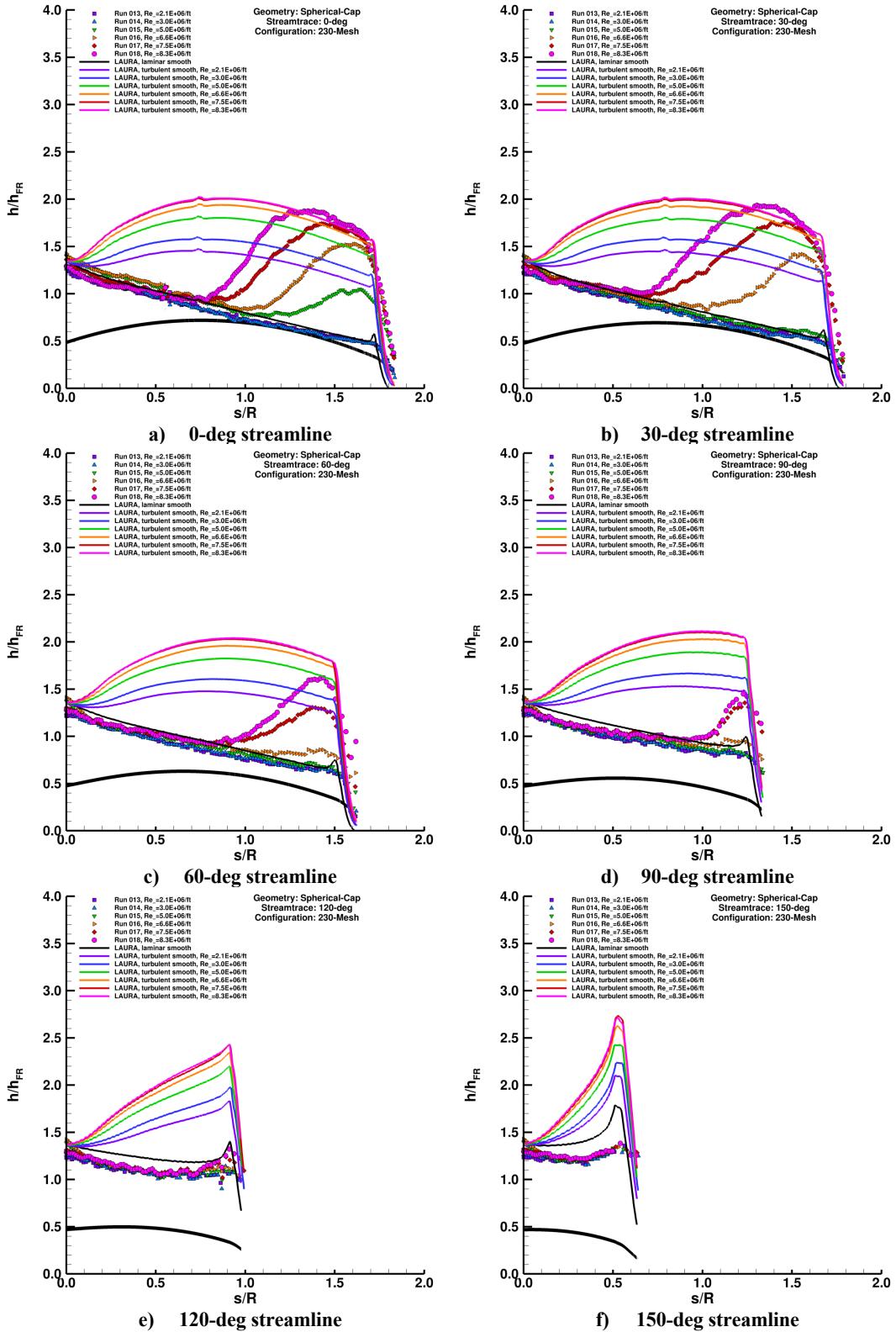
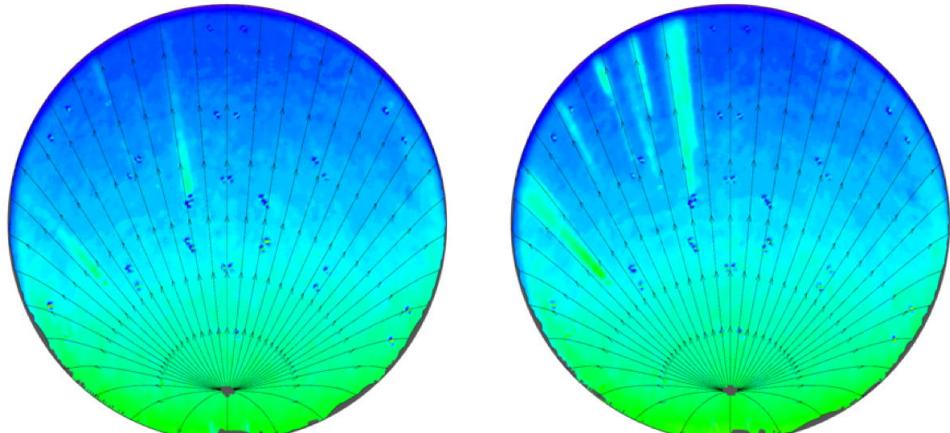
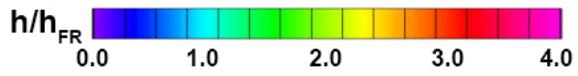
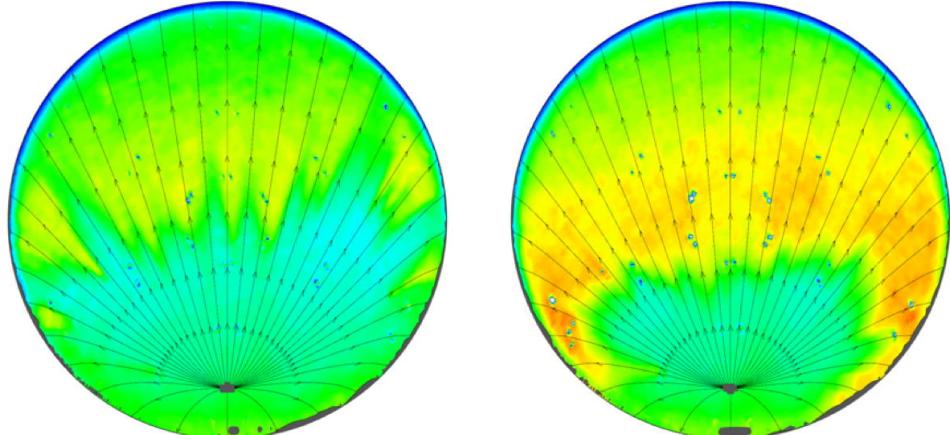


Figure 52. Reynolds Number effects, spherical-cap geometry, 230-mesh model plots.



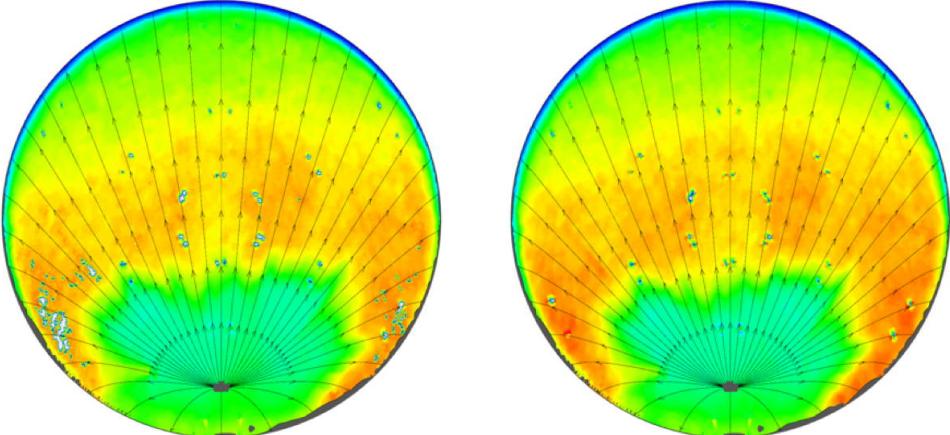
a) Test 7057, Run 19, $Re_\infty = 2.1 \times 10^6/\text{ft}$

b) Test 7057, Run 20, $Re_\infty = 3.0 \times 10^6/\text{ft}$



c) Test 7057, Run 21, $Re_\infty = 5.0 \times 10^6/\text{ft}$

d) Test 7057, Run 22, $Re_\infty = 6.5 \times 10^6/\text{ft}$



e) Test 7057, Run 23, $Re_\infty = 7.2 \times 10^6/\text{ft}$

f) Test 7057, Run 24, $Re_\infty = 8.1 \times 10^6/\text{ft}$

Figure 53. Reynolds Number effects, spherical-cap geometry, 140-mesh model images.

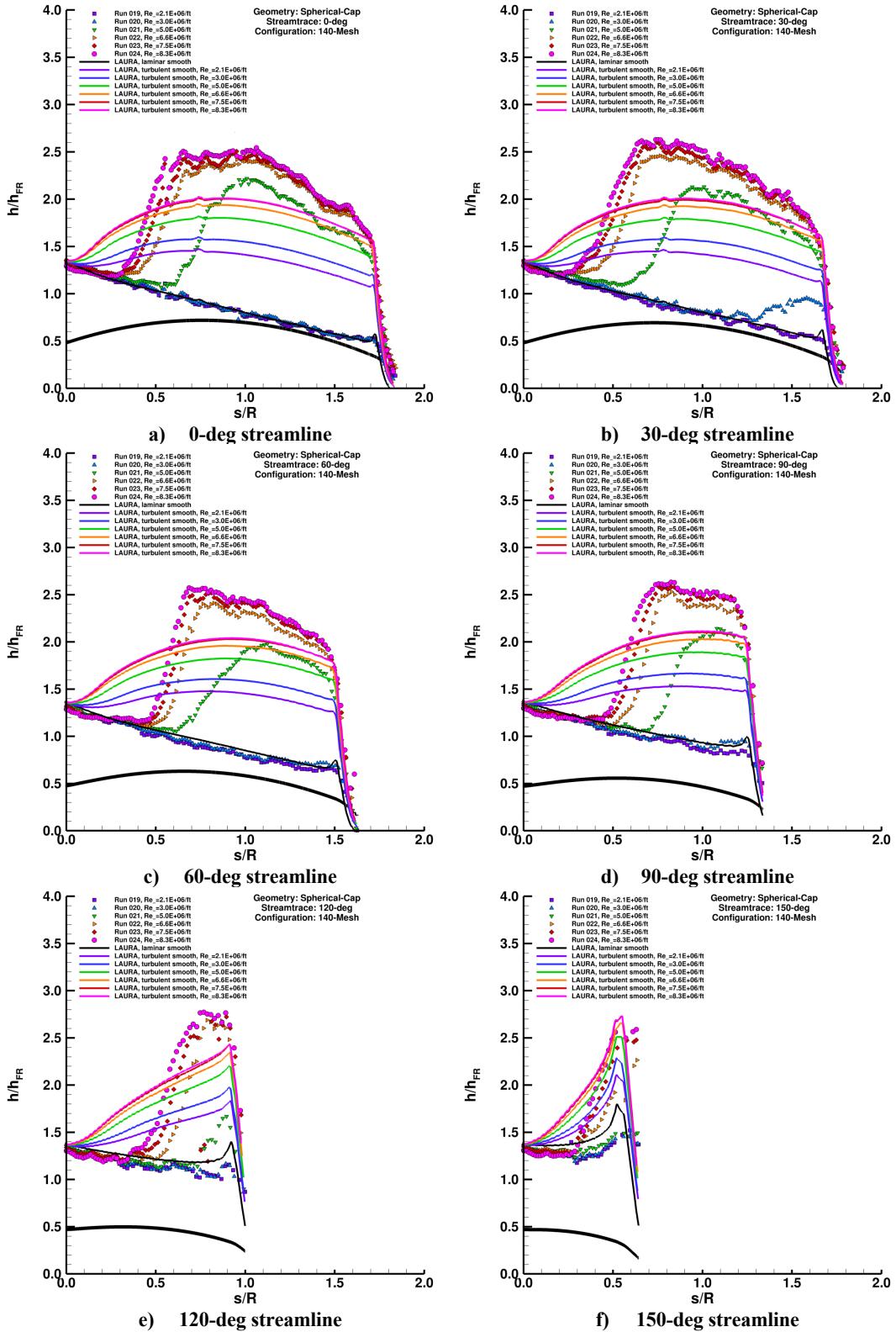


Figure 54. Reynolds Number effects, spherical-cap geometry, 140-mesh model plots.

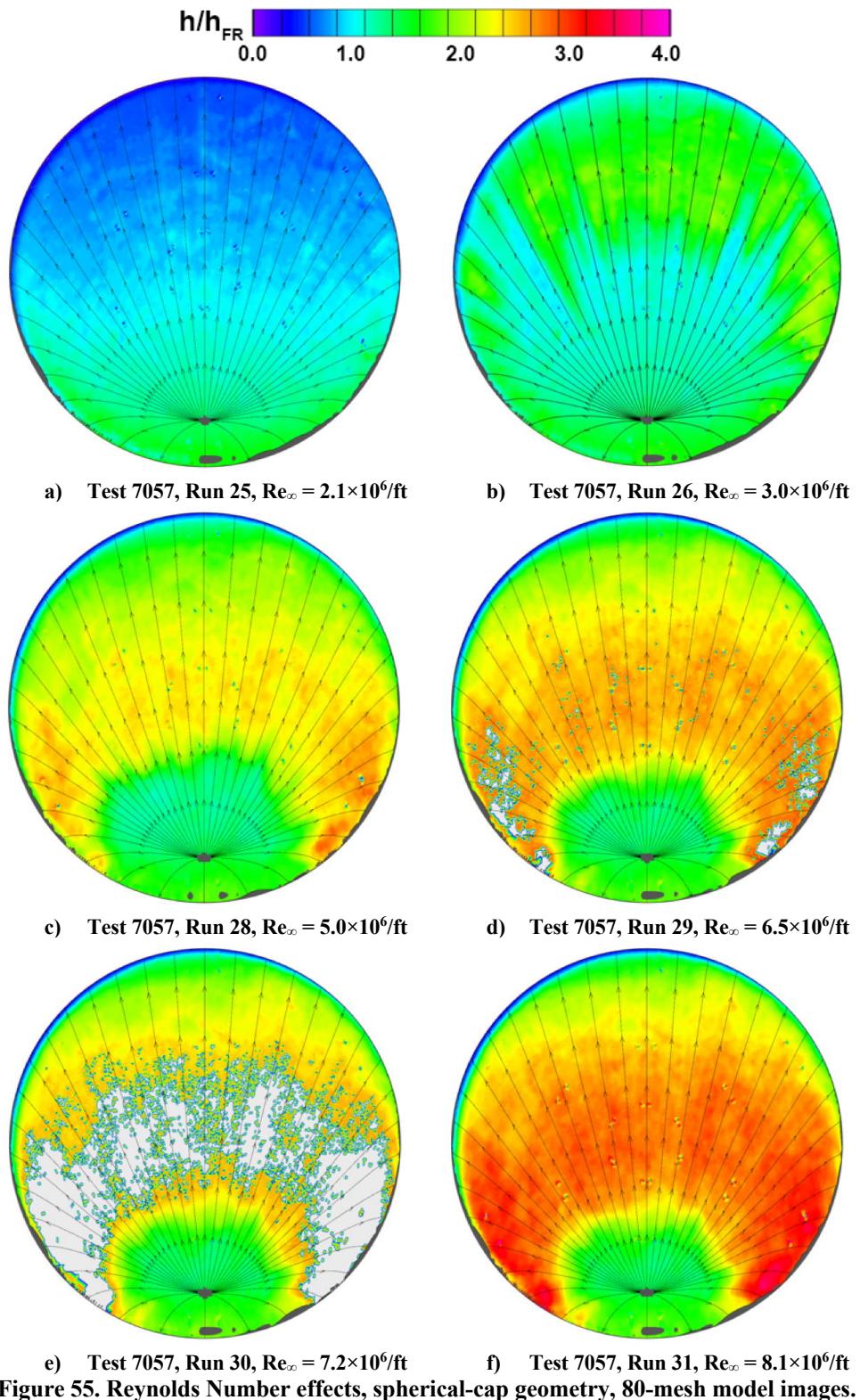


Figure 55. Reynolds Number effects, spherical-cap geometry, 80-mesh model images.

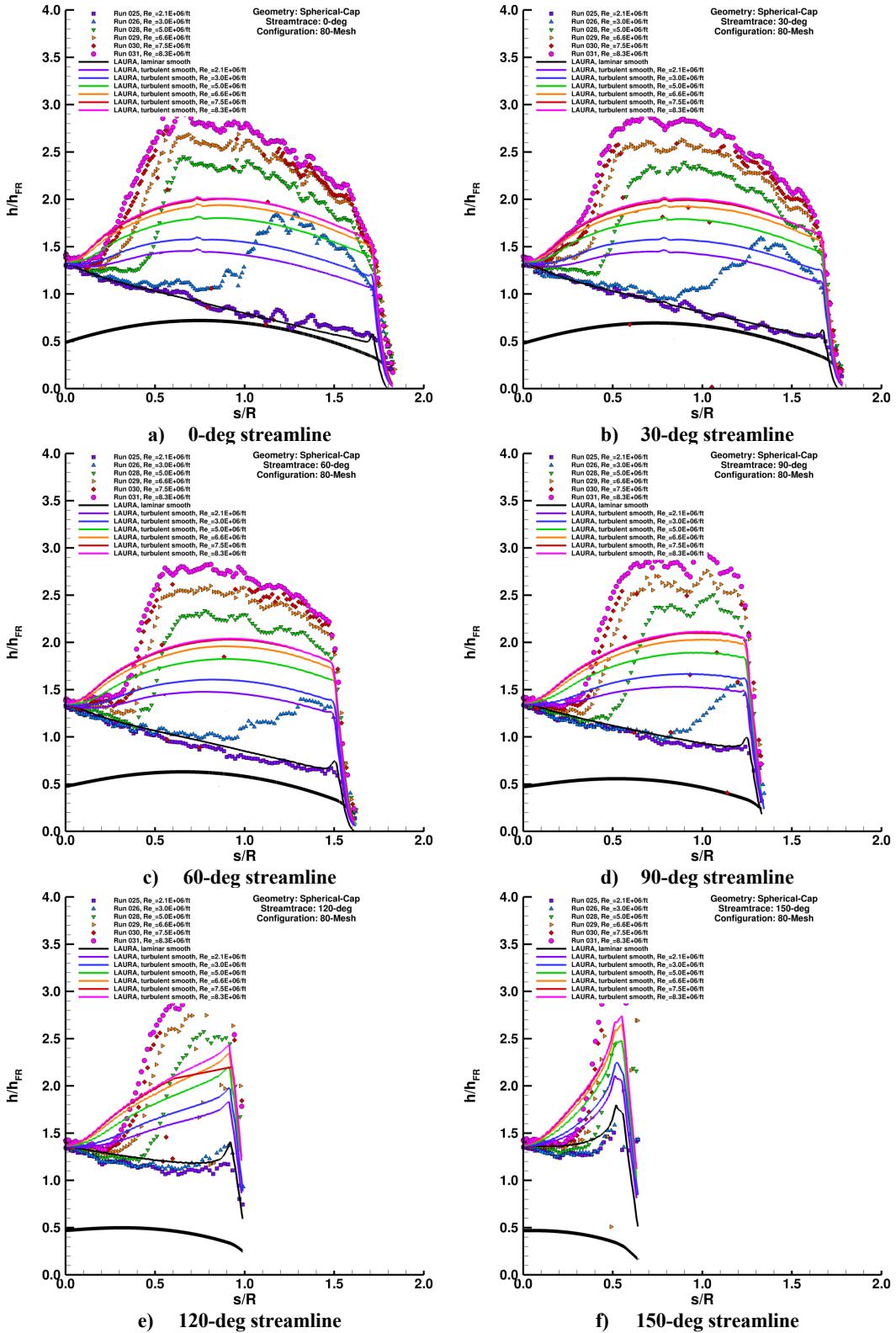


Figure 56. Reynolds Number effects, spherical-cap geometry, 80-mesh model plots.

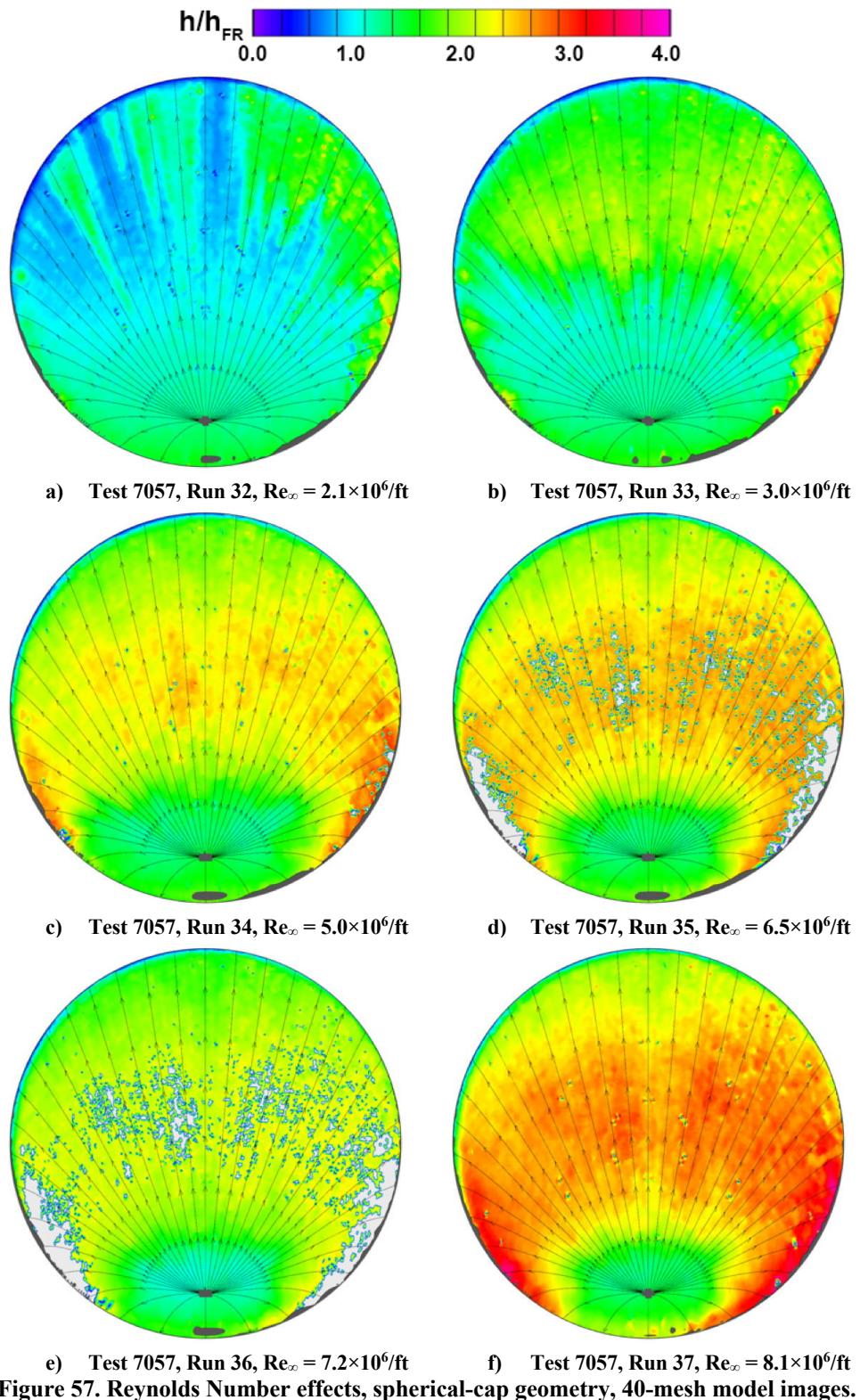


Figure 57. Reynolds Number effects, spherical-cap geometry, 40-mesh model images.

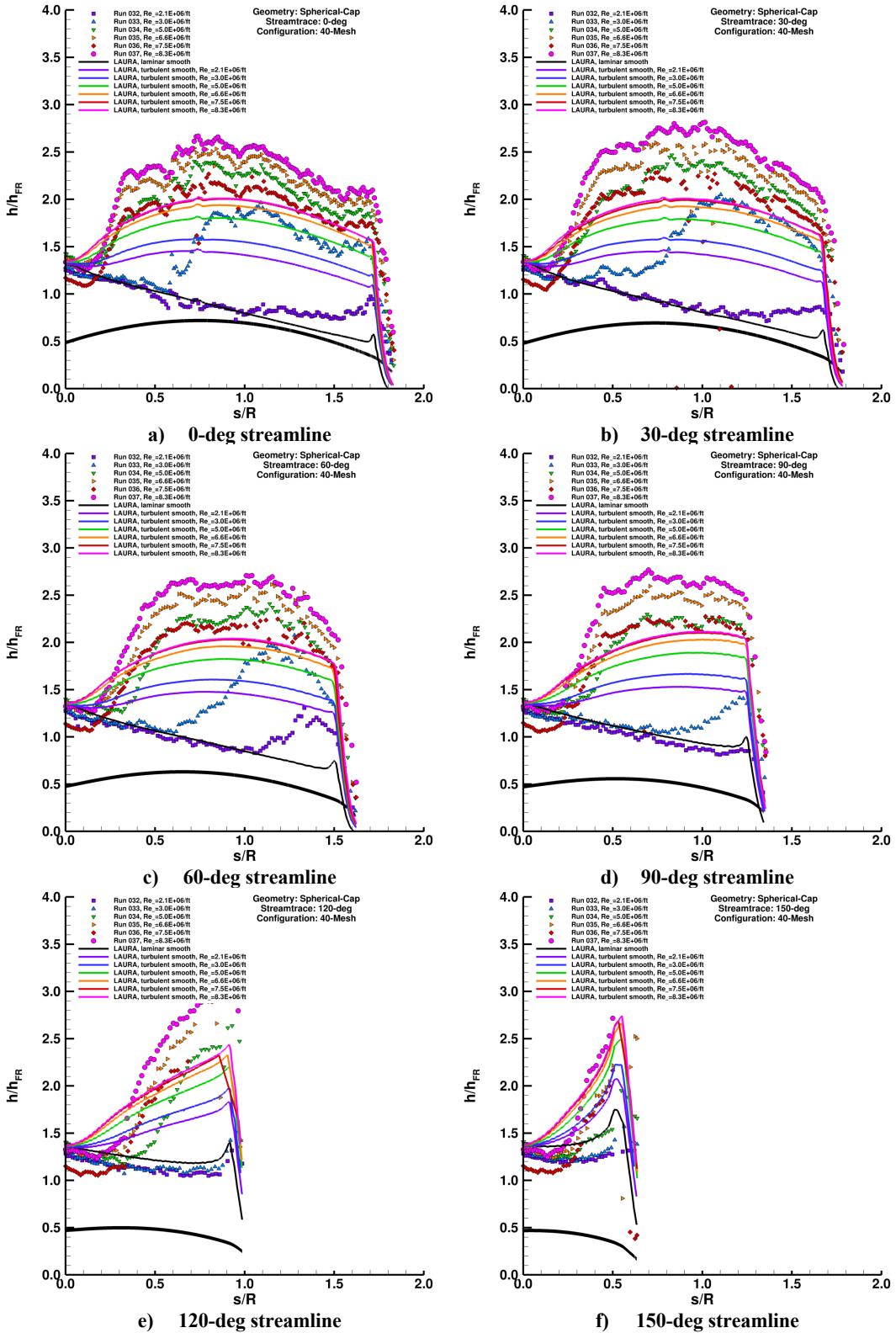


Figure 58. Reynolds Number effects, spherical-cap geometry, 40-mesh model plots.

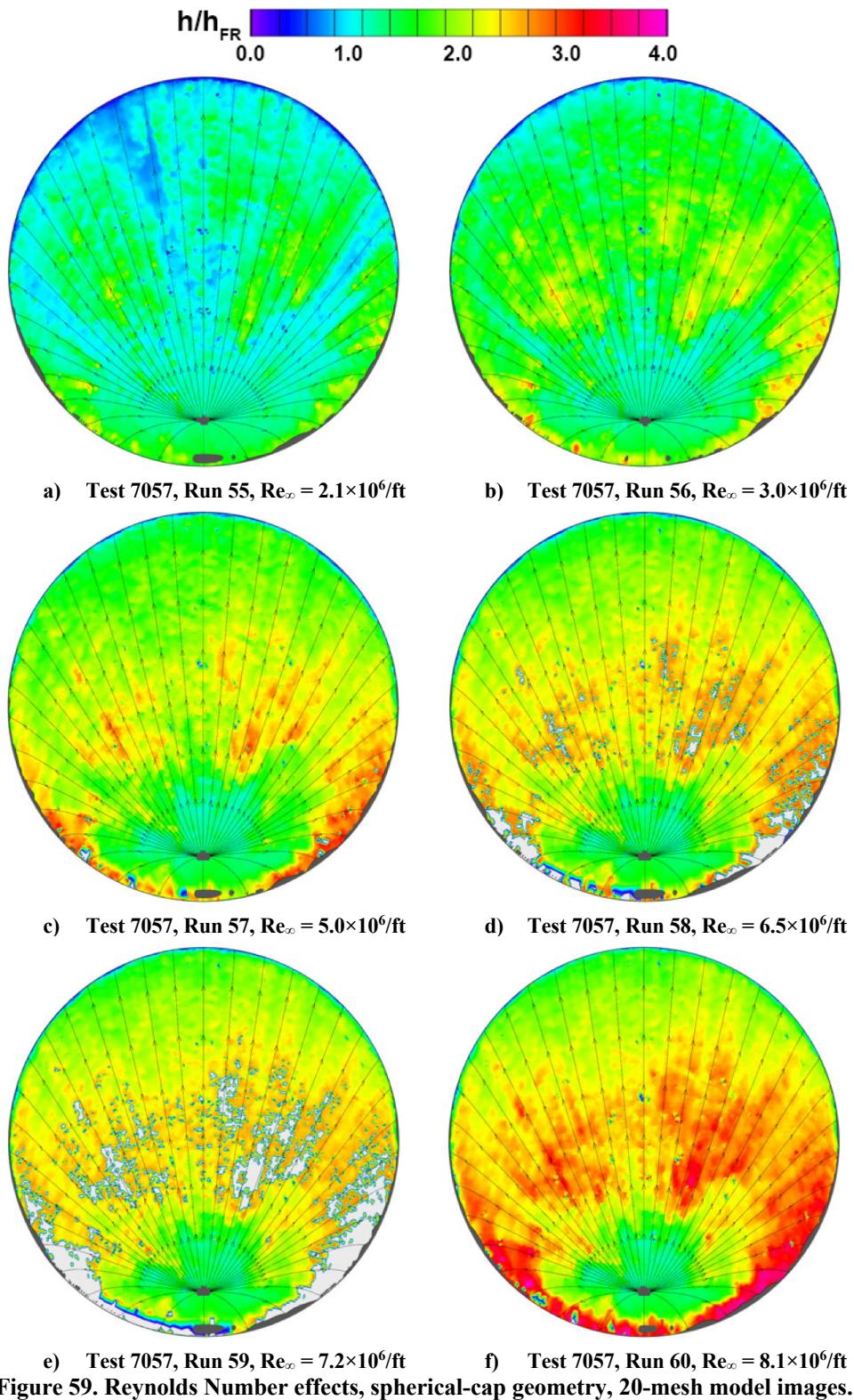


Figure 59. Reynolds Number effects, spherical-cap geometry, 20-mesh model images.

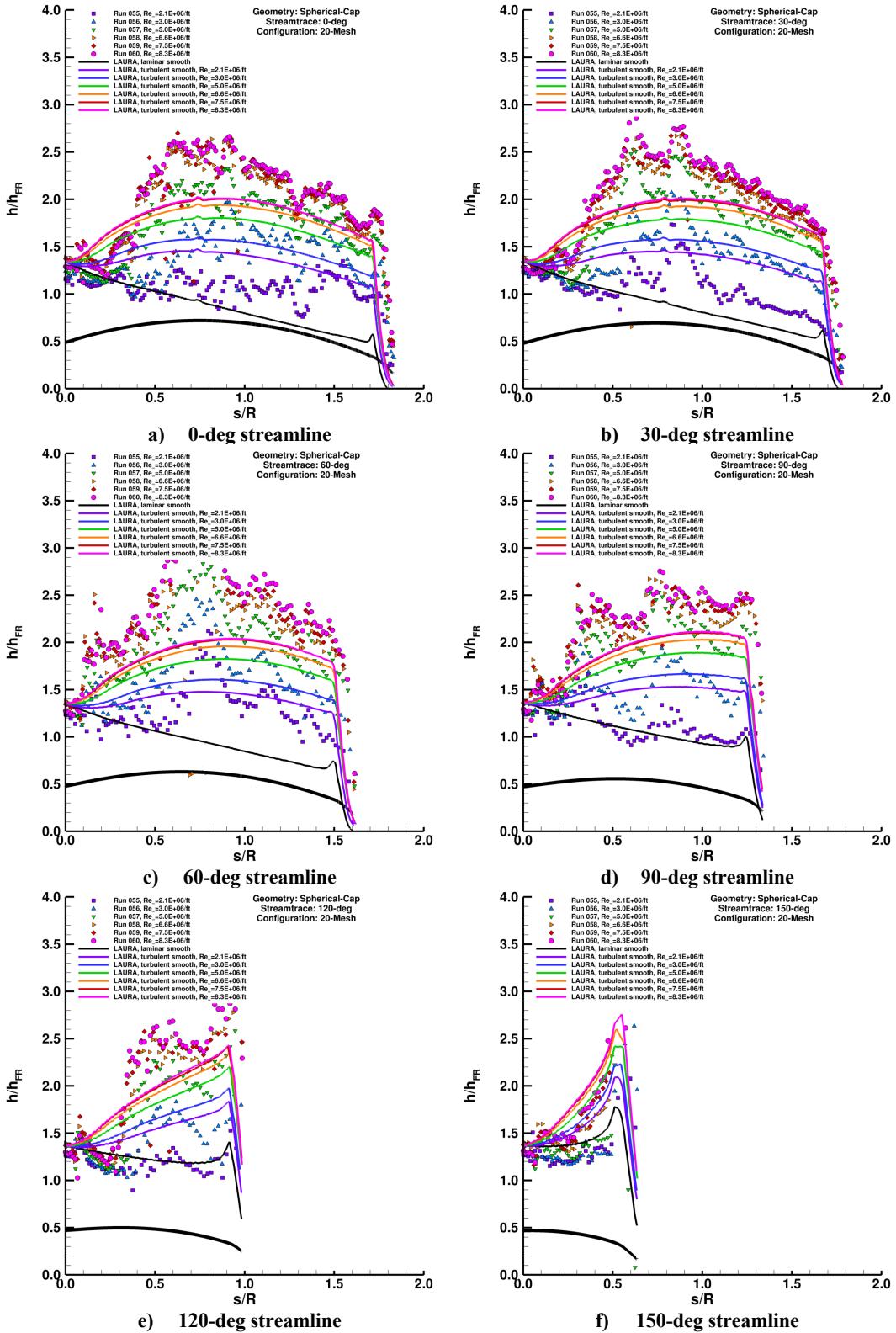


Figure 60. Reynolds Number effects, spherical-cap geometry, 20-mesh model plots.

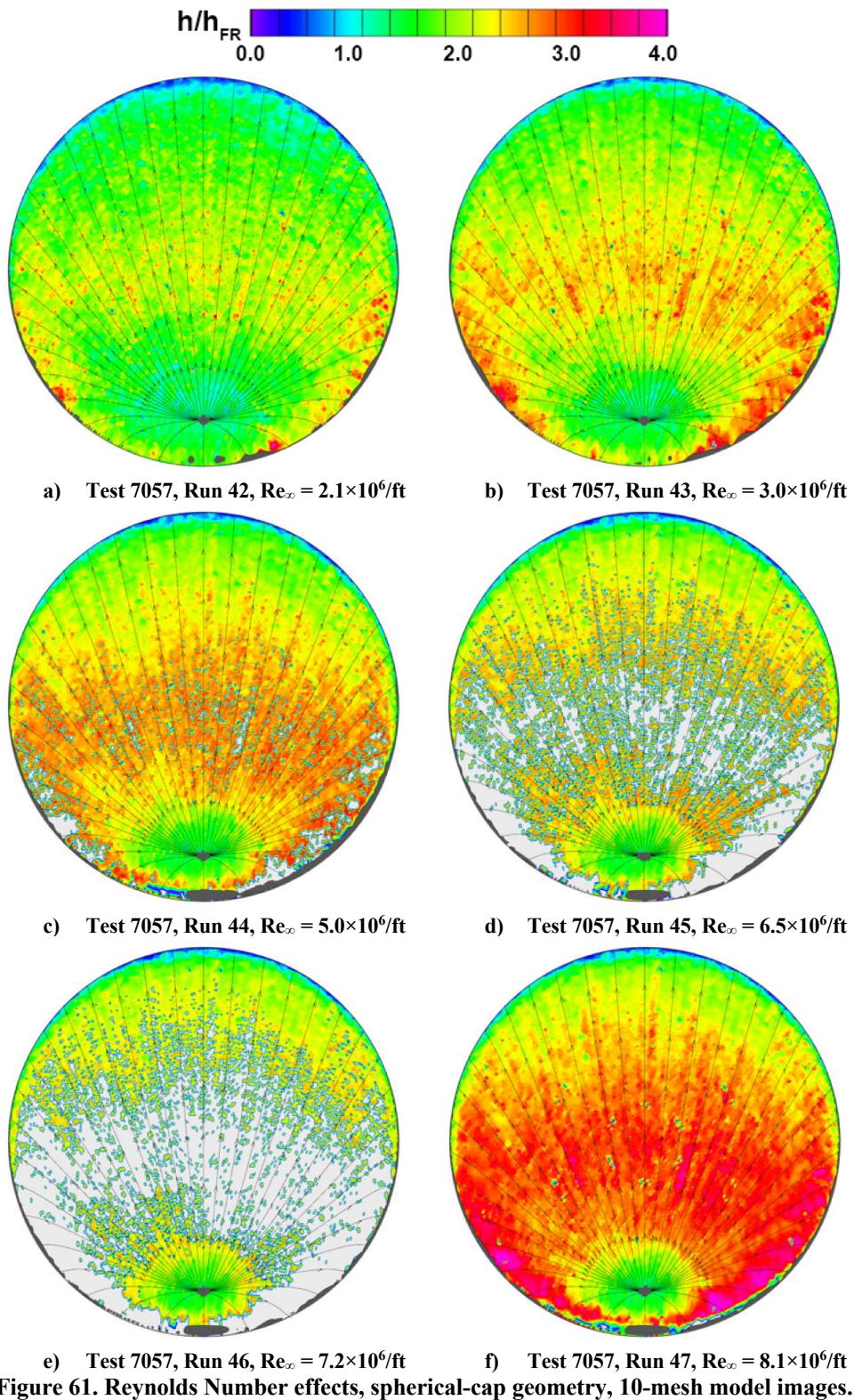


Figure 61. Reynolds Number effects, spherical-cap geometry, 10-mesh model images.

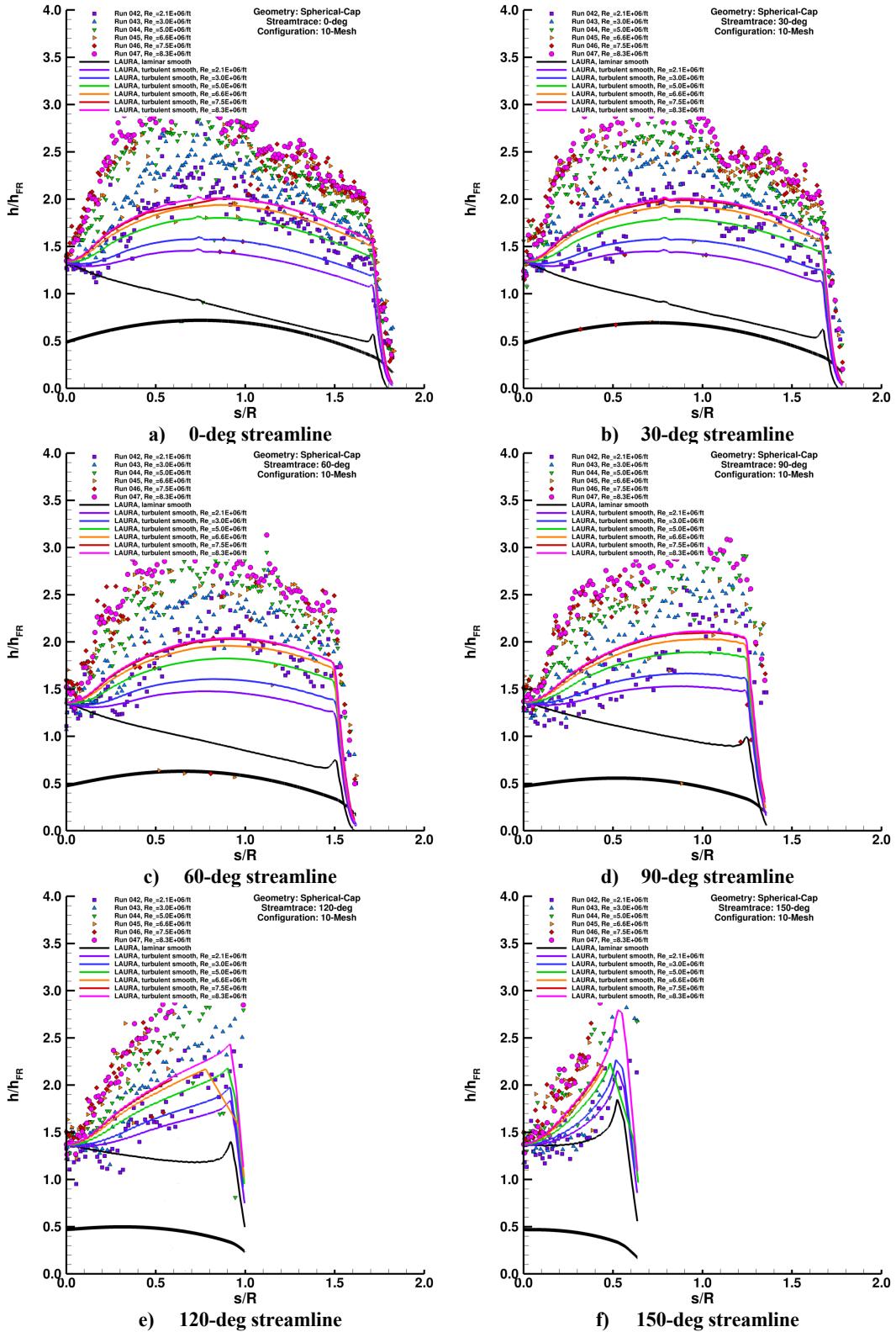


Figure 62. Reynolds Number effects, spherical-cap geometry, 10-mesh model plots.

