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2 **Precision of Spacecraft Doppler Tracking at Low Signal-to-Noise Ratios**

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4 **D. R. Buccino¹, K. Oudrhiri¹, M. Parisi¹, R.S. Park¹, E. Mazarico², P. Tortora^{3,4}, P.**
5 **Withers⁵, A. Genova⁶, M. Zannoni^{3,4}**

6 ¹ Jet Propulsion Laboratory, California Institute of Technology

7 ² Goddard Space Flight Center, National Aeronautics and Space Administration

8 ³ Department of Industrial Engineering, Alma Mater Studiorum - Università di Bologna, Italy

9 ⁴ Centro Interdipartimentale di Ricerca Industriale Aerospaziale, Alma Mater Studiorum -
10 Università di Bologna, Italy

11 ⁵ Boston University, Boston, MA

12 ⁶ Dipartimento di Ingegneria Aerospaziale e Astronautica, Università di Roma "La Sapienza,"
13 Rome, Italy

14
15 Corresponding author: Dustin R. Buccino (Dustin.R.Buccino@jpl.nasa.gov)

16
17 **Key Points:**

- 18 • Spacecraft Doppler measurements can be acquired at low signal-to-noise ratios with
19 open-loop receivers
- 20 • Experiments with Deep Space Network ground systems show Doppler data with a signal-
21 to-noise ratio as low as ~ 4 dB-Hz can be used
- 22 • Typical Doppler measurements at low SNR maintain a line-of-sight velocity precision of
23 < 0.1 mm/s at 60-second integration time
24

25 **Abstract**

26 The signal-to-noise ratio received at Earth is typically larger than 10 dB-Hz for radio tracking of
27 a spacecraft in deep space, allowing a phase-locked loop to execute reliable carrier tracking.
28 Recently, missions have been proposed to utilize low-gain antennas for Doppler tracking where
29 the signal-to-noise ratio may be at the single-digit (dB-Hz) level. In this work, we discuss
30 spacecraft Doppler tracking at these low signal-to-noise ratios through an analysis of thermal
31 noise on the radio link, results from ground testing with the Deep Space Network, and
32 demonstrations with an active spacecraft. We show that by utilizing an open-loop receiver to
33 capture the carrier signal and by applying post-processing techniques, radio data with the signal-
34 to-noise ratio as low as 4 dB-Hz can be used to derive Doppler data with precision sufficient for
35 tracking a spacecraft in deep space.

36

37 **Plain Language Summary**

38 Typical deep-space probes utilize High Gain Antennas for radio tracking that enable large signal-
39 to-noise ratios despite the long distances from Earth. When a Low Gain Antenna is used instead,
40 the signal-to-noise ratio is decreased, resulting in increased thermal noise and questions about the
41 ability to track. This work shows that signals can be tracked with the Deep Space Network's
42 Open Loop Receiver down to 4 dB-Hz with sufficient precision for radio science experiments.

43 **1 Introduction**

44 Radio links between spacecraft in deep space and Earth-based radio telescopes has been
45 used for decades to aid in the navigation of spacecraft, and to investigate planetary atmospheres,
46 ionospheres, surface properties, and interior structure (Asmar et al., 2019). By tracking the
47 carrier frequency of the radio signal sent by a planetary spacecraft, one can measure the Doppler
48 shift caused by the relative motion between the spacecraft and the ground antenna. This
49 measurement, referred to as the Doppler observable or carrier frequency observable, is widely
50 used in spacecraft navigation and radio science experiments.

51 The National Aeronautics and Space Administration's (NASA) Deep Space Network
52 (DSN) is a set of large-aperture (34-meter or 70-meter) ground stations used for communication
53 and navigation of spacecraft beyond the orbit of the Moon. A typical application of the radio link
54 involves a sizable ($>2\text{m}$) High-Gain Antenna (HGA) on the spacecraft transmitting to the large-
55 aperture dishes of the DSN equipped with cryogenically-cooled low-noise amplifiers. This often
56 results in a radio carrier signal-to-noise ratio (SNR) of >10 dB-Hz with typical values ~ 40 dB-
57 Hz. The DSN's closed-loop receiver utilizes a phase-locked loop (PLL) to track the carrier
58 signal, requiring an SNR >10 dB-Hz with a 1 Hz loop-bandwidth to track (O'Dea and Kinman,
59 2019).

60 Recently, missions have proposed using wide-beamwidth Low Gain Antennas (LGAs) to
61 allow tracking a spacecraft's radio carrier and conduct radio science observations that could
62 otherwise not be accomplished with a body-fixed HGA due to pointing constraints onboard the
63 spacecraft (e.g., Mazarico et al., 2021 and Dehant et al., 2020). The resulting SNRs would be in
64 the low single-digits (dB-Hz). This low-SNR precludes the use of the typical closed-loop
65 receiver. Thus, any data collected at these levels would need to rely on the DSN's Open Loop
66 Receiver (OLR) which is used in numerous radio science, radio astronomy, very long baseline
67 interferometry, delta-differential one-way ranging, and engineering applications (Bedrossian &

68 Rogstad, 2021). The OLR records the downconverted antenna voltages as In-phase and
69 Quadrature (IQ) samples. The receiving team would process these IQ samples into typical
70 tracking observables such as the Doppler observable. This article demonstrates the capabilities of
71 the DSN's Open Loop Receiver (OLR) to capture, detect, and process a precise Doppler
72 observable at these low-SNR levels (<10 dB-Hz).

