

Soft X-ray Energy Spectra in the Wide-Field Galactic Disk Area Revealed with HaloSat

2 KAZUKI AMPUKU,¹ IKUYUKI MITSUISHI,¹ KOKI SAKUTA,¹ PHILIP KAARET,^{2,3} DANIEL M. LARocca,³ AND
3 LORELLA ANGELINI⁴

4 ¹*Graduate School of Science, Division of Particle and Astrophysical Science, Nagoya University, Furo-cho, Chikusa-ku, Nagoya, Aichi,*
5 *464-8602, Japan*

6 ²*NASA/Marshall Space Flight Center, Huntsville, AL 35812*

7 ³*University of Iowa Department of Physics and Astronomy, Van Allen Hall, 30 N. Dubuque Street, Iowa City, IA 52242, USA*

8 ⁴*NASA/Goddard Space Flight Center, Greenbelt, MD 20771*

9 ABSTRACT

10 We analyzed data from *HaloSat* observations for five fields in the Galactic disk located far away
11 from the Galactic center ($135^\circ < l < 254^\circ$) to understand the nature of soft X-ray energy emission
12 in the Galactic disk. The fields have 14° diameter and were selected to contain no significant high-
13 flux X-ray sources. All five *HaloSat* soft X-ray energy spectra (0.4–7 keV with energy resolution of
14 < 100 eV below 1 keV) show a possibility of the presence of unresolved high-temperature plasma
15 in the Galactic disk (UHTPGD) with a temperature of 0.8–1.0 keV and an emission measure of $(8$ –
16 $11) \times 10^{-4} \text{ cm}^{-6} \text{ pc}$ in addition to the soft X-ray diffuse background components mainly studied at
17 higher galactic latitudes (solar wind charge exchange emission, local hot bubble, Milky Way halo
18 emission, and the cosmic X-ray background). This suggests that the UHTPGD is present across the
19 whole Galactic disk. We also observed UHTPGD emission in a region with no bright sources in an
20 *XMM-Newton* field contained within one of the *HaloSat* fields. The temperature and emission measure
21 are consistent with those measured with *HaloSat*. Moreover, the stacked spectra of the X-ray point-like
22 sources and NIR-identified point sources such as stars in the *XMM-Newton* field also show a spectral
23 feature similar to the UHTPGD emission. This suggests that the UHTPGD may partly originate from
24 point-like sources such as stars.

25 *Keywords:* X-rays: diffuse background — Galaxy: disk — X-rays: ISM — Galaxy: halo — X-rays:
26 stars

27 1. INTRODUCTION

28 Diffuse background emission across many wavebands
29 provides important astronomical information. Signifi-
30 cant time and effort have been invested to understand
31 the nature of the diffuse backgrounds in the microwave,
32 infrared, and X-ray bands (e.g., [Penzias & Wilson 1965](#);
33 [Hauser & Dwek 2001](#); [Shanks et al. 1991](#)). In the X-ray
34 band, the diffuse background above 2 keV is dominated
35 by the cosmic X-ray background (CXB) that originates
36 mainly from discrete extragalactic sources (e.g., [Hickox](#)
37 [& Markevitch \(2006\)](#)), predominantly AGNs at a wide
38 range of redshifts. The CXB is well described empiri-
39 cally by a powerlaw with a photon index of ~ 1.4 in

40 the 2–10 keV band (e.g., [Vecchi et al. \(1999\)](#); [Kushino](#)
41 [et al. \(2002\)](#); [De Luca & Molendi \(2004\)](#)). However,
42 [McCammon et al. \(2002\)](#) reveals that the CXB is re-
43 sponsible for only ~ 40 % of the total emission in the
44 soft X-ray diffuse background (SXDB) in the ~ 0.4 –
45 keV band. This shows the existence of other compo-
46 nents. Contributors to the SXDB include: emission
47 due to geocoronal and heliospheric solar wind charge
48 exchange (SWCX) processes (e.g., [Snowden et al. 2004](#);
49 [Ezoe et al. 2010](#); [Yoshitake et al. 2013](#); [Kuntz 2019](#)), hot
50 gas in the solar neighborhood referred to as the Local
51 Hot Bubble (LHB) (e.g., [Liu et al. 2017](#); [Farhang et al.](#)
52 [2019](#)), and hot plasma in the Milky Way Halo (MWH)
53 (e.g., [Yoshino et al. 2009](#); [Sakai et al. 2014](#); [Nakashima](#)
54 [et al. 2018](#)).

55 However, the situation changes in the Galactic disk
56 because dense neutral material with $N_{\text{H}} \sim 10^{21} \text{ cm}^{-2}$

57 or more significantly blocks soft X-ray photons below
 58 1 keV arising from the MWH and CXB. Although it
 59 was expected that the SXDB below 1 keV should de-
 60 crease by at least by $\sim 40\%$, the SXDB intensity in the
 61 *ROSAT* R45 band ($\sim 0.44\text{--}1.0$ keV) is only $\sim 20\%$ lower
 62 in the Galactic disk compared to higher Galactic lati-
 63 tudes (McCammion et al. 2002). This suggests the exist-
 64 ence of widely-distributed excess emission that partly
 65 compensates for the absorption of the MWH and CXB.

66 This mystery has been known as the ‘‘M band prob-
 67 lem’’ (McCammion & Sanders 1990; Cox 2005). The
 68 term M band is from an energy band used in the pro-
 69 portional counters in Wisconsin and the Nagoya-Leiden
 70 rocket programs; the energy range is almost equal to the
 71 *ROSAT* R45 band. The origin of the excess emission
 72 has been under debate for a long time. Possible candi-
 73 dates include: hot gas with a temperature of $\sim 3 \times 10^6$
 74 K, dM stars, young expanding superbubbles or super-
 75 nova remnants in low density fields (Nousek et al. 1982;
 76 Sanders et al. 1983; Rosner et al. 1981; Cox 2005).