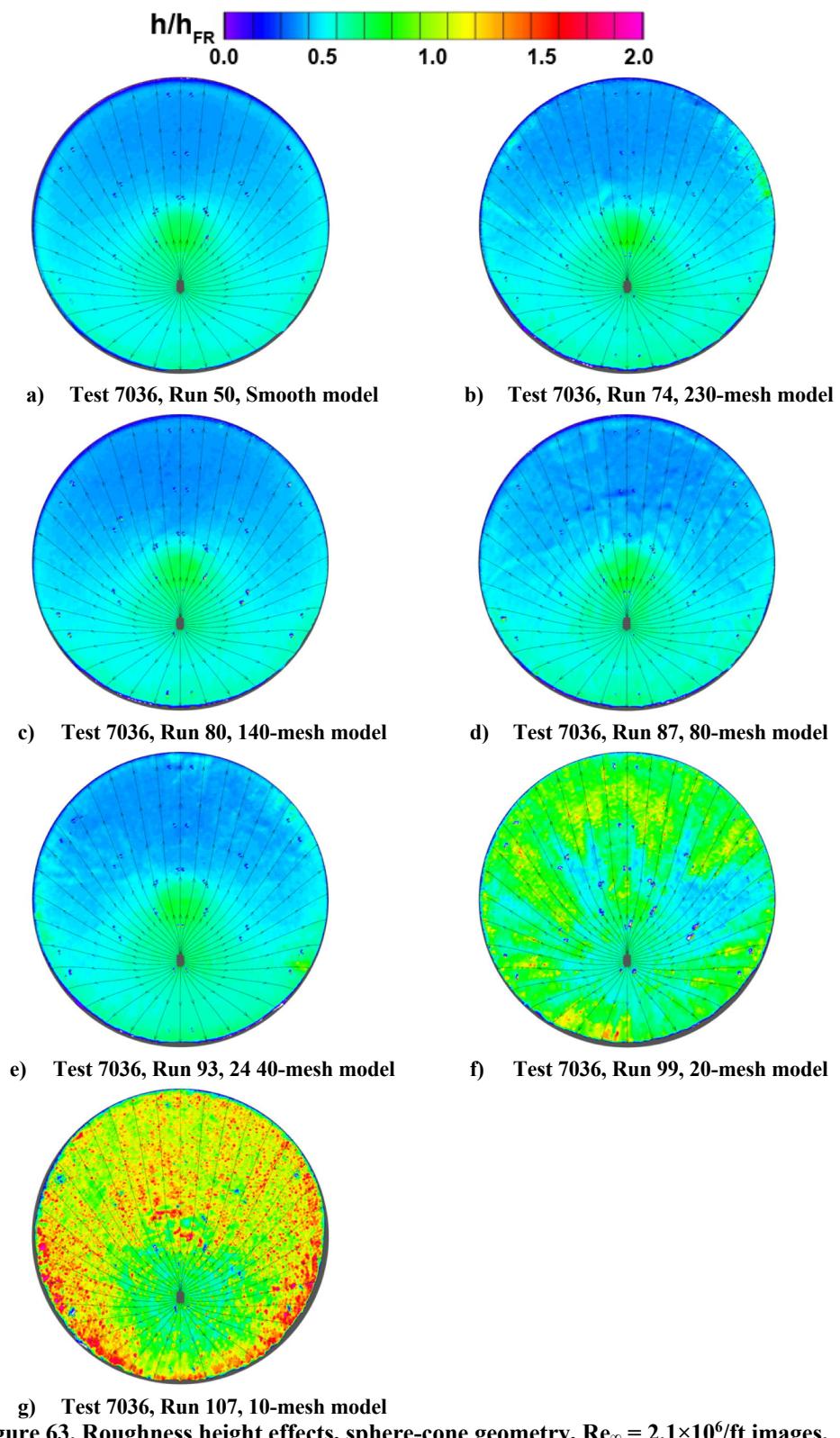


Figure 63. Roughness height effects, sphere-cone geometry, $Re_\infty = 2.1 \times 10^6/\text{ft}$ images.

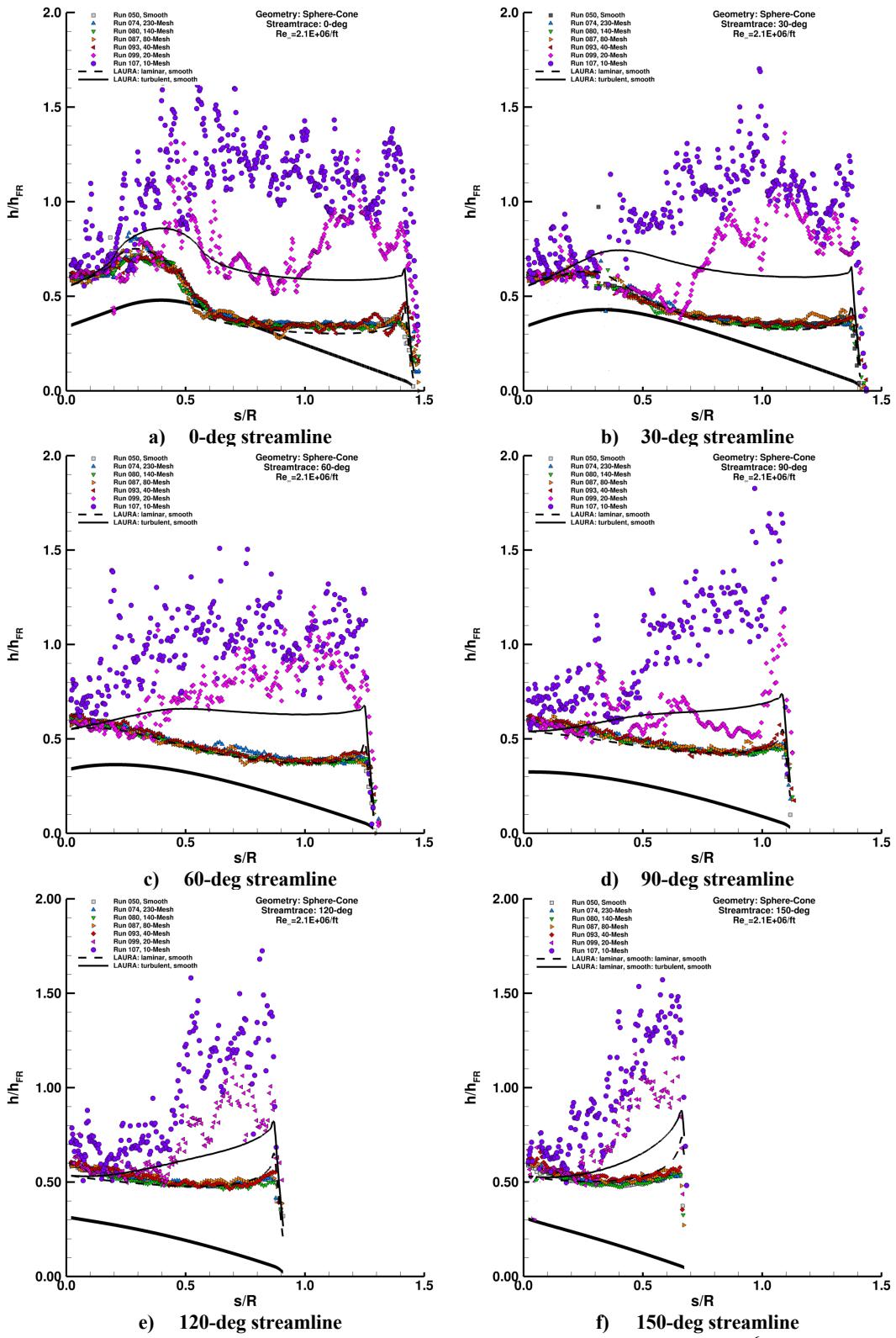


Figure 64. Roughness height effects, sphere-cone geometry, $Re_\infty = 2.1 \times 10^6 / \text{ft}$ plots.

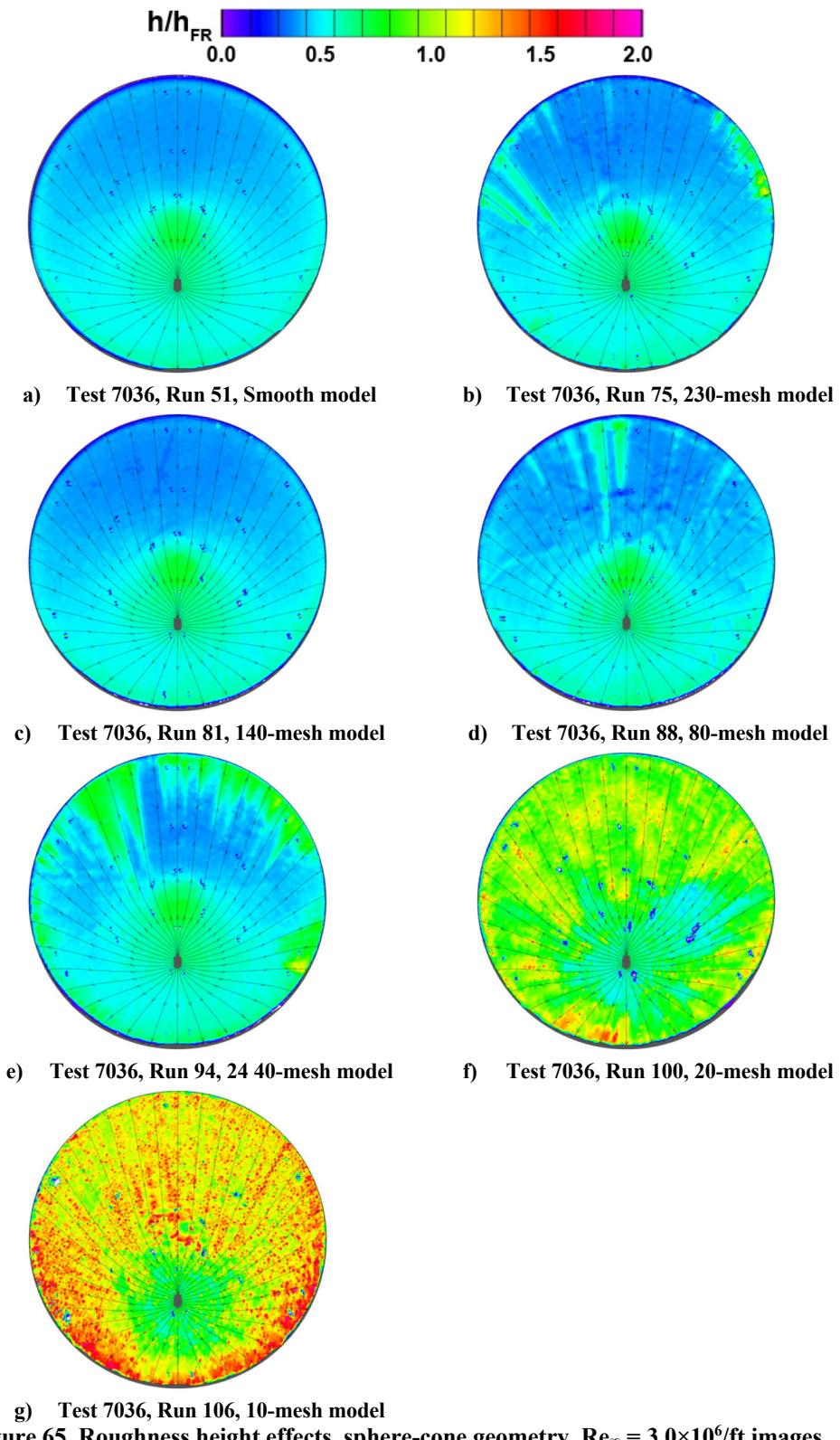


Figure 65. Roughness height effects, sphere-cone geometry, $Re_\infty = 3.0 \times 10^6/\text{ft}$ images.

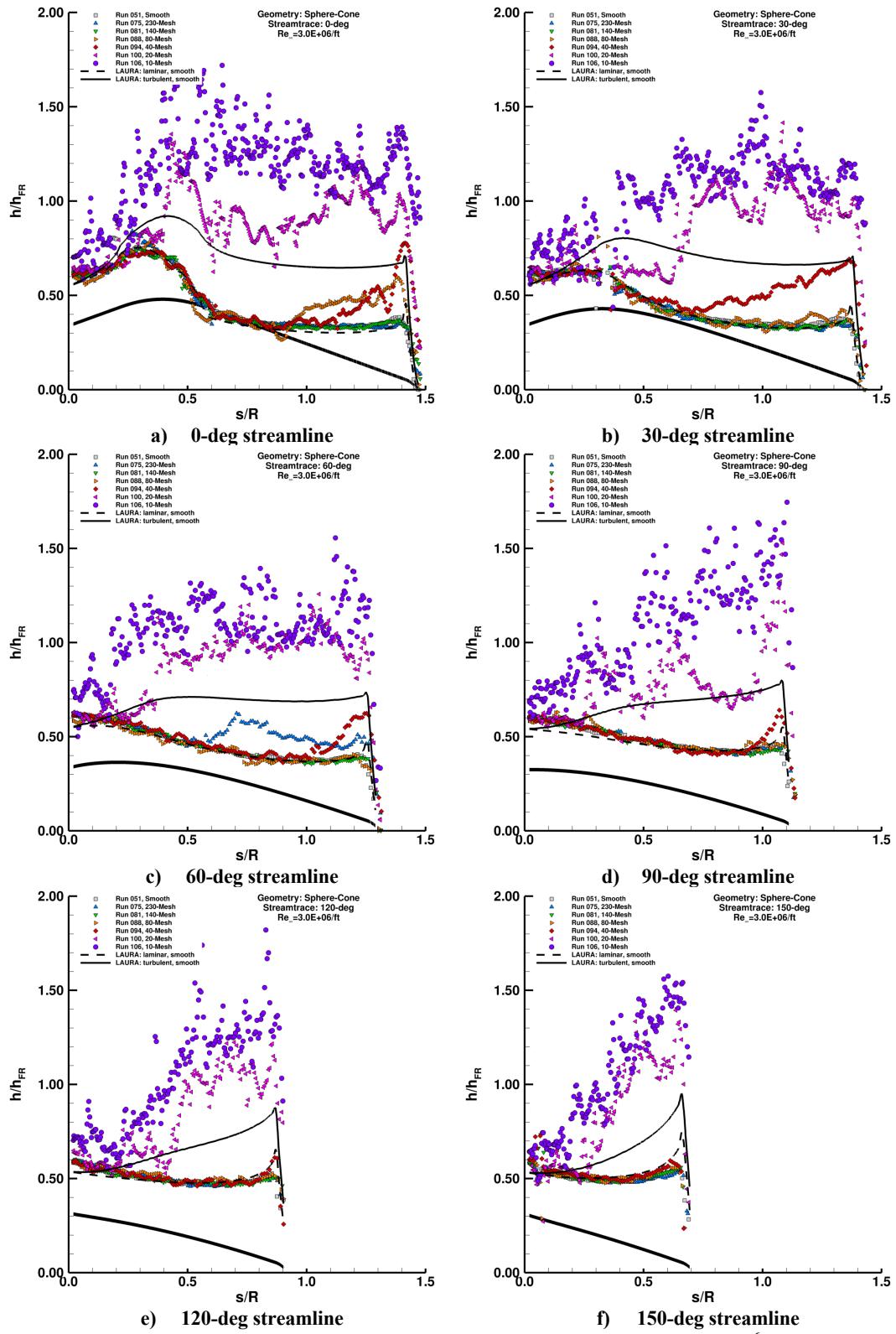


Figure 66. Roughness height effects, sphere-cone geometry, $Re_\infty = 3.0 \times 10^6$ ft plots.

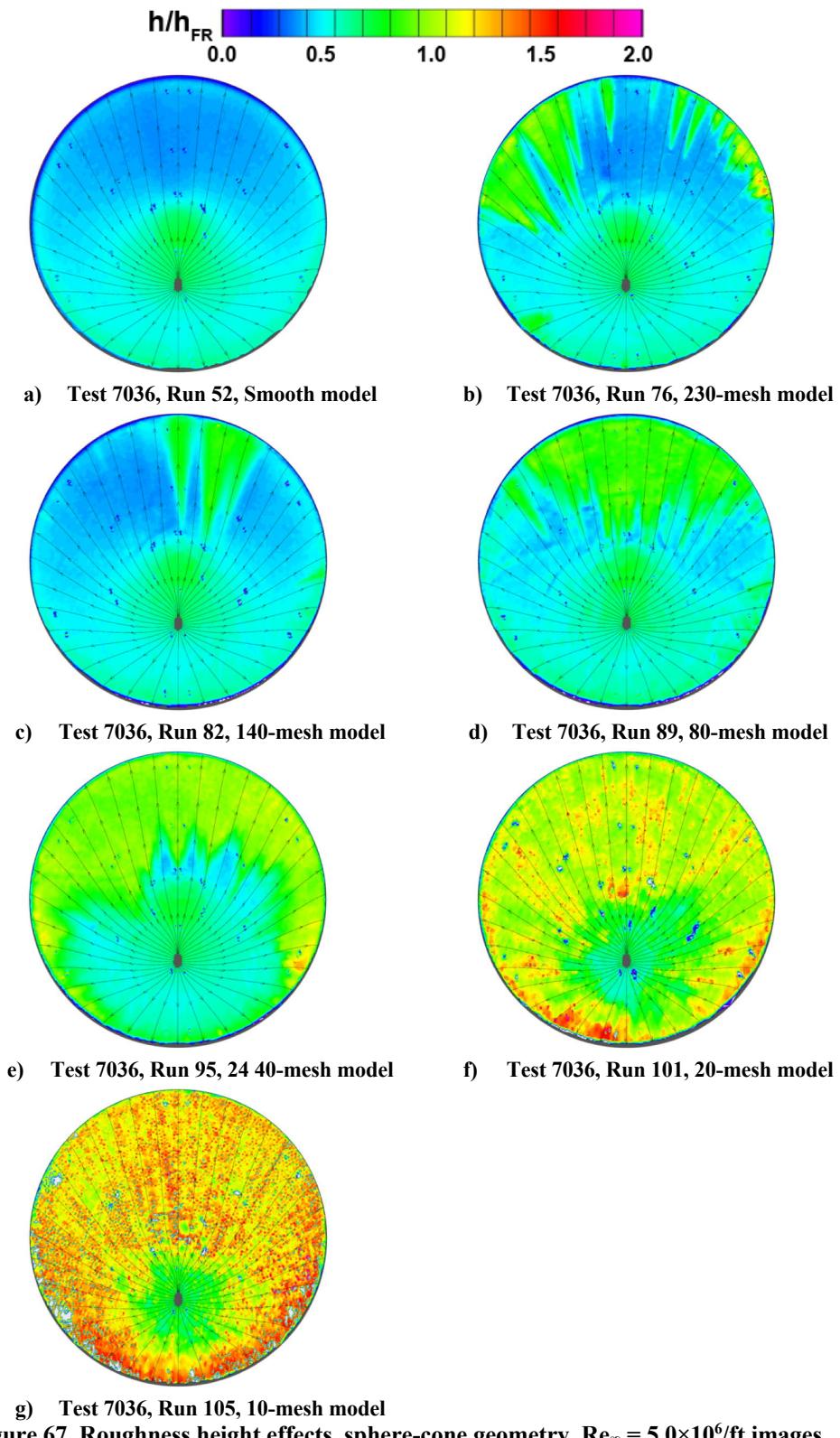


Figure 67. Roughness height effects, sphere-cone geometry, $Re_\infty = 5.0 \times 10^6/\text{ft}$ images.

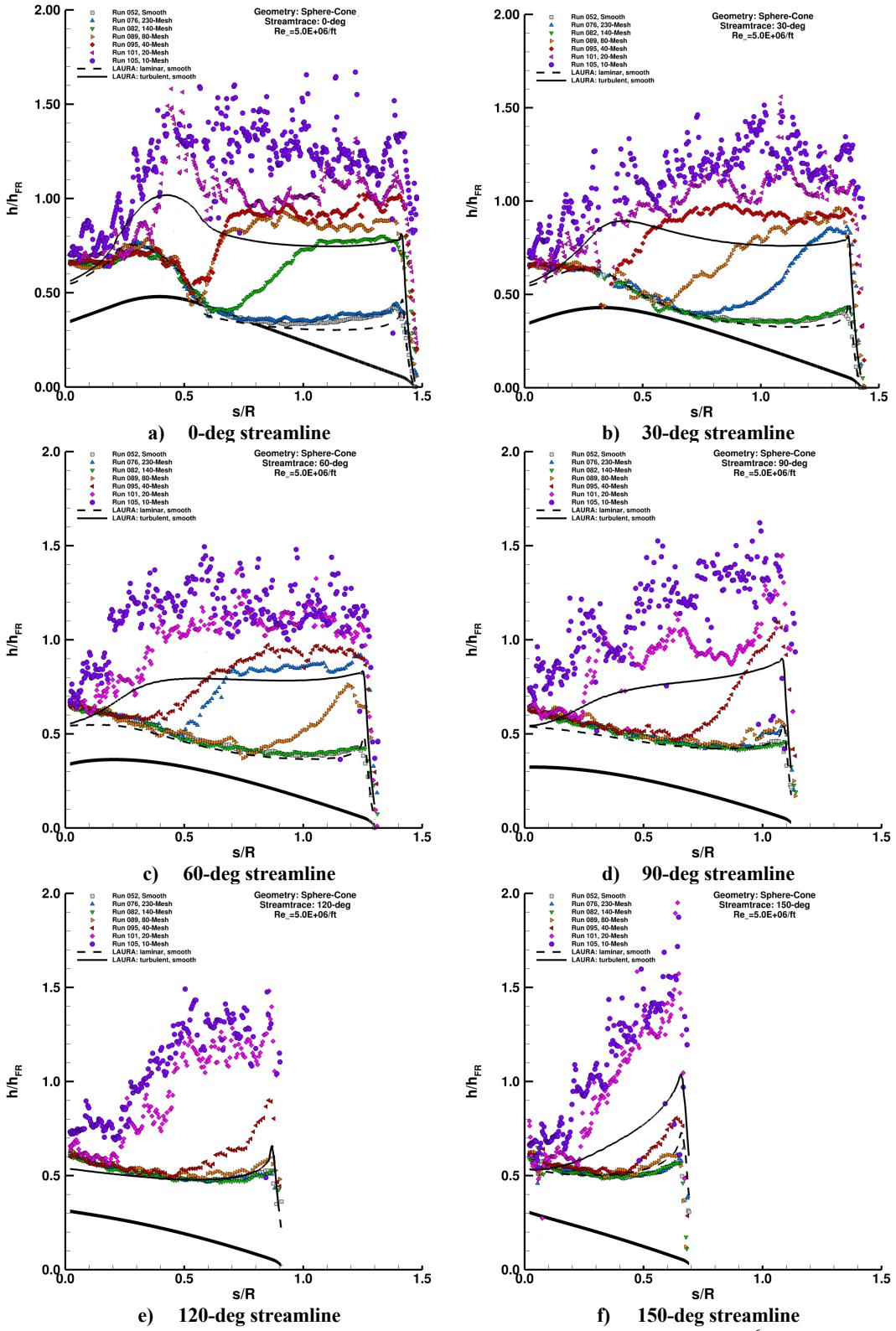


Figure 68. Roughness height effects, sphere-cone geometry, $Re_\infty = 5.0 \times 10^6$ ft plots.

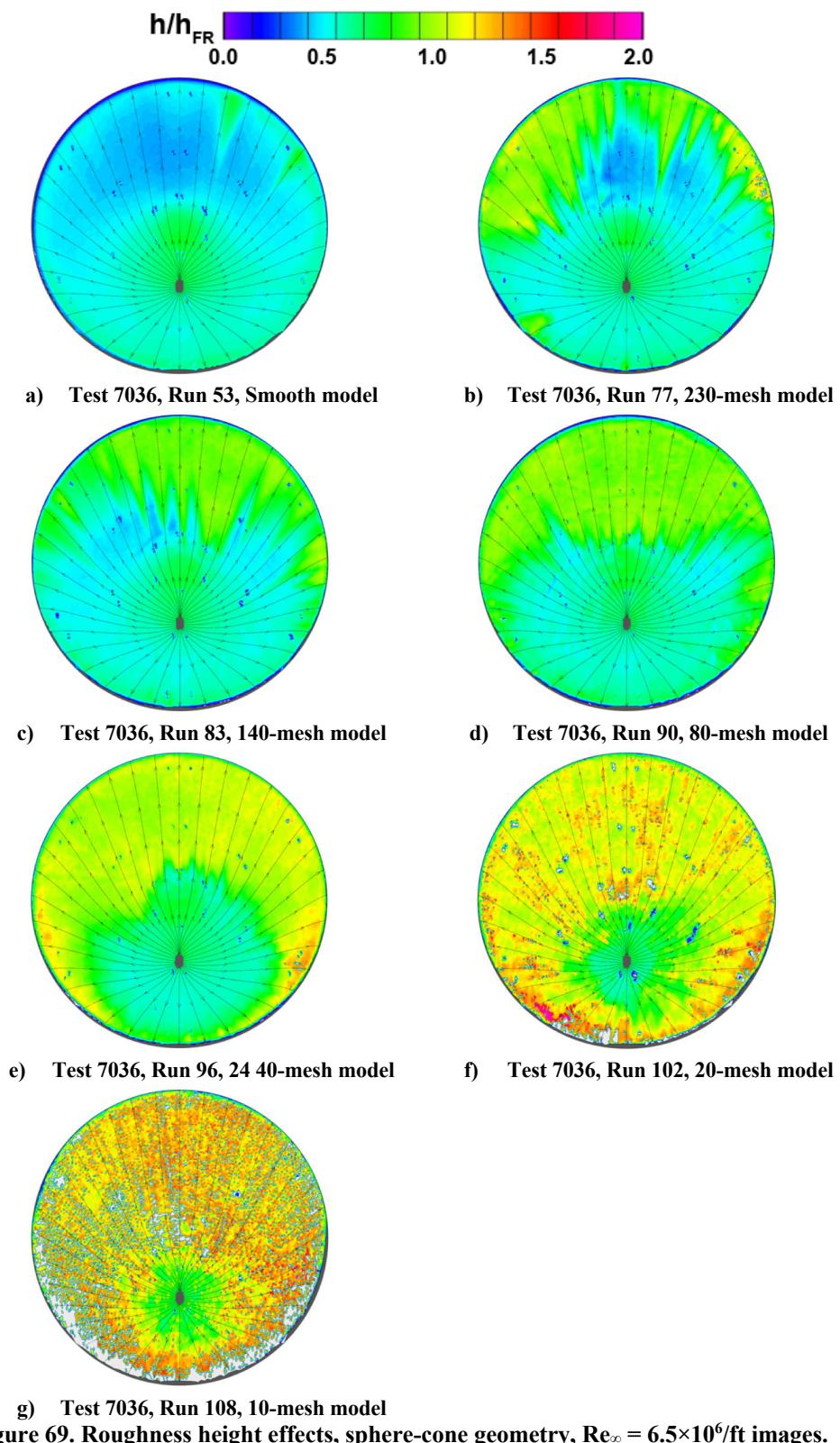


Figure 69. Roughness height effects, sphere-cone geometry, $Re_\infty = 6.5 \times 10^6/\text{ft}$ images.

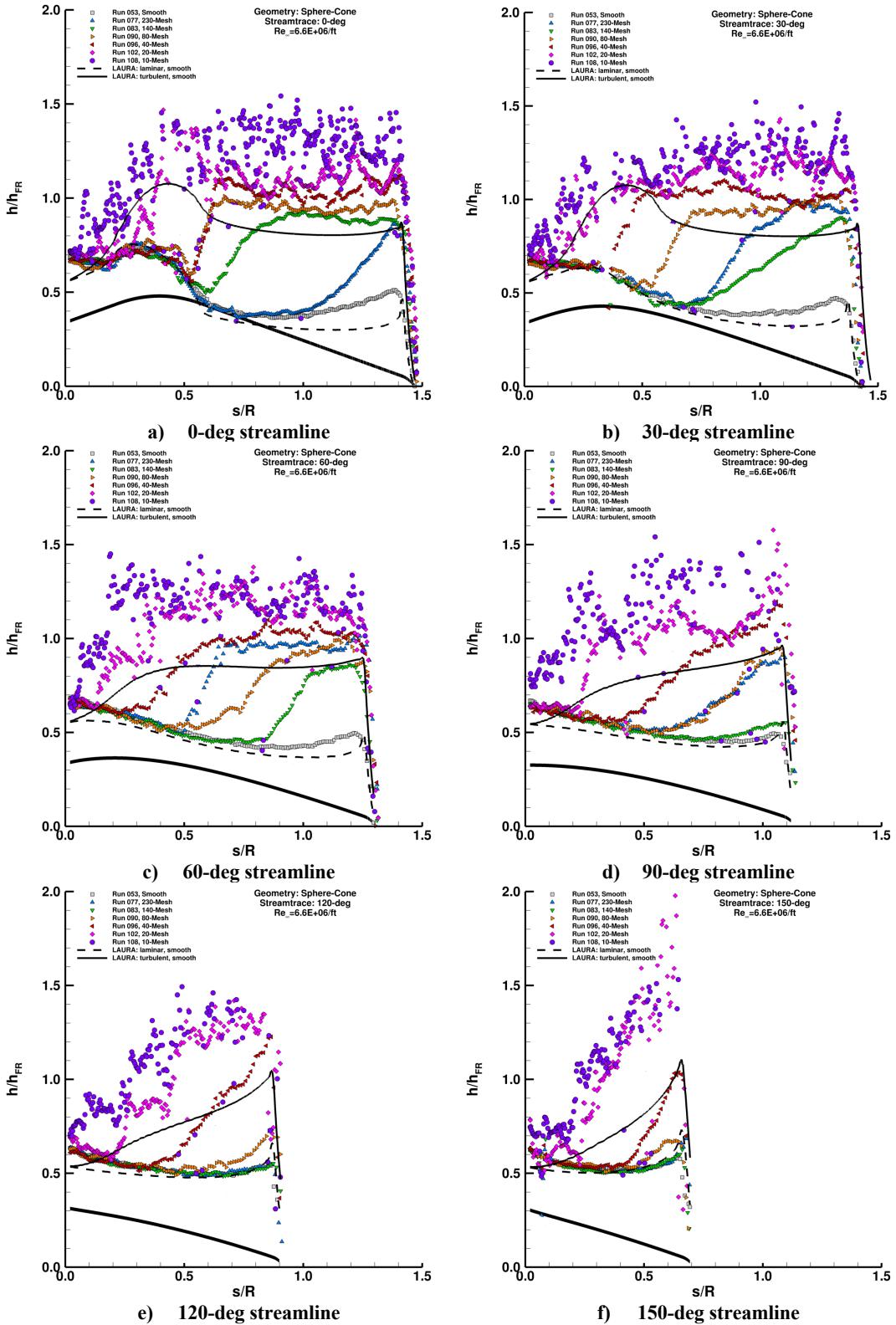


Figure 70. Roughness height effects, sphere-cone geometry, $Re_{\infty} = 6.5 \times 10^6$ ft plots.

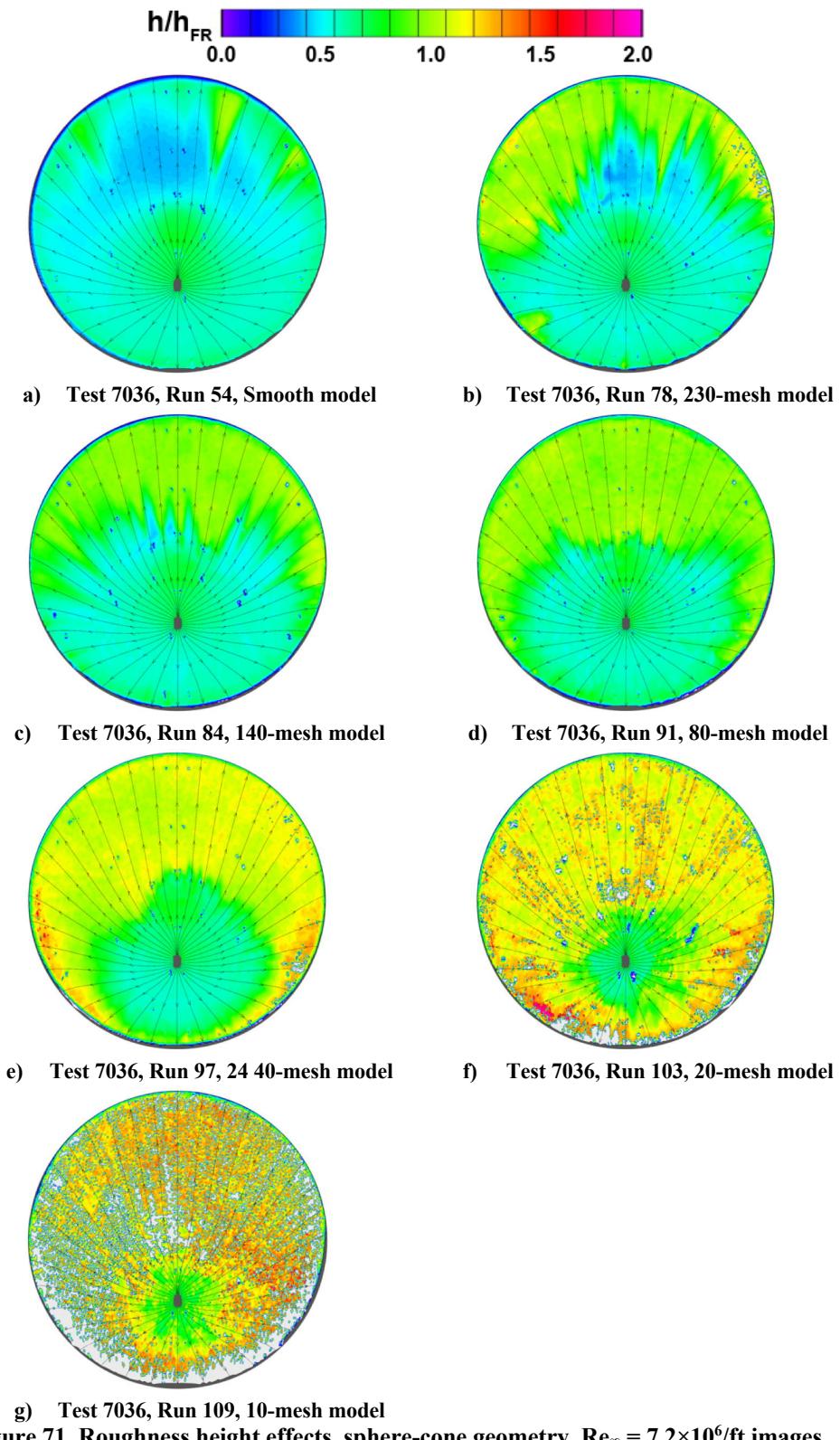


Figure 71. Roughness height effects, sphere-cone geometry, $Re_\infty = 7.2 \times 10^6/\text{ft}$ images.

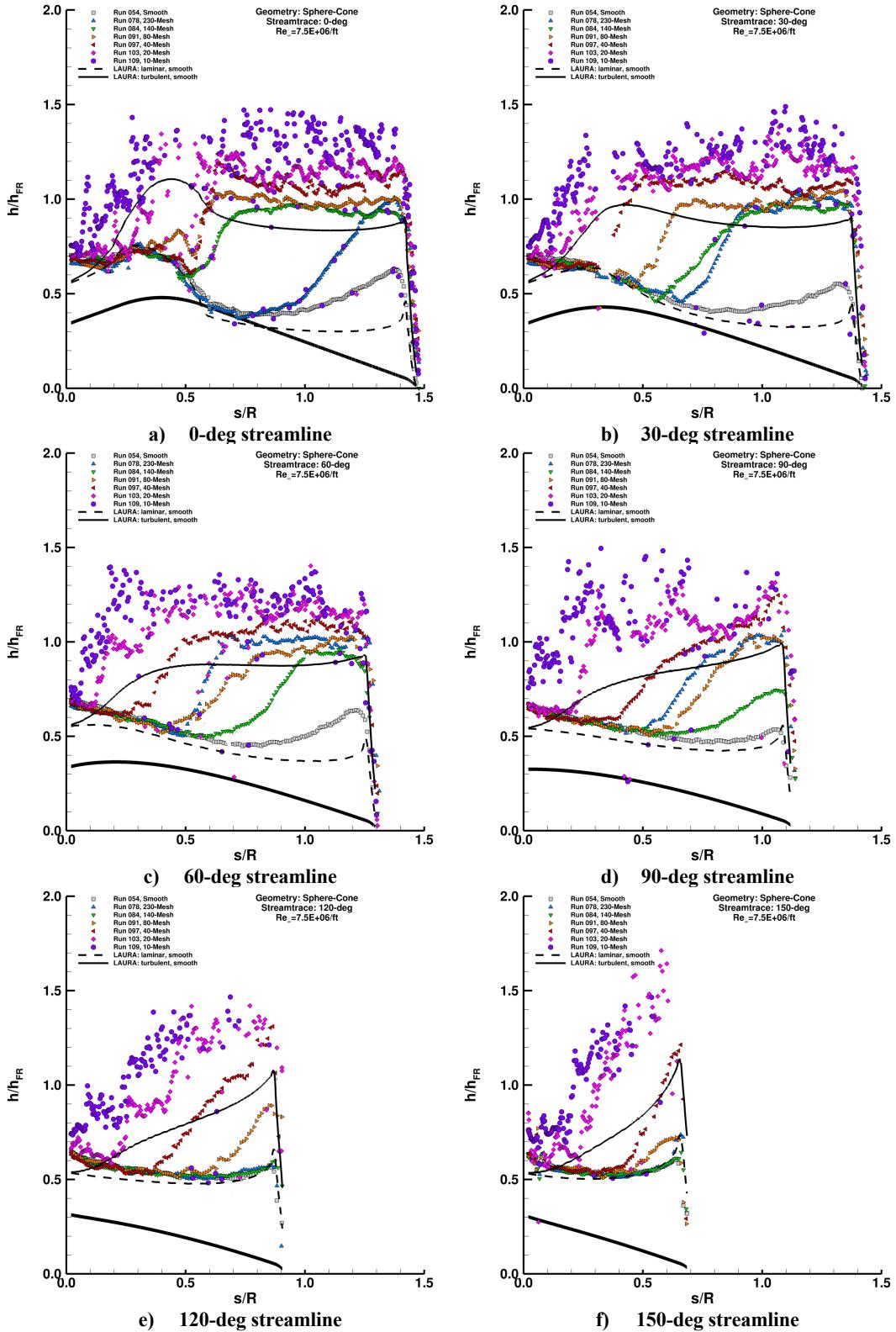


Figure 72. Roughness height effects, sphere-cone geometry, $Re_\infty = 7.2 \times 10^6$ ft plots.

