

1 Detection of extensive equatorial plasma depletions after the 2022 Tongan volcanic
2 eruption from multiple geodetic satellite ranging systems

3

4 Shin-Chan Han (1)*, Hyosub Kil (2), Richard Ray (3), Frank Lemoine (3), Colin Waters (4)

5

6 (1) School of Engineering, University of Newcastle, Callaghan, NSW, Australia

7 (2) Applied Physics Laboratory, The Johns Hopkins University, Maryland, USA

8 (3) Geodesy and Geophysics Laboratory, NASA Goddard Space Flight Center, Maryland, USA

9 (4) School of Information and Physical Sciences, University of Newcastle, Callaghan, NSW,
10 Australia

11

12 *Correspondent: shin-chan.han@newcastle.edu.au

13

14 Submitted to

15 JGR Space Physics

16 27MAR24

17 Revised 22OCT24

18

19

20 **Abstract**

21 We present a number of unique observations of ionospheric anomalies following the
22 Hunga-Tonga Hunga-Ha'apai (HTHH) volcanic eruption on 15 January 2022. All are based on
23 non-dedicated geodetic satellite systems: Global Positioning System tracking of Low Earth
24 Orbit (LEO) CubeSats, intersatellite tracking between two GRACE Follow-On satellites,
25 satellite radar altimeters to the ocean surface, and Doppler radio beacons from ground
26 stations to LEO geodetic satellites. Their observations revealed the development of
27 anomalously large trough-like plasma depletions, along with plasma bubbles, in the
28 equatorial regions of the Pacific and East Asian sectors. Trough-like plasma depletions
29 appeared to be confined within approximately $\pm 20^\circ$ magnetic latitude, accompanied by
30 density enhancements just outside this latitude range. These plasma depletions and
31 enhancements were aligned with the magnetic equator and occurred across broad
32 longitudes. They were detected in regions where atmospheric waves from the HTHH
33 eruption passed through around the time of the sunset terminator. We interpret these
34 phenomena in terms of the E dynamo electric fields driven by atmospheric waves from the
35 eruption. The uplift of the ionosphere beyond satellite altitudes, followed by subsequent
36 plasma diffusion to higher latitudes along magnetic field lines, results in the formation of
37 trough-like plasma depletions around the magnetic equator and density enhancement at
38 higher latitudes. The detection of plasma bubbles in the Asian sector during the non-bubble
39 season (January) is likely associated with the uplift of the ionosphere at the sunset
40 terminator.

41

42 **Plain Language Summary**

43 Satellite-based ranging measurements use radio signals, which are impacted by electrons in
44 the Earth's atmosphere. A correction for electron density must be applied to the ranging
45 measurement prior to its scientific use. We have found that, after the Tonga volcanic
46 eruption on 15 January 2022, the necessary corrections in the equatorial region were much
47 smaller than usual for four geodetic satellites systems including GPS, GRACE Follow-On,
48 satellite altimeters, and DORIS radio beacons. This implies a reduction in electron content
49 in the Earth's ionosphere in the equatorial region. Our observations demonstrate that the
50 spatial and temporal patterns of the reduction are associated with the atmospheric

51 pressure waves generated by the eruption. We suggest electric fields driven by atmospheric
52 waves from the eruption as the source of equatorial electron density depletions.

53

54 **Key points**

55 - Geodetic satellites detected plasma depletions and enhancements at low latitudes after
56 the Tongan volcanic eruption on 15 January 2022.

57 - These phenomena were observed within the region affected by the atmospheric waves
58 propagated from the eruption.

59 - We suggest that E-region dynamo electric fields are a plausible source of the plasma
60 density modulation at low latitudes.

61

62

63 1. Introduction

64

65 The Earth's ionosphere is the region of the upper atmosphere (approximately 80 – 1,000 km
66 altitude) where extreme ultraviolet and x-ray solar radiation produces a significant amount
67 of ions and electrons. The major constituent of the ionosphere above ~200 km altitude is
68 the oxygen ion (O^+). Below that, molecular constituents like nitro oxygen (NO^+) and
69 molecular oxygen ions (O_2^+) dominate. Therefore, the plasma density above 200 km
70 altitude is primarily determined by production, loss, and transportation of O^+ . The loss rate
71 of O^+ is proportional to the number density of nitrogen and oxygen molecules which varies
72 with altitude. Normally, the peak electron density appears at altitudes around 300 km, the
73 result of a tradeoff between solar radiation energy increasing with altitude and atomic
74 oxygen density decreasing with altitude. The maximum electron density varies generally in
75 the range of $10^5 - 10^6$ electrons/cm³ depending on various factors including local time,
76 latitude, longitude, season, and solar cycle (Kelley, 2009). At low and mid latitudes, the
77 ionospheric height is modulated by electric fields and plasma motion along magnetic field
78 lines driven by neutral winds. The modulation of the ionospheric height also affects the
79 latitudinal plasma distribution and the development of plasma instabilities (Hanson and
80 Moffett, 1966; Kelley, 2009; Kil and Lee, 2013).

81

82 Various forms of atmospheric, ionospheric, and oceanic perturbations have been detected
83 after the Hunga Tonga-Hunga Ha'apai (HTHH) volcanic eruption, which occurred on 15
84 January 2022 shortly after 4:00 UTC (e.g., Aa et al., 2022; Matoza et al., 2022; Themens et al.,
85 2022; Wright et al., 2022). These studies reported the propagation of atmospheric waves
86 with different speeds, the generation of traveling ionospheric disturbances (TIDs), the
87 modulation of ionospheric height, and the development of unexpected plasma instabilities.
88 Initially large-scale TIDs (LSTIDs) with 2-3 times the surface acoustic speed were generated
89 by acoustic shock waves from the eruption and followed by their slowing at distant locations,
90 and subsequently medium-scale TIDs (MSTIDs) propagating at the speed of sound were
91 observed for many hours (e.g., Themens et al., 2022; Zhang et al., 2022).

92

93 Two fascinating ionospheric phenomena caused by the volcanic eruption are MSTIDs that
94 propagate with the speed of atmospheric Lamb waves (300-350 m/s) and severe plasma

95 depletions at low latitudes, measured by ground-based Global Navigation Satellite System
96 (GNSS) station networks and space-weather dedicated satellite missions (Aa et al., 2022;
97 Amores et al., 2022; Harding et al., 2022; Zhang et al., 2022). MSTIDs in the dayside are
98 produced by the direct interaction between plasma and neutral particles along the path of
99 atmospheric waves. Ionospheric perturbations are intensified by electric fields at night,
100 producing stronger MSTIDs. One of the characteristics of MSTIDs at night is the formation
101 of their mirror (or conjugate) structures at magnetic conjugate locations, and this
102 characteristic was identified after the HTHH volcanic eruption (Hong et al., 2022; Lin et al.,
103 2022). Total electron content (TEC) anomalies as large as 10 TECU (1 TECU = 10^{16}
104 electrons/m²) travelling with the Lamb waves were also detected at higher altitudes (>550
105 km) from the constellation of Global Positioning System (GPS)-tracked CubeSats (Han et al.,
106 2022 & 2023). Recent modeling studies by Vadas et al. (2023a & 2023b) also suggested TIDs
107 could have been generated by secondary atmospheric gravity waves in contrast to the
108 leakage of the surface Lamb wave.

109

110 The formation of plasma depletions at low latitudes is closely linked to the uplift of the
111 background ionosphere by electric fields and the generation of plasma bubbles (Kil and
112 Paxton, 2006; Lee et al., 2014; Mannucci et al., 2005). Ground events such as volcanic
113 eruptions may launch atmospheric waves that propagate radially away from the epicenter.
114 These waves generate dynamo electric fields in the ionosphere by which the ionospheric
115 height can be lifted or lowered depending on the wind direction. Multiple studies reported
116 the occurrence of the uplift of the ionosphere during the HTHH event (Gasque et al., 2022;
117 Harding et al., 2022; Hong et al., 2022; Huba et al., 2023; Rajesh et al., 2022). The uplift of
118 the ionosphere in the equatorial region results in the transport of equatorial plasma to
119 higher latitudes by the fountain process (Hanson & Moffett, 1966). This process produces
120 an ionization trough near the magnetic equator and ionization crests outside the equatorial
121 region. The uplift of the ionosphere also promotes the development of equatorial
122 ionospheric phenomena known as “plasma bubbles”. Plasma bubbles, characterized by
123 plasma depletions with respect to the background, are produced by the transport of low-
124 density plasma on the bottomside ionosphere to the topside by the generalized Rayleigh-
125 Taylor instability (Huba et al., 2008; Kelley, 2009; Sultan, 1996; Yokoyama et al., 2014).

126 These two processes (uplift and plasma instability) are considered to be the major sources
127 of the plasma depletions at low latitudes (Kil and Lee, 2013; Lee et al., 2014).

128

129 Geodetic satellite ranging systems are sensitive to electromagnetic media. A medium like
130 the ionosphere causes dispersion of electromagnetic waves by altering direction and speed
131 of propagation. The amount of alteration, often quantified as refractivity, depends on the
132 frequency of electromagnetic waves propagating into the electrically charged ionosphere.
133 For precise geodetic ranging, raw ranging measurements (like beat phase or travel time)
134 must be corrected for refraction by the ionosphere (as well as by the troposphere if
135 relevant). Otherwise, ionospheric refraction may introduce tens of meters of systematic
136 delay in, for example, ranging based on GNSS (Seeber, 2003).

137

138 The frequency dependence of refractivity facilitates disentangling ranging and ionospheric
139 interference from measurements, which can be achieved by exploiting multi-frequency
140 waves. This is the fundamental reason for using multi-frequency L-band signals (notably L1
141 and L2) in GNSS. Similarly, the Gravity Recovery And Climate Experiment Follow-On
142 (GRACE-FO) satellites measure inter-satellite distances with Ka- and K-band waves, and this
143 dual frequency is used to remove ionospheric effects on precise phase measurements in the
144 GRACE-FO microwave interferometer (Tapley et al., 2004; Lee et al., 2011). Satellite radar
145 altimeters, such as those onboard the Sentinel-3 satellites, measure the distance between
146 the satellite antenna phase center and the surface of the Earth, and ionospheric refraction is
147 eliminated by using dual-frequency (e.g., Ku- and C-band) radar (Chelton et al., 2001). For
148 the same reason, Doppler Orbitography and Radiopositioning Integrated by Satellite (DORIS)
149 system uses S- and VHF-band microwave signals (Fleury et al., 1991; Willis et al., 2010; Liu et
150 al. 2023).

151

152 This study aims to identify the longitudinal extent of trough-like equatorial plasma
153 depletions caused by the HTHH volcanic eruption. Ground-based and space-weather
154 dedicated satellite observations revealed the development of plasma depletions (Aa et al.,
155 2022; Hong et al., 2022), but these observations have large data gaps over the sea. Given
156 the size of the TEC anomalies reported, other geodetic satellite ranging measurements are
157 likely affected by the HTHH eruption. We first review how each of four geodetic techniques,

158 including GNSS on Low Earth Orbiters, GRACE-FO satellites, satellite altimeters, and DORIS-
 159 equipped satellites, is used to derive an ionospheric correction that is related to the amount
 160 of electrons integrated over different paths of the radio frequency (RF) wave propagation.
 161 Then, we examine the development of large trough-like plasma depletions over extensive
 162 regions using the observations from all geodetic techniques. Finally, we discuss the
 163 generation mechanism of depletions in association with the global atmospheric wave (Lamb
 164 wave) propagation triggered by the HTHH volcanic eruption.

165

166 **2. Geodetic satellite ranging techniques of measuring TEC**

167

168 *2.1 High-low GNSS ranging to low Earth orbit satellites*

169

170 A radio wave refracts in the ionosphere if the frequency is higher than the plasma frequency.
 171 Ignoring the magnetic field, the refractive index (n) of plasma in the ionosphere is
 172 essentially dependent on the plasma frequency (f_p) and the frequency (f) of the radio wave
 173 through $n = \sqrt{1 - f_p^2/f^2}$. As the plasma frequency increases with electron density, the
 174 amount of refraction increases with the electron density and decreases with the square of
 175 radio frequency (e.g., Seeber, 2003; Kelley, 2009). Thus, multi-frequency GNSS radio
 176 measurements can be used to derive the electron density integrated along the signal
 177 propagation (roughly over 20,000 km, although most electrons are concentrated within
 178 1,000 km of the surface).

179

180 By forming a linear combination of L1 and L2 carrier phase counts, $\Phi_{A,1}^i$ and $\Phi_{A,2}^i$, also
 181 known as geometry-free combination, the TEC between GNSS satellite i and receiver A is

182

$$TEC_A^i(t) = \left(\frac{c}{f_{L1}} \Phi_{A,1}^i(t) - \frac{c}{f_{L2}} \Phi_{A,2}^i(t) \right) / \left(\frac{40.3}{f_{L2}^2} - \frac{40.3}{f_{L1}^2} \right) - B_A^i, \#(1)$$

183

184 where c is the speed of light in vacuum, and f_{L1} and f_{L2} are the frequencies of L1 and L2
 185 radio waves, respectively. The approximate constant of 40.3 is derived from values of the
 186 electron charge and mass as well as permittivity that appears in the plasma frequency f_p

187 (Hartman and Leitinger, 1984; Teunissen and Kleusberg, 1998). The numerator is the
188 difference of the (biased) ranges measured with L1 and L2 carrier phases which is due to the
189 ionospheric refraction being dependent on radio frequency. This TEC is also called ‘slant’
190 TEC along the ray path, and must be differentiated from vertical TEC. Equation (1) implies
191 approximately 16 cm path group delay (or phase advance) per 1 TECU (10^{16} electrons/m²)
192 for L1 frequency. It is also noted that the TEC computed with two phase measurements as
193 in equation (1) contains a bias, B_A^i , that is a composite of integer ambiguities, initial
194 fractional phases and differential channel biases (Leick, 1995). In this study, we do not
195 estimate the bias term from dual frequency GNSS data and only focus on relative TEC
196 variations.

197

198 *2.2 Low-low intersatellite ranging between two GRACE Follow-On satellites*

199

200 The GRACE and GRACE-FO microwave interferometers (MWI) measure relative distance
201 changes between two satellites in tandem by exploiting multiple carrier phases generated
202 by two ultra-stable oscillators (USO) and transmitted simultaneously from two transceivers
203 at both satellites. This is called dual one-way ranging (DOWR) and is designed to eliminate
204 low frequency drift errors of USOs in two satellites without explicit modeling of oscillator
205 (clock) error (Thomas, 1999). In addition, the MWI uses two K- and Ka-band radio
206 frequencies to eliminate the first-order effect of ionospheric phase advance at the satellite
207 altitude (Case et al., 2004; Wen et al., 2017). The second-order effect could be more than
208 three orders of magnitude smaller (e.g., Kedar et al., 2003), which may not be insignificant
209 considering the observation of the first-order effect being as large as several millimeters and
210 the precision of MWI ranging being only several microns (e.g., Landerder et al., 2022).
211 However, the second-order effect depends on the signal propagation direction relative to
212 the local magnetic field and, thus, the DOWR combining carrier phases propagating from
213 both directions can nullify greatly the second-order effect. This is an advantage of DOWR
214 although not widely recognized within the community except in a recent paper by Müller et
215 al. (2022).

216

217 Here, we examine how to derive TEC variations between two GRACE-FO satellites from the
218 dual-frequency DOWR MWI by considering only the first order effect (as all higher order

219 effects are negligible). This TEC represents the path-integrated electron density over ~220
 220 km in distance between the two satellites at ~480 km altitude. The biased intersatellite
 221 ranging measurements, R_K and R_{Ka} , can be obtained from the DOWR phases, respectively,
 222 for the K and Ka-bands, as follows (Wu et al., 2004):

$$R_K(t) = \frac{c}{f_K^A + f_K^B} \left(\Phi_K^{A,B}(t) + \Phi_K^{B,A}(t) \right), \#(2)$$

$$R_{Ka}(t) = \frac{c}{f_{Ka}^A + f_{Ka}^B} \left(\Phi_{Ka}^{A,B}(t) + \Phi_{Ka}^{B,A}(t) \right), \#(3)$$

224
 225 where c is the speed of light in vacuum, f_K and f_{Ka} are the K- and Ka-band frequencies,
 226 respectively, separately operated at USOs of satellites A and B (as indicated by the super-
 227 script), and $\Phi_{K\text{ or }Ka}^{A,B}$ and $\Phi_{K\text{ or }Ka}^{B,A}$ represent the one-way phase from satellite A to B and
 228 from B to A, respectively; and these phase measurements are available for both K- and Ka-
 229 bands as indicated in the sub-script.

