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Quantum Algorithms for lattice Φ^4 theory

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Motivation

Quantum Field Theories

- Analytically solvable at weak coupling using perturbative methods
 - Some window into strong coupling via dualities (e.g. AdS/CFT)
- Lattice field theory
 - Non-perturbative formulation
 - Simulating dynamics is difficult (exponential size of Hilbert space)
 - Finite density is difficult (sign problem)



Gate complexity estimates

T gate counts

Gauge Group	Transport Coefficients	Heavy-Ion Collisions
$U(1)$	7.35×10^{17}	7.19×10^{23}
$SU(2)$	2.83×10^{34}	3.71×10^{40}
$SU(3)$	3.01×10^{49}	3.22×10^{55}

(Age of the universe $\sim 10^{17} s$)



Scalar Field Theory

ϕ^4 Hamiltonian

$$H = \int d^3x \mathcal{H} = \int d^3x \left[\overbrace{\frac{1}{2}\pi(\mathbf{x})^2 + \frac{1}{2}(\nabla\varphi(\mathbf{x}))^2 + \frac{1}{2}m^2\varphi(\mathbf{x})^2}^{H_0} + \underbrace{\lambda\varphi(\mathbf{x})^4}_{H_\varphi} \right]$$

Simplest field theory:

- Free theory -- relativistic energy-momentum relationship (Klein-Gordon equation)
- ϕ^4 term simplest self-interaction term bounded below
- Scalar fields in nature: Higgs boson, Inflaton(?)



Occupation Basis

Field expansion

$$\varphi(\mathbf{x}) = \sum_{\mathbf{p} \in \Gamma} \frac{1}{(2\pi)^d} e^{i\mathbf{p} \cdot \mathbf{x}} \sqrt{\frac{1}{2\omega_{\mathbf{p}}}} (a_{\mathbf{p}} + a_{-\mathbf{p}}^\dagger)$$

$$\pi(\mathbf{x}) = \sum_{\mathbf{p} \in \Gamma} \frac{-i}{(2\pi)^d} e^{i\mathbf{p} \cdot \mathbf{x}} \sqrt{\frac{\omega_{\mathbf{p}}}{2}} (a_{\mathbf{p}} - a_{-\mathbf{p}}^\dagger)$$

$$[a_p, a_q^\dagger] = (2\pi)^d \delta(p - q)$$

Free part

$$: H_0 := \frac{1}{(2\pi)^d} \sum_p \omega_p a_p^\dagger a_p$$

Interacting part

$$: H_\varphi := \frac{\lambda}{4!} \frac{1}{(2\pi)^d} \sum_{\vec{p}_1, \vec{p}_2, \vec{p}_3, \vec{p}_4 \in \Gamma} \delta \left(\sum_{i=1}^4 \vec{p}_i \right) \prod_{\vec{p}=\vec{p}_1}^{\vec{p}_4} \frac{1}{\sqrt{2\omega_{\vec{p}}}} : (a_{\vec{p}} + a_{-\vec{p}}^\dagger) :$$

Normal ordering

$$: a_p a_k^\dagger a_q : \equiv a_k^\dagger a_p a_q \quad \text{etc. (vanishing vacuum expectation value)}$$



Occupation Basis

One-hot encoding

$P^d = V$ distinct momentum modes

M distinct momentum states

$$\begin{aligned} |p, n\rangle &= |0_{1,1}, \dots, 0_{p-1,M}, 0_{p,1}, \dots, 0_{p,n-1}, 1_{p,n}, 0_{p,n+1}, \dots, 0_{p,M}, 0_{p+1,1}, \dots, 0_{V,M}\rangle \\ &= \left(\bigotimes_{j=1}^{p-1} |0_{j,1} \dots, 0_{j,M}\rangle \right) \bigotimes |0_{p,1}, \dots, 1_{p,n}, \dots, 0_{p,M}\rangle \bigotimes \left(\bigotimes_{j=p+1}^V |0_{j,1} \dots, 0_{j,M}\rangle \right) \end{aligned}$$