77 Masui et al. (2009) conducted a deep pointed obser-
 78 vation of a $18' \times 18'$ blank sky field located in the Galac-
 79 tic disk ($l, b = (235^\circ, 0^\circ)$) to investigate spectral char-
 80 acteristics of the excess emission. They used *Suzaku*
 81 which had low and stable non-X-ray background espe-
 82 cially in the soft energy band and was thus suitable
 83 for study of low-surface brightness soft X-ray sources
 84 such as the SXDB (Mitsuda et al. 2007). They removed
 85 point-like X-ray sources to extract a spectrum of the dif-
 86 fuse emission and assumed that the MWH component is
 87 completely blocked by dense neutral material associated
 88 with the Galactic disk. They revealed an excess compo-
 89 nent that was well modeled by an unabsorbed, optically-
 90 thin collisionally-ionized plasma in thermal equilibrium
 91 with a temperature of ~ 0.8 keV. In this paper, we de-
 92 fined it as unresolved high-temperature plasma in the
 93 Galactic disk (UHTPGD). The lack of absorption sug-
 94 gests that the excess X-ray photons come from relatively
 95 nearby sources. Masui et al. (2009) considered a large
 96 fraction of the UHTPGD originates from faint dM stars
 97 and their model spectrum for spatially unresolved dM
 98 stars reproduces the observed features although it can-
 99 not compensate entirely the decrease at high Galactic
 100 latitudes due to the small scale height. However, the
 101 results are only for a single field. It is important to test
 102 for the presence of UHTPGD across a wide span of the
 103 Galactic disk.

104 *HaloSat* was a small satellite (CubeSat) designed to
 105 map the entire soft X-ray sky in order to measure the
 106 MWH. It was deployed from the *International Space*
 107 *Station* in 2018 (Kaaret et al. 2019) and re-entered
 108 Earth’s atmosphere in early 2021. *HaloSat* had a large

109 field of view (FoV) of near 100 square degrees and a
 110 moderate energy resolution of < 100 eV below 1 keV
 111 (Kaaret et al. 2019; Zajczyk et al. 2020; LaRocca et al.
 112 2020). It was non-imaging and had no X-ray focusing
 113 optics on board. The large FoV and good soft X-ray sen-
 114 sitivity of *HaloSat* enable us to study diffuse soft X-ray
 115 sources with large angular extent such as the SXDB.

116 In this paper, we analyze *HaloSat* data to investigate
 117 whether UHTPGD emission is found across the extent
 118 of the Galactic disk. We selected five *HaloSat* fields lo-
 119 cated in the outer Galactic disk that avoid bright X-ray
 120 sources enabling analysis of the diffuse emission. The
 121 rest of the paper is organized as follows. Section 2
 122 presents our field selection procedure and data reduc-
 123 tion. Section 3, describes our spectral analysis methods
 124 and results for the *HaloSat* data. In Section 4, we dis-
 125 cuss the origin of the detected high temperature plasma
 126 and its relation to the spectra of point-like sources mea-
 127 sured with *XMM-Newton*. Finally, we summarize our
 128 results in Section 5.

129 2. FIELD SELECTION AND DATA REDUCTION

130 To investigate the X-ray spectral properties of wide-
 131 field areas in the Galactic disk, we searched for *HaloSat*
 132 fields in outer Galactic disk satisfying the following con-
 133 ditions.

- 134 • aim point with $|b| < 2^\circ$ and $|l| > 50^\circ$,
- 135 • exposure times > 20 ksec to have good photon
 136 statistics,
- 137 • containing no X-ray bright objects in the *ROSAT*
 138 all sky survey (Boller et al. 2016) or the
 139 MAXI/SSC all sky catalog (Tomida et al. 2016)
- 140 • containing no extended soft X-ray sources such as
 141 the Cygnus superbubble or Vela supernova rem-
 142 nant (e.g., Bluem et al. 2020; Silich et al. 2020).

143 The final sample contains a total of 5 fields that are
 144 presented in Table 1. Data for the five fields were down-
 145 loaded from the *HaloSat* archive at the HEASARC.
 146 Spectra were obtained for each of the three detectors
 147 following standard screening on the hard (3–7 keV)
 148 and very large event (> 7 keV) count rates as de-
 149 scribed in the *HaloSat* analysis document (see details in
 150 <https://heasarc.gsfc.nasa.gov/docs/halosat/analysis/>).

151 We used HEASoft v6.28 to extract the spec-
 152 tra and XSPEC version 12.11 for the spectral fit-
 153 ting. The response matrix and ancillary response
 154 to apply in the spectral data and as well the re-
 155 sponse for the particle-induced instrumental back-
 156 ground were obtained from the *HaloSat* calibration

Table 1. *HaloSat* and *XMM-Newton* observations

Obs. ID	Field name	Aim point [deg.]		Exposure [ksec]
		(R.A. _{J2000} , Dec. _{J2000})	(<i>l</i> , <i>b</i>)	
<i>HaloSat</i>				
HS0017	HaloSat J0530+3400	(05 30 33.12, +34 00 18.0)	(174.00, 0.00)	32*
HS0019	HaloSat J0320+5715	(03 20 05.52, +57 15 43.2)	(142.00, 0.00)	23*
HS0021	HaloSat J0759-2954	(07 59 59.28, -29 54 14.4)	(247.00, 0.00)	28*
HS0022	HaloSat J0729-1746	(07 29 16.08, -17 46 51.6)	(233.00, 0.00)	43*
HS0023	HaloSat J0700-0430	(07 00 42.00, -04 30 32.4)	(218.00, 0.00)	27*
<i>XMM-Newton</i>				
0500240101	V838 Mon	(07 04 04.85, -03 50 51.1)	(217.80, 1.05)	49 [†]

* The effective observation exposure after data cleaning of the S14 detector.

[†] Net exposure time of the PN detector after the standard data screening.