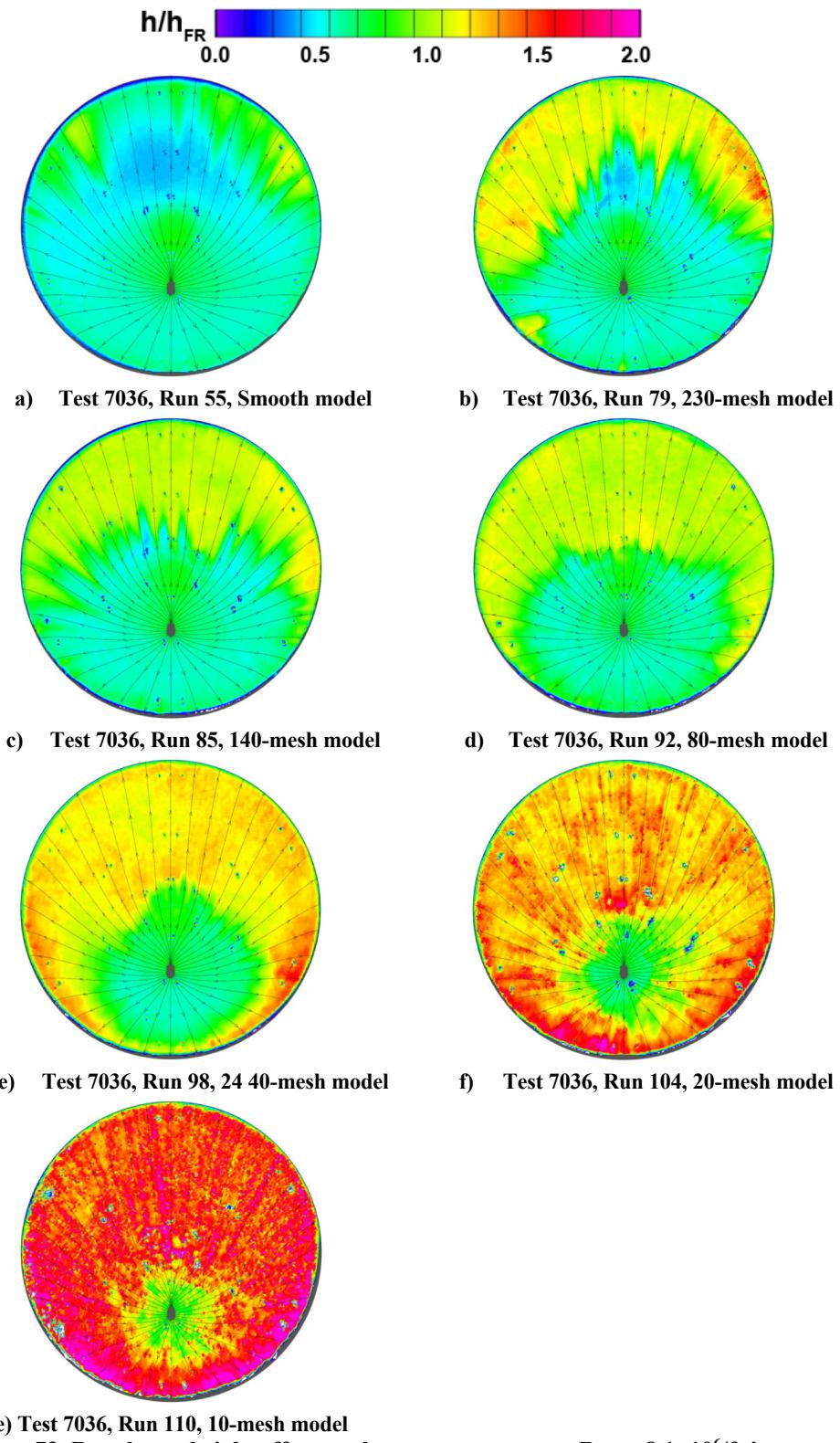


Figure 73. Roughness height effects, sphere-cone geometry, $Re_\infty = 8.1 \times 10^6/\text{ft}$ images.

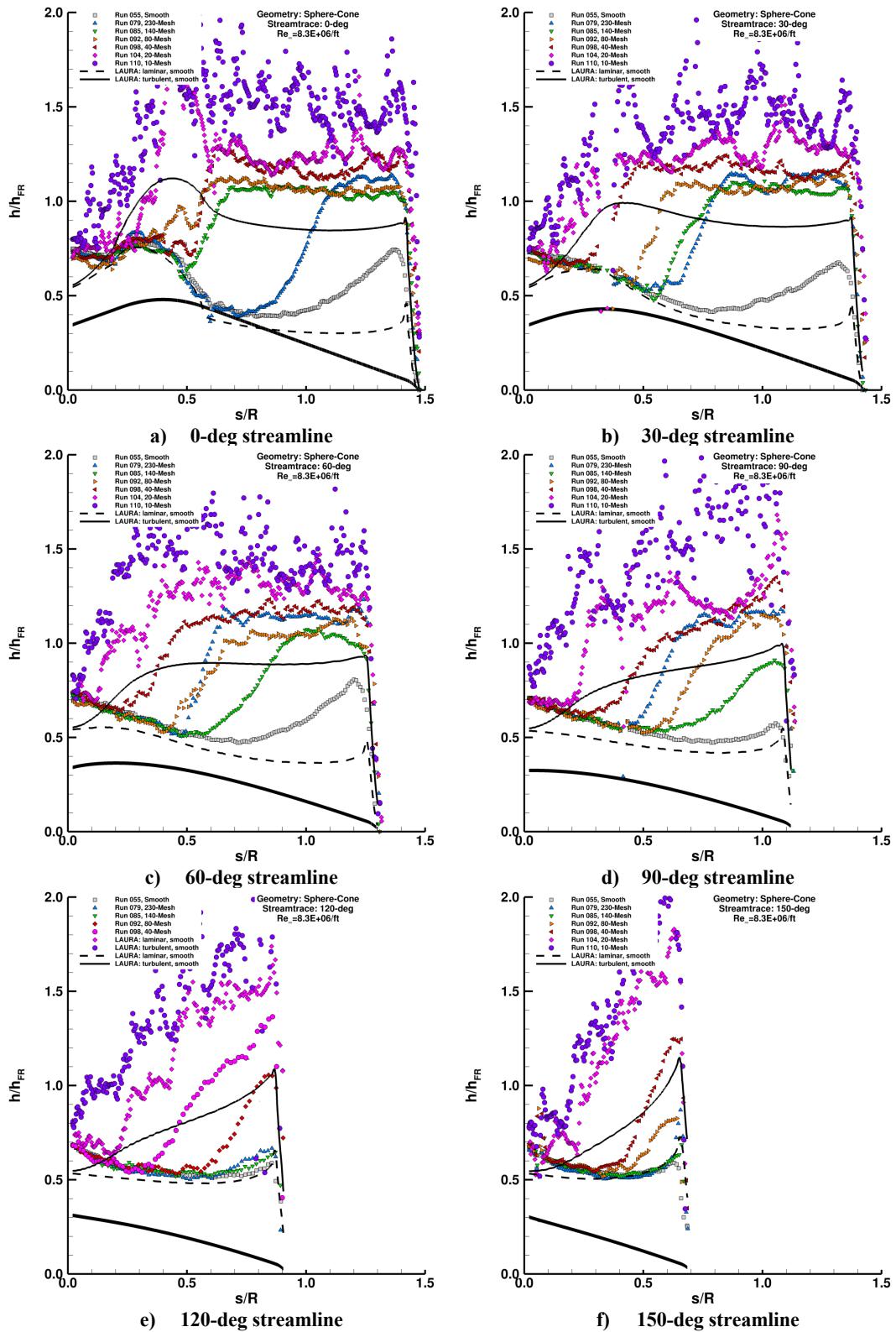


Figure 74. Roughness height effects, sphere-cone geometry, $Re_\infty = 8.1 \times 10^6$ ft plots.

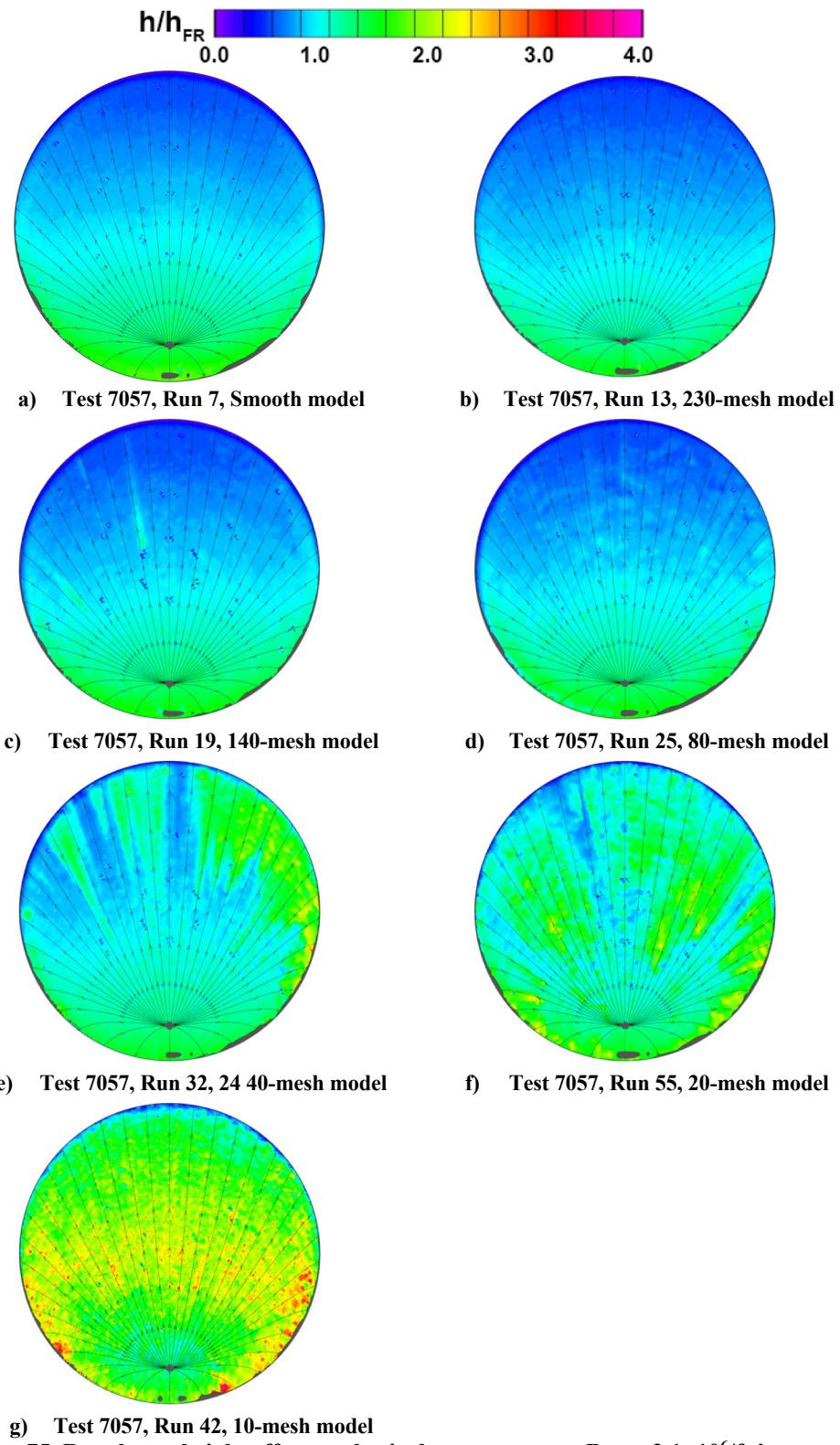


Figure 75. Roughness height effects, spherical-cap geometry, $Re_\infty = 2.1 \times 10^6/\text{ft}$ images.

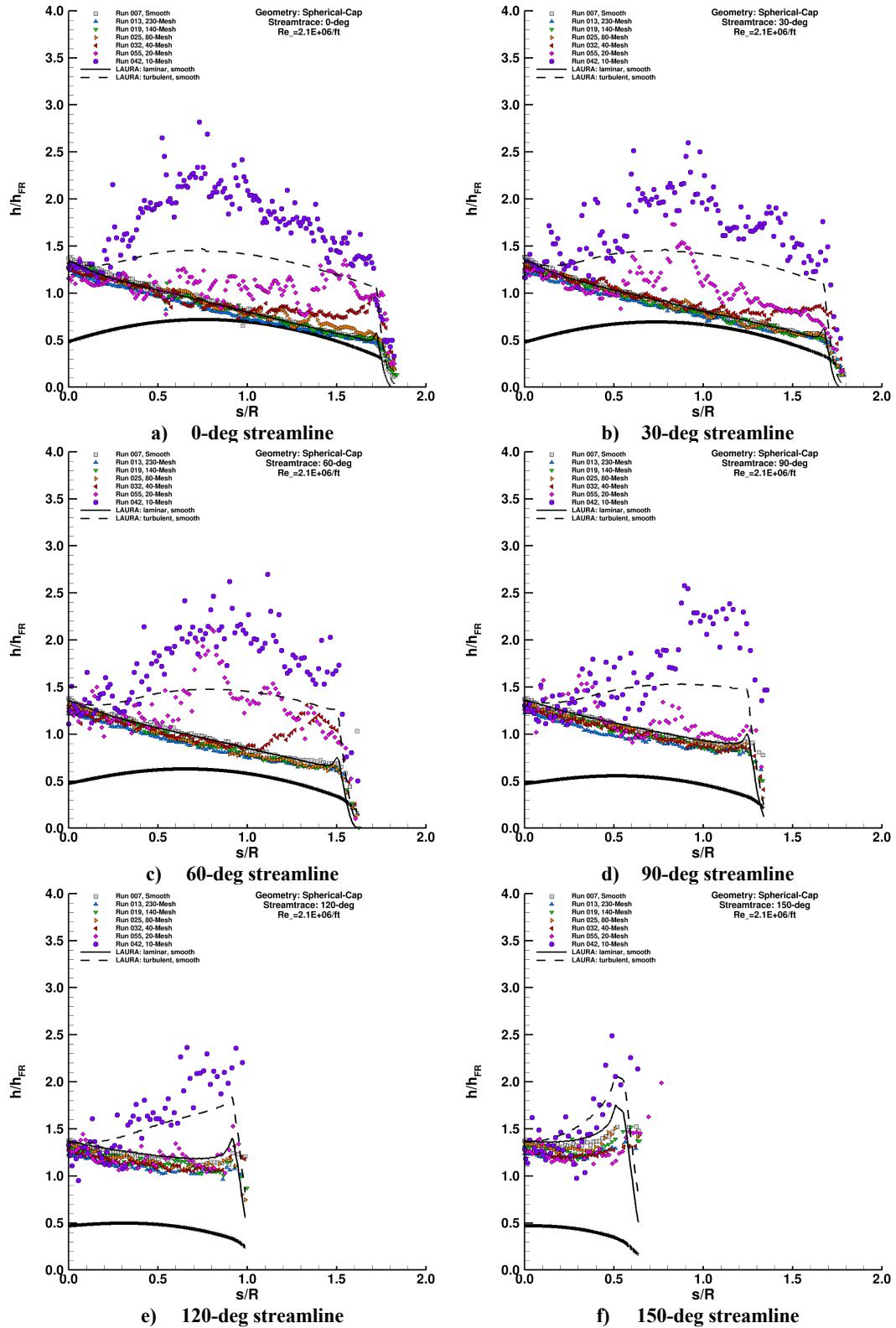


Figure 76. Roughness height effects, spherical-cap geometry, $Re_\infty = 2.1 \times 10^6/\text{ft}$ plots.

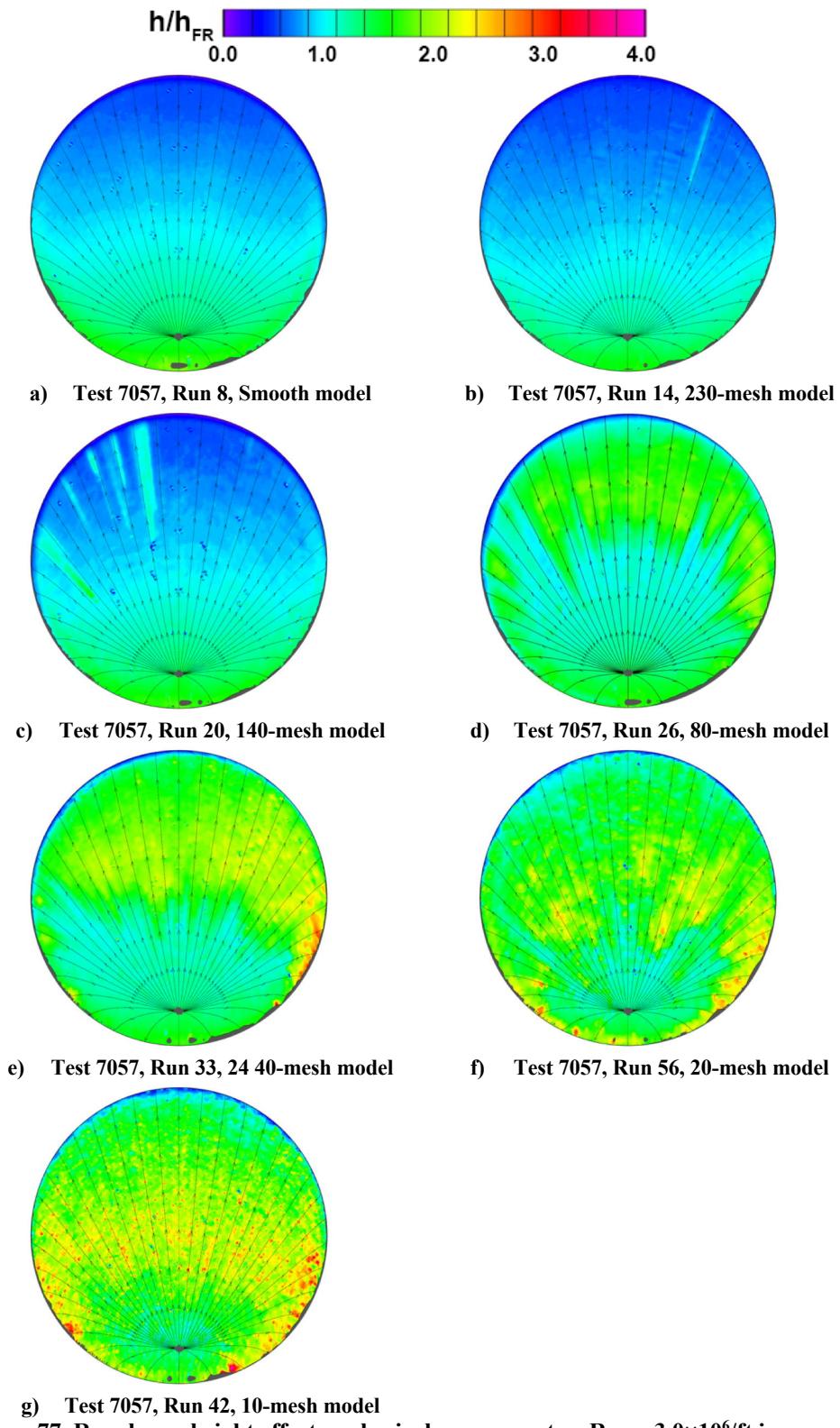


Figure 77. Roughness height effects, spherical-cap geometry, $Re_\infty = 3.0 \times 10^6/\text{ft}$ images.

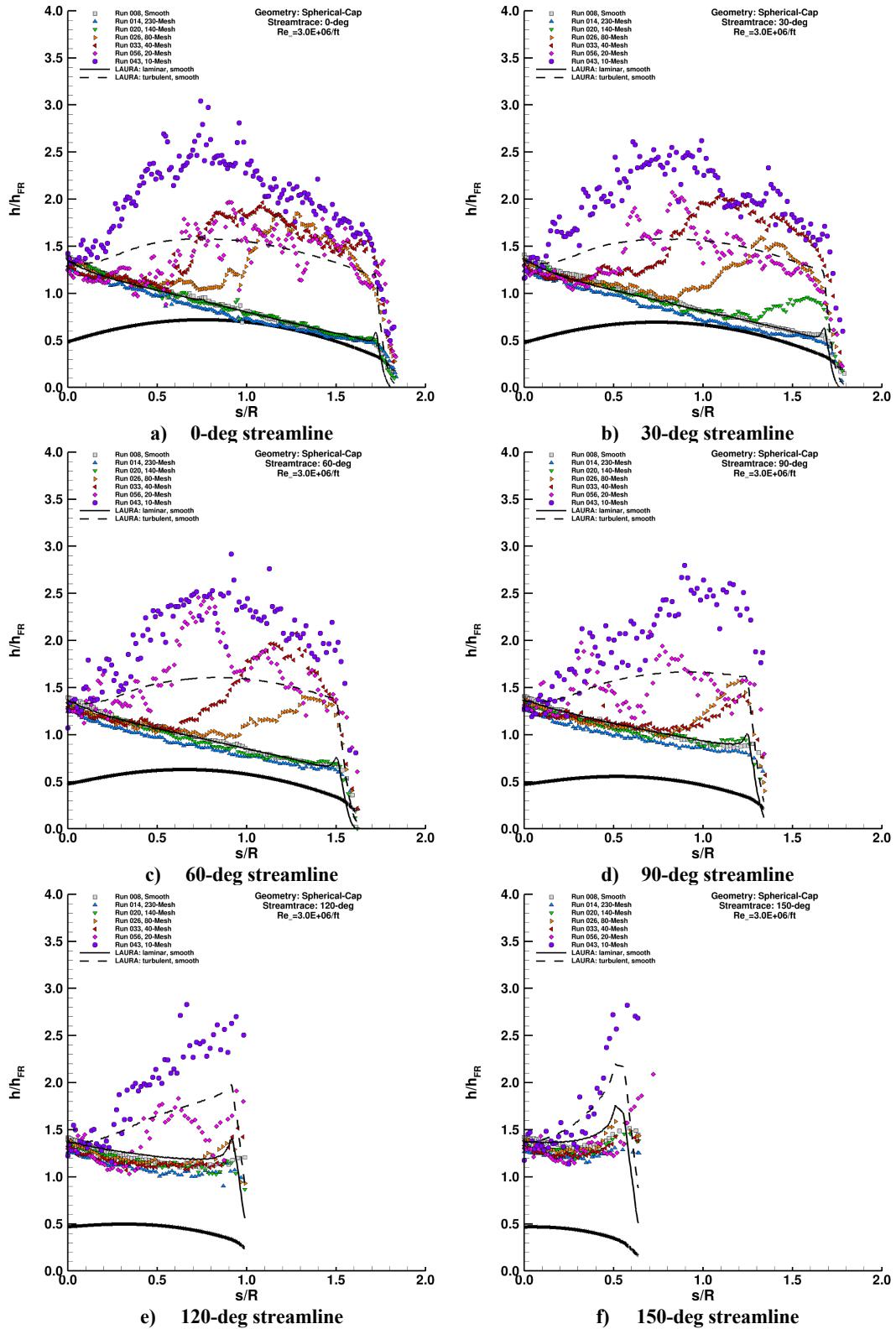


Figure 78. Roughness height effects, spherical-cap geometry, $Re_\infty = 3.0 \times 10^6/\text{ft}$ plots.

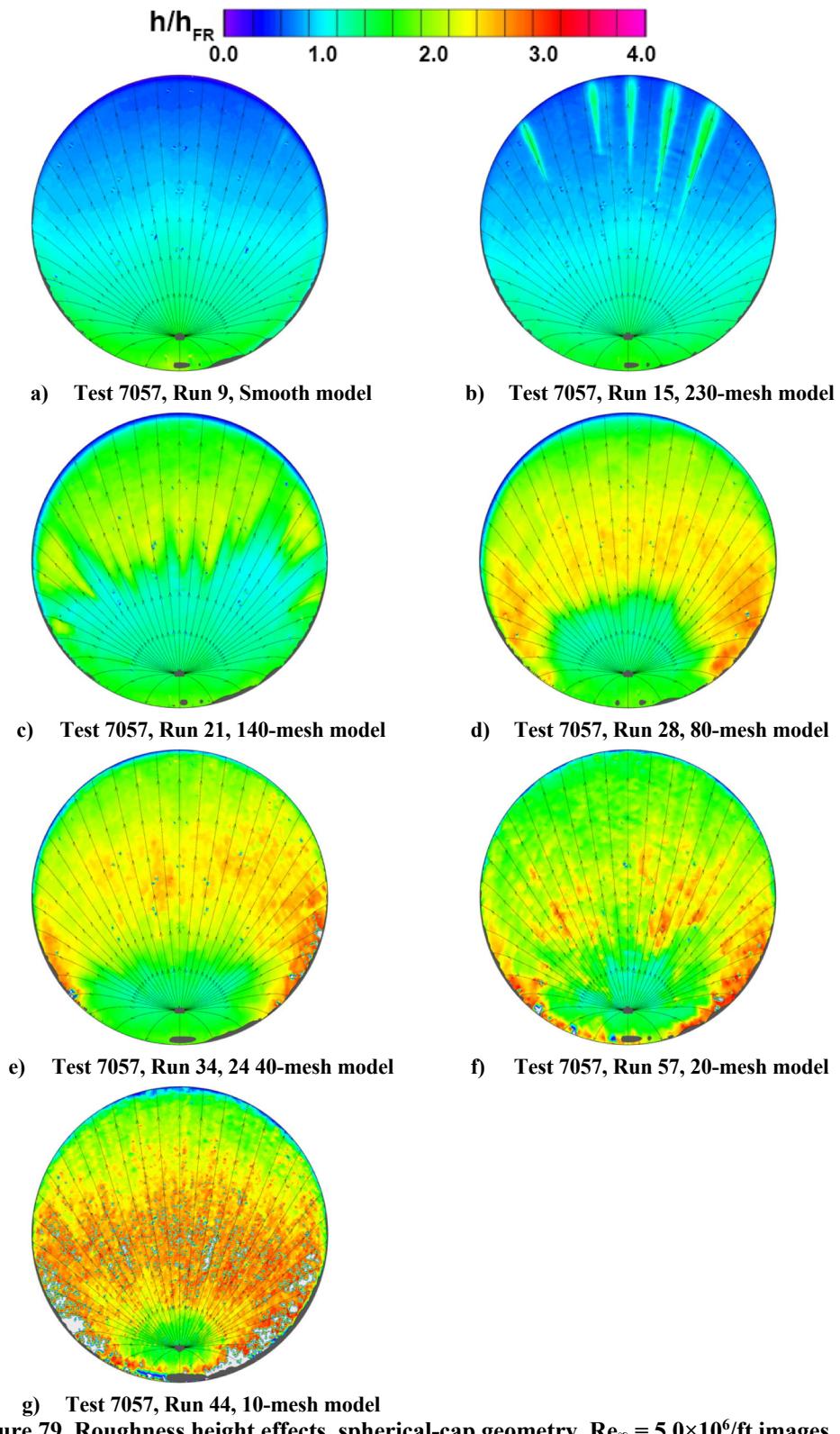


Figure 79. Roughness height effects, spherical-cap geometry, $Re_\infty = 5.0 \times 10^6/\text{ft}$ images.

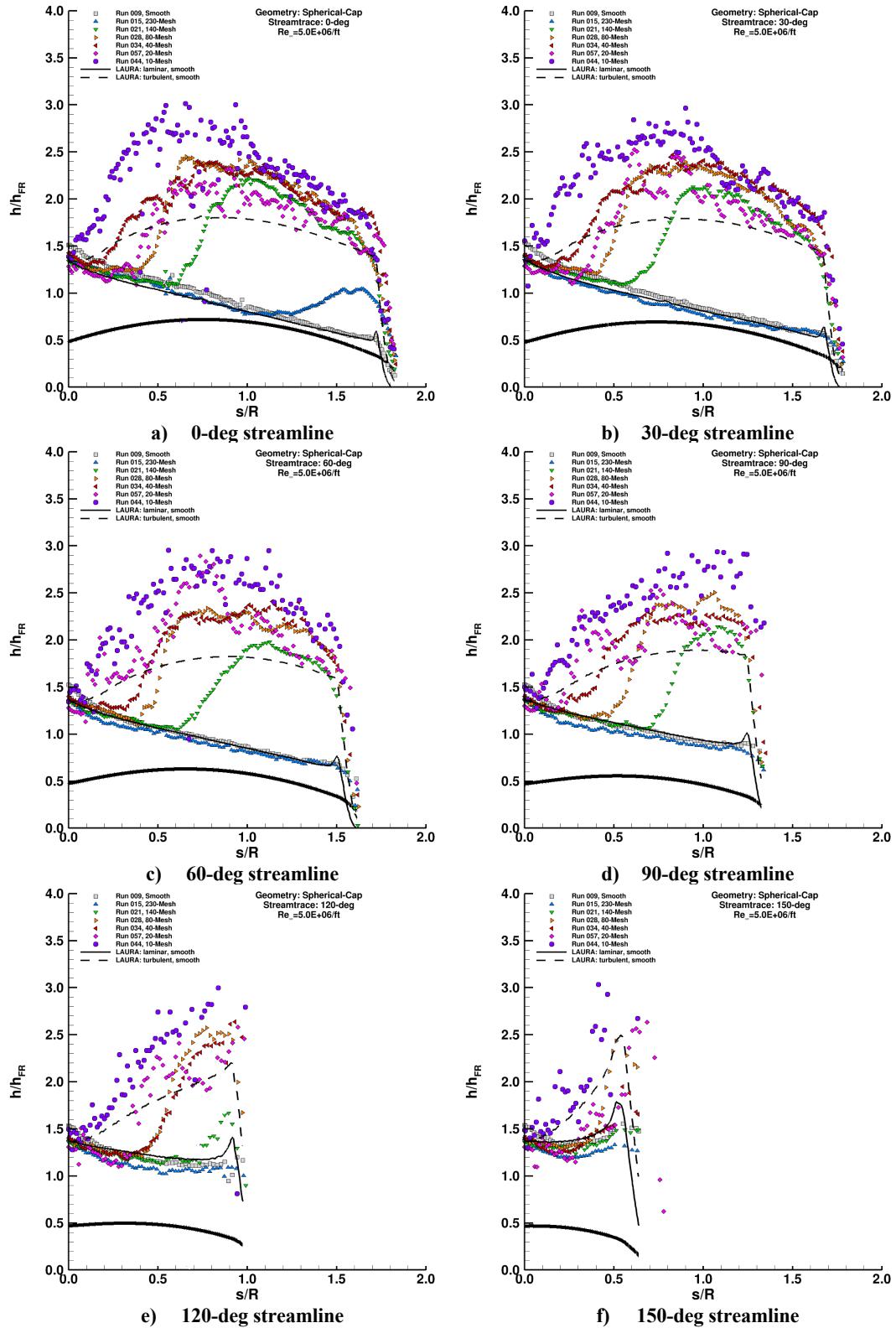


Figure 80. Roughness height effects, spherical-cap geometry, $Re_\infty = 5.0 \times 10^6/\text{ft}$ plots.

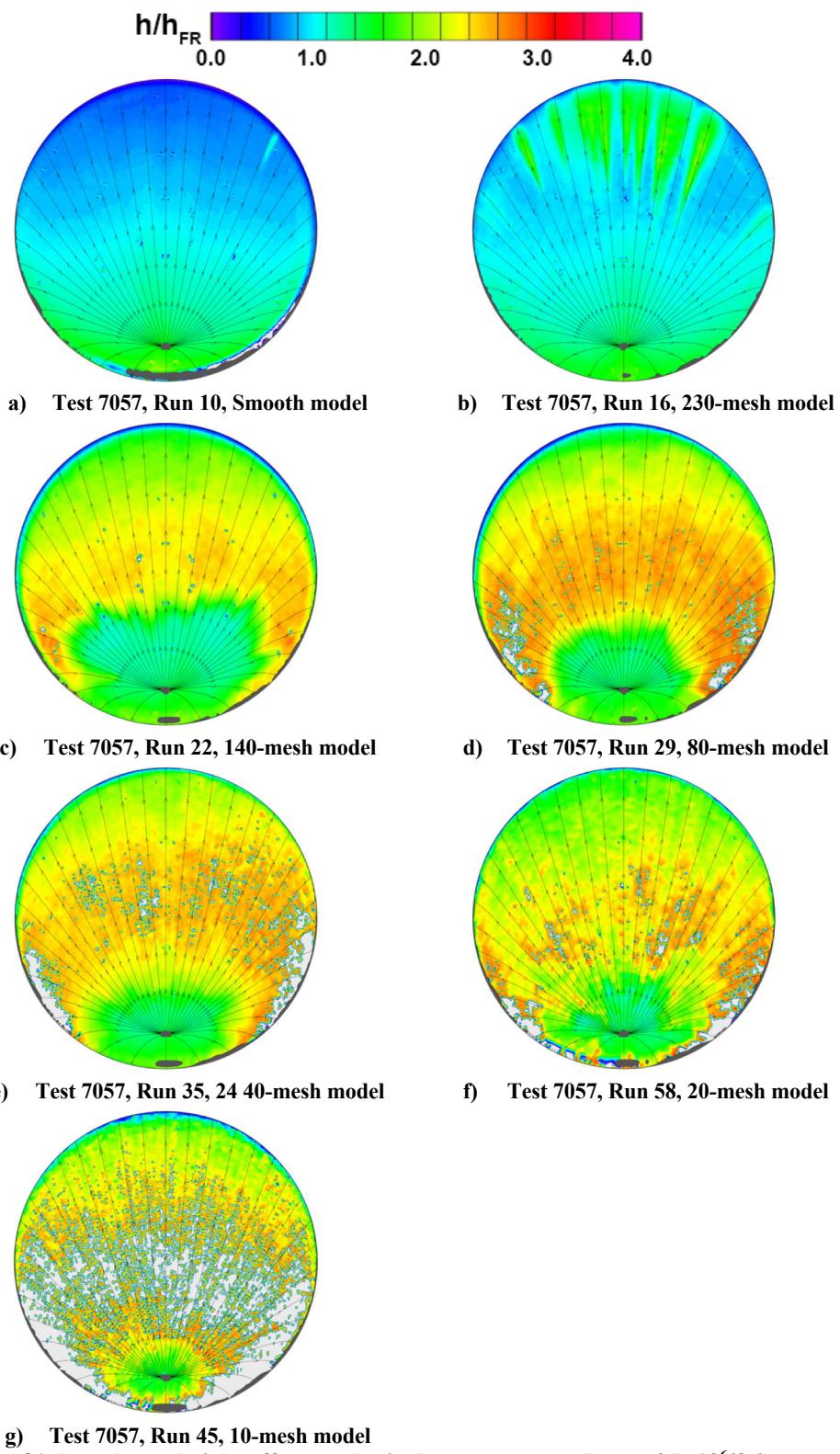


Figure 81. Roughness height effects, spherical-cap geometry, $Re_\infty = 6.5 \times 10^6/\text{ft}$ images.

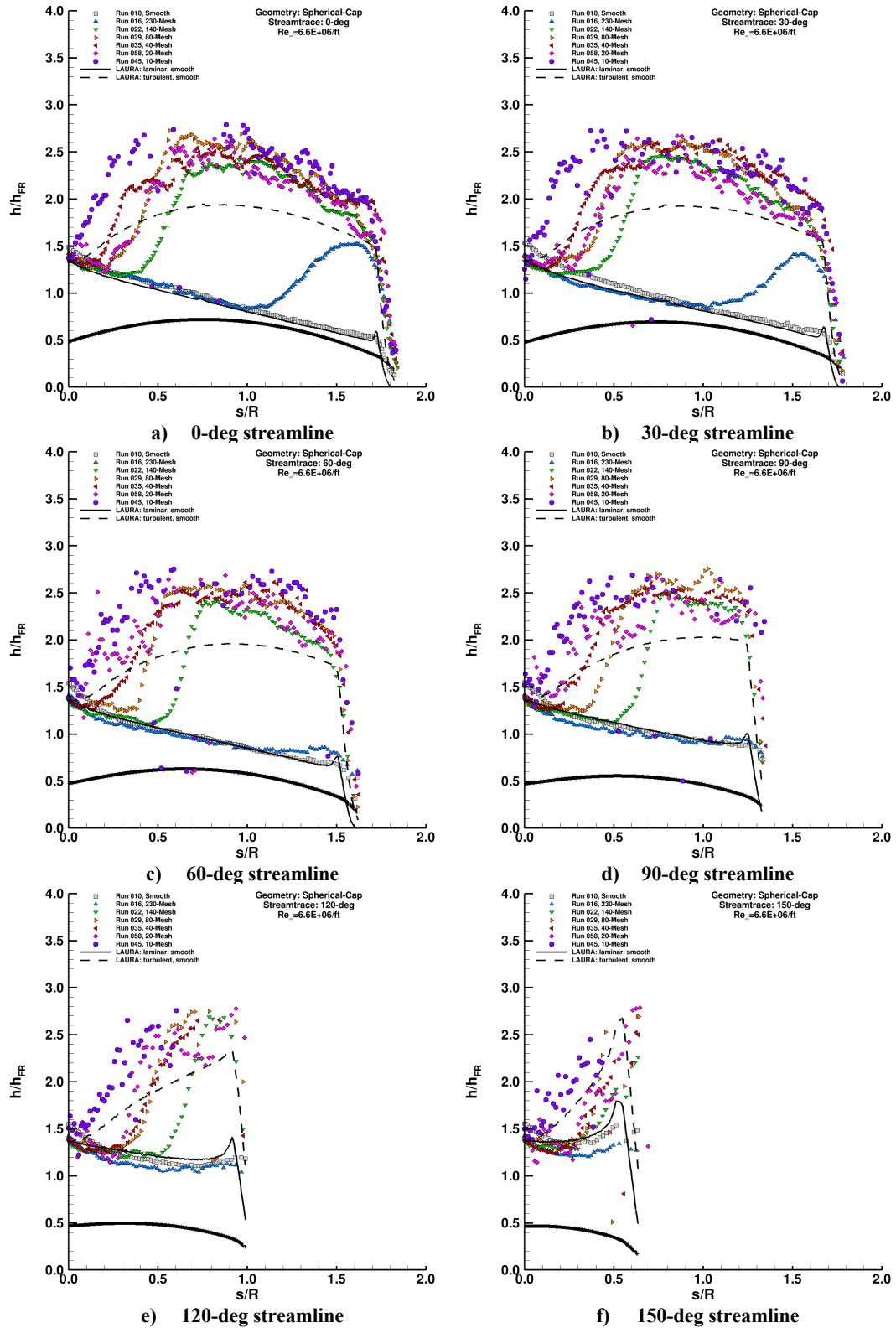


Figure 82. Roughness height effects, spherical-cap geometry, $Re_\infty = 6.5 \times 10^6/\text{ft}$ plots.

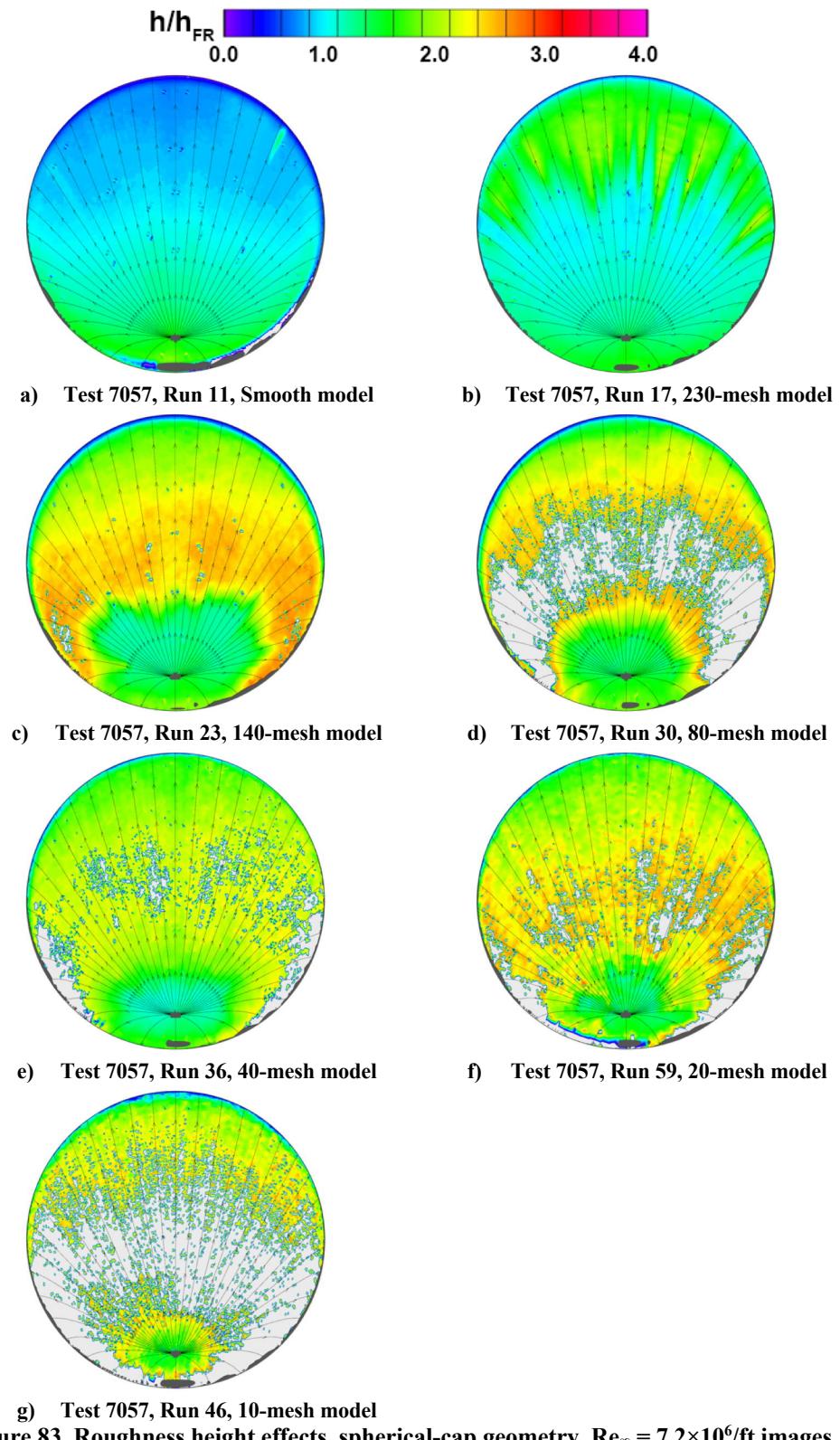


Figure 83. Roughness height effects, spherical-cap geometry, $Re_\infty = 7.2 \times 10^6/\text{ft}$ images.