230
 231 The two (biased) intersatellite ranging measurements, R_K and R_{Ka} , are subject to two
 232 different ionospheric effects proportional to the amount of electrons cumulative along the
 233 ray path (i.e., TEC) between two satellites and to the inverse of frequency squared of the
 234 carrier phases. This leads to ionosphere-free combined intersatellite ranging (with unknown
 235 phase biases) such as

$$R_{KBR}(t) = C_{Ka}R_{Ka}(t) - C_KR_K(t), \#(4)$$

237
 238 where R_{KBR} is the biased range and two constants C_{Ka} and C_K are derived from the K- and
 239 Ka-band frequencies and result exactly in $C_{Ka} = 16/7$ and $C_K = 9/7$ within high precision
 240 (see Case et al., 2004). This R_{KBR} is further corrected for light time and antenna phase
 241 center to derive instantaneous (biased) range that is the fundamental data for the Earth's
 242 gravitational analysis (e.g., Kim, 2000).

243
 244 On the other hand, the (biased) ionospheric path delay (or phase advance) at the Ka-band is
 245 simply computed from the difference between R_{KBR} and R_{Ka} , that is, $R_{KBR} - R_{Ka} =$

246 $C_K(R_{Ka} - R_K)$, and it is also available as one of the Level-1B data products (i.e., “iono_corr”
247 within KBR1B). From this ionospheric range effect, we compute the time series of TEC
248 between two GRACE-FO satellites as

249

$$TEC_{GRFO}(t) = C_K\{R_{Ka}(t) - R_K(t)\}/(40.3/f_{Ka}^2). \#(5)$$

250

251 Equation (5) implies approximately 0.4 mm path delay per 1 TECU for Ka band frequency.
252 Using the KBR1B data record, we examine short-baseline (~220 km) TEC anomalies at the
253 satellite altitude with regard to the volcanic eruption event. This measurement is local and
254 differentiated from GNSS TEC measurement over long baselines between LEO and GNSS
255 satellites. Similar to the GNSS TEC, it is a biased TEC measurements as the K- and Ka-band
256 phase ambiguities are not known.

257

258 *2.3 Nadir satellite altimeter ranging to sea surface*

259

260 Our third geodetic technique for recovering TEC is dual-frequency satellite altimetry. This is
261 a much noisier technique compared with the two ranging types discussed above because it
262 entails ranging off the ocean surface. However, it does allow determination of the absolute
263 TEC (not relative) for the whole ionosphere below the satellite. Unlike the previous phase-
264 based ranging techniques that result in ionospheric phase ‘advance’ and thus measured
265 range shorter than true range, the altimeter is based on radar pulses that are subject to
266 group refractivity and leads to range ‘delay’ in measurement. Other than the sign of the
267 ionospheric effect on range measurement, ionospheric phase advance and group delay can
268 be determined using dual-frequency measurements. Although there were other dual-
269 frequency altimeters flying during January 2022 (e.g., Jason-3), here we use data from the
270 Sentinel-3 satellites as they are sun-synchronous (10:00 and 22:00 local time) and thus
271 repeated passes tend to be affected by (somewhat) similar diurnal variability, thereby
272 simplifying comparisons before and after the HTHH event.

273

274 In January 2022 there were two Sentinel-3 satellites in operation, A and B, flying at an
275 altitude of approximately 815 km. Each satellite has two altimeters: a primary one (Ku band)
276 at 13.6 GHz and a secondary (C band) at 5.41 GHz (this C-band frequency is slightly below

277 the comparable frequency used on the Jason satellite altimeters). The altimeter range
 278 perturbation caused by ionospheric path delay at frequency f is related to the TEC below
 279 the satellite according to (Chelton et al., 2001; Eq. 85):

280

$$\Delta R_{\text{ion}}(f) = \frac{40.3}{f^2} \text{TEC}, \#(6)$$

281

282 in meters (and TEC in electrons/m²), which evaluates to approximately 0.218 cm per 10¹²
 283 electrons/cm² for the K-band frequency, or ~2.18 mm per TECU.

284

285 Let f_k and f_c be the two altimeter frequencies, and let their ratio be denoted

286

$$\delta = f_c/f_k = 0.3985.$$

287

288 The range perturbation ΔR_{ion} at f_k can be determined from the difference between the two
 289 measured ranges $R(f_k)$ and $R(f_c)$ (Chelton et al., 2001; Eq. 96):

290

$$\Delta R_{\text{ion}}(f_k) = a_c [R(f_c) - R(f_k)], \#(7)$$

291

292 where

293

$$a_c = \delta^2/(1 - \delta^2) = 0.189.$$

294

295 A potentially serious complication arises from the fact that the altimeter range is affected by
 296 the ocean sea state, a bias error which itself is frequency dependent. This “sea-state bias” is
 297 caused primarily by wave troughs being better reflectors than wave peaks, so the mean
 298 reflected power is biased toward troughs (e.g., Walsh et al., 1991; Arnold et al., 1995). Sea-
 299 state bias corrections are generally determined separately for each altimeter frequency by
 300 regressing altimeter-based sea-surface heights against wave height (and often wind speed
 301 as well), after an initial ionospheric correction (e.g., Tran et al., 2021). Here we use non-
 302 parametric corrections developed by Tran and her colleagues and supplied by the Sentinel-3

303 project. The final TEC determination then uses (7) after the two ranges, $R(f_k)$ and $R(f_c)$,
304 have been corrected for sea-state bias.

305

306 As noted above, the TEC determined on individual satellite passes is fairly noisy. The
307 standard approach when applying ionospheric corrections to altimeter-based ocean
308 elevations is to smooth raw (typically 1-Hz) ionospheric data along-track over distances of
309 100-150 km (Zlotnicki, 1994). Below we also apply along-track smoothing, but we also show
310 the unsmoothed 1-Hz Sentinel-3 TEC profiles so readers can readily observe the inherent
311 noise levels in these data.

312

313 *2.4 DORIS ground beacons*

314

315 The last geodetic ranging technique of interest is Doppler Orbitography and
316 Radiopositioning Integrated by Satellite (DORIS) (e.g., Willis et al., 2010). DORIS utilizes two
317 RF of 401.25 MHz (VHF) and 2.036 GHz (S-band) continuously emitted from a network of
318 nearly 60 global ground stations (Saunier, 2023). Satellites equipped with DORIS receivers
319 record dual-frequency Doppler shifts that can be used to compute line-of-sight (LOS) range-
320 rate and, thus, changes in the TEC between the ground stations and the satellites (e.g.,
321 Fleury et al., 1991; Bernhardt et al., 2006; Dettmering et al., 2014). Such measured
322 ionospheric group delay or phase advance is corrected for station coordinates and precise
323 orbit determination with DORIS tracking data in real-time or ex post facto for the creation of
324 geophysical products (e.g. Jayles et al., 2015; Lemoine et al., 2016). The ionosphere
325 correction is derived from the DORIS/RINEX tracking data provided by the International
326 DORIS Service (IDS) (Willis et al., 2010). We follow the procedures described by Lemoine et
327 al. (2016) to compute the ionosphere correction for Doppler data. Currently, a number of
328 scientific satellites are equipped with DORIS receivers including Jason-3, Sentinel-3A/3B and
329 Sentinel-6A. The DORIS data cover the layers of the ionosphere up to satellite altitudes as
330 high as 1,300 km, but, in the LOS direction only. This is a sampling geometry that is different
331 from the geodetic methods described earlier, and as such leads to complementary
332 information.

333

334 We used the DORIS/RINEX-derived ionosphere correction given in range-rate every 10
335 seconds and integrated over time to determine the relative (biased) TEC change in LOS
336 between the ground stations and the satellites. Typically, 15 minutes of data were available
337 from each station around the equatorial region between a satellite's rise and set, covering
338 roughly 60 degrees or less in latitude.

339

340 **3. Observational results**

341

342 *3.1 Measurements of plasma depletion by GNSS, GRACE Follow-On, and radar altimeter* 343 *satellites*

344

345 For GNSS TEC variation, we processed GPS L1 and L2 carrier phase measurements every
346 second, as available from multiple patch antennae of the Spire's 3U CubeSats (Nguyen et al.,
347 2020; Han et al., 2023): (1) from the antenna in the zenith direction of the CubeSats, (2)
348 from the antenna pointing in the velocity direction, and (3) from the antenna pointing
349 opposite to the velocity direction. The first antenna receives signals from GPS satellites
350 above the local horizon of the CubeSats (and they are used typically to determine precise
351 positions of the CubeSats). The last two are designed to measure the signals from GPS
352 satellites at and below the local horizon of the CubeSats, with the local elevation range from
353 0° down to -21° . The lower bound of the elevation is dependent on the LEO satellite
354 altitude. These GPS data from two side-looking antennae were used to examine line-of-
355 sight (slant) TEC variations between the occulting CubeSats and GPS satellites covering
356 altitudes from the bottom of the ionosphere (~ 80 km) to the CubeSat altitude (~ 550 km), in
357 contrast to the high-altitude TEC variation above the CubeSats measured by the zenith-point
358 antenna. The TEC measurements below the local horizon were also converted to electron
359 density profiles in altitude assuming spherically stratified distribution of electrons over the
360 regions of the impact points (where the maximum refractions occur) using the radio
361 occultation geometry and Abel inversion (Hajj et al., 2002).

362

363 Figure 1 shows four different cases of GPS and CubeSats, and electron density profiles on
364 three consecutive days of 14, 15 and 16 of January 2022. For each case, the left panel
365 shows the ground tracks of the CubeSat and GPS satellite and location of the HTHH volcano.

366 The Lamb wave front at the average time of GPS measurements (UTC shown on top of each
367 panel) is also shown with large curves of different colors representing a few hPa variations
368 of surface pressure (Amores et al., 2022). In each case, there are ~15 minutes of GPS data
369 measured from the occulting CubeSat while it crosses over the rim until it loses tracking
370 from the GPS satellite. Figure 1 also illustrates the average location of the occultation
371 tangent points (marked with stars in colors next to CubeSat ground tracks) where the
372 maximum bending of the GPS radio ray takes place (Haji et al., 2002). The right panel of
373 each case shows the electron density profiles around the tangent points estimated from
374 GPS slant TEC measurements. Note that the GPS electron density is a relative measurement.
375 In these profiles, the minimum value of GPS estimates was matched to that of the
376 International Reference Ionosphere (IRI) model values (Blitza, 2018).

377

378 In the first case (Figure 1a), GPS measurements took place between the occulting CubeSat
379 FM143 and GPS G04 at UTC 11:24 and local time (LT) 01:24 on 15 and 16 of January 2022.
380 The CubeSat sampled the electron density changes over the area northeast of the volcano
381 near the geomagnetic equator right after the eastward Lamb wave passed the region on
382 Day 15. The GPS data from 16 January identified a typical electron density profile with rapid
383 increase of electron density from 200 to 300 km followed by gradual decrease at higher
384 altitudes, showing the peak layer formed around 280 km. However, the GPS measurements
385 on Day 15 do not show any substantial change in altitude over the entire F region and thus
386 no peak formed likely due to electrons depleted from the F region. There is no GPS data on
387 14 of January from the same CubeSat at the same location.

388

389 For the second case (Figure 1b), CubeSat FM117 measured the electron density over the
390 area northwest of the volcano near the geomagnetic equator at UTC noon (11:48) and LT
391 midnight (23:48) on three different days. Once again, the measurements were taken after
392 the Lamb wave swept westward over the region on Day 15. The GPS measurements have
393 peak values of 2 and 4.5×10^5 electrons/cm³ on Day 14 and 16, respectively, at an altitude
394 of 350 km. However, there was no evident electron density peak formed on Day 15 and
395 there is little electron density variation in altitude from 80 to 550 km relative to variation on
396 other days, again indicating possible depletion of electrons from the entire F region.

397

398 The third case of GPS TEC measurements (Figure 1c) occurred above the South China Sea by
399 CubeSat FM115 and GPS G08 at UTC 14:00 (LT 22:00) on three different days. The
400 westward propagating Lamb wave passed through the region before the TEC measurements
401 were taken on Day 15. The electron density increased drastically from 200 km to 270 km by
402 as much as $6-7 \times 10^5$ electrons/cm³ and gradually decreased at higher altitudes on Day 14
403 and 16. However, the electron densities on Day 15 do not change much across 100 – 550
404 km only varying by 1×10^5 electrons/cm³, possibly due to electron depletion by 6-7 times
405 the background density within the entire F region (>200 km) compared to other days.

406

407 The last example (Figure 1d) is from CubeSat FM119 occulting over west of Peru close to the
408 geomagnetic equator at a similar UTC 14:12 but at local time of early morning around 08:12
409 and east of the volcano. The region was just behind the eastward propagating Lamb wave.
410 The GPS electron density estimates show similar patterns of vertical electron distribution
411 but the peak value of 8×10^5 electrons/cm³ and the peak altitudes of 290 and 370 km on
412 both Day 15 and 16 (note that no CubeSat GPS data are available at this location and this
413 time on Day 14). Unlike the previous cases, no depletion was measured on Day 15,
414 although the observation was made after the Lamb waves swept the region.

415

416 For all examined cases, the IRI model (not shown in Figure 1), that relies solely on solar and
417 geomagnetic indices, was inefficient at capturing any plasma anomaly that originated
418 from the volcanic eruption and its associated atmospheric perturbation. The electron
419 depletion was observed from both west and east sides of the volcano where the
420 atmospheric wave propagated in two opposite directions and from the GPS measurements
421 during the local nighttime on Day 15. The last case of Figure 1 shows no depletion, which is
422 possibly due to the measurements occurring during local morning when plasma begins to
423 build up with increasing sun light. Possible dynamo mechanisms of depleting plasma in the
424 F region are discussed in Section 4.

425

426 We now examine the electron distribution at satellite altitude of 480 km across all latitudes,
427 as measured by the GRACE-FO K/Ka-band ranging system. During January 2022, the angle
428 between the GRACE-FO satellite orbital plane and the geocentric vector to the sun (known
429 as ‘beta angle’) was ~45 degrees. This puts the ascending tracks of the GRACE-FO satellites

430 at a local time of approximately 19:55-20:40 (a half orbital period). This orbital
431 configuration allowed them to sample electron density perturbations while the atmospheric
432 Lamb wave propagated over the Pacific after the volcanic eruption on Day 15. During the
433 local nighttime, electron density perturbations can persist due to the absence of solar
434 photoionization.

435

436 We processed the GRACE-FO intersatellite ranging data for TEC on 14, 15 and 16 of January
437 2022. The GRACE-FO ground tracks shifted eastward roughly by 5 degrees of longitude each
438 day. The measurements over the same local time on three consecutive days is ideal for
439 detecting the anomaly induced by the volcanic eruption. As the Earth rotation rate is ~ 460
440 m/s at the equator and the Lamb wave speed is 300-350 m/s, the GRACE-FO satellites were
441 able to measure the electron density over the following regions: (1) east of the Lamb wave
442 before the wave arrives (i.e., unperturbed state), (2) inside the regime of the Lamb wave
443 perturbation, (3) west of the Lamb wave before the wave arrives (i.e., unperturbed state);
444 all at the same local time of 19:55-20:40.

445

446 Figure 2 shows six cases of GRACE-FO measurements with the left panel of each case
447 showing the ground tracks of the GRACE-FO ascending orbit on three consecutive days
448 along with the Lamb wave fronts (black dashed curves) and the right panel showing the
449 measurements of TEC changes integrated over ~ 220 km baseline between two GRACE-FO
450 satellites at different average UTC but at almost the same average LT of 20:12 – 20:24. As
451 the GRACE-FO TEC measurements are biased (due to K/Ka-band carrier phase biases), the
452 minimum value of TEC measurements on each track was set equal to that of the IRI model
453 prediction. Note that any remaining bias in the data (not corrected by the IRI model) does
454 not affect our analysis of reporting relative changes over the repeat passes across the entire
455 latitudes.