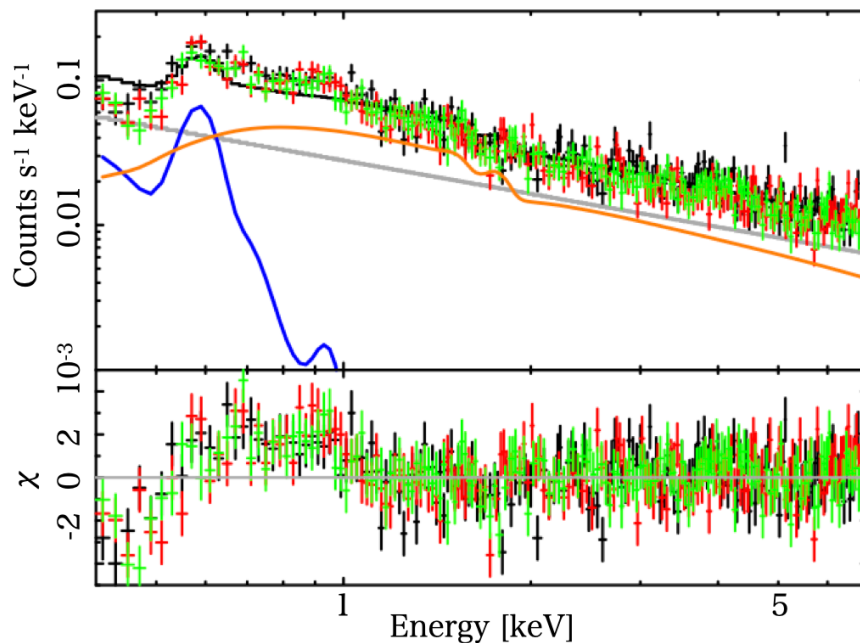


Figure 1. Example of the *HaloSat* spectral fitting results for the simple model (model0) consisting of the LHB (blue) and CXB (orange). The instrumental background (gray) was modelled with as a powerlaw with a diagonal response matrix. Black, red, and green crosses show the observed data (upper) and residuals subtracting the model from the data divided by error bars (lower) for each detector. For simplicity, the best fit models are shown only for one detector.

157 at the HEASARC ([https://heasarc.gsfc.nasa.gov/docs/](https://heasarc.gsfc.nasa.gov/docs/heasarc/caldb/halosat/halosat_caldb.html)
 158 [heasarc/caldb/halosat/halosat_caldb.html](https://heasarc.gsfc.nasa.gov/docs/heasarc/caldb/halosat/halosat_caldb.html)). Through-
 159 out the paper, if not otherwise specified, errors are
 160 quoted at the 90% confidence level and the solar abun-
 161 dance table of Wilms et al. (2000) was utilized.

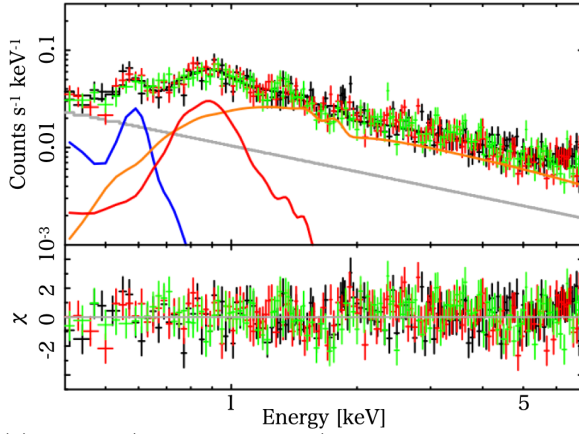
162 3. ANALYSIS AND RESULTS

163 We analyzed the spectra of the five *HaloSat* fields. Our
 164 spectral fitting procedure is based on previous works
 165 modeling the SXDB (e.g., Mitsuishi et al. 2013). Each
 166 of the five *HaloSat* fields were fitted independently. We

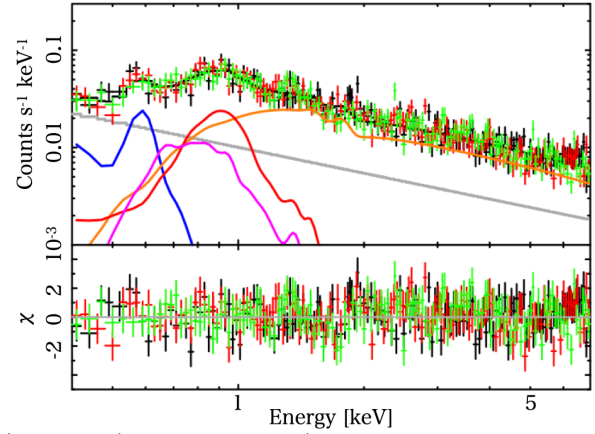
167 fit the spectra for all three detectors simultaneously us-
 168 ing common parameters for the astrophysical emission.

169 The instrumental background was modeled as a pow-
 170 erlaw and allowed to vary between detectors (Silich et al.
 171 2020; Kaaret et al. 2020). The instrumental background
 172 photon index and 1-sigma error were calculated for each
 173 detector for each observation based on the empirical re-
 174 lation described in [https://heasarc.gsfc.nasa.gov/docs/](https://heasarc.gsfc.nasa.gov/docs/halosat/analysis/back20210209.pdf)
 175 [halosat/analysis/back20210209.pdf](https://heasarc.gsfc.nasa.gov/docs/halosat/analysis/back20210209.pdf) obtained using ac-
 176 tual data. In our fitting, the instrumental background
 177 normalization was a free parameter for each detector
 178 while the photon index was allowed to vary within the

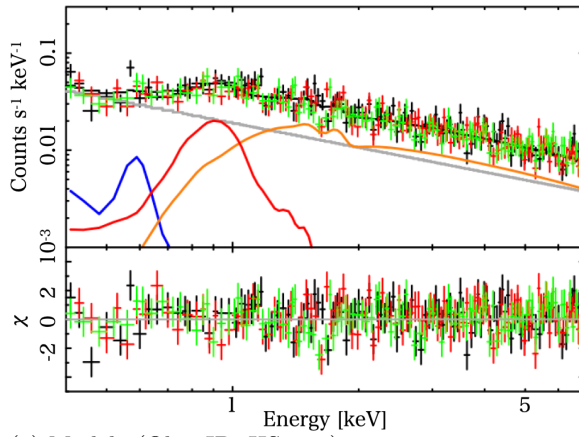
(a) Model1 (Obs. ID: HS0017)



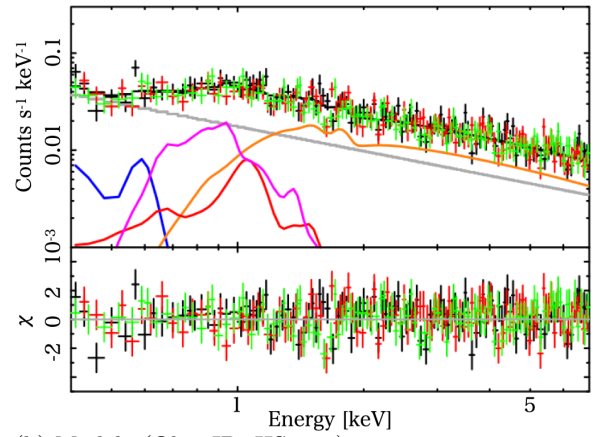
(b) Model2 (Obs. ID: HS0017)