73 **2 Previous Use of Low-SNR Signals**

74 Previously, deep-space missions have conducted radio Doppler tracking using LGAs at
75 low SNR. However, typically these data are only used as the prime data set if either the SNR is
76 large enough to track with the closed-loop receiver reliably, or as a part of a measurement that
77 requires the OLR such as radio occultation or bistatic scattering experiments. LGAs have
78 previously been utilized for gravity science measurements; recent examples include Juno at
79 Jupiter, Cassini at Titan, Dawn at Vesta and Ceres, and MESSENGER at Mercury.

80 Juno's gravity science investigation primarily utilizes dual-frequency links through an
81 HGA, but the mission plan called for closest-approach periods with off-Earth pointing. During
82 these, it was decided to track the spacecraft at a single-frequency with either a Medium Gain
83 Antenna (MGA) or LGA to get additional Doppler tracking to supplement the prime gravity
84 science data. An example of this is Juno's PJ-04 flyby on 07 February 2017. An average SNR of
85 10.7 dB-Hz on the 34-meter DSS-25 antenna was observed. The closed-loop receiver could not
86 track the signal due to the MGA's low SNR and Doppler spin signatures (Buccino et al., 2018b).

87 Cassini conducted one flyby of Titan while utilizing the spacecraft's LGA on 16 March
88 2015 ("T110"). That experiment had been prepared by a previous test, carried out on 15 August
89 2010, which showed the feasibility of gravity measurements using the LGA (Barbaglio et al.,
90 2012). Due to the significant Earth-Saturn distance and low link margin calculated by the
91 telecommunications engineers, in order to acquire a lock on the spacecraft a special maneuver
92 was added to bring the LGA antenna closer to boresight only during uplink acquisition to the
93 spacecraft transponder. After the coherent link was established, the spacecraft transponder
94 retained the lock. During the flyby, the SNR was variable between ~10-20 dB-Hz due to
95 spacecraft pointing and scanning maneuvers required by other instruments (Barbunis, 2016). The
96 data acquired during T110 were successfully used to retrieve the gravity field and ephemerides
97 of Titan (Durante et al., 2019; Lainey et al., 2020).

98 Dawn utilized radio tracking to determine the gravity fields of Vesta and Ceres. The
99 HGA tracking was supplemented with LGA tracking to increase the amount of Doppler data for
100 the gravity field investigation. The SNR on the LGAs was predicted to be ~22-24 dB-Hz at a 70-
101 meter DSN antenna, so that the corresponding Doppler quality was adequate for gravity recovery
102 (Konopoliv et al., 2011).

103 The use of LGAs in radio occultations is rarer. MAVEN has a limited number of
104 opportunities to conduct carrier-only radio occultation observations using the HGA, which
105 requires an Earth-pointed orientation and no telemetry on the downlink signal. The number of
106 observations was increased by a factor of 2-3 by also conducting telemetry-on observations using
107 the LGA in which Earth-pointing requirements are significantly relaxed and telemetry can be
108 included on the downlink signal. Furthermore, as LGA observations are not confined to
109 MAVEN's biweekly dedicated communication periods, multiple observations can be acquired
110 daily in support of investigations of time-variable phenomena in the environment of Mars (Felici

111 et al., 2020). Withers et al. (2020) generated vertical profiles of ionospheric electron density
 112 from these observations and reported average uncertainties of 1900 cm⁻³ and 2700 cm⁻³ for
 113 HGA and LGA observations, respectively. These are proportional to the average uncertainties in
 114 the frequency residual. Galileo also performed gravity and occultation experiments on its LGA
 115 after the deployment failure of the HGA, with limitations (Hinson et al., 1997).

116 **3 Thermal Noise**

117 Doppler noise budgets for deep-space probes are well-understood and quantified (Asmar
 118 et al., 2005; Iess et al., 2014). Typical noise sources for precision Doppler measurements include
 119 instrumentation noise (frequency and timing, ground electronics, antenna mechanical noise,
 120 spacecraft electronics/oscillators, thermal noise) and propagation noise (solar and planetary
 121 plasma, Earth ionosphere, Earth troposphere). Using a coherent two-way link, one where the
 122 downlink signal is phase-coherent with the uplink signal, the frequency reference is provided by
 123 the highly-stable hydrogen maser clocks from the DSN frequency and timing system and
 124 onboard spacecraft oscillator noise is not contributing. This is the type of measurement (coherent
 125 two-way) typically performed for gravity science studies where signal stability is a crucial factor.
 126 In this configuration, dominating noise sources are typically solar plasma and Earth troposphere
 127 scintillation when using HGAs.

128 Thermal noise is present in the radio Doppler measurements for a radio link having a
 129 finite SNR. During coherent tracking, thermal noise occurs at both the spacecraft transponder
 130 and at the DSN receiver. The thermal noise at the spacecraft transponder can be ignored if the
 131 SNR at the spacecraft is larger than the signal-to-noise ratio at the DSN receiver, as is typically
 132 the case due to the larger transmitter power of the DSN (~18 kW) versus the spacecraft (~20-40
 133 W). This assumption should be checked in each application. With this assumption the thermal
 134 noise (σ_f) can be calculated as (derived from O’Dea and Kinman, 2019 or Buccino et al., 2022):

$$135 \quad \sigma_f = \frac{\sqrt{2B_L}}{2\pi T_c \sqrt{P_c/N_0}} \quad (1)$$

136

137 where B_L is the bandwidth of the phase-locked loop tracking the signal, T_c is the count-
 138 time of the measurement, and P_c/N_0 is the carrier signal-to-noise spectral density ratio on a linear
 139 scale. Thus, when the SNR is low, as in the case of using LGAs or a low count-time (in some
 140 radio science applications, such as radio occultations), thermal noise becomes a dominating noise
 141 source in the error budget. Thermal noise is not a function of carrier frequency.