456

457 In the first case of UTC 06:24 (Figure 2a), the GRACE-FO sampled the region outside and east
458 of the Lamb wave while the Lamb wave propagated outward from the HTHH volcano, thus
459 sampling the unperturbed ionosphere. The GRACE-FO TEC measurements do not exhibit
460 any abnormality on 15 of January relative to the measurements on the other days. There

461 are two orders of magnitude variation of TEC across latitude, showing the maximum of
462 several TECU around the geomagnetic equator.

463

464 The GRACE-FO ground tracks shifted westward in the next orbits while the Lamb wave
465 propagated outward from the HTHH volcano and the GRACE-FO satellites sampled the
466 region of atmospheric perturbation at longitudes of 180°, 160°, and 140° E, at UTC 08:00,
467 09:36, and 11:12, respectively (Figures 2b-d). The measurements on 15 January found 1,000
468 times (or even more) reduction in TEC relative to the TEC on the other days. The TEC value
469 of 0.01 TECU over the length of 220 km between two GRACE-FO satellites corresponds to
470 450 electrons/cm³, which is considered to be lower than the typical plasma density at the
471 bottom of the ionosphere (Kelley, 2009). The interpretation of plasma depletions in the
472 equatorial region is not straightforward because multiple processes may contribute to their
473 formation. Irregular density structures are interpreted as signatures of plasma bubbles;
474 however, density depletions approximately 3,000 km wide across the geomagnetic equator
475 are difficult to explain solely by plasma bubbles. The electron density enhancements at
476 18°N and 18°S in geomagnetic latitudes, along with the plasma depletions in the equatorial
477 region, can be understood as a consequence of the fountain effect (Hanson & Moffett,
478 1966). The deep depletions around the magnetic equator are likely caused by the uplift of
479 the F region beyond satellite altitudes.

480

481 The last example of the GRACE-FO measurements took place outside the Lamb wave
482 propagation (Figures 2e and 2f), but west of the Lamb wave (the opposite to the case of
483 Figure 2a). As the Earth rotation was 50% faster than the Lamb wave propagation, the
484 GRACE-FO satellites overtook the Lamb waves from west and east within a few orbital
485 revolutions. Figures 2e and 2f illustrate the measurements west of the Lamb wave front
486 capturing the ionosphere before the Lamb wave arrived. The GRACE-FO measurements
487 look normal, showing the peak electrons near the geomagnetic equator, for three
488 consecutive days just as the case of Figure 2a. No depletion was observed.

489

490 From the observations in Figure 2, we note plasma density enhancements outside the
491 equatorial region, accompanying severe plasma depletions around the magnetic equator
492 (Figures 2b–2d). These enhancements are pronounced in the Northern Hemisphere and are

493 interpreted as ionization crests. The intensities of the ionization crests in the opposite
494 hemispheres can differ due to variations in neutral composition and winds. The coincidence
495 of equatorial plasma depletions and plasma enhancements outside the magnetic equator is
496 also identified from other satellite observations shown below in Figures 3 and 4. As
497 discussed in Section 4, these observations provide insight into the source of equatorial
498 plasma depletions.

499

500 The vertical TEC integrated from the ocean surface to the altimeter satellite altitudes of
501 ~800 km was measured by the nadir-pointing altimeter ranging systems. The Sentinel-3A
502 satellite orbit is sun-synchronous with its ascending track crossing the equator at 22:00 local
503 time. This allowed sampling of the ionosphere in the Pacific during the local nighttime a few
504 hours after the HTHH volcanic eruption. The GRACE-FO satellites also happened to sample
505 the data at a local time of around 20:00. The vertical TEC measurements on three
506 consecutive days (14, 15 and 16 January 2022) are presented along with the ground tracks
507 in Figure 3; the original 1-sec measurements (blue dots) and the 60-sec moving averaged
508 data (magenta curve) are shown. By the times of the satellite passes over the Pacific, the
509 Lamb wave had propagated outward and reached up to 25°N and 45°N on 15 of January,
510 respectively, in Figures 3a and 3b.

511

512 The latitudinal samples of vertical TEC showed a wide region of low density around the
513 geomagnetic equator between 18°S and 18°N, only on Day 15 (Figure 3a). Two different
514 ionosphere models (JPL's Global Ionosphere Maps, GIM; Martire et al., 2024, and NOAA
515 Ionospheric Climatology, NIC09; Scharroo and Smith, 2010) predicted approximately 10
516 TECU around the equatorial region. The measurements from 14 and 16 January are much
517 more consistent with the models than the measurements from 15 January. The second
518 example of vertical TEC is from the subsequent orbit about 90 minutes later but at the same
519 local time (Figure 3b). As in the previous case, the reduction of TEC by as much as 10 TECU
520 is observed on 15 January only. It is also observed that TEC increased up to 30 – 40 TECU
521 around 20 degrees north and south of the geomagnetic equator at the same time of the
522 depletion observed around the geomagnetic equator, similar to the localized increase found
523 by the GRACE-FO observations. Two narrow peaks are more noticeable in Figure 3b. This
524 resembles the intensification of the equatorial ionization anomaly by the fountain effect

525 (Hanson & Moffett, 1966). Note that GIM is a model integrating TEC measurements from a
526 number of global GNSS stations, while NIC09 is a climatological analysis of GIM. Therefore,
527 the HTHH volcanic eruption effect is not expected from the climatological NIC09 model. It is
528 not clear why the GIM results (available through altimeter geophysical data record) do not
529 present the electron depletion and enhancement as much as seen from the altimeter
530 measurements. This is possibly related to the limited spatial resolution and lack of GNSS
531 stations over ocean regions (Martire et al., 2024), which is also apparent in altimeter/GNSS
532 comparisons of the diurnal cycle (Ray, 2020).

533

534 We also examined data from Sentinel-3B which was in the same orbital plane, for the five
535 consecutive orbits on Day 15, along with the Lamb wave propagation model (Figure 4). The
536 first two measurements occurred outside the atmospheric perturbation region while the
537 last three sampled the ionosphere after the Lamb wave had passed, but all are at similar
538 local times. There was clear reduction in TEC observed from the last three cases over the
539 geomagnetic equatorial regions, at least 30 degrees wide in latitude, consistent with the
540 spatial extent of the decrease of horizontal TEC measured by GRACE-FO satellites (Figure 2).
541 Ionization crest features accompanying plasma depletions around the magnetic equator are
542 also visible in Figure 4.

543

544 *3.2. Contrasting electron density measurements across the HTHH Lamb wave*

545

546 We found a unique CubeSat observation that demonstrates the association of equatorial
547 plasma depletions with atmospheric waves caused by the volcanic eruption. As mentioned
548 earlier, the Spire CubeSats were equipped with two side-looking antennae pointing in the
549 orbital velocity direction ('front') and the opposite ('back'). This configuration can be
550 exploited to measure contrasting electron density distribution across the Lamb waves even
551 from a single CubeSat. Figure 5a shows the ground tracks of CubeSat FM119 when it
552 crossed the Lamb wave from west to east. The Lamb wave propagated westward while the
553 low inclination CubeSat orbited eastward, and they met at longitude of $\sim 100^\circ\text{E}$.

554

555 On Day 15, the first half of FM119 measurements (the ground track marked red in Figure 5a)
556 took place southwest of Sumatra Island in the Indian Ocean where the Lamb wave

557 perturbation had not yet reached (i.e., outside of the wave boundary). The back-side
558 antenna of FM119 observed GPS signals from G05 at the local time around 19:24. The
559 electron density profile #2 in Figure 5b was derived from this back-side antenna and G05.
560 The peak density is 4.5×10^5 electrons/cm³ at an altitude of 410 km. No perturbation or
561 depletion was identified in the region. The same CubeSat FM119 started to track and occult
562 with respect to G08 via the front-side antenna measuring electron density changes over the
563 South China Sea where the Lamb wave had just passed. This second half of the ground track
564 is marked blue in Figure 5a, and the electron density profile #1 in Figure 5b is derived from
565 these data along this track. By comparing profiles #1 and #2, we can identify a significant
566 plasma depletion along the second half of the track. This degree of sudden density change
567 cannot be accounted for by the local time difference between the two tracks. A
568 differentiating factor between the two tracks is the Lamb wave boundary. The first half
569 track (red) lies beyond the Lamb wave boundaries, while the second half track (blue) falls
570 within them. This observation indicates that Lamb waves were responsible for the
571 ionospheric plasma depletion.

572

573 For comparison, we examined the data from the same pairs of FM119 and GPS satellites on
574 Day 16. The local times and locations of the observations on Day 16 are almost the same as
575 those on Day 15. As we can see in Figure 5a, the FM119 and GPS satellite ground tracks are
576 close to each other on both days. The electron density profiles over the Indian Ocean (#3)
577 and the South China Sea (#4) on Day 16 are shown in Figure 5b with dashed lines. The
578 electron density profiles on Day 16 do not exhibit any anomalous behavior. Thus, the
579 drastic electron depletion in the entire F region is the phenomenon only over the region
580 disturbed by the atmospheric Lamb wave on the day of the HTHH volcanic eruption. This
581 observation is consistent with the findings from the high-altitude TEC measurements by
582 GRACE-FO. The observations of equatorial plasma depletions over a broad range of
583 longitudes by different geodetic satellites can be explained by either the uplift of the F
584 region above satellite altitude (~500-600 km) or the poleward transport of equatorial
585 plasma as a result of the uplift. This is discussed further in Section 4.

586

587 *3.3. Spatiotemporal variations in ionospheric delay from DORIS radio beacons*

588

589 Turning to our fourth geodetic technique, we examined these time-integrated DORIS data to
590 validate our observations of the electron depletion as detected with the other geodetic
591 satellite measurements. Figure 6 illustrates three ground tracks of Jason-3 on 14, 15 and 16
592 January 2022 at UTC around 9:00 (top panels) and 21:00 (bottom panels), from the
593 descending and ascending tracks, respectively. The color code indicates the “relative”
594 ionospheric delay correction in meter (proportional to TEC), for which time series are shown
595 in the middle panels (n.b. the absolute values have no meaning). The “U-shaped” variation
596 is typical as the ionospheric delay is largely dependent on the elevation of a satellite seen
597 from a ground station (right panels). There are ionospheric delay changes as much as 4 m
598 for normal days in two different time periods in each day. That change was reduced to 50
599 cm around UTC 9:00 on Day 15 (Figure 6a). This is a couple of hours after the atmospheric
600 Lamb wave passed through the region around the station BETB in Betio, Republic of Kiribati.
601 This is a large electron depletion in the ionosphere, a reduction by as much as 10 times the
602 usual. The electrons recovered fully by sunrise 12 hours later, as seen from the bottom
603 panel of the data around UTC 21:00 (Figure 6b). Note that almost the same results were
604 observed from Sentinel-6A (not shown), as (at this time) it shared the same orbit as Jason-3
605 with a time difference of only 30 seconds.

606

607 Since the depletion propagates with the westward Lamb wave during the local nighttime,
608 we examined other DORIS stations further west. Figure 7 shows three cases from two
609 stations in Guam, US (MLAC) and in Manila, Philippines (MANB) with various available
610 satellites including Jason-3 (or Sentinel-6A) and Sentinel-3A. Figure 7a shows the electron
611 depletion around UTC 11:00 on Day 15 only when the Lamb wave front passed the Guam
612 station MLAC but not the Manila station MANB; the ionospheric delay was 1-6 m on the
613 normal days while it got reduced to 20 cm, and became somewhat erratic, on Day 15. The
614 jump in the measurement on Day 14 is due to loss of DORIS tracking (cycle slip) that
615 resulted in a ‘reset’ of phase count. When the two pieces are joined together, the delay was
616 about 1 m on Day 14. The reason for the smaller delay on Day 14, relative to the one on
617 Day 16, even though both passes feature similar elevation change, is not known.

618

619 The measurements from MANB to Jason-3 and Sentinel-6A (Figure 7b) show there is no
620 depletion in electrons and ionospheric delay at UTC 10:30 on Day 15, which is similar to

621 normal days. There are approximately 3-3.5 meters of ionospheric delays observed on all
622 three days (two pieces of the Day 14 measurements need to be combined considering the
623 offset introduced by reset of phase count on Day 14). At the time the DORIS observations
624 took place, the Lamb wave had not yet reached the station and thus the DORIS system must
625 have sampled the unperturbed atmosphere. However, just after the Lamb wave passed the
626 Philippines (MANB), the electron depletion was observed from the MANB station and
627 Sentinel-3A around UTC 14:30, by which time the Lamb wave was approaching India (Figure
628 7c). The DORIS data show that the ionospheric delay correction was reduced to 10-20 cm
629 from typical values of a few meters and that this is the result of electron depletion in the
630 ionosphere. The appearance of such reduction is consistent with the timing of the Lamb
631 wave arrival, once again suggesting that the source of electron depletion is likely dynamo
632 electric fields driven by westward propagating Lamb waves.

633

634 **4. Discussion**

635

636 In this section, we discuss the sources of trough-like equatorial plasma depletions across
637 broad longitudes. Various processes, including neutral composition changes, impulsive
638 shock waves and ionospheric uplifts, can be considered as sources of trough-like equatorial
639 plasma depletions during the HTHH volcanic eruption. The volcanic eruption can cause
640 changes in neutral composition by vertical atmospheric expansion including an increase in
641 water vapor in the F region (Choi et al., 2023). Just as impulsive shock waves generated by
642 earthquakes can cause ionospheric plasma depletions via upwelling of the atmosphere or
643 leading to outward plasma flow (Astafyeva et al., 2019; Shinagawa et al., 2013; Zettergren
644 et al, 2017), acoustic shock waves generated by the HTHH volcanic eruption could generate
645 equatorial plasma depletions (Aa et al., 2022; Astafyeva et al., 2022; He et al., 2023). These
646 processes may affect the plasma density in the local region around the epicenter, but the
647 formation of plasma depletions along the magnetic equator is difficult to explain by these
648 mechanisms. The intensity of these processes would be ordered by the distance from the
649 epicenter rather than by magnetic latitude. Alignment of these processes along the
650 magnetic equator has not yet been demonstrated by observations. Additionally, the
651 ionization crest features accompanying equatorial plasma depletions are not explained by
652 the effects of atmospheric expansion and shock waves.

653

654 The development of strong plasma depletions around the magnetic equator, accompanied
655 by density enhancements outside the equatorial region, is a typical characteristic of the
656 equatorial fountain process. An enhanced fountain process after the HTHH volcanic
657 eruption is supported by the Ionospheric Connection Explorer (ICON) satellite observations
658 of upward plasma motion in the F region (Aa et al., 2022; Gasque et al., 2022). These
659 observations are consistent with the detection of intense plasma bubbles in the Asian and
660 Pacific sectors in association with the upward ionospheric drift (Aa et al., 2022; Hong et al.,
661 2022). The ICON observations of strong neutral winds in the Pacific sector (Harding et al.,
662 2022) and the numerical simulations of dynamo electric fields driven by gravity waves (Huba
663 et al., 2023) further support the role of dynamo electric fields.

664

665 Hong et al. (2022) suggested the uplift of the ionosphere to the west of HTHH by the effect
666 of winds driven by the volcanic eruption. Figure 8a (and Figure 4 of Hong et al., 2022)
667 illustrates this process. The Hall current (J_H), driven by westward winds (U), becomes
668 discontinuous at the terminator due to the day-night difference in Hall conductivity (σ_H).
669 Consequently, negative charges accumulate at the terminator, causing the ionosphere in
670 the sunlit region to be lifted perpendicular to the eastward electric field (E) and magnetic
671 field (B), i.e., $E \times B$ drift. This process allows plasma depletion in the equatorial region to
672 persist at night as the lifted region moves into darkness. The ionospheric uplift in the east
673 of HTHH can also be explained by a similar process by changing the wind direction. In the
674 east, as illustrated in Figure 8b, positive charges accumulate at the terminator by eastward
675 winds resulting in the uplift of the ionosphere in the dark region. The optimal time for the
676 development of electric fields is when the passage of the winds coincides with the local
677 sunset. A downward $E \times B$ drift can also be considered as the cause of the plasma
678 depletion in the equatorial region; chemical reactions with molecular gases at lower
679 altitudes reduce the plasma density. However, the enhancement in the plasma density
680 outside the equatorial region accompanied with equatorial plasma depletions is not
681 explained by the downward $E \times B$ drift.