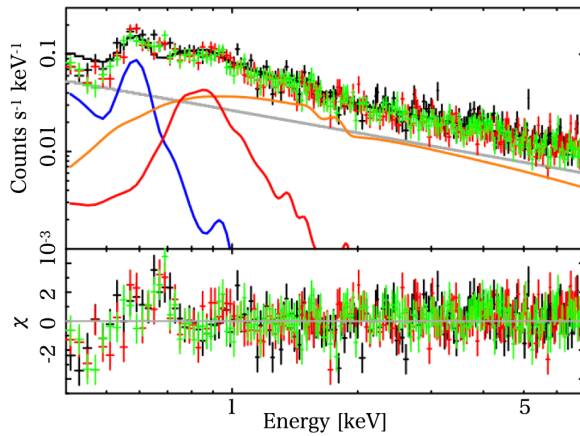
(a) Model1 (Obs. ID: HS0019)



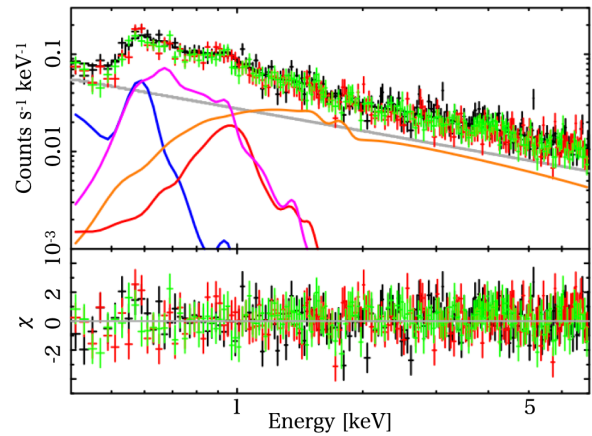
(b) Model2 (Obs. ID: HS0019)



(a) Model1 (Obs. ID: HS0021)



(b) Model2 (Obs. ID: HS0021)



179 1-sigma uncertainty from the empirical relation. In the
180 models below, the instrumental background component
181 is denoted as $powerlaw_{IB}$.

182 As a first step, we adopted a simple model consist-
183 ing of only the SWCX, LHB and CXB components,
184 i.e. assuming that the UHTP emission is negligible and
185 the MWH is completely blocked by the cold dense neu-
186 tral material of the Galactic disk. We used one unab-
187 sorbed optically thin thermal collisionally-ionized equi-
188 librium (CIE) plasma model for the sum of the SWCX
189 and LHB emission and an absorbed powerlaw model
190 for the CXB. We refer to this as Model0 which is
191 defined in XSPEC as $apec_{SWCX+LHB} + tbabs_{Galactic} \times$
192 $powerlaw_{CXB} + powerlaw_{IB}$, where $tbabs_{Galactic}$ gives
193 the absorption along the line-of-sight direction which
194 is mainly associated with the Galactic disk. In this
195 model, the emission measure of the thermal plasma
196 and the Galactic absorption column density were free
197 parameters. We fixed the *apec* redshift to 0, the
198 *apec* abundance to 1, and the powerlaw photon index
199 to 1.4 (Kushino et al. 2002). The plasma tempera-
200 ture was limited to the range 0.097 ± 0.014 keV and
201 the surface brightness of the powerlaw was limited to
202 $(6.38 \pm 0.64) \times 10^{-8}$ erg cm $^{-2}$ s $^{-1}$ sr $^{-1}$ based on previous
203 studies (Kushino et al. 2002; Liu et al. 2017).

204 Fits to this model show significant residuals (see
205 Fig. 1) for all fields except HS0019, especially around
206 1 keV corresponding to Fe L-shell complex. This sug-
207 gests that additional thermal components are needed
208 as found in Masui et al. (2009). Additionally, residu-
209 als in the 0.5–0.7 keV band, corresponding to He- and
210 H- like oxygen emission lines, were found in all of the
211 fields except HS0019, which supports the above sce-
212 nario. There seems to be no significant residuals for
213 HS0019 because the spectrum below ~ 1 keV is rel-
214 atively flat. However, the fitted column density for
215 HS0019 and all the other fields is significantly smaller
216 ($\ll 10^{21}$ cm $^{-2}$) than the expected values based on mul-
217 tiwavelength observations (Kalberla et al. 2005; Zhu
218 et al. 2017), which indicates that the spectrum needs
219 an additional component to make up for soft emis-
220 sion that should be absorbed. Thus, an additional
221 soft component is also needed for HS0019. **Note that**
222 **we considered a heliospheric SWCX scenario for**
223 **the observed residuals around both 0.5–0.7 and**
224 **0.8–1.1 keV by adding multiple gaussian models**
225 **corresponding to oxygen, neon, and iron emis-**
226 **sion lines as shown in Ringuette et al. (2021)**
227 **and Huang et al. (2023). We found that addi-**
228 **tion of SWCX emission can reduce the residuals.**
229 **However the absorption column density remains**
230 **much lower than the predicted values, suggest-**

231 **ing the need for contributions from an additional**
232 **spectral component. We also considered non-**
233 **solar abundances for the LHB component as dis-**
234 **cussed in McCammon et al. (2002) and reached**
235 **the same conclusions.** The normalization of the in-
236 strumental background show typical values and no sig-
237 nificant difference is seen in the three detectors.

238 Therefore, as a next step, we added an unabsorbed
239 thermal plasma component for the UHTPGD emis-
240 sion, that we denote as $apec_{UHTPGD}$. We denote
241 this as Model1 which is defined as: $apec_{SWCX+LHB} +$
242 $tbabs_{Galactic} \times powerlaw_{CXB} + apec_{UHTPGD} +$
243 $powerlaw_{IB}$. This model assumes that all of the hot
244 gas associated with the MWH is blocked completely by
245 the dense neutral material as is the case with Masui
246 et al. (2009). Model1 provides a significantly better fit
247 compared to Model0 including only the LHB and CXB
248 components because the residuals around 1 keV were
249 well fitted by the additional thermal plasma compo-
250 nent. The UHTPGD components have temperatures
251 of 0.7–0.8 keV and emission measure in the range $(8-$
252 $15) \times 10^{-4}$ cm $^{-6}$ pc. The fitted column density became
253 much larger, in the range $(0.1-0.7) \times 10^{22}$ [cm $^{-2}$], which
254 is almost consistent with the expected values for each
255 field. The UHTPGD temperature is consistent with
256 that in Masui et al. (2009) and interestingly, the surface
257 brightness values are consistent with each other within
258 the statistical errors even though the surface bright-
259 ness in Masui et al. (2009) is calculated from only one
260 pointed observation of *Suzaku* with a FoV of $18' \times 18'$.
261 However, residuals around 0.7 keV persist in several
262 fields (HS0021, HS0022, and HS0023). The residuals
263 could alternatively be explained by an LHB component
264 with higher temperature near 0.2 keV if the tempera-
265 ture constraint was not applied. However, the degree of
266 the deviation in the temperature of the LHB in these
267 large fields seems unreasonable considering previous
268 work (Kushino et al. 2002; Liu et al. 2017). We con-
269 firmed that the UHTPGD parameters do not change
270 significantly if a high-temperature LHB is allowed. As
271 another possibility, we tried another fitting model by
272 setting an oxygen abundance free. We confirmed that
273 the results of the temperature and emission measure of
274 the UHTPGD are consistent with each other within the
275 statistical errors even though the resultant parameters
276 are physically unacceptable because the oxygen abun-
277 dance values are too high over 3. These results suggest
278 that one more softer thermal component emitting H-like
279 oxygen emissions such as the hot gas associated with
280 the MWH is needed.

