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SO₂, silicate clouds, but no CH₄ detected in a warm Neptune

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Abstract

WASP-107b is a warm (\sim 740 K) transiting planet with a Neptune-like mass of \sim 30.5 M_{\oplus} and Jupiter-like radius of \sim 0.94 $R_{\rm J}$ [1, 2] whose extended atmosphere is eroding [3]. Previous observations showed evidence for water vapour and a thick high-altitude condensate layer in WASP-107b's atmosphere [4, 5]. Recently, photochemically produced sulphur dioxide (SO₂) was detected in the atmosphere of a hot (\sim 1,200 K) Saturn-mass planet from transmission spectroscopy near 4.05 μ m [6, 7], but for temperatures below \sim 1,000 K sulphur is predicted to preferably form sulphur allotropes instead of SO₂ [8, 9, 10]. Here we report the 9 σ -detection of two fundamental vibration bands of SO₂, at 7.35 μ m and 8.69 μ m, in the transmission spectrum of WASP-107b using the Mid-Infrared Instrument (MIRI) of the JWST. This discovery establishes WASP-107b as the second irradiated exoplanet with confirmed photochemistry, extending the temperature range of exoplanets exhibiting detected photochemistry from \sim 1,200 K down to \sim 740 K. Additionally, our spectral analysis reveals the presence of silicate clouds, which are strongly favoured (\sim 7 σ) over simpler cloud setups. Furthermore, water is detected (\sim 12 σ), but methane is not. These findings provide evidence of disequilibrium chemistry and indicate a dynamically active atmosphere with a super-solar metallicity.

WASP-107b was observed with JWST MIRI on 19-20 January 2023. The SLITLESSPRISM subarray of the low-resolution spectrometer was used, offering a spectral resolution ranging from 30 and 100 over a wavelength span of 4.61 to 11.83 μ m. We performed three independent data reductions using the CASCADe [11], Eureka! [12], and TEATRO packages; see Supplementary Information (SI). Each method extracted 51 spectroscopic light curves. For all channels, we obtained a minimal level of correlated noise in the residuals, consistent with normally distributed noise. The 1σ error displayed a minimum of 80 ppm at $7.5\,\mu$ m. The transmission spectra derived from the different reductions, shown in Figure 1 and tabulated in Extended Data Table 1, are within 3σ agreement and 95% of the points within 2σ ; see SI.

We performed atmospheric retrievals using two independent frameworks, ARCiS [13] and petit-RADTRANS (pRT) [14], including both our JWST data and previous near-infrared (1.121–1.629 μ m) HST data [4]. Free abundance retrievals were run including the following species: H₂O, CO, CO₂, CH₄, C₂H₂, SO₂, SO, H₂S, SiO, HCN, NH₃, and PH₃. The remaining atmosphere consisted of H₂ and He. A variety of cloud models were tested, ranging from cloud-free to more complex models, the latter focusing mostly on silicate clouds (MgSiO₃, SiO₂, and SiO); see SI.

Figure 2 shows the best fit to the data, including main contributions from molecular species and clouds to the spectrum. The figure presents the results based on the CASCADe package, but our conclusions are consistent across the three data reductions. Both retrieval codes detect SO_2 at $\sim 9\sigma$, H_2O at $\sim 12\sigma$, and the presence of high-altitude clouds at $\sim 9\sigma$, with a $\sim 7\sigma$ preference for silicate clouds over more simple cloud setups. We also tentatively detect H_2S ($\sim 4\sigma$), NH_3 ($\sim 2-3\sigma$), and CO ($\sim 2-3\sigma$), although the CO detection relies on the first three spectral points and requires confirmation at shorter wavelengths. CH_4 is not detected, with an upper limit of its volume mixing ratio (VMR) being a few times 10^{-6} . Table 1 presents the detection significance and VMR for each species. Due to differences in cloud structure setups, the absolute VMRs are different between the two retrieval codes. We therefore focus the discussion on the detection significance and relative abundances.

The MIRI data of WASP-107b presents the mid-infrared discovery of SO_2 in an exoplanet atmosphere. Both the ν_1 symmetric stretch and the ν_3 asymmetric stretch vibration bands of SO_2 (with fundamental frequency at $8.69\,\mu\mathrm{m}$ and $7.35\,\mu\mathrm{m}$, respectively [15]) are detected (Figure 1 and Figure 2). Recently, the $(\nu_1 + \nu_3)$ combination band of SO_2 was assigned as carrier of a spectral feature near $4.05\,\mu\mathrm{m}$ detected in the JWST NIRspec spectrum of WASP-39b[6, 16], a hot ($\sim 1,200\,\mathrm{K}$) irradiated Saturn-mass exoplanet. The MIRI detection of SO_2 in WASP-107b extends the range of exoplanet temperatures with detected SO_2 from $\sim 1,200\,\mathrm{K}$ down to $\sim 740\,\mathrm{K}$.