$$a_p^\dagger = \sum_n \sqrt{n+1} (\sigma_n^- \sigma_{n+1}^+)_p$$

$$a_p = \sum_n \sqrt{n+1} (\sigma_n^+ \sigma_{n+1}^-)_p,$$

$$\sigma^+ = \frac{1}{2}(X - iY), \quad \sigma_- = \frac{1}{2}(X + iY)$$



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Occupation Basis

Number operator

$$\hat{a}_p^\dagger \hat{a}_p = \hat{n}_p = \sum_n |1\rangle \langle 1|_{p,n} = \sum_n \left(\frac{I - Z}{2} \right)_{p,n}$$

Free Hamiltonian

$$: H_0 := \frac{1}{(2\pi)^d} \sum_p \omega_p a_p^\dagger a_p = \frac{1}{(2\pi)^d} \sum_{n,p} n \omega_p \left(\frac{I - Z}{2} \right)_{n,p}$$

$e^{-i:H_0:t}$ requires MV many R_Z gates



Occupation Basis

	$\#(c)R_z$	$\# T$	$\# \text{CNOT}$	$\# H$
e^{-iH_0t}	MV	-	-	-
$e^{-iH_1\varphi t}$	$\frac{M^4V^2(V-1)}{48}$	$\frac{M^4V^2(V-1)}{8M^3V^2}$	$\frac{11M^4V^2(V-1)}{20M^3V^2}$	$\frac{M^4V^2(V-1)}{24M^3V^2}$
$e^{-iH_2\varphi t}$	$\frac{M^3V^2}{3}$	$\frac{8M^3V^2}{3}$	$\frac{16M^3V^2}{3}$	$\frac{2M^3V^2}{3}$
$e^{-iH_3\varphi t}$	$2M^2V(V-1)$	$8M^2V(V-1)$	$16M^2V(V-1)$	$3M^2V(V-1)$
$e^{-iH_4\varphi t}$	$3MV$	-	$4MV$	$4MV$

$$p_3 = p_1 + k, \quad p_4 = p_2 - k$$

1) $p_1 \neq p_2, k \neq 0$

3) $p_1 \neq p_2, k = 0$

2) $p_1 = p_2, k \neq 0$

4) $p_1 = p_2, k = 0$



Linear Combination of Unitaries (LCU)

Given

$$A = \sum_{i=1}^L \alpha_i U_i \quad \alpha_i > 0$$

Construct

$$U_{PREP}|\bar{0}\rangle = \frac{1}{\sqrt{|\alpha|_1}} \sum_{i=1}^L \sqrt{\alpha_i} |i\rangle \quad |\alpha|_1 = \sum_{i=1}^L \alpha_i$$