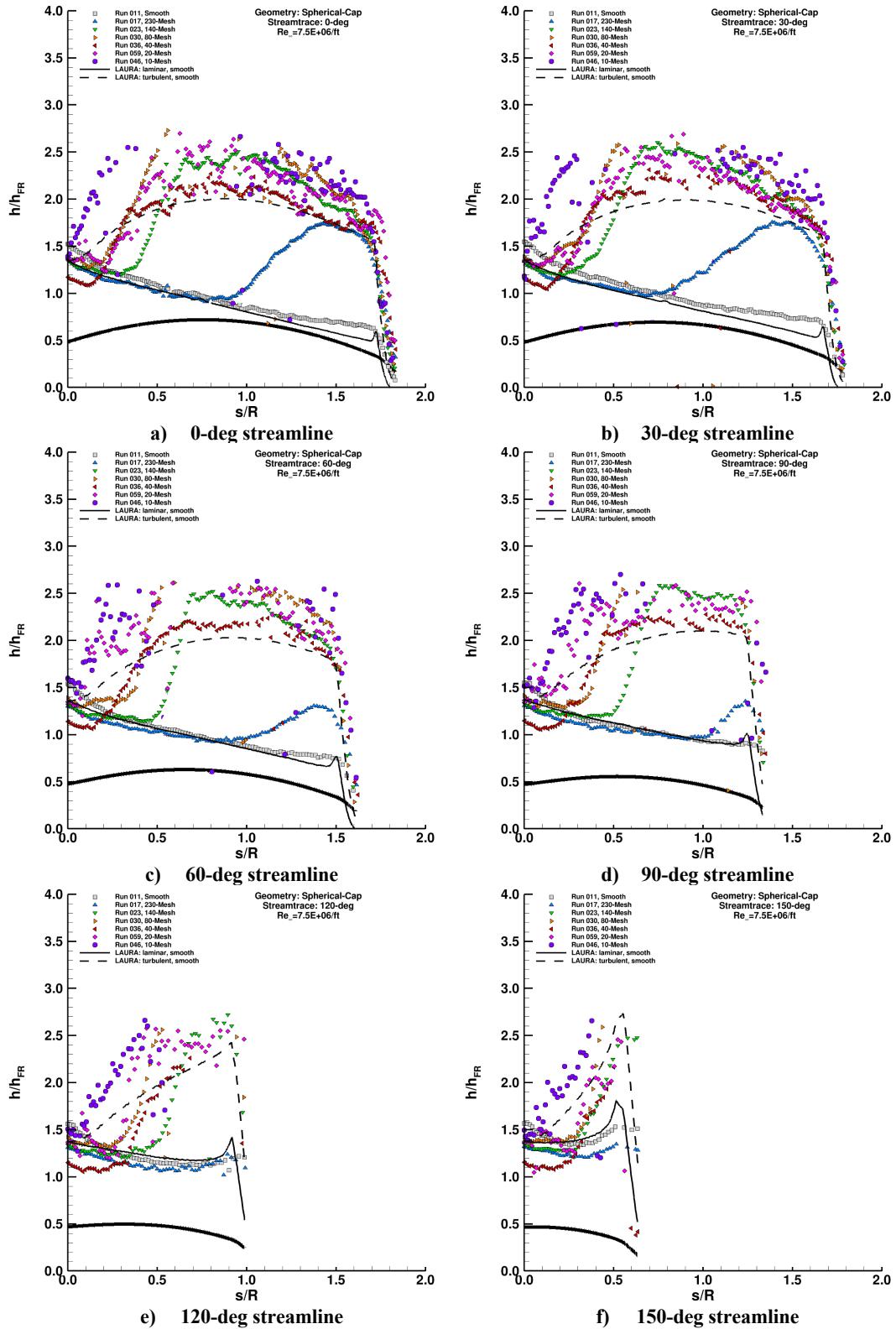


Figure 84. Roughness height effects, spherical-cap geometry, $Re_\infty = 7.2 \times 10^6$ ft plots.

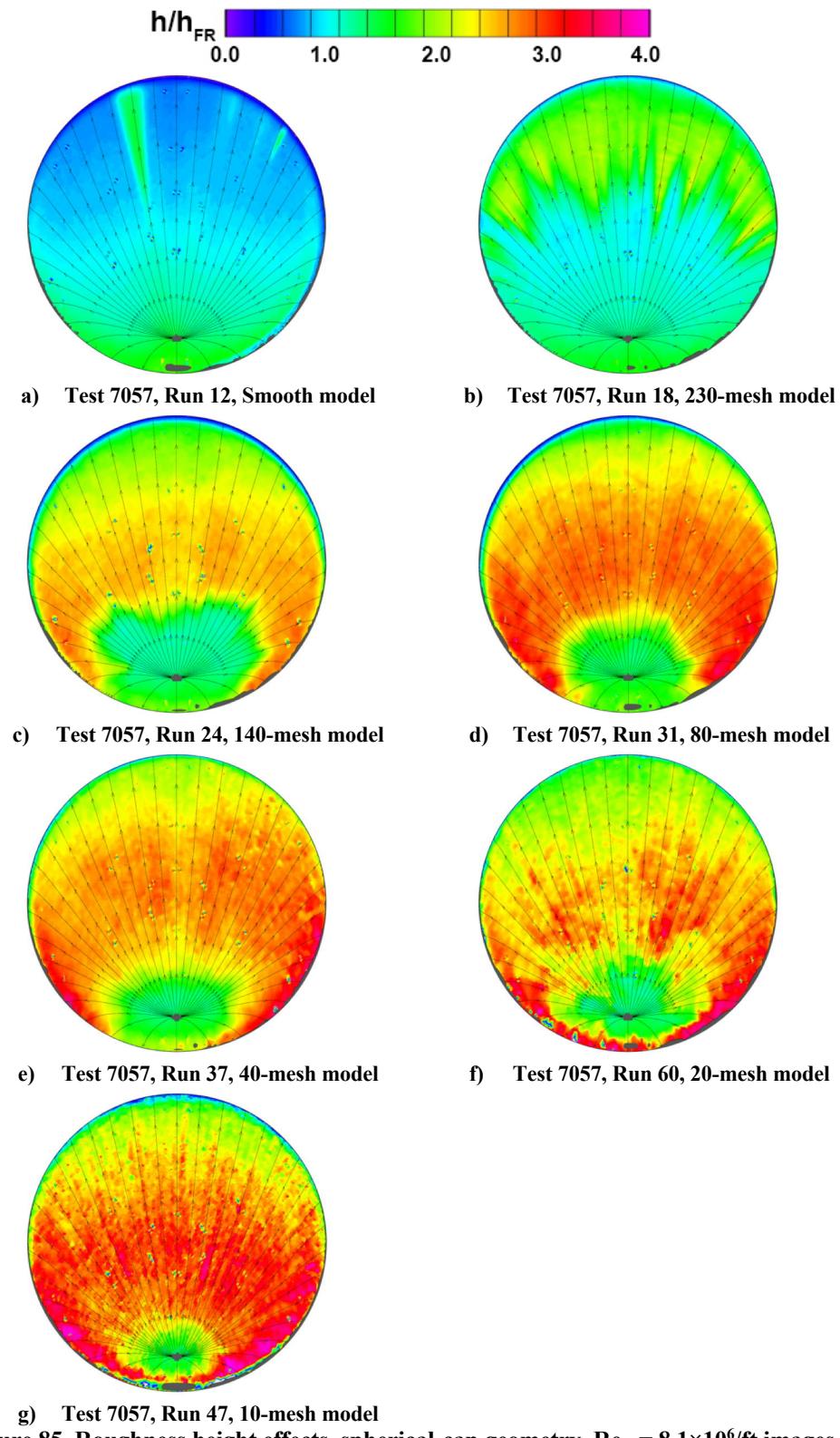


Figure 85. Roughness height effects, spherical-cap geometry, $Re_\infty = 8.1 \times 10^6/\text{ft}$ images.

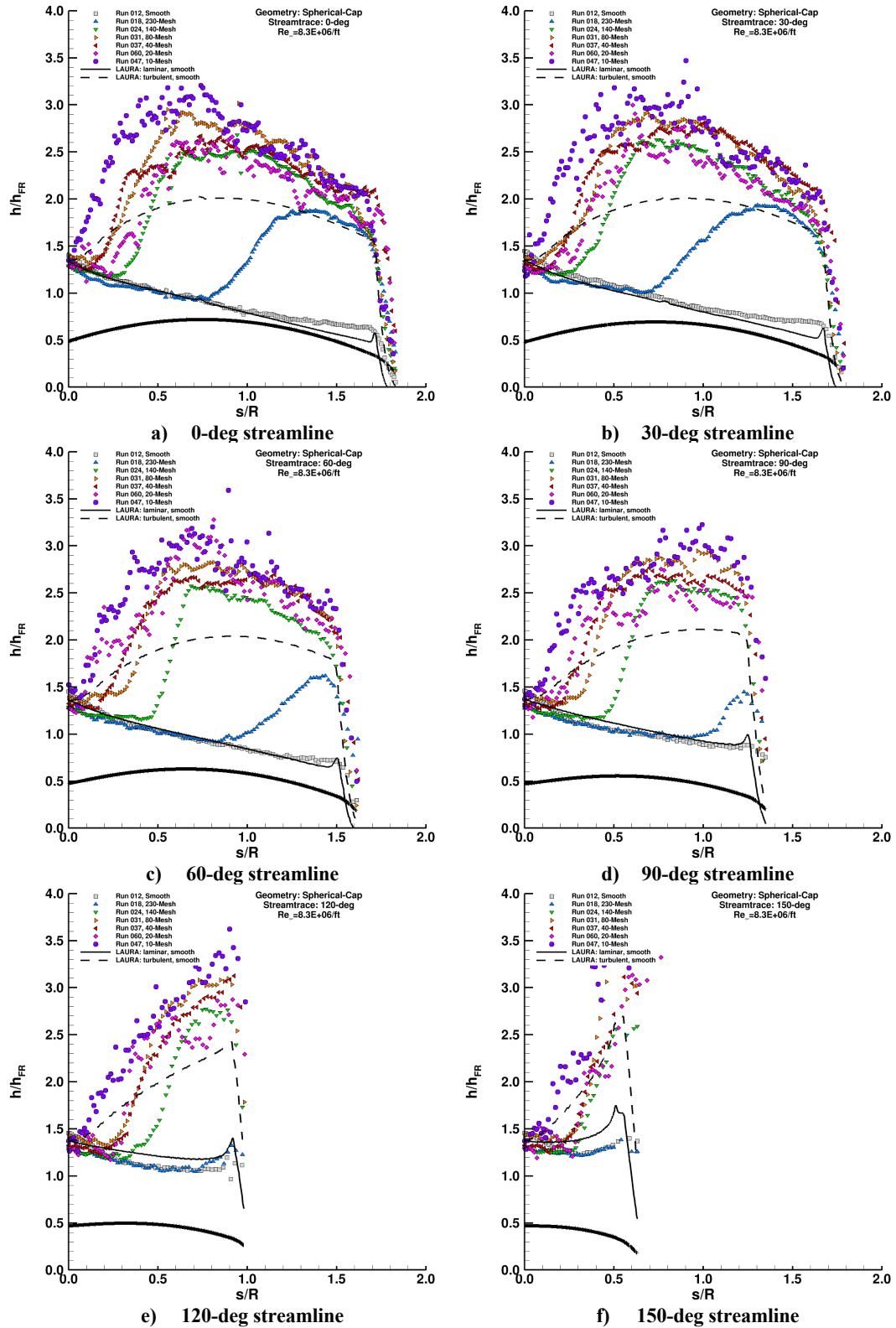
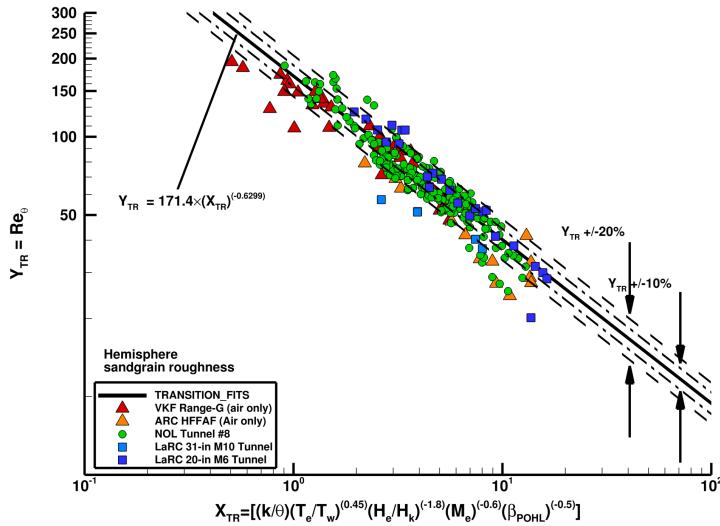
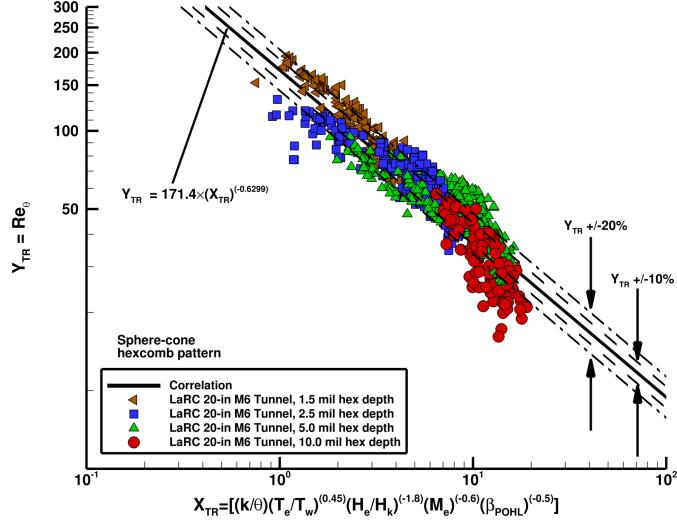


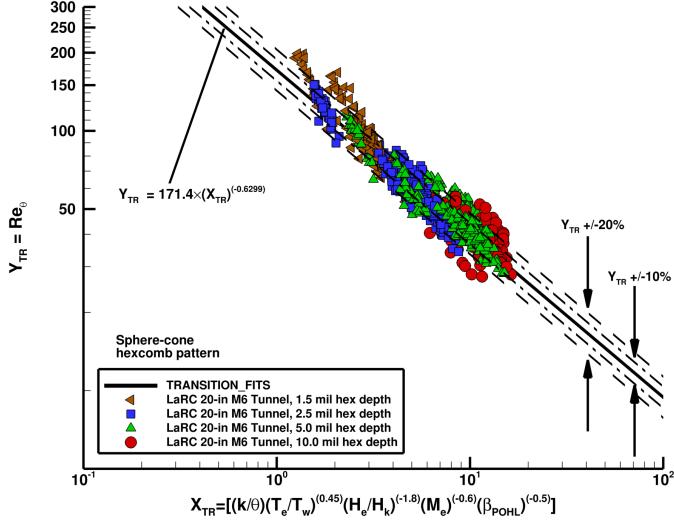
Figure 86. Roughness height effects, spherical-cap geometry, $Re_\infty = 8.1 \times 10^6 \text{ ft}$ plots.



a) Applied to hemisphere geometry with sand-grain roughness

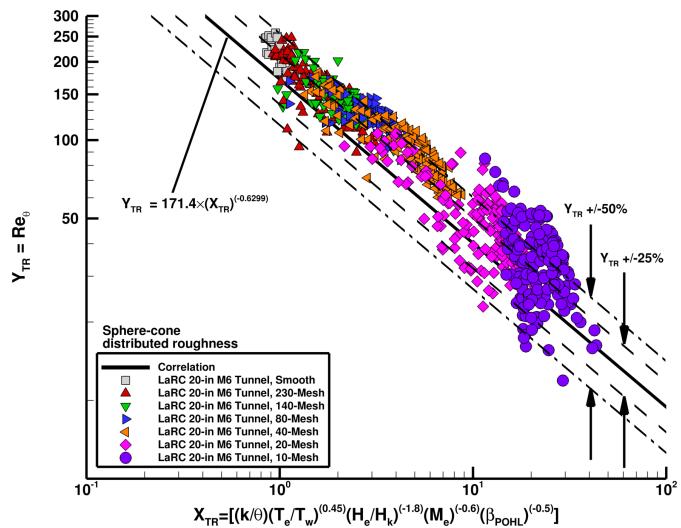


b) Applied to sphere-cone geometry with pattern-hexcomb roughness

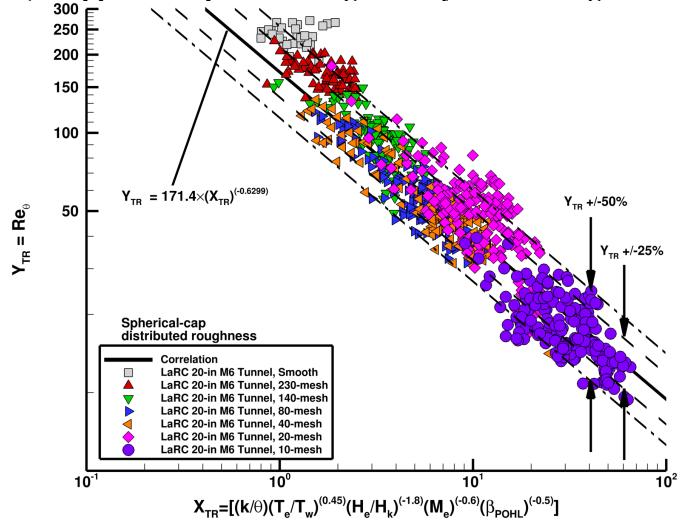


c) Applied to spherical-cap geometry with pattern-hexcomb roughness

Figure 87. Roughness transition correlation applied to prior datasets.



a) Applied to sphere-cone geometry with sand-grain roughness



b) Applied to spherical-cap geometry with sand-grain roughness

Figure 88. Roughness transition correlation applied to current datasets.

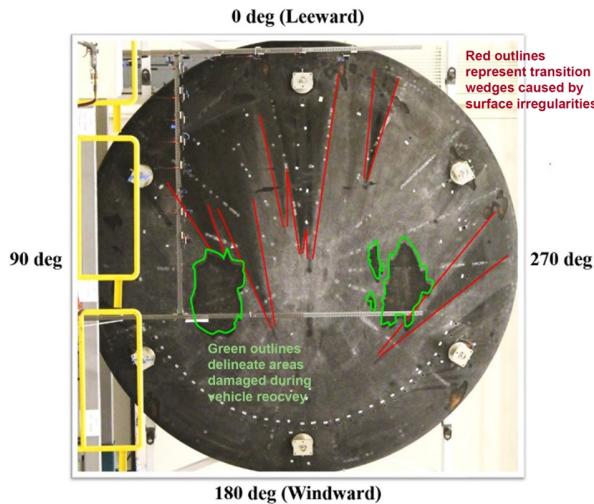


Figure 89. Post-flight recovery picture of Orion EFT-1 heatshield

Appendix A. Sphere-Cone Geometry Global Heating Images

Global heating images for the sphere-cone geometry from Test 7036 in the LAL 20-Inch Mach 6 Air Tunnel are presented in this Appendix in Figure 90 through Figure 131.

At higher Reynolds numbers and/or larger roughness heights, white patches on the images indicate areas where the measured surface temperatures exceed the calibrated range of the phosphor thermography and thus no valid data were obtained.

Boundary-layer edge streamlines determined from laminar, smooth-surface LAURA simulations have been superimposed on the images to illustrate the nature of the flow field.

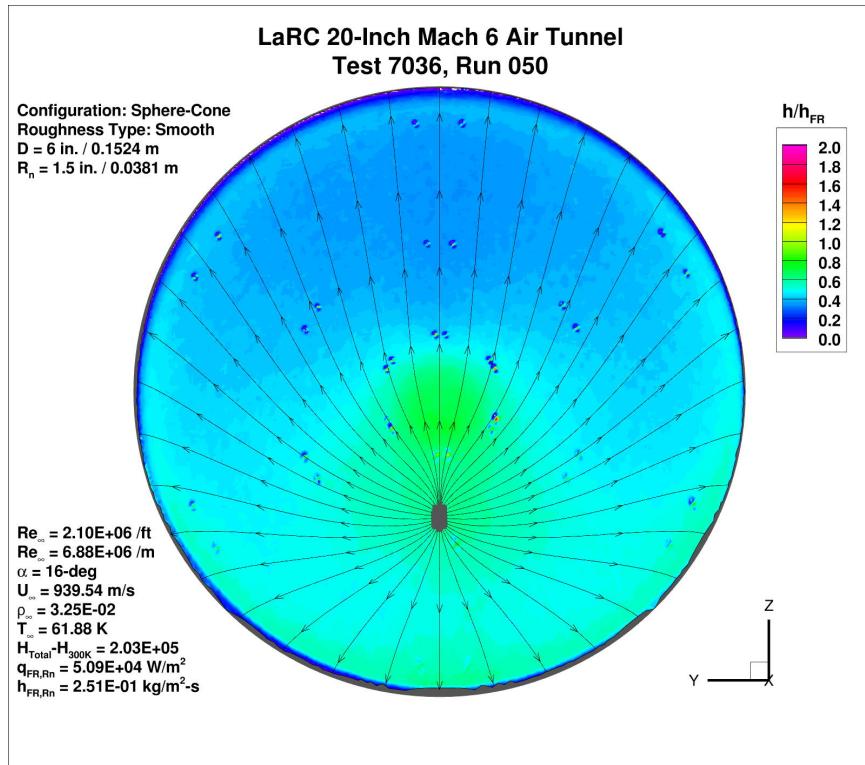


Figure 90. Test 7036, Run 50, $Re_{\infty} = 2.1 \times 10^6$ /ft, sphere-cone, smooth OML.

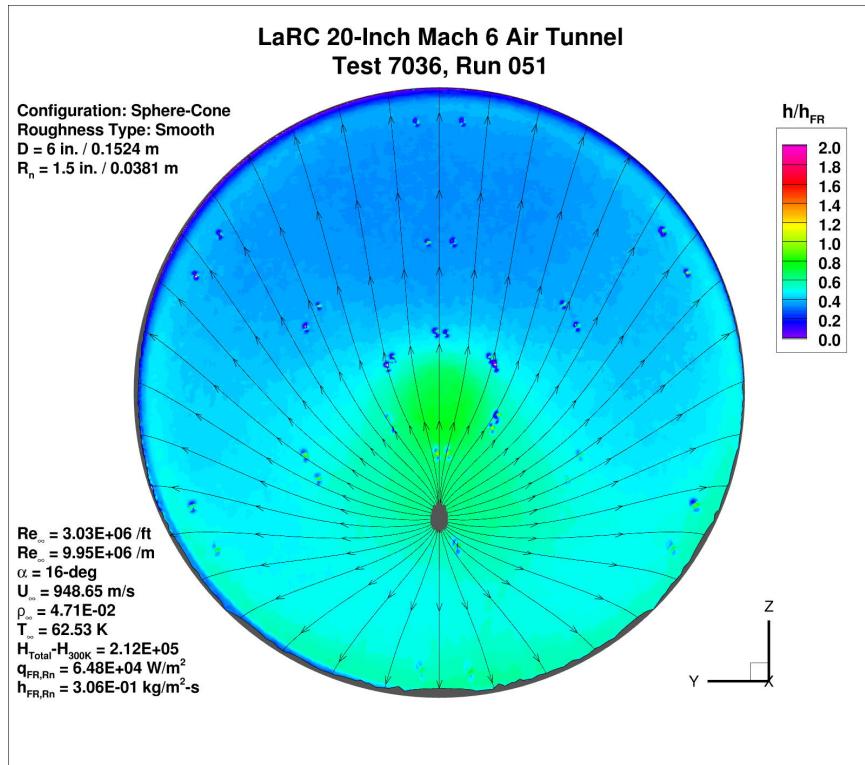


Figure 91. Test 7036, Run 51, $Re_{\infty} = 3.0 \times 10^6$ /ft, sphere-cone, smooth OML.

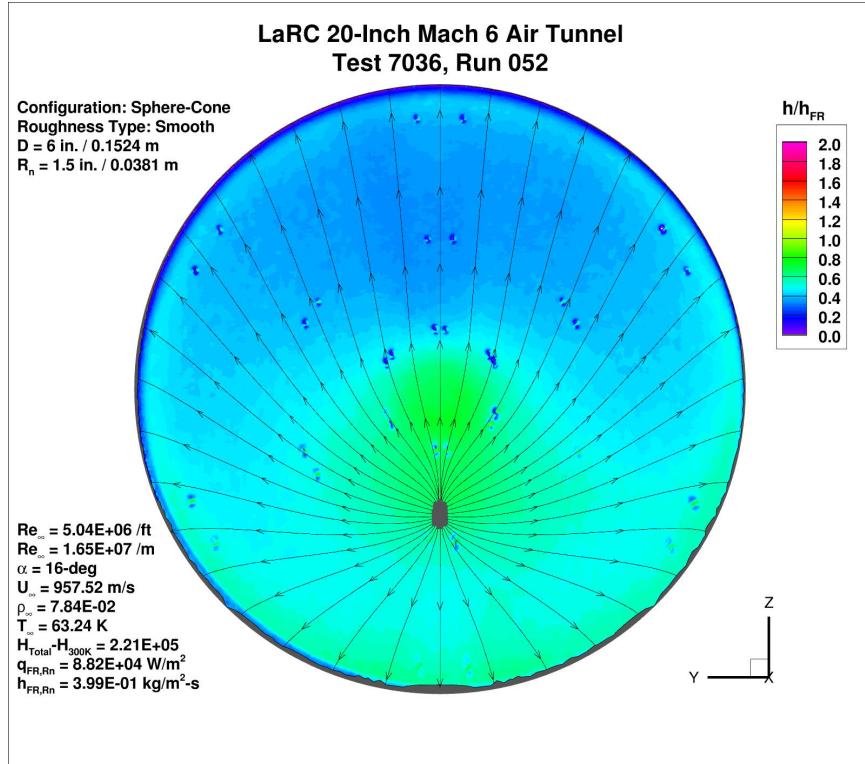


Figure 92. Test 7036, Run 52, $Re_{\infty} = 5.0 \times 10^6$ /ft, sphere-cone, smooth OML.

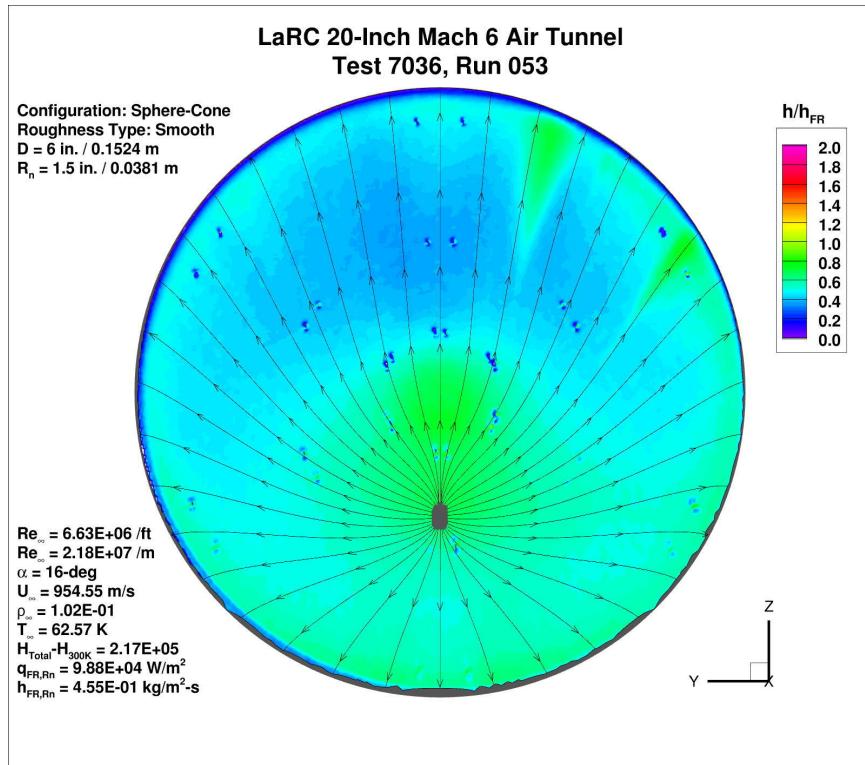


Figure 93. Test 7036, Run 53, $Re_{\infty} = 6.5 \times 10^6$ /ft, sphere-cone, smooth OML.

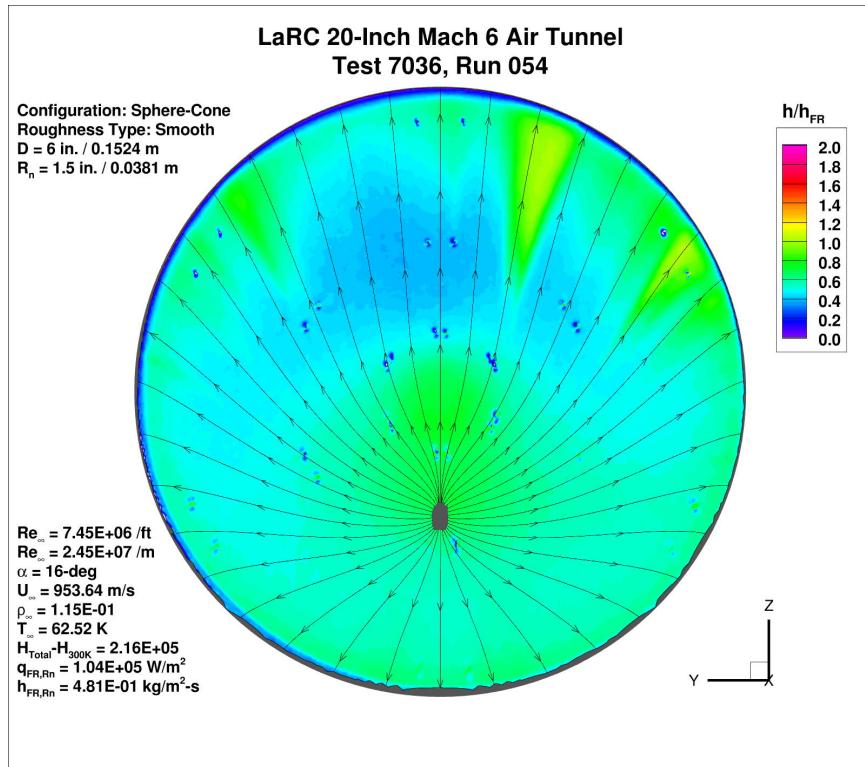


Figure 94. Test 7036, Run 54, $Re_{\infty} = 7.2 \times 10^6$ /ft, sphere-cone, smooth OML.

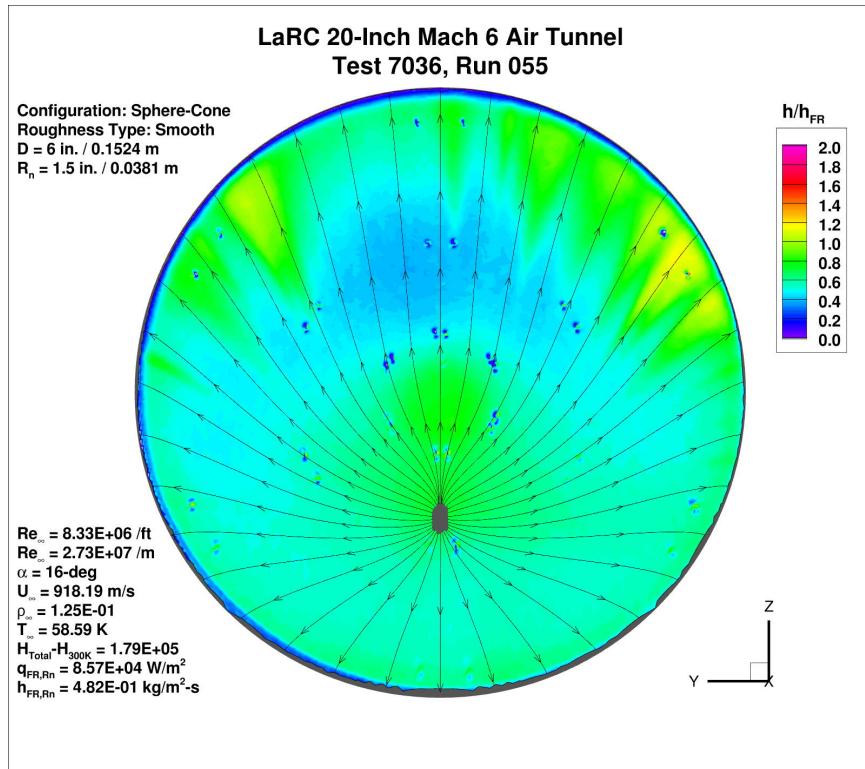


Figure 95. Test 7036, Run 55, $Re_{\infty} = 8.1 \times 10^6$ /ft, sphere-cone, smooth OML.

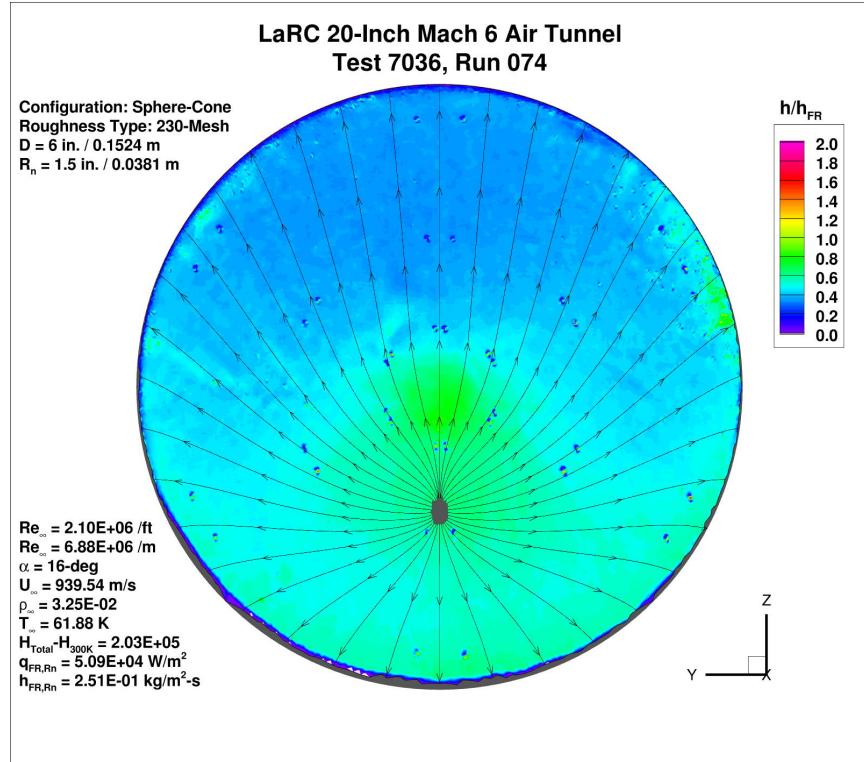


Figure 96. Test 7036, Run 74, $Re_{\infty} = 2.1 \times 10^6/\text{ft}$, sphere-cone, 230-mesh.

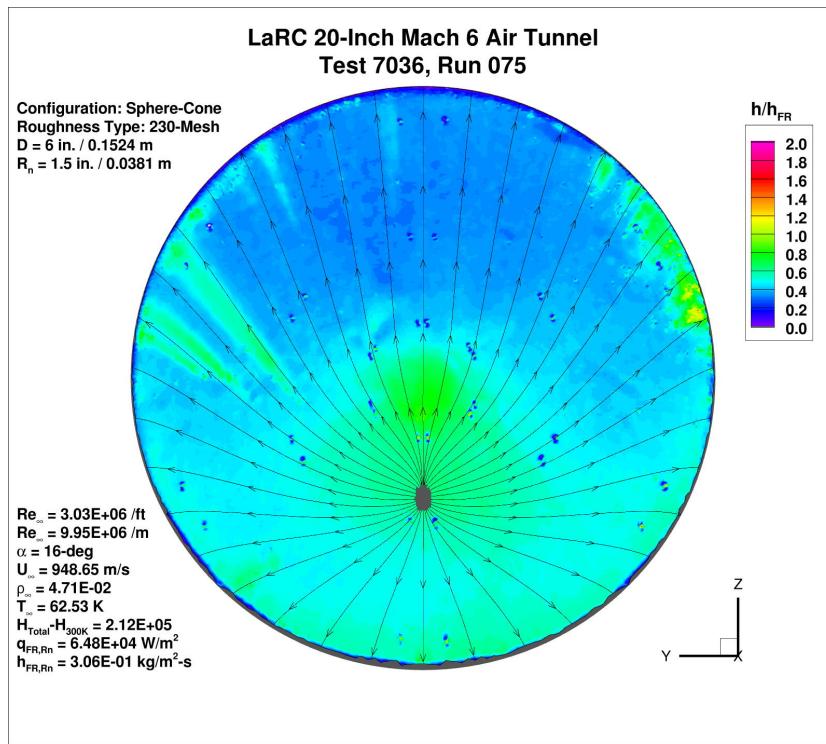


Figure 97. Test 7036, Run 75, $Re_{\infty} = 3.0 \times 10^6/\text{ft}$, sphere-cone, 230-mesh.

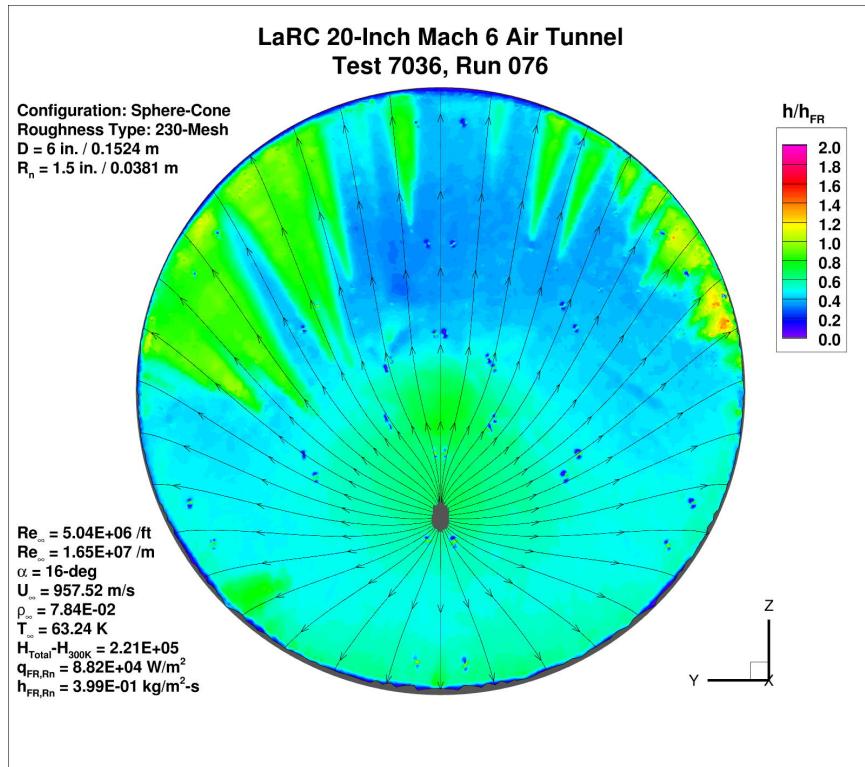


Figure 98. Test 7036, Run 76, $Re_{\infty} = 5.0 \times 10^6/\text{ft}$, sphere-cone, 230-mesh.

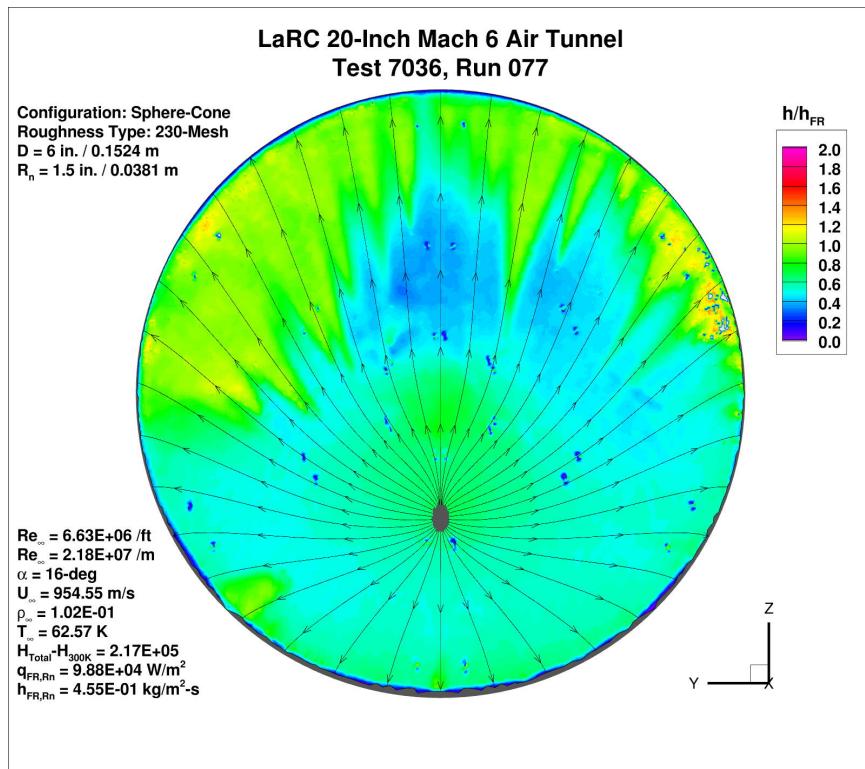


Figure 99. Test 7036, Run 77, $Re_{\infty} = 6.5 \times 10^6/\text{ft}$, sphere-cone, 230-mesh.