142 **4 Open-Loop Data Processing**

143 Various techniques exist to extract Doppler observables from open-loop data, which are
 144 recorded as In-phase and Quadrature samples of the downconverted antenna voltage at the DSN’s
 145 OLR. These include fast-Fourier transforms (FFTs), spectral-optimized FFTs, Doppler-
 146 compensated FFTs, and software PLLs. FFTs are typically used where the signal strength or
 147 frequency is highly variable, such as in radio occultations. PLLs are more desirable for gravity
 148 observations, as they give higher precision measurements than FFTs because the bin width of the
 149 FFT does limit the measurement. However, PLLs require a more stable frequency to maintain
 150 lock. Using a PLL is analogous to measurements that come from the closed-loop receiver.

151 Typically, radio science observations use a spectral-optimized FFT (Paik and Asmar,
 152 2010, Togni et al., 2021) as used in MAVEN occultations for example (Withers et al, 2018), or a
 153 PLL (Densmore, 1988) as used in Juno gravity science (Buccino et al., 2018a). This study
 154 focuses on the utilization of the PLL for precision Doppler measurements. The PLL employed is
 155 a second-order loop of Densmore (1988) integrated into the Radio Science Validation &
 156 Processing (RSVP) toolkit, successor to the STBLTY toolkit used for Galileo radio science
 157 processing (Morabito and Asmar, 1995). The PLL consists of a phase detector and loop filter, a
 158 low pass filter, and an integrator that runs over the open-loop IQ samples. The PLL is
 159 configurable to produce an output phase at a user-specified integration time and loop-bandwidth.
 160 The Doppler observable (f) at time t is then produced as a differenced-phase measurement of
 161 adjacent phase points (ϕ) over the count-time (T_c):

$$162 \quad f(t) = -\frac{\phi\left(t - \frac{T_c}{2}\right) - \phi\left(t + \frac{T_c}{2}\right)}{T_c} \quad (2)$$

163
 164 Three important factors must be considered when running a PLL to produce Doppler
 165 observables, particularly at low SNR: maintain a loop-SNR above the tracking threshold,
 166 minimize static phase error (SPE) in the loop, and optimize the thermal noise by producing
 167 unique phase points. The DSN closed-loop receiver can utilize a second-order loop or third-order
 168 loop. In typical DSN operations a third-order loop is utilized. The second-order loop is analogous
 169 in behavior to the second-order PLL utilized by this study - the same used in previous radio
 170 science investigations, e.g. Juno, Galileo - and thus those constraints may be applied here.

171 Firstly, the loop-SNR must be above the tracking threshold to acquire and maintain lock
 172 on the carrier signal in the recording. The PLL's loop-SNR (ρ_L) tracking threshold is
 173 approximately 7-10 dB. The loop-SNR is defined as the downlink carrier SNR divided by the
 174 loop bandwidth (O'Dea and Kinman, 2019):

$$175 \quad \rho_L = \frac{P_c}{N_0} \cdot \frac{1}{B_L} \quad (3)$$

176
 177 Secondly, to maintain this lock through regions where the frequency is changing rapidly
 178 or is unknown, the SPE must be kept below the lock threshold. In the application of the DSN's
 179 closed-loop tracking second-order loop (O'Dea and Kinman, 2019), the SPE must be kept below
 180 30 degrees and is a function of the loop bandwidth, Doppler rate (α), and Doppler acceleration
 181 (β), which are typically caused by uncertainties in the predicted spacecraft's trajectory:

$$182 \quad SPE(\text{rad}) = \frac{25\pi\alpha}{32 B_L^2} + \left[\frac{25\pi\beta}{32B_L^2} t - \frac{125\pi\beta}{128B_L^3} \right] \quad (4)$$

183
 184 Finally, to produce phase points and utilize all available SNR in the recording, a
 185 constraint is placed that the Doppler integration time is proportional to the loop bandwidth
 186 (Folkner and Border, 2013; Buccino et al. 2022):

$$B_L = \frac{1}{2T_c} \quad (5)$$

187

188

189 Each of these considerations places constraints on the configuration of the PLL. For
 190 example, to track a signal at 4 dB-Hz, the loop bandwidth must be 0.4 Hz to maintain the loop-
 191 SNR tracking threshold (Equation 3). Additionally, to maintain unique phase points and optimize
 192 thermal noise, the integration time should be one point every 1.25 seconds (Equation 5). The
 193 tolerable Doppler rate and acceleration would be limited per Equation 4 by the loop bandwidth
 194 (for 0.4 Hz loop bandwidth, the limit to Doppler rate, assuming $\beta = 0 \text{ Hz/s}^2$ is $\alpha < 0.05 \text{ Hz/s}$)
 195 Residual Doppler rate and Doppler acceleration can be decreased or negated by phase-steering or
 196 counter-rotation of the IQ samples from the OLR using updated models of the spacecraft
 197 trajectory (Buccino et al., 2018a). It is often desirable to have an integration rate of the PLL
 198 divisible by an integer of 60, then 60-second Doppler points can be compressed from the output
 199 of the PLL with existing software. In this case, the loop bandwidth is set to 0.25 Hz and the
 200 integration time is 2 seconds, allowing for a signal at 4 dB-Hz to have a loop SNR above 10 dB
 201 and a divisible-by-60 integration time. The maximum tolerable residual Doppler rate and
 202 acceleration would be reduced (a 0.25 Hz loop bandwidth would limit the Doppler rate $\alpha < 0.013$
 203 Hz/s , assuming $\beta = 0 \text{ Hz/s}^2$).