682

683 The magnitude of electric fields would depend on the duration and magnitude of the wind
684 surge and the temporal distance from sunset. Our observations and previous studies show

685 the development of equatorial plasma depletions in the longitude range of $100^{\circ} - 230^{\circ}\text{E}$.
686 The development of equatorial plasma depletions in this longitude region is seen to be
687 related to the occurrence time of the HTHH volcanic eruption around LT 16:30. Earth's
688 rotation speed at the equator (~ 470 m/s at 100 km altitude) is faster than the propagation
689 speed of Lamb waves (300–350 m/s). In the west of HTHH, the arrival time of the Lamb
690 waves at a local region moves to a later LT as the longitude moves westward. Beyond a
691 certain longitude ($\sim 100^{\circ}\text{E}$), the Lamb waves cannot reach the sunset terminator. Then, the
692 process illustrated in Figure 8a does not occur. The arrival time of the Lamb waves at a local
693 region also moves to a later LT as the longitude moves eastward in the east of HTHH.
694 Therefore, the process illustrated in Figure 8b would be effective in the longitude close to
695 HTHH. In our observations (CubeSat observation at UTC 14:12 in Figure 1d), TEC depletion is
696 absent around 260°E , even though the Lamb waves have already passed the region when
697 the CubeSat observations were made. Since the CubeSat observation was made after
698 sunrise (LT 08:12), plasma depletions would be filled by photoionization. However, we
699 cannot rule out the absence of plasma depletion in the region throughout the night because
700 the region was too far away from the sunset terminator (after midnight) when the Lamb
701 waves passed the region.

702

703 We explained the dynamo electric fields using the Hall currents in the E region, but
704 Pedersen currents and F region dynamo electric fields also contribute to the vertical plasma
705 motion at the sunset terminator. Numerical simulations are required to address the
706 contributions of different processes. Rajesh et al. (2022) explained the enhancement of the
707 upward velocity during the HTHH volcanic eruption in association with the F region dynamo
708 instead of the E region dynamo due to reduced E region conductivity. However, Gasque et
709 al. (2022) attributed the ICON observations of the F-region vertical plasma motion to the E
710 region dynamo because this interpretation is consistent with the propagation speed of
711 observed atmospheric waves. The observation time of the vertical plasma motion by ICON
712 in the Northern Hemisphere was around 17:00 LT (Gasque et al., 2022), before sunset in the
713 E region in the Southern Hemisphere. Therefore, the E region might dominate the E and F
714 region coupling, causing the F region plasma motion. Our illustration in Figure 8 aligns with
715 this interpretation.

716

717 5. Conclusion

718

719 We investigated the ionospheric disturbances in the equatorial region caused by the HTHH
720 volcanic eruption on 15 January 2022 by examining the data from four different geodetic
721 ranging systems: between GPS and CubeSat satellites (L1/L2 bands), between two GRACE
722 Follow-On satellites (K/Ka bands), between altimeter satellites and the ocean surface (C/Ku
723 bands) and between the ground beacon stations and DORIS-equipped satellites (S/VHF
724 bands). These dual-frequency satellite systems were used to measure electron density
725 changes integrated along ray paths in line of sight from transmitters to receivers. Five types
726 of TEC measurements were taken (1) at high altitude above CubeSat satellite altitude of
727 ~550 km, (2) at low altitude below ~550 km, (3) at GRACE Follow-On satellite altitude of
728 ~480 km, (4) for the entire vertical column of ionosphere below the altimeter satellite
729 altitude of 800 km, and (5) in line-of-sight between the ground and satellites above top of
730 the ionosphere. While the CubeSats, altimeter and DORIS satellites sample TEC variations
731 over a wide range of altitude, the GRACE Follow-On satellites sample horizontal TEC
732 variation along the orbit at a fixed altitude.

733

734 The observations from these systems reveal the development of anomalous plasma
735 depletions in the equatorial region across a broad longitude range after the HTHH volcanic
736 eruption. On the basis of our observations, the conditions for the development of the
737 depletions are (1) inside the boundary of the Lamb waves produced by the volcanic eruption
738 and (2) close to the sunset terminator when the Lamb waves pass the region. These
739 conditions facilitate the generation of dynamo electric fields, and by which the ionosphere
740 can be lifted to high altitudes resulting in plasma depletion in the equatorial region.
741 However, the dynamical coupling of the ionosphere with other waves such as secondary
742 gravity waves (Vadas et al., 2023a & 2023b) cannot be ruled out. The size and magnitude of
743 depletion by the HTHH volcanic eruption were anomalously large; for example, at the
744 satellite altitude of 500 km, the reduction of electrons by more than 1,000 times was
745 recorded from the GRACE Follow-On K/Ka band ranging system over the extensive areas of
746 ~3,000 km around the geomagnetic equator. Such magnitudes of depletions are unusual
747 phenomena, indicating that Earth's hazardous surface events can be a source of severe
748 ionospheric disturbances.

749

750 This study demonstrated the use of geodetic (non-space-weather-dedicated) satellite
751 constellations for capturing ionospheric perturbations induced by ground events. These
752 geodetic satellite systems can be cohesively operated and analyzed to support space
753 weather monitoring and prediction and, possibly, activity concerning navigation,
754 communications, and surveillance systems that may be affected by extreme event
755 irregularities in the ionosphere. All four geodetic techniques (GNSS high-low tracking, radio
756 occultation, low-low intersatellite tracking, and Doppler ground beacon) and increasing
757 number of suitable satellites will produce new opportunities to monitor the ionosphere
758 continuously and to study the atmosphere-surface coupled processes including atmospheric
759 and gravity waves caused by earthquakes, tsunamis, extreme weather, and volcanic
760 eruptions.

761

762 **Acknowledgement.** This work was supported by NASA's GRACE and GRACE Follow-On
763 science team project, Australian Research Council Discovery Program projects
764 (DP170100224 & DP240102399), and the projects sponsored by Geoscience Australia,
765 FrontierSI, and NSW Space Research Network (RP220204). HK acknowledges support by
766 NSF-AGS2029840 and NASA-NNH19ZDA001N. We are indebted to two anonymous
767 reviewers who made several valuable suggestions for this work.

768

769 **Open Research.** The Spire Global's GNSS/GPS tracking datasets (aka RINEX files) are publicly
770 available via NASA's Commercial Smallsat Data Acquisition Program from the NASA program
771 website <https://csdap.earthdata.nasa.gov>. The GRACE-FO Level-1B data (KBR1B for K-Band
772 ionospheric correction and GNV1B for orbit) used in this study are publicly available through
773 <https://podaac.jpl.nasa.gov/cloud-datasets>. Interested readers can begin by searching
774 GRACE-FO under Projects and Level-1B under Processing Levels. The dual-frequency
775 satellite altimeter data (including ionosphere correction) are available at
776 <http://rads.tudelft.nl> with associated software at <https://github.com/remkos/rads>. The
777 DORIS/RINEX data are available from the data centers of the International DORIS Service
778 (IDS), at the Institut National de l'Information Géographique et Forestière (IGN-France) and
779 the NASA CDDIS (Noll and Soudarin, 2006; Noll, 2010) as described at [https://ids-](https://ids-doris.org/ids/data-products/tables-of-data-products.html)
780 [doris.org/ids/data-products/tables-of-data-products.html](https://ids-doris.org/ids/data-products/tables-of-data-products.html)

781

782

783 **References**

- 784 Aa, E., Zhang, S., Erickson, P. J., Vierinen, J., Anthea, J., Goncharenko, L. P., Spicher, A., &
785 Rideout, W. C. (2022). Significant equatorial plasma bubbles and global ionospheric
786 disturbances after the 2022 Tonga volcano eruption. *Space Weather*, 20,
787 e2022SW003101.
- 788 Amores, A., Monserrat, S., Marcos, M., Argüeso, D., Villalonga, J., Jordà, G., & Gomis,
789 D.(2022) Numerical simulation of atmospheric Lamb waves generated by the 2022
790 Hunga-Tonga volcanic eruption. *Geophysical Research Letters*, 49, e2022GL098240.
- 791 Arnold, D. V., W. K. Melville, R. H. Stewart, J. A. Kong, W. C. Keller, E. Lamarre (1995),
792 Measurements of electromagnetic bias at Ku and C bands, *J. Geophys. Res.*, 100, 969-980.
- 793 Astafyeva, E., Rolland, L., Lognonné, P., Khelifi, K., & Yahagi, T. (2013). Parameters of seismic
794 source as deduced from 1 Hz ionospheric GPS data: Case-study of the 2011 Tohoku-oki
795 event. *Journal of Geophysical Research: Space Physics*, 118, 5942–5950. [https://doi.org/](https://doi.org/10.1002/jgra.50556)
796 10.1002/jgra.50556
- 797 Astafyeva, E., Maletckii, B., Mikesell, T. D., Munaibari, E., Ravanelli, M., Coisson, P., ... &
798 Rolland, L. (2022). The 15 January 2022 Hunga Tonga eruption history as inferred from
799 ionospheric observations. *Geophysical Research Letters*, 49(10), e2022GL098827.
- 800 Bernhardt, P. A., C. L. Siefring, I. J. Galysh, T. F. Rodillo, D. E. Koch, T. L. MacDonald, M. R.
801 Wilkens, and G. P. Landis (2006), Ionospheric applications of the Scintillation and
802 Tomography Receiver in Space (CITRIS) used with the DORIS radio beacon network, *J.*
803 *Geod.*, 80, 473–485.
- 804 Bilitza, D. (2018), IRI the International Standard for the Ionosphere, *Adv. Radio Sci.*, 16, 1-11,
805 <https://doi.org/10.5194/ars-16-1-2018>
- 806 Case, K.; Gerhard, K.; Wu, S.C. Gravity Recovery and Climate Experiment (GRACE) Level-1B
807 Data Product User Handbook. JPL D-22027. 2010. Available online:
808 https://opendap.jpl.nasa.gov/opendap/allData/grace/docs/Handbook_1B_v1.3.pdf
- 809 Chelton, D. B., J. C. Ries, B. J. Haines, L.-L. Fu, P. S. Callahan (2001), Satellite altimetry, in
810 *Satellite Altimetry and Earth Sciences: A Handbook of Techniques and Applications* (Ed:
811 L.-L. Fu and A. Cazenave), pp. 1-131, Academic Press.
- 812 Choi, J. M., Lin, C. C. H., Rajesh, P. K., Lin, J. T., Chou, M., Kwak, Y. S., & Chen, S. P. (2023).
813 Giant ionospheric density hole near the 2022 Hunga-Tonga volcanic eruption: multi-point
814 satellite observations. *Earth, Planets and Space*, 75(1), 184.

815 Dettmering D, Limberger M, Schmidt M (2014) Using DORIS measurements for modeling the
816 vertical total electron content of the Earth's ionosphere. *J Geod* 88:1131–1143. DOI
817 10.1007/s00190-014-0748-2

818 Fejer, B. G., de Paula, E. R., Gonzalez, S. A., & Woodman, R. F. (1991). Average vertical and
819 zonal F region plasma drifts over Jicamarca. *Journal of Geophysical Research*, 96 (A8),
820 13,901–13, 906. <https://doi.org/10.1029/91JA01171>

821 Fleury R, Foucher F, Lassudrie-Duchesne P (1991) Global TEC measurement capabilities of
822 the DORIS system. *Adv Space Res* 11(10):51–54. doi:10.1016/0273-1177(91)90321-A

823 Gasque, L. C., Wu, Y. J., Harding, B. J., Immel, T. J., & Triplett, C. C. (2022). Rapid volcanic
824 modification of the E-region dynamo: ICON's first glimpse of the Tonga eruption.
825 *Geophysical Research Letters*, 49(18), e2022GL100825.

826 Hajj, G.A., E.R. Kursinski, L.J. Romans, W.I. Bertiger, S.S. Leroy, (2002), A technical
827 description of atmospheric sounding by GPS occultation, *Journal of Atmospheric and*
828 *Solar-Terrestrial Physics*, 64, 451-469.

829 Han, S.-C., McClusky, S., Mikesell, T. D., Tregoning, P., & Sauber, J. (2022). Looking to the sky
830 for better tsunami warnings. *Eos*, 103, 20–23. <https://doi.org/10.1029/2022EO220519>

831 Han, S.-C., McClusky, S., Mikesell, T. D., Rolland, L., Okal, E., & Benson, C. (2023). CubeSat
832 GPS observation of traveling ionospheric disturbances after the 2022 Hunga-Tonga
833 Hunga-Ha'apai volcanic eruption and its potential use for tsunami warning. *Earth and*
834 *Space Science*, 10, e2022EA002586. <https://doi.org/10.1029/2022EA002586>

835 Hanson, W. B., & Moffett, R. J. (1966). Ionization transport effects in the equatorial F region.
836 *Journal of Geophysical Research*, 71(23), 5559–5572.
837 <https://doi.org/10.1029/JZ071i023p05559>

838 Harding, B. J., Wu, Y. J., Alken, P., Yamazaki, Y., Triplett, C. C., Immel, T. J. et al. (2022).
839 Impacts of the January 2022 Tonga Volcanic Eruption on the Ionospheric Dynamo: ICON-
840 MIGHTI and Swarm Observations of Extreme Neutral Winds and Currents. *Geophysical*
841 *Research Letters*, 49(9), <https://doi.org/10.1029/2022gl098577>

842 Hartmann, G., Leitinger, R. (1984): *Range errors due to ionospheric and tropospheric effects*
843 *for signal frequencies above 100 MHz*. *Bull. Geod.* 58: 109–136.

844 He, J., Astafyeva, E., Yue, X., Ding, F., & Maletckii, B. (2023). The giant ionospheric depletion
845 on 15 January 2022 around the Hunga Tonga-Hunga Ha'apai volcanic eruption. *Journal of*
846 *Geophysical Research: Space Physics*, 128(1), e2022JA030984.

847 Hong, J., Kil, H., Lee, W. K., Kwak, Y.-S., Choi, B.-K., & Paxton, L. J. (2022). Detection of
848 different properties of ionospheric perturbations in the vicinity of the Korean Peninsula
849 after the Hunga-Tonga volcanic eruption on 15 January 2022. *Geophysical Research*
850 *Letters*, 49, e2022GL099163. <https://doi.org/10.1029/2022GL099163>

851 Huba, J. D., G. Joyce, and J. Krall (2008), Three-dimensional equatorial spread F modeling,
852 *Geophys. Res. Lett.*, 35, L10102, doi:10.1029/2008GL033509.

853 Huba, J. D., Becker, E., & Vadas, S. L. (2023). Simulation study of the 15 January 2022 Tonga
854 event: Development of super equatorial plasma bubbles. *Geophysical Research Letters*,
855 50(1), e2022GL101185.

856 Immel, T. J., E. Sagawa, S. L. England, S. B. Henderson, M. E. Hagan, S. B. Mende, H. U. Frey,
857 C. M. Swenson, and L. J. Paxton (2006), Control of equatorial ionospheric morphology by
858 atmospheric tides, *Geophys. Res. Lett.*, 33, L15108, doi:10.1029/2006GL026161.

859 Jayles C., Chaveau J.P., Auriol A. (2015). DORIS/DIODE Real-Time Orbit Determination on
860 Board Saral/AltiKa, *Marine Geodesy*, 38, 233–240, doi: 10.1080/01490419.2015.1015695.

861 Kedar, S., G. A. Hajj, B. D. Wilson, and M. B. Heflin, The effect of the second order GPS
862 ionospheric correction on receiver positions, *Geophys. Res. Lett.*, 30(16), 1829,
863 doi:10.1029/2003GL017639, 2003.

864 Kelley, M. C. (2009). *The earth's ionosphere : plasma physics and electrodynamics*, San
865 Diego, CA: Academic Press.

866 Kil, H., and L. J. Paxton (2006), Ionospheric disturbances during the magnetic storm of 15
867 July 2000: Role of the fountain effect and plasma bubbles for the formation of large
868 equatorial plasma density depletions, *J. Geophys. Res.*, 111, A12311,
869 doi:10.1029/2006JA011742.