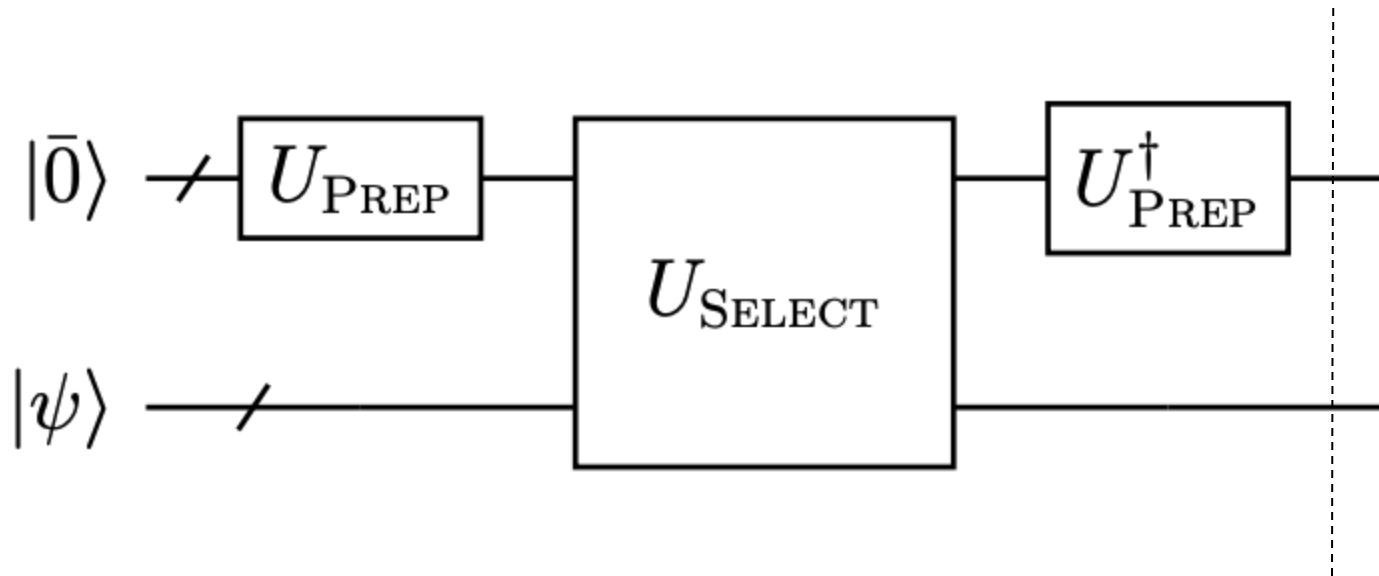
$$U_{SELECT} = \sum_{i=1}^L |i\rangle\langle i| \otimes U_i$$



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Linear Combination of Unitaries (LCU)



Success probability

$$\frac{||A|\psi\rangle||^2}{||\alpha||_1^2}$$

$$\frac{1}{|\alpha|_1} |\bar{0}\rangle \otimes A|\psi\rangle + |\Phi^\perp\rangle$$

$$(|\bar{0}\rangle\langle\bar{0}| \otimes \mathbb{I}) |\Phi^\perp\rangle = 0$$



Scalar Field Theory

ϕ^4 Hamiltonian

$$H = \sum_{\vec{x} \in \Omega} \left[\frac{1}{2} \pi^2(x) + (m^2 + D + 1) \phi^2(\vec{x}) + \frac{\lambda}{4!} \phi^4(\vec{x}) - 2 \sum_{i=1}^D \phi(\vec{x}) \phi(\vec{x} + a \hat{x}_i) \right]$$

Rescaled variables

$$\begin{aligned} \phi &= a^{\frac{D}{2}-1} \Phi, & \pi &= a^{\frac{D}{2}} \Pi, \\ m &= aM, & \lambda &= a^{4-D} \Lambda \end{aligned}$$

Field amplitude basis

$$\hat{\phi}|\phi\rangle = \phi|\phi\rangle$$

$$\hat{\phi}/\Delta\phi = \begin{pmatrix} -k+1 & & & \\ & \ddots & & \\ & & 0 & \\ & & & \ddots \\ & & & & k \end{pmatrix}$$



LCU with unequally sized groups of terms

$$\hat{H} = \alpha_0 \left(U_0^{(0)} + \cdots + U_{N_0-1}^{(0)} \right) + \cdots + \alpha_{M-1} \left(U_0^{(M-1)} + \cdots + U_{N_{M-1}-1}^{(M-1)} \right)$$