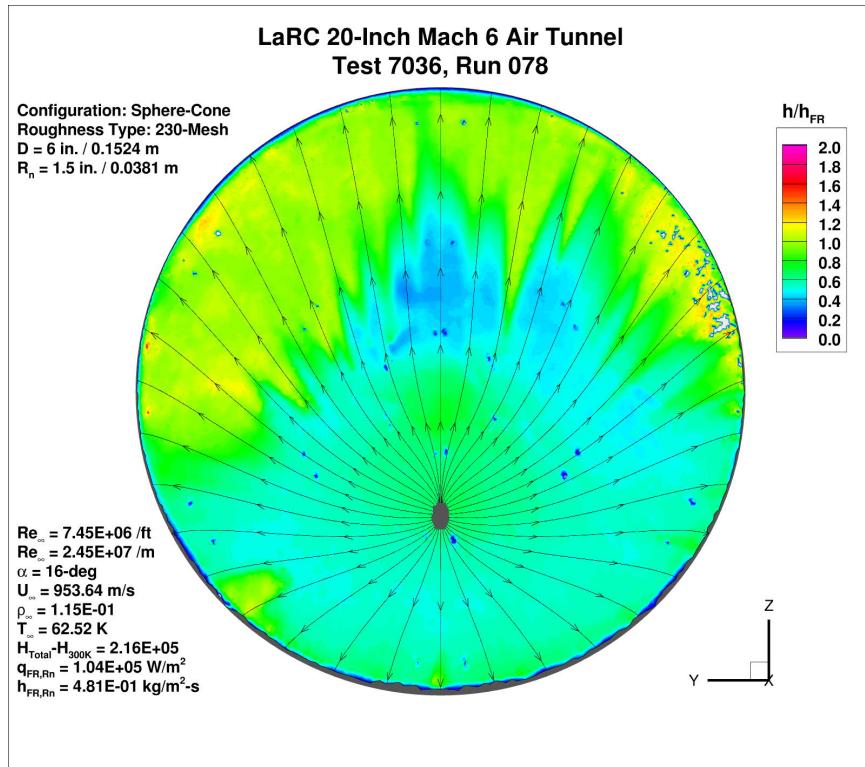


Figure 100. Test 7036, Run 78, $Re_{\infty} = 7.2 \times 10^6$ /ft, sphere-cone, 230-mesh.

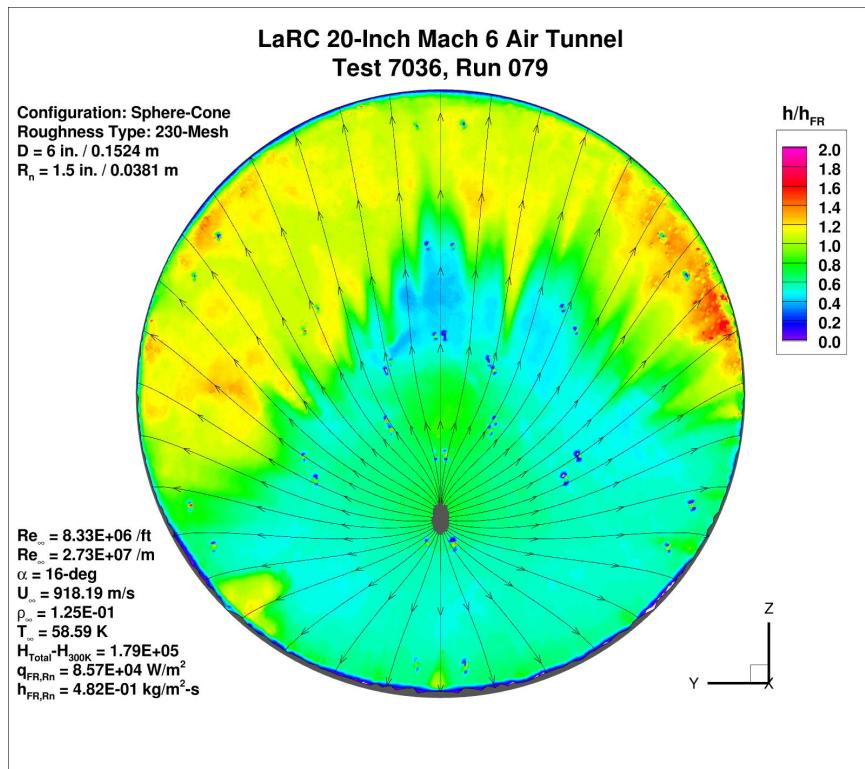


Figure 101. Test 7036, Run 79, $Re_{\infty} = 8.1 \times 10^6$ /ft, sphere-cone, 230-mesh.

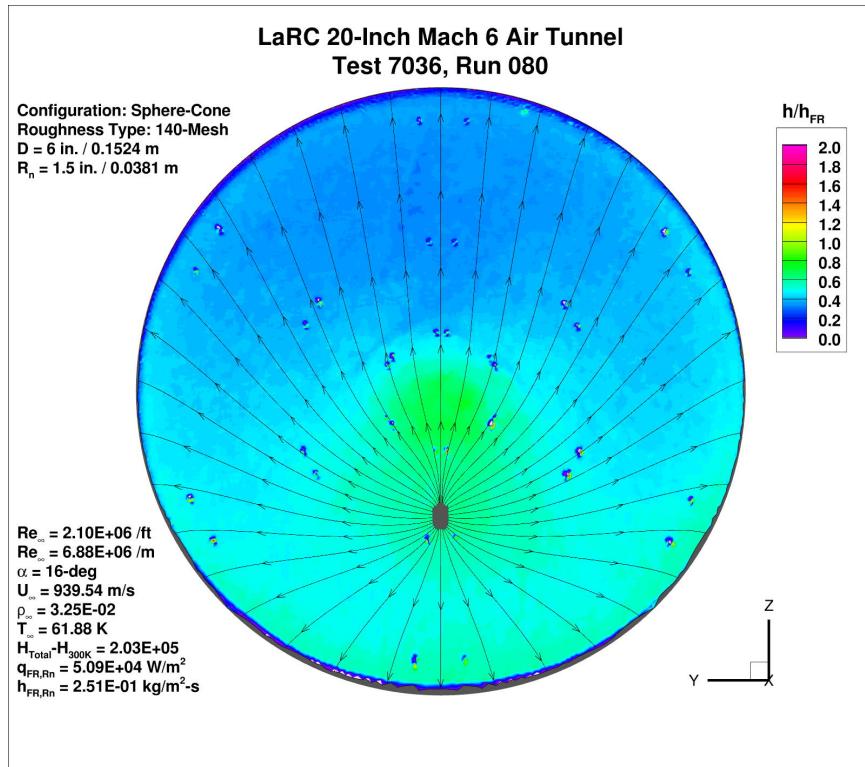


Figure 102. Test 7036, Run 80, $Re_{\infty} = 2.1 \times 10^6/\text{ft}$, sphere-cone, 140-mesh.

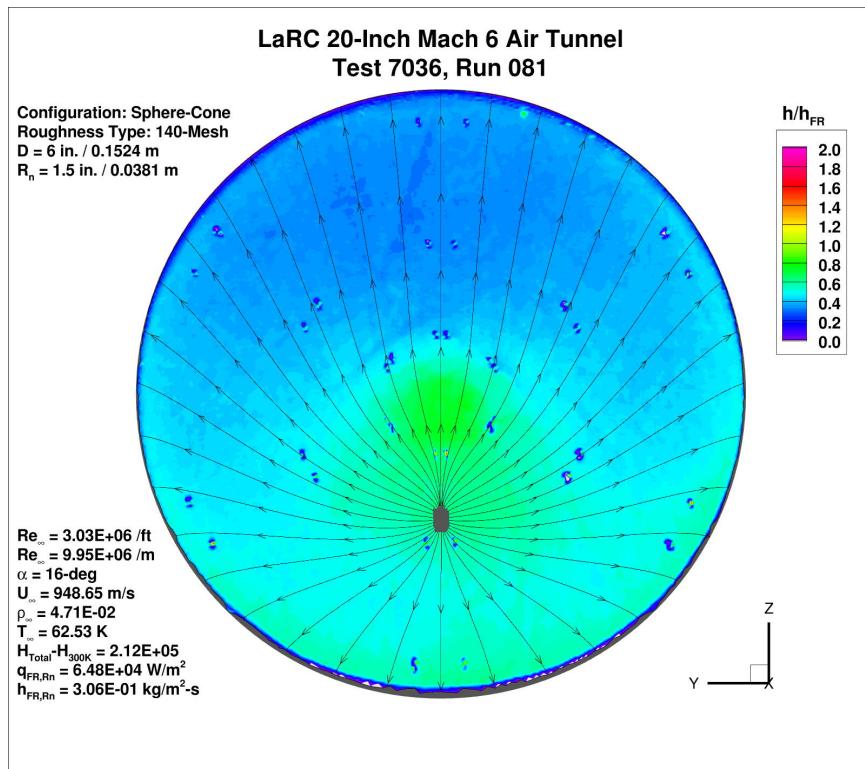


Figure 103. Test 7036, Run 81, $Re_{\infty} = 3.0 \times 10^6/\text{ft}$, sphere-cone, 140-mesh.

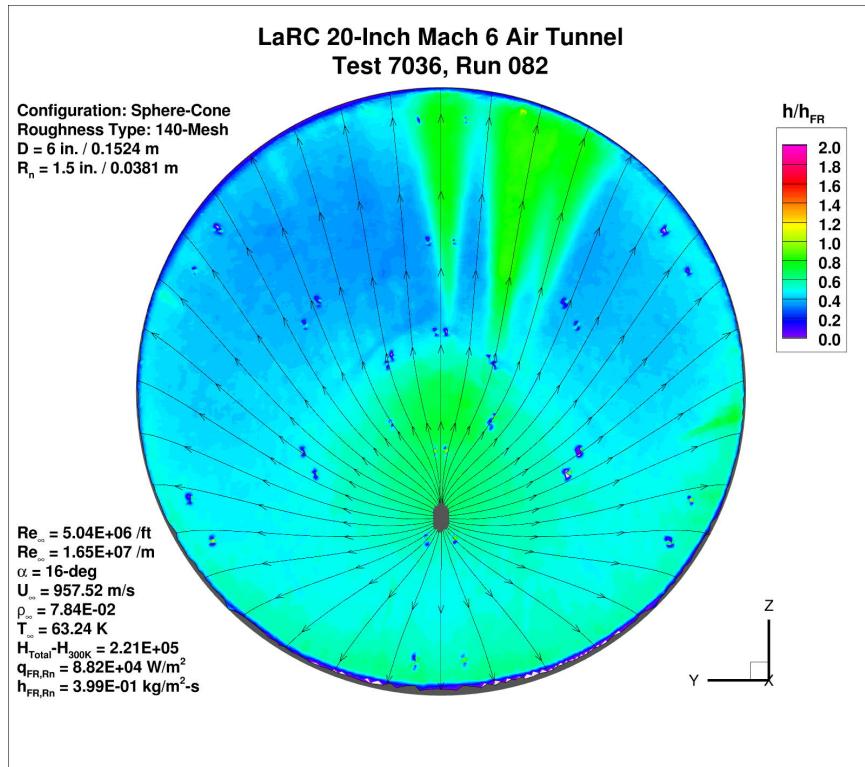


Figure 104. Test 7036, Run 82, $Re_{\infty} = 5.0 \times 10^6$ /ft, sphere-cone, 140-mesh.

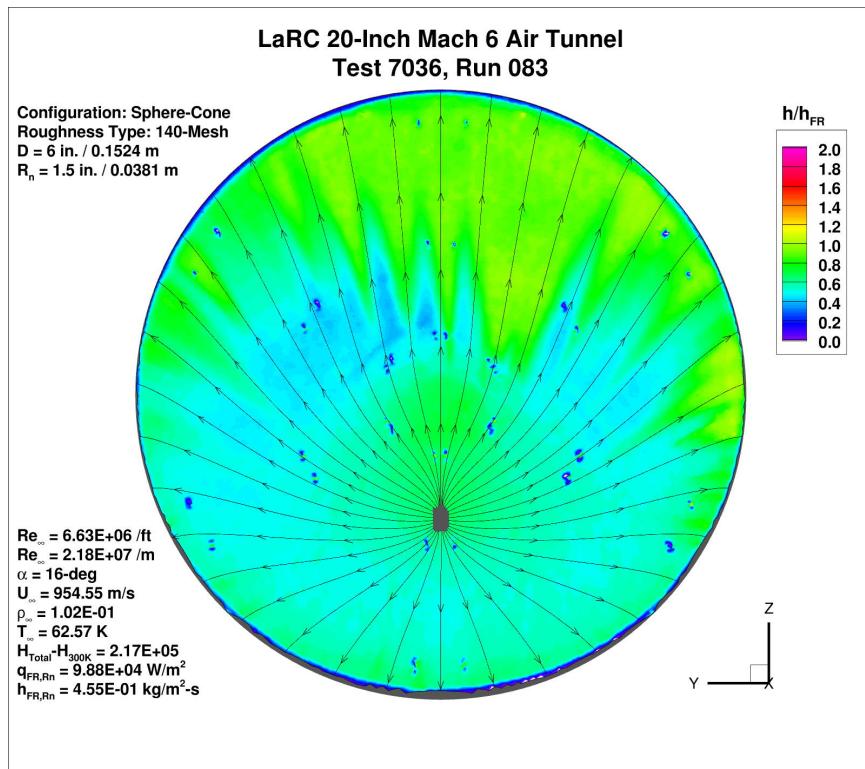


Figure 105. Test 7036, Run 83, $Re_{\infty} = 6.5 \times 10^6$ /ft, sphere-cone, 140-mesh.

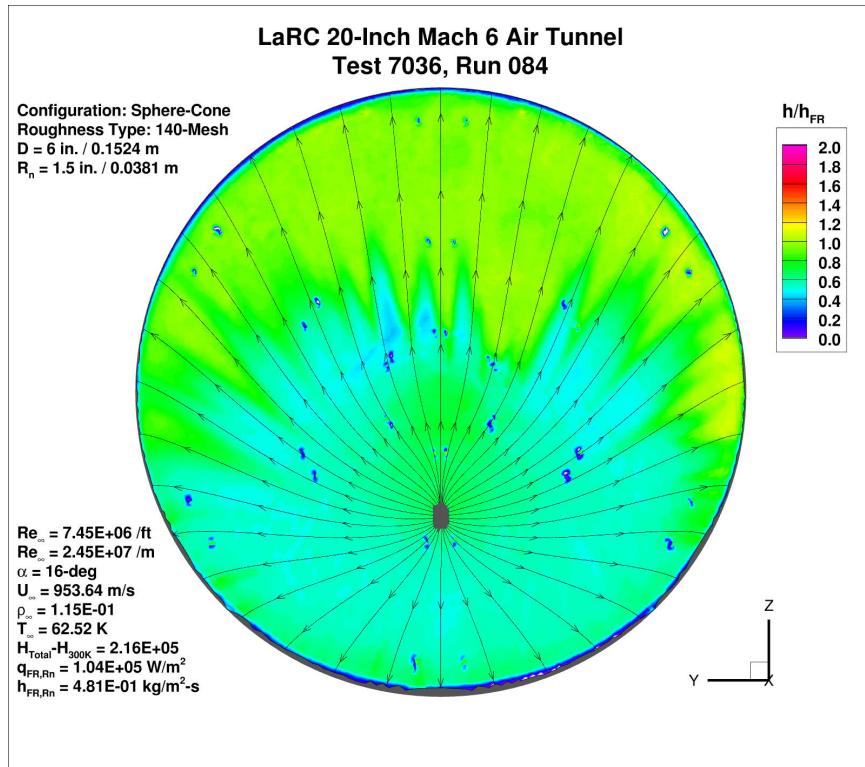


Figure 106. Test 7036, Run 84, $Re_{\infty} = 7.2 \times 10^6$ /ft, sphere-cone, 140-mesh.

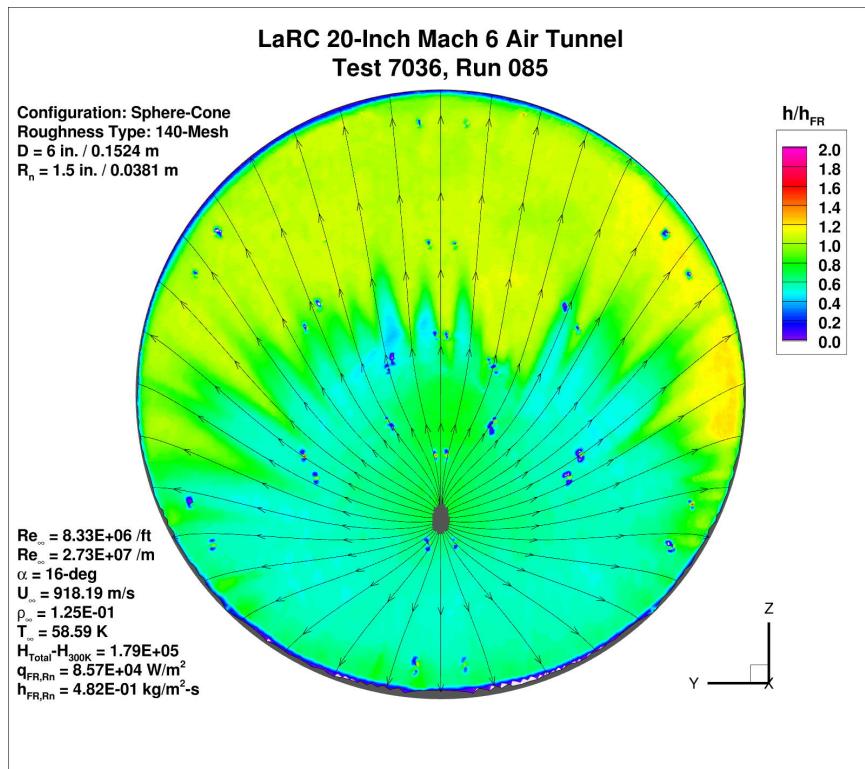


Figure 107. Test 7036, Run 85, $Re_{\infty} = 8.1 \times 10^6$ /ft, sphere-cone, 140-mesh.

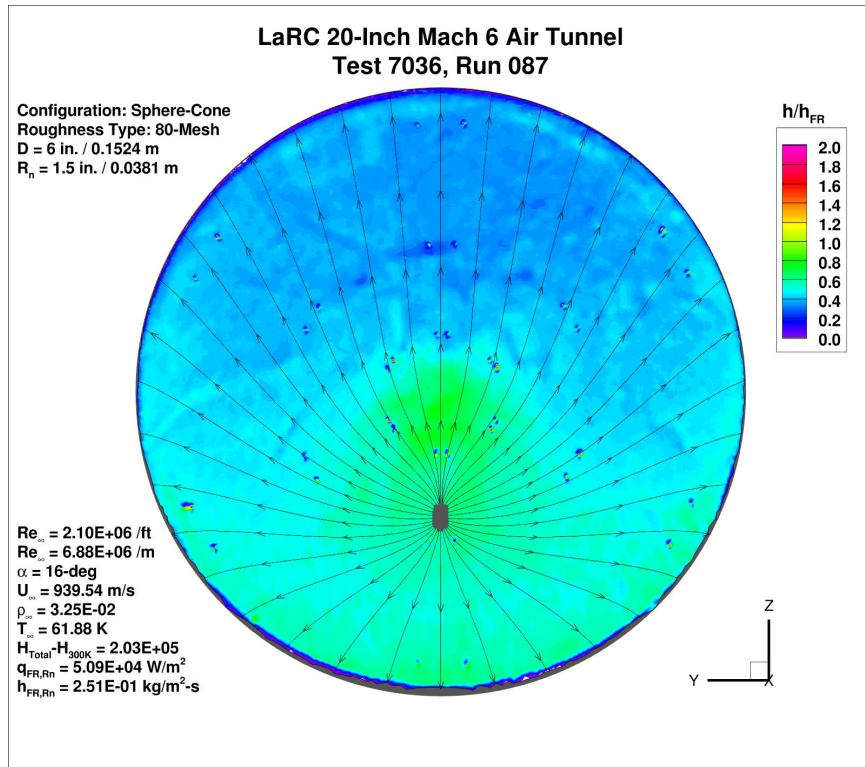


Figure 108. Test 7036, Run 87, $Re_{\infty} = 2.1 \times 10^6/\text{ft}$, sphere-cone, 80-mesh.

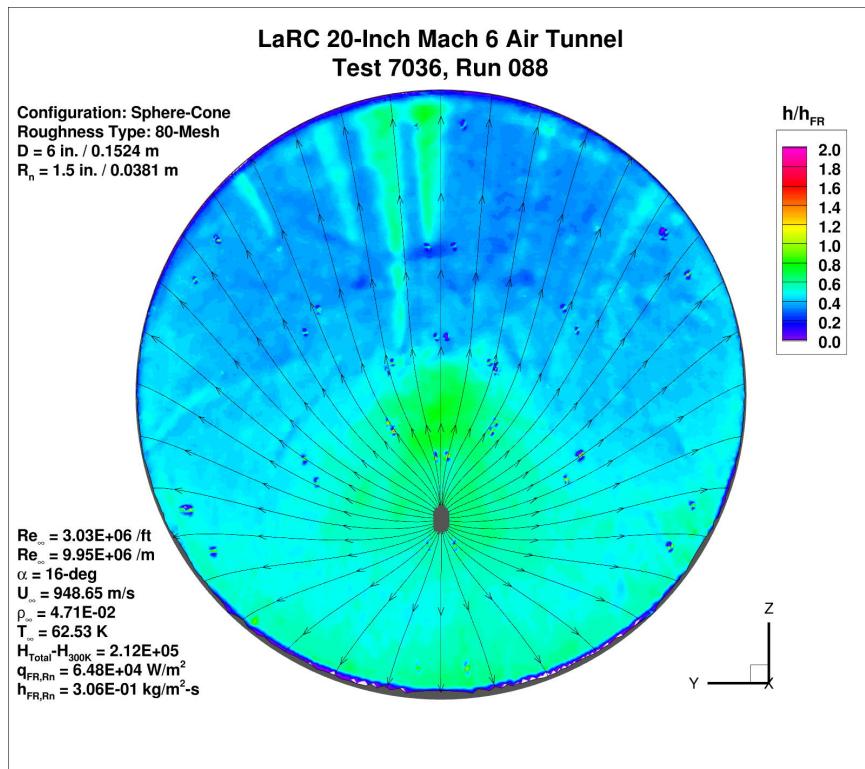


Figure 109. Test 7036, Run 88, $Re_{\infty} = 3.0 \times 10^6/\text{ft}$, sphere-cone, 80-mesh.

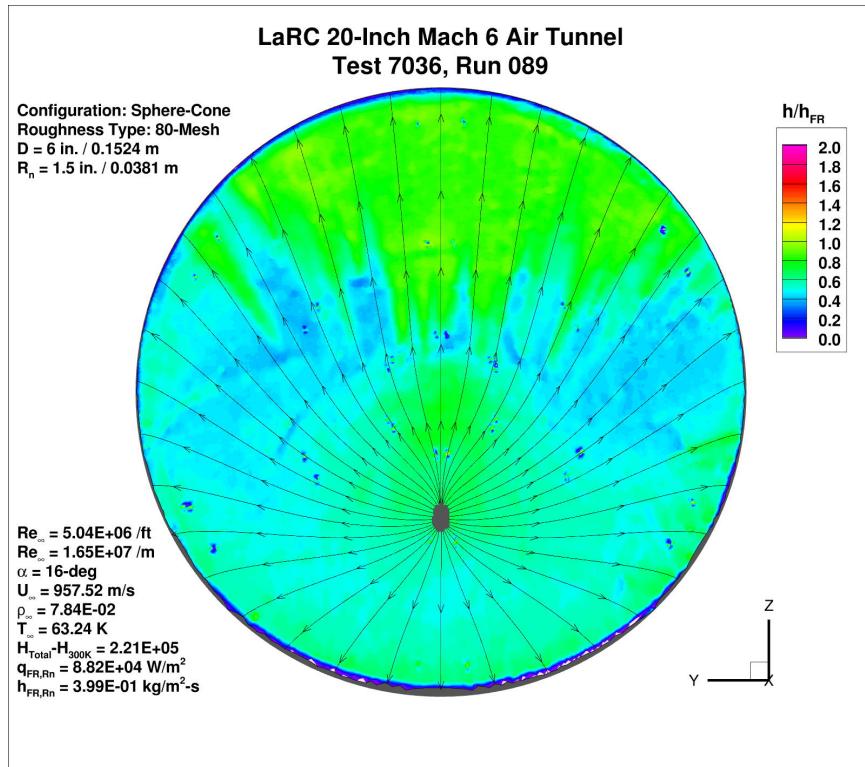


Figure 110. Test 7036, Run 89, $Re_{\infty} = 5.0 \times 10^6 / \text{ft}$, sphere-cone, 80-mesh.

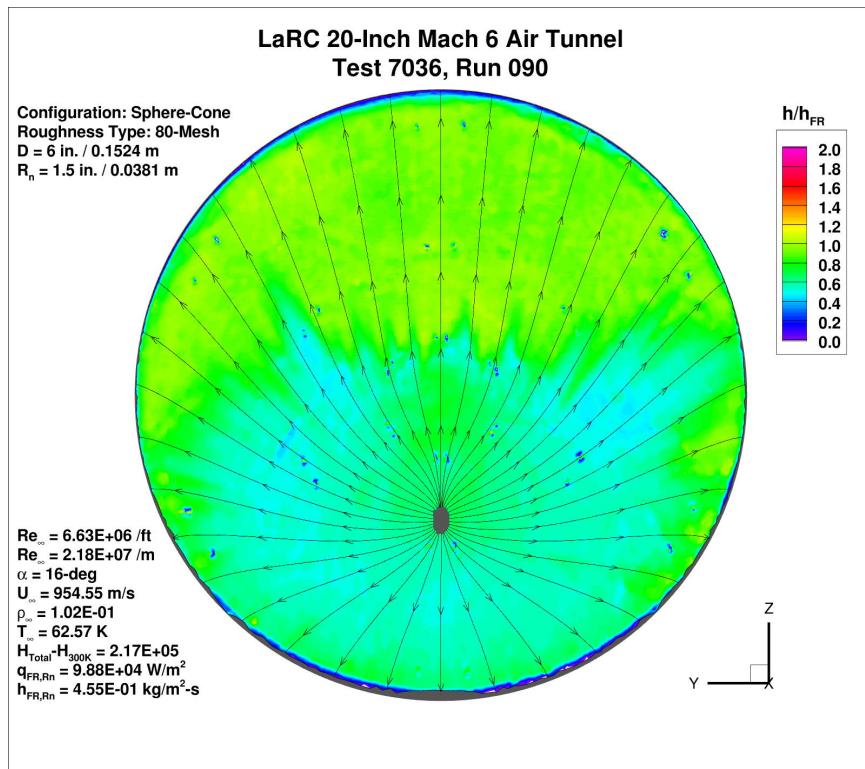


Figure 111. Test 7036, Run 90, $Re_{\infty} = 6.5 \times 10^6 / \text{ft}$, sphere-cone, 80-mesh.

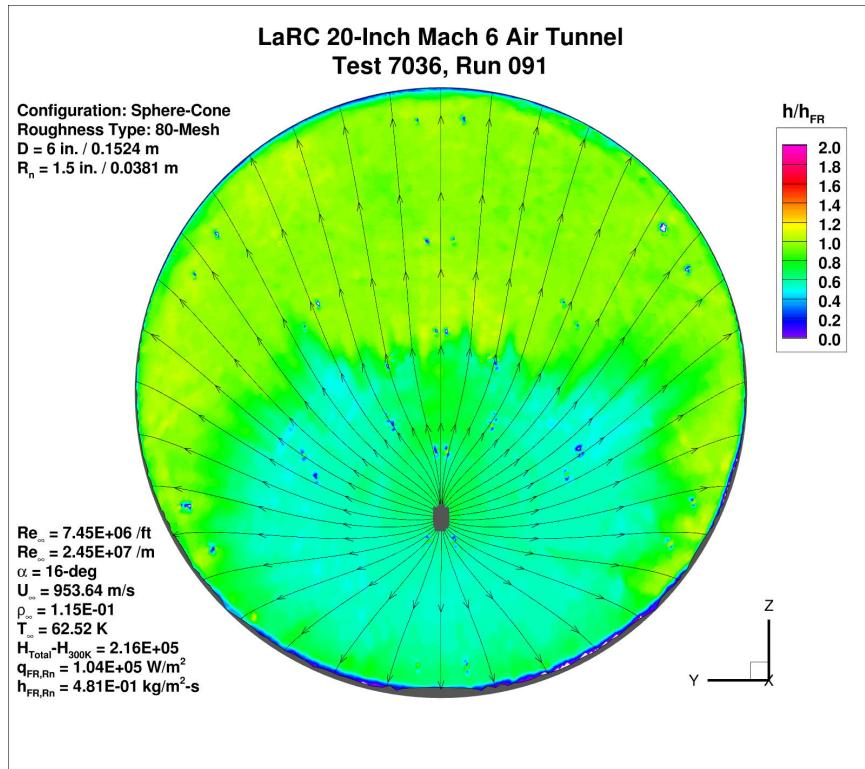


Figure 112. Test 7036, Run 91, $Re_{\infty} = 7.2 \times 10^6 / \text{ft}$, sphere-cone, 80-mesh.

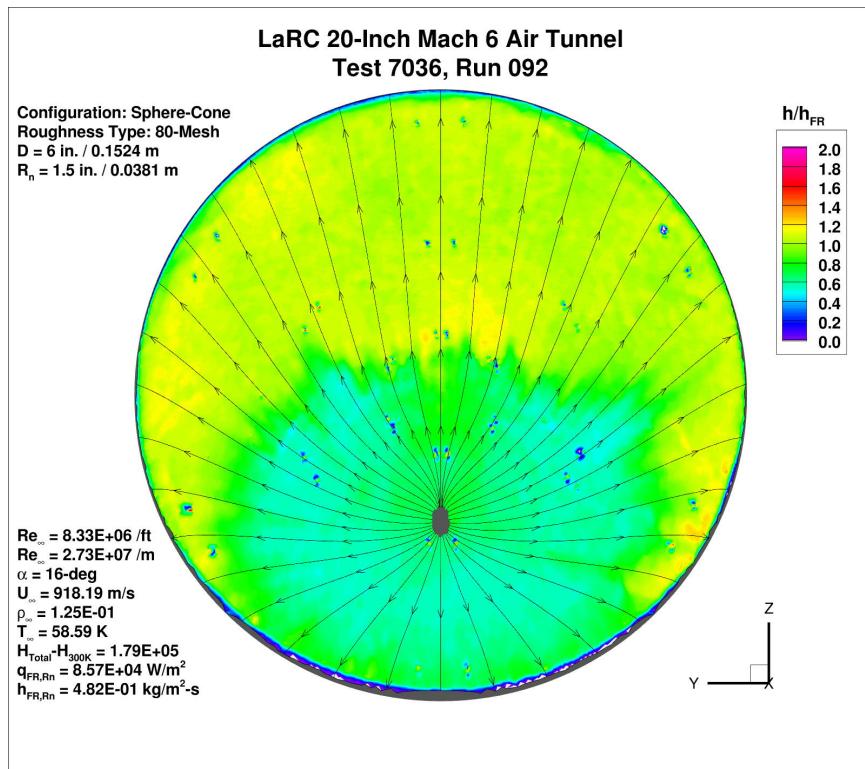


Figure 113. Test 7036, Run 92, $Re_{\infty} = 8.1 \times 10^6 / \text{ft}$, sphere-cone, 80-mesh.

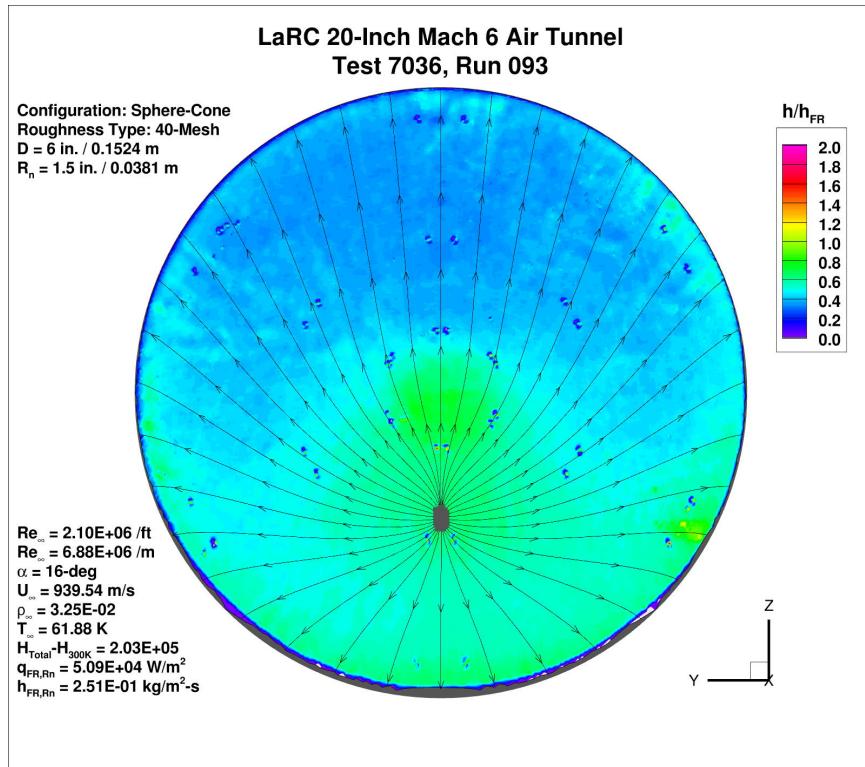


Figure 114. Test 7036, Run 93, $Re_{\infty} = 2.1 \times 10^6/\text{ft}$, sphere-cone, 40-mesh.

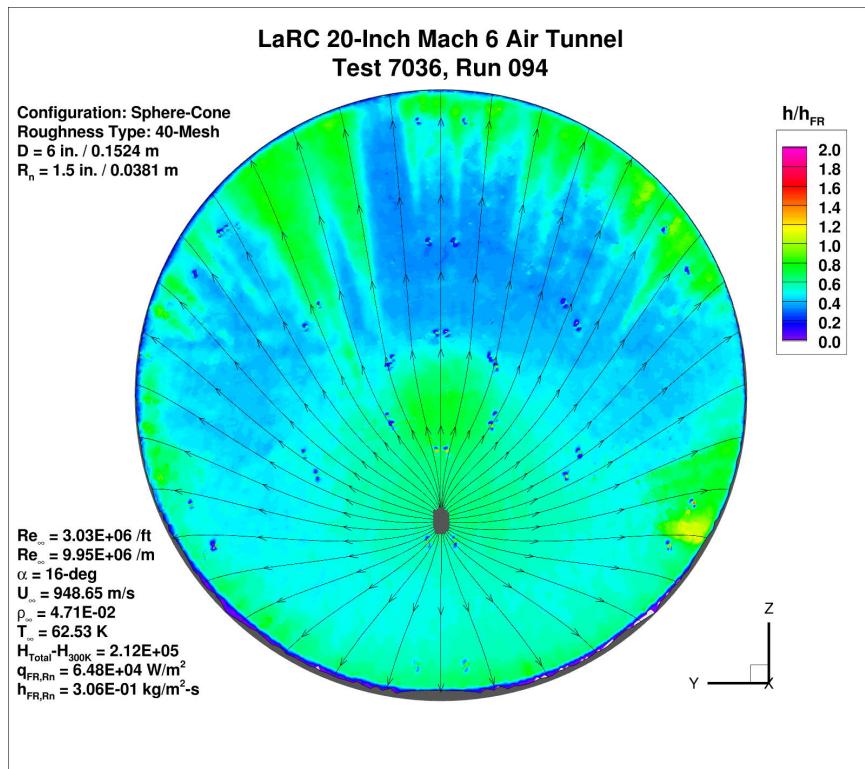


Figure 115. Test 7036, Run 94, $Re_{\infty} = 3.0 \times 10^6/\text{ft}$, sphere-cone, 40-mesh.

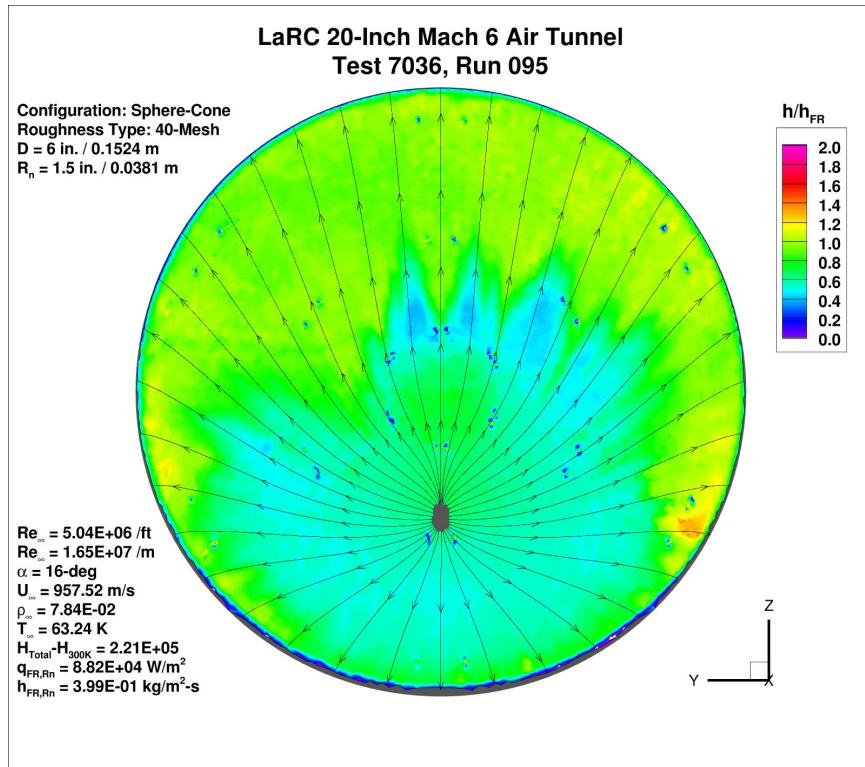


Figure 116. Test 7036, Run 95, $Re_{\infty} = 5.0 \times 10^6 / \text{ft}$, sphere-cone, 40-mesh.

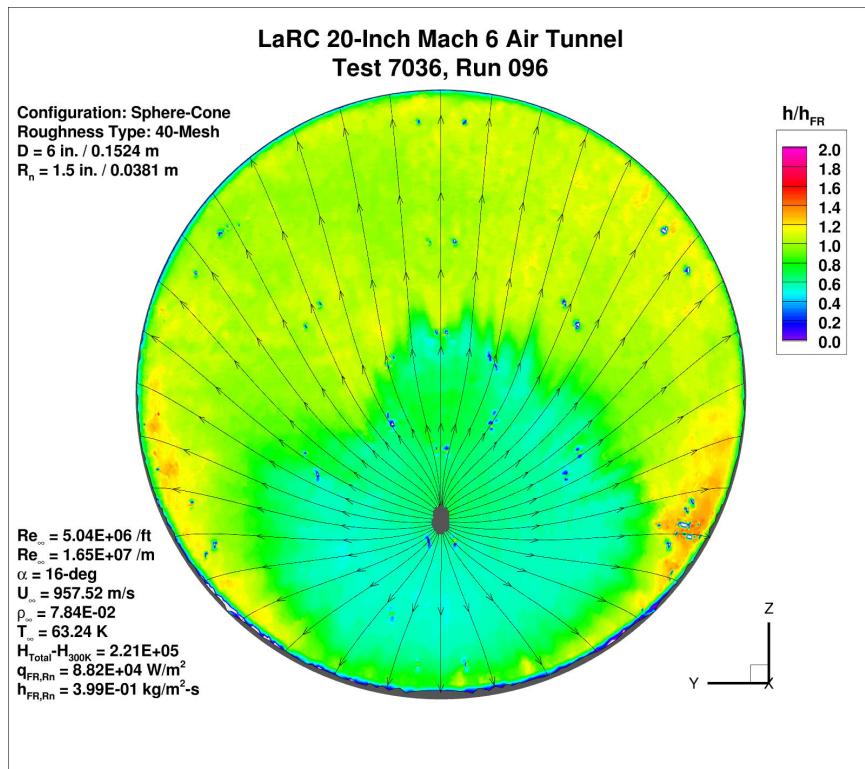


Figure 117. Test 7036, Run 96, $Re_{\infty} = 6.5 \times 10^6 / \text{ft}$, sphere-cone, 40-mesh.

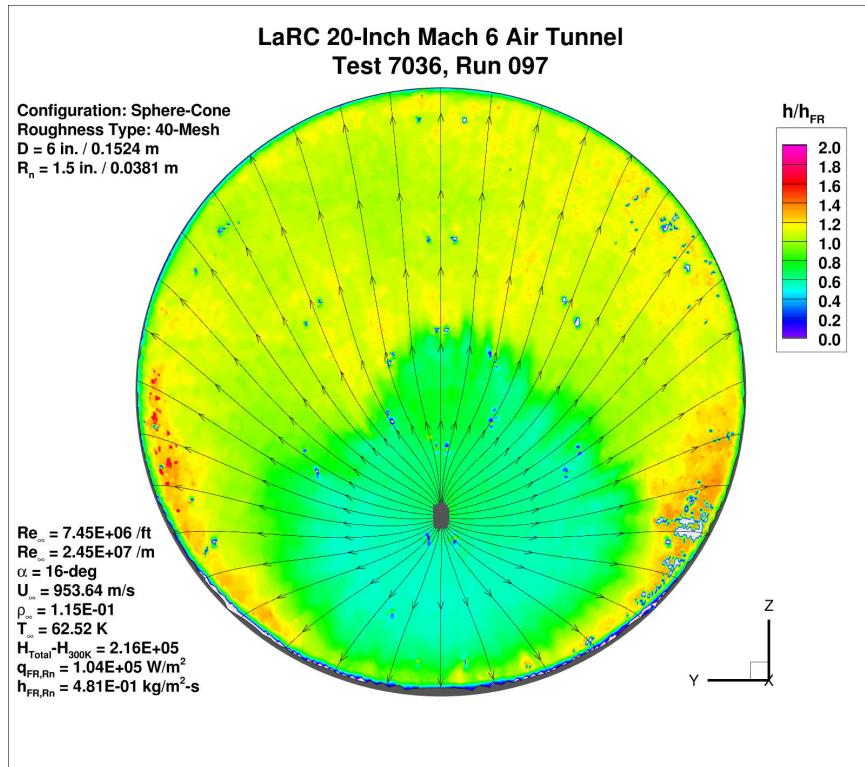


Figure 118. Test 7036, Run 97, $Re_{\infty} = 7.2 \times 10^6/\text{ft}$, sphere-cone, 40-mesh.