281 Thus, we added another thermal plasma component,
282 but with absorption equal to that for the CXB, cor-

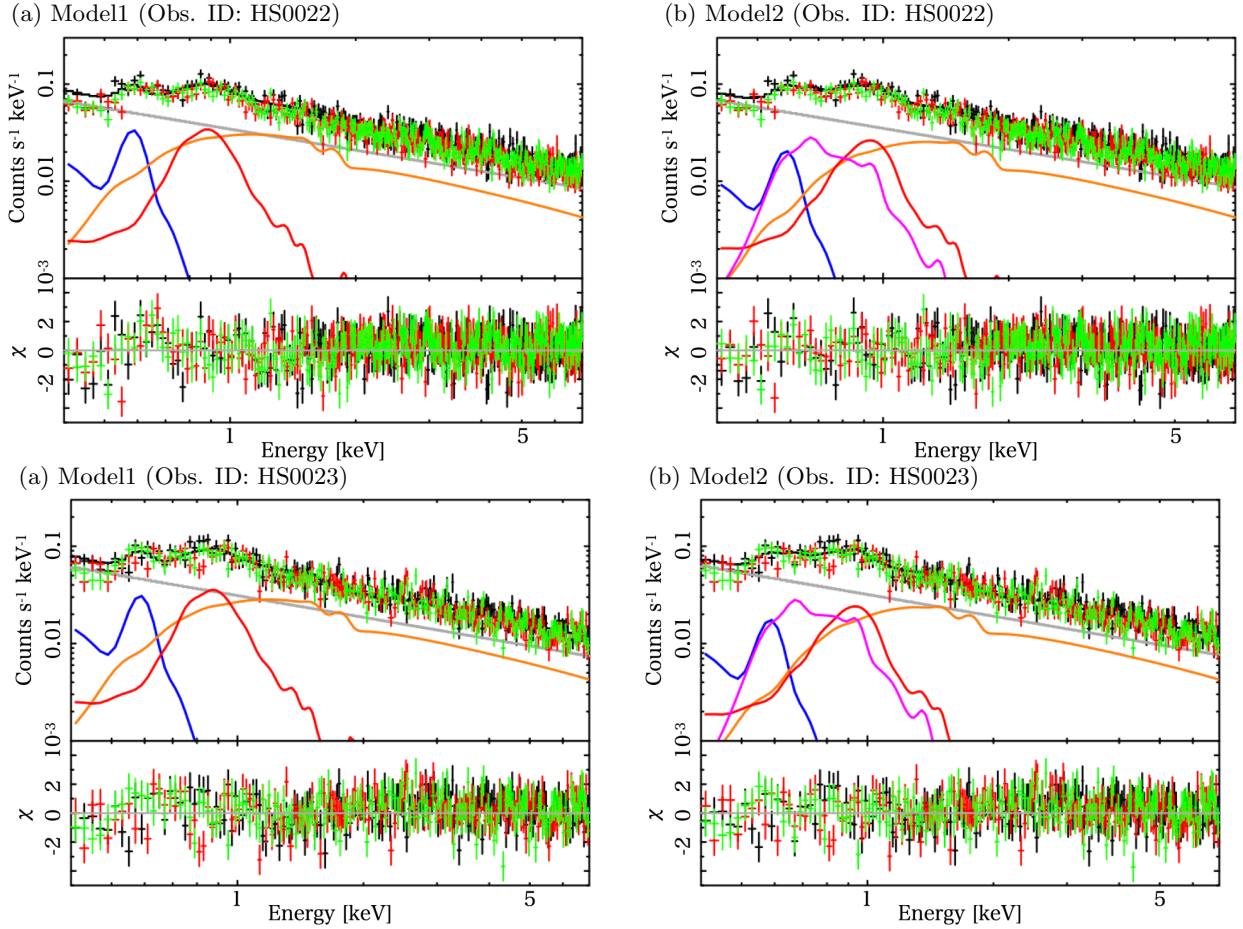


Figure 2. Observed *HaloSat* spectra in the wide fields of the Galactic disk fitted with Model1 (left) and Model2 (right) for the selected five fields. Black, red, and green crosses show the observed data (upper) and residuals subtracting the model from the data divided by error bars (lower) for each detector. Blue and magenta, and orange solid lines correspond to the unabsorbed and absorbed optically thin thermal CIE plasmas corresponding to the LHB and MWH, and the absorbed non-thermal model known as the CXB, respectively. Red and gray lines indicate the excess high temperature plasma and instrumental background models, respectively. For simplicity, the best fit models are shown only for one detector.

Table 2. The results of the model fitting of the five selected *HaloSat* fields.