870 Kil, H., and W. K. Lee (2013), Are plasma bubbles a prerequisite for the formation of broad
871 plasma depletions in the equatorial F region?, *Geophys. Res. Lett.*, 40, 3491–3495,
872 doi:10.1002/grl.50693.

873 Kil, H., Oh, S.-J., Paxton, L. J., & Fang, T.-W. (2009). High-resolution vertical drift model
874 driven from the ROCSAT-1 data. *Journal of Geophysical Research*, 114 (A10).
875 <https://doi.org/10.1029/2009JA014324>

876 Kil, H., Y. S. Kwak, L. J. Paxton, R. R. Meier, and Y. Zhang (2011), O and N₂ disturbances in the
877 F region during the 20 November 2003 storm seen from TIMED/GUVI, *J. Geophys. Res.*,
878 116, A02314, doi:10.1029/2010JA016227.

879 Kim, J. Simulation Study of a Low-Low Satellite-to-Satellite Tracking Mission. Ph.D. Thesis,
880 University of Texas, Austin, TX, USA, 2000.

881 Landerer, F.W.; Flechtner, F.M.; Save, H.; Webb, F.H.; Bandikova, T.; Bertiger, W.I.;
882 Bettadpur, S.V.; Byun, S.H.; Dahle, C.; Dobslaw, H.; et al. Extending the global mass
883 change data record: GRACE Follow-On instrument and science data performance.
884 *Geophys. Res. Lett.* 2020, 47, e2020GL088306.

885 Lee, CK., Han, SC., Bilitza, D. *et al.* Validation of international reference ionosphere models
886 using in situ measurements from GRACE K-band ranging system and CHAMP planar
887 Langmuir probe. *J Geod* **85**, 921–929 (2011). <https://doi.org/10.1007/s00190-011-0442-6>

888 Lee, W. K., H. Kil, Y.-S. Kwak, L. J. Paxton, Y. Zhang, I. Galkin, and I. S. Batista (2014),
889 Equatorial broad plasma depletions associated with the enhanced fountain effect, *J.*
890 *Geophys. Res. Space Physics*, 119, 402–410, doi:10.1002/2013JA019137.

891 Leick, A. (1995): *GPS satellite surveying*. 2nd. Edition, John Wiley, New York.

892 Lemoine F.G., Chinn D.S., Zelensky N.P., Beall J.W., Le Bail K. (2016). The development of the
893 GSFC DORIS contribution to ITRF2014, *Adv. Space Res.*, 58(12), 2520–2542, doi:
894 10.1016/j.asr.2015.12.043.

895 Lemoine J.M., Capdeville H., Soudarin L. (2016). Precise orbit determination and station
896 position estimation using DORIS RINEX data, *Adv. Space Res.*, 58(12), 2677–2690, doi:
897 10.1016/j.asr.2016.06.024.

898 Lin, J., Rajesh, P. K., Lin, C. C. H., Chou, M., Liu, J., Yue, J., Hsiao, T., Tsai, H., Chao, H., & Kung,
899 M. (2022). Rapid Conjugate Appearance of the Giant Ionospheric Lamb Wave Signatures
900 in the Northern Hemisphere After Hunga-Tonga Volcano Eruptions. *Geophysical Research*
901 *Letters*, 49(8), 1–6. <https://doi.org/10.1029/2022GL098222>

902 Liu A., Wang N., Dettmering D., Schmidt M., Wang L., Yuan H. (2023). “Using DORIS data for
903 validating real-time GNSS ionosphere maps”, *Adv. Space Res.*, 72(1), 115-128,
904 <https://doi.org/10.1016/j.asr.2023.01.050>.

905 Mannucci, A. J., Tsurutani, B. T., Iijima, B. A., Komjathy, A., Saito, A., Gonzalez, W. D., et al.
906 (2005). Dayside global ionospheric response to the major interplanetary events of
907 October 29–30, 2003 “Halloween Storms”. *Geophysical Research Letters*, 32, L12S02.
908 <https://doi.org/10.1029/2004GL021467>

909 Martire, L., Runge, T.F., Meng, X. et al. (2024). The JPL-GIM algorithm and products: multi-
910 GNSS high-rate global mapping of total electron content. *J Geod* 98, 44.
911 <https://doi.org/10.1007/s00190-024-01860-3>

912 Matoza, R. S. et al. (2022). Atmospheric waves and global seismoacoustic observations of
913 the January 2022 Hunga eruption, Tonga, *Science*, 377, 95-100.
914 DOI:10.1126/science.abo7063

915 Müller, V.; Hauk, M.; Misfeldt, M.; Müller, L.; Wegener, H.; Yan, Y.; Heinzl, G. Comparing
916 GRACE-FO KBR and LRI Ranging Data with Focus on Carrier Frequency Variations. *Remote*
917 *Sens.* 2022, 14, 4335. <https://doi.org/10.3390/rs14174335>

918 Nguyen, V. A., Nogués-Correig, O., Yuasa, T., Masters, D., & Irisov, V. (2020). Initial GNSS
919 phase altimetry measurements from the spire satellite constellation. *Geophysical*
920 *Research Letters*, 47, e2020GL088308. <https://doi.org/10.1029/2020GL088308>

921 Noll, C.E, Soudarin L. (2006). On-line Resources Supporting the Data, Products, and
922 Information Infrastructure for the International DORIS Service. *J Geodesy*, 80, 419–427,
923 doi: 10.1007/s00190-006-0051-y.

924 Noll, C.E (2010). The crustal dynamics data information system: A resource to support
925 scientific analysis using space geodesy, *Adv. Space Res.*, 45(12), 1421–1440, doi:
926 10.1016/j.asr.2010.01.018.

927 Rajesh, P. K., Lin, C. C. H., Lin, J. T., Lin, C. Y., Liu, J. Y., Matsuo, T., ... & Tsai, H. F. (2022).
928 Extreme poleward expanding super plasma bubbles over Asia-Pacific region triggered by
929 Tonga volcano eruption during the recovery-phase of geomagnetic storm. *Geophysical*
930 *Research Letters*, 49(15), e2022GL099798.

931 Ray, R. D. (2020). Daily harmonics of ionospheric total electron content from satellite
932 altimetry, *J. Atmos. Sol. Terr. Phys.*, 209, 105423, 10.1016/j.jastp.2020.105423

933 Saunier J. (2023). The DORIS network: Advances achieved in the last fifteen years, *Adv.*
934 *Space Res.*, 72, 3–22, doi: 10.1016/j.asr.2022.07.016.

935 Scharroo, R., Smith, W. H. F. (2010), A global positioning system-based climatology for the
936 total electron content in the ionosphere, *J. Geophys. Res.*, 115, A10318, doi:
937 10.1029/2009JA014719

938 Scherliess, L., & Fejer, B. G. (1997). Storm time dependence of equatorial disturbance
939 dynamo zonal electric fields. *Journal of Geophysical Research*, 102(A11),
940 <https://doi.org/10.1029/97JA02165>

941 Scherliess, L., & Fejer, B. G. (1999) Radar and satellite global equatorial F region vertical drift
942 model. *Journal of Geophysical Research*, 104(A4), 6829–6842.
943 <https://doi.org/10.1029/1999JA900025>

944 Seeber, G. (2003): *Satellite Geodesy: Foundations, Methods, and Applications*. 2nd Edition,
945 W. de Gruyter, Berlin-New York.

946 Shinagawa, H., Tsugawa, T., Matsumura, M., Iyemori, T., Saito, A., Maruyama, T., et al.
947 (2013). Two-dimensional simulation of ionospheric variations in the vicinity of the
948 epicenter of the Tohoku-oki earthquake on 11 March 2011. *Geophysical Research Letters*,
949 40, 5009–5013. <https://doi.org/10.1002/2013GL057627>

950 Sultan, P. J. (1996). Linear theory and modeling of the Rayleigh-Taylor instability leading to
951 the occurrence of equatorial spread F. *Journal of Geophysical Research: Space Physics*,
952 101(A12), 26875–26891. <https://doi.org/10.1029/96JA00682>

953 Tapley B.D., Bettadpur S., Watkins M., Reigber C. (2004). The Gravity Recovery and Climate
954 Experiment: mission overview and early results. *Geophys Res Lett* .31(9), L09607,
955 <https://doi.org/10.1029/2004GL019779>.

956 Teunissen, P. J. G., Kleusberg, A. (1998): *GPS for Geodesy*. Second Edition. Springer-Verlag,
957 Berlin-Heidelberg-New York.

958 Themens, D. R., Watson, C., Žagar, N., Vasylkevych, S., Elvidge, S., McCaffrey, A., et al. (2022).
959 Global propagation of ionospheric disturbances associated with the 2022 Tonga volcanic
960 eruption. *Geophysical Research Letters*, 49, e2022GL098158.
961 <https://doi.org/10.1029/2022GL098158>

962 Thomas, J. (1999), An analysis of gravity field estimation based on intersatellite dual one-
963 way biased ranging, JPL Publ. 98-15, pp. 3 – 13, Jet Propul. Lab., Pasadena, Calif., May.

964 Tran, N., D. Vandemark, E. D. Zaron, P. Thibaut, G. Dibarboure, N. Picot (2021), Assessing the
965 effects of sea-state related errors on the precision of high-rate Jason-3 altimeter sea level
966 data, *Adv. Space Res.*, 68, 963-977. DOI: 10.1016/j.asr.2019.11.034

967 Vadas, S. L., Becker, E., Figueiredo, C., Bossert, K., Harding, B. J., & Gasque, L. C. (2023a).
968 Primary and secondary gravity waves and large-scale wind changes generated by the
969 Tonga volcanic eruption on 15 January 2022: Modeling and comparison with ICON-
970 MIGHTI winds. *Journal of Geophysical Research: Space Physics*, 128(2), e2022JA031138.

971 Vadas, S. L., Figueiredo, C., Becker, E., Huba, J. D., Themens, D. R., Hindley, N. P., ... &
972 Bossert, K. (2023b). Traveling ionospheric disturbances induced by the secondary gravity

973 waves from the Tonga eruption on 15 January 2022: Modeling with
974 MESORAC/HIAMCM/SAMI3 and comparison with GPS/TEC and ionosonde data. *Journal*
975 *of Geophysical Research: Space Physics*, 128(6), e2023JA031408.

976 Walsh, E. J., et al. (1991), Frequency dependence of electromagnetic bias in radar altimeter
977 sea surface range measurements, *J. Geophys. Res.*, 96, 20571-20583.

978 Wen, H.; Gerhard, K.; William, B.; Meegyeong, P.; Carly, S.; Felix, L. Gravity Recovery and
979 Climate Experiment (GRACE) Follow-On (GRACE-FO) Level-1 Data Product User Handbook;
980 JPL D-56935. 2019. Available online: [https://podaac-tools.jpl.](https://podaac-tools.jpl.nasa.gov/drive/files/allData/gracefo/docs/GRACE-FO_L1_Handbook.pdf)
981 [nasa.gov/drive/files/allData/gracefo/docs/GRACE-FO_L1_Handbook.pdf](https://podaac-tools.jpl.nasa.gov/drive/files/allData/gracefo/docs/GRACE-FO_L1_Handbook.pdf).

982 Willis, P., et al. (2010) The International DORIS Service (IDS): Toward maturity, *Adv. Space*
983 *Res.*, 45(12), 1408–1420, doi: 10.1016/j.asr.2009.11.018.

984 Wu, S.-C., G. Kruizinga, W. Bertiger (2004), Algorithm Theoretical Basis Document for GRACE
985 Level-1B Data Processing V1.1, Jet Propulsion Laboratory, California Institute of
986 Technology, GRACE 327-741 (JPL D-27672)

987 Wright, C., Hindley, N., Alexander, M. J., Barlow, M., Hoffmann, L., Mitchell, C. et al. (2022).
988 Surface-to-space atmospheric waves from Hunga Tonga–Hunga Ha’apai eruption, *Nature*,
989 609, 741-746. <https://doi.org/10.1038/s41586-022-05012-5>

990 Yokoyama, T., Shinagawa, H., & Jin, H. (2014). Nonlinear growth, bifurcation, and pinching
991 of equatorial plasma bubble simulated by three-dimensional high-resolution bubble
992 model. *Journal of Geophysical Research: Space Physics*, 119, 10,474–10,482.
993 <https://doi.org/10.1002/2014JA02070>

994 Zhang, S. R., Vierinen, J., Aa, E., Goncharenko, L. P., Erickson, P. J., Rideout, W., ... & Spicher,
995 A. (2022). 2022 Tonga volcanic eruption induced global propagation of ionospheric
996 disturbances via Lamb waves. *Frontiers in Astronomy and Space Sciences*, 9, 871275.

997 Zettergren, M. D., Snively, J. B., Komjathy, A., & Verkhoglyadova, O. P. (2017). Nonlinear
998 ionospheric responses to large-amplitude infrasonic-acoustic waves generated by
999 undersea earthquakes. *Journal of Geophysical Research: Space Physics*, 122, 2272–2291.
1000 <https://doi.org/10.1002/2016JA023159>

1001 Zlotnicki, V. (1994). Correlated environmental corrections in TOPEX/POSEIDON, with a note
1002 on ionospheric accuracy. *J. Geophys. Res.*, 99, 24,907-24,914.

1003

1004

1005 **Figure Caption**

1006

1007 **Figure 1.** CubeSat GPS measurements of electron density (N_e) changes. Each panel, (a)
1008 through (d), presents ground tracks of CubeSats and GPS satellites and the atmospheric
1009 Lamb wave front at different times in UTC on the days before, on and after the HTHH
1010 volcanic eruption. It also shows vertical distribution of electron density (N_e) computed from
1011 the slant TEC measurements between CubeSats and GPS satellites. The local time (LT) and
1012 UTC time represent the average epoch of the TEC measurements. The CubeSat and GPS
1013 satellites repeat their ground tracks roughly within a few hundreds of km every day. This
1014 repeated orbital sampling allows us to monitor electron content variations over the same
1015 region at the same time but on different days. The geomagnetic equator is depicted as a
1016 black curve on the map, the location of HTHH volcano is marked with a magenta triangle,
1017 the Lamb wave front is shown in black dashed curve, and lastly, the tangent points are
1018 shown in stars next to the CubeSat ground tracks. Four pairs of CubeSat and GPS satellites
1019 were examined, including FM143 – G04, FM117 – G31, FM115 – G08, and FM119 – G29, to
1020 find changes along the geomagnetic equator over a wide range of longitudes. Three cases
1021 of local nighttime measurements, (a)-(c), found substantial reduction in electron density
1022 after the Lamb wave passage on the day of HTHH volcanic eruption, while the last case (d)
1023 shows no change observed at the local morning time, due possibly to regeneration of
1024 plasma. Three different colors are used for three different days. The sunset terminator at
1025 the equator at the time of the Day 15 measurements is shown in gray dashed line. This
1026 same color and symbol is used throughout all figures in this paper. East of this sunset
1027 terminator is the night side.

1028

1029 **Figure 2.** GRACE Follow-On measurements of changes in TEC integrated over ~ 220 km
1030 between two satellites at the altitude of ~ 480 km, from 6 am to 4 pm in UTC (local time of 8-
1031 9 pm) on 14, 15, and 16 of January 2022. Each panel shows the ground tracks of GRACE
1032 Follow-On satellites and the Lamb wave front (black dashed curves) propagated from the
1033 HTHH volcano (magenta triangle) on 15 of January 2022. The direction of the wave
1034 propagation is shown with small arrows. The geomagnetic equator is marked by a heavy
1035 black curve. The TEC measurements between the two satellites are presented along the
1036 geomagnetic latitude of the midpoint of two satellites and shown in a log scale (as the

1037 change is up to 4 orders of magnitude). Three measurements on three days are compared
1038 to highlight the drastically reduced TEC only on the eruption day and only over the region
1039 after the volcano-induced Lamb wave passage, i.e., panels of (b) through (d). The
1040 measurements outside of the Lamb wave boundary show no change in electron density.
1041 This series of measurements clearly suggest the depletion associated with the Lamb wave.
1042 It is observed that the volcano-induced TEC reduction at the satellite altitude (~480 km) is
1043 from 10 to 0.001 TECU and is as wide as 3,000 km along the geomagnetic equator.

