

1 **Thermospheric Neutral Density Variation during the**
2 **“SpaceX” Storm: Implications from Physics-based**
3 **Whole Geospace Modeling**

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10 **Key Points:**

- 11 • Our model predicts up to $\sim 150\%$ enhancement in neutral density along Starlink
12 orbit during a moderate geomagnetic storm on 3-4 February 2022.
13 • Empirical models tend to underestimate storm effects in thermospheric density
14 enhancement.
15 • The whole geospace model resolves the gradual recovery and wave perturbations
16 missed by empirical models.

Abstract

The Starlink satellites launched on 3 February 2022 were lost before they fully arrived in their designated orbits. The loss was attributed to two moderate geomagnetic storms that occurred consecutively on February 3-4. We investigate the thermospheric neutral mass density variation during these storms with the Multiscale Atmosphere-Geospace Environment (MAGE) model, a first-principles, fully coupled geospace model. Simulated neutral density enhancements are validated by Swarm satellite measurements at the altitude of 400-500 km. Comparison with standalone TIEGCM and empirical NRLMSIS 2.0 and DTM-2012 models suggests better performance by MAGE in predicting the maximum density enhancement and resolving the gradual recovery process. Along the Starlink satellite orbit in the middle thermosphere (~ 200 km altitude), MAGE predicts up to 150% density enhancement near the second storm peak while standalone TIEGCM, NRLMSIS 2.0 and DTM-2012 suggest only $\sim 50\%$ increase. MAGE also suggests altitudinal, longitudinal, and latitudinal variability of storm-time percentage density enhancement due to height dependent Joule heating deposition per unit mass, thermospheric circulation changes, and travelling atmospheric disturbances. This study demonstrates that a moderate storm can cause substantial density enhancement in the middle thermosphere. Thermospheric mass density strongly depends on the strength, timing, and location of high-latitude energy input, which cannot be fully reproduced with empirical models. A physics-based, fully coupled geospace model that can accurately resolve the high-latitude energy input and its variability is critical to modeling the dynamic response of thermospheric neutral density during storm time.

Plain Language Summary

On February 3, 2022, 40 Starlink satellites were launched by the SpaceX Corporation when a moderate geomagnetic storm occurred, followed by another storm on February 4. The storm activities have been regarded as the culprit for the loss the Starlink satellites afterwards. Although strong geomagnetic storms are well-known to be able to increase the neutral atmospheric mass density so as to satellite drag in the thermosphere where many space vehicles are orbiting around the Earth, a not-so-strong storm was not expected to bring such huge impacts based on engineering design evaluation using empirical atmospheric density models. This study compares the performance of a state-of-the-art physics-based, fully coupled whole geospace model and empirical models in predicting the neutral mass density variation in the thermosphere. It turns out that the physics-based model is more accurate in capturing the magnitude of storm enhancement of neutral density. It also resolves the gradual recovery process even though it is not reflected in some geomagnetic indices that are used to drive the empirical models. Using such a first-principles whole geospace model is suggested as a necessary step in future space weather applications.

1 Introduction

The ionosphere-thermosphere (I/T) system is an important atmospheric region where many space vehicles are orbiting around the Earth. The thermospheric density is a critical parameter in determining the drag on those satellites and the resulting stability of the satellites and their orbits, as well as potential for collision with other objects. The variability of neutral density with geomagnetic conditions such as storms can thus significantly affect satellite operations. It is well-known that the I/T system is driven by internal dynamics and external forcing from the sun (e.g. Richards, 2001; L. Liu et al., 2011), magnetosphere (e.g., Fuller-Rowell et al., 1994; J. Forbes et al., 1996; Wang et al., 2008; Zhu et al., 2022), and lower atmosphere (e.g., H.-L. Liu et al., 2013; H.-L. Liu, 2016). Therefore, thermospheric density exhibits complicated temporal variability across a broad range of time scales, including solar cycle, seasonal, diurnal, and geomagnetic storm-time

67 variations, and spatial variations with latitude, longitude, and altitude (e.g., J. M. Forbes,
68 2007; Qian & Solomon, 2012).

69 Geomagnetic storms, especially major ones with Dst minimum less than -100 nT,
70 can greatly disturb the I/T system and thermospheric density (e.g., Fuller-Rowell et al.,
71 1994; J. Forbes et al., 1996). Using measurements from the Satellite Electrostatic Tri-
72 axial Accelerometer experiment in the middle thermosphere (~ 200 km altitude), J. Forbes
73 et al. (1996) found that the thermospheric density increased by 50-70% in the dayside
74 high latitude region when Kp increased from 1-2 to 4-7. Rhoden et al. (2000) also re-
75 ported an increase of density by 134% from Kp 1 to 6. CHAMP and GRACE data have
76 been widely used in the recent two decades to study the upper thermospheric density
77 variation at 400-500 km altitude in the noon and midnight sectors. H. Liu and Lühr (2005)
78 reported a maximum density enhancement of 400-500% during the 29-31 October 2003
79 super storm. Bruinsma et al. (2006) found a global density increase of 300-800% dur-
80 ing the 20-21 November 2003 storm. Lei et al. (2010) reported a storm time density in-
81 crease of 200-400% during the main phase of the 7-9 November 2004 major storm.

82 The storm-time enhancement of thermospheric density is caused by increased mag-
83 netospheric energy input at high latitudes. During storm times, high latitude ionospheric
84 plasma convection is enhanced and causes a substantial increase in Joule heating. Joule
85 heating rate is proportional to ionospheric Pedersen conductivity, which usually max-
86 imizes in the E region and is increased during storms too, due to enhanced auroral pre-
87 cipitation (Campbell, 2003). Auroral precipitation also brings extra energy inputs to the
88 I/T system in the auroral region through ionization, dissociation, and excitation of neu-
89 tral constituents (Qian & Solomon, 2012). Therefore, storm effects on thermospheric den-
90 sity are expected to depend on many factors, including geomagnetic activity level and
91 altitude. Furthermore, disturbances at high latitudes are transmitted to lower latitudes
92 by storm-induced changes in large scale circulation and travelling atmospheric distur-
93 bances. There is typically a few hours time delay in the thermosphere responses between
94 low and high latitudes (e.g., Bruinsma et al., 2006). Consequently, the storm time vari-
95 ation of thermospheric density becomes a global phenomenon. The dependence of ther-
96 mospheric and ionospheric responses on three-dimensional spatial distribution of Joule
97 heating rate and ion drag, which is controlled by the characteristics of a particular storm
98 event, such as the onset time, solar wind condition, type, strength and duration of the
99 event, makes it more challenging to accurately characterize storm time variation of ther-
100 mosphere neutral density.

101 However, satellite and ground-based measurements can only cover certain local lon-
102 gitude/latitude regions, altitudes and local times during a storm. To provide a global
103 coverage and for prediction purposes, various I/T models have been developed and used
104 to investigate thermospheric density distribution and variability. Empirical models use
105 statistical relations with geomagnetic indices and solar fluxes to derive neutral density
106 as a function of latitude, longitude, altitude and time (e.g. Emmert et al., 2021; Hedin,
107 1987; Bowman et al., 2008; Bruinsma, 2013). While empirical models such as NRLM-
108 SIS 2.0 are good at characterizing the average density distribution and longer term vari-
109 ability, they usually have difficulty in resolving storm-time responses as these models may
110 easily miss the dynamic variability of the storm-time I/T system on small spatial and
111 temporal scales (e.g. Rhoden et al., 2000; Bruinsma et al., 2006). On the other hand,
112 recent developments in whole geospace modeling shows great advances in characteriz-
113 ing storm-time large and mesoscale structures such as traveling atmospheric disturbances
114 (Pham et al., 2022). A physics-based global geospace model solves for I/T state from
115 first-principles equations and treats the I/T as a component of the coupled geospace sys-
116 tem.

117 During the past few years of solar minimum, major geomagnetic storms have be-
118 come rare. Storm-time thermospheric density variation and its impacts on satellite drag
119 appear to be less concerning. However, a moderate geomagnetic storm that occurred on

120 3-4 February 2022 was reported to have caused the loss of 40 Starlink satellites of the
 121 SpaceX Corporation orbiting in the middle thermosphere. This storm is hereinafter re-
 122 ferred to as the “SpaceX storm” (Hapgood et al., 2022). Empirical models such as the
 123 NRLMSIS 2.0 and DTM models and standalone ionosphere-thermosphere general cir-
 124 culation models such as TIEGCM predict relatively moderate density enhancement of
 125 about 50% at the Starlink orbit altitude of ~ 200 km, which is less likely to be strong
 126 enough to account for the satellite drag that failed the satellites (e.g., Dang et al., 2022).
 127 In this study, we use a newly developed whole geospace model, the Multiscale Atmosphere-
 128 Geospace Environment (MAGE), to investigate the thermospheric density variations dur-
 129 ing the “SpaceX storm” event. We also compare the performance of MAGE with sev-
 130 eral other models, including a standalone I/T general circulation model (TIEGCM) (Richmond
 131 et al., 1992), the empirical models of NRLMSIS 2.0 (Emmert et al., 2021) and the Drag
 132 Temperature Model (DTM) (Bruinsma, 2013).

133 2 Methodology

134 The MAGE model is a whole geospace model that is being developed with the aim
 135 to explore the magnetosphere-ionosphere-thermosphere as a tightly coupled system. The
 136 current version used here, MAGE 1.0, consists of the Grid Agnostic MHD for Extended
 137 Research Applications (GAMERA) global magnetohydrodynamic (MHD) model of the
 138 magnetosphere (Zhang et al., 2019; Sorathia et al., 2020), the Rice Convection Model
 139 (RCM) of the ring current (Toffoletto et al., 2003), Thermosphere Ionosphere Electro-
 140 dynamics General Circulation Model (TIEGCM) of the upper atmosphere (Richmond
 141 et al., 1992), and the RE-developed Magnetosphere-Ionosphere Coupler/Solver (REMIX)
 142 (Merkin & Lyon, 2010). MAGE is driven by observed solar wind and interplanetary mag-
 143 netic field (IMF) at its upstream boundary. Magnetosphere-driven high latitude iono-
 144 spheric convection and auroral precipitation are used to drive the TIEGCM. Details about
 145 the model configuration can be found in Pham et al. (2022) and Lin et al. (2021).

146 Figure 1 shows the solar wind and IMF conditions and geomagnetic activity indices
 147 of Kp and SymH from February 1 to February 6. A coronal mass ejection arrived at the
 148 Earth by the end of February 1, manifested by abruptly increased solar wind speed and
 149 density and oscillating IMF. IMF B_Z suddenly dropped to -12 nT at 23 UT on Febru-
 150 ary 1 although the southward B_Z lasted for less than an hour. Solar wind speed increased
 151 from ~ 350 km/s to nearly 500 km/s. Solar wind density increased from ~ 4 /cc to more
 152 than 10/cc, resulting in a dynamic pressure enhancement by more than five times. A sud-
 153 den commencement can be identified from the SymH index. IMF B_Z stayed mostly north-
 154 ward on February 2 except a two-hour interval of -6 nT. The IMF smoothly turned south-
 155 ward on February 3 and reached a minimum of -18 nT at 10 UT. The B_Z turning was
 156 accompanied by B_Y rotation from +10 nT to -18 nT. The SymH index showed a typ-
 157 ical main phase feature of rapid decrease (Wanliss & Showalter, 2006) and reached a min-
 158 imum value of -75 nT at ~ 11 UT on February 3. At around 12 UT on February 3, an-
 159 other dynamic pressure impulse arrived at the Earth with northward IMF. The first storm
 160 gradually recovered till 4 UT on February 4 when IMF B_Z turned to -10 nT and stayed
 161 southward for the next two days. A second storm can be seen in the SymH index with
 162 another minimum of -60 nT at 22 UT on February 4. The storm activities were recover-
 163 ing since February 5. In this study, the observed solar wind and IMF conditions (Fig-
 164 ures 1a-1c) were used to drive the MAGE model to investigate the thermospheric mass
 165 density variation during the storm and its potential impacts on the Starlink satellites
 166 which were launched at around 18 UT on February 3 (vertical red dashed line).