204 5 Ground System Demonstration

205 A test was conducted with the DSN to demonstrate the capability to make Doppler
 206 measurements at very low SNR. The DSN's DSS-25 station in Goldstone, California was placed
 207 in the same configuration as is done during a normal tracking pass where coherent Doppler
 208 measurements are taken, where instead of a spacecraft, a device called the "test translator" was
 209 placed at the final stage to provide the phase coherent downlink from the generated uplink signal.
 210 This configuration is shown in Figure 1. The X-band uplink exciter generated a 7170 MHz
 211 signal, derived from the DSN's hydrogen maser-based Frequency and Timing Subsystem (FTS).
 212 The X-band transmitter amplifies this signal through a diplexer. Typically, this signal would then
 213 be transmitted *out the horn* to the desired spacecraft. Instead, the test translator phase-coherently
 214 receives the uplink signal and multiplies it by a turnaround ratio of 880/749, typical of most
 215 deep-space X-band transponders. The test translator is equipped with a precise variable
 216 attenuator, which reduces the signal power to a level typical of reception. Now at a frequency of
 217 approximately 8424 MHz, this signal is sent back through the diplexer and into the
 218 cryogenically-cooled Low-Noise Amplifier (LNA). The signal is downconverted from a radio
 219 frequency (RF) level to an intermediate frequency (IF) level of approximately 300 MHz. The
 220 signal is then digitized through the DSN common platform: IF gain control, IF digitization, and
 221 finally flows into the digital IF switch where the Open-Loop Receiver (OLR) records the signal
 222 with In-Phase and Quadrature complex samples (Bedrossian & Rogstad, 2021).

223 During this test, the variable attenuator of the test translator artificially reduced the signal
 224 power such that the received SNR at the open-loop receiver was between 2 dB-Hz and 12 dB-
 225 Hz, in steps of 2 dB-Hz. Approximately 15 minutes of data were collected at 2 dB-Hz, 6 dB-Hz,
 226 8 dB-Hz, 10 dB-Hz and 12 dB-Hz, and 30 minutes were collected at 4 dB-Hz. The OLR was
 227 configured to sample the signal at a rate of 1 kHz with 16-bit samples.

228 SNR measurements were made with an FFT estimation with 30 seconds of integration
 229 time. The signal was able to be tracked at each targeted SNR to within ± 0.2 dB-Hz, though the
 230 lowest targeted SNR of 2 dB-Hz was, in actuality, measured at 2.6 dB-Hz. Variations from the
 231 target are caused by fluctuations in transmitter power output. The 2.6 dB-Hz minimum led to the
 232 selection of a PLL loop-bandwidth (B_L) of 0.25 Hz. To optimize the thermal noise in the Doppler
 233 measurement, the integration time (T_c) was set to 2 seconds ($T_c = 1/2B_L$). The resulting SNR
 234 measurements and Doppler residuals (observed minus predicted) are shown in Figure 2. The
 235 root-mean-square (RMS) of the Doppler residuals was 34 mHz at 4 dB-Hz and 42 mHz at 2.6
 236 dB-Hz. With standard frequency compression of the Doppler residuals to 60-second integration
 237 time, the resulting RMS was 1.4 mHz at 4 dB-Hz (in units of range-rate, using the conversion
 238 factor $c/2f_c$ where c is the speed of light and f_c is the carrier frequency, is 0.025 mm/s). For
 239 phase-noise dominated measurements, using phase compression over standard frequency
 240 compression will result in higher precision measurements (Buccino et al., 2022).

241 As expected per thermal noise theory (Section 2), as the SNR is reduced the noise in the
 242 Doppler measurement increases. Since the test was conducted on the ground, no propagation
 243 noise or antenna mechanical noise is introduced. Thus, the only remaining noise sources are
 244 thermal noise, DSN instrumental noise, and DSN FTS noise. DSN instrumental and FTS noise is
 245 negligibly small compared to thermal noise in this case. Therefore, the measured noise can be
 246 directly compared with the model for thermal noise in Eq. (1) and is shown in Figure 3. The
 247 measured performance using the digital PLL is comparable to the thermal noise model, and
 248 differences can be explained by the uncertainty in the SNR power measurements (Figure 2).

249 The Allan deviation (Allan, 1966) is a measure of fractional frequency stability,
 250 commonly used in clock and radio tracking performance. Using the 30 minutes of data collected
 251 at 4 dB-Hz in the ground system demonstration, the Allan deviation was computed from the
 252 residual frequencies and is plotted as a function of integration time τ in Figure 4 ($\sim 1 \times 10^{-12}$ at 10-
 253 sec integration time). The Allan deviation (σ_y) is proportional to white phase noise characterized
 254 by a slope of τ^{-1} as expected since this test was conducted in the absence of spacecraft
 255 instrumental noise and propagation noise.