The measured VMR of SO_2 in WASP-107b (see Table 1) is several orders of magnitude higher than expected for chemical equilibrium, which predicts H_2S to be the dominant sulphur-bearing molecule for Neptune-like planets [9]. Disequilibrium processes can drive abundances considerably away from chemical equilibrium, with photochemistry and atmospheric transport being the dominant mechanisms [17]. Notably, the SO_2 feature observed in WASP-39b has been attributed to photochemical processes occurring within its atmosphere [10]. Indeed, UV irradiation initiates the photodissociation of H_2O , yielding H and OH radicals. These OH radicals are key for oxidising sulphur that is liberated from H_2S [10]. However, these models [10] predict that SO_2 would not be detectable using JWST MIRI for a planet with an equilibrium temperature of \sim 740 K. This is in contrast with our detection of SO_2 in the atmosphere of WASP-107b.

To unravel the production paths of SO_2 in WASP-107b, we computed a grid of disequilibrium models (see SI). To ensure an accurate consideration of the upper-atmosphere chemistry, we observed – contemporaneously with the JWST observations – the Near-Ultraviolet (NUV) emission of the host star WASP-107 with Neil Gehrels Swift. Additionally, we reanalysed the 2018 X-ray emission observed with XMM-Newton (see SI). Figure 3 provides evidence that only models incorporating photochemistry in combination with a super-solar metallicity predict a detectable level of SO_2 in WASP-107b. Key disparities from prior predictions [10, 17] stem from UV radiation and gravity. Previous models adopted a gravity of 1,000 cm/s² [10] and 2,140 cm/s² [17], while WASP-107b's gravity is $\sim 260 \, \text{cm/s}^2$. Moreover, the NUV flux is a factor $\sim 200 \, \text{lower}$ for WASP-107b than for WASP-39b, and a factor $\sim 100 - 1,000 \, \text{lower}$ in the FUV (see SI). A lower gravity, an overall decreased UV flux, and a low FUV/NUV ratio collectively contribute to the increased formation of SO_2 ; see Extended Data Figure 1.

The overarching scenario that unfolds is that the initiating pathways for SO_2 formation in WASP-107b are twofold. First, H_2O photodissociation in upper layers at pressures $\lesssim 10^{-5}$ bar generates atomic H and OH radicals, leading to sequential oxidation of sulphur liberated from H_2S . Second, in the pressure range of $\sim 10^{-5}-1$ bar, photolysis of various abundant molecules – beyond just H_2O – provides free atoms and radicals, partially redistributed through eddy diffusion. This initiates a cascade of thermochemical reactions that progressively yield sufficient OH radicals for SO_2 oxidation. Given the fact that a large fraction of these reactions are temperature-independent and barrierless, the equilibrium temperature is not the sole determinant for SO_2 formation. As long as the UV irradiation and FUV/NUV ratio remain moderate and the gravity is low, these processes lead to detectable SO_2 levels even at the low equilibrium temperature of WASP-107b.

The sensitivity to metallicity can be attributed to the larger abundance of sulphur and OH radicals at higher metallicities. At a metallicity of $6\times$ solar, the SO₂ spectral features contribute partially to the 7.8 μ m feature and dominate at $\gtrsim 10\times$ solar (see Extended Data Figure 2). For metallicities $\geq 6\times$ solar, the gas-phase photochemical models predict a SO₂ VMR above $\sim 5\times 10^{-7}$

at pressures between $10^{-7}-10^{-4}$ bar (see Figure 3). This result is corroborated by the retrieval outcomes, which show that the spectral contribution of SO_2 predominantly comes from the region above the high-altitude cloud layer (situated at pressures of a few times 10^{-5} bar; see Extended Data Figure 4). This establishes SO_2 as a key diagnostic for the atmospheric metallicity in exoplanet atmospheres.

At super-solar metallicity, our models also predict a detectable CH₄ feature (see Extended Data Figure 2). The confirmation of the HST CH₄ non-detection [4] in the MIRI wavelength range raises questions regarding the predicted predominance of CH₄ in the atmosphere [18, 19]. One potential reason for the non-detection of CH₄ is that it is hidden by the high-altitude cloud layer. In contrast to SO₂, CH₄ is photochemically destroyed at high altitudes, i.e. above the cloud layer (see Extended Data Figure 3). Another explanation for the non-detection of CH₄ is the influence of the strong irradiation experienced by WASP-107b, which can significantly impact the thermodynamical properties of its atmosphere [20, 21, 22]. WASP-107b is highly inflated, pointing to a deep atmosphere that has significantly higher temperatures [23, 24, 25, 26] than those predicted by conventional models. The latter suggest an intrinsic heat flux ($\sigma T_{\rm int}^4$, with σ the Boltzmann constant and $T_{\rm int}$ the intrinsic temperature) with corresponding $T_{\rm int}$ of 150 K [27]. We computed chemical models with different intrinsic temperatures ranging from 250-600 K to reproduce a hot deep atmosphere (see Extended Data Figure 3). Our results show that the CH₄ molar fraction in the atmospheric layers probed by our MIRI observations is reduced by more than 3 ordersof-magnitude when we increase $T_{\rm int}$ from 250 to 600 K. Interestingly, the SO₂ abundance in the detectable upper atmosphere is almost unaffected since sulphur species do not quench in deep regions [9] (see Extended Data Figure 3). Lowering the C/O ratio from solar (= 0.55) to 0.10 yields an additional decrease of the CH₄ abundance by a factor ~ 15 (see Extended Data Figure 3).