Let $N_{max} = \max\{N_i\}_{i=0}^{M-1}$,

$$U_{PREP} |\bar{0}\rangle = \frac{1}{\sqrt{|\alpha|_1 N_{max}}} \left(\sum_{i=0}^{M-1} \sqrt{\alpha_i N_i} |i\rangle \right) \left(\sum_{m=0}^{N_i-1} |m\rangle \right) \left(\sum_{k=0}^{\frac{N_{max}}{N_i}-1} |k\rangle \right)$$

$$U_{SELECT} = \sum_{i=0}^{M-1} |i\rangle \langle i| \otimes U_{SUB-SELECT_i}$$

$$U_{PREP}^\dagger U_{SELECT} U_{PREP} |\bar{0}\rangle |\psi\rangle = |\bar{0}\rangle \frac{\hat{H}}{|\alpha|_1} |\psi\rangle$$



LCU decompositions

Equal weight LCU

$$\frac{\hat{\phi}}{\Delta\phi} = \frac{1}{2} \sum_{m=0}^{2k-1} U^{(m)}$$

$$U^{(m)} = - \sum_{i=0}^{m-1} |i\rangle\langle i| + \sum_{i=m}^{2k-1} |i\rangle\langle i| = \sum_{i=0}^{2k-1} [2\Theta(i-m) - 1] |i\rangle\langle i|$$

$$\hat{\phi}/\Delta\phi = \begin{pmatrix} -k+1 & & & \\ & \ddots & & \\ & & 0 & \\ & & & \ddots \\ & & & & k \end{pmatrix}$$

Unequal weight LCU^[1]

$$\frac{\hat{\phi}}{\Delta\phi} = \frac{1}{2} \mathbb{I} - \frac{1}{2} \sum_{j=0}^{\log_2 2k-1} 2^j Z_j$$

$O(|\Omega| \log_2^4 k)$ T gates

#qubits $\sim O(|\Omega| \log_2 k)$
(either way)

$$n = \frac{1}{2} \left(\underbrace{+1 \cdots +1}_{k+n} \quad \underbrace{-1 \cdots -1}_{k-n} \right)$$



Equal weight LCU

Use comparator (adapted from Gidney adder^[1])

$$\mathbf{CMP}_{A,B,C} |i\rangle |j\rangle |0\rangle = |i\rangle |j\rangle |j < i\rangle = |i\rangle |j\rangle |\Theta(i - j - 1)\rangle$$

$O(|\Omega| \log_2 k)$ T gates

to simplify

$$\begin{aligned} & U_{PREP_\phi}^\dagger U_{SELECT_\phi} U_{PREP_\phi} |\bar{0}\rangle |\psi\rangle |0\rangle_{anc} \\ &= \left(H^{\otimes m} \mathbf{CMP}^\dagger Z_{anc} \mathbf{CMP} H^{\otimes m} \right) |\bar{0}\rangle |\psi\rangle |0\rangle_{anc} \\ &= \left(|\bar{0}\rangle \frac{\hat{\phi}}{\phi_{max}} |\psi\rangle + |\Phi^\perp\rangle \right) |0\rangle_{anc} \end{aligned}$$

[1] "Halving the cost of quantum addition", Craig Gidney, Quantum 2, 74 (2018)



Equal weight LCU

Other terms in the Hamiltonian

$$\left(\frac{\hat{\phi}}{\Delta\phi}\right)^2 = \frac{1}{2} \sum_{i=0}^{2k^2-1} U^{(i)}, \quad \text{where}$$

$$U^{(i)} = \sum_{j=0}^{2k-1} [2\Theta(k^2 + (k-j-1)^2 - i - 1) - 1] |j\rangle\langle j|$$