256 **6 Demonstration with Spacecraft**

257 Although there have been various uses of low-SNR signals (Section 2), to the best of our
 258 knowledge none have yet reached the low single digits (~ 4 dB-Hz). The ground testing (Section
 259 4) demonstrated the feasibility of such tracking. In order to test with the full link, 34-meter DSN
 260 antennas were scheduled concurrently to a primary spacecraft track on a different antenna (a
 261 configuration called a “shadow track”). The antenna is pointed off of boresight to reduce the
 262 SNR artificially. By taking advantage of the beam pattern of the antenna, the SNR can be
 263 reduced to a targeted quantity for short periods. Over longer durations, mechanical loading on
 264 the antenna changes the antenna pointing, causing the targeted SNR to vary. Nonetheless, this
 265 technique can be applied to prove the feasibility of low-SNR tracking with the full link
 266 configuration, including propagation noise, without changing the spacecraft’s command
 267 sequences. Two such tests were conducted.

268 The first utilized Juno spacecraft (Asmar et al., 2017) tracking in a coherent three-way
 269 mode (DSS-25 was the primary transmit antenna and DSS-35 was the secondary antenna used
 270 for the test). To reduce the SNR from the nominal level of ~ 40 dB-Hz to ~ 4 dB-Hz, a pointing
 271 offset of 71-72 millidegrees was used, placing it near the null between the primary beam and the

272 first side lobe. Two data segments were analyzed: the first with an SNR of 3 dB-Hz and the
273 second with an SNR of 4.8 dB-Hz (averaged over the duration of the data span). The segments
274 were 1 millidegree apart in pointing. Processing the open-loop data using the same configuration
275 of the PLL from Section 4 ($B_L = 0.25$ Hz; $T_c = 2$ sec) could be performed, though the second
276 data segment at ~ 3 dB-Hz had a phase cycle slips where the PLL briefly lost lock (causing a loss
277 of ~ 28 seconds of data out of the 12 minute data span; the Doppler measurements affected by the
278 cycle slips are not used in the statistical calculations that follow). These cycle slips, although
279 unimportant for the statistics reported in this work, may be significant if they occurred during a
280 radio science experiment and may have a loss of information depending on the timing where
281 they occurred. The measured Doppler residual RMS at 2-second integration time (Figure 5) was
282 88 mHz for the first segment at 3 dB-Hz and 50 mHz for the second segment at 4.8 dB-Hz. This
283 differs from the ground-only demonstration due to additional propagation noise sources (e.g.
284 Earth troposphere, Earth ionosphere, and solar plasma). After data compression (frequency
285 averaging) to 60-second integration time, the resulting RMS was 2.5 mHz at 3 dB-Hz and 1.2
286 mHz at 4.8 dB-Hz (the equivalent range-rate RMS is 0.045 mm/s and 0.021 mm/s, respectively).

287 The second utilized the New Horizons spacecraft (Tyler et al., 2009) tracking in a non-
288 coherent one-way mode with DSS-25 receiving (the New Horizons Ultra Stable Oscillator
289 provided the frequency reference). Utilizing an antenna pointing offset of 60 millidegrees
290 resulted in an SNR of 4.5 dB-Hz. Again, with the same PLL configuration ($B_L = 0.25$ Hz; $T_c = 2$
291 sec), the Doppler residual RMS was measured to be 32 mHz at a 2-second integration time
292 (Figure 6). The data collected were less than 5 minutes in duration, yielding less than 5 Doppler
293 points at a 60-second integration time. Therefore 60-second integration statistics are not reported
294 due to the low number of points. Note this is improved over the Juno result (50 mHz at 4.8 dB-
295 Hz), however, this is nearly entirely due to the measurement geometries where Juno was affected
296 by higher levels of plasma noise during data collection.

297 7 Conclusion

298 Doppler observable measurements are used for numerous applications. In this work, we
299 have demonstrated that the same measurements typically taken at high SNR can also be made at
300 low-SNR levels by utilizing the Open-Loop Receiver of the Deep Space Network at the cost of
301 increased thermal noise in the link. Signal processing of open-loop data with a phase-locked loop
302 was used to generate Doppler observable measurements. Ground testing demonstrated that
303 reliable tracking was achieved at an SNR of ~ 4 dB-Hz. Doppler measurement noise was
304 consistent with thermal noise theory. Flight demonstrations with the Juno spacecraft (coherent
305 tracking) and New Horizons (noncoherent 1-way tracking with USO reference) confirmed the
306 tracking ability using a full link configuration with a spacecraft. When signal levels were
307 reduced below 4 dB-Hz with flight demonstration, tracking became less reliable and cycle slips
308 occurred. For the 60-second integration time, Doppler measurement noise was < 0.1 mm/s in units
309 of range-rate. For applications of this work, it is crucial to account for all potential noise sources
310 (e.g., variable solar plasma noise, Earth troposphere, antenna mechanical), as was seen in
311 spacecraft demonstrations where Doppler measurement precision was variable between the two
312 tests. Furthermore, it is important to consider the effect of residual Doppler rate and acceleration,
313 primarily function of how well the spacecraft trajectory is predicted to be known, in applications
314 of low-SNR signal tracking. Lower tracking bandwidths needed to maintain lock also require
315 lower residual Doppler rate and acceleration. Although lower Doppler rate and accelerations may
316 be required to precisely track, the science information contained within the data may still contain

317 useful information for the analysis of gravitational fields, ionosphere and atmosphere via radio
318 occultations, and further improving the spacecraft's trajectory.

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334 **Data Availability Statement**

335 The open-loop test data and processed Doppler observable measurements are provided in a
336 corresponding dataset with this publication on Zenodo (Buccino, 2023).