The presence of silicate clouds in the atmospheres of strongly irradiated planets has been predicted for a long time [28]. Although a hint of silicate emission has been claimed for one single object [29], clear observational evidence has been lacking. Recently, silicate clouds were detected in a young, self-luminous planet at large separation from its host star [30]. The MIRI data of WASP-107b present the detection of the $10 \,\mu m$ Si-O stretching mode of solid silicate particles. The retrieval models favour a high-altitude cloud layer composed of small (sub-micron) amorphous silicate particles. At the pressure levels and temperatures where we find the cloud layer, gas phase SiO can nucleate and condensate efficiently to form small solid state mineral particles [31]. However, according to the traditional picture on cloud formation, these particles would eventually rain down to deeper, hotter layers of the atmosphere, depleting the relatively cold upper atmosphere from gas phase SiO. In the deeper, hotter layers, the particles evaporate, allowing the gas phase material to mix and serve as building blocks for new nucleation/condensation cycles. Thus, we can conclude that the presence of silicate clouds at such high altitudes in a relatively cold part of the atmosphere is another indication of strong mixing from either a hot inner atmosphere (similar to the finding related to CH₄) or potentially from the intensely irradiated hot dayside.

Figure 1: **JWST MIRI transmission spectrum of WASP-107b.** Comparison of the JWST MIRI transmission spectra obtained from the three independent reductions considered in this work (coloured points, with 1- σ error bars). The blue dots show the results from the CASCADe code, the red dots using the Eureka! package and the black dots are from the TEATRO routines. The thick dashed grey line indicates the band-averaged transit depth from the CASCADe analysis at 20,463 ppm and the shaded area the 95% confidence interval of 39 ppm. The right y-axis gives the planetary spectrum in units of atmospheric scale height of the planetary atmosphere, assuming a hydrogen dominated atmosphere. The spectra have a constant wavelength bin width, corresponding to a spectral resolution of 50 at 7.5 μ m. The two horizontal bars indicate the ν_1 symmetric stretch and ν_3 asymmetric stretch vibration bands of SO₂.

Figure 2: Transmission spectrum of WASP-107b with key contributions. The (red and pink) points with 1- σ error bars correspond to the measured JWST MIRI CASCAde transit depths of the spectrophotometric light curves at different wavelengths. The (red and pink) points in the area indicated by the grey band are from HST. The median spectrum model predicted from ARCiS [13] retrievals is shown in blue and from petitradtrans (property) [14] in green (offset by 2,000 ppm for clarity). The shaded regions of the model spectra correspond to the 1-,2-,3- σ credibility envelopes predicted by the retrievals. The bottom part of the figure shows the silicate cloud and molecular contributions for those gases inferred by our analysis of WASP-107b's spectrum.

Figure 3: **Predicted SO**₂ **molar fraction.** Panel (a): Adopted temperature-pressure (T-P) profile for WASP-107b following the analytical equation of ref. [32]; see SI. Panel (b): Predicted SO₂ molar fraction for four values of the metallicity $(1, 2, 6, \text{ and } 10 \, \text{Z}_{\odot}; \text{ in black, green, red, and blue, respectively)}$ for models with (full line) and without (dotted line) the inclusion of photochemistry. Predictions are calculated for a model with an intrinsic temperature of $400 \, \text{K}$, a solar C/O ratio, and a $\log_{10}(K_{zz}, \text{cgs}) = 10$.

Table 1: Outcome of the retrieval analyses. Given are the detection significance and volume mixing ratios (VMR) of the various components in the ARCiS and petitradtrans (prt) retrieval analysis.