1044

1045 **Figure 3.** Sentinel-3A satellite altimeter measurements of vertical TEC changes on 14, 15,
1046 and 16 of January 2022. During UTC approximately 9 to 11 am (local time around 10:30pm),
1047 the satellite flew over the Pacific Ocean and detected ionospheric changes associated with
1048 the volcano-driven Lamb wave (black dashed curve). The map includes the ground tracks as
1049 well as the Lamb wave front at the time of altimeter measurements. The vertical TEC
1050 measurements (blue dots) on three consecutive days are compared with two model outputs
1051 of JPL's Global Ionosphere Maps GIM (red curve) and NOAA's Ionospheric Climatology NIC09
1052 (yellow curve). The smoothed (60-sec moving average) TEC measurements are also shown
1053 for better comparison with the models (magenta curve). On the day of eruption, the
1054 vertical TEC measurements along the satellite orbit decreased as much as 10 TECU relative
1055 to the model predictions and relative to the measurements on the other ('normal') days.
1056 The spatial extent of the TEC depletion is around 3,000 km around the geomagnetic latitude,
1057 similar to the GRACE-FO measurements of high-altitude TEC variation.

1058

1059 **Figure 4.** Sentinel-3B satellite altimeter measurements (bottom panels) of vertical TEC
1060 changes during the five consecutive ascending orbits (top panels) on 15 of January, from
1061 UTC hours from 5 am to 12 pm and Local time around 10 pm. These data allow us to
1062 examine the TEC measurements with the Lamb wave propagation every orbit (~100 min
1063 sampling). The first two cases (first two left columns) present the data outside the Lamb
1064 wave boundary (black dashed curve), while the last three show the measurements inside,
1065 but all are nearly at the same local time. The large depletion is found only in the region
1066 around the geomagnetic equator and only inside of the Lamb wave boundary.

1067

1068 **Figure 5.** CubeSat GPS measurements of electron density changes across the Lamb wave.
1069 This example is similar to Figure 1, but it uses the two GPS antennae pointing in the satellite
1070 velocity direction and its opposite. The same CubeSat FM119 tracked GPS satellite G05
1071 from the backward antenna while it picked up G08 through the front antenna. The data on
1072 two days of 15 and 16 of January were compared (as no data on 14 of January is available).
1073 The pair of FM119 – G05 sampled the electron density outside the Lamb wave above the
1074 Indian Ocean while the other pair of FM119 – G08 measured inside the Lamb wave above
1075 the Southern China Sea. The four vertical electron density profiles on two days and from
1076 two different locations (outside and inside of the wave) are compared in the right panel. It
1077 is clear that drastic reduction of electron density was observed only from the inside region
1078 and only on the eruption day. This highlights excellent samplings of measuring dynamical
1079 ionospheric anomalies associated with the propagating atmospheric perturbation by
1080 exploiting the multi-faceted GPS antennae from a single CubeSat.

1081
1082 **Figure 6.** (a) The left panel shows the descending ground tracks of Jason-3 on 14, 15, and 16
1083 January 2022 around the ground station BETB (marked in magenta inverted triangle). The
1084 color code presents the time-integrated ionospheric delay in range (i.e., biased ionospheric
1085 delay in meter) along the ground track each day. The biased delay and the elevation angle
1086 of the satellite seen from the ground station are shown as a time-series in the middle and
1087 right panels, respectively, for each track for three different days. The delay varies over ~4 m
1088 on normal days, however, there was large reduction in the delay around UTC 9:00 on Day 15,
1089 after the Lamb wave passage. (b) The same as (a), but from the ascending track data 12
1090 hours later. This shows the ionosphere returned to its normal condition. For (b), the Lamb
1091 wave front was around the antipode near the southern tip of Algeria, which is outside of the
1092 map.

1093
1094 **Figure 7.** (a) The same as Figure 6, but between the Sentinel-6A and the ground station
1095 MLAC in Guam, USA. It demonstrates large reduction in the ionospheric delay associated
1096 with the electron depletion after the Lamb wave passage. (b) The case between Jason-3
1097 and MANB (a ground station in Manila, Philippines) between UTC 10:00 – 11:30, prior to the
1098 Lamb wave arrival. (c) The case between Sentinel-3A and MANB after the Lamb wave

1099 passed the region (UTC 14:00-15:00). The reduction in the ionospheric delay was found only
1100 after the Lamb wave.

1101

1102 **Figure 8.** Schematic illustration of the generation of electric fields (a) west and (b) east of
1103 HTHH in association with the volcanic eruption. Winds (\mathbf{U}) from the volcanic eruption
1104 generate westward (a) and eastward (b) Hall currents (\mathbf{J}_H) in the E region, and negative (a)
1105 and positive (b) charges accumulate at the terminator by the day-night difference of the Hall
1106 conductivity (σ_H). The polarization electric fields (\mathbf{E}) in the E region map into the F region,
1107 resulting in the upward $\mathbf{E} \times \mathbf{B}$ plasma drift in sunlight (a) and darkness (b).

1108

Figure 1.

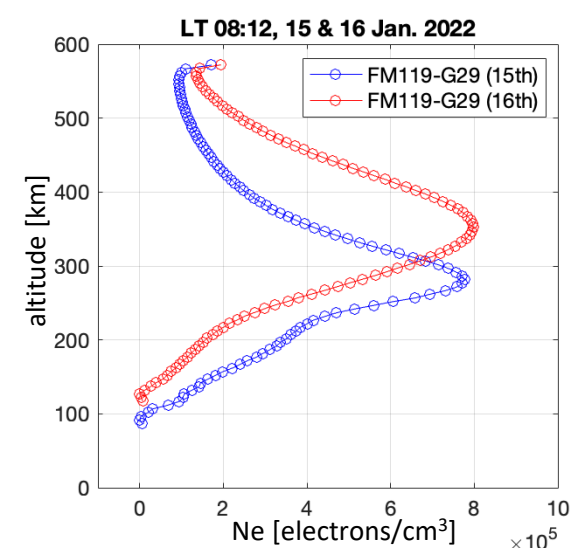
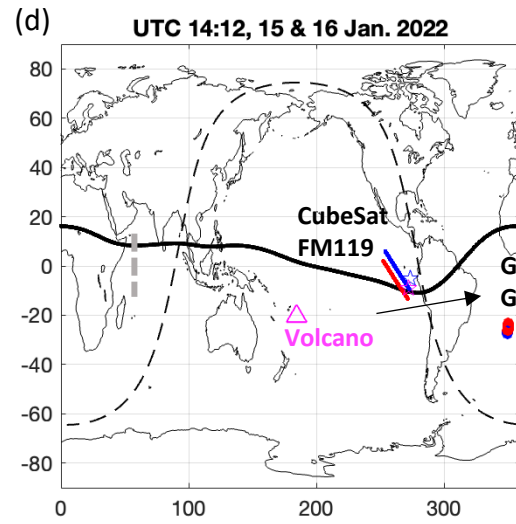
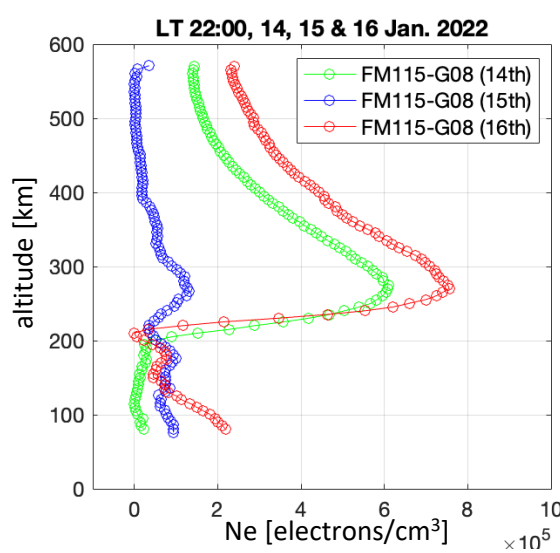
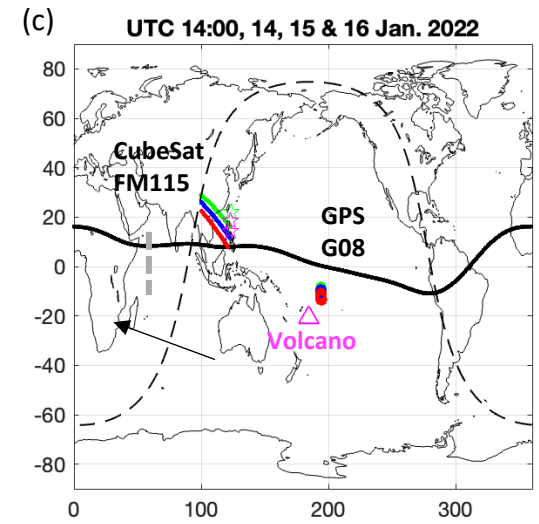
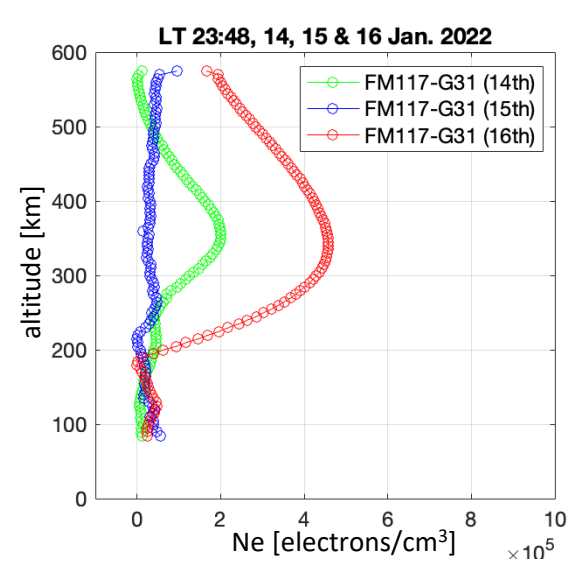
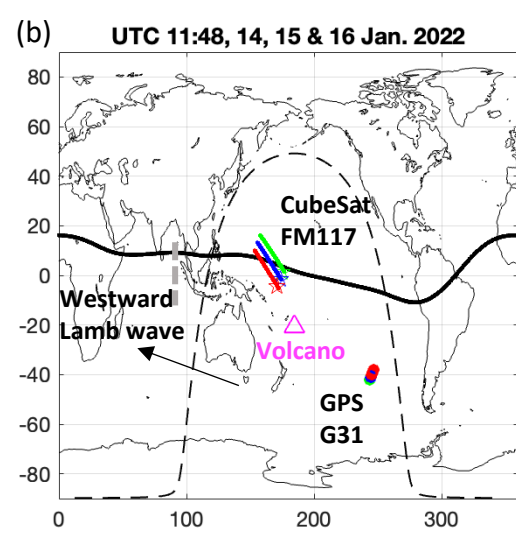
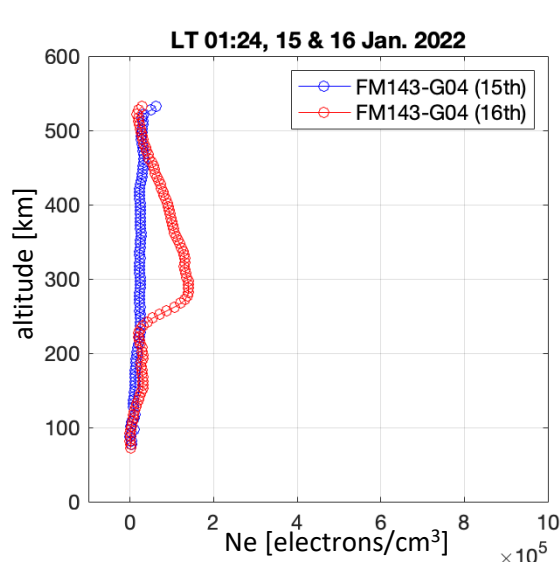
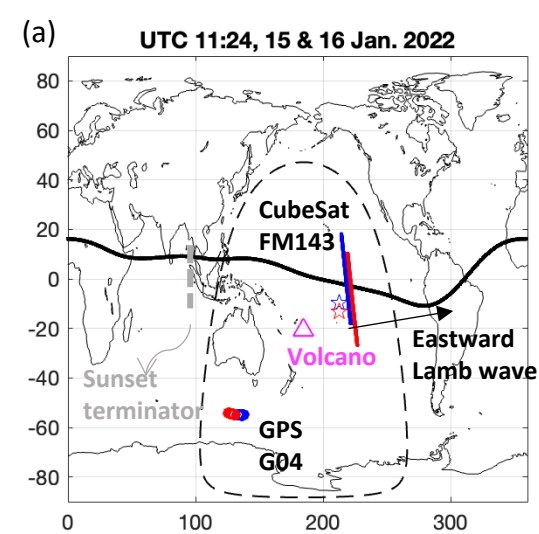


Figure 2.