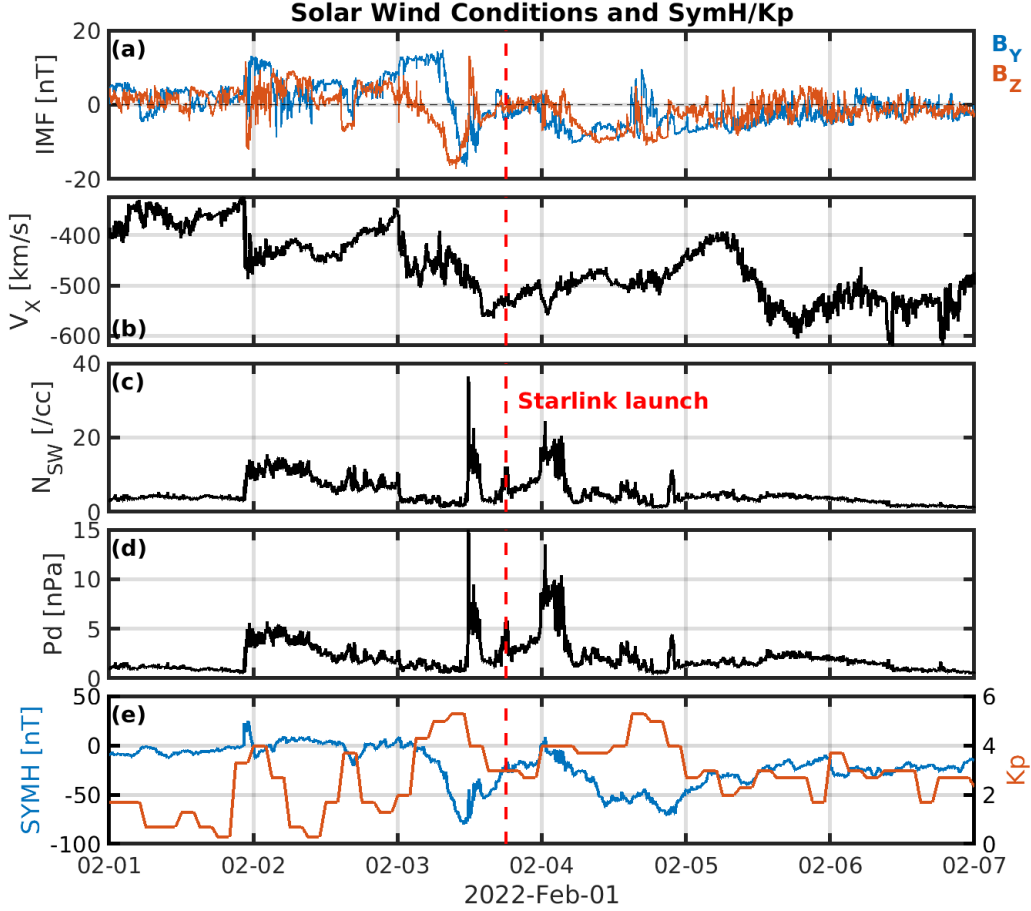


Figure 1. Solar wind/IMF data and SymH/Kp indices during 1-6 February 2022. (a) IMF B_Y (blue) and B_Z (red) in the Solar Magnetic (SM) coordinates. (b) Solar wind velocity V_X . (c) Solar wind density N_{SW} . (d) Solar wind dynamic pressure P_d . (e) SymH (blue) and Kp indices (red). The vertical red dashed line indicates the launch time of Starlink satellites.

3 Results and Discussion

3.1 Neutral density variations at Swarm orbit altitude

We start by analyzing the neutral density variations measured by the Swarm satellites during the SpaceX storm event. This is because the neutral densities derived from the Starlink orbit are not yet publicly available, and when available, will require calibration and comparison against science grade neutral density data. The Swarm satellites A and B provide such science grade neutral densities (van den IJssel et al., 2020). Although the Swarm satellite orbits are higher than the newly-launched Starlink satellites during this time, the densities along the Swarm orbit demonstrate the response of the thermosphere to the storm and provide a baseline to assess the accuracy of the model responses.

Figures 2a-2b show the relative variation of the neutral density $\Delta\rho/\rho_0$ observed by the Swarm A and B satellites. Here ρ_0 is the daily mean density on February 1, the quiet-day reference before the storm. $\Delta\rho = \rho - \rho_0$ is the absolute variation of neutral density from ρ_0 . The relative variation is averaged over each orbit period, i.e., 93 min for

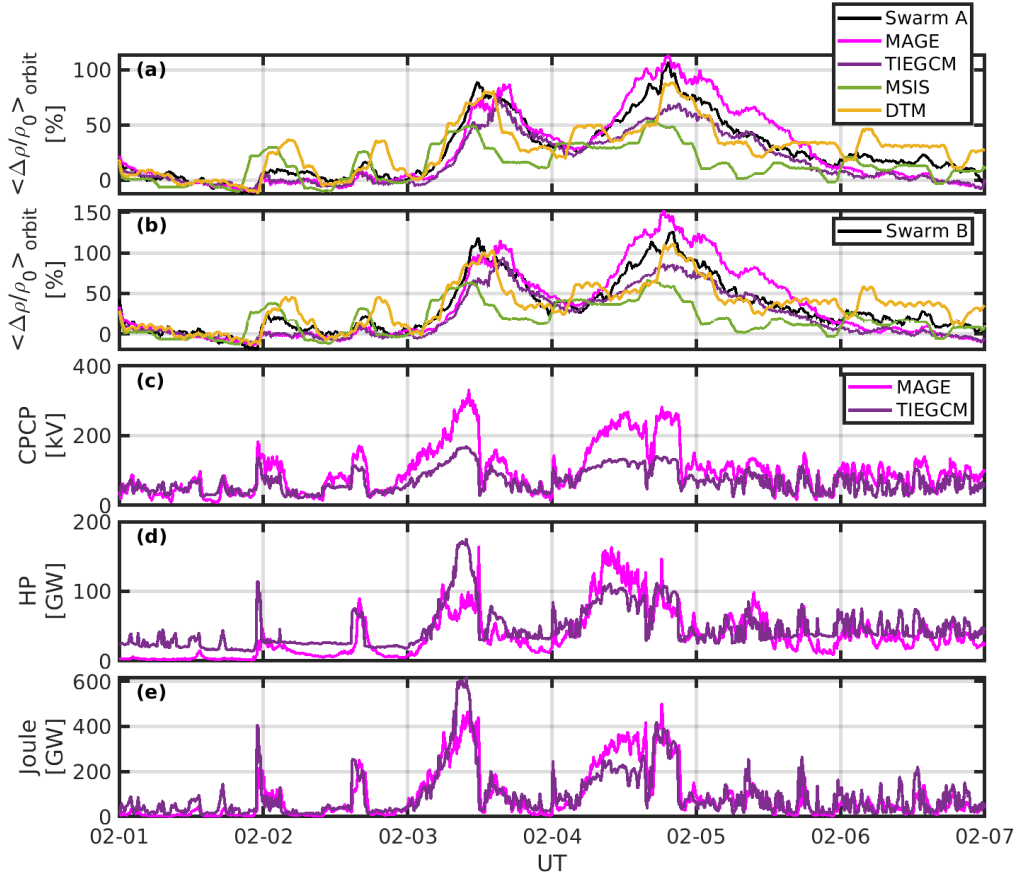


Figure 2. (a-b) Neutral density variations measured by the Swarm satellite (black) and simulated by the MAGE (magenta), standalone TIEGCM driven by the Weimer empirical convection model (purple), NRLMSIS 2.0 (green), and DTM-2012 (yellow) models during the SpaceX storm event. The density variation is normalized by the daily mean value on February 1 and averaged over each orbit. (c-e) Cross-polar cap potential (CPCP), hemisphere power (HP), and integrated high latitude Joule heating rate in the northern hemisphere simulated by the MAGE and standalone TIEGCM.

182 Swarm A and 95 min for Swarm B. Swarm C results are not shown because they are identical to those from Swarm A, which shared the same orbit. Swarm A and B has an average orbit altitude of 440 km and 515 km, respectively.

185 As shown in Figure 2a, Swarm A observed an increase of neutral density starting on February 3, normalized by the daily mean density of February 1. It reached 87% around
 186 12 UT when the first storm reached the SymH minimum. During the recovery phase of
 187 the first storm, the orbit mean density decreased, but it still maintained a perturbation
 188 level that was $\sim 25\%$ higher than the quiet day reference value by 4 UT on February
 189 4. After the second storm started, the mean density at Swarm A orbit increased again
 190 to a peak value of 105% at around 20 UT on February 4. The Swarm A-measured density
 191 then gradually decreased to the pre-storm level during the second recovery phase
 192 by February 7. The density variation observed by Swarm B showed a similar two peak
 193 pattern to Swarm A observations. The peak values of the relative increase are 118% and
 194

195 124% during the two storms, slightly higher than the values observed by Swarm A as
 196 Swarm B orbit is higher.

197 MAGE-simulated neutral density variations are shown by the magenta curves in
 198 Figures 2a and 2b. The neutral density is sampled along the Swarm trajectories in MAGE
 199 outputs, which are saved every 5 minutes. The density variation is also normalized by
 200 the daily mean density of February 1 and averaged over each orbit. It can be seen that
 201 the MAGE model reproduces the relative density variation in quantitative agreement
 202 with Swarm satellite observations. The magnitude of the MAGE-simulated relative den-
 203 sity changes and the temporal variations in the first storm period are close to Swarm ob-
 204 servation except at ~ 16 UT on February 3 when MAGE predicts a perturbation peak
 205 which is not seen in the data. For the second storm, the MAGE results appear to over-
 206 estimate the perturbations around the storm peak at ~ 18 UT as well as in the storm
 207 recovery phase on February 5. Nevertheless, the observed major density enhancements
 208 of $\sim 100\%$ are captured by the MAGE model during both storm peaks on February 3
 209 and February 4. MAGE also reproduces the process of gradual increase and recovery dur-
 210 ing different phases of the two storms.

211 As a comparison, the relative neutral density variations simulated by the standalone
 212 TIEGCM, and empirical NRLMSIS 2.0 and DTM-2012 models are shown with the pur-
 213 ple, green, and yellow curves in Figures 2a and 2b. The standalone TIEGCM simula-
 214 tion is driven by the Weimer (2005) convection model with the same solar wind/IMF con-
 215 ditions for MAGE run. The NRLMSIS 2.0 model is driven by the 3-hourly Ap index,
 216 and daily and 81-day average F10.7 indices. The DTM-2012 model is driven by hourly
 217 Hpo (Yamazaki et al., 2022) and solar F30 indices. It can be seen that the Weimer-driven
 218 TIEGCM simulation underestimated the relative density enhancement during both storms
 219 (purple curves). In particular, the second storm peak on February 4 was only $\sim 65\%$
 220 higher than the quiet-day reference along the Swarm A orbit while the Swarm A data
 221 and MAGE simulation both suggest a $\sim 110\%$ enhancement. DTM-2012 (yellow curves)
 222 predicted comparable density enhancements during the two storms but overestimated
 223 the density enhancement on February 2 and February 6 when the geomagnetic activity
 224 level was relatively low. The model also underestimated density perturbations in the re-
 225 covery phase of the first storm and the initial phase of the second storm. NRLMSIS 2.0
 226 (green curves) underestimated the density enhancement for both storms. Its peak value
 227 is about half of that measured by Swarm. The density enhancement simulated by NRLM-
 228 SIS 2.0 also occurred earlier than Swarm measurements, probably due to the timing un-
 229 certainty of the 3-hour Ap index that was used to drive the model.