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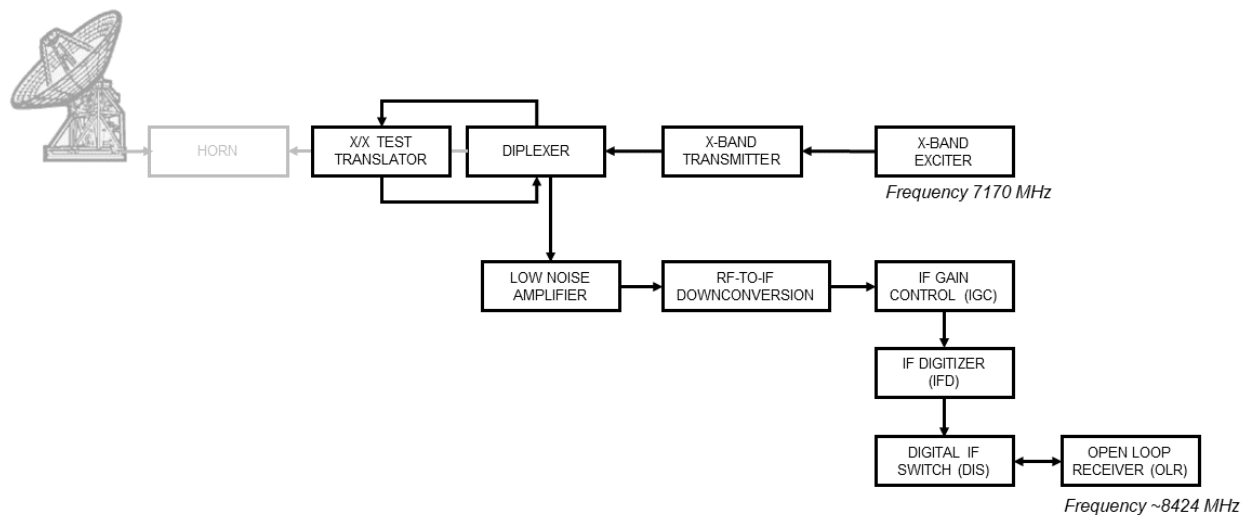
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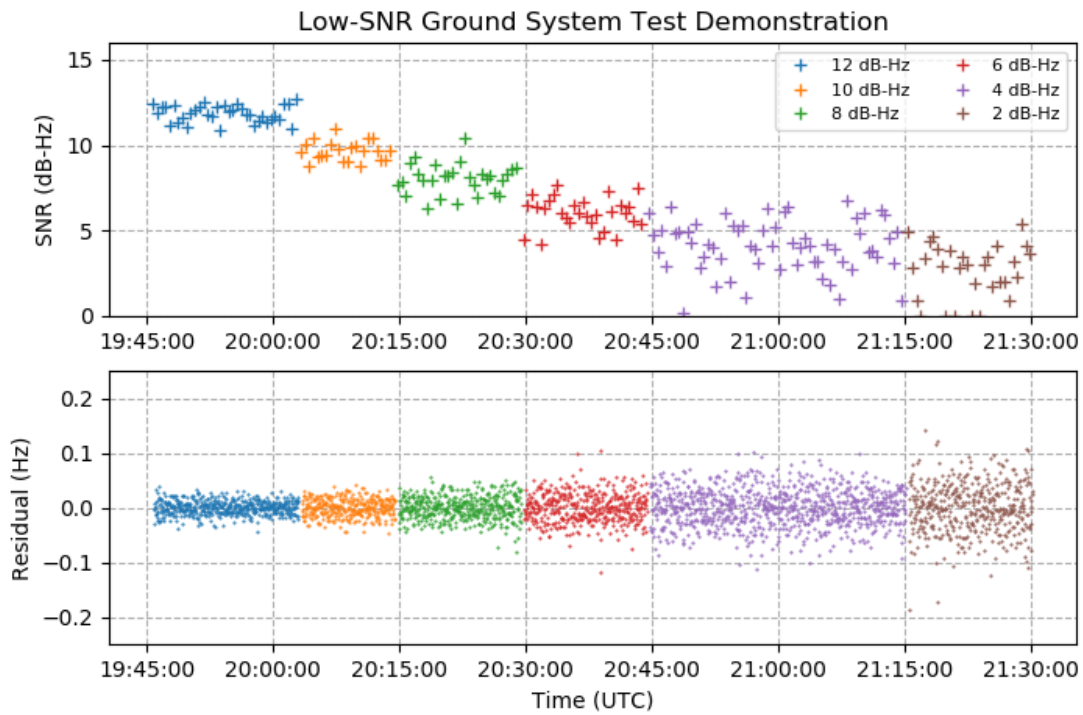


415

416 **Figure 1.** Configuration of the DSS-25 antenna in Goldstone, California to demonstrate Doppler
 417 tracking capabilities at very-low SNR. In lieu of a spacecraft, the test translator device with a
 418 precise variable attenuator was utilized to convert the uplink frequency to a downlink frequency
 419 and adjust the signal power level to the desired levels of 2 dB-Hz to 12 dB-Hz.