Obs. ID	Model	N_{H} [$\times 10^{22} \text{ cm}^{-2}$]	kT_{UHTPGD} [keV]	$\text{EM}_{\text{UHTPGD}}^*$ [$\times 10^{-4} \text{ cm}^{-6} \text{ pc}$]	$\chi^2/\text{d.o.f}$
HS0017	Model 1	$0.31^{+0.11}_{-0.08}$	$0.77^{+0.05}_{-0.05}$	11^{+3}_{-2}	452/399
	Model 2	$0.42^{+0.43}_{-0.12}$	$0.84^{+0.61}_{-0.08}$	$9.6^{+3.2}_{-5.9}$	445/397
HS0019	Model 1	$0.74^{+0.20}_{-0.18}$	$0.84^{+0.11}_{-0.11}$	$8.2^{+2.6}_{-2.6}$	336/316
	Model 2	$0.90^{+0.23}_{-0.18}$	$1.3^{+0.6}_{-0.3}$	$6.7^{+7.6}_{-4.3}$	321/314
HS0021	Model 1	$0.098^{+0.046}_{-0.038}$	$0.67^{+0.06}_{-0.07}$	15^{+2}_{-2}	624/493
	Model 2	$0.32^{+0.15}_{-0.09}$	$0.95^{+0.31}_{-0.13}$	$8.4^{+2.2}_{-2.3}$	494/491
HS0022	Model 1	$0.22^{+0.07}_{-0.06}$	$0.77^{+0.04}_{-0.04}$	13^{+2}_{-2}	790/702
	Model 2	$0.37^{+0.10}_{-0.08}$	$0.88^{+0.08}_{-0.06}$	11^{+2}_{-2}	747/700
HS0023	Model 1	$0.27^{+0.13}_{-0.10}$	$0.75^{+0.06}_{-0.06}$	13^{+3}_{-3}	599/509
	Model 2	$0.46^{+0.17}_{-0.12}$	$0.88^{+0.27}_{-0.09}$	10^{+2}_{-5}	569/507

* Emission Measure integrated over the line of sight, i.e. $\text{EM} = \int n_e n_{\text{H}} dl$.

283 responding to the MWH hot gas. The temperature
 284 was limited to 0.225 ± 0.023 keV, the abundance was
 285 fixed to 0.3 solar, and redshift was fixed to 0, follow-
 286 ing the results obtained in high-latitude regions (Kaaret
 287 et al. 2020). This model, hereafter Model2, is de-
 288 scribed in XSPEC as: $apec_{\text{SWCX+LHB}} + tbabs_{\text{Galactic}} \times$
 289 $(powerlaw_{\text{CXB}} + apec_{\text{MWH}}) + apec_{\text{UHTPGD}} + powerlaw_{\text{IB}}$.
 290 Fitting with Model2 decreases the residuals around
 291 0.7 keV due to the addition of the MHW hot gas. The
 292 goodness of fit improves significantly ($>99.9\%$ in the F-
 293 test) for the fields, HS0021, HS0022, and HS0023, and
 294 somewhat ($<99.9\%$ in the F-test) for the fields, HS0017
 295 and HS0019. The spectral analysis results for Model1
 296 and Model2 are summarized in Table 2 and the spectral
 297 plots and residuals in Figure 2. As for the other models,
 298 the normalization of the instrumental background has
 299 typical values and no significant difference is seen in the
 300 three detectors.

301 The best fitted column density, temperature, and
 302 emission measure of the UHTPGD components are $(0.3-$
 303 $0.9) \times 10^{22}$ [cm^{-2}], 0.8–1.3 keV, and $(7-11) \times 10^{-4}$
 304 pc , respectively. The best fit values of the column
 305 density and temperature are systematically higher than
 306 those of Model1 and the emission measure decreases
 307 as expected with the addition of the MWH compo-
 308 nent. The emission measure of the LHB is consistent
 309 with those of Model1 within the statistical errors ex-
 310 cept HS0021 and the MWH emission measure is consis-
 311 tent amongst the fields within the statistical errors. For
 312 HS0017, all of the results are consistent within the sta-
 313 tistical errors between Model1 and Model2. For HS0019,
 314 the nominal value of the column density increases sig-
 315 nificantly but is consistent with the expected value.

316 In summary, all five HaloSat soft X-ray energy spectra
 317 show a possibility of the presence of UHTPGD with a
 318 temperature of 0.8–1.0 keV in the best fit models and the
 319 parameters are consistent with those found in a smaller
 320 field with *Suzaku* (Masui et al. 2009). These results
 321 suggest that UHTPGD is distributed across much of the
 322 Galactic disk.

323 4. A STELLAR ORIGIN?

324 To investigate the relation of the observed excess high
 325 temperature plasma to point-like X-ray sources includ-
 326 ing stars, we analyzed *XMM-Newton* archival data be-
 327 cause *XMM-Newton* has relatively large FoV and effec-
 328 tive area and good angular resolution which allow us to
 329 study properties of both diffuse and point-like sources.
 330 We searched the *XMM-Newton* archive for observations
 331 within our *HaloSat* fields with long exposure times. We
 332 excluded fields with bright diffuse sources such as clus-
 333 ters of galaxies and supernova remnants. We selected

334 the observation (ObsID: 0500240101) that lies within
 335 the *HaloSat* field HS0023. We use only the PN data
 336 since the PN has the largest effective area. We applied
 337 the standard data screening on, e.g., time filtering to
 338 remove flaring high-background periods (see (e.g., Mit-
 339 suishi et al. 2013)), resulting in a net exposure time after
 340 filtering of 49 ks, long enough to achieve our scientific
 341 goals. Information on the observation is shown in Ta-
 342 ble 1.

343 First, we extracted a spectrum of the entire area of the
 344 observation with a radius of 12.5 arcmin from the aim
 345 point for comparison with the *HaloSat* spectra. Fitting
 346 with Model0, we found the same type of residuals, e.g.
 347 significant residuals around Fe-L complex and H-like
 348 oxygen emission lines, as for the *HaloSat* spectra. The
 349 absorption column density was fixed to be $\sim 0.53 \times 10^{22}$
 350 cm^{-2} based on the Galactic HI survey (Kalberla et al.
 351 2005).

352 Thus, we fitted the *XMM-Newton* spectrum with
 353 Model1 and Model2 and found that the fit improves
 354 significantly from Model0 to Model1 and then from
 355 Model1 and to Model2 similarly to the *HaloSat* results.
 356 The spectrum fitted with Model2 is shown in Fig. 4.
 357 We confirm the presence of high temperature plasma
 358 with a temperature of 1.2 keV and emission measure of
 359 $12 \times 10^{-4} \text{ cm}^{-6} \text{ pc}$ in Model2. The observed temperature
 360 in Model2 is consistent with that for HS0023. The emis-
 361 sion measure in Model2 is also consistent, even though
 362 the FoV of *XMM-Newton* is approximately two orders
 363 of magnitude smaller than that of *HaloSat*.