230 Figures 2c-2e show the strength of high latitude convection pattern and energy in-
 231 puts simulated by the MAGE model and standalone TIEGCM. The convection strength
 232 is indicated by the cross-polar cap potential (CPCP). The hemisphere power (HP) is the
 233 total precipitation energy flux integrated over each hemisphere above 45° magnetic lat-
 234 itude. This latitude range is chosen to cover the large-scale convection driven by the mag-
 235 netospheric dynamics during the storms. During the first storm, the density enhance-
 236 ment was accompanied by greatly enhanced CPCP, HP, and total Joule heating rate. CPCP
 237 reached ~ 300 kV, HP exceeded 100 GW and Joule heating rate was more than 400 GW
 238 at 12 UT on February 3 in both hemispheres. Although CPCP, HP, and Joule heating
 239 all dropped suddenly after the first storm peak at ~ 12 UT due to northward turning
 240 of IMF (Figure 1a), the neutral density remained enhanced and recovered with a much
 241 slower rate. The gradual recovery of neutral density is more evident after the second storm
 242 on February 4. Both the Swarm observed and MAGE simulated neutral density did not
 243 recover to the pre-storm level until February 6 even though the Kp index dropped to be-
 244 low 3 on February 5 (Figure 1e) and drove a fast recovery of the MSIS neutral density.

245 Figure 2c also shows that the TIEGCM CPCP, which is essentially the Weimer (2005)
 246 model, is only 50% of the MAGE CPCP. Compared to the CPCP, the HP and total Joule
 247 heating rate in the Weimer driven TIEGCM are closer to MAGE results and sometimes

248 higher (Figures 2d-2e). Note the CPCP, HP, and Joule are all hemisphere integrated quantities.
 249 Their spatial distribution and temporal evolution of these quantities also play a
 250 critical role in the storm time variation of the ionosphere-thermosphere system. Pham
 251 et al. (2022) found that MAGE simulated neutral density has a better agreement with
 252 CHAMP and GRACE measurements compared to Weimer driven standalone TIEGCM,
 253 even though the MAGE run has a lower Joule.

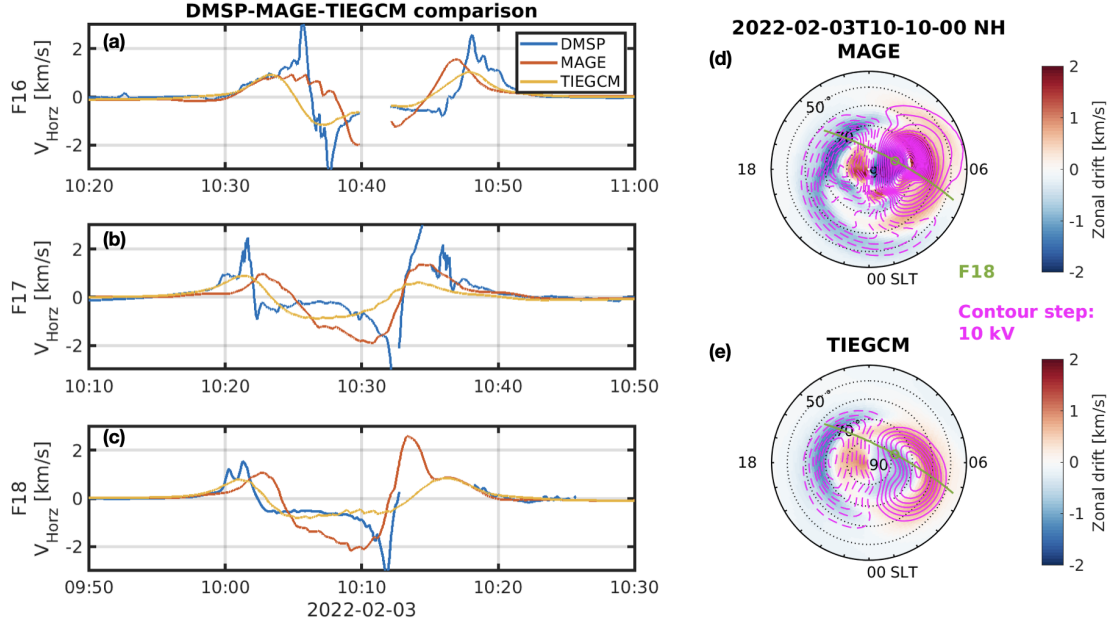


Figure 3. (a-c) Cross-track ion drift velocity measured by DMSP and sampled from MAGE and standalone TIEGCM simulation results. (d-e) Zonal ion drift and convection potential contour simulated by MAGE and standalone TIEGCM.

254 The MAGE simulated high latitude convection is further validated with DMSP measurements
 255 of cross track ion drift. Figure 3 shows the comparison of V_{Horz} measured by
 256 DMSP (blue), and simulated by MAGE (red) and standalone TIEGCM (yellow) at around
 257 10:10 UT on February 3, when the CPCP reached the maximum value. The ion drifts
 258 are smoothed with a moving average over 15s to remove the observed oscillations on a
 259 spatial scale smaller than the grid size, which is 1.25° in both MAGE and standalone
 260 TIEGCM. The V_{Horz} comparison shows that the TIEGCM convection is generally weaker
 261 than DMSP measurements, while MAGE simulation results are more consistent with DMSP
 262 data. Note CPCP cannot be directly measured without a global coverage of the entire
 263 convection pattern. DMSP V_{Horz} can only partially represent CPCP when the trajectory
 264 is close to the convection cell centers, which is roughly the case for DMSP F18 during
 265 the interval around 10:10 UT (see the green curve and magenta contours in Figures
 266 3d-3e). Figures 3d-3e show the zonal ion drift and convection potential contour at 2022-
 267 02-03/10:10 UT. The convection contour is equally spaced by every 10 kV. While the
 268 TIEGCM sampling of V_{Horz} agrees well with DMSP F18 data in the polar cap, its dawn-
 269 side side convection velocity is much weaker than DMSP measurements after 10:10 UT,
 270 suggesting the CPCP from Weimer is indeed underestimated. The MAGE sampling of
 271 V_{Horz} resolves the dawnside convection with a much better agreement, including resolv-
 272 ing the ~ 2 km/s convection flow, although the MAGE results appear to be overesti-
 273 mated in the polar cap. A similar comparison is found for DMSP F16 and F17.

274 In this study, we focus on the large scale storm time variation of thermospheric neu-
 275 tral density, which is more relevant to the accumulated effects on satellite drag. Hence,
 276 we present the orbit averaged relative variations. The absolute values of neutral density
 277 in the model are sometimes higher than Swarm and GRACE Follow-on (FO) measure-
 278 ments, which is not shown here. The difference in absolute values is related to the initial
 279 state of TIEGCM used by MAGE, which shows higher density than Swarm and GRACE-
 280 FO data. However, with the initial offset removed, the relative variation in MAGE sim-
 281 ulation results does show a quantitative agreement with observations in terms of the storm
 282 time neutral density evolution.

283 3.2 Neutral density variations along the Starlink orbit

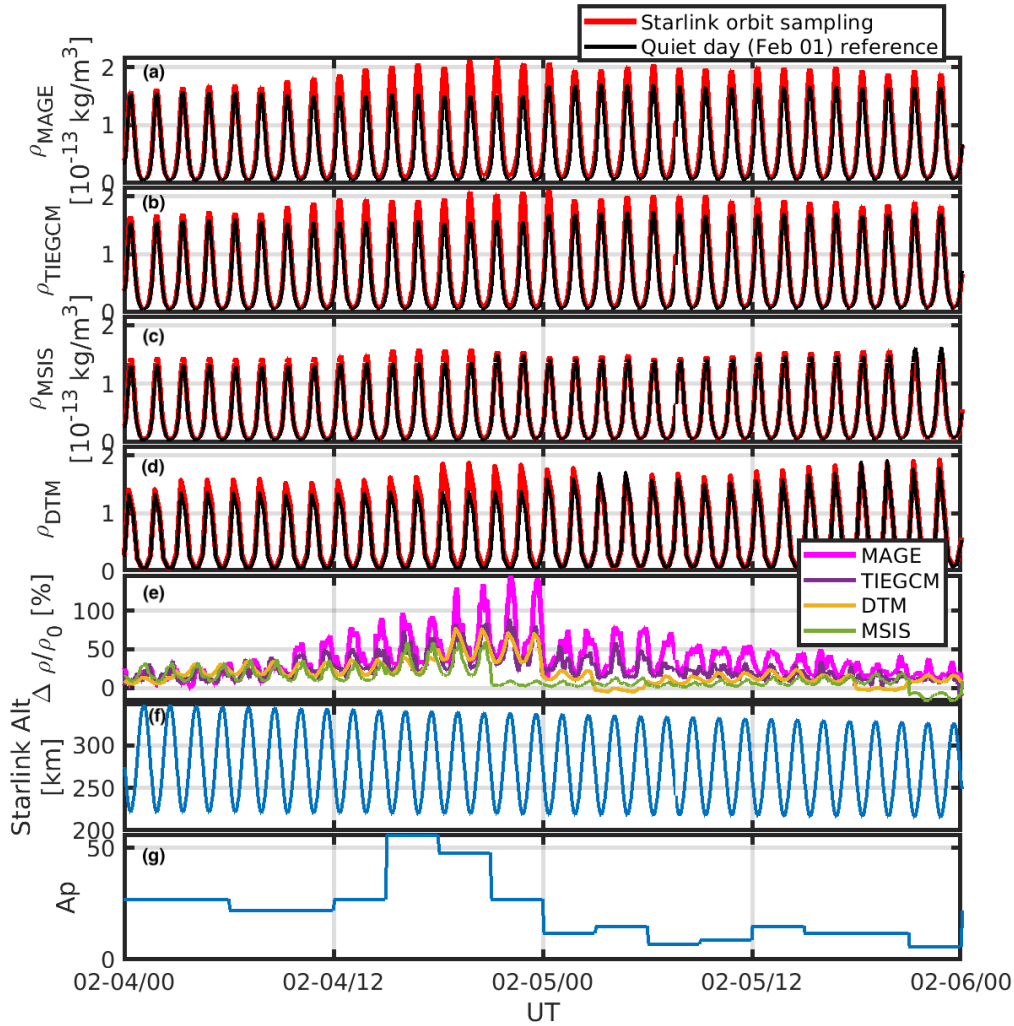


Figure 4. Neutral density variations along the Starlink orbit calculated by (a) MAGE, (b) TIEGCM, (c) DTM-2012, and (d) NRLMSIS 2.0. (e) Relative variations of neutral density based on the values on February 1 at the same UT and location. (f) Starlink altitude. (g) The Ap index used to drive NRLMSIS 2.0.

284 We now turn to an analysis of the modeled neutral density variation along the Star-
 285 link orbit. Since Starlink was not in orbit before February 4, we compare the February

286 4 and 5 density with that on February 1 at the same orbital location to investigate the
 287 storm effects. The Starlink orbit data was derived from publicly available NORAD Two-
 288 Line Elements (TLEs) obtained from CelesTrak (<https://celestrak.org/>). The altitude
 289 of the Starlink orbit ranged from 200 km to 330 km (Figure 4f), involving a density dif-
 290 ference of more than one order of magnitude. Therefore, we do not use orbit average but
 291 use a point-to-point comparison of neutral density. Figures 4a-4d show the neutral den-
 292 sity simulated by MAGE, standalone TIEGCM, DTM-2012, and NRLMSIS 2.0, respec-
 293 tively. The red curves show the density sampled along the Starlink orbit and the black
 294 curves show the density sampled at the same location and UT on February 1.