	Significa	nce $[\sigma]$	$\log_{10}(\mathrm{VMR})$		
Component	ARCiS	pRT	ARCiS	pRT	
$\overline{\mathrm{H_2O}}$	13.2	12.5	$-2.19^{+0.42}_{-0.26}$	$-3.81^{+0.38}_{-0.29}$	
SO_2	8.8	9.1	$-5.03^{+0.33}_{-0.18}$	$-6.72^{+0.30}_{-0.23}$	
H_2S	4.7	3.5	$-2.65^{+0.49}_{-0.38}$	$-3.88^{+0.41}_{-0.33}$	
NH_3	2.3	3.4	$-5.47^{+0.34}_{-0.66}$	$-6.04_{-0.25}^{+0.30}$	
CO	2.8	2.3	$-2.41^{+0.29}_{-0.28}$	$-4.58^{+0.40}_{-0.59}$	
PH_3	$_(b)$	$_(a)$	$-6.29_{-2.24}^{+1.\overline{29}}$	$-7.60^{+0.89}_{-1.39}$	
HCN	$_(b)$	$_(a)$	$-9.26^{+1.73}_{-1.63}$	$-9.19_{-1.15}^{+1.20}$	
C_2H_2	_(b)	$_(a)$	$-9.08^{+1.73}_{-1.73}$	$-9.19^{+1.19}_{-1.14}$	
SiO	_(b)	$_(a)$	$-6.08^{+1.90}_{-3.34}$	$-9.03^{+1.47}_{-1.40}$	
CH_4	_(b)	$_(b)$	$-8.52^{+2.09}_{-2.09}$	$-8.83^{+1.25}_{-1.23}$	
CO_2	_(b)	$_(b)$	$-8.05^{+2.46}_{-2.37}$	$-8.49^{+1.79}_{-1.70}$	
SO	$_(b)$	$_(c)$	$-7.38^{+2.76}_{-2.74}$	1110	
${\bf Silicate~Clouds}$	7.2	7.1			

 $^{^{(}a)}$ included in retrieval, but not tested for detection significance because posterior indicates an upper limit.

(b) not favoured (Bayes factor < 2).

(c) not included.

- Supplementary Information is linked to the online version of the paper at www.nature.com/nature.
- Data availability. The JWST MIRI data presented in this paper are part of the JWST MIRI GTO program (Program identifier (PID) 1280; P.I. P.O. Lagage). The JWST data will be publicly available in the Barbara A. Mikulski Archive for Space Telescopes (MAST; https://archive.stsci.edu/) after January 20, 2024, and can be found either using the program identifier or using the Data Object Identifier (DOI): https://doi.org/10.17909/as3s-x893. The DOI link can also be used to retrieve the publicly available HST WFC3 data used in this paper from MAST. The NUV data are in an online publicly-accessible archive: https://heasarc.gsfc.nasa.gov/cgi-bin/W3Browse/swift.pl. The HST data are available in the MAST archive and can be found using DOI: https://doi.org/10.17909/as3s-x893. The XMM X-ray data archive is available at https://nxsa.esac.esa.int/nxsa-web/#search.

- Code availability. The codes used in this publication to extract, reduce, and analyse the data are as follows. The data reduction pipeline jwst can be found at https://jwst-p ipeline.readthedocs.io/en/latest/. The data analysis codes Eureka!, CASCADe, and TEATRO can be found respectively at https://eurekadocs.readthedocs.io/en/latest/, https://gitlab.com/jbouwman/CASCADe, and https://github.com/ncrouzet/TEATRO. The CASCADe-filtering, and CASCADe-jitter sub-packages can be found, respectively, at https://gitlab.com/jbouwman/CASCADe-filtering.and https://gitlab.com/jbouw man/CASCADe-jitter. The atmospheric model codes used to fit the data can be found at https://www.exoclouds.com/ for the ARCiS-code [13] and at https://petitradtrans. readthedocs.io/en/latest/ for the petitRADTRANS-code [14]. The XSPEC package [33] is available at https://heasarc.gsfc.nasa.gov/xanadu/xspec/. The HEASoft package (including the Swift UVOT tools) is available at https://heasarc.gsfc.nasa.gov/lheaso ft/download.html (version 6.31.1 was used in our work). The NUV Swift-project pipeline is available at https://swift.gsfc.nasa.gov/quicklook/swift_process_overview.html. The VULCAN chemical network can be found at https://github.com/exoclime/VULCAN/b lob/master/thermo/SNCHO_photo_network.txt.
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- Competing Interests. The authors declare that they have no competing financial interests.

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92 Extended Data

Extended Data – Table 1: Measured transit depth of WASP-107b. First and fifth column list the wavelength, and the other columns the transit depths as obtained with Eureka!, CASCADe and TEATRO. Each wavelength bin has a full width of $0.15~\mu m$.