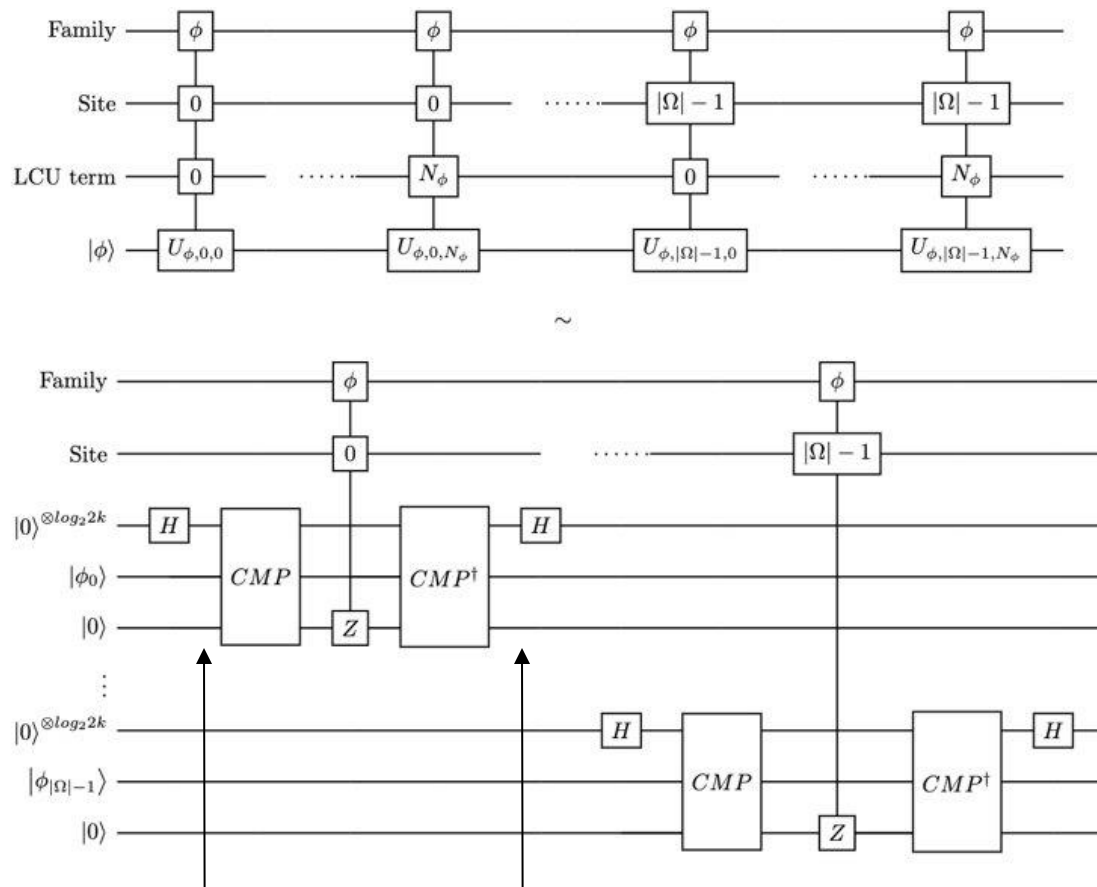
$$\pi^2 = \mathcal{F}^\dagger \phi^2 \mathcal{F}$$

$$\left(\frac{\hat{\phi}}{\Delta\phi}\right)^4 = \frac{1}{2} \sum_{i=0}^{2k^4-1} U^{(i)}, \quad \text{where}$$

$$U^{(i)} = \sum_{j=0}^{2k-1} [2\Theta(k^4 + (k-j-1)^4 - i - 1) - 1] |j\rangle\langle j|$$



Equal weight LCU



$O(|\Omega| \log_2^2 k)$ T gates

Insert $U_{initial}$... and its inverse for ϕ^2, π^2, ϕ^4
(involves arithmetic and Fourier transform)



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T gate complexity

$$\text{Count}(\mathbf{T})_{\phi\phi} \sim O(|\Omega| D \log_2(|\Omega| D k))$$

$$\text{Count}(\mathbf{T})_{\phi^2, \pi^2, \phi^4} \sim O(|\Omega| (\log_2^2 k + \log_2 |\Omega|))$$

$$k = \frac{\phi_{max}}{\Delta\phi} \quad |\Omega| : \text{Lattice size} \quad D : \text{spatial dimensionality}$$



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