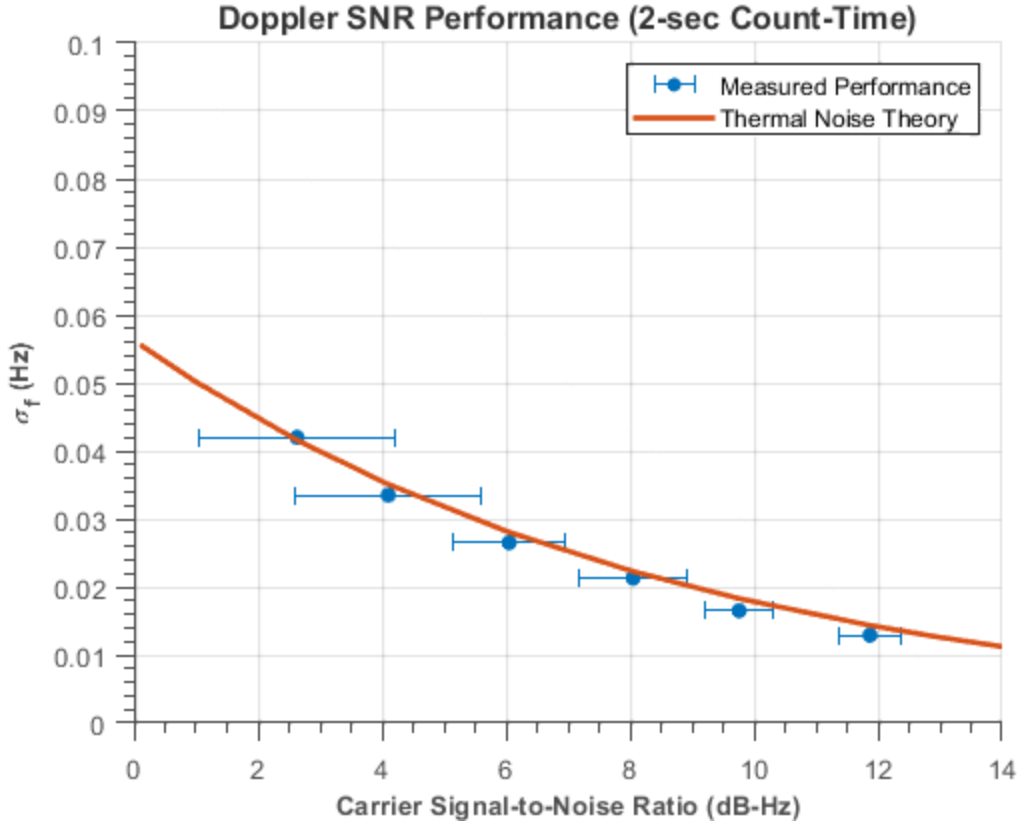
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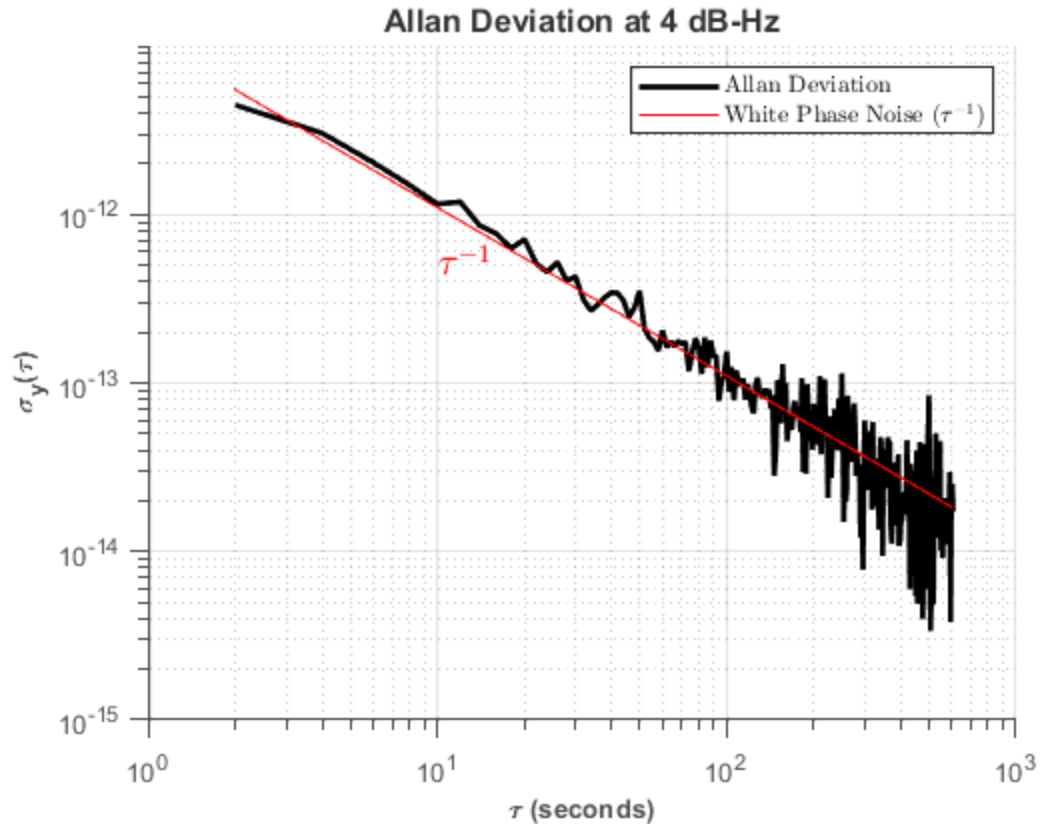
422

423 **Figure 2.** Measured SNR (top) and measured Doppler residuals (bottom) during the ground
424 demonstration of low-SNR tracking capabilities. SNR was measured using an FFT at 30 second
425 integration time and Doppler residuals were estimated using a PLL at 2-second integration time.
426