364 To help understand the contribution of resolved point-
 365 like sources to the observed high temperature plasma, we
 366 investigated a stacked spectrum made of 197 point-like
 367 source spectra detected in the *XMM-Newton* observa-
 368 tion as reported in the catalog (e.g., Webb et al. 2020).
 369 Source and background spectra were found each point-
 370 like source in the catalog. The source regions are circles
 371 with radius of 20 arcseconds while the background re-
 372 gions are annuli with radii of 45 and 60 arcseconds, both
 373 are centered on the source coordinates tabulated in the
 374 catalog. All of the source spectra were summed, the
 375 same done for the background, and finally the stacked
 376 background spectrum was subtracted from the stacked
 377 point-like sources spectrum.

378 As our first step in analysis of the stacked spectrum,
 379 we confirmed that the spectrum was not represented
 380 well only by the non-thermal model corresponding to the
 381 CXB because of large residuals below ~ 1 keV including
 382 highly-ionized oxygen emission lines and the Fe-L com-
 383 plex were found, suggesting that an additional thermal
 384 plasma(s) is needed. We then added unabsorbed opti-
 385 cally thin thermal CIE plasma models until no signifi-

386 cant (99.9%) improvement in the goodness of fit was ob-
 387 served. The preferred model includes two plasma com-
 388 ponents and the CXB.

389 The spectrum with the best fit model and the fit-
 390 ting results are indicated in Figure 4 left and Table 3.
 391 The plasma temperatures are low (~ 0.3 keV) and high
 392 (~ 1 keV). The latter temperature is consistent with that
 393 found for the full observation area while the emission
 394 measure is ~ 50 %. This suggests that the UHTPGD
 395 partly originates from point-like sources. The spectral
 396 shape of the low-temperature plasma is similar to that
 397 of the MWH. Thus, the contribution of the plasma origi-
 398 nated from the point-like sources is included in the emis-
 399 sion measure of the MWH of the entire FoV spectral
 400 analysis. We confirmed the intensity of the CXB de-
 401 creases as expected considering the contribution of un-
 402 resolved point sources.

403 To further study the contribution from stars to
 404 UHTPGD as discussed in Masui et al. (2009), we used
 405 the Two Micron All Sky Survey (2MASS) Point Source
 406 Catalog (PSC) (Skrutskie et al. 2006) because most of
 407 the 2MASS point sources can be considered to be stars
 408 (Cambr esy et al. 2001). Approximately 60% of the X-
 409 ray point-like sources were cross-matched with sources
 410 tabulated in the 2MASS PSC with a matching radius
 411 of $5''$. The cross-matched sources have relatively high
 412 X-ray fluxes compared to the whole set of point-like X-
 413 ray sources in the field. For example, of the 20 brightest
 414 X-ray point-like sources tabulated in the catalog, 18 are
 415 cross matched in the 2MASS PSC. This may suggest
 416 that the observed UHTPGD is mainly due to the con-
 417 tribution of stars.

418 We extracted the stacked spectrum of X-ray point-
 419 like sources cross-matched with the 2MASS PSC and
 420 conducted spectral analysis with the same model as de-
 421 scribed above for the full set of X-ray point-like sources.
 422 There are no significant differences between the temper-
 423 ature and emission measure of the two plasmas in the
 424 stacked spectra of all versus the NIR-identified point-like
 425 sources (see Figure 3 bottom and Table 3), suggesting
 426 that the contribution from stars to the UHTPGD in the
 427 entire FoV of the pointed observation is significant.

428 5. SUMMARY AND CONCLUSIONS

429 The composition of the SXDB in the Galactic disk
 430 is not well known even though it is essential to our
 431 understanding of Galactic structure. The *ROSAT* all
 432 sky survey revealed excess emission in the R45 band
 433 (~ 0.44 – 1.0 keV) widely distributed in the Galactic disk.

434 Masui et al. (2009) showed that this excess has an en-
 435 ergy spectrum consistent with UHTPGD with a temper-
 436 ature of 0.8 keV in one field in the Galactic disk. Using
 437 *HaloSat*, we investigated the soft X-ray energy spectra
 438 of five fields in the outer Galactic disk ($135^\circ < l < 254^\circ$)
 439 selected to exclude bright X-ray sources. Our anal-
 440 ysis shows that all of the spectra require a thermal
 441 plasma component with temperature in the range of
 442 0.8 – 1.0 keV. Thus, for the first time, we present evi-
 443 dence that UHTPGD is pervasive in the outer Galac-
 444 tic disk. The observed temperatures are not signifi-
 445 cantly different from the UHTPGD found previously in
 446 one field (Masui et al. 2009), while the emission mea-
 447 sure of the UHTPGD observed with *HaloSat* varies from
 448 $(8$ – $11) \times 10^{-4}$ cm $^{-6}$ pc and the value is different by a
 449 factor of ~ 1.3 . The spectrum of the total area of an
 450 *XMM-Newton* observation contained within one *HaloSat*
 451 field shows UHTPGD temperature and emission mea-
 452 sure consistent with found with *HaloSat*.

453 To determine whether the observed UHTPGD may
 454 originate from stars, we examined the same *XMM-*
 455 *Newton* observation and constructed stacked spectra of
 456 all X-ray point-like sources detected in the FoV and
 457 of the X-ray sources with infrared counterparts in the
 458 2MASS catalog which are predominantly stars. The
 459 stacked spectra are well fitted with a model including
 460 two thermal plasma components and the CXB. There
 461 are no significant differences between the fit results in
 462 the two thermal plasma components for the different
 463 stacked spectra. The temperature of the highest tem-
 464 perature plasma component is consistent with that ob-
 465 served in spectral analysis for the entire area while the
 466 emission measure is ~ 50 %. This suggests that the
 467 UHTPGD may partly originate from point-like sources
 468 such as stars.

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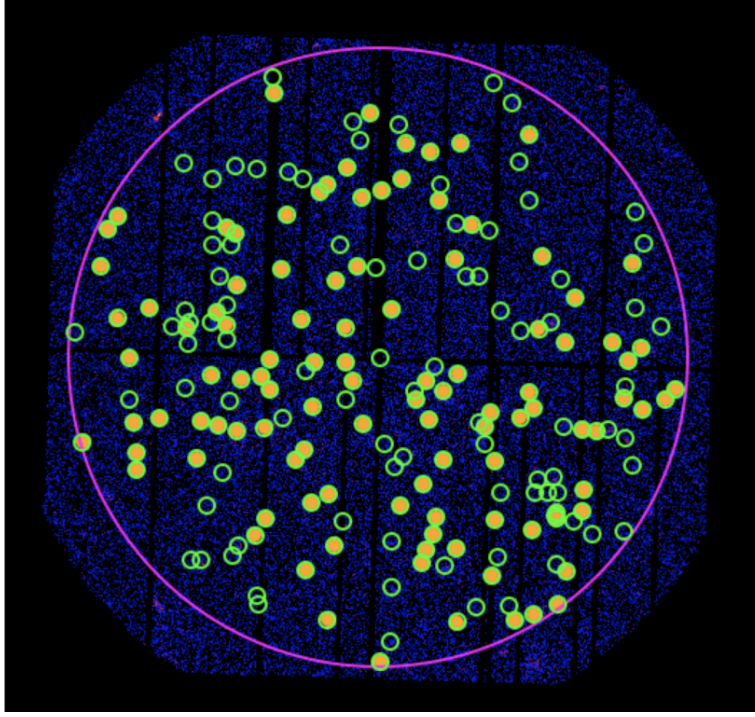