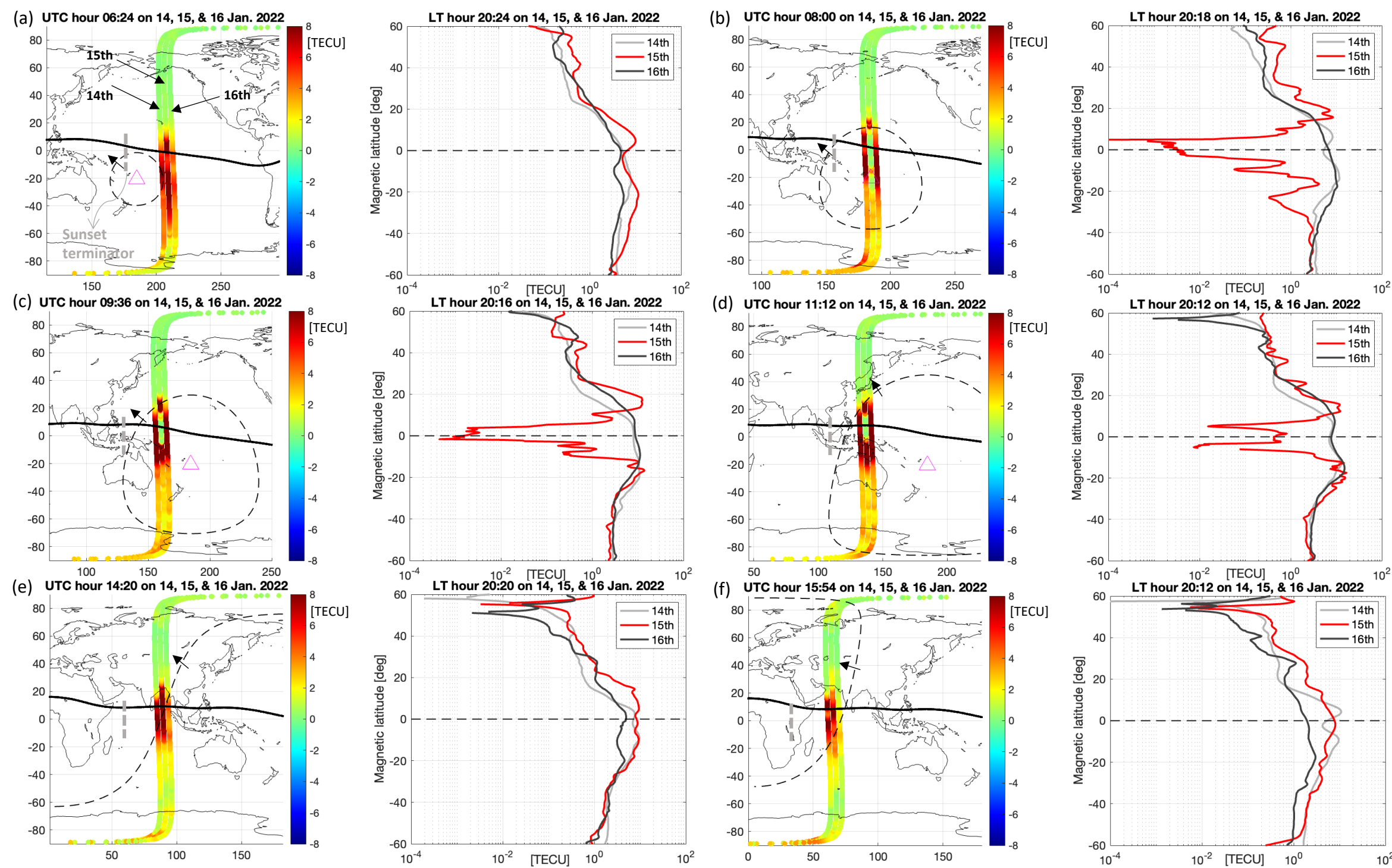
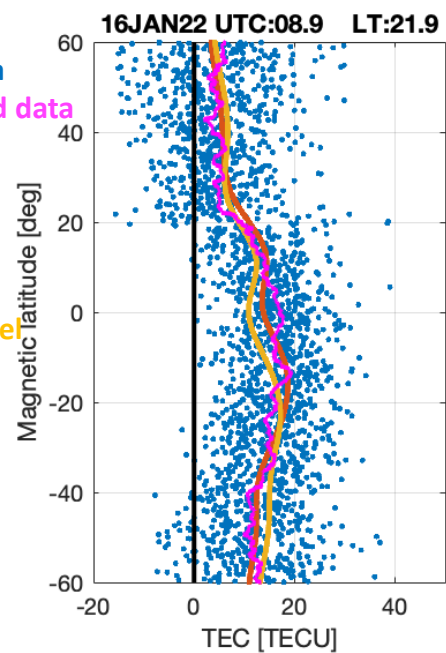
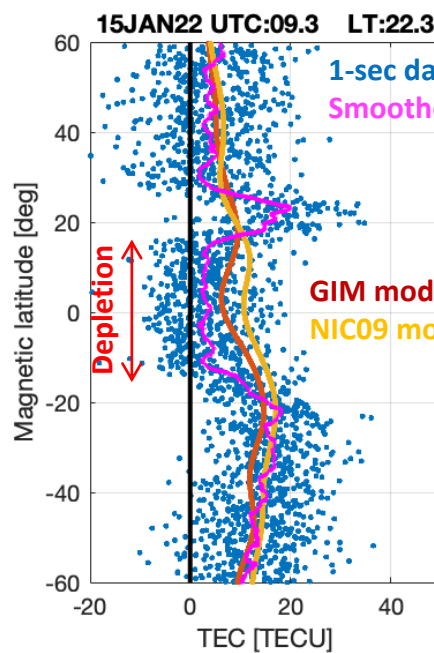
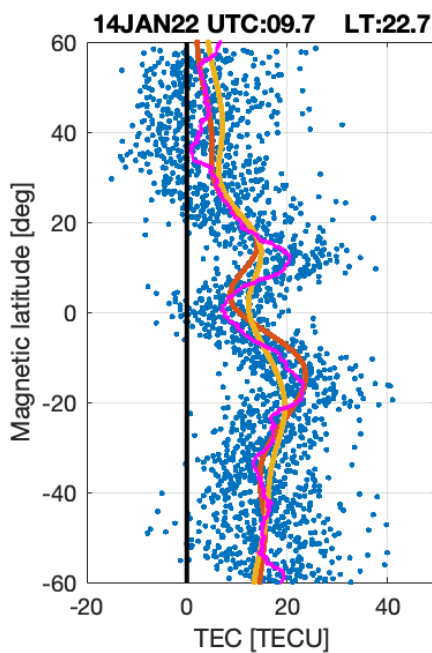
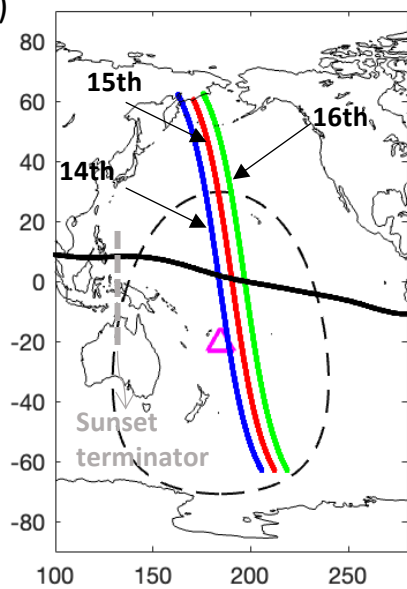


Figure 3.

(a)



(b)

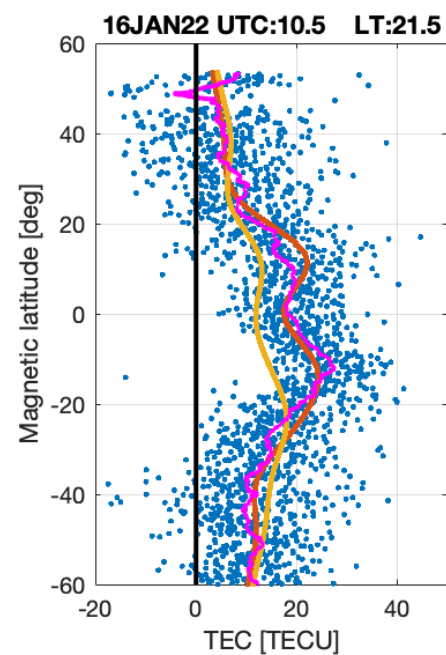
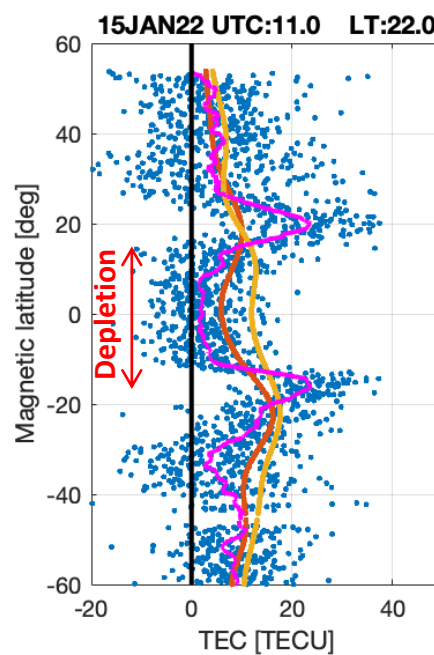
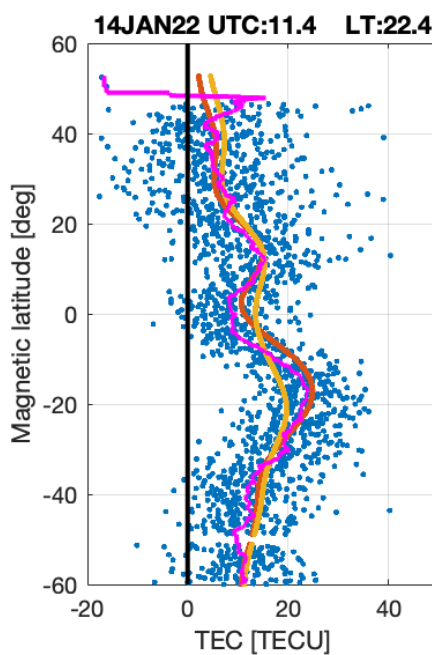
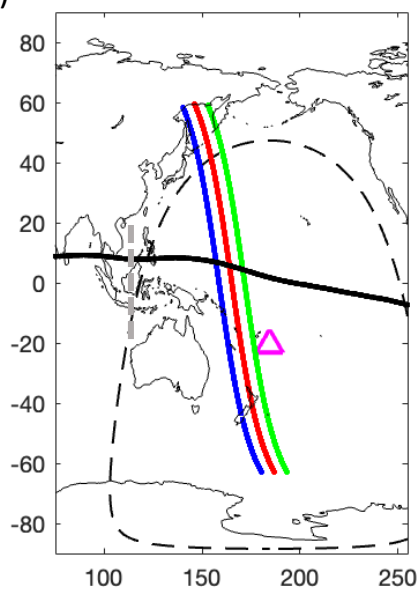


Figure 4.

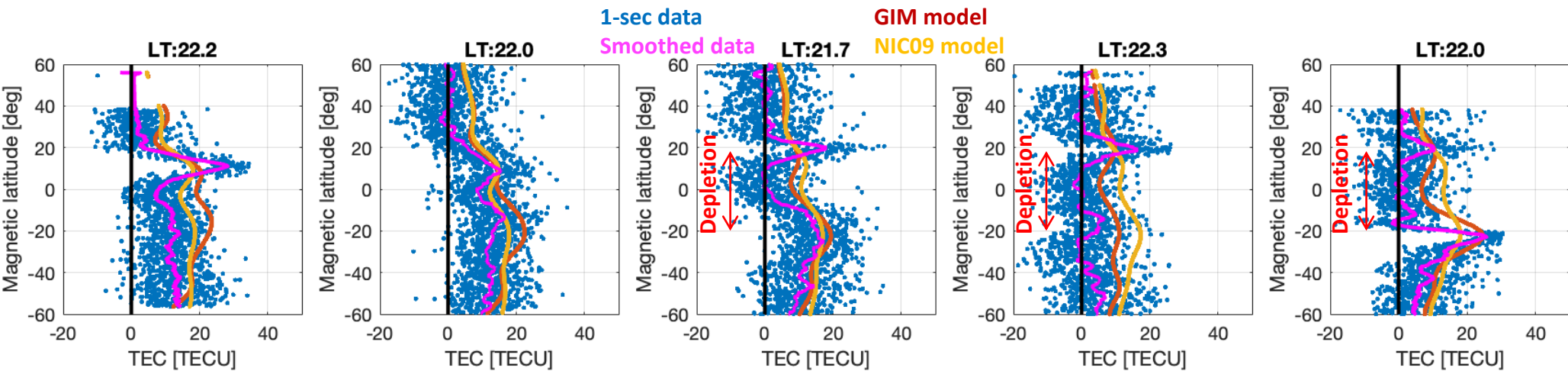
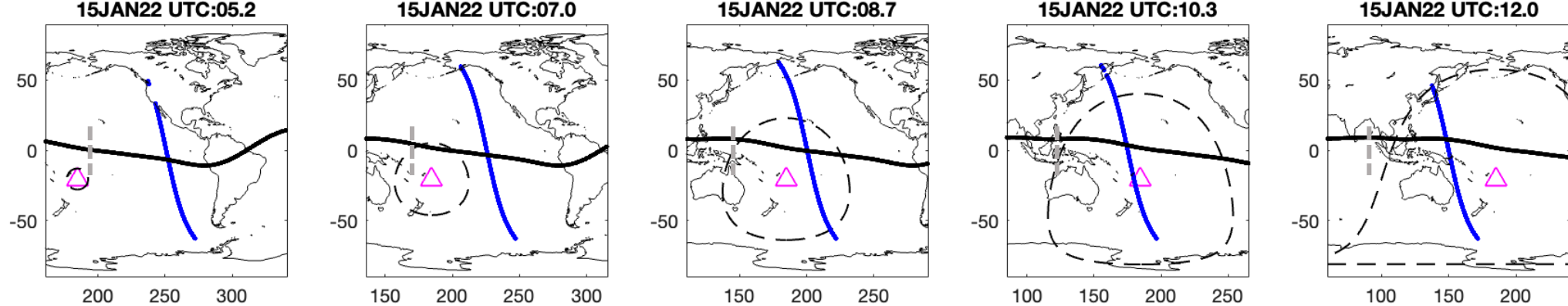
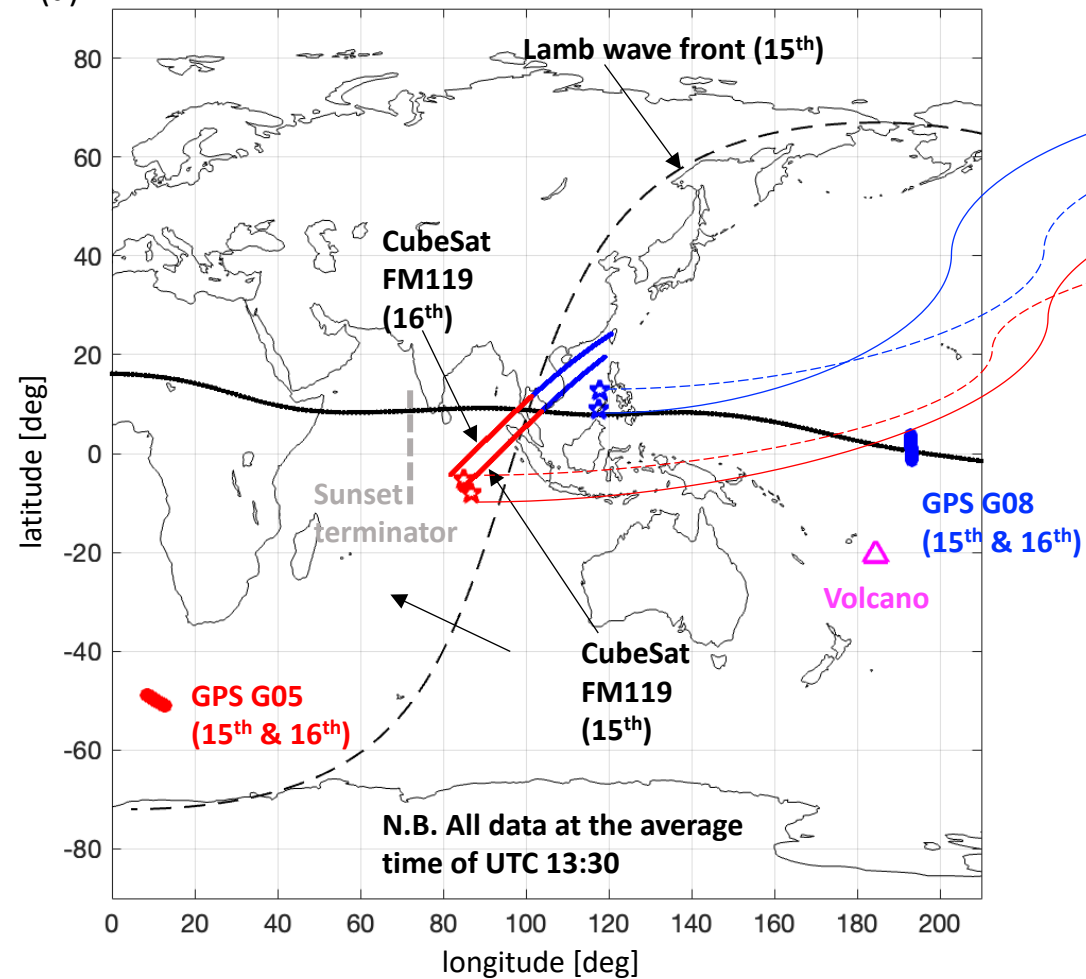


Figure 5.