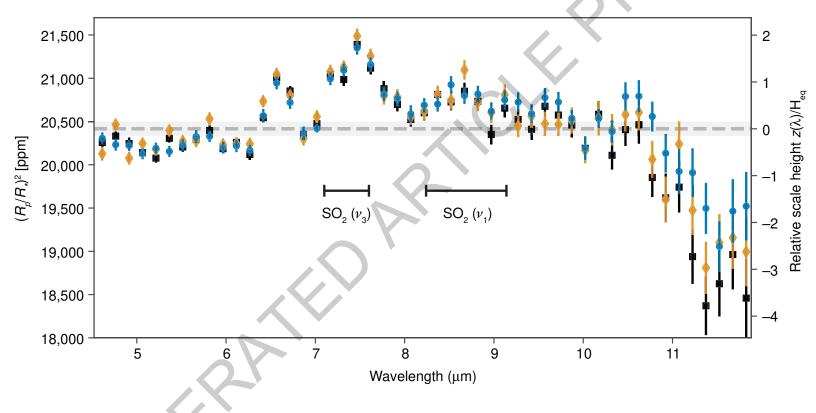
Extended Data – Figure 1: Impact of gravity and UV irradiation on predicted SO₂ molar fraction. The base model (shown in black in each panel) has an intrinsic temperature of 400 K, a solar C/O ratio, a metallicity of $10 \times$ solar, a $\log_{10}(K_{zz}, cgs) = 10$, and uses the SED of HD 85512 [34] – used as a proxy for WASP-107 – as input stellar spectrum (see SI). Panel (a): Temperature-pressure (T-P) profile for a gravity g of $2.6 \,\mathrm{m/s^2}$ (black, purple and brown), $4.3 \,\mathrm{m/s^2}$ (blue), and $10 \,\mathrm{m/s^2}$ (orange). Panel (b): Predicted SO₂ molar fractions for different gravity values with colours corresponding to panel (a). While the black curve uses the HD 85512 SED as input spectrum, the purple curve uses the WASP-39 SED instead (see SI), and the brown curve the HD 85512 SED scaled with a factor 100.

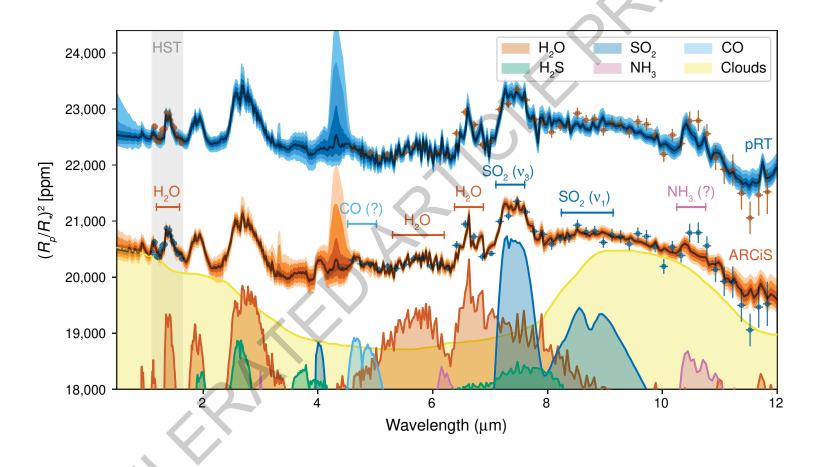
Extended Data – Figure 2: **Predicted transit depth for WASP-107b assuming a cloudless atmosphere.** The four panels display for a metallicity of 1, 2, 6, and $10 \, \mathrm{Z}_{\odot}$, the total transit depth (in black) and the transit depth without relative contributions of SO_2 (light blue), $\mathrm{H}_2\mathrm{O}$ (dark blue), CO_2 (light green), CH_4 (pink) and NH_3 (orange). Predictions are calculated for a model with an intrinsic temperature of $400 \, \mathrm{K}$, a solar C/O ratio, and a $\log_{10}(K_{\mathrm{zz}}, \, \mathrm{cgs}) = 10$.

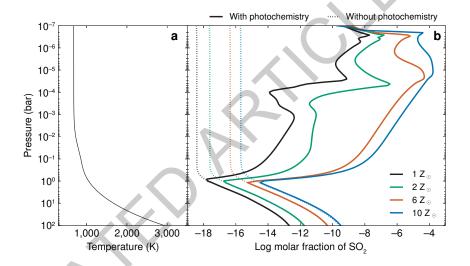
Extended Data – Figure 3: Sensitivity of SO_2 and CH_4 molar fractions to various input parameters. Shown are the impact of the intrinsic temperature, C/O ratio and eddy diffusion coefficient on predicted SO_2 and CH_4 molar fractions. The base model (shown in black in each panel) has an intrinsic temperature of $400 \, \text{K}$, a solar C/O ratio, a metallicity of $10 \times \text{solar}$, and a $\log_{10}(K_{zz}, \text{cgs}) = 10$. Panel (a): Temperature-pressure (T-P) profile for intrinsic temperatures of 250, 400, and $600 \, \text{K}$ (light blue, black, and orange, respectively). Panel (b): Predicted SO_2 (full line) and CH_4 (dotted line) molar fractions for different T-P structures (and hence intrinsic temperatures), with colours corresponding to panel (a). Panel (c): Predicted SO_2 and CH_4 molar fraction for different C/O ratios. Comparison with predictions for a solar C/O ratio (of 0.55, black line) for which either the carbon or oxygen atomic abundance has been adapted (indicated by 'C' or 'O' between parenthesis, respectively). Panel (d): Predicted SO_2 and CH_4 molar fraction for different values of the eddy diffusion coefficient K_{zz} .