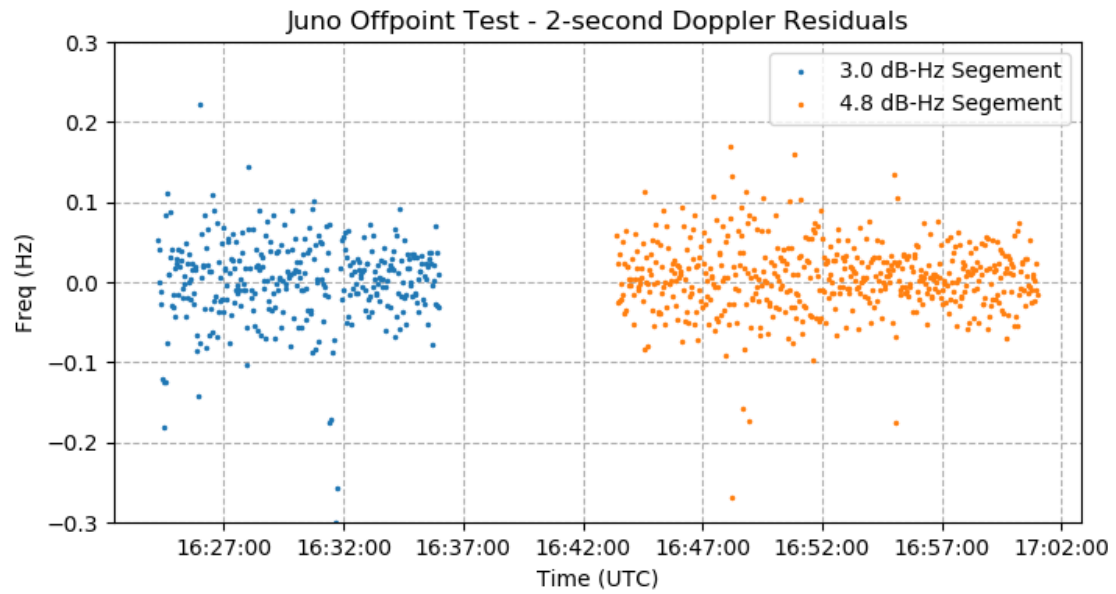
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428 **Figure 3.** Standard deviation of Doppler residuals as function of SNR during the ground system
 429 demonstration of low-SNR tracking capabilities, estimated using a PLL ($B_L = 0.25$ Hz, $T_c = 2$
 430 sec). The error bars are the 1-sigma standard deviation of the power measurement from Figure 2.
 431 The measurements are compared with a model of thermal noise from Eq. (1). The measurements
 432 are taken in the absence of propagation noise and spacecraft instrumental noise.
 433



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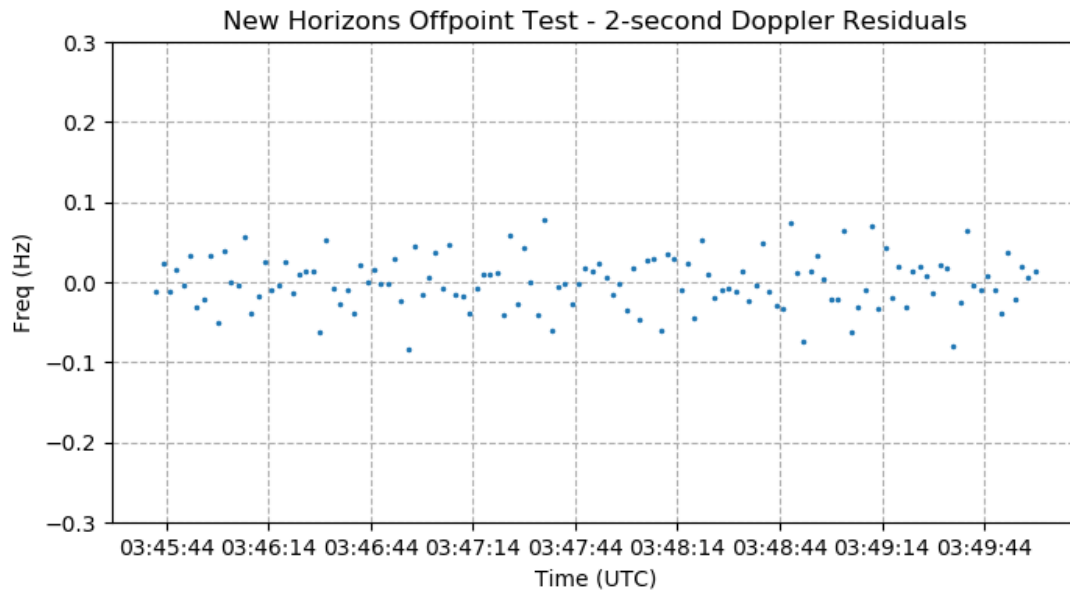
435 **Figure 4.** Allan deviation of the Doppler residuals at an SNR level of 4 dB-Hz from the ground
 436 system demonstration of low-SNR tracking capacities. The measured Allan deviation is
 437 proportional to τ^{-1} , indicating the signal is characterized by white phase noise.
 438



439

440 **Figure 5.** Coherent three-way Doppler residuals from Juno at 2-second integration time while
441 artificially reducing the received SNR by introducing a pointing offset into the DSN antenna.
442 The first segment was collected with an average SNR of 3 dB-Hz (blue) and second segment was
443 collected with an average SNR of 4.8 dB-Hz (orange).

444



445

446 **Figure 6.** Non-coherent one-way Doppler residuals from New Horizons at 2-second integration
447 time while artificially reducing the received SNR by introducing a pointing offset into the DSN
448 antenna. The average SNR during this segment of data was 4.5 dB-Hz.