295 By comparing the red and black curves, all four models show storm time enhance-
 296 ment of neutral density along the Starlink orbit after it was launched. However, differ-
 297 ent models predict vastly different levels of relative density enhancement. To better il-
 298 lustrate the storm effects, we compare the ratio of neutral density to February 1 for MAGE,
 299 TIEGCM, DTM-2012 and NRLMSIS 2.0 in Figure 4e. Despite a periodic variation that
 300 is closely related to the orbit altitude, the MAGE results suggest a density enhancement
 301 by up to 150% during the second storm on February 4 (magenta curve). The orbital max-
 302 imum enhancement dropped to $\sim 75\%$ after 0 UT on February 5 and slowly recovered
 303 to $\sim 20\%$ by the end of February 6. The variation of neutral density in MAGE simu-
 304 lation results is determined by the high latitude energy input, which is shown in Figures
 305 2c-2f. The hemisphere power and integrated Joule heating were greatly enhanced dur-
 306 ing the two storms, which increase the neutral temperature and elevate the atmosphere,
 307 as well as the neutral density along the satellite orbit. In the storm recovery phase when
 308 IMF B_z turned northward, high-latitude energy inputs significantly decreased and the
 309 I/T system began to recover to its undisturbed state before the storms. This recovery
 310 process is slow (Burns et al., 1991) and the density enhancement simulated by MAGE
 311 also takes time to dissipate as shown in Figure 4, as well as for the Swarm densities at
 312 higher altitudes (Figures 2a and 2b).

313 On the other hand, standalone TIEGCM (purple), DTM-2012 (yellow) and NRLM-
 314 SIS 2.0 (green) results only suggest an enhancement of at most 50–75% during the sec-
 315 ond storm on February 4. Particularly, the NRLMSIS 2.0 simulated neutral density en-
 316 hancement dropped to near zero immediately at 21 UT on February 4. The variation
 317 of NRLMSIS 2.0 results can be largely attributed to the 3-hour Ap index that is used
 318 as the input for the model. Note the F10.7 index inputs for all models are daily values
 319 and only varied by less than 5% during these events (between 120 and 126). Figure 4g
 320 shows that the Ap index increased from 27 to 56 at 15 UT and decreased from 48 to 27
 321 at 21 UT, which yields a moderate density enhancement in NRLMSIS 2.0 results (Fig-
 322 ure 4c). DTM-2012 and standalone TIEGCM results show a rapid recovery on Febru-
 323 ary 5, in contrast to the gradual recovery shown in MAGE results. While the Ap index
 324 does suggest a rapid recovery of geomagnetic activity at auroral latitudes on February
 325 5, the SymH index (Figure 1e) and Swarm measurements of mass density (Figures 4a-
 326 4b) reveal that the storm effects were still present on February 5 and it took a longer
 327 time for the thermospheric neutral density to recover. This contrast also highlights the
 328 need to properly represent the recovery phase in order to accurately predict the mass
 329 density variation during the entire process of geomagnetic storms.

330 3.3 Neutral density variation with altitude

331 The density ratios along the Starlink orbit exhibit a periodic pattern, which is mainly
 332 attributed to the altitudinal variation of the orbit. In fact, comparison of the results at
 333 Swarm A, Swarm B, and Starlink orbits also suggests an altitudinal dependence of storm
 334 time neutral density variations. To better understand the altitudinal dependence, we com-
 335 pared the vertical distributions of density enhancement in Figure 5 with MAGE simu-
 336 lation results. The density variation relative to the same location and UT on February
 337 1 is evaluated for all longitudes and latitudes. The average variation over all longitudes

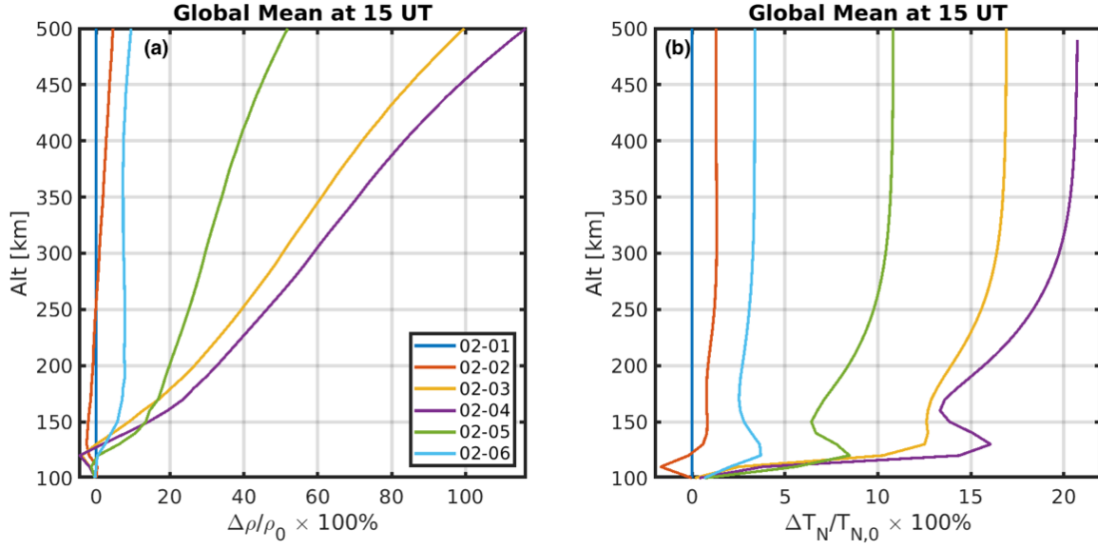


Figure 5. Global mean relative neutral density (a) and neutral temperature (b) variations with altitude at 15 UT on different days.

338 and latitudes is shown in Figure 5a for 15 UT of six consecutive days from February 1
 339 to 6. From February 3 to February 5 (yellow, purple, and green), the percentage den-
 340 sity enhancement increases nearly linearly with altitude. In the altitude range between
 341 200 km and 330 km where Starlink orbit is located, global thermospheric neutral den-
 342 sity is on average enhanced by 40–60% at 15 UT on February 4 and by 20–30% on Febru-
 343 ary 5.

344 Figure 5b shows a similar analysis for neutral temperature variation. Compared
 345 to the quasi-linear dependence of density enhancement on altitude during the storm time,
 346 the neutral temperature enhancement is nearly constant above 200–300 km altitude. Dur-
 347 ing the first storm on February 3, the neutral temperature was globally enhanced by 10%–
 348 17% below 300 km altitude and uniformly enhanced by 17% above 300 km. The tem-
 349 perature enhancement was more than 20% at high altitudes near the peak of the second
 350 storm at 15 UT on February 4.

351 There is a local peak of neutral temperature enhancement at ~ 125 km altitude
 352 during the storm time. The temperature enhancement peak is attributed to the altitu-
 353 dinal distribution of Pedersen conductivity, which has a maximum value around this height
 354 and results in the maximum Joule heating deposition (Campbell, 2003). The storm time
 355 temperature enhancement increases with altitude because the background density de-
 356 creases exponentially while Joule heating rate per unit mass increases with altitude (Huang
 357 et al., 2012; Deng et al., 2011).

358 In Figure 5 we only show the comparison below 500 km altitude to ensure the alti-
 359 tudinal range is covered by MAGE at all longitudes and latitudes for all six days. Note
 360 in MAGE, the TIEGCM grid is defined by pressure levels which vary with altitude/temperature
 361 depending on many factors including solar zenith angle, solar radiation and magneto-
 362 spheric conditions. We show 15 UT of each day because it is near the peak of both two
 363 storms. Other UTs also show similar dependencies on altitude. The density enhance-
 364 ment is contributed mostly by the increased thermospheric temperature and thus a larger
 365 atmospheric scale height, which raises the altitudes of the pressure surfaces. Thermo-
 366 spheric composition also contributes to neutral density changes (Lei et al., 2010). Con-

367 sequently, in the altitude coordinates in which satellite observations are made the neu-
 368 tral density at the same altitudes during the storms is increased.

369 3.4 Neutral density variation with latitude

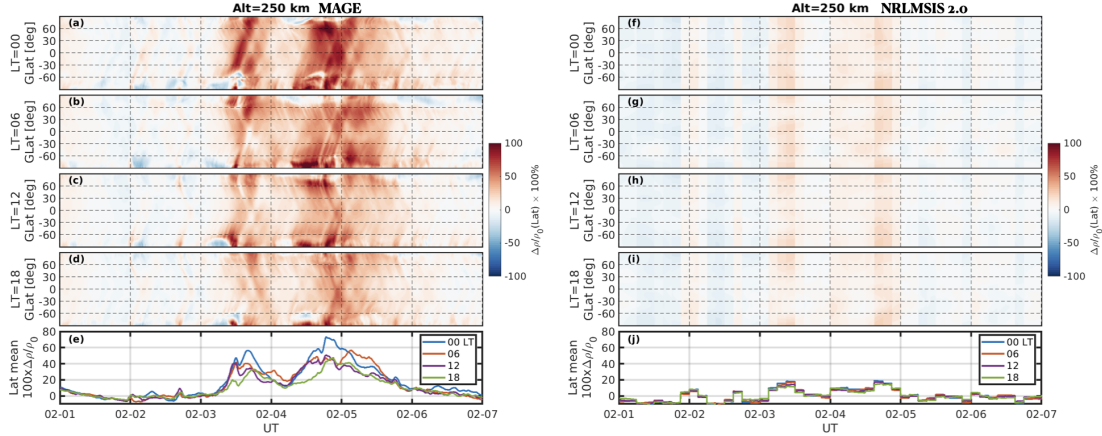


Figure 6. MAGE (left) and NRLMSIS 2.0 (right) simulated neutral density variations at different latitudes and four local times at the altitude of 250 km.

370 While Figure 5a suggests that the global mean density enhancement is up to 40%
 371 and temperature enhancement is up to 18% during the storms at the average Starlink
 372 orbit altitude of 250 km, thermospheric response exhibits substantial spatial and tem-
 373 poral variability. Figure 6 shows the latitudinal and UT dependence of neutral density
 374 variation simulated by MAGE and NRLMSIS 2.0 at four different local times (LTs): mid-
 375 night (00LT), dawn (06LT), dusk (18LT), and noon (12LT). The relative density vari-
 376 ation is normalized by the daily mean density at the same latitude and local time on Feb-
 377 ruary 1.

378 Latitudinally, the relative density variation is generally stronger at high latitudes
 379 than at middle and low latitudes. For instance, at 15 UT on February 4 along 00 LT,
 380 the relative density enhancement is $\sim 80\%$ at 60° geographical latitude and $\sim 50\%$ at
 381 -85° geographical latitude, whereas it is less than 40% within 20° from the equator. The
 382 stronger enhancement at higher latitudes is attributed to the storm time high latitude
 383 energy inputs from the magnetosphere, which are mostly deposited in the auroral zone
 384 and polar cap regions. As shown in Figure 2c, the CPCP is greatly enhanced during the
 385 two storms, indicating a much stronger high latitude convection. On the other hand, the
 386 magnetospheric precipitation is also intensified during storm time as represented by the
 387 hemispheric power (Figure 2d). High-latitude energy inputs are localized and vary with
 388 time (Pham et al., 2022). Therefore, storm-time neutral density changes depend on lo-
 389 cal time. For instance, the latitudinal mean density changes peak at different UTs for
 390 different LTs (Figure 6e). On February 3, the density changes attain the first peak al-
 391 most at the same UTs for all LTs, whereas for the second peak on the same day occurred
 392 first at midnight, followed by peaks at dusk, noon and dawn. Similar LT dependence is
 393 also seen during the second storm on February 4. In addition, the midnight sector has
 394 the largest relative density perturbations while noon experiences the least changes for
 395 both storms.