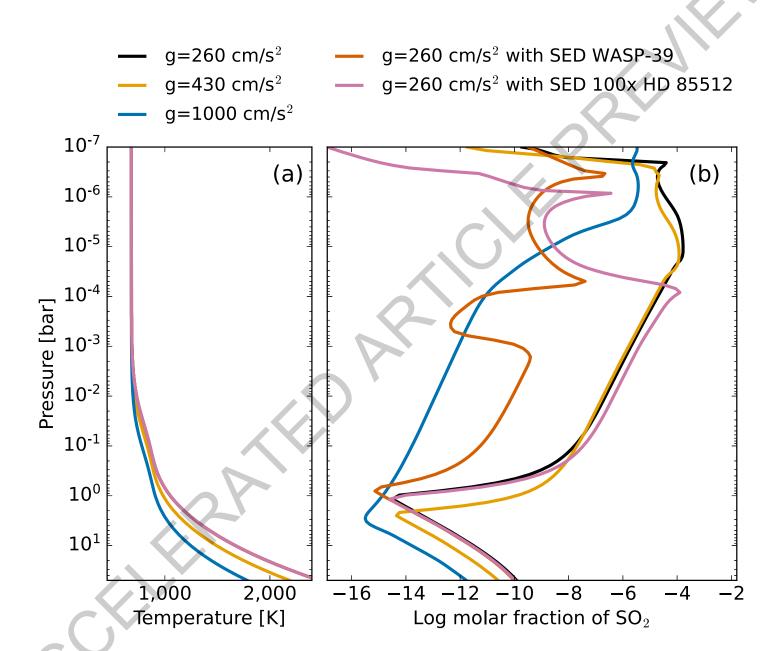
Extended Data – Figure 4: Contribution function for the atmosphere as retrieved by ARCiS. Shown in the contour colours are the relative contributions of various layers in the atmosphere to the transit spectrum as a function of wavelength. The dark horizontal line in this plot is located at the geometrically thin cloud layer, dominating the spectrum at all wavelengths. The spectral variation on top of this baseline predominantly comes from below the cloud layer at the wavelengths where water features are present, while it is dominated by regions above the cloud layer for the spectral features of SO_2 . This is consistent with the pressures where we expect SO_2 to be abundant in the atmosphere from our photochemical modelling (see Figure 3). The molecular contribution to the spectral variation in the $9.5-11~\mu m$ region is significantly fainter than at other wavelengths. This is the region where the silicate feature is most prominent and thus also the spectral variation is dominated by the cloud layer.

Extended Data – Figure 5: **Silicate cloud detection significance.** The detection significance, measured in σ , of the silicate cloud model is presented in comparison to both the cloud-free model (indicated by the blue dashed curve) and the parameterised cloud model (represented by the red solid line) as a function of the maximum wavelength used in the analysis. Even when limiting the analysis to wavelengths below $10 \,\mu\text{m}$, the silicate cloud remains favoured, at a significance of 5.7σ , over the parameterised cloud configuration.