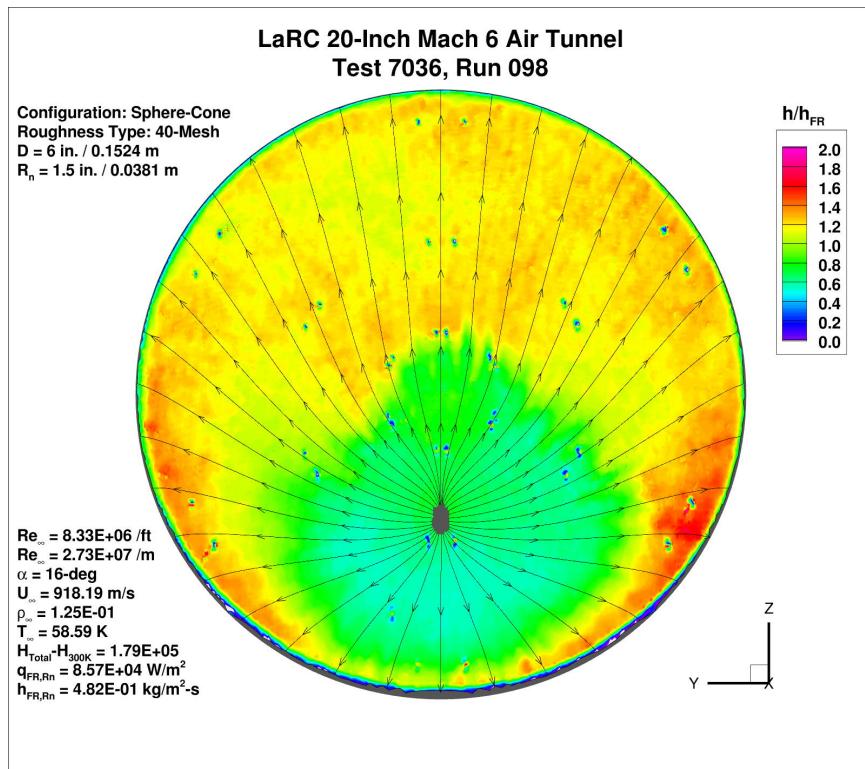


Figure 119. Test 7036, Run 98, $Re_{\infty} = 8.1 \times 10^6/\text{ft}$, sphere-cone, 40-mesh.

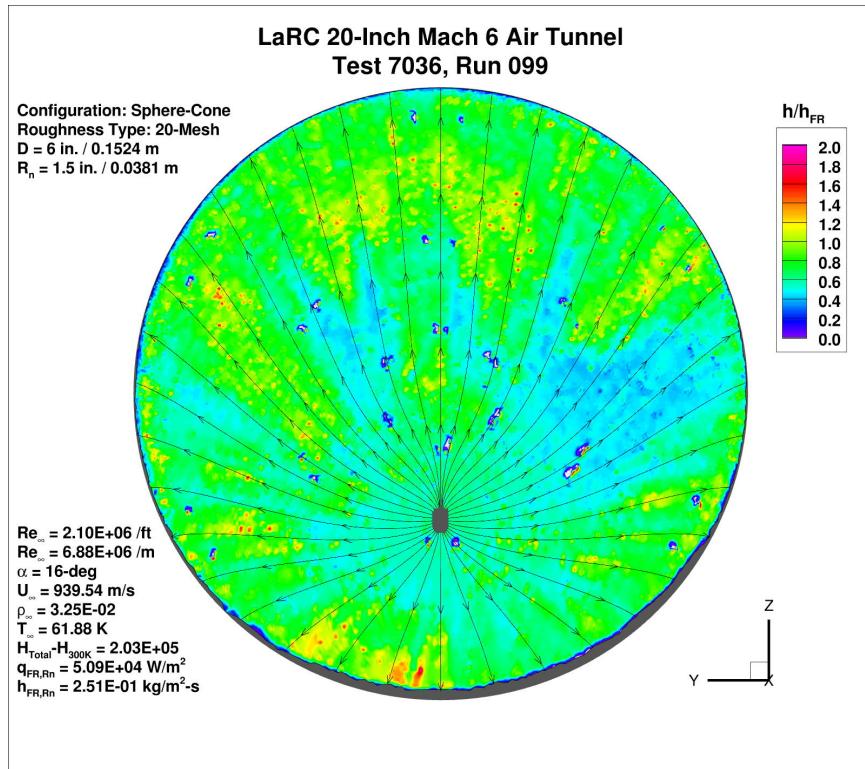


Figure 120. Test 7036, Run 99, $Re_{\infty} = 2.1 \times 10^6$ /ft, sphere-cone, 20-mesh.

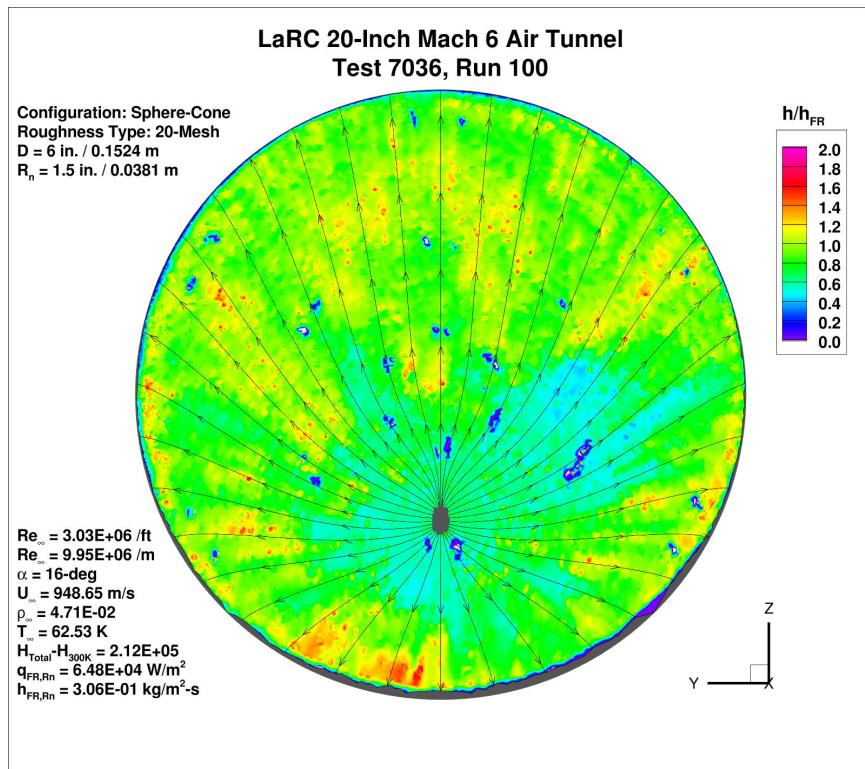


Figure 121. Test 7036, Run 100, $Re_{\infty} = 3.0 \times 10^6$ /ft, sphere-cone, 20-mesh.

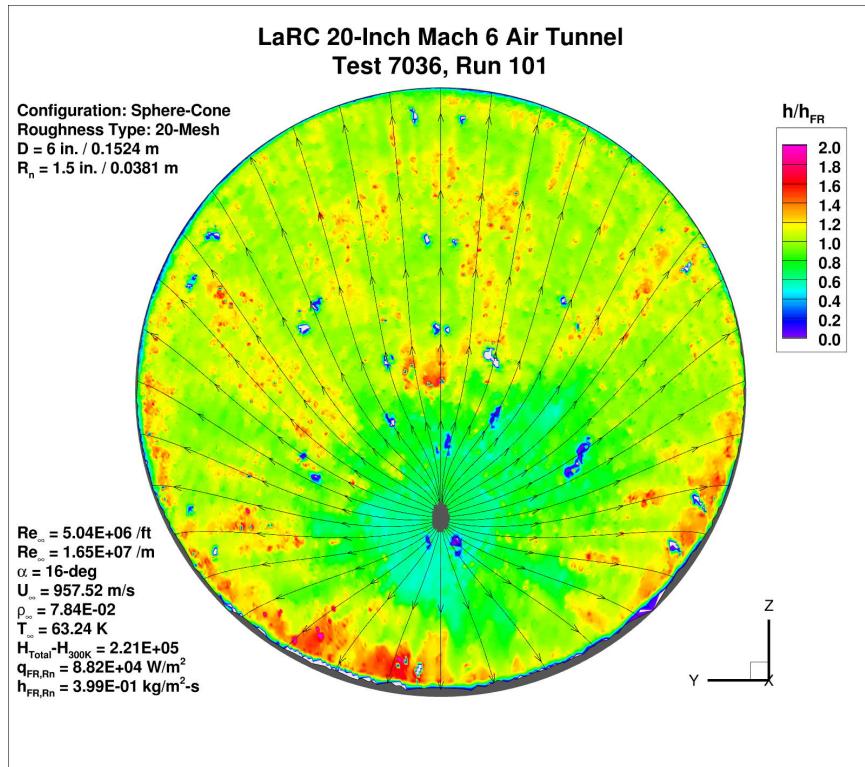


Figure 122. Test 7036, Run 101, $Re_{\infty} = 5.0 \times 10^6$ /ft, sphere-cone, 20-mesh.

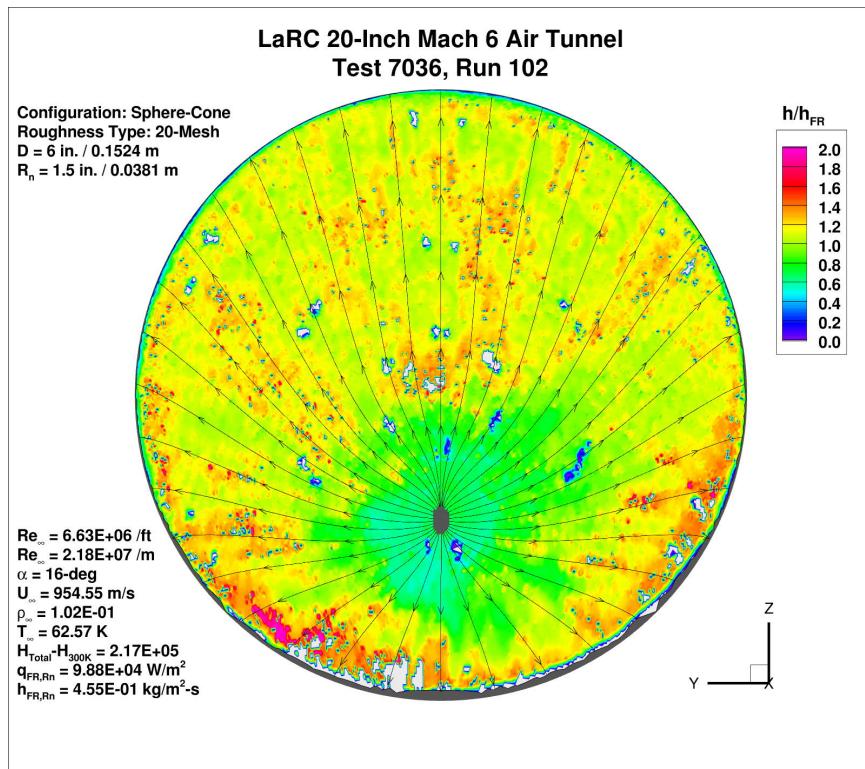


Figure 123. Test 7036, Run 102, $Re_{\infty} = 6.5 \times 10^6$ /ft, sphere-cone, 20-mesh.

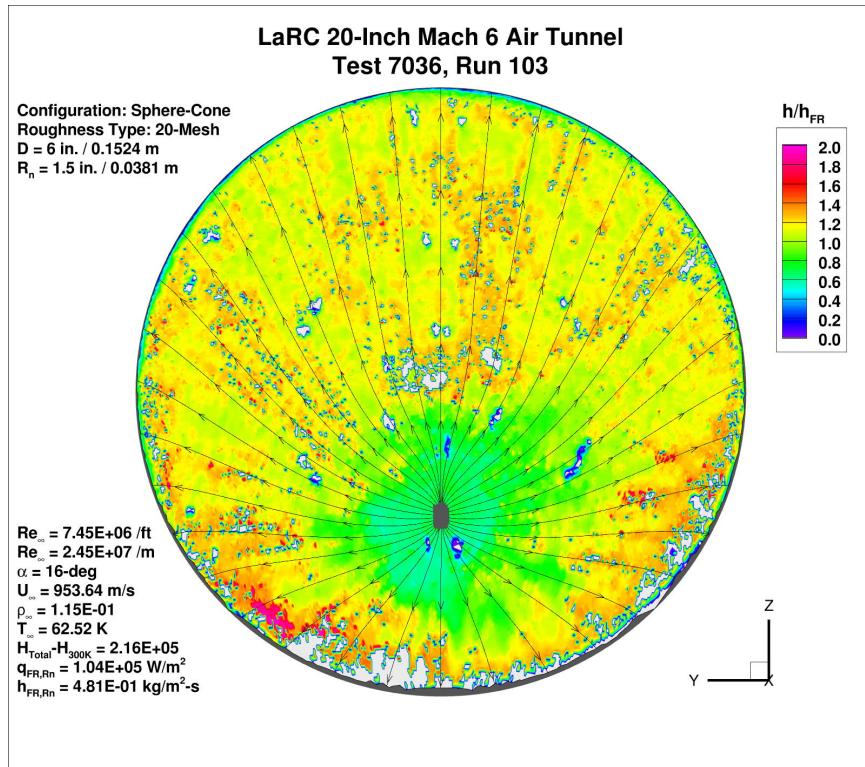


Figure 124. Test 7036, Run 103, Re_∞ = 7.2×10⁶/ft, sphere-cone, 20-mesh.

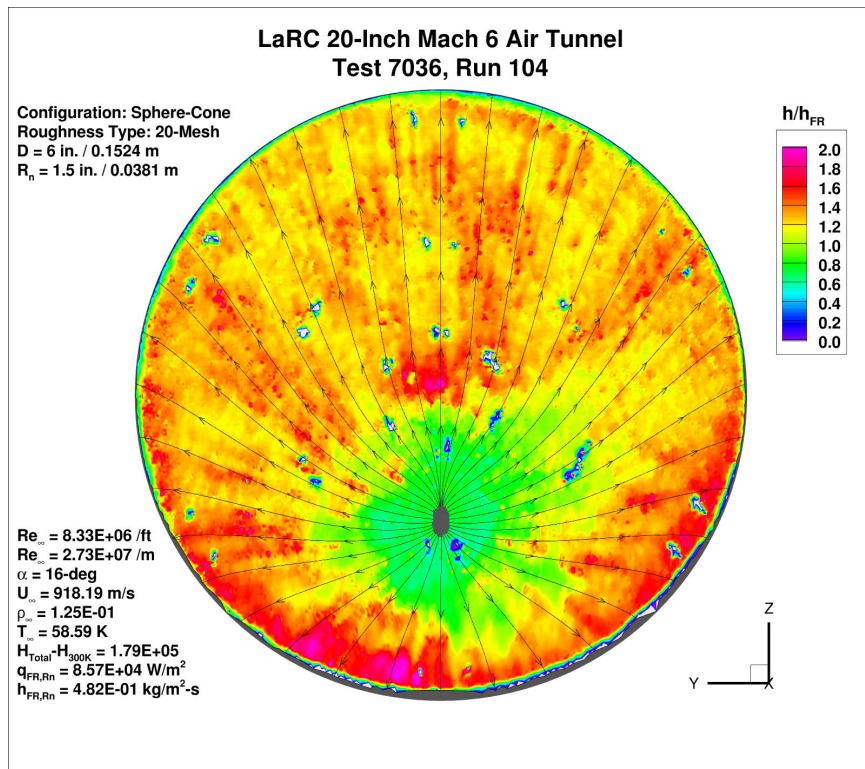


Figure 125. Test 7036, Run 104, Re_∞ = 8.1×10⁶/ft, sphere-cone, 20-mesh.

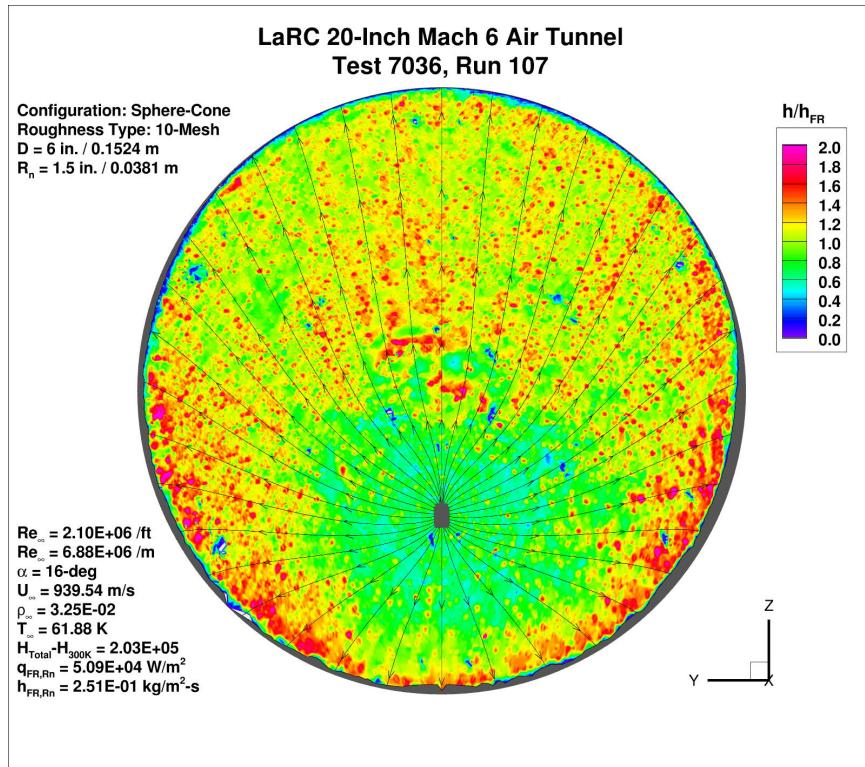


Figure 126. Test 7036, Run 107, $Re_{\infty} = 2.1 \times 10^6$ /ft, sphere-cone, 10-mesh.

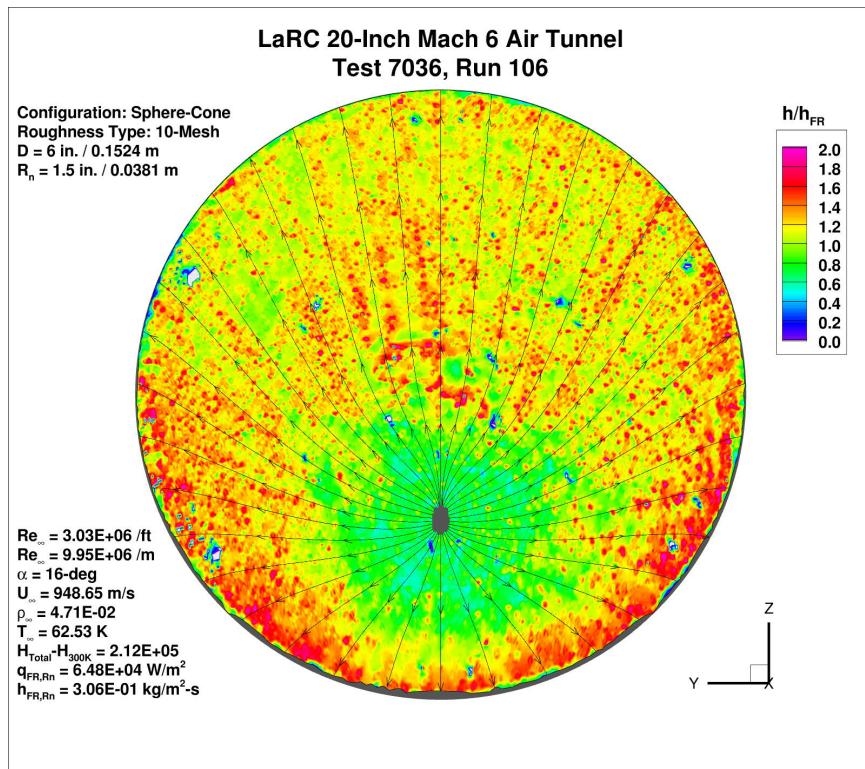


Figure 127. Test 7036, Run 106, $Re_{\infty} = 3.0 \times 10^6$ /ft, sphere-cone, 10-mesh.

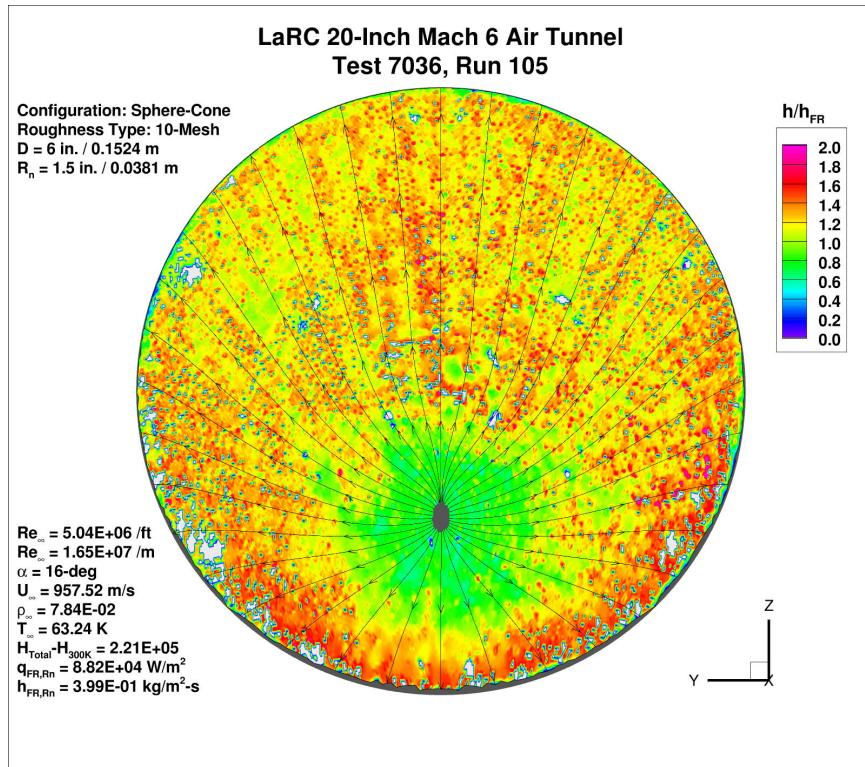


Figure 128. Test 7036, Run 105, Re_∞ = 5.0×10⁶/ft, sphere-cone, 10-mesh.

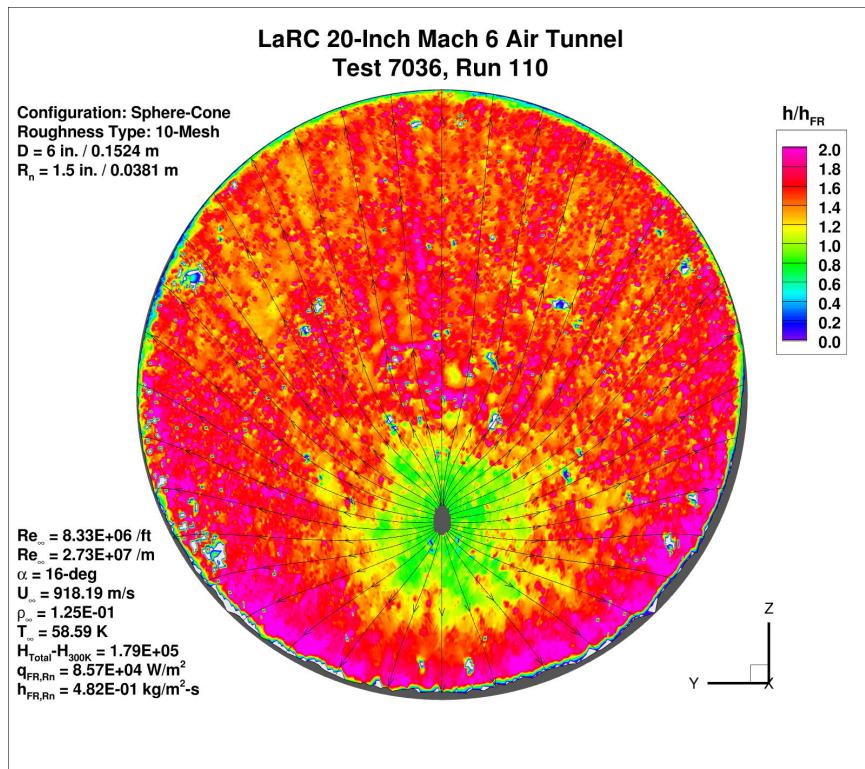


Figure 129. Test 7036, Run 108, Re_∞ = 6.5×10⁶/ft, sphere-cone, 10-mesh.

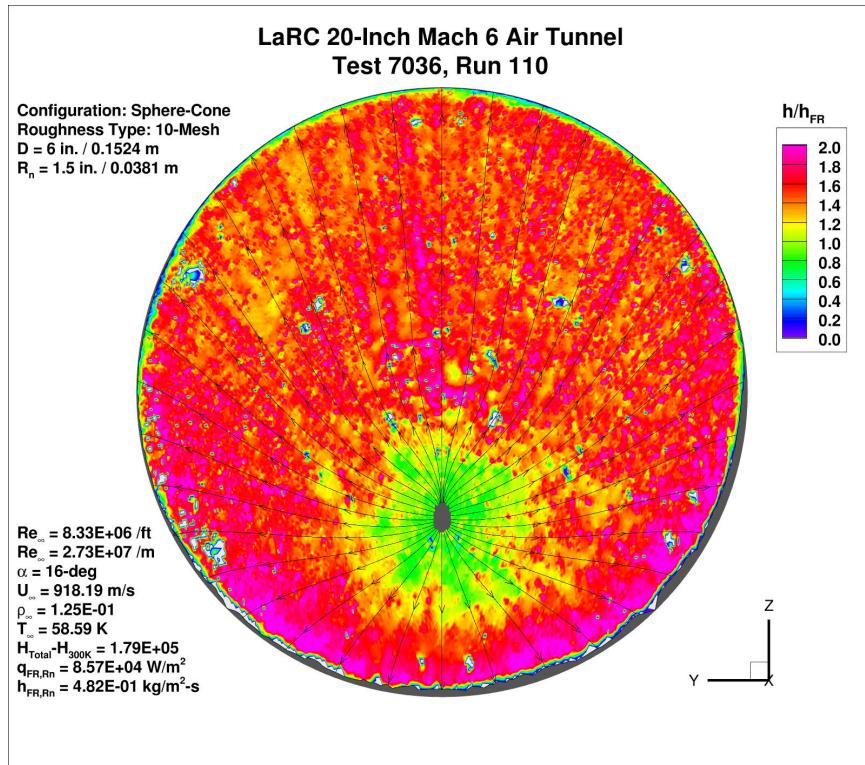


Figure 130. Test 7036, Run 109, $Re_{\infty} = 7.2 \times 10^6 /ft$, sphere-cone, 10-mesh.

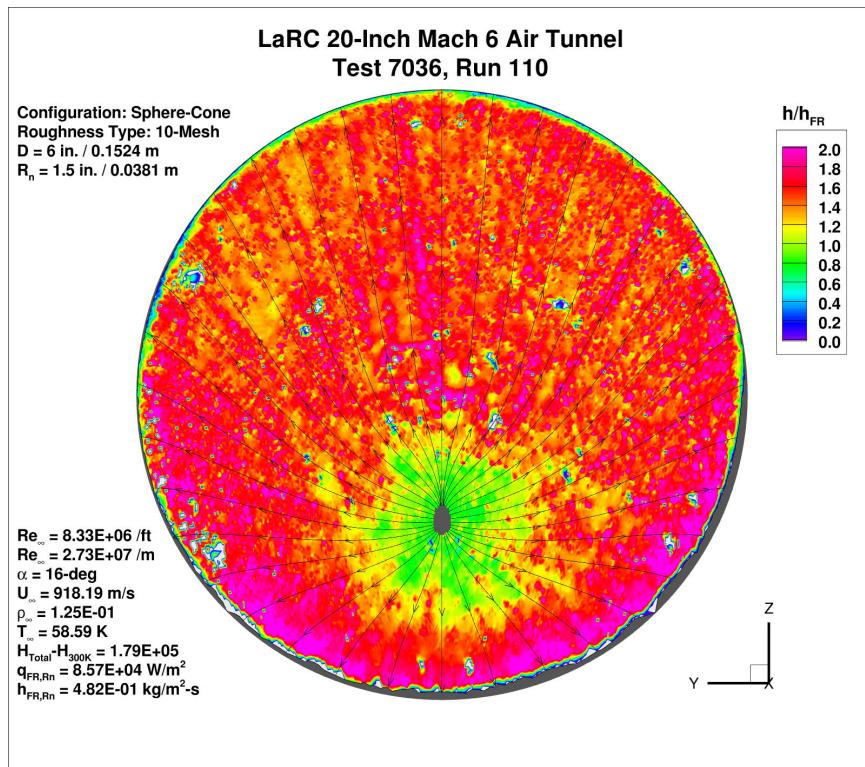


Figure 131. Test 7036, Run 110, $Re_{\infty} = 8.1 \times 10^6 /ft$, sphere-cone, 10-mesh.

Appendix B. Spherical-Cap Geometry Global Heating Images

Global heating images for the sphere-cone geometry from Test 7057 in the LAL 20-Inch Mach 6 Air Tunnel are presented in this Appendix in Figure 132 through Figure 173.

At higher Reynolds numbers and/or larger roughness heights, white patches on the images indicate areas where the measured surface temperatures exceed the calibrated range of the phosphor thermography and thus no valid data were obtained.

Boundary-layer edge streamlines determined from laminar, smooth-surface LAURA simulations have been superimposed on the images to illustrate the nature of the flow field.

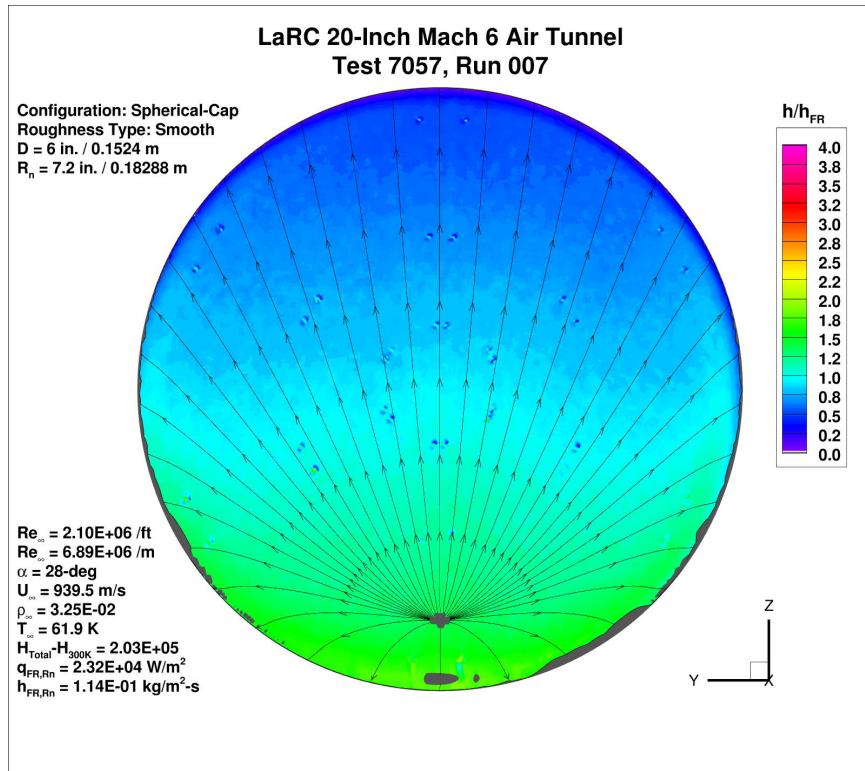


Figure 132. Test 7057, Run 7, Re_∞ = 2.1×10⁶/ft, spherical-cap, smooth OML.

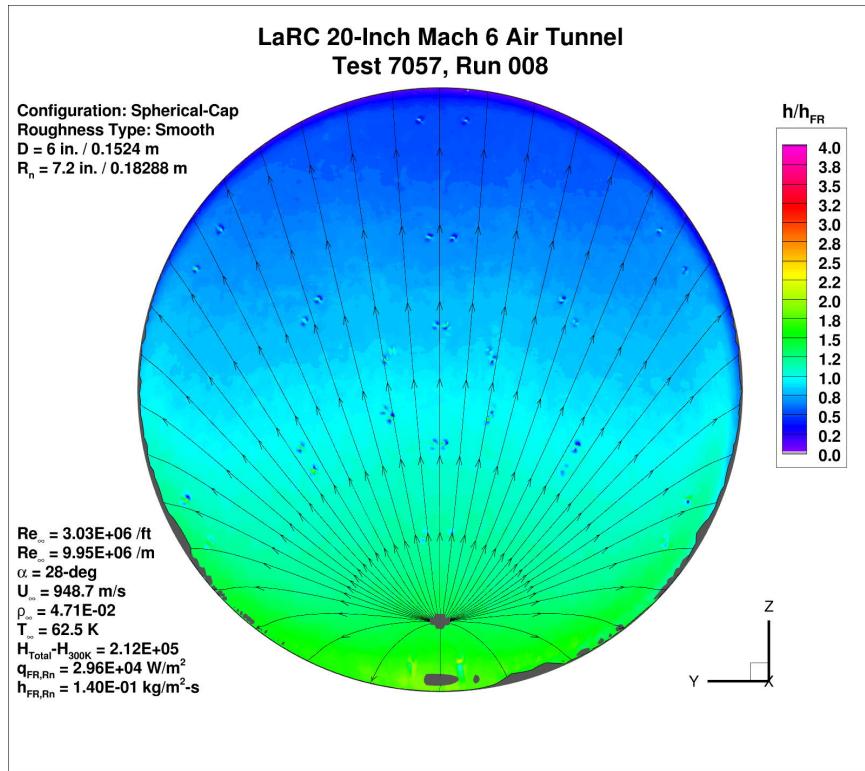


Figure 133. Test 7057, Run 8, Re_∞ = 3.0×10⁶/ft, spherical-cap, smooth OML.

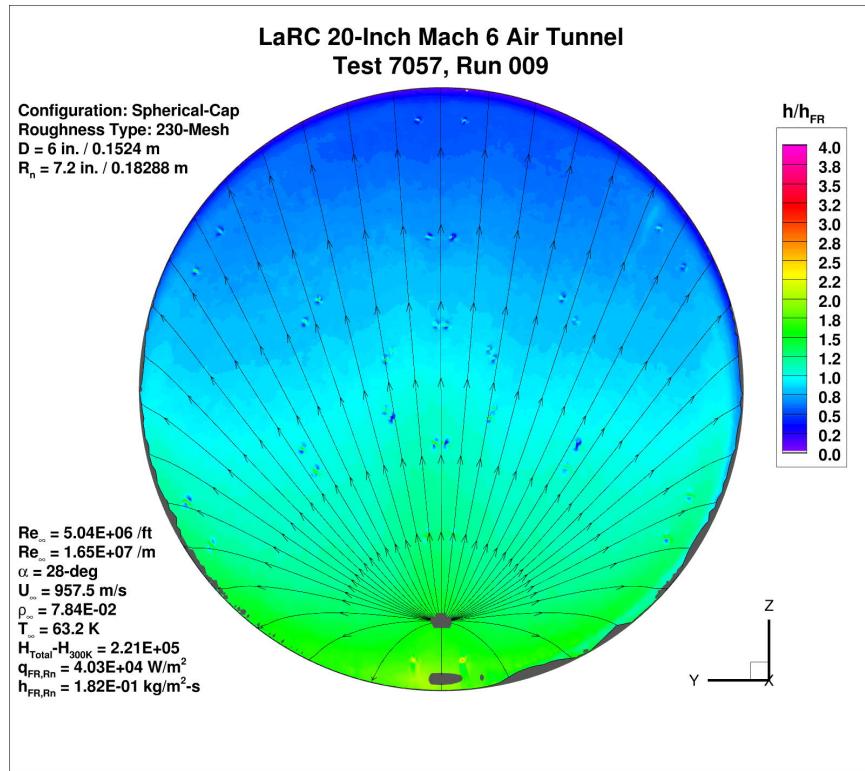


Figure 134. Test 7057, Run 9, $Re_{\infty} = 5.0 \times 10^6/\text{ft}$, spherical-cap, smooth OML.

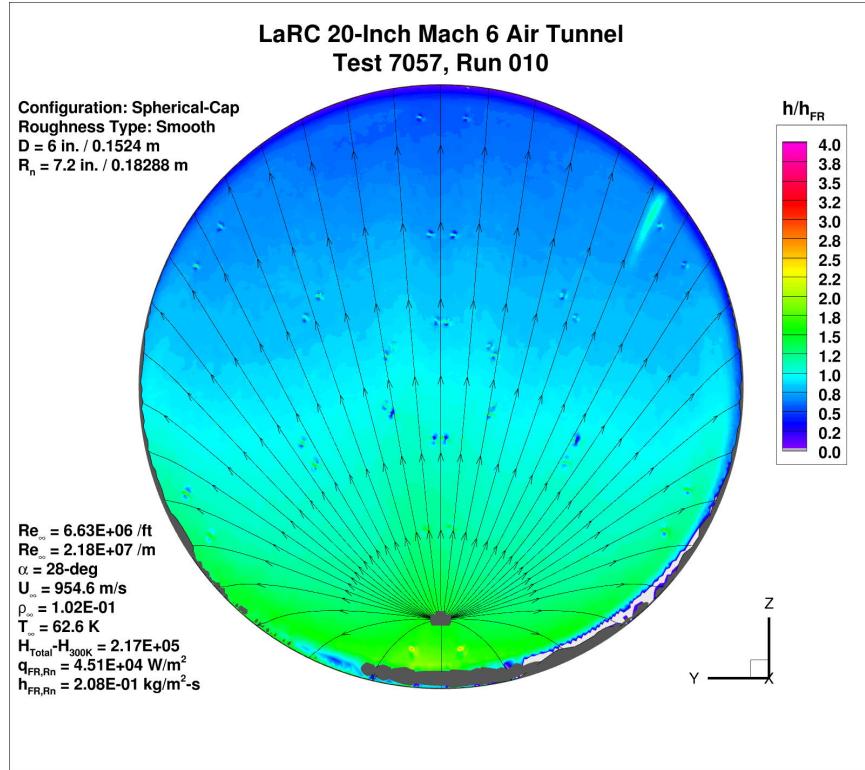


Figure 135. Test 7057, Run 10, $Re_{\infty} = 6.5 \times 10^6/\text{ft}$, spherical-cap, smooth OML.

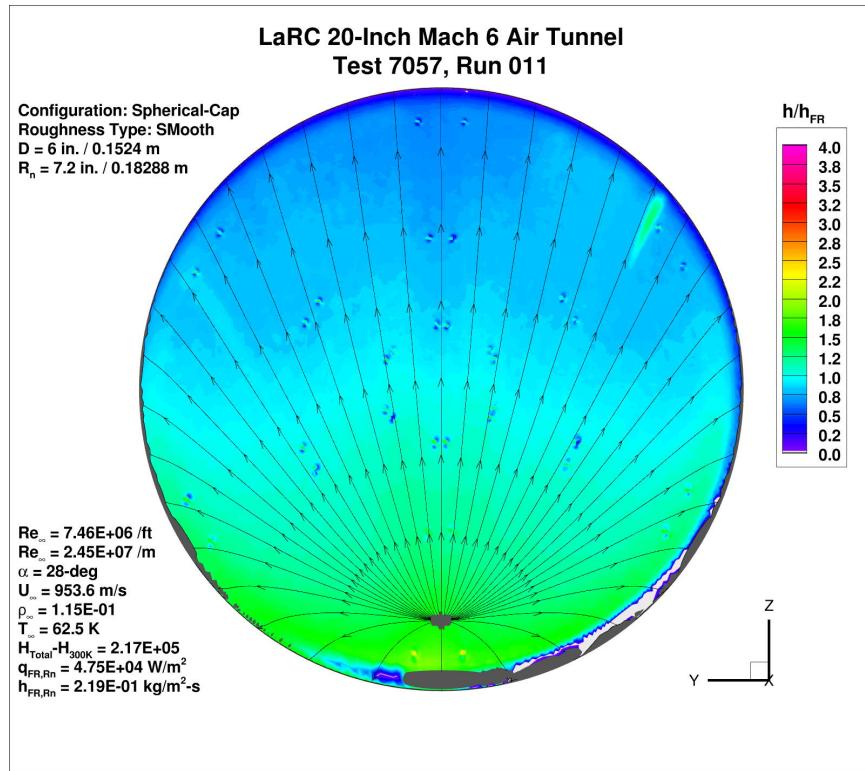


Figure 136. Test 7057, Run 11, $Re_{\infty} = 7.2 \times 10^6 / \text{ft}$, spherical-cap, smooth OML.