396 The right column shows that same latitudinal variations of neutral density from
 397 the NRLMSIS 2.0 model. Neutral density perturbations are noticeably smaller than those
 398 simulated by the MAGE model, as summarized by the latitudinal mean variations shown

399 in Figure 6e for MAGE and Figure 6j for NRLMSIS 2.0. The NRLMSIS 2.0 results do
 400 not show obvious LT dependence either. It also does not have the propagation char-
 401 acter of density perturbations as seen in MAGE results, that is, there is an obvious time
 402 delay between the high latitude and low latitude density variations. This is understand-
 403 able as NRLMSIS 2.0 is an empirical model and it does not include dynamics as the first
 404 principles model MAGE does.

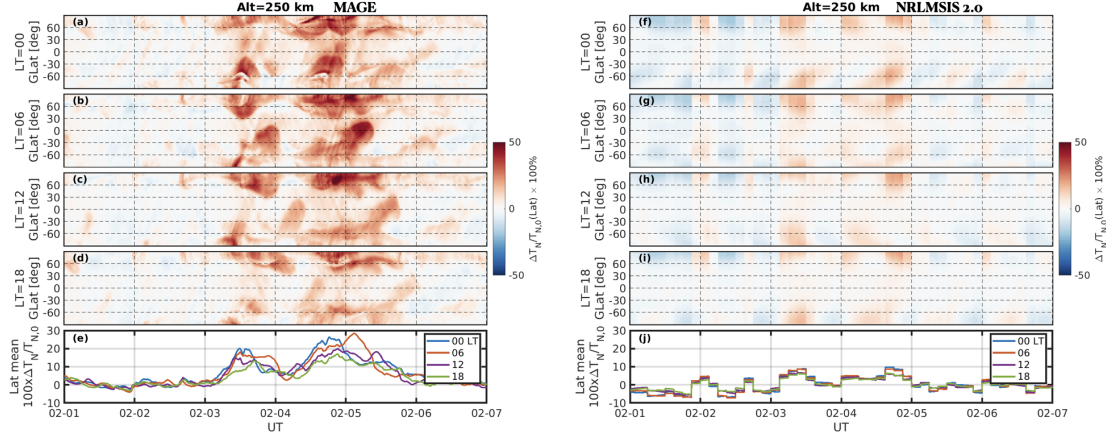


Figure 7. MAGE (left) and NRLMSIS 2.0 (right) simulated neutral temperature variation at different latitudes and four local times at the altitude of 250 km.

405 In addition to density variations, we also analyzed the responses of neutral temper-
 406 ature during this event. Figures 7a-7d show the keograms of neutral temperature sam-
 407 pled at the same altitude and four local times as in Figure 6. The neutral temperature
 408 response shows similar latitudinal and local time dependence, i.e., stronger enhance-
 409 ment at high latitudes than low latitudes during the storms and the nightside undergoes a larger
 410 relative neutral temperature variation. However, the temperature variation does not fully
 411 resemble that of the neutral density, implying other processes such as composition changes
 412 may have contributed to the density variations (e.g., Bruinsma et al., 2006).

413 The MAGE keograms of density (Figure 6) and temperature (Figure 7) variation
 414 also show a time delay of the density and temperature enhancement at different latitudes,
 415 which is known to result from traveling atmospheric disturbance (TAD) launched by the
 416 increased energy inputs at high latitudes (e.g., Fuller-Rowell et al., 1994; Pham et al.,
 417 2022). The comparison between Figures 6(a-e) and Figures 6(f-i) illustrates the contrast
 418 between a physics-based whole geospace model and an empirical I/T model. The MAGE
 419 model shows clear propagation effects from high to low latitudes which is caused by TADs,
 420 whereas the NRLMSIS 2.0 model results hardly show such wave or propagation effects.
 421 Again, as mentioned earlier, the first-principles MAGE model solves the dynamics of the
 422 I/T system and can resolve large and mesoscale structures like TADs. In addition, as
 423 a fully coupled geospace model, MAGE provides a more accurate physical specification
 424 of the magnetospheric forcing on the I/T system, including the location and temporal
 425 variation of high-latitude energy inputs which is critical to accurately simulate the I/T
 426 responses such as TADs to storms (Pham et al., 2022).

427 The comparison among four local times, midnight, dawn, noon, and dusk, in Fig-
 428 ures 6a-6d reveals that neutral density and temperature variations are nonuniform at dif-
 429 ferent local times. Figure 6e shows the latitudinal average of the relative density vari-
 430 ations for the four local times. MAGE simulation results indicate that the nightside ther-
 431 mosphere undergoes the largest relative variation in neutral density while the dayside

432 shows the least relative variation across all latitudes on average. The day/night asym-
433 metry is consistent with the observational analysis by J. Forbes et al. (1996) and has to
434 do with the stronger wind surge on the nightside. Increased Joule heating during the storm
435 time forces a global wind surge which undergoes less dissipation by ion drag at low and
436 middle latitudes on the nightside than on the dayside (e.g., J. Forbes et al., 1996). Fur-
437 thermore, the poleward pressure gradient produced by solar heating works against the
438 equatorward propagation of wind and temperature perturbations during the daytime (Wang
439 et al., 2008).

440 4 Summary

441 In this study, we used the MAGE model, a fully coupled whole geospace model,
442 to investigate the neutral density variation during the “SpaceX storm” that occurred on
443 3-4 February 2022. MAGE-simulated storm-time density enhancements are in quanti-
444 tative agreement with Swarm satellite measurements in the upper thermosphere. Com-
445 parison with the standalone TIEGCM and empirical models such as DTM-2012 and NRLM-
446 SIS 2.0 shows an overall better agreement of the physics-based model in simulating neu-
447 tral density perturbations during the storm event. The MAGE model also resolves the
448 recovery phase as a gradual process rather than an instant change suggested by empir-
449 ical models driven with the 3-hour Ap index. The better performance of MAGE simu-
450 lation results is attributed to its capability to dynamically characterize the high-latitude
451 energy inputs as a result of magnetospheric forcing.

452 We further evaluated the density enhancement along the Starlink satellite orbit in
453 the middle thermosphere near 250 km, where the dynamics are rich but least understood,
454 and observational data are sparse. Despite a moderate storm with a SymH minimum of
455 only -75 nT, the storm time density enhancement simulated by MAGE is up to 150%
456 along the Starlink orbit during the second storm on February 4, much stronger than stan-
457 dalone TIEGCM and empirical model results. The greatly enhanced neutral density is
458 more likely to have caused strong satellite drag that finally failed the Starlink satellites.
459 As we are entering the ascending phase of the current solar cycle, the geomagnetic ac-
460 tivity level is expected to increase within the next decade. Reliable observational and
461 numerical capabilities, including those enabled by the fully coupled geospace models like
462 MAGE, to characterize and resolve storm effects on neutral density at all altitudes from
463 middle to upper thermosphere, as well as on other critical I/T parameters, are becom-
464 ing increasingly important in space weather applications.

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474 and Information Systems Laboratory (CISL).

475 Data Availability Statement

476 The OMNI data are available at <https://cdaweb.gsfc.nasa.gov/index.html/>.
477 The Swarm neutral density data is obtained from <http://thermosphere.tudelft.nl>.
478 The simulation data supporting the analysis of this study is archived at this repository:
479 <https://doi.org/10.5065/axkv-me49>.

480

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Figure1.

Solar Wind Conditions and SymH/Kp

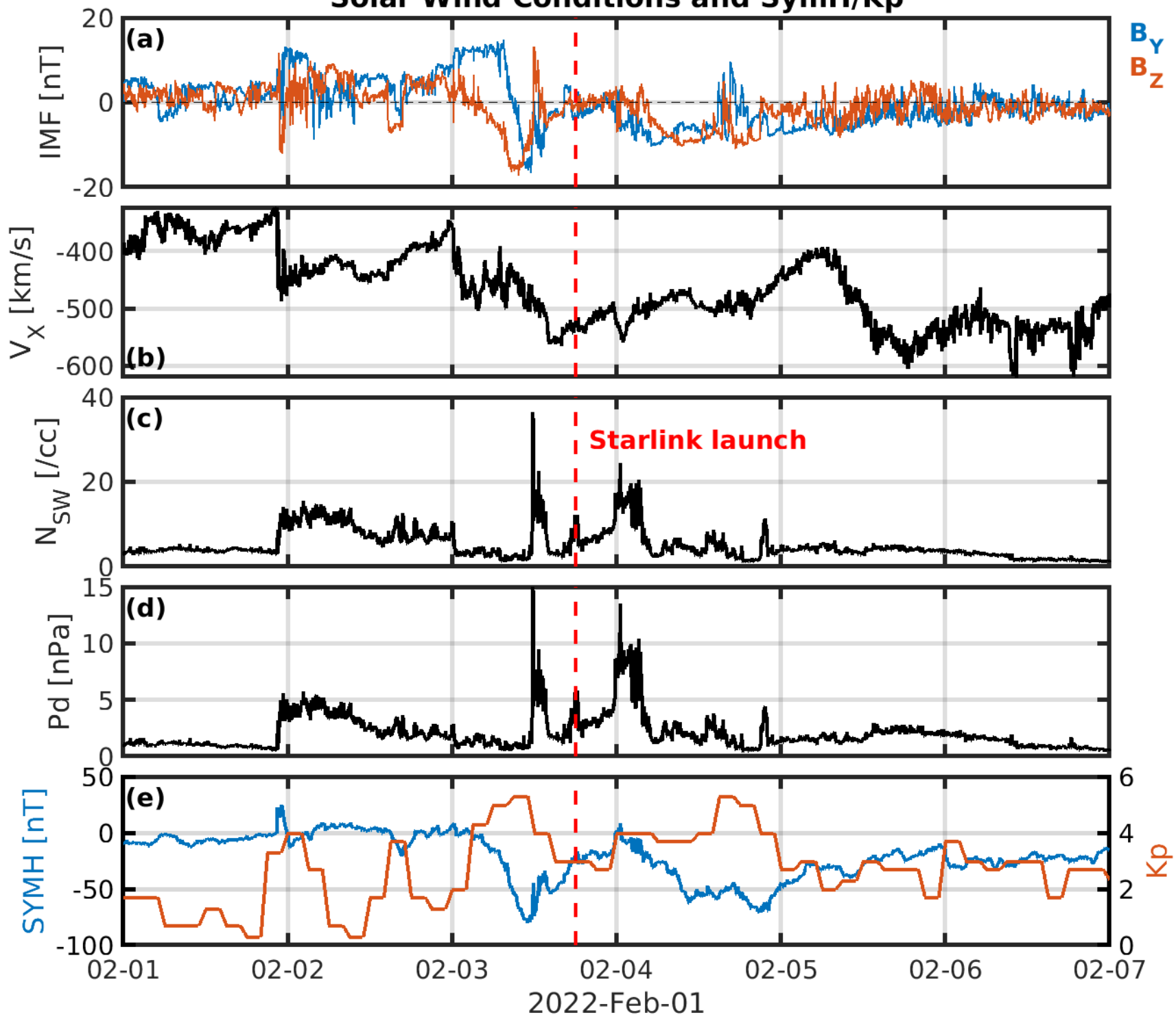


Figure2.