(a)



(b)

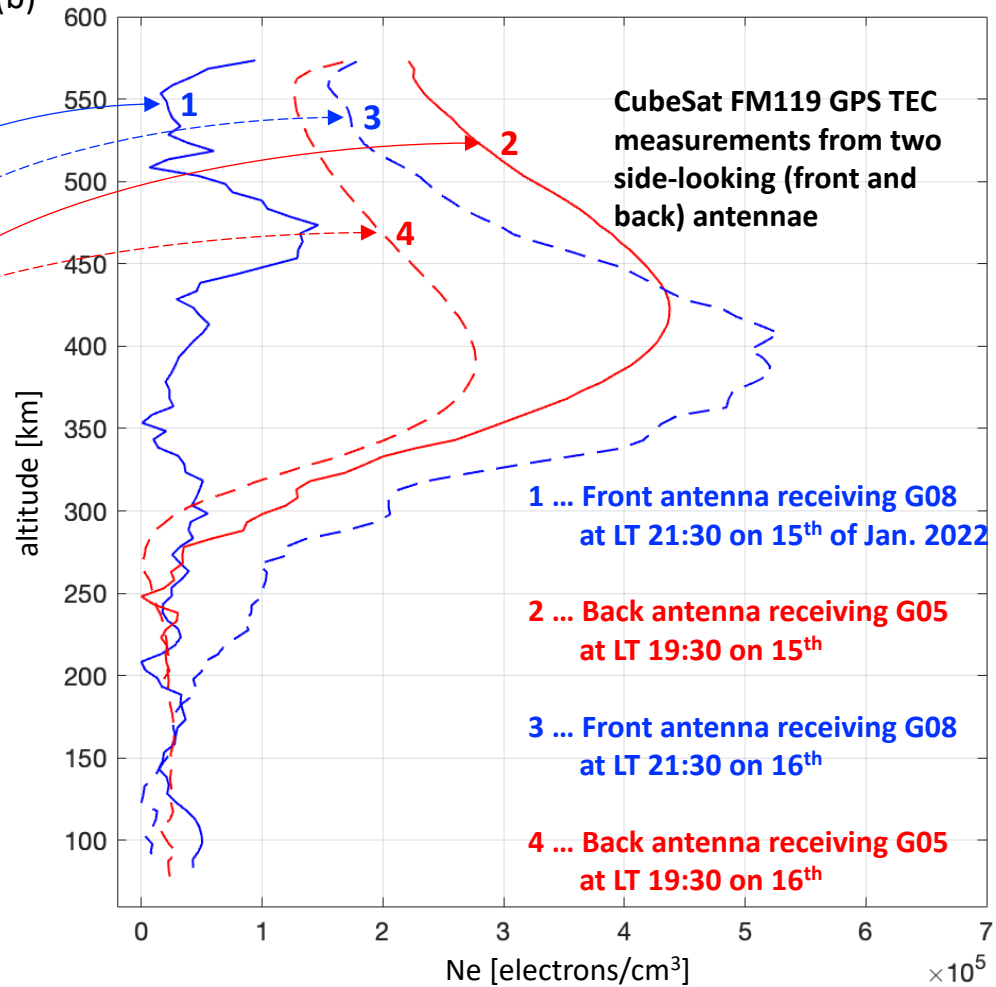
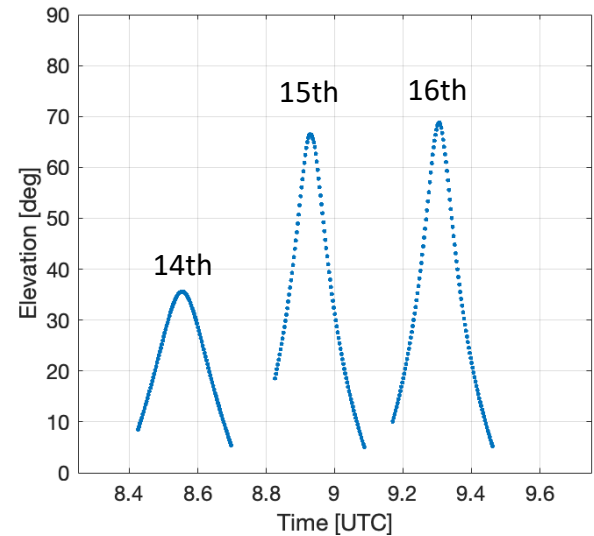
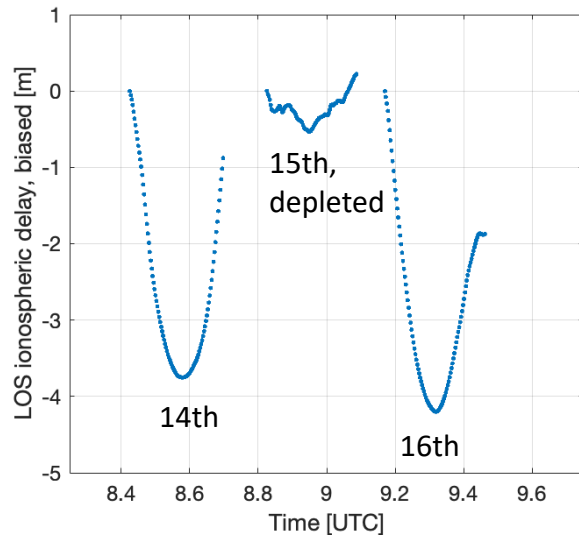
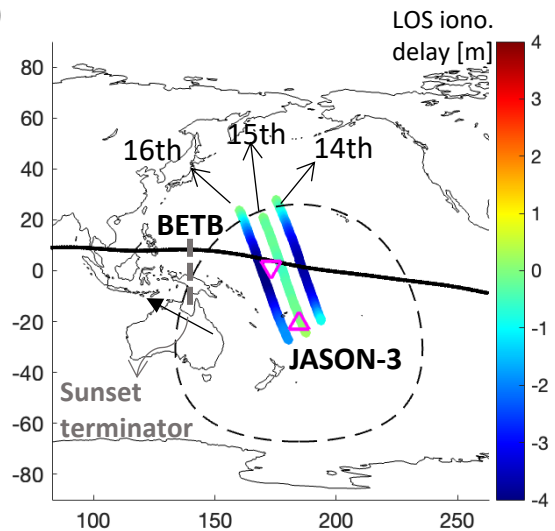


Figure 6.

(a)



(b)

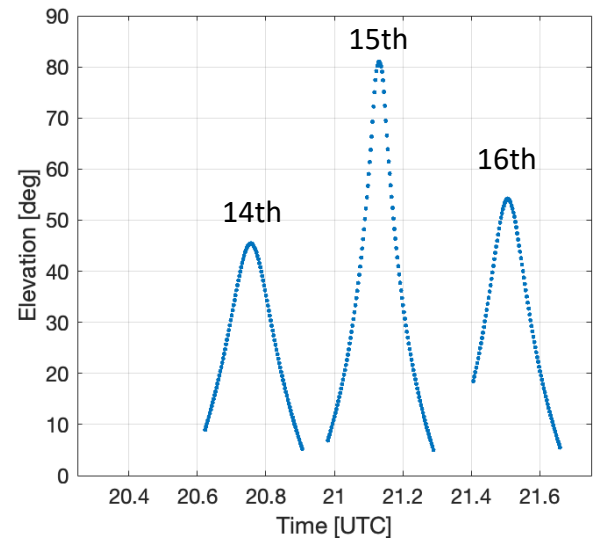
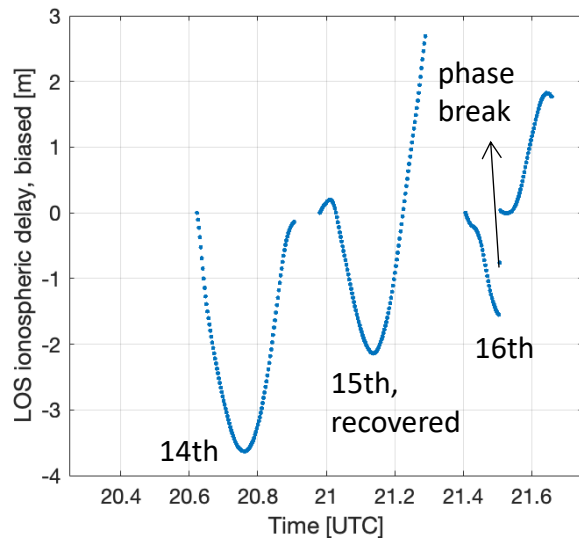
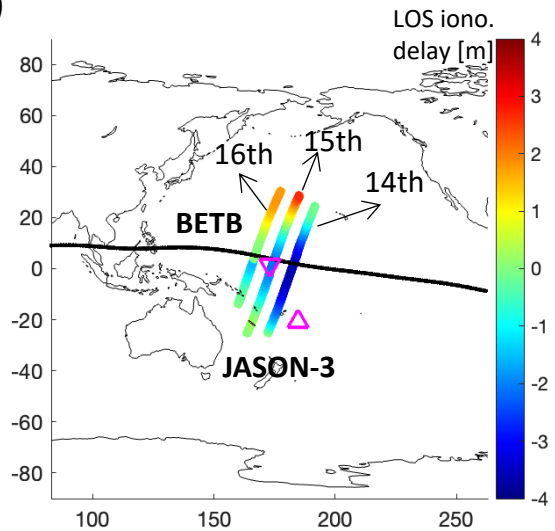


Figure 7.

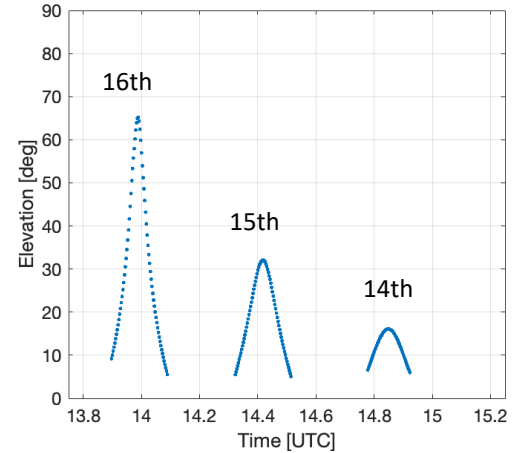
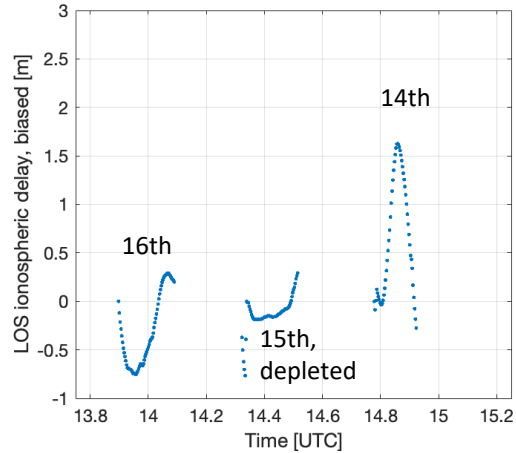
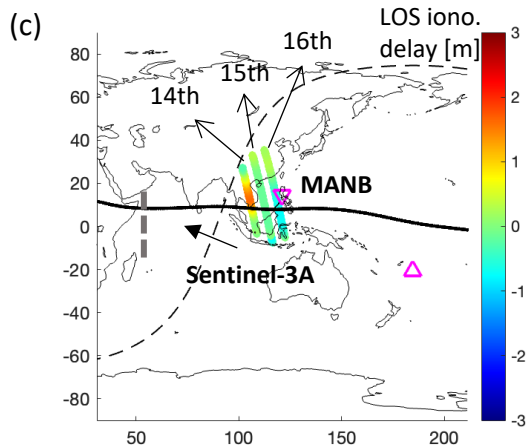
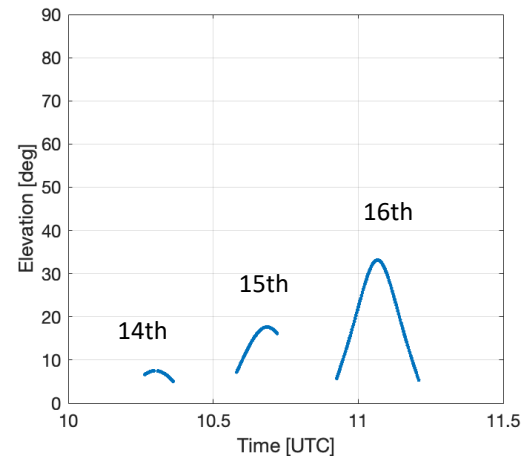
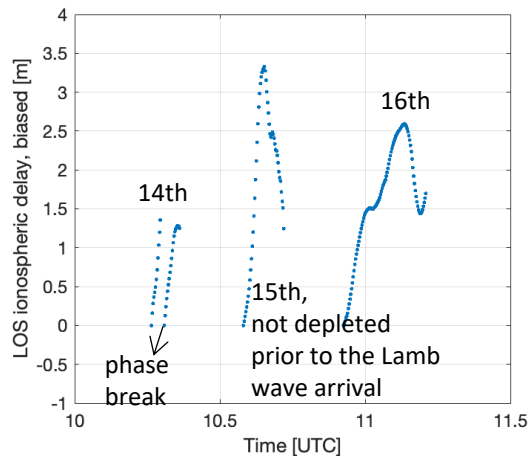
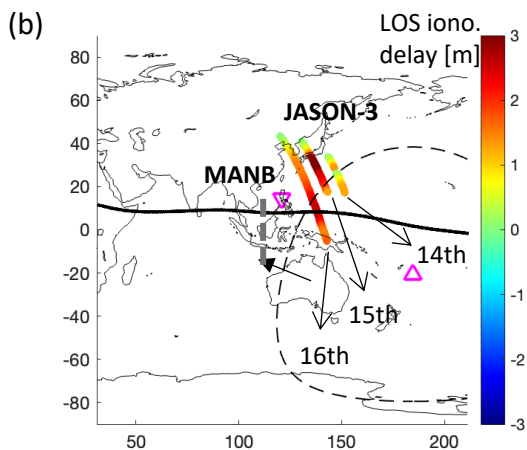
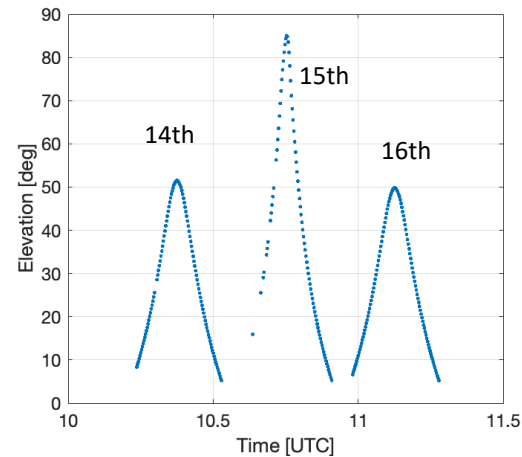
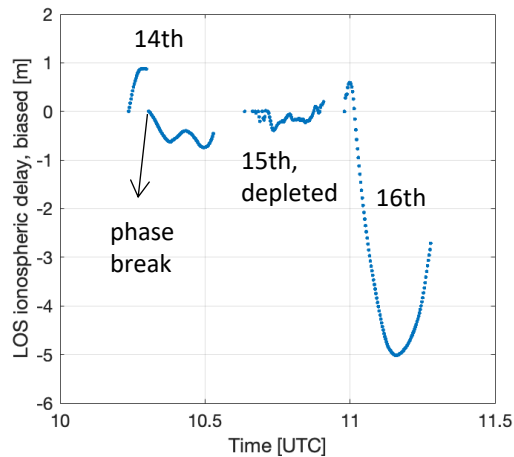
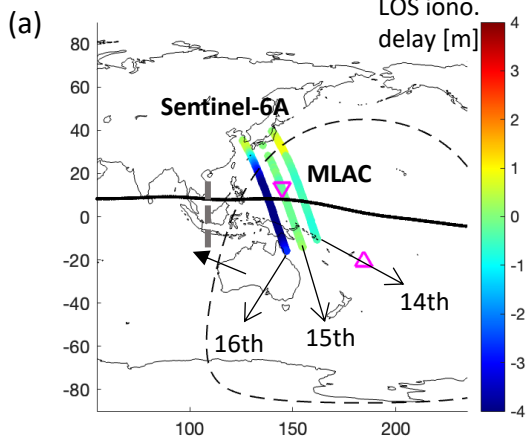
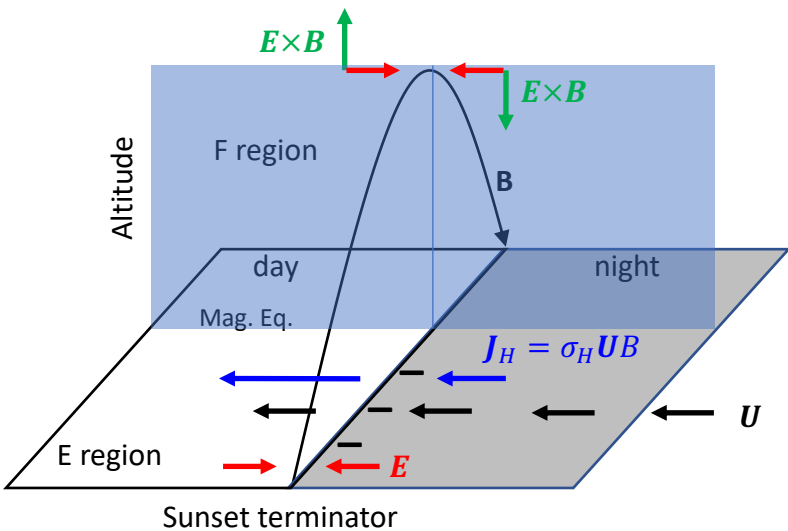


Figure 8.

(a) West



Tonga

(b) East