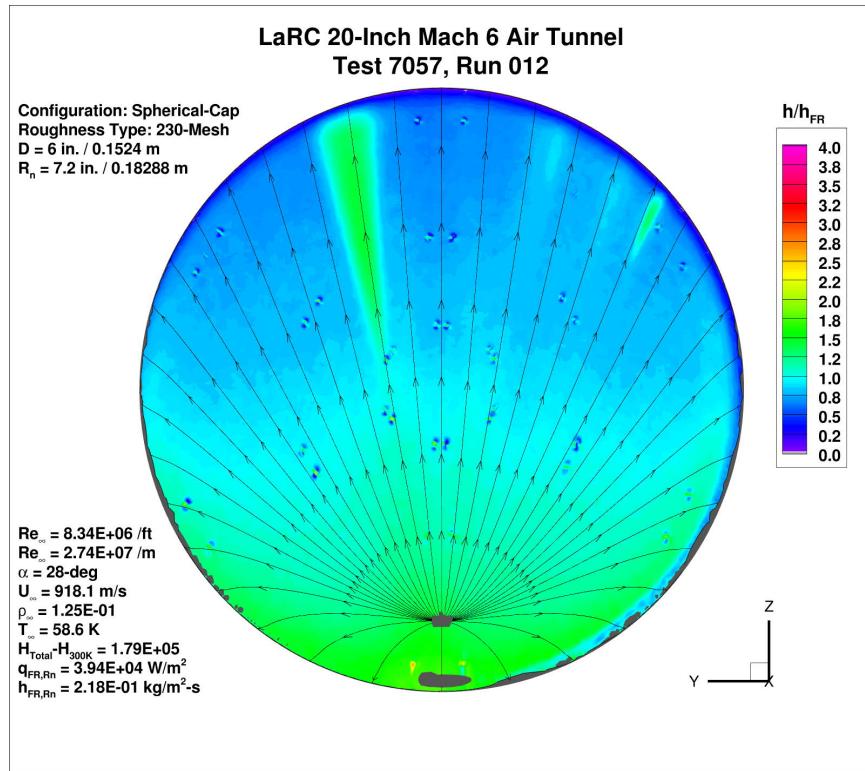


Figure 137. Test 7057, Run 12, $Re_{\infty} = 8.1 \times 10^6 / \text{ft}$, spherical-cap, smooth OML.

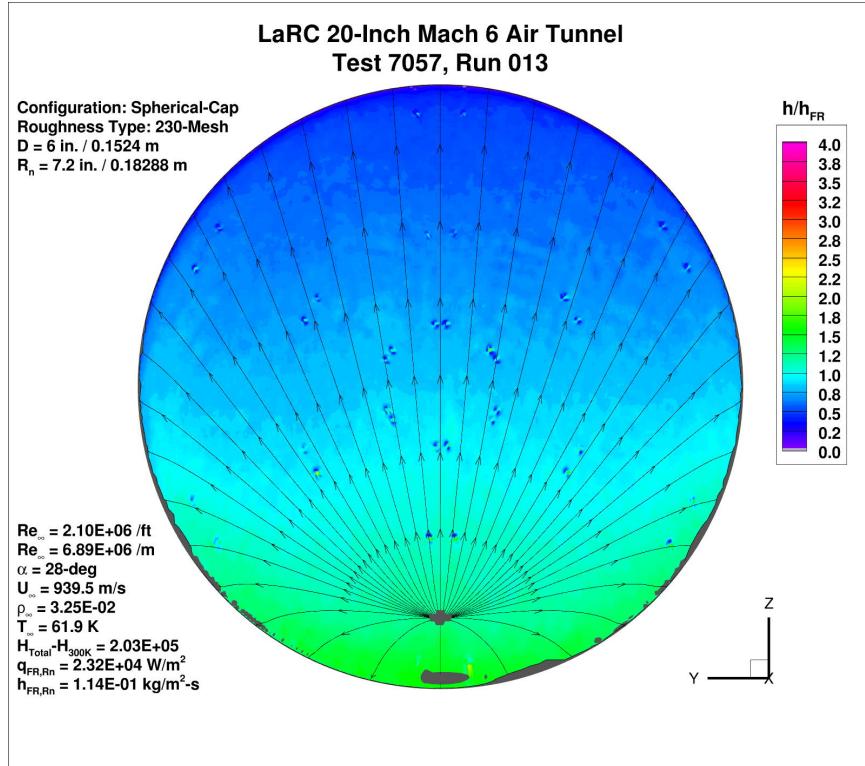


Figure 138. Test 7057, Run 13, $Re_{\infty} = 2.1 \times 10^6 / \text{ft}$, spherical-cap, 230-mesh.

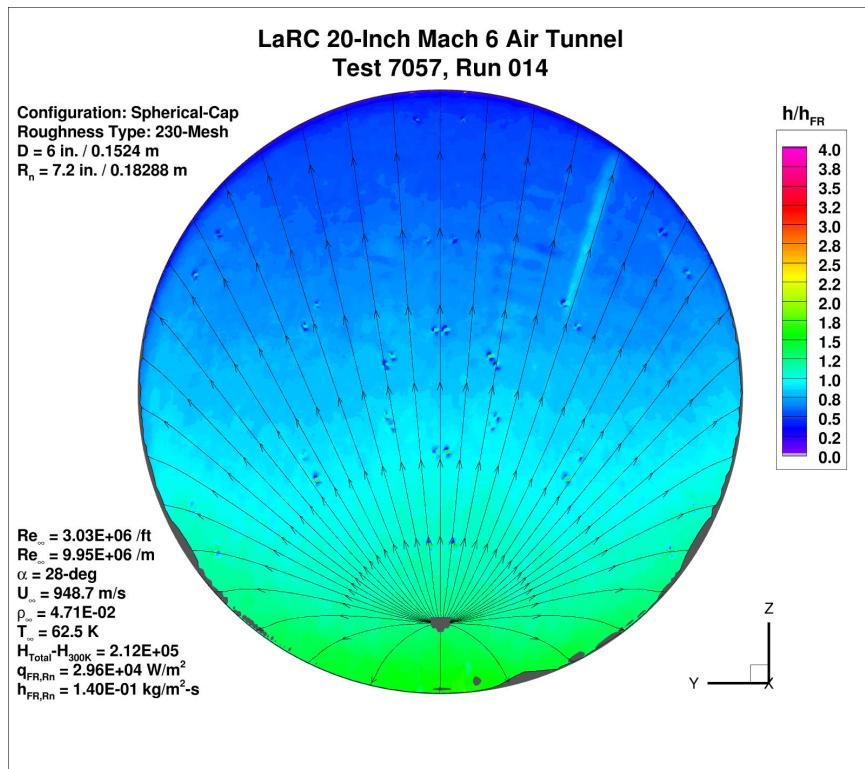


Figure 139. Test 7057, Run 14, $Re_{\infty} = 3.0 \times 10^6 / \text{ft}$, spherical-cap, 230-mesh.

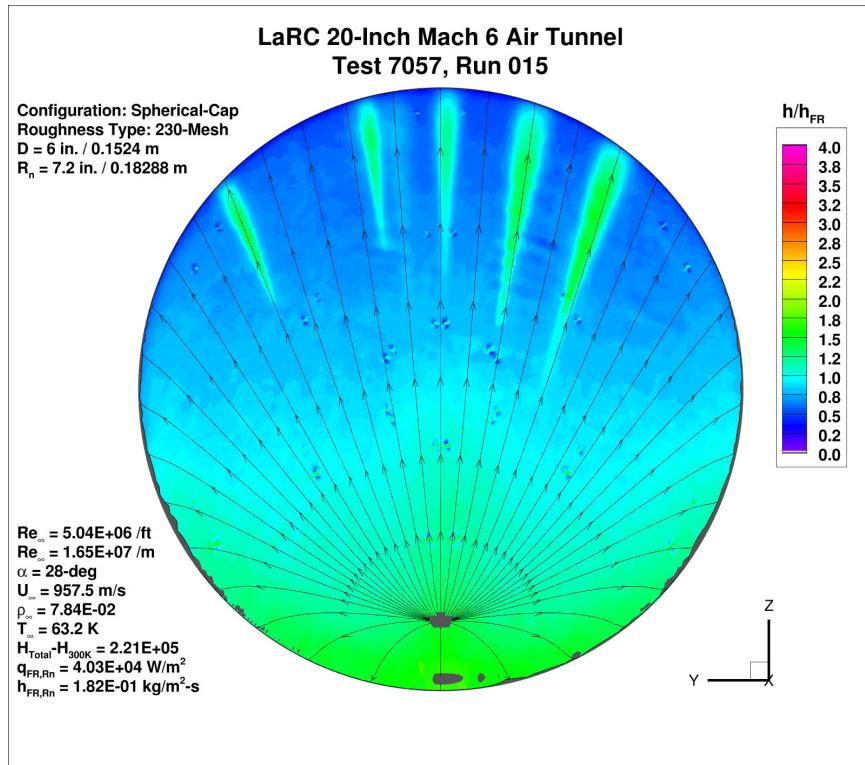


Figure 140. Test 7057, Run 15, $Re_{\infty} = 5.0 \times 10^6 / \text{ft}$, spherical-cap, 230-mesh.

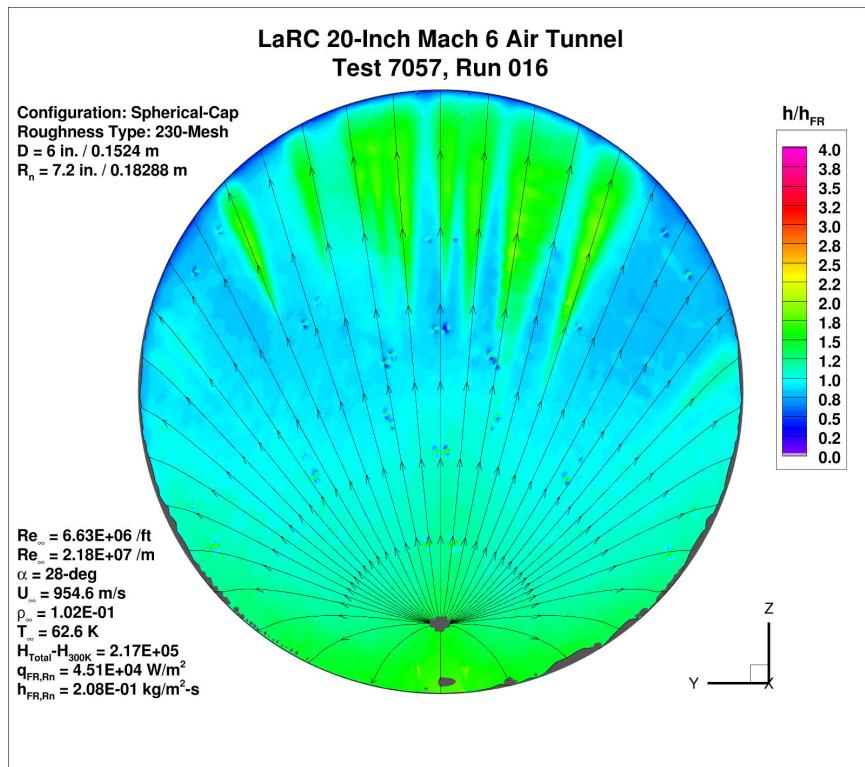


Figure 141. Test 7057, Run 16, $Re_{\infty} = 6.5 \times 10^6 / \text{ft}$, spherical-cap, 230-mesh.

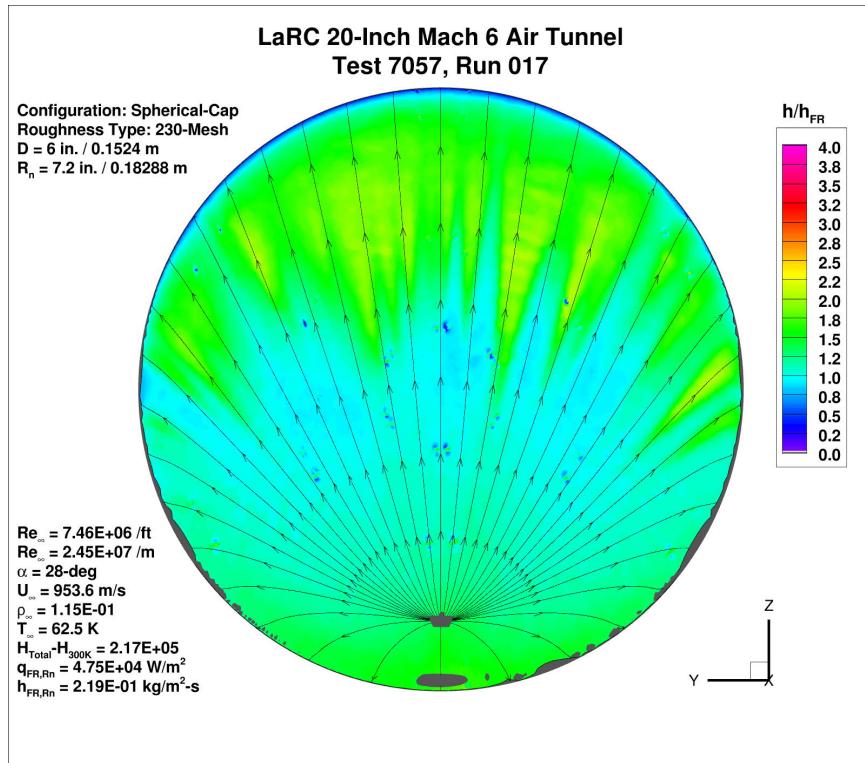


Figure 142. Test 7057, Run 17, Re_∞ = 7.2×10⁶/ft, spherical-cap, 230-mesh.

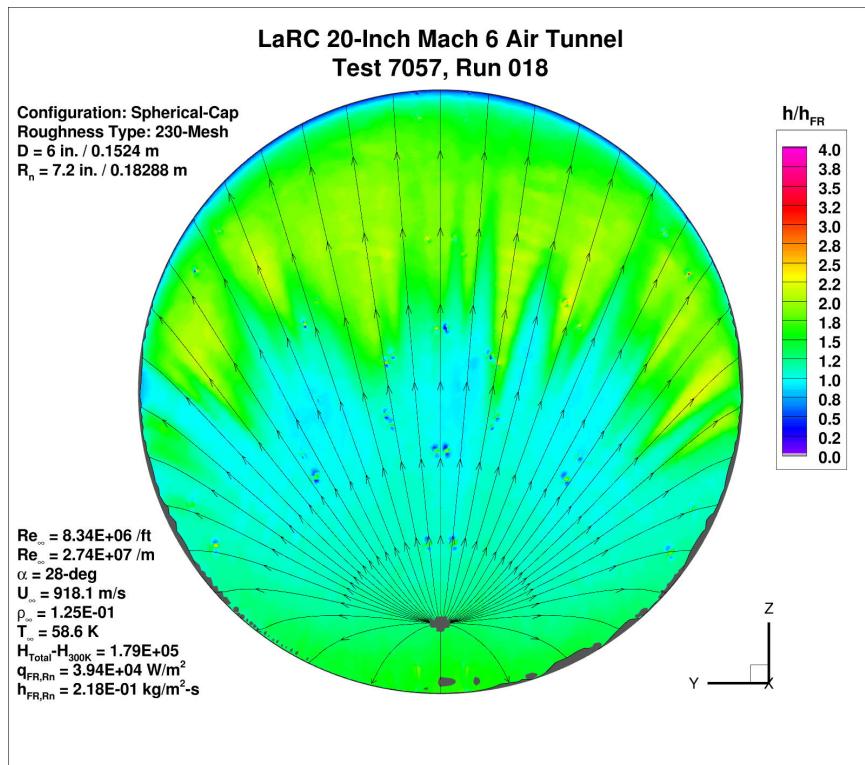


Figure 143. Test 7057, Run 18, Re_∞ = 8.1×10⁶/ft, spherical-cap, 230-mesh.

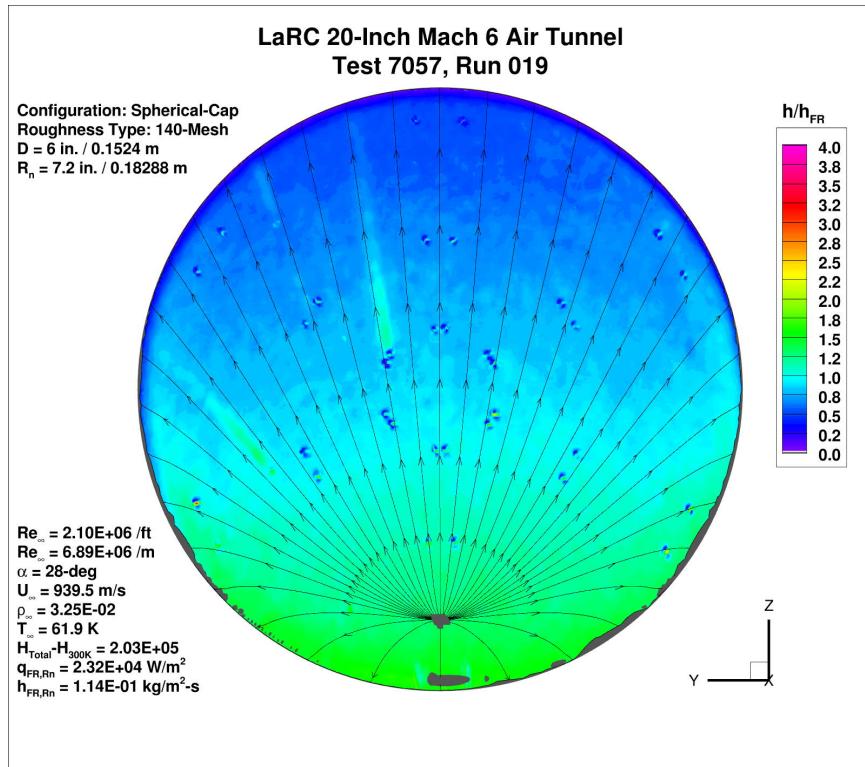


Figure 144. Test 7057, Run 19, Re_∞ = 2.1×10⁶/ft, spherical-cap, 140-mesh.

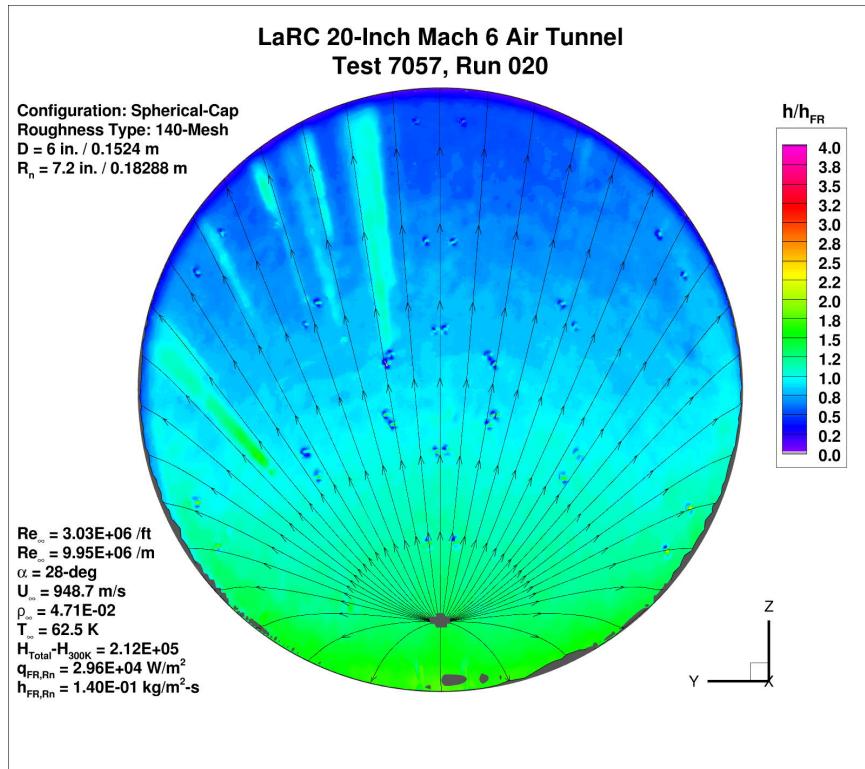


Figure 145. Test 7057, Run 20, Re_∞ = 3.0×10⁶/ft, spherical-cap, 140-mesh.

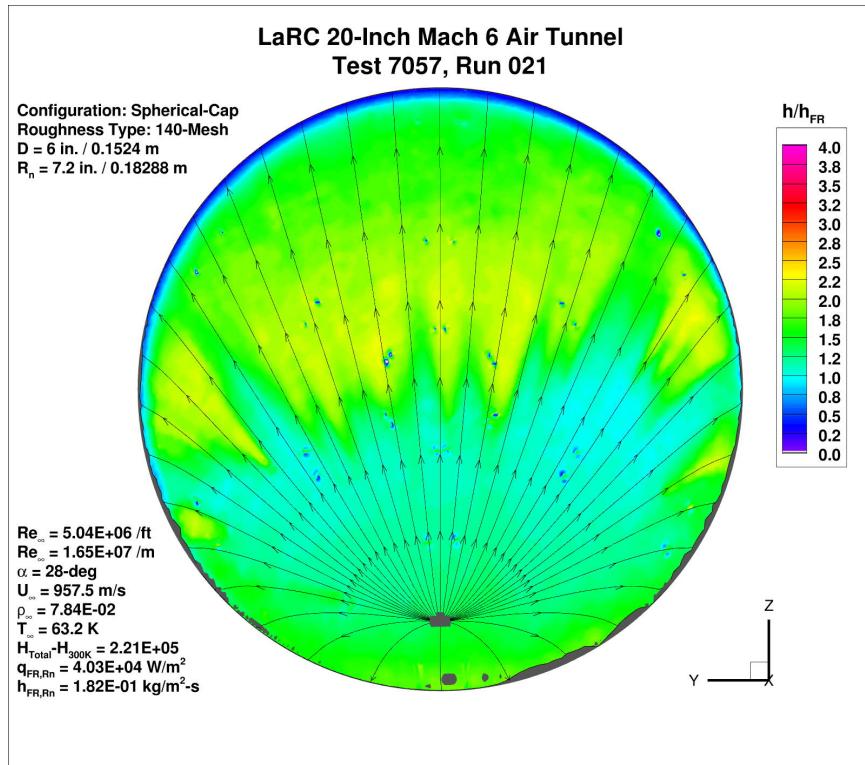


Figure 146. Test 7057, Run 21, $Re_{\infty} = 5.0 \times 10^6 / \text{ft}$, spherical-cap, 140-mesh.

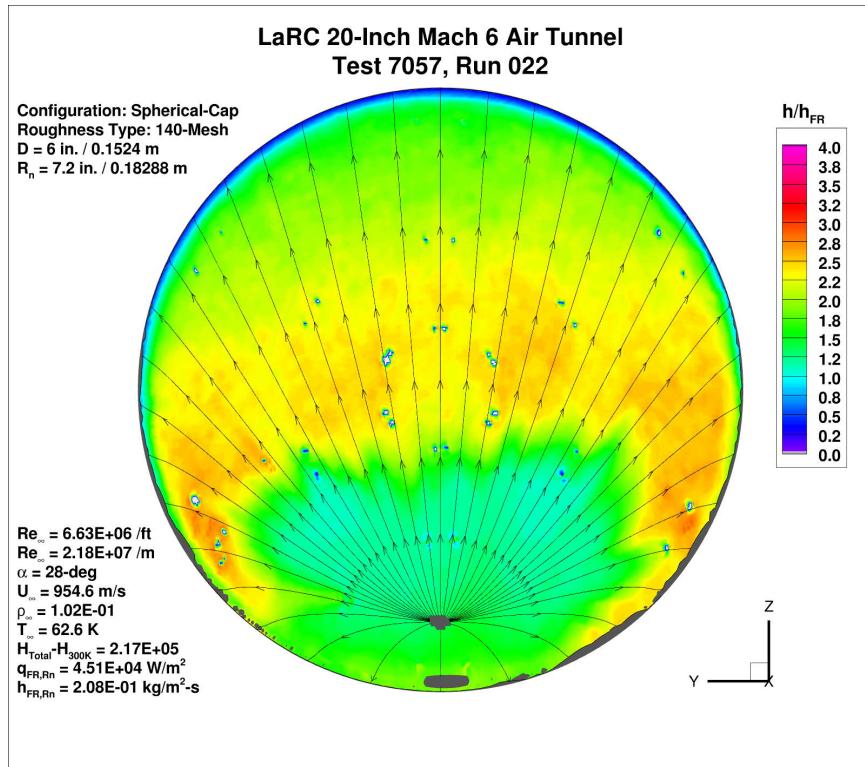


Figure 147. Test 7057, Run 22, $Re_{\infty} = 6.5 \times 10^6 / \text{ft}$, spherical-cap, 140-mesh.

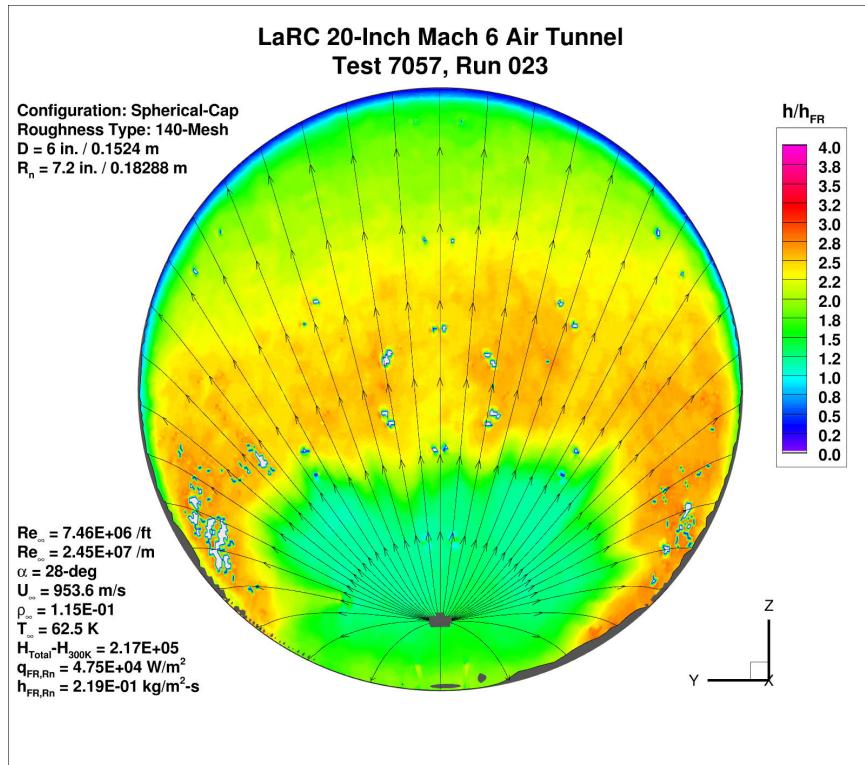


Figure 148. Test 7057, Run 23, Re_∞ = 7.2×10⁶/ft, spherical-cap, 140-mesh.

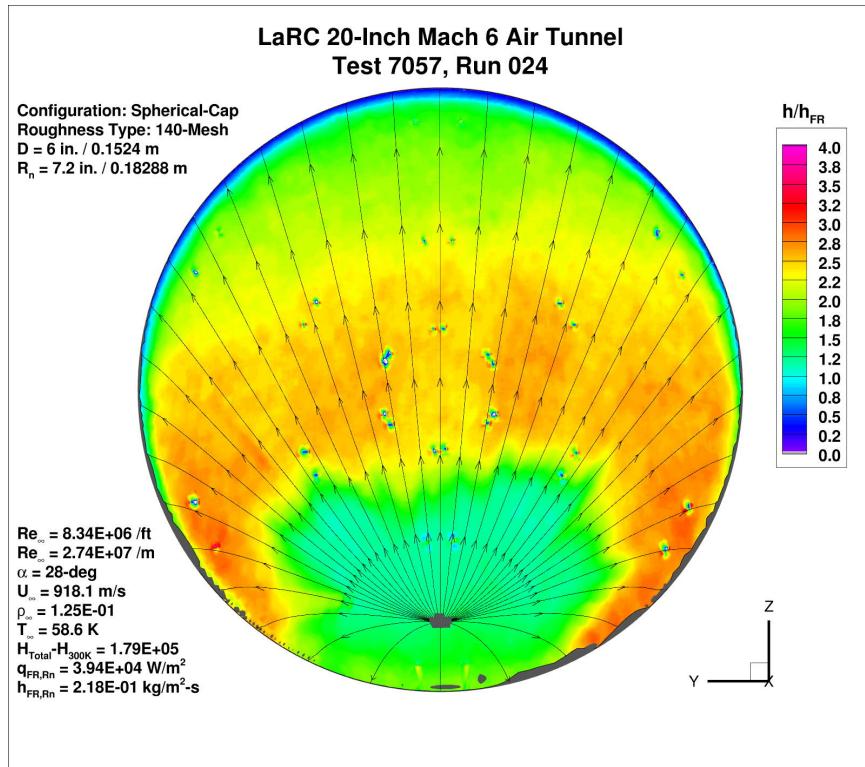


Figure 149. Test 7057, Run 24, Re_∞ = 8.1×10⁶/ft, spherical-cap, 140-mesh.

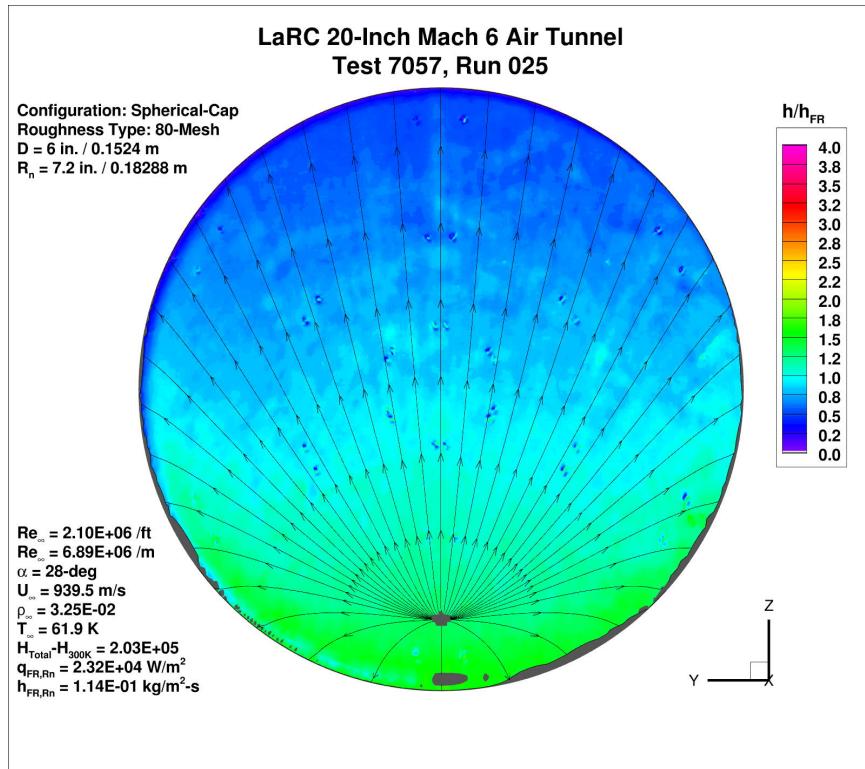


Figure 150. Test 7057, Run 25, $Re_{\infty} = 2.1 \times 10^6/\text{ft}$, spherical-cap, 80-mesh.

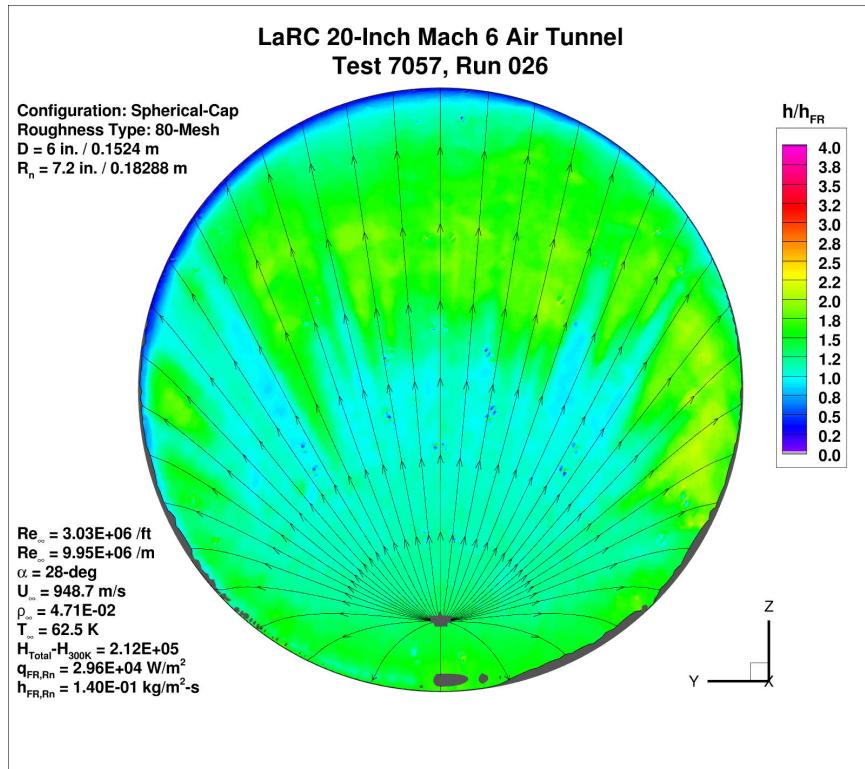


Figure 151. Test 7057, Run 26, $Re_{\infty} = 3.0 \times 10^6/\text{ft}$, spherical-cap, 80-mesh.

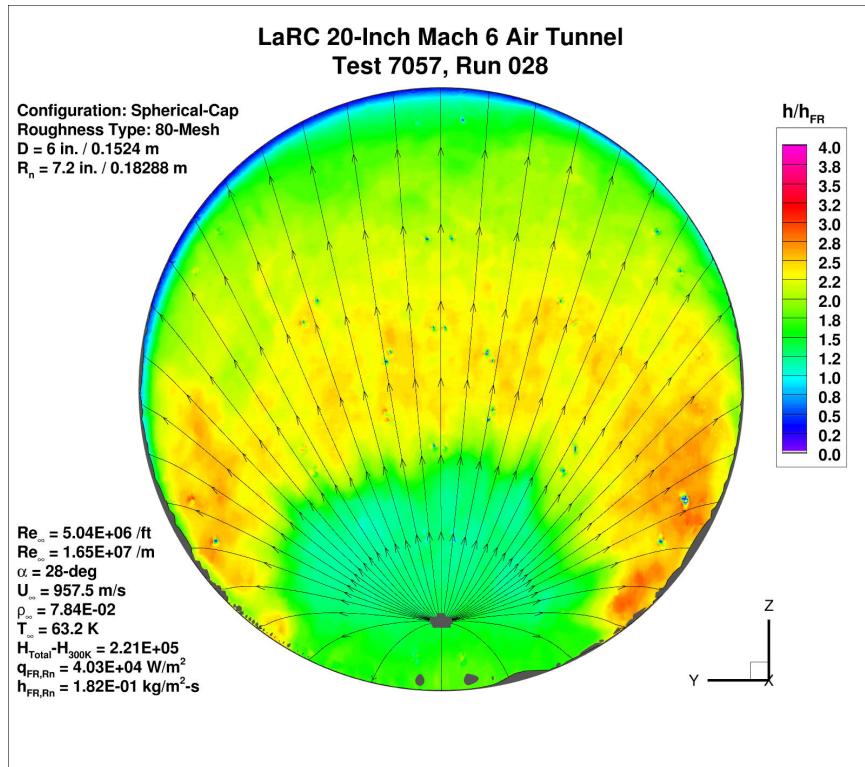


Figure 152. Test 7057, Run 28, $Re_{\infty} = 5.0 \times 10^6/\text{ft}$, spherical-cap, 80-mesh.

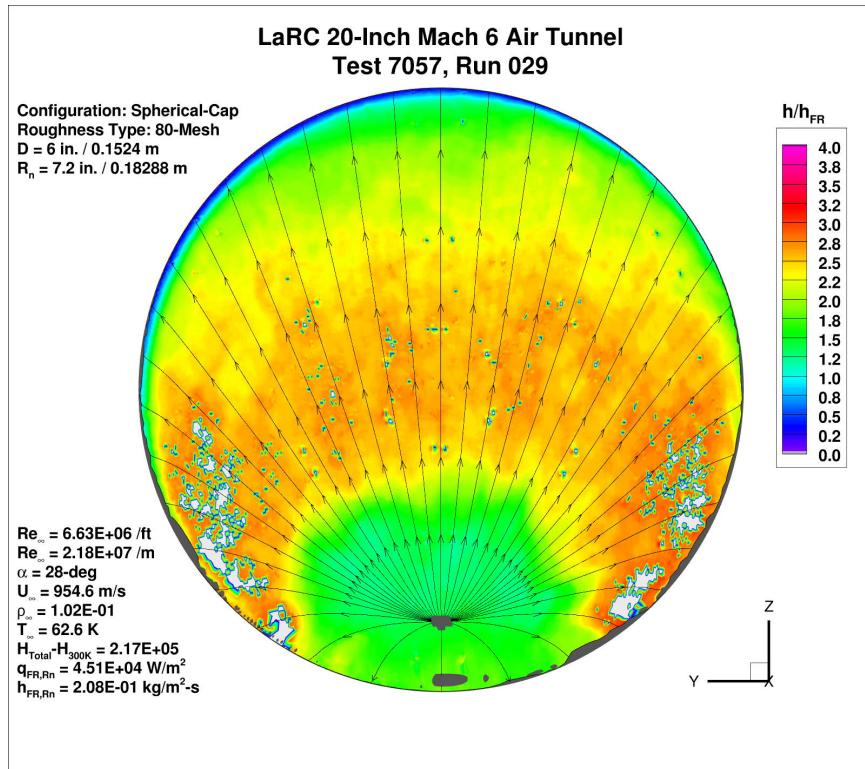


Figure 153. Test 7057, Run 29, $Re_{\infty} = 6.5 \times 10^6/\text{ft}$, spherical-cap, 80-mesh.

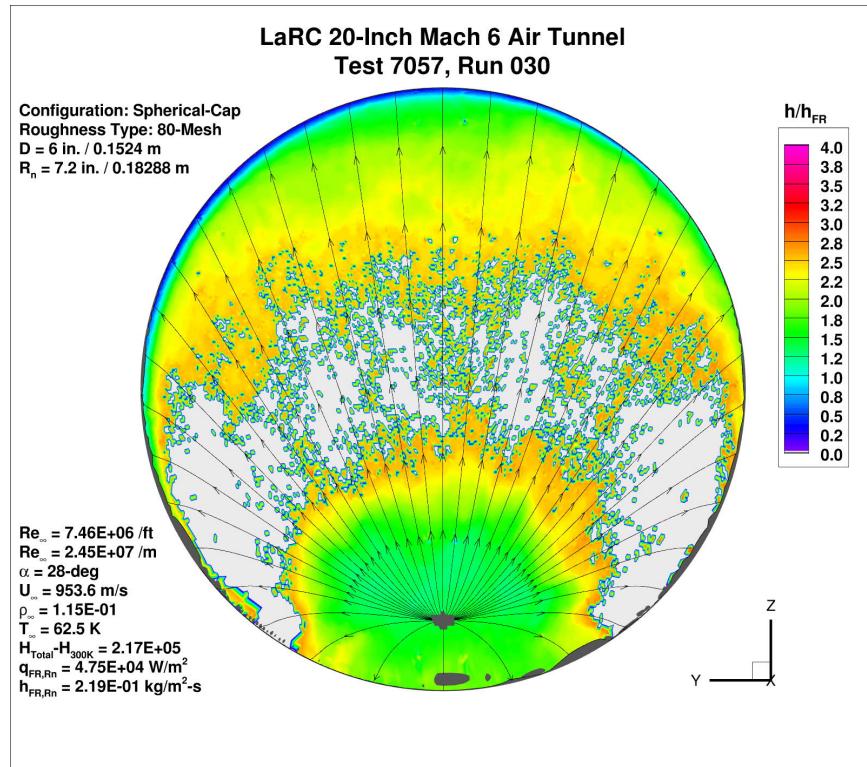


Figure 154. Test 7057, Run 30, $Re_{\infty} = 7.2 \times 10^6 / \text{ft}$, spherical-cap, 80-mesh.