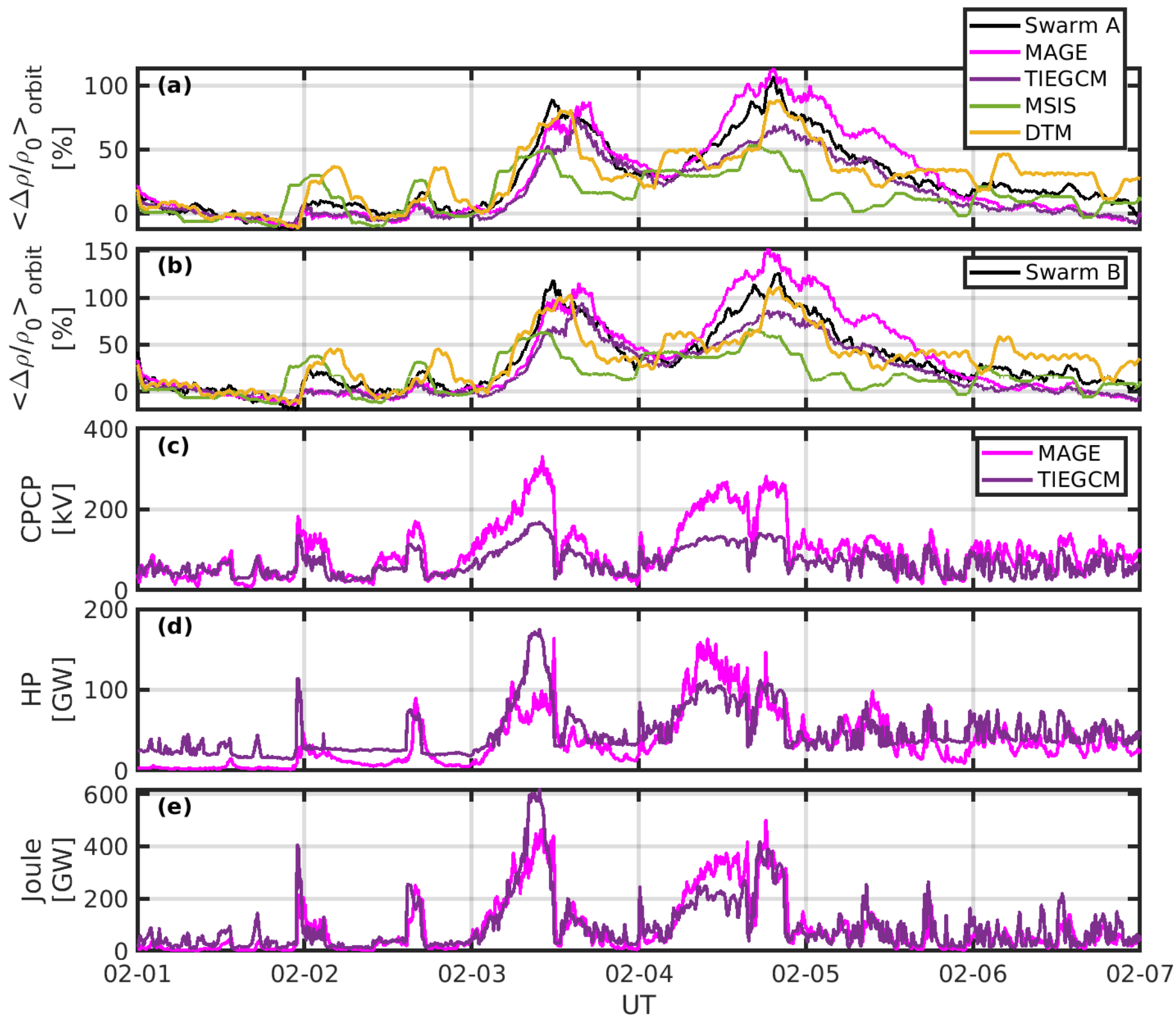
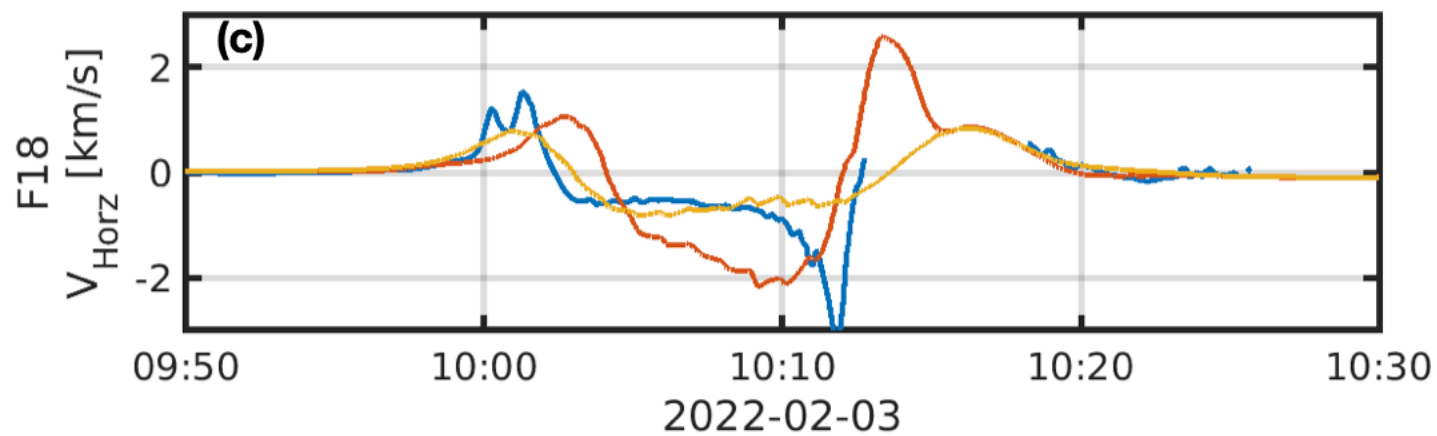
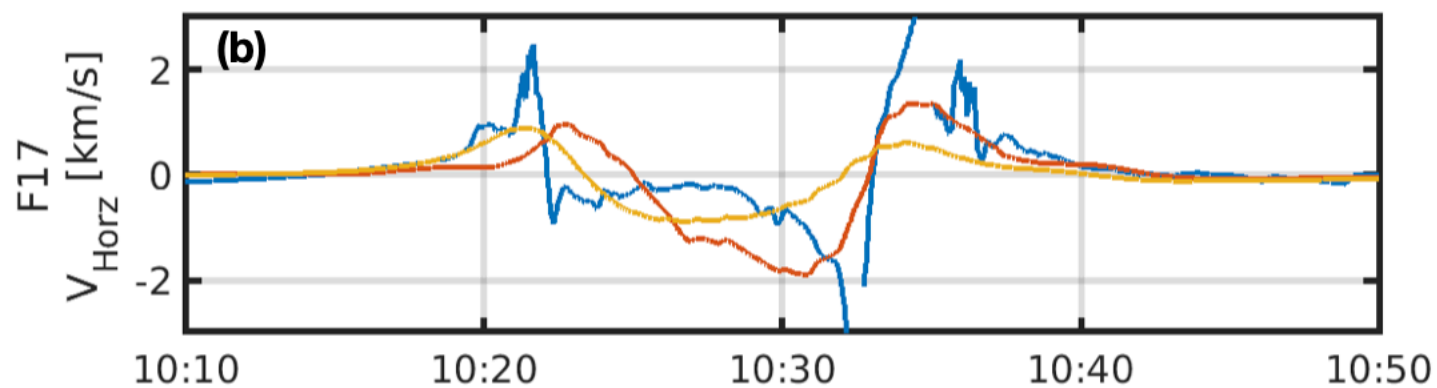
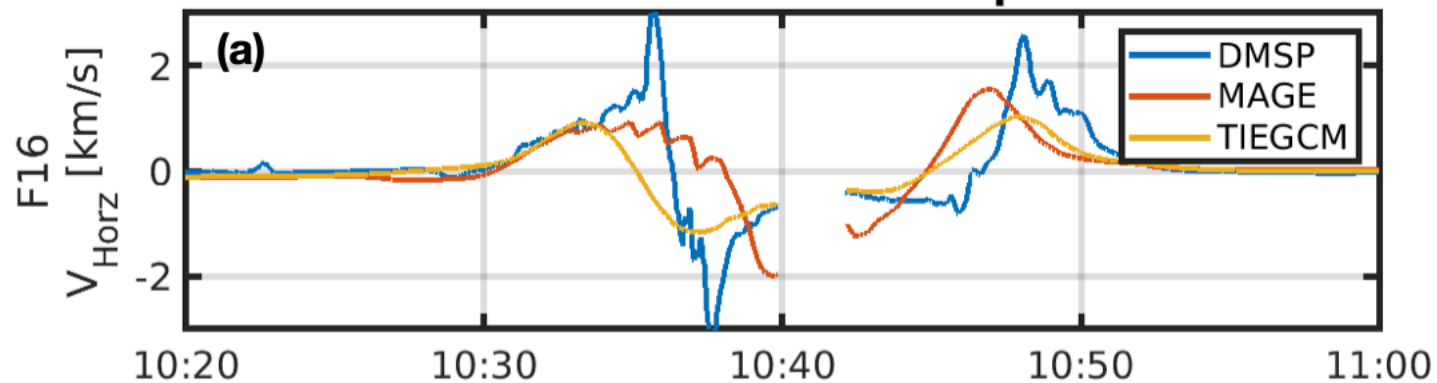


Figure3.

DMSP-MAGE-TIEGCM comparison



2022-02-03T10-10-00 NH

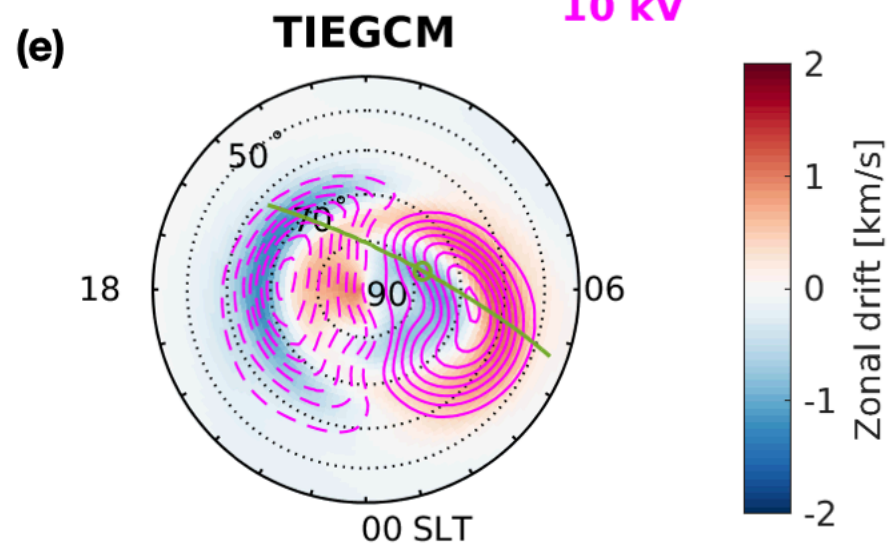
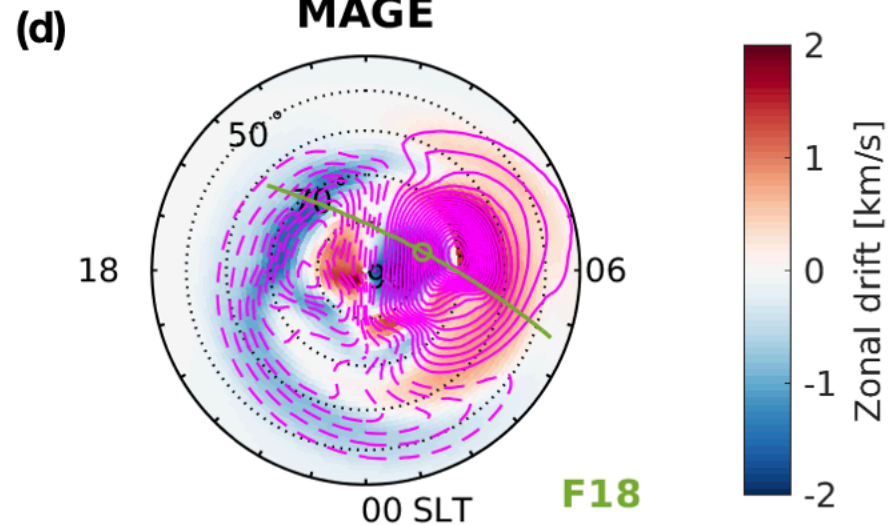


Figure4.