Figure 3. 0.4–5.0 keV *XMM-Newton* PN image with X-ray point-like sources (green circles). The cross-matched X-ray point-like sources with 2MASS (filled orange circle) is also shown. Spectral analysis was conducted within a circle with a radius of 12.5 arcmin (magenta).

Table 3. The results of the model fitting for a pointed *XMM-Newton* observation.

	kT_{low} [keV]	EM_{low}^* [$\times 10^{-4} \text{ cm}^{-6} \text{ pc}$]	kT_{high} [keV]	$\text{EM}_{\text{high}}^*$ [$\times 10^{-4} \text{ cm}^{-6} \text{ pc}$]	$\chi^2/\text{d.o.f}$
FoV			$1.2^{+0.1}_{-0.1}$	12^{+4}_{-3}	366/317
All point source	$0.28^{+0.03}_{-0.03}$	$5.8^{+0.8}_{-0.9}$	$0.99^{+0.08}_{-0.07}$	$5.8^{+0.8}_{-0.8}$	152/145
2MASS point source	$0.31^{+0.03}_{-0.03}$	$4.9^{+0.6}_{-0.6}$	$1.0^{+0.1}_{-0.1}$	$5.8^{+0.6}_{-0.6}$	130/87

* Emission Measure integrated over the line of sight, i.e. $\text{EM} = \int n_e n_{\text{H}} dl$.

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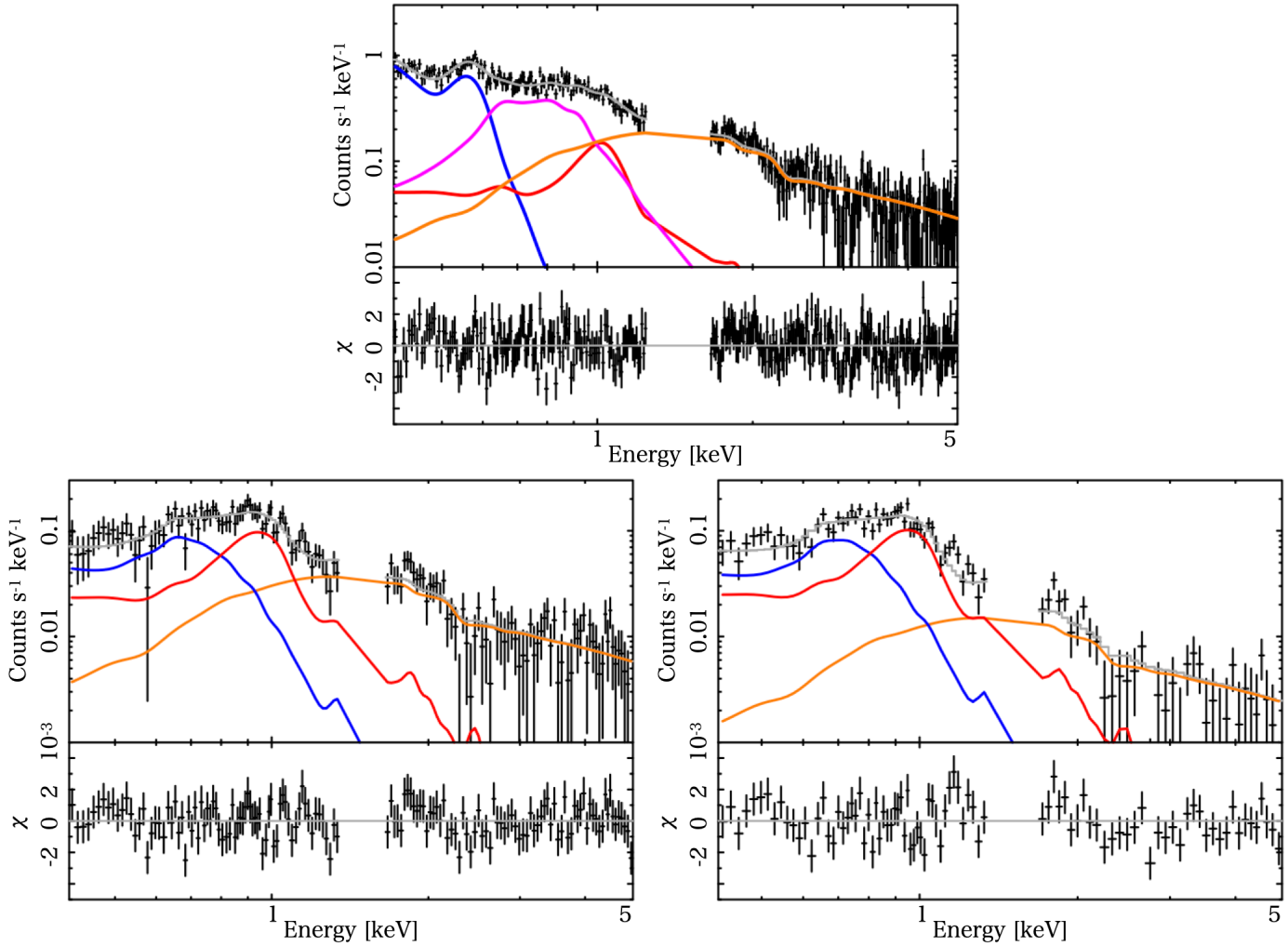


Figure 4. Spectrum of an *XMM-Newton* observation obtained from the entire area (top) and stacked spectra of all of the X-ray point-like sources (bottom left) and cross-matched X-ray point sources with 2MASS point sources (bottom right). Top: blue, magenta, red, and orange lines correspond to the SWCX+LHB, MWH, UHTPGD and the CXB, respectively. Bottom: blue, red, and orange lines show the low- and high-temperature plasmas and the CXB, respectively.