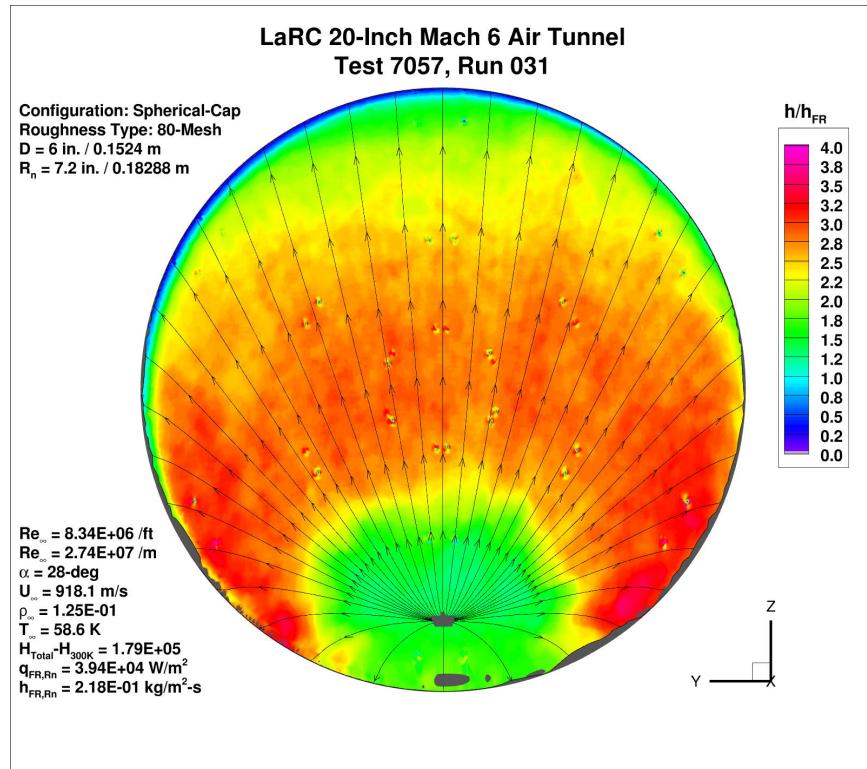


Figure 155. Test 7057, Run 31, $Re_{\infty} = 8.1 \times 10^6 / \text{ft}$, spherical-cap, 80-mesh.

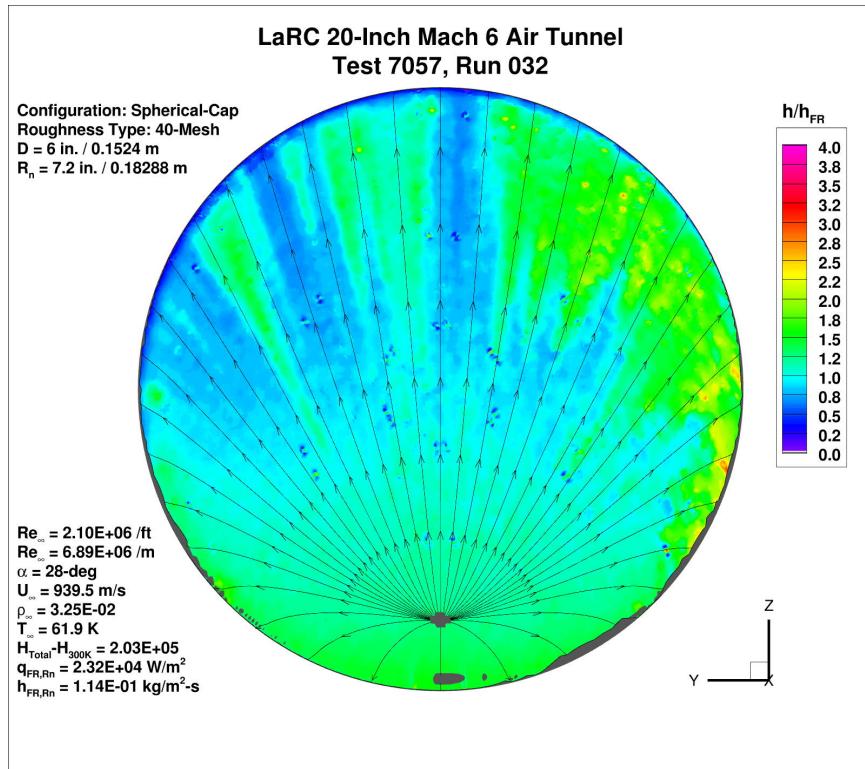


Figure 156. Test 7057, Run 32, $Re_{\infty} = 2.1 \times 10^6/\text{ft}$, spherical-cap, 40-mesh.

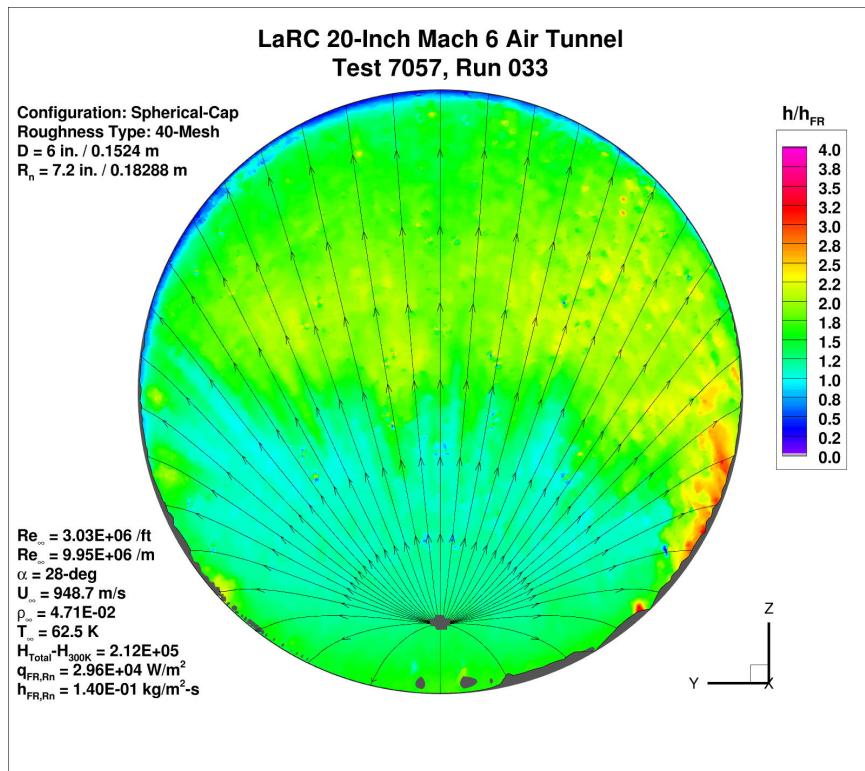


Figure 157. Test 7057, Run 33, $Re_{\infty} = 3.0 \times 10^6/\text{ft}$, spherical-cap, 40-mesh.

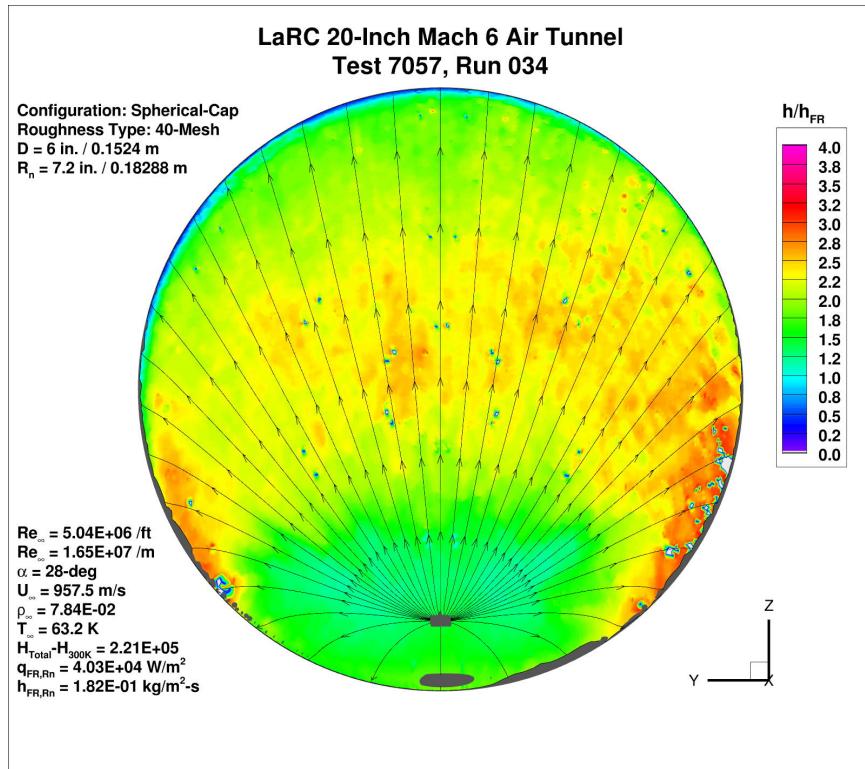


Figure 158. Test 7057, Run 34, $Re_{\infty} = 5.0 \times 10^6/\text{ft}$, spherical-cap, 40-mesh.

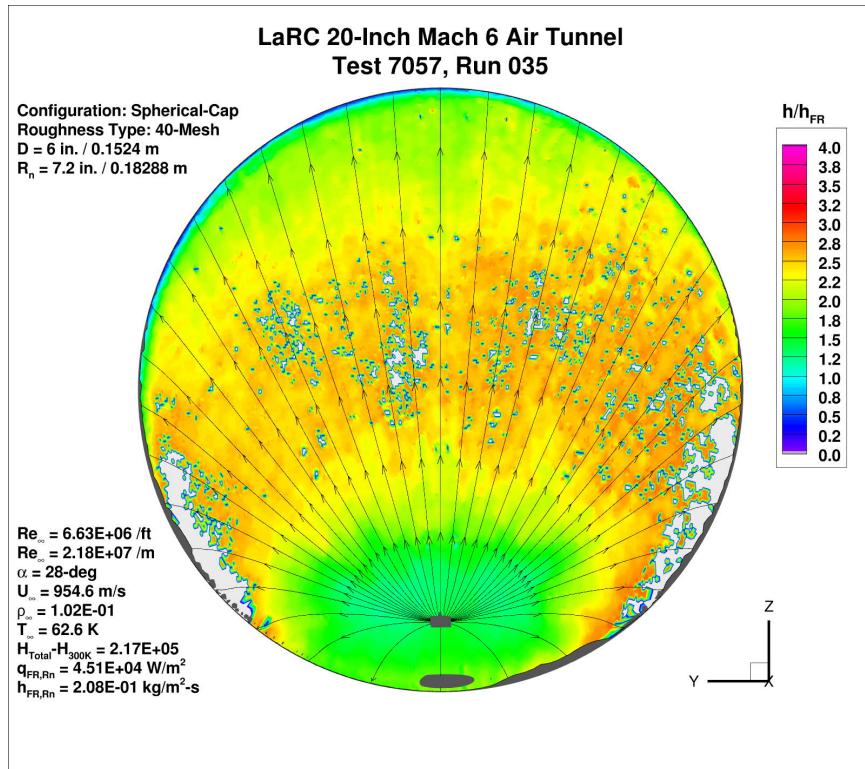


Figure 159. Test 7057, Run 35, $Re_{\infty} = 6.5 \times 10^6/\text{ft}$, spherical-cap, 40-mesh.

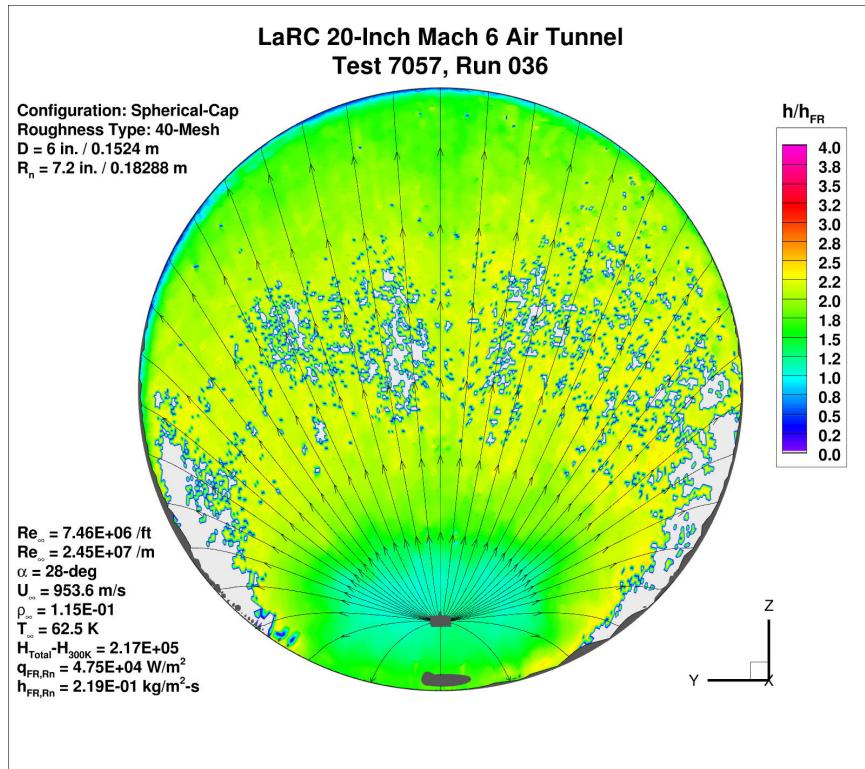


Figure 160. Test 7057, Run 36, $Re_{\infty} = 7.2 \times 10^6 / \text{ft}$, spherical-cap, 40-mesh.

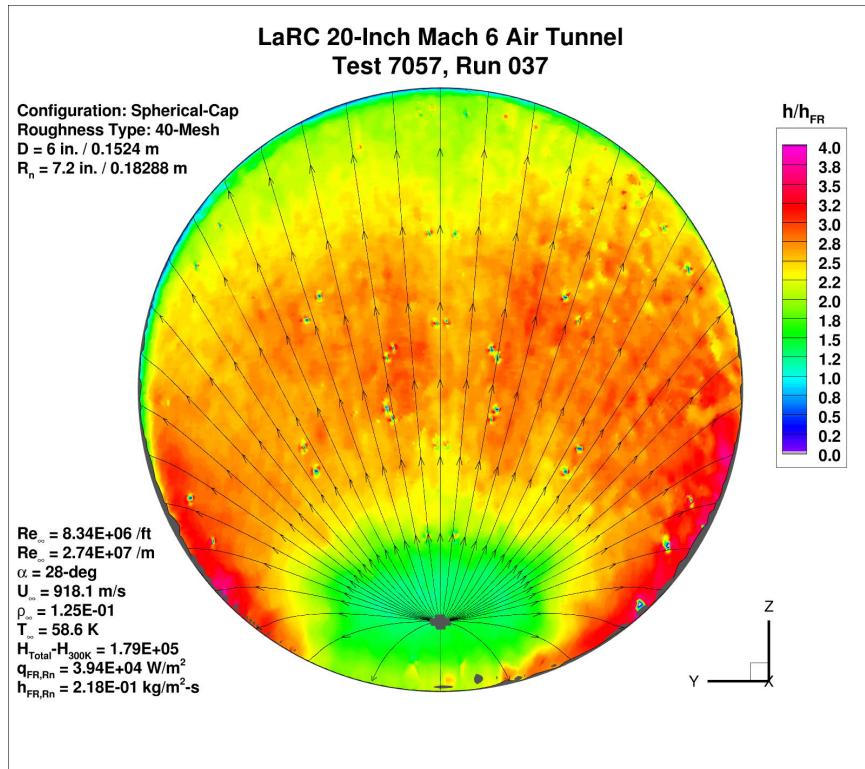


Figure 161. Test 7057, Run 37, $Re_{\infty} = 8.1 \times 10^6 / \text{ft}$, spherical-cap, 40-mesh.

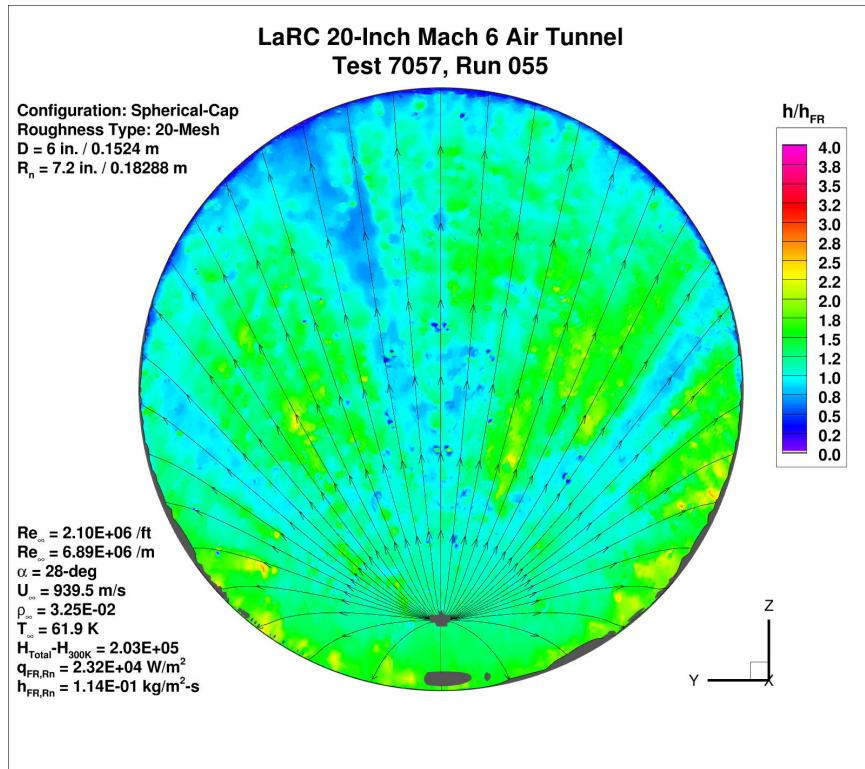


Figure 162. Test 7057, Run 55, Re_∞ = 2.1×10⁶/ft, spherical-cap, 20-mesh.

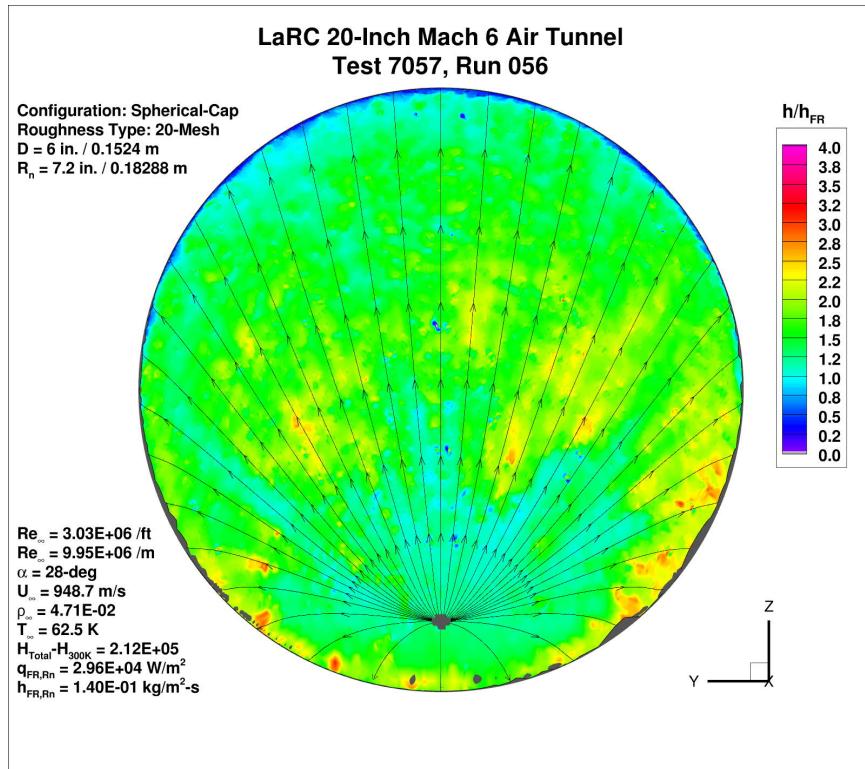


Figure 163. Test 7057, Run 56, Re_∞ = 3.0×10⁶/ft, spherical-cap, 20-mesh.

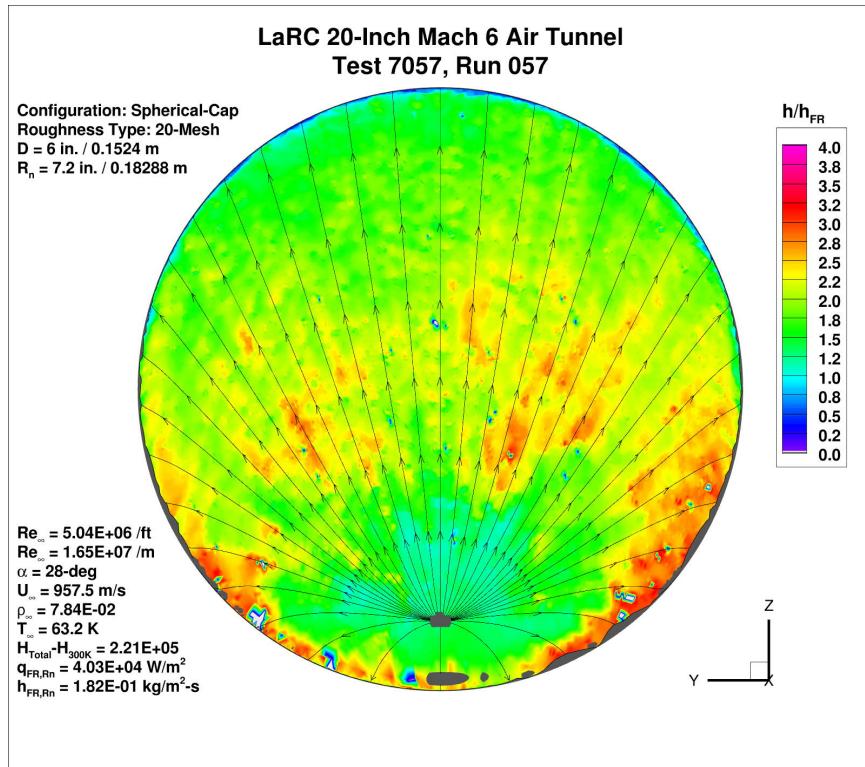


Figure 164. Test 7057, Run 57, $Re_{\infty} = 5.0 \times 10^6/\text{ft}$, spherical-cap, 20-mesh.

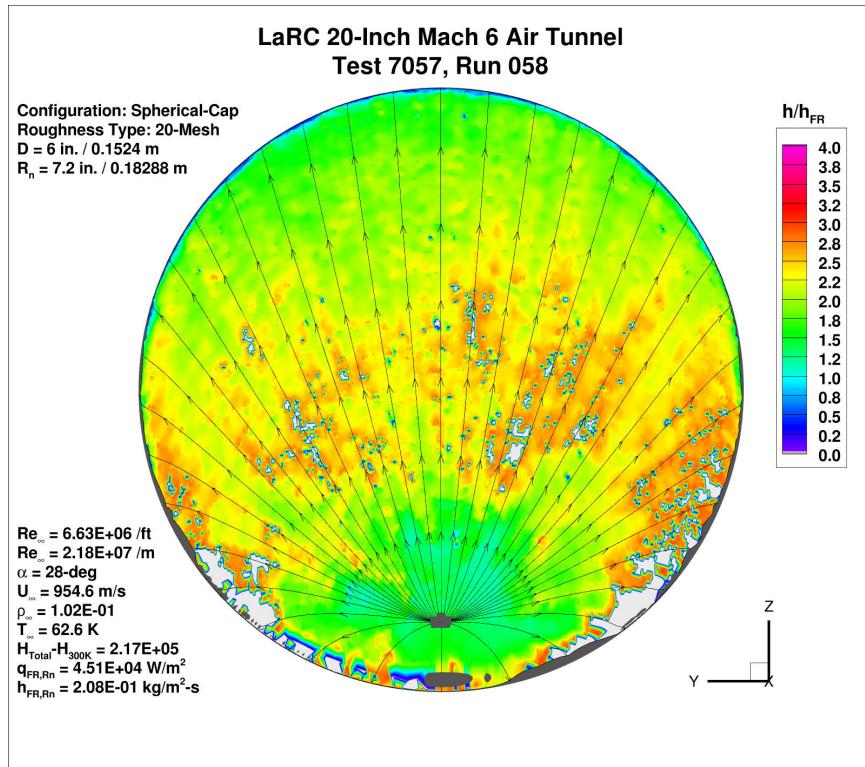


Figure 165. Test 7057, Run 58, $Re_{\infty} = 6.5 \times 10^6/\text{ft}$, spherical-cap, 20-mesh.

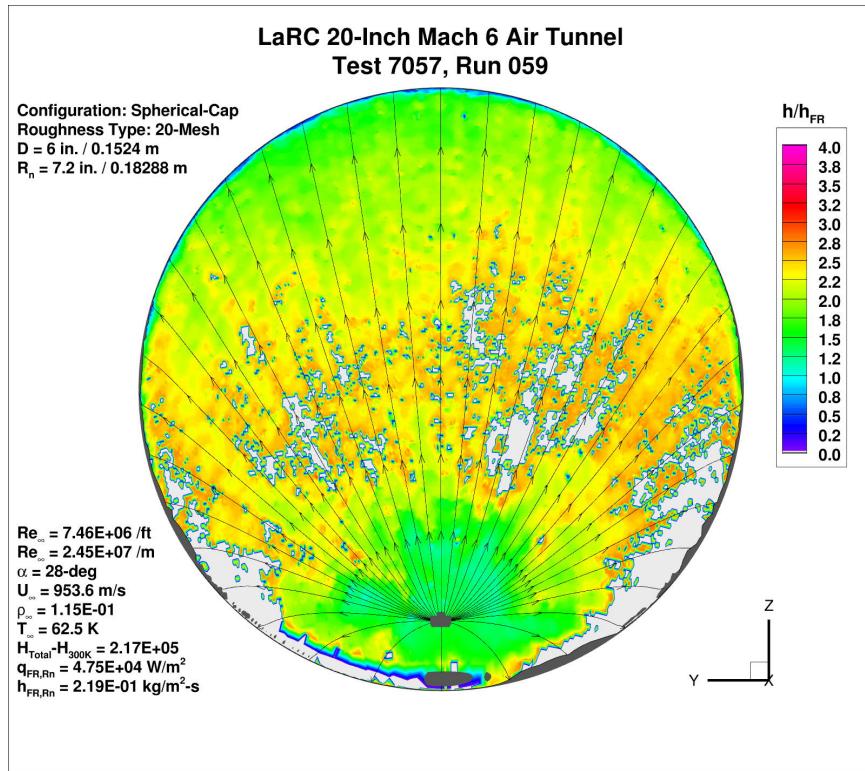


Figure 166. Test 7057, Run 59, Re_∞ = 7.2×10⁶/ft, spherical-cap, 20-mesh.

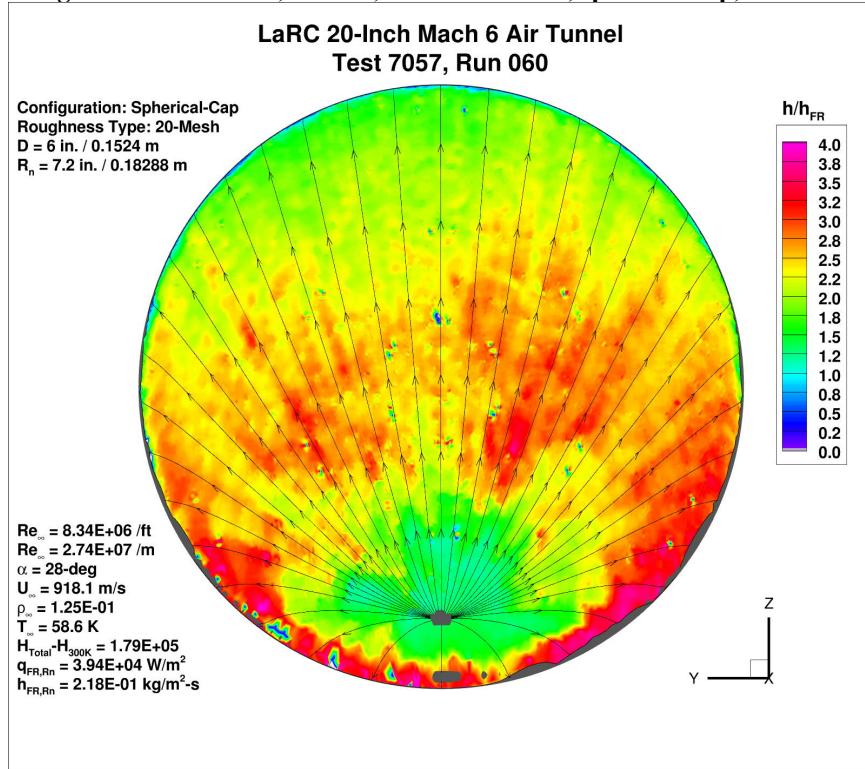


Figure 167. Test 7057, Run 60, Re_∞ = 8.1×10⁶/ft, spherical-cap, 20-mesh.

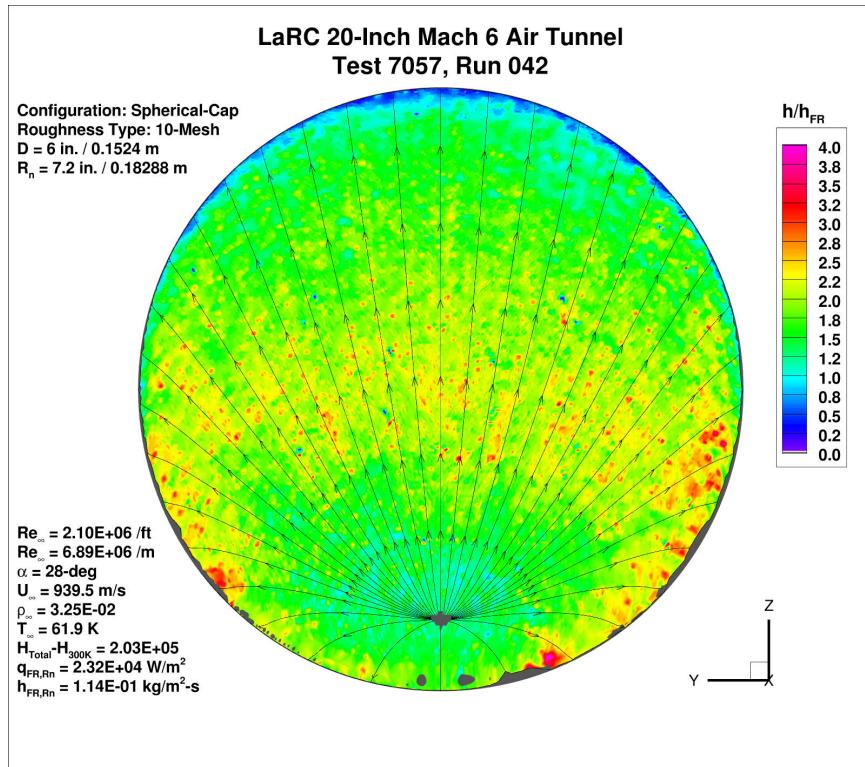


Figure 168. Test 7057, Run 42, $Re_{\infty} = 2.1 \times 10^6 / \text{ft}$, spherical-cap, 10-mesh.

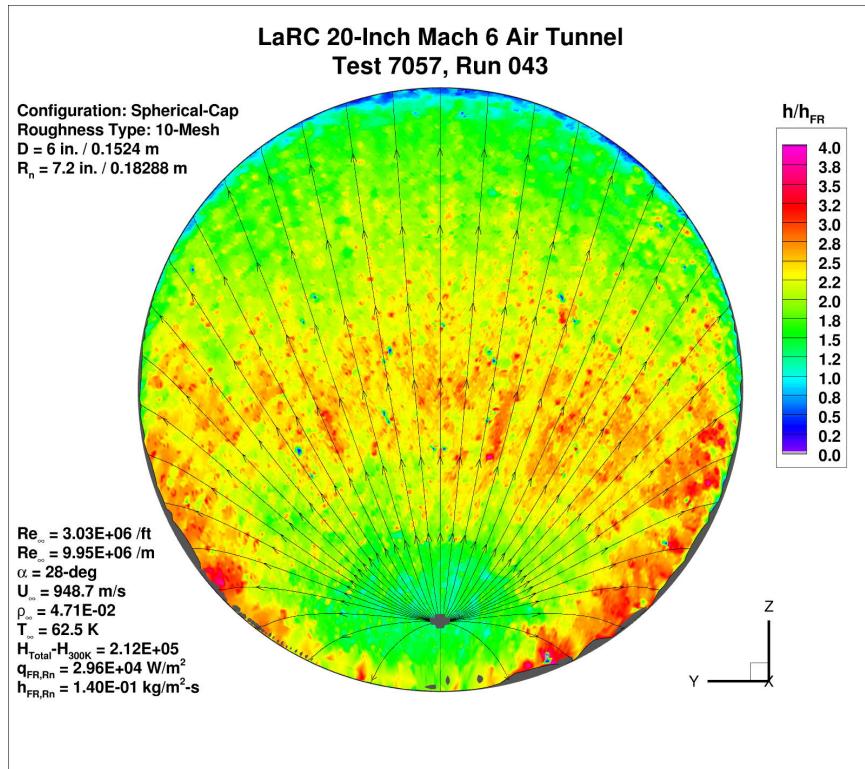


Figure 169. Test 7057, Run 43, $Re_{\infty} = 3.0 \times 10^6 / \text{ft}$, spherical-cap, 10-mesh.

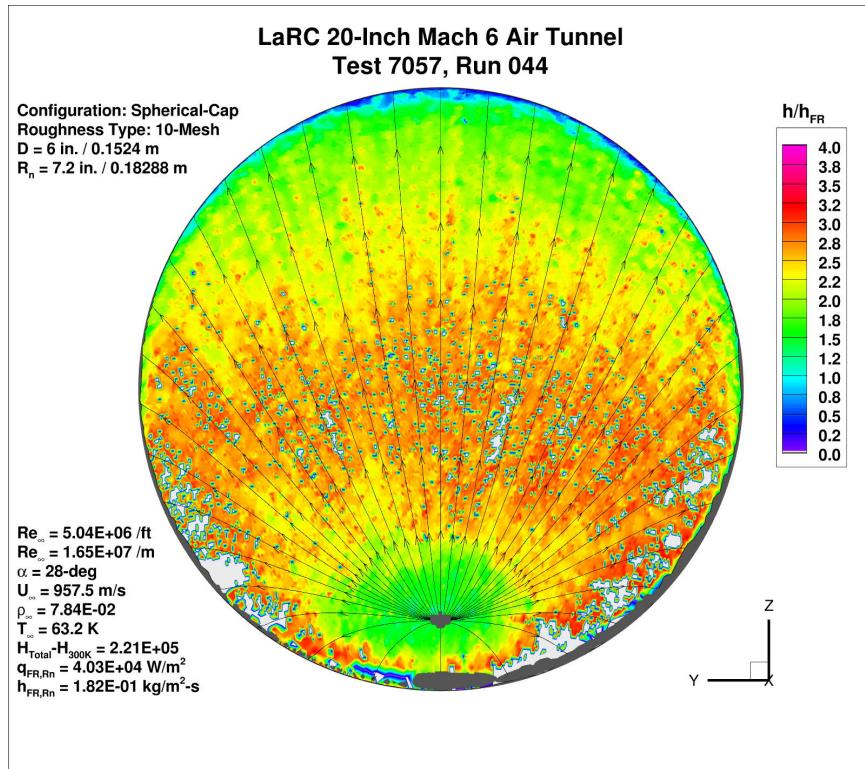


Figure 170. Test 7057, Run 44, Re_∞ = 5.0×10⁶/ft, spherical-cap, 10-mesh.

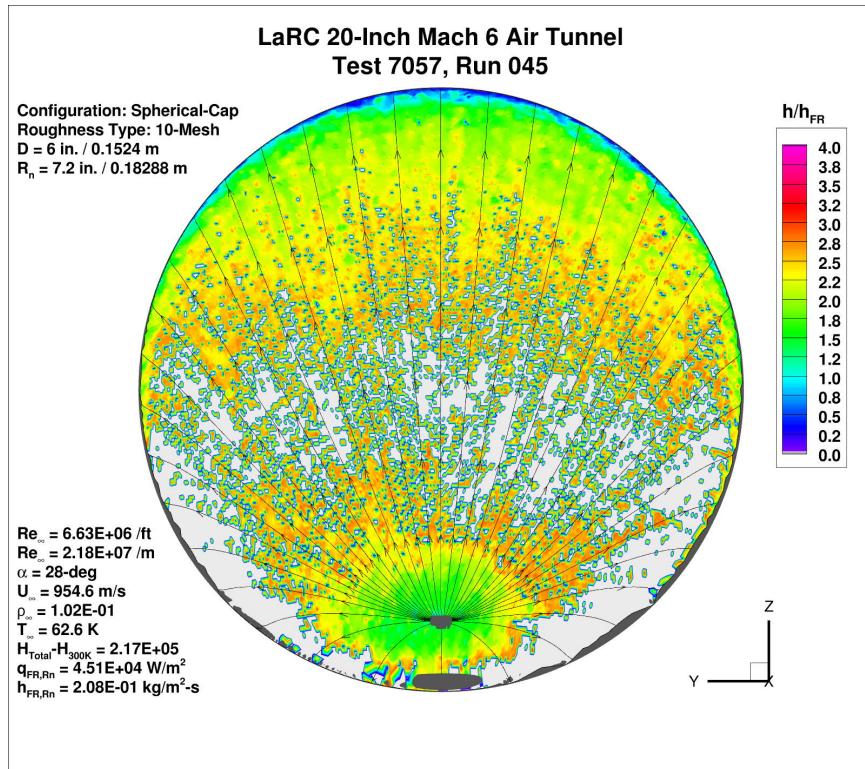


Figure 171. Test 7057, Run 45, Re_∞ = 6.5×10⁶/ft, spherical-cap, 10-mesh.

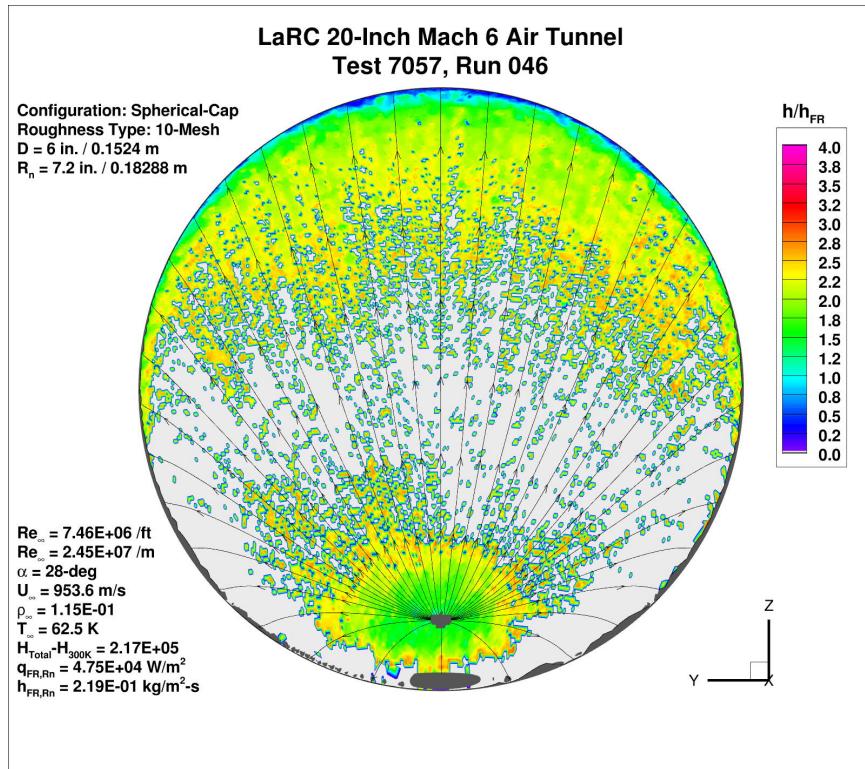


Figure 172. Test 7057, Run 46, Re_∞ = 7.2×10⁶/ft, spherical-cap, 10-mesh.

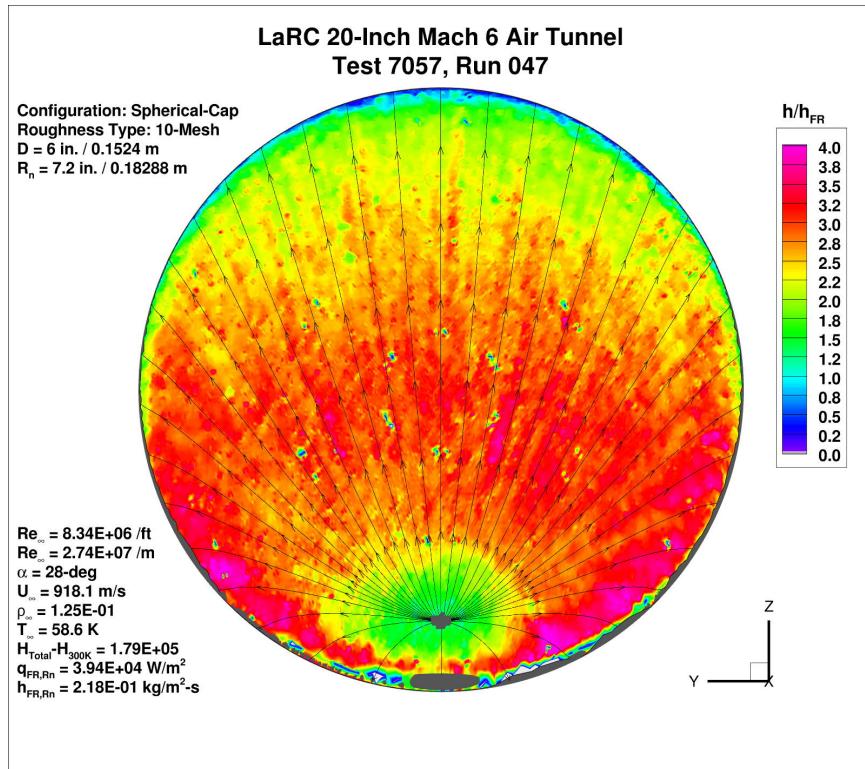


Figure 173. Test 7057, Run 47, Re_∞ = 8.1×10⁶/ft, spherical-cap, 10-mesh.