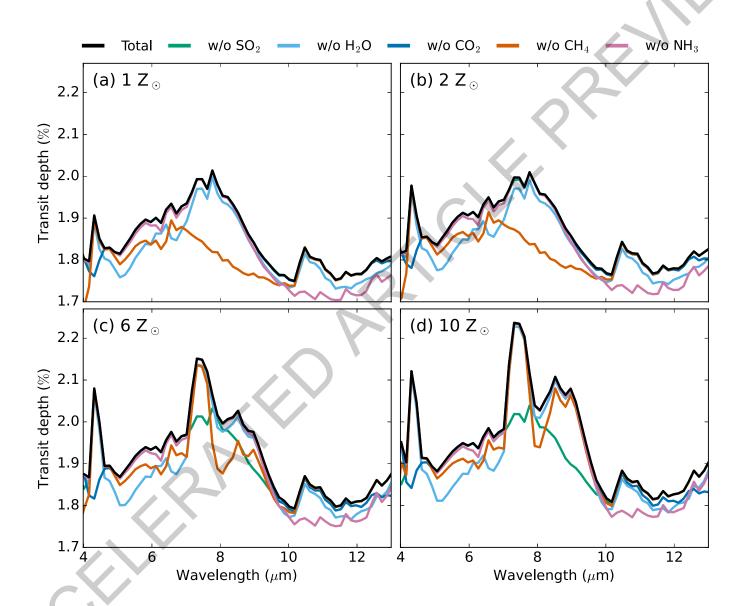




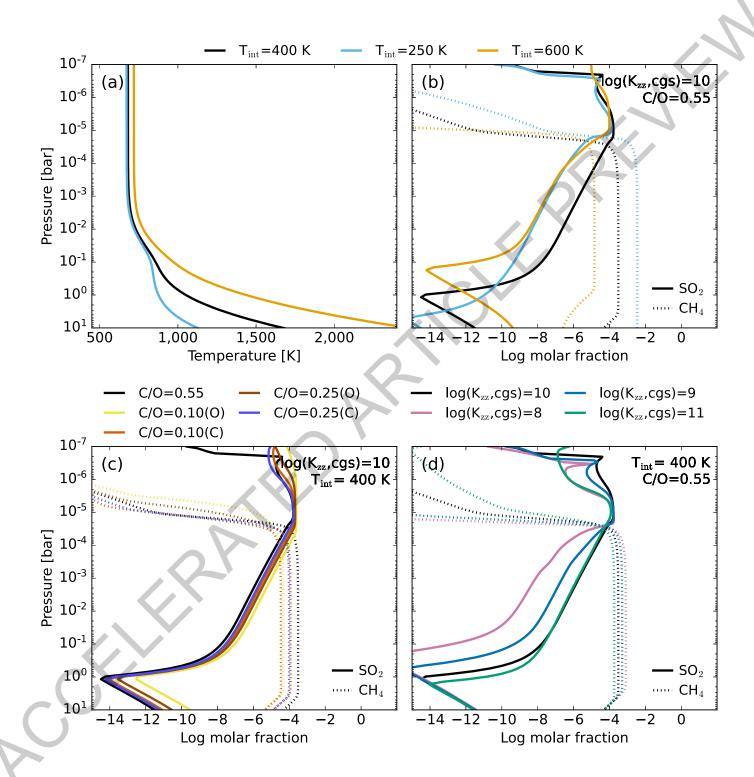




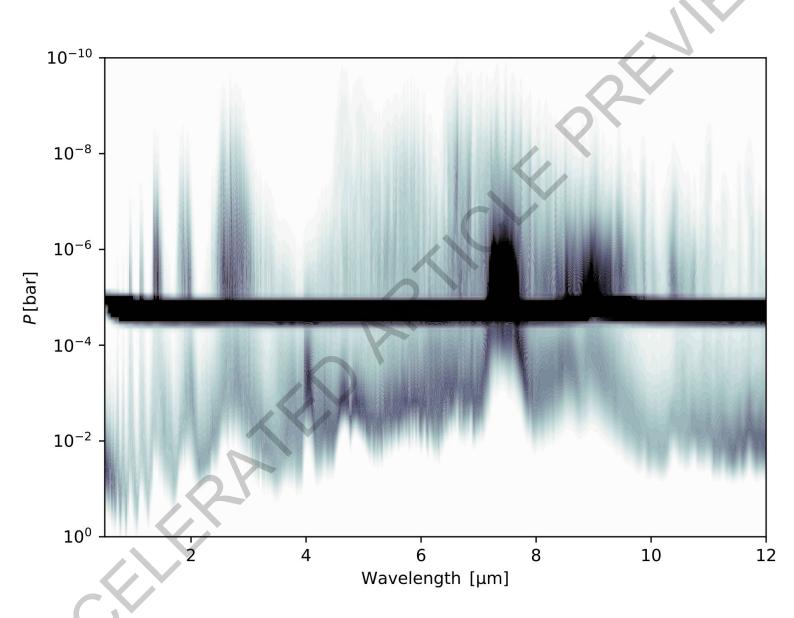
Extended Data Fig. 1



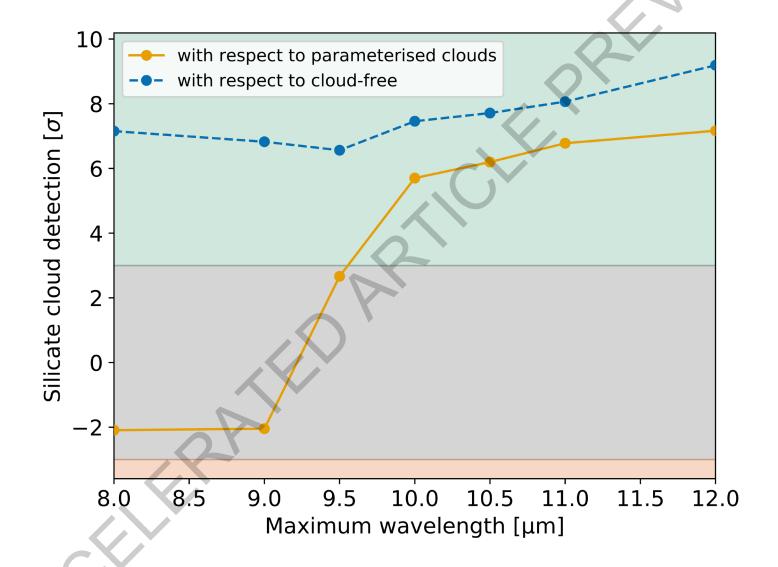
Extended Data Fig. 2



Extended Data Fig. 3



Extended Data Fig. 4



Extended Data Fig. 5

λ		Transit depth		λ		Transit depth	
	Eureka!	CASCAde	TEATRO		Eureka!	CASCADe	TEATRO
(μm)	(ppm)	(ppm)	(ppm)	(μm)	(ppm)	(ppm)	(ppm)
4.61	$20,\!126\pm80$	$20,\!305\pm69$	$20,\!257 \pm 73$	8.37	$20,813 \pm 100$	$20,701 \pm 87$	$20,811 \pm 93$
4.76	$20,464 \pm 60$	$20,233 \pm 70$	$20,331 \pm 52$	8.52	$20{,}745\pm106$	$20,925 \pm 99$	$20{,}727 \pm 98$
4.91	$20,075 \pm 73$	$20,223 \pm 63$	$20,245 \pm 70$	8.67	$21,095 \pm 112$	$20,800 \pm 96$	$20,850 \pm 99$
5.06	$20,245 \pm 60$	$20,128 \pm 64$	$20,139 \pm 53$	8.82	$20,725 \pm 108$	$20,818 \pm 96$	$20,733 \pm 95$
5.21	$20,178 \pm 60$	$20,192 \pm 64$	$20,077 \pm 52$	8.97	$20,603 \pm 118$	$20,620 \pm 100$	$20,350 \pm 112$
5.36	$20,396 \pm 55$	$20,154 \pm 63$	$20,306 \pm 46$	9.12	$20,807 \pm 118$	$20,750 \pm 114$	$20,654 \pm 107$
5.51	$20,272 \pm 69$	$20,227 \pm 70$	$20,206 \pm 58$	9.27	$20,447 \pm 125$	$20,724 \pm 114$	$20,523 \pm 114$
5.67	$20,274 \pm 65$	$20,326 \pm 72$	$20,292 \pm 53$	9.42	$20,562 \pm 132$	$20,590 \pm 110$	$20,411 \pm 122$
5.82	$20,530 \pm 76$	$20,329 \pm 67$	$20,396 \pm 65$	9.57	$20,475 \pm 137$	$20,775 \pm 111$	$20,676 \pm 127$
5.97	$20,234 \pm 64$	$20,196 \pm 66$	$20,186 \pm 55$	9.72	$20,466 \pm 127$	$20,725 \pm 117$	$20,574 \pm 122$