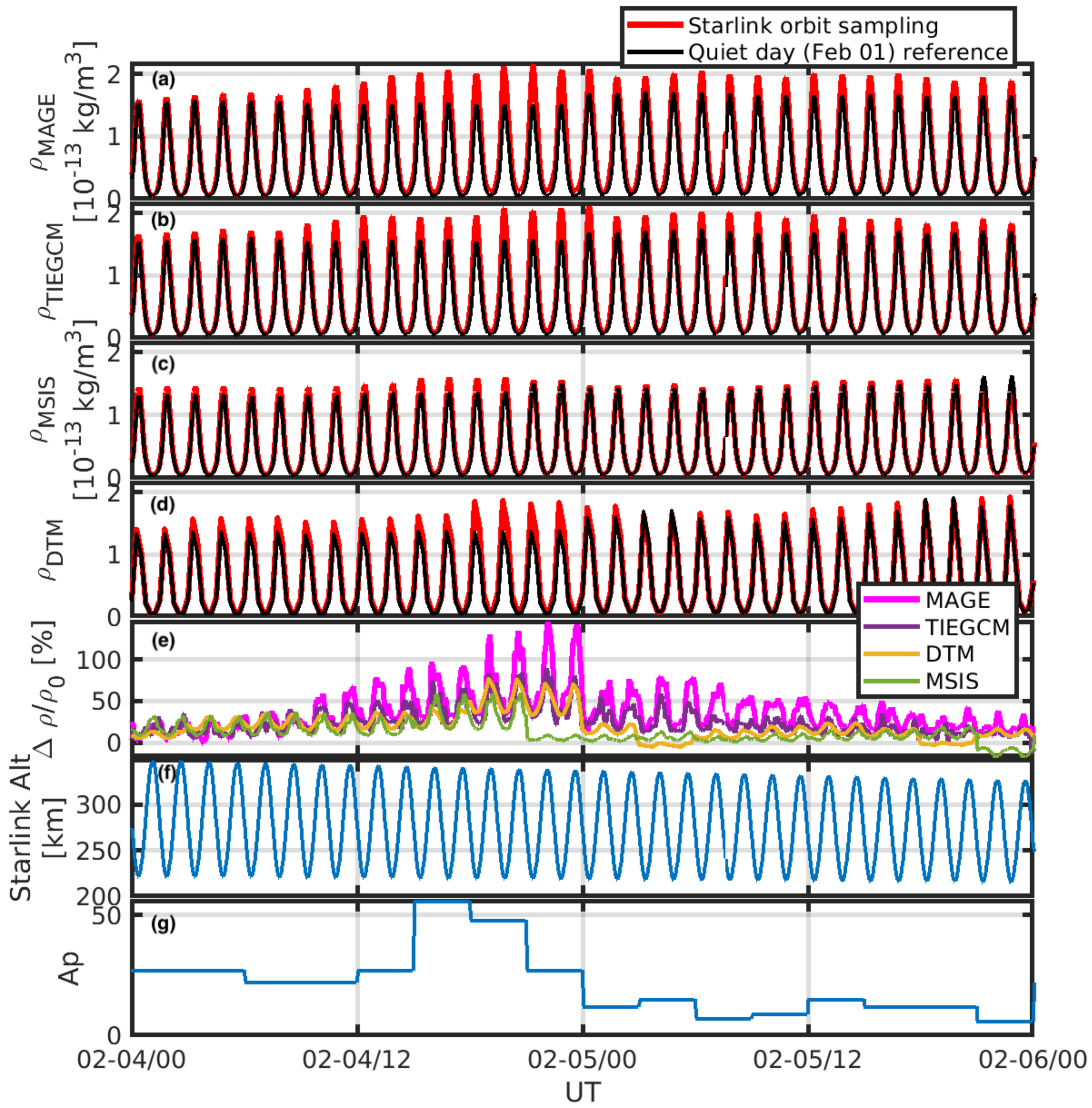
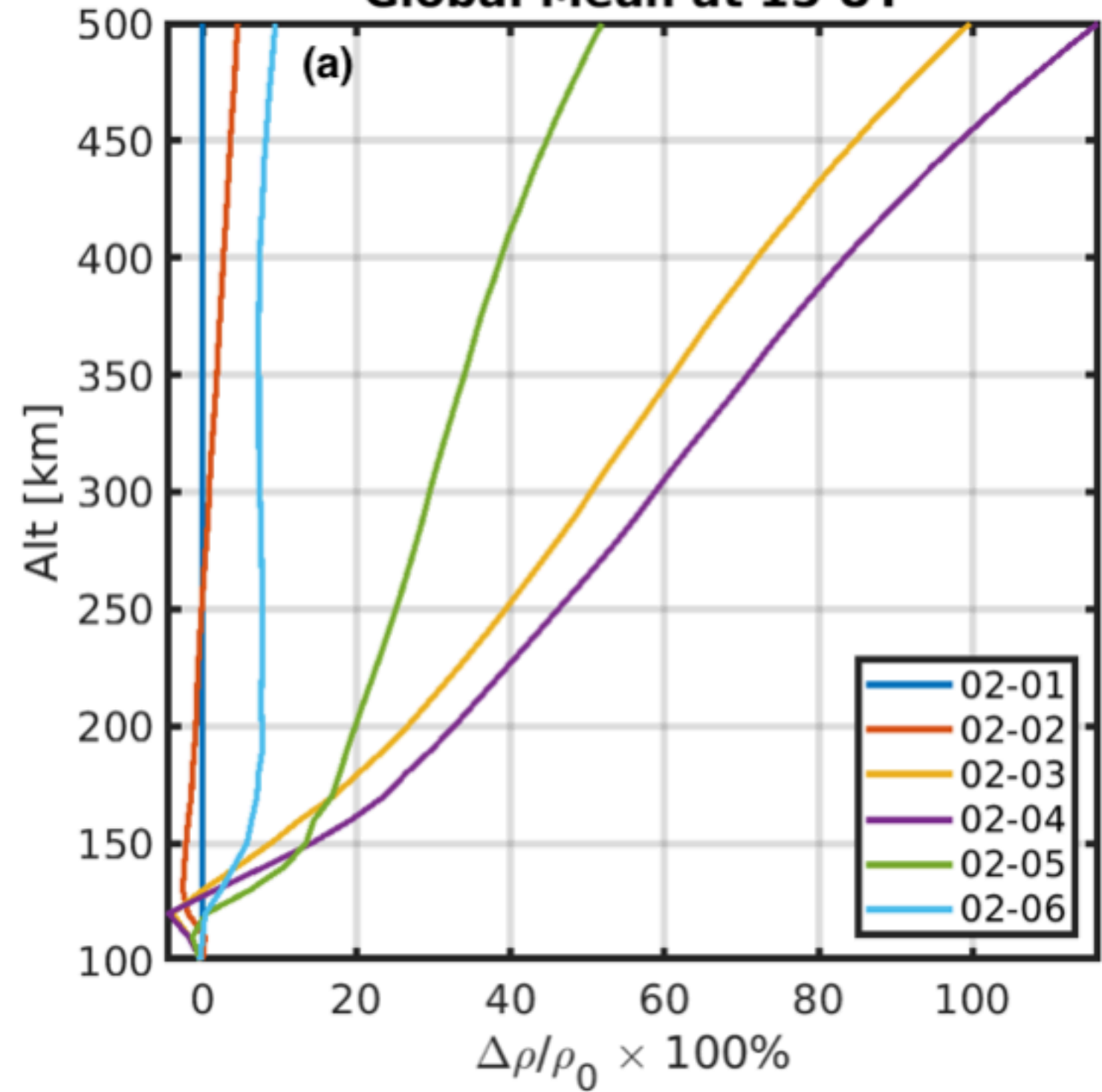


Figure5.