6.12	$20,243 \pm 68$	$20,217 \pm 72$	$20,258 \pm 59$	9.87	$20,512 \pm 149$	$20,537 \pm 94$	$20,460 \pm 141$
6.27	$20,240 \pm 66$	$20,166 \pm 64$	$20,118 \pm 61$	10.02	$20,166 \pm 148$	$20,190 \pm 134$	$20,190 \pm 139$
6.42	$20,733 \pm 66$	$20,567 \pm 78$	$20,544 \pm 55$	10.17	$20,539 \pm 174$	$20,535 \pm 110$	$20,584 \pm 154$
6.57	$21,045 \pm 77$	$20,945 \pm 73$	$21,006 \pm 65$	10.32	$20,404 \pm 186$	$20,384 \pm 125$	$20,109 \pm 165$
6.72	$20,799 \pm 67$	$20,720 \pm 74$	$20,853 \pm 58$	10.47	$20,580 \pm 192$	$20,789 \pm 164$	$20,406 \pm 189$
6.87	$20,301 \pm 79$	$20,364 \pm 74$	$20,351 \pm 73$	10.62	$20,606 \pm 213$	$20,792 \pm 186$	$20,462 \pm 219$
7.02	$20,555 \pm 72$	$20,420 \pm 47$	$20,475 \pm 65$	10.77	$20,062 \pm 208$	$20,558 \pm 172$	$19,853 \pm 226$
7.17	$21,076 \pm 78$	$20,993 \pm 73$	$21,044 \pm 69$	10.92	$19,600 \pm 266$	$20,133 \pm 223$	$19,618 \pm 282$
7.32	$21,128 \pm 76$	$21,093 \pm 77$	$20,985 \pm 71$	11.07	$20,239 \pm 262$	$19,923 \pm 260$	$19,742 \pm 294$
7.47	$21,486 \pm 88$	$21,351 \pm 79$	$21,390 \pm 77$	11.22	$19,474 \pm 278$	$19,909 \pm 282$	$18,938 \pm 316$
7.62	$21,257 \pm 83$	$21,160 \pm 83$	$21,117 \pm 73$	11.37	$18,810 \pm 295$	$19,496 \pm 295$	$18,371 \pm 339$
	$20,797 \pm 100$		$20,880 \pm 85$				$18,625 \pm 378$
	$20,782 \pm 93$		$20,697 \pm 81$,	•	$18,963 \pm 404$
	$20,554 \pm 90$	$20,587 \pm 101$	$20,524 \pm 83$		$18,994 \pm 397$		
	$20,\!617\pm98$	$20,688 \pm 82$	$20,601 \pm 89$		_	_	_