Global Mean at 15 UT



Global Mean at 15 UT

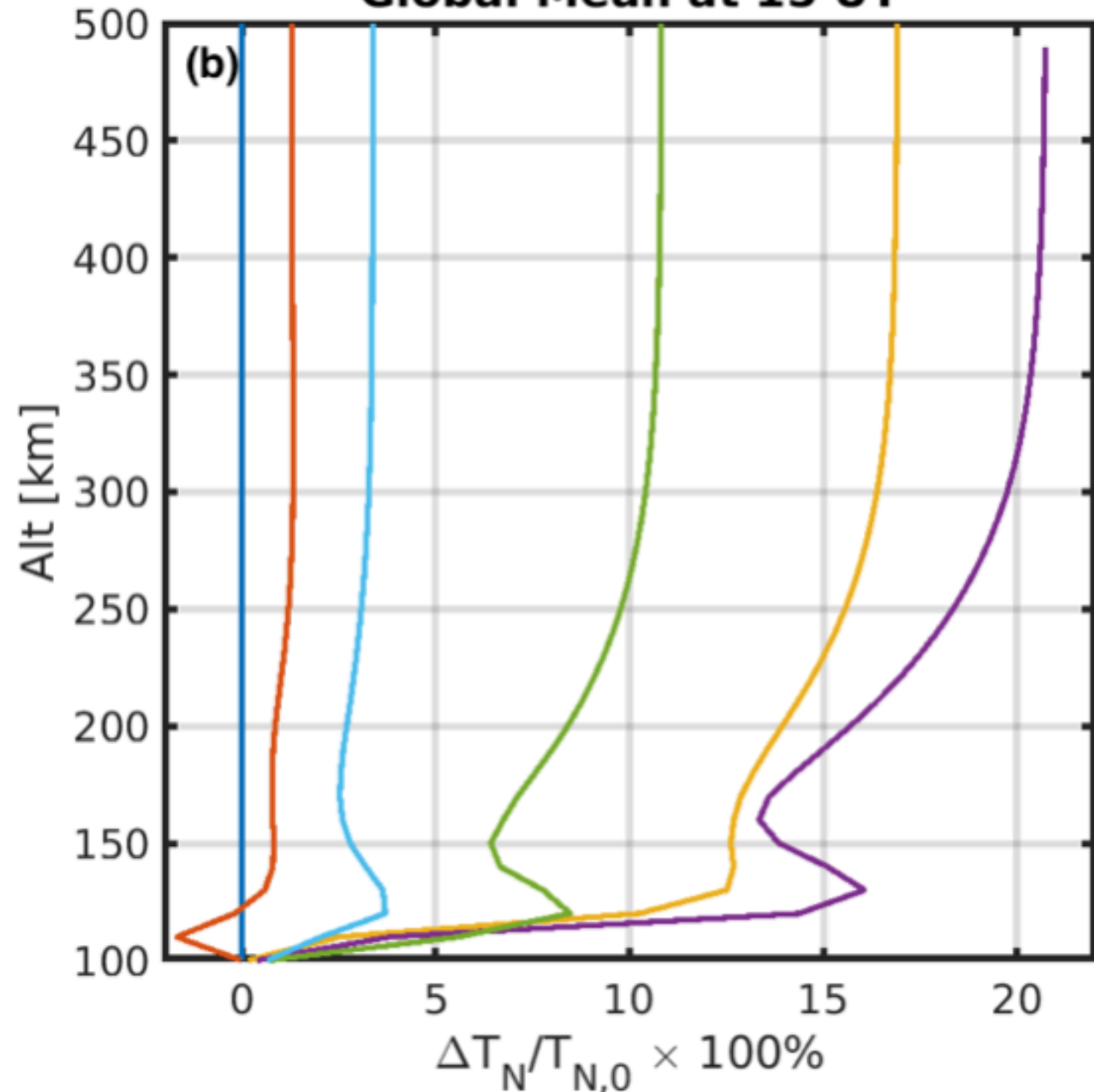
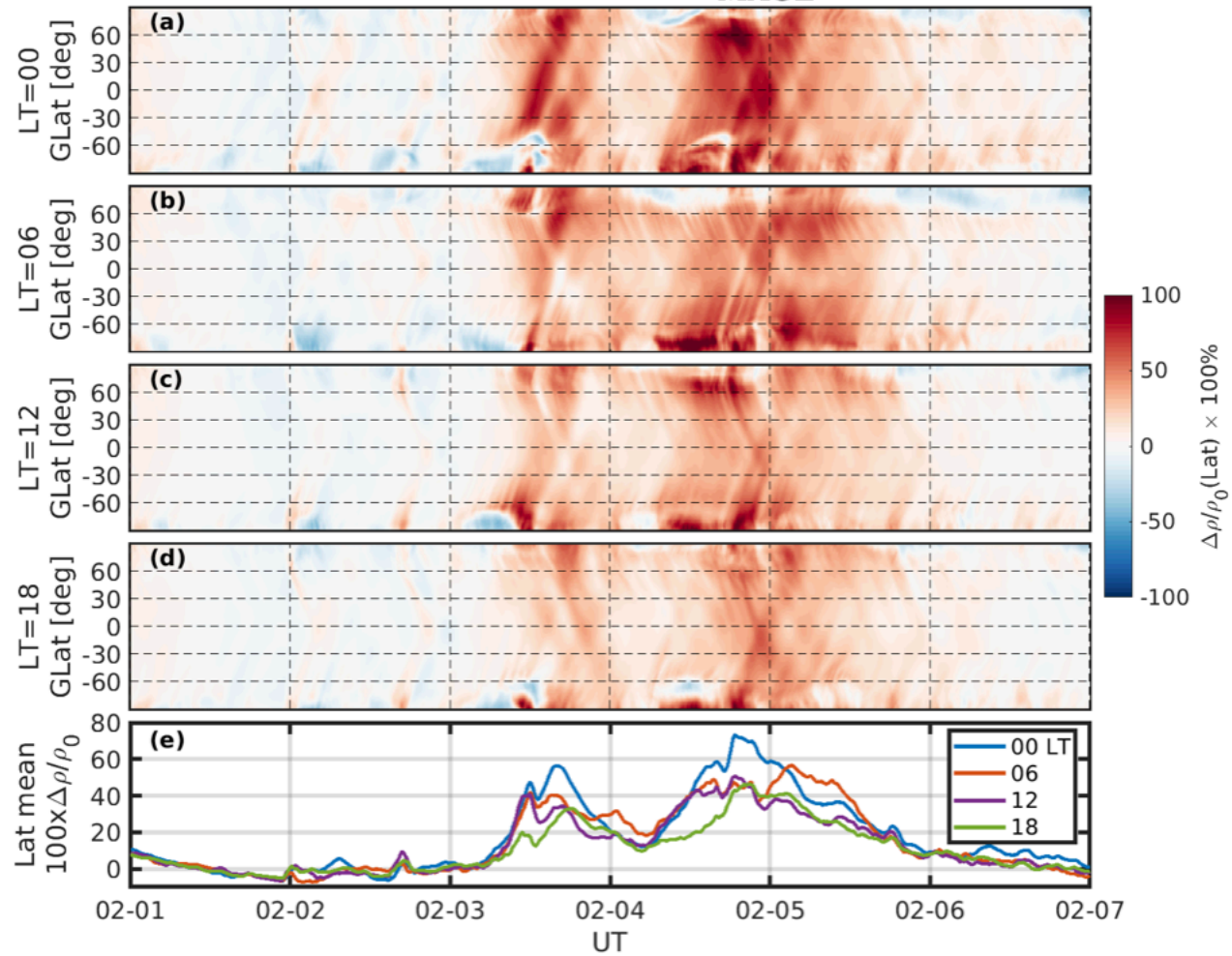


Figure6.

Alt=250 km MAGE



Alt=250 km NRLMSIS 2.0

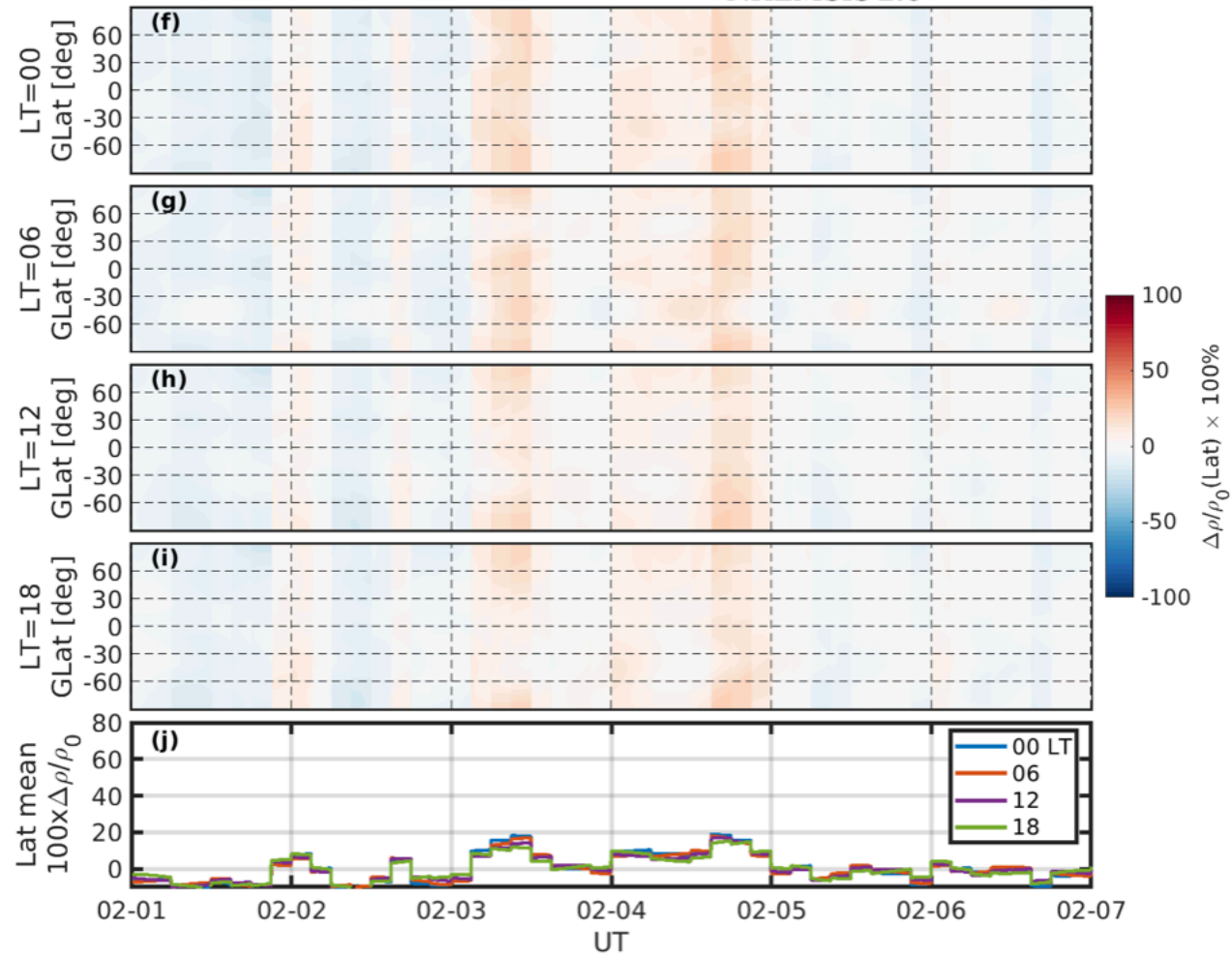
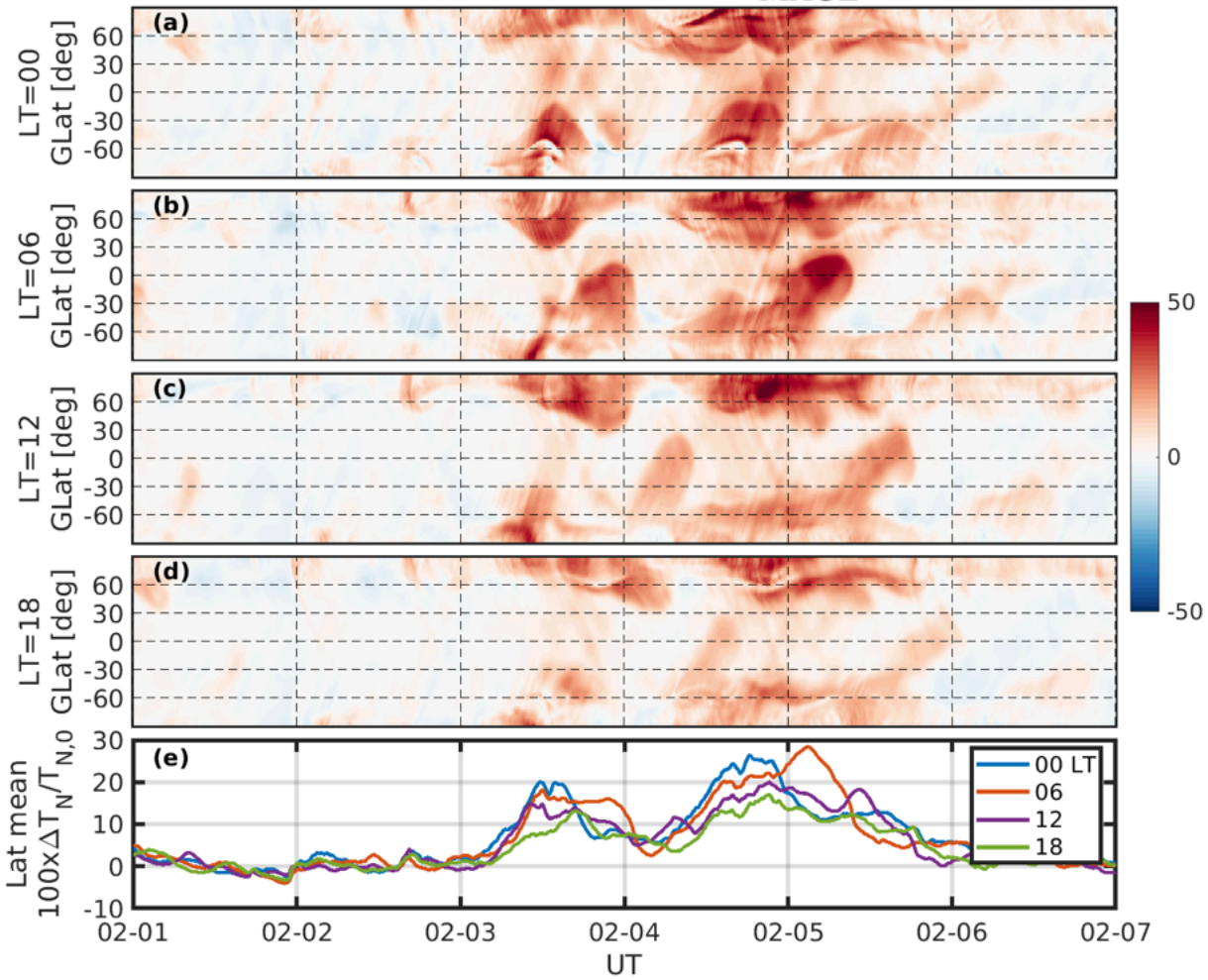


Figure7.

Alt=250 km MAGE



Alt=250 km NRLMSIS 2.0

